



中国科学院近代物理研究所
Institute of Modern Physics, Chinese Academy of Sciences



Observation of the Antimatter Hypernucleus ${}_{\Lambda}^4\overline{H}$

Qiu, Hao (仇浩)

Institute of Modern Physics, CAS

for the STAR Collaboration



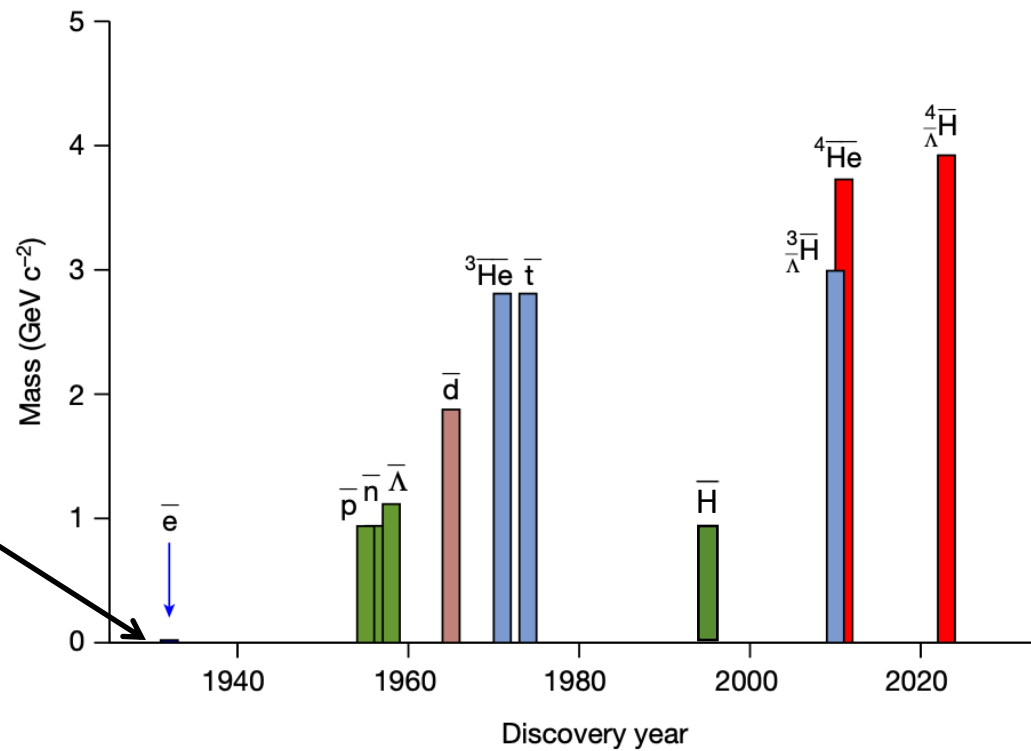
U.S. DEPARTMENT OF
ENERGY

Office of
Science

Introduction: History of Anti-matter Discovery



Dirac, P.A.M.,
The Quantum Theory of the Electron.
Proc. Roy. Soc. Lond. A 117, 610 (1928).



Introduction: History of Anti-matter Discovery



Carl D. Anderson

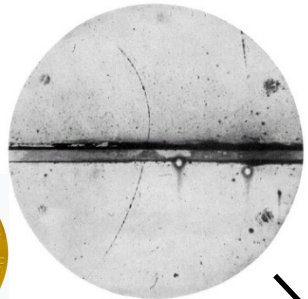
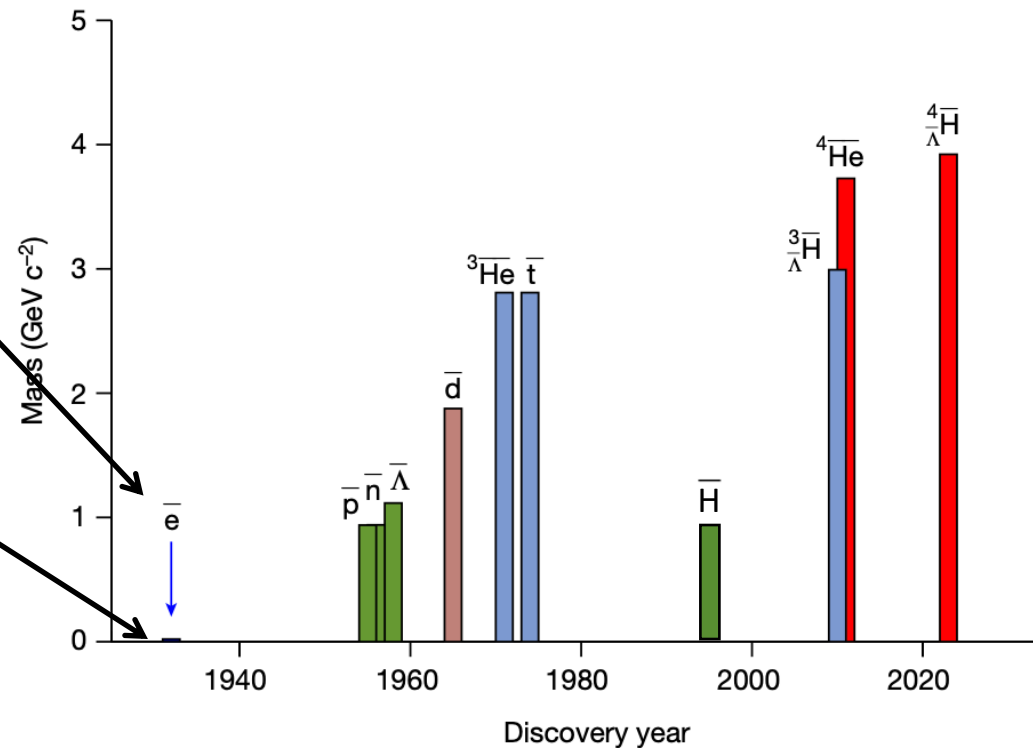


Fig. 1. A 65 million volt positron (β^+) (β^+ reversed) passing through a 0.5 cm lead plate and emerging as a 22 million volt positron (β^+) (β^+ reversed). The length of this latter path is four ten times greater than the possible length of a proton path of this curvature.

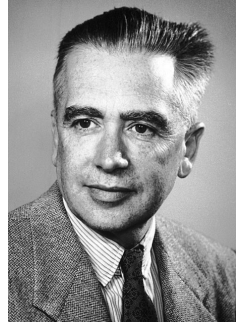


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Introduction: History of Anti-matter Discovery



Emilio Segrè



Owen Chamberlain

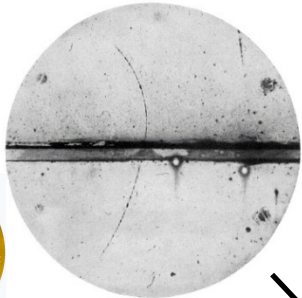
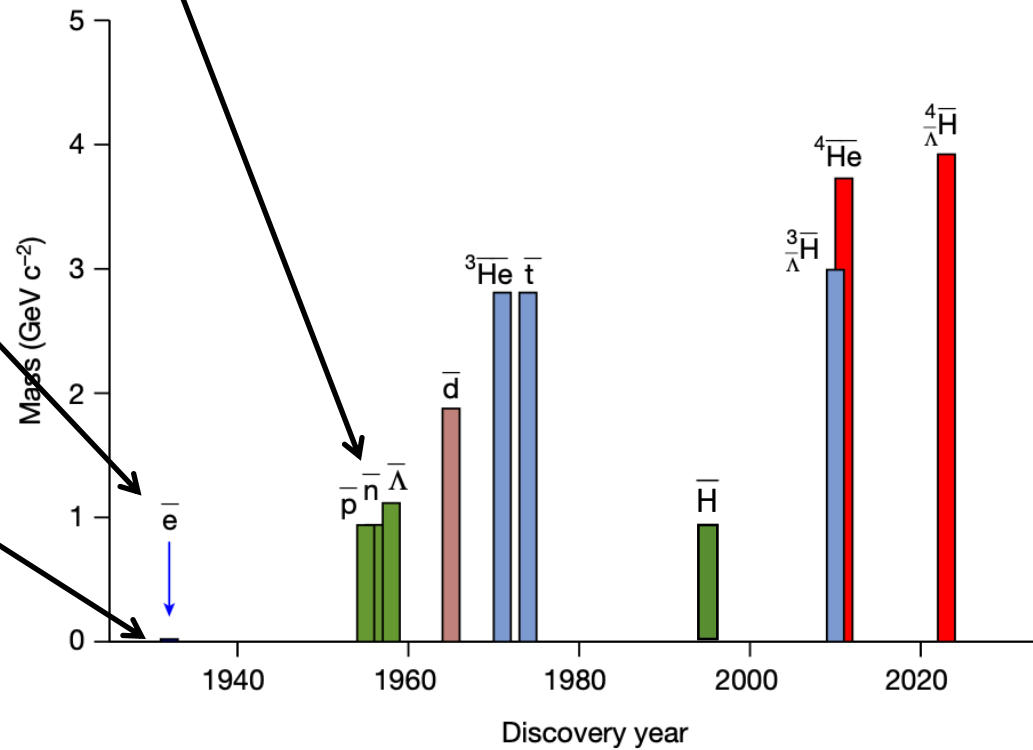


Fig. 1. A 65 million volt positron (β^+) (β^+ general) passing through a 0.1 cm lead plate and emerging as a 22 million volt positron (β^+) (β^+ general). The length of this latter path is about ten times greater than the possible length of a positron path of this curvature.



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Introduction: History of Anti-matter Discovery

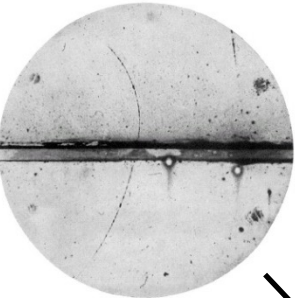
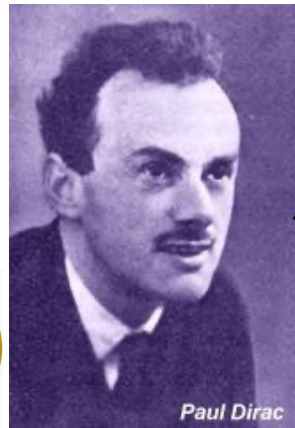
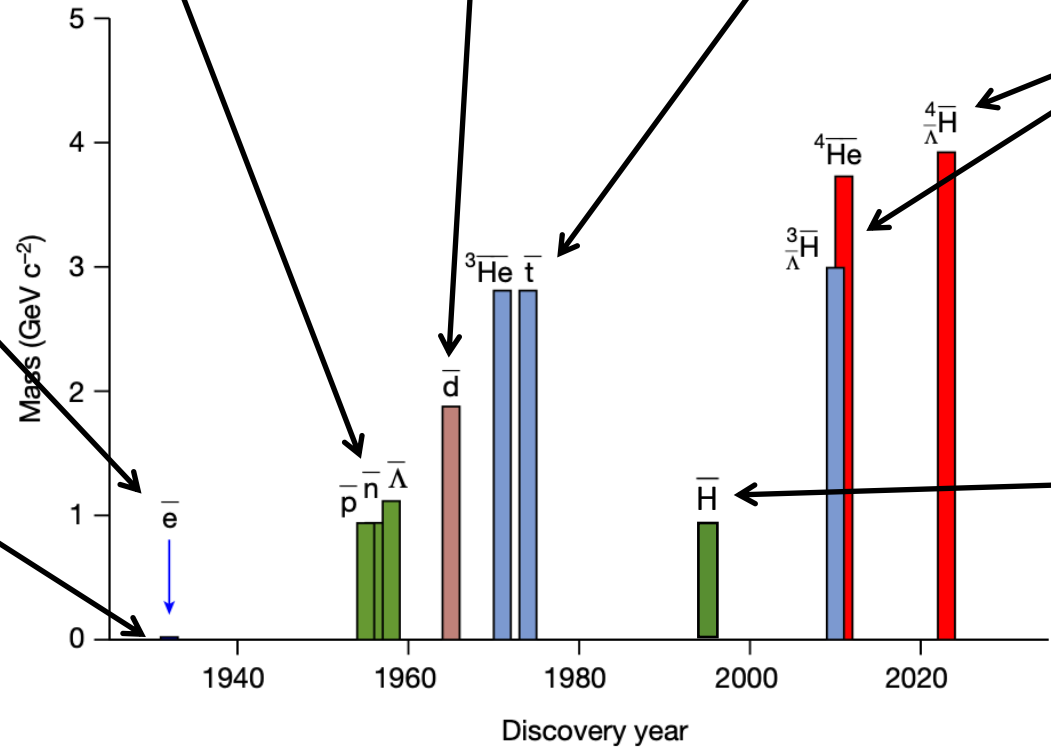


Fig. 1. A 65 million volt positron (β^+) (β^+ general) passing through a 0.1 cm lead plate and emerging as a 22 million volt positron (β^+) (β^+ general). The length of this layer plate is four ten times greater than the possible length of a positron path of this curvature.

