

Canadian Institute of  
Nuclear Physics  
Institut canadien de  
physique nucléaire

# Dilepton emission in heavy ion collisions and equilibrium of QCD matter

Xiang-Yu Wu

In collaboration with Lipei Du, Charles Gale and Sangyong Jeon

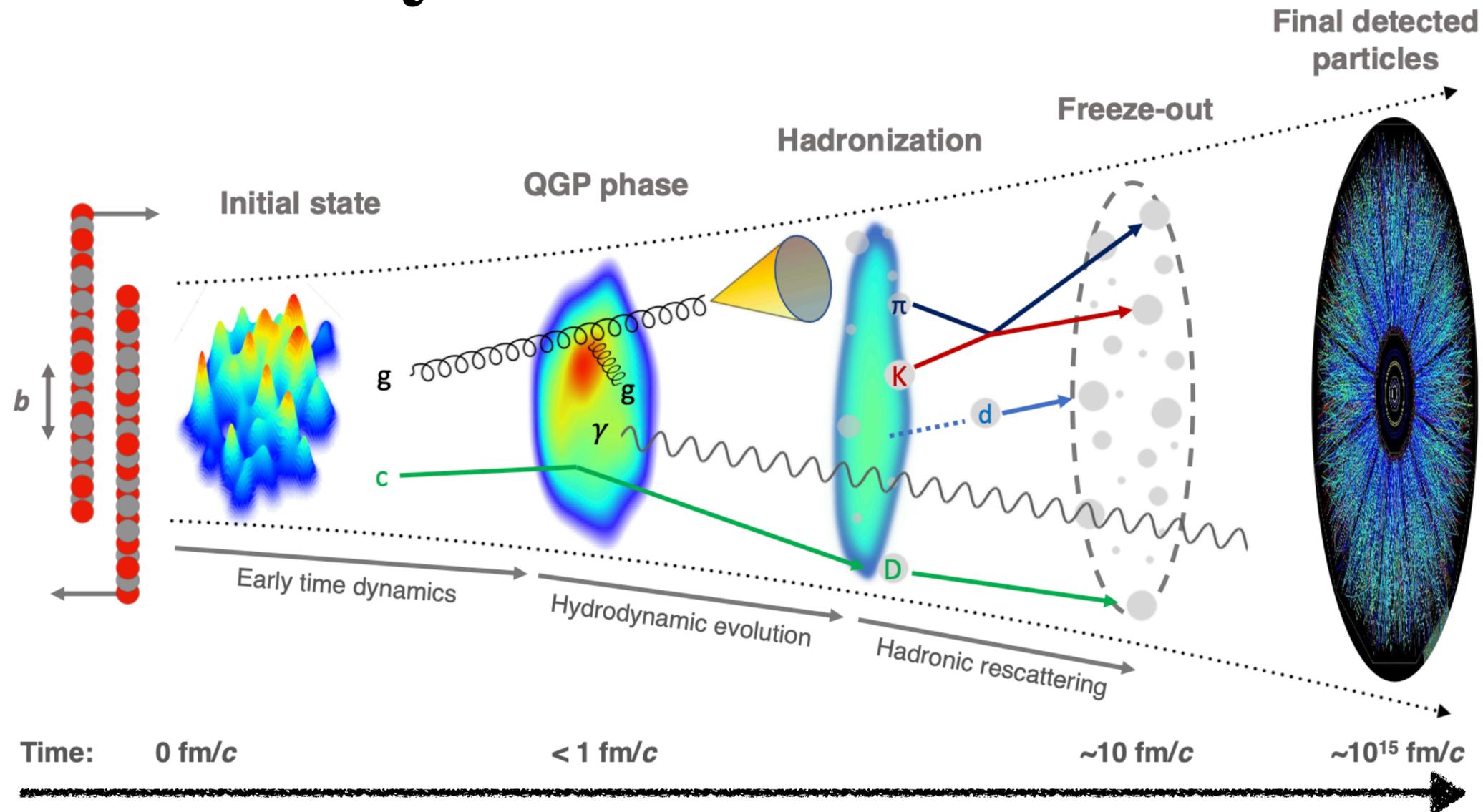
McGill University

Based on e-Print: 2407.04156 [nucl-th]

Hard Probe 2024, Nagasaki, Japan  
September 24, 2024



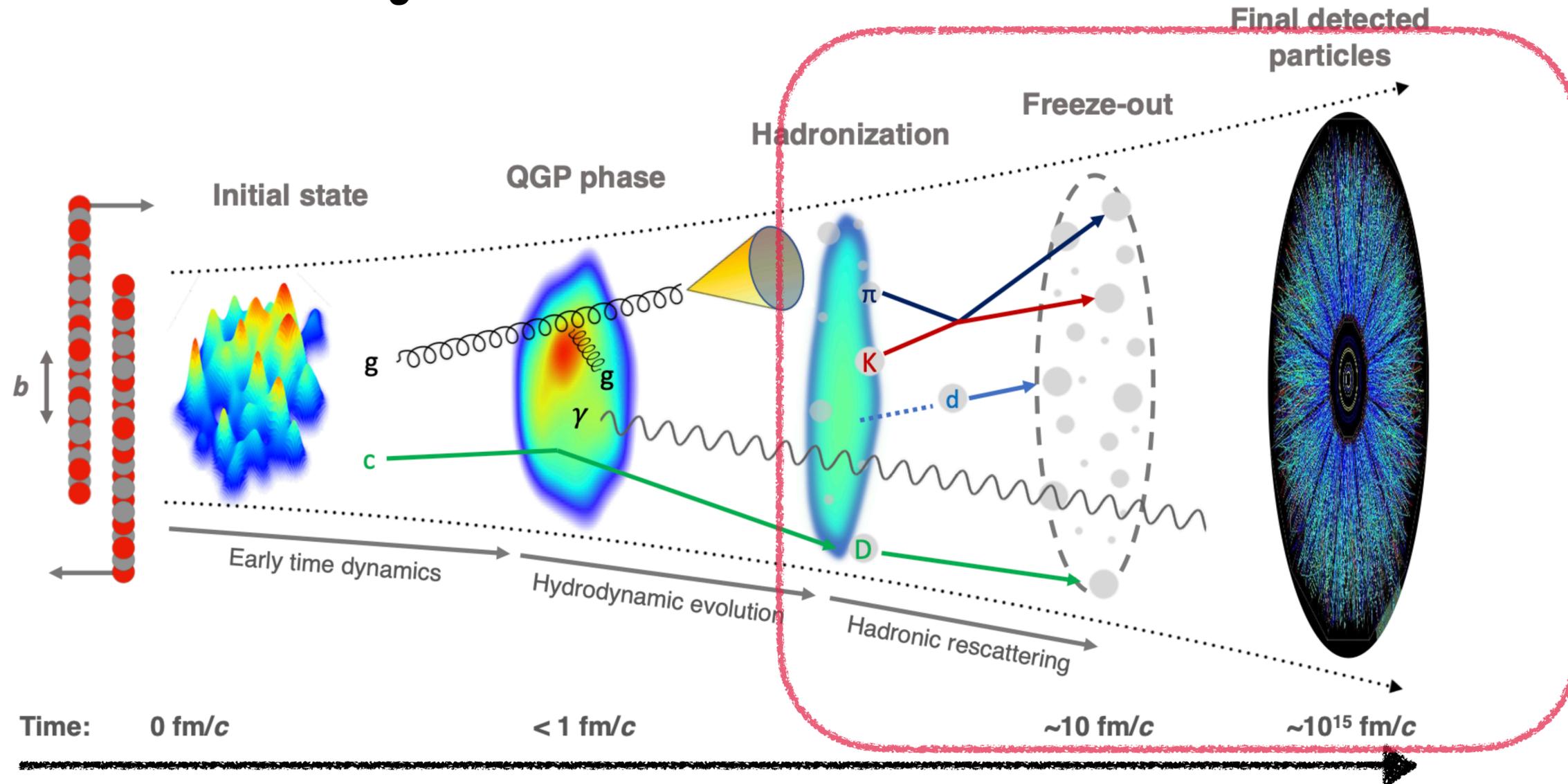
# Relativistic Heavy Ion Collision



Goal: explore and study the properties of QGP.

[ALICE, Eur.Phys.J.C 84 (2024) 8, 813]

# Relativistic Heavy Ion Collision



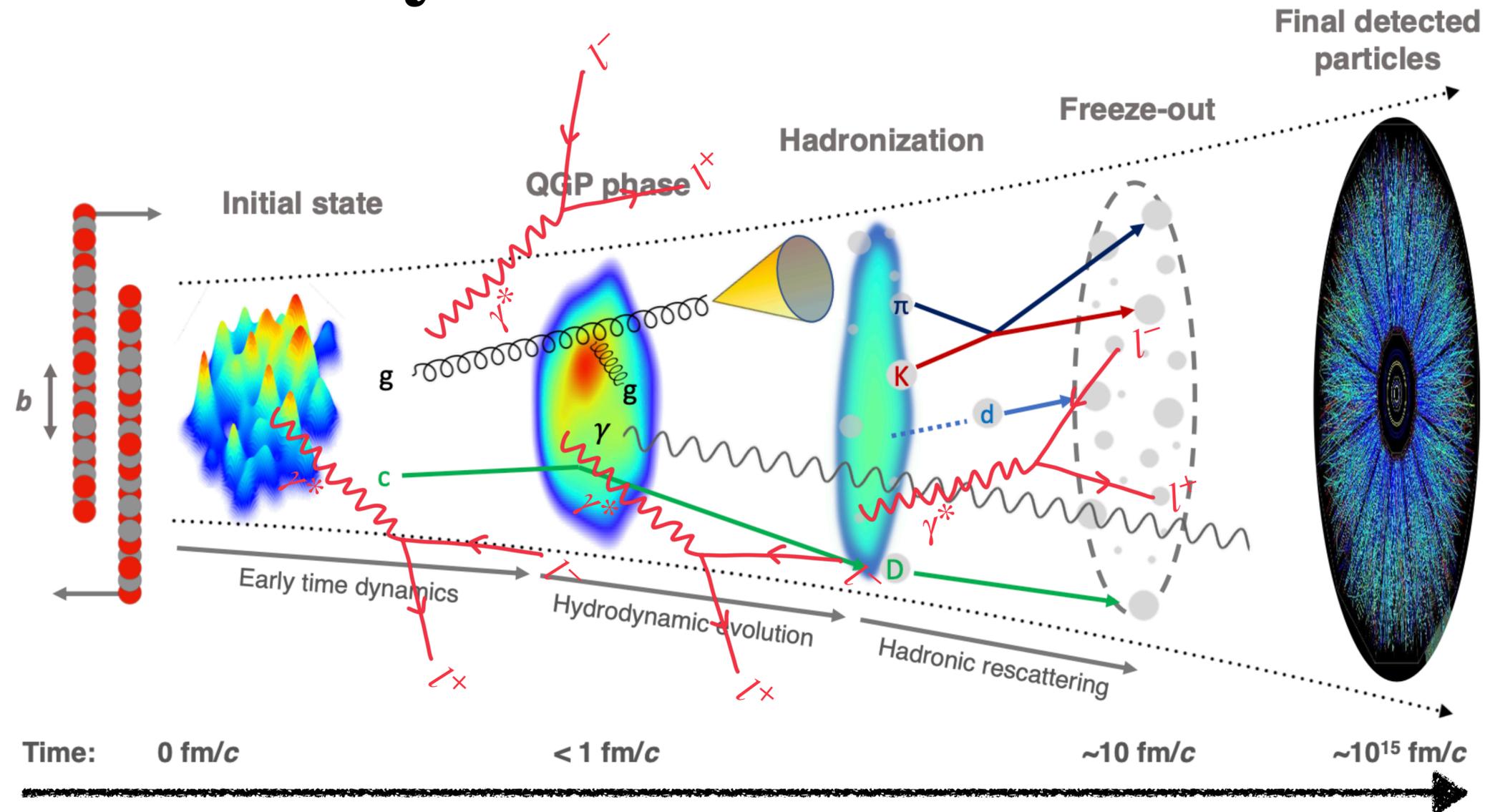
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Probes:

- Soft hadrons:  $T < T_{\text{frz}}$ , indirect probe.

# Relativistic Heavy Ion Collision



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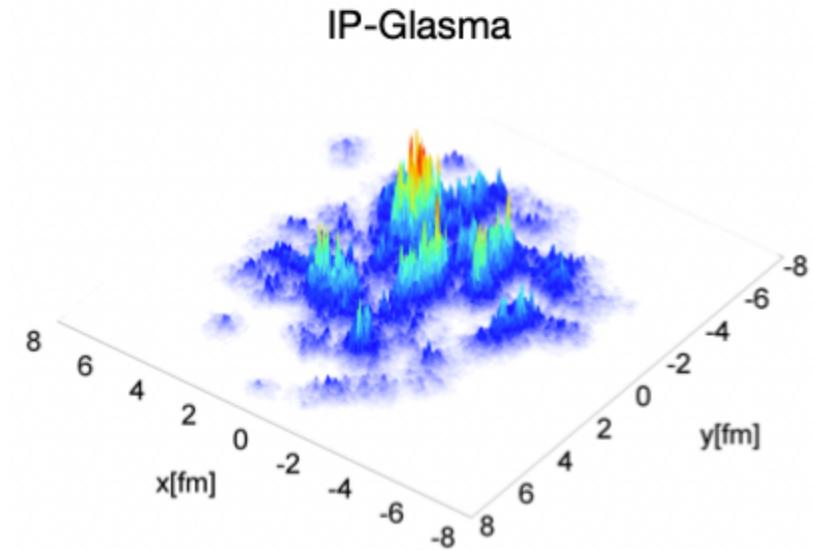
- Soft hadrons:  $T < T_{\text{frz}}$ , indirect probe.
- EM probes ( $\gamma/l\bar{l}$ ):  $T > T_{\text{frz}}$ , direct probe.

# iEBE-MUSIC Framework

$\tau = 0^+$  fm

Initial condition: IP-Glasma

- Color glass condensates.
- Evolved by Yang-Mills equations.



