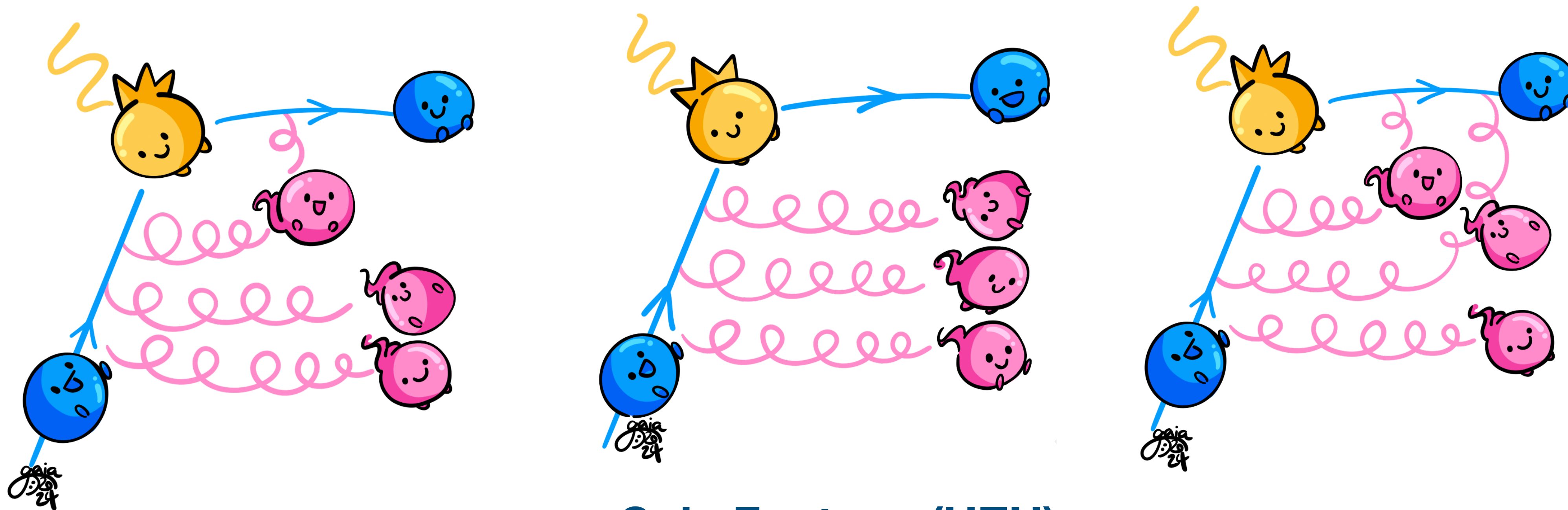


3 loops & 4 cuts:

towards N3LO RRR antenna functions



Gaia Fontana (UZH)

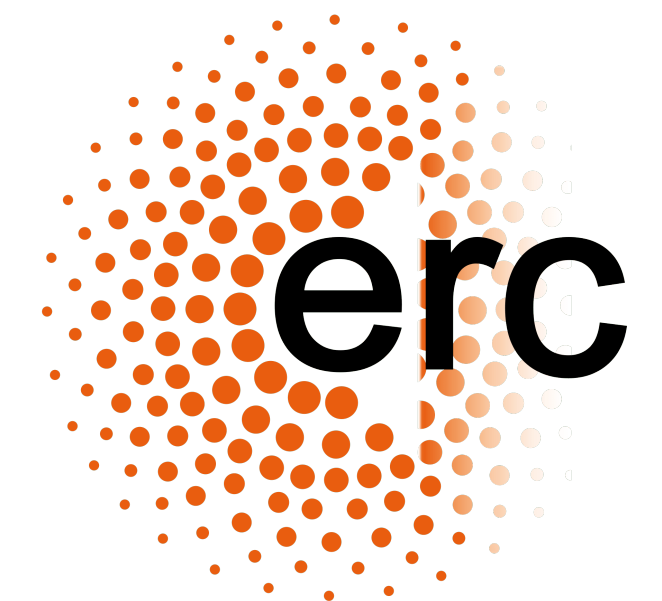
In collaboration with Thomas Gehrmann & Kay Schönwald

Based on JHEP 03 (2024) 159 & upcoming works

35th Rencontres de Blois
Blois, 22/10/2024



Universität
Zürich^{UZH}



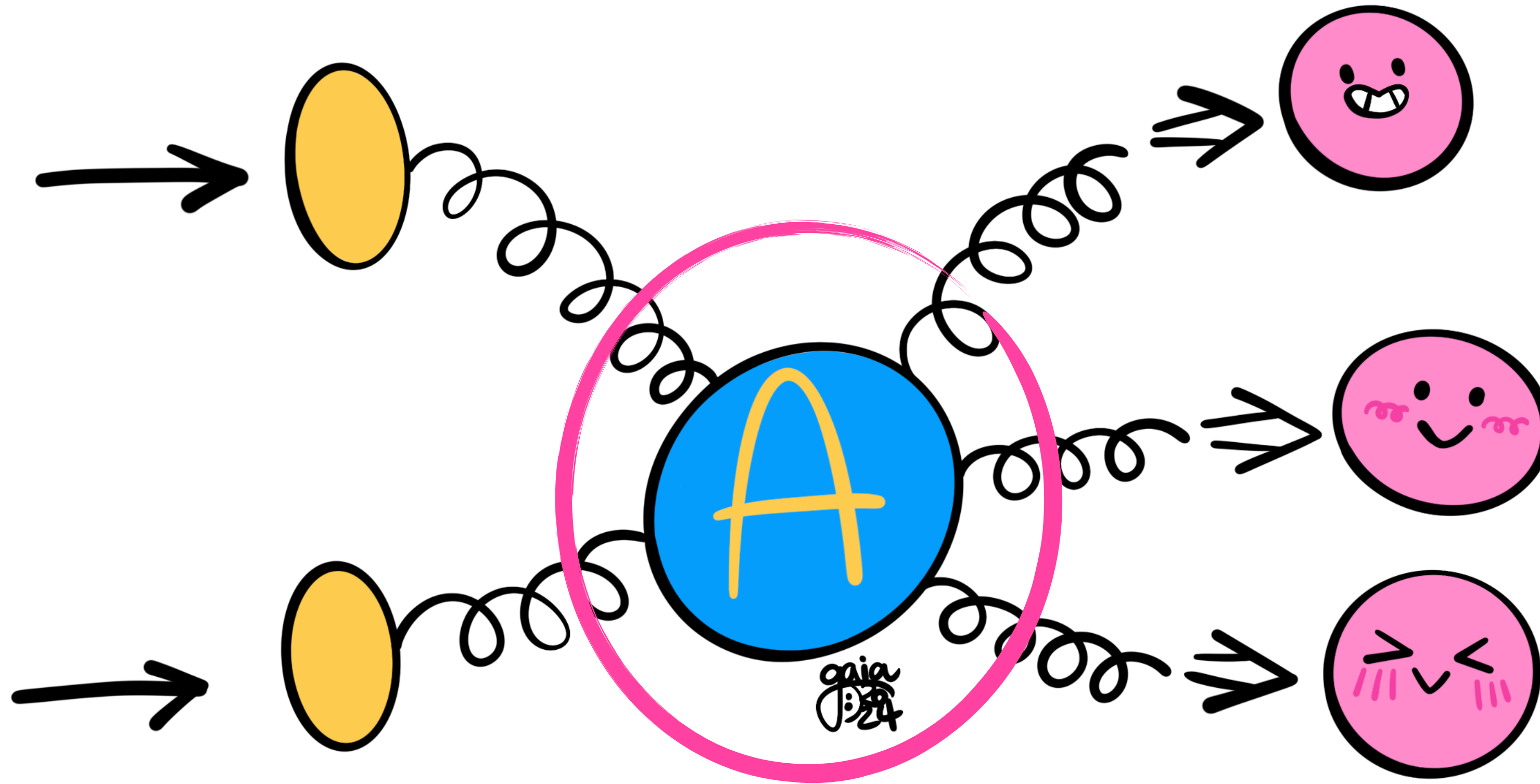
Why precision?

