

Hamiltonian Truncation and Effective Field Theory

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University of Geneva

based on work with Tim Cohen, Rachel Houtz, Markus Luty and Dorian Wenzel
arXiv: 2110.08273 and work in progress

Landscape of quantum field theory

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- QFT is nature's language

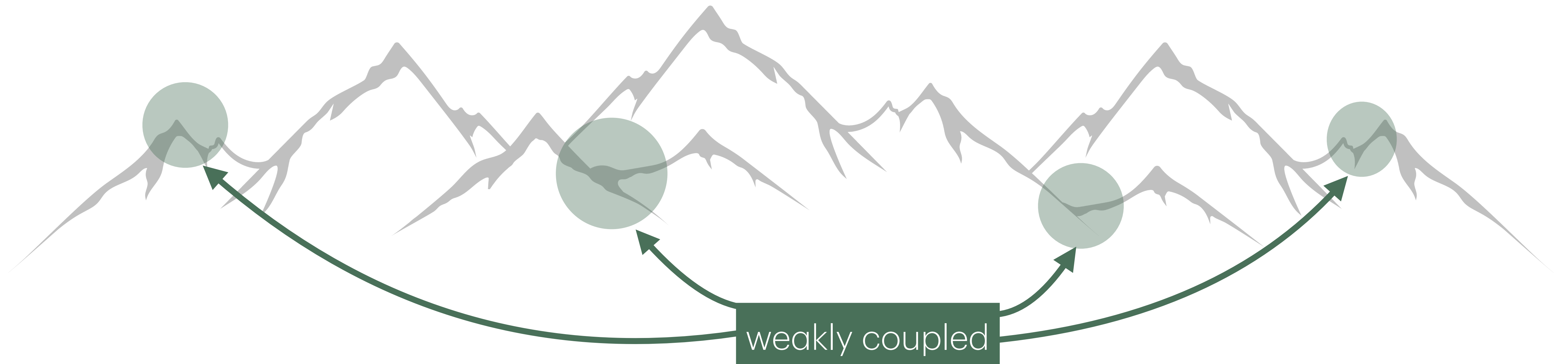
Landscape of quantum field theory

- QFT is nature's language
- vast landscape of possible QFTs



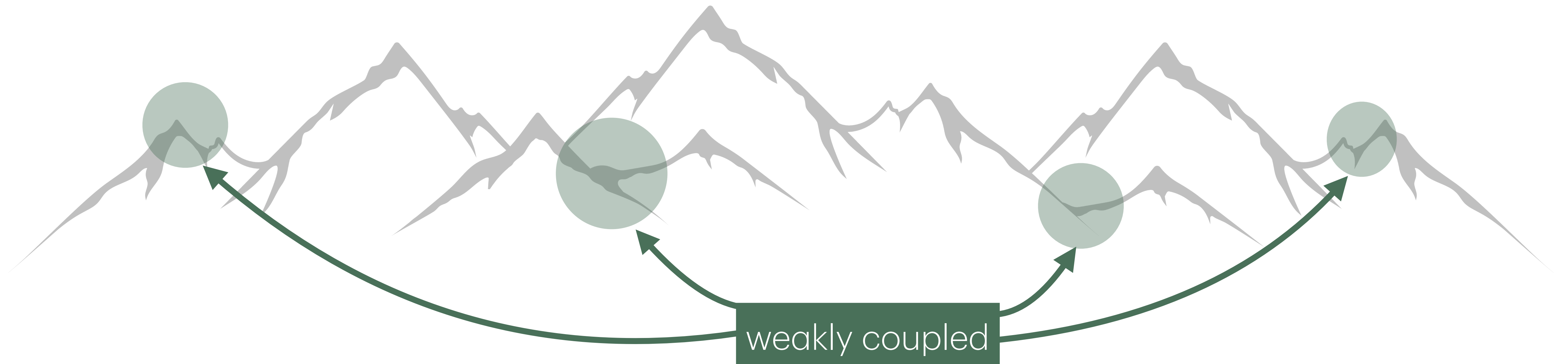
Landscape of quantum field theory

- QFT is nature's language
- vast landscape of possible QFTs
- powerful tools for weak coupling



Landscape of quantum field theory

- QFT is nature's language
- vast landscape of possible QFTs
- powerful tools for weak coupling
- need **new equipment** to explore whole landscape



QFT at strong coupling



weak coupling

perturbation theory



strong coupling

QFT at strong coupling

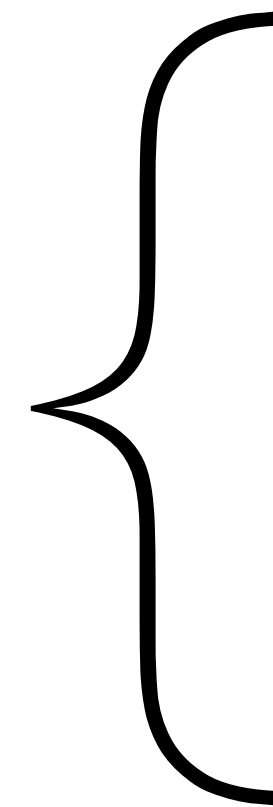


weak coupling

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strong coupling



Lattice methods

Conformal bootstrap

Hamiltonian truncation



QFT at strong coupling



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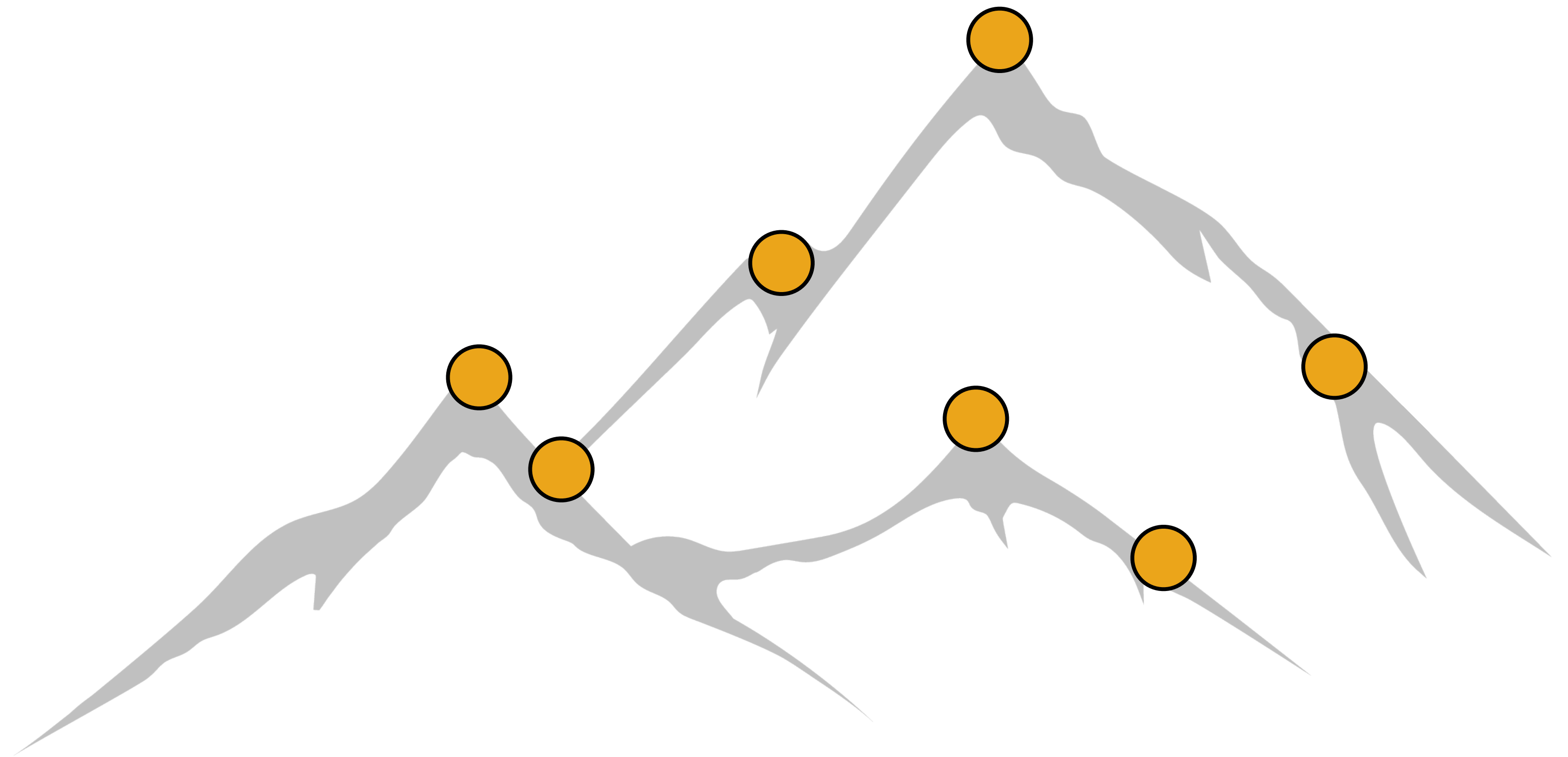


QFT = fixed points + flows



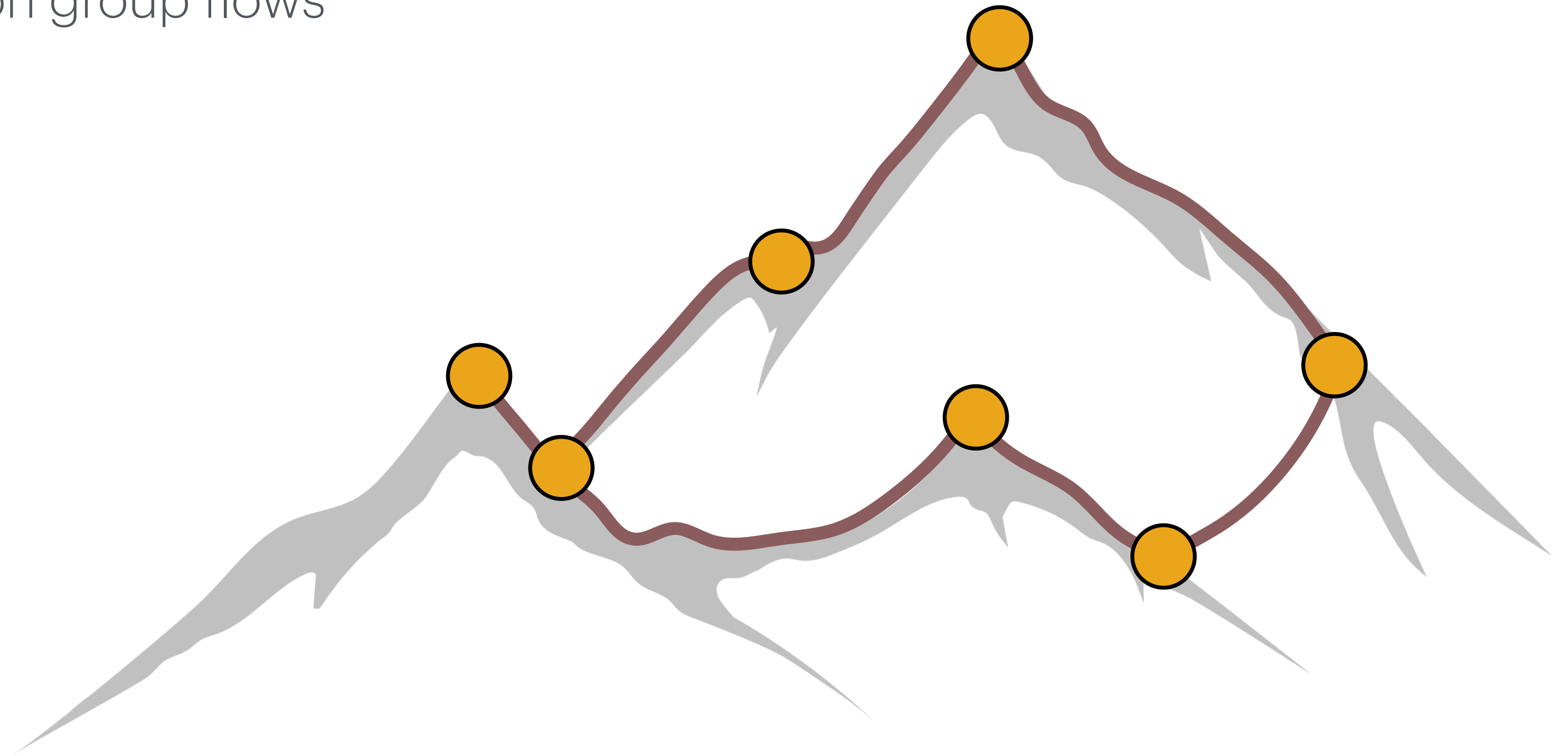
QFT = fixed points + flows

- special points in landscape = fixed points



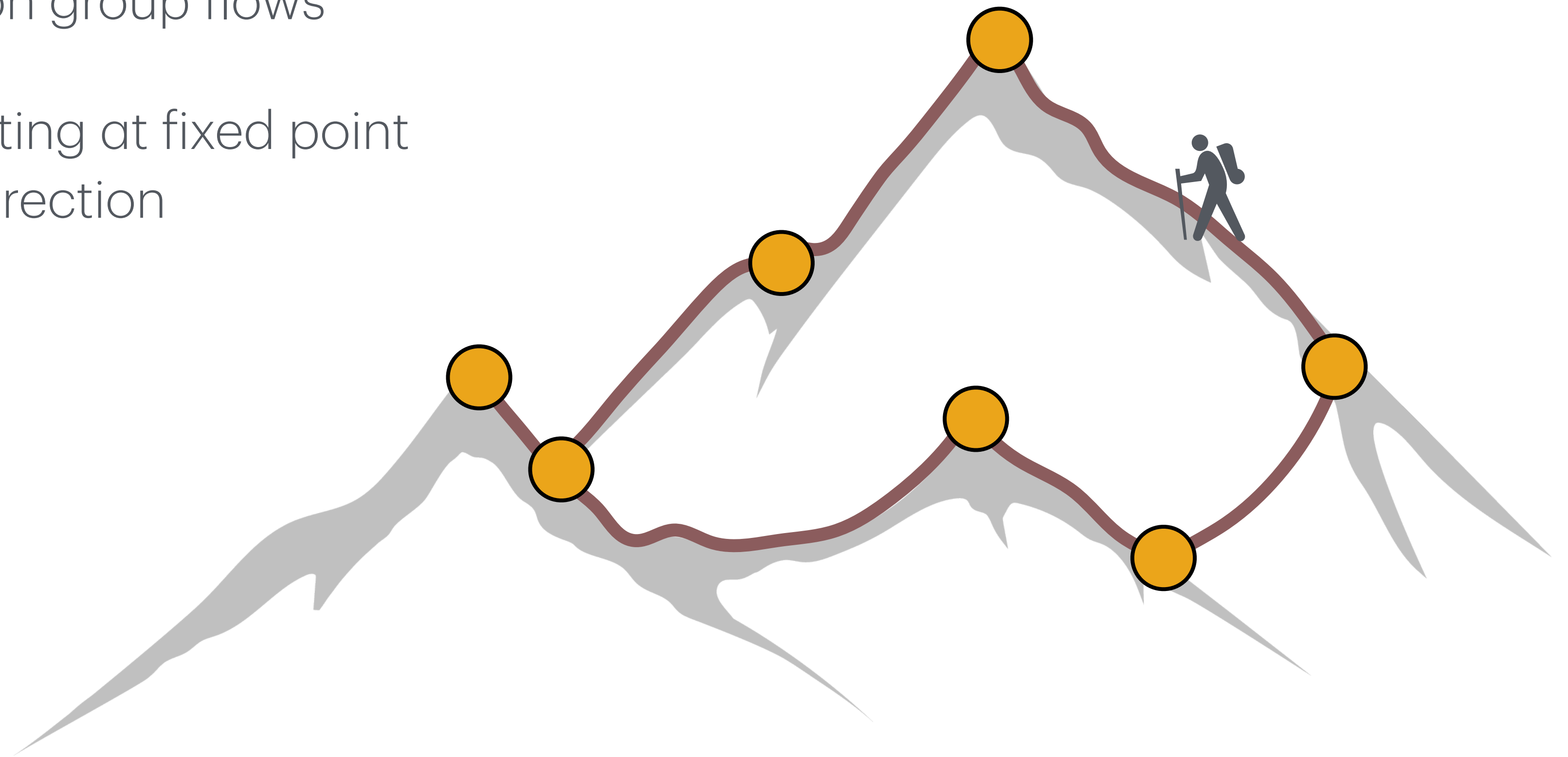
QFT = fixed points + flows

- special points in landscape = fixed points
- connected by renormalization group flows



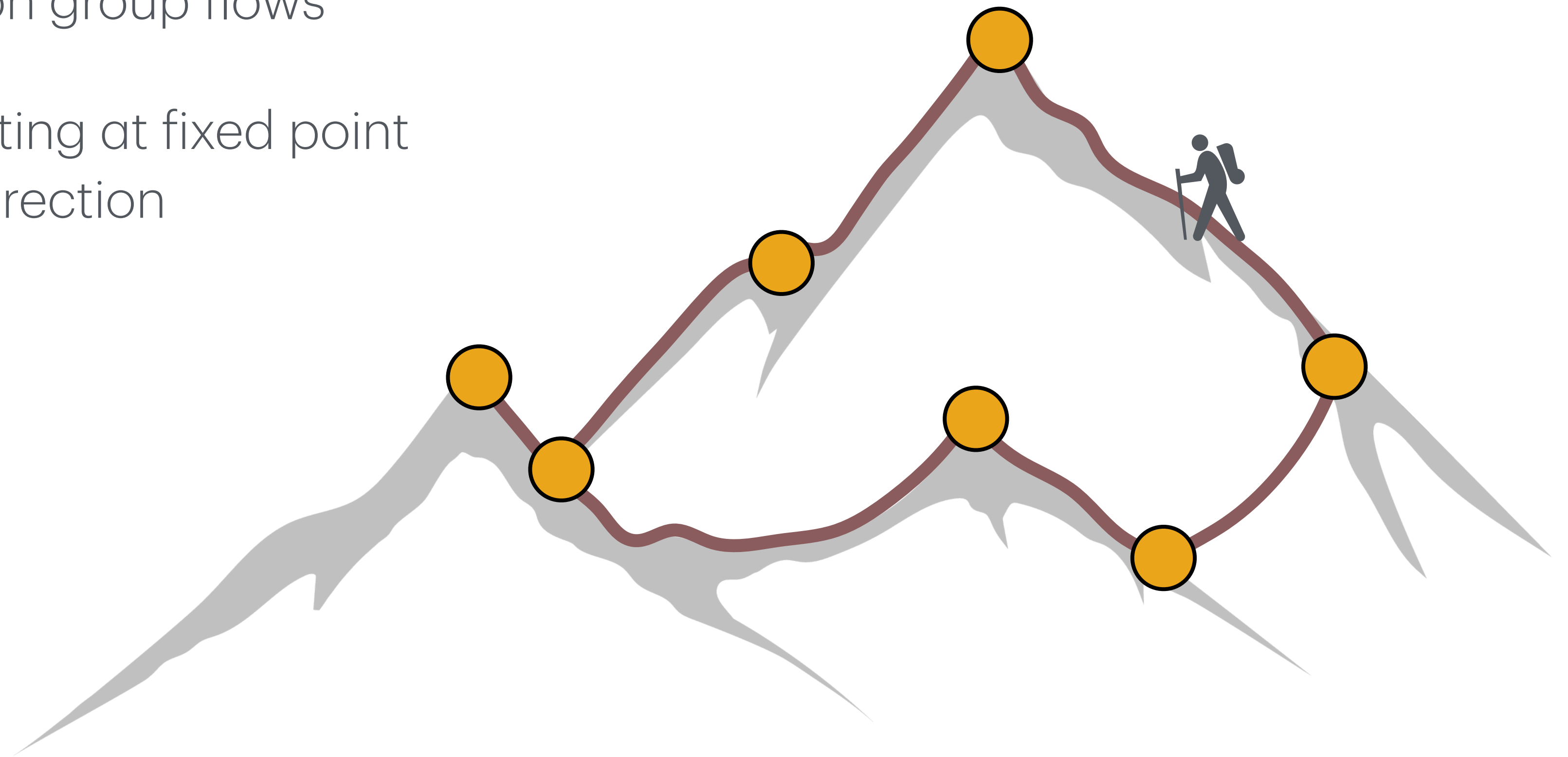
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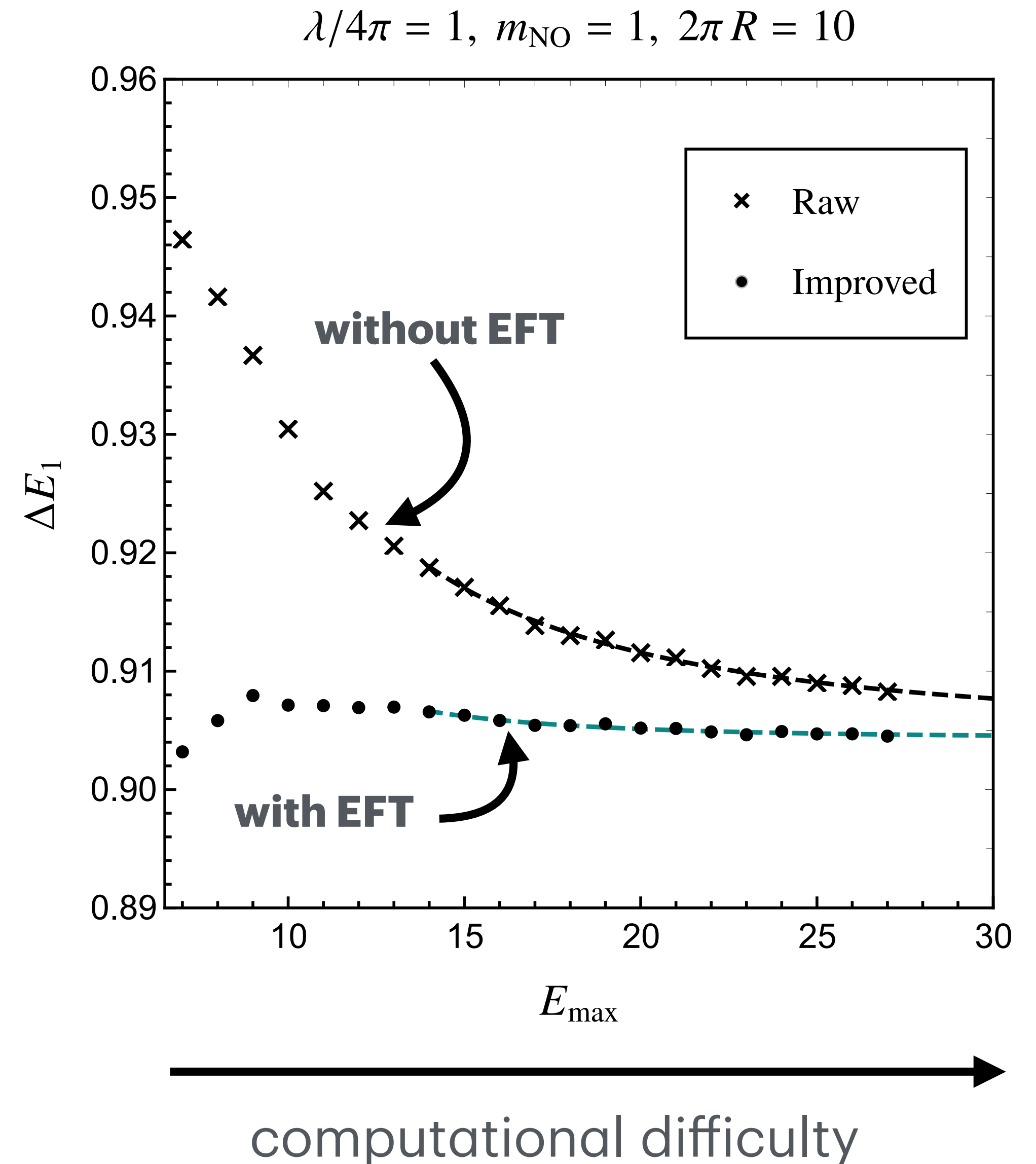
- special points in landscape = fixed points
- connected by renormalization group flows
- can think of **any** QFT as starting at fixed point and flowing in a particular direction
- captures intuition
 - universality
 - relevant vs. irrelevant



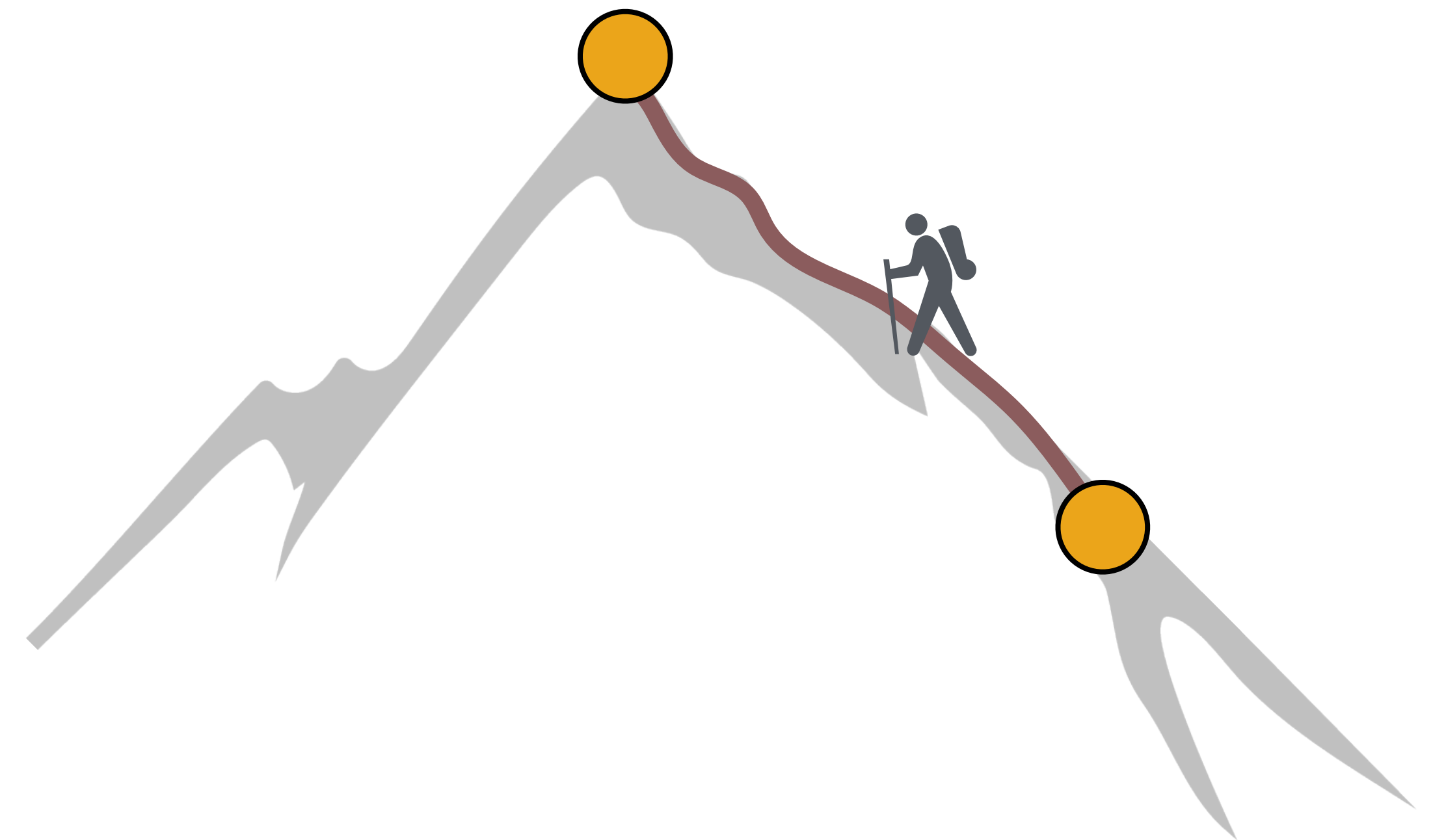
Hamiltonian truncation solidifies this conceptual picture

Punchline

- Hamiltonian truncation = non-perturbative method for computing observables in strongly coupled QFTs
- effective field theory = powerful tool for compensating for ignorance
- effective field theory techniques **drastically improve** convergence in Hamiltonian truncation calculations

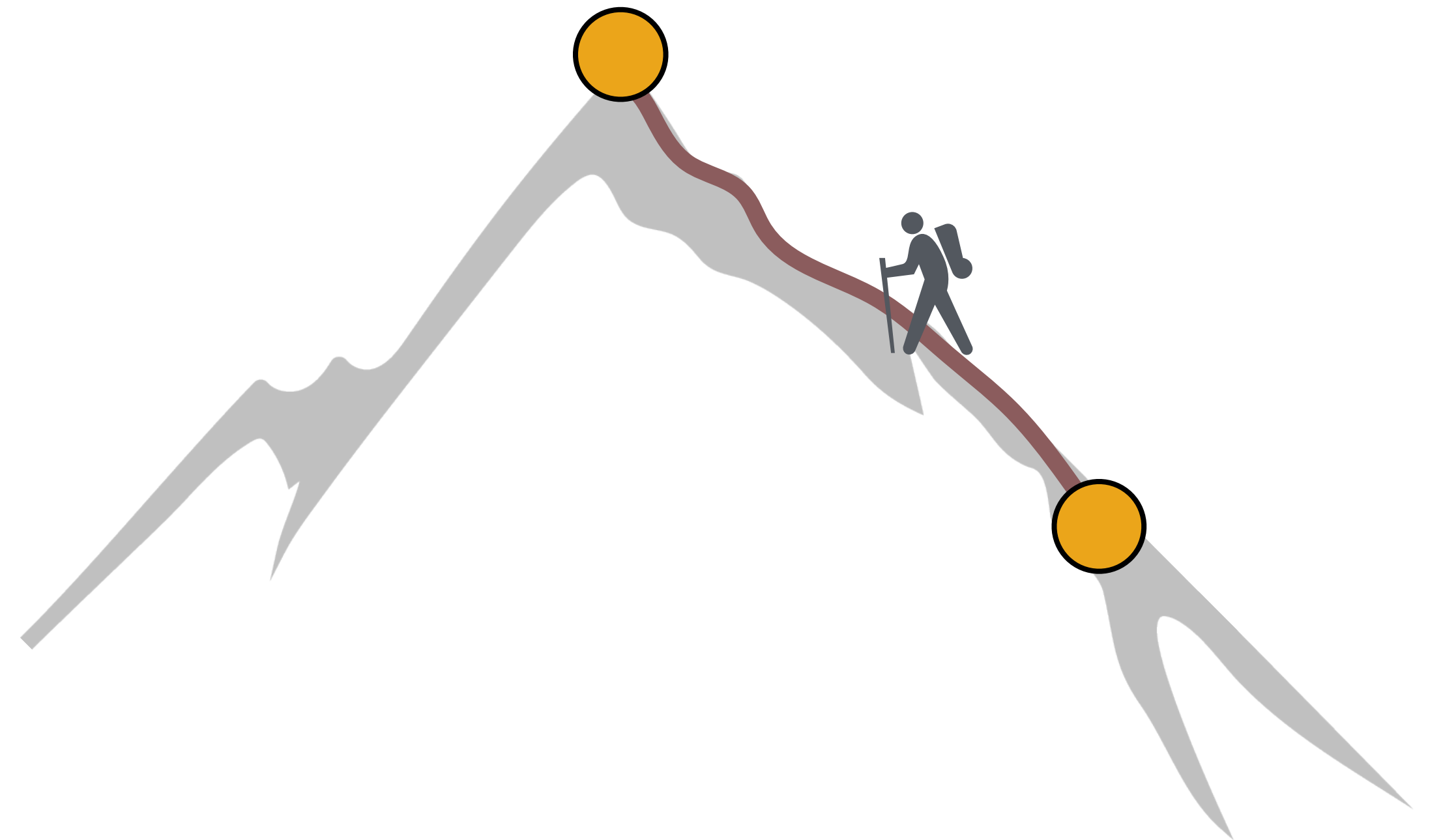


Hamiltonian truncation



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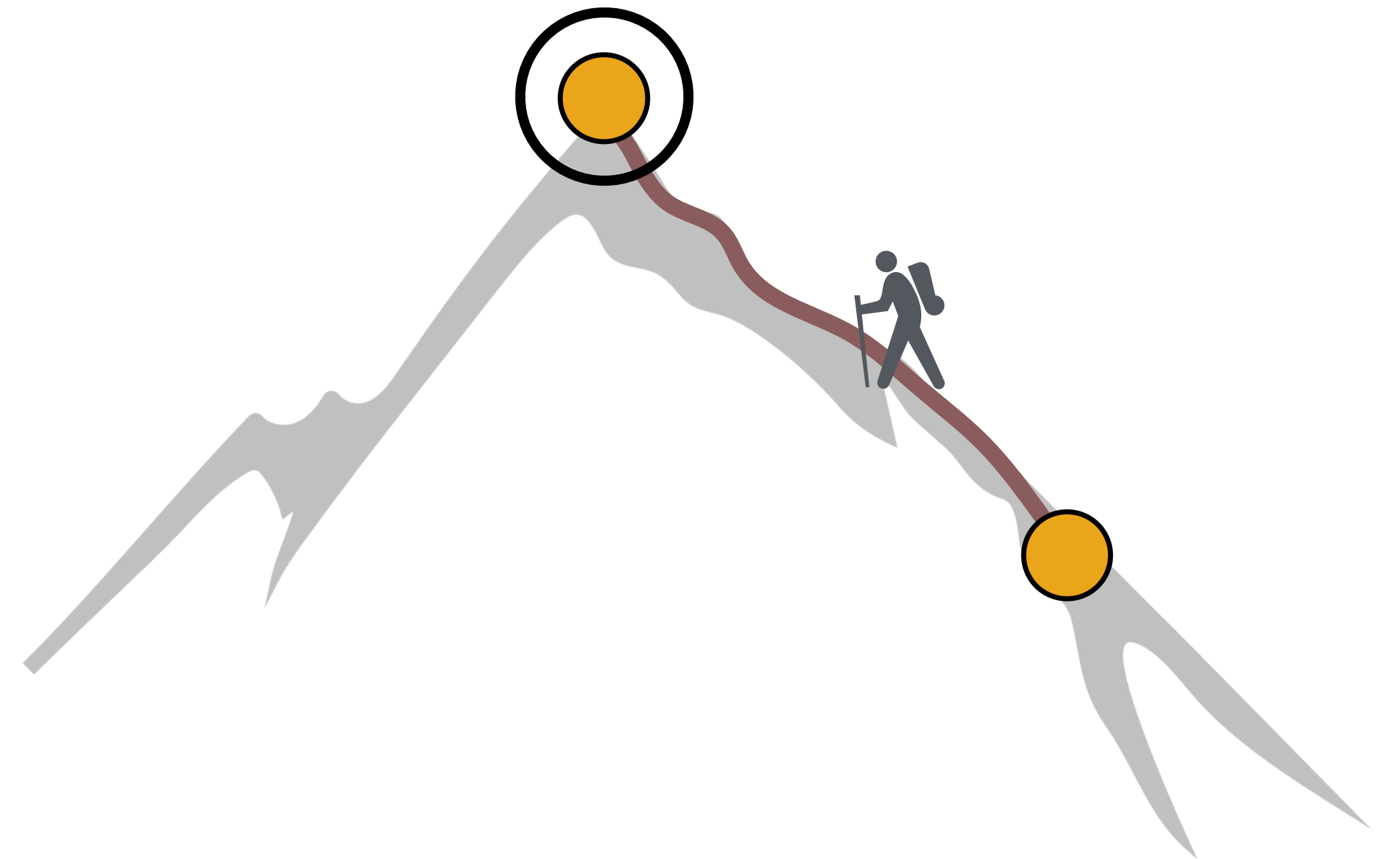
$$H = H_0 + V$$



