

Bundesministerium für Bildung und Forschung



tu technische universität



$b \rightarrow c \ell \nu$ global Wilson coefficent fit with LHCb and Belle II

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> **Challenges in Semileptonic B decays** Vienna 23-27 September 2024

Outline

1 Introduction

(2) LFU tests at LHCb and Belle II





Introduction

- Semileptonic B decays are ideal for testing Lepton Flavour Universality(LFU)
- R(D) and R(D*) measurements report a 3.1 σ discrepancy with the Standard Model

$$R(D^*) = \frac{BR(B^0 \to D^* \tau \nu)}{BR(B^0 \to D^* \mu \nu)}$$

- Measurements from LHCb, Belle and Babar avaialable
- R(D*) measurements from LHCb use $\tau \rightarrow \mu\nu\nu$ or $\tau \rightarrow \pi^{+}\pi^{-}\pi^{+}\nu$



The goal:

Measure potential New Physics scenarios contributions

LFU tests in semileptonic *B* decays at LHCb

Muonic τ decay

 $R(D^{*+})$ Run I (2015) $0.336 \pm 0.027 (stat) \pm 0.030 (syst)$

 $R(D^0)$ vs $R(D^*)$ Run I (2023) $R(D^*)=0.281\pm0.018(stat)\pm0.024(syst)$ $R(D^0) = 0.441 \pm 0.060(stat) \pm 0.066(syst)$ > PRL 131 111802

Hadronic τ decay $R(D^{*+})$ Run I (2018) $0.291 \pm 0.019(stat) \pm 0.026(syst) \pm$ 0.013(ext)▶ PRL 120 171802

 $R(D^{*+})$ part Run II (2023) $0.247 \pm 0.015(stat) \pm 0.015(svst) \pm$ 0.012(ext) PRD 108 012018

 $R(D^+)$ vs $R(D^{*+})$ Run II (2024) $R(\Lambda_c^+)$ Run I (2022) $R(D^{*+})=0.402\pm0.081(stat)\pm0.085(syst)$ 0.242 $\pm0.026(stat)\pm0.040(syst)\pm$ $R(D^+)=0.249\pm0.043(stat)\pm0.047(syst)$ 0.059(ext) \rightarrow PRL 128 191803

▶ arXiv:2406.03387

▶ PRL 120 121801

 $R(J/\psi)$ Run I (2018) 0.71±0.17(stat)±0.18(syst)

 D^* polarisation (2023) $0.242 \pm 0.026(stat) \pm 0.040(syst) \pm$ 0.059(ext) LHCb-PAPER-2023-020

Muonic vs hadronic tau decays at LHCb

Muonic τ decay



- High statistics sample
- $R(D^*)$ directly measured
- Multiple missing neutrinos
- Precise background modelling

Hadronic τ decay



- High purity sample: allowed by being able to fully reconstruct the τ vertex
- $R(D^*)$ needs external input
- Low statistics

Complementary analyses that provide independent result

LFU tests in semileptonic *B* decays at Belle II



LFU tests in semileptonic *B* decays at Belle (II)

Belle

Had tag: R(D) vs R(D*) (2015) R(D)=0.375±0.064(stat)±0.026(syst) R(D*)=0.293±0.038(stat)±0.015(syst) ▶ PRD 92 072014

Had tag: P_{τ} and $R(D^*)$ (2018) $P_{\tau}(D^*) = -0.38 \pm 0.51(stat) \pm \frac{0.21}{0.16}(syst)$ $R(D^*) = 0.270 \pm 0.035(stat) \pm \frac{0.028}{0.025}(syst)$ • PRD 97 012004

SL tag: R(D) vs $R(D^*)$ (2020) $R(D)=0.307\pm0.037(stat)\pm0.016(syst)$ $R(D^*)=0.283\pm0.018(stat)\pm0.014(syst)$

• PRD 92 072014

Belle II

Had tag: R(X) (2024) $R(X)=0.228\pm0.016(stat)\pm0.036(syst)$

PRL 132 211804

Had tag: $R(D^*)$ (2024) $R(D^*)=0.262\pm_{0.039}^{0.041}(stat)\pm_{0.032}^{0.035}(syst)$

$R(D^*)$ measurement strategy

$$R(D^*) = rac{N_{sig}}{N_{norm}} rac{\epsilon_{norm}}{\epsilon_{signal}}$$

- Signal/norm yield ratio determined from fit to data
- Signal/norm efficiency ratio determined from simulation: assuming SM scenarios
- Form factors varying in the templates or a systematic uncertainty assigned (depending on analysis setup)
- Largest systematic uncertainties coming from the form factor parameterisation, background modelling and simulation sample size
- Depending on the choice of τ decay different backgrounds to model

New Physics

• We can use operators with unknown coupling constants and write them in an effective Hamiltonian



- C_i^{NP} are the Wilson coefficients that describe the NP effects
- *O_i* are effective operators that can be of a vector, scalar or tensor type