[Schenke, Tribedy & Venugopalan, PRL 108, 252301 (2012)]  
[Schenke, Tribedy & Venugopalan, PRC 86, 034908 (2012)]

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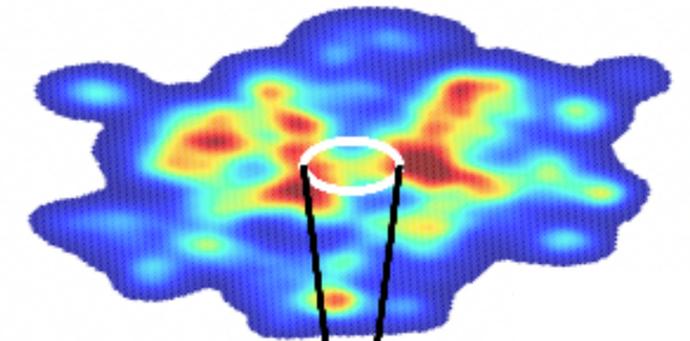
$\tau = 0.1$  fm

Pre-equilibrium stage: K $\phi$ MP $\phi$ ST

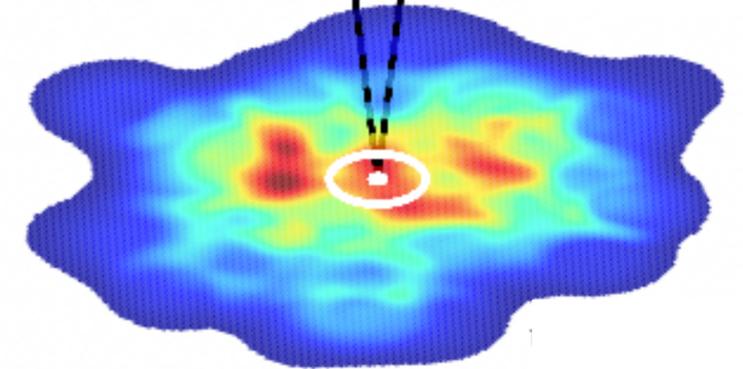
$$T^{\mu\nu}(\tau, \mathbf{x}) = \bar{T}^{\mu\nu}(\tau, \mathbf{x}) + \delta T^{\mu\nu}(\tau, \mathbf{x})$$

- Background and perturbation.
- Kinetic theory scaling curve and linear response.

$\tau = 0.1$ fm



$\tau = 0.8$ fm



[Kurkela, Mazeliauskas, Paquet, Schlichting & Teaney, PRC 99, 034910 (2019)]  
[Kurkela, Mazeliauskas, Paquet, Schlichting & Teaney, PRL122, 122302 (2019)]

# iEBE-MUSIC Framework

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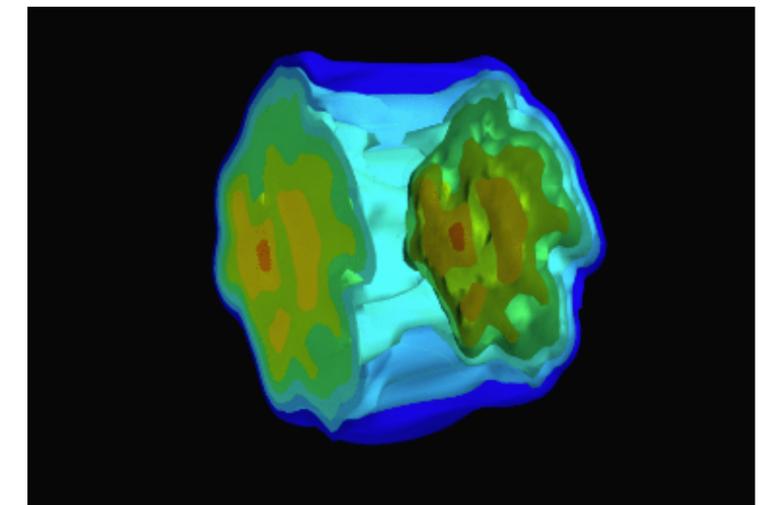
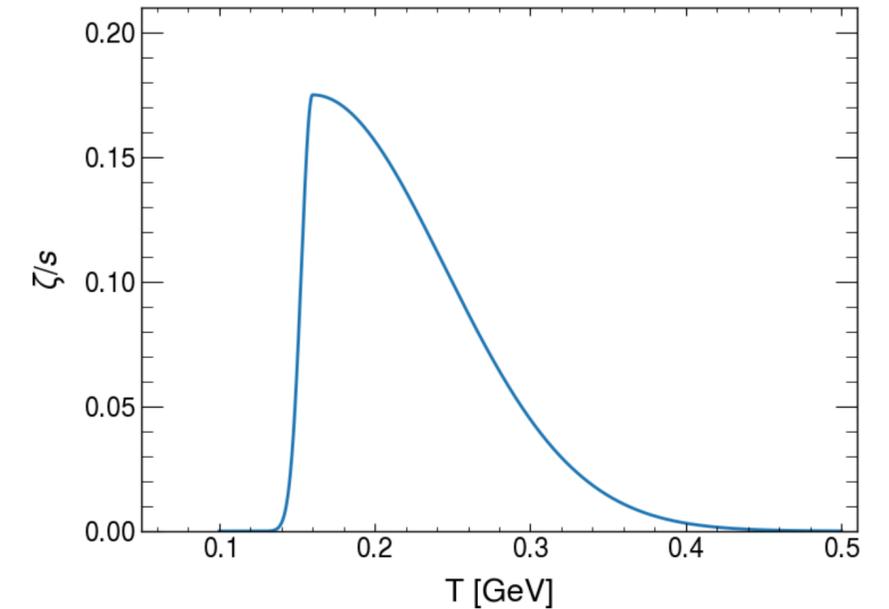
- Background and perturbations.
- Kinetic theory scaling curve and linear response.

$\tau = 0.8$  fm

Hydrodynamic evolution: MUSIC

- (2+1)-D event-by-event 2<sup>nd</sup> viscous hydrodynamics model.
- Shear viscosity  $\eta/s = 0.12$  and bulk viscosity  $\zeta/s(T)$ .
- HotQCD+hadron gas equation of states.

$\tau \approx 10$  fm



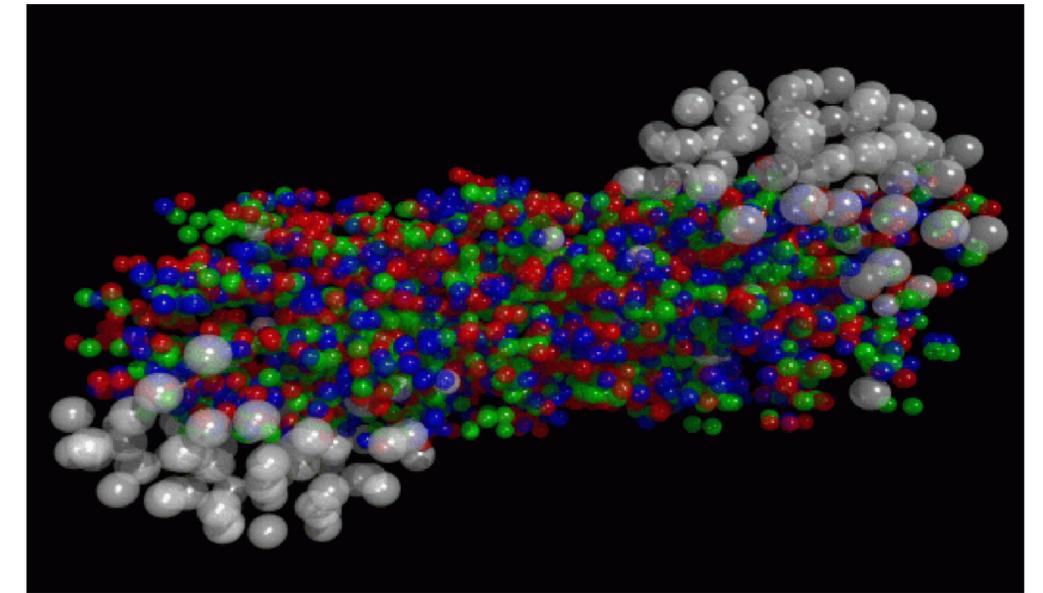
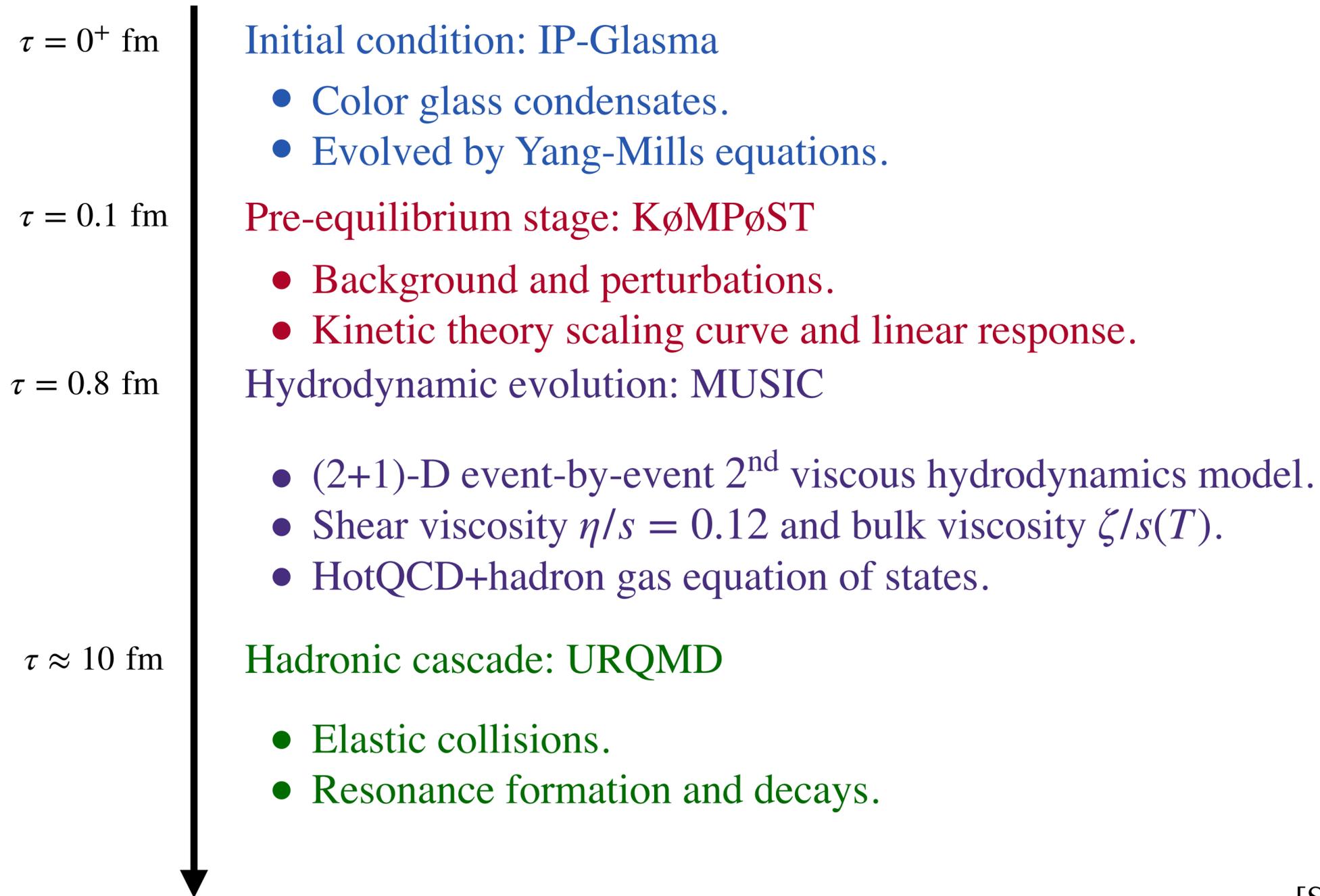
[Paquet, Shen, Denicol, Luzum, Schenke, Jeon & Gale, PRC 93, 044906 (2016)]

[Schenke, Jeon & Gale, PRC 82, 014903 (2010)]

[Schenke, Jeon & Gale, PRL 106, 042301 (2011)]

