Generalised Parton Distributions (GPDs) and Transverse Momentum-dependent Distributions (TMDs) at the Electron-Ion Collider

Daria Sokhan
University of Glasgow, Scotland, UK

Forward Physics and QCD at the LHC and EIC
Bad Honnef, Germany – 24th October 2023

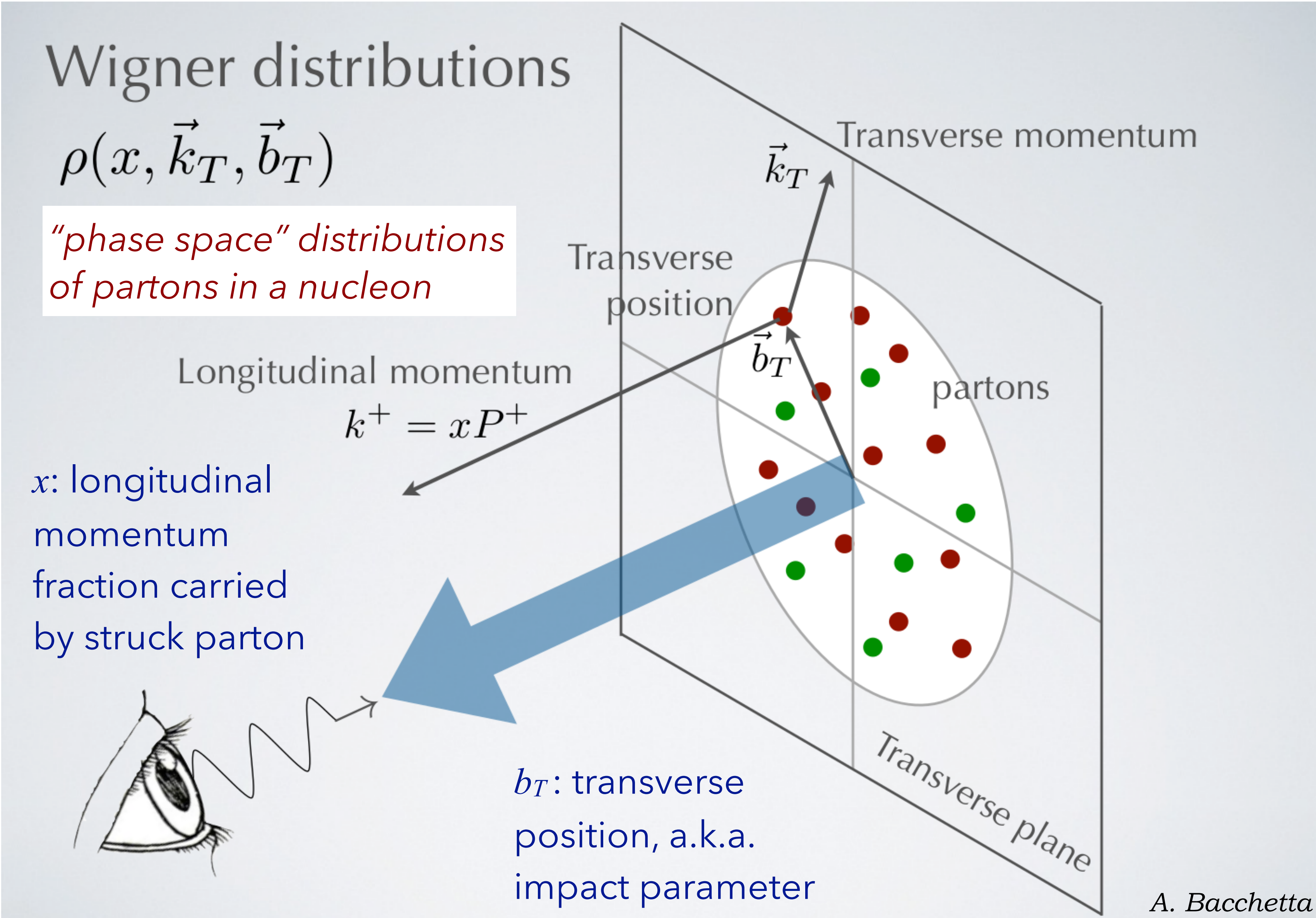
Generalised Parton Distributions (GPDs) and Transverse Momentum-dependent Distributions (TMDs) at the Electron-Ion Collider

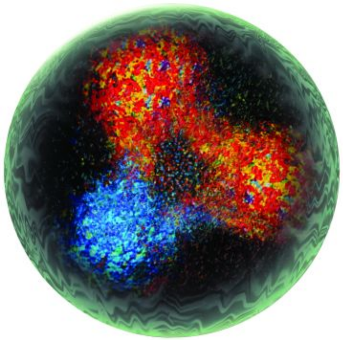
Daria Sokhan
University of Glasgow, Scotland, UK

Plots and studies presented are from many contributors to the Exclusive, Diffractive & Tagging WG of the Yellow Report, ATHENA, ECCE and ePIC
(not all acknowledged personally)

Forward Physics and QCD at the LHC and EIC
Bad Honnef, Germany – 24th October 2023

A constructivist view of the nucleon





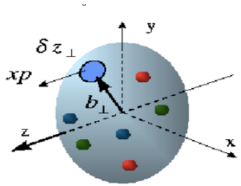
Wigner function: full phase space parton distribution of the nucleon

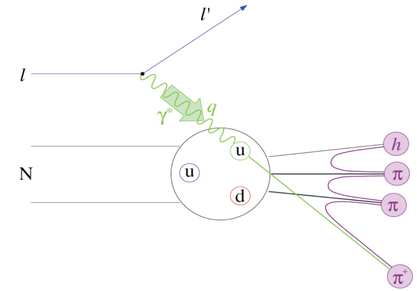
Generalised Transverse Momentum Distributions (GTMDs)

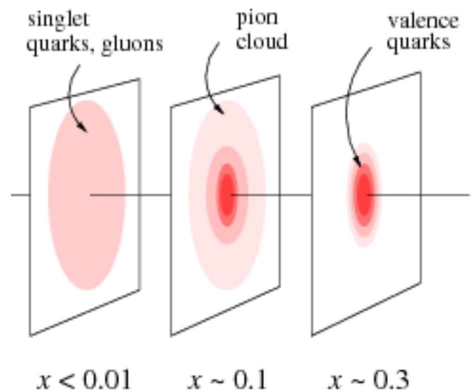
Possible access via exclusive di-jet production or exclusive π⁰-production at high Q^2.

$\int d^2k_T$

$\int d^2b_T$





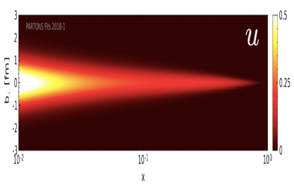


Generalised Parton Distributions (GPDs)

Exclusive processes

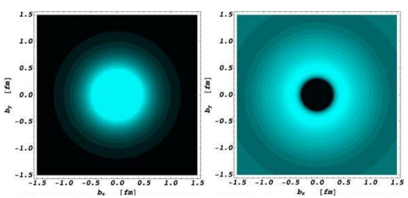
Transverse Momentum-Dependent distributions (TMDs)

Semi-inclusive DIS (SIDIS)



$\int dx$

$\int d^2k_T$

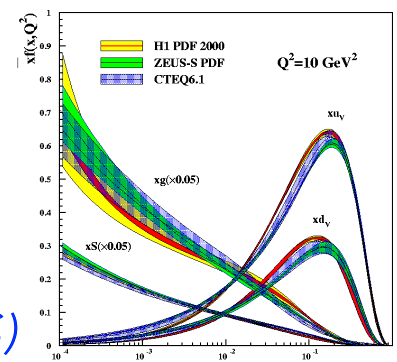


Form Factors

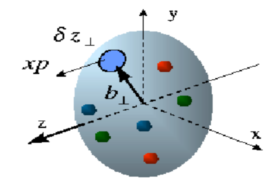
Elastic scattering

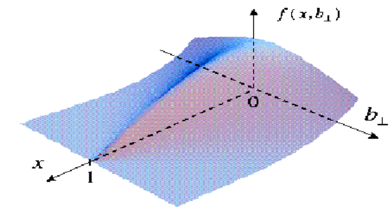
Parton Distribution Functions (PDFs)

Deep Inelastic Scattering (DIS)



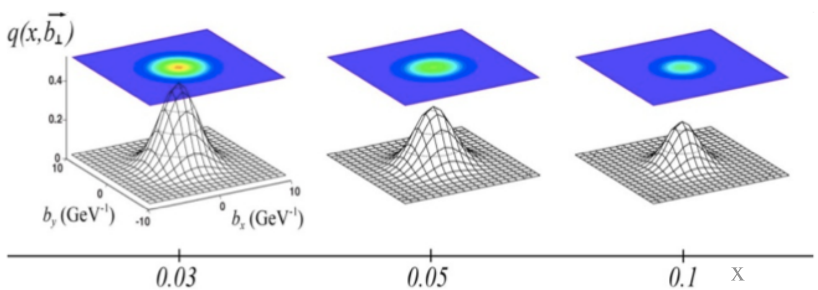
Generalised Parton Distributions





- proposed by Müller (1994), Radyushkin, Ji (1997).
- can be interpreted as relating, in the infinite momentum frame, transverse position of partons (impact parameter $b_\perp$) to longitudinal momentum fraction (x).

* **Tomography** of the nucleon: transverse spatial distributions of quarks and gluons in longitudinal momentum space.



* Information on the orbital angular momentum contribution to nucleon spin: **the spin puzzle**.

$$J_N = \frac{1}{2} = \frac{1}{2}\Sigma_q + L_q + J_g$$

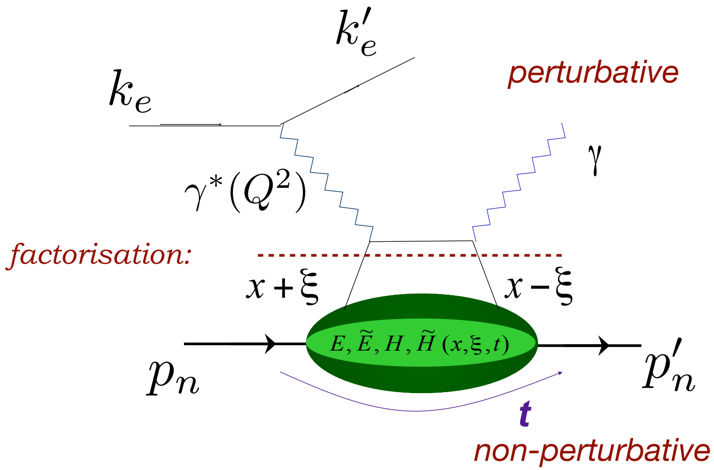
Ji's relation:

$$J^q = \frac{1}{2} - J^g$$
$$= \frac{1}{2}\int_{-1}^{1} x dx \left\{H^q(x,\xi,0) + E^q(x,\xi,0)\right\}$$

* Indirect access to mechanical properties of the nucleon: possibilities of extracting **pressure distributions** within the nucleon.

* Combine with TMDs to access **spin-orbit correlations** of quarks and gluons, study non-perturbative interactions of partons.

Deeply Virtual Compton scattering



"Handbag" diagram

$$Q^2 = -(\mathbf{k} - \mathbf{k}')^2 \qquad t = (\mathbf{p}'_\mathbf{n} - \mathbf{p}_\mathbf{n})^2$$

Bjorken variable: $x_B = \dfrac{Q^2}{2\mathbf{p}_\mathbf{n} \cdot \mathbf{q}}$

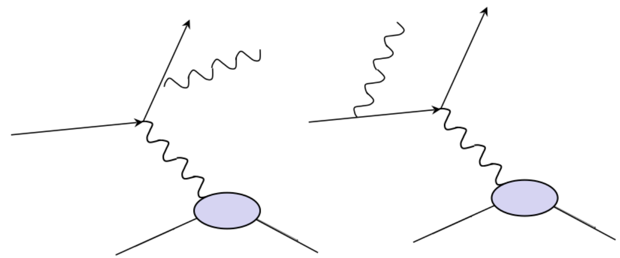
$x \pm \xi$ *longitudinal momentum fractions of the struck parton*

Skewness: $\xi \cong \dfrac{x_B}{2 - x_B}$

- At high exchanged Q^2 and low t sensitivity to four parton helicity-conserving, chiral-even GPDs:

$$E^q, \tilde{E}^q, H^q, \tilde{H}^q(x, \xi, t)$$

- Experimentally, measure DVCS, Bethe-Heitler and their interference:



$$d\sigma \propto |T_{DVCS}|^2 + |T_{BH}|^2 + T_{BH}T^*_{DVCS} + T_{DVCS}T^*_{BH}$$

- Observables are parametrised in terms of Compton Form Factors (CFFs): complex functions where $\mathcal{Re}$ parts are integrals of GPDs over x and $\mathcal{Im}$ parts are GPDs at $x = \pm\xi$

Experimental access to GPDs

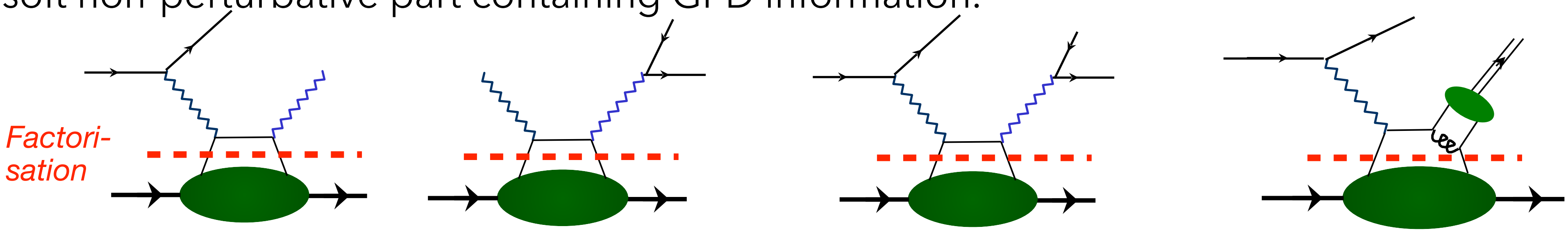
Accessible via CFFs and structure functions in *exclusive* processes (all final state particles are determined), *eg*:

Exclusive @ ePIC: talk by Alex Jentsch (Wed)

- Deeply Virtual Compton Scattering (DVCS)
- Time-like Compton Scattering (TCS)
- Hard Exclusive Meson Production (HEMP) – a.k.a. Deeply Virtual Meson Production (DVMP)
- Double DVCS
- Certain diffractive processes, eg: diffractive ρ-production with the emission of a meson or virtual photon from the nucleon
- Hard exclusive production of a meson-photon or photon-photon pair
- Charged-current meson production, eg: $ep \to \nu_e \pi^- p$

See EIC Yellow Report for details

Relies on *factorisation* of the process amplitude into a hard, perturbative part and the soft non-perturbative part containing GPD information.



DVCS — *Virtual photon space-like*

TCS — *Virtual photon time-like*

DDVCS — *One time-like, one space-like virtual photon*

DVMP — *Virtual photon space-like*

Electron-Ion Collider

World's first polarized electron-proton/light ion and electron-Nucleus collider.

For e-N collisions at the EIC:

- ✓ Polarized beams: e, p, d/^{3}He
- ✓ e beam 5 - 18 GeV
- ✓ Luminosity L_{ep} ~ 10^{33-34} cm^{-2}s^{-1}
- ✓ 30 - 140 GeV variable CoM

EIC: talk by Silvia dalla Torre (Mon)

For e-A collisions at the EIC:

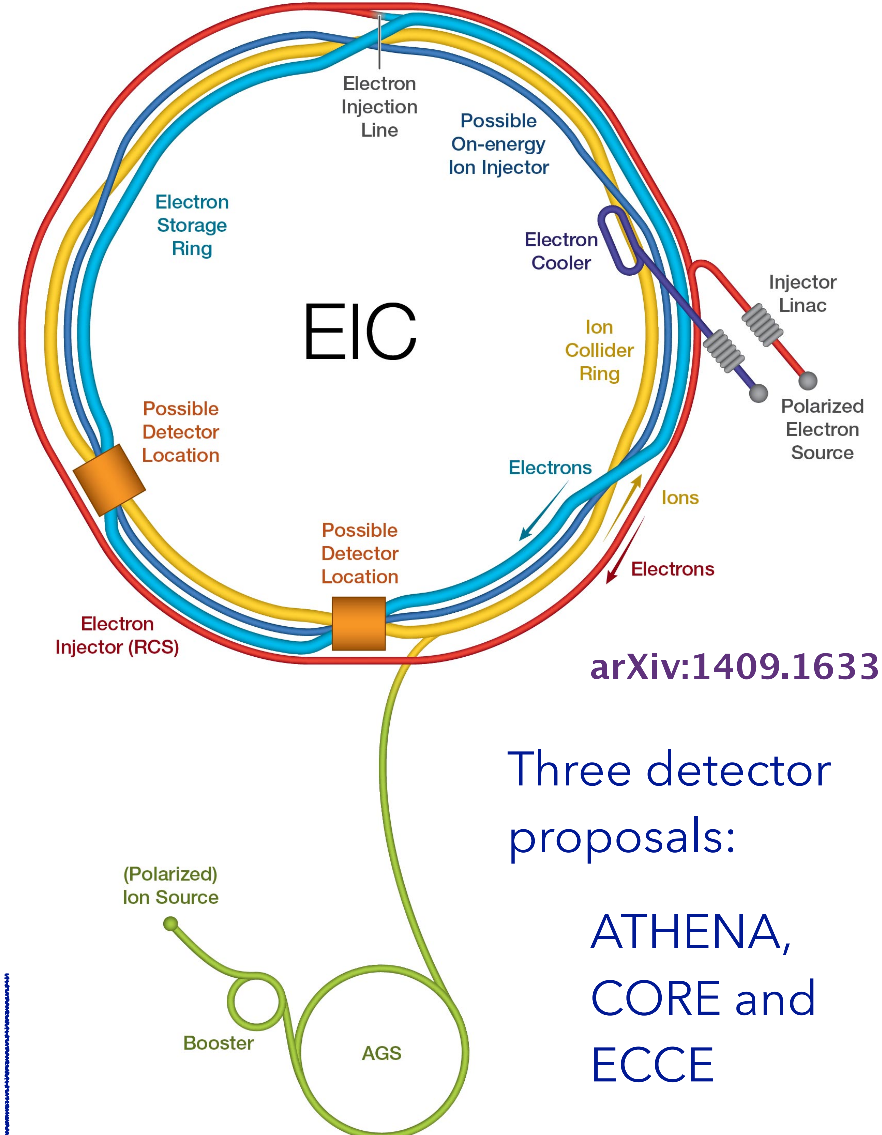
- ✓ Wide range of nuclei
- ✓ Luminosity per nucleon same as e-p
- ✓ Variable centre of mass energy

