



U. S. DEPARTMENT OF  
**ENERGY**

Office of  
Science



Creighton  
UNIVERSITY

# Studying photonuclear cross sections with UPCs in ALICE

Simone Ragoni for the ALICE Collaboration  
Creighton University, USA

# Outline

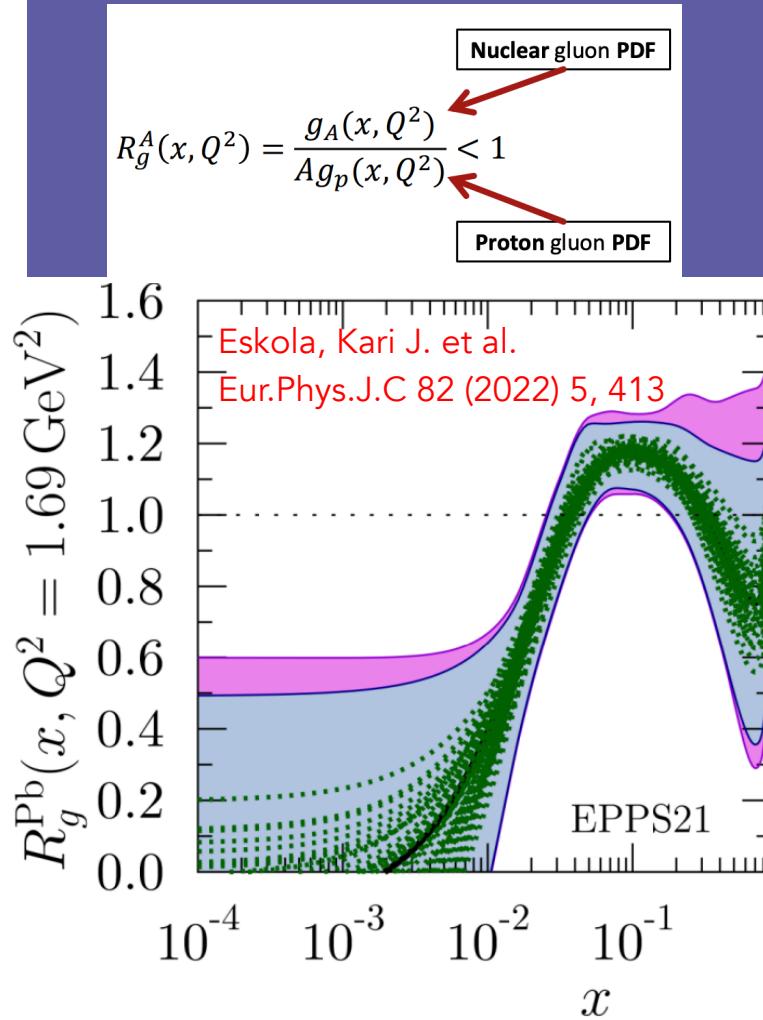
- Physics of ultra-peripheral collisions (UPC)
- The ALICE detector
- Measurements of the energy dependence of the photonuclear cross sections
  - Rapidity-differential measurements
  - Peripheral photoproduction
  - Neutron emission

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# Physics of ultra-peripheral collisions (UPC)

Studying nuclear shadowing

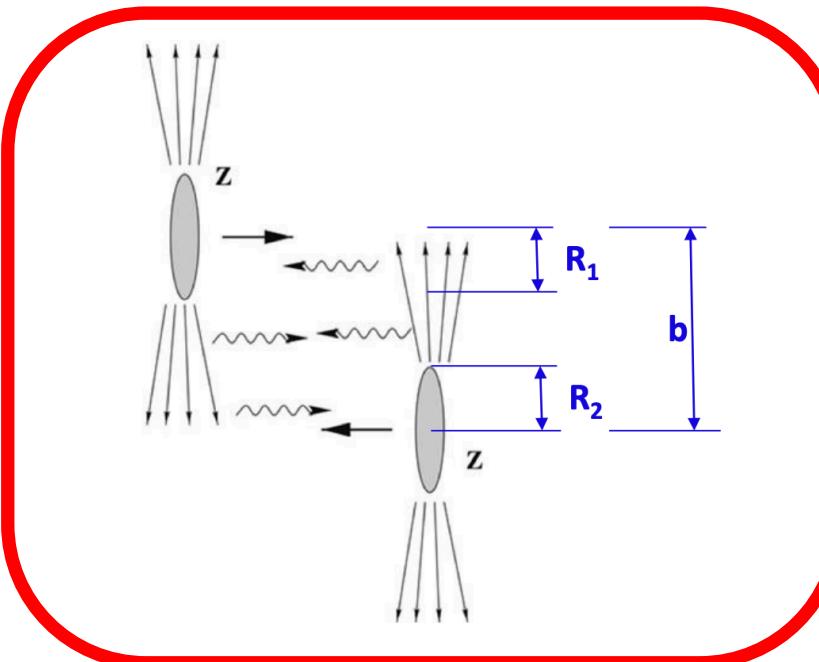


LO

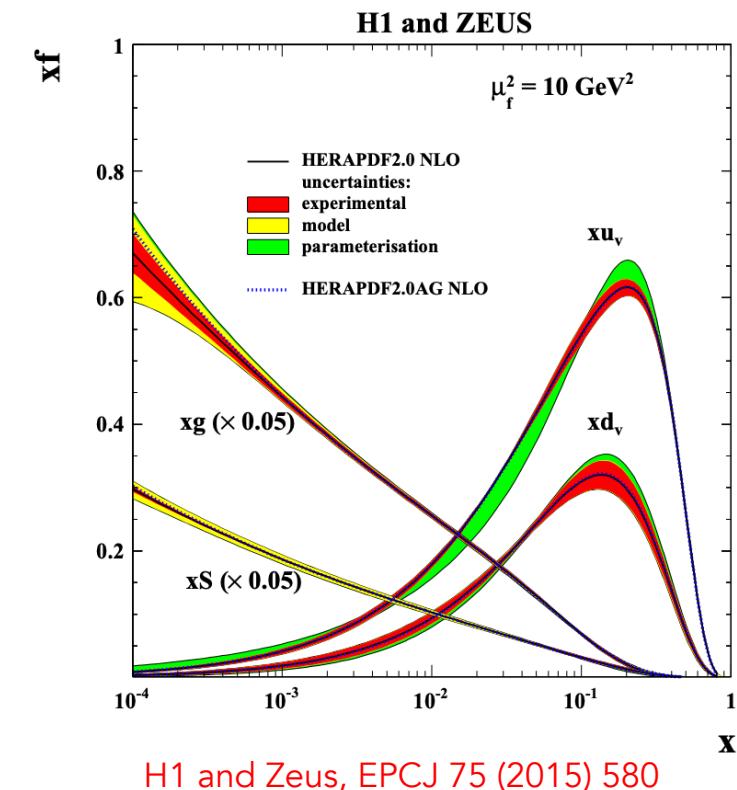
$$\frac{d\sigma^T(\gamma p \rightarrow J/\Psi + p)}{dt} = \frac{|M|^2}{16\pi s^2}$$

$$= [F_N^{2G}(t)]^2 \frac{\alpha_s^2 \Gamma_{ee}^J m_J^3}{3\alpha_{e.m.}} \pi^3 \left[ \bar{x} G(\bar{x}, \bar{q}^2) \frac{2\bar{q}^2 - |q_t^J|^2}{(2\bar{q}^2)^3} \right]^2$$

Ryskin: Z. Phys. C 57, 89-92 (1993)



Looking for gluon saturation



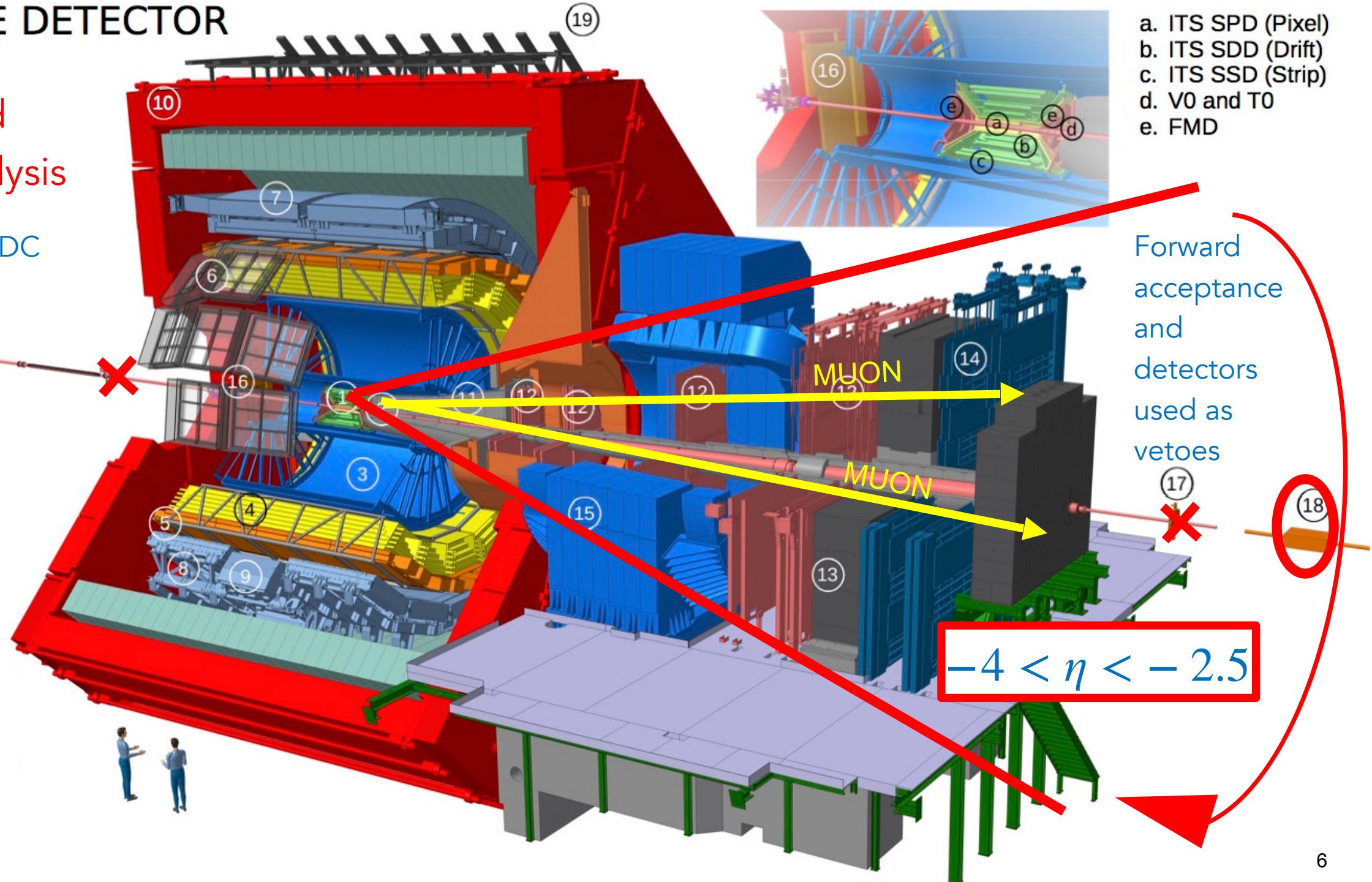
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# THE ALICE DETECTOR

