

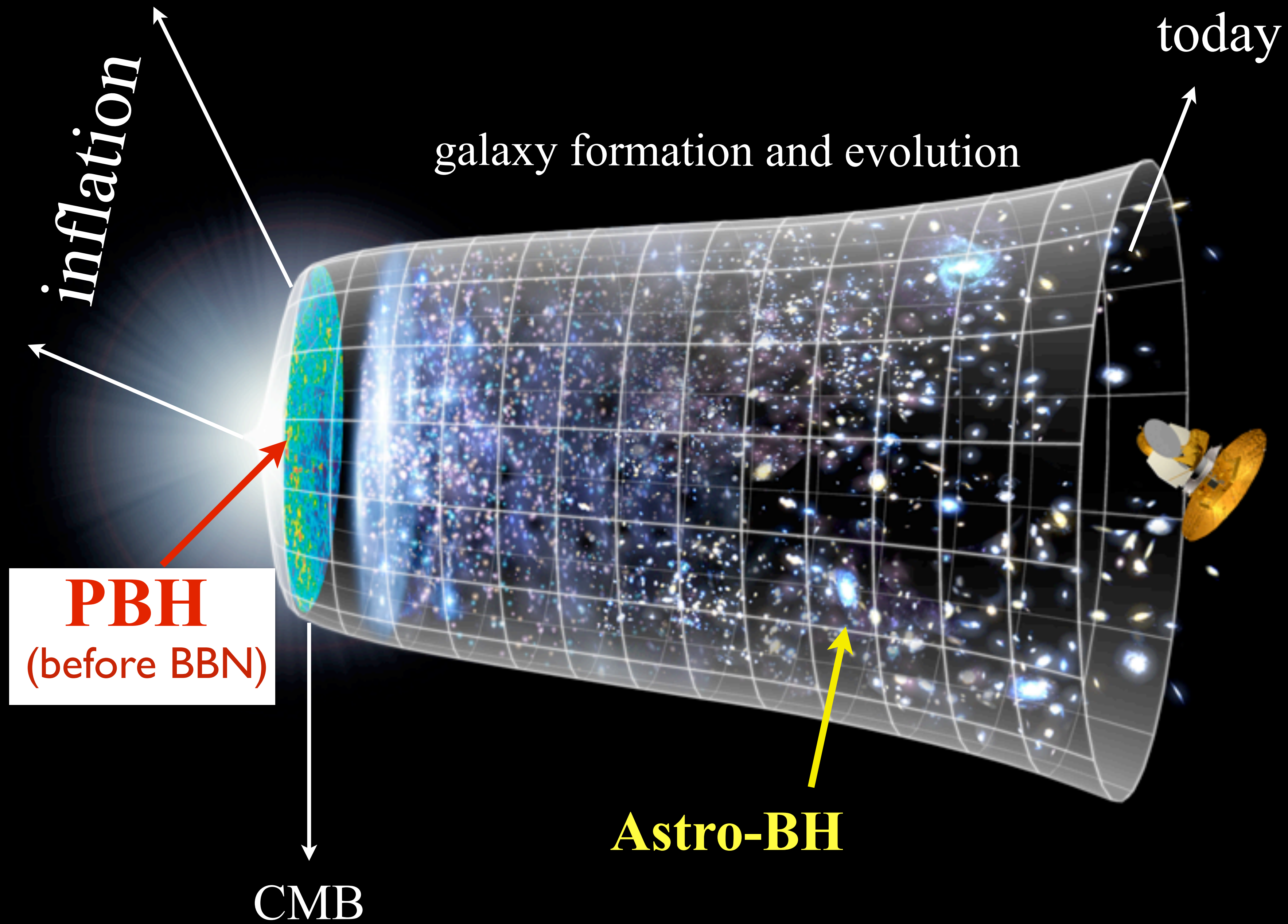
Primordial Black Holes

Guillermo Ballesteros



25/10/2024

35th Rencontres de Blois on Particle Physics and Cosmology



PBH
(before BBN)

CMB

Astro-BH

Primordial black holes

Zeldovich and Novikov 1967, Hawking 1971, Hawking & Carr 1974, ...

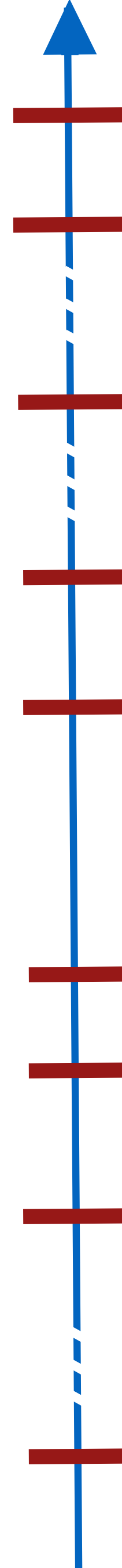
- **Dark matter candidate**
- Reheat the universe after inflation
- Catalysers for particle dark matter or other BSM production
- Baryogenesis
- Seeds of supermassive black holes
- Possible BH merger outliers in LIGO/Virgo/Kagra
- Possible BH microlensing events
- PTA indications of a stochastic GW background
- Window into the early universe

Dark Matters

(pre-BBN) PBHs



$6 \cdot 10^{38}$ GeV



$1 M_{\odot} = 2 \times 10^{30}$ kg

$5 \times 10^{-19} M_{\odot} \sim$ evaporation limit

WIMPS



1 TeV

1 GeV

$10^{-54} M_{\odot}$

WDM

1 keV

1 eV

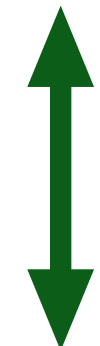
$10^{-68} M_{\odot}$

Axion and Co.



1 meV

Fuzzy DM



10^{-22} eV

$10^{-90} M_{\odot}$

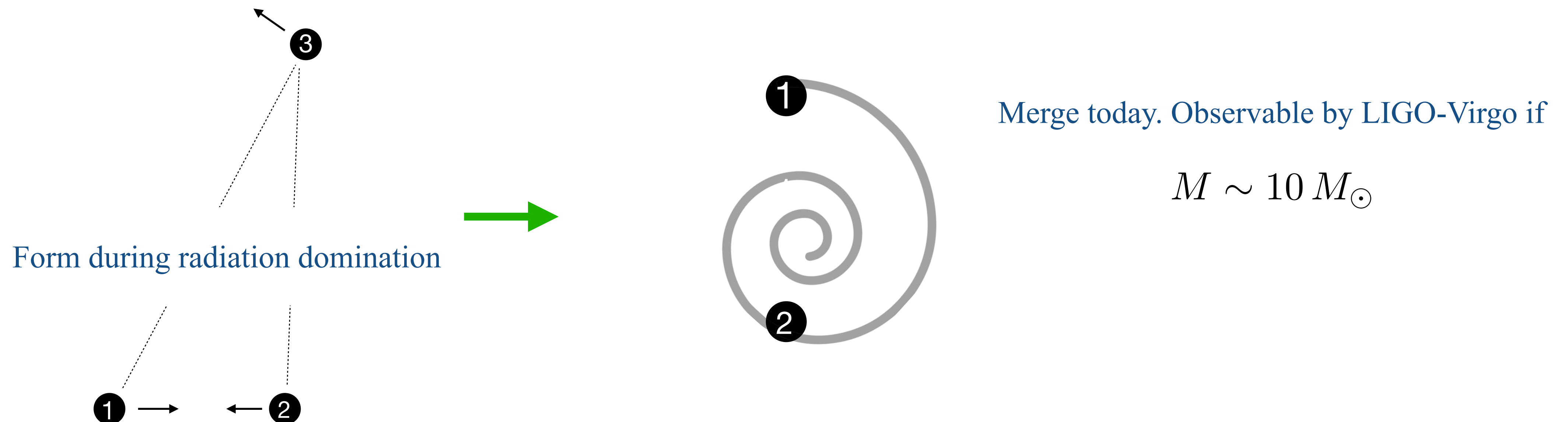
PBHs and other dark matter candidates

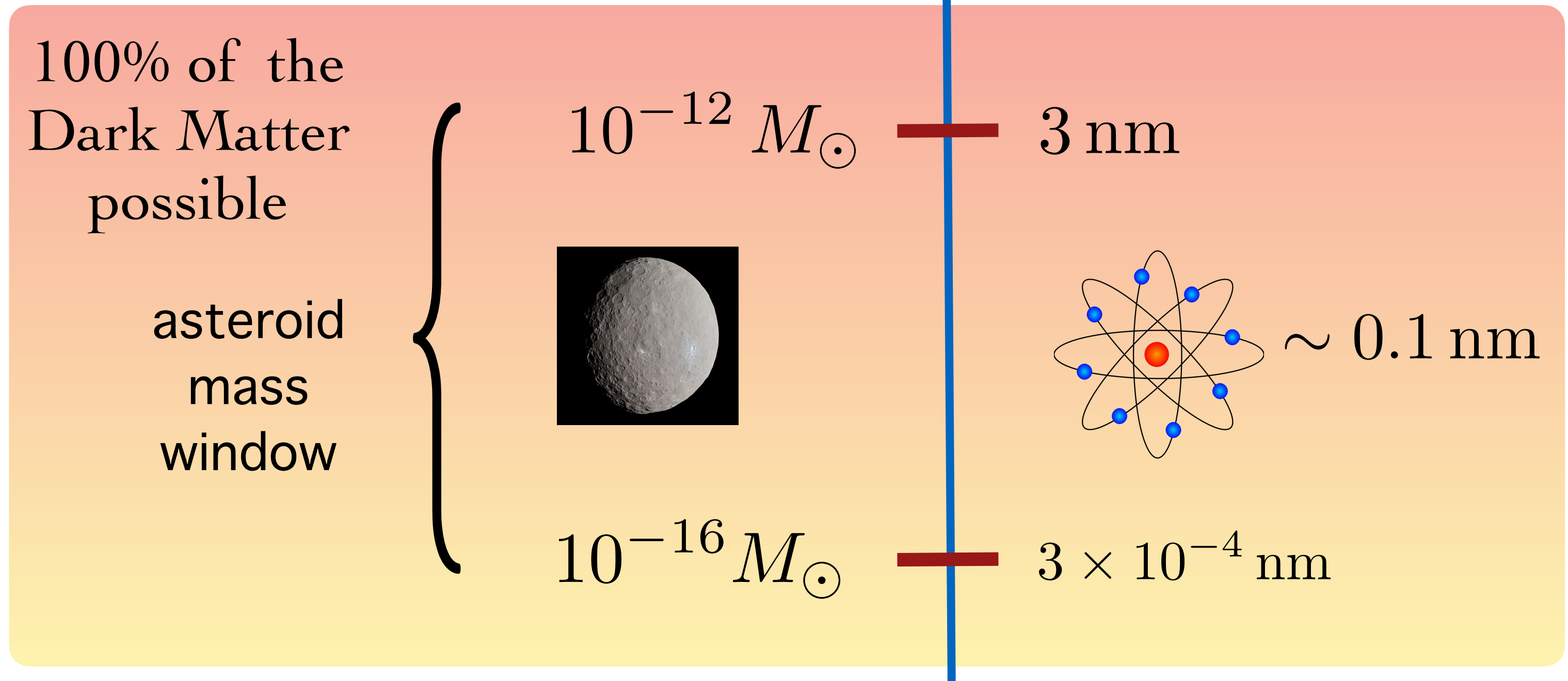
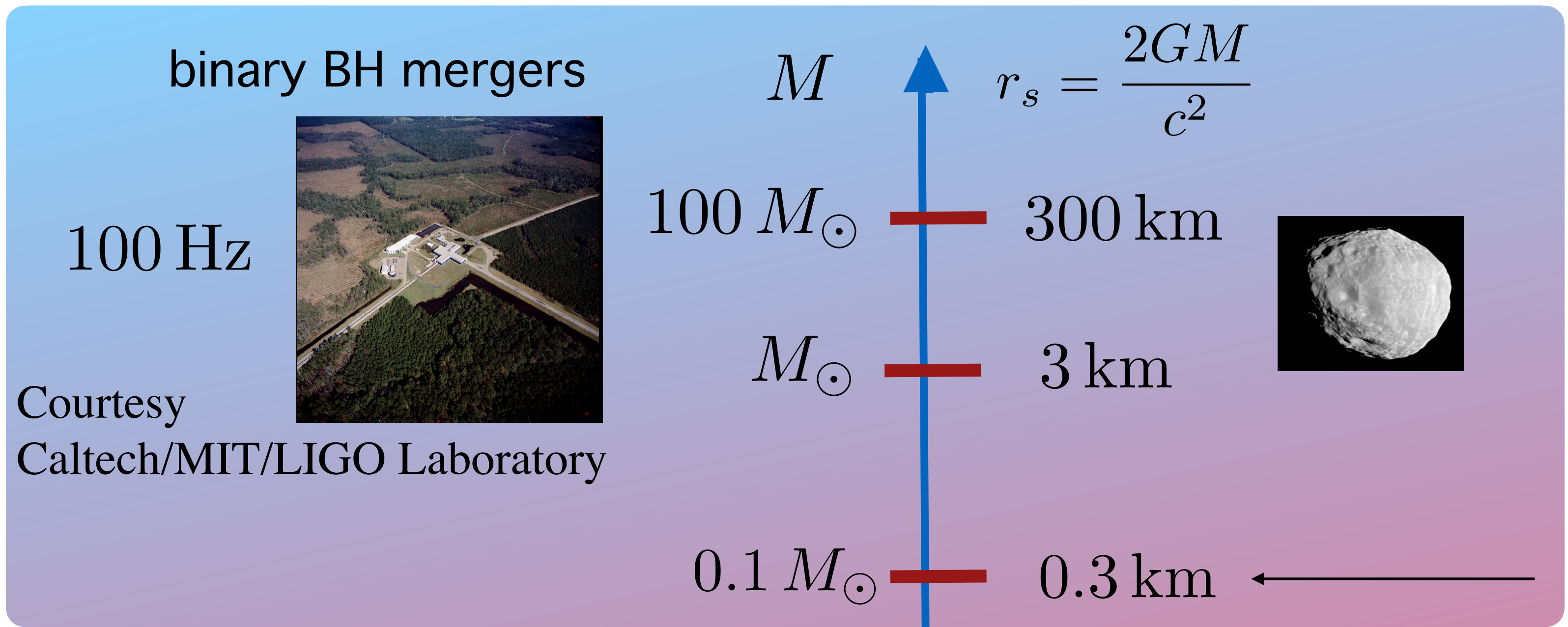
- * WIMPs: possibly largely incompatible (absence of annihilation signals)

(e.g. Lacki+ 2018, Adamek+2019, Gaggero+2019)

