

Triple-gauge couplings in LHC diboson production: a SMEFT view from every angle

In collaboration with Giovanni Pelliccioli and Eleni Vryonidou, [[arXiv:2405.19083](https://arxiv.org/abs/2405.19083)]

Hesham El Faham
The University of Manchester



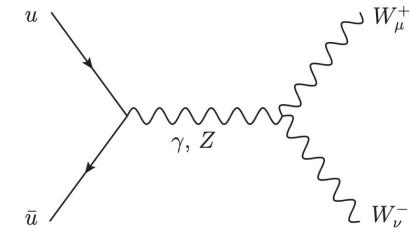
Motivation for diboson at the LHC

- Important probe for EWSB
- Fully leptonic diboson production → **relatively clean signature at the LHC**
- With Run 3 and HL-LHC → **promising for precision and differential measurements**
- Irreducible background for Higgs analyses

At LO, production is dominated by quark-initial states and gluon-initiated ones are loop-induced
→ at **NLO in QCD, mixed channel opens up with enhancement from gluon luminosity**

On diboson in the SMEFT

- Dominating quark-initiated channel is sensitive to dim-6 TGC
- At NLO QCD, **sensitivity to TGC is non-trivial and depends on phase-space setups**
- Dim-6 TGCs **non-trivially correlate with Vqq-induced ones** [Grojean et al. \[1810.05149\]](#)
- Linear **suppression is expected for $2 \rightarrow 2$ due to helicity selection rules** [Azatov et al. \[1607.05236\]](#)
- A priori, one can not neglect dim-8 SMEFT insertions e.g. [Degrande et al. \[2303.10493\]](#)
→ **dim-8 effects are not expected to alter the power induced by *purely* dim-6 TGC quadratic contributions** [Corbett et al. \[2304.03305\]](#)

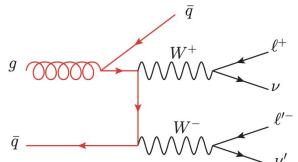
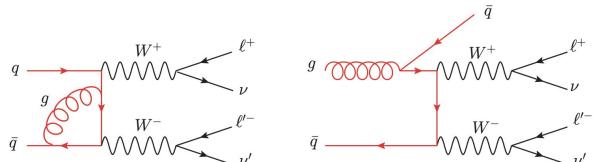


Goal

- Purely CP-even and CP-odd SMEFT coefficients in the Warsaw basis [Grzadkowski et al. \[1008.4884\]](#)

$$\epsilon_{ijk} W_{\mu\nu}^i W^{j,\nu\rho} W_\rho^{k,\mu}, \quad \epsilon_{ijk} \tilde{W}_{\mu\nu}^i W^{j,\nu\rho} W_\rho^{k,\mu} \longleftrightarrow \lambda_z = -c_W \frac{v}{\Lambda^2} \frac{3}{2} g, \quad \tilde{\lambda}_z = -c_{\tilde{W}} \frac{v}{\Lambda^2} \frac{3}{2} g$$

- Full NLO in QCD, including the complete off-shell effects and spin correlations



→ Impact of different phase-space setups?

→ EFT effects on angular coefficients and observables?

→ Impact of NLO QCD?

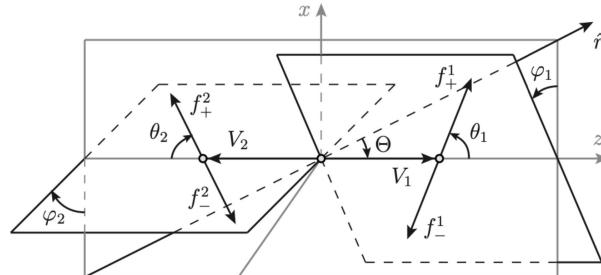
Diboson analysis features

- Z couples \sim equally to left and right-hand fermions \rightarrow **can not identify helicities of final states**
- W couples to left-hand fermions \rightarrow **but neutrino reconstruction is problematic**

Interference suppression at $2 \rightarrow 2$ is lifted at $2 \rightarrow 3$ or $2 \rightarrow 4$

