



McGill

# Improving our understanding of $B \rightarrow D\pi\ell\nu$ & $B \rightarrow \pi\pi\ell\nu$ decays

**Raynette van Tonder**

raynette.vantonder@mail.mcgill.ca

**In collaboration with:**

E. Gustafson, F. Herren, B. Kubis,  
R. Van de Water & M. Wagman





# $B \rightarrow X\ell\nu$ modelling & composition

---

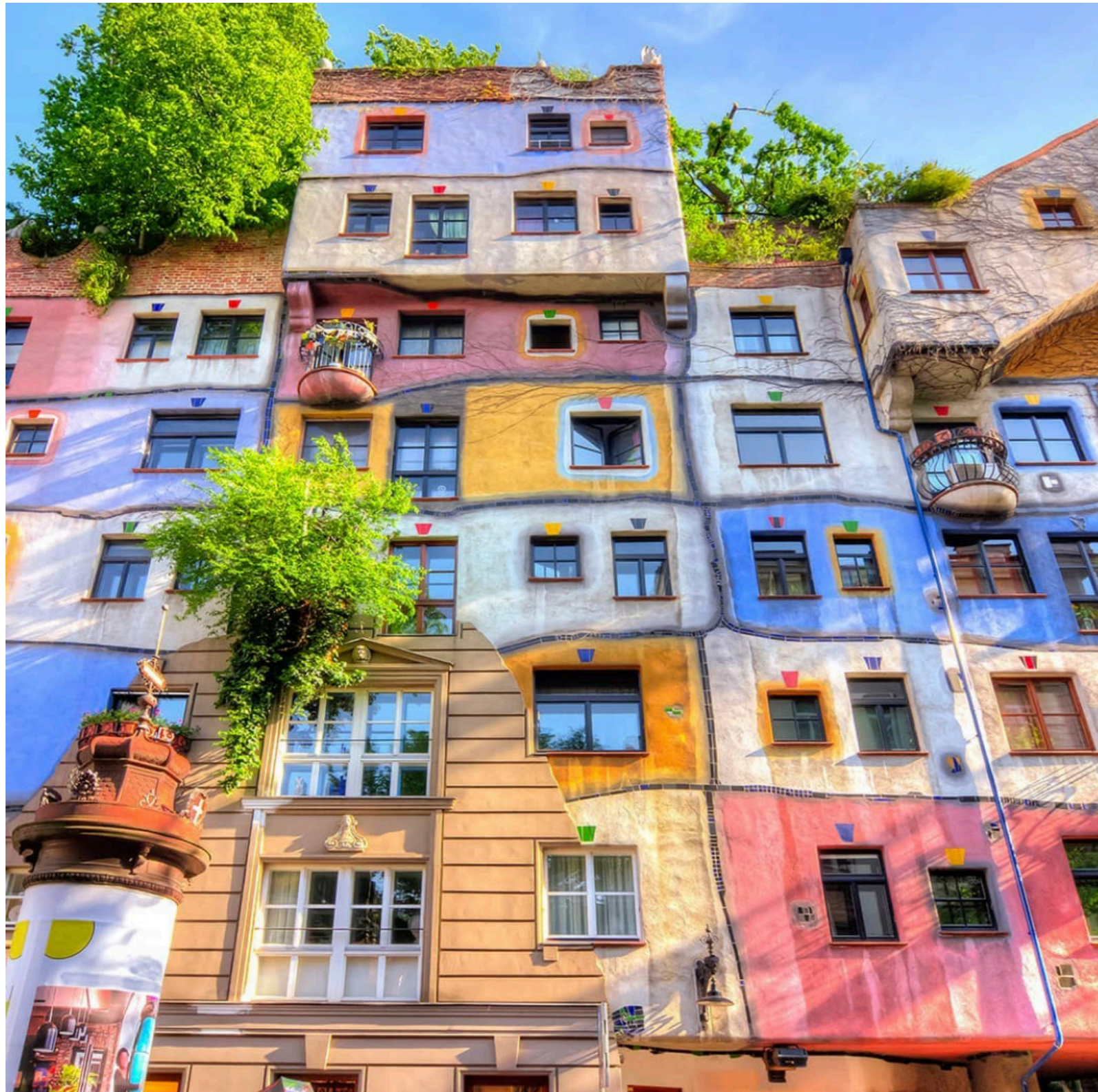


$B \rightarrow X\ell\nu$  modelling is like the Hundertwasserhaus:



# $B \rightarrow X_{\ell v}$ modelling & composition

---



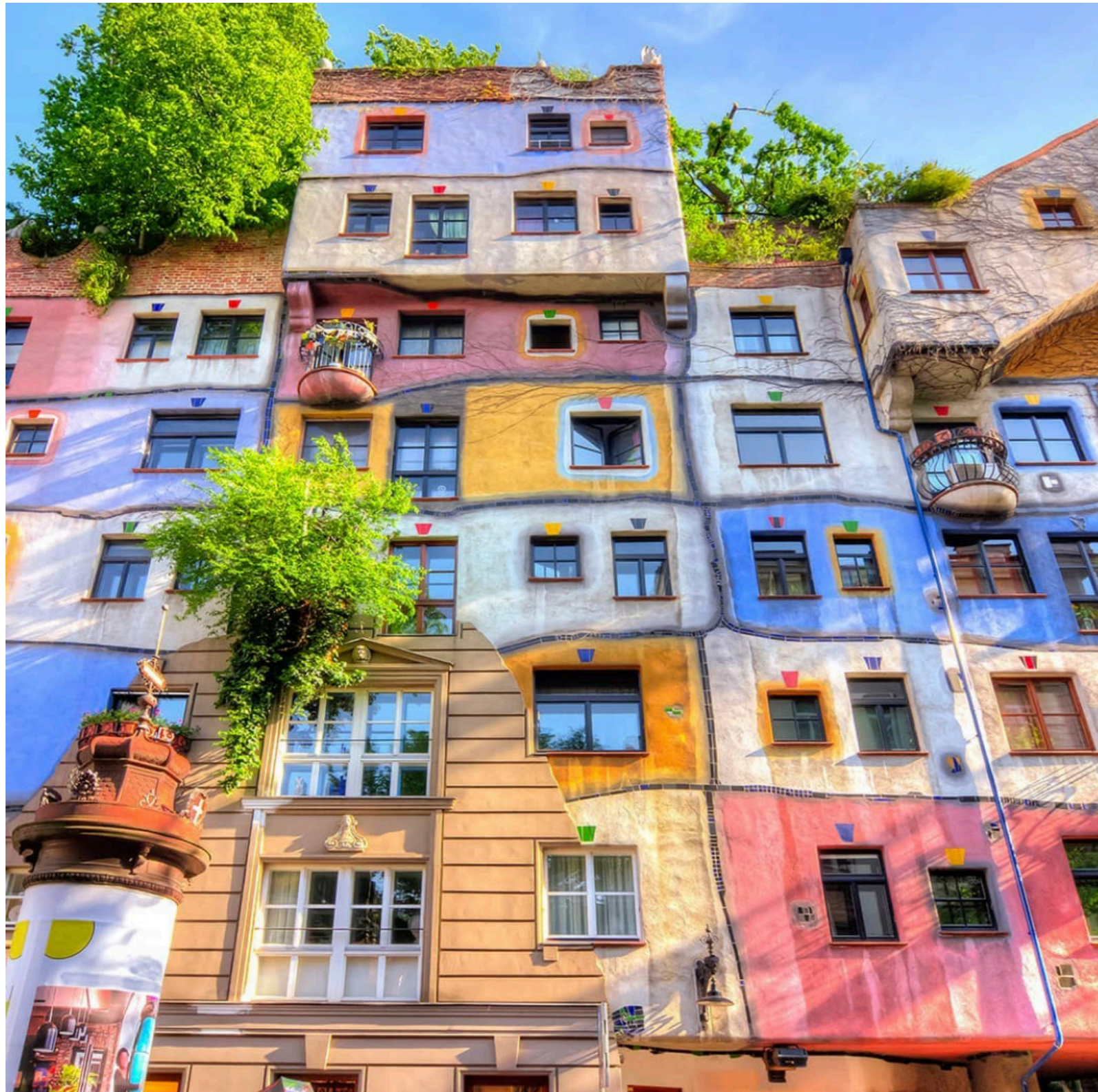
$B \rightarrow X_{\ell v}$  modelling is like the Hundertwasserhaus:

Each individual process is an important building block with its own characteristic shape and style...



# $B \rightarrow X_{\ell v}$ modelling & composition

---



$B \rightarrow X_{\ell v}$  modelling is like the Hundertwasserhaus:

Each individual process is an important building block with its own characteristic shape and style...

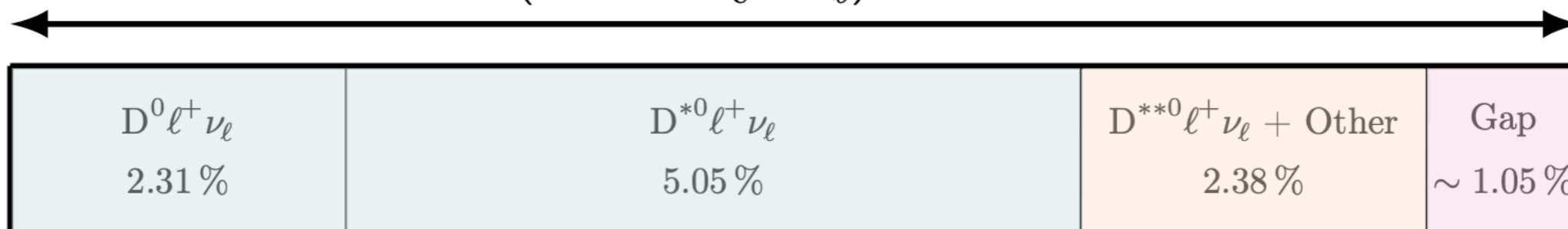
Together they form a wonky, yet beautiful, piece of architecture.



# Some blocks are still missing...

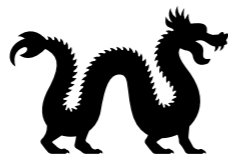
A **leading systematic** for many analyses (not just semileptonic):

$$\mathcal{B}(B^+ \rightarrow X_c^0 \ell^+ \nu_\ell) \approx 10.79\%$$



...or is it even bigger?

Decay	$\mathcal{B}(B^+)$	$\mathcal{B}(B^0)$
$B \rightarrow D \ell^+ \nu_\ell$	$(2.4 \pm 0.1) \times 10^{-2}$	$(2.2 \pm 0.1) \times 10^{-2}$
$B \rightarrow D^* \ell^+ \nu_\ell$	$(5.5 \pm 0.1) \times 10^{-2}$	$(5.1 \pm 0.1) \times 10^{-2}$
$B \rightarrow D_1 \ell^+ \nu_\ell$	$(6.6 \pm 0.1) \times 10^{-3}$	$(6.2 \pm 0.1) \times 10^{-3}$
$B \rightarrow D_2^* \ell^+ \nu_\ell$	$(2.9 \pm 0.3) \times 10^{-3}$	$(2.7 \pm 0.3) \times 10^{-3}$
$B \rightarrow D_0^* \ell^+ \nu_\ell$	$(4.2 \pm 0.8) \times 10^{-3}$	$(3.9 \pm 0.7) \times 10^{-3}$
$B \rightarrow D_1' \ell^+ \nu_\ell$	$(4.2 \pm 0.9) \times 10^{-3}$	$(3.9 \pm 0.8) \times 10^{-3}$
$B \rightarrow D \pi \pi \ell^+ \nu_\ell$	$(0.6 \pm 0.9) \times 10^{-3}$	$(0.6 \pm 0.9) \times 10^{-3}$
$B \rightarrow D^* \pi \pi \ell^+ \nu_\ell$	$(2.2 \pm 1.0) \times 10^{-3}$	$(2.0 \pm 1.0) \times 10^{-3}$
$B \rightarrow X_c \ell \nu_\ell$	$(10.8 \pm 0.4) \times 10^{-2}$	$(10.1 \pm 0.4) \times 10^{-2}$



Fairly well known.



Broad states based on 3 measurements. (BaBar, Belle, DELPHI)



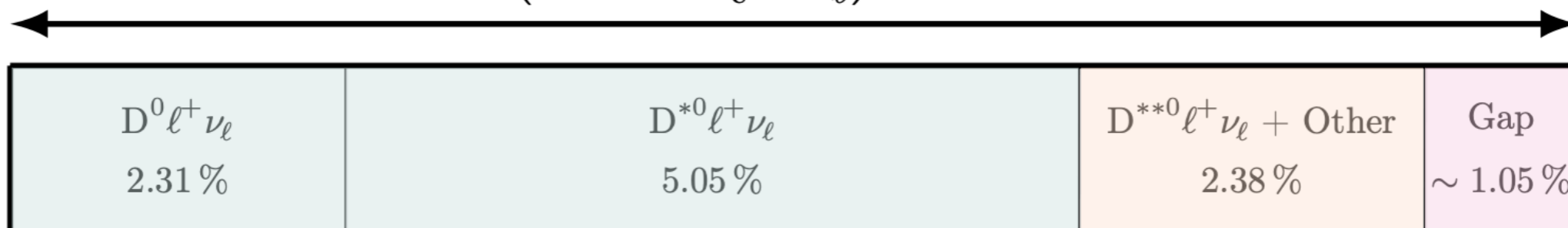
Some hints from BaBar & recent Belle result.



# Some blocks are still missing...

A **leading systematic** for many analyses (not just semileptonic):

$$\mathcal{B}(B^+ \rightarrow X_c^0 \ell^+ \nu_\ell) \approx 10.79\%$$



...or is it even bigger?

Decay	$\mathcal{B}(B^+)$	$\mathcal{B}(B^0)$
$B \rightarrow D \ell^+ \nu_\ell$	$(2.4 \pm 0.1) \times 10^{-2}$	$(2.2 \pm 0.1) \times 10^{-2}$
$B \rightarrow D^* \ell^+ \nu_\ell$	$(5.5 \pm 0.1) \times 10^{-2}$	$(5.1 \pm 0.1) \times 10^{-2}$
$B \rightarrow D_1 \ell^+ \nu_\ell$	$(6.6 \pm 0.1) \times 10^{-3}$	$(6.2 \pm 0.1) \times 10^{-3}$
$B \rightarrow D_2^* \ell^+ \nu_\ell$	$(2.9 \pm 0.3) \times 10^{-3}$	$(2.7 \pm 0.3) \times 10^{-3}$
$B \rightarrow D_0^* \ell^+ \nu_\ell$	$(4.2 \pm 0.8) \times 10^{-3}$	$(3.9 \pm 0.7) \times 10^{-3}$
$B \rightarrow D_1' \ell^+ \nu_\ell$	$(4.2 \pm 0.9) \times 10^{-3}$	$(3.9 \pm 0.8) \times 10^{-3}$
$B \rightarrow D \pi \pi \ell^+ \nu_\ell$	$(0.6 \pm 0.9) \times 10^{-3}$	$(0.6 \pm 0.9) \times 10^{-3}$
$B \rightarrow D^* \pi \pi \ell^+ \nu_\ell$	$(2.2 \pm 1.0) \times 10^{-3}$	$(2.0 \pm 1.0) \times 10^{-3}$
$B \rightarrow D \eta \ell^+ \nu_\ell$	$(4.0 \pm 4.0) \times 10^{-3}$	$(4.0 \pm 4.0) \times 10^{-3}$
$B \rightarrow D^* \eta \ell^+ \nu_\ell$	$(4.0 \pm 4.0) \times 10^{-3}$	$(4.0 \pm 4.0) \times 10^{-3}$
$B \rightarrow X_c \ell \nu_\ell$	$(10.8 \pm 0.4) \times 10^{-2}$	$(10.1 \pm 0.4) \times 10^{-2}$

Fairly well known.

Broad states based on 3 measurements. (BaBar, Belle, DELPHI)

Some hints from BaBar & recent Belle result.

Fill the gap with current "best guess".



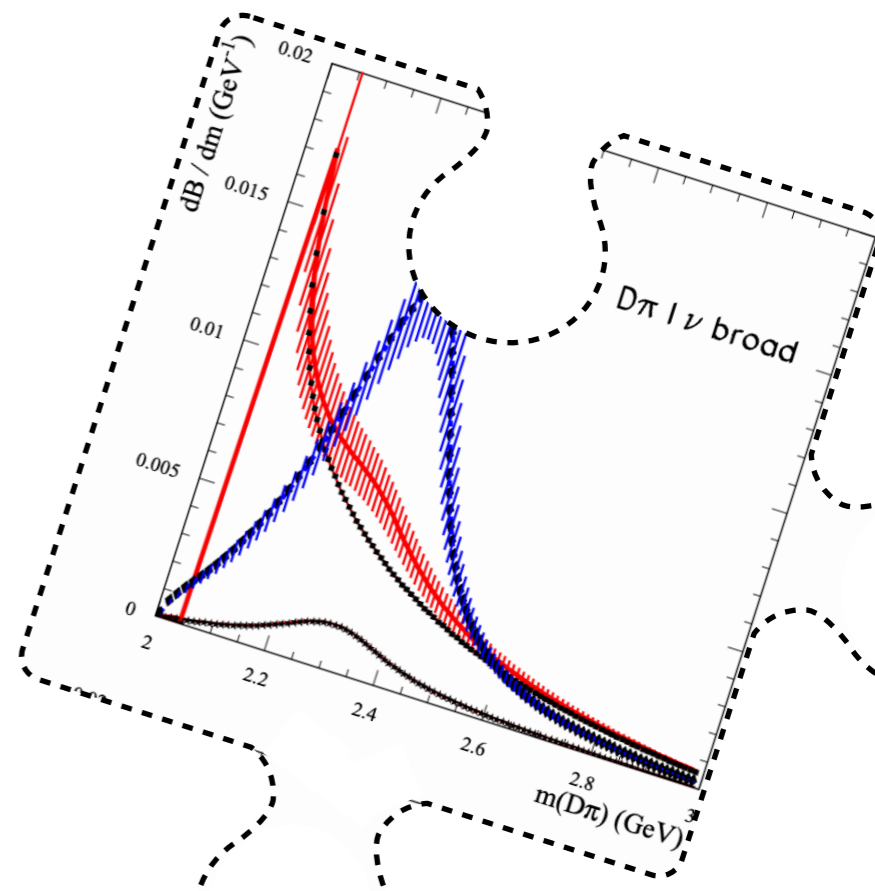
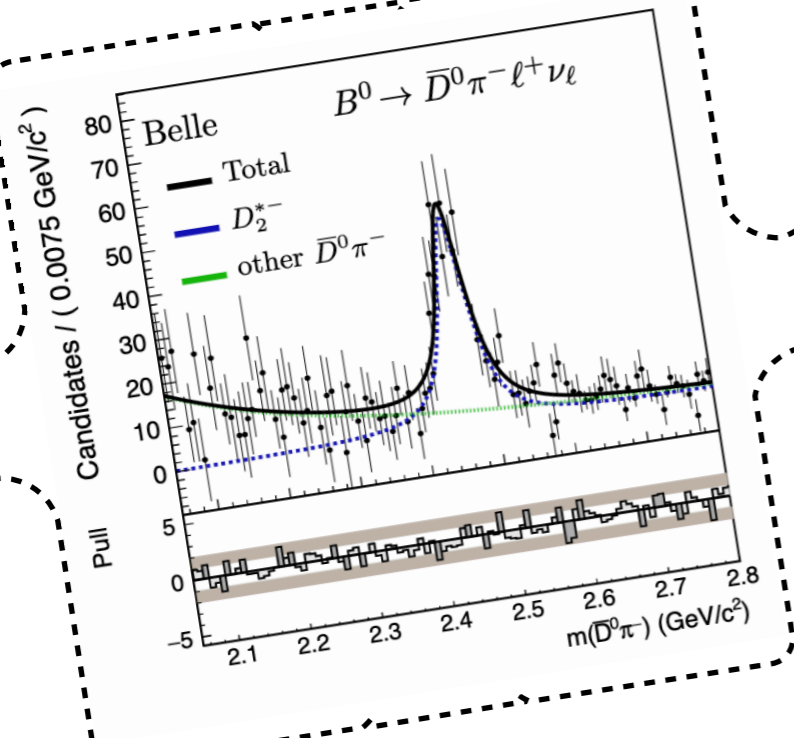
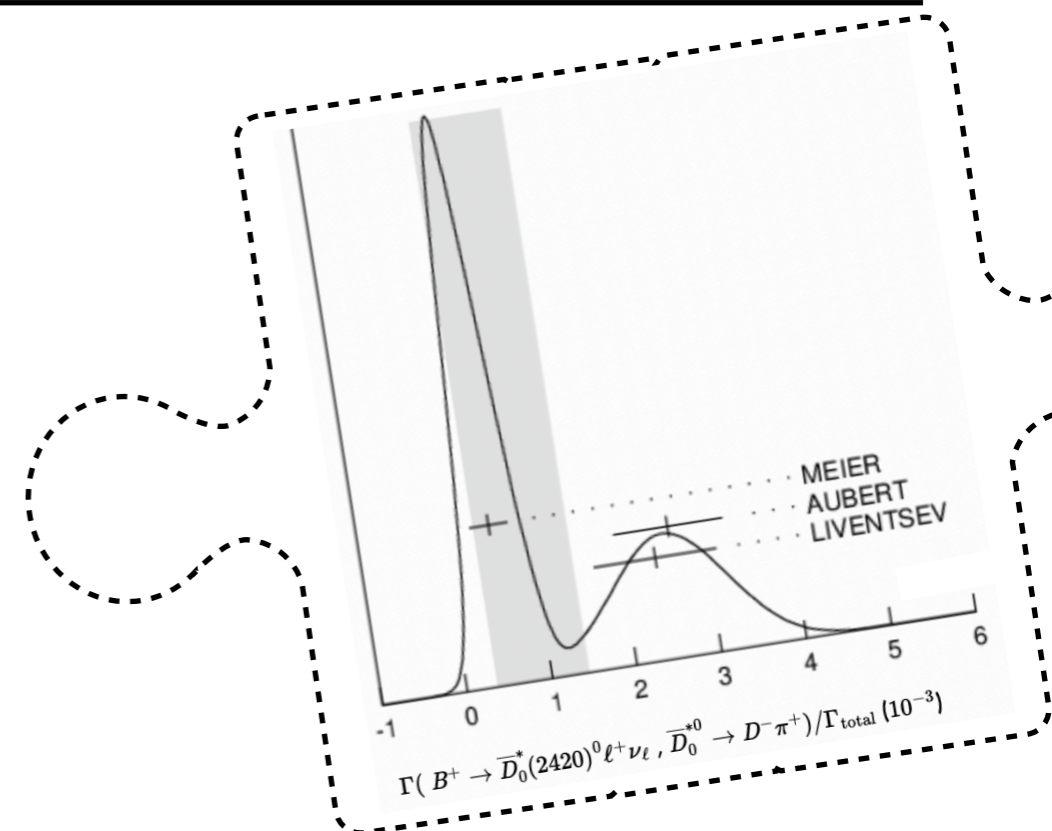
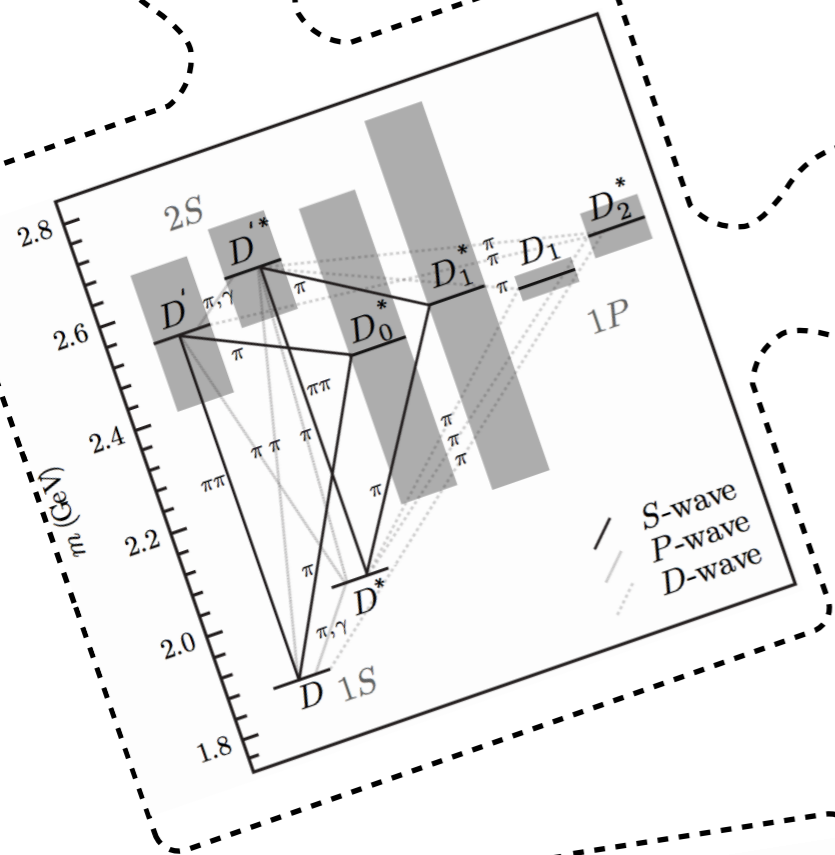




