



Topological defects during cosmological phase transitions

Simone Blasi
DESY Hamburg

Based on:

SB, Mariotti [2203.16450], PRL

SB, Jinno, Konstandin, Rubira, Stomberg [2302.06952], JCAP

Agrawal, **SB**, Mariotti, Nee [2312.06749]

SB, Mariotti, [2405.08060]



EuCAPT

Introduction

Higgs mechanism + Hot Big Bang = Cosmological phase transitions

Key to address SM open questions: e.g. matter/antimatter asymmetry, dark matter...

Aftermath of phase transitions **directly observable in gravitational waves**

QCD and EWPT are not first order in the SM: need for **new particles** or **new symmetries**

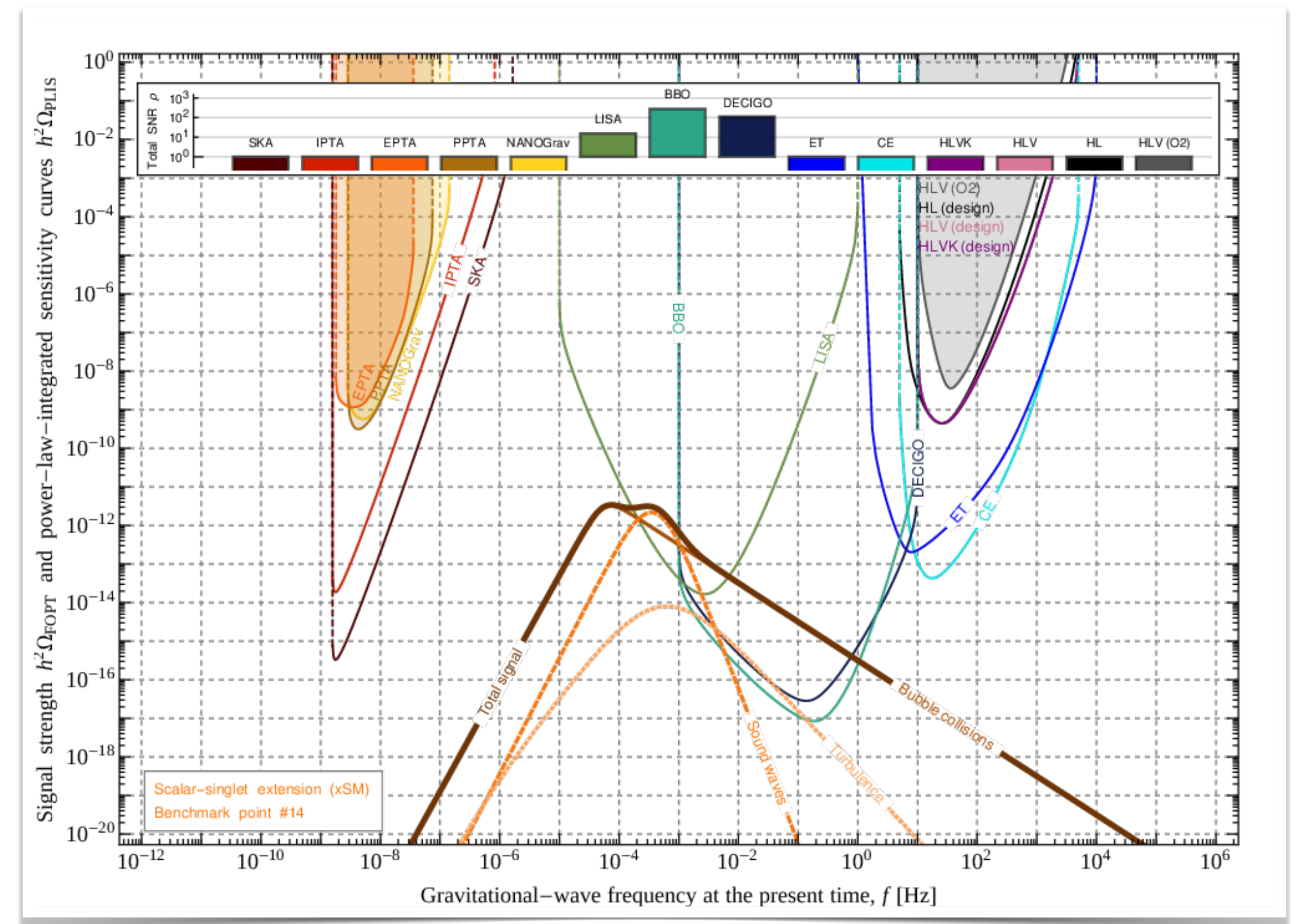
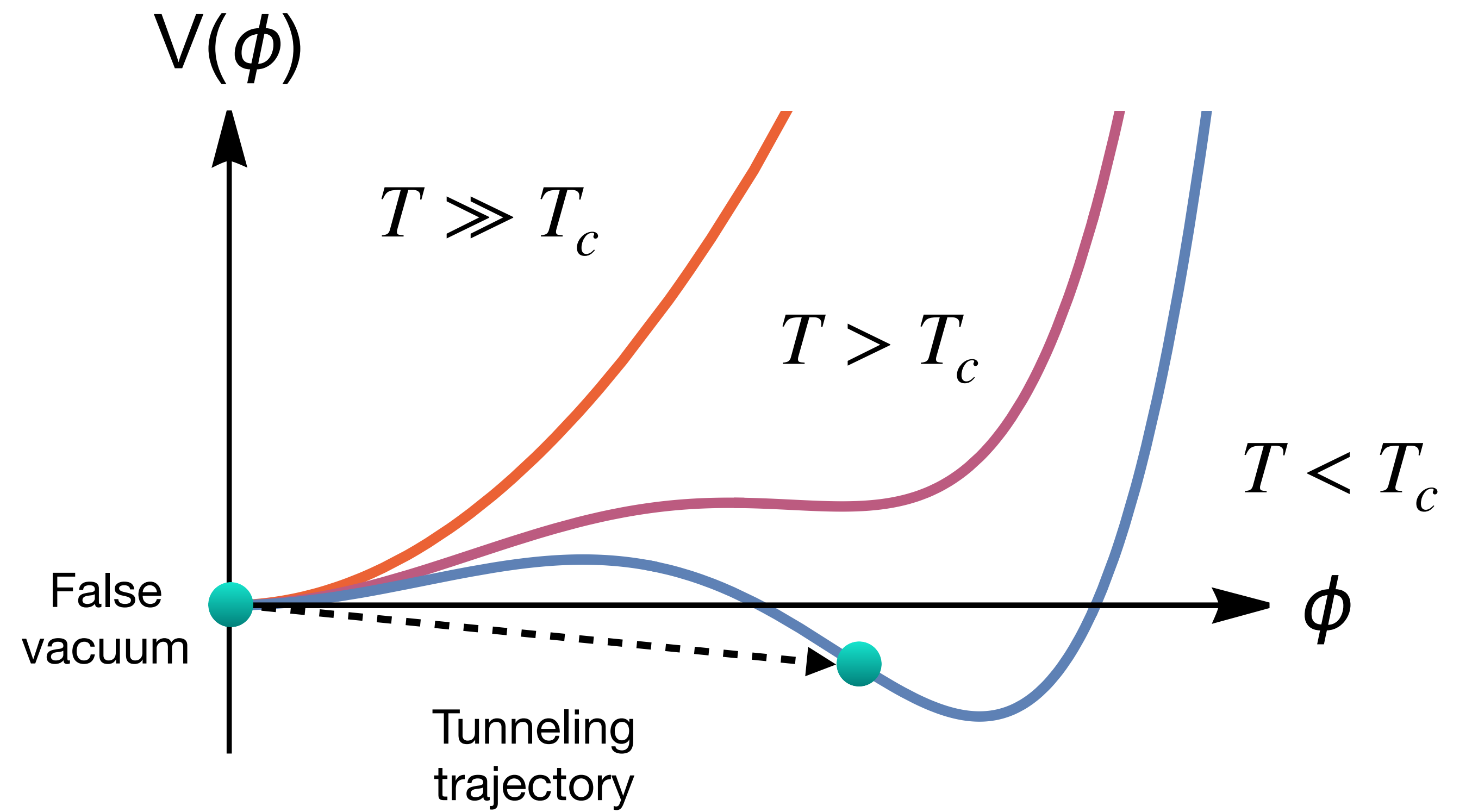


Fig. from Schmitz [2002.04615] JHEP

Nucleation theory

Coleman 1977 (PRD)
Callan, Coleman 1977 (PRD)
Linde 1983 (NPB)



Nucleation theory

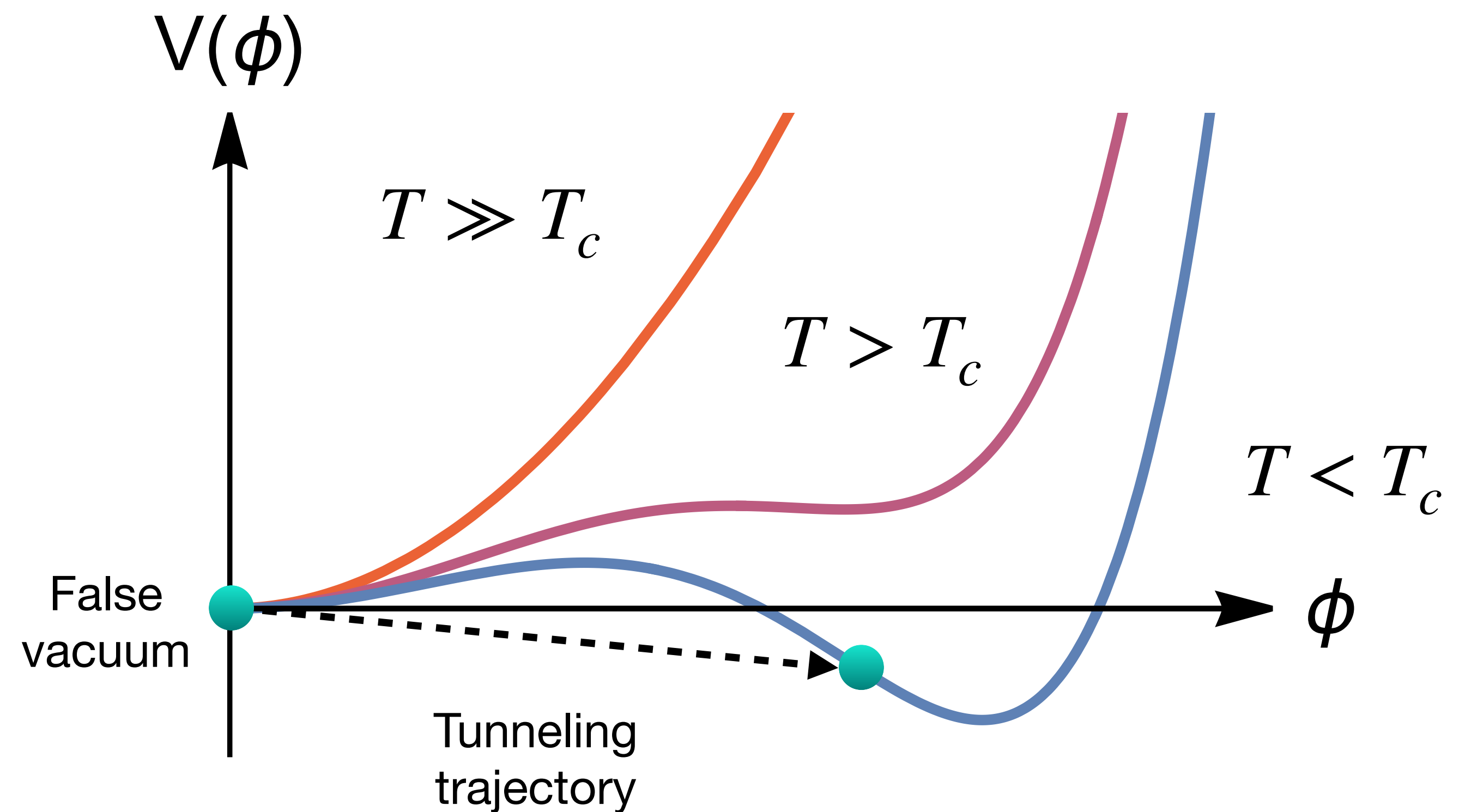
- Assume thermal fluctuations in homogeneous spacetime:

$$\phi(\mathbf{x}, \tau) = \phi(r), \quad r = |\mathbf{x}|$$

- Tunneling rate per unit volume given by O(3) action S_3/T

$$\gamma_V \sim T^4 \exp(-S_3/T)$$

Coleman 1977 (PRD)
 Callan, Coleman 1977 (PRD)
 Linde 1983 (NPB)



Nucleation theory

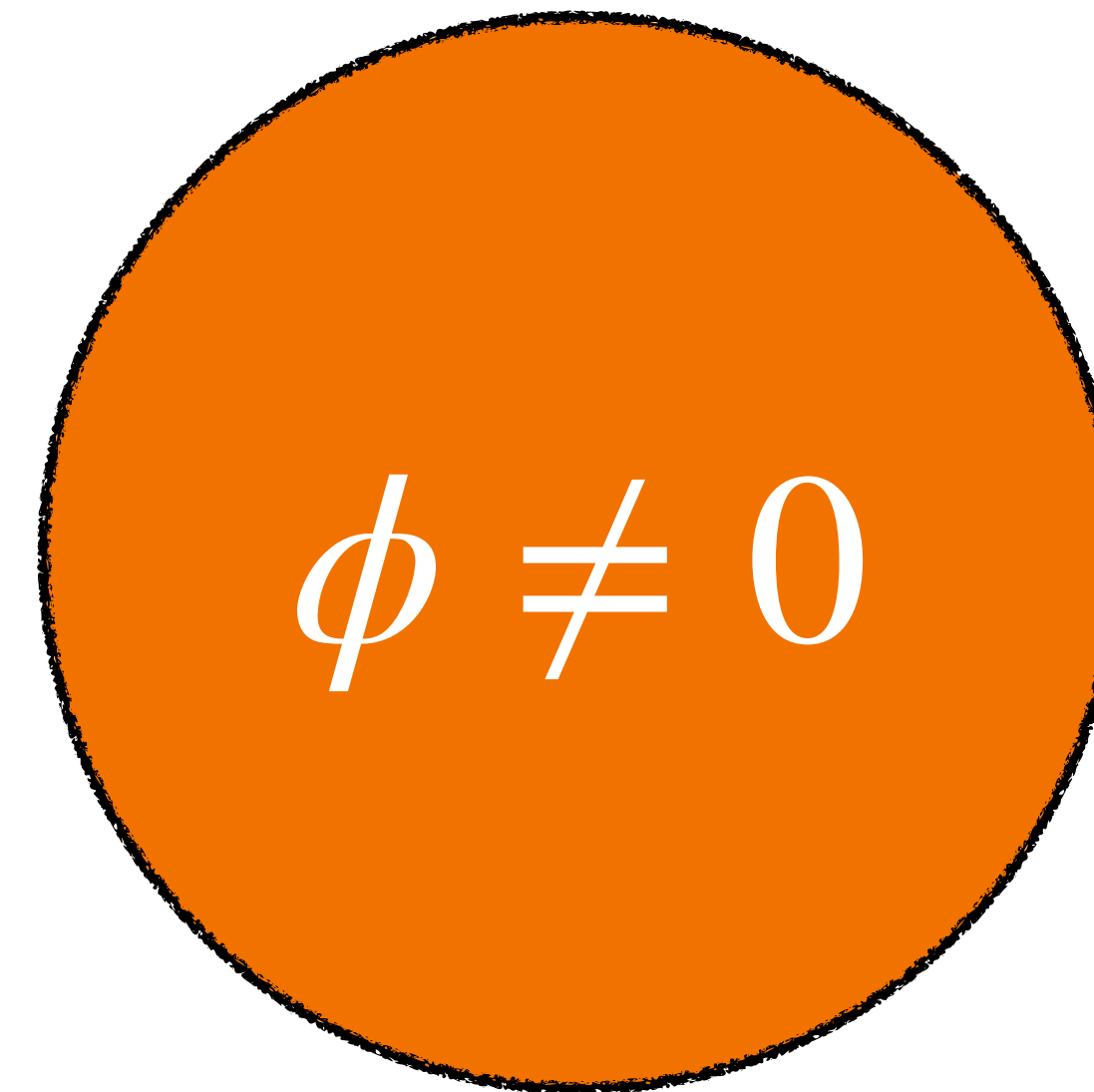
Coleman 1977 (PRD)
Callan, Coleman 1977 (PRD)
Linde 1983 (NPB)

- Assume thermal fluctuations in homogeneous spacetime:

$$\phi(\mathbf{x}, \tau) = \phi(r), \quad r = |\mathbf{x}|$$

- Tunneling rate per unit volume given by O(3) action S_3/T

$$\gamma_V \sim T^4 \exp(-S_3/T)$$



$$\phi = 0$$

What about impurities?

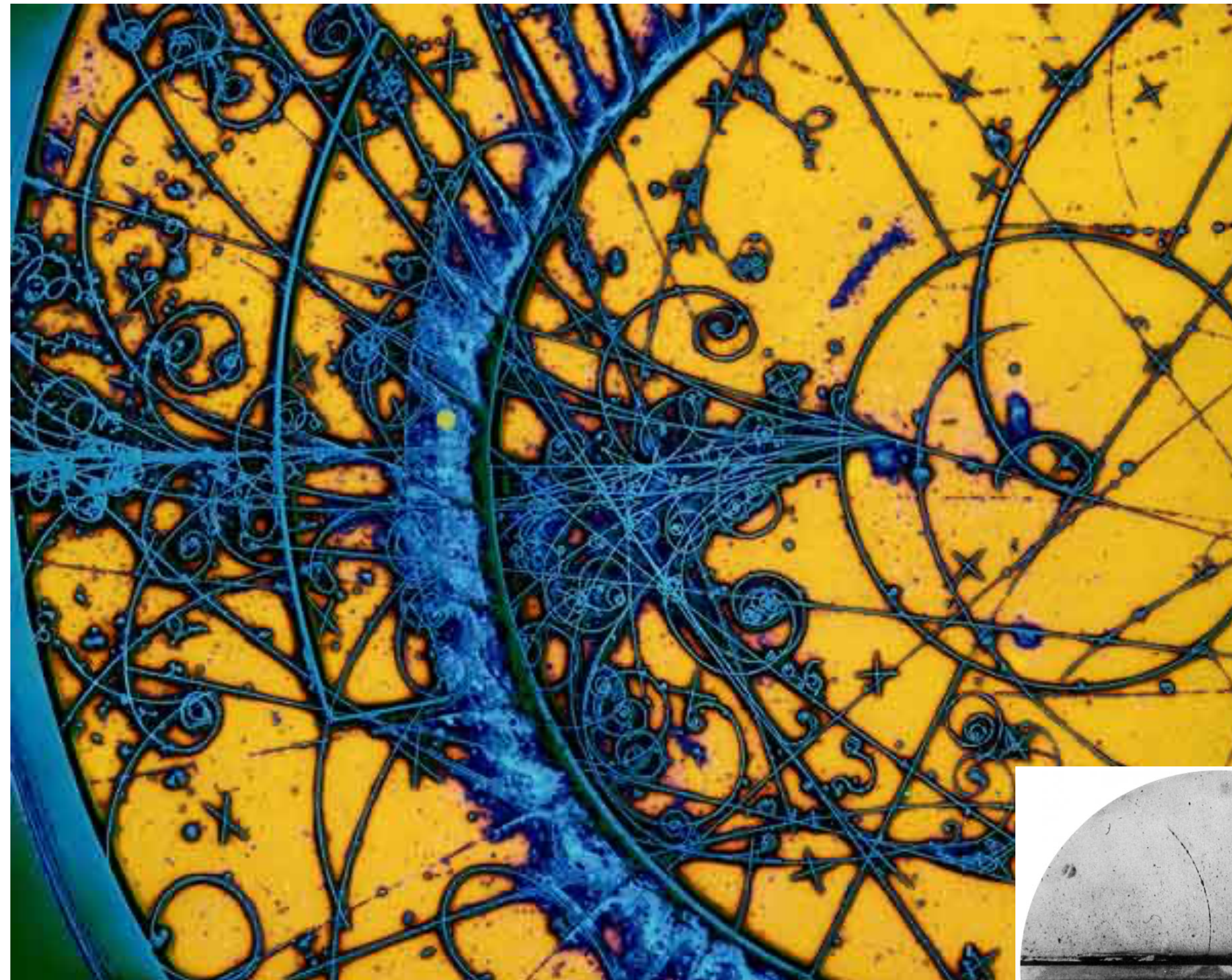
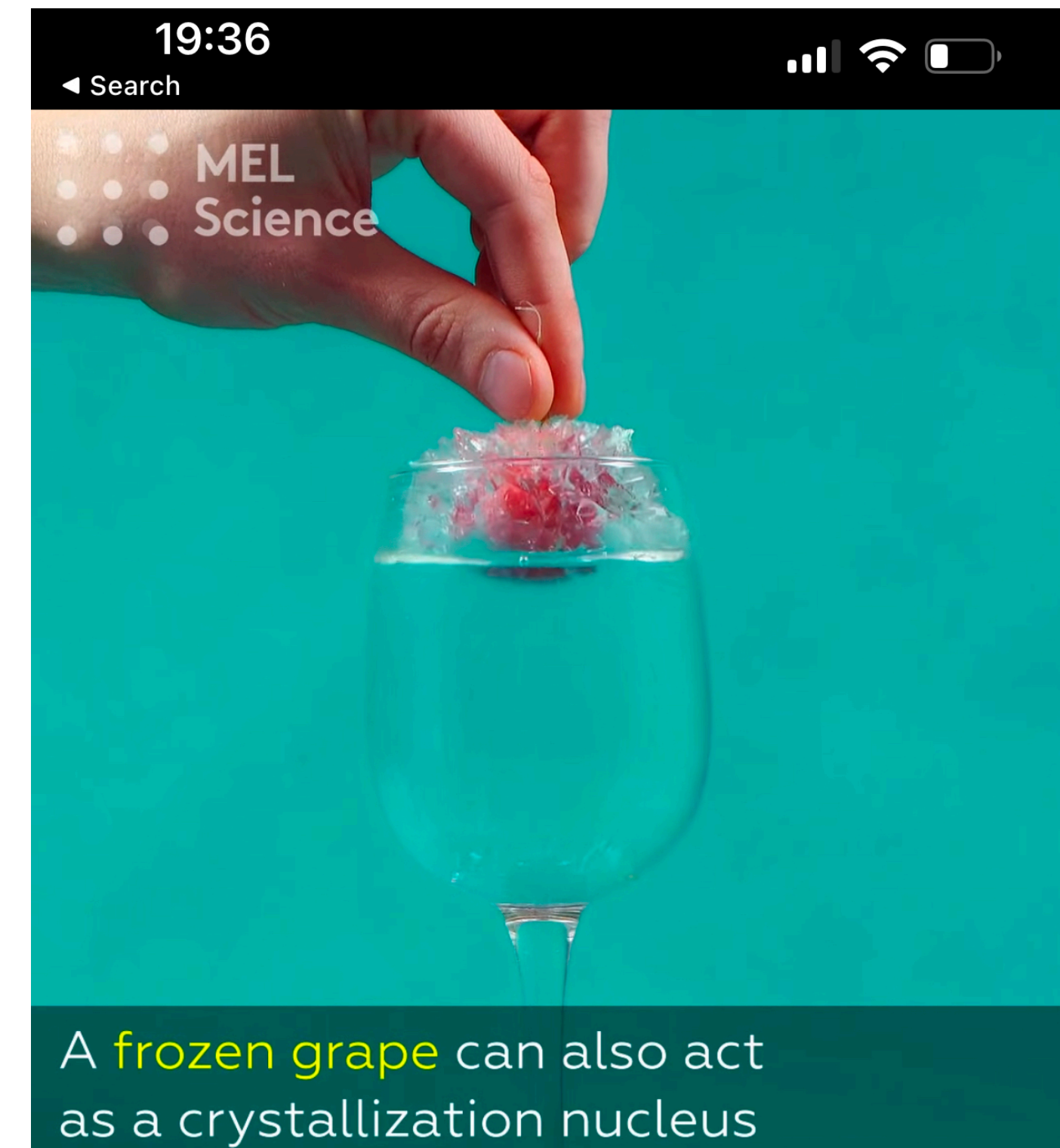
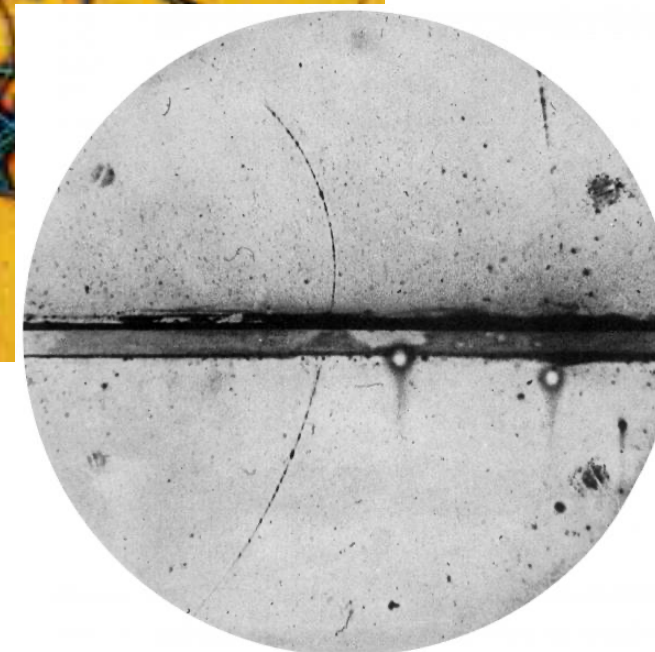


