

*Unraveling **Dark Matter** with novel collider signatures in Type-I 2HDM+a*

Roadmap of Dark Matter models for Run 3

Ilia Kalaitzidou

Spyros Argyropoulos

Ulrich Haisch

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Motivation

Open question: Dark Matter nature

- **Initial simplified DM models:** Add a singlet acting as mediator between the visible and dark sectors [arXiv:1506.03116](https://arxiv.org/abs/1506.03116)
- **Problem:** Unitarity is be violated: Interactions between DM mediator and SM fermions are not gauge invariant [arXiv:1510.02110](https://arxiv.org/abs/1510.02110)
- **Solution:** Extend the SM Higgs sector



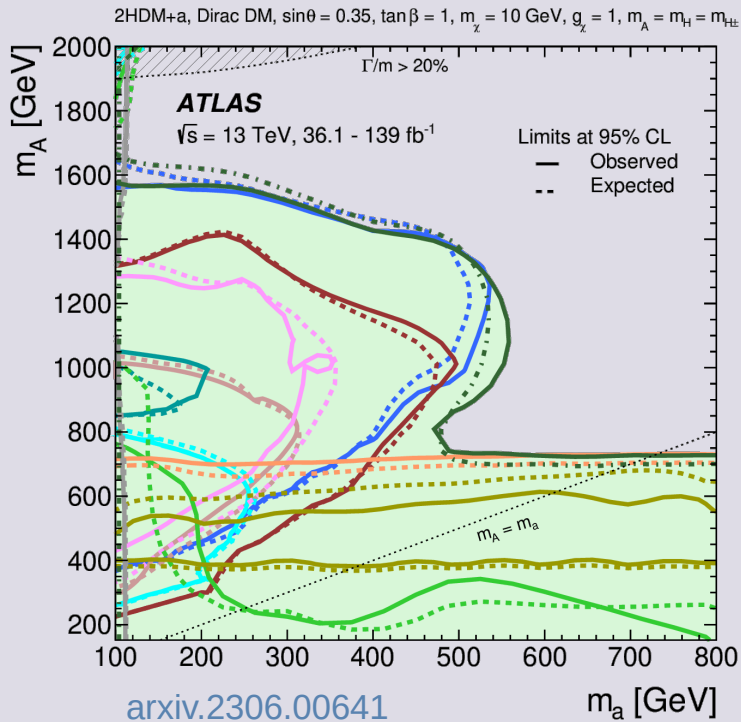
Two Higgs Doublet Model with an additional pseudoscalar DM mediator 2HDM+a: simplest gauge-invariant and renormalisable extension of the simplified pseudoscalar DM model

→ New channels for particle interaction → More distinctive collider signatures

Previous searches

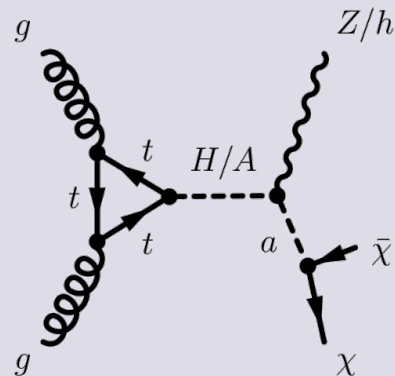
Exclusion m_A - m_a plane

No mass hierarchy: $m_A = m_H = m_{H^\pm}$

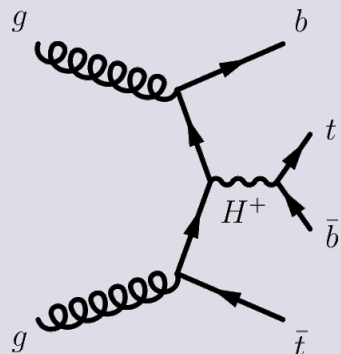


- $E_T^{\text{miss}} + h(b\bar{b})$, 139 fb $^{-1}$
 JHEP 11 (2021) 209
- $E_T^{\text{miss}} + h(\tau\tau)$, 139 fb $^{-1}$
 arXiv:2305.12938
- $E_T^{\text{miss}} + h(\gamma\gamma)$, 139 fb $^{-1}$
 JHEP 10 (2021) 13
- $E_T^{\text{miss}} + Z(l\bar{l})$, 139 fb $^{-1}$
 PLB 829 (2022) 137066
- $E_T^{\text{miss}} + Z(q\bar{q})$, 36.1 fb $^{-1}$
 JHEP 10 (2018) 180
- $E_T^{\text{miss}} + tW$, 139 fb $^{-1}$
 arXiv:2211.13138
- $E_T^{\text{miss}} + j$, 139 fb $^{-1}$
 PRD 103 (2021) 112006
- $tbH^\pm(tb)$, 139 fb $^{-1}$
 JHEP 06 (2021) 145
- $t\bar{t}t$, 139 fb $^{-1}$
 arXiv:2211.01136
- $h \rightarrow \text{invisible}$, 139 fb $^{-1}$
 arxiv:2301.10731
- **Combination**
 $E_T^{\text{miss}} + h(b\bar{b})$, $E_T^{\text{miss}} + Z(l\bar{l})$, $tbH^\pm(tb)$

- $E_T^{\text{miss}} + h(b\bar{b})$ and $E_T^{\text{miss}} + Z(l\bar{l})$ dominate the sensitivity



- tbH^\pm provides complementary sensitivity

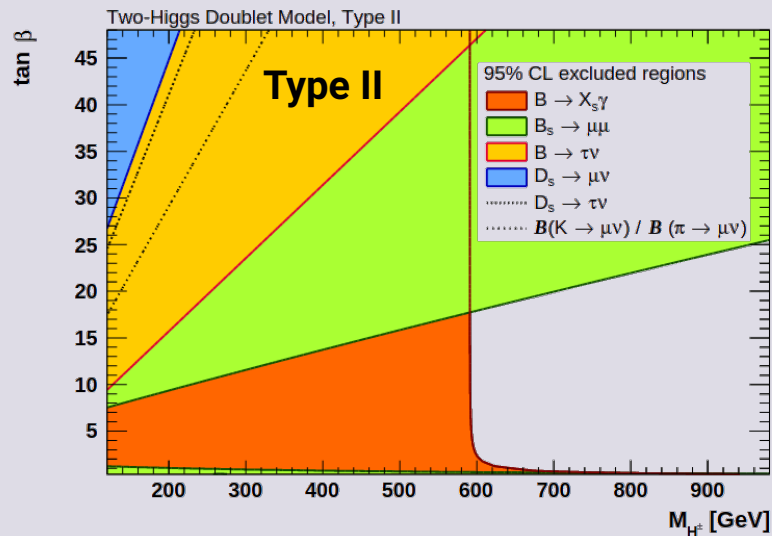
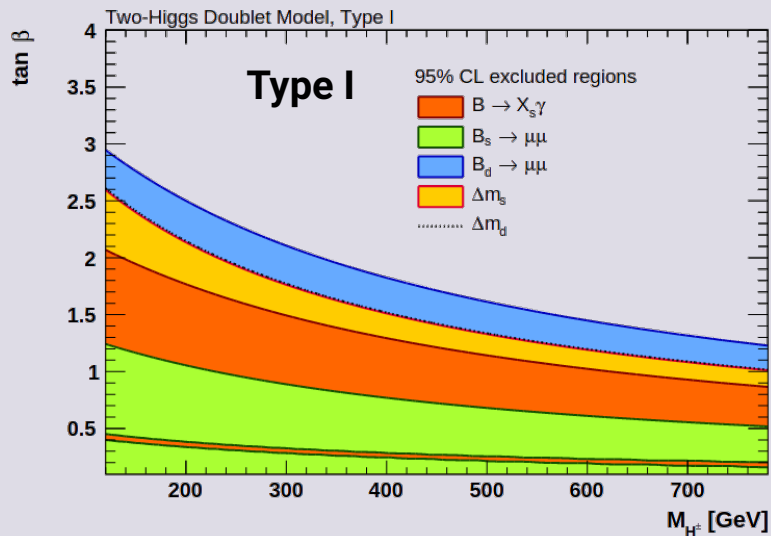


2HDM Type I

Why Type I?

All previous LHC searches consider only Type II Yukawa sector

- Constraints from flavour physics on charged Higgs mass are very weak in Type I → allow lower H^\pm masses

