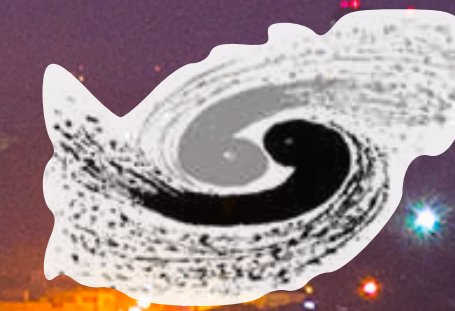


LHC and beyond

João Guimarães da Costa



中国科学院高能物理研究所

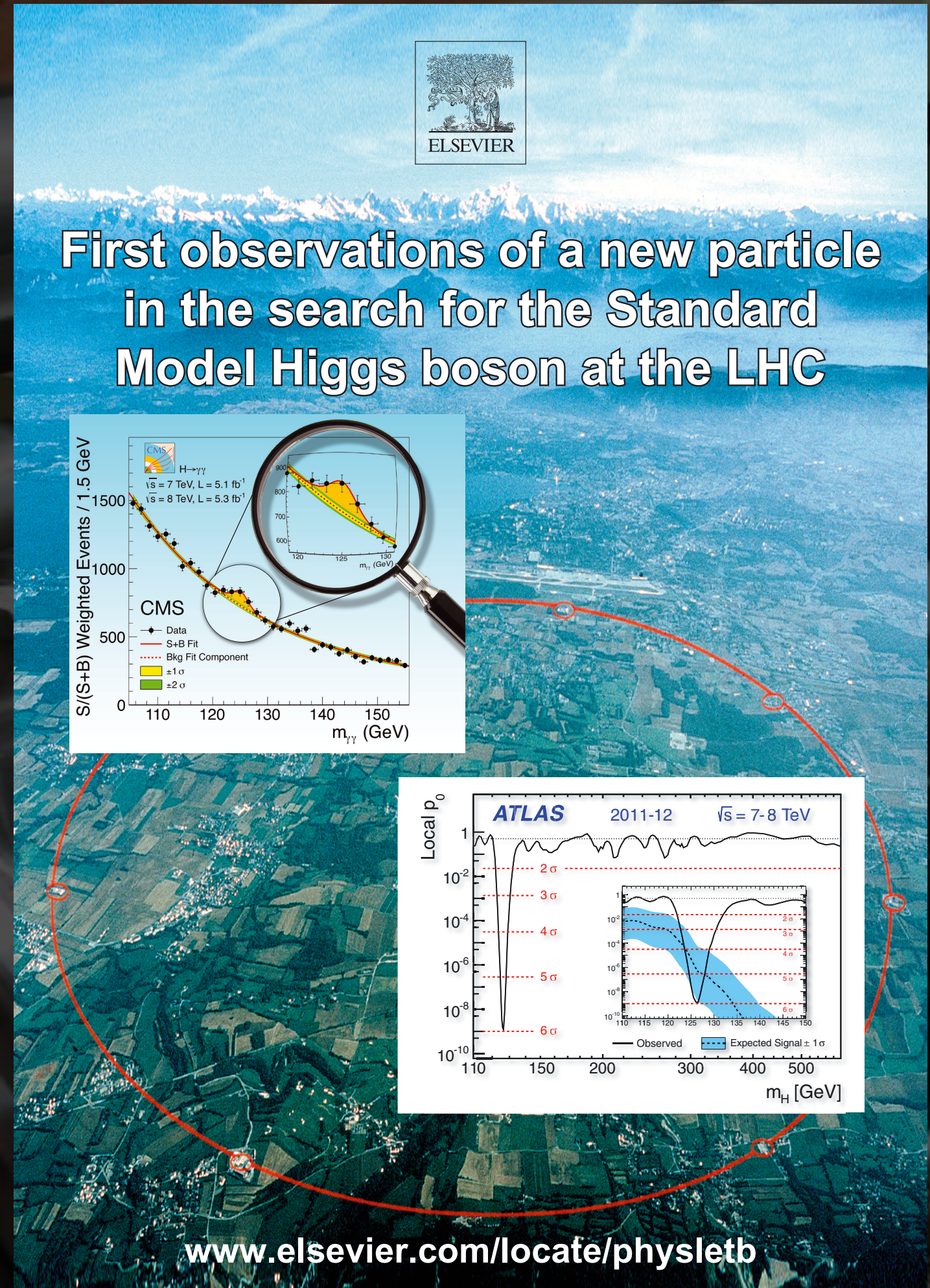
*Institute of High Energy Physics
Chinese Academy of Sciences*

Nakhon
Pathom
Thailand
June 23, 2024

The Higgs Boson Discovery at LHC

Predicted in 1964, discovered in **2012!** 48 year hunting!

ATLAS & CMS Observation



2013 Nobel Prize



An effort by thousands of scientists and engineers from all over the world



François Englert and Peter Higgs

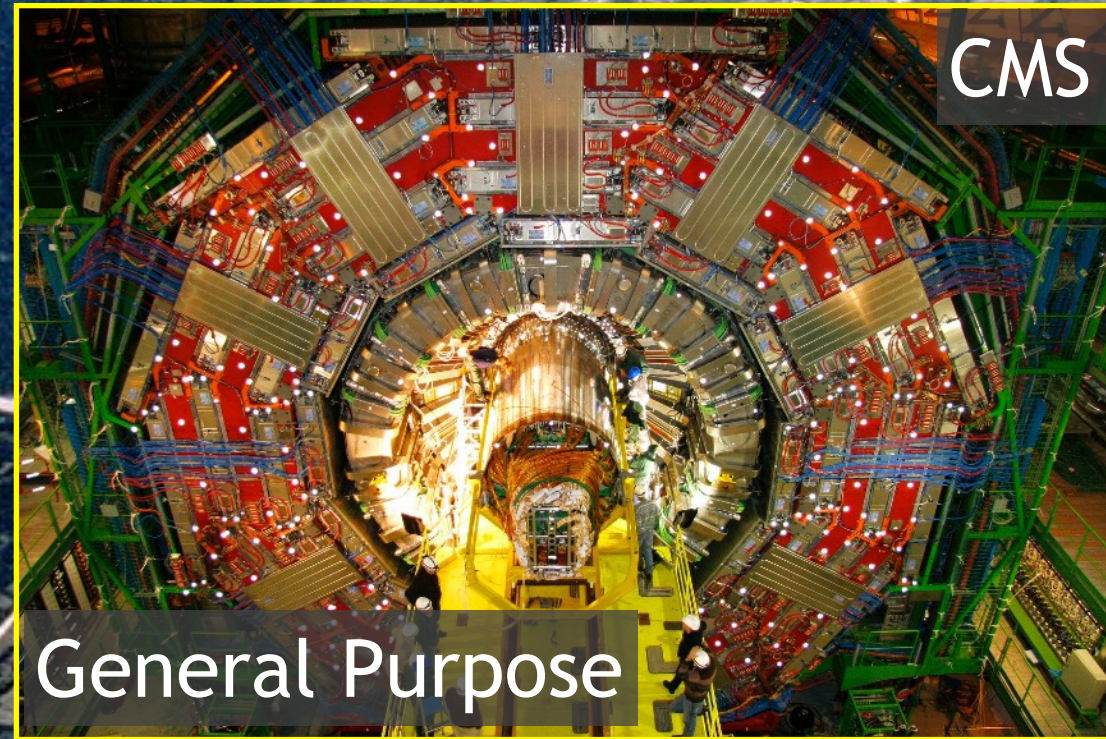
Huge impact to humanity

Technology
Cultural
International Collaboration

**How did we get there
and
what are we doing now?**

Large Hadron Collider proton-proton collisions Center of mass energy: 7-8-14 TeV

Lake Geneva

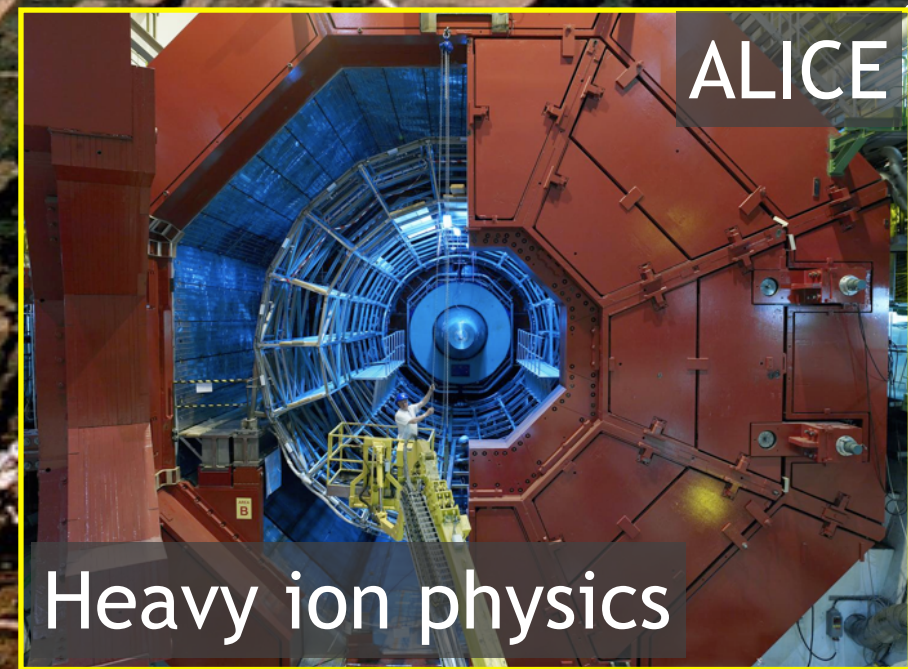


LHC ring:
27 km circumference

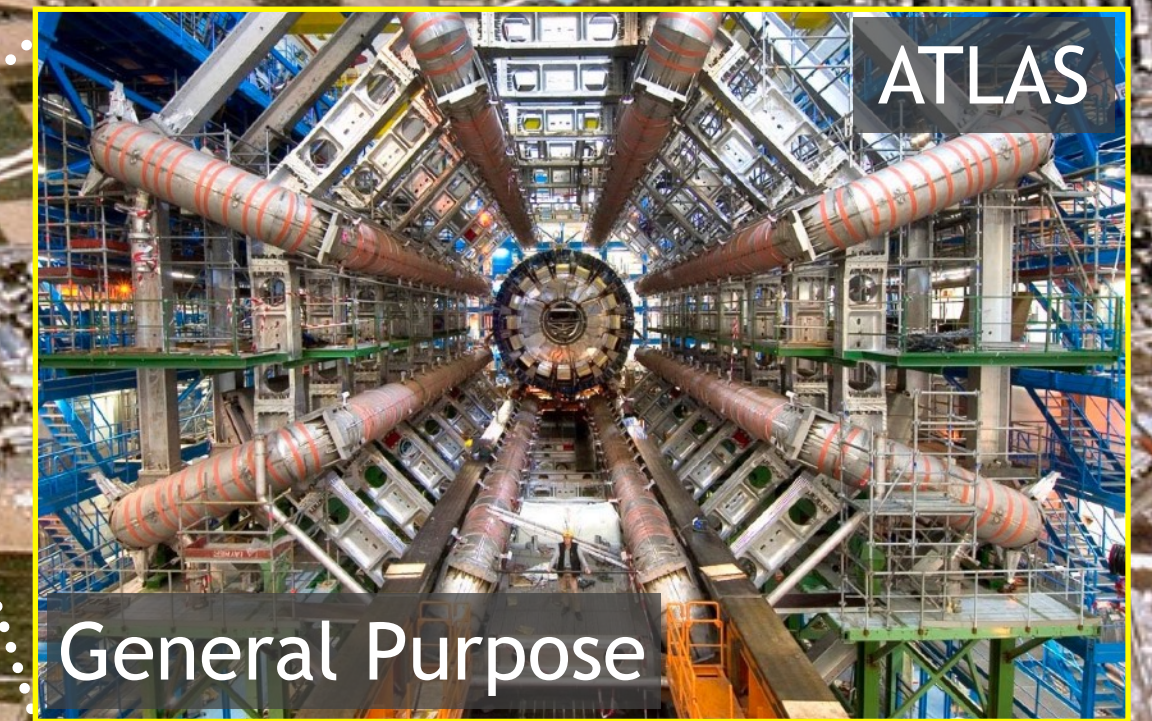


SPS ring

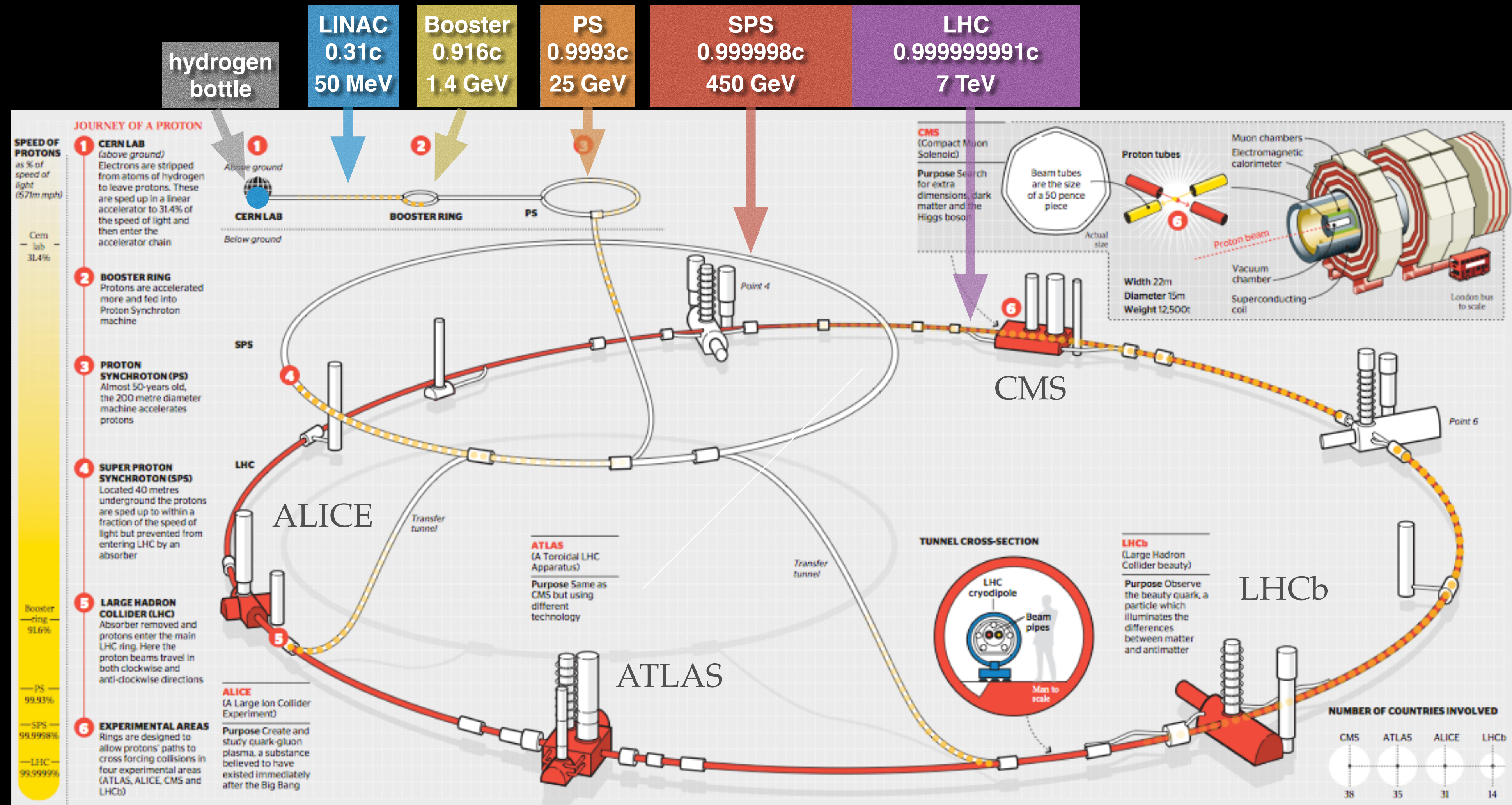
Airport



CERN



The LHC accelerator complex

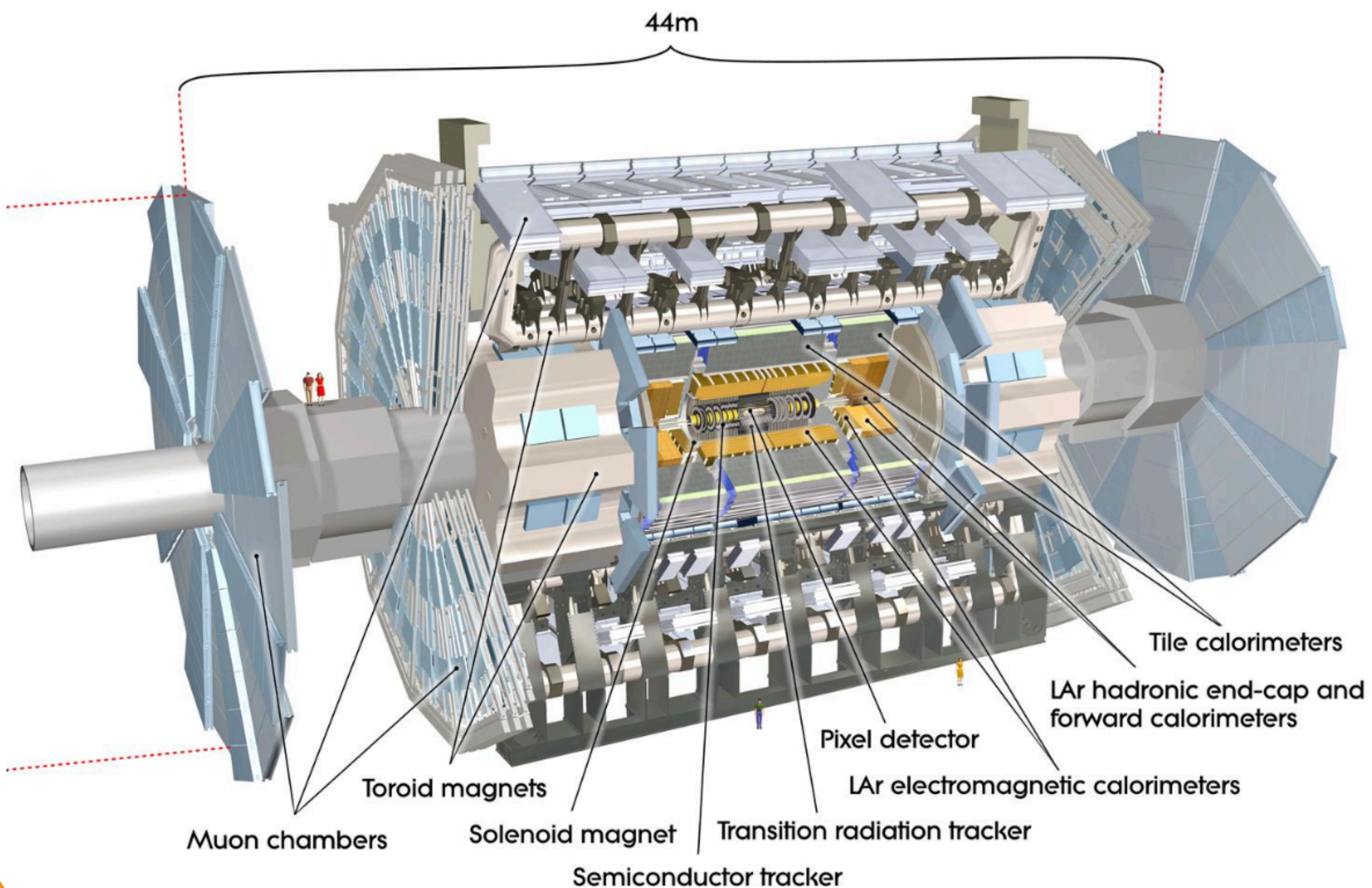


Depth ~ 50-175 m

LHC Experiments: General Purpose Detectors

ATLAS

- 25 m diameter and 44 m length
- Over 7000 tons
- O(100) million readout channels

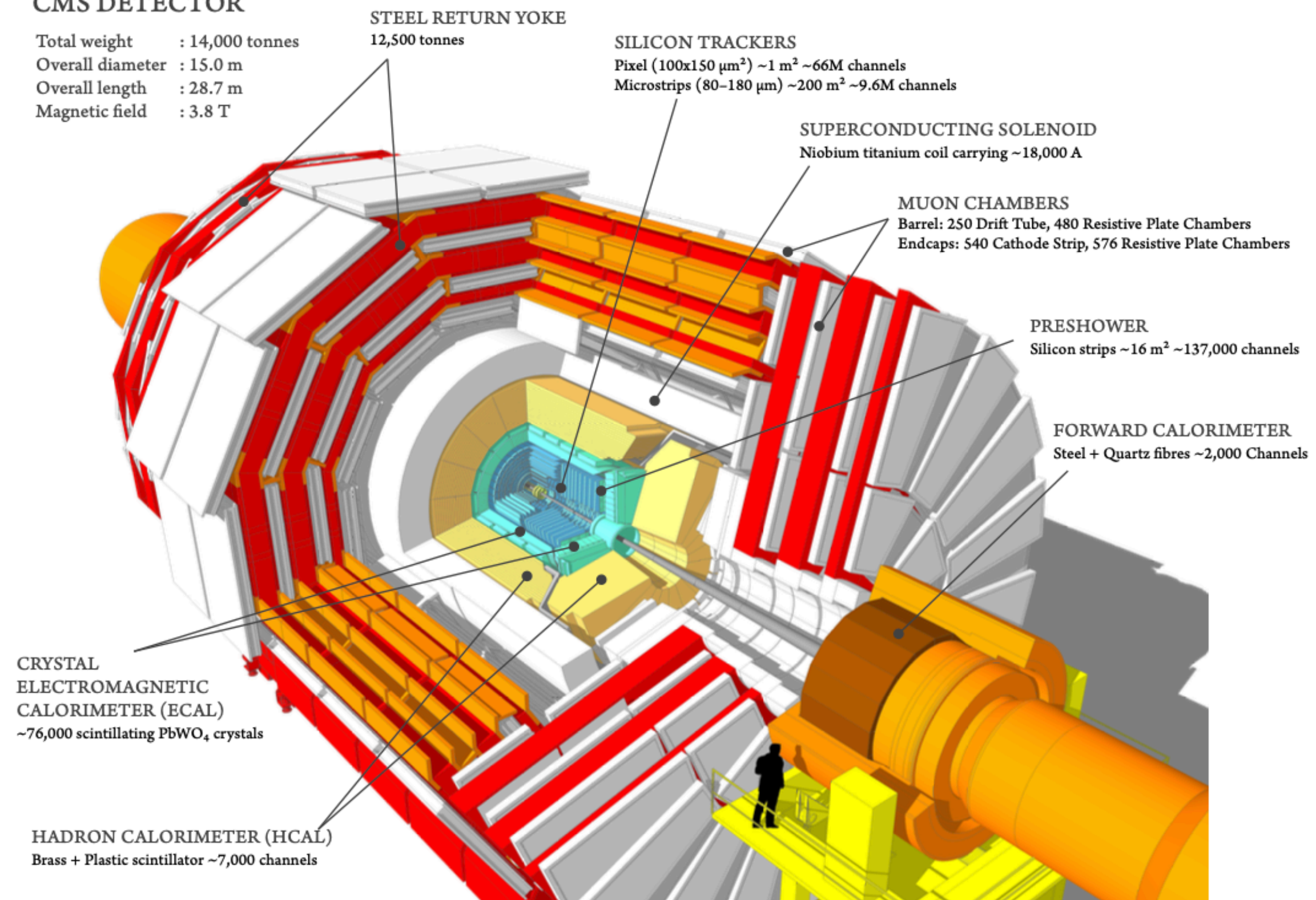


CMS

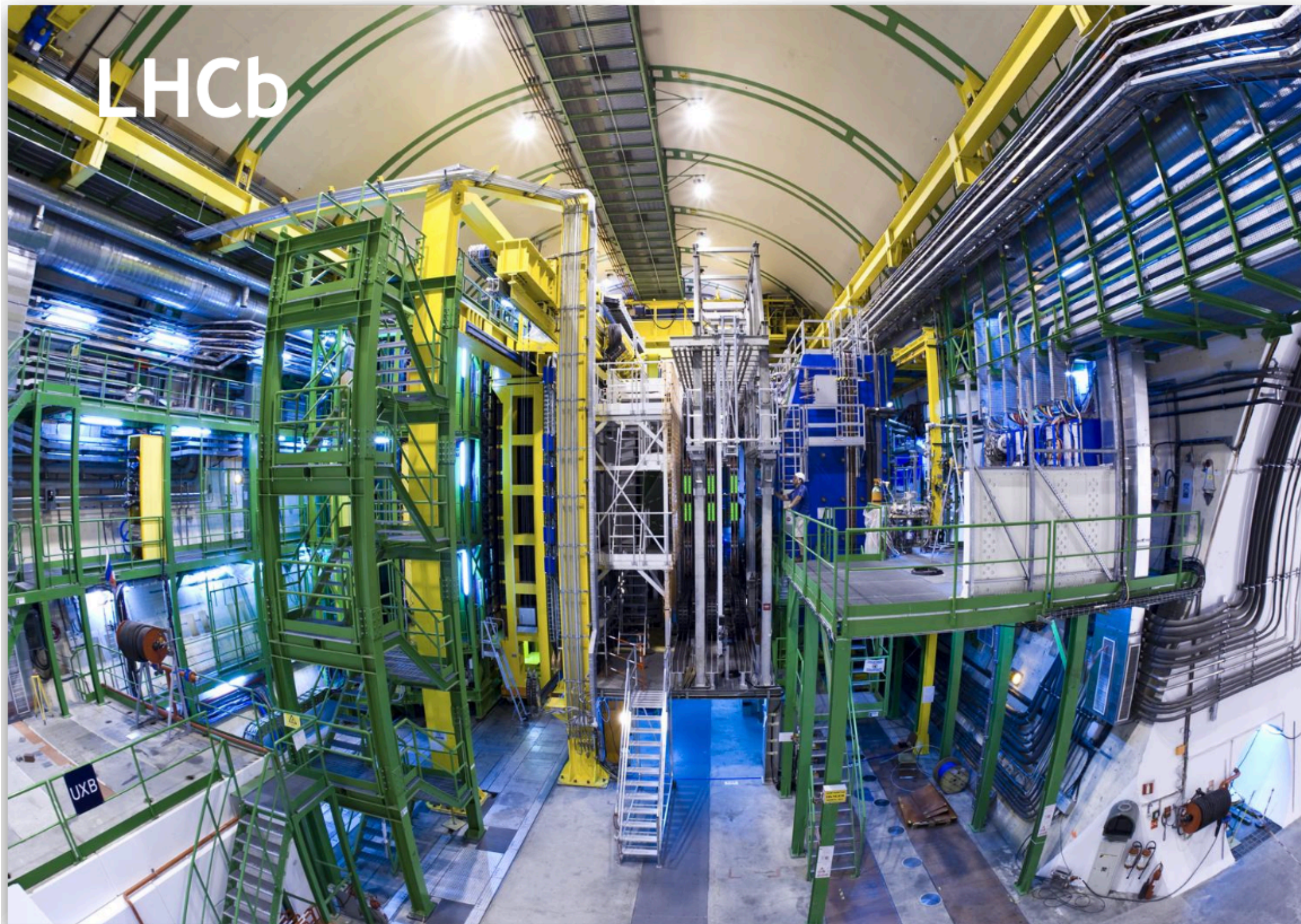
- 15 m diameter and 21 m length
- 14000 tons
- O(75) million readout channels

CMS DETECTOR

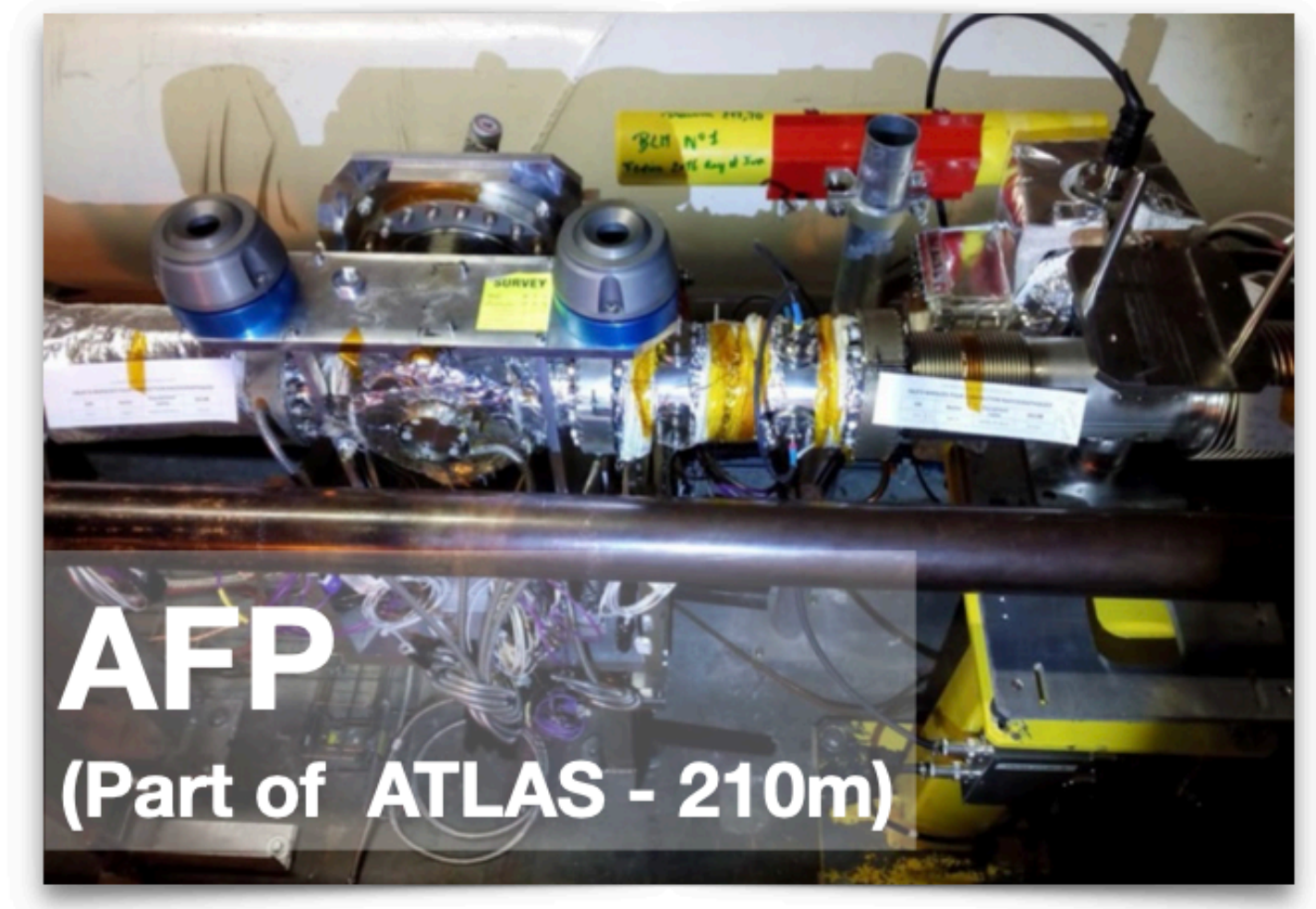
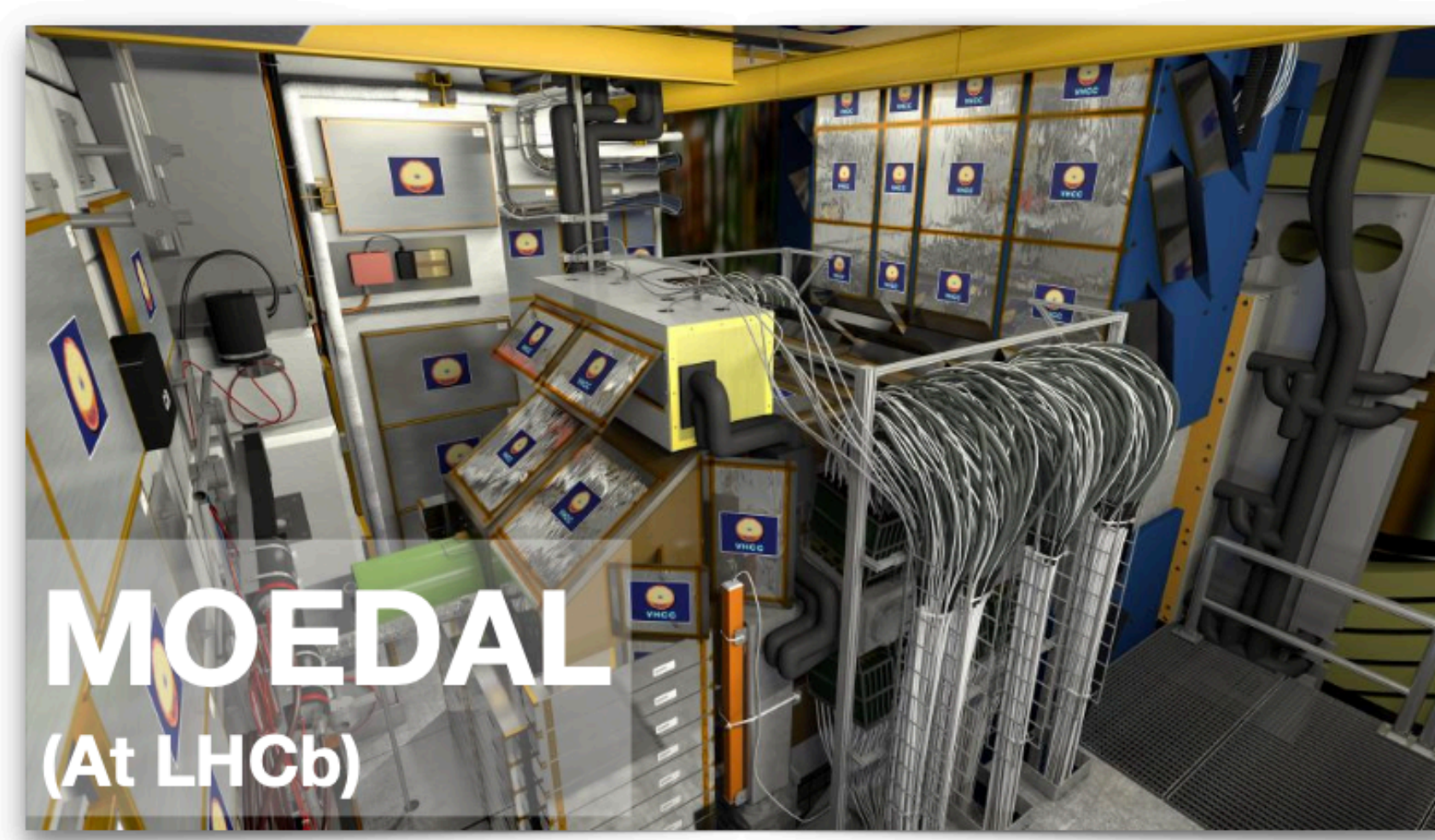
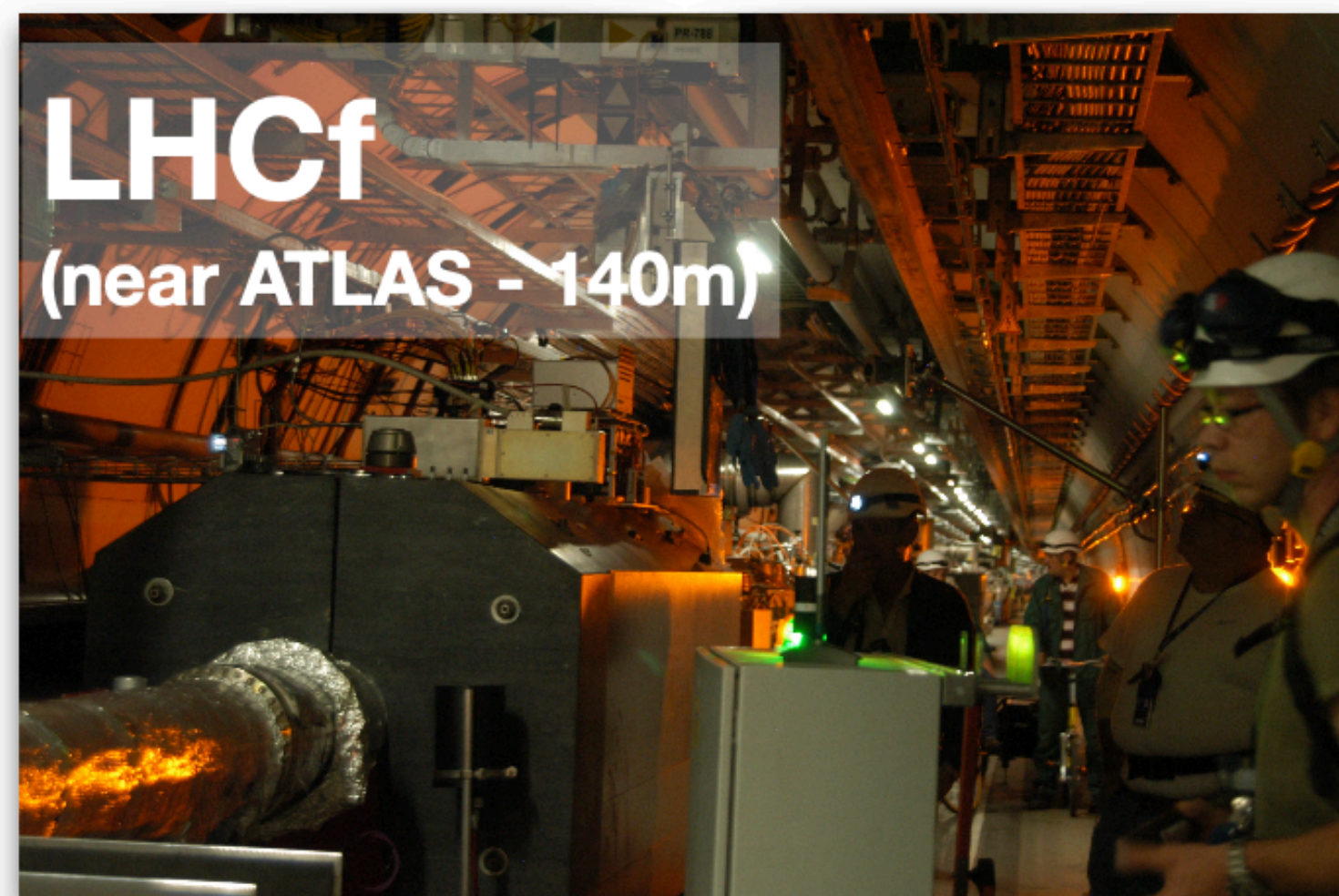
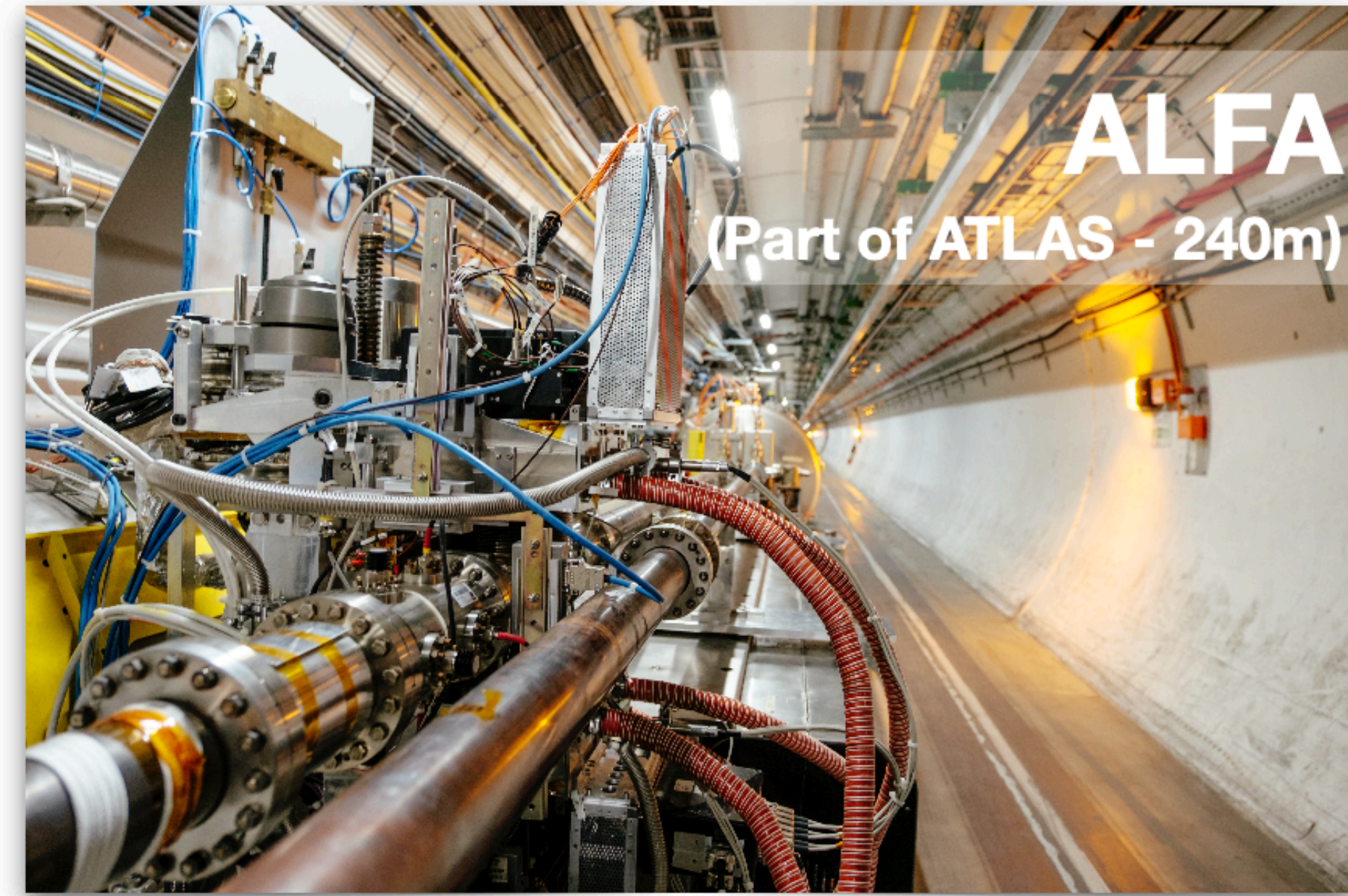
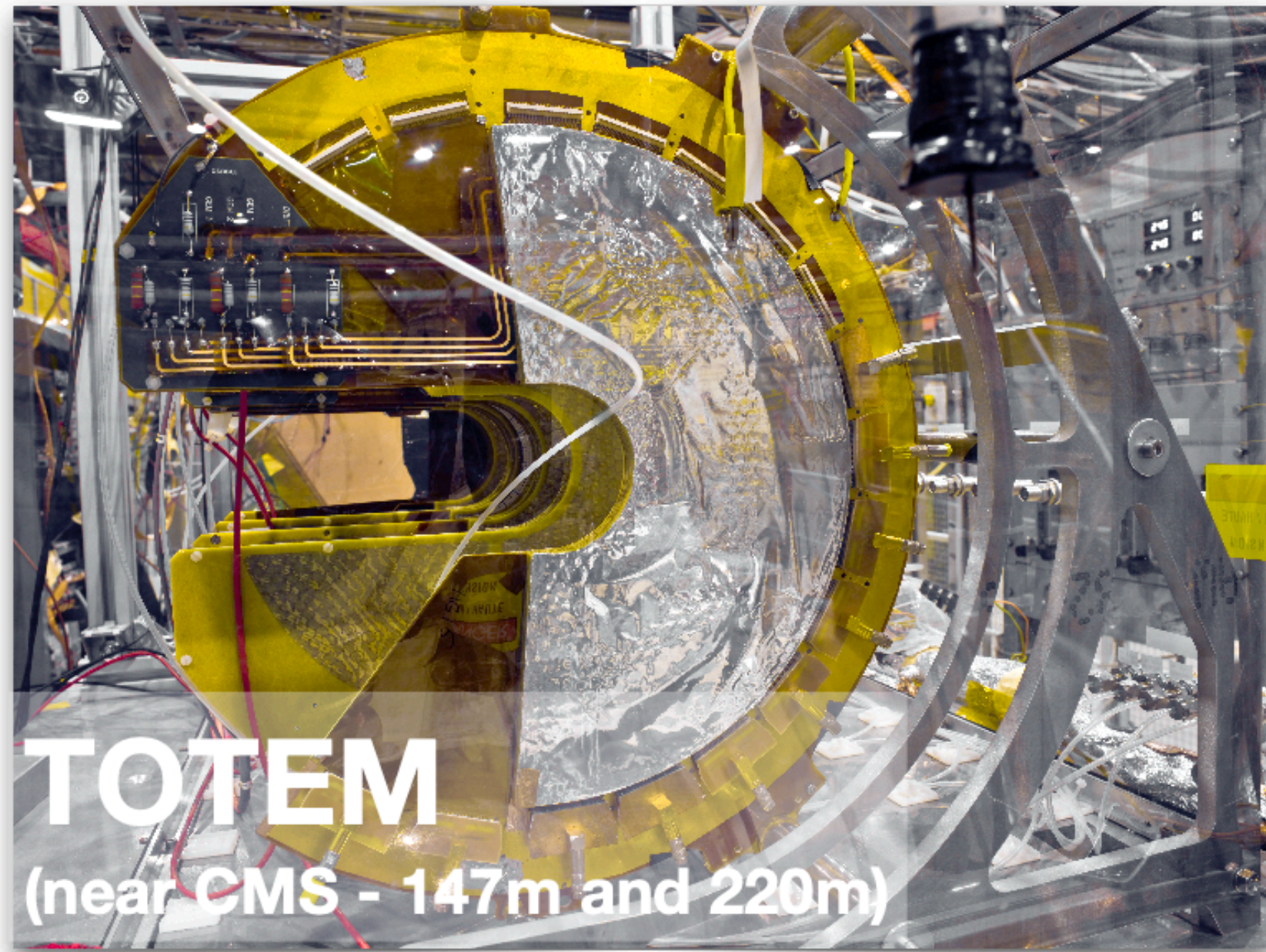
Total weight : 14,000 tonnes
 Overall diameter : 15.0 m
 Overall length : 28.7 m
 Magnetic field : 3.8 T



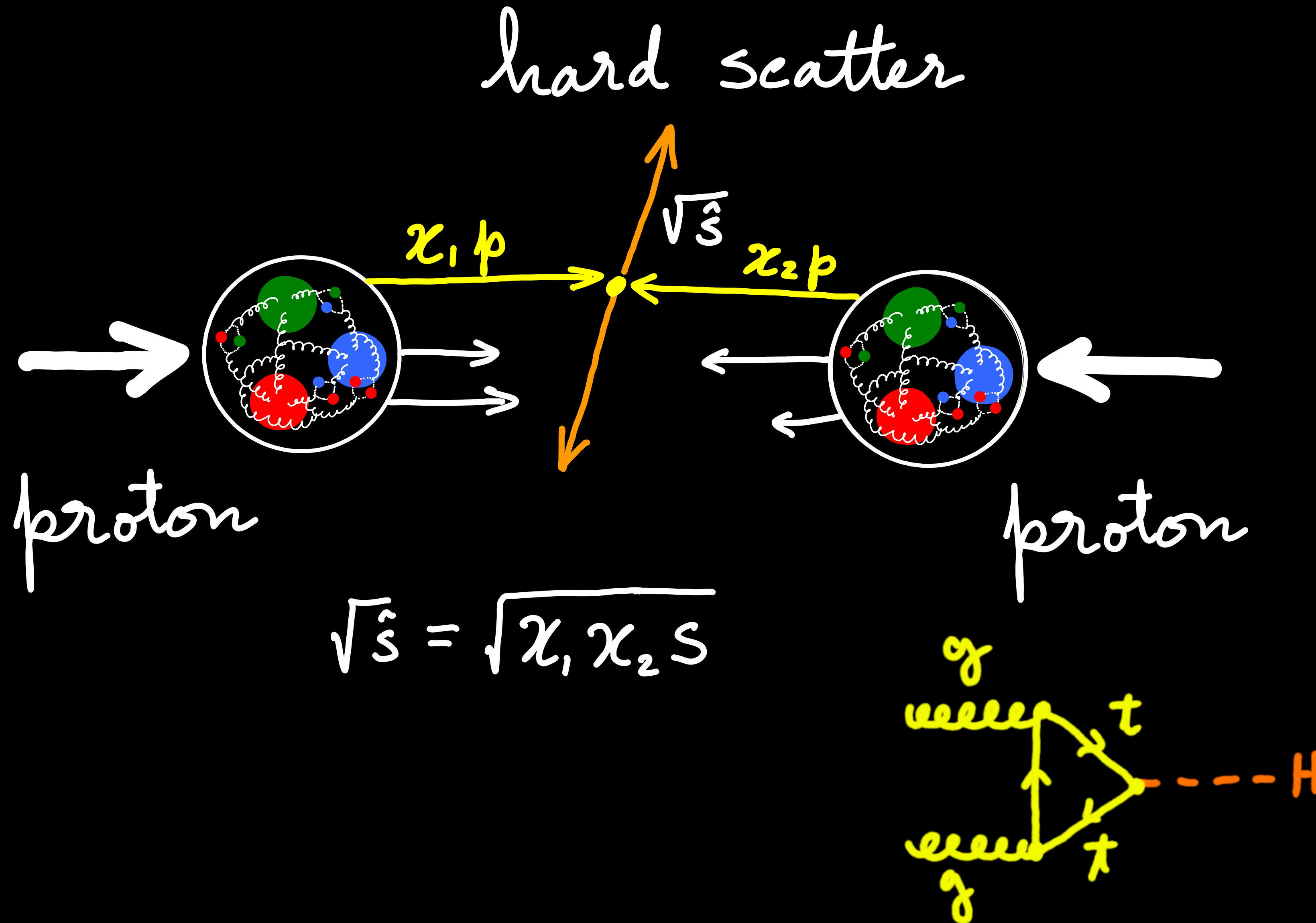
LHC Experiments: Specialized Detectors



Other even more specialized experiments



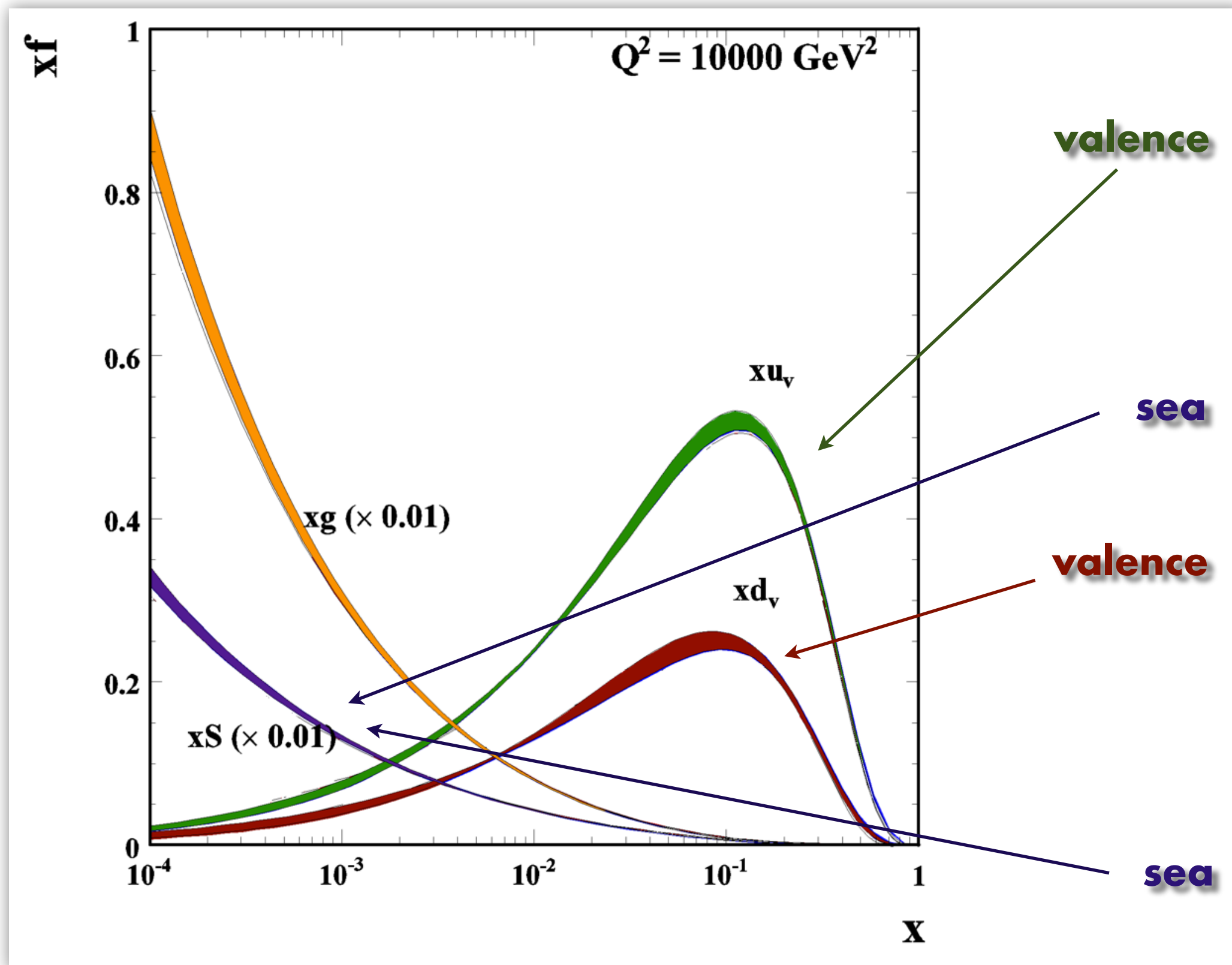
Look into the pp collisions...



Parton Distribution Functions (PDF)



- PDFs give the probability to find a parton with a momentum fraction of x
- PDFs are not calculable, but measured in experiments



PDFs Sum rules

Momentum sum rule

$$\sum_i \int_0^1 dx x f_i(x, Q^2) = 1$$

Flavour conservation sum rules

$$\int_0^1 (f_u(x, Q^2) - f_{\bar{u}}(x, Q^2)) dx = 2$$

$$\int_0^1 (f_d(x, Q^2) - f_{\bar{d}}(x, Q^2)) dx = 1$$

$$\int_0^1 (f_s(x, Q^2) - f_{\bar{s}}(x, Q^2)) dx = 0$$

- In a proton, quarks heavier than u, v are only available in the sea

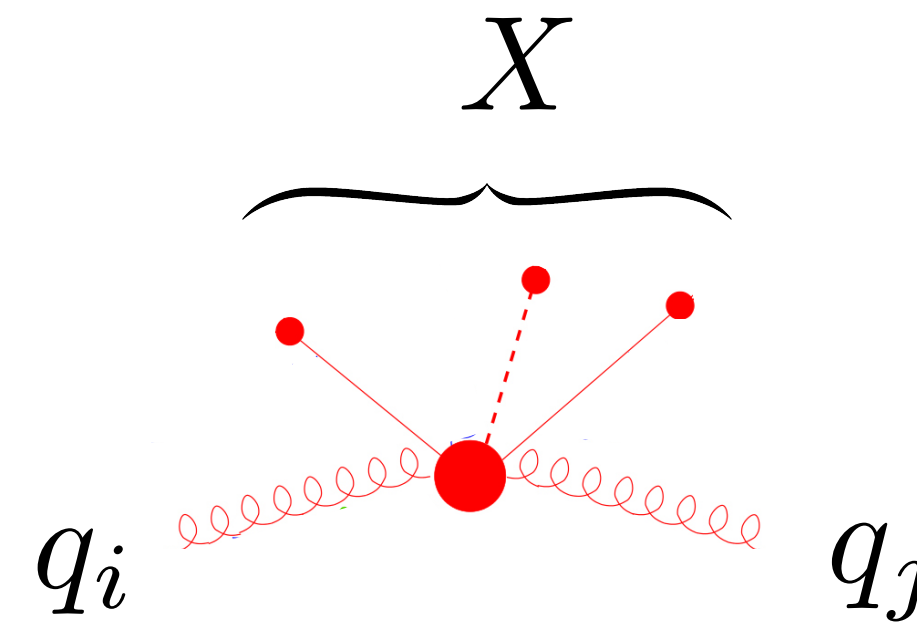
Hadronic Collisions Phenomenology

Truth-level event

$$a(q_i q_j \rightarrow X) = a_0 \times (1 + \alpha_S a_1 + \alpha_S^2 a_2 + \alpha_S^3 a_3 + \dots)$$

$$\alpha_S \sim 0.1$$

α_S - strong coupling



Hard scattering

Perturbative regime (with large QCD corrections):

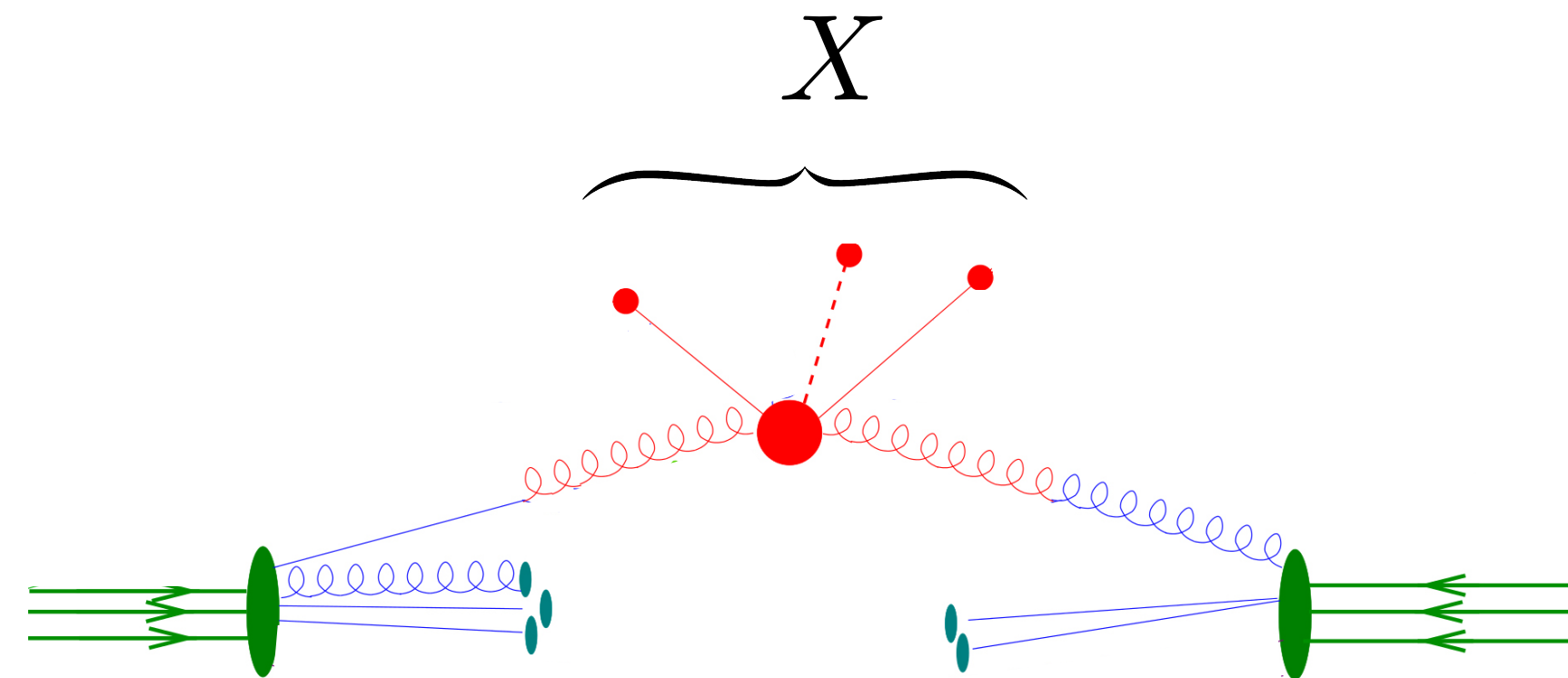
- NLO (now mostly automated)
- NNLO (state-of-the-art that has become standard)
- N3LO possible (with large impact on Higgs results)

Hadronic Collisions Phenomenology

Truth-level event

$$\sigma_W = \sum_q \int dx_1 dx_2 f_q(x_1) f_{\bar{q}}(x_2) \times \hat{\sigma}_{q\bar{q}}$$

total x-sec parton distribution functions parton x-sec



Parton Distribution Functions

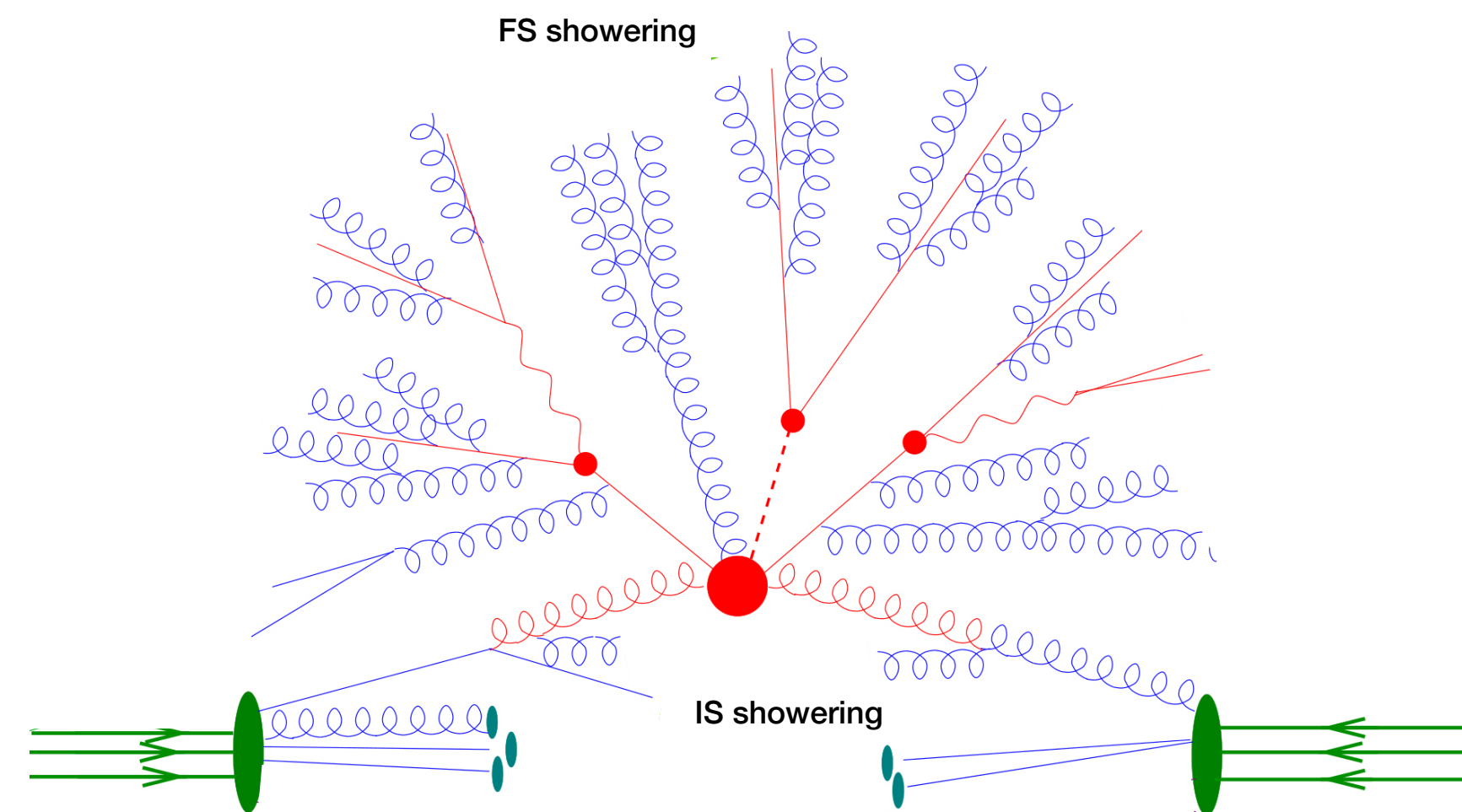
Hard scattering

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Hadronic Collisions Phenomenology

