



## **Hadron Spectroscopy**

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# There are four interactions !

- It all started with the big bang! → Gravity governed by General Relativity (<u>it was good</u>!)
- Let there be light: and there was light! → Electromagnetic and weak interactions governed by Electroweak theory (<u>it was good</u>!)
- Let there be quarks and gluons! 

   strong interaction governed by QCD (it was good at short distance only!)

 Yes, let's study the strong interaction at long distance — nonperturbative part of QCD!







## The study of hadron spectrum



Due to the limited time, in these two days I just focus on the charmonium, bottomonium and XYZ states.





## **BEPCII and BESIII**

1.2. 1<sup>st</sup> I.R. Experi. hall 3. Power Station of ring mag. and computer center 5. 2<sup>nd</sup> I.R. Experi. hall 4. RF Station 6. Tunnel of storage ring 7. Tunnel of Trans. line 8. Tunnel of linac 9. Klystron gallery 10. Nuclear phy. Experi. hall 11. Power sta. of trans. line 12. East hall for S.R. experi. 13. West hall for S.R. experi 14. Computer center BESI  $\sqrt{s} = 2 \sim 4.95 \text{ GeV}$ Peak luminosity:  $1.02 \times 10^{33} \text{ cm}^{-1} \text{s}^{-1}$ 

 $10 \times 10^9 J/\psi$  events 2.7 ×  $10^9 \psi(3686)$  events 16 fb<sup>-1</sup>  $\psi(3770)$  events World largest  $J/\psi$ ,  $\psi(3686)$ , and  $\psi(3770)$  data samples on resonance









## **BESIII data samples**







BESIII has collected rich datasets in the XYZ region  $\sqrt{s} > 3.8$  GeV with integrated luminosity of around 22 fb<sup>-1</sup>.







# **KEKB and Belle**

Peak luminosity:  $2.11 \times 10^{34} \text{ cm}^{-1} \text{s}^{-1}$ Integrated luminosity (~980 fb<sup>-1</sup> in total):  $\Upsilon(5S)$ : 121 fb<sup>-1</sup>,  $\Upsilon(4S)$ : 711 fb<sup>-1</sup>,  $\Upsilon(3S)$ : 3 fb<sup>-1</sup>,  $\Upsilon(2S)$ : 25 fb<sup>-1</sup>,  $\Upsilon(1S)$ : 6 fb<sup>-1</sup>, continuum: 90 fb<sup>-1</sup>















### The LHC as a Beauty and Charm factory



Proton-Proton Collisions at  $\sqrt{s} = 13$ TeV ~ 20 000  $b\bar{b}$  pairs per second, x 20 of  $c\bar{c}$ pairs

High B-baryon production fraction

 $B^{+}: B^{0}: B^{0}_{s}: \Lambda^{0}_{b}$  $(u\overline{b}) (d\overline{b}) (s\overline{b}) (udb)$ 4: 4: 1: 2







## LHCb detector and performance







## Data process workflow at HEP experiment









## Charmonium States















**Burton Richter** 

Samuel C.C.Ting

- All charmonia below charm threshold
- All n=1 charmonia well known and measured
- Mass difference not large (<710 MeV), so not many channels
- Big transition rates
- Study since 1974 !











## **Charmonium Spectrum**

(1.1)

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#### Charmonium: The model

E. Eichten,\* K. Gottfried, T. Kinoshita, K. D. Lane,\* and T.-M. Yan<sup>†</sup> Laboratory of Nuclear Studies, Cornell University, Ithaca, New York 14853 (Received 9 February 1978)

A comprehensive treatment of the charmonium model of the  $\psi$  family is presented. The model's basic assumption is a flavor-symmetric instantaneous effective interaction between quark color densities. This interaction describes both quark-antiquark binding and pair creation, and thereby provides a unified approach for energies below and above the threshold for charmed-meson production. If coupling to decay channels is ignored, one obtains the "naive" model wherein the dynamics is completely described by a single charmed-quark pair. A detailed description of this "naive" model is presented for the case where the instantaneous potential is a superposition of a linear and Coulombic term. A far more realistic picture is attained by incorporating those terms in the interaction that couple charmed quarks to light quarks. The coupled-channel formalism needed for this purpose is fully described. Formulas are given for the inclusive  $e^+e^-$  cross section and for  $e^+e^-$  annihilation into specific charmed-meson pairs. The influence of closed decay channels on  $\psi$  states below charm threshold is investigated, with particular attention to leptonic and radiative widths.



color gauge interaction leads to forces that are so strong at large distances that quarks are permanently confined in color-neutral bound statesthe mesons and baryons. We also adopt this assumption.

Secondly, the large masses of the  $\psi$  resonances and charmed mesons lead to the assumption that the charmed quarks are so heavy that they may be treated nonrelativistically.<sup>4</sup> No one has yet succeeded in calculating the effective form of the interquark forces from quantum chromodynamics,<sup>16</sup> even in the nonrelativistic limit. To fill this gap we postulate that in this limit many of the gross features of the potential between the charmed quarks can be simulated by the potential

 $V(r)=-\frac{\kappa}{r}+\frac{r}{a^2}.$ 

#### **Cornell potential**

• At short distance Cornell model works pretty well  $V(r) = -4\alpha_s/3r+kr$ 



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 $n^{(2S+1)}$ 

Since quark is a spin=1/2 fermion, then for a system of quark and anti-quark the total spin should be 0 or 1.  $P=(-1)^{L+1}$ ,  $C=(-1)^{L+S}$ .

The possible spin-parity quantum numbers  $J^{PC}$  for conventional mesons ( $L \leq 3$ ).

 $1^{+-} \mid 0^{++}, 1^{++}, 2^{++}$ 

S

0

1

3

0

 $0^{-+}$ 

n radial quantum number S total spin of c & cbar L orbital angular momentum  $L = 0, 1, 2 \dots$  correspond to S, P, D,  $\dots$  $\mathbf{J} = \mathbf{S} + \mathbf{L}$  $P = (-1)^{L+1}$  parity  $C = (-1)^{L+S}$  charge conj.







## The quarkonium system

- When distance becomes larger
  - Theory 1: let there be screened potential
  - Theory 2: let there be hybrids with excited gluons
  - Theory 3: let there be tetraquark states
  - Theory 4: let there be meson molecules
  - Theory 5: let there be cusps
  - Theory 6: let there be final state interaction
  - Theory 7: let there be coupled-channel effect
  - Theory 8: let there be mixing
  - Theory 9: let there be mixture of all these effects
  - Theories ...

• The world is not that good!



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# What to measure ?

- Masses and widths of the charmonia
- Transition rates
- Multipole amplitudes (helicity amplitudes)
- Mass distributions, intermediate states
- Relations between similar/different modes
- Search for undetected modes
- C-violation, P-violation, CP-violation transitions as a probe of physics beyond SM and/or new physics



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$J/\psi(1S)$	$I^G(J^{PC})$ = $0^-(1^{})$			
$\psi/\psi(1S)$ MASS		$3096.900\pm0.006$ MeV	~	
$J/\psi(1S)$ width		$92.6\pm1.7$ keV (S = 1.1)	~	<ul> <li>They can be produced</li> </ul>
$\psi(2S)$ See the Revi	$I^G(J^{PC})$ = $0^-(1^{})$ iew on `` $\psi(2S)$ and $\chi_c$ branching ratios'' before the $\chi_{c0}(1P)$ Listings.			by e <sup>+</sup> e <sup>-</sup> annihilation directly
			<ul> <li>Expand all sections</li> </ul>	• Other lower
$\psi(2S)$ MASS		$3686.10\pm0.06$ MeV (S = 5.9)	~	charmonium states can be
$m_{\psi(2S)} - m_{J/\psi(1S)}$		$589.188\pm0.028$ MeV	~	produced by their
$\psi(2S)$ WIDTH		$294\pm8$ keV	~	transitions
$\psi(3770)$	$I^G(J^{PC})$ = $0^-(1^{})$			
$\psi(3770)$ MASS (MeV)		$3773.7\pm0.4$ MeV (S = 1.4)	~	
$m_{\psi(3770)} - m_{\psi(2S)}$		$87.6\pm0.4$ MeV (S = 1.4)	~	
$\psi(3770)$ WIDTH		$27.2\pm1.0$ MeV	~	







- E1 dominant transitions rates between ψ' and χ, χ and J/ψ were well
   measured.
- M1 transition between  $\psi'$ and  $\eta_c$ , J/ $\psi$  and  $\eta_c$  were measured with big uncertainties.
- Strong and EM transitions between  $\psi$ ' and  $J/\psi$ :
  - $\pi^+\pi^-$ ,  $\pi^0\pi^0$ : rates, mass distribution, isospin test, multipoles,  $\sigma$  pole, CPV
  - π<sup>0</sup>,η: Isospin violation strength, quark mass
  - EM:  $\pi^+\pi^-\pi^0$ ,  $\pi^0\pi^0\pi^0$ ?
  - Strong: 2(π<sup>+</sup>π<sup>-</sup>), π<sup>+</sup>π<sup>-</sup> π<sup>0</sup>π<sup>0</sup>, 4π<sup>0</sup>? PS small, how much?







# Why study transition ?

- Largest  $\psi$  decay modes (experimentally interesting)
- Understand how charm and anti-charm quarks interact (detailed information on the potential between cc-bar)
- Multipole amplitudes --- S-D mixing in  $\psi$ ' and  $\psi$ '' ( $\psi$ '' charmless decays)
- Channels with low momentum pions --- does chiral theory work?
- Shed light on  $\psi$  hadronic decays and radiative decays (eg. "12% rule")
- Chance to study  $\eta_{c_1} h_c$  and  $\eta_c$ ' more
- Search for rare and forbidden transitions





hadrons

hadrons

hadrons



(~82%) • Transitions  $MM^{\sim}$ • Hadronic transitions (~54%) • Radiative transitions (~28%) mm)mm • Leptonic decays (~ 2%)  $\mathbf{m}$ (~15%) • Hadronic decays • Strong decays (~13%)  $MM^{\epsilon}$ • EM decays (~ 2%) MNNN • Radiative decays (~1%) • Rare decays and beyond SM (<<1%)

 $\psi(2S)$  decays





## Hadronic decays: The "12% rule"

 $\boldsymbol{\Gamma}_{h} = |M_{h}|^{2} |\boldsymbol{\Psi}(0)|^{2}$ 

 $= (2/9\pi)(\pi^2 - 9)^{\frac{5}{18}}\alpha_s^{3}(\frac{4}{3}\alpha_s)^3 m_{\mathcal{O}'}.$ 

The leptonic width via one photon into  $\bar{l}l$  is

 $\Gamma_{l} = |M_{l}|^{2} |\Psi(0)|^{2} = \frac{1}{2} (\frac{2}{3}\alpha)^{2} (\frac{4}{3}\alpha_{s})^{3} m_{\mathcal{C}'},$ 

where  $\alpha \approx \frac{1}{137}$ . Although separately these calculations are not trustworthy, the ratio

 $\frac{\Gamma_{l}}{\Gamma_{h}} = \frac{\frac{2}{9}\alpha^{2}}{(2/9\pi)(\pi^{2} - 9)5/\alpha_{s}^{3}}$ 

is independent of wave-function effects.

M. Appelquist and H. D. Politzer, PRL34, 43 (1975)



(5)



This is the famous (or notorious):

"12% rule".