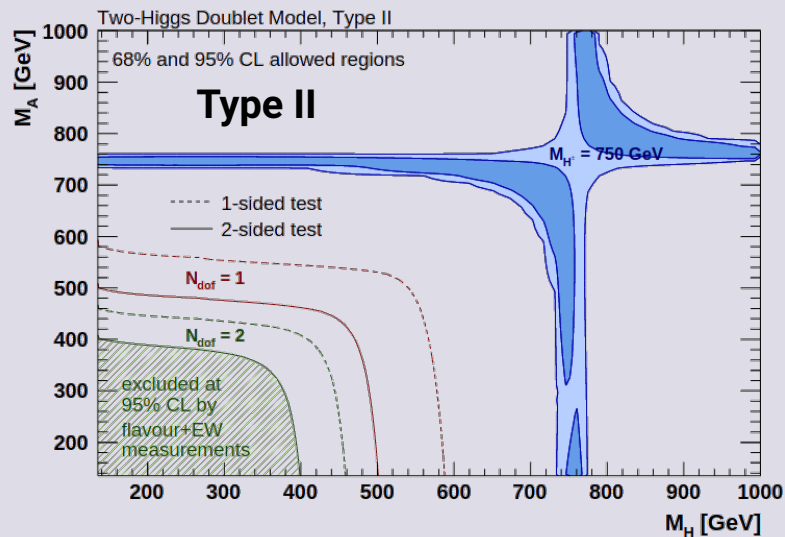
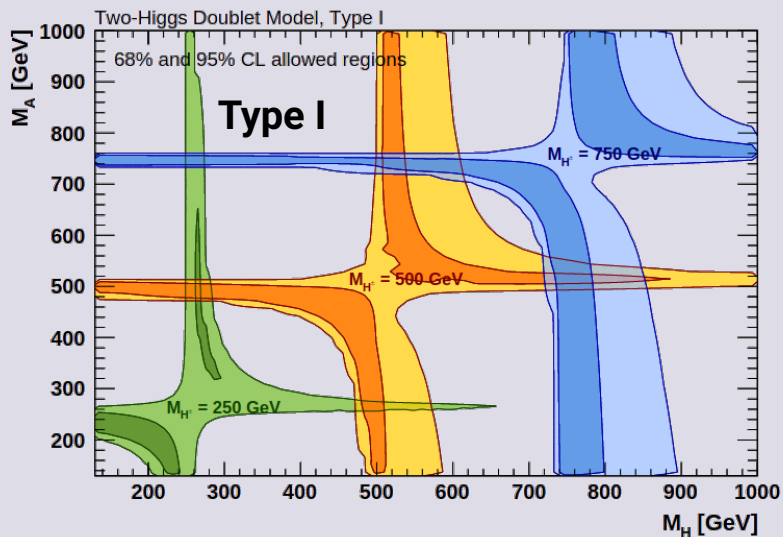


arxiv.1803.01853

2HDM Type I

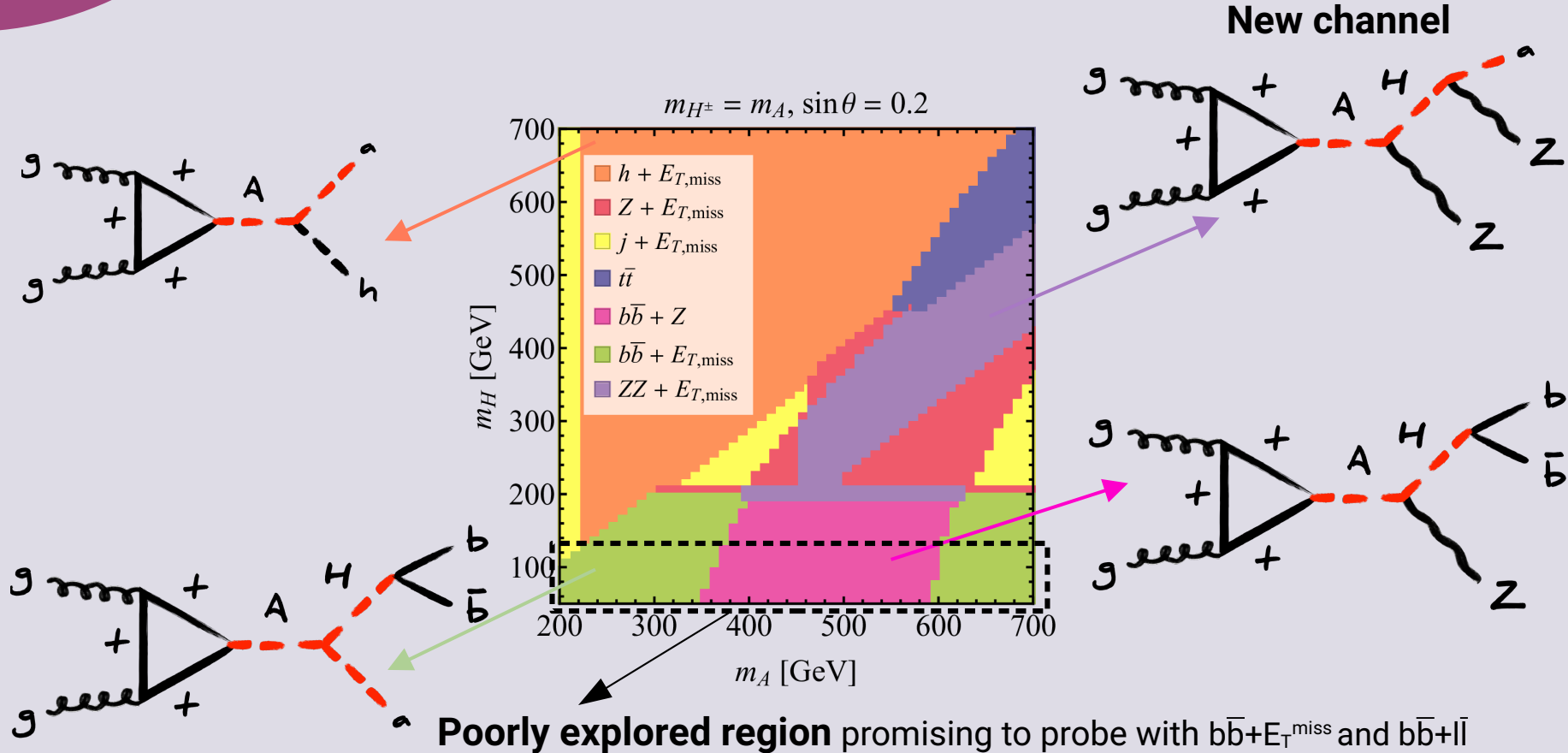
Why Type I?

- H^\pm should be close to the mass of A or H: Allow smaller m_{H^\pm} → Smaller allowed masses for A/H → *Explore masses below the SM Higgs mass*



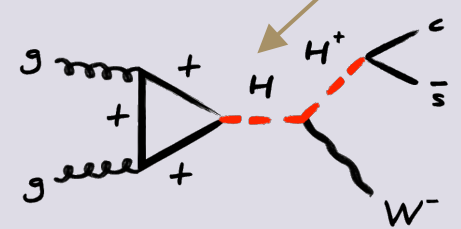
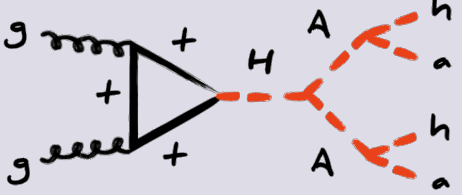
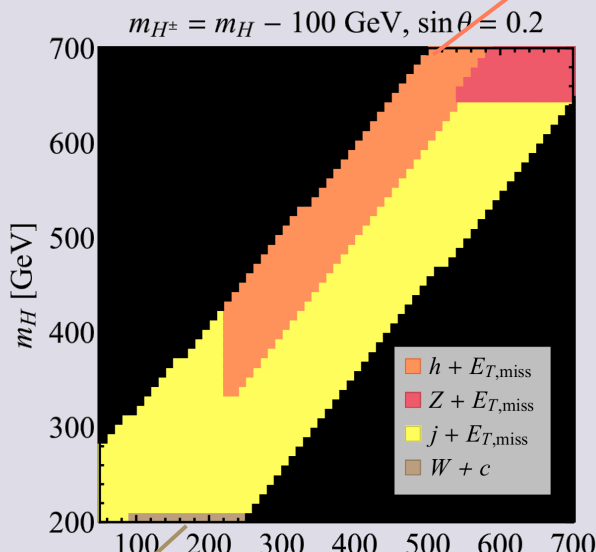
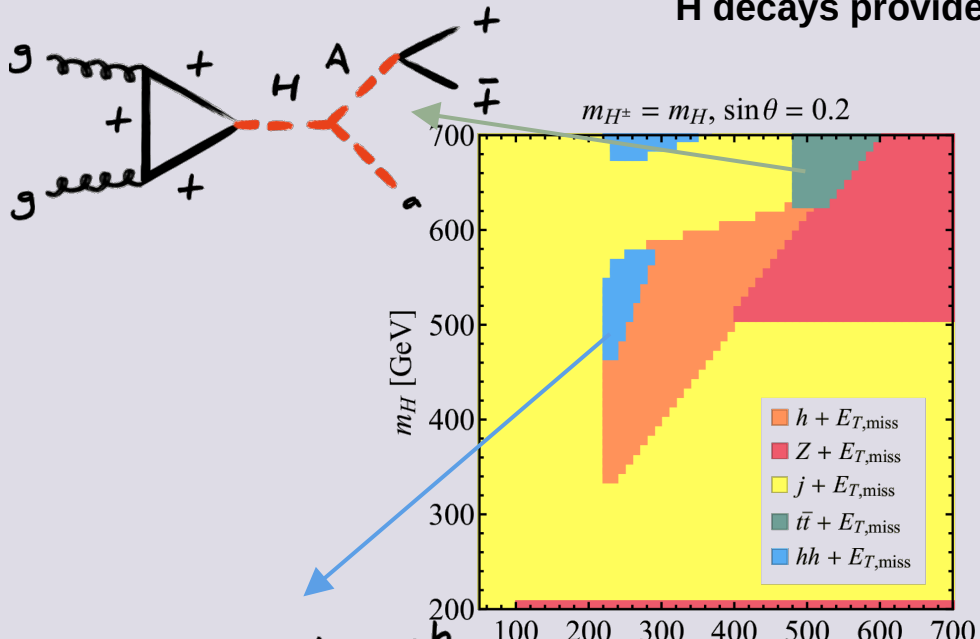
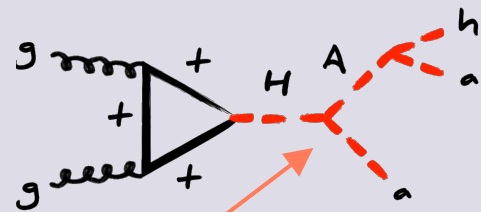
arxiv.1803.01853

A-boson dominant decays



H-boson dominant decays

H decays provide new channels



In the following we focus on $b\bar{b} + E_{T,miss}$ and $l\bar{l} + b\bar{b}$ produced with an A-boson resonance

