

Illuminating Quarkonium Production Inclusive quarkonium photoproduction via UPC at the LHC

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Forward Physics and QCD at the LHC and EIC anr^{® STR}®NG

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 $\sqrt{s_{NN}} = 8.16 \text{ TeV} \rightarrow W_{\gamma N}^{max} \approx 1.5 \text{ TeV}; \sqrt{s_{NN}} = 13 \text{ TeV} \rightarrow W_{\gamma N}^{max} \approx 5 \text{ TeV}$

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- **How do we isolate photon interactions at the LHC? ultra-peripheral collisions**

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Motivation: inclusive quarkonium photoproduction at the LHC

- Owing to the large mass of its constituents, quarkonia are in principle the simplest QCD bound states
	- however, production mechanism remains an open question; Colour Singlet Model Colour Octet Mechanism and NRQCD Colour Evaporation Model
- **Photoproduction** is generally **simpler** than hadroproduction
- Measuring inclusive photoproduction presents the opportunity to constrain the production mechanism
- \bullet In advance of the EIC, p-Pb is the ideal system at the LHC since...
	- **no ambiguity** as to which beam particle emits the photon $[p \cdot p \cdot \text{or } Pb \cdot Pb]$
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- \bullet How do photoproduction vs. hadroproduction contributions compare?

- **Hadroproduction contribution is larger than photoproduction;** $\sigma_{\text{had.}} \gg \sigma_{\text{photo.}}$
- In order to make a measurement we must be able t[o k](#page-7-0)i[ll t](#page-9-0)[h](#page-8-0)[e](#page-7-0) h[a](#page-9-0)[dr](#page-0-0)[o](#page-37-0)[pr](#page-38-0)[od](#page-0-0)[u](#page-37-0)[ct](#page-38-0)[io](#page-0-0)[n](#page-37-0) $\text{Im} \rho \propto \rho$

Signal vs. background in detector acceptance

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Must impose cuts to enhance signal with respect to background!

Killing background: searching for the intact photon emitter with ZDC

• Far-forward detectors close to beam-pipe; used to classify centrality

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Minimum bias data (≥7 GeV in forward calorimeter) CMS 10.1088/1748-0221/16/05/P05008

- **•** Can resolve single to few neutron emissions
- **•** All of the signal is in the 0-neutron bump [signal with neutron emission is negligible]
- $\bullet \geq 1$ -neutron region is all background
- **Efficiency** (ϵ) for detecting 1n is $> 98\%$ cms, 2102.06640; ALICE, 1203.2436

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	- maximally 2% of 1n events look like 0n events

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Killing background: rapidity gaps

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 $W_{\gamma p}$: to know the collision energy

z : discriminant variable for quarkonium production mechanism (singlet vs. octet) and allows us to control the resolved-photon contribution

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KRAMER, hep-ph/016120

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- At the LHC the scattered photon-emitter is in the beam-pipe and cannot be measured. To learn about the photon energy must examine the final-state system.
	- In the exclusive case this is simple; detected particle gives the photon energy
	- This is not true for the inclusive case... how well can we reconstruct the final state?

KRAMER, hep-ph/016120

z-reconstruction depends on; (i) position of the detectors; (ii) kinematics of the event.

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Summary and outlook

- Quarkonium offer the chance to examine the interplay between perturbative and non-perturbative regime of the strong interaction
- The LHC can be used as a photon-nucleon collider
	- measuring inclusive J/ψ photoproduction at the LHC appears feasible which is complementary to existing HERA measurements
- \bullet In J/ψ photoproduction events in Pbp collisions
	- in CMS, ATLAS and ALICE the ZDC is sufficient to suppress background events
	- in LHCb a combination of gap and HeRSCheL based cuts are likely sufficient to suppress background
- \bullet The $\Delta \eta$ value at which the cut is placed allows for control over statistics and purity
- \bullet Both z and $W_{\gamma\rho}$ reconstruction appear possible with varying resolution which will allow control of the resolved contribution and offer the possibility to constrain the quarkonium production mechanism.