Brookhaven National Lab selected as the site

Dedicated studies of EIC physics and design:

EIC White Paper, Eur. Phy. J. A 52, 9 (2016)

EIC Yellow Report, Nuc. Phys. A 1026, 122447 (2022)



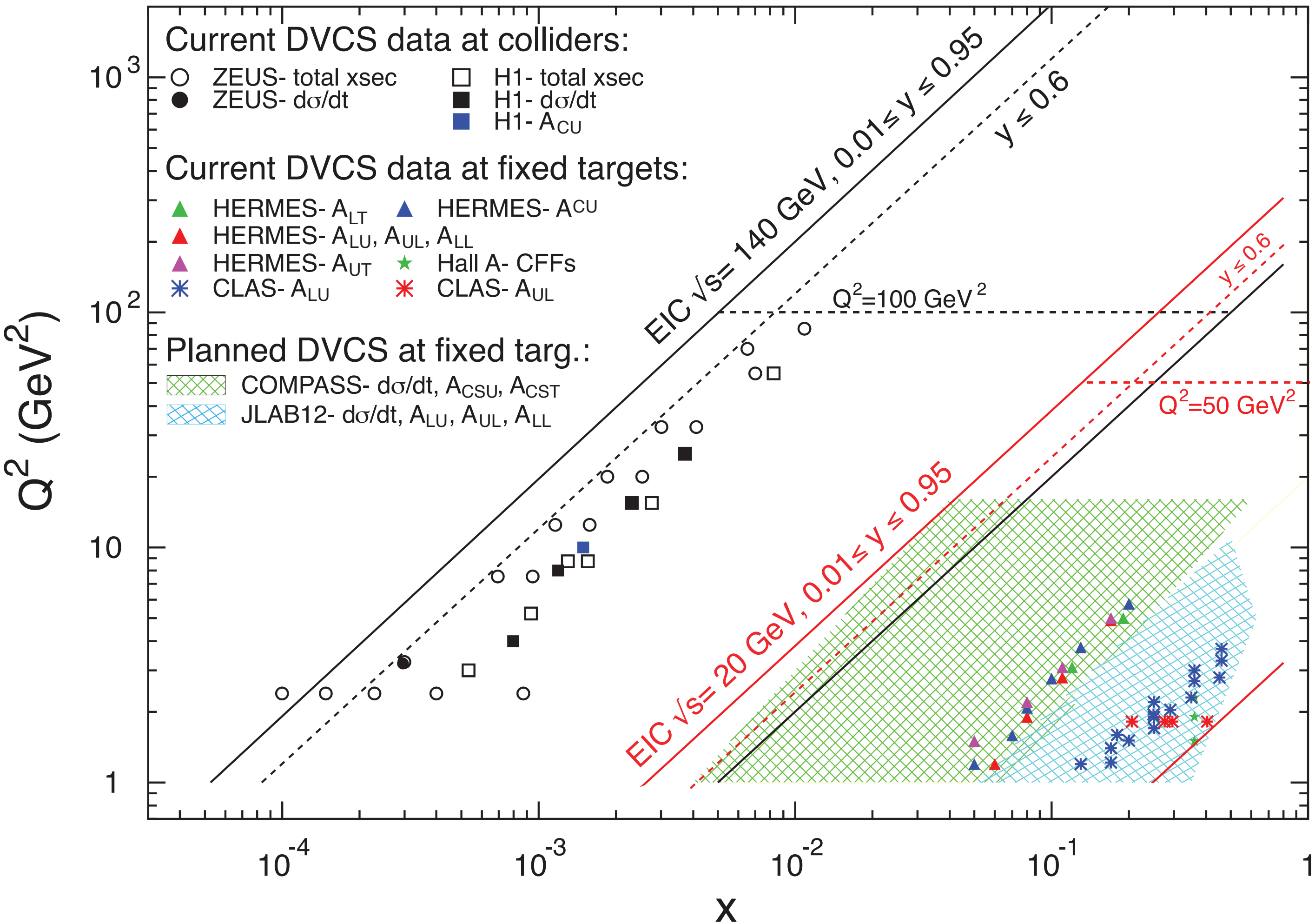
arXiv:1409.1633

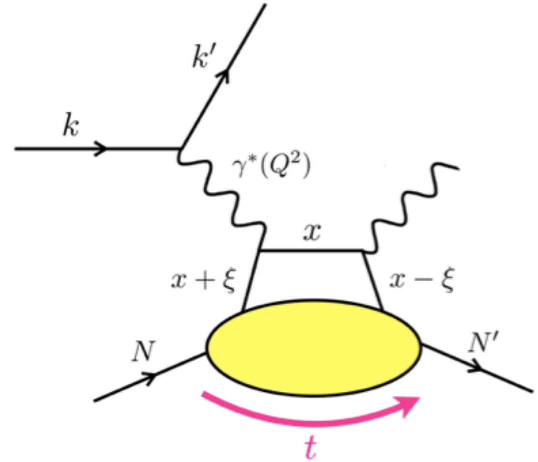
Three detector proposals:

ATHENA, CORE and ECCE

Merged to form the new "Detector 1" collaboration: ePIC

EIC kinematic reach: DVCS





$$Q^2 = -(\mathbf{k} - \mathbf{k}')^2$$

y: inelasticity

https://arxiv.org/abs/2103.05419

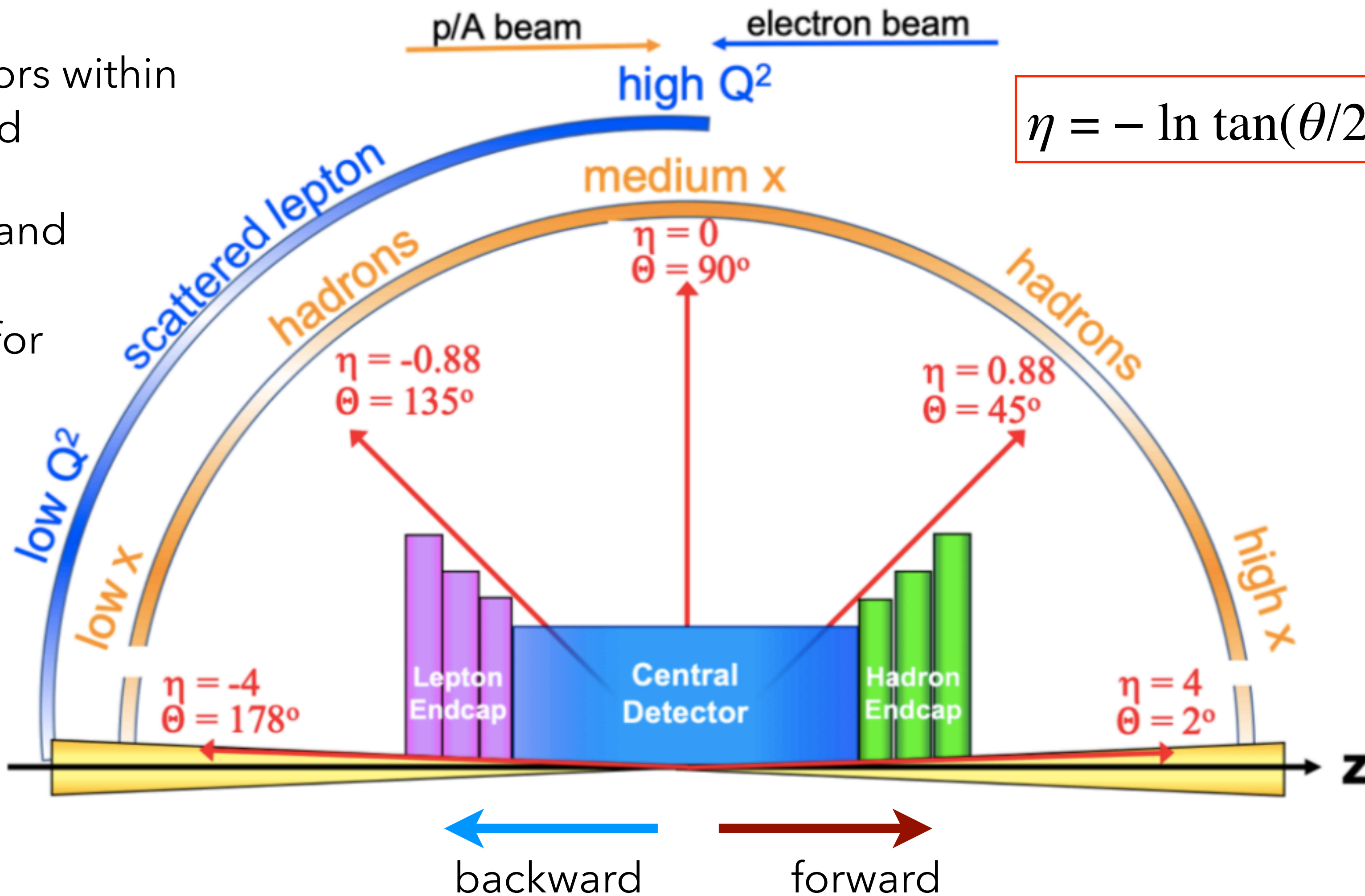
Detector configuration

Very asymmetric beams

Hermetic detectors within a central solenoid

Very far-forward and far-backward instrumentation for lowest scattering angles.

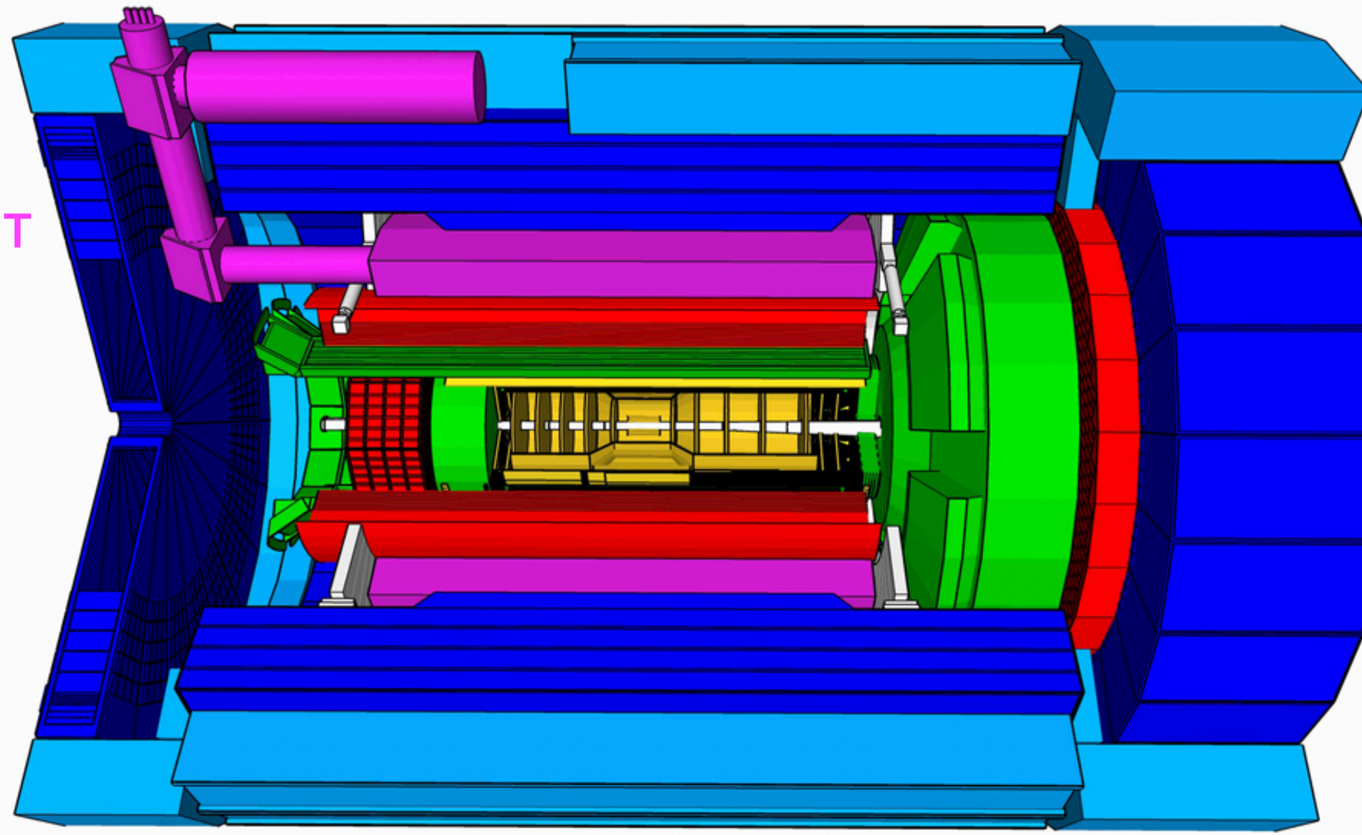
$$\eta = -\ln\tan(\theta/2)$$



EIC: talk by Silvia dalla Torre (Mon)

The ePIC detector

Electron-Proton and -Ion Collider detector

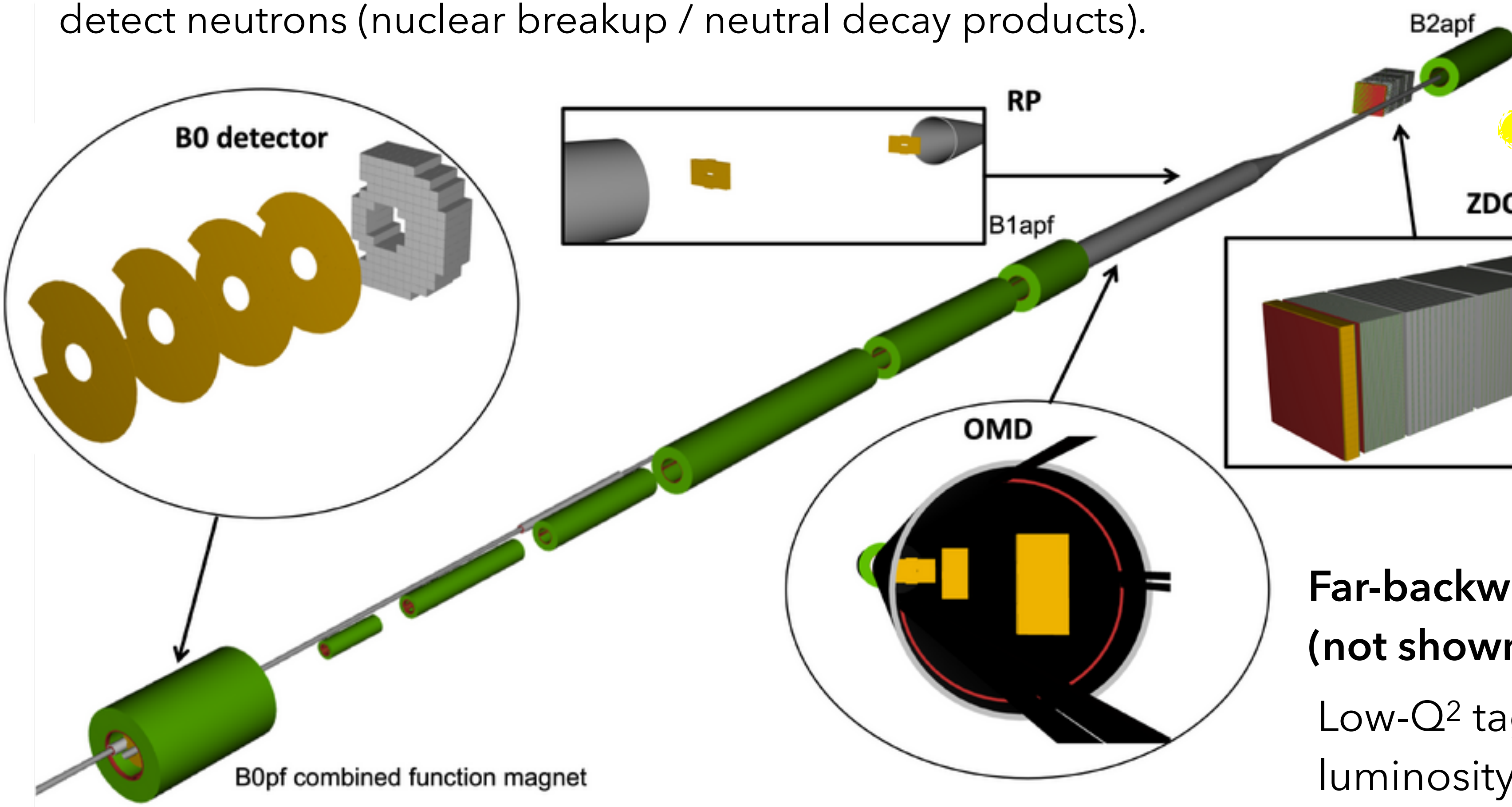


Result of the merging of ECCE and ATHENA collaborations.

Far-forward and -backward detectors

Far-forward detectors: Far from interaction point, very low angles.

Roman Pots inside the beam pipe, Off-Momentum Detector (OMD) for nuclear break-up and B0 tracker for larger angles, large acceptance Zero degree Calorimeter (ZDC) to detect neutrons (nuclear breakup / neutral decay products).



