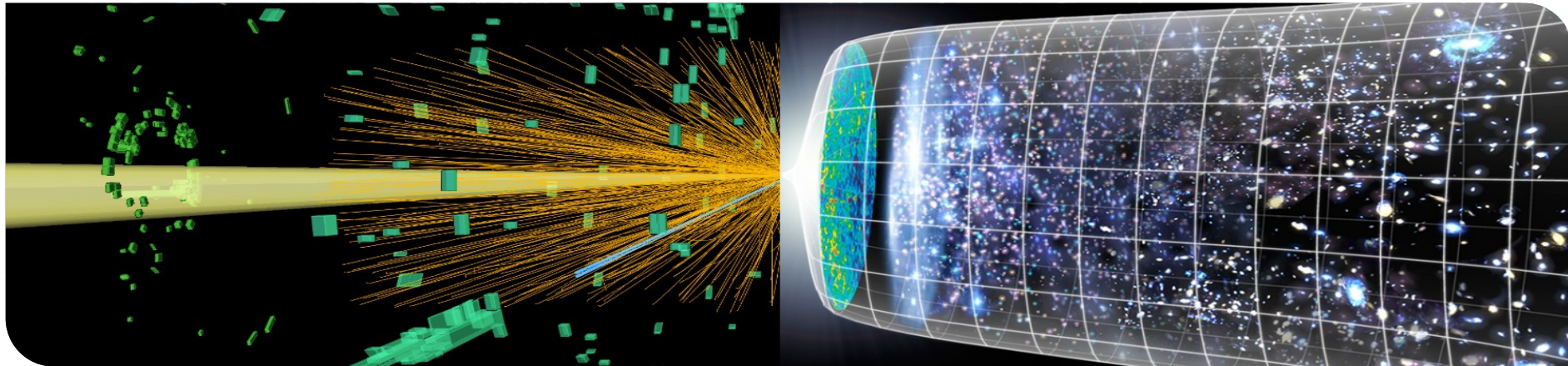


Strongly-interacting dark sectors and dark matter relic density: an overview

Felix Kahlhoefer

Roadmap of Dark Matter Models for Run 3

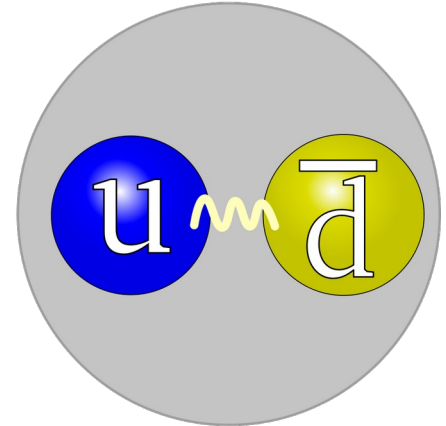


Strongly-interacting dark sectors

- Focus of this talk: QCD-like dark sectors (mostly SU(N) theories with $N_f < 3 N_c$)
- At high energies: dark sector **contains dark gluons and dark quarks:**

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu}^a F^{\mu\nu a} + \bar{q}_d i \not{D} q_d - \bar{q}_d M_q q_d$$

- Quark masses small or comparable to confinement scale Λ_d
- At low energies: Confinement into dark mesons and baryons
- In case of particle-antiparticle asymmetry, dark baryons could be DM, otherwise they annihilate away



Cline & Perron, arXiv:2204.00033

Stable dark pions

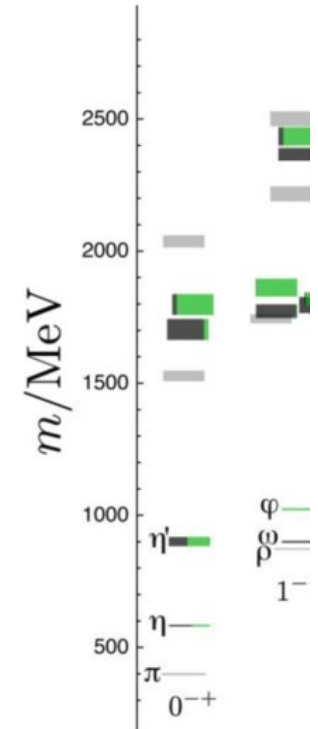
- In principle, all dark mesons could be unstable (as in QCD) → no DM candidate
- More appealing: Stabilisation mechanism
 - Option 1: Abelian $U(1)'$ gauge symmetry
 - Different charges for different dark quarks (e.g. +1 for up-type and -1 for down-type quarks)
 - Dark mesons of type $u\bar{d}$ or $d\bar{u}$ carry $U(1)'$ charge
 - Lightest such mesons (dark pions) must be stable
 - Option 2: Global discrete symmetry (G-parity)
 - Also neutral dark pions can be stabilised

