



Exotic cousins of the proton: tetraquarks, pentaquarks and hadronic molecules

Marek Karliner
Tel Aviv University
Joint work with Jon Rosner

35-th Recontres de Blois, Oct 23 2024

Outline

- quarks are fundamental building blocks of protons, neutrons and all hadrons
- all quarks are equal, but heavy quarks are more equal than others

new combinations with heavy quarks, incl. exotics:

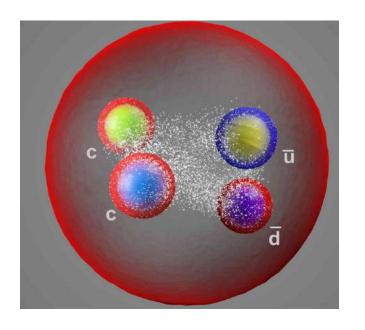
- newly discovered T_{cc}^+ tetraquark = $(cc\bar{u}\bar{d})$
- ullet stable bbar uar d tetraquark
- hadronic molecules, esp. LHCb pentaquarks ^{6 by the latest count:}
 3 nonstrange & 3 strange
- "like a new layer in the periodic table"

∃ robust experimental evidence for multiquark states, a.k.a. exotic hadrons with heavy Q

- non $\bar{q}q'$ mesons, e.g. $\bar{Q}Q\bar{q}q$, $QQ\bar{q}\bar{q}$ Q=c,b q=u,d,s
- non qq'q'' baryons, e.g. $\bar{Q}Qqq'q''$ two key questions:
- which additional exotics should we expect?
- how are quarks organized inside them?



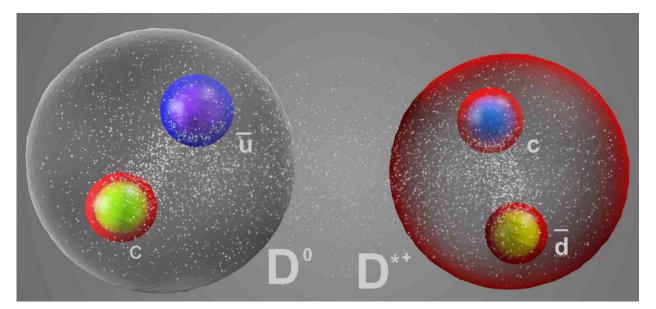
tightly-bound tetraquark



each quark sees the color charges of all other quarks

or

hadronic molecule?



two color singlets interacting by light meson x-change



CERN-EP-2021-165 LHCb-PAPER-2021-031 September 2, 2021

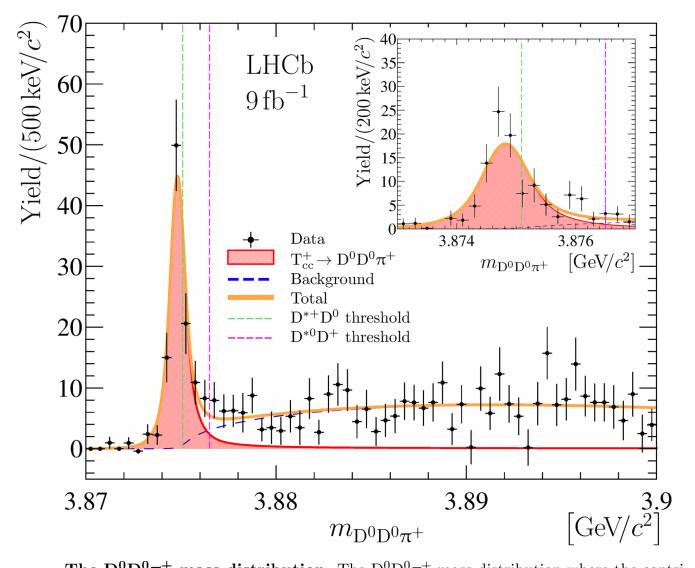
Observation of an exotic narrow doubly charmed tetraquark

LHCb collaboration[†]

Abstract

Conventional hadronic matter consists of baryons and mesons made of three quarks and quark-antiquark pairs, respectively. The observation of a new type of hadronic state, a doubly charmed tetraquark containing two charm quarks, an anti-u and an anti-d quark, is reported using data collected by the LHCb experiment at the Large Hadron Collider. This exotic state with a mass of about $3875 \,\mathrm{MeV}/c^2$ manifests itself as a narrow peak in the mass spectrum of $\mathrm{D^0D^0}\pi^+$ mesons just below the $\mathrm{D^{*+}D^0}$ mass threshold. The near-threshold mass together with a strikingly narrow width reveals the resonance nature of the state.

2 0 <u>o</u> 2 Sep 2021 \mathfrak{C} [hep-ex] arXiv:2109.01038v1 Phys. Rev. Lett Commun Nature



The $D^0D^0\pi^+$ mass distribution. The $D^0D^0\pi^+$ mass distribution where the contribution of the non- D^0 background has been statistically subtracted. The result of the fit described in the text is overlaid.

Table 1: Signal yield, N, Breit-Wigner mass relative to $D^{*+}D^{0}$ mass threshold, $\delta m_{\rm BW}$, and width, $\Gamma_{\rm BW}$, obtained from the fit to the $D^{0}D^{0}\pi^{+}$ mass spectrum. The uncertainties are statistical only.

Parameter	Value		
$N \ \delta m_{ m BW} \ \Gamma_{ m BW}$	117 ± 16 $-273 \pm 61 \text{ keV}/c^2$ $410 \pm 165 \text{ keV}$	@ 4.3 σ	

$$\delta m_{\rm pole} = -360 \pm 40^{+4}_{-0} \text{ keV/}c^2,$$

$$\Gamma_{\rm pole} = 48 \pm 2^{+0}_{-14} \text{ keV},$$

$$[M(D^{*0}) + M(D^+)] - [M(D^{*+}) + M(D^0)] = 1.4 \text{ MeV} \gg \Gamma(T_{cc}^+)$$

so $T_{cc}^+ \iff D^{*+}D^0$, with very little $D^{*0}D^+$

TH predictions for T_{cc}^+ mass, I = 0, $J^P = 1^+$

Theory predictions for the mass of the ground isoscalar $J^P = 1^+$ cc $\overline{u}d$ tetraquark T_{cc}^+ state [44–75]. Masses are shown relative to the $D^{*+}D^0$ mass threshold.

important digression: X(3872) observed by CMS in heavy ion collisions

PHYSICAL REVIEW LETTERS 128, 032001 (2022)

Evidence for X(3872) in Pb-Pb Collisions and Studies of its Prompt Production at $\sqrt{s_{NN}} = 5.02 \text{ TeV}$

A. M. Sirunyan *et al.* CMS Collaboration

(Received 25 February 2021; revised 2 September 2021; accepted 22 December 2021; published 19 January 2022)

The first evidence for X(3872) production in relativistic heavy ion collisions is reported. The X(3872) production is studied in lead-lead (Pb-Pb) collisions at a center-of-mass energy of $\sqrt{s_{\rm NN}}=5.02$ TeV per nucleon pair, using the decay chain X(3872) $\rightarrow J/\psi\pi^+\pi^- \rightarrow \mu^+\mu^-\pi^+\pi^-$. The data were recorded with the CMS detector in 2018 and correspond to an integrated luminosity of 1.7 nb⁻¹. The measurement is performed in the rapidity and transverse momentum ranges |y|<1.6 and $15< p_{\rm T}<50$ GeV/c. The significance of the inclusive X(3872) signal is 4.2 standard deviations. The prompt X(3872) to ψ 2S yield ratio is found to be $\rho^{\rm Pb-Pb}=1.08\pm0.49({\rm stat})\pm0.52({\rm syst})$, to be compared with typical values of 0.1 for pp collisions. This result provides a unique experimental input to theoretical models of the X(3872) production mechanism, and of the nature of this exotic state.