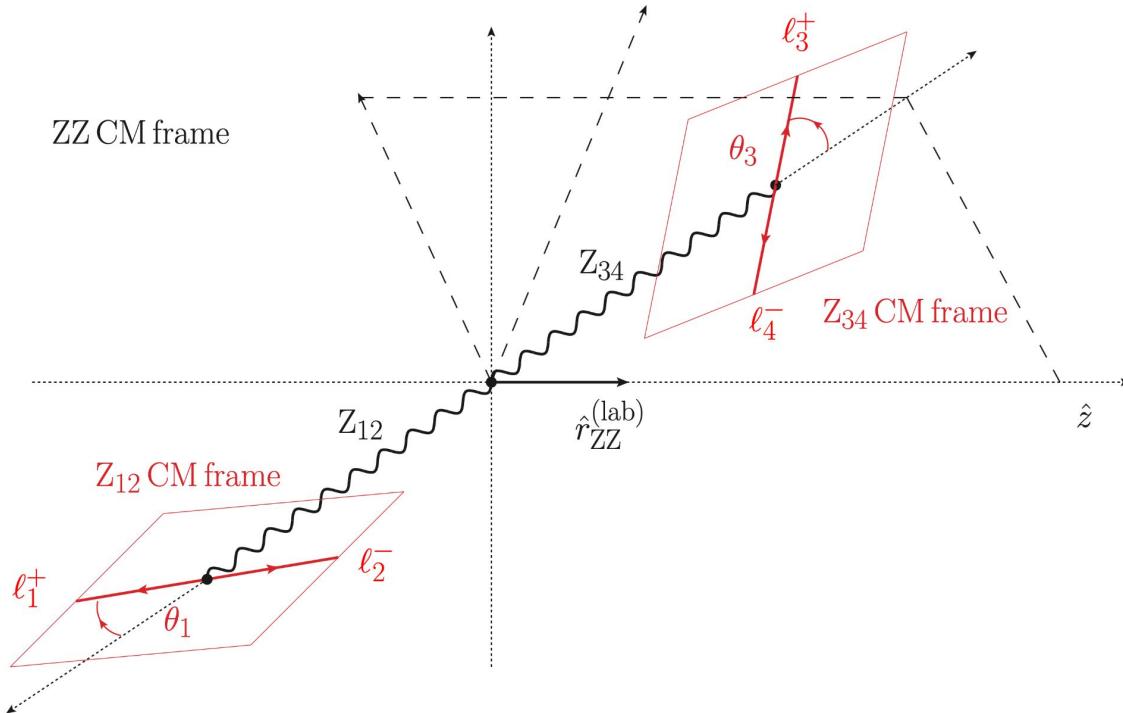
See also talks by Matteo and Marion

\rightarrow **the angle spanned by the decay products and/or real radiation ‘restores’ the interference**



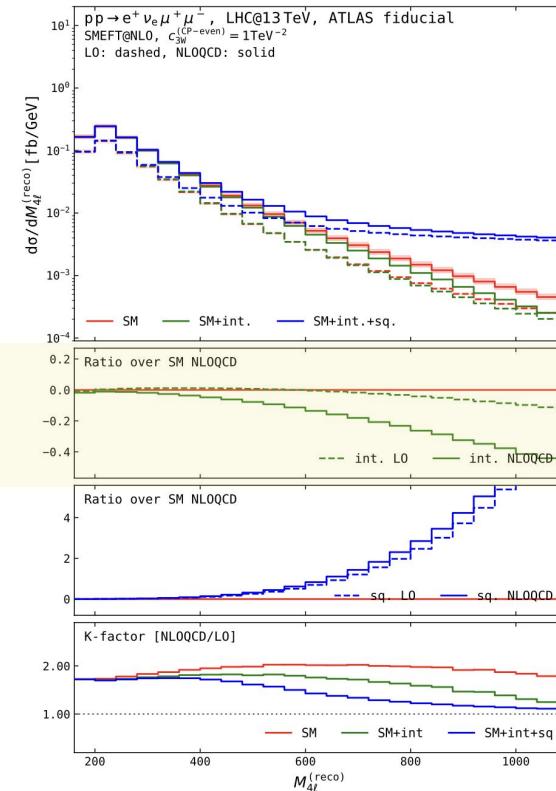
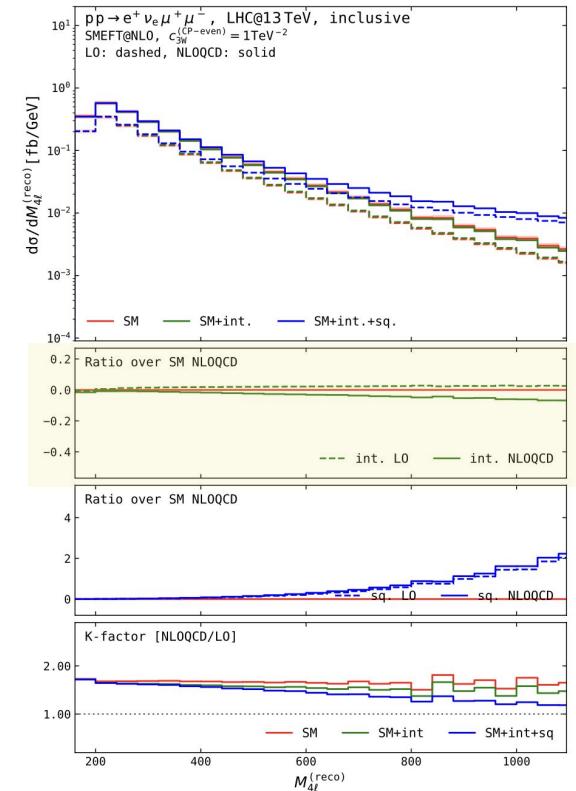
$$\frac{d\sigma_{\text{int}}(q\bar{q} \rightarrow WZ \rightarrow 4\psi)}{d\phi_Z d\phi_W} \propto \cos(2\phi_Z) + \cos(2\phi_W)$$