- **Precision physics** as
 - test of the Standard model
 - gate to new physics
- **High-Lumi upgrade of LHC** :
 - theory and experiments must have comparable uncertainties
 - needed: %-level accuracy:
 - perturbation theory @ **NNLO** and often **N³LO**



TOM GAULD for NEW SCIENTIST

Recipe for a theoretical prediction



Many ingredients

- PDFs to describe the proton structure
- **Hard scattering**
- Radiation and evolution to hadronic states

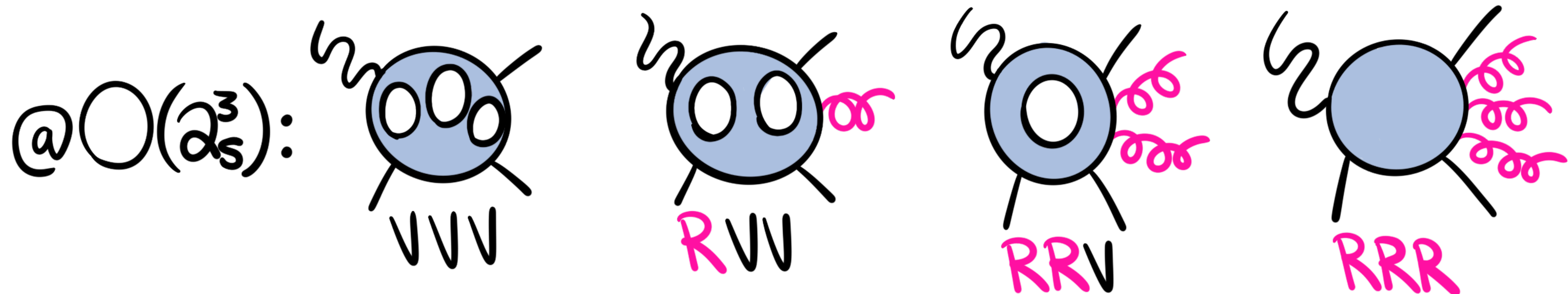
Hard Scattering

Looking @ QCD corrections:

$$d\sigma = d\sigma_{LO} + \alpha_S d\sigma_{NLO} + \alpha_S^2 d\sigma_{NNLO} + \alpha_S^3 d\sigma_{N3LO} + \dots$$

Perturbative series in the strong coupling

Beyond LO: contributions from diagrams with increasing loops and legs

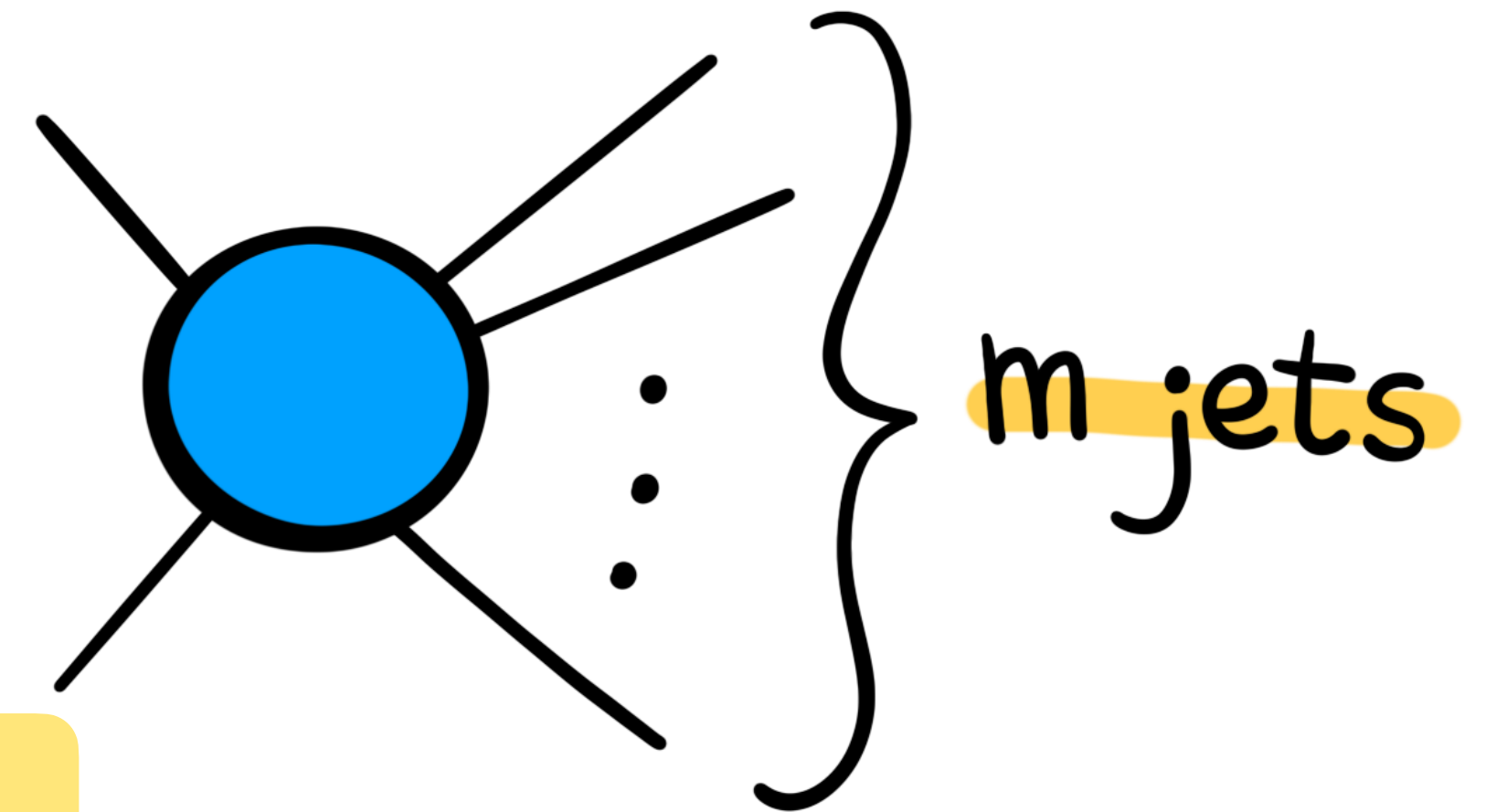


Real corrections!

$$d\sigma = d\sigma_{LO} + \alpha_S d\sigma_{NLO} + \alpha_S^2 d\sigma_{NNLO} + \alpha_S^3 d\sigma_{N3LO} + \dots$$

@ LO

$$d\sigma_{LO} = \int_{d\Phi_m} d\sigma_{Born}$$



@ NLO

$$d\sigma_{NLO} = \int_{d\Phi_{m+1}} d\sigma_{NLO}^R + \int_{d\Phi_m} d\sigma_{NLO}^V$$

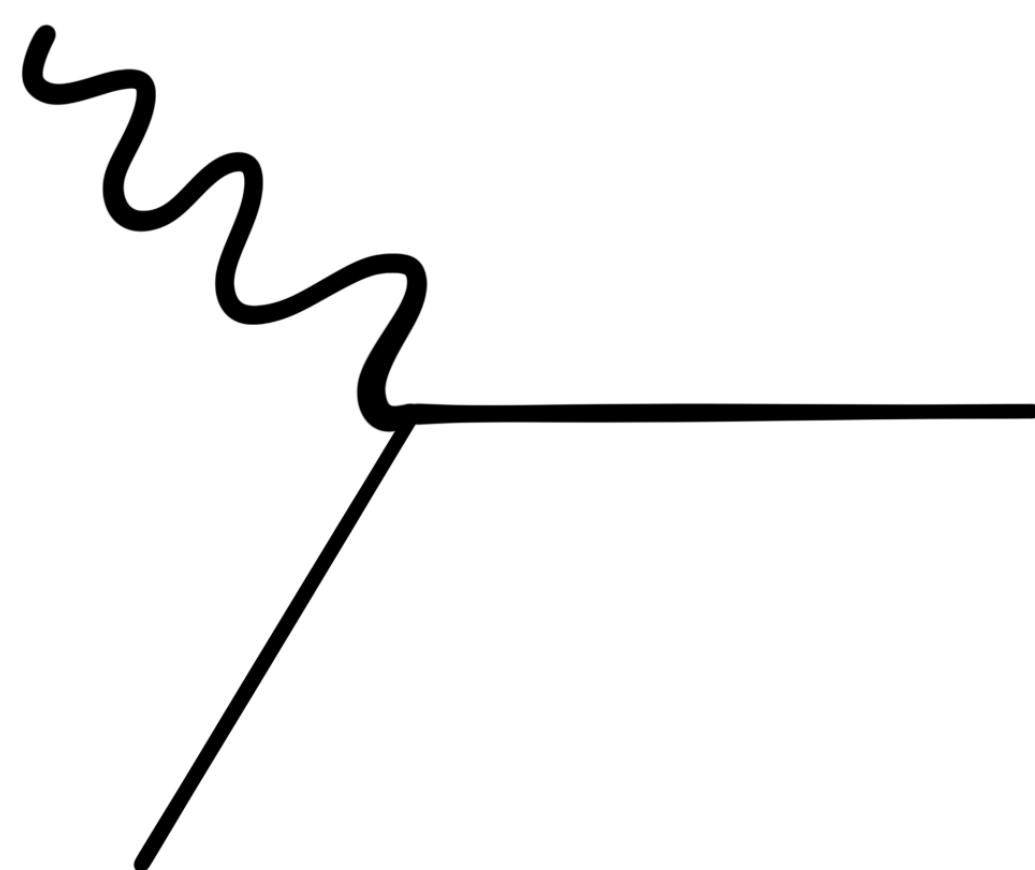
@ NNLO

$$d\sigma_{NNLO} = \int_{d\Phi_{m+2}} d\sigma_{NNLO}^{RR} + \int_{d\Phi_{m+1}} d\sigma_{NNLO}^{RV} + \int_{d\Phi_m} d\sigma_{NNLO}^{VV}$$