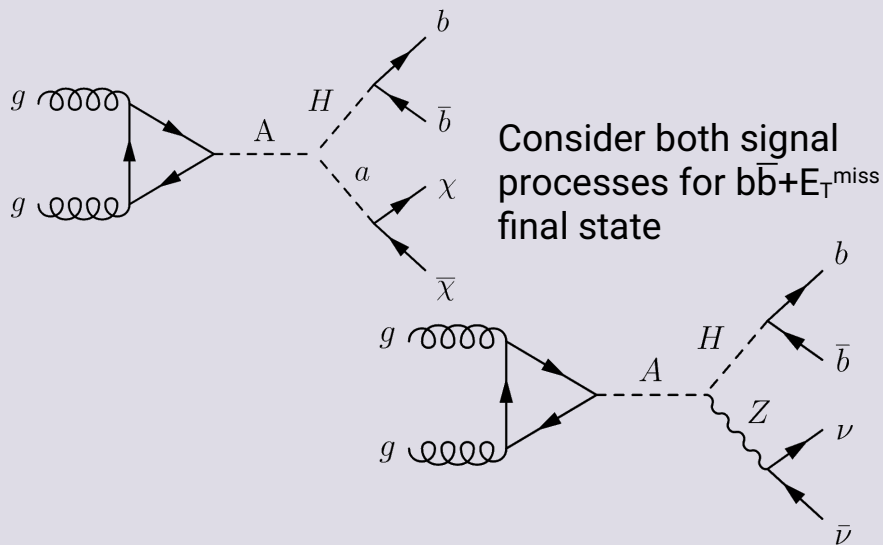
$b\bar{b} + E_T^{\text{miss}}$ signature

Requirements

- **0 leptons, exactly 2 b-jets, $E_T^{\text{miss}} > 150\text{GeV}$**

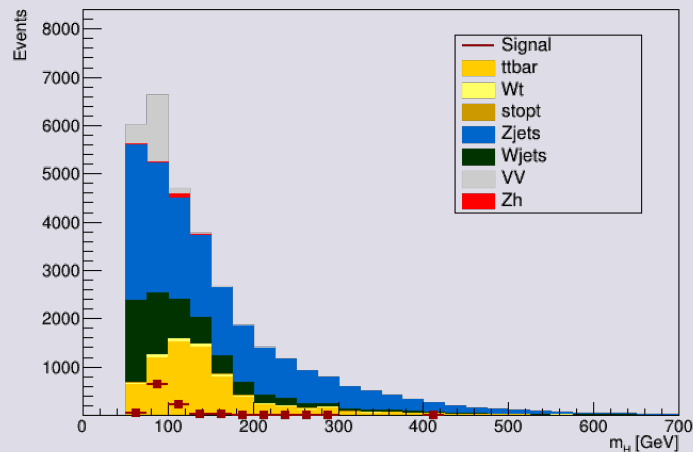
Reconstruction

- **a or Z:** Missing transverse momentum
- **H candidate:** 2 b-jets
- **Transverse mass for A candidate:** $H + E_T^{\text{miss}}$



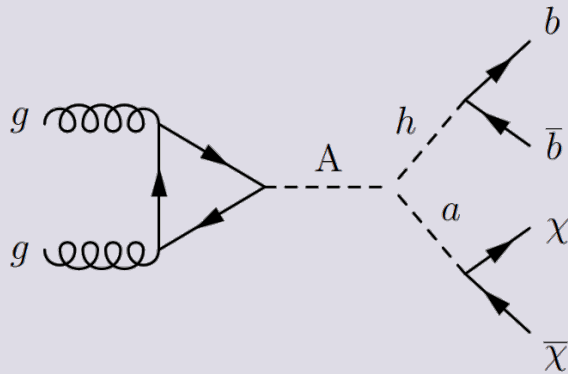
Apply cuts to improve sensitivity, similar to the ATLAS $A \rightarrow Z(\nu\nu)H(bb)$ analysis [arXiv:2311.04033](https://arxiv.org/abs/2311.04033)

Calculate sensitivity with m_H

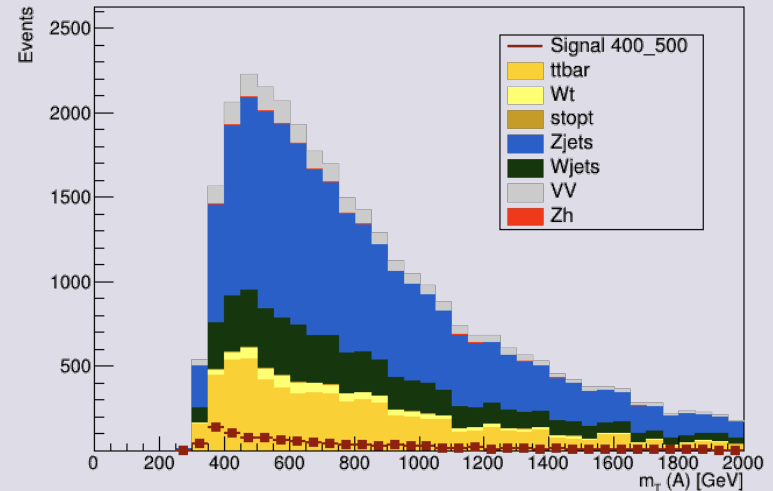


$h + E_T^{\text{miss}}$ signature

- Complementary exclusion for the phase space where $A \rightarrow Ha$ decay is not kinematically allowed
- Same cuts and reconstruction as $b\bar{b} + E_T^{\text{miss}}$

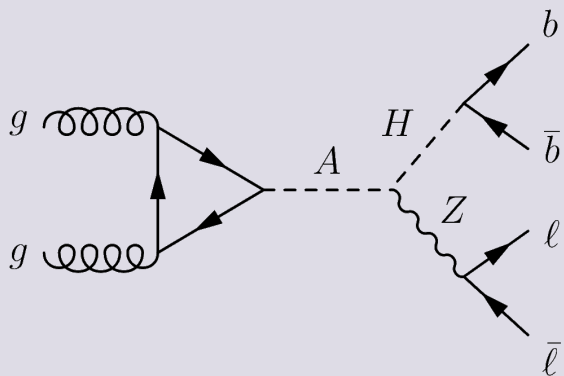


Calculate sensitivity with $m_T(A)$



$b\bar{b}+l\bar{l}$ signature

- Previous $A \rightarrow ZH \rightarrow llbb$ analyses both in [ATLAS](#) and [CMS](#) cover $m_{b\bar{b}}$ above 125 GeV

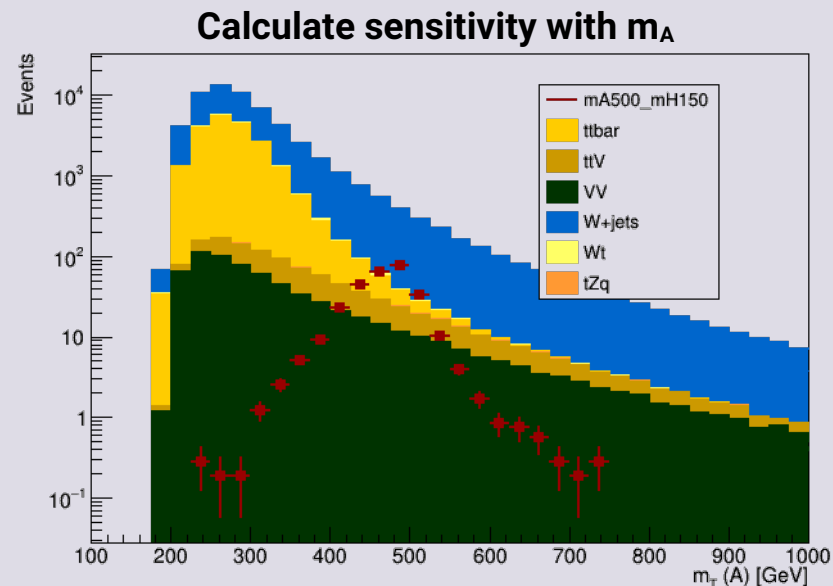


- **Z**: OSSF lepton-pair ($\mu^+\mu^-/e^+e^-$)
- **H candidate**: 2 b-jets
- **A candidate**: H+Z

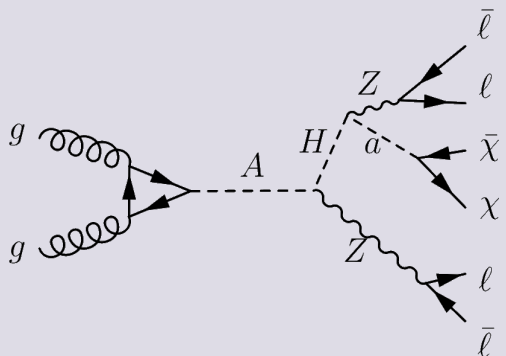
Cuts inspired from the ATLAS $A \rightarrow ZH \rightarrow llbb$ analysis

[arXiv:2011.05639](#)

- Apply m_{bb} window: $0.85m_H - 20 < m_{bb} < m_H + 20 \rightarrow$ Increase sensitivity



