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Semi-invisible $h \rightarrow Z X$ Higgs decays:

Extended scalar windows into dark matter



Based on: Aguilar-Saavedra, Cano, Cerdeño, JMN, *PRD106 115023 (2206.01214)*

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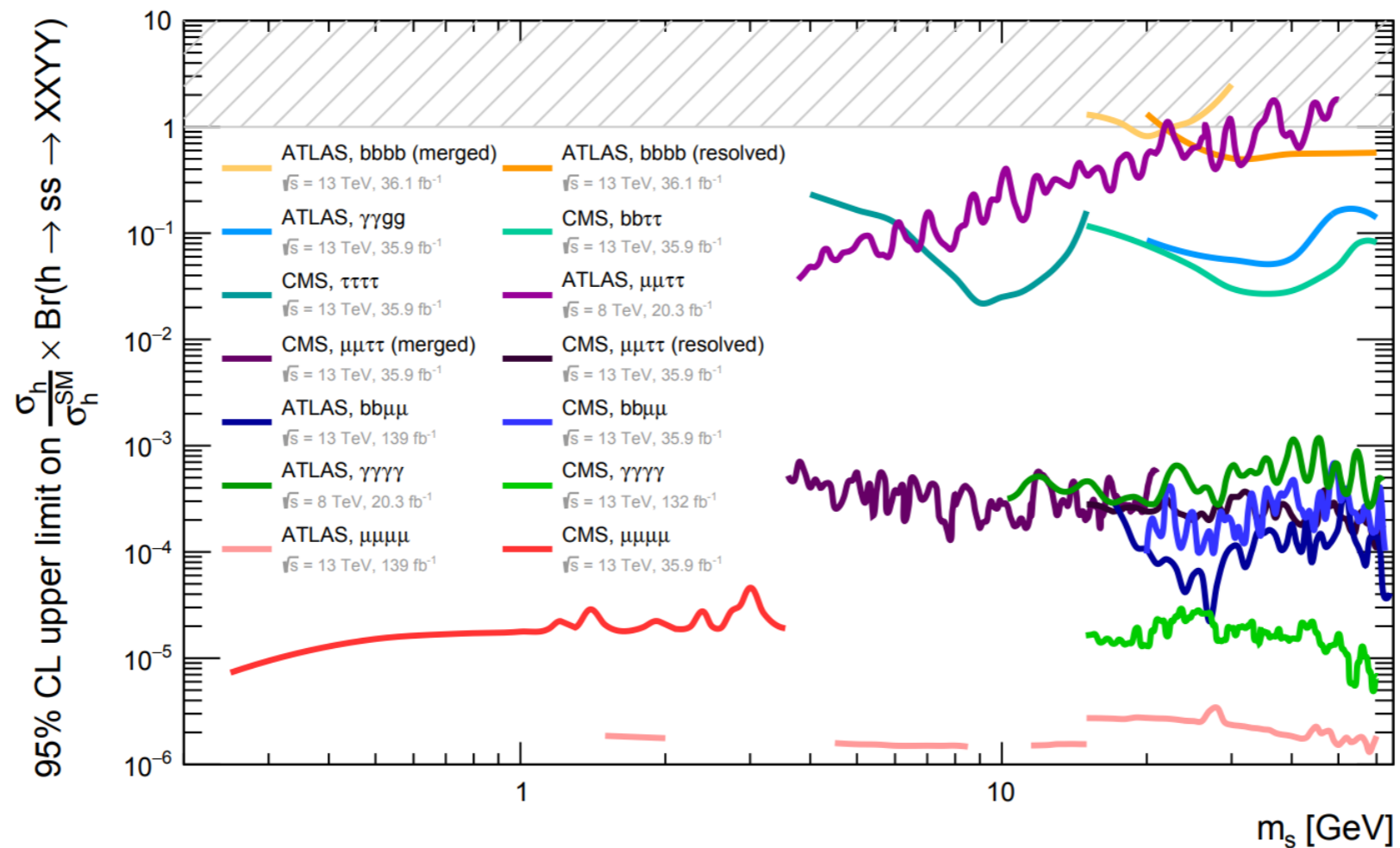
Why semi-invisible (exotic) Higgs boson decays?

Exotic Higgs decays

○ Higgs \rightarrow visible

e.g. $h \rightarrow ss \rightarrow XXYY$

“The Classic”



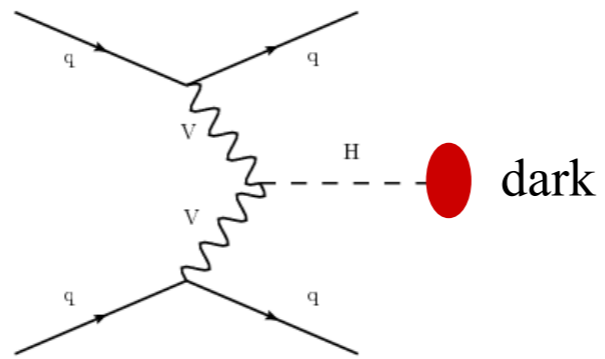
M. Cepeda et al. ([arXiv:2111.12751](https://arxiv.org/abs/2111.12751))

State-of-Art Exotic Higgs decay review
(focused on experiment)

Exotic Higgs decays

○ Higgs \rightarrow visible

○ Higgs \rightarrow invisible (**dark sector**) 



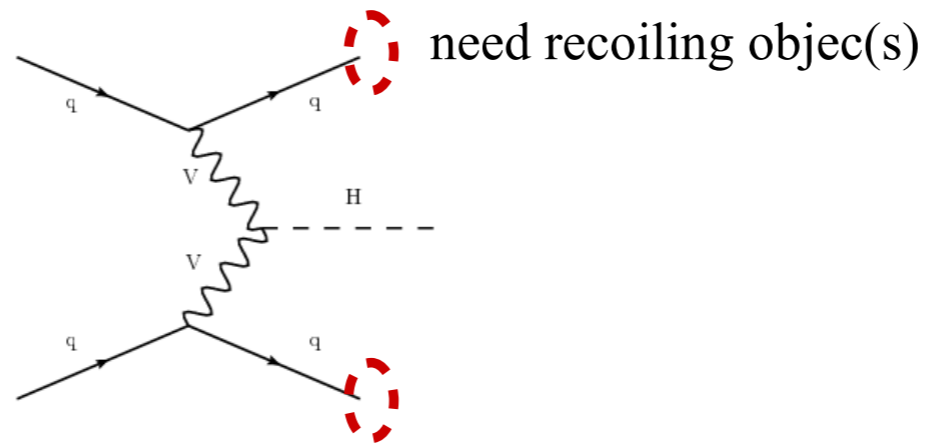
“The Classic”
(Dark Ed.)



Exotic Higgs decays

○ Higgs \rightarrow visible

○ Higgs \rightarrow invisible (dark sector)



Exotic Higgs decays

○ Higgs \rightarrow visible

○ Higgs \rightarrow invisible (dark sector)

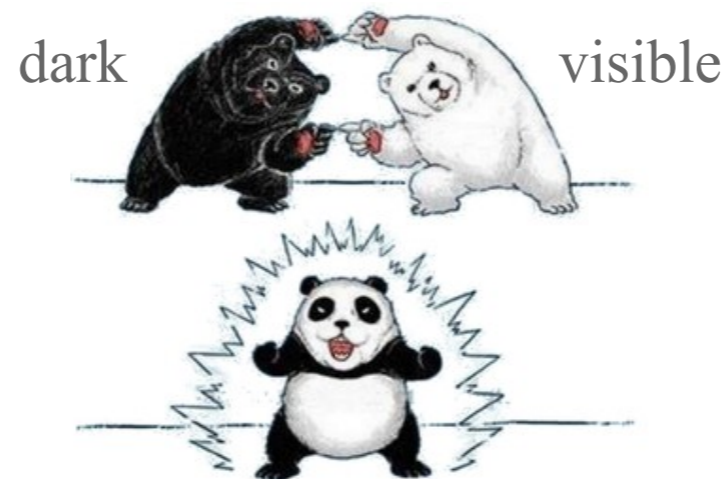
○ Higgs \rightarrow **semi**-invisible 

(visible part allows to access properties of dark sector)

Exotic Higgs decays

- Higgs \rightarrow visible
- Higgs \rightarrow invisible (dark sector)
- Higgs \rightarrow **semi**-invisible

(visible part allows to access properties of dark sector)



Critical information on Higgs interactions with dark sector

Exotic Higgs decays

○ Higgs \rightarrow visible

○ Higgs \rightarrow invisible

Decay	Mode	Reference	\sqrt{s} (TeV)	$\int \mathcal{L}$ (fb $^{-1}$)	m (GeV)	Interpretations
$h \rightarrow ss/aa/vv$						
$eeee$ (r)	ggF	CMS [79]	13	137	4-8, 11.5-62.5	SM+v, SM+ALP
(r)	ggF	ATLAS [80]	13	139	15-60	SM+s, SM+v
$ee\mu\mu$ (r)	ggF	CMS [79]	13	137	4-8, 11.5-62.5	SM+v, SM+ALP
(r)	ggF	ATLAS [80]	13	139	15-60	SM+v
$\mu\mu\mu\mu$ (m)	ggF	D0 [81]	1.96	4.2	0.2143-3	SM+s, SM+v
(r)	ggF	CMS [78]	13	35.9	0.25-8.5	SM+s, dark SUSY
(r)	ggF	CMS [79]	13	137	4-8, 11.5-60	SM+v, SM+ALP
(m/r)	ggF	ATLAS [80]	13	139	1.2-2, 4.4-8, 12-60	SM+s, SM+v
$\mu\mu\tau\tau$ (m/r)	ggF	D0 [81]	1.96	4.2	3.6-19	SM+s
(m/r)	ggF	ATLAS [82]	8	20.3	3.7-50	SM+s
(m/r)	ggF	CMS [83]	13	35.9	3.6-21	SM+s
(r)	ggF	CMS [84]	13	35.9	15-62.5	SM+s
$\tau\tau\tau\tau$ (m)	ggF	CMS [77]	13	35.9	4-15	SM+s
$bb\mu\mu$ (r)	ggF	ATLAS [85]	13	139	18-60	SM+s
(r)	ggF	CMS [86]	13	35.9	20-62.5	SM+s
$bb\tau\tau$ (r)	ggF	CMS [87]	13	35.9	15-60	SM+s
$bbbb$ (m)	Zh	ATLAS [88]	13	36.1	15-30	SM+s
(r)	Wh/Zh	ATLAS [76]	13	36.1	20-60	SM+s
$\gamma\gamma\gamma\gamma$ (r)	ggF	ATLAS [89]	8	20.3	10-62	SM+s
(r)	ggF	CMS [90]	13	132	15-60	SM+s
$\gamma\gamma gg$ (r)	VBF	ATLAS [91]	13	36.7	20-60	SM+s
$h \rightarrow Za/Zv$						
gg (m)	ggF	ATLAS [92]	13	139	0.5-4	SM+s
ss (m)	ggF	ATLAS [92]	13	139	1.5-3	SM+s
ee (r)	ggF	CMS [79]	13	137	4-8, 11.5-35	SM+v
(r)	ggF	ATLAS [80]	13	139	15-55	SM+v
$\mu\mu$ (r)	ggF	CMS [79]	13	137	4-8, 11.5-35	SM+v
(r)	ggF	ATLAS [80]	13	139	15-30/15-55	SM+s, SM+v