DOI: 10.1103/PhysRevLett.128.032001

The production cross section is much larger, due to multi-parton events, but the huge combinatorial background has been a major challenge. This is a proof that this challenge can be dealt with, at least in some cases,

Prompt production of X(3872) in Pb-Pb collisions. \Longrightarrow what about T_{cc}^+ ?

hadrons w. heavy quarks are much simpler:

heavy quarks almost static

ullet smaller spin-dep. interaction $\propto 1/m_Q$

key to accurate prediction of b quark baryons

- Phenomenological approach
- Identify eff. d.o.f. & their interactions
- Extract model parameters from exp
- Then use them to make predictions

apply the toolbox to

doubly-heavy baryons , e.g. ccu and doubly-heavy tetraquarks, e.g. $cc\bar{u}\bar{d}$ in both heavy cc diquark 3_c^* coupled to a light 3_c doubly-heavy baryons non-exotic, must exist

⇒ excellent testing ground for the toolbox MK & JR, PRD 90, 094007(2014)

(Q) (Q) = QQQ'

doubly heavy baryons: mass predictions

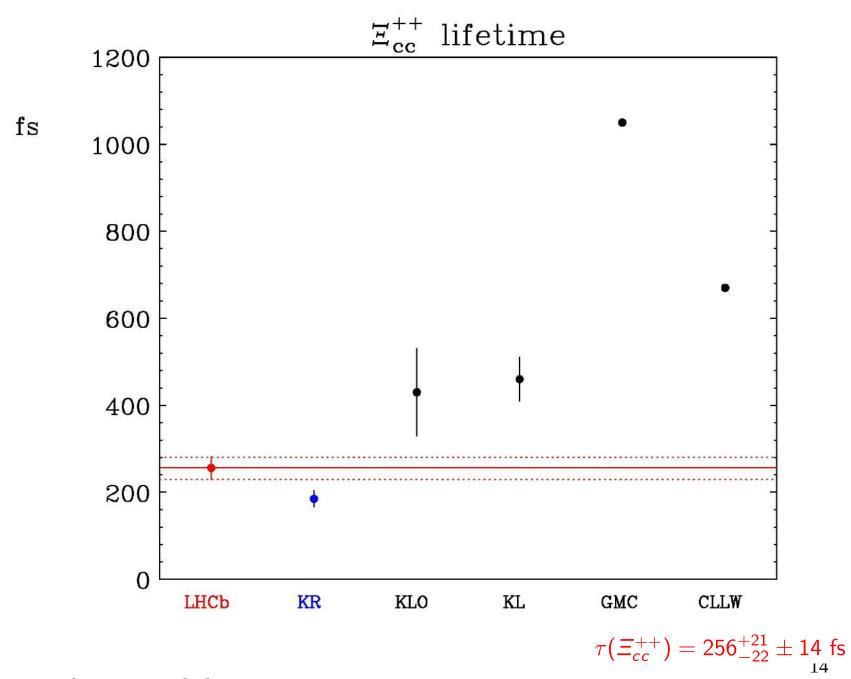
MK & JR, Phys. Rev. D90, 094007 (2014)

TABLE XVIII. Summary of our mass predictions (in MeV) for lowest-lying baryons with two heavy quarks. States without a star have J = 1/2; states with a star are their J = 3/2 hyperfine partners. The quark q can be either u or d. The square or curved brackets around cq denote coupling to spin 0 or 1.

State	Quark content	M(J=1/2)	M(J=3/2)
$\Xi_{cc}^{(*)}$	ccq	3627 ± 12	3690 ± 12
$\Xi_{bc}^{(*)}$	b[cq]	6914 ± 13	6969 ± 14
Ξ'_{bc}	b(cq)	6933 ± 12	* * *
$\Xi'_{bc} \ \Xi_{bb}^{(*)}$	bbq	10162 ± 12	10184 ± 12

LHCb: 3621.6 ± 0.4

PRL 119,112001, (2017)



M. Karliner Exotic hadrons.

Blois 23/10/2024

ccq mass calculation

sum of:

- \bullet 2 m_c
- *V_{cc}* in 3^{*}_c
- V_{HF}(cc)
- $V_{HF}(cq)$
- \bullet m_q

ccq mass calculation

sum of:

- \bullet 2 m_c

- $V_{HF}(cq)$
- \mathbf{m}_{q}

Effective masses

in mesons:

$$m_u^m = m_d^m = m_q^m = 310$$
 MeV, $m_c^m = 1663.3$ MeV

in baryons:

$$m_u^b = m_d^{\ b} = m_q^{\ b} = 363 \ {
m MeV}, \ m_c^{\ b} = 1710.5 \ {
m MeV}$$

V(cc) from $V(c\bar{c})$:

$$\bar{M}(c\bar{c}:1S) \equiv [3M(J/\psi) + M(\eta_c)]/4 = 3068.6 \text{ MeV}$$

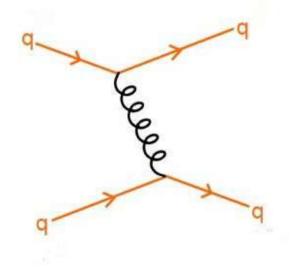
$$V(c\bar{c}) = \bar{M}(c\bar{c}:1S) - 2m_c^m = -258.0 \text{ MeV}.$$

$$V(cc) = \frac{1}{2}V(c\bar{c}) = -129.0 \text{ MeV}.$$

in weak coupling follows from color algebra in 1gx

here a *dynamical assumption*:

V(cc) and $V(c\bar{c})$ factorize into color×space



gluon exchange by 2 quarks

$V_{HF}(cc)$ from $V_{HF}(c\bar{c})$:

$$V_{HF}(cc) = \frac{a_{cc}}{m_c^2}$$

$$V_{HF}(c\bar{c}) = M(J/\psi) - M(\eta_c) = 113.2 \text{ MeV} = \frac{4a_{\bar{c}c}}{m_c^2}$$

assume
$$a_{cc} = \frac{1}{2}a_{c\bar{c}}$$
,

$$\Rightarrow \frac{a_{cc}}{m_c^2} = 1/2 \cdot \frac{M(J/\psi) - M(\eta_c)}{4} = 14.2 \text{ MeV}$$

Contributions to Ξ_{cc} mass

Contribution	Value (MeV)
$2m_{c}^{b} + m_{q}^{b}$	3783.9
cc binding	-129.0
$a_{cc}/(m_c^b)^2$	14.2
$-4a/m_q^b m_c^b$	-42.4
Totaĺ	3627 ± 12

