Two-loop Feynman integrals for top-quark pair plus jet production

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Status of multi-scale two-loop QCD corrections



Feynman Integrals for Planar Topologies



Outlook

Motivation

High-Luminosity LHC Plan



Experimental precision ~ $\mathcal{O}(1\%)$ for many observables

NNLO QCD Corrections required to reduce theoretical uncertainty

ttj production can be used to extract numerical value of top-quark mass m_{t} 53

$$\mathscr{R}(m_t^R,\rho_s) = \frac{1}{\sigma_{t\bar{t}j}} \frac{d\sigma_{t\bar{t}j}}{d\rho_s} (N)$$

Current theoretical predictions at NLO in QCD

- Theoretical uncertainty $\delta \sim O(10\%)$
- Higher-order corrections might be crucial to reduce uncertainties



$$\rho_s = \frac{2m_0}{m_{t\bar{t}j}}$$

[Alioli, Fernandez, Fuster, Irles, Moch, Uwer '13]

[Dittmaier, Uwer, Weinzierl '07] [Melnikov,Schulze '10]

[Alioli, Fuster, Garzelli, Gavardi, Irles, Melini '22]

Analytic structure of two-loop five-point processes with internal massive propagators

Status of two-loop multi-scale computations

Gehrmannm, Henn, Lo Presti, Abreu, Dixon, Herrmann, Page, Zeng, Chicherin, Wasser, Zhang, Zoia, Sotnikov, Ita, Moriello, Tschernow, Canko, Papadopoulos, Syrrakos, Badger, Brønnum-Hansen, Hartanto, Peraro, Dormans, Febres Cordero, Heinrich, Pascual, Chawdhry, Mitov, Poncelet, Czakon, Agarwal, Buccioni, von Manteuffel, Tancredi, Kryś, Kallweit, Wiesemann, Marcoli, Moodie, Popescu Catani, Devoto, Grazzini, Mazzitelli, Savoini,...

AN AN A RAIN TO BE AN AND A RAIN AND THE AND	
$pp \rightarrow jjj$	Leading Color
$pp \rightarrow \gamma \gamma j$	Leading Color (Planar)
$pp \rightarrow \gamma \gamma \gamma$	Leading Color (Planar)
$pp \rightarrow \gamma \gamma j$	
$gg \to \gamma\gamma g$	
$pp \rightarrow Wb\bar{b}$	Leading Color
$pp \rightarrow Wjj$	Leading Color
$pp \rightarrow Zjj$	Leading Color (Planar)
$pp \rightarrow W\gamma j$	Leading Color (Planar)
$pp \rightarrow Hb\bar{b}$	Leading Color
$pp \rightarrow t\bar{t}H$	Two-Loop Approx.



Adapted from Sotnikov



Amplitude Computation Pipeline





IBP reduction and amplitude reconstruction performed exploiting Finite Fields method

[von Manteuffel, Schabinger '14] [Peraro '16]

Master Integrals Computation Status



Master Integrals Computation Status

ttj [Badger,MB,Chaubey,Marzucca '22]

One Massive Internal Propagator

Two Massive External Particles



ttH [Febres Corder, Figueiredo, Kraus, Page, Reina '23]

Two Massive Internal Propagators

Three Massive External Particles







Scattering Kinematics:

$$p_1^2 = p_2^2 = m_t^2 \quad p_3^2 = p_4^2 = p_5^2 = 0$$



- **One Internal Massive propagator**
- Topology described by 88 MIs
- **Differential Equations:**

Canonical basis

Alphabet



AMFLOW: High-precision numerical boundary [Liu,Ma,Wang '18] conditions [Liu,Ma '23]

DiffExp: Semi-Analytic solution with generalised power series expansion



Contributes to the VV planar leading color

alised [Moriello '19] [Hidding '20]

Strategy of the Computation





Analytic Information: Canonical Basis, Alphabet

A First Look at Planar sector of five-point two-loop scattering amplitude with massive propagators

Two-Loop Planar Feynman Integrals for ttj

Massive propagator increases complexity of calculation

	Planar Massless	Planar one-mass	Planar Etj	
Number of MIs	61	74-86	88-123	
Length of Alphabet	31	38-49	71	
Number of roots	1		6	

However, the structure of the Canonical Basis and of the Alphabet are similar....