Figure: Bubble chamber



A supercool experiment

82K views 3 yr ago ...more

MEL Chemistry 13.1K

Subscribe

1.1K



Share

Download

Save

**MONOPOLE AND VORTEX DISSOCIATION AND DECAY OF THE
FALSE VACUUM**

Paul Joseph STEINHARDT

Lyman Laboratory of Physics, Harvard University, Cambridge, Massachusetts 02138, USA

Received 17 February 1981

“If monopole (or vortex) solutions exist for a metastable or false vacuum, **a finite density of monopoles (or vortices) can act as impurity sites that trigger inhomogeneous nucleation** and decay of the false vacuum.”

See also:

“Impurities in the early Universe”, Hosotani (1982)

“Cosmic separation of Phases”, Witten (1984)

The nature of impurities

- Compact objects and gravitational effects

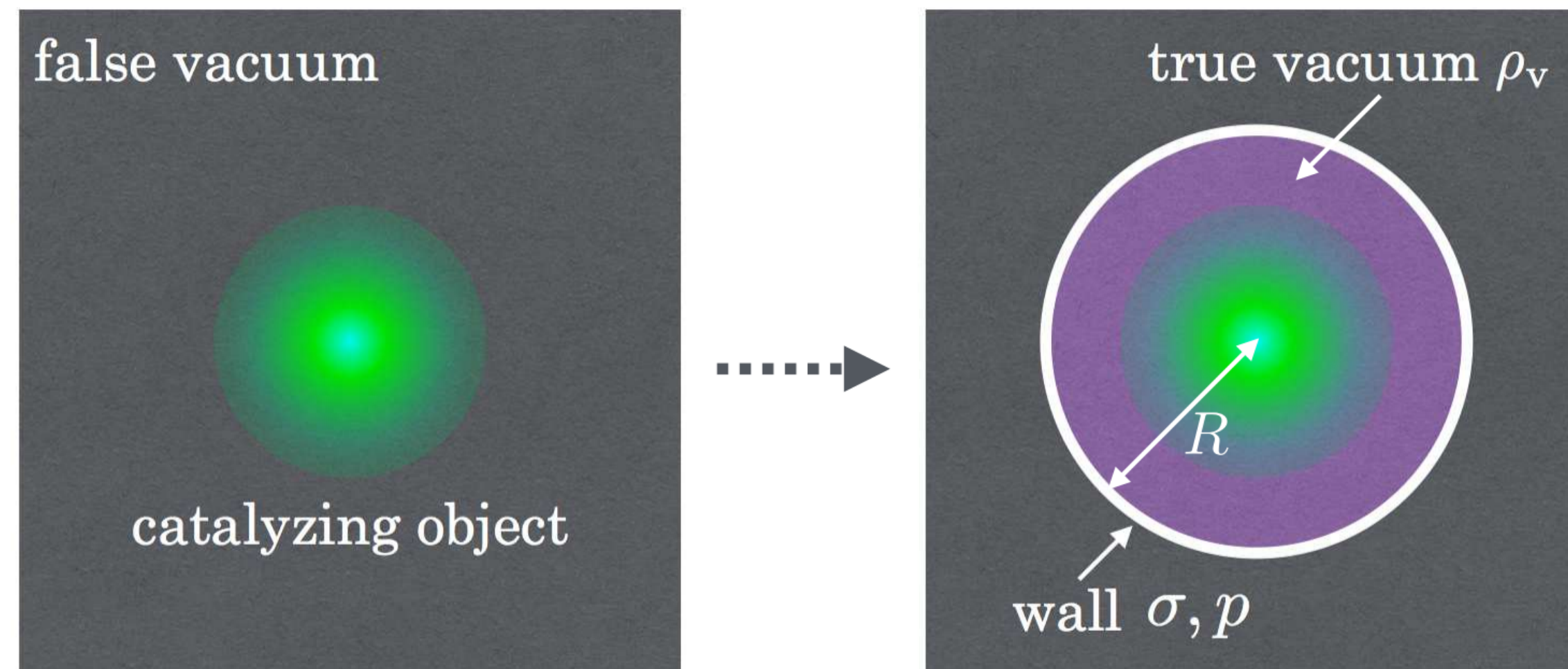


Fig. from Oshita, Yamada, Yamaguchi [1808.01382], PLB

- Primordial density fluctuations

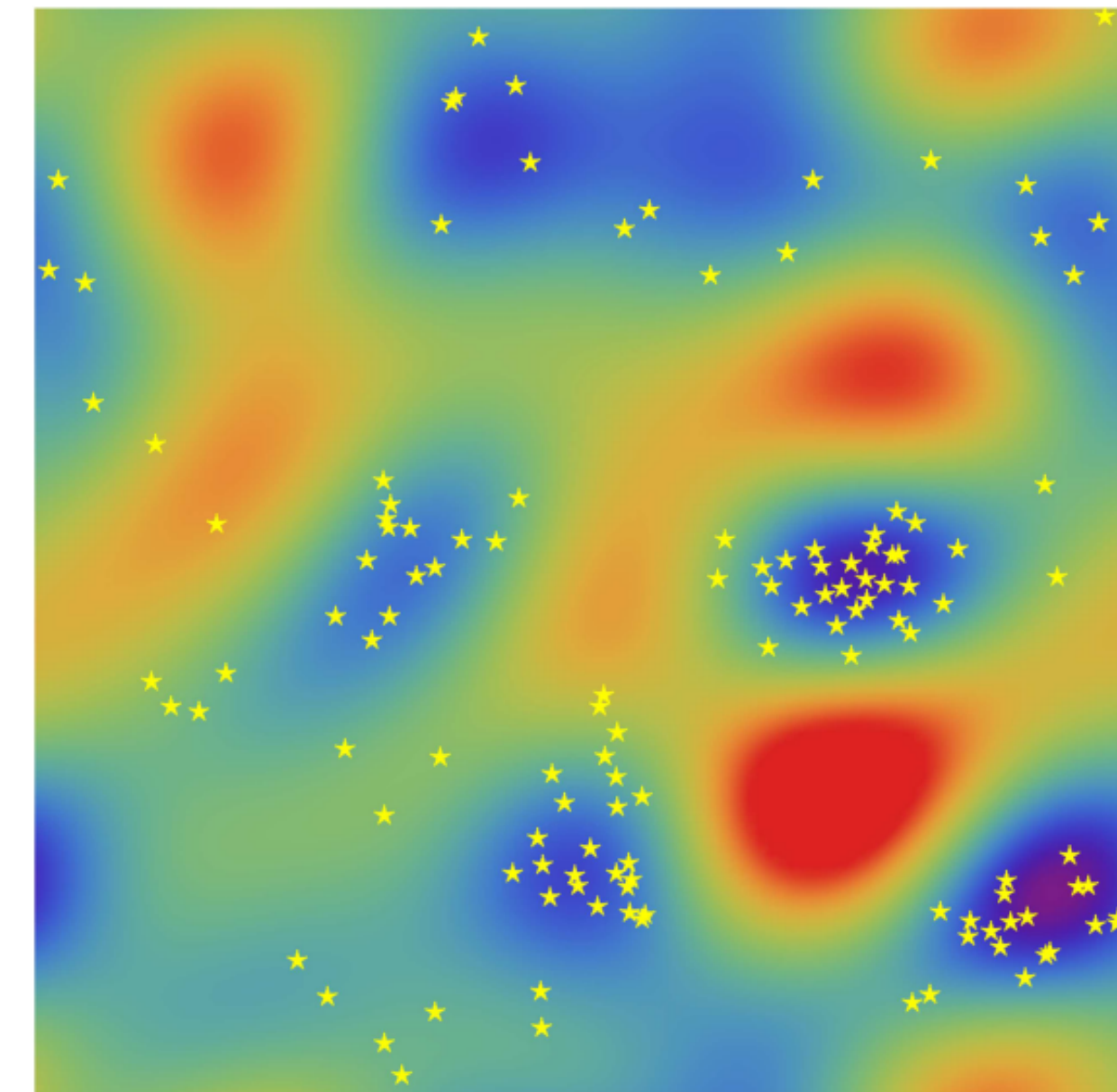


Fig. from Jinno, Konstandin, Rubira, van de Vis, [2108.11947], JCAP

The nature of impurities

- Topological defects (this talk)

Domain walls

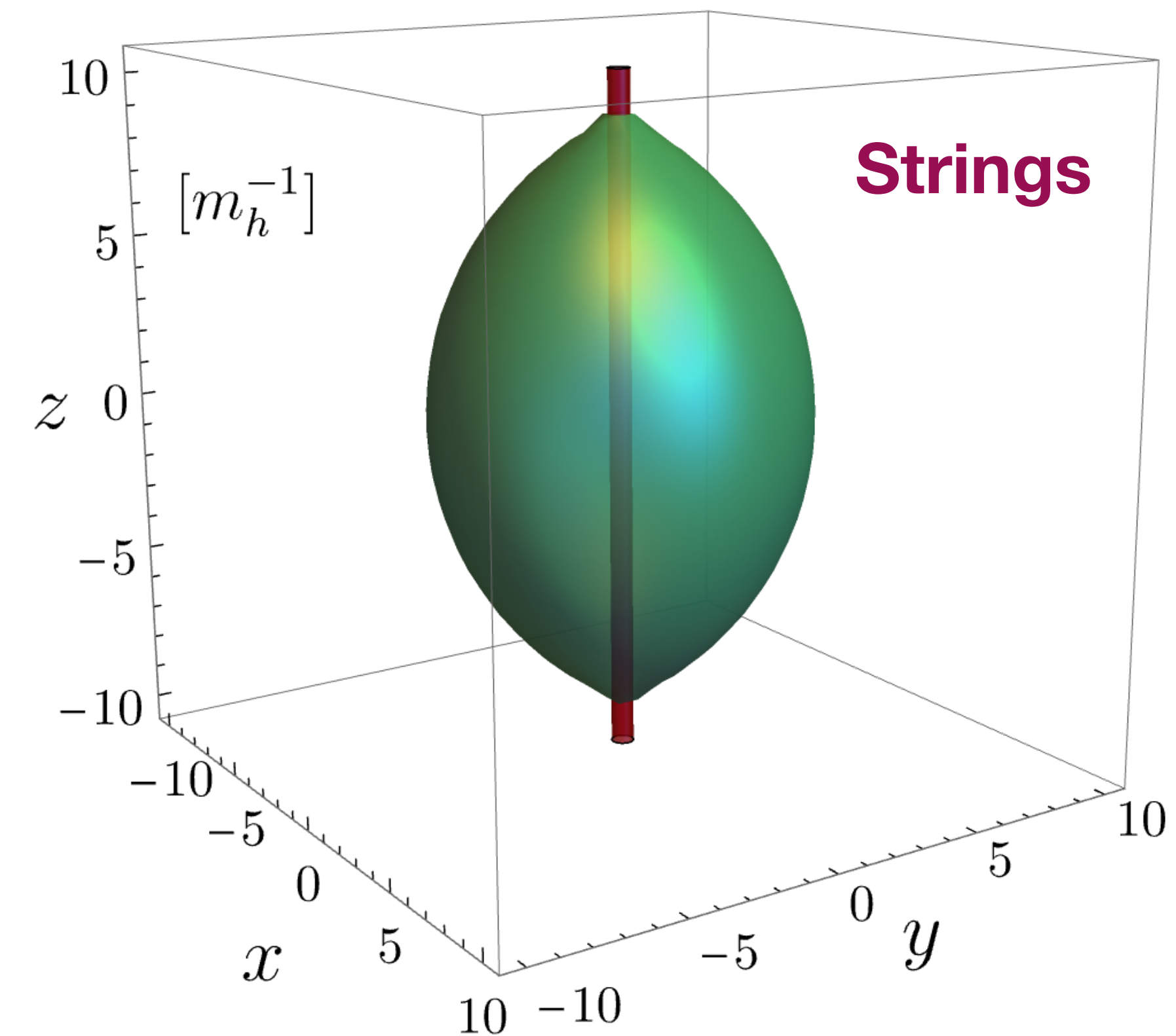
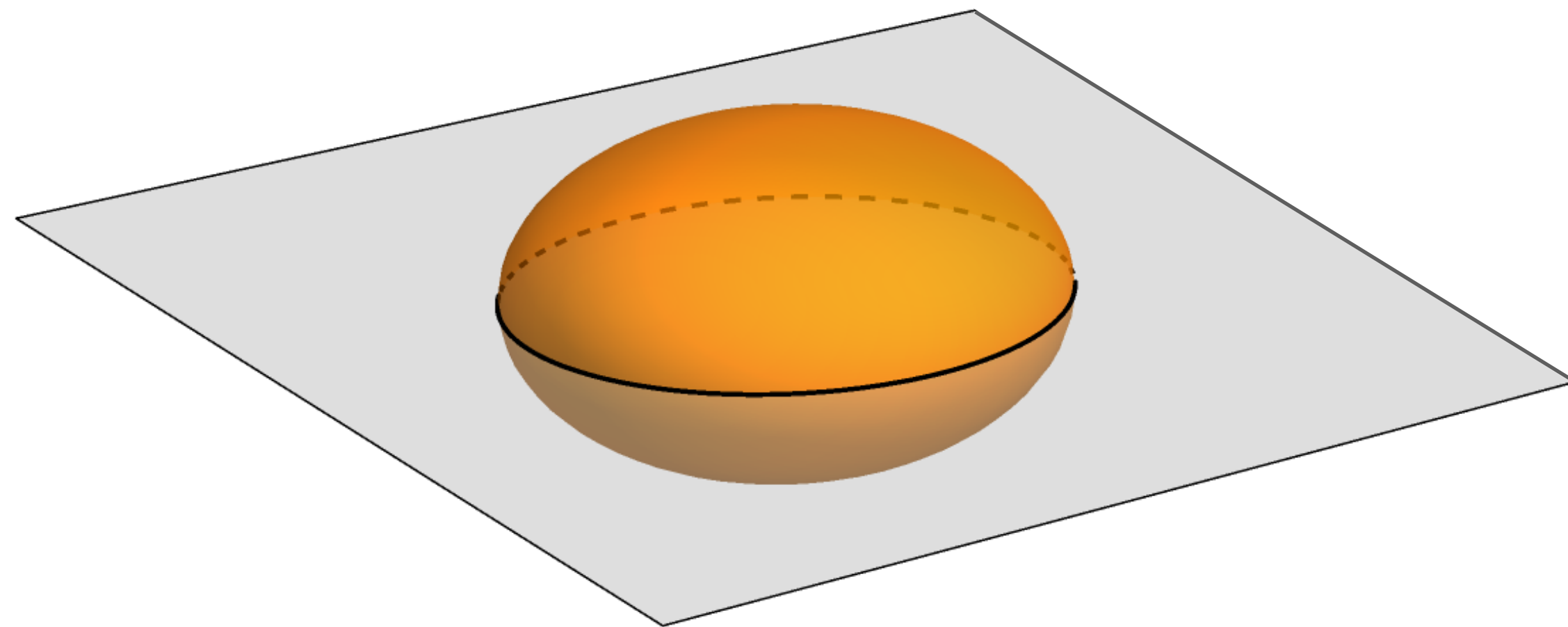


Fig. From Agrawal, **SB**, Mariotti, Nee [2312.06749]

Fig. From **SB**, Mariotti, [2405.08060]

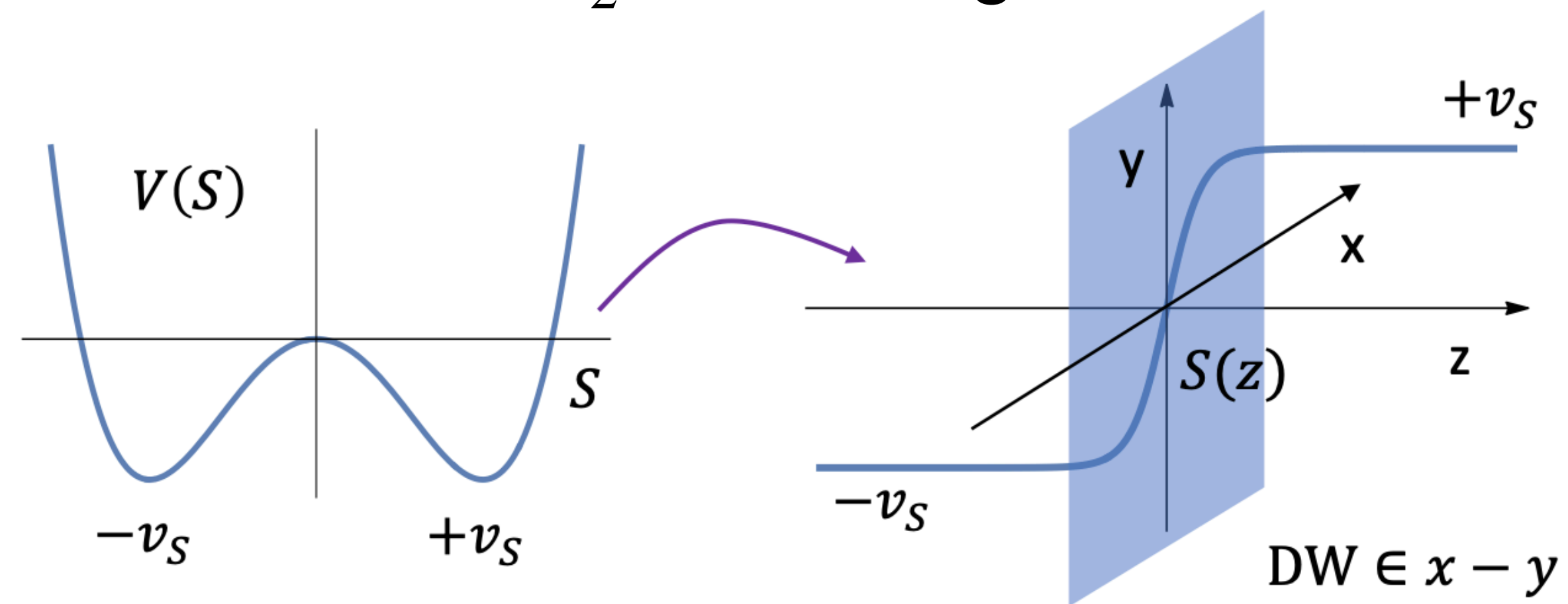
Topological classification

Defect	Dimension	Homotopy	Mass
Domain walls	2	$\pi_0(\mathcal{M})$	σL^2
Strings	1	$\pi_1(\mathcal{M})$	μL