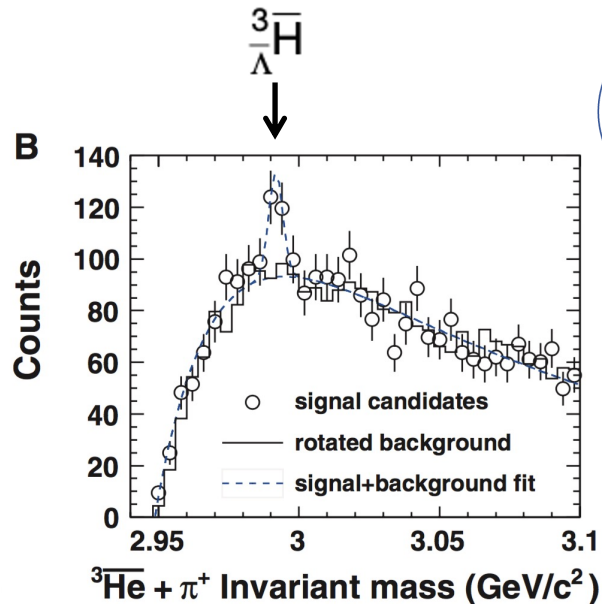
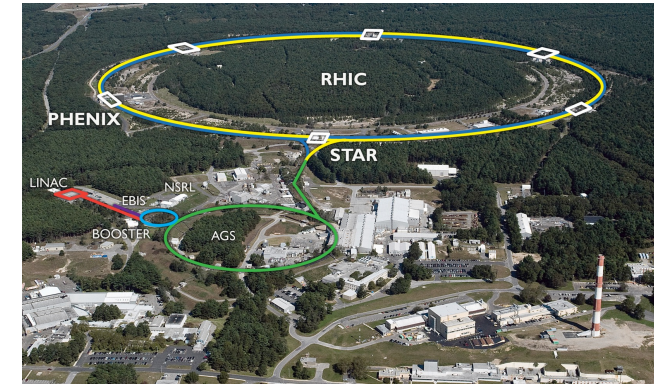
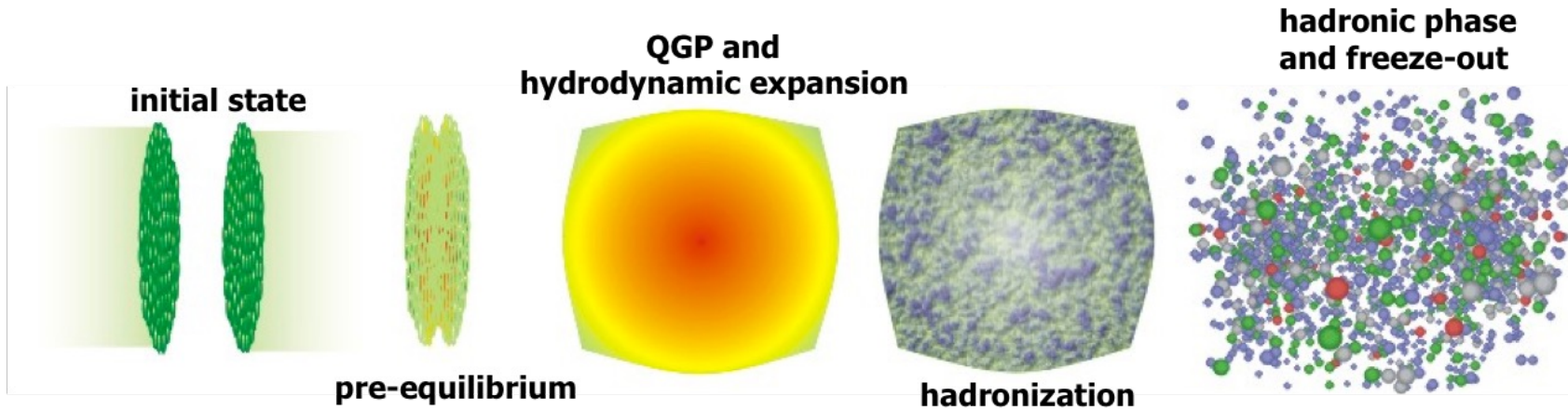


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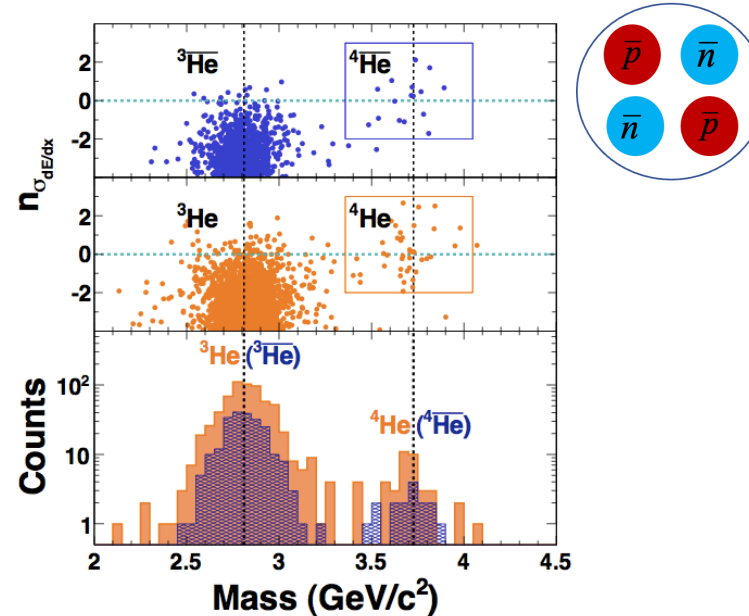




Introduction: History of Anti-matter Discovery

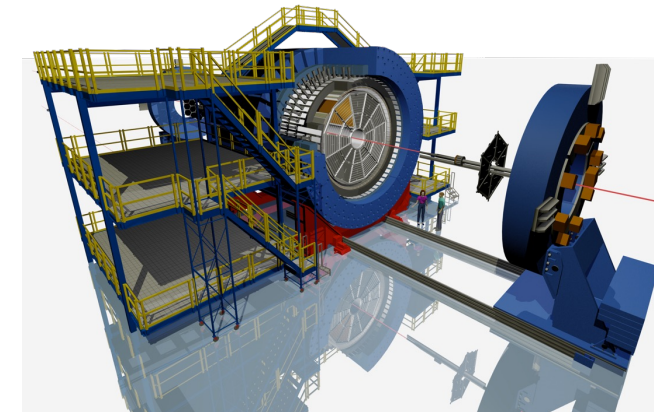


Science 328, 58 (2010)



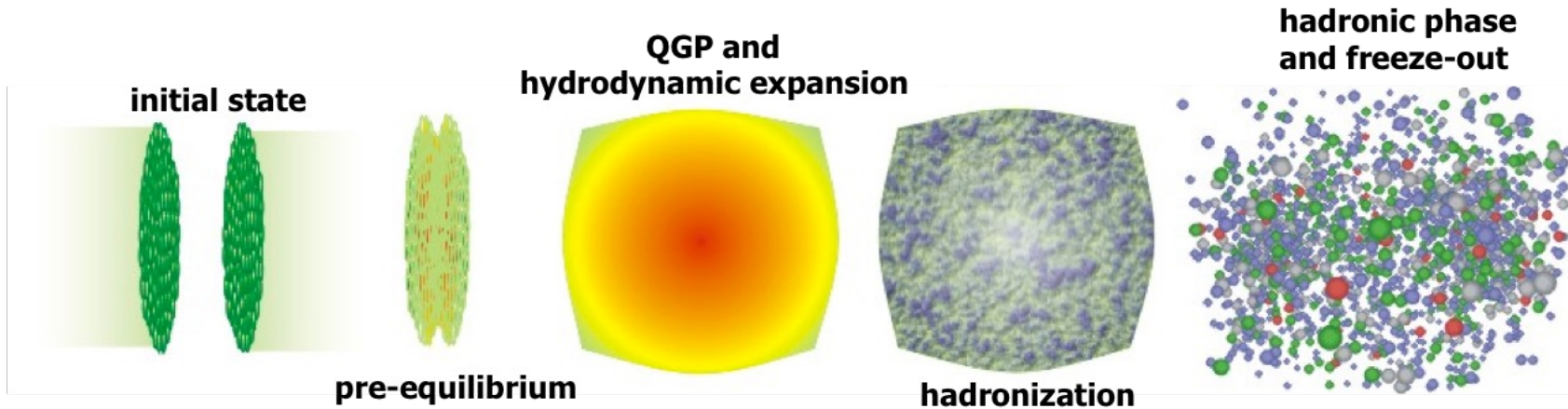
Nature 473, 353 (2011)

RHIC



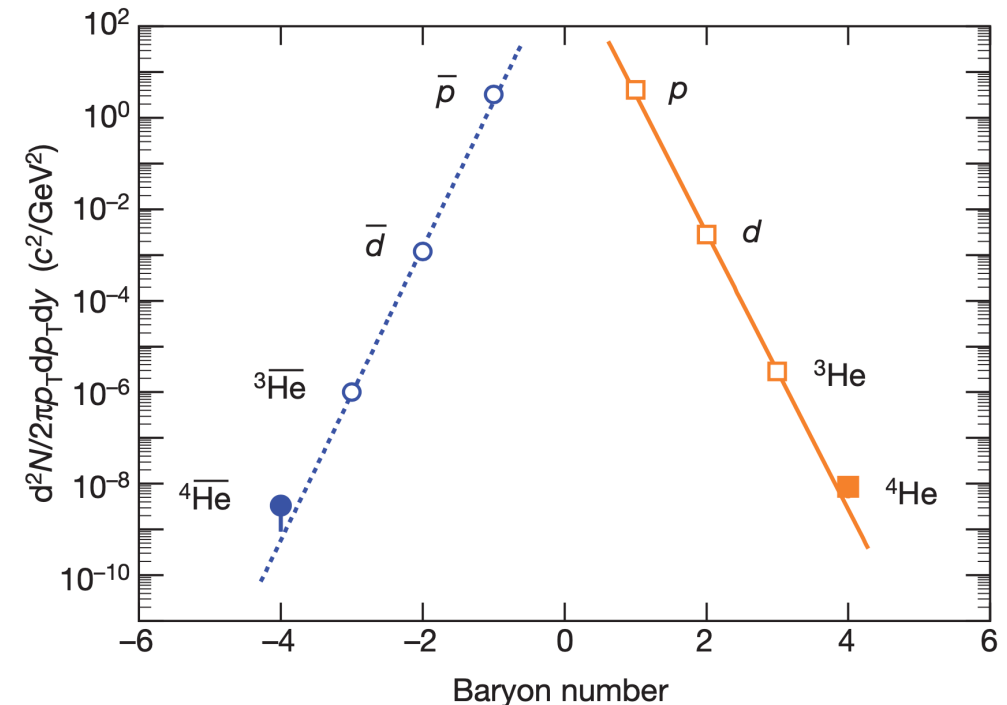
STAR

Introduction: Anti(hyper)nucleus Production

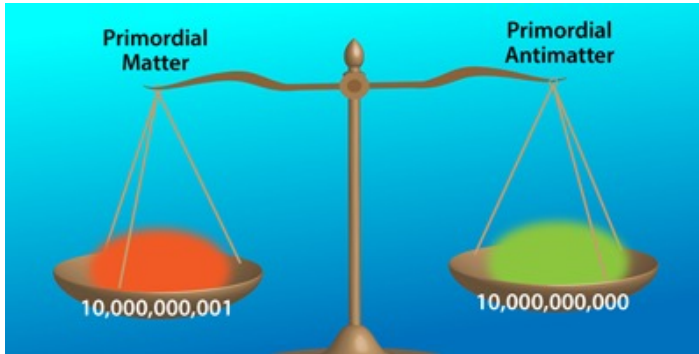


Nature 473, 353 (2011)