- * axions: no problem

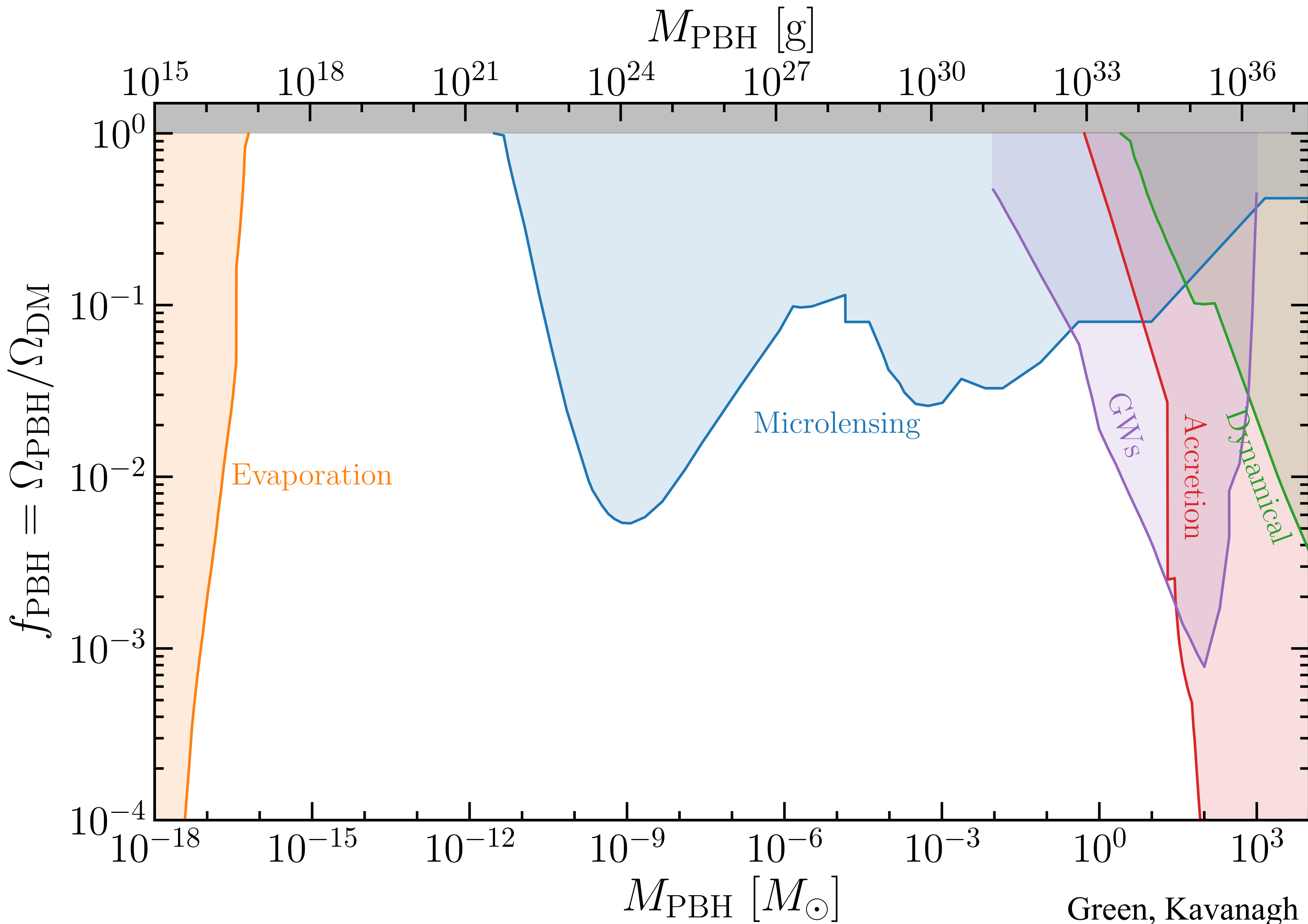
- 2016: LIGO: merger with **~30 Solar masses BH**
- Did **LIGO** detect (**THE**) dark matter? (Bird et al 2016...)
- Most likely **NO**. Or at least, most likely not the main component of it.
- Two mechanisms to form PBHs binaries:
 - Late Universe (in dark matter halos)
 - Early Universe (before matter radiation equality). **Nakamura et al 1997**
- **MAX 0.1% – 1%** of the DM in the range **1-100 Solar Masses**





Possible Stochastic GW backgrounds from inflation

0.03 Hz – 3 Hz
e.g. LISA



Evaporation by Hawking radiation

$$T = \frac{\hbar c^2}{8\pi G k_B M} = 6 \times 10^{-8} \frac{M_\odot}{M} K$$

Bound for $M \lesssim 10^{17} g \simeq 5 \times 10^{-17} M_\odot$

Carr+ 2009

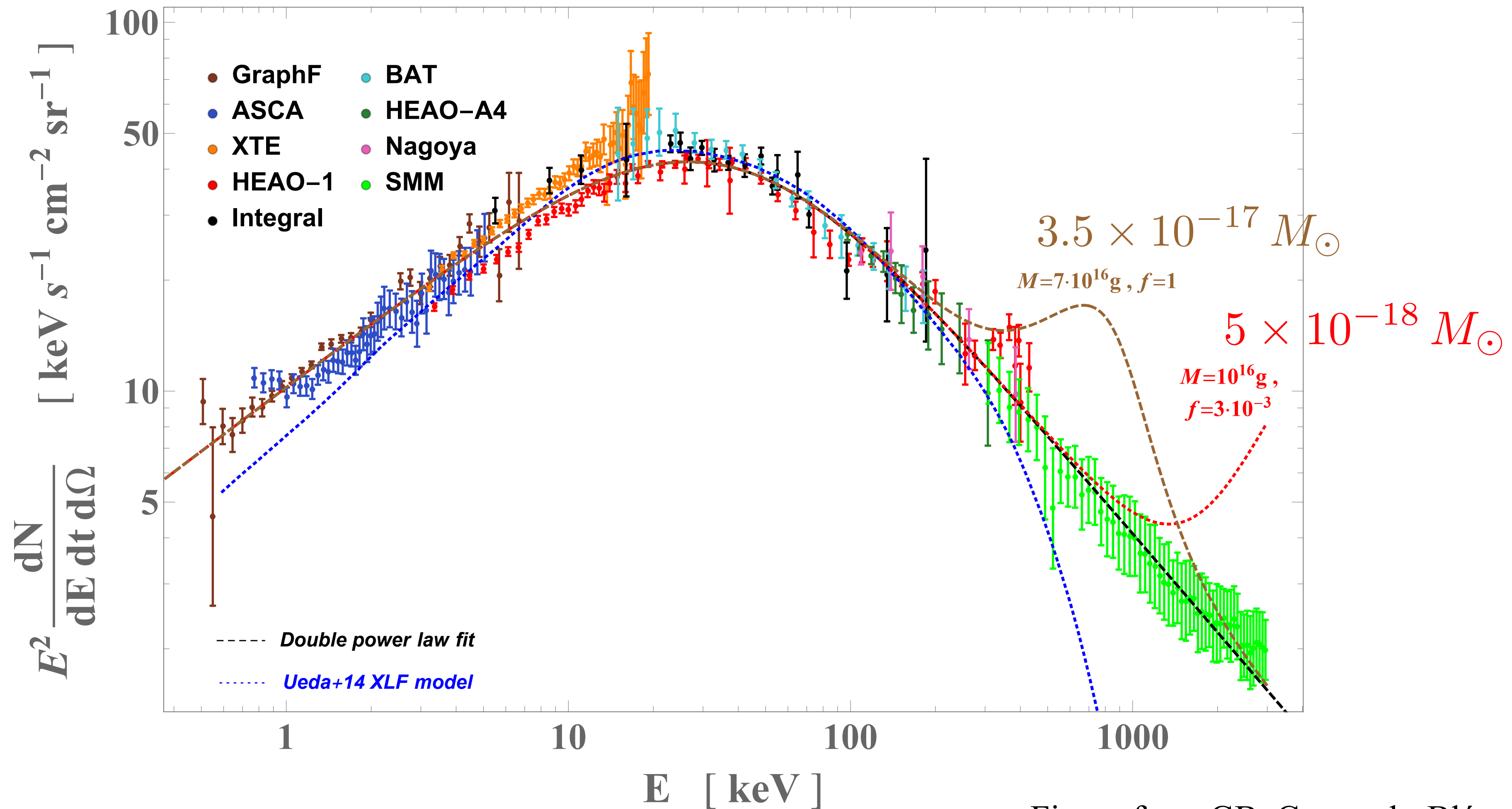


Figure from GB, Coronado-Blázquez, Gaggero, 2019

Situation before ~2019

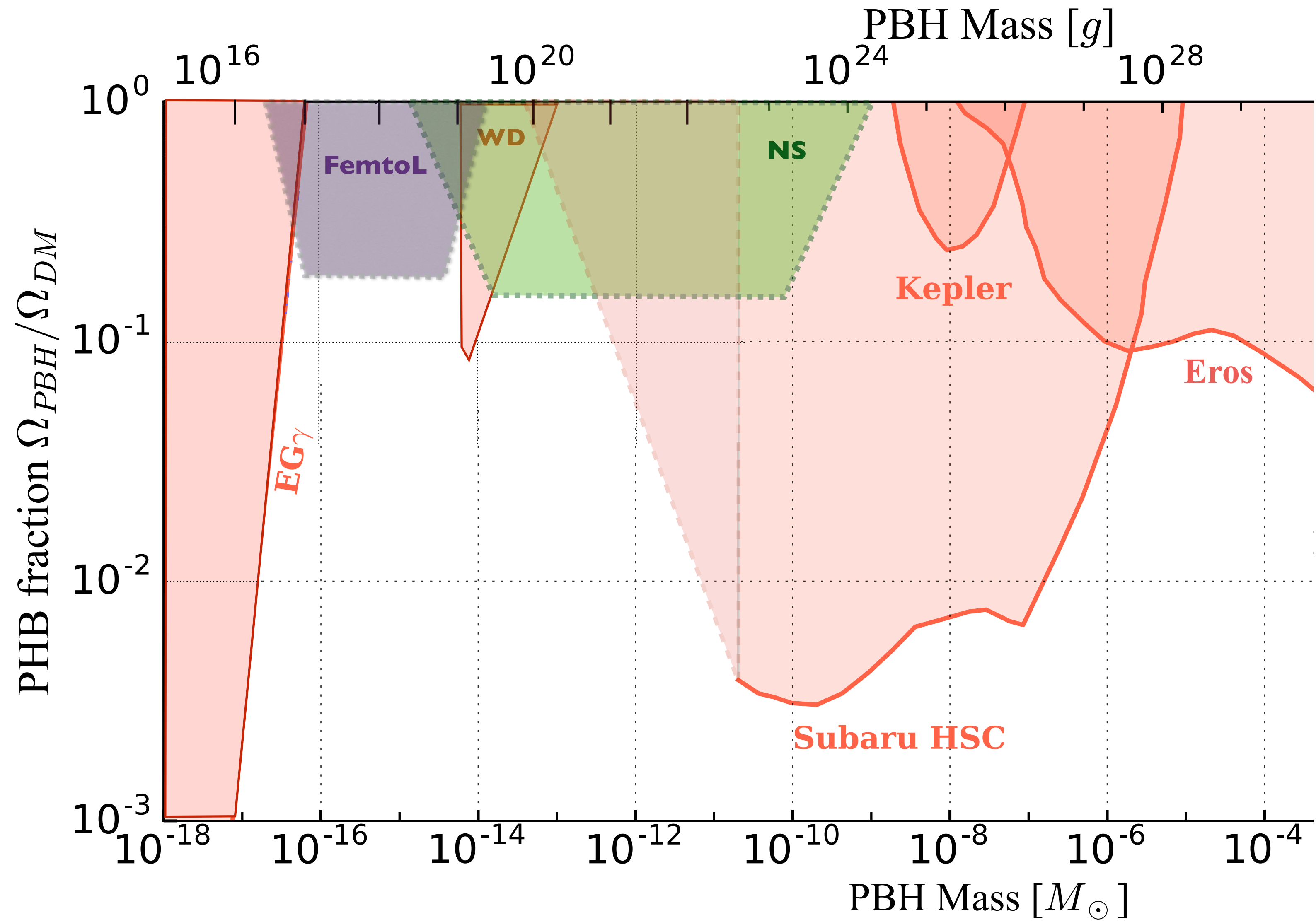


Figure from Katz et al. 2018
(modified)

See also Montero-Camacho et al. 2019
Smyth et al. 2019

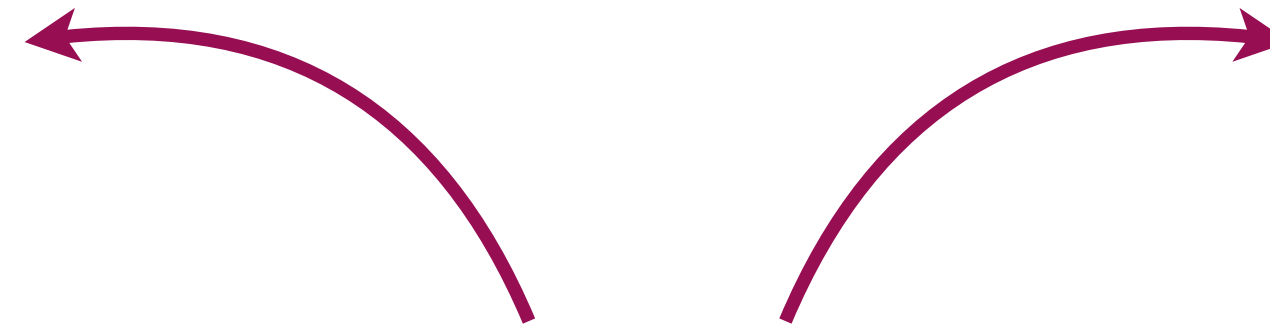
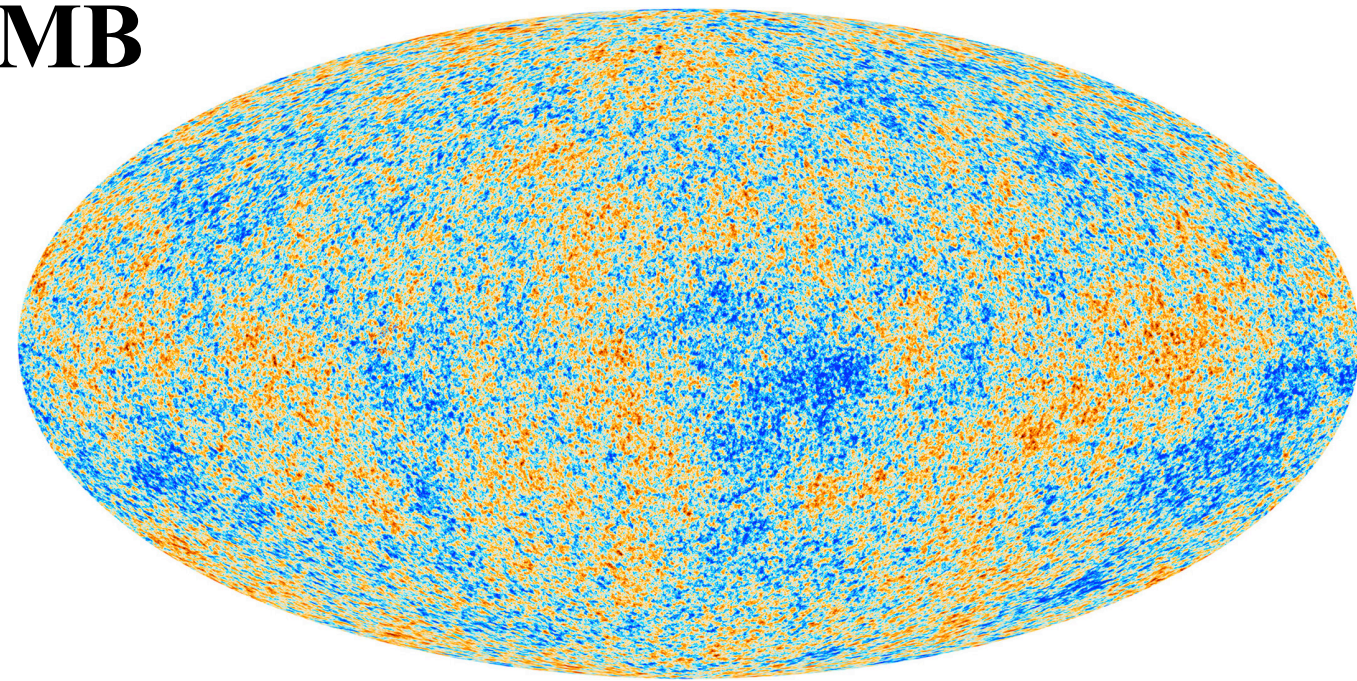
Mechanisms for PBH formation

- Large primordial fluctuations (inflation)
- Reheating after inflation
- 1st order phase transitions
- Collapse of topological defects
- False vacuum decay
- ...