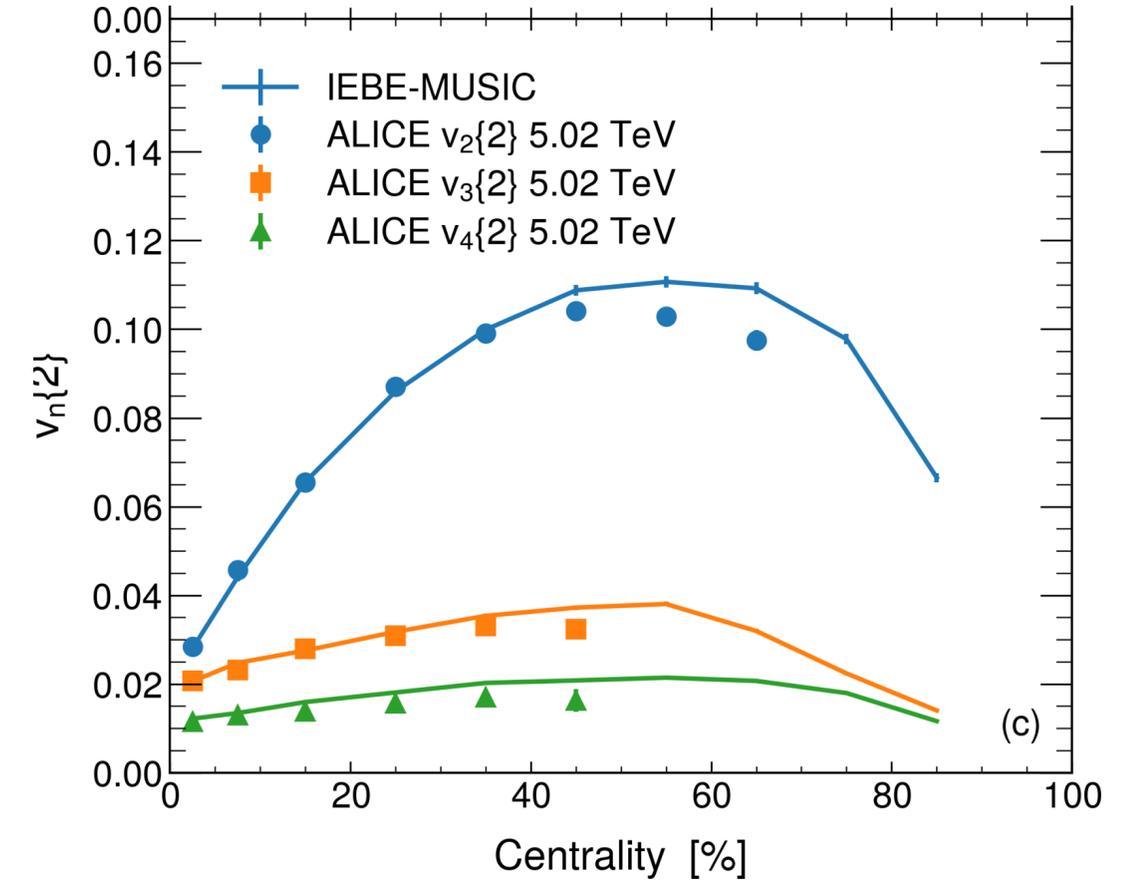
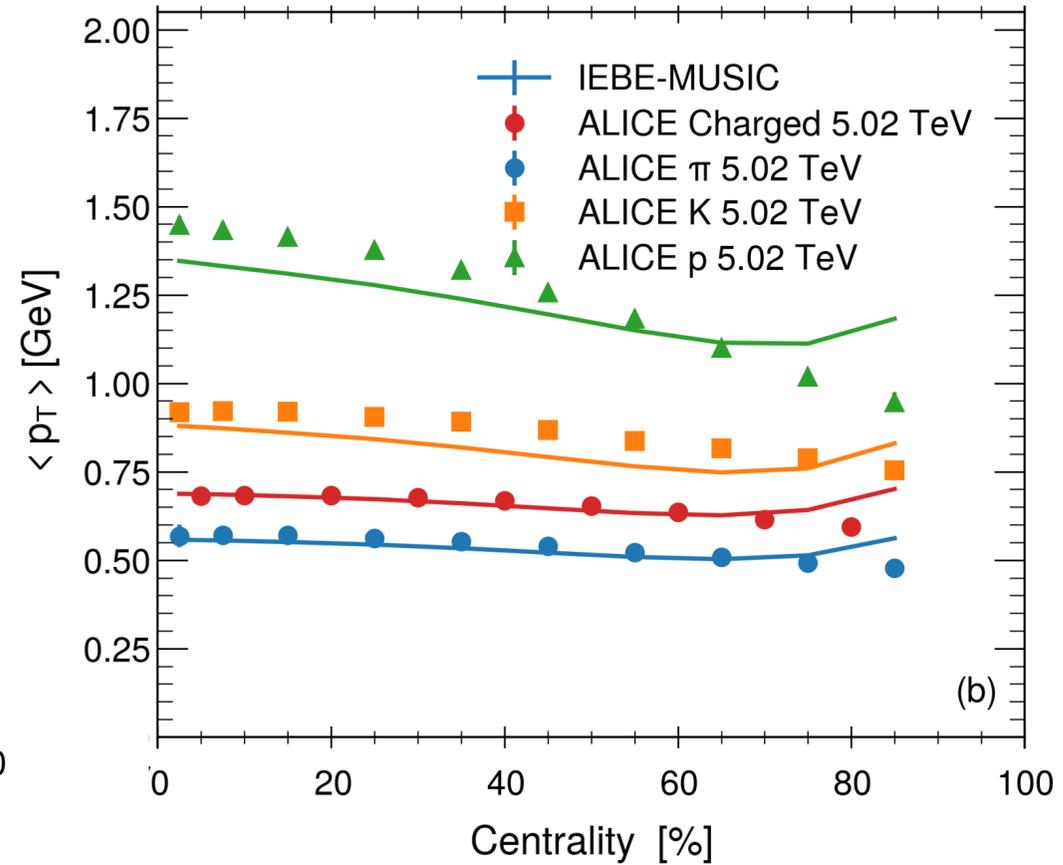
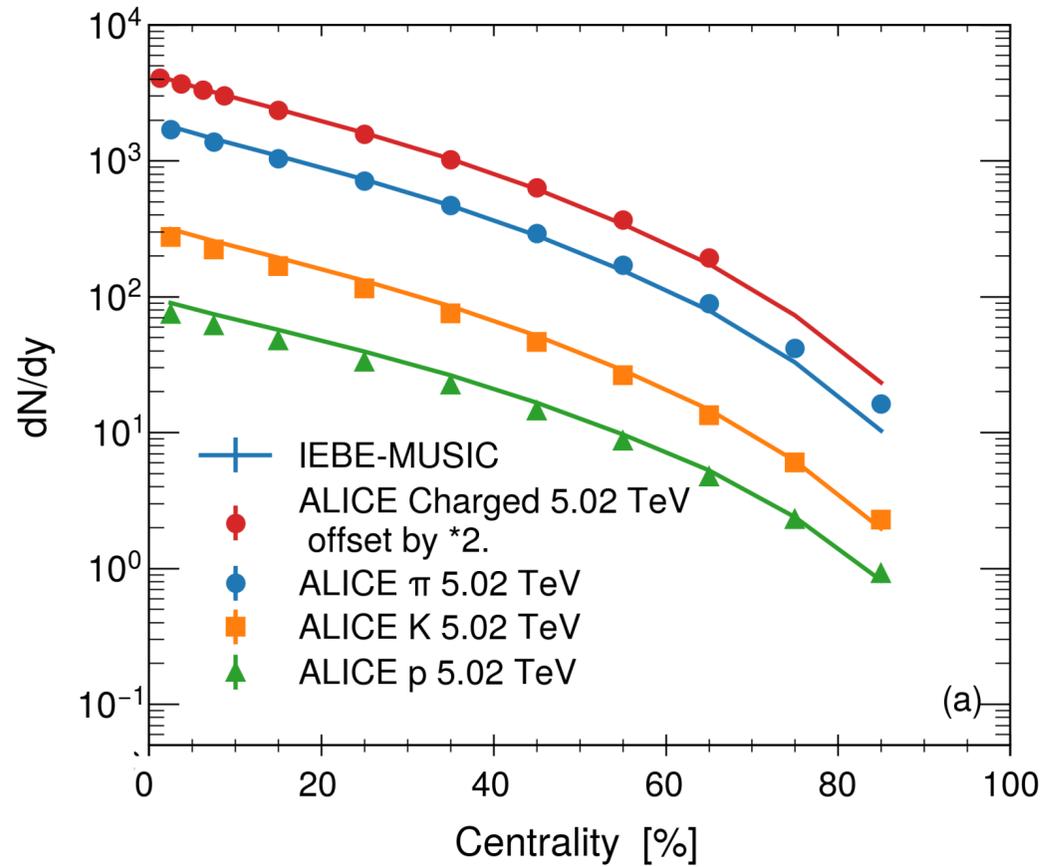
[Gale, Paquet, Schenke & Shen, PRC 105, 014909 (2022)]

# iEBE-MUSIC Framework



[S. A. Bass et al, Prog. Part. Nucl. Phys. 41, 255–369 (1998)]

# Model Calibration



- Hydrodynamic results agree well with hadron multiplicity in all centrality bins.
- Small discrepancy in mean transverse momentum  $\langle p_T \rangle$  and anisotropic flow  $v_n\{2}$  originates from vanishing bulk viscosity in IP-Glasma and K $\emptyset$ MP $\emptyset$ ST.
- The discrepancy increases in peripheral collision due to shorter lifetime in hydrodynamic phase.

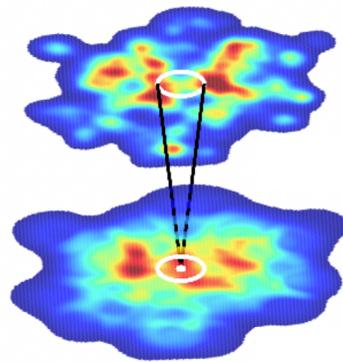
# Dilepton Production

$\tau = 0^+$  fm

Initial condition: IP-Glasma

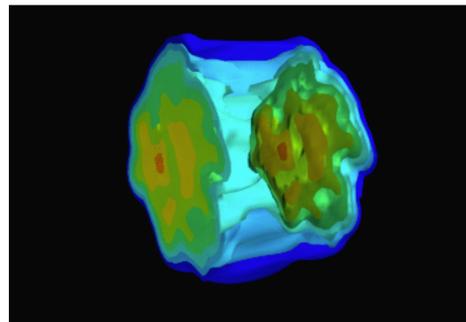
$\tau = 0.1$  fm

Pre-equilibrium stage: K $\emptyset$ MP $\emptyset$ ST



$\tau = 0.8$  fm

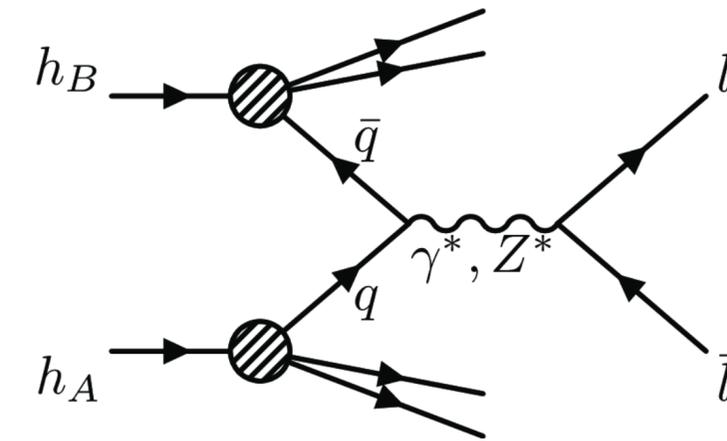
Hydrodynamic evolution: MUSIC



$\tau \approx 10$  fm

Hadronic cascade: URQMD

Drell-Yan dilepton



DYTurbo package: NLO order in p+p collision.

In A+A collisions

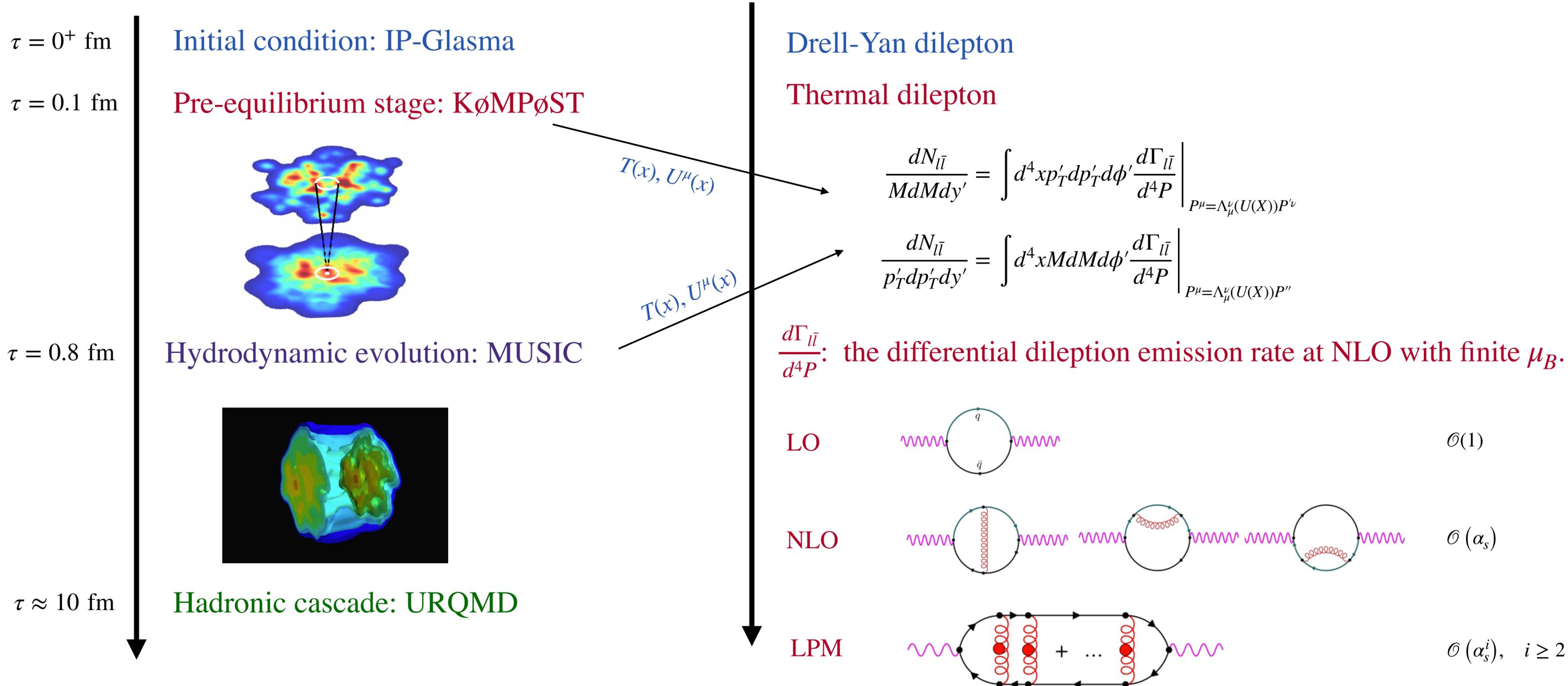
$$\frac{dN_{ee}^{DY}}{dMdy} = \frac{d\sigma_{ee}^{DY;pp}}{dMdy} \cdot \frac{N_{coll}}{\sigma_{in}^{pp}}$$

nPDF: EPPS16nlo-CT14nlo-Pb208.

Thickness function: MC-Glauber.