$\psi(3$	770)	$I^G(J^{PC})$ = $0^-(1^{})$		$\psi^{n}(c\bar{c})$ $\bar{c}$	$ \begin{array}{c} D^{0} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \\ \overline{u} \\ \end{array} \right) \\ \underbrace{u}_{\overline{D}^{0}} \left( \begin{array}{c} \end{array}{u}_{\overline{D}^{0}} \left( \begin{array}{c} $	$D^{+}) (D_{s}^{+})$ $\overline{d}) (\overline{s})$ $d) (s)$ $D^{-}) (D_{\overline{s}}^{-})$
$\psi(3770)$	) MASS (MeV)		$3773.7\pm0.4$ MeV	(S = 1.4)		~
$m_{\psi(3770)}$	$)^{-}m_{\psi(2S)}$		$87.6\pm0.4$ MeV (S	= 1.4)		~
$\psi(3770)$	) WIDTH		$27.2\pm1.0$ MeV			~
$\psi(3770)$	D) DECAY MOD	DES		•	Expand all o	decays
Mode			Fraction ( $\Gamma_i$ / $\Gamma$ )	Scale Factor/ Conf. Level	P(MeV/c)	
$\Gamma_1$	$D\overline{D}$		$(93^{+8}_{-9})\%$	S=2.0	287	~
$\Gamma_2$	$D^0\overline{D}^0$		$(52^{+4}_{-5})\%$	S=2.0	287	~
$\Gamma_3$	$D^+D^-$		$(41\pm4)\%$	S=2.0	254	~





With the singly tagged D sample, we can do some absolute measurements and search for some new decay modes of D mesons



Charmed Mesons: The existence of charmed mesons was predicted by Bjorken and Glashow [Phys. Lett. 11, 255(1964)]. A charmed meson is a bound state of c quark and one of the light antiquarks.





$$\eta_c(1S)$$
  $I^{G(J^{PC})}=0^+(0^{-+})$ 

$\eta_c(1S)$ MASS				$2983.9\pm0.4$ MeV (S =	1.2)	~
$\eta_c(1S)$ WIDTH				$32.0\pm0.7$ MeV		~
						Y

$$\eta_c(2S)$$
  ${}^{I^G(J^{PC})=0^+(0^{-+})}$ 

Quantum numbers are quark model predictions.



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### $\eta_c$ Produced from $\gamma\gamma$ -fusion









### First Observation of $\eta_c'$

### The search for $\eta'_c$ has a long and checkered history:







### To be clear about $\eta'_c$ width is important!







### Parameters of $\eta_c$

#### Puzzle: the mass and width of $\eta_c$ determined from different production mechanisms are very different









## The decays of $\eta_c'$

#### Not much $\eta'_c$ decays are observed:









## The decays of $\eta_c'$

#### Not much $\eta'_c$ decays are observed:







## $\eta_{c2}(1D)$ : spin-singlet low-lying D-wave charmonium

Mass: in the range of 3.80 to 3.88 GeV/c<sup>2</sup>, and lies between  $D\overline{D}$  and  $D^*\overline{D}$  threshold [PRD 72, 054026 (2005), PRD 79, 094004 (2009)].

Width: very narrow

**Decay mode**: branching fraction of  $\eta_{c2}(1D) \rightarrow \gamma h_c(1P)$  is large (> 50%) [PRD 80, 014001 (2009)].

The  $\eta_{c2}(1D)$  was searched in *B* decays and  $e^+e^-$  annihilations by Belle, but no signal was found [JHEP 05, 034 (2020), PRD 104, 012012 (2021)].









 $h_c(1P)$  $I^G(J^{PC})$  =  $0^-(1^{+-})$ 

Quantum numbers are quark model prediction,  $^{C}=$  established by  $\eta_{c}\gamma$  decay.

	$h_c(1P$	) MASS	$3525.37\pm0.14$ MeV	$3525.37\pm0.14$ MeV (S = 1.2)					
	$h_c(1P$	) WIDTH	$0.78\pm0.28$ MeV		~				
	$h_c(1P$	) PARTIAL WIDTHS							
	$h_c(1F)$	P) DECAY MODES		<ul> <li>Expand all d</li> </ul>					
	Mode		Fraction ( $\Gamma_i$ / $\Gamma$ )	Scale Factor/ Conf. Level	P(MeV/c)				
	$\Gamma_1$	$J/\psi(1S)\pi^0$	$< 5  imes 10^{-4}$	CL=90%	382	~			
	$\Gamma_2$	$J/\psi(1S)\pi\pi$	not seen		312	~			
	$\Gamma_3$	$J/\psi(1S)\pi^+\pi^-$	$< 2.7  imes 10^{-3}$	CL=90%	305	~			
▼ R	adiative	decays							
$\Gamma_{23}$		$\gamma\eta$	$(4.7\pm2.1) imes10$	-4		1720	~		
$\Gamma_{24}$		$\gamma \eta^{\prime}(958)$	$(1.5\pm0.4) imes10$	-3		1633	~		
$\Gamma_{25}$		$\gamma\eta_c(1S)$	$(57\pm5)\%$			500	~		

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### The importance of $h_c$ in hadron spectroscopy

- According to QCD potential model, the hyperfine splitting between spin triplet and spin singlet is determined by the vector component and coupled channel effects.
- The hyperfine splitting between spin triplet and spin singlet of *P*-wave charmonia is defined as  $\Delta M_{\rm HF}(1P) \equiv \langle M(\chi_{cJ}) \rangle M(h_c)$ . QCD predicts  $\Delta M_{\rm HF}(1P) = 0$ .
- Before the discovery of  $h_c$ , the spin-weighted average mass of  $\chi_{cJ} \langle M(\chi_{cJ}) \rangle = 3525.4 \pm 0.1 \text{ MeV}/c^2$  [PDG 2004].









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#### Further measurements of $h_c$

In 2008, CLEO Collaboration measured the mass of  $h_c$  and a product of branching fractions  $\mathcal{B}(\psi(2S) \rightarrow \pi^0 h_c) \times \mathcal{B}(h_c \rightarrow \gamma \eta_c)$ more precisely, with 24.5 × 10<sup>6</sup>  $\psi(2S)$  events (8 times of 2005) [PRL 101, 182003 (2008)].



- In 2010, BESIII Collaboration measured the  $\mathcal{B}(\psi(2S) \rightarrow \mathcal{B})$
- $\pi^0 h_c$ ) and  $\mathcal{B}(h_c \to \gamma \eta_c)$ , and determined the upper limit
- $\Gamma(h_c) < 1.44$  MeV at the 90% confidence level [PRL 104, 132002 (2010)].
- In 2022, BESIII Collaboration updated the  $h_c$  mass (3525.32 ± 0.06 ± 0.15) MeV/ $c^2$  and width (0.78<sup>+0.27</sup><sub>-0.24</sub> ± 0.12) MeV [PRD 106, 072007 (2022)]. Thus the  $\Delta M_{HF}(1P) = 0$  is consistent between experiment and theory.









## Decay modes of $h_c$ (I) – radiative decay and hadronic transition

- Except h<sub>c</sub> → γη<sub>c</sub>, a dominant decay mode of h<sub>c</sub> with the BF of 57%, more decay modes have been investigated.
- In 2016, BESIII Collaboration reported two radiative decay modes of  $h_c$ ,  $h_c \rightarrow \gamma \eta'$  (8.4 $\sigma$ ) and  $h_c \rightarrow \gamma \eta$  (4.0 $\sigma$ ) [PRL **116**, 251802 (2016)].



- Hadronic transitions of  $h_c \rightarrow \pi \pi J/\psi$  and  $h_c \rightarrow \pi^0 J/\psi$  have been predicted theoretically.
- In 2018, BESIII Collaboration searched for the process  $h_c \rightarrow \pi^+\pi^- J/\psi$ , and the upper limit was determined to be 2.7 × 10<sup>-3</sup> (90% confidence) [PRD **97**, 052008 (2018)].
- In 2022, BESIII Collaboration searched for the process  $h_c \rightarrow \pi^0 J/\psi$  via the reaction  $e^+e^- \rightarrow \pi^+\pi^-h_c$  with the data samples between 4.189 and 4.437 GeV [JHEP **2022**, 3 (2022)]. Q: Why not use  $\psi(2S) \rightarrow \pi^0 h_c \rightarrow \pi^0(\pi^0 J/\psi)$ ?







#### Decay modes of $h_c$ (II) – hadronic decay

- Theoretical predictions of the BFs of  $h_c \rightarrow light hadrons$  vary a lot in different QCD potential models.
- Since 2019, many hadronic decays of  $h_c$  have been observed by BESIII Collaboration.

















# Bottomonium States



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### **Bottomonium Spectrum**



 $n^{(2S+1)}I$ n radial quantum number S total spin of c & cbar L orbital angular momentum  $L = 0, 1, 2 \dots$  correspond to S, P, D,  $\dots$  $\mathbf{J} = \mathbf{S} + \mathbf{L}$  $P = (-1)^{L+1}$  parity  $C = (-1)^{L+S}$  charge conj.

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## **Bottomonium Spectrum**

Four main ways to access Bottomonia:

- Direct production from  $e^+e^-$ :  $J^{PC} = 1^{--}$ :  $\Upsilon(nS)$
- ISR production:  $J^{PC} = 1^{--}$ :  $\Upsilon(nS)$
- Hadronic transitions from  $\Upsilon(nS)$  through  $\eta$ ,  $\pi\pi$ , etc

$$J^{PC} = 0^{-+}, 1^{--}, 1^{+-} \dots : \Upsilon(nS), \eta_b(nS), h_b(nS), \dots$$

• Radiative transitions from  $\Upsilon(nS)$ 

$$\mathsf{J}^{\mathsf{PC}} = 0^{-+}, 0^{++}, 1^{++}, 2^{++}: \eta_{\mathrm{b}}(\mathrm{nS}), \chi_{\mathrm{b}}$$







# First observation of bottomonium: $\Upsilon(1S)$

The first bottomonium resonance, i.e.  $\Upsilon(1S)$ , was discovered in 1977 in the bombardment of a beam of high energy protons to a stationary nuclear target [PRL 39, 252 (1977), PRL 39, 1240 (1977)].









### Masses, widths, and dominant decay modes of $\Upsilon(1S)$ , $\Upsilon(2S)$ , and $\Upsilon(3S)$

Parameters	Υ(1S)	Υ(2S)	Y(3S)	
Mass (MeV/c <sup>2</sup> )	9460.40±0.10	10023.4 <u>+</u> 0.5	10355.1 <u>+</u> 0.5	These numbers
Width (keV)	54.02 <u>+</u> 1.25	31.98 <u>+</u> 2.63	20.32 <u>+</u> 1.85	are from the latest
$\mathcal{B}(\Upsilon \to \text{ggg}) \ (\%)$	81.7 <u>±</u> 0.7	58.8 <u>+</u> 1.2	35.7 <u>±</u> 2.6	T DG2023.
$\mathcal{B}(\Upsilon \rightarrow \gamma gg) \ (\%)$	2.2 <u>±</u> 0.6	1.87 <u>+</u> 0.28	1.1 <u>+</u> 0.2	

- Below the BB threshold, Y(1S, 2S, 3S) decay via this Okubo-Zweig-lizuka (OZI) suppressed way, thus leading to a narrow natural width.
- They decay into three gluons (ggg) or two gluons plus a photon (γgg) with large branching fractions, providing an entry to many potential final states, including glueballs, light Higgs bosons, and states made of light quarks.

Recent measurement of  $\mathcal{B}(\Upsilon \rightarrow \gamma gg)$  from CLEO [PRD 74, 012003 (2006)]

Photon energy spectra ( $X_{\gamma} = p_{\gamma}/E_{beam}$ ):







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#### "80% rule" for $\Upsilon(1S)$ and $\Upsilon(2S)$ decays?

From the pQCD calculations [PRL 34, 43 (1975)]:

- The 12% rule for charmonium decay
- The 80% rule for bottomonium decay

$$Q_{\psi} = \frac{\mathcal{B}_{\psi' \to \text{hadrons}}}{\mathcal{B}_{J/\psi \to \text{hadrons}}} = \frac{\mathcal{B}_{\psi' \to e^+ e^-}}{\mathcal{B}_{J/\psi \to e^+ e^-}} \approx 12\%$$
$$Q_{\Upsilon} = \frac{\mathcal{B}_{\Upsilon(2S) \to \text{hadrons}}}{\mathcal{B}_{\Upsilon(1S) \to \text{hadrons}}} = \frac{\mathcal{B}_{\Upsilon(2S) \to e^+ e^-}}{\mathcal{B}_{\Upsilon(1S) \to e^+ e^-}} = 0.80 \pm 0.08$$

The 12% rule was found to be severely violated for  $\rho\pi$  and other Vector–Pseudoscalar and Vector–Tensor final states [PRL 51, 963 (1983), PRD 69, 072001 (2004)]. This is the so-called " $\rho\pi$  puzzle".

So, how about bottomonium decays?













# Hadronic transitions among $\Upsilon(\mathrm{n}S)$





This ratio is slightly lower than the value calculated by the QCD multipole expansion method in Ref. [Front. Phys. 1, 19 (2006)].







#### **First Observation of** $\chi_{bJ}(1P)$ **:**

#### $1^{3}P_{J}$ states, observed later than $\chi_{b}(2P)$ .









#### **First Observation of** $\chi_{bJ}(2P)$ **:**







#### **First Observation of** $\chi_b(3P)$

[PRD 36, 3401 (1987); PRD 38, 279 (1988)]

It is predicted that the  $\chi_b(3P)$  has an average mass of ~10.52 GeV, with hyperfine mass splitting of 10~20 MeV.