Truth-level event



Parton shower

QCD radiation showers not calculable exactly - Approximated in the soft and collinear regime — Pythia/Herwig MC

Parton Distribution Functions

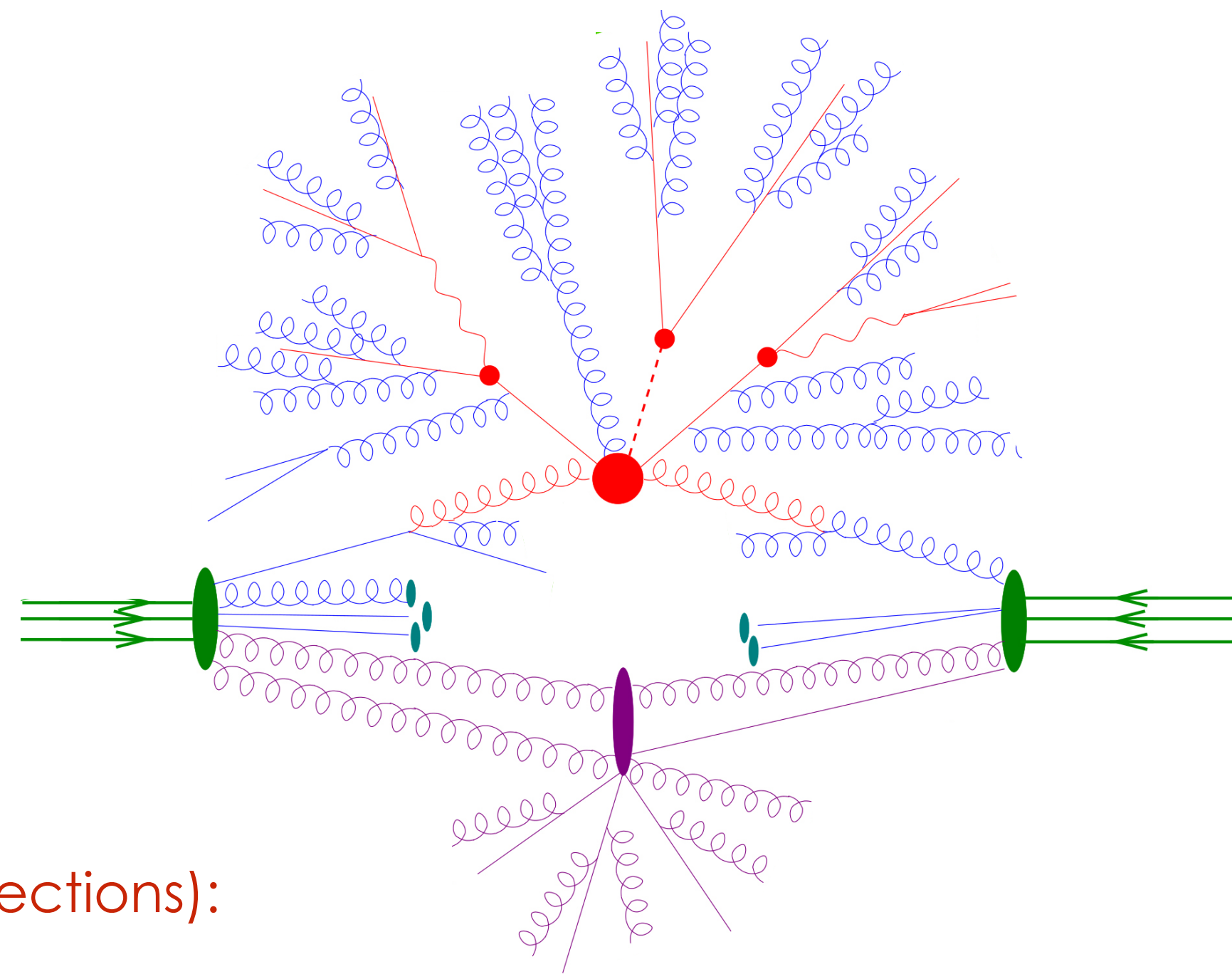
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Hadronic Collisions Phenomenology

Truth-level event



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Parton shower

QCD radiation showers not calculable exactly - Approximated in the soft and collinear regime — Pythia/Herwig MC

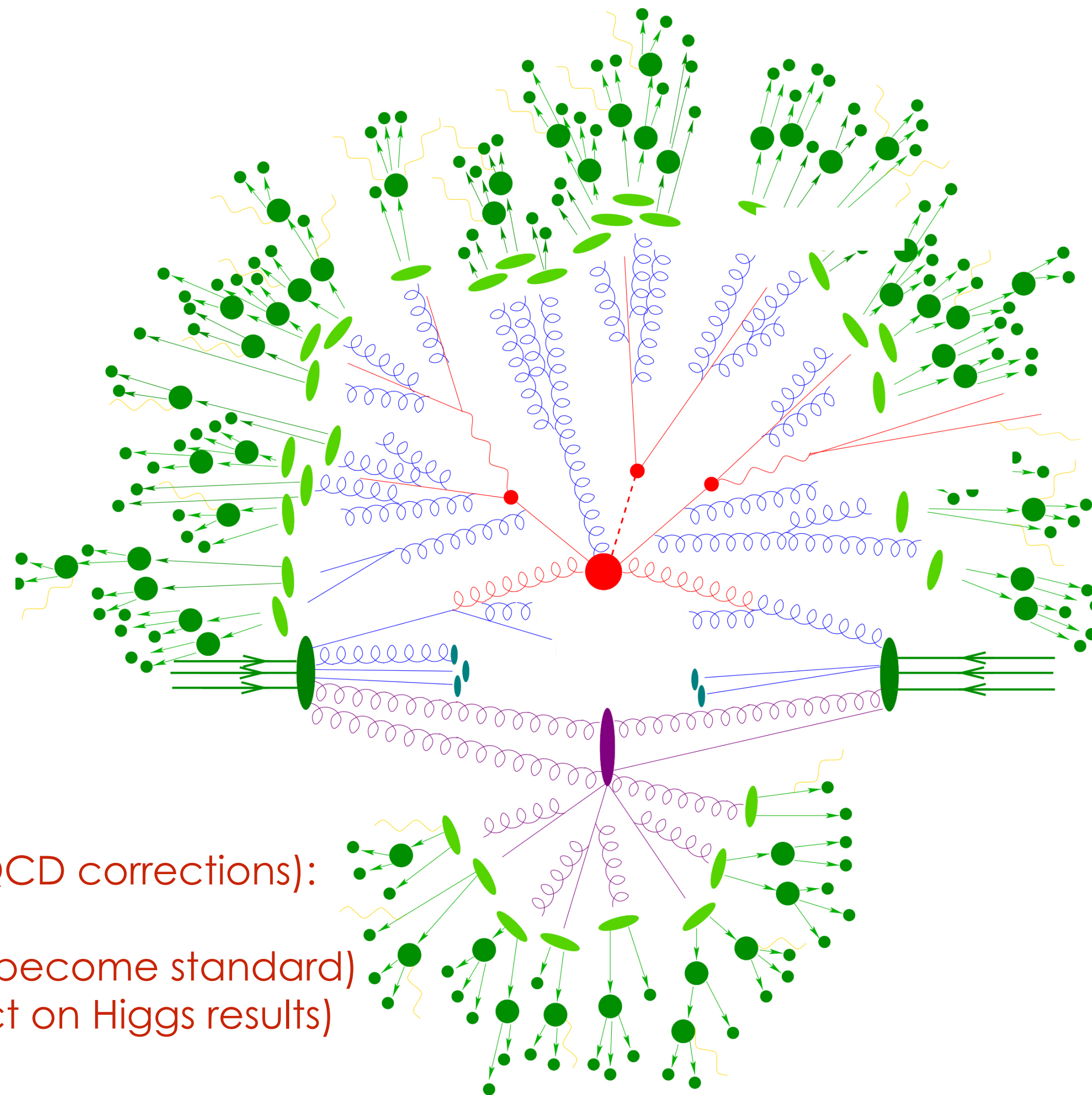
Parton Distribution Functions

Multiple parton interactions

Additional interactions between partons of the same protons

Hadronic Collisions Phenomenology

Truth-level event



Hadronization

Pythia, Herwig, Sherpa MC

Parton shower

QCD radiation showers not calculable exactly - Approximated in the soft and collinear regime — Pythia/Herwig MC

Parton Distribution Functions

Hard scattering

Perturbative regime (with large QCD corrections):

- NLO (now mostly automated)
- NNLO (state-of-the-art that has become standard)
- N3LO possible (with large impact on Higgs results)

Multiple parton interactions

Additional interactions between partons of the same protons

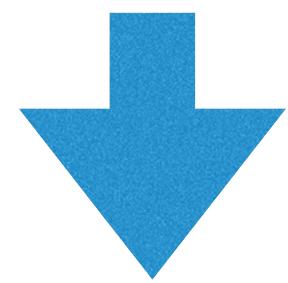
Picture by Frank Krauss

Hadronic Collisions Phenomenology

Truth-level event

Next

- Hadron decays
- Additional collisions from other protons (pile-up)

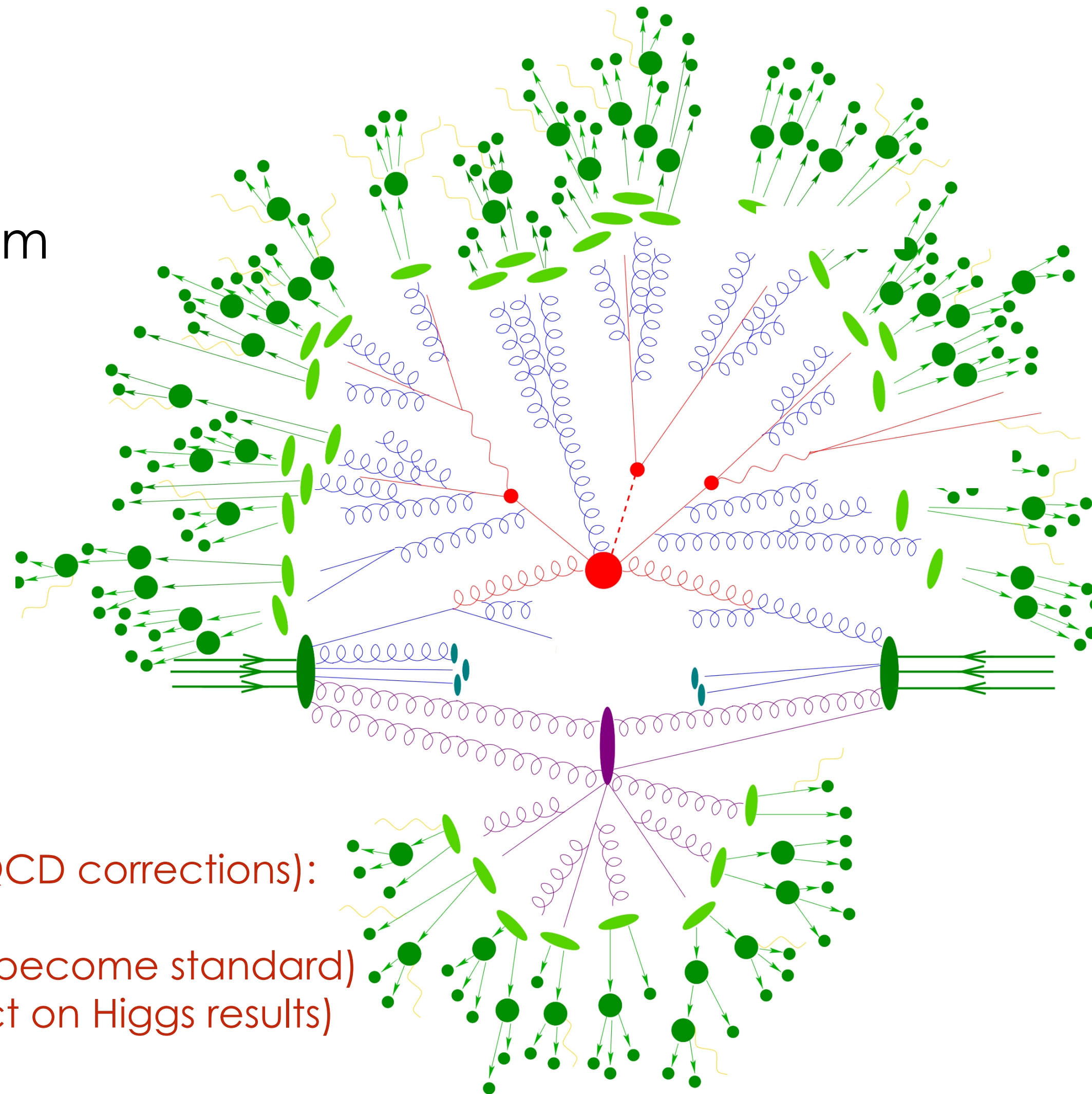


Reconstruction

Hard scattering

Perturbative regime (with large QCD corrections):

- NLO (now mostly automated)
- NNLO (state-of-the-art that has become standard)
- N3LO possible (with large impact on Higgs results)



Hadronization

Pythia, Herwig, Sherpa MC

Parton shower

QCD radiation showers not calculable exactly - Approximated in the soft and collinear regime — Pythia/Herwig MC

Parton Distribution Functions

Multiple parton interactions

Additional interactions between partons of the same protons

Picture by Frank Krauss

Luminosity

Number of collisions that can be produced in a detector per cm² and per second

Number of protons per bunch
 1.2×10^{11}

Number of bunches
1400

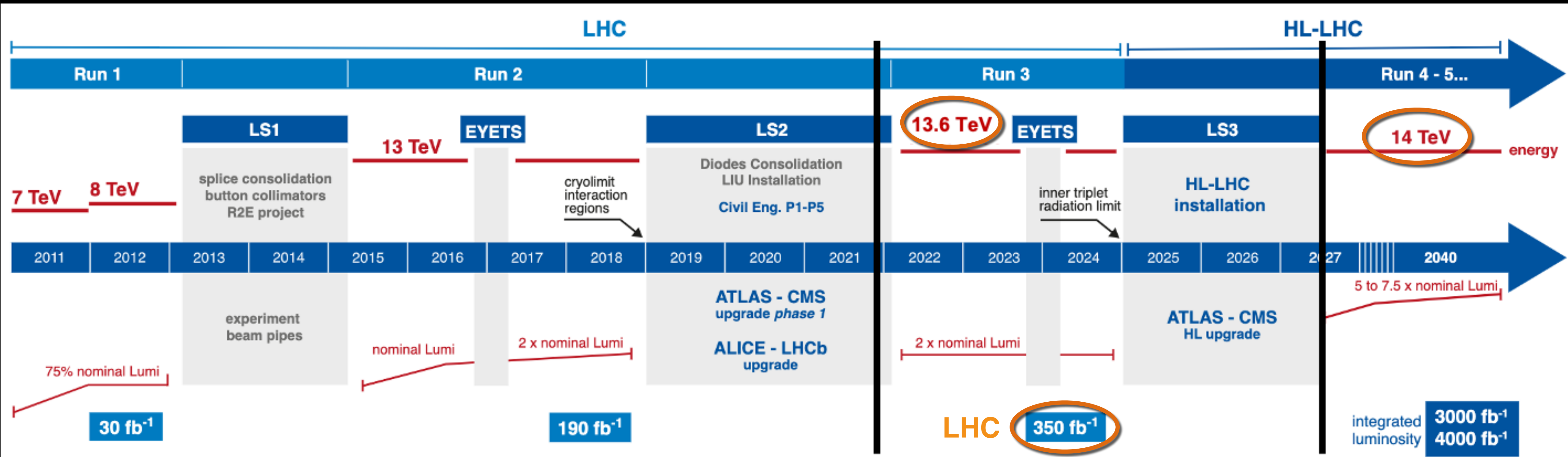
Number of turns per second
11245

$$\mathcal{L}_p = \frac{N^2 K_b f}{4\pi \sigma_x \sigma_y}$$

Beam size at interaction point (IP)
~20-25 μm radius

$$N_{ev} = \int \mathcal{L}_p dt \times \sigma(pp \rightarrow X)$$

High-Luminosity LHC Plan



We are here!
moving towards this target

LHCb and ALICE plan
major upgrades in LS4

CEPC startup?

HL-LHC goal is deliver 3000 fb⁻¹ in 10 years

- Implies integrated luminosity of 250-300 fb⁻¹ per year
- Requires peak luminosities of $5-7 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$ while using luminosity leveling (3-5 hours at peak luminosity)

Design for "ultimate" performance $7.5 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$ and 4000 fb⁻¹

Hadron collisions - Luminosity

$$\mathcal{L} = \frac{N_p^2 f}{4\pi\sigma_x\sigma_y}$$

$$\rightarrow \frac{N_p^2 f k_b}{4\pi\sigma_x\sigma_y}$$

$$\rightarrow \frac{N_p^2 f k_b \gamma}{4\pi\beta^* \epsilon_N}$$

$$\rightarrow \frac{N_p^2 f k_b \gamma}{4\pi\beta^* \epsilon_N} F$$

Luminosity for several bunches kb in trains

Beam size expressed in beam parameters, dependence in beam energy!

Correction due to geometrical effects (e.g. beam crossing angle)

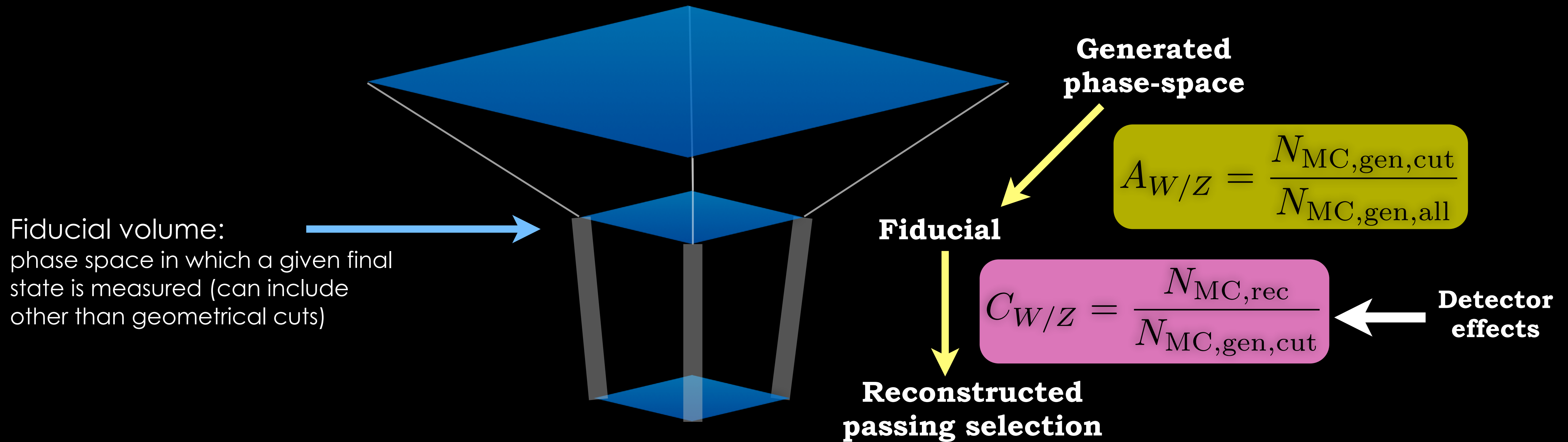
Luminosity grows with energy!

Emittance (fixed at injection)

Beta* Focussing

Parameter	2010	2011	2012	2016	2017	2018	Nominal	HL-LHC
CoM Energy	7 TeV	7 TeV	8 TeV	13 TeV	13 TeV	13 TeV	14 TeV	14 TeV
N_p	$1.1 \cdot 10^{11}$	$1.4 \cdot 10^{11}$	$1.6 \cdot 10^{11}$	$1.2 \cdot 10^{11}$	$1.2 \cdot 10^{11}$	$1.2 \cdot 10^{11}$	$1.15 \cdot 10^{11}$	$2.2 \cdot 10^{11}$
Bunches k	368	1380	1380	2300	2450	2500	2808	2760
Spacing	150 ns	50 ns	50 ns	25 ns	25 ns	25 ns	25 ns	25ns
ϵ (mm rad)	2.4-4	1.9-2.3	2.5	2.6	2.3	2.6	3.75	2.5
β^* (m)	3.5	1.5-1	0.6	0.4	0.3-0.4	0.4	0.55	0.15
L ($\text{cm}^{-2}\text{s}^{-1}$)	$2 \cdot 10^{32}$	$3.3 \cdot 10^{33}$	$\sim 7 \cdot 10^{33}$	$1.5 \cdot 10^{33}$	$2.0 \cdot 10^{34}$	$2 \cdot 10^{34}$	10^{34}	$8 \cdot 10^{34}$
PU	~2	~10	~30	~30	~50	~50	~25	~130

Fiducial volume and acceptance



Fiducial cross section

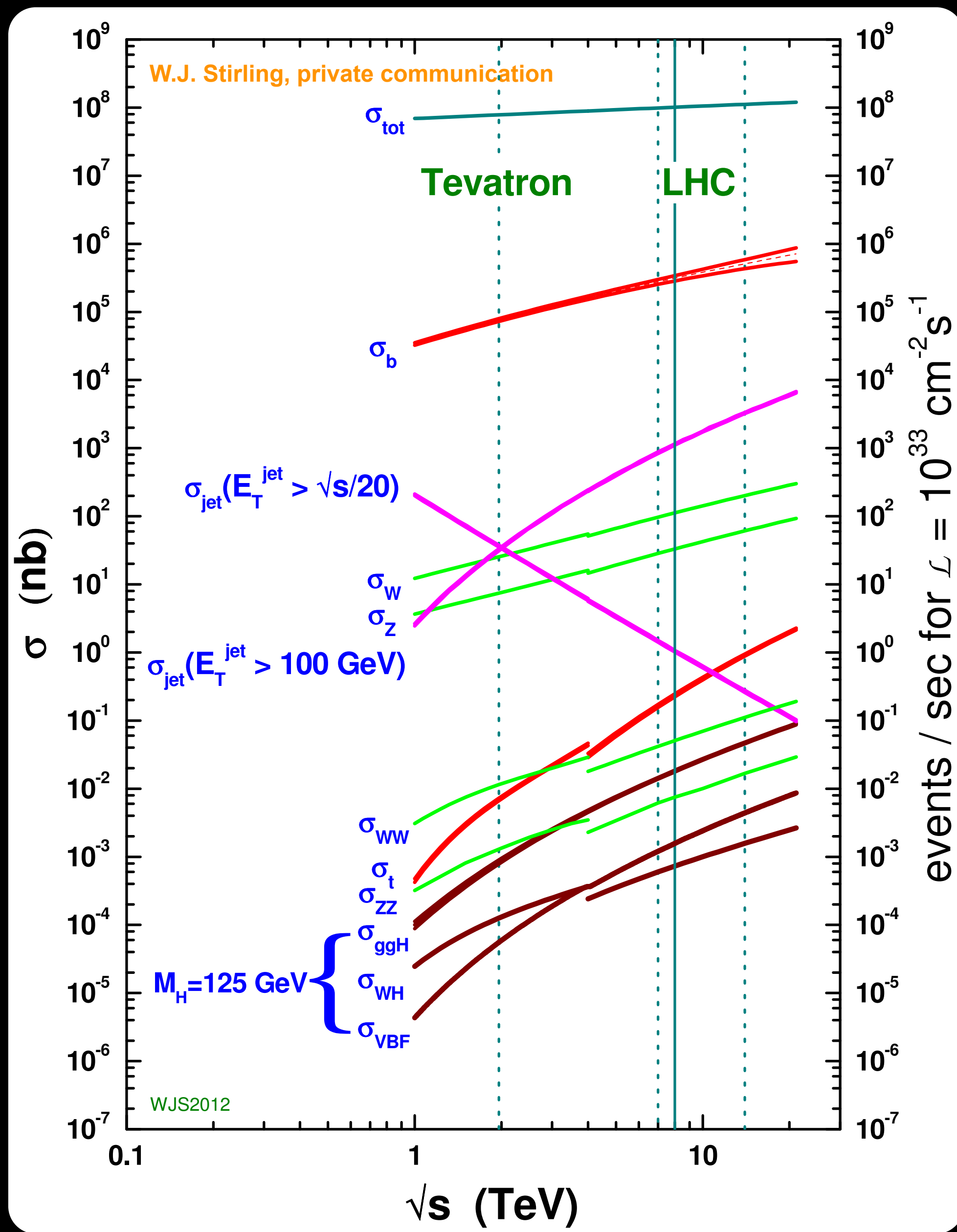
$$\sigma_{fid} = \frac{N - B}{C_{W/Z} \cdot L_{int}}$$

No theoretical uncertainty from extrapolation outside experimental acceptance

Total cross section

$$\sigma_{tot} = \sigma_{W/Z} \times BR(W/Z \rightarrow l\nu/\ell\ell) = \frac{\sigma_{fid}}{A_{W/Z}}$$

The total cross section at hadron collider



Total cross section: 100 mb

Nominal LHC Luminosity

$$2 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$$

Number of events produced:

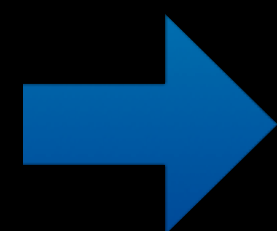
$$\sigma \times \mathcal{L} =$$

$$100 \times 10^{-27} (\text{cm}^2) \times 2 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$$

$$\sim 2 \times 10^9 \text{ events/sec}$$

Production rates at Hadron Colliders

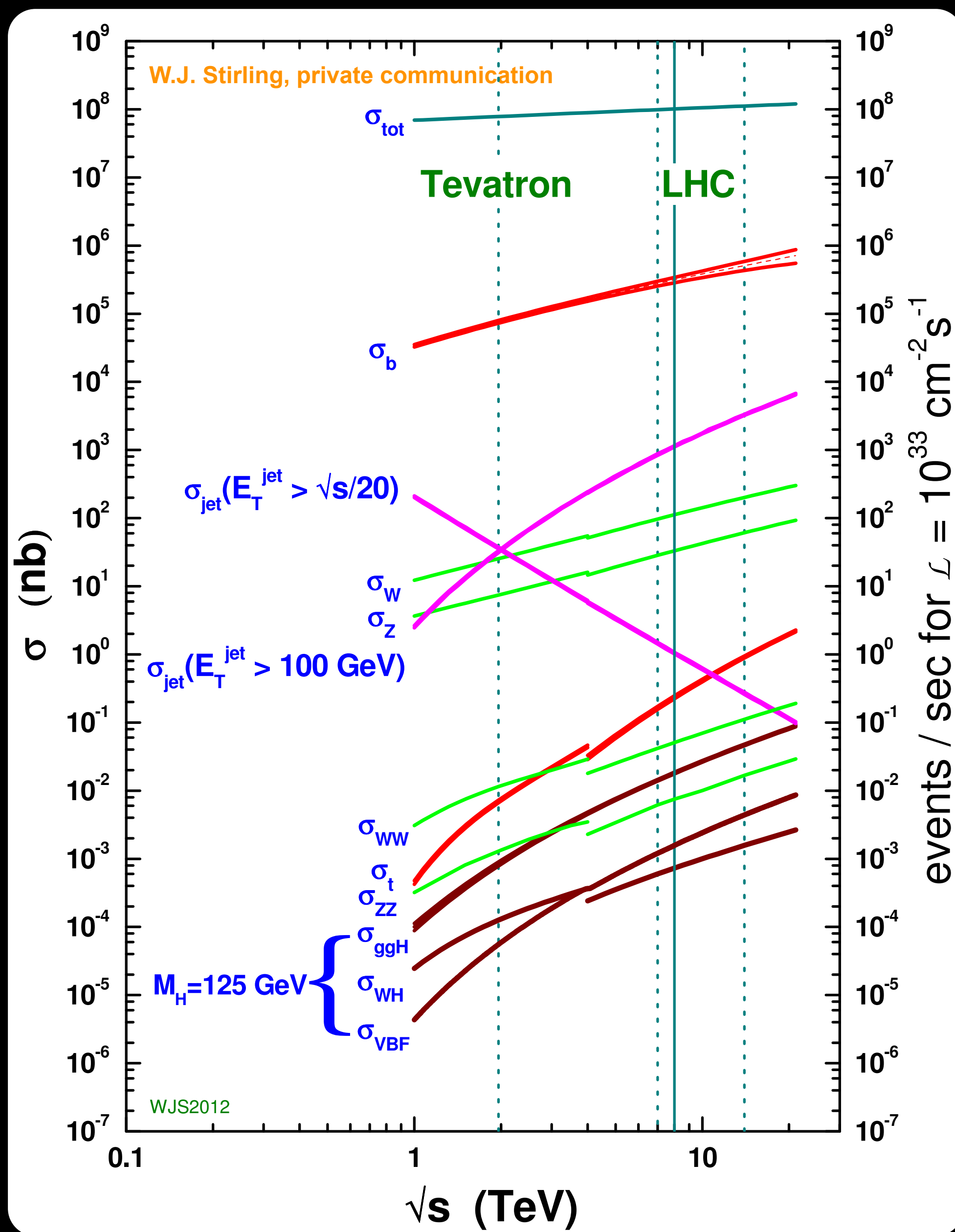
Tevatron timeline



(*) **W: 1988**
Z: 1988

Top: 1994
WW: 2005
WZ: 2007
ZZ: 2008

(*) Discovered at CERN in 1983 by UA1 and UA2



LHC timeline

W: May 2010
Z: Jun 2010

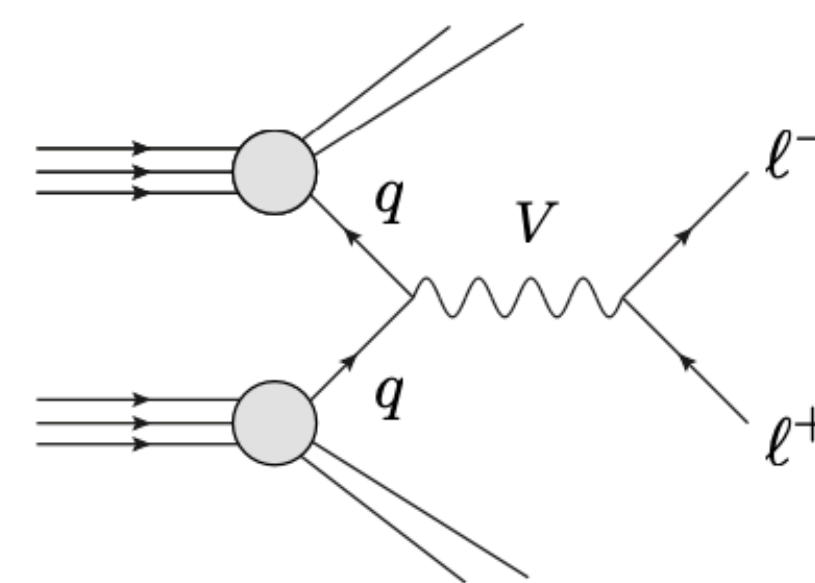
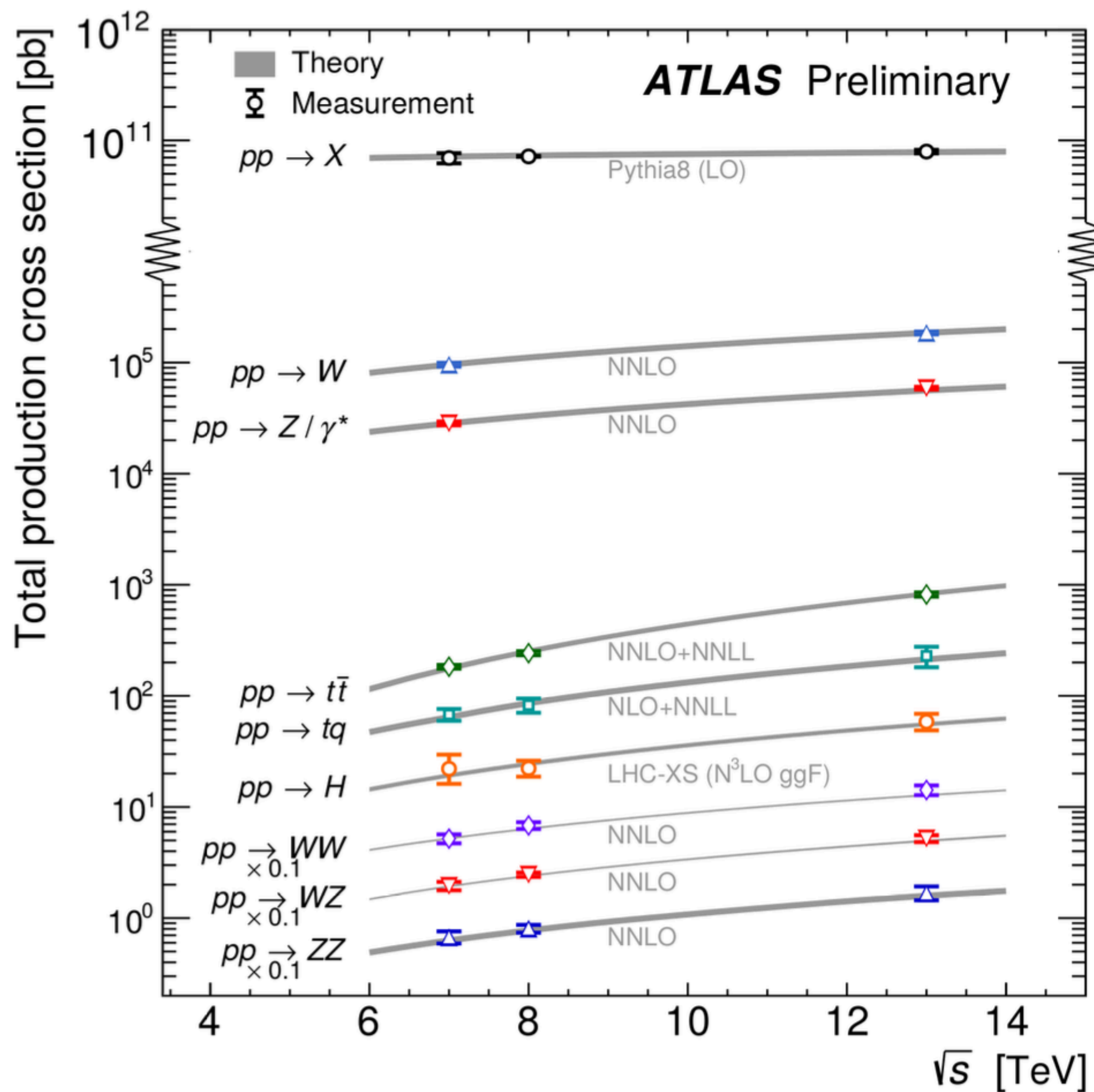
Top: Jul 2010

WW: Dec 2010
WZ: Mar 2011
ZZ: Jul 2011

H: July 4, 2012

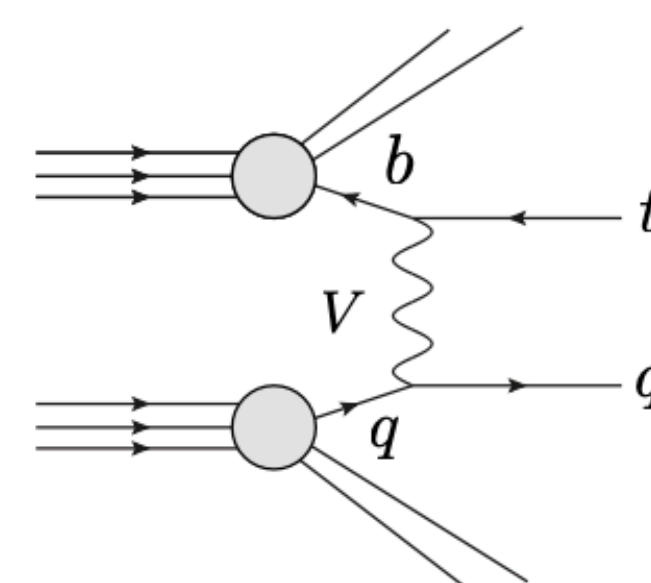


Total cross sections at LHC

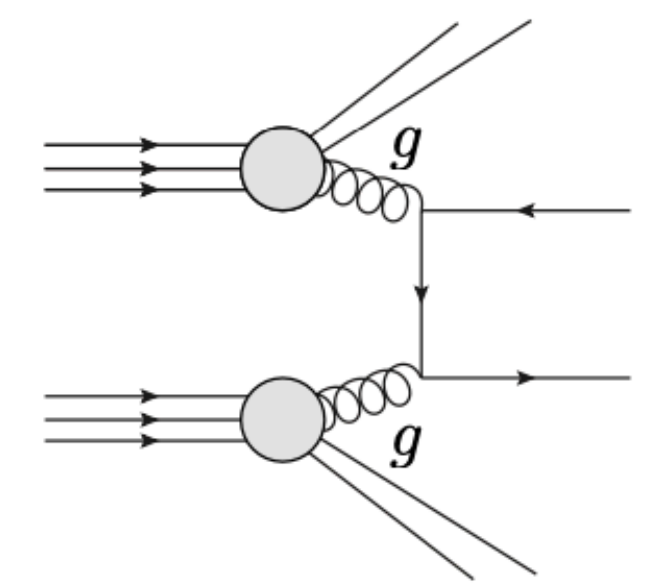


Vector boson production (often referred to as Drell Yan).

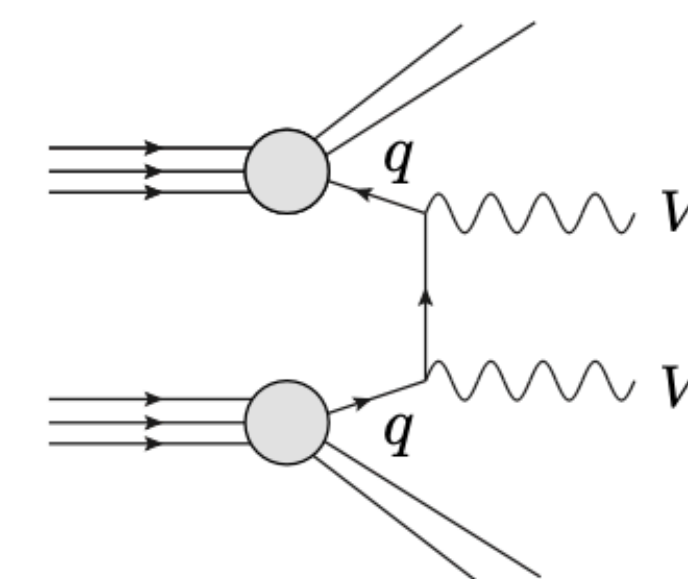
LEP ~4 M Z per experiment
LHC ~100 M (leptonic) / exp.
(for 100 fb^{-1})



Single top production
 $tq \sim 200 \text{ pb}$



Top pair production
 $t\bar{t} \sim 1 \text{ nb}$

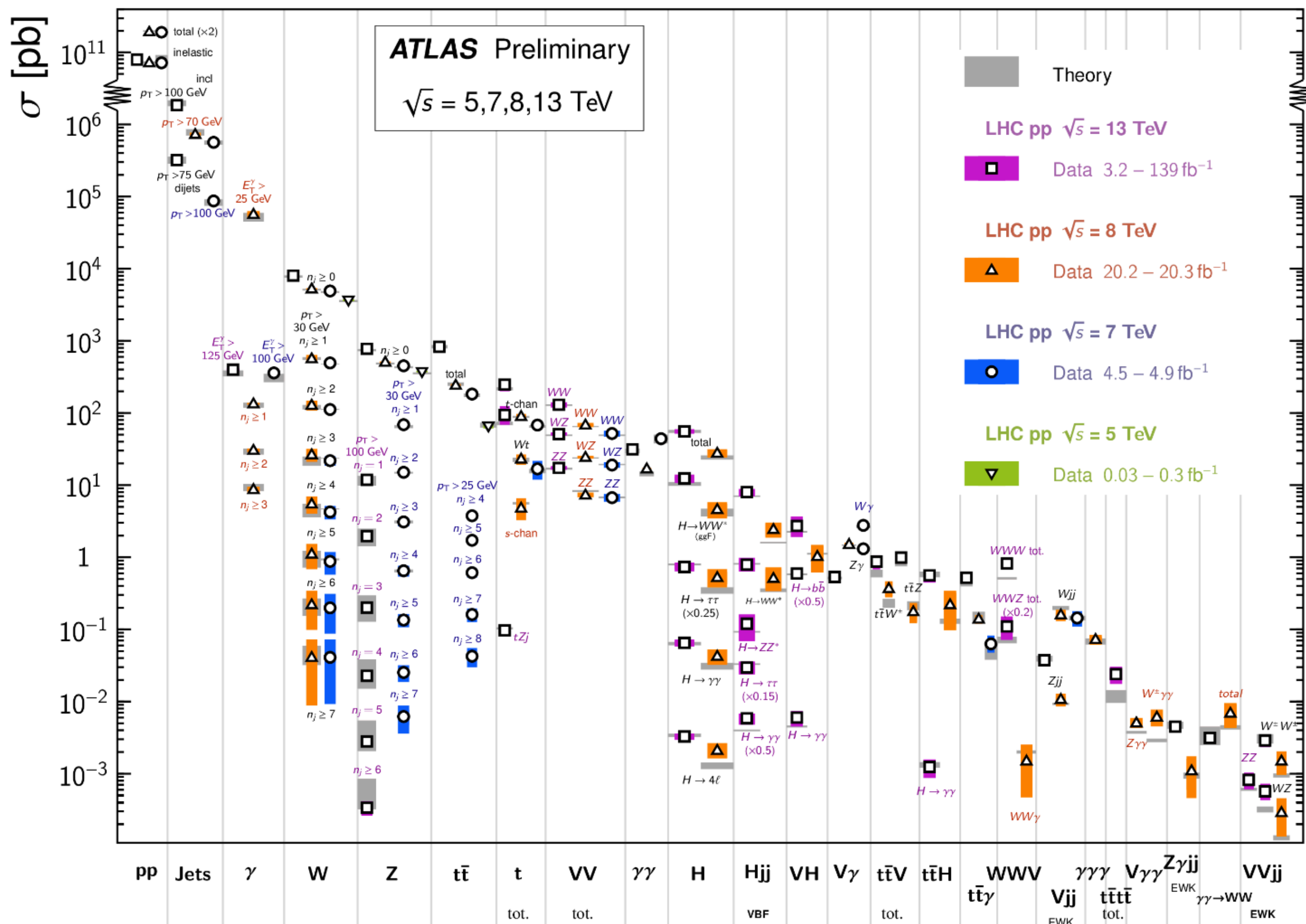


Diboson production
 $WW \sim 100 \text{ pb}$
 $ZZ \sim 20 \text{ pb}$

Fiducial cross sections at LHC

Standard Model Production Cross Section Measurements

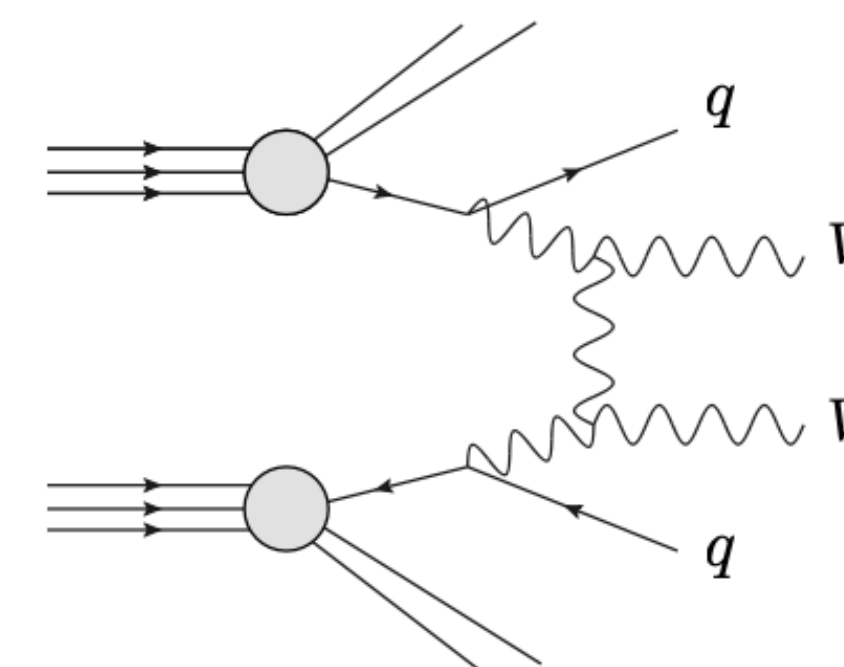
Status: February 2022



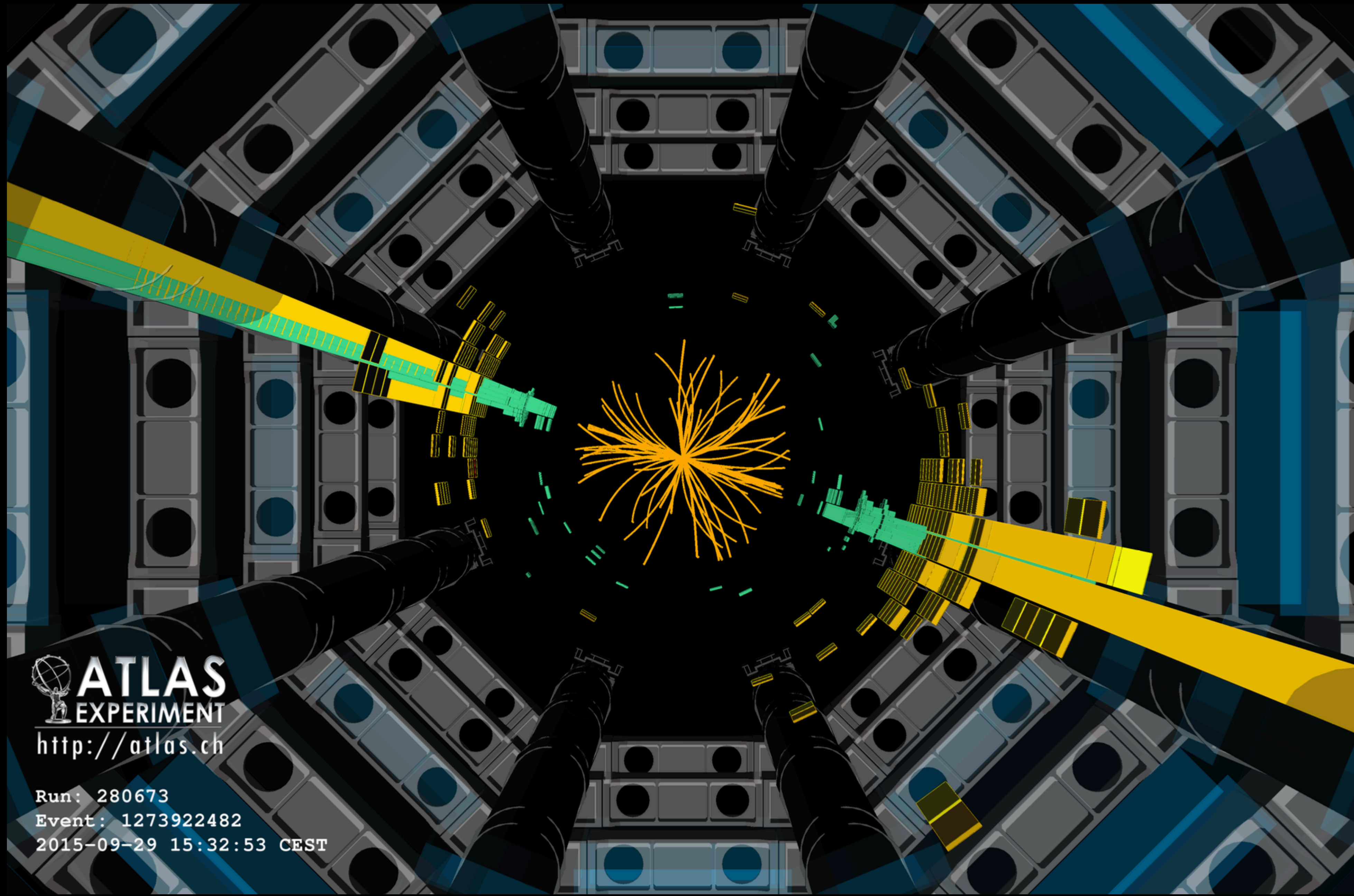
Very large number of fiducial cross section measurement made at the LHC

Down to processes as rare as three boson production

Now measured VBS diboson production (VVjj)



Jet cross sections

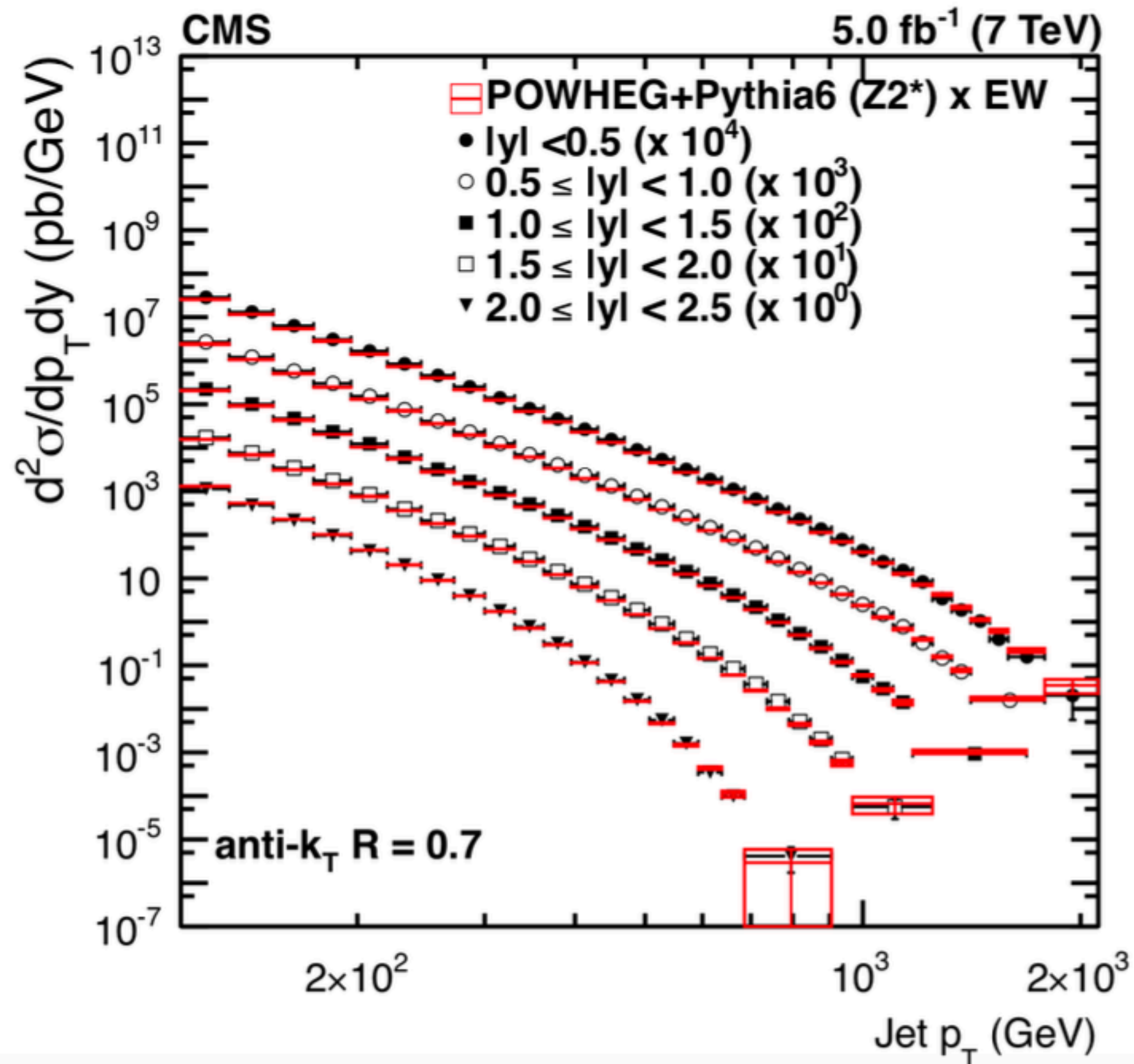


 **ATLAS**
EXPERIMENT
<http://atlas.ch>

Run: 280673
Event: 1273922482
2015-09-29 15:32:53 CEST

Differential jet production cross sections

Example: Double differential jet cross section measurement

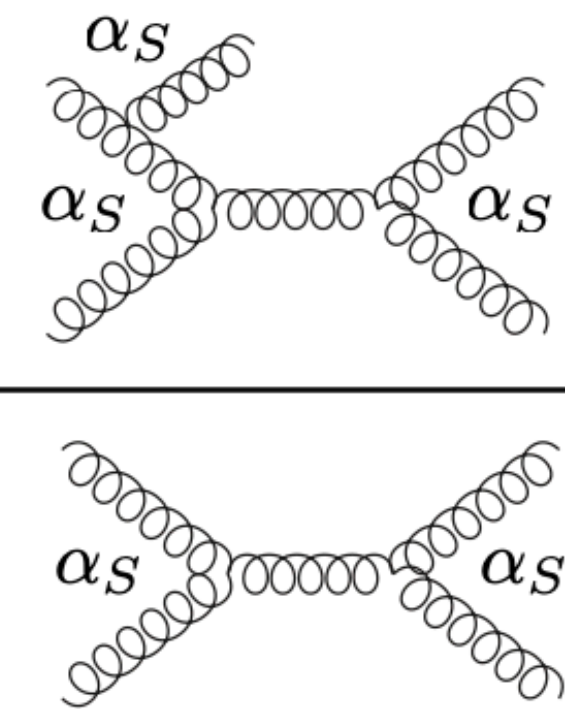


Count number of events in reconstructed bins of $(\Delta p_T, \Delta y)$

Unfold to truth particle jet quantities taking into account reconstruction and trigger efficiencies and unfolding resolution matrix

$$\frac{d^2\sigma}{dp_T dy} = \frac{1}{\epsilon \mathcal{L}} \frac{N_j}{\Delta p_T \Delta y}$$

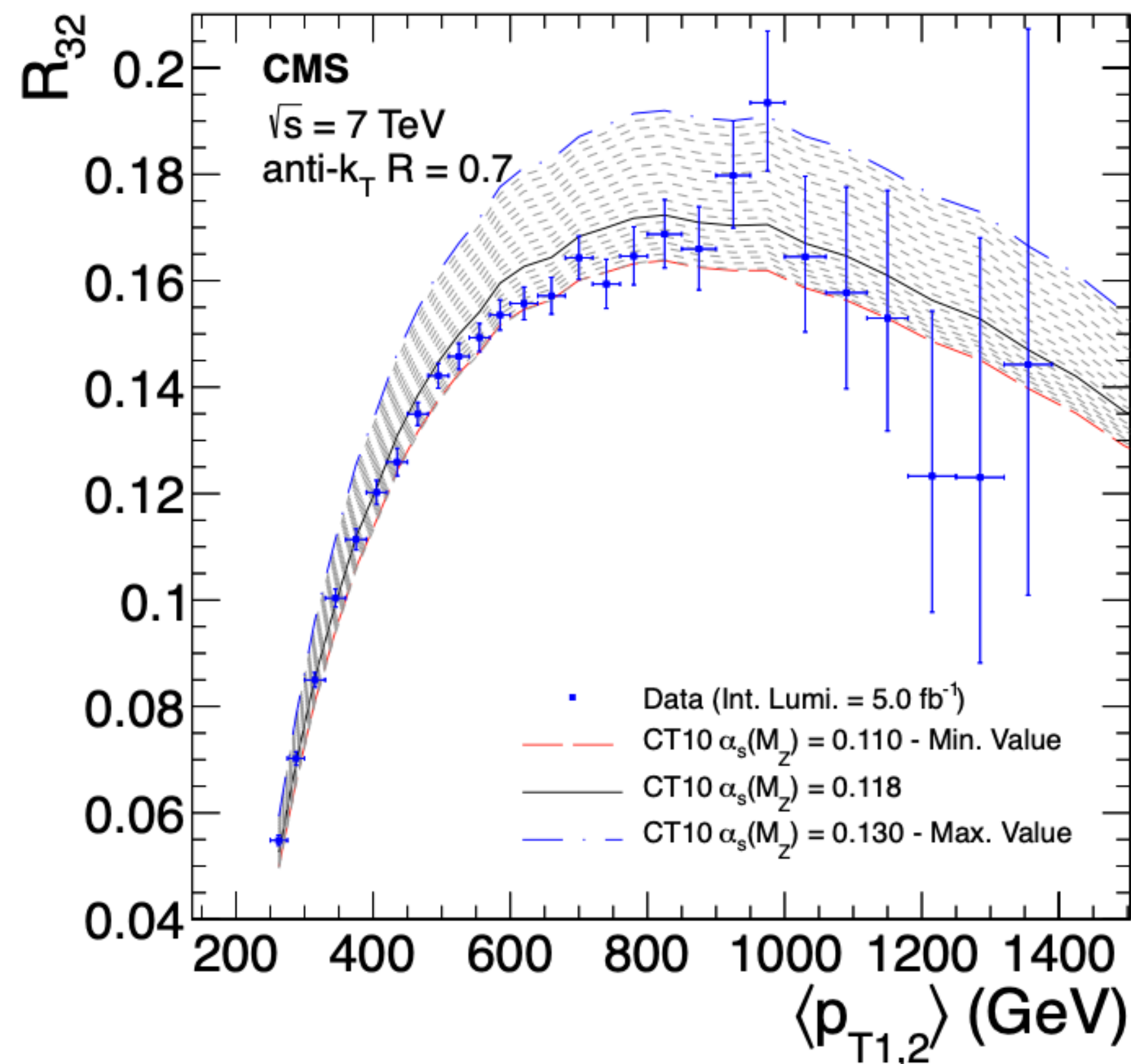
Ratios of differential jet production cross sections

$$R_{3/2} = \frac{\sigma_{3-jets}}{\sigma_{2-jets}} = \frac{\text{[3-jet diagram]} \propto \alpha_S^3}{\text{[2-jet diagram]} \propto \alpha_S^2} \propto \alpha_S$$


Measurement of the ratio of cross sections

- **R3/2** is the ratio of inclusive 3-jet to inclusive 2-jet cross sections as a function of the average pT of the two leading jets:

$$\langle p_{T1,2} \rangle$$

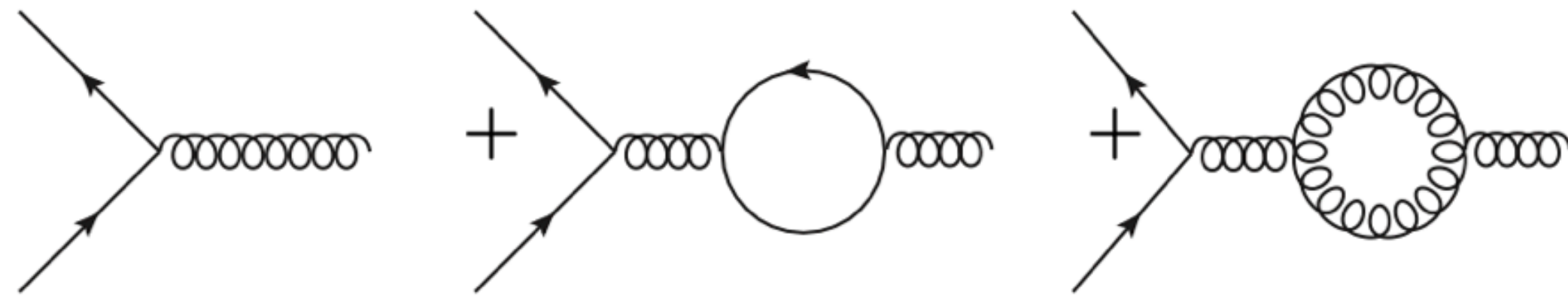


Using ratios can improve precision by cancelling systematic uncertainties.

- Again interesting dependence in strong coupling constant.
- Main experimental uncertainty partly canceled but still dominant.
- With the transverse momenta of all the jets can infer the energy scale of the process.