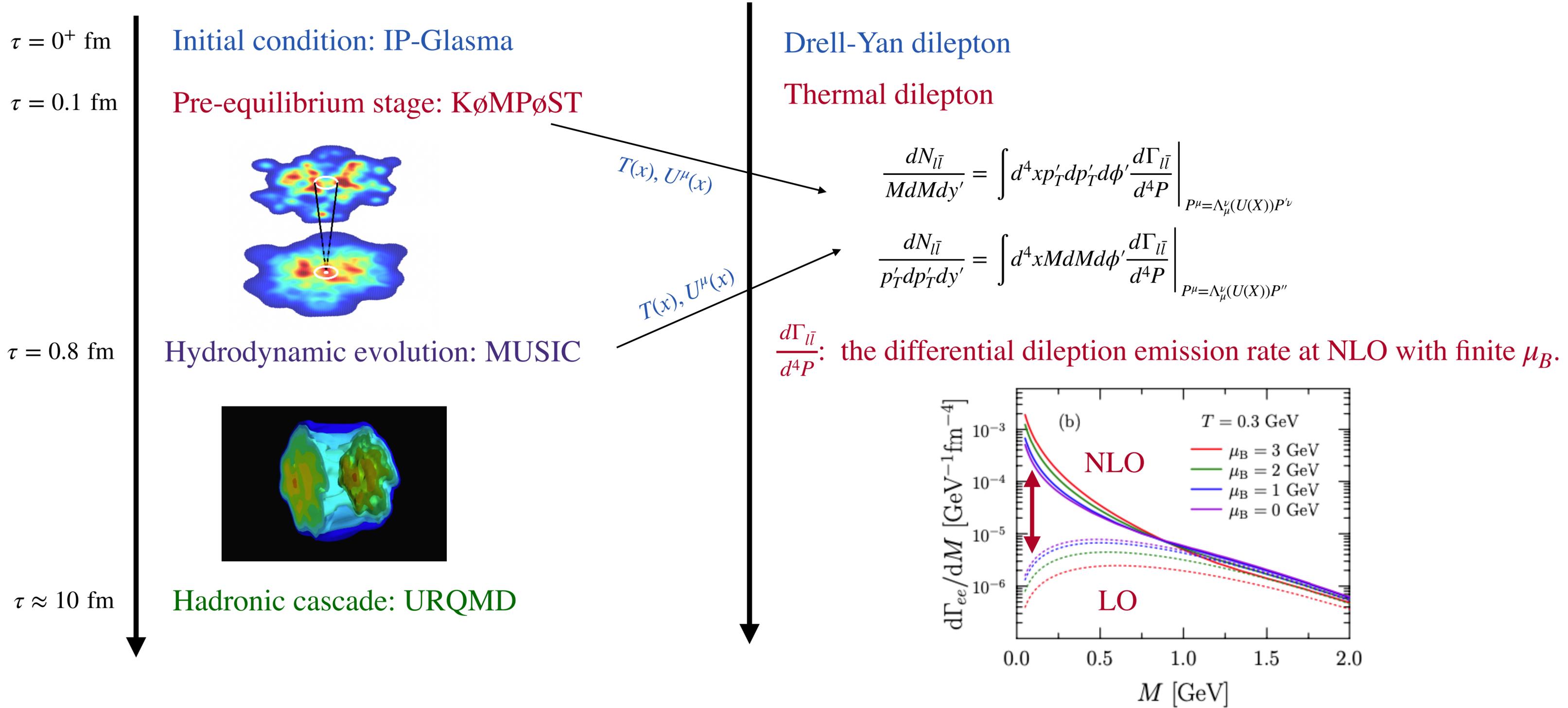
[Camarda et al., EPJC 80, 251 (2020)]  
 [Eskola, Paakkinen, Paukkunen & Salgado, EPJC 77, 163 (2017)]  
 [Loizides, Kamin & Enterria, PRC 97, 054910 (2018)]

# Dilepton Production



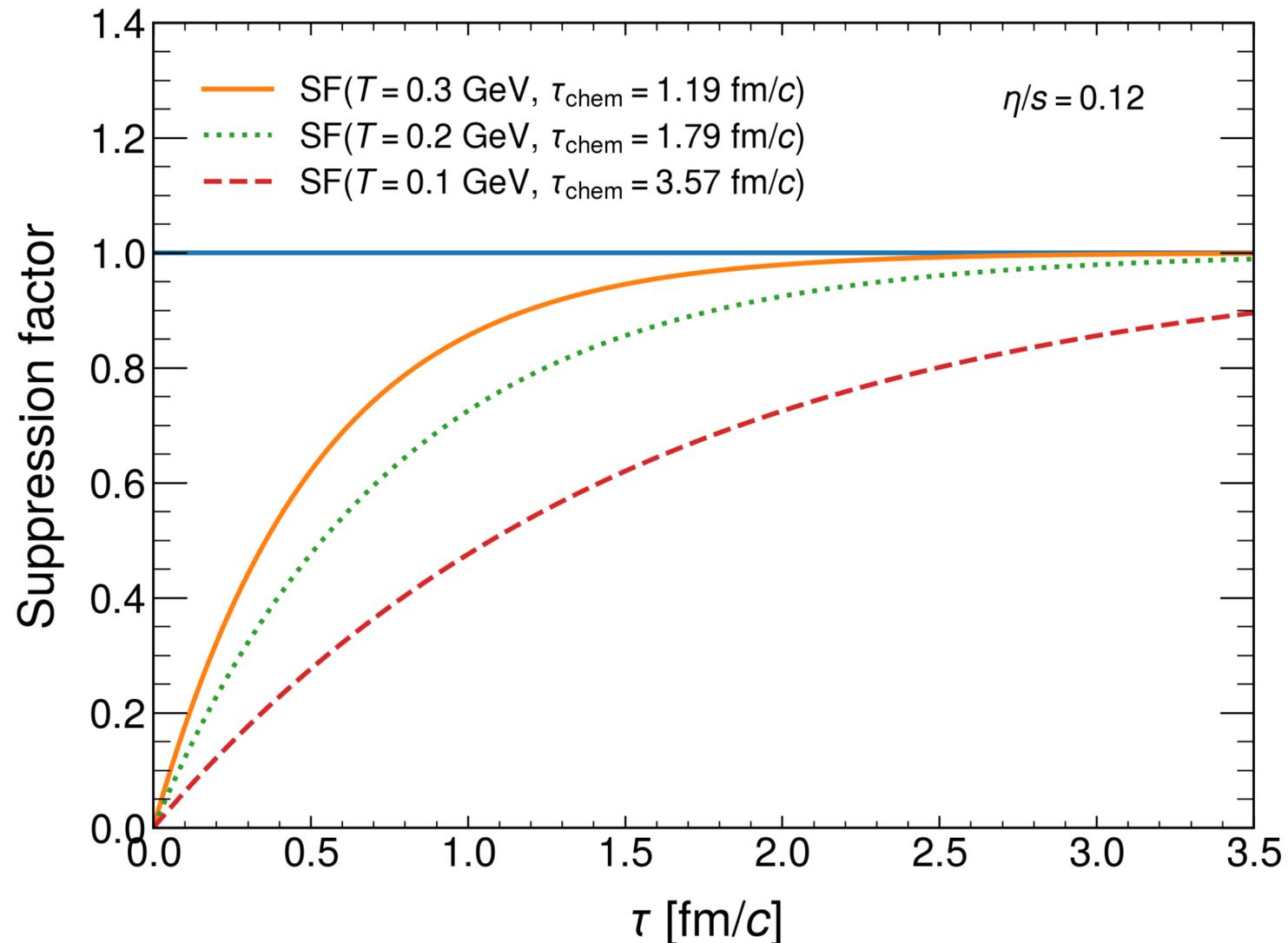
[Churchill, Du, Gale, Jackson & Jeon, PRC 109, 044915 (2024), PRL 132, 172301 (2024)]

# Dilepton Production



# Chemical Equilibrium

- The gluon fields dominate in the early stage.
- Quarks and anti-quarks are produced via  $gg \rightarrow q\bar{q}$  and  $g \rightarrow q\bar{q}$ .
- The effective suppression factor  $SF(T, \tau)$  is introduced in relation to chemical relaxation time  $\tau_R = 4\pi\eta/Ts$ .

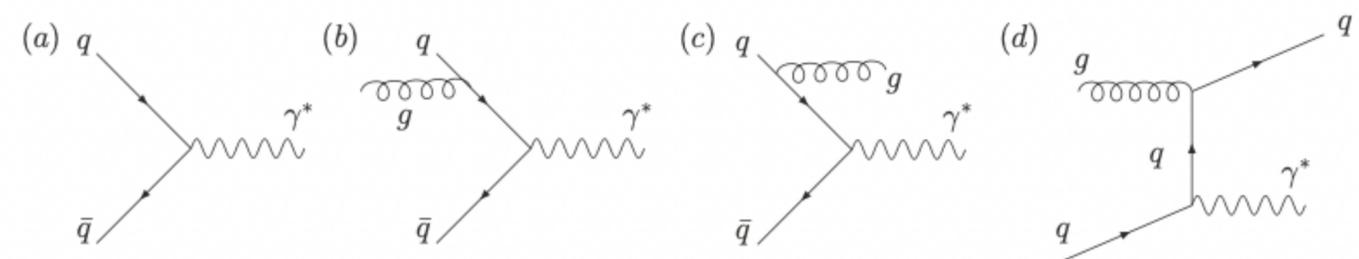


Systems with higher temperatures reach chemical equilibrium faster than those at lower temperatures.

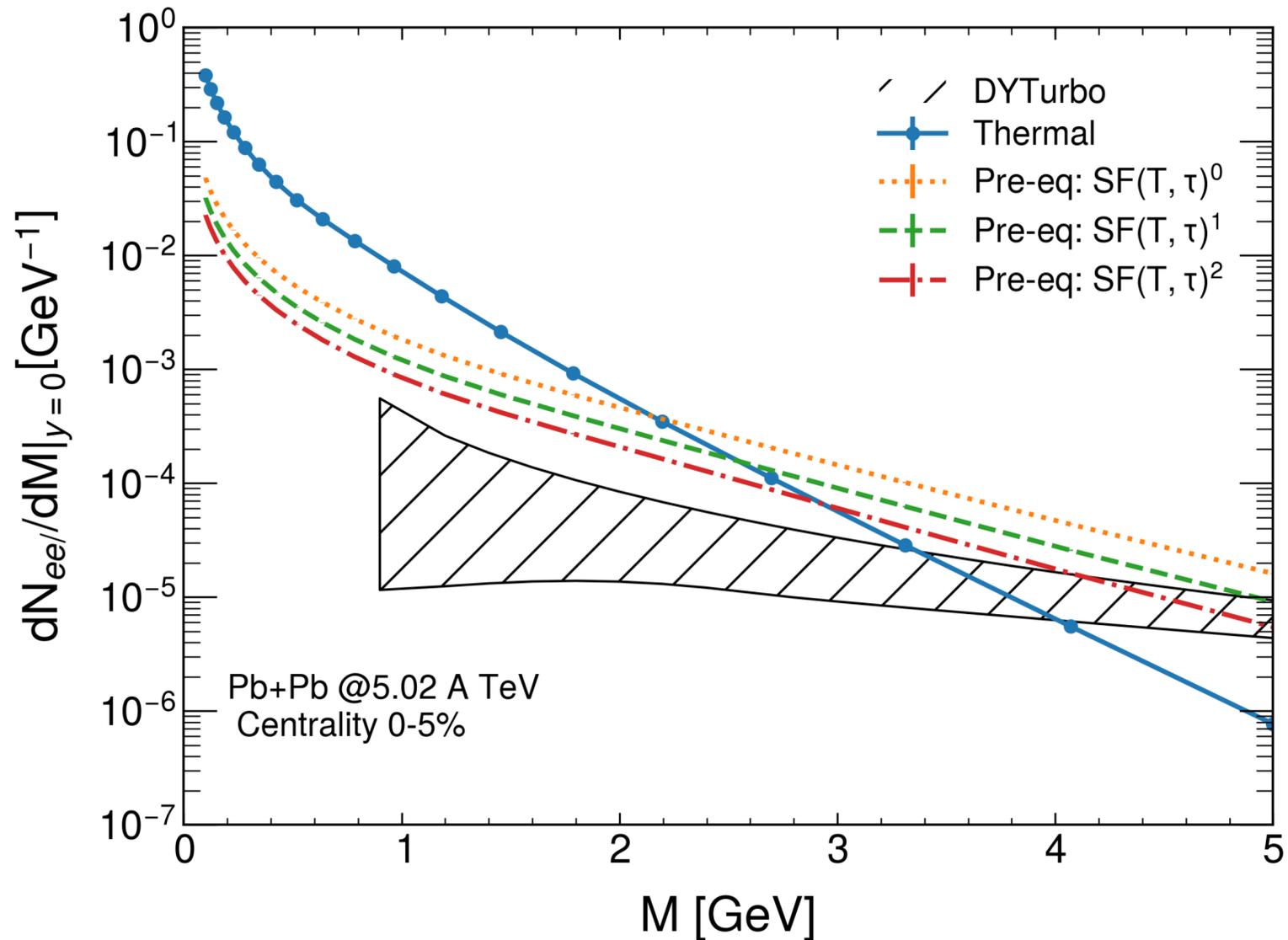
$$\frac{dN_{l\bar{l}}}{MdMdy'} = \int dV SF(T, \tau)^n \frac{d\Gamma_{l\bar{l}}}{d^4P} \Bigg|_{P^\mu = \Lambda_\mu^\nu(U(X))P^\nu} p'_T dp'_T d\phi'$$

$$\frac{dN_{l\bar{l}}}{p'_T dp'_T dy'} = \int dV SF(T, \tau)^n \frac{d\Gamma_{l\bar{l}}}{d^4P} \Bigg|_{P^\mu = \Lambda_\mu^\nu(U(X))P''} MdM d\phi'$$

Here  $n = 0, 1, 2$ .