Forward  
 $J/\psi$  analysis

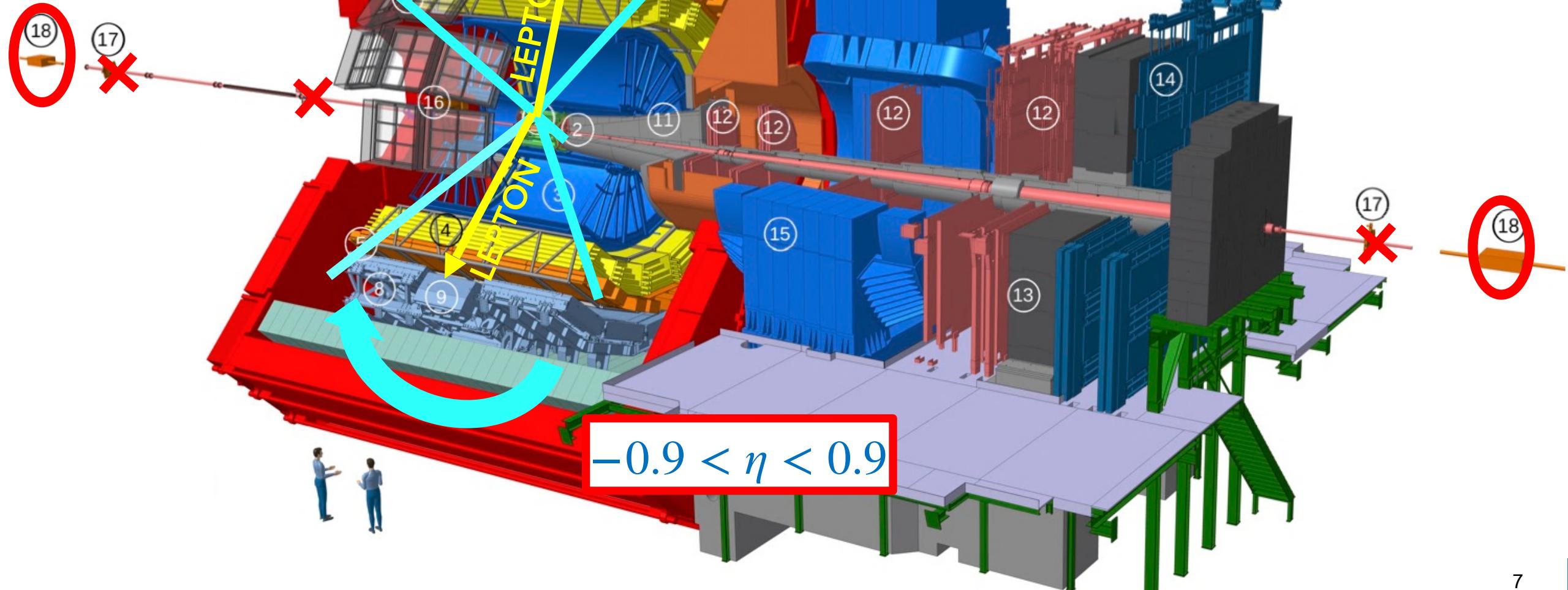
Vetoes and ZDC



# THE ALICE DETECTOR

Midrapidity  
 $J/\psi$  analysis

Vetoes and ZDC

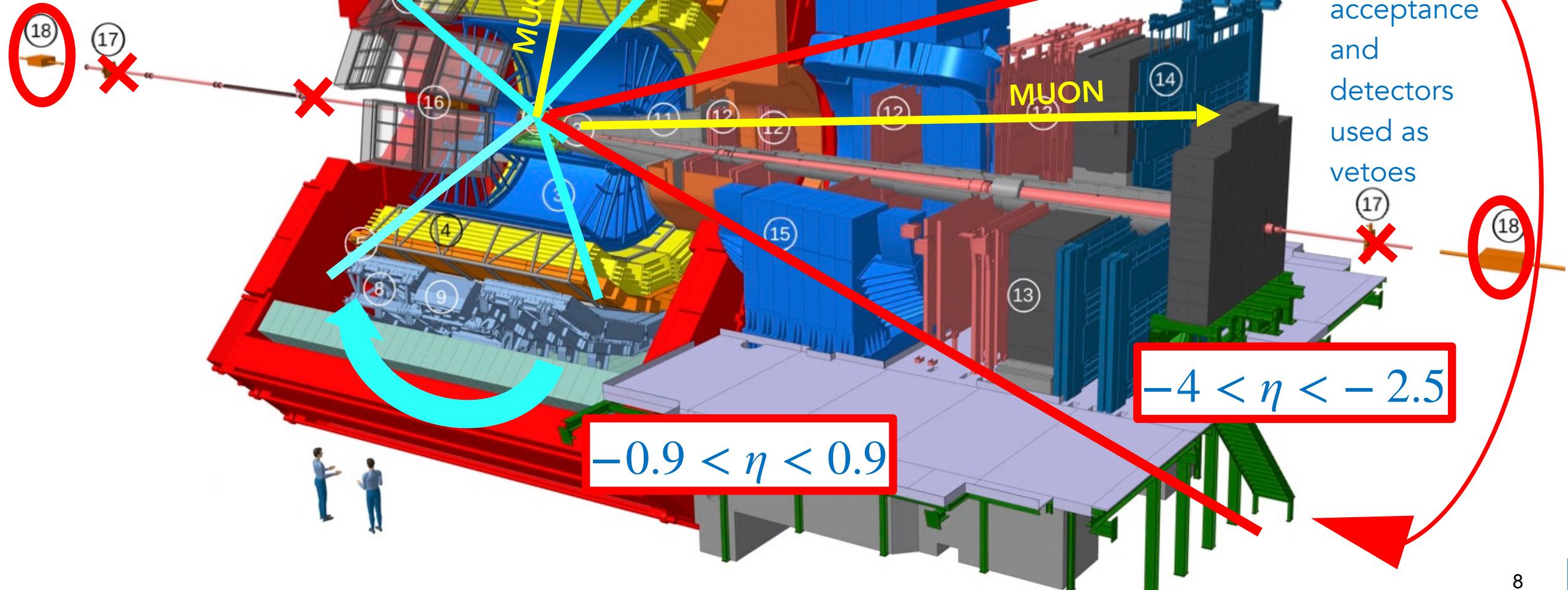


- a. ITS SPD (Pixel)
- b. ITS SDD (Drift)
- c. ITS SSD (Strip)
- d. V0 and T0
- e. FMD

# THE ALICE DETECTOR

Semiforward  
configuration

Vetoes and ZDC



# Outline

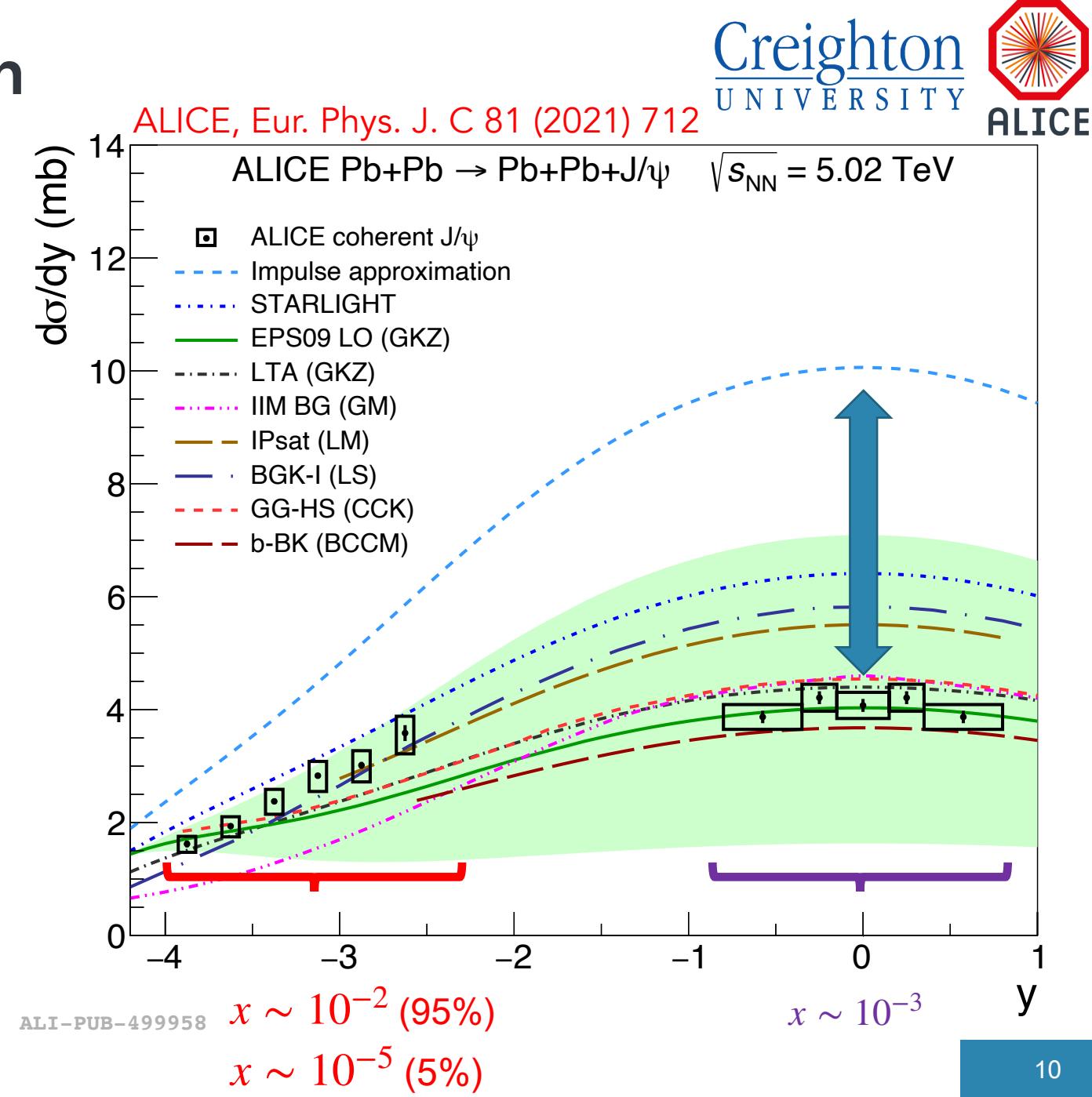
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# Coherent J/ $\psi$ cross section

- ALICE data exhibit moderate nuclear shadowing
- Nuclear suppression factor

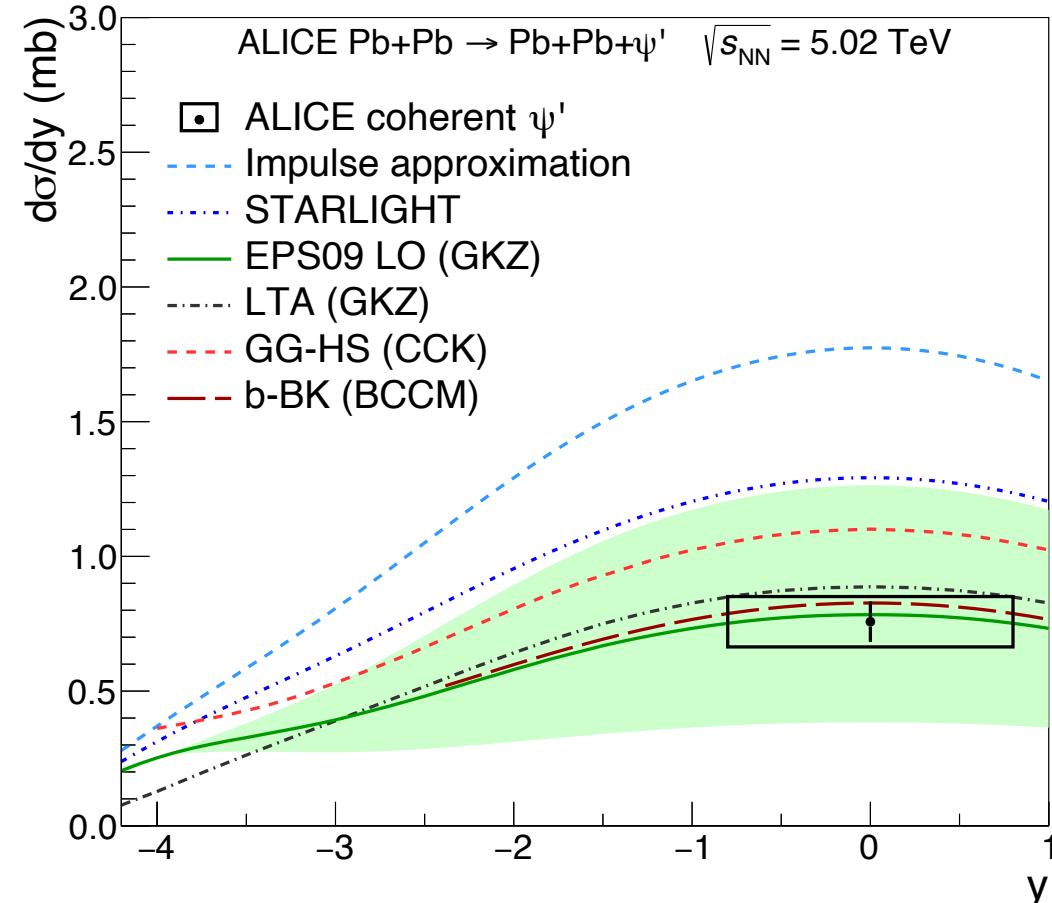
$$S_{\text{Pb}}(y \sim 0) = \sqrt{\frac{d\sigma}{dy}_{\text{data}} / \frac{d\sigma}{dy}_{IA}} = 0.64 \pm 0.04$$

- IA = impulse approximation (no nuclear effects)
- $S(W_{\gamma p})$  - nuclear suppression factor - provides a way to test the consistency of the data with the available nuclear and nucleon PDFs and to measure the nuclear shadowing factor

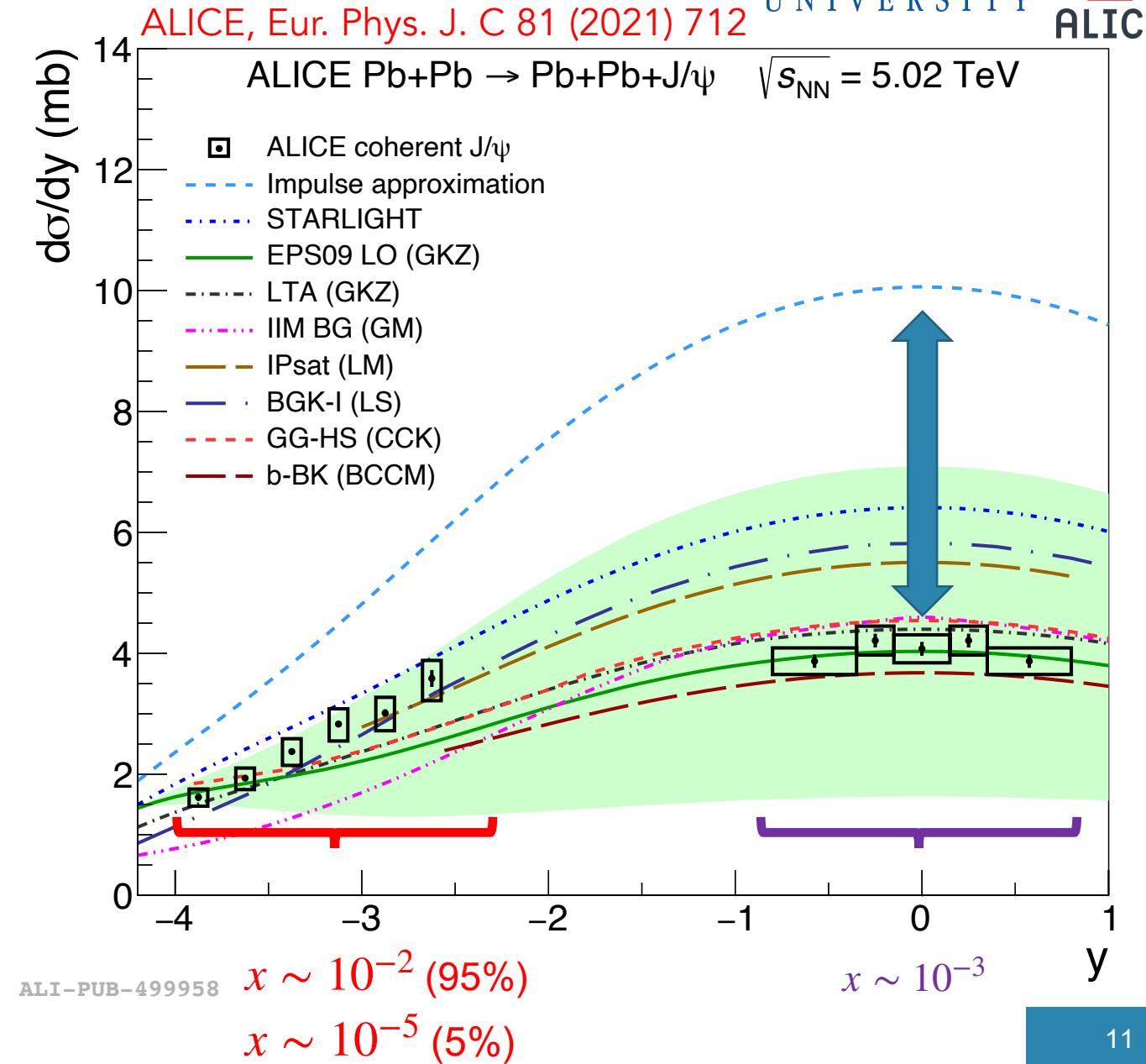


# Coherent J/ $\psi$ and $\psi'$ cross section

$$S_{\text{Pb}}(y \sim 0, \psi') = \sqrt{\frac{d\sigma}{dy}_{\text{data}} / \frac{d\sigma}{dy}_{IA}} = 0.66 \pm 0.6$$