→

u	\bar{u}	d	\bar{d}	π^0
$-\bar{d}$	$-d$	\bar{u}	u	$-\pi^0$

Describing dark mesons

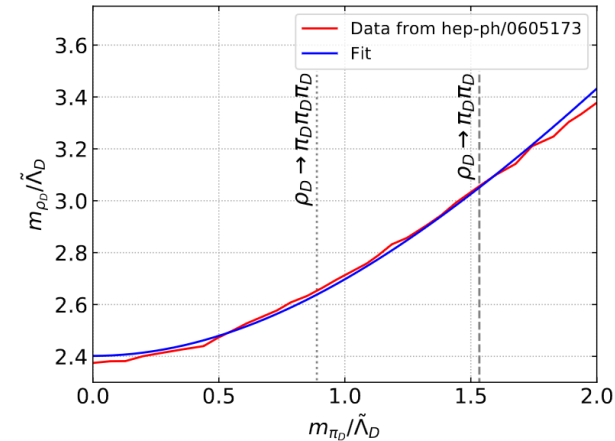
- Low-energy theory features many different states
 - Dark pions (Pseudo-Goldstone bosons of chiral symmetry breaking)
 - Dark rho mesons (spin-1)
 - ...
- Apparently many free parameters:
 - Masses of various dark mesons (dark pions, dark rho mesons, ...)
 - Interactions between them
 - Interactions with SM particles
- Great complexity! Need some guidance...



Crucial connections

- But high energy theory has only 3 parameters (for given number of colours and flavours):
 - Dark quark mass (assuming degenerate states)
 - Dark gauge coupling
 - Coupling to mediator (e.g. U(1)' gauge boson)
- Low-energy pheno determined by one dimensionful parameter (e.g. m_π) and one dimensionless ratio (e.g. m_π/Λ_d)
 - Dark rho meson mass determined from lattice simulations
 - Coupling between ρ_d and π_d determined by KSRF relation
 - Couplings to mediator follow from group theory

Dark showers Snowmass report,
arXiv:2203.09503



What constrains the dark pion mass?

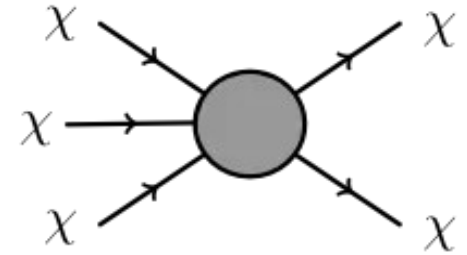
- Need input from astrophysics and cosmology
- Bullet Cluster excludes DM self-interaction cross sections larger than $\sim 2 \text{ cm}^2/\text{g}$ ($= 4 \text{ barn}/\text{GeV}$)
- Dark pions have $\sigma_c = \frac{3}{64\pi} \frac{m_\pi^2}{f_\pi^4} (1 + \mathcal{O}(N_{f_D}^{-2}))$
- Consider sizable dark quark masses, such that m_π and f_π are not too different
→ Dark pion masses below $\sim 100 \text{ MeV}$ excluded



Dark pion relic density calculation

- Conceptually simplest way for dark pions to evolve in the early universe:
Number-changing processes ($3\pi_d \rightarrow 2\pi_d$) induced by the Wess-Zumino-Witten term

$$\frac{2N_d}{15\pi^2 f_\pi^5} \epsilon^{\mu\nu\rho\sigma} \text{Tr} (\pi \partial_\mu \pi \partial_\nu \pi \partial_\rho \pi \partial_\sigma \pi)$$



- Rest mass converted into kinetic energy
- Kinetic energy transferred to SM via scattering
- Relic density determined by freeze-out of number-changing processes

→ **SIMP mechanism**

Hochberg et al., arXiv:1402.5143

Miracle or no miracle?