Measurement of the Strong Coupling Constant α_s

Crucial dependence of the strong coupling constant with the energy scale predicted (1973):



$$\alpha_S^{-1}(Q^2) = \alpha_S^{-1}(Q_0^2) \left[1 + \frac{33 - 2n_f}{12\pi} \alpha_S(Q_0^2) \ln \left(\frac{Q^2}{Q_0^2} \right) \right]$$

Asymptotic freedom: QCD is perturbative at high energies

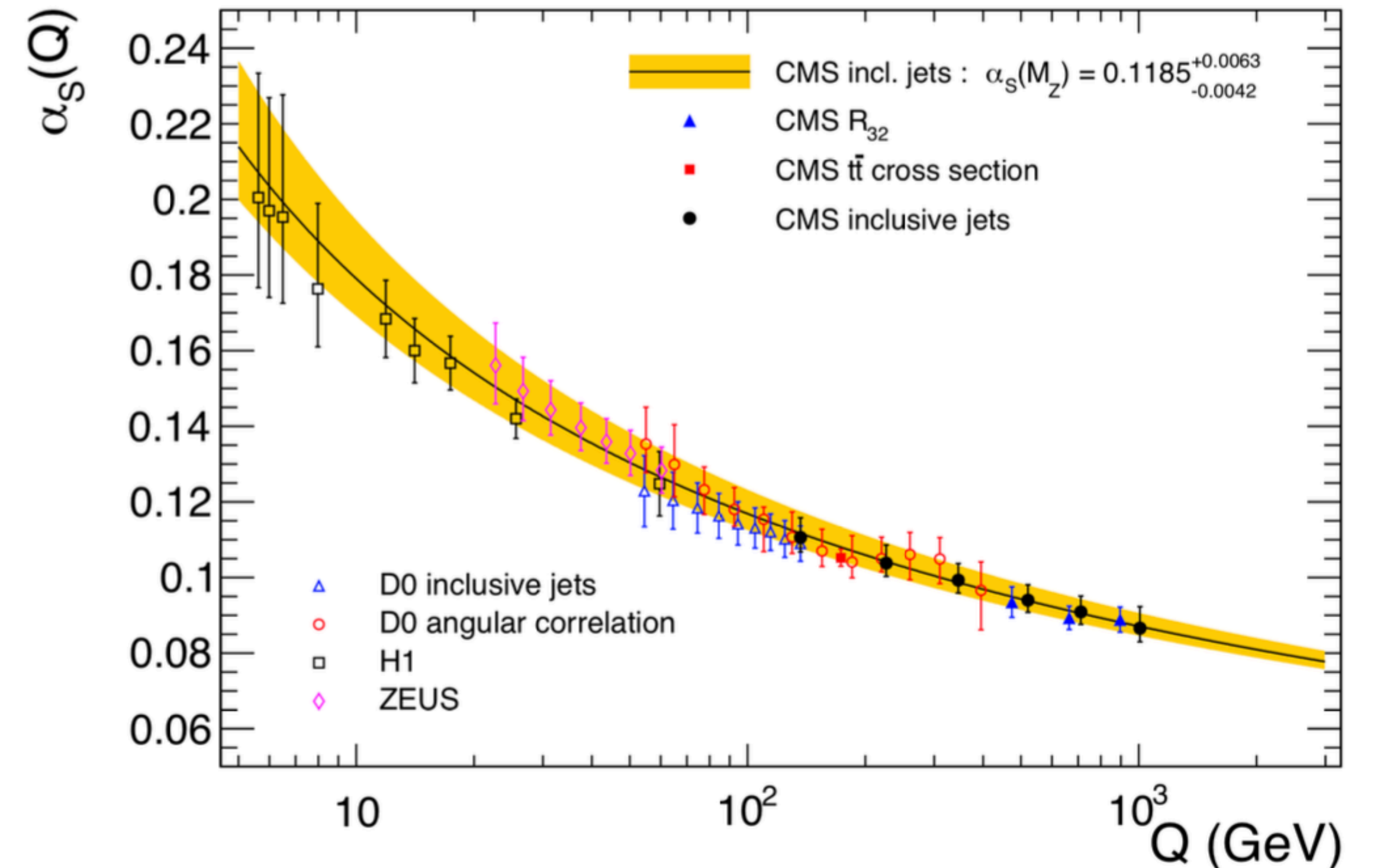


2004

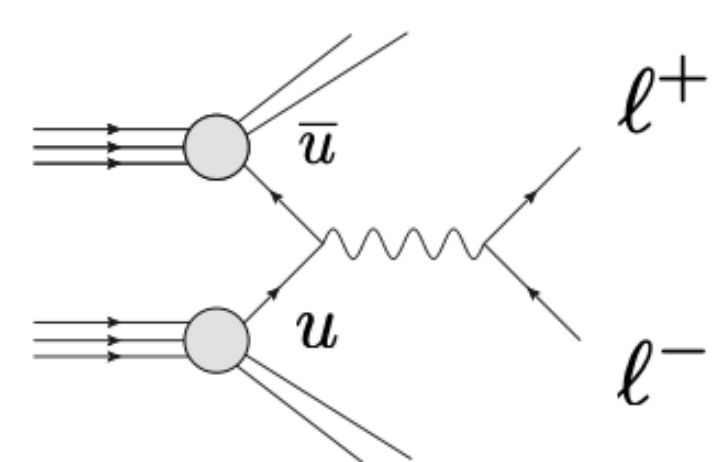
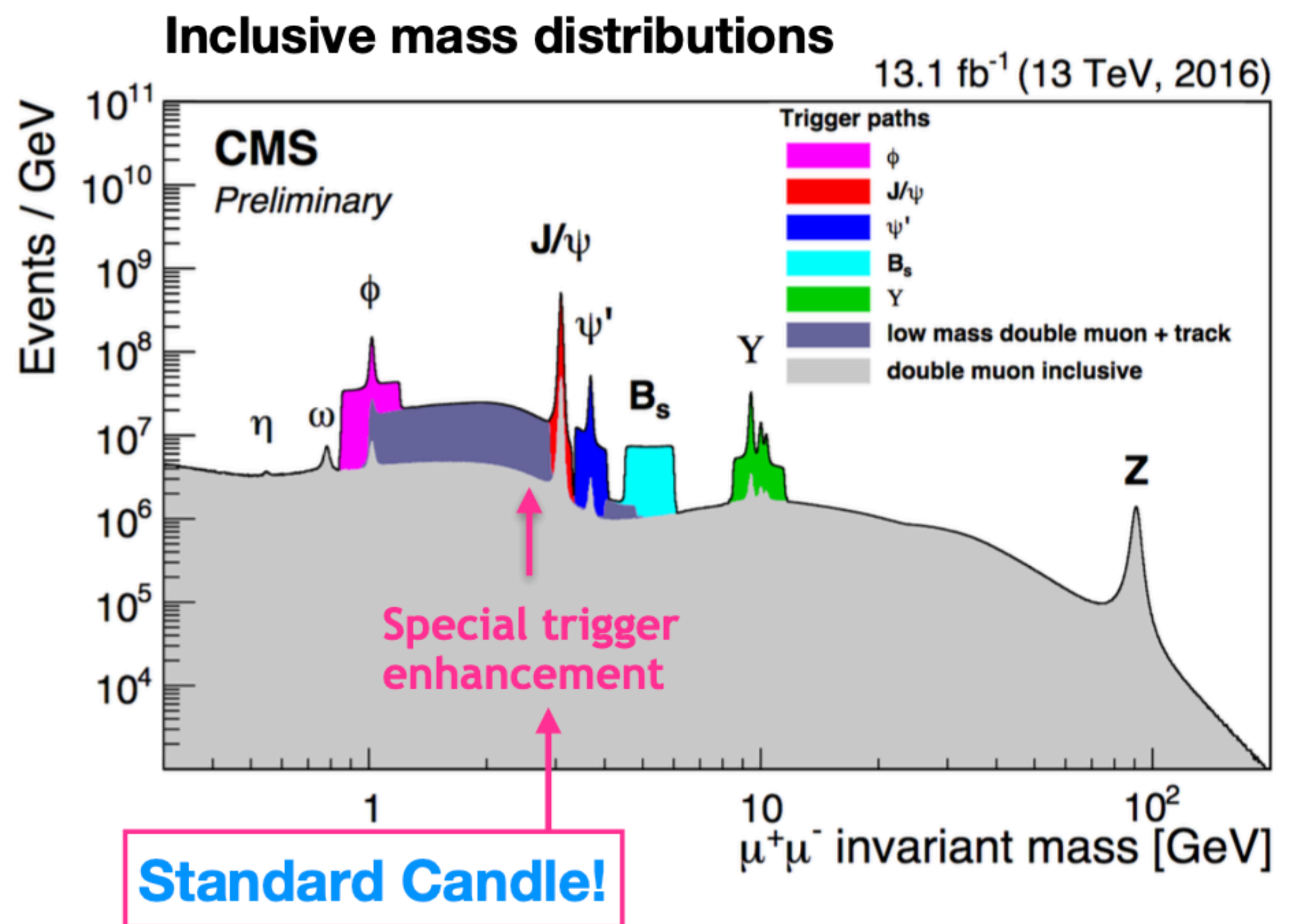
Asymptotic freedom and QCD

David Gross, David Politzer, Frank Wilczek

From the measurements of jet cross sections and their ratios, the strong coupling constant can be measured at the highest energy scales!

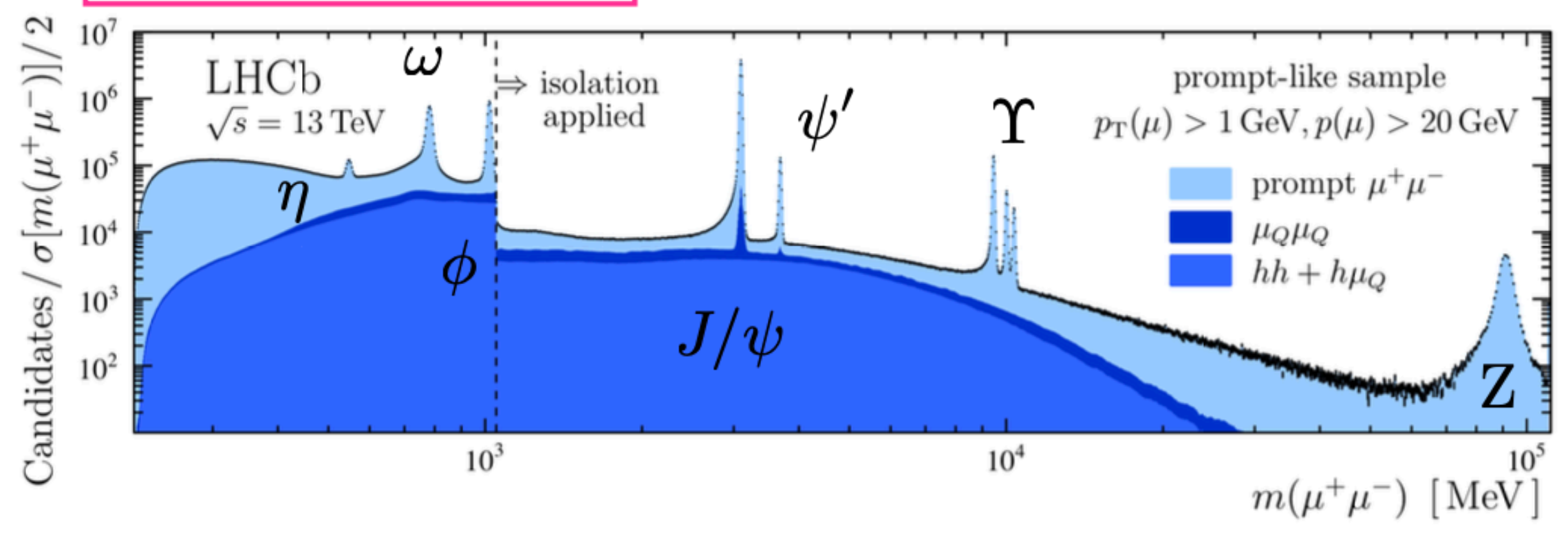
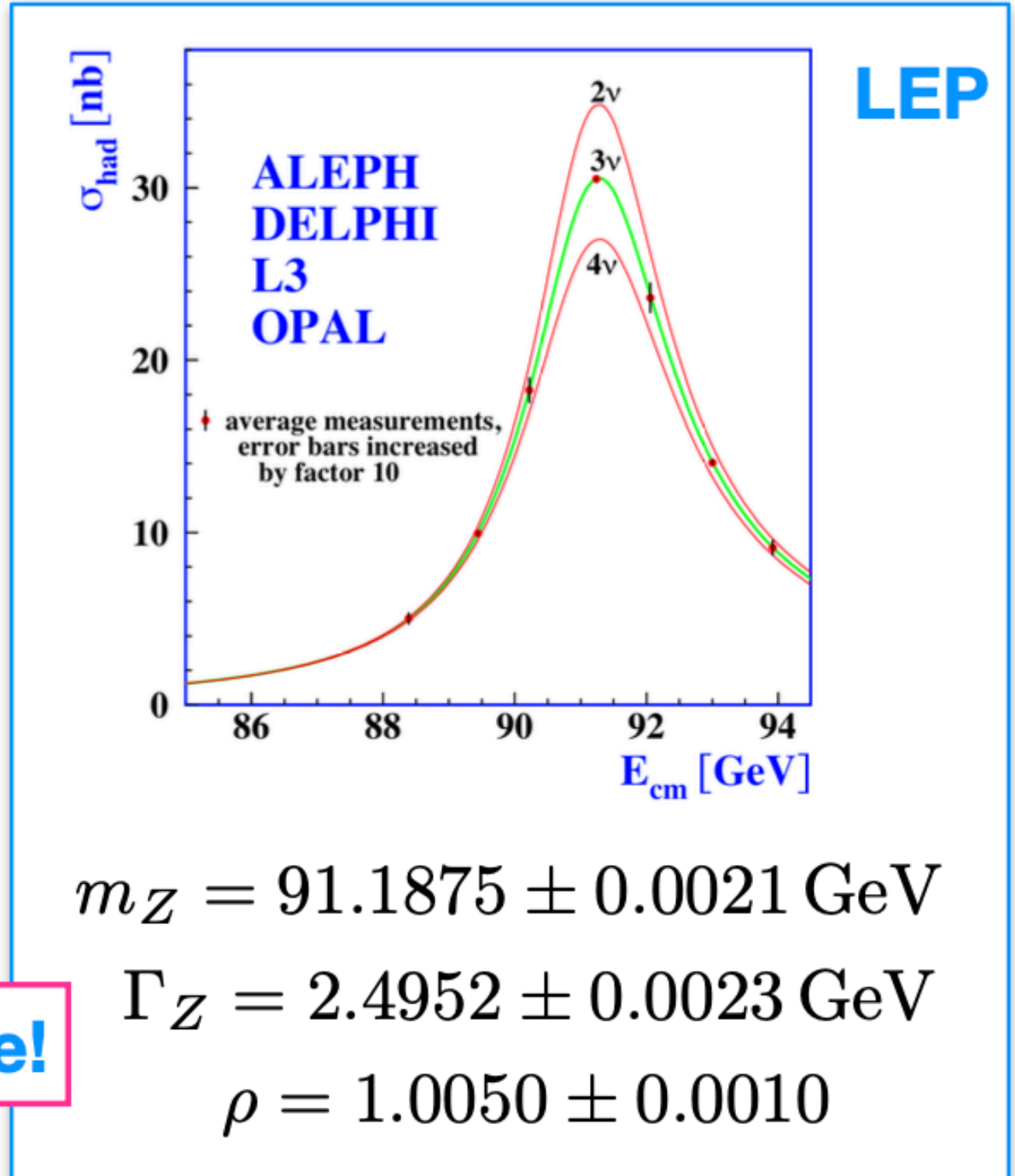


The di-lepton mass spectrum at LHC



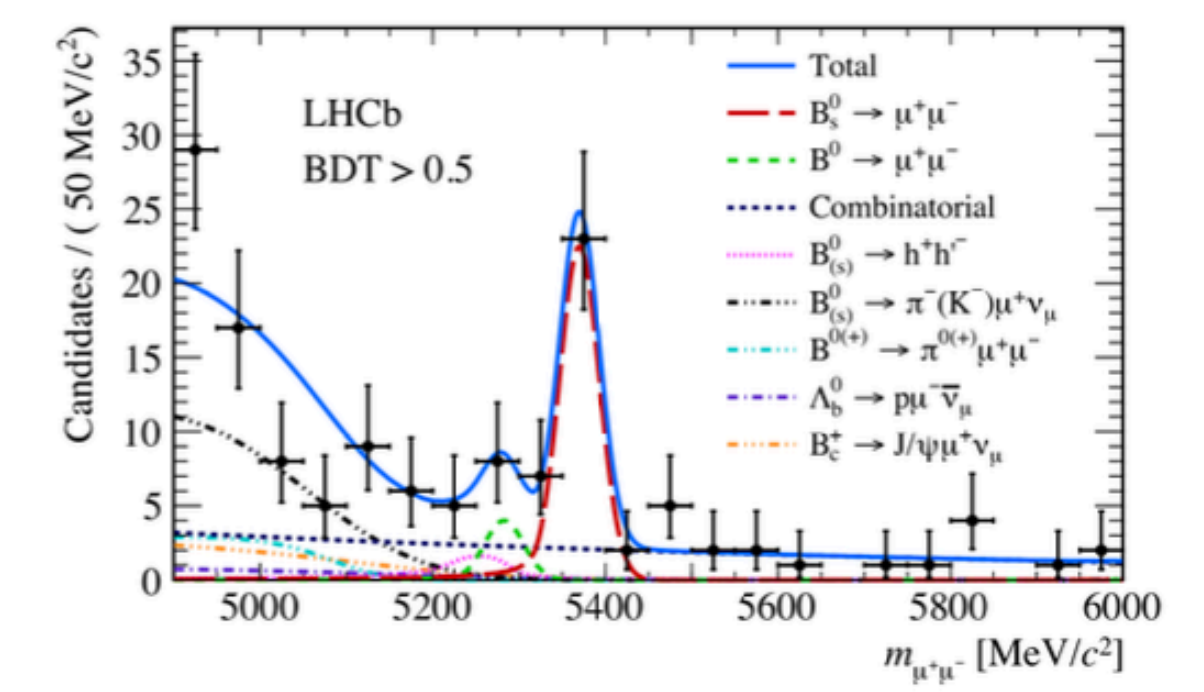
Z, J/Psi and Upsilon in electrons and muons are extremely important standard candles for calibration.

Standard Candle!

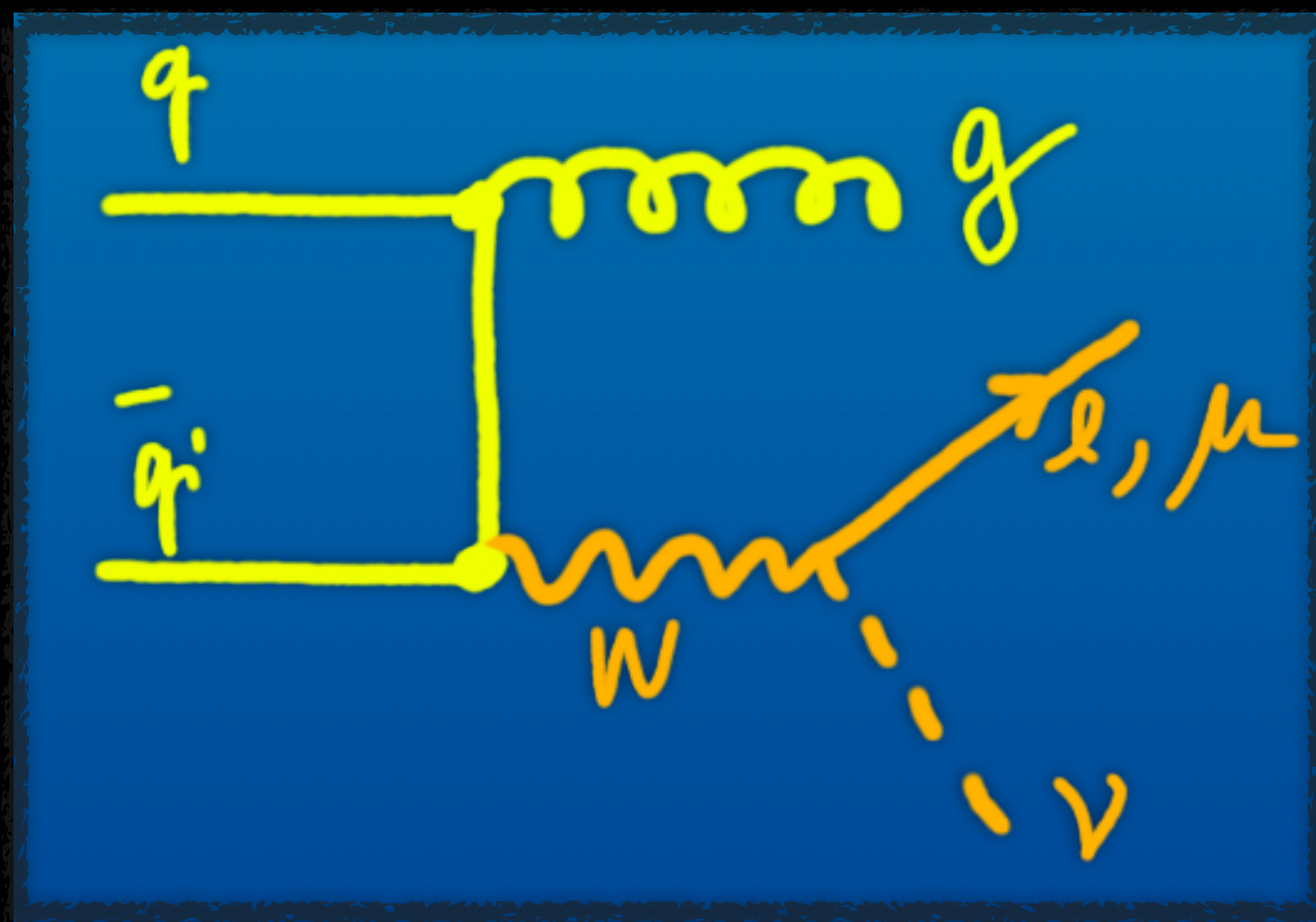


An exclusive analysis scrutinising the Bs mass region

$$\text{Br}(B_s^0 \rightarrow \mu^+\mu^-) = (3.65 \pm 0.23) \times 10^{-9}$$



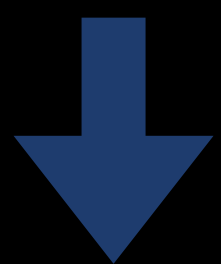
The challenge of the W Mass Measurement



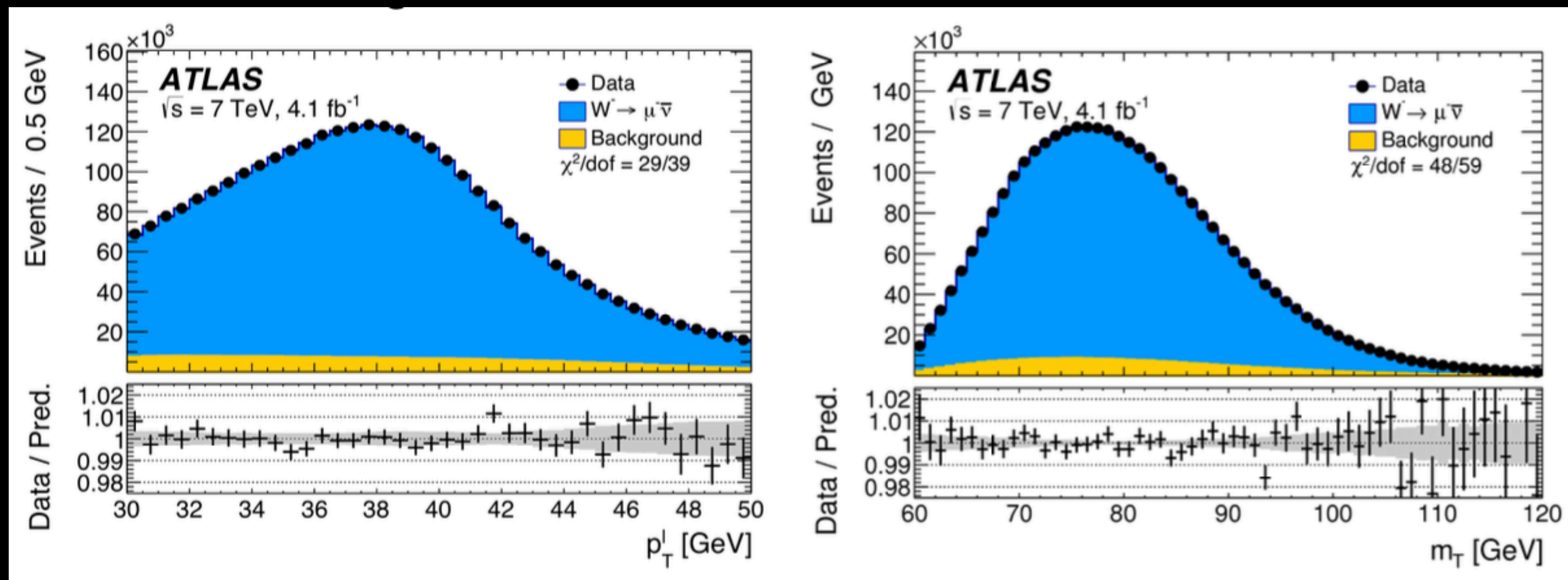
$$m_T = \sqrt{2 p_T^l p_T^{miss} (1 - \cos\Delta\varphi)}$$

Strategy: fit two kinematic distributions in several categories

Cannot measure **W mass** directly from its decay particles because the **neutrino** is not measured directly in the detector



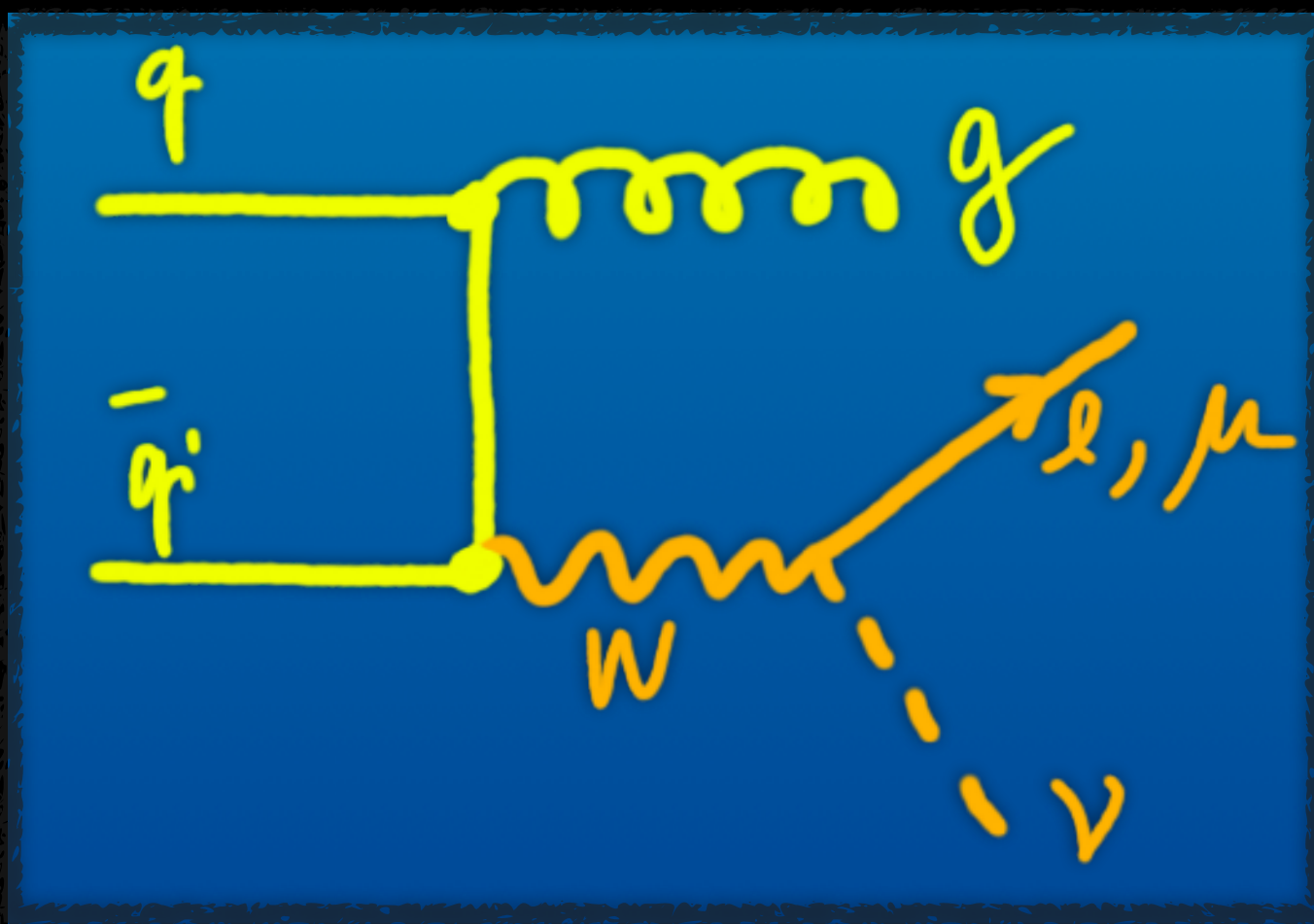
Measure **W recoil** precisely



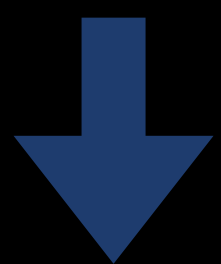
p_T^l : Clean energy measurement, but more sensitive to the modelling of the W transverse momentum

m_T : Less sensitive to modelling but more difficult to reconstruct because of the missing transverse energy

The challenge of the W Mass Measurement



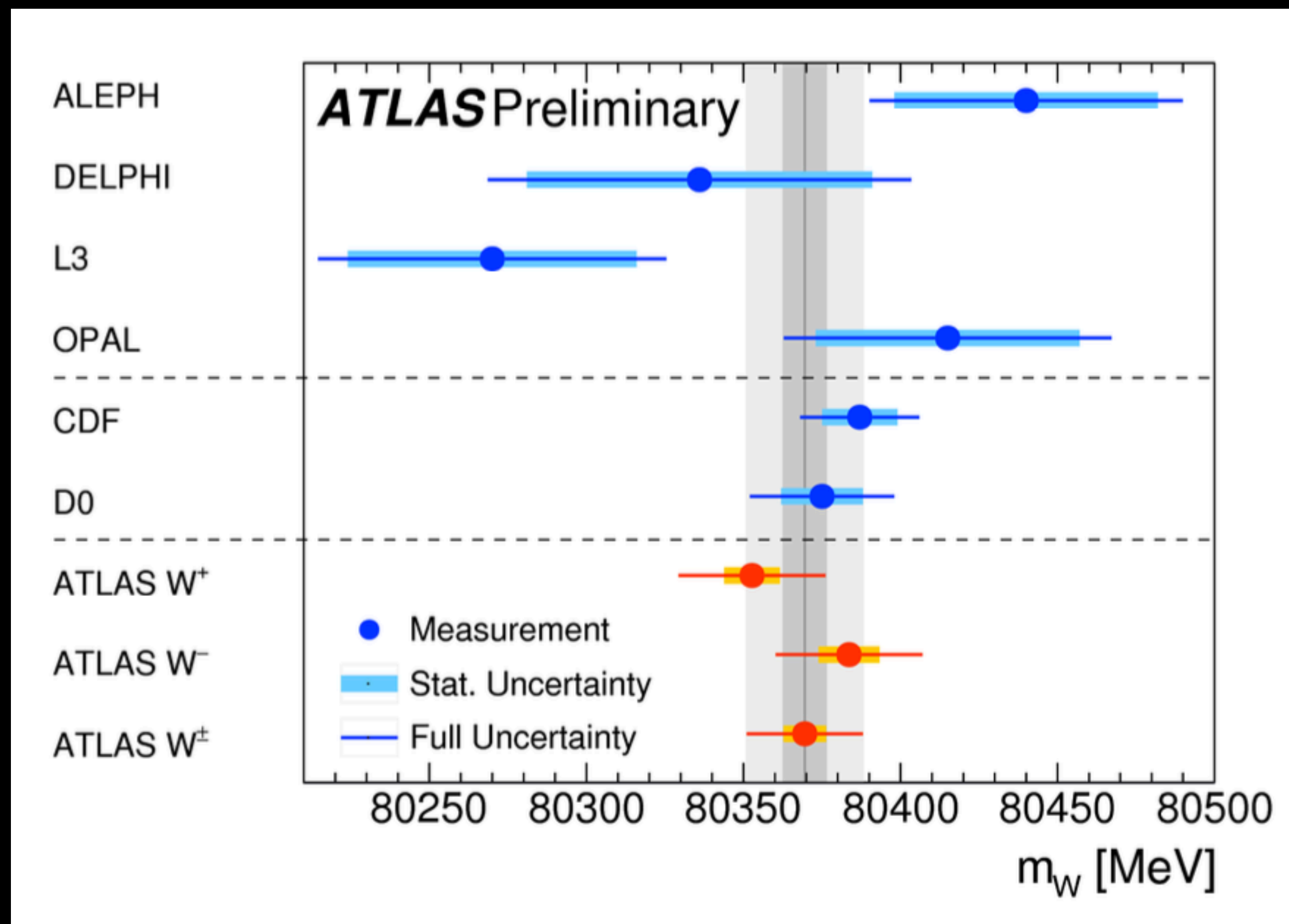
Cannot measure **W mass** directly from its decay particles because the **neutrino** is not measured directly in the detector



Measure **W recoil** precisely

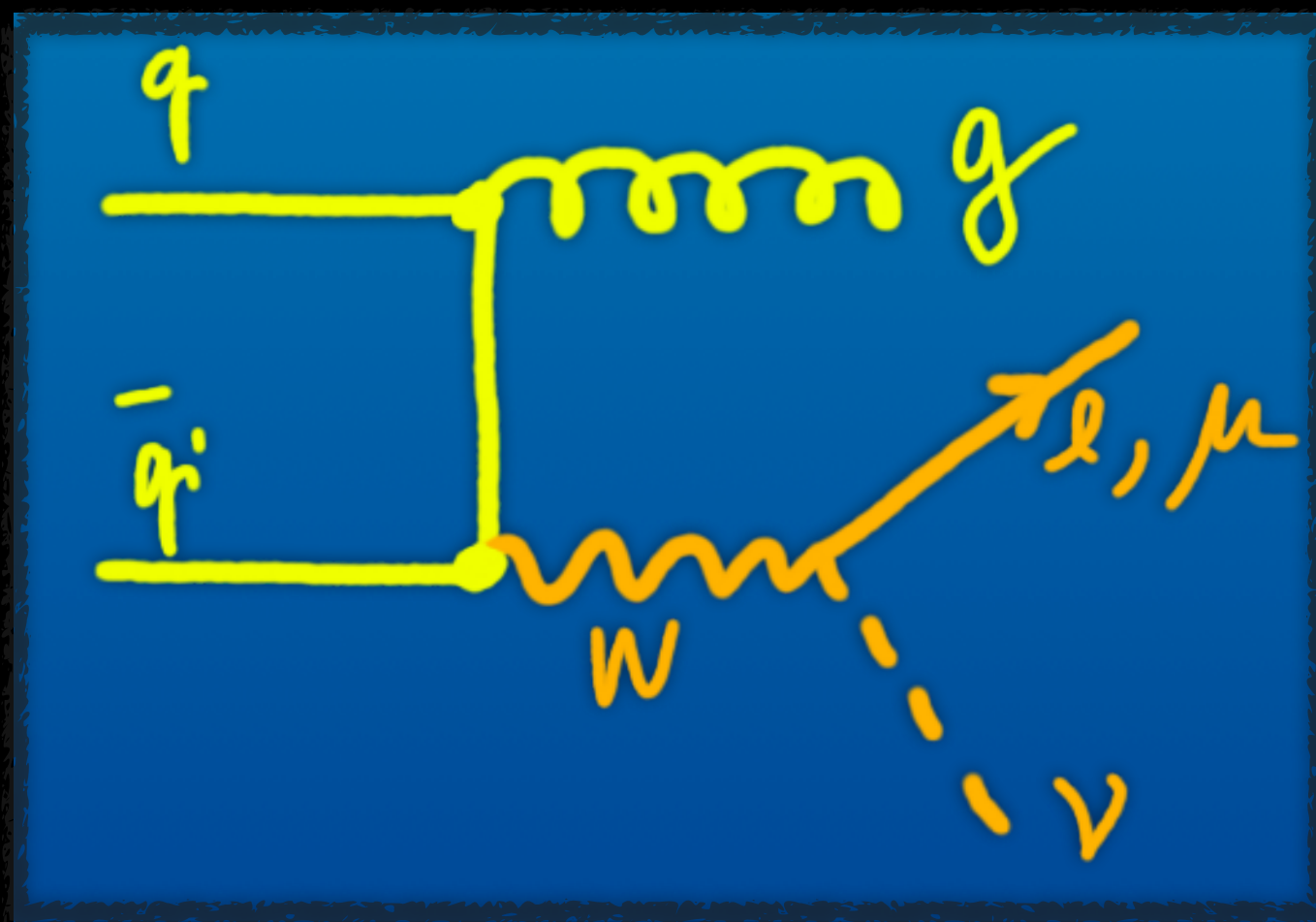
$$m_T = \sqrt{2 p_T^l p_T^{miss} (1 - \cos \Delta \varphi)}$$

Strategy: fit two kinematic distributions in several categories



$m_W = 80369.5 \pm 18.5$ MeV (large component from modeling)

The challenge of the W Mass Measurement



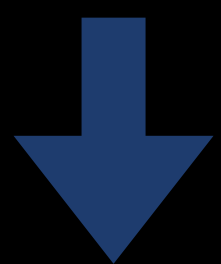
$$m_T = \sqrt{2 p_T^l p_T^{miss} (1 - \cos\Delta\varphi)}$$

Strategy: fit two kinematic distributions in several categories

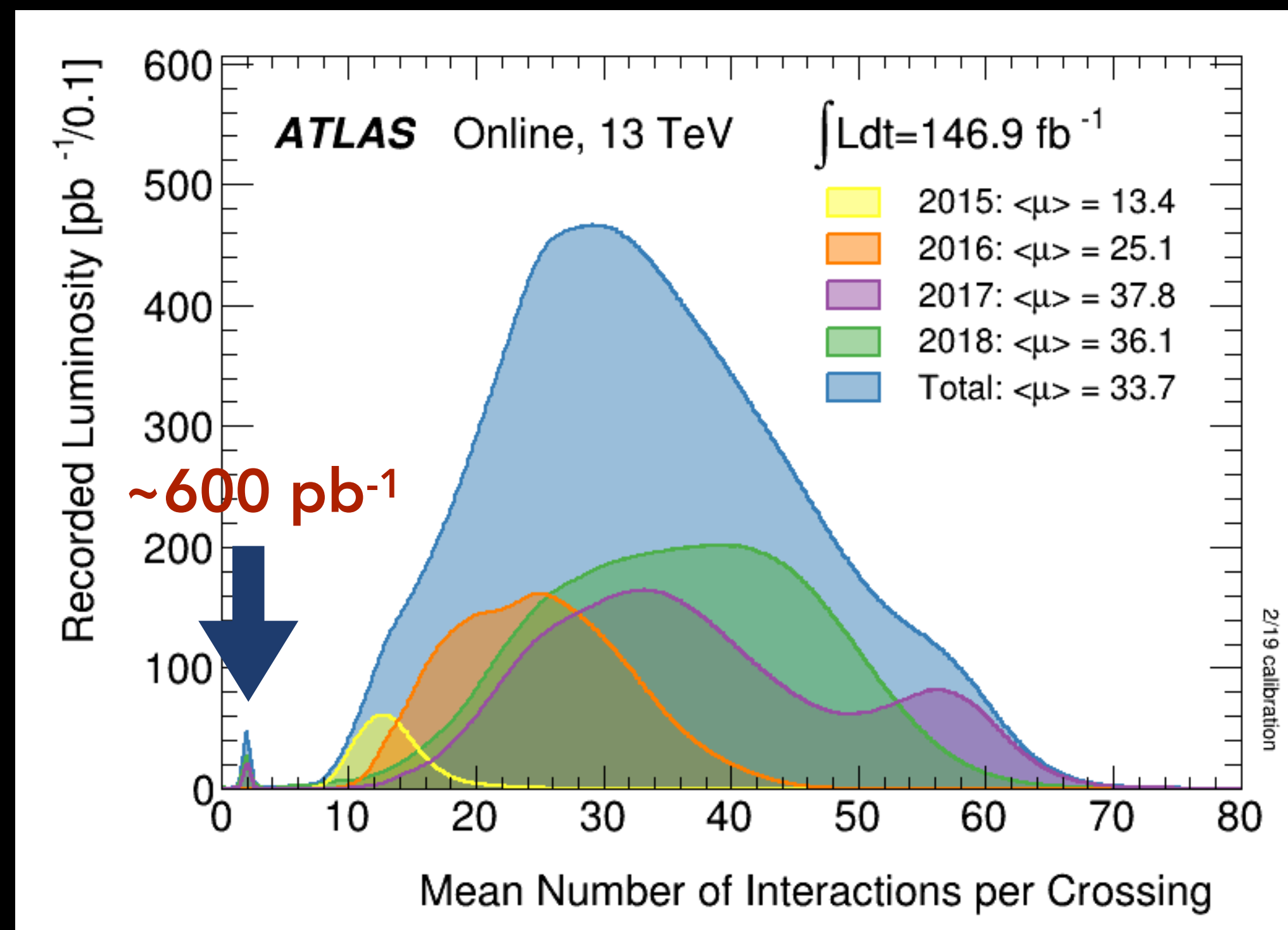
For further improvements:

Special precision-data runs with low pile-up events

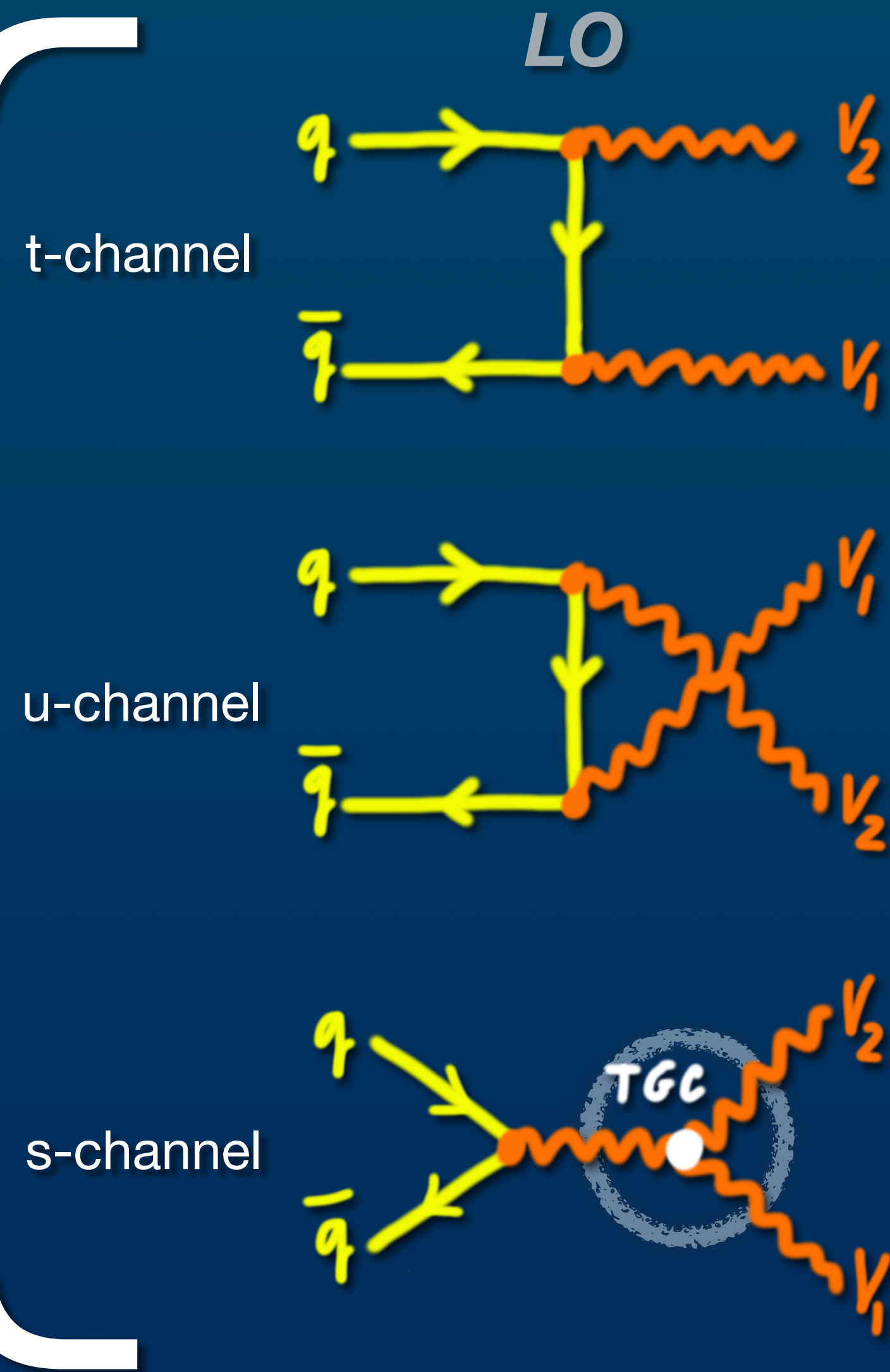
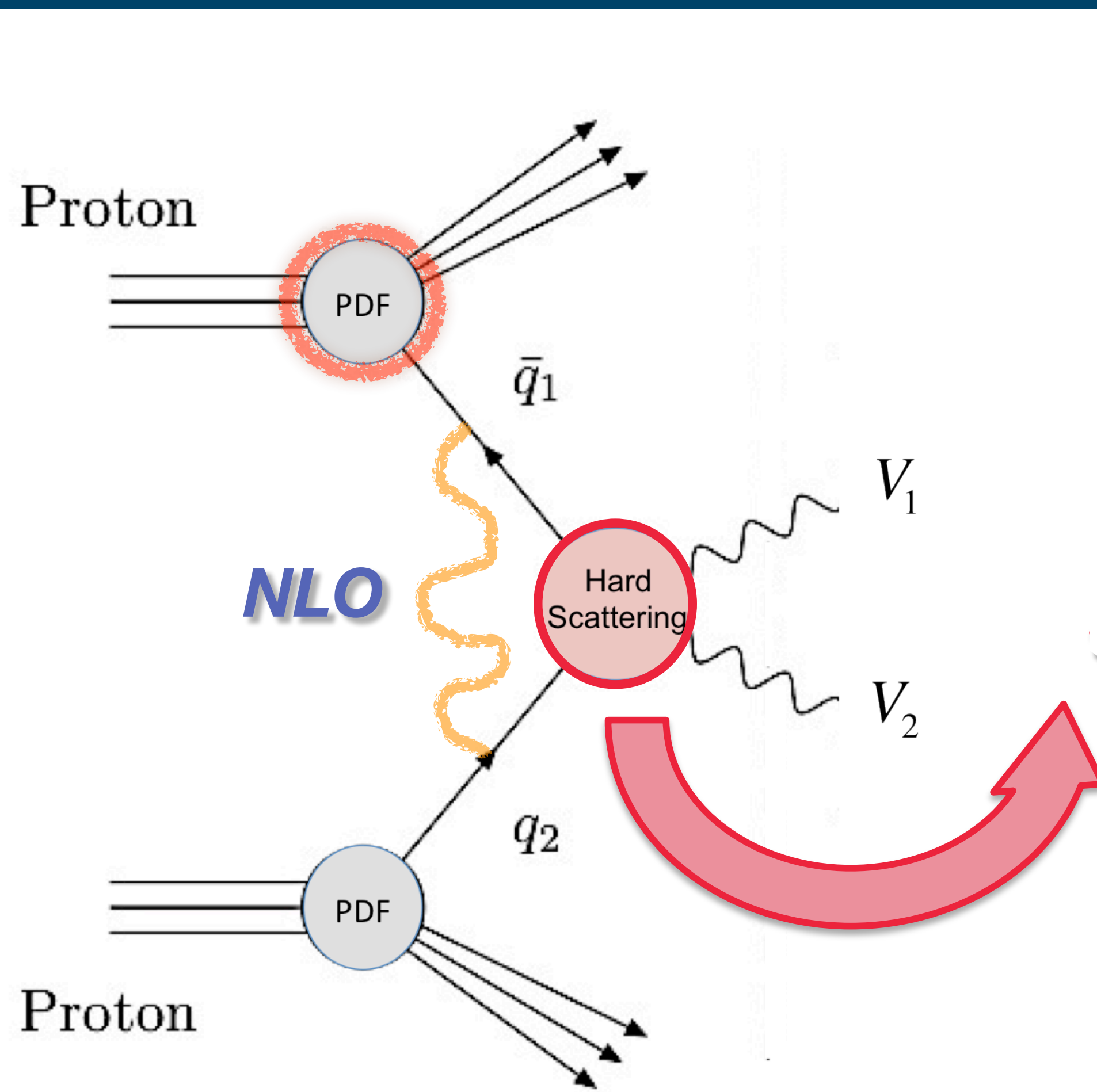
Cannot measure **W** mass directly from its decay particles because the **neutrino** is not measured directly in the detector



Measure **W** recoil precisely



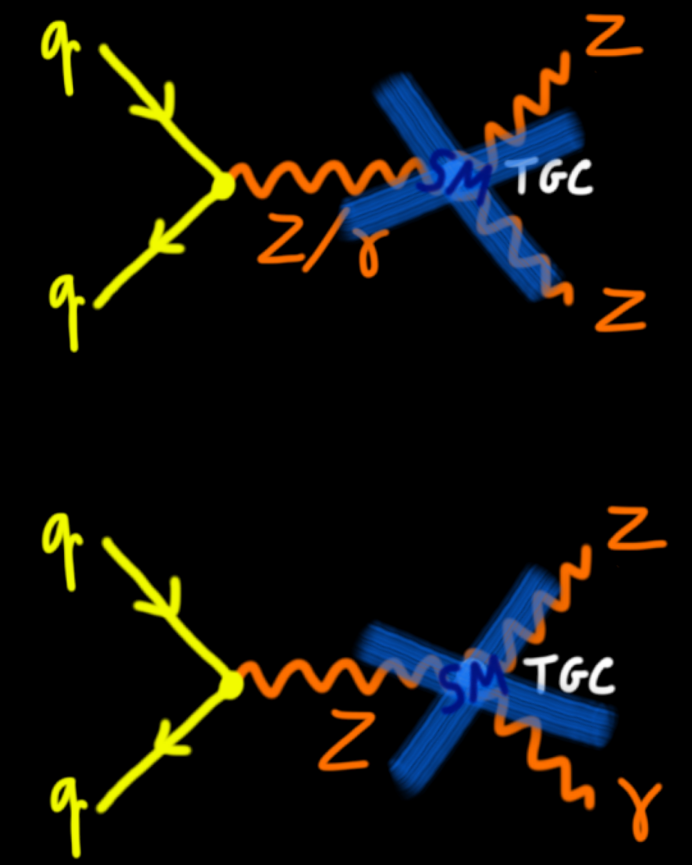
Diboson production at the LHC



Final states:

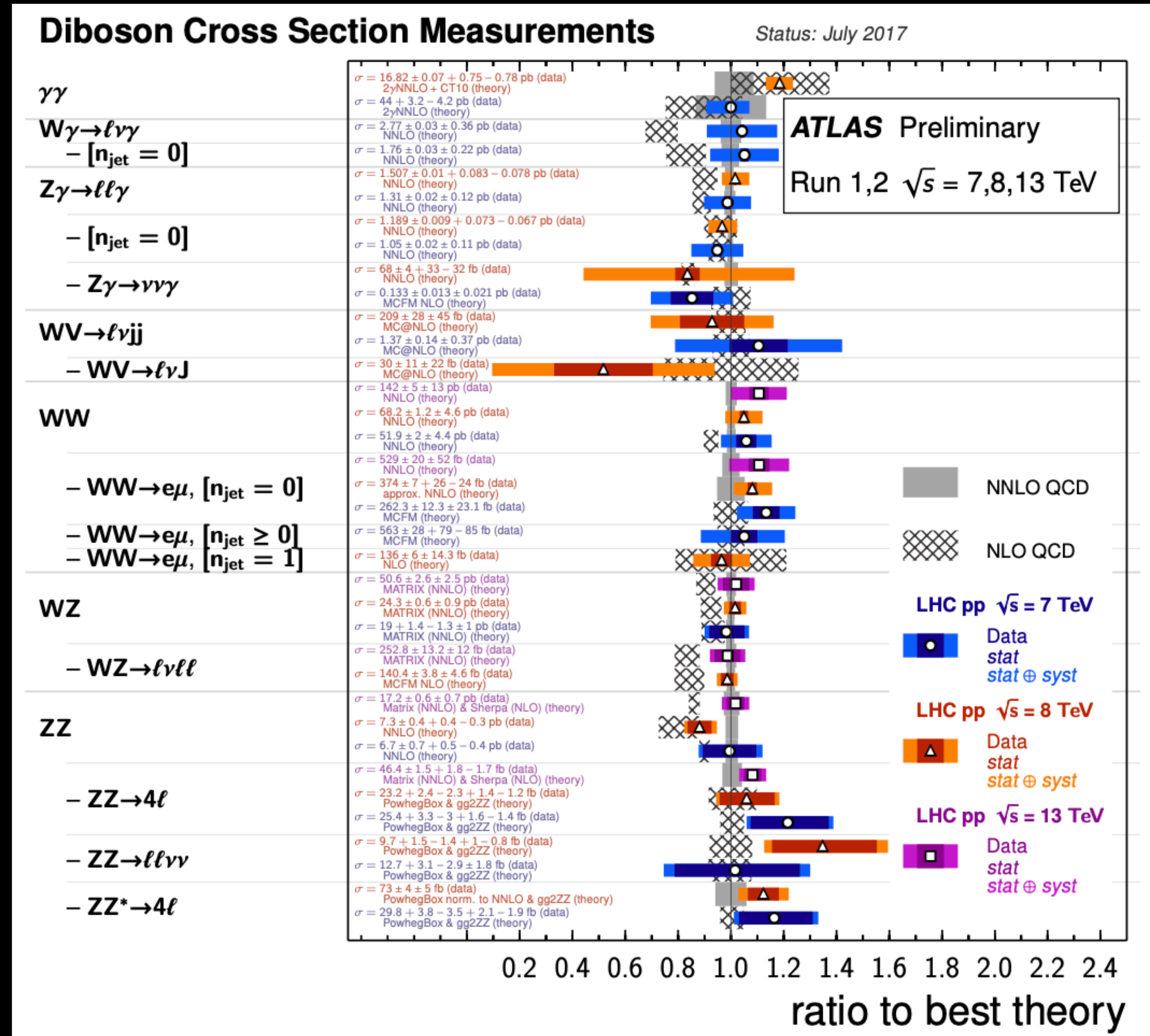
- WW**
- WZ**
- ZZ**
- $W\gamma$**
- $Z\gamma$**

Anomalous Triple Gauge Couplings

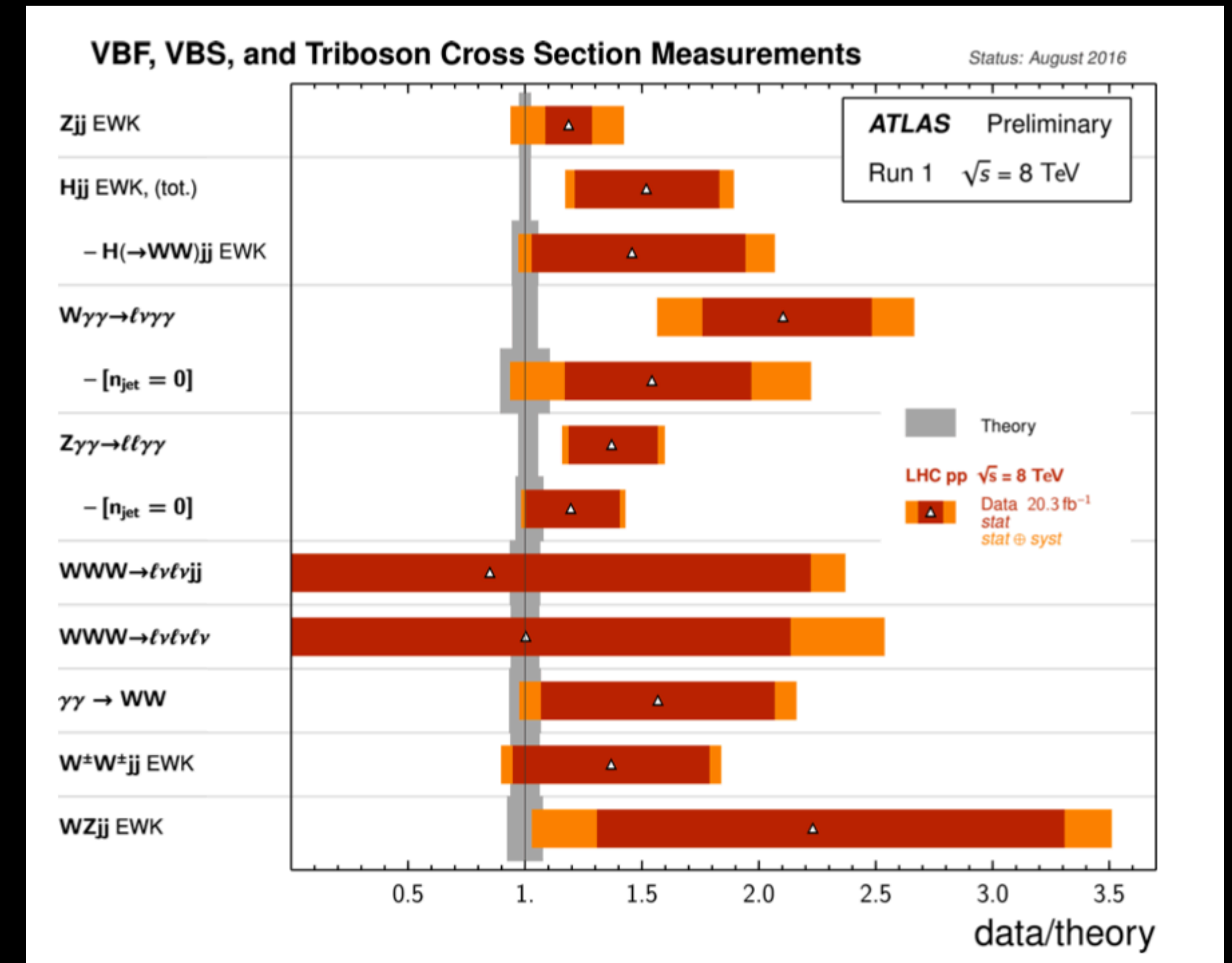


Di- (and Tri-) boson measurements at the LHC

Large number of processes measured and used to constrain anomalous gauge couplings



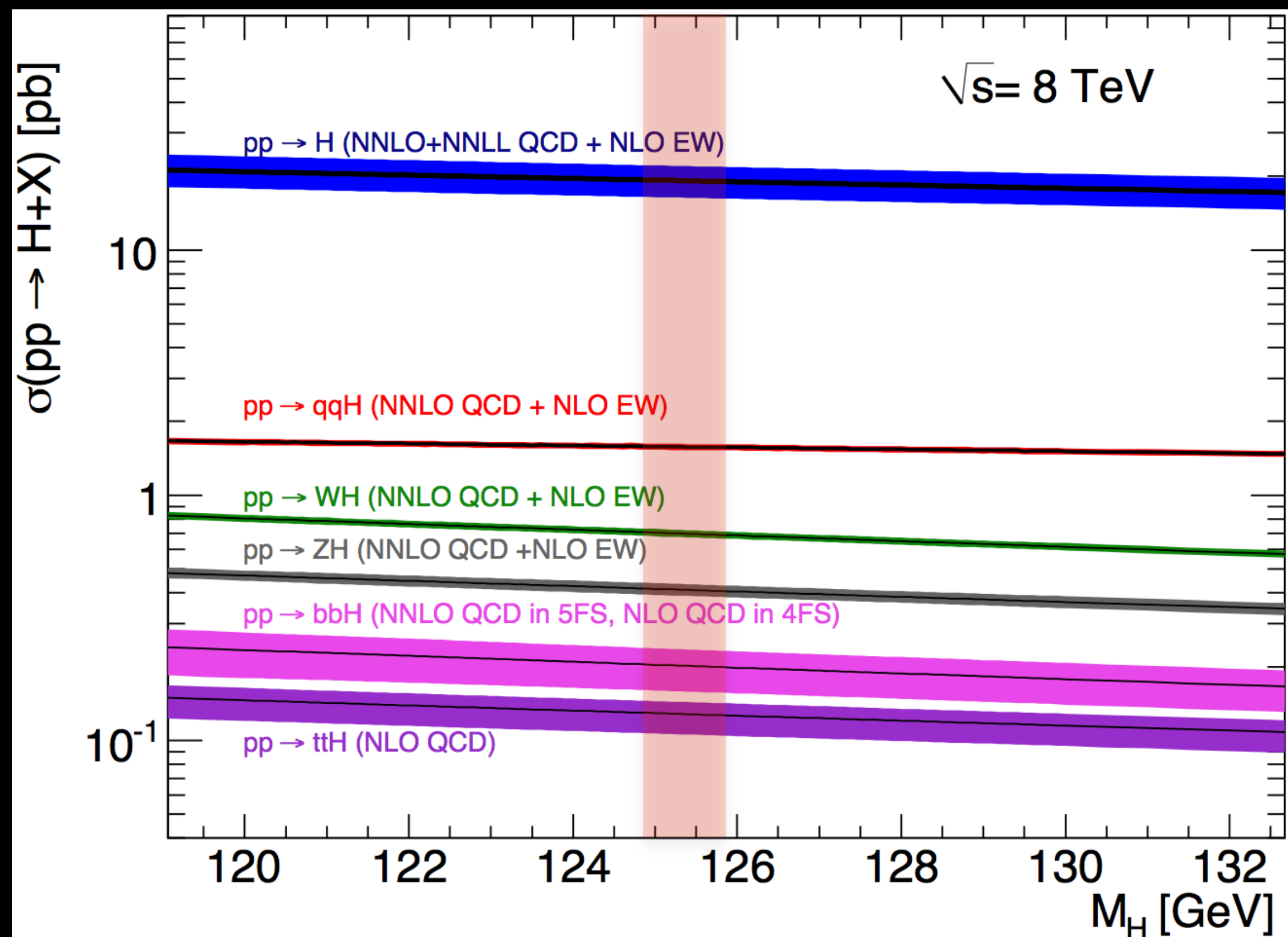
Even more rare processes...