@ N3LO

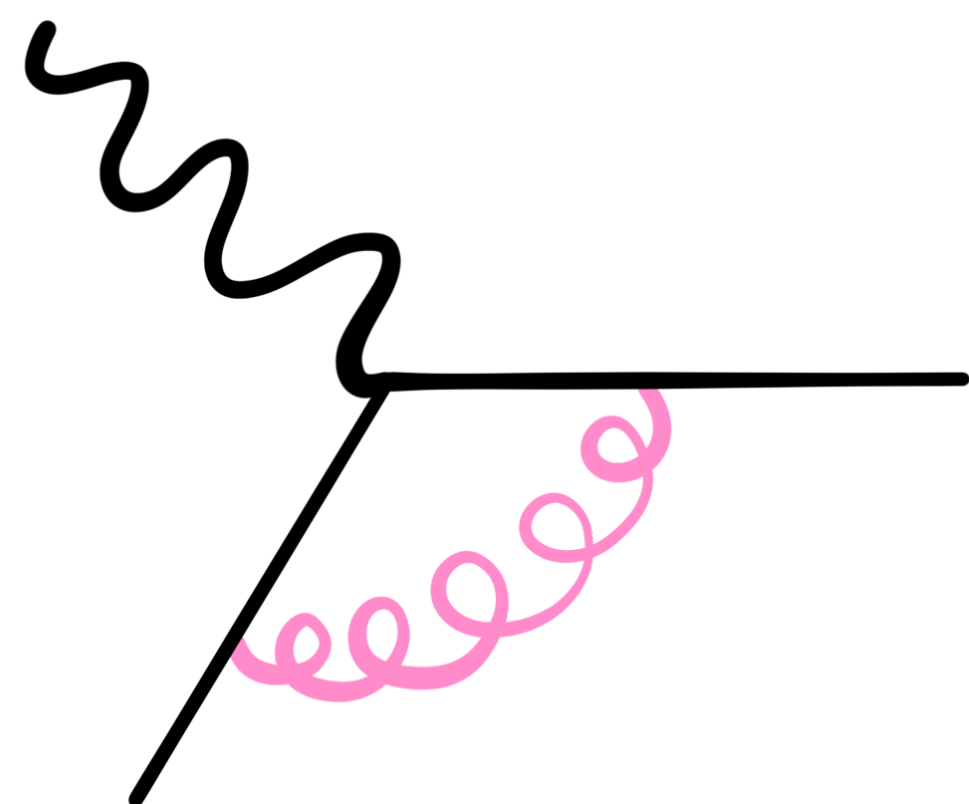
$$d\sigma_{N3LO} = \int_{d\Phi_{m+3}} d\sigma_{N3LO}^{RRR} + \int_{d\Phi_{m+2}} d\sigma_{N3LO}^{RRV} + \int_{d\Phi_{m+1}} d\sigma_{N3LO}^{RVV} + \int_{d\Phi_m} d\sigma_{N3LO}^{VVV}$$

@ LO :

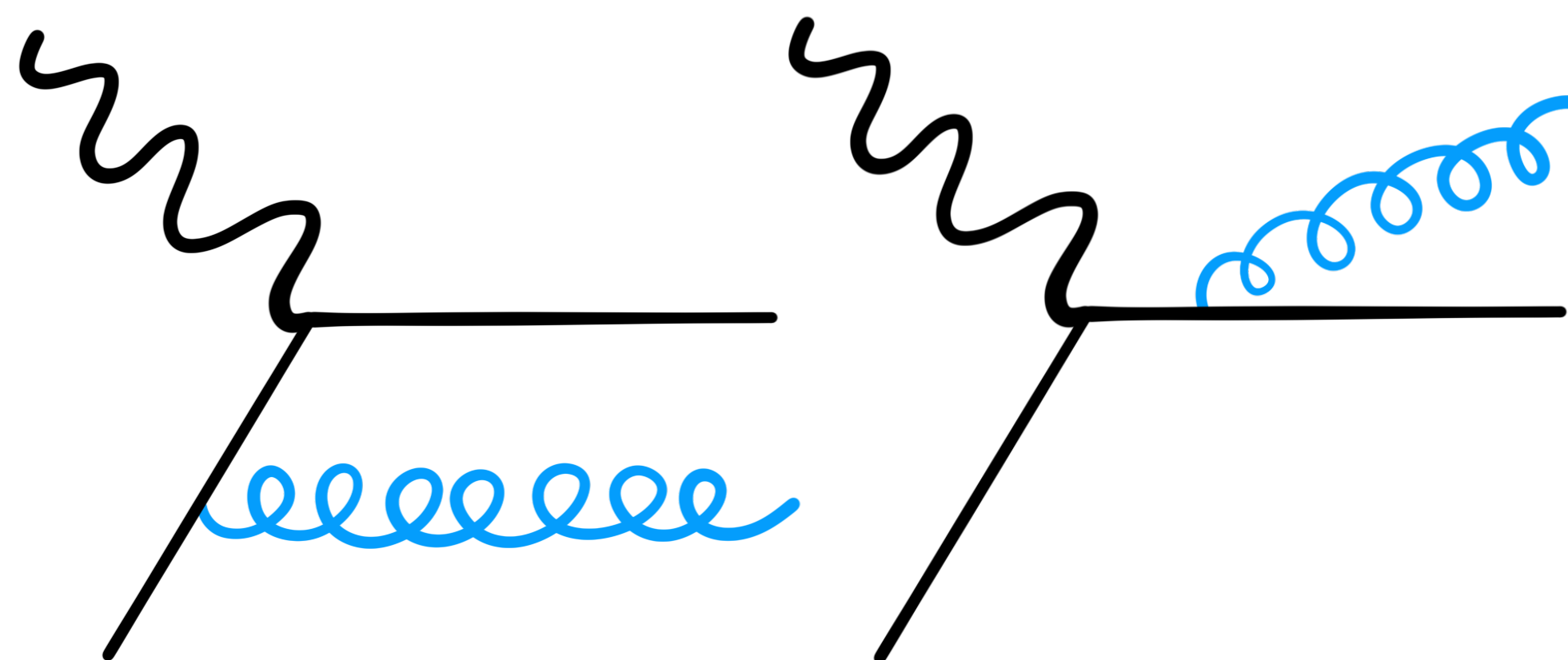


$$\gamma q \rightarrow q$$

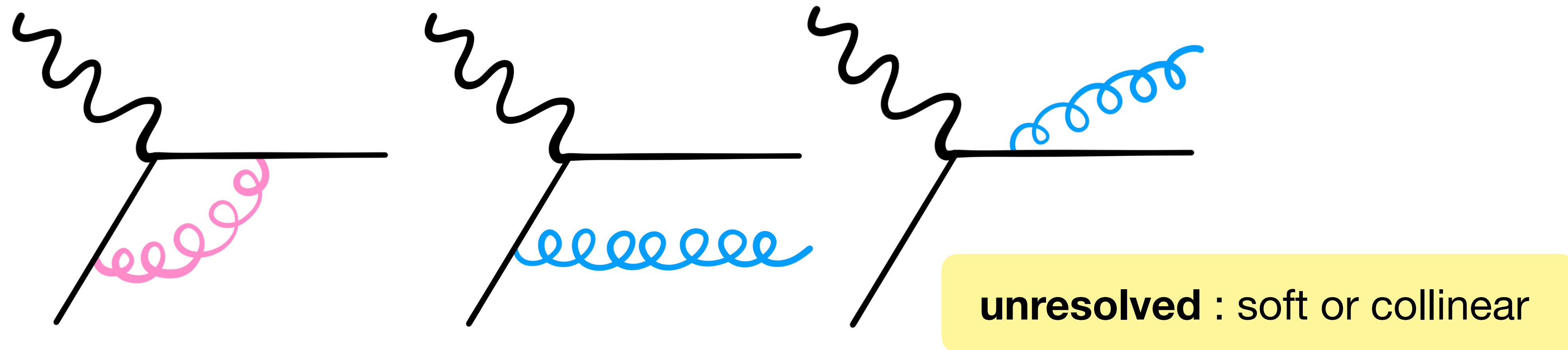
@ NLO corrections: all contributions must be taken into account



