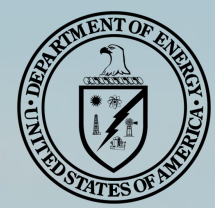


Supported in part by:



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Jets at STAR

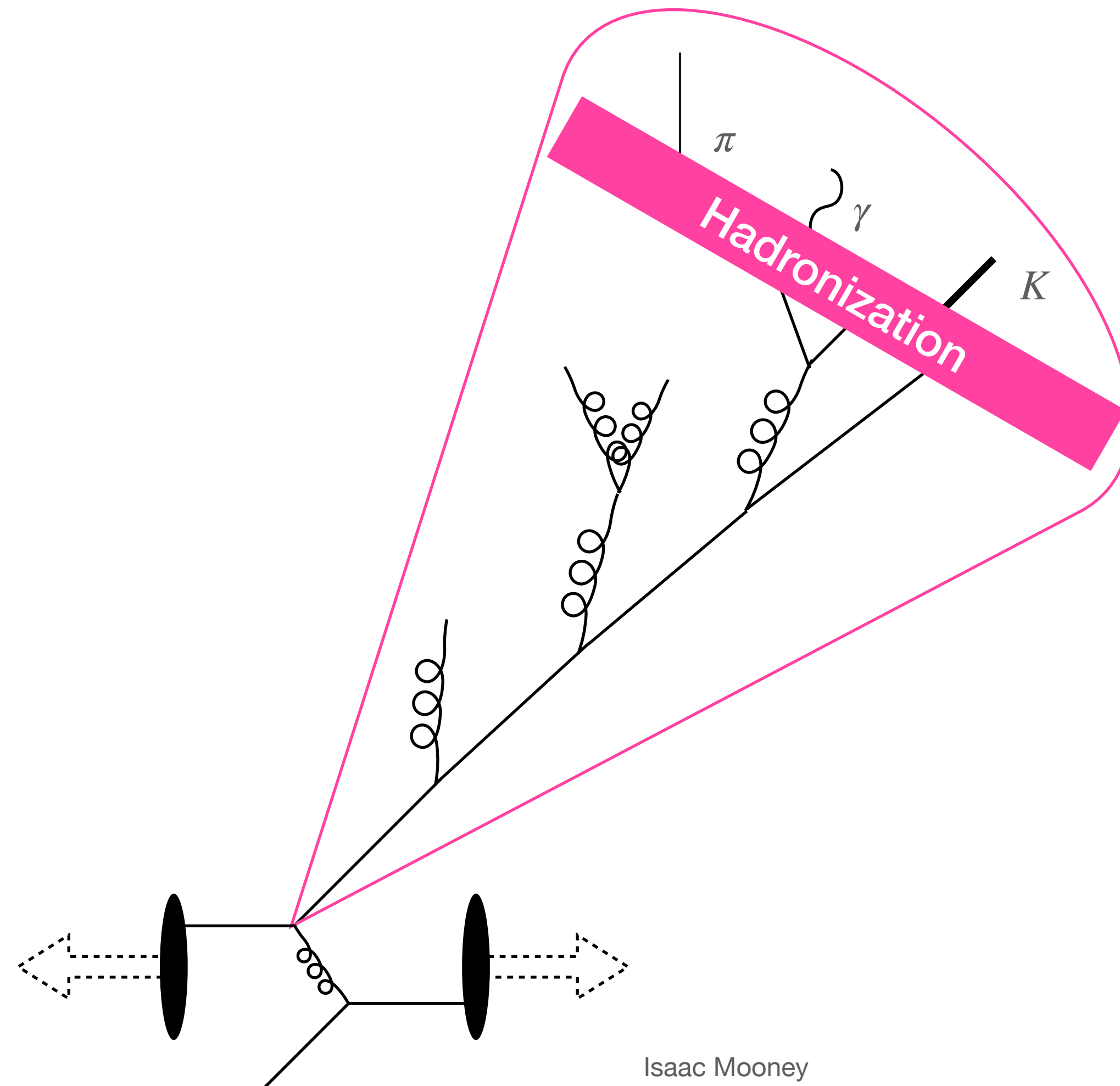
39th Winter Workshop on Nuclear Dynamics

Jackson, Wyoming
February 12, 2024

Isaac Mooney (Yale University, BNL) for the STAR Collaboration isaac.mooney@yale.edu

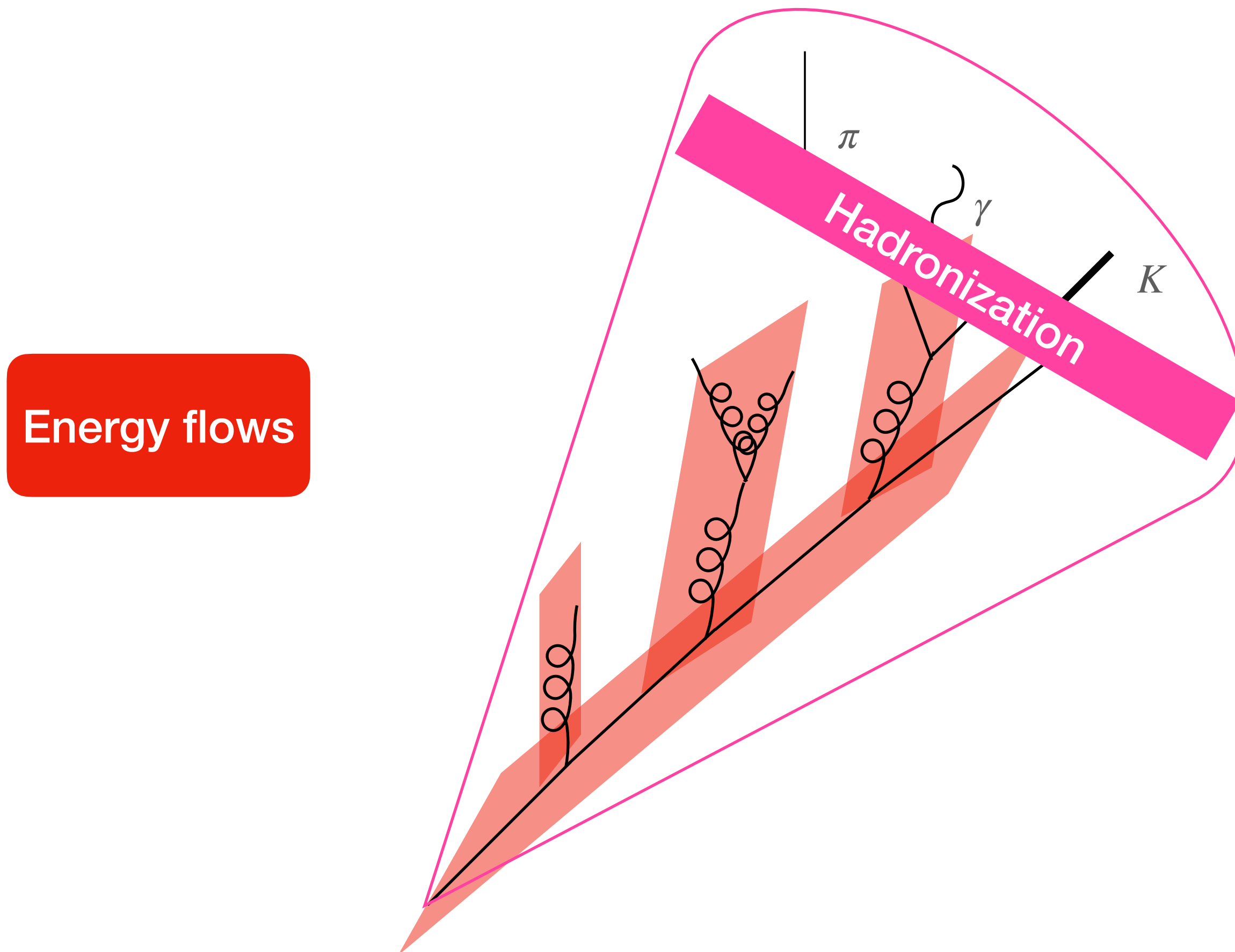
How to understand jet evolution in vacuum

Two ways: the How and the What



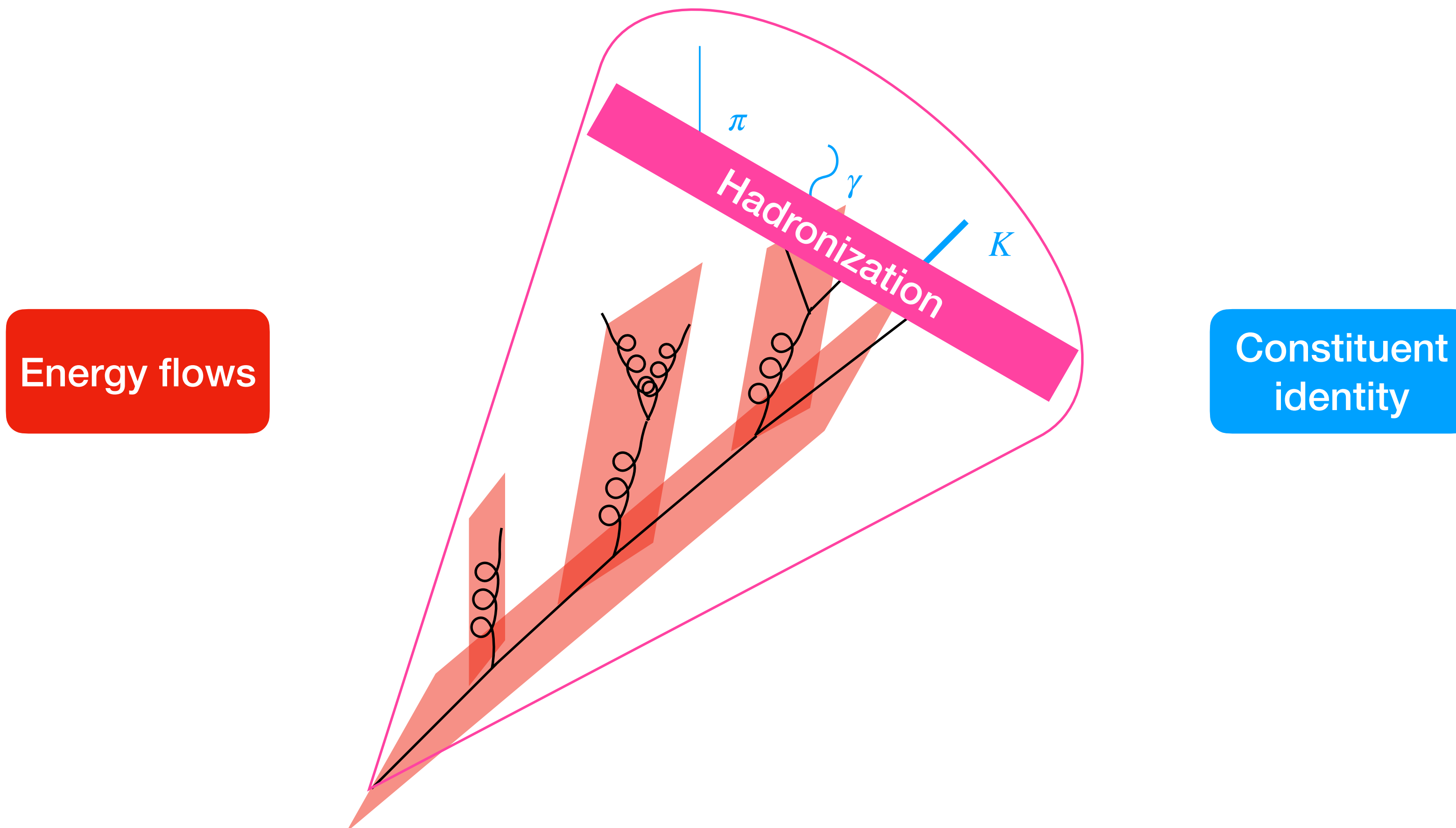
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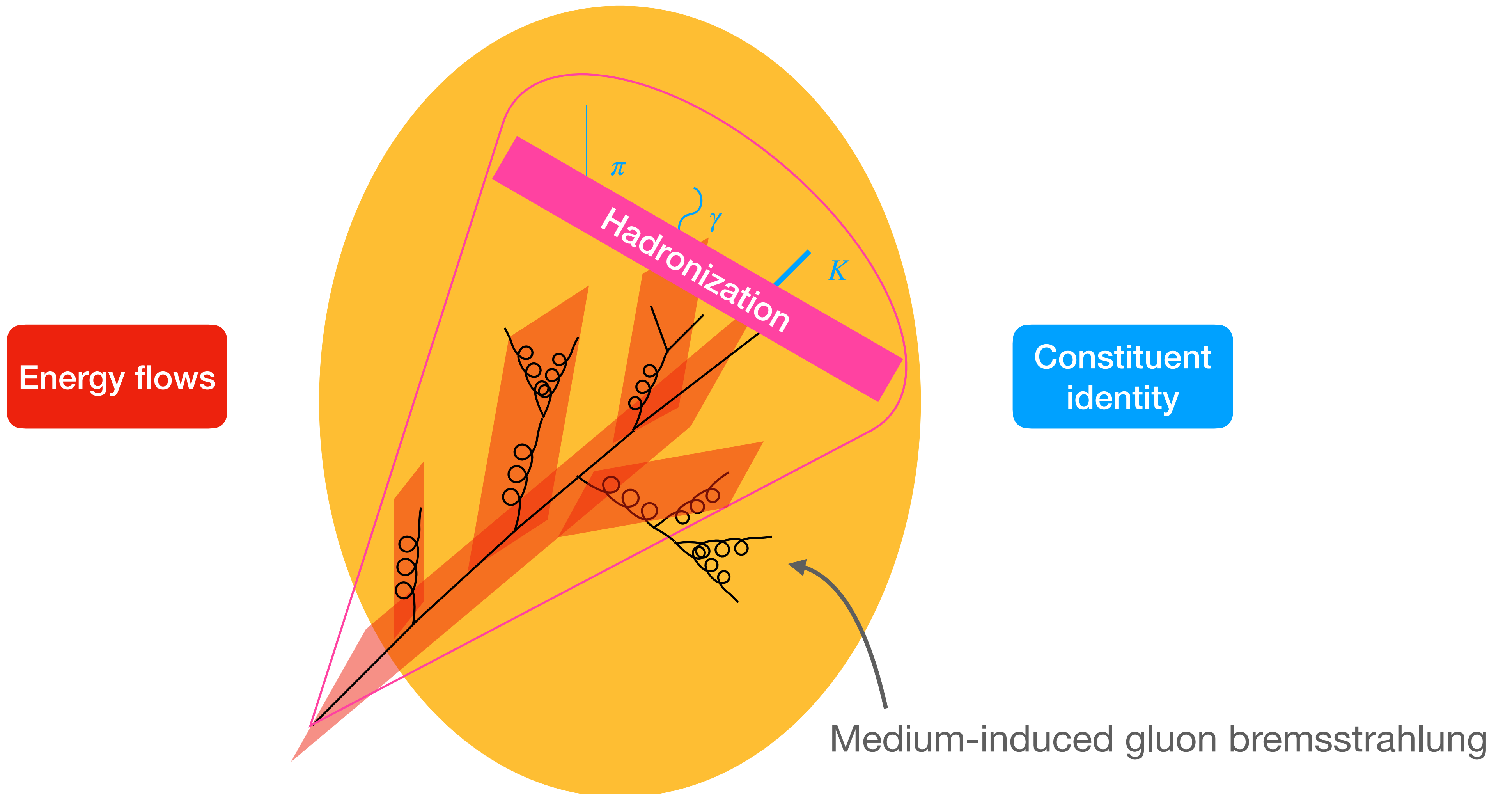
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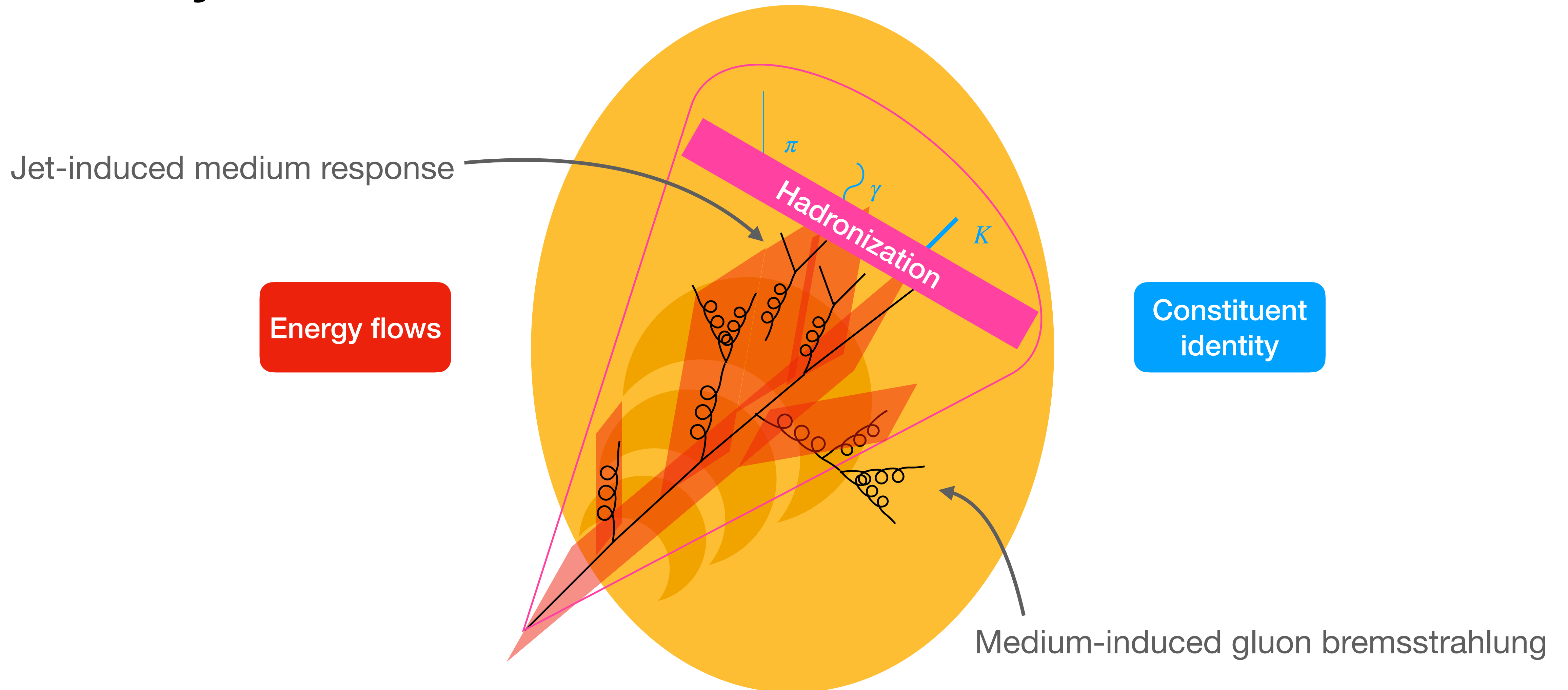
How to understand jet evolution in media

