

# Spin-1 Thermal Targets for Dark Matter **Searches at Beam Dump and Fixed Target** Experiments

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# **Sub-GeV Dark Matter**

- Direct detection experiments are becoming more sensitive, with **no signal yet**.
- Instead of being GeV TeV scale like the canonical WIMP, dark matter could have **sub-GeV** mass, **evading nuclear recoil direct detection** experiments.
- Sub-GeV dark matter is a plausible candidate which can account for the **thermal** relic abundance in the universe.
- Lee-Weinberg bound of electroweak mediators is circumvented by adding a

# **Spin-1 Dark Matter**

We **extend** the existing studies on sub-GeV dark matter at fixed target experiments, which include spin-0 and spin-1/2 dark matter models, to include spin-1 dark matter. We consider **simplified spin-1 models** which are directly comparable to the previously studied models, and an **ultraviolet complete model** based on a non-abelian gauge group where dark matter is a spin-1 Strongly Interacting Massive Particle (SIMP).

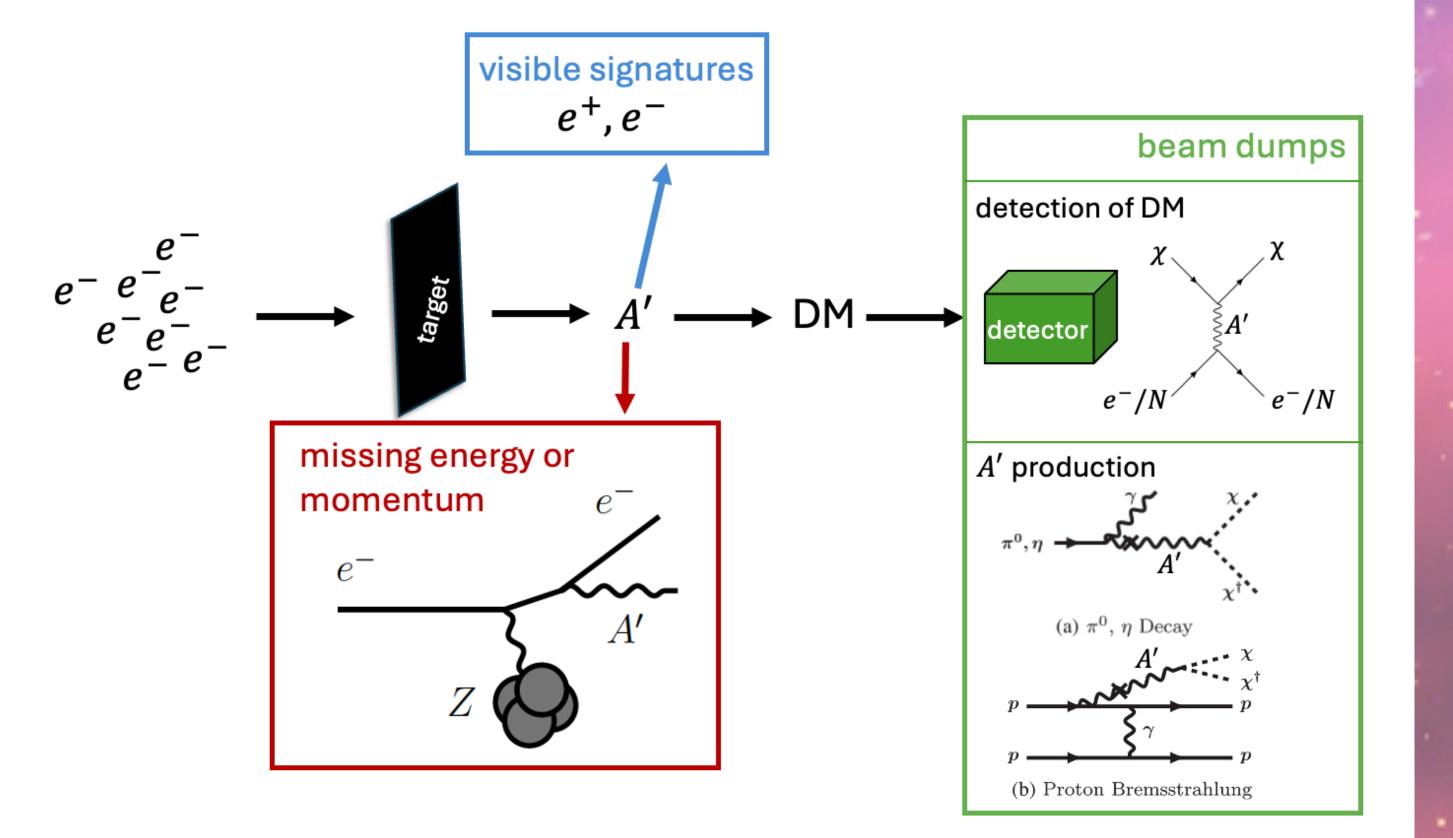
#### Simplified Spin-1 Dark Matter Models

new mediator, for example the **dark photon** A'.

