

A Geometrical Formalism of Functional Matching

Xu-Xiang Li 黎旭翔 (University of Utah)

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University of Bologna

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Based on work in preparation collaborating with Xiaochuan Lu and Zhengkang “Kevin” Zhang

Field Space Geometry in Higgs Sector

- SM Higgs doublet under different parametrization

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} -\phi^2 - i\phi^1 \\ \phi^4 + i\phi^3 \end{pmatrix}$$

$$\Delta\mathcal{L} = \frac{1}{2} \delta_{ij} (\partial_\mu \phi^i) (\partial^\mu \phi^j) - V(|\phi|^2)$$

$$\Delta\mathcal{L} = (\partial_\mu H)^\dagger (\partial^\mu H) - V(H^\dagger H)$$

$$H = U(x) \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + h \end{pmatrix}$$

$$U(x) = \exp\left(-\frac{i\sigma^a \pi^a}{v}\right)$$

$$\Delta\mathcal{L} = \frac{1}{2} (\partial_\mu h)^2 + \frac{1}{2} \left(1 + \frac{h}{v}\right)^2 g_{ab}(\pi) (\partial_\mu \pi^a) (\partial^\mu \pi^b) - V(h)$$

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Cartesian Coordinates

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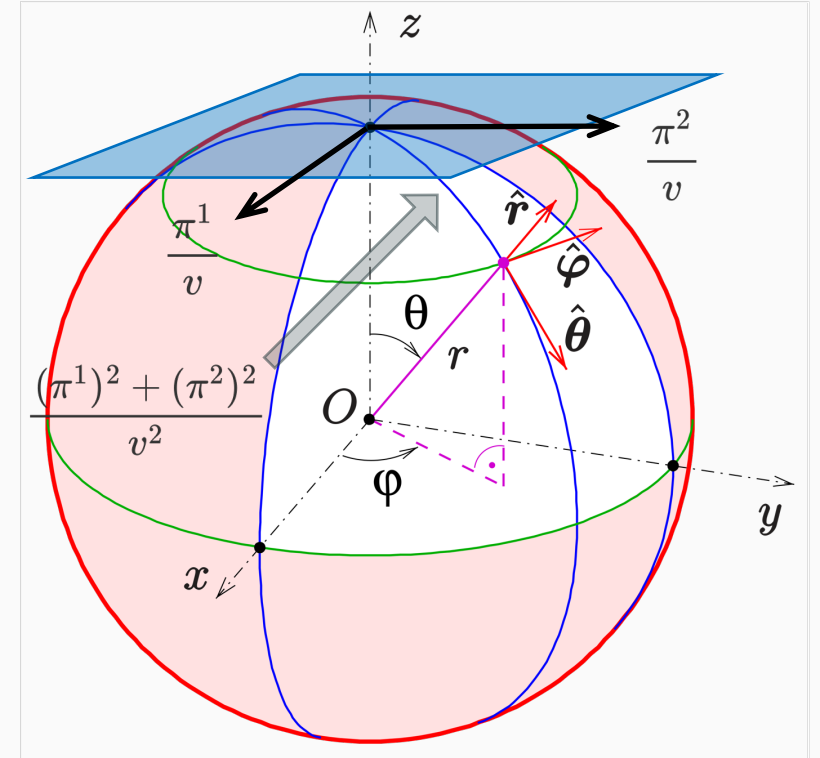
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Spherical Coordinates