Hamiltonian truncation

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diagonalizable
(e.g. fixed point)

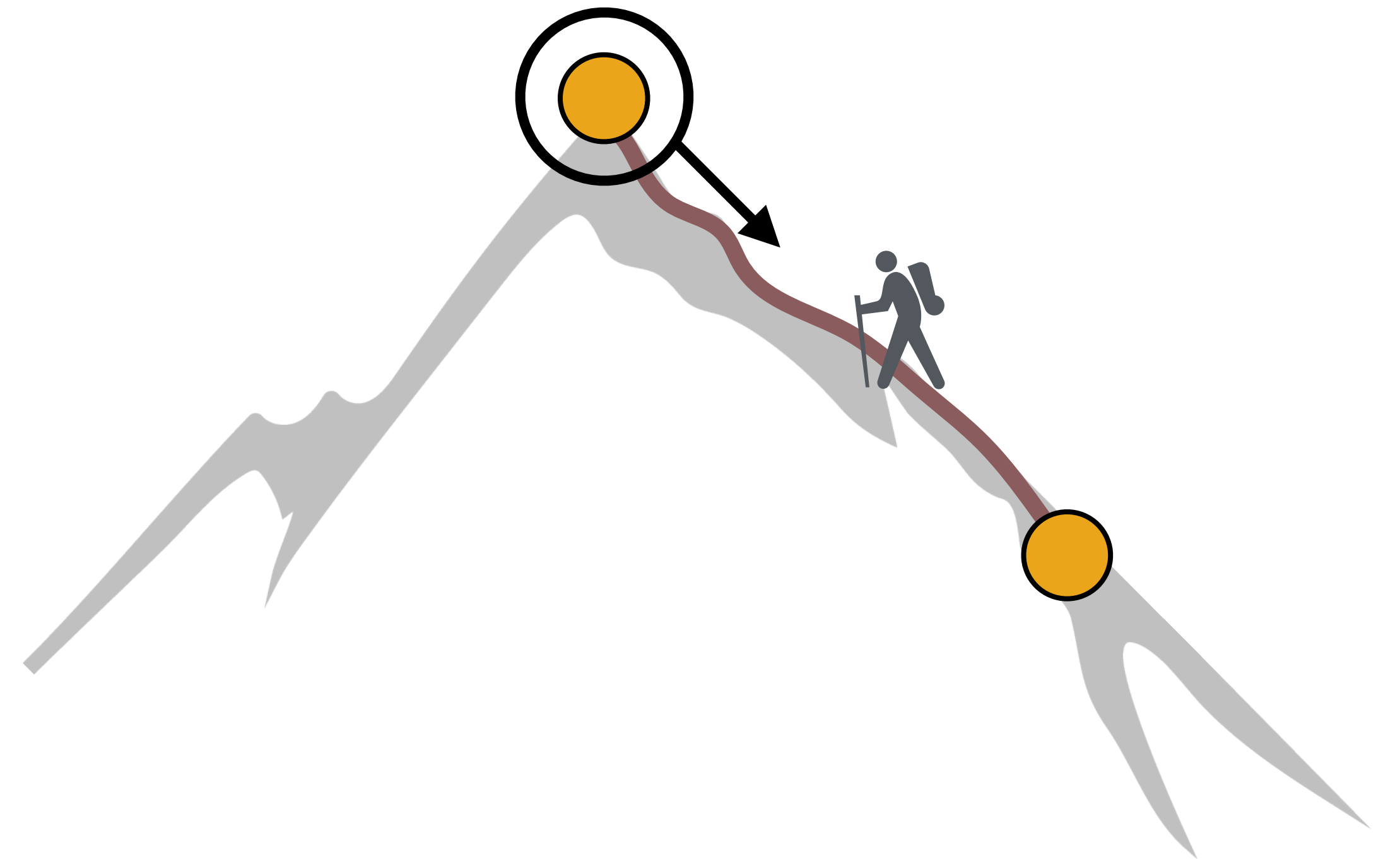


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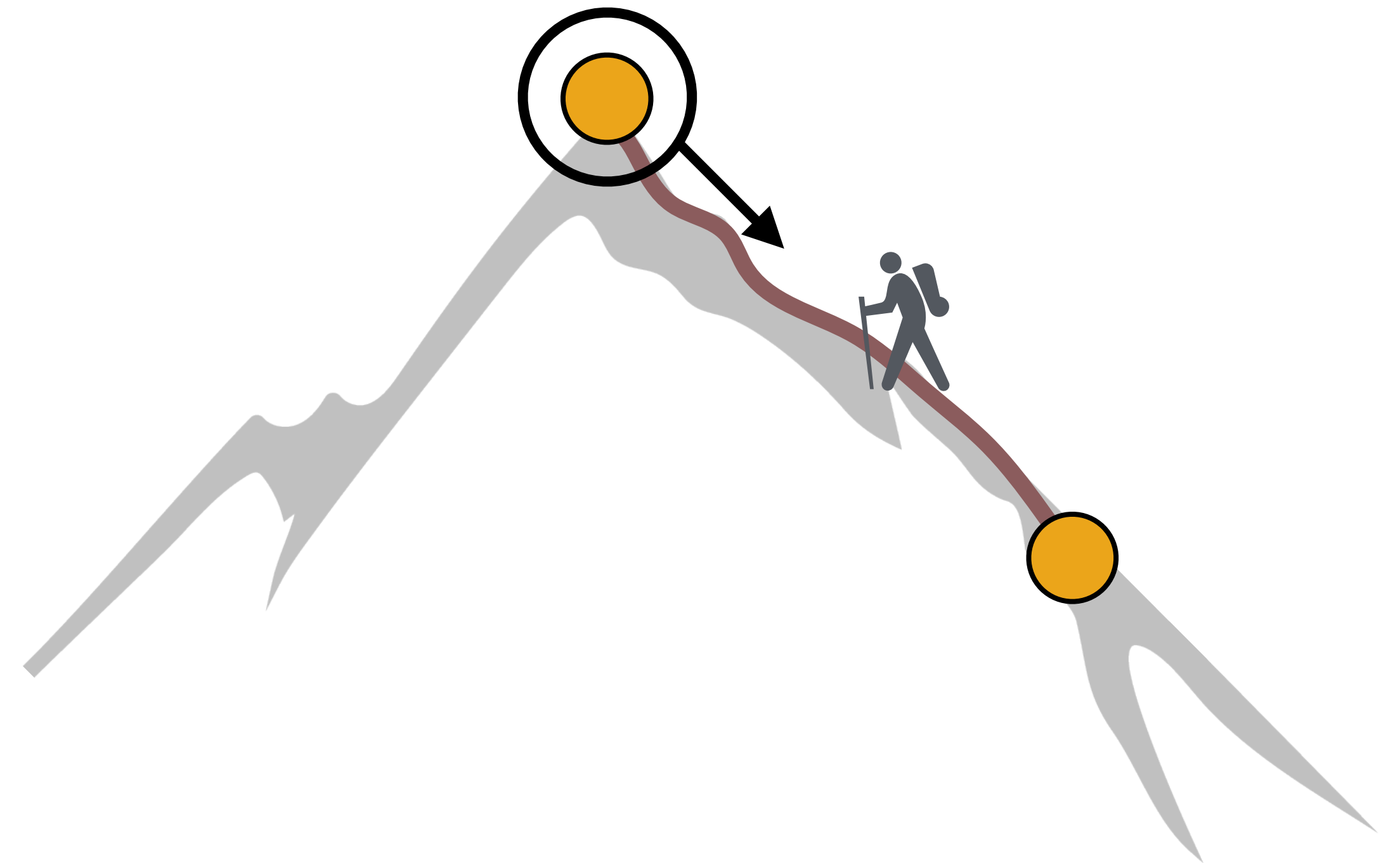


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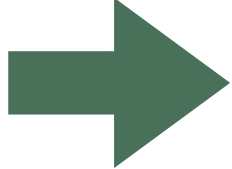
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can't diagonalize full H → truncate!

Hamiltonian truncation in practice

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- discretize  finite volume

- lots of approaches: DLCQ (Pauli et al '85), TCSA (Yurov, '90), Massive Fock Space (Brooks et al '84, Rychkov et al '14), LCT (Katz et al '16), RCMPS (Tilloy '21), ...

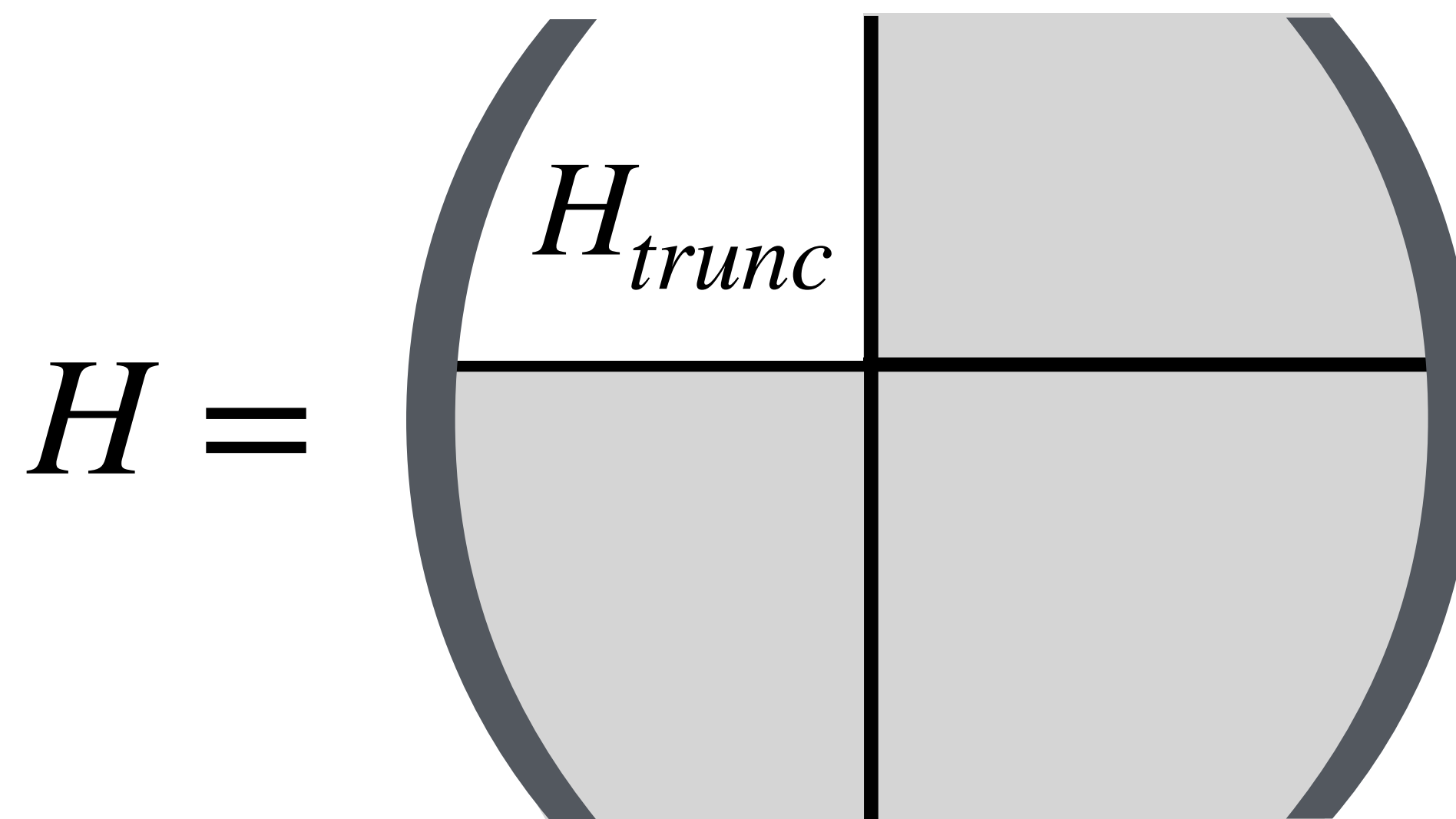
Hamiltonian truncation in practice

- discretize \rightarrow finite volume
- truncate \rightarrow separate Hamiltonian

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(small in UV, big in IR)



$$H_0 |i\rangle = E_i |i\rangle$$

$$E_i \leq E_{max}$$

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Hamiltonian truncation in practice

- discretize \rightarrow finite volume
- truncate \rightarrow separate Hamiltonian
- diagonalize \rightarrow spectrum (finite volume)

$$H = \begin{array}{|c|c|} \hline H_{trunc} & \\ \hline & \\ \hline \end{array}$$

$$H = H_0 + V$$

diagonalizable
(e.g. fixed point)

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Example: anharmonic oscillator

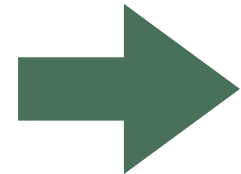
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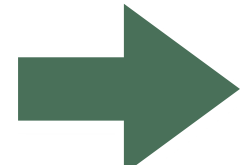
$|0\rangle, |1\rangle, |2\rangle, \dots$

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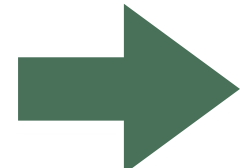
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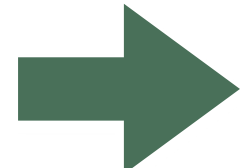
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$$\langle m | V | n \rangle = \frac{\lambda}{4} \langle m | (a + a^\dagger)^4 | n \rangle$$

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- diagonalize!

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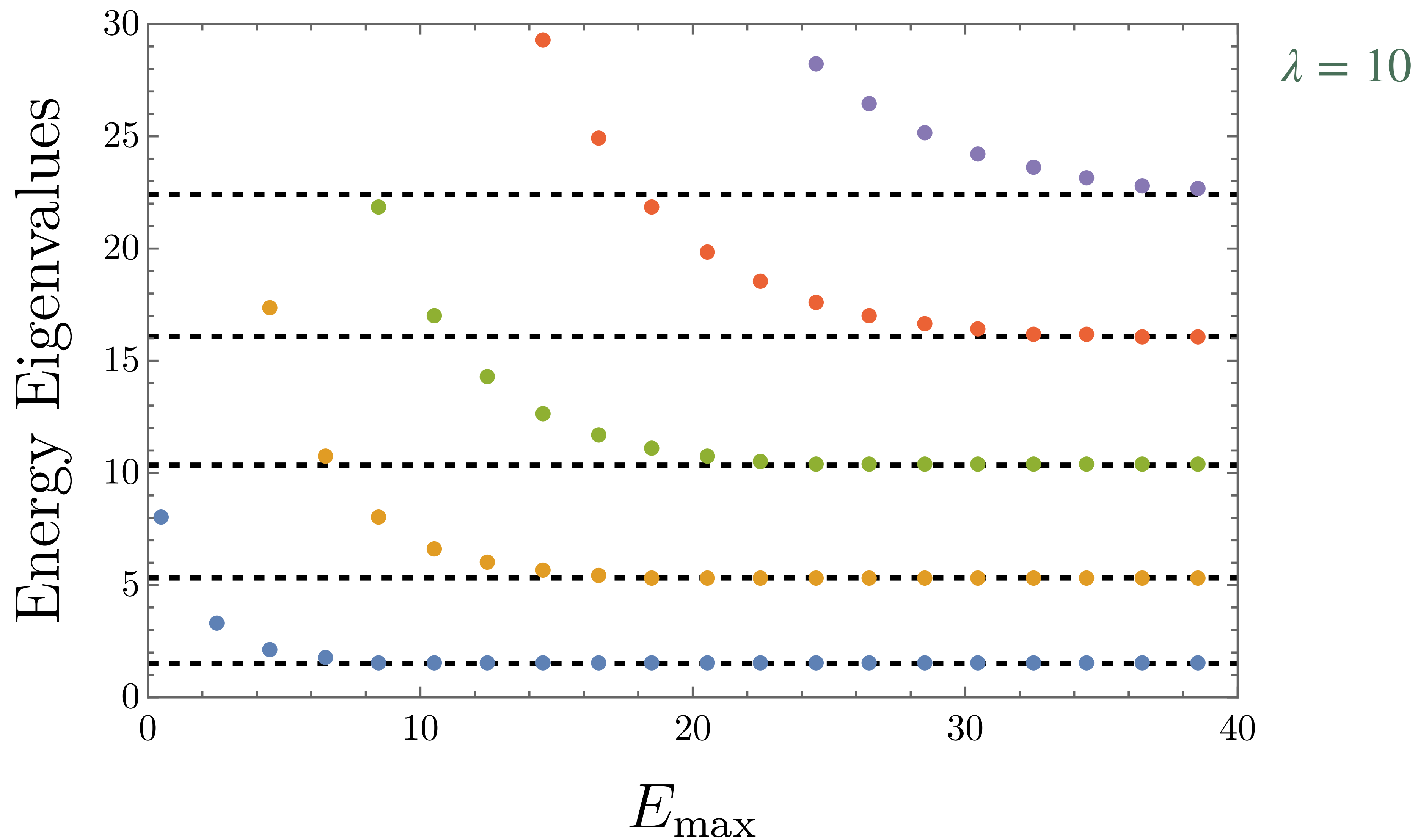
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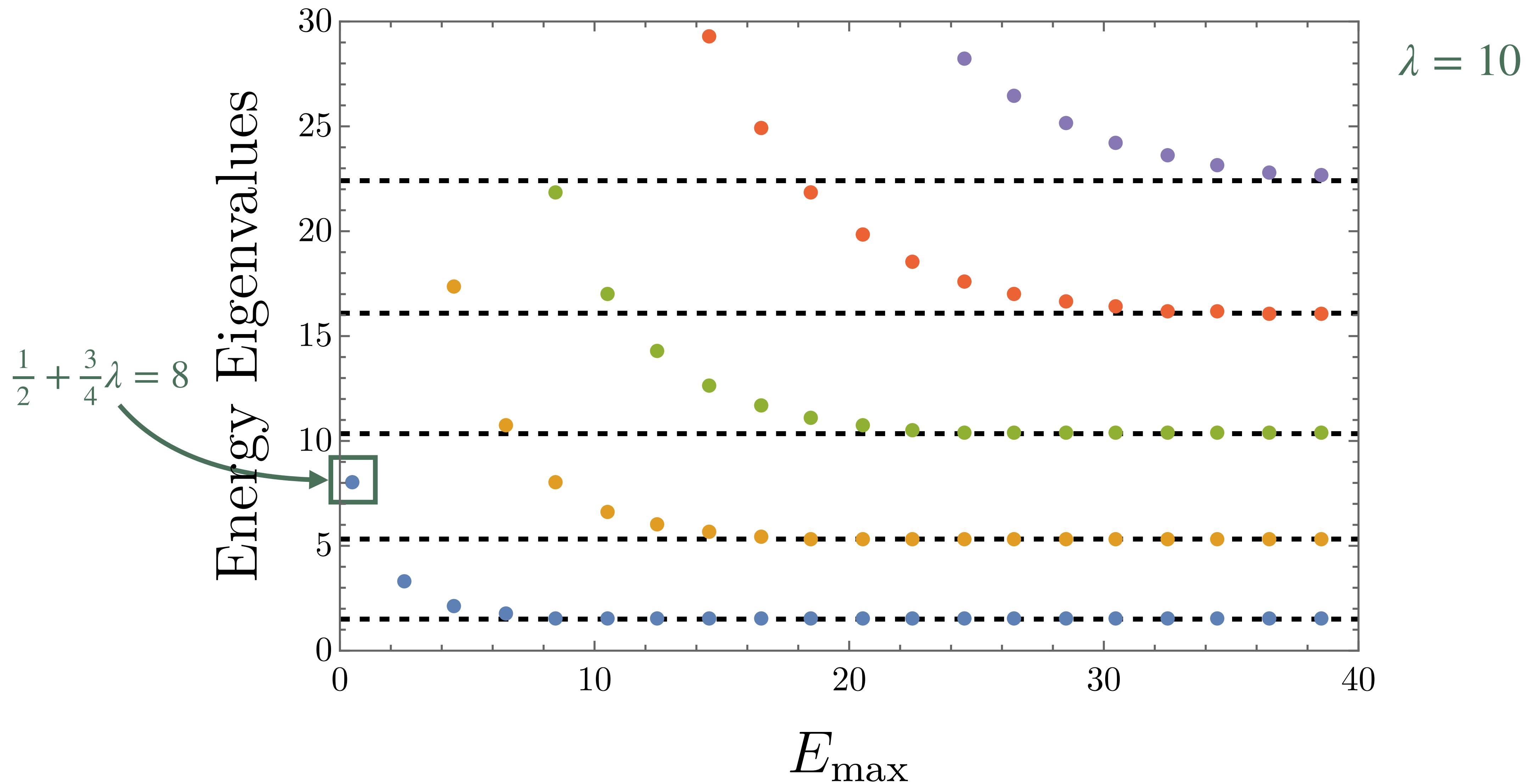
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- diagonalize \rightarrow eigenvalues: $\frac{3}{2} + \frac{21}{4}\lambda \pm \sqrt{1 + \frac{9}{4}\lambda(11\lambda + 4)}$

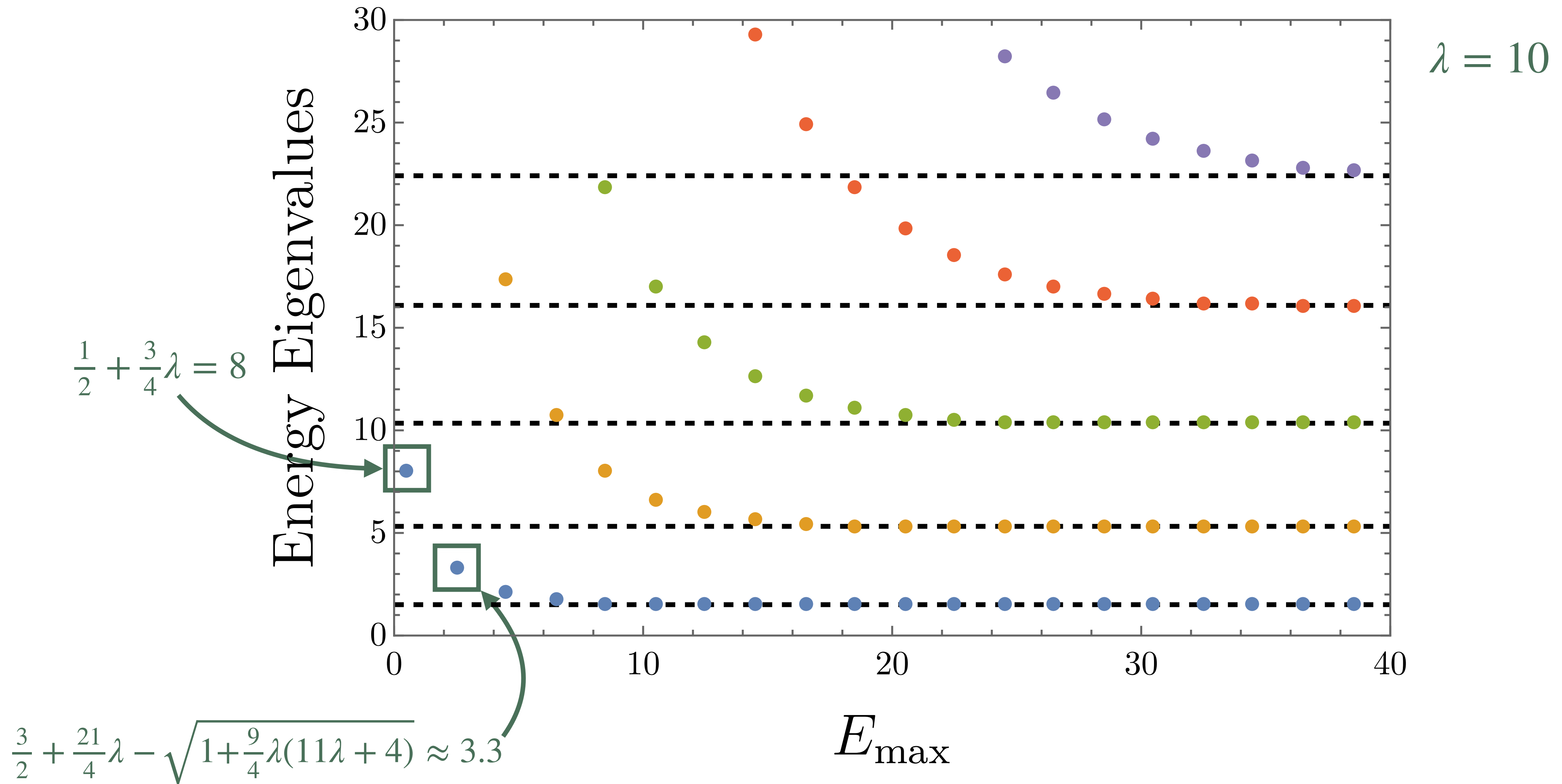
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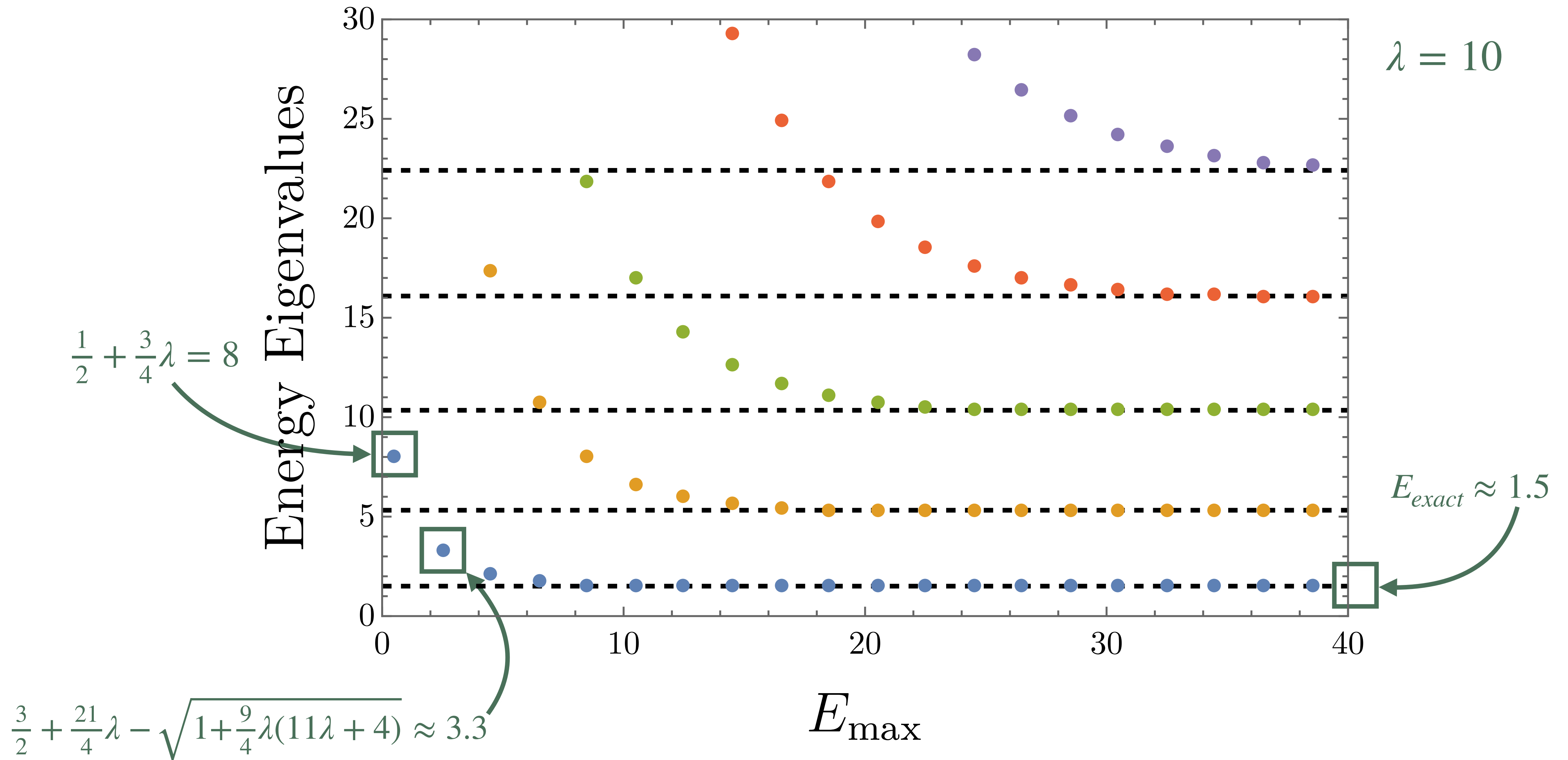
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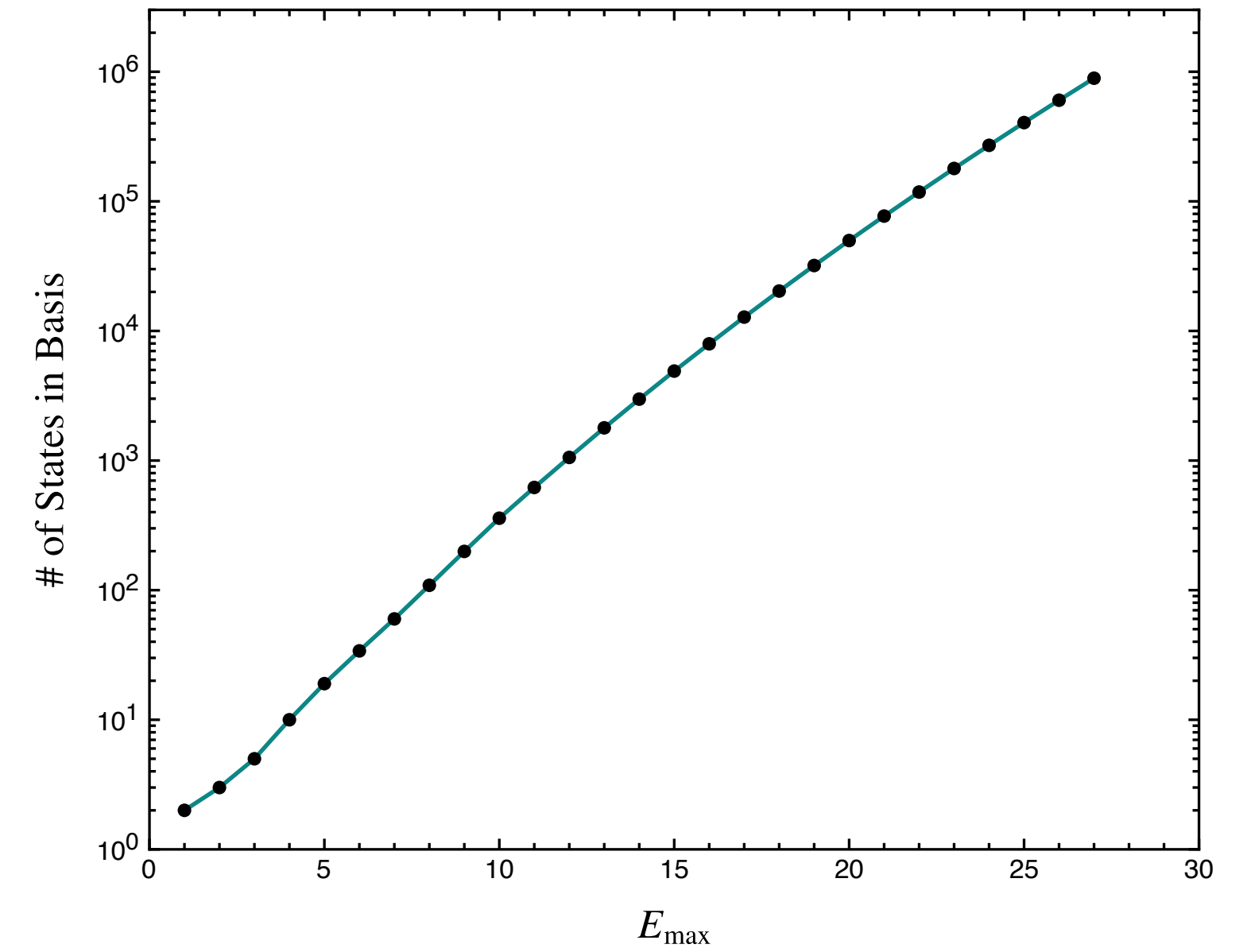
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- direct access to **dynamics** ($i\partial_t\Psi = H\Psi$)

Just one potential issue

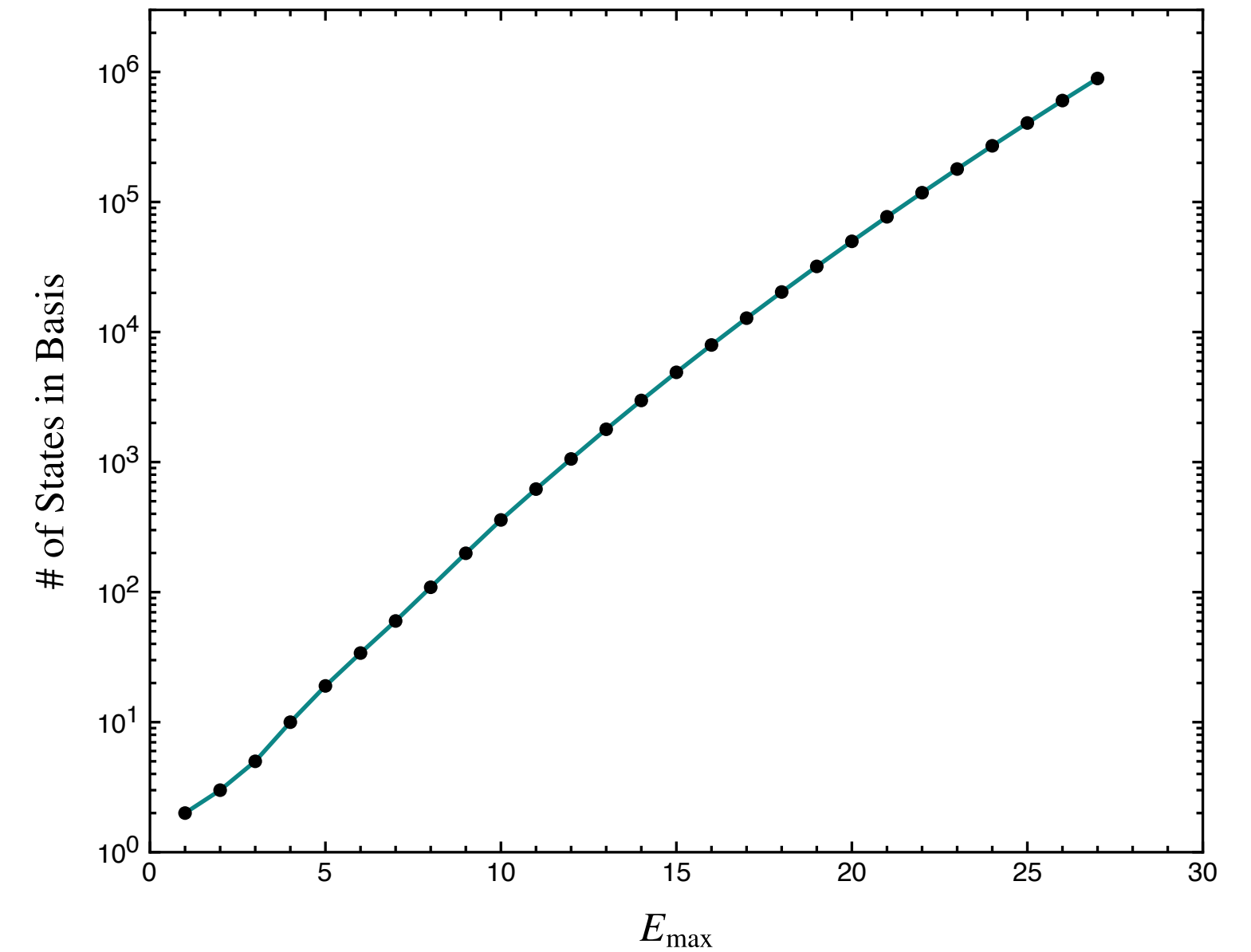
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- add corrections to account for effects from states outside truncated Hilbert space (“integrate out”)

