

CP violation in gluon-induced diboson production

Marion Thomas

University of Manchester

In collaboration with Eleni Vryonidou, in preparation

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CP violation in the SMEFT

CP-odd component of Higgs interactions not ruled out by measurements



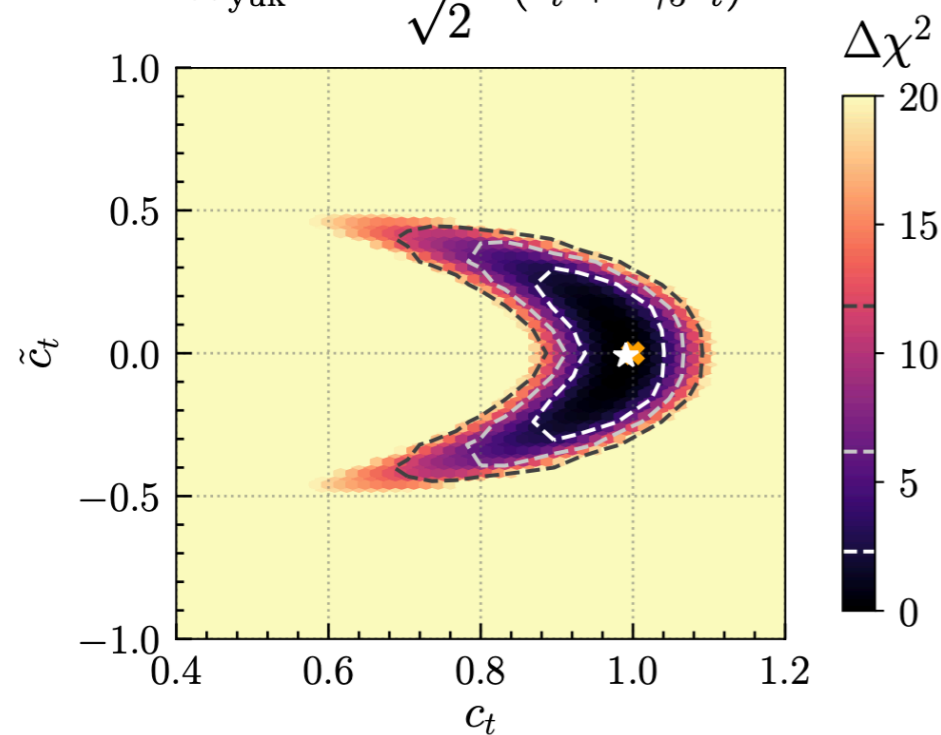
CP-mixing angle in top Yukawa in $t\bar{t}H$ and tH production [\[arXiv:2303.05974\]](#)



CP-odd contribution to $t\bar{t}H$ in $H \rightarrow \gamma\gamma$ [\[arXiv:2003.10866\]](#)

+ many more

$$\mathcal{L}_{\text{yuk}} = -\frac{y_t^{\text{SM}}}{\sqrt{2}} \bar{t} (c_t + i\gamma_5 \tilde{c}_t) t H$$



Bahl et al. in arXiv:2007.08542

EFT studies of Higgs interactions at one-loop, eg. with Higgs Characterisation framework [\[arXiv:1306.6464\]](#)
[\[arXiv:1407.5089\]](#)

→ So far focus on CP-violation in HVV and Hff interactions

What about a more general treatment?

→ We extend the SMEFT@NLO UFO to include dim-6 CP-odd SMEFT operators entering gluon-induced diboson production and study their impact on kinematic distributions.

CP-violating SMEFT operators

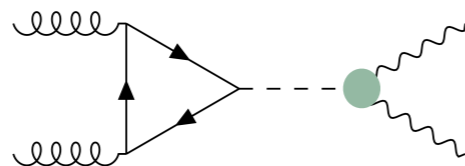
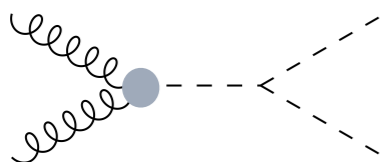
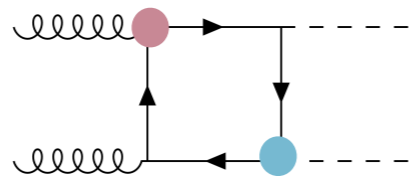
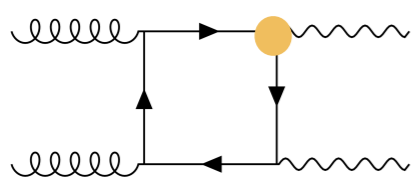
Warsaw basis of dim-6 SMEFT operators.
 Flavour symmetry: $U(2)_q \times U(3)_d \times U(2)_u$

Hermitian operators

$$c_i = \text{RE}c_i$$

$$\mathcal{O}_{\varphi\tilde{G}} \quad c_{\varphi\tilde{G}} \quad \left(\varphi^\dagger\varphi - \frac{v^2}{2}\right) \tilde{G}_A^{\mu\nu} G_{\mu\nu}^A$$

$$\mathcal{O}_{\varphi\tilde{W}} \quad c_{\varphi\tilde{W}} \quad \left(\varphi^\dagger\varphi - \frac{v^2}{2}\right) \tilde{W}_I^{\mu\nu} W_{\mu\nu}^I$$