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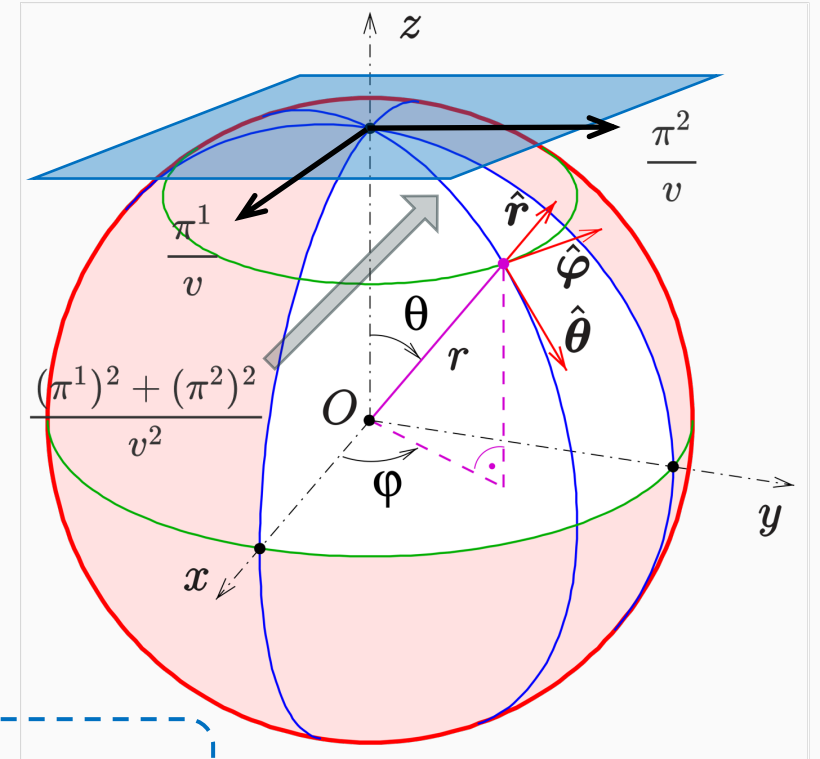
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Spherical Coordinates



Field Space Geometry in General Scalar Sector

- A general scalar theory with two derivatives

$$\mathcal{L} = \frac{1}{2} g_{ij}(\phi) (\partial_\mu \phi^i) (\partial^\mu \phi^j) - V(\phi)$$

- Under non-derivative field redefinition

$$\phi^i \rightarrow \phi'^i = \phi'^i(\phi)$$

$$(\partial_\mu \phi^i) \rightarrow (\partial_\mu \phi'^i) = \frac{\partial \phi'^i}{\partial \phi^j} (\partial_\mu \phi^j)$$

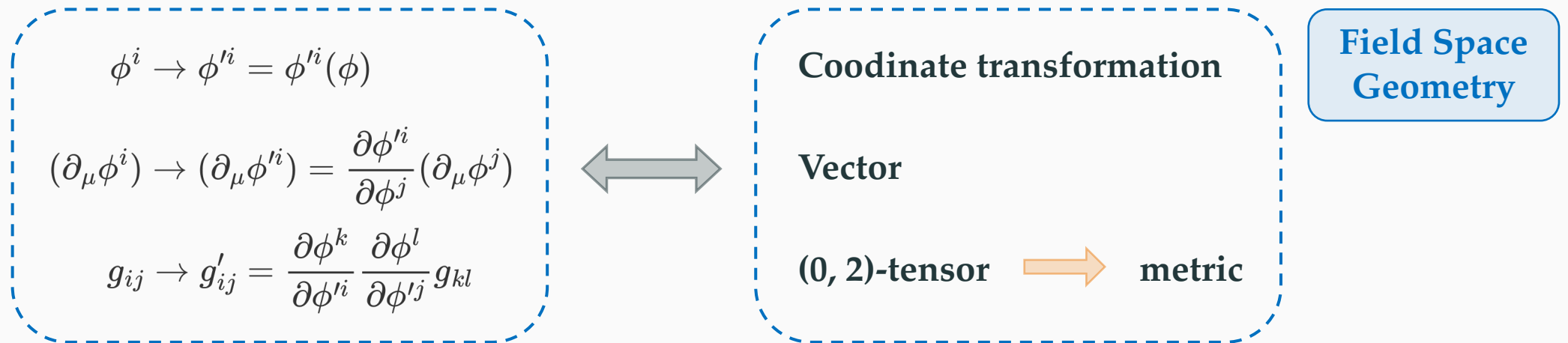
$$g_{ij} \rightarrow g'_{ij} = \frac{\partial \phi^k}{\partial \phi'^i} \frac{\partial \phi^l}{\partial \phi'^j} g_{kl}$$