- Anti(hyper)nuclei binding energy \sim several MeV / nucleon
- QGP temperature \sim several hundred MeV
- \Rightarrow They are produced by coalescence of antibaryons in the last stage of the collision

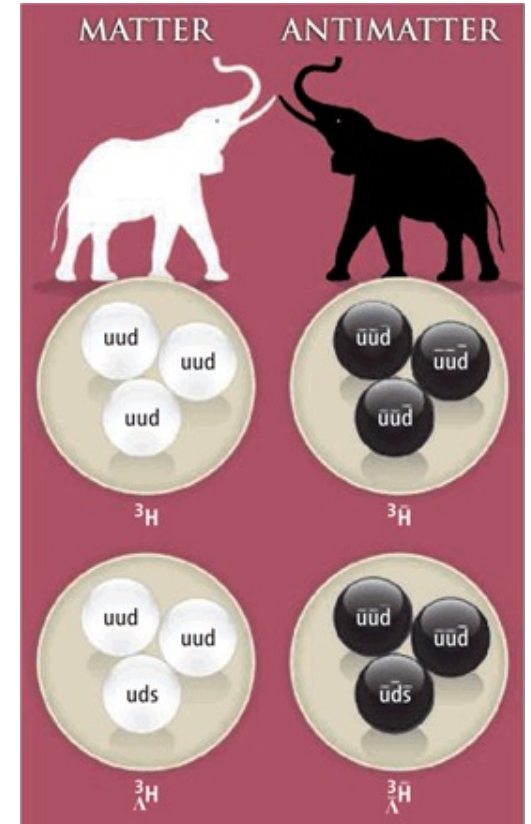


Introduction: Matter-antimatter Asymmetry in the Universe

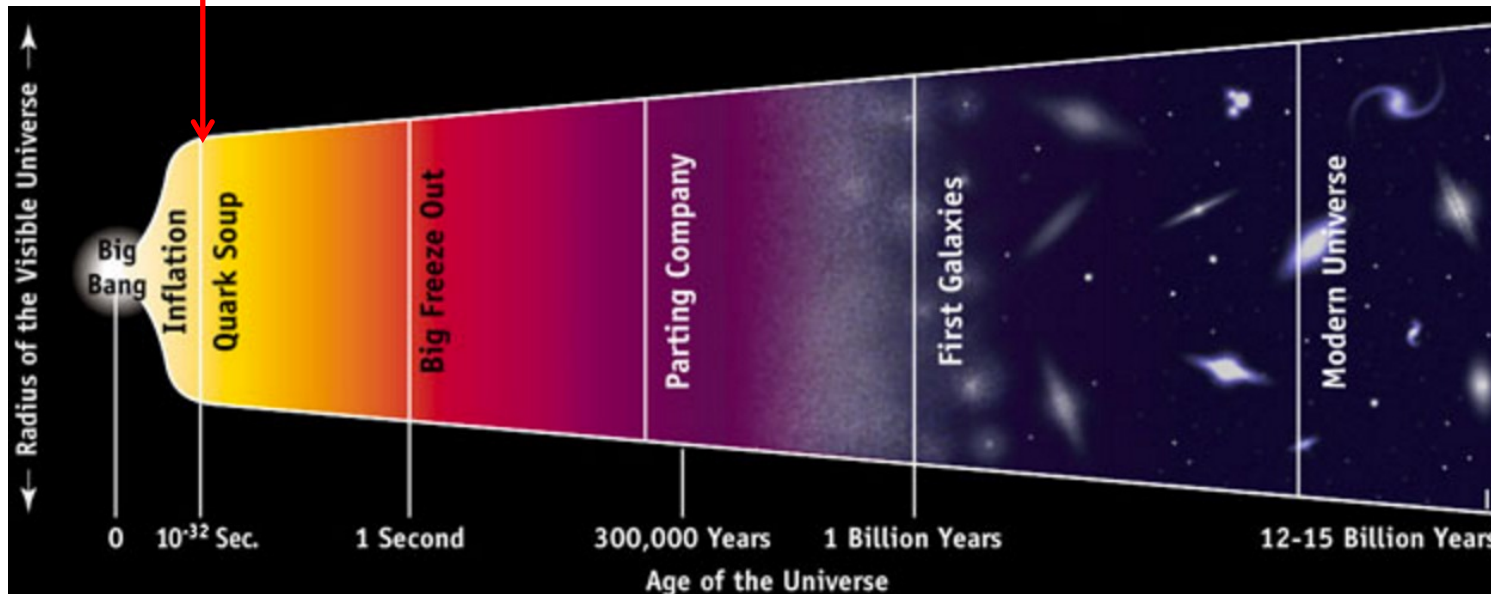


?

- Matter-antimatter asymmetry in early universe is the precondition for the existence of the matter world today
- The source of this asymmetry is still not clear

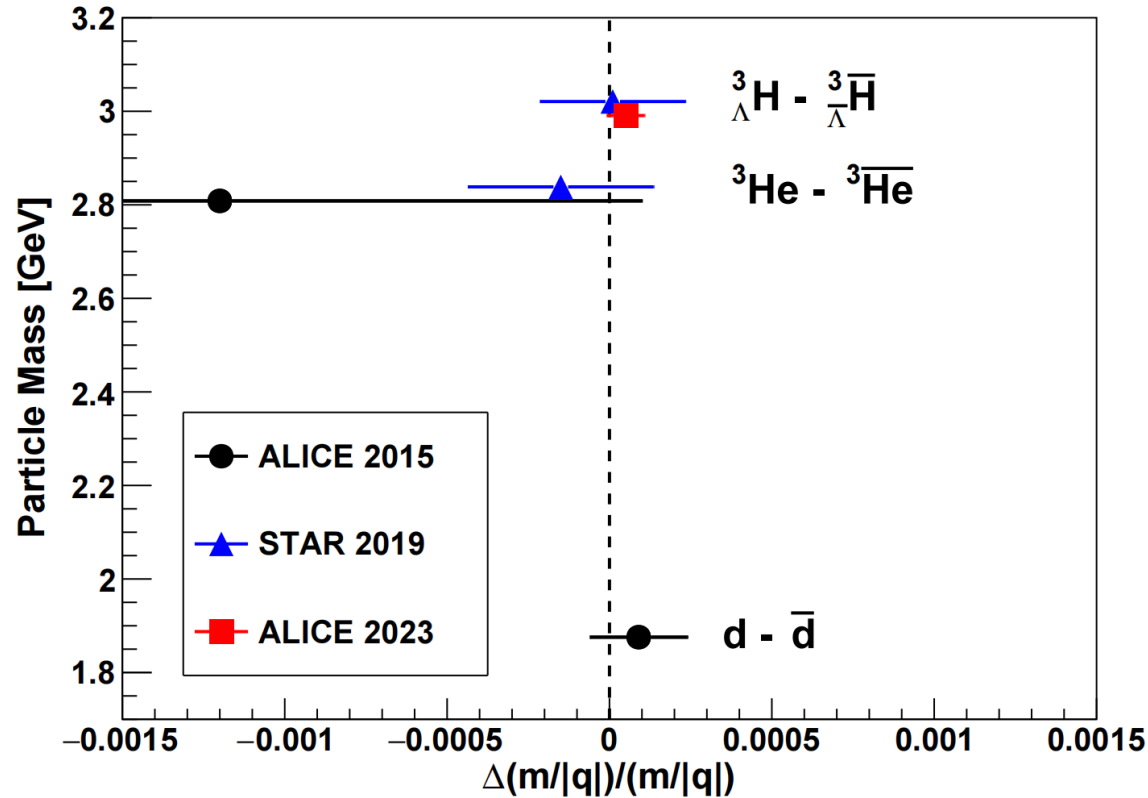


CPT theorem =>
Symmetry of matter-
antimatter properties





Introduction: CPT symmetry test with heavy-ion collisions



$$\frac{\tau_{{}^3_{\Lambda}\text{H}} - \tau_{{}^3_{\Lambda}\bar{\text{H}}}}{\tau_{{}^3_{\Lambda}\text{H}}} = [3 \pm 7(\text{stat}) \pm 4(\text{syst})] \times 10^{-2}$$

Nature Phys. 11 (2015) 10, 811-814

Nature Phys. VOL 16, April 2020, 409–412

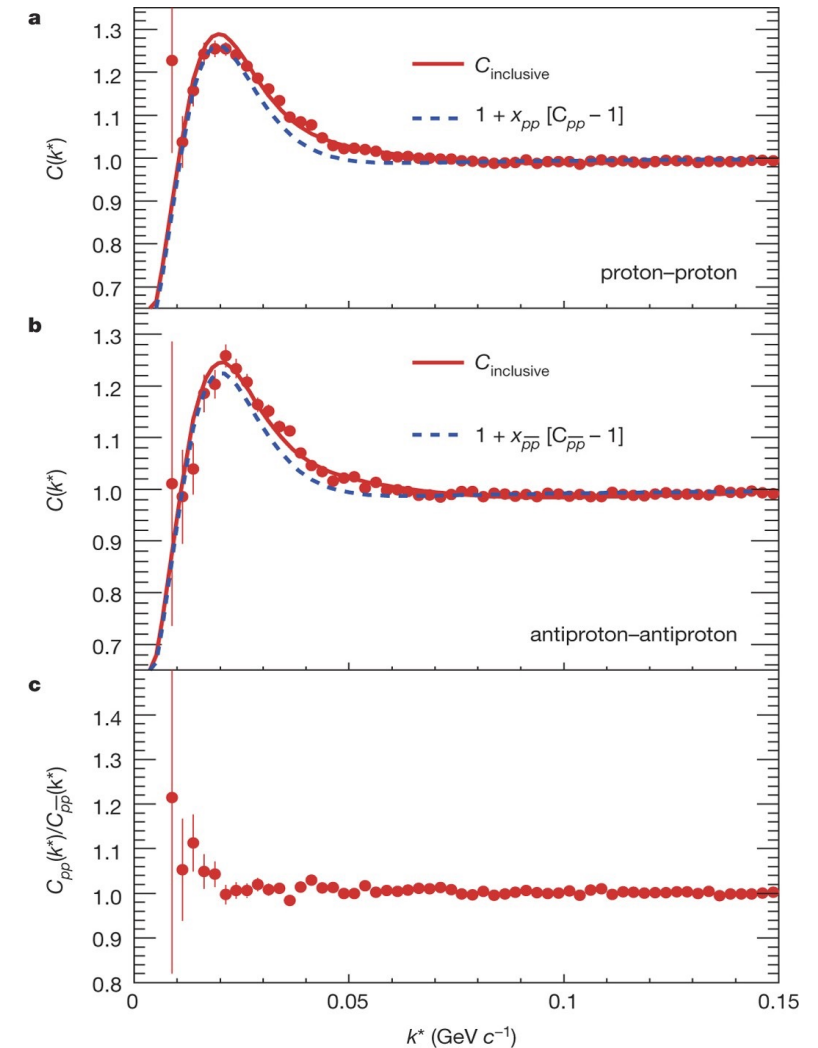
PHYSICAL REVIEW LETTERS 131, 102302 (2023)

- No significant mass or binding energy difference between d & \bar{d} , ${}^3\text{He}$ & ${}^3\bar{\text{He}}$, ${}^3_{\Lambda}\text{H}$ and ${}^3_{\Lambda}\bar{\text{H}}$
- No significant lifetime difference between ${}^3_{\Lambda}\text{H}$ and ${}^3_{\Lambda}\bar{\text{H}}$