#### Hyperfine Mass Splitting of $\chi_b(3P)$

[PRD 36, 3401 (1987); PRD 38, 279 (1988)]

It is predicted that the  $\chi_b(3P)$  has an average mass of ~10.52 GeV, with hyperfine mass splitting of 10~20 MeV.









#### **Decays of** $\chi_b$ :

The	bra	nching	fraction	$\overline{s}$ of	the kn	own de	cay mo	odes	pen q	uest
are	far	below	100%,	in	which	$\chi_b \to \gamma$	$\Upsilon(1S)$	are	The	moi
don	ninar	nt.						I	four	

states	Known channels	$\chi_b \to \gamma \Upsilon(nS)$
$\chi_{b0}(1P)$	1.94%	1.94%
$\chi_{b1}(1P)$	52.7%	35.2%
$\chi_{b2}(1P)$	<26.0%	18.0%
$\chi_{b0}(2P)$	1.38%	1.38%
$\chi_{b1}(2P)$	38.4%	29.6%
$\chi_{b2}(2P)$	<27.9%	25.5%
$\chi_{b1}(3P)$	-	Only observed mode
$\chi_{b2}(3P)$	-	Only observed mode

Open question:
$\square$ The more hadronic decay modes of $\chi_b$ are needed to be
found.
• Prob the physics of soft gluon emission and hadronization.
Study the non-perturbative QCD
□ Can exotic states be found?
• Glueball?
• XYZ via $h\chi_{b0}$ final states or in $\chi_{b0}$ decays daughters?
□ Fragmentation function: part of the great blueprint for QCD
• Where is $\chi_{b0}(3P)$ ?
Detect new physics?
• Invisible decay?
Supersymmetric quarkonia transition?





	$\Upsilon(4$	$IS)$ $I^G(J^{PC})$ = 0 $^-(1^{})$				
			also known as $\Upsilon(10580)$	1		
	22(4.0)				3 fac	tories
	T(4S)	MASS	$10579.4 \pm 1.2$ MeV		·	
	1 (45)	WIDTH	$20.3\pm2.3$ MeV			
	$\Upsilon(4S)$	DECAY MODES		•	Expand all dec	cays
	Mode		Fraction ( $\Gamma_i$ / $\Gamma$ )	Scale Factor/ Conf. Level	P(MeV/c)	
_	$\Gamma_1$	$B\overline{B}$	> 96%	CL=95%	326	~
	$\Gamma_2$	$B^+B^-$	$(51.4 \pm 0.6)\%$		331	~
	$\Gamma_3$	$D_s^+$ anything $+$ c.c.	$(17.8 \pm 2.6)\%$			~
	$\Gamma_4$	$B^0\overline{B}^0$	$(48.6 \pm 0.6)\%$		326	~
	$\Gamma_5$	$J/\psi K^0_S+$ ( $J/\psi$ , $\eta_c)$ $K^0_S$	$< 4  imes 10^{-7}$	CL=90%		~
	$\Gamma_6$	non- $B\overline{B}$	< 4%	CL=95%		<b>~</b> 55















$\Upsilon(1)$	$I020)$ $I^{G}(J^{PC}) = 0^{-}(1^{-})$	-)			
$\Upsilon(11020$	) MASS		$11000\pm4$ MeV		
$\Upsilon(11020$	) WIDTH		$24^{+8}_{-6}$ MeV		
$\Upsilon(1102)$	0) DECAY MODES				
Mode			Fraction ( $\Gamma_i$ / $\Gamma$ )	Scale Factor/ Conf. Level	P(MeV/c)
$\Gamma_1$	$e^+e^-$		$(5.4^{+1.9}_{-2.1}) imes 10^{-6}$		5500
$\Gamma_2$	$\Upsilon(1S)\pi^+\pi^-$				1408
$\Gamma_3$	$\Upsilon(2S)\pi^+\pi^-$				894
$\Gamma_4$	$\Upsilon(3S)\pi^+\pi^-$				564
$\Gamma_5$	$\chi_{bJ}(1P)\pi^+\pi^-\pi^0$		$(9^{+9}_{-8}) imes 10^{-3}$		1007
$\Gamma_6$	$\chi_{b1}(1P)\pi^+\pi^-\pi^0$		seen		975
$\Gamma_7$	$\chi_{b2}(1P)\pi^+\pi^-\pi^0$		seen		956

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# Exotic States

Note: In the light meson energy range exotic states overlap with conventional states; in the charmonium/bottomonium states the density is lower and also the overlap, but it is easier.





而篤志 切問而近思



# Hadrons: normal & multiquarks (exotic)

- Quark model: hadrons are composed from 2 (meson) quarks or 3 (baryon) quarks
- QCD does not forbid hadrons with  $N_{quarks} \neq 2, 3$ 
  - Glueball: N<sub>quarks</sub> = 0 (gg, ggg, ...)
    Hybrid: N<sub>quarks</sub> = 2 (or more) +
  - Multiquark state: N<sub>quarks</sub> > 3
  - Molecule: bound state of more than 2 hadrons







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Volume 8, number 3

PHYSICS LETTERS

1 February 1964

Multiquark states have been discussed since the 1<sup>st</sup> page of the quark model

A SCHEMATIC MODEL OF BARYONS AND MESONS \*

M. GELL-MANN California Institute of Technology, Pasadena, California

Received 4 January 1964

If we assume that the strong interactions of baryons and mesons are correctly described in terms of the broken "eightfold way" 1-3), we are tempted to look for some fundamental explanation of the situation. A highly promised approach is the purely dynamical "bootstrap" model for all the strongly interacting particles within which one may try to de rive isotopic spin and strangeness conservation and broken eightfold symmetry from soft-consistency alone 4). Of course, with only arong interactions, the orientation of the asymmetry in the unitary space cannot be specified; one hopes that in some way the selection of specific components of the Fspin by electromagnetism and the weak interactions determines the choice of isotopic spin and hypercharge directions.

Even if we consider the scattering amplitudes of strongly interacting particles on the mass shell only and treat the matrix elements of the weak, electromagnetic, and gravitational interactions by means

M. Gell-Mann, Phys. Lett. 8, 214 (1964)



ber  $n_t - n_t$  would be zero for all known baryons and mesons. The metripteresting example of such a model is one in which the triplet has spin  $\frac{1}{2}$  and z = 1 of the four particles d<sup>-</sup>, s<sup>-</sup>, u<sup>0</sup> and b<sup>0</sup> exhibit a parallel with the leptons.

A simpler and more elegant scheme can be constructed if we allow non-integral values for the charges. We can dispense entirely with the basic baryon b if we assign to the triplet t the following properties: spin  $\frac{1}{2}$ ,  $z = -\frac{1}{3}$ , and baryon number  $\frac{1}{3}$ . We then refer to the members  $u^3$ ,  $d^{-\frac{1}{3}}$ , and  $s^{-\frac{1}{3}}$  of the triplet as "quarks" 6) q and the members of the anti-triplet as anti-quarks  $\bar{q}$ . Baryons can now be constructed from quarks by using the combinations (q q q),  $(q q \bar{q} \bar{q})$ , etc., while mesons are made out of  $(q \bar{q})$ ,  $(q q \bar{q} \bar{q})$ , etc. It is assuming that the lowest baryon configuration (q q q) gives just the representations 1, 8, and 10 that have been observed, while the lowest meson configuration  $(q \bar{q})$  similarly gives just 1 and 8. Gell-Mann in his quark model paper has mentioned "exotic states" since 1964. After that, many experiments focused on finding exotic hadrons.







# A bit history on exotics hunting

- "The absence of exotics is one of the most obvious features of QCD" R. L. Jaffe, 2005
- Deuteron  $\rightarrow$  H state,  $\Omega^{-}\Omega^{-}$  bound state, ...
- No solid signature of glueballs
- Pentaquark state appeared and disappeared
  ("The story of pentaquark shows how poorly we understand QCD" F. Wilczek, 2005)
- There are lots of new states from low to high mass in various experiments! Are they normal or exotic?



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# XYZ states



# Success=X+Y+Z

#### Classification:

- $Q\overline{Q}q\overline{q}$ X: Neutral,  $J^{PC} \neq 1^{--}$ ; Y: Neutral,  $J^{PC} = 1^{--}$ ; Z: Charged
- $Q\overline{Q}qqq: P_c^+$
- Study of exotic hadrons can
  - provide new insights into internal structure and dynamics of hadrons
  - act as a unique probe to nonperturbative behavior of QCD

- Quarkonium:  $q\overline{q}$ , the simplest system of a hadron.
- Below DD/BB thresholds both charmonium and bottomonium are successful stories of QCD.
- But there are many exotic states observed in the past decade, and they are hard to fit in the two families.

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# "X Y Z" - the beginning



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## X States

Due to the time limitation, it is impossible to cover all of the XYZ states in this lecture. I select some typical results.







# $\chi_{c1}(3872)$ $I^{G}(J^{PC}) = 0^{+}(1^{++})$

#### also known as X(3872)

This state shows properties different from a conventional  $q\bar{q}$  state. A candidate for an exotic structure. See the review on non-  $q\bar{q}$  states. First observed by CHOI 2003 in  $B \to K\pi^+\pi^- J/\psi(1S)$  decays as a narrow peak in the invariant mass distribution of the  $\pi^+\pi^- J/\psi(1S)$  final state. Isovector hypothesis excluded by AUBERT 2005B and CHOI 2011 . AAIJ 2013Q perform a full five-dimensional amplitude analysis of the angular correlations between the decay products in  $B^+ \to \chi_{c1}(3872)K^+$  decays, where  $\chi_{c1}(3872) \to J/\psi\pi^+\pi^-$  and  $J/\psi \to \mu^+\mu^-$ , which unambiguously gives the  $J^{PC} = 1^{++}$  assignment under the assumption that the  $\pi^+\pi^-$  and  $J/\psi$  are in an S-wave. AAIJ 2015AO extend this analysis with more data to limit D-wave contributions to < 4% at 95% CL. See the review on ``Spectroscopy of Mesons Containing Two Heavy Quarks."



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# What is the X(3872)?

- Mass: Very close to D<sup>0</sup>D<sup>\*0</sup> threshold
- Width: Very narrow, 1.19 ±0.21 MeV [LHCb, PRD102, 092005; JHEP (2008) 123]
- J<sup>PC</sup>=1<sup>++</sup>
- Production
  - in pp/pp collision rate similar to charmonia
  - In B decays KX similar to cc, K\*X smaller than cc
  - Y(4260)→γ+X(3872)
- Decay BR: open charm ~ 50%, charmonium~O(%)
- Nature (very likely exotic)
  - Loosely D<sup>0</sup>D<sup>\*0</sup> bound state (like deuteron)?
  - Mixture of excited  $\chi_{c1}$  and  $\overline{D}^0 D^{*0}$  bound state?



 $M(\pi\pi J/\psi) - M(J/\psi)$  [GeV]





5 MeV

Candidates/













## How to understand X(3872)?



 $D^0-\overline{D}^{*0}$  molecule?

Lots of literature about this

Impossible to produce such an fragile extended object in prompt high energy hadron colliders at the rates reported by CDF & CMS  $\Gamma_{\text{"tot"}} \approx 15 \Gamma(X(3872) \rightarrow \pi^+ \pi^- J/\psi)$   $\Gamma(X(3872) \rightarrow \pi^+ \pi^- J/\psi) < 80 \text{ keV}$ 

QCD diquark-diantiquark? Maiani et al. PRD 71, 014028 (2005)

Predicts partner states (e.g., a nearby state with u→d) that have vet be seen. no charged partners of the X(3872)

no nearby neutral X(3872) partners





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# First determination of $B(B^{\pm} \rightarrow X(3872)K^{\pm})$

- The determination of the B(B<sup>±</sup>→X(3872)K<sup>±</sup>) leads to B(X(3872)→J/ψπ<sup>+</sup>π<sup>-</sup>), bringing useful information regarding the complex nature of the X(3872).
- The original tetraquark model [PRD **71**, 014028 (2005)] predicts it to be about 50%. Various molecular models [PRD **72**, 054022 (2005); PRD **69**, 054008 (2004)] predict it to be ≤10%.



