

Hard probes and nuclear PDFs

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Hard Probes 2024
26 September 2024

Hard probes and the need for nuclear PDFs

“Hard probes, e.g. high- p_T hadrons and heavy quarks, are formed in the early stages of the collisions via initial **hard scattering** of high-energy **partons** (quarks and gluons). Their production mechanism is well understood by perturbative QCD. Thus they serve as “calibrated probes” injected in the QGP medium.”

Akiba, Prog. Theor. Exp. Phys. (2015) 03A105

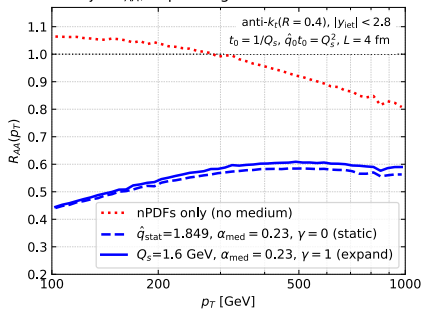
But in order to have these probes *well calibrated*, we need to know the distributions of the initial-state partons in the colliding nuclei precisely enough. I.e. need to study the **nuclear PDFs** (nPDFs).

cf. Apolinário et al., Prog. Part. Nucl. Phys. 127 (2022) 103990

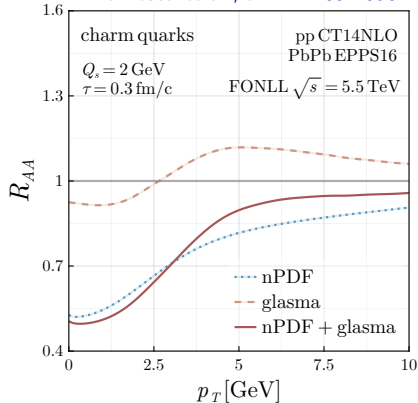
Pablos, EPJ Web of Conferences 296 (2024) 01028

Caucal et al., JHEP 04 (2021) 209

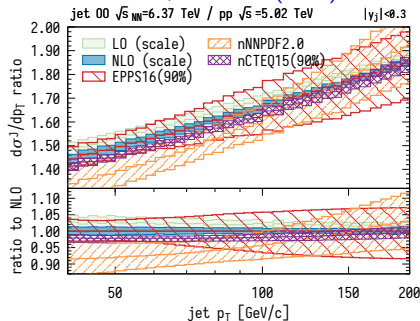
jet R_{AA} , expanding v. static medium



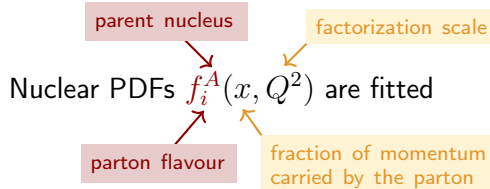
Avramescu et al., arXiv:2409.10564



Brewer et al., PRD 105 (2022) 074040



Nuclear PDFs from global analyses



to inclusive hard cross section data

- use $\{e, \nu, \pi, p\} + A$ collisions to avoid hot-QCD effects
- rely only to the QCD collinear factorization
- use model-agnostic *parametrisations* of nuclear effects as a function of x

Use statistical inference, minimize:

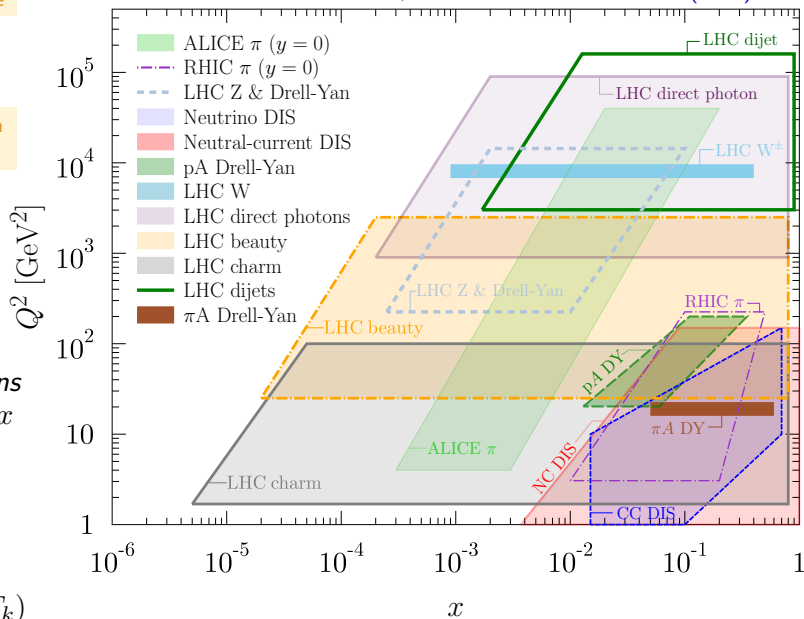
Sum over data sets

$$\chi_{\text{tot}}^2 = \sum_k (D_k - T_k)^T C_k^{-1} (D_k - T_k)$$

↓ data ↑ theory ↑ cov.

Data correlations important!

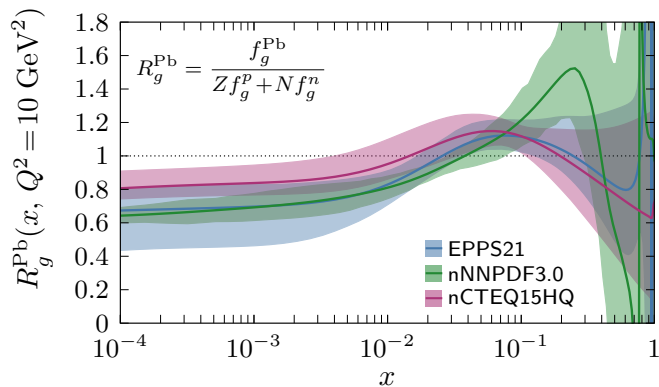
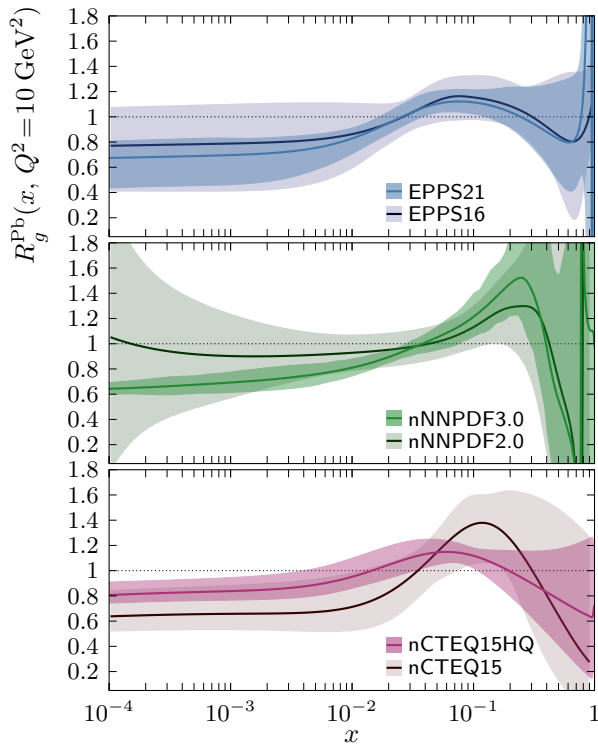
New review: Klasen & Paukkunen, Ann. Rev. Nucl. Part. Sci. 74 (2024) 1–41



Summary of recent nPDF global fits

| | KSASG20 | TUJU21 | EPPS21 | nNNPDF3.0 | nCTEQ15HQ* |
|---------------------------------|-----------------|-----------------|--------------------|--------------------|-------------------------|
| Order in α_s | NLO & NNLO | NLO & NNLO | NLO | NLO | NLO |
| IA NC DIS | ✓ | ✓ | ✓ | ✓ | ✓ |
| νA CC DIS | ✓ | ✓ | ✓ | ✓ | |
| pA DY | ✓ | | ✓ | ✓ | ✓ |
| πA DY | | | ✓ | | |
| RHIC dAu π^0, π^\pm | | | ✓ | | ✓ |
| LHC pPb π^0, π^\pm, K^\pm | | | | | ✓ |
| LHC pPb dijets | | | ✓ | ✓ | |
| LHC pPb HF | | | ✓ ^{GMVFN} | ✓ ^{FO+PS} | ✓ ^{ME fitting} |
| LHC pPb W,Z | | ✓ | ✓ | ✓ | ✓ |
| LHC pPb γ | | | | ✓ | |
| Q, W cut in DIS | 1.3, 0.0 GeV | 1.87, 3.5 GeV | 1.3, 1.8 GeV | 1.87, 3.5 GeV | 2.0, 3.5 GeV |
| p_T cut in inc.-h, HF | N/A | N/A | 3.0, 3.0 GeV | N/A, 0.0 GeV | 3.0, 3.0 GeV |
| Data points | 4353 | 2410 | 2077 | 2188 | 1484 |
| Free parameters | 18 | 16 | 24 | 256 | 19 |
| Error analysis | Hessian | Hessian | Hessian | Monte Carlo | Hessian |
| Free-proton PDFs | CT18 | own fit | CT18A | ~NNPDF4.0 | ~CTEQ6M |
| Free-proton corr. | no | no | yes | yes | no |
| HF treatment | FONLL | FONLL | S-ACOT | FONLL | S-ACOT |
| Indep. flavours | 3 | 4 | 6 | 6 | 5 |
| Reference | PRD 104, 034010 | PRD 105, 094031 | EPJC 82, 413 | EPJC 82, 507 | PRD 105, 114043 |