Introduction: CPT symmetry test with heavy-ion collisions

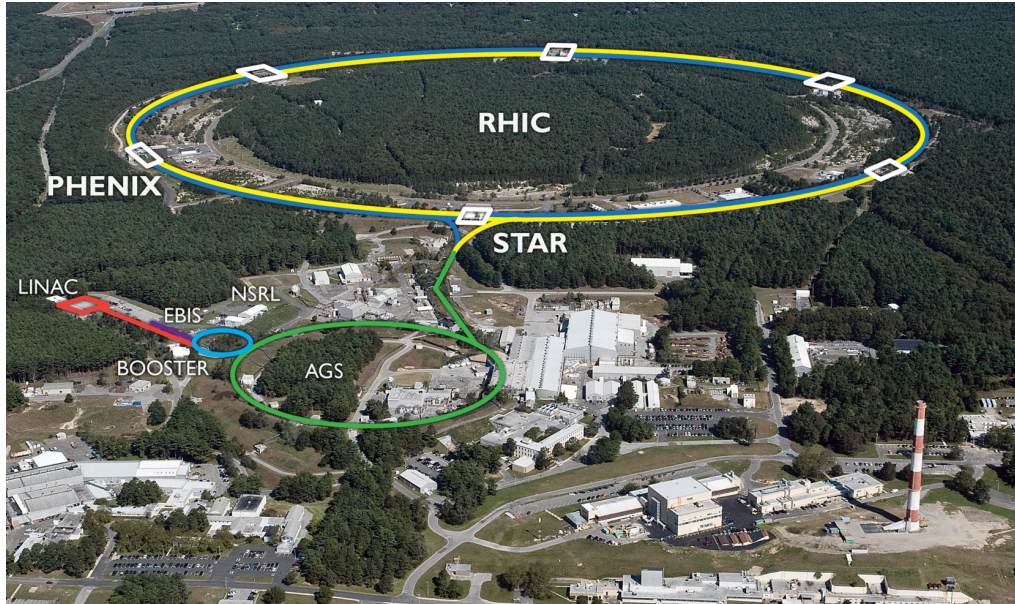
- No difference between p-p and $\bar{p}\text{-}\bar{p}$ correlation functions
 - \Rightarrow No difference between p-p and $\bar{p}\text{-}\bar{p}$ interactions



Nature volume 527, pages345–348 (2015)

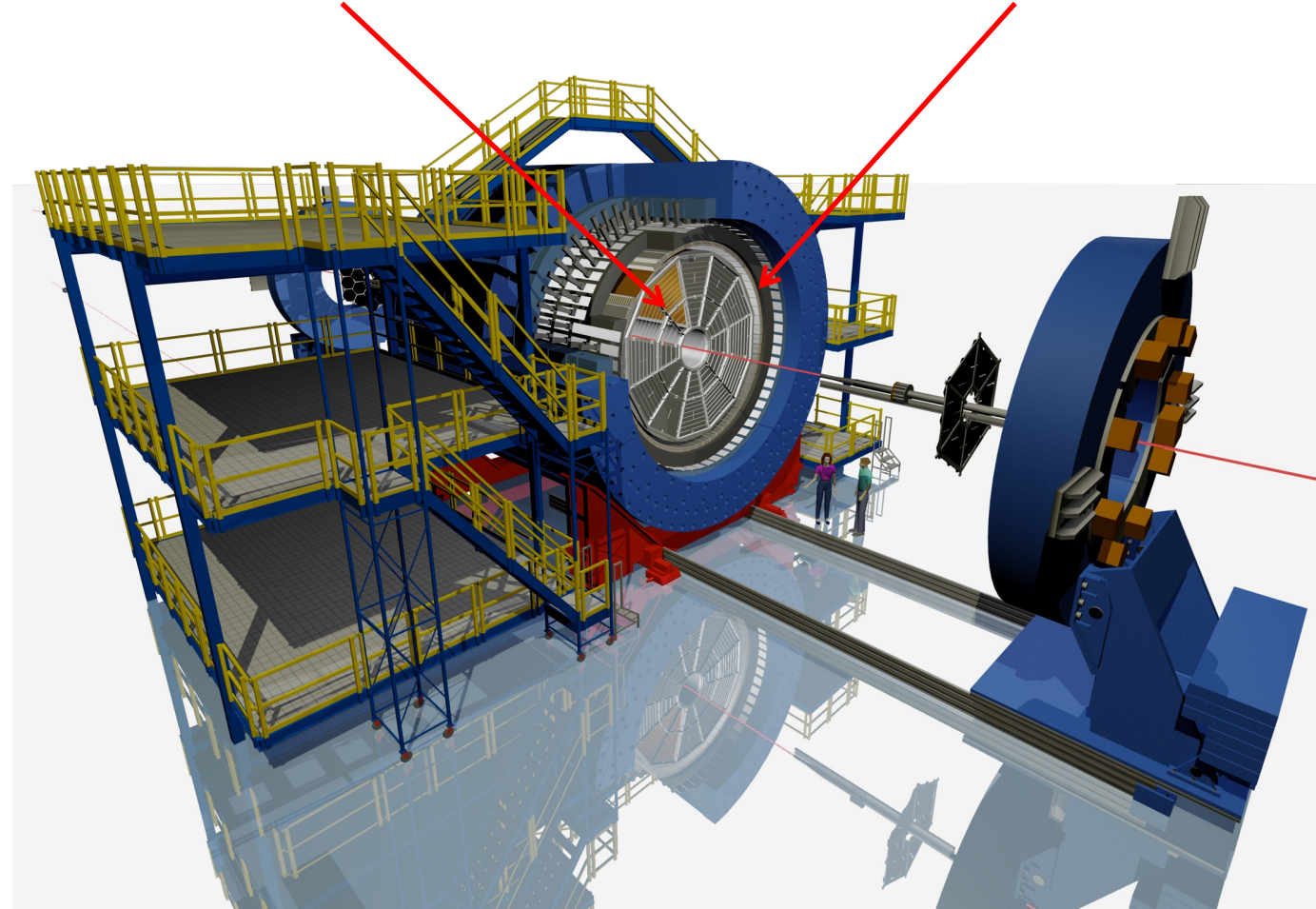


RHIC-STAR



Time Projection Chamber
• tracking $\Rightarrow p/Z, dE/dx$

Time-Of-Flight detector
• TOF + $p/Z + L \Rightarrow m^2/Z^2$





Data Sets

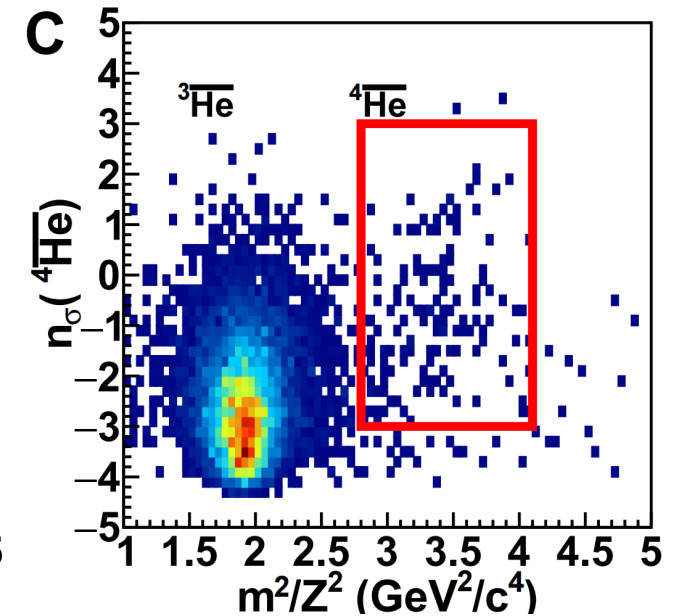
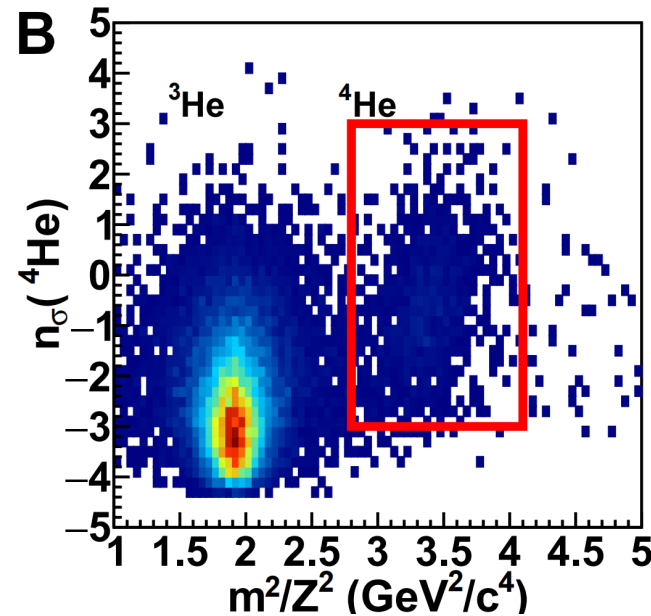
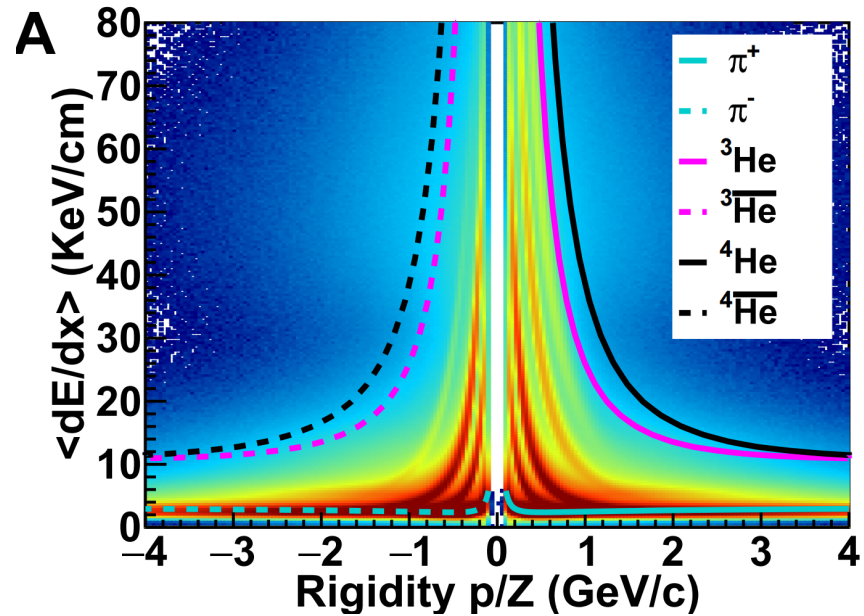
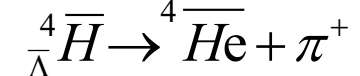
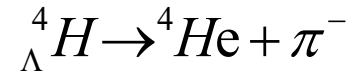
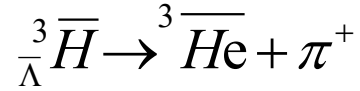
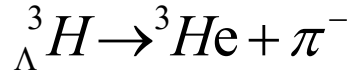
data set	year	N events
AuAu@200 GeV	2010	~606 M
AuAu@200 GeV	2011	~626 M
UU@193GeV	2012	~512 M
ZrZr+RuRu(Isobar)@200GeV	2018	~4.7 B

Trigger:

- Minimum bias trigger
- Central trigger
- Electromagnetic and hadronic triggers
-