ZZ+E_T^{miss} signature

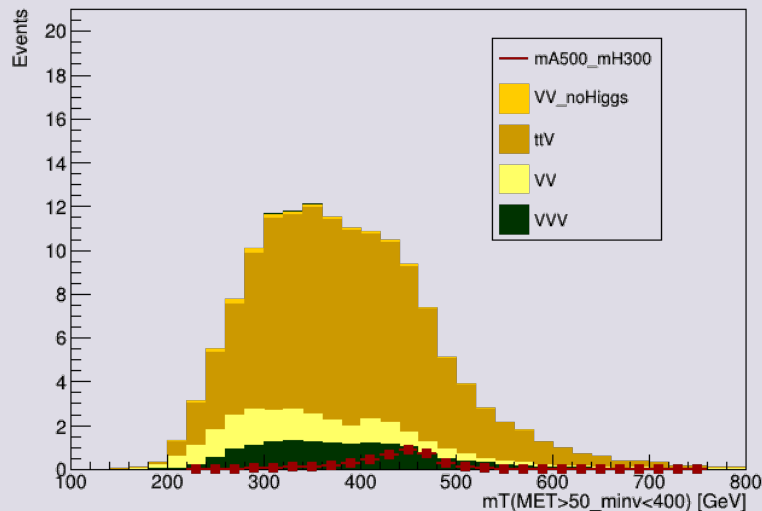


Cuts inspired from the ATLAS $l\bar{l}l\bar{l}+MET$ analysis [arxiv.2107.00404](https://arxiv.org/abs/2107.00404)

- Exactly four leptons
- $E_{T}^{miss} > 50$ GeV
- $m(4l) < 400$ GeV

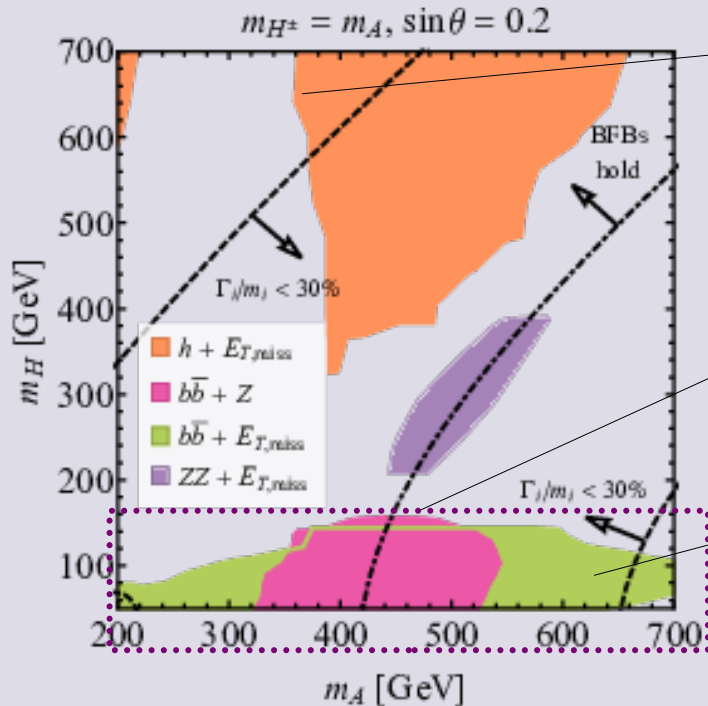
- Transverse mass for **A** candidate: $m^{inv}(4l)+E_{T}^{miss}$

Calculate sensitivity with $m_T(A)$



Results

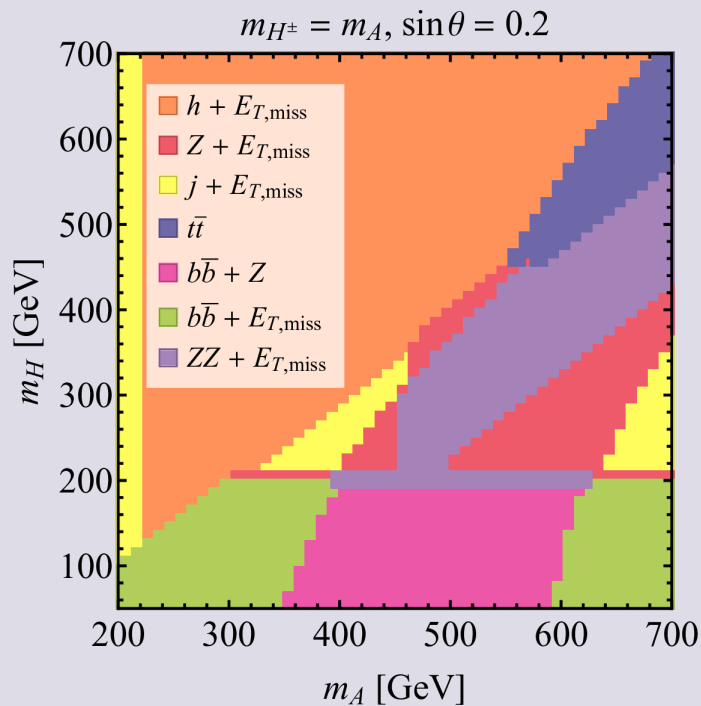
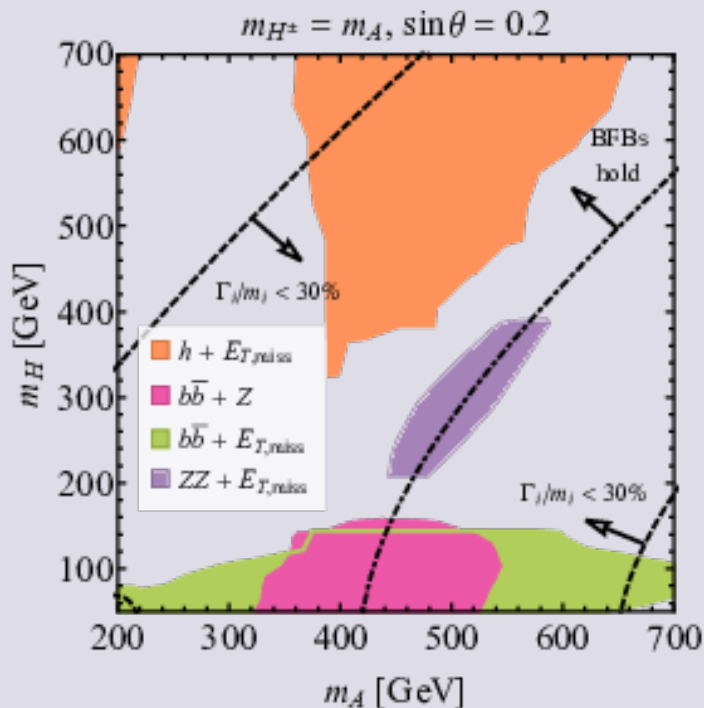
- Exclusion shown for Run2 luminosity except $ZZ+E_T^{\text{miss}}$ shown for Run2 + Run3



- Minimum requirement on E_T^{miss} sets a lower bound on $m_A \gtrsim 375$ GeV
- This restriction could be overcome by employing new techniques like **new b-jet triggers**
- $\text{Br}(H \rightarrow b\bar{b})$ drops drastically when $H \rightarrow aa/Za$ becomes kinematically possible
- Expand exclusion to masses **below** the mass of the **SM Higgs boson** → Particularly interesting to investigate observed small excess at ~ 95 GeV

Results

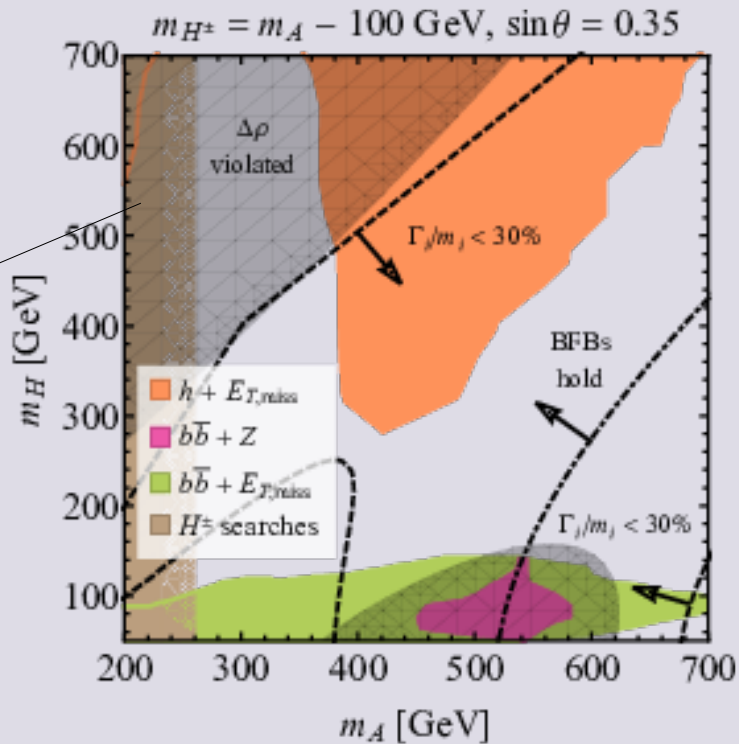
- Empty areas in the parameter region can be covered by further signatures such as $Z+E_T^{\text{miss}}$, $j+E_T^{\text{miss}}$, $t\bar{t}$



Results

- Allowing larger mass splitting $\Delta m = m_A - m_{H^\pm} = 100 \text{ GeV} \rightarrow A \rightarrow H^+W^-$ becomes kinematically possible

- **Extra constraints** through $A \rightarrow H^+W^-$
- Stronger restrictions on allowed parameter space



Conclusion

- **2HDM+a of Type I is not yet explored** → Leads to promising **new signatures**
- **Goal:** New benchmarks of uncovered final states → **New analyses with Run3 data**
- **New decay channels:** $A \rightarrow a H(bb)$, $A \rightarrow Z H(aZ)$
 $H \rightarrow a A(tt)$, $H \rightarrow H^+ W^-$
- $b\bar{b} + E_T^{\text{miss}}$ and $l\bar{l}b\bar{b}$ expand exclusion to masses **below the SM Higgs mass**

