



#### **Flavor hierarchy of parton energy loss in quark-gluon plasma from a Bayesian analysis**

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#### Jets are versatile probes of QGP



**Jet quenching: energy loss, broadening, substructure, medium response**

## Evidences for jet quenching



#### General framework for jet quenching study



**pQCD** factorization: Large-p<sub>T</sub> processes may be factorized into long-distance pieces in **terms of PDF & FF, and short-distance parts describing hard interactions of partons.**

#### General framework for jet quenching study



## Jet-medium interaction



**Bjorken 1982; Bratten, Thoma 1991; Thoma, Gyulassy, 1991; Mustafa, Thoma 2005; Peigne, Peshier, 2006; Djordjevic, 2006; Wicks et al (DGLV), 2007; GYQ et al (AMY), 2008; …**

**BDMPS-Z**: Baier-Dokshitzer-Mueller-Peigne-Schiff-Zakharov **ASW**: Amesto-Salgado-Wiedemann **AMY**: Arnold-Moore-Yaffe (& Caron-Huot, Gale) **GLV**: Gyulassy-Levai-Vitev (& Djordjevic, Heinz) **HT**: Wang-Guo (& Zhang, Wang, Majumder & GYQ, Zhang, Hou)

## Flavor hierachy of parton energy loss



**He, Luo, Wang, Zhu, PRC 2015; Cao, Luo, GYQ, Wang, PRC 2016 ; PLB 2018; etc.**

#### Flavor hierarchy of high  $p_T$  hadron suppression



Combination of NLO-pQCD + LBT + Hydro can explain the flavor hierarchy of  $R_{AA}$ .

**W. J. Xing, GYQ, S. Cao, H. Xing, PLB 2022**

#### Constrain E-Loss using data-driven method

**The theoretical framework:**

$$
\frac{d\sigma_{\rm pp \to hX}}{dp_{\rm T}^h} = \sum_{j} \int dp_{\rm T}^j dz \frac{d\hat{\sigma}_{\rm pp \to jX}}{dp_{\rm T}^j} (p_{\rm T}^j) D_{j \to h}(z) \delta \left( p_{\rm T}^h - z p_{\rm T}^j \right)
$$
\n
$$
\frac{1}{\langle N_{\rm coll} \rangle} \frac{d\sigma_{\rm AA \to hX}}{dp_{\rm T}^h} = \sum_{j} \int dp_{\rm T}^j dx dz \frac{d\hat{\sigma}_{\rm p'p' \to jX}}{dp_{\rm T}^j} (p_{\rm T}^j) W_{\rm AA}(x) D_{j \to h}(z) \delta \left( p_{\rm T}^h - z (p_{\rm T}^j - x \langle \Delta p_{\rm T}^j \rangle) \right)
$$

⟨∆ ⟩**is the average energy loss for parton j, is the energy loss distribution**  with  $x = \Delta p_T^j / \langle \Delta p_T^j \rangle$ .

#### Some works on extracting  $\langle \Delta p_T \rangle$  and  $W_{AA}(x)$

- **F. Arleo, PRL 2017**
	- Take  $W_{AA}(x)$  from BDMPS medium-induced gluon spectrum, & extract parton  $\langle \Delta p_T^{} \rangle$  from  $R_{AA}^{}$  data on  $h^\pm$ ,  $D$  ,  $J/\psi.$
- **He, Pang, Wang, PRL 2019**
	- Use a general ansatz of jet  $W_{AA}(x)$ , and extract the flavor-averaged jet  $\langle \Delta p_T \rangle$  and  $W_{AA}(x)$  from single inclusive jet & y-jet data.
- **Zhang, Liao, GYQ, Wang, Xing, Sci. Bull. 2023** 
	- Extract gluon & charm quark  $\langle \Delta p_T \rangle$  and  $W_{AA}(x)$  for from  $J/\psi R_{AA}$ data.
- **This work (Xing, Cao, GYQ, PLB 2024) :** 
	- Perform a simultaneous analysis on **parton** ∆ for all parton species (g, q, c, and b) from light & heavy flavor hadron  $R_{AA}$  data.



### Details about the analysis

• **The formula for hadron production in AA collisions:**

$$
\frac{1}{\langle N_\text{coll}\rangle}\frac{d\sigma_\text{AA\to hX}}{dp_\text{T}^h} = \sum_j \int dp_\text{T}^j dxdz \frac{d\hat{\sigma}_\text{p'p'\to jX}}{dp_\text{T}^j} (p_\text{T}^j) W_\text{AA}(x) D_{j\to h}(z) \delta\left(p_\text{T}^h - z(p_\text{T}^j - x \langle \Delta p_\text{T}^j \rangle) \right)
$$

• **Parameterize**  $p_T$ -dependence of  $\langle \Delta p_T \rangle$  for gluons  $(g)$ , light quarks  $(g)$ , charm **quarks () and bottom quarks() as:**

$$
\left\langle \Delta p_{\rm T}^{j}\right\rangle =C_{j}\beta_{g}p_{\rm T}^{\gamma}{\rm log}(p_{\rm T})
$$

 $\bm{C}_{{\bm g}} = \bm{1}$  and  $\bm{C}_{{\bm q}}$ ,  $\bm{C}_{{\bm c}}$ ,  $\bm{C}_{{\bm b}}$  represents the  $\langle \Delta \bm{p}_{\bm T} \rangle$  ratio relative to gluon's.