Two ways: the How and the What



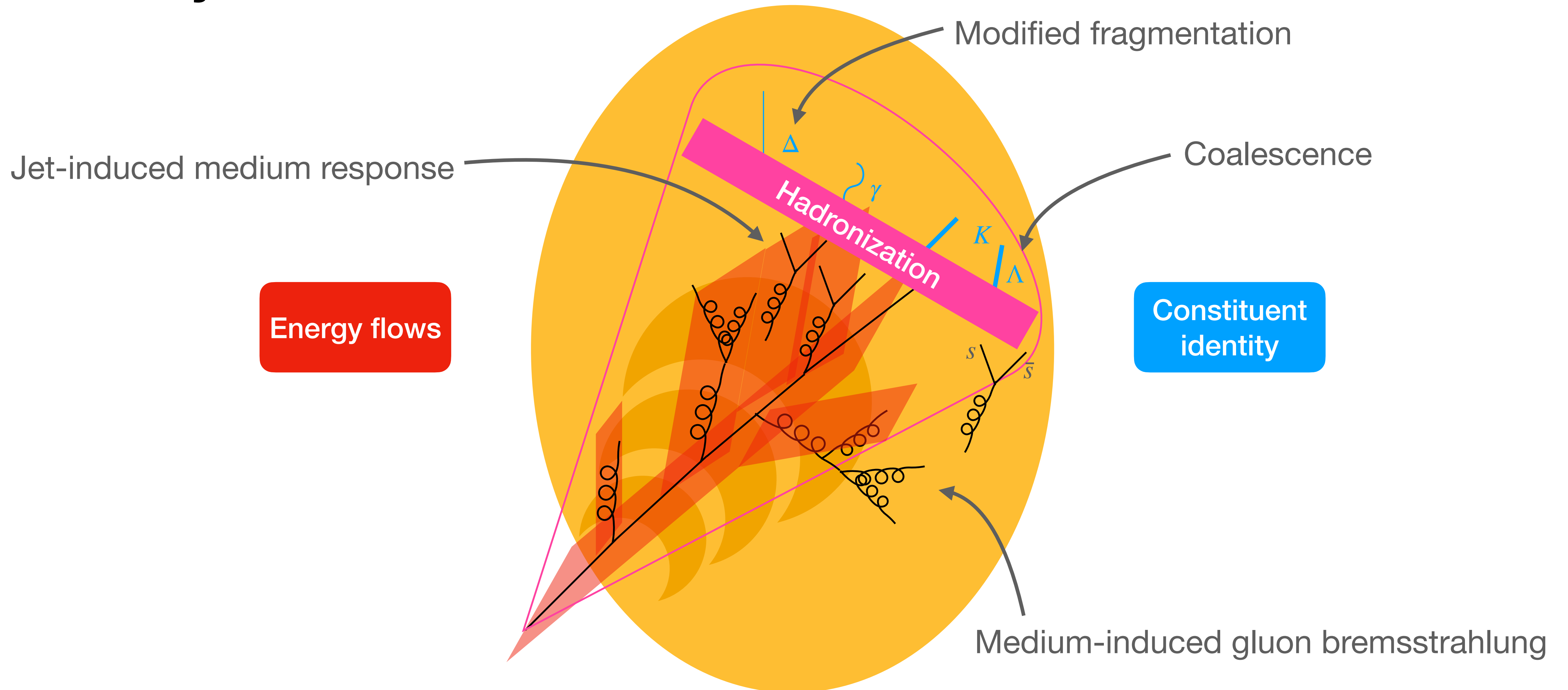
How to understand jet evolution in media

Two ways: the How and the What



How to understand jet evolution in media

Two ways: the How and the What



Solenoidal Tracker at RHIC (STAR)

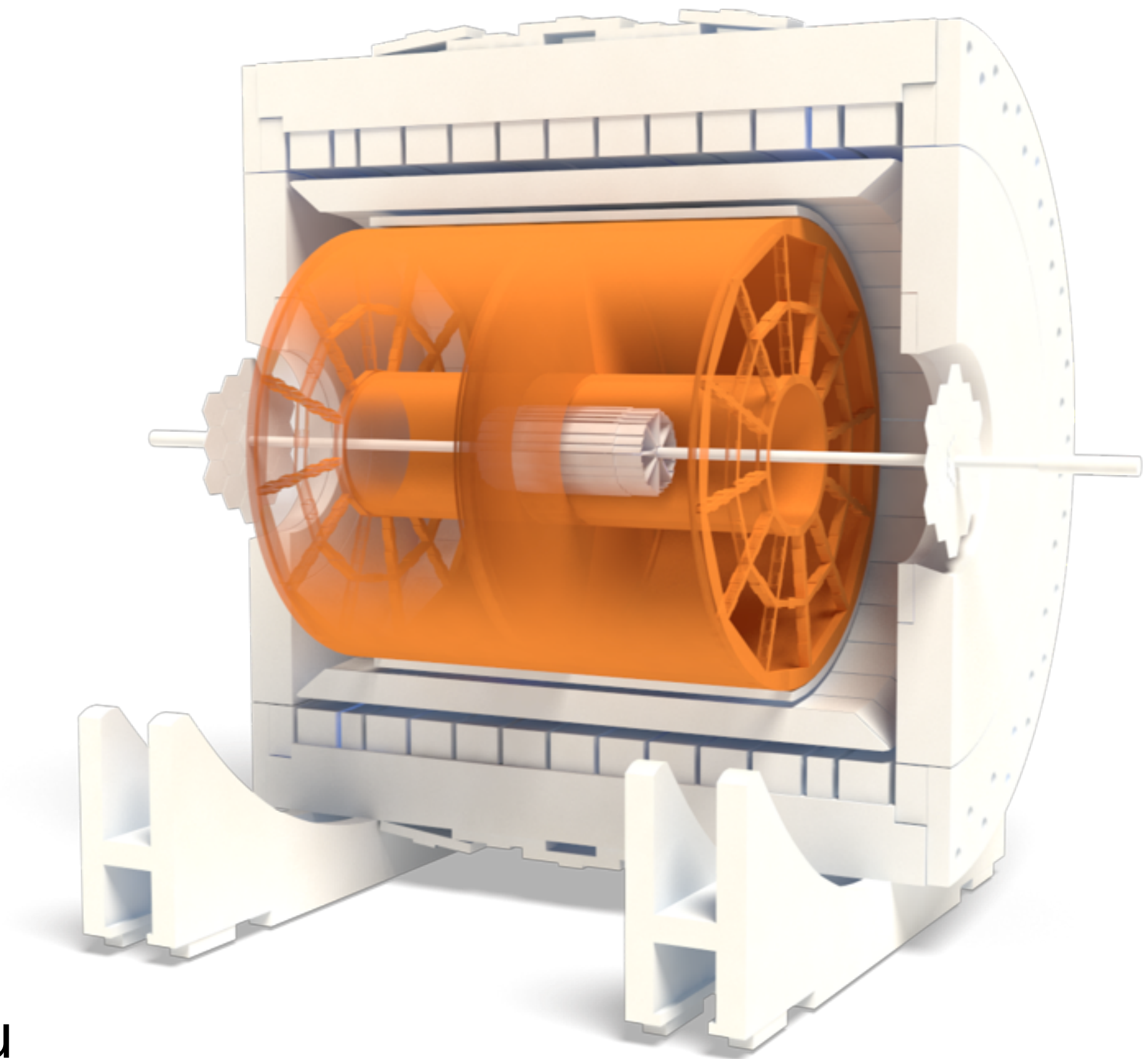
Main subdetectors

Relativistic Heavy Ion Collider (RHIC)
 collides $p+p$, $p+Au$, $O+O$, $Zr+Zr$, $Ru+Ru$, $Au+Au$, etc.
 beams at $\sqrt{s_{NN}} = 200$ GeV, etc.

Time Projection Chamber (TPC) [$|\eta| < 1$]:
 momenta of charged tracks + centrality

Barrel Electromagnetic Calorimeter (BEMC) [$|\eta| < 1$]:
 neutral energy deposits + provides online trigger
 (Jet Patch: $E_T^{patch} > 7.4$ GeV, High Tower: $E_T > 4.2$ GeV)

Inner Beam-Beam Counter (iBBC) [$3.4 < |\eta| < 5.0$]:
 forward detector,
 east/Au-going side activity used as centrality proxy in $p+Au$



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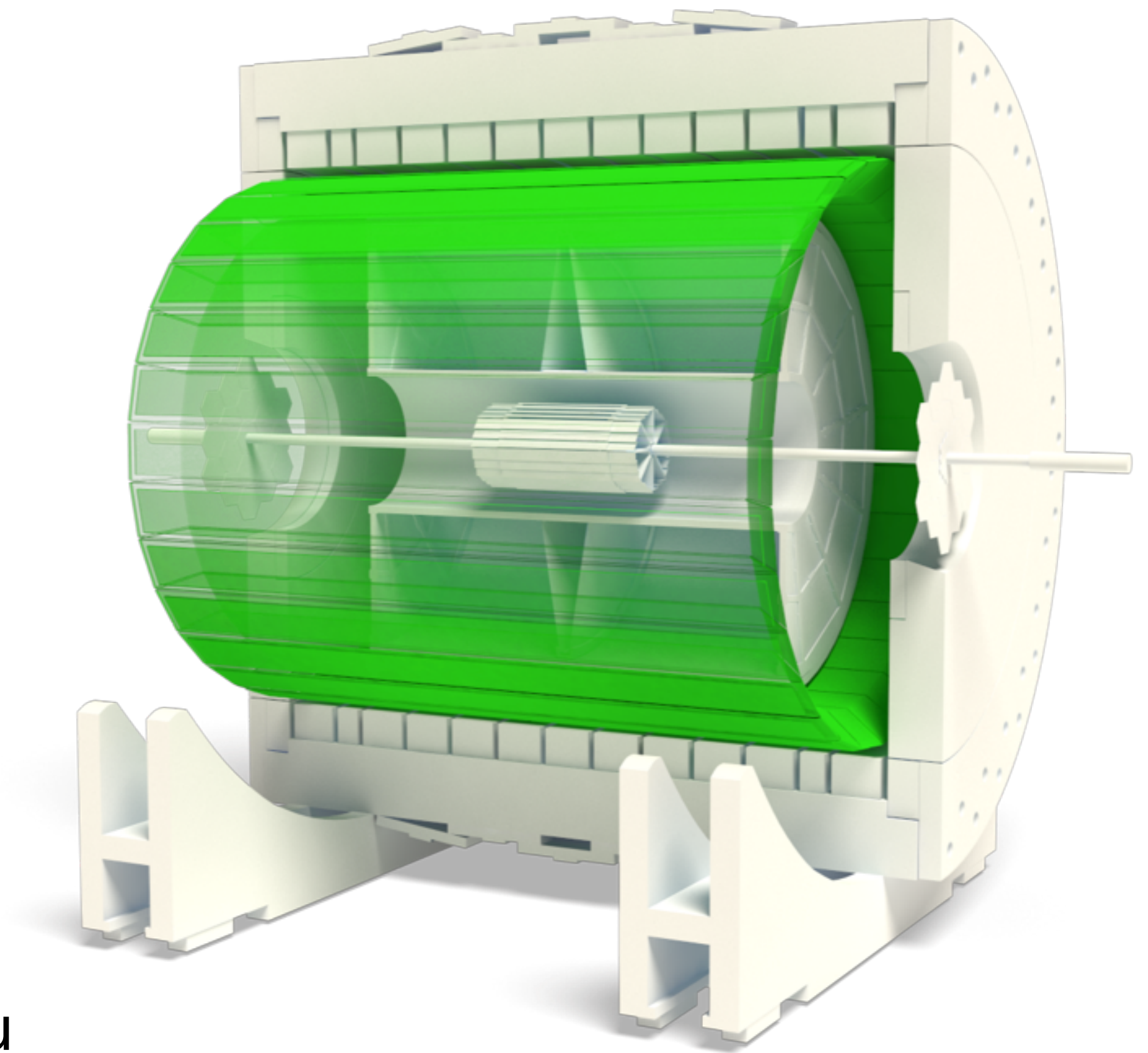
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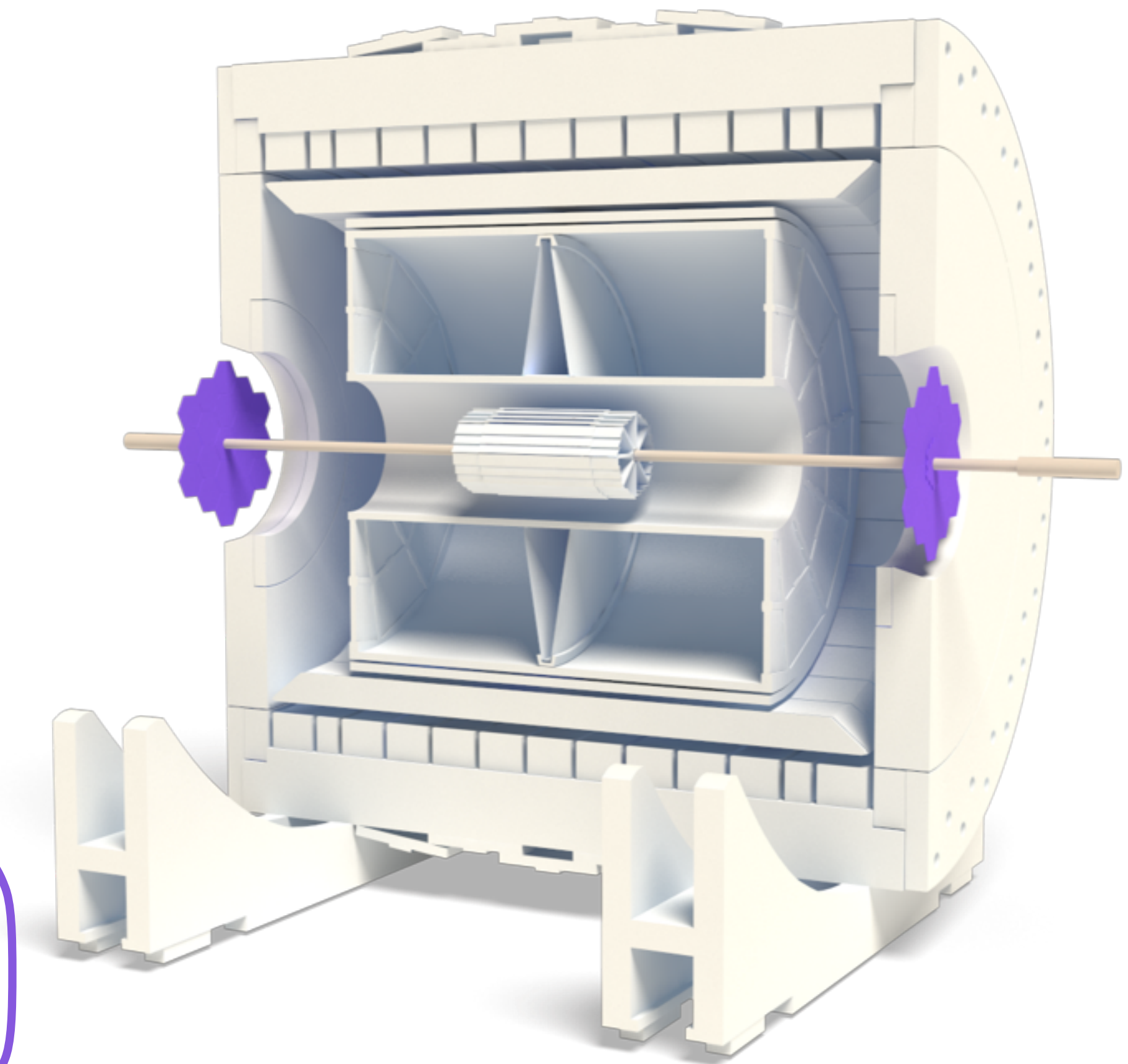
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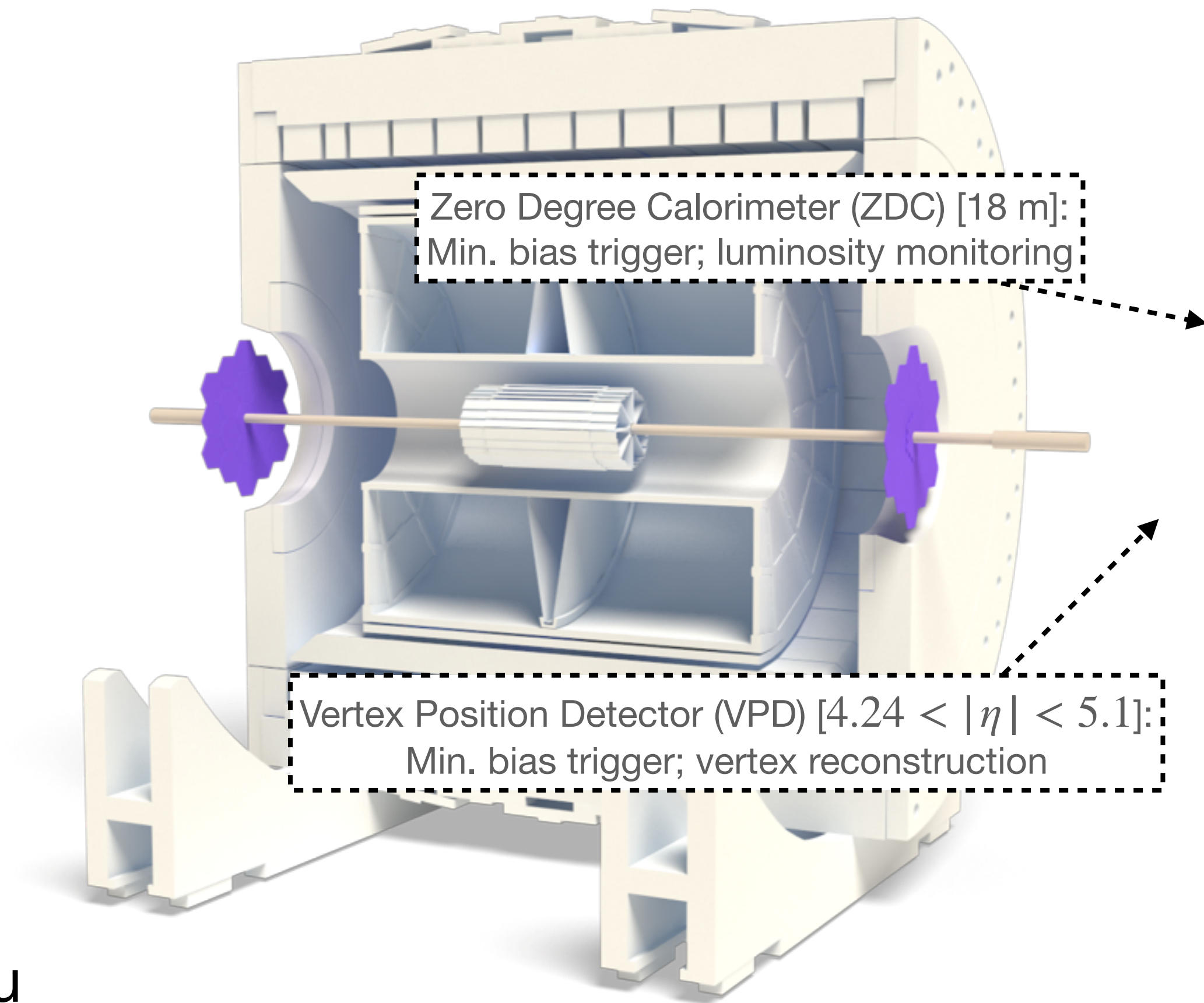
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Precision QCD; exploring the Lund plane
with *multi-dimensional jet substructure*