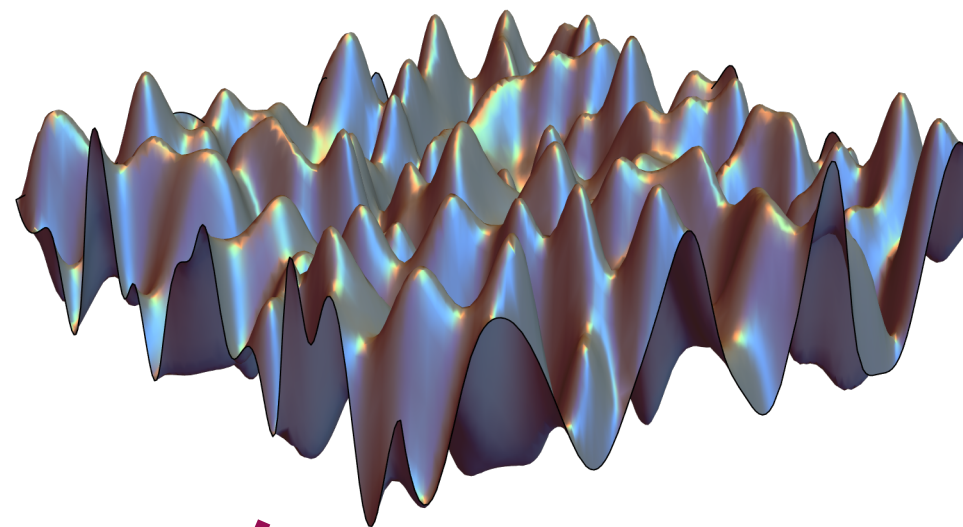
PBH dark matter (from inflation)

Credit: ESA/Planck

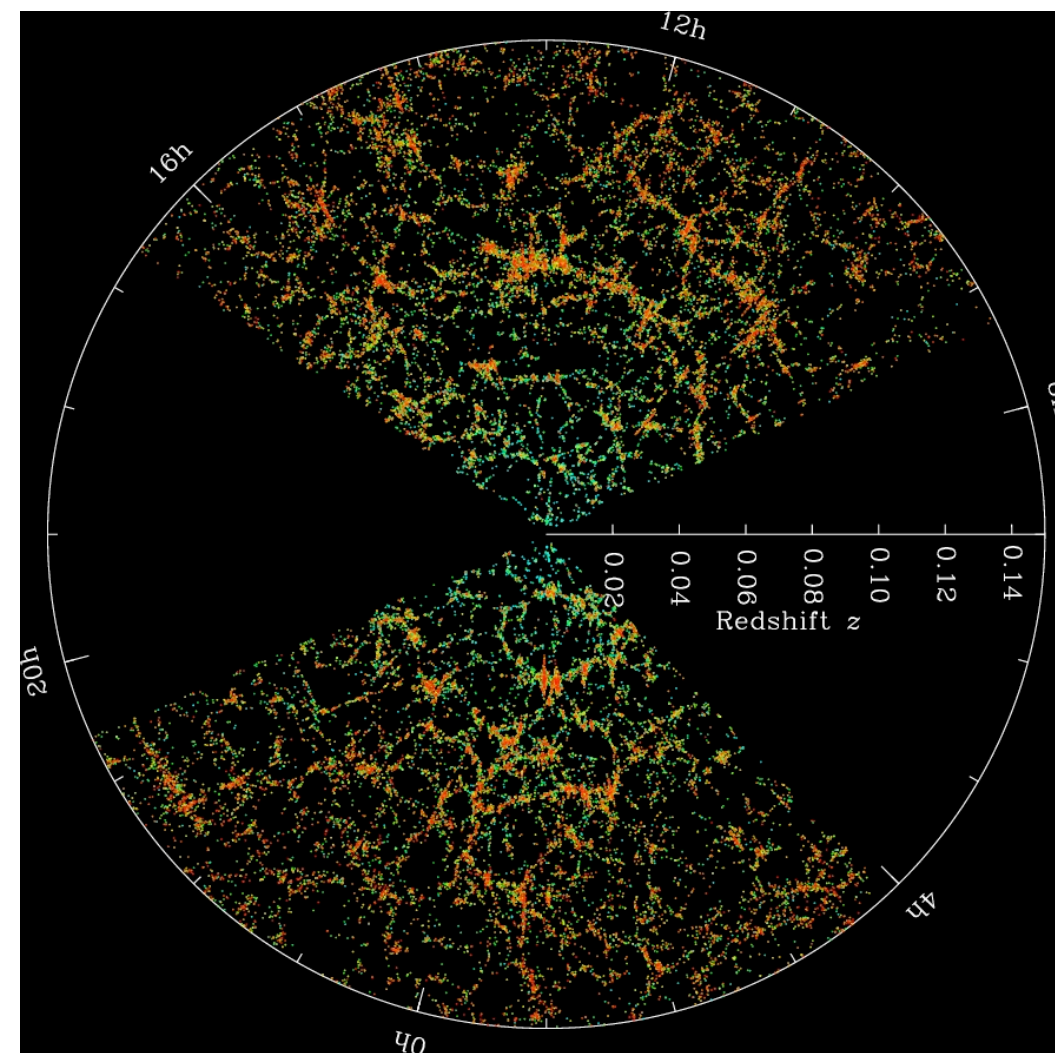
CMB



Quantum fluctuations
of space-time



LSS



Credit: M Blanton
and SDSS

Milky Way halo structure

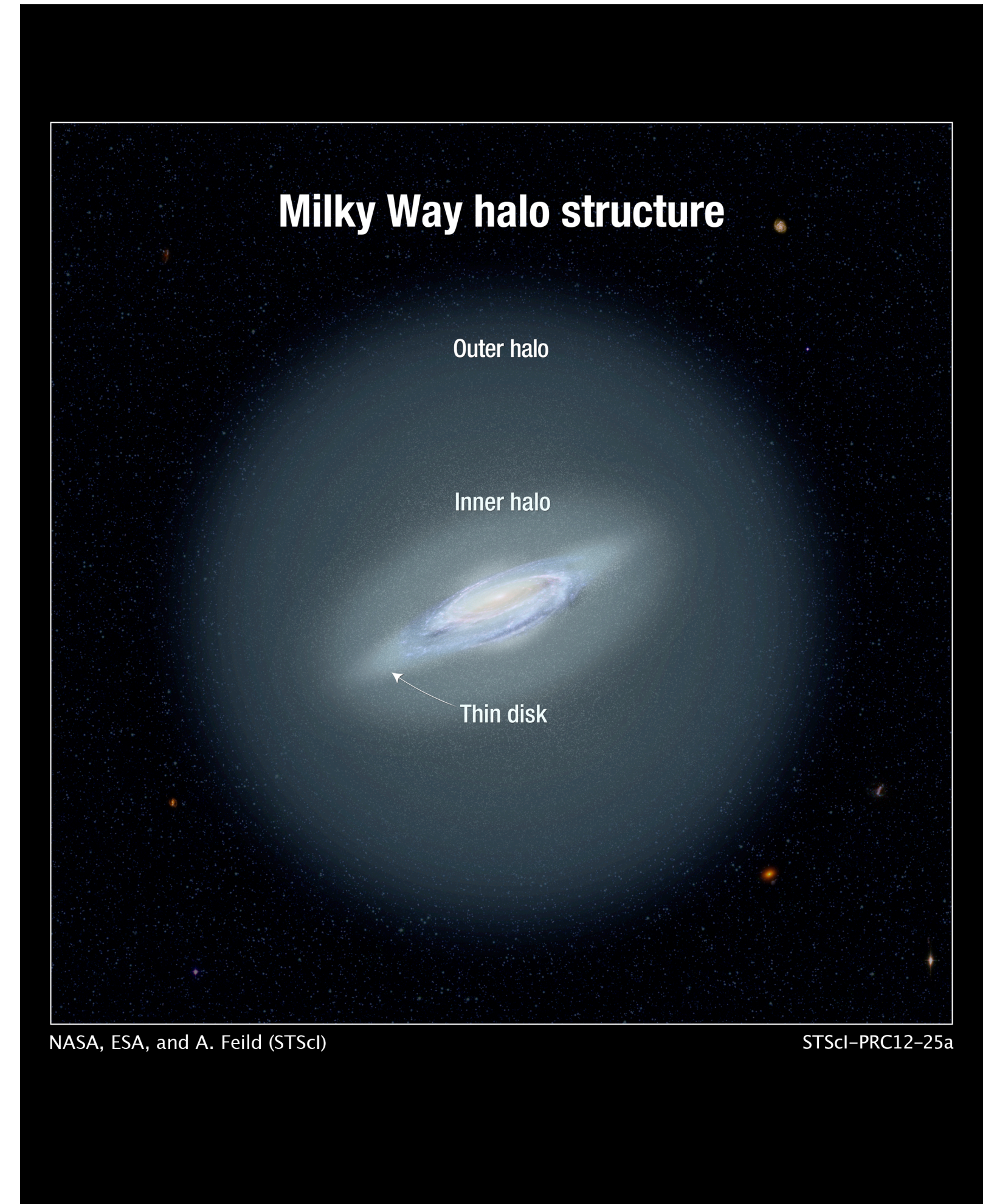
Outer halo

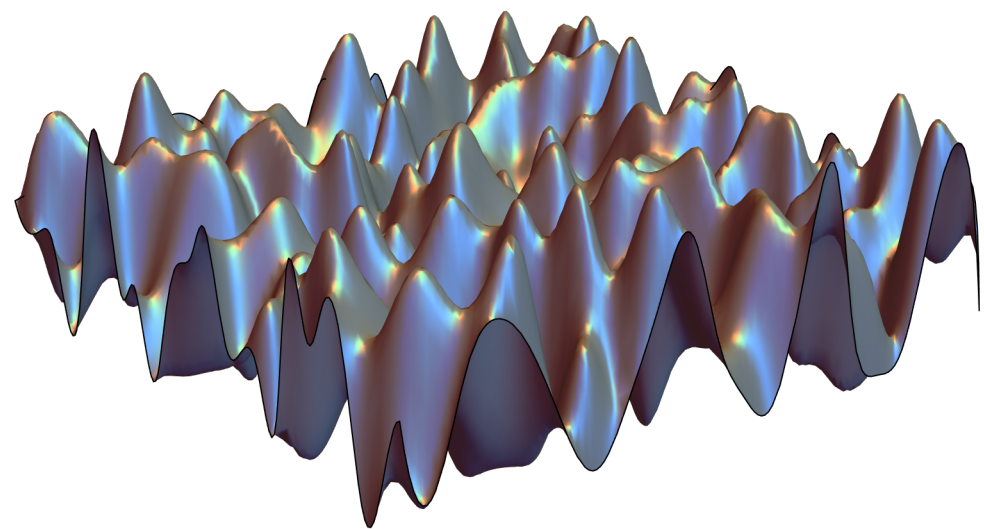
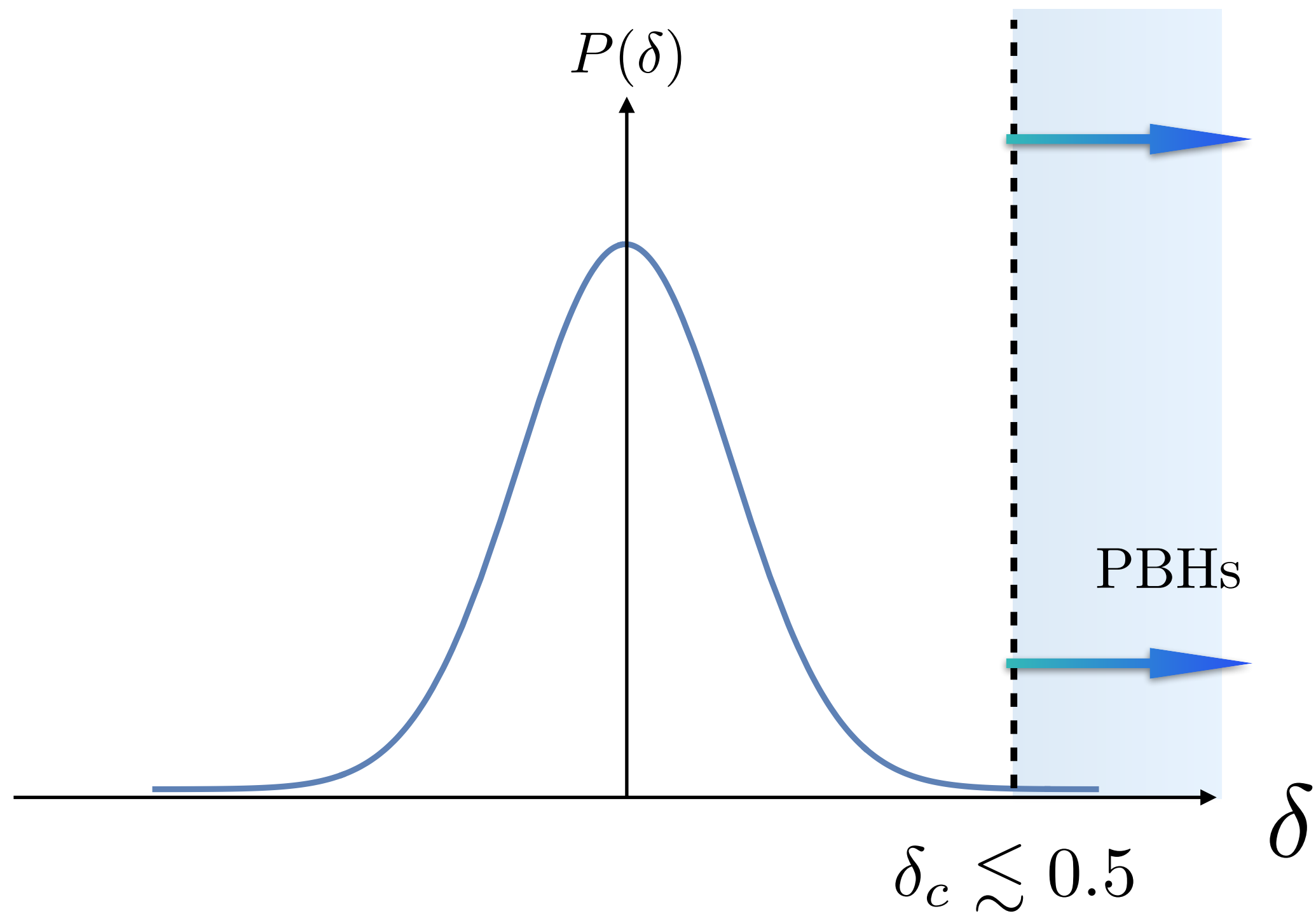
Inner halo

