

THE THREE LOOP AMPLITUDE FOR V+jet PRODUCTION

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based on 2307.15405, 2306.10170 and 2301.10849,
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Summary

1. Introduction

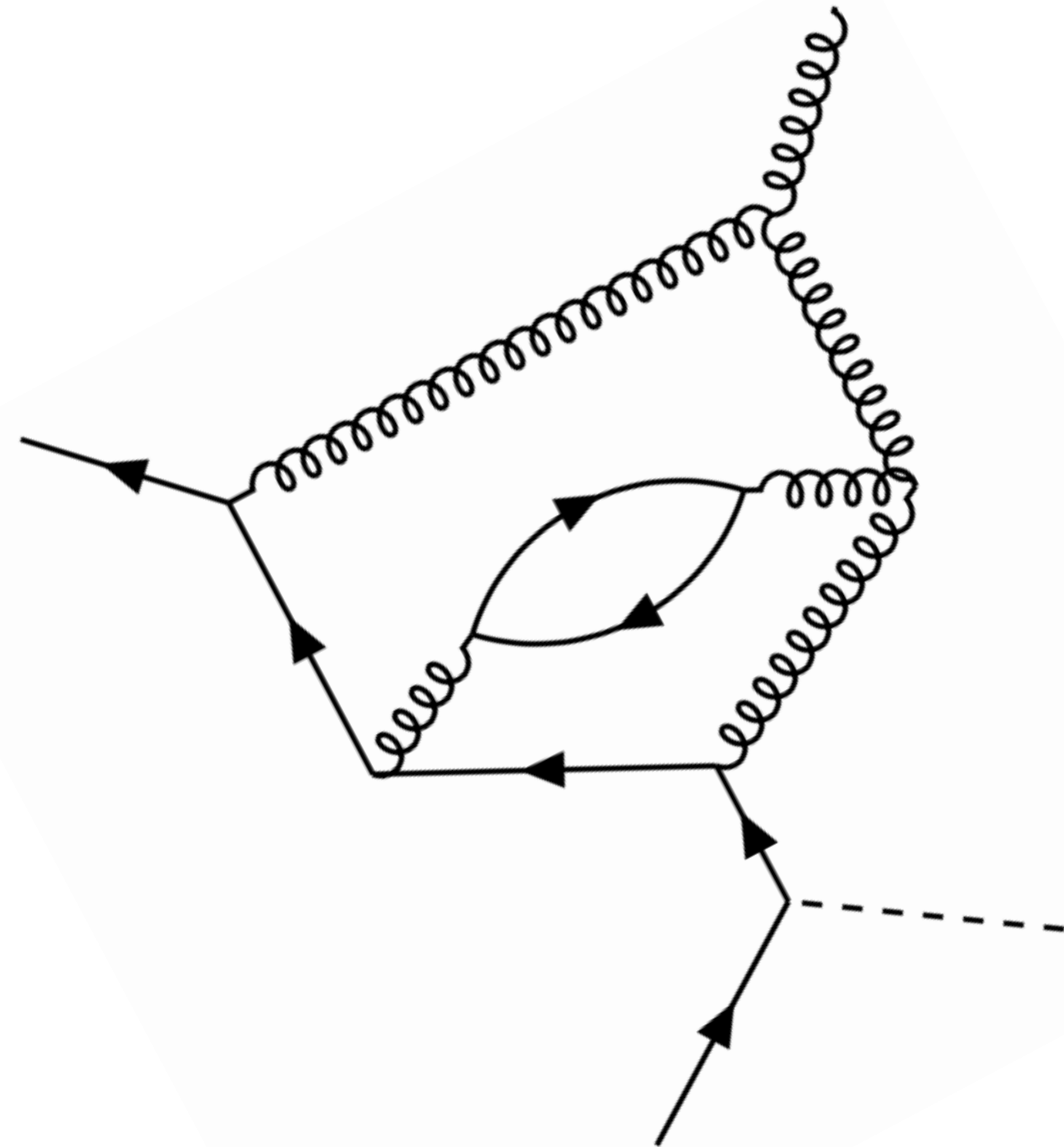
- Do we need N3LO for V+jet production?
- State of the Art in theory prediction

2. The Calculation

- Tensor decomposition
- The workflow
- Results

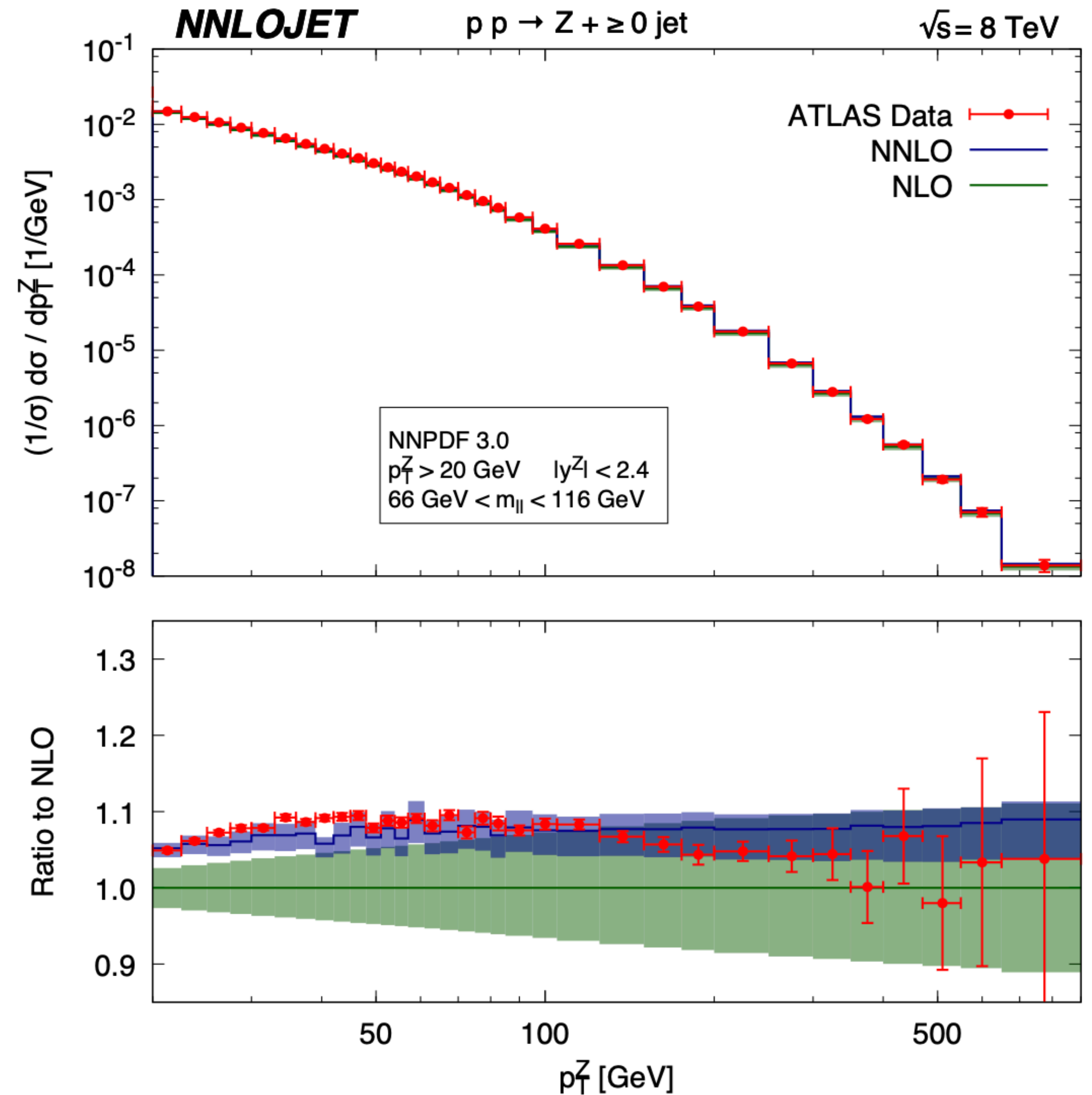
3. What's next?