$B \rightarrow D\pi\ell v$

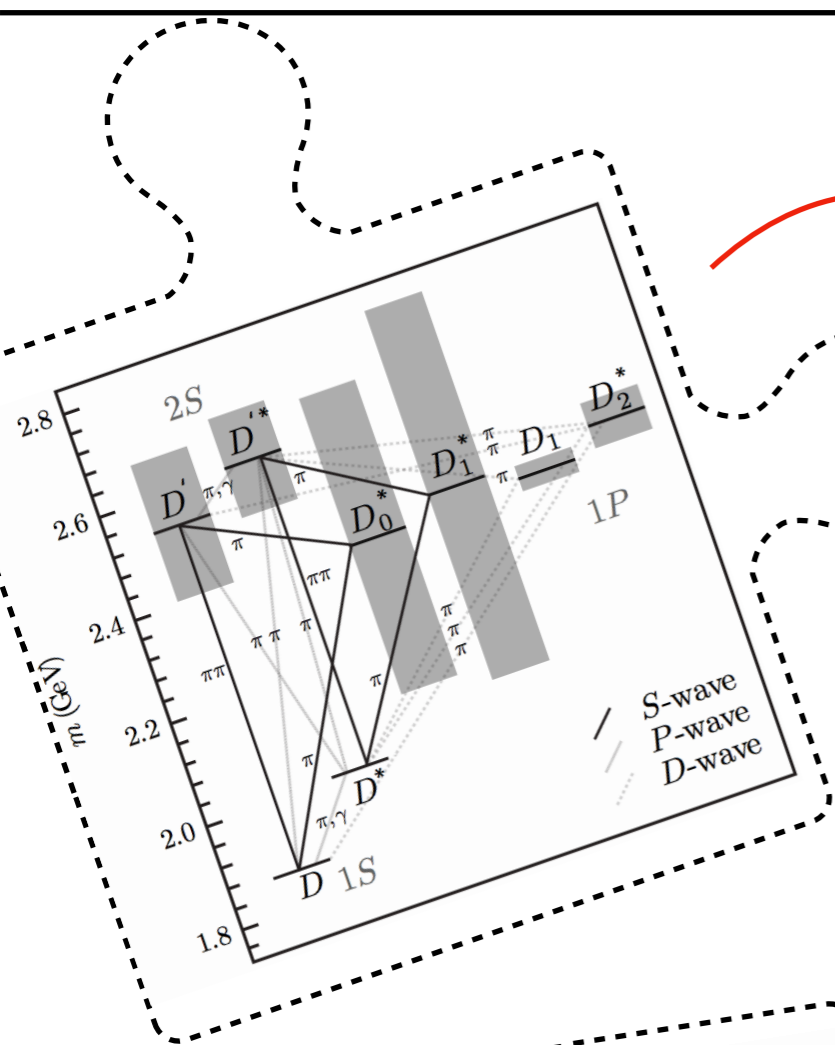


# So what do we know about $B \rightarrow D\pi\ell\nu$ ?

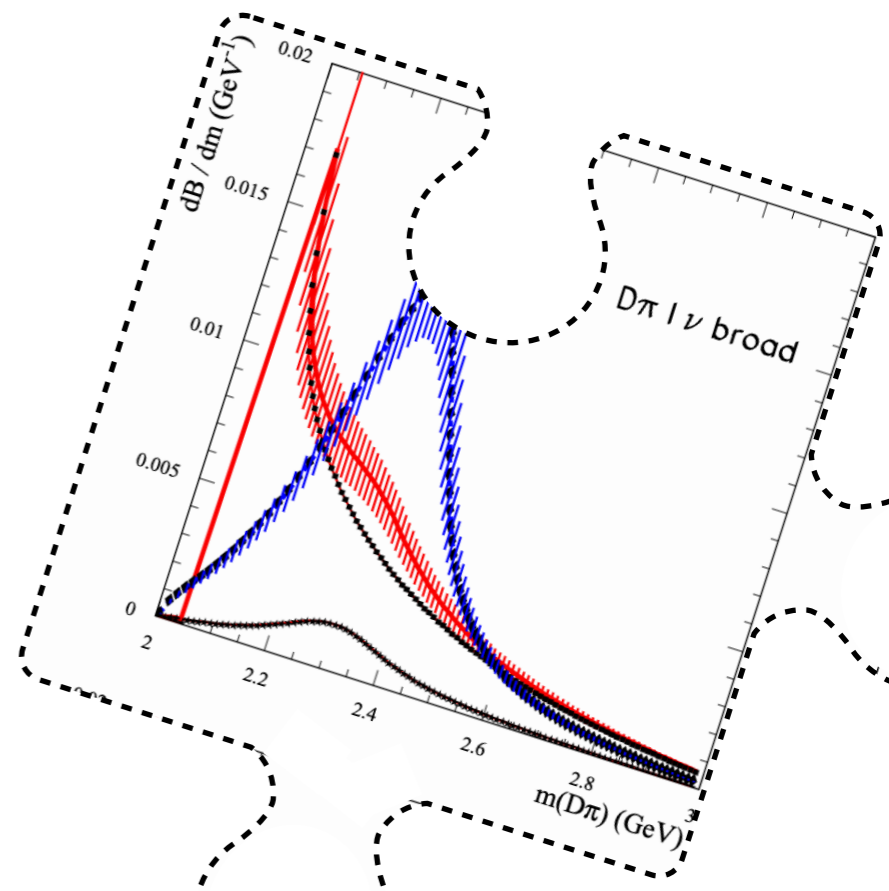
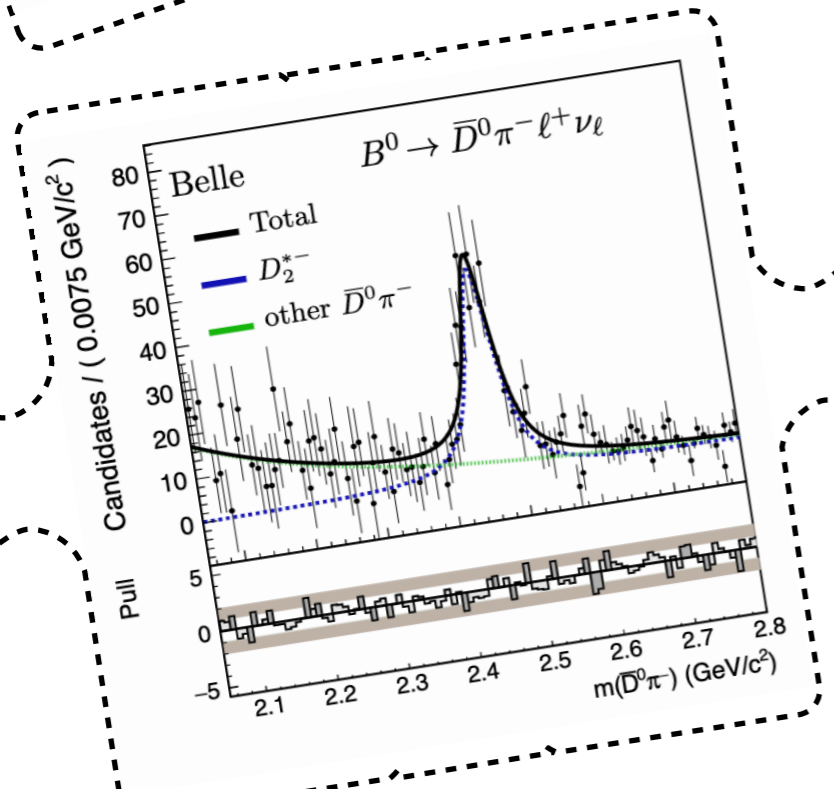
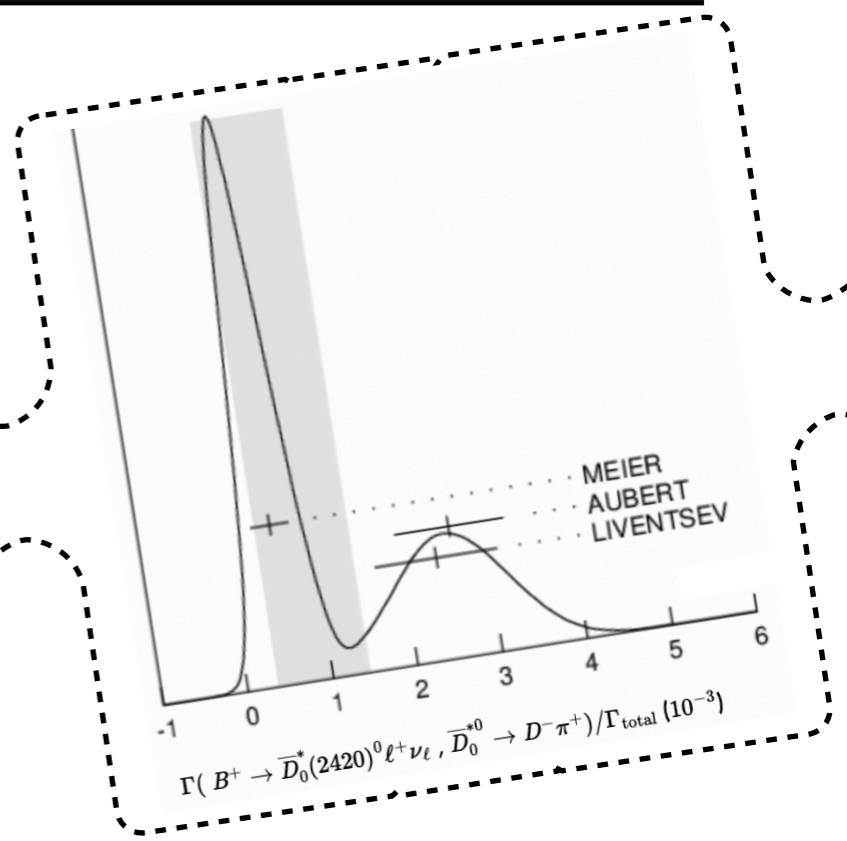




# So what do we know about $B \rightarrow D\pi\ell\nu$ ?

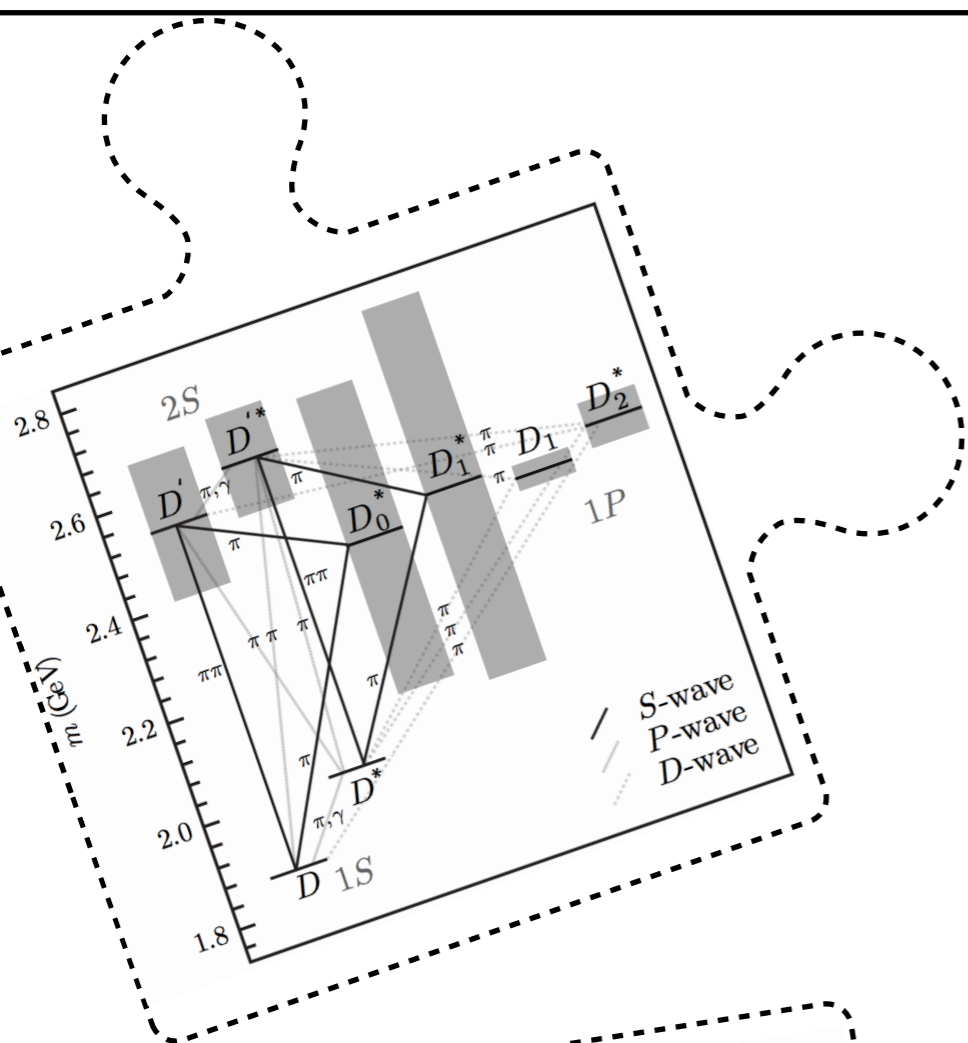


Decay via  $D^*$ ,  $D_2^*$  and  $D_0^*$

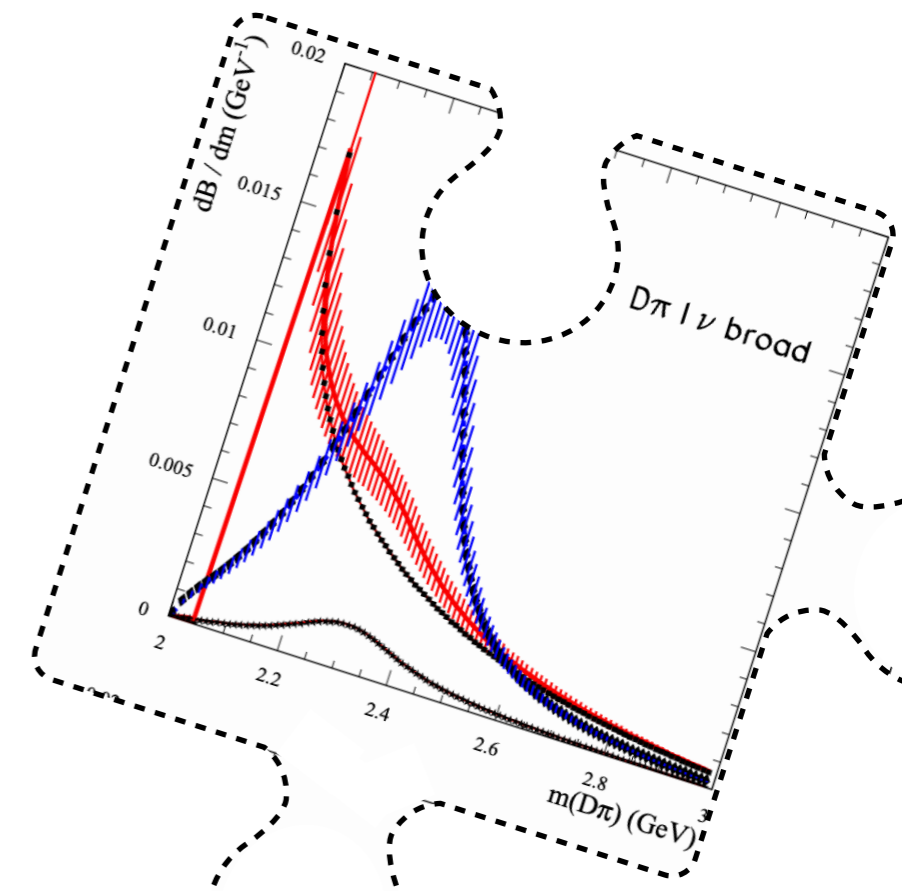
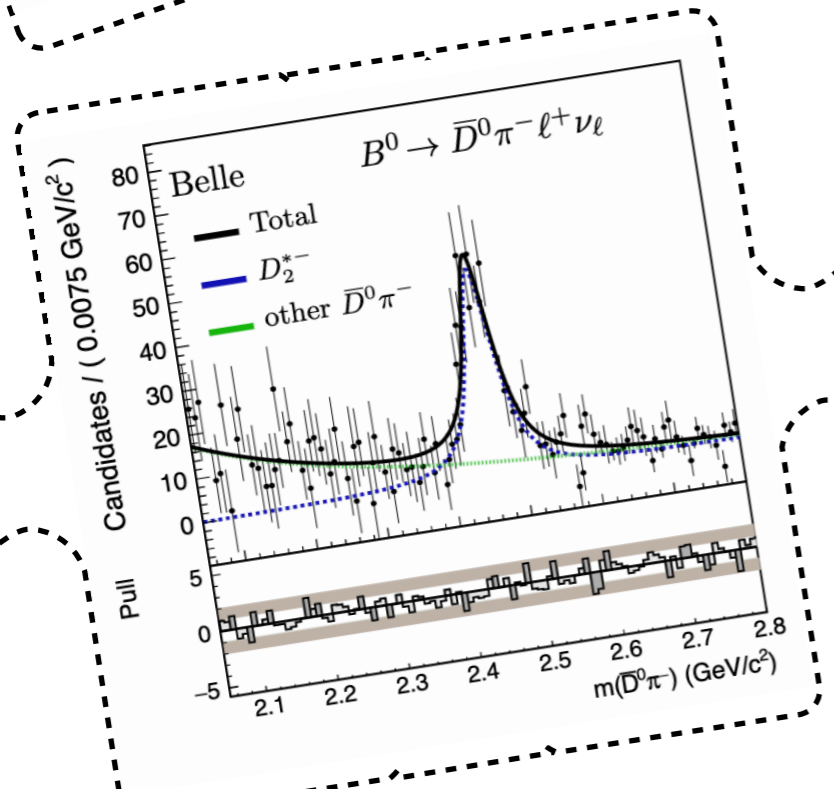
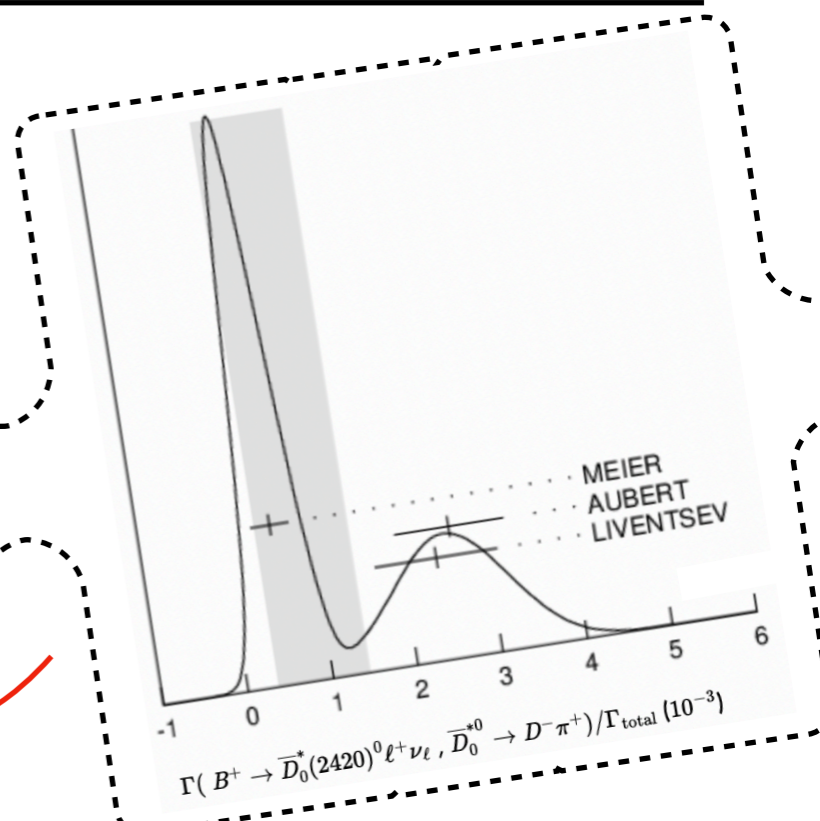