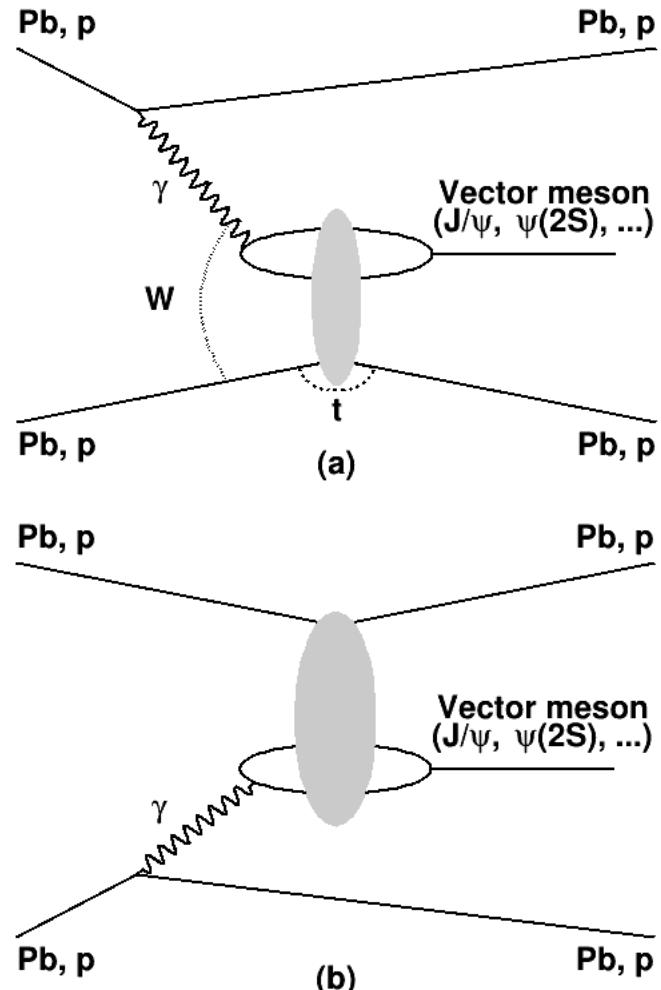
ALI-PUB-499963



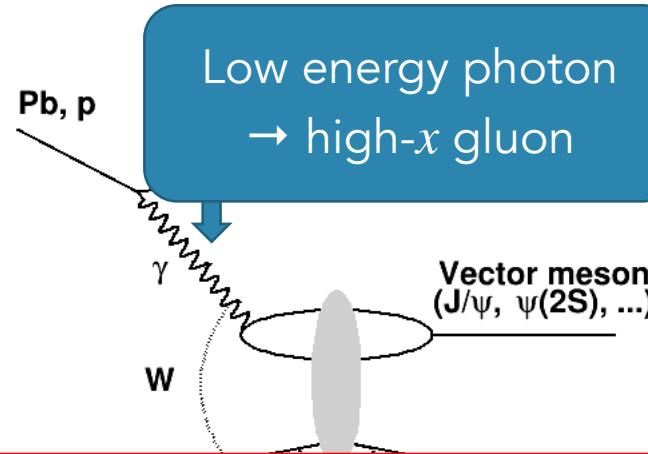
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# Measuring the energy dependence

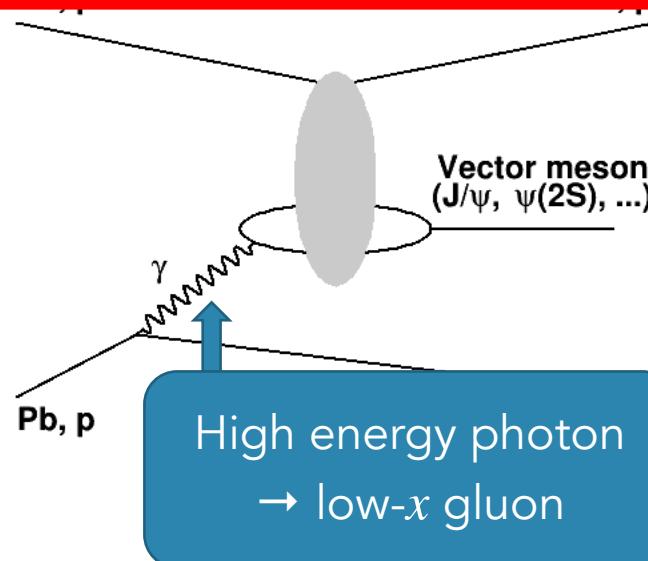


# Measuring the energy dependence



$$x = \frac{M_{J/\psi}^2}{W_{\gamma Pb,n}^2}$$

$$\frac{d\sigma}{dy} = n(\gamma, +y)\sigma_{\gamma Pb}(+y) + n(\gamma, -y)\sigma_{\gamma Pb}(-y)$$



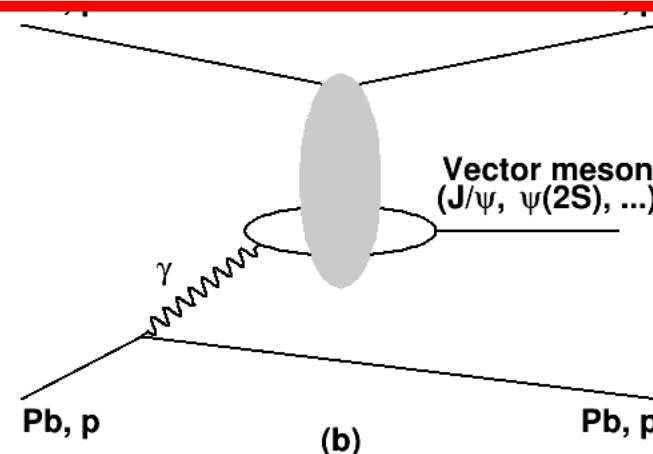
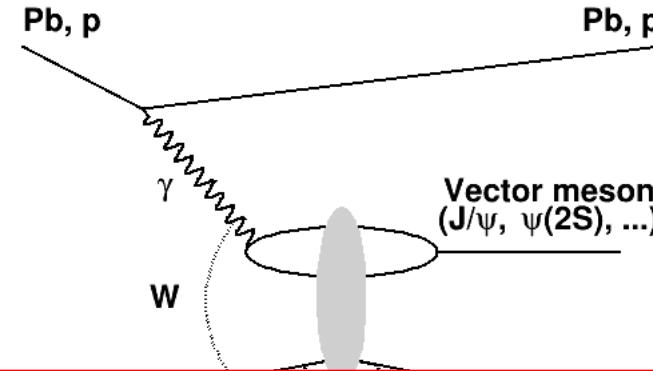
# Measuring the energy dependence

Coherent J/ $\psi$  at midrapidity:

- UPC cross sections can be linked directly to the photonuclear cross section

$$\frac{d\sigma}{dy} = 2n(\gamma)\sigma_{\gamma Pb}$$

$$\frac{d\sigma}{dy} = n(\gamma, +y)\sigma_{\gamma Pb}(+y) + n(\gamma, -y)\sigma_{\gamma Pb}(-y)$$



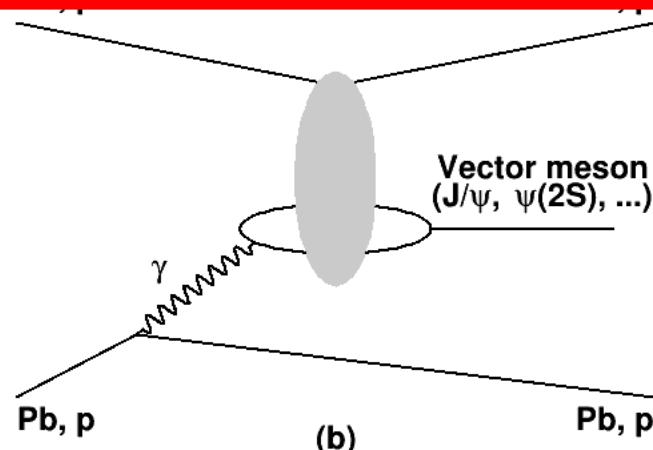
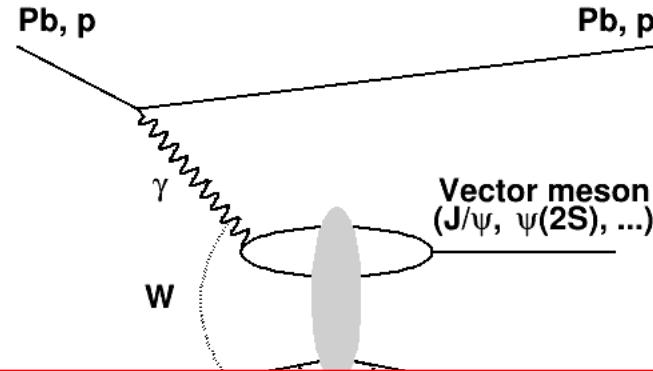
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Coherent J/ $\psi$  at forward rapidity:

- 95% of the cross section originates from the low-energy photon

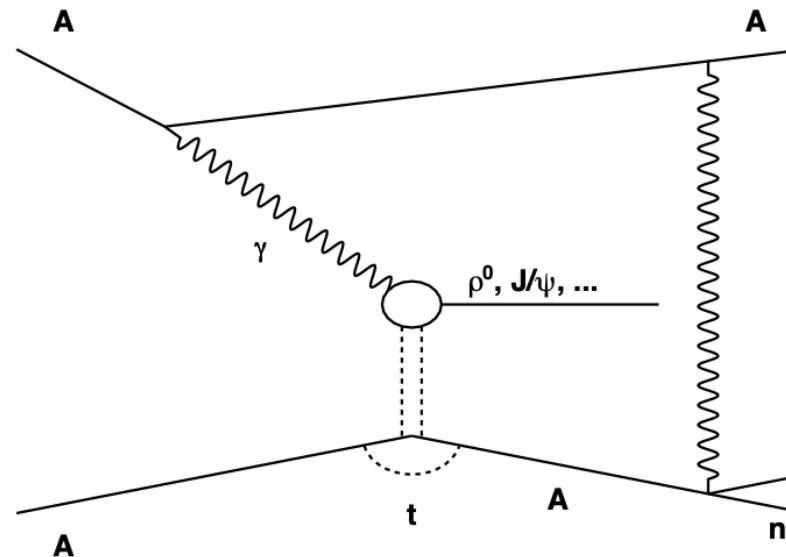
$$\frac{d\sigma}{dy} \simeq n(\gamma)\sigma_{\gamma Pb}$$

But that remaining 5% is very interesting

# Techniques for the photon direction ambiguity

## Neutron emission:

- $x = \frac{M_{\text{VM}}}{\sqrt{s_{\text{NN}}}} \cdot e^{\pm y}$
- Ambiguity due to sign in the rapidity of the photon emitter  $\rightarrow 10^{-2}, 10^{-5}$
- Additional photon exchanges may lead to neutron emission

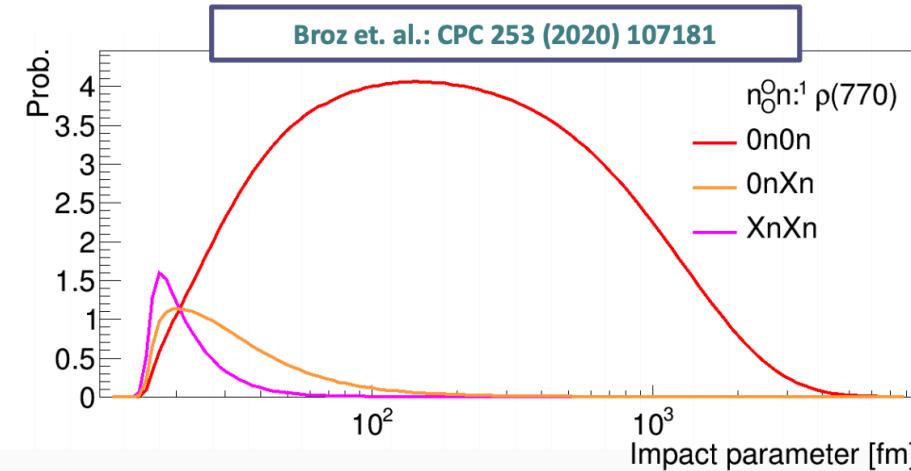


- Using the neutron ZDCs on the A and C side to detect the neutrons!
- E.g. 0N0N: no neutrons on either ZDCs
- E.g. 0NXN: neutrons only on one side

# Techniques for the photon direction ambiguity

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- Additional photon exchanges may lead to neutron emission

$$\frac{d\sigma_{PbPb}^{0N0N}}{dy} = n_{0N0N}(\gamma, +y) \cdot \sigma_{\gamma Pb}(+y) + n_{0N0N}(\gamma, -y) \cdot \sigma_{\gamma Pb}(-y)$$
$$\frac{d\sigma_{PbPb}^{0NXN}}{dy} = n_{0NXN}(\gamma, +y) \cdot \sigma_{\gamma Pb}(+y) + n_{0NXN}(\gamma, -y) \cdot \sigma_{\gamma Pb}(-y)$$

Guzey et al.,  
Eur.Phys.J.C 74 (2014) 7, 2942

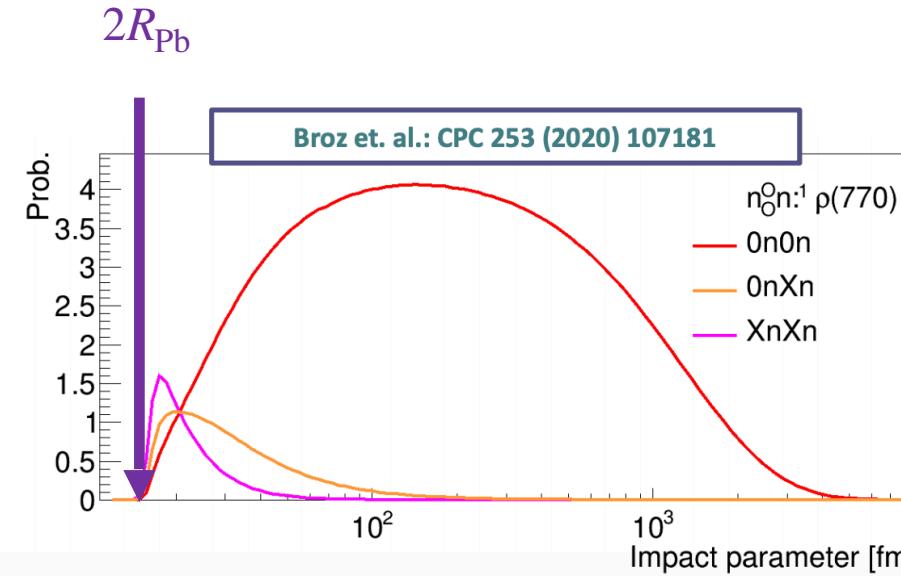
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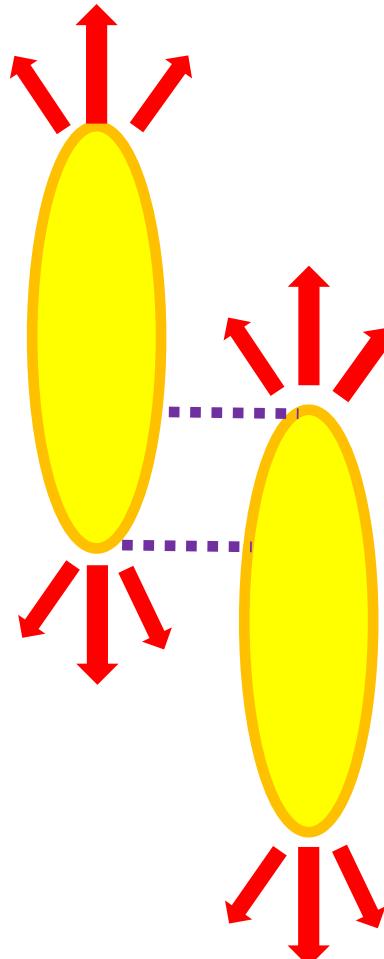
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# Techniques for the photon direction ambiguity

Peripheral  
photoproduction:

- $b < R_1 + R_2$
- Hadronic interactions + photoproduction



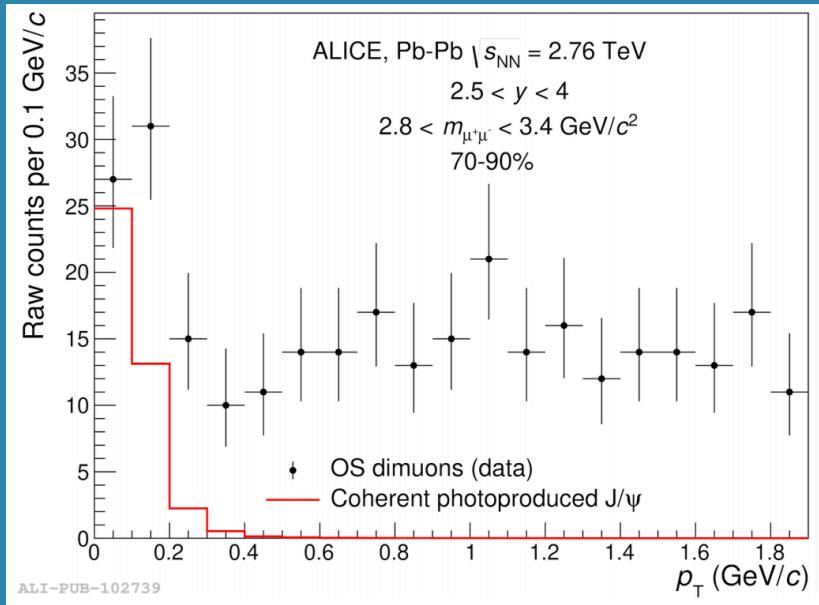
- Simultaneously use UPC and peripheral results to get rid of the ambiguities!

$$\frac{d\sigma_{PbPb}^P}{dy} = n_P(\gamma, +y) \cdot \sigma_{\gamma Pb}(+y) + n_P(\gamma, -y) \cdot \sigma_{\gamma Pb}(-y)$$
$$\frac{d\sigma_{PbPb}^U}{dy} = n_U(\gamma, +y) \cdot \sigma_{\gamma Pb}(+y) + n_U(\gamma, -y) \cdot \sigma_{\gamma Pb}(-y)$$

J.G. Contreras PRC 96 (2017) 015203

# Peripheral J/ $\psi$ photoproduction using ALICE Run 1 data

- First observed with Run 1 data by ALICE, then with Run 2 statistics by both ALICE and LHCb, and by STAR
- First extraction of the photonuclear cross sections with Run 1 data

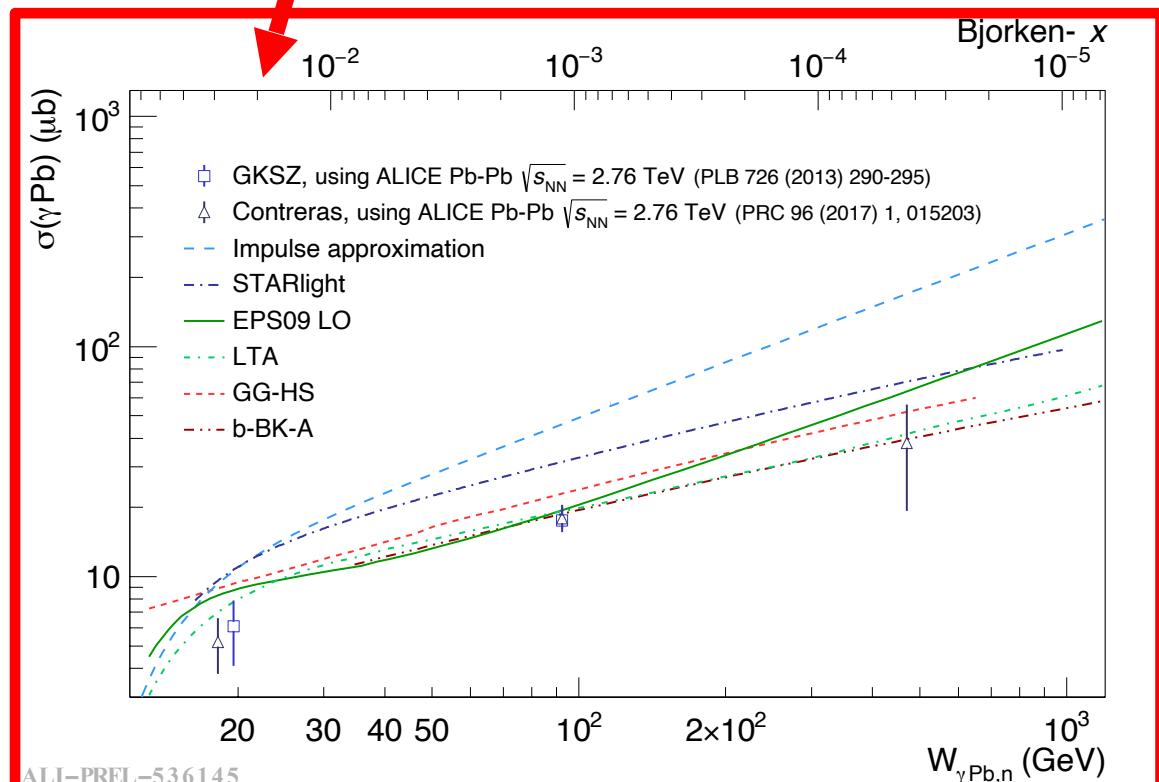


ALICE, PRL 116 (2016), 222301 LHCb, Phys.Rev.C 105 (2022) 3, L032201  
ALICE, arXiv:2204.10684 STAR, Phys.Rev.Lett. 123 (2019) 132302

Peripheral results → 
$$\frac{d\sigma_{PbPb}^P}{dy} = n_P(\gamma, +y) \cdot \sigma_{\gamma Pb}(+y) + n_P(\gamma, -y) \cdot \sigma_{\gamma Pb}(-y)$$

UPC results → 
$$\frac{d\sigma_{PbPb}^U}{dy} = n_U(\gamma, +y) \cdot \sigma_{\gamma Pb}(+y) + n_U(\gamma, -y) \cdot \sigma_{\gamma Pb}(-y)$$

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# Coherent J/ $\psi$ with neutron emission

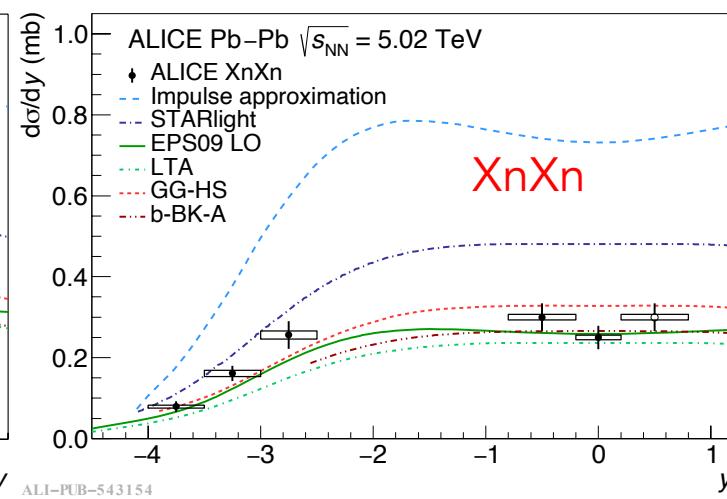
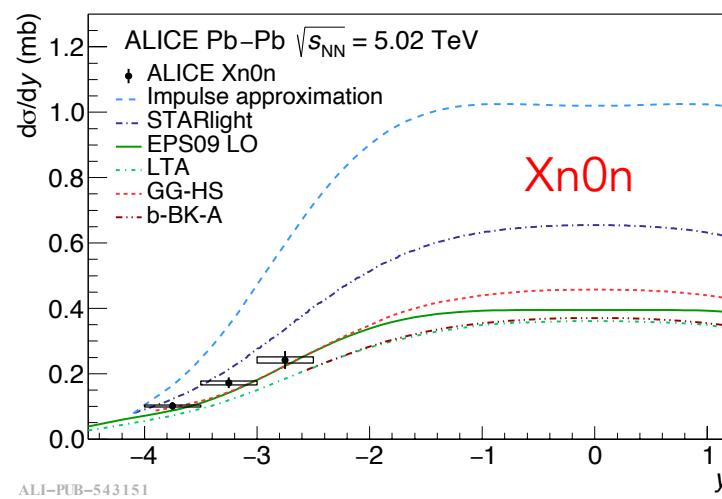
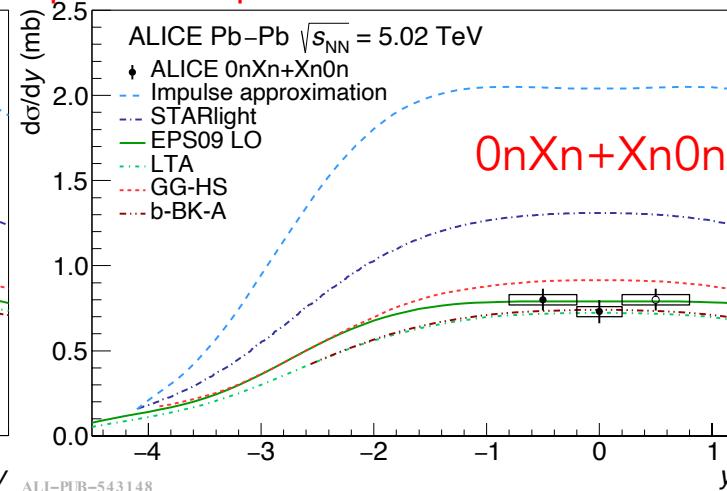
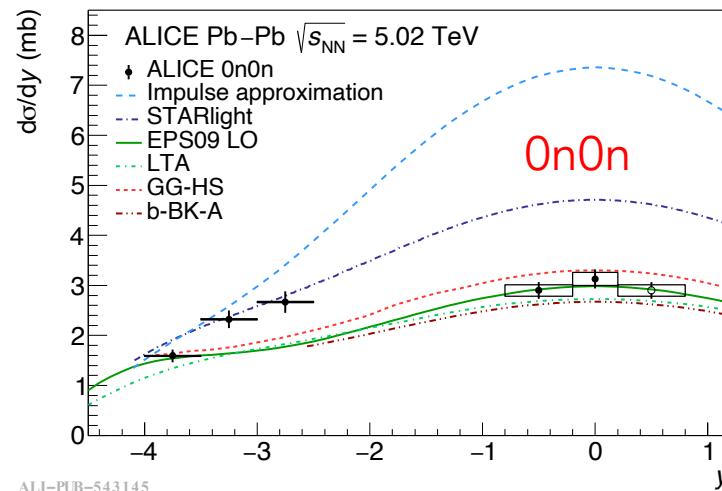
$$\frac{d\sigma_{PbPb}^{0N0N}}{dy} = n_{0N0N}(\gamma, +y) \cdot \sigma_{\gamma Pb}(+y) + n_{0N0N}(\gamma, -y) \cdot \sigma_{\gamma Pb}(-y)$$

$$\frac{d\sigma_{PbPb}^{0NXN}}{dy} = n_{0NXN}(\gamma, +y) \cdot \sigma_{\gamma Pb}(+y) + n_{0NXN}(\gamma, -y) \cdot \sigma_{\gamma Pb}(-y)$$



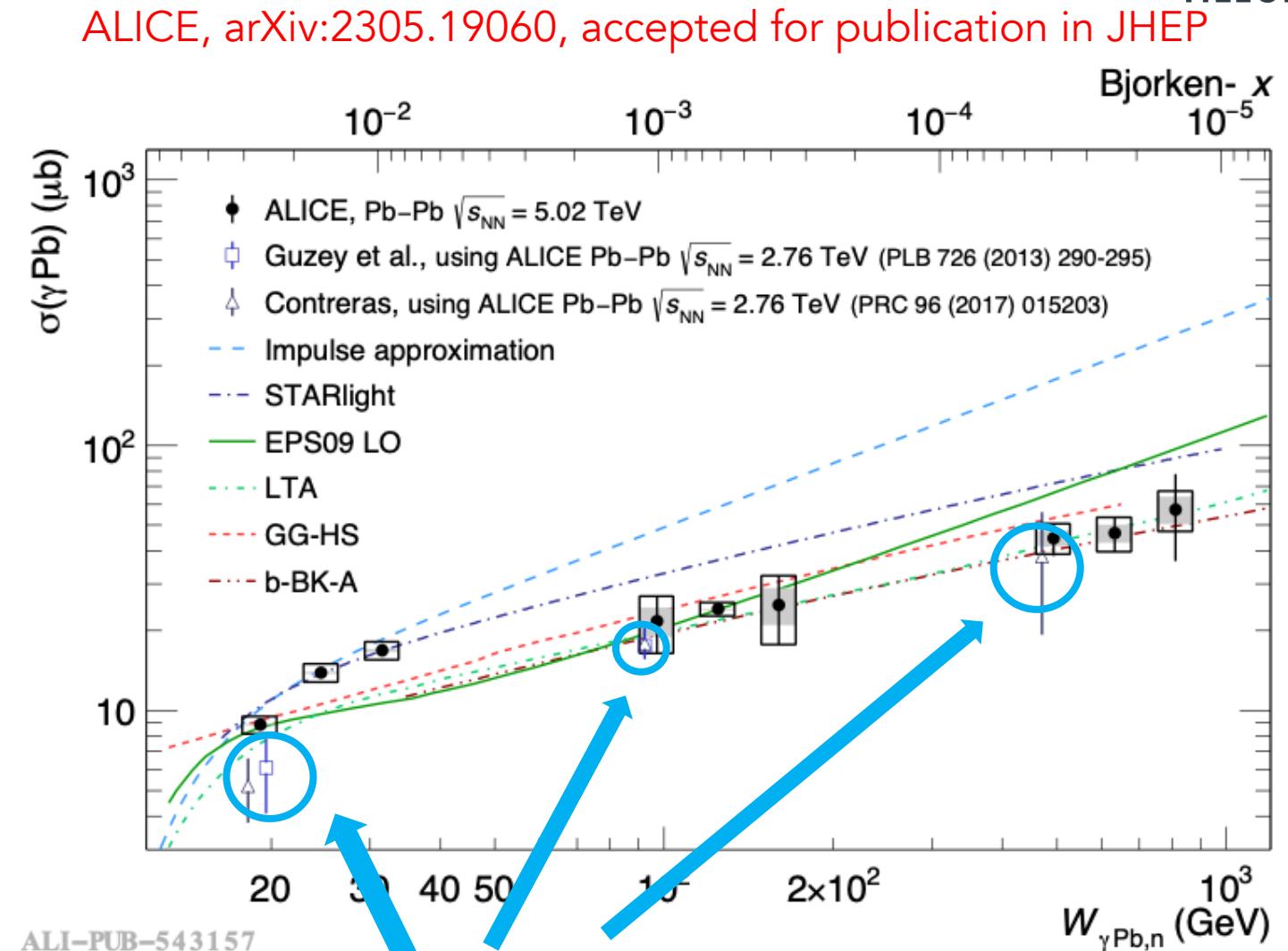
$y$	$n_\gamma(0n0n)$	$n_\gamma(0nXn+Xn0n)$	$n_\gamma(XnXn)$	$\sigma_{\gamma Pb}^{\text{IA}} (\mu\text{b})$
$3.5 < y < 4$	178.51	18.18	6.34	10
$3 < y < 3.5$	162.99	18.19	6.34	14
$2.5 < y < 3$	147.46	18.19	6.34	19
$0.2 < y < 0.8$	77.88	17.88	6.33	48
$-0.2 < y < 0.2$	62.86	17.47	6.27	58
$-0.8 < y < -0.2$	48.31	16.75	6.18	71
$-3 < y < -2.5$	3.91	4.97	2.78	176
$-3.5 < y < -3$	1.22	2.15	1.42	215
$-4 < y < -3.5$	0.26	0.61	0.48	262