- Amazingly, the well-known literature result for the SIMP cross section is wrong by a factor of 3

- Reason: Naive non-relativistic limit leads to an expression that is not Lorentz-invariant

Kamada et al., arXiv:2210.0139;

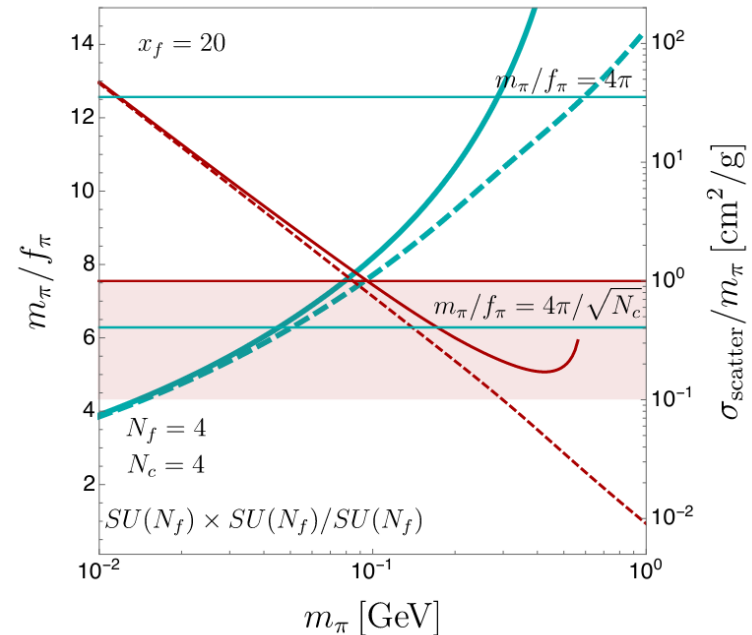
- (N)NLO corrections can be large → lattice needed!

Hansen et al., arXiv:1507.01590; Dengler et al., arXiv:2405.06506

- General conclusion: SIMP mechanism requires dark pion masses of 10–100 MeV

- Tension with Bullet Cluster bound
- Requires $m_\pi / f_\pi \sim 4\pi$ (close to perturbative bound)

Cyan: m_π/f_π required by relic density
Red: Corresponding cross section

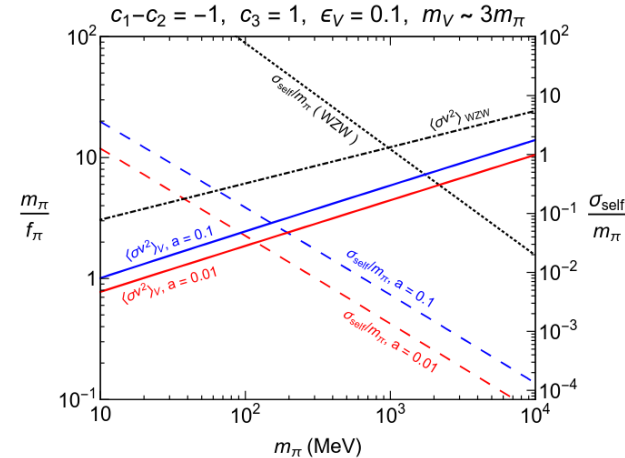
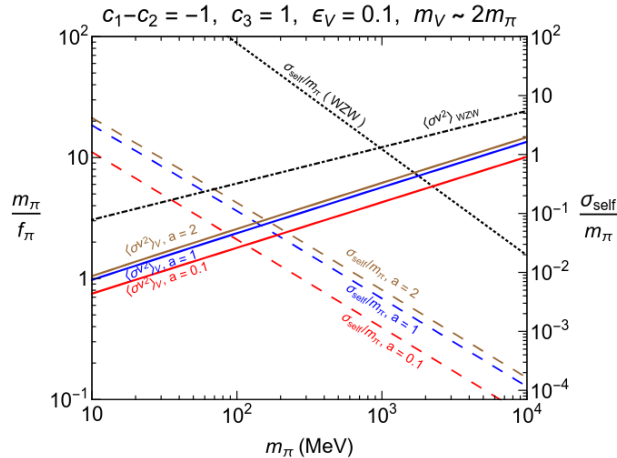
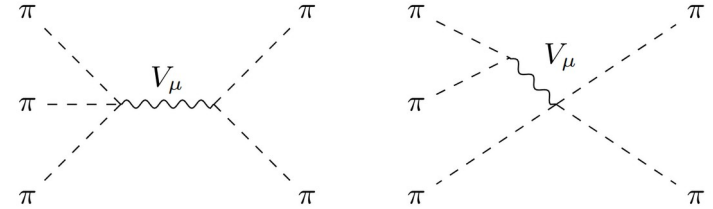