Field Space Geometry in General Scalar Sector

- A general scalar theory with two derivatives

$$\mathcal{L} = \frac{1}{2} g_{ij}(\phi) (\partial_\mu \phi^i) (\partial^\mu \phi^j) - V(\phi)$$

- Under non-derivative field redefinition



- SM Higgs under Unitary representation:

$$g_{ij} = \begin{pmatrix} (1 + \frac{h}{v})^2 g_{ab}(\pi) & 0 \\ 0 & 1 \end{pmatrix} \rightarrow R_{ijkl} = 0 \quad \text{Flat space} \rightarrow \text{Renormalizable}$$

Alonso, Jenkins, Manohar [arXiv: 1511.00724, 1605.03602]
 Craig, Lee [arXiv: 2307.15742]
 Alminawi, Brivio, Davighi [arXiv: 2308.00017]

Partition Function & Effective Action

- Partition function

$$Z[J] = \int \mathcal{D}\phi \sqrt{\text{Det } g(\phi)} e^{i(S[\phi] + J\phi)}$$

- Tree level: Stationary point

$$\left. \frac{\delta S}{\delta \phi} \right|_{\phi=\bar{\phi}} + J = 0$$

- One-loop level: Gaussian approximation

$$\phi = \bar{\phi} + \phi'$$

$$S[\phi] + J\phi = S[\bar{\phi}] + J\bar{\phi} + \frac{1}{2} \phi' \left. \frac{\delta^2 S}{\delta \phi^2} \right|_{\phi=\bar{\phi}} \phi' + O(\phi'^3)$$

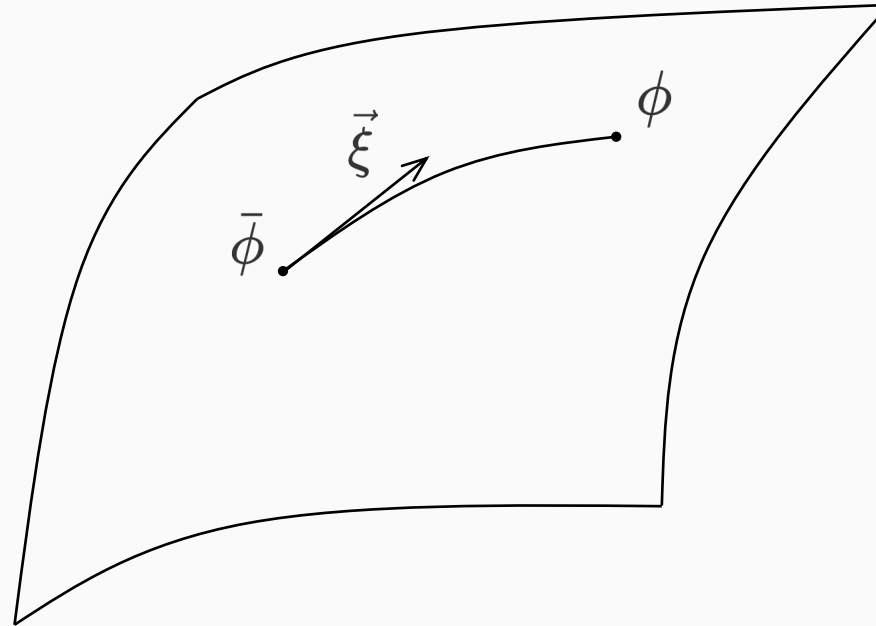
$$-i \log Z[J] = S[\bar{\phi}] + J\bar{\phi} + i \frac{1}{2} \text{Tr} \log \left(g^{-1} \frac{\delta^2 S}{\delta \phi^2} \right) \Big|_{\phi=\bar{\phi}} + \dots$$

$$\Gamma^{[0]}[\varphi] = S[\varphi]$$

$$\Gamma^{[1]}[\varphi] = i \frac{1}{2} \text{Tr} \log \left(g^{-1} \frac{\delta^2 S}{\delta \phi^2} \right) \Big|_{\phi=\varphi}$$