Non-hermitian operators

$$c_i = \text{RE}c_i + i \text{IM}c_i$$

$$\mathcal{O}_{tG} \quad c_{tG} \quad ig_s (\bar{Q}\sigma^{\mu\nu} T_A t) \tilde{\varphi} G_{\mu\nu}^A + \text{h.c.}$$

$$\mathcal{O}_{t\varphi} \quad c_{t\varphi} \quad \left(\varphi^\dagger\varphi - \frac{v^2}{2}\right) \bar{Q} t \tilde{\varphi} + \text{h.c.}$$

$$\mathcal{O}_{tW} \quad c_{tW} \quad (\bar{Q}\sigma^{\mu\nu} \tau_I t) \tilde{\varphi} W_{\mu\nu}^I + \text{h.c.}$$

$$\mathcal{O}_{tB} \quad c_{tB} \quad (\bar{Q}\sigma^{\mu\nu} t) \tilde{\varphi} B_{\mu\nu} + \text{h.c.}$$

$$\mathcal{O}_{tZ} \quad c_{tZ} \quad -\sin\theta_W c_{tB} + \cos\theta_W c_{tW}$$

$$\tilde{X}_{\mu\nu} = \frac{1}{2}\epsilon_{\mu\nu\rho\sigma} X^{\rho\sigma}$$

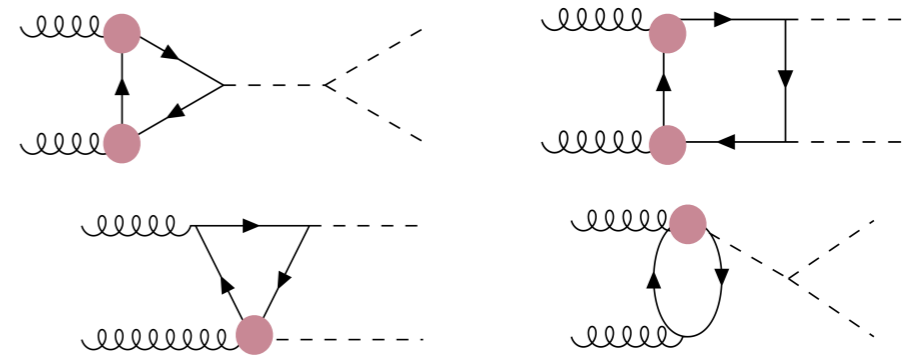
Adding CP-odd operators in SMEFT@NLO

3 ingredients to
add to UFO

Degrande, Durieux, Maltoni, Mimasu,
Vryonidou, Zhang in arXiv:2008.11743

Example: $\text{Im}c_{tG}$ in $gg \rightarrow HH$

$$\mathcal{O}_{tG} \quad c_{tG} \quad ig_s (\bar{Q}\sigma^{\mu\nu} T_A t) \tilde{\varphi} G_{\mu\nu}^A + \text{h.c.}$$



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1. Feynman rules

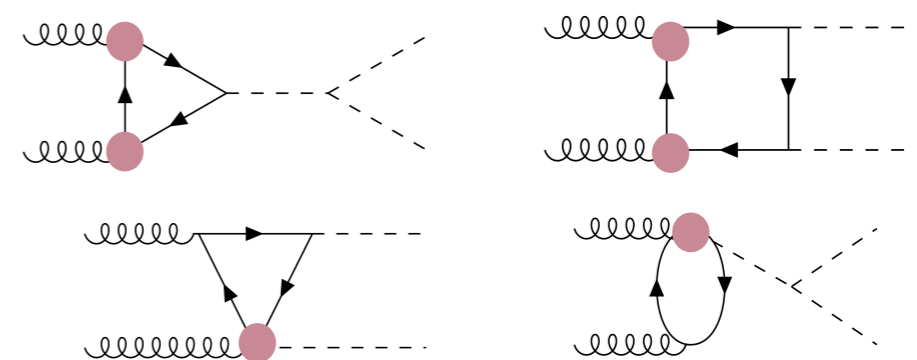
Read from the SMEFT Lagrangian

Checked against SMEFTsim implementation

Brivio in arXiv:2012.11343

Example: $\text{Im}c_{tG}$ in $gg \rightarrow HH$

$$\mathcal{O}_{tG} \quad c_{tG} \quad ig_s (\bar{Q} \sigma^{\mu\nu} T_A t) \tilde{\varphi} G_{\mu\nu}^A + \text{h.c.}$$



$$\begin{array}{c} t \\ \nearrow \\ \bullet \\ \nwarrow \\ \bar{t} \end{array}
 \begin{array}{c} \text{gluon } g, p_1, \mu \end{array}
 = -\frac{\text{Im}c_{tG} v g_s}{\sqrt{2} \Lambda^2} [\gamma^\mu, \not{p}_1] \gamma_5$$

$$\begin{array}{c} t \\ \nearrow \\ \bullet \\ \nwarrow \\ \bar{t} \end{array}
 \begin{array}{c} \text{gluon } g, p_1, \mu \\ \text{Higgs } H \end{array}
 = -\frac{\text{Im}c_{tG} g_s}{\sqrt{2} \Lambda^2} [\gamma^\mu, \not{p}_1] \gamma_5$$

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Example: $\text{IM}c_{tG}$ in $gg \rightarrow HH$