Thin disk

NASA, ESA, and A. Feild (STScI)

STScI-PRC12-25a





Mass:

$$M \sim 10^{-14} \left(\frac{10^{13} \text{ Mpc}^{-1}}{k} \right)^2 M_{\odot}$$

$$N_e \simeq 18 - \frac{1}{2} \log \frac{M}{M_{\odot}}$$

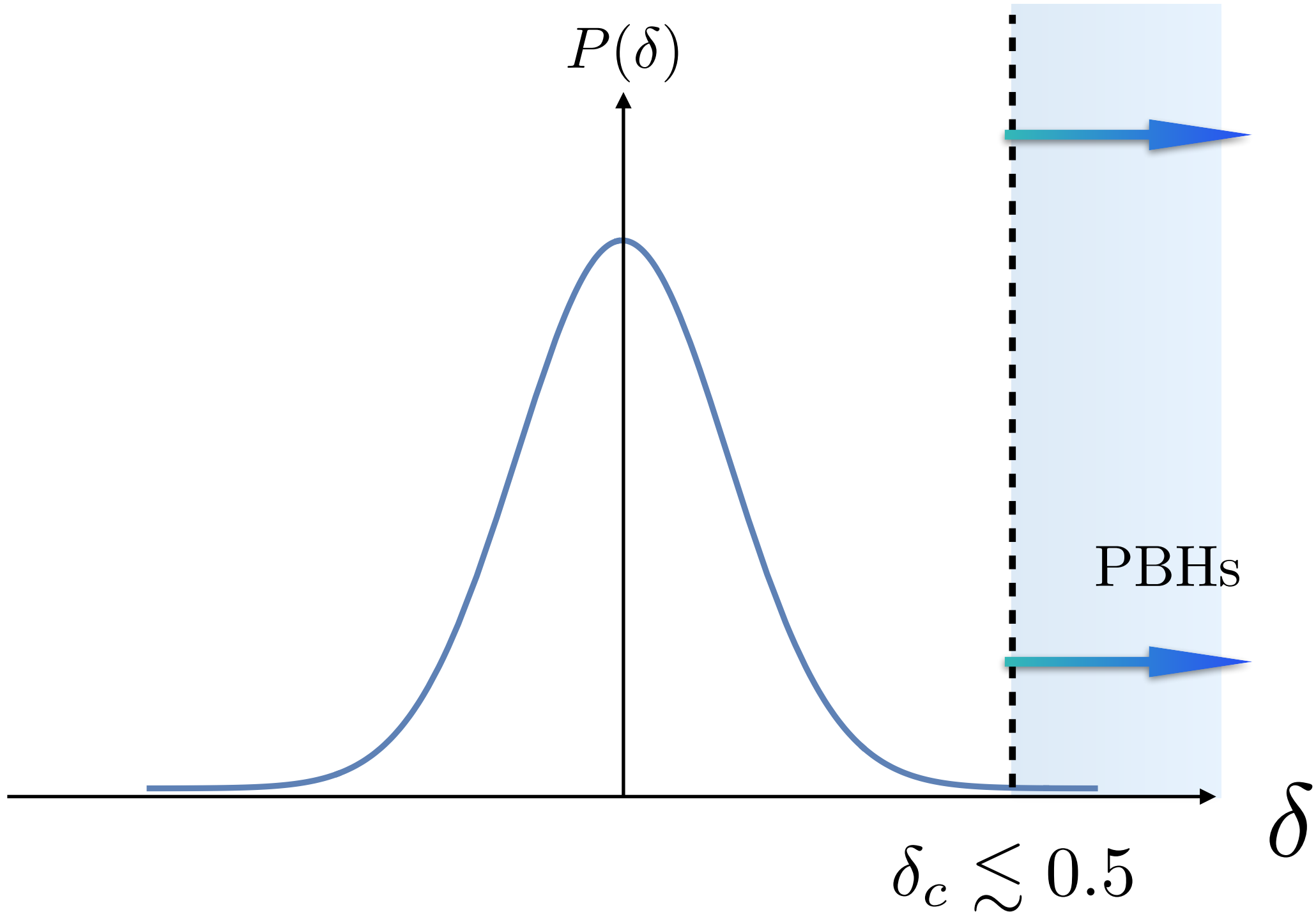
(for PBHs formed during radiation domination)

Abundance (naive Gaussian estimate):

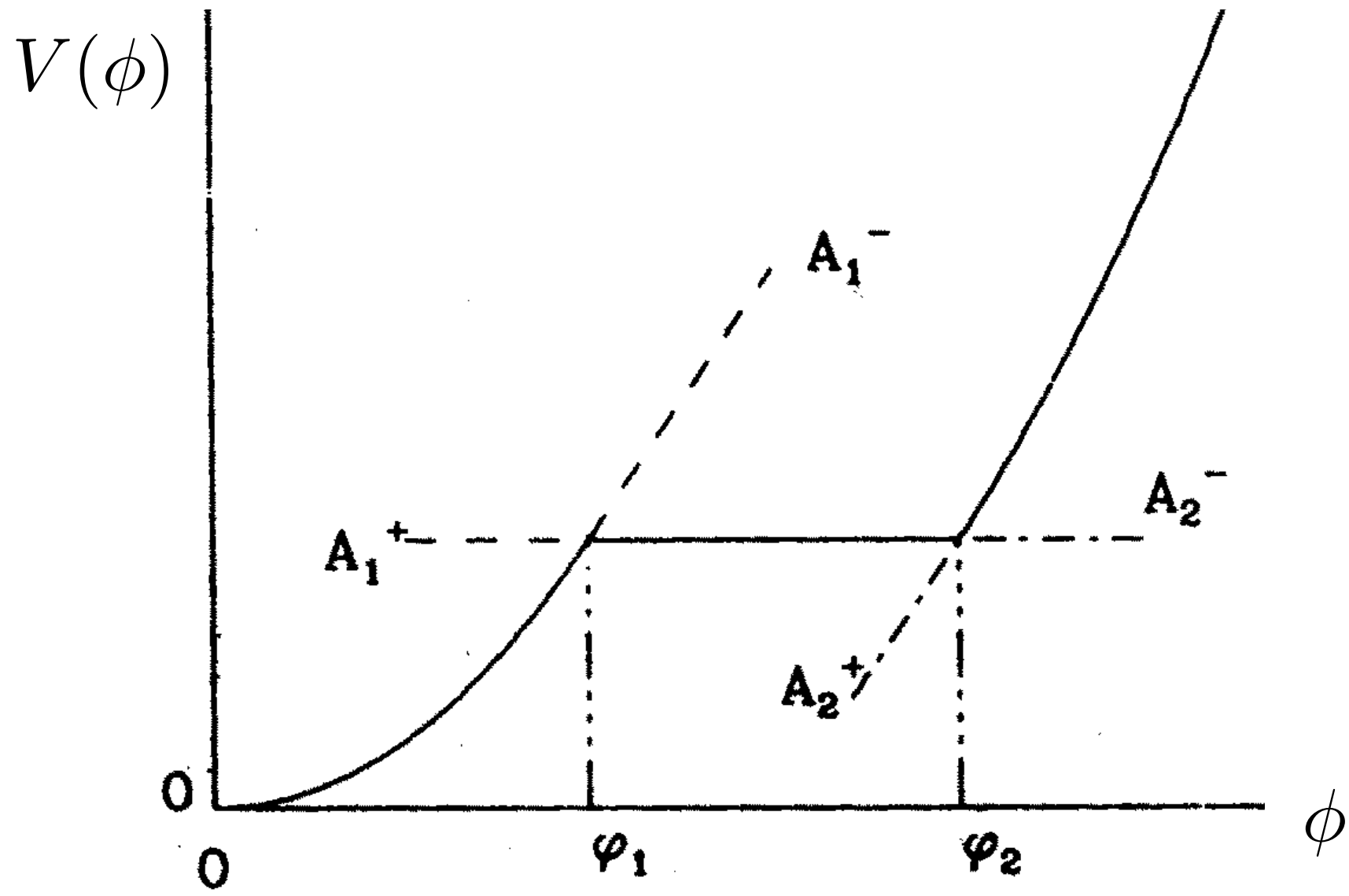
$$f_{\text{PBH}} = \frac{\Omega_{\text{PBH}}}{\Omega_{\text{DM}}} \propto \int_{\delta_c}^{\infty} \exp\left(-\frac{\delta^2}{2\sigma^2}\right) d\delta$$

$$\sigma \sim \mathcal{P}_{\mathcal{R}} \sim 10^{-2} \quad \Rightarrow \quad \frac{\Omega_{\text{PBH}}}{\Omega_{\text{DM}}} \sim 1$$

Inflation and primordial black holes as dark matter

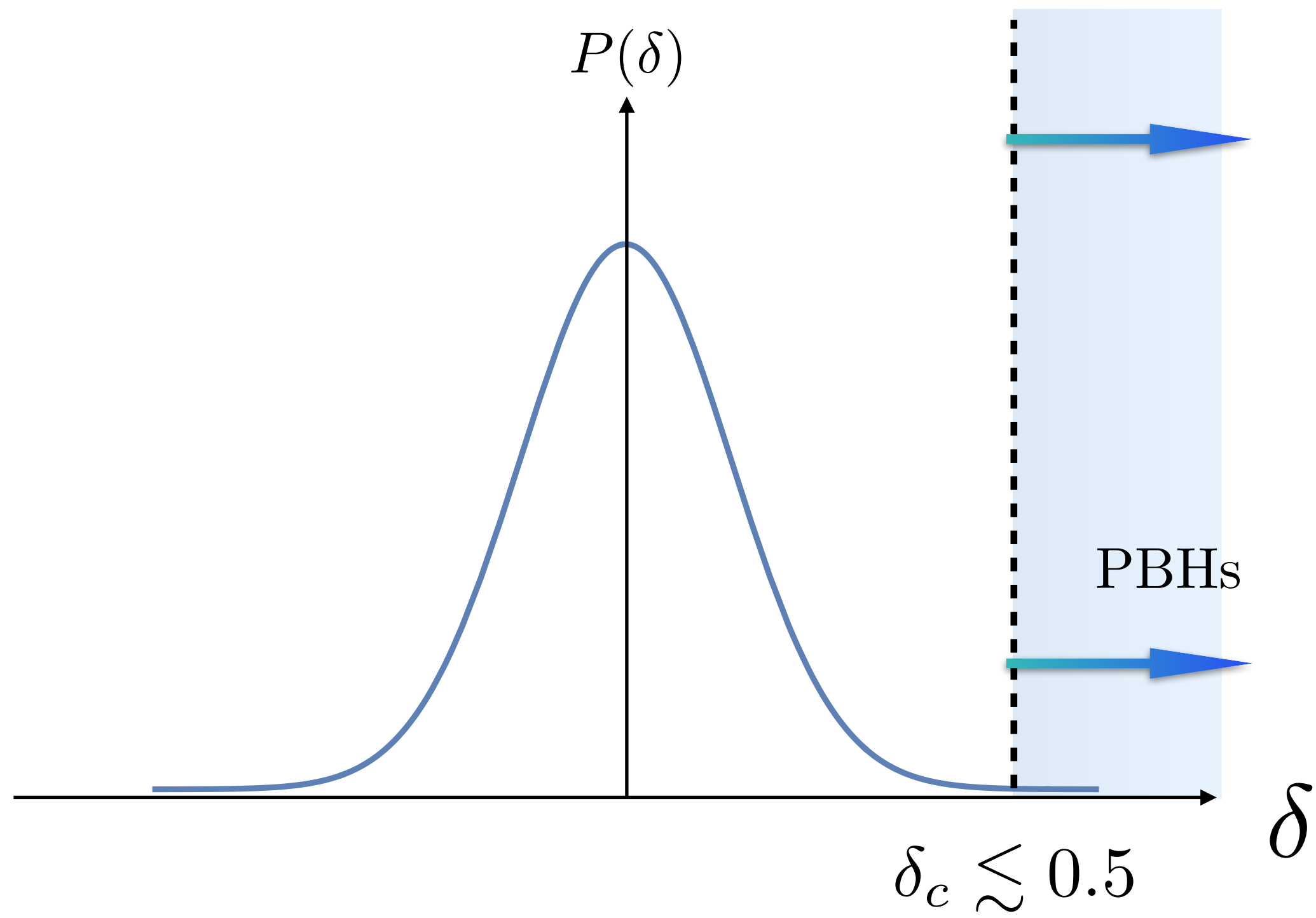


Ivanov, Naselsky, Novikov. 1994



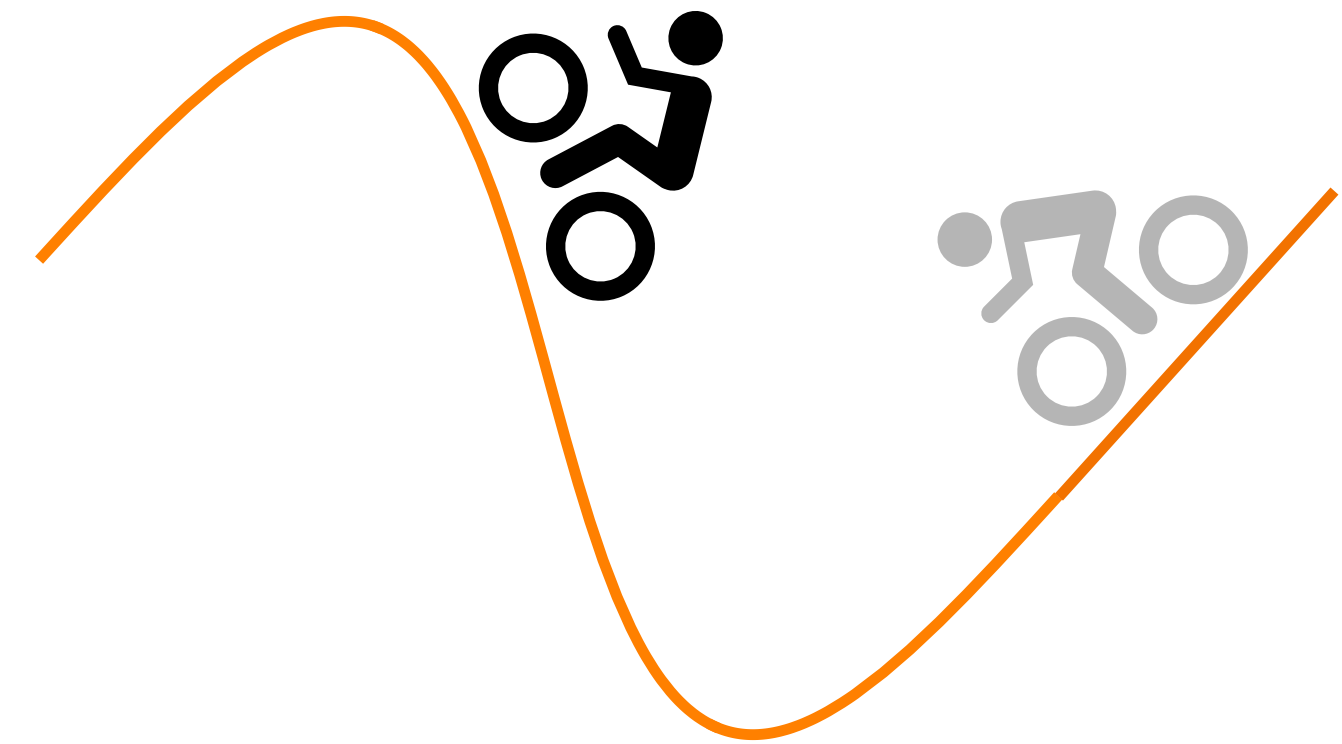
$$\mathcal{P}_{\mathcal{R}} \sim \left(\frac{H}{m_P}\right)^2 \left(\frac{H}{\dot{\phi}}\right)^2 \sim \frac{1}{m_P^2} \left(\frac{V}{V'}\right)^2 \frac{V}{m_P^4}$$

$$V(\phi) = \begin{cases} V_0 + A_+(\phi - \phi_0) & \text{for } \phi > \phi_0 \\ V_0 + A_-(\phi - \phi_0) & \text{for } \phi < \phi_0 \end{cases}$$