*see also [A. Kusina, Mon 3:00 PM](#) for preliminary nCTEQ24/25 results



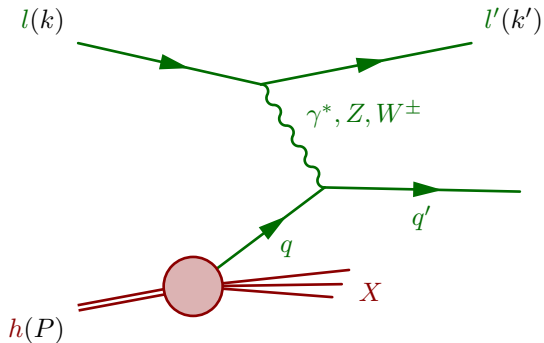
All major global nPDF fits find significant reduction in gluon uncertainties when including LHC data

Constraints driven by dijets & heavy-flavour, but also Ws and light mesons carry sensitivity

Differences between sets due to methodological and data-selection choices

Probes with leptonic final states

Deep inelastic scattering (DIS)

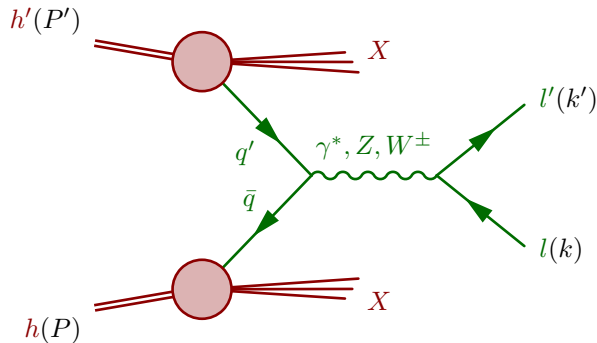


For the photon-mediated case:

$$\frac{d^2\sigma^{\text{DIS}}}{dx dQ^2} = \frac{d^2\hat{\sigma}}{dx dQ^2} \sum_{i \in \{q, \bar{q}\}} e_i^2 f_i^h(x, Q^2) + \text{NLO corrections}$$

$$\left. \begin{aligned} Q^2 &= -(k - k')^2 \\ x &= \frac{Q^2}{2P \cdot (k - k')} \end{aligned} \right\} \leftarrow \begin{array}{l} \text{access scale and momentum-} \\ \text{fraction dependence through} \\ \text{external kinematics} \end{array}$$

Drell-Yan (DY)



The photon-mediated case:

$$\frac{d^2\sigma^{\text{DY}}}{dy dM^2} = \frac{4\pi\alpha_{\text{e.m.}}^2}{9M^4} \sum_{i \in \{q, \bar{q}\}} e_i^2 x_1 x_2 f_i^h(x_1, M^2) f_i^{h'}(x_2, M^2) + \text{NLO corrections}$$

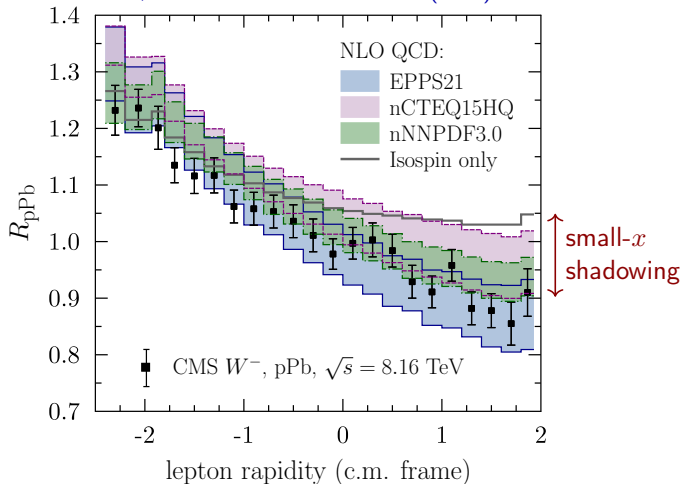
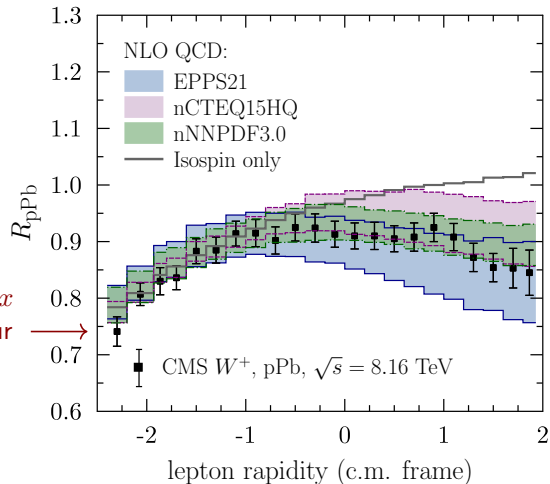
$$M^2 = (k + k')^2 = x_1 x_2 s$$

$$y = \frac{1}{2} \log \frac{(k_0 + k'_0) + (k_3 + k'_3)}{(k_0 + k'_0) - (k_3 + k'_3)} = \frac{1}{2} \log \frac{x_1}{x_2}$$

W bosons in pPb at 8.16 TeV

pPb data from: CMS Collaboration, Phys. Lett. B 800 (2020) 135048
pp baseline: CMS Collaboration, Eur. Phys. J. C 76 (2016) 469

Klasen & Paukkunen, Ann. Rev. Nucl. Part. Sci. 74 (2024) 1–41



Run-2 W boson data included in practically all recent nPDF fits:

nCTEQ15HQ
nNNPDF3.0
TUJU21

← Use absolute pPb cross sections

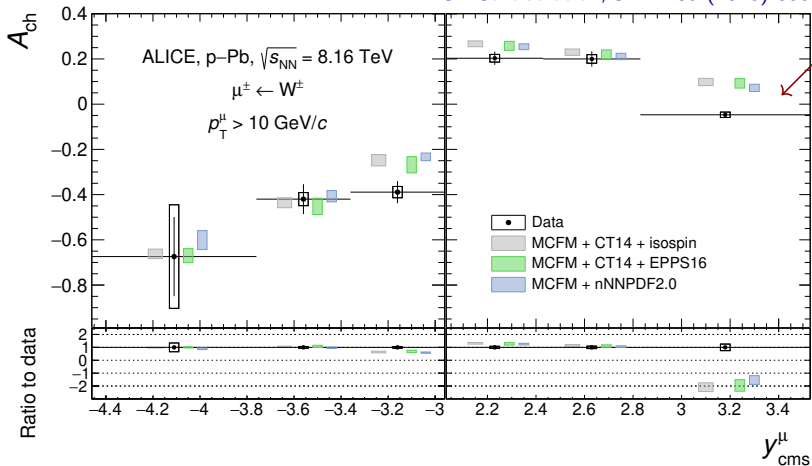
EPPS21

← Use R_{pPb} to cancel free-proton PDFs

$$R_{pPb} = \frac{\sigma^{pPb}(8.16 \text{ TeV})}{\sigma^{pp}(8.0 \text{ TeV})}$$

W boson charge asymmetry in pPb at 8.16 TeV

ALICE Collaboration, JHEP 05 (2023) 036



large data-theory difference
in a forward bin either due to:

strongly flavour dependent shadowing
(not supported by CMS R_{pPb})

large- x proton-PDF flavour asymmetry
(should be tested)

Charge asymmetry very sensitive to free-proton baseline

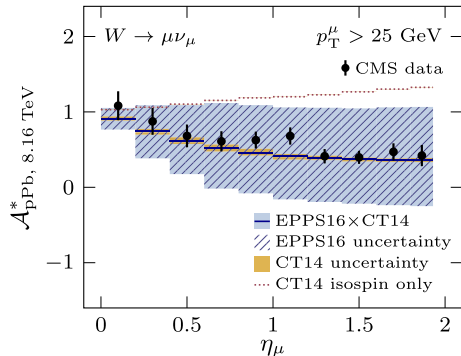
Paukkunen & Salgado, JHEP 03 (2011) 071

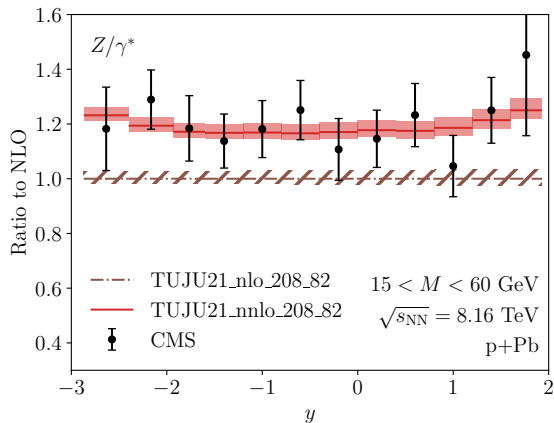
To probe nuclear modifications, use instead:

$$A_{pPb}^* = \frac{d\sigma_{pPb}^{W^+}/d\eta_{\mu}|_{\eta_{\mu}} - d\sigma_{pPb}^{W^-}/d\eta_{\mu}|_{-\eta_{\mu}}}{d\sigma_{pPb}^{W^+}/d\eta_{\mu}|_{-\eta_{\mu}} - d\sigma_{pPb}^{W^-}/d\eta_{\mu}|_{\eta_{\mu}}}$$