Covariant Expansion of Action

- Action = Function(al) on field space



$$S(\phi) \sim S(\bar{\phi}) + \xi^i \nabla_i S(\bar{\phi}) + \frac{1}{2} \xi^i \xi^j \nabla_i \nabla_j S(\bar{\phi}) + \dots$$

Covariant Expansion of Action

- Action = Function(al) on field space

Covariant expansion along geodesic from $\bar{\phi}$ to ϕ

$$\frac{\delta S}{\delta \phi^i} \rightarrow \frac{\nabla S}{\nabla \phi^i} = -g_{ij}(\phi) [\mathcal{D}_\mu(\partial^\mu \phi)]^j - \nabla_i V(\phi)$$

$$\frac{\delta^2 S}{\delta \phi^i \delta \phi^j} \rightarrow \frac{\nabla^2 S}{\nabla \phi^i \nabla \phi^j} = -[g_{ij}(\phi) \mathcal{D}^2 + R_{ijkl}(\partial_\mu \phi^k)(\partial^\mu \phi^l) + \nabla_i \nabla_j V(\phi)]_x \delta^4(x - y)$$

Gaillard, [Nucl. Phys. B 268, 669-692 (1986)]
 Alonso, Jenkins, Manohar [arXiv: 1605.03602]
 Alonso, West [arXiv:2207.02050]

- \Rightarrow Covariant 1-loop Effective action

- Divergent part \Rightarrow RGE

Alonso, Jenkins, Manohar [arXiv: 1511.00724]
 Alonso, Kanshin, Saa [arXiv: 1710.06848]
 Assi, Helset, Manohar, Pagès, Shen [arXiv:2307.03187]
 Jenkins, Manohar, Naterop, Pagès [arXiv:2308.06315]

- Hard region (div + finite) \Rightarrow EFT matching!

Fuentes-Martin, Portoles, Ruiz-Femenia [arXiv: 1607.02142]
 Zhang [arXiv: 1610.00710]

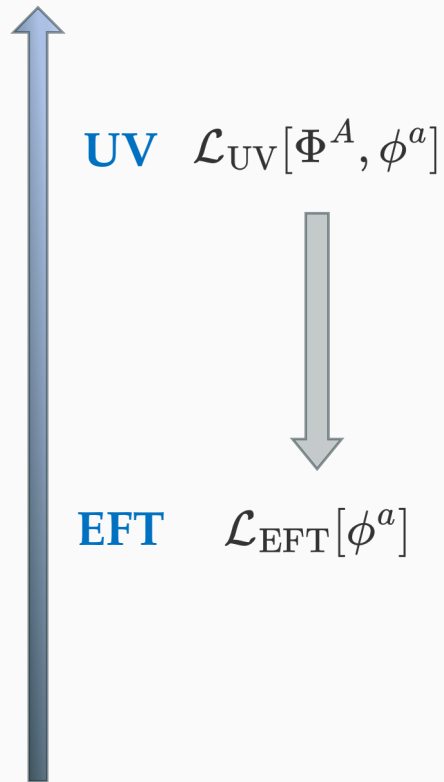
$$\mathcal{D}_\mu = \partial_\mu + (\partial_\mu \phi)^k \Gamma_{kj}^i(\phi) \sim (\partial_\mu \phi^i) \nabla_i$$

