Differential π^0 and photon modification in d+Au collisions [based on nucl-th/2404.17660]



Dennis V. Perepelitsa

23 September 2024

12th International Conference on Hard and EM Probes of **High-Energy Nuclear Collisions**

PHYSICAL REVIEW C 110, L011901 (2024)

Letter

Contribution to differential π^0 and γ_{dir} modification in small systems from color fluctuation effects

Dennis V. Perepelitsa[®]

Department of Physics, University of Colorado Boulder, Boulder, Colorado 80309, USA

(Received 6 May 2024; accepted 20 June 2024; published 1 July 2024)

A major complication in the search for jet quenching in proton- or deuteron-nucleus collision systems is the presence of physical effects which influence the experimental determination of collision centrality in the presence of a hard process. For example, in the proton color fluctuation picture, protons with a large Bjorken-x ($x \ge 0.1$) parton interact more weakly with the nucleons in the nucleus, leading to a smaller (larger) than expected yield in large (small) activity events. A recent measurement by the PHENIX Collaboration compared the yield of neutral pion and direct photon production in d + Au collisions, under the argument that the photon yields correct for such biases, and the difference between the two species is thus attributable to final-state effects (i.e., jet quenching).





Measurement of π^0/γ modification in *d*+Au

T. Hachiya, PHENIX Overview, nucl-ex/2303.12899



- PHENIX measures the double ratio $R_{\rm dAu}(\pi^0)/R_{\rm dAu}(\gamma_{\rm dir})$ in small systems
 - Argument: the normalization to the number of photons takes out all centrality bias effects, and any modification is a final-state effect
- For centrality averaged events, $R_{\rm dAu}(\pi^0)/R_{\rm dAu}(\gamma_{\rm dir}) = 0.92 \pm 0.02 \text{ (stat)} \pm 0.15 \text{ (syst)}$
- In the most central events, $R_{\rm dAu}(\pi^0)/R_{\rm dAu}(\gamma_{\rm dir}) = 0.77 \pm 0.03 \text{ (stat)} \pm 0.13 \text{ (syst)}$



The PHENIX results would require $\delta p_{\rm T}/p_{\rm T} \approx$ 1-2% in 0-5% *d*+Au collisions

At LHC energies, limit of $\delta p_T/p_T < 1.4\%$ at 90% CL (in 0-20% ZDC-selected *p*+Pb events)

3

Well-known "multiplicity" bias

- Increased event activity in the presence of a hard-scattering
 - overestimate hard scattering yields in high-multiplicity events, underestimate them in low-multiplicity events
- This particular bias does not have a strong process (γ_{dir} vs. π^0) or kinematic dependence
 - the PHENIX strategy likely eliminates this particular bias 🔽





ALICE, PRC 91 (2015) 064905 + much work by PHENIX, Steinberg, Morsch, Loizides, others

A different bias for extremely high- x_p process

- For hard-scattering processes with Bjorken-*x* in the proton, $x_p \gtrsim 0.1$:
 - data at RHIC and LHC: there is an anti-correlation between increasing x_p and the backward multiplicity (i.e. in the nucleus-direction)
 - ➡ i.e. relative movement of events from "central" to "peripheral" category
 - \rightarrow observed indirectly as an $R_{CP} < 1$



rare high- x_p proton

M. Alvioli et al PRD 98 (2018) 071502(R)

- One compelling explanation: proton **color fluctuations**
 - \rightarrow protons with a large- x_p parton are quite different than "average" protons, and strike fewer nucleons in the nucleus







Evidence for x_p-UE^{backward} anti-correlation





Strong UE-dependence of the yield of jets, increasing with more forward (higherx_p) selections

Strongly suppressed $R_{\rm CP}$ for jets, increasing with *p*_T



Evidence for x_p-UE^{backward} anti-correlation

- Recent ATLAS measurement of the R_{CP} for dijet events
 - event-level parton kinematics (x_p , x_A , Q^2) estimated via the dijet kinematics
- Modification is (mostly) a universal function just of x_p
 - strongly supportive of color fluctuation picture







Quantifying color fluctuation effects

PHYSICAL REVIEW D 98. 071502(R) (2018)

Rapid Communications

Global analysis of color fluctuation effects in proton– and deuteron–nucleus collisions at RHIC and the LHC

M. Alvioli,¹ L. Frankfurt,^{2,3} D. V. Perepelitsa,⁴ and M. Strikman³

¹Consiglio Nazionale delle Ricerche, Istituto di Ricerca per la Protezione Idrogeologica, via Madonna Alta 126, I-06128 Perugia, Italy ²Tel Aviv University, Tel Aviv, Israel ³104 Davey Lab, The Pennsylvania State University, University Park, Pennsylvania 16803, USA ⁴University of Colorado, Boulder, Colorado 80309 USA



- An analysis of earlier ATLAS & PHENIX data quantitatively extracted the magnitude of the proton's "shrinking"
 - The effective nucleon-nucleon crosssection for a proton depleted by a high x_p parton is $\lambda(x_p) \times \sigma_{NN}$
 - $\rightarrow \lambda(x_p) < 1$ and decreases quickly with x_p
 - I use the results of this Color Fluctuation Model (CFM) with no changes to the parameters or introduction of additional physics





x_p ranges probed by PHENIX measurement



- MC simulation, matching specific PHENIX acceptance & kinematics
- For direct photons, typical x_p values probed are $x_p \sim p_T^{\gamma}/E_{\rm beam} \sim 0.11$
- However, **neutral pions** carry only a fraction of the fragmenting parton's $p_{\rm T}$
 - → $\langle x_p \rangle \approx 0.17$, with long tail to large x_p values (28% from $x_p > 0.2$)
 - they will incur stronger CFM effects