# So what do we know about $B \rightarrow D\pi\ell\nu$ ?

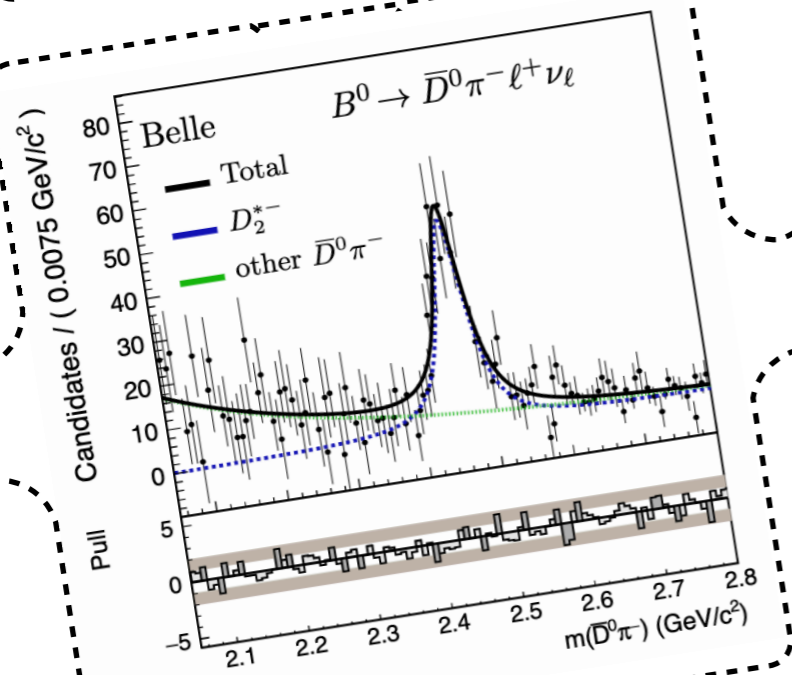
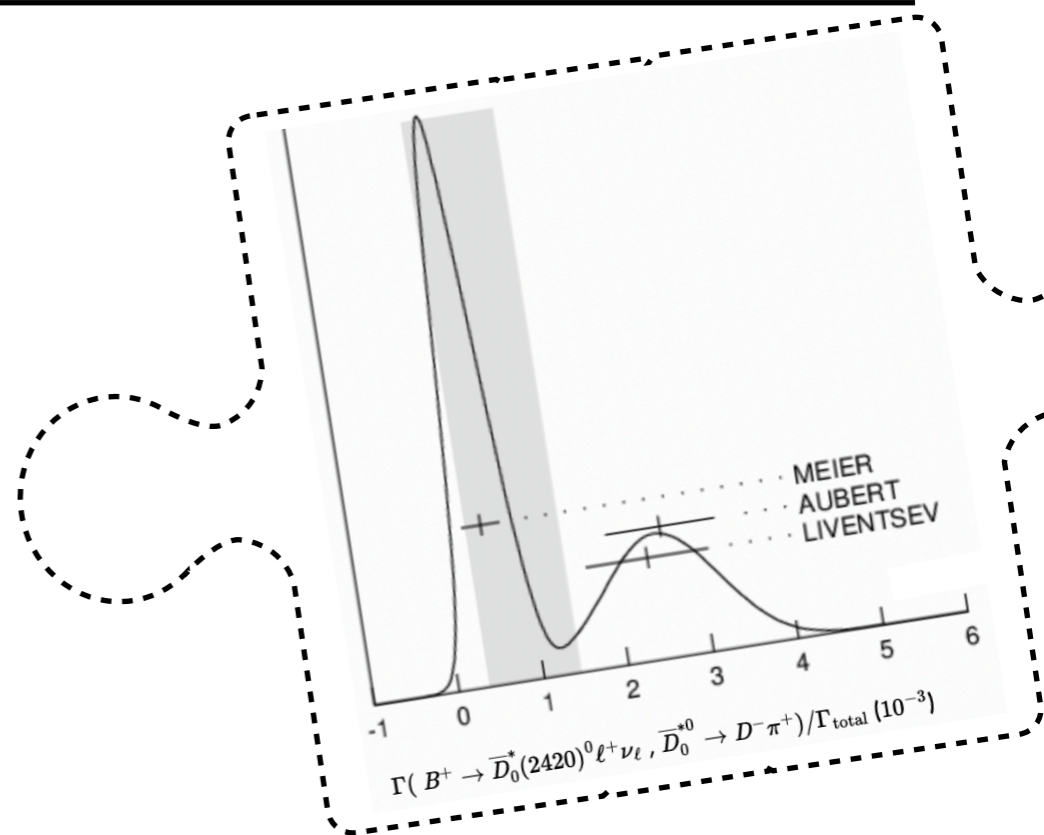
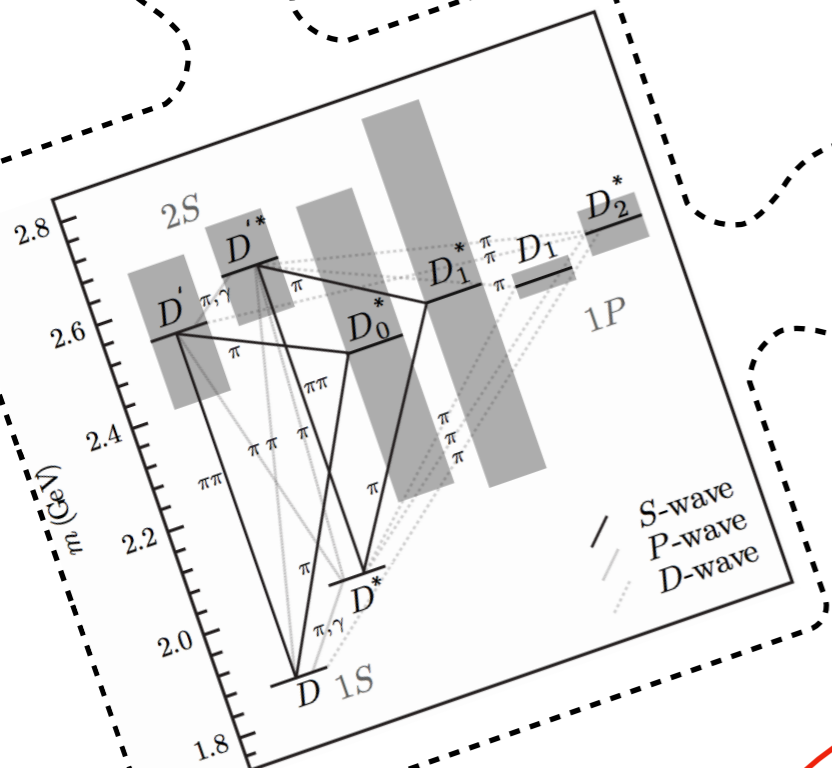


Lower  $B \rightarrow D_0^*$  BF from Belle

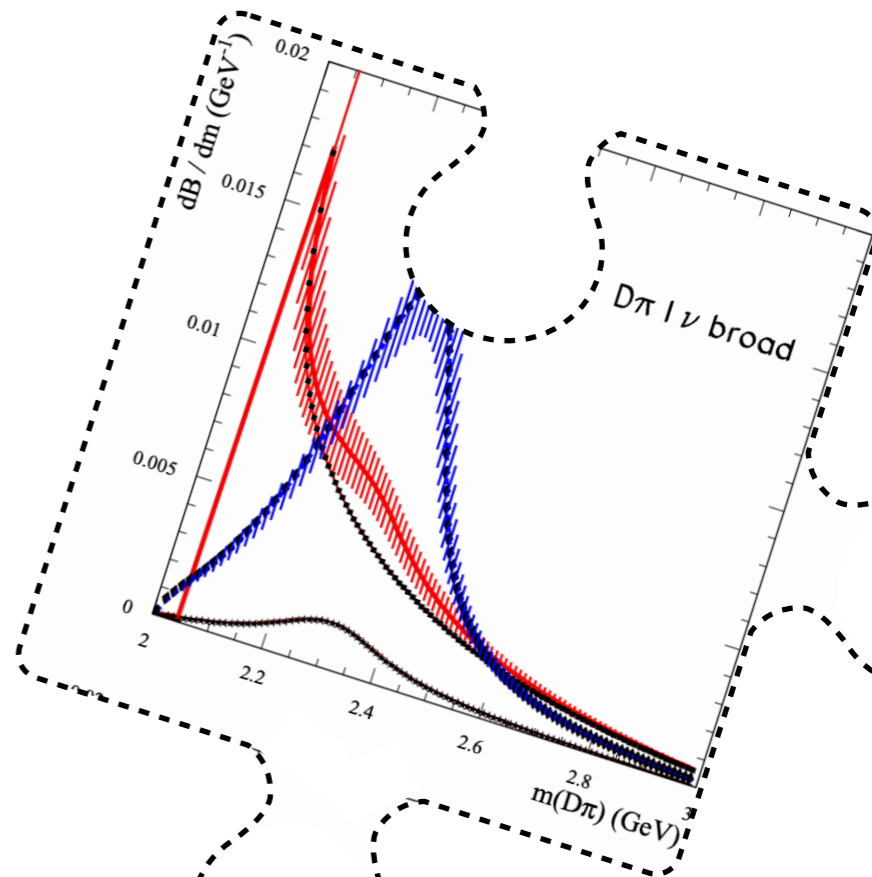




# So what do we know about $B \rightarrow D\pi\ell\nu$ ?

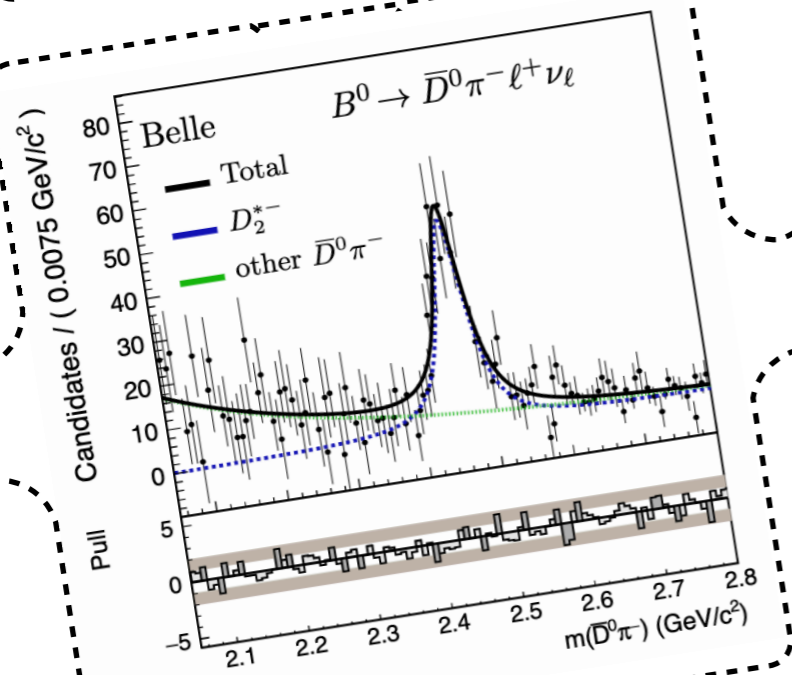
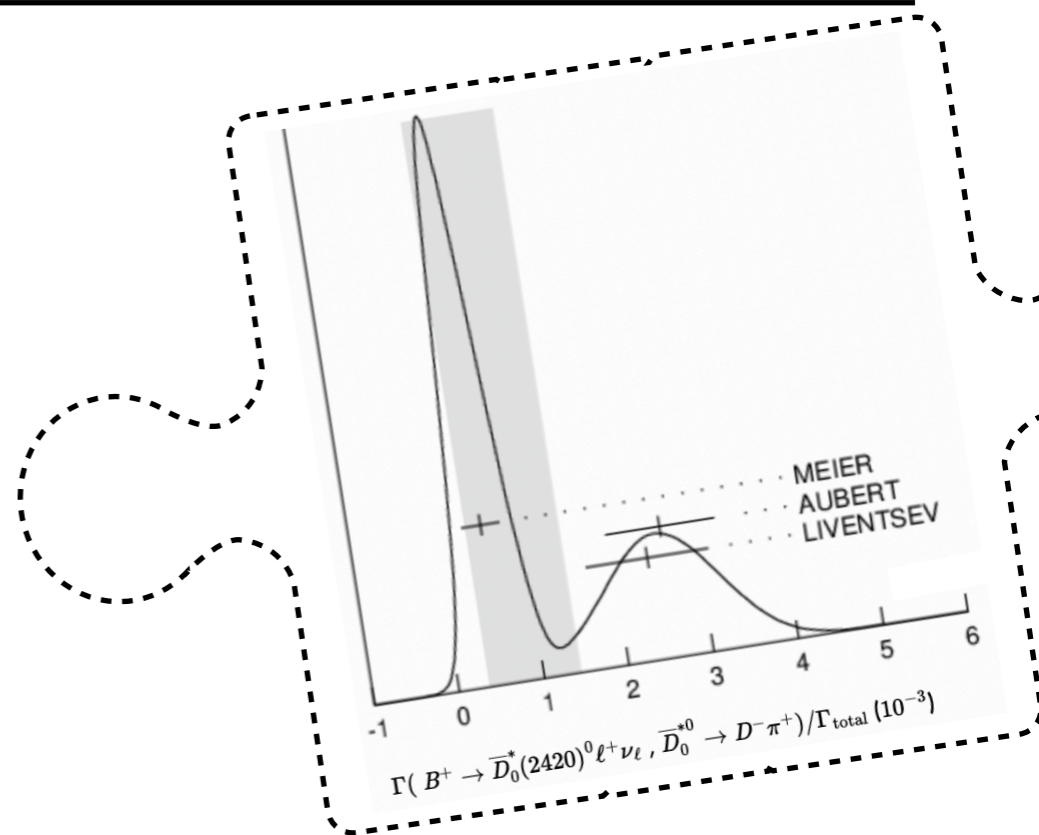
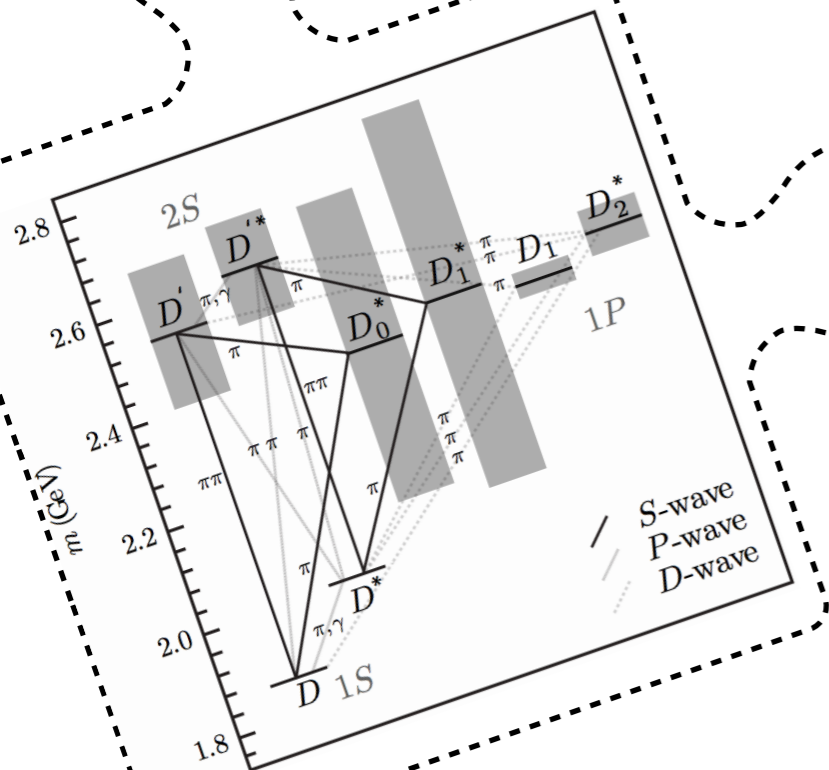