Separating p-QCD and np-QCD
with *energy correlators*

Energy flows

Path-length dependence of jet energy loss in medium
with *jet anisotropies (with respect to event plane)*

Energy-density dependence of jet energy loss in medium;
angular distribution of radiation in quenched jets
with *inclusive/semi-inclusive jet & high- p_T hadron yields*

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Jet substructure

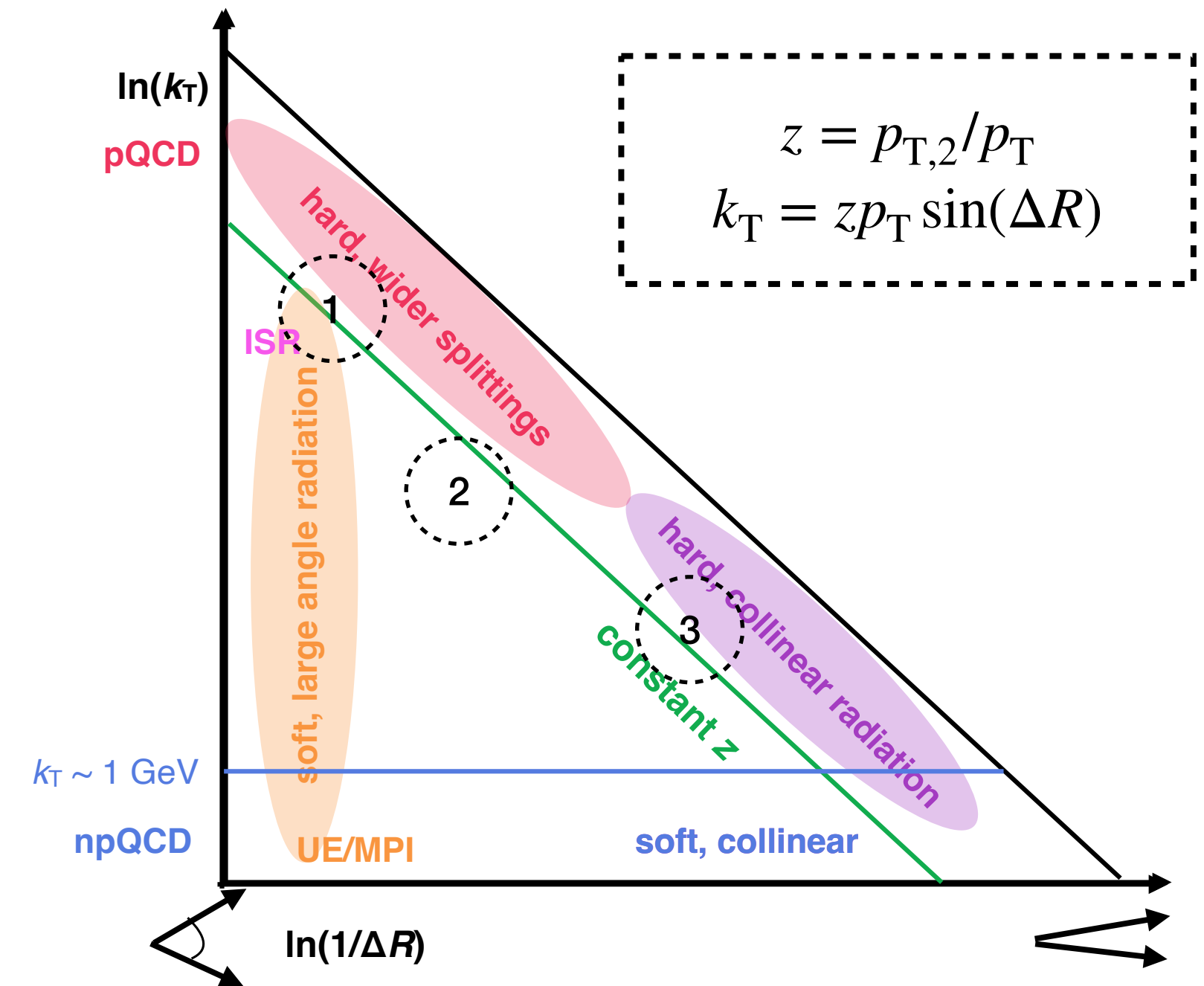
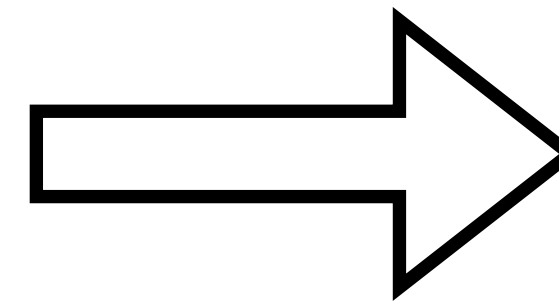
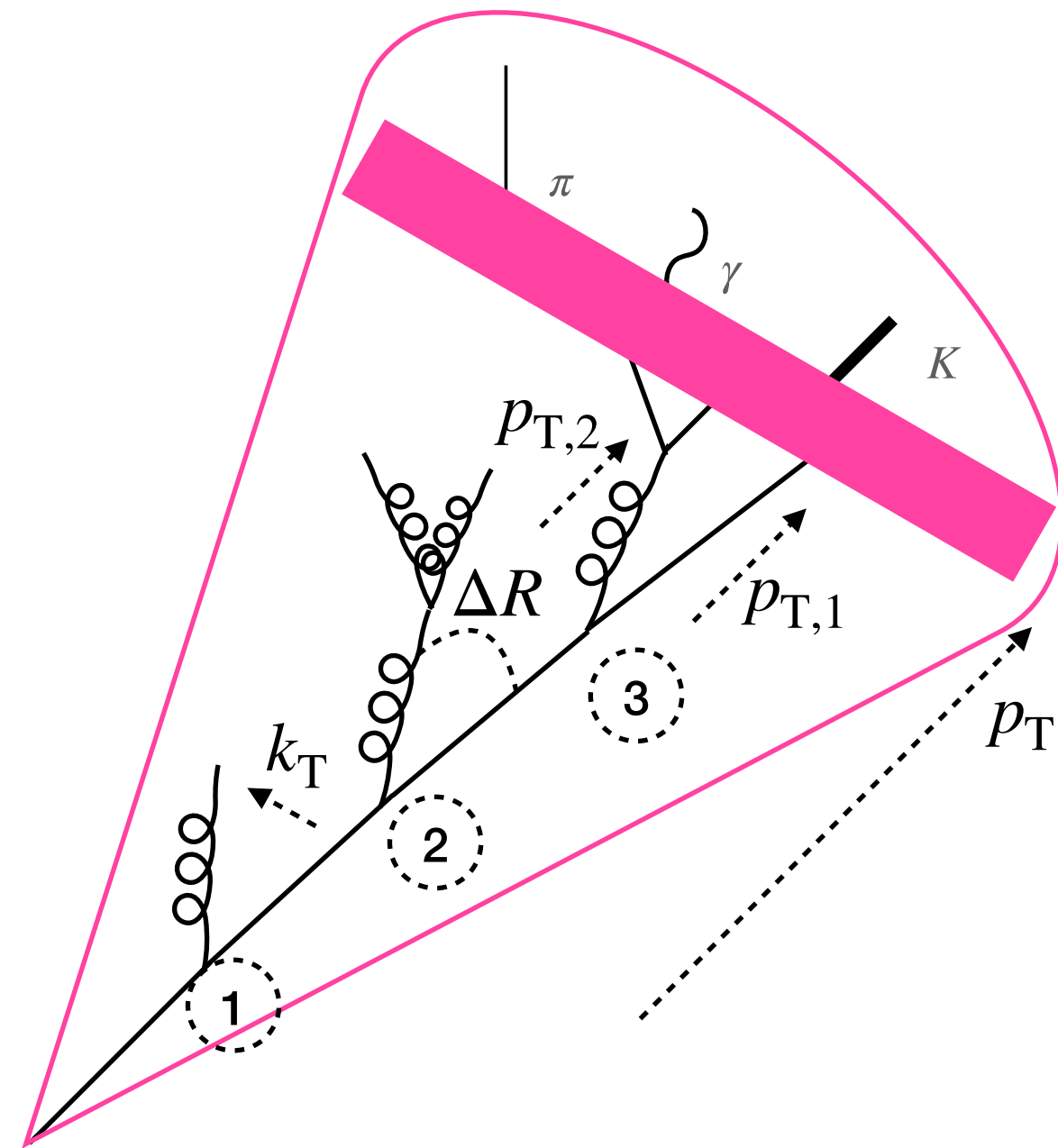
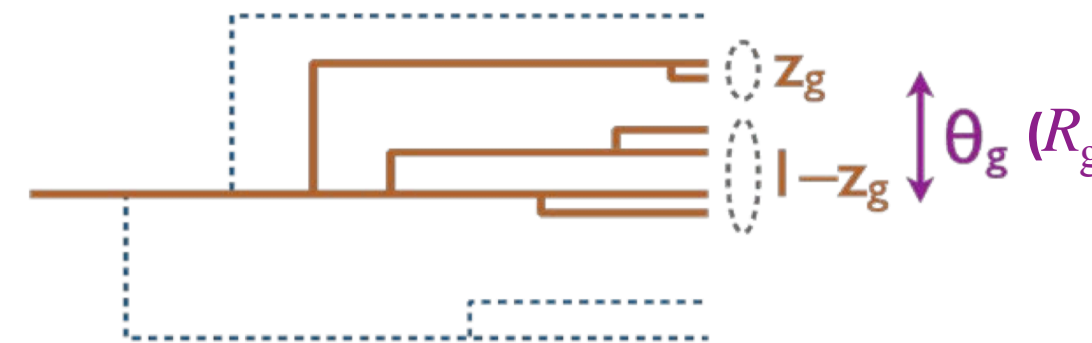


Image: Laura Havener, modified from Andrews et al., [J.Phys.G 47 \(2020\) 6, 065102](https://arxiv.org/abs/1908.07554)

$$\frac{\min(p_{T,i}, p_{T,j})}{p_{T,i} + p_{T,j}} > z_{\text{cut}} \left(\frac{\Delta R_{ij}}{R} \right)^\beta$$

$$z_g = \frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}}$$

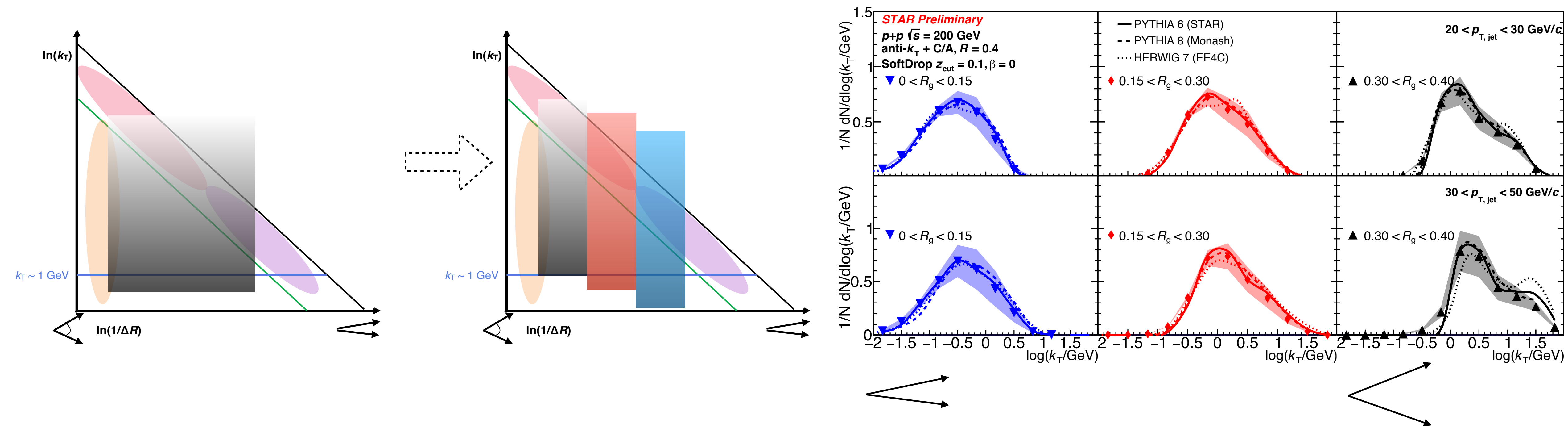


$$M_g = \left| \sum_{i \in J_g} p_i \right|$$

Image: Larkoski, Marzani, Thaler, Xue, [PRL 119 \(2017\) 13, 132003](https://arxiv.org/abs/1603.04914)

