

# Exploring light flavor hadronization in hard and soft events with event shape classifiers in small collision systems at the LHC with ALICE

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## sQGP in small collision systems?



✓ Strangeness enhancement in high-multiplicity (HM) small collisions systems



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### **Selection bias**



 $\checkmark$  Selection bias towards hard pp collisions  $\rightarrow$  selecting multiplicity classes and measuring particle spectra in the same pseudorapidity

- ✓ We need new observables to isolate events with specific topologic characteristics
  - Transverse spherocity  $S_{\Omega}^{p_{T}=1}$
  - Relative transverse activity  $R_{\rm T}$
  - Charged-particle flattenicity  $1 \rho$







## The ALICE detector in Run 1 and Run 2

111111111

### TPC ( $|\eta| < 0.9$ ): tracking, PID (d*E*/d*x*)

TOF ( $|\eta| < 0.9$ ): PID (particle velocity)

### ITS ( $|\eta| < 0.9$ ): tracking, vertexing, pile up rejection, SPD tracklets estimator





Transverse spherocity  $S_{O}^{p_{T}=1}$  measurement



# Transverse spherocity $S_{\Omega}^{p_{\rm T}=1}$

Event-by-event selection quantified the topology in the azimuthal plane

$$S_{O}^{p_{T}=1} = \frac{\pi^{2}}{4} \min_{\hat{n}} \left( \frac{\sum_{i} |\hat{p}_{T,i} \times \hat{n}|}{N_{\text{trks}}} \right)^{2}$$

- $\hat{p}_{T,i}$  is the transverse momentum unit vector for a charged particle
- $N_{\rm trks}$  is the number of charged particles in a given event
- $\hat{n}$  is the unit vector that minimizes  $S_{\Omega}^{p_{\rm T}=1}$
- ✓  $S_{O}^{p_{T}=1} \rightarrow 0$ : jet-like event, particle production driven by hard physics
- ✓  $S_{O}^{p_{T}=1} \rightarrow 1$ : isotropic event, particle production driven by soft physics

![](_page_7_Figure_8.jpeg)

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![](_page_7_Picture_12.jpeg)

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![](_page_8_Figure_8.jpeg)

![](_page_8_Figure_9.jpeg)

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![](_page_9_Figure_8.jpeg)

![](_page_9_Figure_9.jpeg)

# -differential average transverse momentum $\langle p_{\rm T} \rangle$

![](_page_10_Figure_1.jpeg)

![](_page_10_Picture_5.jpeg)

![](_page_10_Figure_6.jpeg)

![](_page_10_Figure_7.jpeg)

![](_page_10_Figure_8.jpeg)

## -differential average transverse momentum $\langle p_{\rm T} \rangle$

![](_page_11_Figure_1.jpeg)

![](_page_11_Figure_2.jpeg)

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# Integrated yield as a function of $S_{\Omega}^{p_{\rm T}=1}$

- Strangeness enhancement in isotropic events
- Strangeness-based ordering
- Proton remains mostly unaffected (S=0)

![](_page_12_Figure_7.jpeg)

EXALICE, JHEP 05 (2024) 184 ALI-PUB-574667

![](_page_12_Picture_9.jpeg)

![](_page_12_Picture_10.jpeg)

# Integrated yield as a function of $S_{\Omega}^{p_{\rm T}=1}$

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PYTHIA 8.2 predictions

- Ropes tune qualitatively reproduce the trends, but not the strangeness ordering
- Monash tune cannot capture the trends

![](_page_13_Figure_10.jpeg)

EXALICE, JHEP 05 (2024) 184 ALI-PUB-574667

![](_page_13_Picture_12.jpeg)

![](_page_13_Picture_13.jpeg)

# Integrated yield as a function of $S_{0}^{p_{T}=1}$

- Strangeness enhancement in isotropic events
- Strangeness-based ordering
- Proton remains mostly unaffected (S=0)

PYTHIA 8.2 predictions

- Ropes tune qualitatively reproduce the trends, but not the strangeness ordering
- Monash tune cannot capture the trends
- EPOS-LHC qualitatively reproduced the trends, but not the strangeness ordering
- Herwig 7.2 cannot capture the trends

![](_page_14_Figure_12.jpeg)

ALI-PUB-574672 📚 ALICE, JHEP 05 (2024) 184

![](_page_14_Picture_15.jpeg)