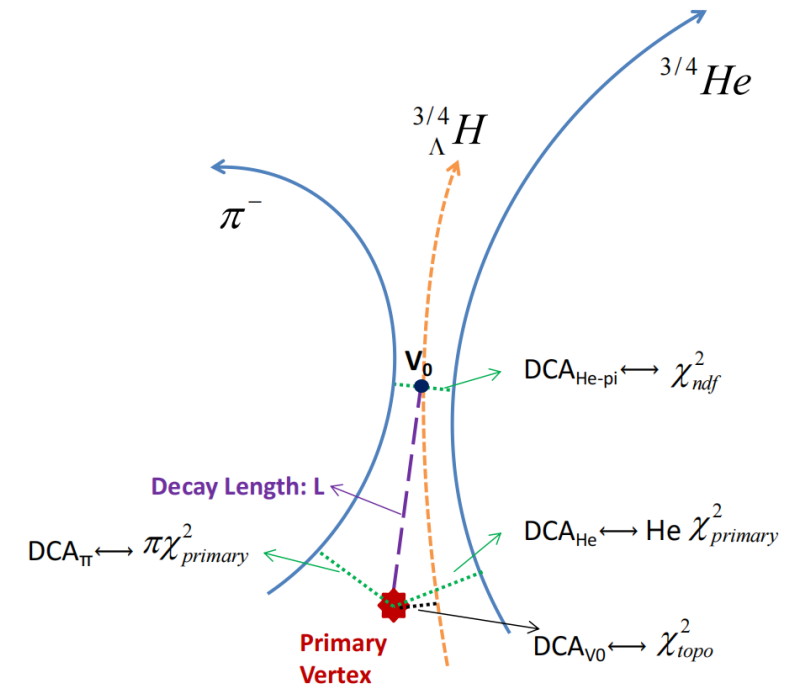
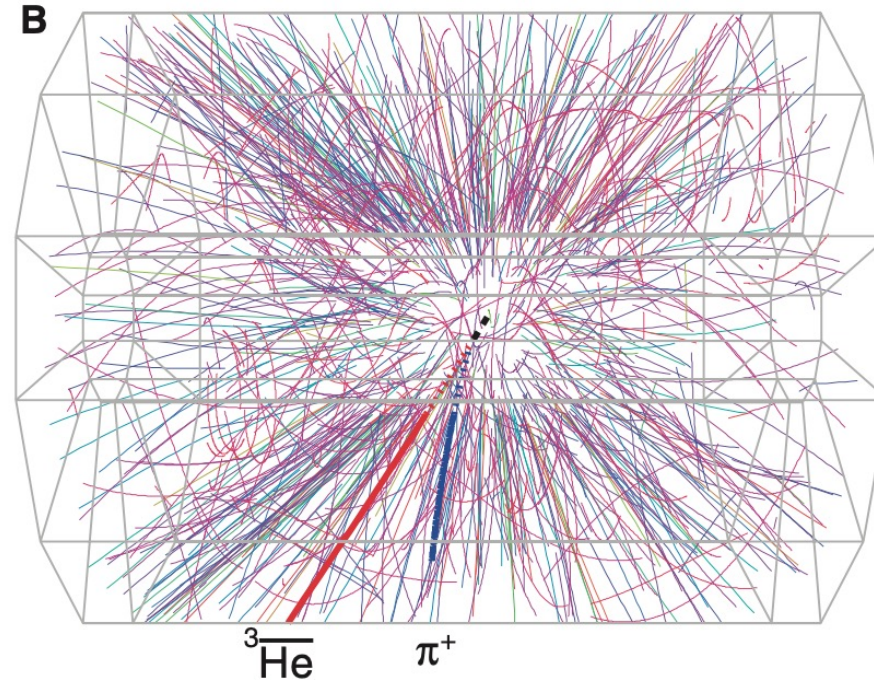
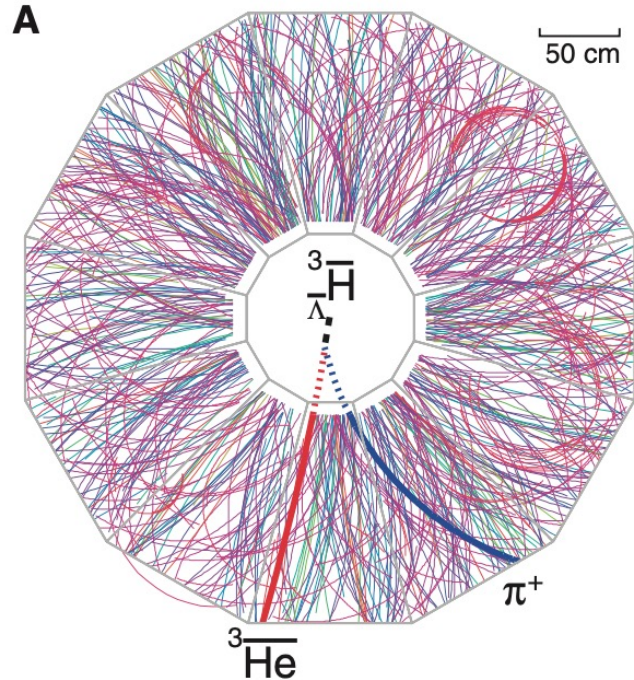
- Use as many triggers as possible to find signal and measure lifetime
- Use minimum bias trigger for production yield ratios measurement

Decay Channels & Daughter Particle Identification



- ${}^3\text{He}$ PID: ($Z < 0 \parallel p > 2.$) && $|n_{\sigma}({}^3\text{He})| < 3$ && (if TOF matched, $1 < M^2/Z^2 < 3$);
- ${}^4\text{He}$ PID: ($Z < 0 \parallel p > 2.$) && $|n_{\sigma}({}^4\text{He})| < 3$ && ($|n_{\sigma}({}^3\text{He})| > 3.5 \parallel 2.8 < M^2/Z^2 < 4.1$);
- π PID: $|n_{\sigma_{\pi}}| < 3$;

Decay Vertex Reconstruction



Science 328, 58 (2010)

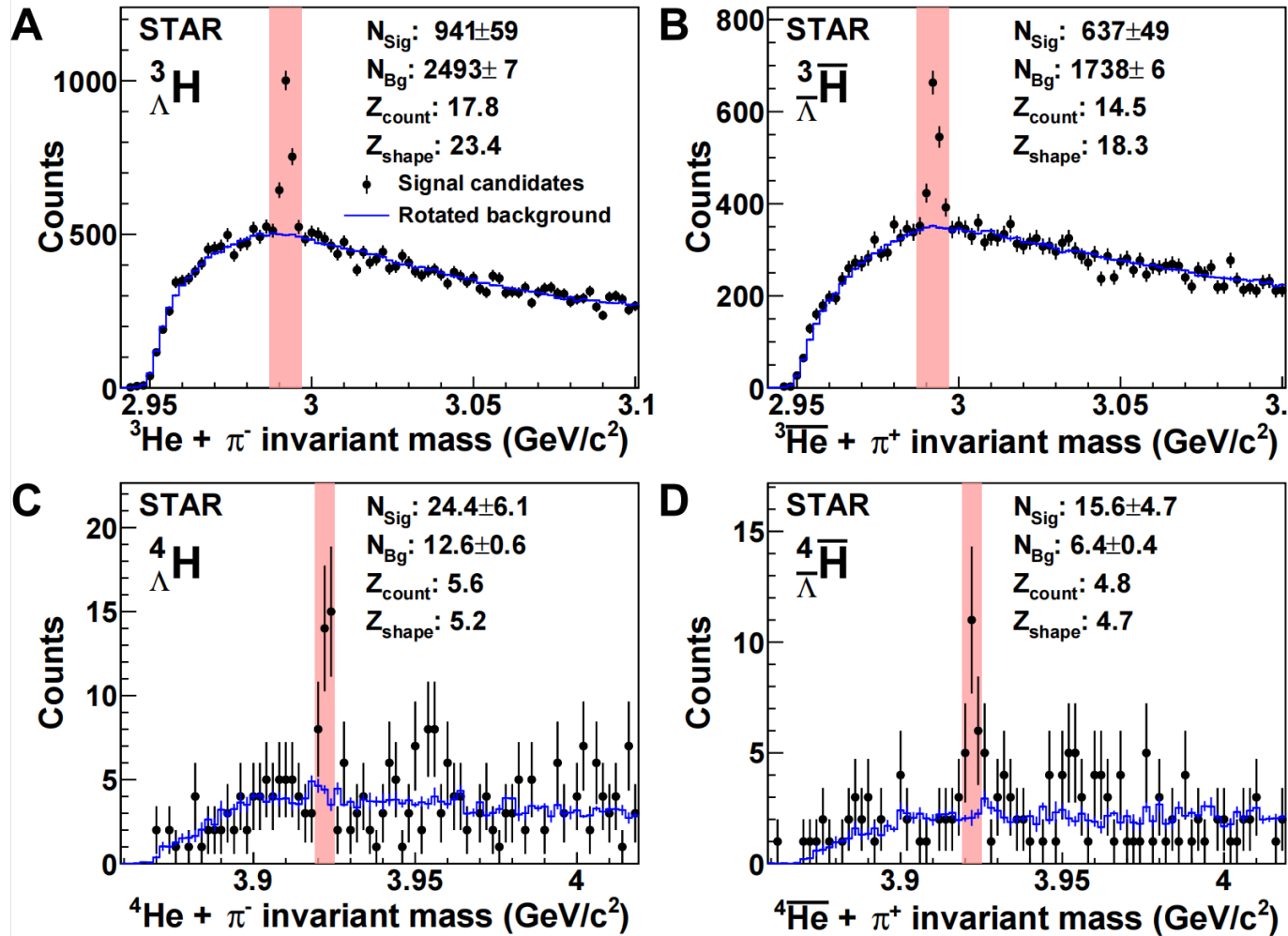
S. Gorbunov and I. Kisel, CBM-SOFT-note-2007-003, 7 May 2007

M. Zyzak, "Online selection of short-lived particles on many-core computer architectures in the CBM experiment at FAIR," Dissertation thesis, Goethe University of Frankfurt, 2016

- KF(Kalman Filter) Particle package for decay vertex reconstruction
- Topology cuts obtained by optimizing $\frac{3}{\Lambda} \bar{H}$ significance
 - blind for ${}^4_{\Lambda} H$ and ${}^4_{\Lambda} \bar{H}$

Particle	$\chi^2_{\text{prim He}}$	$\chi^2_{\text{prim } \pi}$	χ^2_{ndf}	χ^2_{topo}	L/dL	L	He DCA
${}^3_{\Lambda} H$ & ${}^4_{\Lambda} H$	<2000	>10	<5	<2	>3.5	>3.4cm	<1cm
${}^3_{\Lambda} \bar{H}$ & ${}^4_{\Lambda} \bar{H}$	<2000	>10	<5	<3	>3.5	>3.4cm	-