Exclusive @ EIC: talk by Alex Jentsch (Wed)

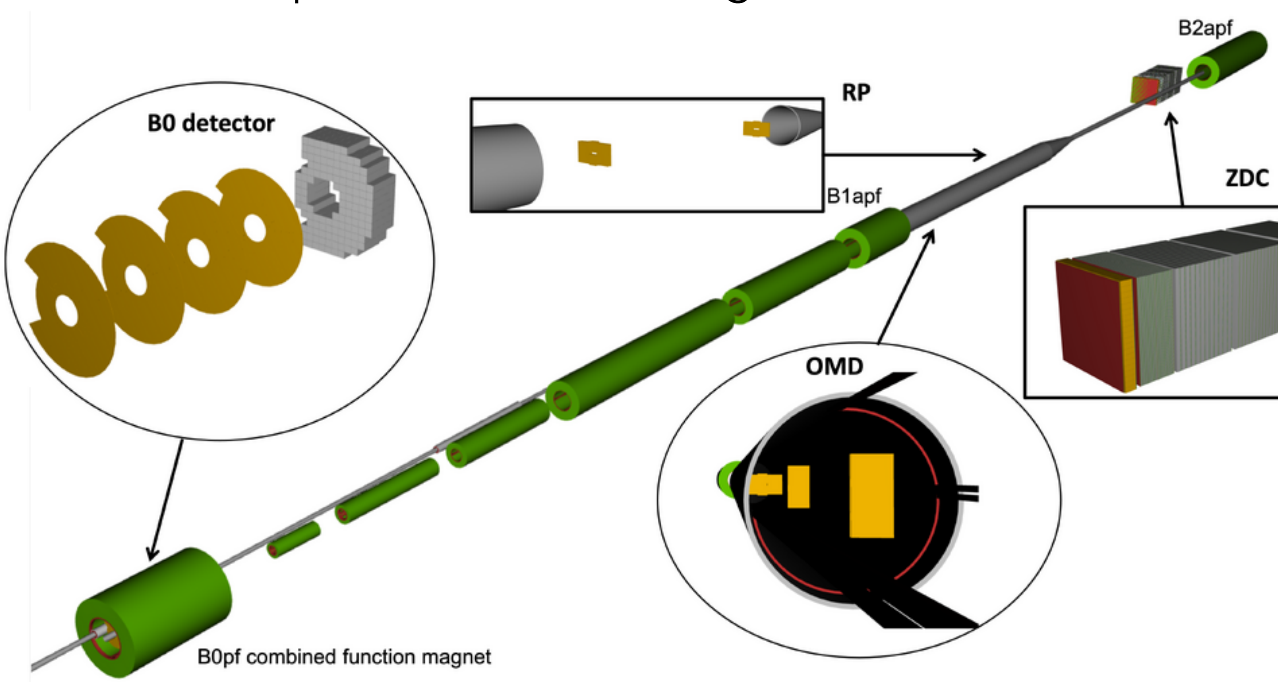
Far-backward detectors (not shown):

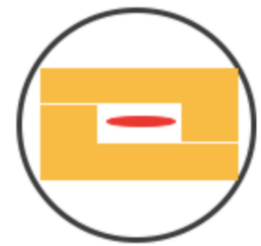
Low-Q^2 tagger, luminosity monitor

Recoil protons in *ep*

- The impact parameter information in many exclusive processes is encoded in *t*, via a Fourier Transform. Require accurate measurement of *t* from as close to zero as possible and across a wide range in *ep and e(light-A)* collisions.

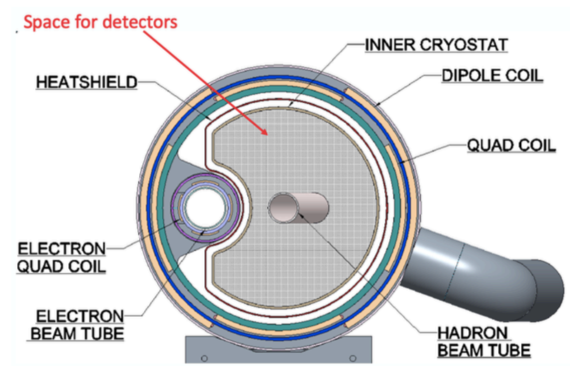
- Scattered protons / light ions detected in Roman Pots (for the lowest values of *t*) and in the B0 spectrometers (for higher values).





Hadron beam-pipe and Roman Pots in cross-section

$$t = (p' - p)^2$$

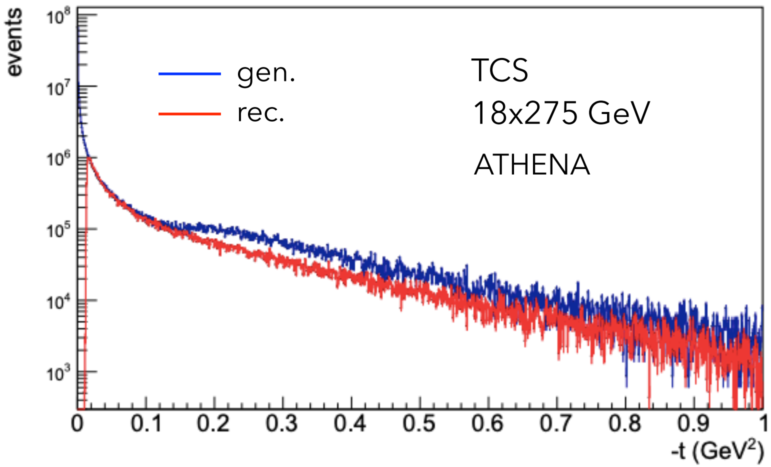


EIC Yellow Report

B0 spectrometer configuration

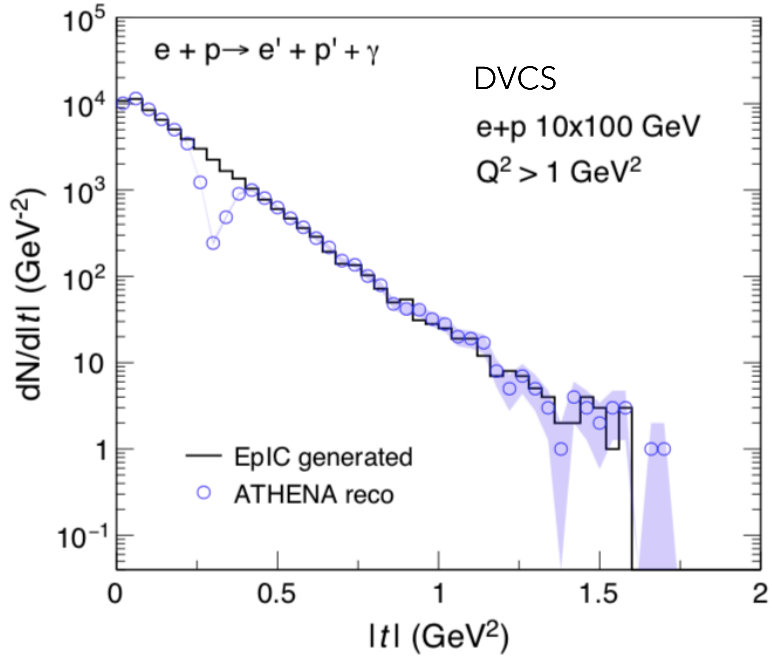
Recoil protons in *ep*

- Scattered protons detected in Roman Pots (for the lowest values of t) and in the B0 spectrometers (for higher values).



Light ions bend less and the t-distribution drops faster: detection entirely in the Roman Pots.

Can study coherent DVCS on He-4: spin-0, so coherent amplitude parametrised by one GPD. Also deuteron (but spin-1, too many GPDs!)

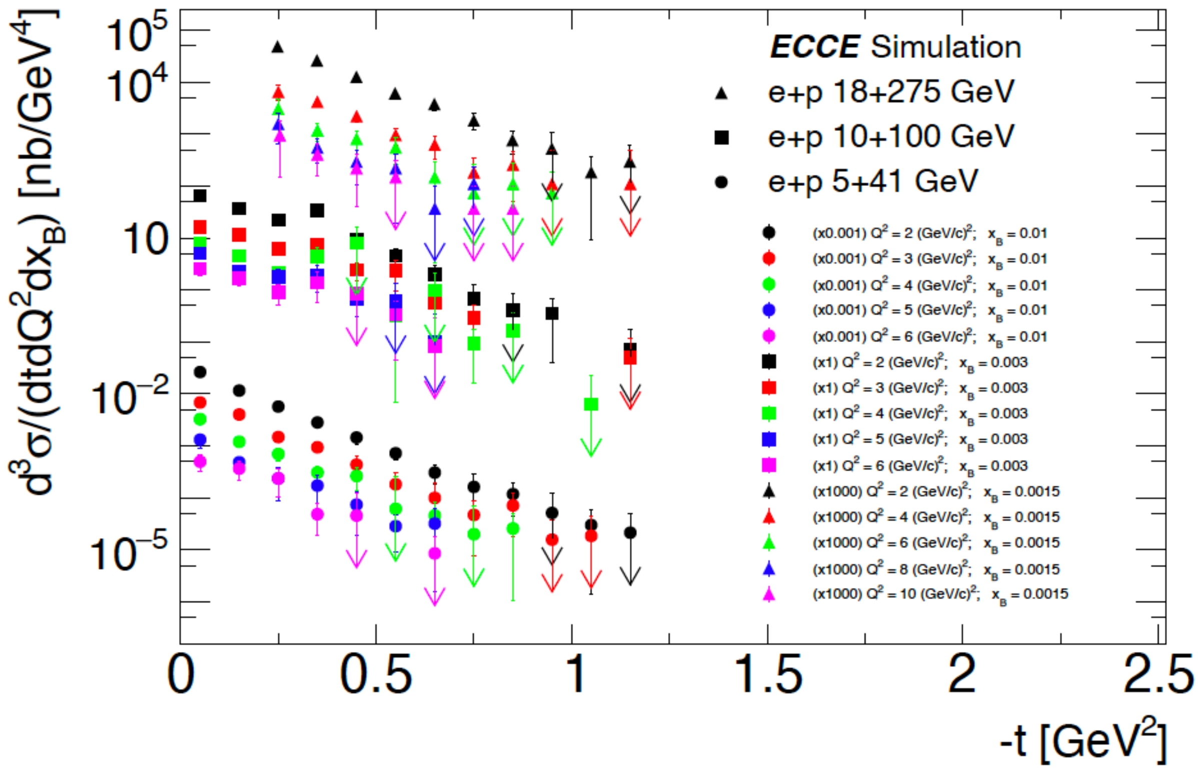


Dip in t-distribution is due to a gap between Roman Pots and B0 tracker: intrinsic to interaction region (IR). Gap position depends on proton beam-energy.

Plots: Kong Tu (BNL), D. S.

Coherent DVCS at the EIC

- DVCS on the proton:



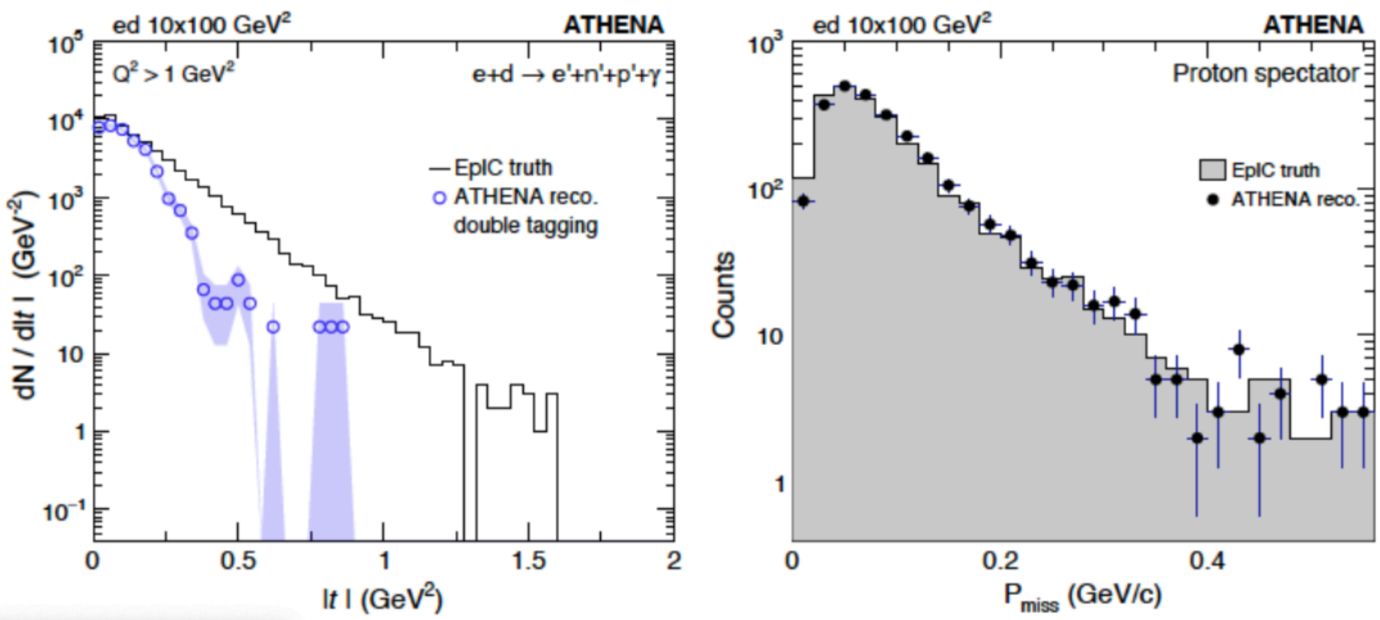
Detection of electron and photon in the central detector

Multi-dimensional binning: strong constraints on extraction of Compton Form Factors.

Plots: I. Korover (MIT)

DVCS in ed

- Enables measurement on neutron in deuterium: quark-flavour separation of GPDs, sensitivity to other Compton Form Factors (eg: Im E in DVCS beam-spin asymmetry on neutron vs Im H in BSA on proton).
- Both the spectator proton and the scattered neutron tagged in the measurement.
- Spectator proton is used to determine initial neutron momentum, to enable reconstruction of *t*:



- Scattered neutron tagged in ZDC: loss of *t*-acceptance at high *t* is due to limitations of ZDC acceptance. Can obtain better *t* acceptance from electron - photon.

Imaging light nuclei

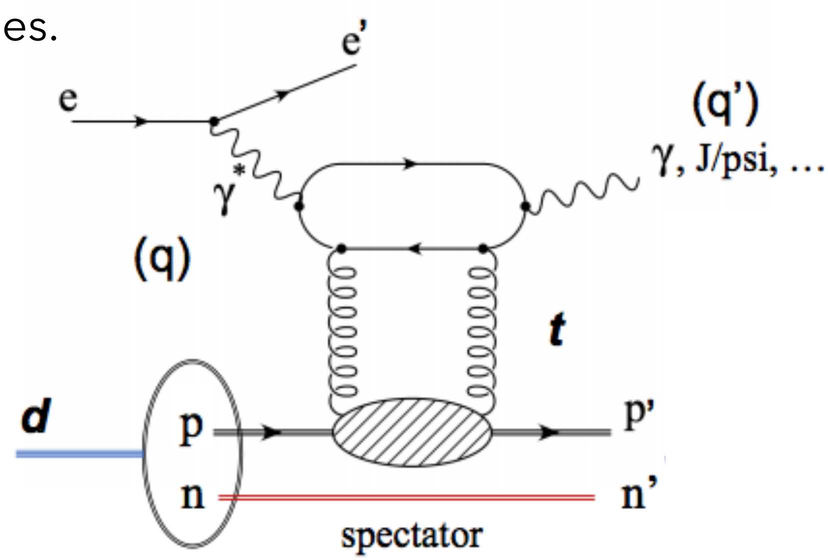
Coherent DVCS on light nuclei requires their intact detection and provides access to nuclear GPDs: imaging of partons in a nuclear medium.

Deuteron: spin-1. Many more GPDs at leading twist – theoretically well-described, experimentally almost untested. Challenge to detect *d* at low-*t*: use veto of deuteron breakup.

^{3}He: spin-1/2. DVCS amplitude has same GPD decomposition as for nucleon, binding energy larger than for deuteron – ideal to look for onset of nuclear effects.
Polarised neutron – possibility for completely new studies.

^{4}He: spin-0. Only one leading-twist GPD! Fully bound nucleus – access to medium-modification effects. Studies with TOPEG generator (*G. Penman, Glasgow*)

Incoherent DVCS (or meson-production): scatter from the nucleon, tag the process by detecting the spectator recoil → access to measurements on a quasi-free neutron.

