

A new framework on global analysis of fragmentation functions





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- based on 2305.14620 with ChongYang Liu, XiaoMin Shen, Bin Zhou and 2401.02781 with ChongYang Liu, XiaoMin Shen, HongXi Xing, YuXiang Zhao
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Outline

♦ 1. Introductions

◆ 2. Fragmentation functions and QCD factorization

◆ 3. A new global analysis on FFs to light charged hadrons

♦ 4. Outlook and Summary

QCD at its 50 years

simulations

running coupling constant



[PDG 2022]

• Quantum Chromodynamics is a beautiful theory with enormous success, rich phenomena and powerful predictions; yet still more to come after 50 years with future facilities and developments of lattice

three-point energy correlator inside jet

$$E3C = \frac{d\sigma^{[3]}}{dx_L} = \sum_{i,j,k}^n \int d\sigma \frac{E_i E_j E_k}{E^3} \delta(x_L - \max(\Delta R_{i,j}, \Delta R_{i,k}, \Delta R_{j,k}))$$

$$CMS Preliminary 36.3 \text{ fb}^{-1} (13 \text{ TeV})$$
Free hadron Confinement Perturbative
$$Data$$

[SMP-22-015]

Observables at colliders

color-neutral hadrons appearing in both the initial and final states

inclusive cross sections at pp collisions



Standard Model Total Production Cross Section Measurements

• QCD predictions to observables rely on both perturbative calculations describing interactions of highenergy quarks and gluons, as well as non-perturbative inputs and hadronization corrections since only





Jet charge and flavor-tagging

fragmentation functions or from models implemented in MC generators



u/d quark separation [2103.09649]

+ Jet charge is a typical observable requiring knowledge on transition of parton to hadrons, especially the distribution of electric charges to hadrons; can be calculated in QCD from first principle based on



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Single inclusive hadron production

e.g., from single-inclusive annihilation (SIA), semi-inclusive DIS (SIDIS), pp collisions

full description of final state hadrons



rely on model and PS at the lowest order

◆ In its simplest form, fragmentation functions (FFs) describe number density of the identified hadron wrt the fraction of momentum of the initial parton it carries, as measured in single inclusive hadron production,

single inclusive hadron production/observable

[1607.02521]



$$\frac{1}{\sigma_{\text{tot}}} \frac{d\sigma^{e^+e^- \to hX}}{dz} = F^h(z, Q^2), \quad z = \frac{2E_h}{\sqrt{s}}$$

exp. definition of unpolarized collinear FFs

other forms: polarized FFs, TMD FFs, di-hadron FFs

QCD collinear factorization



$$\frac{1}{\sigma_{\text{tot}}} \frac{d\sigma^{e^+e^- \to hX}}{dz} = \frac{1}{\sum_q e_q^2} \left(2F_1^h(z, Q^2) + F_L^h(z, Q^2) \right)$$

$$2F_1^h(z,Q^2) = \sum_q e_q^2 \left(D_1^{h/q}(z,Q^2) + \frac{\alpha_s(Q^2)}{2\pi} \left(C_1^q \otimes D_1^{h/q} + C_1^g \otimes D_1^{h/g} \right)(z,Q^2) \right)$$

$$\frac{d^3 \sigma^{\ell p \to \ell h X}}{dx \, dy \, dz} = \frac{2\pi \alpha_{\rm em}^2}{Q^2} \left(\frac{1 + (1 - y)^2)}{y} \, 2F_1^h(x, z, Q^2) + \frac{2(1 - y)}{y} \, F_L^h(x, z, Q^2) \right)$$

$$2F_1^h(x, z, Q^2) = \sum_q e_q^2 \left(f_1^{q/p} D_1^{h/q} + \frac{\alpha_s(Q^2)}{2\pi} \left(f_1^{q/p} \otimes C_1^{qq} \otimes D_1^{h/q} + f_1^{q/p} \otimes C_1^{qq} \otimes D_1^{h/q} \right) \right],$$

+ QCD collinear factorization ensures universal separation of long-distance and short-distance contributions in high energy scatterings involving initial/final state hadrons, and enables predictions on cross sections