New Physics measurement strategies

- Model the New Physics effects in the fitting template
 - Hammer usage crucial
- More sensitivity to NP using the shape of the angles as fitting variables: see <a>talk by L.Grillo at this workshop
- Measure **directly Wilson coefficients** : flexible to use NP model or be model independent
- Measure angular coefficients
- Currently no published analyses, several in progress at LHCb

HAMMER - Helicity Amplitude Module for Matrix Element Reweightingt

 Tool that weights a MC sample from the generation amplitude to a new desired one
 Hammer
 arXiv:2002.00020v2

Theoretical approach

The decay rate $(B \rightarrow X l \nu_l)$:

$$\frac{d\Gamma}{dq^2} = \frac{G_F^2}{(2\pi)^3} V_{ij}^2 \frac{(q^2 - m_I^2)^2 p_X}{12m_B^2 q^2} (H_+^2(q^2) + H_-^2(q^2) + H_0^2(q^2)(1 + \frac{m_I^2}{2q^2}) + \frac{3}{2} \frac{m_I^2}{qq^2} H_s^2(q^2))$$

where $H_i(q^2)$ are the helicity amplitudes.

Reweighting to New Physics scenarios, e.g. by adding extra scalar, vector or tensor couplings can be done with the weight vector for each event calculated as:

$$\omega_i = \frac{\Gamma_{old}}{\Gamma_{new}} \frac{d^n \Gamma_{new}/dx}{d^n \Gamma_{old}/dx}$$

where Γ_{old} is the decay rate for the model implemented in the Monte Carlo and Γ_{new} is the the decay rate for updated model

Wilson coefficient fit strategy

- Use Hammer to directly fit the Wilson coefficients
- Model the signal and background(where possible) template with Hammer
- Allows for floating the Wilson coefficients and as well the form factor parameters in the fit
- Produce multidimensional templates with a choice of kinematic variables
- Easy to fit for specific New Physics model
- Provide sensitivity of each NP operator
- BLPR form factor parameterisation in use (calculations by Bernlochner et al. arXiv:1703.05330v4), using both the leading and $O(\Lambda_{QCD}/m_b)$ and subleading Isgur-Wise function)

Combine several analyses to measure WCs?

Project proposal

- Combine statistical power of several channels and observables from LHCb and Belle II
- Similar measurements available within LHCb and Belle II
- A starting point for this could be already published analyses

Leveraging LHCb and Belle II



- Global $b
 ightarrow c \ell
 u$ fit that avoids biases and SM assumptions
- Share WCs between templates of different measurements and (where possible) background models and systematics

Combined WCs analysis strategy

• A scan over the WCs is performed: combined multiple fits in different phase space regions

Profile likelihood method

 $\chi^2(\alpha) = -2 \ln L(\alpha)$ where α is the phase space.

With the confidence interval as

$$1 - CL = rac{1}{\sqrt{2}\Gamma(1/2)} \int_{\delta\chi^2}^{\infty} e^{-1/2} t^{-1/2} dt$$



 Different regions of phase space can be scanned separately to paralelize the scan

WC template shapes at LHCb

 $TqLIL - [\bar{c}\chi_L^T \gamma^{\mu} P_L b] [\bar{l}\lambda_L^V \gamma_{\mu} P_L \nu]: \text{ plots for } B^0 \to D^* \tau \nu \text{ from HAMMER}$

0.5 0.45 0.25 0.4 0.35 0.2 0.3 0.15 0.25 0.2 0.1 0.15 0.1 0.05 0.05 6 10 $m^2_{miss} \, [GeV^2]$ q2 [GeV2] 0.28 0.5 0.26 0.45 α LIL = 0.0 0.24 0.4 aLL = 0.1T qLIL = 0.2 0.35 0.22 T aLIL = 0.30.3 $T_qLIL = 0.5$ 0.2 0.25 0.18 0.2 0.16 0.15 0.14 0.1 -0.50.5 -0.50 0.5 $^{-1}$ $\cos(\theta_{i})$ $\cos(\theta_{\downarrow})$

WC template shapes at Belle II

Shape change for selected New Physics in p_{ℓ}^* and m_{miss}^2



▶ arXiv:2002.00020v2

First proof of principle

- LHCb input: Toy MC with $B^0 \rightarrow D^* \tau \nu$ (with $\tau \rightarrow \mu \nu \nu$)
- Belle II input: Toy MC with $B^0 \rightarrow D^* \tau \nu$ (with $\tau \rightarrow \pi / \rho \nu$)





- The combination (dashed pink contour) demonstrates improvement in sensitivity
- No backgrounds or systematic uncertainties considered in the plots

LHCb toy studies: WC with $\tau \rightarrow \mu \nu \nu$

- **Asimov** scan is applied in signal+control configuration to study the sensitivity
- Each scan considers one of the WCs floating in the fit (the Re and Im part)
- The SM contribution is shared to parametrize the yields



- Less sensitive to the scalar WC as compared to vector and tensor
- Im part results with small sensitivity in all WCs

Other background contributions have been excluded in these preliminary studies

LHCb toy studies: WC with $\tau \rightarrow 3\pi\nu$

- Asimov scan is applied to study the sensitivity
- Each scan considers one of the WCs floating in the fit (the Re and Im part)
- A background (non-Hammer reweighted) is added \rightarrow yield not shared parameter