![](_page_14_Picture_16.jpeg)

# Relative transverse activity $R_{\rm T}$ measurement

![](_page_15_Picture_1.jpeg)

## Relative transverse activity $R_{\rm T}$

Event-by-event selection based on the underlying-event (UE) activity in the midrapidity interval (UE refers to everything that does not come from the main hard partonic scattering)

- $R_{\rm T} = N_{\rm ch}^{\rm TS} / \langle N_{\rm ch}^{\rm TS} \rangle$
- $N_{\rm ch}^{\rm TS}$  is the charged-particle multiplicity in the transverse region (TS)
- $\langle N_{ch}^{TS} \rangle$  is the average multiplicity over all events in TS

![](_page_16_Figure_5.jpeg)

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![](_page_16_Figure_7.jpeg)

might be attributed to the initial- and final-state radiations

### $p_{\rm T}$ -spectra as a function of $R_{\rm T}$

#### ➢ ALICE, JHEP 01 (2024) 056

![](_page_17_Figure_2.jpeg)

Toward and Away regions:

- For  $p_{\rm T}$  < 4 GeV/c, the  $p_{\rm T}$  spectra is  $R_{\rm T}$  dependent
- For  $p_{\rm T}$  > 4 GeV/c, the spectral shape is almost  $R_{\rm T}$  independent

![](_page_17_Picture_7.jpeg)

![](_page_17_Figure_8.jpeg)

### $p_{\rm T}$ -spectra as a function of $R_{\rm T}$

▲ ALICE, JHEP 01 (2024) 056

![](_page_18_Figure_2.jpeg)

Toward and Away regions:

- For  $p_{\rm T}$  < 4 GeV/c, the  $p_{\rm T}$  spectra is  $R_{\rm T}$ dependent
- For  $p_{\rm T}$  > 4 GeV/*c*, the spectral shape is almost  $R_{\rm T}$  independent

Transverse region:

- A  $p_{\rm T}$ -hardening with increasing  $R_{\rm T}$ , due to autocorrelation effects

Phys. Rev. D 104 (2021) 016017

![](_page_18_Picture_10.jpeg)

![](_page_18_Picture_11.jpeg)

![](_page_18_Picture_12.jpeg)

### $p_{\rm T}$ -spectra as a function of $R_{\rm T}$

#### ► ALICE, JHEP 01 (2024) 056

![](_page_19_Figure_2.jpeg)

Toward and Away regions:

- For  $p_{\rm T}$  < 4 GeV/*c*, the  $p_{\rm T}$  spectra is  $R_{\rm T}$ dependent
- For  $p_{\rm T}$  > 4 GeV/*c*, the spectral shape is almost  $R_{\rm T}$  independent

Transverse region:

- A  $p_{\rm T}$ -hardening with increasing  $R_{\rm T}$ , due to autocorrelation effects
  - Phys. Rev. D 104 (2021) 016017

MC predictions:

- PYTHIA 8.2 describes data better than EPOS LHC

![](_page_19_Picture_12.jpeg)

![](_page_19_Picture_13.jpeg)

![](_page_19_Picture_14.jpeg)

## $\langle p_{\rm T} \rangle$ as a function of $R_{\rm T}$

![](_page_20_Figure_1.jpeg)

**ALI-PUB-567949** 

Low  $R_{\rm T}$ : the  $\langle p_{\rm T} \rangle$  is independent of collision system for  $R_{\rm T} \approx 0$  as jet dominates at low  $R_{\rm T}$ 

#### ▲ ALICE, JHEP 01 (2024) 056

![](_page_20_Picture_6.jpeg)

![](_page_20_Picture_7.jpeg)

## $\langle p_{\rm T} \rangle$ as a function of $R_{\rm T}$

![](_page_21_Figure_1.jpeg)

High  $R_{\rm T}$ :

- Toward and Away: the  $\langle p_{\rm T} \rangle$  is nearly flat for  $R_{\rm T} > 1$ , and exhibits a system size ordering

#### EXALICE, JHEP 01 (2024) 056

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![](_page_21_Picture_7.jpeg)

![](_page_21_Picture_8.jpeg)

## $\langle p_{\rm T} \rangle$ as a function of $R_{\rm T}$

![](_page_22_Figure_1.jpeg)

**ALI-PUB-567949** 

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- Transverse: the  $\langle p_{\rm T} \rangle$  is increasing with increasing  $R_{\rm T}$

