A theoretical overview on *t*-channel models and their phenomenology

Luca Panizzi



Roadmap of Dark Matter models for Run 3 - CERN 13-17 May 2024

Motivation

A white paper on t-channel scenarios is being written Dark Matter via *t*-channel Production A Report of the LHC Dark Matter Working Group

Joint effort TH-EXP to provide guidelines and benchmarks for new analysis during Run 3 and future upgrades

More than 50 authors involved

Study of scenarios based on the schematic interaction



Why is this important?



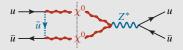
Complementary to s-channel

t-channel

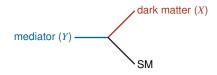
s-channel

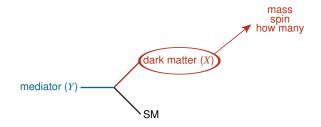
mediator always heavier than DM mediator can also be lighter than DM odd mediators allowed in interactions

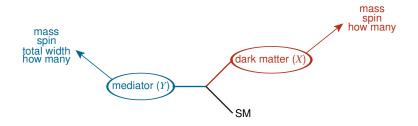
But interferences can happen in non-minimal/full models...

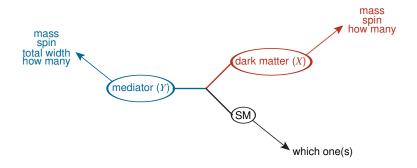


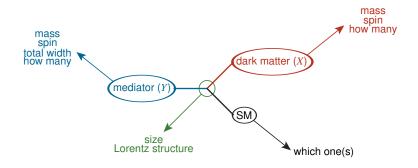
coloured mediators interesting at a hadron collider

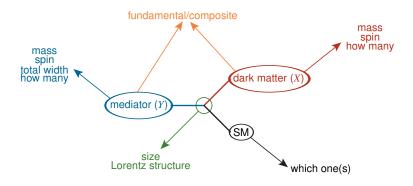


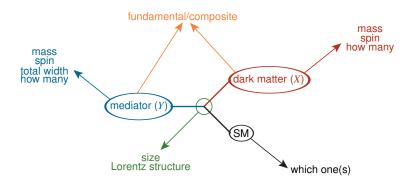










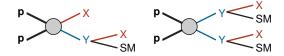


Depending on the possibilities:

- Can we observe a signal? And how?
- How does cosmology constrain the parameters?
- How do we reinterpret results?
- Can we define benchmarks for LHC to cover the widest range of possibilities?

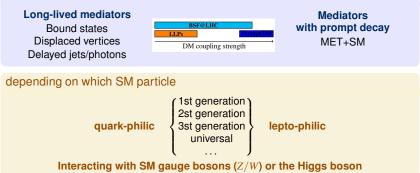
Which signatures





Not all processes might be possible at tree-level

depending on coupling or mass splitting



This talk: quark-philic scenarios with prompt-decay mediators

Classification of simplified scenarios

Real DM				Complex DM					
Mediator spin					Mediator spin				
		0	1/2	1			0	1/2	1
DM	0	×	F3S	×	DM	0	×	F3C	×
spin	1/2	S3M	×	to be done	spin	1/2	S3D	×	to be done
	1	×	F3V	×		1	×	F3W	×

Examples of theories which can be described by these simplified models

S3M	SUSY: squarks+neutralino (Majorana fermion)
S3D	Right-handed neutrino portals with extended scalar sectors
F3S	UED: KK quark partners + KK photon (real scalar)
F3C	SUSY: sleptons+sneutrinos (not aware of quark-philic models)
F3V	?
F3W	FPVDM: vector-like quark + vector DM (non-abelian gauge boson)

Complex DM scenarios excluded by cosmology for interactions with light quarks Is it true also for non-minimal models? Is it true also for bottom and top?

Numerical models

Simplified models suitable for performing MC simulations at NLO in QCD and testing against cosmological observables

Coloured mediators

DMSimpt : A general framework for t-channel dark matter models at NLO in QCD

Contact Information

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See ParXiv:2001.05024 [hep-ph].