Possible ways forward

- Assume resonant enhancement of $3 \rightarrow 2$ cross section through exchange of on-shell dark meson

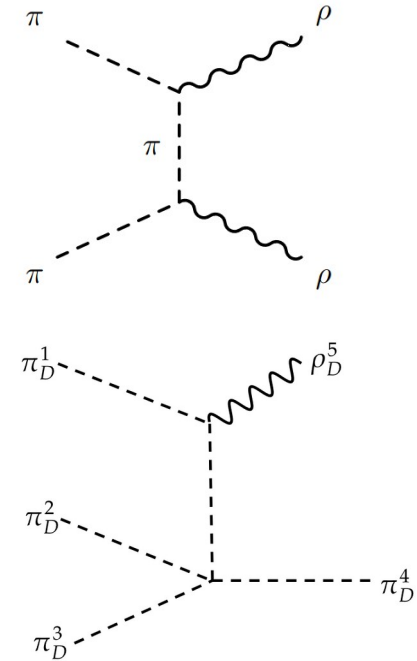
Choi et al., arXiv:1801.07726

- Requires tuning but relaxes Bullet Cluster constraint



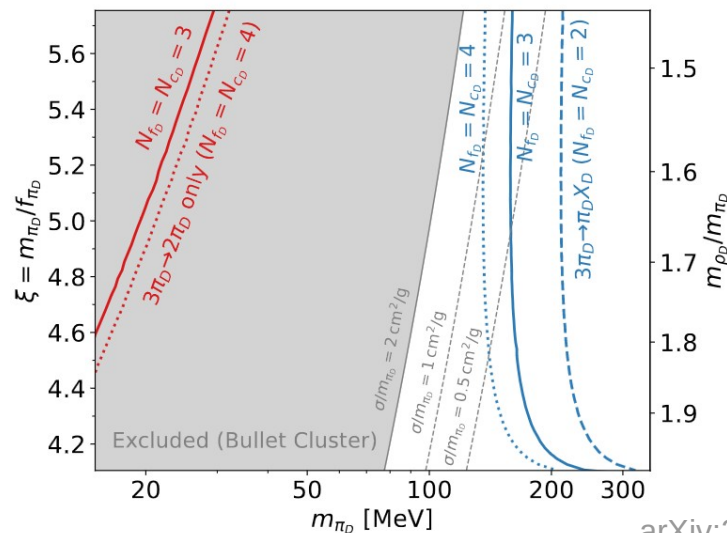
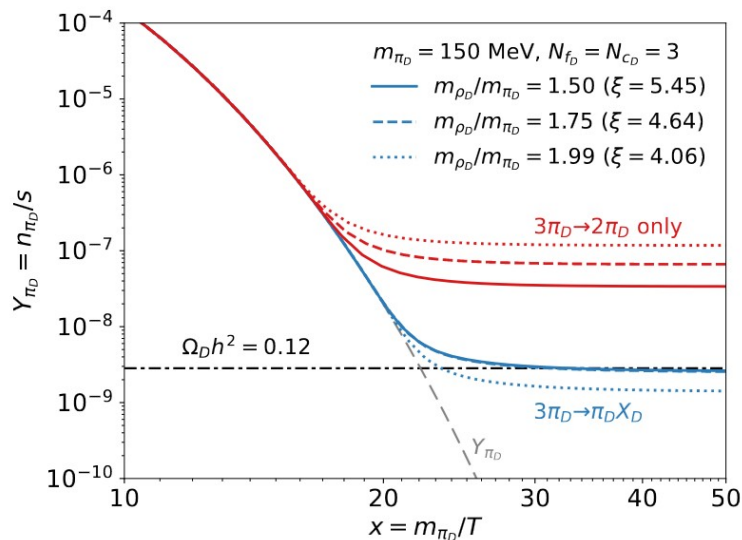
Dark rho mesons in the final state

- $2\pi_d \rightarrow \pi_d + \rho_d$ Berlin et al., arXiv:1801.05805
 - Requires anomalous interaction and violates G-parity
- $2\pi_d \rightarrow 2\rho_d$ Bernreuther, FK et al., arXiv:1907.04346
 - Requires non-zero kinetic energy in the initial state
 - Annihilation rate exponentially suppressed for small temperatures
 - Requires very small $\pi_d - \rho_d$ mass difference (problem with chPT validity)
- $3\pi_d \rightarrow \pi_d + \rho_d$ Bernreuther, FK et al., arXiv:2311.17157
 - Kinematically allowed in the non-relativistic limit for $m_\rho < 2 m_\pi$
 - No momentum suppression (unlike SIMP process)
 - Resonant enhancement for m_ρ close to $2 m_\pi$