#### ▲ ALICE, JHEP 01 (2024) 056

- Toward and Away: the  $\langle p_{\rm T} \rangle$  is nearly flat for  $R_{\rm T} > 1$ , and exhibits a system size ordering

![](_page_22_Picture_9.jpeg)

![](_page_22_Picture_10.jpeg)

# Charged-particle flattenicity 1- $\rho$ measurement

![](_page_23_Picture_1.jpeg)

![](_page_23_Picture_2.jpeg)

## Charged-particle flattenicity 1-p

Event-by-event selection based on the relative standard deviation of the multiplicity measured in the 64 V0 channels

$$\rho = \frac{\sqrt{\sum_{i=1}^{64} \left( N_{\rm ch}^{\rm cell,i} - \langle N_{\rm ch}^{\rm cell} \rangle \right) / N_{\rm cell}^2}}{\langle N_{\rm ch}^{\rm cell} \rangle}$$

- $N_{\rm ch}^{\rm cell,i}$  is the particle multiplicity in the *i*-th cell
- $\langle N_{\rm ch}^{\rm cell} \rangle$  is the average multiplicity over the all 64 cells per event

- ✓ large flattenicity  $1 \rho \rightarrow 0$ : jet-like events, with low  $N_{\rm mpi}$
- ✓ small flattenicity  $1 \rho \rightarrow 1$ : isotropic events, with high  $N_{\rm mpi}$

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![](_page_24_Figure_8.jpeg)

## $\mathcal{Q}_{\mathrm{pp}}$ as a function of $p_{\mathrm{T}}$

![](_page_25_Figure_1.jpeg)

Intermediate  $p_{\rm T}$ : a bump structure develops with increasing multiplicity -----

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![](_page_25_Picture_5.jpeg)

## $\mathcal{Q}_{\mathrm{pp}}$ as a function of $p_{\mathrm{T}}$

![](_page_26_Figure_1.jpeg)

- Intermediate  $p_{\rm T}$ : a bump structure develops with increasing multiplicity -
- High  $p_{\rm T}$ :  $Q_{\rm pp}$  approaches unity

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![](_page_26_Picture_6.jpeg)

![](_page_27_Figure_3.jpeg)

- Intermediate  $p_{\rm T}$ : a bump structure develops with increasing multiplicity
- High  $p_{\rm T}$ :  $Q_{\rm pp}$  approaches unity

# $Q_{pp}$ : data vs MC models

![](_page_28_Figure_1.jpeg)

- PYTHIA 8 w/o CR: a nearly flat  $Q_{pp}$  as a function of  $p_{\rm T}$
- PYTHIA 8 w CR: overall the best description of data
- —

EPOS LHC: overestimates and underestimates  $Q_{pp}$  at intermediate and high  $p_{T}$  values, respectively

![](_page_28_Picture_7.jpeg)

### Particle ratios: flattenicity vs multiplicity

![](_page_29_Figure_1.jpeg)

The particle ratios as a function of flattenicity exhibit a steeper increase with multiplicity than those as a function of VOM

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### ➢ arXiv:2407.20037

Flattenicity (1-p) only: flattenicity selection based on unbiased events (0-100% VOM)

0-1% VOM:

flattenicity selection based on high-multiplicity events (0-1% VOM)

![](_page_29_Picture_10.jpeg)

## Summary

- 1. Transverse spherocity  $S_{\Omega}^{p_{T}=1}$ 
  - $S_{O}^{p_{T}=1}$  can be used to select strangeness enhanced/suppressed events
  - Strangeness enhancement in high-multiplicity pp collisions is a feature of the isotropic events -
- 2. Relative transverse activity  $R_{\rm T}$ 
  - $R_{\rm T} > 3$  emerges bias towards multi-jet topologies
  - $p_{\rm T}$ -spectra are hardening with increasing  $R_{\rm T}$  due to autocorrelation effects
- 3. Charged particle flattenicity  $1 \rho$  can be used to select pp collisions with large number of MPI and small bias than the VOM multiplicity estimator

![](_page_30_Picture_10.jpeg)

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Outlook: measurement of  $\phi$  meson production in and out of jets for pp collisions at  $\sqrt{s} = 13.6$  TeV is ongoing

![](_page_31_Figure_12.jpeg)

![](_page_31_Picture_13.jpeg)