

Diagrammatic Analysis of Charmless Three-Body B Decays

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Abstract

The diagrammatic method has been used in the past to extract, among other things the weak phase γ using observables in $B \rightarrow K\pi$ decays. We generalize the diagrammatic decomposition method to more-than-two-body decays, and systematically apply this method to the charmless decays of B mesons to three pseudoscalars. The method of contractions is used to demonstrate the existence of relations between tree and electroweak-penguin diagrams in $B \rightarrow PPP$ decays. Together with the use of diagrammatic decomposition and Dalitz-plot analyses, these relations can be used to cleanly extract the weak phase Γ from $B \rightarrow K\pi\pi$ decays.

Accounting for Symmetry

- ▶ Isospin (SU(3)) symmetry requires treating pions (and kaons) as identical particles
- ▶ Bose-Einstein symmetry requires symmetrical wave functions
- ▶ To account for this, diagrams will come in **different versions**
- ▶ Each version corresponds to a representation of permutation group S2 (for isospin) or S3 (for SU(3))
- ▶ Comparing predictions to experiment will thus require extracting the components of amplitudes with the correct symmetries
- ▶ **This can be done with Dalitz-plot amplitude analyses**

Applying Diagrams to $B \rightarrow PPP$

The dominant contribution to $B \rightarrow PPP$ decays comes from the following diagrams :