Belle reports large falling component

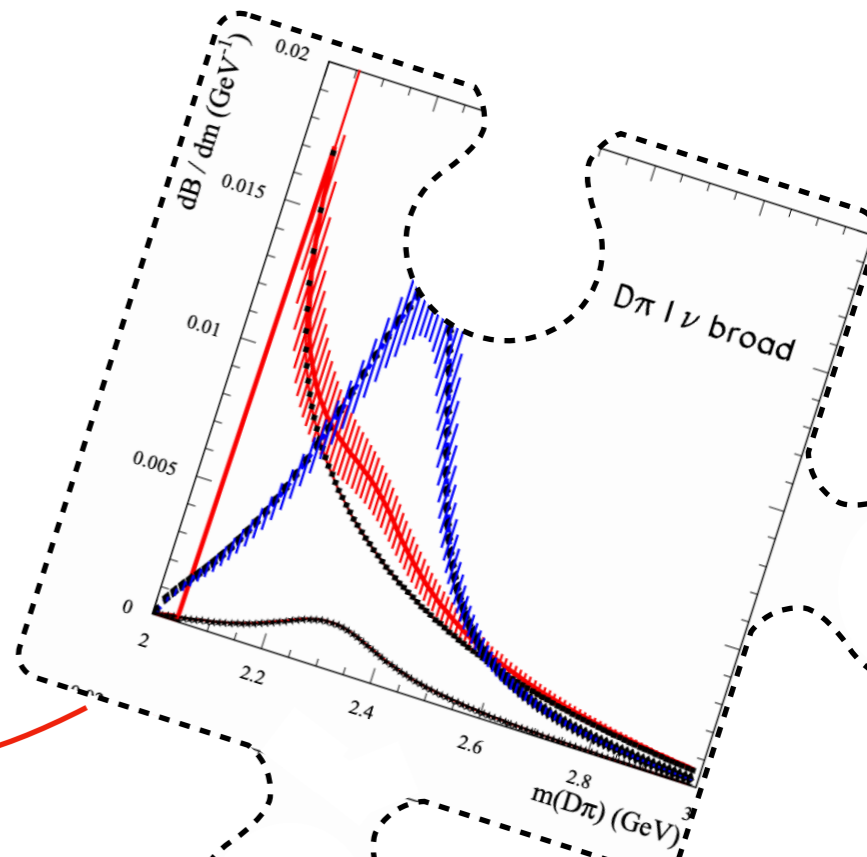




# So what do we know about $B \rightarrow D\pi\ell\nu$ ?

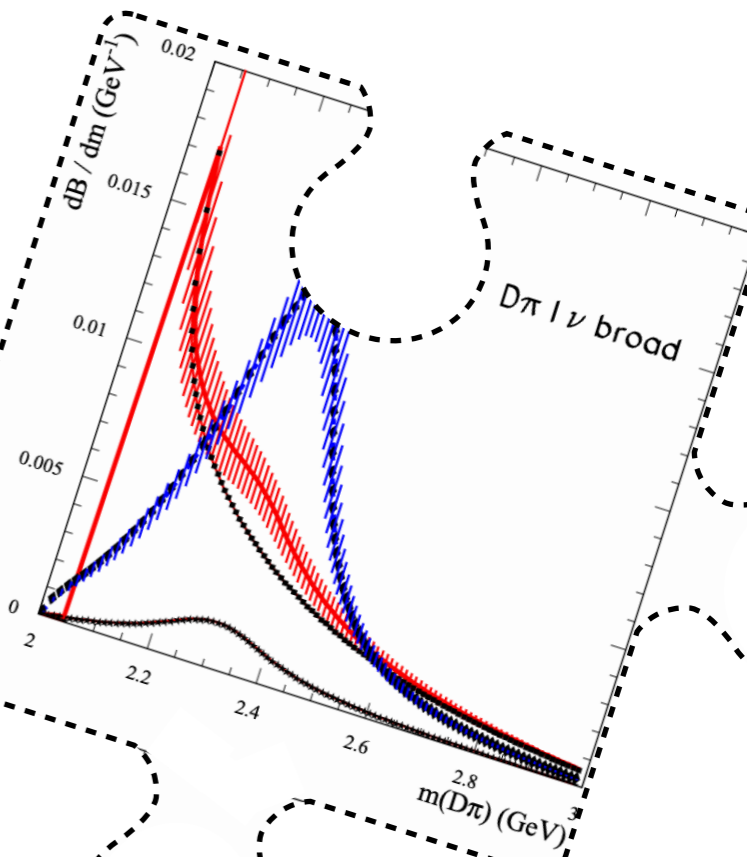
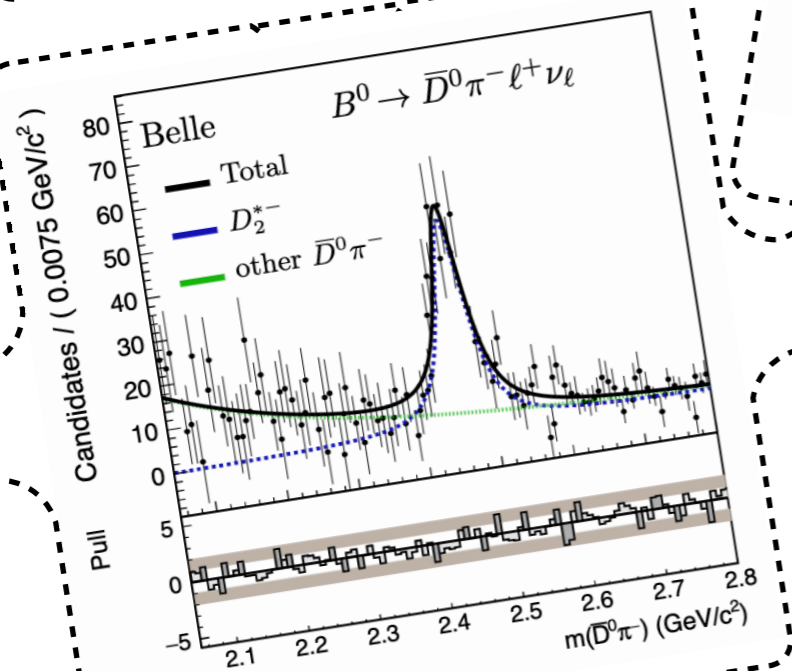
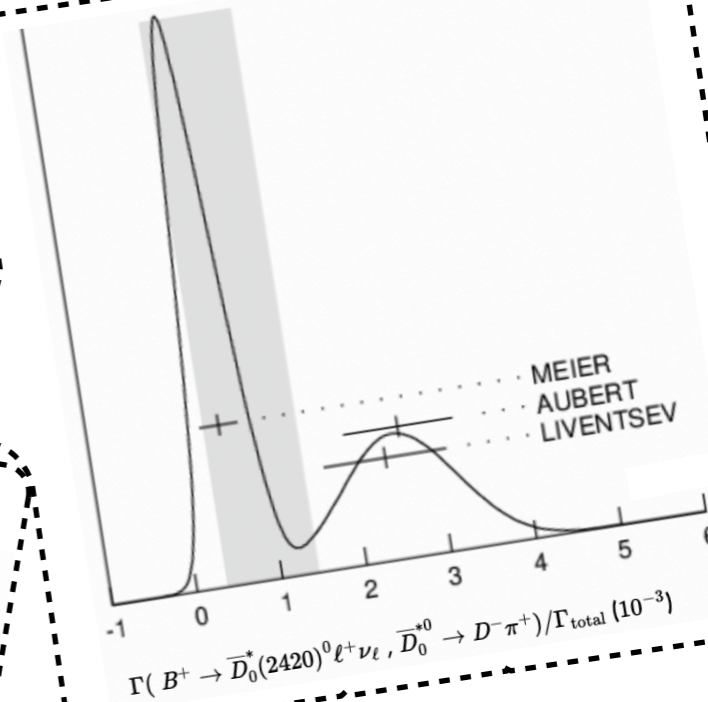
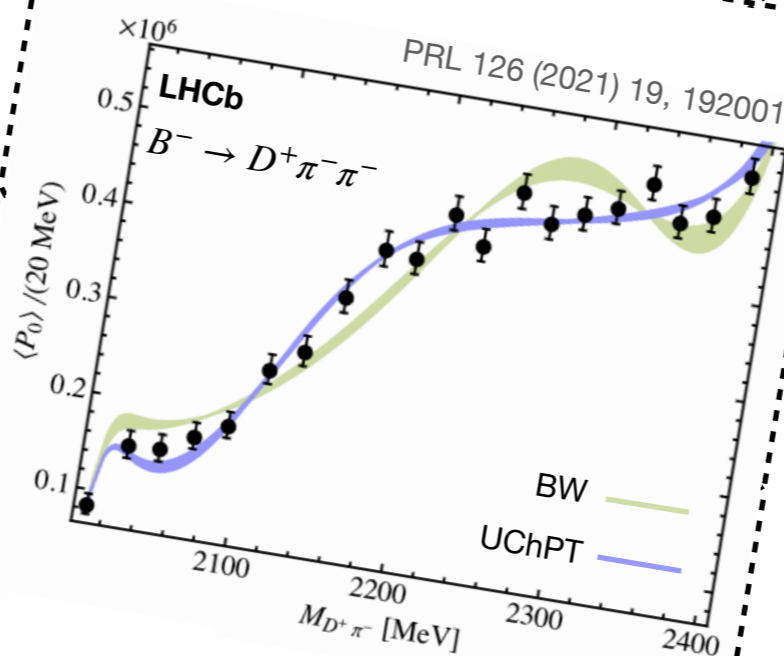
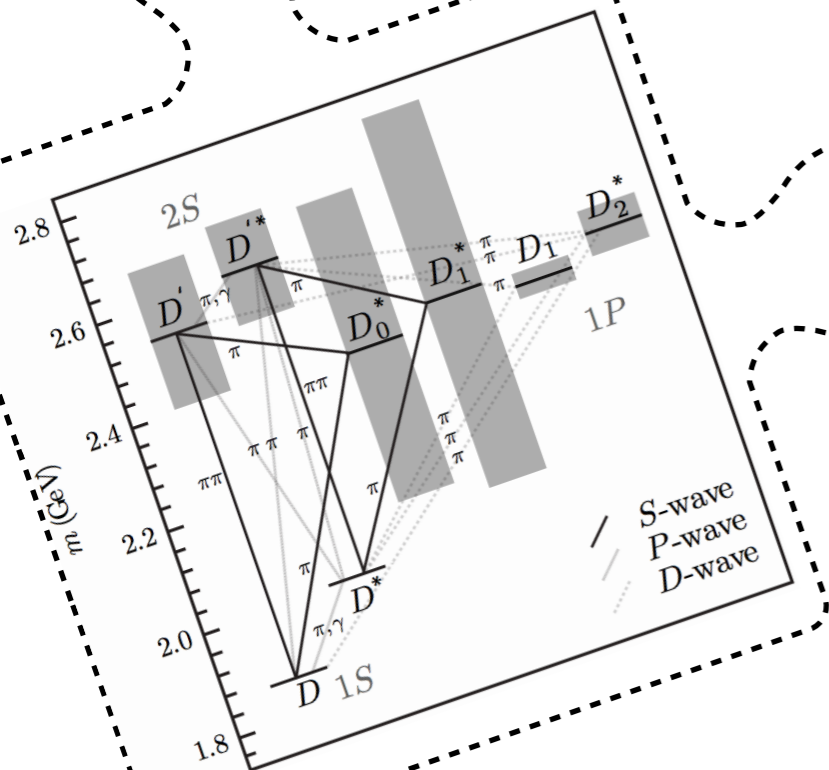


Are we missing the  $D^*$  tail?



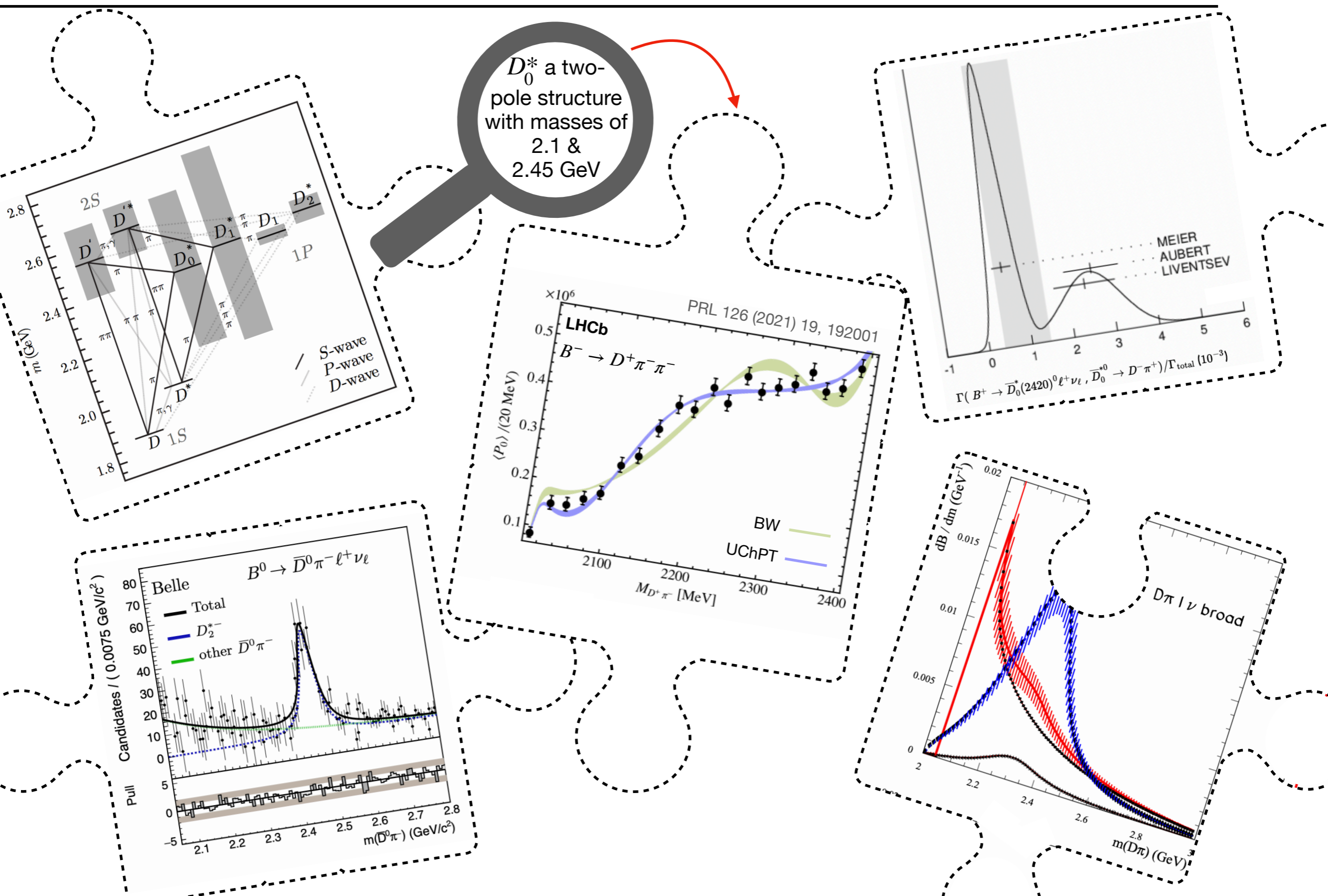


# So what do we know about $B \rightarrow D\pi\ell\nu$ ?





# So what do we know about $B \rightarrow D\pi\ell\nu$ ?





# A “how-to” for $B \rightarrow D\pi\ell\nu$

---

$$\vec{f}(q^2, M_{D\pi}^2) = \Omega(M_{D\pi}^2) \vec{P}(q^2)$$



# A “how-to” for $B \rightarrow D\pi\ell\nu$

---

## Step #1: Coupled channel treatment

**Numerically solve** the integral equation for the Omnès matrix:

$$\text{Im } \Omega(s + i\epsilon) = \frac{1}{\pi} \int_{s_{\text{thr}}}^{\infty} \frac{T^*(s') \Sigma(s') \Omega(s')}{s' - s - i\epsilon} ds'$$

$$\vec{f}(q^2, M_{D\pi}^2) = \Omega(M_{D\pi}^2) \vec{P}(q^2)$$



# A “how-to” for $B \rightarrow D\pi\ell\nu$

## Step #1: Coupled channel treatment

**Numerically solve** the integral equation for the Omnès matrix:

$$\text{Im } \Omega(s + i\epsilon) = \frac{1}{\pi} \int_{s_{\text{thr}}}^{\infty} \frac{T^*(s') \Sigma(s') \Omega(s')}{s' - s - i\epsilon} ds'$$



$$\vec{f}(q^2, M_{D\pi}^2) = \Omega(M_{D\pi}^2) \vec{P}(q^2)$$

The Omnès Matrix describes the **interactions** between **final state hadrons** and the **lineshapes of resonances**:

- $T$  is the T matrix and  $\Sigma$  contains the relevant **phase-space factors**
- Allows **simultaneous extraction** of:  
 $B \rightarrow D\pi\ell\nu$ ,  $B \rightarrow D_s K\ell\nu$  and  $B \rightarrow D\eta\ell\nu$



# A “how-to” for $B \rightarrow D\pi\ell\nu$

## Step #1: Coupled channel treatment

**Numerically solve** the integral equation for the Omnès matrix:

$$\text{Im } \Omega(s + i\epsilon) = \frac{1}{\pi} \int_{s_{\text{thr}}}^{\infty} \frac{T^*(s') \Sigma(s') \Omega(s')}{s' - s - i\epsilon} ds'$$

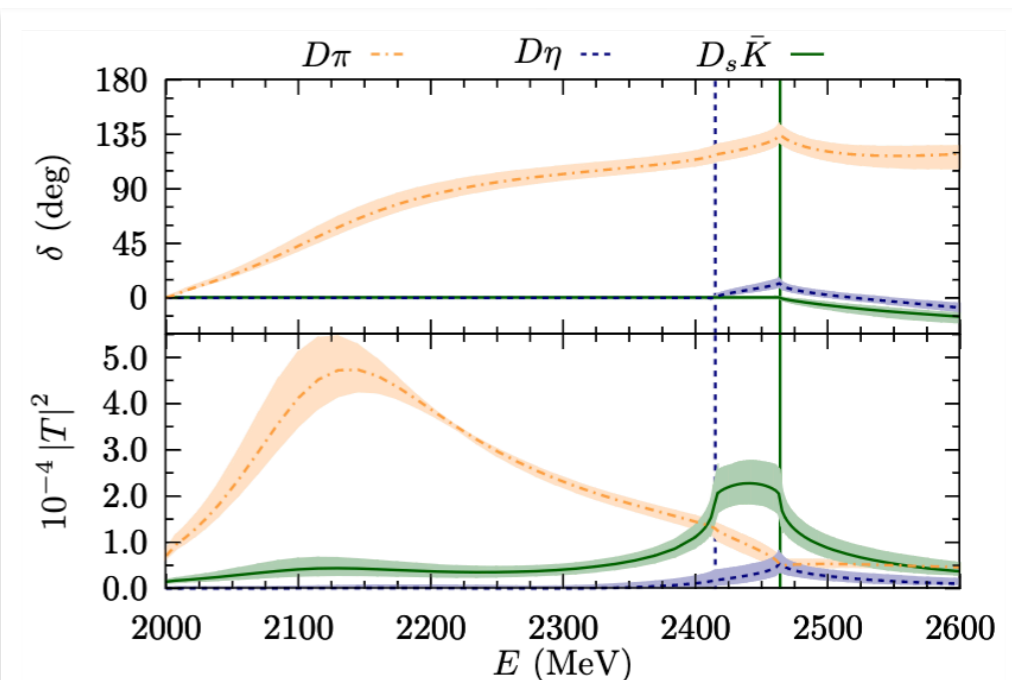
$$\vec{f}(q^2, M_{D\pi}^2) = \Omega(M_{D\pi}^2) \vec{P}(q^2)$$

PLB 767 (2017) 465-469

Hadron Spectrum: JHEP 10 (2016) 011

The Omnès Matrix describes the **interactions** between **final state hadrons** and the **lineshapes of resonances**:

- $T$  is the T matrix and  $\Sigma$  contains the relevant **phase-space factors**
- Allows **simultaneous extraction** of:  
 $B \rightarrow D\pi\ell\nu$ ,  $B \rightarrow D_s K\ell\nu$  and  $B \rightarrow D\eta\ell\nu$





# A “how-to” for $B \rightarrow D\pi\ell\nu$

---

## Step #1: Coupled channel treatment

**Numerically solve** the integral equation for the Omnès matrix.

## Step #2: Unitarity bounds

Generalize BGL unitarity bounds to **multi-hadron final states**.

$$\vec{f}(q^2, M_{D\pi}^2) = \Omega(M_{D\pi}^2) \vec{P}(q^2)$$

# A “how-to” for $B \rightarrow D\pi\ell\nu$

## Step #1: Coupled channel treatment

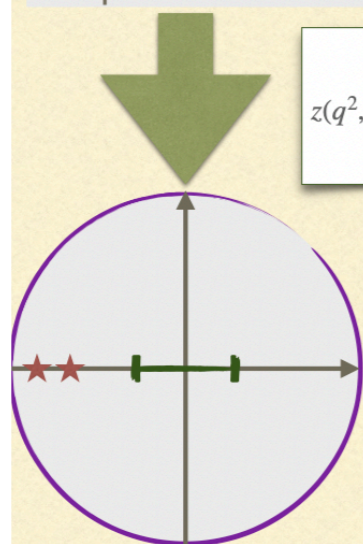
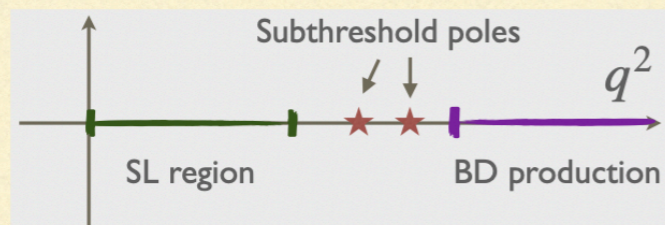
**Numerically solve** the integral equation for the Omnès matrix.

## Step #2: Unitarity bounds

Generalize BGL unitarity bounds to **multi-hadron final states**.

See Florian Herren's [talk](#).

## Theoretical fundamentals: Unitarity bounds



$$z(q^2, q_0^2) = \frac{\sqrt{q_+^2 - q^2} - \sqrt{q_+^2 - q_0^2}}{\sqrt{q_+^2 - q^2} + \sqrt{q_+^2 - q_0^2}}$$

$$1 \geq \frac{1}{2\pi i} \oint \frac{dz}{z} |B(z)\Phi(z)f(z)|^2$$

$$f(z) = \frac{1}{\Phi(z)B(z)} \sum_{i=0}^{\infty} a_i z^i \quad 1 \geq \sum_{i=0}^{\infty} |a_i|^2$$

- Mapping  $q^2$  to the dimensionless variable  $z$  transforms integration region to unit circle
- In this form it is evident that our FFs live in the Hardy space  $H^2$
- Insert Blaschke products to get rid of subthreshold poles and zeroes in kinematic factors
- Series expansion (or orthogonal polynomials)
- Semileptonic region:  $|z| < 1$

**Ingredient 2: Convergent expansion**



# A “how-to” for $B \rightarrow D\pi\ell\nu$

## Step #1: Coupled channel treatment

**Numerically solve** the integral equation for the Omnès matrix.

## Step #2: Unitarity bounds

Generalize BGL unitarity bounds to **multi-hadron final states**.  
See Florian Herren’s talk.

$$\vec{f}(q^2, M_{D\pi}^2) = \Omega(M_{D\pi}^2) \vec{P}(q^2)$$

## Step #3: Fit to $M_{D\pi}$ -spectrum

- Latest Belle results: PRD 107 (2023) 9, 092003
- **Combined fit** to both charged modes.
- We do not include data **above 2.55 GeV**;
  - To **avoid influence** from unknown higher resonances.
- Use **PDG averages** for  $D_2^*$  mass and width.

# A “how-to” for $B \rightarrow D\pi\ell\nu$

## Step #1: Coupled channel treatment

**Numerically solve** the integral equation for the Omnès matrix.

## Step #2: Unitarity bounds

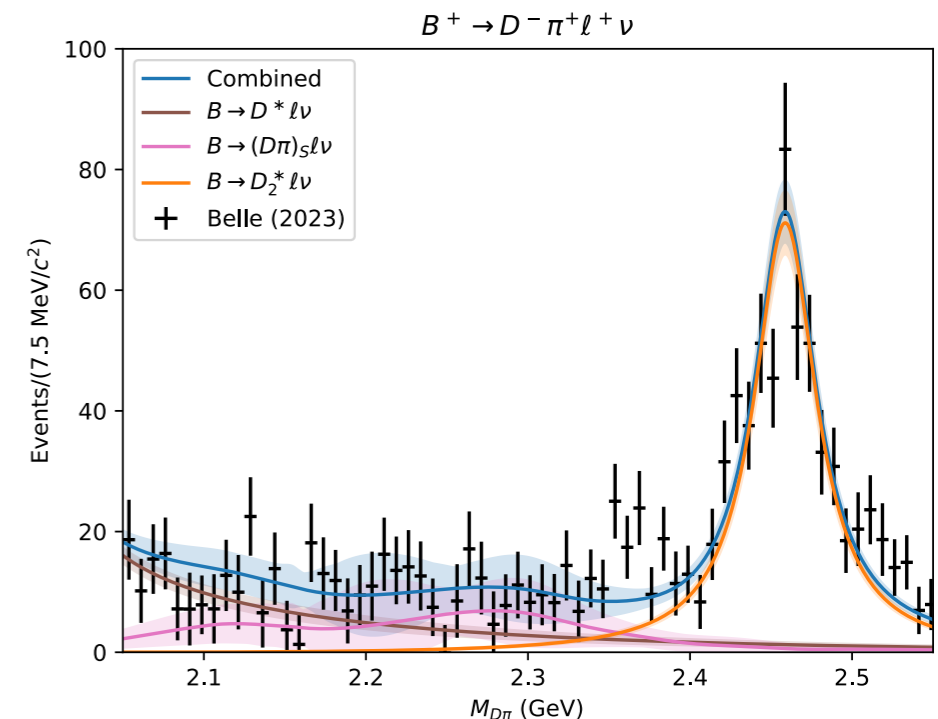
Generalize BGL unitarity bounds to **multi-hadron final states**.  
See Florian Herren’s talk.

$$\vec{f}(q^2, M_{D\pi}^2) = \Omega(M_{D\pi}^2) \vec{P}(q^2)$$

**Excellent agreement with data:**  
 $\chi^2/\text{dof} = 1.0$  (134)

## Step #3: Fit to $M_{D\pi}$ -spectrum

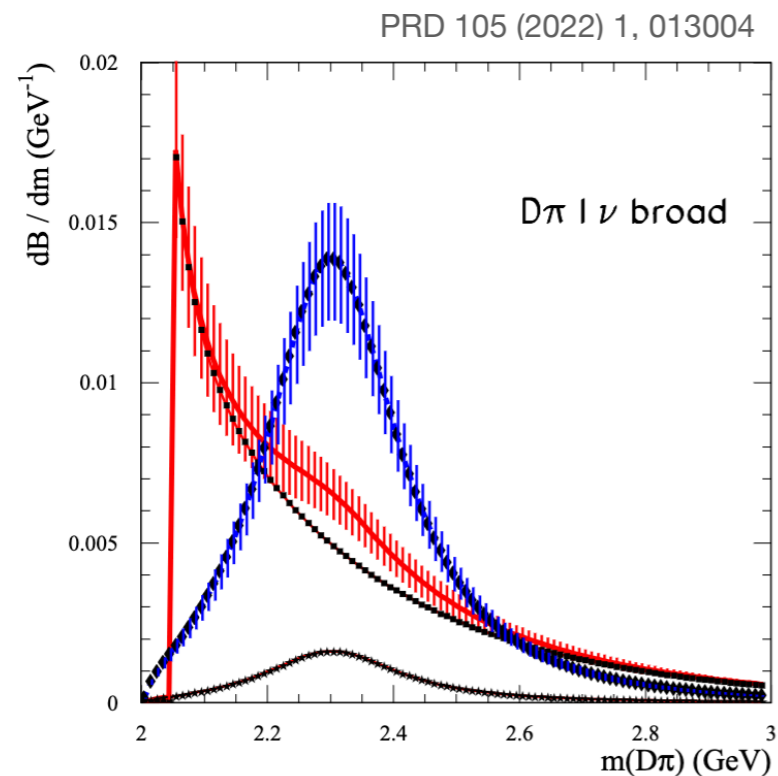
- Latest Belle results: PRD 107 (2023) 9, 092003
- **Combined fit** to both charged modes.
- We do not include data **above 2.55 GeV**;
  - To **avoid influence** from unknown higher resonances.
- Use **PDG averages** for  $D_2^*$  mass and width.