ALICE, arXiv:2305.19060, accepted for publication in JHEP



# Coherent J/ $\psi$ with neutron emission

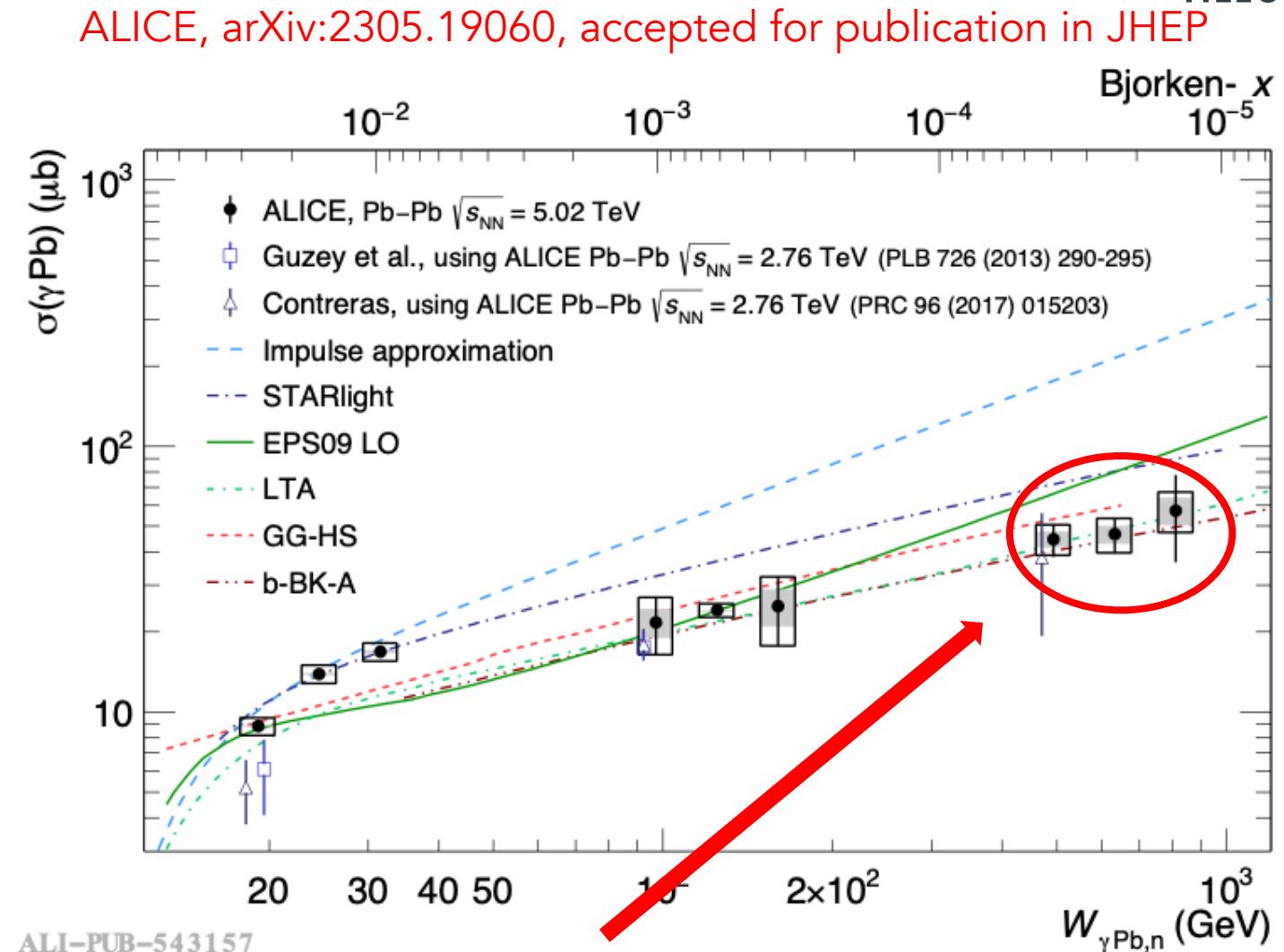
- First measurement of the energy dependence of the photonuclear cross section down to Bjorken-  $x \sim 10^{-5}$ !
- At low- $x$  data favours both saturation and shadowing models



Good agreement between Run 1 peripheral  
and Run 2 neutron emission measurements!

# Coherent J/ $\psi$ with neutron emission

- First measurement of the energy dependence of the photonuclear cross section down to Bjorken-  $x \sim 10^{-5}$ !
- At low- $x$  data favours both saturation and shadowing models
- New Run 2 results probe unprecedented Bjorken- $x$  region like no other LHC experiment!



Neutron emission extends the range in energy being explored by about 300 GeV!

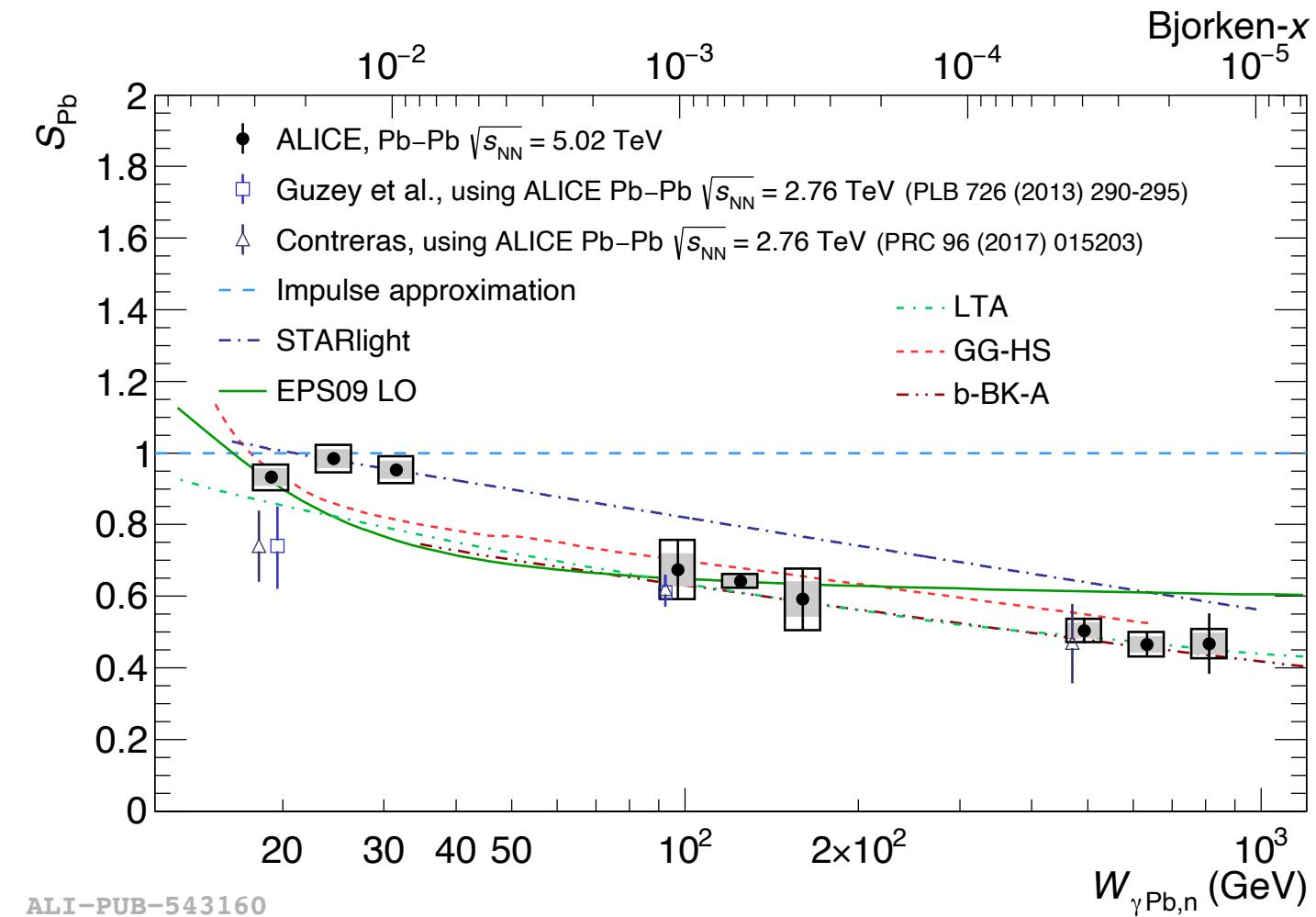
# Coherent J/ $\psi$ with neutron emission

ALICE, arXiv:2305.19060, accepted for publication in JHEP

- First measurement of the nuclear suppression factor at Bjorken- $x \sim 10^{-5}$ !

$$S_{\text{Pb}}(y) = \sqrt{\frac{d\sigma}{dy}_{\text{data}} / \frac{d\sigma}{dy}_{IA}}$$

- At low- $x$  data favours both saturation and shadowing models
- Additional theoretical uncertainty from impulse approximation → dominates at low energies

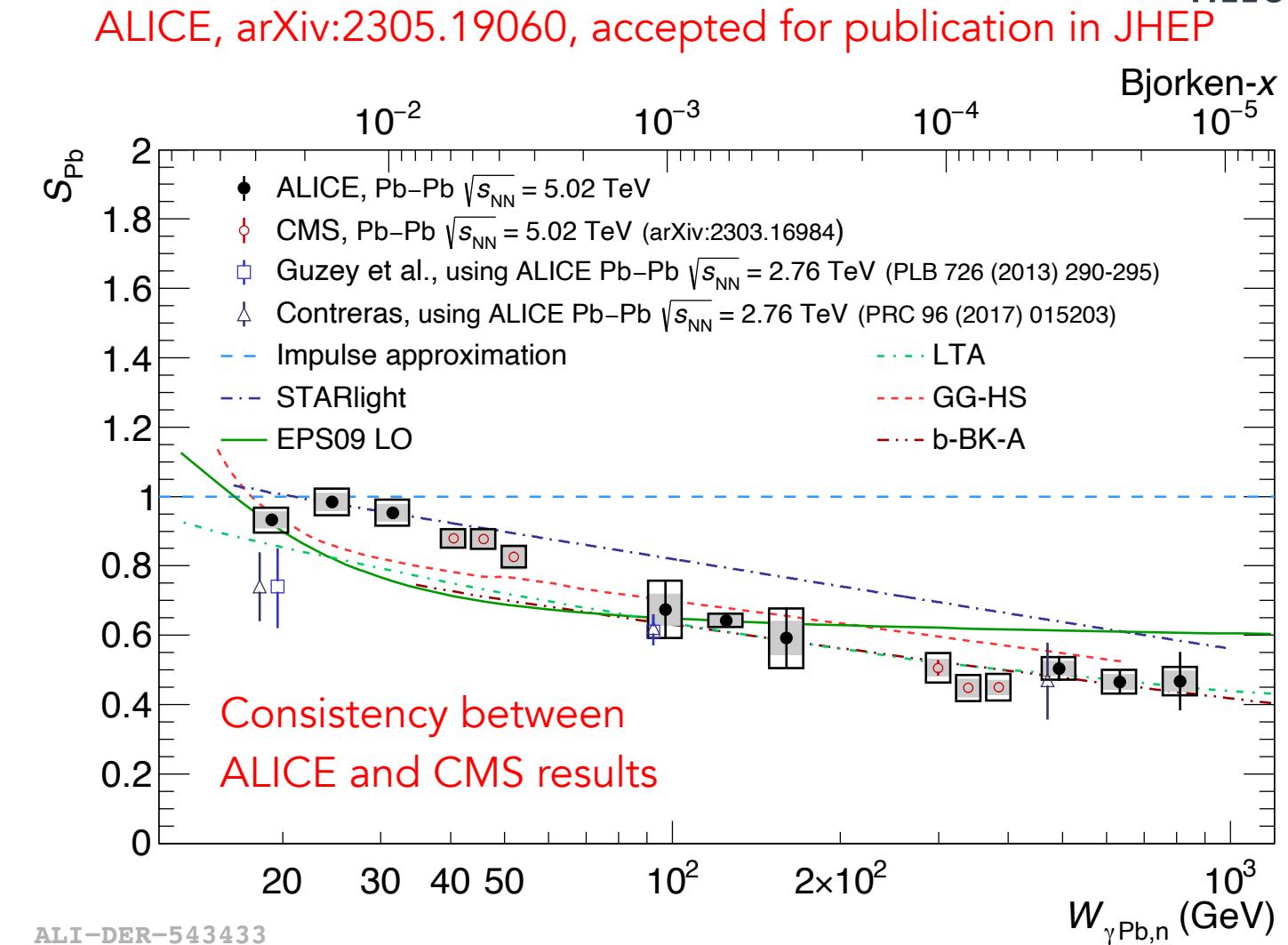


# Coherent J/ $\psi$ with neutron emission

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# ALICE in Run 3 and 4

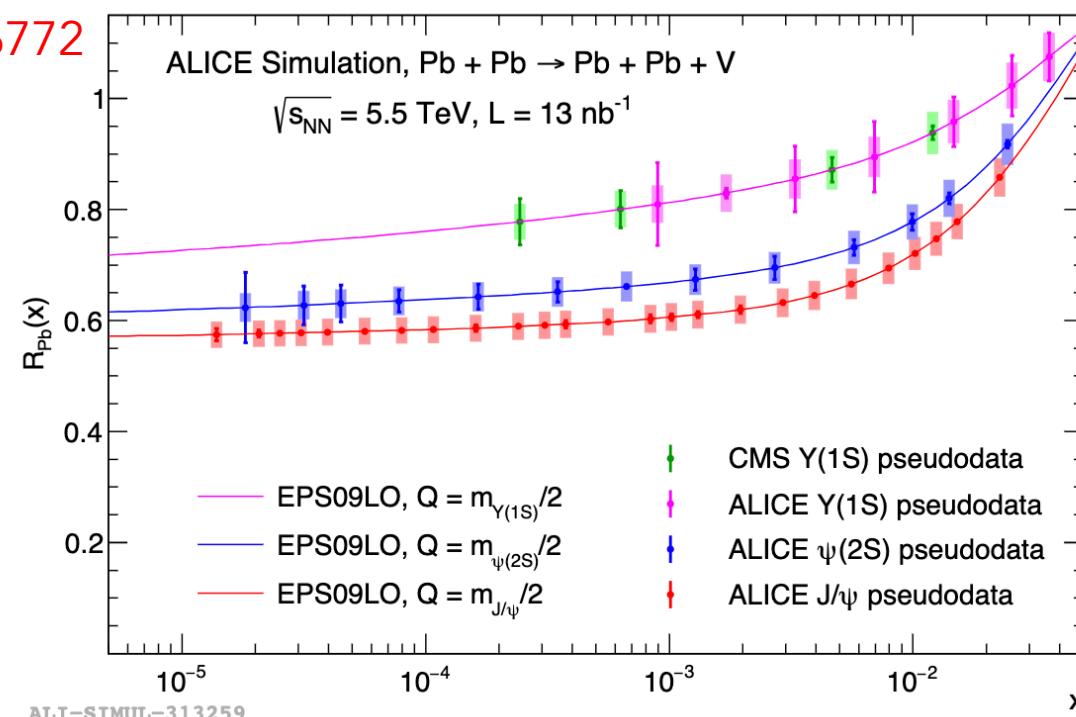
- Significant increase in integrated lumi from  $1 \text{ nb}^{-1}$  for Run 2 to  $13 \text{ nb}^{-1}$  for Run 3 and Run 4 together
- Great increase in statistics with continuous readout
- Uncertainties for nuclear suppression factor expected to be at the level of 4%
  - Nuclear shadowing studied as a function of  $x$  and  $Q^2$
- New measurements e.g. bottomonium states
- MFT in Run 3, FoCal in Run 4!