- Full amplitude, non planar 3L master integrals, ...



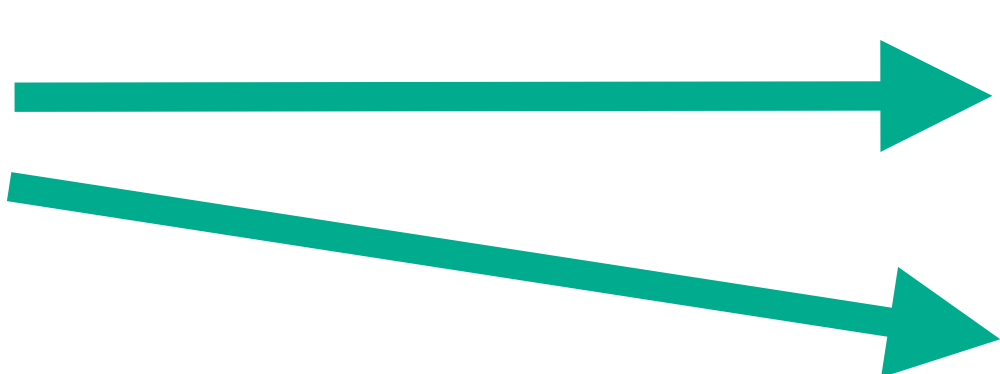
Do we need N3LO for V+jet?

- Determination of α_s and constraint on PDFs
- Measures of differential distribution for Z production extremely accurate $\mathcal{O}(< 1\%)$ [CMS, ATLAS]: high rate and clear leptonic signature
- NNLO for V+jet obtained via antenna subtraction [Gehrmann - De Ridder, Gehrmann, Glover, et al., 2016], N-jettiness [Boughezal, Liu, Petriello, 2016]
 - $Z - p_T$ distribution, leading jet distribution, Angular coefficients \rightarrow details of gauge boson and its polarisation states



State of the Art in theory predictions

$$d\sigma_{pp \rightarrow X} = \sum_{i,j} \int \int dx_1 dx_2 f_{p_1,i}(x_1, Q_F) f_{p_2,j}(x_2, Q_F) \times d\hat{\sigma}_{ij \rightarrow X}(x_1 P_1, x_2 P_2, Q_F) + \mathcal{O}((\Lambda_{QCD}/Q)^p)$$

- Among the ingredients to the $\mathcal{O}(1\%)$ accuracy is the precise determination of partonic cross section
- Perturbative expansion 
 - Multiloop and higher multiplicity amplitudes
 - Suitable subtraction procedure (Antenna, Slicing, ...)

On the amplitudes side

- 4 point massless master integrals [Henn et al. '20]
 - 4 points 3L massless amps. ($gggg, qqgg, 4q, qq\gamma\gamma, \dots$) [Caola, Chakraborty, Gambuti, von Manteuffel, Tancredi, '21-'22]
- ➔ 4 point with one external mass @3L is one of the next things to do

State of the Art in theory predictions: new results!

Perturbative expansion for V + jet amplitude: the N3LO

$$d\sigma_{q\bar{q}\rightarrow Vj} = d\sigma_{q\bar{q}\rightarrow Vj}^{LO} + \left(\frac{\alpha_s}{2\pi}\right) d\sigma_{q\bar{q}\rightarrow Vj}^{NLO} + \left(\frac{\alpha_s}{2\pi}\right)^2 d\sigma_{q\bar{q}\rightarrow Vj}^{NNLO} + \mathcal{O}(\alpha_s^3)$$