1. Feynman rules

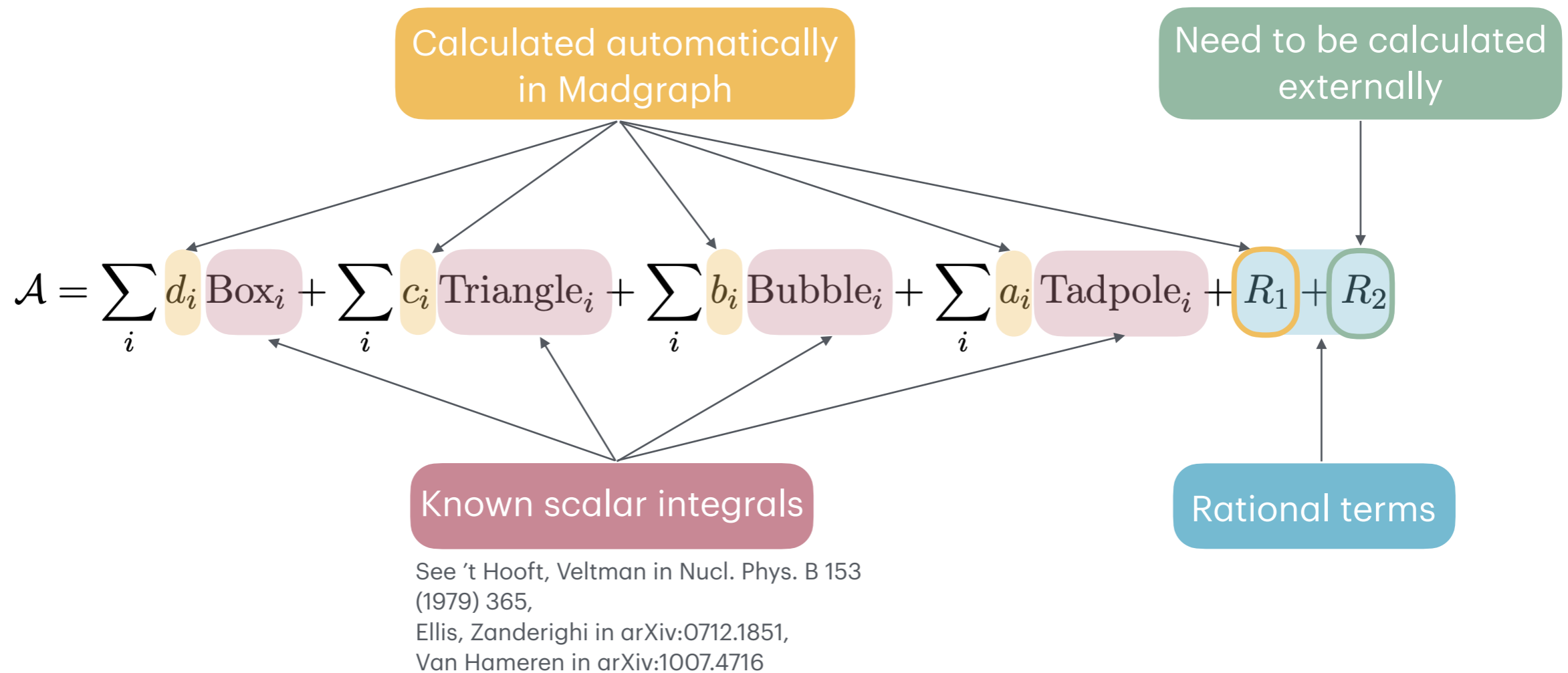
Read from the SMEFT Lagrangian

2. Rational Terms R_2

Calculated from one-loop irreducible diagrams

What are the rational terms?

Implementation of one-loop QCD calculations in MadLoop relies on Ossola-Papadopoulos-Pittau (OPP) reduction method.



Ossola, Papadopoulos, Pittau in arXiv:0609007, 0711.3596, 0802.1876
Hirschi et al. in arXiv:1103.0621

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Example: $\text{IM}c_{tG}$ in $gg \rightarrow HH$

Need underlying $ggH R_2$

$$p_1, a, \alpha \quad p_2, b, \beta \quad H = \frac{\text{IM}c_{tG}}{\Lambda^2} \frac{\sqrt{2} i g_s^2 m_t}{3\pi^2} \epsilon^{\alpha\beta\gamma\delta} p_1^\gamma p_2^\delta \delta_{ab}$$

One-loop irreducible diagrams for $gg \rightarrow HH$

$$p_1, a, \alpha \quad p_2, b, \beta \quad H \quad H = \frac{\text{IM}c_{tG}}{\Lambda^2} \frac{i g_s^2 m_t}{\sqrt{2}\pi^2 v} \epsilon^{\alpha\beta\gamma\delta} p_1^\gamma p_2^\delta \delta_{ab}$$

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1. Feynman rules

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2. Rational Terms R_2

Calculated from one-loop irreducible diagrams

3. UV Counterterms

Checked against the Renormalisation
Group Evolution (RGE)

Alonso, Jenkins, Manohar and Trott in arXiv:1308.2627,
1310.4838, 1312.2014

Example: $\text{Im}c_{tG}$ in $gg \rightarrow HH$

Renormalised with $\mathcal{O}_{\phi\tilde{G}}$



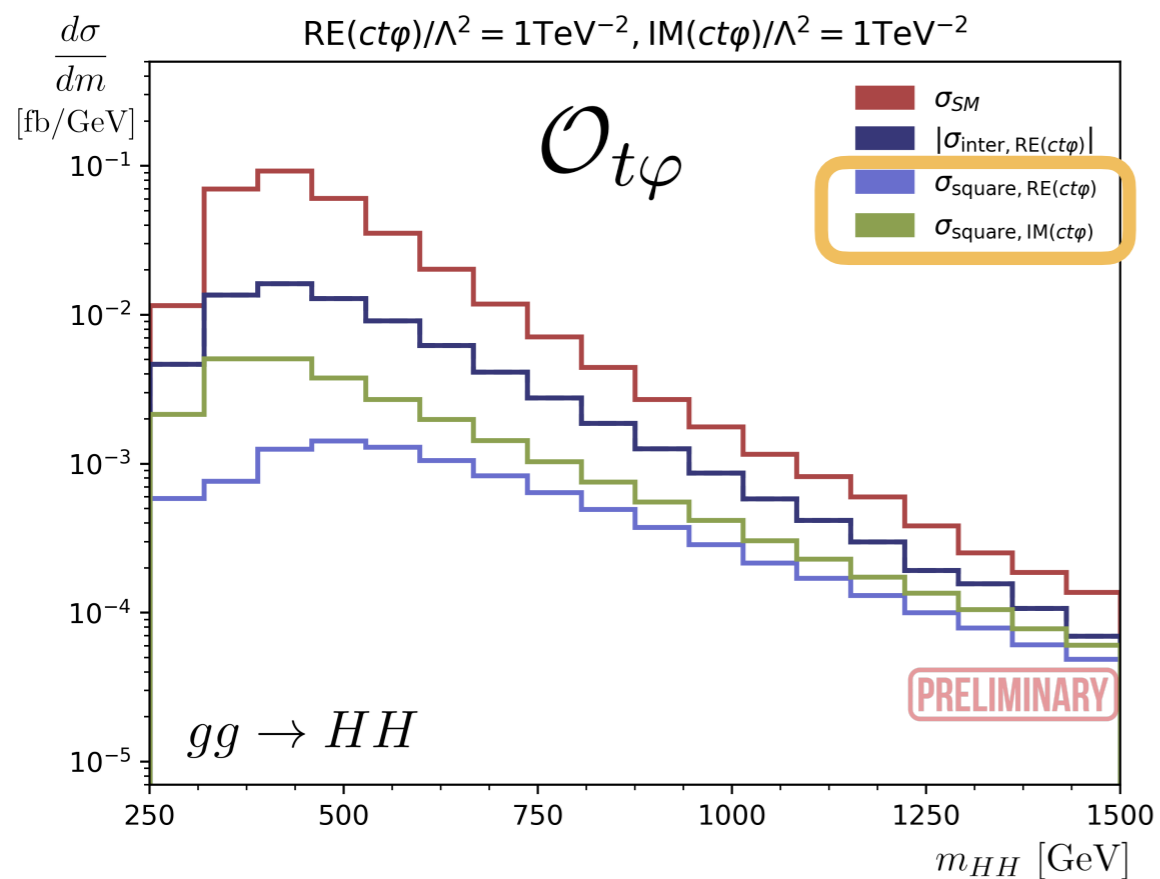
ggH and $ggHH$ UV counterterms:

$$\begin{array}{c} p_1, a, \alpha \\ \text{wavy line} \\ \bullet \\ \text{wavy line} \\ p_2, b, \beta \end{array} \text{---} H = \frac{\text{Im}c_{tG}}{\Lambda^2} \frac{i g_s^2 m_t}{\sqrt{2}\pi^2} \left(\frac{1}{\epsilon} - \log \left(\frac{\mu_{EFT}^2}{\mu_R^2} \right) \right) \epsilon^{p_1 p_2 \alpha \beta} \delta_{ab}$$

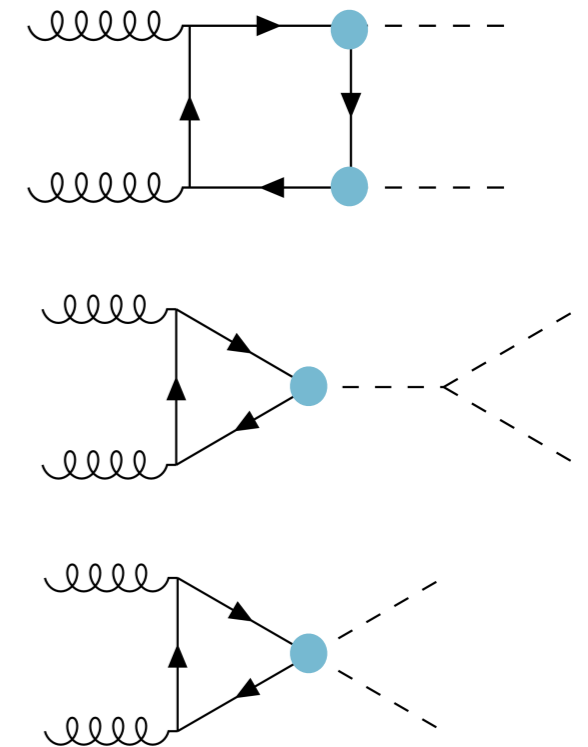
$$\begin{array}{c} p_1, a, \alpha \\ \text{wavy line} \\ \bullet \\ \text{wavy line} \\ p_2, b, \beta \end{array} \begin{array}{c} H \\ \text{dashed line} \\ \bullet \\ \text{dashed line} \\ H \end{array} = \frac{\text{Im}c_{tG}}{\Lambda^2} \frac{i g_s^2 m_t}{\sqrt{2}\pi^2 v} \left(\frac{1}{\epsilon} - \log \left(\frac{\mu_{EFT}^2}{\mu_R^2} \right) \right) \epsilon^{p_1 p_2 \alpha \beta} \delta_{ab}$$