Novel collider signatures in the type-I 2HDM+a model

[arxiv.2404.05704](https://arxiv.org/abs/2404.05704)

Spyros Argyropoulos^a Ulrich Haisch^b and Iliia Kalaitzidou^a

^aPhysikalisches Institut, Universität Freiburg,
Hermann-Herder Str. 3a, 79104 Freiburg, Germany Freiburg, Germany

^bMax Planck Institute for Physics,
Föhringer Ring 6, 80805 München, Germany

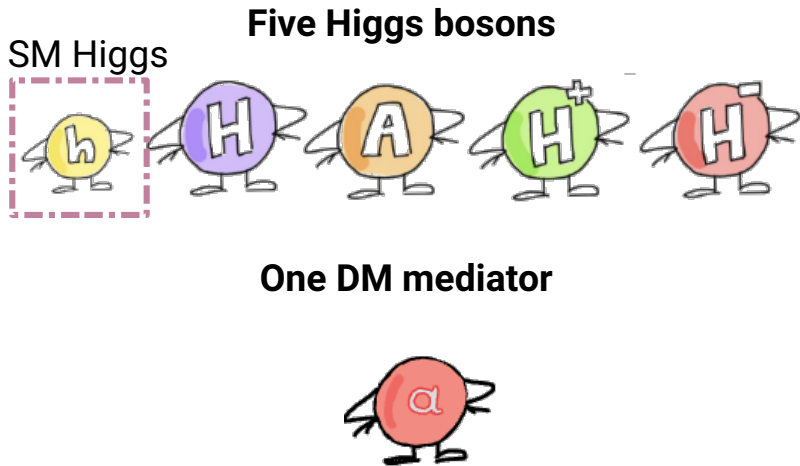
Thank you!



Back-up slides

2HDM+a theory

- Two Higgs doublets H_1, H_2 , one pseudoscalar singlet P



Mixing angles:

$\alpha \rightarrow$ Mixing of CP-even states ($H \leftrightarrow h$)

$$\beta \rightarrow \tan \beta \equiv \frac{v_2}{v_1}$$

$\theta \rightarrow$ Mixing of CP-odd states ($A \leftrightarrow a$)

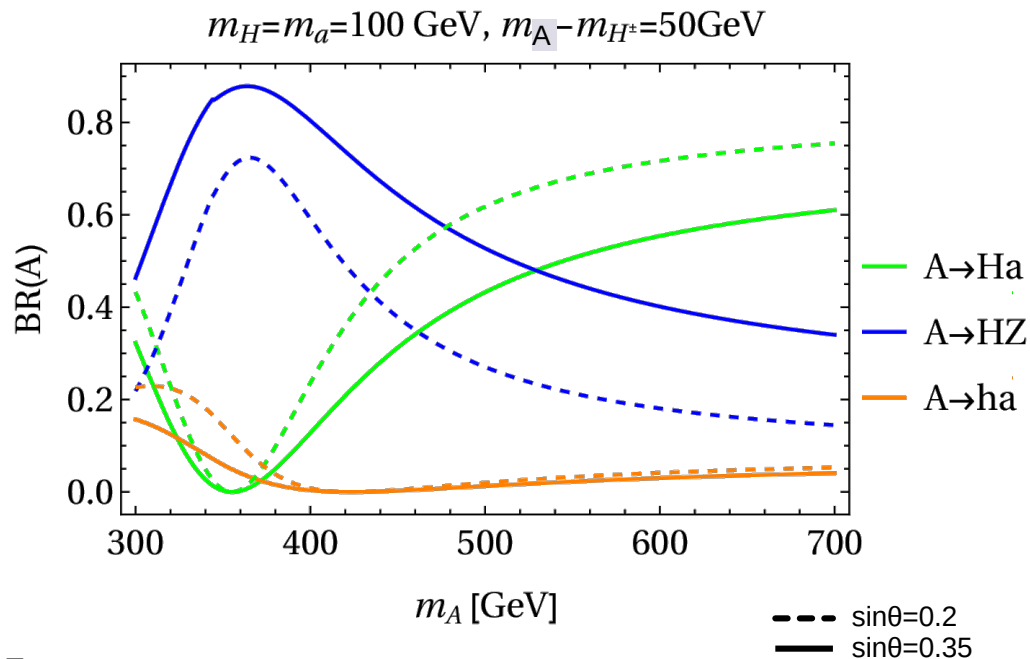
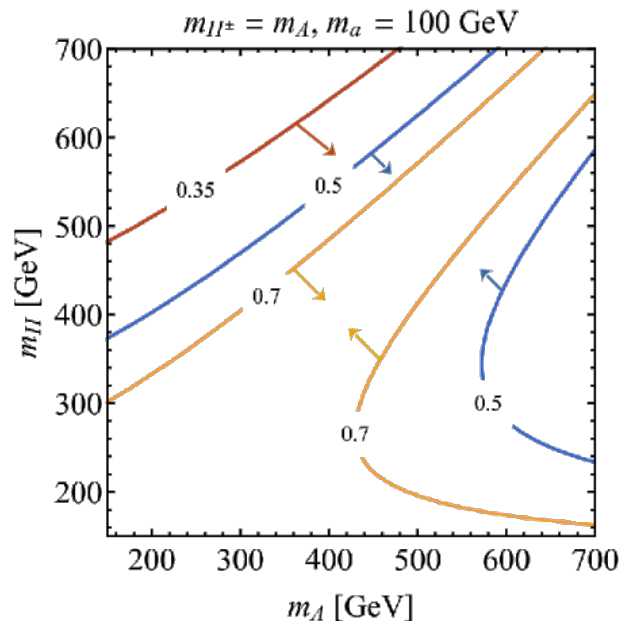
Couplings

	u-type	d-type	leptons	g_A^u	g_A^d
Type I	H_2	H_2	H_2	$1/\tan\beta$	$-1/\tan\beta$
Type II	H_2	H_1	H_1	$1/\tan\beta$	$\tan\beta$

- Masses: $m_A, m_H, m_{H^\pm}, m_a, m_\chi$ (DM mass)

Constraints on mixing angle θ

- Constraints on $\sin\theta$ from EW precision observables
- $\sin\theta$ choice affects A branching ratio



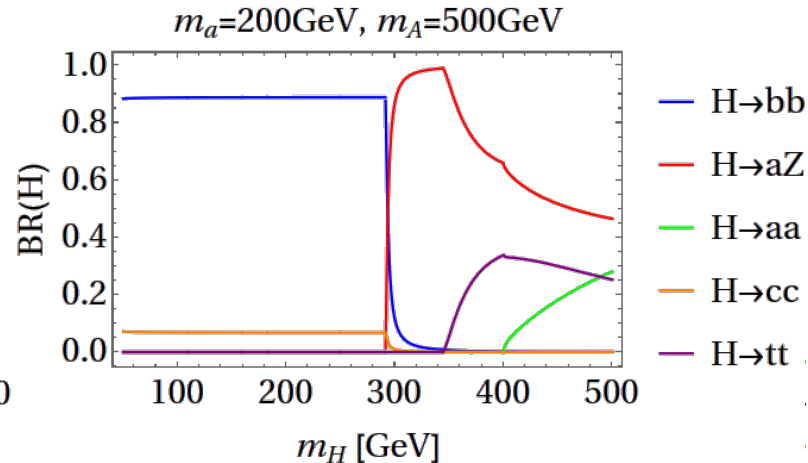
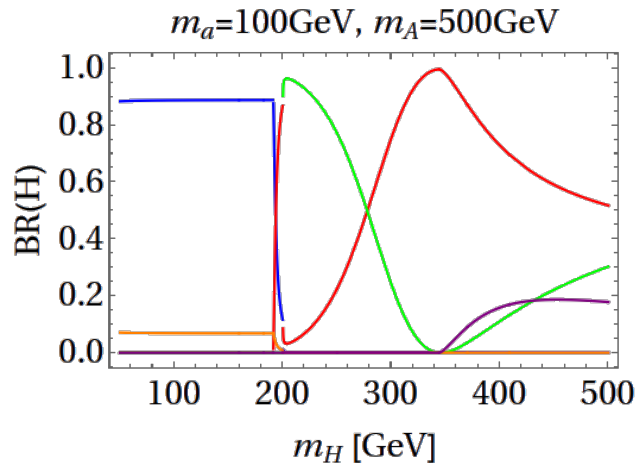
We choose **small $\sin\theta$** values of **0.2** and **0.35**

Constraints on Model Parameters

- Measurements for $h \rightarrow$ invisible constrain $m_a \gtrsim 100$ GeV

m_a choice affects H branching ratios

We choose:
 $m_a = 100$ GeV



- SM Higgs boson couplings \rightarrow tight constraints on $\cos(\beta-\alpha)$: We choose $\cos(\beta-\alpha)=0$ (alignment limit where $h \rightarrow$ SM Higgs)

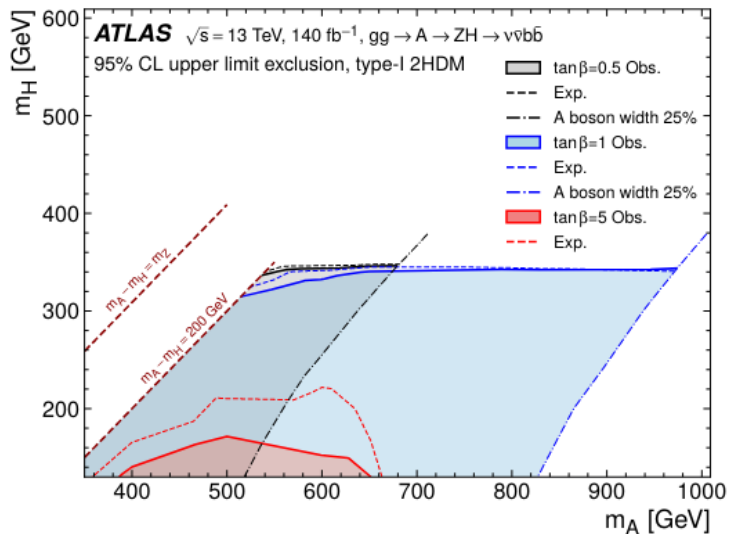
Additional constraints:

- Decay widths of BSM (pseudo)scalars should remain small \rightarrow narrow width approximation
- Scalar potential should be bounded from below (BFB) to assure stability

$\bar{b}b + E_T^{\text{miss}}$ signature

ATLAS $A \rightarrow Z(\nu\nu)H(bb)$ analysis

arXiv:2311.04033

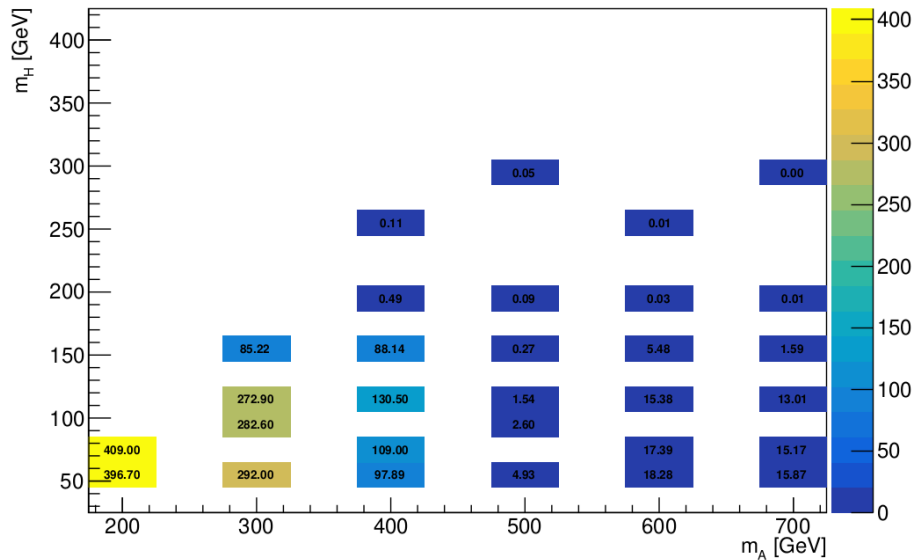


Requirement	Regions			
	2L (CR)	$e\mu$ (CR)	1L (VR)	\emptyset L
				Hlo / Hhi (CR)
Number of jets	2–5			
Number of b -jets	≥ 2			
$m(b\bar{b})$	> 50 GeV			
Number of τ^{had}	0			
$p_T(V)$	> 150 GeV			
$\min_i \Delta\phi(\vec{E}_T^{\text{miss}}, \vec{p}_i^{\text{jet}})$	$> \pi/10$			
$\Delta R(b_1, b_2)$	< 3.3 (2 b -jets)			
	< 3.5 (≥ 3 b -jets)			
Number of leptons	2		1	0
Lepton flavour	$ee/\mu\mu$	$e\mu$	e/μ	-
$p_T(\ell_1)$	> 27 GeV			-
$ m_Z^{\text{cand}} - m_Z $	< 10 GeV		-	
S_{MET}	< 5	-	> 3	> 10
$m_{\text{top}}^{\text{near}}$	-			> 180 GeV
$m_{\text{top}}^{\text{far}}$	-			> 200 GeV
$ m(b\bar{b}) - m_H $	-		$> 0.2 \cdot m_H$	$< 0.2 \cdot m_H$

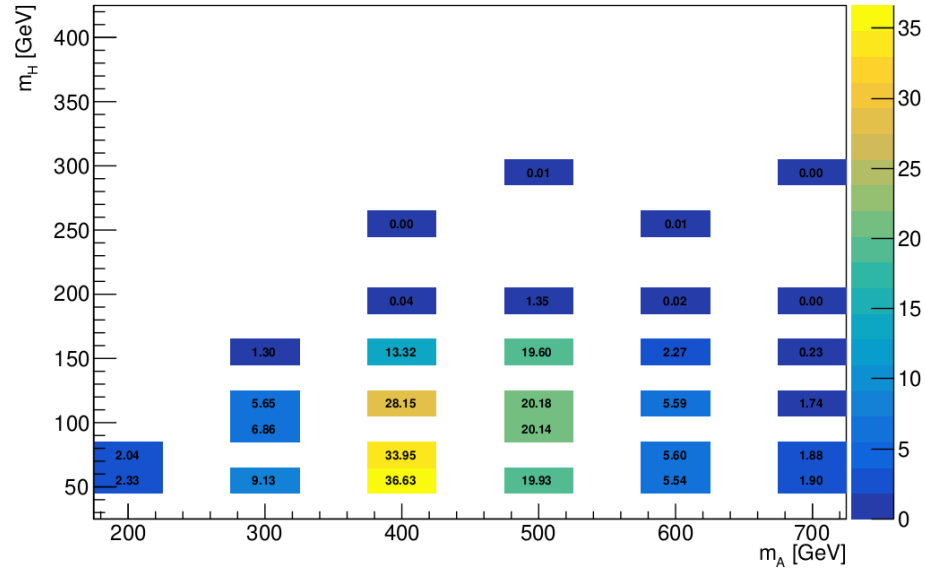
$b\bar{b} + E_T^{\text{miss}}$ signature

Cross-sections in fb

$\sigma(\text{ggA}) \times \text{Br}(\text{A} \rightarrow \text{aH}) \times \text{Br}(\text{a} \rightarrow \chi\chi) \times \text{Br}(\text{H} \rightarrow \text{bb})$



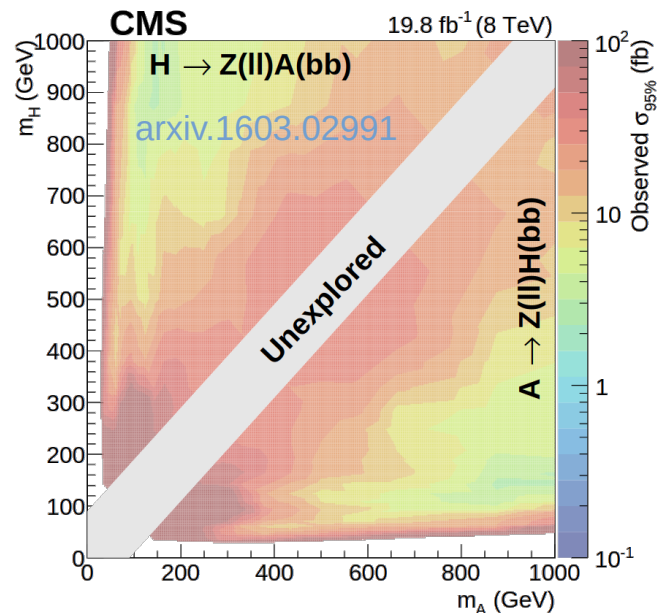
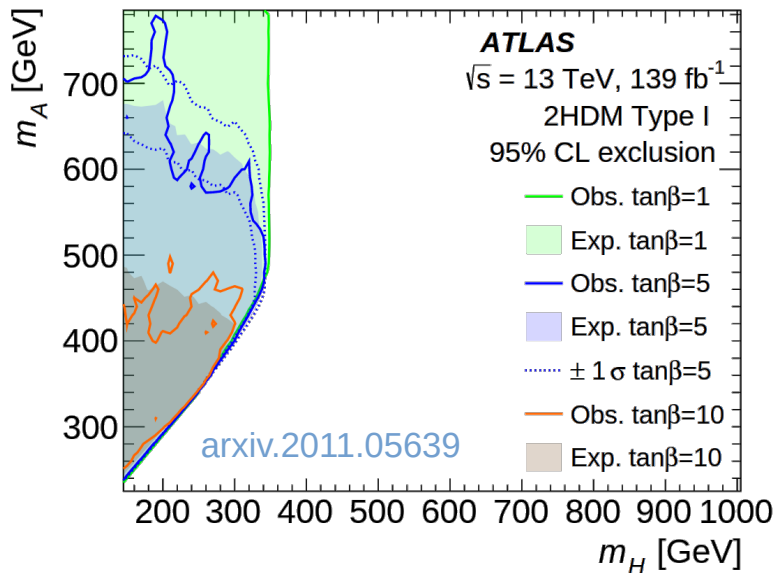
$\sigma(\text{ggA}) \times \text{Br}(\text{A} \rightarrow \text{ZH}) \times \text{Br}(\text{Z} \rightarrow \nu\nu) \times \text{Br}(\text{H} \rightarrow \text{bb})$



$b\bar{b} + \ell\bar{\ell}$ signature

Previous analyses

- Previous $A \rightarrow ZH \rightarrow \ell\bar{\ell} b\bar{b}$ analyses both in **ATLAS** and **CMS** cover $m_{b\bar{b}}$ above 125 GeV



$b\bar{b} + \ell\bar{\ell}$ signature

- $A \rightarrow ZH \rightarrow \ell\bar{\ell}b\bar{b}$ in ATLAS
arxiv.2011.05639

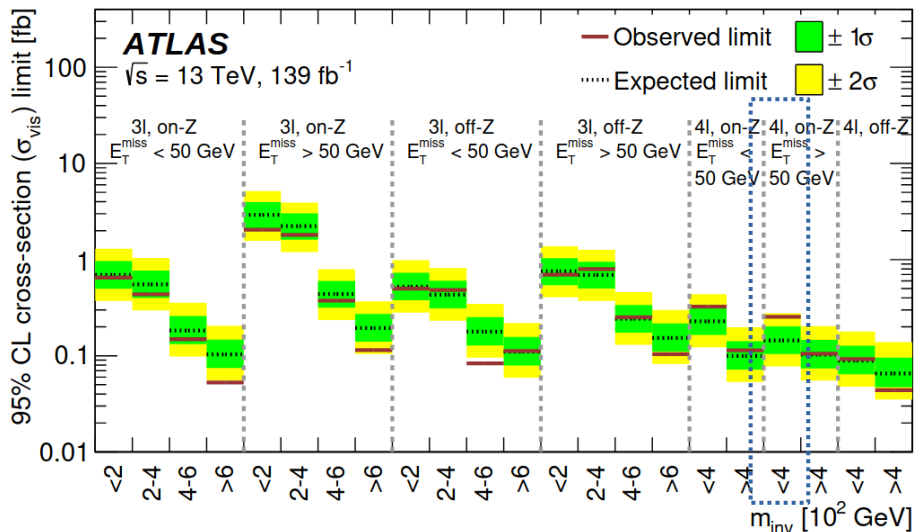
Table 1: Summary of the event selection for signal and control regions in the $A \rightarrow ZH \rightarrow \ell\bar{\ell}b\bar{b}$ channel.