Meson	$\sigma$	Central 1 Total	Forward 1 Total 1
$\rho \rightarrow \pi^+ \pi^-$	5.2b	5.5 B	4.9 B
$\rho' \rightarrow \pi^+ \pi^- \pi^+ \pi^-$	730 mb	210 M	190 M
$\phi \rightarrow K^+ K^-$	0.22b	82 M	15 M
$J/\psi \rightarrow \mu^+ \mu^-$	1.0 mb	1.1 M	600 K
$\psi(2S) \rightarrow \mu^+ \mu^-$	30 $\mu b$	35 K	19 K
$Y(1S) \rightarrow \mu^+ \mu^-$	2.0 $\mu b$	2.8 K	880

CERN Yellow Rep. Monogr. 7  
(2019) 1159-1410, arXiv

1812.06772

$|y| < 0.9$     $2.5 < |y| < 4$



# Summary

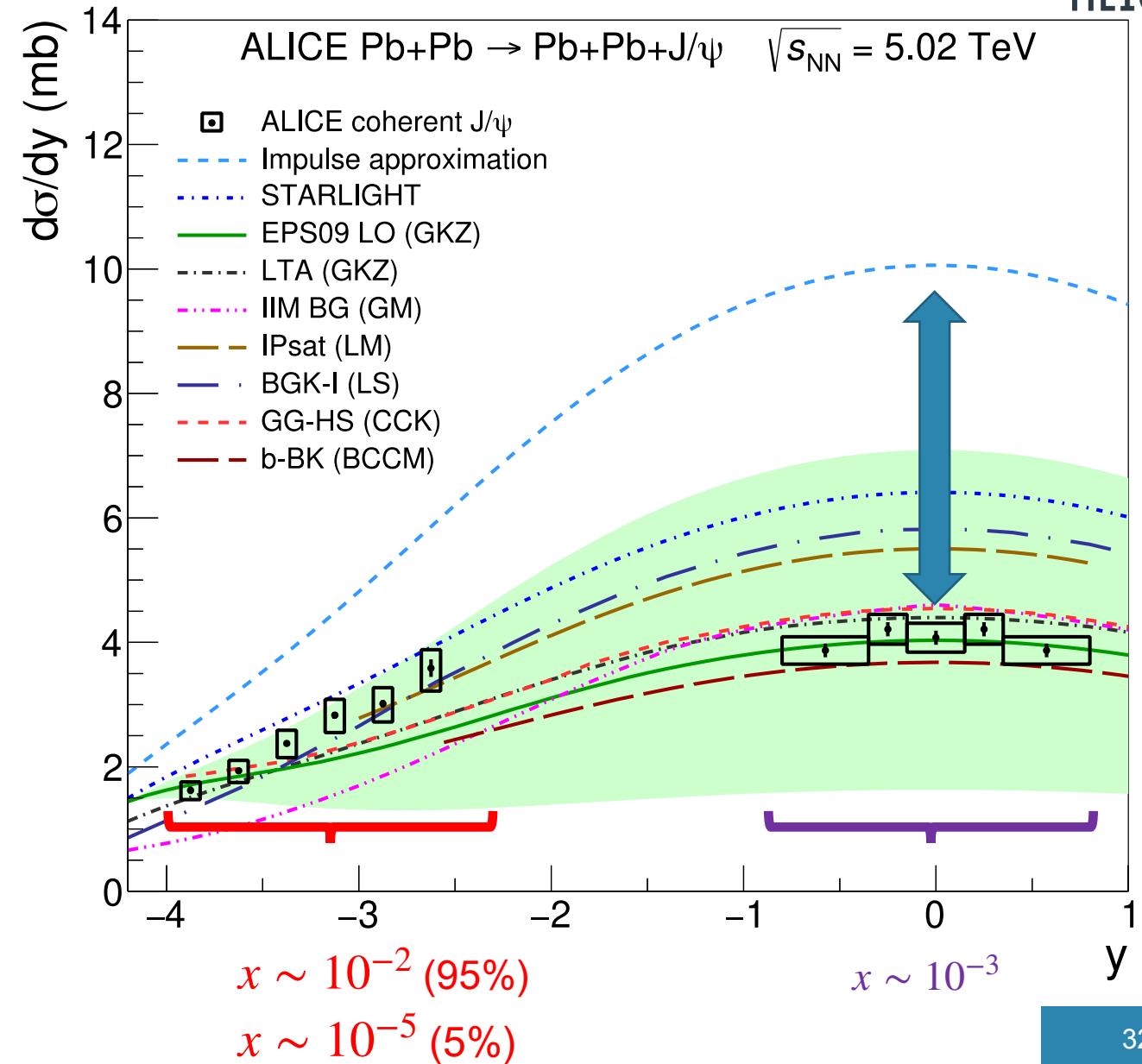
- First measurement of the nuclear suppression factor at Bjorken- $x \sim 10^{-5}$ ! Obtained with the neutron emission technique
- At high- $x$ , models fail to describe the data
- At low- $x$ , nuclear shadowing and saturation consistent with the data
- Peripheral and neutron emission techniques in good agreement



# Backup slides

# Coherent J/ $\psi$ cross section

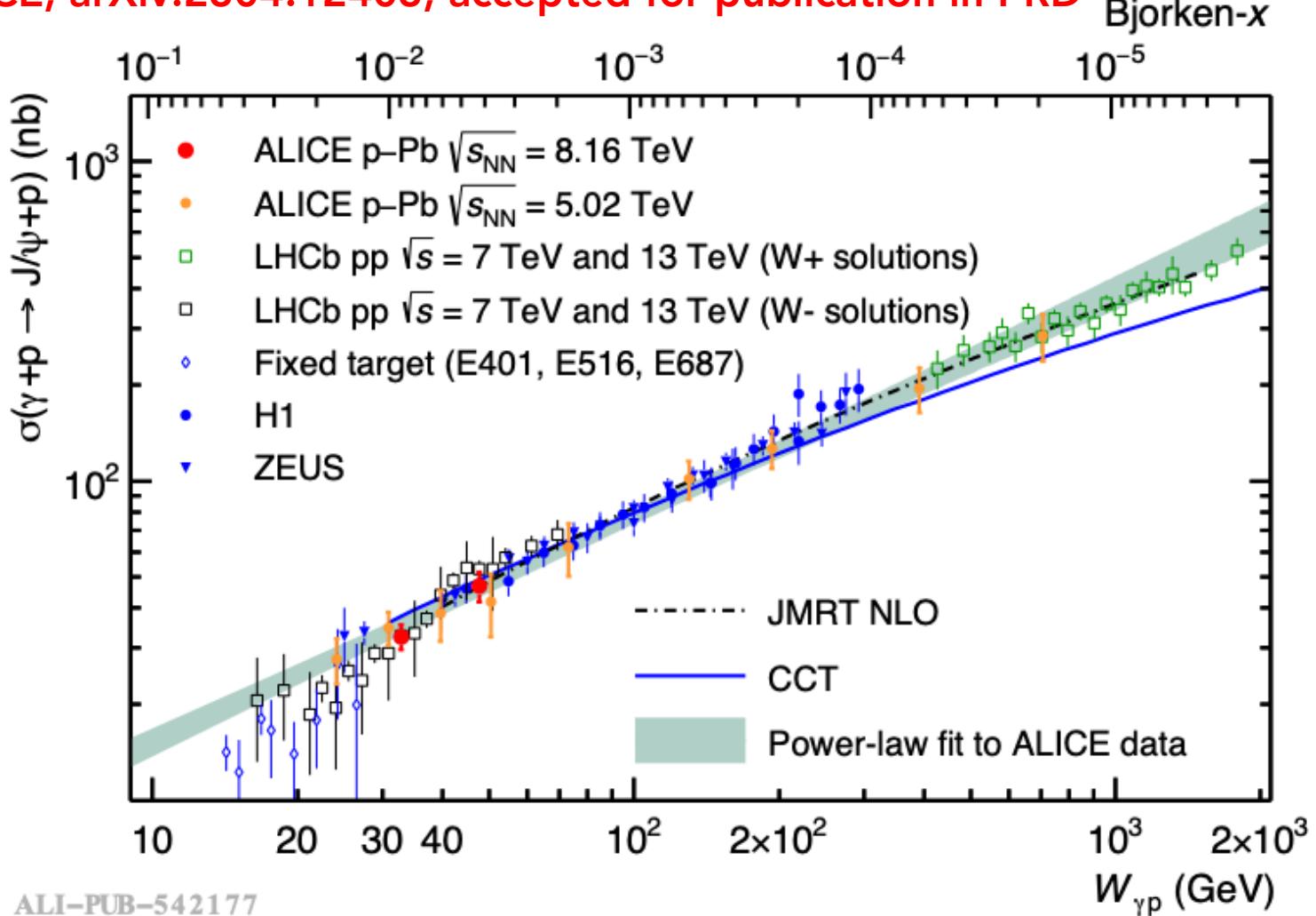
- Impulse approximation: Photoproduction data from protons, does not include nuclear effects except coherence
- STARlight: Photoproduction data from protons + Vector Meson Dominance model, includes multiple scattering but no gluon shadowing [Klein, Nystrand et al: Comput. Phys. Commun. 212 (2017) 258]
- EPS09: parametrization of nuclear shadowing data [Guzey, Kryshen, Zhalov, PRC93 (2016) 055206]
- LTA: nuclear shadowing [Guzey, Kryshen, Zhalov, PRC93 (2016) 055206]
- IIM BG, IPsat, BGK-I: Color dipol-based approaches [Goncalves, Machado et al.: PRC 90 (2014) 015203, JPG 42 (2015) 105001], [T. Lappi, H. Mäntysaari, PRC 83 (2011) 065202; 87 (2013) 032201]
- GG-HS: Color dipole + hot spots [Cepila, Contreras et al. PRC97 (2018) 024901]
- LS: Color dipole model [Luszczak, Schafer: PRC 99, 044905 (2019)]
- b-BK: Color dipole + Balitsky-Kovchegov equation



# Exclusive J/ $\psi$ in p-Pb UPC

ALICE, arXiv:2304.12403, accepted for publication in PRD

- JMRT NLO: DGLAP formalism with main NLO contributions [Jones et al. JHEP 11 (2013) 085, J. Phys. G 44 (2017) 03LT01, Phys. Rev. D 97 (2018) 116013]
- CCT: colour dipole + energy dependent hot spot model [Cepila et al., Phys. Lett. B766 (2017) 186–191]

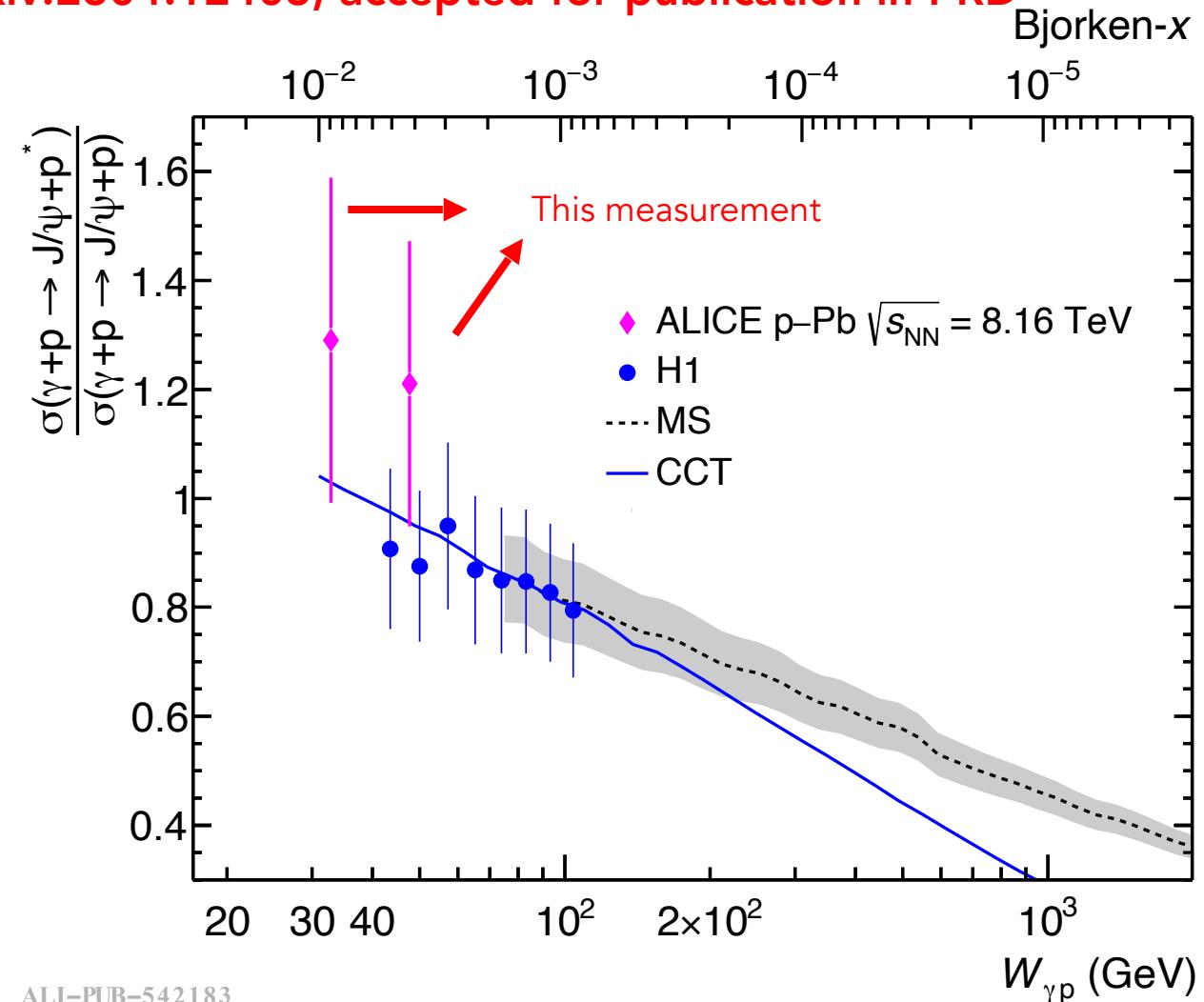


Eur. Phys. J. C (2019) 79: 402 (ALICE midrapidity and semiforward)  
Phys. Rev. Lett. 113 no. 23, (2014) 232504 (ALICE forward )