Integrate over the loop momentum k



Radiation of an extra gluon



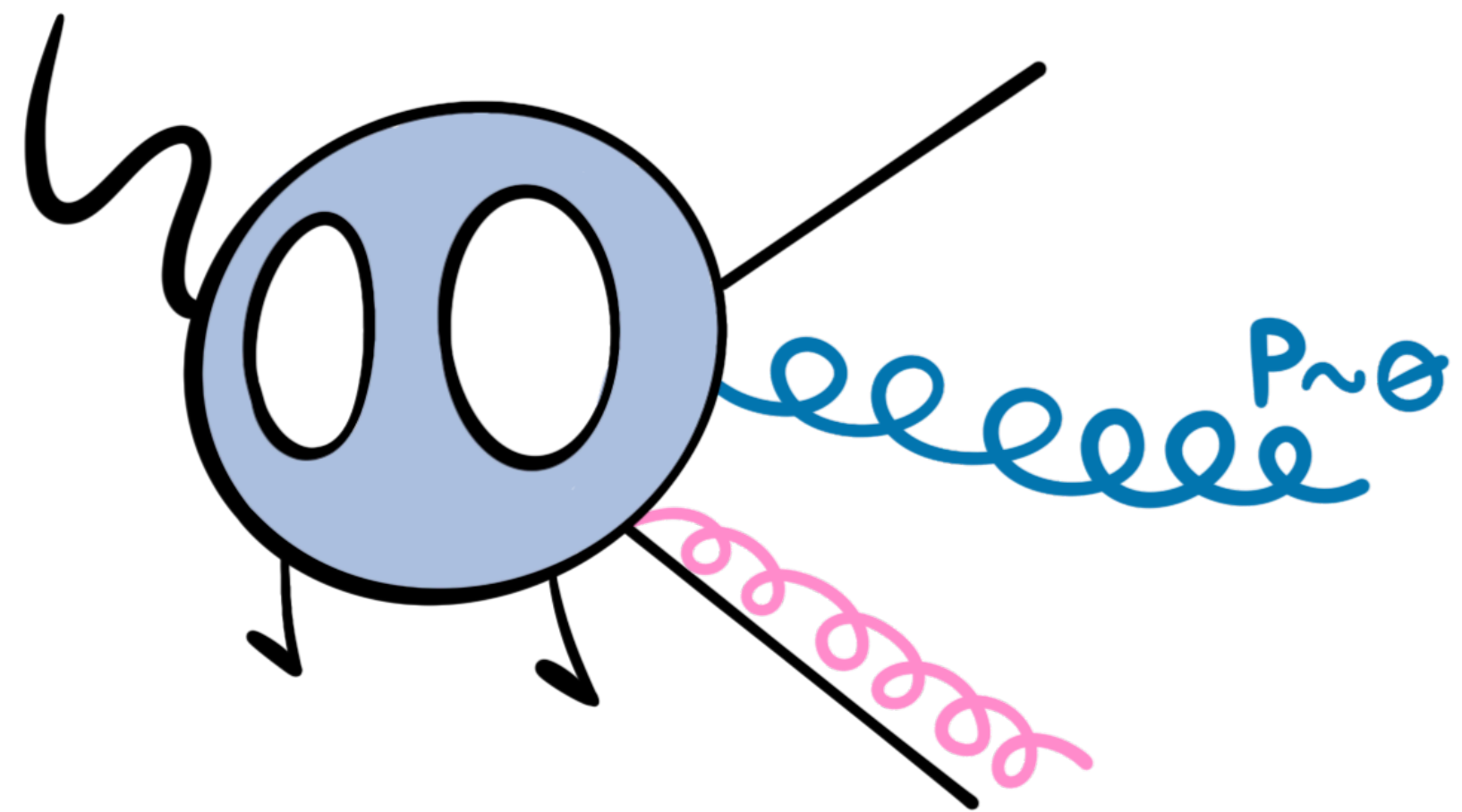
the extra gluon can become unresolved :
must add real radiation!

KLN thm

[Kinoshita 1962 ; Lee, Nauenberg 1964]

finiteness when summing over all unresolved configurations

- Separate pieces are IR-divergent:
 - **Explicit** poles in ϵ after **loop** integration
 - **Implicit** divergencies from **real** radiation



**How do we deal with
these divergencies?**

NLO example

⊖ Hard to solve analytically ⊖

$$d\sigma_{NLO} = \int_{d\Phi_{m+1}} d\sigma_{NLO}^R + \int_{d\Phi_m} d\sigma_{NLO}^V$$

- finite

$$\int_{d\Phi_{m+1}} d\sigma_{NLO}^R$$

∞

+

$$\int_{d\Phi_m} d\sigma_{NLO}^V$$

∞

NLO example

⊖ Hard to solve analytically ⊖

$$d\sigma_{NLO} = \int_{d\Phi_{m+1}} d\sigma_{NLO}^R + \int_{d\Phi_m} d\sigma_{NLO}^V$$

• finite

$$\int_{d\Phi_{m+1}} d\sigma_{NLO}^R - \int_{d\Phi_{m+1}} d\sigma_{NLO}^S$$

• finite

+

$$\int_{d\Phi_m} d\sigma_{NLO}^V + \int_{d\Phi_{m+1}} d\sigma_{NLO}^S$$

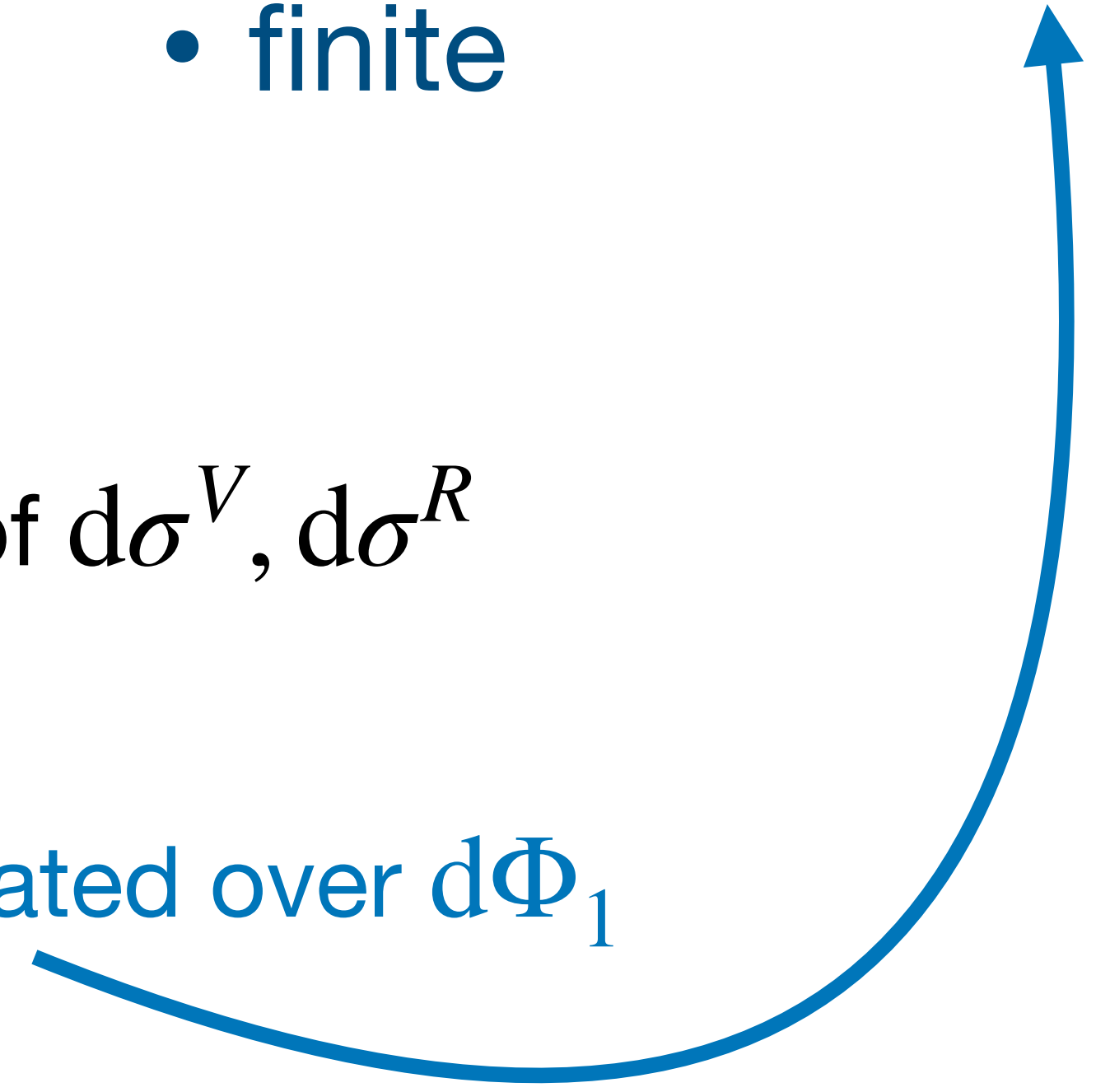
• finite

Subtraction Schemes

$$d\sigma_{NLO} = \int_{d\Phi_{m+1}} (d\sigma_{NLO}^R - d\sigma_{NLO}^S) + \int_{d\Phi_m} \left[d\sigma_{NLO}^V + \int_1 d\sigma_{NLO}^S \right]$$