Model Description and FeynRules Implementation

We extend the Standard Model by a dark matter candidate X and a coloured mediator Y. The model includ or bosonic dark matter) or 0 (fermionic dark matter). The model Lagrangian is given by

 $\mathcal{L} = \mathcal{L}_{\mathrm{SM}} + \mathcal{L}_{\mathrm{kin}} + \mathcal{L}_{F}(\chi) + \mathcal{L}_{F}(\bar{\chi}) + \mathcal{L}_{S}(S) + \mathcal{L}_{S}(\bar{S}) + \mathcal{L}_{V}(V) + \mathcal{L}_{V}(\bar{V}) \; .$

The first term consists in the Standard Model Lagrangian, the second one includes gauge-invariant kinetic Dirac fermion, Majorana fermion, complex scalar, real scalar, complex vector and real vector dark matter,

$$\begin{split} \mathcal{L}_F(X) &= \left[\lambda_{\mathbf{Q}} X Q_L \varphi_{\mathbf{Q}}^{\dagger} + \lambda_{\mathbf{u}} X u_R \varphi_{\mathbf{u}}^{\dagger} + \lambda_{\mathbf{d}} X d_R \varphi_{\mathbf{d}}^{\dagger} + \mathbf{h.c.} \right], \\ \mathcal{L}_S(X) &= \left[\lambda_{\mathbf{Q}} \bar{\psi}_{\mathbf{Q}} Q_L X + \lambda_{\mathbf{u}} \bar{\psi}_{\mathbf{u}} u_R X + \lambda_{\mathbf{d}} \bar{\psi}_{\mathbf{u}} d_R X + \mathbf{h.c.} \right], \\ \mathcal{L}_V(X) &= \left[\lambda_{\mathbf{Q}} \bar{\psi}_{\mathbf{Q}} \gamma^{\mu} X_{\mu} Q_L + \lambda_{\mathbf{u}} \bar{\psi}_{\mathbf{u}} \gamma^{\mu} X_{\mu} u_R + \lambda_{\mathbf{d}} \bar{\psi}_{\mathbf{u}} \gamma^{\mu} A_{\mu} d_R + \mathbf{h.c.} \right], \end{split}$$

where ϕ and ψ consists in coloured scalar and fermionic mediators.

http://feynrules.irmp.ucl.ac.be/wiki/DMsimpt

	Spin		
Mediator	0	1/2	
Dark matter	1/2	0 or 1	

- DM real or complex
- Couplings with any SM quark
- Restrictions to select representations or coupling hierarchies (only one generation, universal couplings...)

C. Arina, B. Fuks and L. Mantani, Eur. Phys. J. C 80 (2020) no.5, 409, [arXiv:2001.05024 [hep-ph]].

Other models available for specific problems (leptophilic DM, multi-component DM...) A unified model will also be released

We need to provide useful information for both TH and EXP community

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 Accurate kinematical description of the signal LO vs NLO

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Accurate kinematical description of the signal
 I O vs NLO

\rightarrow	LO vs NLO	

Process	LO	NLO
XX	$\begin{array}{c} q \longrightarrow X \\ \overline{q} \longrightarrow X \\ \overline{q} \longrightarrow X \end{array}$	$\begin{array}{c} q \\ g \\ \overline{q} \\ \overline$
ХҮ	$\begin{array}{c} q \longrightarrow X \\ Y \\ g \\ g$	
YY	8 0000000 ···· Y Y g 0000000 ···· Y	

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Process	LO	NLO
ХХ	$\begin{array}{c} q \longrightarrow X \\ Y & \\ \overline{q} \longrightarrow X \end{array}$	$\begin{array}{c} q \rightarrow & X \\ g \downarrow & Y \\ \overline{q} \rightarrow & X \\ \overline{q} \rightarrow & X \\ \end{array} \begin{array}{c} \chi \\ g \downarrow & \chi \\ \chi \\ g \\ \overline{q} \rightarrow & X \\ \chi \\ g \\ \overline{q} \rightarrow & \chi \\ \chi \\ g \\ \overline{q} \rightarrow & \chi \rightarrow \\ q$
ХҮ	$\begin{array}{c} q \longrightarrow & X \\ Y & \\ g & \\ g & \\ g & \\ g & \\ y & \\ y$	$\begin{array}{c} q & X \\ g & Y \\$
YY	g 20000000	q coorsococo Y g coorsoco Y g y Y Y corroros g g coorsococo Y g seesee

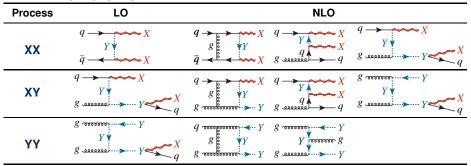
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Accurate kinematical description of the signal
 LO vs NLO