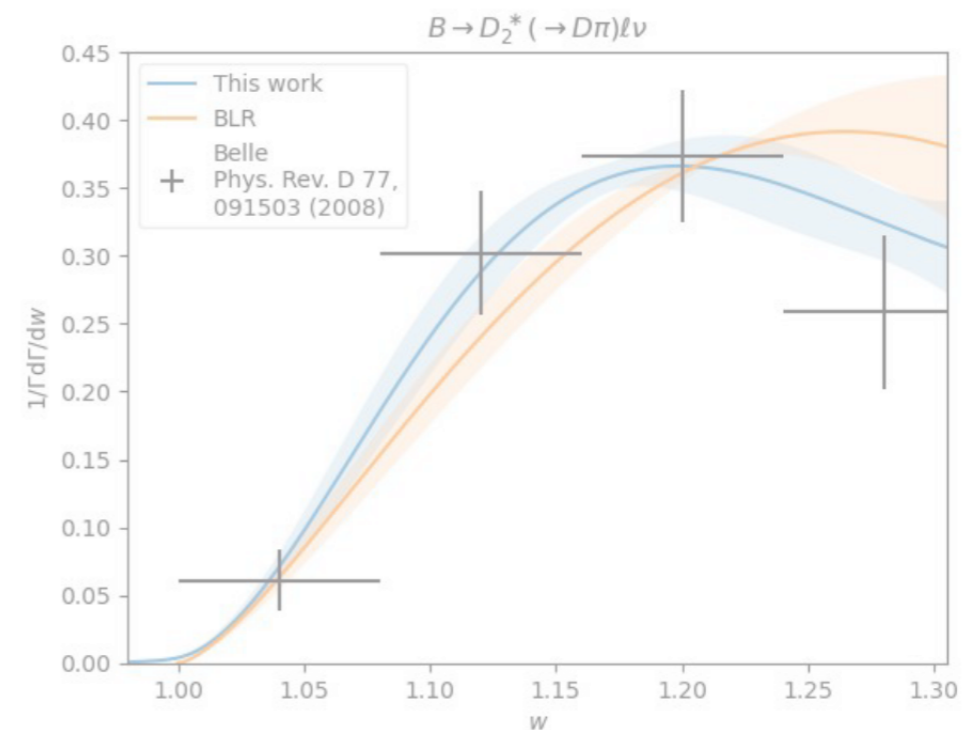
# Additional constraints

- Recently, it was pointed out that virtual  $D^*$  contributions should be taken into account in semileptonic decays.



- We introduce Blatt-Weisskopf damping factors and include  $r_{BW}$  as fit parameter.
- Use FNAL/MILC  $D^*$  FFs and fit after integrating over the  $D\pi$  invariant mass.

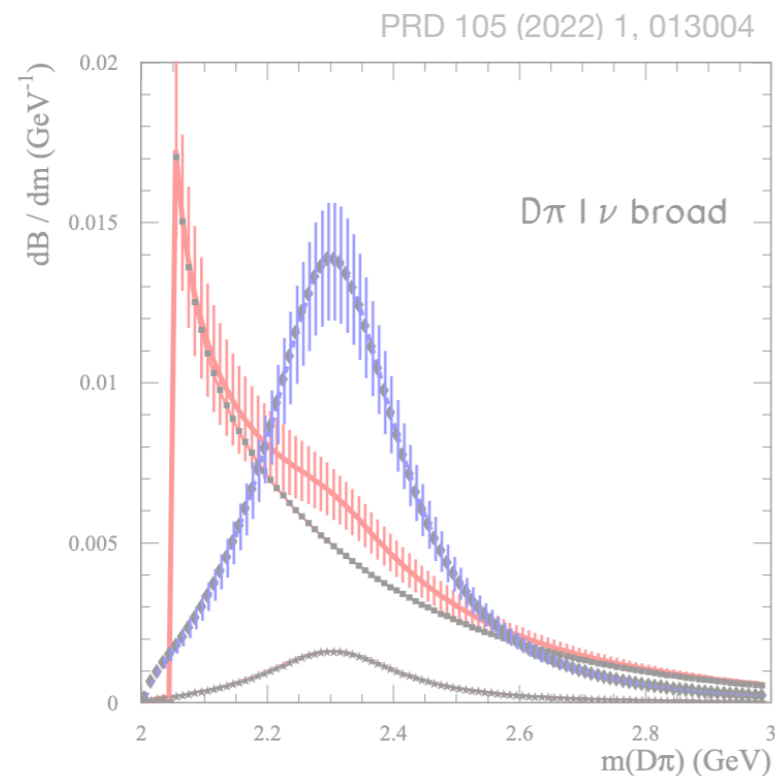
- The  $D_2^*$  FFs are fitted to the spectra measured by Belle with loose constraints on:
  - $B \rightarrow D_2^*(\rightarrow D\pi)\ell\nu$  decay rate
  - $B \rightarrow D_2^*(\rightarrow D\pi)\pi/K$  BFs



- Uncertainties could be decreased by implementing the HQET constraints present in the LLSW parametrization.

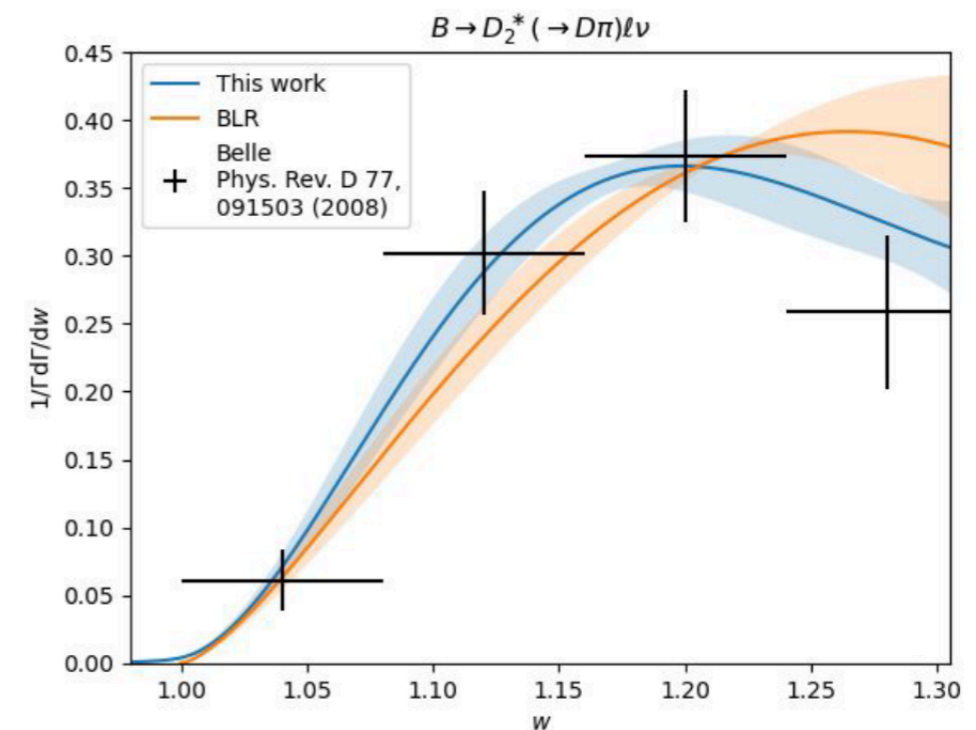
# Additional constraints

- Recently, it was pointed out that virtual  $D^*$  contributions should be taken into account in semileptonic decays.



- We introduce Blatt-Weisskopf damping factors and include  $r_{BW}$  as fit parameter.
- Use FNAL/MILC  $D^*$  FFs and fit after integrating over the  $D\pi$  invariant mass.

- The  $D_2^*$  FFs are fitted to the spectra measured by Belle with loose constraints on:
  - $B \rightarrow D_2^*( \rightarrow D\pi)\ell\nu$  decay rate
  - $B \rightarrow D_2^*( \rightarrow D\pi)\pi/K$  BFs

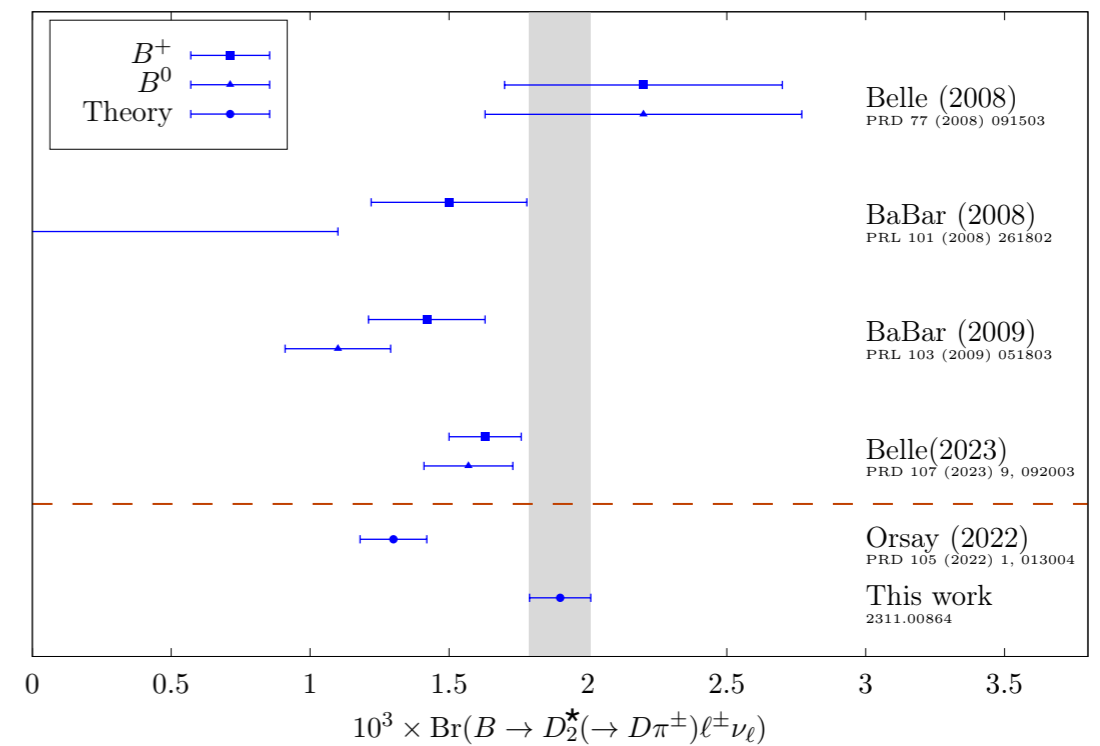
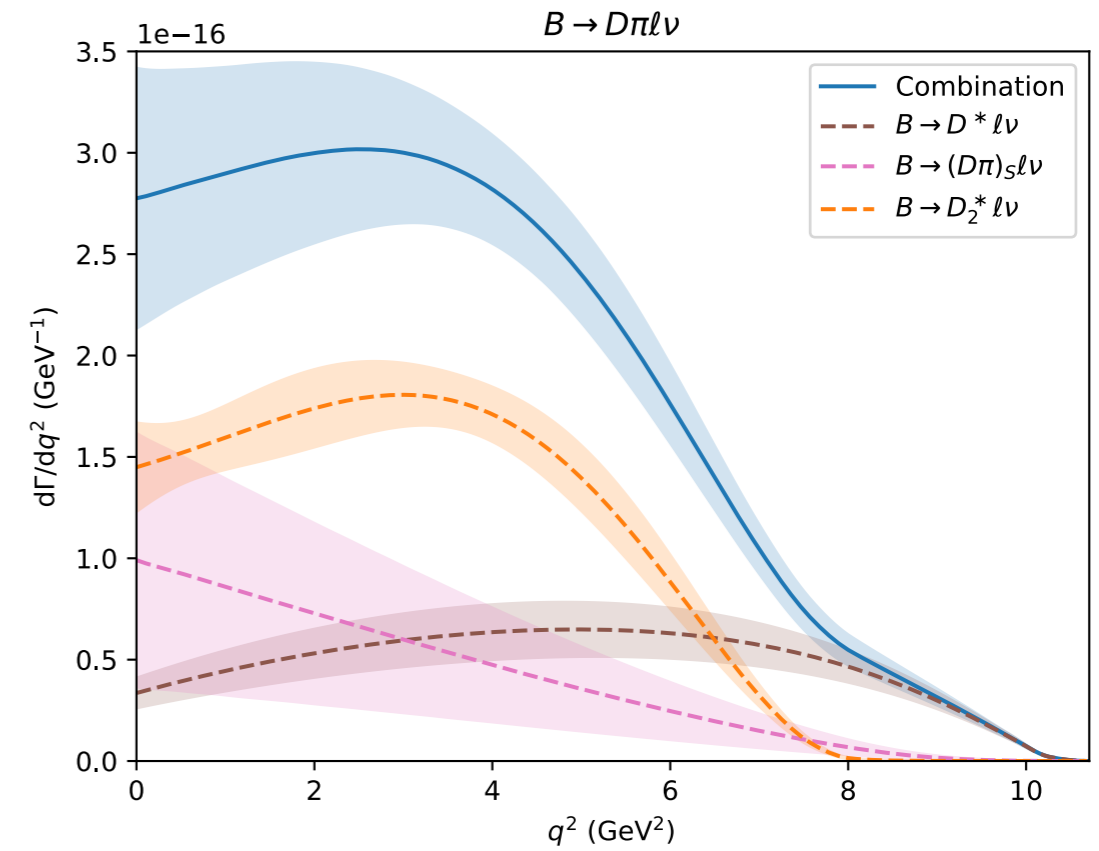


- Uncertainties could be decreased by implementing the HQET constraints present in the LLSW parametrization.



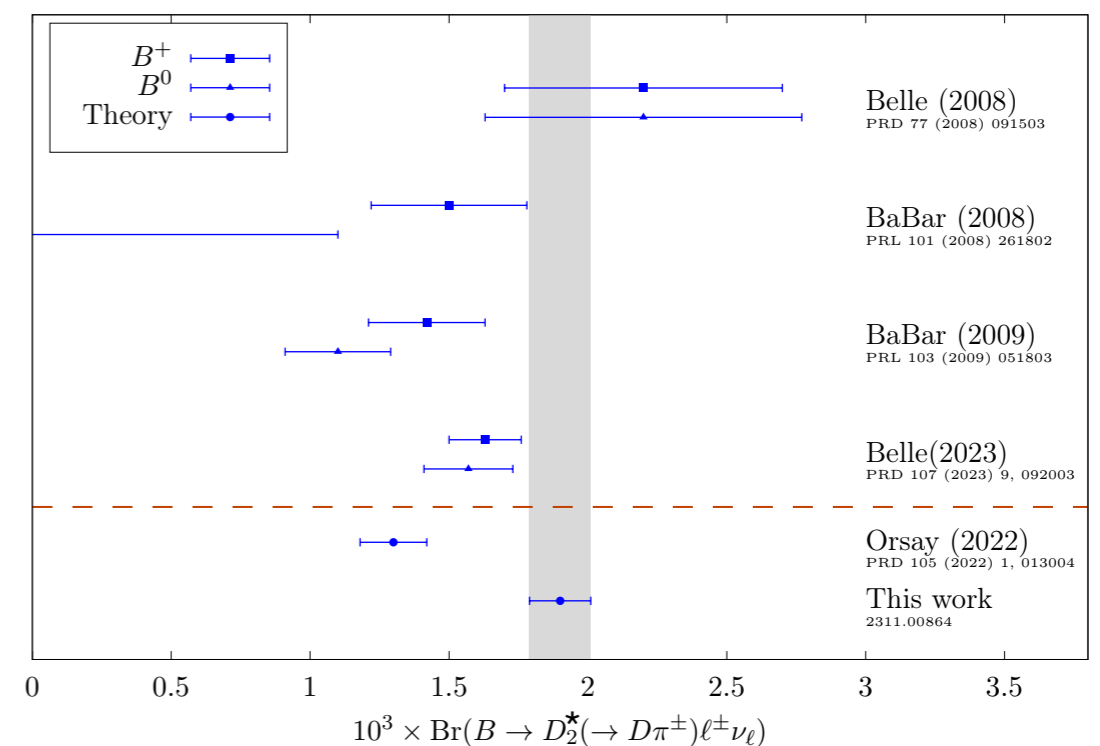
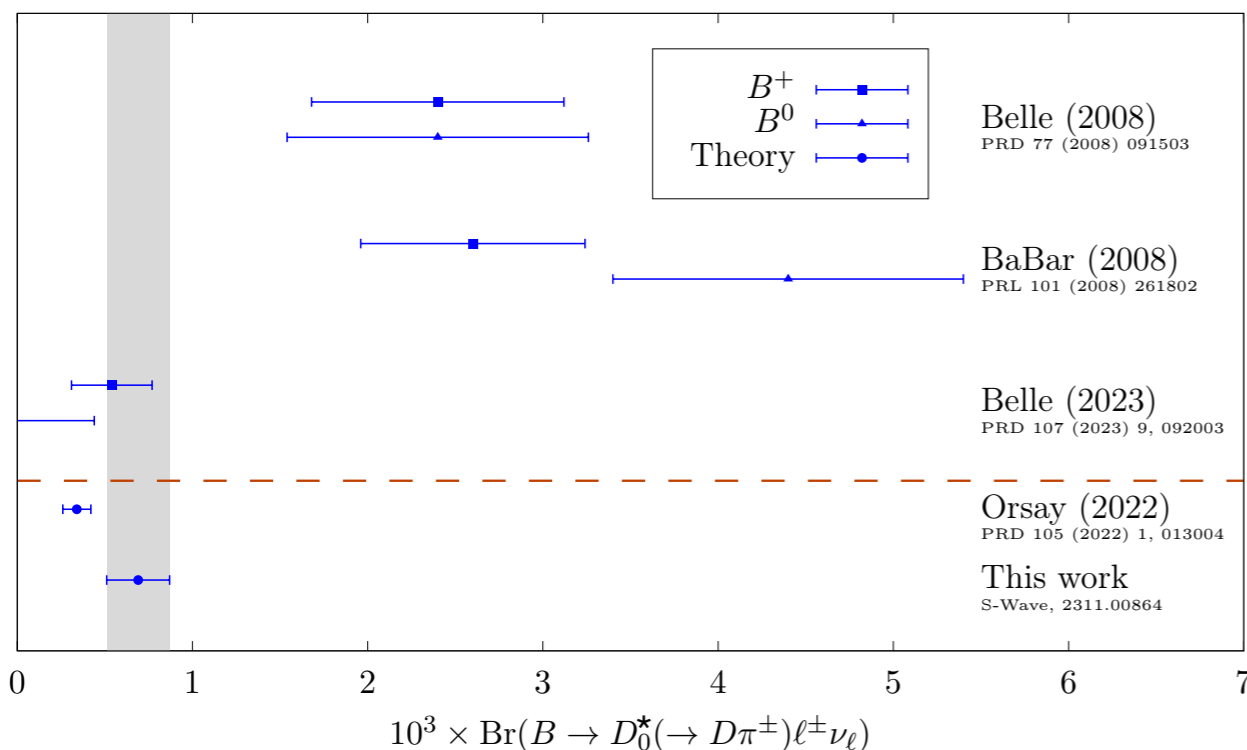
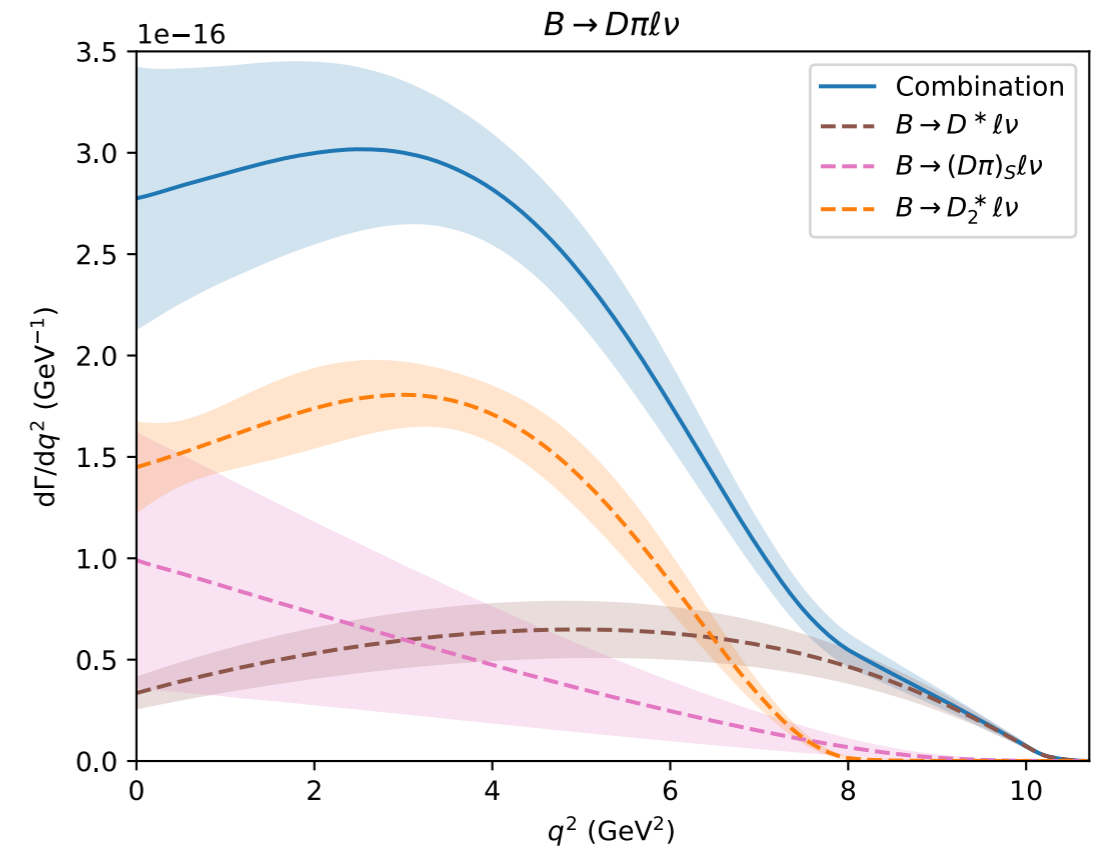
# Results & predictions

- We find a **significantly larger  $D_2^*$  yield** than quoted by the PDG.