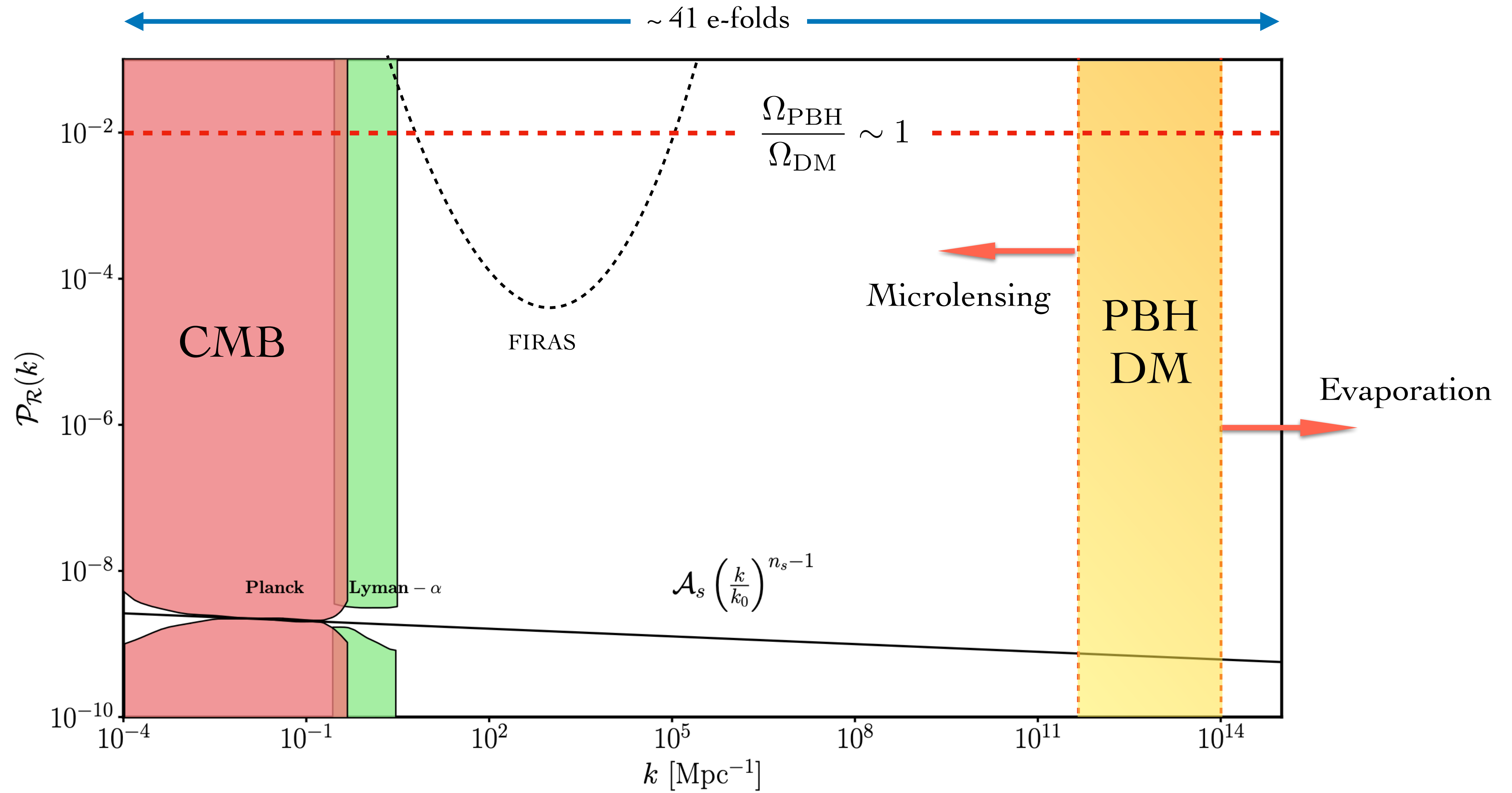
Ultra Slow-Roll (USR)

$$\ddot{\phi} + 3H\dot{\phi} \simeq 0$$



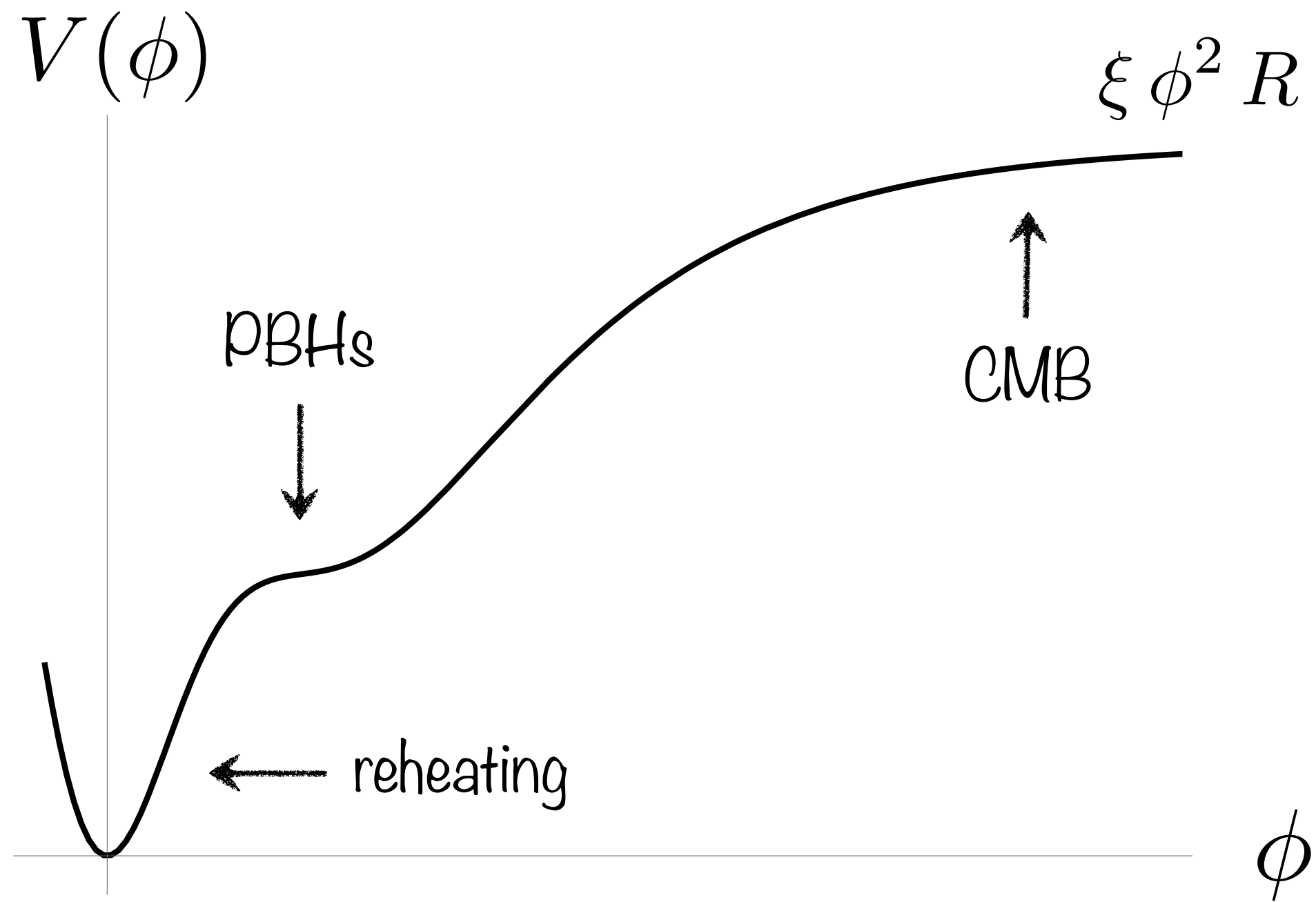
$$\mathcal{P}_{\mathcal{R}} \sim \left(\frac{H}{m_P}\right)^2 \left(\frac{H}{\dot{\phi}}\right)^2 \sim \frac{1}{m_P^2} \left(\frac{V}{V'}\right)^2 \frac{V}{m_P^4}$$

$$\eta \sim \frac{\ddot{\phi}}{\dot{\phi} H} \sim -3$$



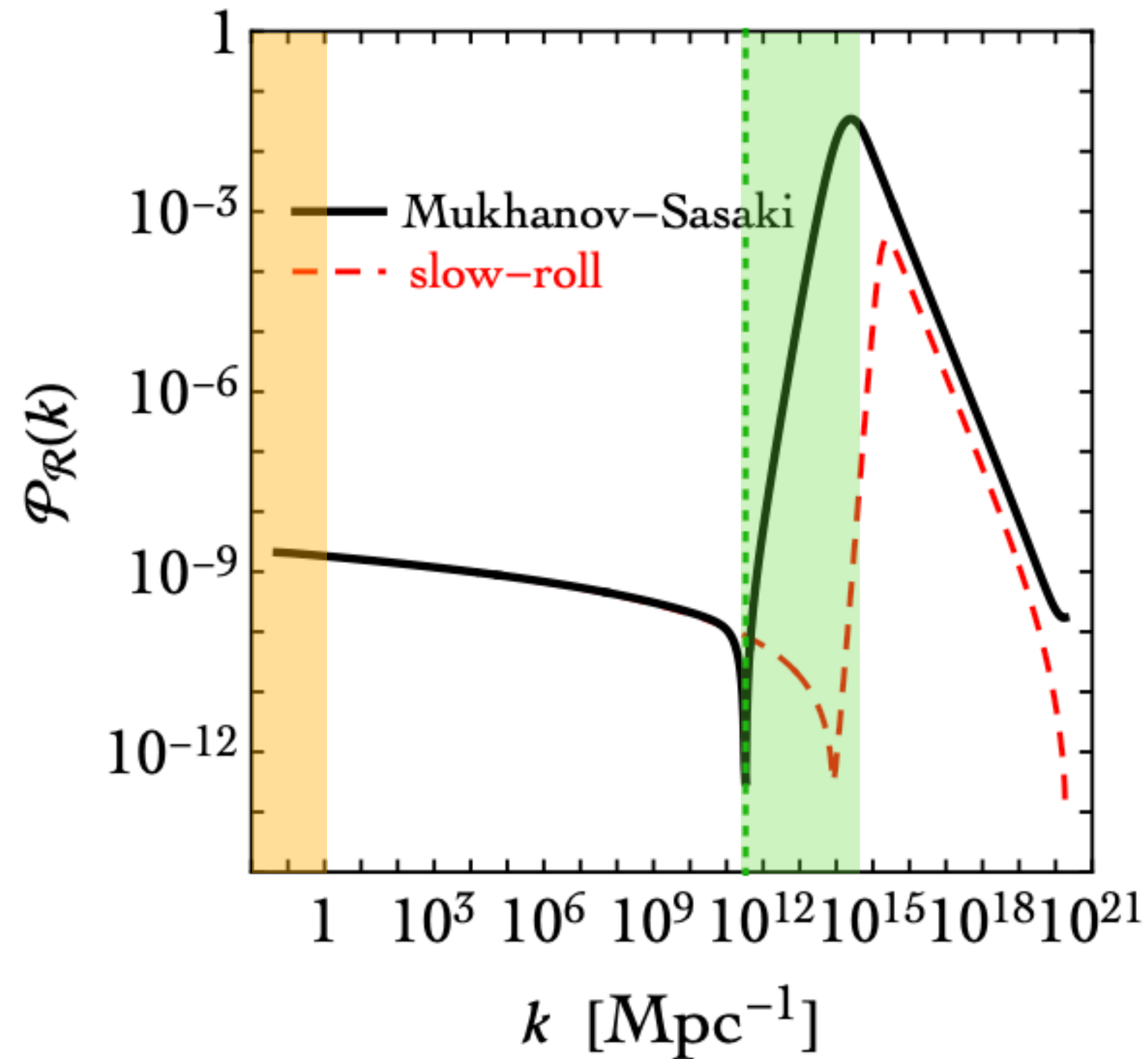
Franco-Abellán, EUCAPT symposium 2023 (modified)

+ enough inflation & successful reheating



$$V = \sum_{n=2}^{4+\epsilon \mathcal{O}(5)} a_n \phi^n$$

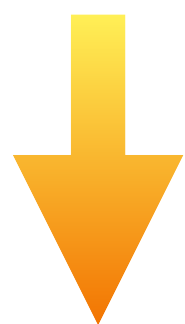
$$\frac{V(\phi)}{\phi^4} = \lambda(\phi_0) + \frac{1}{2} \beta_\lambda(\phi_0) \log \frac{\phi^2}{\phi_0^2} + \frac{1}{8} \beta'_\lambda(\phi_0) \left(\log \frac{\phi^2}{\phi_0^2} \right)^2 + \dots$$