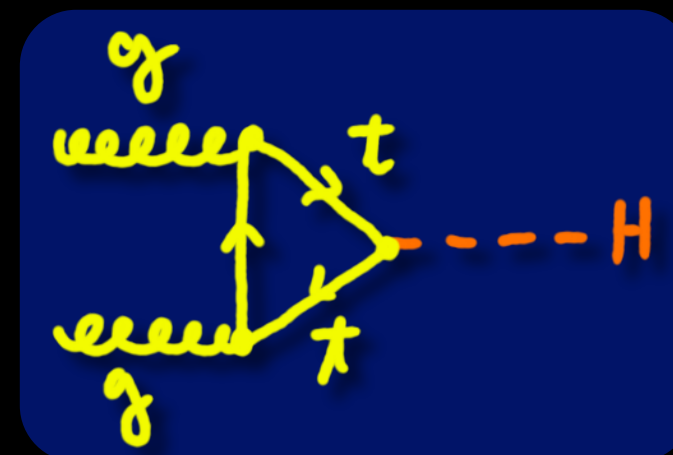
Exposing the need for NNLO QCD calculations

Higgs

Higgs Production at LHC ($m_H = 125 \text{ GeV}$)

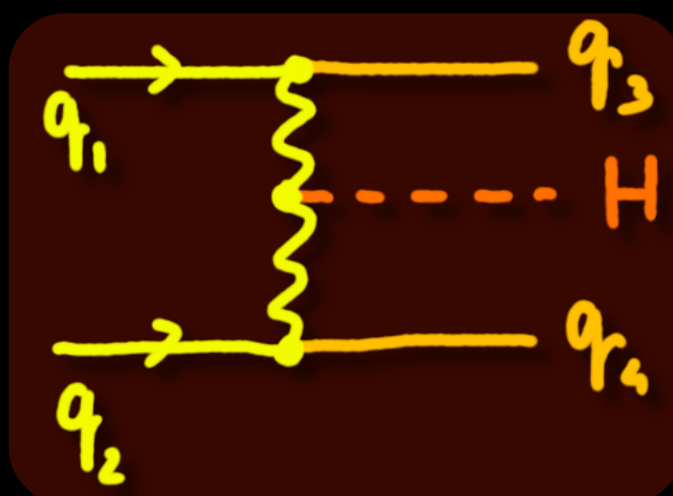


Earlier cross section calculations assumed a full range of Higgs masses



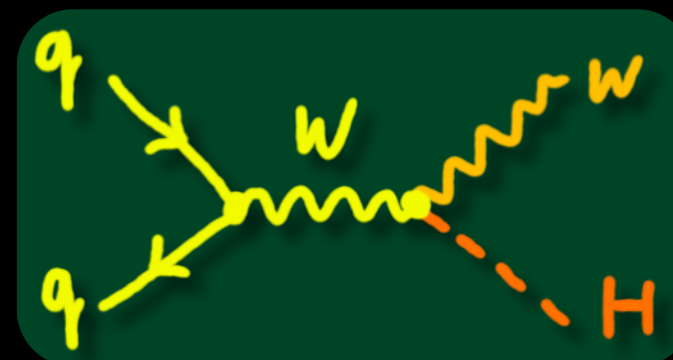
Gluon Fusion (ggF)

NNLO+NNLL QCD
+ NLO EWK



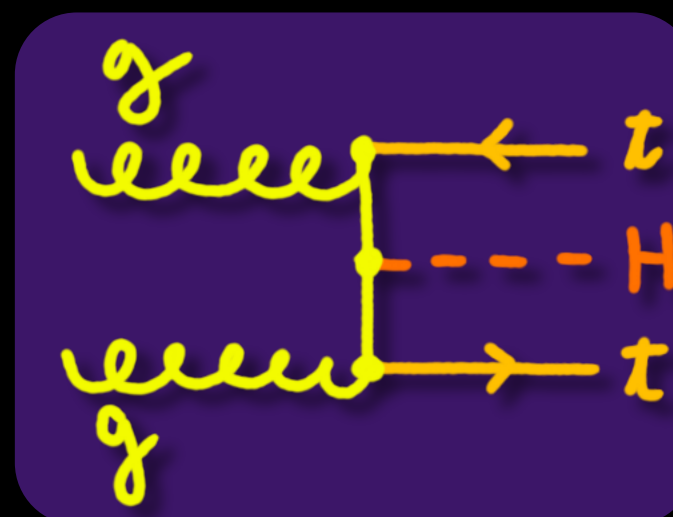
Vector Boson Fusion (VBF)

NNLO QCD
+ NLO EWK



Associate Production (VH)

NNLO QCD
+ NLO EWK



Associate Production (ttH)

NLO QCD

Production Rates
in Run 2 (13 TeV)
for 150 fb^{-1}

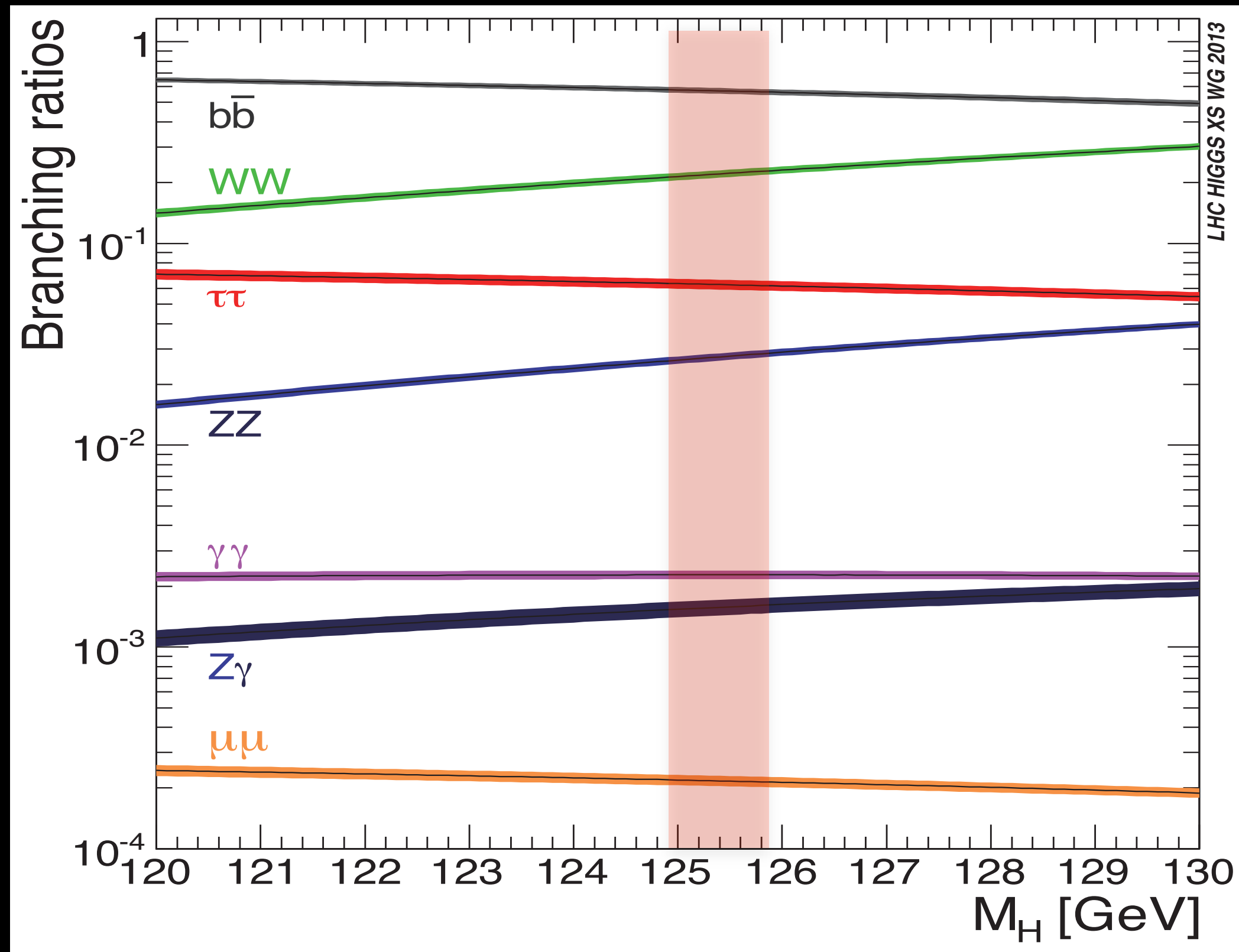
~8M events

~600k events

~400k events

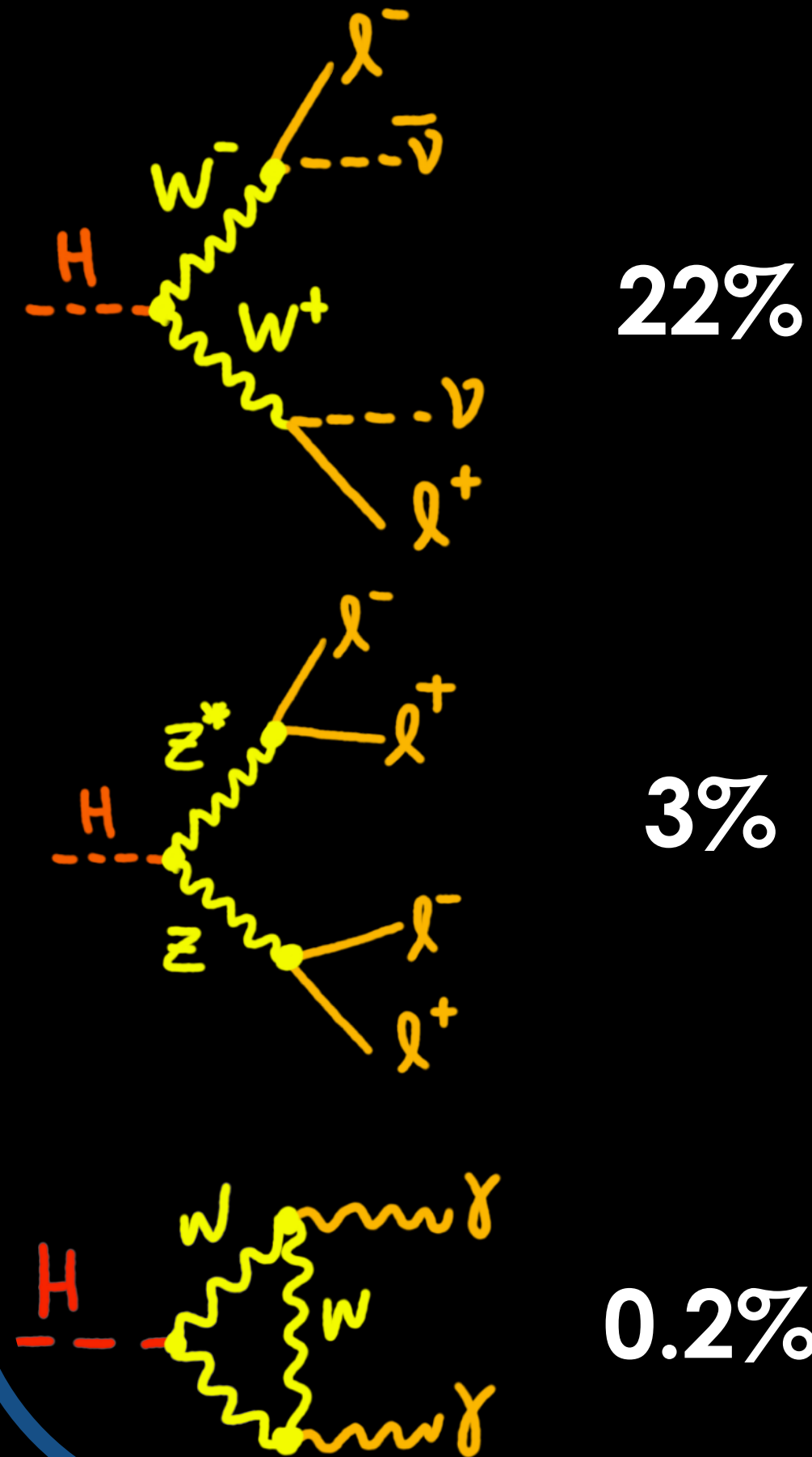
~80k events

Higgs Decay Branching Fractions ($m_H = 125 \text{ GeV}$)



Possible to test Higgs sector with many decay modes

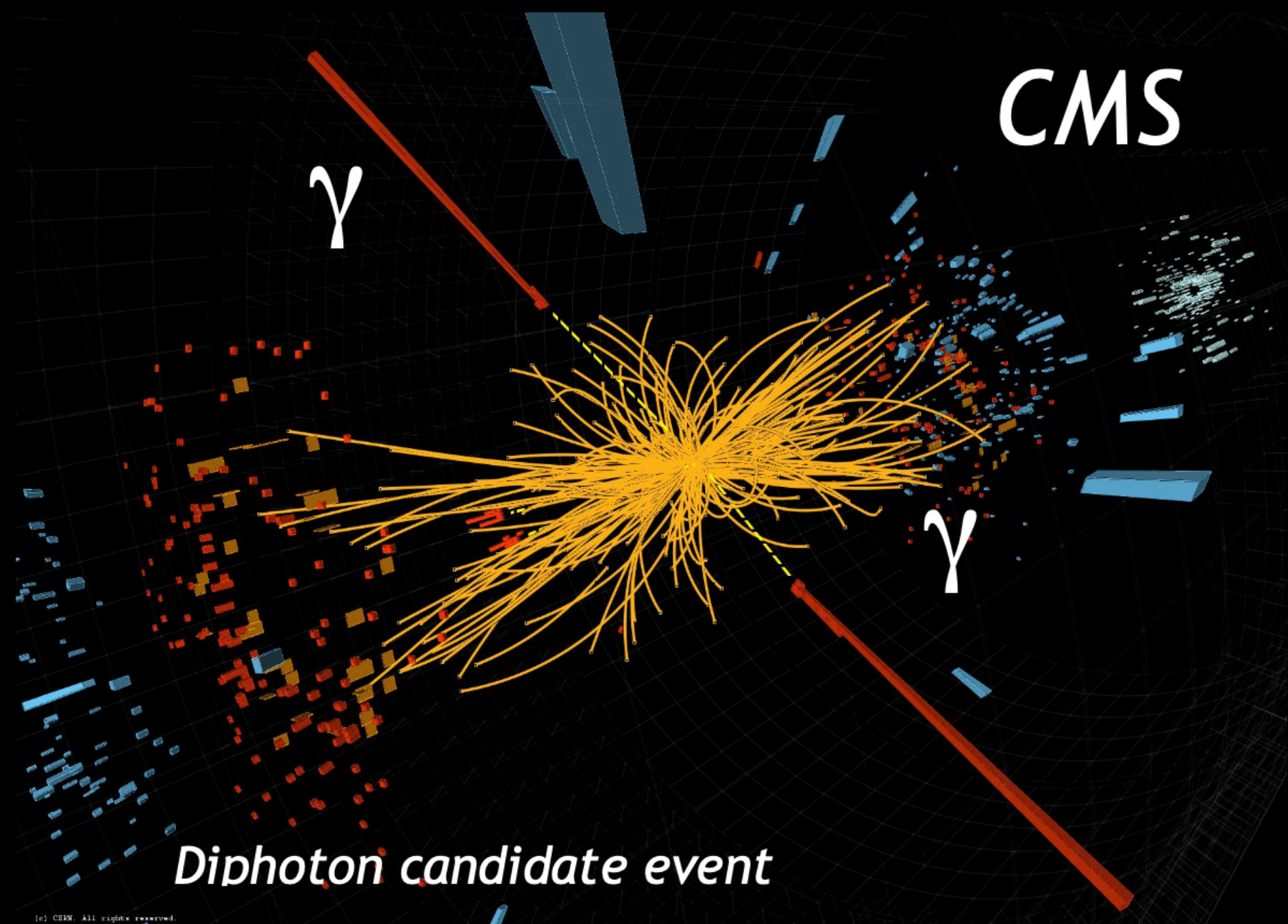
Boson decays



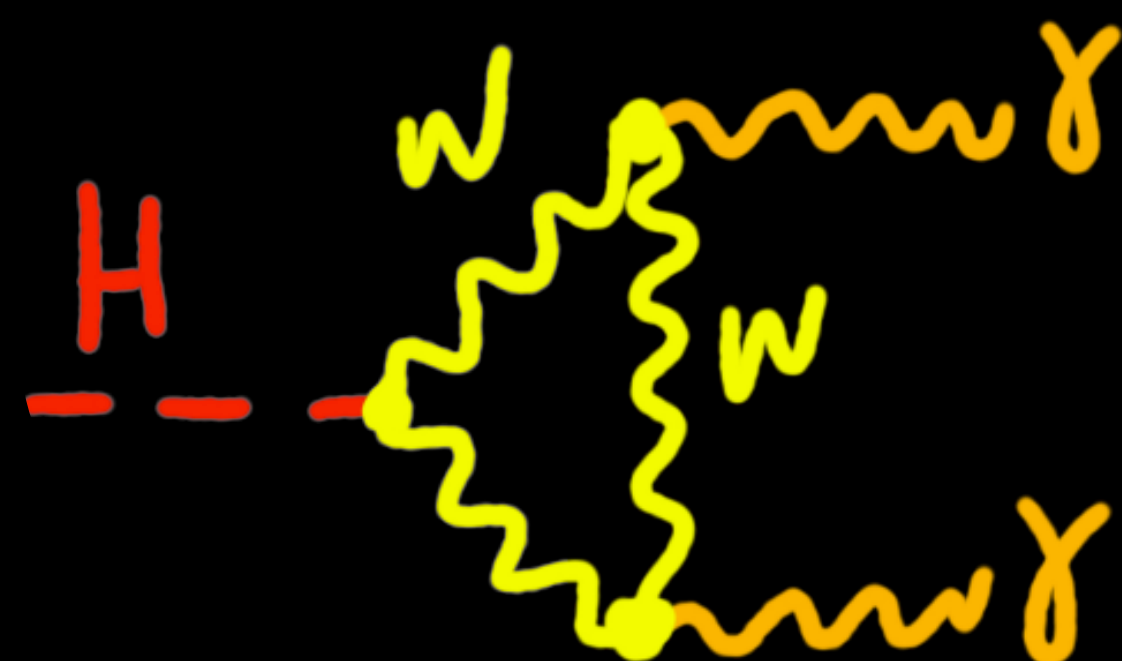
Fermion decays



Discovery channel: 2 photons

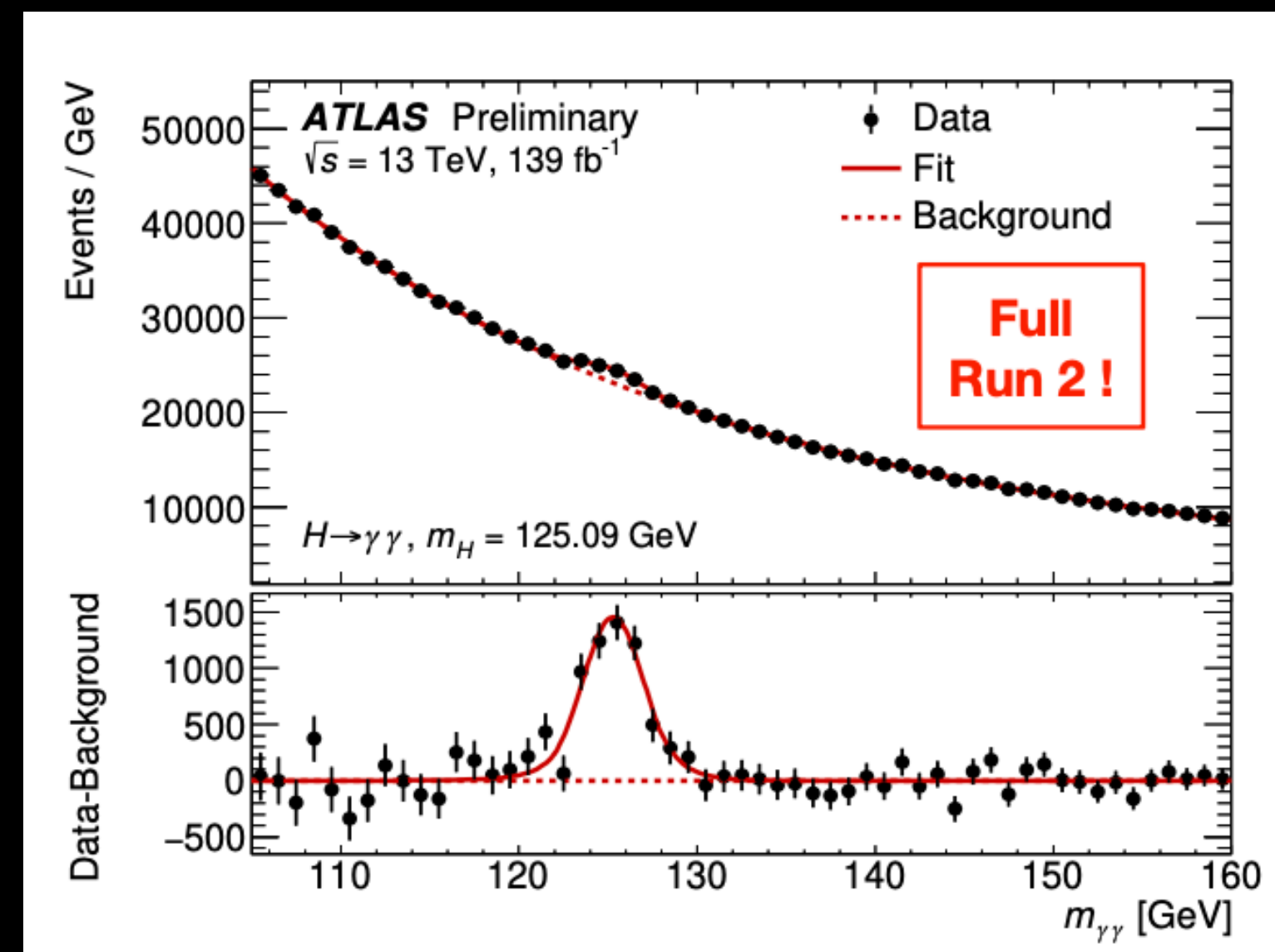
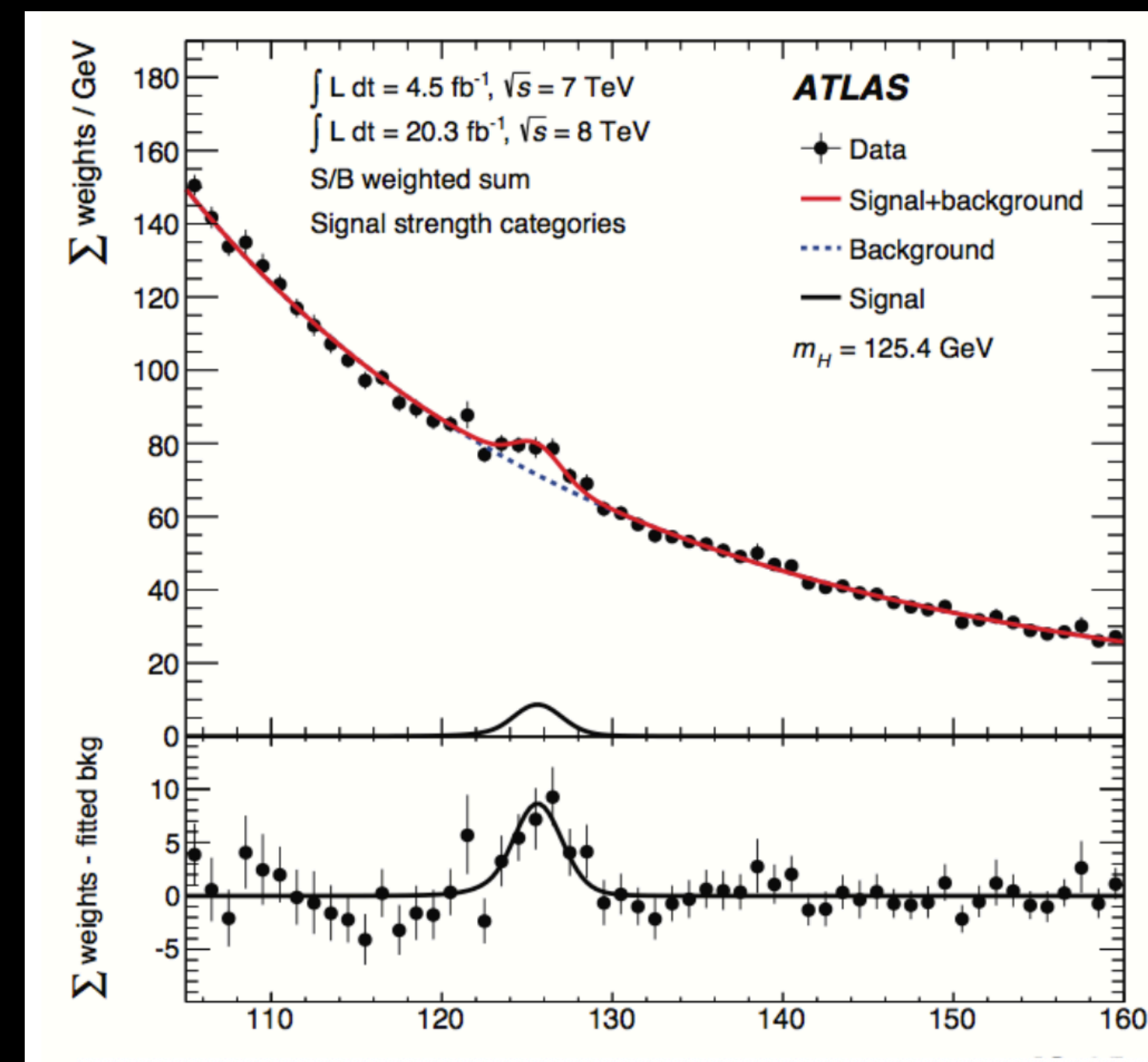


High mass resolution $O(1\%)$ allowing data driven estimate of background in the sidebands

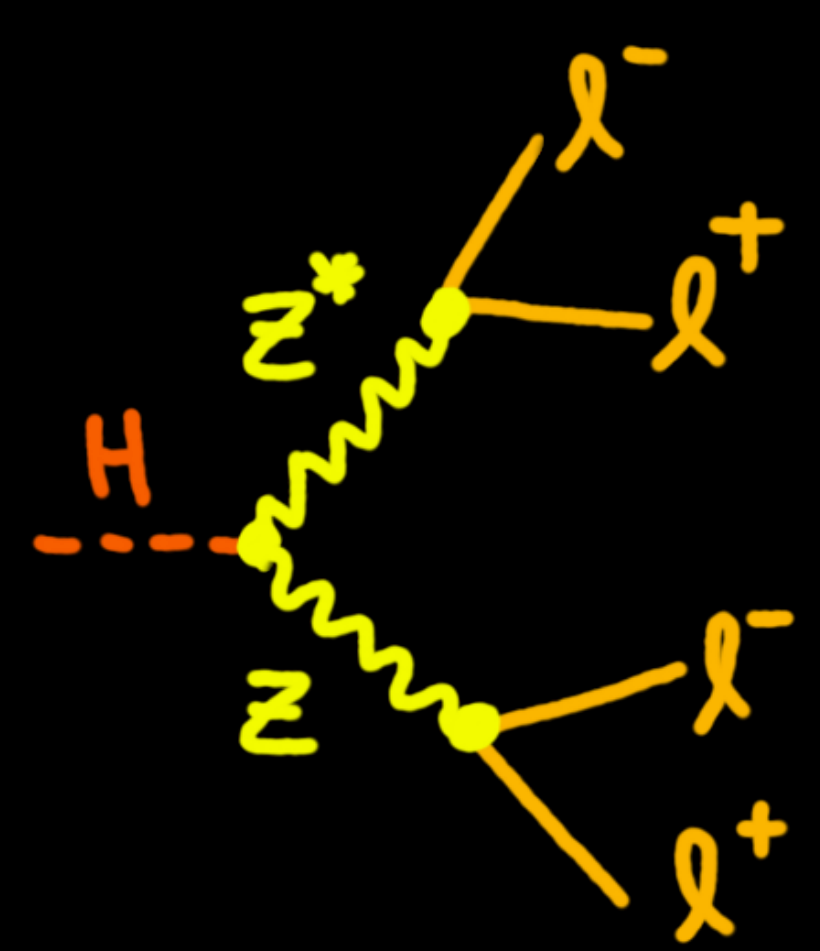
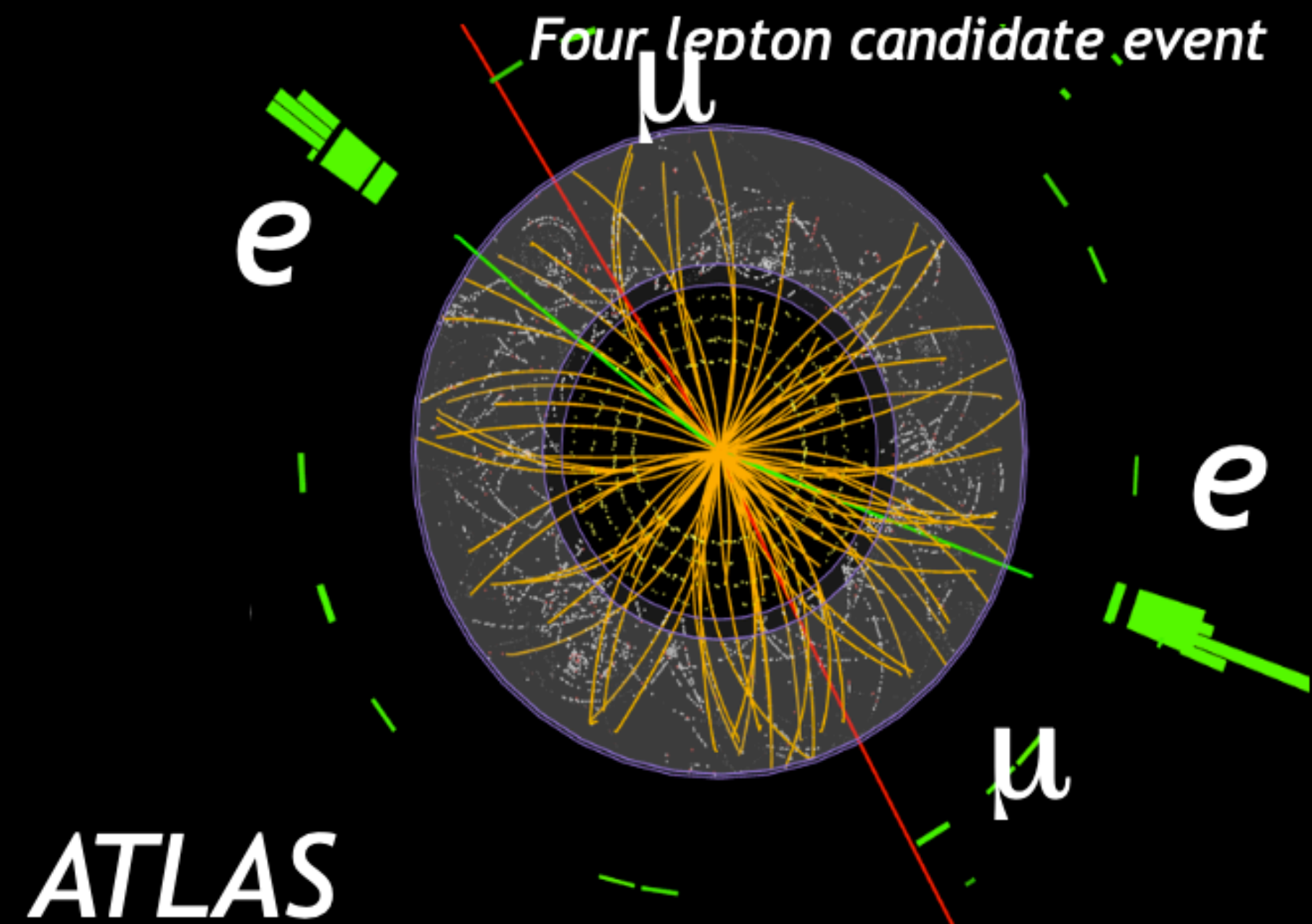


Observation implies that the bump does not originate from spin 1 particle: Landau-Yang theorem

Sensitive to new physics



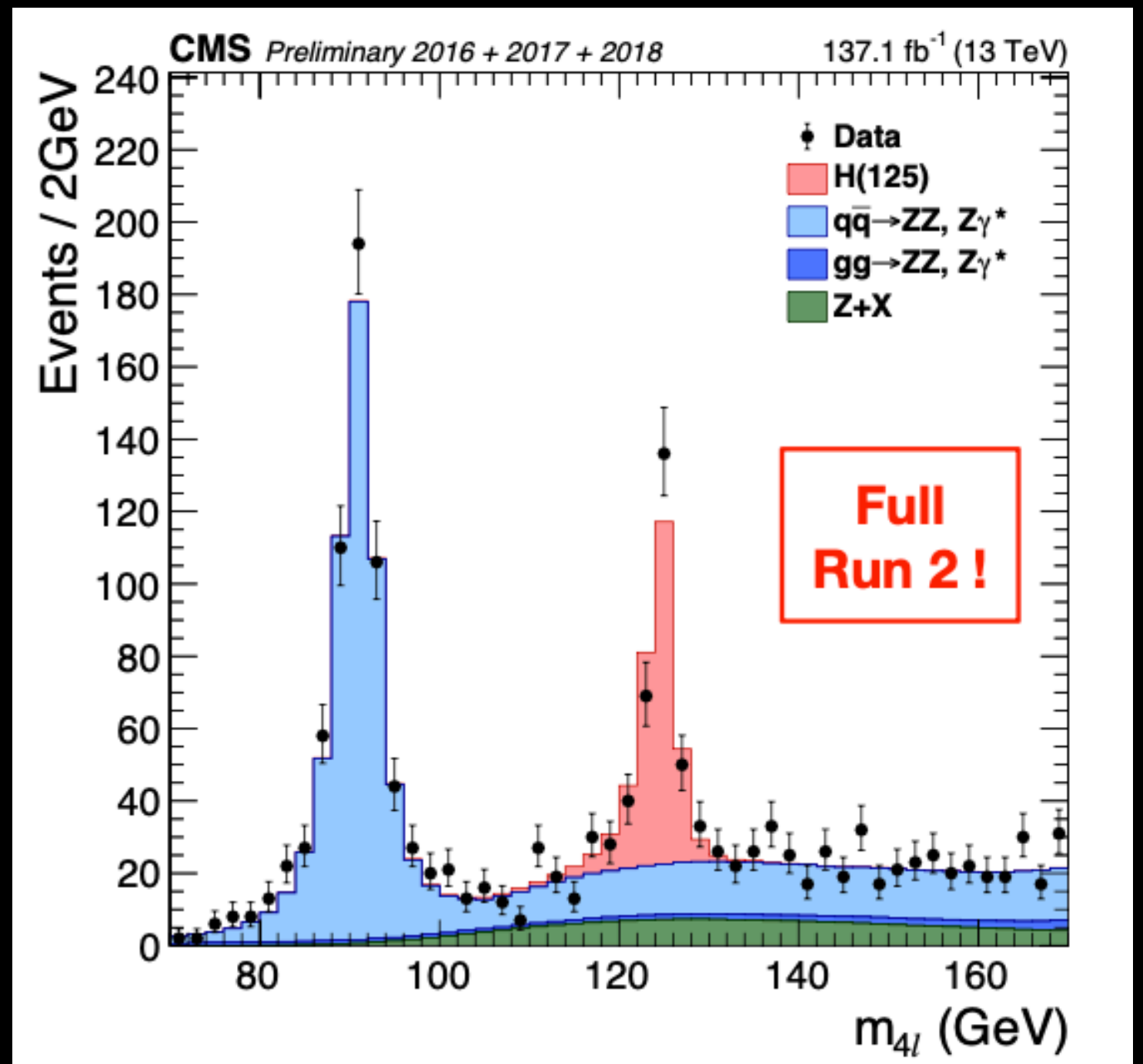
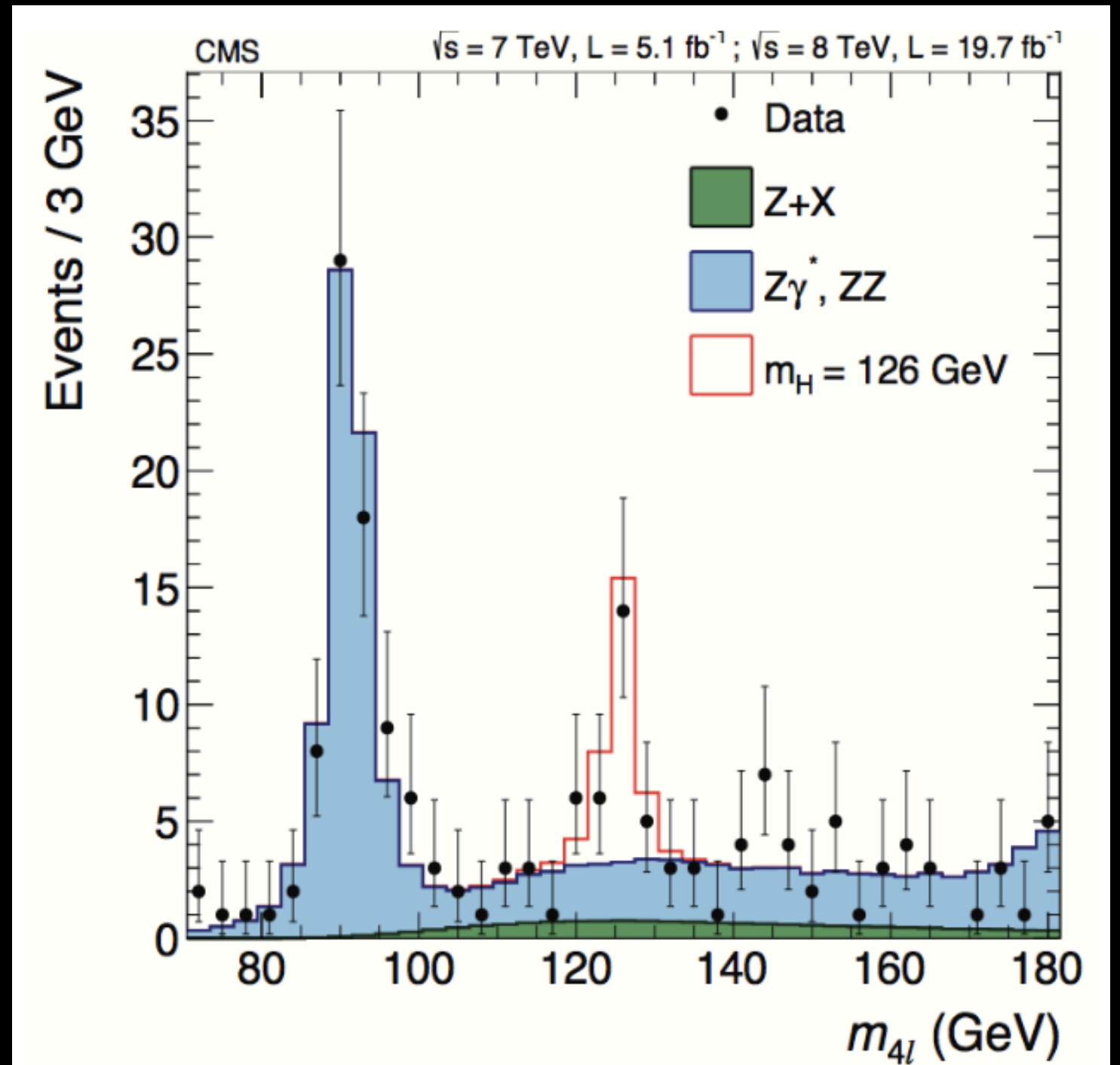
Discovery channel: 4 leptons



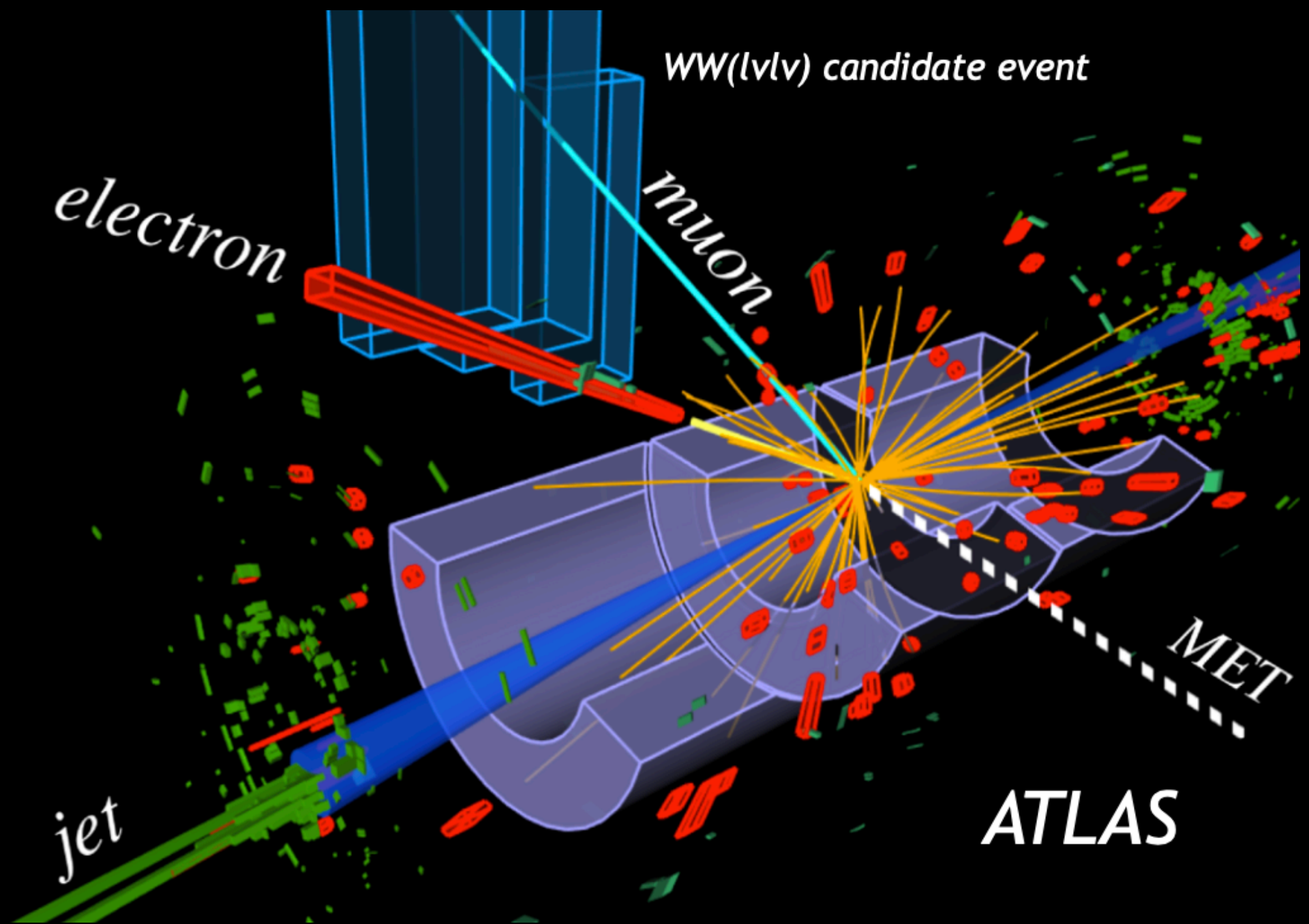
Very high S/B ratio

Backgrounds can be estimated from MC

Polarization can be reconstructed



Discovery channel: WW

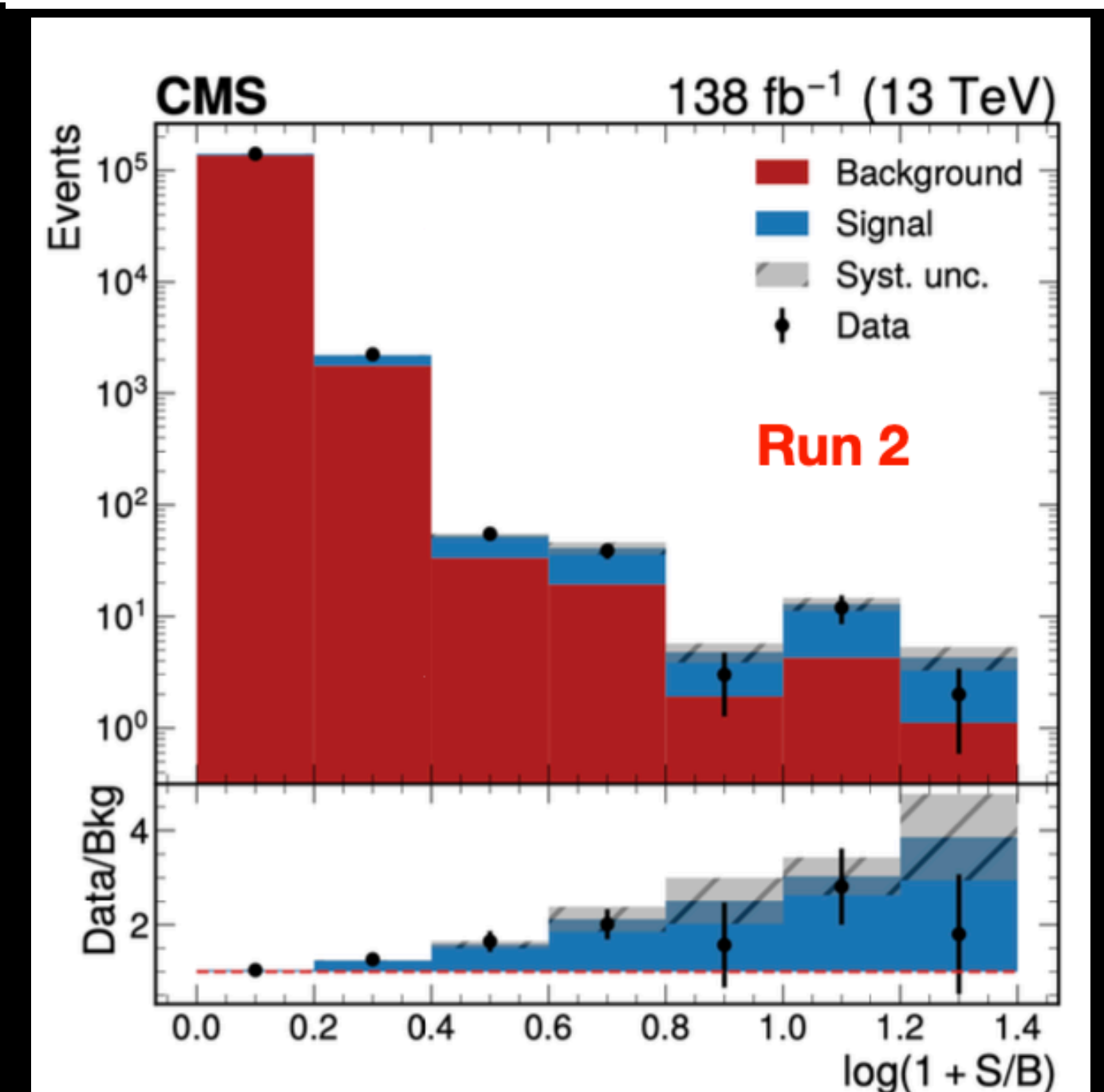
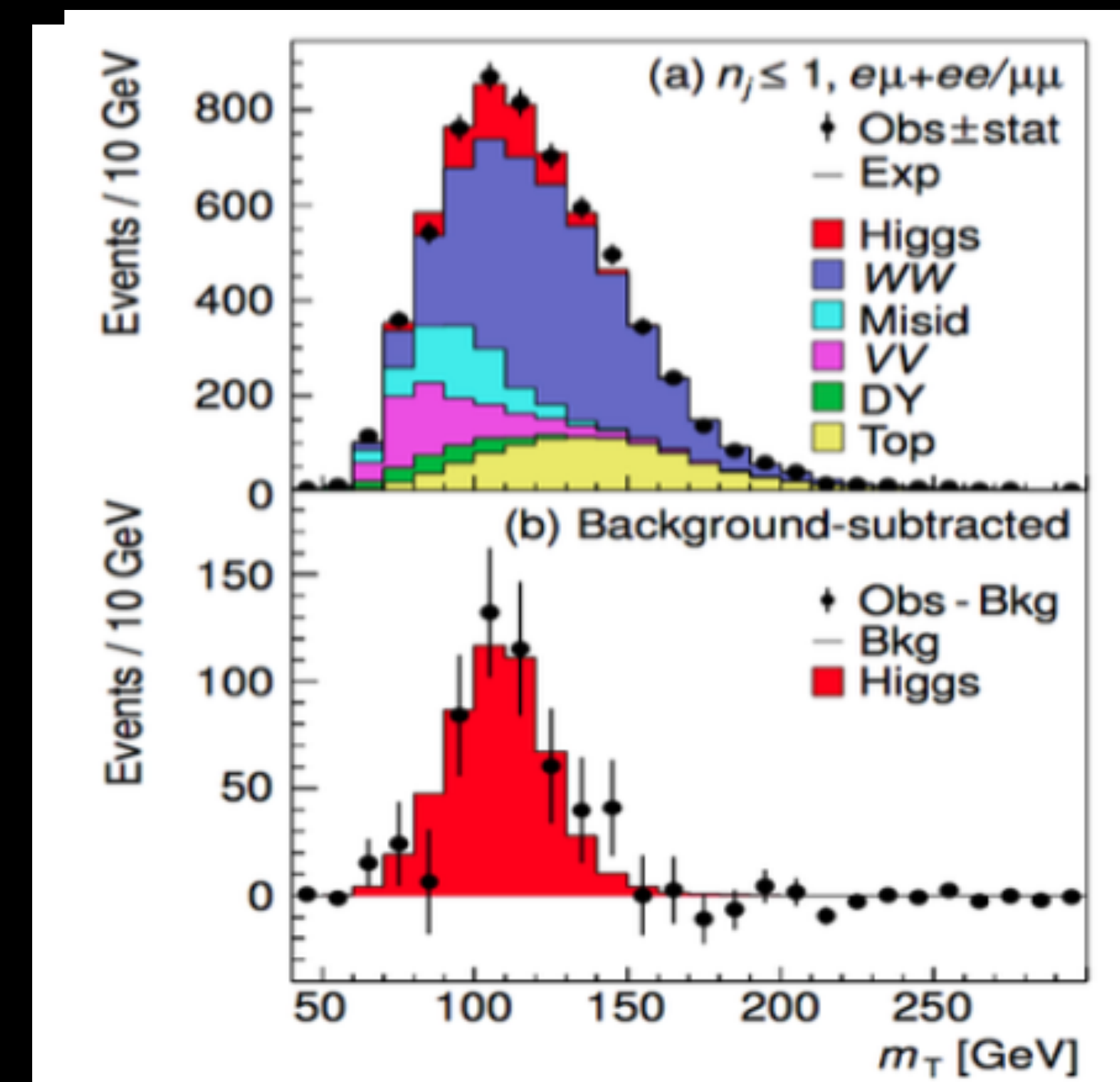
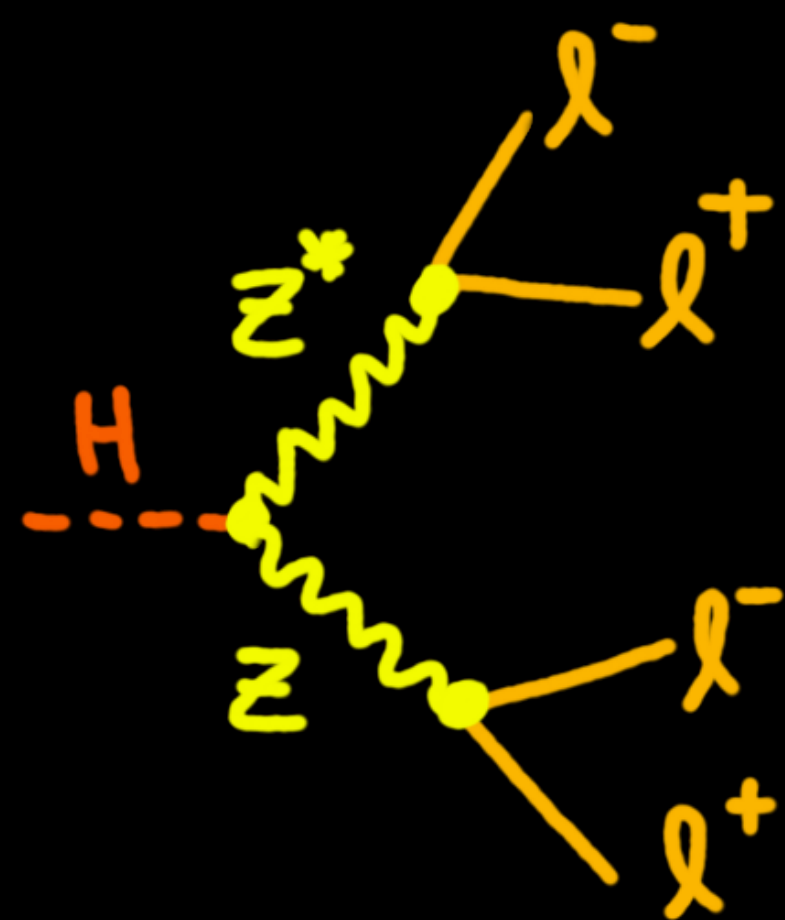


Large event rate

but

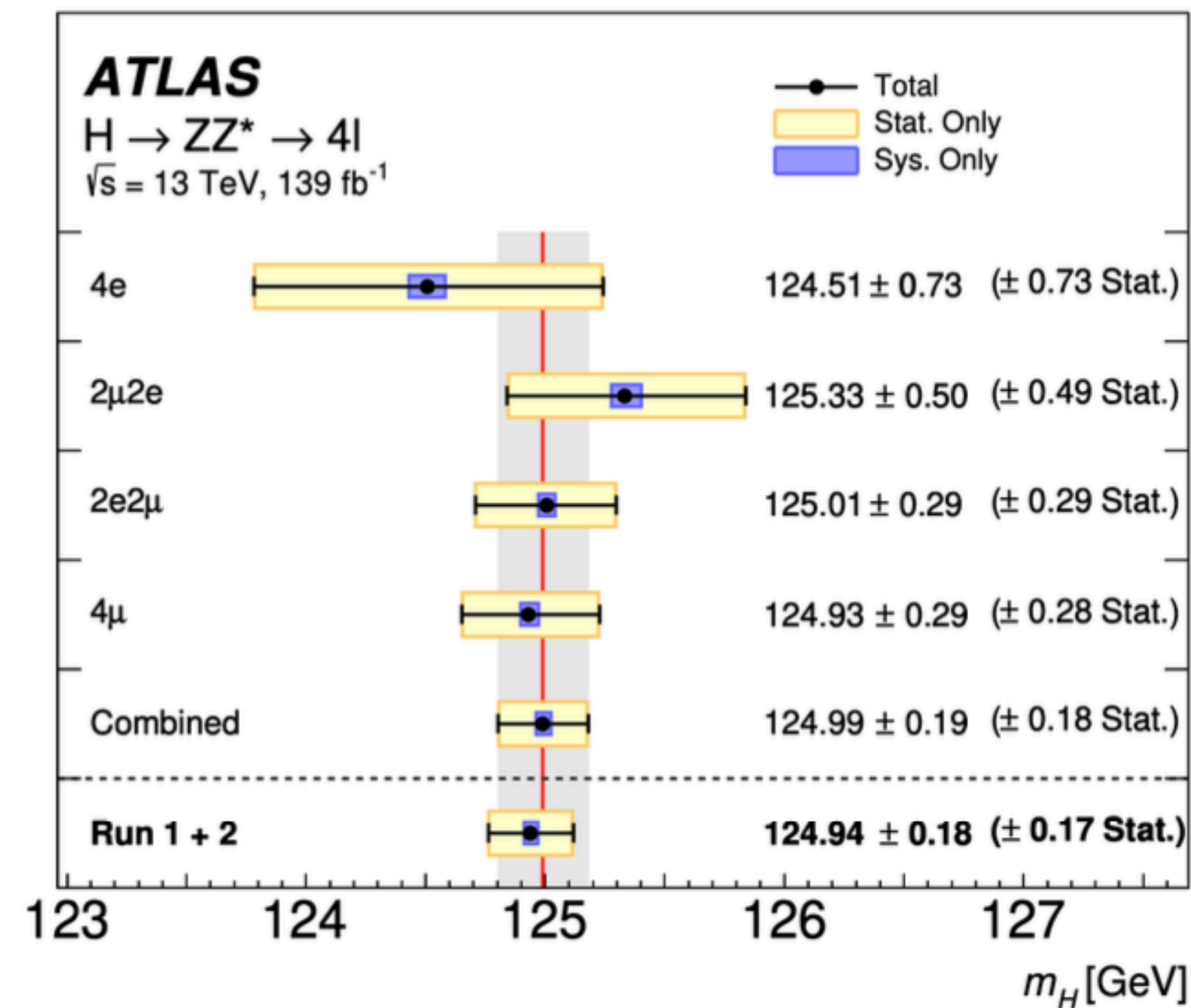
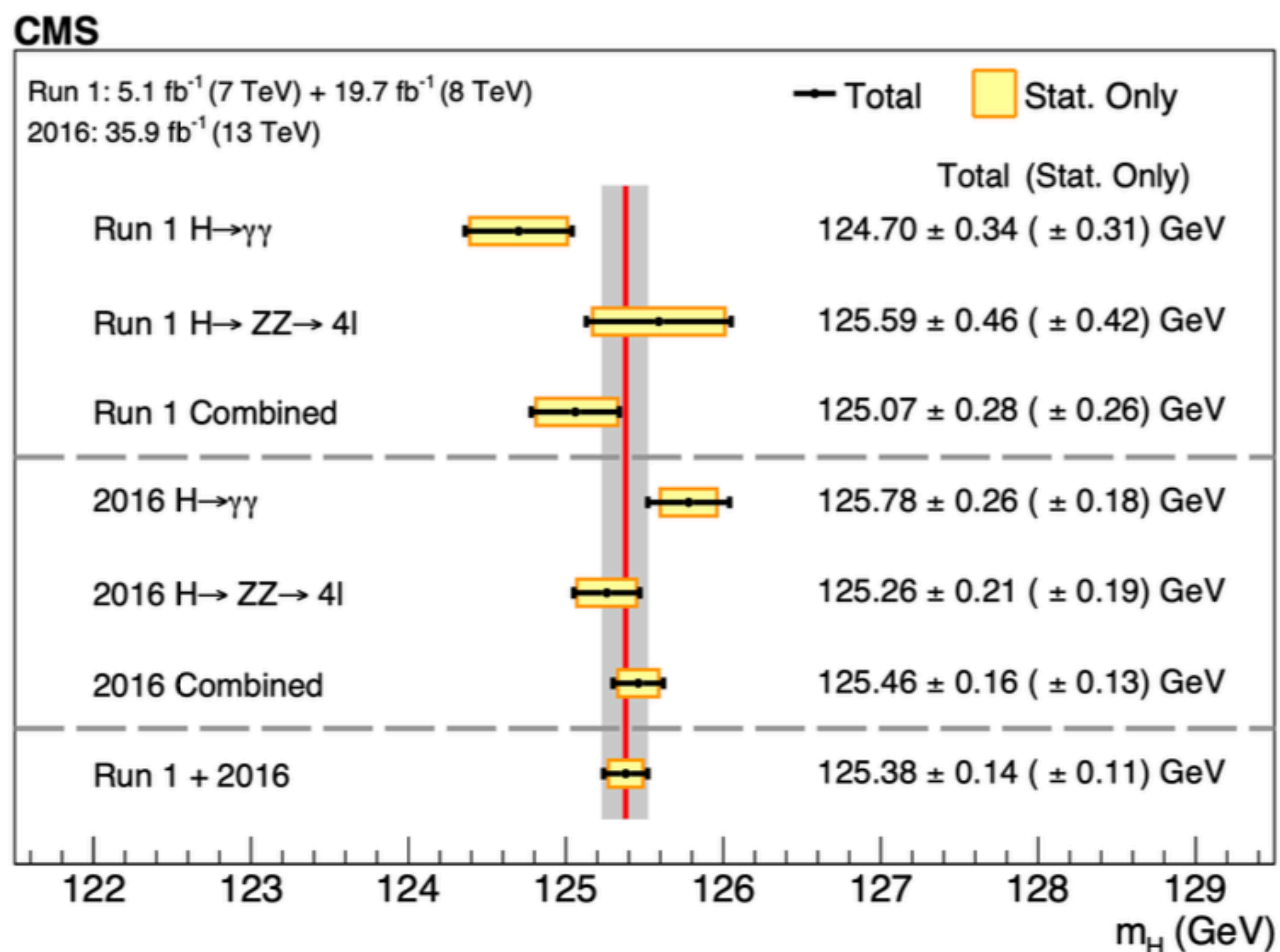
Large background from SM WW and top quark production

Mass resolution spoiled by the presence of neutrinos



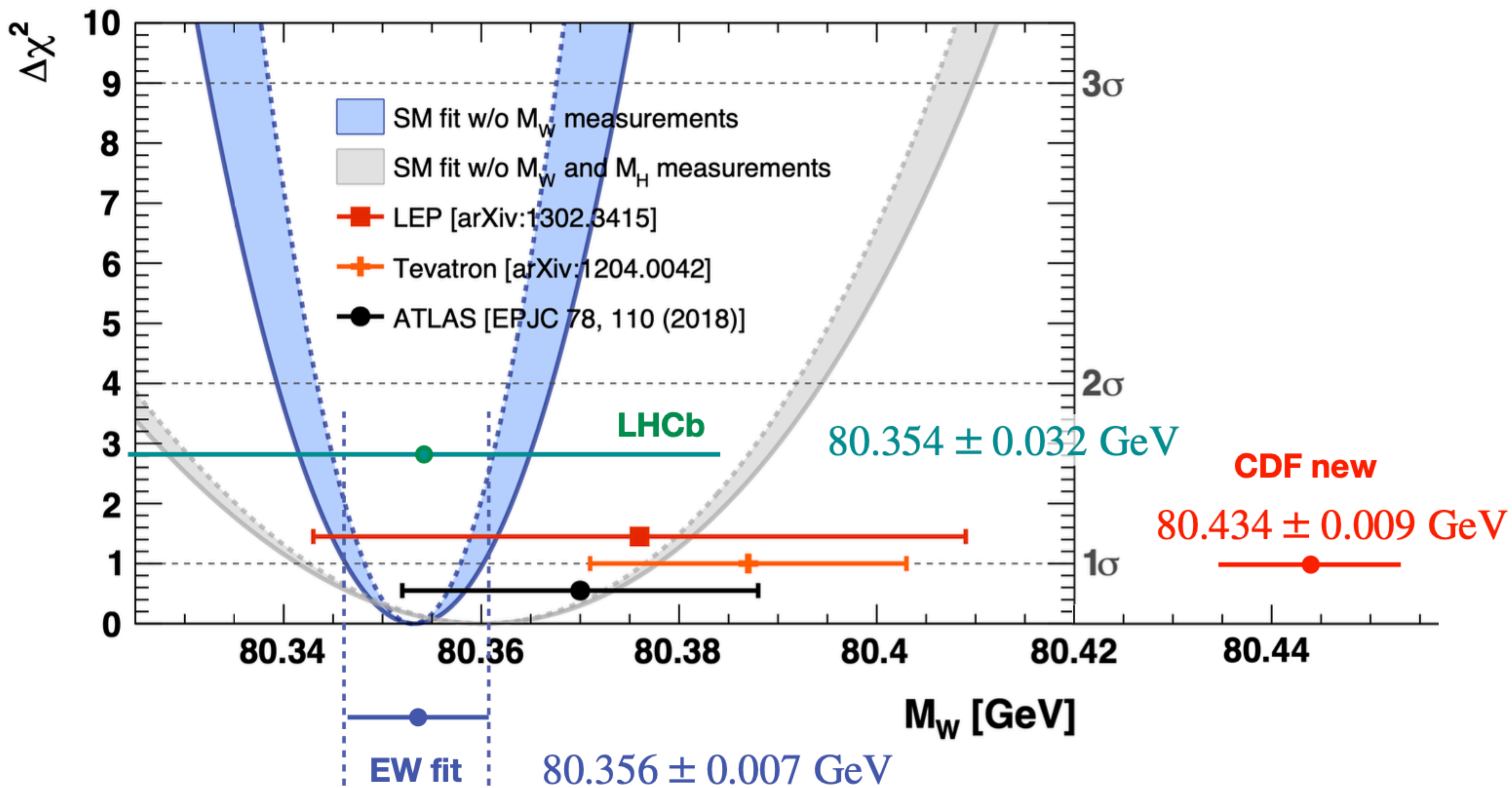
Higgs boson mass

Measurement done exclusively in the di-photon and 4-leptons channels



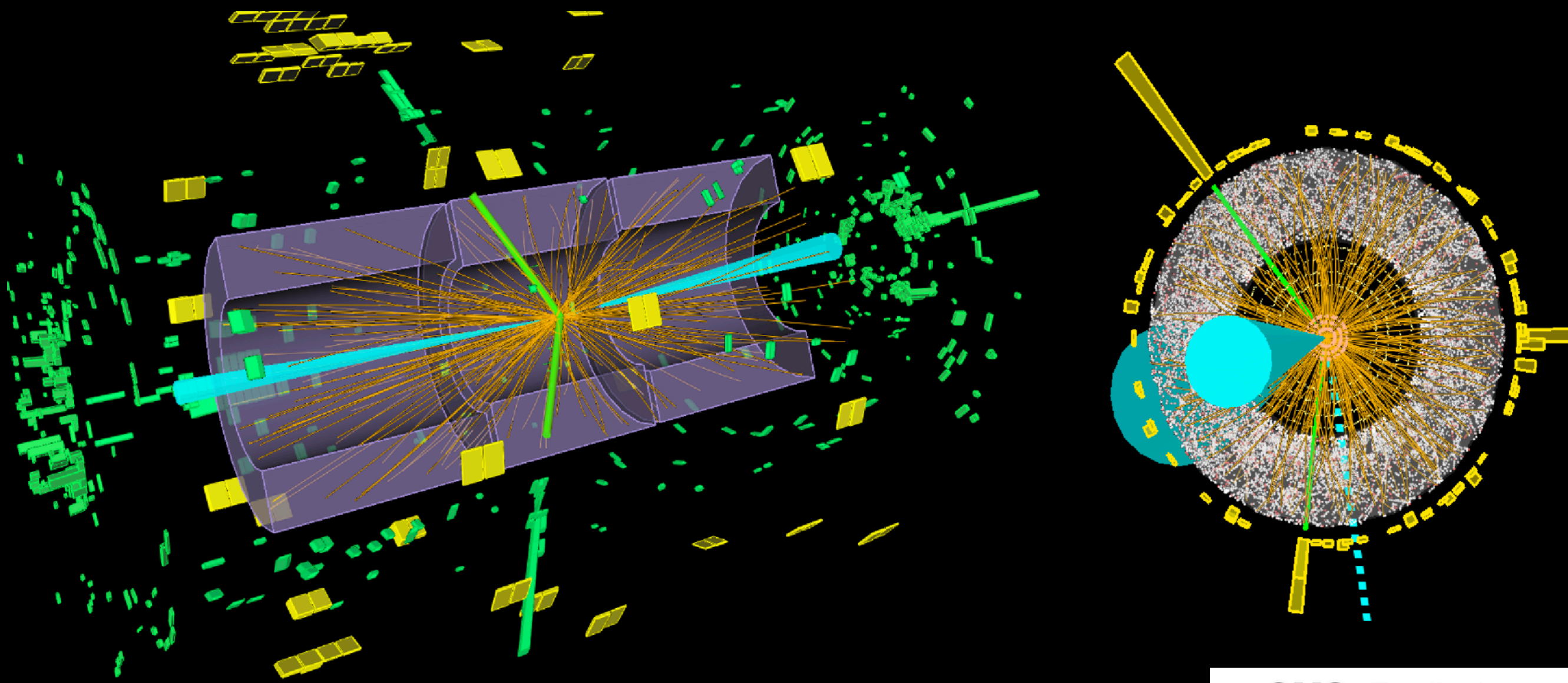
0.11% uncertainty

Precision measurements and predictions

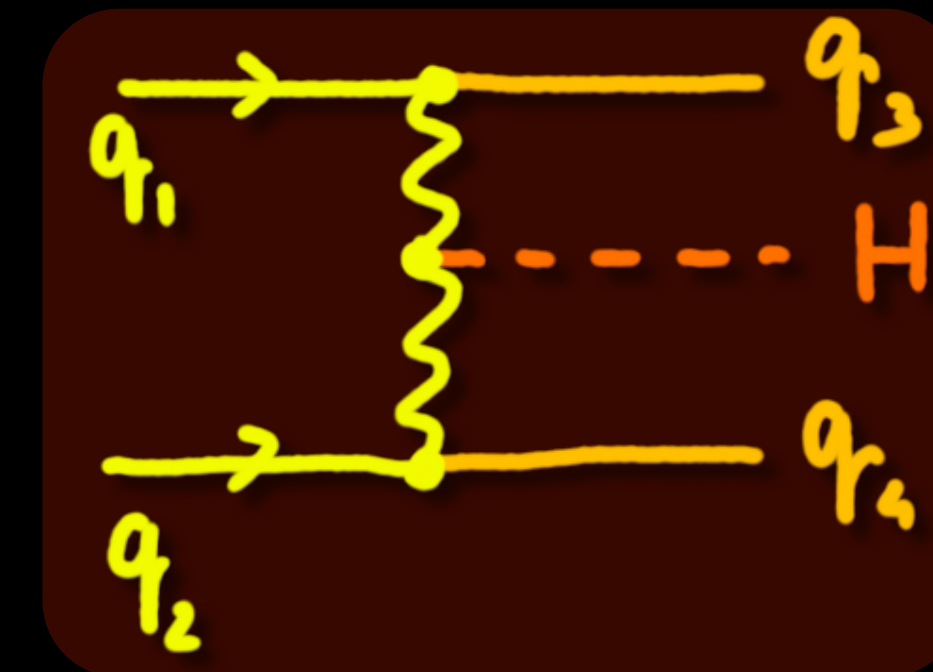


The knowledge of the Higgs mass has large impact on the precision of indirect measurements!