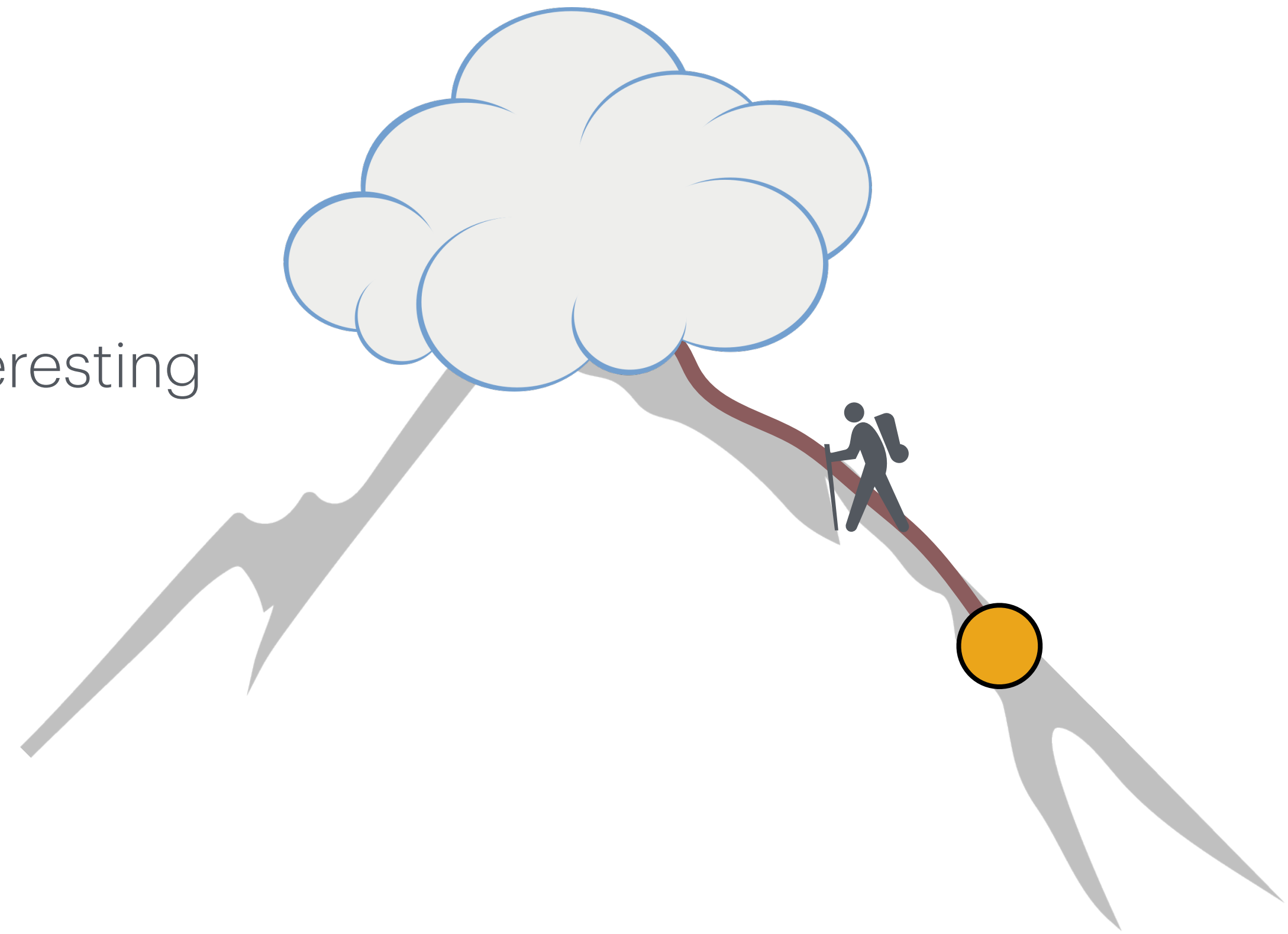


$$H = \left(\begin{array}{c|c} H_{trunc} & \\ \hline & \end{array} \right) + \left(\begin{array}{c|c} H_{corr} & \\ \hline & \end{array} \right)$$

- similar approaches: Feverati et al '06, Hogervorst et al '14, Elias-Miro et al '17, ...

Effective field theory

- don't need to know everything to say something interesting

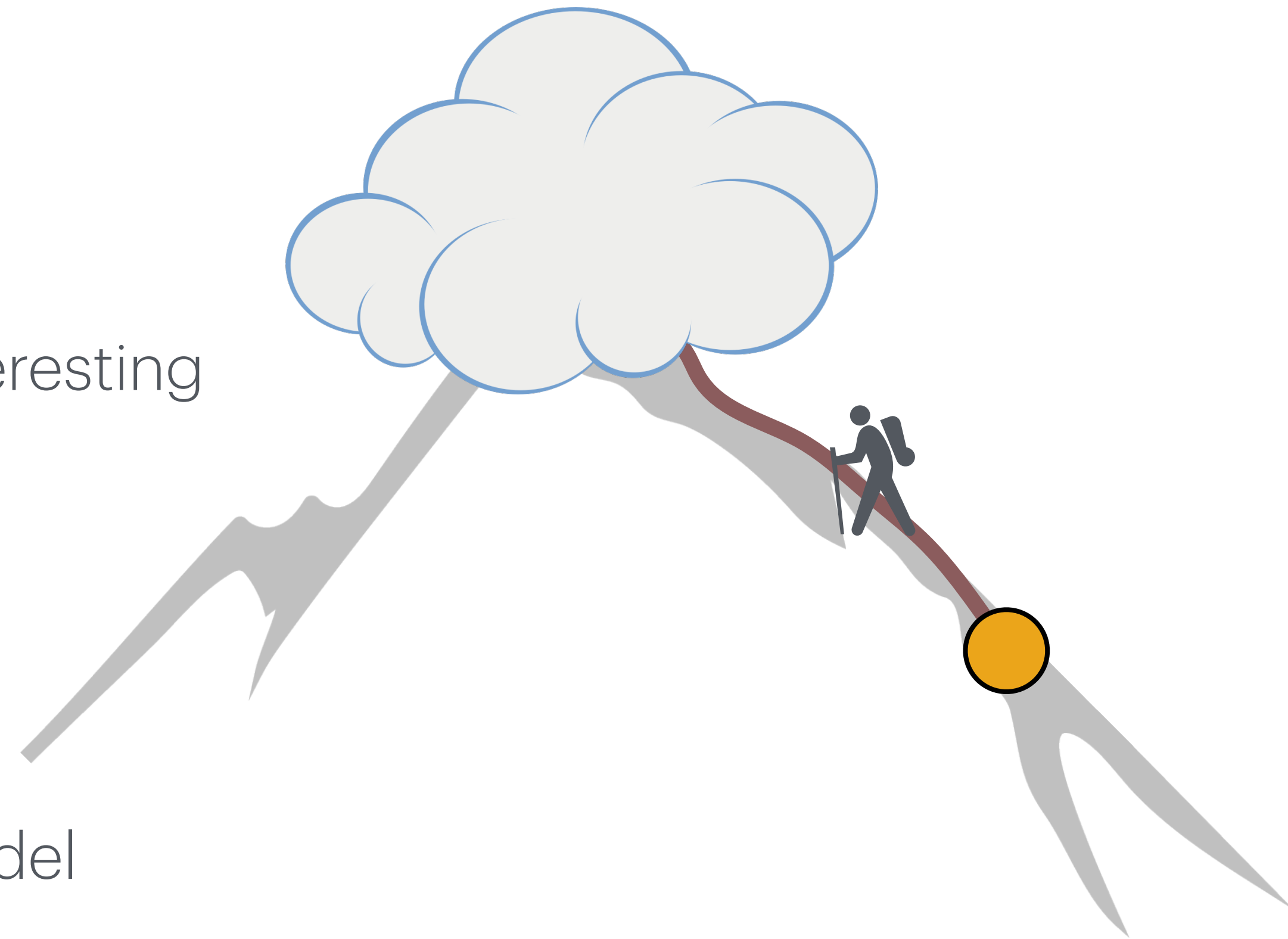


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Construct EFT:

- capture **relevant** physics with simplest possible model



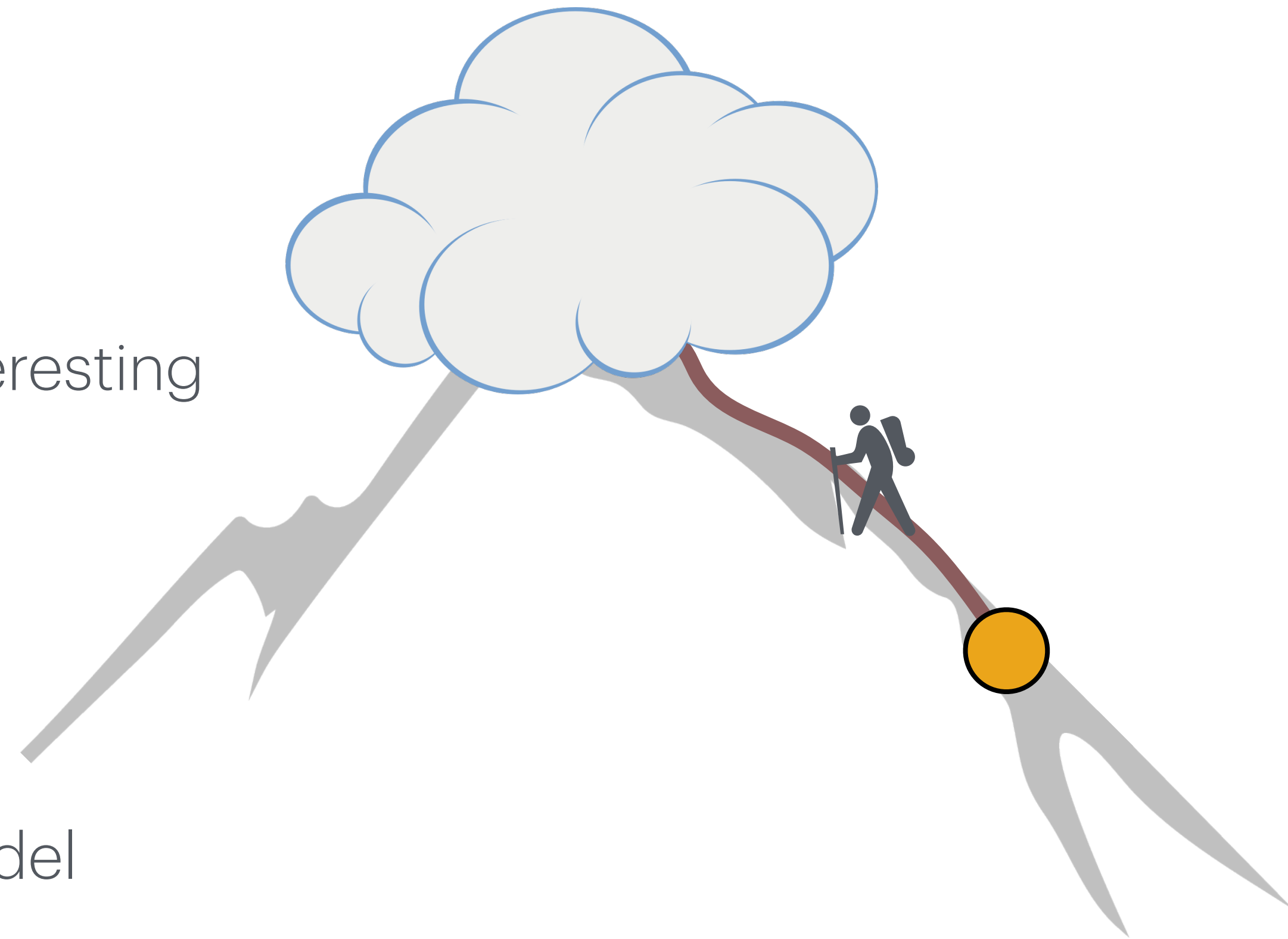
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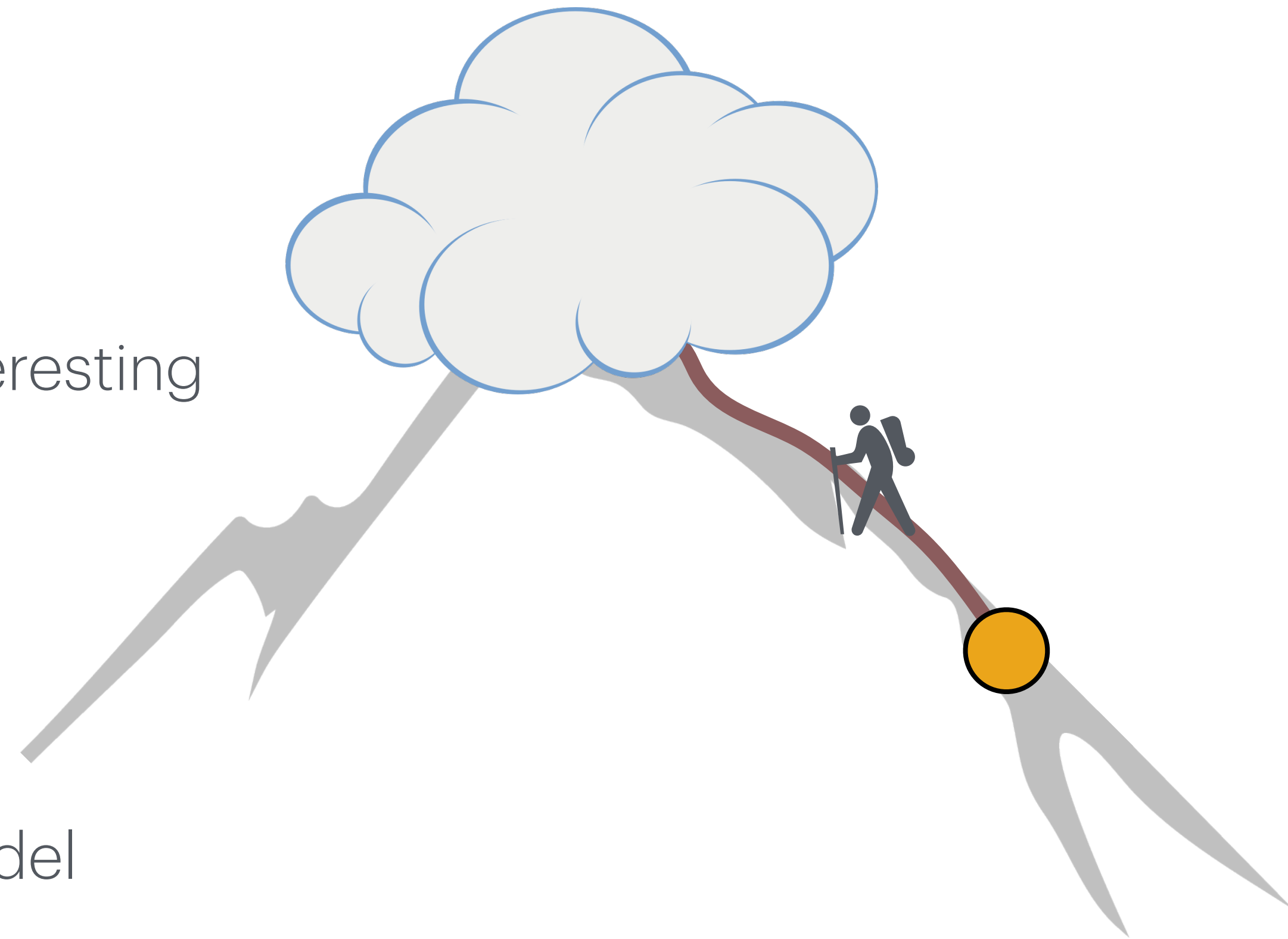
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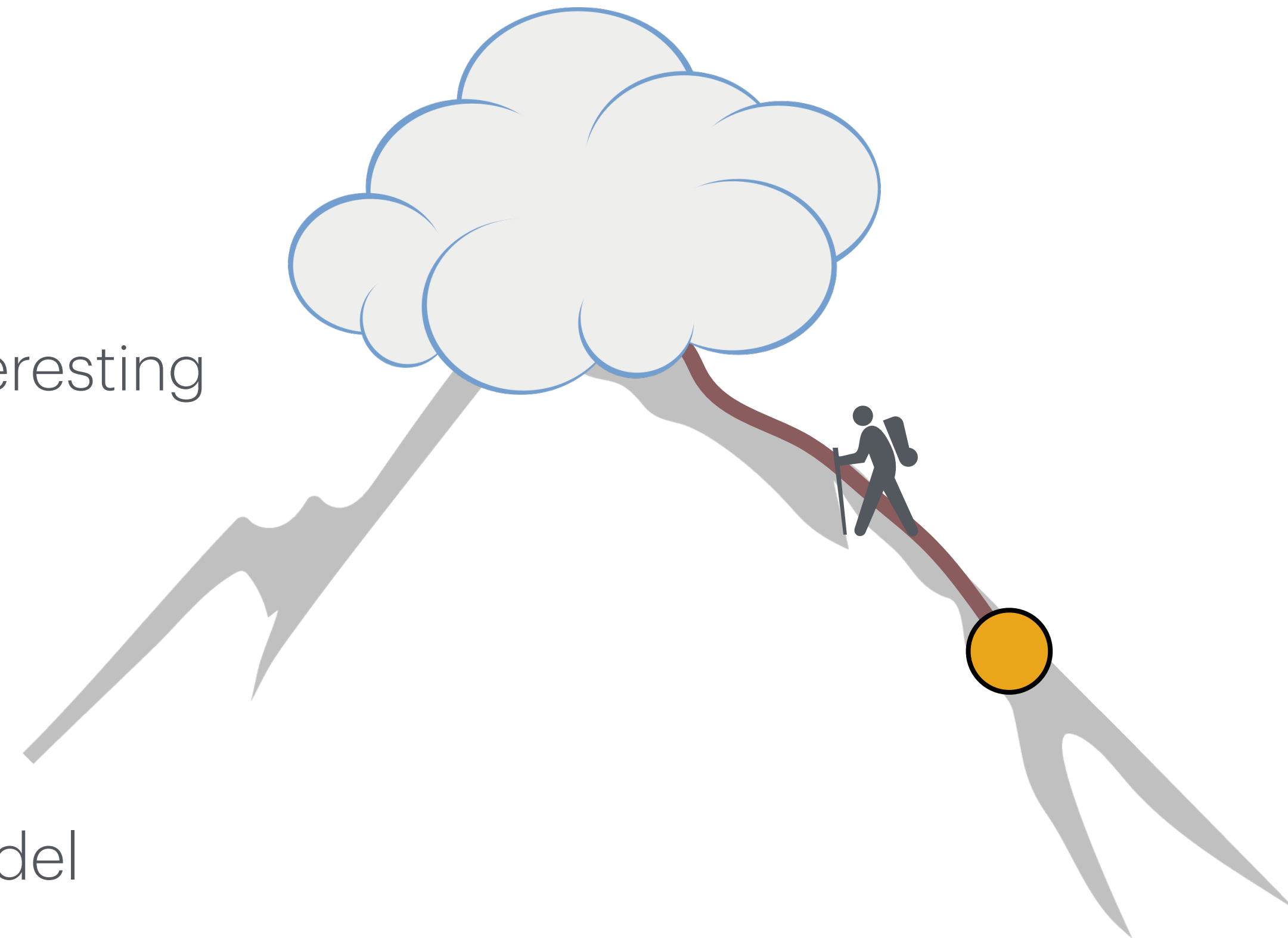
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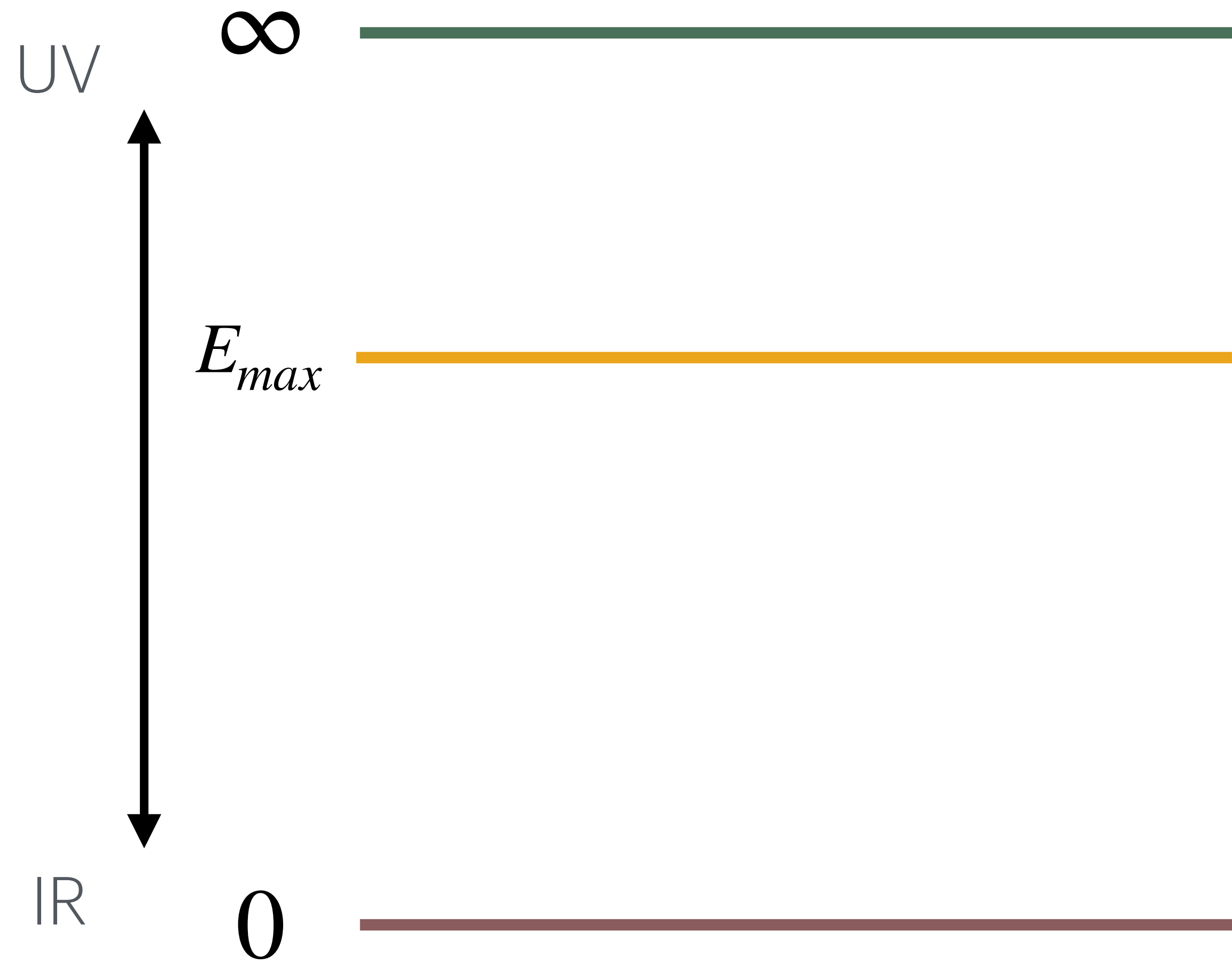
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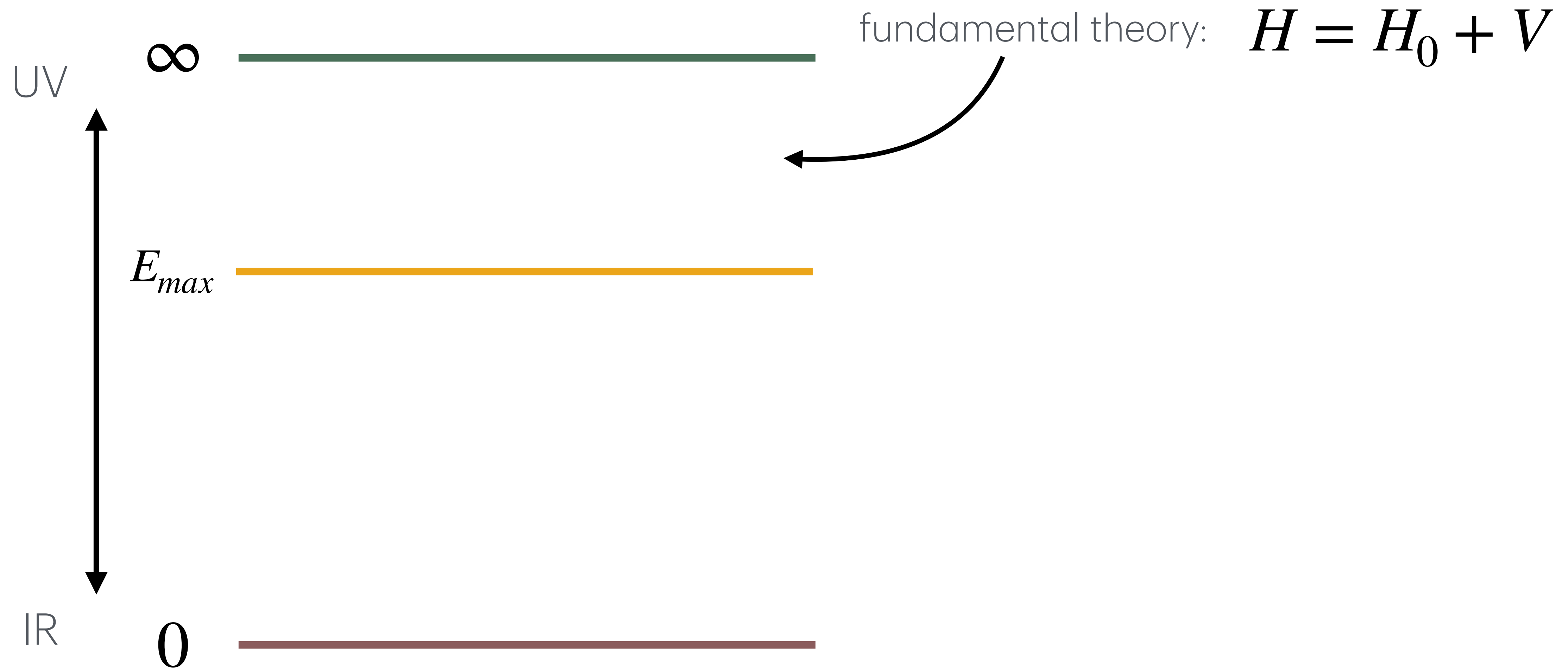
hallmarks

- separation of scales
- consistent power counting
- written in terms of local operators