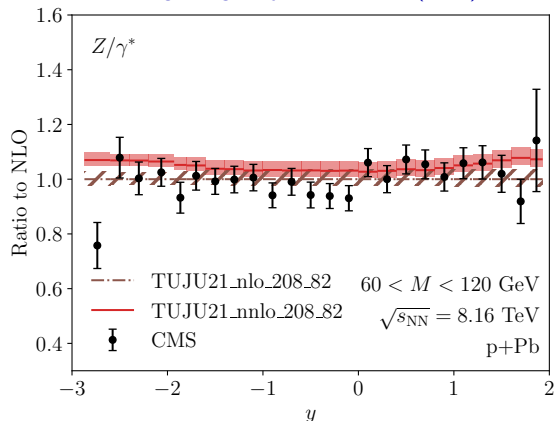
Eskola, PP, Paukkunen, Salgado, Eur. Phys. J. C 82 (2022) 271

Eskola, PP, Paukkunen, Salgado,
Eur. Phys. J. C 82 (2022) 271





[Helenius, Walt, Vogelsang, Phys. Rev. D 105 \(2022\) 094031](#)



Low-mass DY ($15 < M < 60$ GeV): First clear evidence for the need of NNLO nPDFs

[Helenius, Walt, Vogelsang, Phys. Rev. D 105 \(2022\) 094031](#)

Z-peak region ($60 < M < 120$ GeV): Poor χ^2 in all nPDF analyses due to data fluctuations

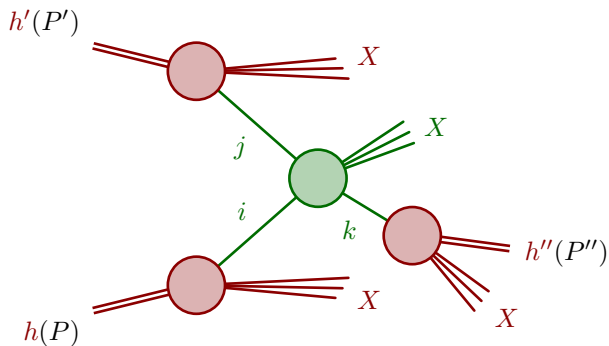
[Eskola, PP, Paukkunen, Salgado, EPJC 82 \(2022\) 413](#)
[Abdul Khalek et al., EPJC 82 \(2022\) 507](#)

Also ALICE & LHCb 8.16 TeV data available

[ALICE Collaboration, JHEP 09 \(2020\) 076](#)
[LHCb Collaboration, JHEP 06 \(2023\) 022](#)

Hadroproduction of hadronic final states

Hadron-production

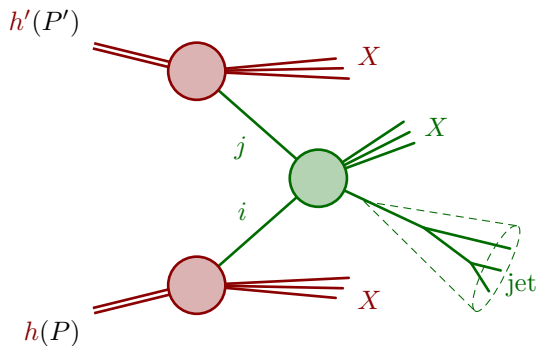


$$\sigma^{h+h' \rightarrow h''+X} = \sum_{i,j,k \in \{q,\bar{q},g\}} f_i^h \otimes f_j^{h'} \otimes \hat{\sigma}^{ij \rightarrow k+X} \otimes D_k^{h''}$$

Account for the hadronization effects with the *parton to hadron fragmentation functions* $D_k^{h''}$

→ a source of uncertainty for PDF fits

Jet-production



$$\sigma^{h+h' \rightarrow \text{jet}+X} = f_{\text{NP}} \sum_{i,j \in \{q,\bar{q},g\}} f_i^h \otimes f_j^{h'} \otimes \hat{\sigma}^{ij \rightarrow \text{jet}+X}$$

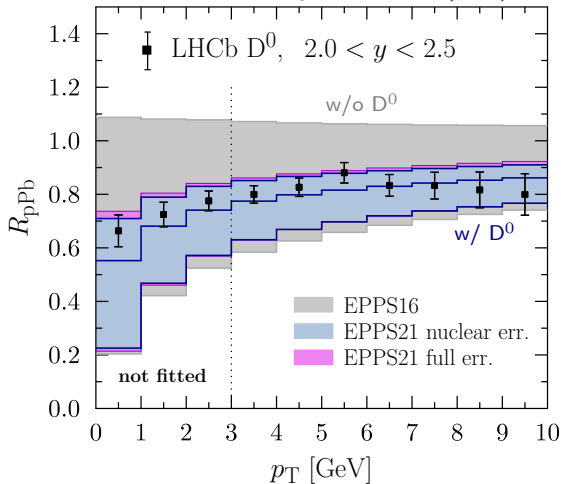
Instead of fragmentation functions:

- need an IR-safe definition of a jet
- non-perturbative corrections f_{NP}

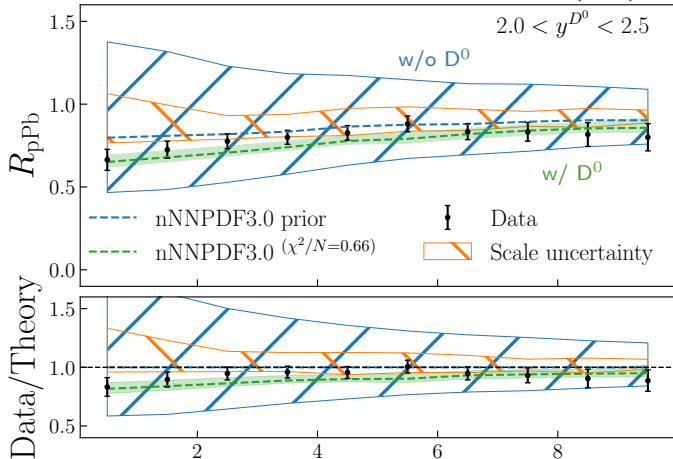
D^0 s in EPPS21 and nNNPDF3.0

data from: LHCb Collaboration, JHEP 10 (2017) 090

Eskola, PP, Paukkunen, Salgado, EPJC 82 (2022) 413



Abdul Khalek et al., EPJC 82 (2022) 507



nNNPDF3.0 with POWHEG+PYTHIA finds a large **scale uncertainty** in $R_{pPb} \rightarrow$ fit only forward data

Abdul Khalek et al., EPJC 82 (2022) 507

EPPS21 uses S-ACOT- m_T GM-VFNS \rightarrow scale uncertainty small except at low p_T

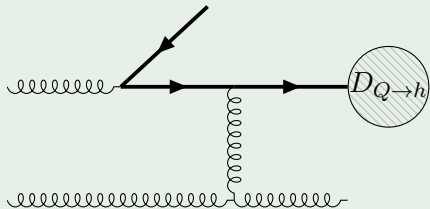
Helenius & Paukkunen, JHEP 05 (2018) 196
Eskola, Helenius, PP, Paukkunen, JHEP 05 (2020) 037

Heavy-flavour production mass schemes

FFNS

In *fixed flavour number scheme*,
valid at small p_T , heavy quarks are produced
only at the matrix element level

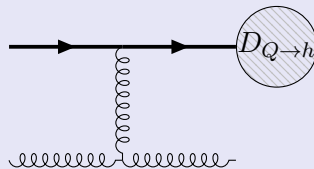
Contains $\mathcal{O}(m)$ and $\log(p_T/m)$ terms



ZM-VFNS

In *zero-mass variable flavour number scheme*,
valid at large p_T , heavy quarks are treated as
massless particles produced also in ISR/FSR

Resums $\log(p_T/m)$ but ignores $\mathcal{O}(m)$ terms



- subtraction term +

GM-VFNS

A *general-mass variable flavour number scheme* combines the two by supplementing subtraction terms
to prevent double counting of the resummed splittings, **valid at all** p_T

Resums $\log(p_T/m)$ and includes $\mathcal{O}(m)$ terms in the FFNS matrix elements

Important: includes also **gluon-to-HF fragmentation** – large contribution to the cross section!