Topological classification

Defect	Dimension	Homotopy	Mass
Domain walls	2	$\pi_0(\mathcal{M})$	σL^2
Strings	1	$\pi_1(\mathcal{M})$	μL

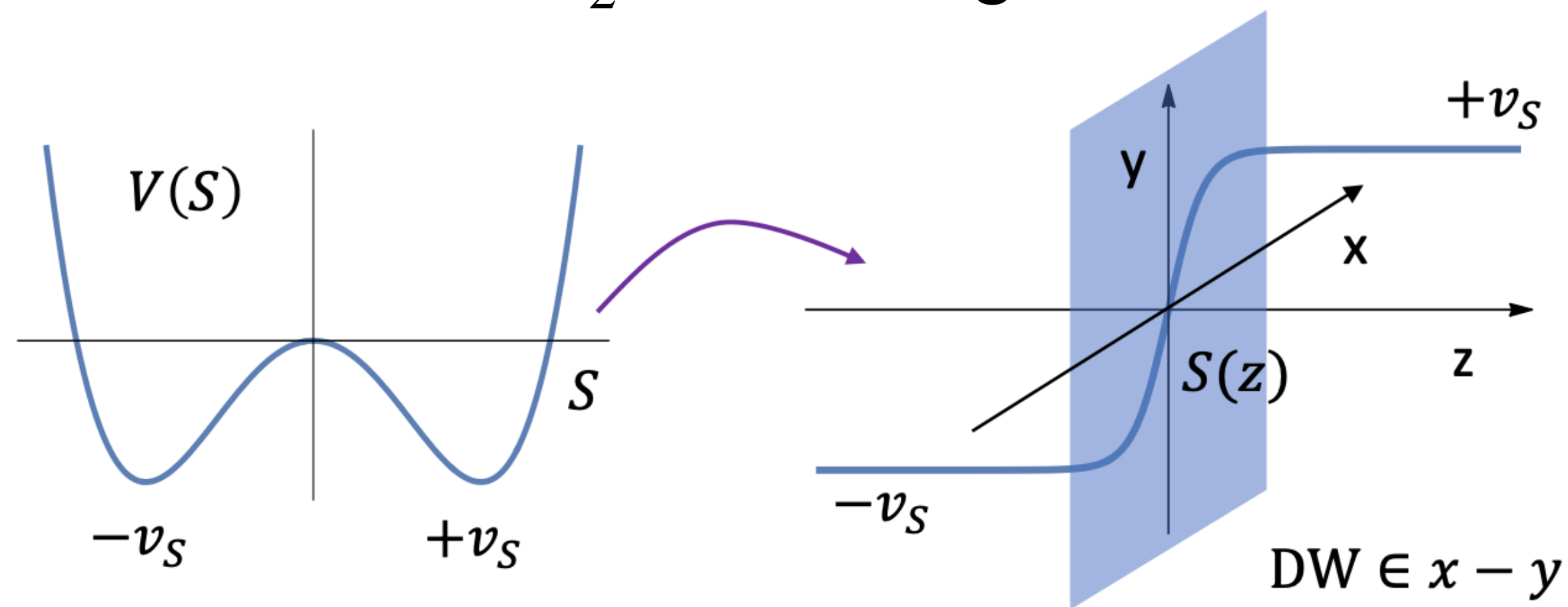
$\mathbb{Z}_2 \rightarrow$ nothing



Topological classification

Defect	Dimension	Homotopy	Mass
Domain walls	2	$\pi_0(\mathcal{M})$	σL^2
Strings	1	$\pi_1(\mathcal{M})$	μL

$\mathbb{Z}_2 \rightarrow$ nothing



$U(1) \rightarrow$ nothing

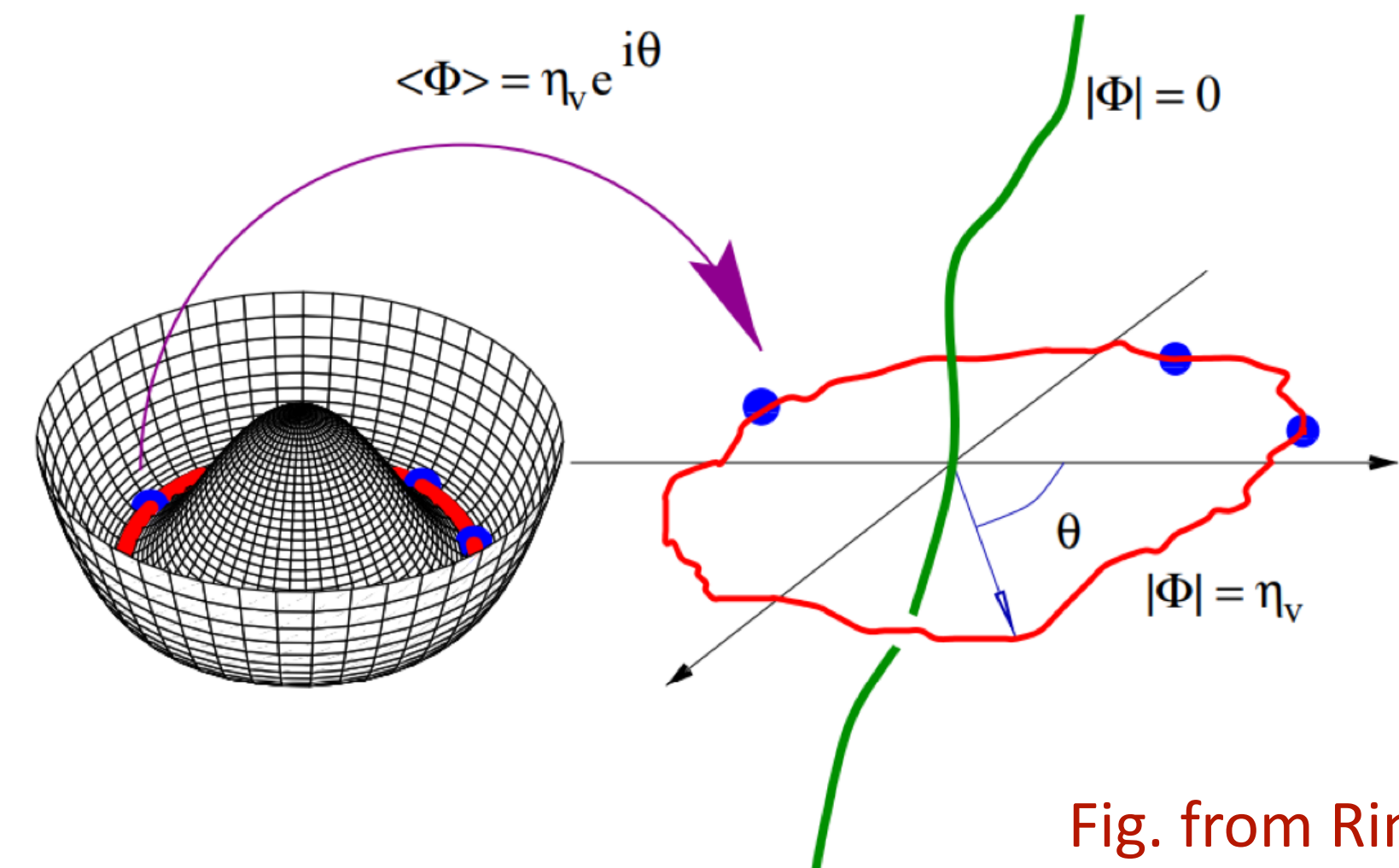
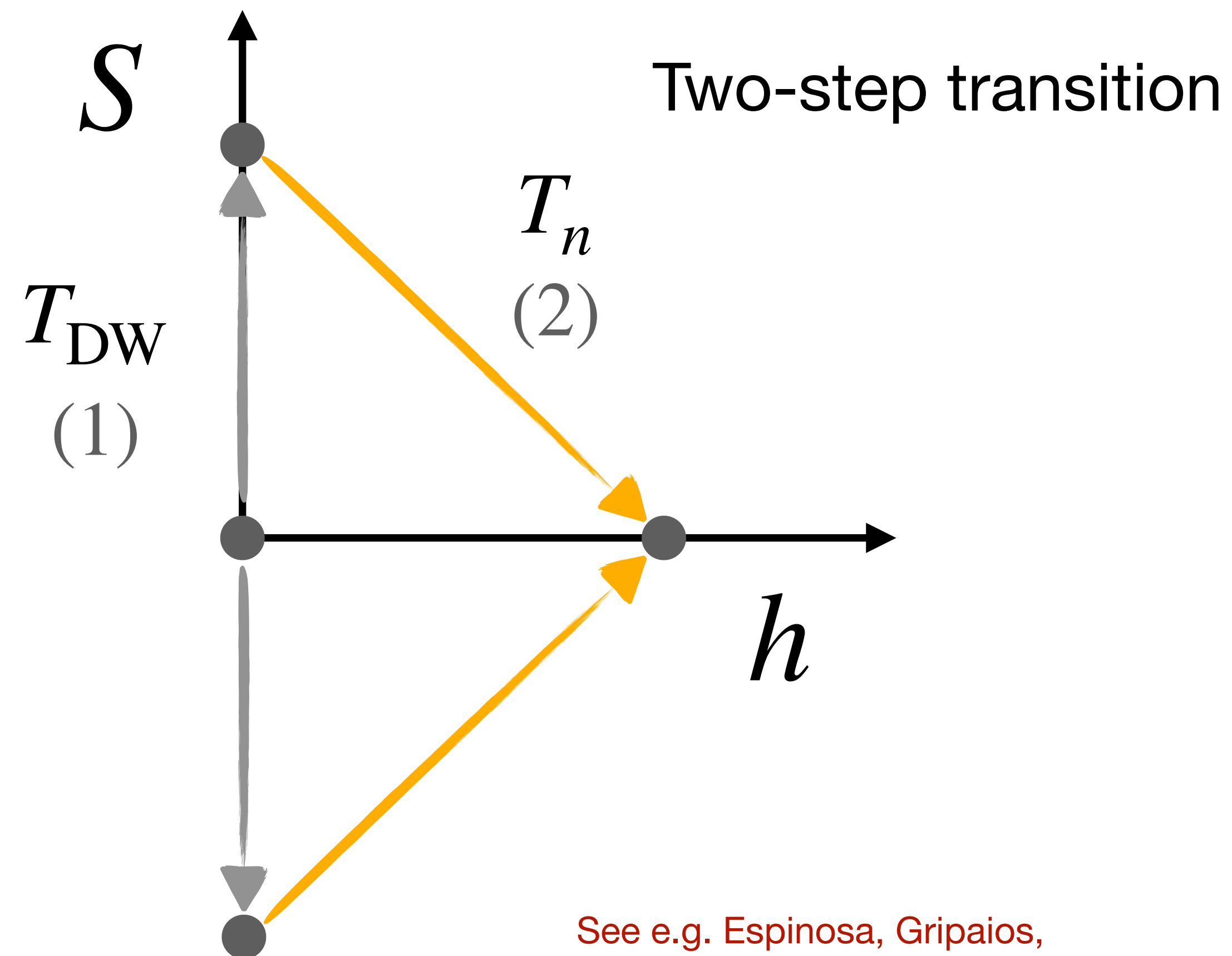


Fig. from Ringeval 2010

Electroweak phase transition with a singlet

- SM + scalar singlet with $\mathbb{Z}_2 : S \rightarrow -S$

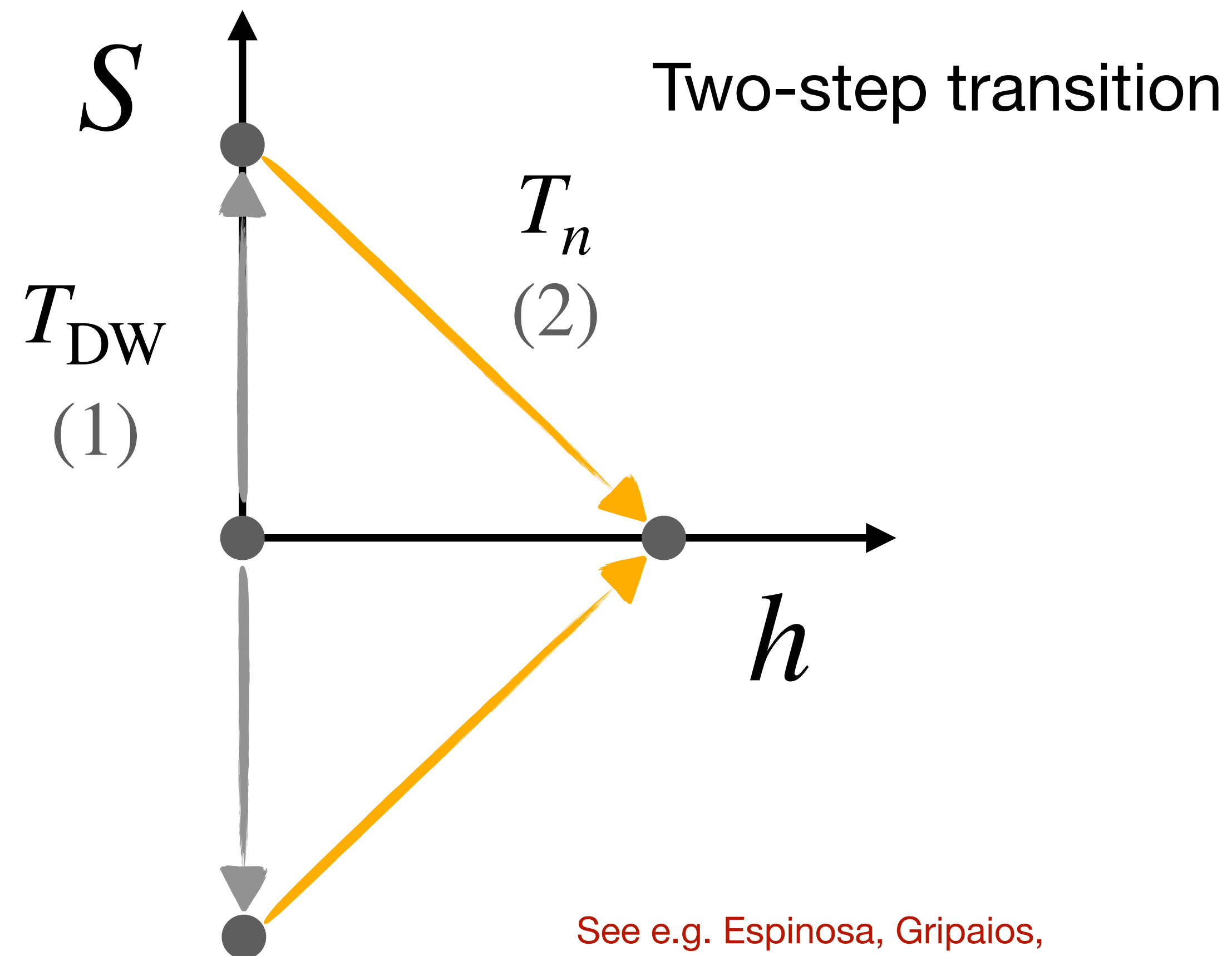


See e.g. Espinosa, Gripaios,
Konstandin, Riva [1110.2876]
JCAP

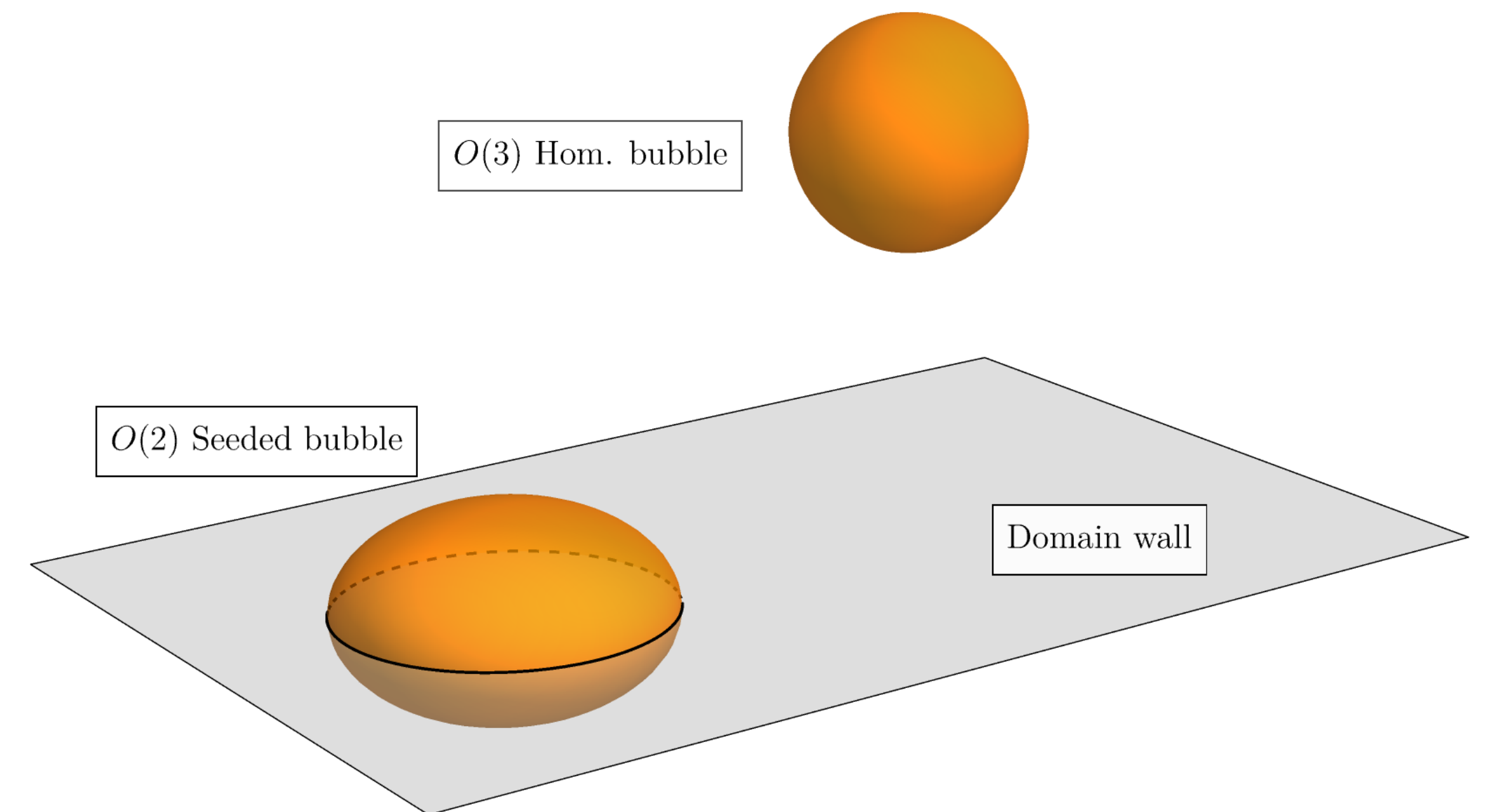
Electroweak phase transition with a singlet

- SM + scalar singlet with $\mathbb{Z}_2 : S \rightarrow -S$

- Competition between homogenous and seeded nucleation for 2nd step:



See e.g. Espinosa, Gripaios, Konstandin, Riva [1110.2876] JCAP



SB, Mariotti [2203.16450], PRL
Agrawal, **SB**, Mariotti, Nee [2312.06749]

Electroweak phase transition with a singlet

- Seeded tunneling is faster whenever there is a two—step transition
- New parameter space becomes viable thanks to the walls
- Phenomenology of the phase transition is drastically changed

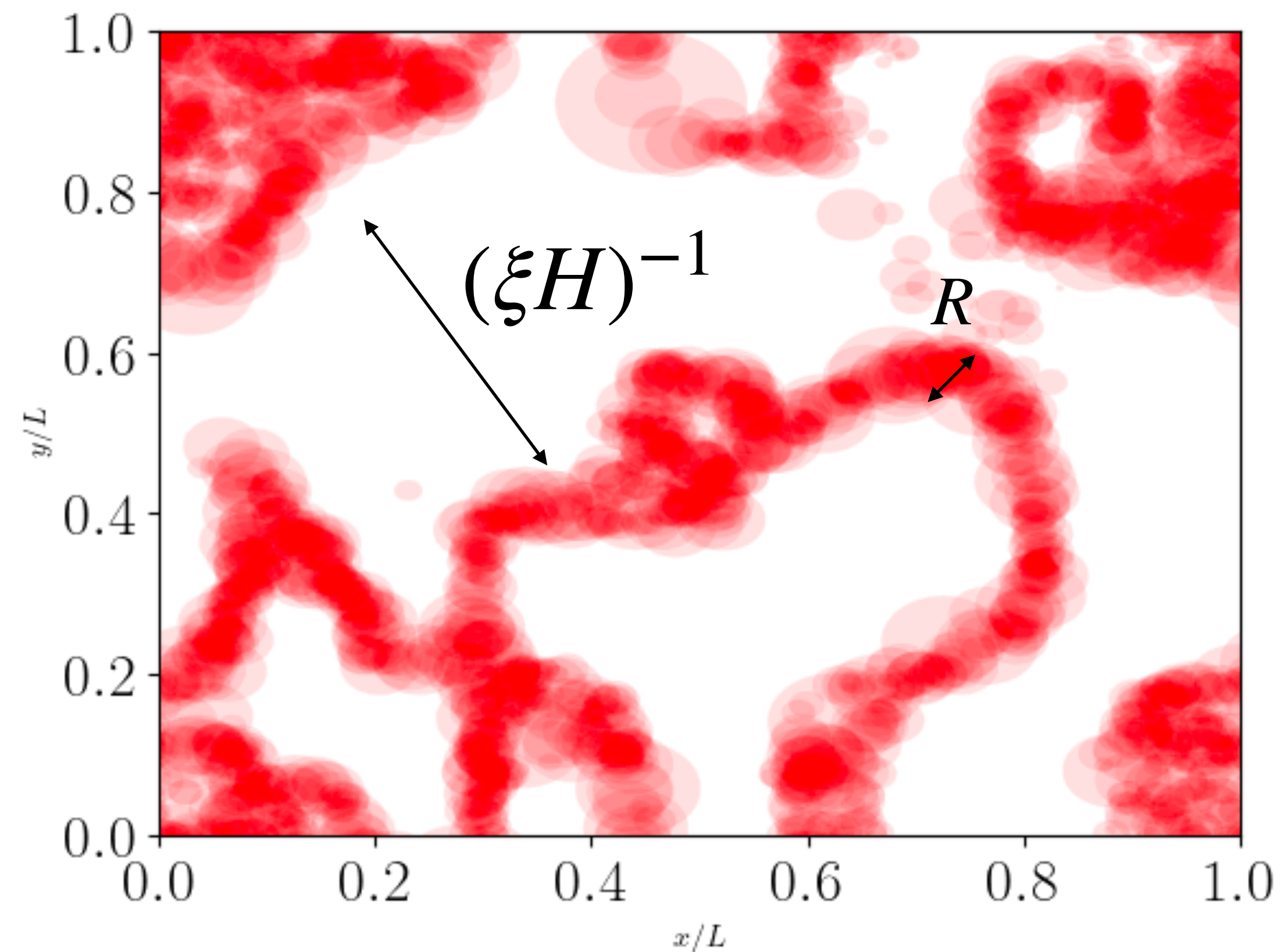
Electroweak phase transition with a singlet