✓ Tree Level V + 4j

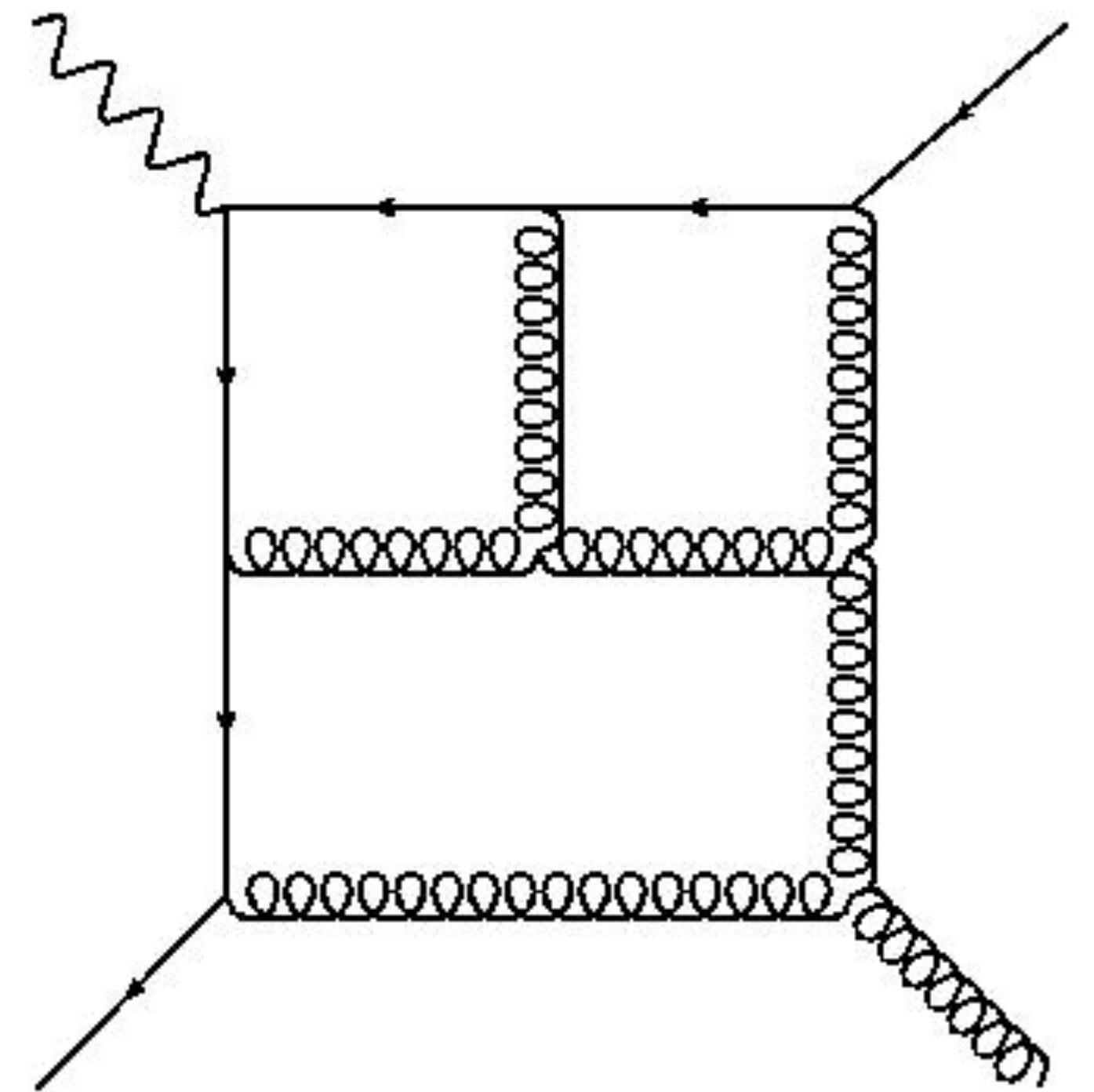
✓ 1- Loop V + 3j

▶ 2- Loop V + 2j, all integrals [Abreu, Chicherin, Ita, Page, Sotnikov, Tschernow, Zoia, 2023], and leading color amps. [Abreu, Cordero, Ita, Klinkert, Page, Sotnikov, 2021]

★ 3- Loop V + j is work in progress, **RESULTS!**

1. 2- Loop V + j at transcendental weight six, including axial-vector coupling, done!

2. 3- Loop V + j in leading color approximation, done!



Computing amplitudes: tensor decomposition

- We are interested in helicity amplitudes
 - We work with 4-dim external states ('t Hooft-Veltman scheme) and dim-reg for the loops
 - Write a decomposition based on external particles and symmetries

Use gauge symmetry, gluons transversality and kinematics to constrain the tensor structures



$$\mathcal{A} = \sum_{i=1}^N F_i T_i$$

- Define projectors to form factors $P_j = \sum_{i=1}^N c_{(j)}^i T_i^\dagger$, such that $\sum_{pol} P_j \mathcal{A} = F_j$, see [Peraro, Tancredi, 2021]
- 24 tensor structures for $Zggg$ and 12 for $Zgq\bar{q}$
- Helicity amplitude projectors as linear combinations of the form factor projectors

Computational setup

1. Diagrams generated with QGRAF [Nogueira, 1993]
2. Feynman rules, colour algebra and projection onto form factors is performed in FORM [Vermaseren, 1999]

► The resulting integrand is a combination of scalar integrals

$$F = \sum_{l=1}^{N_{int}} r_l(s_{ij}, d) Int_l$$

3. Integral **reduction** performed with KIRA [Klappert, Lange, Maierhöfer, Usovitsch, 2021] and REDUZE [Studerus, von Matteuffel, 2012]

4. Insert reductions and **solution**, UV renorm., IR subtr., analytic continuation

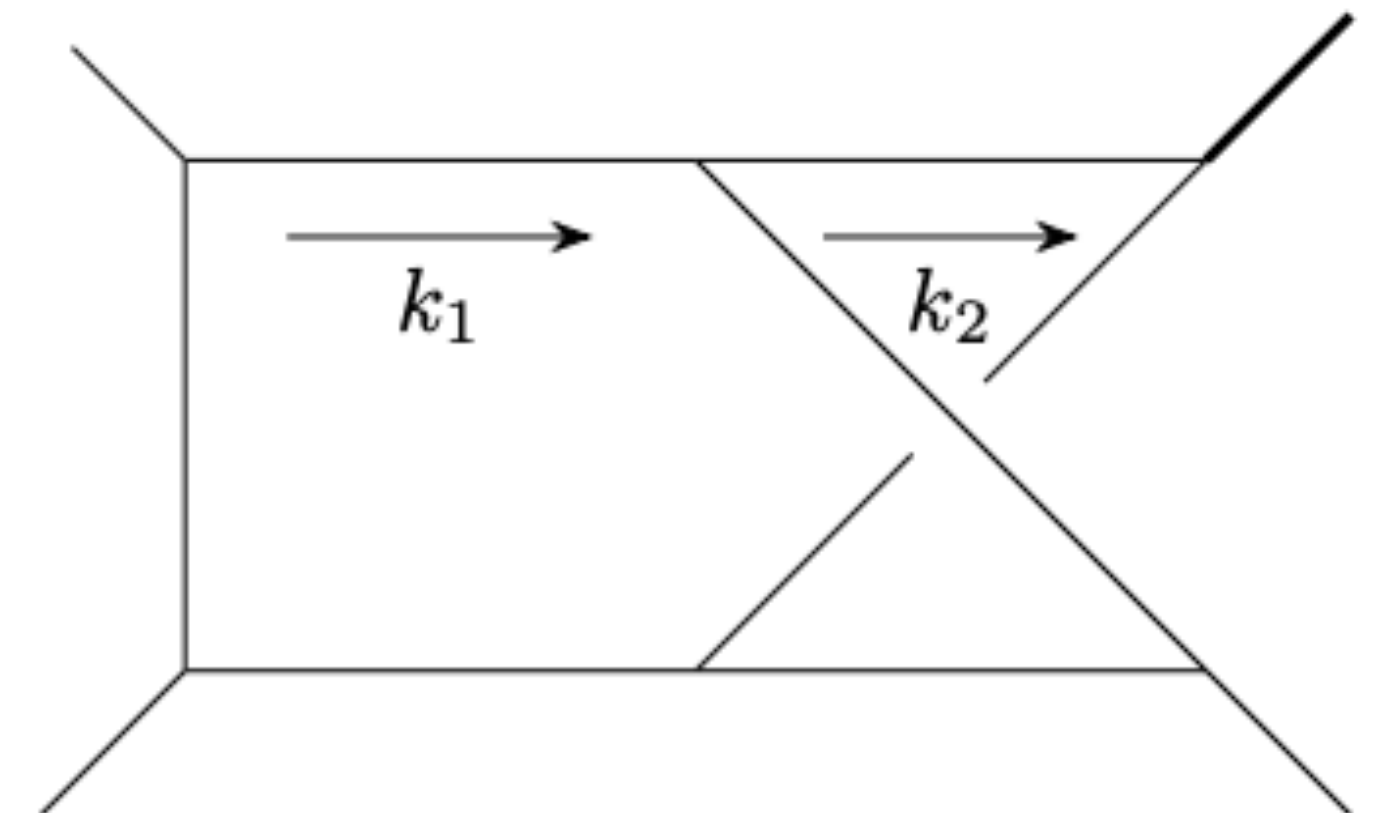
At 3L non negligible **analytic** and **computational** complexity.