• finite • finite

- Add and subtract the same quantity $d\sigma^S$
 - Mimics singular behaviour in IR-limits of $d\sigma^V, d\sigma^R$
 - Makes the integrals individually finite
 - Simple enough to be analytically integrated over $d\Phi_1$



Antenna Subtraction scheme

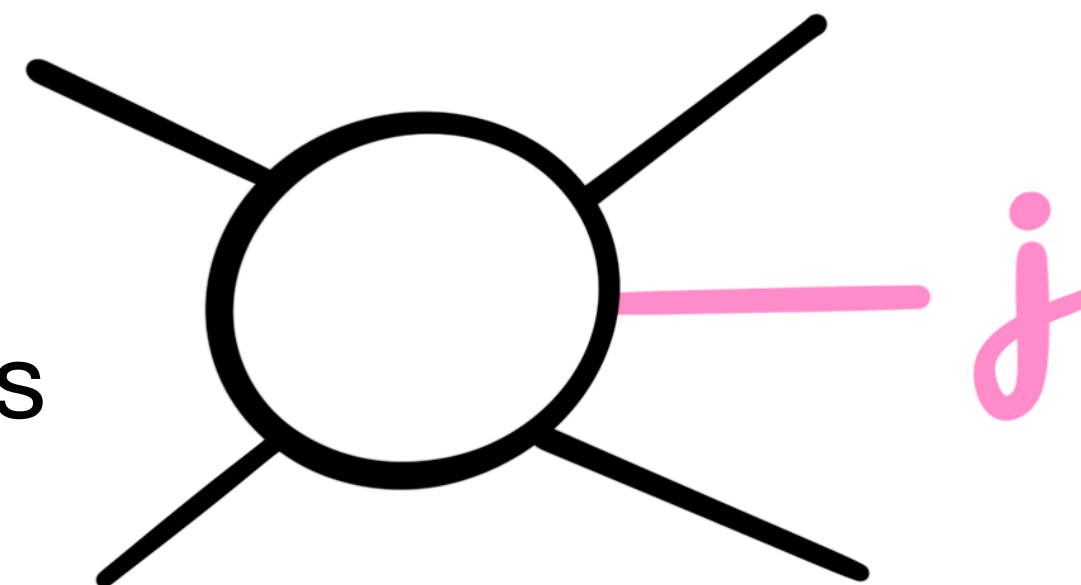
[Gehrmann-De Ridder, Gehrmann, Glover (2007)]

• Antenna functions

- Built from simple matrix elements
- Mimic the divergent behaviour in singular limits
- Can be easily integrated over phase space

$$d\sigma_{NLO}^S \sim X_{2+ \text{ extra radiation}}^\ell \tilde{M}^\ell_{\text{hard partons}} J_m$$

- Exploit factorisation of matrix elements in IR limits



What can happen?

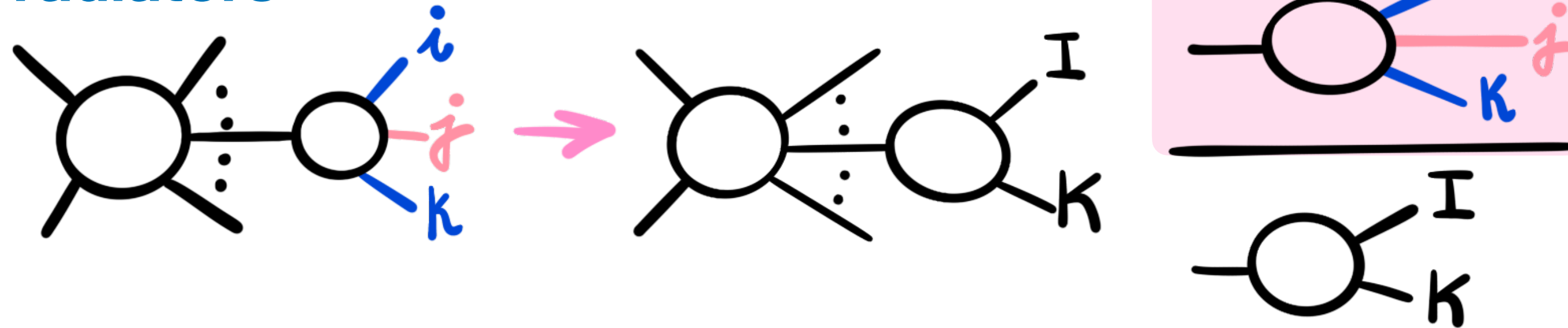
Matrix element with an extra radiation j



IR limit factorization

$$j \parallel i, \quad j \parallel k, \quad j \text{ soft}$$

Final state hard radiators

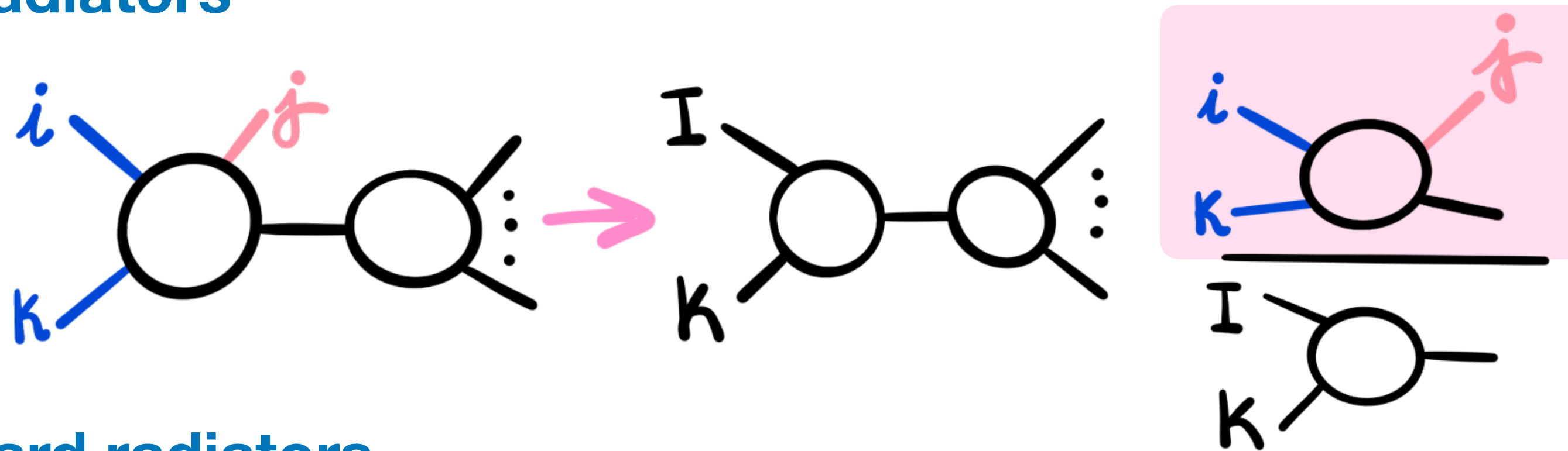


Annihilation into
hadrons

Solved! :)

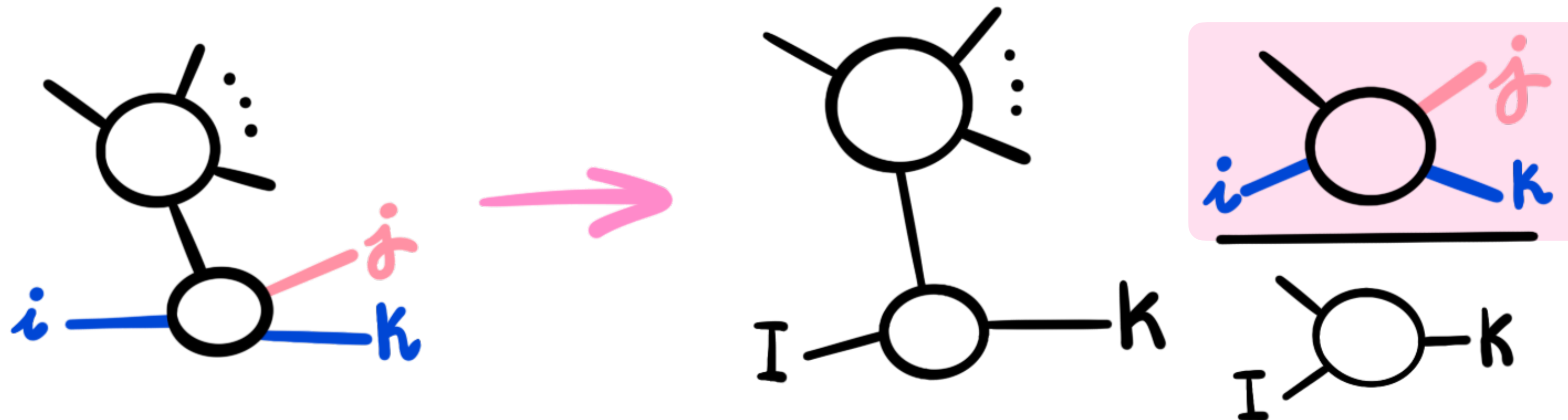
[Chen, Jakubčík, Marcoli,
Stagnitto '23]