The ± 12 MeV error estimate from ave. error for Qqq baryons

The same theoretical toolbox that led to the accurate Ξ_{cc} mass prediction now predicts

a stable, deeply bound $bb\bar{u}\bar{d}$ tetraquark,

215 MeV below BB^* threshold the first manifestly exotic stable hadron

Editors' Suggestion

PRL 119, 202001 (2017)

PHYSICAL REVIEW LETTERS

week ending 17 NOVEMBER 2017



Discovery of the Doubly Charmed Ξ_{cc} Baryon Implies a Stable $bb\bar{u}\bar{d}$ Tetraquark

Marek Karliner^{1,*} and Jonathan L. Rosner^{2,†}

¹School of Physics and Astronomy, Raymond and Beverly Sackler Faculty of Exact Sciences, Tel Aviv University, Tel Aviv 69978, Israel ²Enrico Fermi Institute and Department of Physics, University of Chicago, 5620 South Ellis Avenue, Chicago, Illinois 60637, USA (Received 28 July 2017; published 15 November 2017)

Recently, the LHCb Collaboration discovered the first doubly charmed baryon $\Xi_{cc}^{++} = ccu$ at 3621.40 ± 0.78 MeV, very close to our theoretical prediction. We use the same methods to predict a doubly bottom tetraquark $T(bb\bar{u}\bar{d})$ with $J^P = 1^+$ at 10389 ± 12 MeV, 215 MeV below the $B^-\bar{B}^{*0}$ threshold and 170 MeV below the threshold for decay to $B^-\bar{B}^0\gamma$. The $T(bb\bar{u}\bar{d})$ is therefore stable under strong and electromagnetic interactions and can only decay weakly, the first exotic hadron with such a property. On the other hand, the mass of $T(cc\bar{u}\bar{d})$ with $J^P = 1^+$ is predicted to be 3882 ± 12 MeV, 7 MeV above the D^0D^{*+} threshold and 148 MeV above the $D^0D^+\gamma$ threshold. $T(bc\bar{u}\bar{d})$ with $J^P = 0^+$ is predicted at 7134 \pm 13 MeV, 11 MeV below the \bar{B}^0D^0 threshold. Our precision is not sufficient to determine whether $bc\bar{u}\bar{d}$ is actually above or below the threshold. It could manifest itself as a narrow resonance just at threshold.

DOI: 10.1103/PhysRevLett.119.202001

Calculation of tetraquark bbūd mass

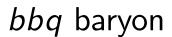
build on accuracy of the Ξ_{cc} mass prediction

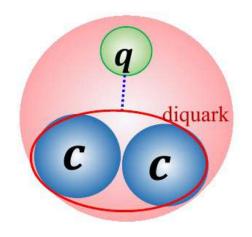
$$V(bb) = \frac{1}{2}V(\bar{b}b)$$

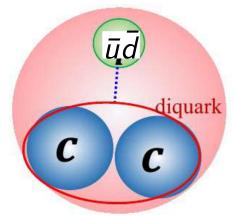
to obtain lowest possible mass, assume:

- bbūd̄ in S-wave
- $\bar{u}\bar{d}$: $\mathbf{3_c}$ "good" antidiq., S=0, I=0 (it's the lightest one)
- \Rightarrow bb must be $\overline{3}_{c}$; Fermi stats: spin 1 $(bb)_{S=1} (\bar{u}\bar{d})_{S=0} \Rightarrow J^{P} = 1^{+}$.
- \Rightarrow $(bb)(\bar{u}\bar{d})$ very similar to bbq baryon:

$$q \leftrightarrow (\bar{u}\bar{d})$$







Ξ_{cc} discovery \Rightarrow quantitative validation

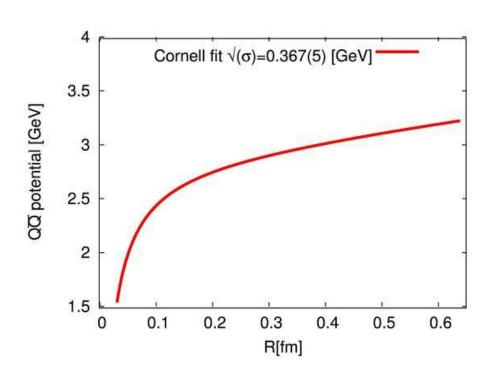
qualitatively
$$E_{binding} \sim \alpha_s^2 M_Q$$

so for
$$M_Q o \infty$$

 $QQ\bar{u}\bar{d}$ must be bound

Contributions to mass of $(bb\bar u\bar d)$ Tq with $J^P=1^+$

Contribution	Value (MeV)
$2m_b^b$	10087.0
$2m_q^b$	726.0
$a_{bb}/(m_{b}^{b})^{2}$	7.8
$-3a/(m_{q}^{b})^{2}$	-150.0
bb binding	-281.4
Total	10389.4 ± 12

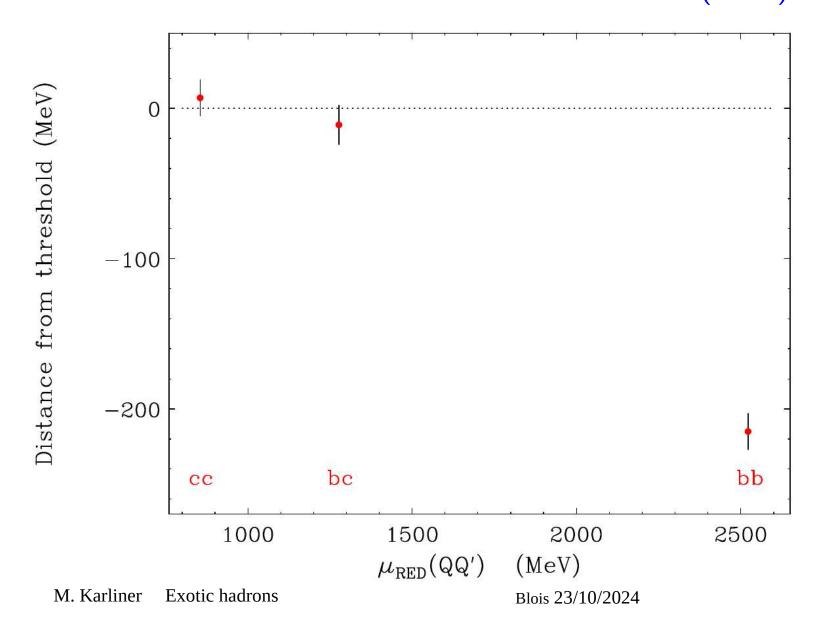


$$T(bb\bar{u}\bar{d})$$
:
 $m_b \approx 5 \text{ GeV}$
 $\Rightarrow R(bb) \sim 0.2 \text{ fm}$
 $V(r) = -\frac{\alpha_s(r)}{r} + \sigma r$
 $\Rightarrow B(bb) \approx -280 \text{ MeV}$
tightly bound, but $\bar{3}_c$,
so cannot disangage from $\bar{u}\bar{d}$

The channel $T_{bb} \to BB^*$ is kinematically closed because in BB^* the two b quarks are far from each other and the v. large bb binding energy is lost

$\Rightarrow T_{bb}$ is stable against strong decay

Distance of the $QQ'\bar{u}\bar{d}$ Tq masses from the relevant two-meson thresholds (MeV).