The 2 loop amplitudes at $\mathcal{O}(\epsilon^2)$: master integrals and UV renormalisation

- 4 point with one external mass @2L [Gehrmann, Remiddi 2001]
- ✓ We provide a canonical basis [Henn, 2012], using the package DLogBasis [Henn, Mistlberger, 2020]
- ✓ Solution in terms of GHPLs to $\mathcal{O}(\epsilon^2)$ using PolyLogTools (not at weight 6!) [Duhr, Dulat, 2019]
Alphabet $\{y, z, 1 - y, 1 - z, y + z, 1 - y - z\}$, where $y = s_{13}/m_V^2$, $z = s_{23}/m_V^2$.
- The axial-vector amplitude needs an **extra finite renormalisation** when working in the Larin scheme, $\gamma^\mu \gamma^5 \rightarrow \frac{1}{6} \epsilon^{\mu\nu\rho\sigma} \gamma_\nu \gamma_\rho \gamma_\sigma$, in order to restore the axial Ward Identities [Larin, 1993]

- We checked the renormalisation against the axial anomaly relation

$$\partial_\mu \langle J_5^\mu(x) \mathcal{O}(x_1, \dots, x_n) \rangle = \frac{\alpha_s}{8\pi} \langle G_{\mu\nu}^a \tilde{G}_{\mu\nu}^a(x) \mathcal{O}(x_1, \dots, x_n) \rangle$$



The 3 loop leading color for $Vgq\bar{q}$

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$$\begin{aligned}\mathcal{A}^{(3)} = & N^3 \Omega_1^{(3)} + N \Omega_2^{(3)} + \frac{1}{N} \Omega_3^{(3)} + \frac{1}{N^3} \Omega_4^{(3)} + N_f N^2 \Omega_5^{(3)} + N_f \Omega_6^{(3)} + \frac{N_f}{N^2} \Omega_7^{(3)} \\ & + N_f^2 N \Omega_8^{(3)} + \frac{N_f^2}{N} \Omega_9^{(3)} + N_f^3 \Omega_{10}^{(3)} + N_{f,V} N^2 \Omega_{11}^{(3)} + N_{f,V} \Omega_{12}^{(3)} \\ & + \frac{N_{f,V}}{N^2} \Omega_{13}^{(3)} + N_f N_{f,V} N \Omega_{14}^{(3)} + \frac{N_f N_{f,V}}{N} \Omega_{15}^{(3)}, + N_f^2 N_{f,V} \Omega_{16}^{(3)}\end{aligned}$$

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 \end{aligned}$$

- Since $N_F \approx N_C$, we consider $N_F^2 N_C, N_C^2 N_F, N_C^3$ \longrightarrow only planar contributions!
- Planar topologies at 3 loop computed [Canko, Syrrakos, '22]
- Alphabet the same as at two loop $\{y, z, 1 - y, 1 - z, y + z, 1 - y - z\}$

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- Unreduced contribution $F = \sum_{l=1}^{N_{int}} r_l(s_{ij}, d) Int_l$ with $N_{Int} \approx 100k$

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KIRA

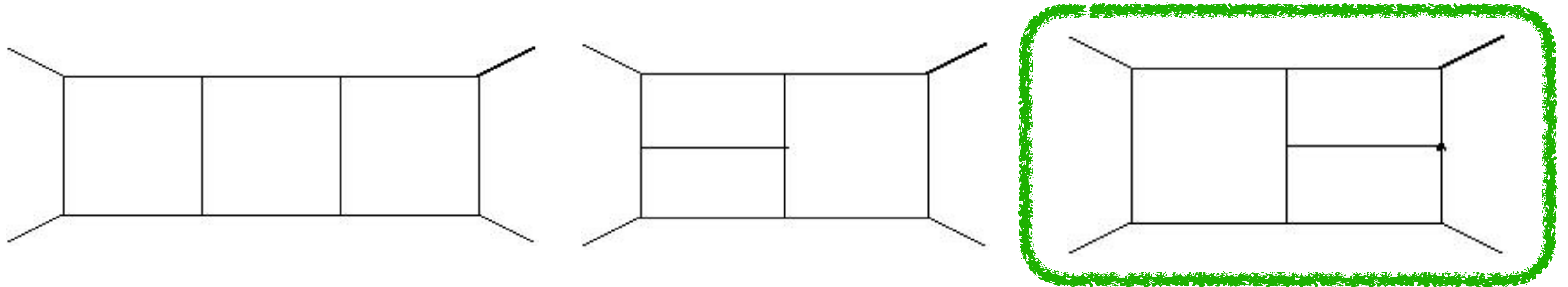


Reduction to masters of family with 10+5 (ISP) = 15 propagators with ints up to rank 5: **expensive!**

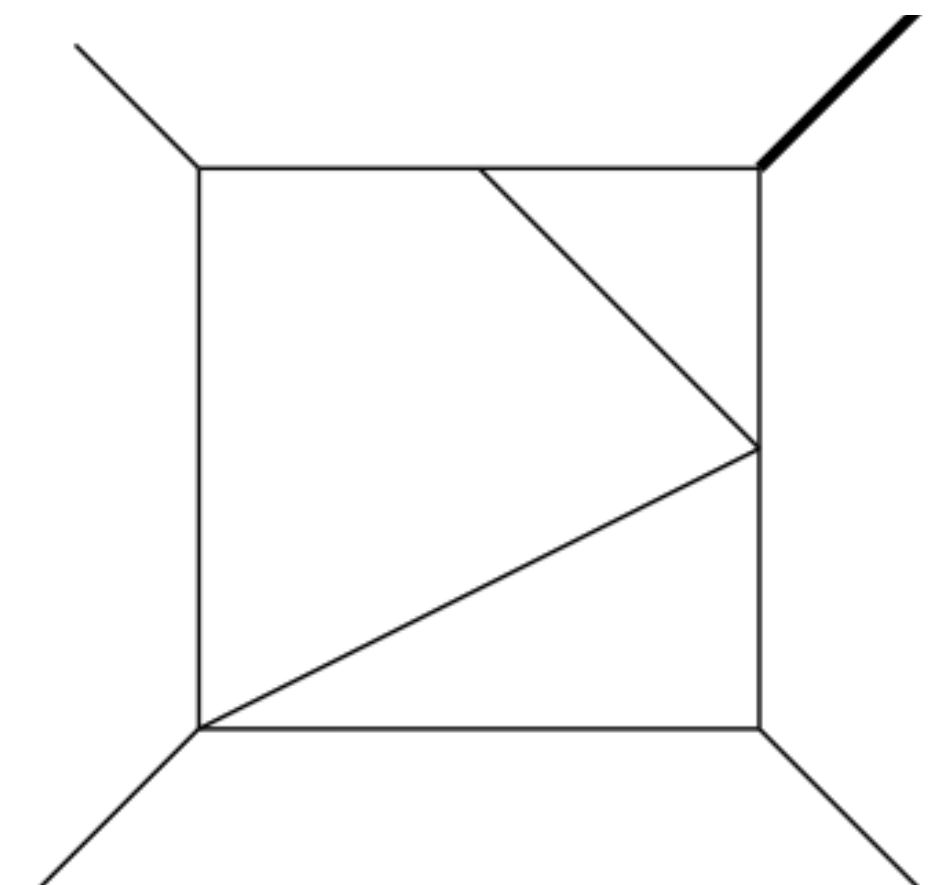
- Reduction in terms of 291 master integral (mod crossings)
- The insertion of the reduction and master solutions both analytically and with finite field reconstruction using FiniteFlow [Peraro, 2019]
 - We reconstructed only the coefficients of the appearing GHPLs