# Results & predictions

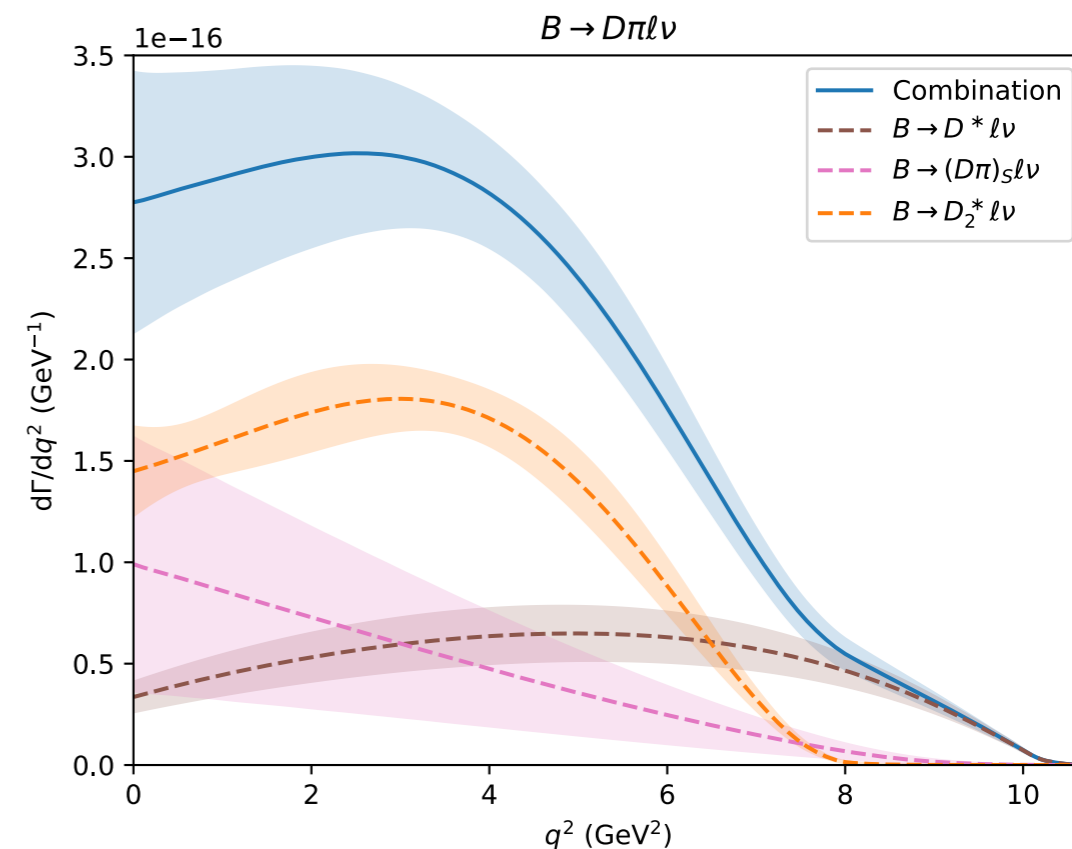
- We find a **significantly larger  $D_2^*$  yield** than quoted by the PDG.
- Our  $D^*$  and S-wave contributions **drop off faster** than the falling exponential used by Belle.



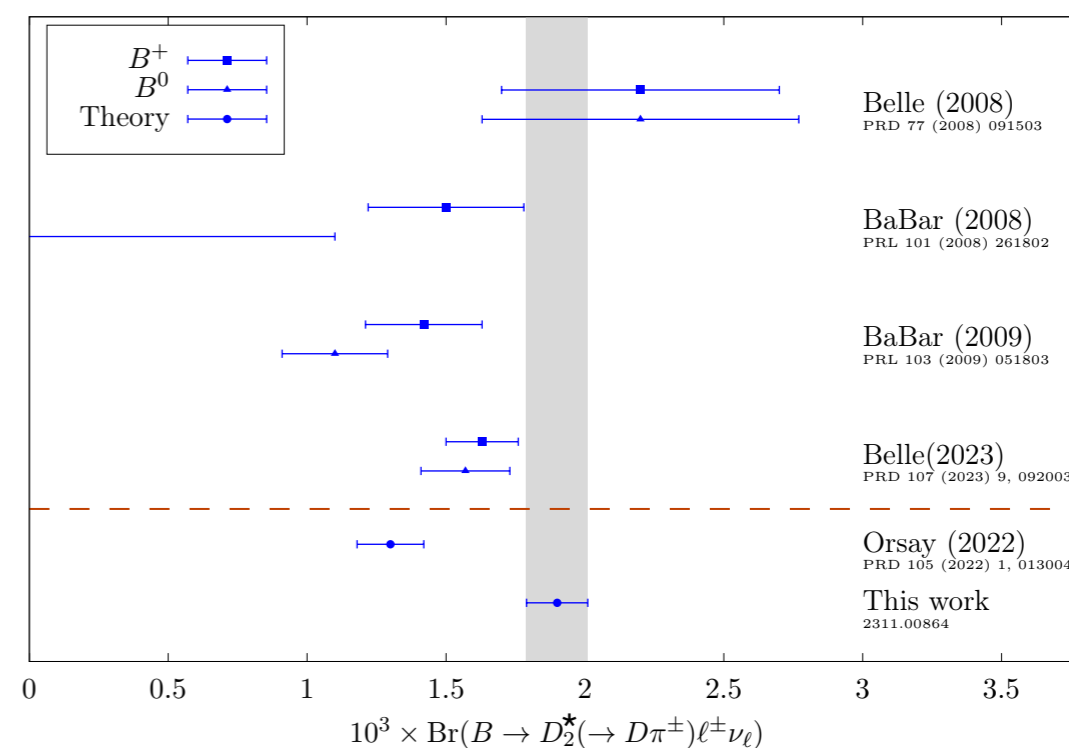
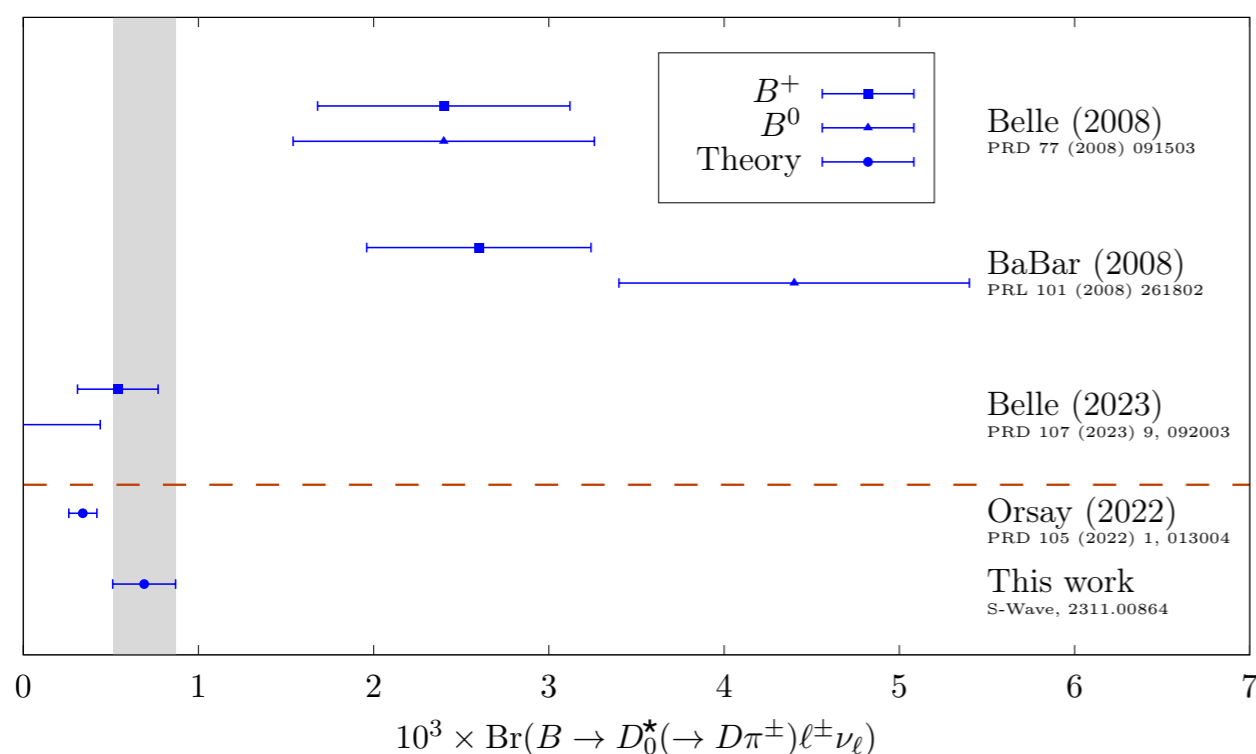


# Results & predictions

- We find a **significantly larger  $D_2^*$  yield** than quoted by the PDG.
- Our  $D^*$  and S-wave contributions **drop off faster** than the falling exponential used by Belle.
- S-Wave  $B \rightarrow D\eta\ell\nu$  decays **cannot account for the gap** between the inclusive BF and the sum of exclusive decays.
- By **heavy quark symmetry**  $B \rightarrow D^*\eta\ell\nu$  decays will also be subdominant.



Prediction:  $\text{Br}(B \rightarrow D\eta\ell\bar{\nu}_\ell) = (1.9 \pm 1.7) \times 10^{-5}$







$B \rightarrow \pi\pi\ell\nu$



# $|V_{ub}|$ and the $B \rightarrow \rho \ell \nu$ conundrum

- Recent Belle II results shown at Moriond EW 2024 report lower  $|V_{ub}|$  value from  $B \rightarrow \rho \ell \nu$  decays than from  $B \rightarrow \pi \ell \nu$ .

## Simultaneous measurements of $B^0 \rightarrow \pi^- \ell^+ \nu$ , $B^+ \rightarrow \rho^0 \ell^+ \nu$



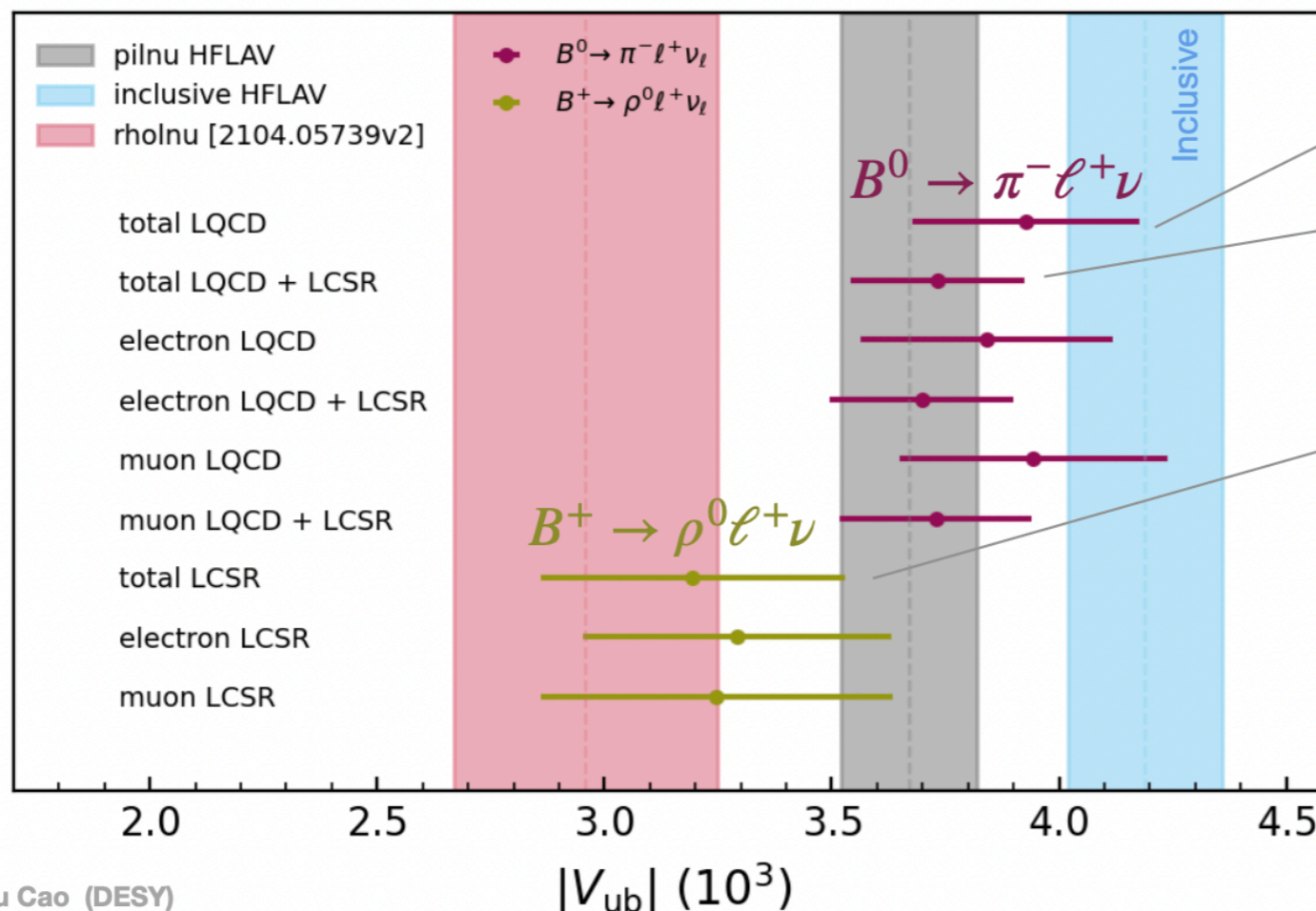
Preliminary

**NEW!!**

- Further split into  $e$  and  $\mu$  modes to provide cross check
- Additional stability tests done by removing higher/lower  $q^2$  bins

Preliminary

arXiv:2407.17403



$$|V_{ub}|_{B \rightarrow \pi \ell \nu \ell} = (3.93 \pm 0.09 \pm 0.13 \pm 0.19) \times 10^{-3}$$

LQCD

stat    syst    theo

$$|V_{ub}|_{B \rightarrow \pi \ell \nu \ell} = (3.73 \pm 0.07 \pm 0.07 \pm 0.16) \times 10^{-3}$$

LQCD+LCSR

$$|V_{ub}|_{B \rightarrow \rho \ell \nu \ell} = (3.19 \pm 0.12 \pm 0.17 \pm 0.26) \times 10^{-3}$$

LCSR

- Leading systematic unc. are the modelling of continuum and non-resonant  $B \rightarrow X_u \ell \nu$  decays
- Overall **theoretical** uncertainty dominating



# $|V_{ub}|$ and the $B \rightarrow \rho \ell \nu$ conundrum

- Recent Belle II results shown at Moriond EW 2024 report lower  $|V_{ub}|$  value from  $B \rightarrow \rho \ell \nu$  decays than from  $B \rightarrow \pi \ell \nu$ .

## Simultaneous measurements of $B^0 \rightarrow \pi^- \ell^+ \nu$ , $B^+ \rightarrow \rho^0 \ell^+ \nu$



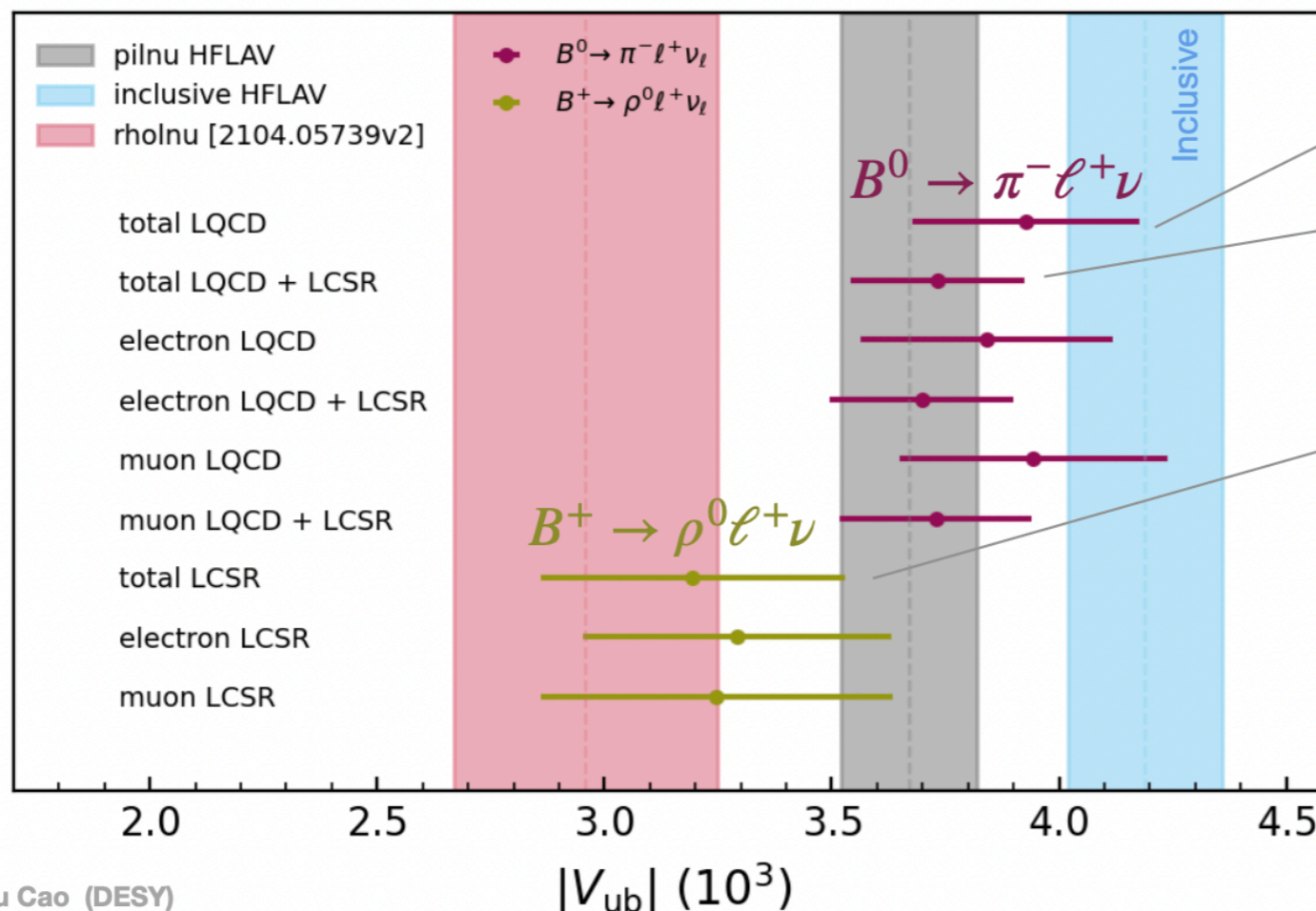
Preliminary

**NEW!!**

- Further split into  $e$  and  $\mu$  modes to provide cross check
- Additional stability tests done by removing higher/lower  $q^2$  bins

Preliminary

arXiv:2407.17403



$$|V_{ub}|_{B \rightarrow \pi \ell \nu \ell} = (3.93 \pm 0.09 \pm 0.13 \pm 0.19) \times 10^{-3}$$

stat    syst    theo

LQCD

$$|V_{ub}|_{B \rightarrow \pi \ell \nu \ell} = (3.73 \pm 0.07 \pm 0.07 \pm 0.16) \times 10^{-3}$$

LQCD+LCSR

$$|V_{ub}|_{B \rightarrow \rho \ell \nu \ell} = (3.19 \pm 0.12 \pm 0.17 \pm 0.26) \times 10^{-3}$$

LCSR

- Leading systematic unc. are the modelling of continuum and non-resonant  $B \rightarrow X_u \ell \nu$  decays
- Overall **theoretical** uncertainty dominating





# What do we have to work with?

---

- We collect **measurements** from many different sources for the **lineshapes**:
  - P-wave (the dominant contribution),
  - D-wave (the sub-dominant contribution),
  - last, but not least, the S-wave.