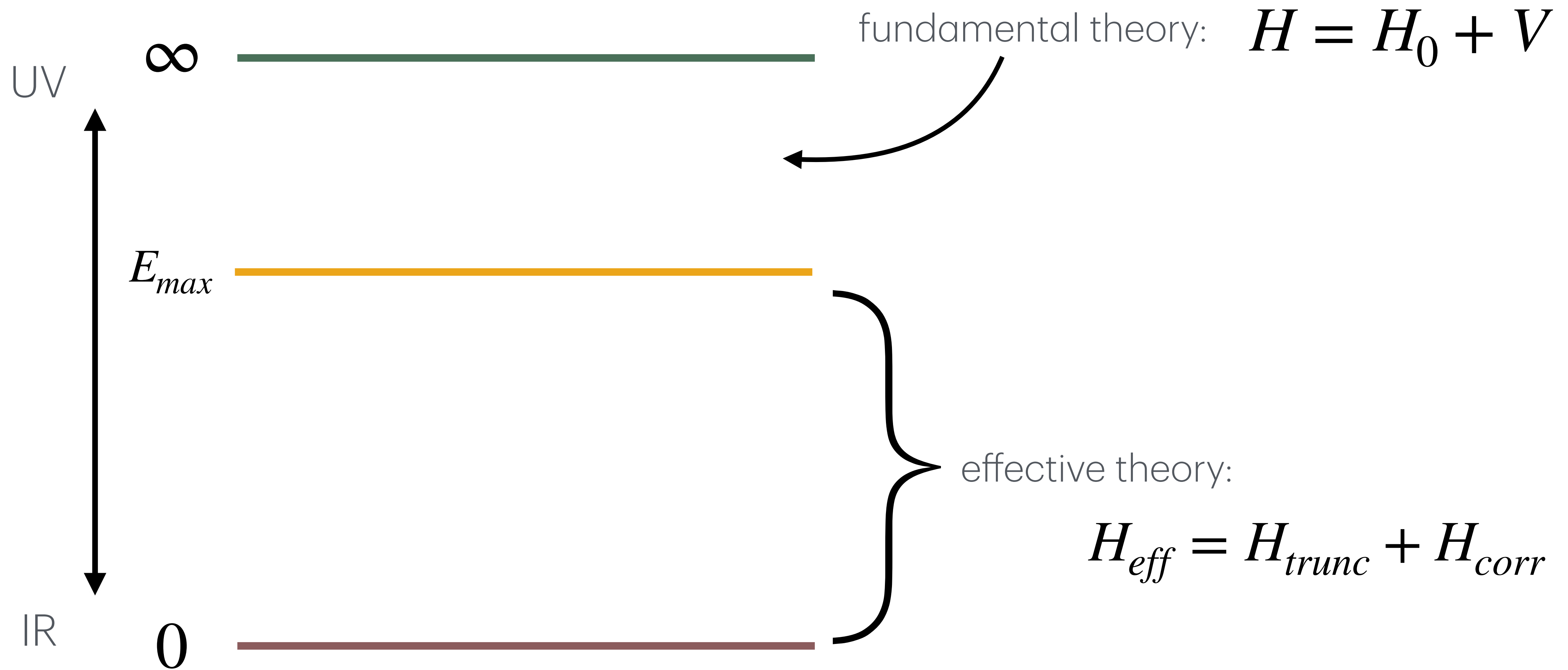
Effective field theory for truncation



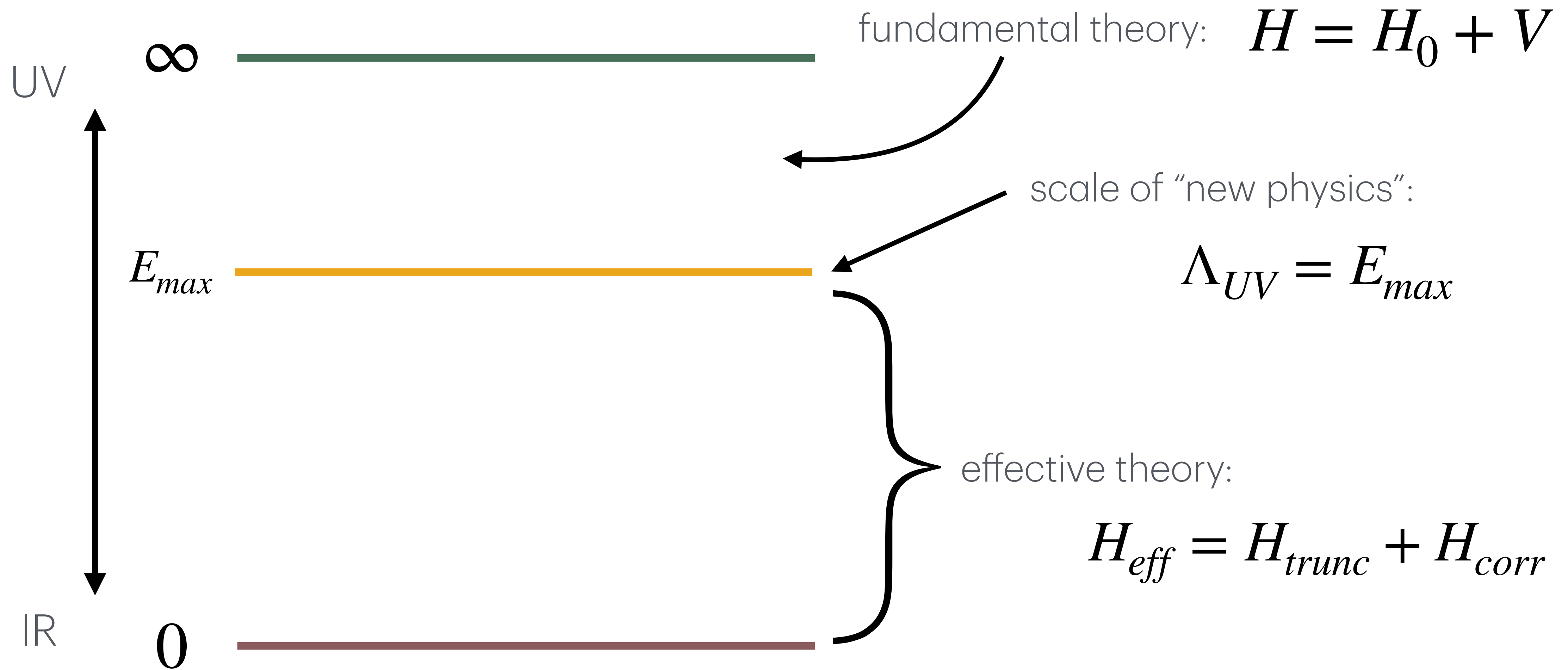
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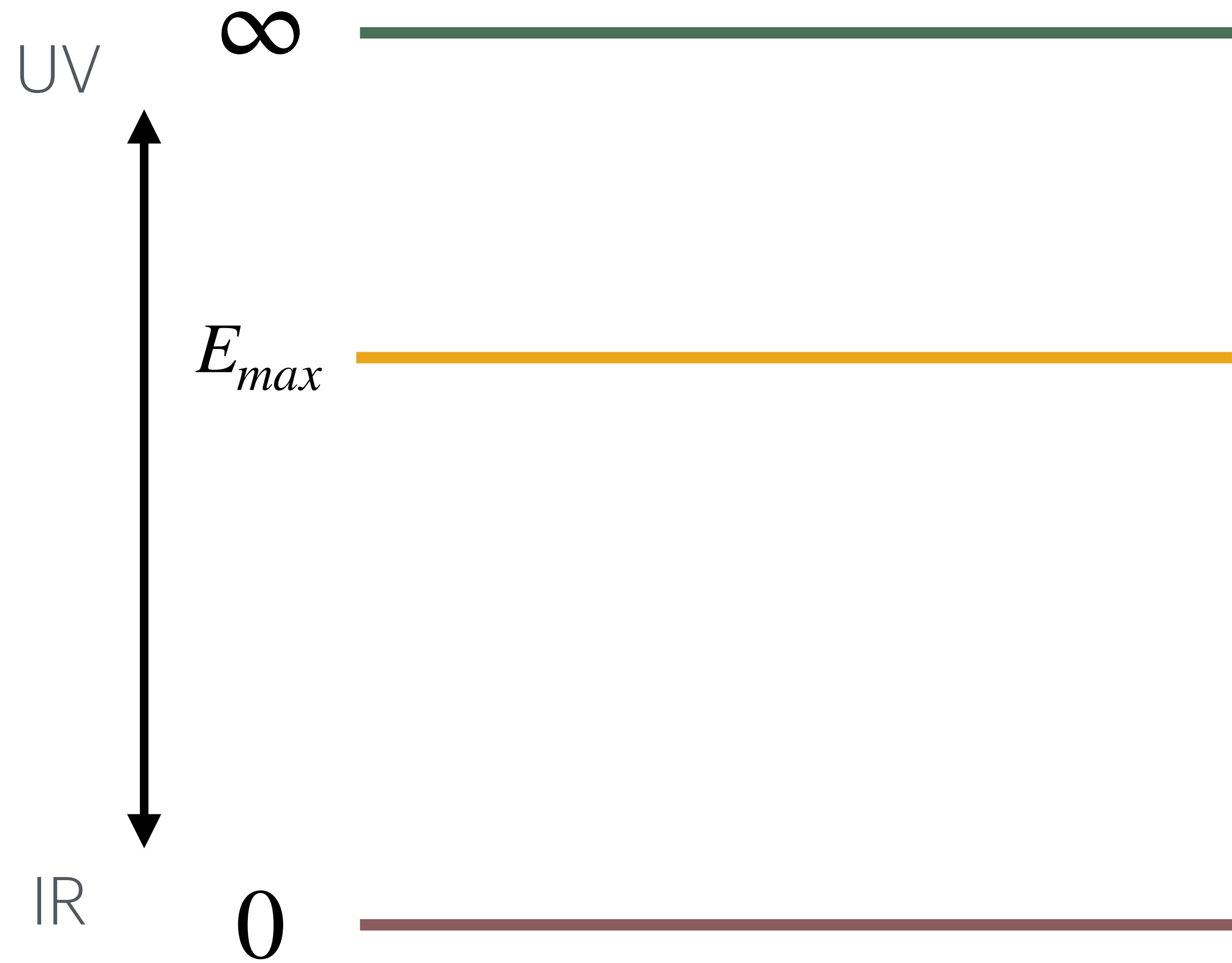
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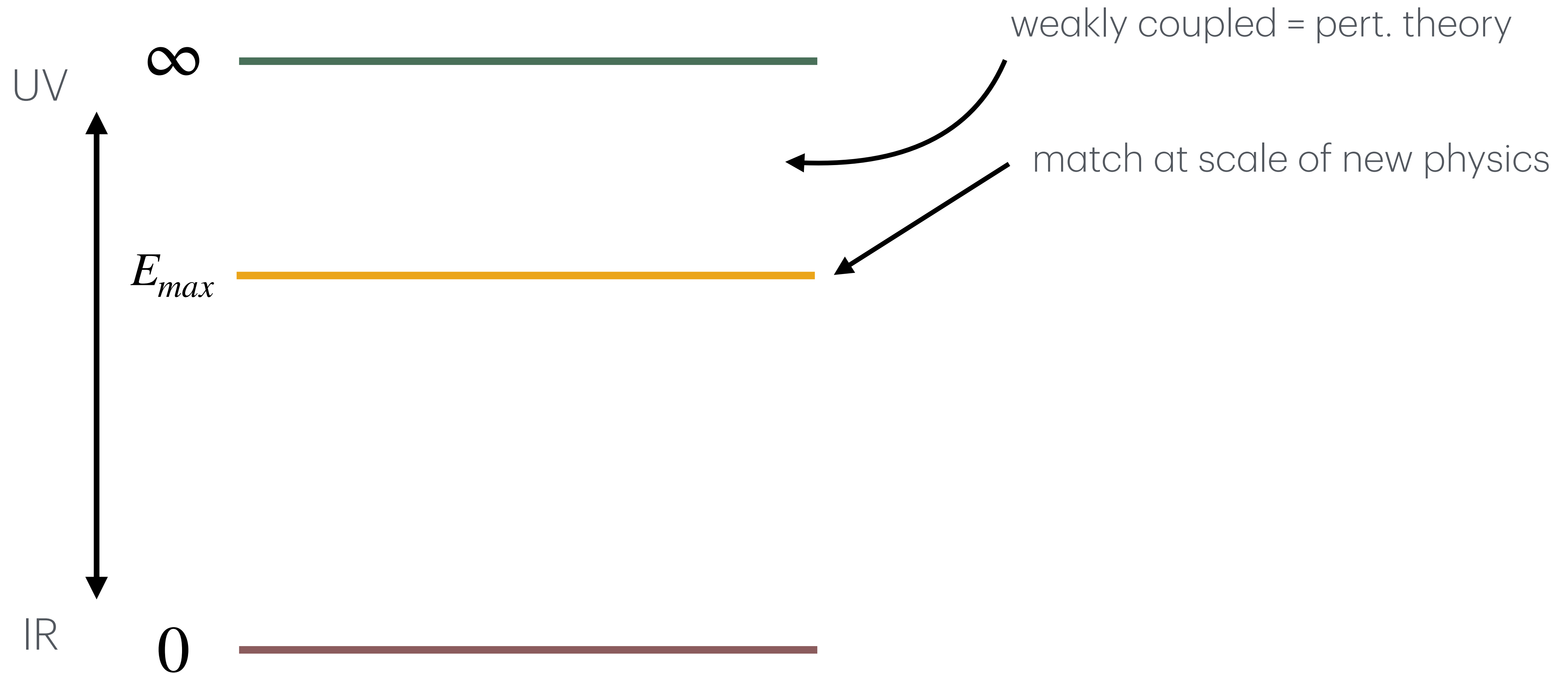
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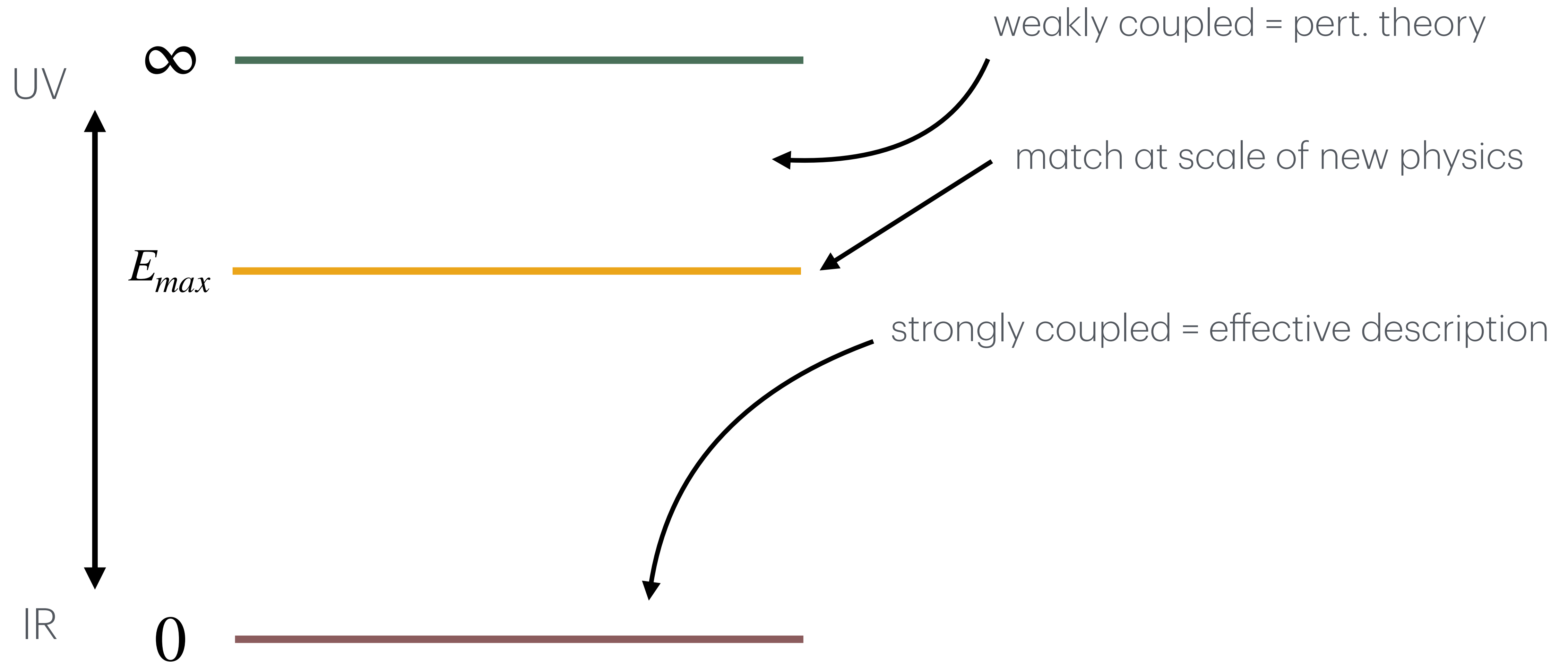
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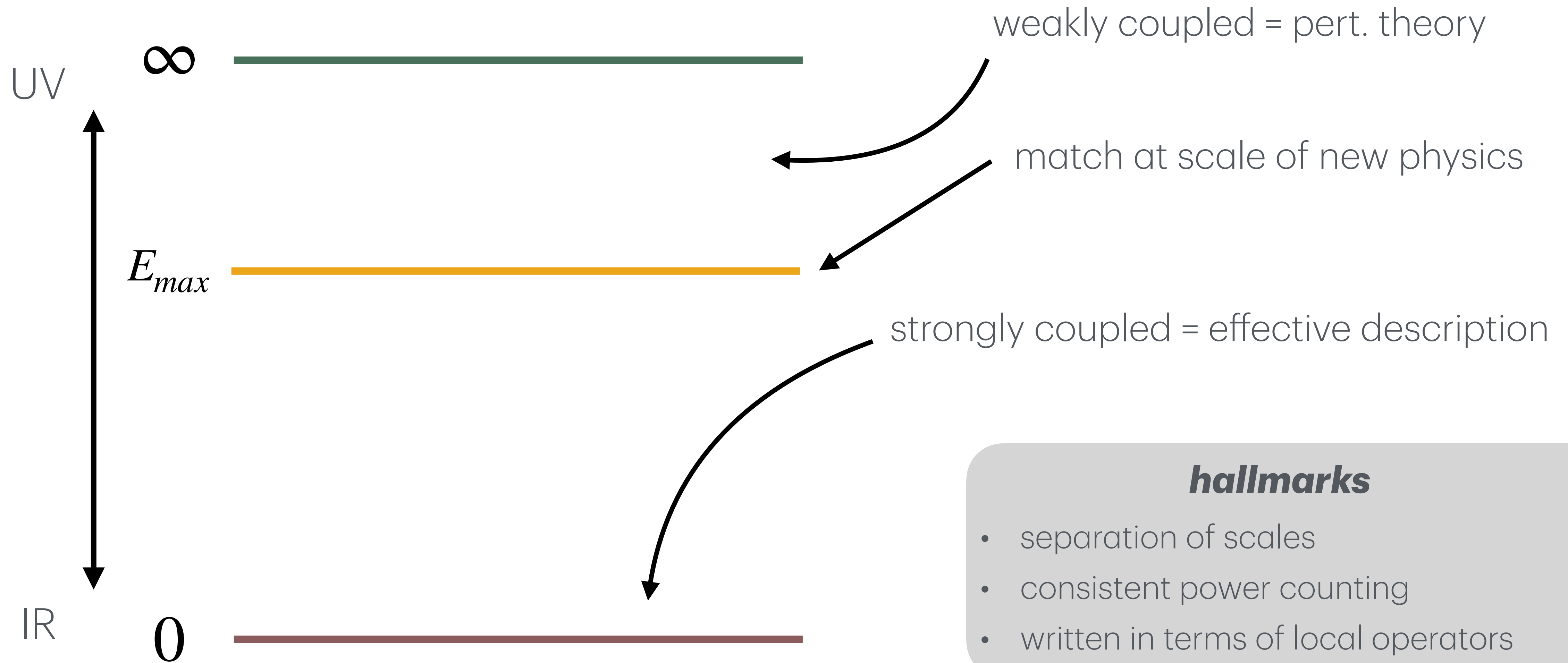
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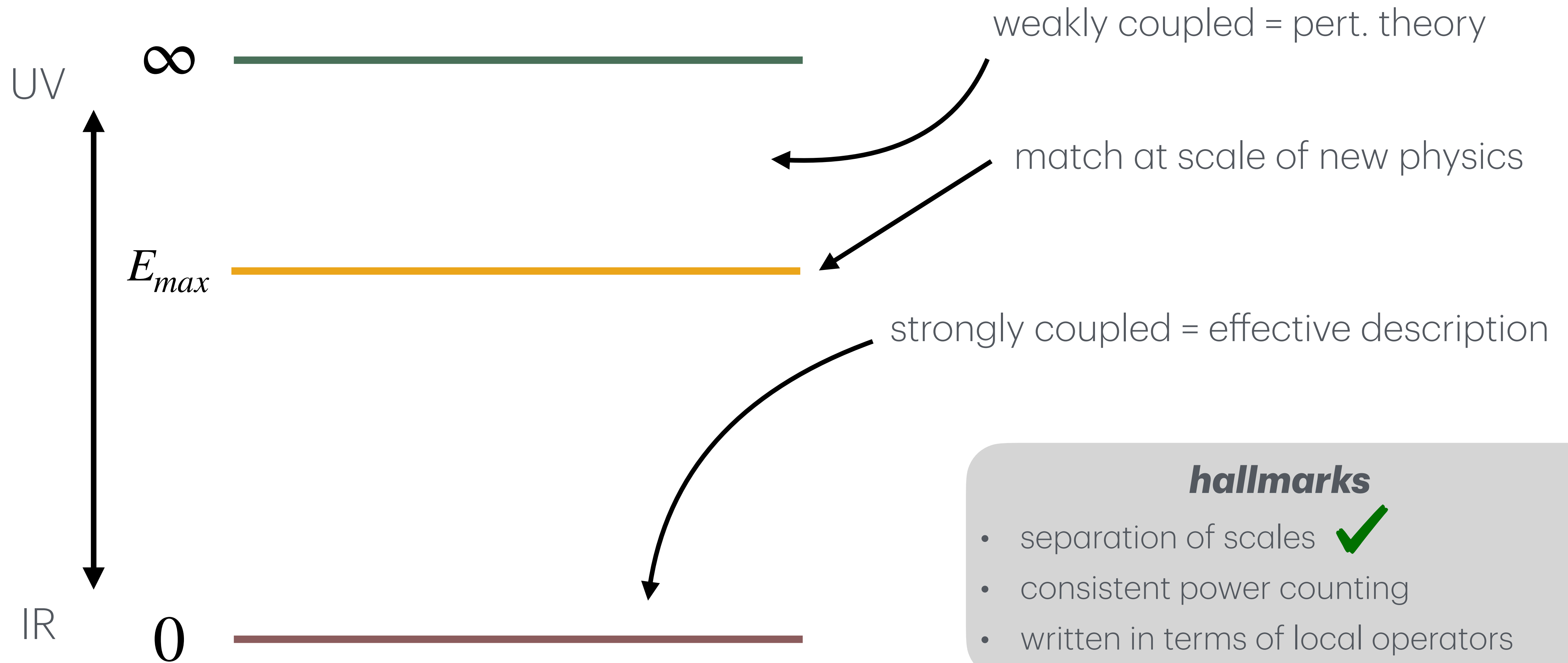
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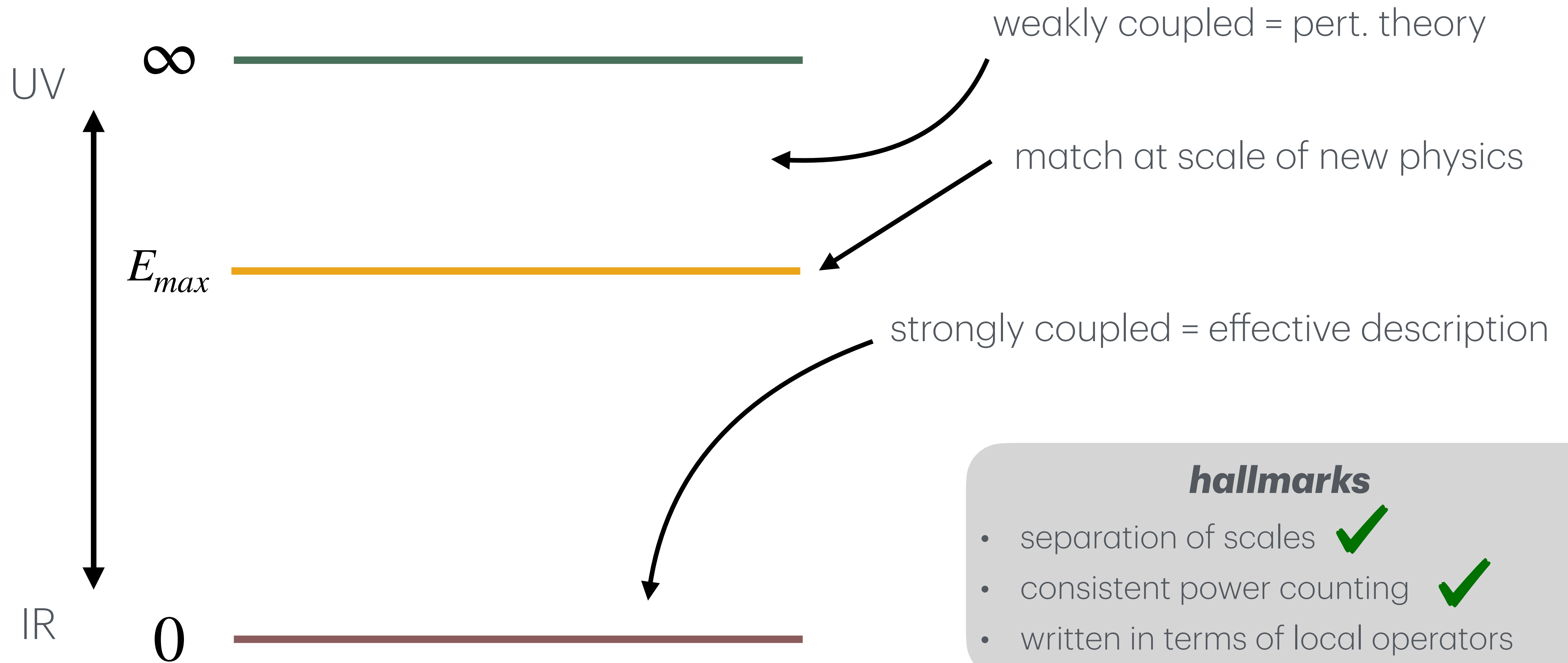
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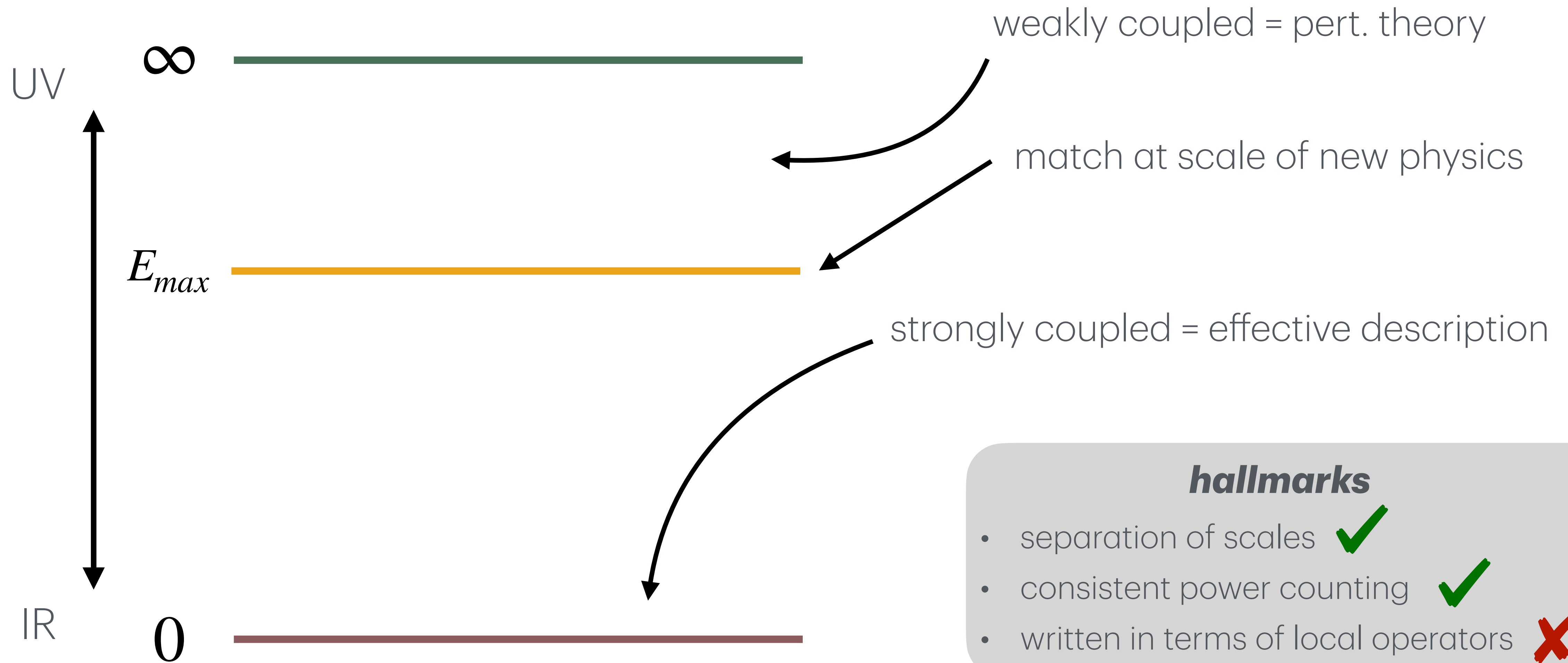
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What to match?

full theory: $H|\Psi_i\rangle = \mathcal{E}_i|\Psi_i\rangle$

$$\mathcal{E}_i = E_i + \mathcal{E}_{1i} + \mathcal{E}_{2i} + \dots$$

$H_0|i\rangle = E_i|i\rangle$

$\mathcal{E}_{ni} = \mathcal{O}(V^n)$

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The diagram consists of two curved arrows. The first arrow starts from the E_i term in the energy expansion and points to the $H_0|i\rangle = E_i|i\rangle$ equation. The second arrow starts from the \mathcal{E}_{ni} term and points to the $\mathcal{E}_{ni} = \mathcal{O}(V^n)$ equation.

effective theory: $H_{eff} = H_0 + H_1 + H_2 + \dots$

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effective theory: $H_{eff} = \underbrace{H_0 + H_1}_{H_{trunc}} + \underbrace{H_2 + \dots}_{H_{corr}}$ $H_n = \mathcal{O}(V^n)$

spectrum? $\mathcal{E}_{1i} = \langle i|H_1|i\rangle = \langle i|V|i\rangle$ only fixes diagonal elements

eigenvectors? $\langle \Psi_f|i\rangle \rightarrow \langle f|H_1|i\rangle = \langle f|V|i\rangle$ uniquely fixes H_{eff}

Example Theory 1: 2D $\lambda\phi^4$

$$\mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m^2\phi^2 - \frac{\lambda}{4!}\phi^4$$

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$$\omega_k = \sqrt{m^2 + k^2/R^2}$$

- strongly relevant interaction
- \mathbb{Z}_2 ($\phi \rightarrow -\phi$) symmetry broken at strong coupling = phase transition
 - in universality class of Ising Model: good check

Power Counting for 2D $\lambda\phi^4$

$$[\lambda] = 2, \quad [\phi] = 0$$

$$H_n \simeq \frac{\lambda^n}{E_{max}^{2n-2}} \int dx \text{ (dimensionless)}$$

$$H_2 \simeq \frac{\lambda^2}{E_{max}^2} \int dx \left(\phi^2 + \phi^4 + \mathcal{O}\left(\frac{E_f}{E_{max}}, \frac{R^{-1}}{E_{max}}\right) \right)$$

$$H_3 \simeq \frac{\lambda^3}{E_{max}^4} \int dx \text{ (dimensionless)}$$

Power Counting for 2D $\lambda\phi^4$

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What we expect:

$$H_n \simeq \frac{\lambda^n}{E_{max}^{2n-2}} \int dx \text{ (dimensionless)}$$

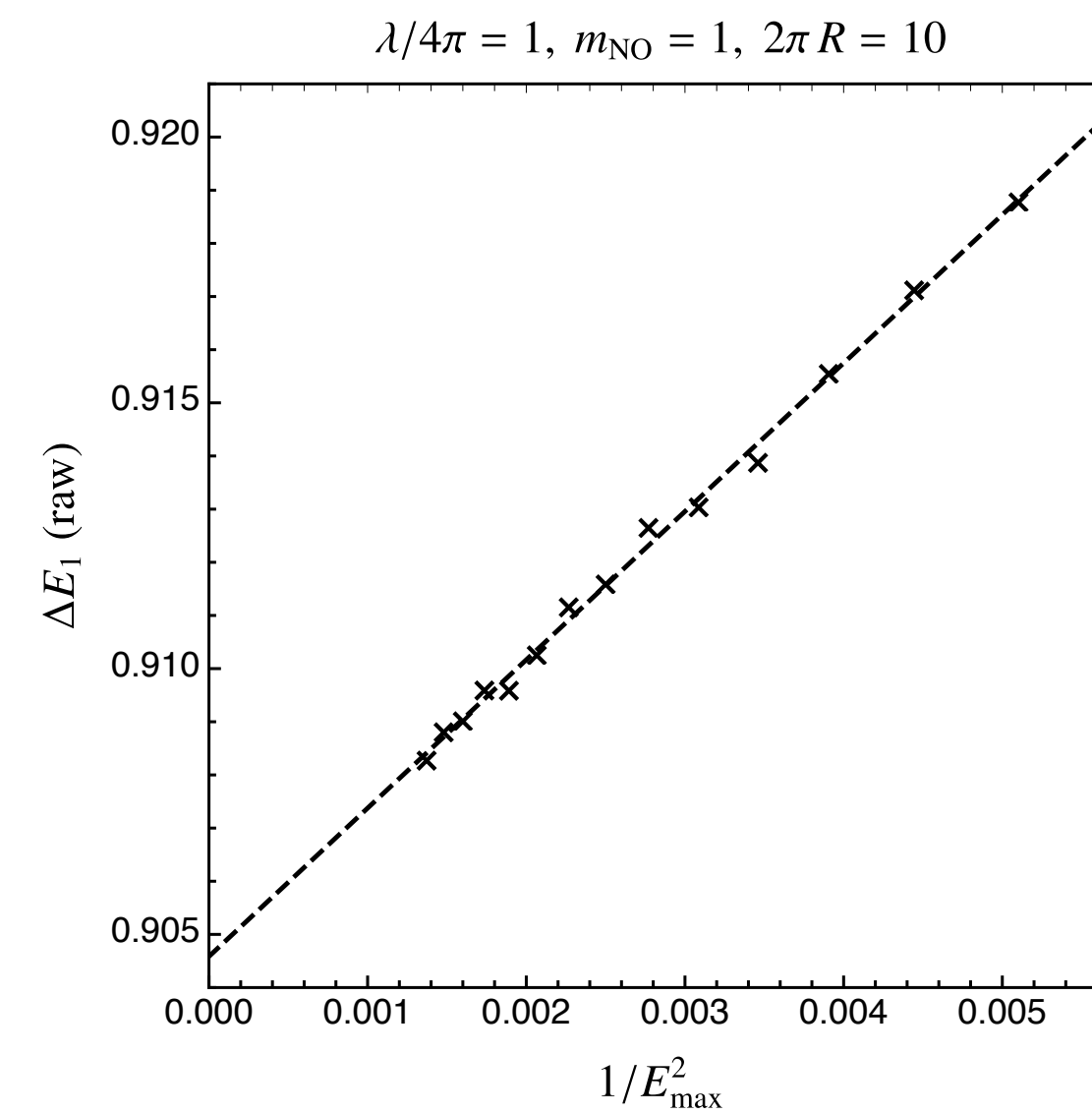
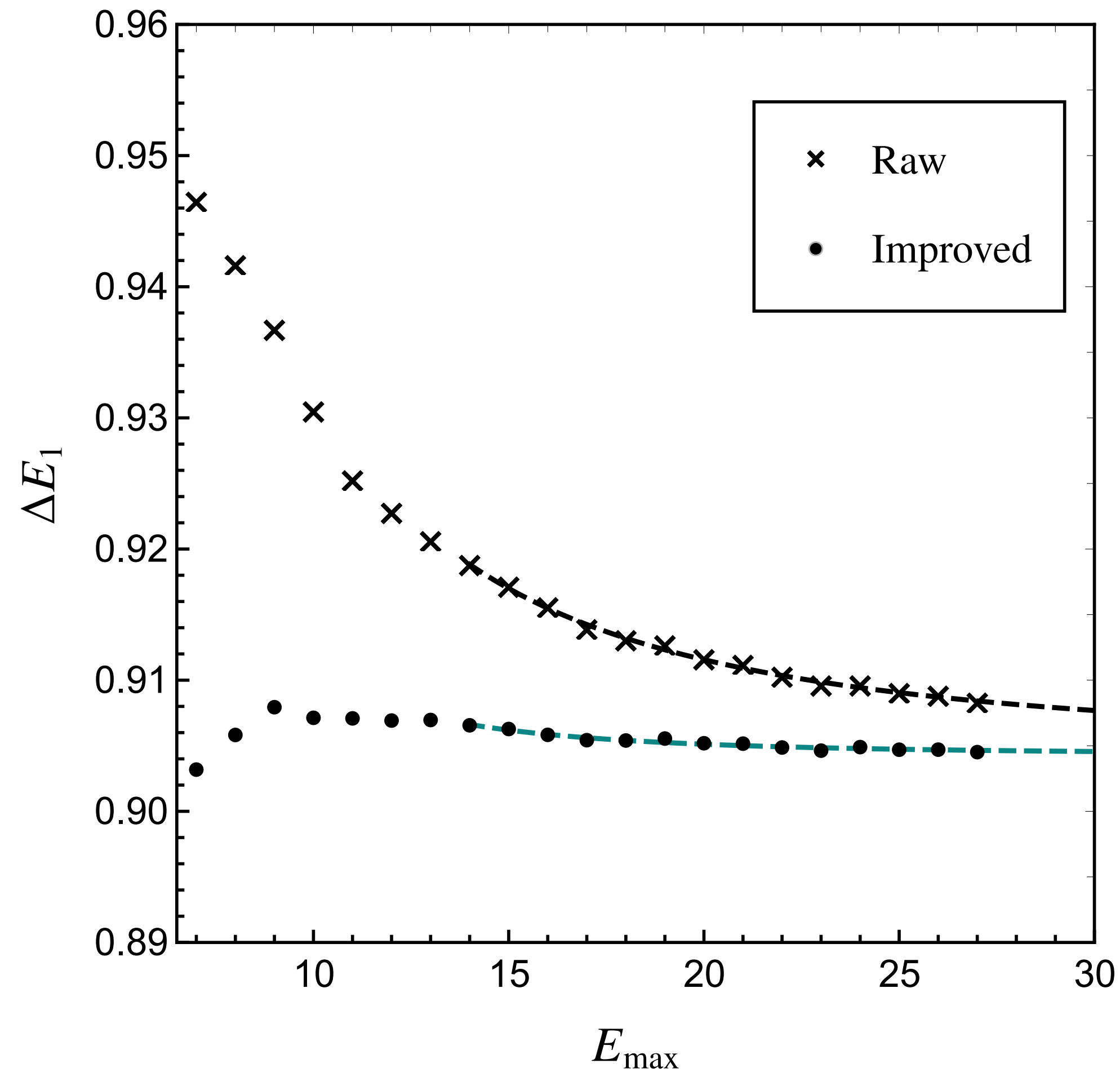
$$H_2 \simeq \frac{\lambda^2}{E_{max}^2} \int dx \left(\phi^2 + \phi^4 + \mathcal{O}\left(\frac{E_f}{E_{max}}, \frac{R^{-1}}{E_{max}}\right) \right)$$

$$H_3 \simeq \frac{\lambda^3}{E_{max}^4} \int dx \text{ (dimensionless)}$$

- error $\sim 1/E_{max}^2$ for raw truncation
- error $\sim 1/E_{max}^3$ after including corrections
- phase transition:
 - \mathbb{Z}_2 symmetry breaking at critical coupling
 - 2D Ising model