# Dilepton $M$ spectrum

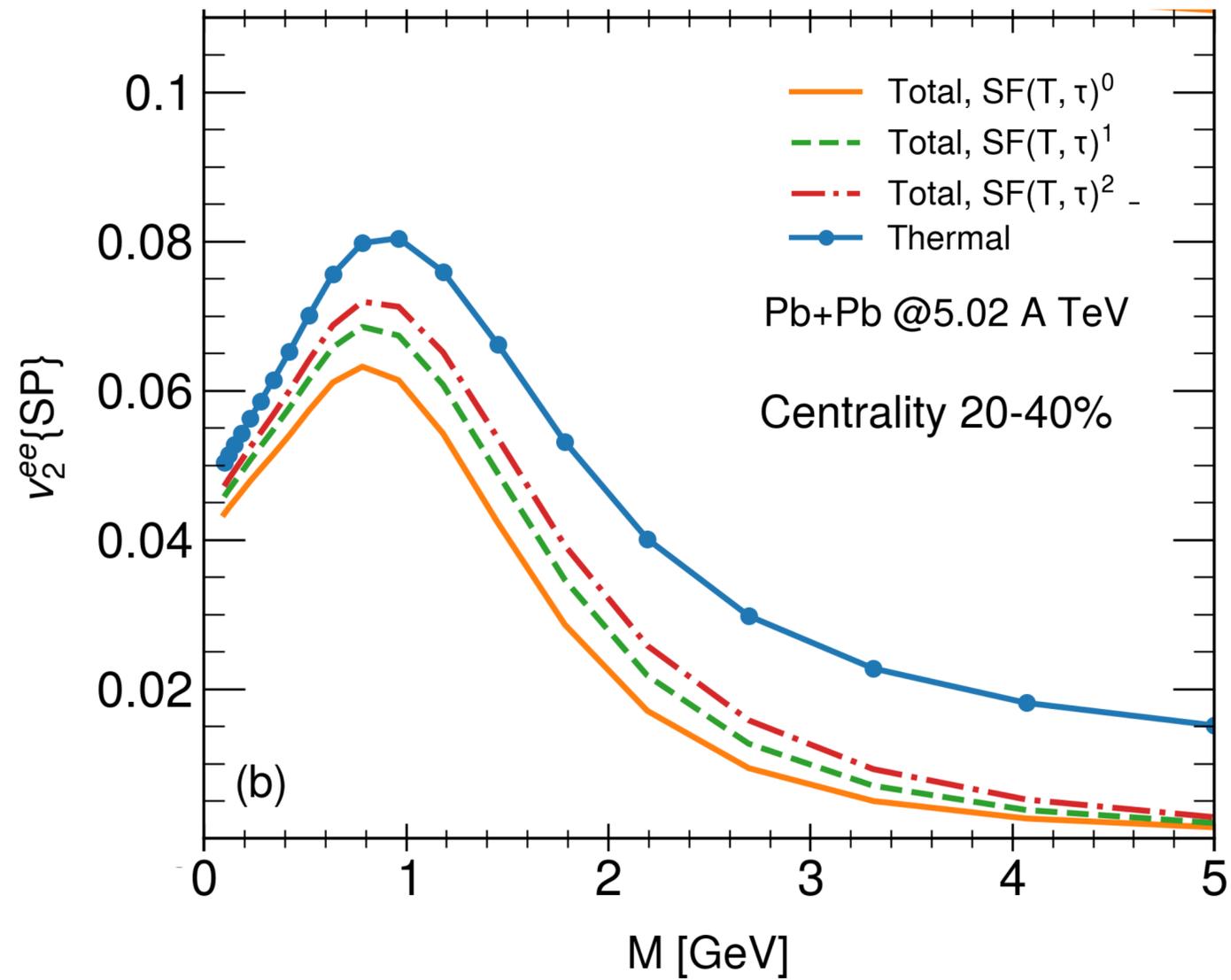


- Thermal dilepton production dominates in the low-mass region (LMR).
- The pre-equilibrium dileptons are the dominant source over thermal dileptons in the intermediate mass region (IMR).
- Drell-Yan dileptons are consistently smaller than the pre-equilibrium contribution until the high-mass region (HMR).
- Dilepton production during the pre-equilibrium phase will be sensitive to the power of suppression factor for chemical equilibrium.

Oscar's Talk@Tue 9:00AM.

[Oscar Garcia-Montero, Philip Plaschke, Sören Schlichting, e-Print: 2403.04846 [hep-ph]]

# Dilepton $v_n^{ee}\{SP\}(M)$

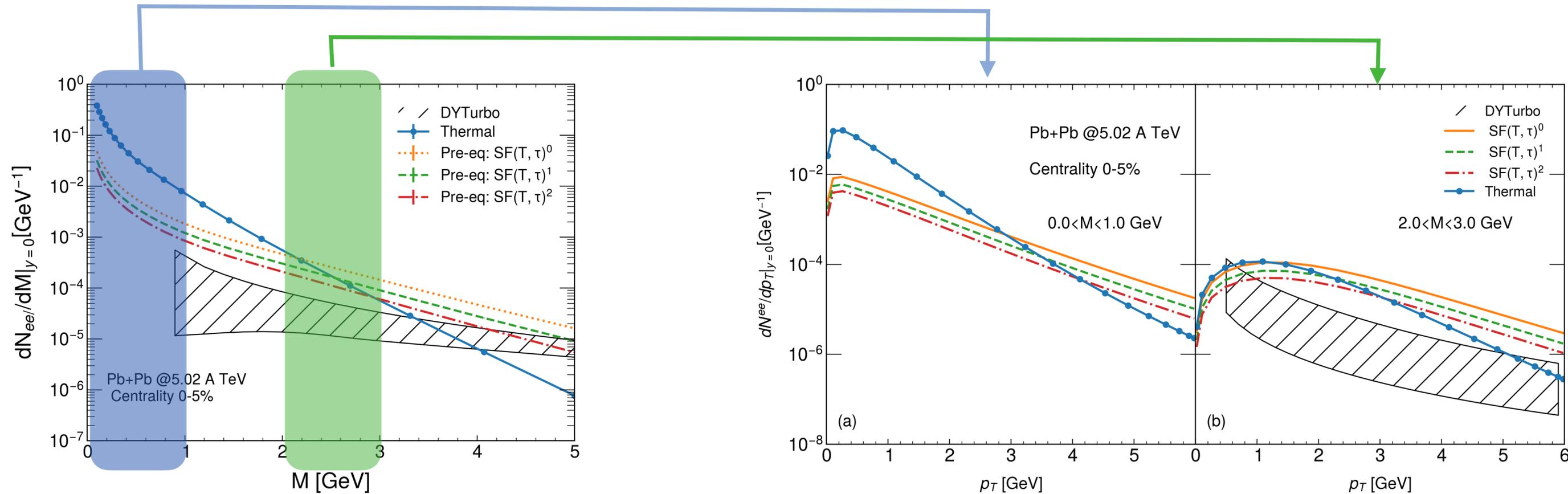


- The scalar-product (SP) method:

$$v_n^{\bar{l}l}\{SP\}(X) = \frac{\left\langle v_n^{\bar{l}l}(X) v_n^h \cos \left\{ n \left[ \Psi_n^{\bar{l}l}(X) - \Psi_n^h \right] \right\} \right\rangle}{\sqrt{\left\langle (v_n^h)^2 \right\rangle}}$$

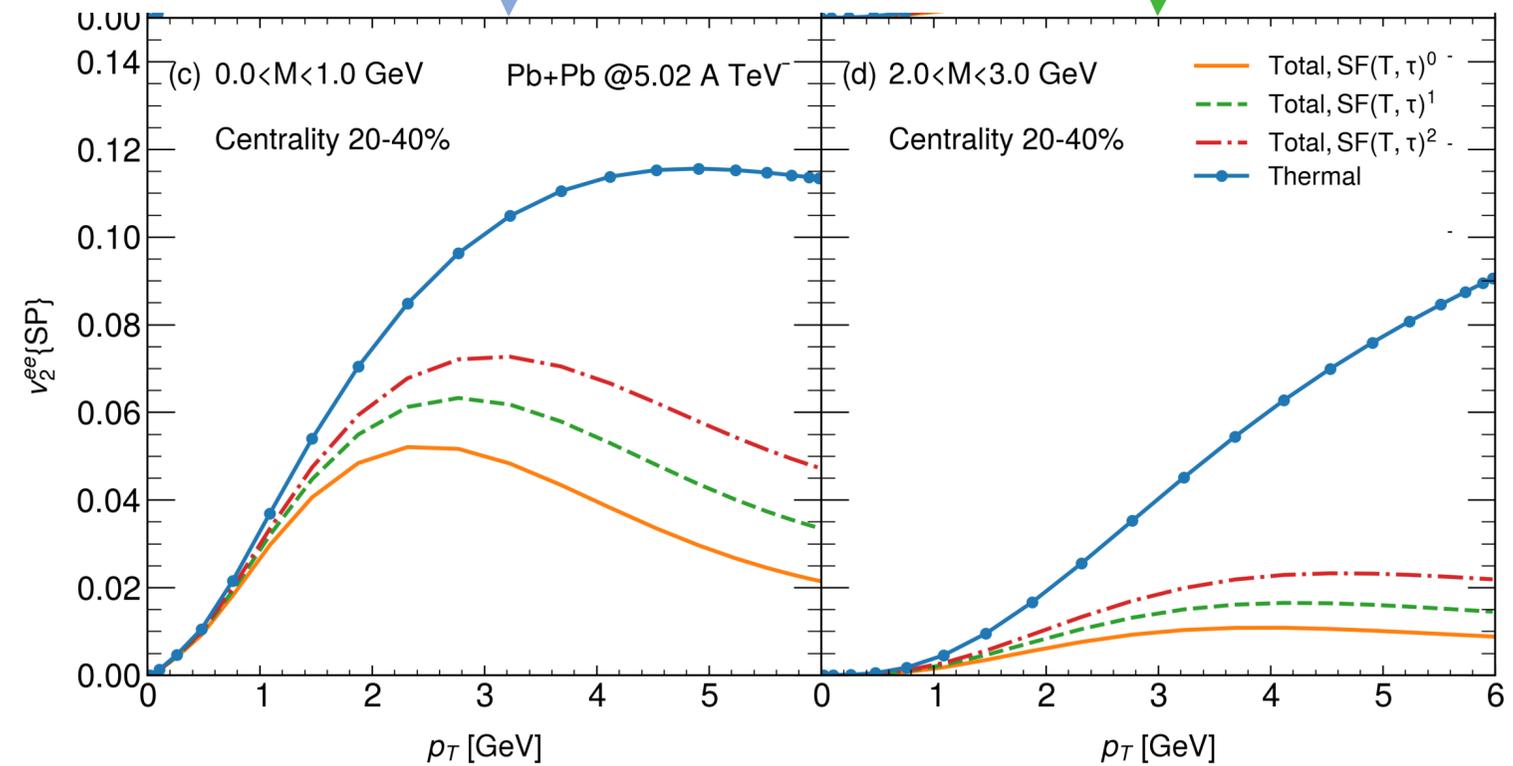
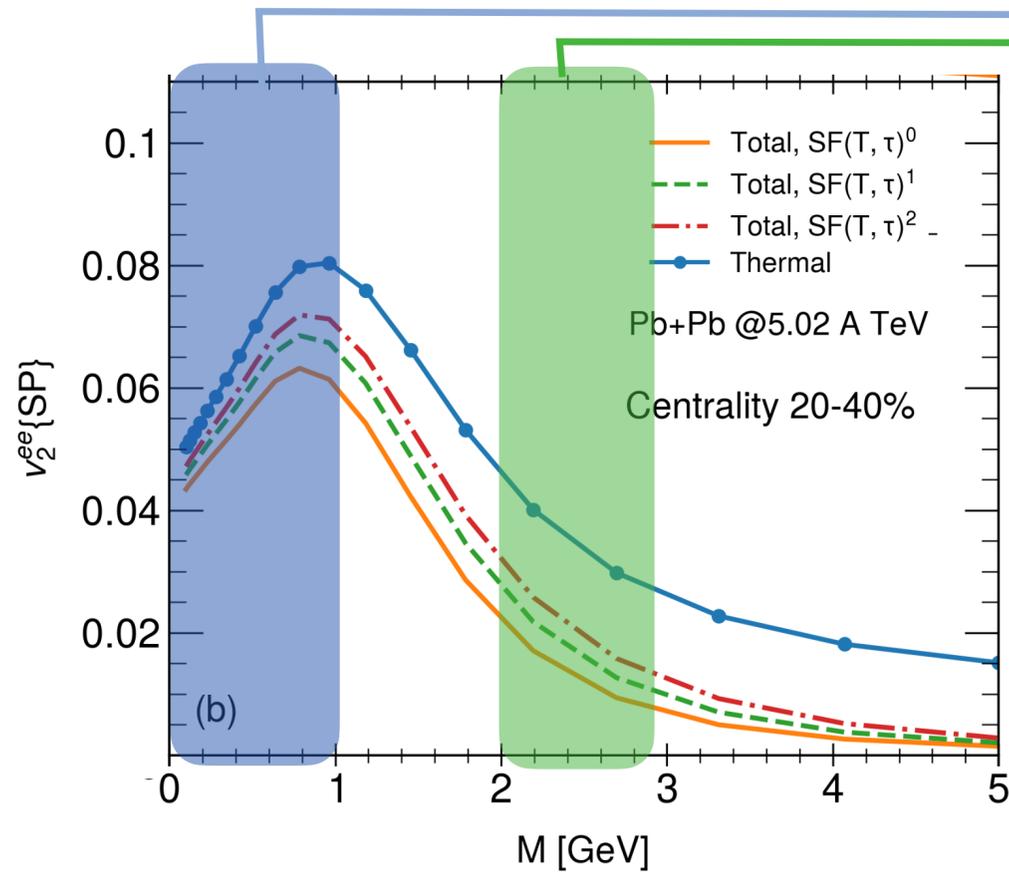
- The total  $v_2^{ee}(M)$  will be suppressed due to the almost zero momentum anisotropy in the pre-equilibrium stage.
- For the higher order of the suppression factor, the total dilepton flow  $v_2^{ee}(M)$  moves closer to that of thermal dileptons.

# Dilepton $p_T$ spectra



- By carefully selecting the invariant mass window, it is possible to reveal the pre-equilibrium stage.
- $0 < M < 1$  GeV:
  - ▶ the thermal dilepton production dominates at low  $p_T$  region.
  - ▶ dilepton from the pre-equilibrium stage becomes dominant from intermediate  $p_T$  region.
- $2 < M < 3$  GeV:
  - ▶ the pre-equilibrium dilepton is comparable to thermal dilepton production in the low  $p_T$  region.
  - ▶ the pre-equilibrium dilepton consistently exceeds both thermal dilepton and Drell-Yan dilepton production.