Growing amplitudes and tail effects

Modified top-Higgs interactions in $gg \rightarrow HH$

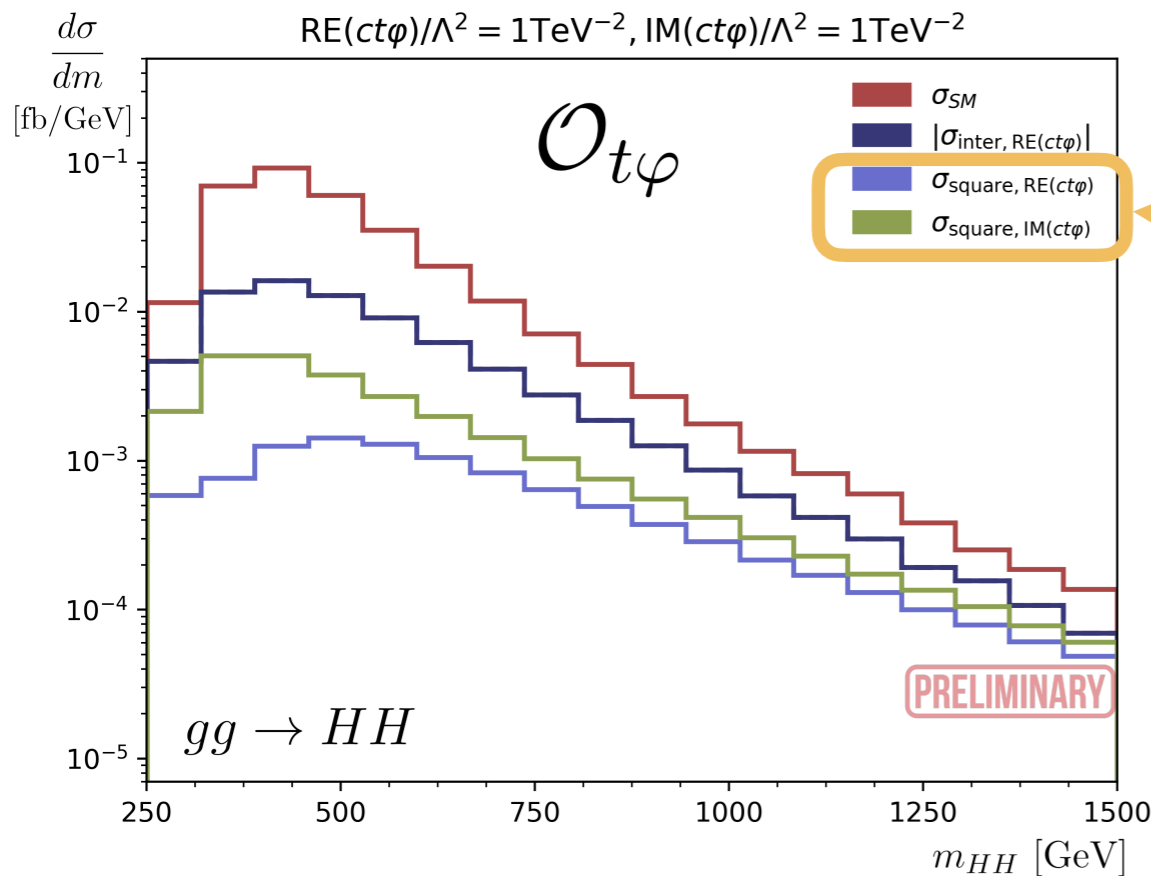


Converge at high energies

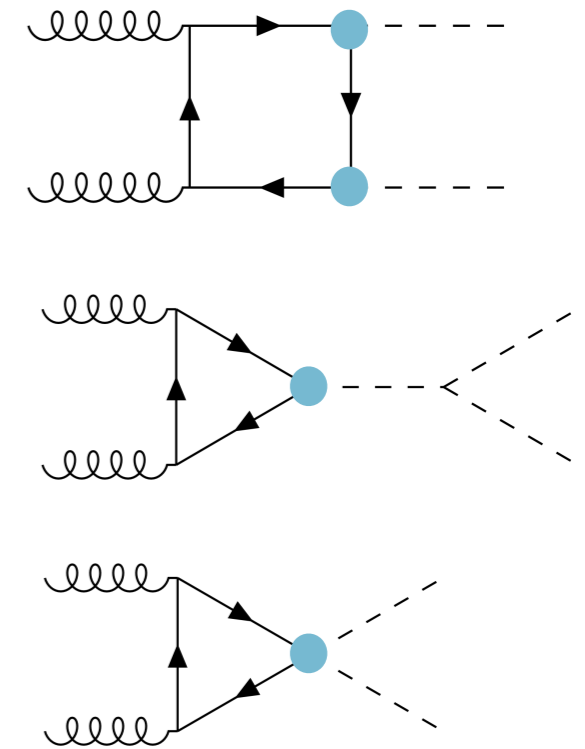


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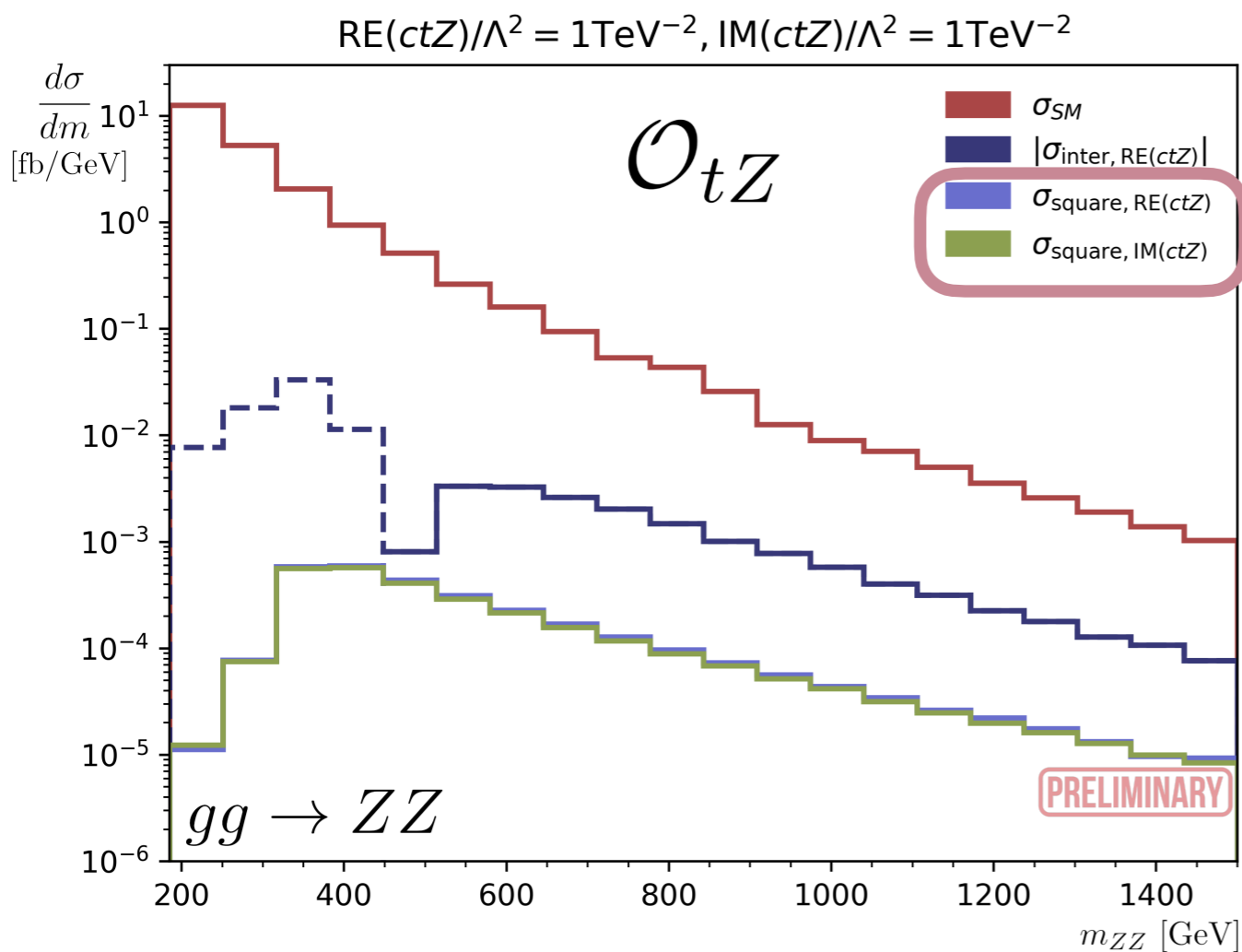
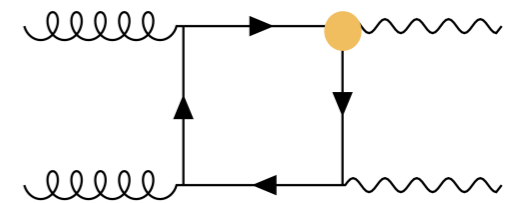