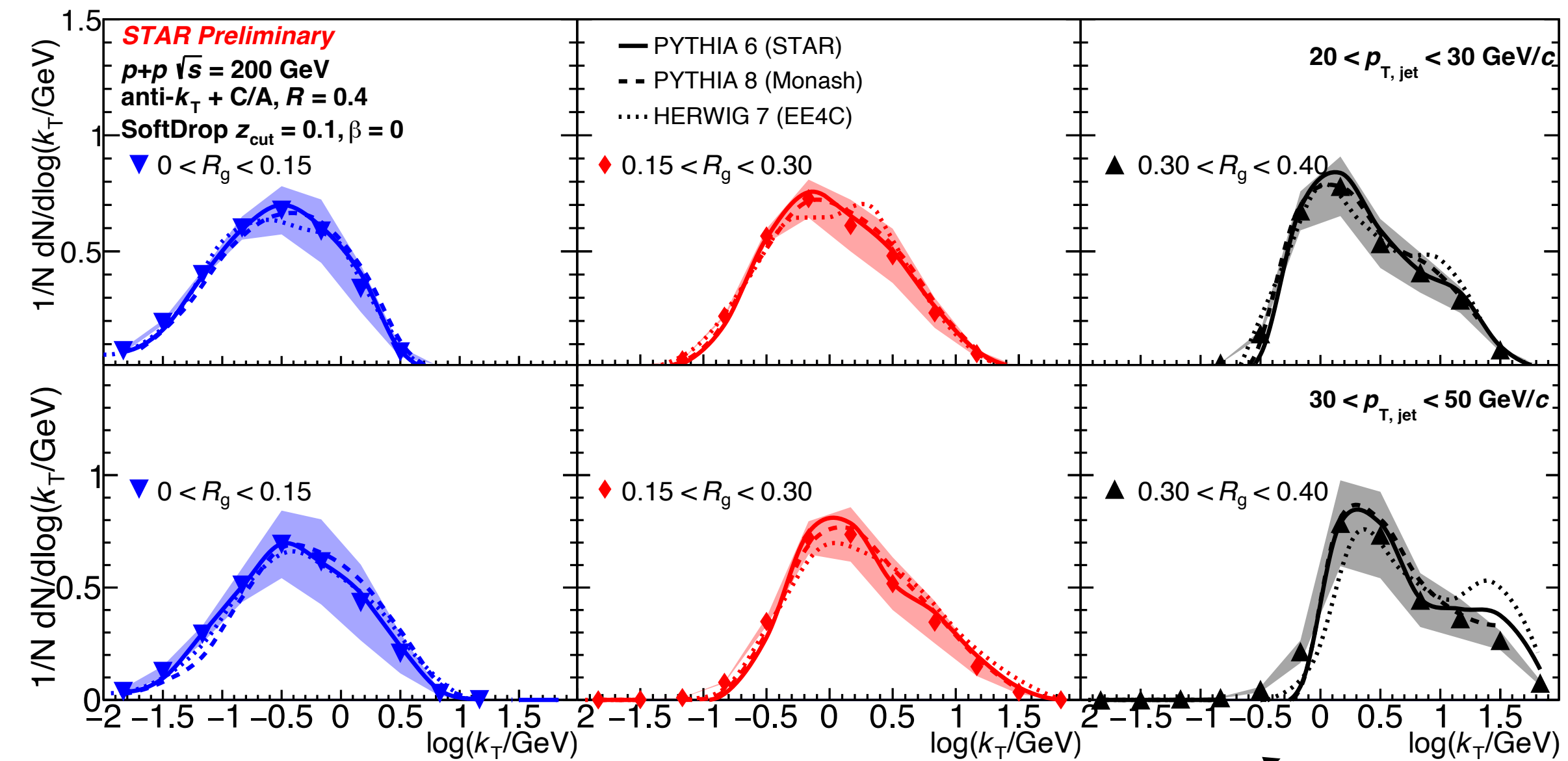
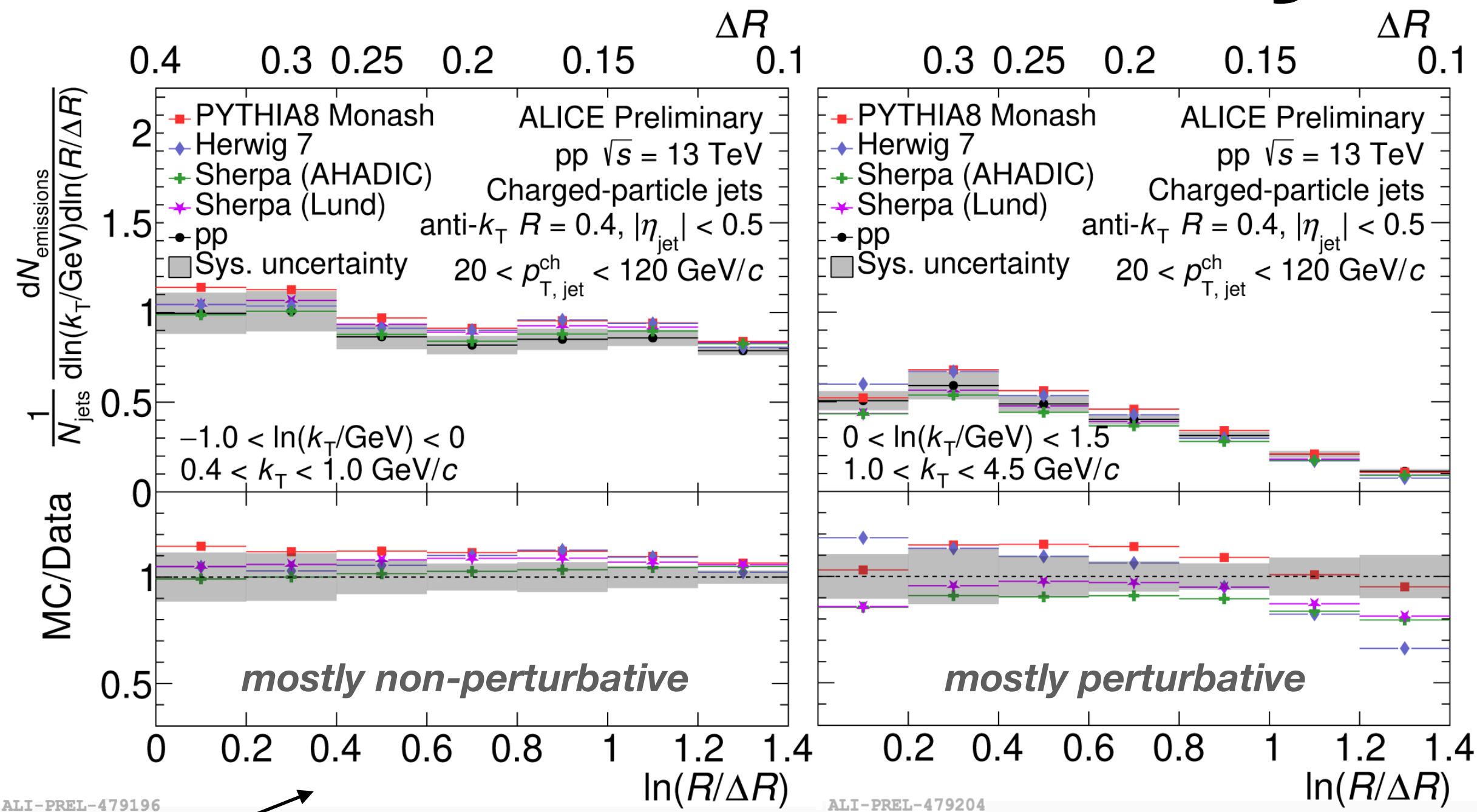
SoftDrop¹ grooming: reduce soft non-perturbative contribution
 → better theoretical control

Multi-dimensional jet substructure



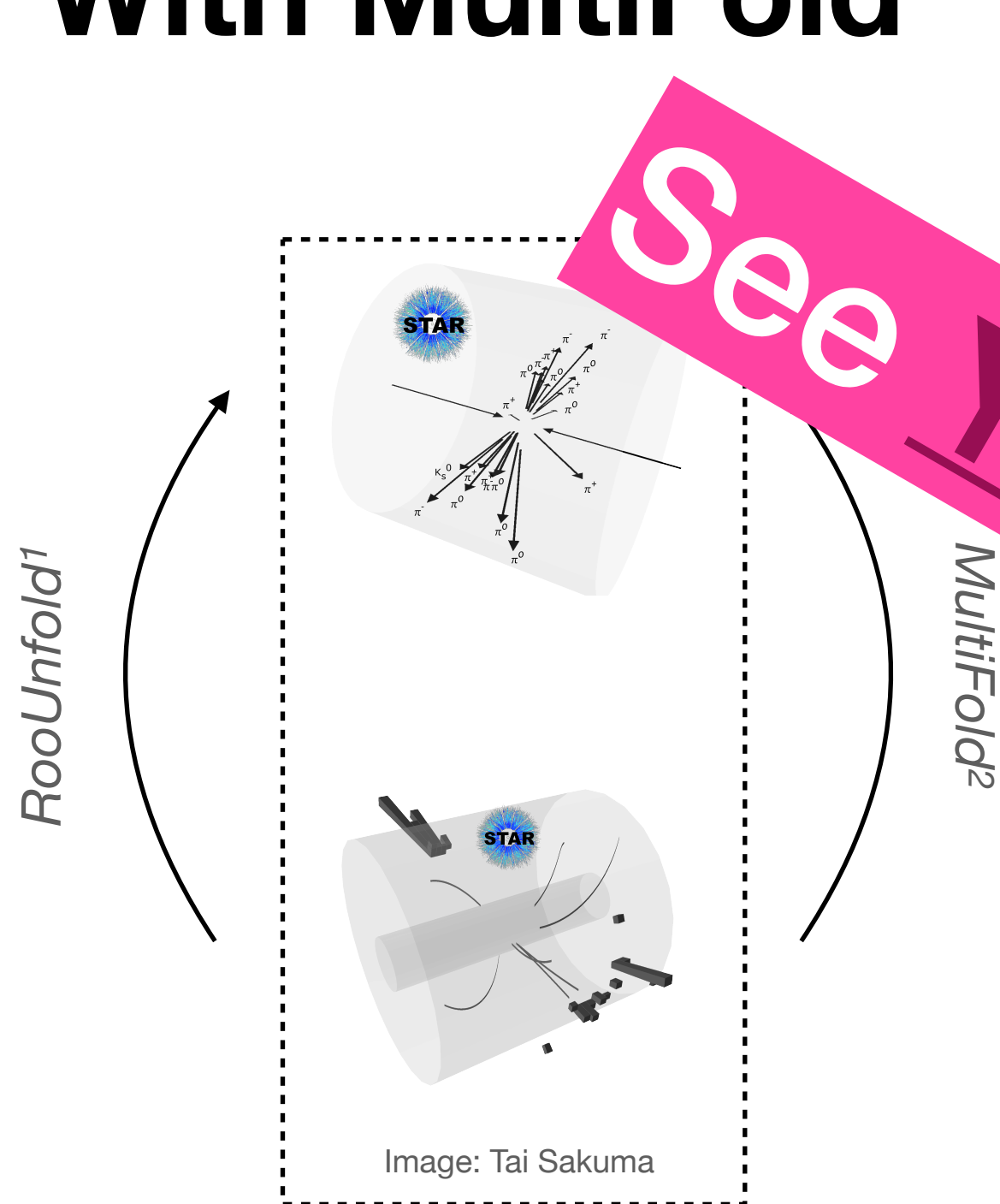
- Now able to make slices in the Lund Plane → more stringent tests of Monte Carlo (MC) models
- Observe: wider splits are harder. MCs in good agreement.

Multi-dimensional jet substructure

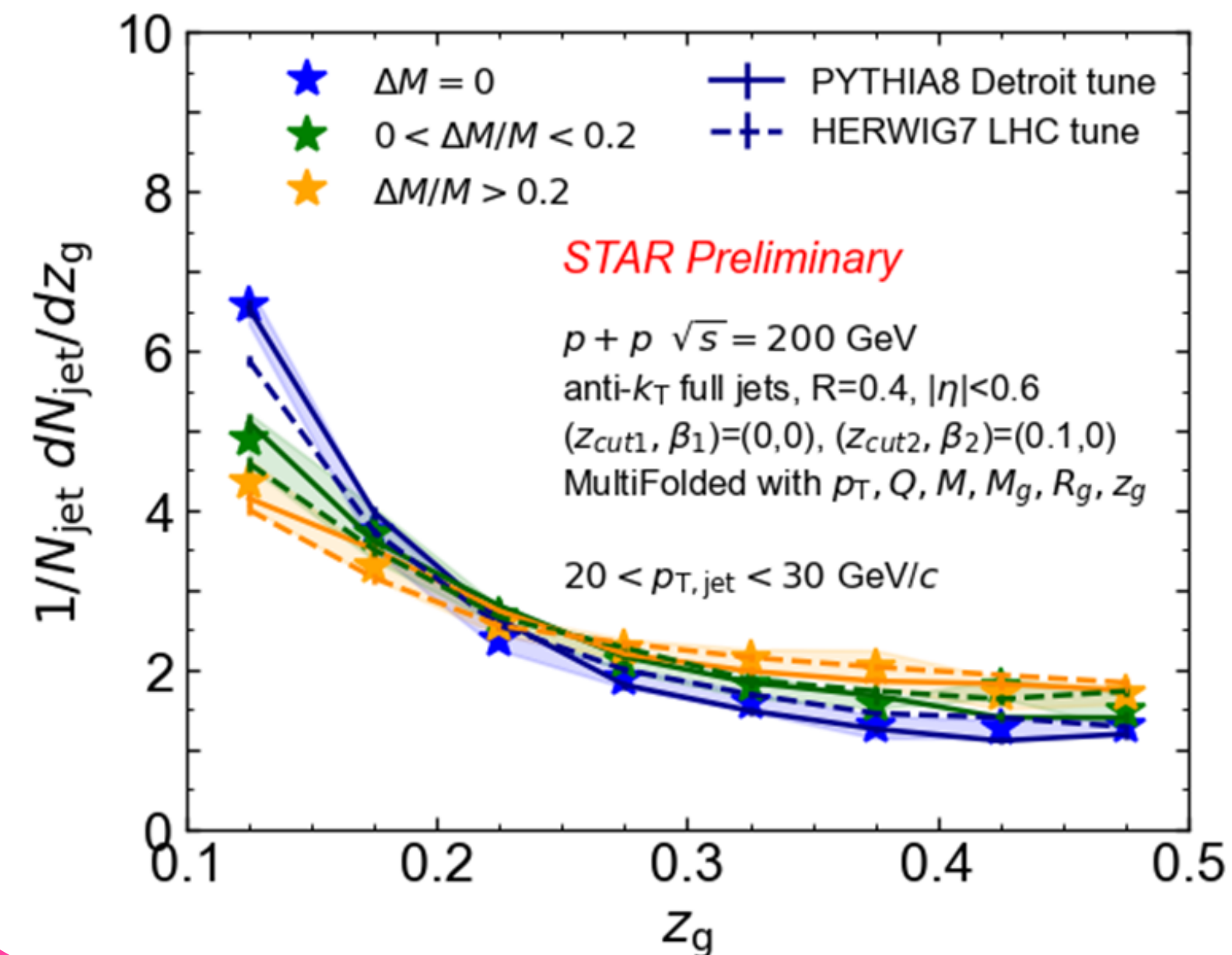
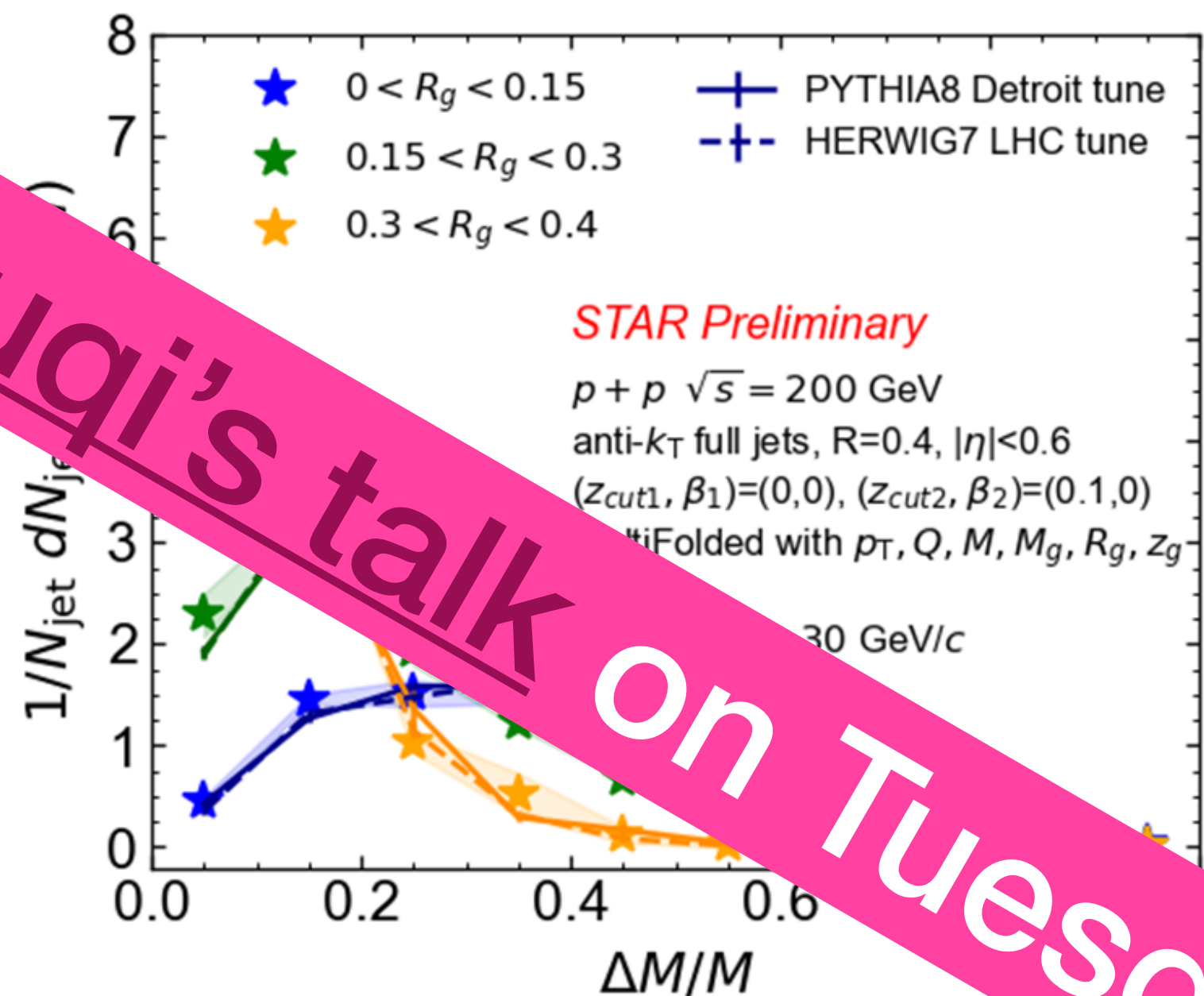


- Now able to make slices in the Lund Plane \rightarrow more stringent tests of Monte Carlo (MC) models
- Observe similarly in ALICE: high- k_T splits are wider. But tension with models for narrow splits with high k_T

N -dimensional observables With MultiFold



$$\Delta M = M - M_g$$

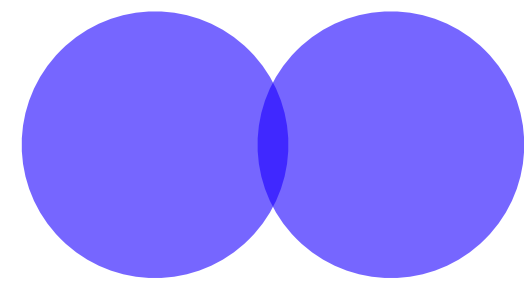


- First application of MultiFold at RHIC
- Consistent with angular ordering + kinematic constraint in early and late time splittings

See Youqi's talk on Tuesday at 11 AM!