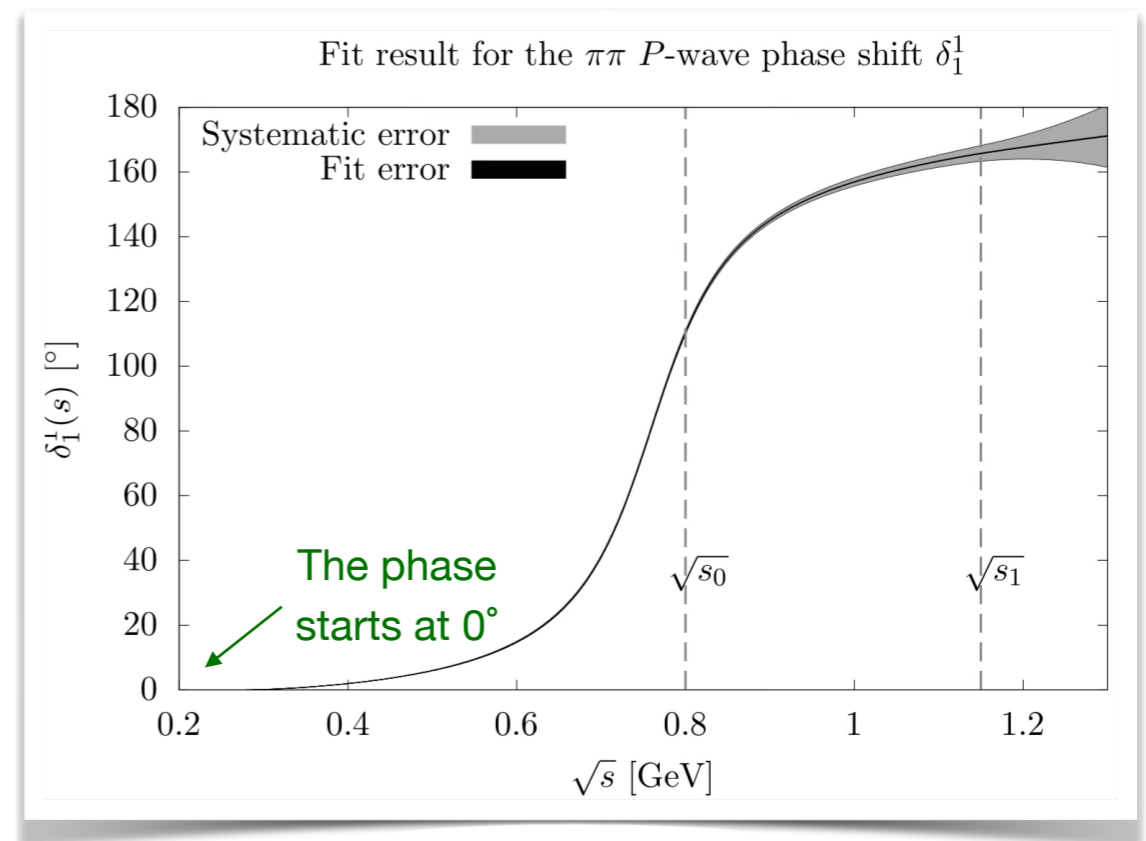
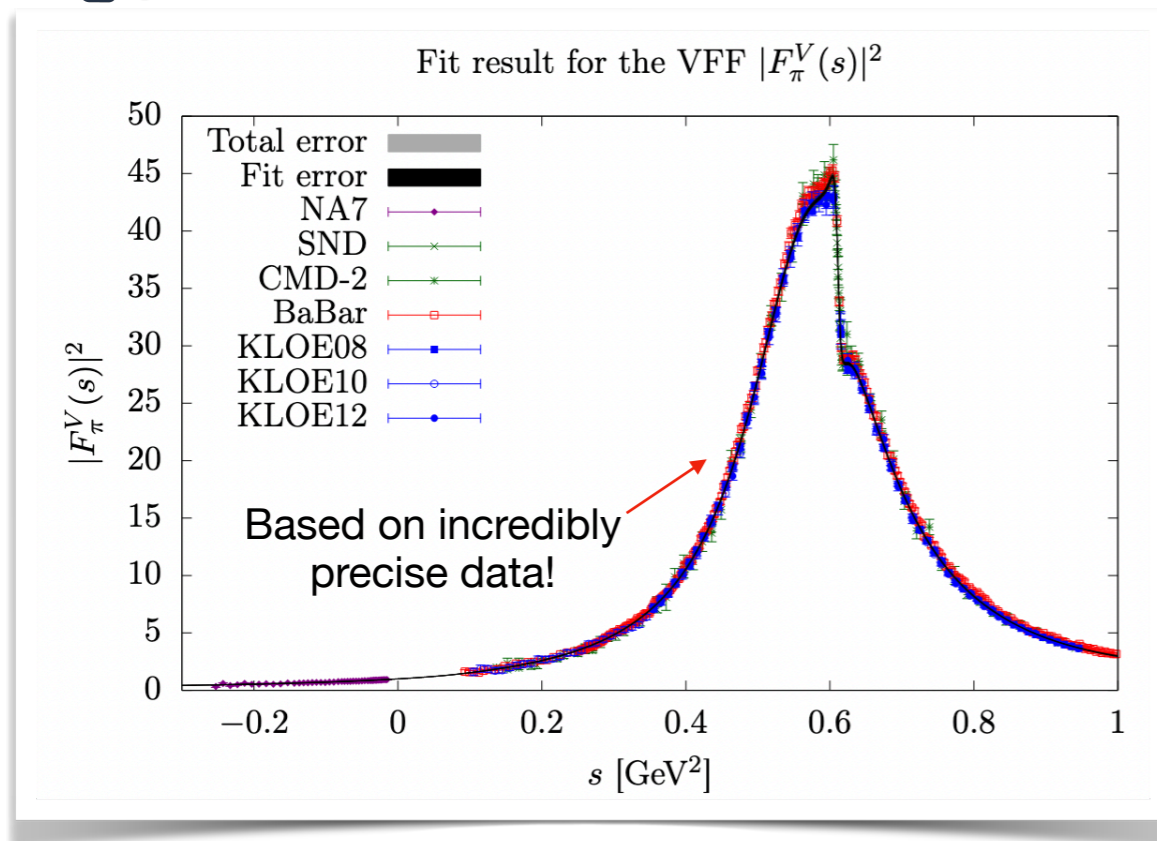
# What do we have to work with?

- We collect **measurements** from many different sources for the **lineshapes**:
  - **P-wave (the dominant contribution)**,
  - D-wave (the sub-dominant contribution),
  - last, but not least, the S-wave.

Data from  $e^+e^- \rightarrow \pi^+\pi^-$  production provides a high precision determination of the P-wave phase.



For single-channel problems, this is everything we need for the Omnès factor.

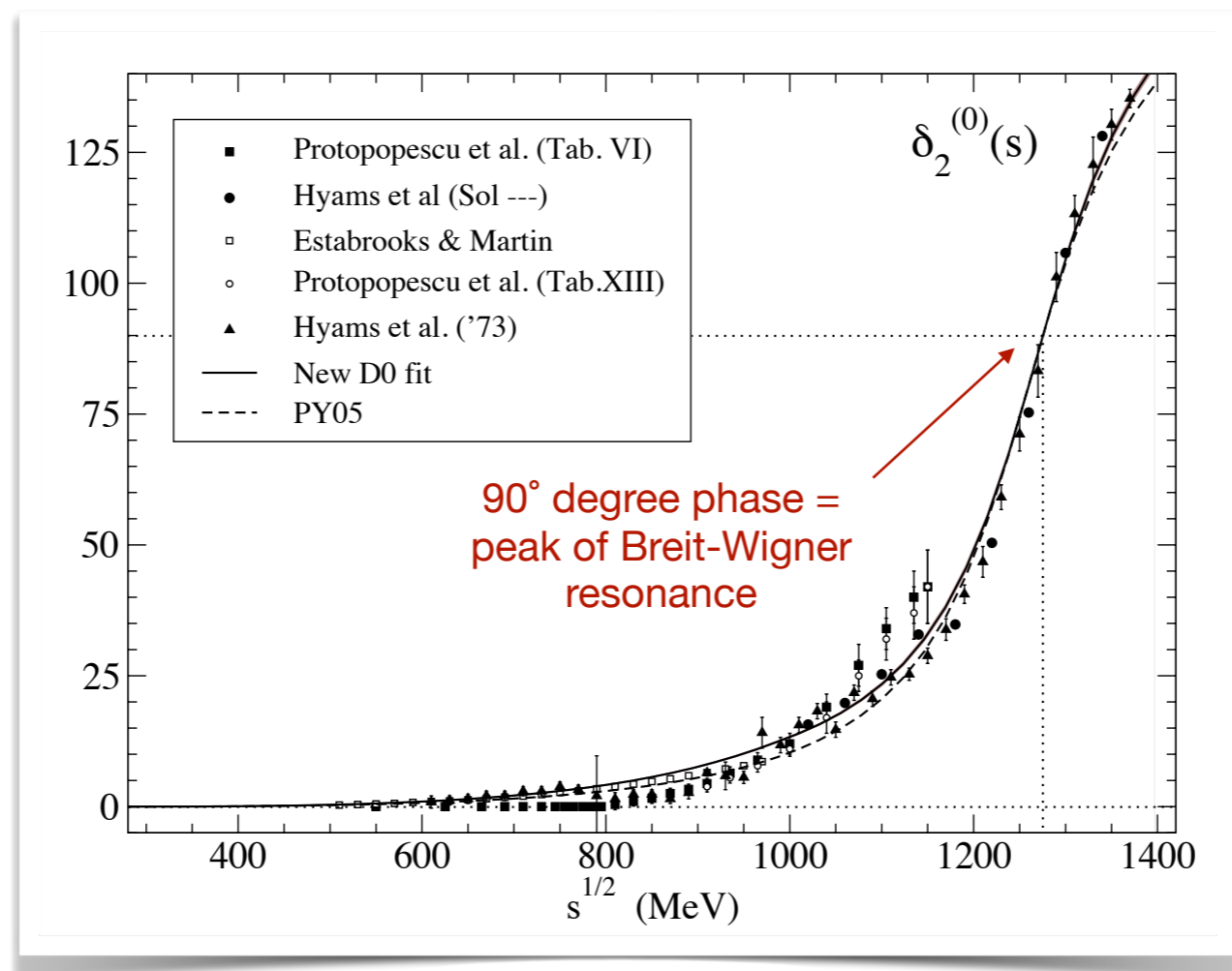




# What do we have to work with?

- We collect **measurements** from many different sources for the **lineshapes**:
  - P-wave (the dominant contribution),
  - **D-wave (the sub-dominant contribution),**
  - last, but not least, the S-wave.

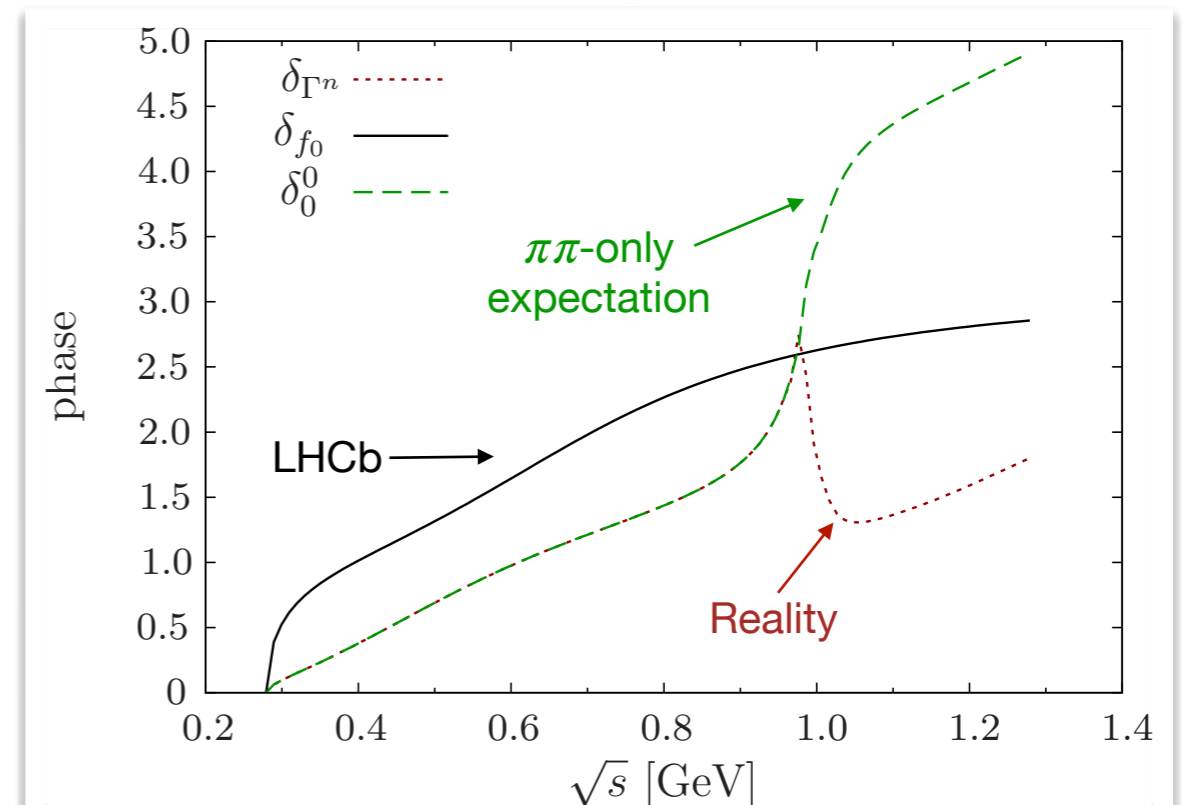
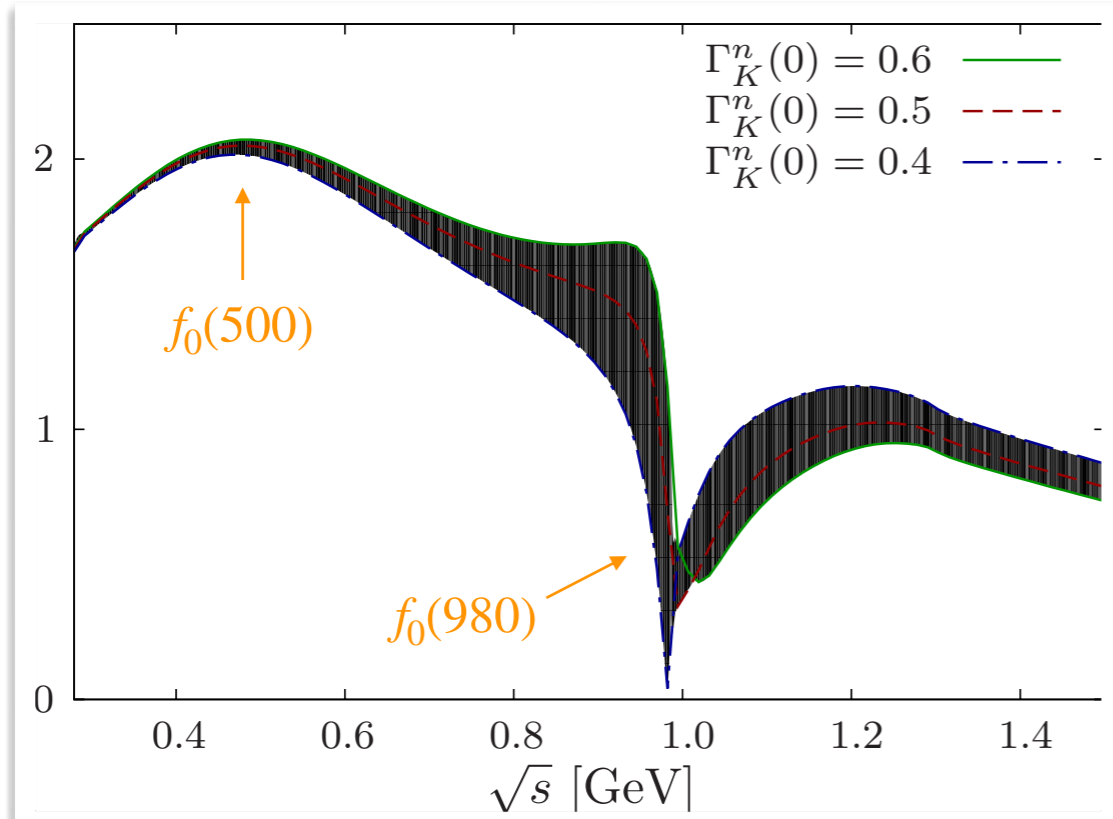
Data from  $\pi^+\pi^- \rightarrow \pi^+\pi^-$  scattering gives the necessary information to describe the  $f_2(1270)$  resonance.



# What do we have to work with?

- We collect **measurements** from many different sources for the **lineshapes**:
  - P-wave (the dominant contribution),
  - D-wave (the sub-dominant contribution),
  - **last, but not least, the S-wave.**

A coupled-channel analysis of  $\bar{B}_{d/s}^0 \rightarrow J/\psi\pi\pi$  decays provide the last piece of the puzzle.

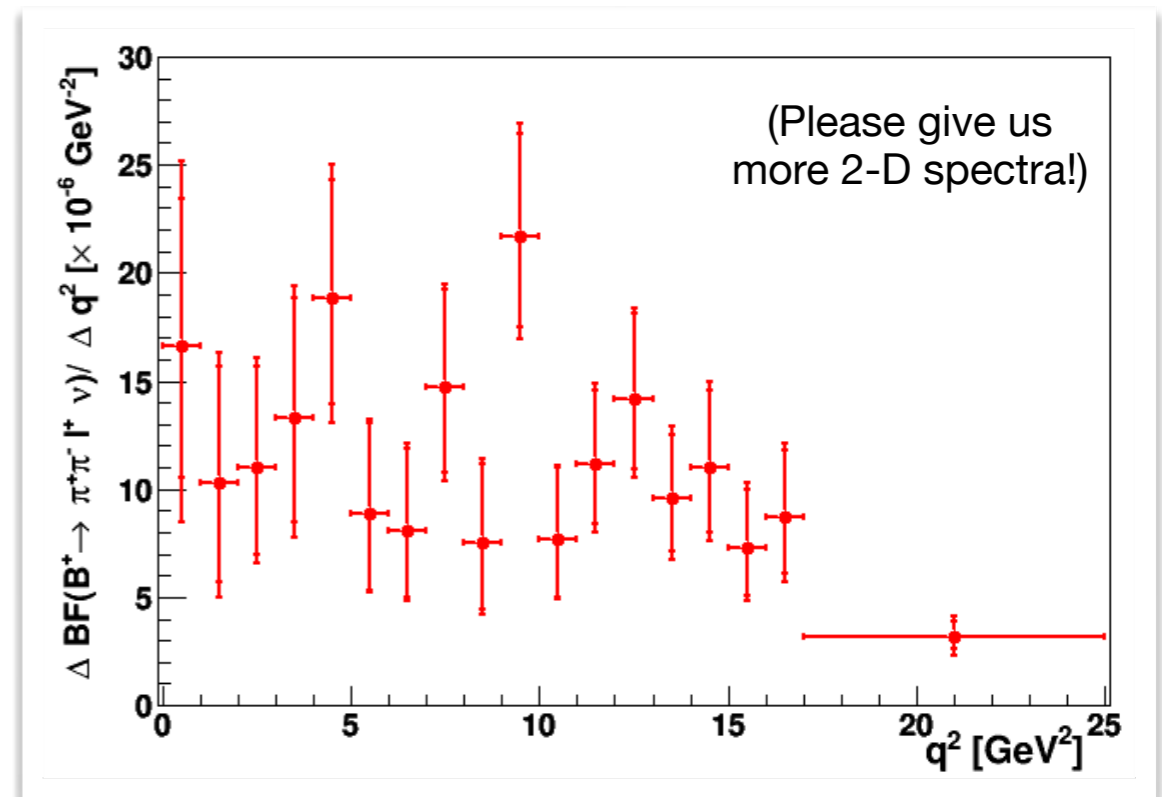
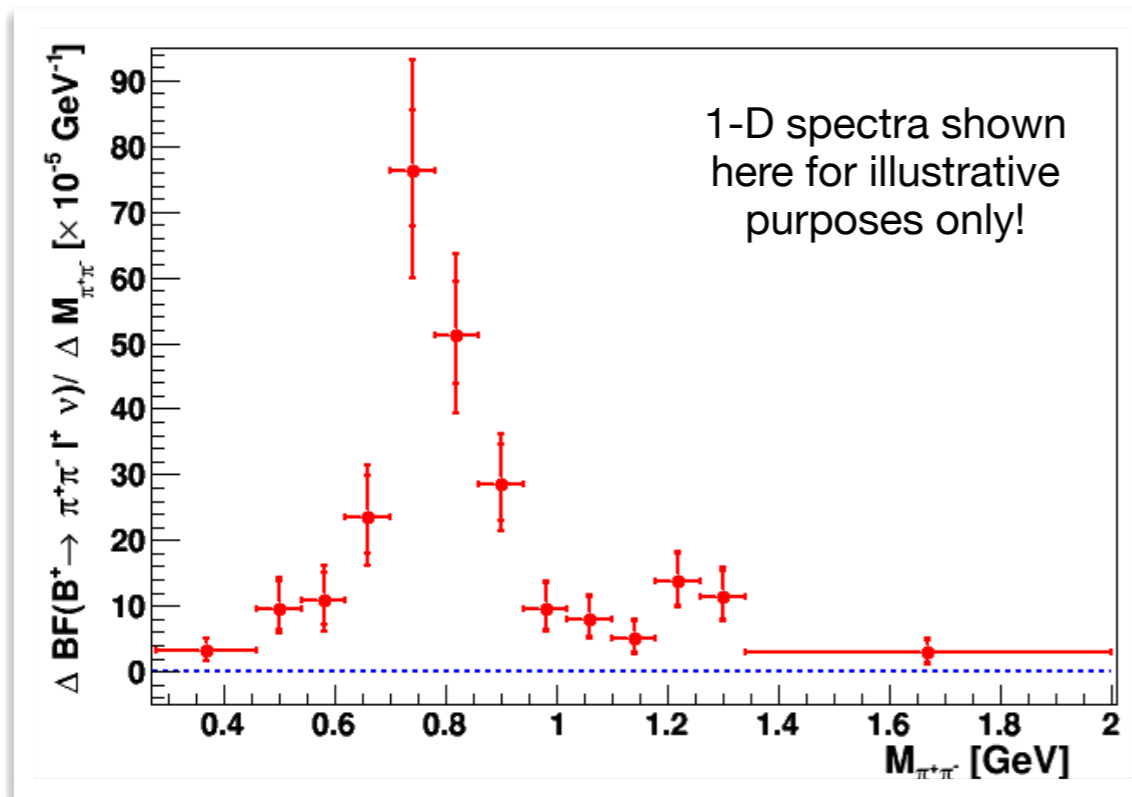




# What do we have to work with?

- We collect **measurements** from many different sources for the **lineshapes**:
  - P-wave (the dominant contribution),
  - D-wave (the sub-dominant contribution),
  - last, but not least, the S-wave.