Single-electron or single-muon trigger		
Exactly 2 leptons (e or μ) ($p_T > 7$ GeV) with the leading one having $p_T > 27$ GeV		
Opposite electric charge for $\mu\mu$ pairs; $80 \text{ GeV} < m_{\ell\ell}, e\mu < 100 \text{ GeV}$, $\ell = e, \mu$		
At least 2 b -jets ($p_T > 20$ GeV) with one of them having $p_T > 45$ GeV		
$E_T^{\text{miss}}/\sqrt{H_T} < 3.5 \text{ GeV}^{1/2}$, $\sqrt{\Sigma p_T^2}/m_{\ell\ell b\bar{b}} > 0.4$		
	$n_b = 2$ category	$n_b \geq 3$ category
	Exactly 2 b -tagged jets	At least 3 b -tagged jets
Signal region	ee or $\mu\mu$ pair $0.85 \cdot m_H - 20 \text{ GeV} < m_{b\bar{b}} < m_H + 20 \text{ GeV}$	ee or $\mu\mu$ pair $0.85 \cdot m_H - 25 \text{ GeV} < m_{b\bar{b}} < m_H + 50 \text{ GeV}$
Z+jets control region	ee or $\mu\mu$ pair $m_{b\bar{b}} < 0.85 \cdot m_H - 20 \text{ GeV}$ or $m_{b\bar{b}} > m_H + 20 \text{ GeV}$	ee or $\mu\mu$ pair $m_{b\bar{b}} < 0.85 \cdot m_H - 25 \text{ GeV}$ or $m_{b\bar{b}} > m_H + 50 \text{ GeV}$
Top control region	$e\mu$ pair $0.85 \cdot m_H - 20 \text{ GeV} < m_{b\bar{b}} < m_H + 20 \text{ GeV}$	$e\mu$ pair $0.85 \cdot m_H - 25 \text{ GeV} < m_{b\bar{b}} < m_H + 50 \text{ GeV}$

ZZ+E_T^{miss} signature

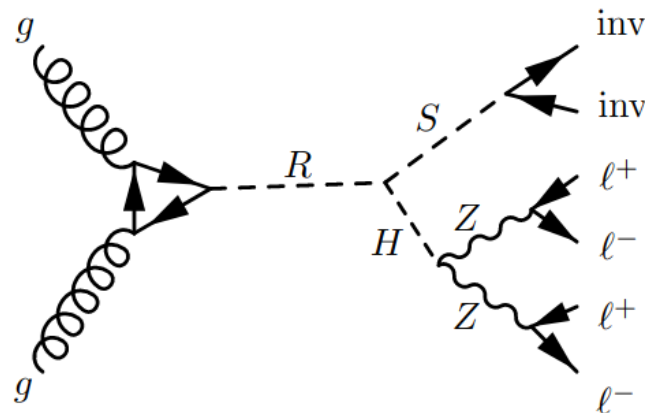
Selection from [arxiv.2107.00404](https://arxiv.org/abs/2107.00404)

- Exactly four leptons
- $|m_Z - 91.2| < 10$ GeV
- $p_T(l) > 25$ GeV
- $E_T^{\text{miss}} > 50$ GeV
- $m(4l) < 400$ GeV



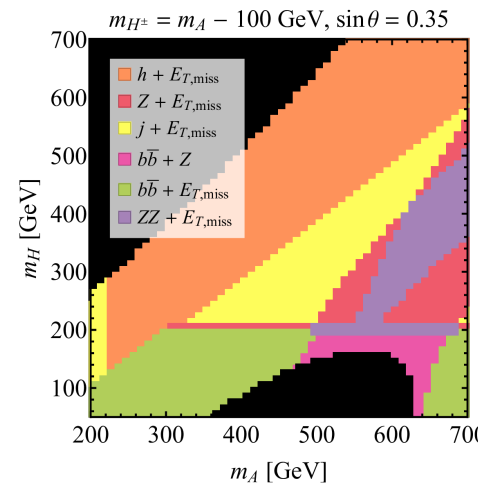
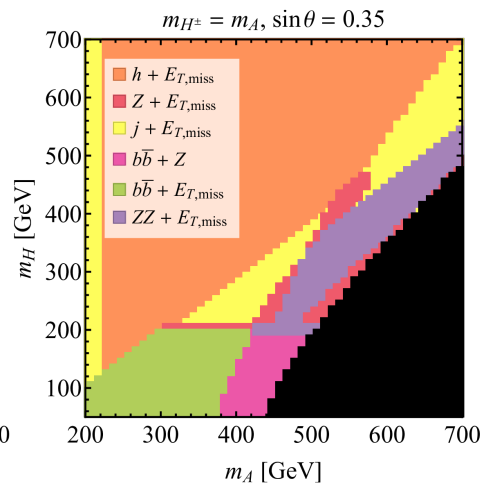
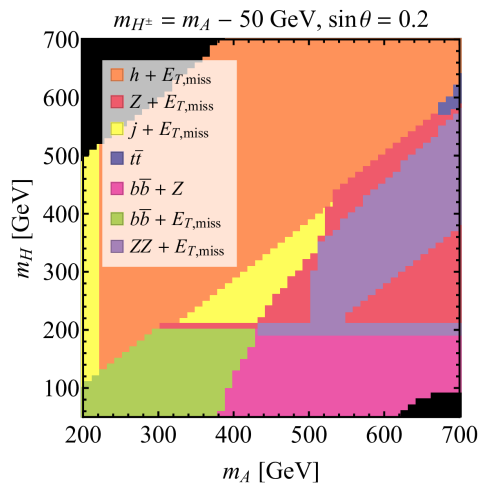
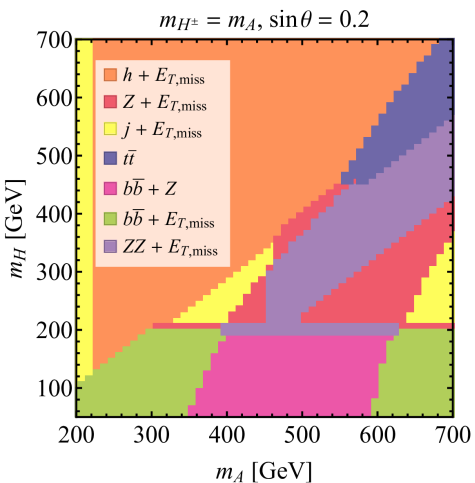
[arxiv.2401.04742](https://arxiv.org/abs/2401.04742)

Additional ATLAS 4l+MET search with 4l coming from H resonance



Benchmark scenarios

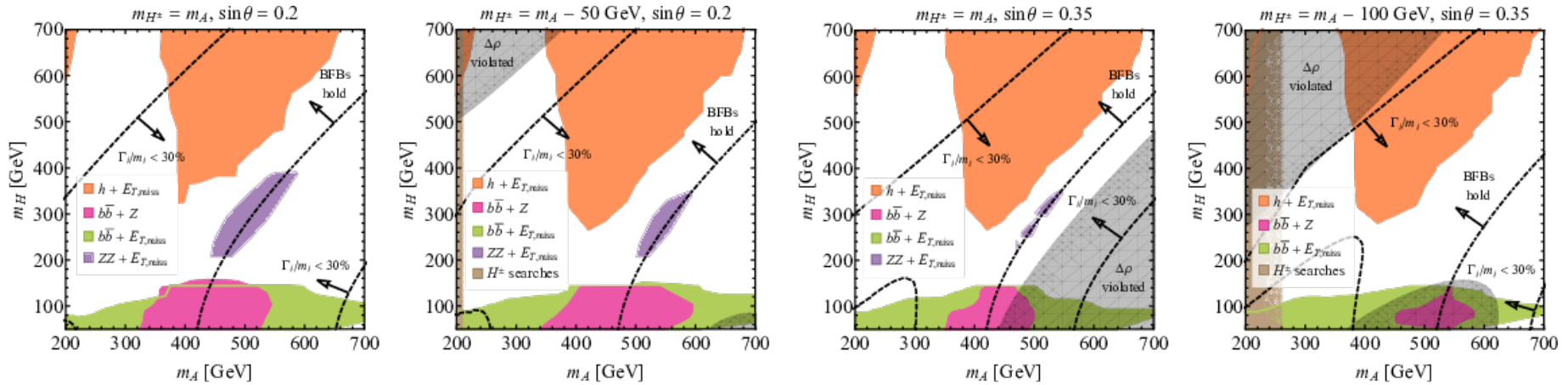
Study four benchmark points for different $\Delta m = m_A - m_{H^\pm}$ and $\tan\beta$