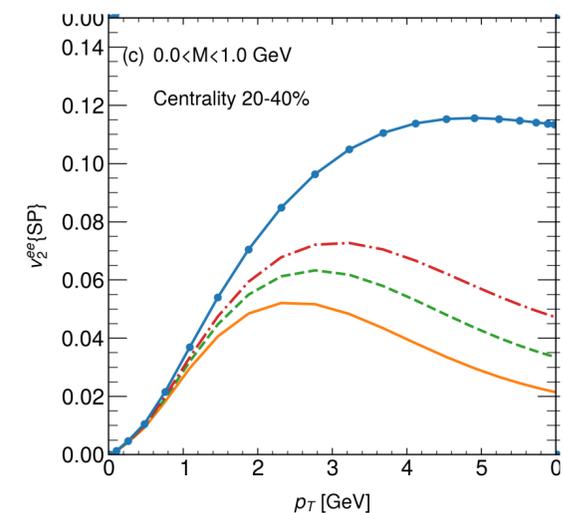
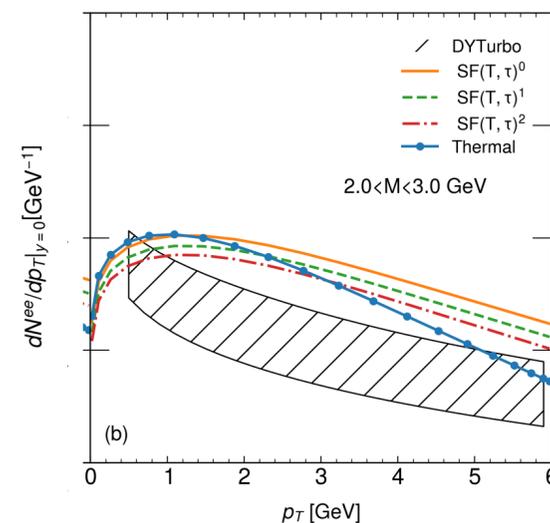
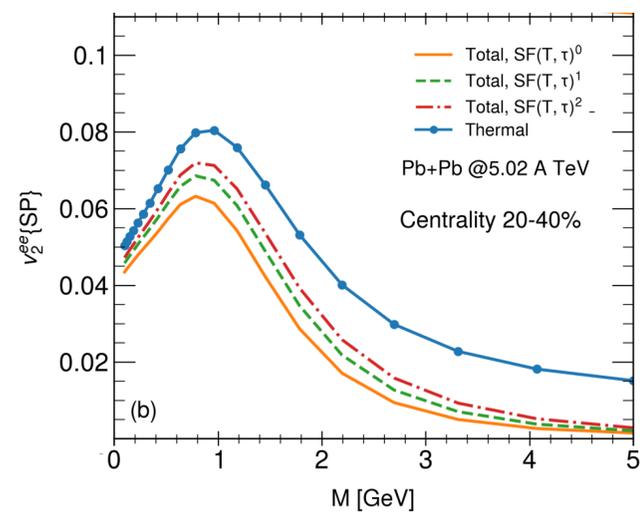
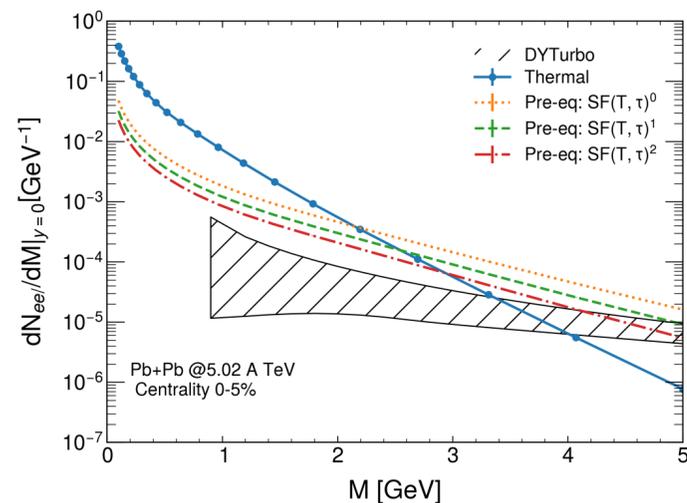
# Dilepton $v_2^{ee}\{SP\}(p_T)$



- Total elliptic flow  $v_2^{ee}(p_T)$  is highly sensitive to the effects of chemical equilibrium.
- $0 < M < 1$  GeV:
  - ▶ after considering the pre-equilibrium stage, the total dilepton flow  $v_2^{ee}(p_T)$  is significantly suppressed.
- $2 < M < 3$  GeV:
  - ▶ the total dilepton elliptic flow  $v_2^{ee}(p_T)$  shows a noticeable reduction compared to that in the IMR.

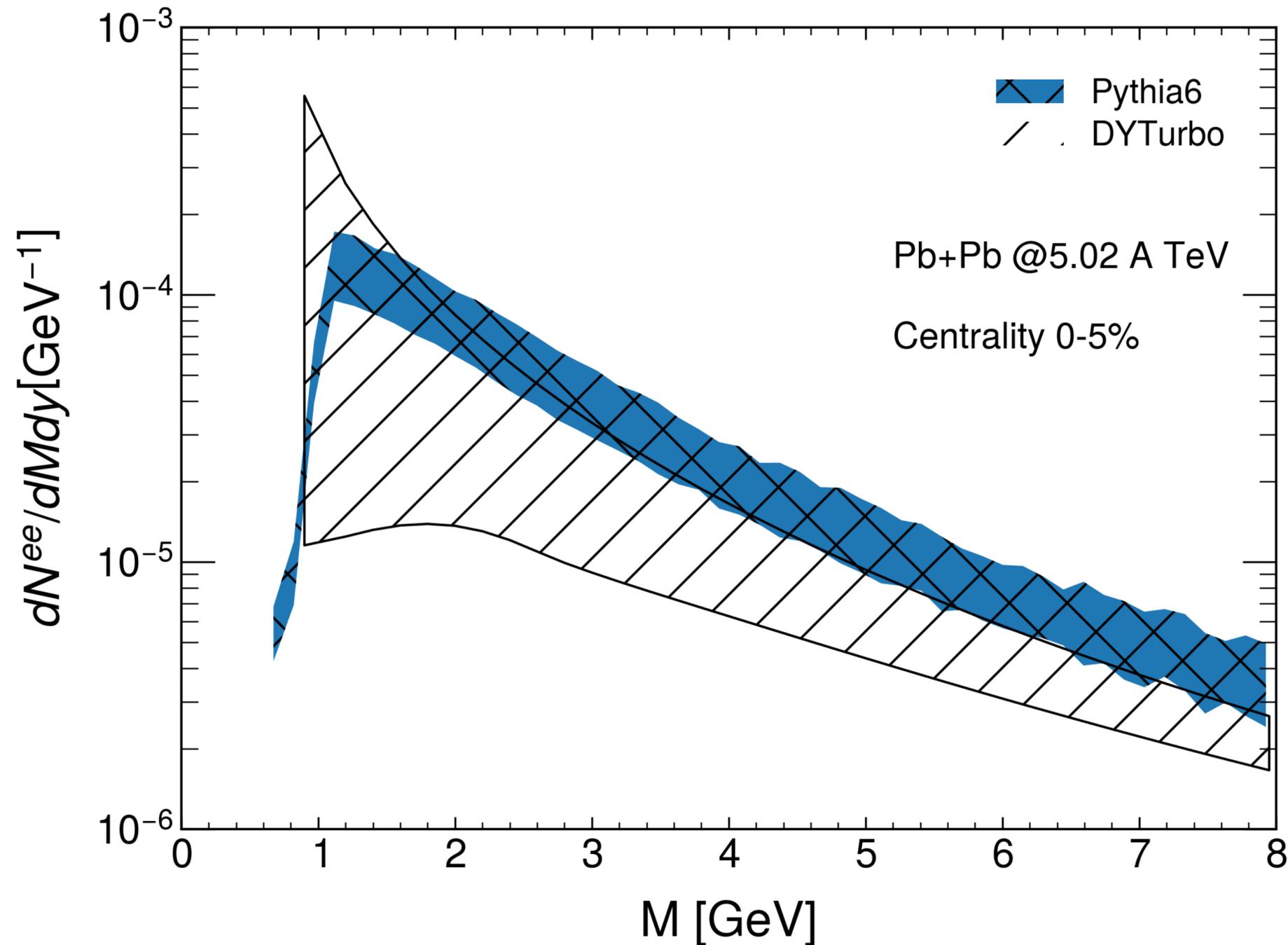
# Summary

- Dilepton observables offer an opportunity to explore quark production of the pre-equilibrium stage.
- Dilepton flow could help constrain the effects of chemical equilibrium during the pre-equilibrium stage in the future.
- An analysis of thermal dilepton production and anisotropic flow based on NLO order emission rate is performed in Pb+Pb collisions at a LHC collision energy of  $\sqrt{s_{NN}} = 5.02$  TeV.
- The pre-equilibrium dilepton contribution dominates in the IMR, exceeding both Drell-Yan and thermal dilepton production.
- Dilepton production is slightly affected by chemical equilibrium.



# Backup

# DYTurbo vs Pythia 6



Pythia 6:

PDF CTEQ5L

Bands: K factor, effectively account for higher order corrections.

DYTurbo:

Bands:

NLO, factorization and renormalization scales.

Upper limits:  $2m_{ll}$ , lower limit:  $0.5m_{ll}$ .

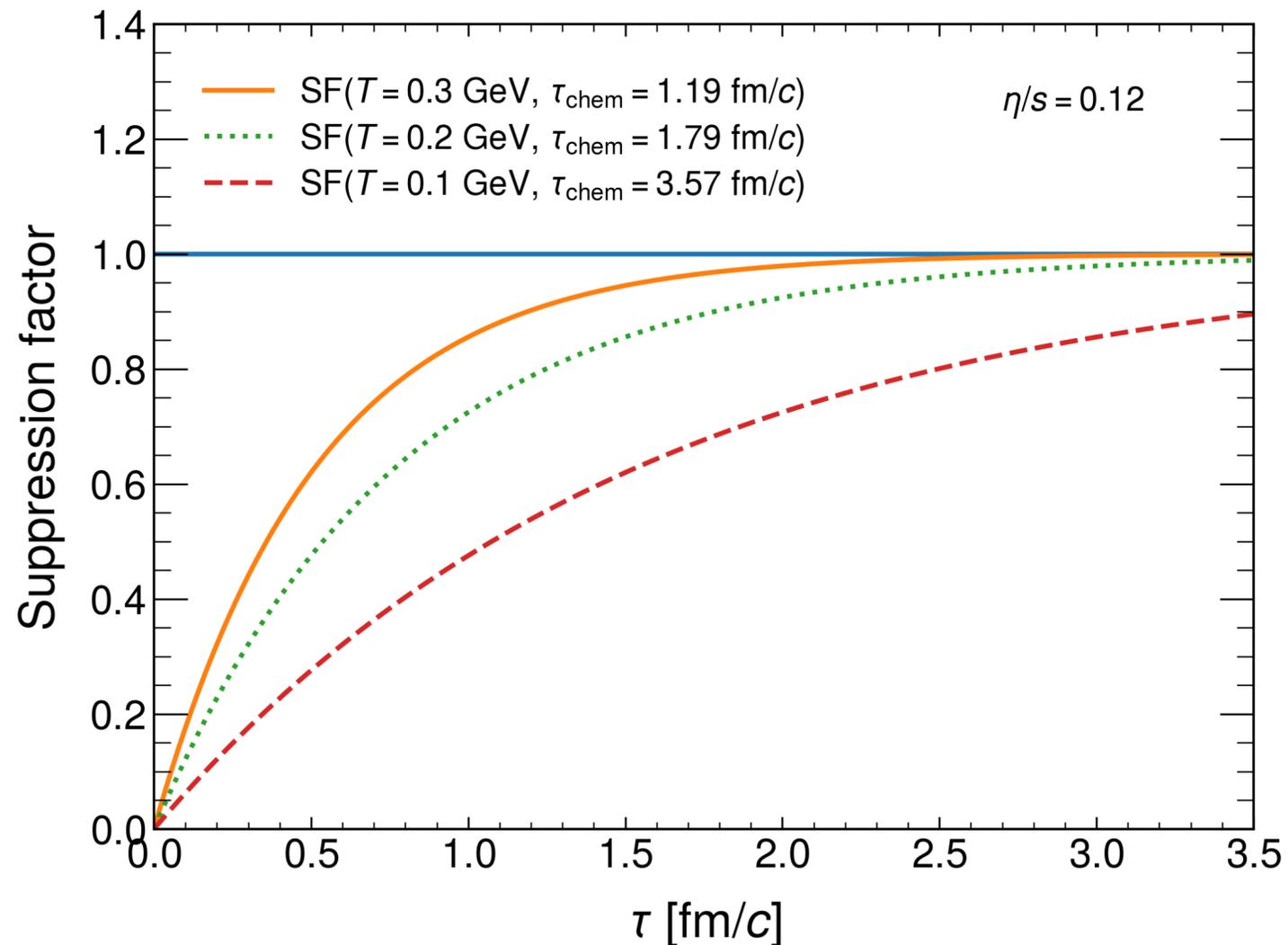
Parton distribution functions and fragmentation functions limits scale  $Q \sim 1\text{GeV}$ .

# Chemical Equilibrium

Introduce an effective suppression factor (SF) to include the process of the establishment of chemical equilibrium

$$\text{SF}(T, \tau) = 1 - e^{-A \frac{\tau}{\tau_R(T)}}$$

with chemical relaxation time  $\tau_R = 4\pi\eta/Ts$

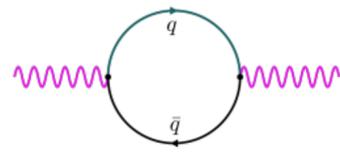


The chemical equilibrium is found to be  $\tau_{\text{chem}} = 1.2\tau_R$ . The constant A can be determined by requiring  $\text{SF}(\tau_{\text{chem}}) = 0.9$ .