Tetraquark production

$$\sigma(pp \to T(bb\bar{u}\bar{d}) + X \lesssim \sigma(pp \to \Xi_{bb} + X)$$

same bottleneck: $\sigma(pp \to \{bb\} + X)$

hadronization:

$$\{bb\}
ightarrow \{bb\}q
ightarrow \{bb\} ar{u}ar{d}
ight\} P(ar{u}ar{d}) \lesssim P(q)$$

 $\{bb\}
ightarrow \{bb\}ar{u}ar{d}
ight\} ar{\mathbf{3}}_c \qquad \mathbf{3}_c$

LHCb observed $ccu = \Xi_{cc}^{++}$

$$\sigma(pp \to \Xi_{bb} + X) = (b/c)^2 \cdot \sigma(pp \to \Xi_{cc} + X)$$

 $\Rightarrow \Xi_{bb}$ and $T(bb\bar{u}\bar{d})$ accessible, near thr. \rightarrow v. narrow with much more $\int \mathcal{L} dt$

$$T(cc\bar{u}d)$$

near thr. \rightarrow v. narrow
accessible
now: D^0D^{*+} , etc.

Inclusive signature of either bbq or $bb\bar{q}\bar{q}$: displaced B_c

T. Gershon & A. Poluektov JHEP 1901 (2019) 019

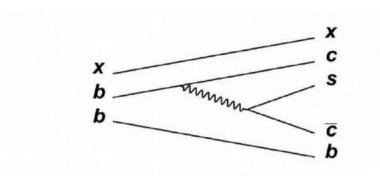
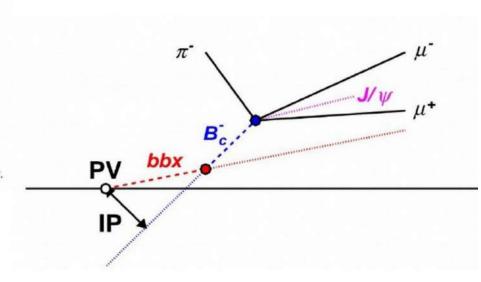


Diagram for production of a B_c^- meson from a double beauty hadron decay.



 $\mathcal{O}(1\%)$ of all B_c -s @LHC come from bbx

- major enhancement of eff. bbx rate
- bbq or bbūd?

incl. $\sigma(bbx)$: heavy ions $\gg pp$

 \Rightarrow displaced B_c @ALICE & RHIC!

crude estimate of bbūd lifetime

$$M_{initial} = M(bb\bar{u}\bar{d}) = 10,389.4 \text{ MeV}$$

$$M_{final} = M(\bar{B}) + M(D) = 7,144.5 \text{ MeV},$$

$$W^{-*}
ightarrow e \bar{\nu}_e$$
, $\mu \bar{\nu}_\mu$, $\tau \bar{\nu}_\tau$, 3 colors of $\bar{u}d$ and $\bar{c}s$,

a kinematic suppression factor

$$F(x) = 1 - 8x + 8x^3 - x^4 + 12x^2 \ln(1/x) ,$$

$$x \equiv \{ [M(\bar{B}) + M(D)] / M(bb\bar{u}\bar{d}) \}^2,$$

 $|V_{cb}| = 0.04$, factor of 2 to count each decaying b quark.

$$\Rightarrow \Gamma(bb\bar{u}\bar{d}) = \frac{18 G_F^2 M (bb\bar{u}\bar{d})^5}{192\pi^3} F(x) |V_{cb}|^2 = 17.9 \times 10^{-13} \text{ GeV} ,$$

$$\tau(bb\bar{u}\bar{d}) = 367 \text{ fs.}$$

bbūd decay channels

(a) "standard process" $bb\bar{u}\bar{d} \to cb\bar{u}\bar{d} + W^{*-}$.

$$(bbar uar d) o D^0ar B^0\pi^-$$
 , $D^+B^-\pi^-$

$$(bb\bar u\bar d) o J/\psi K^-ar B^0$$
, $J/\psi ar K^0 B^-$.

$$(bbar uar d) o\Omega_{bc}\,ar p$$
, $\Omega_{bc}\,ar\Lambda_c$, $\Xi_{bc}^0\,ar p$, $\Xi_{bc}^0\,ar\Lambda_c$

In addition, a rare process where both $b \rightarrow c\bar{c}s$,

$$(bb\bar{u}\bar{d}) \rightarrow J/\psi J/\psi K^- \bar{K}^0$$
.

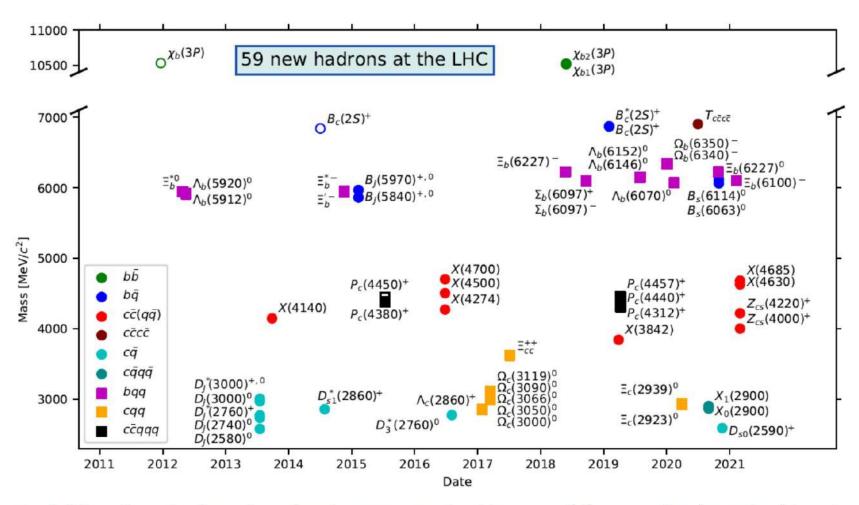
striking signature: $2J/\psi$ -s from same 2ndary vertex

(b) The W-exchange bar d o car u

e.g.
$$(bb\bar{u}\bar{d}) \rightarrow D^0B^-$$
.