Initial state hard radiators



Drell-Yan

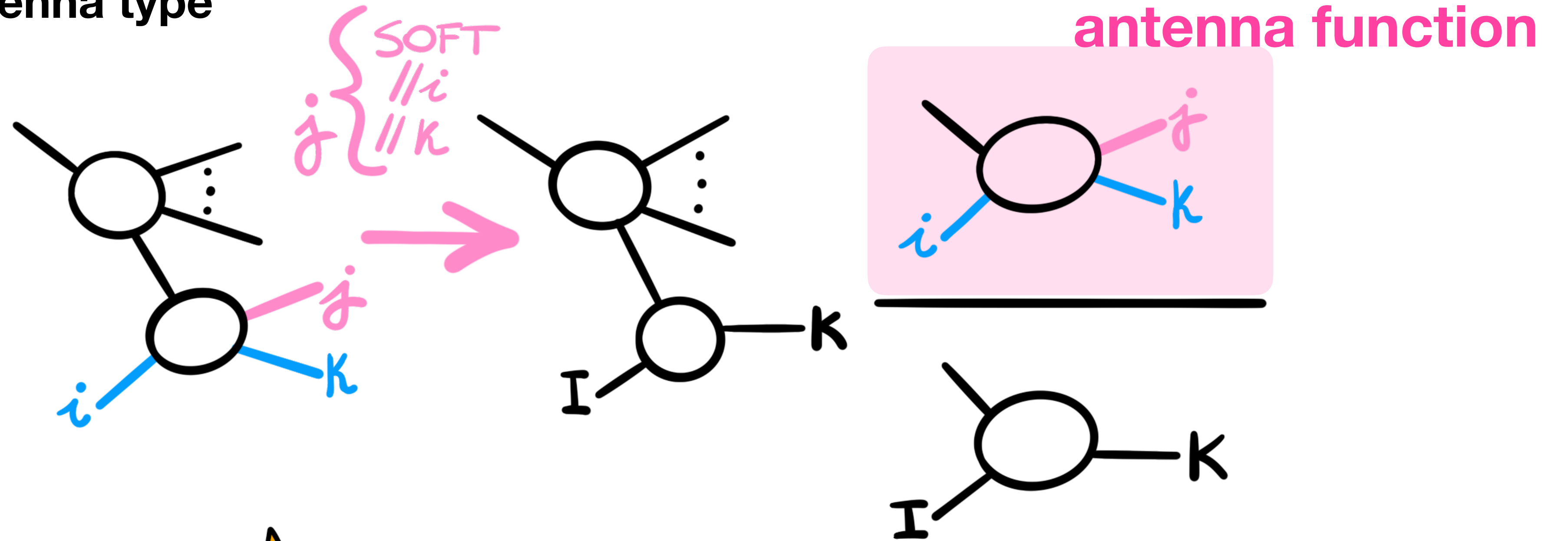
Initial-final state hard radiators



Deep Inelastic
Scattering

• Antenna functions

Focus: initial-final antenna type



off-shell current



initial-state parton



extra-radiation parton

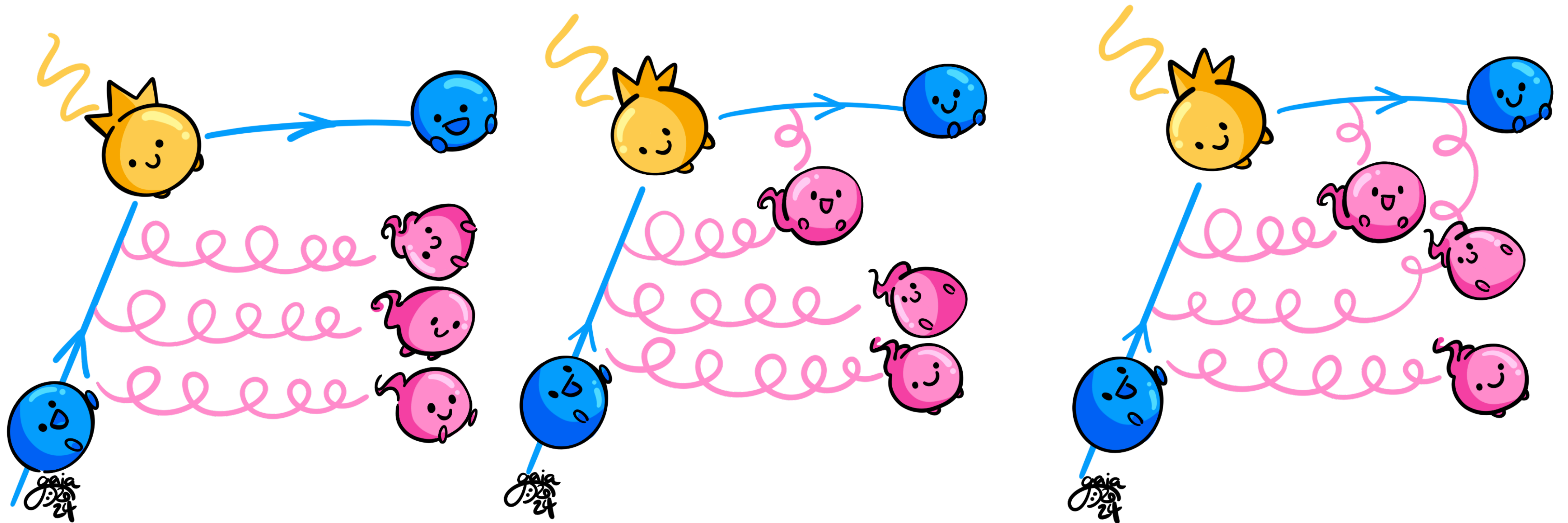


final-state parton



⇒ DIS-like process

Towards N3LO RRR antennae

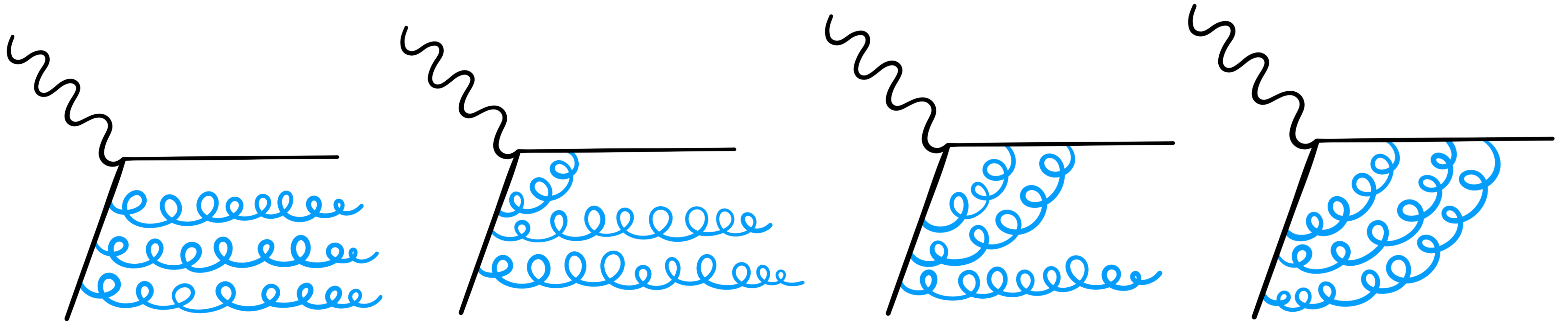


Deep Inelastic Scattering N3LO

Representative contributions at order α_s^3

$$q_1 + q_2 \rightarrow p_1 + p_2 + (p_3) + (p_4)$$

$$q_2^2 = -Q^2 < 0, \quad q_1^2 = 0, \quad p_i^2 = 0, \quad i = 1, 2, 3, 4$$



RRR

2 → 4

RRV

2 → 3

RVV

2 → 2

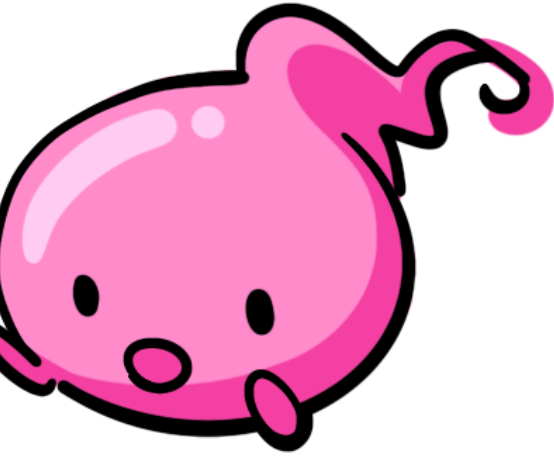
VVV

2 → 1

Workflow

Phase space integral

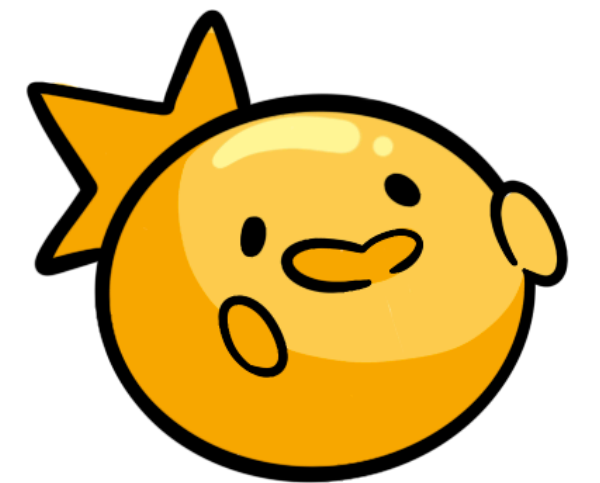
Reverse
unitarity



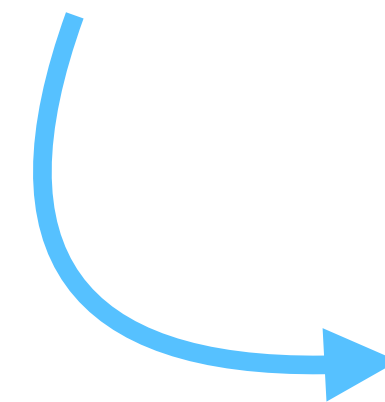