$h \rightarrow E_T^{\text{miss}}$					
Decay	Mode	Reference	\sqrt{s} (TeV)	$\int \mathcal{L}$ (fb $^{-1}$)	Br(H \rightarrow Inv) UL
E_T^{miss}	VBF	ATLAS [117]	13	139	0.13 (0.13)
	VBF	CMS [118]	13	138	0.17 (0.11)
	$Z(\ell)h$	ATLAS [111]	13	139	0.18 (0.18)
	$Z(\ell)h$	CMS [112]	13	137	0.29 (0.25)
	ggF	ATLAS [119]	13	139	0.34 (0.39)
	ggF, $V(qq)h$	CMS [120]	13	137	0.278 (0.253)
	tth	ATLAS [110]	13	139	0.40 (0.36)
	tth	CMS [121]	13	35.9	0.46 (0.48)
	Combination	ATLAS [110]	7, 8, 13	4.7+20.3+139	0.11 (0.11)
	Combination	CMS [122]	7, 8, 13	4.9+19.7+38.2	0.19 (0.15)

M. Cepeda et al. ([arXiv:2111.12751](https://arxiv.org/abs/2111.12751))

**Widely explored experimentally
(& theoretically!)**

(many Higgs production & decay modes covered)

Exotic Higgs decays

- Higgs \rightarrow visible
- Higgs \rightarrow invisible
- **Higgs \rightarrow semi-invisible** poorly explored so far...

Th.

Englert, Spannowsky, Wymant, Phys.Lett.B 718 (2012), 538

$h \rightarrow aa$ (jets + MET)

Petersson, Romagnoni, Torre, JHEP 10 (2012), 016

$h \rightarrow \gamma + \text{MET}$

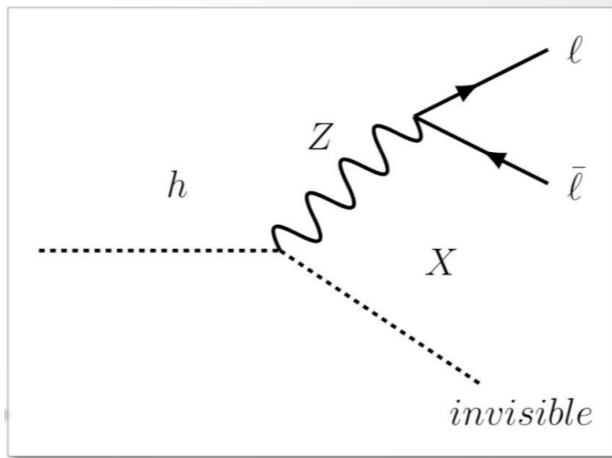
*Reviews by Curtin et al. (1312.4992),
Cepeda et al. (2111.12751) cover few
more channels/models:*

$bb + \text{MET}, \tau\tau + \text{MET}, \gamma\gamma + \text{MET}...$

Exp.

$h \rightarrow s/v + E_T^{\text{miss}}$					
Decay	Mode	Reference	\sqrt{s} (TeV)	$\int \mathcal{L}$ (fb $^{-1}$)	Interpretations
$E_T^{\text{miss}} + \gamma$	VBF	CMS [113]	13	130	SM+v
	VBF	ATLAS [114]	13	139	SM+v
	Zh	CMS [109]	13	137	SM+v
	ggF, Zh	CMS [115]	8	19.4	Other
$E_T^{\text{miss}} + bb$	Zh	ATLAS [116]	13	139	NMSSM

BSM Higgs boson decay $h \rightarrow ZX$ ($X \rightarrow invisible$)

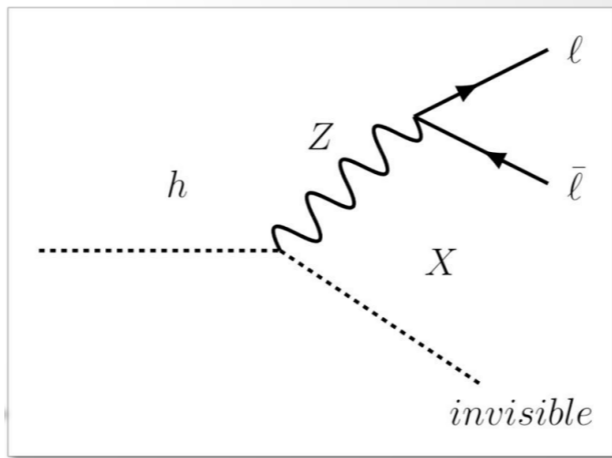


Aguilar-Saavedra, Cano, Cerdeño, No, 2206.01214

Decay channel previously unexplored in literature...

(... LHC $h \rightarrow ZX$ searches with visible X ($\ell\ell$ or $\gamma\gamma$) exist)

BSM Higgs boson decay $h \rightarrow ZX$ ($X \rightarrow invisible$)

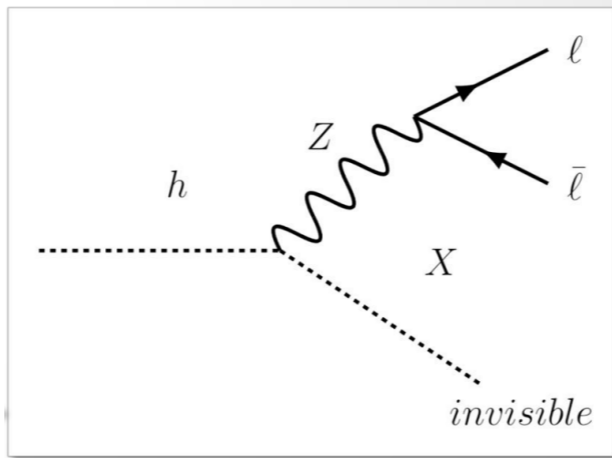


Aguilar-Saavedra, Cano, Cerdeño, No, 2206.01214

Decay channel previously unexplored in literature...

- Model-independent sensitivity analysis for **HL-LHC**
- BSM model interpretations:
 - Extended Higgs sector portal to DM
 - Axion-like particles (ALPs)

BSM Higgs boson decay $h \rightarrow ZX$ ($X \rightarrow invisible$)



Aguilar-Saavedra, Cano, Cerdeño, No, 2206.01214

Decay channel previously unexplored in literature...

- Model-independent sensitivity analysis for HL-LHC
- BSM model interpretations:

**Extended Higgs sector portal to DM
Axion-like particles (ALPs)**

$$h \rightarrow Z (\ell\ell) + invisible$$

○ Already present in SM: $h \rightarrow ZZ^* \rightarrow \ell\ell \nu\nu$

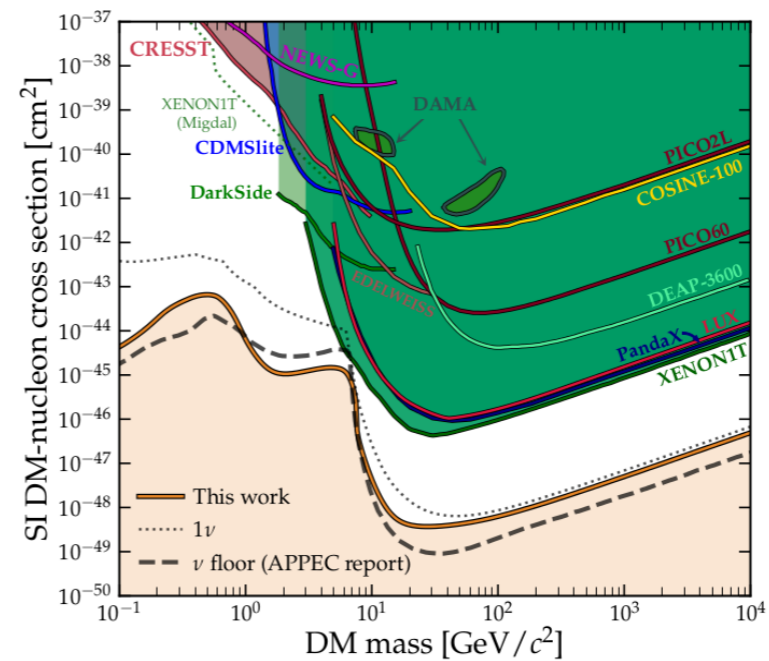
Irreducible SM background for
BSM $h \rightarrow Z (\ell\ell) + invisible$

$$h \rightarrow Z (\ell\ell) + \textit{invisible}$$

○ Already present in SM: $h \rightarrow ZZ^* \rightarrow \ell\ell \nu\nu$

“SM Higgs **neutrino floor**” for
BSM $h \rightarrow Z (\ell\ell) + \textit{invisible}$

“in **analogy** to the neutrino floor/fog
(from coherent neutrino-nucleus
scattering) in DM direct detection
experiments”