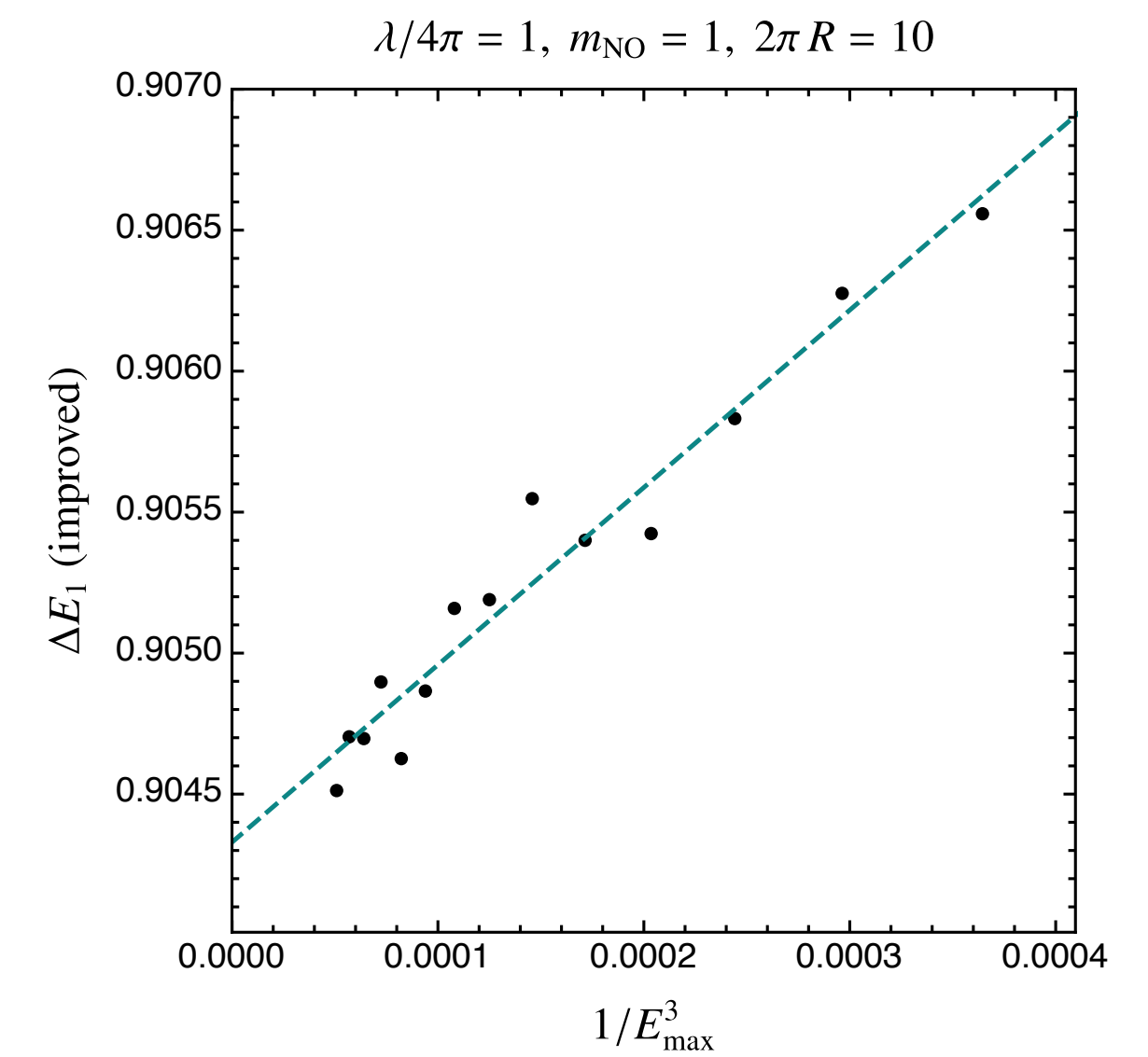
Scaling raw vs. improved

$$\lambda/4\pi = 1, m_{\text{NO}} = 1, 2\pi R = 10$$



without EFT

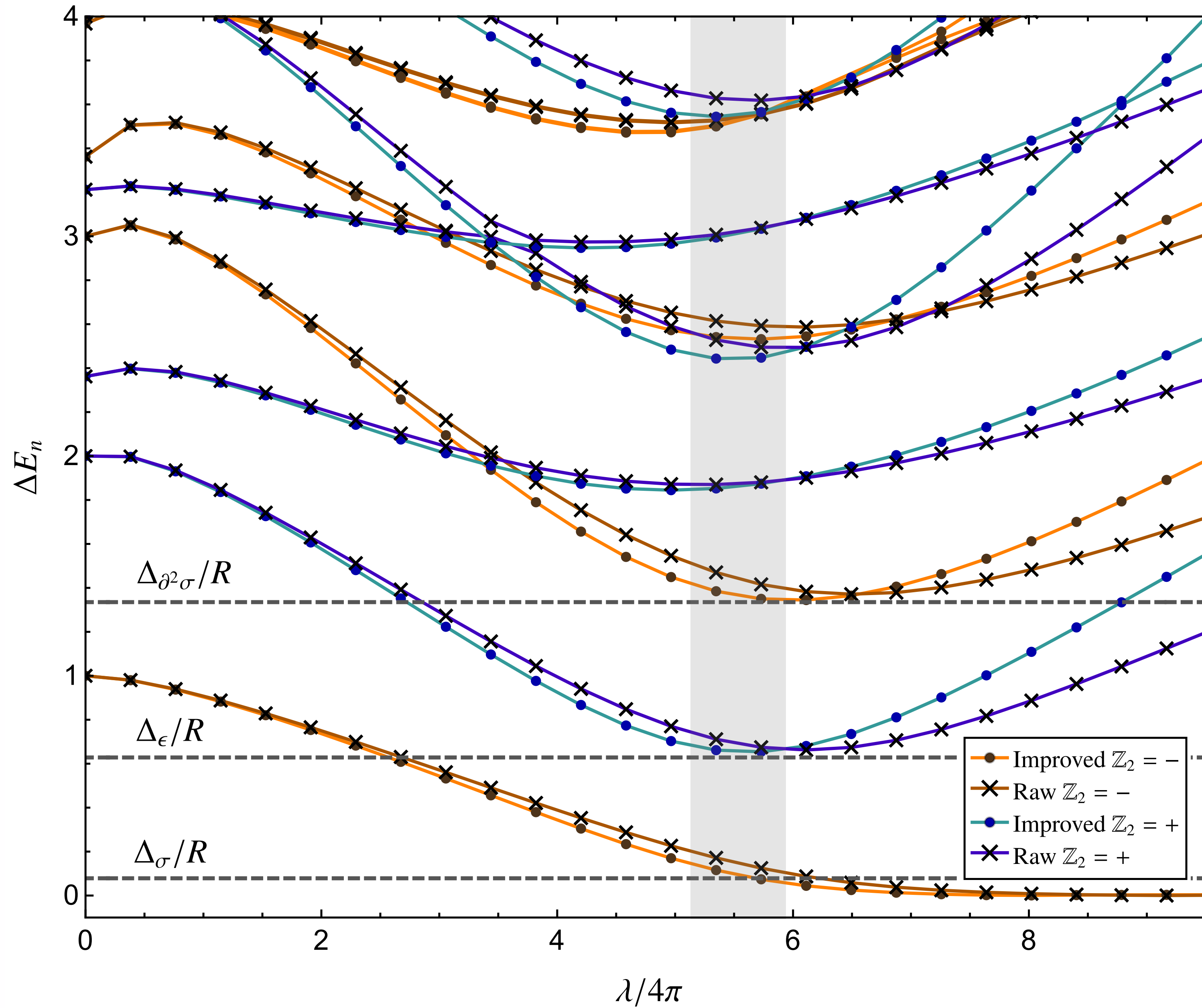
$$\sim 1/E_{\text{max}}^2$$



with EFT

$$\sim 1/E_{\text{max}}^3$$

$$E_{\max} = 27, m_{\text{NO}} = 1, 2\pi R = 10$$



What about 4D QCD?

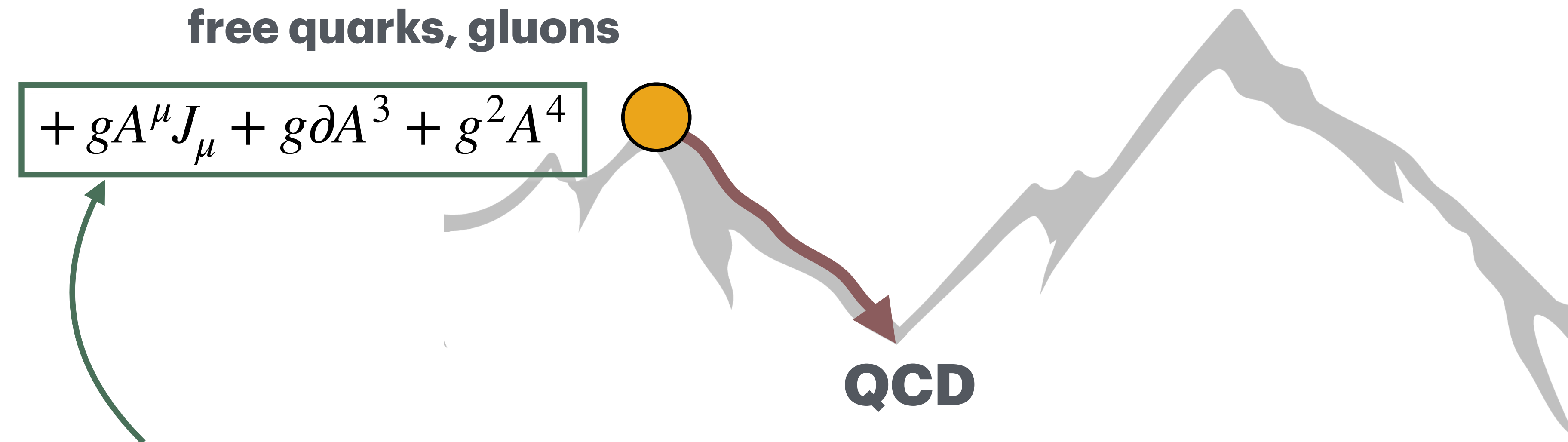


What about 4D QCD?

free quarks, gluons



What about 4D QCD?



- marginal (convergence unknown)
- not gauge-invariant

What about 4D QCD?

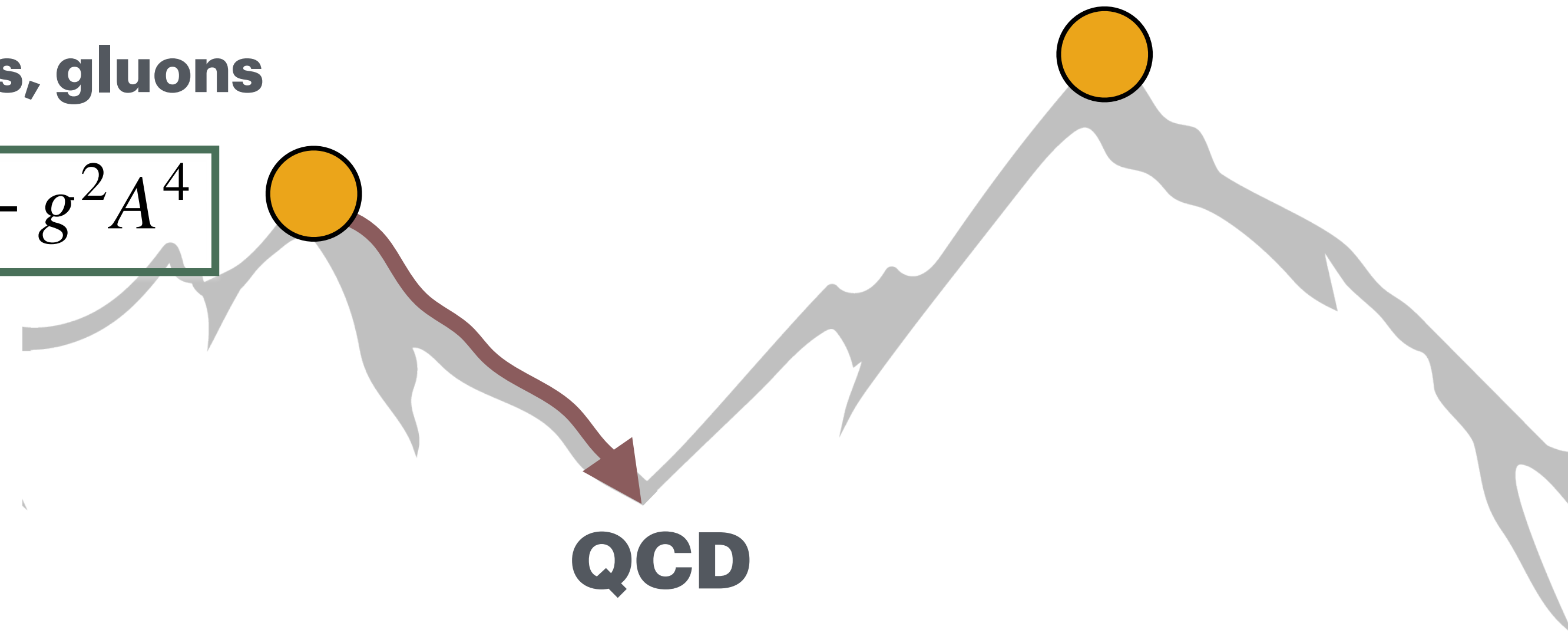
Banks-Zaks (e.g. $SU(3), N_f = 16$)

free quarks, gluons

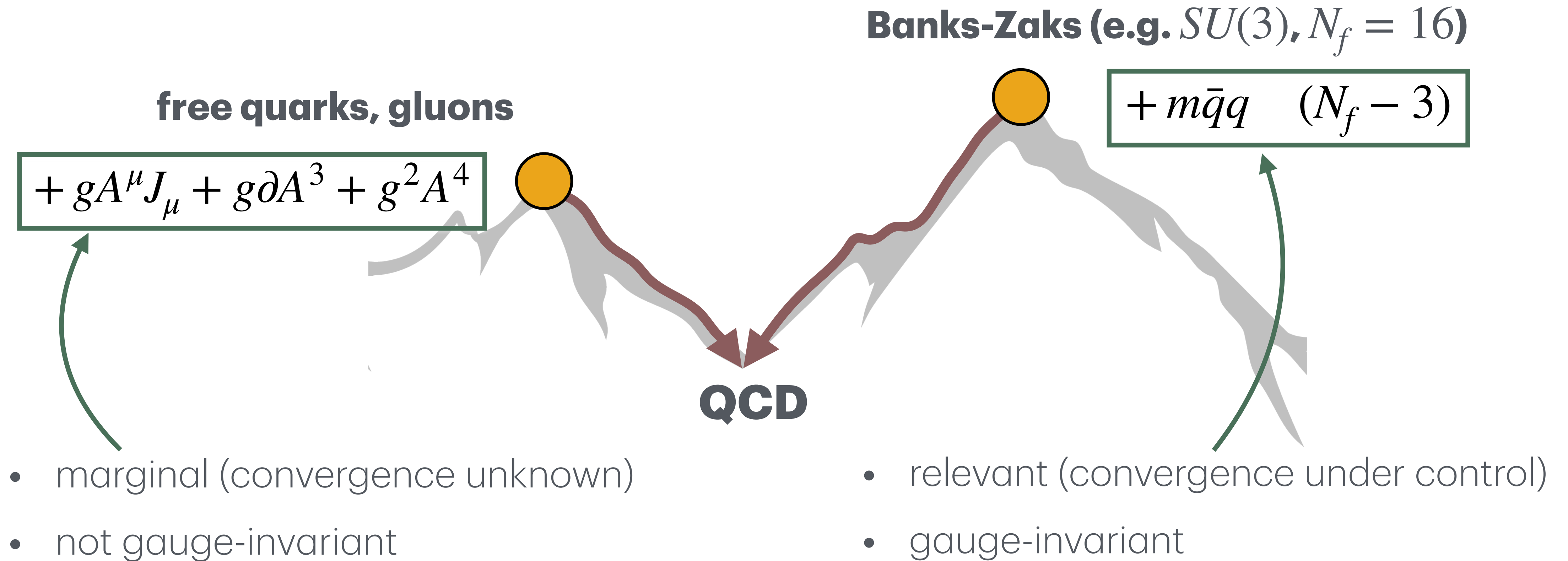
$$+ gA^\mu J_\mu + g\partial A^3 + g^2 A^4$$

QCD

- marginal (convergence unknown)
- not gauge-invariant



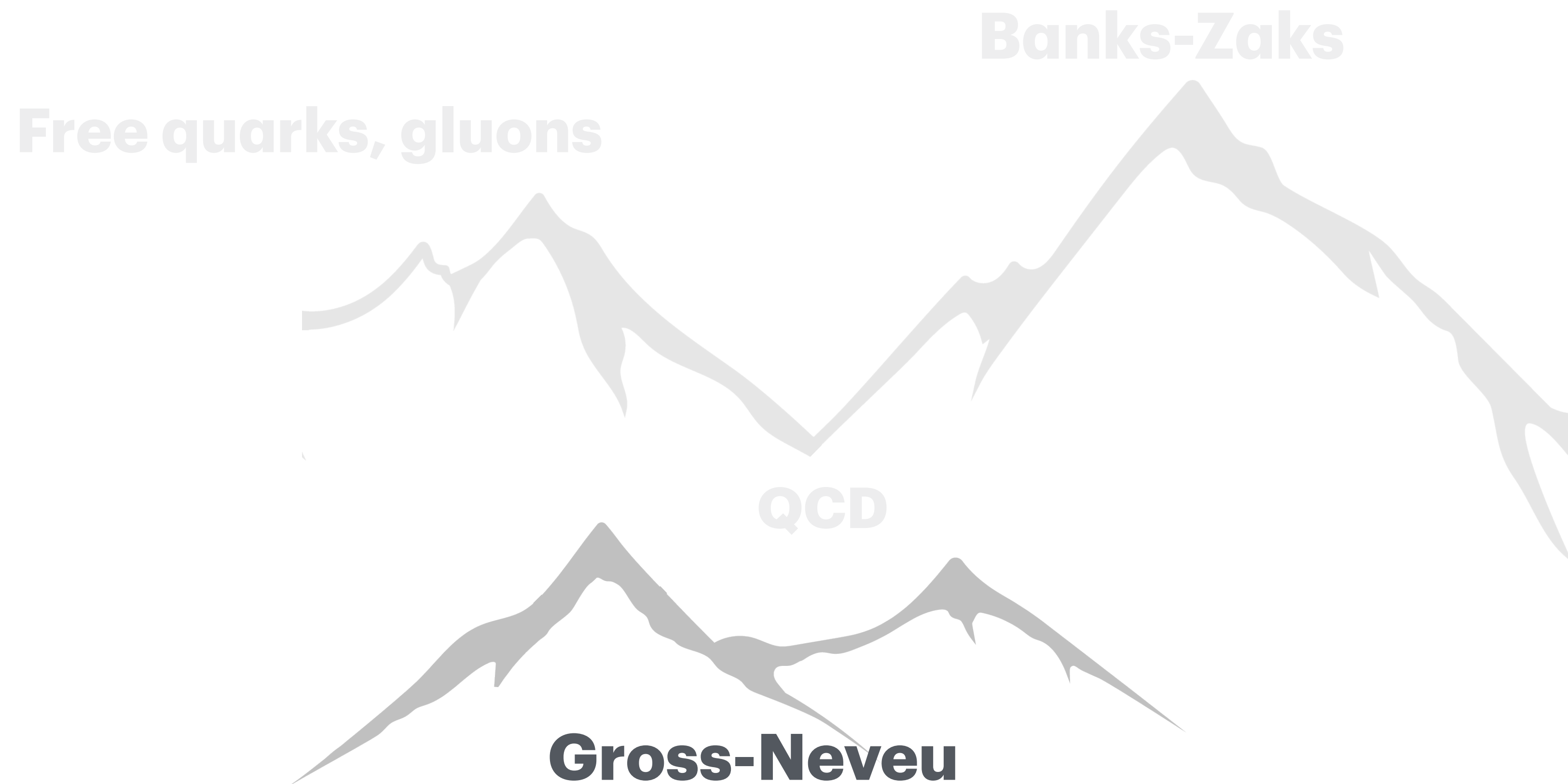
What about 4D QCD?



First step: 2D Gross-Neveu



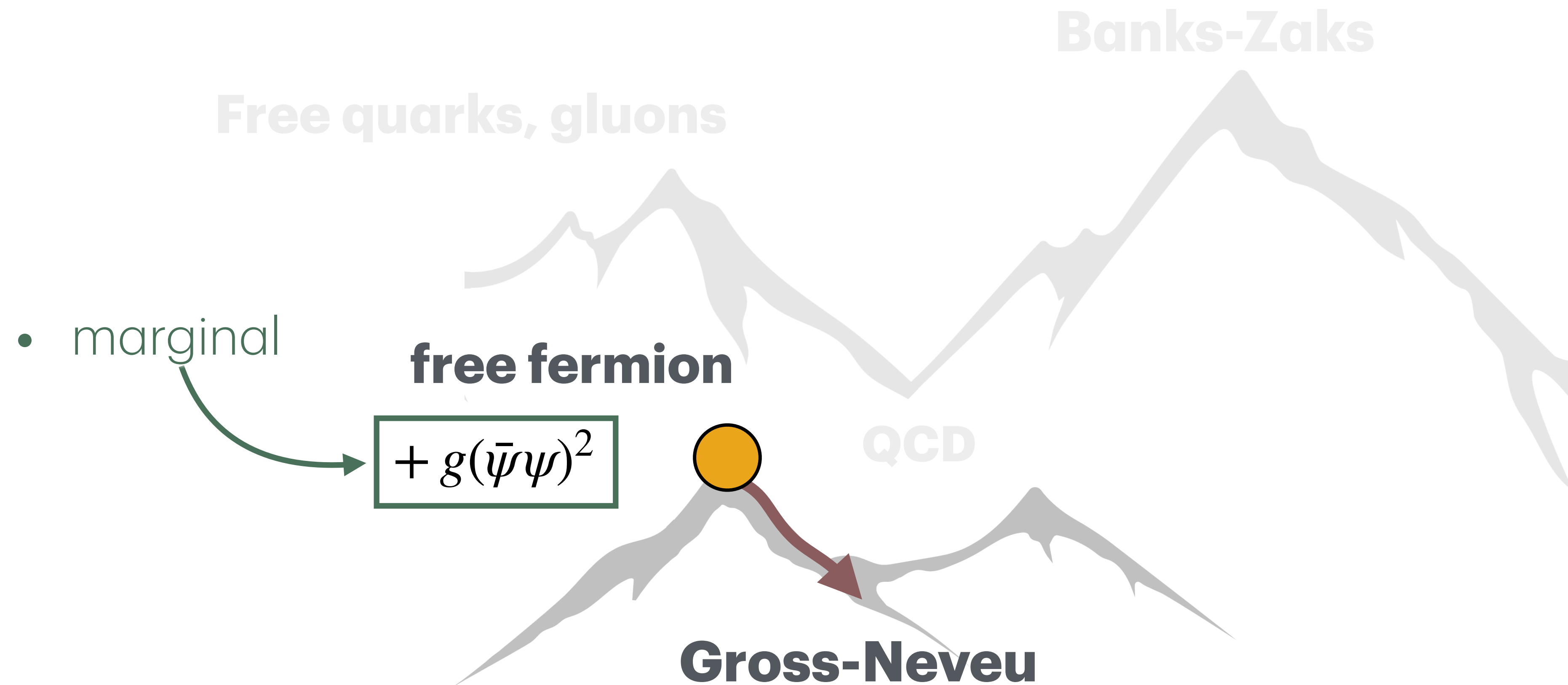
First step: 2D Gross-Neveu



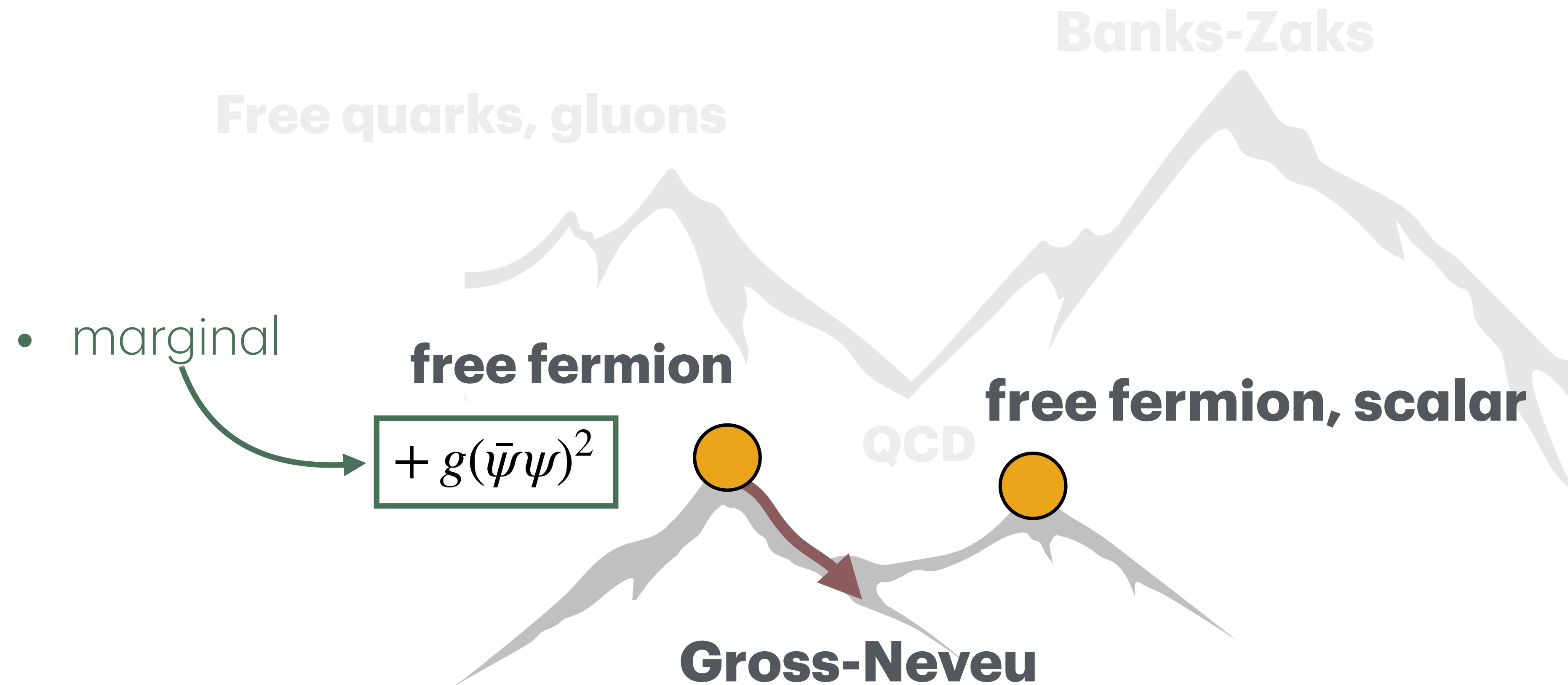
First step: 2D Gross-Neveu



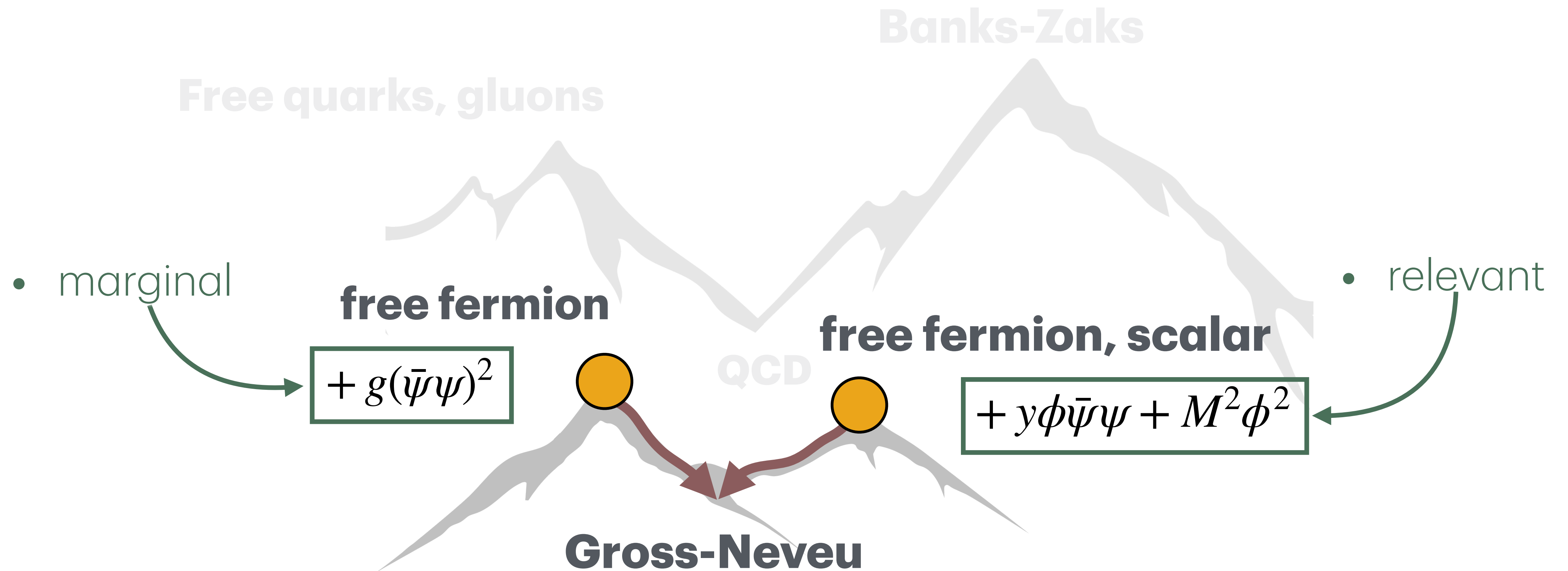
First step: 2D Gross-Neveu



First step: 2D Gross-Neveu



First step: 2D Gross-Neveu



Simpler testing ground on way to QCD

Example Theory 2: 2D Yukawa

$$\mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m_\phi^2\phi^2 + \bar{\psi}(i\gamma^\mu\partial_\mu - m_\psi)\psi + y\phi\bar{\psi}\psi$$

$$\phi(x, t) = \phi(x + 2\pi R, t)$$

$$\psi(x, t) = \psi(x + 2\pi R, t)$$

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H_0

V

$$\phi(x, t) = \phi(x + 2\pi R, t)$$

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H_0

V

$$\phi(x, t) = \phi(x + 2\pi R, t)$$

$$\psi(x, t) = \psi(x + 2\pi R, t)$$

- now with **real** UV divergences at $\mathcal{O}(y^2)$
 - regulate with local counterterms
 - expect finite after including H_2 (DeLouche et al '23)
- charge conjugation symmetry

Power Counting for 2D Yukawa

$$[y] = 1, \quad [\phi] = 0, \quad [\psi] = 1/2$$

$$H_2 \simeq y^2 \int dx \left(\ln(E_{max}) + \phi^2 \ln(E_{max}) + \frac{1}{E_{max}} \bar{\psi} \psi + \frac{H_0}{E_{max}} \phi^2 + \dots \right)$$

$$H_3 \simeq \frac{y^3}{E_{max}} \int dx \left(\frac{1}{E_{max}} \phi \bar{\psi} \psi + \dots \right)$$

$$H_4 \simeq \frac{y^4}{E_{max}^2} \int dx \left(1 + \phi^2 + \frac{1}{E_{max}} \bar{\psi} \psi + \dots \right)$$

Power Counting for 2D Yukawa

$$[y] = 1, \quad [\phi] = 0, \quad [\psi] = 1/2$$

What we expect:

$$H_2 \simeq y^2 \int dx \left(\ln(E_{max}) + \phi^2 \ln(E_{max}) + \frac{1}{E_{max}} \bar{\psi} \psi + \frac{H_0}{E_{max}} \phi^2 + \dots \right)$$

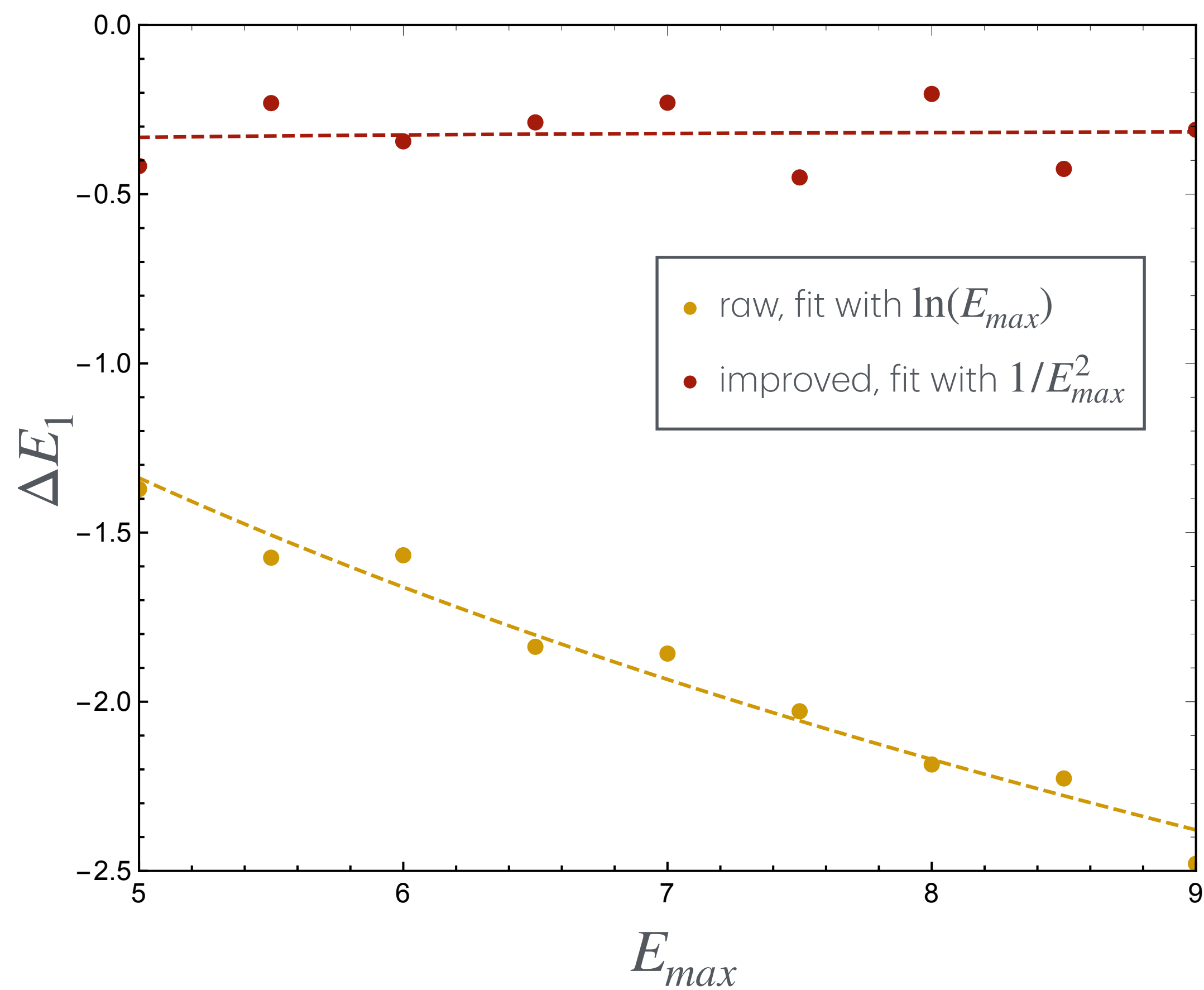
$$H_3 \simeq \frac{y^3}{E_{max}} \int dx \left(\frac{1}{E_{max}} \phi \bar{\psi} \psi + \dots \right)$$

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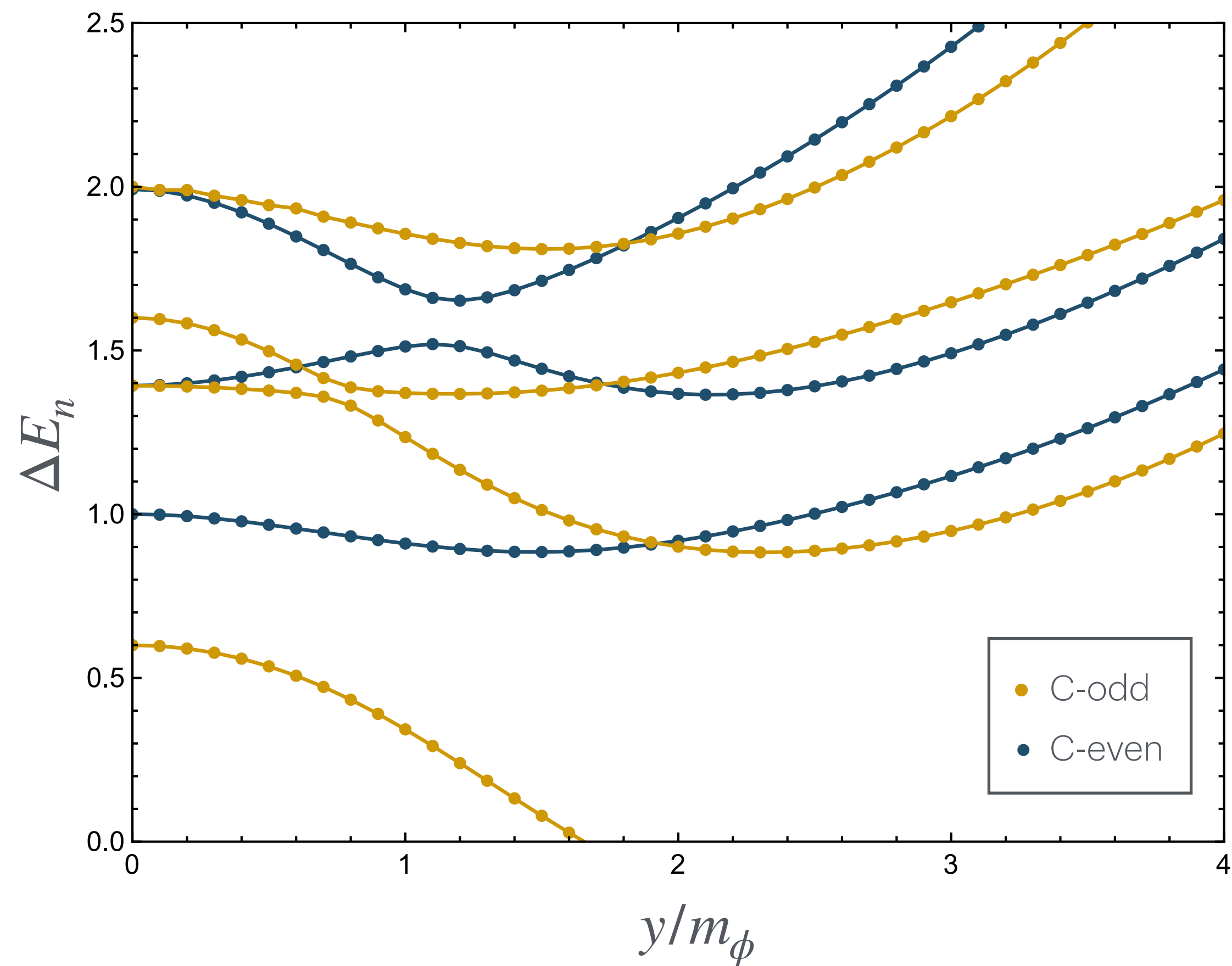
- diverge like $\ln(E_{max})$ with no correction
- error $\sim 1/E_{max}$ including local counterterms
- error $\sim 1/E_{max}^2$ after including corrections