Helicity coordinate system



Impact of NLO QCD and selection cuts

interference ‘restored’ through selection cuts



Inclusive (left)

→ Real NLO radiation restores the suppressed LO SMEFT interference

Fiducial (right)

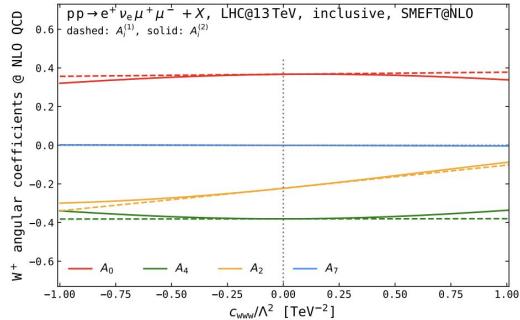
→ The interference restoration is already manifest at LO due to the modulation from the cuts

→ Non-trivial K-factors

Inclusive angular coefficients

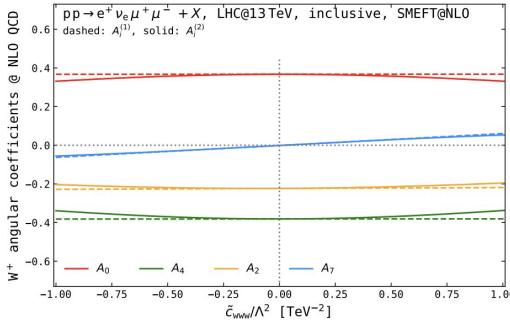
Inclusive setup

CP-even



(a)

CP-odd



(b)

At the linear-level (dashed),

→ polarisation fractions, **A0** and **A4**, are barely distorted by CP-even and unaffected by the CP-odd modifications

$$A_0 = 2 f_0, \quad A_4 = 2 c_{LR} (f_R - f_L)$$

At the quadratic-level (solid),

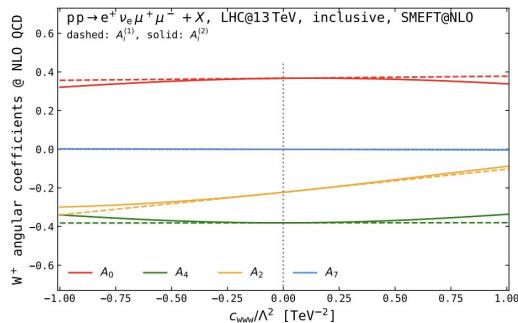
→ **right handed** and **longitudinal** fractions of the W are modified for CP-even and CP-odd