A data-driven approach – nCTEQ15HQ

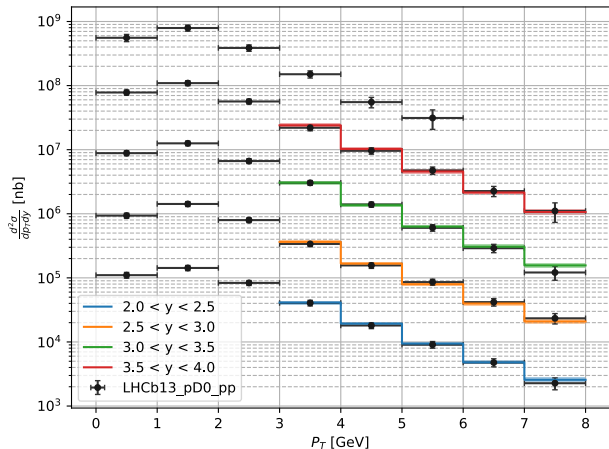
Duwentäster et al.,
PRD 105 (2022) 114043

nCTEQ15HQ uses a data-driven approach

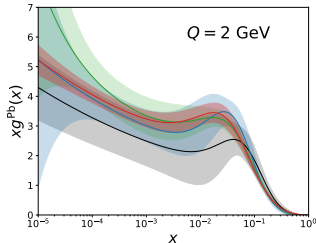
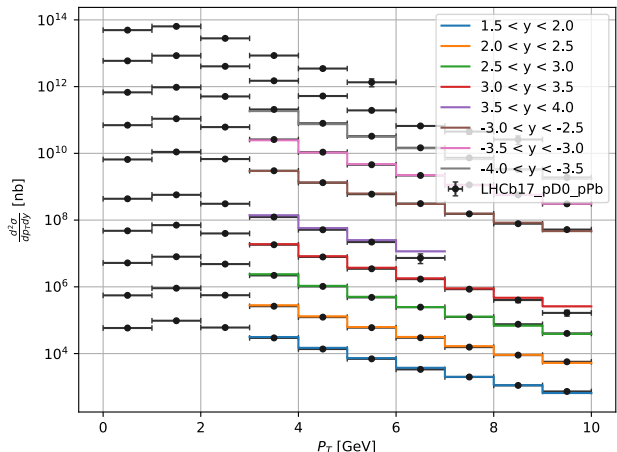
Lansberg & Shao, EPJC 77 (2017) 1
Kusina et al., PRL 121 (2018) 052004

to fit the D^0 and J/ψ data:

1. Fit the matrix elements to pp data...
(assume $2 \rightarrow 2$ kinematics, gg IS only)



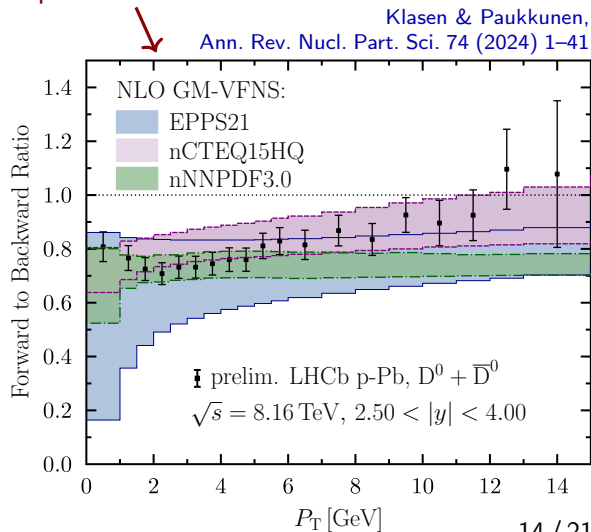
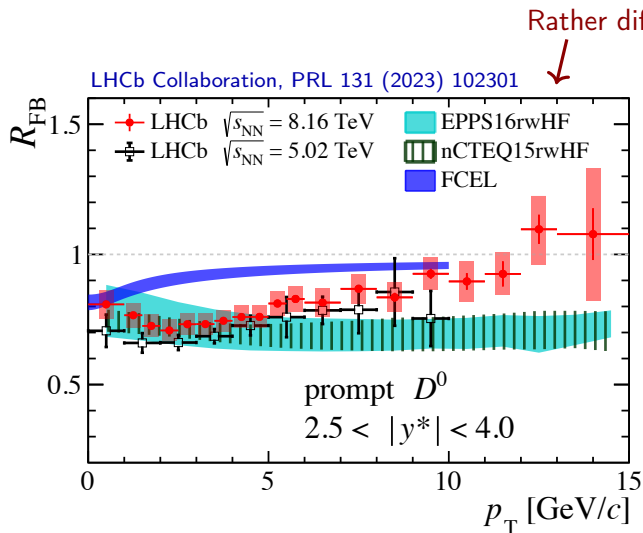
2. ... use the fitted matrix elements to fit nuclear PDFs with pPb data

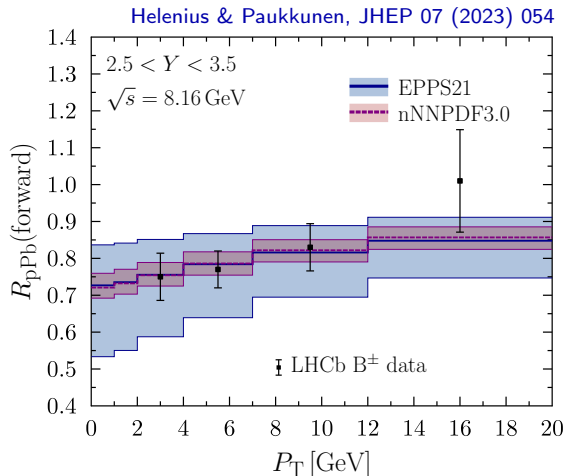
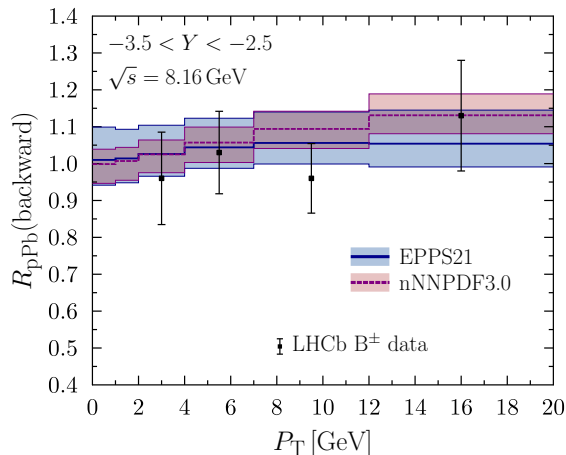


D^0 s in pPb at 8.16 TeV

New LHCb measurement at 8.16 TeV
initially claimed to be in tension with nPDFs
(not included in the nPDF analyses yet)

Not only probing nPDFs but also testing
production and interaction mechanism!
(Here HELAC vs. S-ACOT- m_T vs. FCEL)





B-meson production theoretically clean due to high b -quark mass, but scale-variation (\sim higher order) uncertainties can still be relevant in GM-VFNS at NLO towards low- p_T

Helenius & Paukkunen, JHEP 07 (2023) 054

LHCb data in agreement with S-ACOT- m_T using EPPS21 and nNNPDF3.0 nPDFs

→ Need more statistics for strong constraints

Neutral pions in pPb at 8.16 TeV

Forward π^0 s agree with D^0 -constrained nPDFs, but at backward rapidities this agreement seems to break down!

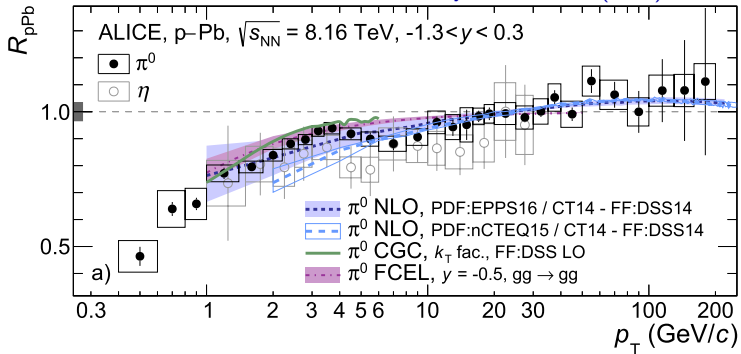
Do I see a bump also in midrapidity data at 3–4 GeV?