The 3 loop leading color for $Vgq\bar{q}$: the adjacency condition

- It was noticed in $\mathcal{N} = 4$ super-YM that the letters $\{1 - x, 1 - y, 1 - z\}$ never appear next to each other in the symbol [Dixon, McLeod, Wilhelm, 2021] \longrightarrow adjacency conjecture
- In the case considered, it holds up to two loop at the level of the integrals

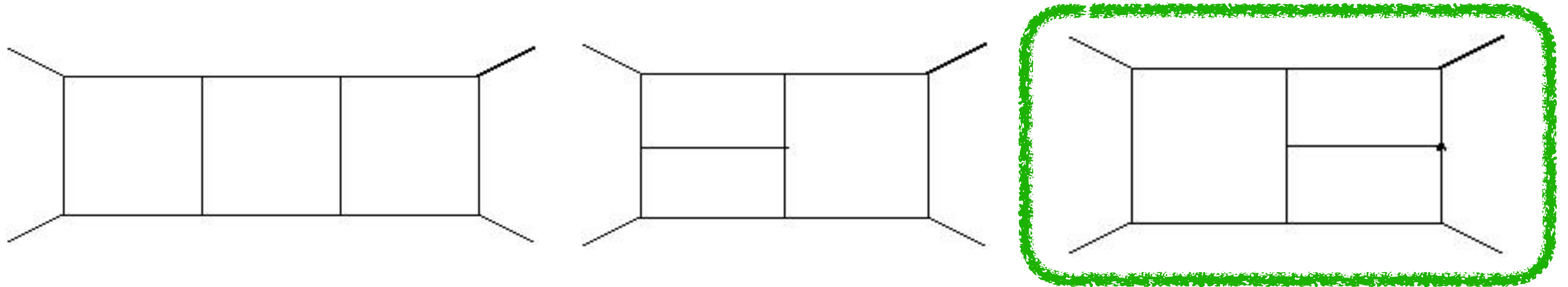


- Together with the evaluation of the first non-planar topologies, the adjacency conjecture was found to be violated in one of the planar topologies at 3L [Henn, Lim, Bobadilla, 2023]



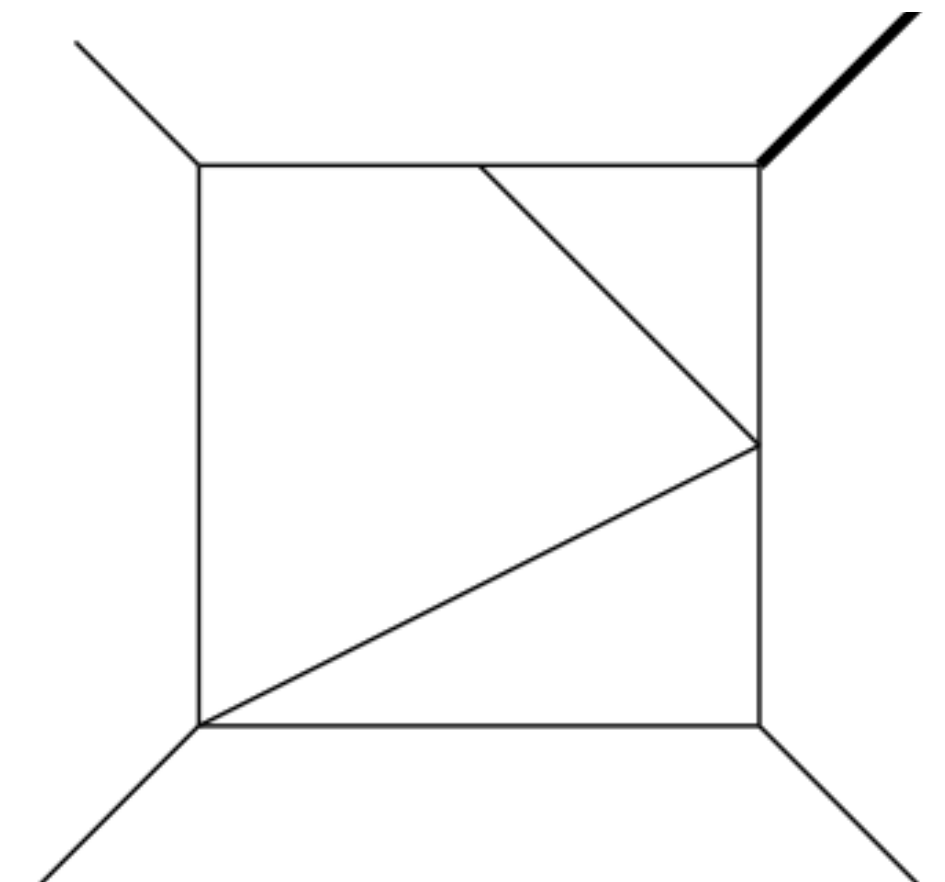
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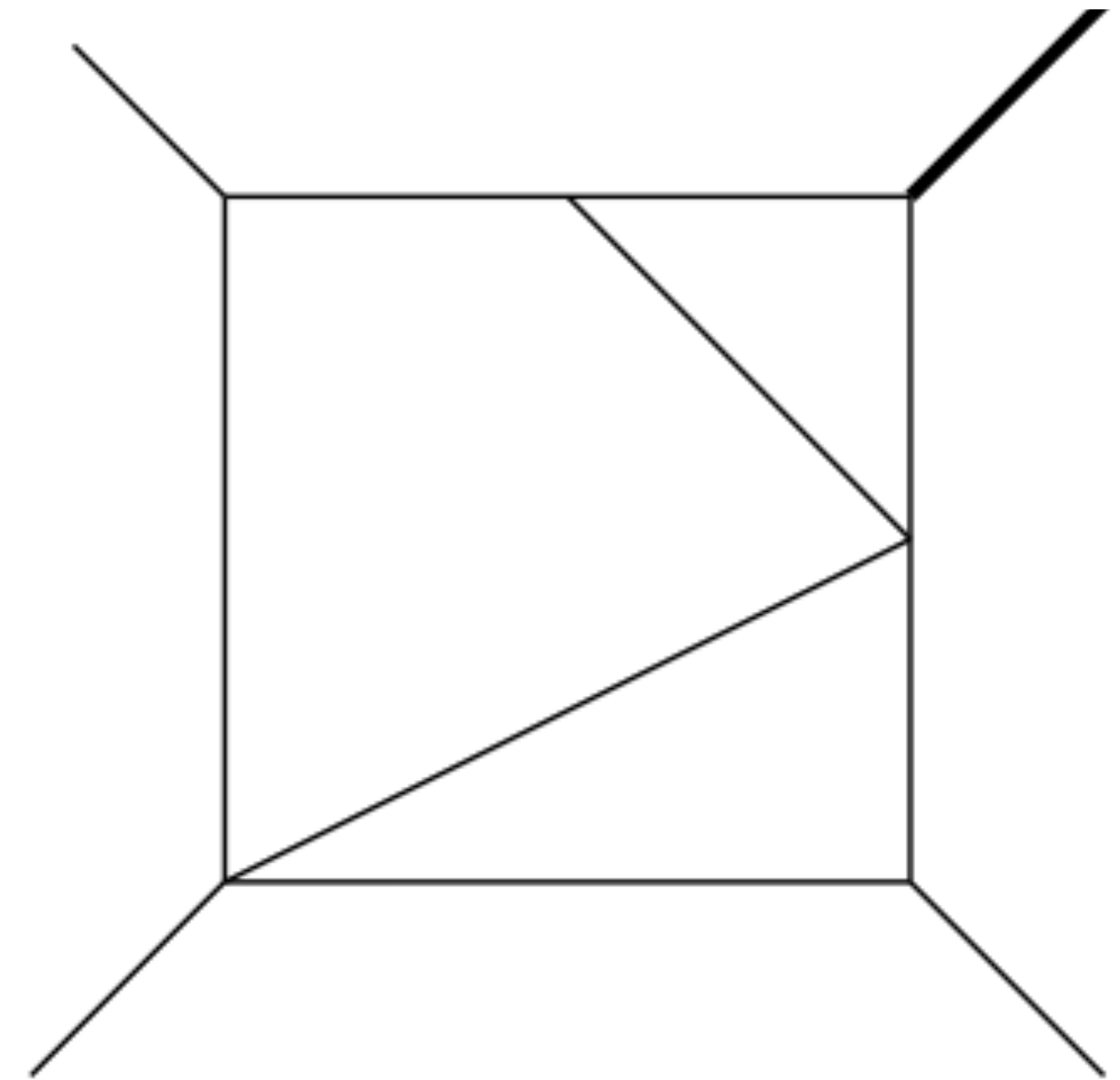
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?
 $\dots \otimes 1 - x \otimes 1 - y \otimes \dots$
?