Covariant derivative on field space

Geometrical Functional Matching

- Top-down: find EFT from UV theory

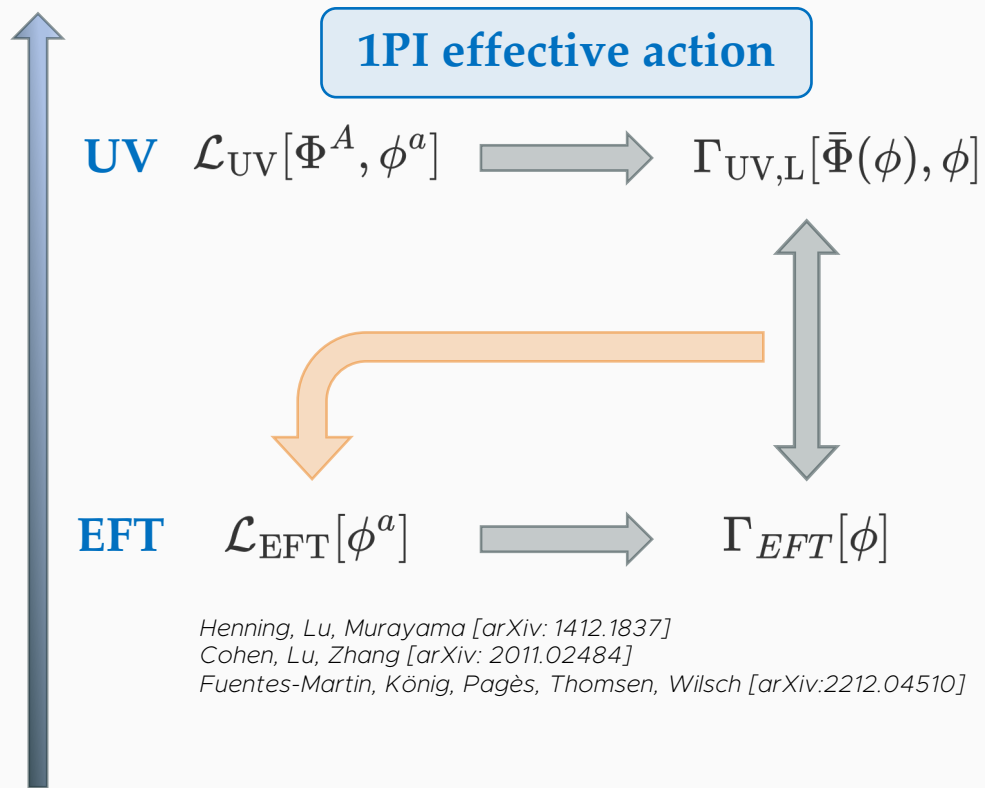
Energy



Geometrical Functional Matching

- Top-down: find EFT from UV theory

Energy



Henning, Lu, Murayama [arXiv: 1412.1837]

Cohen, Lu, Zhang [arXiv: 2011.02484]

Fuentes-Martin, König, Pagès, Thomsen, Wilsch [arXiv:2212.04510]

- Tree level:

$$\int d^4x \mathcal{L}_{EFT}^{[0]}[\phi] = \int d^4x \mathcal{L}_{UV}[\bar{\Phi}(\phi), \phi]$$

- One-loop level:

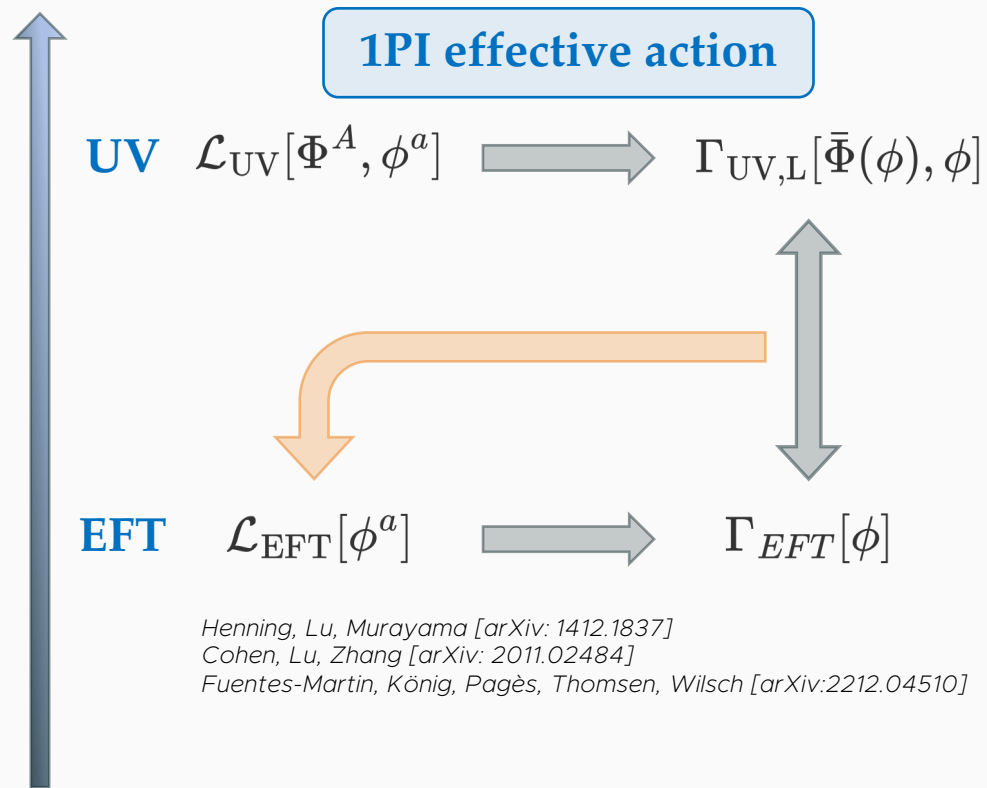
Covariant!

$$\int d^4x \mathcal{L}_{EFT}^{[1]}[\phi] = i \frac{1}{2} \int d^4x \int \frac{d^4p}{(2\pi)^4} \text{tr} \log \left(-g^{-1} \nabla^2 S_{UV} \right) \Big|_{i\partial_\mu \rightarrow i\partial_\mu - p_\mu, \Phi = \bar{\Phi}[\phi]}$$

Geometrical Functional Matching

- Top-down: find EFT from UV theory

Energy



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- Tree level:

$$\int d^4x \mathcal{L}_{EFT}^{[0]}[\phi] = \int d^4x \mathcal{L}_{UV}[\bar{\Phi}(\phi), \phi] \longleftarrow \text{Heavy EOM}$$

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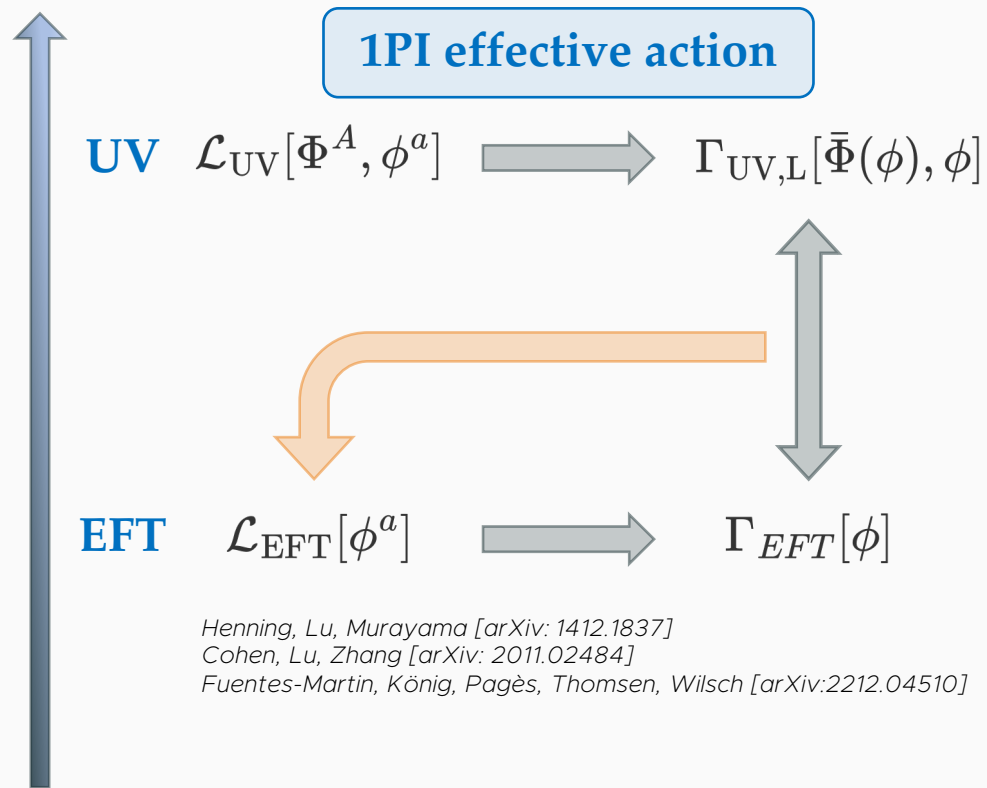
Take hard region