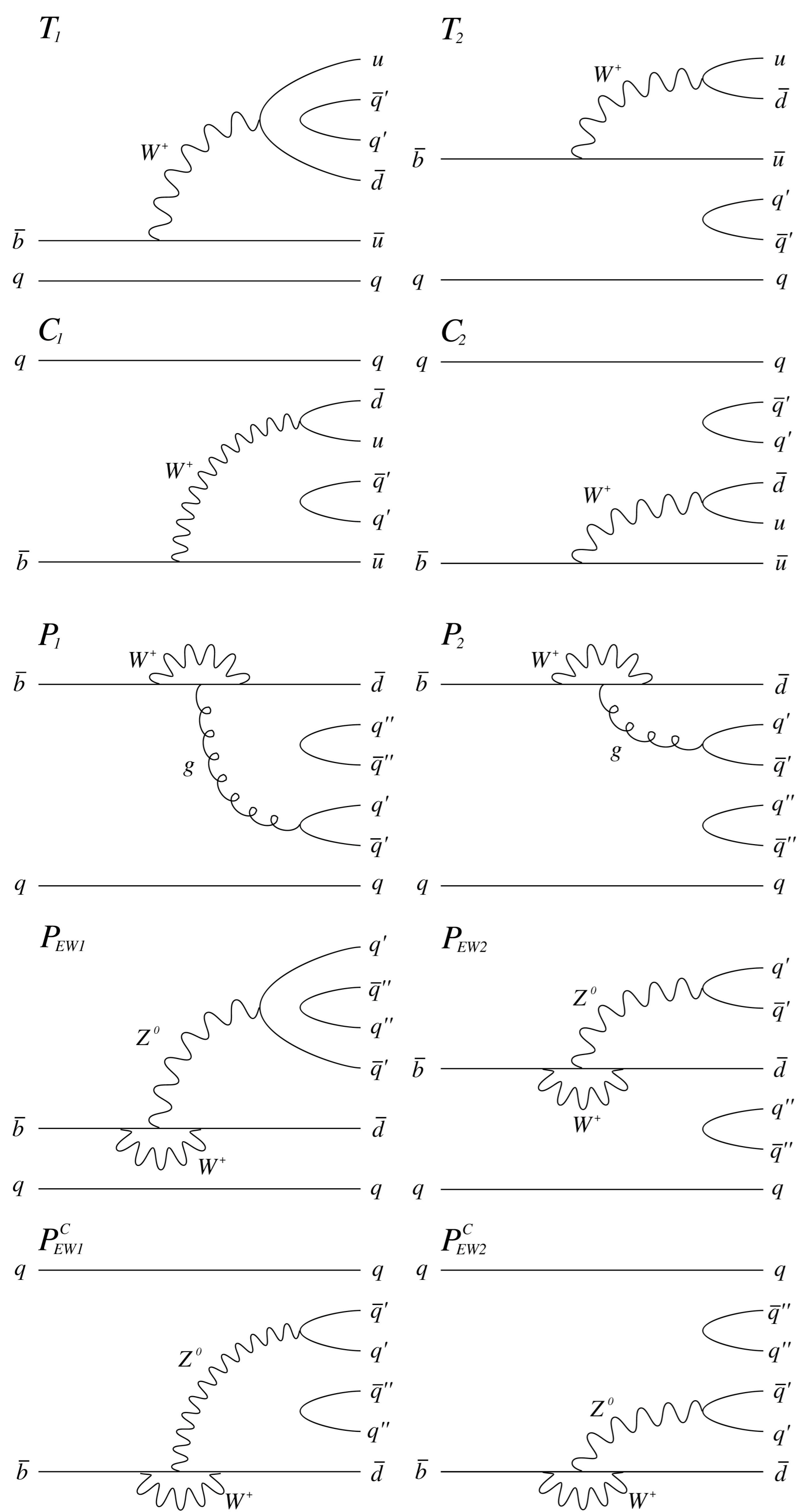


FIGURE: Diagrams for a $B \rightarrow \pi\pi\pi$ decay

- ▶ Other diagrams actively involving the spectator quark are expected to be suppressed by a factor of f_B/m_B
- ▶ Isospin symmetry establishes the equivalence of diagrams within a given type of decay. The different types are :
 - ▶ $B \rightarrow K\pi\pi$
 - ▶ $B \rightarrow KK\bar{K}$
 - ▶ $B \rightarrow K\bar{K}\pi$
 - ▶ $B \rightarrow \pi\pi\pi$

An Important Example : $B \rightarrow (K\pi\pi)_{sym}$

$$\begin{aligned} \sqrt{2}A(B^+ \rightarrow K^0\pi^+\pi^0)_{sym} &= -T'_a e^{i\gamma} - T'_b e^{i\gamma} + P'_{EW,a} + P'_{EW,b}, \\ A(B^0_d \rightarrow K^0\pi^+\pi^-)_{sym} &= -T'_a e^{i\gamma} - P'_a e^{i\gamma} + P'_b, \\ \sqrt{2}A(B^0_d \rightarrow K^0\pi^0\pi^0)_{sym} &= -T'_b e^{i\gamma} + P'_a e^{i\gamma} - P'_b + P'_{EW,a} + P'_{EW,b}, \\ A(B^+ \rightarrow K^+\pi^+\pi^-)_{sym} &= -P'_a e^{i\gamma} + P'_b - P'_{EW,a}, \\ \sqrt{2}A(B^+ \rightarrow K^+\pi^0\pi^0)_{sym} &= T'_a e^{i\gamma} + T'_b e^{i\gamma} + P'_a e^{i\gamma} - P'_b - P'_{EW,b}, \\ \sqrt{2}A(B^0_d \rightarrow K^+\pi^0\pi^-)_{sym} &= T'_a e^{i\gamma} + T'_b e^{i\gamma} - P'_{EW,a} - P'_{EW,b}, \end{aligned}$$

Extracting Symmetrical Amplitudes

- ▶ Dalitz-plots can be used to extract the amplitude components with the desired symmetries
- ▶ Isobar model : $\mathcal{M} = \sum a_j F_j(x, y)$
- ▶ a_j coefficients can be fitted to the measured Dalitz-plot
- ▶ $f(x, y)_{sym} = \frac{1}{\sqrt{2}} (f(x, y) + f(y, x))$
- ▶ Representations of S3 can also be extracted