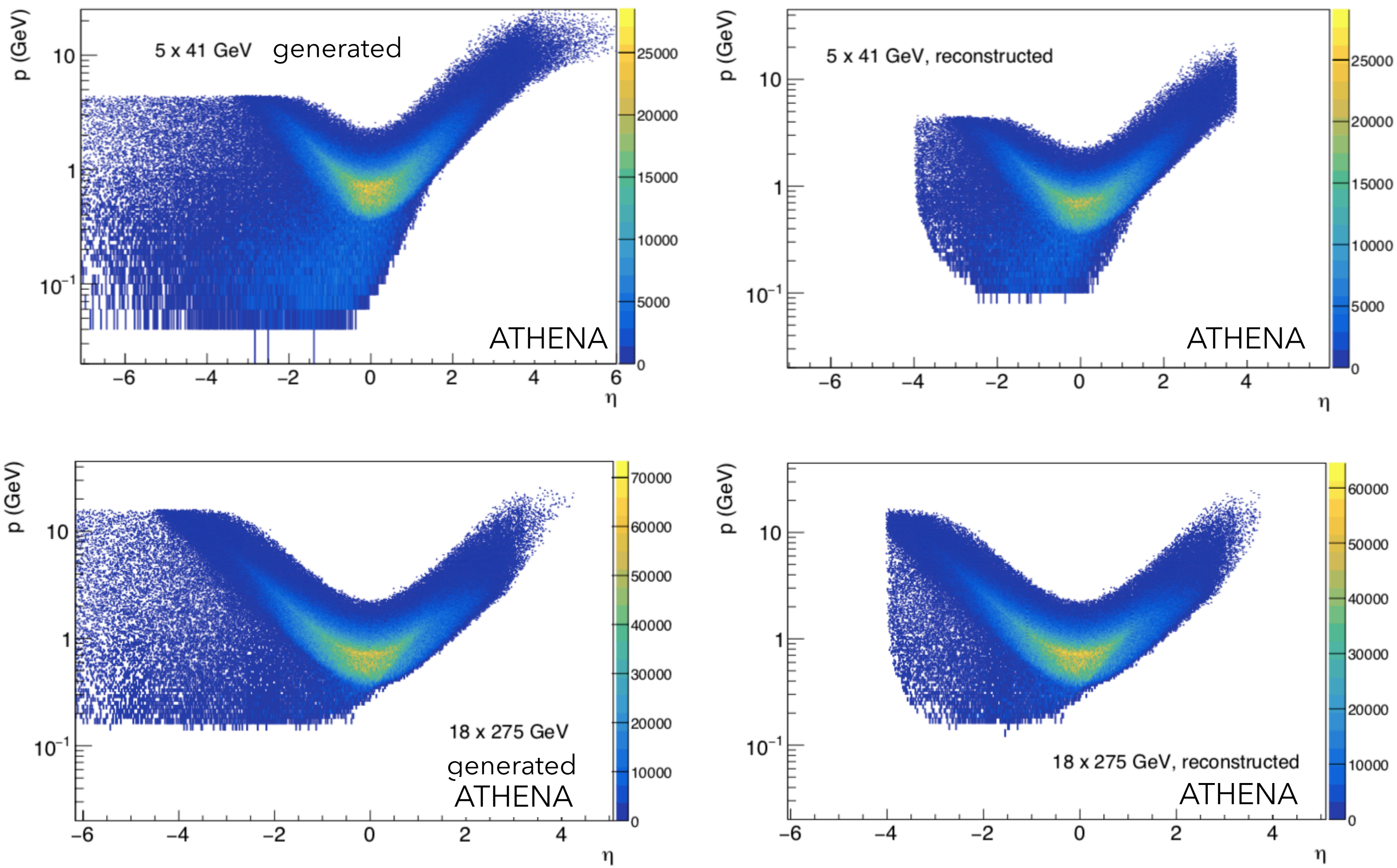


K. Tu, A. Jentsch

Flavour-decomposition, sensitivity to different GPDs…

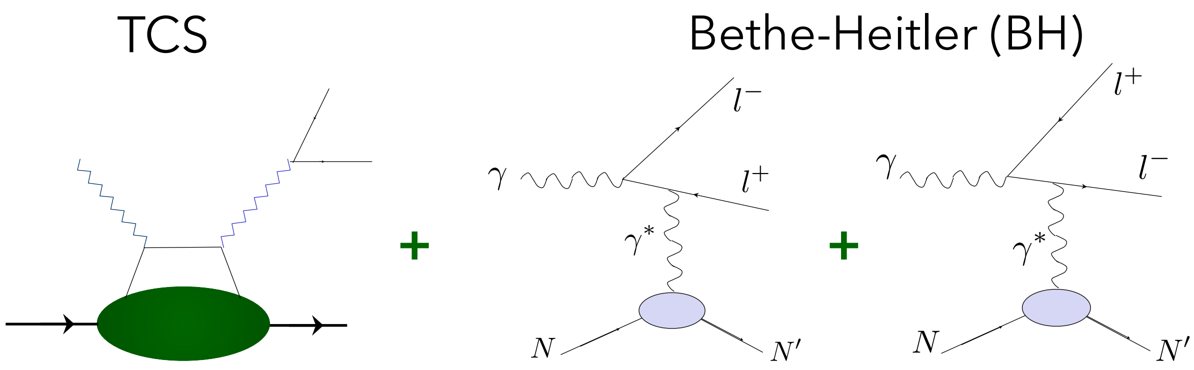
TCS e^+e^-

* TCS-produced virtual photon decays into e^+e^- pairs at mid-rapidity: need excellent acceptance in central region (barrel + end-caps), as scattered electron will in general be reconstructed though missing mass and momentum (low-Q^2 tagger can provide direct detection only in a part of the phase space). Also need excellent PID to suppress $\pi^+\pi^-$ background.



TCS (+ BH + Int) e+e-: observables

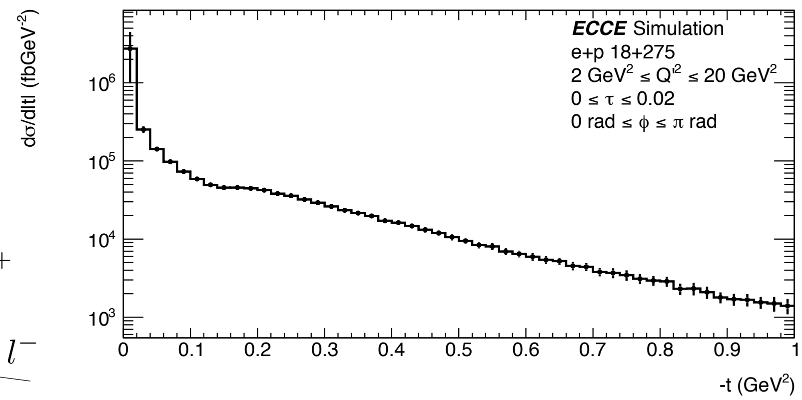
- Pure TCS cross-section is dominated by a factor of ~100 by Bethe-Heitler (BH): extract TCS signal from the BH-TCS interference.

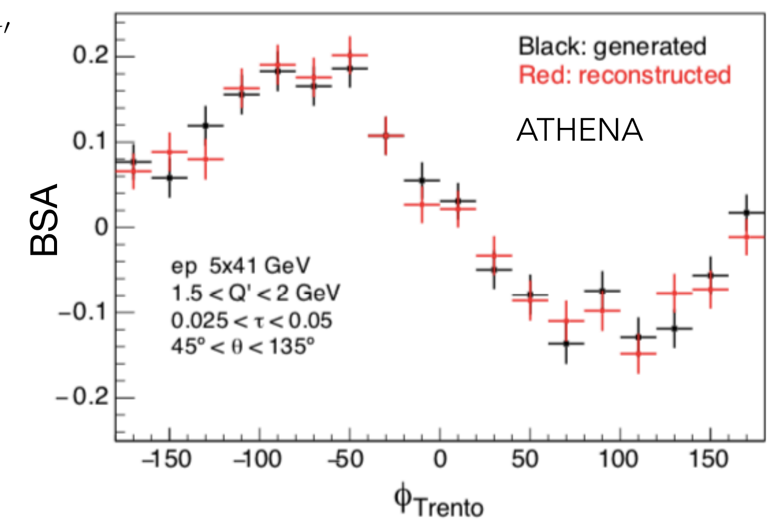


$$\sigma(\gamma p \to p' e^+ e^-) = \sigma_{BH} + \sigma_{TCS} + \sigma_{INT}$$

- Sensitivity to Interference term in single-spin asymmetries: beam-spin (BSA), target-spin.

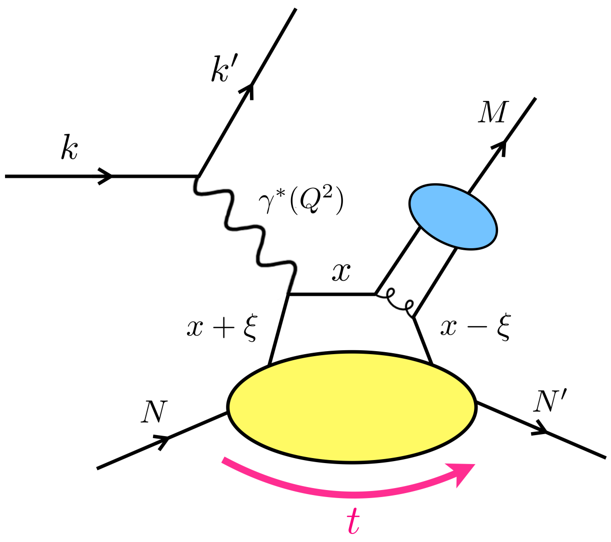
- Studied with the EpIC generator using the PARTONS framework.





Plots: Kayleigh Gates (Glasgow), D.S.

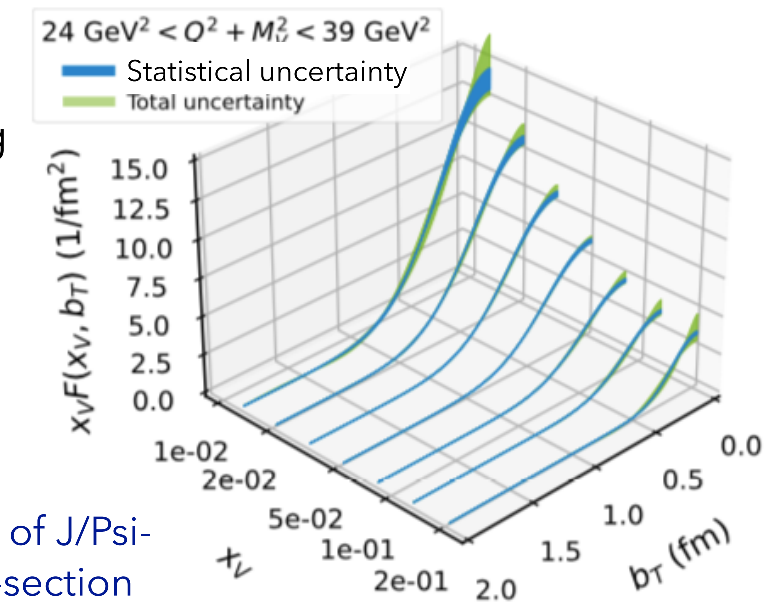
GPDs through meson-production



- Hard exclusive electro-production of vector mesons gives access to gluon GPDs, particularly clean in heavy mesons: J/Ψ and Υ

Hard scale in the scattering given by: $Q^2 + M_v^2$

Hence: $x_v = \frac{Q^2+M_v^2}{2\mathbf{p}\cdot\mathbf{q}}$

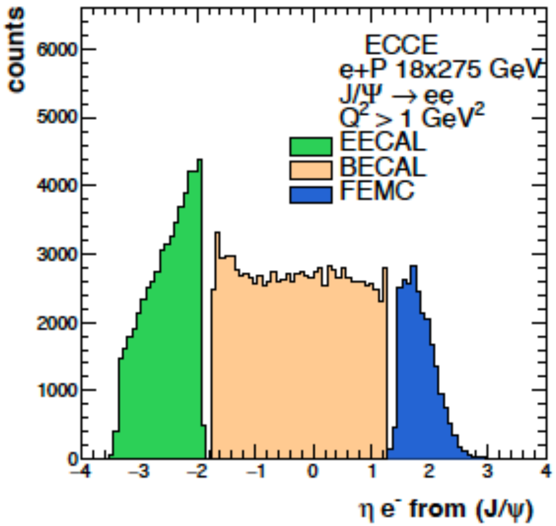


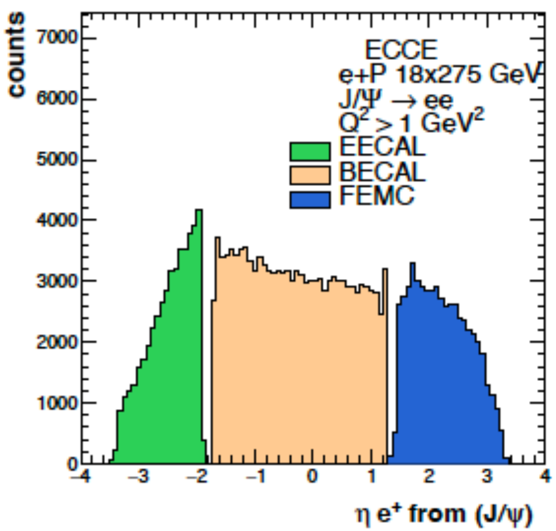
Fourier transform of J/Psi-production cross-section

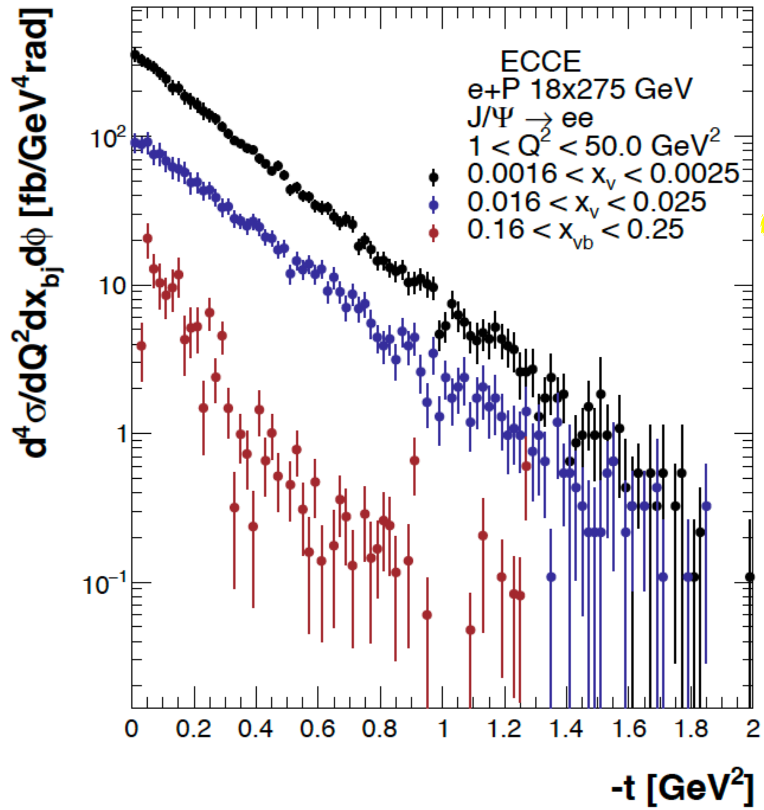
- Light vector-meson production additionally enables flavour-decomposition of GPDs.

- Light pseudo-scalar meson production gives, at high Q^2, access to parity-odd GPDs: $\tilde{H}$, $\tilde{E}$ and at low Q^2 to chiral-odd, transversity GPDs which are not accessible at leading-twist in DVCS processes.

J/Psi production







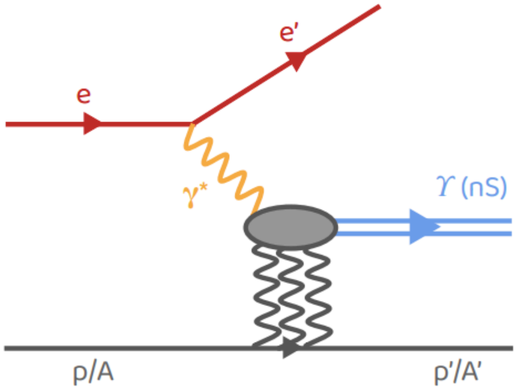
Photoprod. of quarkonium: talk by Vadim Guzey (Wed)

- Excellent acceptance coverage for J/Psi decay leptons
- Multi-dimensional binning

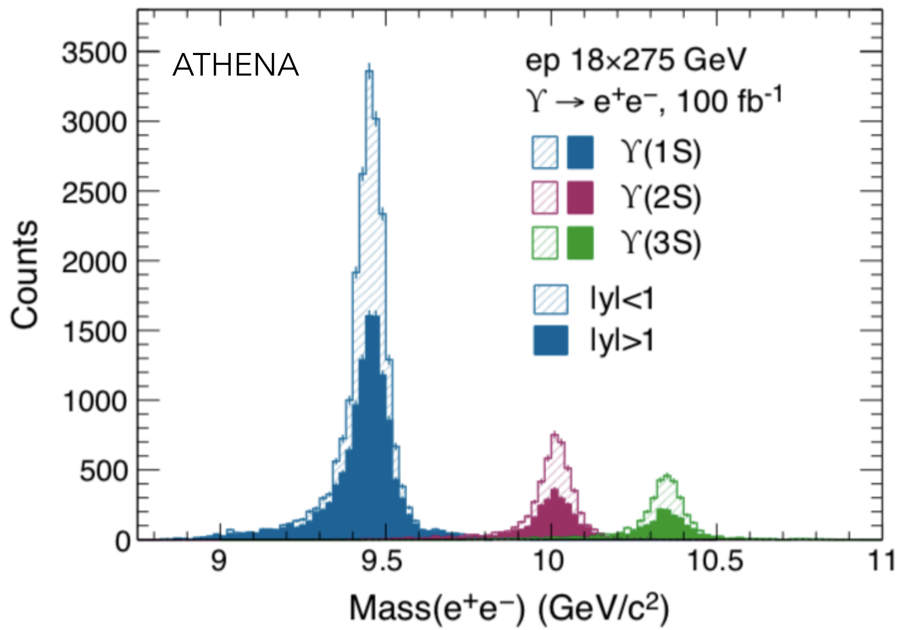
Plots: N. Santiesteban

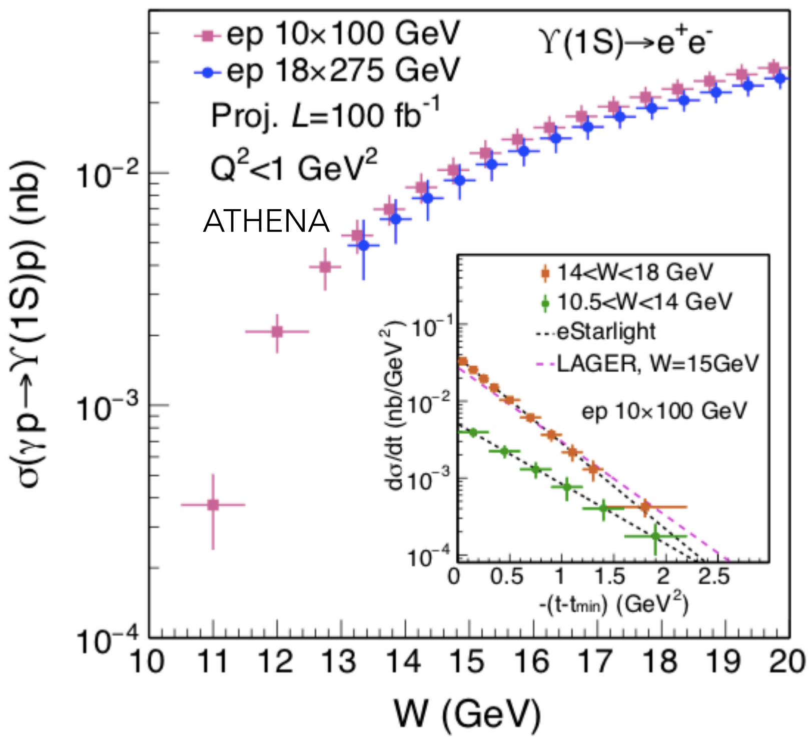
Upsilon-production

- Sensitivity to gluon distributions, information on colour correlations, upsilon-proton scattering lengths, possibly saturation. Near-threshold production: little-known, twist-4 effects contribute significantly.