.. negligible effect on the left handed one

Inclusive angular coefficients

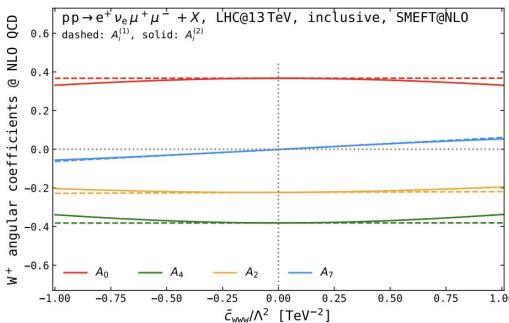
Inclusive setup

CP-even



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CP-odd



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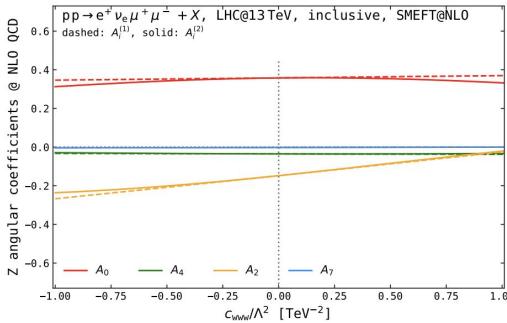
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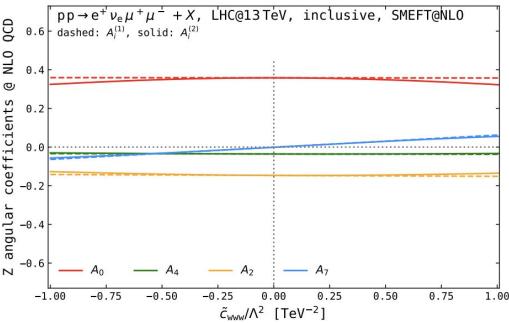
At the quadratic-level (solid),

→ right handed and longitudinal fractions of the W are modified for CP-even and CP-odd

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(c)



(d)

At the linear- and quadratic-levels,

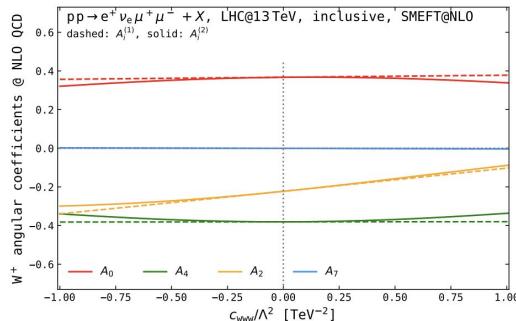
→ Longitudinal fraction of the Z behaves similarly to W

→ The very small absolute value of **A4** for the Z manifest the left-right balance which is not altered by the EFT

Inclusive angular coefficients

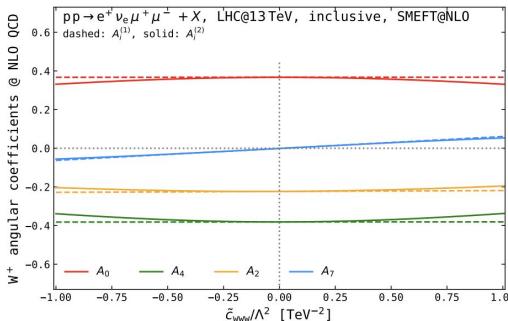
Inclusive setup

CP-even



(a)

CP-odd



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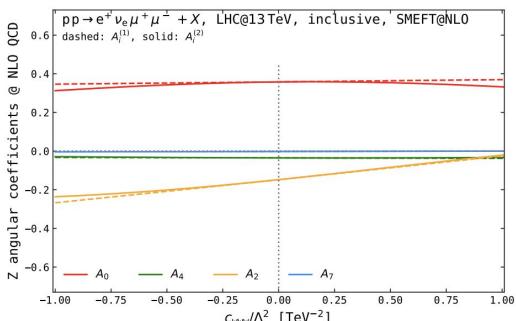
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At the quadratic-level (solid),

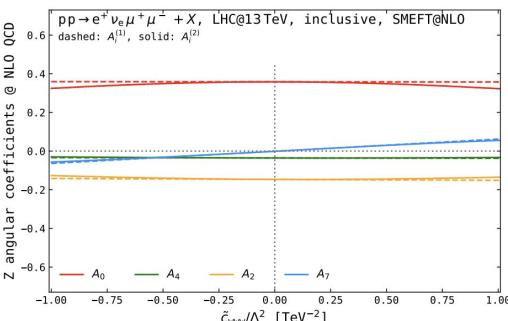
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.. negligible effect on the left handed one

Z



(c)



(d)