- Seeded tunneling is faster whenever there is a two—step transition
- New parameter space becomes viable thanks to the walls
- Phenomenology of the phase transition is drastically changed

- Previous studies can still apply if explicit Z_2 breaking is introduced, implications need to be taken into account consistently

Gravitational waves from seeded bubbles

- Domain wall network mimicked by Ising model



- Spectrum shifted to IR with enhanced amplitude

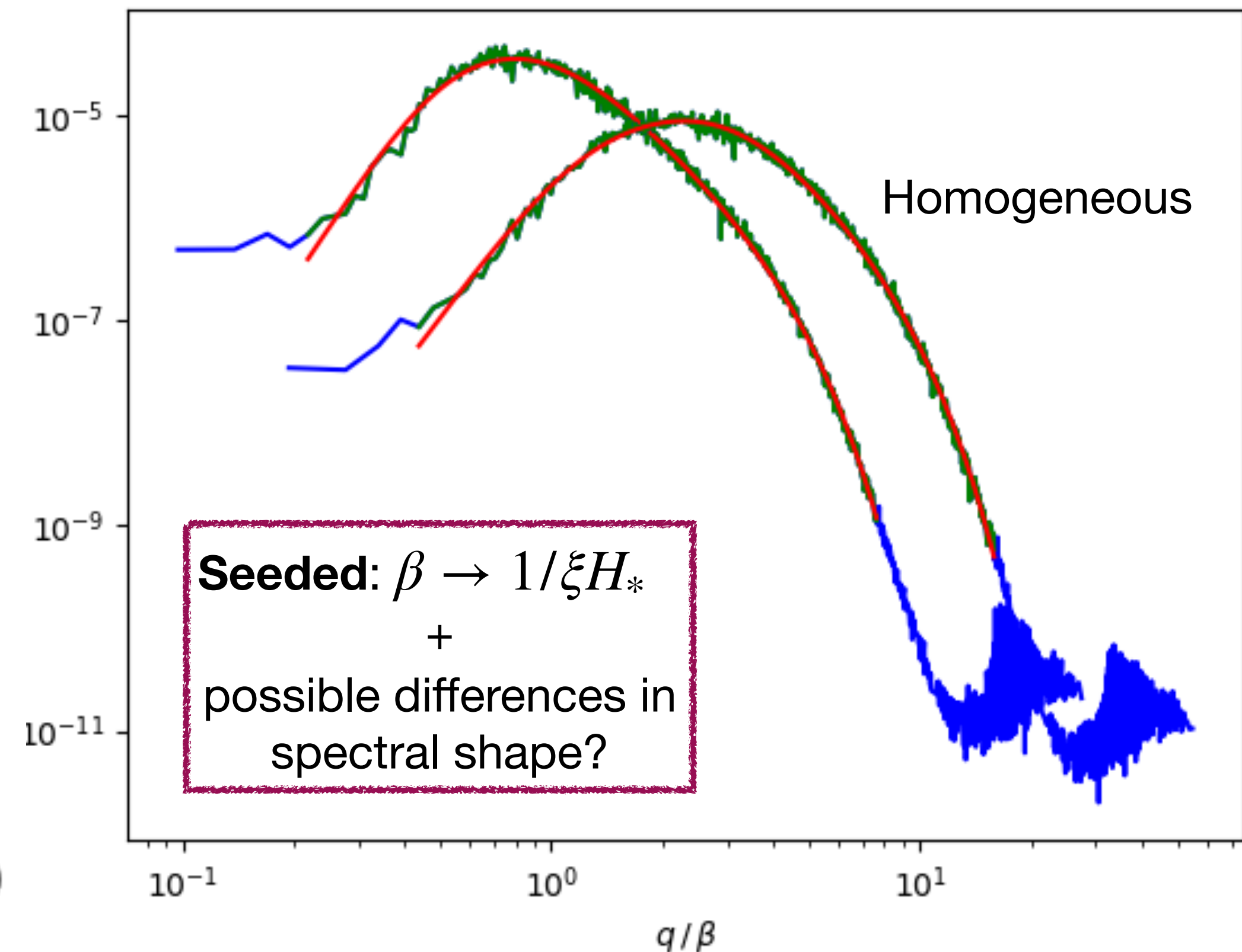
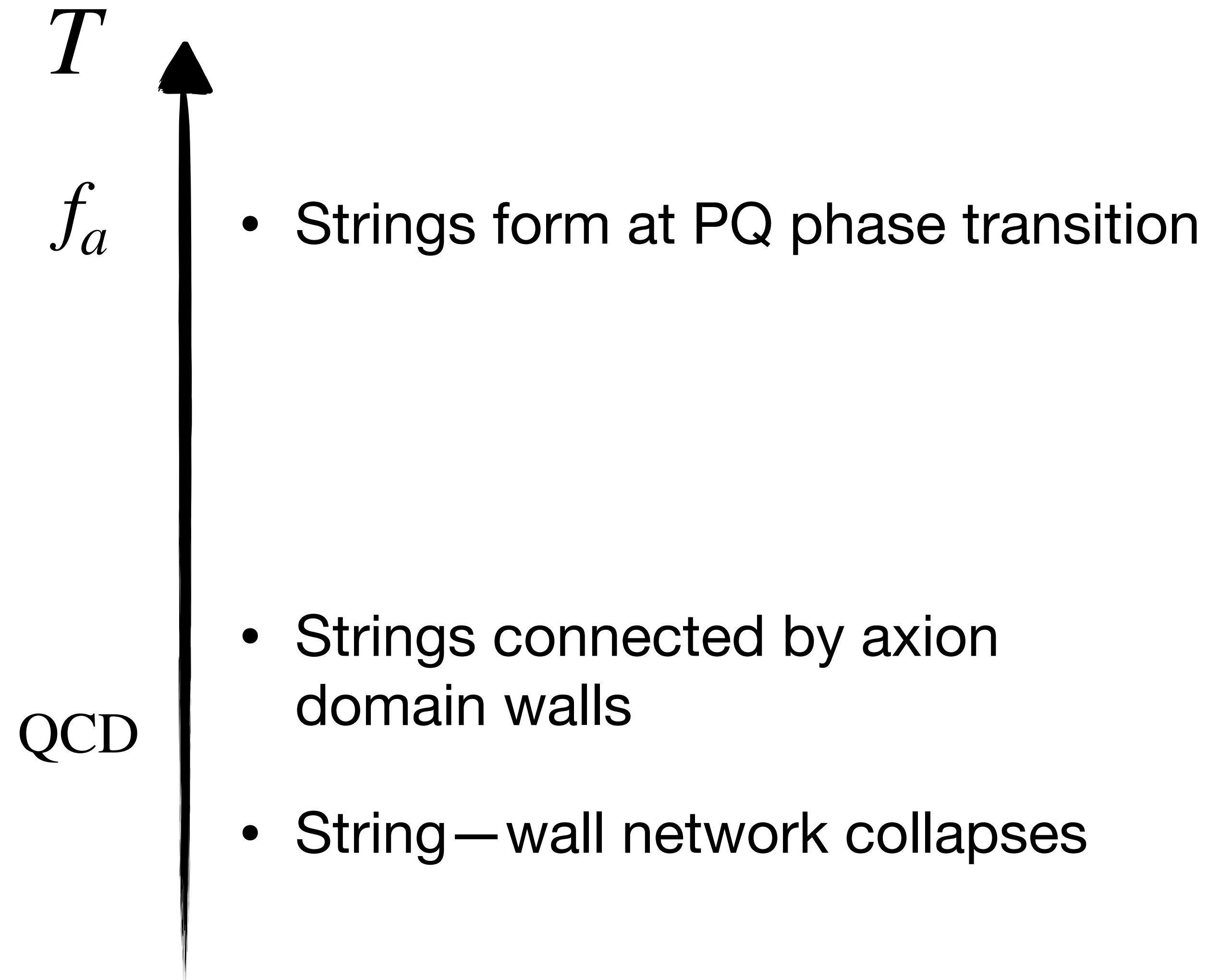


Fig. from **SB**, Jinno, Konstandin, Rubira, Stomberg [2302.06952] JCAP

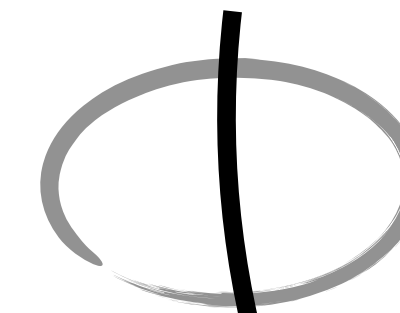
What about other defects?

SB, Mariotti [2405.08060]

QCD axion strings

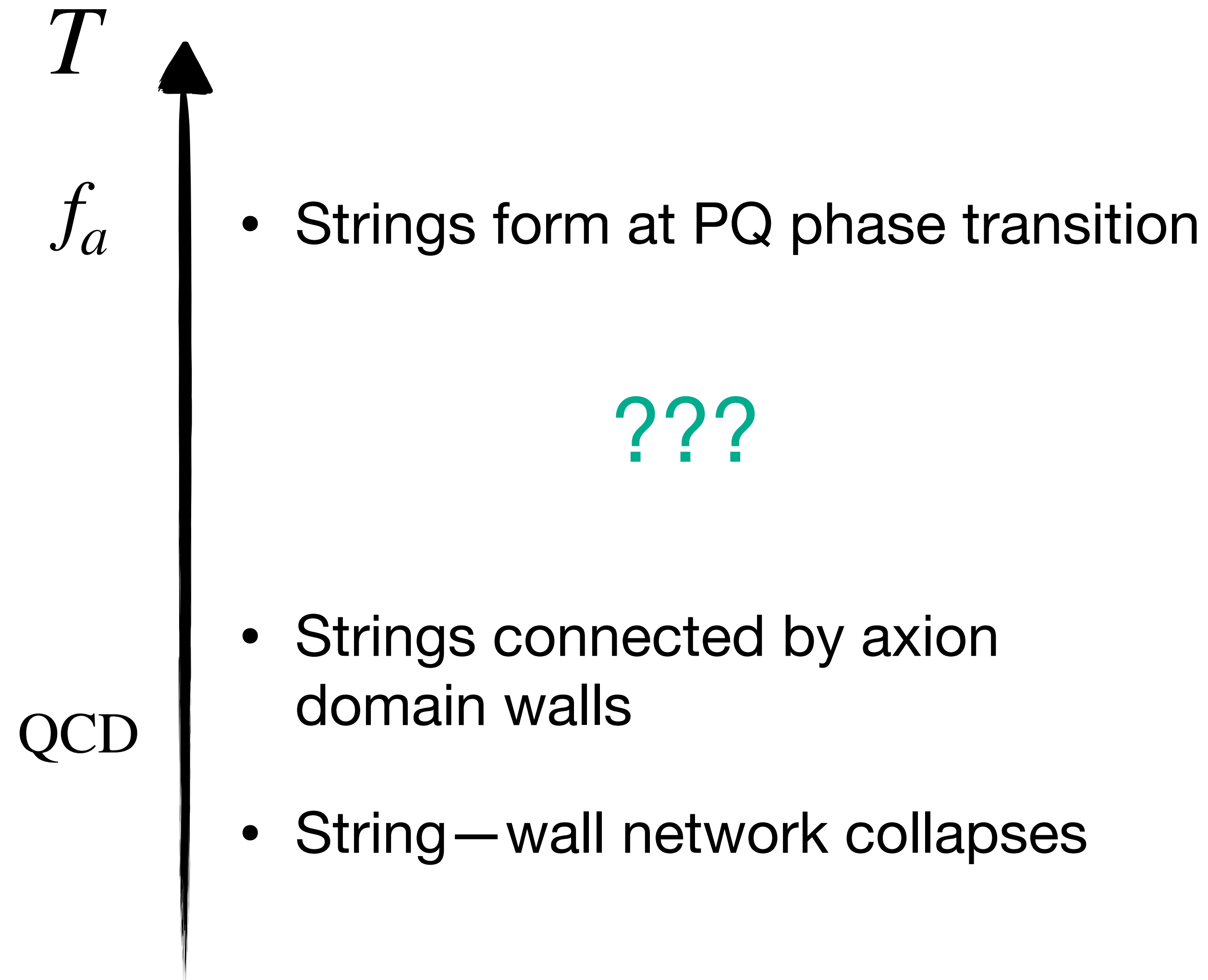


$$\alpha(\theta) : 0 \rightarrow 2\pi$$

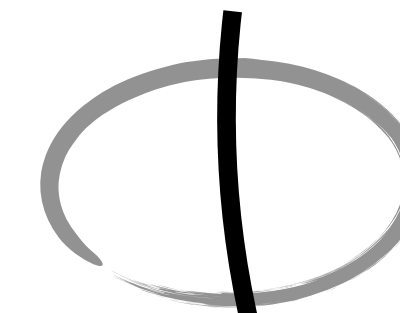


$$\delta \approx m_\rho^{-1}$$

QCD axion strings



$$\alpha(\theta) : 0 \rightarrow 2\pi$$

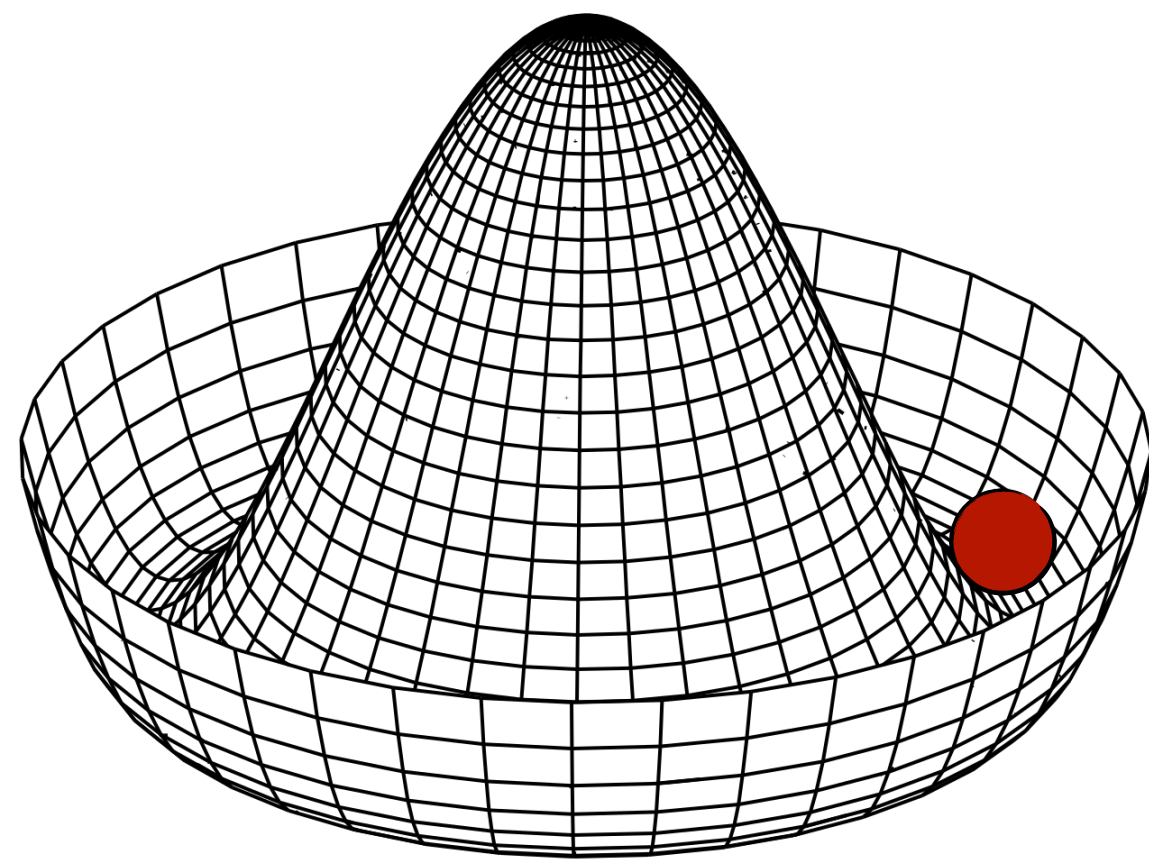


$$\delta \approx m_\rho^{-1}$$

QCD axion strings

- Potential for PQ field

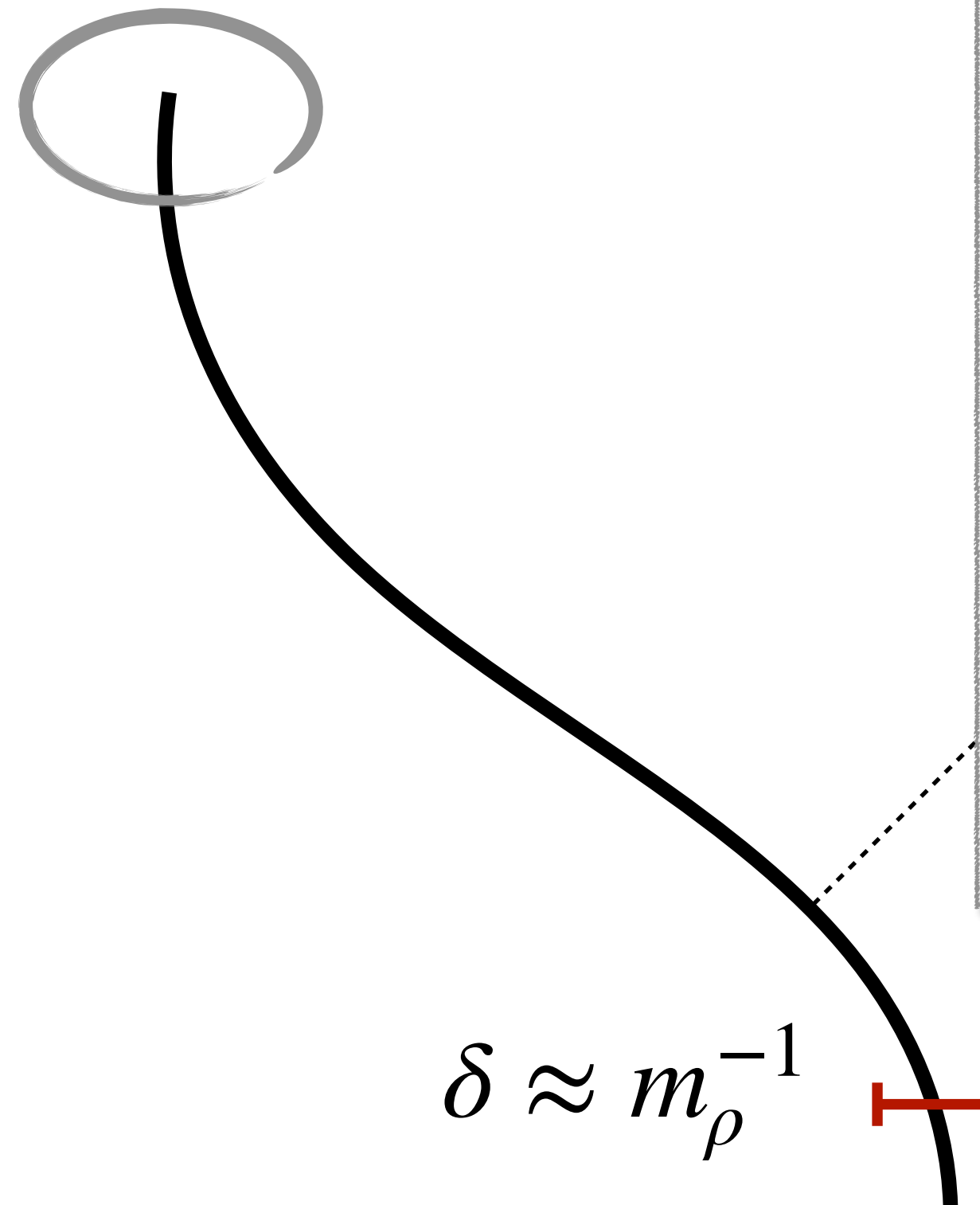
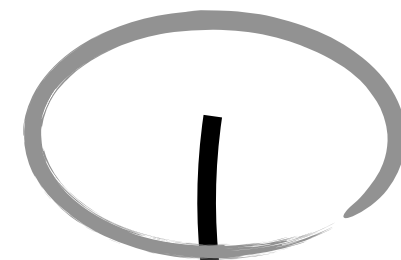
$$\Phi = \rho e^{i\alpha}$$