Similar/larger enhancements seen in charged hadrons

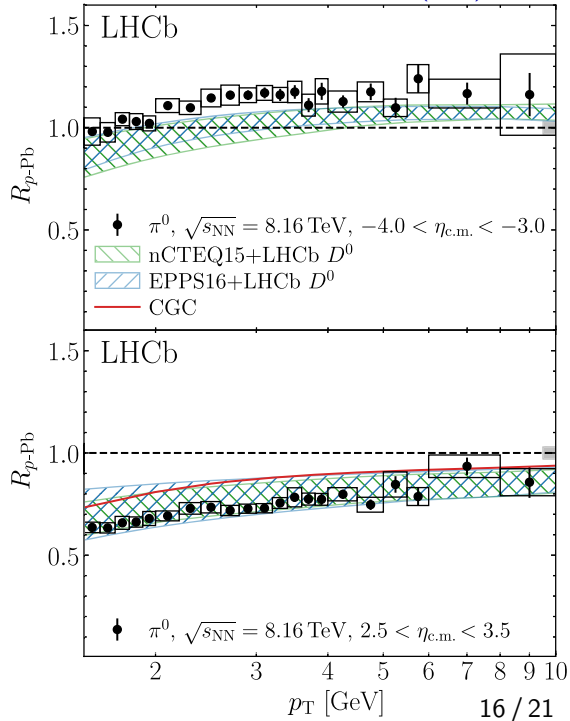
PHENIX Collaboration, PRC 101 (2020) 034910
PHENIX Collaboration, PRC 105 (2022) 064902
LHCb Collaboration, PRL 128 (2022) 142004

→ Traditional $p_T > 3$ GeV cut probably too small

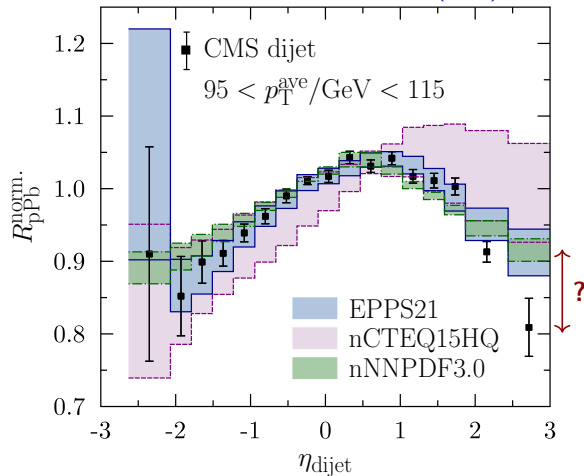
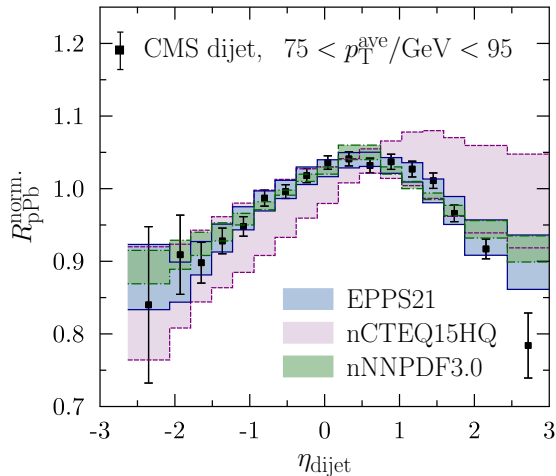
ALICE Collaboration, Phys. Lett. B 827 (2022) 136943



LHCb Collaboration, PRL 131 (2023) 042302



Klasen & Paukkunen, Ann. Rev. Nucl. Part. Sci. 74 (2024) 1–41



EPPS21 and nNNPDF3.0 include CMS dijets, given as self-normalized ratio to cancel hadronization effects:

$$R_{\text{pPb}}^{\text{norm.}} = \frac{d^2\sigma^{\text{pPb}}/dp_T^{\text{ave}}d\eta_{\text{dijet}}}{d\sigma^{\text{pPb}}/dp_T^{\text{ave}}} \bigg/ \frac{d^2\sigma^{\text{pp}}/dp_T^{\text{ave}}d\eta_{\text{dijet}}}{d\sigma^{\text{pp}}/dp_T^{\text{ave}}}$$

Inability to fit forward data due to the missing (induced) data correlations? Or NNLO / NP effects?

Dijets in pPb at 8.16 TeV

New measurement of dijets at 8.16 TeV from CMS

Preliminary results with rapidity spectra in 16 bins of p_T from 50 to 500 GeV

Note: the rapidity spectra carry non-vanishing proton-PDF dependence and scale uncertainty at NLO

Eskola, PP, Paukkunen, EPJC 79 (2019) 511

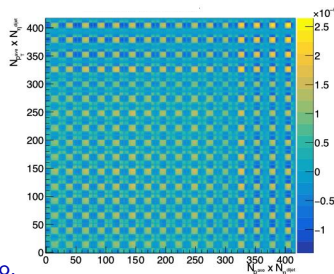
→ Cure with forward-to-backward ratio?

Eskola, Paukkunen, Salgado, JHEP 10 (2013) 213

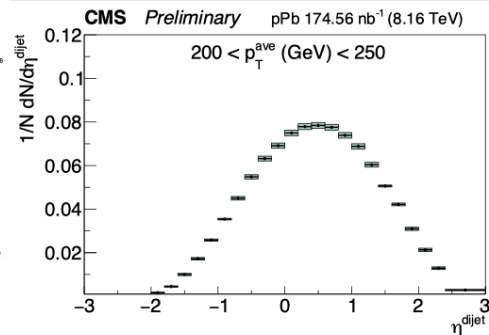
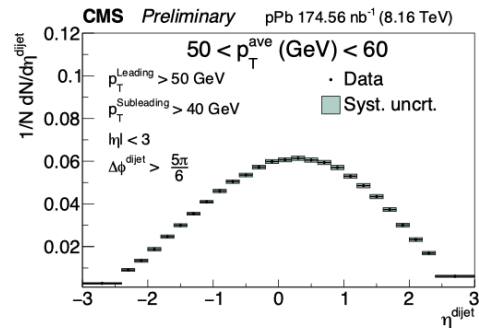
Systematic uncertainties important

- Correlations significant
- ✓ Covariance matrix expected to be provided
- N.B. For self-normalized quantities even statistical uncertainties become correlated

Eskola, PP, Paukkunen, Salgado, EPJC 82 (2022) 271



G. Nigmatkulov, Mon 2:00 PM



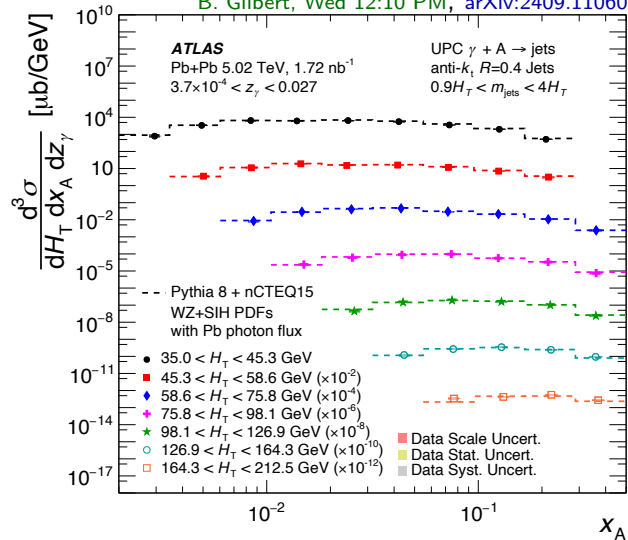
Novel probes: Inclusive dijets in UPCs

Dijet photoproduction in UPCs has been promoted as a probe of nuclear PDFs

Strikman, Vogt & White, PRL 96 (2006) 082001

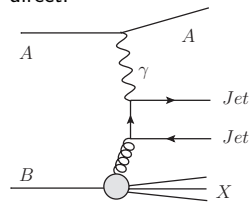
ATLAS measurement now final!

B. Gilbert, Wed 12:10 PM, arXiv:2409.11060

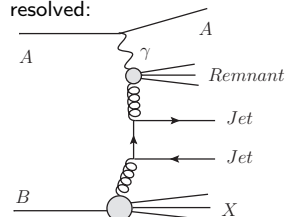


Guzey & Klasen, PRC 99 (2019) 065202

direct:



resolved:



Triple differential in

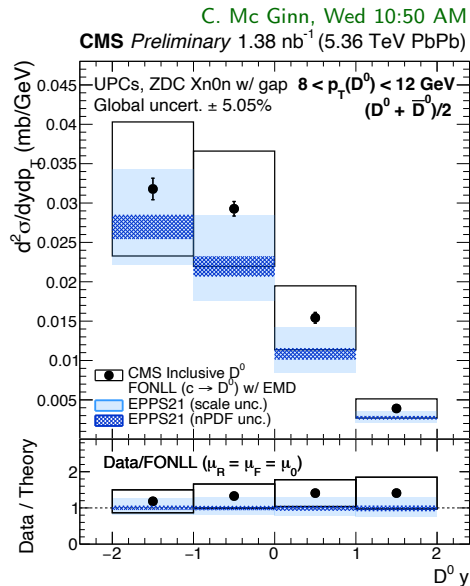
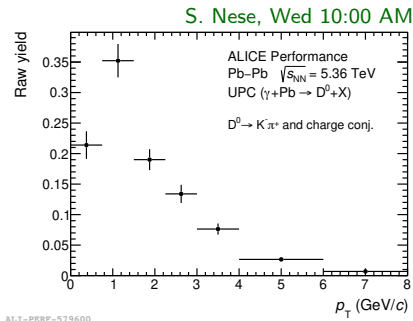
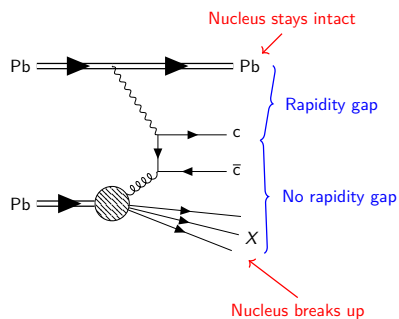
$$H_T = \sum_{i \in \text{jets}} p_{T,i}, \quad z_\gamma = \frac{M_{\text{jets}}}{\sqrt{s_{\text{NN}}}} e^{+y_{\text{jets}}},$$

$$x_A = \frac{M_{\text{jets}}}{\sqrt{s_{\text{NN}}}} e^{-y_{\text{jets}}}$$

Note: transverse-plane collision geometry gets resolved at large z_γ!