Black: Constraints from EW precision measurements and decay widths $> 30\%$

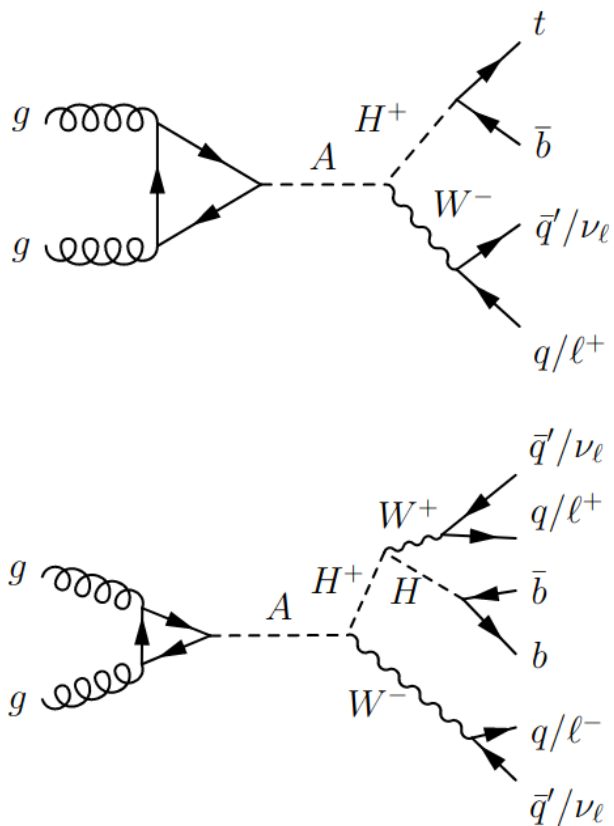
Exclusion

- Show four different benchmark points for different $\Delta m = m_A - m_{H^\pm}$ and $\sin\theta$



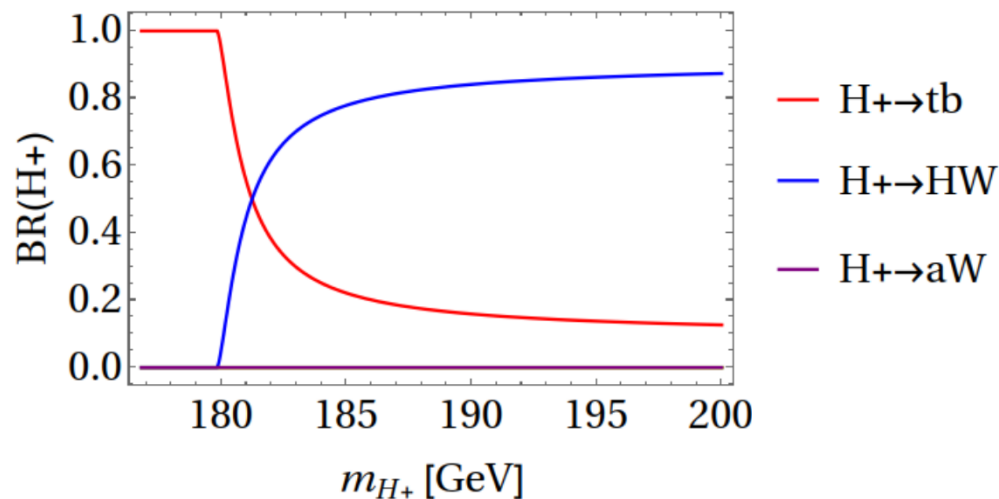
- $\Delta\rho$ violated: Constraints from electroweak precision observables
- BFBs hold: Scalar potential is bounded from below
- $\Gamma_i/m_i < 30\%$: Decay widths of scalars should remain small

Signatures with charged Higgs



- Allowing **larger mass splitting** between m_A and m_{H^\pm} makes further **new unexplored signal signatures** kinematically possible such as $A \rightarrow H^+W^-$

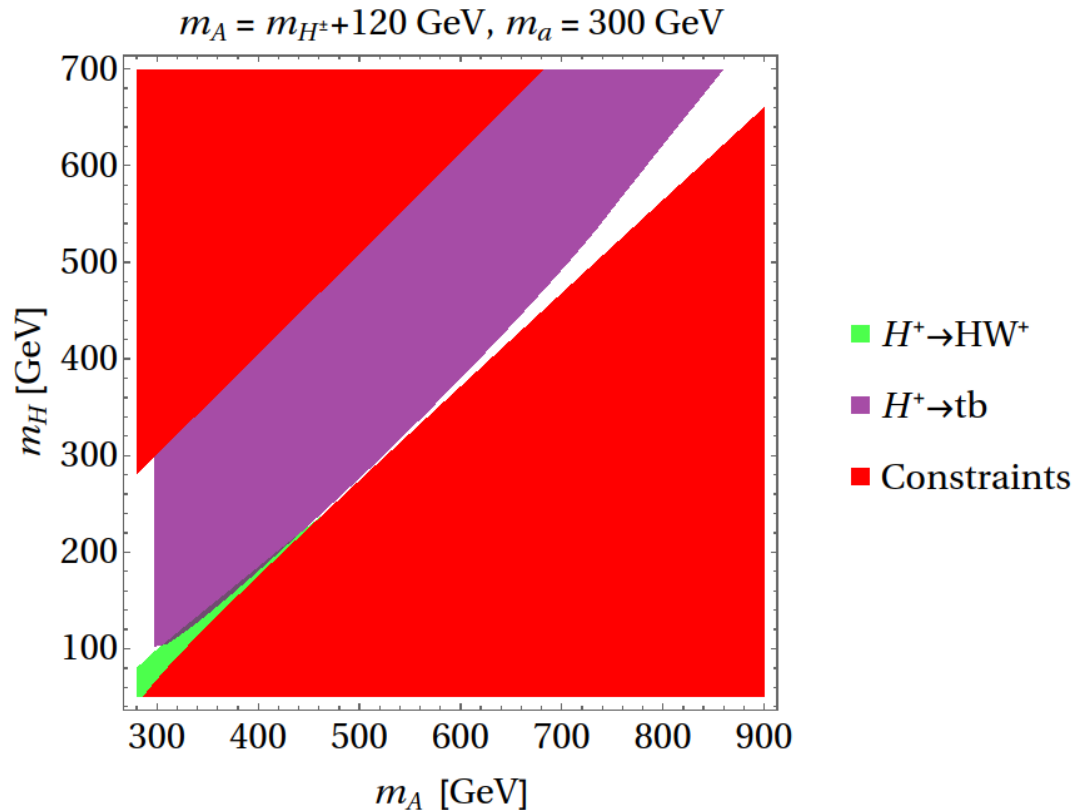
$$m_{H^\pm} = m_A - 120, \quad m_H = 100 \text{ GeV}, \quad m_a = 600 \text{ GeV}$$



Signatures with charged Higgs

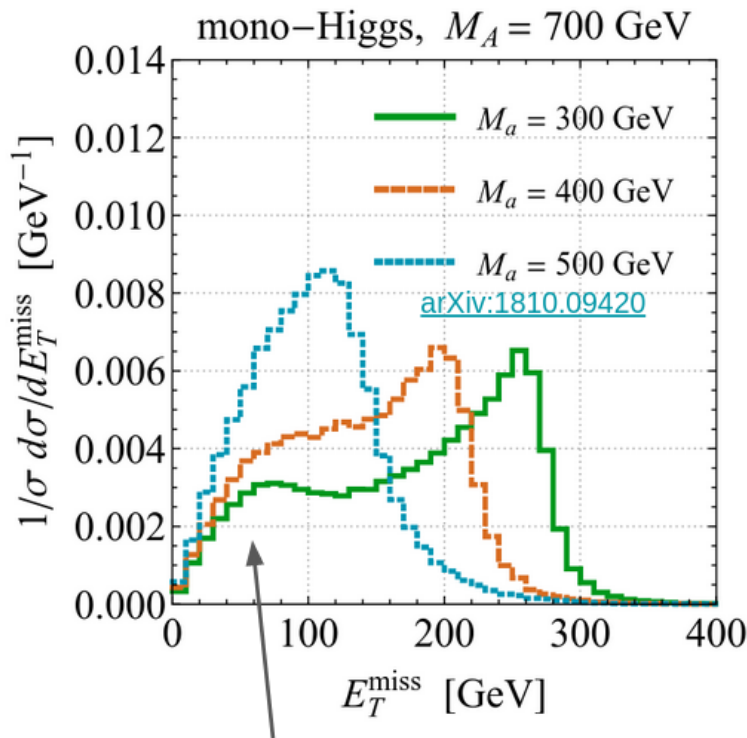
The $A \rightarrow H^+W^-$ decay

- **No previous $A \rightarrow H^+W^-$ analysis**
- Only a small region (bottom left corner) is sensitive for the $H^\pm \rightarrow W^\pm H$ decay
- Larger region where the $H^\pm \rightarrow t\bar{b}$ decay is important
- Both of them give a final state not previously explored



$b\bar{b} + E_T^{\text{miss}}$ signature

Impact of Box Diagrams



$$E_{T,\text{max}}^{\text{miss}} \approx \frac{\lambda^{\frac{1}{2}}(M_A, M_{h,H}, M_a)}{2M_A}$$

with

$$\lambda(M_A, M_{h,H}, M_a) = (M_A^2 - M_{h,H}^2 - M_a^2)^2 - 4M_{h,H}^2 M_a^2$$

Box diagrams become important for large mass difference between a and A (decay $A \rightarrow h_{\text{SM}} a(xdx)$)