$T(bb\bar{u}\bar{d})$ Summary

- stable, deeply bound $bb\bar{u}\bar{d}$ tetraquark
- $J^P = 1^+$, $M(bb\bar{u}\bar{d}) = 10389 \pm 12 \; {
 m MeV}$
- 215 MeV below BB* threshold
- first manifesty exotic stable hadron
- $(bb\bar{u}\bar{d}) o \bar{B}D\pi^-$, $J/\psi \bar{K}\bar{B}$, $J/\psi J/\psi K^- \bar{K}^0$, $D^0 B^-$
- bb $\bar{u}\bar{d}$ cousins $\begin{cases} \bullet & (bc\bar{u}\bar{d}) \colon J^P = 0^+, \text{ borderline bound} \\ 7134 \pm 13 \text{ MeV, } 11 \text{ MeV below } \bar{B}^0D^0 \\ \bullet & (cc\bar{u}\bar{d}) \colon J^P = 1^+, \text{ observed: } 3875 \text{ MeV,} \end{cases}$
 - just $\mathcal{O}(300)$ keV below D^0D^{*+} , $\Gamma\ll 1$ MeV



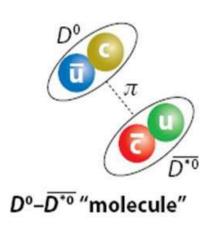
The full list of new hadrons found at the LHC, organised by year of discovery (horizontal axis) and particle mass (vertical axis). The colours and shapes denote the quark content of these states. (Image: LHCb/CERN)

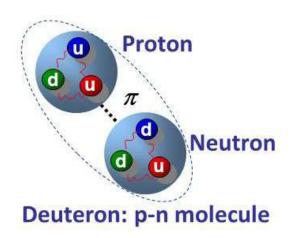
5 narrow exotic states close to meson-meson thresholds

state	mass MeV	width MeV	$ar{Q}Q$ decay	phase space	nearby threshold	ΔE MeV
			mode	MeV		
X(3872)	3872	< 1.2	$J/\psi\pi^+\pi^-$	495	$ar{D}D^*$	< 1
$Z_b(10610)$	10608	21	γ_π	1008	$ar{B}B^*$	2 ± 2
$Z_b(10650)$	10651	10	γ_π	1051	$ar{B}^*B^*$	2 ± 2
$Z_c(3900)$	3900	24 - 46	$J/\psi\pi$	663	$ar{D}D^*$	24
$Z_c(4020)$	4020	8 - 25	$J/\psi\pi$	783	$ar{D}^*D^*$	6
×					ŪD	
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- masses and widths approximate
- quarkonium decays mode listed have max phase space
- offset from threshold for orientation only, v. sensitive to exact mass

Hadronic molecules: deuteron-like

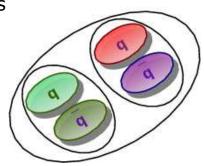




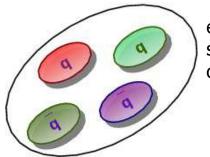
Tetraquarks: same 4 quarks, but tightly bound:

Hadronic Molecule

two color singlets attract through residual forces



Tetraquark



each quark sees color charges of all the other quarks Belle, PRL 116, 212001 (2016):

$$rac{arGamma(Z_b(10610) o ar{B}B^*)}{arGamma(Z_b(10610) o arGamma(1S)\pi)} pprox rac{86\%}{0.54\%} = \mathcal{O}(100)$$

despite 1000 MeV of phase space for $\Upsilon(1S)\pi$ vs few MeV for $\bar{B}B^*$!

overlap of Z_b wave function with $\Upsilon\pi$ dramatically smaller than with $\bar{B}B^*$

similiarly

$$\frac{\Gamma(X(3872) o \bar{D}D^*)}{\Gamma(X(3872) o J/\psi \pi^+ \pi^-)} = 9.1^{+3.4}_{-2.0}$$

$$\frac{\Gamma(Z_c(3885) o \bar{D}D^*)}{\Gamma(Z_c(3885) o J/\psi\pi)} = 6.2 \pm 1.1 \pm 2.7$$

4 pieces of experimental evidence in support of molecular interpretation of Z_Q and X(3872):

- 1. masses near thresholds and J^P of S-wave
- 2. narrow width despite very large phase space
- 3. BR(fall apart mode) \gg BR(quarkonium + X)
- 4. no states which require binding through 3 pseudoscalar coupling

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the binding mechanism can in principle

apply to any two heavy hadrons

which couple to isospin

and are heavy enough,

be they mesons or baryons

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doubly-heavy hadronic molecules: most likely candidates with $Q\bar{Q}'$, Q=c, b, $\bar{Q}'=\bar{c}$, \bar{b} :

$$D\bar{D}^*$$
, $D^*\bar{D}^*$, D^*B^* , $\bar{B}B^*$, \bar{B}^*B^* ,

$$\Sigma_c \bar{D}^*$$
, $\Sigma_c B^*$, $\Sigma_b \bar{D}^*$, $\Sigma_b B^*$, the lightest of new kind

$$\Sigma_c \bar{\Sigma}_c$$
, $\Sigma_c \bar{\Lambda}_c$, $\Sigma_c \bar{\Lambda}_b$, $\Sigma_b \bar{\Sigma}_b$, $\Sigma_b \bar{\Lambda}_b$, and $\Sigma_b \bar{\Lambda}_c$.

 $c\bar{c}$ and $b\bar{b}$ states decay strongly to $\bar{c}c$ or $\bar{b}b$ and π -(s) $b\bar{c}$ and $c\bar{b}$ states decay strongly to B_c^\pm and π -(s)

QQ' candidates – dibaryons

$$\Sigma_c \Sigma_c$$
, $\Sigma_c \Lambda_c$, $\Sigma_c \Lambda_b$, $\Sigma_b \Sigma_b$, $\Sigma_b \Lambda_b$, and $\Sigma_b \Lambda_c$.

Thresholds for $Q\bar{Q}'$ molecular states

Channel	Minimum	Minimal quark	Threshold	Example of
	isospin	content ^{a,b}	$(MeV)^c$	decay mode
$\overline{Dar{D}^*}$	0	с̄сq̄q	3875.8	$J\!/\psi\pi\pi$
$D^*ar{D}^*$	0	$car{c}qar{q}$	4017.2	$J\!/\psi\pi\pi$
D^*B^*	0	$car{b}qar{q}$	7333.8	$B_c^+\pi\pi$
$ar{\mathcal{B}}\mathcal{B}^*$	0	$bar{b}qar{q}$	10604.6	$\Upsilon(\mathit{nS})\pi\pi$
$ar{B}^*B^*$	0	$bar{b}qar{q}$	10650.4	$\Upsilon(n S)\pi\pi$
$\Sigma_car{D}^*$	1/2	c̄c̄qqq′	4462.4	$J\!/\psi$ p
$\Sigma_c B^*$	1/2	cБqqq′	7779.5	$B_c^+ p$
$\mathcal{\Sigma}_bar{\mathcal{D}}^*$	1/2	b̄cqqq′	7823.0	$B_c^- p$
$\Sigma_b B^*$	1/2	$bar{b}qqq'$	11139.6	$\Upsilon(nS)p$
$\Sigma_car{\Lambda}_c$	1	c̄cqq'ū̄d̄	4740.3	$J\!/\psi\pi$
$\Sigma_car{\Sigma_c}$	0	$car{c}qq'ar{q}ar{q}'$	4907.6	$J\!/\!\psi\pi\pi$
$\Sigma_car{\Lambda}_b$	1	cБqq′ ūd̄	8073.3^{d}	$B_c^+\pi$
$\Sigma_bar{\Lambda}_c$	1	b̄cqq'ū̄d	8100.9^{d}	$B_c^-\pi$
$\Sigma_bar{\Lambda}_b$	1	bБ̄qq′ ū̄d̄	11433.9	$\Upsilon(n{\cal S})\pi$
$\Sigma_b ar{\Sigma}_b$	0	$bar{b}qq'ar{q}ar{q}'$	11628.8	$\Upsilon(nS)\pi\pi$

^algnoring annihilation of quarks.

^cBased on isospin-averaged masses.

^bPlus other charge states when $I \neq 0$.