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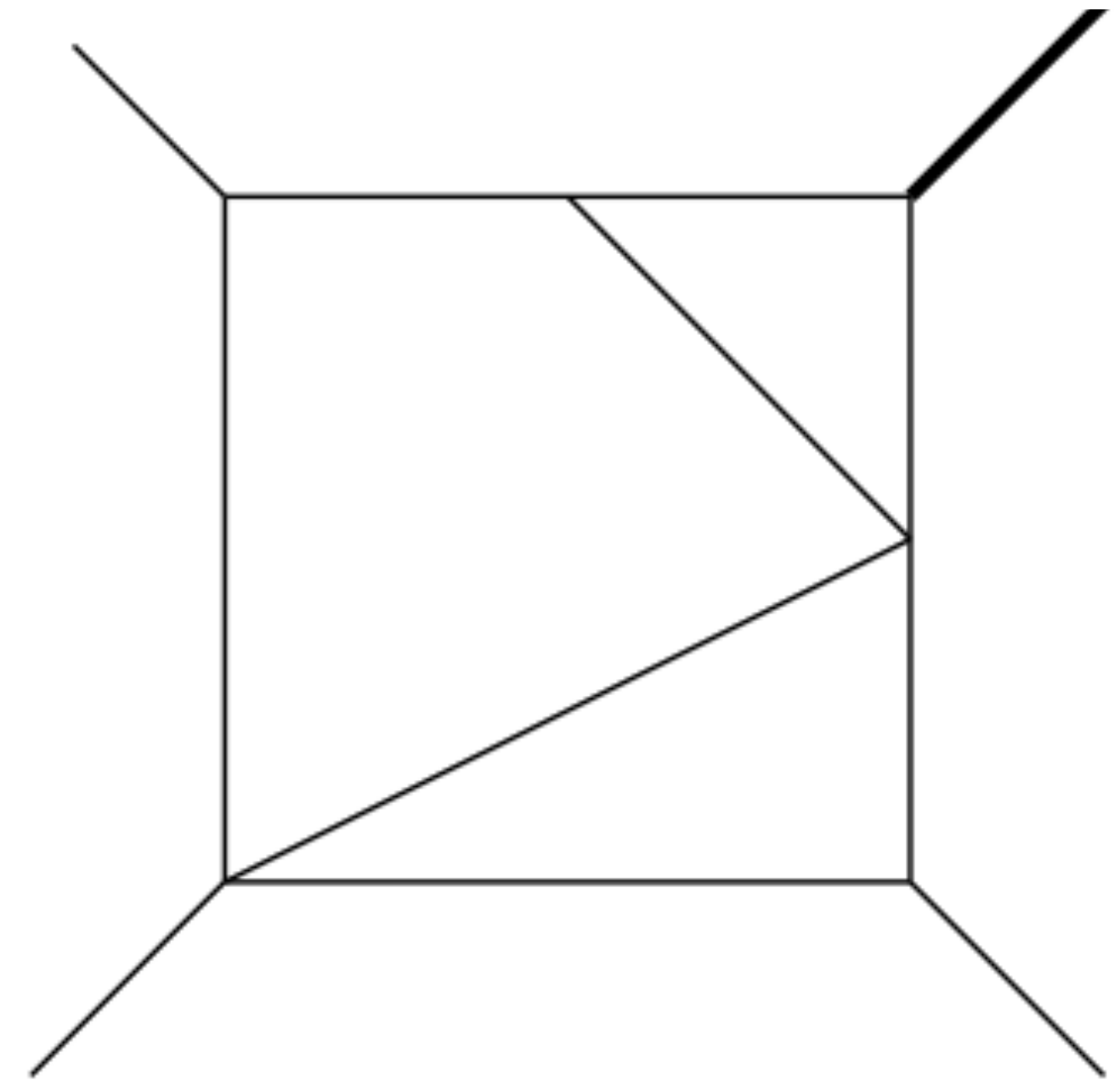
- $d\vec{I} = \epsilon \sum_i A_i d \log \alpha_i \vec{I}$ where $\alpha_i \in \{y, z, 1 - y, 1 - z, y + z, \dots\}$
 - In term of the diff equation matrices, the condition which is violated is: $A_i \cdot A_j = 0$ whenever $\alpha_i, \alpha_j \in \{1 - x, 1 - y, 1 - z\}$