O' Hare, 2109.03116

$$h \rightarrow Z (\ell\ell) + \textit{invisible}$$

○ Already present in SM: $h \rightarrow ZZ^* \rightarrow \ell\ell \nu\nu$

SM “Higgs neutrino floor” for
BSM $h \rightarrow Z (\ell\ell) + \textit{invisible}$

Target sensitivity for HL-LHC

$$\mathbf{Br}_{h \rightarrow Z + \textit{invisible}} \approx \mathbf{0.0053}$$

$$h \rightarrow Z (\ell\ell) + \text{invisible}$$

○ Which Higgs production mode @LHC?

ggF ($pp \rightarrow 2\ell + \text{MET}$) **X** Too much background

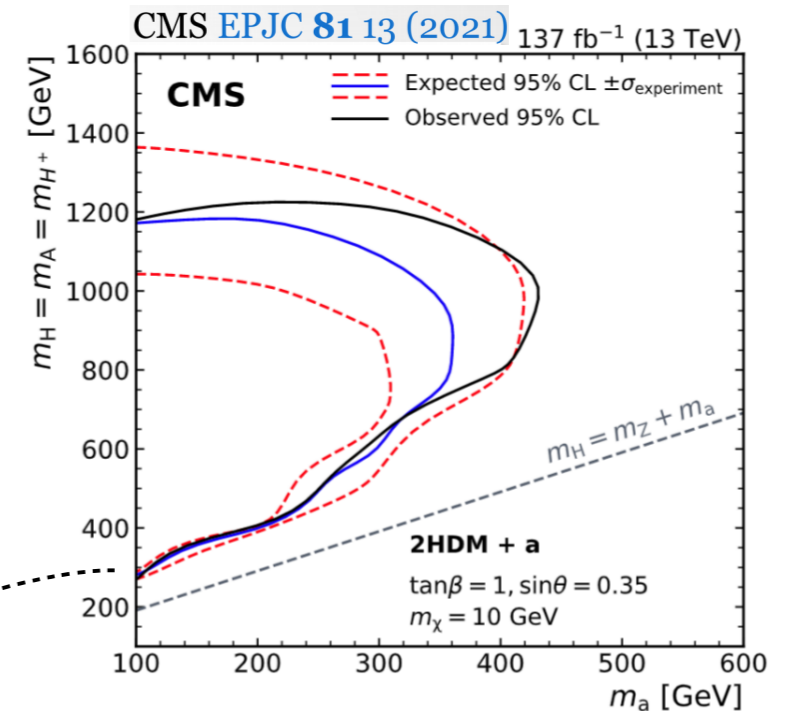
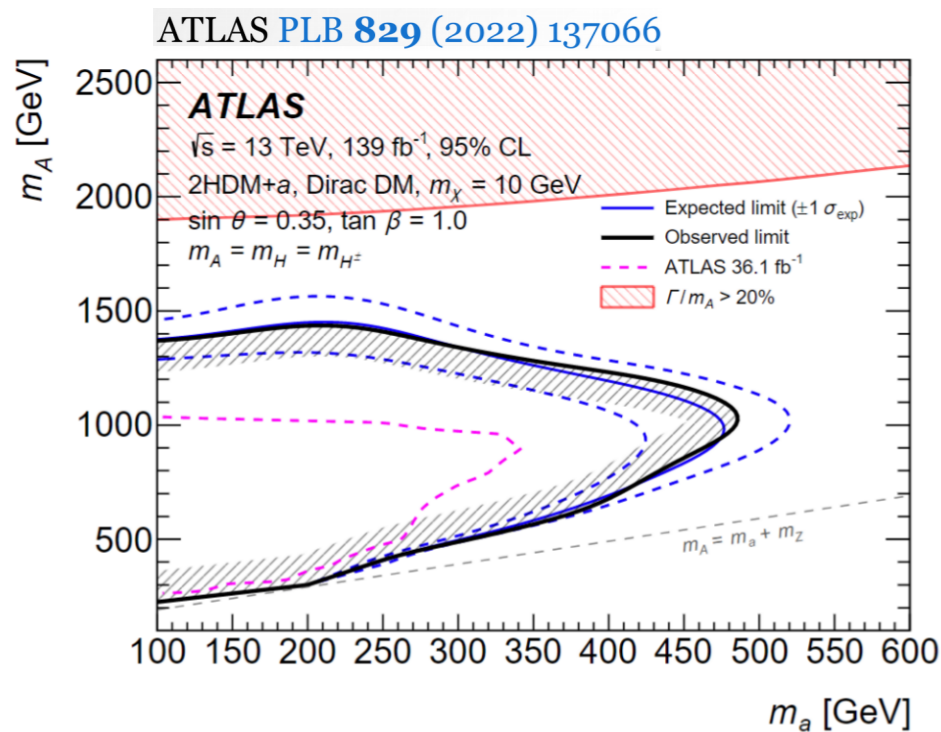
$h \rightarrow Z (\ell\ell) + invisible$

○ Which Higgs production mode @LHC?

ggF ($pp \rightarrow 2\ell + MET$) ✗ Too much background

e.g. Seen explicitly in ATLAS/CMS searches for “mono-Z”

ggF $\rightarrow H, H \rightarrow Z + a$ (invisible)



$E_T^{\text{miss}} > 90 \text{ GeV}$
 $m_T > 200 \text{ GeV}$

insensitive to $m_H = m_h = 125 \text{ GeV}$

$$h \rightarrow Z (\ell\ell) + \text{invisible}$$

○ Which Higgs production mode @LHC?

ggF ($pp \rightarrow 2\ell + \text{MET}$) ✗ Too much background

VBF ($pp \rightarrow 2\ell + \text{MET} + 2j$) ✗ Too much background?

Zh ($pp \rightarrow 4\ell + \text{MET}$) ✓ very clean

$$pp \rightarrow Zh, h \rightarrow Z(\ell\ell) + \text{invisible}$$

Our initial event selection:

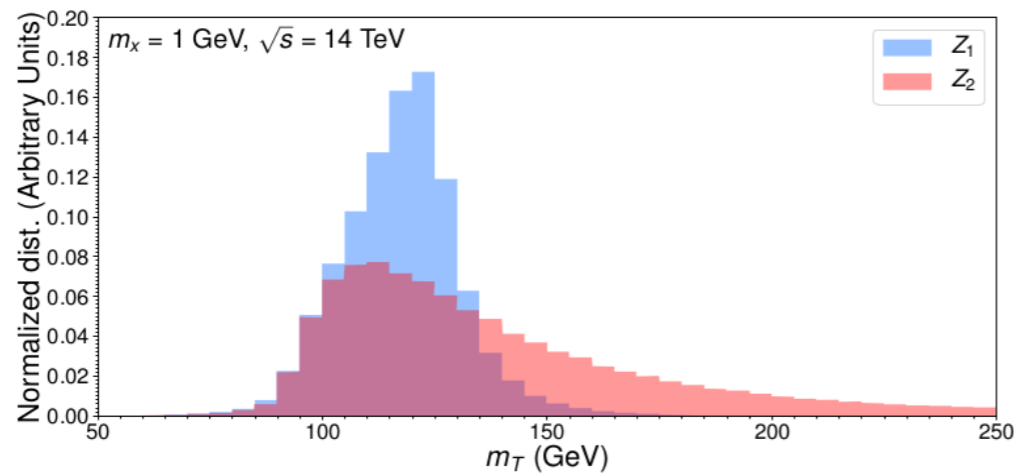
- Two (SF,OS) lepton pairs
- Pass 1-, 2- or 3-lepton ATLAS (2018) Triggers
- Veto high- p_T jets ($p_T > 50$ GeV)
- If multiple lepton combinations in event, minimize Δ

$$\Delta^2 = m_Z^{-2} [(m_{\ell\ell_1} - m_Z)^2 + (m_{\ell\ell_2} - m_Z)^2]$$

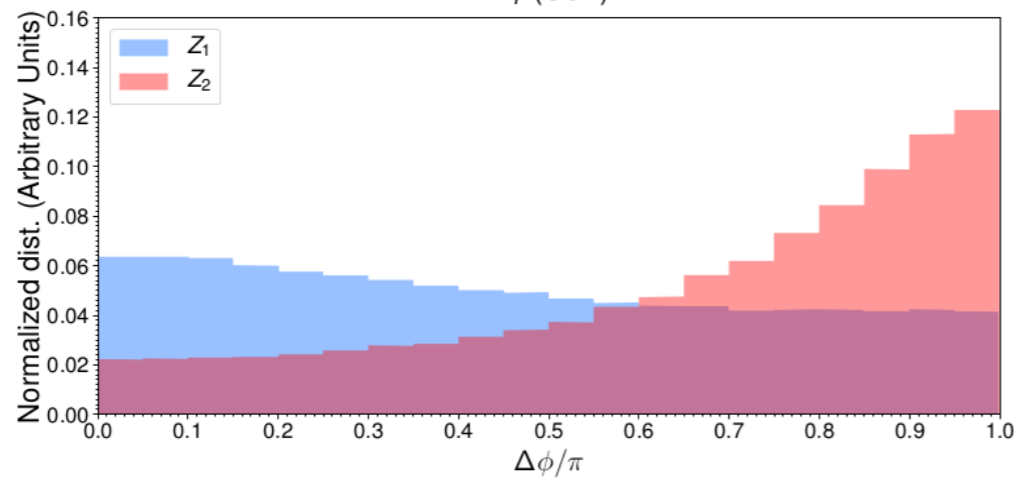
$$pp \rightarrow Z h, h \rightarrow Z (\ell\ell) + \text{invisible}$$

- ❶ Small Cross-section: $\sigma(pp \rightarrow 4\ell + \text{inv}) \simeq 3.8 \text{ fb} \times \text{BR}(h \rightarrow Z + \text{inv})$ (@13 TeV LHC)
- ❷ How to choose Z boson “from Higgs decay”? (Z_1)

Discriminating kinematic variables



$$M_T^2 = \left(\sqrt{M_Z^2 + |\vec{p}_T^Z|^2} + \cancel{E}_T \right)^2 - \left| \vec{p}_T^Z + \vec{\cancel{E}}_T \right|^2$$