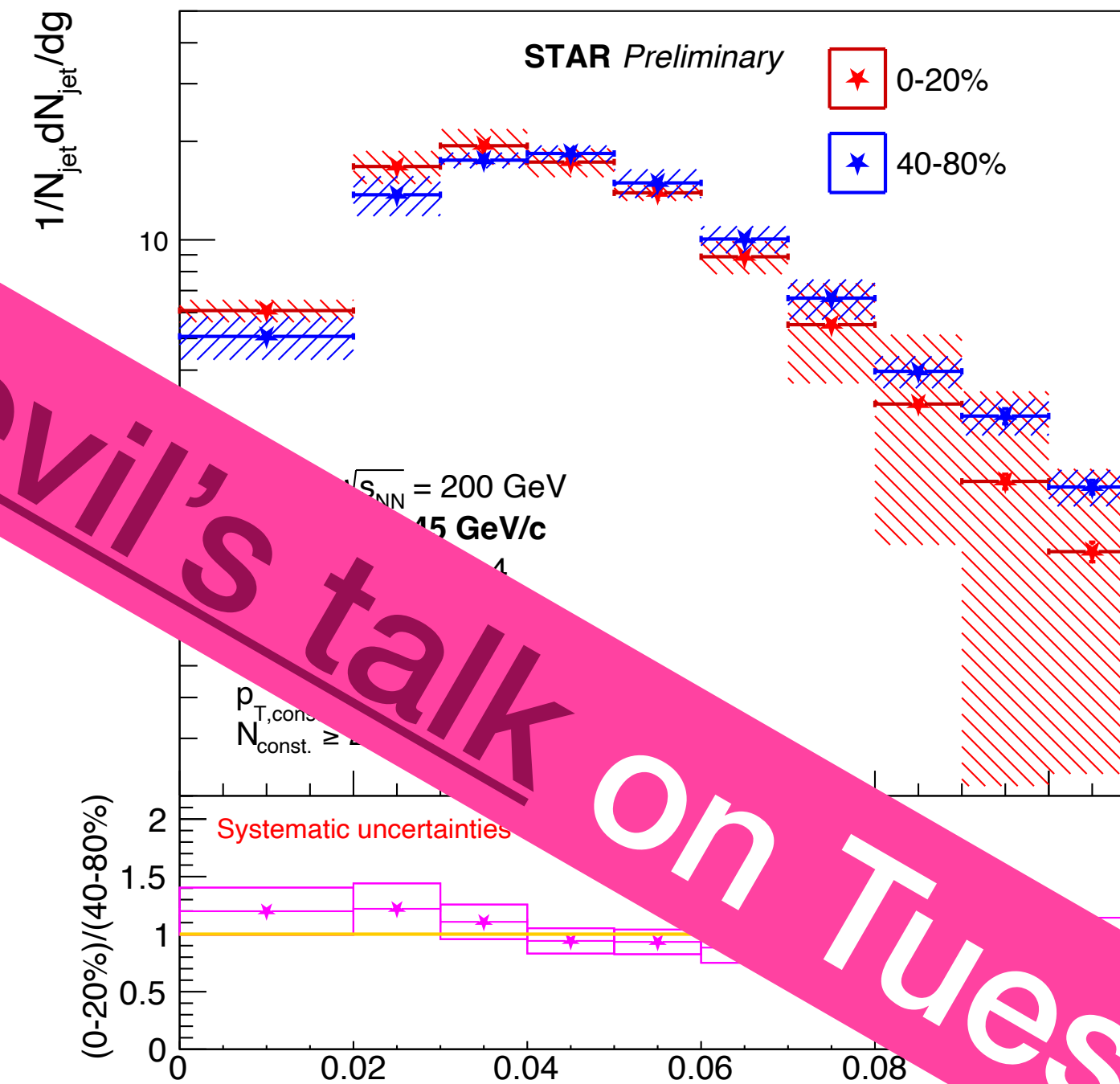
Generalized angularities in AA

With MultiFold

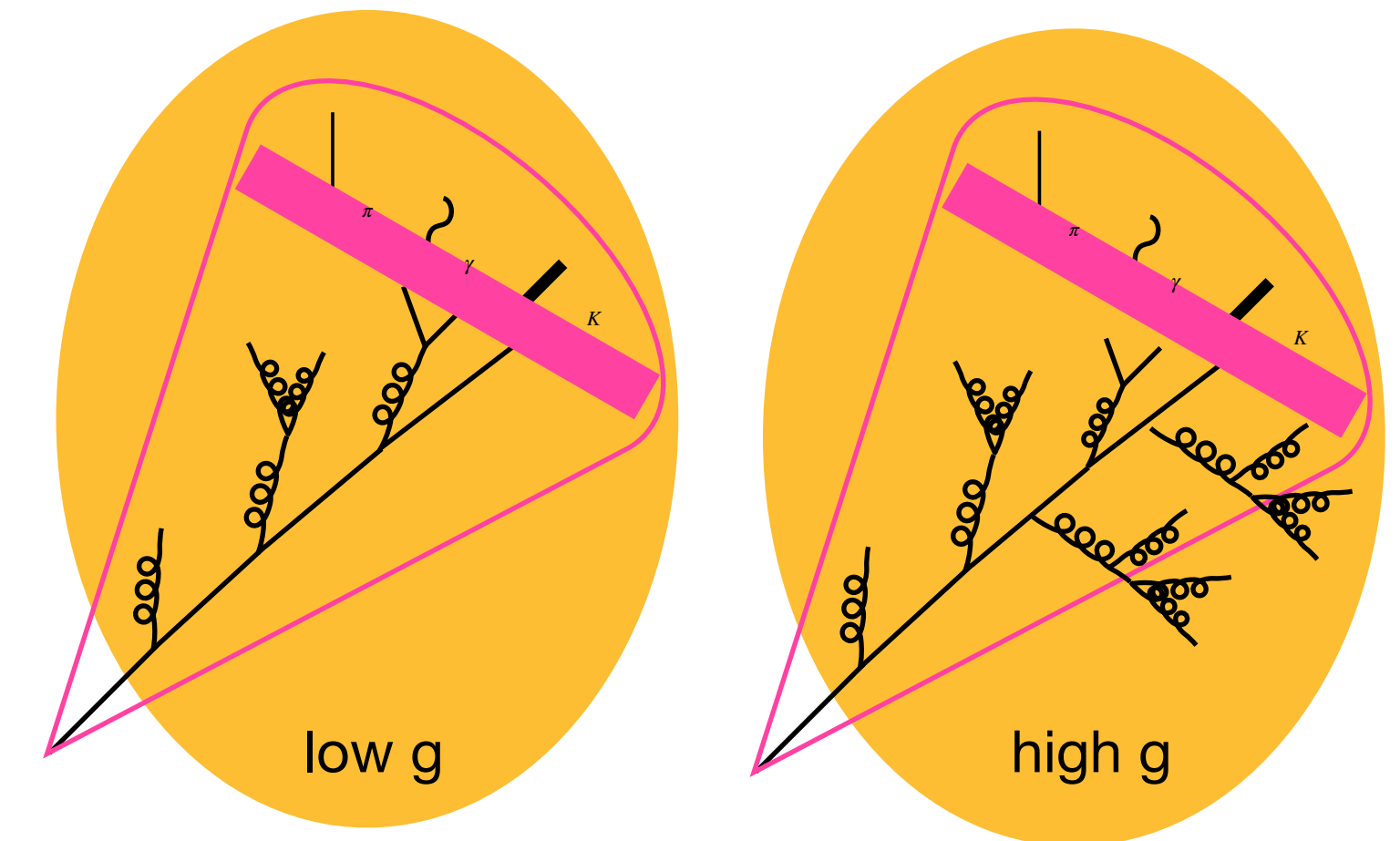


$$\lambda_{\beta}^{\kappa} = \sum_{\text{cons} \in \text{jet}} \left(\frac{p_{T,\text{cons}}}{p_{T,\text{jet}}} \right)^{\kappa} \Delta R (\text{cons, jet})^{\beta}$$

$$g = \lambda_1^1 = \frac{\sum_{\text{trk} \in \text{jet}} p_{T,\text{trk}} \Delta R}{p_{T,\text{jet}}}$$



Data corrected using MultiFold in 7D



- Generalized angularities allow tunable contributions to momentum, angular scales in IRC safe way
- With conservative systematic uncertainties *in biased pop.*, g in peripheral and central collisions are consistent

See Sevil's talk on Tuesday at 11:30 AM!

Precision QCD; exploring the Lund plane
with *multi-dimensional jet substructure*

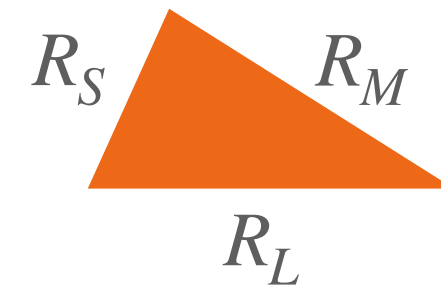
Separating p-QCD and np-QCD
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Path-length dependence of jet energy loss in medium
with *jet anisotropies (with respect to event plane)*

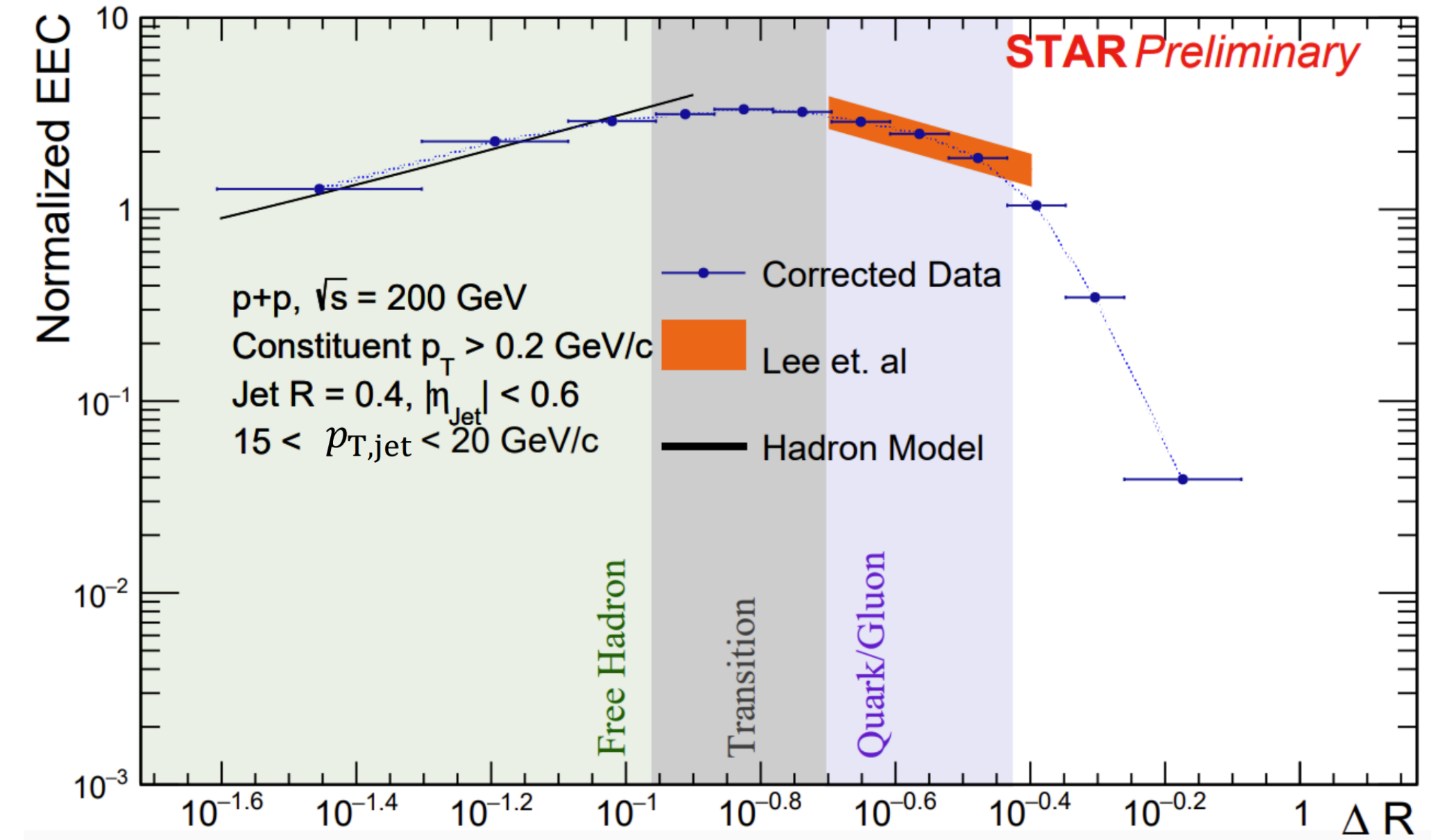
Energy-density dependence of jet energy loss in medium;
angular distribution of radiation in quenched jets
with *inclusive/semi-inclusive jet & high- p_T hadron yields*

Energy correlators



$$\text{ENC}(R_L) = \left(\prod_{k=1}^N \int d\Omega_{\vec{n}_k} \right) \delta(R_L - \Delta \hat{R}_L) \frac{1}{(E_{\text{jet}})^N} \left\langle \mathcal{E}(\vec{n}_1) \mathcal{E}(\vec{n}_2) \dots \mathcal{E}(\vec{n}_N) \right\rangle^{1,2,3}$$

$$\text{EEC}(\Delta R) = \frac{1}{\mathcal{O}} \frac{d\mathcal{O}}{d(\Delta R)}, \quad \mathcal{O} = \sum_{\text{jets}} \sum_{i \neq j} \frac{E_i E_j}{p_{T,\text{jet}}^2}$$