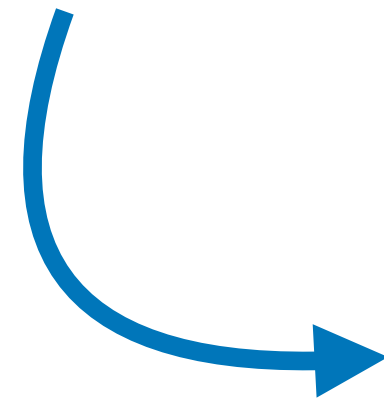
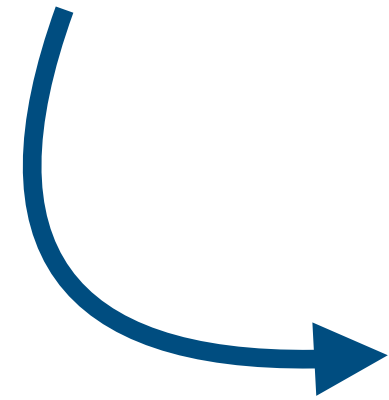
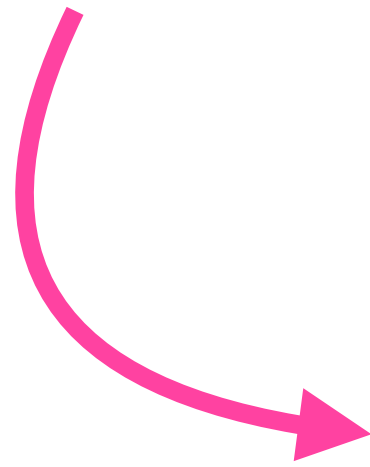
Reduction to
master integrals



DE & canonical form



Boundaries



Reverse Unitarity

[Anastasiou, Melnikov 2002]

phase space \rightarrow (cut) loops

$$-2\pi i \delta^+(p_i^+) = \frac{1}{p_i^2 + i0^+} - \frac{1}{p_i^2 - i0^+} = \frac{1}{[p_i^2]_{cut}}$$

$$I_{RRR} = \int d\Phi_4 (2\pi)^D \delta^D \left(q_1 + q_2 - \sum_{i=1}^4 p_i \right) \prod_j \frac{1}{D_j^{\alpha_j}}$$

Notice!

$$d\Phi_n = \prod_{i=1}^n \frac{d^d p_i}{(2\pi)^d} \delta^+(p_i^2)$$

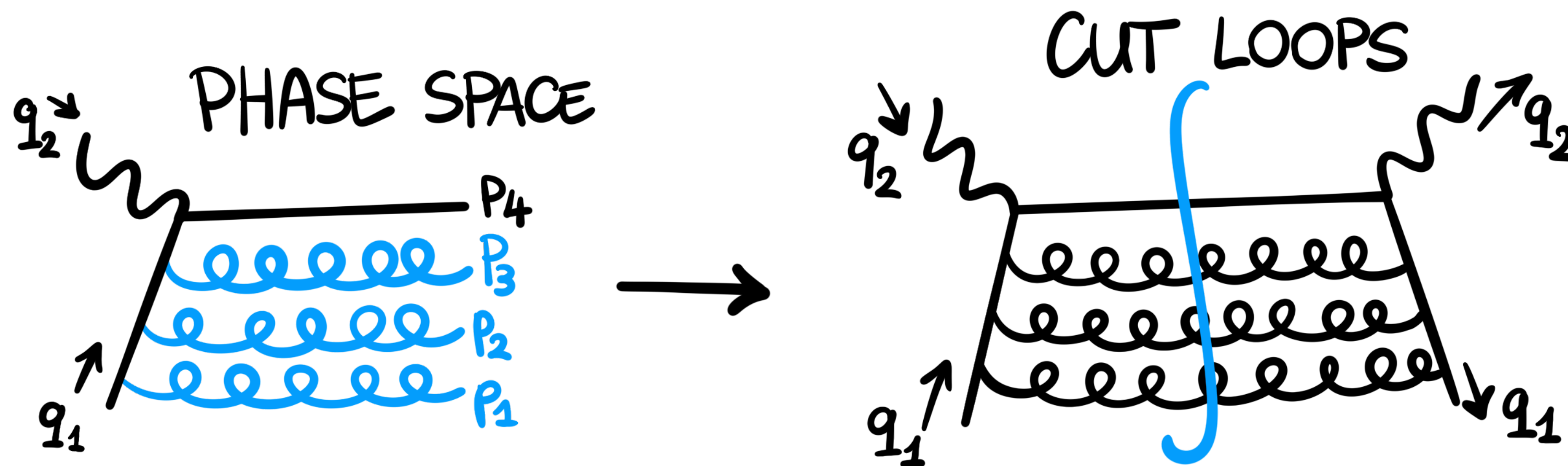
$$I_{RRR} = \int \frac{d^D p_1}{(2\pi)^D} \int \frac{d^D p_2}{(2\pi)^D} \int \frac{d^D p_3}{(2\pi)^D} \frac{1}{[p_1^2]_{cut}} \frac{1}{[p_2^2]_{cut}} \frac{1}{[p_3^2]_{cut}} \frac{1}{[p_4^2]_{cut}} \prod_j \frac{1}{D_j^{\alpha_j}}$$