Covariant!

Geometrical Functional Matching

- Top-down: find EFT from UV theory

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Take hard region

~~Covariant!~~
 only covariance of light components (submanifold)

Geometrical Functional Matching

$$\mathcal{L}_{EFT}^{[1]}[\phi] = i \frac{1}{2} \int \frac{d^4 p}{(2\pi)^4} \text{tr} \log \left[(\mathcal{D}_\mu + ip_\mu)^2 + R^i{}_{kjl} (\partial_\mu \varphi^k) (\partial^\mu \varphi^l) + \nabla^i \nabla_j V \right] \Big|_{\Phi = \bar{\Phi}[\phi]}$$

- Divide into covariant quantities of submanifold

The diagram illustrates the decomposition of the covariant derivative \mathcal{D}_μ and the Riemann tensor R_{abcd} into intrinsic and extrinsic curvature components. The covariant derivative is shown as a 2x2 block matrix:

$$\mathcal{D}_\mu = \begin{pmatrix} \mathcal{D}_\mu^L & K_\mu^{LH} \\ K_\mu^{HL} & \mathcal{D}_\mu^H \end{pmatrix}$$

The Riemann tensor is decomposed as:

$$R_{abcd} = \hat{R}_{abcd} + K_{ac} K_{bd} - K_{ad} K_{bc}$$

Annotations: A green box labeled "extrinsic curvature" points to the K terms in the \mathcal{D}_μ matrix and the $K_{ac} K_{bd} - K_{ad} K_{bc}$ terms in the Riemann tensor decomposition. An orange box labeled "intrinsic curvature" points to the \hat{R}_{abcd} term in the Riemann tensor decomposition.

- Take hard region: separate hard scale from potential

The diagram shows the separation of a hard scale from the potential in the EFT Lagrangian. The potential term is decomposed as:

$$\nabla^i \nabla_j V \supset \mathbf{M}^2 = \begin{pmatrix} 0 & 0 \\ 0 & M^2 \end{pmatrix} \longrightarrow \mathcal{L}_{EFT}^{[1]}[\phi] = i \frac{1}{2} \int \frac{d^4 p}{(2\pi)^4} \text{tr} \log \left[(\mathcal{D}'_\mu + K_\mu + \underbrace{(ip_\mu)^2 + \mathbf{M}^2}_{\text{hard scale}} + U \right] \Big|_{\Phi = \bar{\Phi}[\phi]}$$

A blue box labeled "hard scale" points to the $(ip_\mu)^2 + \mathbf{M}^2$ term in the EFT Lagrangian.