$\lambda_{g_1}, \lambda_{g_2}, \lambda_{H_1}, \lambda_{H_2}$	(+, +, 0, 0)	(+, -, 0, 0)
$\text{RE}c_{t\phi}$	$\frac{m_t v g_s^2 \delta_{ab}}{32\pi^2} \left[-3 \left(\log\left(\frac{s}{m_t^2}\right) - i\pi \right)^2 + 20 \right]$	s^0
$\text{IM}c_{t\phi}$	$\frac{m_t v g_s^2 \delta_{ab}}{32\pi^2} \left[-3i \left(\log\left(\frac{s}{m_t^2}\right) - i\pi \right)^2 \right]$	0

High energy behaviour of the helicity amplitudes
 $\rightarrow \text{RE}c_{t\phi}$ and $\text{IM}c_{t\phi}$ dominated by (+, +, 0, 0)

Growing amplitudes and tail effects

Modified top-Z interactions in $gg \rightarrow ZZ$



Total IM/RE ct_Z amplitudes dominated by:

$$\begin{aligned} &A(+, +, +, -) \\ &A(+, +, -, -) \\ &A(+, -, -, -) \end{aligned}$$

Log growth with energy

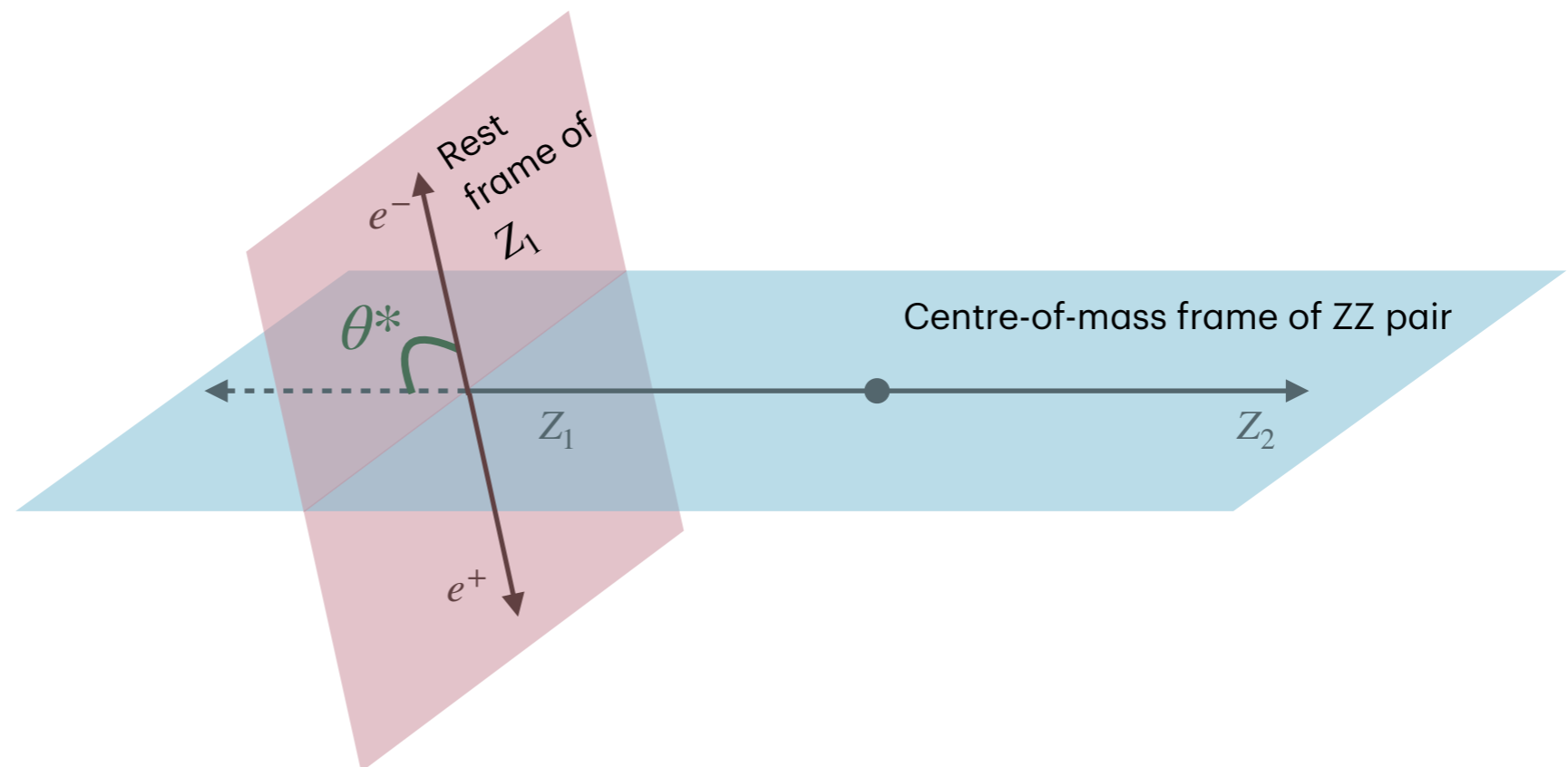
Same magnitudes for $\sqrt{s} > 300$ GeV
 \rightarrow overlapping σ_{square} distributions

Distinguishing CP-odd contributions with angular observables

Consider $gg \rightarrow ZZ \rightarrow e^+e^-\mu^+\mu^-$

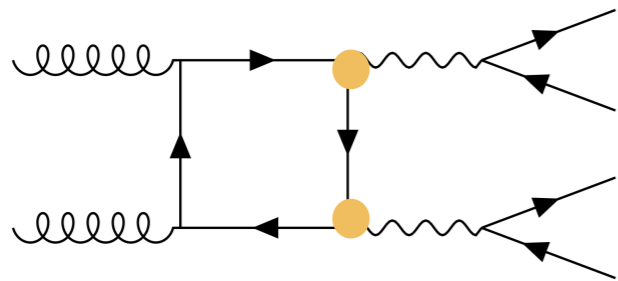
No kinematic cuts except on-shell Z bosons

θ^* : angle between Z boson momentum in ZZ pair centre-of-mass frame and momentum of e^- in Z rest frame

