

# Exploring light flavor hadronization in hard and soft events with event shape classifiers in small collision systems at the LHC with ALICE

Feng Fan, for the ALICE Collaboration Central China Normal University





## sQGP in small collision systems?



✓ Strangeness enhancement in high-multiplicity (HM) small collisions systems



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### **Selection bias**



 $\checkmark$  Selection bias towards hard pp collisions  $\rightarrow$  selecting multiplicity classes and measuring particle spectra in the same pseudorapidity

- ✓ We need new observables to isolate events with specific topologic characteristics
  - Transverse spherocity  $S_{\Omega}^{p_{T}=1}$
  - Relative transverse activity  $R_{\rm T}$
  - Charged-particle flattenicity  $1 \rho$







## The ALICE detector in Run 1 and Run 2

111111111

### TPC ( $|\eta| < 0.9$ ): tracking, PID (d*E*/d*x*)

TOF ( $|\eta| < 0.9$ ): PID (particle velocity)

### ITS ( $|\eta| < 0.9$ ): tracking, vertexing, pile up rejection, SPD tracklets estimator





Transverse spherocity  $S_{O}^{p_{T}=1}$  measurement



# Transverse spherocity $S_{\Omega}^{p_{\rm T}=1}$

Event-by-event selection quantified the topology in the azimuthal plane

$$S_{O}^{p_{T}=1} = \frac{\pi^{2}}{4} \min_{\hat{n}} \left( \frac{\sum_{i} |\hat{p}_{T,i} \times \hat{n}|}{N_{\text{trks}}} \right)^{2}$$

- $\hat{p}_{T,i}$  is the transverse momentum unit vector for a charged particle
- $N_{\rm trks}$  is the number of charged particles in a given event
- $\hat{n}$  is the unit vector that minimizes  $S_{\Omega}^{p_{\rm T}=1}$
- ✓  $S_{O}^{p_{T}=1} \rightarrow 0$ : jet-like event, particle production driven by hard physics
- ✓  $S_{O}^{p_{T}=1} \rightarrow 1$ : isotropic event, particle production driven by soft physics



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# -differential average transverse momentum $\langle p_{\rm T} \rangle$

![](_page_10_Figure_1.jpeg)

![](_page_10_Picture_5.jpeg)

![](_page_10_Figure_6.jpeg)

![](_page_10_Figure_7.jpeg)

![](_page_10_Figure_8.jpeg)

## -differential average transverse momentum $\langle p_{\rm T} \rangle$

![](_page_11_Figure_1.jpeg)

![](_page_11_Figure_2.jpeg)

5

# Integrated yield as a function of $S_{\Omega}^{p_{\rm T}=1}$

- Strangeness enhancement in isotropic events
- Strangeness-based ordering
- Proton remains mostly unaffected (S=0)

![](_page_12_Figure_7.jpeg)

EXALICE, JHEP 05 (2024) 184 ALI-PUB-574667

![](_page_12_Picture_9.jpeg)

![](_page_12_Picture_10.jpeg)

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PYTHIA 8.2 predictions

- Ropes tune qualitatively reproduce the trends, but not the strangeness ordering
- Monash tune cannot capture the trends

![](_page_13_Figure_10.jpeg)

EXALICE, JHEP 05 (2024) 184 ALI-PUB-574667

![](_page_13_Picture_12.jpeg)

![](_page_13_Picture_13.jpeg)

# Integrated yield as a function of $S_{0}^{p_{T}=1}$

- Strangeness enhancement in isotropic events
- Strangeness-based ordering
- Proton remains mostly unaffected (S=0)

PYTHIA 8.2 predictions

- Ropes tune qualitatively reproduce the trends, but not the strangeness ordering
- Monash tune cannot capture the trends
- EPOS-LHC qualitatively reproduced the trends, but not the strangeness ordering
- Herwig 7.2 cannot capture the trends

![](_page_14_Figure_12.jpeg)

ALI-PUB-574672 📚 ALICE, JHEP 05 (2024) 184

![](_page_14_Picture_15.jpeg)

![](_page_14_Picture_16.jpeg)

# Relative transverse activity $R_{\rm T}$ measurement

![](_page_15_Picture_1.jpeg)

## Relative transverse activity $R_{\rm T}$

Event-by-event selection based on the underlying-event (UE) activity in the midrapidity interval (UE refers to everything that does not come from the main hard partonic scattering)