BaBar, 424 fb<sup>-1</sup>, PRL 124, 152001 (2020)

- Increase signal efficiency by a factor of 3 by retaining all B tag candidates instead of the best one.
- There is  $3\sigma$  evidence of the decay  $B^{\pm} \rightarrow X(3872)K^{\pm}$ , detected for the first time using this recoil technique.

•  $B(B^{\pm} \rightarrow X(3872)K^{\pm}) = (2.1 \pm 0.6 \pm 0.3) \times 10^{-4}$ 





#### Absolute branching fractions of X(3872) decays

- Globally analyzing the measurements by BESIII, Belle, Babar, LHCb
- The absolute branching fractions of X(3872) are free parameters in the fitting

$$\chi^{2}(x) = \sum_{i=1}^{25} \frac{(x_{i} - x)^{2}}{\sigma_{i}^{2}},$$

- Statistical uncertainties are dominant for most measurements.
- Possible correlation between the systematics of different measurements in an experiments is neglected.

C.H.Li, C.Z.Yuan, Phys.Rev. D100 (2019) 094003

Index $(i)$	Parameters	Values	Experiments
	$X(3872) \to \pi^+\pi^- J/\psi$	$(\times 10^{-6})$	
1	$B^+ \rightarrow X(3872)K^+$	$8.61 \pm 0.82 \pm 0.52$	Belle [14]
2		$8.4\pm1.5\pm0.7$	BaBar $[15]$
3	$B^0 \rightarrow X(3872)K^0$	$4.3\pm1.2\pm0.4$	Belle [14]
4		$3.5\pm1.9\pm0.4$	BaBar $[15]$
	$X(3872) \rightarrow \gamma J/\psi$	$(\times 10^{-6})$	
5	$B^+ \to X(3872)K^+$	$1.78^{+0.48}_{-0.44} \pm 0.12$	Belle [22]
6		$2.8\pm0.8\pm0.1$	BaBar $[23]$
7	$B^0 \to X(3872)K^0$	$1.24^{+0.76}_{-0.61} \pm 0.11$	Belle $[22]$
8		$2.6\pm1.8\pm0.2$	BaBar $[23]$
	$X(3872) \rightarrow \gamma \psi(3686)$	$(\times 10^{-6})$	
9	$B^+ \to X(3872)K^+$	$0.83^{+1.98}_{-1.83} \pm 0.44$	Belle [22]
10		$9.5\pm2.7\pm0.6$	BaBar $[23]$
11	$B^0 \rightarrow X(3872)K^0$	$1.12^{+3.57}_{-2.90} \pm 0.57$	Belle $[22]$
12		$11.4\pm5.5\pm1.0$	BaBar $[23]$
	$X(3872) \to D^{*0}\bar{D}^0 + c.c.$	$(\times 10^{-4})$	
13	$B^+ \to X(3872)K^+$	$0.77 \pm 0.16 \pm 0.10$	Belle [16]
14		$1.67 \pm 0.36 \pm 0.47$	BaBar $[17]$
15	$B^0 \to X(3872)K^0$	$0.97 \pm 0.46 \pm 0.13$	Belle $[16]$
16		$2.22 \pm 1.05 \pm 0.42$	BaBar $[17]$
	$X(3872) \rightarrow \omega J/\psi$	$(\times 10^{-6})$	
17	$B^+ \to X(3872)K^+$	$6 \pm 2 \pm 1$	BaBar [18]
18	$B^0 \rightarrow X(3872)K^0$	$6\pm3\pm1$	BaBar $[18]$
	Ratios		
19	$\frac{\mathcal{B}(X(3872) \rightarrow \gamma J/\psi)}{\mathcal{B}(X(3872) \rightarrow \pi^+\pi^- J/\psi)}$	$0.79 \pm 0.28$	BESIII [19]
20	$\frac{\mathcal{B}(X(3872) \rightarrow D^{*0}\bar{D}^0 + c.c.)}{\mathcal{B}(X(3872) \rightarrow \pi^+\pi^- J/\psi)}$	$14.81\pm3.80$	BESIII $[19]$
21	$\frac{\mathcal{B}(X(3872) \to \omega J/\psi)}{\mathcal{B}(X(3872) \to \pi^+ \pi^- J/\psi)}$	$1.6^{+0.4}_{-0.3}\pm0.2$	BESIII $[20]$
22	$\frac{\mathcal{B}(X(3872) \to \pi^0 \chi_{c1})}{\mathcal{B}(X(3872) \to \pi^+ \pi^- J/\psi)}$	$0.88^{+0.33}_{-0.27}\pm0.10$	BESIII [21]
23	$\frac{\mathcal{B}(X(3872) \to \gamma \psi(3686))}{\mathcal{B}(X(3872) \to \gamma J/\psi)}$	$2.46 \pm 0.64 \pm 0.29$	LHCb $[24]$
	$B^+ \to X(3872)K^+$	$(\times 10^{-4})$	
24		$2.1\pm0.6\pm0.3$	BaBar $[27]$
25		$1.2\pm1.1\pm0.1$	Belle $[26]$








#### Absolute branching fractions of X(3872) decays

Parameter	index Decay mode	Branching fraction
1	$X(3872) \to \pi^+\pi^- J/\psi$	$(4.1^{+1.9}_{-1.1})\%$
2	$X(3872) \to D^{*0}\bar{D}^0 + c.c$	. $(52.4^{+25.3}_{-14.3})\%$
3	$X(3872) \rightarrow \gamma J/\psi$	$(1.1^{+0.6}_{-0.3})\%$
4	$X(3872) \rightarrow \gamma \psi(3686)$	$(2.4^{+1.3}_{-0.8})\%$
5	$X(3872) \to \pi^0 \chi_{c1}$	$(3.6^{+2.2}_{-1.6})\%$
6	$X(3872) \rightarrow \omega J/\psi$	$(4.4^{+2.3}_{-1.3})\%$
7	$B^+ \to X(3872)K^+$	$(1.9 \pm 0.6) \times 10^{-4}$
8	$B^0 \rightarrow X(3872)K^0$	$(1.1^{+0.5}_{-0.4}) \times 10^{-4}$
	$X(3872) \rightarrow \text{unknown}$	$(31.9^{+18.1}_{-31.5})\%$

- $X(3872) \rightarrow \pi^+\pi^- J/\psi \sim (4.1^{+1.9}_{-1.1})\%$
- $X(3872) \rightarrow D^0 D^{*0} \sim (52.4^{+25.3}_{-14.3})\%$
- Unknown decay ~  $(31.9^{+18.1}_{-31.5})\%$

C.H.Li, C.Z.Yuan, Phys.Rev. D100 (2019) 094003

- Statistical uncertainties are dominant.
- At Belle II, we need improve the measurements related with X(3872) decays







# Probably a mixture of DD<sup>\*</sup> & a cc<sup>\*</sup> core<sup>\*</sup>













#### *Hints before the discovery of* $X(3872) \rightarrow J/\psi \pi^+ \pi^-$

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CDF internal. 1994

x<sup>2</sup>/ndf 31.40 Observation of a narrow charmonium-like state in exclusive  $B^\pm o K^\pm \pi^+\pi^- J/\psi$ #1  $-109.4 \pm$ 0.9881 34.57± 0.2665 P2 decays P3  $95.55 \pm$ 5.546 3.685 ± 0.1631E-03 P4 Belle Collaboration • S.K. Choi (Gyeongsang Natl. U.) et al. (Sep, 2003) P5  $0.3827E - 02 \pm 0.1723E - 03$ Published in: Phys.Rev.Lett. 91 (2003) 262001 • e-Print: hep-ex/0309032 [hep-ex] reference search ि pdf DOI [→ cite links Ն\_ՆՆՆՆ 20 3.95 3.85 3.9 3.7 3.75 3.8 3.6 2016 W.K.H. Panofsky Prize in Experimental Particle Physics Recipient M 2.5Mev c10+minpi+dr.4 **BaBar internal, 2003** Stephen L Olsen  $B^+ \rightarrow J/\psi \pi^+ \pi^- K^+$ Institute for Basic Science AWG meeting June 2003 v(2S) motivation: background to 0.640 Citation: 0.173  $171.0 \pm$ J/w K<sub>1</sub>; test factorization... -23.00 ± 0.4698E-01 "For leadership in the BaBar and Belle Experiments, which established the violation of 3.684 ± 0.4084E-03 0.6907E-02 ± 0.4281E-0 CP symmetry in B-meson decay, and furthered our understanding of quark mixing and 15.88 ± 6.060 3.866 ± 0.3190E 02  $Mass = 3866 \pm 6$ quantum chromodynamics." 0.6671E-02 ± 0.1946E-02 **Background: v**(3836) Stephen Lars Olsen received a B.S. from the City College of New York in 1963 and a Ph.D. in physics from the University of Wisconsin in 1970. He is currently an Emeritus Research Fellow at the Center for Underground Physics of the Institute for Basic Science in Korea. His research has concentrated mostly on studies of heavy guarks and their associated hadrons using CLEO at Cornell, AMY and Belle experiments at KEK in Japan, and the BES experiments at IHEP in Beijing. He currently participates in the KIMS dark matter and AMoRE neutrinoless double beta decay searches at the Yangyang Underground Laboratory in Korea. Olsen was an Alfred P. Sloan Fellow (1972-1977), a John Simon Guggenheim Fellow (1986-1987), a Japan  $M(J/\psi \pi^+\pi^-)$  (GeV) Society for the Promotion of Science Fellow (1987-1988). He was awarded the University of Hawaii Regents From BaBar B-Factory Symposium (C. Hearty) What can we learn from this story? Medal for Excellence inResearch in 2002 and was designated as a University of Wisconsin Distinguished Alumni in 2007. He was elected Fellow of the APS in 1984 http://www-conf.slac.stanford.edu/b-factory-symposium/talks.asp

#### E705, PRD 50, 4258 (1994)

E705 saw  $\psi(3836)$  (2<sup>--)</sup>in 1994, 3.836  $\pm$  0.013 GeV PRL 115 011803, PRL 111 032001



**CDF saw a hint in 1994**, unpublished BaBar saw a hint in 2003, unpublished

Both CDF and Babar spotted hints of *X*(3872) *before its discovery!* 





								_
$\chi_{c1}(4274)$ ${}^{I^G(J^{PC})=0^+(1^{++})}$			X(4630)	$I^G(J^{PC})$ = $0^+(?^{?+})$				
was $X(4274)$ This state shows properties different from a conventional $q\overline{q}$ state. A candidate for an exotic structure. See the review on non- $q\overline{q}$ states. Seen by AAIJ 2017C in $B^+ \to \chi_{c1}K^+$ , $\chi_{c1} \to J/\psi\phi$ using an amplitude analysis of $B^+ \to J/\psi\phi K^+$ with a significance (accounting for systematic uncertainties) of 6.0 $\sigma$ .			This state sh "Heavy Non- $ ightarrow J/\psi\phi K^+$ with a signifi	This state shows properties different from a conventional $q\overline{q}$ state. A candidate for an exotic structure. See the review on "Heavy Non- $q\overline{q}$ Mesons." Seen by AAIJ 2021E in $B^+ \to X(4630)K^+$ with $X(4630) \to J/\psi\phi$ using an amplitude analysis of $B^+$ $\to J/\psi\phi K^+$ with a significance (accounting for systematic uncertainties) of 5.5 $\sigma$ . The $J^P = 1^-$ assignment is favored over $2^-$ with a significance of 3 $\sigma$ and other assignments are disfavored by more than 5 $\sigma$ .				
			X(4630) MASS		$4626^{+24}_{-110}$ MeV			~
$\chi_{c1}(4274)$ MASS	$4286^{+8}_{-9}$ MeV (S = 1.7)		✓ X(4630) WIDTH		$174^{+140}_{-80}$ MeV			~
$\chi_{c1}(4274)$ width	$51\pm7$ MeV		✓ X(4630) DECAY MO	DDES				
$\chi_{c1}(4274)$ decay modes	Sca	le Factor/	Mode		Fraction ( $\Gamma_i$ / $\Gamma$ )	Scale Factor/ Conf. Level	P(MeV/c)	
Mode	Fraction $(\Gamma_i / \Gamma)$	onf. Level P(MeV/c)	$\Gamma_1 \qquad J/\psi \phi$		seen		943	~
$\chi_{c1}(4685)$ $I^{G}(J^{PC}) = 0^{+}(1^{++})$	All are from an am	plitude ana	$\frac{1}{\chi_{c0}} (4700)$	► $K^+ \phi J / \psi$ )) $I^G(J^{PC}) = 0^+(0^{++})$				
This state shows properties different from a con "Heavy Non- $q \overline{q}$ Mesons." Seen by AAIJ 2021E in $B^+  o J/\psi \phi K^+$ with a significance (accounting high significance.	nventional $q\overline{q}$ state. A candidate for an exotic struc $B^+ o\chi_{c1}(4685)K^+$ with $\chi_{c1}(4685) o J/\psi\phi$ us for systematic uncertainties) of 15 $\sigma$ . The $J^P=1$	ture. See the review on ing an amplitude analysis of <sup>+</sup> assignment is favored with	This state $q \overline{q}$ states significar	shows properties different from a convent . Seen by AAIJ 2017C in $B^+  o \chi_{c0} K^+$ , $\chi_{c0}$ .ce (accounting for systematic uncertainties	was $X(4700)$ tional $q\overline{q}$ state. A candidate for an exo $ ightarrow J/\psi\phi$ using an amplitude analysis s) of 5.6 $\sigma$ .	tic structure. See the of $B^+  o J/\psi \phi K^+$	e review on nor with a	n-
$\chi_{c1}(4685)$ MASS	$4684^{+15}_{-17}{ m MeV}$		✓		4004+16			
$\chi_{c1}(4685)$ width	$126\pm40$ MeV		✓ 10 $\chi_{c0}(4700)$ MASS		$4094_{-5}^{-3}$ MeV			~
$\chi_{c1}(4685)$ decay modes			$\chi_{c0}(4700)$ WIDTH		$\delta t_{\pm 10}$ MeV			~
Mode	Fraction ( $\Gamma_i / \Gamma$ )	cale Factor/ Conf. Level P(MeV/c)	$\chi_{c0}(4700)$ deca	Y MODES		Carla Fratan	,	
$\Gamma_1$ $J/\psi\phi$	seen	1002	✓ Mode		Fraction ( $\Gamma_i$ / $\Gamma$ )	Conf. Level	P(MeV/c	c)
			$\Gamma_1 = J/\psi \phi$	11 100 1-000	seen	- 10-01 ( ) ( )	1011	76
					捕避而背	まれの	雨沂	Ð