Reverse Unitarity

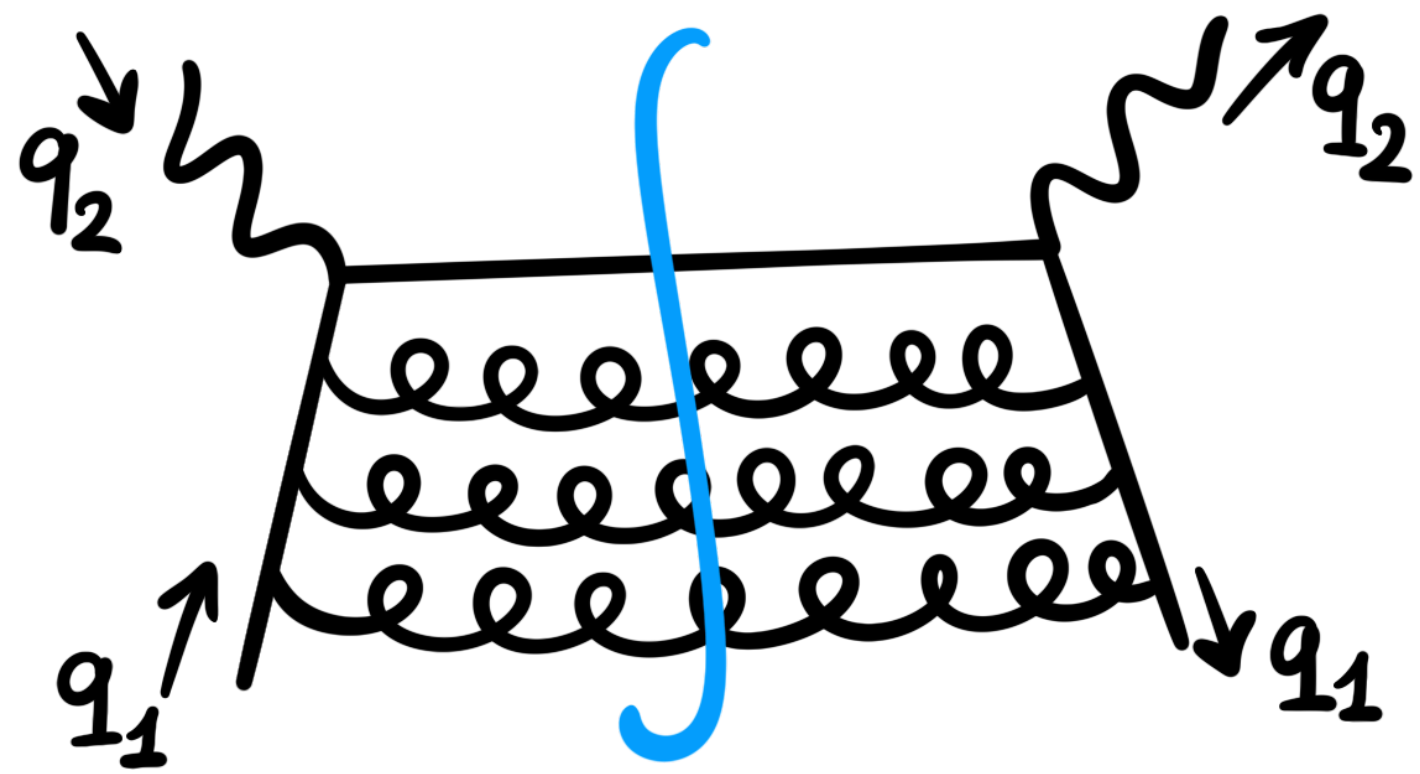
phase space \rightarrow (cut) loops

$$-2\pi i \delta^+(p_i^+) = \frac{1}{p_i^2 + i0^+} - \frac{1}{p_i^2 - i0^+} = \frac{1}{[p_i^2]_{cut}}$$

Diagrammatically



$$\int d\Phi_4 (2\pi)^D \delta^D(P+q - \sum_{i=1}^4 p_i) \prod_j \frac{1}{D_j^{a_j}} \rightarrow \int \prod_{i=1}^3 \frac{d^D p_i}{(2\pi)^D} \frac{1}{p_1^2 p_2^2 p_3^2 p_4^2} \prod_j \frac{1}{D_j^{a_j}}$$



This DIS amplitude contains a lot of integrals $\{I_j\}$

→ how to make things better?

Reduction to Master integrals

[Chetyrkin, Tkachov '81; Laporta 2000]

Reduction into a basis of linearly independent **master integrals**

$$\{g_j\} \subset \{I_j\}$$

$$I_j = \sum_k c_{jk} g_k$$

rational coefficients
master integrals

modulo identities:

- Integration By Parts
- Lorentz Invariance
- symmetry relations

DE for Feynman integrals

[Barucchi, Ponzano '73; Kotikov '91; Bern, Dixon, Kosower '94; Gehrmann, Remiddi 2000]

Derivative of MI with respect to external invariants

$$\partial_z g_i = \sum_j a_{ij} I_j$$

Use IBP relations to rewrite RHS

$$I_j = \sum_k c_{jk} g_k$$

Obtain a system of first order DE for the MI!

$$\partial_z g_i = \sum_{jk} a_{ij} c_{jk} g_k$$

$$\partial_z \vec{g} = A(\epsilon, z) \cdot \vec{g}$$

How to solve a differential equation:

- Generic solution
- Boundary condition

Rewrite the DE in **canonical form** [Henn 2013]: solution in terms of iterated integrals

$$\partial_z \vec{g} = \epsilon A^\star(z) \cdot \vec{g}$$

Boundary Conditions

- Consistency conditions
 - Finding relations between boundaries
- Evaluation in some kinematic limit
 - Fix the remaining ones



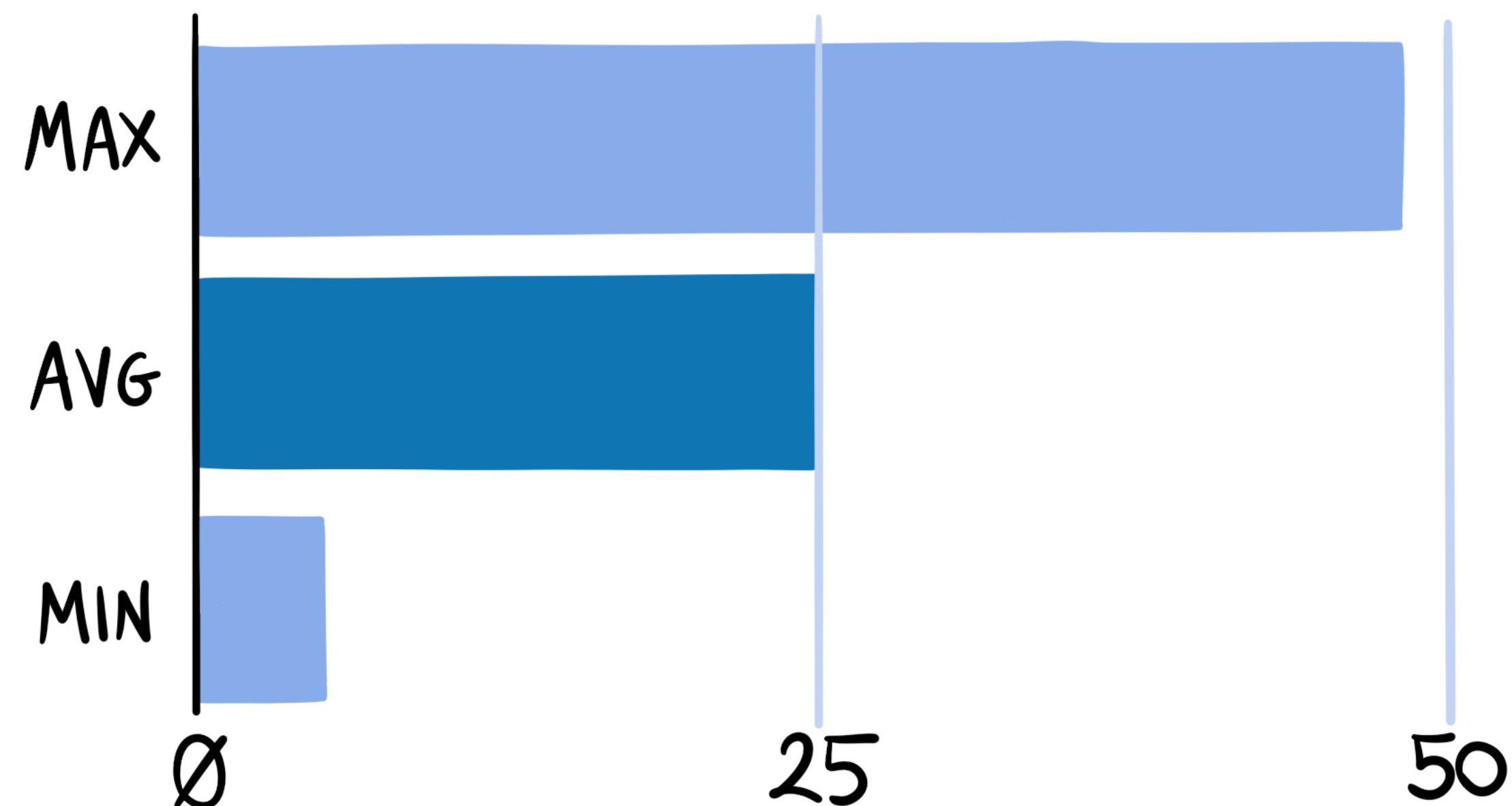
Getting to know the RRR families

Physical 4-cuts of the 3 loop inclusive DIS amplitude

$$I_{RRR} = \int \frac{d^D p_1}{(2\pi)^D} \int \frac{d^D p_2}{(2\pi)^D} \int \frac{d^D p_3}{(2\pi)^D} \frac{1}{[p_1^2]_{cut}} \frac{1}{[p_2^2]_{cut}} \frac{1}{[p_3^2]_{cut}} \frac{1}{[p_4^2]_{cut}} \prod_j \frac{1}{D_j^{\alpha_j}}$$

- 65 families
- Few number of MI for each family → **4 cuts**
- **Total: 1620 MIs** (No symmetries between families included)

#MI PER FAMILY



- DE matrix M is a function of $M(z, \epsilon)$ [Lee 2020]
⇒ **playground for automatic tools! eg LIBRA**
- Analysis of their parametric representation to simplify the DE & get to a canonical form

[Henn 2013]

Results

Canonical DE for all the families ✓

We can find a generic solution



Boundary conditions →
SOON

[Liu, Ma 2022]

- Numerical evaluation with **AMF Low @ 200 digits** (~80% done...) & **PSLQ**
- Constraints from symmetry relations between the families
- Calculation of the amplitude → which boundaries are actually needed

Outlook:

- Extend calculation to RVV and RRV layers
- Ultimate goal: obtaining the full set of integrated initial-final antennae

Thank you for your attention!