- FFs/PDFs, reveal inner structure of hadrons; nonperturbative (NP) origin, universal, e.g. DIS vs. pp collisions; fitted from data
- * runnings of FFs/PDFs with μ_D/μ_f are governed by the DGLAP equation

unpolarized collinear FFs, operator definition

$$D_{1}^{h/q}(z) = \frac{z}{4} \sum_{X} \int \frac{d\xi^{+}}{2\pi} e^{ik^{-}\xi^{+}} \operatorname{Tr} \left[\langle 0 | \mathcal{W}(\infty^{+}, \xi^{+}) \psi_{q}(\xi^{+}, 0^{-}, \vec{0}_{T}) | P_{h}, S_{h}; X | \chi_{q}(0^{+}, 0^{-}, \vec{0}_{T}) \mathcal{W}(0^{+}, \infty^{+}) | 0 \rangle \gamma^{-} \right].$$

$$\frac{d}{d\ln\mu^2} D_1^{h/i}(z,\mu^2) = \frac{\alpha_s(\mu^2)}{2\pi} \sum_j \int_z^1 \frac{du}{u} P_{ji}(u,\alpha_s(\mu^2)) D_1^{h/j}\left(\frac{z}{u},\mu^2\right)$$

[Collins, Soper, Sterman]

 $X\rangle$

FMNLO (fragmentation at NLO in QCD)

- accurate to NLO in QCD
 - automation of fragmentation calculations from arbitrary hard processes at NLO, within SM and BSMs via MG5_aMC@NLO
 - fast convolution algorithms of partonic cross sections with FFs without repeating the time consuming MC integrations
 - future goal/generalizations: transverse observables, NNLO corrections





2023.05: **FMNLOv1.0** first release of FMNL0 interfaced with MG5_aMC@NL0.

+ FMNLO is a new program for automated and fast calculations of fragmentation cross sections of arbitrary processes. It is based on a hybrid scheme of phase-space slicing method and local subtraction method,

https://fmnlo.sjtu.edu.cn/~fmnlo/

QCD inclusive dijets at LHC

[JG+, 2305.14620]

Z-boson tagged jet



Global data and phenomenological analysis

phenomenological FFs from global analysis at NLO/NNLO in QCD

single incl. production $\overleftarrow{\mathfrak{G}}^{\mathfrak{f}_{10}}$ unidentified $\overset{20}{\mathsf{r}}_{\mathsf{T}}^{\mathsf{r}} < 30 \, \text{GeV}$ Jet fragmentation to $\overleftarrow{\mathfrak{G}}^{\mathfrak{f}_{10}}$ with the second se charged hadrons (LHCb) charged hadrons (SIA & SIDIS)



◆ Measurements are available from colliders SLAC, LEP, HERA, RHIC, LHC and fixed-target HERMES, COMPASS experiments for various charged hadrons as well neutral hadrons; several major groups provide

 $50 < p_{\rm T}^{\rm jet} < 100 \, {\rm GeV}$

global analysis

* $m_{ajp_{T}} g_{T} g_$

mostly done¹⁰ at NLO in QCD since exact NNLO₁ coefficient functions not known for SIDI^{Sythia}⁸ pp

* different determination can be quite different derestor sedection of data sets as well as the ory treatments, not converge as well as the case of PDF fits

 10^{-2}

[1607.02521 for a review]

 10^{-1}

 \mathcal{Z}

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A new global analysis of FFs

parametrization of FFs to charged pion/kaon/ proton at initial scale (Q=5 GeV):

$$zD_{i}^{h}(z,Q_{0}) = z^{\alpha_{i}^{h}}(1-z)^{\beta_{i}^{h}} \exp\left(\sum_{n=0}^{m} a_{i,n}^{h}(\sqrt{z})^{n}\right)$$

parton-to- π^+	favored	α	β	a_0	a_1	a	$^{\prime}2$	d.o.f.
u	Y							5
$d \simeq u$	Y	_	_		_	-	_	1
$\bar{u} = d$	N					2	x	4
$s = \bar{s} \simeq \bar{u}$	N	_				2	X	3
$c = \bar{c}$	N					Σ	X	4
$b = \overline{b}$	N					2	x	4
g	N		F					4

parton-to- K^+	favored	lpha	$ \beta $	a_0	$ a_1 $	a_2	d.o.f.
u	Y					X	4
$\overline{s} \simeq u$	Y	-	-		-	X	1
$\bar{u} = d = d = s$	Ν					X	4
$c = \overline{c}$	Ν					X	4
$b = \overline{b}$	Ν					x	4
g	Ν		F			X	3

parton-to-p	favored	lpha	β	a_0	a_1	a_2	d.o.f.
u = 2d	Y					х	4
$\bar{u} = \bar{d} = s = \bar{s}$	Ν				x	x	3
$c = \overline{c}$	Ν					x	4
$b = \overline{b}$	Ν					x	4
g	N		F			х	3