Color fluctuation effect for γ_{dir} vs. π^{\vee}



- I apply the CFM to these x_p distributions and the PHENIX Glauber+NBD centrality model
 - The CFM has no physics to change the minimum-bias R_{dAu} — all results are only relative enhancements/suppressions
- Note the stronger effect for pions (higher $< x_p >$) than direct photons (lower $< x_p >$)
 - These only partially cancel the PHENIX π^0/γ_{dir} observable is sensitive to this x_p -dependent effect





Benchmarking: direct photon R_{dAm}



- Compare the CFM to PHENIX direct photon R_{dAm} data vs. centrality
 - The data has a modest enhancement at low $N_{\rm coll}$, suppression at high $N_{\rm coll}$, and min-bias $R_{\rm dAu} \approx 1$
- The Color Fluctuation Model predicts a modification pattern exactly like this
 - The predicted modifications are modest since x_p is not too high





Comparison to π^0/γ_{dir} **data**



- Compare the CFM to the PHENIX $R_{\rm dAu}(\pi^0)/R_{\rm dAu}(\gamma_{\rm dir})$ vs. $N_{\rm coll}$ results
 - ➡ The dashed red line is the CFM assuming all min-bias $R_{dAu} = 1$
 - But the PHENIX measurement in MB has $R_{\rm dAu}(\pi^0)/R_{\rm dAu}(\gamma_{\rm dir}) = 0.92 \pm 0.02 \text{ (stat)} \pm 0.15 \text{ (syst)}$
- The solid red line is the CFM with x0.90 MB normalization to better match what is observed in data
 - \rightarrow Great agreement vs. N_{coll} without any need for final-state physics effects!



How to take out the color fluctuation effect?

Change projectile size



- Final-state energy loss should increase from p+Au to d+Au to ³He+Au, etc.
- The impact of color fluctuations in one nucleon is diluted as one increases the projectile size

 \rightarrow Choose matching $\langle x_p \rangle$ selections



- Example ranges of $p_{\rm T}^{\gamma}$ and $p_{\rm T}^{\pi^0}$ with similar sampled x_p distributions
- Use γ 's and full jets in, SPHENIX p+Au



Spectator neutrons are much less biased



Study by ATLAS presented at HP'24 on sensitivity of FCal vs. ZDC to large- x_p selections

Major bias to the **backward multiplicity** from a high- x_p process

Bias to the spectator neutrons from a high- x_p process is smaller by a factor of 6!

Downside: loss of sensitivity to very central events



Conclusion

- Some centrality bias effects are processor kinematics-dependent
- The PHENIX π^0/γ_{dir} ratio is moderately sensitive to color fluctuation effects, due to the different x_p distributions
 - the data are quantitatively reproduced by a previous Color Fluctuation Model with no new parameters or retuning
 - no additional final-state effects needed to explain the PHENIX data

Letter

PHYSICAL REVIEW C **110**, L011901 (2024)

Contribution to differential π^0 and γ_{dir} modification in small systems from color fluctuation effects

Dennis V. Perepelitsa[®]*

Department of Physics, University of Colorado Boulder, Boulder, Colorado 80309, USA

(Received 6 May 2024; accepted 20 June 2024; published 1 July 2024)

A major complication in the search for jet quenching in proton- or deuteron-nucleus collision systems is the presence of physical effects which influence the experimental determination of collision centrality in the presence of a hard process. For example, in the proton color fluctuation picture, protons with a large Bjorken-x ($x \ge 0.1$) parton interact more weakly with the nucleons in the nucleus, leading to a smaller (larger) than expected yield in large (small) activity events. A recent measurement by the PHENIX Collaboration compared the yield of neutral pion and direct photon production in d + Au collisions, under the argument that the photon yields correct for such biases, and the difference between the two species is thus attributable to final-state effects (i.e., jet quenching). The main finding suggests a significant degree of jet quenching for hard processes in small systems. In this paper, I argue that the particular photon and pion events selected by PHENIX arise from proton configurations with significantly different Bjorken-x distributions, and thus are subject to different magnitudes of modification in the color fluctuation model. Using the results of a previous global analysis of data from the BNL Relativistic Heavy Ion Collider (RHIC) and the CERN Large Hadron Collider (LHC), I show that potentially all of the pion-to-photon difference in PHENIX data can be described by a proton color fluctuation picture at a quantitative level before any additional physics from final-state effects is required. This finding reconciles the interpretation of the PHENIX measurement with others at RHIC and LHC, which have found no observable evidence for jet quenching in small systems.

DOI: 10.1103/PhysRevC.110.L011901

More information: DVP, PRC 110 (2024) L011901 nucl-th/2404.17660



Is there a x_p -UE^{backward} anti-correlation just in *p*+*p* collisions?

➡ Not really.

ATLAS, PLB 756 (2016) 10



 $-\infty \leftarrow \eta$

 $\eta \rightarrow +\infty$