- Less sensitive in scalar WC as compared with vector and tensor
- Similar Im and Re WC sensitivity

NOTE: Difference in Hammer reweighting between muonic and hadronic τ : here we reweight the τ vertex with RCT FF parameterization

LHCb toy studies: LHCb combination

- **Complementary** measurements from $\tau \rightarrow 3\pi\nu$ and $\tau \rightarrow \mu\nu\nu$
- NP and SM Wilson Coefficients shared among the two modes
- The **muonic mode** constraints the Re part of the WCs (higher statistics)
- The **hadronic** mode constraints the Im part of the WCs (τ vertex weight)



Not a final sensitivity:

- Hammer reweighting to be aligned between samples
- Full background model to be added

What we want to provide

- If there is NP the global WC fit could provide consistent picture among observables
- Not possible to interpret consistently our current $R(D^{(*)})$ measurements: they contain SM templates that could result with biases

How will this proposal evolve

- 1. Each experiment should **continue to perform SM measurements** and pursue their own physics programme
- 2. Concentrate at the LHCb or Belle combination only at the moment
- 3. Combine the Belle II and LHCb results

Summary and next steps

- We are working on a **combined WCs** fit in $b \rightarrow c \tau \nu$ decays using several LHCb + Belle II analyses
- Currently setting up the framework and studying sensitivity of the WCs
- The framework allows to **combine** multiple analysis
- Preliminary sensitivity studies here done with hadronic and muonic τ modes from LHCb

Next steps:

- Full description of the background in the toy studies needs to be added
- Carry on working with simulation to perform a full sensitivity study
- The framework can be used to setup a combined measurement with Belle once all the inputs are confirmed

BACKUP

BGL and **CLN** form factor parameterisation

$$\frac{d\Gamma(B \to D^*\ell\nu)}{dq^2} = \frac{G_F^2 |V_{cb}^2| |\eta_{EW}|^2 |\vec{p}|^2 q^2}{96\pi^3 m_{B^0}^2} \left(1 - \frac{m_\ell^2}{q^2}\right) \times \left[\left(|\mathsf{H}_+|^2 + |\mathsf{H}_-|^2 + |\mathsf{H}_0|^2\right) \left(1 - \frac{m_\ell}{2q^2}\right) + \frac{3}{2} \frac{m_\ell^2}{q^2} |\mathsf{H}_t|^2 \right]$$

Helicity amplitudes in BGL () PRL 74 4603 (1995)

1 N

$$\begin{split} f(z) &= \frac{1}{P_{1*}(z)\phi_{f}(z)}\sum_{n=0}^{N}b_{n}z^{n} \\ m_{p*} & H_{\pm}(\omega) = f(\omega) \mp m_{B}m_{D*}\sqrt{\omega^{2}-1}g(\omega) \\ \omega(q^{2}) &= \frac{m_{B}^{2}+m_{D*}^{2}-q^{2}}{2m_{D*}m_{B}} & H_{0}(\omega) = \frac{F_{1}(\omega)}{\sqrt{q^{2}}} \\ H_{1}(\omega) &= m_{B}(\frac{r\omega^{2}-1}{1+r^{2}-2r\omega}F_{2}(\omega)) \\ H_{1}(z) &= \frac{1}{P_{1-}(z)\phi_{F_{1}}(z)}\sum_{n=0}^{N}a_{n}z^{n} \\ F_{1}(z) &= \frac{1}{P_{0-}(z)\phi_{F_{1}}(z)}\sum_{n=0}^{N}a_{n}z^{n} \\ \end{split}$$

Helicity amplitudes in CLN (Nucl. Phys. B 530 1 (1998)

$$\begin{split} H_{\pm}(\omega) &= m_{B}\sqrt{(r)}(\omega+1)h_{A_{1}}(\omega)\Big[1\mp\sqrt{\frac{\omega-1}{\omega+1}}R_{1}(\omega)\Big] & h_{A_{1}}(\omega) = h_{A_{1}}(1)[1-8\rho_{ps}^{2}z(\omega)+(53\rho_{Ds}^{2}-15)z^{2}(\omega)-(231\rho_{Ds}^{2}-91)z^{3}(\omega)] \\ H_{0}(\omega) &= m_{B}\sqrt{(r)}(\omega+1)\frac{1-r}{\sqrt{(q^{2})}}h_{A_{1}}(\omega)\Big[1+\frac{\omega-1}{1-r}(1-R_{2}(\omega))\Big] & R_{1}(\omega) = R_{1}(1)-0.12(\omega-1)+0.05(\omega-1)^{2} \\ R_{2}(\omega) &= R_{2}(1)+0.11(\omega-1)+0.06(\omega-1)^{2} \\ R_{0}(\omega) &= R_{0}(1)-0.11(\omega-1)+0.01(\omega-1)^{2} \end{split}$$

WC template shapes at Belle II

Change of selection efficiency of $B\to D^*\tau\nu$ vs New Physics contribution