(GeV<sup>2</sup>/c<sup>4</sup>)

2.5





#### The history/story of X(4140)/Y(4140)

CDF—PRL102:242002 (2009)



Mod.Phys.Lett. A32 (2017), 1750139

10 MeV/c<sup>2</sup>

andidates



X(4140) (renamed), mass-4.14 GeV, width—15 MeV This is the first unexpected particle discovered by Tevatron! Possible second state: mass—4.27 GeV, width—30 MeV Experienced a long road for confirmation!

- Necessarily exotic since it is narrow and above the DsDs threshold
- [*csc̄s*] tetraquark ?
- Hint of a second structure: X(4274)





#### Belle: Confirm or refute? (2009, 2010)



-B factories suffer from low pt track inefficiency -Belle cannot confirm or deny the existence of Y(4140) Belle spotted another possible new state in the same final state but from a different production: X(4350) needs to be confirmed at Belle II with larger data samples.

#### LHCb: contests CDF report (2011)



LHCb Versus CDF: Two Punches In The Face!

By Tommaso Dorigo | July 27th 2011 05:48 AM | 10 comments | 🖴 Print | 🖂 E-mail | Track

result. Note that, as reported in the figure, if the CDF signal were as estimated by CDF, LHCb would

have been able to fit 39+-9+-6 events. The Y(4140) is on very shaky ground at the moment, and the

new PDG will likely change its status in the particle zoo... This is punch number 1.

 $\chi\chi \rightarrow \phi J/\psi$ 

4.8

X(4350)

4.6

 $M(\phi J/\psi)$  (GeV/c<sup>2</sup>)

4.4





# 1958

#### **Result from CMS (2012)**





significance greater than 5σ, confirms the existence of Y(4140) for the first time from another experiment

 evidence for a second structure in the same mass spectrum D0 provides the second independent confirmation of Y(4140) with 3.1 $\sigma$  significance







 $M(J/\psi \phi)$  [GeV] significance: >> 5 $\sigma$ Mass and width are consistent with their previous measurements and CDF/CMS D0 provides additional confirmation from a different production





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No significance for both structures though there are hints

BaBar provides useful information even though there is no significant signals Three events @4.15 GeV

BES sets limits, cannot compare because it is from a different process

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#### A short summary before LHCb enters

[PRD 95 (2017) 012002]

Year	Experiment	$B \to J\!/\psi  \phi K$	X(4140) peak			
	luminosity	yield	Mass [MeV]	Width $[MeV]$	Sign.	Fraction $\%$
2008	CDF $2.7 \text{ fb}^{-1}$ [1]	$58\pm10$	$4143.0 {\pm} 2.9 {\pm} 1.2$	$11.7^{+8.3}_{-5.0}{\pm}3.7$	$3.8\sigma$	
2009	$Belle \ [22]$	$325\pm21$	4143.0 fixed	11.7 fixed	$1.9\sigma$	
2011	$CDF \ 6.0 \ fb^{-1} \ [29]$	$115\pm12$	$4143.4 {}^{+2.9}_{-3.0} {\pm} 0.6$	$15.3^{+10.4}_{-6.1}{\pm}2.5$	$5.0\sigma$	$14.9 \pm 3.9 \pm 2.4$
2011	LHCb $0.37 \text{ fb}^{-1}$ [21]	$346\pm20$	4143.4 fixed	15.3 fixed	$1.4\sigma$	< 7 @ 90%CL
2013	CMS 5.2 fb $^{-1}$ [25]	$2480 \pm 160$	$4148.0 {\pm} 2.4 {\pm} 6.3$	$28 \ {}^{+15}_{-11} \ \pm 19$	$5.0\sigma$	$10{\pm}3$ (stat.)
2013	D0 10.4 fb $^{-1}$ [26]	$215\pm37$	$4159.0 {\pm} 4.3 {\pm} 6.6$	$19.9 {\pm} 12.6 {}^{+1.0}_{-8.0}$	$3.0\sigma$	$21{\pm}8{\pm}4$
2014	$BaBar \ [24]$	$189\pm14$	4143.4 fixed	15.3 fixed	$1.6\sigma$	< 13.3 @ 90%CL
2015	D0 10.4 fb $^{-1}$ [27]	$p \bar{p}  ightarrow J / \psi \phi$	$4152.5{\pm}1.7_{-5.4}^{+6.2}$	$16.3 {\pm} 5.6 {\pm} 11.4$	$4.7\sigma$ (5.7	$(\sigma)$
Average			$4147.1 {\pm} 2.4$	$15.7{\pm}6.3$		
						81







# **Results from LHCb (2016)**

LHCb, PRL 118 (2017), 022003; PRD 95 (2017), 012002



- No light quark (u,d) components
- Cannot exchange pion— $J/\phi$  or  $\phi$  has no isospin
- Cannot exchange photon--pion— $J/\phi$  or  $\phi$  has no charge
- A case of more general tetra-quark dynamics
- New important piece to the exotic meson family

▶ LHCb re-confirmed both X(4140) and X(4274), Observed X(4500) and X(4700)

- LHCb found two additional resonances in the same mass spectrum
- ▶ This is 7 years after the first report from CDF
- Waing for Belle II larger data samples: signals should be more cleaner







# **Updated Results from LHCb (2021)**



lew states: 2	Z <sub>cs</sub> (4000), X( X(4150) < 50	4685) > 15	σ; <i>Z<sub>cs</sub></i> (422	0), <i>X</i> (4630)
Contribution	Significance $[\times \sigma]$	$M_0[{ m MeV}]$	$\Gamma_0[{\rm MeV}]$	$\mathrm{FF}\left[\% ight]$
$X(2^{-})$				
X(4150)	4.8(8.7)	$4146 \pm 18 \pm 33$	$135 \pm 28  {}^{+ 59}_{- 30}$	$2.0\pm0.5^{+0.8}_{-1.0}$
$X(1^{-})$				
X(4630)	5.5(5.7)	$4626 \pm 16^{+18}_{-110}$	$174 \pm 27  {}^{+ 134}_{- 73}$	$2.6 \pm 0.5 ^{+2.9}_{-1.5}$
All $X(0^+)$	Stat.(Syst. inc	luded)		$20 \pm 5 {}^{+14}_{-7}$
X(4500)	20 (20)	$4474 \pm 3 \pm 3$	$77\pm6{}^{+10}_{-8}$	$5.6 \pm 0.7 {}^{+2.4}_{-0.6}$
X(4700)	17 (18)	$4694 \pm 4  {}^{+ 16}_{- 3}$	$87\pm8{}^{+16}_{-6}$	$8.9 \pm 1.2  {}^{+4.9}_{-1.4}$
$\mathrm{NR}_{J/\psi\phi}$	4.8(5.7)			$28 \pm 8 {}^{+19}_{-11}$
All $X(1^+)$				$26 \pm 3 {+ 8 \atop -10}$
X(4140)	13 (16)	$4118 \pm 11 {}^{+19}_{-36}$	$162 \pm 21  {}^{+ 24}_{- 49}$	$17 \pm 3  {}^{+ 19}_{- 6}$
X(4274)	18 (18)	$4294 \pm 4^{+3}_{-6}$	$53 \pm 5 \pm 5$	$2.8 \pm 0.5  {}^{+ 0.8}_{- 0.4}$
X(4685)	15 (15)	$4684 \pm 7^{+13}_{-16}$	$126 \pm 15 {}^{+37}_{-41}$	$7.2 \pm 1.0 {}^{+4.0}_{-2.0}$
All $Z_{cs}(1^+)$				$25 \pm 5 {}^{+11}_{-12}$
$Z_{cs}(4000)$	15 (16)	$4003 \pm 6 {}^{+}_{-}{}^{4}_{14}$	$131 \pm 15 \pm 26$	$9.4 \pm 2.1 \pm 3.4$
$Z_{cs}(4220)$	5.9(8.4)	$4216 \pm 24  {}^{+43}_{-30}$	$233 \pm 52  {}^{+ 97}_{- 73}$	$10 \pm 4^{+10}_{-7}$







## **Updated Results from LHCb (2021)**

- For X(4140), no evidence of a narrow threshold resonance at  $J/\psi\phi$  in our data
- 4 new  $J/\psi K^+$  and  $J/\psi \phi$  structures observed in  $B^+ \rightarrow J/\psi \phi K^+$  decays with 6 times data and much clean environment
  - A 1<sup>+</sup>  $Z_{cs}(4000)^+ \rightarrow J/\psi K^+$  observed for 1<sup>st</sup> time, significance > 15 $\sigma$
  - A broad  $Z_{cs}(4220)^+ > 5\sigma$
  - A new 1<sup>+</sup> X(4685) is > 15 $\sigma$ , and new X(4630) > 5 $\sigma$
  - 4 X states previously observed are confirmed, and  $J^P$  determined with higher significances

CMS should update their results on this channel with a (much) larger data sample, and more sophisticated analysis technique, than previously.





Are *X*(4700) and *X*(4740) the

same state? Further amplitude



## The story of $J/\psi \phi$ system is not finished yet !

#### Study of $B_s^0 \rightarrow J/\psi \pi^+ \pi^- K^+ K^-$ decays [JHEP 02 (2021) 024]

χ<sub>c1</sub>(3872) and J/ψφ structures can be studied in this decay
 Production rate measurements can shed light on the nature of exotic states









#### **Other productions for charmonium-like states**









Y States









### $e^+e^- \rightarrow \pi^+\pi^- J/\psi$ cross section : Y(4260)



X(4260)	$I^{G}(J^{PC}) = ?^{?}(1^{})$			
 X(4260) MASS	4251 ± 9	AVERAGE		
X(4260) WIDTH	120 ±12	AVERAGE		

#### BABAR PRL95,142001(2005)



also known as Y(4230); was  $\psi(4260)$ 

The original  $\psi(4260)$  (also known as Y(4260)) was observed by AUBERT,B 2005I as a peak in the energy dependence of the  $e^+ e^- \rightarrow \pi^+ \pi^- J/\psi$  cross section and was confirmed by HE 2006B, YUAN 2007, LEES 2012AC, and LIU 2013B in the same process. A higher-statistics analysis by ABLIKIM 2017B revealed an asymmetry in the cross section and resulted in a shift of the peak position to a lower mass. The  $\psi(4260)$  was therefore renamed  $\psi(4230)$ . The energy-dependent cross sections for  $e^+e^-$  to other channels also exhibit peaks in the same mass region. The parameters corresponding to those peaks are also listed here, but the number of states in this region remains to be determined. For details see the review on "Spectroscopy of mesons containing two heavy quarks."