# Exclusive and dissociative J/ $\psi$ in p-Pb

ALICE, arXiv:2304.12403, accepted for publication in PRD

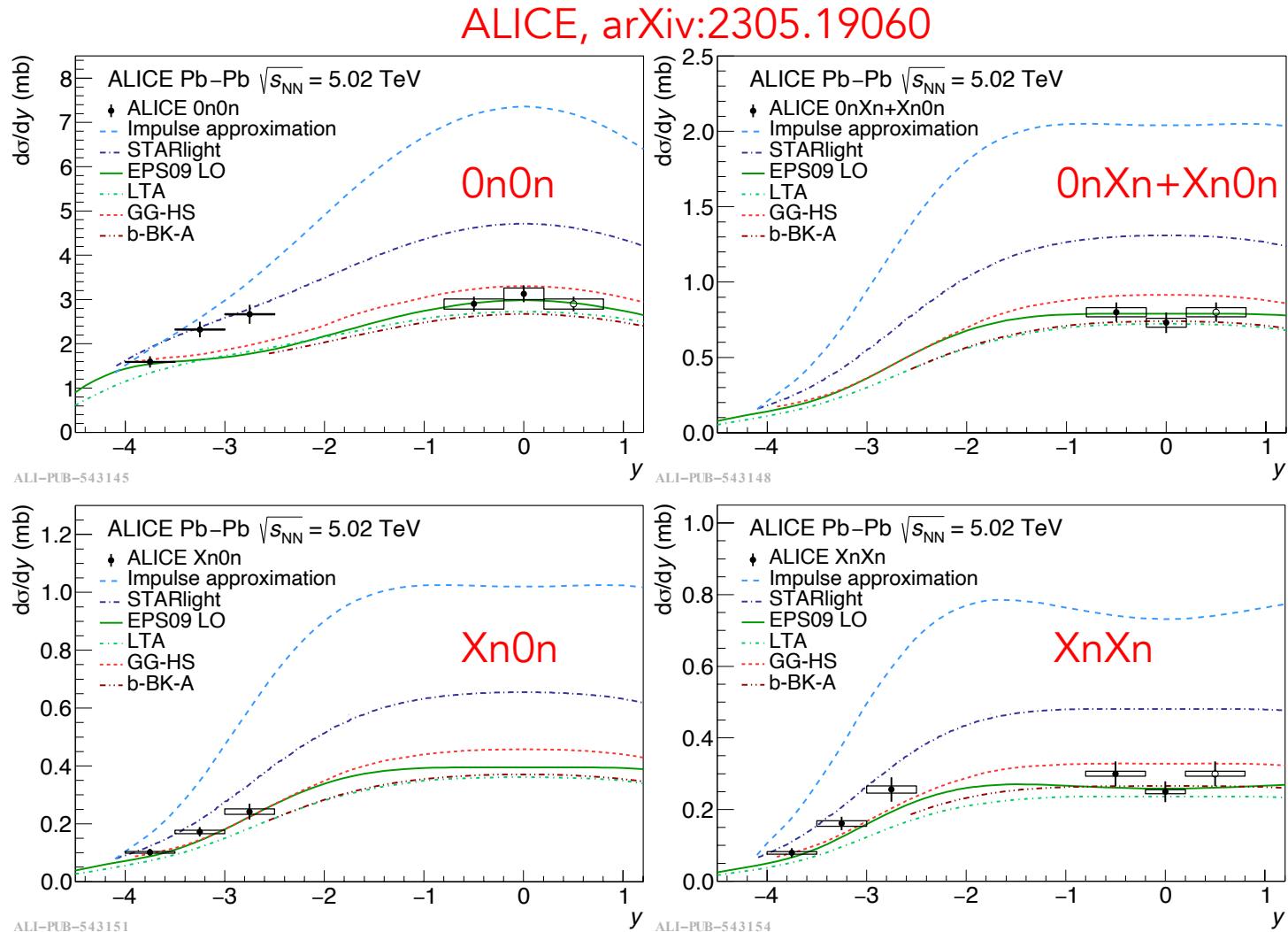
- CCT: colour dipole + energy dependent hot spot model [Cepila et al., Phys. Lett. B766 (2017) 186–191]
- MS: JIMWLK + constrains from fits to H1 data [Mäntysaari, Schenke, Phys. Rev. D 98 (2018) 034013]



ALI-PUB-542183

# Coherent J/ $\psi$ with neutron emission

- Event migrations: missed detection of a neutron (i.e. XN0N->0N0N), and neutrons from pile-up events
- $\epsilon_A = \epsilon_C = 0.933 \pm 0.017$
- $p_A = 0.0237 \pm 0.0005$
- Events are collected for pile-up, with  $\epsilon_{pu} = 0.920 \pm 0.002$  and  $0.962 \pm 0.001$  for mid and forward rapidity respectively
- Further correction for charged particles produced at beam rapidities by the dissociation of either nuclei,  $\epsilon_{emd}$



Systematics relevant to the extraction  
of the UPC cross section, i.e.  $d\sigma/dy$

# Coherent J/ $\psi$ with neutron emission

- Photon flux used as weights in a  $\chi^2$  minimisation:

$$\chi_{\text{exp}}^2(\vec{m}, \vec{b}) = \sum_i \frac{\left(m^i - \sum_j \gamma_j^i m^i b_j - \mu^i\right)^2}{\delta_{i,\text{stat}}^2 \mu^i \left(m^i - \sum_j \gamma_j^i m^i b_j + (\delta_{i,\text{uncor}} m^i)^2\right)} + \sum_j b_j^2.$$

- Two types of uncertainty on the photon flux:
  - Total flux, obtained by varying the nuclear radius according to nuclear skin measurements
  - Fraction of the flux to each neutron class

$$\frac{d\sigma_{PbPb}^{0N0N}}{dy} = n_{0N0N}(\gamma, +y) \cdot \sigma_{\gamma Pb}(+y) + n_{0N0N}(\gamma, -y) \cdot \sigma_{\gamma Pb}(-y)$$

$$\frac{d\sigma_{PbPb}^{0NXN}}{dy} = n_{0NXN}(\gamma, +y) \cdot \sigma_{\gamma Pb}(+y) + n_{0NXN}(\gamma, -y) \cdot \sigma_{\gamma Pb}(-y)$$

Large difference in weight of the two solutions for the photon fluxes!



$y$	$n_\gamma(0n0n)$	$n_\gamma(0nXn+Xn0n)$	$n_\gamma(XnXn)$	$\sigma_{\gamma Pb}^{\text{IA}} (\mu\text{b})$
$3.5 < y < 4$	178.51	18.18	6.34	10
$3 < y < 3.5$	162.99	18.19	6.34	14
$2.5 < y < 3$	147.46	18.19	6.34	19
$0.2 < y < 0.8$	77.88	17.88	6.33	48
$-0.2 < y < 0.2$	62.86	17.47	6.27	58
$-0.8 < y < -0.2$	48.31	16.75	6.18	71
$-3 < y < -2.5$	3.91	4.97	2.78	176
$-3.5 < y < -3$	1.22	2.15	1.42	215
$-4 < y < -3.5$	0.26	0.61	0.48	262

# Coherent J/ $\psi$ with neutron emission

- Photon flux used as weights in a  $\chi^2$  minimisation:

$$\chi_{\text{exp}}^2(\vec{m}, \vec{b}) = \sum_i \frac{\left(m^i - \sum_j \gamma_j^i m^i b_j - \mu^i\right)^2}{\delta_{i,\text{stat}}^2 \mu^i \left(m^i - \sum_j \gamma_j^i m^i b_j + (\delta_{i,\text{uncor}} m^i)^2\right)} + \sum_j b_j^2.$$

- Two types of uncertainty on the photon flux:
  - Total flux, obtained by varying the nuclear radius according to nuclear skin measurements
  - Fraction of the flux to each neutron class

$$\frac{d\sigma_{PbPb}^{0N0N}}{dy} = n_{0N0N}(\gamma, +y) \cdot \sigma_{\gamma Pb}(+y) + n_{0N0N}(\gamma, -y) \cdot \sigma_{\gamma Pb}(-y)$$

$$\frac{d\sigma_{PbPb}^{0NXN}}{dy} = n_{0NXN}(\gamma, +y) \cdot \sigma_{\gamma Pb}(+y) + n_{0NXN}(\gamma, -y) \cdot \sigma_{\gamma Pb}(-y)$$

Large difference in weight of the two solutions for the photon fluxes!

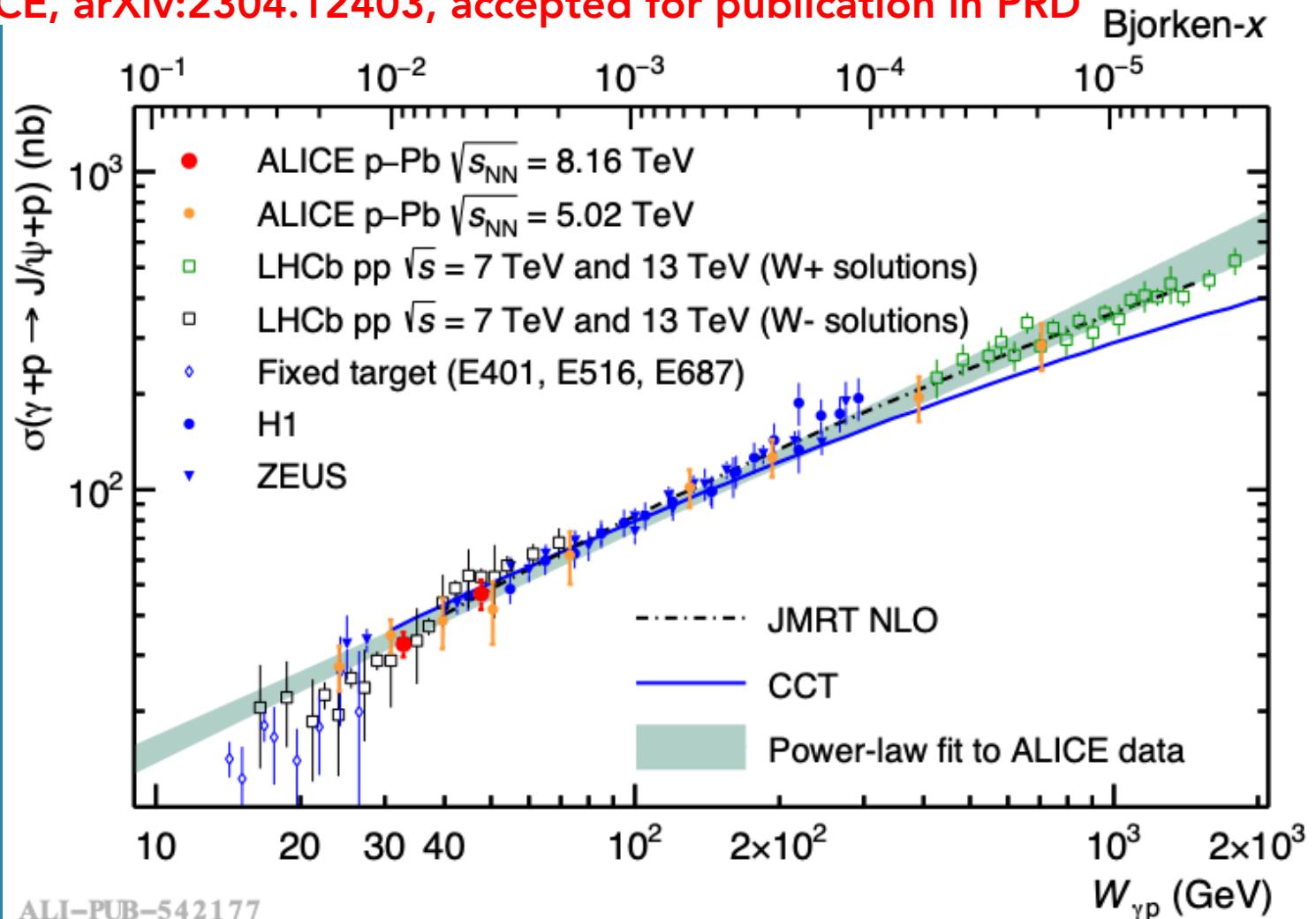


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# Exclusive J/ $\psi$ in p-Pb UPC

ALICE, arXiv:2304.12403, accepted for publication in PRD

- $x = e^{\pm|y|} M_{J/\psi} / 2E_p$
- Probing Bjorken- $x \sim 10^{-5}$  with ALICE data
- power-law growth of cross sections observed -> no change in the behaviour of the gluon PDF in the proton between HERA and LHC energies
- ALICE points: forward, semiforward and midrapidity configurations
  - Forward: two muons in the spectrometer
  - Semiforward: one in the spectrometer, one in the central barrel
  - Midrapidity: two muons/ electrons in the central barrel