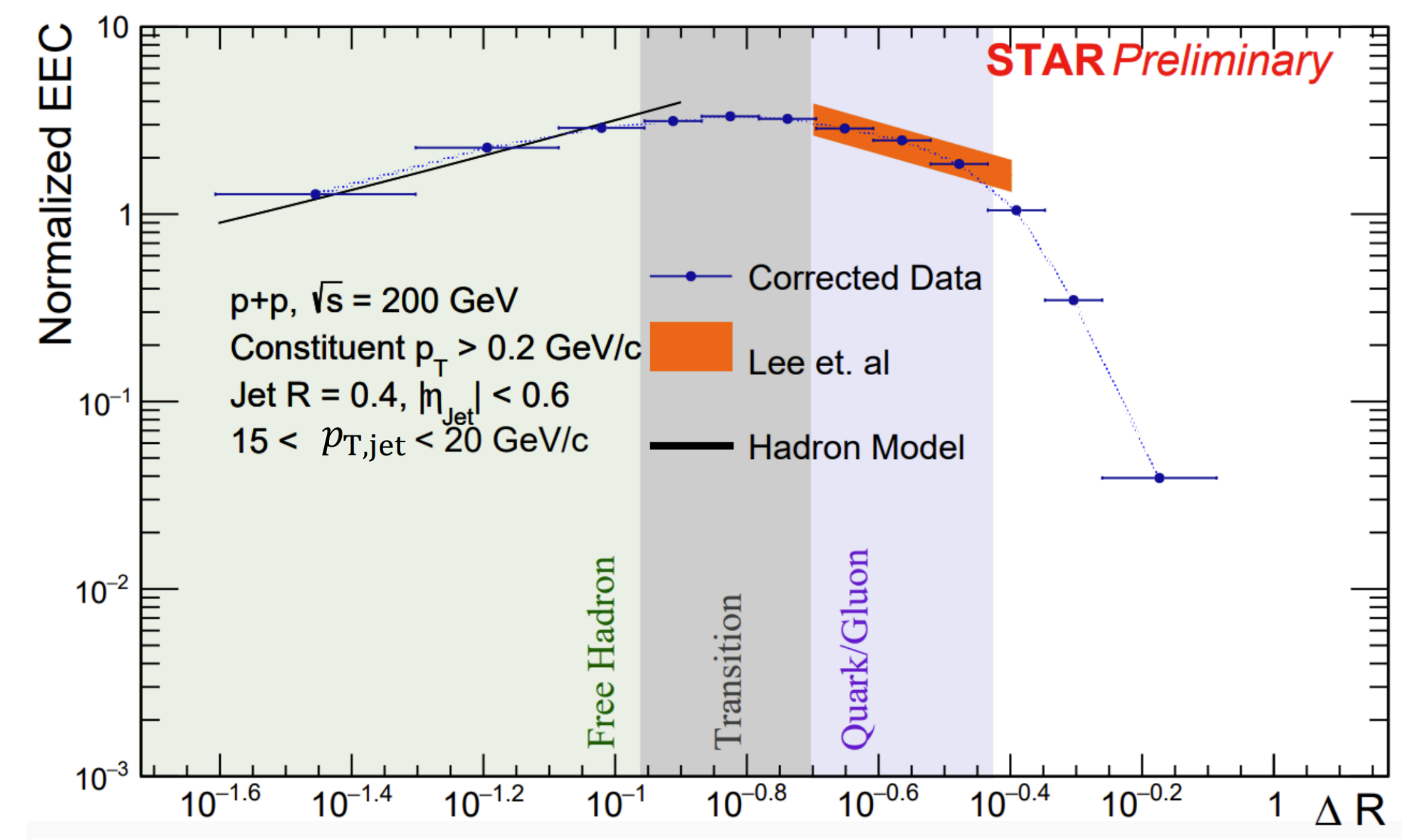
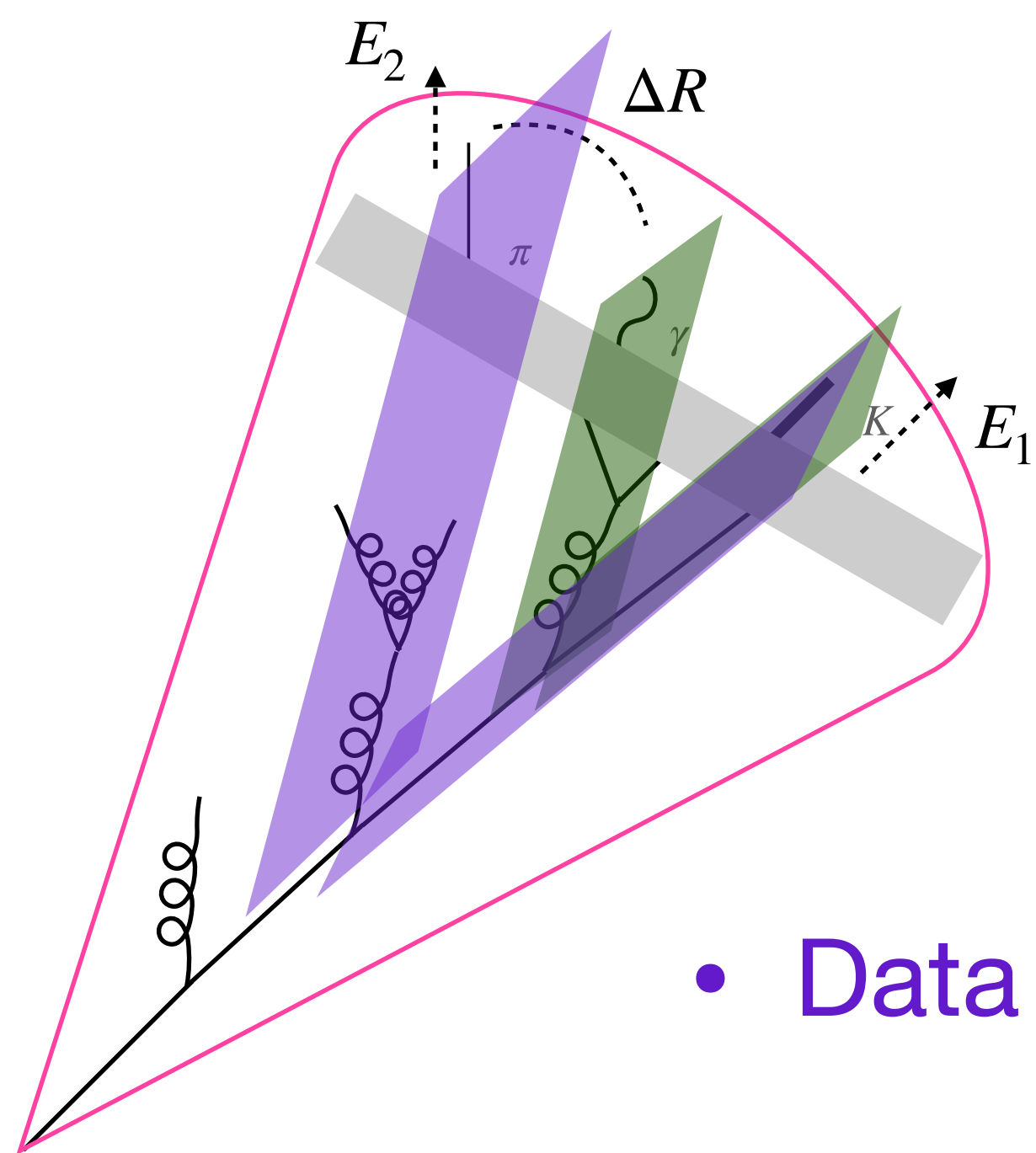


- Change in scaling when virtuality $\sim p_T R_L \sim \Lambda_{\text{QCD}}$ so $R_L^{\text{transition}} \propto 1/p_T$
- No need to recluster or remove npQCD contributions
- Simple scaling in the **hadronic** and **partonic** regimes

Energy correlators

$$\text{ENC}(R_L) = \left(\prod_{k=1}^N \int d\Omega_{\vec{n}_k} \right) \delta(R_L - \Delta \hat{R}_L) \frac{1}{(E_{\text{jet}})^N} \langle \mathcal{E}(\vec{n}_1) \mathcal{E}(\vec{n}_2) \dots \mathcal{E}(\vec{n}_N) \rangle^{1,2,3}$$

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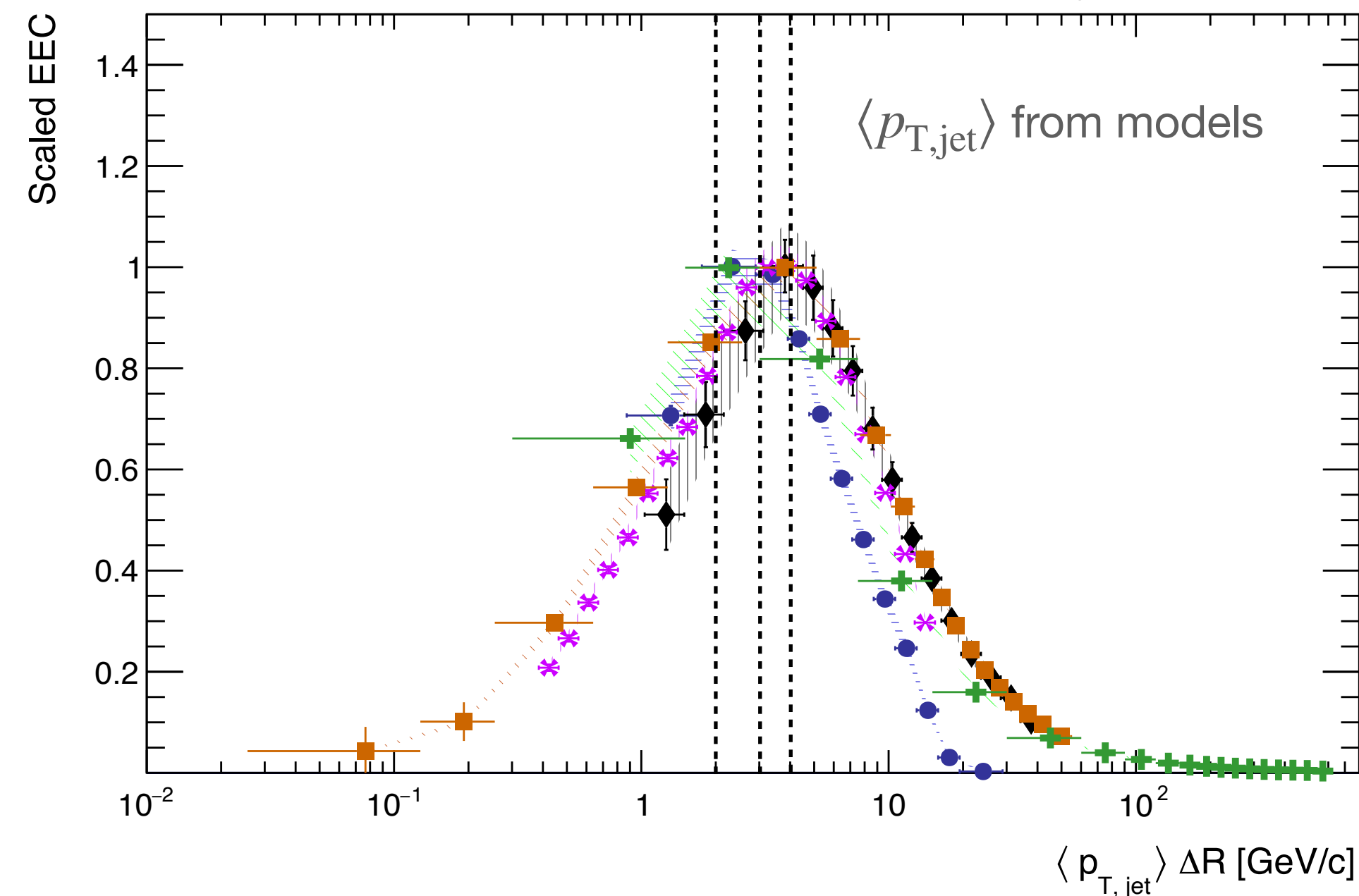


- Data agree well with NLL pQCD calculation (& MC model, not shown)
- Data agree well with model assuming non-interacting hadrons

Energy correlators

$$\text{ENC}(R_L) = \left(\prod_{k=1}^N \int d\Omega_{\vec{n}_k} \right) \delta(R_L - \Delta\hat{R}_L) \frac{1}{(E_{\text{jet}})^N} \langle \mathcal{E}(\vec{n}_1) \mathcal{E}(\vec{n}_2) \dots \mathcal{E}(\vec{n}_N) \rangle^{1,2,3}$$

- STAR Preliminary: $\sqrt{s} = 200$ GeV $30 < \text{Full Jet } p_T < 50$ GeV/c
- *— ALICE Preliminary: $\sqrt{s} = 5.02$ TeV, $20 < \text{Charged Jet } p_T < 40$ GeV/c
- ◆— ALICE Preliminary: $\sqrt{s} = 13$ TeV, $60 < \text{Charged Jet } p_T < 80$ GeV/c
- CMS Preliminary: $\sqrt{s} = 13$ TeV $97 < \text{Full Jet } p_T < 220$ GeV/c
- +— CMS Preliminary: $\sqrt{s} = 13$ TeV, $1410 < \text{Full Jet } p_T < 1784$ GeV/c



CMS: Lu, Boost '23
ALICE: Fan, Quark Matter '23

$$\text{EEC}(\Delta R) = \frac{1}{\mathcal{O}} \frac{d\mathcal{O}}{d(\Delta R)}, \quad \mathcal{O} = \sum_{\text{jets}} \sum_{i \neq j} \frac{E_i E_j}{p_{T,\text{jet}}^2}$$

- Testing universality of transition region by comparing to LHC data:
- ~ 2 orders of magnitude in \sqrt{s} and $p_{T,\text{jet}}$ from STAR \rightarrow ALICE \rightarrow CMS, transition $\sim 2 - 4$ GeV/c
- STAR more similar to CMS high- p_T (high- x) jets than ALICE or CMS low- p_T jets — **q vs. g** differences

Precision QCD; exploring the Lund plane
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Separating p-QCD and np-QCD
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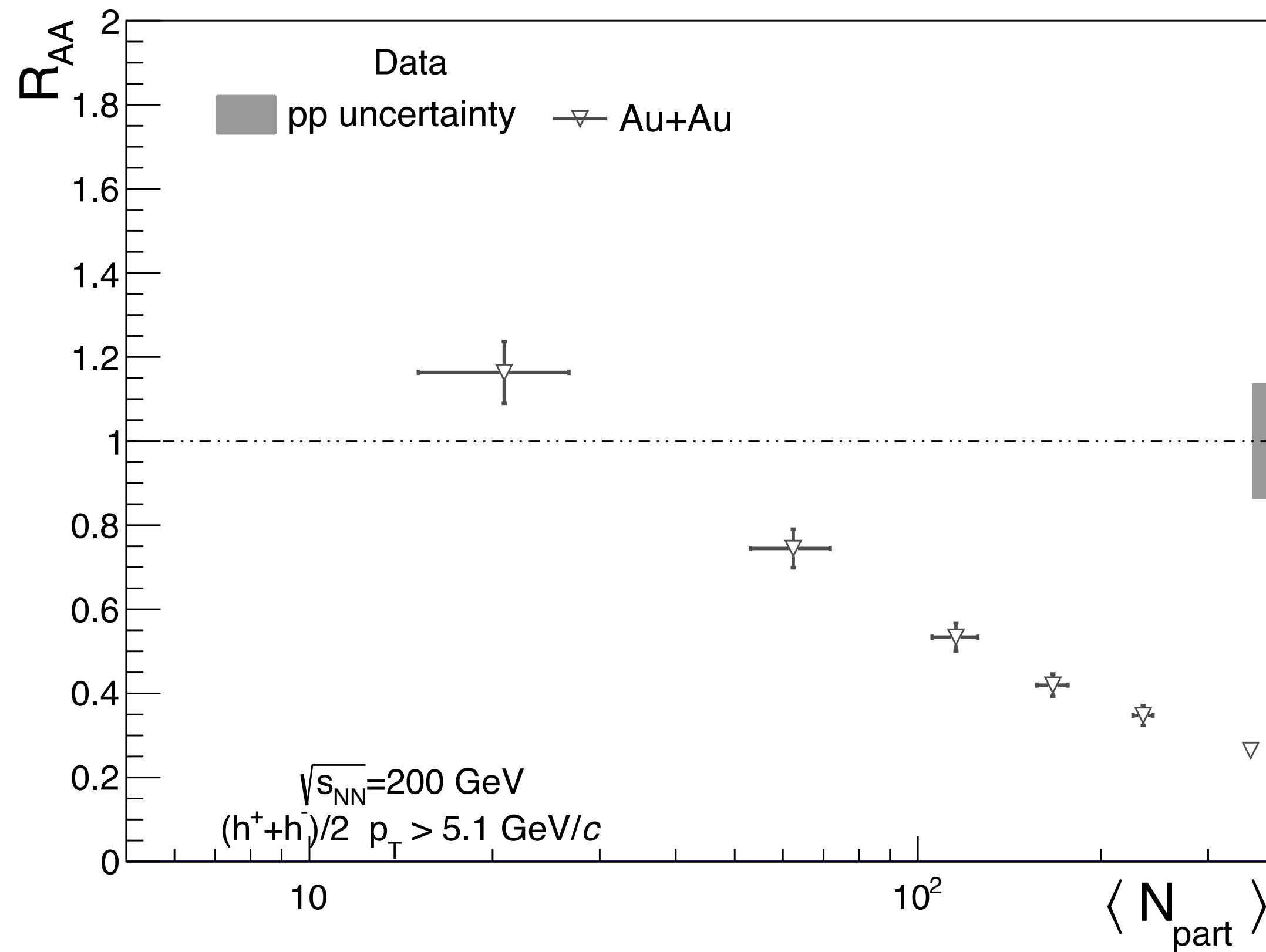
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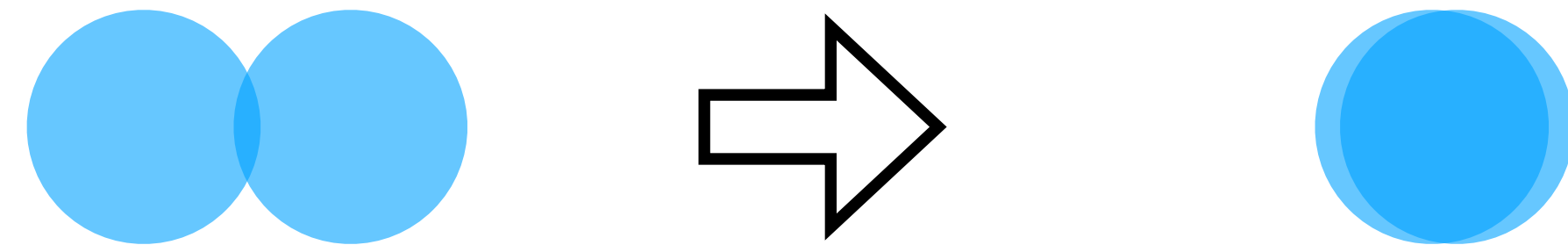
Energy-density dependence of jet energy loss in medium;
angular distribution of radiation in quenched jets
with *inclusive/semi-inclusive jet & high- p_T hadron yields*

Inclusive yield modification

$$R_{AA} = \frac{1}{N_{ev}^{AA}} \frac{d^2 N^{AA} / d\eta dp_T}{\langle T_{AA} \rangle d^2 \sigma^{NN} / d\eta dp_T}, \quad \langle T_{AA} \rangle = \langle N_{coll} \rangle / \sigma_{inel}^{NN}$$