Signals



Nature 2024, DOI: 10.1038/s41586-024-07823-0

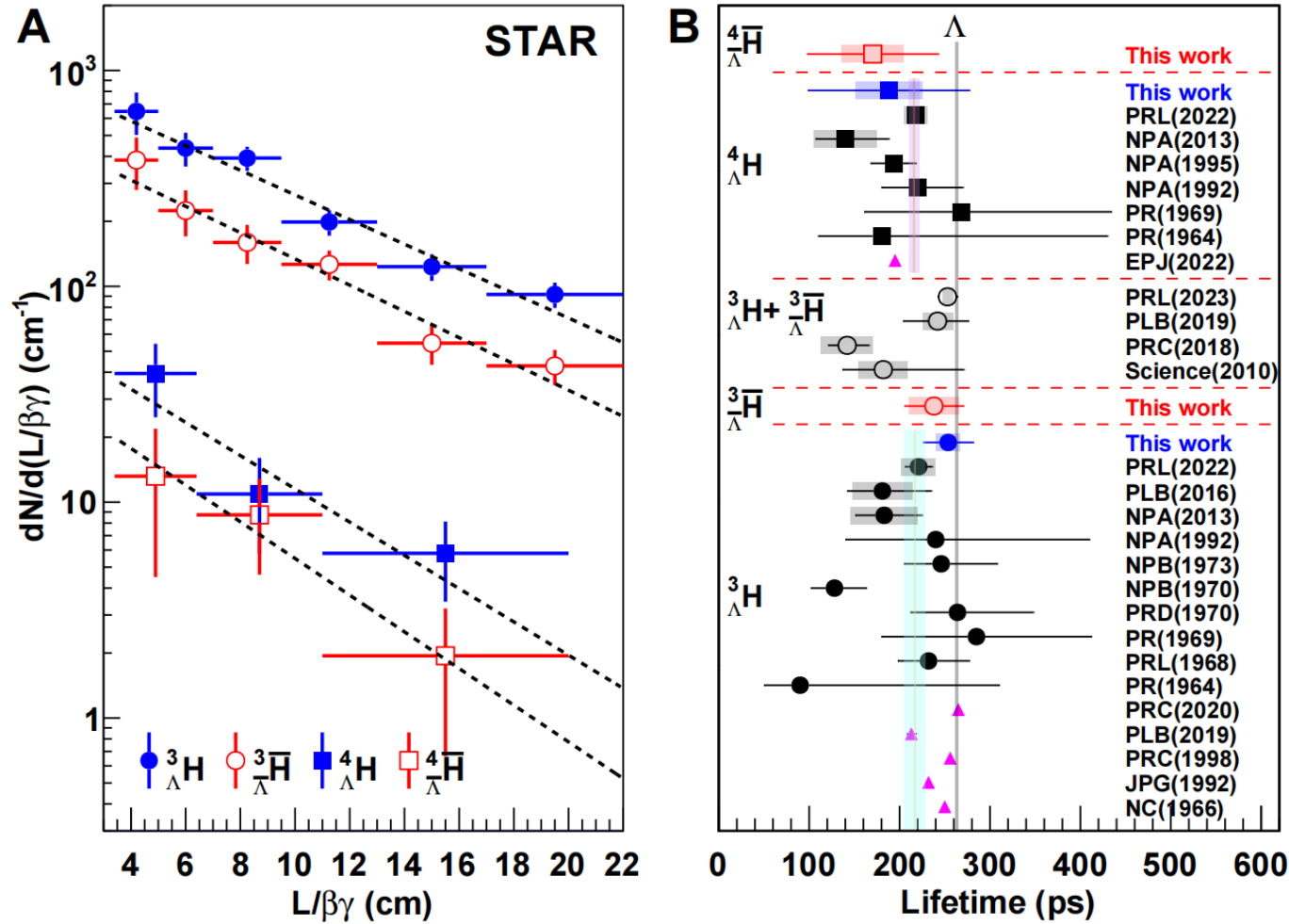
- Background invariant-mass distributions obtained by rotating the (anti)He daughter track before reconstructing the decay vertex

$$Z_{\text{count}} = \sqrt{2 \left[(N_{\text{Sig}} + N_{\text{Bg}}) \ln \left(1 + \frac{N_{\text{Sig}}}{N_{\text{Bg}}} \right) - N_{\text{Sig}} \right]}$$

- Z_{shape} obtained with `RooStats()::AsymptoticCalculator()` assuming pure background vs. background + Gaussian signal
- 15.6 ${}^4\bar{\text{H}}$ signal candidates
- Significances $Z_{\text{count}} = 4.8$, $Z_{\text{shape}} = 4.7$

The heaviest antihypernucleus observed

Lifetime Measurements & CPT Symmetry Test



- Efficiency corrected
- Fit with exponential function: $N(t) = N_0 e^{-t/\beta\gamma\tau}$
- Our results consistent with previous average

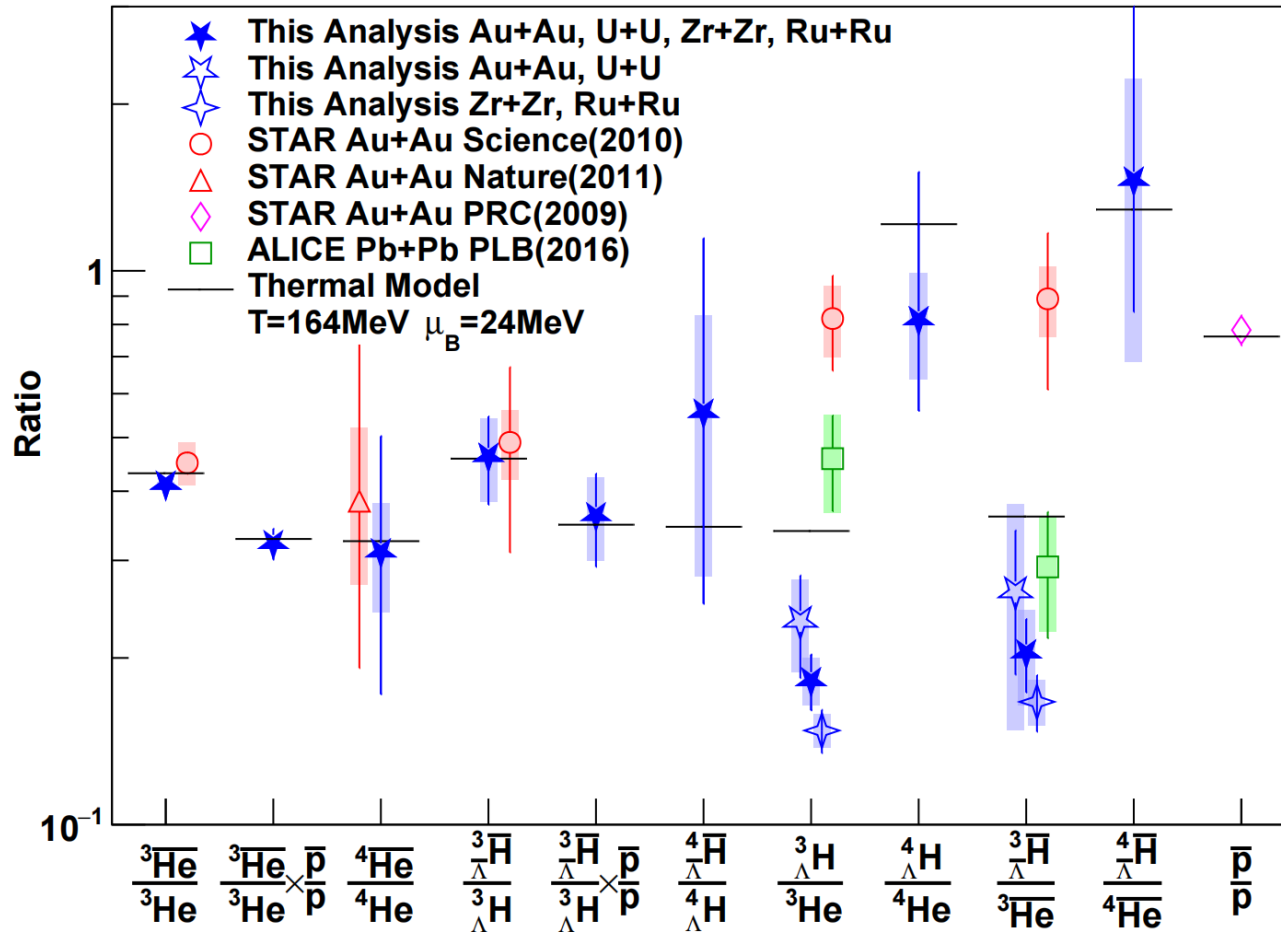
$$\tau_{{}^3_{\Lambda}\text{H}} - \tau_{{}^3_{\Lambda}\bar{\text{H}}} = 16 \pm 43(\text{stat.}) \pm 20(\text{sys.}) \text{ ps}$$

$$\tau_{{}^4_{\Lambda}\text{H}} - \tau_{{}^4_{\Lambda}\bar{\text{H}}} = 18 \pm 115(\text{stat.}) \pm 46(\text{sys.}) \text{ ps}$$

- No lifetime difference between antihypernuclei and their corresponding hypernuclei within uncertainties

Nature 2024, DOI: 10.1038/s41586-024-07823-0

Yield Ratios

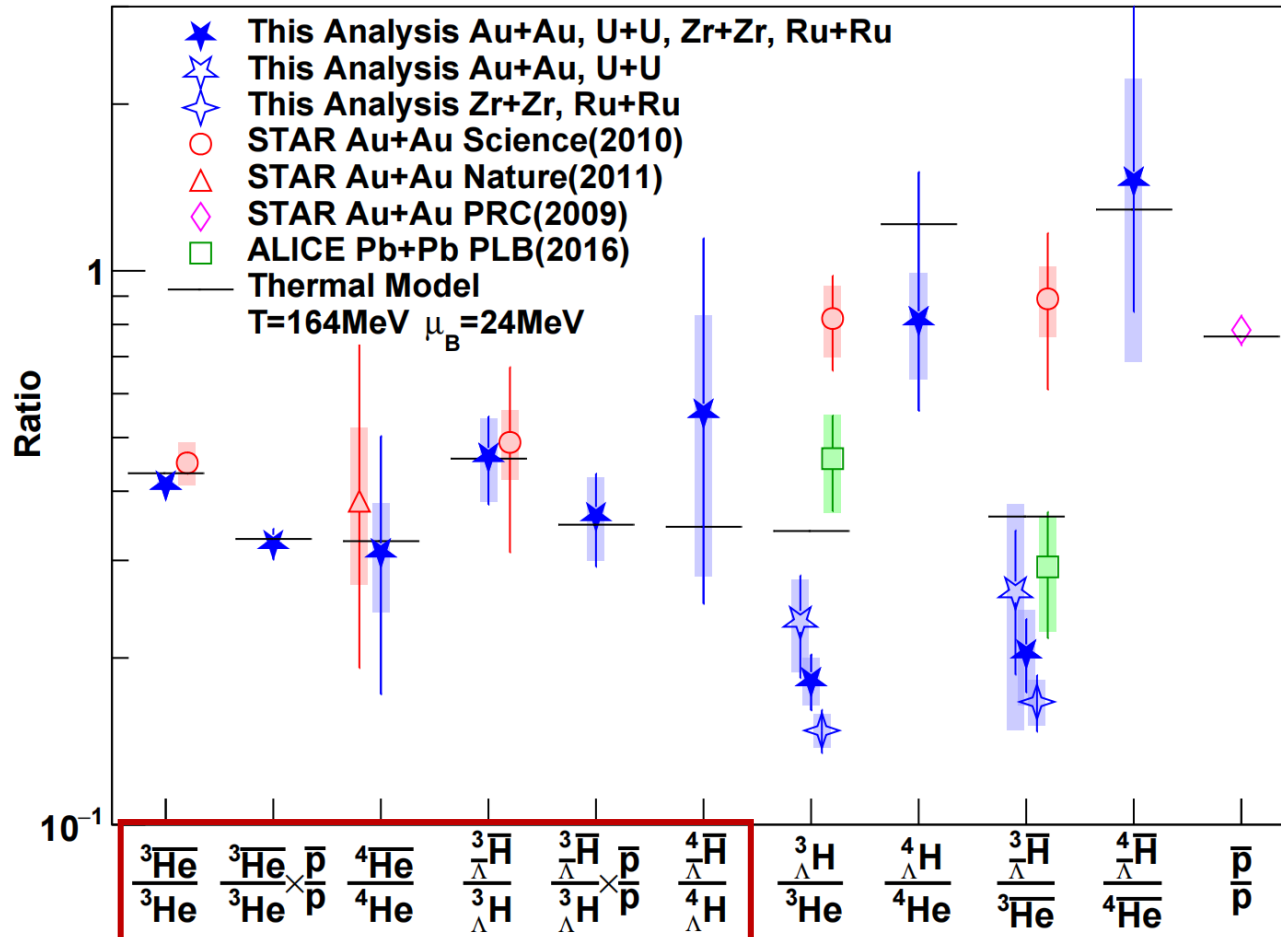


- Branch fraction assumed:
 - 25% for ${}^3_{\Lambda}H \rightarrow {}^3He + \pi^-$
 - 50% for ${}^4_{\Lambda}H \rightarrow {}^4He + \pi^-$
- Phase space: $0.7 < p_T/M < 1.5$, $|y| < 0.7$
- $\frac{{}^3_{\Lambda}H}{{}^3He}$ & $\frac{{}^3_{\Lambda}\bar{H}}{{}^3\bar{He}}$ ratios measured in large and small collision systems separately to have a fair comparison with previous measurements
- Our results are consistent with previous results, except that the $\frac{{}^3_{\Lambda}H}{{}^3He}$ & $\frac{{}^3_{\Lambda}\bar{H}}{{}^3\bar{He}}$ ratios are lower than Science 2010 results by 2.8 & 1.9 σ