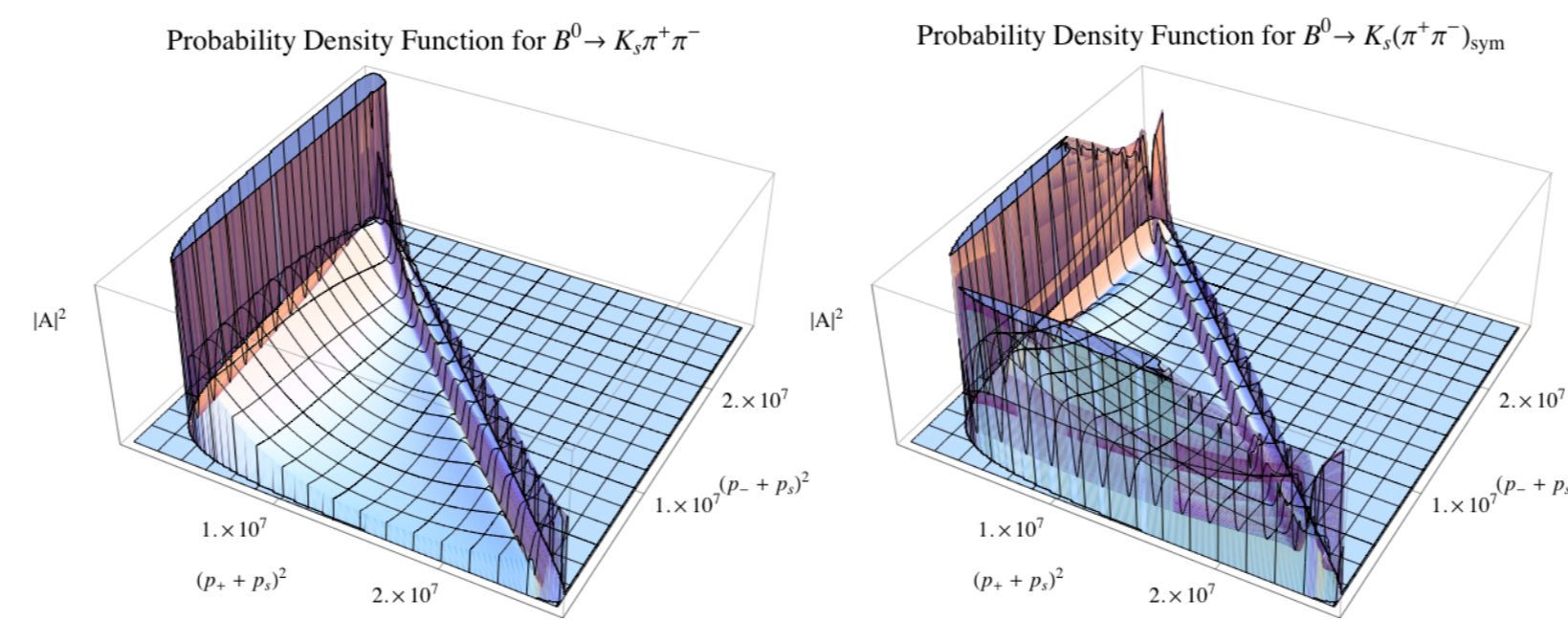


FIGURE: Probability Density Functions for $B^0_d \rightarrow K_S\pi^+\pi^-$ and for the symmetric part of the amplitude

Contractions and EWP-Tree Relations

- ▶ The effective weak hamiltonian :

$$H_{eff} = \frac{G_F}{\sqrt{2}} \sum_{q=d,s} \left(\sum_{p=u,c} \lambda_p^{(q)} (c_1(\mu) O_1^p(\mu) + c_2(\mu) O_2^p(\mu)) - \lambda_t^{(q)} \sum_{i=3}^{10} c_i(\mu) O_i(\mu) \right)$$
- ▶ Matrix elements of four-quark operators can be expressed in terms of contractions
- ▶ Contractions can be related to diagrams :