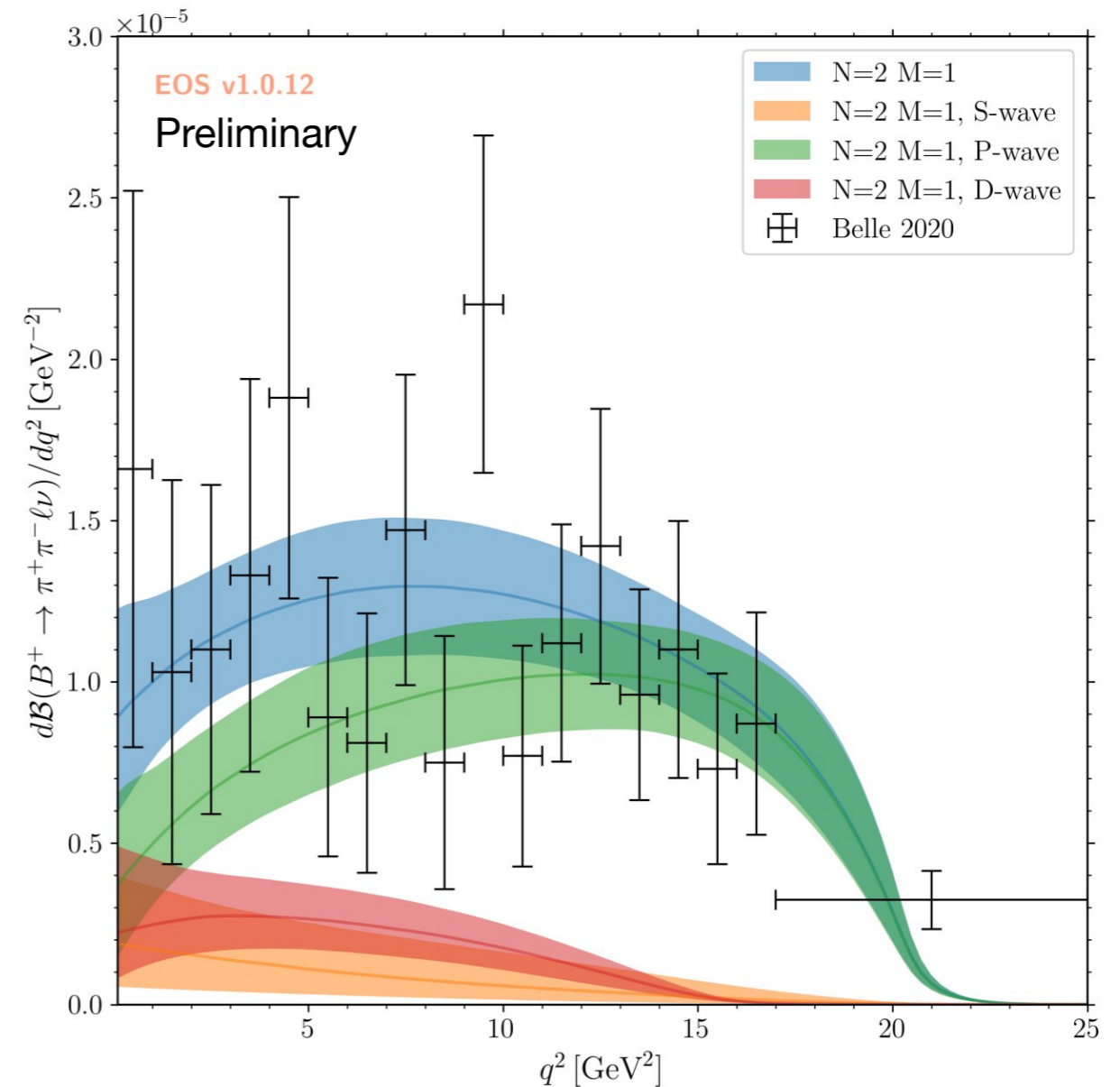
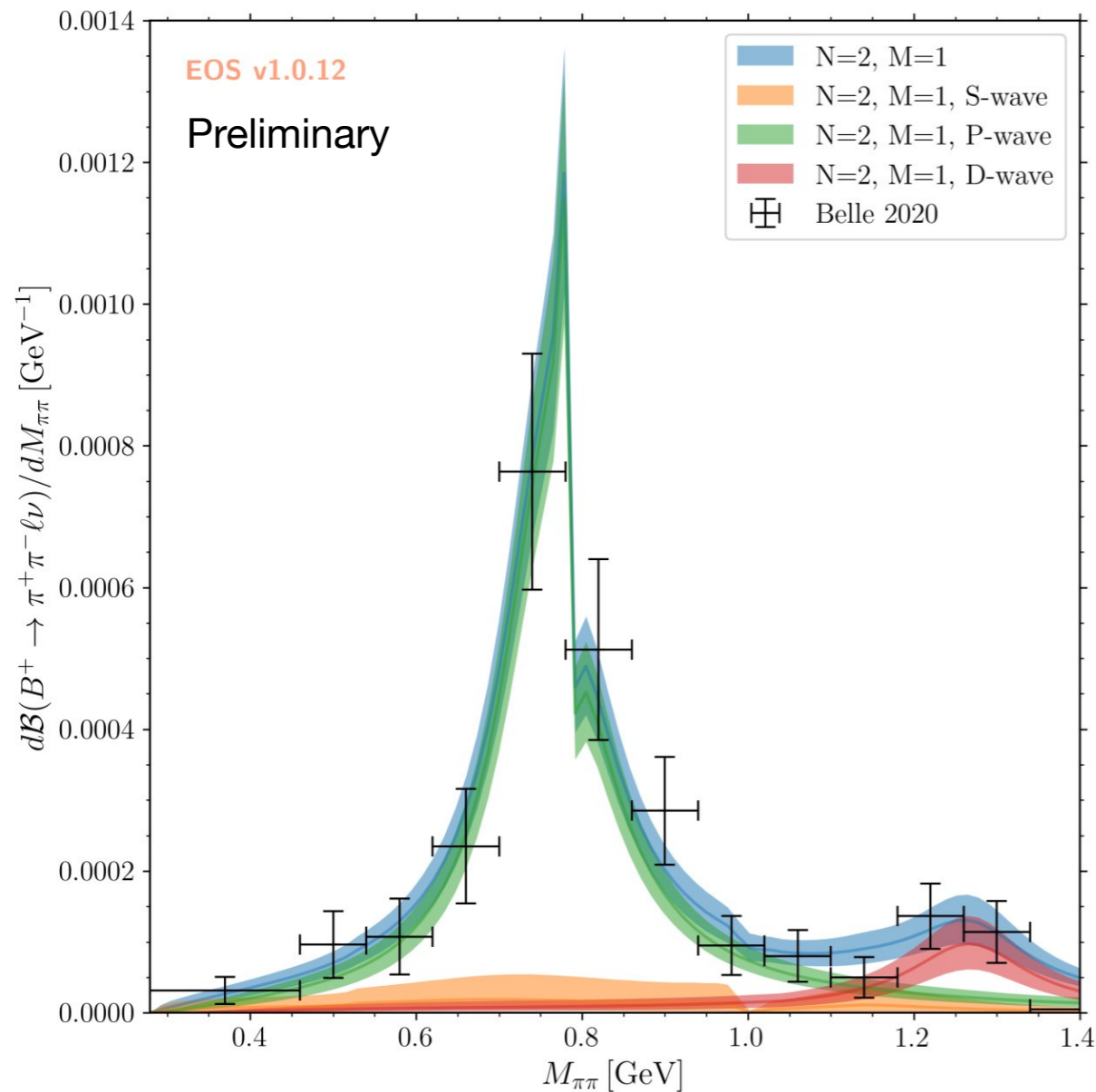
A recent Belle measurement of  $B^+ \rightarrow \pi^+ \pi^- \ell^+ \nu$  decays provides 2-D correlated spectra of  $q^2$  &  $M_{\pi^+ \pi^-}$ .  
We fit to the 2-D spectra using the [EOS package](#).



# Preliminary results

Paper in progress!

- We provide a complete description:
  - Lineshapes, form factors, and correlated uncertainties.
- Currently working on refining the unitarity bounds by including additional processes.
- The final results will be available in EOS for direct interfacing with future analyses.





# Conclusion & Outlook

---

- We developed a **powerful framework** allowing for the **complete description** of processes with 2 hadrons in the final state,
- We **improve the modelling** of many processes of interest to ongoing experimental analyses:

# Conclusion & Outlook

---

- We developed a **powerful framework** allowing for the **complete description** of processes with 2 hadrons in the final state,
- We **improve the modelling** of many processes of interest to ongoing experimental analyses:

$$B \rightarrow D^{**} \ell \nu$$

- $D_0^*$  and friends are not what we think they are. We need more spectra (not BFs) to disentangle different resonances.
- The  $D^*$  tail should be correctly handled in MC!





# Conclusion & Outlook

---

- We developed a **powerful framework** allowing for the **complete description** of processes with 2 hadrons in the final state,
- We **improve the modelling** of many processes of interest to ongoing experimental analyses:

$$B \rightarrow D^{**}\ell\nu$$

- $D_0^*$  and friends are not what we think they are. We need more spectra (not BFs) to disentangle different resonances.
- The  $D^*$  tail should be correctly handled in MC!



$$B \rightarrow D^{(*)}\eta\ell\nu \quad \& \quad B \rightarrow X_c\ell\nu$$

- We rule out the current “best guess” used to fill the semileptonic gap.  
  
(So stop saying it could fill the gap, prove it with direct experimental evidence.)



# Conclusion & Outlook

- We developed a **powerful framework** allowing for the **complete description** of processes with 2 hadrons in the final state,
- We **improve the modelling** of many processes of interest to ongoing experimental analyses:

$$B \rightarrow D^{**}\ell\nu$$

- $D_0^*$  and friends are not what we think they are. We need more spectra (not BFs) to disentangle different resonances.
- The  $D^*$  tail should be correctly handled in MC!



$$B \rightarrow D^{(*)}\eta\ell\nu \quad \& \quad B \rightarrow X_c\ell\nu$$

- We rule out the current “best guess” used to fill the semileptonic gap.

(So stop saying it could fill the gap, prove it with direct experimental evidence.)



$$B \rightarrow \rho\ell\nu \quad \& \quad B \rightarrow X_u\ell\nu$$

- We provide lineshapes and form factors for all  $B \rightarrow \pi\pi\ell\nu$  decays.
- Looking forward to upcoming lattice results.
- It would be interesting to see LCSR calculations using our parametrization!





# Conclusion & Outlook

- We developed a **powerful framework** allowing for the **complete description** of processes with 2 hadrons in the final state,
- We **improve the modelling** of many processes of interest to ongoing experimental analyses:

$$B \rightarrow D^{**}\ell\nu$$

- $D_0^*$  and friends are not what we think they are. We need more spectra (not BF's) to disentangle different resonances.
- The  $D^*$  tail should be correctly handled in MC!



$$B \rightarrow D^{(*)}\eta\ell\nu \quad \& \quad B \rightarrow X_c\ell\nu$$

- We rule out the current “best guess” used to fill the semileptonic gap.  
(So stop saying it could fill the gap, prove it with direct experimental evidence.)



$$B \rightarrow \rho\ell\nu \quad \& \quad B \rightarrow X_u\ell\nu$$

- We provide lineshapes and form factors for all  $B \rightarrow \pi\pi\ell\nu$  decays.
- Looking forward to upcoming lattice results.
- It would be interesting to see LCSR calculations using our parametrization!



$$B_s \rightarrow DK\ell\nu \quad \& \quad B \rightarrow ??\ell\nu$$

- Next on the menu:  $B_s \rightarrow DK\ell\nu$  can be treated similarly to  $B \rightarrow D\pi\ell\nu$ .
- We're always open to suggestions and fruitful collaborations!







Thank you for your attention!



# A tale of two 'gap' models

## Model 1:

Equidistribution of all final state particles in phase space

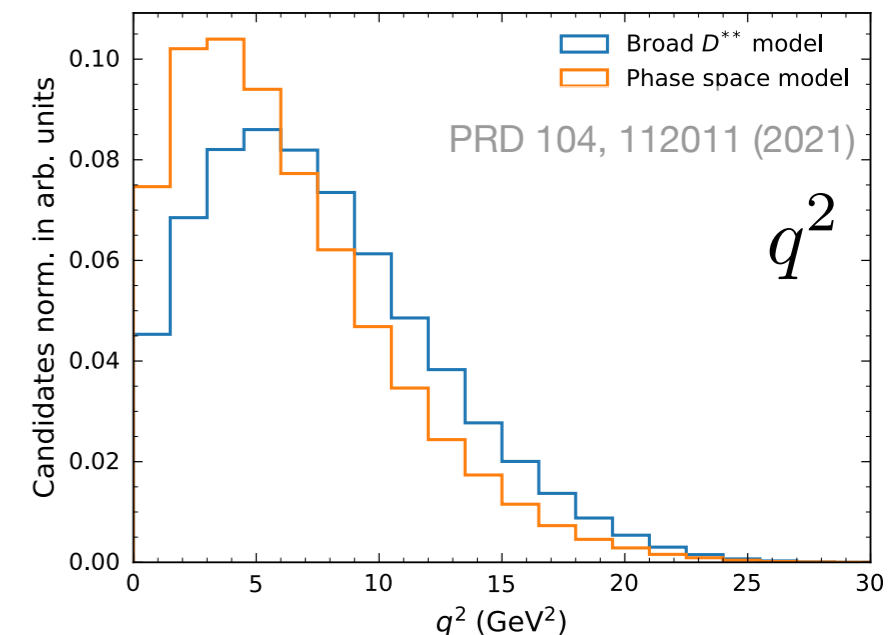
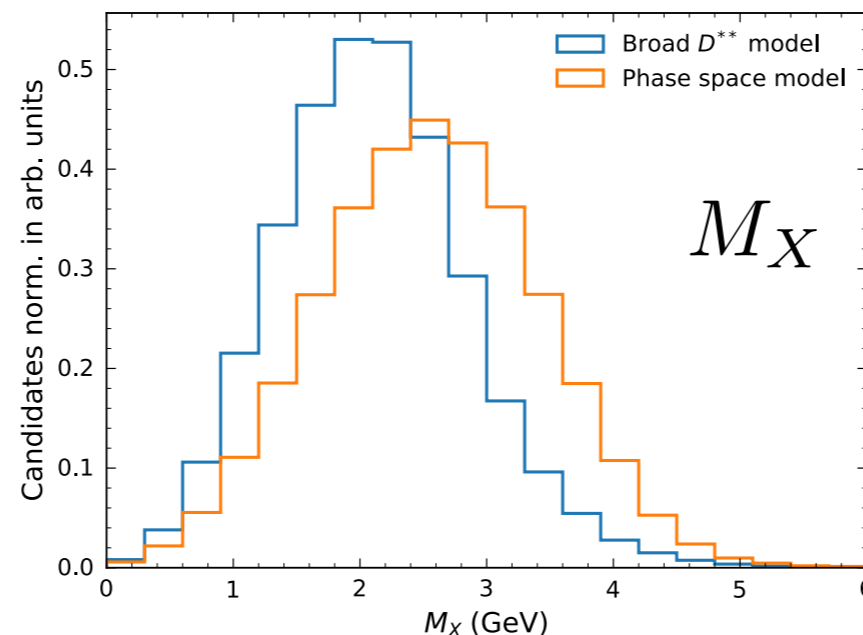
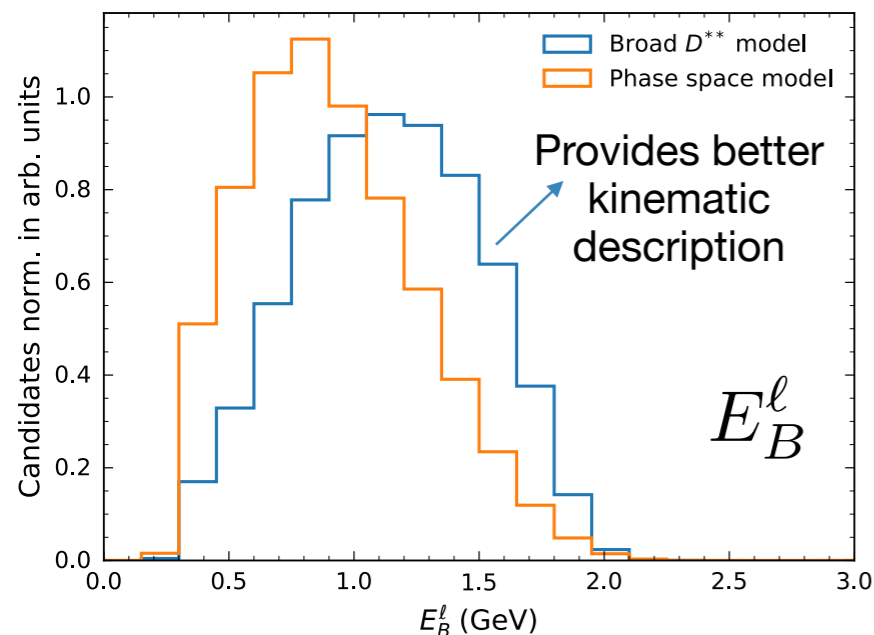
Decay	$\mathcal{B}(B^+)$	$\mathcal{B}(B^0)$
$B \rightarrow D \ell^+ \nu_\ell$	$(2.4 \pm 0.1) \times 10^{-2}$	$(2.2 \pm 0.1) \times 10^{-2}$
$B \rightarrow D^* \ell^+ \nu_\ell$	$(5.5 \pm 0.1) \times 10^{-2}$	$(5.1 \pm 0.1) \times 10^{-2}$
$B \rightarrow D_1 \ell^+ \nu_\ell$	$(6.6 \pm 0.1) \times 10^{-3}$	$(6.2 \pm 0.1) \times 10^{-3}$
$B \rightarrow D_2^* \ell^+ \nu_\ell$	$(2.9 \pm 0.3) \times 10^{-3}$	$(2.7 \pm 0.3) \times 10^{-3}$
$B \rightarrow D_0^* \ell^+ \nu_\ell$	$(4.2 \pm 0.8) \times 10^{-3}$	$(3.9 \pm 0.7) \times 10^{-3}$
$B \rightarrow D_1' \ell^+ \nu_\ell$	$(4.2 \pm 0.9) \times 10^{-3}$	$(3.9 \pm 0.8) \times 10^{-3}$
$B \rightarrow D \pi \pi \ell^+ \nu_\ell$	$(0.6 \pm 0.9) \times 10^{-3}$	$(0.6 \pm 0.9) \times 10^{-3}$
$B \rightarrow D^* \pi \pi \ell^+ \nu_\ell$	$(2.2 \pm 1.0) \times 10^{-3}$	$(2.0 \pm 1.0) \times 10^{-3}$
$B \rightarrow D \eta \ell^+ \nu_\ell$	$(4.0 \pm 4.0) \times 10^{-3}$	$(4.0 \pm 4.0) \times 10^{-3}$
$B \rightarrow D^* \eta \ell^+ \nu_\ell$	$(4.0 \pm 4.0) \times 10^{-3}$	$(4.0 \pm 4.0) \times 10^{-3}$
$B \rightarrow X_c \ell \nu_\ell$	$(10.8 \pm 0.4) \times 10^{-2}$	$(10.1 \pm 0.4) \times 10^{-2}$

## Model 2:

Decay via intermediate broad  $D^{**}$  state

Decay	$\mathcal{B}(B^+)$	$\mathcal{B}(B^0)$
$B \rightarrow D_0^* \ell^+ \nu_\ell$ ( $\hookrightarrow D \pi \pi$ )	$(0.03 \pm 0.03) \times 10^{-2}$	$(0.03 \pm 0.03) \times 10^{-2}$
$B \rightarrow D_1^* \ell^+ \nu_\ell$ ( $\hookrightarrow D \pi \pi$ )	$(0.03 \pm 0.03) \times 10^{-2}$	$(0.03 \pm 0.03) \times 10^{-2}$
$B \rightarrow D_0^* \pi \pi \ell^+ \nu_\ell$ ( $\hookrightarrow D^* \pi \pi$ )	$(0.108 \pm 0.051) \times 10^{-2}$	$(0.101 \pm 0.048) \times 10^{-2}$
$B \rightarrow D_1^* \pi \pi \ell^+ \nu_\ell$ ( $\hookrightarrow D^* \pi \pi$ )	$(0.108 \pm 0.051) \times 10^{-2}$	$(0.101 \pm 0.048) \times 10^{-2}$
$B \rightarrow D_0^* \ell^+ \nu_\ell$ ( $\hookrightarrow D \eta$ )	$(0.396 \pm 0.396) \times 10^{-2}$	$(0.399 \pm 0.399) \times 10^{-2}$
$B \rightarrow D_1^* \ell^+ \nu_\ell$ ( $\hookrightarrow D^* \eta$ )	$(0.396 \pm 0.396) \times 10^{-2}$	$(0.399 \pm 0.399) \times 10^{-2}$

(Assign 100% BR uncertainty in systematics covariance matrix)





# The $\omega$ phase-shift

- Since the phases are known for both the  $\rho$  and  $\omega$ , we follow standard procedures based on the works of Leutwyler to treat the  $\rho$ - $\omega$  interference.