Process	LO	NLO
хх	$\begin{array}{c} q \longrightarrow X \\ \hline q \longrightarrow X \\ \hline q \longrightarrow X \end{array}$	$\begin{array}{c} q \xrightarrow{\qquad \ \ } X \\ g \xrightarrow{\qquad \ \ } Y \\ \overline{q} \xrightarrow{\qquad \ \ } X \\ \overline{q} \xrightarrow{\qquad \ \ } X \\ \overline{q} \xrightarrow{\qquad \ \ } X \\ g \xrightarrow{\qquad \ \ } y \xrightarrow{\qquad \ \ } q \\ \overline{q} \xrightarrow{\qquad \ \ } X \\ g \xrightarrow{\qquad \ \ } y \xrightarrow{\qquad \ } q \\ \overline{q} \xrightarrow{\qquad \ \ } y \xrightarrow{\qquad \ } q \\ \overline{q} \xrightarrow{\qquad \ \ } y \xrightarrow{\qquad \ } q \\ \overline{q} \xrightarrow{\qquad \ \ } y \xrightarrow{\qquad \ } q \\ \overline{q} \xrightarrow{\qquad \ \ } y \xrightarrow{\qquad \ } y \xrightarrow{\qquad \ } y \xrightarrow{\qquad \ } y \qquad \ $
ХҮ	$\begin{array}{c} q \longrightarrow & X \\ Y & \\ g & \\ g$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$
YY	g	q correspondent Y g correspondent Y g y Y Y respondent S g correspondent Y g correspondent Y

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Accurate kinematical description of the signal
 LO vs NLO



Double-counting between real emission and tree-level processes Removed through suitable algorythm in MadGraph (MadSTR)

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- Accurate kinematical description of the signal
 - LO vs NLO
 - \longrightarrow beware of limitations: narrow width approximation $\Gamma_Y \ll m_Y$

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 - ----- recasts using publicly available codes in MadAnalysis 5
 - → is there any model-independent conclusion we can make?

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Widest possible reinterpretation potential

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Provide public models and simulated data for (at least) Run 3 studies
 Writing easy-to-use tools to map simplified model parameters to any theory



Database of simulated samples and recast data under construction (not public yet)

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People involved in the analysis

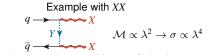
E. Bagnaschi, F. Benoit, A. Cagnotta, A. Desai, B. Fuks, O. Iorio, L. Munoz S. Manohar Dogra, A. Moreno, Y. Sheng (limitations mostly due to accessibility to HPC resources)

Relevance of the different processes

Master equation to reconstruct signal for any flavour hypothesis

$$\begin{split} \sigma_{\text{Tot}}^{e\!f\!f}(M_Y, M_X, \lambda) &= \lambda^0 \ \hat{\sigma}_{Y\bar{Y}_{QCD}}(M_Y) \quad \epsilon_{Y\bar{Y}_{QCD}}(M_Y, M_X) \\ &+ \lambda^4 \ \hat{\sigma}_{YY_t}(M_Y, M_X) \ \epsilon_{YY_t}(M_Y, M_X) \\ &+ \lambda^4 \ \hat{\sigma}_{Y\bar{Y}_t}(M_Y, M_X) \ \epsilon_{Y\bar{Y}_t}(M_Y, M_X) \\ &+ \lambda^2 \ \hat{\sigma}_{Y\bar{Y}_t}(M_Y, M_X) \ \epsilon_{Y\bar{Y}_t}(M_Y, M_X) \\ &+ \lambda^2 \ \hat{\sigma}_{XX}(M_Y, M_X) \ \epsilon_{XX}(M_Y, M_X) \\ &+ \lambda^2 \ \hat{\sigma}_{XY}(M_Y, M_X) \ \epsilon_{XY}(M_Y, M_X) \\ &+ \lambda^2 \ \hat{\sigma}_{XY}(M_Y, M_X) \ \epsilon_{XY}(M_Y, M_X) \end{split}$$

 $\hat{\sigma}$ are the cross-sections after factorizing the new coupling ϵ are the efficiencies associated with a given experimental signal region



The kinematic properties are driven **only** by the masses λ just **rescales** the cross-sections without affecting the shape of distributions