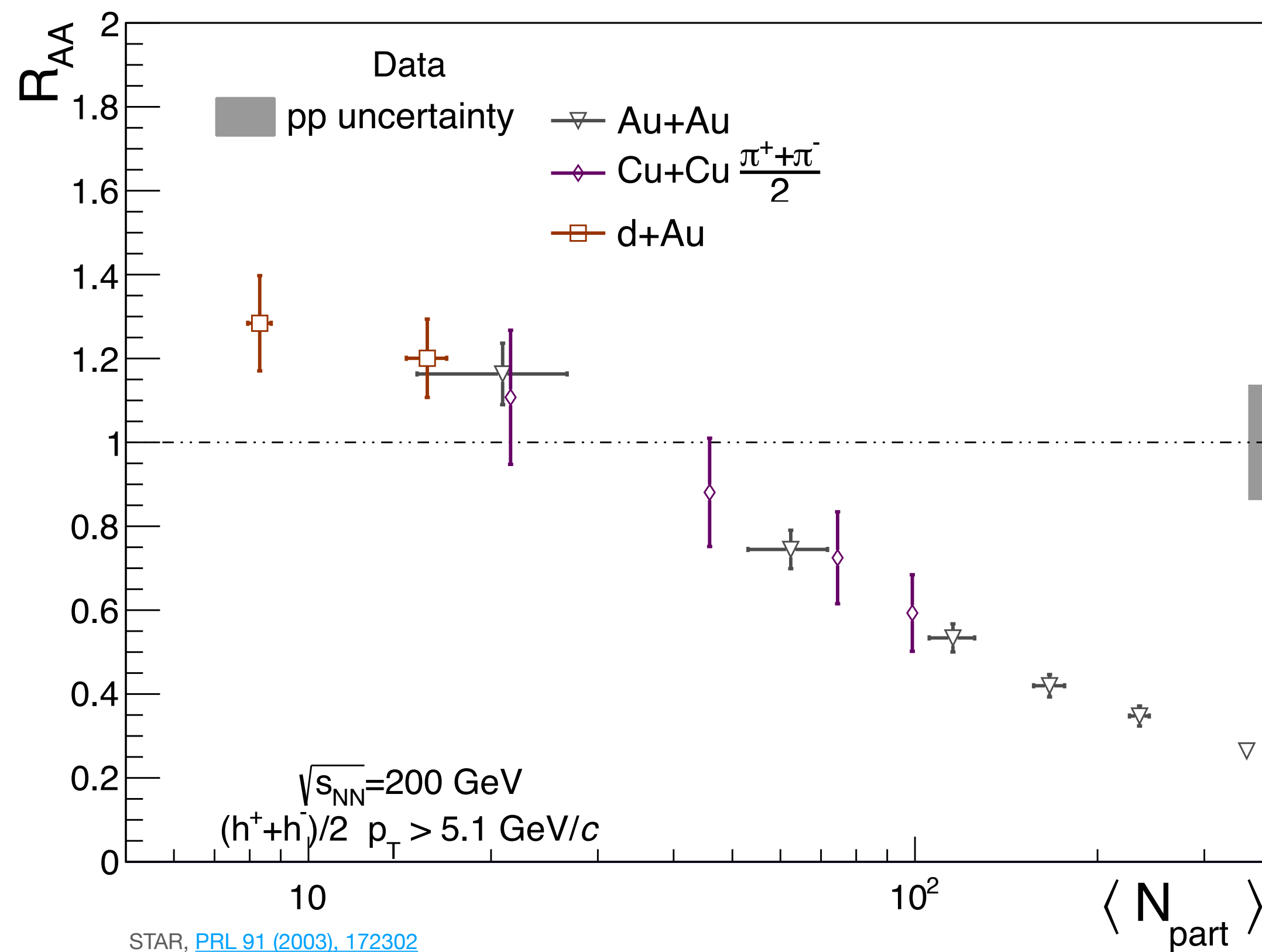
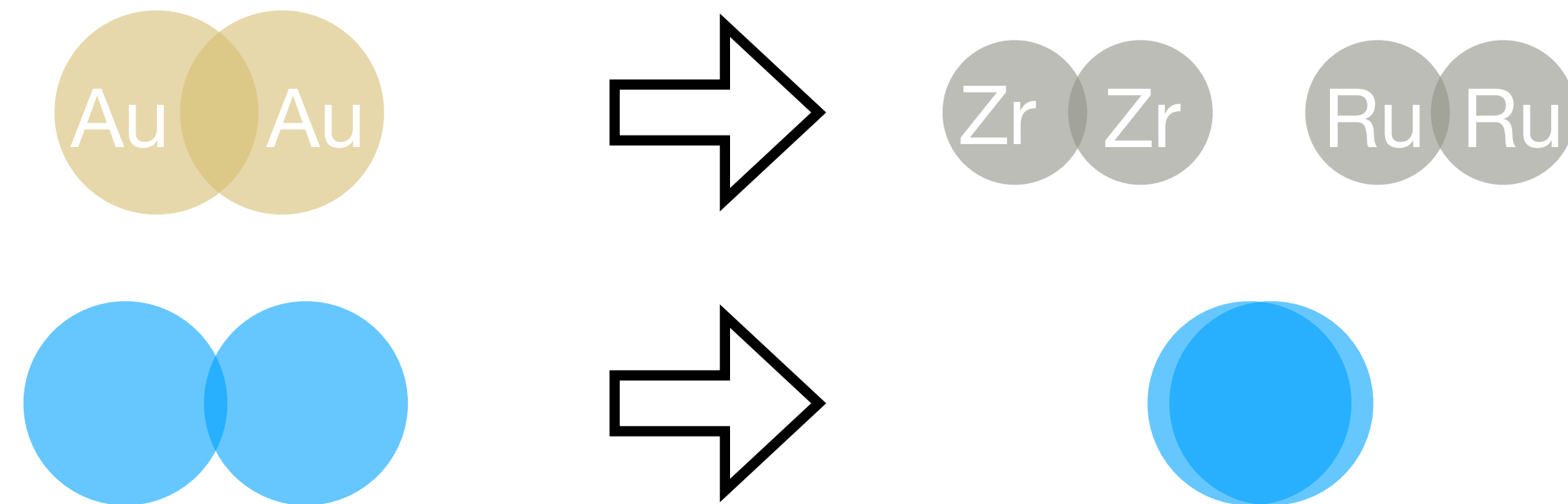
STAR, [PRL 91 \(2003\), 172302](https://arxiv.org/abs/1703.07501)



- Suppression strongly increases with $\langle N_{part} \rangle$

Inclusive yield modification

$$R_{AA} = \frac{1}{N_{ev}^{AA}} \frac{d^2 N^{AA} / d\eta dp_T}{\langle T_{AA} \rangle d^2 \sigma^{NN} / d\eta dp_T}, \quad \langle T_{AA} \rangle = \langle N_{coll} \rangle / \sigma_{inel}^{NN}$$

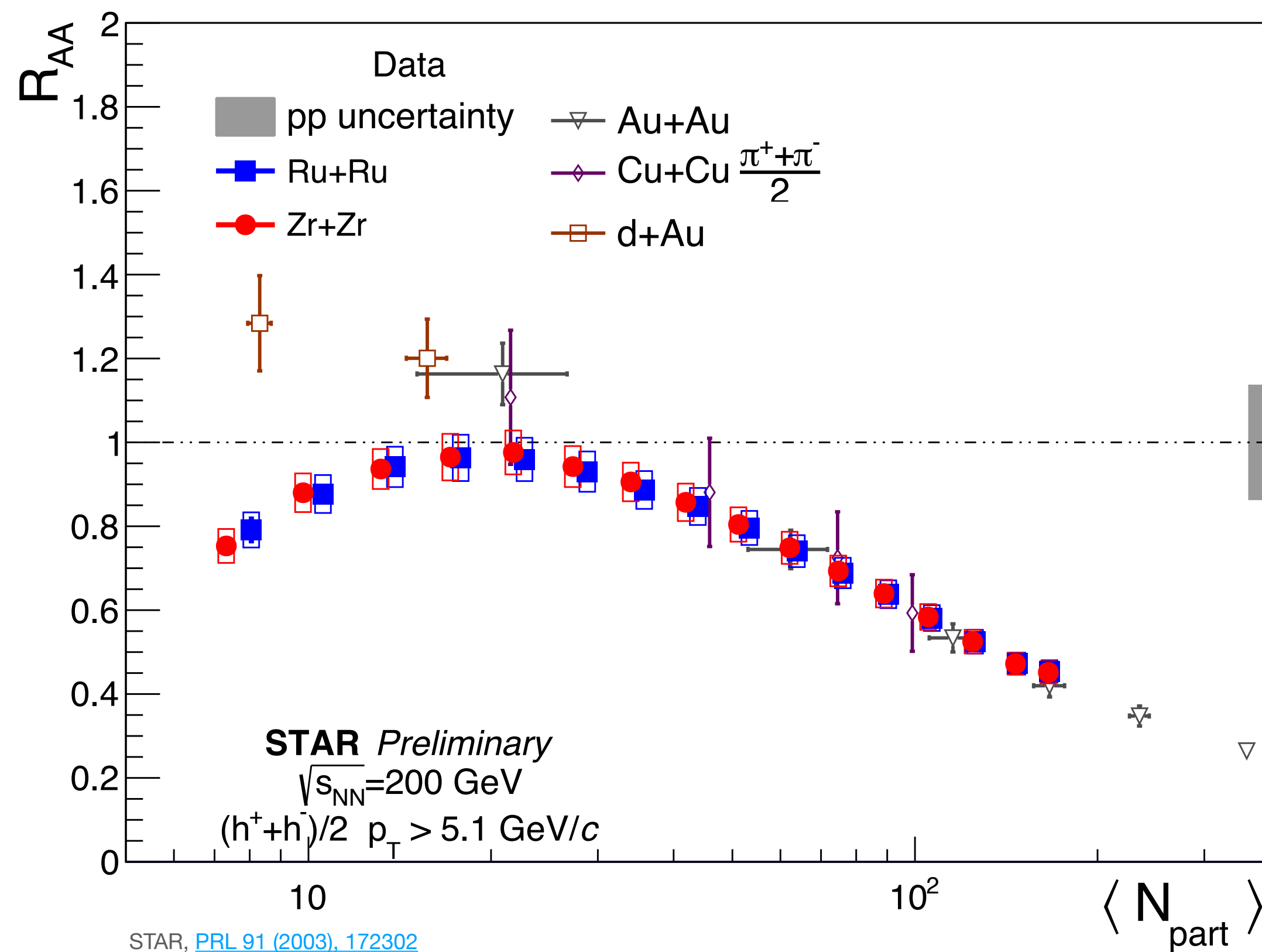


STAR, [PRL 91 \(2003\), 172302](#)
 STAR, [PRL 91 \(2003\), 072304](#)
 STAR, [PRC 81 \(2010\), 054907](#)

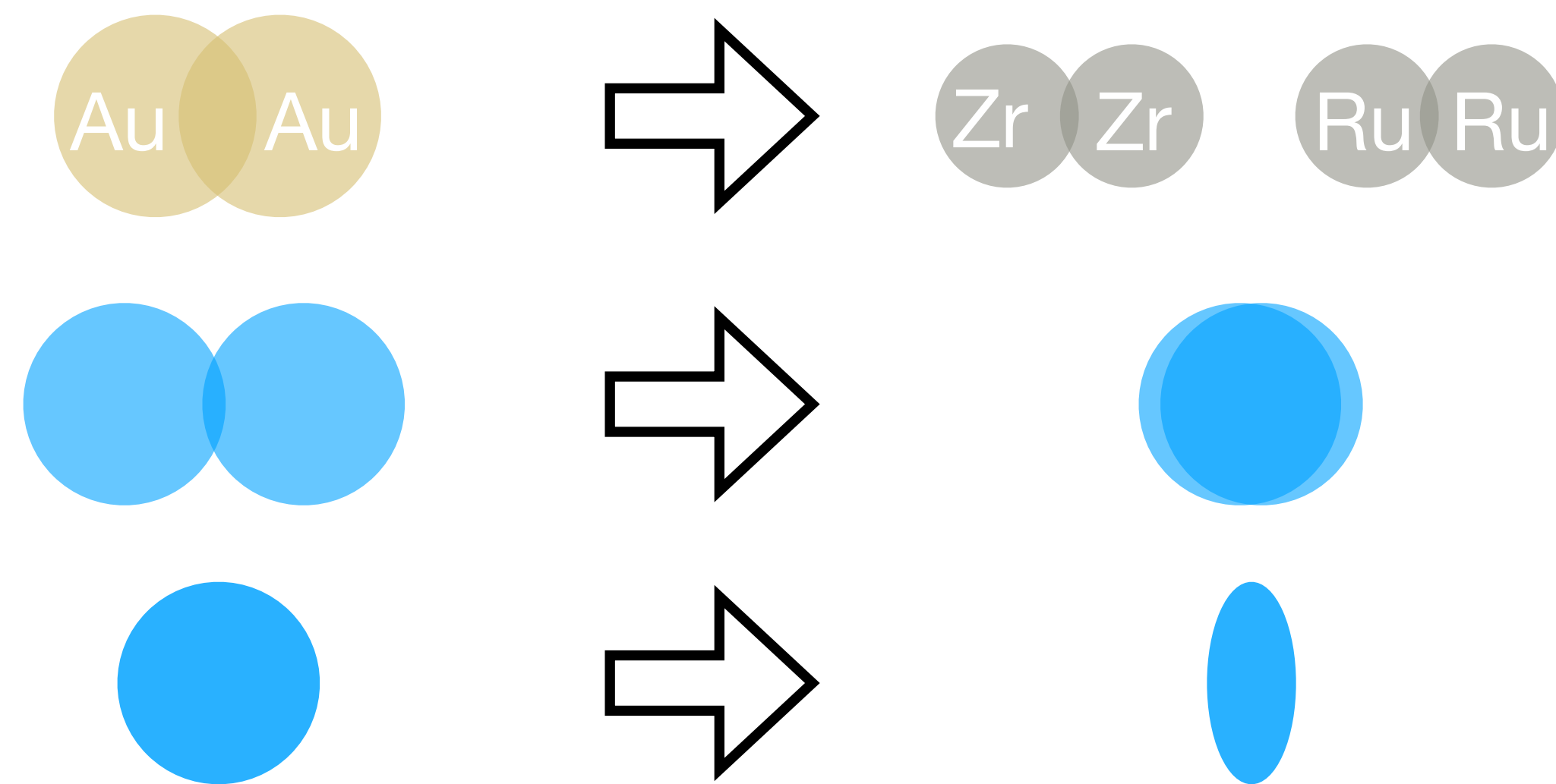
- R_{AA} falls with $\langle N_{part} \rangle$ independent of collision species (system size)

Inclusive yield modification

$$R_{AA} = \frac{1}{N_{ev}^{AA} \langle T_{AA} \rangle} \frac{d^2 N^{AA} / d\eta dp_T}{d^2 \sigma^{NN} / d\eta dp_T}, \quad \langle T_{AA} \rangle = \langle N_{coll} \rangle / \sigma_{inel}^{NN}$$



STAR, [PRL 91 \(2003\), 172302](#)
 STAR, [PRL 91 \(2003\), 072304](#)
 STAR, [PRC 81 \(2010\), 054907](#)

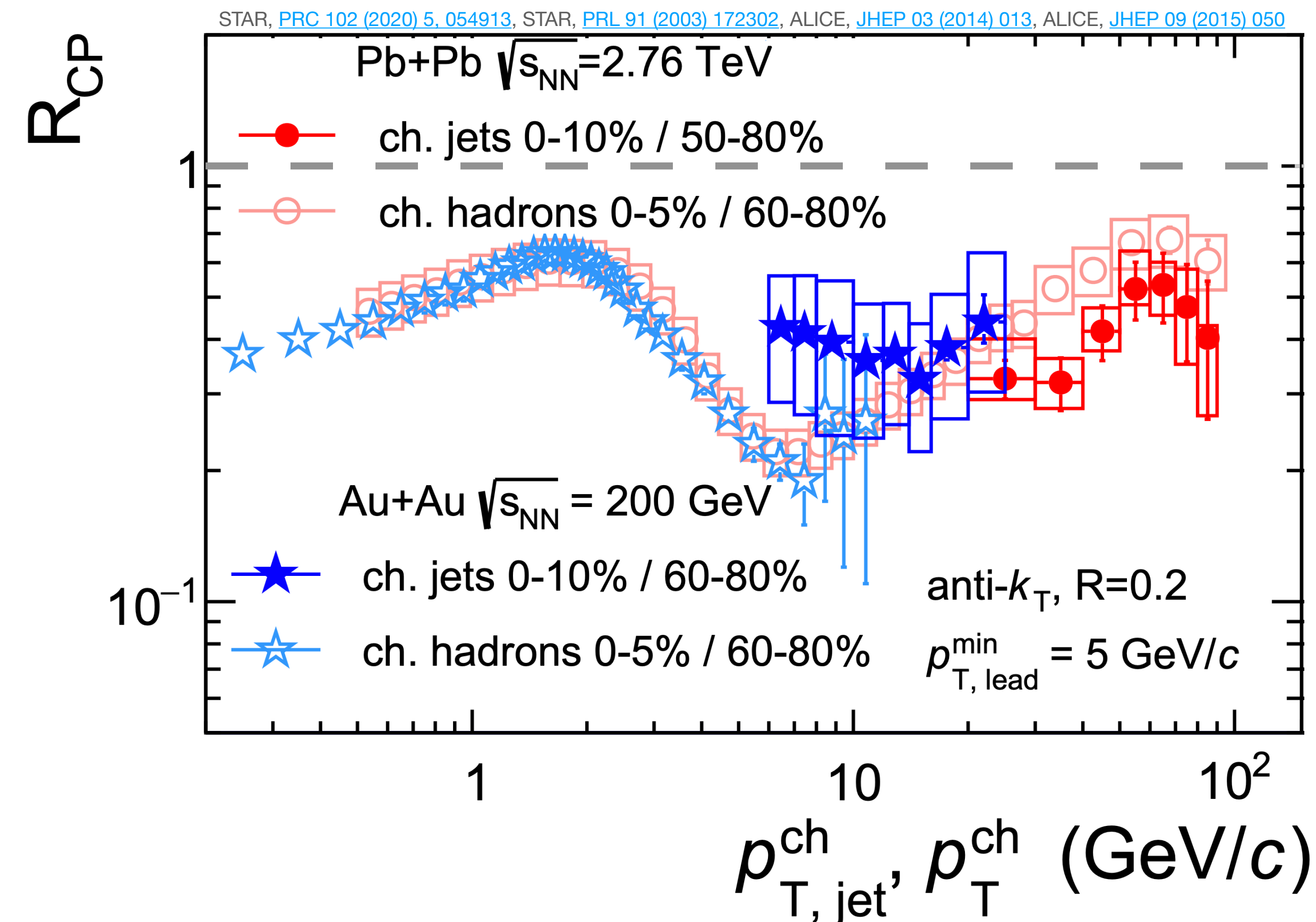


- R_{AA} falls with $\langle N_{part} \rangle$ independent of collision species (system size), above ~ 20
- Later: for given $\langle N_{part} \rangle$, how does geometry influence E -loss?