^dThresholds differ by 27.6 MeV.

New Exotic Meson and Baryon Resonances from Doubly Heavy Hadronic Molecules

Marek Karliner^{1,*} and Jonathan L. Rosner^{2,†}

¹School of Physics and Astronomy, Raymond and Beverly Sackler Faculty of Exact Sciences,

Tel Aviv University, Tel Aviv 69978, Israel

²Enrico Fermi Institute and Department of Physics, University of Chicago, 5620 S. Ellis Avenue,

Chicago, Illinois 60637, USA

(Received 13 July 2015; published 14 September 2015)

We predict several new exotic doubly heavy hadronic resonances, inferring from the observed exotic bottomoniumlike and charmoniumlike narrow states X(3872), $Z_b(10610)$, $Z_b(10650)$, $Z_c(3900)$, and $Z_c(4020/4025)$. We interpret the binding mechanism as mostly molecularlike isospin-exchange attraction between two heavy-light mesons in a relative S-wave state. We then generalize it to other systems containing two heavy hadrons which can couple through isospin exchange. The new predicted states include resonances in meson-meson, meson-baryon, baryon-baryon, and baryon-antibaryon channels. These include those giving rise to final states involving a heavy quark Q = c, b and antiquark $\bar{Q}' = \bar{c}$, \bar{b} , namely, $D\bar{D}^*$, $D^*\bar{D}^*$, D^*B^* , $\bar{B}B^*$, \bar{B}^*B^* , \bar{S}^*B^* , $\bar{S}^*C_c\bar{S}^*$, \bar{S}^*C

DOI: 10.1103/PhysRevLett,115.122001 PACS numbers: 14.20.Pt, 12.39.Hg, 12.39.Jh, 14.40.Rt

PRL 115, 072001 (2015)

Selected for a Viewpoint in Physics PHYSICAL REVIEW LETTERS

week ending 14 AUGUST 2015

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Observation of $J/\psi p$ Resonances Consistent with Pentaquark States in $\Lambda_b^0 \to J/\psi K^- p$ Decays

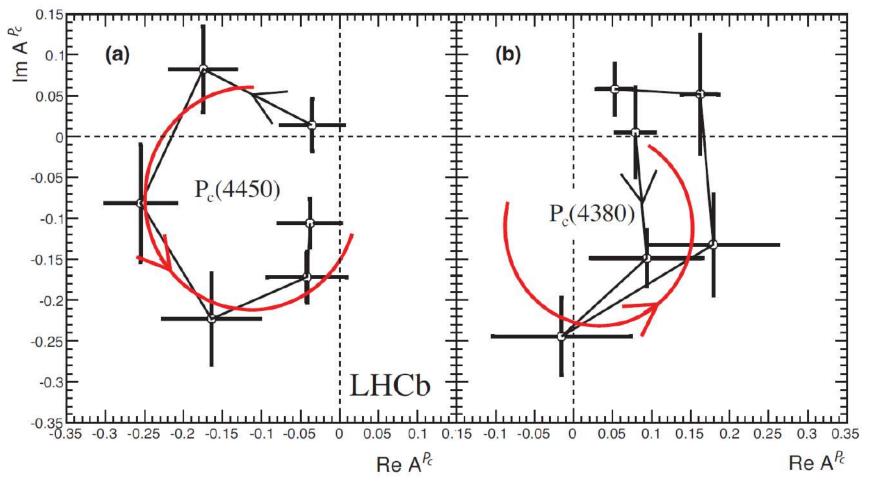
R. Aaij et al.*

(LHCb Collaboration)

(Received 13 July 2015; published 12 August 2015)

Observations of exotic structures in the $J/\psi p$ channel, which we refer to as charmonium-pentaquark states, in $\Lambda_b^0 \to J/\psi K^- p$ decays are presented. The data sample corresponds to an integrated luminosity of 3 fb⁻¹ acquired with the LHCb detector from 7 and 8 TeV pp collisions. An amplitude analysis of the three-body final state reproduces the two-body mass and angular distributions. To obtain a satisfactory fit of the structures seen in the $J/\psi p$ mass spectrum, it is necessary to include two Breit-Wigner amplitudes that each describe a resonant state. The significance of each of these resonances is more than 9 standard deviations. One has a mass of $4380 \pm 8 \pm 29$ MeV and a width of $205 \pm 18 \pm 86$ MeV, while the second is narrower, with a mass of $4449.8 \pm 1.7 \pm 2.5$ MeV and a width of $39 \pm 5 \pm 19$ MeV. The preferred J^P assignments are of opposite parity, with one state having spin 3/2 and the other 5/2.

DOI: 10.1103/PhysRevLett.115.072001 PACS numbers: 14.40.Pq, 13.25.Gv

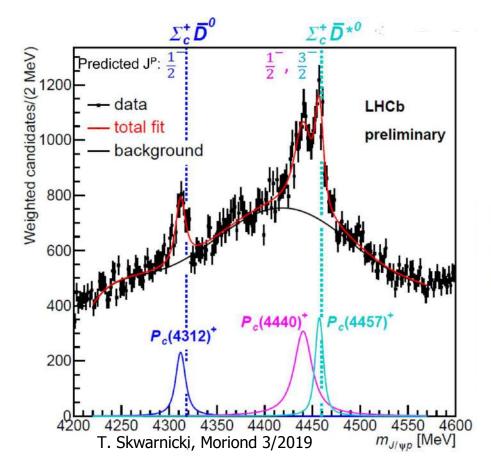


 $P_c(4450)$: predicted, narrow: $\Gamma=39\pm5\pm19$, 10 MeV from $\Sigma_c\bar{D}^*$ threshold perfect Argand plot: a molecule

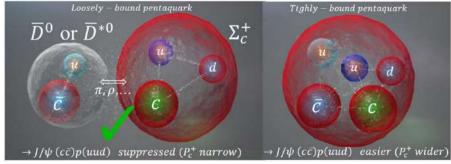
 $P_c(4380)$: not predicted, wide: $\Gamma = 205 \pm 18 \pm 86$ MeV, Argand plot not resonance-like 777

$P_c(4450)$ might be just the first of many "heavy deuterons"

as of 2015



The near-threshold masses and the narrow widths of $P_c(4312)^+$, $P_c(4440)^+$ and $P_c(4457)^+$ favor "molecular" pentaquarks with meson-baryon substructure!



observe all 3 S-wave states:

$$\Sigma_c \bar{D}; \quad J^P = \frac{1}{2}^-,$$

$$\Sigma_c \bar{D}^*$$
; $J^P = \frac{1}{2}^-$, $\frac{3}{2}^-$

M. Karliner Exotic hadrons

for $Q \to \infty$ 4 more *S*-wave states:

$$\Sigma_c^* \bar{D}; \quad J^P = \frac{3}{2}^-$$

$$\Sigma_c^* \bar{D}^*; J^P = \frac{1}{2}^-, \frac{3}{2}^-, \frac{5}{2}^-$$

but $\Gamma(\Sigma_c^* o \Lambda_c \pi) \approx 15$ MeV...

two v. different types of exotics:

 $Q\bar{Q}q\bar{q}$

 $QQ\bar{q}\bar{q}$

e.g.