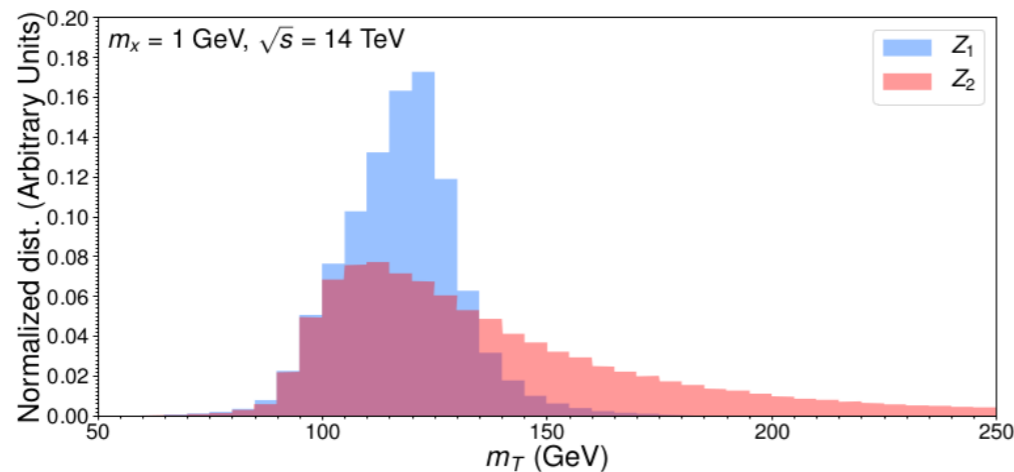


$$\Delta\phi_{\ell\ell, \cancel{E}_T}$$

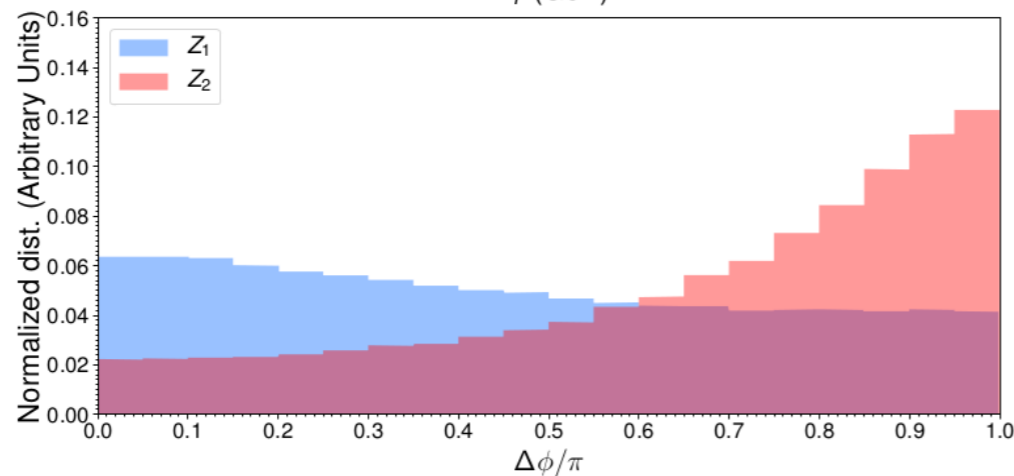
$pp \rightarrow Z h, h \rightarrow Z (\ell\ell) + invisible$

- ❶ Small Cross-section: $\sigma(pp \rightarrow 4\ell + inv) \simeq 3.8 \text{ fb} \times \text{BR}(h \rightarrow Z + inv)$ (@13 TeV LHC)
- ❷ How to choose Z boson “from Higgs decay”? (Z_1)

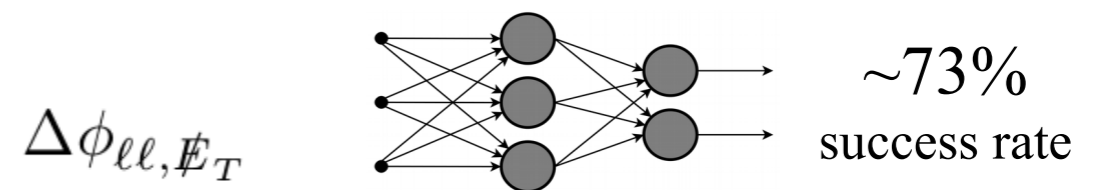
Discriminating kinematic variables



$$M_T^2 = \left(\sqrt{M_Z^2 + |\vec{p}_T^Z|^2} + \cancel{E}_T \right)^2 - \left| \vec{p}_T^Z + \vec{\cancel{E}}_T \right|^2$$



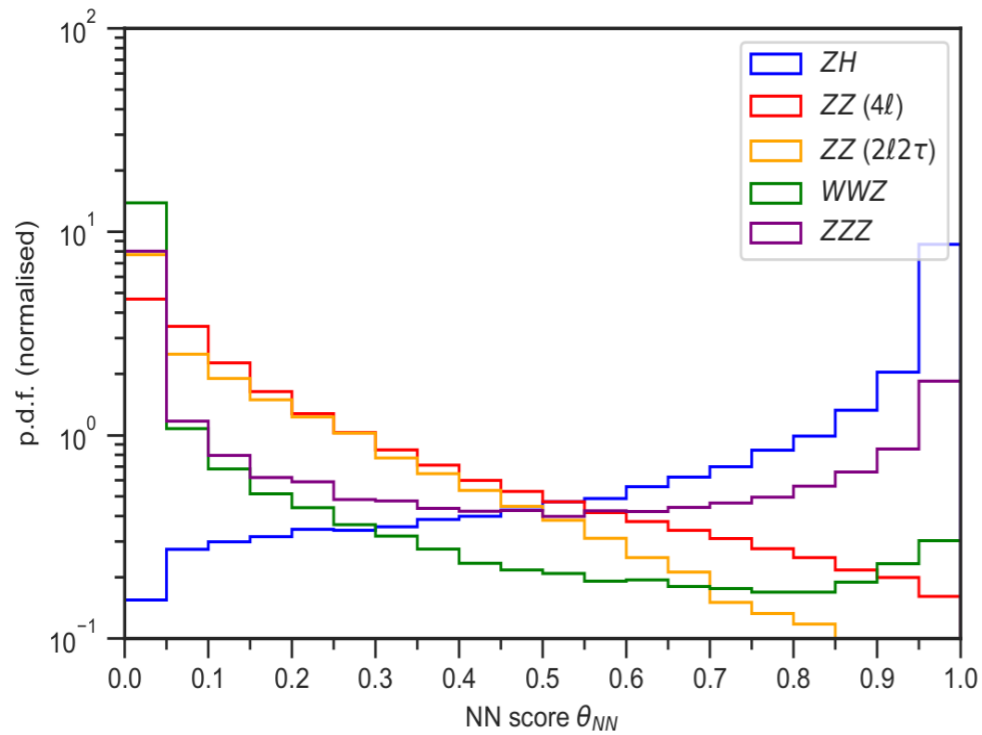
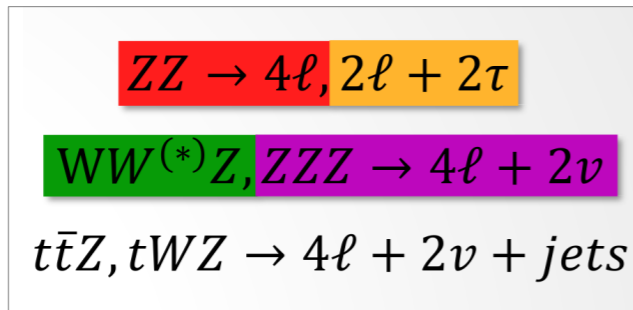
We build a Neural Network to select Z_1 based on kinematics (2 variables)



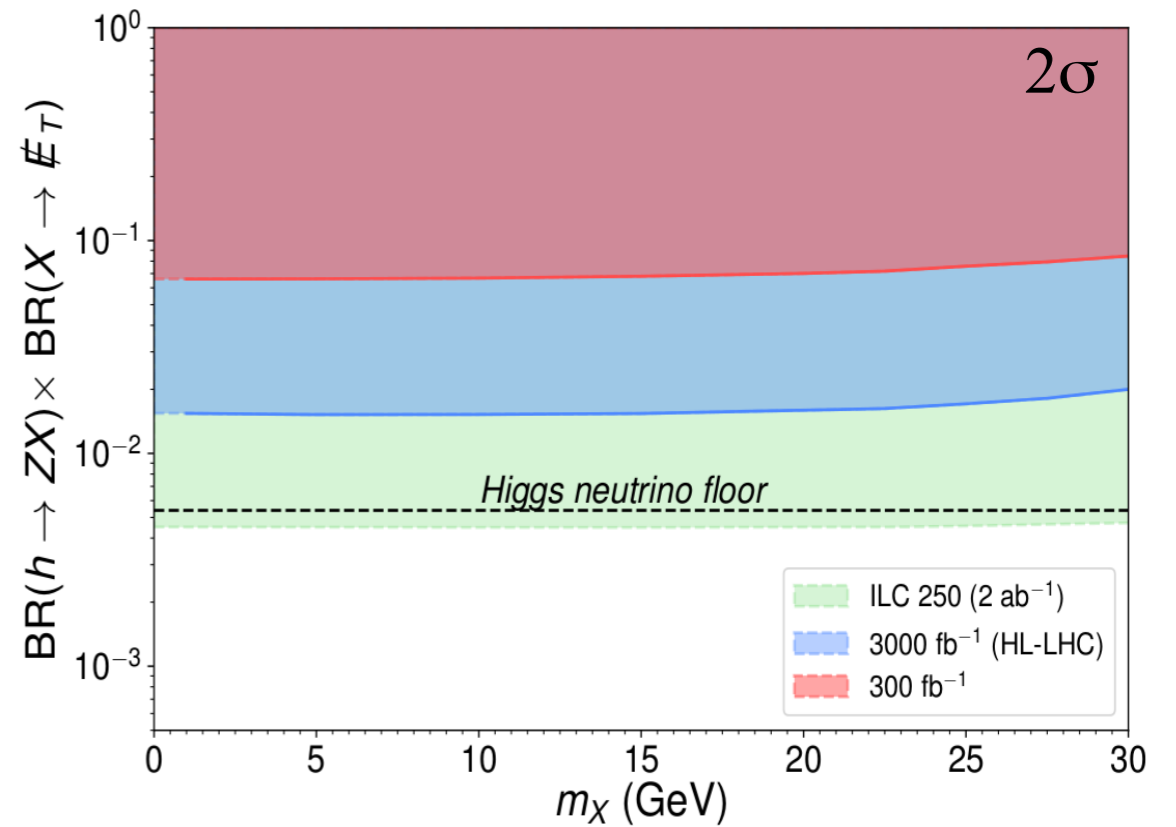
$$pp \rightarrow Zh, h \rightarrow Z(\ell\ell) + \text{invisible}$$