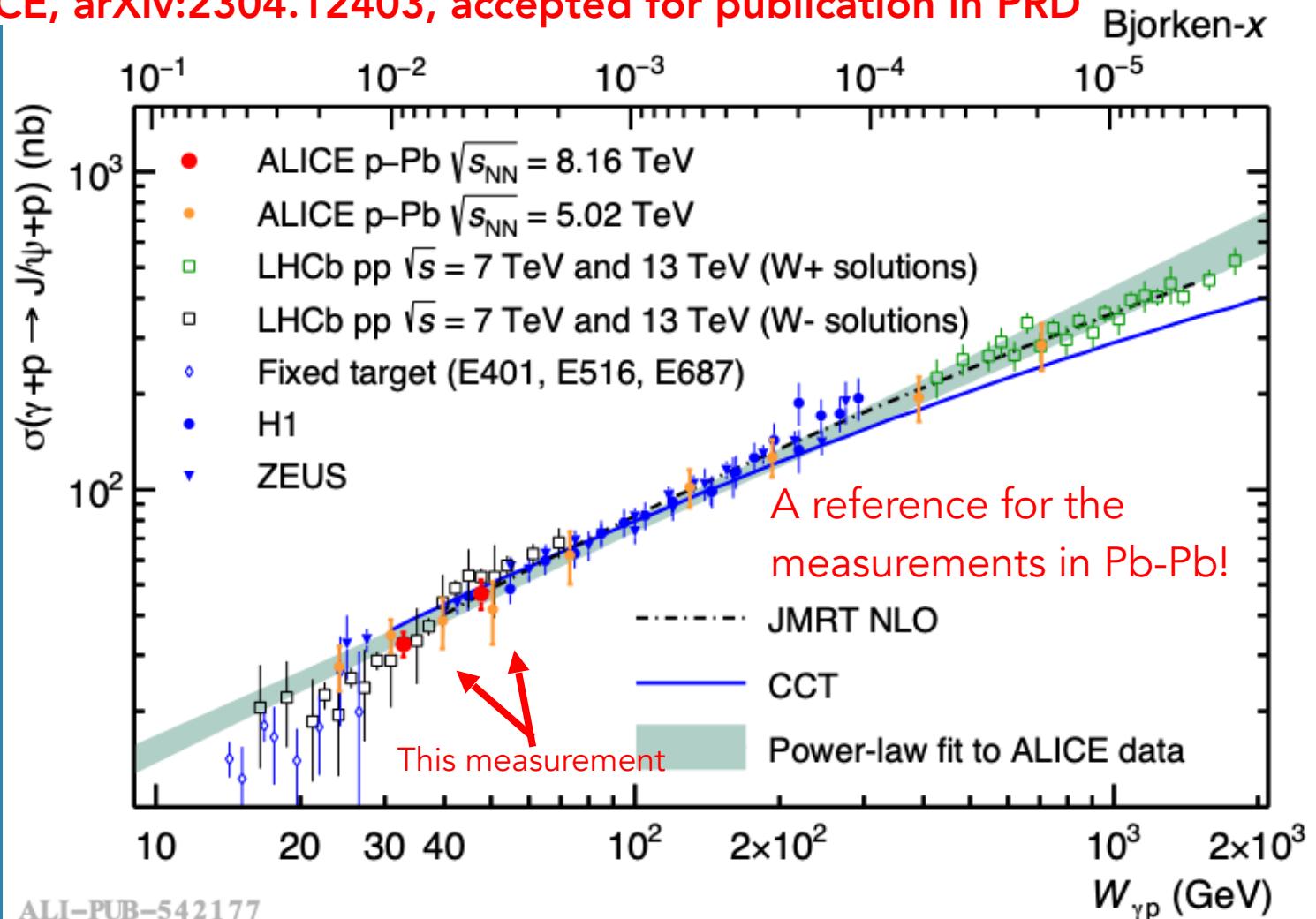


Eur. Phys. J. C (2019) 79: 402 (ALICE midrapidity and semiforward)  
Phys. Rev. Lett. 113 no. 23, (2014) 232504 (ALICE forward )

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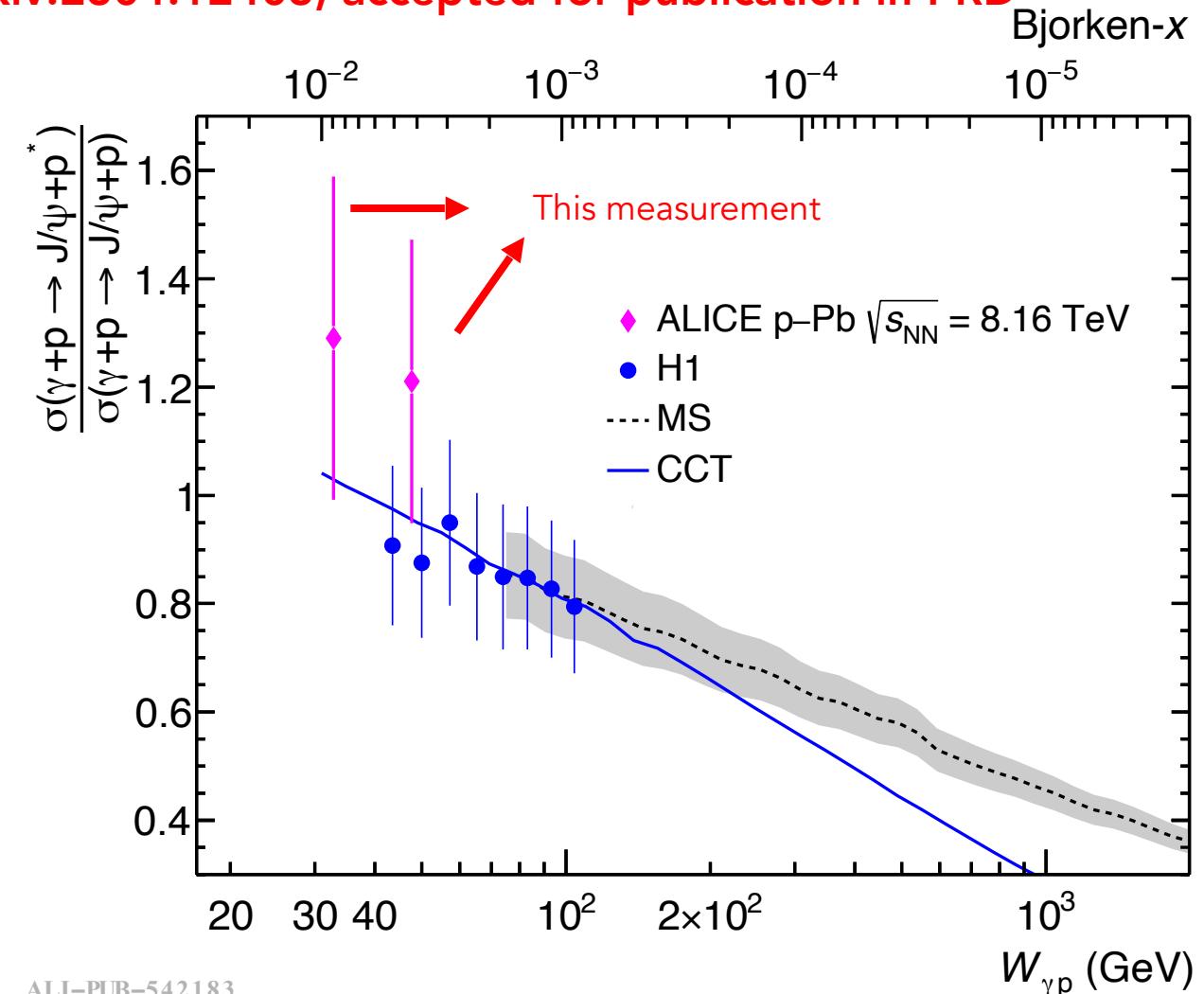
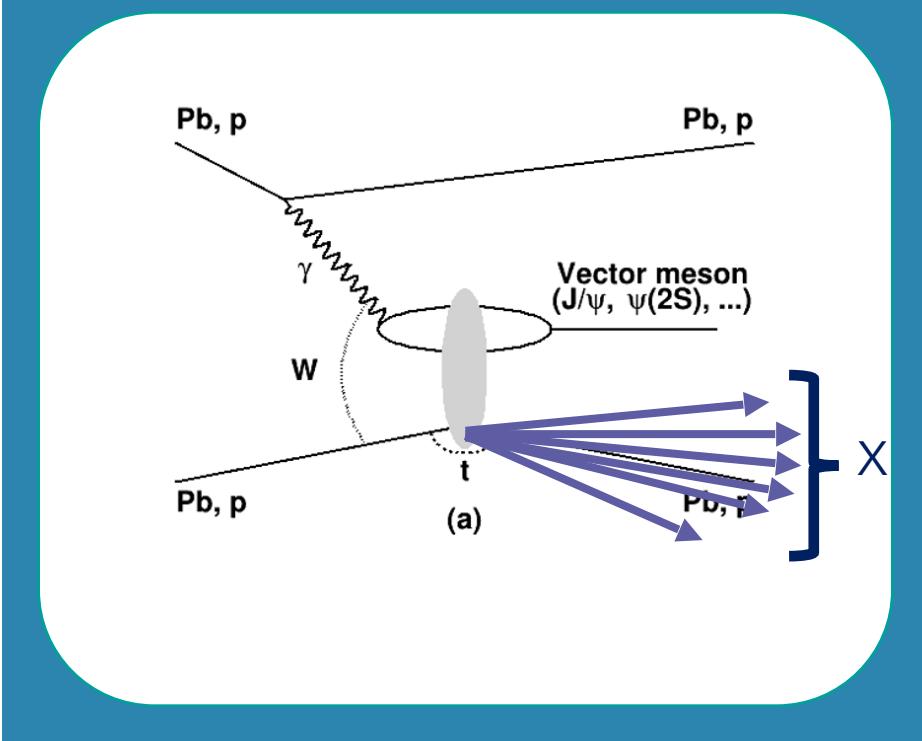


Eur. Phys. J. C (2019) 79: 402 (ALICE midrapidity and semiforward)  
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# Exclusive and dissociative J/ $\psi$ in p-Pb

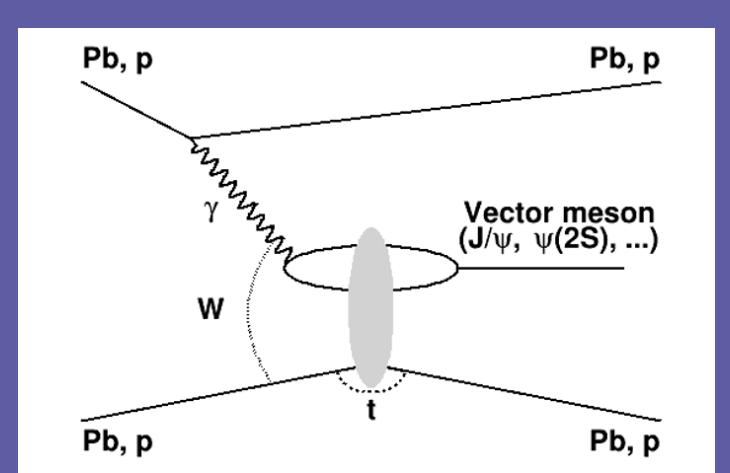
ALICE, arXiv:2304.12403, accepted for publication in PRD

- First result at the LHC of the measurement of dissociative J/ $\psi$

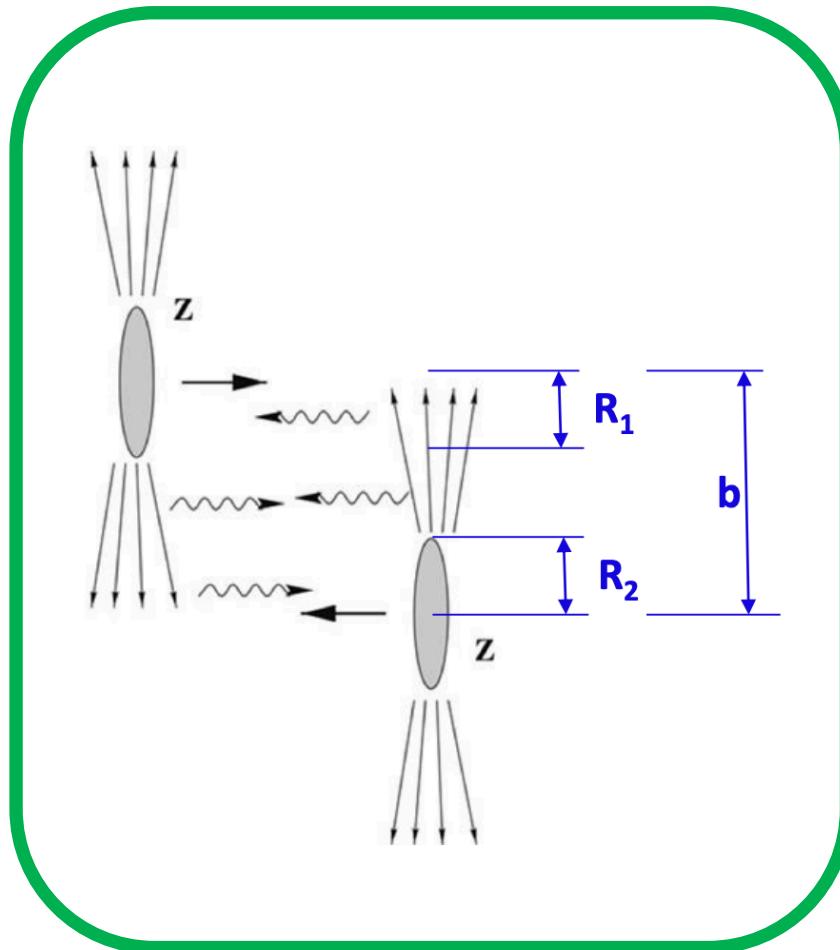


ALI-PUB-542183

# Physics of ultra-peripheral collisions (UPC)

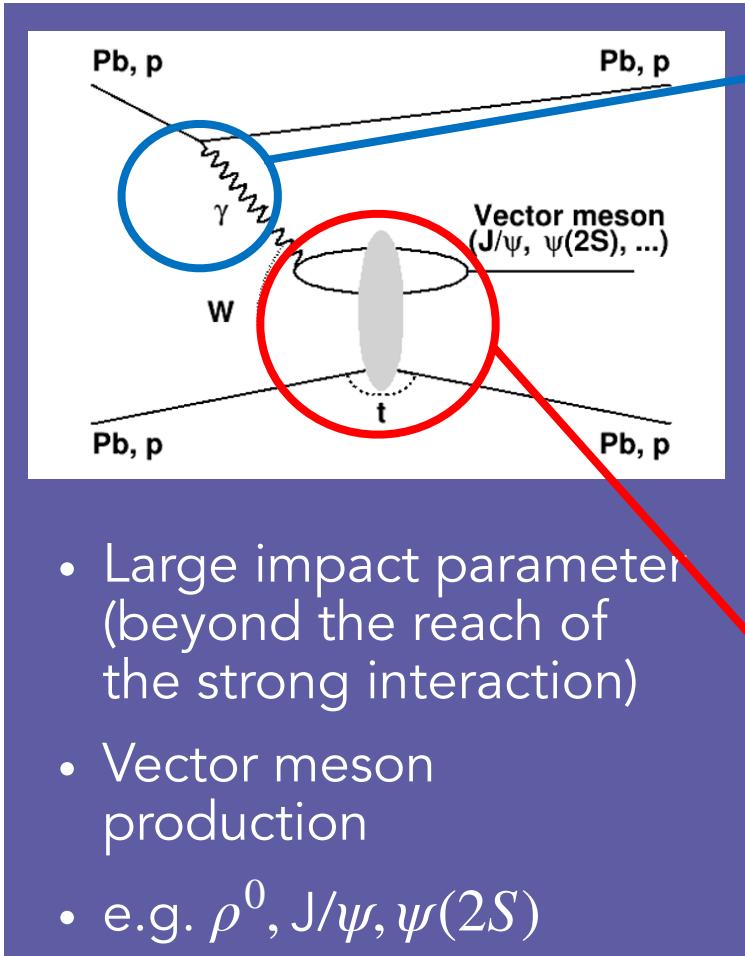


- Large impact parameter (beyond the reach of the strong interaction)
- Vector meson production
- e.g.  $\rho^0$ , J/ $\psi$ ,  $\psi(2S)$



- High photon flux
- Described in Weizsäcker-Williams approach of quasireal photons
- Flux proportional to  $Z^2$

# Introduction to ultra-peripheral collisions (UPC)



Only QED involved at this vertex!

$|t|$  the square of the momentum transferred between the incoming and outgoing target nucleus

$W$  the centre-of-mass energy of the photon-target system

Hard scale ensured by high mass states i.e.  $J/\psi, \psi(2S)$

Semi-hard scale for  $\rho^0$

- *Coherent photoproduction:* photon couples with the **entire nucleus**
- *Incoherent photoproduction:* photon couples with a **single nucleon** only
- Different average  $p_T$  of the vector mesons for the two processes