Higgs observation in tau channel

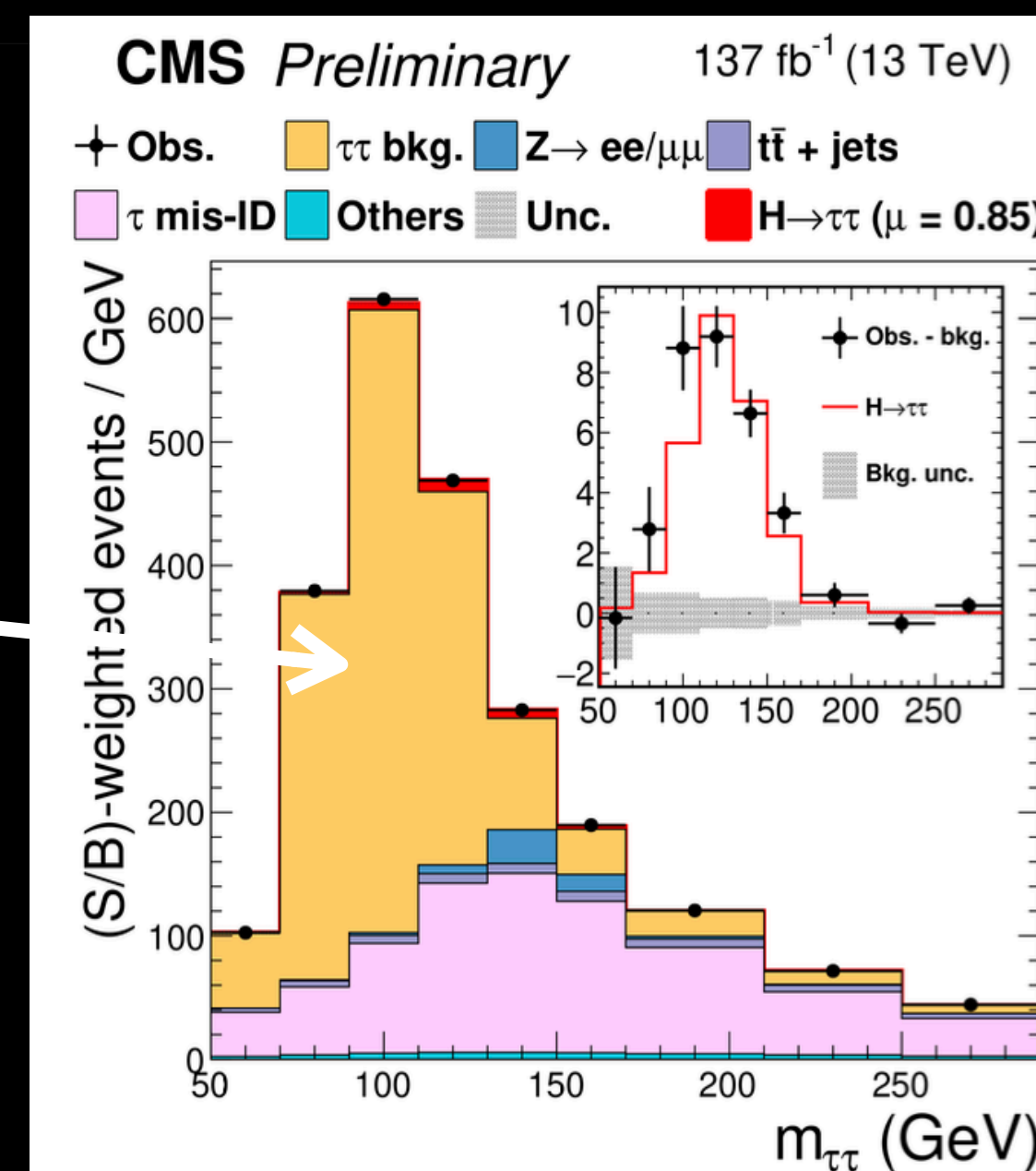


VBF process

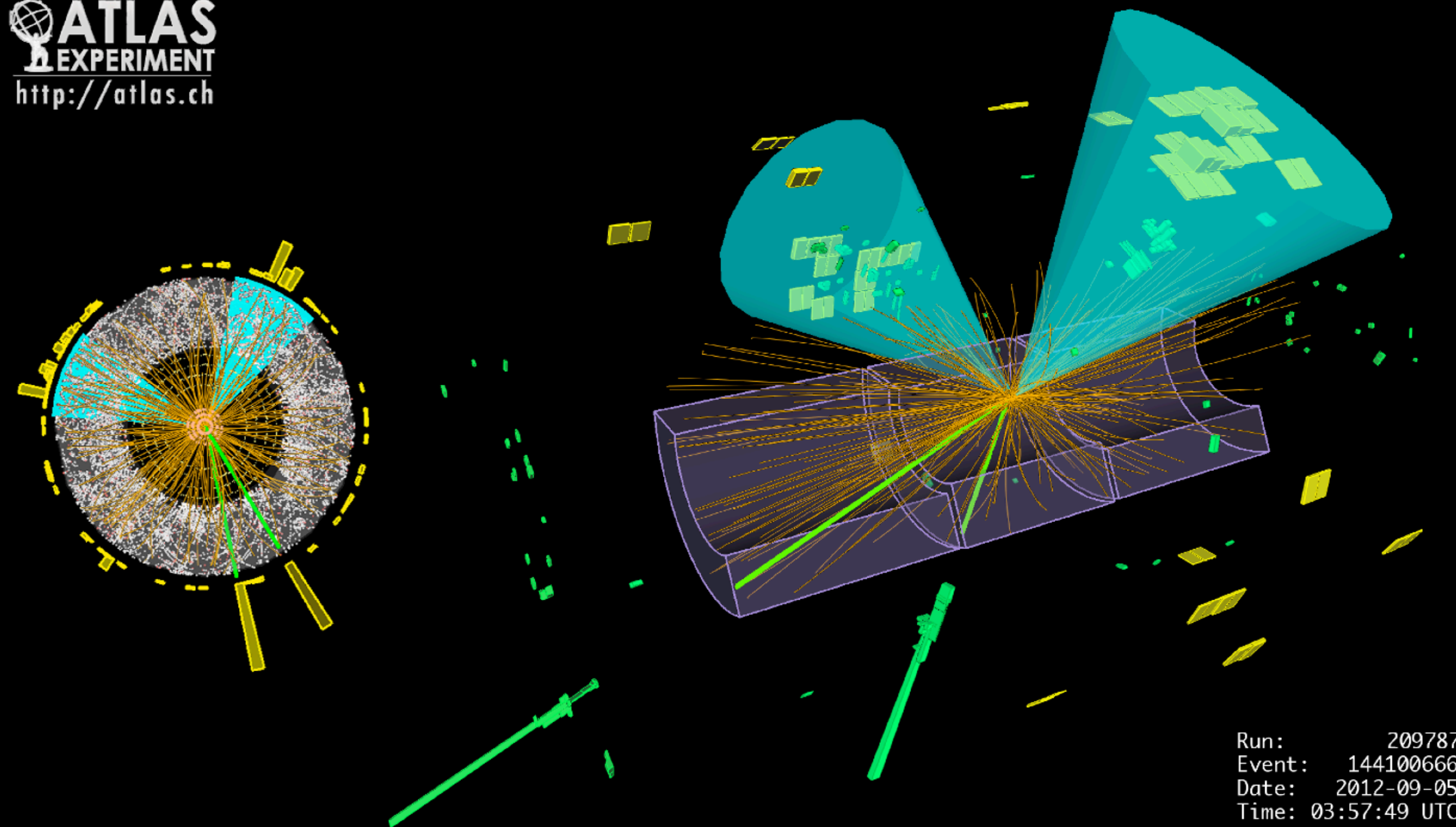


with two forward jets

background
from Z production
plus 2 jets

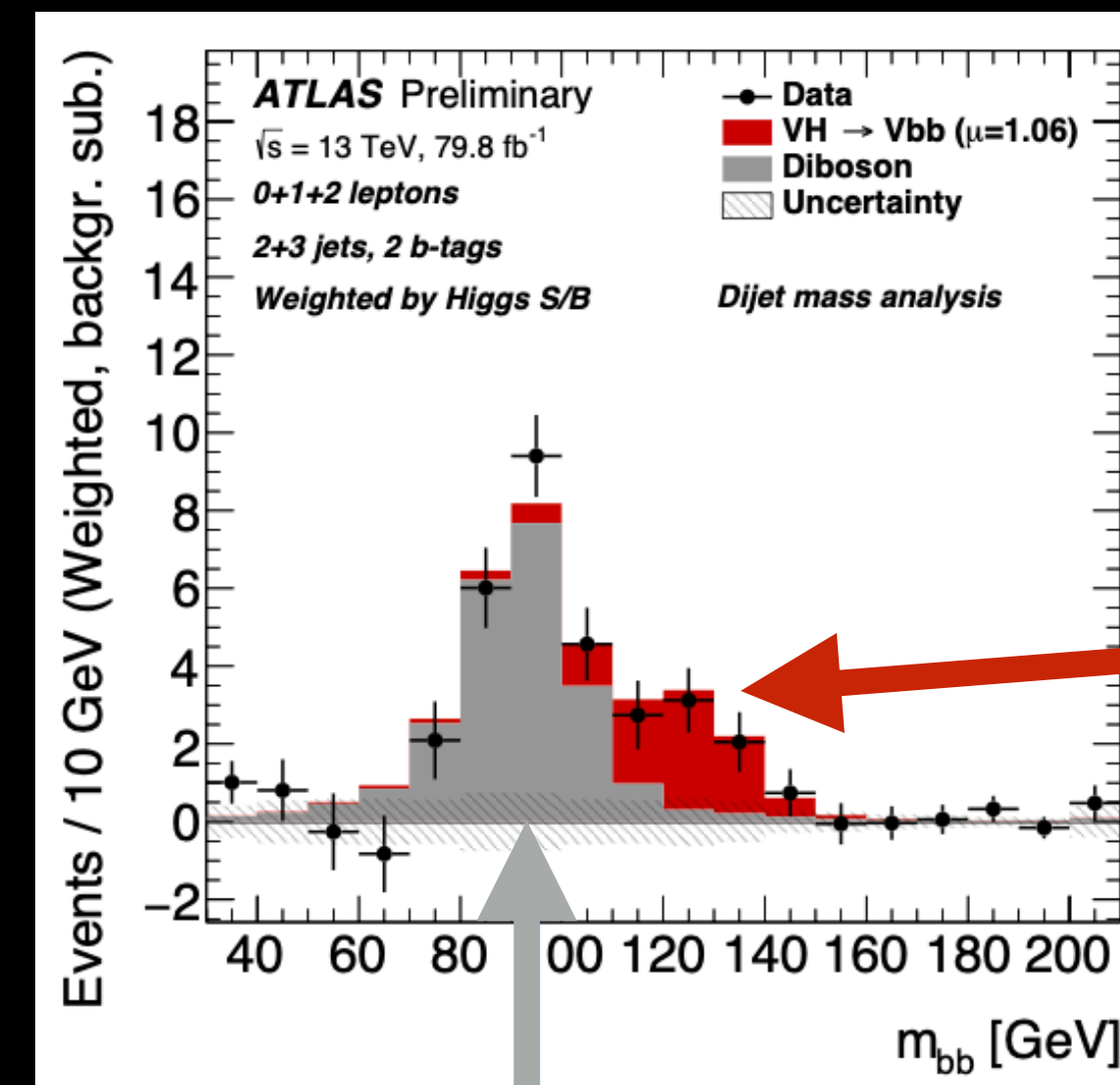
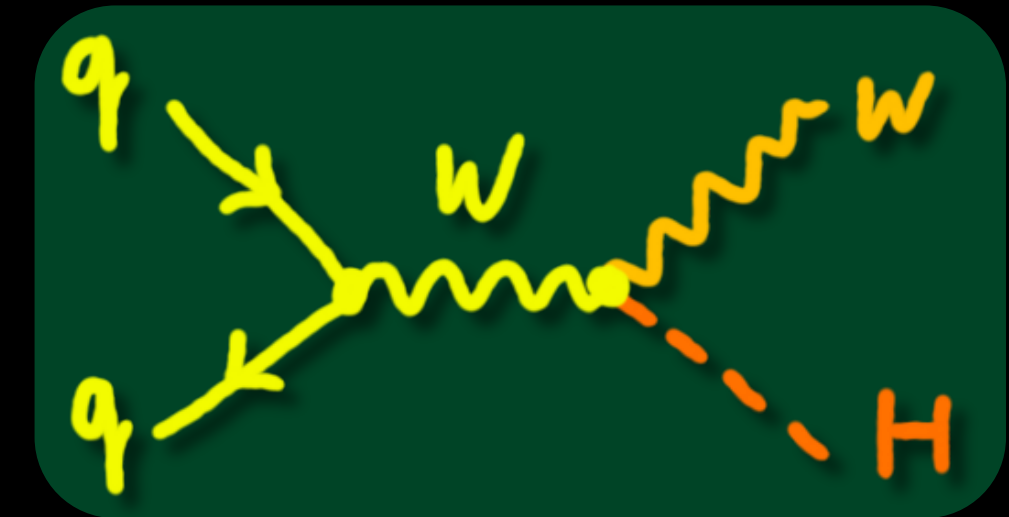


Higgs observation with b-quark decays



Run: 209787
Event: 144100666
Date: 2012-09-05
Time: 03:57:49 UTC

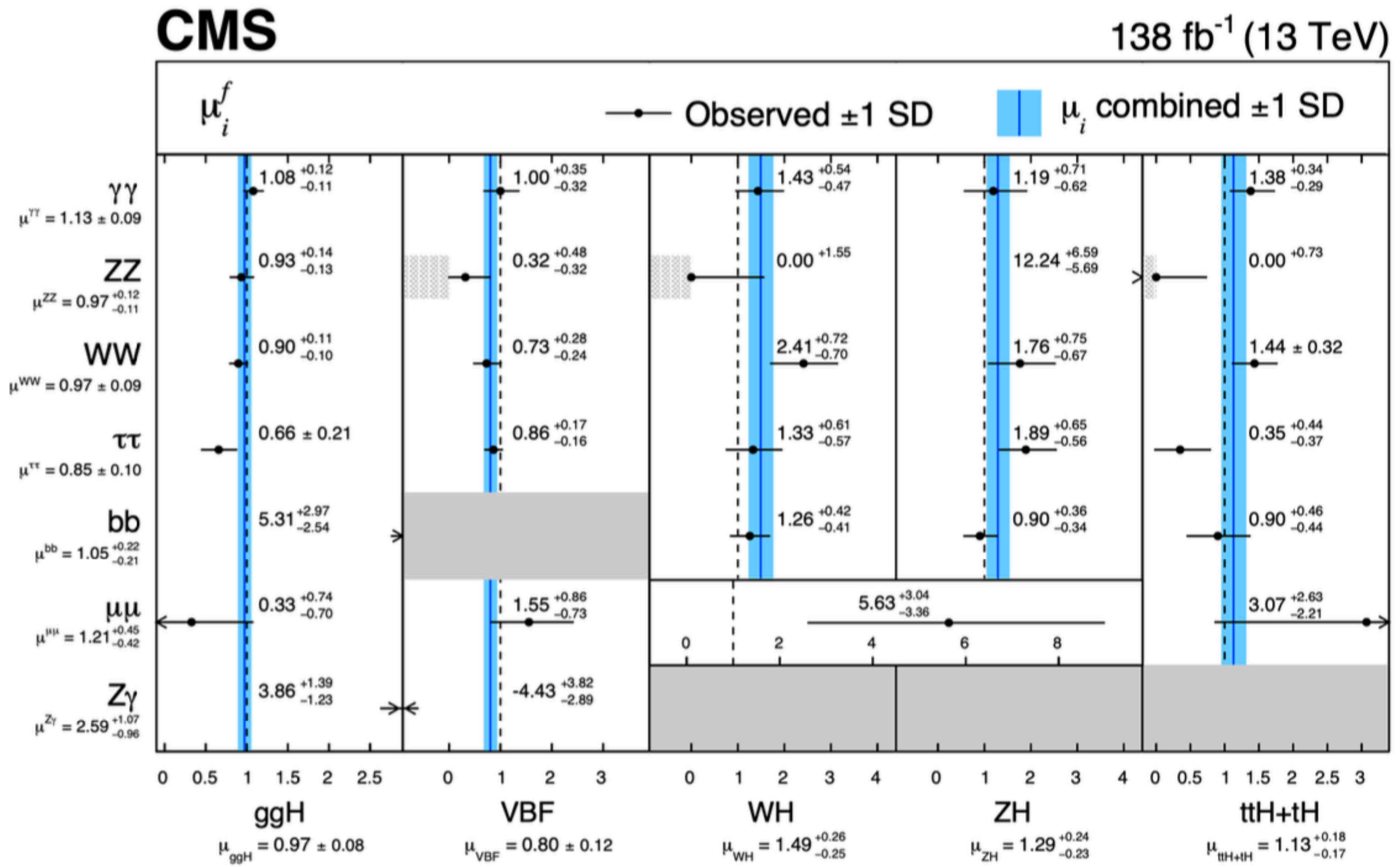
W/Z associate process



Higgs into
b-quarks

Z decays into b-quarks

Overview of the Higgs Boson exploration



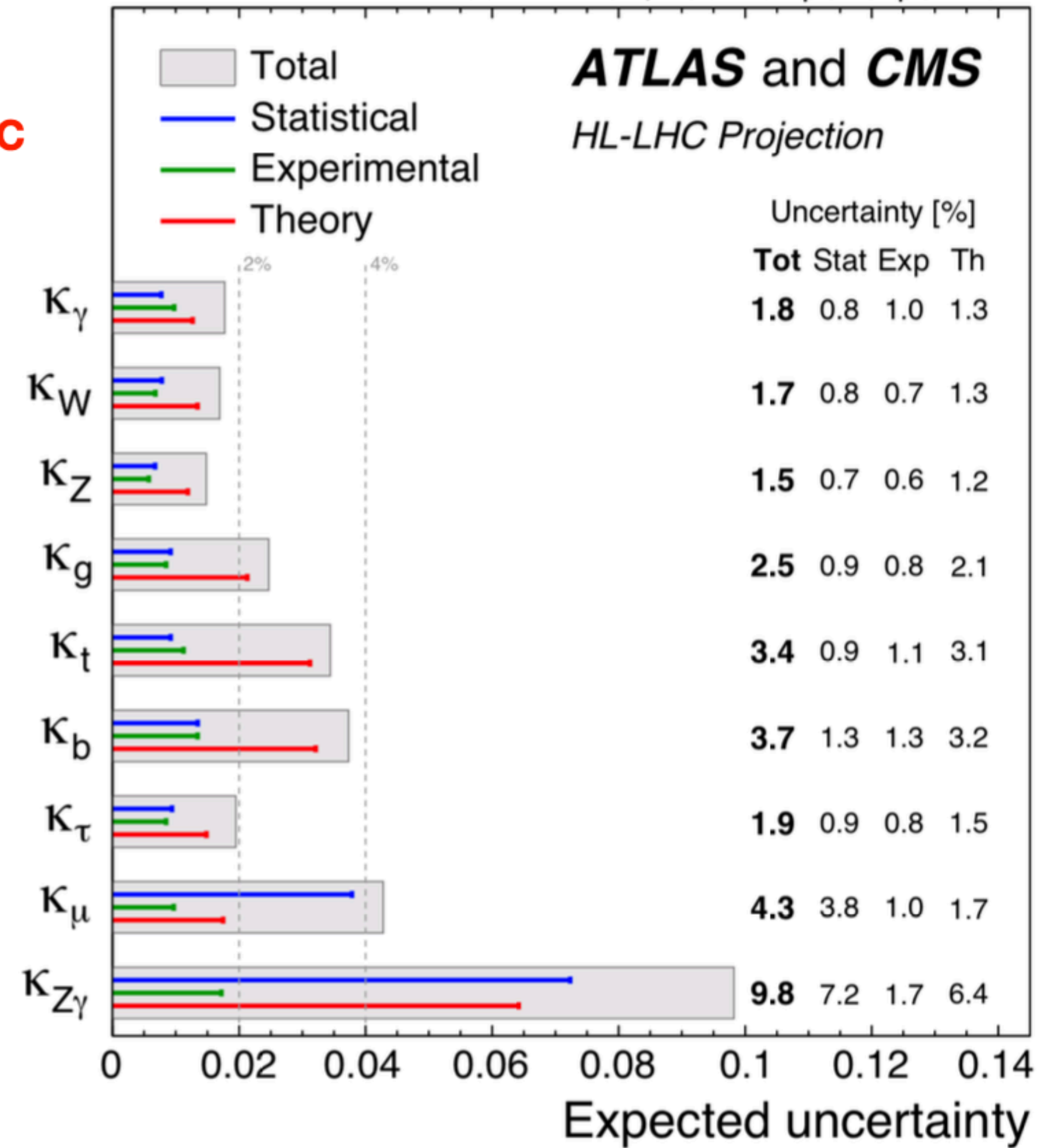
Higgs Boson Coupling Precision at LHC and HL-LHC

$\sqrt{s} = 14 \text{ TeV}, 3000 \text{ fb}^{-1}$ per experiment

	ATLAS - CMS Run 1 combination	ATLAS Run 2	CMS Run 2	Current precision	HL-LHC
κ_γ	13%	1.04 ± 0.06	1.10 ± 0.08	6%	1.8%
κ_W	11%	1.05 ± 0.06	1.02 ± 0.08	6%	1.7%
κ_Z	11%	0.99 ± 0.06	1.04 ± 0.07	6%	1.5%
κ_g	14%	0.95 ± 0.07	0.92 ± 0.08	7%	2.5%
κ_t	30%	0.94 ± 0.11	1.01 ± 0.11	11%	3.4%
κ_b	26%	0.89 ± 0.11	0.99 ± 0.16	11%	3.7%
κ_τ	15%	0.93 ± 0.07	0.92 ± 0.08	8%	1.9%
κ_μ	-	$1.06^{+0.25}_{-0.30}$	1.12 ± 0.21	20%	4.3%
$\kappa_{Z\gamma}$	-	$1.38^{0.31}_{-0.36}$	1.65 ± 0.34	30%	9.8%
B_{inv}		< 11 %	< 16 %	11%	2.5%

Nature 607, 52-59 (2022)

Nature 607, 60-68 (2022)

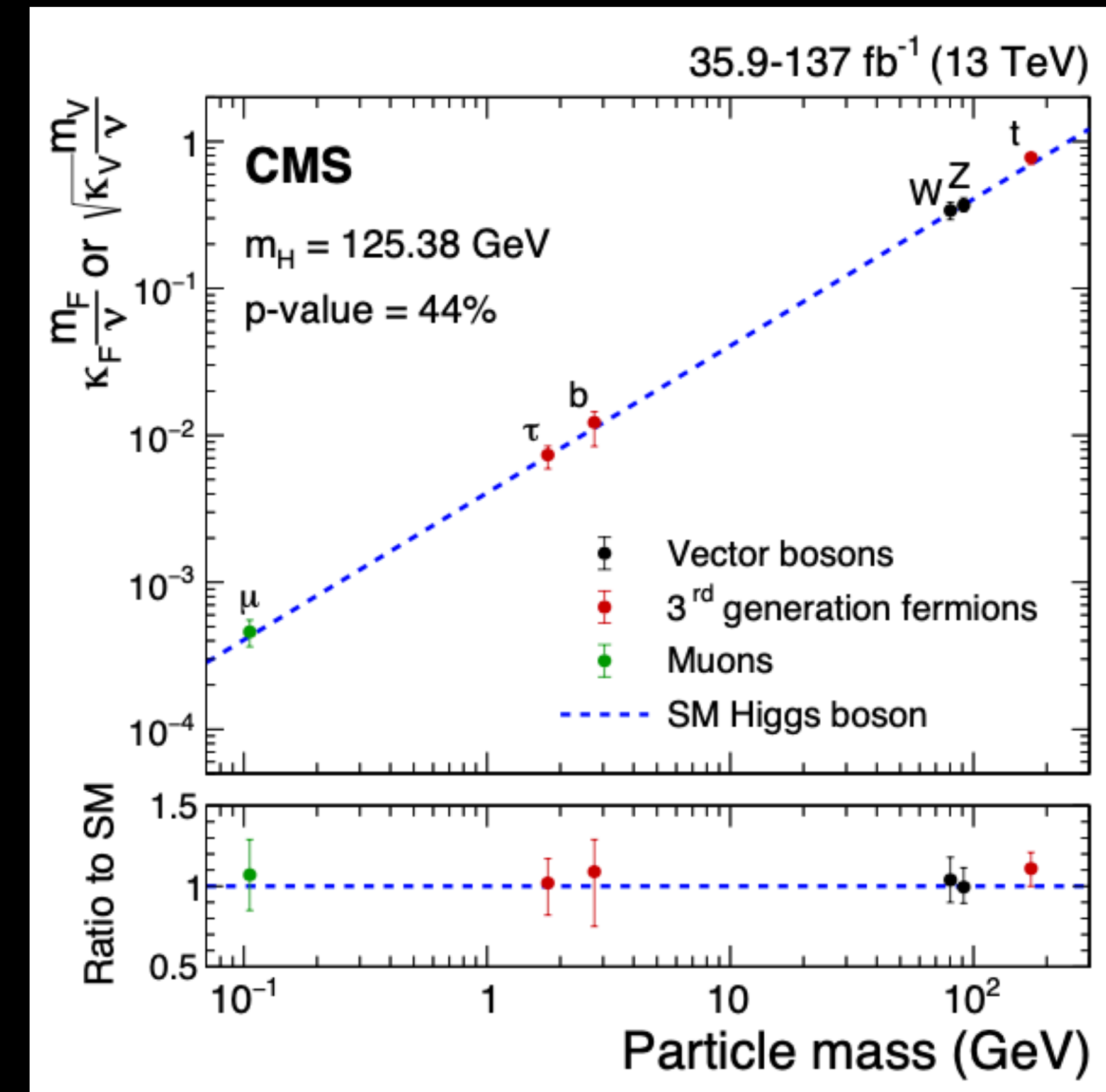


TH Uncertainties dominant (assumed to be 1/2 of Run 2)

What we know about the Higgs

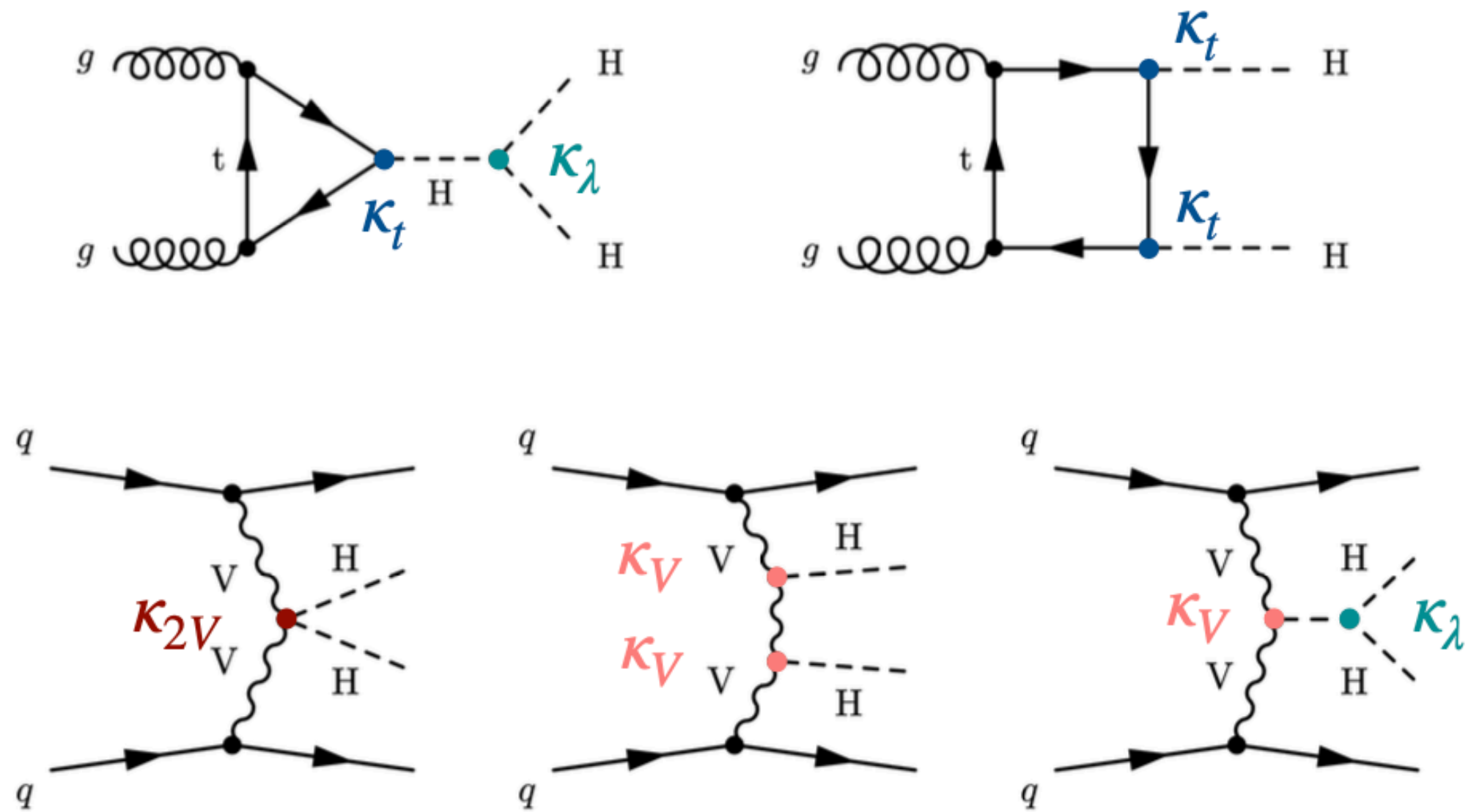
- Gives mass to the **W** and **Z** bosons
- Gives mass to **3rd generation fermions** via Yukawa interaction
- Evidence it couples to the **2nd generation fermions** as expected
- Has **spin 0**
- Not composite (down to 10^{-19} m)

JHEP01(2021)148



Di-Higgs Production and Higgs Self Coupling

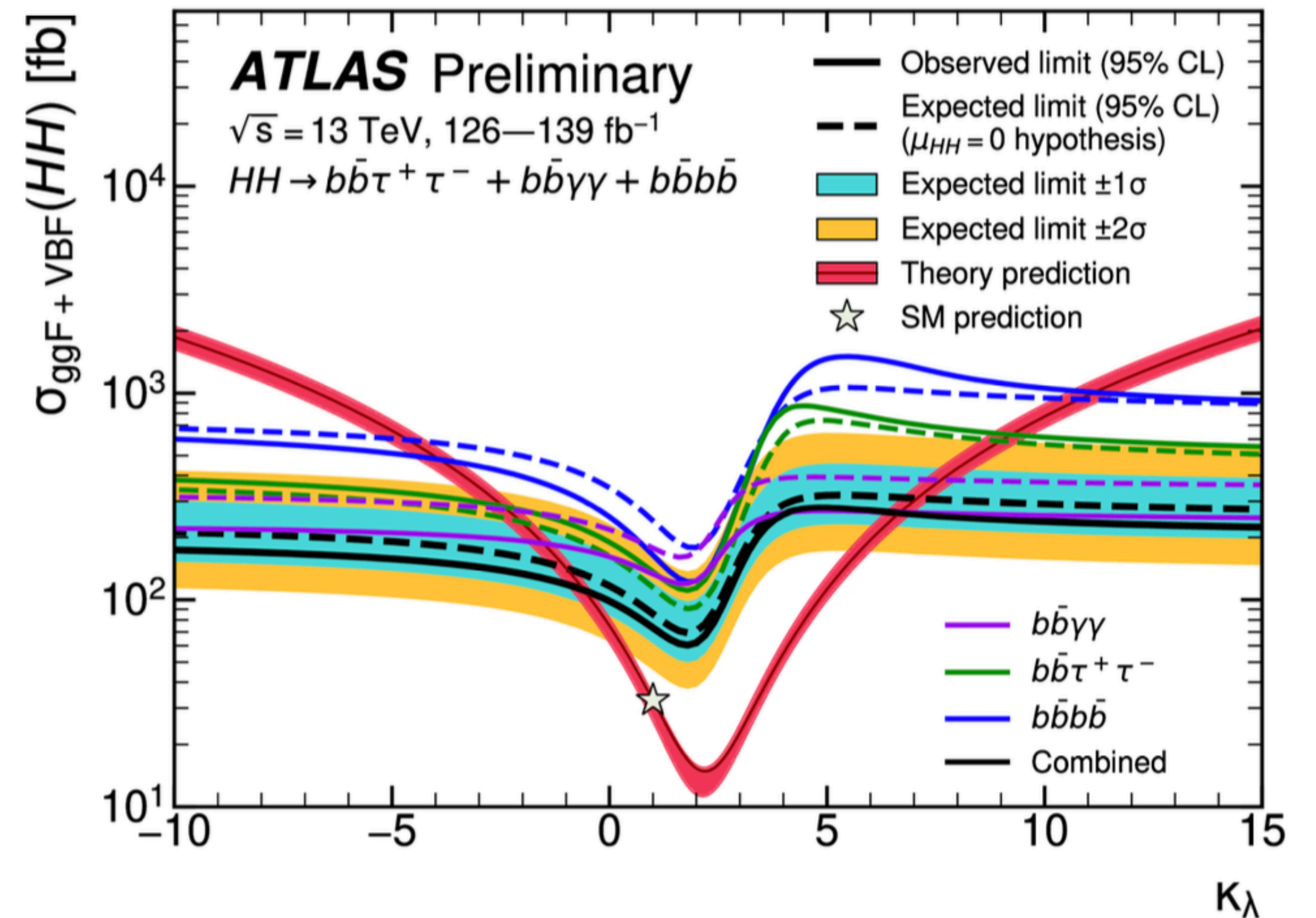
Higgs pair production through gluon fusion and VBF



Cross section $\sim 1000\times$ smaller than Higgs production still

Expect $\sim 100k$ events produced at HL-LHC

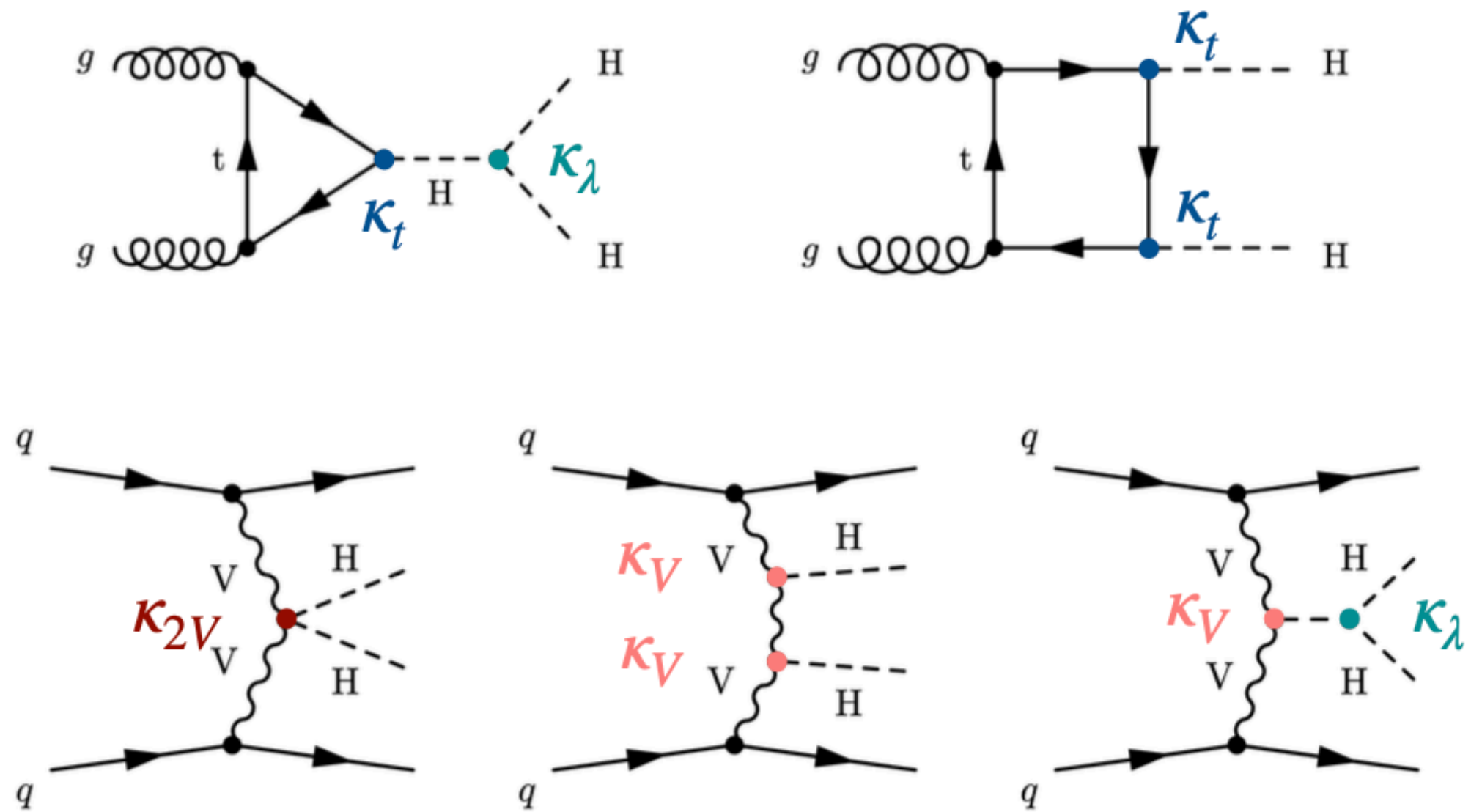
Multiple channels being investigated to maximize sensitivity: $bb, \gamma\gamma, \tau\tau, WW$



ATLAS
 Observed $-0.4 < \kappa_\lambda < 6.3$
 Expected $-1.9 < \kappa_\lambda < 7.5$

Di-Higgs Production and Higgs Self Coupling

Higgs pair production through gluon fusion and VBF

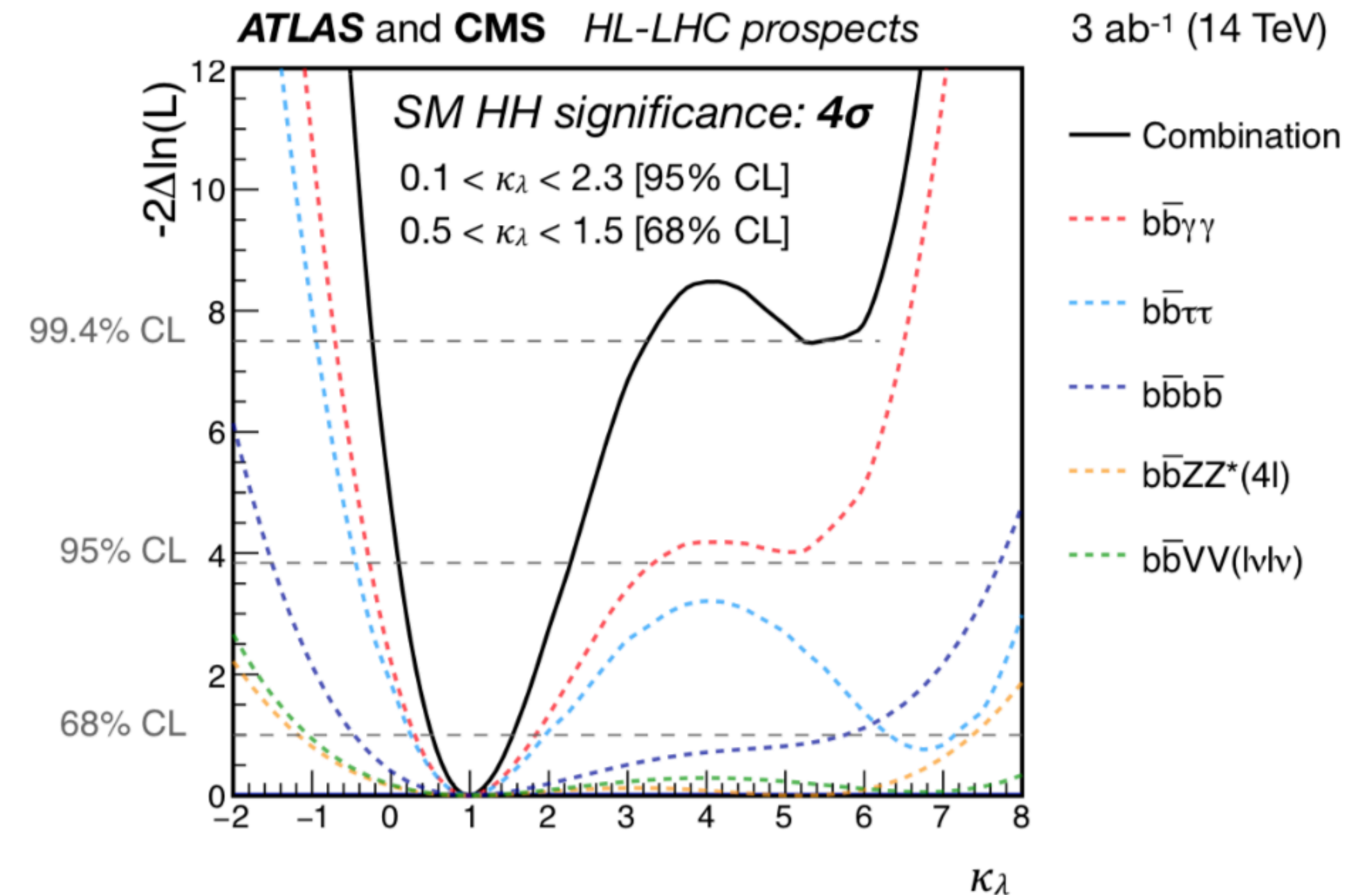


Cross section $\sim 1000\times$ smaller than Higgs production still

Expect $\sim 100k$ events produced at HL-LHC

Multiple channels being investigated to maximize sensitivity: $bb, \gamma\gamma, \tau\tau, WW$

Prospects for HL-LHC



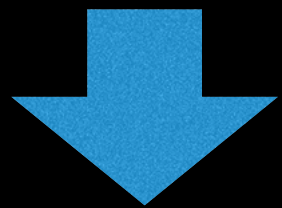
Possible observation of HH signal at 5σ constraint on the Higgs self coupling of $0.5 < \kappa_\lambda < 1.5$

Searches

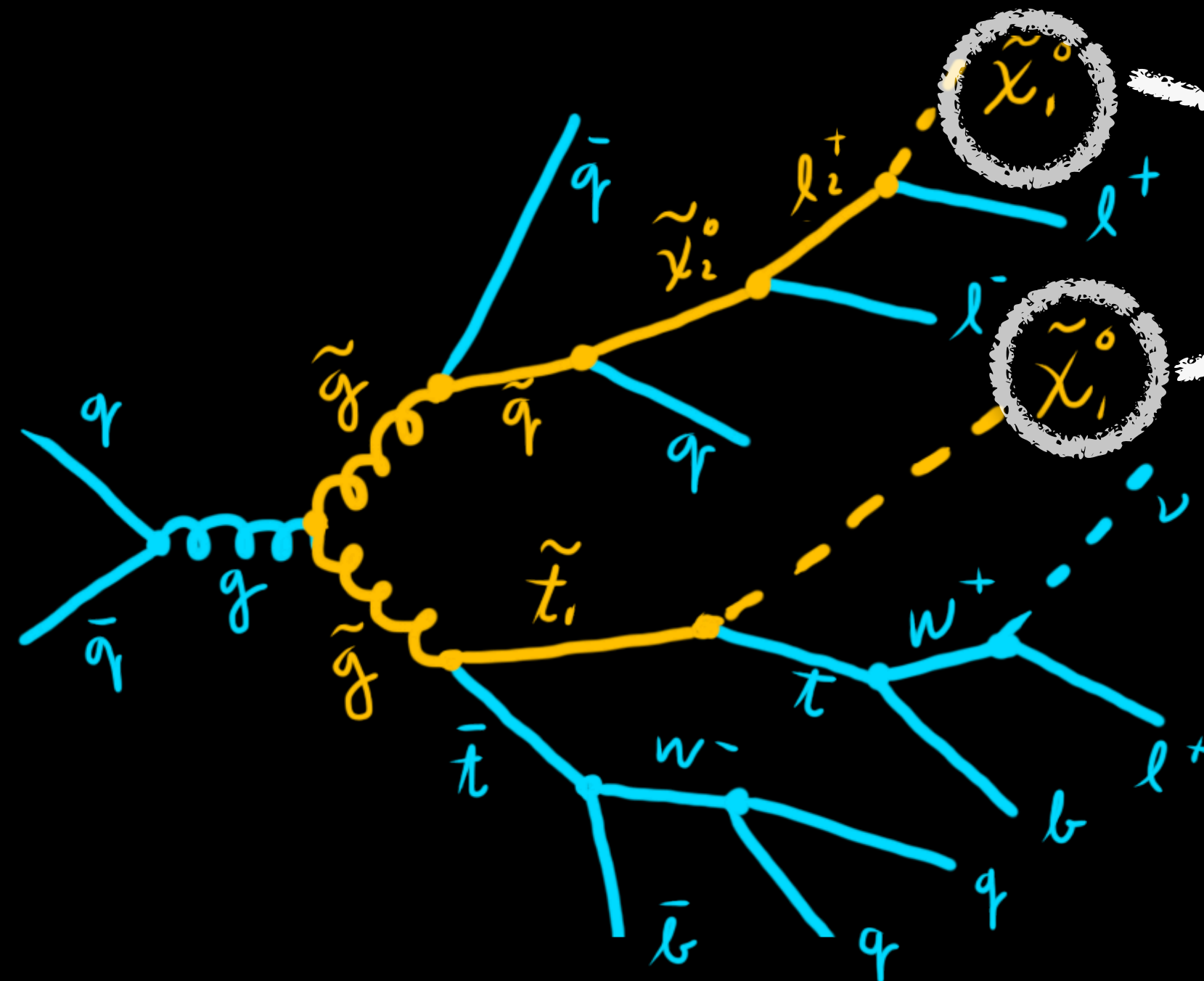
SUPERSYMMETRY (SUSY)

- **A favorite, ideal, long-time candidate to explain most of questions raised**
 - Symmetry between bosons and fermions, provides dark matter candidate, provides unification of forces, no fine-tuning required for Higgs mass
- **Unfortunately, no evidence for SUSY found yet**

Strong SUSY production



Large cross section



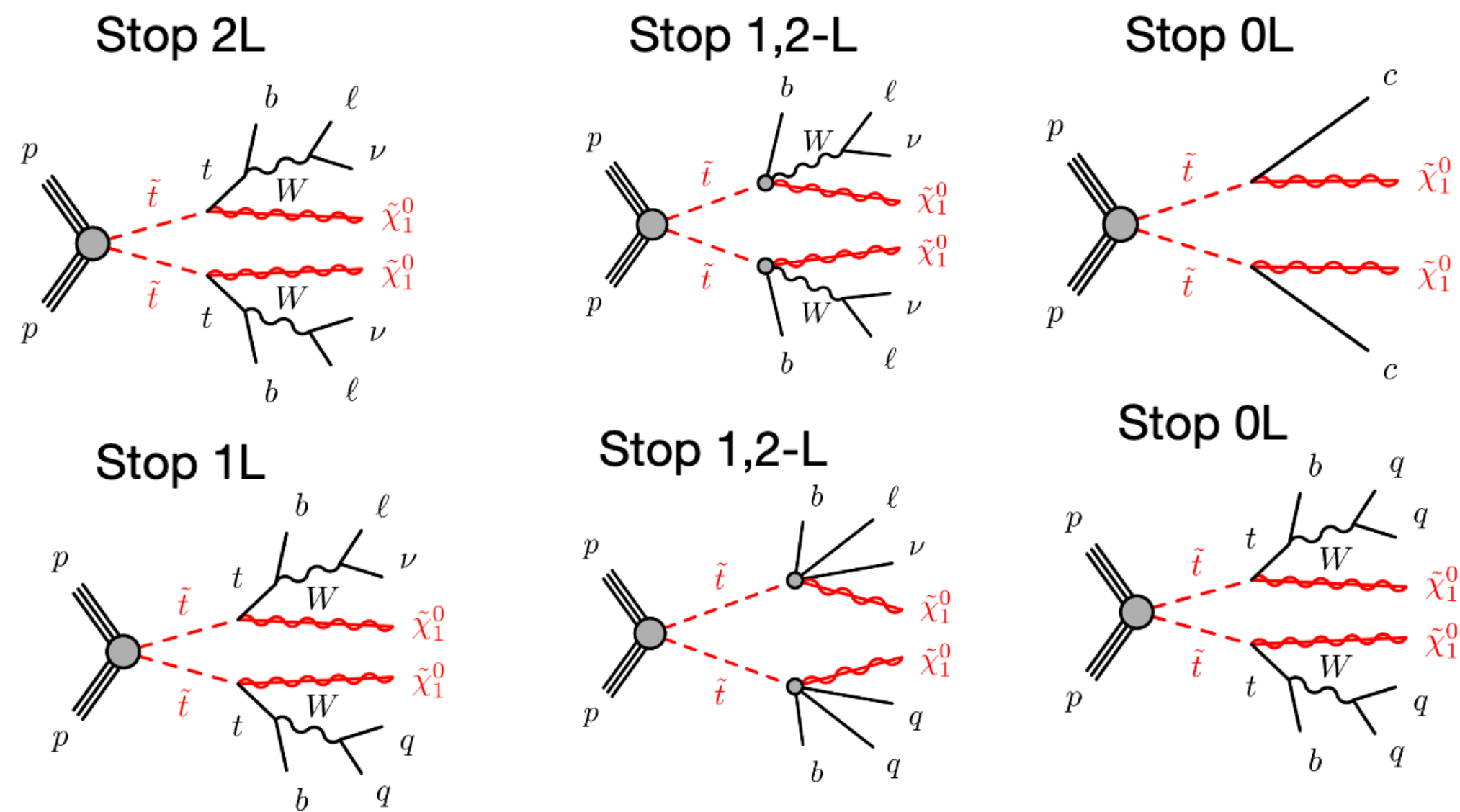
Dark Matter Candidates

(Assuming R-Parity Conservation)

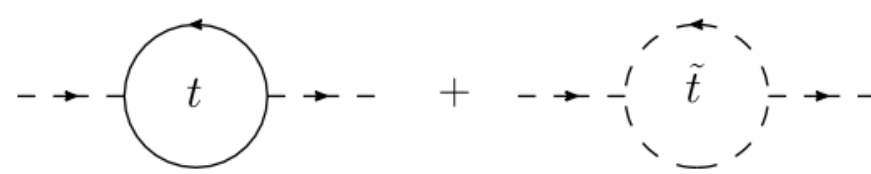
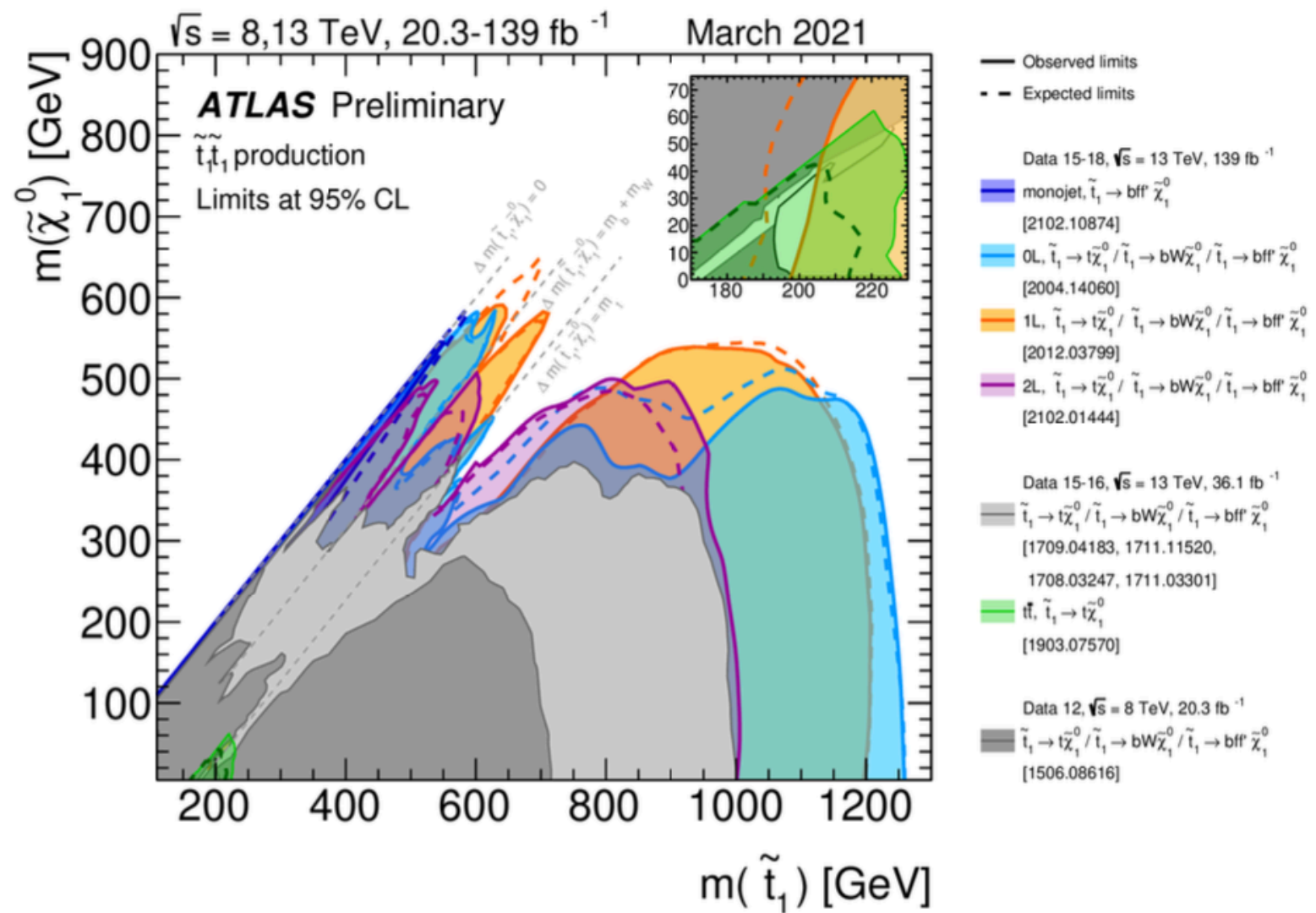
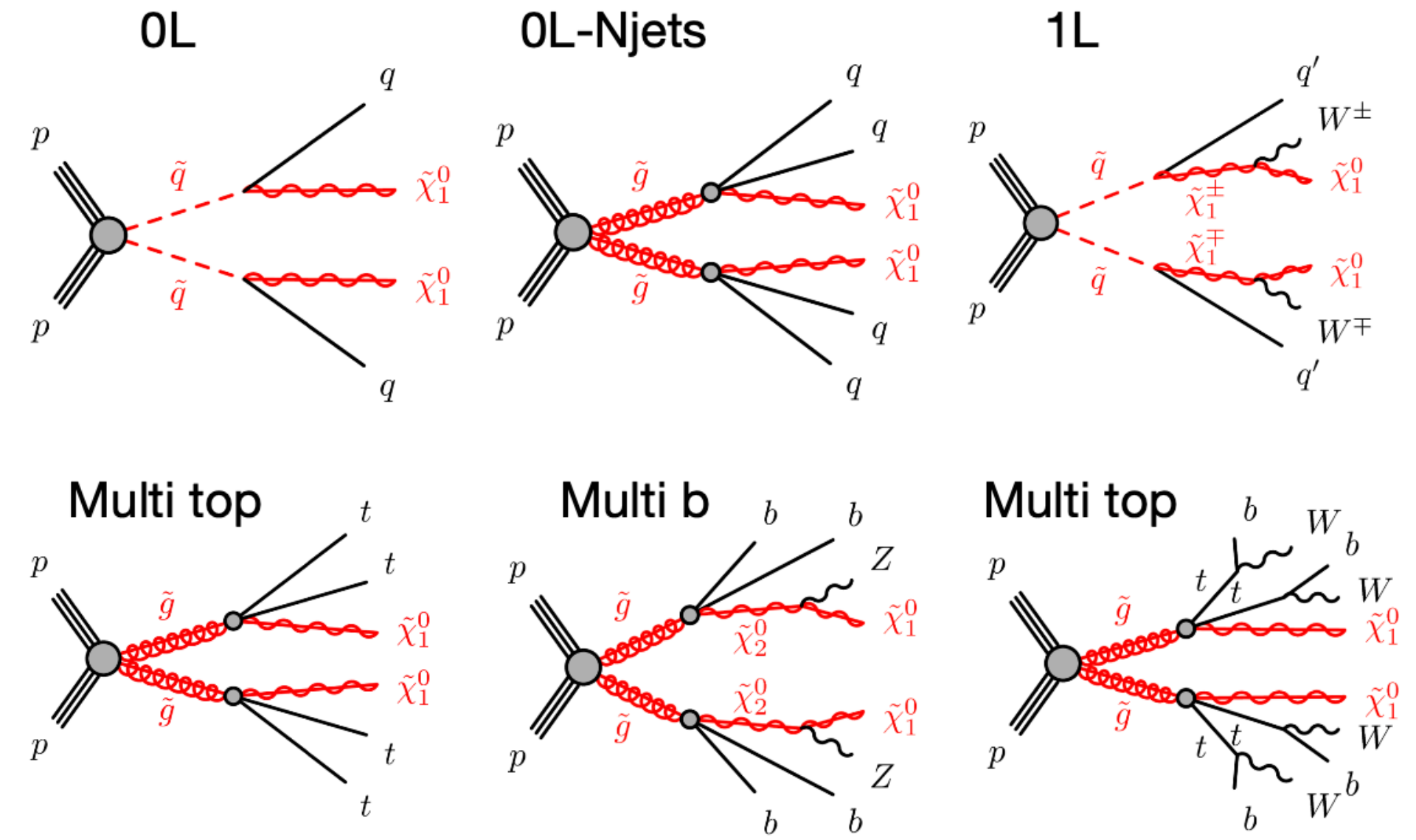
Mass gluino > 2 TeV (depending on models)

Searches for Natural and Strongly Produced SUSY

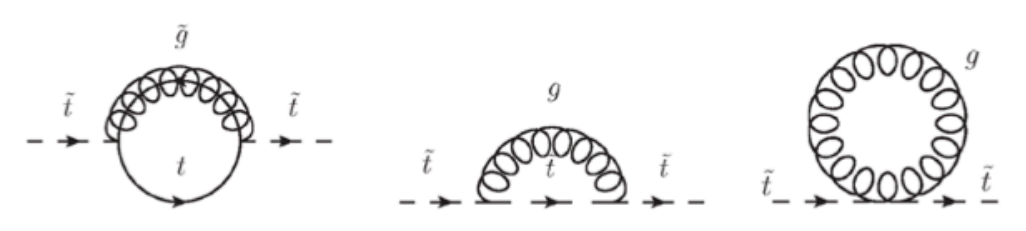
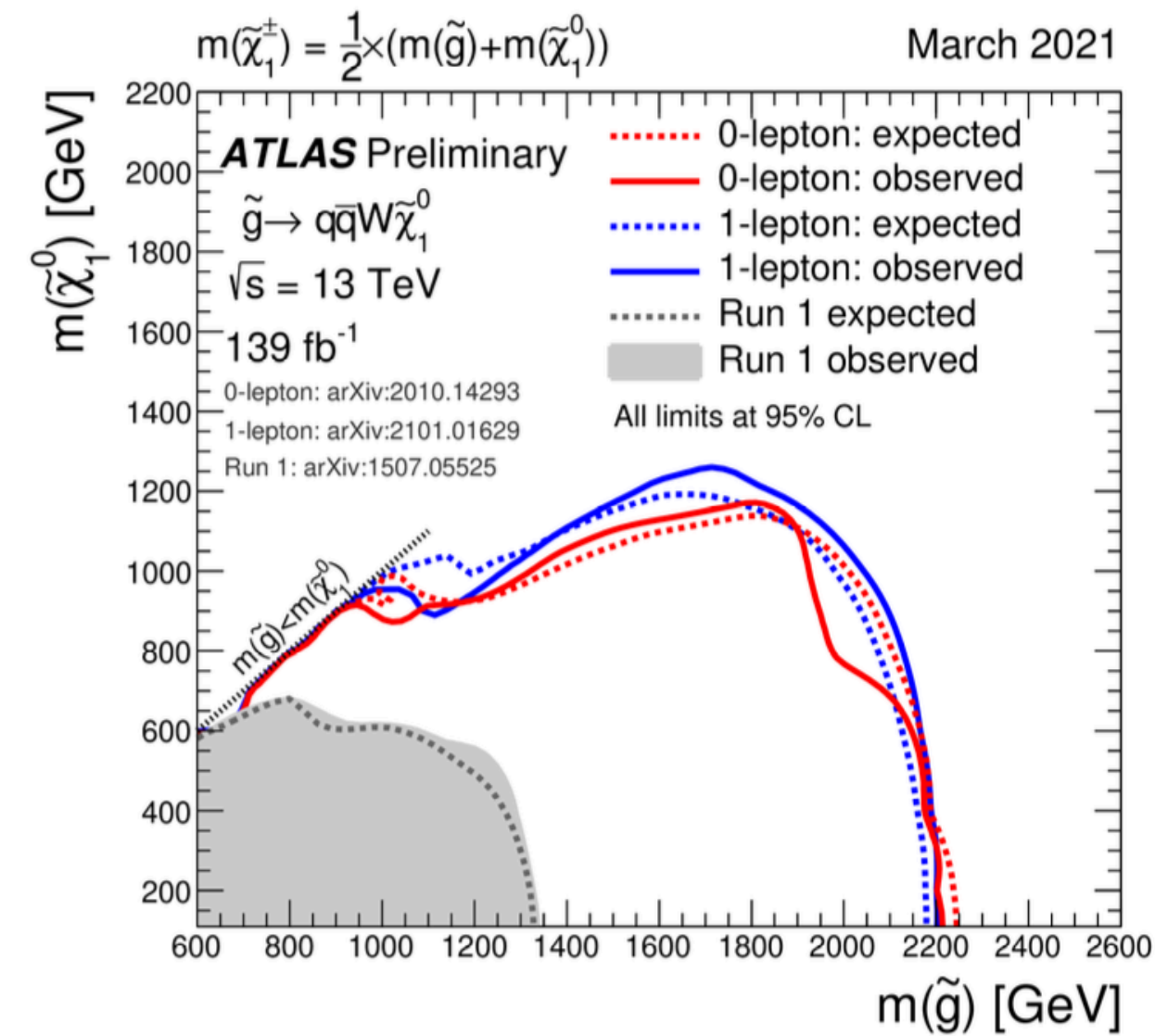
Stop



Squarks and gluinos



Not so natural SUSY:
 Stops > 1 TeV ~Tuning of factor 20, but these exclusions are under specific conditions, and there are unexcluded corridors.