Relevance of the different processes

Master equation to reconstruct signal for any flavour hypothesis

$$\begin{split} & \sigma^{e \mathcal{T}}_{\text{Tot}}(M_Y, M_X, \lambda) = \lambda^0 \; \hat{\sigma}_{Y \bar{Y}_{QCD}}(M_Y) \quad \epsilon_{Y \bar{Y}_{QCD}}(M_Y, M_X) \\ & + \lambda^4 \; \hat{\sigma}_{Y Y_t}(M_Y, M_X) \; \epsilon_{Y Y_t}(M_Y, M_X) \\ & + \lambda^4 \; \hat{\sigma}_{Y \bar{Y}_t}(M_Y, M_X) \; \epsilon_{Y \bar{Y}_t}(M_Y, M_X) \\ & + \lambda^2 \; \hat{\sigma}_{Y \bar{Y}_t}(M_Y, M_X) \; \epsilon_{Y \bar{Y}_t}(M_Y, M_X) \\ & + \lambda^2 \; \hat{\sigma}_{X \bar{Y}_t}(M_Y, M_X) \; \epsilon_{X X}(M_Y, M_X) \\ & + \lambda^2 \; \hat{\sigma}_{X X}(M_Y, M_X) \; \epsilon_{X X}(M_Y, M_X) \\ & + \lambda^2 \; \hat{\sigma}_{X Y}(M_Y, M_X) \; \epsilon_{X Y}(M_Y, M_X) \end{split}$$

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Channel	Scaling	Key features	
XX	$\chi \chi = \lambda^4$ strong dependence on coupling, but requires emission of visible objects		
XY	λ^2	phase-space advantage: potentially competitive	
$Y\bar{Y}_{QCD}$	1	only depends on masses, baseline contribution	
$\begin{array}{c} YY_t \\ Y\bar{Y}_t \\ \bar{Y}\bar{Y}_t \end{array}$	λ^4	enhanced by PDFs for <i>u</i> and <i>d</i> but present only for real DM strong dependence on coupling and interferes with QCD always PDF suppressed and present only for real DM	
$Y\overline{Y}_i$	λ^2	unphysical by itself, but potentially negative contribution	

• up and down \longrightarrow large PDF enhancement for YY_t , unique to these two quarks

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- top \longrightarrow final states with leptons from its decay, limited number of processes: XX (but only at one-loop) and YY_{QCD}

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- top final states with leptons from its decay, limited number of processes: XX (but only at one-loop) and YY_{OCD}

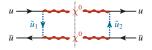
Possibility to combine individual result to describe universal scenarios

 $\mathcal{L} \sim \lambda Y_f X q_f$ with same λ for each q_f

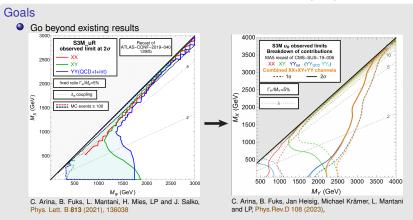
Actually, results can be recombined in almost any way Simulated samples can also be recycled using appropriate weights

Potential to reconstruct complex models with multiple mediators or DM candidates

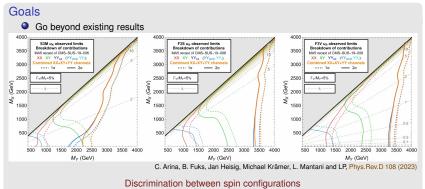
Missing some interference contributions at the moment

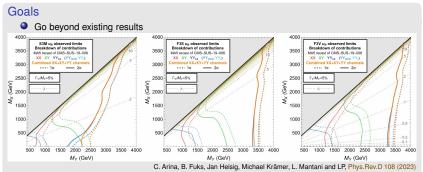


Current results



Combination of all channels, relevance of NLO corrections and interference effects





Discrimination between spin configurations

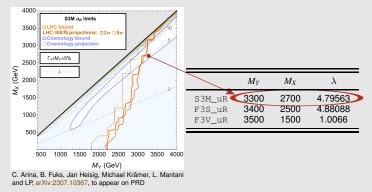
For fixed Γ_Y/m_Y bounds poorly depend on m_X , especially for fermion Y, even if λ decreases.