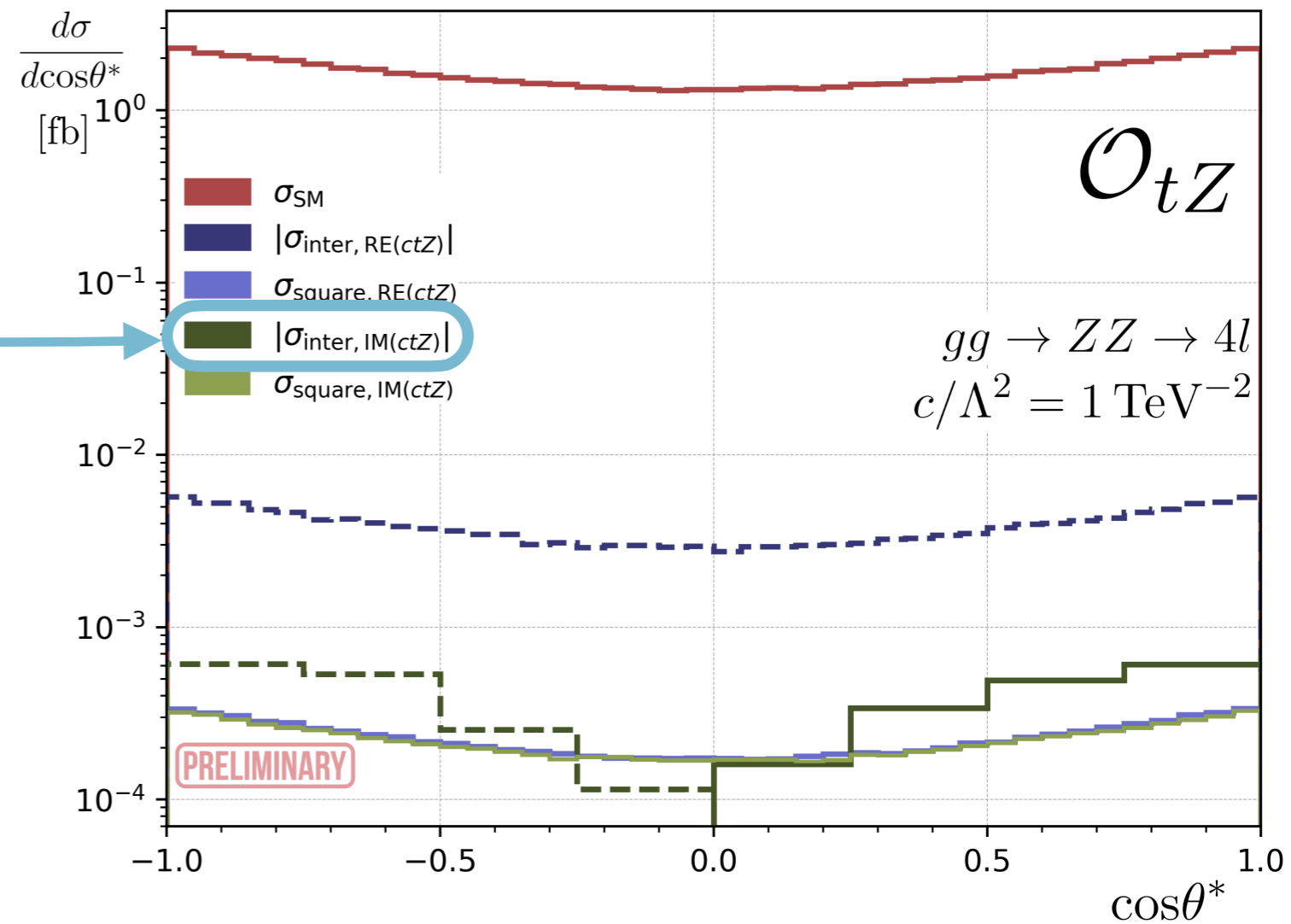


Distinguishing CP-odd contributions with angular observables

$\text{IM}ctZ$ interference is non-zero and odd around $\cos\theta^* = 0$



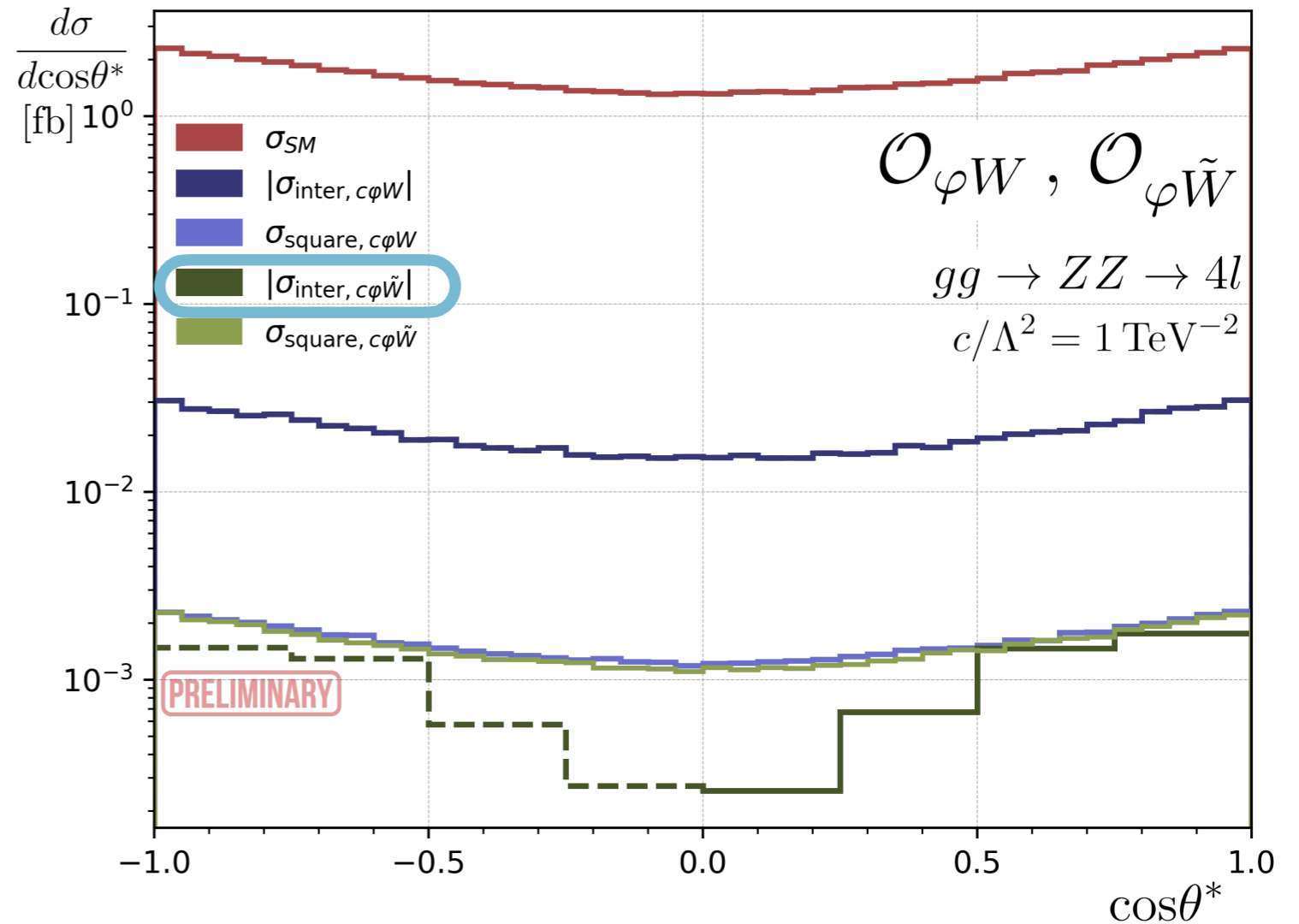
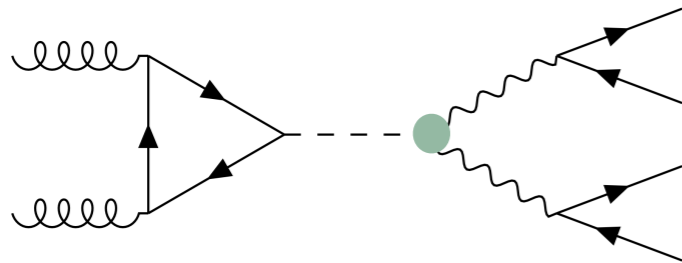
All other distributions are even around $\cos\theta^* = 0$



Dashed line = negative interference

Distinguishing CP-odd contributions with angular observables

Same behaviour for bosonic operators that enter in the tree part, for eg. $\mathcal{O}_{\varphi W}$ and $\mathcal{O}_{\varphi \tilde{W}}$



Dashed line = negative interference

Conclusion

- Study of CP-violation in the SMEFT at one-loop is only beginning.