(HL-)LHC sensitivity:

SM backgrounds



Multivariate (NN) analysis



$BR(h \rightarrow ZX) \sim 1\text{-}2\%$ reachable @HL-LHC

BSM Models

Axion-like particles (ALPs)

- ALP may have:

ALP - Higgs interactions

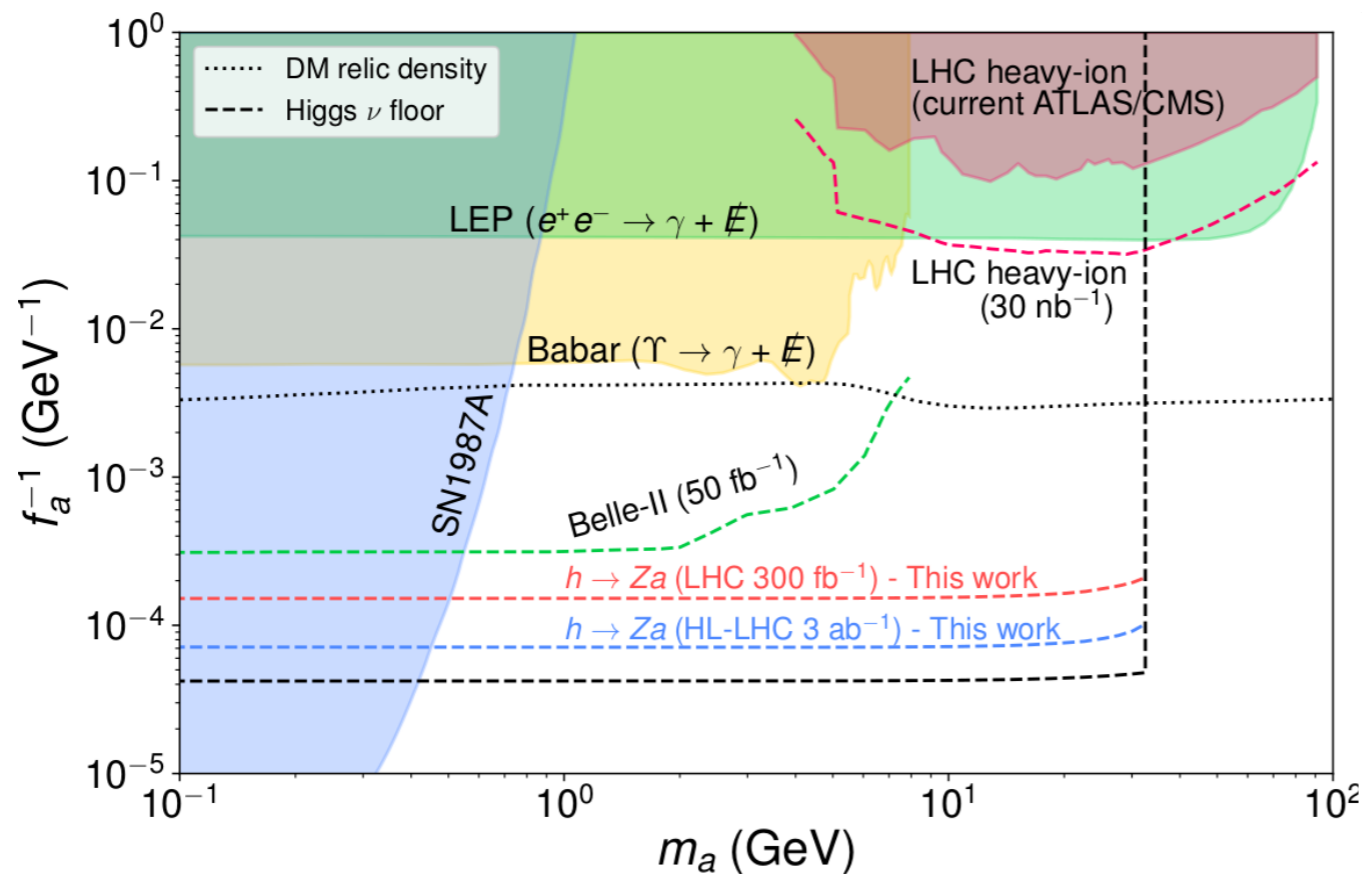
Brivio, Gavela, Merlo, Mimasu, No, del Rey, Sanz, *EJPC* 77 (2017) 8, 572
 Bauer, Neubert, Thamm, *JHEP* 12 (2017), 044



ALP - Dark sector interactions

Dolan, Ferber, Hearty, Kahlhoefer, Schmidt-Hoberg, *JHEP* 12 (2017) 094

(dark decay of ALP) “ALP portal to DM”



$$(y_\chi = 1, c_{aZh} = 1, m_\chi = 0.45 m_a)$$

BSM ALP setup

$$\Gamma(h \rightarrow Za) = (m_h^3 / 16\pi f_a^2) c_{aZh}^2 \lambda^{3/2}$$

$$\text{ALP-DM: } y_\chi \bar{\chi} \gamma^\mu \gamma^5 \chi \partial_\mu a / f_a$$

$$\text{ALP-Photons: } c_{a\gamma\gamma} / f_a a F^{\mu\nu} \tilde{F}_{\mu\nu}$$

$$c_{a\gamma\gamma} \sim \alpha_{EM}$$

BSM Models

2HDM + a \rightarrow “Pseudoscalar portal to DM”

$$\begin{aligned} V = & V_{2\text{HDM}} + \frac{\mu_{a_0}^2}{2} a_0^2 + \frac{\lambda_a}{4} a_0^4 + i \kappa a_0 H_1^\dagger H_2 + \text{h.c.} \\ & + \lambda_{a1} a_0^2 |H_1|^2 + \lambda_{a2} a_0^2 |H_2|^2 + m_\chi \bar{\chi} \chi \\ & + y_\chi a_0 \bar{\chi} i \gamma^5 \chi \end{aligned}$$

$$\begin{aligned} V_{2\text{HDM}} = & \mu_1^2 |H_1|^2 + \mu_2^2 |H_2|^2 - \mu^2 [H_1^\dagger H_2 + \text{h.c.}] \\ & + \frac{\lambda_1}{2} |H_1|^4 + \frac{\lambda_2}{2} |H_2|^4 + \lambda_3 |H_1|^2 |H_2|^2 \\ & + \lambda_4 |H_1^\dagger H_2|^2 + \frac{\lambda_5}{2} [(H_1^\dagger H_2)^2 + \text{h.c.}] \end{aligned}$$

BSM Models

2HDM + a \rightarrow “Pseudoscalar portal to DM”

Pseudoscalar mediator

$$\begin{aligned}
 V = & V_{2\text{HDM}} + \frac{\mu_{a_0}^2}{2} a_0^2 + \frac{\lambda_a}{4} a_0^4 + i \kappa a_0 H_1^\dagger H_2 + \text{h.c.} \\
 & + \lambda_{a_1} a_0^2 |H_1|^2 + \lambda_{a_2} a_0^2 |H_2|^2 + m_\chi \bar{\chi} \chi \\
 & + y_\chi a_0 \bar{\chi} i \gamma^5 \chi
 \end{aligned}$$

Dirac Fermion DM

$$\begin{aligned}
 V_{2\text{HDM}} = & \mu_1^2 |H_1|^2 + \mu_2^2 |H_2|^2 - \mu^2 [H_1^\dagger H_2 + \text{h.c.}] \\
 & + \frac{\lambda_1}{2} |H_1|^4 + \frac{\lambda_2}{2} |H_2|^4 + \lambda_3 |H_1|^2 |H_2|^2 \\
 & + \lambda_4 |H_1^\dagger H_2|^2 + \frac{\lambda_5}{2} [(H_1^\dagger H_2)^2 + \text{h.c.}]
 \end{aligned}$$

BSM Models

2HDM + a \rightarrow “Pseudoscalar portal to DM”

Singlet-doublet mixing

$$V = V_{2\text{HDM}} + \frac{\mu_{a_0}^2}{2} a_0^2 + \frac{\lambda_a}{4} a_0^4 + i \kappa a_0 H_1^\dagger H_2 + \text{h.c.}$$

$$+ \lambda_{a1} a_0^2 |H_1|^2 + \lambda_{a2} a_0^2 |H_2|^2 + m_\chi \bar{\chi} \chi$$

$$+ y_\chi a_0 \bar{\chi} i \gamma^5 \chi$$

Portal coupling

$$V_{2\text{HDM}} = \mu_1^2 |H_1|^2 + \mu_2^2 |H_2|^2 - \mu^2 [H_1^\dagger H_2 + \text{h.c.}]$$

$$+ \frac{\lambda_1}{2} |H_1|^4 + \frac{\lambda_2}{2} |H_2|^4 + \lambda_3 |H_1|^2 |H_2|^2$$

$$+ \lambda_4 |H_1^\dagger H_2|^2 + \frac{\lambda_5}{2} [(H_1^\dagger H_2)^2 + \text{h.c.}]$$

BSM Models

2HDM + a \rightarrow “Pseudoscalar portal to DM”

$$V = V_{2\text{HDM}} + \frac{\mu_{a_0}^2}{2} a_0^2 + \frac{\lambda_a}{4} a_0^4 + i \kappa a_0 H_1^\dagger H_2 + \text{h.c.} \\ + \lambda_{a1} a_0^2 |H_1|^2 + \lambda_{a2} a_0^2 |H_2|^2 + m_\chi \bar{\chi} \chi \\ + y_\chi a_0 \bar{\chi} i \gamma^5 \chi$$

$$V_{2\text{HDM}} = \mu_1^2 |H_1|^2 + \mu_2^2 |H_2|^2 - \mu^2 [H_1^\dagger H_2 + \text{h.c.}] \\ + \frac{\lambda_1}{2} |H_1|^4 + \frac{\lambda_2}{2} |H_2|^4 + \lambda_3 |H_1|^2 |H_2|^2 \\ + \lambda_4 |H_1^\dagger H_2|^2 + \frac{\lambda_5}{2} [(H_1^\dagger H_2)^2 + \text{h.c.}]$$

Why?