$$\begin{split} & \text{S3M_uR}: \begin{cases} \mathcal{M}_{uu}^2 \propto \lambda^4 \frac{tu-M_r^2}{(t-M_r^2)^2} \\ \mathcal{M}_{u\bar{u}}^2 \propto \lambda^4 \frac{sM_X^2}{(t-M_X^2)^2} \\ \end{cases} \\ & \text{F3S_uR}: \mathcal{M}_{uu}^2 = \mathcal{M}_{u\bar{u}}^2 \propto \lambda^4 \frac{(t-M_r^2)^2}{(t-M_X^2)^2} \\ & \text{F3V_uR}: \begin{cases} \mathcal{M}_{uu}^2 \propto \lambda^4 \frac{[2M_X^2(t-M_r^2)+M_Y^2(t+2u-3M_r^2)]^2}{M_X^2(t-M_X^2)} \\ \mathcal{M}_{u\bar{u}}^2 \propto \lambda^4 \frac{[2M_X^2(M_Y^2-t)+M_Y^2(s-u+M_r^2)]^2}{M_{u\bar{u}}^2(t-M_Y^2)^2} \end{cases} \end{cases}$$

- **F3S** and **F3V**: the YY_t all amplitudes become independent of m_X for $m_X \rightarrow 0$
- S3M: the uū-initiated amplitude decreases with m_x and leaves only the m_x-independent uu-initiated one for low DM masses

Goals

- Go beyond existing results
- Identify benchmarks allowed by LHC and cosmology observables

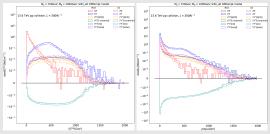


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	M_Y	M_X	λ
S3M_uR	3300	2700	4.79563
F3S_uR	3400	2500	4.88088
F3V_uR	3500	1500	1.0066

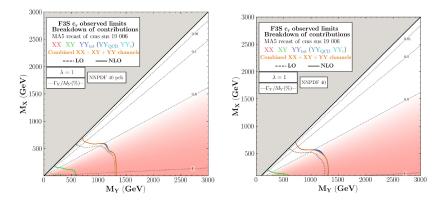
Kinematical studies for subsequent analyses (preliminary, courtesy of A. Desai)



Do we need differential *K*-factors or can we just apply a constant one? Which process dominates in which region, and how to emphasize its significance?

Interaction with the charm quark

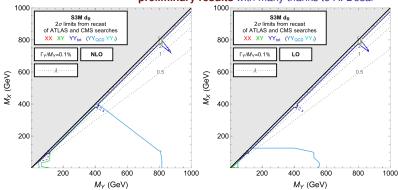
preliminary results courtesy of F. Benoit and L. Munoz



- Exploring difference by using perturbative/intrinsic charm PDFs
- Results in the red area have large mediator width: care must be taken

Interaction with down-type quarks

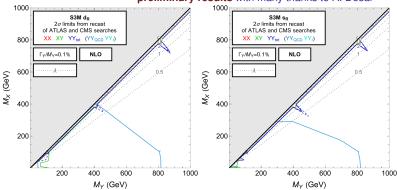
preliminary results with many thanks to A. Desai



- Simulations done with fixed $\Gamma_Y/m_Y = 0.1\%$: small λ overall, baseline QCD pair production bound dominates.
- Large impact of NLO corrections

Interaction with down-type quarks

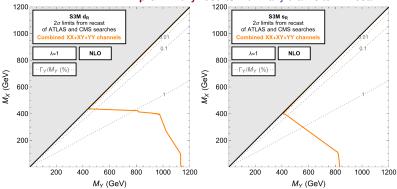
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- Consistent among different flavours (grid to be refined)

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- Large impact of NLO corrections
- Consistent among different flavours (grid to be refined)
- Rescaling for λ = 1: PDF-enhancement for down, increasing constraints



Systematic study of t-channel scenarios and their phenomenology at LHC

Analysis with Run 2 data

focus on kinematically relevant parameters (masses, spins) accurate description of the processes and determination of current bounds recombination of samples to determine bounds for different hypotheses



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Perspectives for Run 3

Production of a database of samples for different *t*-channels processes Publication of numerical models and analysis tools to reproduce and analyse data



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Wishlist

Combine with *s*-channel Include lepto-philic DM scenarios Include flavoured DM scenarios Include interferences for non-minimal scenarios

Further pheno and experimental input will be needed in due course