

Dimensionally reduced EFTs for cosmological phase transitions

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Based on [2406.02667], by:

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Thermally induced field phase transitions



Fig. 1: Temperature evolution of scalar potential



Fig. 2: O(3) symmetric bounce solution



[Caprini et al. - 1512.06239]

As the transition rate grows, in a first-order PT (FOPT):

- 1) Bubbles of true vacuum nucleate and expand in a hot plasma
- 2) Bubble fronts collide
- 3) Sound waves
- 4) Turbulence



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Out of equilibrium



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Probe

Answers for electroweak baryogenesis

Out of equilibrium



From PTs to GWs

How do we connect a QFT model to these GW spectra?



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Inverse duration of PT

$$\frac{\beta}{H_*} = T_* \frac{\mathrm{d}S_3[\varphi_c]}{\mathrm{d}T} \bigg|_{T_*}$$

Strength parameter

$$\alpha \approx -0.03 \frac{V_3(\varphi_T)}{T_*^3}$$

Terminal bubble wall velocity

A whole different story...

[Laurent, Cline - 2204.13120]



What is the state of the art?



- Several studies of FOPTs in BSM extensions and in the SMEFT:
 - N⁴LO Higgs potential at finite T [Ekstedt et al. 2
 - EWPT in dim-6 SMEFT
 - EWPT in dim-6 XSM
 - $\circ \quad \text{EWPT in } \Sigma \text{SM}$

- [Ekstedt *et al.* 2405.18349] [Camargo-Molina *et al.* - 2103.14022] [Oikonomou *et al.* - 2403.01591] [Niemi *et al.* - 1802.10500]
- Tools for PT-related computations:
 - CosmoTransitions
 - FindBounce
 - DRalgo

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• BubbleDet

[Wainwright - 1109.4189] [Guada *et al.* - 2002.00881] [Ekstedt *et al.* - 2205.08815] [Ekstedt *et al.* - 2308.15652] Effective operators with derivatives **not included**

- Theoretical uncertainties persist from the particle physics side:
 - LISA [Caprini et al. 1910.13125]



See earlier talk by M. Chala

UV theory in 4D Minkowski:

At temperature T, in Euclidean space:

$$\begin{aligned} \mathscr{L}_{4} &= \frac{1}{2} (\partial_{\mu} \phi)^{2} - \frac{1}{2} m^{2} \phi^{2} - \kappa \phi^{3} - \lambda \phi^{4} + \overline{\Psi} i \partial \overline{\Psi} - g \phi \overline{\Psi} \Psi \quad \text{[Gould, Xie - 2310.02308]} \\ S_{0} &= (-i) \frac{1}{T} \int d^{3} \mathbf{x} \sum_{n=-\infty}^{\infty} \left[\frac{1}{2} (\partial_{i} \varphi_{n})^{2} + \frac{1}{2} \left(m^{2} + (2\pi nT)^{2} \right) \varphi_{n}^{2} \right. \\ &+ \overline{\psi}_{n} \partial \psi_{n} + \left(2\pi \left(n + \frac{1}{2} \right) T \right)^{2} \overline{\psi}_{n} \psi_{n} \right] \end{aligned}$$





[Kajantie et al. - 9508379] [Criado - 1901.03501]

We find an off-shell basis of operators up to (4D) dim-8 with the help of *BasisGen* and match at 1-loop (fermion only loops)¹:

$$\begin{aligned} \mathscr{L}_{3} &= \frac{1}{2} K_{3} (\partial \varphi)^{2} + \frac{1}{2} m_{3}^{2} \varphi^{2} + \kappa_{3} \varphi^{3} + \lambda_{3} \varphi^{4} \\ &+ \alpha_{61} \varphi^{6} + \beta_{61} \partial^{2} \varphi \partial^{2} \varphi + \beta_{62} \varphi^{3} \partial^{2} \varphi \\ &+ \alpha_{81} \varphi^{8} + \alpha_{82} \varphi^{2} \partial_{\mu} \partial_{\nu} \varphi \partial^{\mu} \partial^{\nu} \varphi + \beta_{81} \varphi \partial^{6} \varphi + \beta_{82} \varphi^{3} \partial^{4} \varphi + \beta_{83} \varphi^{2} \partial^{2} \varphi \partial^{2} \varphi + \beta_{84} \varphi^{5} \partial^{2} \varphi \end{aligned}$$

¹ Scalar loop contributions to effective operators are subleading (see Appendix A in [2406.02667]).



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$$+ \alpha_{61}\varphi^{6} + \beta_{61}\partial^{2}\varphi\partial^{2}\varphi + \beta_{62}\varphi^{3}\partial^{2}\varphi$$

$$+ \alpha_{81}\varphi^{8} + \alpha_{82}\varphi^{2}\partial_{\mu}\partial_{\nu}\varphi\partial^{\mu}\partial^{\nu}\varphi + \beta_{81}\varphi\partial^{6}\varphi + \beta_{82}\varphi^{3}\partial^{4}\varphi + \beta_{83}\varphi^{2}\partial^{2}\varphi\partial^{2}\varphi + \beta_{84}\varphi^{5}\partial^{2}\varphi$$

$$dim-6$$

$$dim-8$$

$$Control EFT validity$$

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3D EFT approach

Computing relevant PT magnitudes





Fig. 5: PT magnitudes in two models, with and w/o effective operators



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Two observations

Eff. ops. **allow for PTs in a wider range** of values of the Yukawa Including eff. ops. yields **very different estimations** at large Yukawas



Fig. 6: GW power spectra in two models, with and w/o effective operators

(*) Generated with **PTPlot** [Caprini et al. - 1910.13125]



Fig. 6: GW power spectra in two models, with and w/o effective operators

Including eff. ops. can change the peak amplitude and frequency by **one order of magnitude**



We have learned

- Strong FOPTs occur in regions of parameter space **close the limit of validity** of the 3D EFT
- Higher dimensional effective operators are thus relevant in the estimation of PT-related magnitudes (order of magnitude differences in GW spectra)
- **Bounce solutions** can be obtained perturbatively in the presence of higher derivative terms

Future work

- Application to well-motivated BSM models
- Adding loop corrections to scalar zero mode

Thank you for your attention!

¡Gracias por vuestra atención!