Results

- Including the process $3\pi_d \rightarrow \pi_d + \rho_d$ enhances annihilation rate by orders of magnitude
- Freeze-out delayed, larger masses viable, Bullet Cluster constraint can be evaded



arXiv:2311.17157

Phenomenological implications

- For $m_\rho < 2 m_\pi$, dark rho mesons cannot decay into pairs of dark pions
- But ρ^0 mesons must couple to SM particles for thermalisation

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- For $m_\rho < 2 m_\pi$, dark rho mesons cannot decay into pairs of dark pions
- But ρ^0 mesons must couple to SM particles for thermalisation
- Simple example: Assume U(1)' gauge boson Z' with kinetic mixing κ
 - At energies below Z' mass: Effective interaction between dark quarks and SM fermions

$$\mathcal{L}_{\text{eff}} \supset \frac{1}{\Lambda^2} \sum_f q_f \bar{f} \gamma^\mu f \bar{q}_d \gamma_\mu q_d \quad \Lambda = \frac{m_{Z'}}{\sqrt{\kappa e e_d}}$$

- Dark ρ^0 mesons have the same quantum numbers as the Z' gauge boson

- Mixing leads to effective coupling $\mathcal{L}_{\text{eff}} \supset \frac{2}{g} \frac{m_{\rho_d}^2}{\Lambda^2} \rho_d^{0\mu} \sum_f q_f \bar{f} \gamma_\mu f$

Bernreuther, FK et al., arXiv:2203.08824

Phenomenological implications

- For $m_\rho < 2 m_\pi$, dark rho mesons cannot decay into pairs of dark pions
- But ρ^0 mesons must couple to SM particles for thermalisation
 - Prediction: ρ^0 mesons must decay into SM states
 - Semi-visible jets!
- Couplings of ρ^0 to SM particles can be very small
 - Displaced vertices or emerging jets
 - No guarantee for sizable production cross section
- Preferred mass range for dark mesons (well) below 1 GeV
 - Dark shower substructure quite similar to QCD except for fraction of invisible particles

Similar mechanisms with other light states

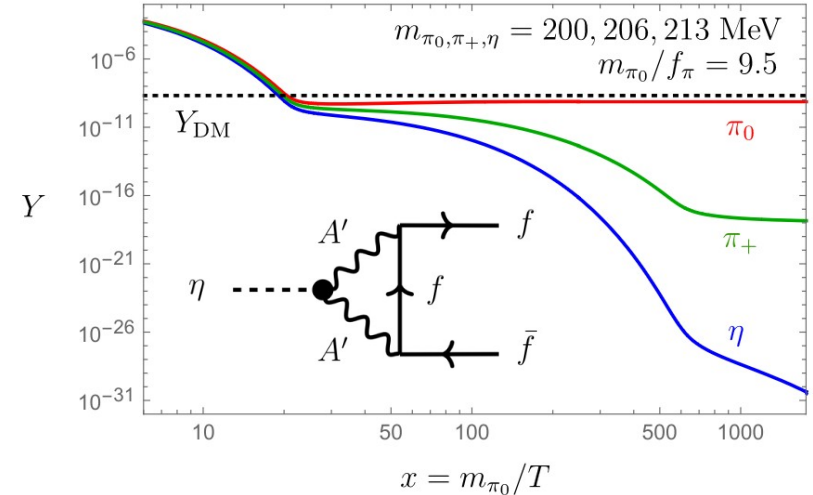
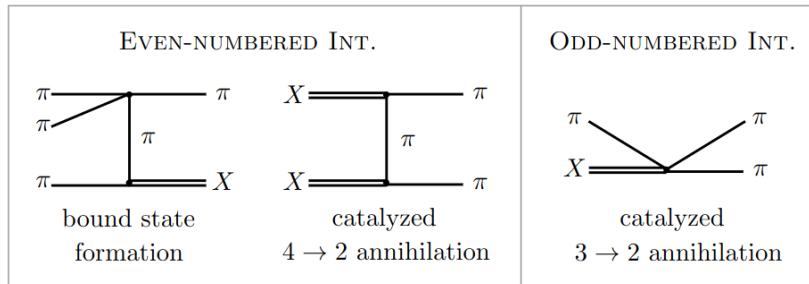
Dark eta mesons