 $Z_b(10610)$

 $\bar{B}B^*$

molecule

 $T(bb\bar{u}\bar{d})$

tightly-bound tetraquark

why is it so?

Exotics with $\overline{Q}Q$ vs. $\overline{Q}Q$: very different

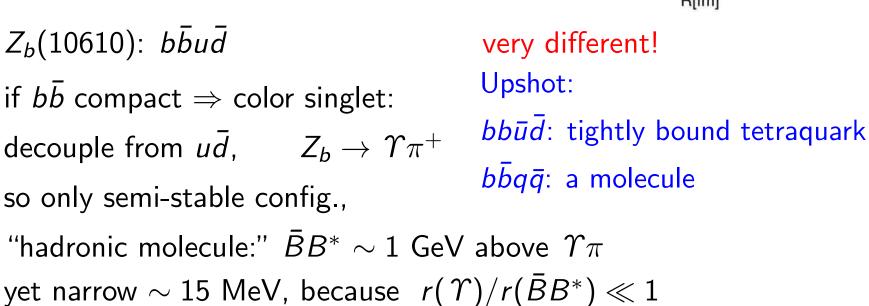
 $V(\bar{Q}Q) = 2V(QQ)$, hundreds of MeV

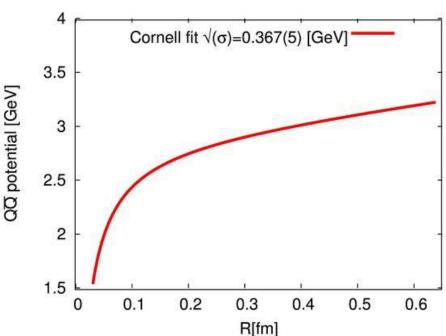
but only if $\overline{Q}Q$ color singlet

- $\Rightarrow \bar{Q}Q$ can immediately hadronize as quarkonium
- \Rightarrow exotics: \overline{Q} in one hadron and \overline{Q} in the other
- ⇒ deuteron-like "hadronic molecules"
- vs. QQ never a color singlet,
- ⇒ tightly bound exotics, tetraquarks

$$T(bb\bar{u}\bar{d})$$
:
 $m_b \approx 5 \text{ GeV}$
 $\Rightarrow R(bb) \sim 0.2 \text{ fm}$
 $V(r) = -\frac{\alpha_s(r)}{r} + \sigma r$
 $\Rightarrow B(bb) \approx -280 \text{ MeV}$
tightly bound, but $\bar{3}_c$,
so cannot disangage from $\bar{u}\bar{d}$

 $Z_b(10610)$: $b\bar{b}u\bar{d}$ if $b\bar{b}$ compact \Rightarrow color singlet: decouple from $u\bar{d}$, $Z_b \to \Upsilon \pi^+$ so only semi-stable config.,





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recent news from LHCb: new strange Pq-s

 $J/\psi \Lambda$ resonances in

$$B^- o J/\psi \Lambda \bar{p}, \quad \Xi_b^- o J/\psi \Lambda K^-$$

⇒ new "molecular" pentaquarks:

$$(car csud)pprox eta_c^0(csd)ar D^{*0}(ar cu) o J\!/\!\psi \Lambda$$

vs.
$$(car c u u d) pprox \Sigma_c^+(c u d) ar D^{*0}(ar c u) o J/\psi \, p$$

LHCb arXiv:2012.10380, Sci. Bull. 66, 1278-1287 (2021)

LHC seminar "Particle Zoo 2.0: New tetra- and pentaquarks at LHCb", July 5, 2022, https://indico.cern.ch/event/1176505/ and LHCb-PAPER-2022-031, in preparation.

$\Xi_c \bar{D}^{(*)}$ molecules $\Longrightarrow \Xi'_c \bar{D}^{(*)}$ molecules

PHYSICAL REVIEW D 106, 036024 (2022)

New strange pentaquarks

Marek Karliner^{1,*} and Jonathan L. Rosner^{2,†}

¹School of Physics and Astronomy, Raymond and Beverly Sackler Faculty of Exact Sciences,

Tel Aviv University, Tel Aviv 69978, Israel

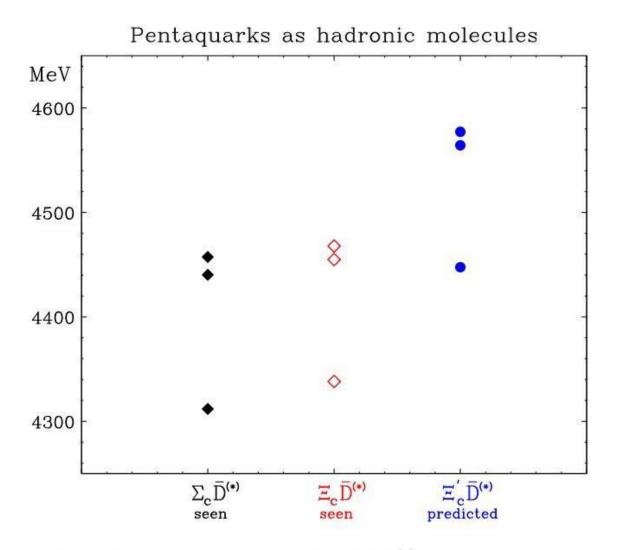
²Enrico Fermi Institute and Department of Physics, University of Chicago,

5620 South Ellis Avenue, Chicago, Illinois 60637, USA

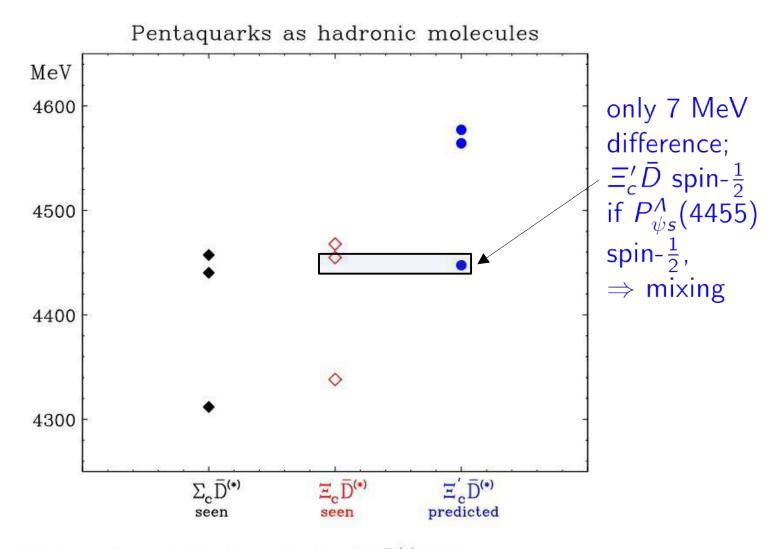
(Received 25 July 2022; accepted 8 August 2022; published 25 August 2022)

The new strange pentaquarks observed by LHCb are very likely hadronic molecules consisting of $\Xi_c\bar{D}$ and $\Xi_c\bar{D}^*$. We discuss the experimental evidence supporting this conclusion, pointing out the similarities and differences with the $P_c(4312)$, $P_c(4440)$ and $P_c(4457)$ pentaquarks in the nonstrange sector. The latter clearly are hadronic molecules consisting of $\Sigma_c\bar{D}$ and $\Sigma_c\bar{D}^*$. Following this line of thought, we predict three additional strange pentaquarks consisting of $\Xi_c'\bar{D}$ and $\Xi_c'\bar{D}^*$. The masses of these states are expected to be shifted upward by $M(\Xi_c') = M(\Xi_c) \approx 110$ MeV with respect to the corresponding known strange pentaquarks.