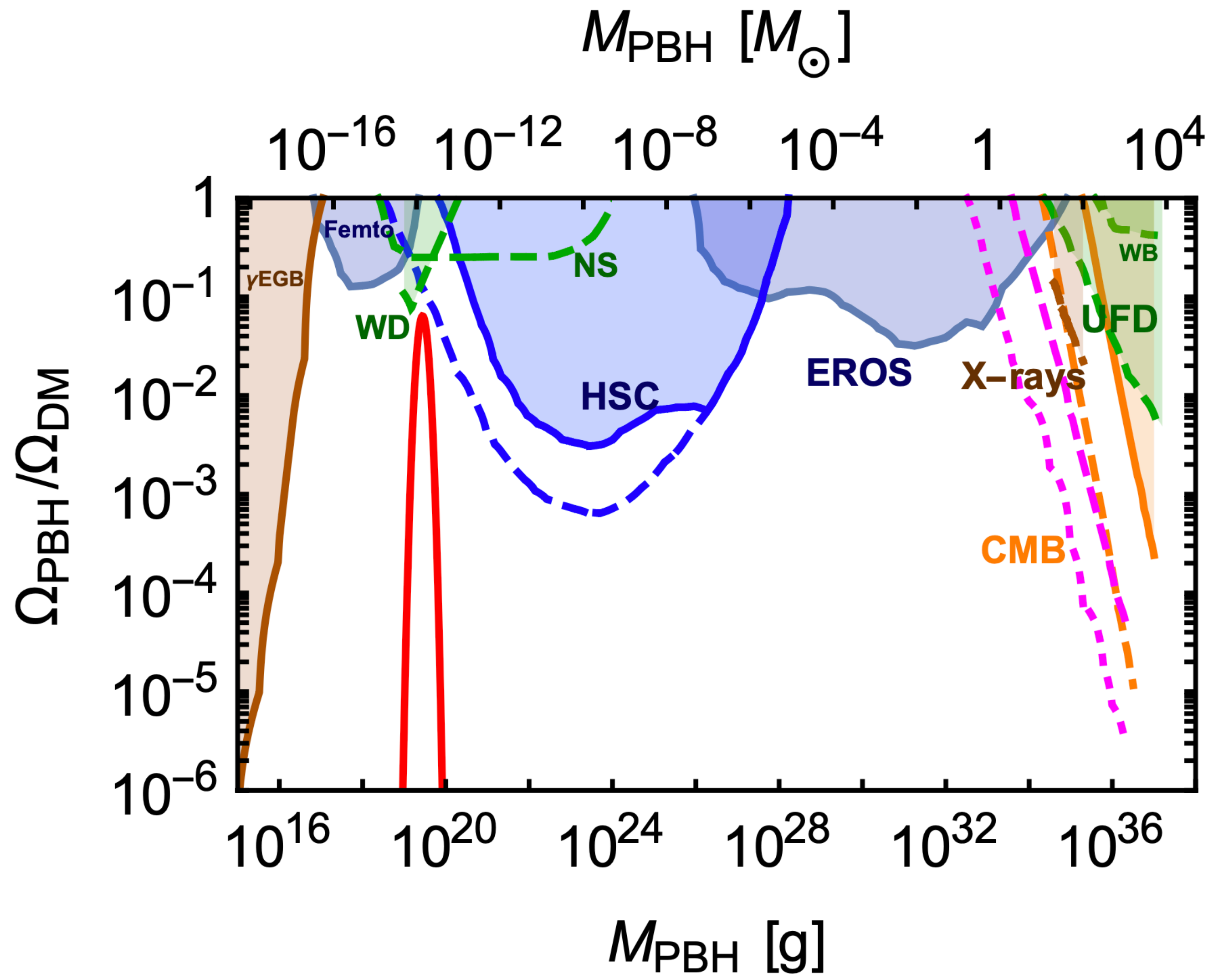
GB, Taoso 2017

GB, Rey, Taoso, Urbano 2020

- Enough inflation
- Agreement with the CMB
(n_s within 3σ)
- $\frac{\Omega_{\text{PBH}}}{\Omega_{\text{DM}}} \sim 1$ (Gaussian, RD)



$$10^{-16} M_{\odot} \leftrightarrow 10^{-12} M_{\odot}$$

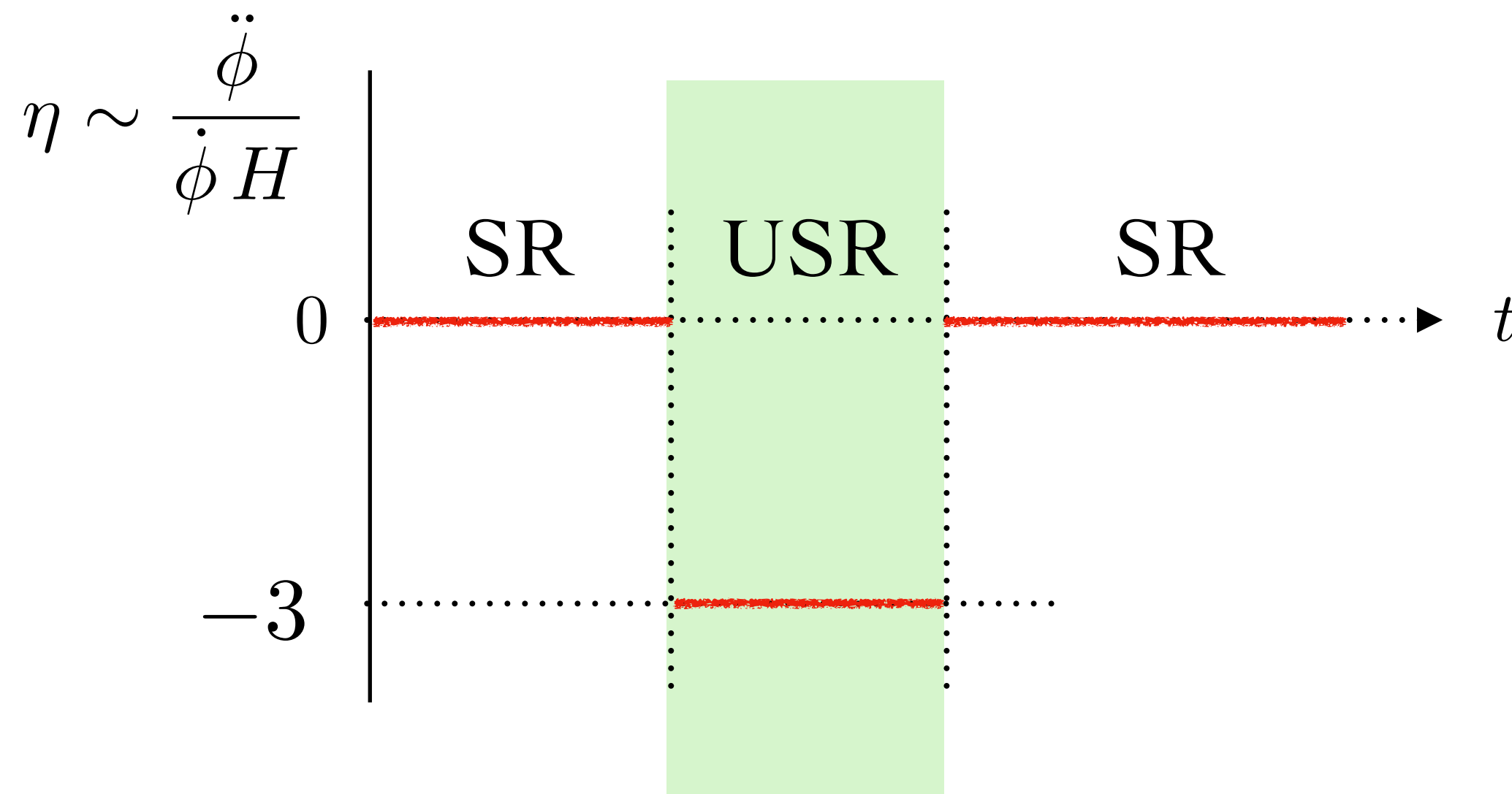


Breakdown of perturbation theory in USR inflation?

Claim: *a large enough tree-level primordial spectrum for PBH DM implies perturbation theory breaks at CMB scales.*

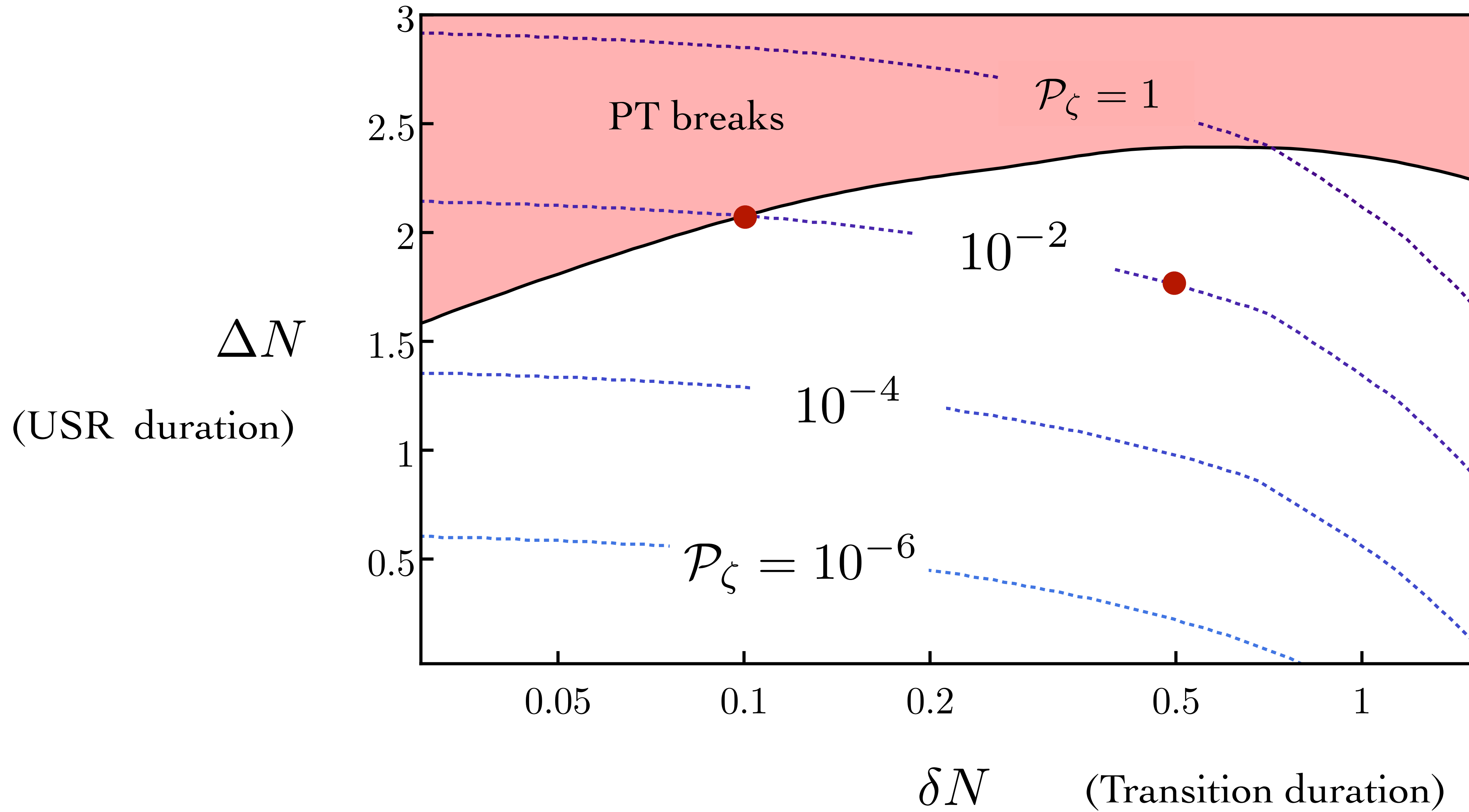
Kristiano and Yokoyama, 2022 & 2023

Toy model: SR \rightarrow USR \rightarrow SR (with sharp transitions)



$$\mathcal{P}_\zeta \ll \frac{1}{(\Delta\eta)^2} \simeq 0.03$$

(for perturbation theory to hold)



Other mechanisms to form PBH

- * Single-field inflation other than USR
- * Transient Dissipation during inflation

* Single-field inflation other than USR

$$\mathcal{S} = \int d^4x M \frac{a^3 \epsilon}{c_s^2} \left(\dot{\mathcal{R}}^2 - \frac{c_s^2}{a^2} |\nabla \mathcal{R}|^2 \right) \quad \mathcal{R} \simeq C_1 + C_2 \int \frac{c_s^2}{a^3 M^2 \epsilon H} dN$$

Rate of change of ϵ

$$\frac{d\mathcal{R}}{dN} \propto \exp \left[- \int (3 + \epsilon - 2\eta - 2s + \mu) dN \right]$$

Rate of change of c_s

Rate of change of M

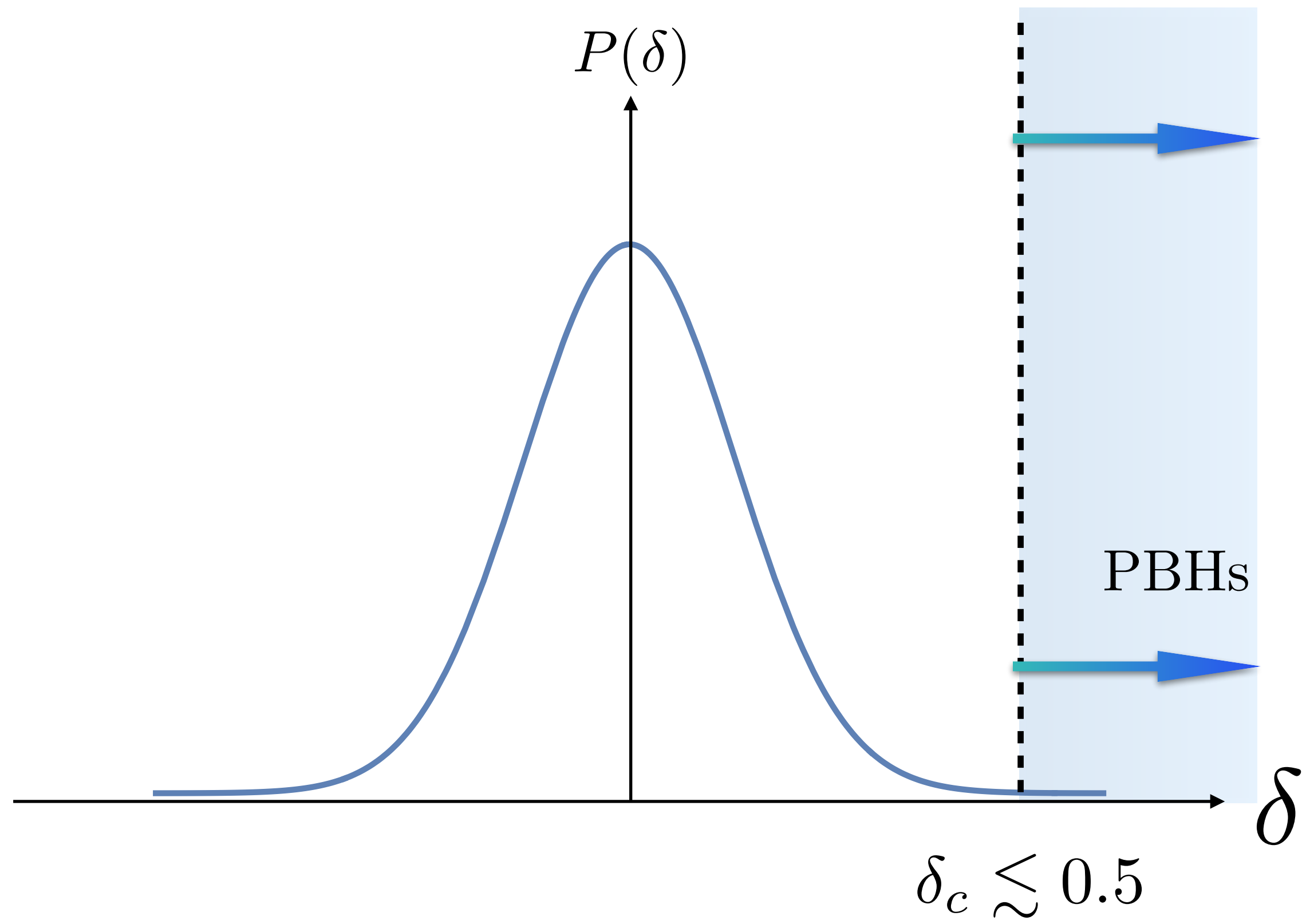
✳ (Transient) Dissipation during inflation