Eskola, Guzey, Helenius, PP, Paukkunen, arXiv:2404.09731

Novel probes: Inclusive UPC charmed meson photoproduction



ALICE plans to measure inclusive D^0 , D^\pm , D^* and J/ψ photoproduction

CMS preliminary UPC D^0 data in agreement with FONLL predictions within scale and systematic uncertainties

→ Important cross check w.r.t. D^0 hadroproduction data

Inclusive UPC processes emerging as new nPDF probes!

Nuclear PDFs are being constrained by an increasing amount of LHC data

- Recent global fits include few thousand data points on a variety of processes
 - Collinear factorization works in pA across a large phase space in x, Q^2
 - Nuclear gluon content better constrained than ever before
- Data to theory agreement is starting to become sensitive to NNLO effects (low-mass DY), treatment of heavy-flavour production (D^0), imposed p_T cuts...
 - Can aim at over-constraining nuclear PDFs to find the limits where additional effects (saturation, cold nuclear matter energy loss, hadronization nonuniversality) become significant
 - Need both precision data and calculations
- *Inclusive* UPC processes emerging as new nPDF probes!

Thank you!

... and my apologies I could not cover every nPDF related analysis available

Backup

Collinear factorization in perturbative QCD

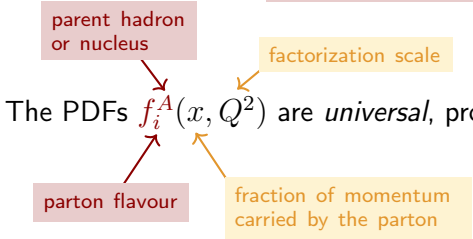
The cross section for producing an inclusive final state $k + X$ can be described as a convolution of...

... Coefficient Functions $\hat{d}^{ij \rightarrow k+X'}$ which are calculable from perturbative QCD...

$$d\sigma^{AB \rightarrow k+X}(Q^2) \stackrel{Q \gg \Lambda_{\text{QCD}}}{=} \sum_{i,j,X'} f_i^A(Q^2) \otimes \hat{d}^{ij \rightarrow k+X'}(Q^2) \otimes f_j^B(Q^2) + \mathcal{O}(1/Q^2)$$

... and Parton Distribution Functions f_i^A, f_j^B which contain long-range physics and cannot be obtained by perturbative means...

... plus "Higher Twist" corrections which are suppressed at high enough momentum scale $Q \gg \Lambda_{\text{QCD}}$

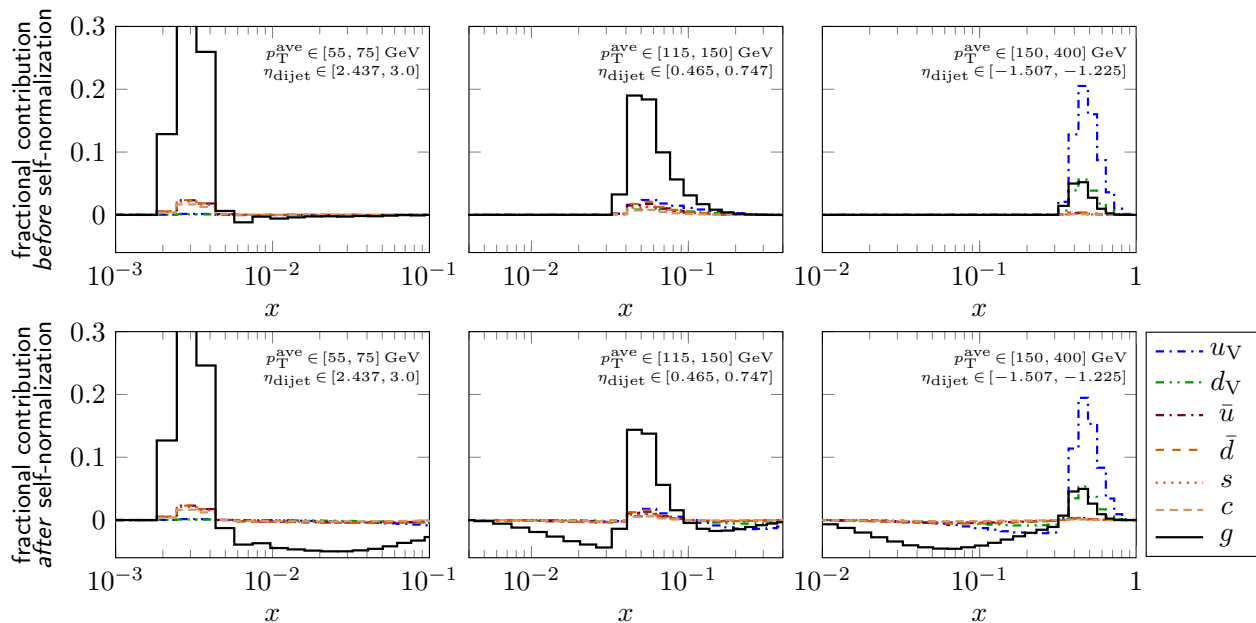


and obey the DGLAP equations $Q^2 \frac{\partial f_i^A}{\partial Q^2} = \sum_j P_{ij} \otimes f_j^A$

splitting functions

Mellin conv.

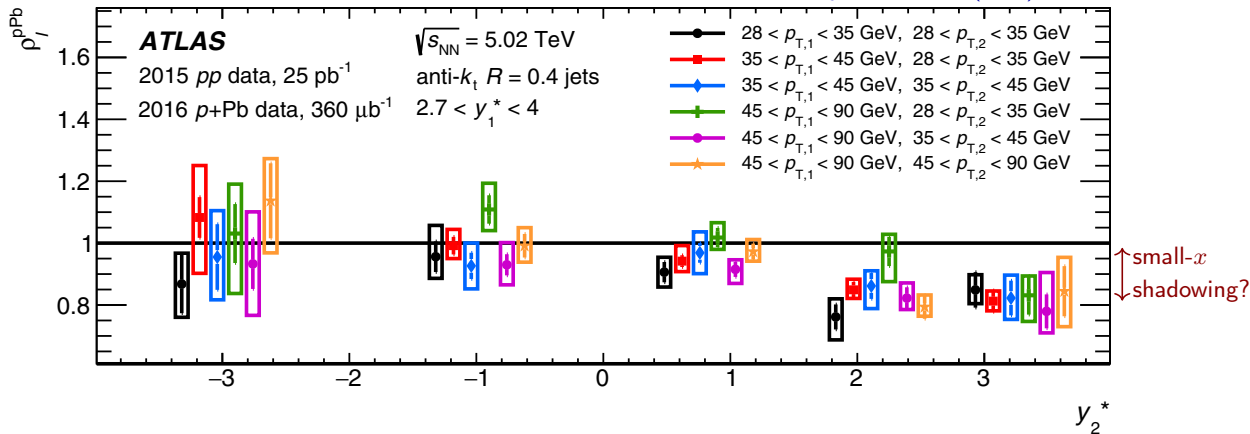
Dijets in pPb at 5.02 TeV – x -dependence at NLO pQCD



Self-normalization introduces additional (anti-)correlation between data points and across probed x !

Dijets in pPb at 5.02 TeV – per-trigger yields

ATLAS Collaboration, Phys. Rev. C 100 (2019) 034903



Nuclear modification of per-trigger yields, practically a ratio of 2-jet over 1-jet ratios:

$$\rho_I^{p\text{Pb}} = \frac{d^4\sigma^{p\text{Pb}}_{2\text{-jet}}/dp_{T,1}dy_1^*dp_{T,2}dy_2^*}{\sigma^{p\text{Pb}}_{1\text{-jet}}} \bigg/ \frac{d^4\sigma^{\text{pp}}_{2\text{-jet}}/dp_{T,1}dy_1^*dp_{T,2}dy_2^*}{\sigma^{\text{pp}}_{1\text{-jet}}}$$

Forward suppression indicating small- x shadowing?

