Beam-beam effects in EIC electron coolers

S. Seletskiy (for C-AD cooling group) Brookhaven National Laboratory

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Beam-beam in Electron Ion Collider (EIC)

- There is the **beam-beam effect from the colliding beams**
- Also, EIC employs cooling of hadrons, both at the injection energy (pre-cooler) and at the top energy (high energy cooler).
- In the coolers bunches of electrons co-travel with hadrons. The focusing effect from electrons on hadrons (and vice-versa for a proposed ring cooler) is non-negligible.
- There is the **beam-beam effect in the pre-cooler**
- There is the **beam-beam effect in the high energy cooler**

Why does Electron Ion Collider need coolers?

• We need to achieve and maintain optimal hadron emittances at the collision energy

Table 1 Main parameters for the electron-proton operation in high divergence operation mode.

- We must cool a hadron bunch at injection energy ($\gamma \approx 25$) to obtain the required emittances
- We must cool hadrons at top energy to counteract an IBS-driven heating $\frac{3}{3}$

EIC layout

- Two coolers are needed:
	- The pre-cooler (injection energy cooler) – a single pass electron cooler
	- The high energy cooler one of the possible options for it is to use a ring electron cooler (multi -pass electron cooler)
- Both coolers must share the same cooling section ($L \approx 170$ m)

How electron cooling works

- Mix a hadron bunch with an electron bunch traveling with the same velocity and let the two bunches travel together over some length (for us $L = 170$ m). You will notice that the emittances of hadrons get reduced. Why?
- In the co-moving frame, the mixture of two bunches looks like a mixture of two gases – a gas of ions and a gas of electrons
- Electrons are much lighter; hence, the gas of electrons is much colder than the gas of ions. Heat transfer \equiv Electron Cooling

Beam-beam effect in electron coolers

• Consider two co-traveling bunches of type 1 and type 2 particles. Assume that the bunches are oppositely charged and assume that both bunches are circularly symmetric. Then the space charge effect from bunch 1 on a particle of bunch 2 is:

$$
x'' = -\frac{2}{\gamma^3} \cdot \frac{I_{1(slice)}}{I_A} \cdot Z_2 \cdot \frac{m_e}{m_2} \cdot \frac{\left(1 - e^{-\frac{x^2 + y^2}{2\sigma_1^2}}\right)}{x^2 + y^2} x
$$

• In linear approximation:

$$
x^{\prime\prime}=-\frac{1}{\gamma^3}\cdot\frac{I_{1(slice)}}{I_A}\cdot Z_2\cdot\frac{m_e}{m_2}\cdot\frac{x}{\sigma_1^2}
$$

• The maximum tune shift:

$$
\max(\Delta \nu_{12}) = \frac{1}{4\pi\gamma^3} \cdot \frac{I_{1(peak)}}{I_A} \cdot Z_2 \cdot \frac{m_e}{m_2} \cdot \frac{L_{CS}\langle \beta_{2(CS)} \rangle}{\sigma_1^2} \qquad 0 \qquad \frac{5}{-0.2} \qquad -0.1 \qquad 0.0 \qquad 0.1 \qquad 0.2
$$

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15

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[A]

- The bunches are co-traveling; hence, the space charge effect is reduced by γ^2 , yet the cooling section length L_{CS} is 170 m.
- On a single pass through the CS each particle interacts with a slice of the other bunch. A nonuniform longitudinal distribution will cause periodic modulation of the turn-by-turn focusing kick.

 \mathcal{X}

 \mathcal{Y}

 $\frac{1}{2}$ 1(slice) - bunch 1

Beam-beam effect in EIC pre-cooler

 $\Delta v_{ep} = 1.3 \cdot 10^{-3}$

- The Ring Electron Cooler (REC) must counteract IBS-driven heating of protons at the collision energy ($\gamma =$ 293). The required cooling times for 275 GeV protons are 2 hrs horizontal and 3 hrs longitudinal.
- The REC reutilizes electron bunches for several million turns. Damping wigglers are used in the REC to counteract electrons' emittance growth from IBS, BBS and quantum excitations.
- **Since there are two rings to consider (HSR and REC), both the electron-proton and proton-electron beambeam effects become important.**

Beam-beam effects in REC

 $\xi_p = 12 \cdot 10^{-3}$

$$
\Delta\nu_{pe(x,y)} = \frac{I_p \beta_{eCS(x,y)} L_{CS}}{2\pi I_a \gamma^3 \sigma_{p(x,y)} (\sigma_{px} + \sigma_{py})} \longrightarrow \boxed{\frac{\Delta v_{pe(x)} = 0.04}{\Delta v_{pe(y)} = 0.1}}
$$