Background: $\ddot{\phi} + (3H + \mathbf{\Gamma})\dot{\phi} + V' = 0$

Fluctuations (schematically): $\delta\ddot{\phi}_{\mathbf{k}} + (3H + \Gamma)\delta\dot{\phi}_{\mathbf{k}} + \left(\frac{k^2}{a^2} + \dot{\phi}\Gamma_{\phi}\right)\delta\phi_{\mathbf{k}} \propto \sqrt{\frac{2\Gamma T}{a^3}}\xi_{\mathbf{k}}(t)$

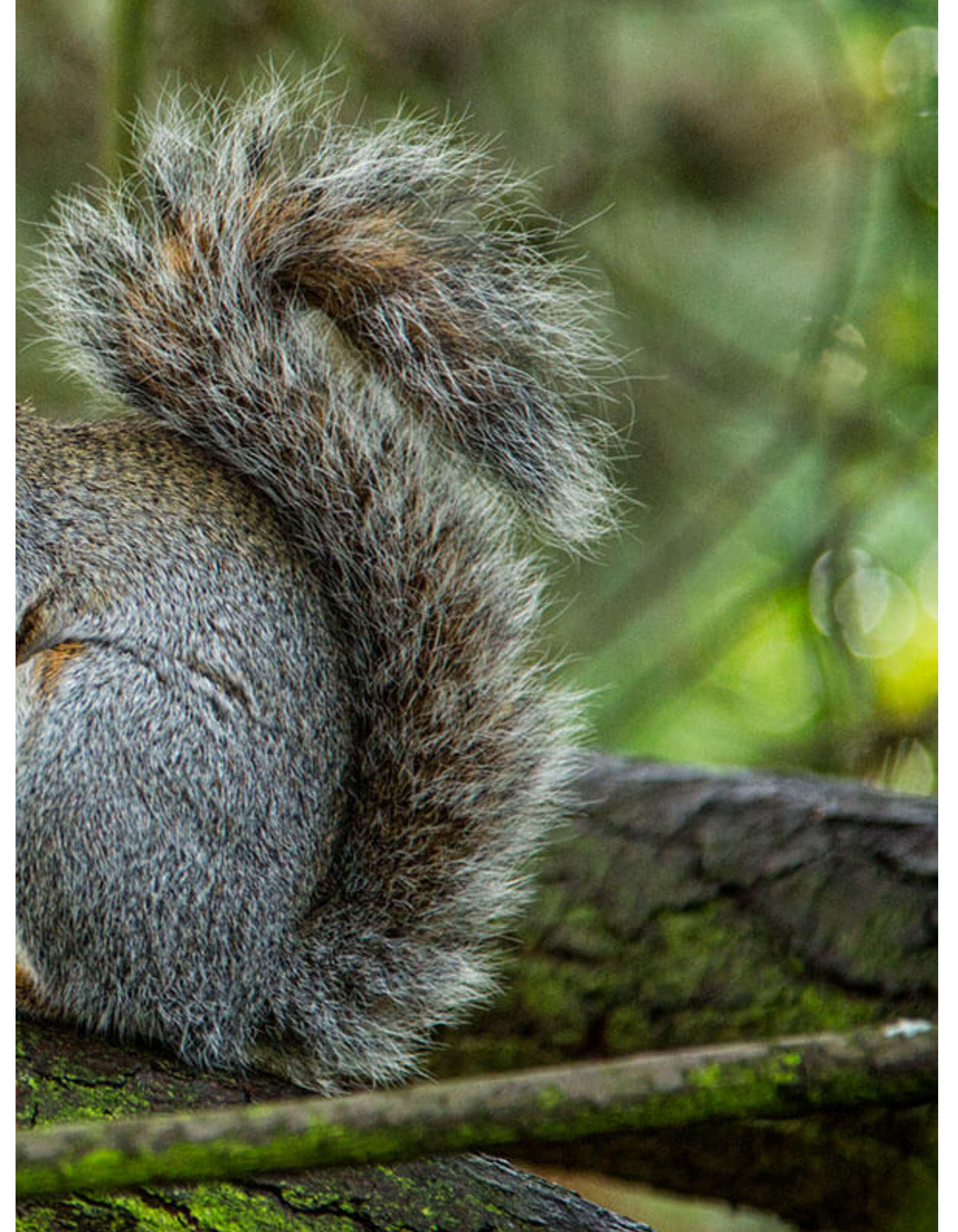
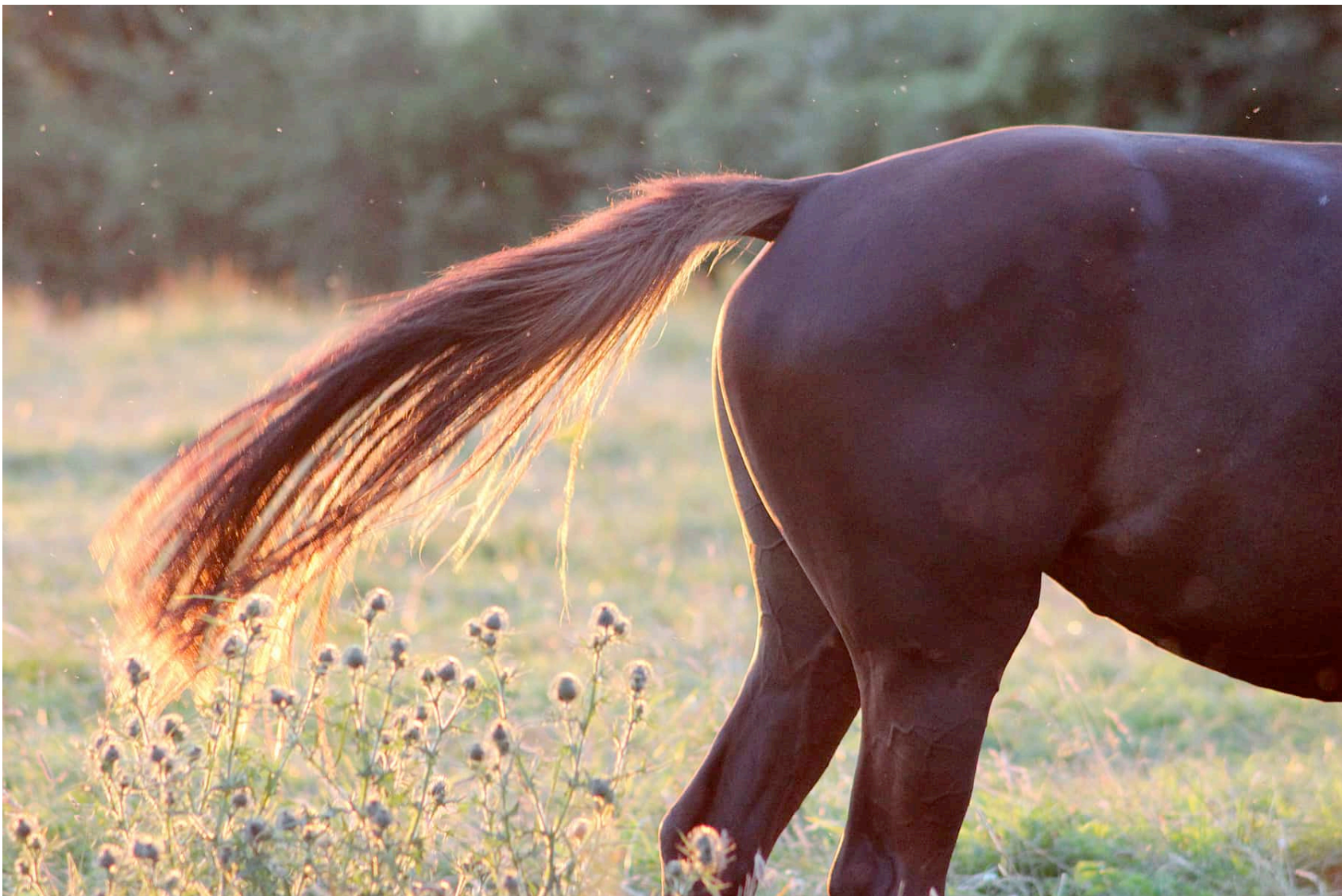
(Stochastic thermal noise)

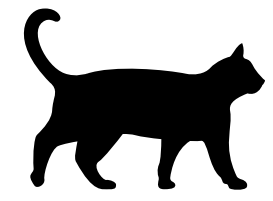
The problem of the abundance



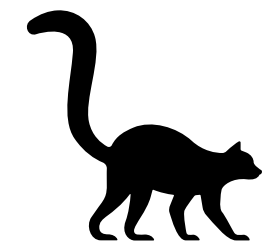
$$\frac{\Omega_{\text{PBH}}}{\Omega_{\text{DM}}} = \int_{\delta_c}^{\infty} f(\delta) d\delta$$

How does the tail of the PDF look like?





The relation between ζ and δ is non-linear

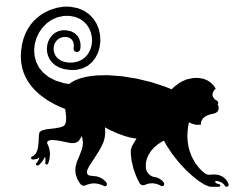


ζ is, in general, intrinsically non-gaussian



The abundance may require more than an integral

$$f(\delta) \sim \exp\left(-\frac{\zeta}{2\mathcal{P}_\zeta} + \frac{\langle\zeta\zeta\zeta\rangle}{\mathcal{P}_\zeta^3}\zeta^3 + \dots\right) \quad \text{Small corrections for small } \zeta$$



Several indications of non-gaussian tails, for large ζ

Non-linear saddle point for ζ^4

Celoria, Creminelli, Tambalo, Yingcharoenrat 2021

USR { Stochastic inflation, numerically
Stochastic δN formalism

Figueroa, Raatikainen, Rasanen, Tomberg 2020

Pattison, Vennin, Wands, Assadullahi 2021

$$\frac{\dot{\phi}}{H} \gg \frac{H}{2\pi} \quad \longrightarrow \quad \mathcal{P}_\zeta \ll 1 \quad \text{Quantum diffusion?}$$

Non-Gaussian tails in the PDF of curvature perturbations arise in USR inflation without invoking stochastic inflation.

GB, Konstandin, Pérez Rodríguez, Pierre, Rey 2024

Gravitational wave signatures of PBH DM

- * Second order induced gravitational waves
- * Gravitational wave emission from PBH mergers

* Second-order induced gravitational waves

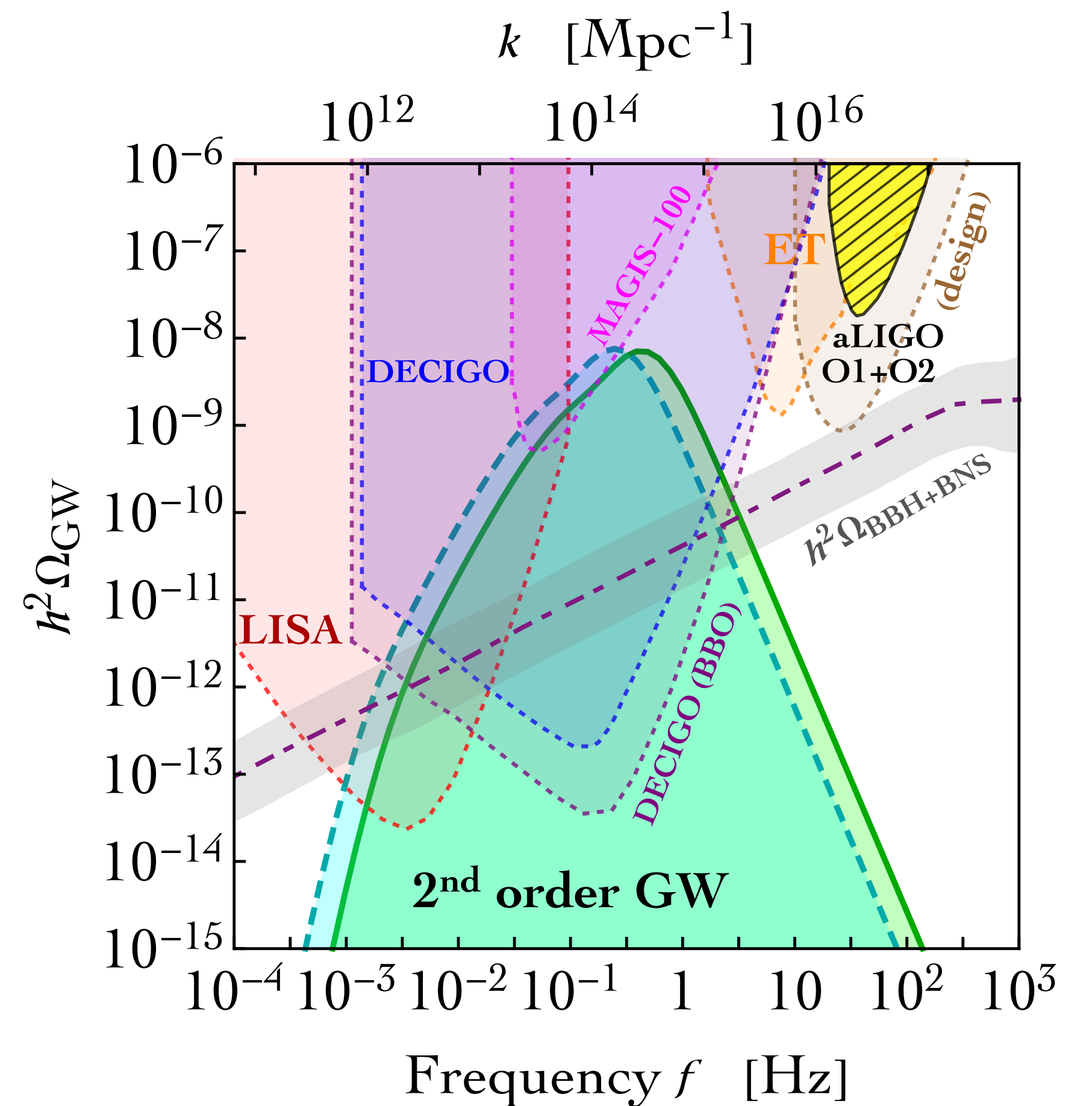
$$h''_{ij} + 2\frac{H'}{h} h_{ij} - \nabla^2 h_{ij} = S_{ij} \sim \partial_i \mathcal{R} \partial_j \mathcal{R} \quad \Rightarrow \quad \Omega_{\text{GW}} \sim \mathcal{P}_h \sim (\mathcal{P}_{\mathcal{R}})^2$$

$$\left(\frac{M_{\text{PBH}}}{10^{17} \text{ g}}\right)^{-1/2} \simeq \frac{k}{2 \cdot 10^{14} \text{ Mpc}^{-1}} \simeq \frac{f}{0.3 \text{ Hz}}$$