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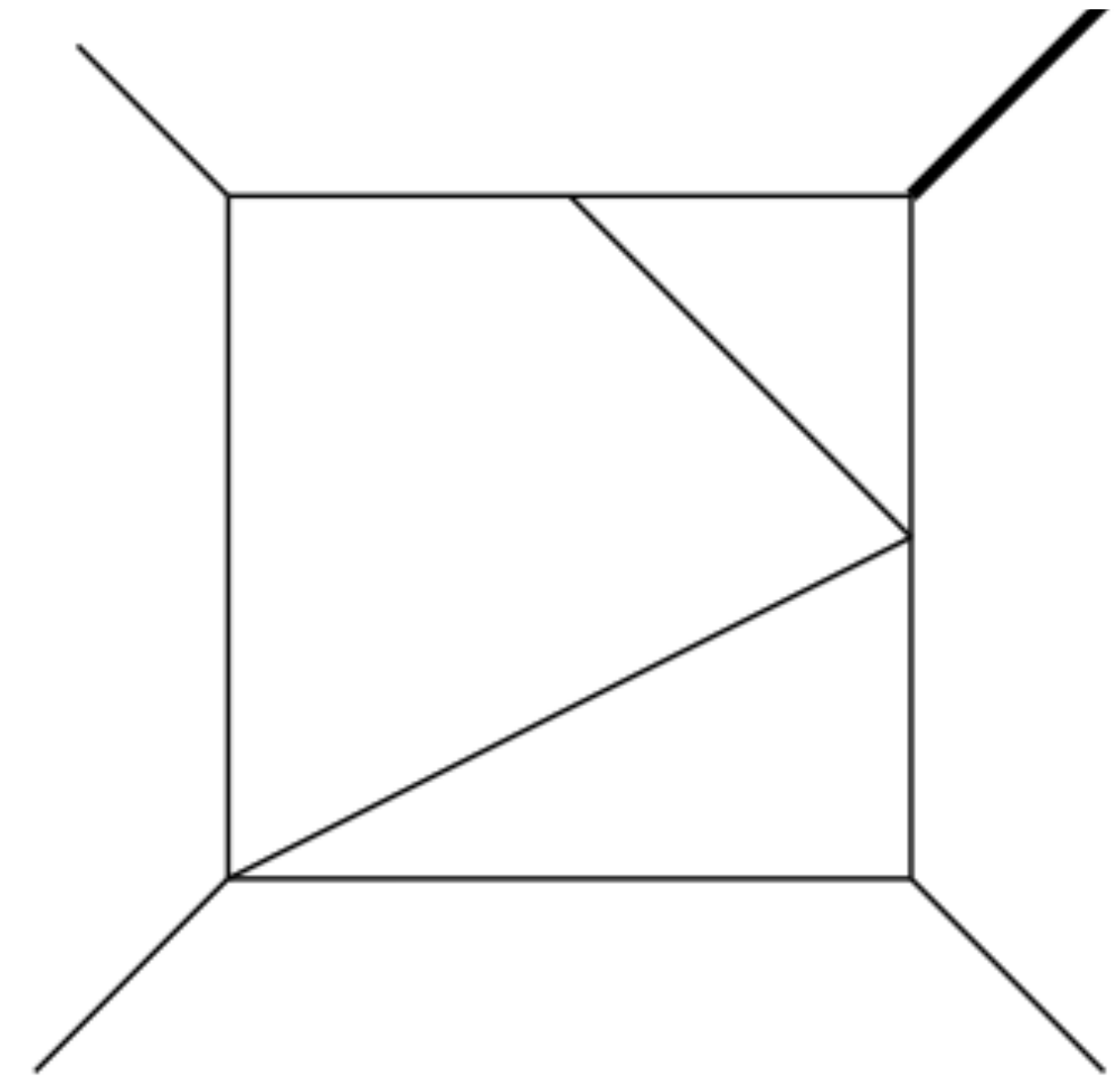
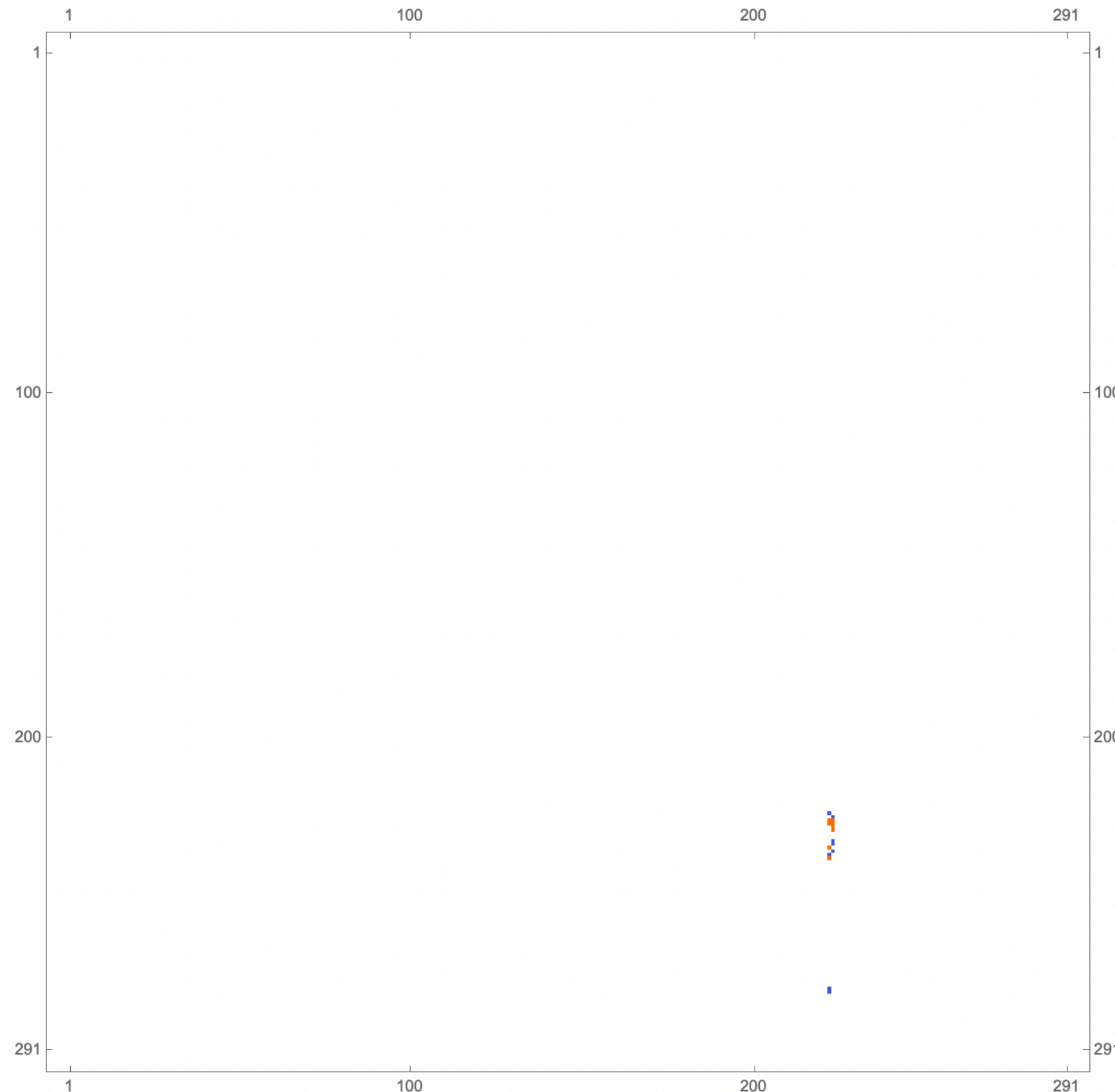
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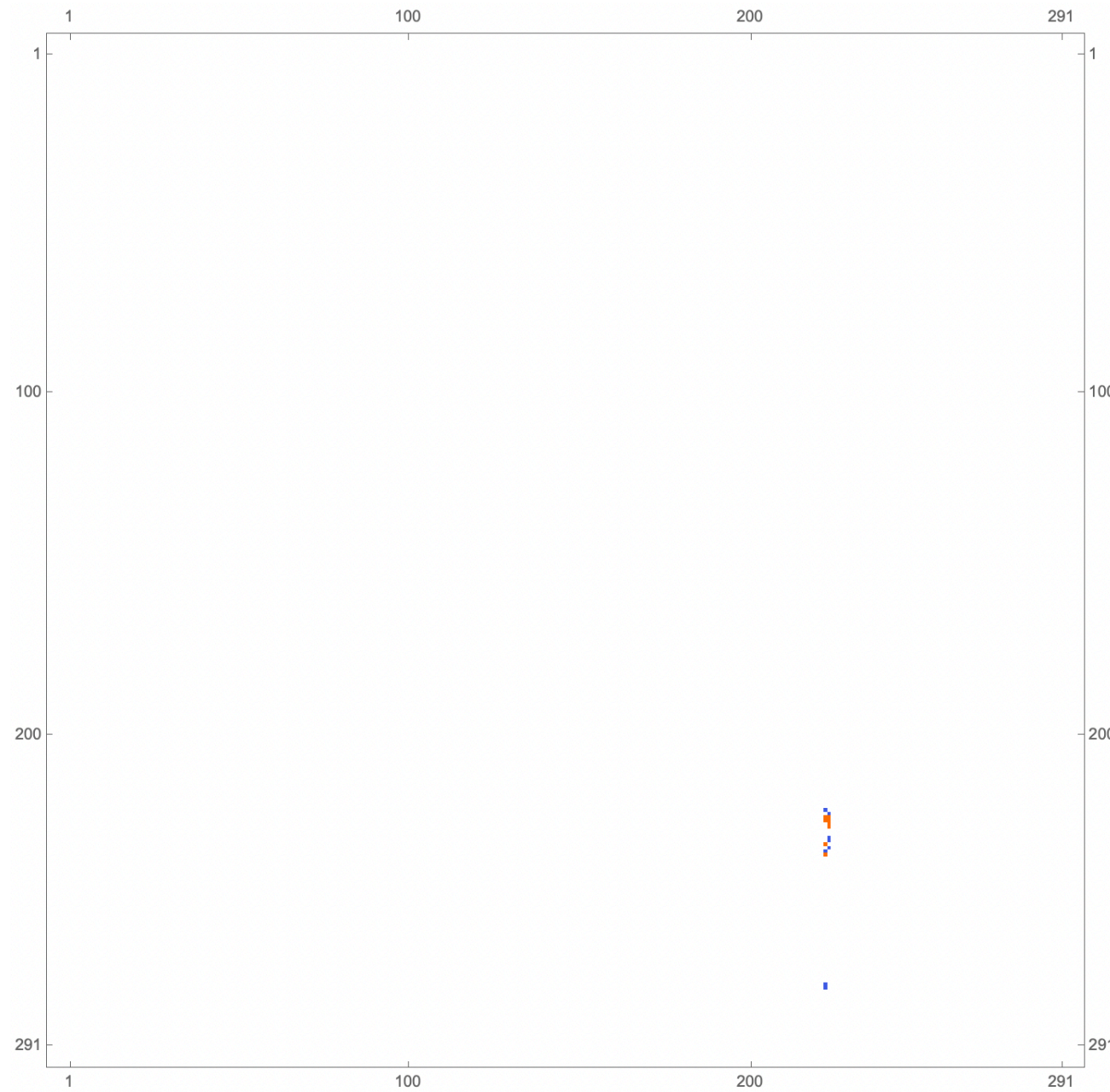
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$$\dots \otimes 1-x \otimes 1-y \otimes \dots$$