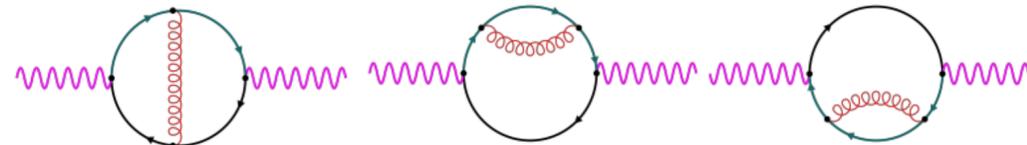
[Kurkela & Mazeliauskas, PRL 122, 142301 (2019), PRD 99, 054018 (2019)]

# How to combine NLO and LPM

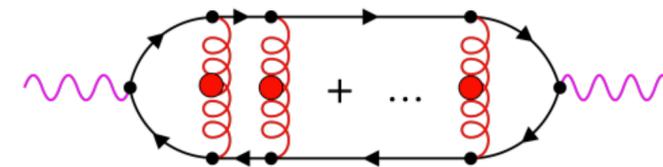
$$\rho_i|_{\text{NLO}}^{\text{resummed}} \equiv \rho_i|_{1\text{-loop}}^{\text{strict}} + \rho_i|_{2\text{-loop}}^{\text{strict}} + \left( \rho_i|_{\text{LPM}}^{\text{full}} - \rho_i|_{\text{LPM}}^{\text{expanded}} \right)$$



**LO**  
 $\mathcal{O}(1)$

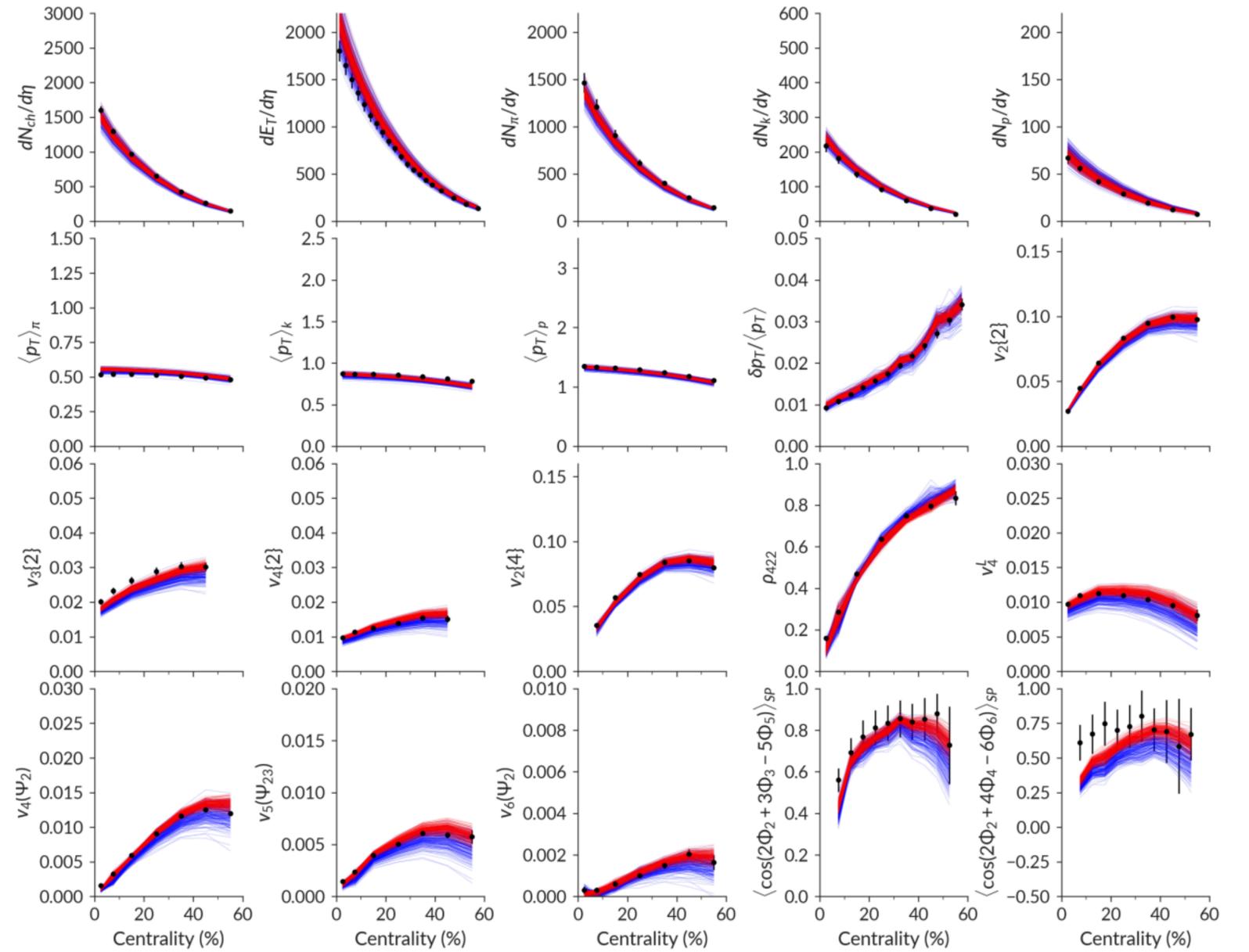
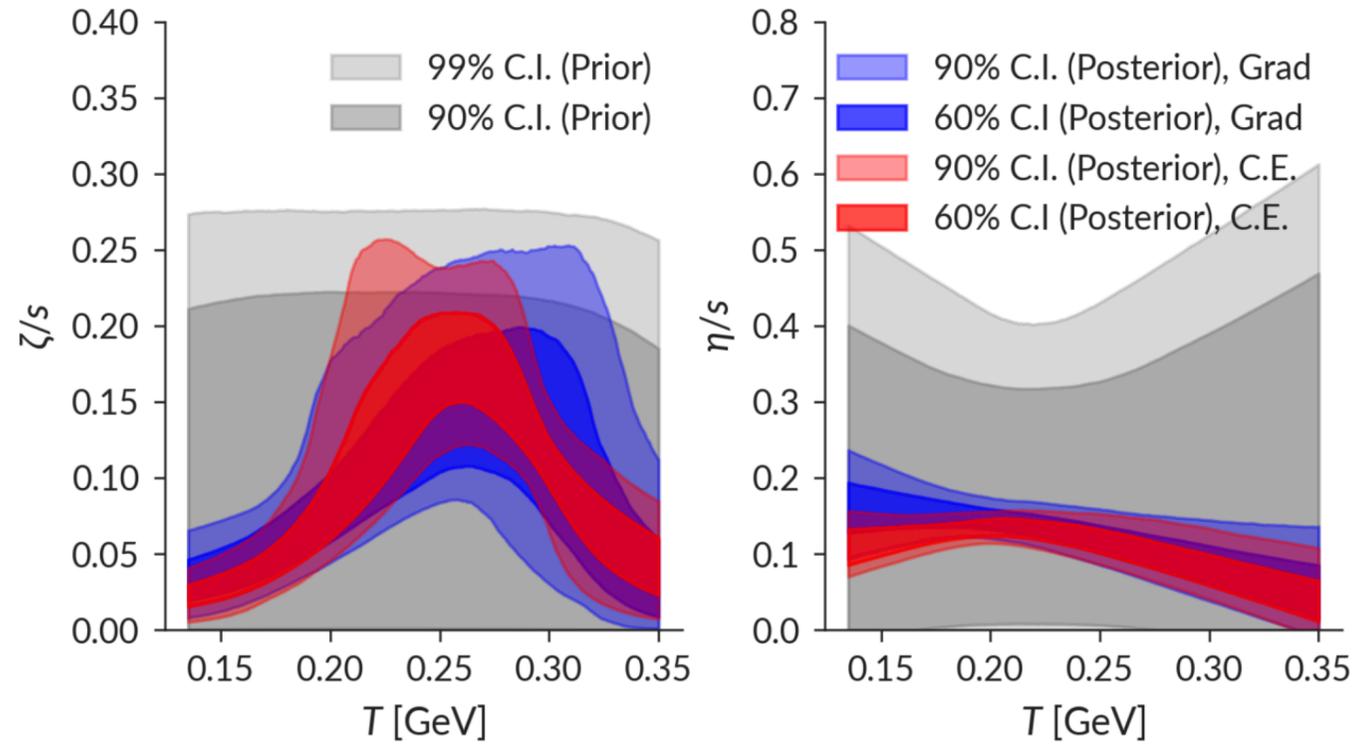


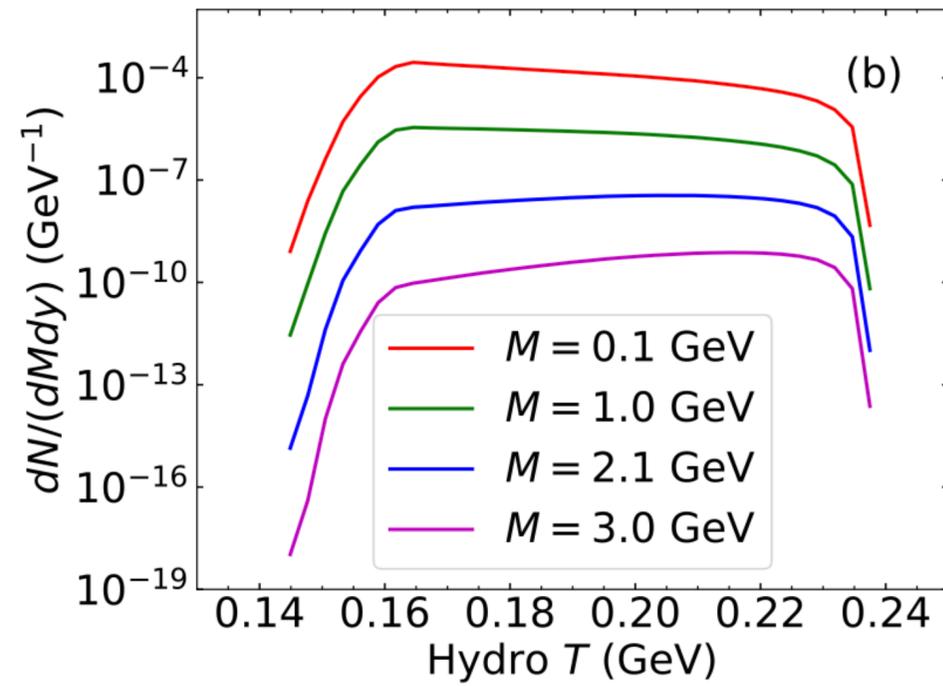
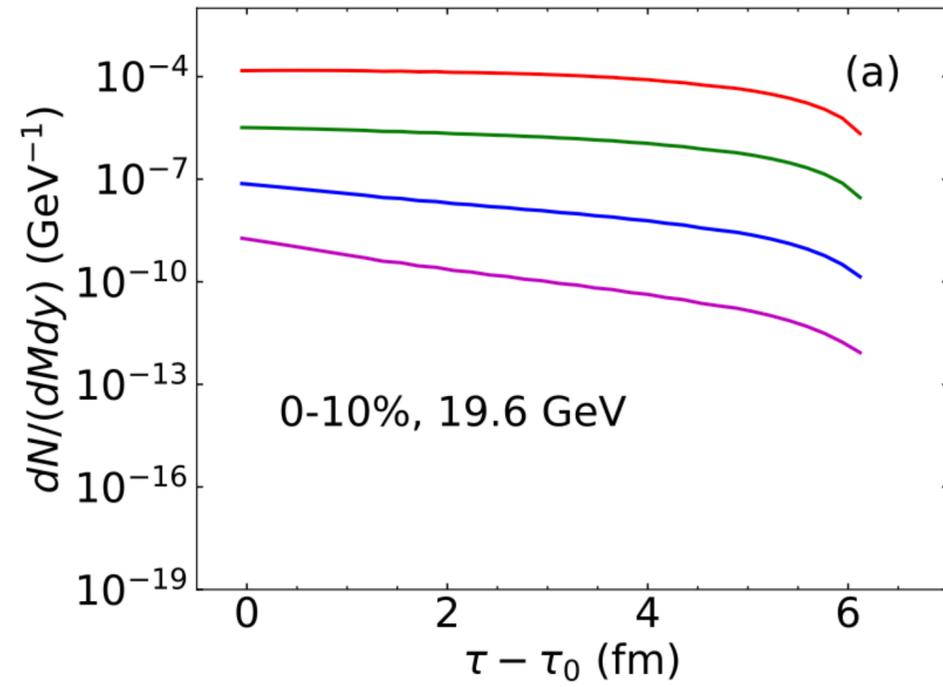
**NLO**  
 $\mathcal{O}(\alpha_s)$