Novel probes: Exclusive UPC J/ψ photoproduction in collinear factorization

First phenomenological implementation of the NLO corrections

Ivanov et al., EPJC 34 (2004) 297

Jones et al., J. Phys. G 43 (2016) 035002

in ultraperipheral Pb+Pb

Eskola et al., PRC 106 (2022) 035202

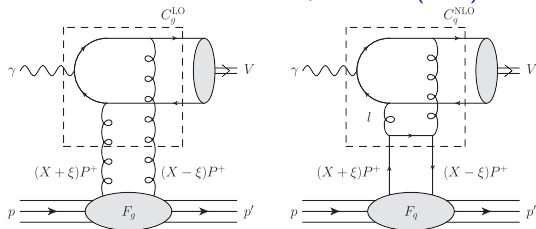
Exclusive process

→ need a mapping between GPDs and traditional PDFs

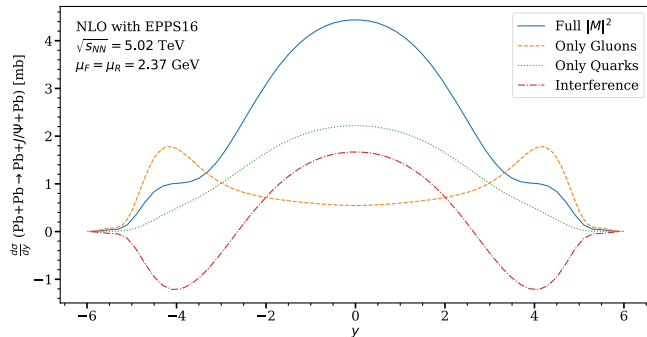
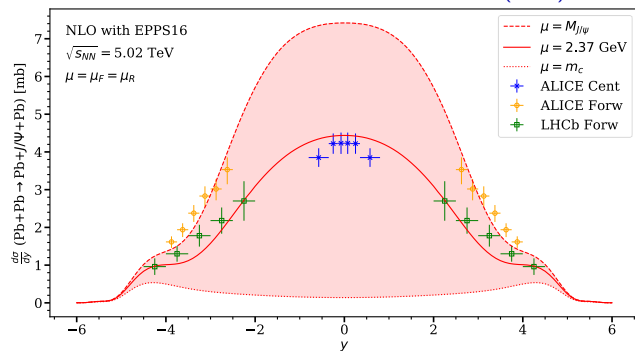
Large scale uncertainty

→ perturbative convergence?

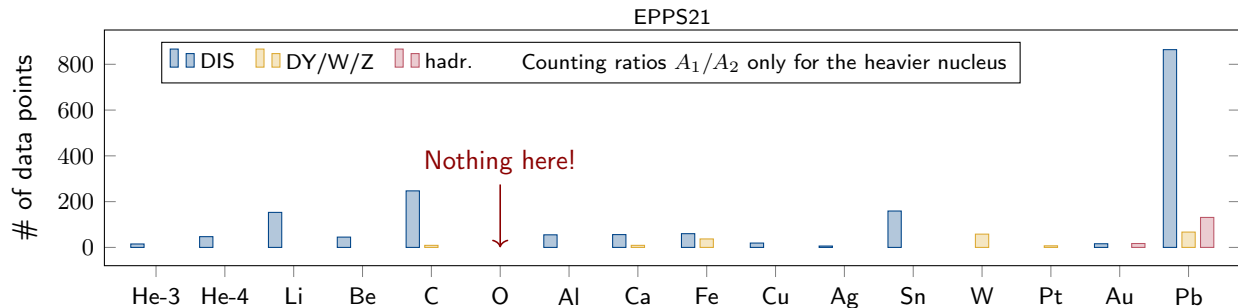
Flett et al., PRD 101 (2020) 094011



Eskola et al., PRC 106 (2022) 035202



Data availability w.r.t. A



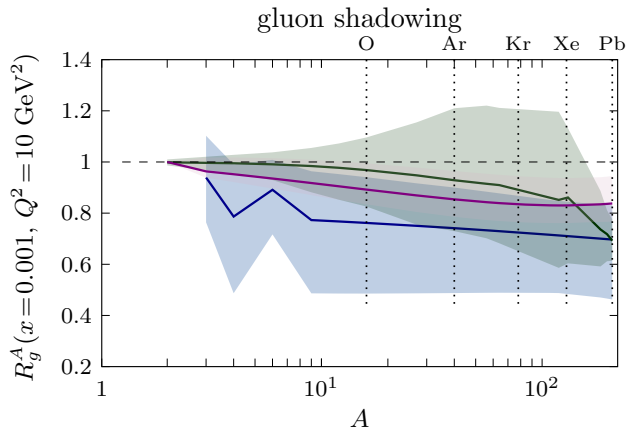
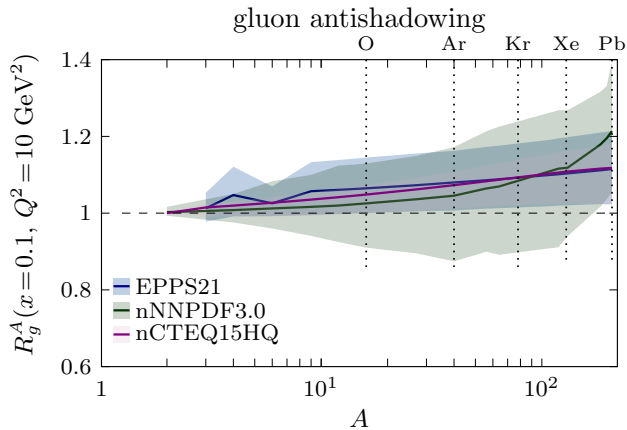
$\sim 50\%$ of the data points are for Pb!

- 😊 Good coverage of DIS measurements for different A (but only fixed target!)
- 😐 DY data more scarce, but OK A coverage
- 😞 Hadronic observables available only for heavy nuclei!

Light-ion runs at LHC could:

- Complement other light-nuclei DY data with W and Z production (strangeness!)
- Give first direct constraints (e.g. dijets, D-mesons) on light-nuclei small- x gluon distributions!

A-dependence of nuclear modifications

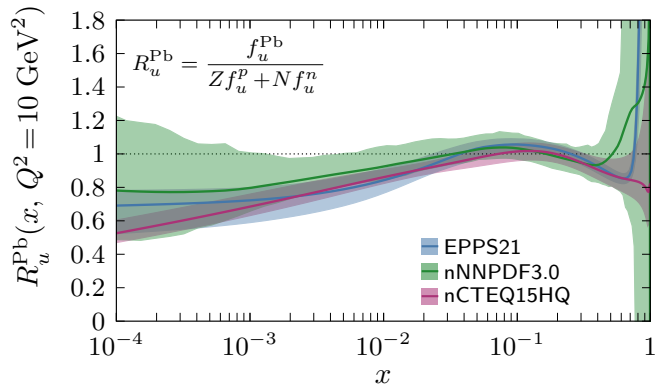
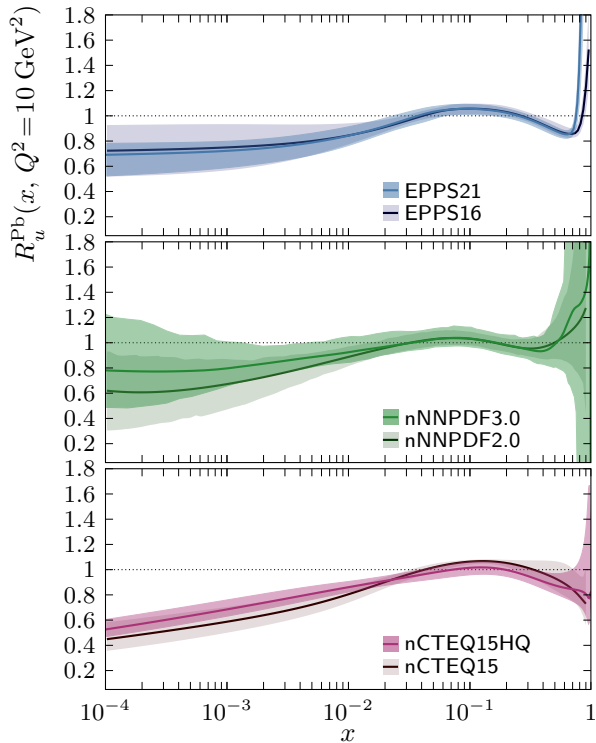


A-dependence of gluon PDFs not well constrained by data!