 $\Delta\bm{\nu_{ep}}_{(\bm{\mathcal{y}})}=\bm{1}.\,\bm{7}\cdot\bm{10}^{-4}$

Tracking electrons in REC

• Considering the central slice of a p-bunch (max I_p):

$$
\delta \theta_x = \frac{K_p}{2} L_{CS} x I_x; \quad \delta \theta_y = \frac{K_p}{2} L_{CS} y I_y; \quad K_p = \frac{2I_p}{I_A \gamma^3 \beta^3}
$$
(1)

$$
I_x = \int_0^\infty \frac{\exp\left(-\frac{x^2}{2(\sigma_x^2 + q)} - \frac{y^2}{2(\sigma_y^2 + q)}\right)}{(\sigma_x^2 + q)^{\frac{3}{2}} (\sigma_y^2 + q)^{\frac{1}{2}}} dq; \quad I_y = \int_0^\infty \frac{\exp\left(-\frac{x^2}{2(\sigma_x^2 + q)} - \frac{y^2}{2(\sigma_y^2 + q)}\right)}{(\sigma_x^2 + q)^{\frac{1}{2}} (\sigma_y^2 + q)^{\frac{3}{2}}} dq
$$

$$
\begin{pmatrix} x_{n+1} \\ x'_{n+1} \end{pmatrix} = M_x \begin{pmatrix} x_n \\ x'_n + \delta \theta_x (x_n, y_n) \end{pmatrix} \qquad \qquad \begin{pmatrix} y_{n+1} \\ y'_{n+1} \end{pmatrix} = M_y \begin{pmatrix} y_n \\ y'_n + \delta \theta_y (x_n, y_n) \end{pmatrix}
$$

REC working point choice (I) We started with $v_x = 0.23$, $v_y = 0.21$

tracking 100k particles

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REC working point choice (II) *tracking 100k particles*

Adjusting the working point to $v_x = 0$. 13, $v_y = 0$. 09, allowed us to keep good emittance, and we got a halo under control.

 x/σ_{x}

Relative trajectory stability (I)

• To estimate requirements to a stability of e- and p- trajectories in the CS with respect to each other we simulate kicks in Eq. (1) as:

$$
\delta\theta_x = \frac{K_p}{2} L_{CS}(x + \Delta x)I_x; \quad \delta\theta_y = \frac{K_p}{2} L_{CS}(y + \Delta y)I_y
$$

where displacements Δx and Δy are varying on each turn and are randomly distributed with σ_x , σ_y

• A quick estimate (overestimate) for the resulting emittance growth can be obtained from a "random walk-emittance dilution" model assuming a linear p-e kick:

$$
\varepsilon = \varepsilon_0 + \frac{\pi}{8} \beta_{e(CS)} \sigma_{\theta(noise)}^2 \cdot N_{turns}
$$

$$
\sigma_{\theta(noise)} = \frac{4\pi \Delta v_{pe}}{\beta_{eCS}} \sigma_{(noise)}
$$

Relative trajectory stability (II)

An analytic formula (with a fudge-factor n=0.4) is a good match for heating times obtained from tracking

$$
\sigma_{(noise)} = \sqrt{\frac{\varepsilon_e \beta_e T_{rev}}{2\pi^3 \Delta v_{pe}^2 \cdot \tau \cdot n}}
$$

 n accounts for the difference between an actual non-linear space charge and a model linear SC and depends on parameters of both bunches.

If we want the noise driven heating time (τ) to be an order of magnitude smaller than the transverse damping time, then the requirement to noise is $\sigma_{(noise)} \leq$ $5 \mu m$

Conclusion

- It is proposed to use electron coolers at EIC both at the injection and at the top energy to achieve and maintain the required hadron parameters.
- In electron coolers the electron and hadron beams are co-traveling in the common section with the same velocity. Although the beam-beam effect is suppressed by γ^2 , the cooling section can be long, and the resulting effect can be substantial.
- We started considering the effect of proton focusing on electron's dynamics in the EIC Ring Electron Cooler
	- A "single slice" p-e focusing was included into simulations
	- Requirements to stability of beam trajectories with respect to each other were derived
	- We are working on including longitudinal dynamics into the simulations
	- Studying of beam-beam in combination with the self-SC, IBS and radiation damping effects is planned