• Establishing a new framework on global analysis of fragmentation functions to identified charged hadrons, including charged pion, kaon and proton, using most recent data from SIA, SIDIS, and pp collisions

- * a joint determination of FFs to charged pion, kaon and proton at NLO in QCD (63 parameters) including estimation of uncertainties with Hessian sets
- * apply a strong selection criteria on the kinematics of fragmentation processes to ensure validity of LT factorization and perturbative calculations (z>0.01 and E_h/ $p_{T,h}>4$ GeV)
- including theory uncertainties (residual) scale variations) into the covariance matrix
- suse fast interpolation techniques for calculations of cross sections which largely increase efficiency of the global fit

[JG+, 2401.02781 and work in preparation]

Selection of data

productions, due to the development of FMNLO

LHC measurements for hadron inside jet measurements (jet fragmentation)

exp.	$\sqrt{s}(\text{TeV})$	luminosity	hadrons	final states	R_{j}	cuts for jets/hadron	observable	$N_{ m pt}$
ATLAS[5]	5.02	25 pb^{-1}	h^{\pm}	$\gamma + j$	0.4	$\Delta \phi_{j,\gamma} > \frac{7\pi}{8}$	$rac{1}{N_{ m jet}}rac{dN_{ m ch}}{dp_{T,h}}$	6
CMS[6]	5.02	27.4 pb^{-1}	h^{\pm}	$\gamma + j$	0.3	$\Delta \phi_{j,\gamma} > \frac{7\pi}{8}, \Delta R_{h,j} < R_j$	$rac{1}{N_{ m jet}}rac{dN_{ m ch}}{d\xi}$	4
ATLAS[7]	5.02	260 pb^{-1}	h^{\pm}	Z + h	no jet	$\Delta \phi_{h,Z} > \frac{3}{4}\pi$	$\frac{1}{n_Z} \frac{dN_{\rm ch}}{dp_{T,h}}$	9
CMS[8]	5.02	320 pb^{-1}	h^{\pm}	Z + h	no jet	$\Delta \phi_{h,Z} > \frac{7}{8}\pi$	$\frac{1}{n_Z} \frac{dN_{\rm ch}}{dp_{T,h}}$	11
LHCb[9]	13	$1.64 \ {\rm fb}^{-1}$	π, K, p	Z+j	0.5	$\Delta \phi_{j,\gamma} > \frac{7\pi}{8}, \Delta R_{h,j} < R_j$	$\frac{1}{n_Z} \frac{dN_{\rm ch}}{d\zeta}$	20
ATLAS[10]	5.02	25 pb^{-1}	h^{\pm}	inc. jet	0.4	-	$rac{1}{N_{ m jet}}rac{dN_{ m ch}}{d\zeta}$	63
ATLAS[11]	7	36 pb^{-1}	h^{\pm}	inc. jet	0.6	$\Delta R_{h,j} < R_j$	$rac{1}{N_{ m jet}}rac{dN_{ m ch}}{d\zeta}$	103
ATLAS[12]	13	$33 { m ~fb}^{-1}$	h^{\pm}	dijet	0.4	$p_T^{\text{lead}}/p_T^{\text{sublead}} < 1.5$	$\frac{1}{N_{\text{iet}}} \frac{dN_{\text{ch}}}{d\zeta}$	280

- * LHC measurements on hadron inside jet provide essential inputs for u/d/g flavor separation with wide kinematic coverages, both in energy scale Q and in momentum fraction z
- ♦ In dijets or inclusive jets production, low p_T and central (high p_T and forward) jets are mostly initiated by g(u-quark); Z or photon tagged jets are more likely from u/d quarks