At the linear- and quadratic-levels,

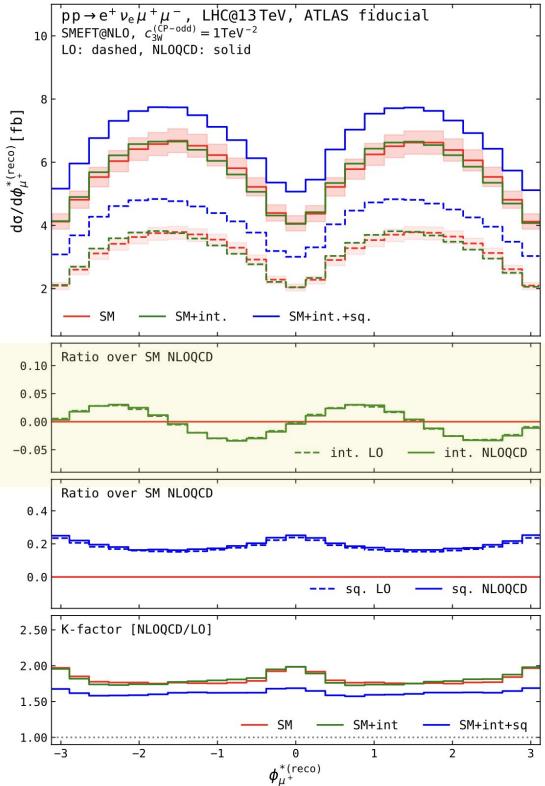
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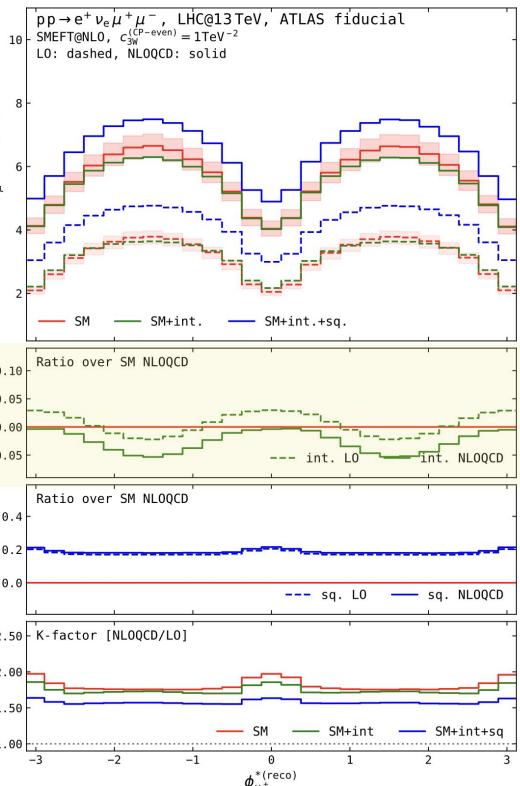
A7 is parity odd sensitive

Differential angular observables

CP-odd



CP-even



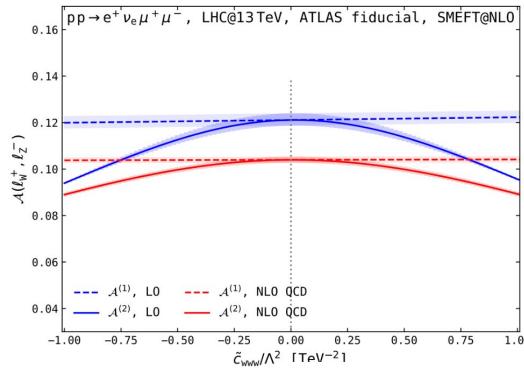
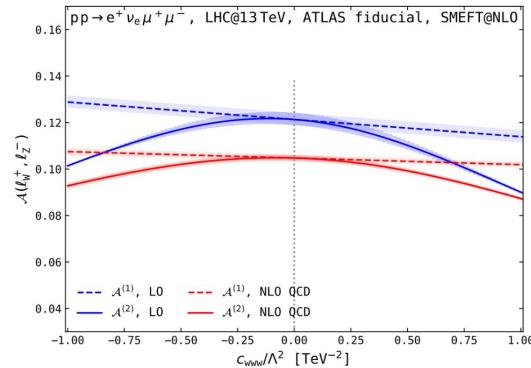
Azimuthal variables are good probes for CP-properties

→ Interference modulation maps the CP-property of TGC

→ Distortion due to selection cuts and neutrino reconstruction relative to SM is mild (inclusive setup not shown here)

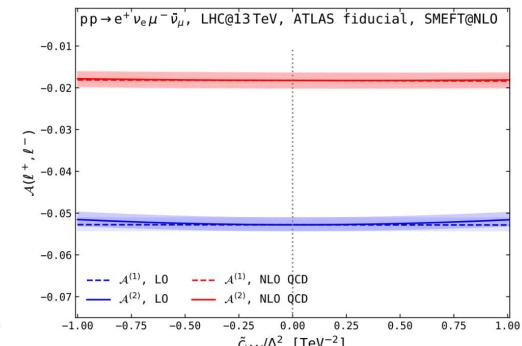
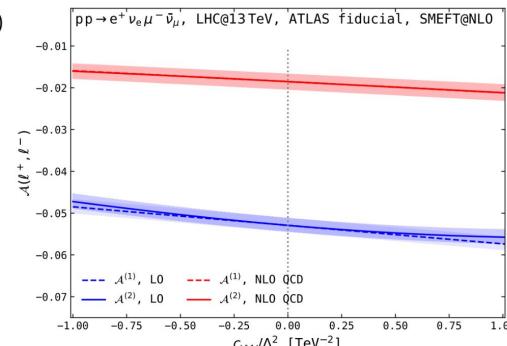
Boost asymmetries

$$\mathcal{A}(i, j) = \frac{d\sigma(|y_i| > |y_j|) - d\sigma(|y_i| < |y_j|)}{d\sigma(|y_i| > |y_j|) + d\sigma(|y_i| < |y_j|)}$$