- $R_{\rm T} = N_{\rm ch}^{\rm TS} / \langle N_{\rm ch}^{\rm TS} \rangle$
- $N_{\rm ch}^{\rm TS}$  is the charged-particle multiplicity in the transverse region (TS)
- $\langle N_{ch}^{TS} \rangle$  is the average multiplicity over all events in TS

![](_page_16_Figure_5.jpeg)

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![](_page_16_Figure_7.jpeg)

might be attributed to the initial- and final-state radiations

### $p_{\rm T}$ -spectra as a function of $R_{\rm T}$

#### ➢ ALICE, JHEP 01 (2024) 056

![](_page_17_Figure_2.jpeg)

Toward and Away regions:

- For  $p_{\rm T}$  < 4 GeV/c, the  $p_{\rm T}$  spectra is  $R_{\rm T}$  dependent
- For  $p_{\rm T}$  > 4 GeV/c, the spectral shape is almost  $R_{\rm T}$  independent

![](_page_17_Picture_7.jpeg)

![](_page_17_Figure_8.jpeg)

### $p_{\rm T}$ -spectra as a function of $R_{\rm T}$

▲ ALICE, JHEP 01 (2024) 056

![](_page_18_Figure_2.jpeg)

Toward and Away regions:

- For  $p_{\rm T}$  < 4 GeV/c, the  $p_{\rm T}$  spectra is  $R_{\rm T}$ dependent
- For  $p_{\rm T}$  > 4 GeV/*c*, the spectral shape is almost  $R_{\rm T}$  independent

Transverse region:

- A  $p_{\rm T}$ -hardening with increasing  $R_{\rm T}$ , due to autocorrelation effects

Phys. Rev. D 104 (2021) 016017

![](_page_18_Picture_10.jpeg)

![](_page_18_Picture_11.jpeg)

![](_page_18_Picture_12.jpeg)

### $p_{\rm T}$ -spectra as a function of $R_{\rm T}$

#### ► ALICE, JHEP 01 (2024) 056

![](_page_19_Figure_2.jpeg)

Toward and Away regions:

- For  $p_{\rm T}$  < 4 GeV/*c*, the  $p_{\rm T}$  spectra is  $R_{\rm T}$ dependent
- For  $p_{\rm T}$  > 4 GeV/*c*, the spectral shape is almost  $R_{\rm T}$  independent

Transverse region:

- A  $p_{\rm T}$ -hardening with increasing  $R_{\rm T}$ , due to autocorrelation effects
  - Phys. Rev. D 104 (2021) 016017

MC predictions:

- PYTHIA 8.2 describes data better than EPOS LHC

![](_page_19_Picture_12.jpeg)

![](_page_19_Picture_13.jpeg)

![](_page_19_Picture_14.jpeg)

## $\langle p_{\rm T} \rangle$ as a function of $R_{\rm T}$

![](_page_20_Figure_1.jpeg)

**ALI-PUB-567949** 

Low  $R_{\rm T}$ : the  $\langle p_{\rm T} \rangle$  is independent of collision system for  $R_{\rm T} \approx 0$  as jet dominates at low  $R_{\rm T}$ 

#### ▲ ALICE, JHEP 01 (2024) 056

![](_page_20_Picture_6.jpeg)

![](_page_20_Picture_7.jpeg)

## $\langle p_{\rm T} \rangle$ as a function of $R_{\rm T}$

![](_page_21_Figure_1.jpeg)

High  $R_{\rm T}$ :

- Toward and Away: the  $\langle p_{\rm T} \rangle$  is nearly flat for  $R_{\rm T} > 1$ , and exhibits a system size ordering

#### EXALICE, JHEP 01 (2024) 056

Low  $R_{\rm T}$ : the  $\langle p_{\rm T} \rangle$  is independent of collision system for  $R_{\rm T} \approx 0$  as jet dominates at low  $R_{\rm T}$ 

![](_page_21_Picture_7.jpeg)

![](_page_21_Picture_8.jpeg)

## $\langle p_{\rm T} \rangle$ as a function of $R_{\rm T}$

![](_page_22_Figure_1.jpeg)

**ALI-PUB-567949** 

Low  $R_{\rm T}$ : the  $\langle p_{\rm T} \rangle$  is independent of collision system for  $R_{\rm T} \approx 0$  as jet dominates at low  $R_{\rm T}$ High  $R_{\rm T}$ :