• For the first time the jet fragmentation data from LHC have been incorporated into the global analysis of FFs to light charged hadrons, including from processes of incl. jet, dijet, Z or photon tagged jet

kinematic/flavor coverage (LO) for ATLAS jet fragmentation





Selection of data

production in SIDIS from HERA and COMPASS, for identified or unidentified charged hadrons

incl. hadron production at RHIC and LHC (pp)

exp.	$\sqrt{s_{NN}}$ (TeV)	# events (million)	$p_{T,h}$	hadrons	observable	$N_{\rm pt}$
ALICE[1]	13	40-60(pp)	[2, 20] GeV	π, K, p, K_S^0	$K/\pi, p/\pi, K_S^0/\pi$	49
ALICE[1]	7	150(pp)	[3, 20] GeV	π, K, p	13TeV/7TeV for π, K, p	37
ALICE[2]	5.02	120(pp)	[2, 20] GeV	π, K, p	$K/\pi, p/\pi$	34
ALICE[3]	2.76	40(pp), 15(Pb-Pb)	[2, 20] GeV	π, K, p	$K/\pi, p/\pi$	27
STAR[4]	0.2	14(pp)	[3, 15] GeV	π, K, p, K_S^0	$K/\pi, p/\pi^+, \bar{p}/\pi^-, K_S^0/\pi, \pi^-/\pi^+, K^-/K^+$	60

incl. hadron production at Z-pole (SIA)

exp.	\sqrt{s}	$lum.(n_Z)$	final states	hadrons	$N_{\rm pt}$
OPAL[19]	m_Z	780 000	$Z \to q\bar{q}$	π^{\pm}, K^{\pm}	20
ALEPH	m_Z	520 000	$Z \to q\bar{q}$	$\pi^{\pm}, K^{\pm}, p(\bar{p})$	42
DELPHI	m_Z	$1 \ 400 \ 000$	$Z \to q\bar{q}$	$\pi^{\pm}, K^{\pm}, p(\bar{p})$	39
-	-	-	$Z \to b\overline{b}$	$\pi^{\pm}, K^{\pm}, p(\bar{p})$	39
SLD	m_Z	400 000	$Z \to q\bar{q}$	$\pi^{\pm}, K^{\pm}, p(\bar{p})$	66
-	-	-	$Z \to b\overline{b}$	$\pi^{\pm}, K^{\pm}, p(\bar{p})$	66
-	-	-	$Z \to c\bar{c}$	$\pi^{\pm}, K^{\pm}, p(\bar{p})$	66
TASSO	34GeV	$77 { m pb}^{-1}$	inc. had.	$\pi^{\pm}, K^{\pm}, p(\bar{p})$	3
TASSO	44GeV	34 pb^{-1}	inc. had.	π^{\pm},π^{0}	5
TPC	29GeV	$70 {\rm \ pb^{-1}}$	inc. had.	π^{\pm}, K^{\pm}	12
OPAL	201.7GeV	433 pb^{-1}	inc. had.	h^{\pm}	17
DELPHI	189GeV	157.7 pb^{-1}	inc. had.	$\pi^{\pm}, K^{\pm}, p(\bar{p})$	9

• Other data include ratios of inclusive production rates of different hadrons measured in pp collisions, single incl. hadron production from SIA (w/wo heavy-flavor tagging) mostly at Z-pole, and incl. hadron

incl. hadron production at HERA and COMPASS (SIDIS)

exp.	$\sqrt{s}(\text{GeV})$	luminosity	kinematic cuts	hadrons	obs	$N_{\rm pt}$
H1[13]	318	44 pb^{-1}	$Q^2 \in [175, 20000] \text{ GeV}^2$	h^{\pm}	$D \equiv \frac{1}{N} \frac{dn_{h\pm}}{dx_{p}}$	16
H1[14]	318	44 pb^{-1}	$Q^2 \in [175,8000] \text{ GeV}^2$	h^{\pm}	$A \equiv \frac{D^+ - D^-}{D^+ + D^-}$	14
ZEUS[15]	300,318	440 pb^{-1}	$Q^2 \in [160, 40960] \text{ GeV}^2$	h^{\pm}	D	32
COMPASS[16, 17]	17.3	540 pb^{-1}	$x \in [0.14, 0.4], y \in [0.3, 0, 5]$	π, K, h	$\frac{dM^{h}}{dz}$	124
COMPASS[18]	17.3	-	$x \in [0.14, 0.4], y \in [0.3, 0, 5]$	π, K, p	$\frac{dM^{h}}{dz}$	97