WW

→ Mild EFT effects on asymmetries



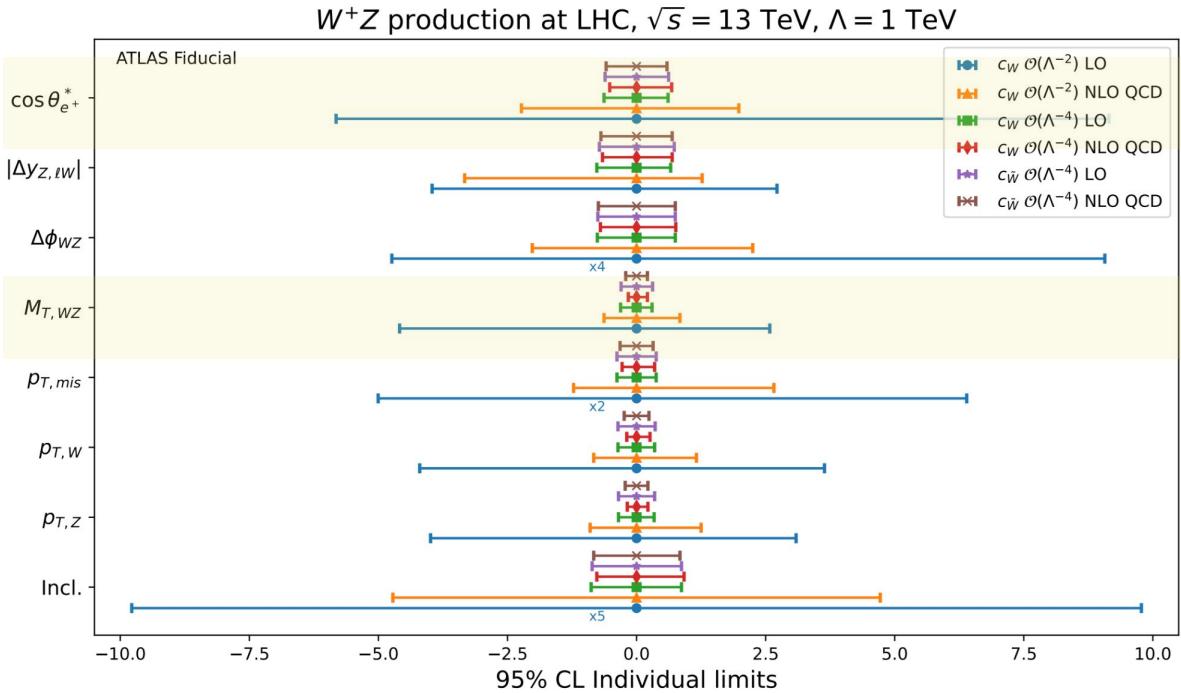
Differential measurements of boost asymmetries might be promising

WZ

→ CP-even (left) modifies the right-handed fraction at linear and quadratic-levels

→ CP-odd (right) linear effects are negligible

Impact of NLO QCD on SMEFT WZ



- dimensionful observables are the most stringent
- polar decay angle is competitive in the dimensionless class
- **NLO QCD** is key in constraining SMEFT

Conclusions

- NLO QCD is key in diboson production; constraining SMEFT, resurrecting $2 \rightarrow 2$ suppressed interference, non trivial k-factors
- The angle spanned by decay products as well as selection cuts have significant impact on the interference behavior
- Angular observables are good probes for TGC CP-properties
- Differential leptonic boost asymmetries might be promising in constraining SMEFT