- Surprisingly, the conjecture is **satisfied** at the level of the amplitude
- Because of the propagation of the conjecture-violating symbols in the upper sectors, it is not easy to trace the cancellation

Results

- Canonical basis for the for 2L and extension up to $\mathcal{O}(\epsilon^2)$
- Solution in terms GHPLs with simple alphabet $\{y, z, 1 - y, 1 - z, y + z, 1 - y - z\}$, fast numerical evaluation
- Completed the missing vector-axial calculation and 3L leading color for $V + jet$
- UV renormalised and IR subtracted helicity amplitudes analytically continued to all physical regions:
 $V \rightarrow gq\bar{q}, q\bar{q} \rightarrow Vg, qg \rightarrow Vq, \bar{q}g \rightarrow V\bar{q}$
- Results have been check numerically at 1L with OpenLoops2 [Buccioni et al., 2019], and against available literature at 2L

The Work in Progress: non planar ints and the Higgs amplitude

In collaboration with Johannes Henn, Jungwon Lim, William J. Torres Bobadilla

★ *TODO list: preliminary news*

- $H + jet$ Feynman diagram representation is combinatorially more complex
 - ▶ 5334 diagrams for $Zgq\bar{q}$ @3L **vs** 35827 diagrams for $Hggg$ @3L
 - ▶ Moderate swell of intermediate expressions ~ 6 Gb for $H + jet$ leading color, ~ 2 mln integrals, with non planar integrals up to rank 6
- 5 non planar top sectors contributing to the leading color of $H + j$
 - ▶ The canonical basis for them has a richer alphabet: quadratic letters and a square root
 - ▶ Ansatz approach and different representations (Baikov, Baikov loop-by-loop, 4-dim param.) to find $dlog$ candidates



Thank You!