Nature 2024, DOI: 10.1038/s41586-024-07823-0
 Science 328, 58 (2010)
 Nature 473, 353–356 (2011)

Phys. Rev. Lett. 97, 152301
 Phys. Lett. B 754 (2016) 360
 Phys. Lett. B 697.3 (2011)

Yield Ratios



$${}^4\overline{\text{He}}/{}^4\text{He} \sim {}^3\overline{\text{He}}/{}^3\text{He} \times \bar{p}/p$$

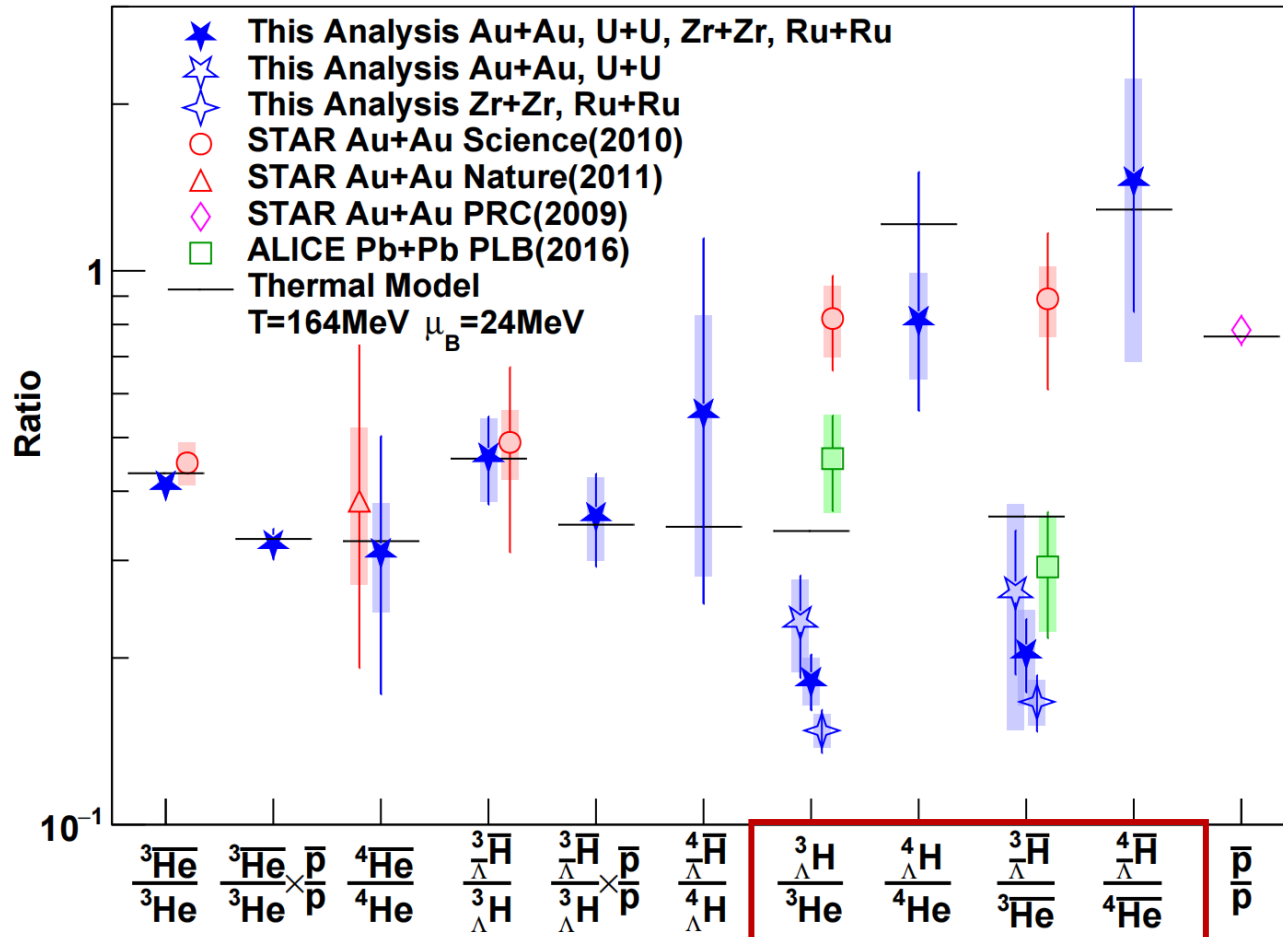
$$\frac{{}^4\overline{\text{H}}}{{}^4\text{H}} \sim \frac{{}^3\overline{\text{H}}}{{}^3\text{H}} \times \bar{p}/p$$

- Consistent with expectation of coalescence picture
- Consistent with thermal model predictions

Nature 2024, DOI: 10.1038/s41586-024-07823-0
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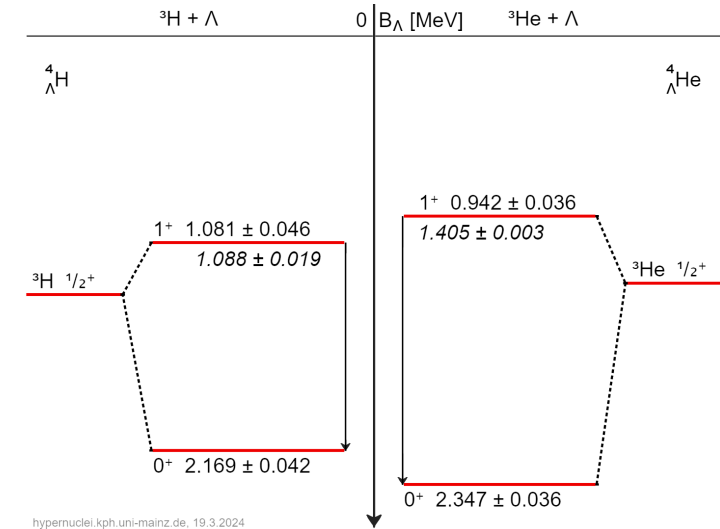
Yield Ratios



$$\frac{^4_\Lambda\text{H}}{^4\text{He}} \sim 4 \times \frac{^3_\Lambda\text{H}}{^3\text{He}}$$

$$\frac{^4_\Lambda\text{Hbar}}{^4\text{Hebar}} \sim 4 \times \frac{^3_\Lambda\text{Hbar}}{^3\text{Hebar}}$$

- Factor 4 due to spin-1 excited states of $^4_\Lambda\text{H}$ & $^4_\Lambda\text{Hbar}$



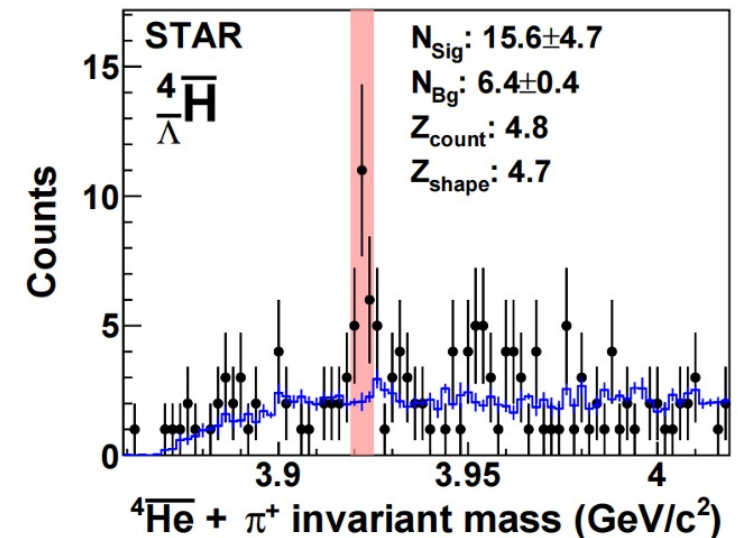
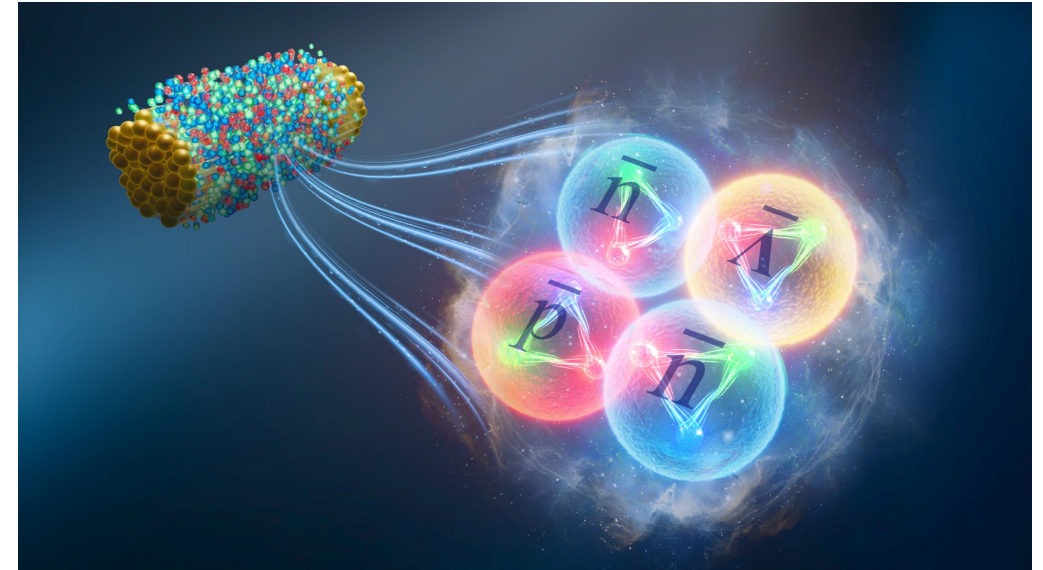
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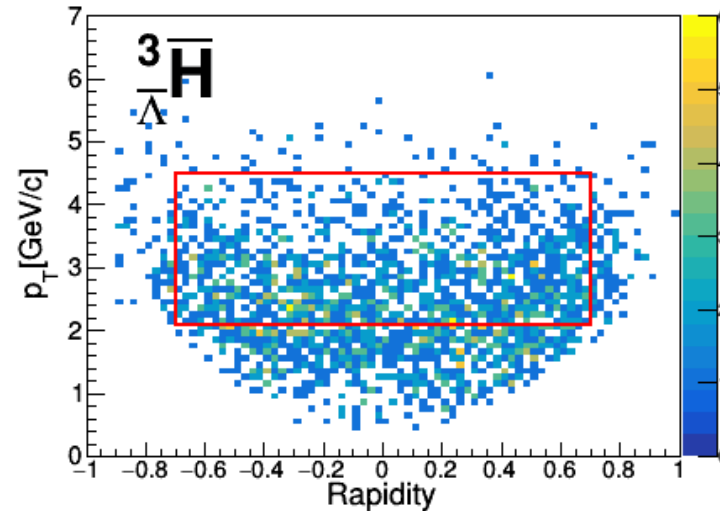
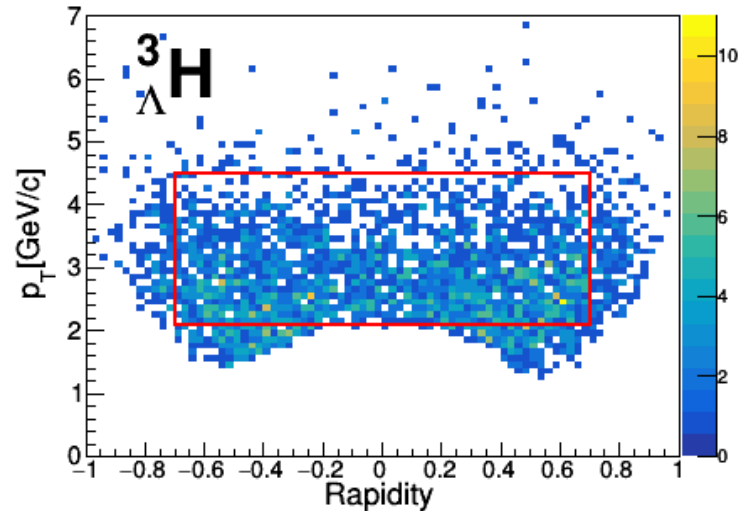
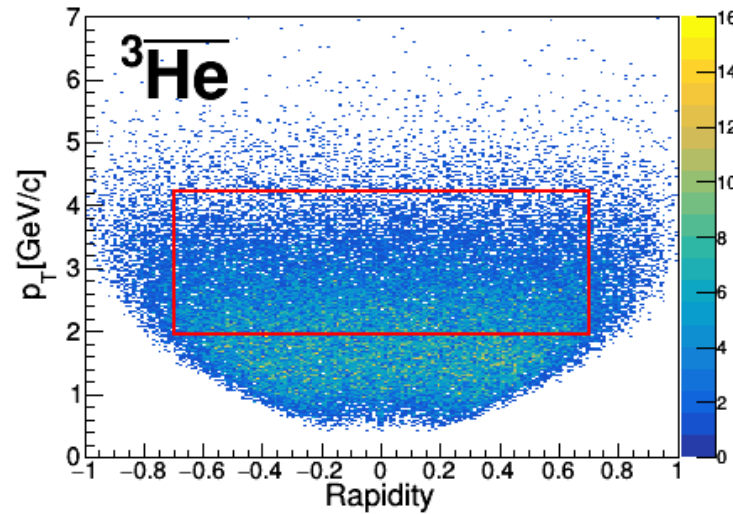
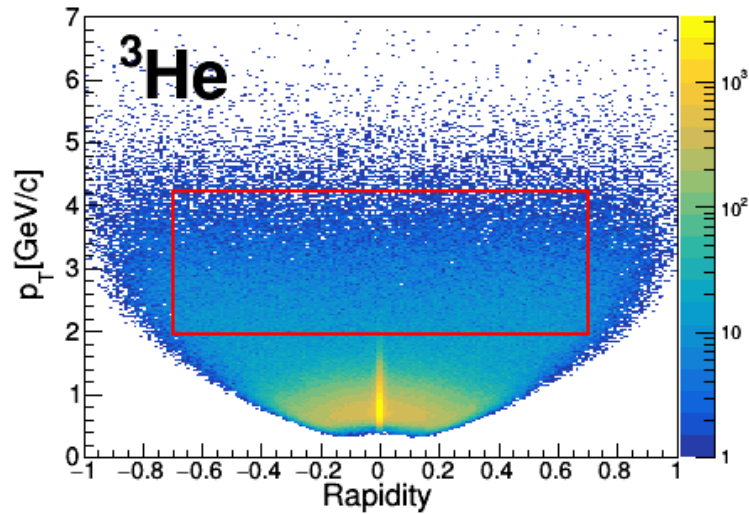
Summary