Geometrical Functional Matching

$$\mathcal{L}_{EFT}^{[1]}[\phi] = i \frac{1}{2} \int \frac{d^4 p}{(2\pi)^4} \text{tr} \log \left[(\mathcal{D}'_\mu + K_\mu + ip_\mu)^2 + \mathbf{M}^2 + U \right] \Big|_{\Phi = \bar{\Phi}[\phi]}$$

- **Geometrical Covariant Derivative Expansion: keep the subcovariance when evaluating the momentum integral**

$$\mathcal{D}_\mu = \partial_\mu + (\partial_\mu \phi)^k \Gamma_{kj}^i(\phi) \longleftrightarrow D_\mu = \partial_\mu - ig A_\mu^a T^a$$

use CDE trick

Gaillard, [Nucl. Phys. B 268, 669-692 (1986)]
 Henning, Lu, Murayama [arXiv: 1412.1837]
 Cohen, Lu, Zhang [arXiv: 2011.02484]
 Fuentes-Martin, König, Pagès, Thomsen, Wilsch [arXiv:2212.04510]

- **Result: effective Lagrangian built by geometric quantities**

$$\mathcal{L}_{\text{eff}} \supset -\frac{i}{2} \int_{\text{hard}} \frac{d^4 p}{(2\pi)^4} \text{tr} \sum_{n=1}^{\infty} \frac{1}{n} \left\{ \sum_{m=0}^{\infty} \left[\frac{1}{p^2 - \mathbf{M}^2} \left(\left\{ p^\mu, \tilde{\mathcal{G}}'_{\mu\nu} \partial_p^\nu \right\} - \left(\tilde{\mathcal{G}}'_{\mu\nu} \partial_p^\nu \right)^2 \right) \right]^m \frac{1}{p^2 - \mathbf{M}^2} \left(\tilde{K}_\mu \tilde{K}^\mu + i \left\{ p_\mu - \tilde{\mathcal{G}}'_{\mu\nu} \partial_p^\nu, \tilde{K}^\mu \right\} \right) \right\}^n$$

Example: 'Non'linear Sigma Model

- Reparametrized linear sigma model:

$$\mathcal{L}_{UV} = \frac{1}{2}(\partial_\mu h)^2 + \frac{1}{2} \left(1 + \frac{h}{v}\right)^2 \left[g_{ab}(\pi) (\partial_\mu \pi^a) (\partial^\mu \pi^b) \right] - \frac{1}{4} \lambda (h^2 + 2hv)^2$$

- h : radial, π : spherical, field space is flat. g the metric on π sphere.
- $O(N+1)$ symmetry breaking \Rightarrow massive h
- Integrate out $h \Rightarrow$ EFT of π (up to dim-8 under canonical dimension counting)

The diagram illustrates the decomposition of the EFT Lagrangian into three parts: potential, extrinsic curvature, and intrinsic curvature. Arrows point from the labels to the corresponding terms in the equation below.

$$\begin{aligned} \mathcal{L}_{EFT}^{[1]}[\pi] &= \frac{1}{2} A(M^2) \left[U \right] + K_\mu \left[K^\mu \right]_{hh} + \frac{1}{d} B(0, M^2, 0) (K^\mu \left[\mathcal{G}_{\mu\nu} \right] K^\nu)_{hh} + \dots \\ &= \frac{1}{16\pi^2} \left[\frac{1}{2} \lambda (-6L + 7) \hat{g} + \frac{\lambda^2}{9M^4} (-102L + 67) \hat{g}^2 + \frac{\lambda^2}{9M^4} (12L - 34) (\partial_\mu \pi \cdot \partial_\nu \pi)^2 \right] \end{aligned}$$

Summary

- **Geometrical Functional Matching: keep geometric covariance during EFT matching**
 - Geometry on field space helps expand action covariantly
 - Hard region expansion breaks covariance, **only covariance on submanifold survives**
 - **Geometric covariant derivative expansion: preserve the covariance**
 - EFT built by geometric quantities, e.g. extrinsic/intrinsic curvature

Thank you!