Canonical basis

[Henn '13]

 $d\vec{f}(\vec{x},\epsilon) = \epsilon \, dA(\vec{x}) \, \vec{f}(\vec{x},\epsilon)$

Rational Letters

$$(\Delta_i)^2$$

 $\operatorname{tr}(ij\cdots k) = \operatorname{tr}(\gamma_{\mu}p_{i}^{\mu}\gamma_{\nu}p_{j}^{\nu}\cdots\gamma_{\rho}p_{k}^{\rho})$

Similar structure to five-point massless and one off-shell leg cases

$$(x), A(\vec{x}) = \sum_{i=1}^{71} c_i \log(w_i(\vec{x})))$$

$$Algebraic$$

$$Letters$$

$$(\Omega_i(a, b)) \qquad (\underline{\mathrm{tr}_+(i\cdots j)}) \qquad (\tilde{\Omega}_i(a, b, c))$$

$$(\Omega(a, b) := \frac{a + \sqrt{b}}{a - \sqrt{b}} \qquad \mathrm{tr}_\pm(i\cdots j) = \frac{1}{2} \operatorname{tr}((1 \pm \gamma_5)\gamma_\mu p_i^\mu \cdots \gamma_4)$$

$$(\tilde{\Omega}(a, b, c) := \frac{(a + \sqrt{b} + \sqrt{c})(a - \sqrt{b} - \sqrt{c})}{(a + \sqrt{b} - \sqrt{c})(a - \sqrt{b} + \sqrt{c})}$$







Canonical basis depends on the set of square-roots:

$$\beta = \sqrt{1 - \frac{4m_t^2}{s_{12}}} \qquad \Delta_1 = \sqrt{\det G(p_{23}, p_1)} \qquad \Delta_2 = \sqrt{\det G(p_{15}, p_2)} \qquad \Delta_3 = \sqrt{1 - \frac{4s_{45}m_t^2}{(s_{12} + s_{23} - m_t^2)^2}} \qquad \Delta_4 = \sqrt{1 + \frac{4s_{34}s_{45}m_t^2}{s_{12}(s_{15} - s_{23})^2}}$$

 $tr_5 = 4\sqrt{\det G(p_3, p_4, p_5, p_1)}$





★ 32 cores workstation



We exploit Finite Fields method as implemented in FiniteFlow to obtain analytic DEQs [Peraro '19] Reconstruction Degrees time 53/57 Non-UT ~3 Weeks 15/15 Quasi-UT ~ 2 Hours

$$\overline{O} \qquad G_{ij}(\vec{v}) = v_i \cdot v_j$$

Differential Equations with Finite Fields



$$(\vec{x}) \widehat{A}^{(1)}(\vec{x}) N^{-1}(\vec{x}) \widehat{I}(\vec{x}, \epsilon)$$

 $V_{ij}(\vec{x})$ Diagonal matrix with
square-roots

Building Canonical Basis of MIs





We exploit emerging patterns for five-point processes in the construction of the UT basis



Given the complexity of the kinematics automated approaches are difficult to apply

Emerging Patterns

five-point MIs

Scalar integrals with numerators and local integrand insertions



Local integrand:
$$\mu_{ij} = -k_i^{[-2\epsilon]} \cdot k_j^{[-2\epsilon]}$$
 Numerators: $\mathcal{N} = (k_i + q)^2$



More in depth analysis performed exploiting Leading Singularities, Loop-by-Loop...



A cooking recipe for Canonical Basis



Candidates choice inspired by similar processes Fast pipeline to generate and tests possible candidates integrals for Canonical Basis

Generalised Power Series Evaluation: A proof of concept

We exploit the Generalised Power Series method as implemented in DiffExp

Series Solution around singular points of DEQs

$$\vec{f}(t,\epsilon) = \sum_{k=0}^{\infty} \epsilon^k \sum_{i=0}^{N-1} \rho_i(t) \vec{f}_i^{(k)}(t), \quad \rho(t) = \begin{cases} 1, & t \in \left[t_i - r_i, t_i + r_i\right) \\ 0, & t \notin \left[t_i - r_i, t_i + r_i\right) \end{cases}, \qquad \vec{f}_i^{(k)}(t) = \sum_{l_1=0}^{\infty} \sum_{l_2=0}^{N_{i,k}} c_k^{(i,l_1,l_2)} \left(t - t_i\right)^{\frac{l_1}{2}} \log(t - t_i)^{l_2} \left(t - t_i\right)^{\frac{l_1}{2}} \log(t - t_i)^{l_2} \left(t - t_i\right)^{\frac{l_1}{2}} \log(t - t_i)^{\frac{l_2}{2}} \log(t - t_i)^{\frac{$$

- Numerical boundary values obtained with AMFlow
- Numerical evaluation of MIs in whole phase-space
- Proof of Concept for this computation
- **Optimisation could lead to phenomenological applications**

ttj Leading Colour Planar Amplitude



- Numerical implementation of all the MIs through generalised power series expansion
- - Optimisation of power serie method for phenomenological application

Numerical reconstruction and evaluation of leading colour planar amplitude for benchmark points

ttj Leading Colour Planar Amplitude

Preliminary results



- **111 MIs for TopoC**
- Differential Equations for TopoC in epsilon factorised form
- New analytic structure: MIs basis contains nested square roots normalisation,...





Summary

- First computation for two-loop five-point Feynman Integrals with Internal Massive Propagators



UT structure follows emerging pattern amongst five-point processes

Outlook



Computation of all the MIs relevant for the NNLO ttj production in the planar leading color limit

Efficient numerical evaluation: optimisation of generalised power series method for phenomenological applications

MIs computation for two-loop planar topology represents the first ingredient for a NNLO QCD corrections to ttj



Thank you for your attention!