DM-SM interactions via pseudoscalar mediator avoid stringent bounds from DM direct detection experiments (e.g. XENON 1T)

Simplest renormalizable model: **2HDM + a**

Ipek, McKeen, Nelson, PRD 90 (2014), 055021

No, PRD 93 (2016), 031701

Goncalves, Machado, No, PRD 95 (2017), 055027

Bauer, Haisch, Kahlhoefer, JHEP 05 (2017), 138

[LHC DM WG Benchmark Model]

Abe et al, Phys. Dark. Univ. 27 (2020), 100351

BSM Models

2HDM + a \rightarrow “Pseudoscalar portal to DM”

$$V = V_{2\text{HDM}} + \frac{\mu_{a_0}^2}{2} a_0^2 + \frac{\lambda_a}{4} a_0^4 + i \kappa a_0 H_1^\dagger H_2 + \text{h.c.} \\ + \lambda_{a1} a_0^2 |H_1|^2 + \lambda_{a2} a_0^2 |H_2|^2 + m_\chi \bar{\chi} \chi \\ + y_\chi a_0 \bar{\chi} i \gamma^5 \chi$$

$$V_{2\text{HDM}} = \mu_1^2 |H_1|^2 + \mu_2^2 |H_2|^2 - \mu^2 [H_1^\dagger H_2 + \text{h.c.}] \\ + \frac{\lambda_1}{2} |H_1|^4 + \frac{\lambda_2}{2} |H_2|^4 + \lambda_3 |H_1|^2 |H_2|^2 \\ + \lambda_4 |H_1^\dagger H_2|^2 + \frac{\lambda_5}{2} [(H_1^\dagger H_2)^2 + \text{h.c.}]$$

$$\Gamma(h \rightarrow Za) = \frac{1}{16\pi} \sin^2 \theta \cos^2 (\beta - \alpha) \frac{m_h^3}{v^2} \lambda^{3/2}$$

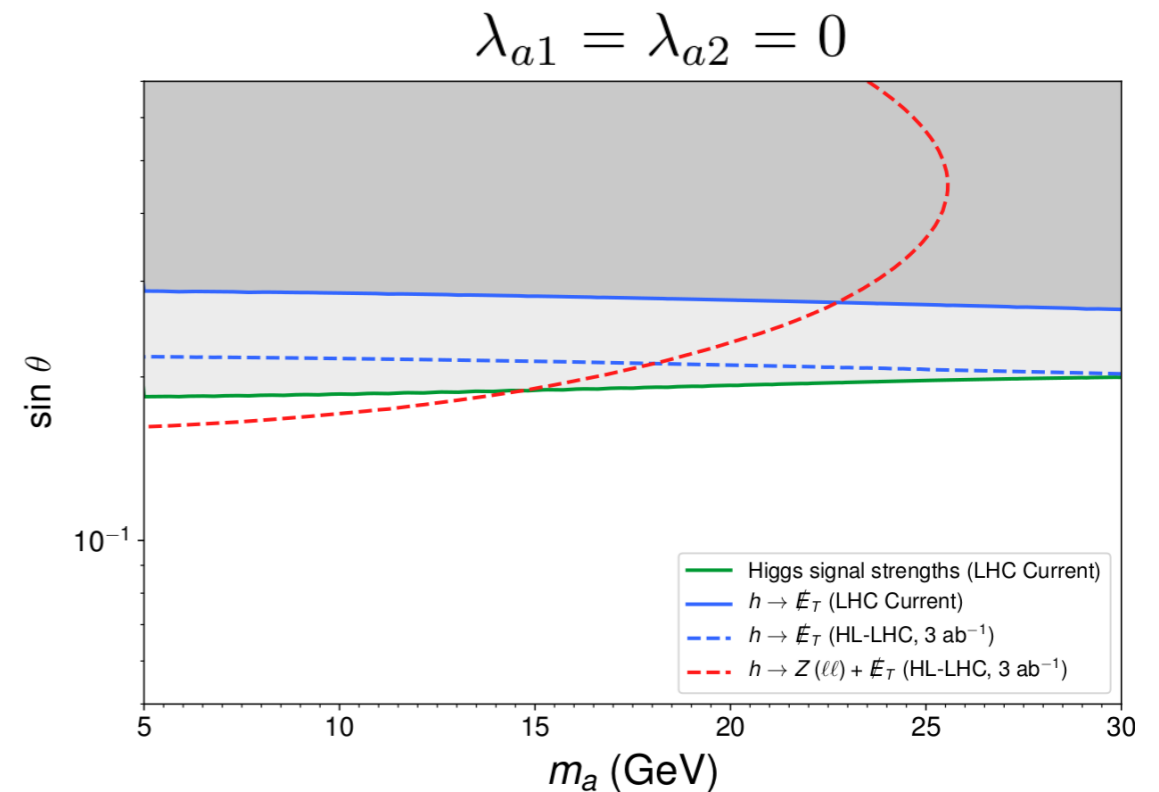
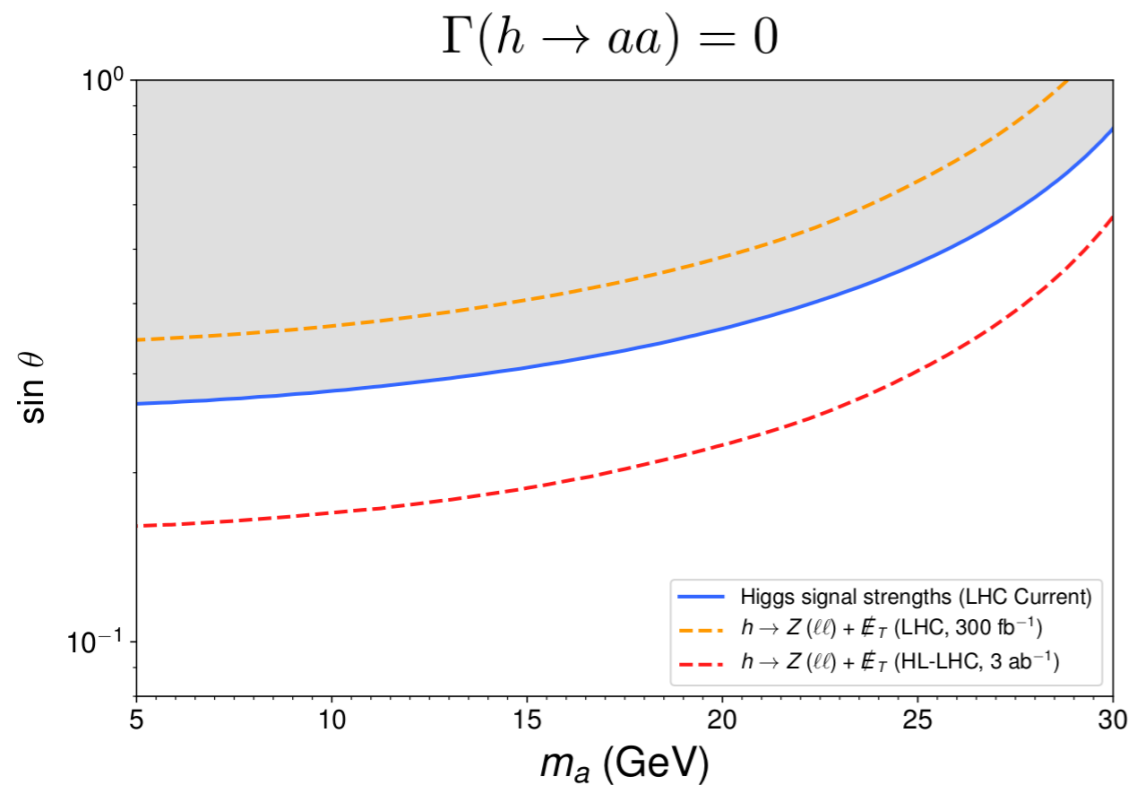
(decay present away from 2HDM alignment)

BSM Models

2HDM + a \rightarrow “Pseudoscalar portal to DM”

$$M = 600 \text{ GeV}, m_{H_0} = m_{H^\pm} = m_{A_0} = 700 \text{ GeV}$$

$$\tan \beta = 6, c_{\beta-\alpha} = 0.2, m_\chi = 0.45 m_a$$



Competitive/complementary to probes via:

- Higgs signal strength measurements *(indirect)*
- Higgs \rightarrow Invisible ($h \rightarrow aa$) *(direct)*

Summary

- Semi-dark Higgs decays: key info on Higgs portal (to dark sector) physics

- $h \rightarrow ZX$ ($X \rightarrow \cancel{E}_T$) so far unexplored at LHC

- Present in SM ($\text{BR}_{\text{SM}} \approx 0.0053$)

- HL-LHC can probe $\text{BR}_{\text{BSM}} \approx 0.01 - 0.02$ (?)

BSM Target sensitivity for Colliders
(Higgs ν floor)

- Powerful constraints on BSM scenarios:

- Extended scalar sectors (2HDM a)

- ALPs (ALP portal to dark sector)

- Dark photons...