Preliminary plots

$$y/m_\phi = \sqrt{4\pi}, m_\phi = 1, m_\psi = 0.3, 2\pi R = 10$$

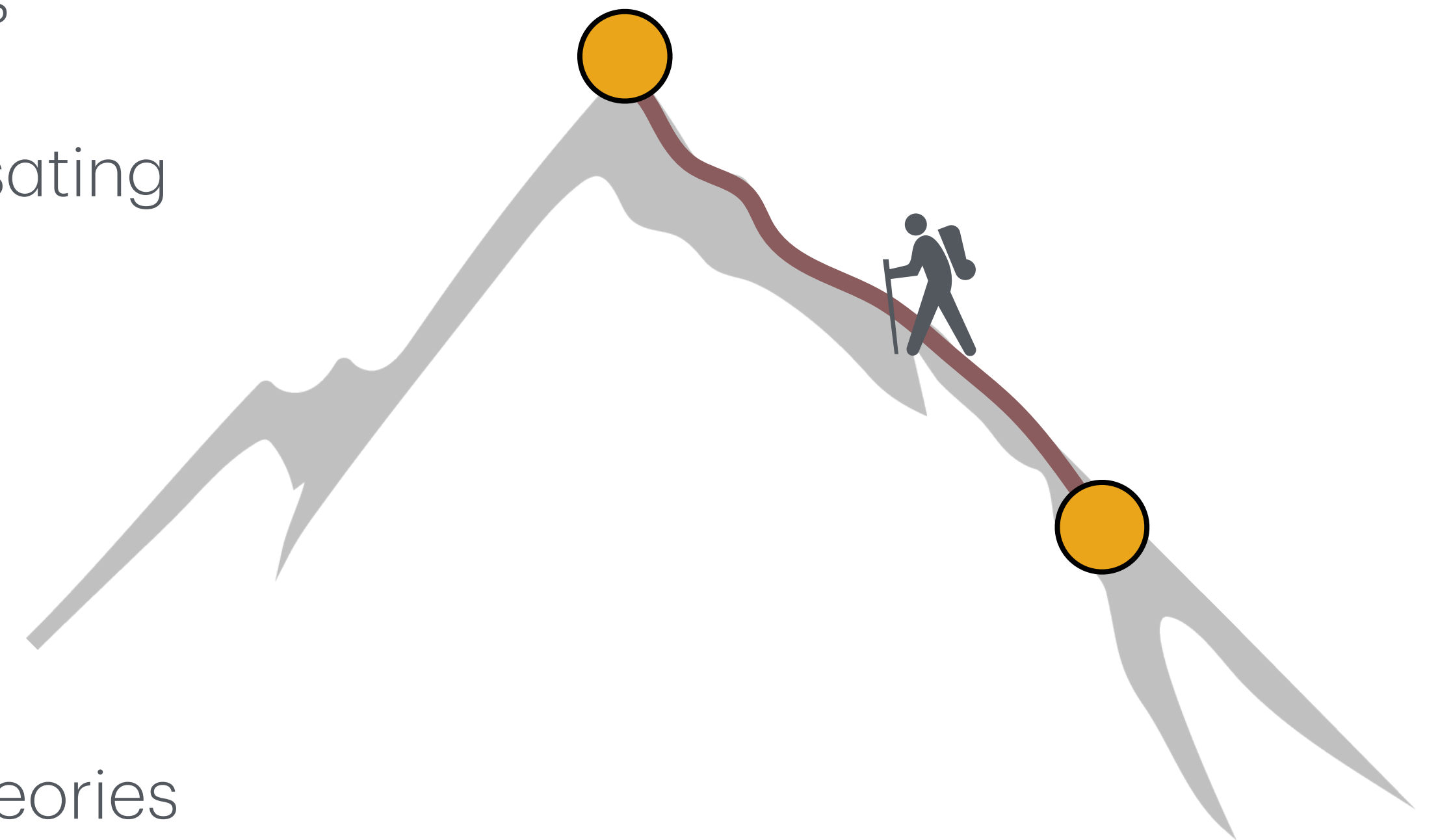


$$E_{max} = 8.5, m_\phi = 1, m_\psi = 0.3, 2\pi R = 10$$



Summary

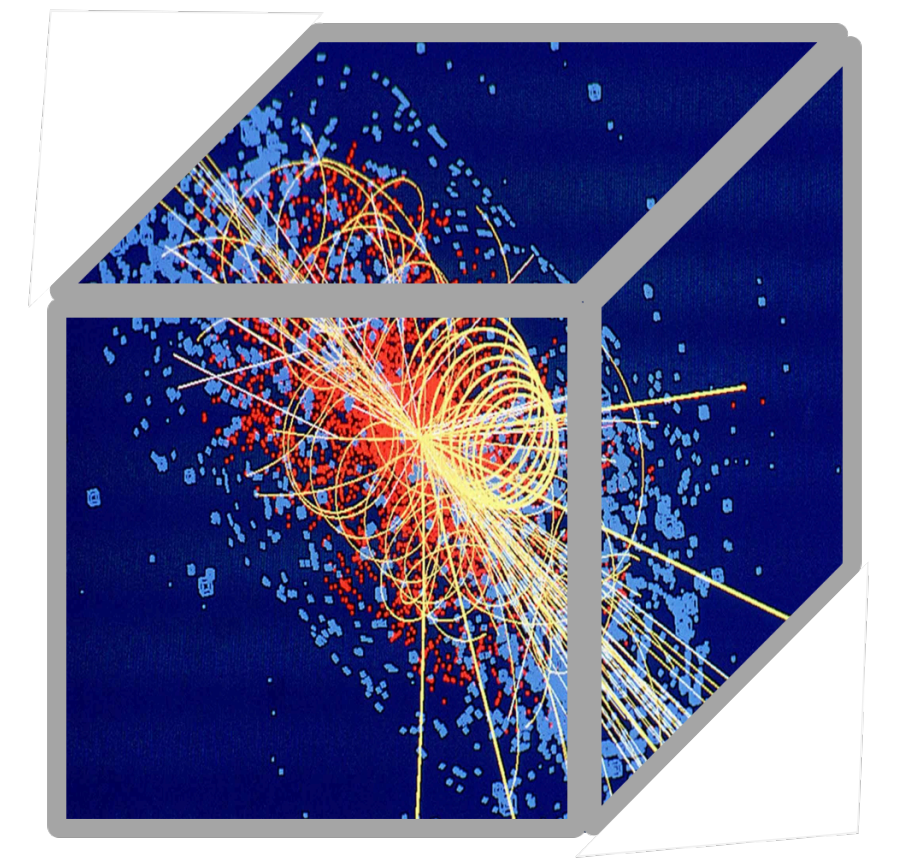
- Hamiltonian truncation = non-perturbative method for computing observables in strongly coupled QFTs
- effective field theory = powerful tool for compensating for ignorance
- effective field theory techniques **systematically improve** convergence in Hamiltonian truncation calculations
- **crucial** for higher dimensions, more complex theories



Future directions

- **improve this method** (next order, include more UV divergences)
 - work in progress with Rachel Houtz
- look at **new observables** (Wilson loops, entanglement entropy, energy correlators)
 - finite volume S-matrix, work in progress with Carl Beadler, Francesco Riva, Matthew Walters
- move to **higher dimensions**
- include **new fields** (gauge bosons)
- **curved spacetime** (cosmology, S-matrix)
- **connection** to other non-perturbative methods
- ... you tell me!

$$S_{fi} = \langle \Psi_f | \Psi_i \rangle ?$$



Landscape of quantum field theory

Landscape of quantum field theory



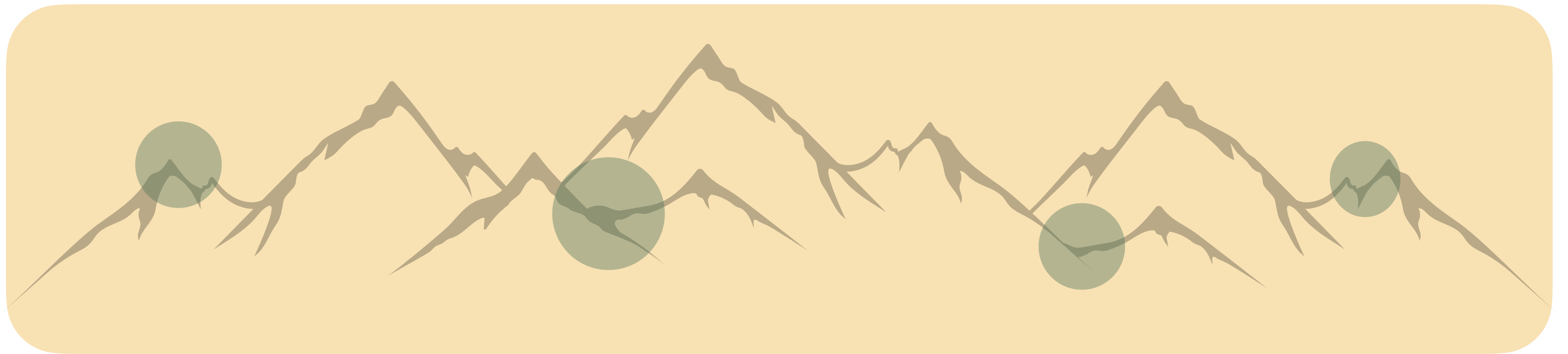
Landscape of quantum field theory



Landscape of quantum field theory

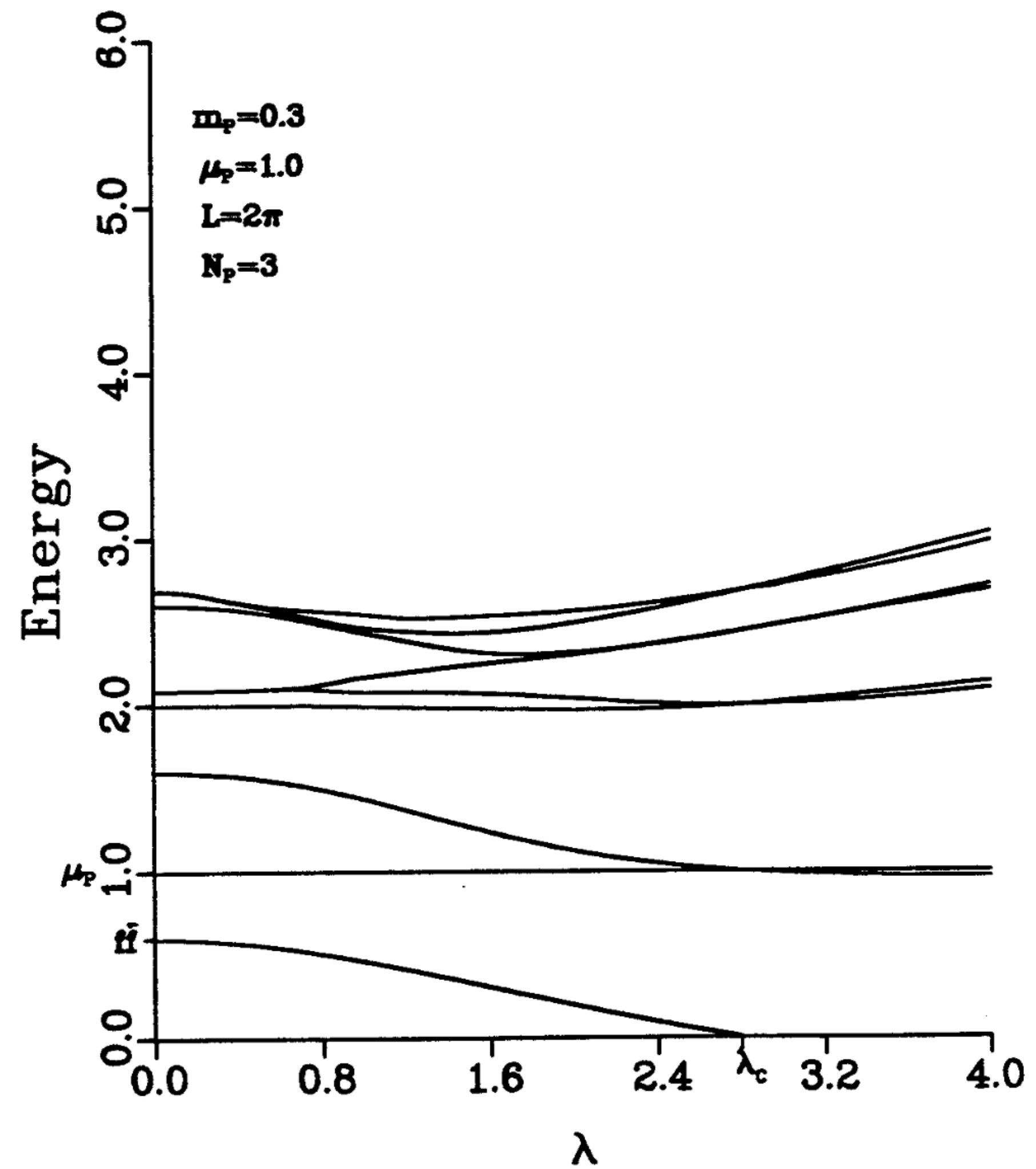
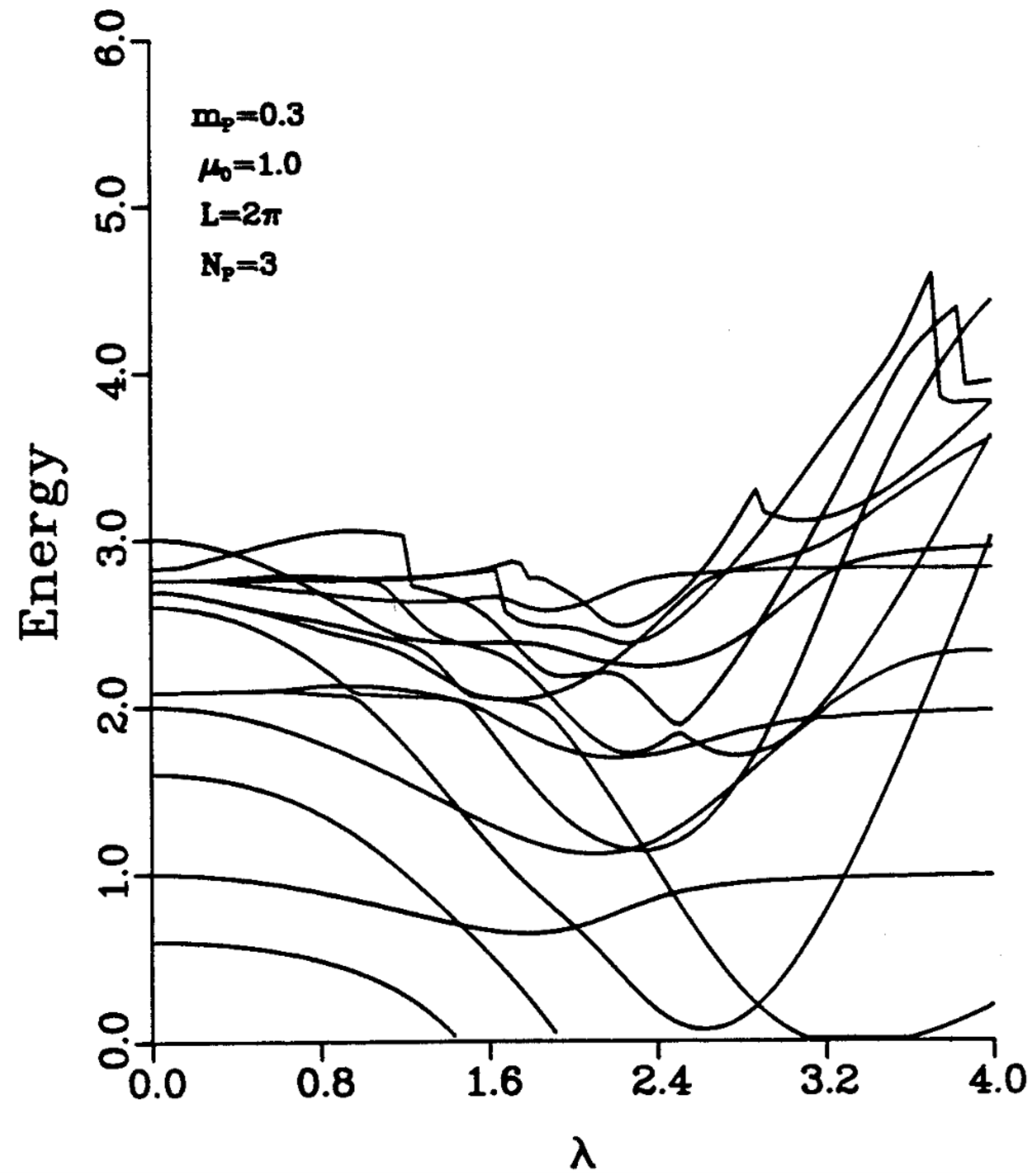


Landscape of quantum field theory

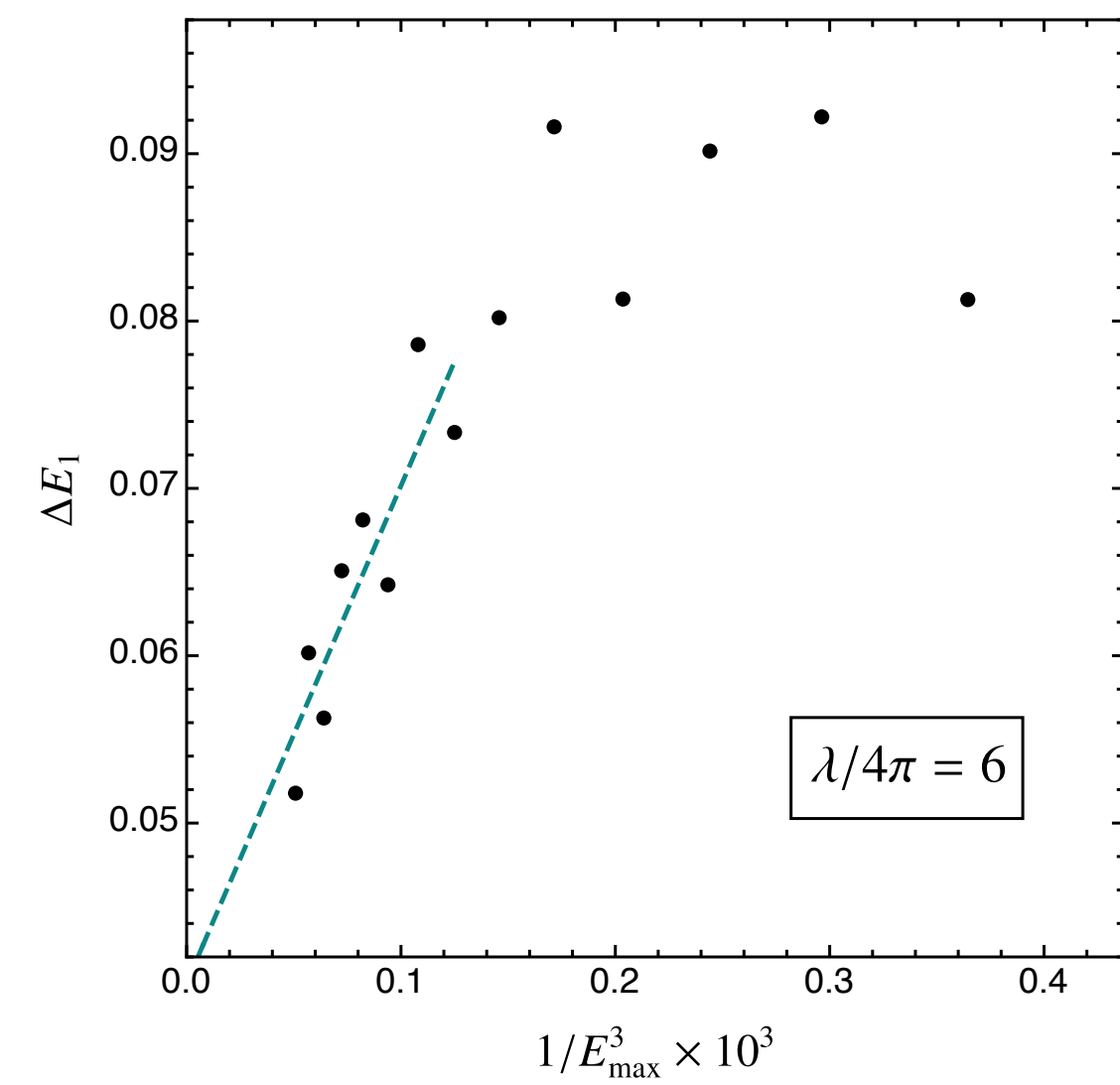
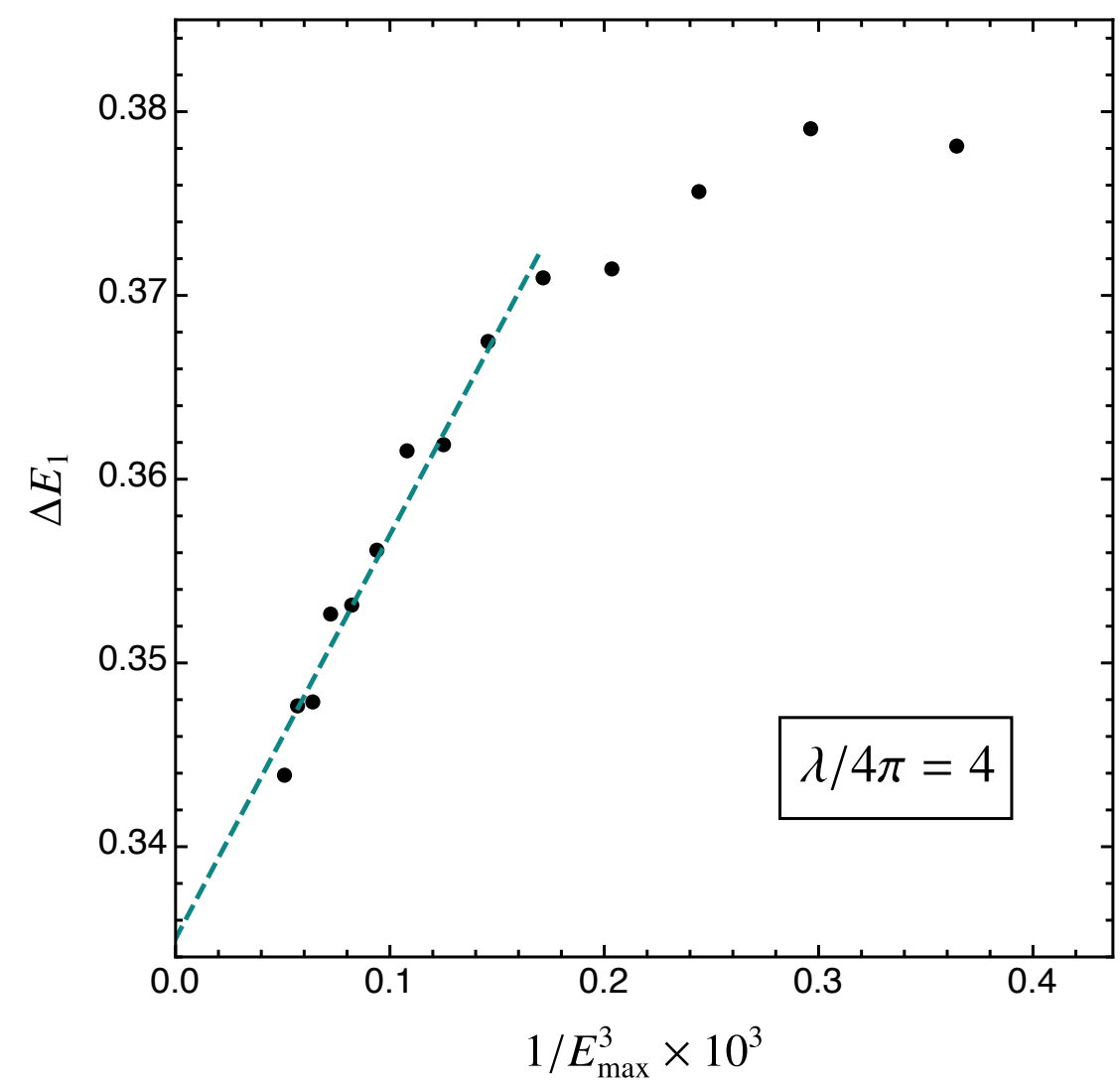
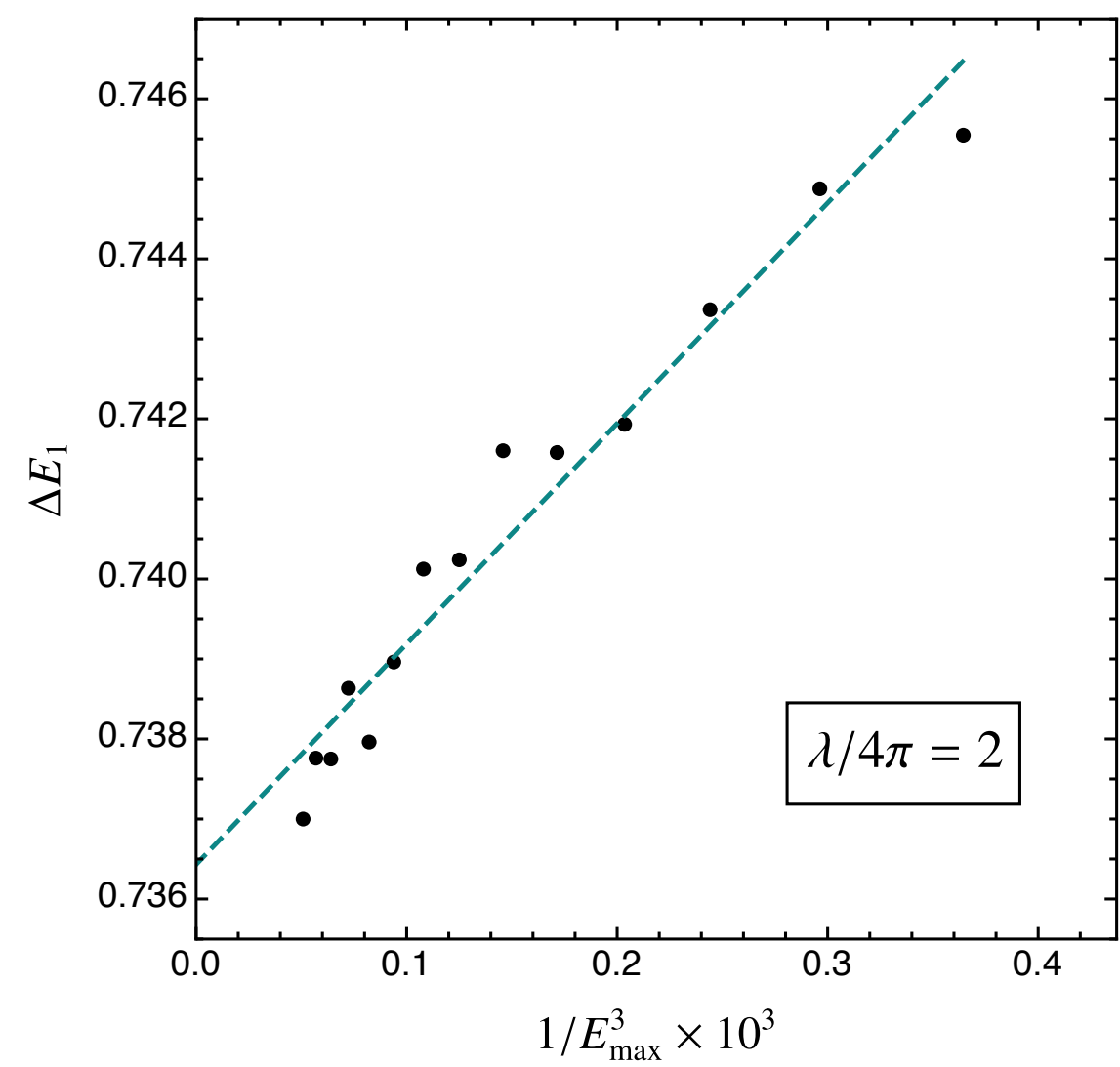
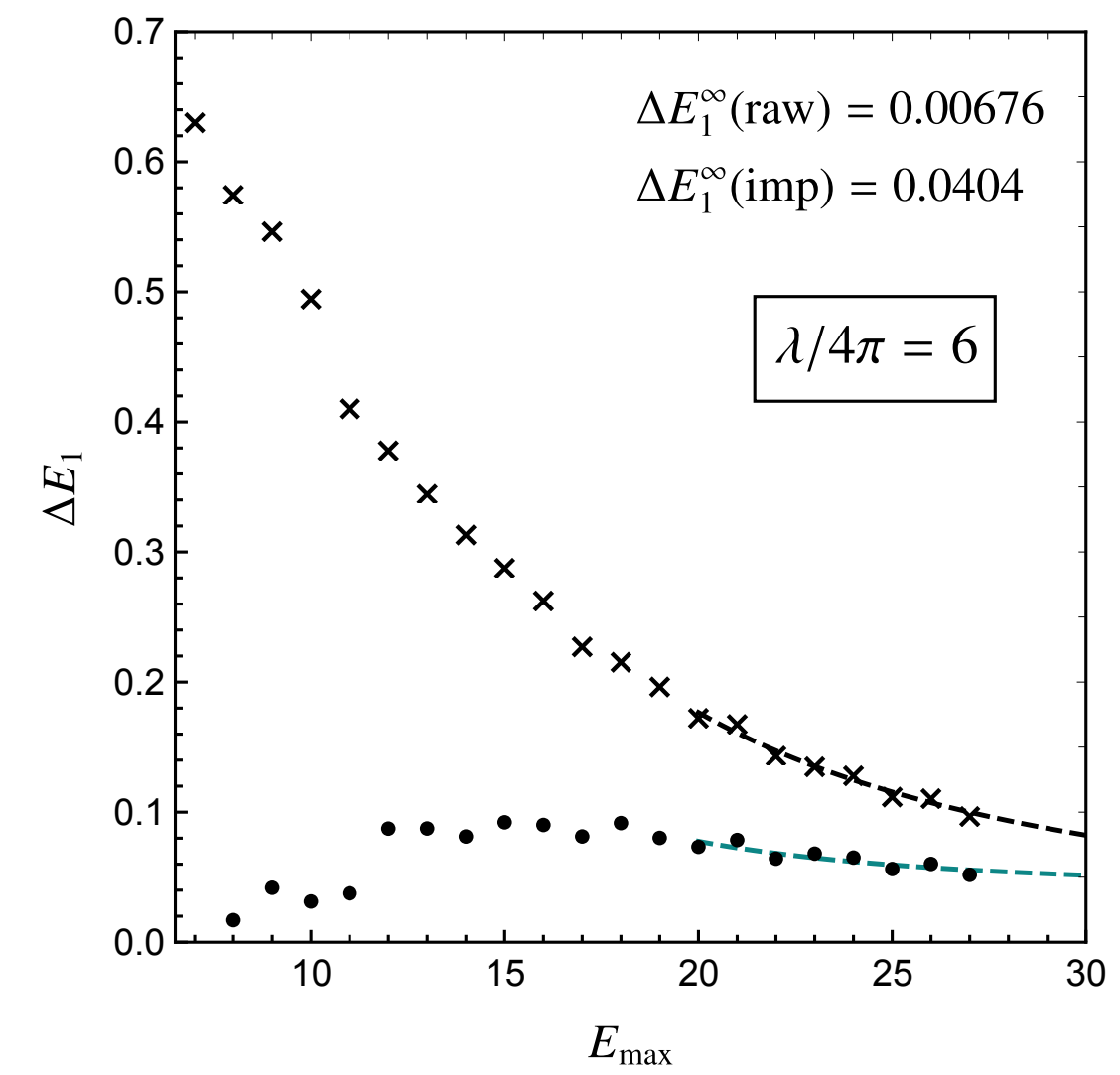
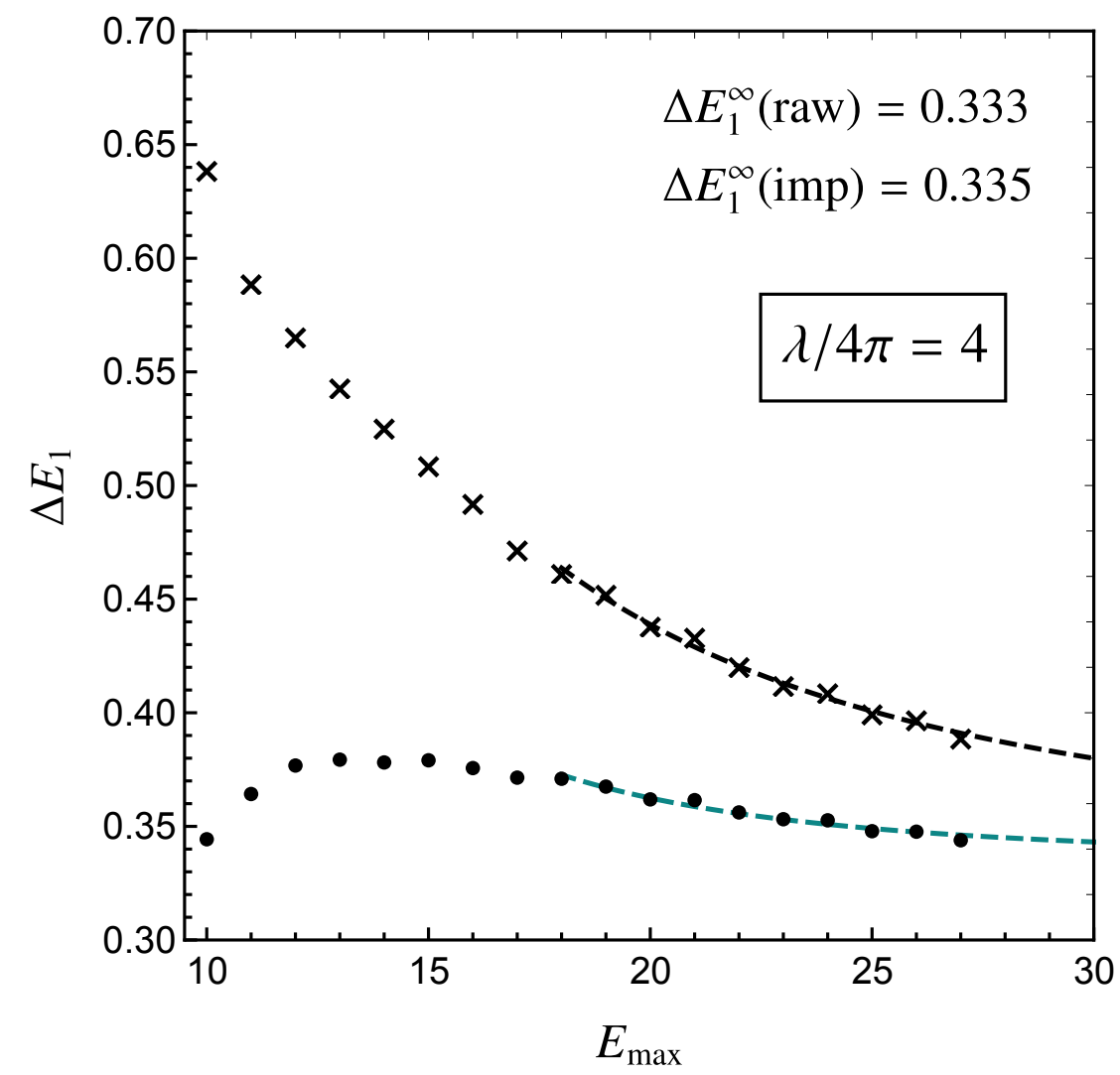
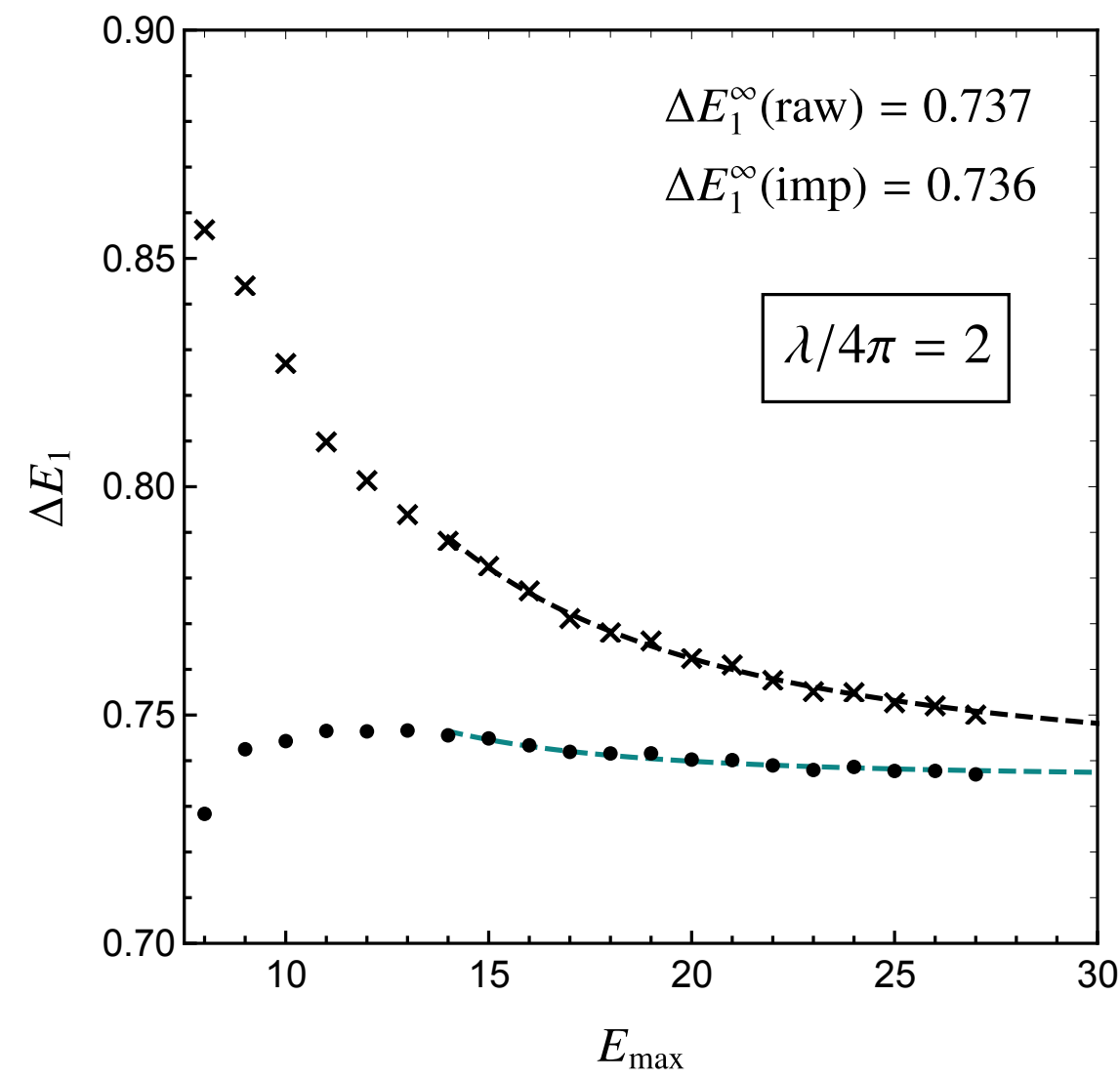