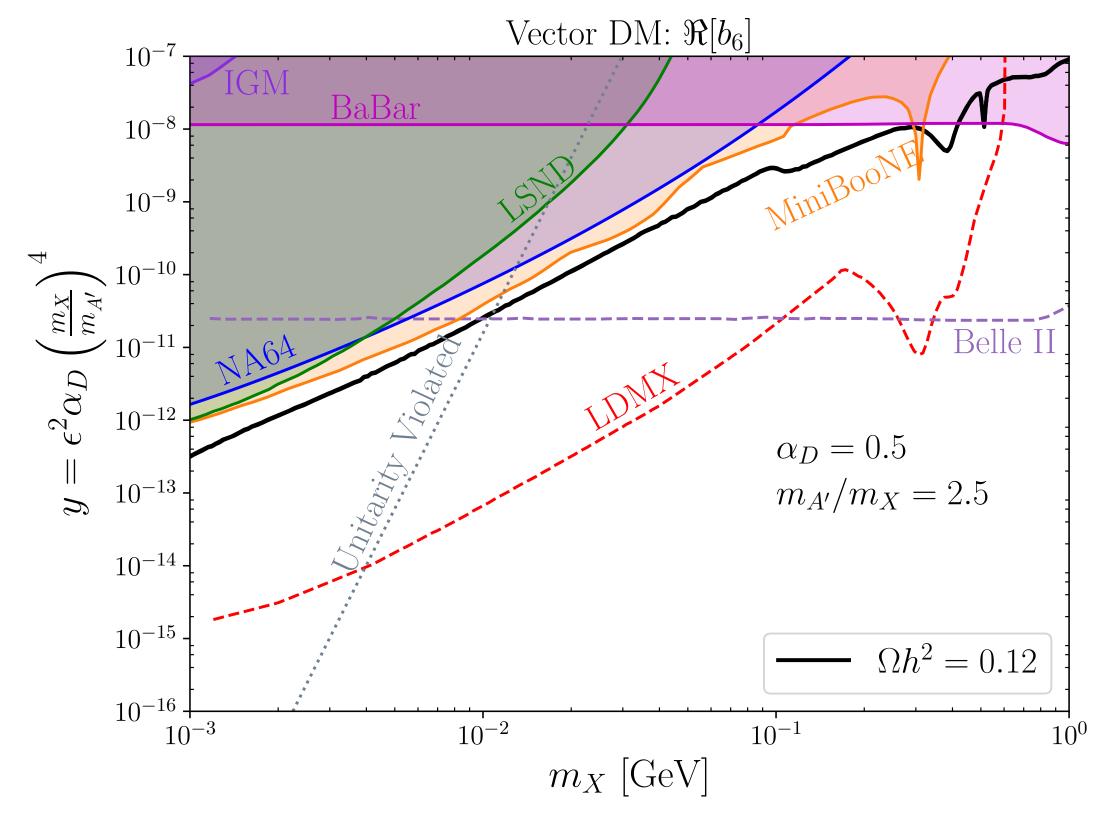
Direct detection via electron recoils and accelerator based experiments work complimentary to one another in probing sub-GeV dark matter.

# **Fixed Target Experiments**

- Accelerator experiments involving a **fixed target** and a **beam** give some of the leading constraints and future sensitivity.
- An electron (or proton) beam is incident on a fixed target leading to interactions between the target nuclei and beam particles, with the goal to produce dark sector particles such as dark photons A'.



- The SM is extended by a single dark matter particle, X and a mediator particle, A'.  $\mathscr{L} = -\left[ib_5 X^{\dagger}_{\nu} \partial_{\mu} X^{\nu} A^{\prime \mu} + b_6 X^{\dagger}_{\mu} \partial^{\mu} X_{\nu} A^{\prime \nu} + h.c.\right]$  $- \left[ b_7 \epsilon_{\mu\nu\rho\sigma} \left( X^{\dagger\mu} \partial^{\nu} X^{\rho} \right) A^{\prime\sigma} + h.c. \right]$  $-e\epsilon A'_{\mu}\bar{f}\gamma^{\mu}f,$
- Certain simplified models are **highly constrained** by **beam dump experiments**, and/or **energy** injection into the CMB, while others are the first to be probed by LDMX.
- Unitarity bounds are strong in these models due to energy dependent vertices and polarization vectors for dark matter.