$$\begin{aligned} C'_1 &= \langle \bar{q}_1 q_2 \bar{q}_3 q_4 \bar{q}_5 q_6 | \bar{b} q_7 \bar{q}_8 q_9 | \bar{b} q_{10} \rangle, & T'_1 &= \langle \bar{q}_1 q_2 \bar{q}_3 q_4 \bar{q}_5 q_6 | \bar{b} q_7 \bar{q}_8 q_9 | \bar{b} q_{10} \rangle, \\ C'_2 &= \langle \bar{q}_1 q_2 \bar{q}_3 q_4 \bar{q}_5 q_6 | \bar{b} q_7 \bar{q}_8 q_9 | \bar{b} q_{10} \rangle, & T'_2 &= \langle \bar{q}_1 q_2 \bar{q}_3 q_4 \bar{q}_5 q_6 | \bar{b} q_7 \bar{q}_8 q_9 | \bar{b} q_{10} \rangle, \\ P'_{EW1} &= \langle \bar{q}_1 q_2 \bar{q}_3 q_4 \bar{q}_5 q_6 | \bar{b} q_7 \bar{q}_8 q_9 | \bar{b} q_{10} \rangle, & P'_{EW1} &= \langle \bar{q}_1 q_2 \bar{q}_3 q_4 \bar{q}_5 q_6 | \bar{b} q_7 \bar{q}_8 q_9 | \bar{b} q_{10} \rangle, \\ P'_{EW2} &= \langle \bar{q}_1 q_2 \bar{q}_3 q_4 \bar{q}_5 q_6 | \bar{b} q_7 \bar{q}_8 q_9 | \bar{b} q_{10} \rangle, & P'_{EW2} &= \langle \bar{q}_1 q_2 \bar{q}_3 q_4 \bar{q}_5 q_6 | \bar{b} q_7 \bar{q}_8 q_9 | \bar{b} q_{10} \rangle. \end{aligned}$$

- ▶ Under SU(3) symmetry, this allows us to obtain relations between the **fully symmetric** diagrams :

$$\begin{aligned} P'_{EW1} &= \frac{3|\lambda_t^{(s)}|}{2|\lambda_b^{(s)}|} c_9 + c_{10} T'_1, & P'_{EW2} &= \frac{3|\lambda_t^{(s)}|}{2|\lambda_b^{(s)}|} c_9 + c_{10} T'_2, \\ P'_{EW1} &= -\frac{3|\lambda_t^{(s)}|}{2|\lambda_b^{(s)}|} c_9 + c_{10} C'_1, & P'_{EW2} &= -\frac{3|\lambda_t^{(s)}|}{2|\lambda_b^{(s)}|} c_9 + c_{10} C'_2. \end{aligned}$$

- ▶ SU(3)-breaking errors are $O(30\%)$, but since trees and EWPs are sub-leading effects, the net error is expected to be $O(5\%)$

Extracting γ from $B \rightarrow K\pi\pi$

- ▶ The $B \rightarrow K\pi\pi$ sector contains **11** (independent, measurable) observables :
 - ▶ $\Gamma(B^+ \rightarrow K^+\pi^+\pi^-)$, $\Gamma(B^+ \rightarrow K^+\pi^0\pi^0)$,
 - ▶ $\Gamma(B^0 \rightarrow K^+\pi^0\pi^-)$, $\Gamma(B^0 \rightarrow K^0\pi^+\pi^-)$,
 - ▶ $\Gamma(B^0 \rightarrow K^0\pi^0\pi^0)$
- ▶ The five direct CP asymmetries
- ▶ The indirect CP asymmetry of $B^0 \rightarrow K^0\pi^+\pi^-$
- ▶ Using the fully symmetric amplitudes and the EWP-tree relations, the diagrammatic decomposition gives **10** theoretical parameters
- ▶ γ **can thus be extracted** by fitting the parameters to the observables

New Relations Between Amplitudes

- ▶ Neglecting annihilation and exchange diagrams reveals new relations between amplitudes
- ▶ These relations are invisible to the exact Wigner-Eckart analysis
- ▶ These are further tests of the SM which can be made once the symmetric and antisymmetric scenarios are experimentally distinguished

References

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