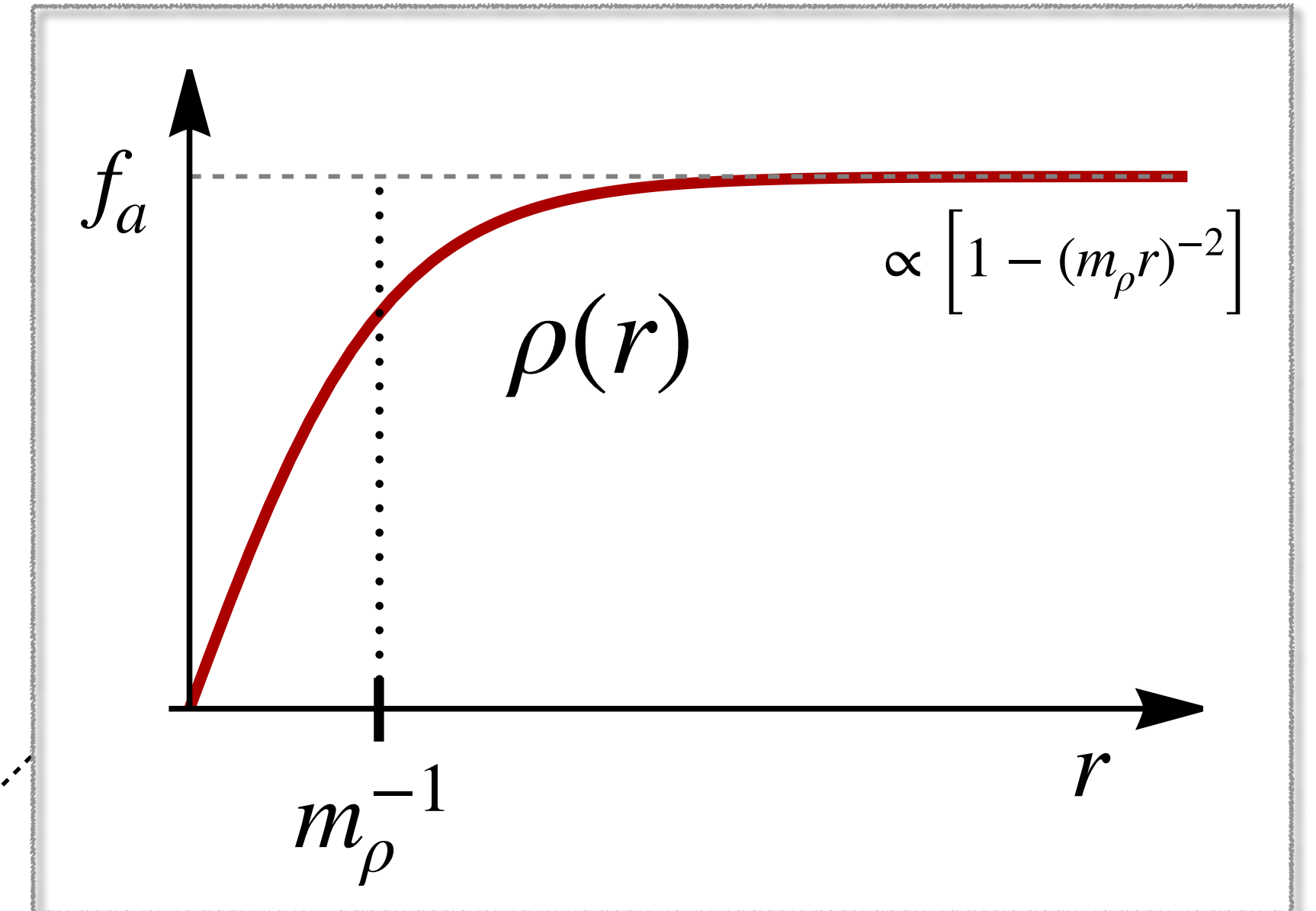
$$V_{\text{PQ}}(\Phi)$$

- Global string solution

$$\alpha(\theta) : 0 \rightarrow 2\pi$$



$$\delta \approx m_\rho^{-1}$$



QCD axion strings

- Relevant contribution to dark matter abundance
- Source of gravitational waves

Hiramatsu et al. [1012.5502] PRD
Gorghetto, Hardy, Villadoro
[1806.04677] JHEP; [2007.04990] SciPost

Gorghetto, Hardy, Nicolaescu [2101.11007] JCAP
Baeza-Ballesteros, Copeland, Figueroa, Lizarraga [2308.08456]

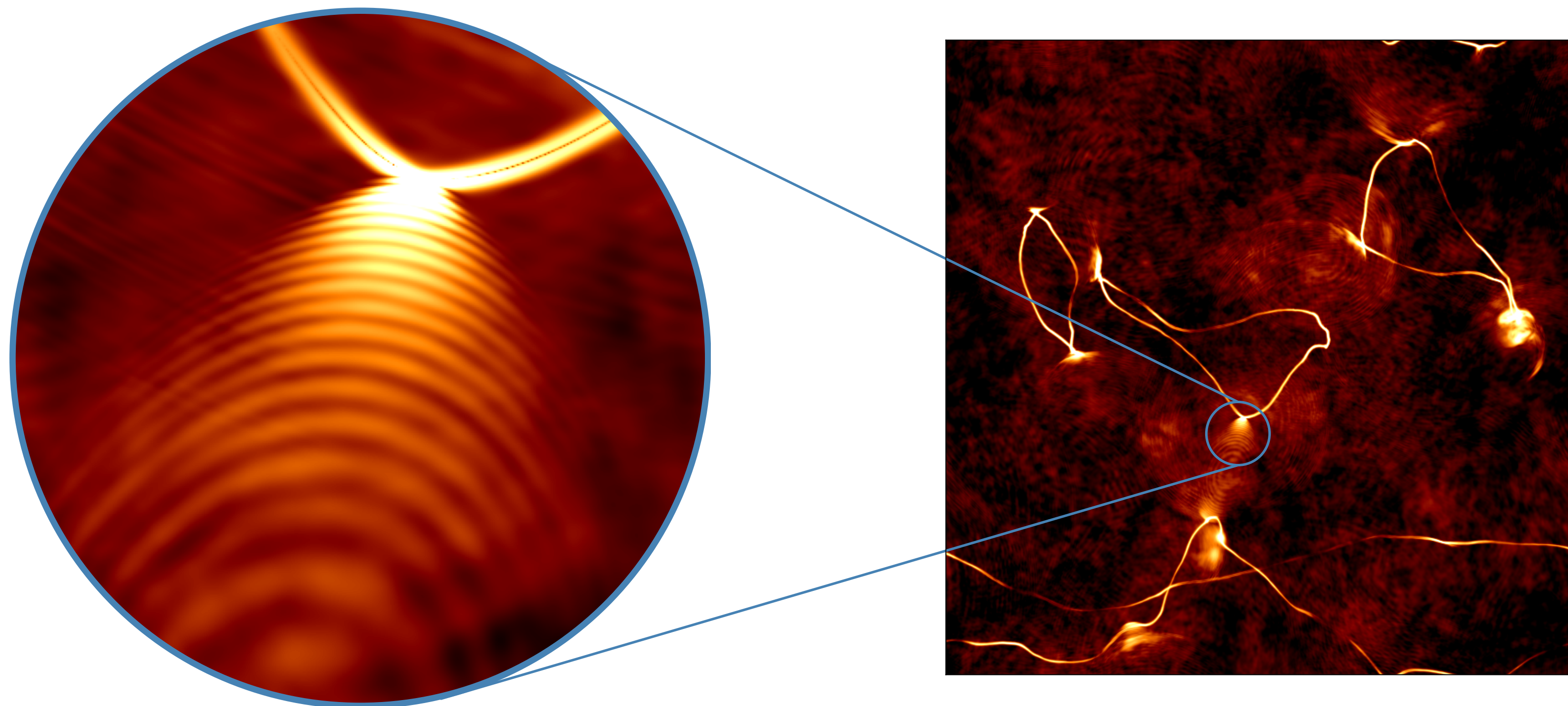


Fig. from Benabou, Buschmann,
Kumar, Park, Said [2308.01334], PRD

Axion strings during the EWPT

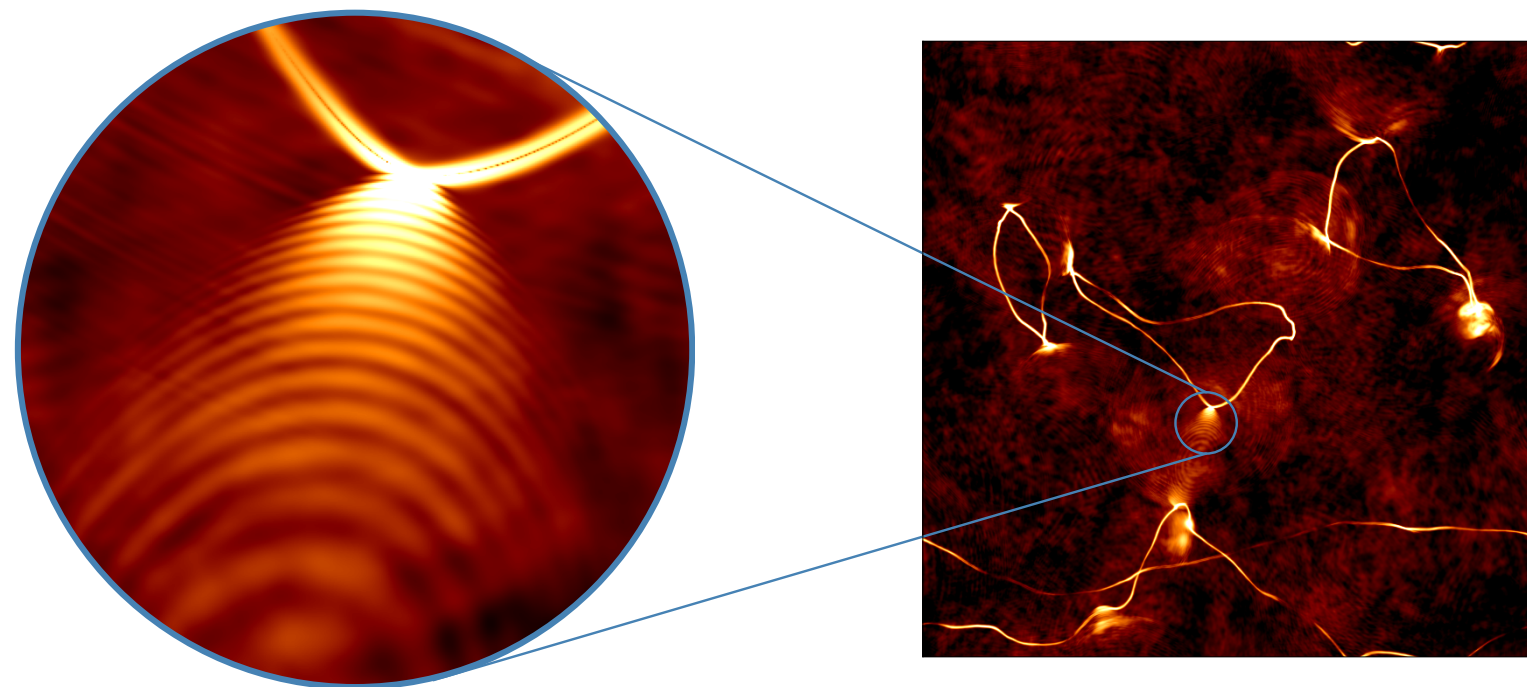
- Consider the minimal KSVZ axion model with a Higgs portal:

$$\mathcal{V} = V_{\text{PQ}}(|\Phi|) + V_{\text{EW}}(|\mathcal{H}|; T) + \kappa \left(|\Phi|^2 - \frac{f_a^2}{2} \right) \left(|\mathcal{H}|^2 - \frac{v^2}{2} \right)$$

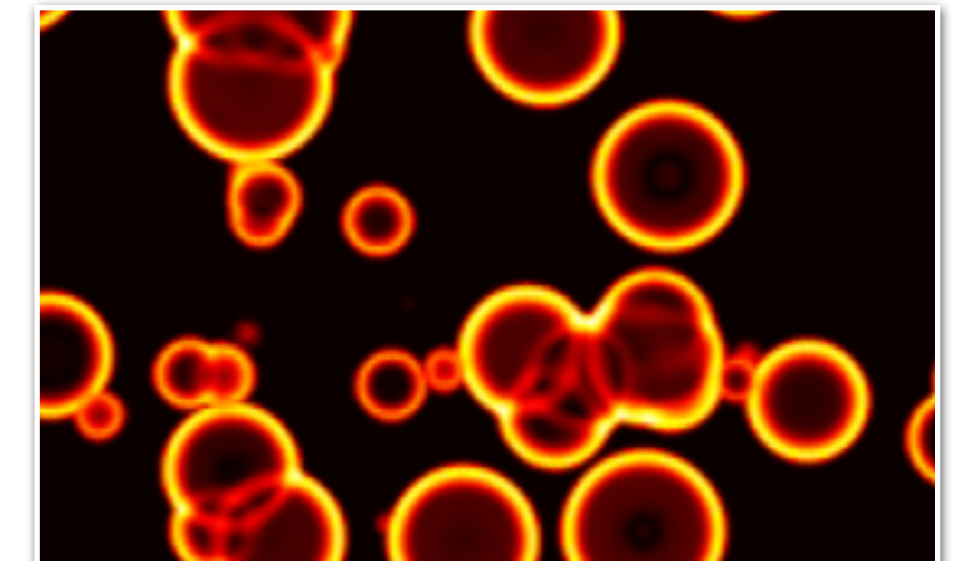
Axion strings during the EWPT

- Consider the minimal KSVZ axion model with a Higgs portal:

$$\mathcal{V} = V_{\text{PQ}}(|\Phi|) + V_{\text{EW}}(|\mathcal{H}|; T) + \kappa \left(|\Phi|^2 - \frac{f_a^2}{2} \right) \left(|\mathcal{H}|^2 - \frac{v^2}{2} \right)$$

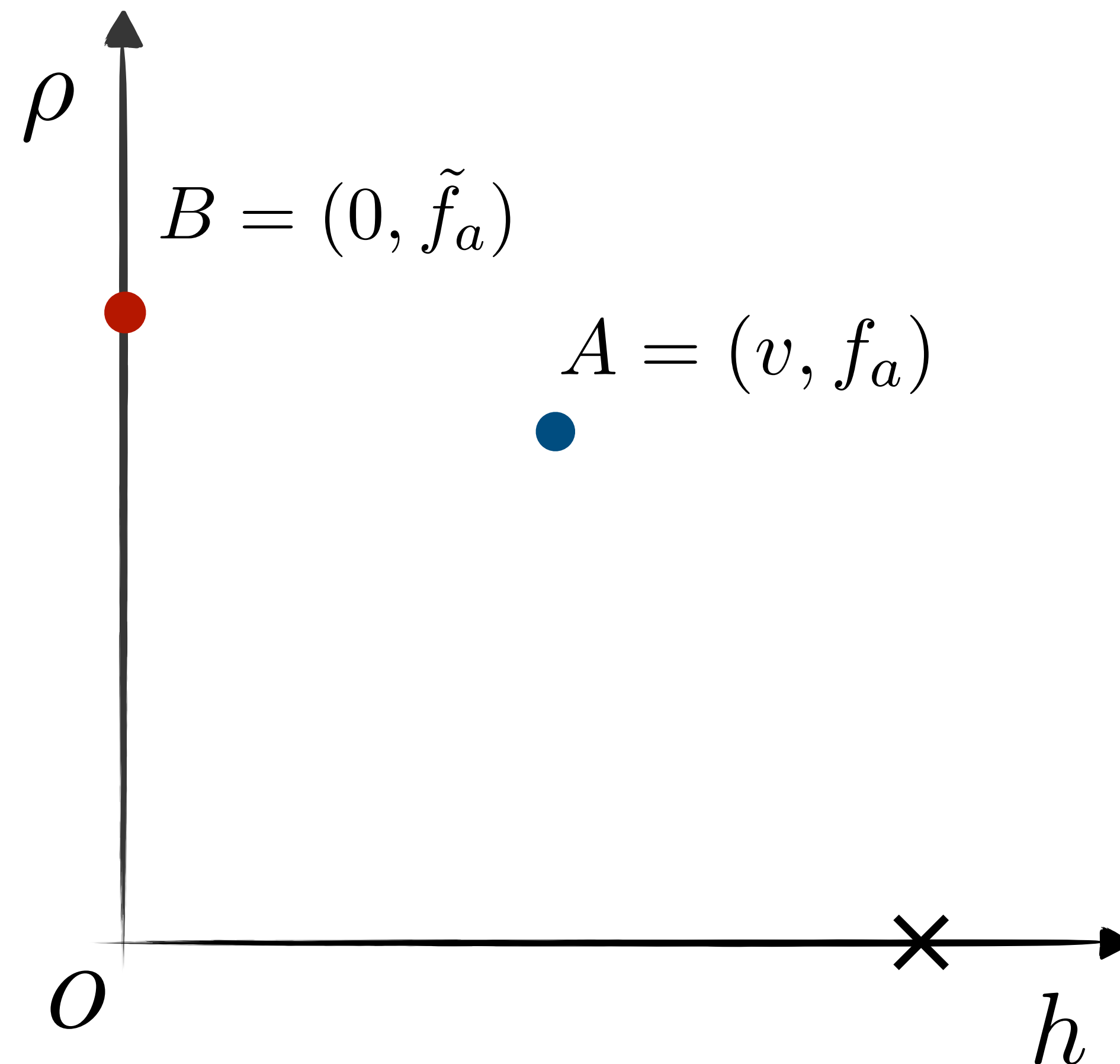


.....
How does this affect electroweak
symmetry breaking?
.....