Thank you



KEEP
CALM
AND
BACKUP
YOUR
WORK

$$e^+e^- \rightarrow Zh \quad h \rightarrow Z(\ell\ell) + \text{invisible} \quad @ILC!$$

○ Lepton colliders: Ideally suited for semi-dark Higgs decay searches

c.o.m. frame = lab frame
(knowledge of longitudinal momentum)

○ $e^+e^- \rightarrow Zh$ ILC $\sqrt{s} = 250$ GeV

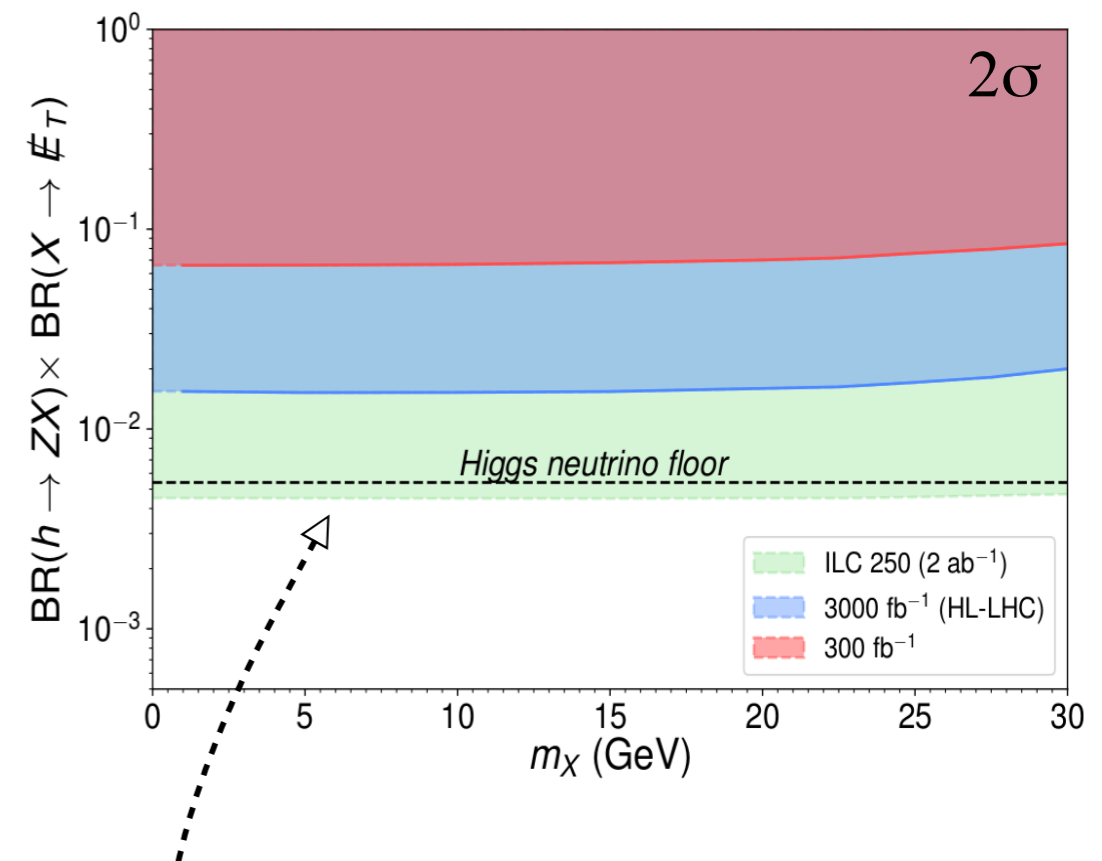
→ Higgs recoil mass for $Z_{1,2}$ ID

$$M_{\text{reco}}^2 = s + m_{Z_2}^2 - 2E_{Z_2}\sqrt{s}$$

→ (semi-visible) Higgs invariant mass for signal discrimination

$$(m_{Z_1}^{\text{miss}})^2 = \left(\sqrt{m_{Z_1}^2 + p_{Z_1}^2} + \cancel{E}\right)^2 - \left|\vec{p}_{Z_1} + \vec{\cancel{E}}\right|^2$$

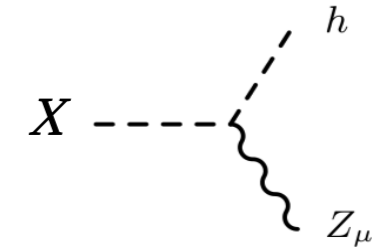
(No need for multivariate analysis)



$BR(h \rightarrow ZX) < 0.5\% \text{ reachable (95\% C.L.)}$

What about existing data? (LEP)

○ Searches for $e^+e^- \rightarrow Z^* \rightarrow h + X$ (*invisible*)



$$\exists h \rightarrow Z + X \text{ decay} \Rightarrow m_X < 35 \text{ GeV}$$

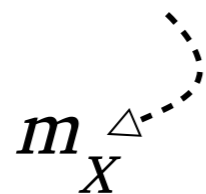
$e^+e^- \rightarrow h + X$ within LEP2 kinematic reach!

However, LEP $h + \text{invisible}$ searches target $h + Z(\nu\nu)$

e.g. OPAL

[ArXiv/hep-ex/0209078](https://arxiv.org/abs/hep-ex/0209078)

$$50 \text{ GeV} < M_{\text{miss}} < 130 \text{ GeV}$$



...insensitive to $m_X < 50 \text{ GeV}$

EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH

CERN-EP/2002-059
23rd July 2002

Search for the Standard Model Higgs Boson
with the OPAL Detector at LEP

The OPAL Collaboration

Abstract

This paper summarises the search for the Standard Model Higgs boson in e^+e^- collisions at centre-of-mass energies up to 209 GeV performed by the OPAL Collaboration at LEP. The consistency of the data with the background hypothesis and various Higgs boson mass hypotheses is examined. No indication of a signal is found in the data and a lower bound of 112.7 GeV/ c^2 is obtained on the mass of the Standard Model Higgs boson at the 95% CL.

arXiv:hep-ex/0209078v1 27 Sep 2002

No LEP constraints

BSM Models

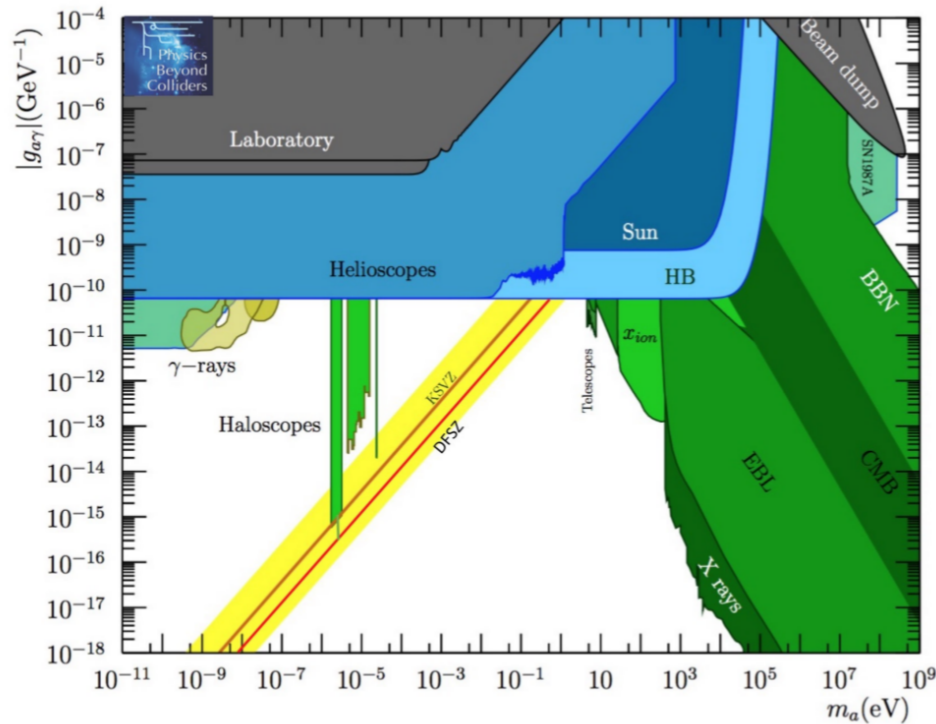
Axion-like particles (ALPs)

- **QCD Axion** = hypothetical particle from the PQ solution to **strong-CP problem**

$$\frac{1}{4}g_{agg} a G\tilde{G} + \frac{1}{4}g_{a\gamma\gamma} a F\tilde{F}$$

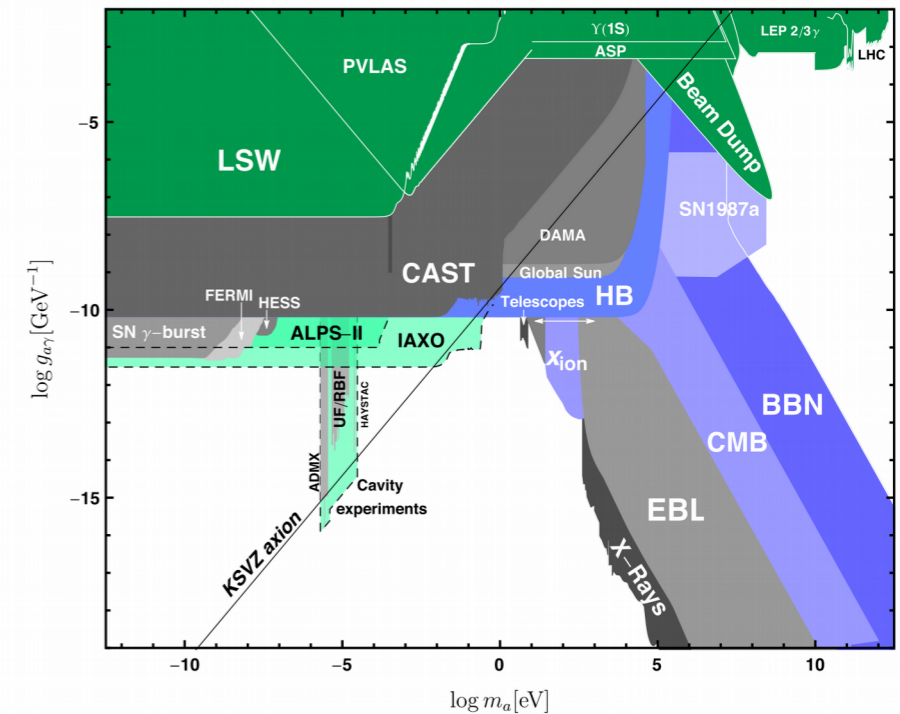
$$g_{agg} = -\frac{1}{2\pi f_a} \alpha_s$$

$$g_{a\gamma\gamma} = -\frac{1}{2\pi f_a} \alpha_{em} \left(\frac{E}{N} - \frac{2m_u + 4m_d}{3m_u + m_d} \right)$$