- Consider quark flavours with different masses
- Dark eta meson heavier than dark pions
- Generally unstable against decays into SM

Katz et al., arXiv:2006.15148

$X = [\pi\pi]$ bound state (at least in $Sp(4)$)

Chu et al., arXiv:2401.12283

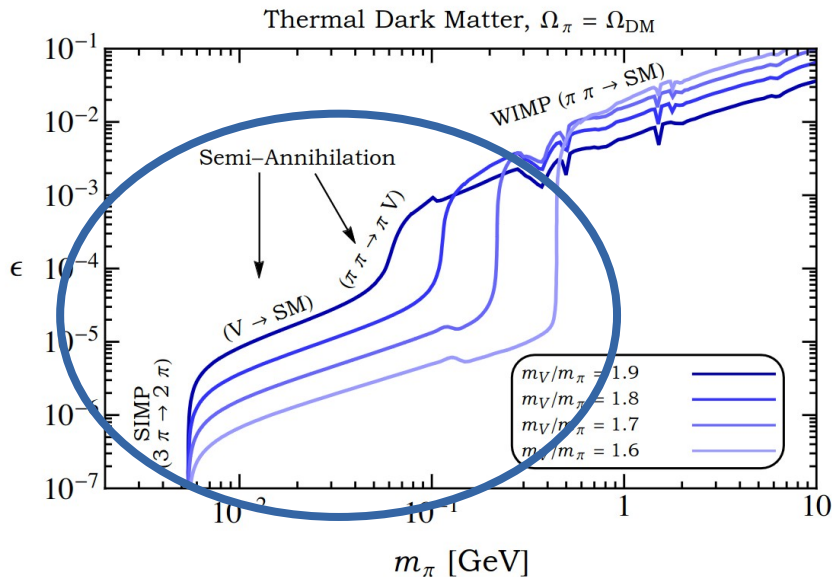


- Dilaton χ arising from breaking of conformal symmetry (not QCD-like any more!)

Appelquist et al., arXiv:2404.07601

SIMPs as WIMPs?

- All scenarios considered so far: relic density requirement places tight constraints on the dark pion mass, but only very weak constraints on the SM couplings

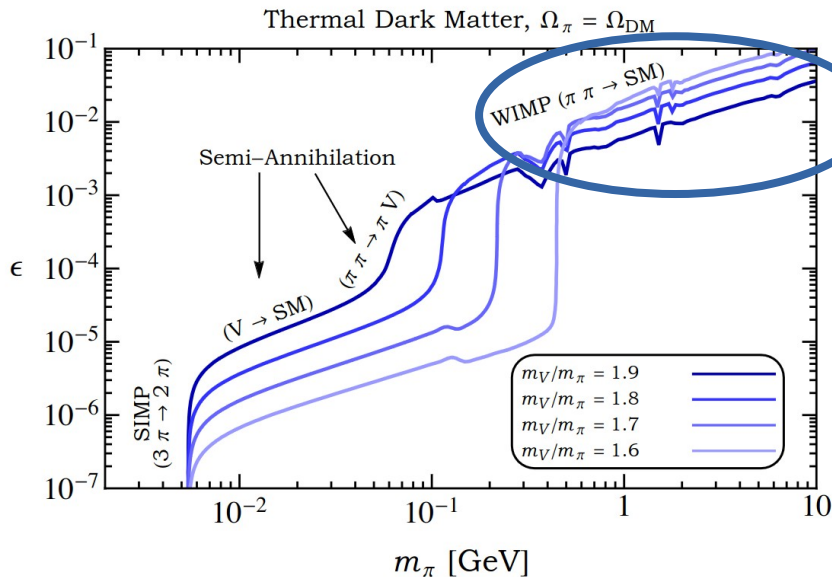


- Couplings vary by many orders of magnitude, while masses vary only slightly
- For the LHC we would prefer the opposite!