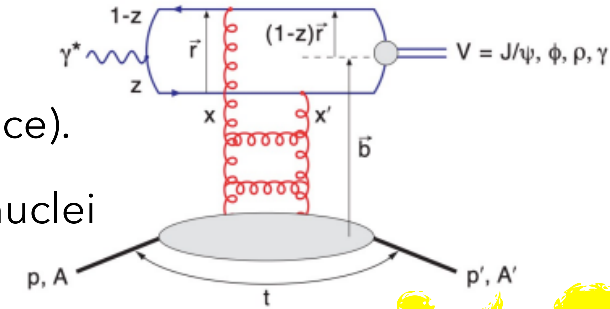
Photoproduction ($Q^2 < 1$ GeV^2) and electroproduction ($Q^2 > 1$ GeV^2).





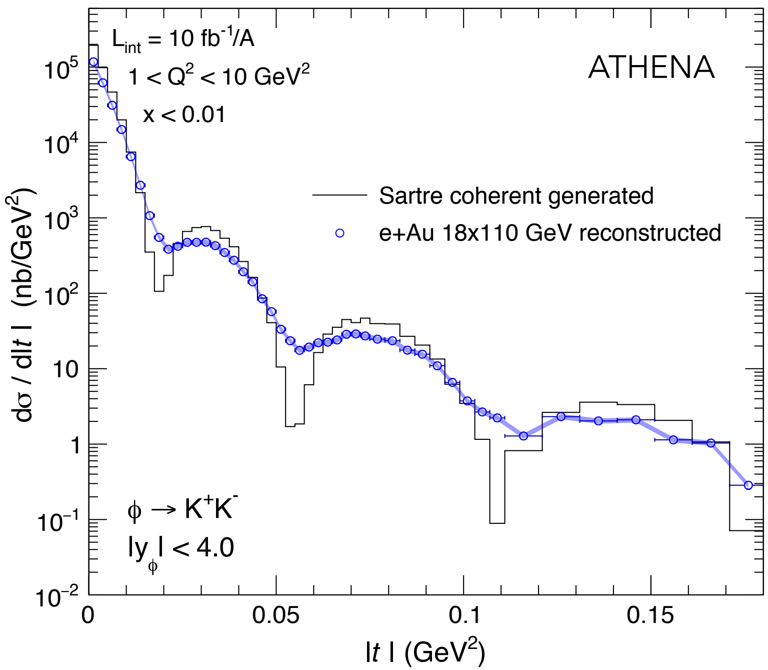
Plots: Yuangjing Ji (LBNL)

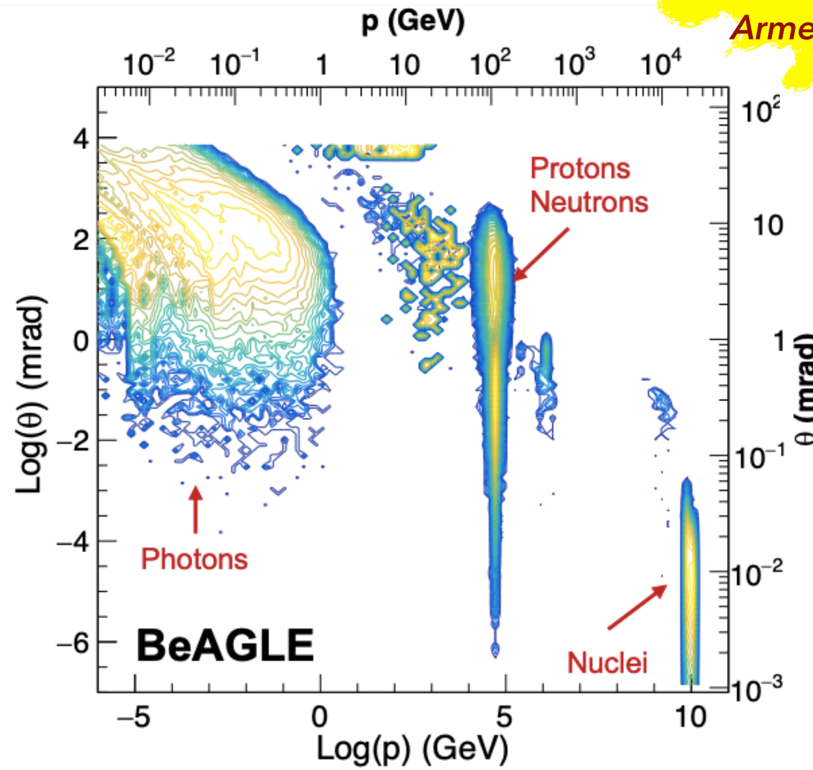
Coherent VM production in eA



- Gluon distributions in nuclei and a probe of saturation (in Q^2-depedence).
- Detector challenge: reconstruct *t* from leptons and mesons, not from nuclei (these escape undetected): resolution is crucial to identify *t* minima.
- Suppression of incoherent background by vetoing nuclear break-up in Far-Forward detectors.

CGC: talk by Nestor Armesto (Tues)

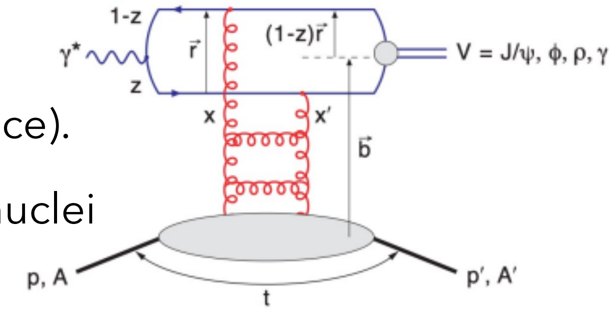




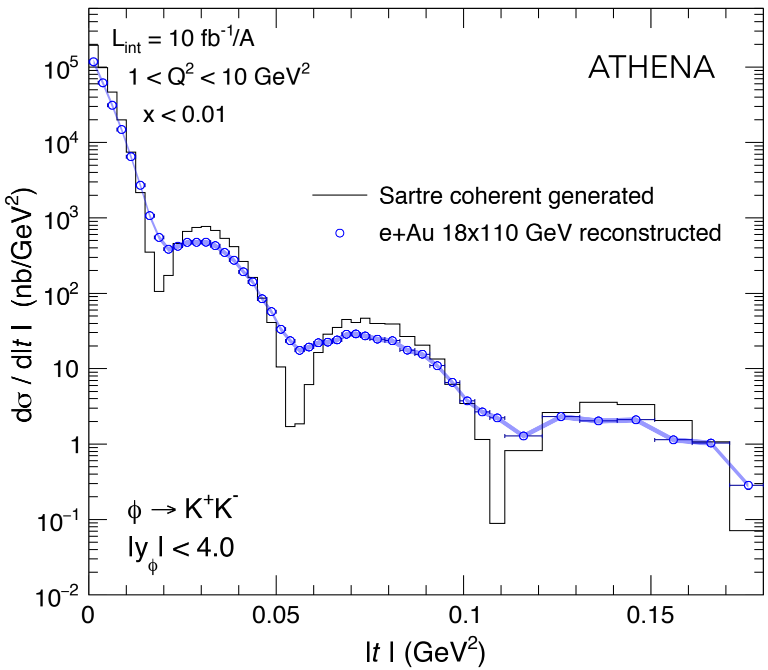
Incoherent backgrounds

Plot: Kong Tu (BNL)

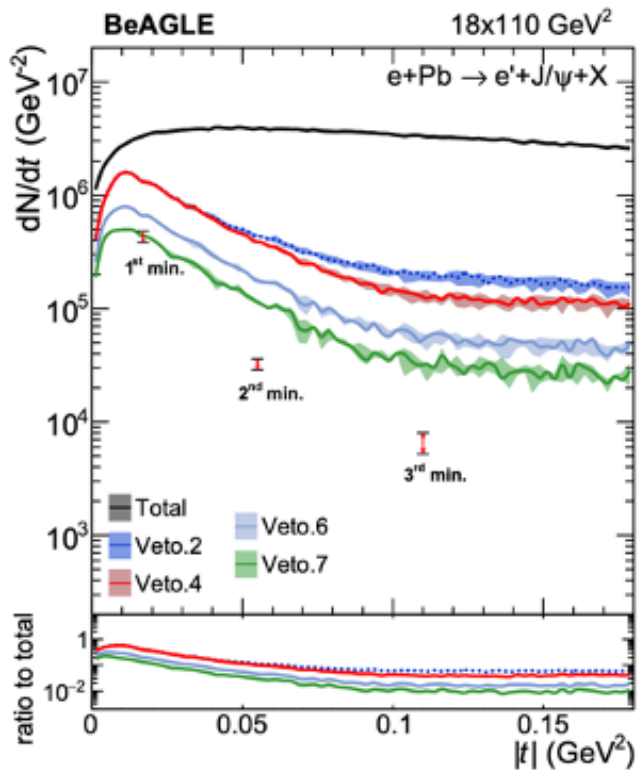
Coherent VM production in eA



- Gluon distributions in nuclei and a probe of saturation (in Q^2-depedence).
- Detector challenge: reconstruct *t* from leptons and mesons, not from nuclei (these escape undetected): resolution is crucial to identify *t* minima.
- Suppression of incoherent background by vetoing nuclear break-up in Far-Forward detectors.



Plot: Kong Tu (BNL)

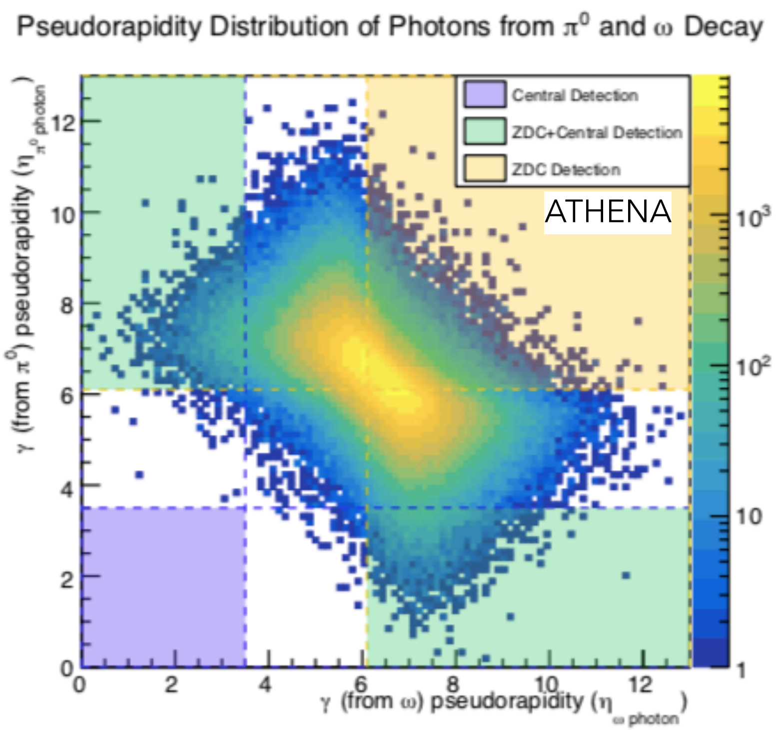


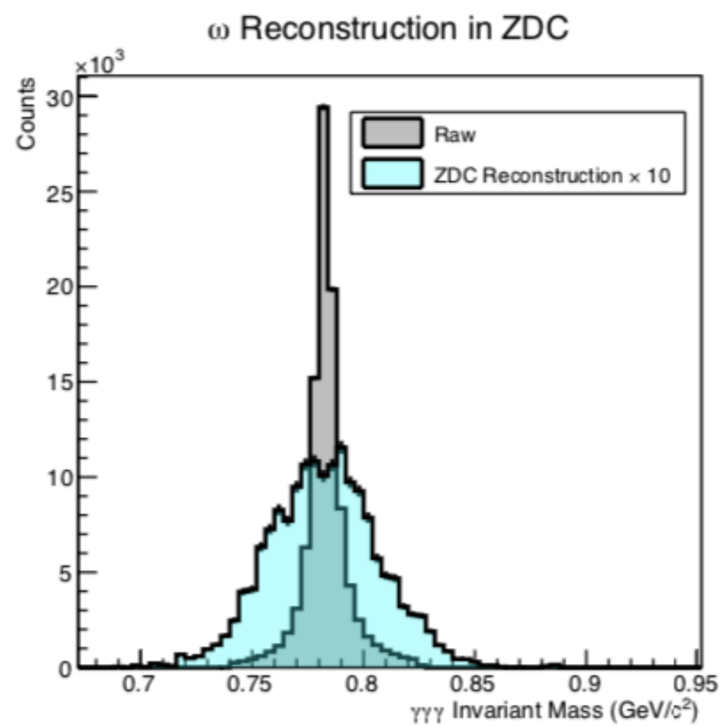
Phys. Rev. D ***104****, 114030*

Backward production of omega

- Similar to normal meson photo-production, but proton at mid-rapidity and meson goes forward with high momentum: u-channel.
- Sensitivity to Transition Distribution Amplitudes (TDA).
- Proton (a few hundred MeV) detected in central detector. Photons from meson decay detected in a combination of central and ZDC.

$$\omega \to \pi^0 \gamma$$





Plots: Zachary Sweger (UC Davis), A. Jentsch (BNL)

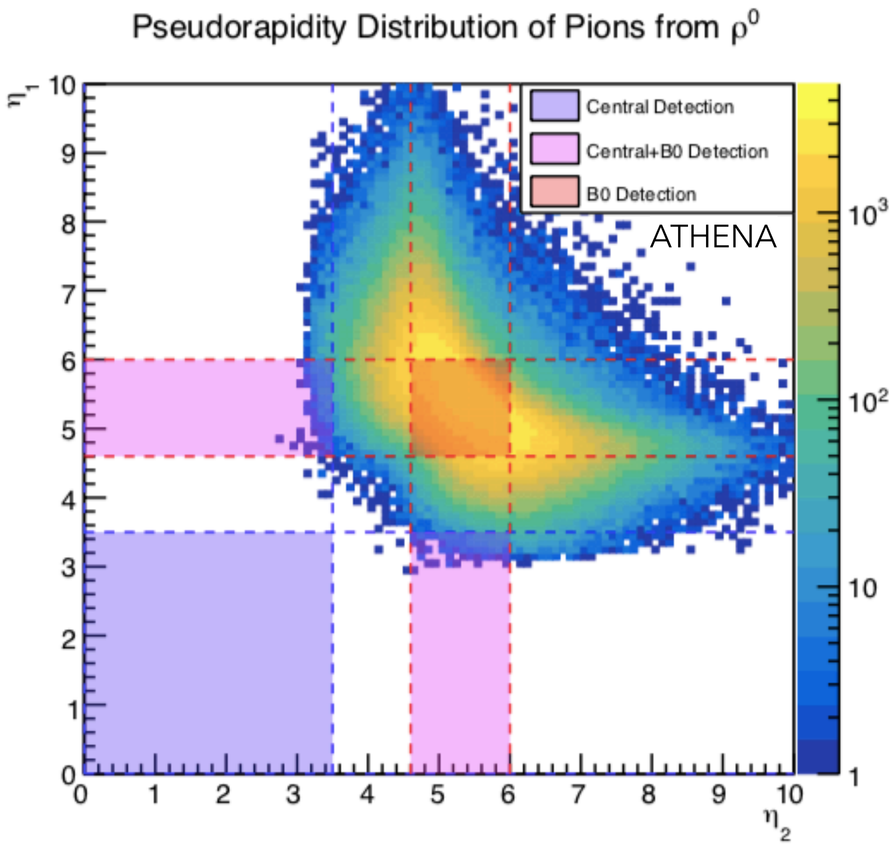
18 x 275 GeV, $Q^2 < 1$ GeV^2

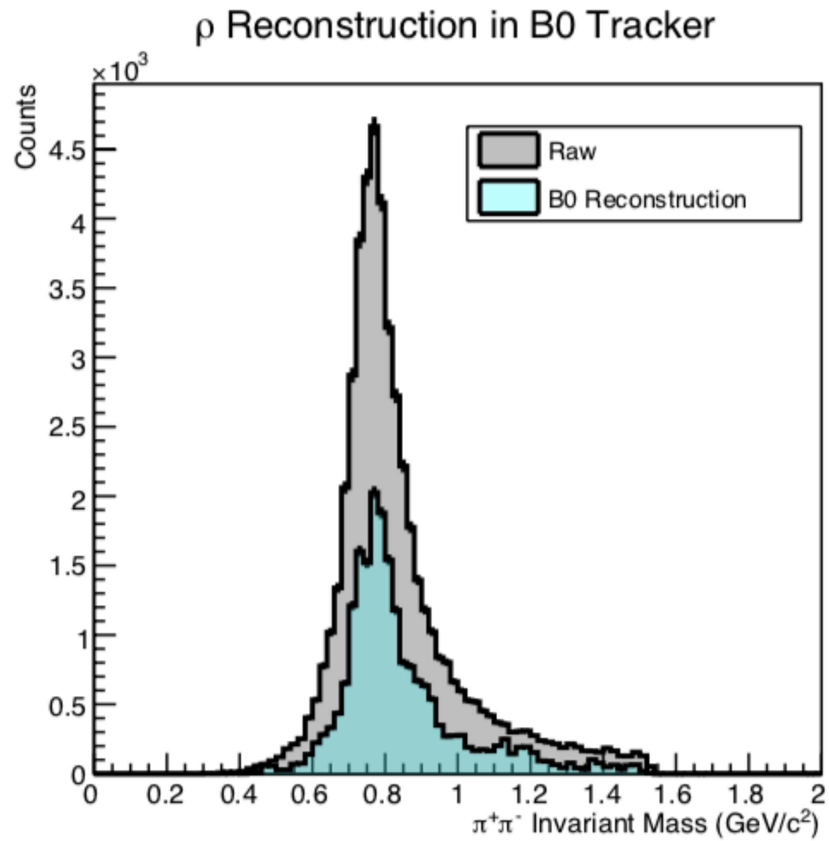
Backward production of rho

- Charged mesons in the u-channel production can also be reconstructed using Central detector, ZDC and B0 tracker.