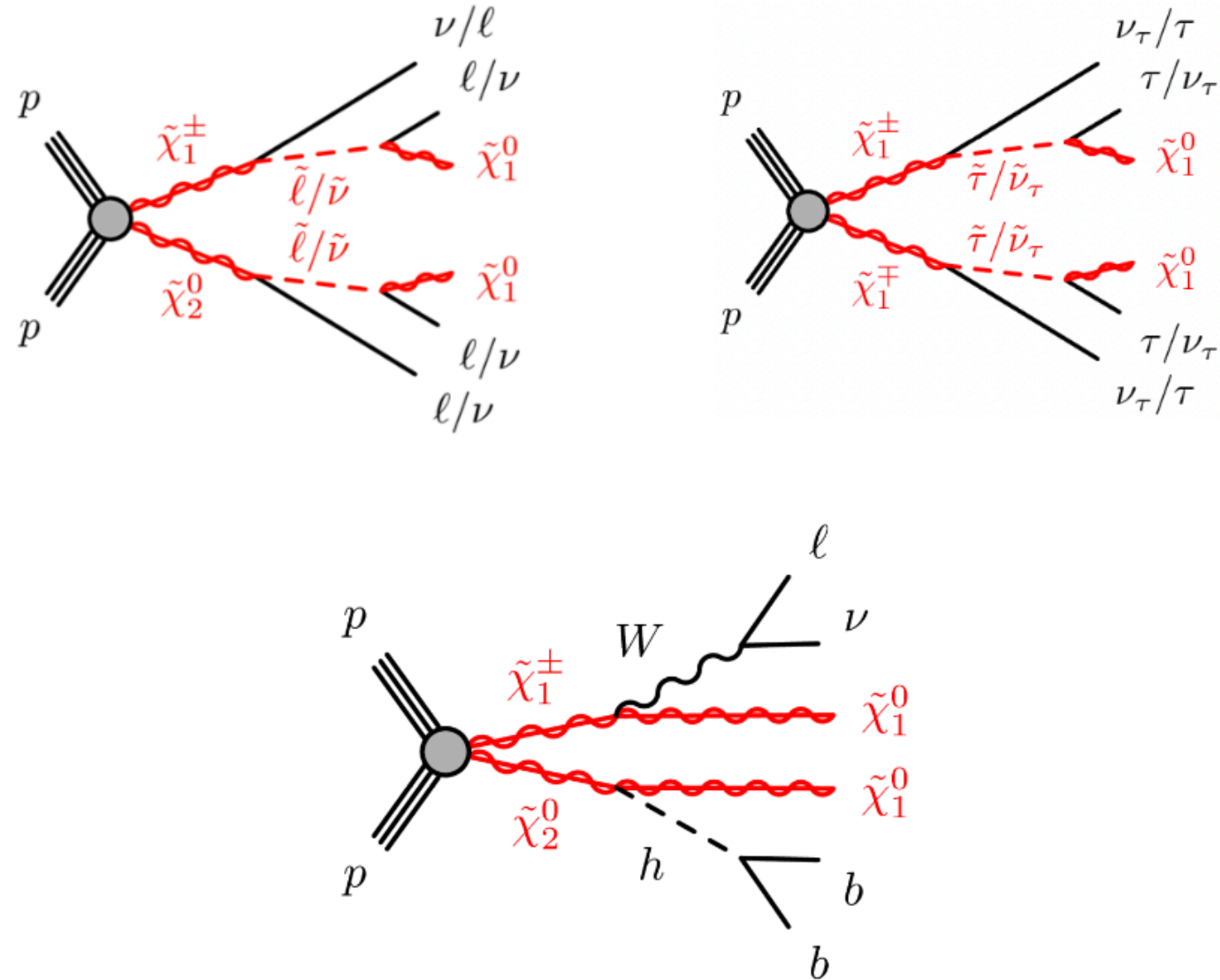


Stop also a scalar requires light gluinos to be light enough: for gluinos > 2 TeV ~tuning of Factor of 30

Searches for Weakly Coupled and Complex Scenarios

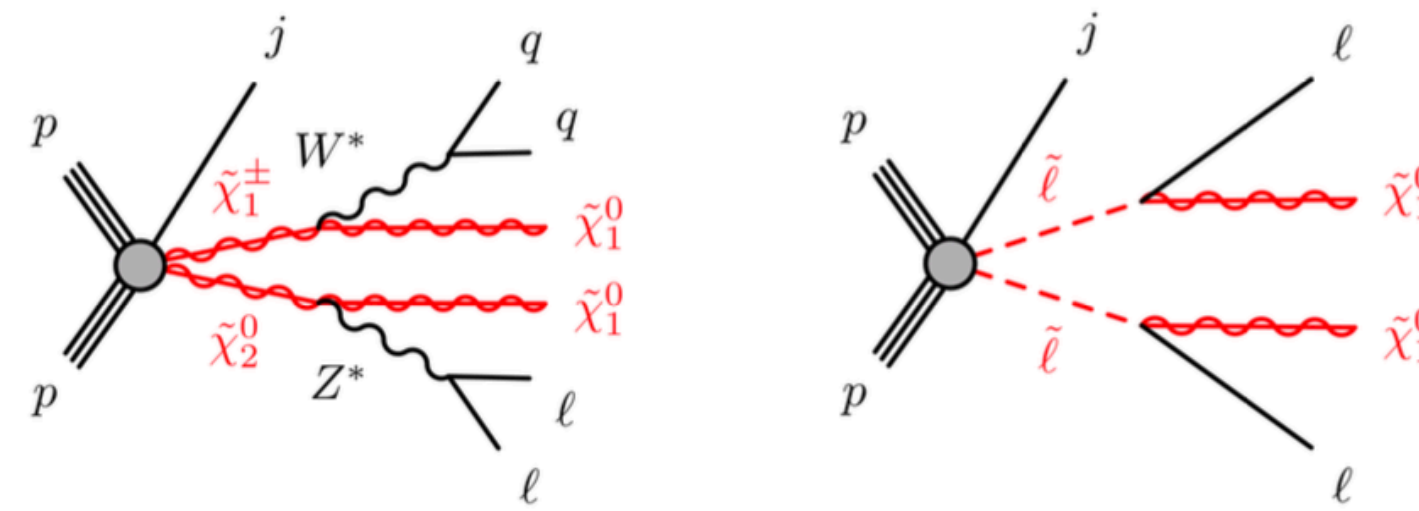
Weakly produced SUSY

Weak production of charginos, neutralinos and sleptons



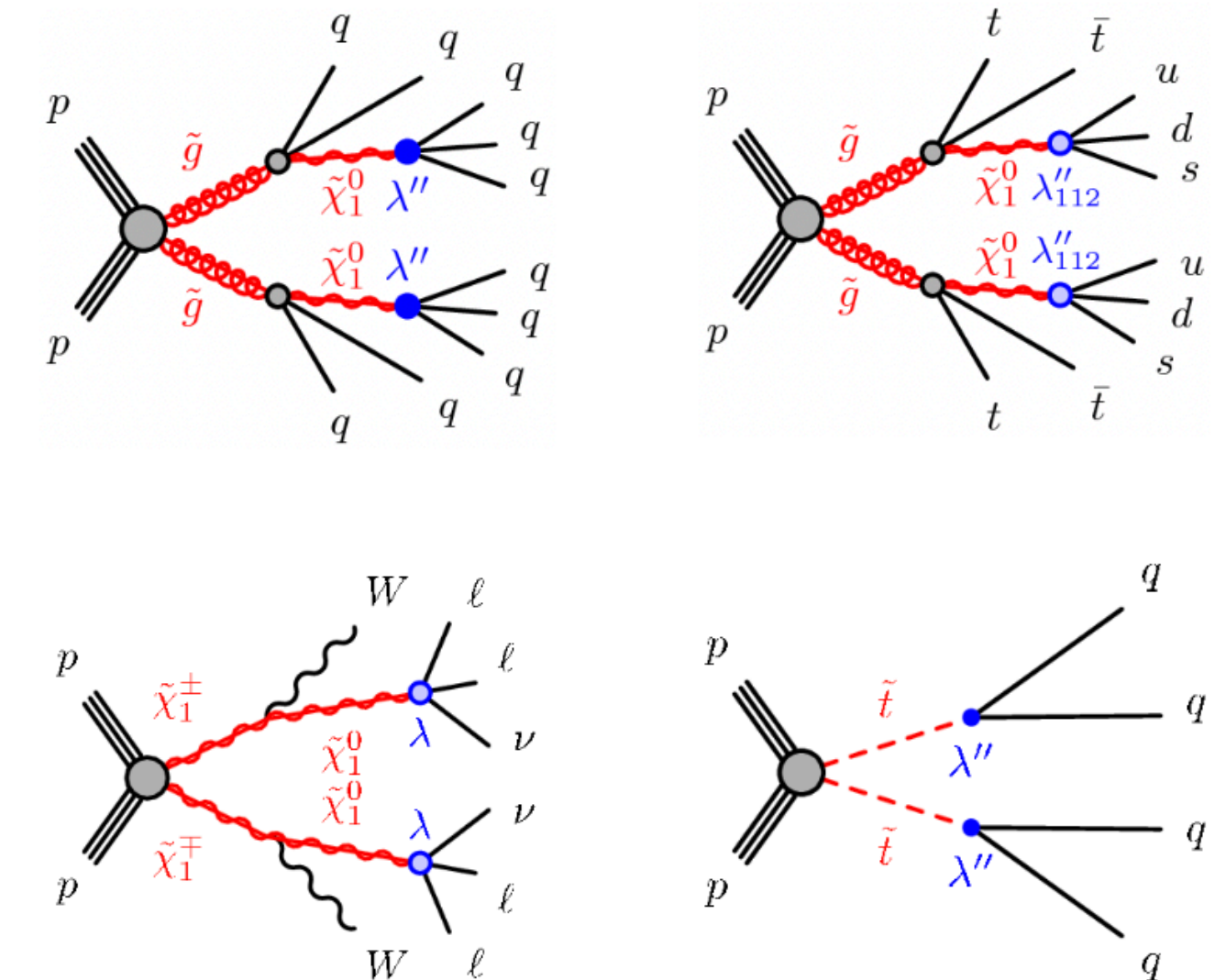
1 to 4 leptons (including taus) in the final state. Including decays to electroweak bosons.

Weak production in compressed scenarios



Scenarios where the charginos, neutralinos or sleptons are close to mass degenerate with the lightest SUSY particle (LSP).

R-Parity violating SUSY



Resulting in topologies without LSP in the final state and therefore no MET.

Rare or complex signatures \rightarrow more difficult to observe

SUSY Searches Outcomes:

Caveats to be aware of:

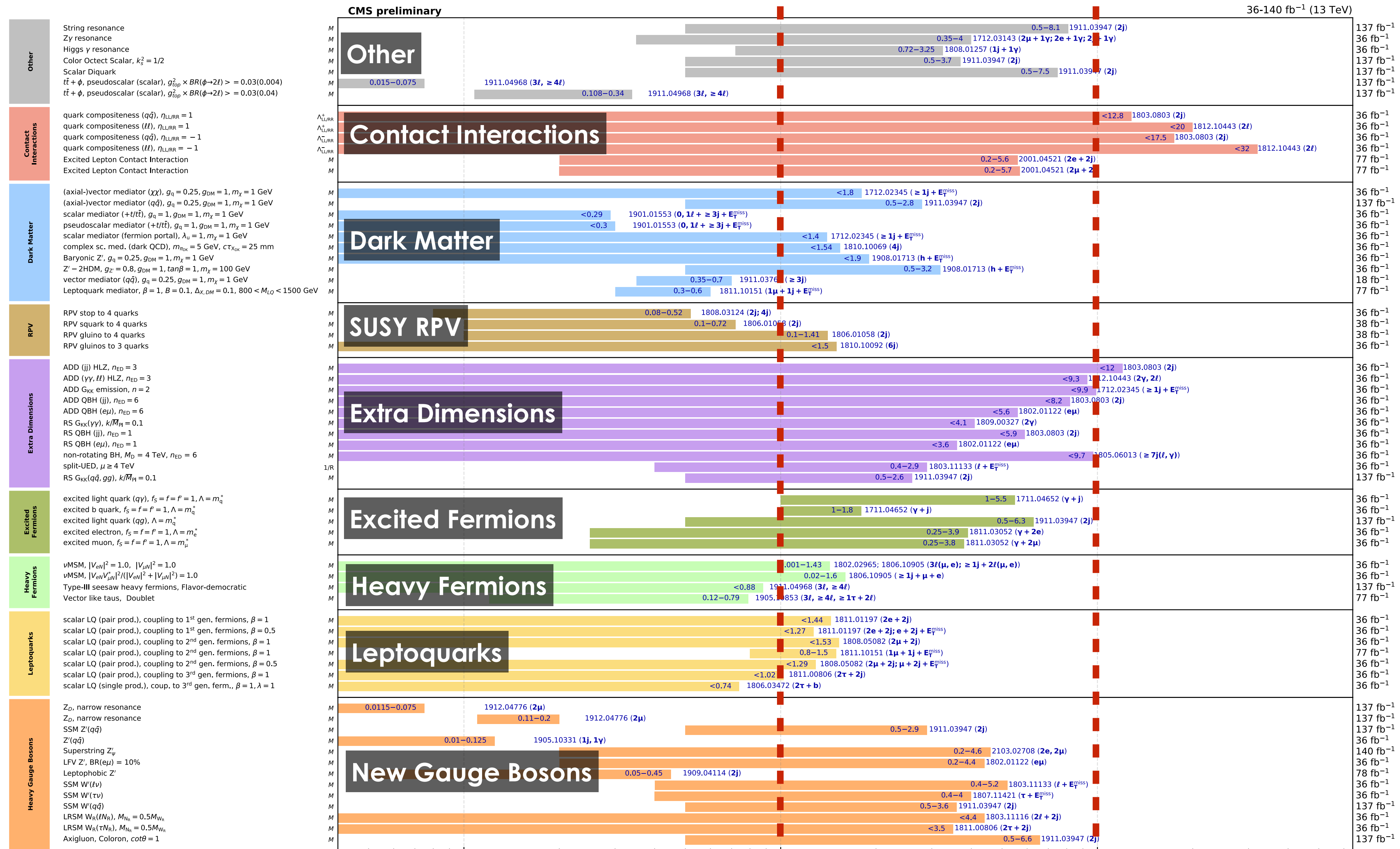
- Searches typically assume Branching Fractions of 100% into a specific channel
- Many searches assume mass degeneracy between various SUSY particles, e.g., squarks of different generation
- Interpretation is usually done via simplified model framework, not in the full model

Limits might be less stringent than they seem to be

DIRECT NEW PHYSICS SEARCHES FROM LHC

Observed exclusion limits at 95%

Overview of CMS EXO results



Selection of observed exclusion limits at 95% C.L. (theory uncertainties are not included).

0.1 TeV

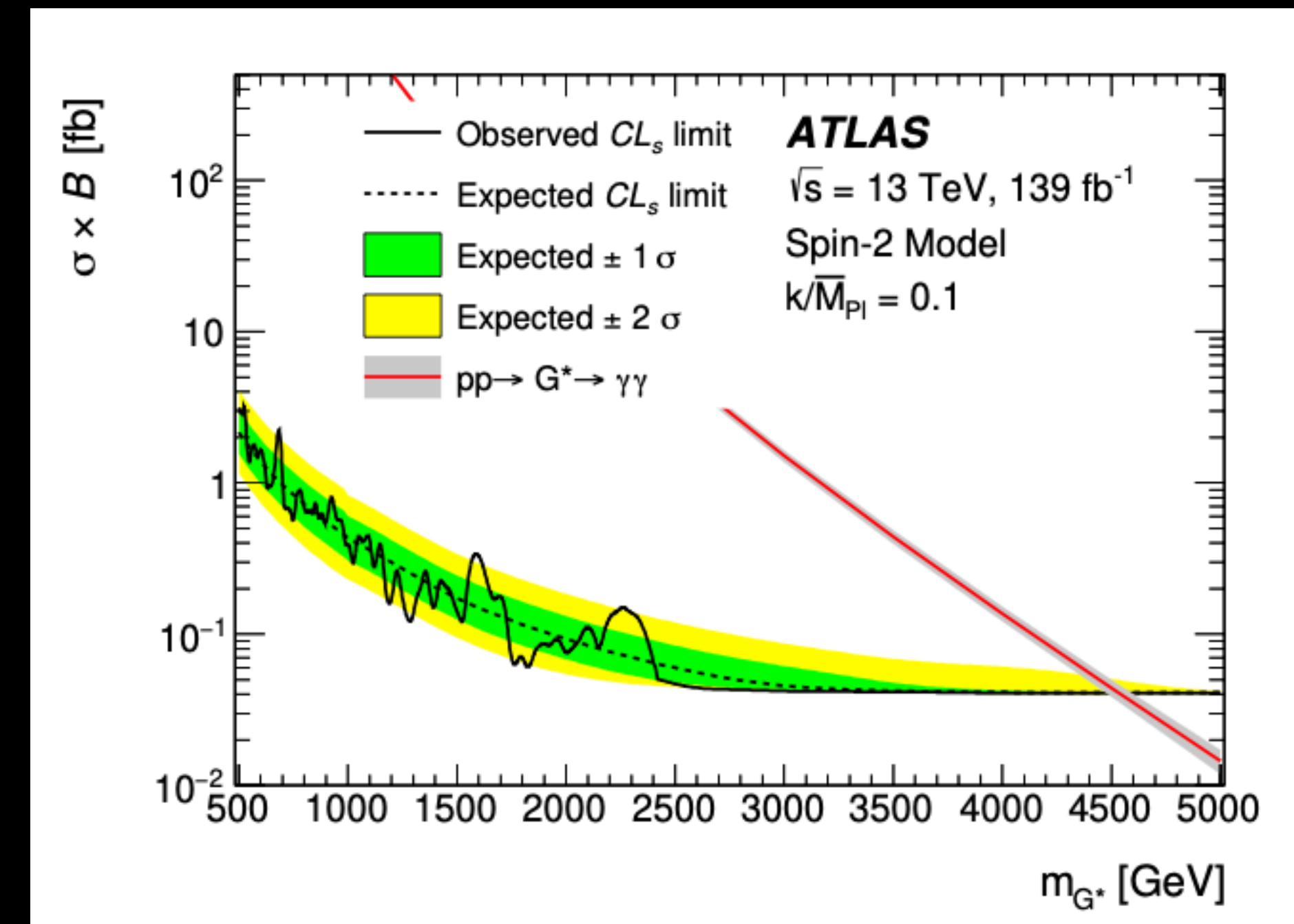
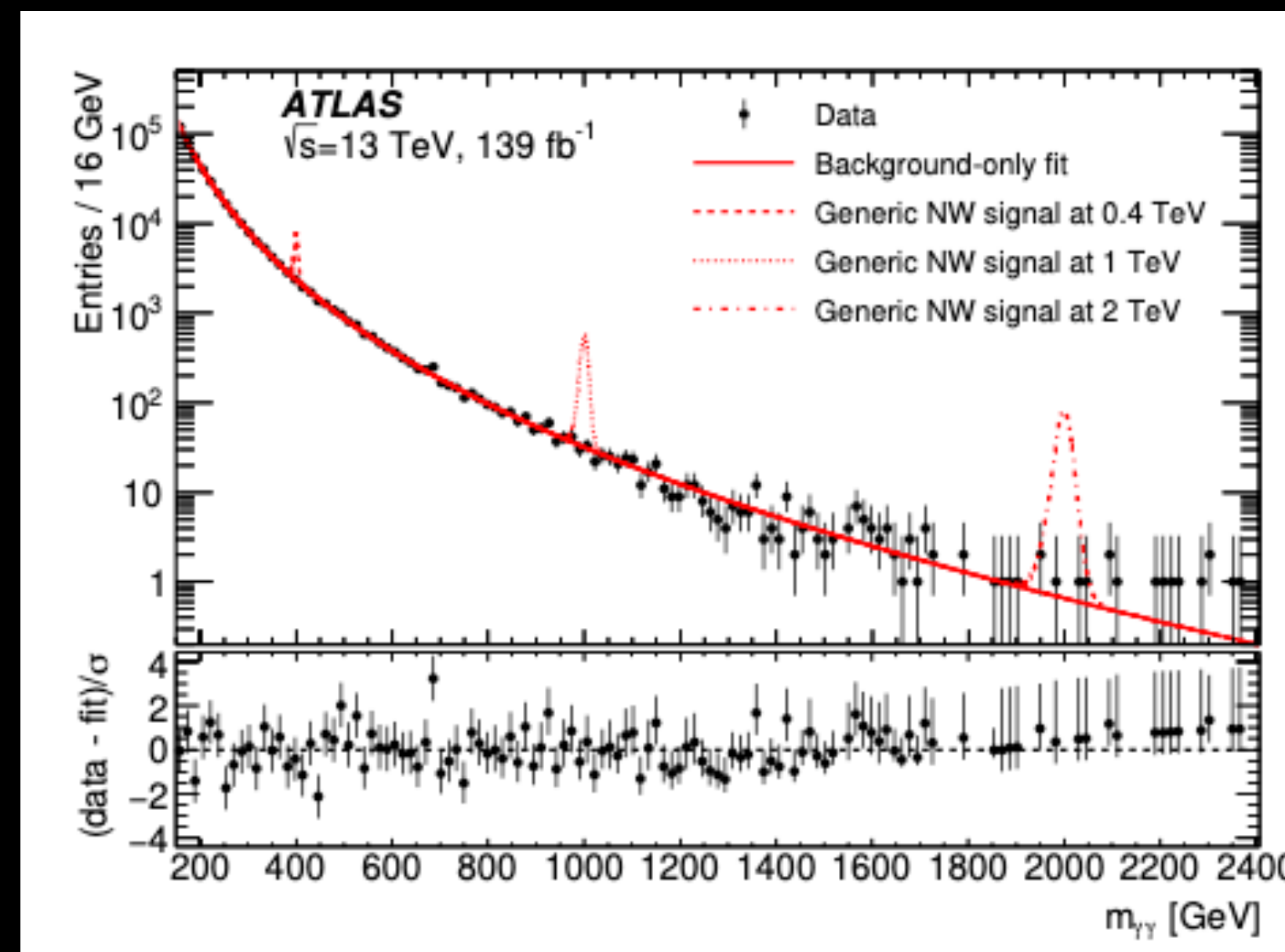
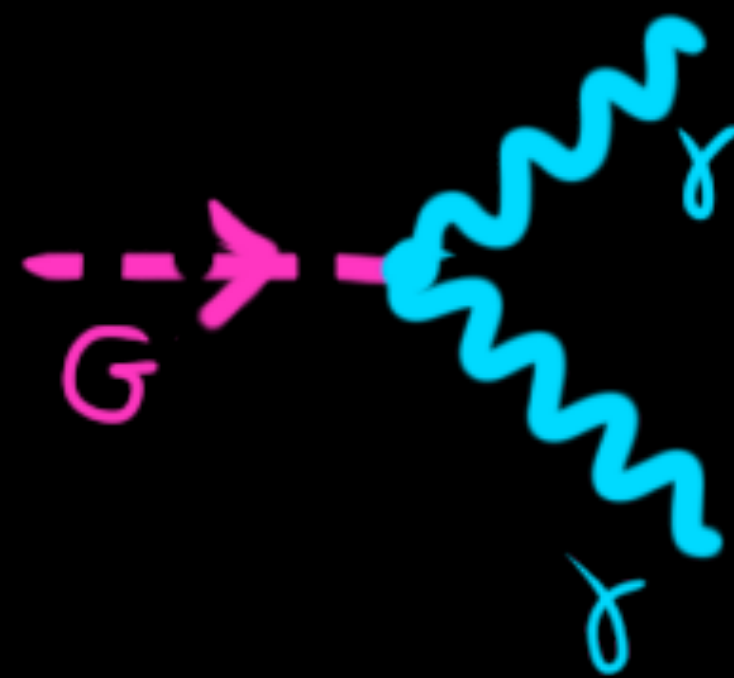
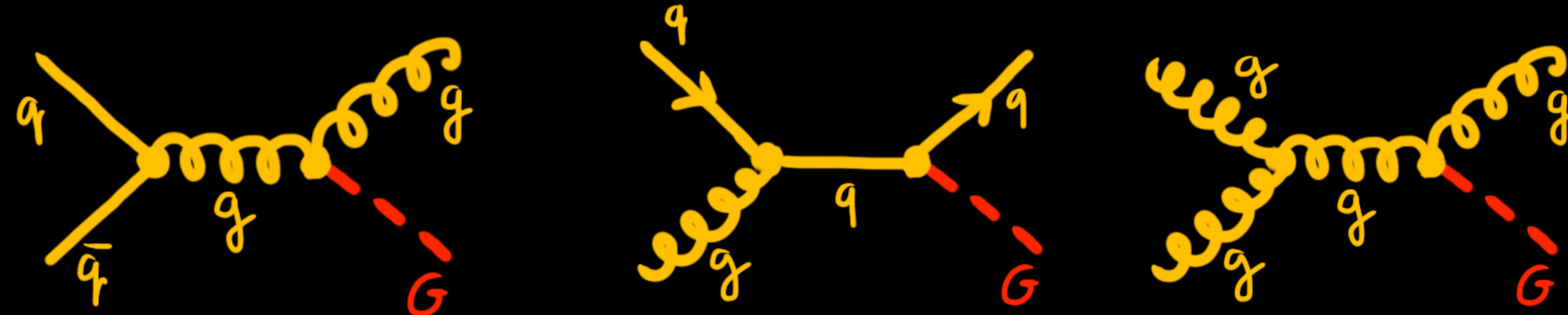
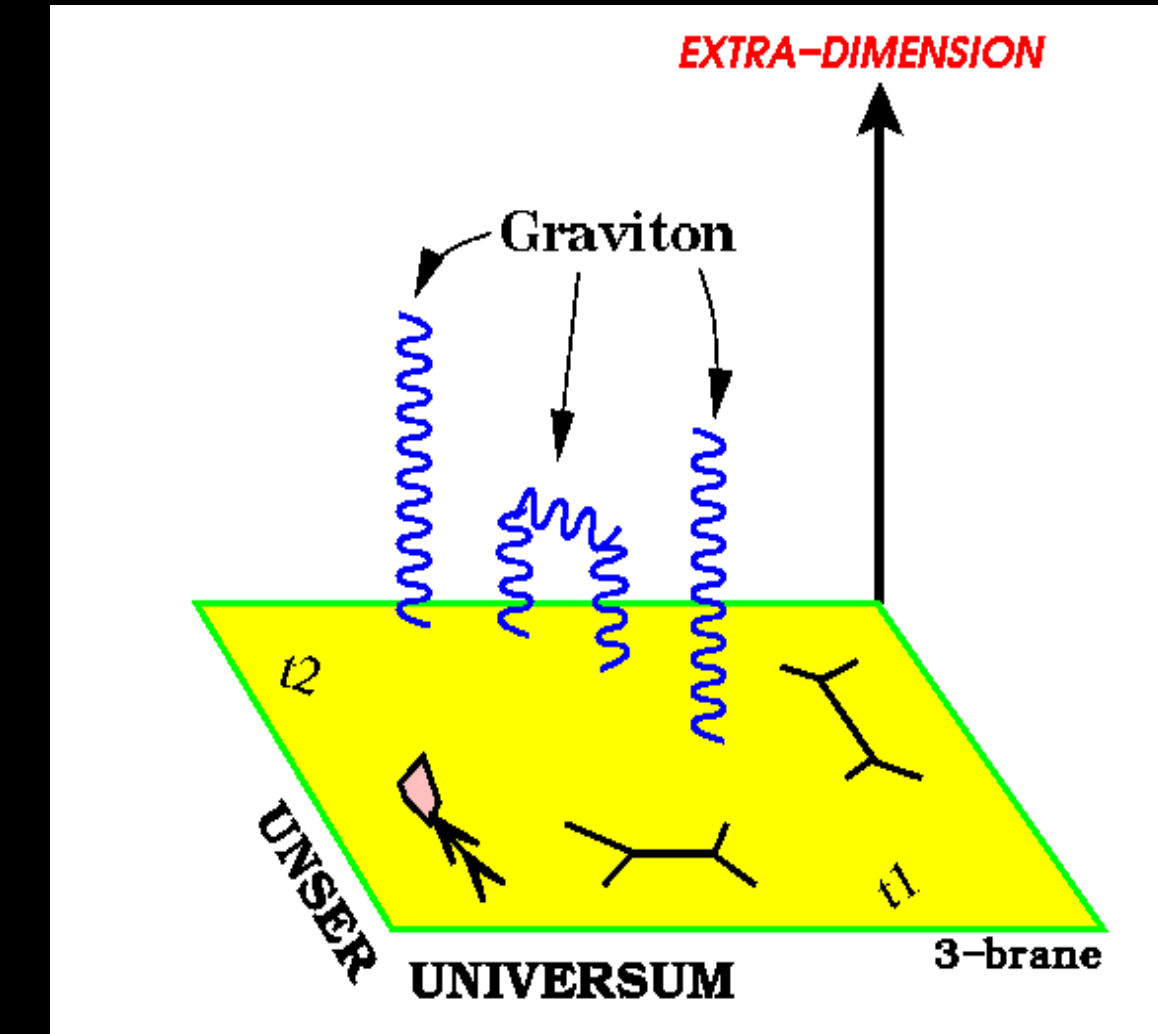
1 TeV

10 TeV

Moriond 2021

Extra Dimensions - an example of bump hunting

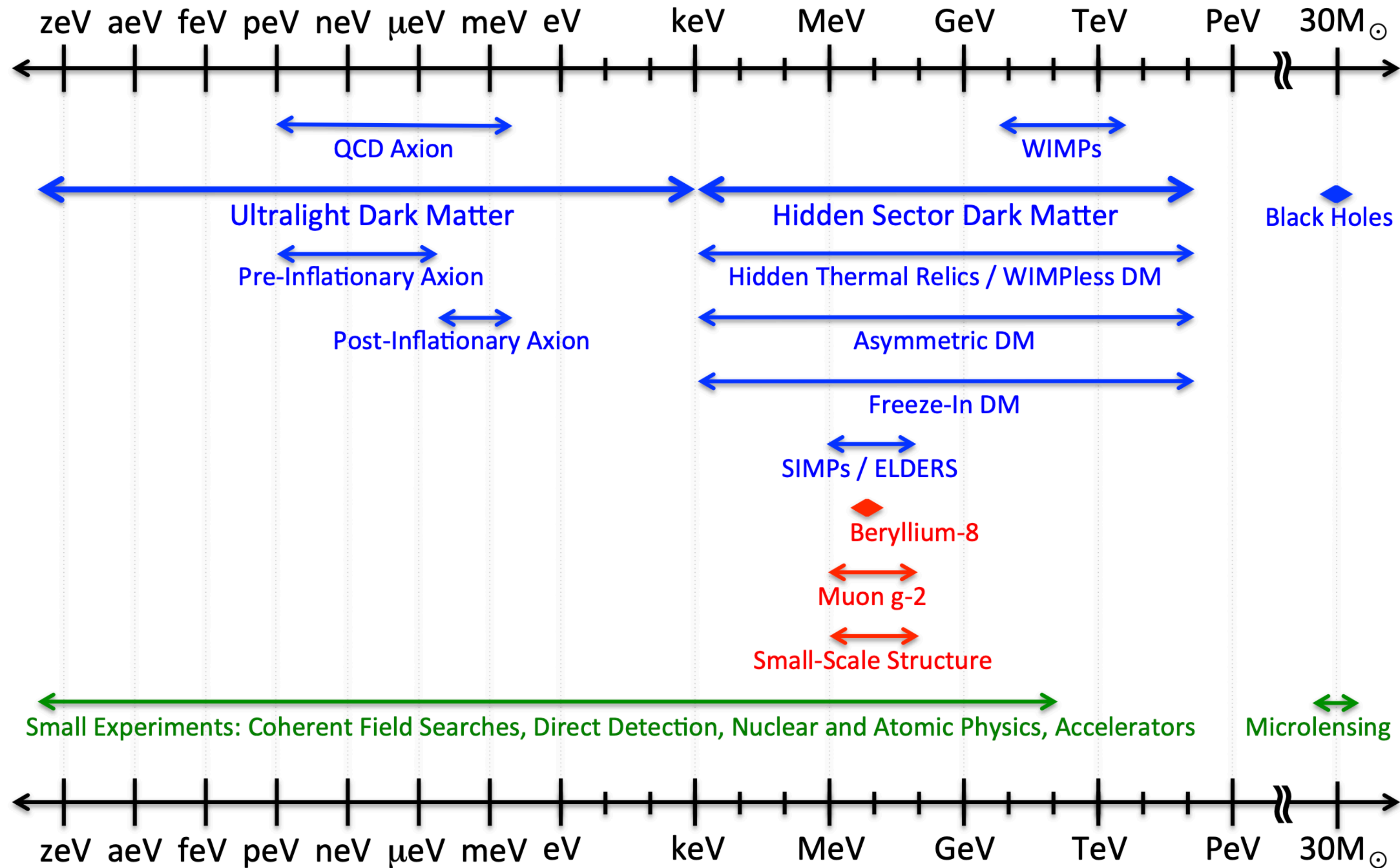
- Hierarchy problem solved by bringing down Planck scale
- n extra dimensions, compactified at radius R
- SM confined to brane in a higher dimensional space
- Only gravity can access extra dimensions



Dark Matter

Dark Matter Landscape

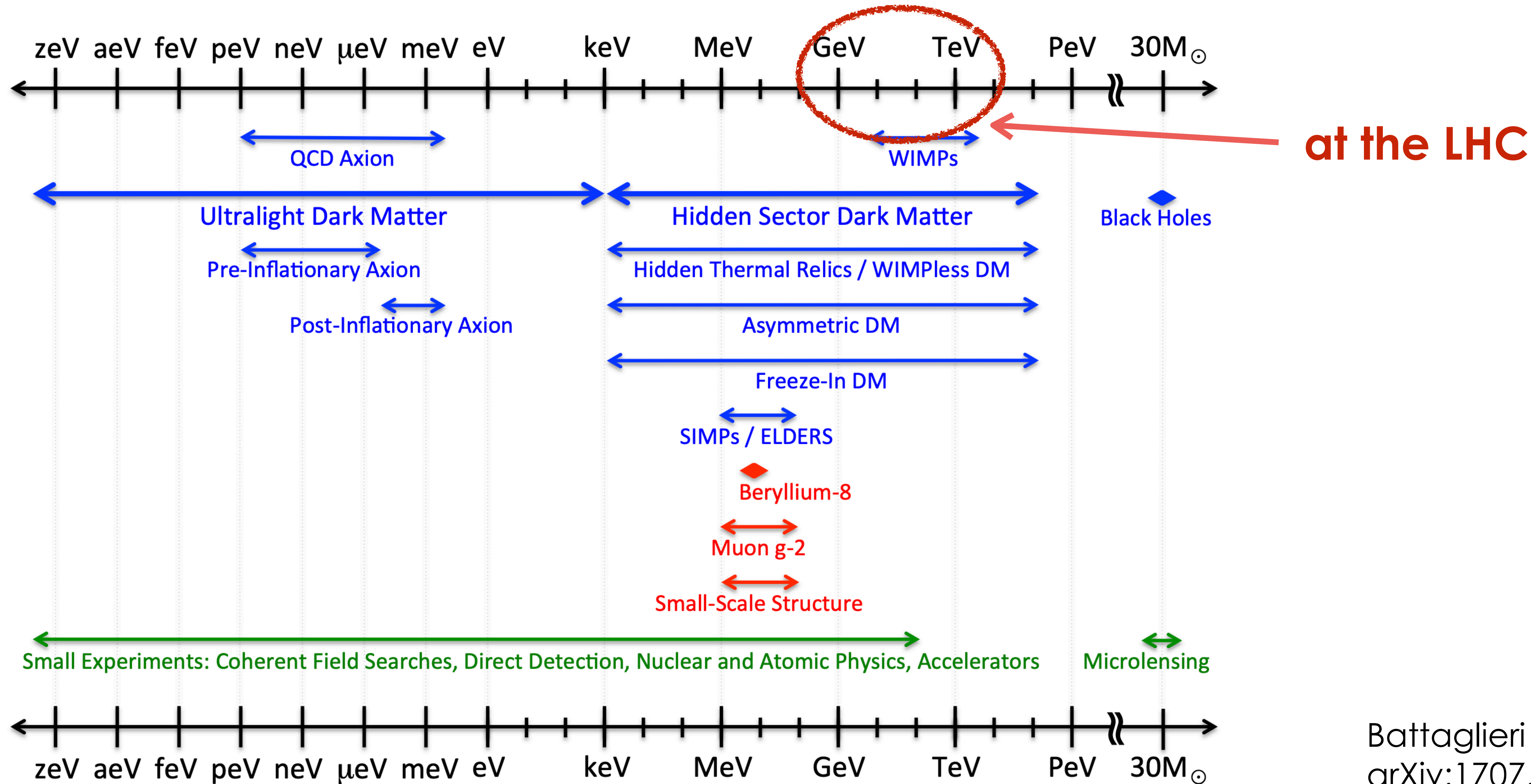
Dark Sector Candidates, Anomalies, and Search Techniques



Battaglieri et al.,
arXiv:1707.04591

Dark Matter Landscape

Dark Sector Candidates, Anomalies, and Search Techniques



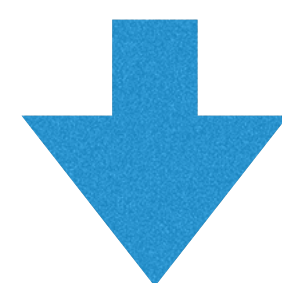
Battaglieri et al.,
arXiv:1707.04591

Dark Matter Searches at Colliders

There are three main approaches to detect DM:

- Dark Matter-nucleon scattering (direct detection)
- Annihilation (indirect detection)
- Pair production at colliders

All three processes are just topological permutations of the same Feynman diagram

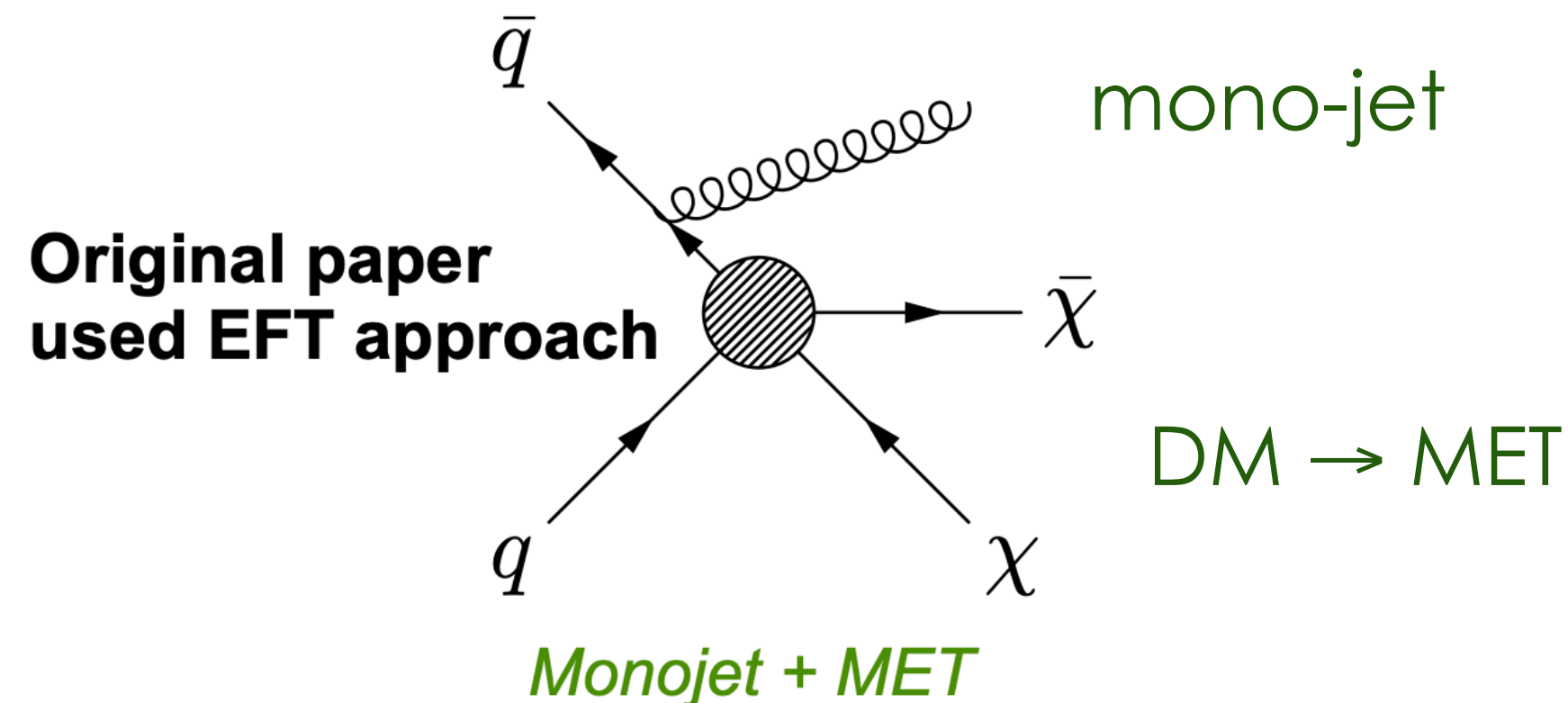
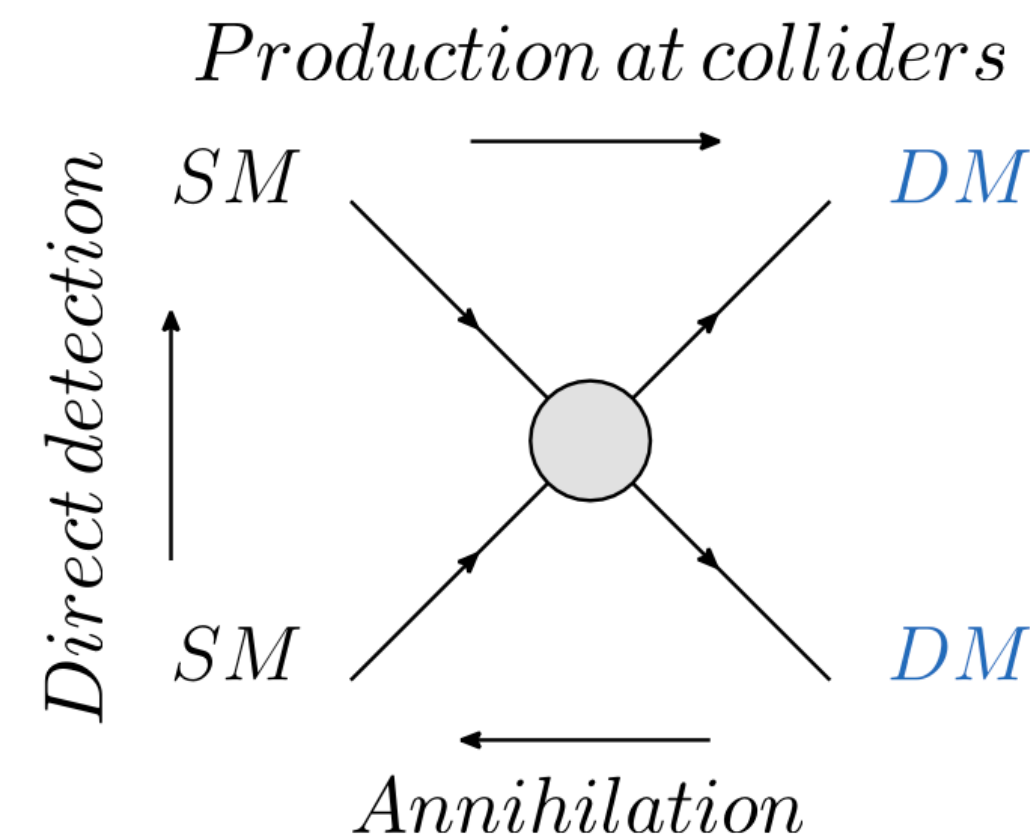


Leads to complementarity of searches

How to trigger on a pair of DM particles in a collider

Use Initial State Radiation (jet, γ , W, Z, ...)

Mono-everything analyses:
mono-jet, mono-photon, mono-top,
mono-V, mono-Higgs...



Beltran, Hooper, Kolb, Krusberg, and Tait,
"Maverick Dark Matter at Colliders" JHEP 09 (2010) 037

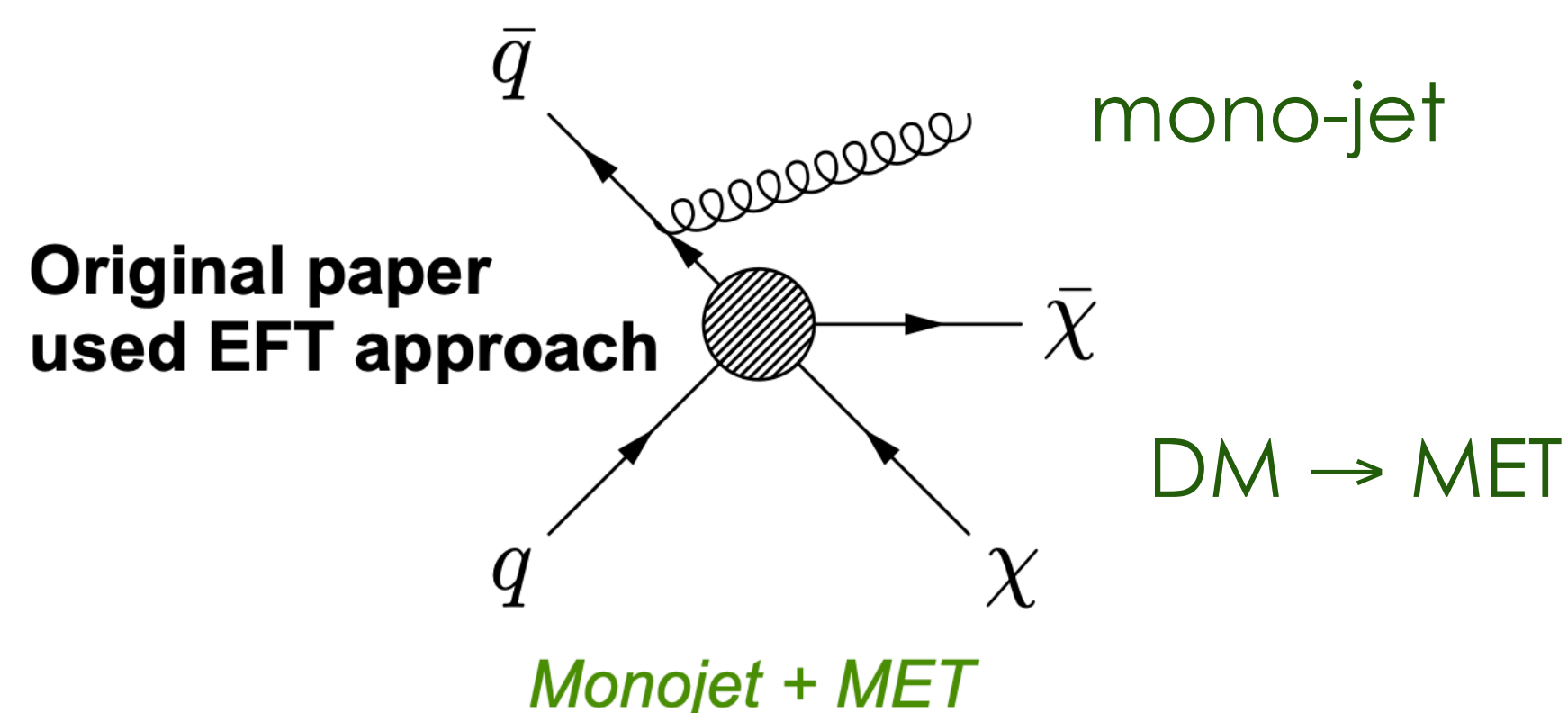
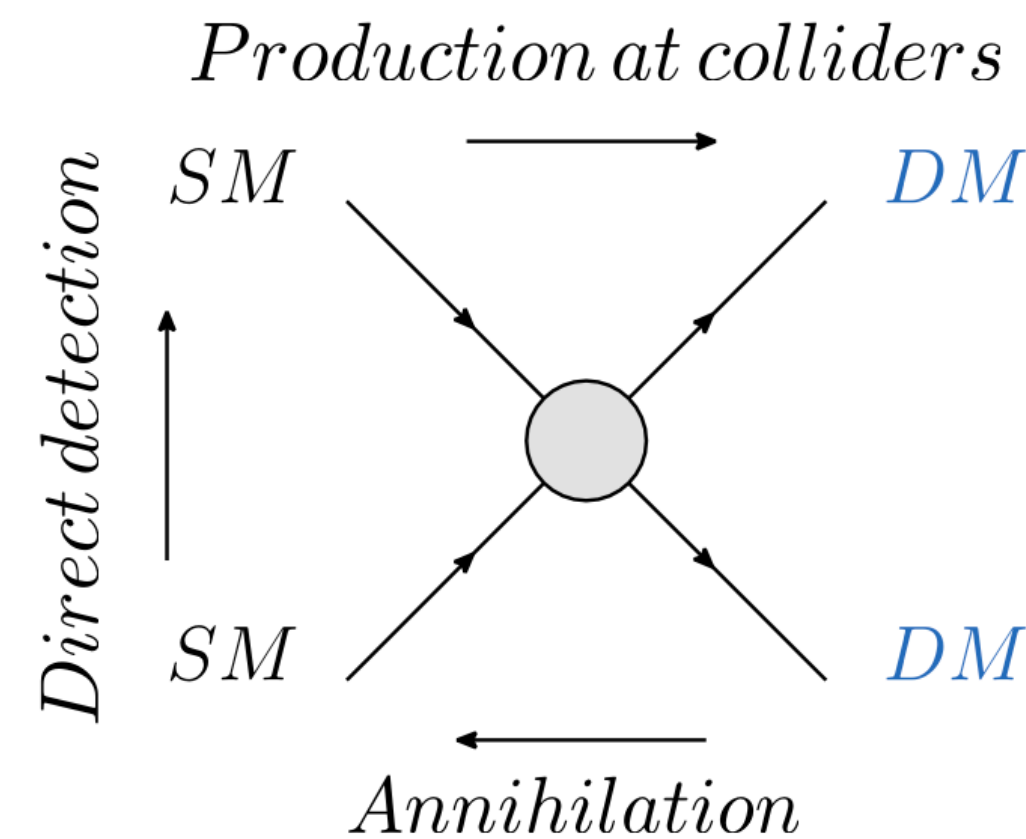
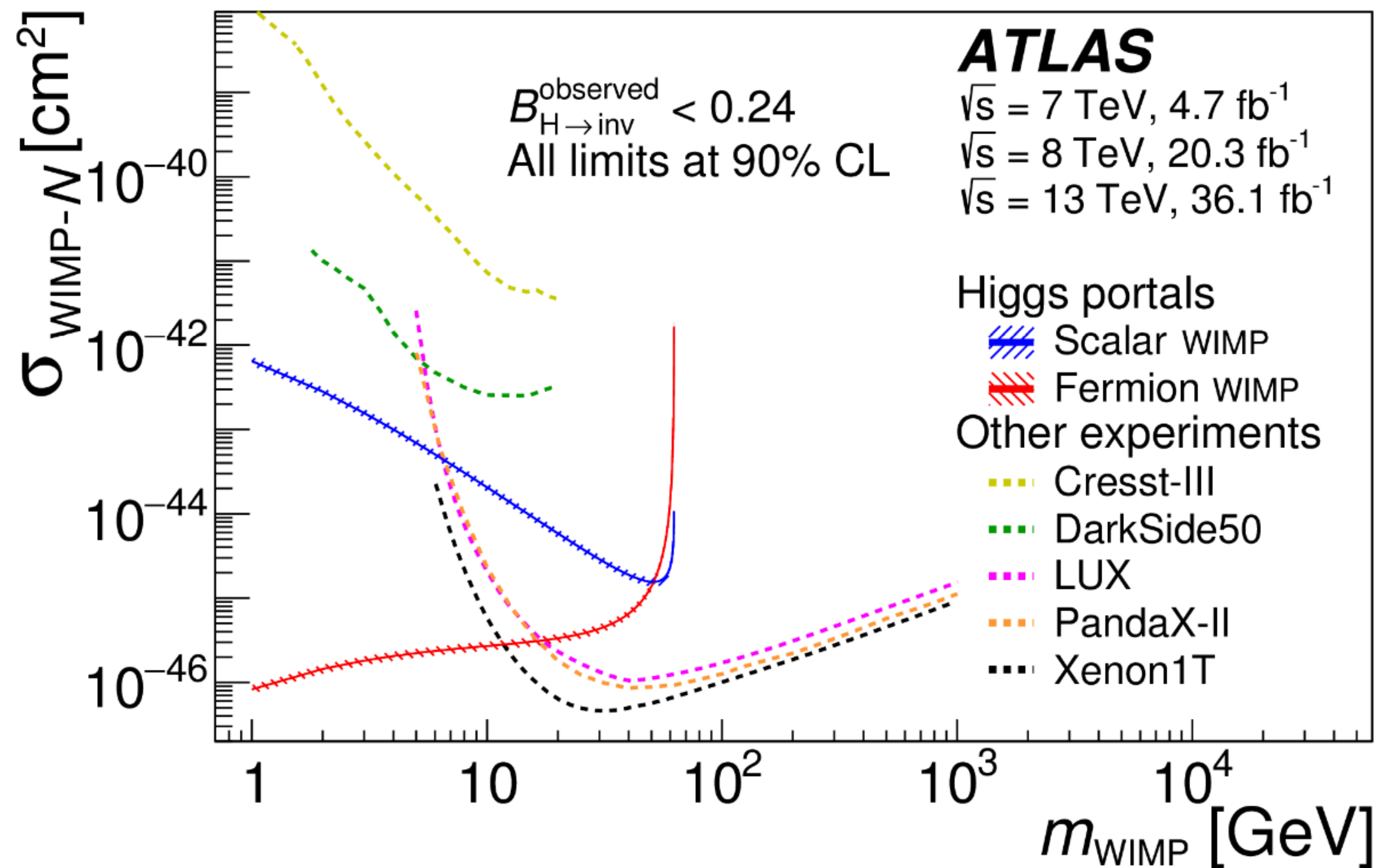
1707 GeV

Monojet event

MET: 1735 GeV

Run: 302393
Event: 738941529
2016-06-20 07:26:47 CEST

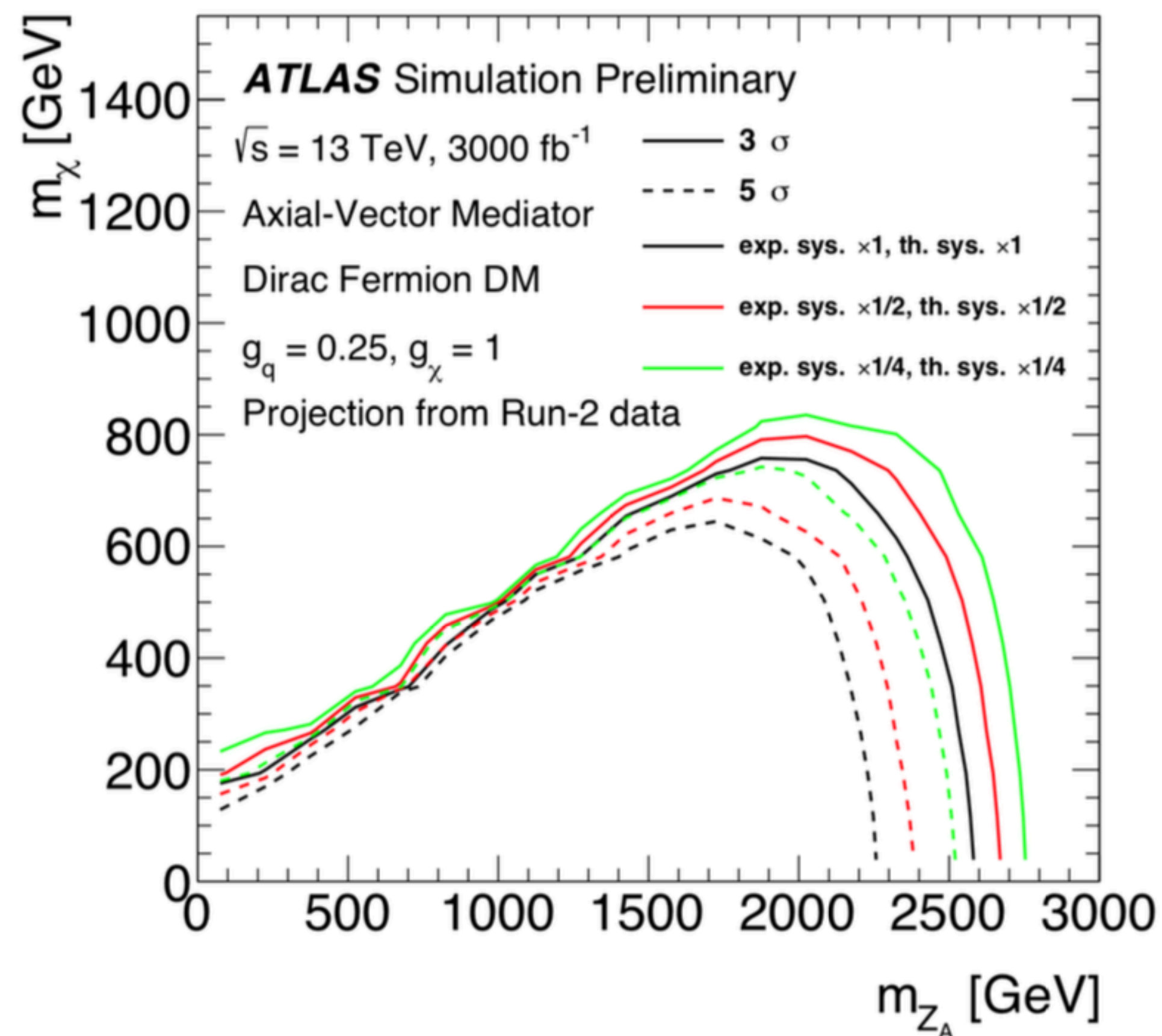
Dark Matter Searches at Colliders: Higgs \rightarrow invisible



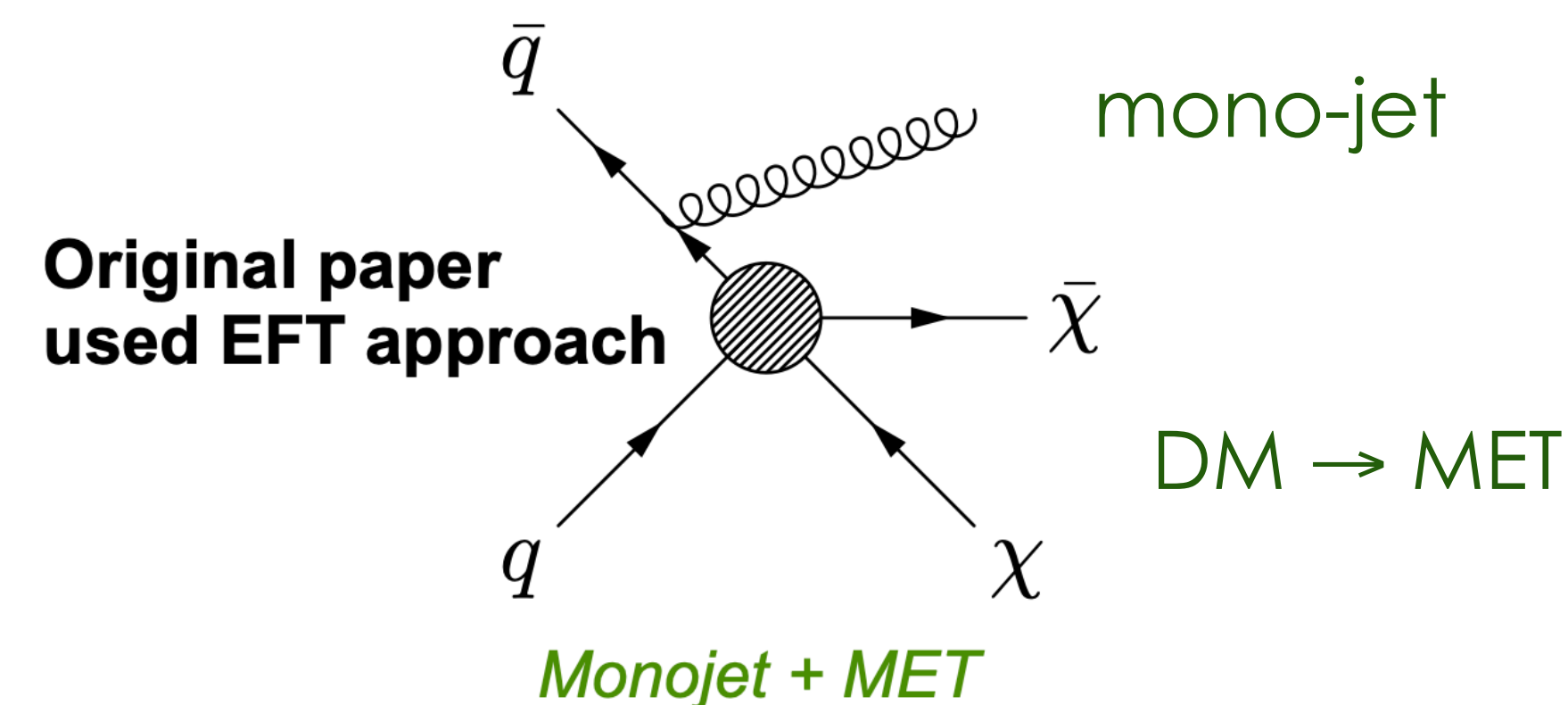
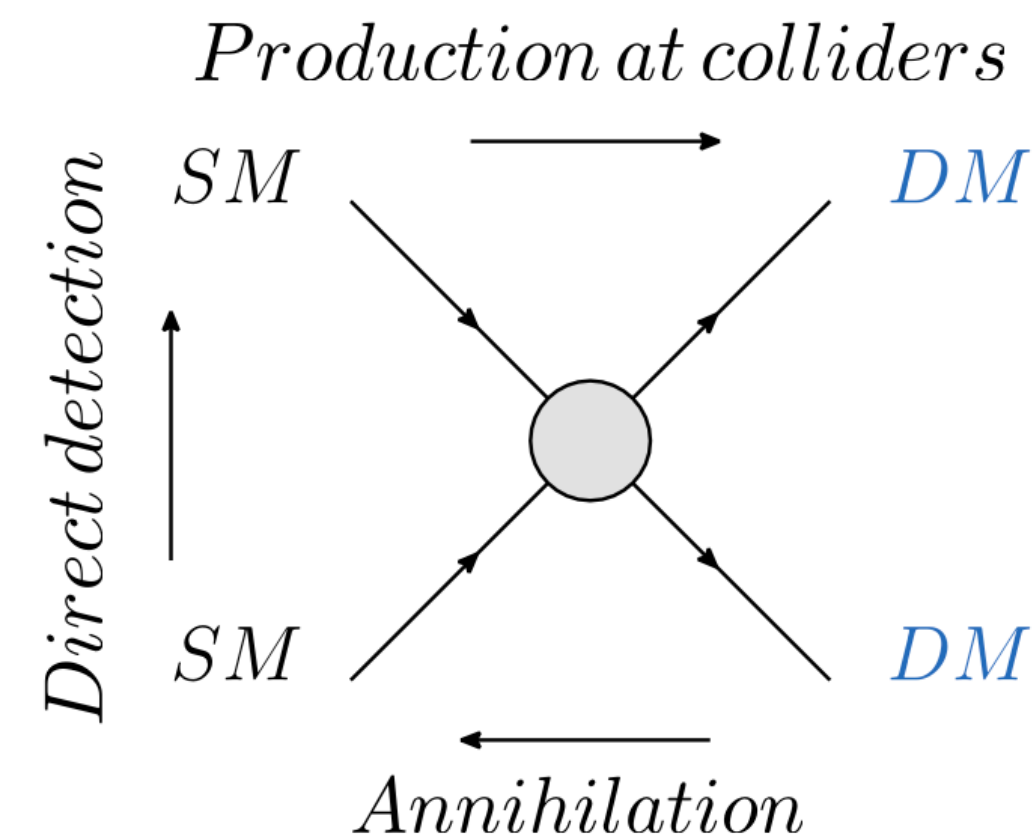
Beltran, Hooper, Kolb, Krusberg, and Tait,
 "Maverick Dark Matter at Colliders" JHEP 09 (2010) 037