Backup

On polarisation fractions and angular terms

2-body decay rate of V boson + projections on spherical harmonics up to rank 2

→ inclusive angular coefficients and polarisation fractions

$$\frac{1}{\sigma} \frac{d\sigma}{d \cos \theta^* d\phi^*} = \frac{3}{16\pi} \left[1 + \cos^2 \theta^* + A_0 \frac{1 - 3 \cos^2 \theta^*}{2} + A_1 \sin 2\theta^* \cos \phi^* \right.$$

azimuthal integral

$$\begin{aligned} &+ \frac{1}{2} A_2 \sin^2 \theta^* \cos 2\phi^* + A_3 \sin \theta^* \cos \phi^* + A_4 \cos \theta^* \\ &\left. + A_5 \sin \theta^* \sin \phi^* + A_6 \sin 2\theta^* \sin \phi^* + A_7 \sin^2 \theta^* \sin 2\phi^* \right]$$

$$\frac{1}{\sigma} \frac{d\sigma}{d \cos \theta^*} = \frac{3}{8} \left[2 f_0 \sin^2 \theta^* \right.$$

→

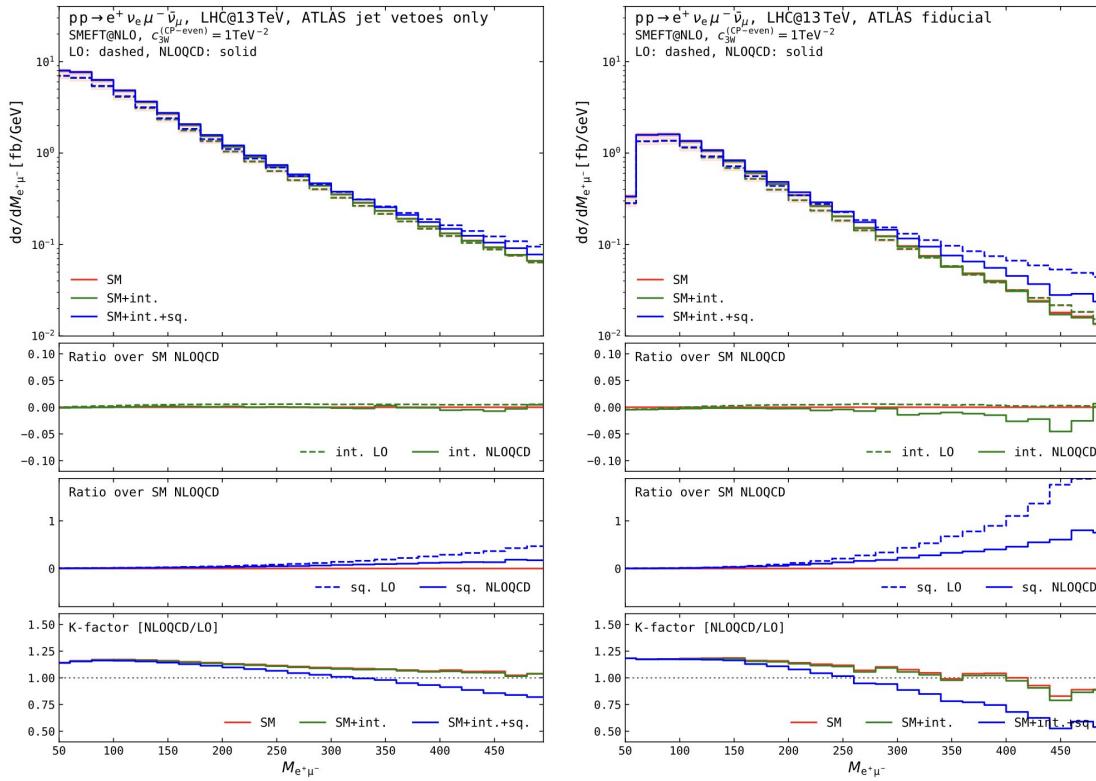
$$\begin{aligned} &+ f_L (1 + \cos^2 \theta^* - 2 c_{LR} \cos \theta^*) \\ &\left. + f_R (1 + \cos^2 \theta^* + 2 c_{LR} \cos \theta^*) \right] \end{aligned}$$

A_i coefficients modulate an angular term

→ underly the dynamics of the production and decay process, the polarisation states of the particles, and possible interference effects