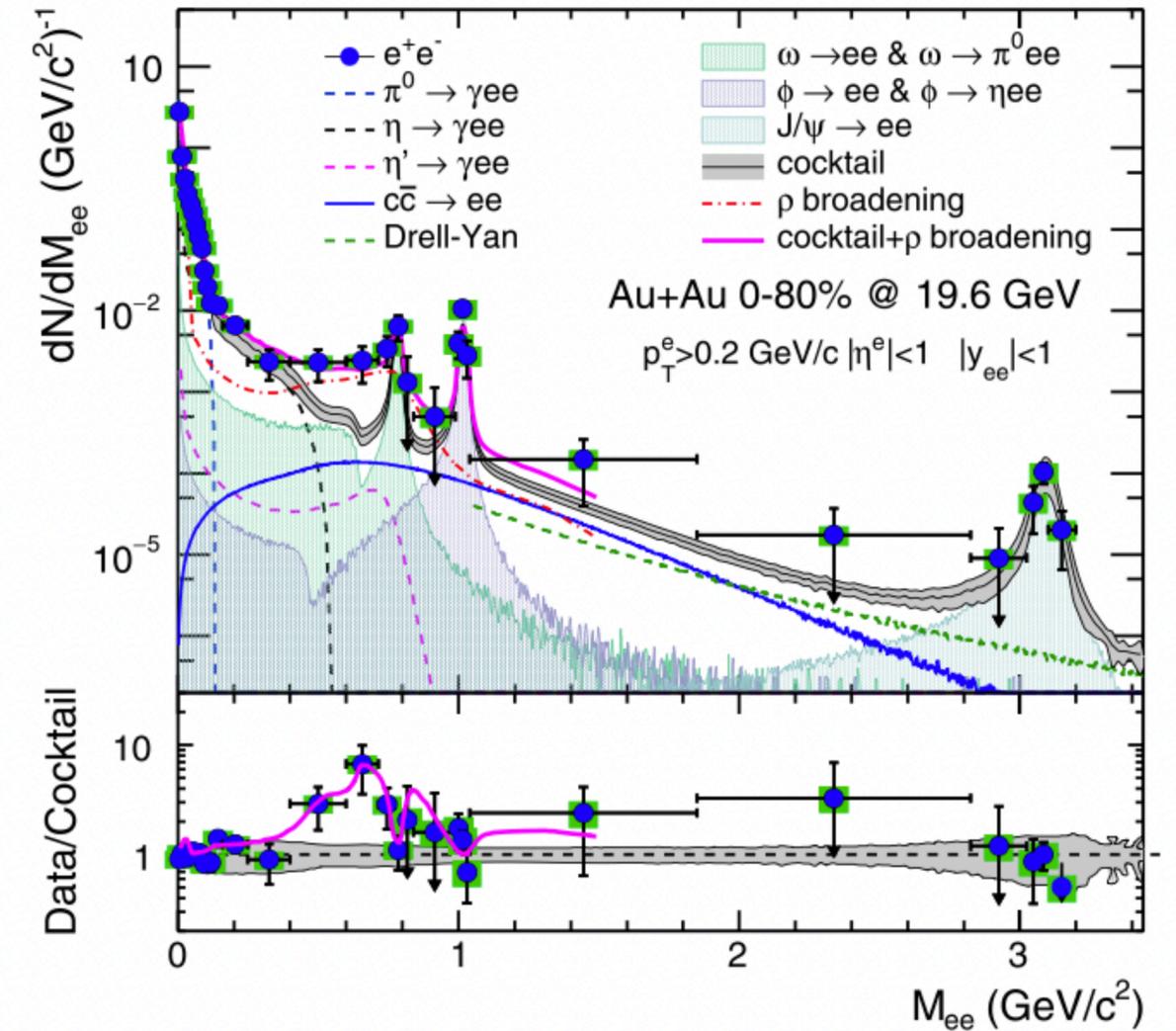
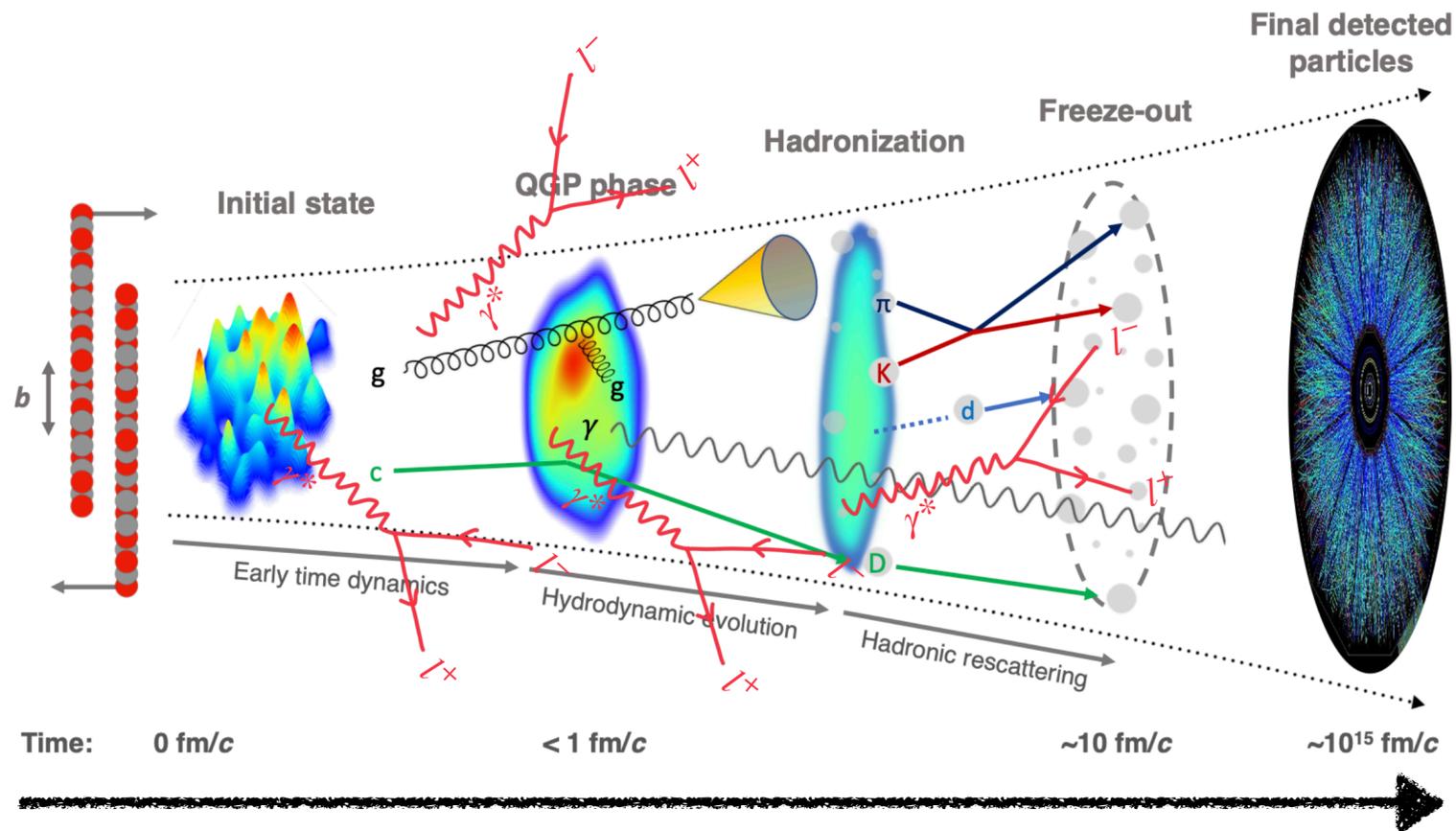
**LPM**  
 $\mathcal{O}(\alpha_s^i), i \geq 2$

# Model





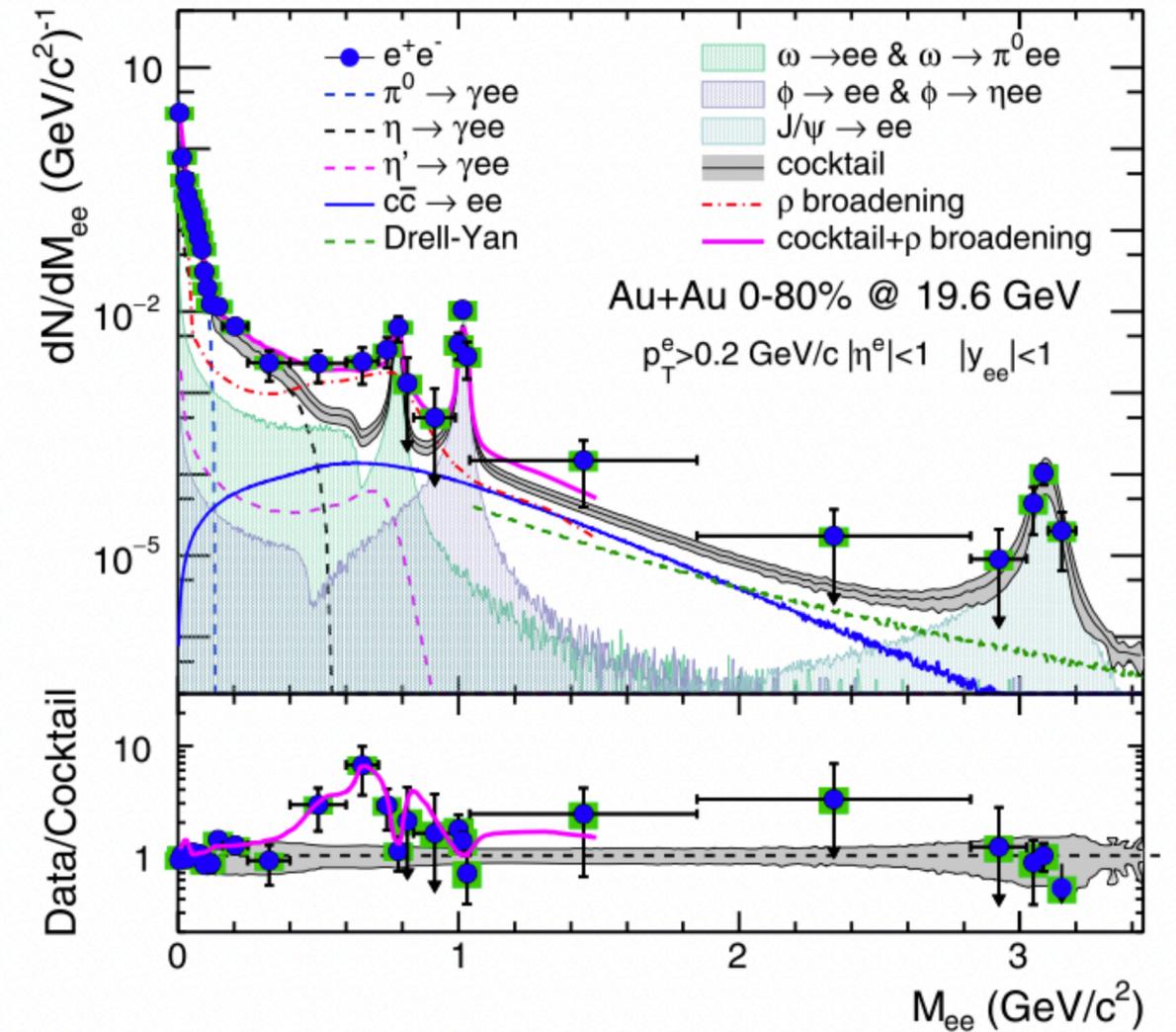
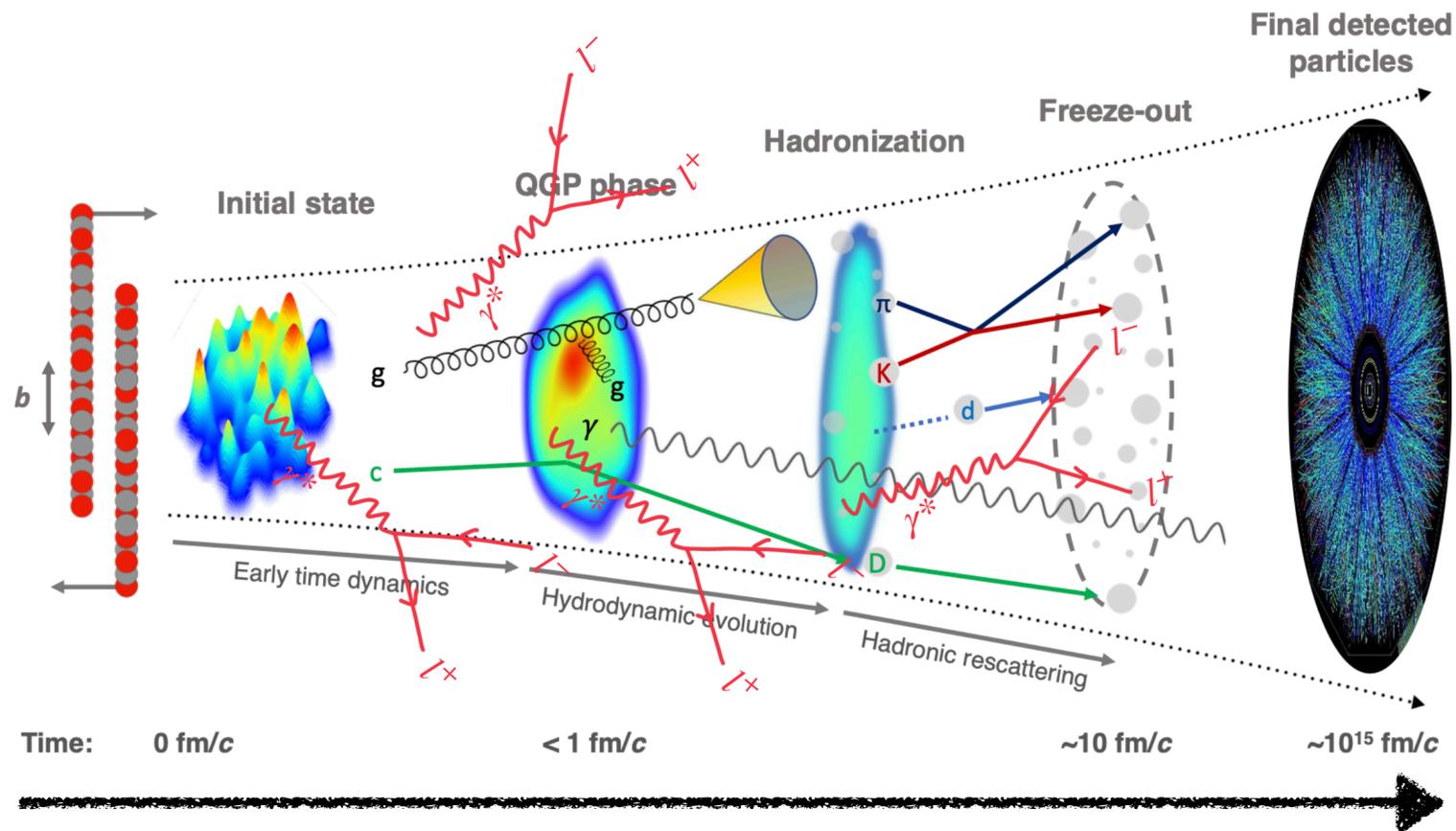
# A blessing and a curse



Dilepton from different stages are a blend together from experimental side.

- The low-mass region (LMR,  $0 < M < 1.1$  GeV): the decays of light mesons.
- The intermediate mass region (IMR,  $1.1 < M < 3$  GeV): thermal dilepton and decays of open heavy flavor.
- The high-mass region (HMR,  $M > 3$  GeV): Drell-Yan process and quarkonium decays.

# A blessing and a curse



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