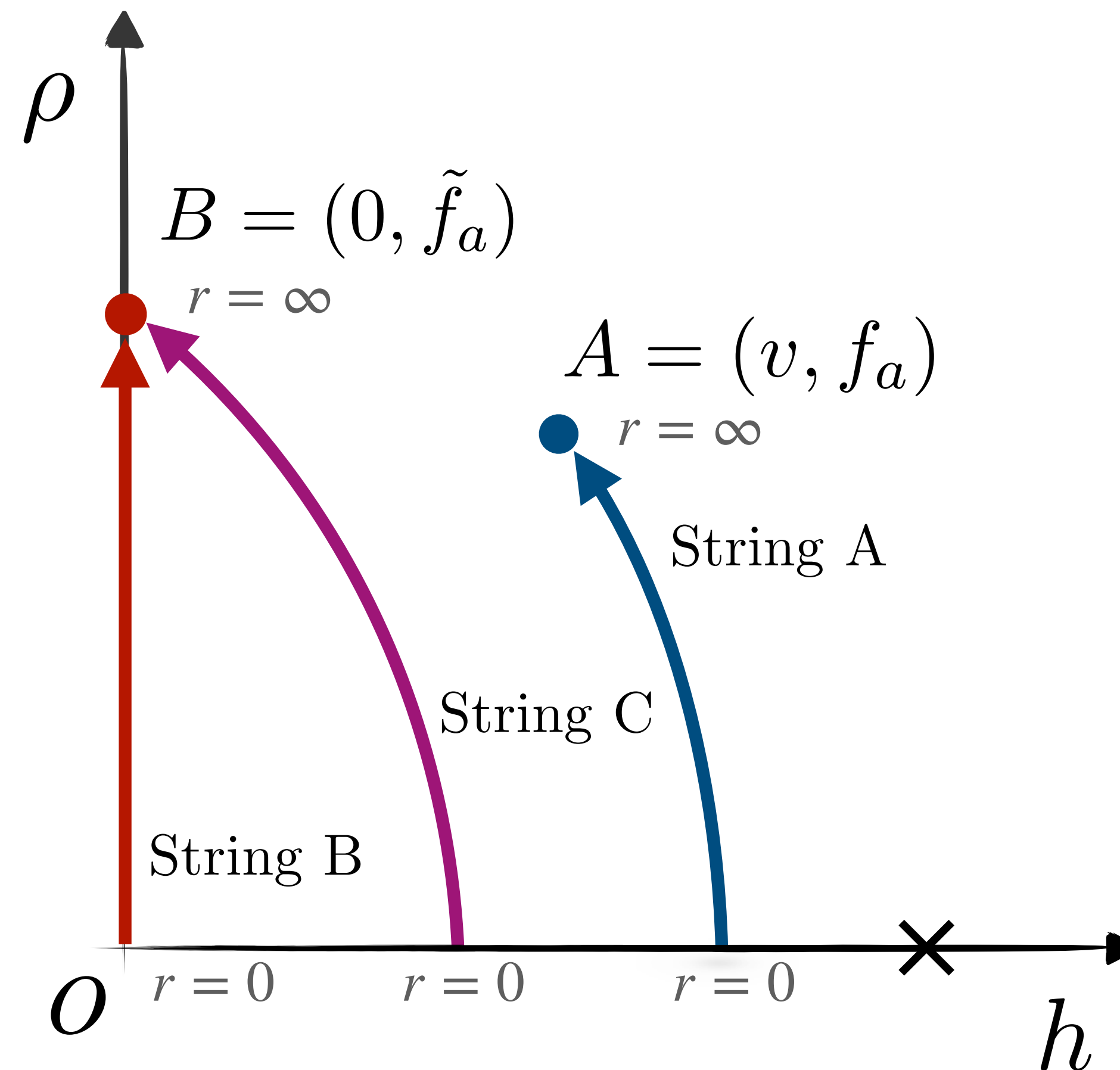
Axion strings during the EWPT

- Relevant points in field space



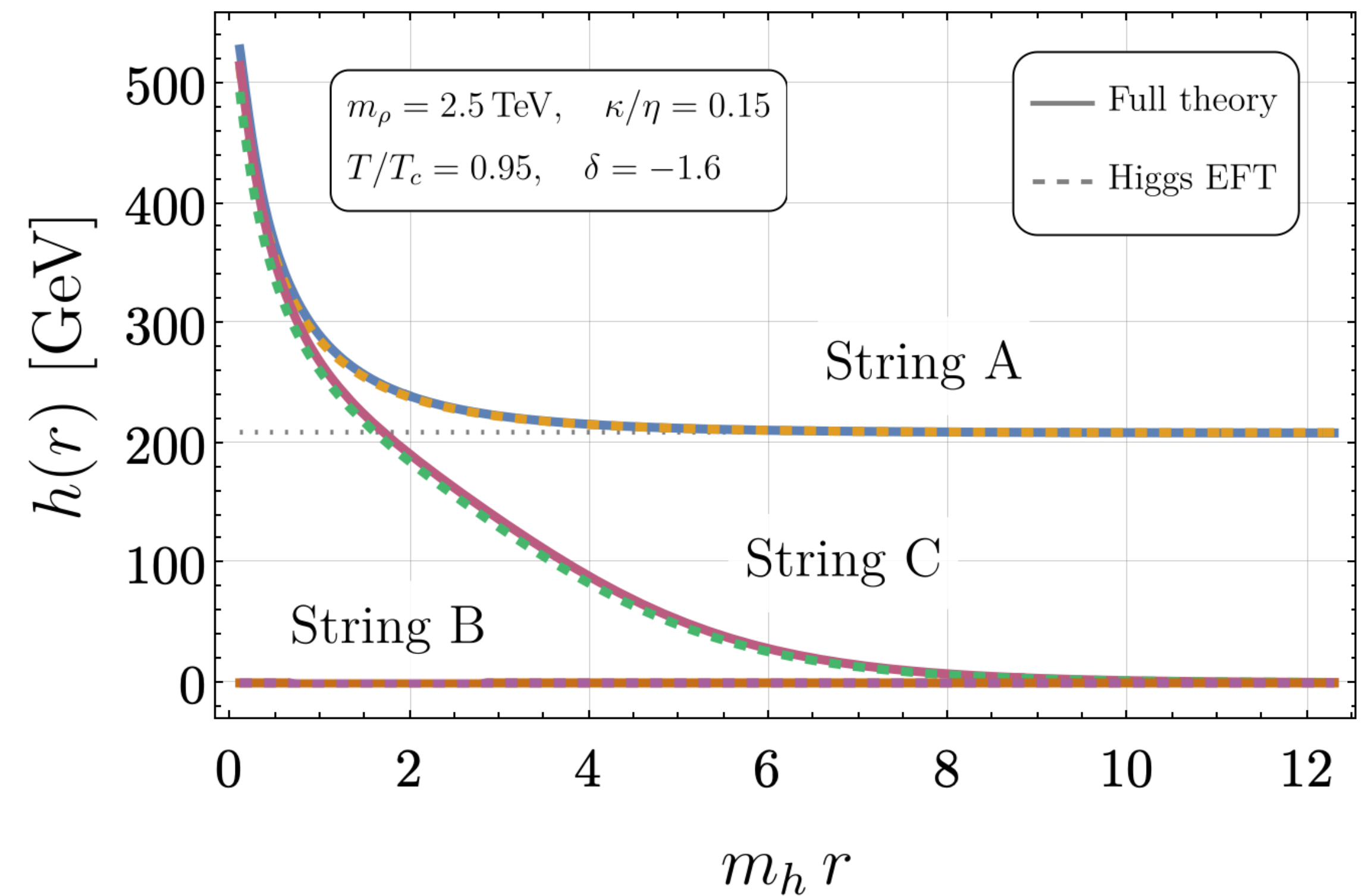
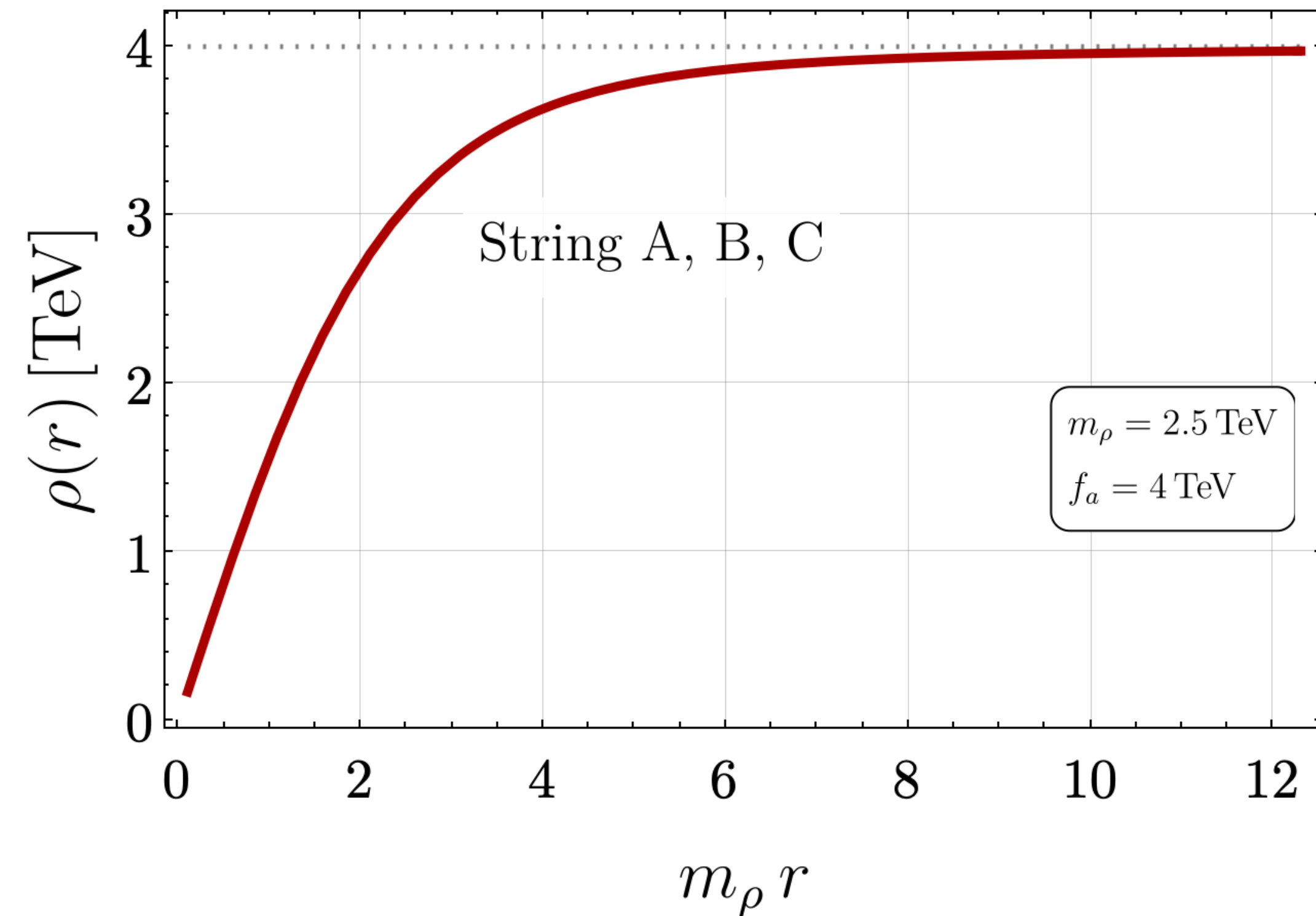
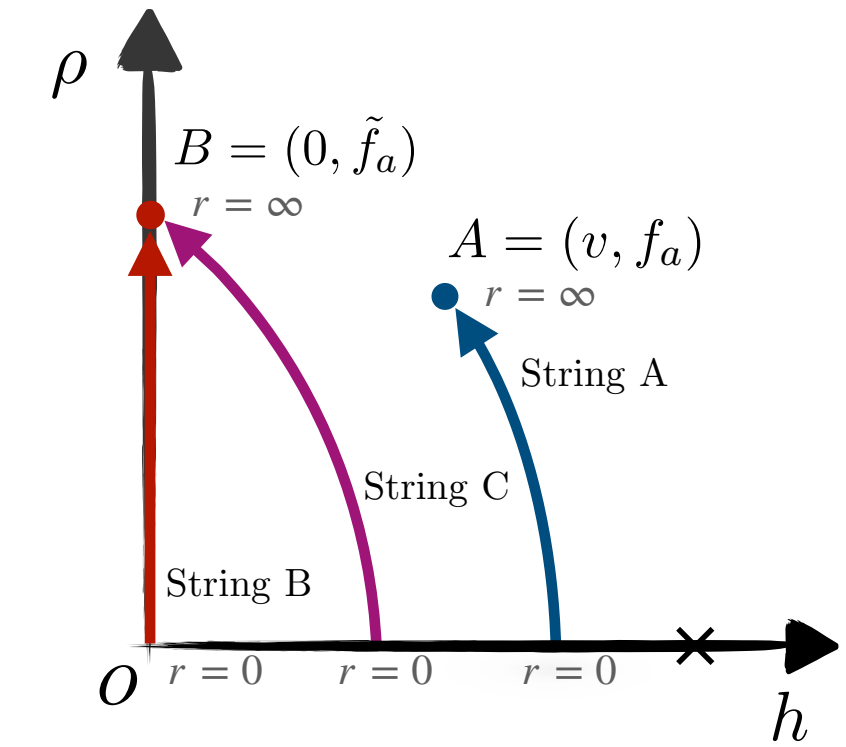
Axion strings during the EWPT

- Relevant points in field space and possible string solutions:



Axion strings during the EWPT

- Relevant points in field space and possible string solutions:



Axion strings during the EWPT

- Large hierarchy between the mass of the Higgs and the PQ radial mode

Axion strings during the EWPT

- Large hierarchy between the mass of the Higgs and the PQ radial mode
- Physics captured by electroweak-scale EFT (SM + axion or ALP):

$$S_{\text{EFT}}[h] = \int d^4x \left\{ \frac{1}{2} (\partial_\mu h)^2 - V_{\text{EW}}(h) - \frac{1}{2} \frac{\kappa}{\eta} (\partial_\mu \alpha)^2 h^2 + \pi \frac{\kappa}{\eta} C(\epsilon) \delta^{(2)}(r - \epsilon) h^2 \right\}$$

Axion strings during the EWPT

- Large hierarchy between the mass of the Higgs and the PQ radial mode
- Physics captured by electroweak-scale EFT (SM + axion or ALP):

$$S_{\text{EFT}}[h] = \int d^4x \left\{ \frac{1}{2} (\partial_\mu h)^2 - V_{\text{EW}}(h) - \frac{1}{2} \frac{\kappa}{\eta} (\partial_\mu \alpha)^2 h^2 + \pi \frac{\kappa}{\eta} C(\epsilon) \delta^{(2)}(r - \epsilon) h^2 \right\}$$

Axion strings during the EWPT

- Large hierarchy between the mass of the Higgs and the PQ radial mode
- Physics captured by electroweak-scale EFT (SM + axion or ALP):

$$S_{\text{EFT}}[h] = \int d^4x \left\{ \frac{1}{2} (\partial_\mu h)^2 - V_{\text{EW}}(h) - \frac{1}{2} \frac{\kappa}{\eta} (\partial_\mu \alpha)^2 h^2 + \pi \frac{\kappa}{\eta} C(\epsilon) \delta^{(2)}(r - \epsilon) h^2 \right\}$$

- Explicit UV scale:

$$\epsilon \sim 1/m_\rho$$

- Axion-Higgs portal, in the string background:

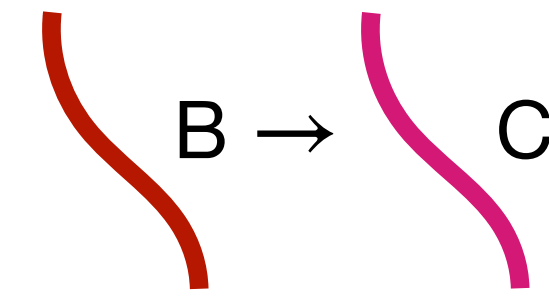
$$\alpha = \theta \Rightarrow \partial_\mu \alpha \sim 1/r$$

- δ -potential imposes UV matching condition:

$$\epsilon h'(\epsilon) = -C(\epsilon) \frac{\kappa}{\eta} h(\epsilon)$$

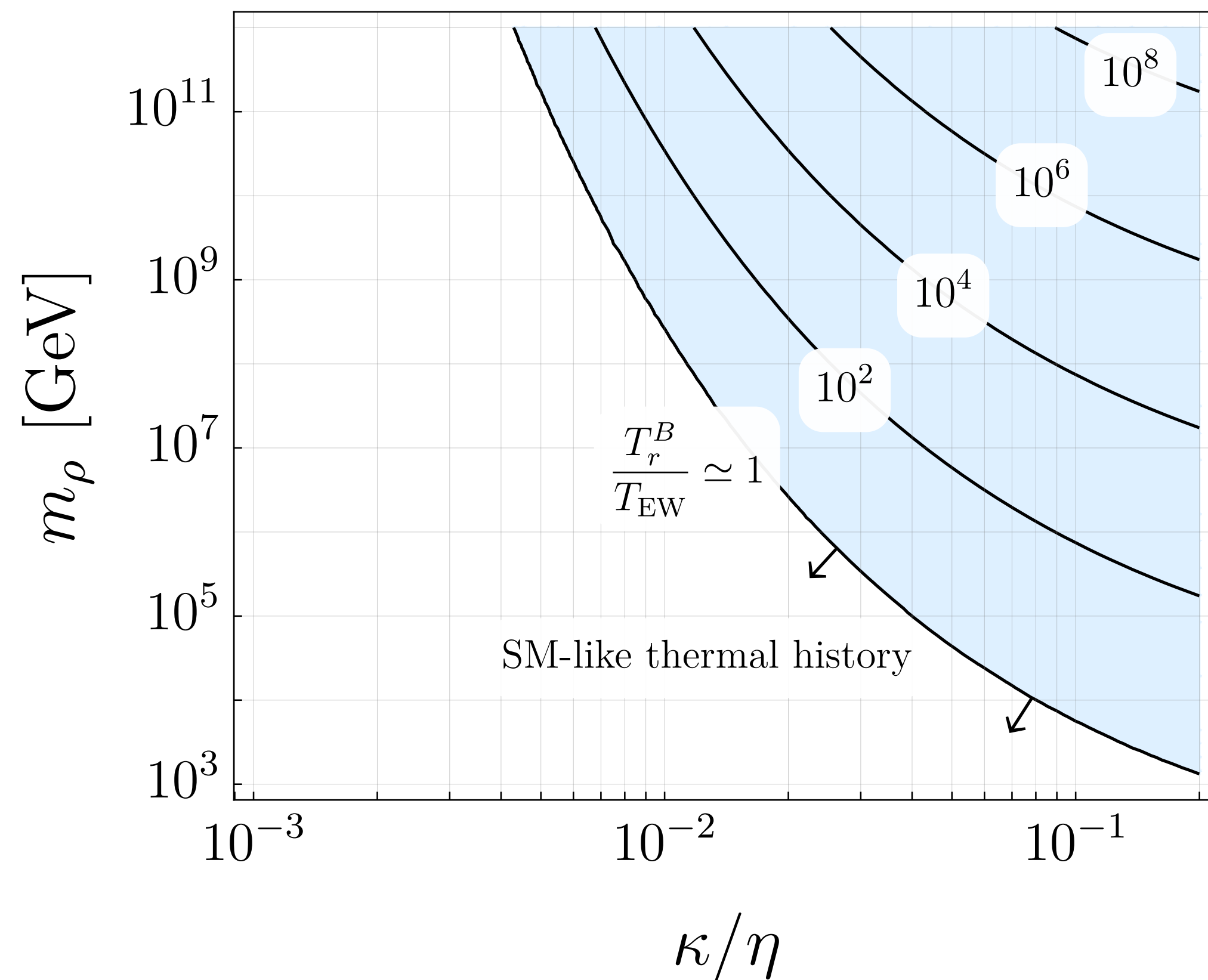
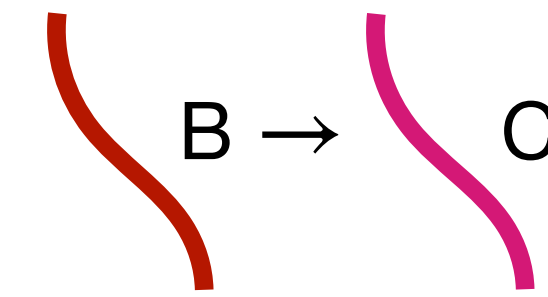
SM + PQ

- Consider the minimal Higgs potential of the SM
- Standard axion string develops a Higgs core at $T = T_r^B$:



SM + PQ

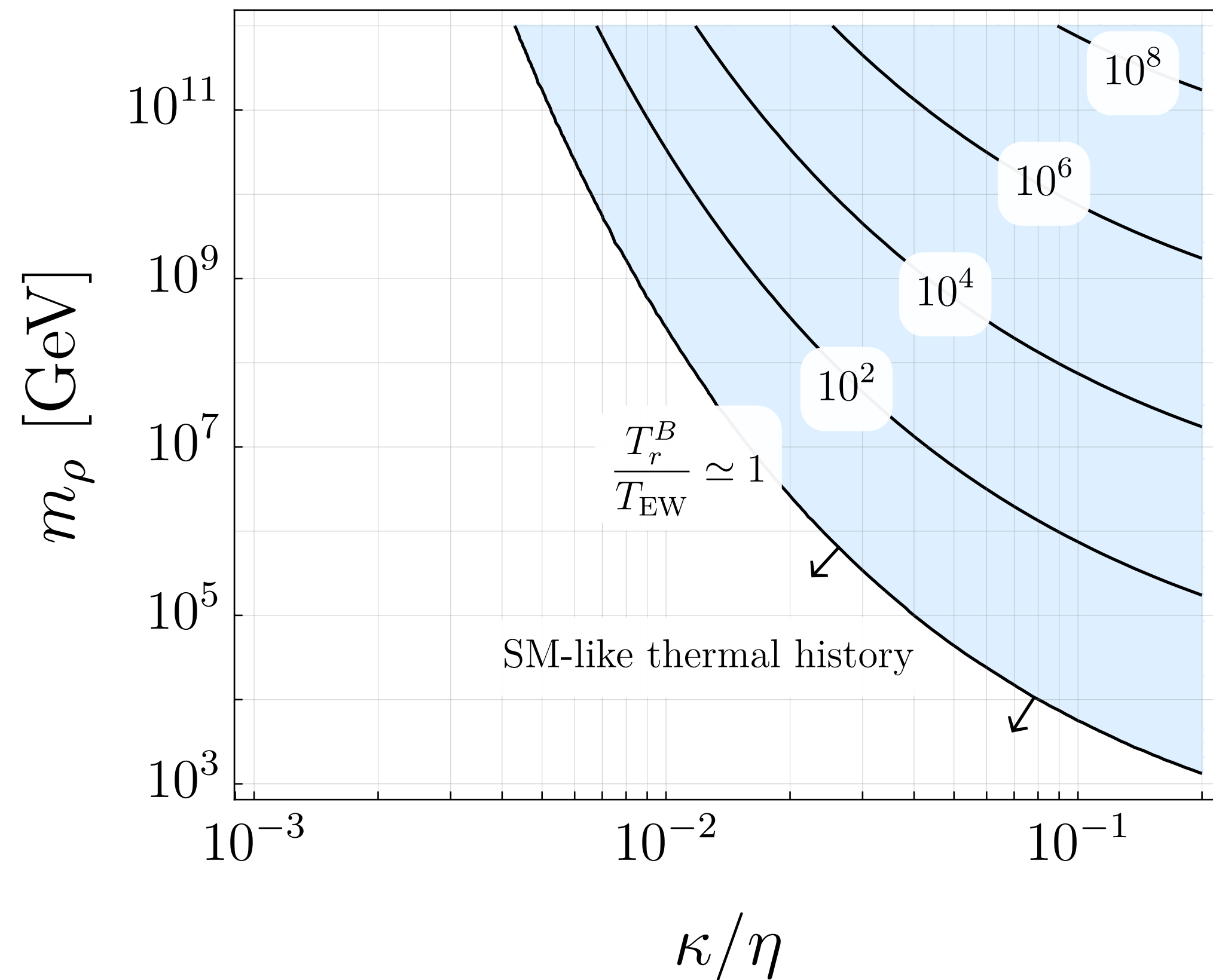
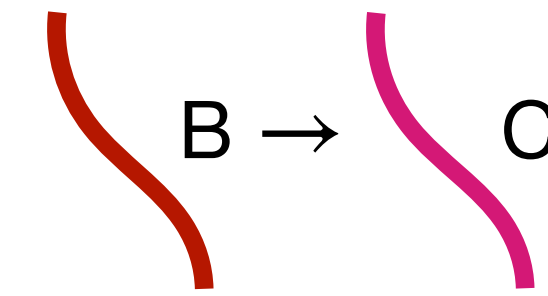
- Consider the minimal Higgs potential of the SM
- Standard axion string develops a Higgs core at $T = T_r^B$:



$$\omega^2(T_r^B) = V''_{EW}(0; T_r^B) - \frac{1}{2}m_\rho^2 f(\kappa/\eta) = 0$$

SM + PQ

- Consider the minimal Higgs potential of the SM
- Standard axion string develops a Higgs core at $T = T_r^B$:



$$\omega^2(T_r^B) = V''_{\text{EW}}(0; T_r^B) - \frac{1}{2} m_\rho^2 f(\kappa/\eta) = 0$$