DOI: 10.1103/PhysRevD.106.036024



Pentaquarks as hadronic molecules. $\Sigma_c \bar{D}^{(*)}$ states are denoted by black diamonds, $\Xi_c \bar{D}^{(*)}$ states by open red diamonds and $\Xi_c' \bar{D}^{(*)}$ states by blue circles.



Pentaquarks as hadronic molecules. $\Sigma_c \bar{D}^{(*)}$ states are denoted by black diamonds, $\Xi_c \bar{D}^{(*)}$ states by open red diamonds and $\Xi_c' \bar{D}^{(*)}$ states by blue circles.

Sci.Bull. 65 (2020) 23, 1983

arXiv:2006.16957

Liupan An, \exists here at FTCF2024

LHCb, June 2020:

- a narrow resonance decaying into two J/ψ -s
- quark content cccc, confirmed by CMS
- $M \approx 6.9 \text{ GeV}$: X(6900)
- tetraquark-like
- ullet ~ 700 MeV above $J\!/\!\psi\,J\!/\!\psi$ threshold
 - \Rightarrow probably a 2*S* excited $cc\bar{c}\bar{c}$ state
- first exotic containing both QQ and $\bar{Q}Q$
- \Longrightarrow bottom analogue: $\Upsilon \Upsilon$ at $\gtrsim 19.4$ GeV ?
- stimulating challenge for EXP and TH

Interpretation of structure in $di-J/\psi$ sprectrum

- structure in LHCb di $-J/\psi$ spectrum around 6.9 and 7.2 GeV
- interpreted in terms of $J^{PC}=0^{++}$ (cc)- $(\bar{c}\bar{c})$ Tq resonances
- Tq masses from recently confirmed string-junction picture
- main peak around 6.9 GeV likely dominated by the $0^{++}(2S)$, radial exc. of (cc)- $(\bar{c}\bar{c})$ Tq, predicted at 6.871 ± 0.025 GeV
- dip around 6.75 GeV: opening of S-wave di- χ_{c0} channel
- dip around 7.2 GeV: opening of di- $\eta_c(2S)$ & $\Xi_{cc}\bar{\Xi}_{cc}$ channels?
- low-mass structure appears to require broad resonance consistent with predicted $0^{++}(1S)$ at $6191.5 \pm 25 \text{MeV}$.
- Implications for $bb\bar{b}\bar{b}$ tetraquarks

recent venture into exotics with light quarks only

Possible mixing of a diquark-antidiquark with a $p\bar{p}$ hadronic molecule

Marek Karliner a† and Jonathan L. Rosner b‡

^a School of Physics and Astronomy Tel Aviv University, Tel Aviv 69978, Israel

^b Enrico Fermi Institute and Department of Physics University of Chicago, 5620 S. Ellis Avenue, Chicago, IL 60637, USA

ABSTRACT

We discuss the possibility that the two nearby resonances observed by BESIII partially below the $p\bar{p}$ threshold might be due to mixing between two metastable states with the same $J^{PC}=0^{-+}$ quantum numbers, but rather different internal structure. One is a $p\bar{p}$ hadronic molecule and the other a bound state of a light-quark diquark and an antidiquark, both with spin 1 and isospin 0, a composite color antitriplet and triplet, respectively. The doubling of resonances, one of which may be interpreted as a hadronic molecule, while the other arises from $q\bar{q}$ annihilation in a state with vacuum quantum numbers may be a more general feature than the specific case considered here.

mixing of diquark-antidiquark with a $p\bar{p}$ H.M. ?

• BESIII: two close resonances in $J/\psi \to \gamma \, 3(\pi^+\pi^-)$

$$X(1840)$$
: $M = 1842.2 \pm 4.2^{+7.1}_{-2.6}$ MeV, $\Gamma = 83 \pm 14 \pm 11$ MeV, $J^{PC} = 0^{-+}$

$$X(1880)$$
: $M=1882.1\pm1.7\pm0.7$ MeV, $\Gamma=30.7\pm5.5\pm2.4$ MeV, $J^{PC}=0^{-+}$ width includes ~ 10 MeV below $p\bar{p}$ threshold if $p\bar{p}$

- ullet postulate: 2nd state a P-wave narrow diquark-antidquark state, both S=1
- $p\bar{p} \Longleftrightarrow \{qq\}_{S=1}$ - $\{\bar{q}\bar{q}\}_{S=1}$ via $q\bar{q}$ annihilation in 3P_0 flavor singlet vacuum q. numbers
- $M(\{qq\}_{S=1} \{\bar{q}\bar{q}\}_{S=1}) = 1886$ MeV, (quark model estimate)
- possible analogues:
 - $f_0(980)$: $K\bar{K}$ molecule and mixture of $(u\bar{u}+d\bar{d})/\sqrt{2}$ and $s\bar{s}$ in 3P_0
 - $-D_{s0}(2317)$: $c\bar{s}$ in 3P_0 ; expect DK^+ molecule near 2359 MeV threshold
 - $-\chi_{c1}(3872)$: $\bar{D}D^*$ with admixture of $c\bar{c}$ in 3P_1

SUMMARY

- narrow $cc\bar{u}\bar{d}$ tetraquark discovered by LHCb
- doubly charmed baryon found exactly where predicted $\equiv_{cc}^{++}(ccu) \Rightarrow (bcq), (bbq)$
- stable $bb\bar{u}\bar{d}$ (?also $bc\bar{u}\bar{d}$) tetraquark: LHCb!
- narrow exotics with $Q\bar{Q}$: "heavy deuterons"/hadronic molecules $\bar{D}D^*$, \bar{D}^*D^* , $\bar{B}B^*$, \bar{B}^*B^* , $\Sigma_c\bar{D}^*(S=\frac{1}{2},\frac{3}{2})$, $\Sigma_c\bar{D}(S=\frac{1}{2})$; $\gamma p \to J/\psi p$? $\Xi_c\bar{D}^{(*)}$ seen, expect 3 additional $\Xi_c'\bar{D}^{(*)}$ states $\Sigma_c B^*$, $\Sigma_b \bar{D}^*$, $\Sigma_b B^*$, D^*B^* , ...
- D^+K^- res. $\Leftrightarrow cs\bar{u}\bar{d}$ Tq w. string junction \mathcal{C} ; $bs\bar{u}\bar{d}=\bar{B}^0K^-$?
- $J/\psi J/\psi$ res. \Leftrightarrow excited $cc\bar{c}\bar{c}$ Tq, probably 2S, $J/\psi \Upsilon$, $\Upsilon \Upsilon$?
- light quark sector: possible mixture of $p\bar{p}$ and $\{qq\}_{S=1}$ - $\{\bar{q}\bar{q}\}_{S=1}$, etc.

both kinds of exotic hadrons exist in Nature exciting new spectroscopy awaiting discovery