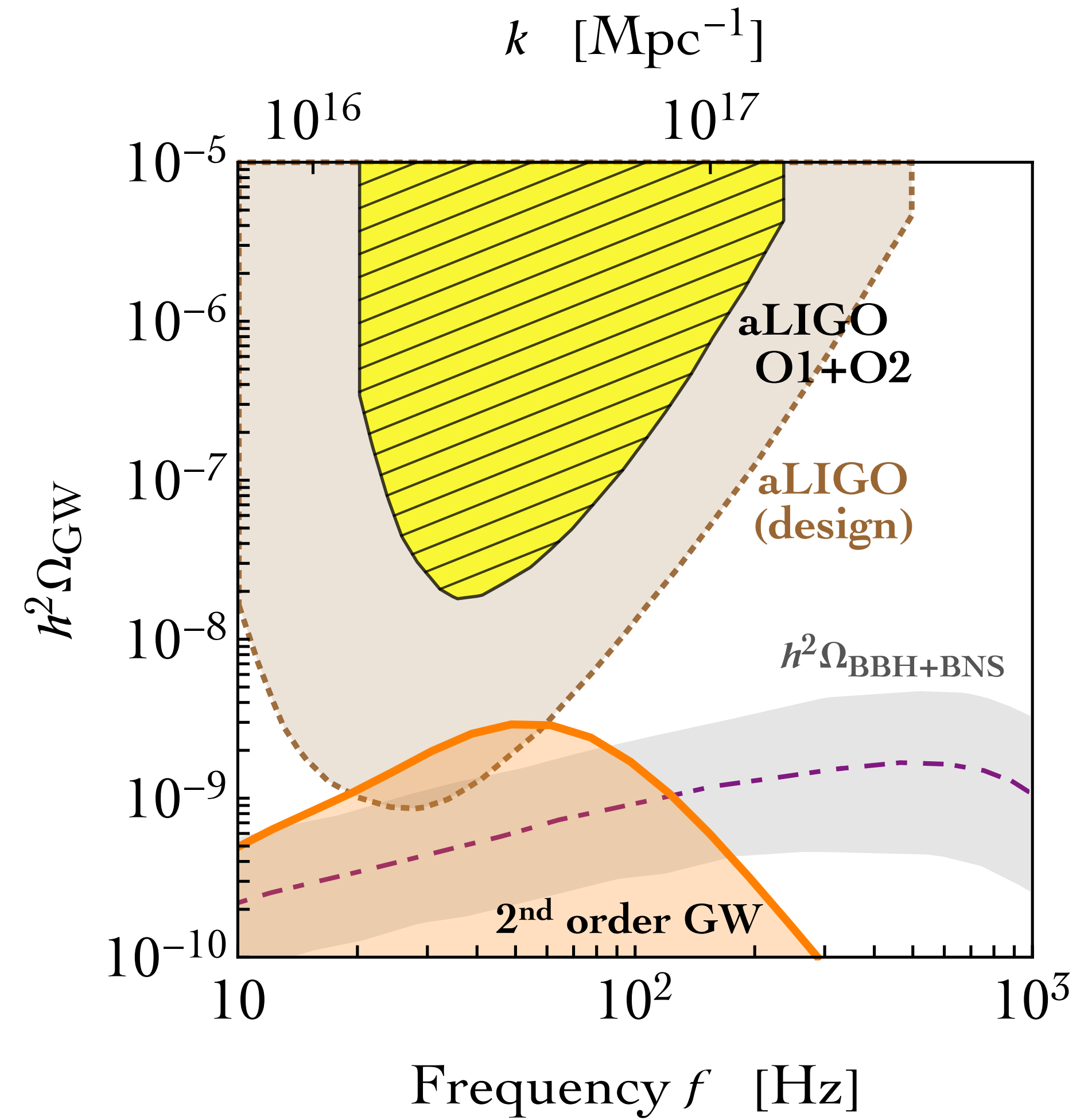
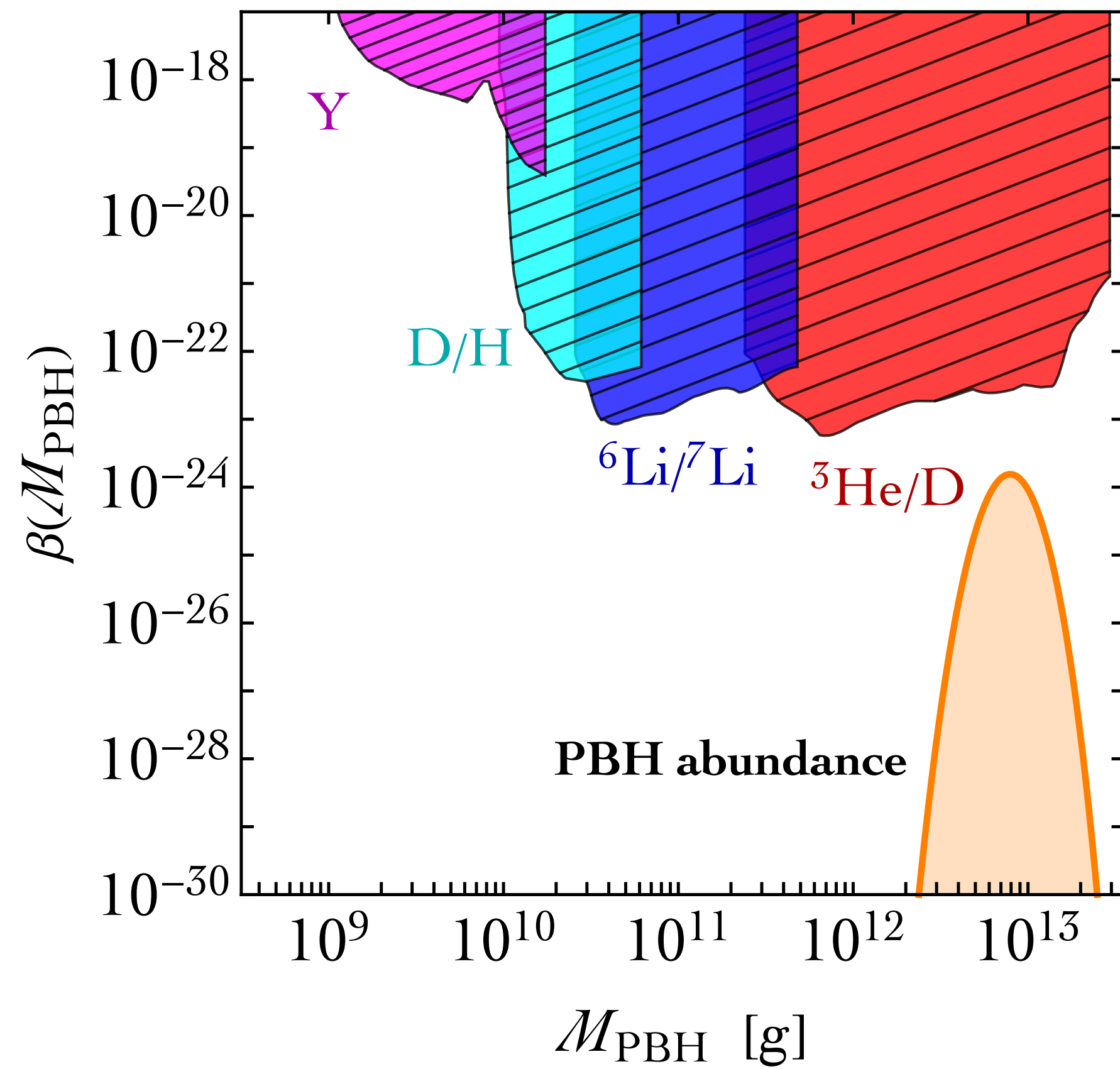
* LISA: for PBH DM

* LIGO/Virgo/KAGRA: if PBHs already evaporated

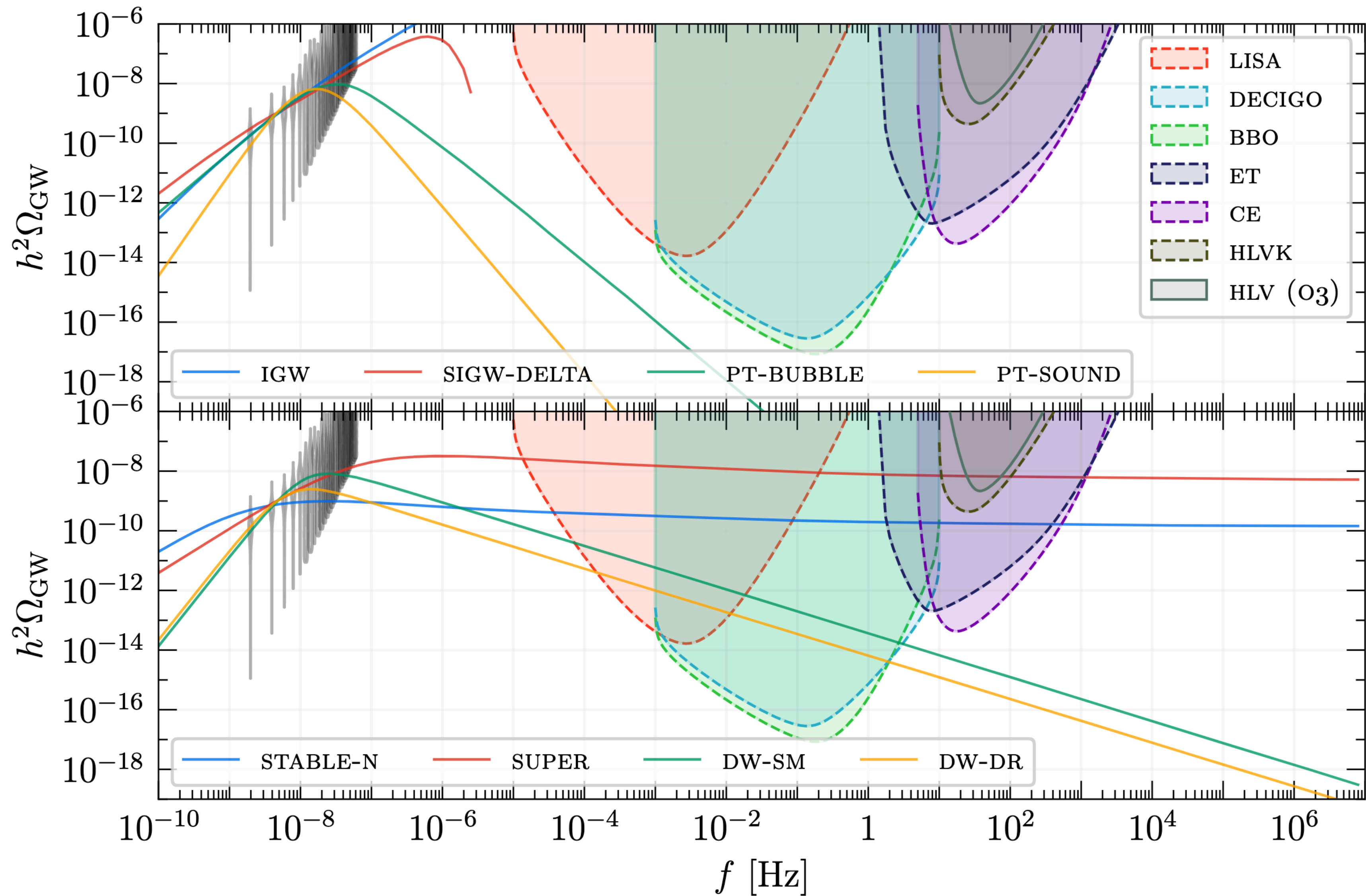
* PTA: if PBHs of ~ 0.1 Solar masses



✳ Second-order induced gravitational waves



* Second-order induced gravitational waves



PTA
 $\sim \text{NHz} \sim 10 M_{\odot}$

* High frequency GW from PBH DM mergers

No known astrophysical sources at GHz

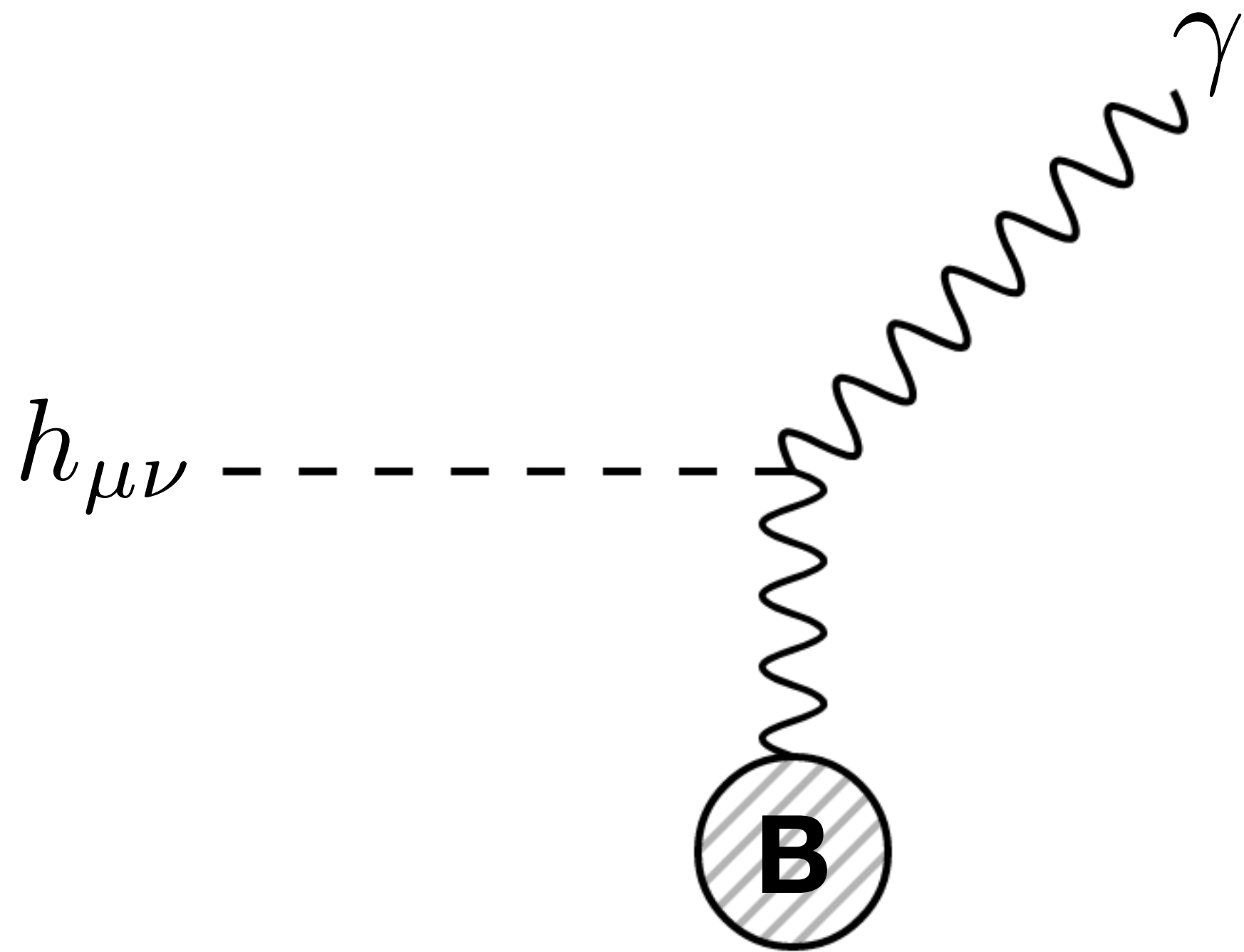
Asteroid mass PBH mergers. GW frequency \gtrsim GHz

Small characteristic strain $\sim 10^{-25}$

- Very strong clustering
- Enhanced local DM density for ~ 10 Kpc

✦ High frequency GW from PBH DM mergers

Inverse Gertsenshtein effect. Connection with axion physics



ALPS II. DESY / Heiner Müller-Elsner

Summary

- Asteroid-mass PBHs are a strong contender to explain the DM.
- PBH from inflation. Interesting for phenomenology and theory playground
- GW: indirect probes of PBH DM. Technical challenges.
- Goal: detect subsolar BH with lensing, GW interferometry, etc.