Berlin et al., arXiv:1801.05805

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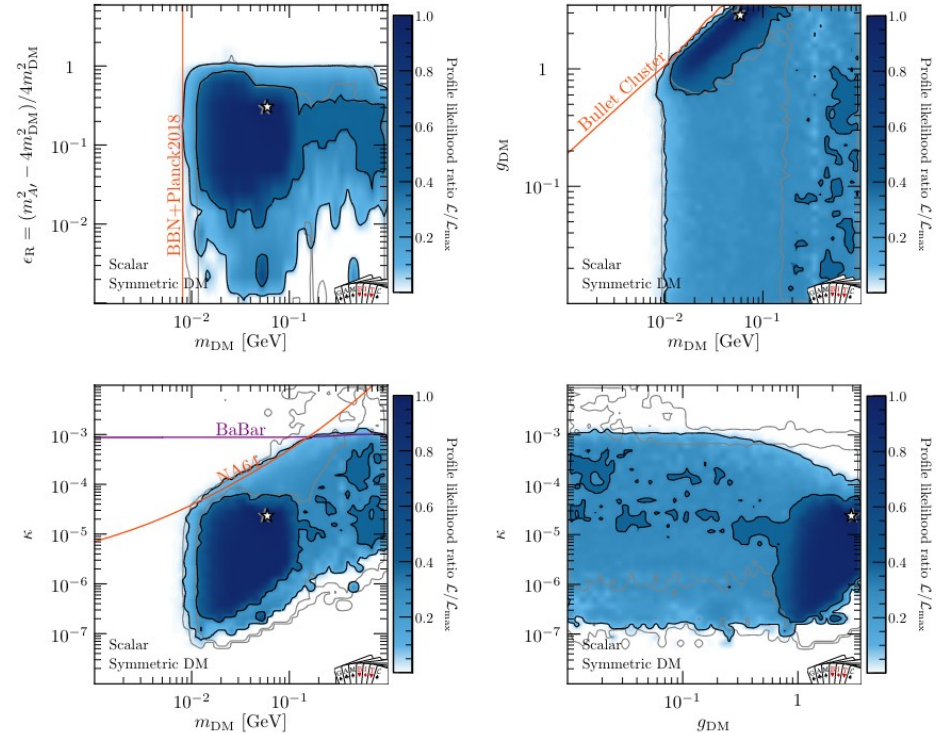


- Couplings vary by many orders of magnitude, while masses vary only slightly
- For the LHC we would prefer the opposite!
- Consider direct annihilations of dark pions into SM particles

Berlin et al., arXiv:1801.05805

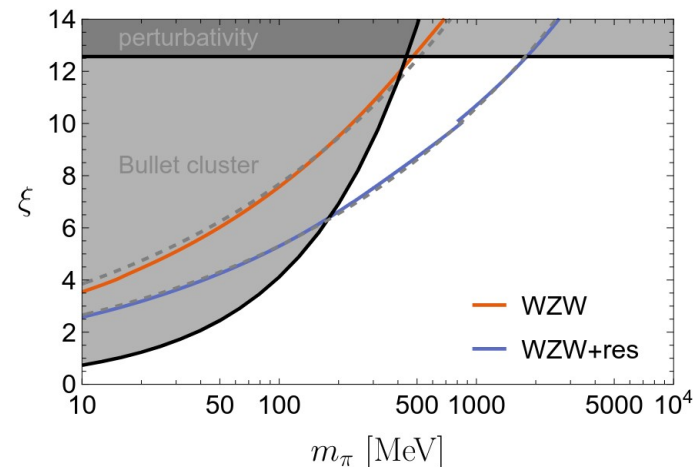
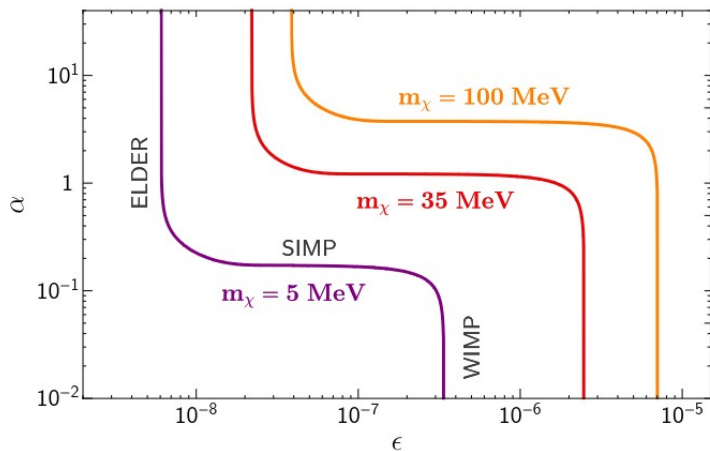
Dark pions as GeV-scale thermal DM

- More flexibility in terms of mass, clearer target in terms of couplings
- Many relevant constraints from missing energy searches, rare meson decays, direct detection etc.
- Viable parameter space for complex scalar DM identified in upcoming global fit
Balan, Gray, FK et al., in preparation
- But: Not all constraints apply to dark pions
- So far no dedicated study (?)