- 15.6 ${}^4_{\Lambda}\bar{H}$ signal candidates observed, with a significance of 4.7 σ
- Lifetimes of (anti)hypernuclei compared
 - $\tau_{\Lambda}{}^3H \approx \tau_{\Lambda}{}^3\bar{H}$, $\tau_{\Lambda}{}^4H \approx \tau_{\Lambda}{}^4\bar{H}$
 - Confirming CPT symmetry
- Various (anti)particle production yield ratios presented
 - ${}^4\bar{He}/{}^4He \sim {}^3\bar{He}/{}^3He \times \bar{p}/p$
 - ${}^4_{\Lambda}\bar{H}/{}^4_{\Lambda}H \sim {}^3_{\Lambda}\bar{H}/{}^3_{\Lambda}H \times \bar{p}/p$
 - ${}^4H/{}^4He \sim 4 \times {}^3H/{}^3He$
 - ${}^4_{\Lambda}\bar{H}/{}^4\bar{He} \sim 4 \times {}^3_{\Lambda}\bar{H}/{}^3\bar{He}$
 - Consistent with coalescence picture and thermal model



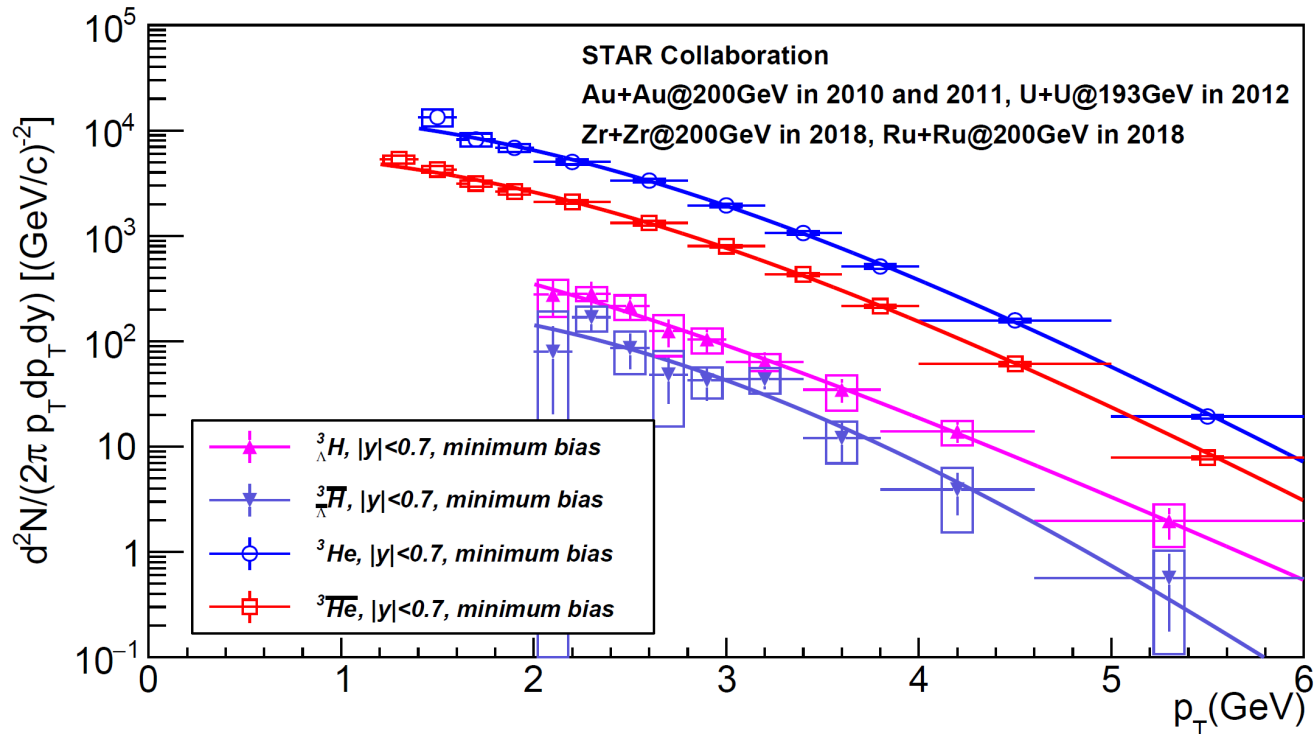
Thanks ☺

Back Up: Yield Ratios Measurement - Phase Space



- Yield measurement in phase space region : $0.7 < p_T/M < 1.5$, $|\text{rapidity}| < 0.7$

Back Up: Yield Ratios Measurement - A = 3 Particles



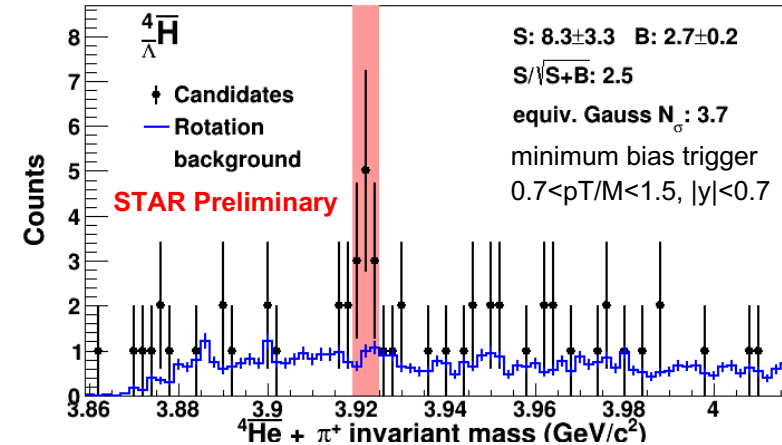
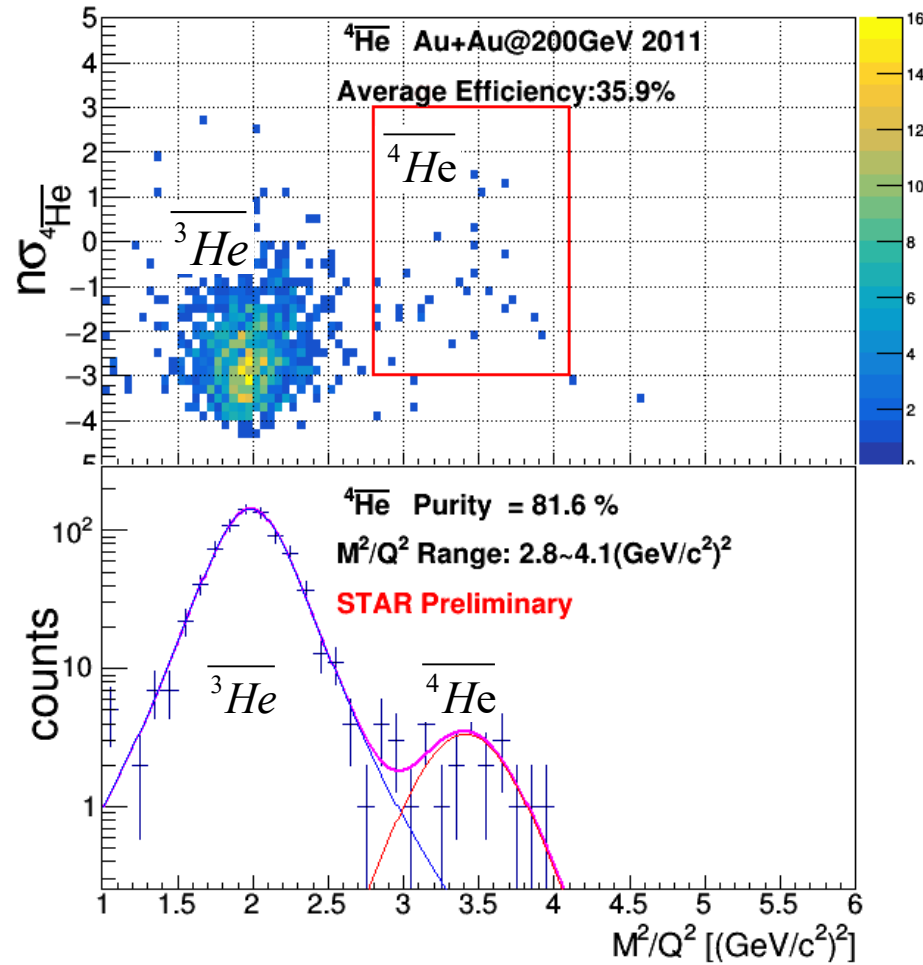
- ${}^3\text{He}$, ${}^3\overline{\text{He}}$, ${}^3\text{H}$ and ${}^3\overline{\text{H}}$: Yields are obtained by integrating over the measured p_T spectrum.

Blast Wave function fit:

$$\frac{1}{2\pi p_T} \frac{d^2 N}{dp_T dy} \propto \int_0^R r dr m_0 I_0\left(\frac{p_T \sinh \rho}{T}\right) K_1\left(\frac{m_T \cosh \rho}{T}\right)$$

• [Physical Review C Volume48, Number5, 1993](#)

Back Up: Yield Ratios Measurement - A = 4 Particles



- For A = 4 particles, the yields are too low to obtain a p_T spectrum.
- An average efficiency is obtained for the whole measured p_T range, assuming Blast Wave functional shape with the same T and β as those of A = 3 particles.