$$\rho \rightarrow \pi^+\pi^-$$

10 x 100 GeV, $Q^2 < 1$ GeV^2





Plots: Zachary Sweger (UC Davis), A. Jentsch (BNL)

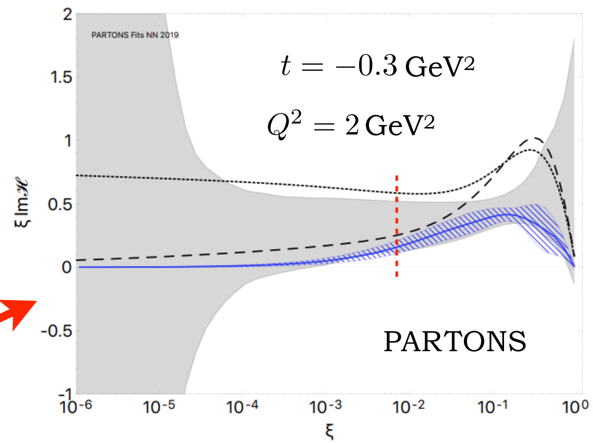
https://link.aps.org/doi/10.1103/PhysRevC.106.015204

Towards nucleon tomography

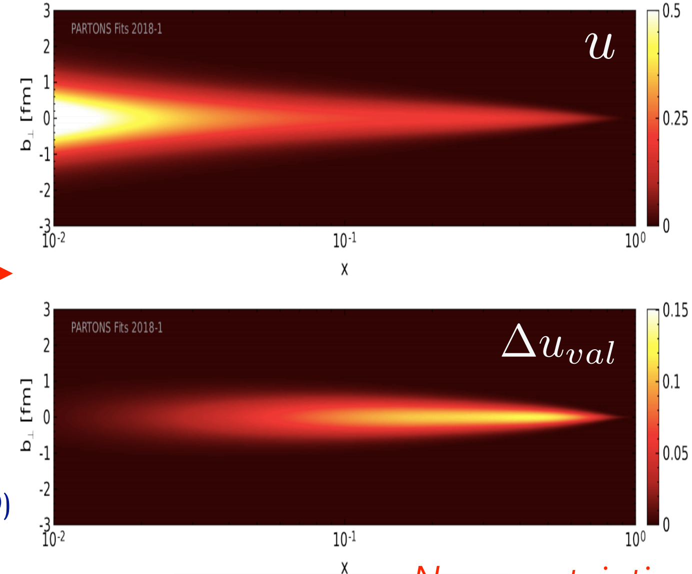
Ongoing imaging efforts on available world-data, strongest constraints in the valence region:

Uncertainties in the extraction of CFFs translate into uncertainties in spatial distributions.

PARTONS global fit with neural networks to minimise model-dependence in the extraction of CFFs.



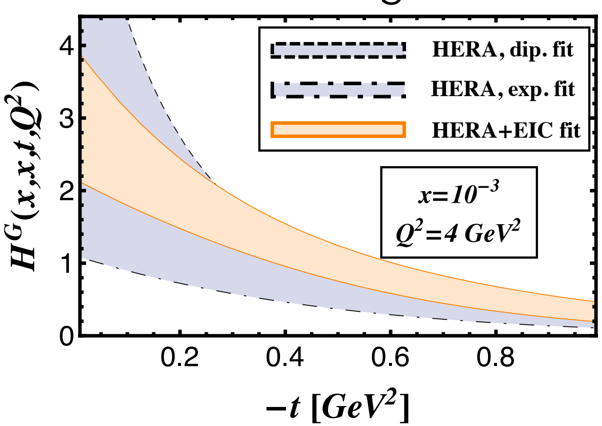
H. Moutarde *et al.*, Eur. Phys. J C79, 614 (2019)



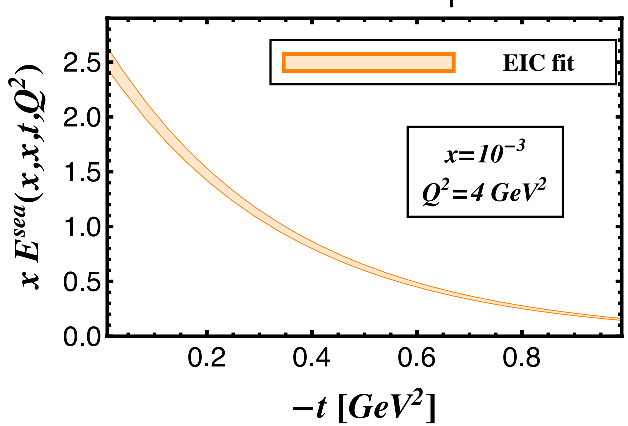
No uncertainties shown

Anticipated constraints from EIC on GPDs H and E:

GPD H for gluons



GPD E for sea quarks

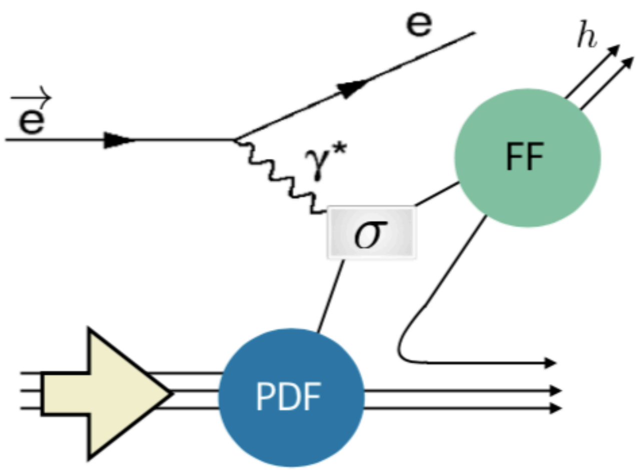


Measurements at EIC will provide significant constraints at low-x and enable extraction of as-yet unknown GPDs.

https://arxiv.org/abs/2103.05419

Transverse Momentum-Dependent Distributions

Can define TMDs for the nucleon in terms of x and k_T: 3D momentum distributions.
In SIDIS, structure functions are convolutions of TMDs with Fragmentation Functions (FF), defined in terms of z (parton's longitudinal momentum fraction in the hadron) and k_T: describe hadronisation of outgoing parton.



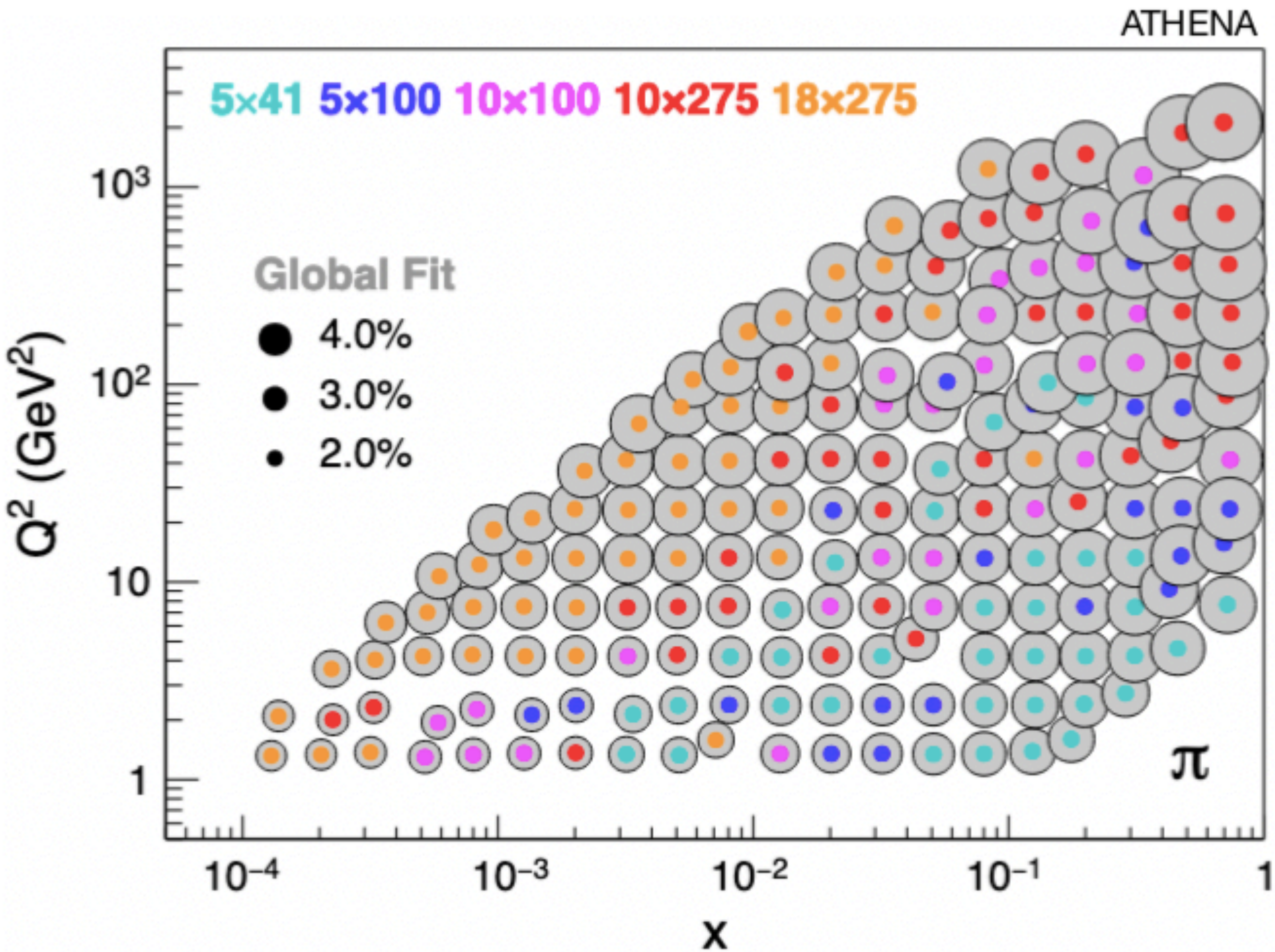
Formalism relies on TMD factorisation: applicable in certain regimes.

TMDs: talk by Valerio Bertone (Tues)

TMDs at play in combinations of quark and nucleon polarisations (at leading twist):

Quark \ Nucleon	Unpol.	Long.	Trans.
Unpol.	f_1 =		$f_{1T}^\perp$ = − Sivers
Long.		g_{1L} = − Helicity	g_{1T} = − Worm-gear
Trans.	$h_1^\perp$ = − Boer-Mulders	$h_{1L}^\perp$ = − Worm-gear	Transversity h_{1T} = − $h_{1T}^\perp$ = − Pretzelocity

Spin-independent TMD f_1



Fit: A. Bacchetta *et al.*
JHEP 06 (2017) 081,
JHEP 06 (2019) 051 (*erratum*)

Experimental uncertainties dominated by assumed:

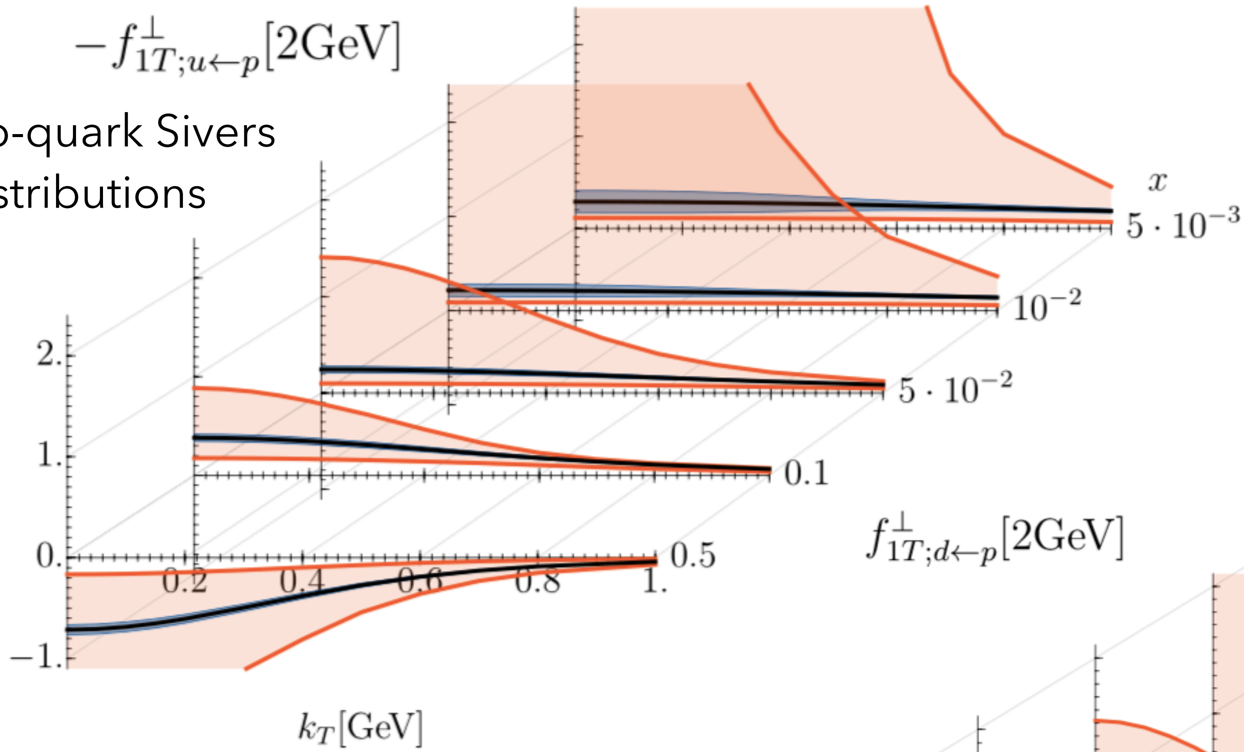
- 3% point-by-point uncorrelated uncertainty,
- 3% scale systematic

Theory uncertainties dominated by TMD evolution.

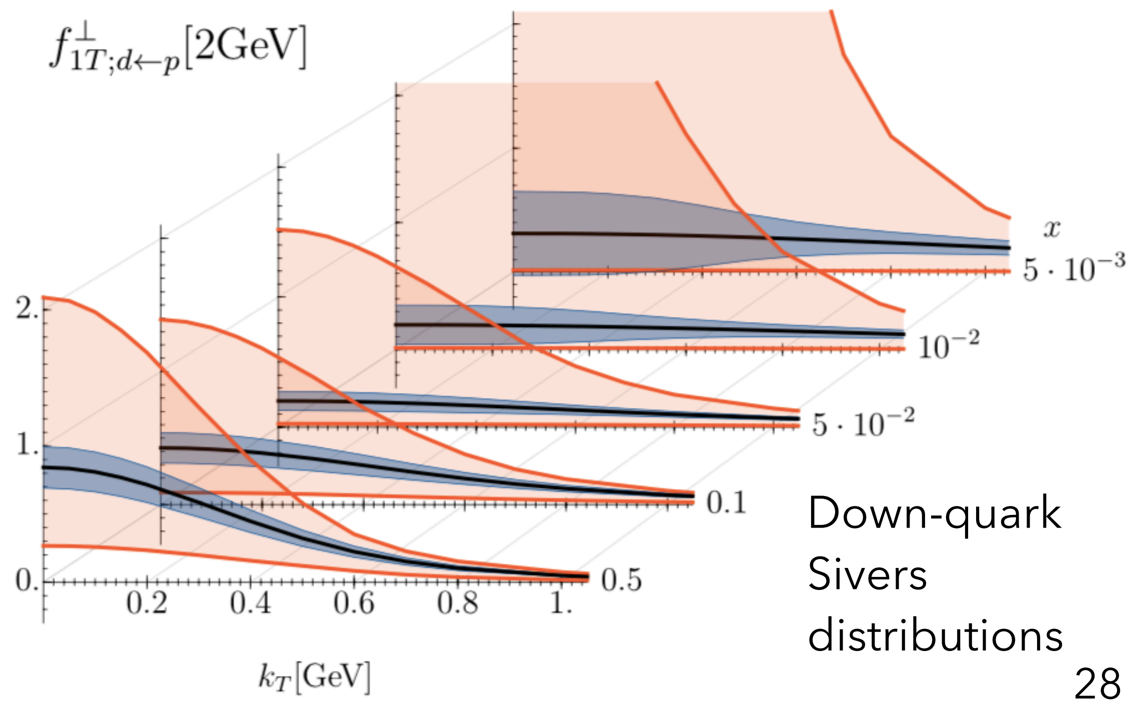
Large lever arm in Q^2 over a wide *x*-range: constraints on TMD evolution.