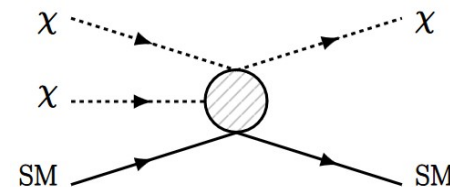
Combination of different interactions

- Combine $3 \rightarrow 2$ process and (resonantly enhanced) annihilation into SM Braat & Postma, arXiv:2301.04513
- Relic density set by thermal decoupling of elastic scattering (ELDERS) Kuflik et al., arXiv:1512.04545



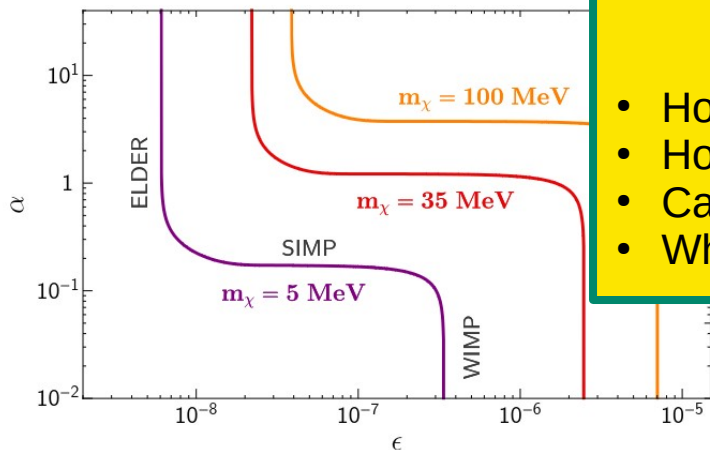
- Consider number-changing processes involving SM particles

Smirnov & Beacom, arXiv:2002.04038



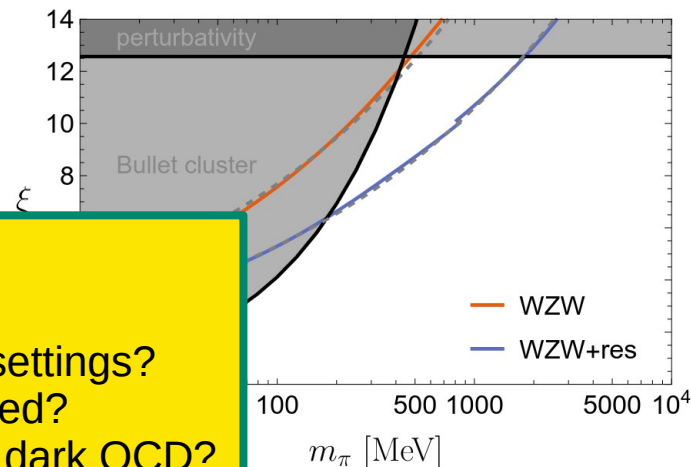
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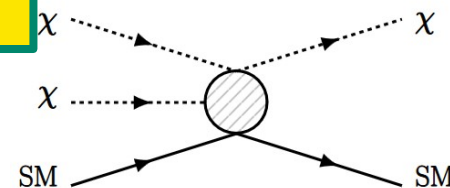
Open questions

- How predictive are these settings?
- How much tuning is required?
- Can they be embedded in dark QCD?
- What is the Bullet Cluster constraint?



Consider number-changing processes involving SM particles

Smirnov & Beacom, arXiv:2002.04038



Conclusions

- QCD-like strongly interacting DM has exciting implications for cosmology and colliders
- Conventional SIMP mechanism in tension with Bullet Cluster constraints
- Many possible modifications using additional states in the dark sector or new interactions
- Simple solution: Annihilations into dark rho mesons with $m_\rho < 2m_\pi$
 - Dark rho mesons must decay visibly into SM particles
- Interactions between dark and visible sector largely unconstrained
- Interesting to think about alternative relic density mechanisms with tighter connection
 - WIMP-like annihilations? Resonant annihilations? Co-SIMP mechanism?