Impact of NLO QCD and selection cuts

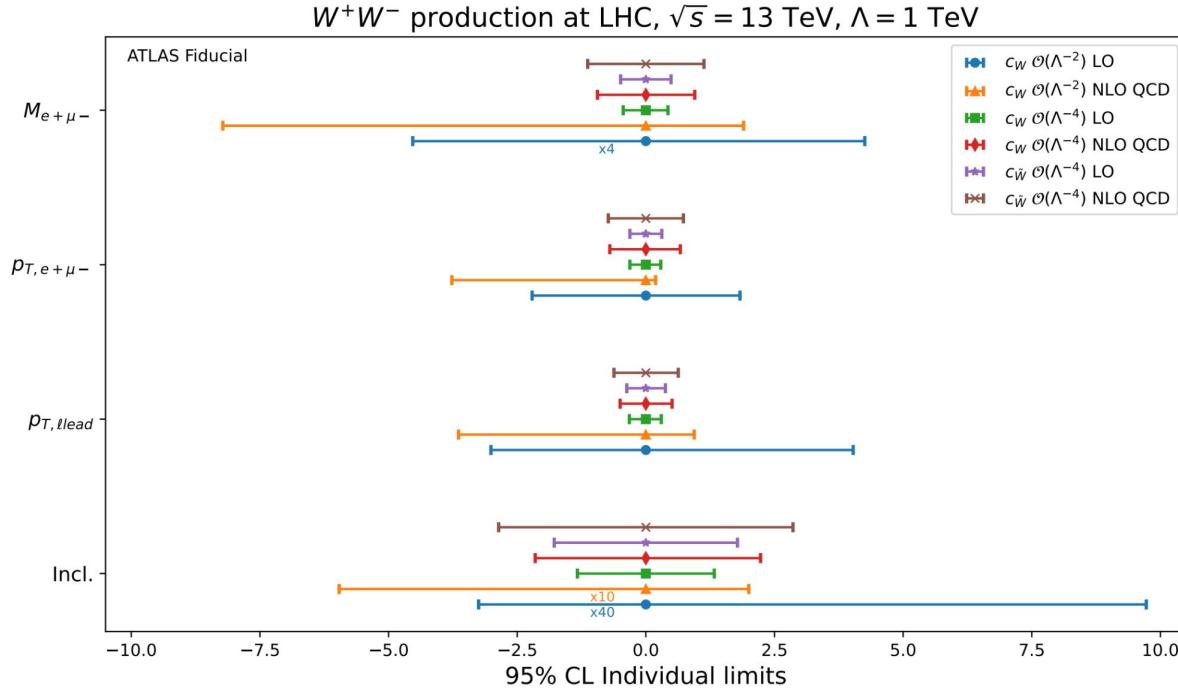
WW



→ Selection cuts still enhances
the interferences

→ WW is less-sensitive to TGC
than WZ

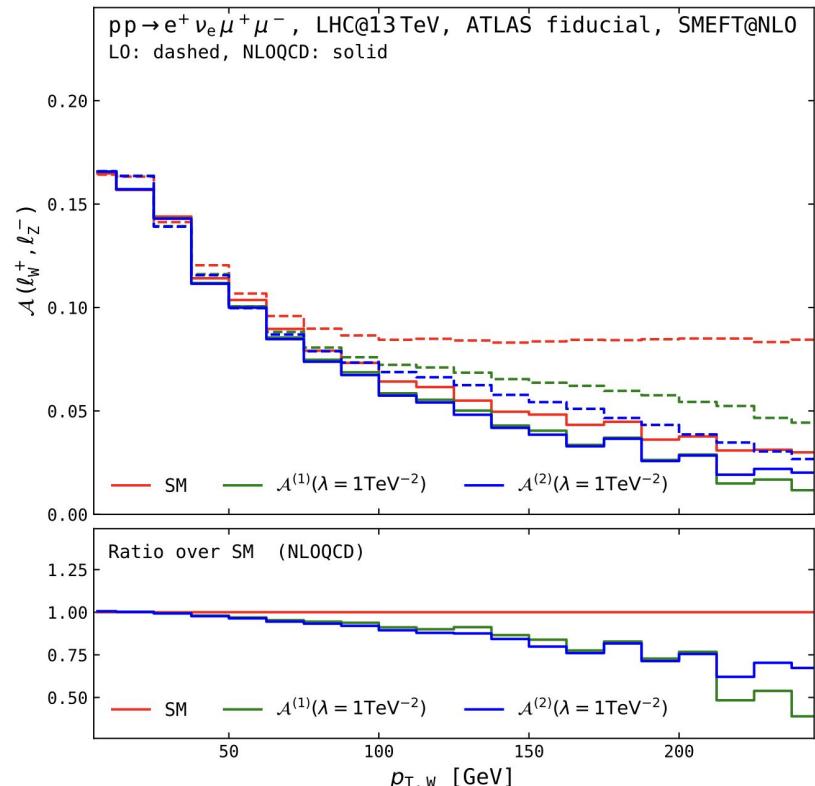
Impact of NLO QCD on SMEFT \rightarrow WW



→ similar conclusions to the WZ case
 → the different NLO QCD behaviour
 is manifest

Boost asymmetries

CP-even



CP-odd