Thank You!

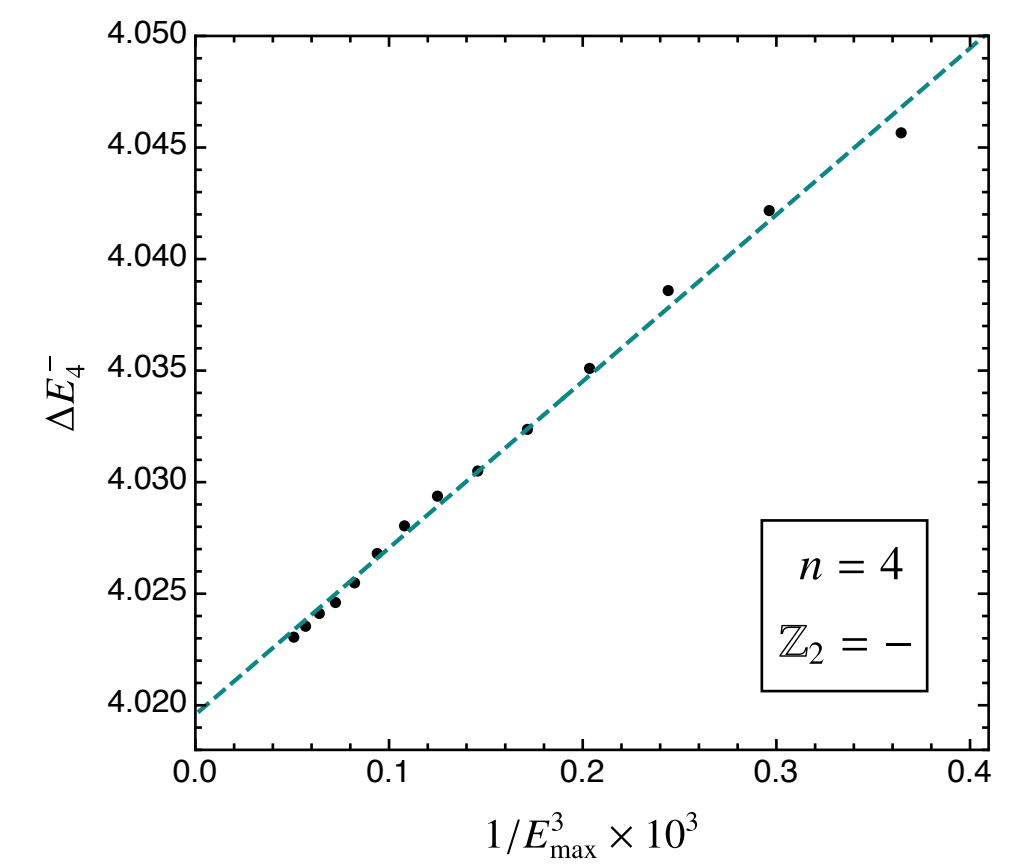
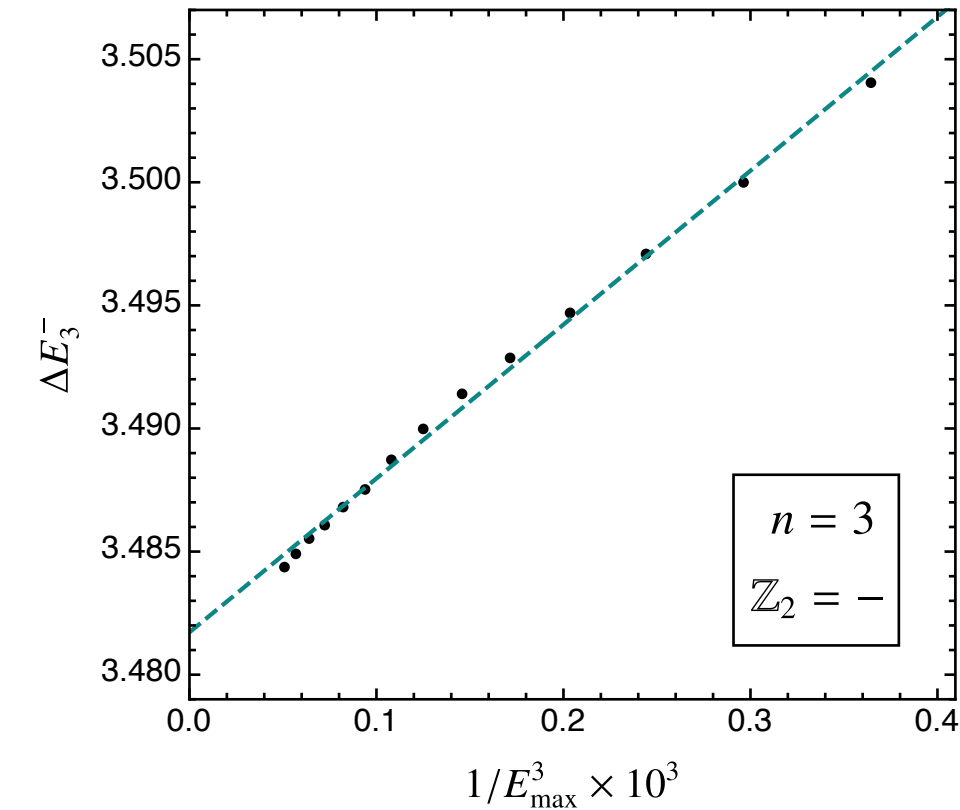
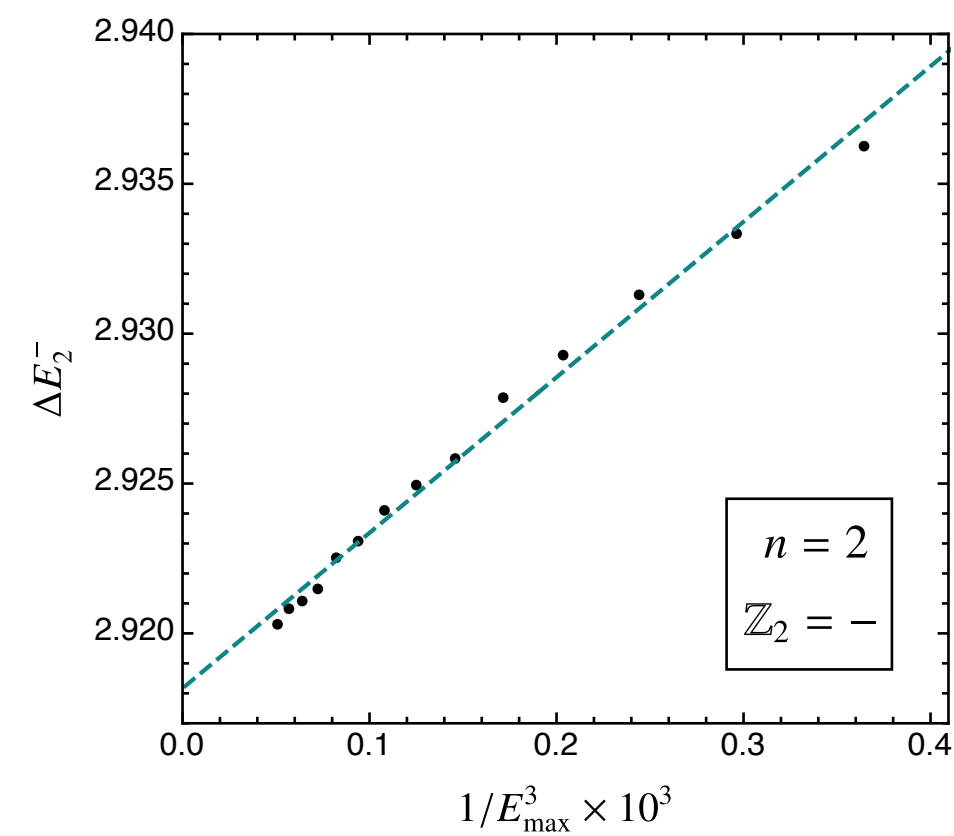
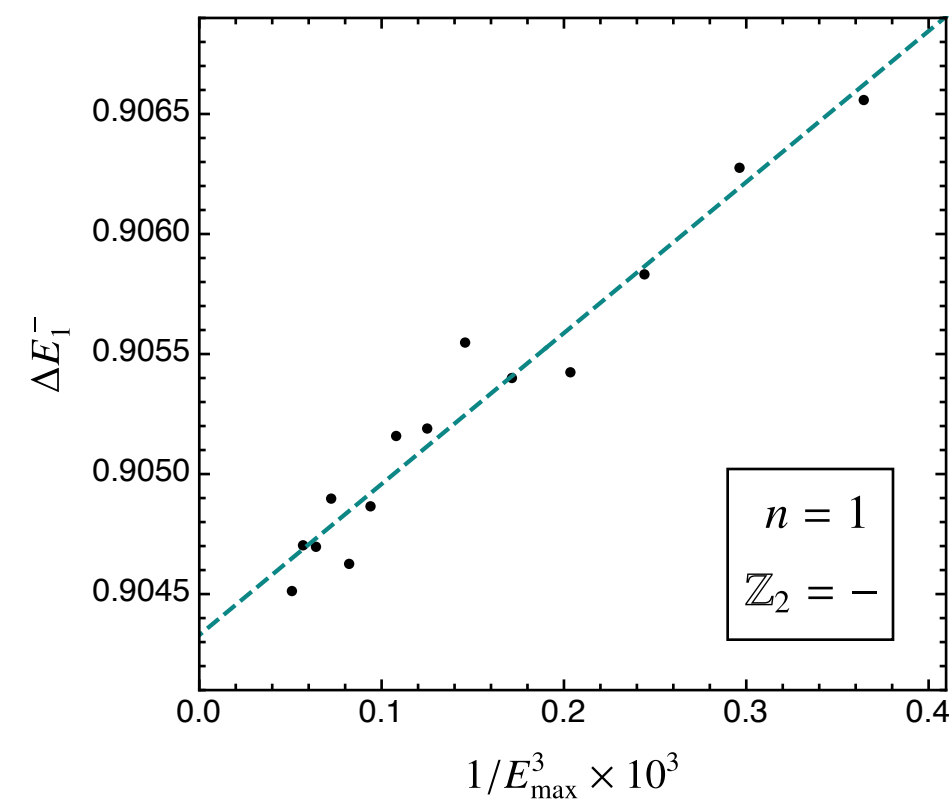
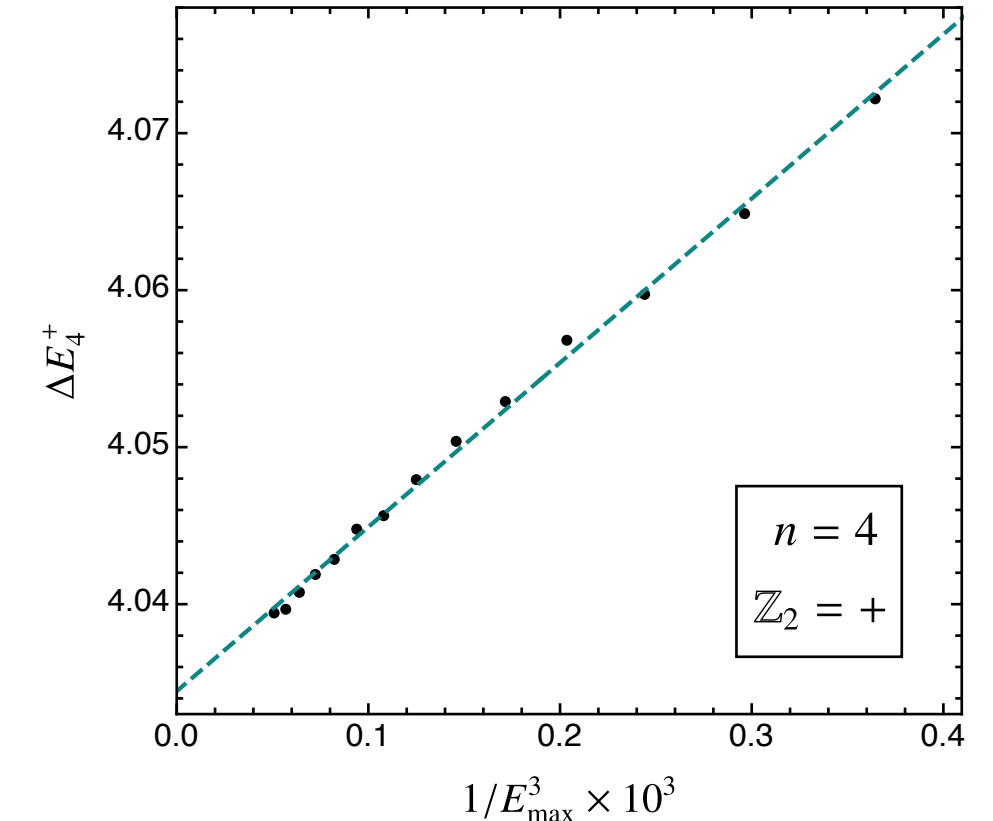
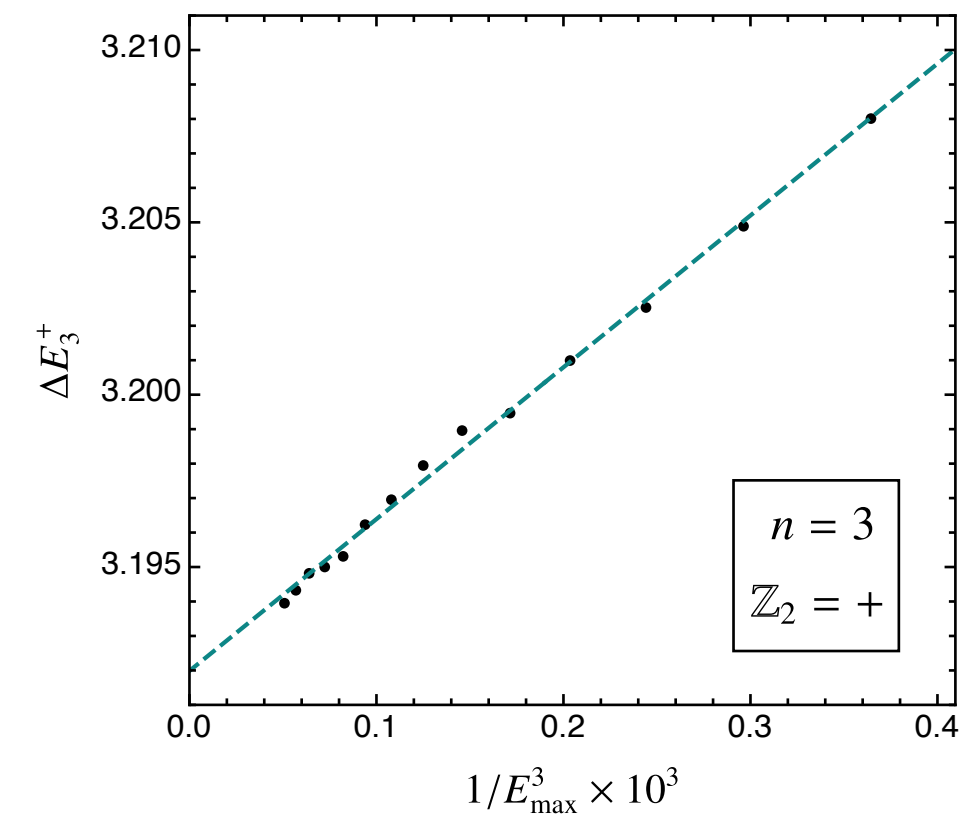
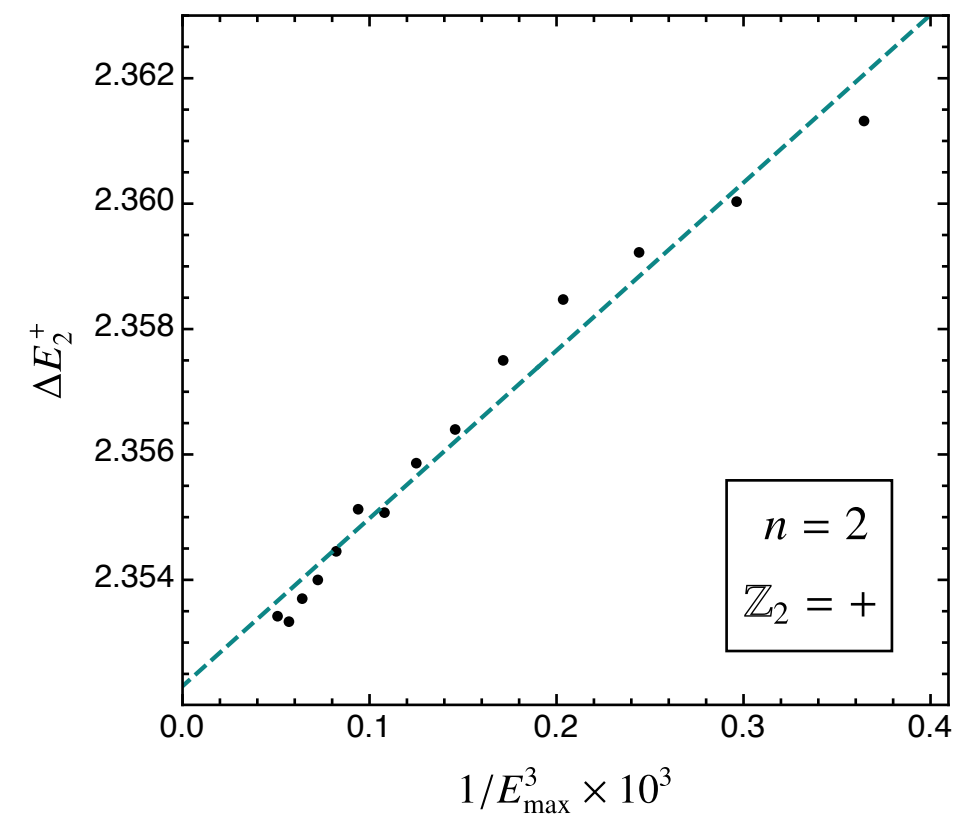
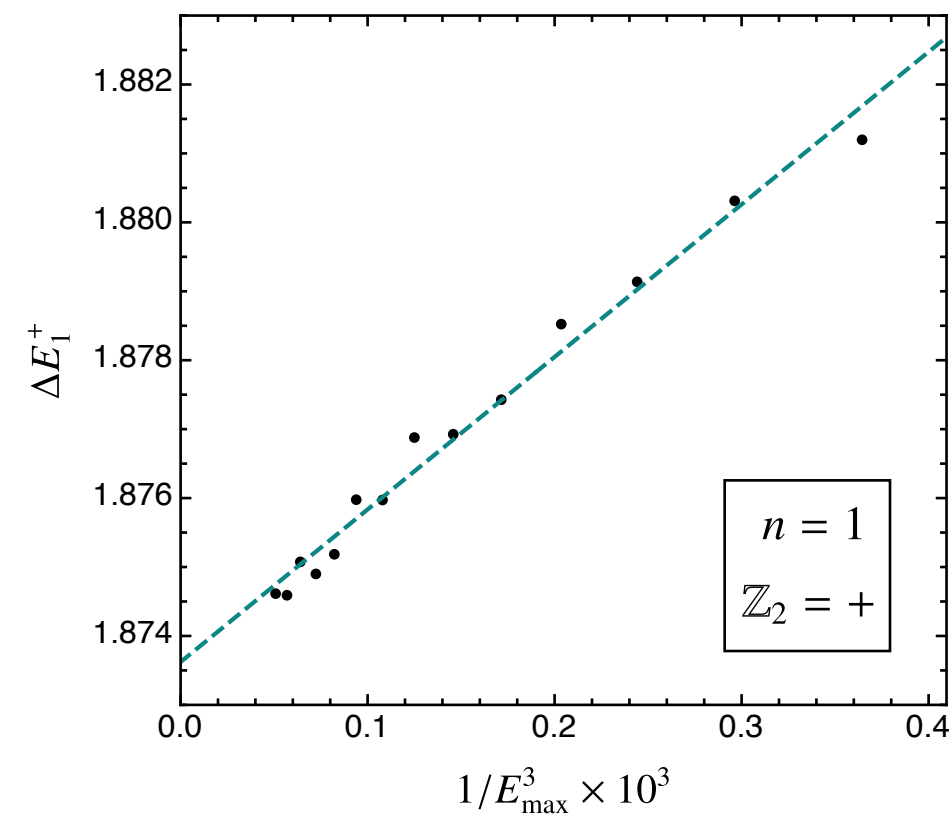
Old paper