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Backup

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Set-up: tuning in p_T

No single model can simultaneously describe the photo- and hadroproduction data

- A correct description of the p_T description require combining different contributions, including NLO ones: not available in existing MC codes
- In order to be accurate in our MC event distribution, we tune leading order colour singlet and colour octet spectra to data in other colliding systems and extrapolate by changing the photon flux and PDFs
- The objective of our MC simulation is not about rate predictions but the characterisation of inclusive events

 $E|E \cap Q \cap Q$

Comput.Phys.Commun. 184 (2013) 2562-2570

- Use HELAC-Onia to generate MC samples [in the NRQCD framework]
- Use MC samples to model the signal and background
	- Signal $[\gamma g \rightarrow J/\psi({}^3S_1^1)g]$ and $[\gamma g \rightarrow J/\psi({}^1S_0^8)g]$
	- Background $[gg \to J/\psi(^3S_1^1)g$] and $[gg \to J/\psi(^3S_1^8)g$]
- Tune to data:
	- · photoproduction signal H1 ep 320 GeV data 10.1140/epjc/s10052-010-1376-5; 10.1007/s10052-002-1009-8
	- hadroproduction background LHCb 5 TeV pp data 10.1007/JHEP11(2021)181
- **•** Use PYTHIA to shower partonic events
- Characterise the inclusive signal and background using showered events

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Tune MC to rapidity integrated data (LHCb data @ 5 TeV). Assumptions:

2 Tuning is \sqrt{s} independant

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- Octet is reduced \bullet
- Singlet is increased \bullet

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Validation 1: tune vs. y-diff. data @ 5 TeV.

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Validation 1: tune vs. y-diff. data @ 5 TeV. 0 2 4 6 8 10 12 14 p_T [GeV] 10^{1} \vdots 10^2 $\frac{1}{2}$ 10^{3} $10⁴$ $\overline{}$ 10^5 10 $6₁$ $10'$ \leftarrow dN/ dp T[diM uon/ GeV] .. Octet - Singlet + Data $2 < v < 2.5x1$ 2.5<y<3.0x10 3.0<y<3.5x100 3.5<y<4.0x1000 4.0<y<4.5x10000

Tuning: photoproduction signal

Tune MC to HERA data @ $\sqrt{s} = 320$ GeV;

- **60** $< W_{\gamma p} < 240 \text{ GeV}$
- $0.3 < z < 0.9$

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NOTE: no tuning factor for octet in $0 < p_T < 1$ GeV as cross section is divergent. However, tuning factors can be computed using distributions from PYTHIA where events are smeared into the $0 < p_T < 1$ GeV region. 4 D F $E \cap Q$

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Per event z reconstruction

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\bullet\ \Delta z=z_{true}-z_{meas.}<0
$$

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- **.** Different contributions separated using experimental cuts ...
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We propose inclusive photoproduction is measured at the LHC; opportunity to extend p_T - & $W_{\gamma p}$ -reach, capture a variety of quarkonium species & improve statistical accuracy of existing data $E|E \cap Q$

Difference between models lies in the hadronisation description.

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- **Q** Colour Singlet Model
	- Heavy quarks $(Q\bar{Q})$ produced in the hard scattering have the same quantum numbers as the final quarkonium (Q) .
	- **NO gluon emissions** during hadronisation
	- $\sigma(Q) = \sigma(Q\bar{Q}) \times \langle \mathcal{O}^{\mathcal{Q}} \rangle$

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- **2** NRQCD and Colour Octet Mechanism
	- Higher Fock states (n) can contribute.
	- Each Fock state has a different hadronisation probability.
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- **3 Colour Evaporation Model**
	- Quantum numbers of $Q\bar{Q}$ decorrelated from Q .
	- \bullet Only the invariant mass of the $Q\bar{Q}$ is constrained.
	- **•** Semi-soft gluon emissions during hadronisation
	- $\sigma({\cal Q}) = \int \frac{\sigma(Q\bar{Q})}{d m_{Q\bar{Q}}}dm_{Q\bar{Q}}$