• The parton energy loss distribution  $W_{AA}(x)$  is taken as:

$$
W_{\text{AA}}(x) = \frac{\alpha^{\alpha} x^{\alpha - 1} e^{-\alpha x}}{\Gamma(\alpha)}
$$

• The parameter set  $\theta = (\beta_g, C_q, C_c, C_b, \gamma, \alpha)$  is to be calibrated.

# Bayesian analysis



### Posterior distributions of parameters





**The energy loss parameters for jet-medium interaction can be well constrained by the Bayesian analysis.**

**Reducing experimental data error bars can improve the precision of the extracted parameters.** 

# Prior and posterior R<sub>AA</sub>



**Our calibrated model calculation provides a simultaneous description on**  $R_{AA}$  **data** of charged hadrons, D mesons and B-decayed  $J/\psi$  measured by CMS.

## Flavor hierarchy of parton energy loss



**Direct extraction of the flavor dependence of parton energy loss in QGP from data. Provides a stringent test of pQCD calculation of parton-medium interaction.**

**W. J. Xing, S. Cao, GYQ, PLB 2024**

# Summary

- **Based on a NLO-pQCD calculation of parton production, a general ansatz for parton E-loss function & parton FF, we calculate R<sub>AA</sub> for both heavy & light flavor** hadrons over a wide  $p_T$  range.
- **Using a Bayesian model-to-data analysis, we perform first simultaneous extraction of E-loss of gluons, light quarks, charm quarks & bottom quarks inside the QGP.**
- **The extracted parton E-loss inside the QGP exhibits a clear flavor hierarchy:**   $\langle \Delta E_a \rangle > \langle \Delta E_a \rangle$  ~ $\langle \Delta E_c \rangle > \langle \Delta E_b \rangle$ , consistent with the pQCD expectation.
- **More data and more precise data can improve the precision for the extracted parton E-loss, providing better constraint on theoretical models.**





#### The 9th International Symposium on Heavy Flavor **Production in Hadron and Nuclear Collisions** (HF-HNC 2024)

#### Guangzhou, China, December 6-11, 2024

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## Linear Boltzmann Transport (LBT) Model

- $p_1 \cdot \partial f_1(x_1, p_1) = E_1 C[f_1]$ • **Boltzmann equation:**  $\Gamma_{12\rightarrow 34} = \frac{\gamma_2}{2E_1} \int \frac{d^3p_2}{(2\pi)^3 2E_2} \int \frac{d^3p_3}{(2\pi)^3 2E_2} \int \frac{d^3p_4}{(2\pi)^3 2E_4}$ • **Elastic collisions:**  $\times f_2(\vec{p}_2) \left[1 \pm f_3(\vec{p}_1 - \vec{k})\right] \left[1 \pm f_4(\vec{p}_2 + \vec{k})\right]$  $\times (2\pi)^4 \delta^{(4)}(p_1+p_2-p_3-p_4) |\mathcal{M}_{12\rightarrow 34}|^2$  $= 1 - e^{-\Gamma_{el}\Delta t}$  $t_{el} = 1 - e^{-\Gamma_{el}\Delta t}$  $P_{el} = 1 - e^{-\Gamma_{el}\Delta t}$ Matrix elements taken from LO pQCD
- **Inelastic collisions:**

$$
\langle N_g \rangle = \ \mathbf{\Gamma}_g \mathbf{\Delta}\, t \ = \Delta t \int dx dk_\perp^2 \frac{dN_g}{dx dk_\perp^2 dt}.
$$

Medium-induced radiation spectra taken from HT: Guo, Wang PRL 2000; Zhang, Wang, Wang, PRL 2004; Zhang, Hou, GYQ, PRC 2019; Zhang, GYQ, Wang, PRD 2019.

• **Elastic + Inelastic:**  $el$  inel  $el$  inel  $P_{tot} = 1 - e^{-\Gamma_{tot} \Delta t} = P_{el} + P_{inel} - P_{el} P_{inel}$ 

 $P_{inel} = 1 - e^{-\langle N_g \rangle}$ 

**He, Luo, Wang, Zhu, PRC 2015; Cao, Luo, GYQ, Wang, PRC 2016, PLB 2018; etc.**

# Closure test





**Start with a pre-set point**  $\theta_0$  **to** calculate  $R_{AA}$  of hadrons, which serves **as the middle points of pseudo-data (the error bars are taken from the experiments data).**

**Confirm that the posterior distributions from Baysian analysis do**  agree with the pre-set value  $\theta_0$ .

**Halving the error bars of the pseudo-data can improve the precision of the extracted parameters.**