EIC constraints on Sivers

Up-quark Sivers distributions



Orange: current uncertainties. Blue: EIC constraints from SIDIS pion and kaon pseudo-data, with ECCE pseudo-data.

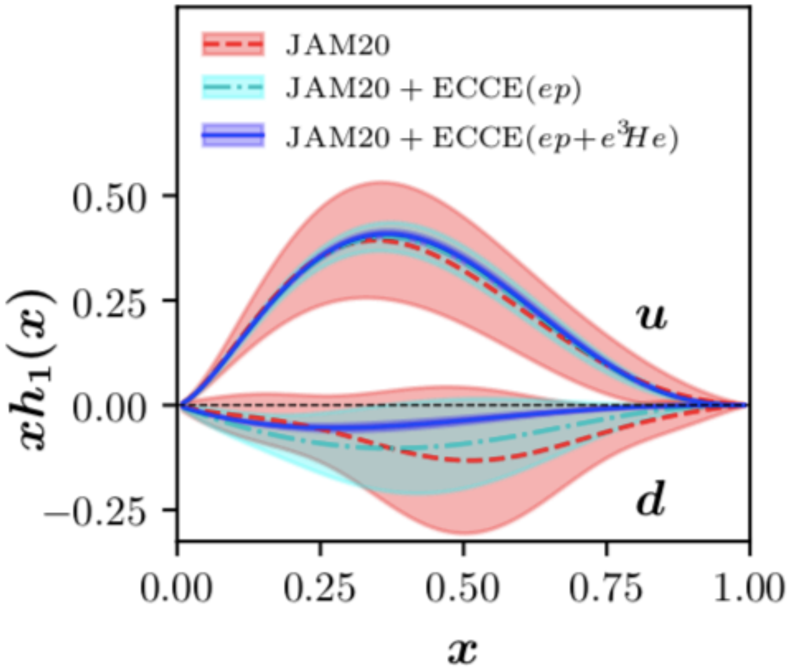


Transverse-momentum dependence of Sivers function obtained by measuring p_T-dependence of SIDIS cross-sections.

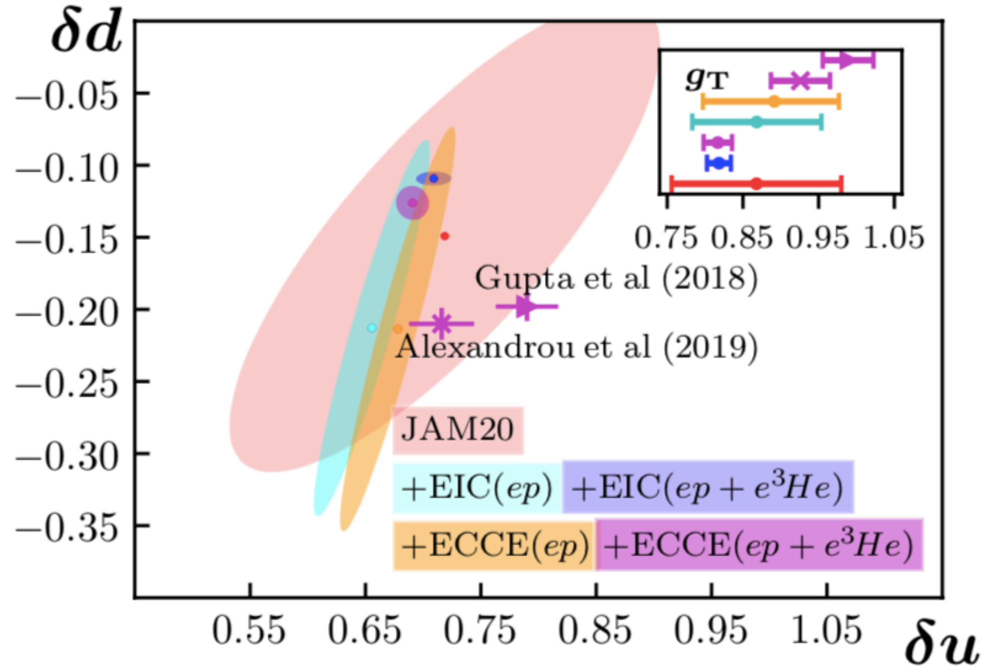
Down-quark Sivers distributions

arXiv:2207.10890v1 [hep-ex]

Transversity and tensor charges



Transversity distributions for up and down quarks



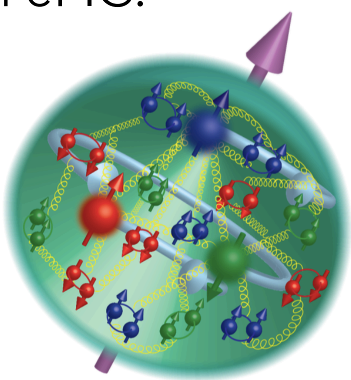
Up, down, and iso-vector combination tensor charges

arXiv:2207.10890v1 [hep-ex]

Study with ECCE detector design

Summary

- Electron-Ion Collider to be built at Brookhaven National Laboratory, start operation ~2032.
- Large range of CoM energies (30 - 140 GeV), high luminosity ($10^{33\text{-}34}$ $cm^{-2}s^{-1}$): high precision measurements across wide swathes of phase space from the gluon sea to the valence quark region.
- Design of the first detector being finalised this year: the ePIC collaboration.
- Hermeticity, tracking, PID, neutral particle detection. Focus on the far-forward region – excellent reconstruction of scattered protons and light ions at the smallest angles. Detection of neutral particles at low angles.
- A range of exclusive and semi-inclusive processes accessible with ePIC: sensitivity to a range of GPDs and TMDs in low-x.
- Join us! http://www.eicug.org/

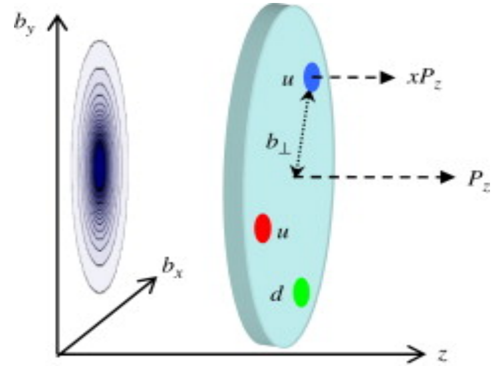


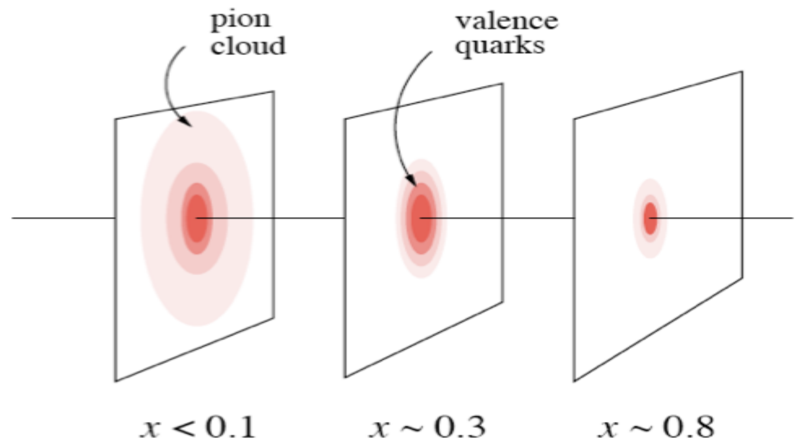

Thank you!
Any questions?

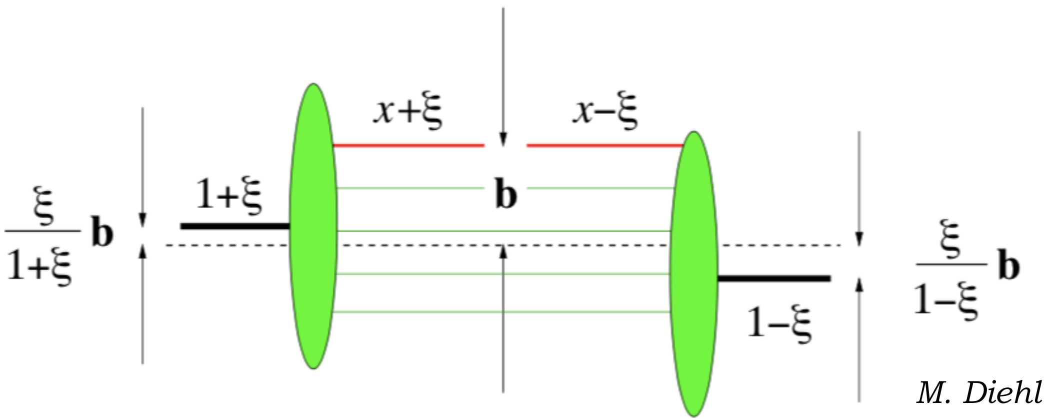

Back-up

Nucleon Tomography from GPDs

At a fixed Q^2, x_B and $\xi=0$ slope of GPD with t is related, via a Fourier Transform, to the transverse spatial distribution.







Formally, the radial separation, **b**, *between the struck parton and the centre of momentum of the remaining spectators.*

Experimentally, can fit the t-dependence of structure functions (from meson-production) or Compton Form Factors (from DVCS/TCS) with an exponential:

eg: $$\frac{d\sigma_U}{dt} = Ae^{Bt}$$

Spin and pressure in the nucleon

- GPDs also provide indirect access to mechanical properties of the nucleon (encoded in the gravitational form factors, GFFs, of the energy-momentum tensor).

  X. D. Ji, PR**D 55**, 7114-7125 (1997)
  M. Polyakov, PL**B 555**, 57-62 (2016)

- Three scalar GFFs, functions of *t*: encode pressure and shear forces ($d_1(t)$), mass ($M_2(t)$) and angular momentum distributions ($J(t)$).

- Can be related to GPDs via sum rules: $\int x\,[H(x,\xi,t)+E(x,\xi,t)]\,\mathrm{d}x = 2J(t)$

$$\int xH(x,\xi,t)\mathrm{d}x = M_2(t) + \frac{4}{5}\xi^2 d_1(t)$$

(Ji's relation)

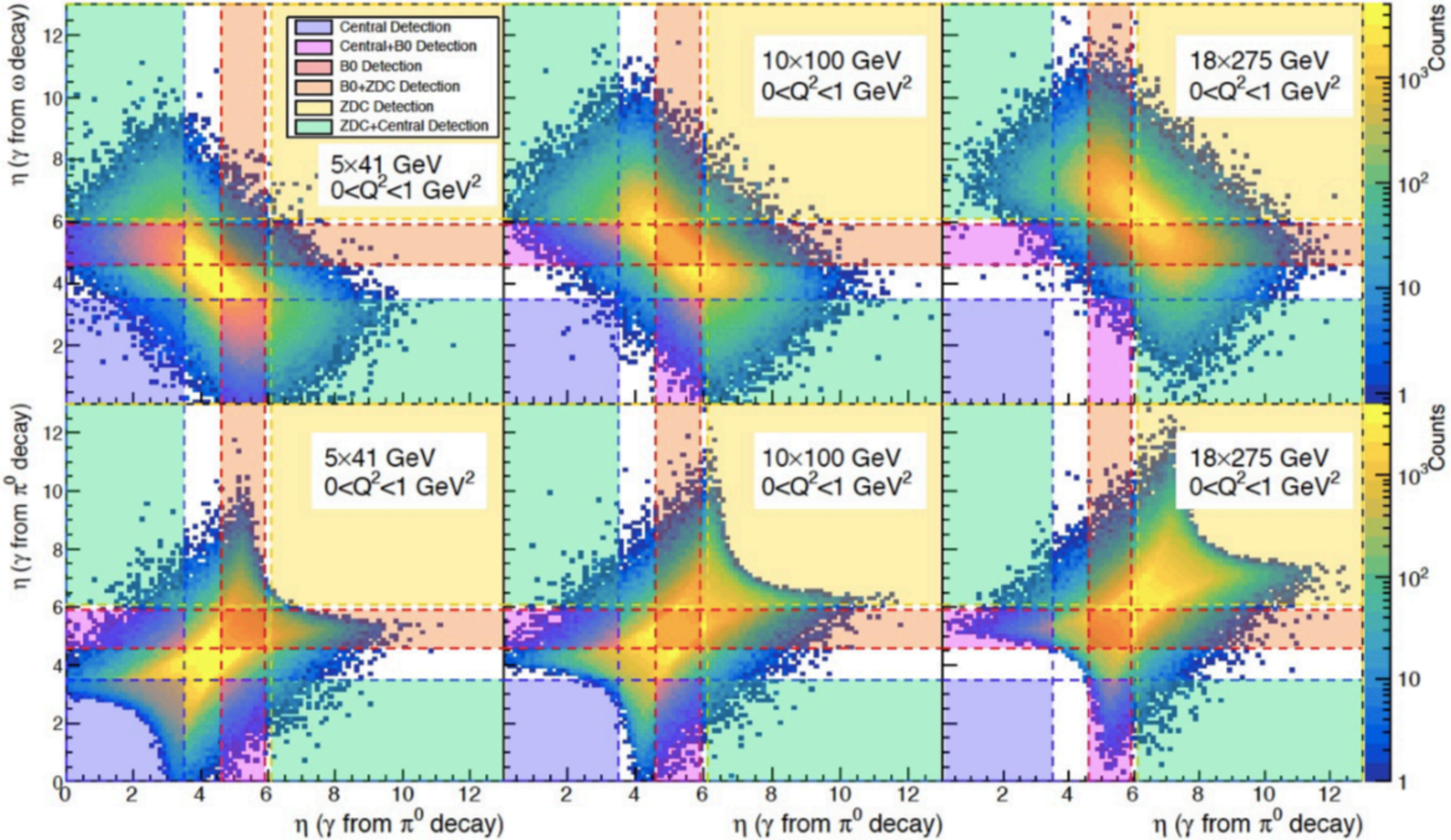
$$J_N = \frac{1}{2} = \frac{1}{2}\Sigma_q + L_q + J_g$$

- $d_1(t)$ (D-term) "last unknown global property of the nucleon" – can be accessed via the $\mathcal{Re}$ and $\mathcal{Im}$ $\mathcal{H}$:

Dispersion relation: $\mathrm{Re}\mathcal{H}(\xi,t) = \int_{-1}^{1}\left(\frac{1}{\xi-x} - \frac{1}{\xi+x}\right)\mathrm{Im}\mathcal{H}(\xi,t)\;dx + \Delta(t).$

Assuming double-distribution parametrisation: $\Delta(t) \propto d_1(t)$

Backward production of omega



Proton beam energy	ω eff. cent.+ZDC	ω eff. cent.+B0+ZDC
41 GeV	1.4%	18%
100 GeV	1.3%	41%
275 GeV	6%	63%

Plots: Zachary Sweger (UC Davis)

Spin and pressure

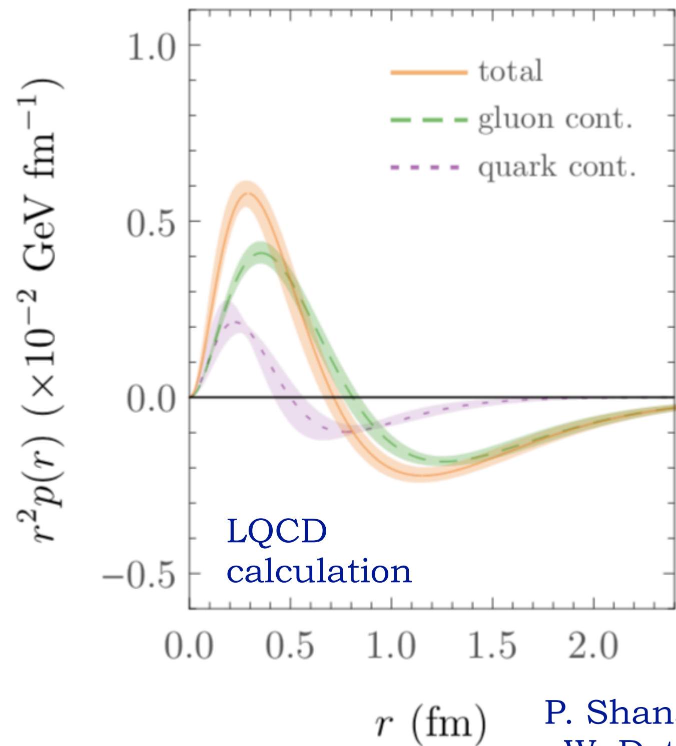
- GPDs provide indirect access to mechanical properties of the nucleon (encoded in the gravitational form factors, GFFs, of the energy-momentum tensor).

  X. D. Ji, PR**D 55**, 7114-7125 (1997) M. Polyakov, PL**B 555**, 57-62 (2003)

- Four GFFs, functions of *t*, of which three are related to moments of GPDs: they encode pressure and shear forces ($d_1(t)$), mass ($M_2(t)$) and angular momentum distributions ($J(t)$):

$$\int x\,[H(x,\xi,t)+E(x,\xi,t)]\,\mathrm{d}x = 2J(t)$$

$$\int xH(x,\xi,t)\mathrm{d}x = M_2(t)+\frac{4}{5}\xi^2 d_1(t)$$



P. Shanahan, W. Detmold, PRL 122,072003 (2019)

- The D-term: "last unknown global property of the nucleon" -- can be related to spatial distribution of shear forces and pressure within the nucleon.