Dark Matter Searches at Colliders

Mono-jets prospects at HL-LHC



Reach TeV range, but model dependent

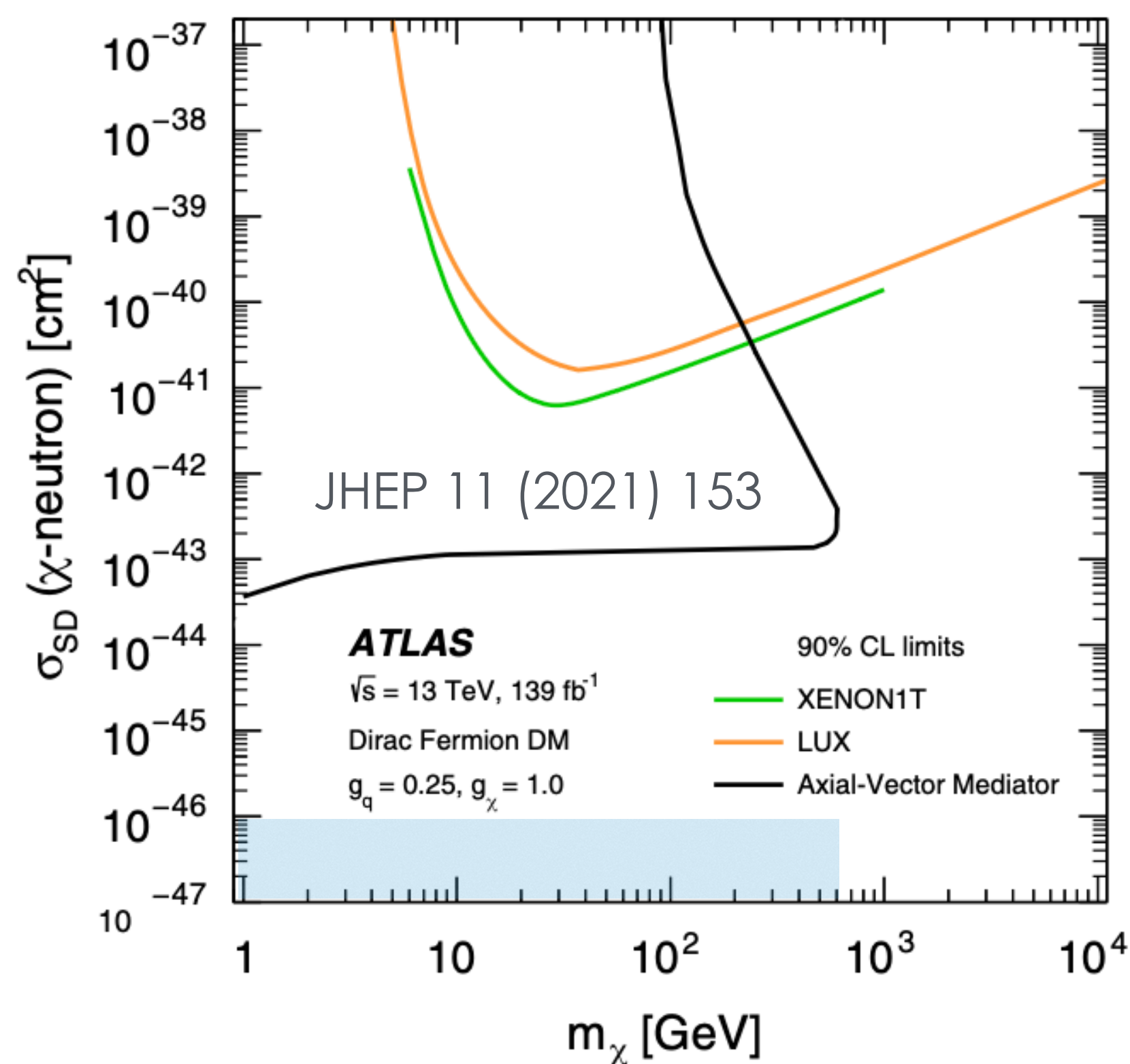


Beltran, Hooper, Kolb, Krusberg, and Tait,
 "Maverick Dark Matter at Colliders" JHEP 09 (2010) 037

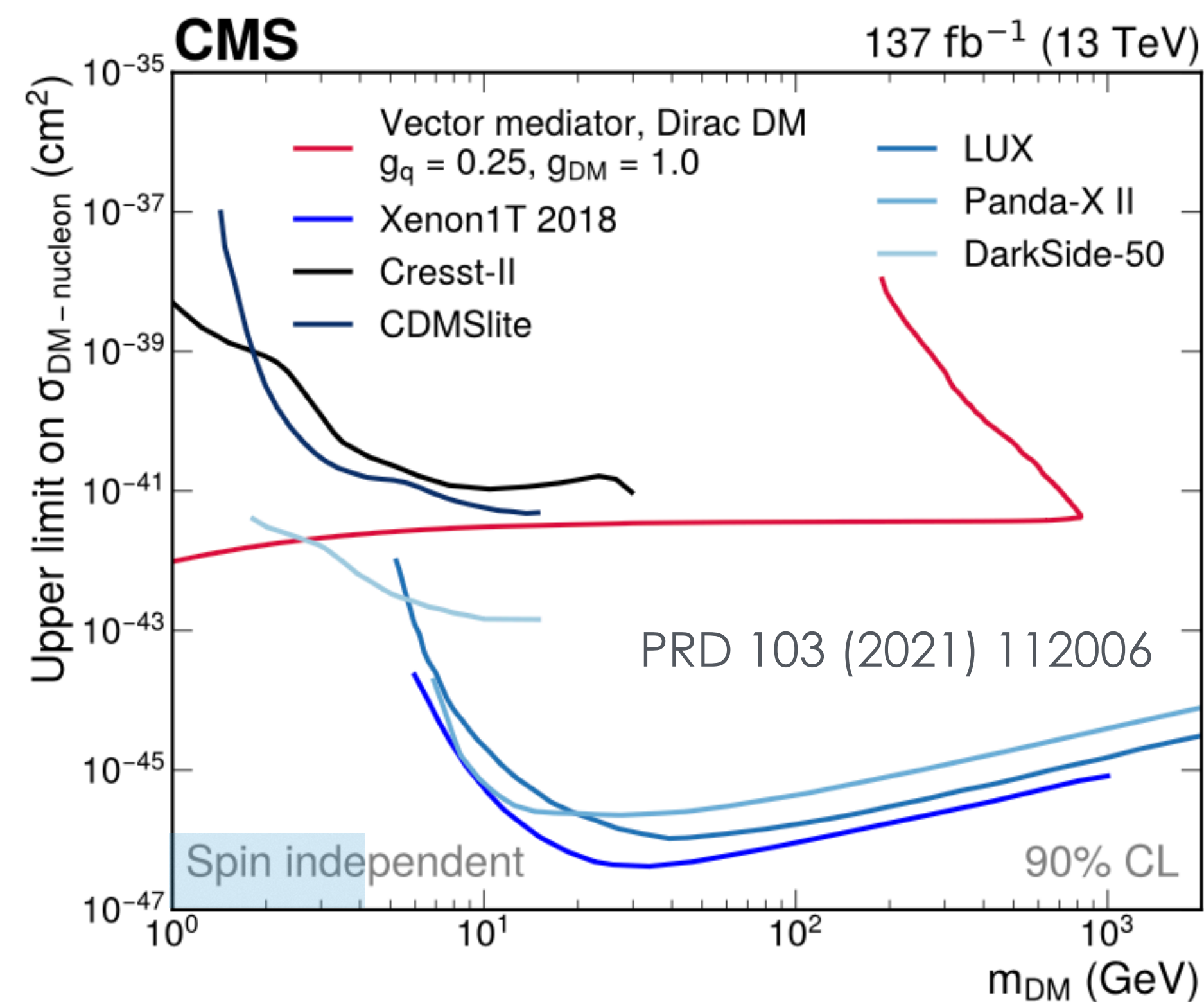
Dark Matter Searches at Colliders

Complementarity to Direct Detection

Collider experiments competitive w/ direct detection ones



Spin Dependent (axial-vector mediator)
 up to $m_{\text{DM}} \sim 500 \text{ GeV}$



Spin Independent (vector mediator)
 for very light $m_{\text{DM}} < \sim 5 \text{ GeV}$

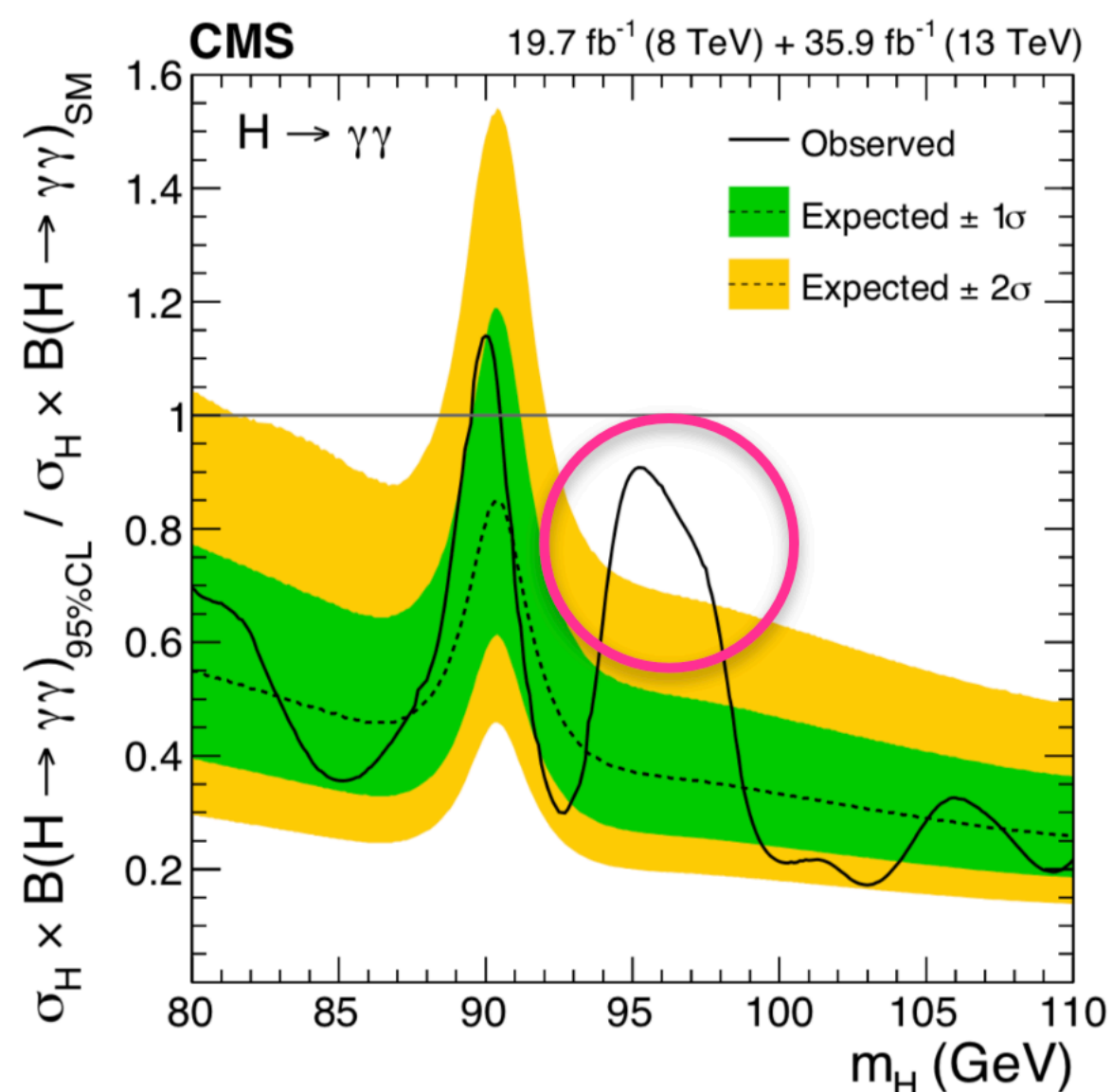
Do anomalies exist in the LHC data?

Anomalies do exist... a few examples

Several 3σ local excesses...

$$a \rightarrow \gamma\gamma$$

Local (global) significance at 95.3 GeV : 2.9σ (1.5σ)

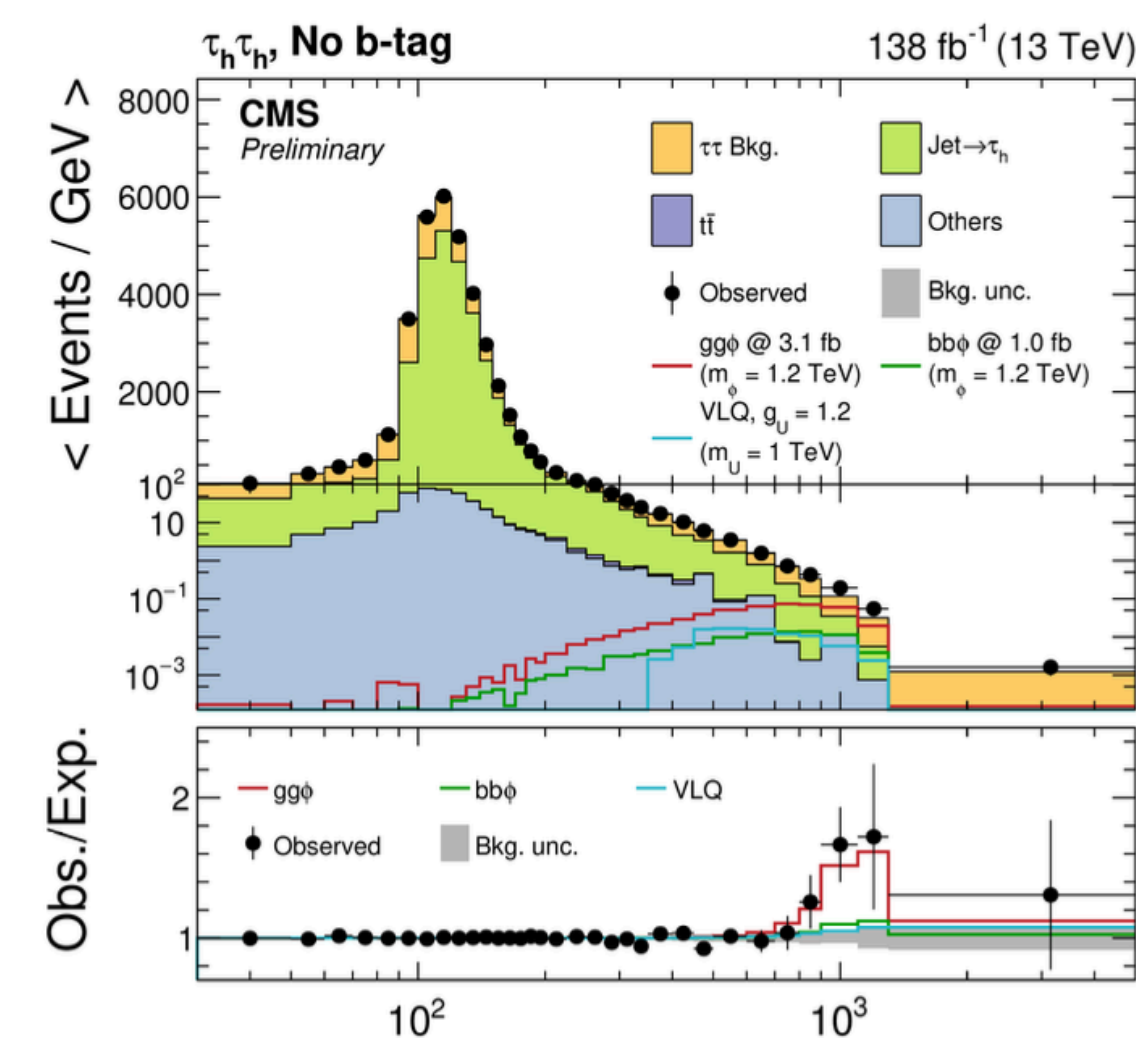
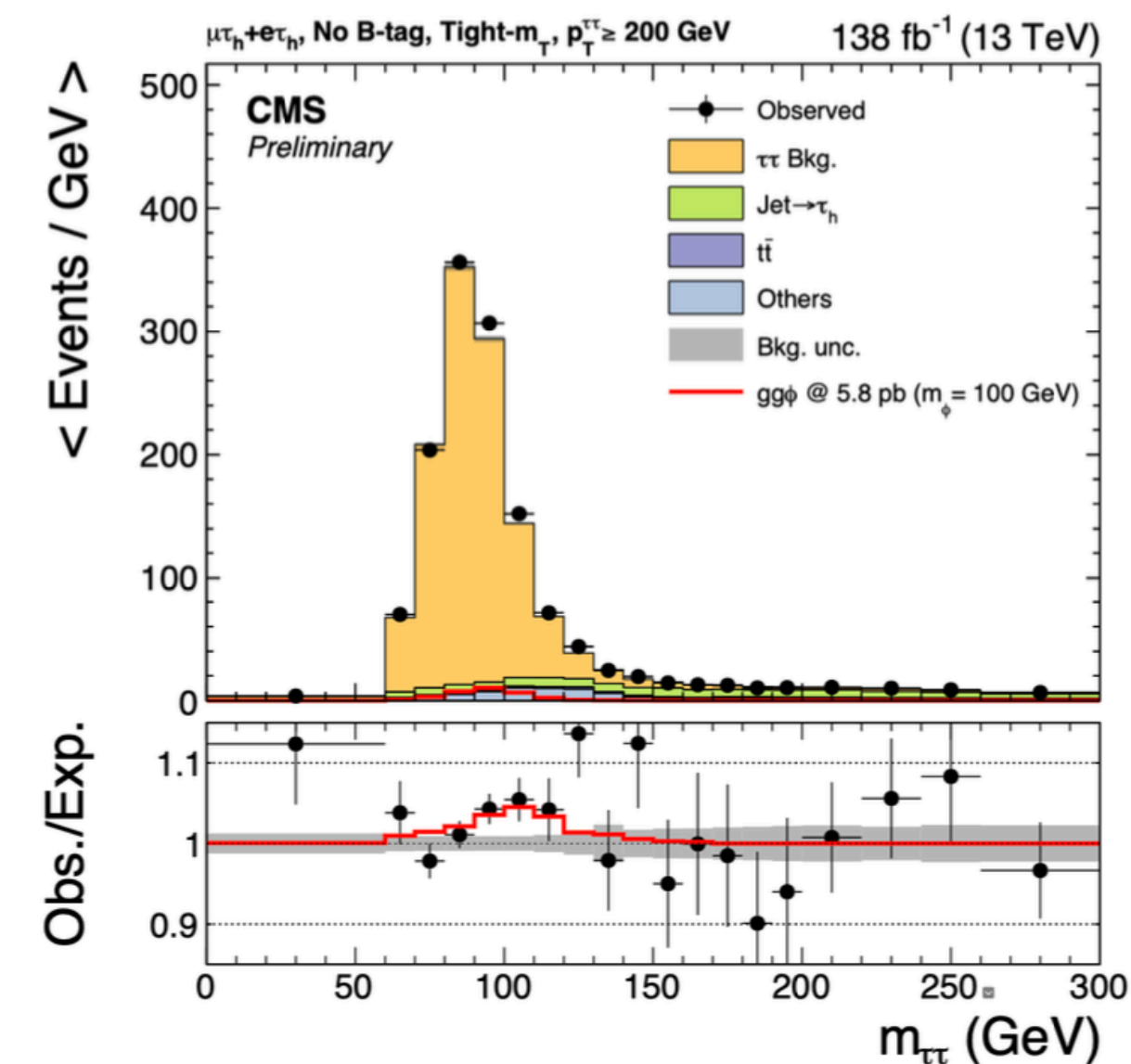


$$\phi \rightarrow \tau^+ \tau^-$$

Local (global) significance at 95 TeV : 3.1σ (2.7σ)

$$\phi \rightarrow \tau^+ \tau^-$$

Local (global) significance at 1.2 TeV : 2.8σ (2.4σ)



Warning:

With such large number of searches statical fluctuations are expected and there could be systematic effects!

NO NEW PHYSICS OBSERVED

→ LOOK FURTHER/DEEPER

Explore new tools and techniques

Long-lived particles (LLPs)

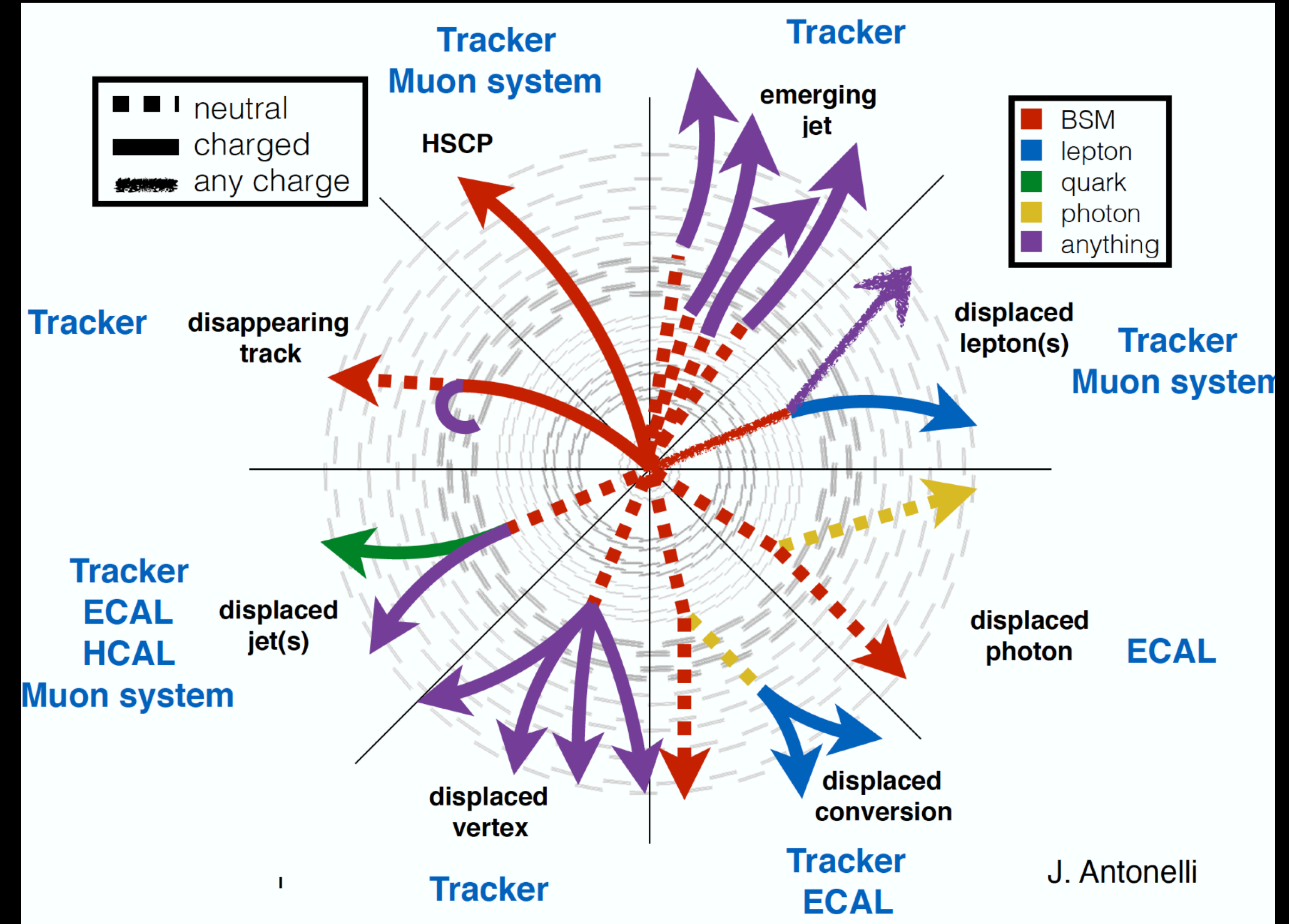
Long-lived particles (LLPs) are common in BSM

- Small phase space for decay (e.g. Split SUSY)
- Small couplings to SM particles
 - Suppressed (e.g. Higgs/gauge portal to Dark Sector)
 - Forbidden by symmetry (SUSY R-parity)

Could have escaped observation so far

Challenges:

- Often requires new triggers
- Exotic detector signatures (requiring new tools)
- Non-standard backgrounds



ATLAS Long-lived Particle Searches* - 95% CL Exclusion

Status: July 2021

ATLAS Preliminary

$\int \mathcal{L} dt = (18.4 - 139) \text{ fb}^{-1}$

$\sqrt{s} = 8, 13 \text{ TeV}$

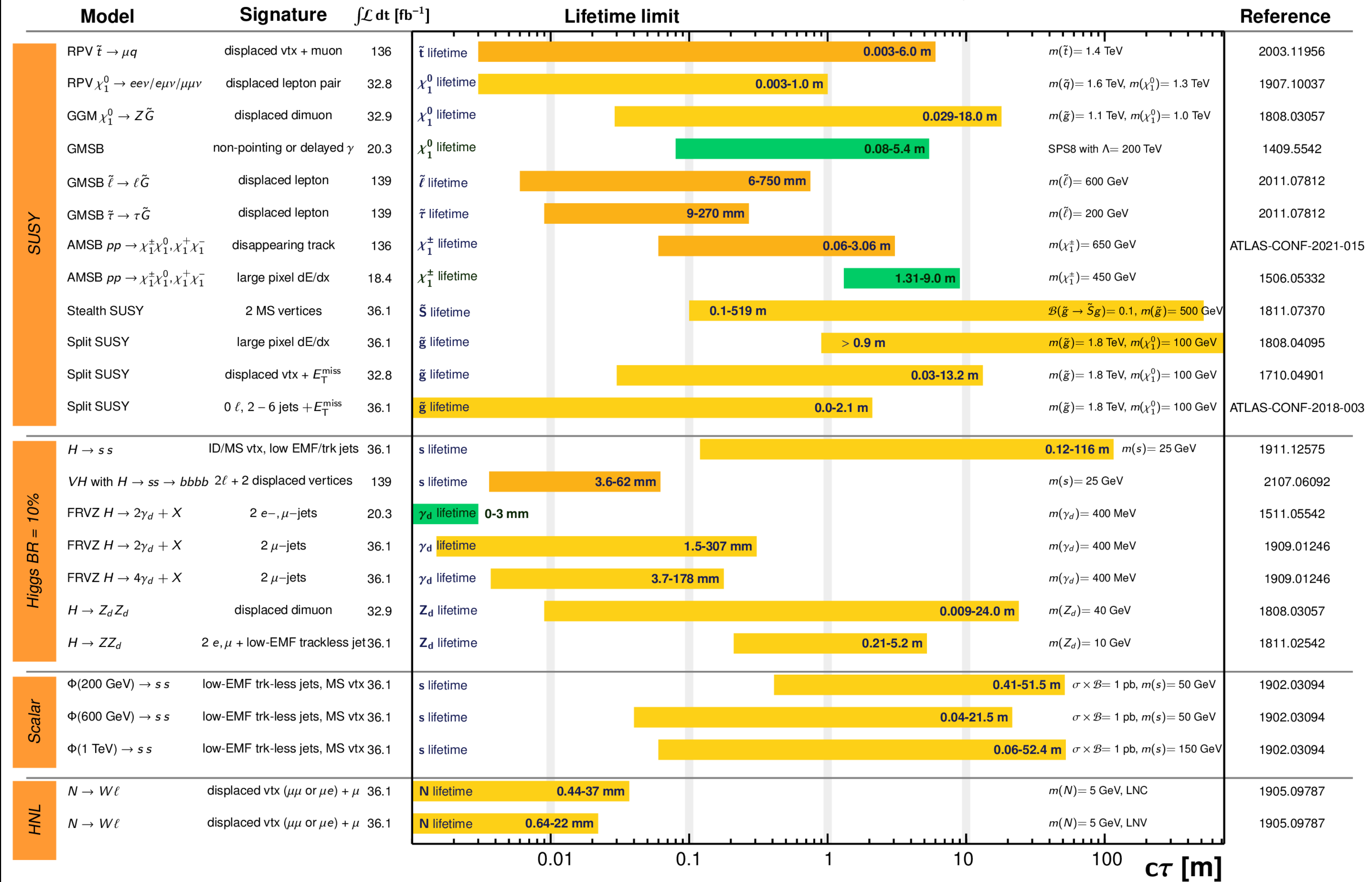
Displaced leptons

Disappearing track

Large pixel dE/dx

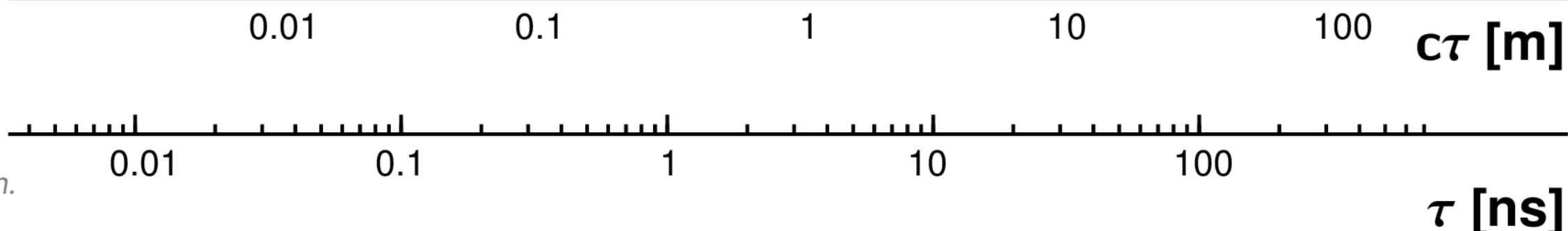
Track-less jets

Displaced vertices



$\sqrt{s} = 8 \text{ TeV}$ $\sqrt{s} = 13 \text{ TeV}$ partial data $\sqrt{s} = 13 \text{ TeV}$ full data

*Only a selection of the available lifetime limits is shown.



NO NEW PHYSICS OBSERVED

→ LOOK FURTHER/DEEPER

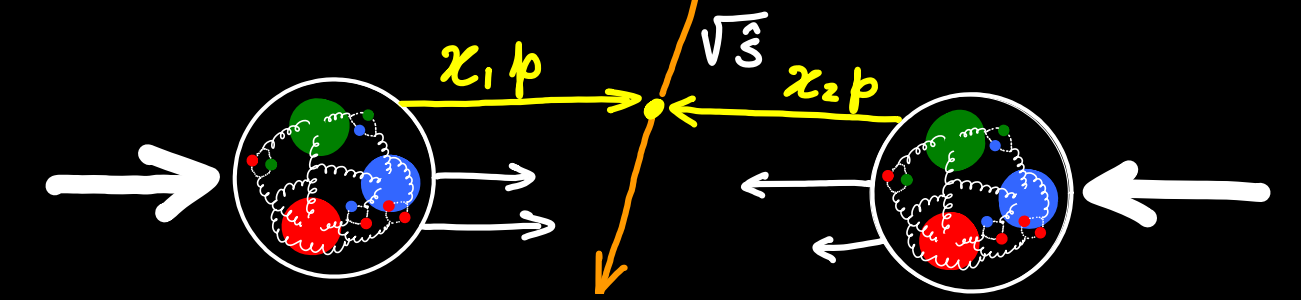
Explore new tools and techniques

The highest **energy** possible

The highest **luminosity** possible

As **low backgrounds** as possible

THE LARGE HADRON COLLIDER



LHC and experiments
being upgraded

HL-LHC starting in 2029: 3000-4000 fb⁻¹

Center of mass energy: 14 TeV

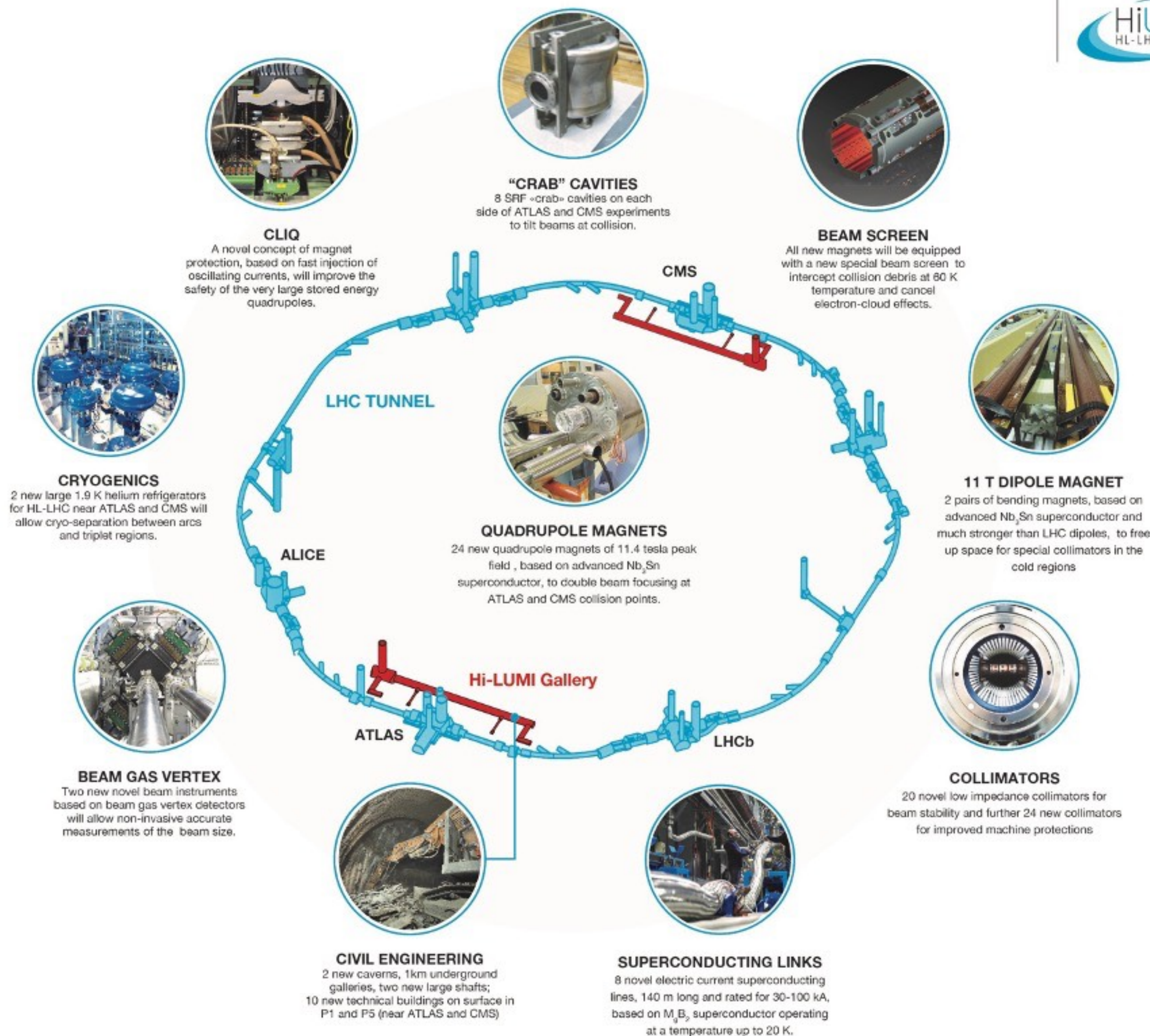


HL-LHC Upgrade

~1.2 km of accelerator will be upgraded



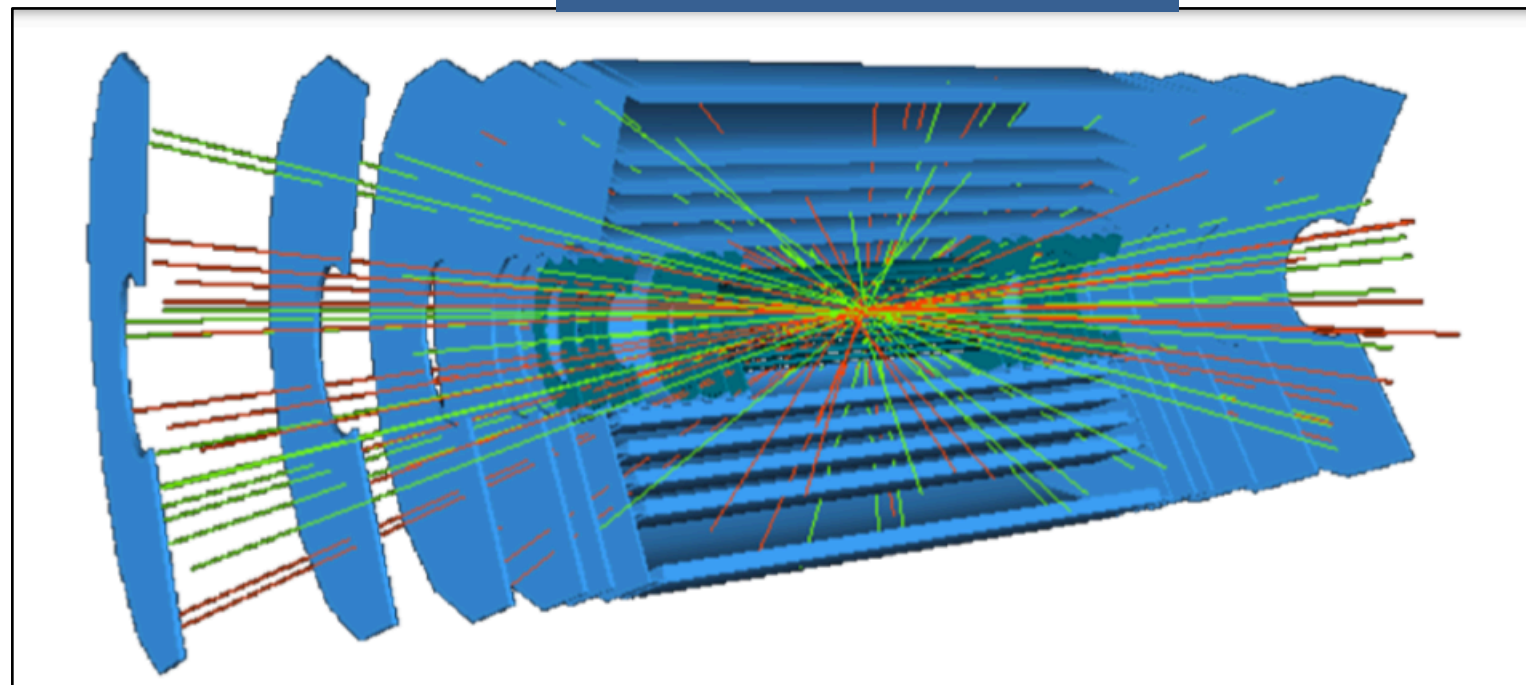
CERN 2019



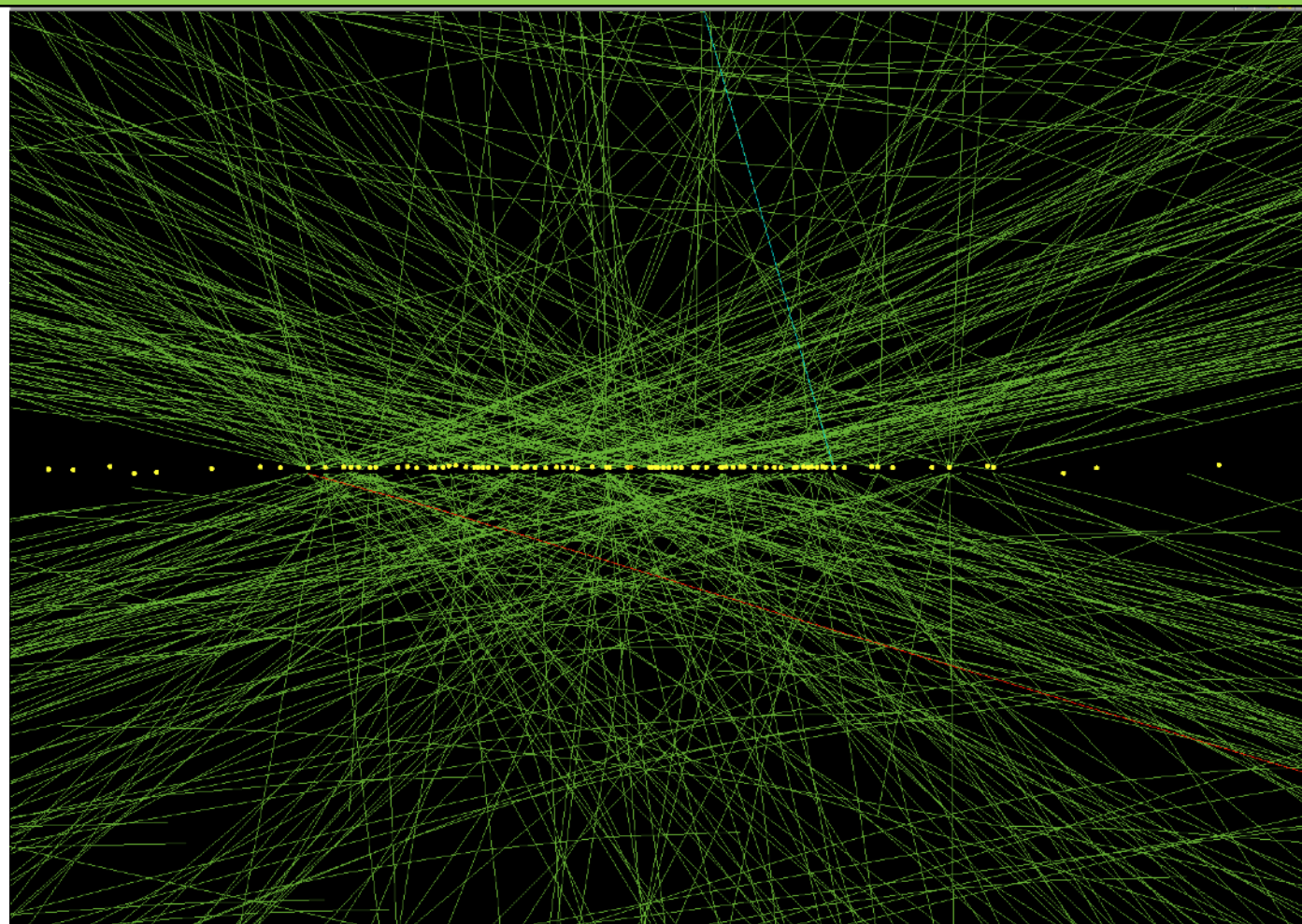
Challenging Phase-II Upgrades of ATLAS and CMS

Higher peak luminosity and larger pile-up (from ~ 30 to 140-200 events/x-ing) require: increased radiation hardness and granularity, dedicated (timing) detectors, larger bandwidth, faster and more granular readout electronics, improved triggers and higher redundancy, to provide similar or better performance than current detectors (including trigger thresholds) in much harsher HL-LHC environment

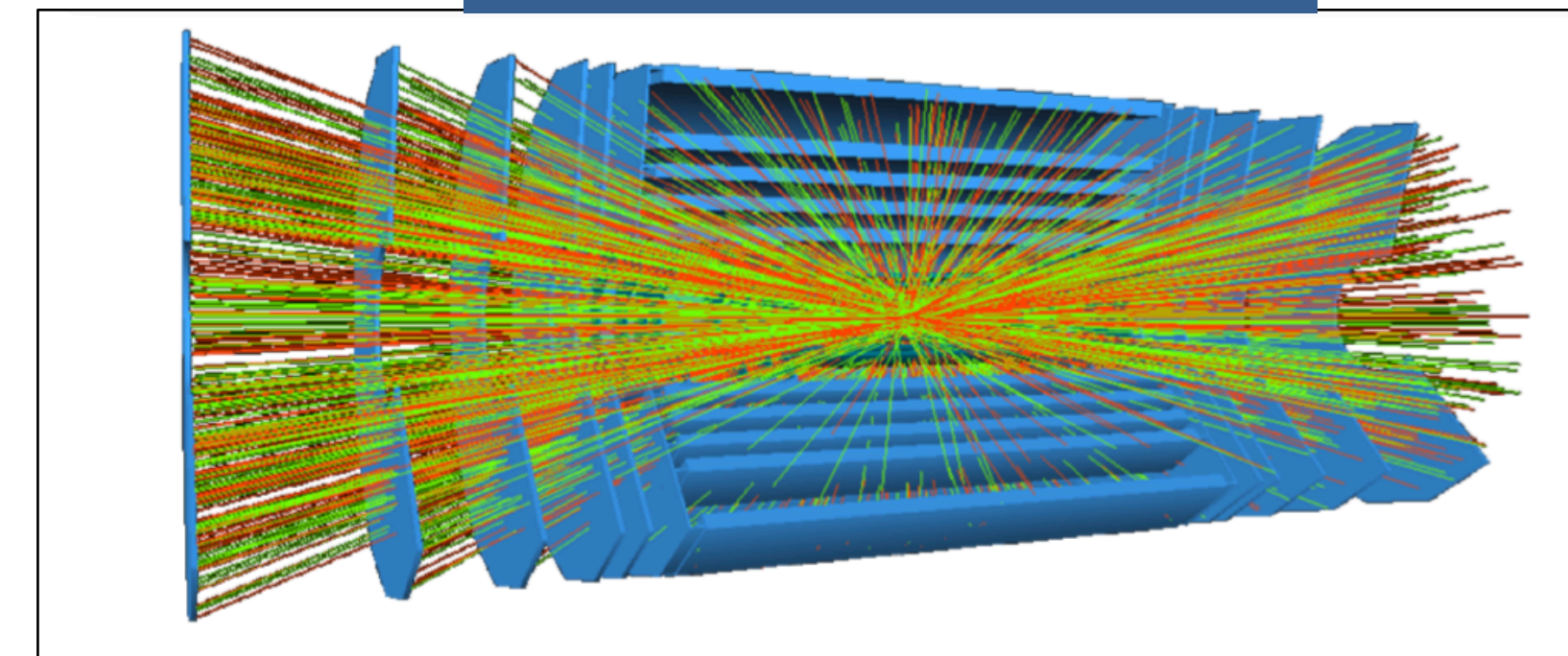
LHC: ~ 30 evts/x-ing



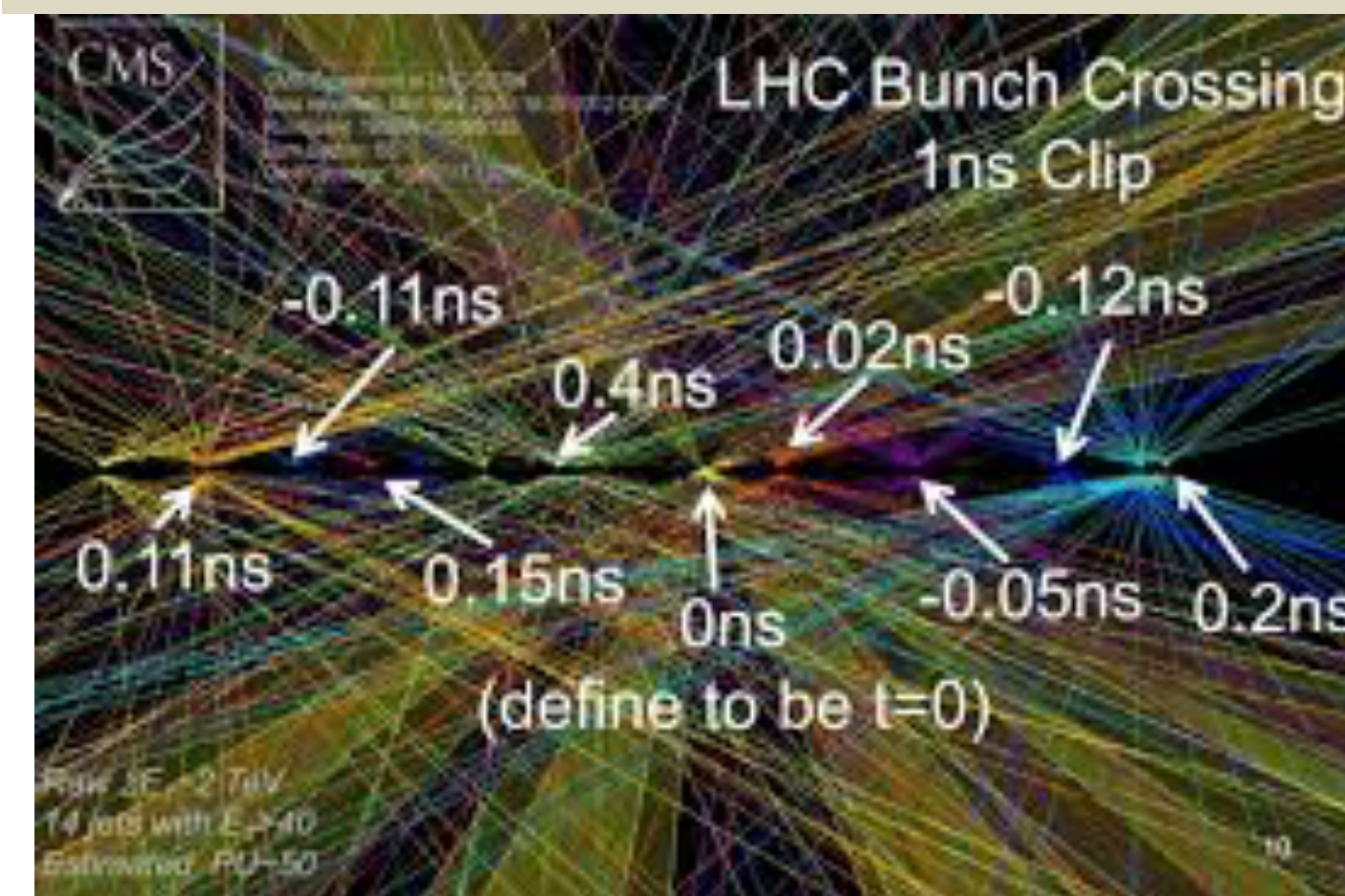
Event with 78 reconstructed vertices (CMS Run 2 data)
Note: ~ 20 expected when detectors were designed



HL-LHC: $\sim 140-200$ evts/x-ing

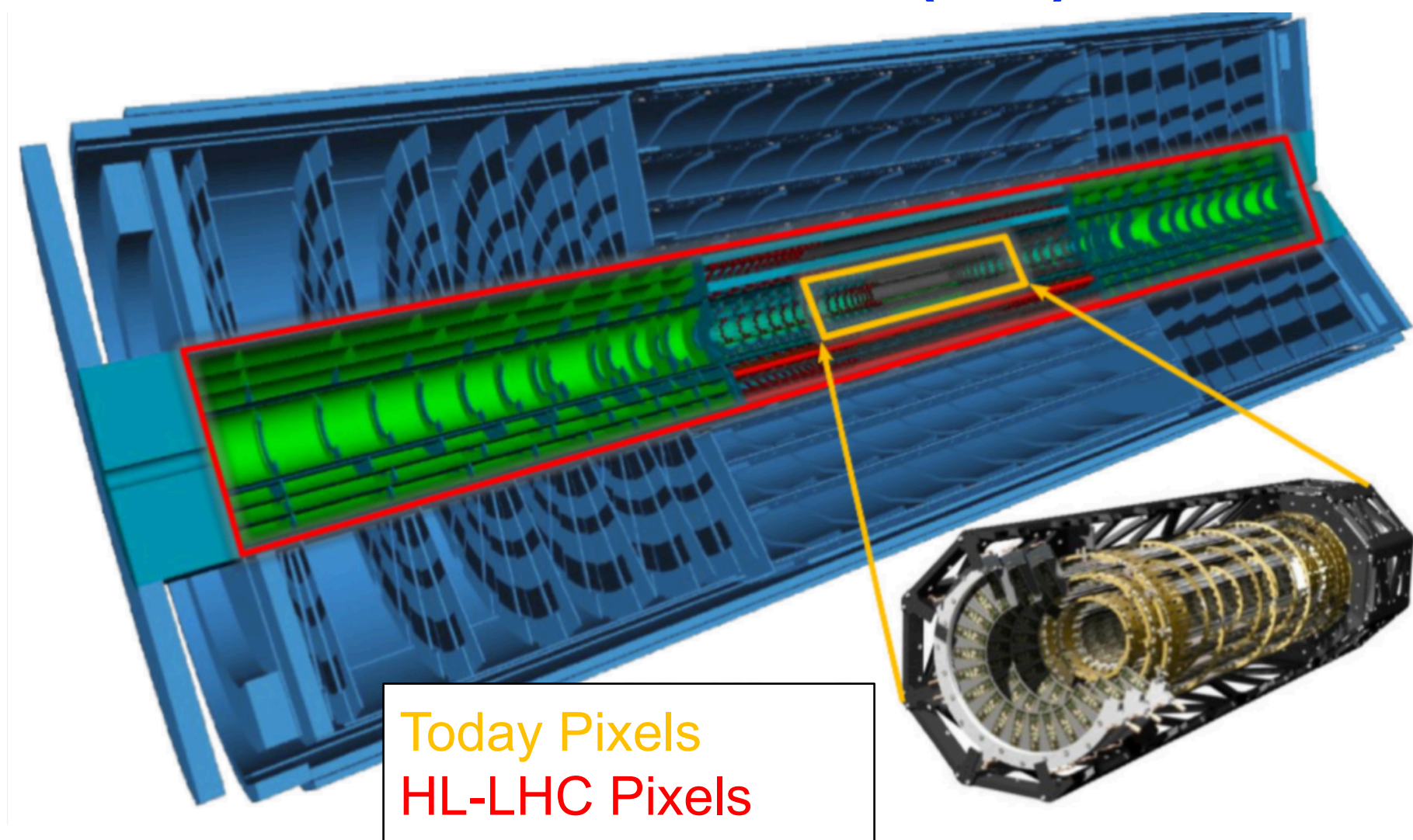


New timing detectors with resolution ~ 30 ps in both experiments will allow 4-dimensional identification of primary vertex



Challenging Phase-II Upgrades of ATLAS and CMS

ATLAS tracker (ITk)



$$|\eta| < 4$$

Low mass, rad hard

Barrel: 5 pixel + 4 strip layers

End-cap: up to 23 pixel + 6 strip rings

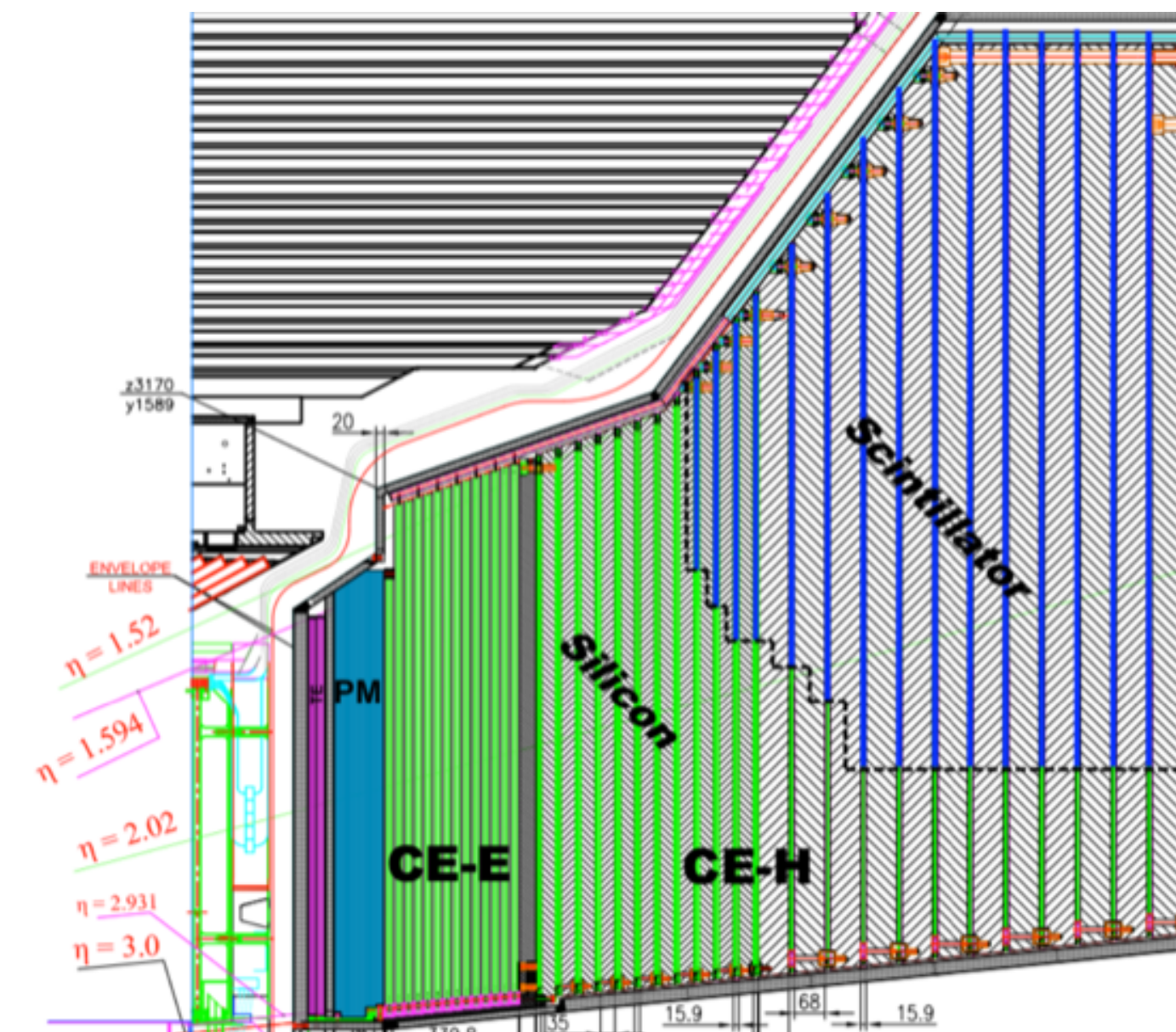
Pixel size: $25 \times 100 \mu\text{m}^2$ and $50 \times 50 \mu\text{m}^2$

Strip size (barrel): $\sim 75 \mu\text{m} \times 24\text{-}42 \text{ mm}$

Total Si area: $\sim 180 \text{ m}^2$

Total # of channels: ~ 5 billion (50 x today)

CMS end-cap calorimeter (HGCal)



$$1.5 < |\eta| < 3$$

Unprecedented lat. and long. segmentation (ILC-type)

Time resolution $\sim 30 \text{ ps}$

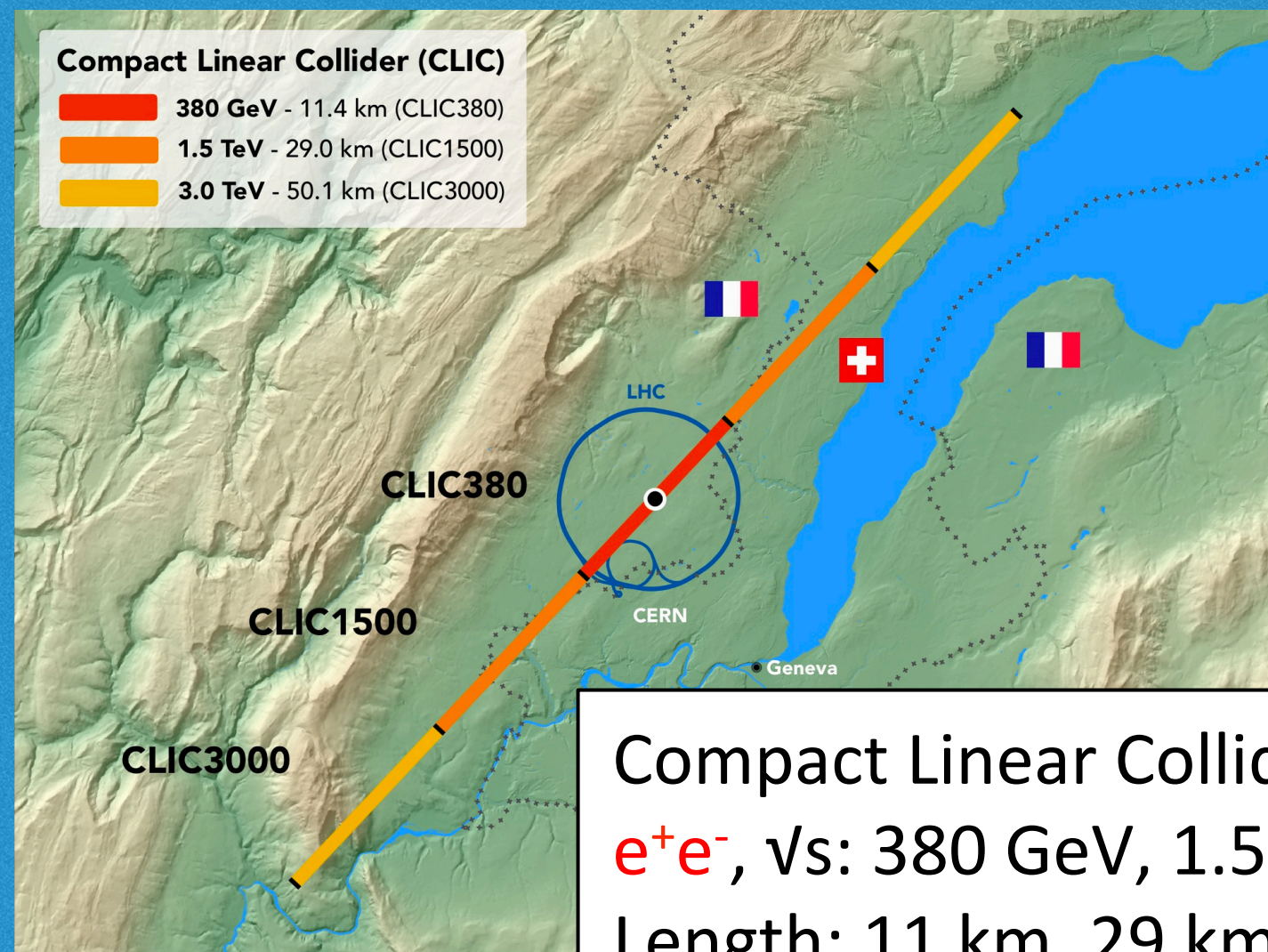
EM part (CE-E): Si pads, Cu/CuW/Pb absorber
26 layers, $25 X_0$