- Having data for even one additional nucleus would help interpolating the effect for others (but note that A -dependence is not necessarily smooth or even monotonous)
- nPDFs a major source of uncertainty for testing existence of QGP in small systems

Huss et al., PRL 126 (2021) 192301
Brewer et al., PRD 105 (2022) 074040

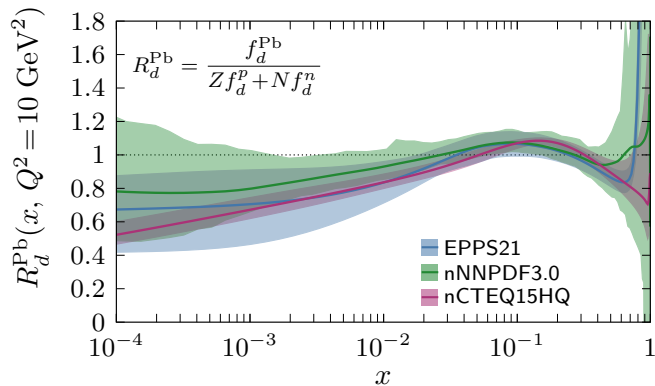
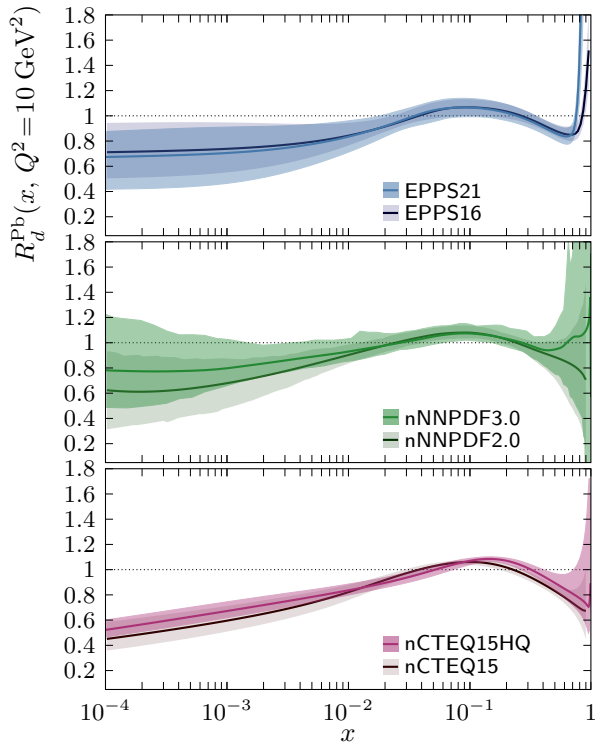
Impact on nPDFs – up



Valence region constrained by fixed-target DIS data

Uncertainties grow towards small x due to lack of collider DIS data

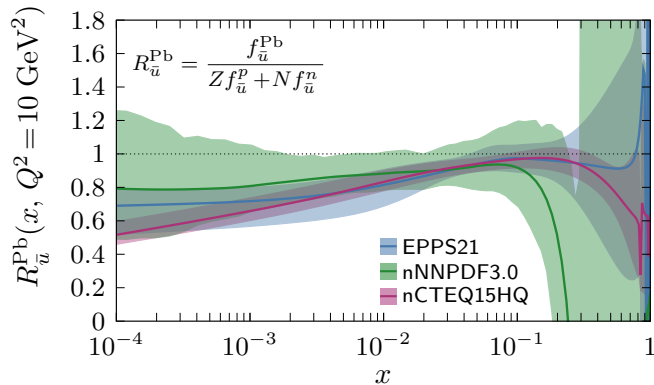
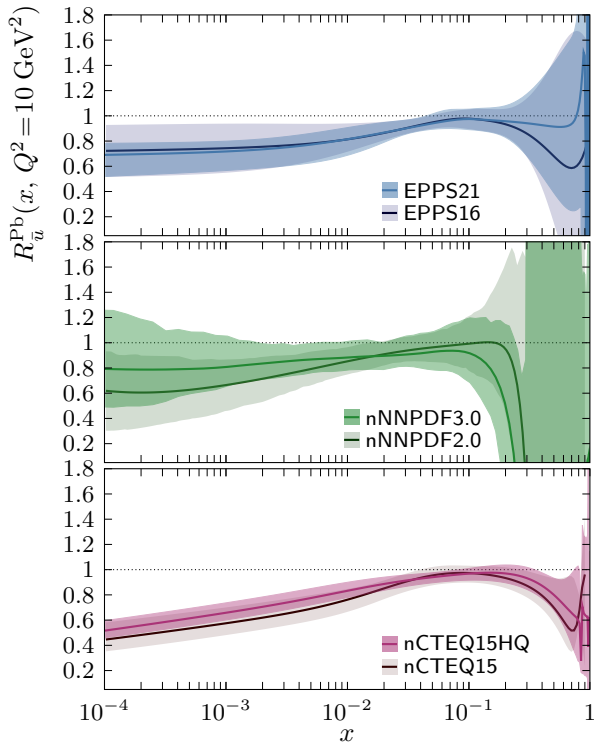
Impact on nPDFs – down



Valence region constrained by fixed-target DIS data

Uncertainties grow towards small x due to lack of collider DIS data

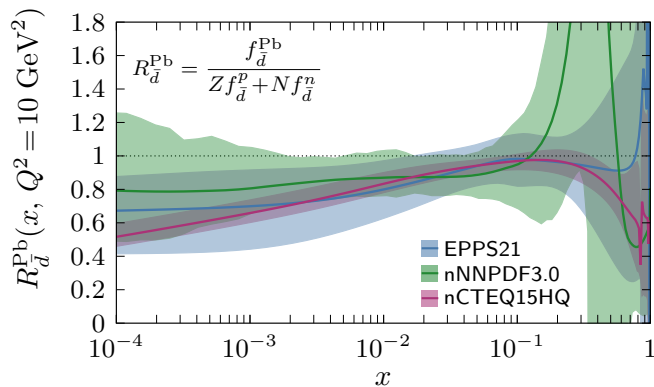
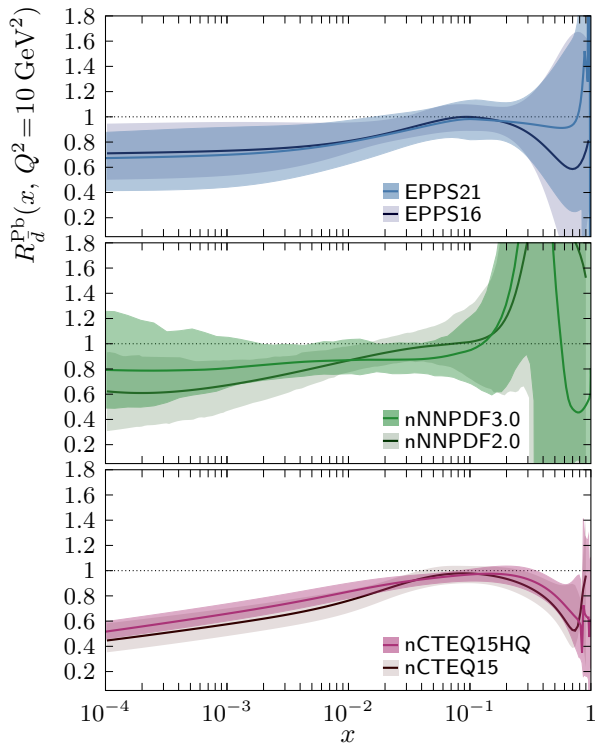
Impact on nPDFs – *anti-up*



Valence / sea quark separation dominated by fixed-target DY data

nNNPDF3.0 has larger uncertainties due to not including these data

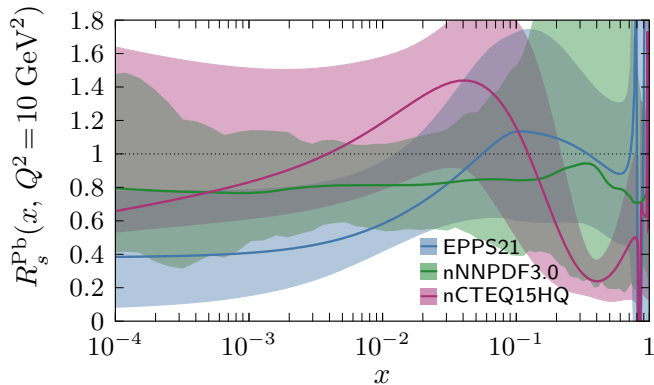
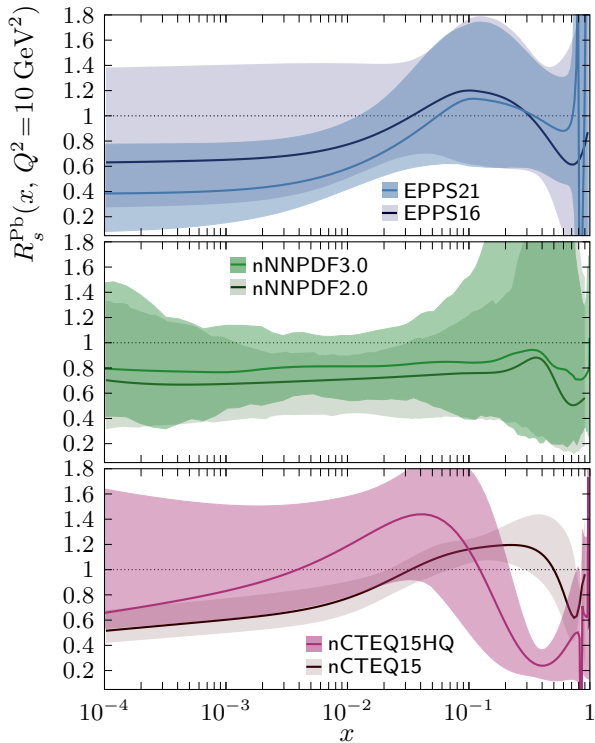
Impact on nPDFs – *anti-down*



Valence / sea quark separation dominated by fixed-target DY data

nNNPDF3.0 has larger uncertainties due to not including these data

Impact on nPDFs – *strange*



Strangeness poorly known in lack of direct constraints

Dimuon process in neutrino-DIS could be used for improved constraints

[Helenius, Paukkunen, Yrjänheikki, JHEP 09 \(2024\) 043](#)