- Possibilities of "imaging" spatial distributions of angular momentum: C. Lorcé, M. Montovani, B. Pasquini, PL**B 776**, 38-47 (2018)

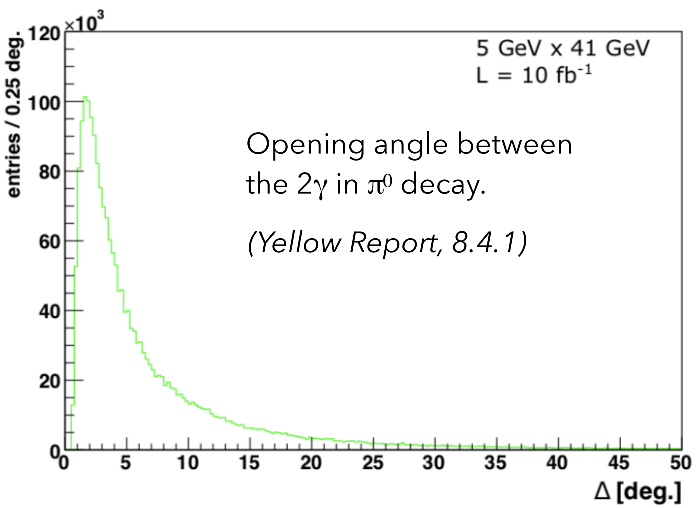
Studies for proposals:	ATHENA	ECCE
* DVCS in ep	EpIC	MILOU3D
* DVCS (incoherent) in ed	EpIC	
* DVCS on He-4		TOPEG
* TCS in ep	EpIC	EpIC
* J/Psi in ep		lAger, eSTARlight
* J/Psi in eA		lAger, eSTARlight
* Φ in eAu/Pb	SARTRE, BeAGLE	SARTRE, BeAGLE
* Υ(1S, 2S, 3S) in ep	eSTARlight, lAger	
* u-channel: ω, ρ in ep	eSTARlight	
* X,Y Ψ(2S) in ep -> J/Ψ $\pi^+\pi^-$ p	elSpectro	elSpectro
* Pion Form Factors		*
* Pion Structure Functions		*
* A^n_1 (He-3 double tagging)		*

Coherent DVCS at the EIC

Plots: I. Korover (MIT)

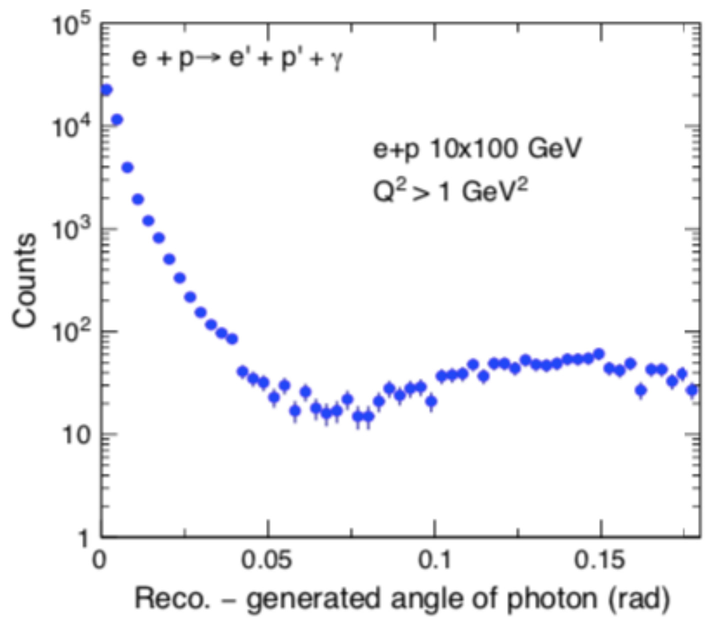
Main background from meson-production of π^0 which decays into 2γ pairs:

- Minimise risk of missing one photon: practically hermetic calorimeter coverage.
- Good calorimeter resolution to ensure photon clusters don't merge.



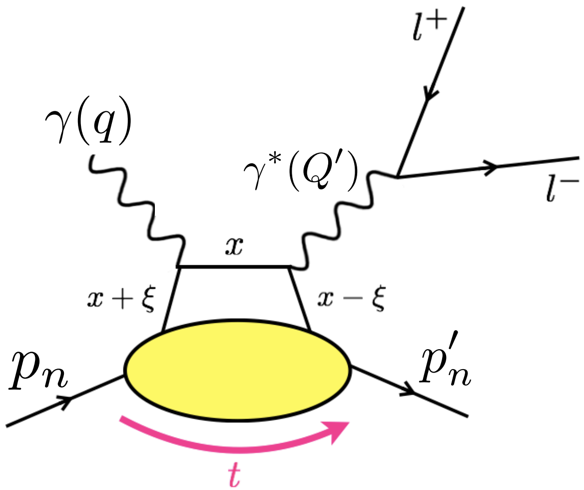
Opening angle between the 2γ in π^0 decay.

(Yellow Report, 8.4.1)



Difference between generated and reconstructed DVCS photon mainly < 0.17mrad (1deg): smallest opening angle for π^0 decay.

Timelike Compton Scattering



- Time-reversal process of DVCS: parametrised in terms of same Compton Form Factors (their complex conjugates).

- Verification of GPD universality.
- Another route to access the D-term.

- Factorisation ensured by hard scale of γ* virtuality:

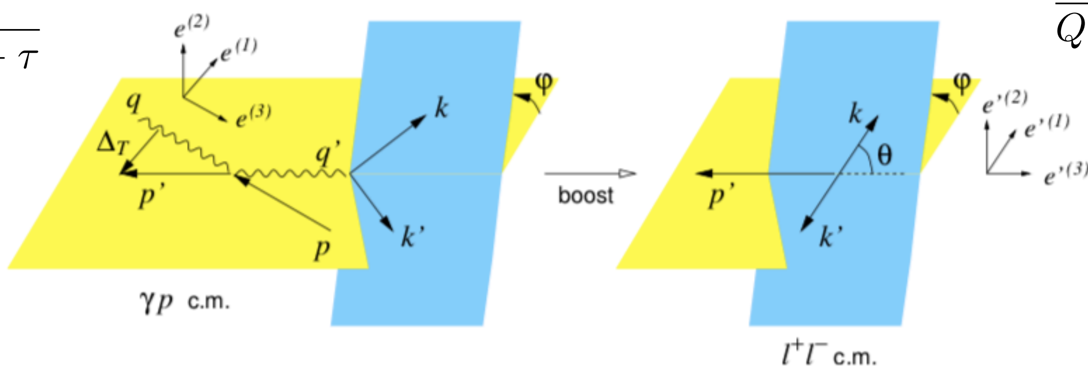
$$\frac{t}{Q'^2} << 1$$

$$Q' = l^+ + l^- \qquad \xi = \frac{\tau}{2-\tau}$$

$$s = (q + p_n)^2$$

$$\tau = \frac{Q'^2}{s - m_p^2}$$

θ: angle between l^+ and scattered proton in lepton CMS



- Measurements establish dependence on Q^2, x, t and φ (angle between leptonic and hadronic planes).

TCS observables

- Unpolarised cross-sections:

  sensitive to $\mathcal{R}e\ \mathcal{H}$.

$$\frac{d^4\sigma_{INT}}{dQ'^2 dt d\Omega} = A\frac{1+\cos^2\theta}{\sin\theta}[\cos\phi\ \mathrm{Re}\tilde{M}^{--} - \nu\cdot\sin\phi\ \mathrm{Im}\tilde{M}^{--}]$$

$$\tilde{M}^{--} = \left[F_1\mathcal{H} - \xi(F_1+F_2)\tilde{\mathcal{H}} - \frac{t}{4m_p^2}F_2\mathcal{E}\right]$$

suppressed

- Circularly-polarised photon cross-section: access to $\mathcal{I}m\ \mathcal{H}$.
- More promising observables: asymmetries and cross-section ratios.
- Photon-polarisation (beam-spin) asymmetry:

$$A_{\odot U} = \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-}$$

access to $\mathcal{I}m\ \mathcal{H}$

- Forward - backward asymmetry:

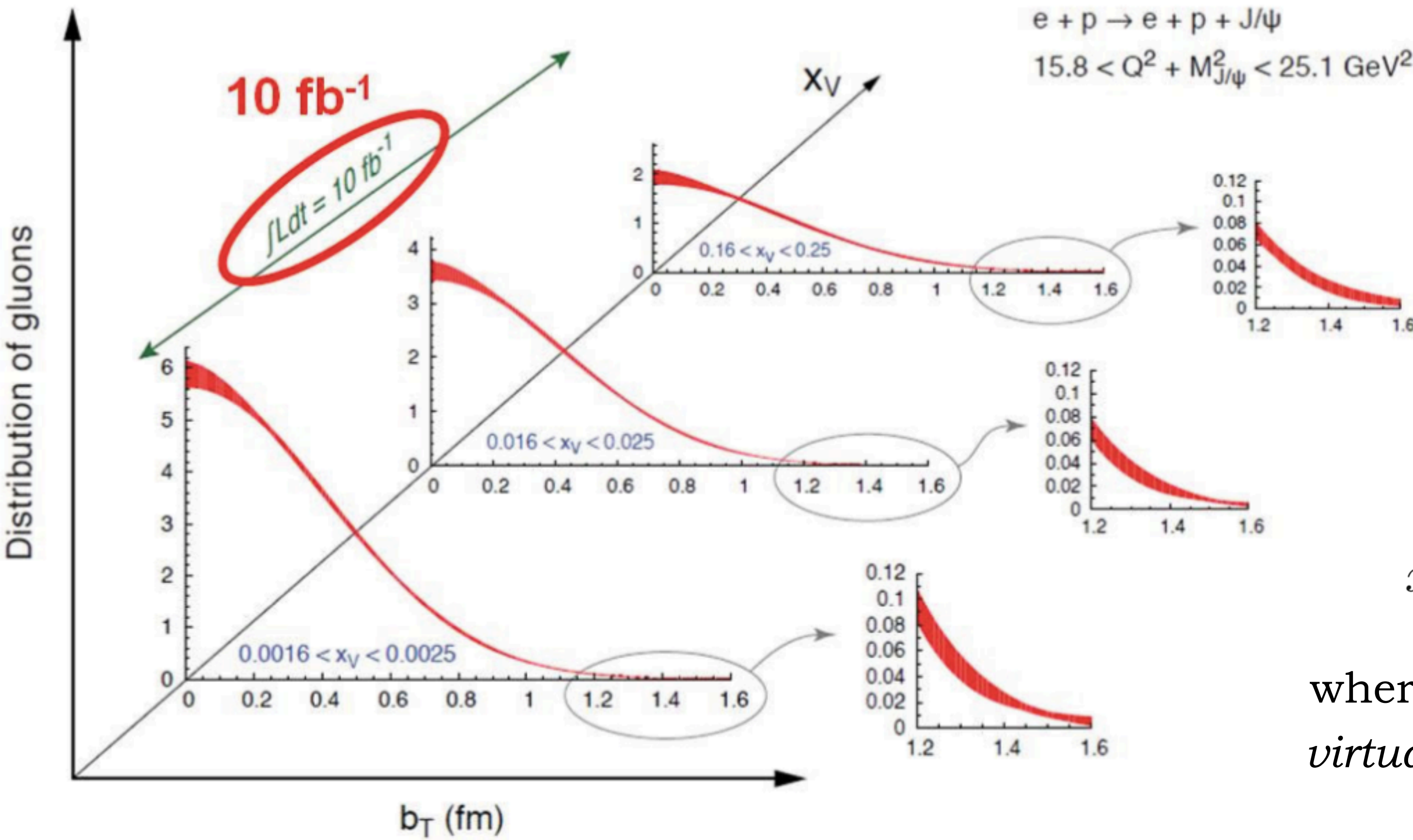
$$A_{FB}(\theta,\phi) = \frac{d\sigma(\theta,\phi) - d\sigma(180^\circ-\theta, 180^\circ+\phi)}{d\sigma(\theta,\phi) + d\sigma(180^\circ-\theta, 180^\circ+\phi)}$$

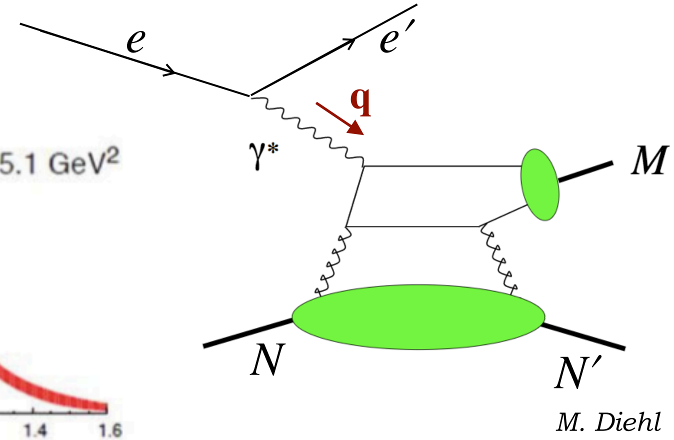
access to $\mathcal{R}e\ \mathcal{H}$

P. Chatagnon et al. (CLAS): arXiv:2108.11746 [hep-ex], *accepted (PRL)*

Nucleon tomography: imaging glue

- Gluon GPDs can be accessed through deeply virtual meson production (DVMP), eg: J/Ψ
- Access to spatial distributions of gluons at different longitudinal momentum fractions:





Gluon momentum fraction related to:

$$x_V = x_B(1 + M^2_{J/\Psi}/Q^2)$$

where $Q^2 = -\mathbf{q}^2 = -(\mathbf{p}_e - \mathbf{p}_e')^2$

virtuality of exchanged photon

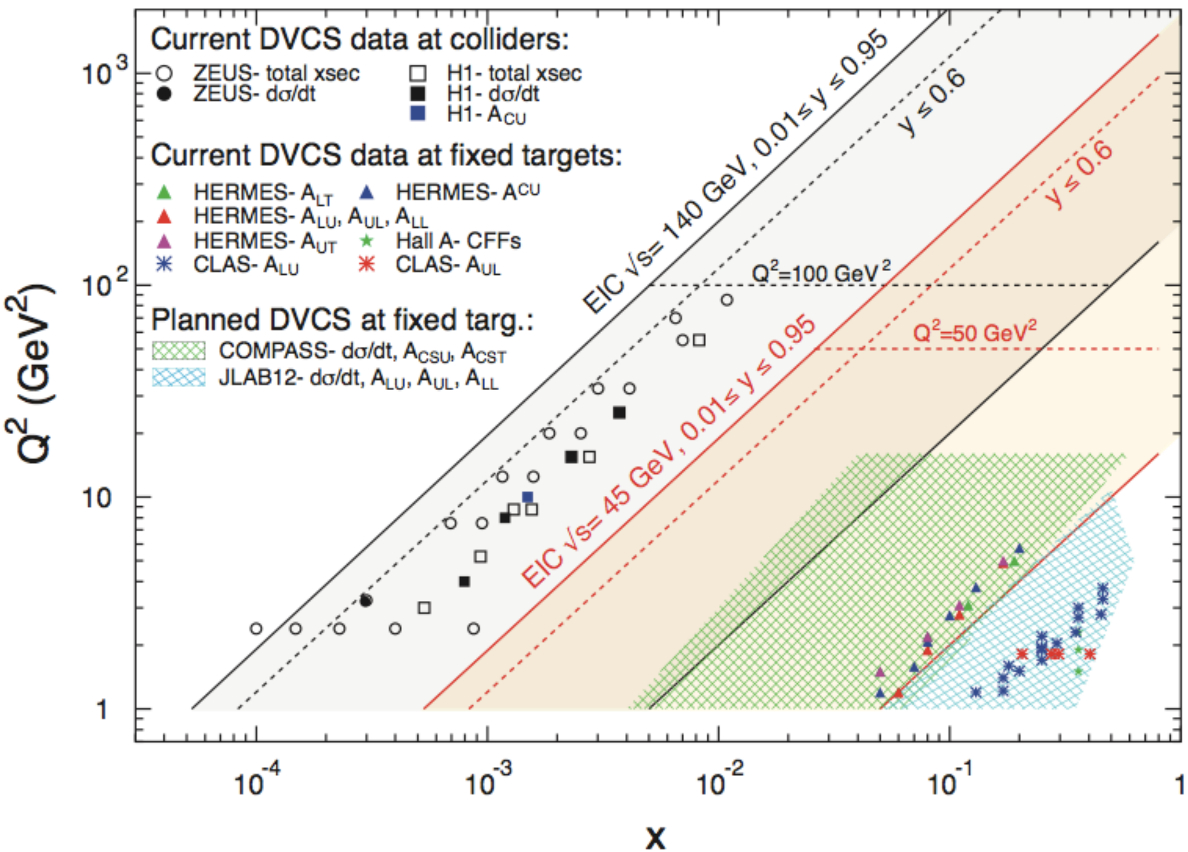
Bjorken variable $x_B = \dfrac{Q^2}{2\mathbf{p}_n \cdot \mathbf{q}}$

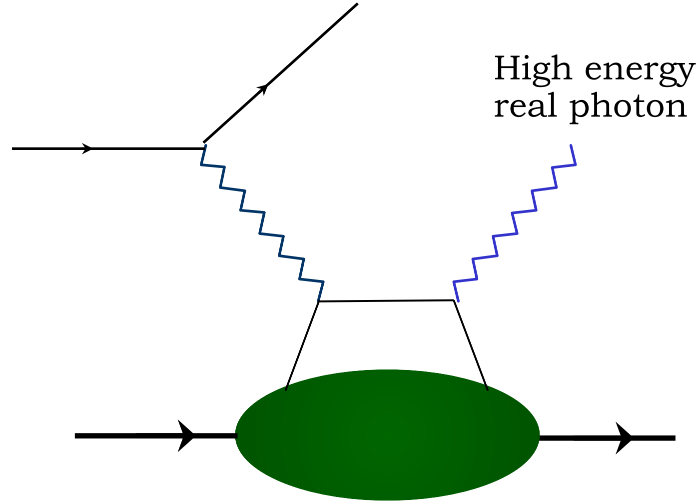
EIC White Paper, *Eur. Phy. J. A 52, 9 (2016)*

Nucleon tomography: imaging quarks

- Quark GPDs are accessible in a related process: Deeply Virtual Compton Scattering (DVCS)

DVCS kinematic reach at the EIC:





- 3D images of sea quark and gluon distributions from exclusive reactions: DVCS and DVMP.

EIC White Paper, *Eur. Phy. J. A 52, 9 (2016)*