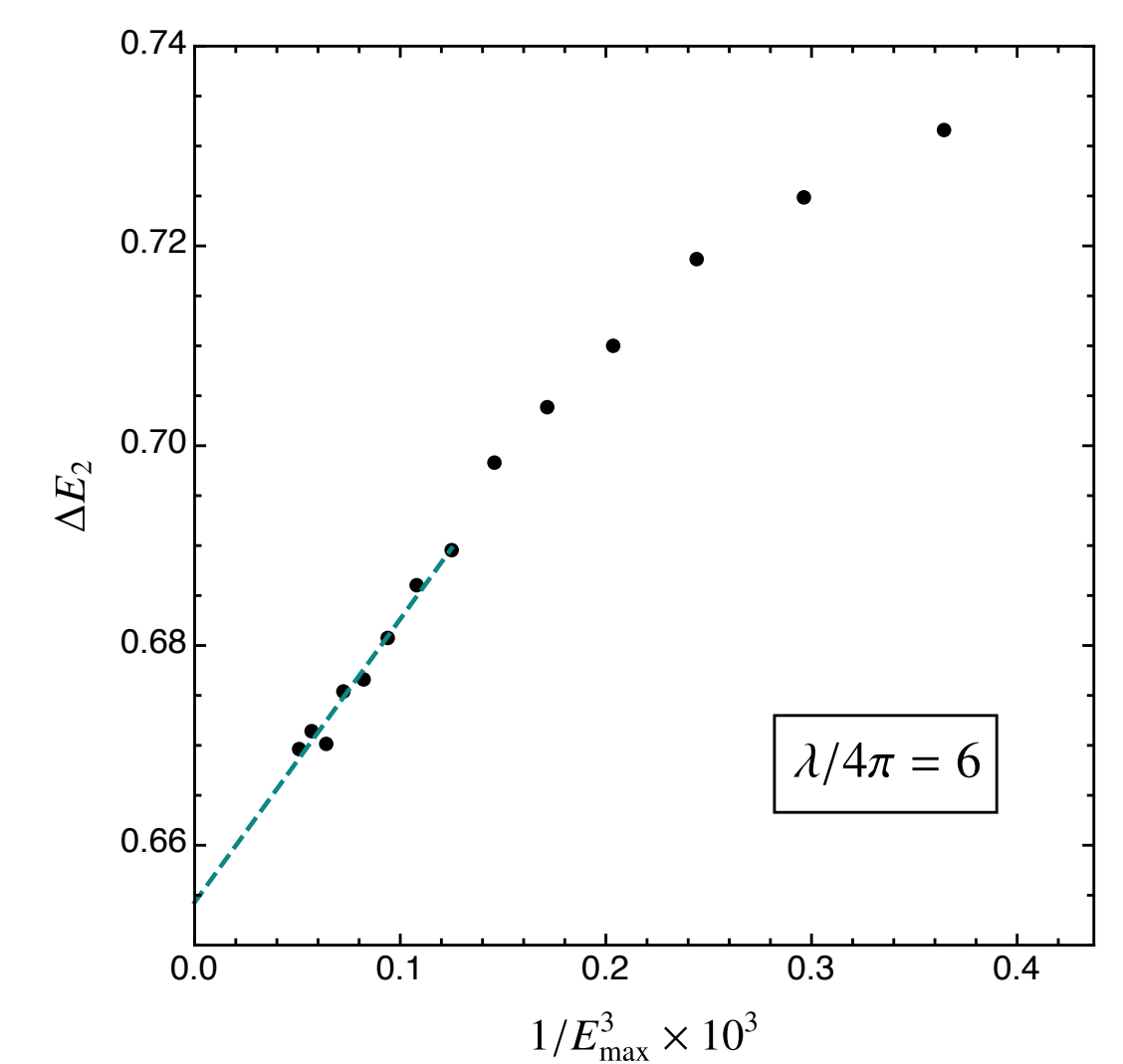
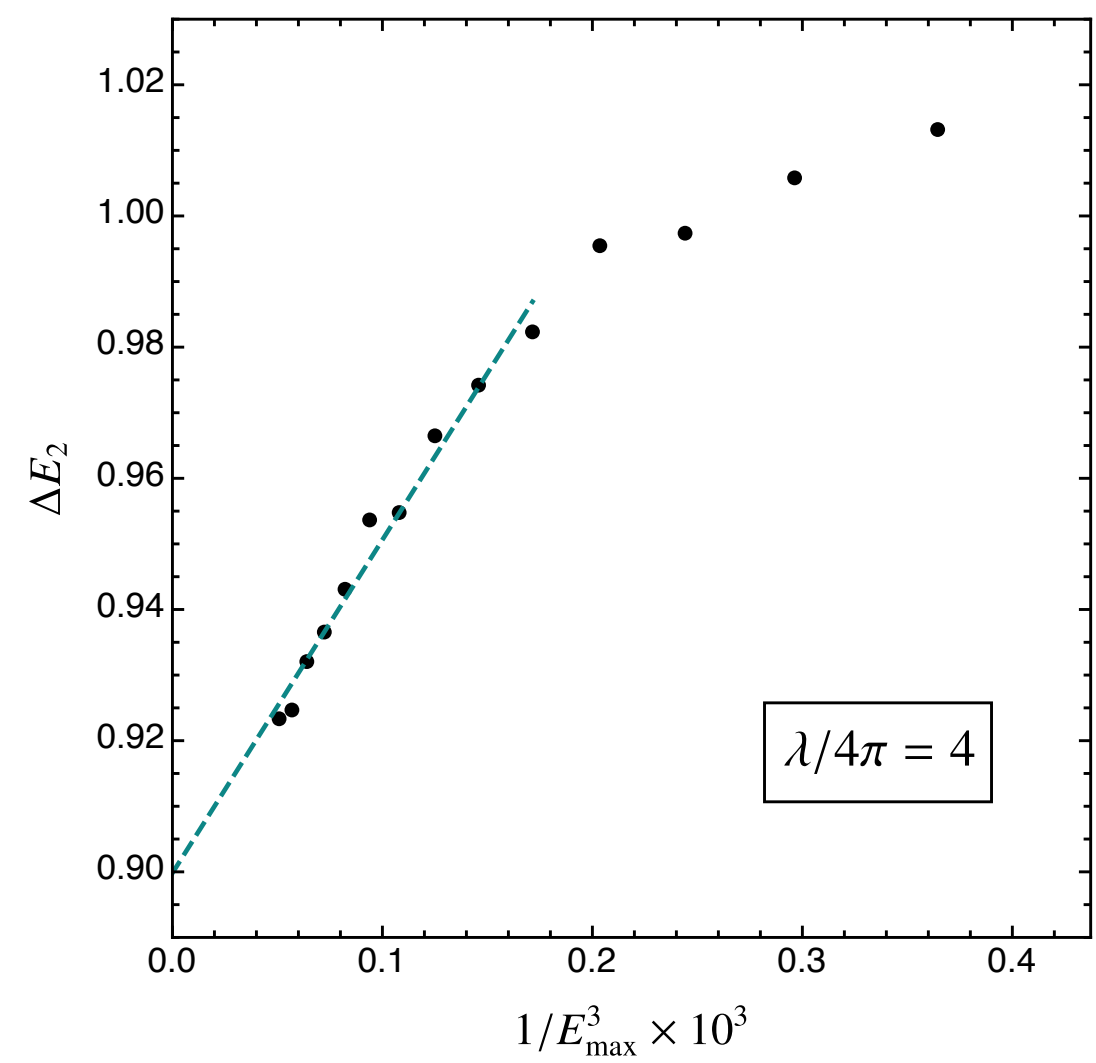
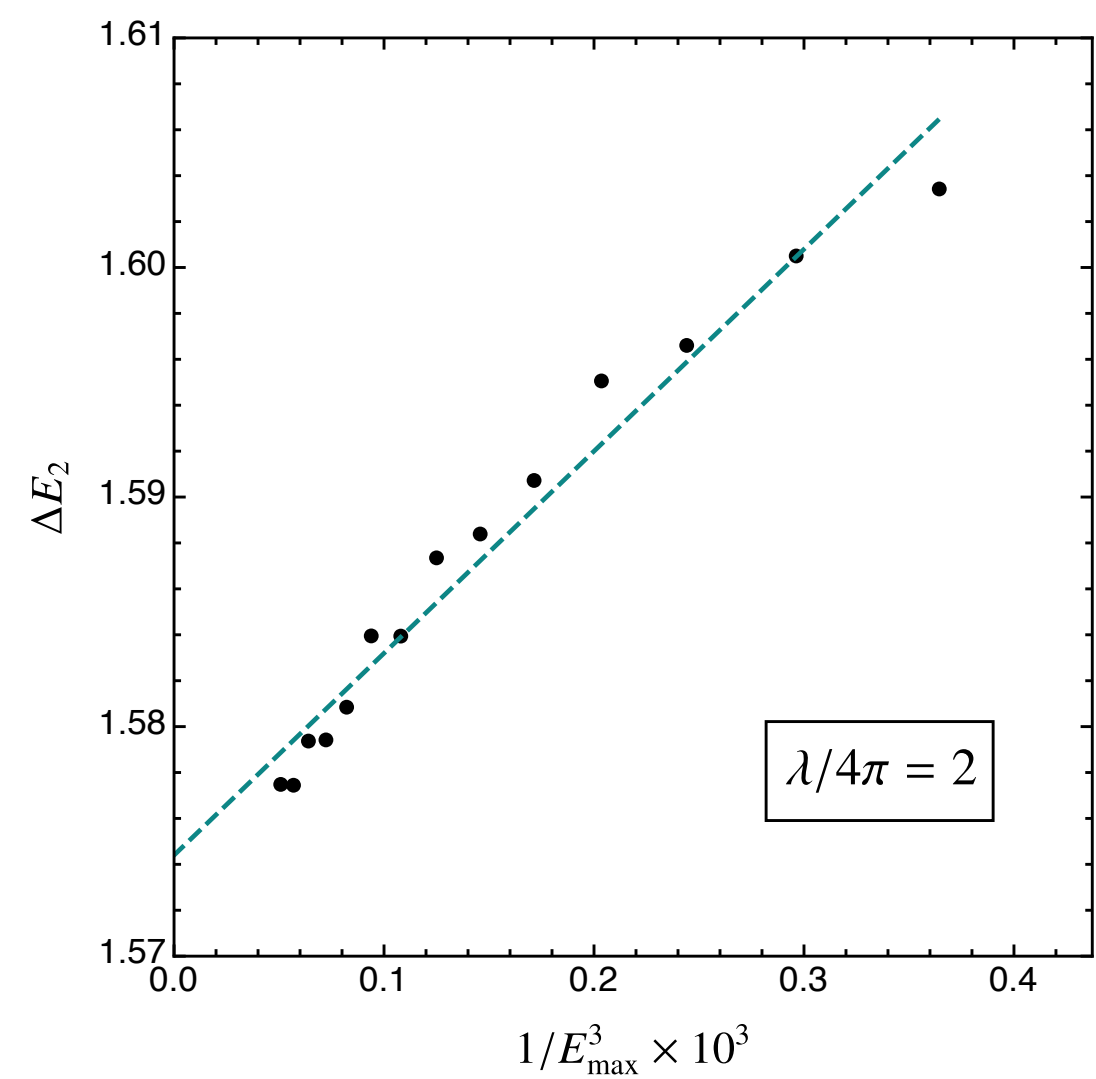
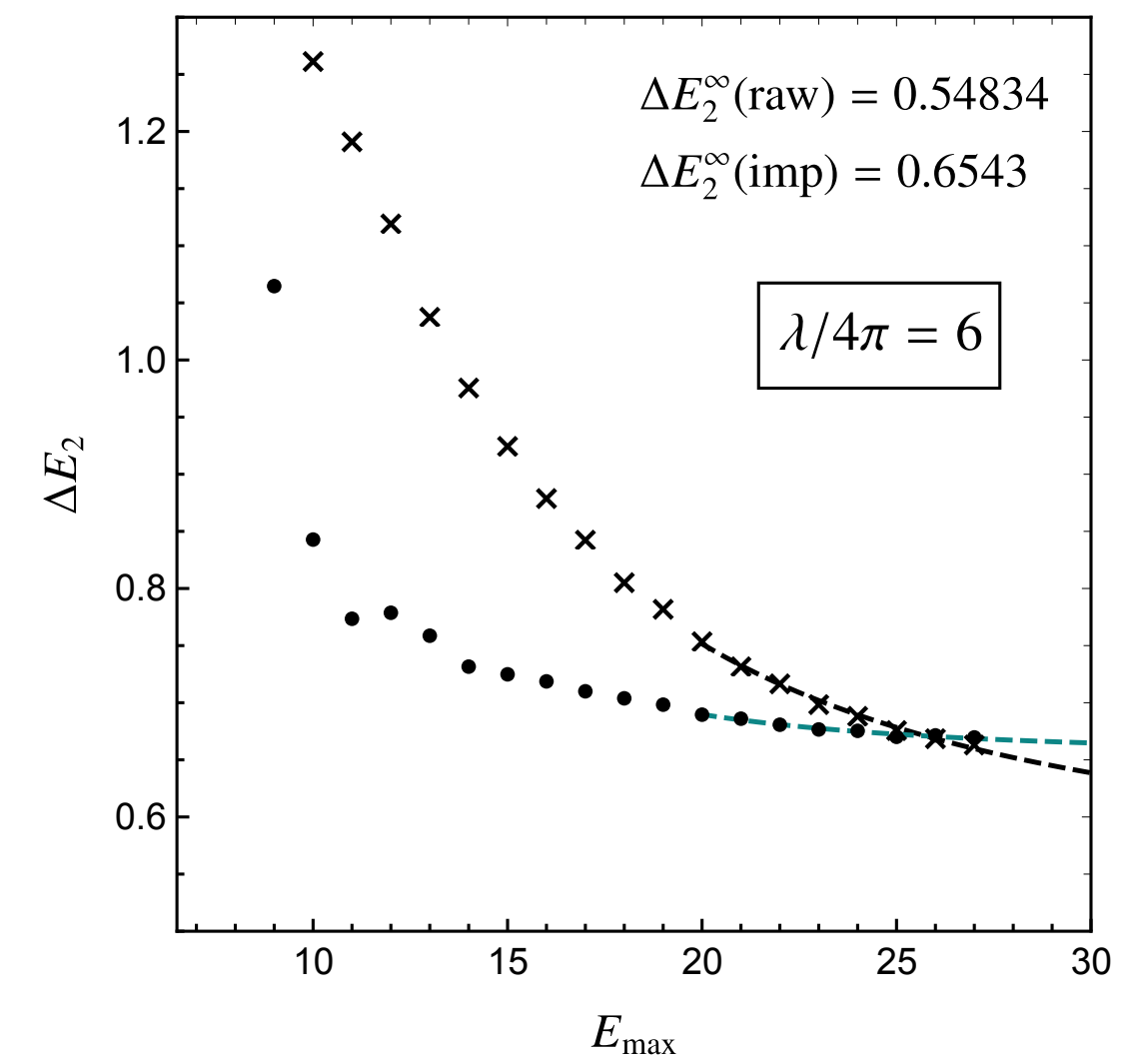
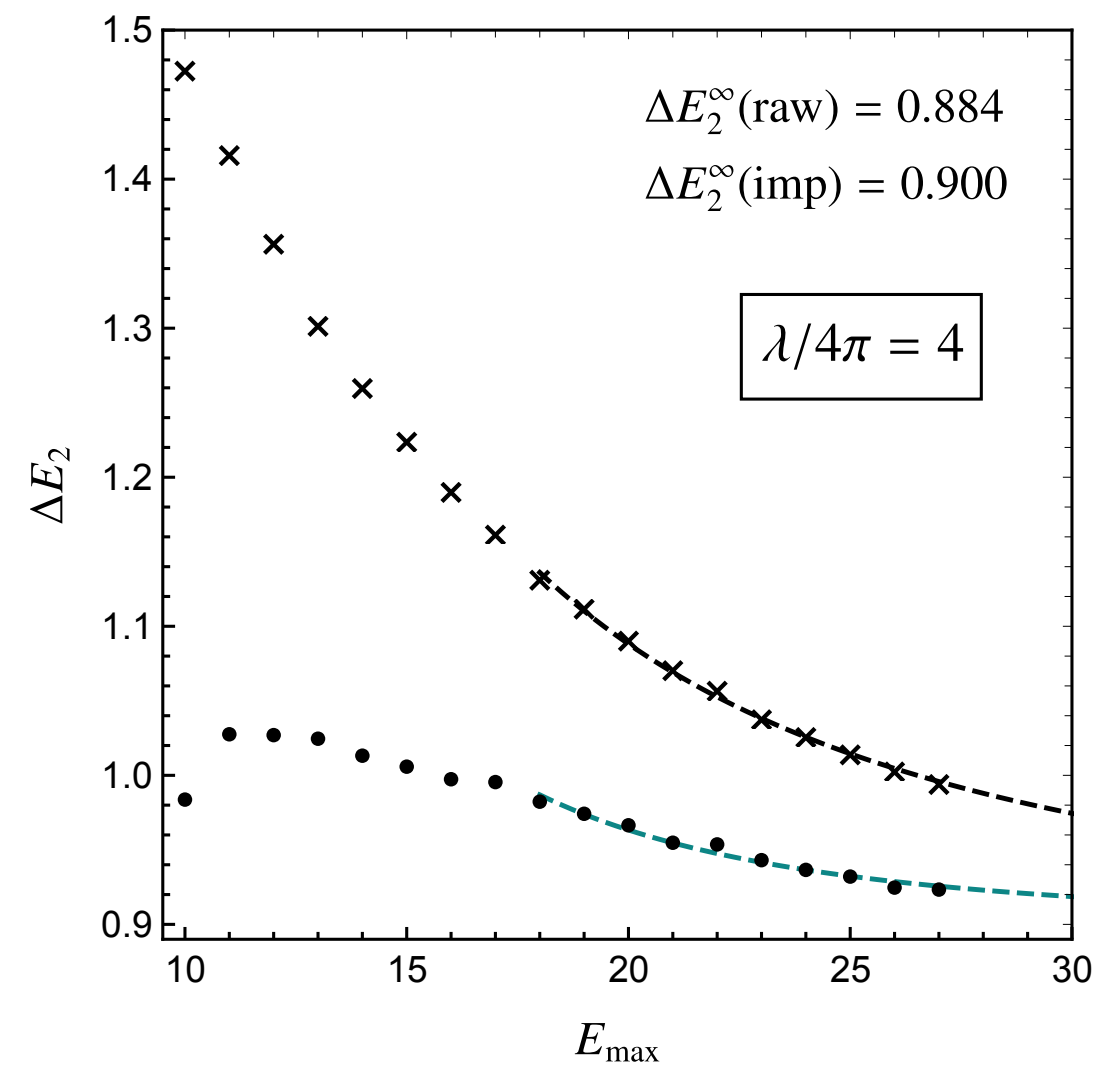
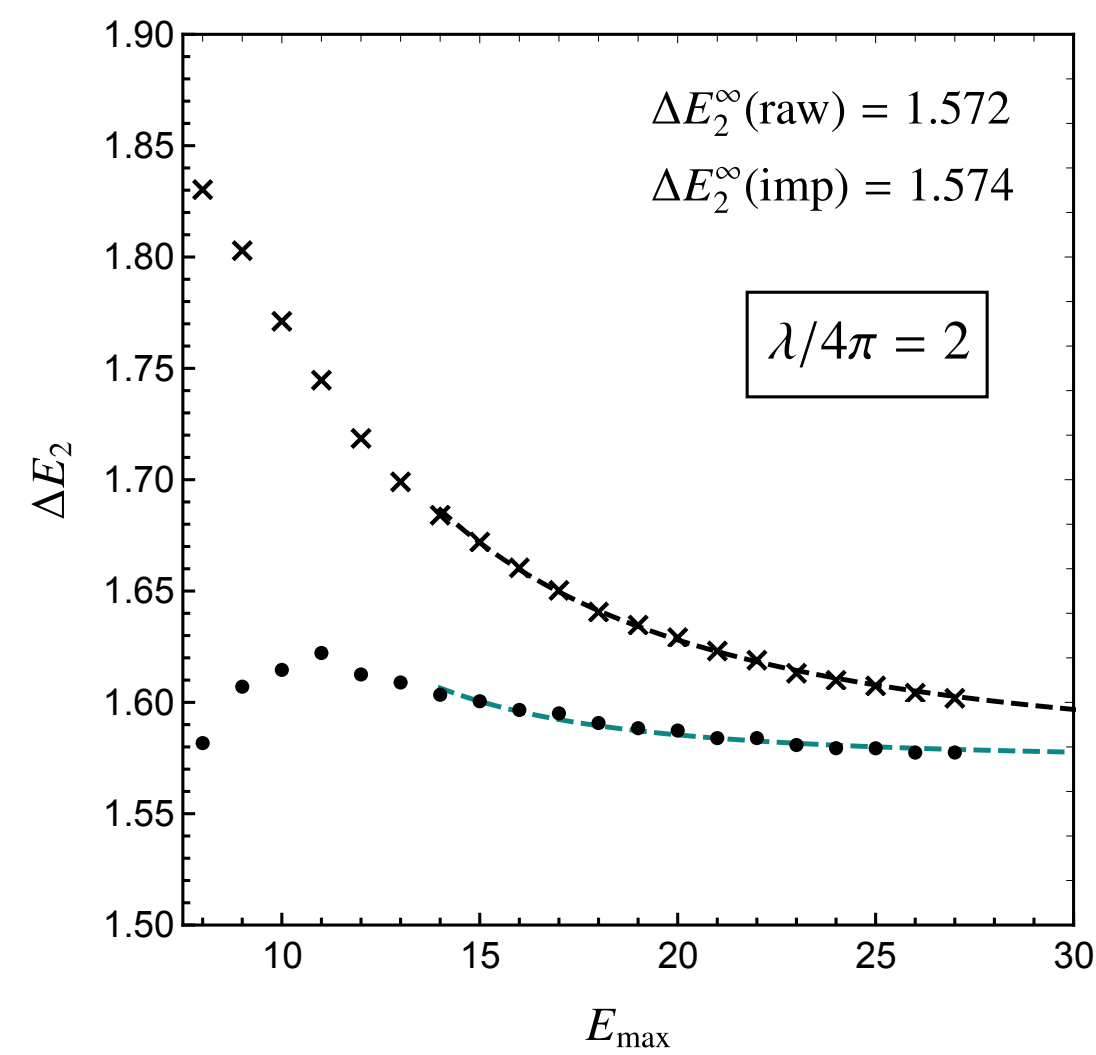
Higher coupling



Higher Excited States



Backup: Higher coupling ΔE_2



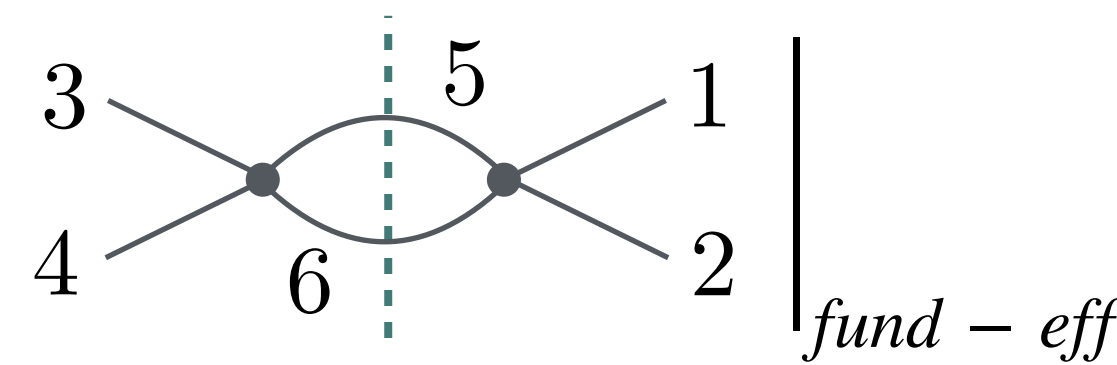
Example Theory 1: 2D $\lambda\phi^4$

$$\langle f | H_2 | i \rangle = \sum_{\alpha \neq i}^{\geq} \frac{\langle f | V | \alpha \rangle \langle \alpha | V | i \rangle}{E_f - E_\alpha}$$

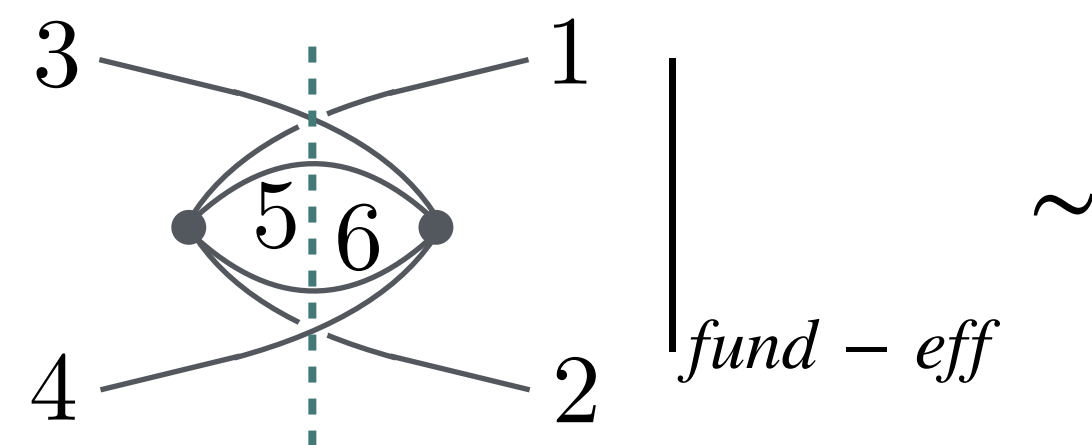
$$\omega_i > 0$$

Example Theory 1: 2D $\lambda\phi^4$

$$\langle f | H_2 | i \rangle = \sum_{\alpha \neq i}^> \frac{\langle f | V | \alpha \rangle \langle \alpha | V | i \rangle}{E_f - E_\alpha}$$



$$\sim \lambda^2 \sum_{5,6} \frac{\Theta(E_\alpha - E_{max})}{\omega_5 \omega_6 (\omega_3 + \omega_4 - \omega_5 - \omega_6)}$$



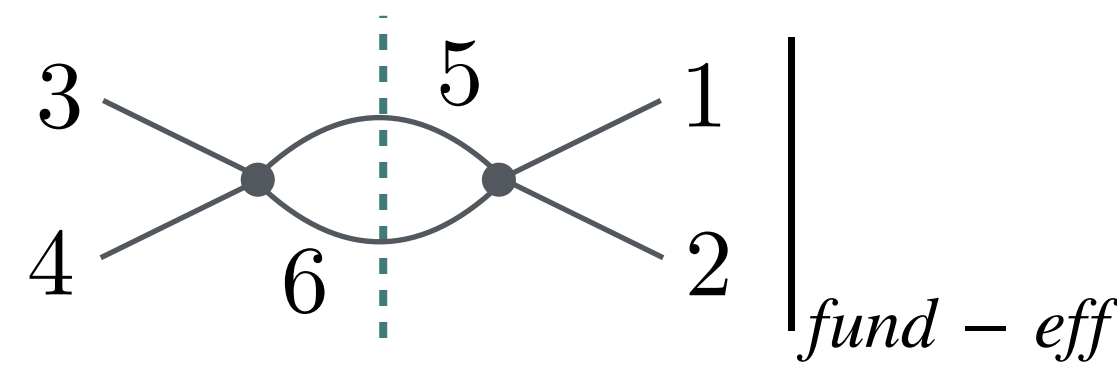
$$\sim \lambda^2 \sum_{5,6} \frac{\Theta(E_\alpha - E_{max})}{\omega_5 \omega_6 (-\omega_1 - \omega_2 - \omega_5 - \omega_6)}$$

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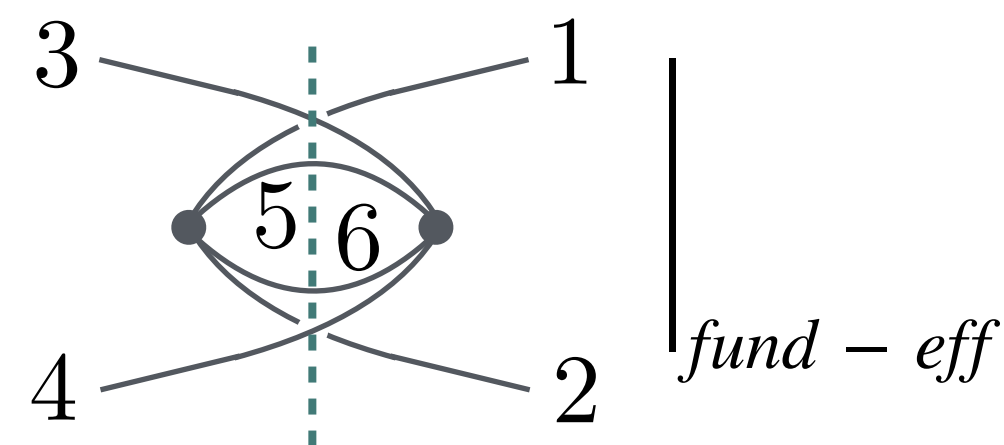
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$$\sim \lambda^2 \sum_{5,6} \frac{\Theta(E_\alpha - E_{max})}{\omega_5 \omega_6 (\omega_3 + \omega_4 - \omega_5 - \omega_6)}$$

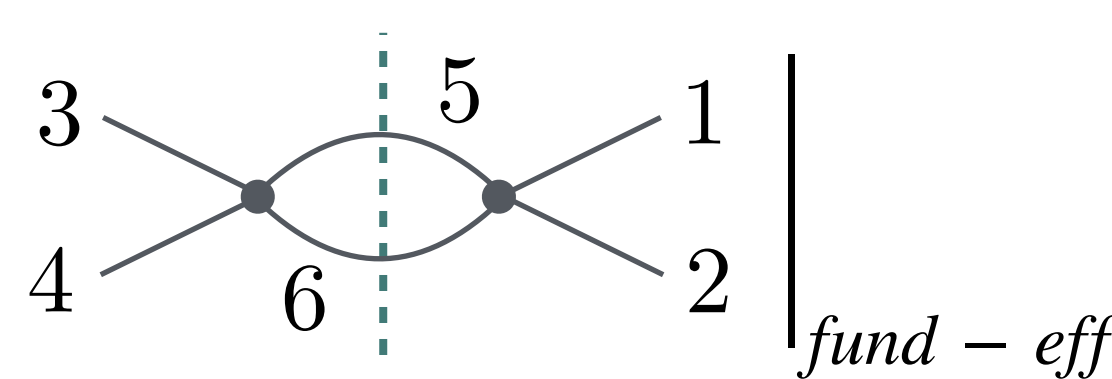


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Example Theory 1: 2D $\lambda\phi^4$

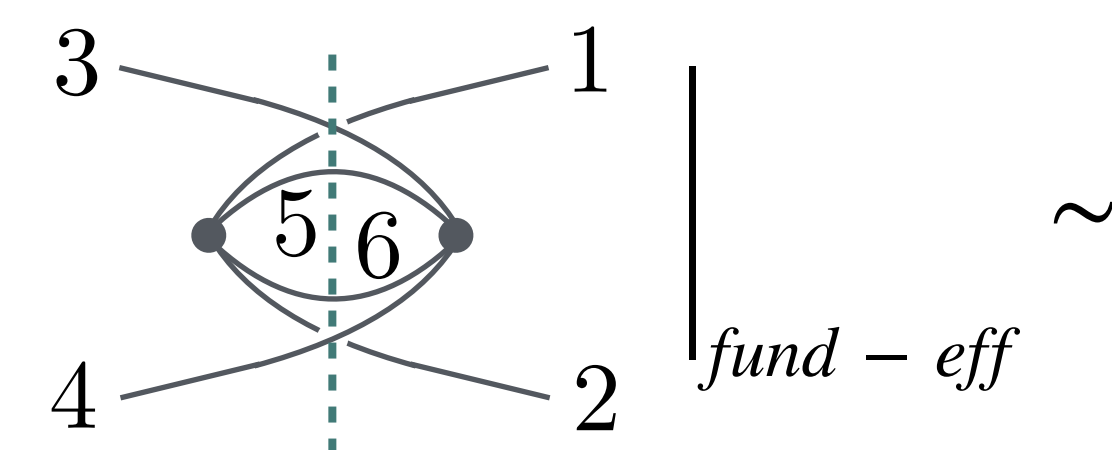
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$$E_\alpha = E_f - (\omega_3 + \omega_4 - \omega_5 - \omega_6)$$

$$\sim \lambda^2 \sum_{5,6} \frac{\Theta(E_\alpha - E_{max})}{\omega_5 \omega_6 (\omega_3 + \omega_4 - \omega_5 - \omega_6)}$$

fund - eff



$$E_\alpha = E_f - (-\omega_1 - \omega_2 - \omega_5 - \omega_6)$$

$$\sim \lambda^2 \sum_{5,6} \frac{\Theta(E_\alpha - E_{max})}{\omega_5 \omega_6 (-\omega_1 - \omega_2 - \omega_5 - \omega_6)}$$

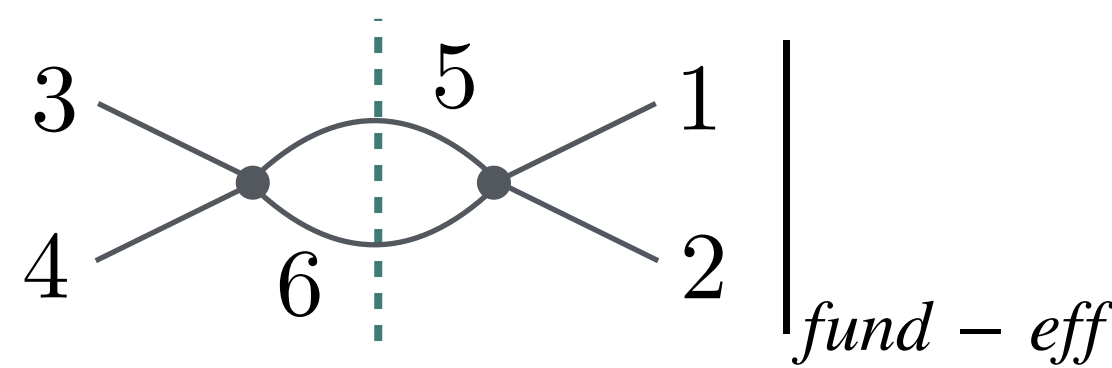
fund - eff

$$\omega_i > 0$$

Example Theory 1: 2D $\lambda\phi^4$

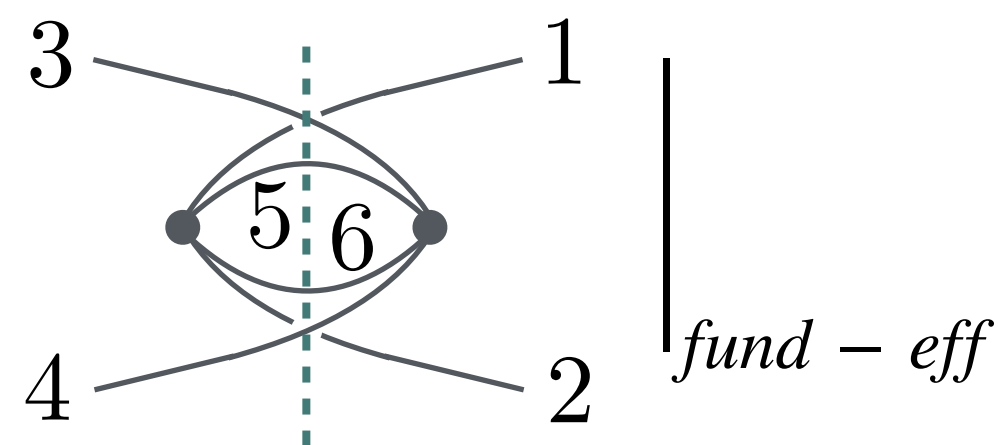
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$$E_\alpha = E_f - (\omega_3 + \omega_4 - \omega_5 - \omega_6)$$

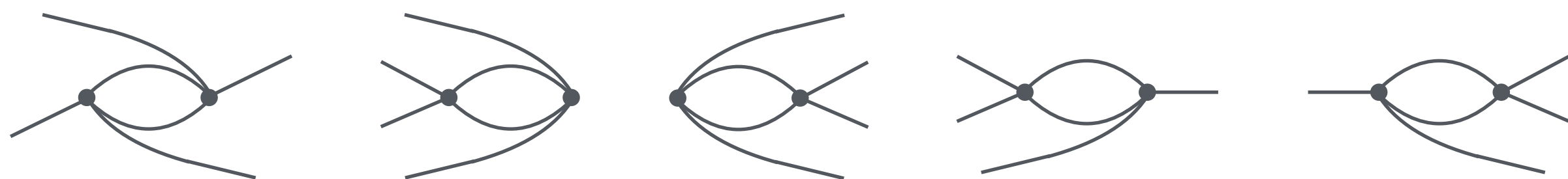


$$\sim \lambda^2 \sum_{5,6} \frac{\Theta(E_\alpha - E_{max})}{\omega_5 \omega_6 (\omega_3 + \omega_4 - \omega_5 - \omega_6)}$$

$$E_\alpha = E_f - (-\omega_1 - \omega_2 - \omega_5 - \omega_6)$$



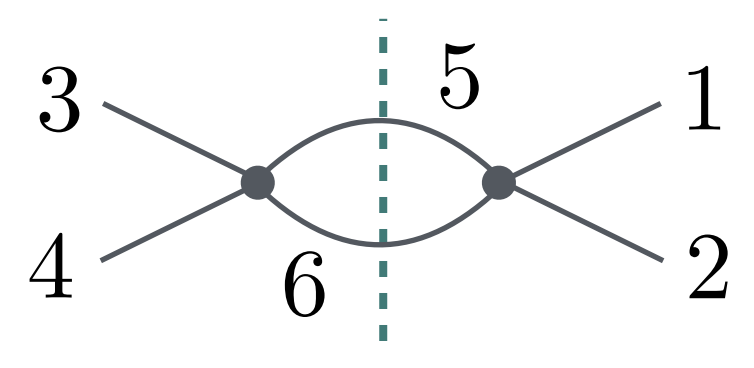
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Example Theory 1: 2D $\lambda\phi^4$

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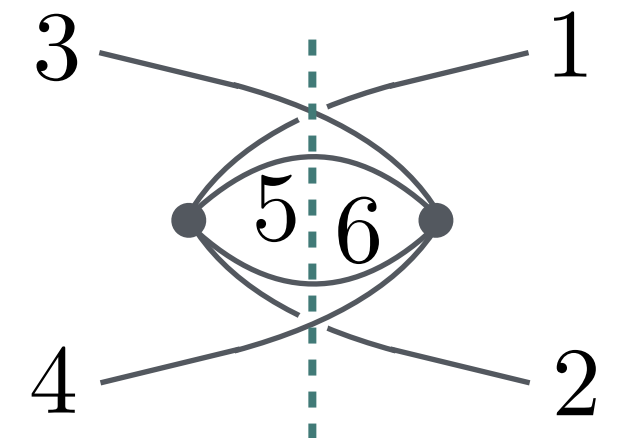
$$E_\alpha = E_f - (\omega_3 + \omega_4 - \omega_5 - \omega_6)$$

$$\sim \lambda^2 \sum_{5,6} \frac{\Theta(E_\alpha - E_{max})}{\omega_5 \omega_6 (\omega_3 + \omega_4 - \omega_5 - \omega_6)}$$

fund - eff

local approximation

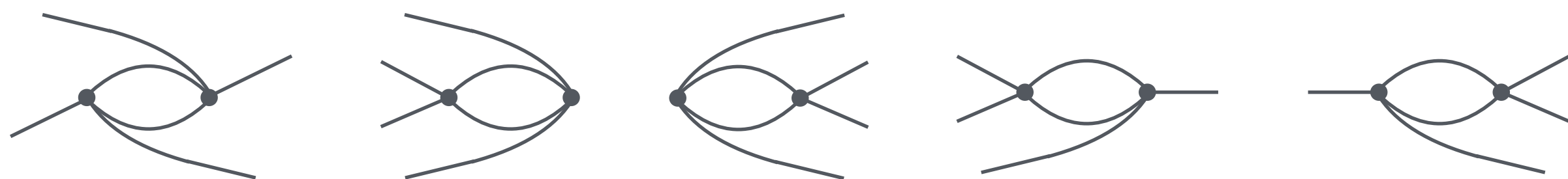
$$\omega_{1,2,3,4} \lesssim E_{i,f} \ll E_{max}$$



$$E_\alpha = E_f - (-\omega_1 - \omega_2 - \omega_5 - \omega_6)$$

$$\sim \lambda^2 \sum_{5,6} \frac{\Theta(E_\alpha - E_{max})}{\omega_5 \omega_6 (-\omega_1 - \omega_2 - \omega_5 - \omega_6)}$$

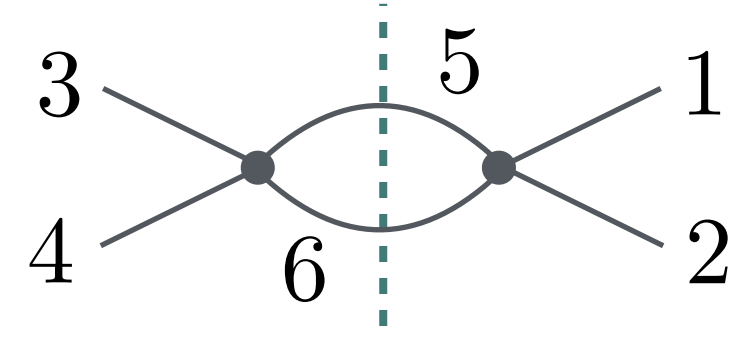
fund - eff



$$\omega_i > 0$$

Example Theory 1: 2D $\lambda\phi^4$

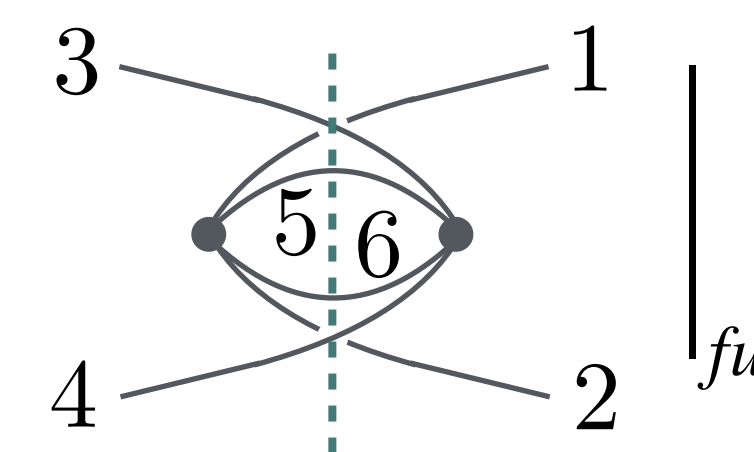
$$\langle f | H_2 | i \rangle = \sum_{\alpha \neq i}^> \frac{\langle f | V | \alpha \rangle \langle \alpha | V | i \rangle}{E_f - E_\alpha}$$



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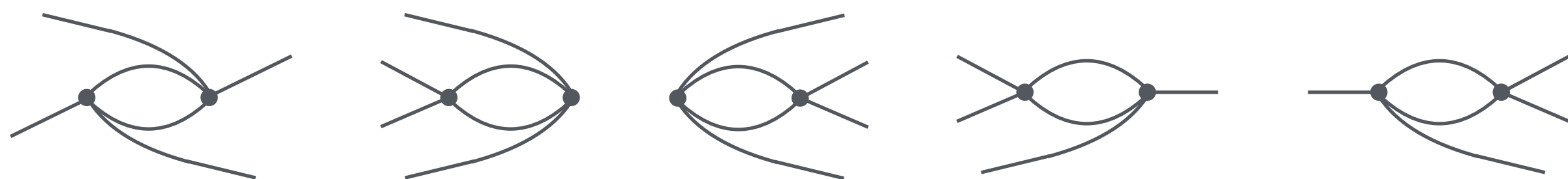
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local approximation

$$\omega_{1,2,3,4} \lesssim E_{i,f} \ll E_{max}$$

$$\sim \lambda^2 \sum_k \frac{\Theta(2\omega_k - E_{max})}{\omega_k^2 (-2\omega_k)} \int dx \phi^4$$



$$\omega_i > 0$$

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