- Transverse: the  $\langle p_{\rm T} \rangle$  is increasing with increasing  $R_{\rm T}$

#### ▲ ALICE, JHEP 01 (2024) 056

- Toward and Away: the  $\langle p_{\rm T} \rangle$  is nearly flat for  $R_{\rm T} > 1$ , and exhibits a system size ordering

![](_page_22_Picture_9.jpeg)

![](_page_22_Picture_10.jpeg)

# Charged-particle flattenicity 1- $\rho$ measurement

![](_page_23_Picture_1.jpeg)

![](_page_23_Picture_2.jpeg)

## Charged-particle flattenicity 1-p

Event-by-event selection based on the relative standard deviation of the multiplicity measured in the 64 V0 channels

$$\rho = \frac{\sqrt{\sum_{i=1}^{64} \left( N_{\rm ch}^{\rm cell,i} - \langle N_{\rm ch}^{\rm cell} \rangle \right) / N_{\rm cell}^2}}{\langle N_{\rm ch}^{\rm cell} \rangle}$$

- $N_{\rm ch}^{\rm cell,i}$  is the particle multiplicity in the *i*-th cell
- $\langle N_{\rm ch}^{\rm cell} \rangle$  is the average multiplicity over the all 64 cells per event

- ✓ large flattenicity  $1 \rho \rightarrow 0$ : jet-like events, with low  $N_{\rm mpi}$
- ✓ small flattenicity  $1 \rho \rightarrow 1$ : isotropic events, with high  $N_{\rm mpi}$

10

![](_page_24_Figure_8.jpeg)

## $\mathcal{Q}_{\mathrm{pp}}$ as a function of $p_{\mathrm{T}}$

![](_page_25_Figure_1.jpeg)

Intermediate  $p_{\rm T}$ : a bump structure develops with increasing multiplicity -----

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![](_page_25_Picture_5.jpeg)

## $\mathcal{Q}_{\mathrm{pp}}$ as a function of $p_{\mathrm{T}}$

![](_page_26_Figure_1.jpeg)

- Intermediate  $p_{\rm T}$ : a bump structure develops with increasing multiplicity -
- High  $p_{\rm T}$ :  $Q_{\rm pp}$  approaches unity

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![](_page_26_Picture_6.jpeg)

![](_page_27_Figure_3.jpeg)

- Intermediate  $p_{\rm T}$ : a bump structure develops with increasing multiplicity
- High  $p_{\rm T}$ :  $Q_{\rm pp}$  approaches unity

# $Q_{pp}$ : data vs MC models

![](_page_28_Figure_1.jpeg)

- PYTHIA 8 w/o CR: a nearly flat  $Q_{pp}$  as a function of  $p_{\rm T}$
- PYTHIA 8 w CR: overall the best description of data
- —

EPOS LHC: overestimates and underestimates  $Q_{pp}$  at intermediate and high  $p_{T}$  values, respectively

![](_page_28_Picture_7.jpeg)

### Particle ratios: flattenicity vs multiplicity

![](_page_29_Figure_1.jpeg)

The particle ratios as a function of flattenicity exhibit a steeper increase with multiplicity than those as a function of VOM

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### ➢ arXiv:2407.20037

Flattenicity (1-p) only: flattenicity selection based on unbiased events (0-100% VOM)

0-1% VOM:

flattenicity selection based on high-multiplicity events (0-1% VOM)

![](_page_29_Picture_10.jpeg)

## Summary

- 1. Transverse spherocity  $S_{\Omega}^{p_{T}=1}$ 
  - $S_{O}^{p_{T}=1}$  can be used to select strangeness enhanced/suppressed events
  - Strangeness enhancement in high-multiplicity pp collisions is a feature of the isotropic events -
- 2. Relative transverse activity  $R_{\rm T}$ 
  - $R_{\rm T} > 3$  emerges bias towards multi-jet topologies
  - $p_{\rm T}$ -spectra are hardening with increasing  $R_{\rm T}$  due to autocorrelation effects
- 3. Charged particle flattenicity  $1 \rho$  can be used to select pp collisions with large number of MPI and small bias than the VOM multiplicity estimator

![](_page_30_Picture_10.jpeg)

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Outlook: measurement of  $\phi$  meson production in and out of jets for pp collisions at  $\sqrt{s} = 13.6$  TeV is ongoing

![](_page_31_Figure_12.jpeg)

![](_page_31_Picture_13.jpeg)