determination of Hessian uncertainties

overall agreement: **X**² breakdown to sub-groups for the best-fit

Experiments	N_{pt}	χ^2	χ^2/N_{pt}
ATLAS jets [†]	446	350.8	0.79
ATLAS Z/γ +jet [†]	15	31.8	2.12
CMS Z/γ +jet [†]	15	17.3	1.15
LHCb Z +jet	20	30.6	1.53
ALICE inc. hadron	147	150.6	1.02
STAR inc. hadron	60	42.2	0.70
pp sum	703	623.3	0.89
TASSO	8	7.0	0.88
TPC	12	11.6	0.97
OPAL	20	16.3	0.81
OPAL (202 GeV) †	17	24.2	1.42
ALEPH	42	31.4	0.75
DELPHI	78	36.4	0.47
DELPHI (189 GeV)	9	15.3	1.70
SLD	198	211.6	1.07
SIA sum	384	353.8	0.92
H1 '	16	12.5	0.78
H1 (asy.) †	14	12.2	0.87
ZEUS †	32	65.5	2.05
COMPASS $(06I)$	124	107.3	0.87
COMPASS $(16p)$	97	56.8	0.59
SIDIS sum	283	254.4	0.90
Global total	1370	1231.5	0.90

the global data sets (1370 points in total) are found, χ^2/N well below 1; individual agreements to the 138 sub-datasets are also tested, motivating usage of a tolerance $\Delta \chi^2 \sim 2$ in



FFs to light charged hadrons



• We arrive at a best-fit of the charged pion, kaon and proton FFs together with 126 Hessian error FFs, two for each of the eigenvector direction; FFs are generally well constrained in the region with z~0.1-07

- ◆ our results show an uncertainty of 3%, 4% and 8% for FFs of gluon to pion at z=0.05, 0.1 and 0.3, respectively
- similarly an uncertainty of 4%, 4% and 7% for FFs of u-quark to pion, kaon and proton at z=0.3, respectively
- FFs of heavy-quarks are well constrained for z between 0.1~0.5 due to the tagged SIA events at Z-pole measurements
- ✤ a preference for larger FFs of s quark to pion possibly due to decays of short-lived strange hadrons
- high precision of gluon FFs is mostly due to the data of jet fragmentation from the LHC

FFs to light charged hadrons

clarifications



FFs (charge-summed) vs. momentum fraction

• Our new extractions on FFs are compared to previous determinations from other groups (DSS and NNFF) for the charge-summed pion, kaon and proton; large discrepancies are found and will need further

- We find general agreement between ours and DSS for FFs of u and d quarks to pion, and of s quark to kaon
- however, large discrepancies are found for FFs to protons and for FFs of gluon to all three charged hadrons
- NNFFs show larger uncertainties in general and can even become negative in some kinematic regions
- Intervalue future works involving coordinations from different groups will be needed for clarifications on discrepancies

Test of sum rules

are tested with the extracted FFs and find consistency

momentum sum rule:
$$\sum_{h} \int_{0}^{1} dz z D_{i}^{h}(z, Q) = 1$$

with finite cutoff: $\langle z \rangle_{i}^{h} = \int_{z_{min}}^{1} dz z D_{i}^{h}(z, Q)$