 $F = \Omega$

From p to Pb in the HeRSCheL region

- \bullet The background is modelled by generating pA events with HELAC-Onia and passing them through PYTHIA; PYTHIA reads these as pp events.
- In a pp collision $N_{coll.} = 1$; whereas in a pA collision there are many more nucleons and therefore it is possible to have $N_{coll.} > 1$ [typically modelled using Glauber-type models].
- Using minimum bias events generated by PYTHIA, one can obtain a probability distribution for the number of charged tracks in the HeRSCheL region. [bottom left]
- To model the HeRSCheL signal using the PYTHIA events (i.e., converting pp to pA) events are randomly assigned a centrality class and then assigned $N_{coll.}$ based on ALICE results. [bottom centre arXiv:1605.05680]
- For a given event, the total number of charged tracks in the HeRSCheL region is given by throwing $i = 1, \ldots, N_{coll.} - 1$ points into the probability distribution, and summing over $N_{coll.}$.
- The transformation from pp to pA HeRSCheL distribution. [bottom right]

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Rapidity-differential gap distributions in LHCb pPb

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Rapidity-differential gap distributions in LHCb Pbp

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In a given kinematic region, the percentage error on z-reconstruction at one standard deviation.

Note: $Δz/z = (z - z_{exp.})/z < 0$.

 $F = \Omega Q$

Diffractive production

- Colourless exchange
- Only CSM contributes
- exclusive: only J/ψ decay products

Inclusive production

- Hard final state gluon
- Resolved vs. direct contribution
- Test production mechanism
- Probe gluon PDF

Lightcone four-vector representation

1 Choose two vectors along an axis such that,

$$
\eta^{\pm} \cdot \eta^{\pm} = 0 \quad \& \quad \eta^{\mp} \cdot \eta^{\pm} = 2. \tag{1}
$$

² A particle's four-momentum can be written as,

$$
p = (E, p_x, p_y, p_z) = [P^+, P^-, \mathbf{p}].
$$
 (2)

3 The scalar product of two four-momenta is given as,

$$
p \cdot q = \frac{1}{2} \left(P^+ Q^- + P^- Q^+ \right) - \mathbf{p} \cdot \mathbf{q}.
$$
 (3)

 \bullet If ρ lies along the vector η^- , then the scalar product reduces to,

$$
p \cdot q = \frac{1}{2} \left(P^- Q^+ \right). \tag{4}
$$

5 Consider some massless particle q,

- If q lies on the vector η^+ : $p \cdot q$ is maximised $\rightarrow p \cdot q = A$.
- If q is perpendicular to the vectors η^{\pm} : $p \cdot q = A/2$.
- If [q](#page-67-0) lies on the vector η^- : $p \cdot q$ $p \cdot q$ is minimis[ed](#page-65-0) $\rightarrow p_{\hat{\sigma}} q = 0$ $\rightarrow p_{\hat{\sigma}} q = 0$ $\rightarrow p_{\hat{\sigma}} q = 0$.

NLO inclusive J/ψ photoproduction at HERA

arXiv:2107.13434

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ATLAS UPC dijet Study

ATLAS-CONF-2022-021

- Pb-Pb @ $\sqrt{s_{NN}} = 5.02 \text{ TeV}$
	- 0 n X n requirement [$\varepsilon_{ZDC} <$ 1 TeV]
	- $\sum_{\gamma} \Delta \eta$ requirement [instead of $\Delta \eta_\gamma^{edge}$]
		- \bullet Include resolved photon in analysis
		- What is the effect of higher order corrections on choice of gap definition?

Event topology (experimental)

Luminosity targets taken from LHC programme coordination meeting; pPb and PbPb targets are for Run 3 and 4 and pp targets are for Run 3 only.

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