\downarrow 2D mass \downarrow 4D mass \downarrow Δm_h^2

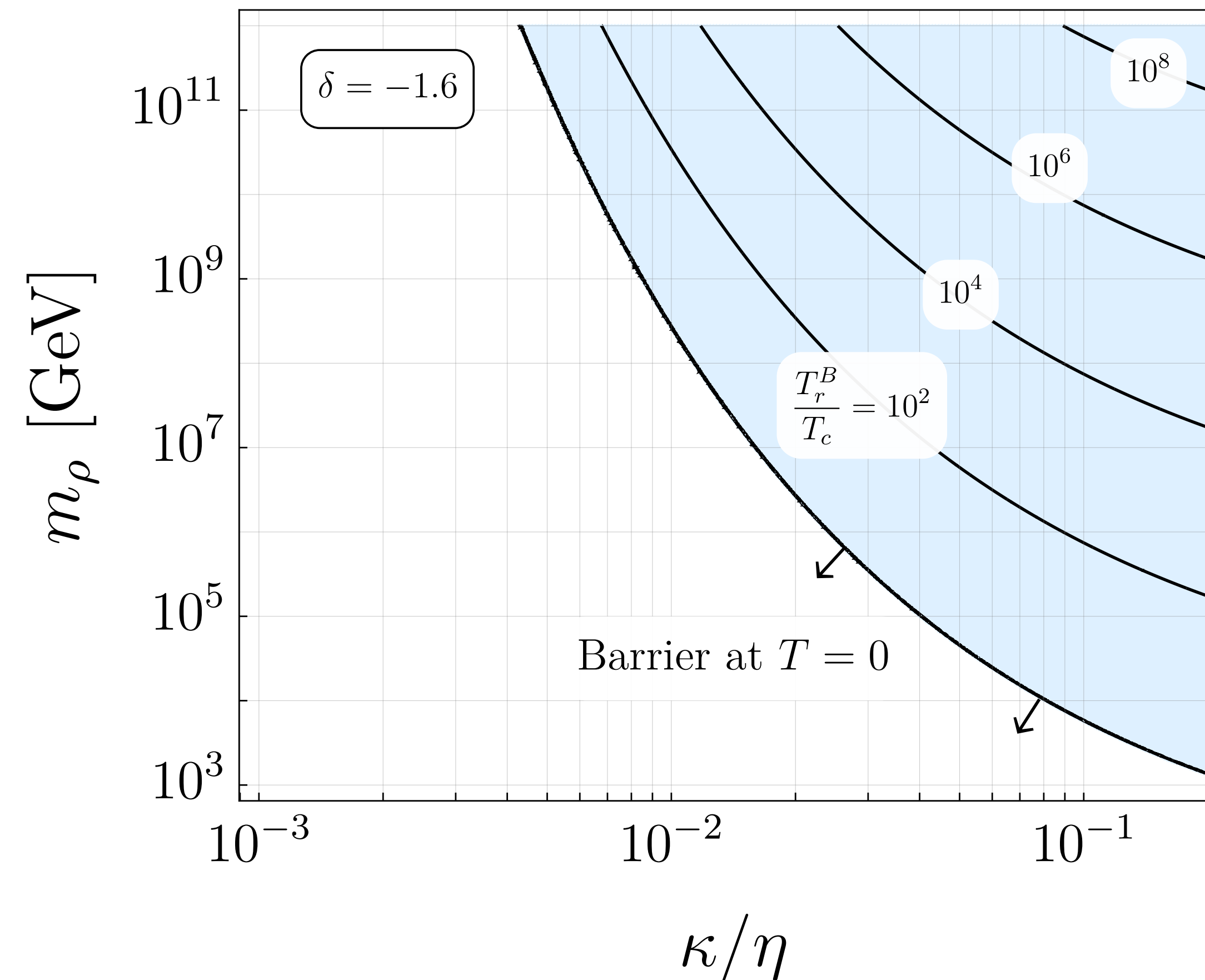
$$f(\kappa/\eta) = \exp \left\{ -\frac{\pi}{\sqrt{\kappa/\eta}} - \gamma_E + 2C(\epsilon) \right\}$$

First order EWPT + PQ

- Consider BSM model with a first order (hom.) EWPT $B \rightarrow A$ which fails to nucleate: transition *still possible* around the axion strings

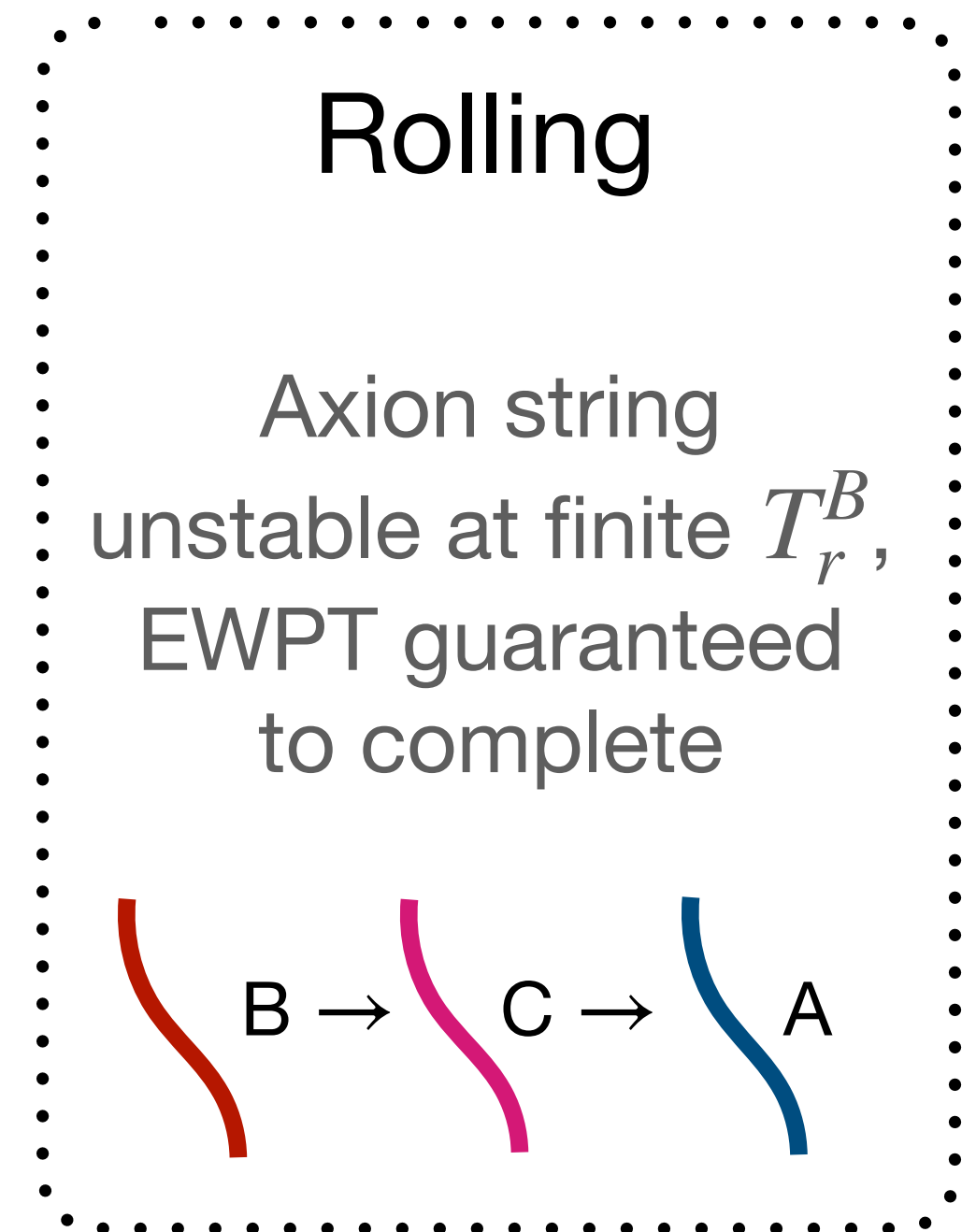
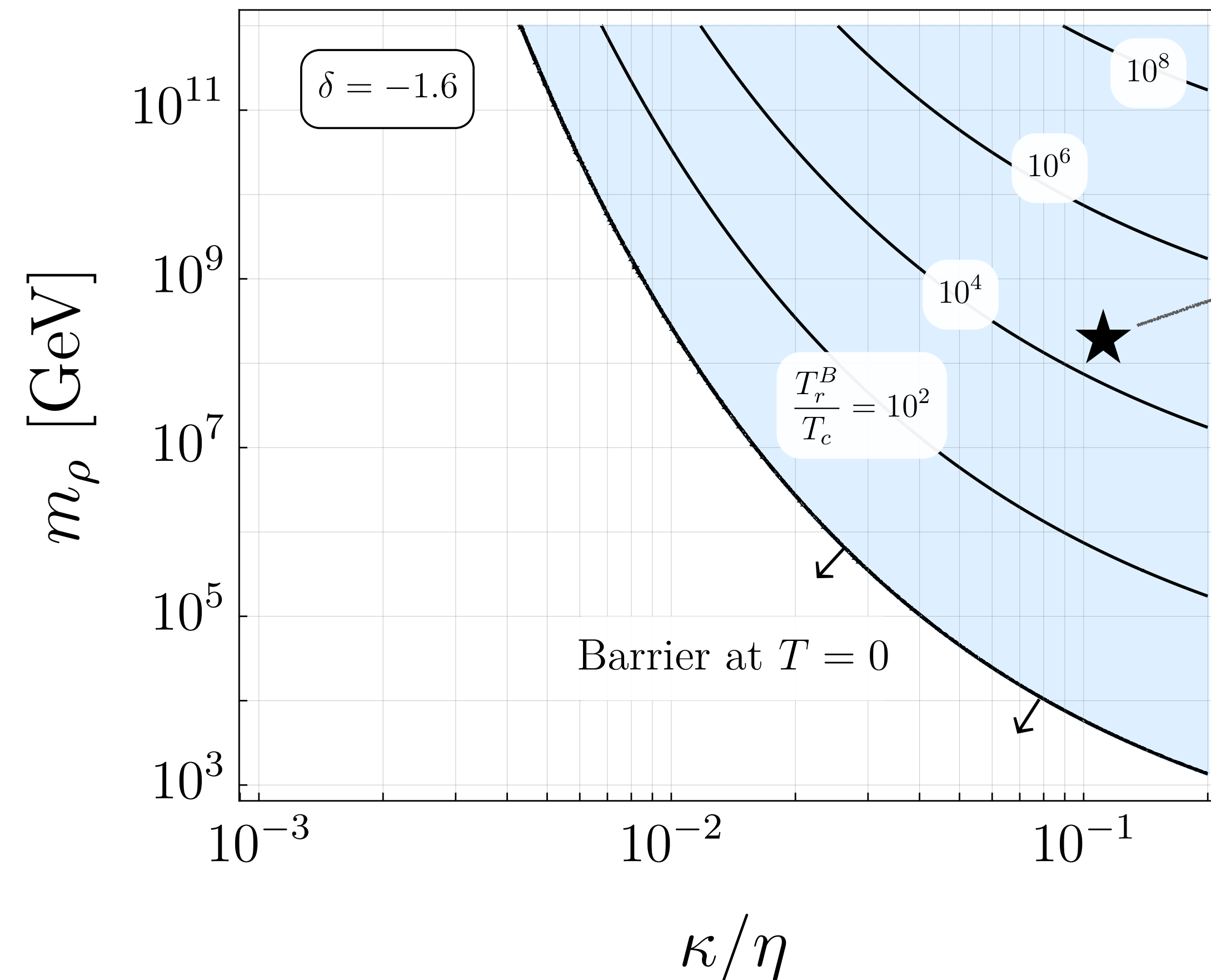
First order EWPT + PQ

- Consider BSM model with a first order (hom.) EWPT $B \rightarrow A$ which fails to nucleate: transition *still possible* around the axion strings



First order EWPT + PQ

- Consider BSM model with a first order (hom.) EWPT $B \rightarrow A$ which fails to nucleate: transition *still possible* around the axion strings

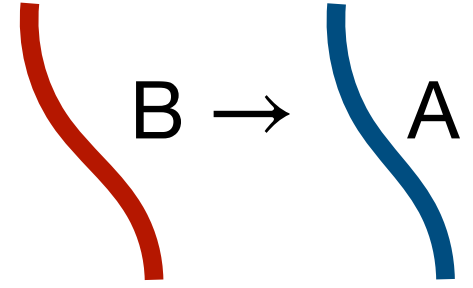
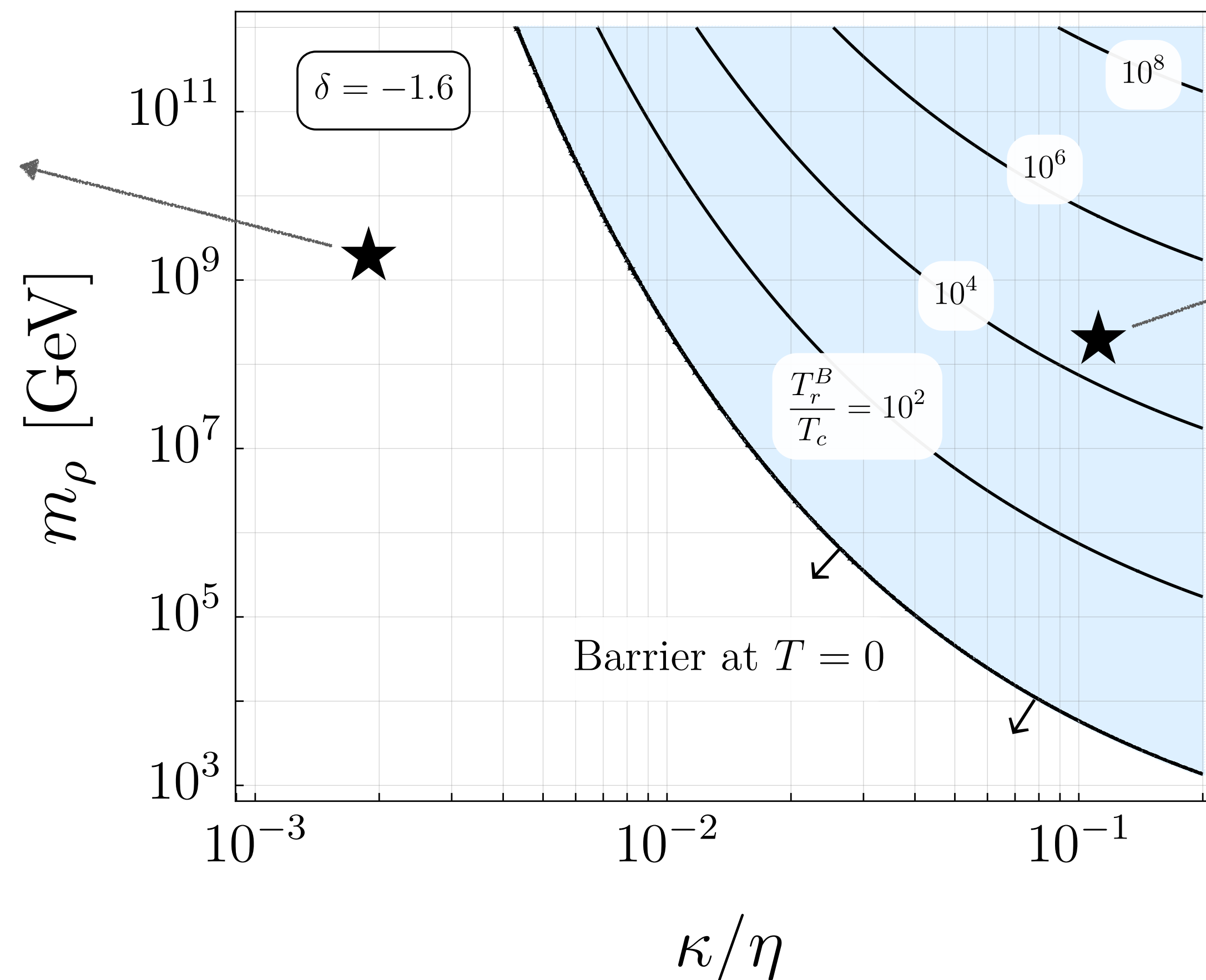


First order EWPT + PQ

- Consider BSM model with a first order (hom.) EWPT $B \rightarrow A$ which fails to nucleate: transition *still possible* around the axion strings

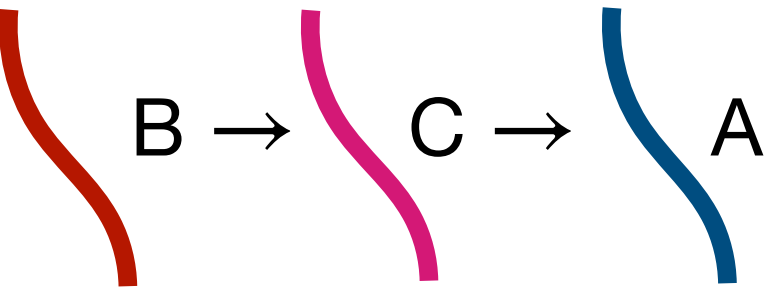
Seeded tunneling

Axion string metastable at $T = 0$: EWPT may complete if fast enough rate

Rolling

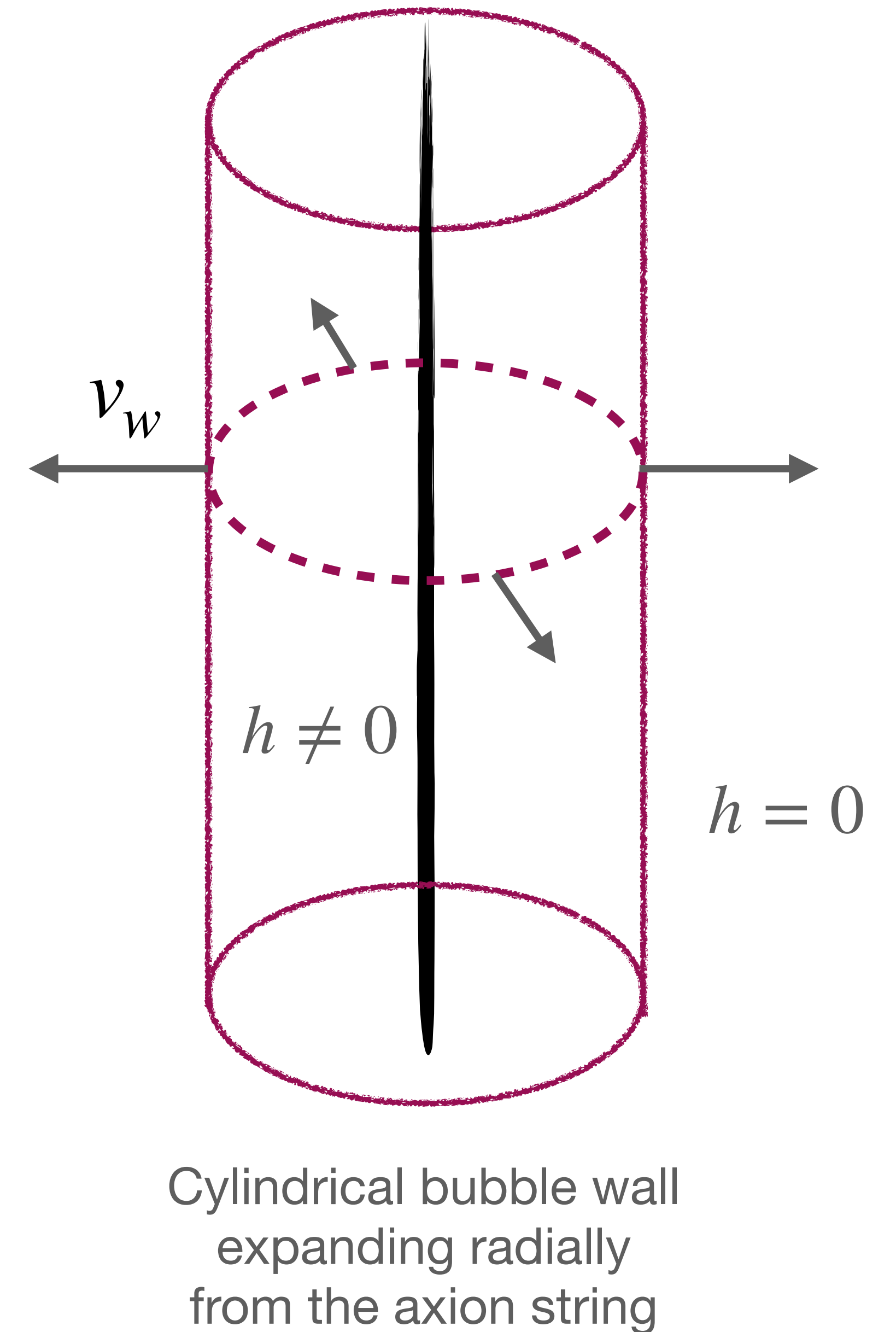
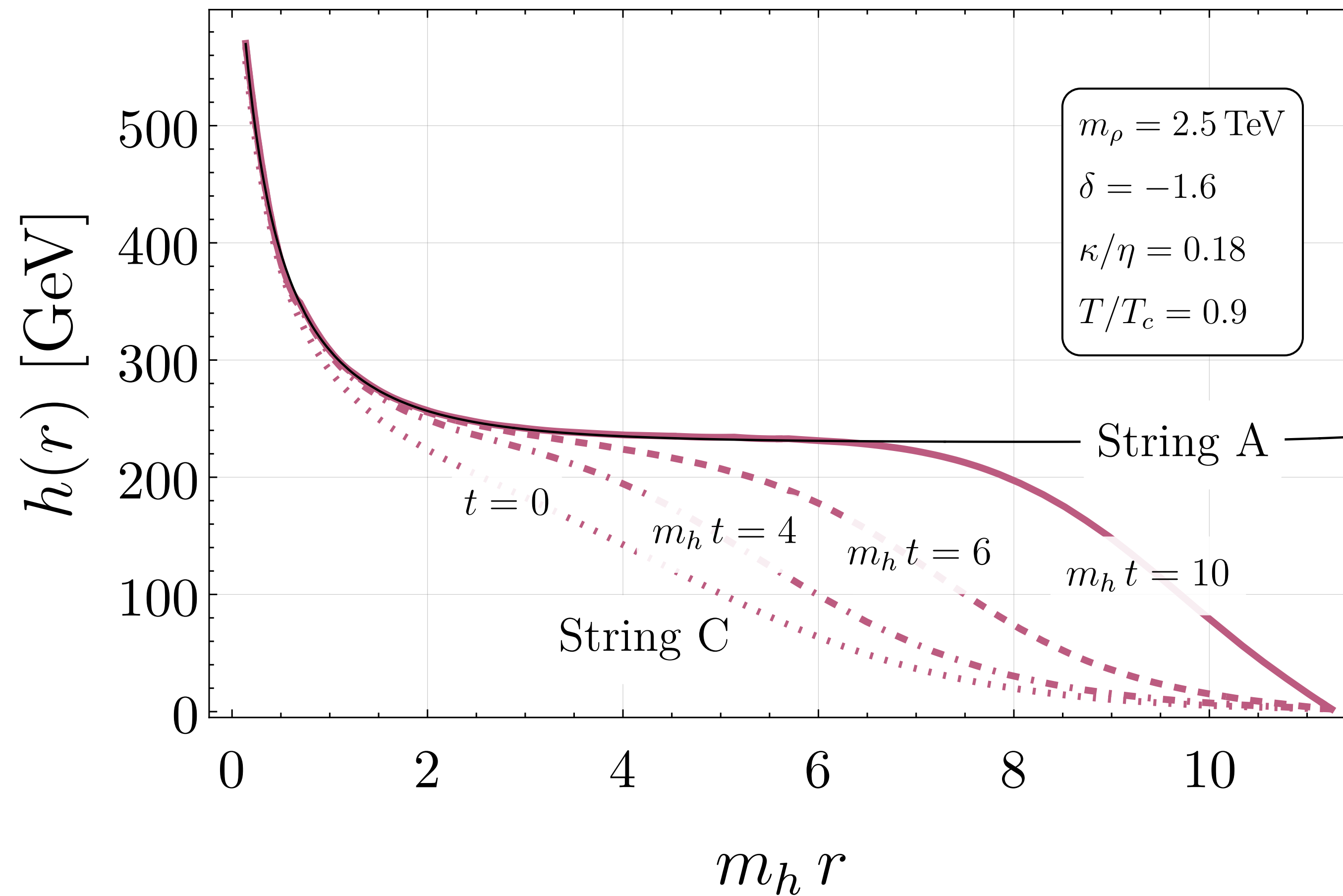
Axion string unstable at finite T_r^B , EWPT guaranteed to complete