#### Missing Energy/Momentum

• We consider the future experimental reach of **LDMX** and the current experiment NA64, where the recoil electron  $p_T$  and/or E is measured to search for signals from dark photon bremsstrahlung.

### Proton Beam Dumps

• The absence of significant signal over background at **LSND** and **MiniBooNE** has lead to the constraints considered here. Dark photons can be produced through meson decays, and in the case of MiniBooNE (where the beam energy) is sufficient) through dark bremsstrahlung.

# **Thermal Targets**

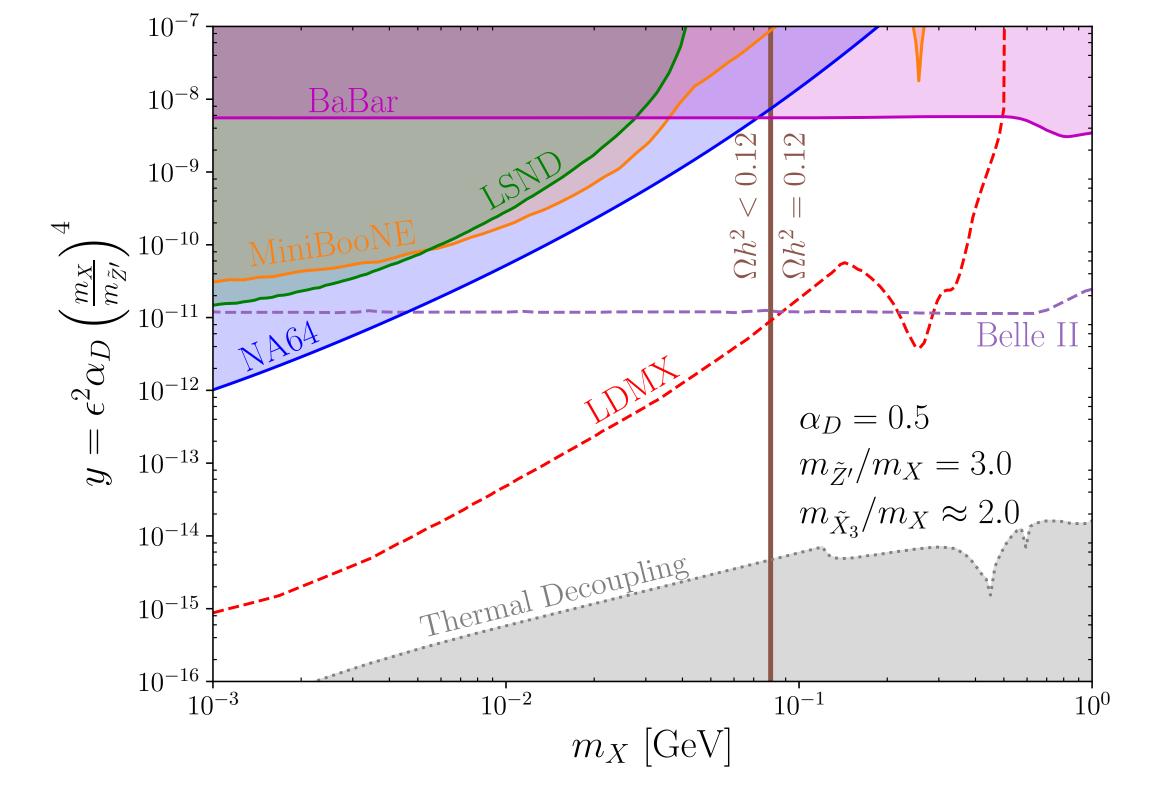
Fixed target experiments search for signatures which are also compatible with **relic** density observations of the CMB from Planck, namely thermal targets. Here we include the relic targets of various dark matter models (including vector/spin-1 dark matter which is from this work) with current and future experimental limits [LDMX and Belle II]:

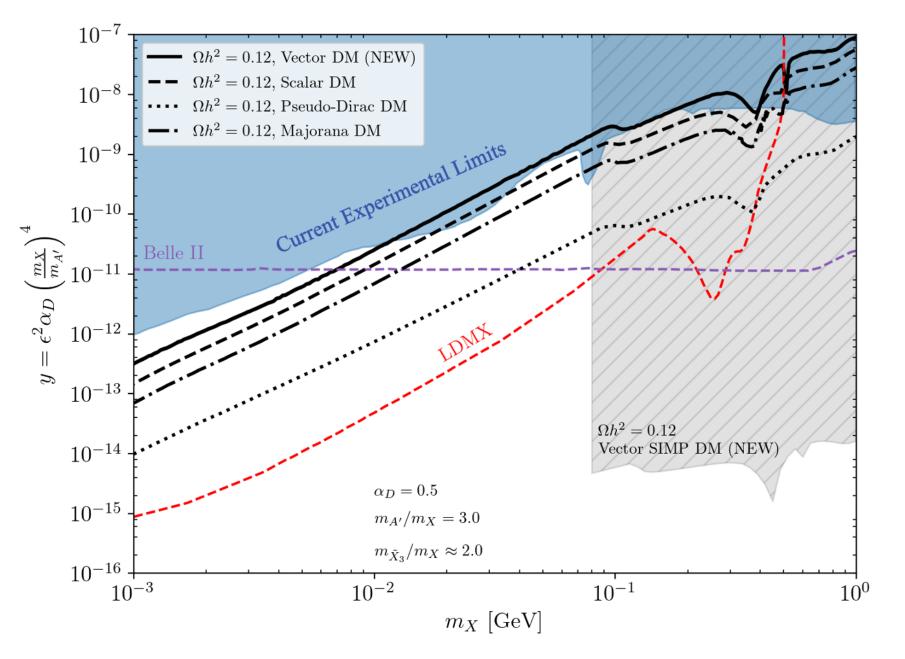
#### SIMP Spin-1 Dark Matter

- The SM is extended by a dark  $SU_X(2) \times U_{Z'}(1)$
- We get an extended Higgs sector, and two new mediators, A' and  $\tilde{X}_3$
- Not subject to strong energy dependent *unitarity bounds*.

$$\begin{aligned} \mathscr{L} &= -ig_X \cos \theta'_X \left[ \left( \partial^{\mu} X^{\nu} - \partial^{\nu} X^{\mu} \right) X^{\dagger}_{\mu} \tilde{X}_{3,\nu} - \left( \partial^{\mu} X^{\nu\dagger} - \partial^{\nu} X^{\mu\dagger} \right) X_{\mu} \tilde{X}_{3,\nu} \right. \\ &\left. + X_{\mu} X^{\dagger}_{\nu} \left( \partial^{\mu} \tilde{X}^{\nu}_{3} - \partial^{\nu} \tilde{X}^{\mu}_{3} \right) \right] \\ &\left. - ig_X \sin \theta'_X \left[ \left( \partial^{\mu} X^{\nu} - \partial^{\nu} X^{\mu} \right) X^{\dagger}_{\mu} \tilde{Z}'_{\nu} - \left( \partial^{\mu} X^{\nu\dagger} - \partial^{\nu} X^{\mu\dagger} \right) X_{\mu} \tilde{Z}'_{\nu} \right. \\ &\left. + X_{\mu} X^{\dagger}_{\nu} \left( \partial^{\mu} \tilde{Z}'^{\nu} - \partial^{\nu} \tilde{Z}'^{\mu} \right) \right] \\ &\left. - e\varepsilon \cos(\theta'_X) \, \tilde{Z}'_{\mu} \, \bar{f} \gamma^{\mu} f + e\varepsilon \sin(\theta'_X) \, \tilde{X}_{3\mu} \, \bar{f} \gamma^{\mu} f \end{aligned}$$

- Relic target is **independent of**  $\epsilon$  since relic density is dominantly set by  $3 \rightarrow 2$  or  $2 \rightarrow 2$ annihilations to  $X_3$ .
- Weaker beam dump limits due to cancellations between diagrams for  $e^{-}/N$  dark matter scattering.





• LDMX Full Luminosity with  $1.6 \times 10^{15}$  EOT and a 8 GeV electron beam.

## Summary

- We have extended the current landscape of sub-GeV DM models considered in the context of **fixed target experiments**,
- to include spin-1 sub-GeV DM [with  $m_{A'} > 2m_{DM}$ ]. Simplified models UV complete SIMP model

Spin-1 dark matter are the first models to be probed at upcoming LDMX!

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