- We implemented the dim-6 CP-odd operators entering in gluon-induced diboson production in Madgraph (\rightarrow calculation of rational terms and UV counterterms)
- In $gg \rightarrow ZZ \rightarrow 4l$ angular observables distinguish between CP-even and CP-odd contributions

 More is under study, stay tuned!

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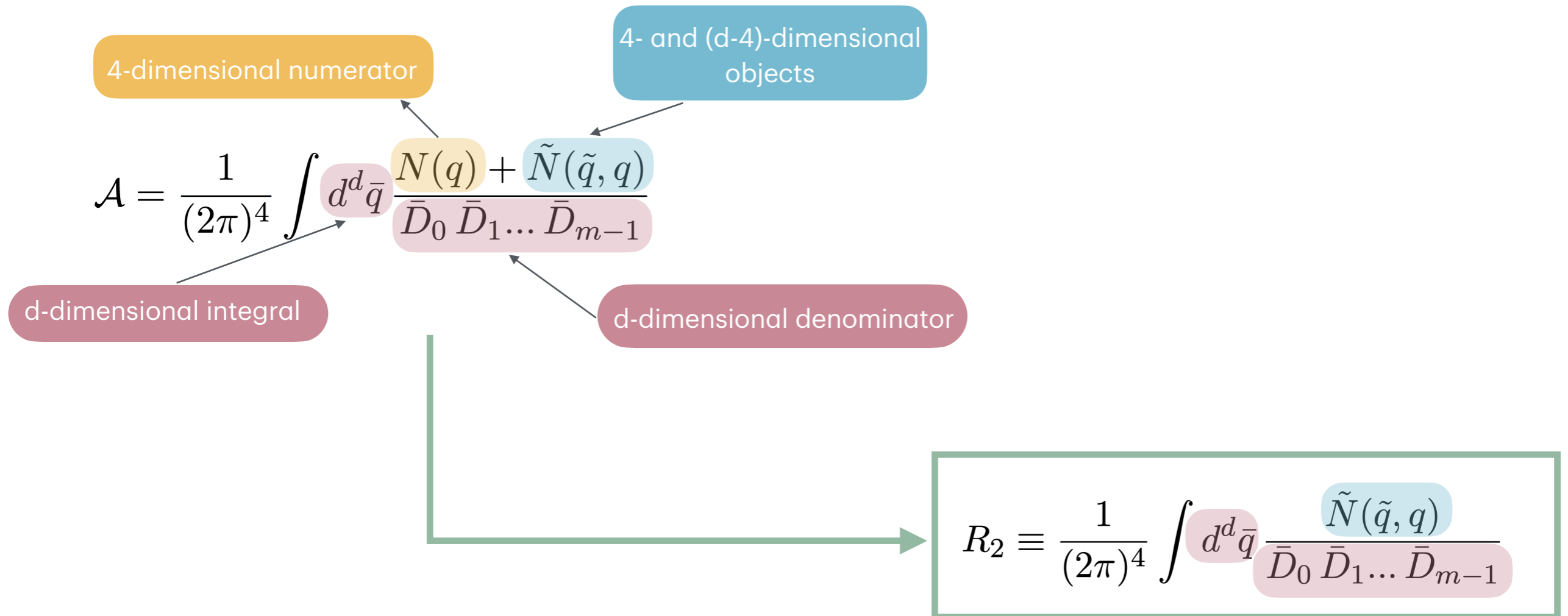
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Thank you !

What are the rational terms?

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In d-dimension, write m-point amplitude for 1-loop irreducible diagrams:



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