PBC CERN Report, arXiv:1901.09966

$$m_a^2 \simeq \frac{f_\pi^2 m_\pi^2}{f_a^2} \frac{m_u m_d}{(m_u + m_d)^2}$$



MADMAX Collaboration, Eur. Phys. J. C79 (2019), 3, 186

- **ALPs** = particle with properties similar to Axion, yet not strict QCD Axion mass-coupling relation

(may not solve strong-CP problem)

(general pseudo-Goldstone bosons)

BSM Models

Axion-like particles (ALPs)

- ALP may have:

ALP - Higgs interactions

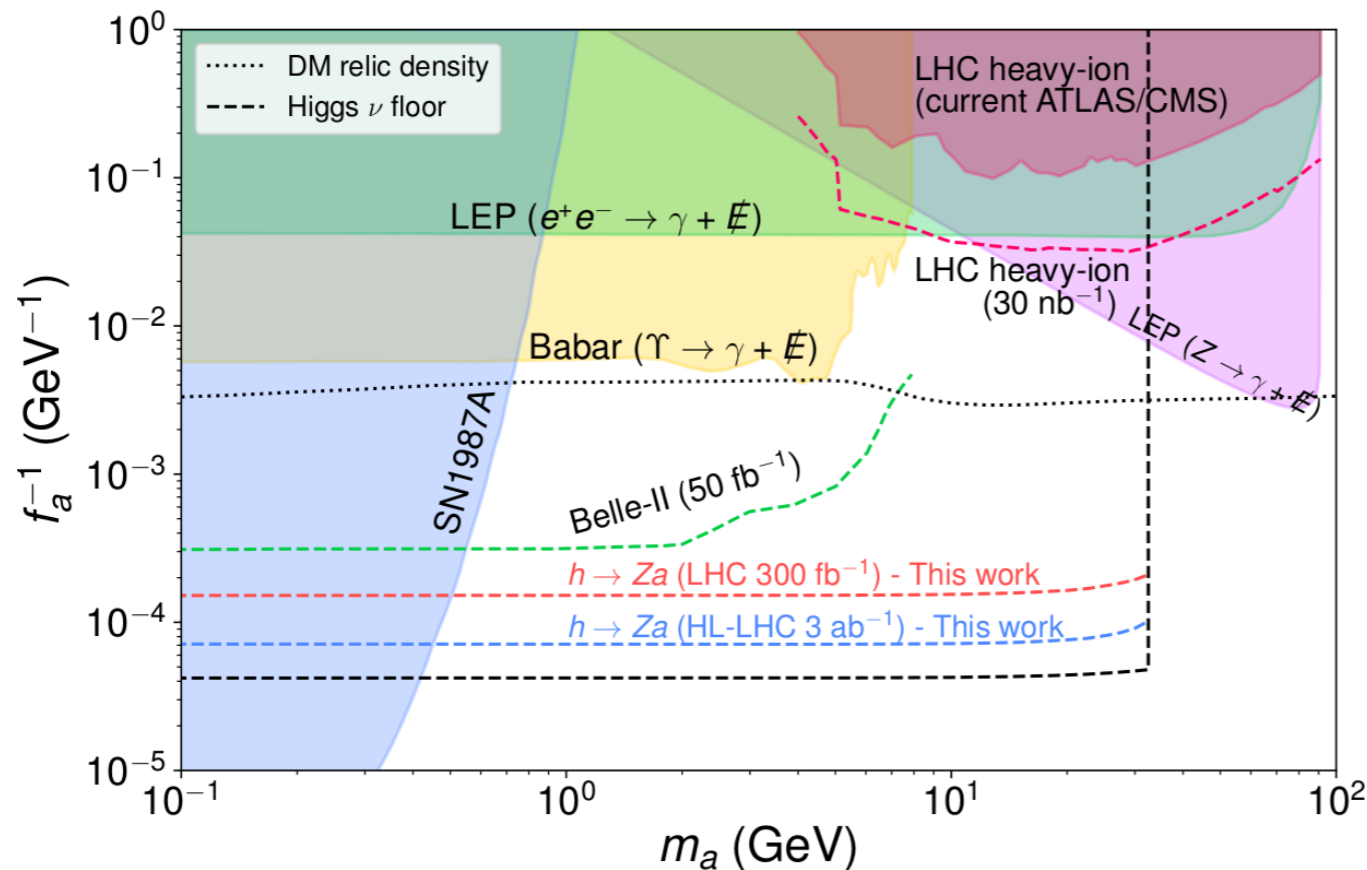
Brivio, Gavela, Merlo, Mimasu, No, del Rey, Sanz, *EJPC* 77 (2017) 8, 572
 Bauer, Neubert, Thamm, *JHEP* 12 (2017), 044



ALP - Dark sector interactions

Dolan, Ferber, Hearty, Kahlhoefer, Schmidt-Hoberg, *JHEP* 12 (2017) 094

(dark decay of ALP) “ALP portal to DM”



$$(y_\chi = 1, c_{aZh} = 1, m_\chi = 0.45 m_a)$$

BSM ALP setup

$$\Gamma(h \rightarrow Za) = (m_h^3 / 16\pi f_a^2) c_{aZh}^2 \lambda^{3/2}$$

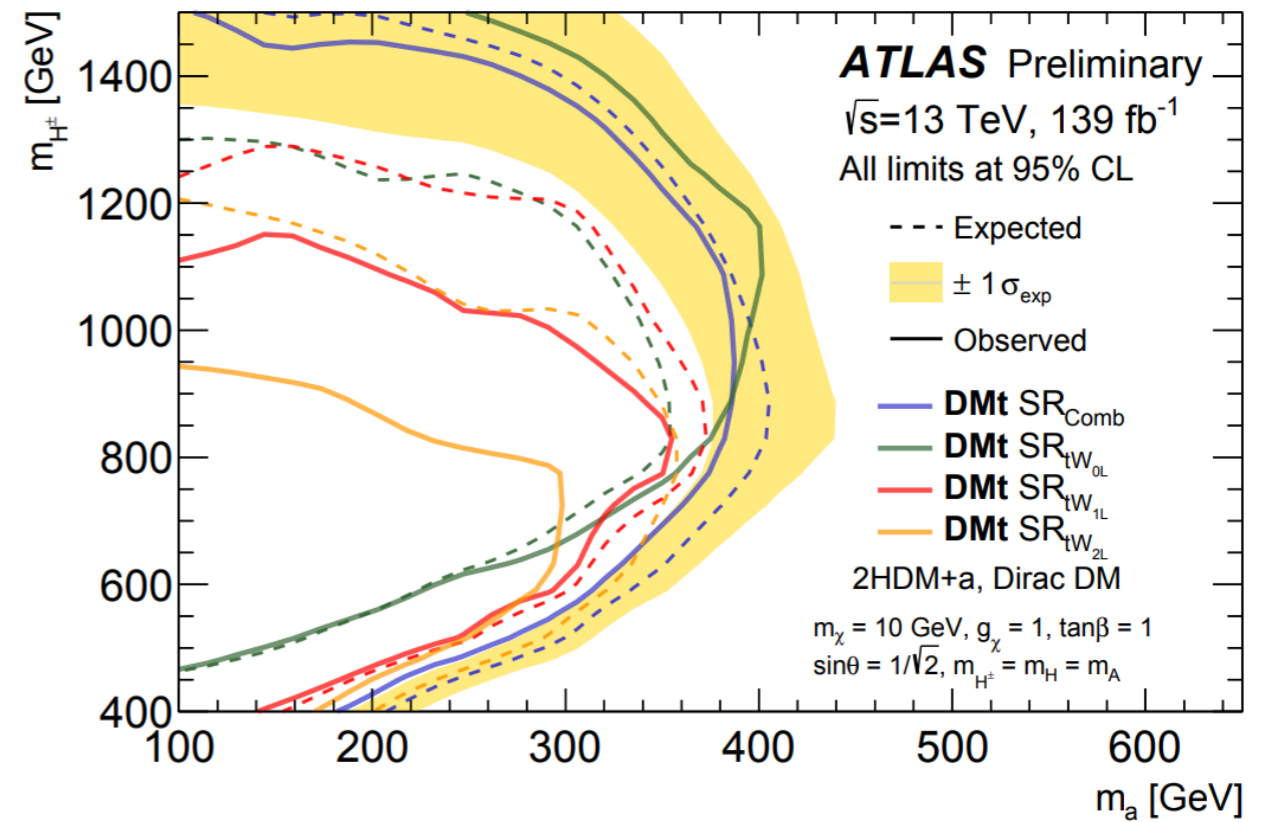
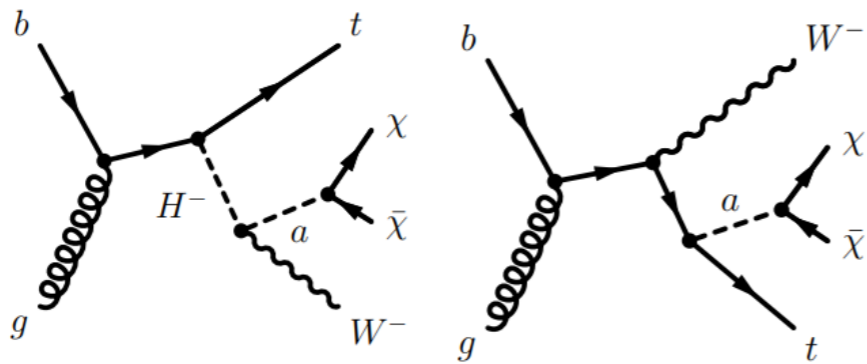
$$\text{ALP-DM: } y_\chi \bar{\chi} \gamma^\mu \gamma^5 \chi \partial_\mu a / f_a$$

$$\text{ALP-Photons: } c_{aBB} / f_a a B^{\mu\nu} \tilde{B}_{\mu\nu}$$

BSM Models

2HDM + a [LHC DM WG Benchmark Model]

e.g.



ATLAS-CONF-22-012