#### $Y(4260) \rightarrow \pi^+\pi^- J/\psi$ confirmed by Belle







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5

0

4.5

 $M(\pi^+\pi^-\psi(2S))$  (GeV/c<sup>2</sup>)





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5.5

Y(4008) Y(4260)

Y(4360) Y(4660)



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### $Y(4260) \rightarrow Y(4230) + Y(4320)$









# Y(4260): mass $\rightarrow$ lower & width $\rightarrow$ narrower





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 $E_{CM}^{4.3}(GeV)$ 

4.4

4.5

<sup>4</sup><u>6</u> 94





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- How big is the influence of the  $D_1\overline{D}$  threshold?
  - $_{\circ}$  Why Y(4320) only seen in  $e^+e^- \rightarrow J/\psi\pi\pi$ ?







# What is the Y(4260)?

The Y(4260) mass is lower and width narrower than previously thought

"Y(4260)" → Y(4220)?

If it is a  $D\overline{D}_1(2420)$  molecule:

B.E.  $\approx$  66 MeV  $\leftarrow$  too large?? "affinity" to  $D\overline{D}_{1}(2420)$  should be high

```
If it is a c\overline{c}-gluon hybrid:
```

its mass is ~65 MeV below current ( $m_{\pi} \approx 400$  MeV) LQCD predictions  $\leftarrow$  not so bad? "affinity" to  $D\overline{D}_{0}(2400)$  should be high

2012 LQCD calc. (m<sub>π</sub>≈400 MeV):

pre-2017: too high by ~35 MeV

post-2017: too high by ~65 MeV

Had. Spectr. Collab. JHEP07, 126

Maiani et al. PRD89,114010 If it is a QCD diquark–diantiquark tetraquark: it should have Isospin- &  $SU_{F}(3)$ -multiplet partner states  $\leftarrow$  not seen

If it is hadrocharmonium:

decays to non-J/ $\psi(h_c)$  charmonium states should be suppressed  $\leftarrow$  they aren't

BESIII is well suited to further investigate this intriguing puzzle  $\leftarrow a''Y(4260)''$  factory



Dubynskiy & Voloshin, PLB 666, 344 Li & Voloshin, Mod. Phys. Lett. A29, 1450060











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√s, GeV



Exclusive cross sections contribution B to the total cross section R

Blue: R-measurement Red: Cross section measurements



r

D\*D\*

√s, GeV















After we have measured all the  $e^+e^-$  annihilation cross sections, what do we do to get the resonant parameters of the vector charmonium(-like) states?

















#### Interpretation of the Y(10753)

- D-wave bottomonium
  - B. Chen, A.L. Zhang, J. He, arXiv:1910.06065, Bottomonium spectrum in the relativistic flux tube model (3D)
  - Q. Li, M.S. Liu, Q.F. Lü, L.C. Gui, X.H. Zhong, arXiv:1905.10344, Canonical interpretation of Y(10750) and Y(10860) in the Y family (4D)
- $\overline{B}^{(*)}B^{(*)}$  dynamically generated pole
  - P. Bicudo, M. Cardoso, N. Cardoso, M. Wagner, arXiv:1910.04827, Bottomonium resonances with I=0 from lattice QCD correlation functions with static and light quarks
- Hybrid
  - J. T. Castellà, arXiv:1908.05179, Spin Structure of heavy-quark hybrids
- Tetraquark state
  - A. Ali, L. Maiani, A. Y. Parkhomenko, W. Wang, arXiv:1910.07671, Interpretation of Yb (10753) as a tetraquark and its production mechanism
  - Z.G. Wang, arXiv:1905.06610, Vector hidden-bottom tetraquark candidate: Y(10750)







# Study of properties of $\Upsilon(10753)$

- Largest bottomonium data sample at Belle and Belle II
- In Nov. 2021, Belle II collected ~20/fb of unique scan data at energies near 10.75 GeV
  - Fill the gaps in Belle Scan data
  - Physics goal is to understand the nature of  $\Upsilon(10753)$



#### More analyses are ongoing

0

0

0

- $\circ \qquad \Upsilon(10753) \to K^+ K^- \Upsilon(nS)$ 
  - $\Upsilon(10753) \rightarrow \eta(\eta')\Upsilon(nS)$

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 $\Upsilon(10753) \rightarrow \gamma X_b, X_b$  $\rightarrow \pi \pi \chi_{bJ}, \pi \pi \Upsilon(nS)$ 

etc...







# Observation of $\Upsilon(10753) \rightarrow \omega \chi_{bJ}$

PRL 130, 091902 (2023)



$\Gamma_{ee}\mathcal{B}_{f}$	Solution I (constructive interference)	Solution II (destructive interference)	
$\Gamma_{ee}\mathcal{B}(\Upsilon(10753) \rightarrow \omega \chi_{b1})$	(0.63±0.39±0.20) eV	(2.01±0.38±0.76) eV	
$\Gamma_{ee}\mathcal{B}(\Upsilon(10753) \rightarrow \omega \chi_{b2})$	(0.53±0.46±0.15) eV	(1.32±0.44±0.55) eV	

The e<sup>+</sup>e<sup>-</sup>  $\rightarrow \omega \chi_{bJ}$  (J = 1, 2) cross sections peak at Y(10753). Fit cross section with function:  $\sigma_{e^+e^- \rightarrow \omega \chi_{b1}}(\sqrt{s}) = |\sqrt{PS_2(\sqrt{s})} + BW(\sqrt{s})e^{i\varphi}|^2, BW(\sqrt{s})$  $= \frac{\sqrt{12\pi\Gamma_{ee}B_f\Gamma}}{s - M^2 + iM\Gamma} \sqrt{\frac{PS_2(\sqrt{s})}{PS_2(M)}}$ 

M and  $\Gamma$  of  $\Upsilon(10753)$  are fixed according to Ref. [JHEP 10, 220(2019)].

1.  $\sigma(e^+e^- \rightarrow \omega \chi_{b1})/\sigma(e^+e^- \rightarrow \omega \chi_{b2})=1.3\pm0.6$  at 10.745 GeV, contradicts the expectation for a pure D-wave bottomonium state of 15 [PLB 738, 172 (2014)] 2. There is also a 1.8 $\sigma$  difference with the prediction for a S-D-mixed state of 0.2 [PRD 104, 034036 (2021)]

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Z States



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 $Z_c(4430)^{\pm}$  exist or not ?





"For the fit ... equivalent to the Belle analysis...we obtain mass & width values that are consistent with theirs,... but only  $\sim$ 1.9s from zero; fixing mass and width increases this to only  $\sim$ 3.1s."

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## LHCb 4-dim analysis of $B \rightarrow K^+\pi^-\psi'$





PRD 78, 072004 (2008)









- Dalitz-plot analysis of  $\underline{B}^0 \rightarrow \chi_{c1} \pi^+ K^- \chi_{c1} \rightarrow J/\psi \gamma$  with 657M BB
- Dalitz plot models: known  $K^* \rightarrow K\pi$  only

K\*'s + one  $Z \rightarrow \chi_{c1} \pi^{\pm}$ 

K\*'s + two Z<sup> $\pm$ </sup> states  $\Rightarrow$  favored by data











### **BaBar doesn't see significant** $Z^{\pm} \rightarrow \chi_{c1} \pi^{\pm}$

PRD85, 052003 (2012)

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$$\begin{split} \mathcal{B}(\bar{B}^{0} \to Z_{1}(4050)^{+}K^{-}) &\times \mathcal{B}(Z_{1}(4050)^{+} \\ &\to \chi_{c1}\pi^{+}) < 1.8 \times 10^{-5}, \\ \hline & \text{Belle: } (3.0^{+1.5} - 0.8^{+3.7} - 1.6) \times 10^{-5} \\ \mathcal{B}(\bar{B}^{0} \to Z_{2}(4250)^{+}K^{-}) &\times \mathcal{B}(Z_{2}(4250)^{+} \\ &\to \chi_{c1}\pi^{+}) < 4.0 \times 10^{-5}, \end{split}$$

Belle: (4.0<sup>+2.3</sup>-0.9<sup>+19.7</sup>-0.5)x10<sup>-5</sup>

"We find that it is possible to obtain a good description of our data without the need for additional resonances in the  $\chi_{c1}\pi$  system."









- 1. Almost full Belle data sample used: Lum=967 fb<sup>-1</sup> data.
- 2. Using ISR photon non-tagged method, Y(4260) was observed significantly.
- 3.  $4.15 \le M(p^+p^-J/y) \le 4.45$  GeV to select Y(4260) resonance.
- 4. Dalitz plot also shows structures.

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## Zc(3900)<sup>±</sup> from Belle



- 1. S-Wave BW, p\*q phase space factor, efficiency applied, to fit  $M_{max}(\pi^{\pm}J/\psi)$  distribution
- 2. Belle observed 689 events, with 139 background.
- 3. M=(3894.5 $\pm$ 6.6 $\pm$ 4.5) MeV;  $\Gamma$ =(63 $\pm$ 24 $\pm$ 26) MeV.
- 4. Significance:5.2σ.
  - Comment: Since Zc(3900) is charged and can decay into  $\pi J/\psi$ , it must have at least four quarks.

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## BESIII + Belle + CLEO's data





1. CLEO's data at 4.17 GeV by K. Seth. 2. M=3885±5 MeV,  $\Gamma$ =34±13 MeV. 3. Significance:  $6\sigma$   $M(Z_c(4430))-M(Z_c(3900)) = 589 \pm 30 \text{ MeV}$  $M(\psi') - M(J/\psi) = 589 \text{ MeV}$ 

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## Z<sub>c</sub>(3900) State (I=1)





Evidence with  $3.7\sigma$  by using CLEO-c data

 $M = 3894.8 \pm 2.3 \pm 3.2 \text{ MeV/c}^2$ ,  $\Gamma = 29.6 \pm 8.2 \pm 8.2 \text{ MeV}$ 

>An iso-spin triplet is established!







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#### Observation of the $Z_c(3900)$ — a charged charmoniumlike structure —





#### QUARK SOUP

Researchers at colliders in China and Japan have succeeded in making exotic matter comprising four quarks, but are still debating whether the fleeting particles are meson pairs or true tetraquarks.



#### PARTICLE PHYSICS

### Quark quartet opens fresh vista on matter

First particle containing four quarks is confirmed.

#### **BY DEVIN POWELL**

 $P_{\text{that may have existed in the first hot} \\ \text{moments after the Big Bang. Arcanely} \\ \text{called } Z_{c}(3900), \text{it is the first confirmed particle made of four quarks, the building blocks} \\ \text{of much of the Universe's matter.}$ 

Until now, observed particles made of quarks have contained only three quarks (such as protons and neutrons) or two quarks (such as the pions and kaons found in cosmic rays). Although no law of physics precludes larger

antimatter counterparts, positrons. These crashes have one-thousandth the energy of those at the world's most powerful accelerator, the Large Hadron Collider (LHC) at CERN near Geneva, Switzerland, but they are still energetic enough to mimic conditions in the early Universe. Collision rates at KEK are

"They have clear evidence of a particle with four quarks." more than twice those at the LHC, and they occasionally give birth to rare particles not found in nature today —

Question: Zc has been confirmed. How about Zs ? How to search for it ?















## **Spin-parity hypothesis with likelihood test**

Amplitude model

helicity amplitude, covariant tensor amplitude

- null and alternative hypothesis
  - Null: spin=J, alternative: spin  $\neq J$
- test with data events
  - Minimize log-likelihood function
- check with angular distributions, moment analysis and invariant mass lineshape
- significance test:
  - likelihood ratio or ToyMC ensemble (avoid look elsewhere effect)

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$$Z_c(3900)$$
  $I^G(J^{PC}) = 1^+(1^{+-})$ 

#### was X(3900)

Properties incompatible with a  $q\overline{q}$  structure (exotic state). See the review on non-  $q\overline{q}$  states. Charged  $Z_c(3900)$  seen as a peak in the invariant mass distribution of the  $J/\psi\pi^{\pm}$  system by BES III (ABLIKIM 2013T) in  $e^+ e^- \rightarrow \pi^+\pi^- J/\psi$  at c.m. energy of 4.26 GeV and by radiative return from  $e^+e^-$  collisions at  $\sqrt{s}$  from 9.46 to 10.86 GeV at Belle (LIU 2013B). Partial wave analysis of ABLIKIM 2017J determines  $J^P = 1^+$  with more than 7  $\sigma$  significance. Neutral  $Z_c(3900)$  seen in the  $J/\psi\pi^0$  invariant mass distribution in  $e^+ e^- \rightarrow \pi^0\pi^0 J/\psi$  at c.m. energies of 4.23, 4.26, and 4.36 GeV by BES III (ABLIKIM 2015U) and at 4.17 GeV by XIAO 2013A. Peaks in ( $D\overline{D}^*$ )<sup>0,±</sup> reported by BES III (ABLIKIM 2014A, ABLIKIM 2015AB) are assumed to be related.