Rolling

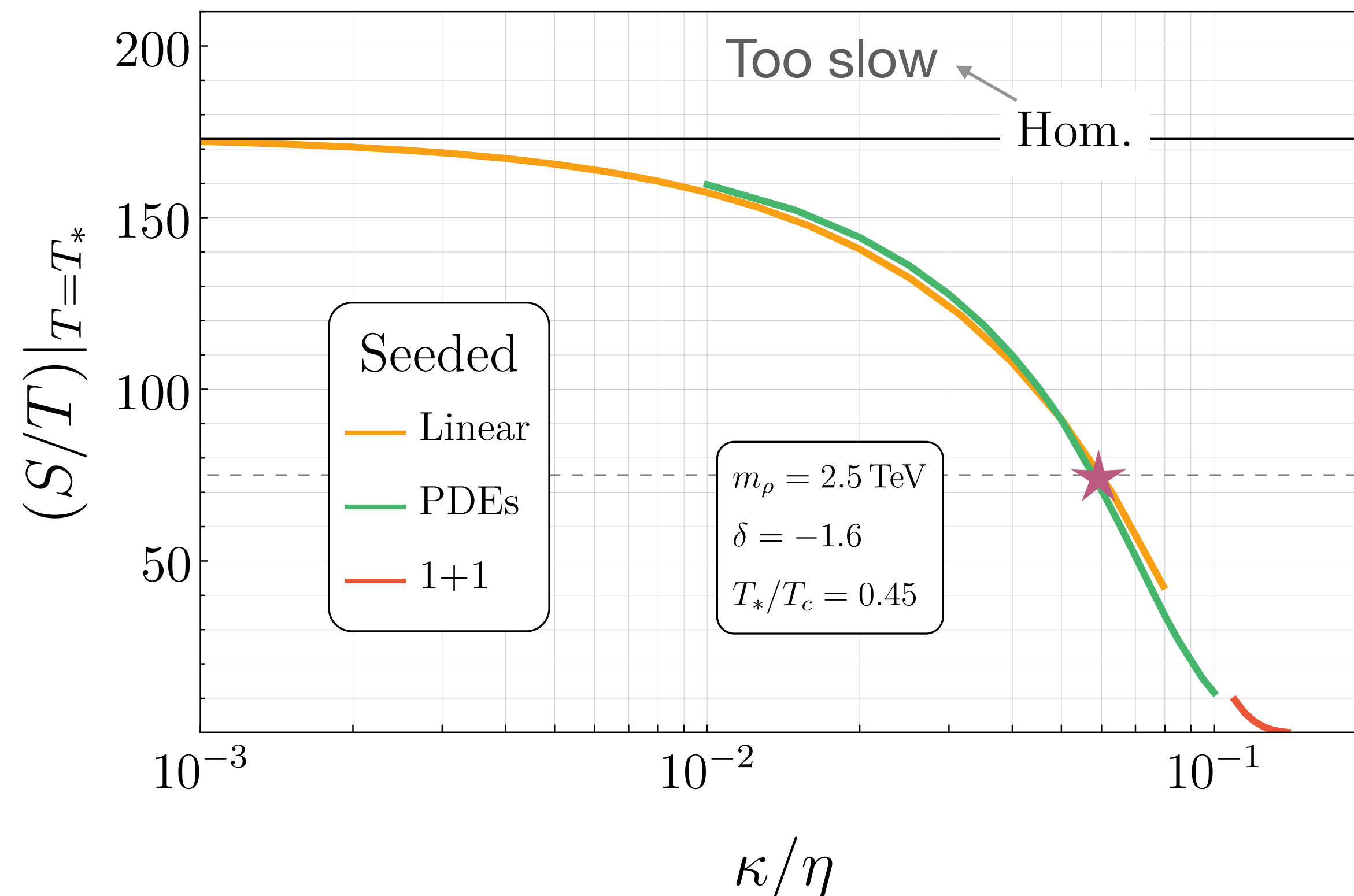
See also Yajnik, PRD (1986)

- Classical evolution from string C to string A at $T \approx T_c$



Seeded tunneling

- Bubble of broken electroweak symmetry nucleated around the axion string with a certain (exponential) rate S/T :



- **Linear** \sim thin wall approximation:

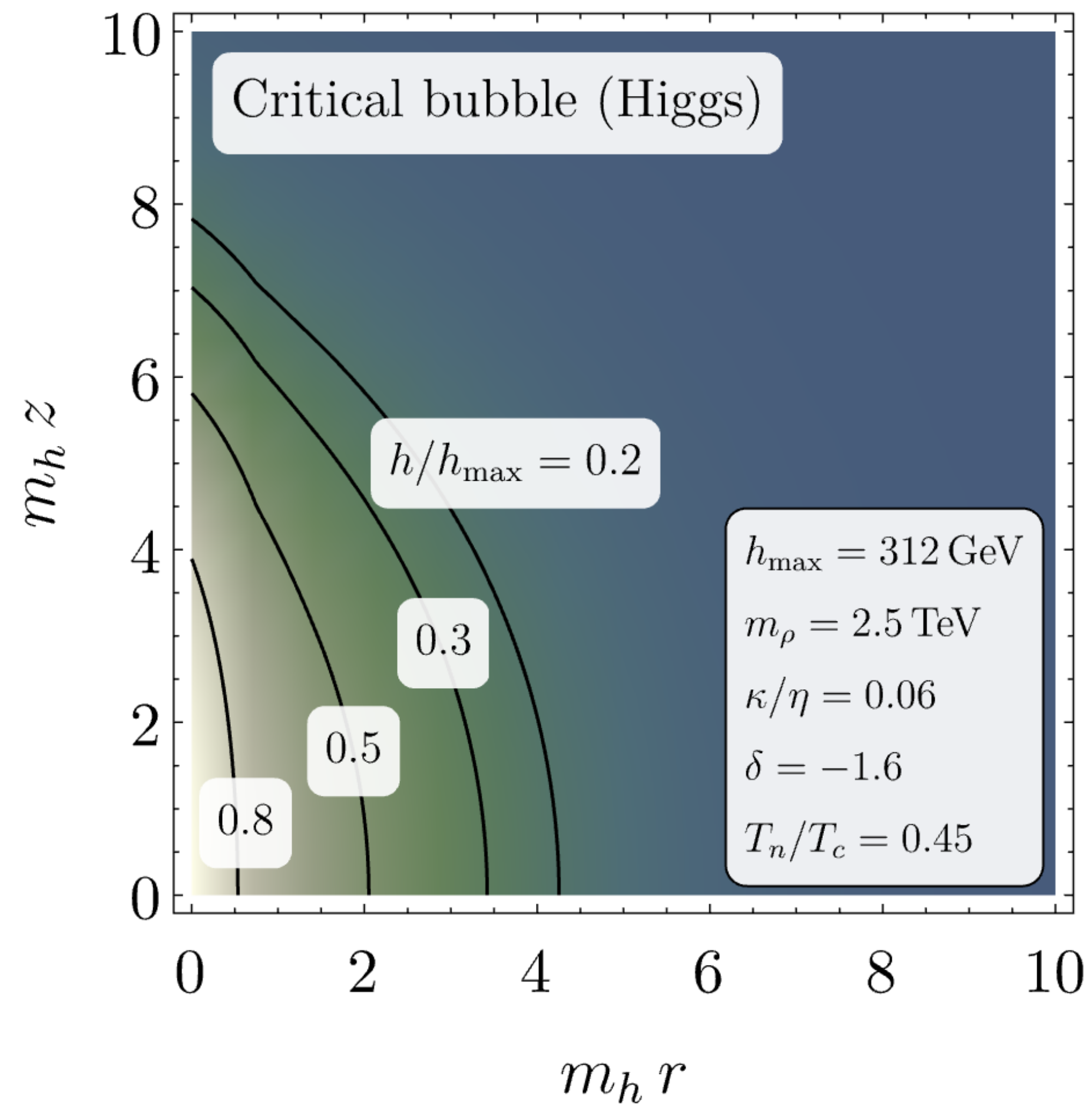
$$\delta S_{\text{TW}} = -2\pi R \frac{\kappa}{\eta} \log(R m_\rho) h_r^2(0) \equiv 2\pi R \Delta\mu_{\text{eff}}$$

- **EFT** on the string for the lightest mode:

$$S_{1+1}[h_0] = \int dz dt \left\{ \frac{1}{2} (\partial_\mu h_0)^2 - \tilde{V}(h_0) \right\}$$

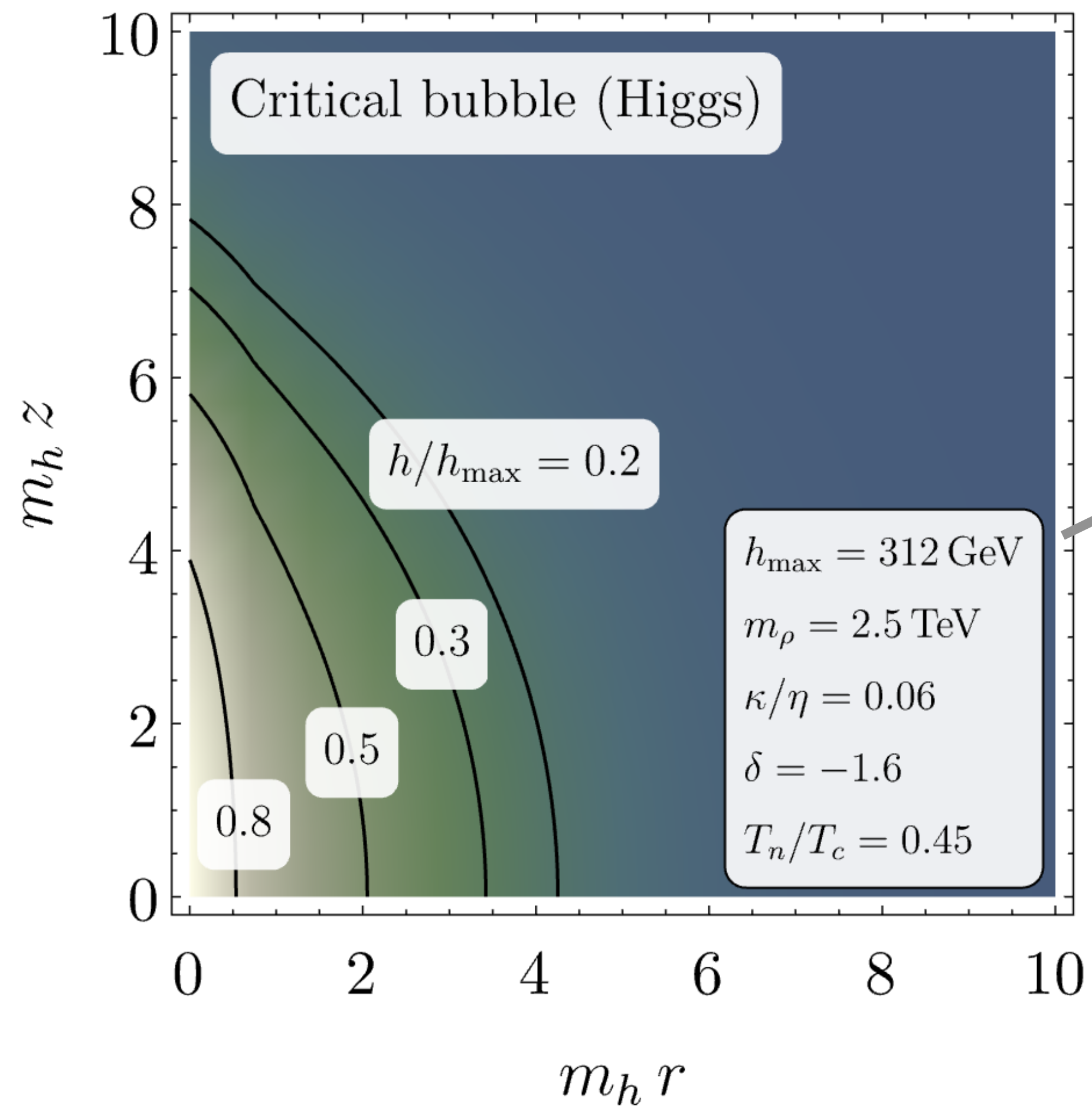
Seeded tunneling

- Profile of the critical bubble: ★



Seeded tunneling

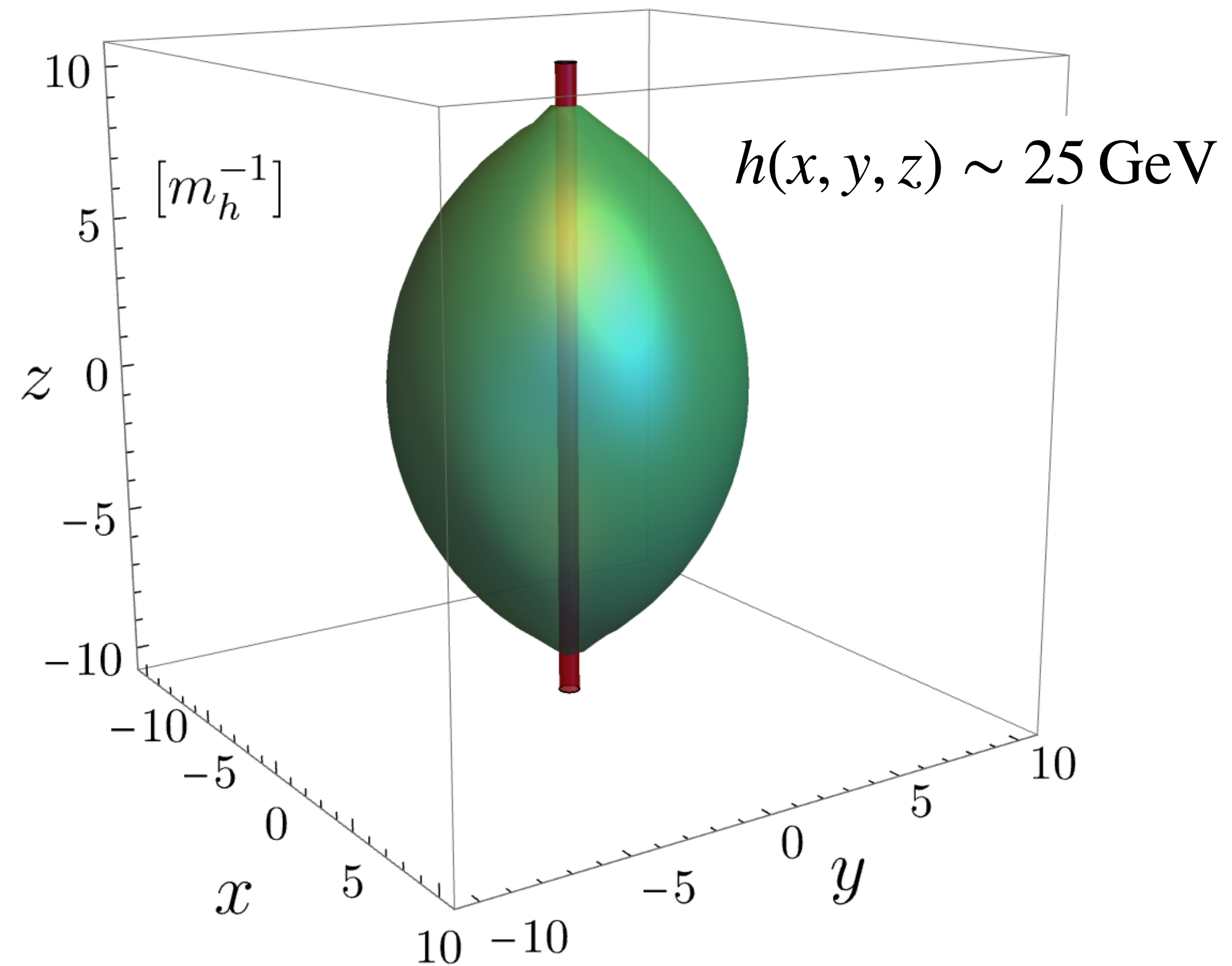
- Profile of the critical bubble: ★



Release point is larger than the VEV: the bubble is partially reconstructing string A

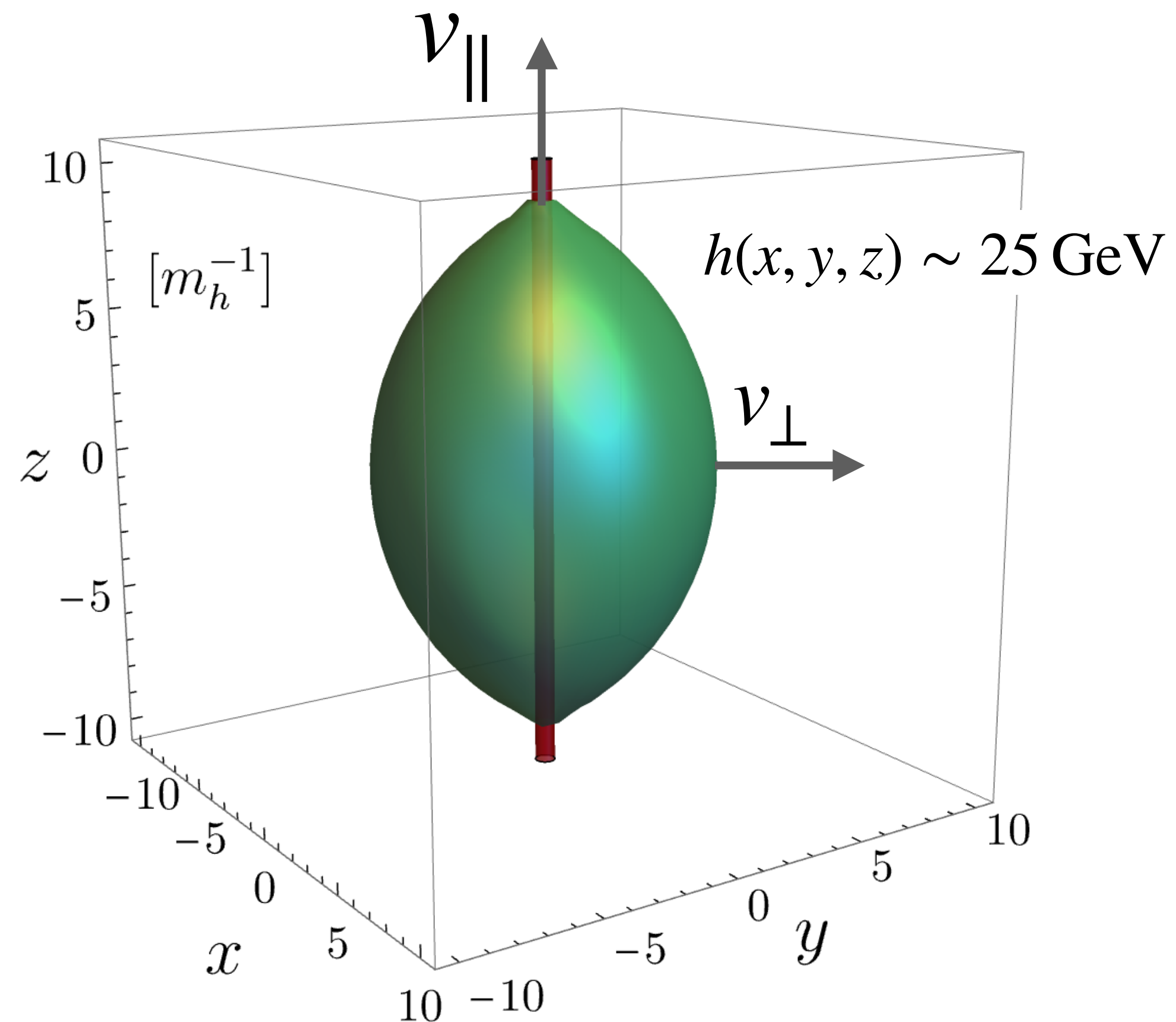
Seeded tunneling

- Profile of the critical bubble: ★



Seeded tunneling

- Profile of the critical bubble: ★



- Percolation as interplay between seeded nucleation rate and density of defects
- Is nucleation rate the same on loops?
- Different velocities parallel or orthogonal to the string?
- Gravitational wave emission before collision (non-spherical bubbles, same for *rolling* case)

Summary

- The presence of impurities in the early Universe can strongly affect the way a phase transition proceeds
- The xSM with Z_2 symmetry is arguably the simplest (and complete) example for a seeded EWPT
- Other defects can exist at the time of the EWPT: dedicated study of QCD axion strings in KSVZ model with Higgs portal
- Pheno aspects of seeded phase transitions: percolation, slow transitions, expansion of non-spherical bubbles, features in the GW signal?

Summary

- The presence of impurities in the early Universe can strongly affect the way a phase transition proceeds
- The xSM with Z_2 symmetry is arguably the simplest (and complete) example for a seeded EWPT
- Other defects can exist at the time of the EWPT: dedicated study of QCD axion strings in KSVZ model with Higgs portal
- Pheno aspects of seeded phase transitions: percolation, slow transitions, expansion of non-spherical bubbles, features in the GW signal?

Thank you!