mom.	g(z > 0.01)	u(z > 0.01)	d(z > 0.01)	s(z > 0.088)
π^+	$0.200^{+0.008}_{-0.008}$	$0.262^{+0.017}_{-0.016}$	$0.128^{+0.020}_{-0.019}$	$0.161^{+0.013}_{-0.013}$
K^+	$0.018^{+0.004}_{-0.003}$	$0.058^{+0.005}_{-0.004}$	$0.019^{+0.004}_{-0.004}$	$0.015\substack{+0.002\\-0.002}$
p	$0.035^{+0.006}_{-0.005}$	$0.044^{+0.004}_{-0.004}$	$0.022\substack{+0.002\\-0.002}$	$0.015\substack{+0.002\\-0.002}$
π^{-}	$0.200\substack{+0.008\\-0.008}$	$0.128^{+0.020}_{-0.019}$	$0.299^{+0.054}_{-0.049}$	$0.161^{+0.013}_{-0.013}$
K^{-}	$0.018^{+0.004}_{-0.003}$	$0.019^{+0.004}_{-0.004}$	$0.019^{+0.004}_{-0.004}$	$0.205\substack{+0.014\\-0.013}$
\bar{p}	$0.035^{+0.006}_{-0.005}$	$0.019^{+0.003}_{-0.003}$	$0.019^{+0.003}_{-0.003}$	$0.015^{+0.002}_{-0.002}$
Sum	$0.507\substack{+0.014 \\ -0.013}$	$0.531^{+0.015}_{-0.013}$	$0.506\substack{+0.042\\-0.037}$	$0.572^{+0.029}_{-0.028}$

momentum carried by individual/all light charged hadrons at Q=5 GeV

+ FFs have the interpretation of number densities of hadrons and satisfy various fundamental sum rules as derived from first principle, including momentum sum rule, charge sum rule, etc.; momentum sum rules



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Opportunities with CEPC

hadron production, sufficiently for precision determination of fragmentation functions alone

Particl e	E _{c.m.} (GeV)	Year s	SR Power (MW)	Lumi. /IP (10 ³⁴ cm ⁻² s ⁻¹)	Integrated Lumi. /yr (ab ⁻¹ , 2 IPs)	Total Integrated L (ab ⁻¹ , 2 IPs)	Total no. of events
Н*	240	10	50	8.3	2.2	21.6	4.3 × 10 ⁶
			30	5	1.3	13	2.6 × 10 ⁶
Z		0	50	192**	50	100	4.1 × 10 ¹²
	91	2	30	115**	30	60	2.5 × 10 ¹²
W	100	_	50	26.7	6.9	6.9	2.1 × 10 ⁸
	160	1	30	16	4.2	4.2	1.3 × 10 ⁸
tī	360	5	50	0.8	0.2	1.0	0.6 × 10 ⁶
	000	Ŭ	30	0.5	0.13	0.65	0.4 × 10 ⁶

CEPC operation plans

+ Future run of Circular Electron Positron Collider (CEPC) will collect high-quality and diverse data on

- * producing qqbar samples at various energies with high statistics, important for u/d separation
- * separated into bins of different polar angles for additional flavor and charge separation
- heavy-quark enriched samples and gluon samples from Higgs hadronic decays
- In the second W-boson production with hadronic decays

[a projection study on FFs from CEPC data alone is in preparation]



Nuclear physics: hadron supp

 Apart from study of QCD, hadron production are thus important for studying properties of Q_{a} match, espectant, jet damperation encode

medium (QGP) induced radiations



$$\tilde{D}_{h/d}(z_d, \mu^2, \Delta E_d) = (1 - e^{-\langle N_g^d \rangle}) \left[\frac{z_d'}{z_d} D_{h/d}(z_d', \mu^2) + \langle N_g^d \rangle \frac{z_g'}{z_d} D_{h/g}(z_g', \mu^2) \right] + e^{-\langle N_g^d \rangle} D_{h/d}(z_d, \mu^2) (13)$$

modifications to FFs as functions of jet

[2208.14419]



incl. hadron production: Pb-Pb vs pp collisions

[1910.07678]



Summary

- production cross sections in high energy scattering from first principle of QCD
- and capability for arbitrary hard processes
- FFs are much improved and large discrepancies are found wrt. previous determinations
- hadrons and medium modified FFs, projections for future e+e- colliders, etc.

Fragmentation functions (FFs) are essential non-perturbative inputs for precision calculations of hadron

FMNLO a new program for automated and fast calculations of fragmentation processes at NLO in QCD is now publicly available, which is desirable for global analysis of FFs providing much improved efficiency

• We perform a new global analysis of FFs to identified charged hadrons, including charged pion, kaon and proton, at NLO in QCD, using most recent data from SIA, SIDIS, and pp collisions; constraints on gluon

◆ Ongoing developments include extensions to higher orders in QCD (NNLO), studies of FFs to neutral

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Thank you for your attention!

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