$Z_c(3900)$ mass		$3887.1\pm2.6$ MeV (S =	1.7)	~
$Z_c(3900)$ WIDTH		$28.4\pm2.6$ MeV		~
	1. 2.	M. Ablikim, et al. (BESIII), Phys.Rev.Lett.,110, 252001 (2013). M. Ablikim, et al. (BESIII), Phys.Rev.Lett.,119, 072001 (2017).		119
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Dalitz plot and mass spectrum







# 系

## Spin and parity measurement of Zc(3900)







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### Angular distribution for the Zc(3900) J<sup>P</sup> assumption

1. $e^+e^- \rightarrow \pi^{\pm}Z_c^{\mp}$	F	
	$\sin \theta_0 \qquad (J^p = 0^-)$	
dN	$1 + \alpha_0 \cos^2 \theta_0 \ (J^P = 1^+)$	$\pi^{\mp}$ $( heta_1, \phi_1)$
	$-\propto \left\{ 1 + \cos^2 \theta_0 \qquad (J^P = 1^-) \right\}$	$\times$
$a\cos\theta_0$	$1 + \alpha_0 \cos^2 \theta_0 \ (J^P = 2^-)$	$Z^+ \sum J/\psi$
	$\left(1 + \cos^2 \theta_0 \qquad (J^P = 2^+)\right)$	$Z_{c}$ $(\theta_{0}, \phi_{0})$
2. $Z_c^{\mp} \rightarrow \pi^{\mp} I/\psi$		$e^+$ $e^-$
	$(1  (I^p = 0^-))$	
	$1 + \alpha_1 \cos^2 \theta_1 \ (I^P = 1^+)$	$\pi^+$
$\frac{dN}{dN} \propto $	$1 + \cos^2 \theta_1$ $(I^P = 1^-)$	
$d\cos heta_1$	$1 + \alpha_1 \cos^2 \theta_1 + \alpha_2 \cos^4 \theta_1$ $(I^P = 2^-)$	
	$1 - 3\cos^2\theta_1 + 4\cos^4\theta_1$ $(I^P = 2^+)$	Chin. Phys. Lett. 33, 061401 (2016)
		122





# 1958

### Angular distribution for the Zc(3900) J<sup>P</sup> assumption



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BESI

## Observation of $Z_{cs}(3985)$ —first $Z_c$ with a strange quark

•  $e^+e^- \to K^+(D_s^-D^{*0} + D_s^{*-}D^0)$  PRL 126, 102001 (2021)

- 3.7fb<sup>-1</sup> data at 4628, 4640, 4660, 4680, and 4700
- Partial reconstruction of the process, tag K and  $D_s^-$
- $D_s^-$  reconstructed with  $K^+K^-\pi^-$  [ $\phi\pi$  or  $K^*K$ ] and  $K_s^0K^-$  e<sup>+</sup>



• Both decay modes can survive the selection

3 D.

- Combinatorial background described by wrong sign (WS) events
- Absolute contribution in signal region determined from a fit to  $RM(K^+D_s^-)$





**BES**III

## Observation of $Z_{cs}(3985)$ —first $Z_c$ with a strange quark



PRL 126, 102001 (2021)

- Assume  $J^P=1^+$
- Simultaneous fit to five data samples
- Signal component:

$$\left|\frac{\sqrt{q}}{M^2 - m_0^2 + im_0(f\Gamma_1(M) + (1 - f)\Gamma_2(M))}\right|$$

a.n.

f = 0.5 represents the fraction of the two decay modes

• Pole position:

 $m = 3982.5^{+1.8}_{-2.6} \pm 2.1 \text{MeV}/c^2$   $\Gamma = 12.8^{+5.3}_{-4.4} \pm 3.0 \text{MeV}$ 

- Significance:  $5.3\sigma$
- At least four quarks  $c\overline{csu}$

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# **E**

## X, Y, Z particles are correlated









**Resonant structure of \Upsilon(5S) \rightarrow (b\bar{b})\pi^+\pi^-**









## $Z_b^{\pm} \rightarrow$ Open Beauty



Assuming that  $Z_b$ decays are saturated by the  $\Upsilon(nS)\pi$ ,  $h_b(mP)\pi$ and  $B^{(*)}B^*$  channels, branching fractions are here:

Channel	Fraction, %			
	$Z_b(10610)$	$Z_b(10650)$		
$\Upsilon(1S)\pi^+$	$0.60 \pm 0.17 \pm 0.07$	$0.17 \pm 0.06 \pm 0.02$		
$\Upsilon(2S)\pi^+$	$4.05 \pm 0.81 \pm 0.58$	$1.38 \pm 0.45 \pm 0.21$		
$\Upsilon(3S)\pi^+$	$2.40 \pm 0.58 \pm 0.36$	$1.62 \pm 0.50 \pm 0.24$		
$h_b(1P)\pi^+$	$4.26 \pm 1.28 \pm 1.10$	$9.23 \pm 2.88 \pm 2.28$		
$h_b(2\mathrm{P})\pi^+$	$6.08 \pm 2.15 \pm 1.63$	$17.0 \pm 3.74 \pm 4.1$		
$B^+ \bar{B}^{*0} + \bar{B}^0 B^{*+}$	$82.6 \pm 2.9 \pm 2.3$	-		
$B^{*+}\bar{B}^{*0}$		$70.6\pm4.9\pm4.4$		

Model-0 :  $Z_b(10650)$  only Model-1:  $Z_b(10610)$  + Non-res. Model-2:  $Z_b(10610)$  +  $Z_b(10650)$ with interference Model-3: Non-resonance

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## X(5568) – puzzle ?









## Pc States



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## Open the pentaquark door: LHCb observation in 2015

- Two  $J/\psi p$  resonant structures are revealed by a full 6D amplitude analysis
  - $P_c(4450)^+$   $\leftarrow$  the prominent peak
  - $P_c(4380)^+$   $\leftarrow$  required to obtain a good fit to the data
  - Consistent with **pentaquarks** with minimal quark content of  $uudc\overline{c}$

26k  $\Lambda_b$  signals PRL 115 (2015) 072001 (most cited paper at LHCb so far)





	$P_{c}(4380)^{\pm}$	$P_{c}(4450)^{\pm}$	
Mass (MeV)	$4380\pm8\pm29$	$4449.8 \pm 1.7 \pm 2.5$	
Width (MeV)	$205\pm18\pm86$	$39\pm5\pm19$	
Fit Fraction (%)	$8.4\pm0.7\pm4.2$	$4.1\pm0.5\pm1.1$	

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## Lots of open questions

 $\Sigma_c^+$ 

- To interpret the nature of  $P_c$ , more studies are needed
  - Inner structures?
  - More states, SU(3) partners?
  - $J^P$ , mode decay modes, production mechanism ...?













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## Fine structures from update

- Run1+Run2, x10  $\Lambda_b^0 \rightarrow J/\psi p K^-$  yield
  - Inclusion of Run 2 data (x 5)
  - Improved data selection (x 2)
- $P_c(4312)^+$  is observed
- $P_c(4450)^+$  peak structure is an overlap of two narrower states,  $P_c(4440)^+$  and  $P_c(4457)^+$
- Their near-threshold masses **favor** the predicted "molecular" pentaquarks with meson-baryon substructure, but **other hypotheses are not ruled out**

State	$M \;[\mathrm{MeV}\;]$	$\Gamma \;[\mathrm{MeV}\;]$	(95%  CL)	$\mathcal{R}~[\%]$
$P_c(4312)^+$	$4311.9 \pm 0.7^{+6.8}_{-0.6}$	$9.8 \pm 2.7^{+}_{-} \stackrel{3.7}{_{-}}{_{-}}$	(< 27)	$0.30 \pm 0.07^{+0.34}_{-0.09}$
$P_c(4440)^+$	$4440.3 \pm 1.3^{+4.1}_{-4.7}$	$20.6 \pm 4.9^{+8.7}_{-10.1}$	(< 49)	$1.11 \pm 0.33^{+0.22}_{-0.10}$
$P_c(4457)^+$	$4457.3 \pm 0.6^{+4.1}_{-1.7}$	$6.4 \pm 2.0^{+}_{-} \stackrel{5.7}{_{-}1.9}$	(< 20)	$0.53 \pm 0.16^{+0.15}_{-0.13}$





1D  $m_{J/\psi p}$  is fitted, ongoing amplitude analysis is in advanced stage







## 1<sup>st</sup> observation of $\Lambda_h^0 \to \eta_c p K^-$

- $\eta_c p$  final state is very sensitive to  $1/2^- P_c$ , where  $\eta_c p$  is in S-wave
- If  $P_c(4312)^+$  is  $\Sigma_c \overline{D}$  molecule, predicted  $\underline{\mathcal{B}(P_c(4312)^+ \rightarrow \eta_c p)}_{\sim 3}$ [PRD 100 (2019) 034020, 100 (2019) 074007, 102 (2020) 036012]  $\mathcal{B}(P_c(4312)^+ \rightarrow J/\psi p)$

Swapped protons

5700

 $m(p\overline{p}pK)$  [MeV/ $c^2$ ]

5750

5650

• LHCb run2 data (5.5 fb<sup>-1</sup>) using  $\eta_c \rightarrow p\bar{p}$ 

5500

5550

5600

• Fit 2D mass spectrum to confirm the existence





2800

2900

3000

3100

 $m(p\overline{p})$  [MeV/ $c^2$ ]

3200

 $P_c(4312)^+$  production fraction in  $\Lambda_h^0 \to \eta_c p K^$ is  $\sim 3\%$  (predicted)



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### Search for $P_c^+$ in $\eta_c p$ system [PRD 102 (2020) 112012]

• Check background-subtracted  $\eta_c p$  mass spectrum

No significant  $P_c(4312)^+$  contribution (~2s)

 $P_c^+$  production fraction obtained  $R(P_c(4312)^+) < 24\% @ 95\%$  C. L.

much larger than the predicted value 3% (no conclusion yet)

• Need run3+4 data, amplitude fit can be performed







## Search for pentaquarks via open charm

- Prompt production with 32 final states
  - $\Lambda_c^+ \overline{D}, \Lambda_c^+ \overline{D}^*, \Lambda_c^+ \pi \overline{D}, \Sigma_c^{(*)} \overline{D}^{(*)} \text{ and } \Lambda_c^+ D, \Lambda_c^+ D^*, \Lambda_c^+ \pi D, \Sigma_c^{(*)} D^{(*)}$
- Scan to search for pentaquarks with narrow width (0-15 MeV)
- No significant narrow peak is found for all the modes
- Upper limits are set on the production rates related to  $\Lambda_c^+$  $R = \frac{N_{P_c}}{N_c} \times \frac{\varepsilon_{\Lambda_c^+}}{N_c} \longrightarrow \frac{\sigma(P_c) \times \mathcal{B}(P_c \to \Lambda_c^+ D(\pi)) \times \mathcal{B}(D)}{\sigma(P_c \to \Lambda_c^+ D(\pi)) \times \mathcal{B}(D)}$