HAD part (CE-H): Si and scintillator, steel absorber
21 layers, 8.5λ

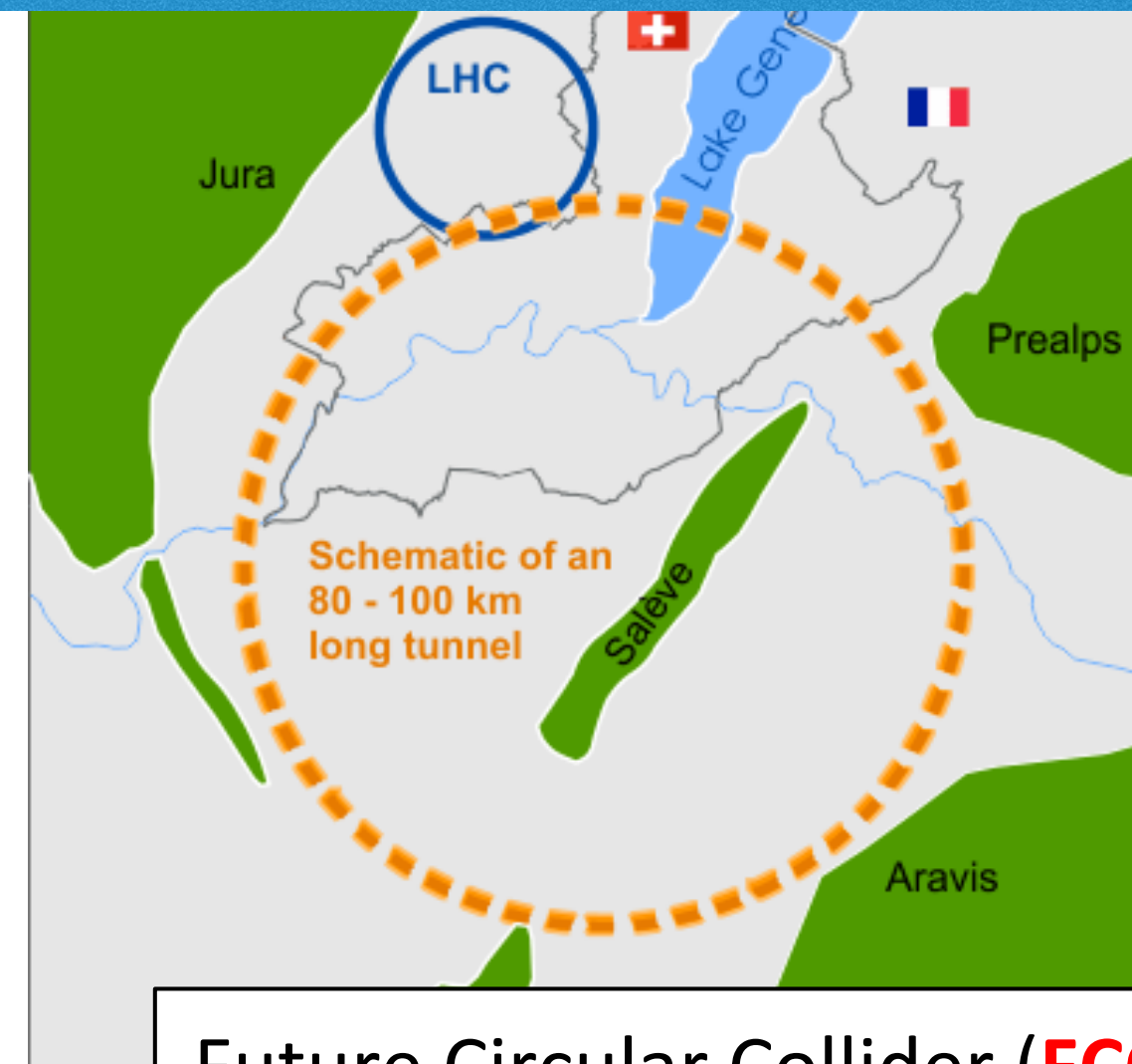
$\sim 600 \text{ m}^2$ of Si pads ($0.5\text{-}1 \text{ cm}^2$)

10^6 channels

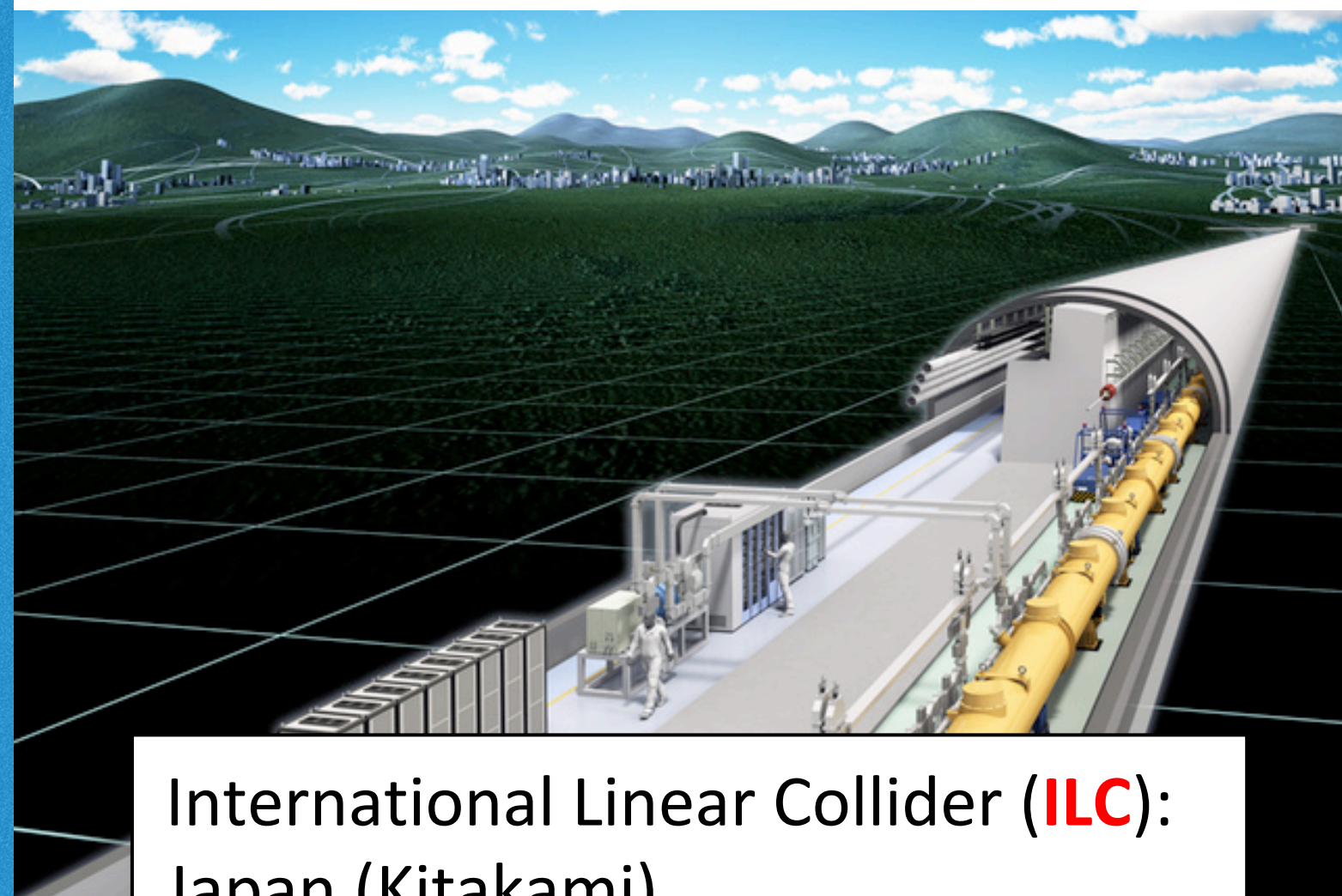
New High-Energy Collider Projects



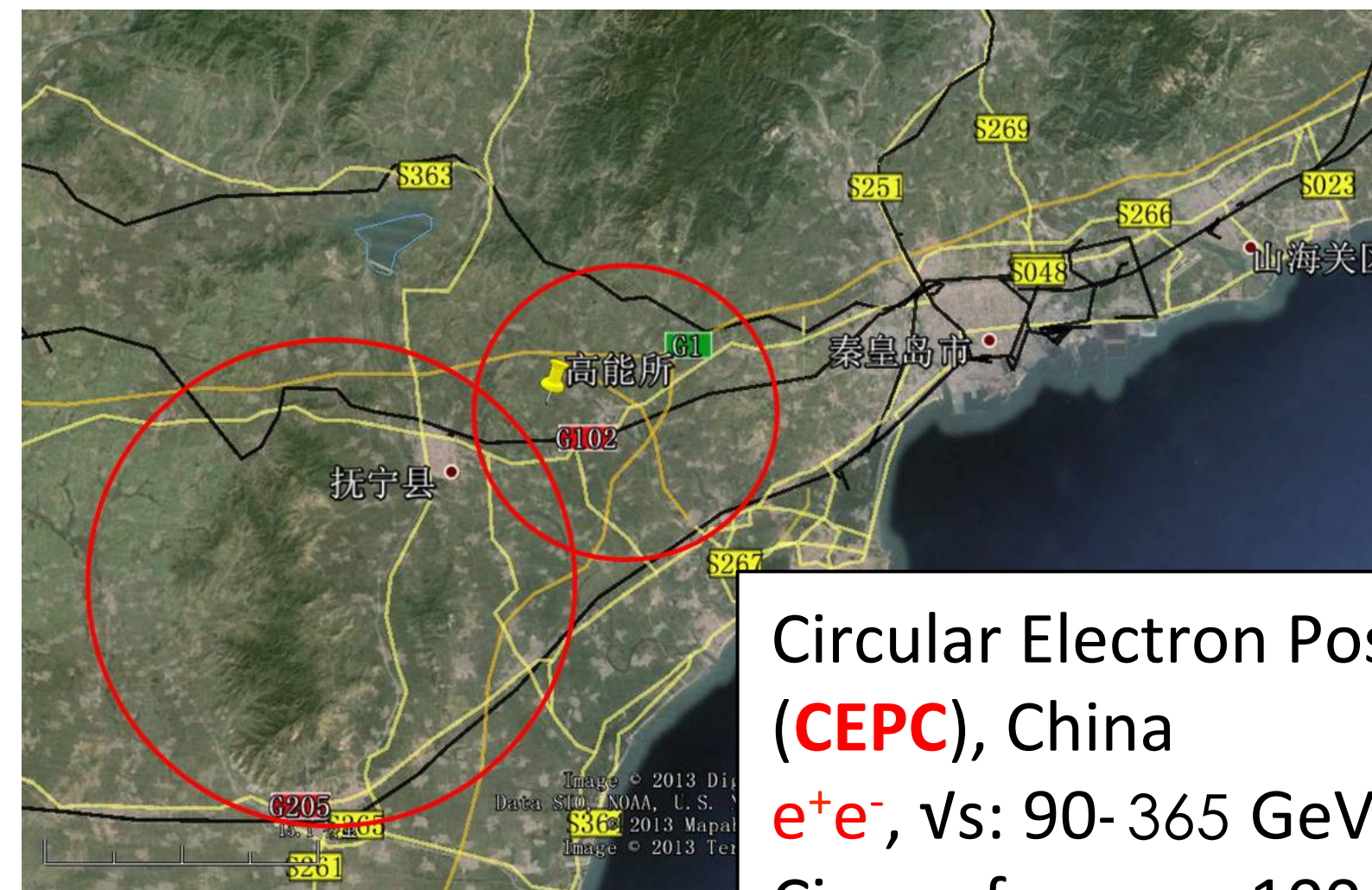
Compact Linear Collider (**CLIC**): CERN
 e^+e^- , vs: 380 GeV, 1.5 TeV, 3 TeV
 Length: 11 km, 29 km, 50 km



Future Circular Collider (**FCC-ee**): CERN
 e^+e^- , vs: 90 - 350 (365) GeV; FCC-hh pp
 Circumference: 97.75 km

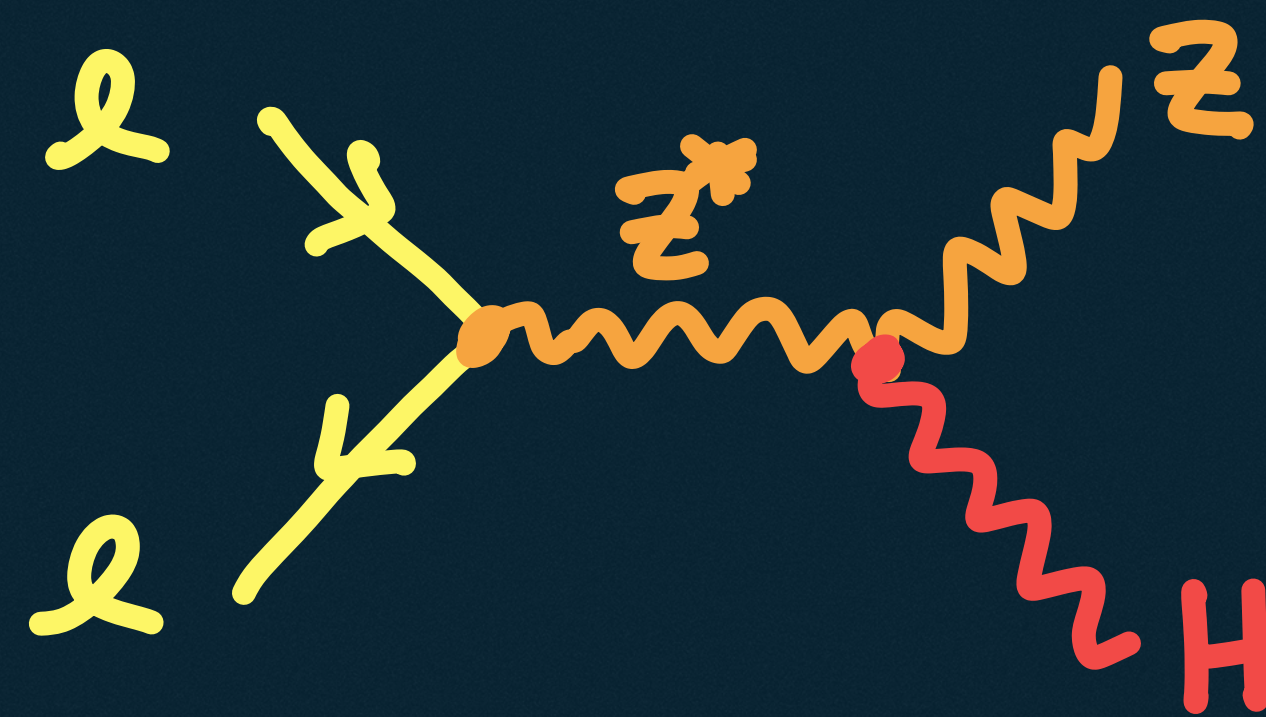
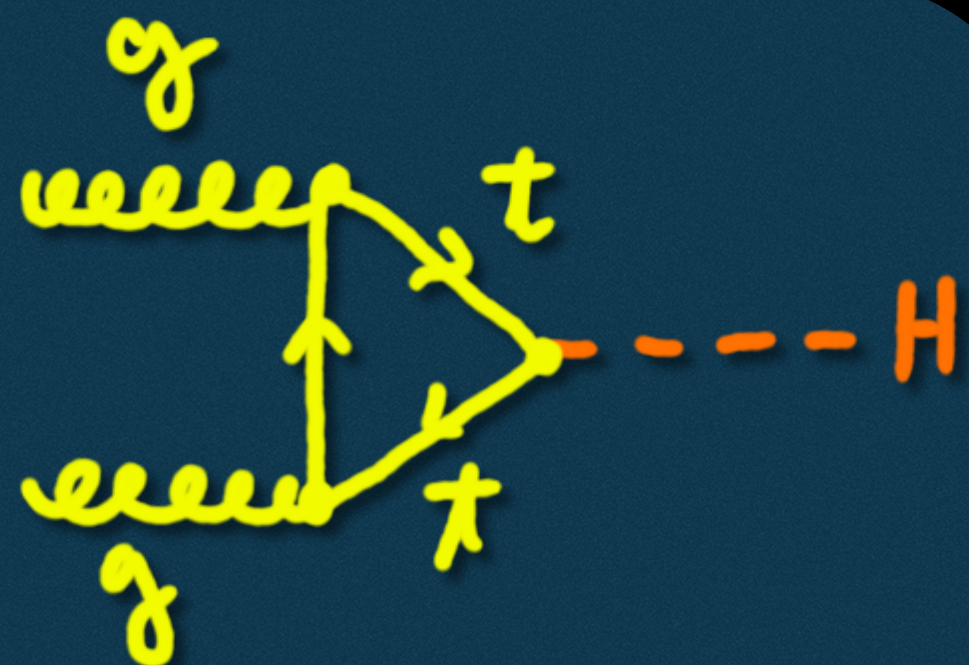
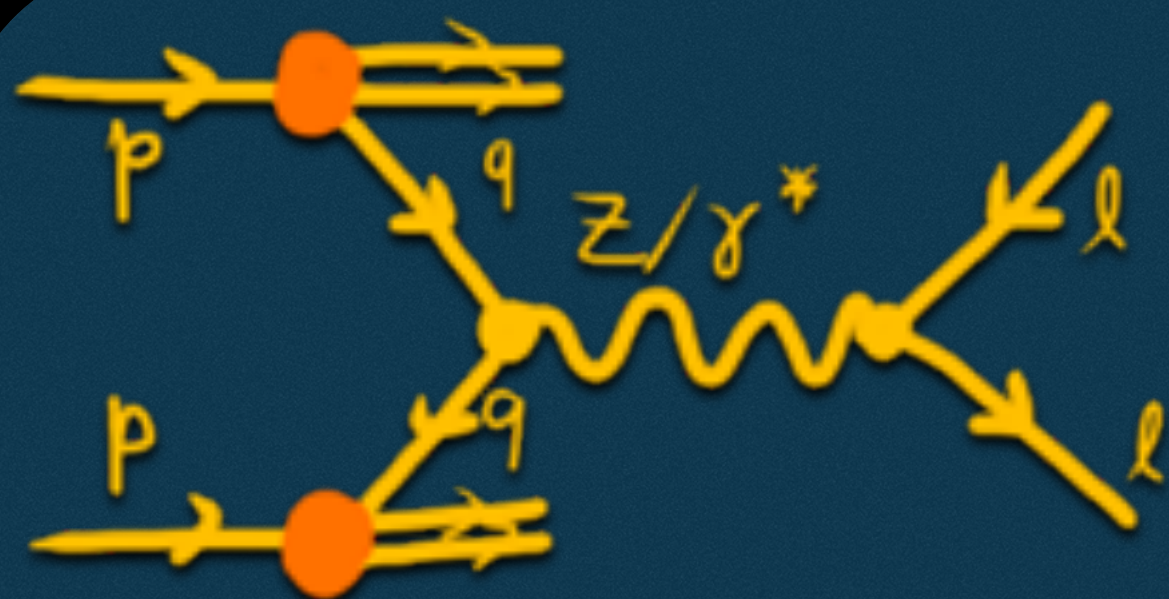


International Linear Collider (**ILC**):
 Japan (Kitakami)
 e^+e^- , vs: 250 - 500 GeV (1 TeV)
 Length: 17 km, 31 km (50 km)



Circular Electron Positron Collider (**CEPC**), China
 e^+e^- , vs: 90- 365 GeV; SPPC pp,
 Circumference: 100 km

Hadron versus lepton colliders



1. Proton are compound objects

- Initial state unknown (particle and momentum)
- Limits achievable precision

2. High rates of QCD background

- Complex triggers
- High levels of radiation
 - Detector design focus on radiation hardness of many sub-detectors

$$S/B \sim 10^{-10}$$

3. Very high-energy circular colliders feasible

1. Electrons are point-like particles

- Initial state well-defined (particle, energy, polarization...)
- High-precision measurements

2. Clean experimental environment

- No (less) need for triggers
- Lower levels of radiation

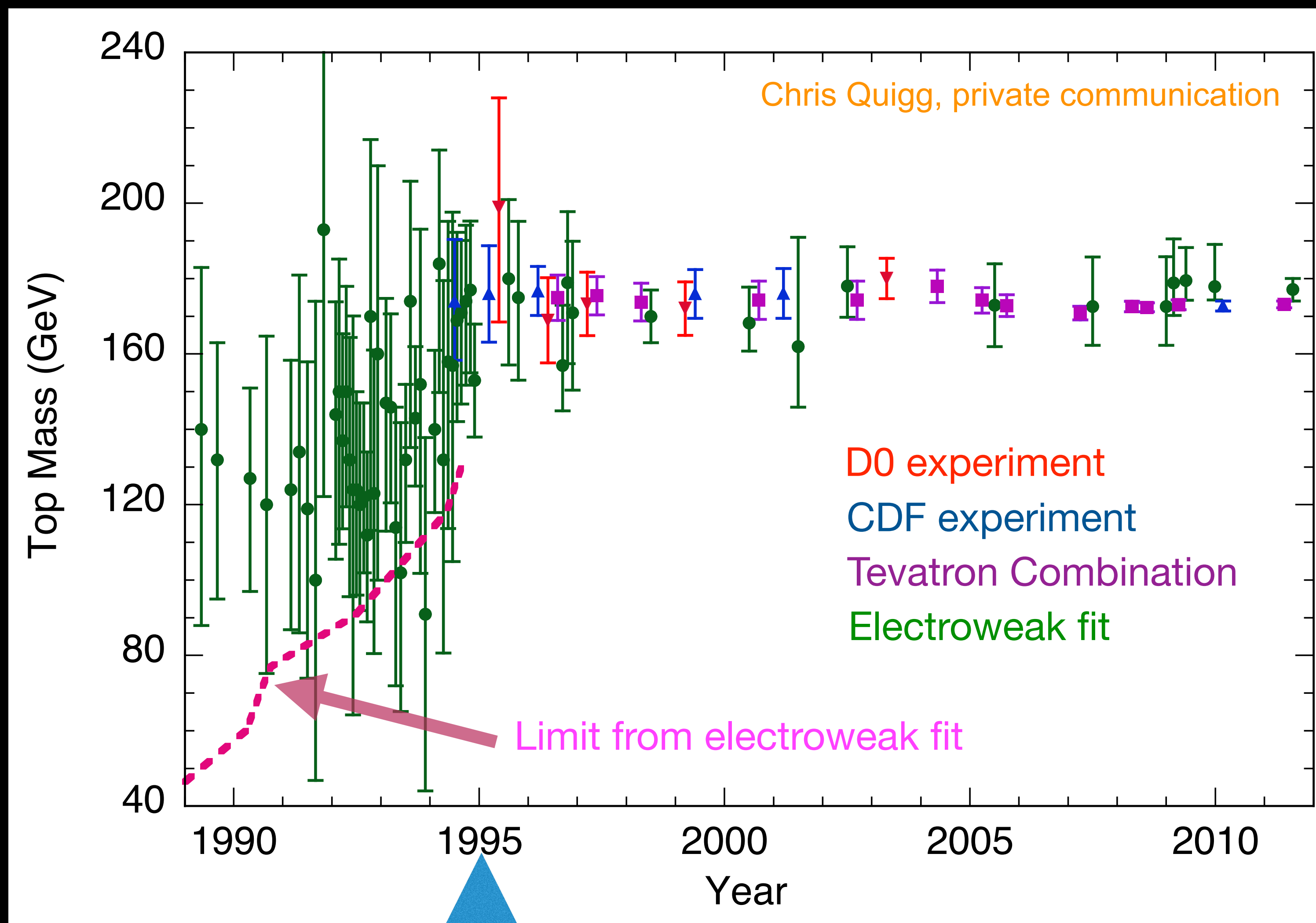
$$S/B \sim 10^{-3}$$

3. Very high-energies require linear colliders

After the Higgs boson discovery, no other new physics found
Need to also pursue outstanding precision

- **PRECISION** physics can play a key role -

Top Mass Prediction from Precision Electroweak data



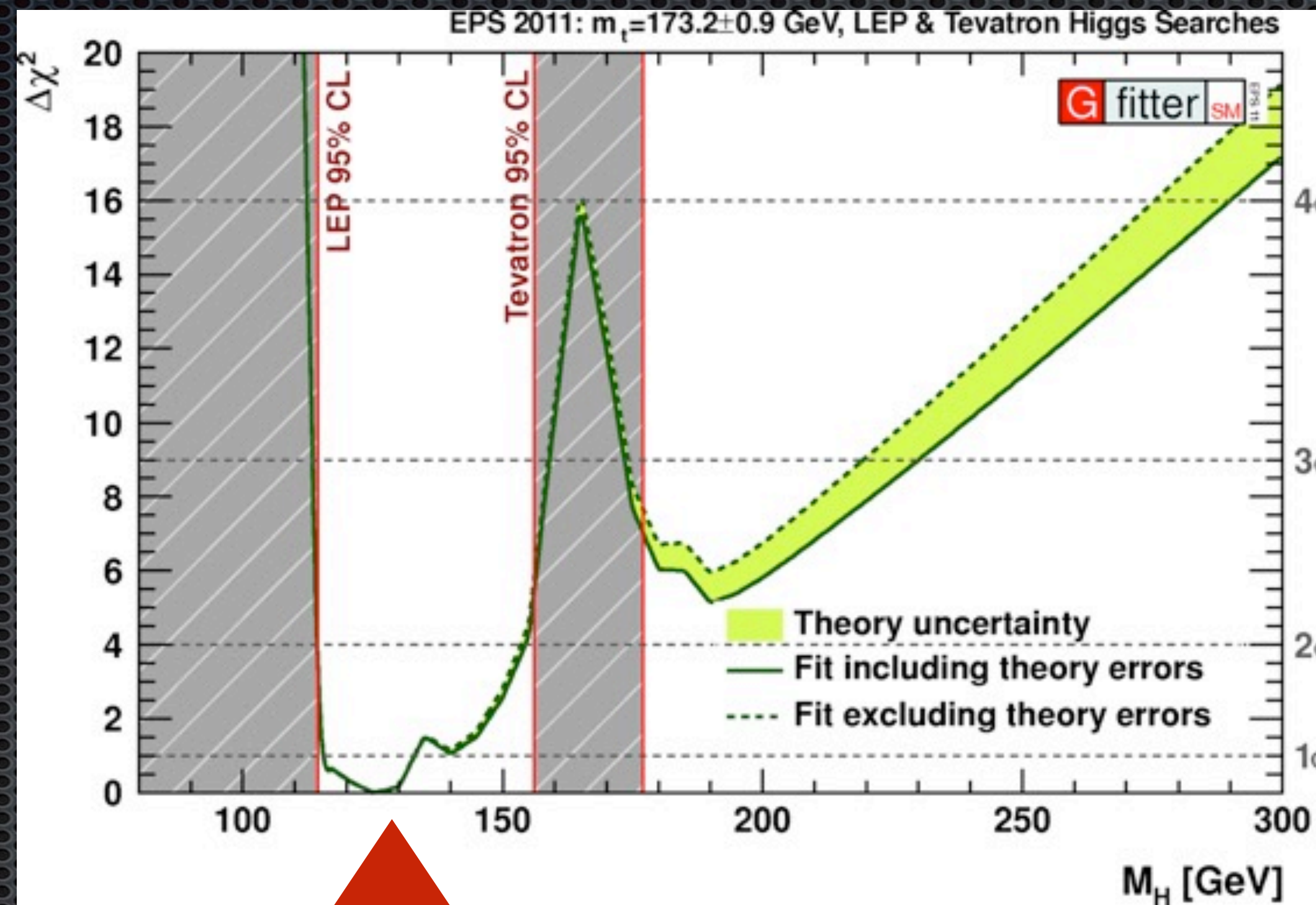
$M_{\text{top}} = 175 \rightarrow 173 \text{ GeV}$

World average:
 $m_{\text{top}} = 173.1 \pm 0.6 \text{ GeV}$
(0.35%)

Top discovery at Tevatron

Overnight update

- Updated with EPS'01 results
- Excludes direct searches from ATLAS and CMS from EPS



Standard Fit

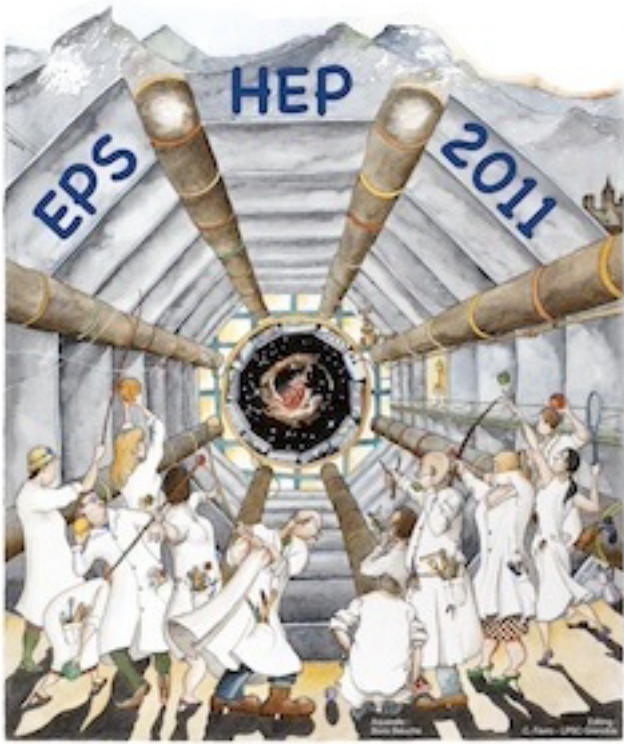
m_H (minimum) = 94.5 GeV, Range $m_H = [71, 124]$, $m_H < 166.5$ GeV @ 95%

Complete Fit

m_H (minimum) = 125.2 GeV, Range $m_H = [116, 133]$, $m_H < 153.9$ GeV @ 95%

Thanks to Matthias Schott from the GFitter group

WARNING:
Slide from
2011



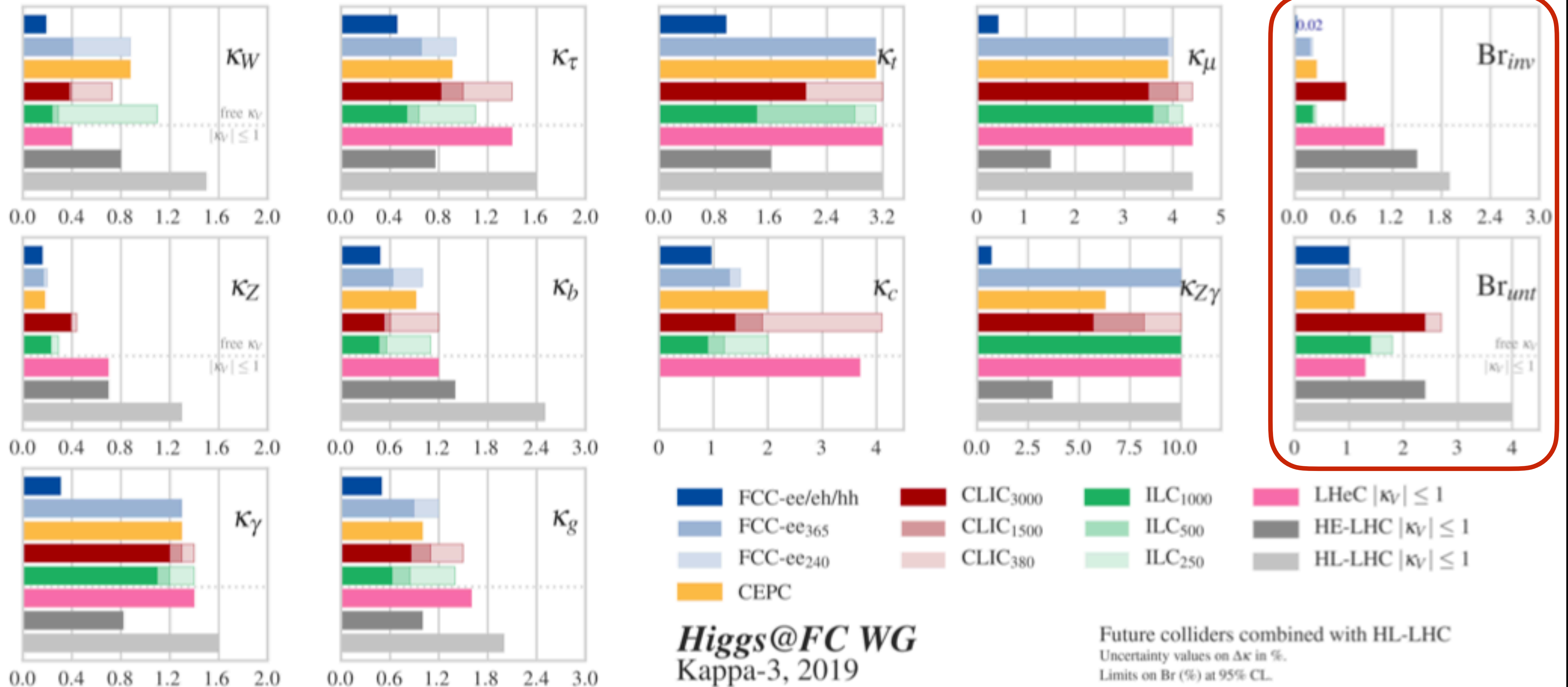
July 21-27

PANIC 2011, July 28, 2011

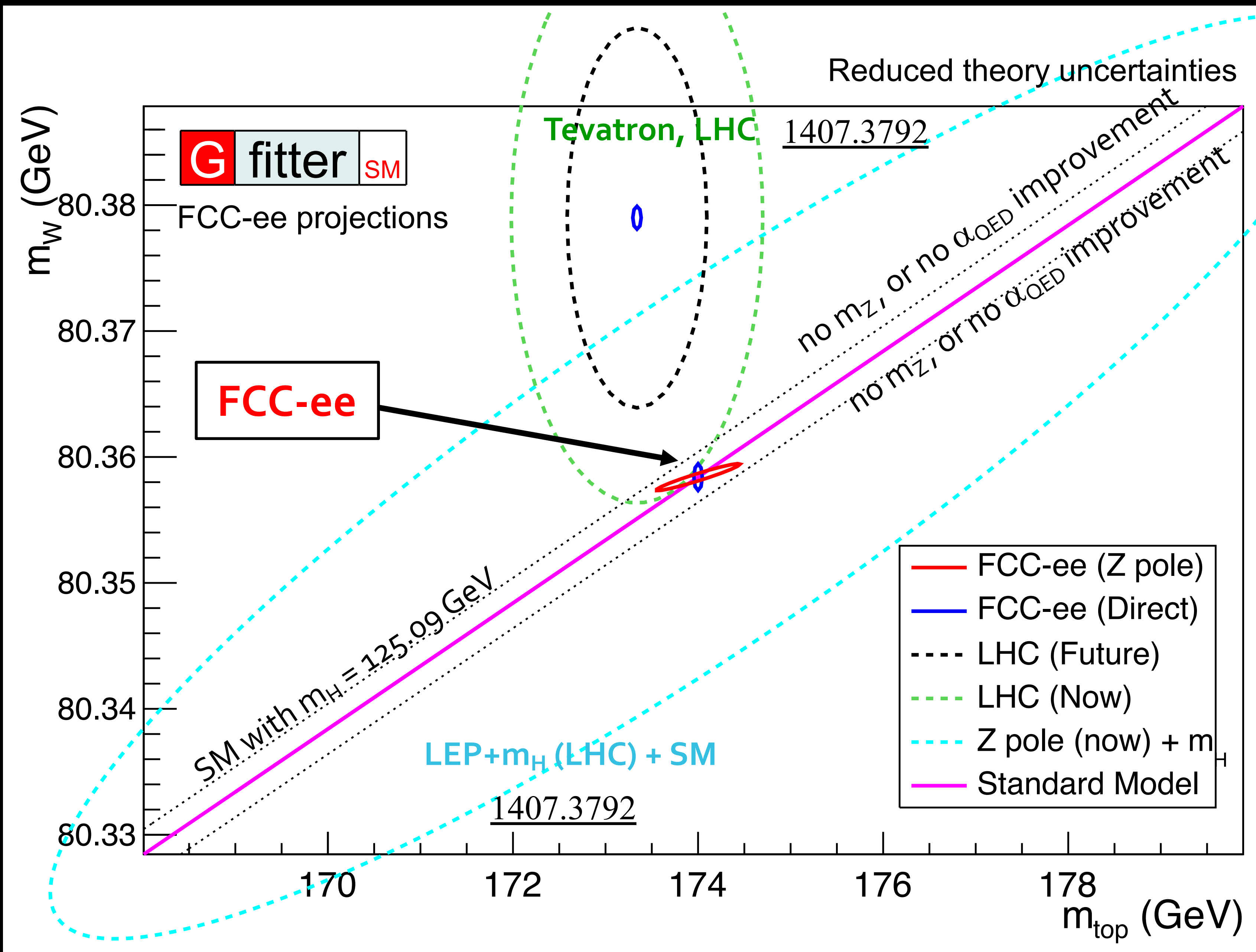
Higgs coupling measurement at future colliders

$$(\sigma \cdot \text{BR})(i \rightarrow H \rightarrow f) = \frac{\sigma_i^{\text{SM}} \kappa_i^2 \cdot \Gamma_f^{\text{SM}} \kappa_f^2}{\Gamma_H^{\text{SM}} \kappa_H^2} \rightarrow \mu_i^f \equiv \frac{\sigma \cdot \text{BR}}{\sigma_{\text{SM}} \cdot \text{BR}_{\text{SM}}} = \frac{\kappa_i^2 \cdot \kappa_f^2}{\kappa_H^2}$$

95% CL limits



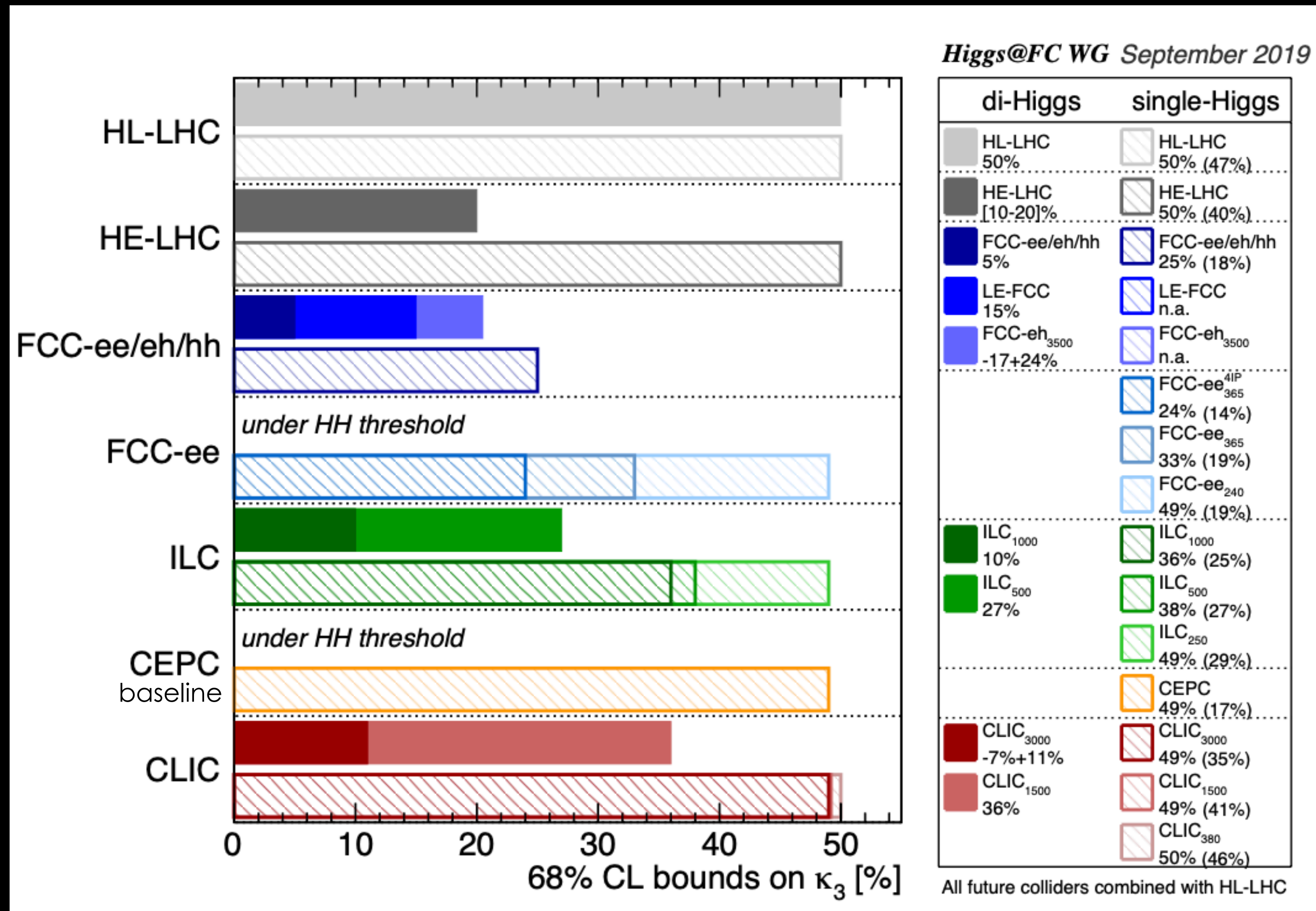
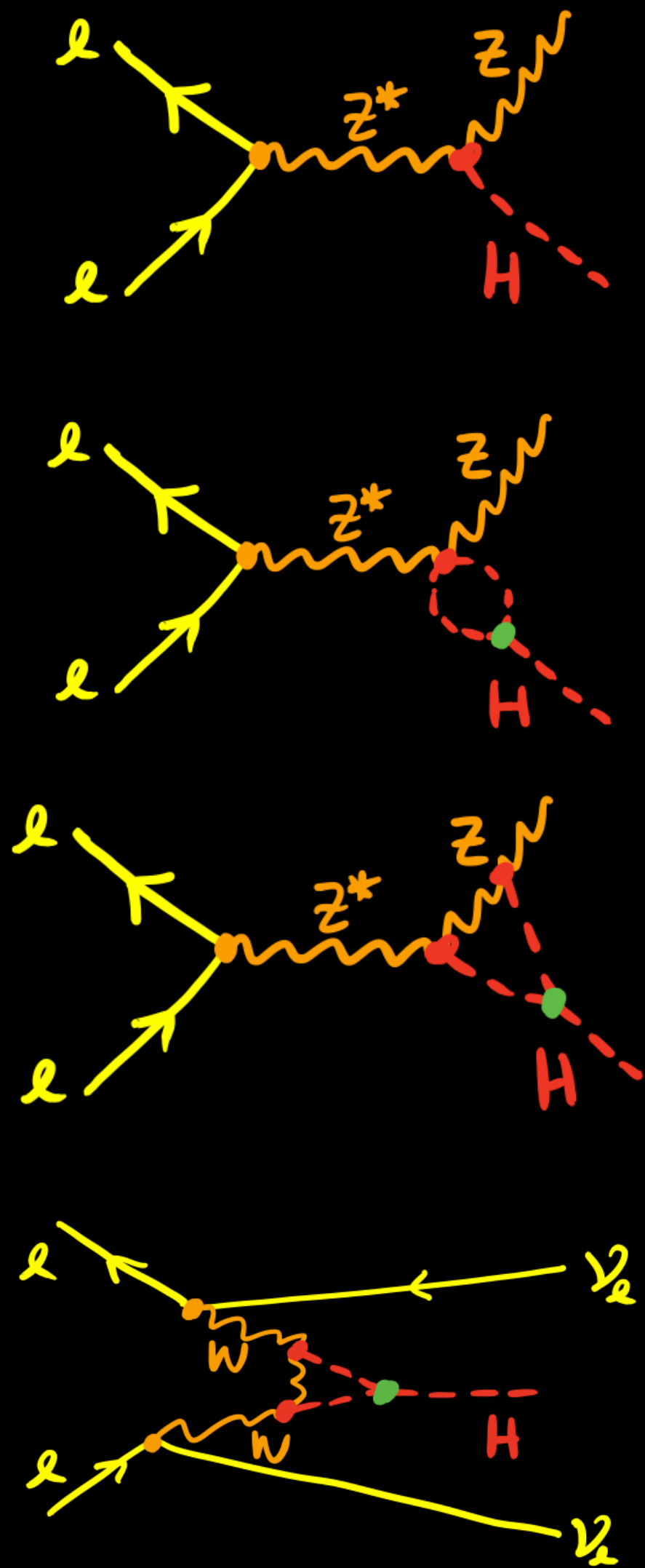
Combination of electroweak measurements



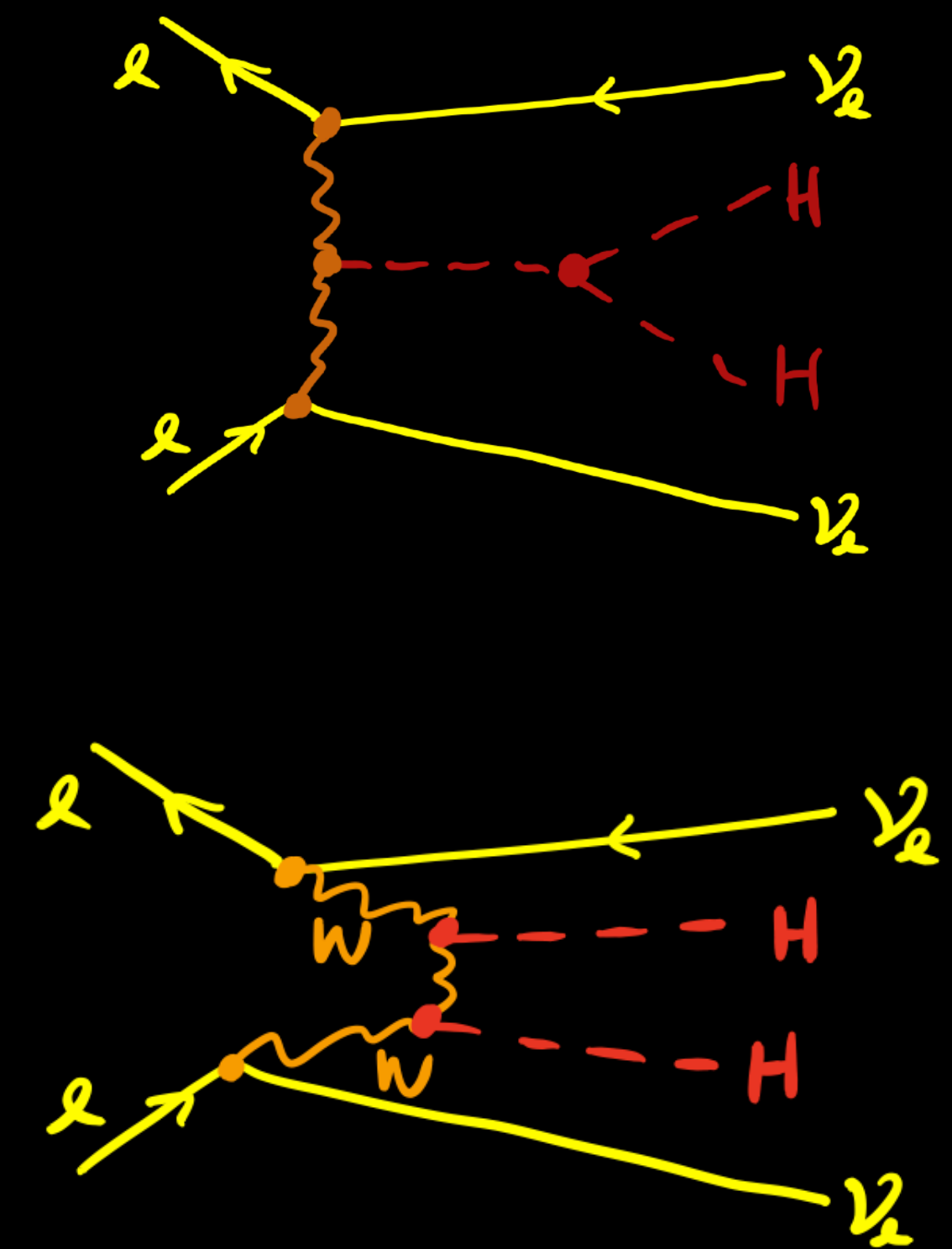
Precision of theory predictions needs to improve to take advantage of this experimental precision for full sensitivity to new physics

Higgs self-coupling at lepton colliders

Global effective-field-theory analysis can assess Higgs trilinear self-coupling



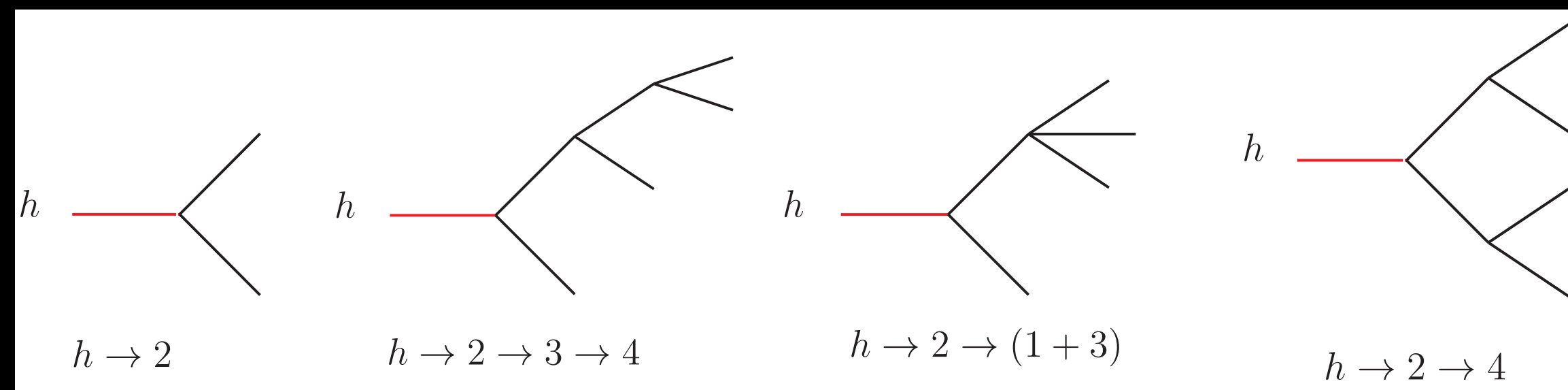
Direct production at ILC and CLIC most important



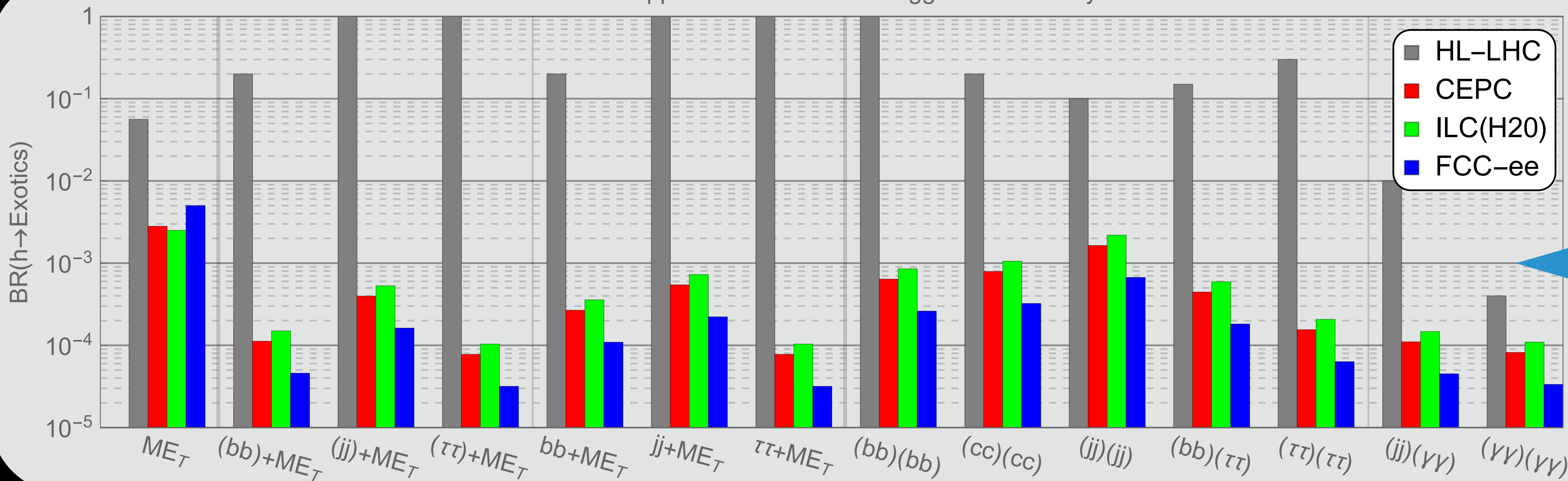
BSM Physics through Exotic Higgs Decays

General search for BSM

e^+e^- collider better than HL-LHC for MET+hadronic activity final states



95% C.L. upper limit on selected Higgs Exotic Decay BR



Final remarks

The standard Model is well established but many questions remain answered

LHC, including HL-LHC, to provide still 20-30x more data
Hence, small discrepancies can still turn into discoveries

Higgs self coupling measurement at HL-LHC is a major milestone

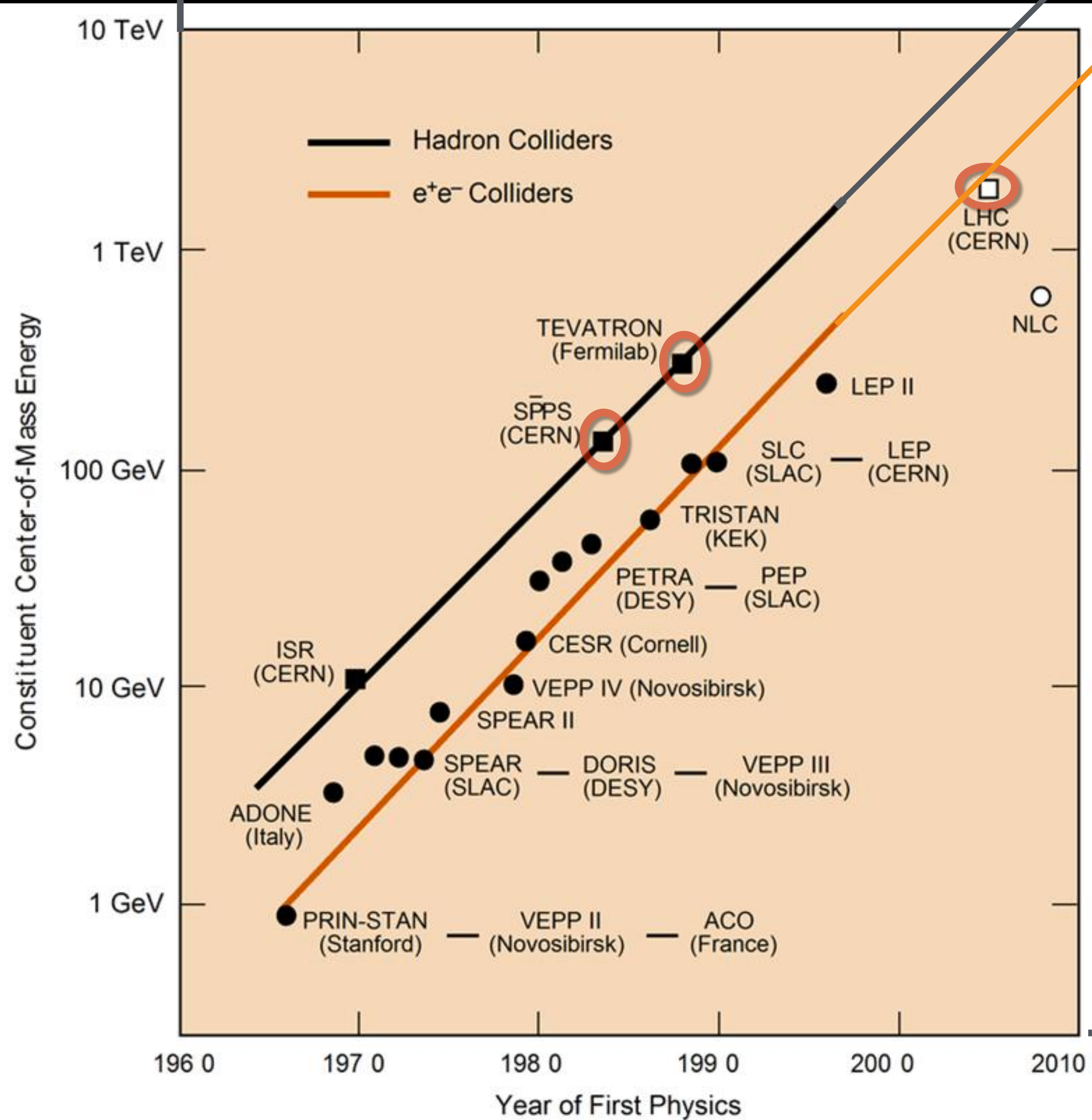
Planning for the next future accelerators has started

(see Tatsuya Tanaka's lecture tomorrow for details)

Expect new exciting sensitivity for BSM physics and potential discoveries

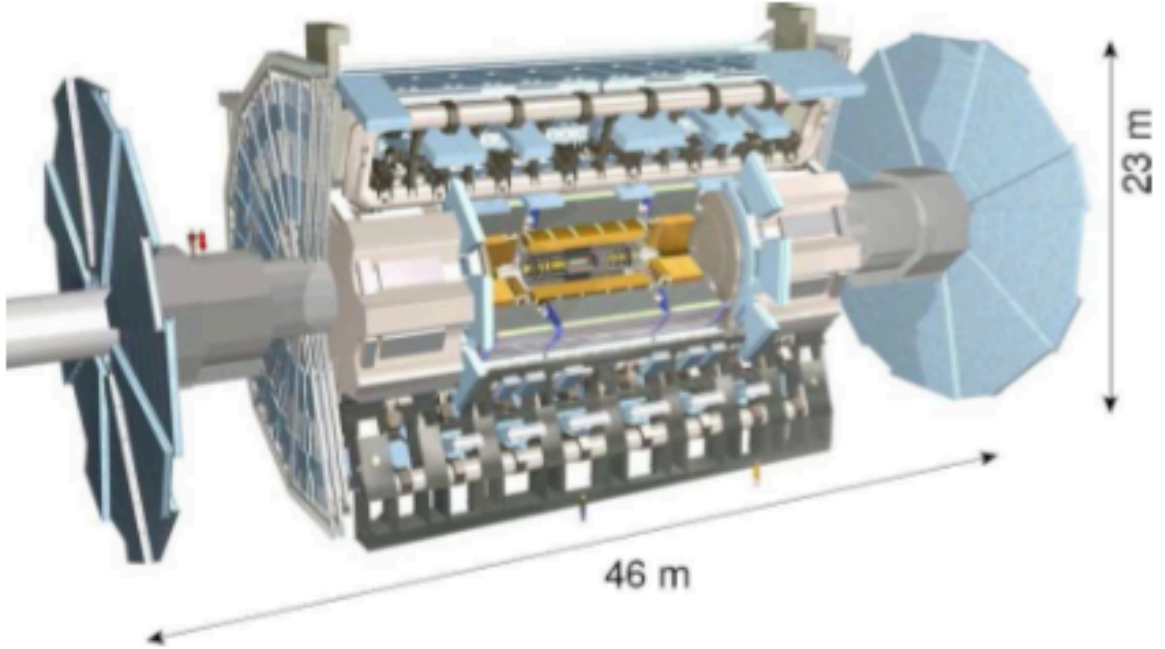
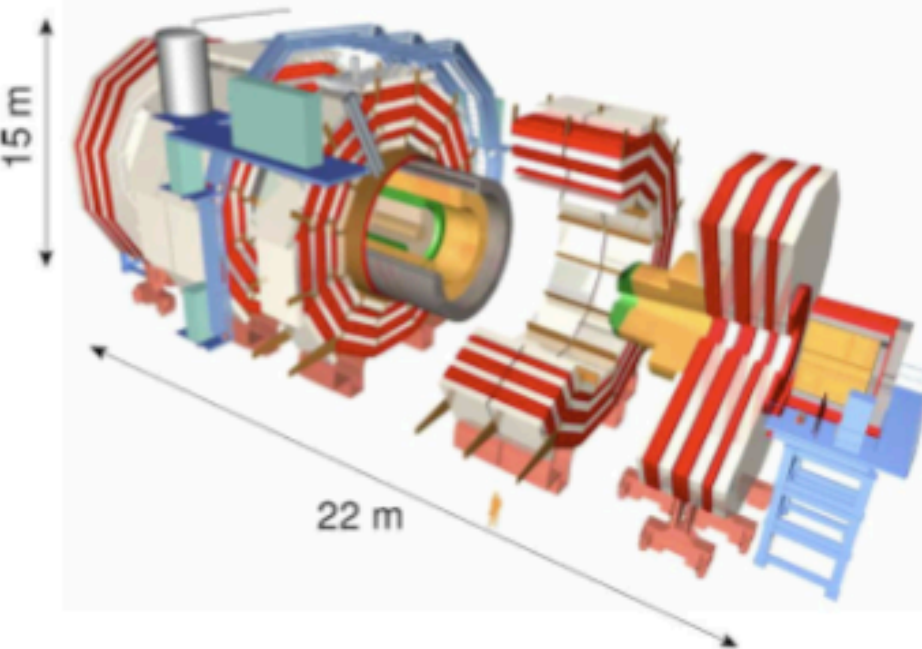
Higgs will continue to play an essential role

Explore synergies of electron and hadron colliders to maximize physics potential



2040 GeV

ATLAS and CMS in more detail

Sub System	ATLAS	CMS
Design		
Magnet(s)	Solenoid (within EM Calo) 2T 3 Air-core Toroids	Solenoid 3.8T Calorimeters Inside
Inner Tracking	Pixels, Si-strips, TRT PID w/ TRT and dE/dx $\sigma_{p_T}/p_T \sim 5 \times 10^{-4} p_T \oplus 0.01$	Pixels and Si-strips PID w/ dE/dx $\sigma_{p_T}/p_T \sim 1.5 \times 10^{-4} p_T \oplus 0.005$
EM Calorimeter	Lead-Larg Sampling w/ longitudinal segmentation $\sigma_E/E \sim 10\%/\sqrt{E} \oplus 0.007$	Lead-Tungstate Crys. Homogeneous w/o longitudinal segmentation $\sigma_E/E \sim 3\%/\sqrt{E} \oplus 0.5\%$
Hadronic Calorimeter	Fe-Scint. & Cu-Larg (fwd) $\gtrsim 11\lambda_0$ $\sigma_E/E \sim 50\%/\sqrt{E} \oplus 0.03$	Brass-scint. $\gtrsim 7\lambda_0$ & Tail Catcher $\sigma_E/E \sim 100\%/\sqrt{E} \oplus 0.05$
Muon Spectrometer System Acc. ATLAS 2.7 & CMS 2.4	Instrumented Air Core (std. alone) $\sigma_{p_T}/p_T \sim 4\%$ (at 50 GeV) $\sim 11\%$ (at 1 TeV)	Instrumented Iron return yoke $\sigma_{p_T}/p_T \sim 1\%$ (at 50 GeV) $\sim 10\%$ (at 1 TeV)

How bad does it look?

2012: $Z \rightarrow \mu\mu$ event