	$n = N_{\Lambda^+} \sim \varepsilon_{P_C}$				$\sigma(\Lambda_c^+)$			
	Decay Mode	Signific Local	ance $(\sigma)$ Global	Corresponding Mass (MeV/c <sup>2</sup> )	Signal Yield	Upper Lin 90% CL	nit (×10 <sup>-3</sup> ) 95% CL	
	$\Lambda_c^+ \overline{D}{}^0$	2.85	1.01	349	$\textbf{46.8} \pm \textbf{23.4}$	1.16	1.21	
	$\Lambda_c^+ D^{*-}$	2.32	0.00	365	$\textbf{15.0} \pm \textbf{10.3}$	2.16	2.39	
	$\Lambda_{c}^{+}\pi^{+}D^{-}$	2.82	0.99	225	$\textbf{68.6} \pm \textbf{13.3}$	1.95	2.40	T a war a st
	$\Sigma_{c}^{0}\overline{D}^{0}$	1.90	0.00	65	47+42	1.02	1 15	_ Largest
I	$\Lambda_{c}^{+}\pi^{-}\overline{D}^{0}$	3.86	2.56	45	$\textbf{60.1} \pm \textbf{25.9}$	1.40	1.70	aignificance
1	$\Sigma_c^0 \bar{D}^-$	2.03	0.00	261	$7.0\pm~2.6$	0.71	0.89	
	$\Lambda_{m{c}}^+\pi^- D^-$	3.67	2.35	249	$\textbf{82.8} \pm \textbf{14.3}$	2.23	2.67	
	$\Lambda_{c}^{+}\pi^{-}D^{*-}$	2.31	0.00	409	$\textbf{23.6} \pm \textbf{23.0}$	2.79	3.28	
	$\Sigma^{*++}_{c} D^{*-}$	1.74	0.00	453	$3.3\pm~2.4$	1.24	1.43	
	$\Sigma_c^{*0} D^-$	1.86	0.00	109	$\textbf{10.7} \pm \textbf{29.1}$	1.32	1.59	
	$\Lambda_c^+ D^+$	2.52	0.59	169	$14.9\pm9.6$	1.34	1.50	
	$\Lambda_c^+ \pi^+ D^0$	3.21	1.72	45	$\textbf{24.8} \pm \textbf{39.3}$	0.98	1.18	
	$\Lambda_c^+ \pi^+ D^{*+}$	3.37	1.99	165	$13.8\pm3.5$	0.97	1.22	
	$\Lambda_c^+ \pi^- D^{*+}$	2.70	0.58	73	$\textbf{5.8} \pm \textbf{71.3}$	1.70	1.94	
	$\Sigma_c^{*++} D^0$	2.11	0.00	113	$3.9\pm~2.8$	0.87	0.99	
	$\Sigma_c^{*0}D^+$	2.18	0.00	69	$4.7\pm4.6$	1.13	1.32	
1.2	HALL BURNER IN THE CASE OF THE PARTY OF THE	11		C. ma		CP1994		

#### [LHCb-PAPER-2023-018]







# **系**

## Evidence of $P_{cs}$ in $J/\psi \Lambda$

- SU(3) partner  $P_{cs}$  is predicted, and suggested to search for in  $\Xi_b^- \rightarrow J/\psi \Lambda K^-$ [JJ Wu PRL 105 (2010) 232001; HX Chen PRC 93(2016) 064203]
- Amplitude analysis with improved helicity formalism
  - $P_{cs}(4459)^0$  found, significance >3.1s





Mass is about 19 MeV below  $\Xi_c^0 \overline{D}^{*0}$  threshold



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## $P_{cs}$ in $B^- \to J/\psi \Lambda \overline{p}$

[PRL 131 (2023) 031901]

- Can search for pentaquark both in  $J/\psi p \& J/\psi \Lambda$ 
  - Limited range:  $m(J/\psi p) < 4.16 \text{ GeV}, m(J/\psi \Lambda) < 4.34 \text{ GeV}$ , Cover thresholds of  $\Lambda_c^+ \overline{D}$  and  $\overline{\mathcal{Z}_c D}$
- A new pentaquark with strangeness  $P^{\Lambda}_{\psi s}(4338)^0$  (*cc̄sud*) observed in the *B*<sup>-</sup>



- $\rightarrow J/\psi \Lambda \bar{p}$  decay
  - At  $\mathbf{Z}_c^+ \mathbf{D}^-$  threshold
  - $m = 4338.2 \pm 0.7 \pm 0.4$  MeV
  - $\Gamma = 7.0 \pm 1.2 \pm 1.3 \text{ MeV}$
  - Fit fraction =  $(12.5 \pm 0.7 \pm 1.9)\%$
  - $J^P = (1/2)^-$  preferred,



















## $di-J/\psi$ States



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### X(4c) states at ATLAS --All-charm Tetra-quarks



• In the di- $J/\psi$  channel, two signal models are tested:

- Model A: three interfering signal peaks; Model B: two signal peaks
- The peak around 6.9 GeV is consistent with the LHCb observed X(6900) (arXiv:2006.16957), with significance far above  $5\sigma$





### X(4c) states at ATLAS



• Model  $\beta$ : only one signal peak

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### X(4c) structures in di-Jpsi channel at CMS



- Interference model: Phys. Rev. Lett. 132 (2024) 111901
  - Signal: interference between BW1, BW2, BW3
  - Background: BW0 + NRSPS + NRDPS

	$BW_1$	BW <sub>2</sub>	BW <sub>3</sub>
m (MeV)	$6638^{+43+16}_{-38-31}$	$6847^{+44+48}_{-28-20}$	$7134_{-25-15}^{+48+41}$
$\Gamma$ (MeV)	$440^{+230+110}_{-200-240}$	$191_{-49-17}^{+66+25}$	$97^{+40+29}_{-29-26}$

CMS found 3 significant  $J/\psi J/\psi$  structures using Run II data

- BW2 consistent with X(6900) reported by LHCb [Sci. Bull. 65, 1983 (2020)]
- Two new structures named as X(6600) [>5σ], X(7300) [4.7σ]

A family of structures which are candidates for all-charm tetra-quarks

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### Comparison with some theoretical calculations







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# XYZ Summary



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35 new hadrons were found at Belle. **10 of these are "exotic"** and cannot be explained in the conventional quark model while the nature of 8 of them are still under investigation. The remaining 17 states are consistent with the quark model. Measurements of all these states will provide critical insights for QCD.



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11.0 11.0 0<sup>*χ<sub>b</sub>*(3*P*)</sup>  $\bullet_{\chi_{b1}(3P)}^{\chi_{b2}(3P)}$ 10.5 10.5 23 new exotic hadrons at the LHC 72 new hadrons at the LHC 7.5-7.5 • T<sub>ov</sub>(6900)  $O^{B_c(2S)^+}$ -ww(6900) 7.0 -7.0 B<sup>\*</sup><sub>c</sub>(25)<sup>+</sup>  $B_{c}(2S)^{+}$ *T<sub>φφ</sub>*(6600)  $\Omega_{b}(6350)^{-1}$ 6.5 6.5  $\Lambda_b(6152)^0 \Omega_b(6340)$ E<sub>b</sub>(6227)<sup>0</sup> = E<sub>b</sub>(6327) Ξ<sub>b</sub>(6227)<sup>-</sup>  $\Lambda_{b}(6146)^{0}$  $\Xi_b(6095)^0$  $\Xi_b(6087)^0$  $\Xi_b(5945)^0 \wedge_b(5920)^0 \wedge_b(5912)^0$  $\Xi_b(5955)^- B_J(5970)^{+,0}$ 6.0 6.0 E<sub>b</sub>(6100) Ξ<sub>b</sub>(5935)- B<sub>j</sub>(5840)<sup>+,0</sup>  $\Sigma_{b}(6097)^{+}$  $\Lambda_b(6070)^0 B_s^*(6114)^0$ Σ<sub>b</sub>(6097)<sup>-</sup> B.\* (6063)0 X(4700) X(4500) X(4685) X(4630) bb X(4700) X(4685) X(4630) P\_v^N(4457)+  $P_{\psi}^{N}(4450)^{+}$ *X*(4500)  $P_{w}^{N}(4457)^{+}$  $P_w^N(4450)^+$ bą P<sub>w</sub><sup>N</sup>(4380)<sup>4</sup>  $P_{\varphi}^{N}(4440)^{+} P_{\varphi}^{N}(4312)^{+}$ P^{(4338)<sup>0</sup> X(4274)  $P_{\psi}^{N}(4440)^{+} P_{\psi}^{N}(4312)^{+}$ P^{(4338)0  $T_{ys1}(4220)^+$  $T_{ys1}^{\theta}(4000)^+$ X(4274)  $T_{\psi s1}(4220)^+$ X(4140) T<sup>e</sup><sub>9/51</sub>(4000)<sup>c</sup> cc̄(qq̄) X(4140) P\_m^N(4380)  $T_{\psi s1}^{\theta}(4000)^{+}$ \_ψ₃(3842)  $T^{\theta}_{\psi s1}(4000)$ ψ<sub>3</sub>(3842) ccą́ą Tcc(3875)+ X(3960) (3875)+ X(3960) cc(qq) Ξ;;; cēcē 3.5 3.5 Ω<sub>c</sub>(3327)<sup>0</sup> ccą̃ą сą  $\Omega_{c}(3119)^{0}$  $\Omega_c(3185)^0$ D,\*(3000)+,0  $\Omega_c(3090)^0$  $\Omega_c(3066)^0$ ■ T<sub>cs0</sub>(2900)<sup>0</sup> T<sub>cs1</sub>(2900)<sup>0</sup> 3.0 cēcē cāqā T<sup>a</sup><sub>cš0</sub>(2900)<sup>++</sup> T<sup>a</sup><sub>cš0</sub>(2900)<sup>0</sup> ٠ D<sup>\*</sup><sub>s1</sub>(2860)<sup>+</sup> Ξ<sub>c</sub>(2939)<sup>0</sup> 3.0 D<sub>1</sub>(3000)<sup>0</sup>  $T_{cs0}^{(2900)^0} \\ T_{cs1}^{(2900)^0}$  $\Lambda_{c}(2860)^{+}$  $T^{a}_{c\bar{s}0}(2900)^{++}$  $\Omega_c(3050)^0$  $\Omega_c(3000)^0$ bqq cāgā D<sub>1</sub>\*(2760)+ Ξ<sub>c</sub>(2923)<sup>0</sup> T<sup>a</sup><sub>c50</sub>(2900)<sup>0</sup> ٠ D<sub>1</sub>(2740)<sup>0</sup>  $D_3^*(2760)$ 2.5 cqq 2.5 D<sub>s0</sub>(2590)<sup>+</sup> cēqqq  $D_{1}(2580)^{0}$ cēqqq 2.0 2.0 2017 2018 2019 2020 2021 2011 2012 2013 2014 2015 2016 2022 2023 2011 2012 2013 2014 2015 2016 2017 2018 2019 2020 2021 2022 2023 Date of arXiv submission patrick.koppenburg@cern.ch 2023-08-16 Date of arXiv submission patrick.koppenburg@cern.ch 2023-08-16 Particle "Zoo" again ! 150

72 new hadrons were found at LHC. 23 of these are "exotic"















## Too many models !

<complex-block>

- Theory 1: screened potential
- Theory 2: hybrids with excited gluons
- Theory 3: tetraquark states
- Theory 4: meson molecules
- Theory 5: cusps effect
- Theory 6: final state interaction
- Theory 7: coupled-channel effect
- Theory 8: mixing of normal quarkonium and exotics
- Theory 9: mixture of all these effects
- Theories ... We need clear features to identify exotic hadronic states !

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## We found more questions to answer, works to do

- In the Experiments sector
  - Search for flavor analog exotic states  $(Z_s, X_b, ...)$
  - Confirm marginal states (X(3940), Y(4008), Z<sub>1</sub>(4050), X(4160), Z<sub>2</sub>(4250), X(4350)....)
  - Search for missing charmonium/bottomonium states ( $\eta_{c2}$ ,  $h_c(2P)$  ... )
  - Are there excited  $Z_c$  states and  $Z_{cs}$  states  $[D^*D_s \text{ or } DD_s^*]$ ?
  - Search for flavor analogs of the  $P_{cs}$  ( $P_s$ , ...)
  - Search for quantum number partners of XYZ states
  - Precise measurements of relative strength to different final states
  - Check more di-charmonium systems or di-bottomonium systems
  - Correlation between charm production & charmonium transitions?
  - Make experimental results more accessible for subsequent interpretation (publish Dalitz plot in text format, supply also efficiency curve ...)
  - Publish upper limits for negative searches

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## We found more questions to answer, works to do

- In the Theory sector
  - Study exclusive e<sup>+</sup>e<sup>-</sup> cross sections using better coupled-channel formalism
  - Give differences in key physical quantities to distinguish between different interpretations (molecule, hybrid, tetraquark state, ...)
  - Improve parameterizations of the data (when appropriate and beneficial, experimentalists and theorists directly work together)
  - theorists, when possible, to publish complete functional forms









### More data, more surprises, more opportunities



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