Inclusive yield modification

$$R_{CP} = \frac{N_{ev}^P \langle T_{AA,P} \rangle d^2N^C/d\eta dp_T}{N_{ev}^C \langle T_{AA,C} \rangle d^2N^P/d\eta dp_T}, \quad \langle T_{AA} \rangle = \langle N_{coll} \rangle / \sigma_{inel}^{NN}$$

- Jet R_{AA} consistent with hadron R_{AA}
- Strong suppression across p_T
- RHIC and LHC jets already have kinematic overlap
- Similar quenching?

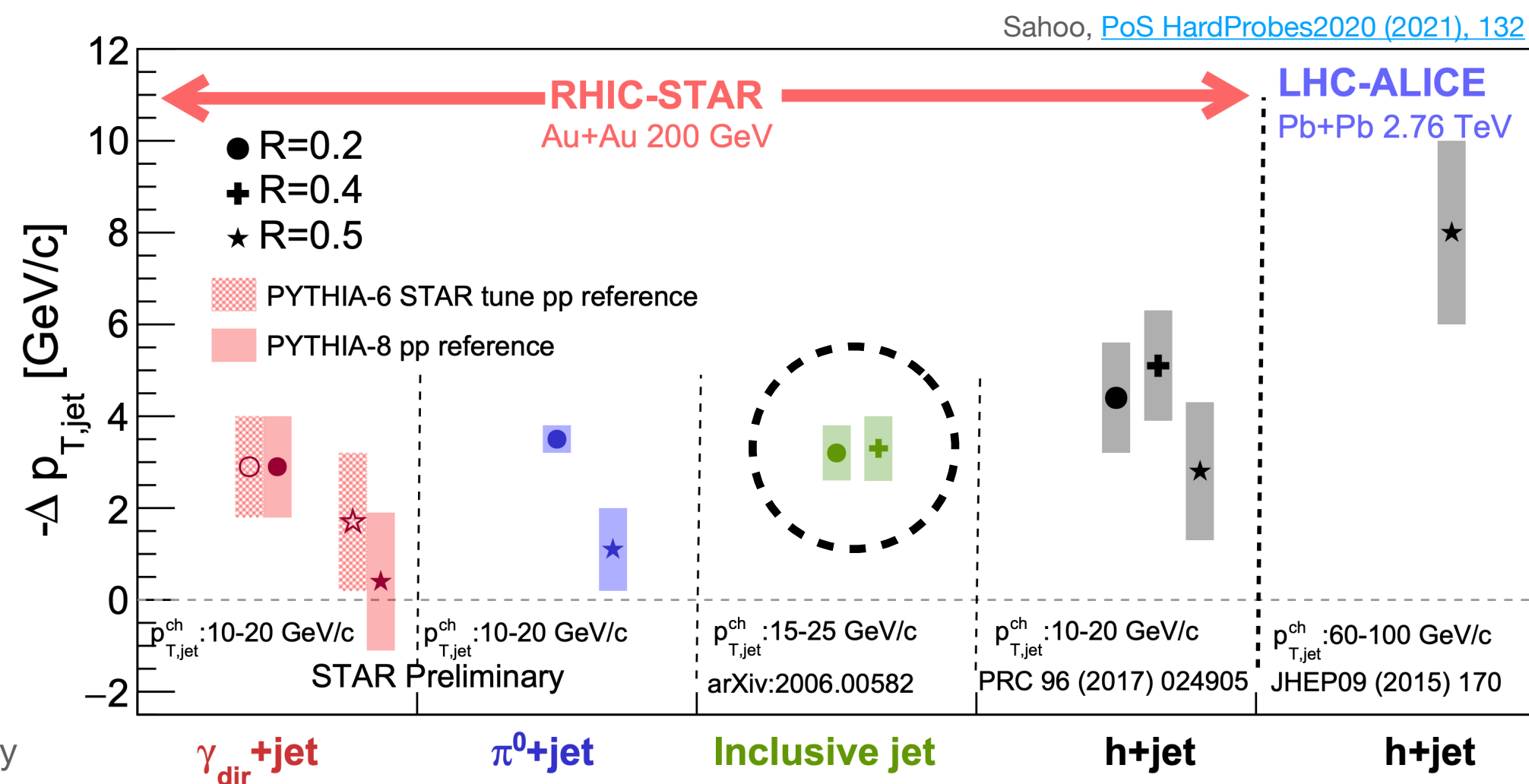
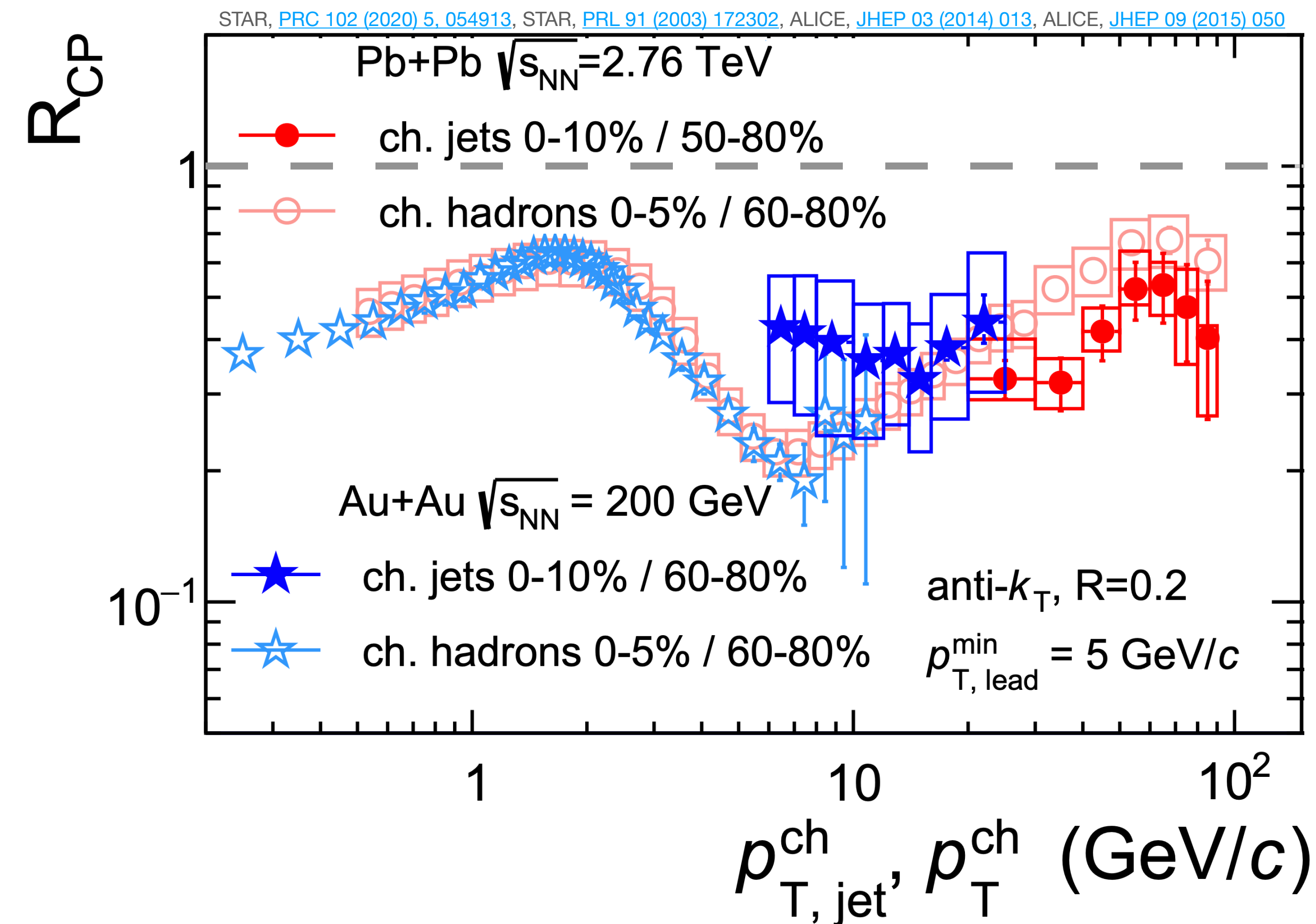


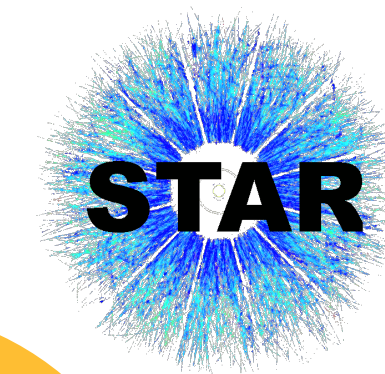
Inclusive yield modification

$$R_{CP} = \frac{N_{ev}^P \langle T_{AA,P} \rangle \frac{d^2N^C}{d\eta dp_T}}{N_{ev}^C \langle T_{AA,C} \rangle \frac{d^2N^P}{d\eta dp_T}}, \quad \langle T_{AA} \rangle = \langle N_{coll} \rangle / \sigma_{inel}^{NN}$$

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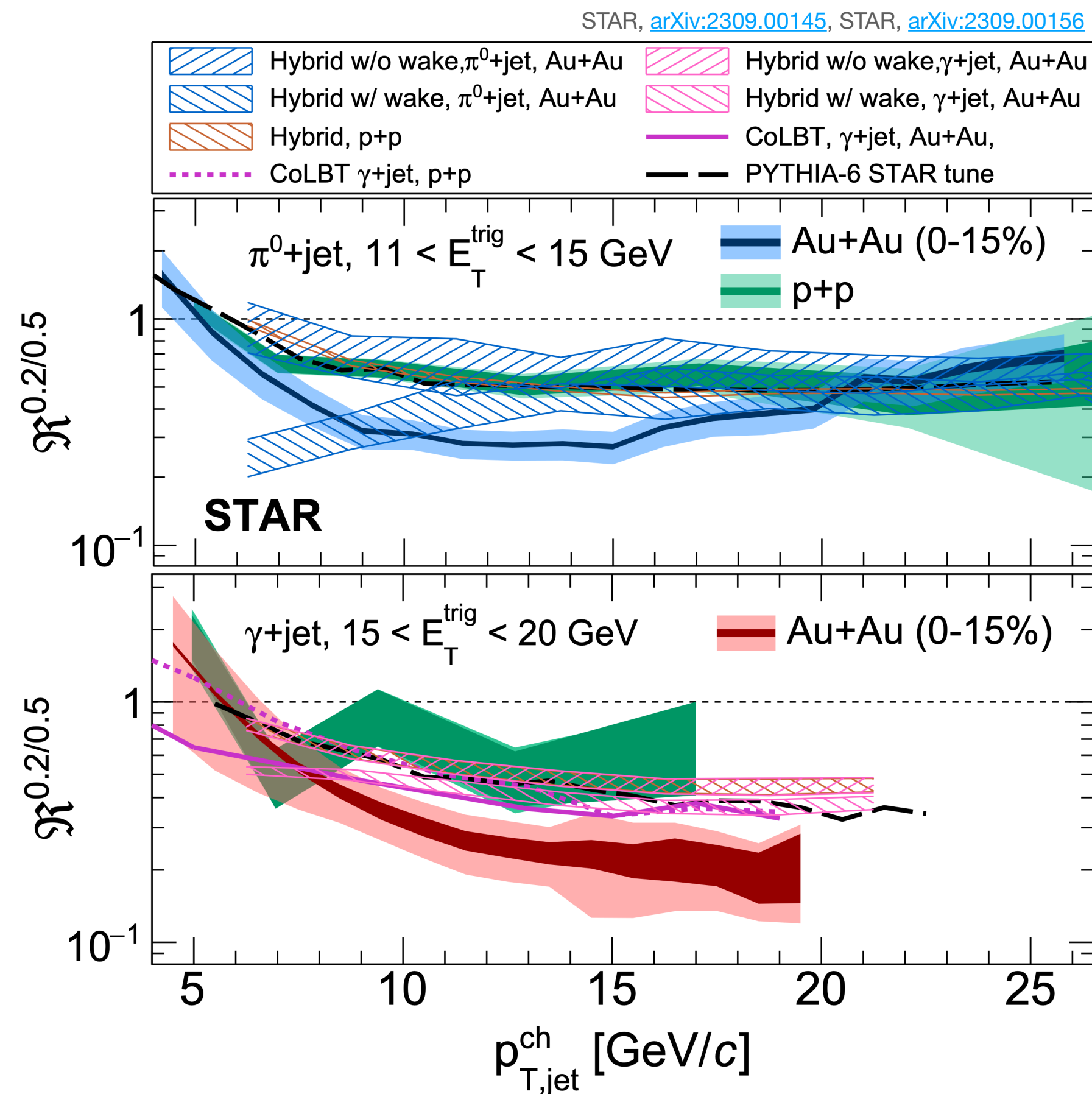
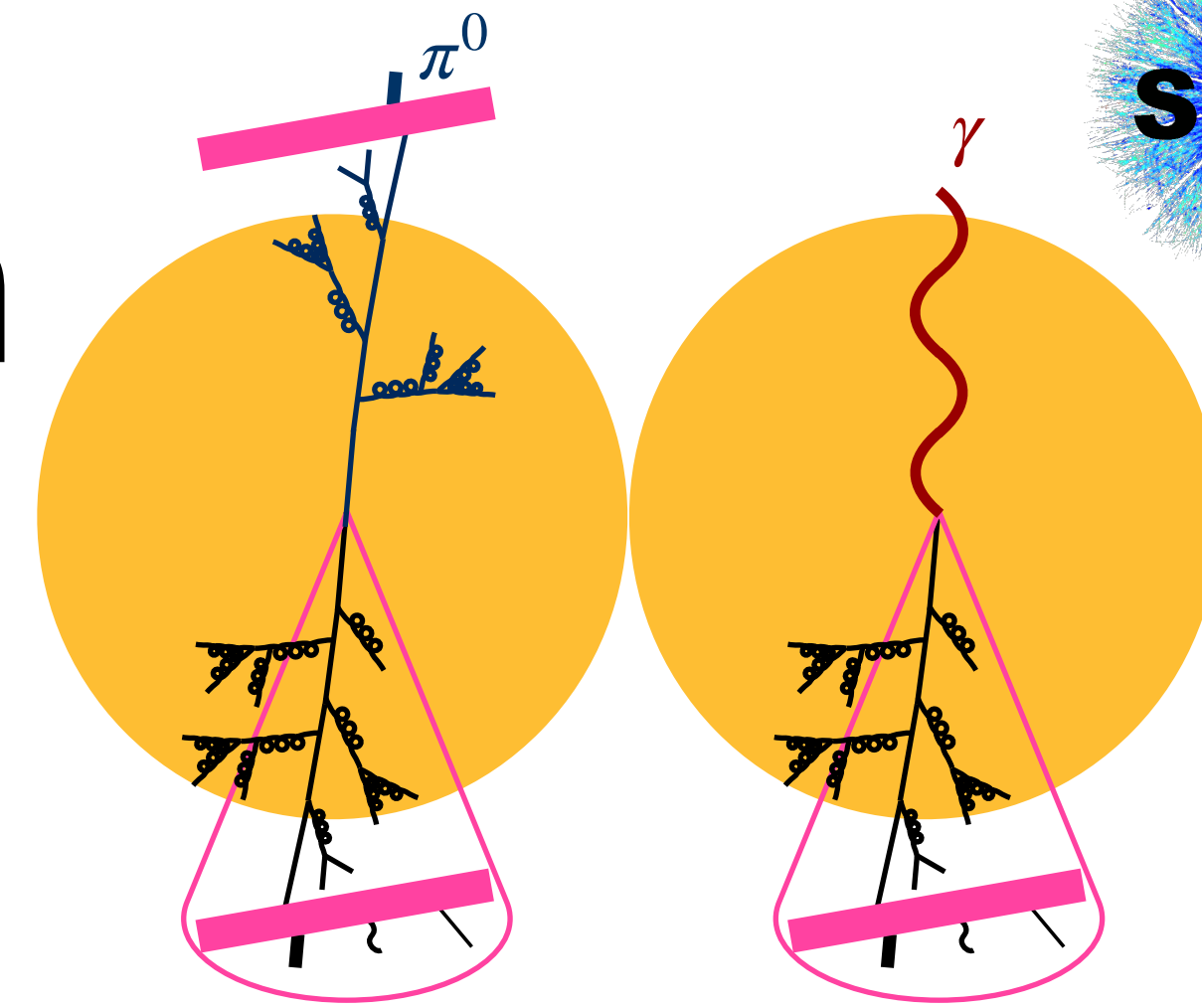
- Similar quenching?
Absolute, smaller. Relative, *larger!*



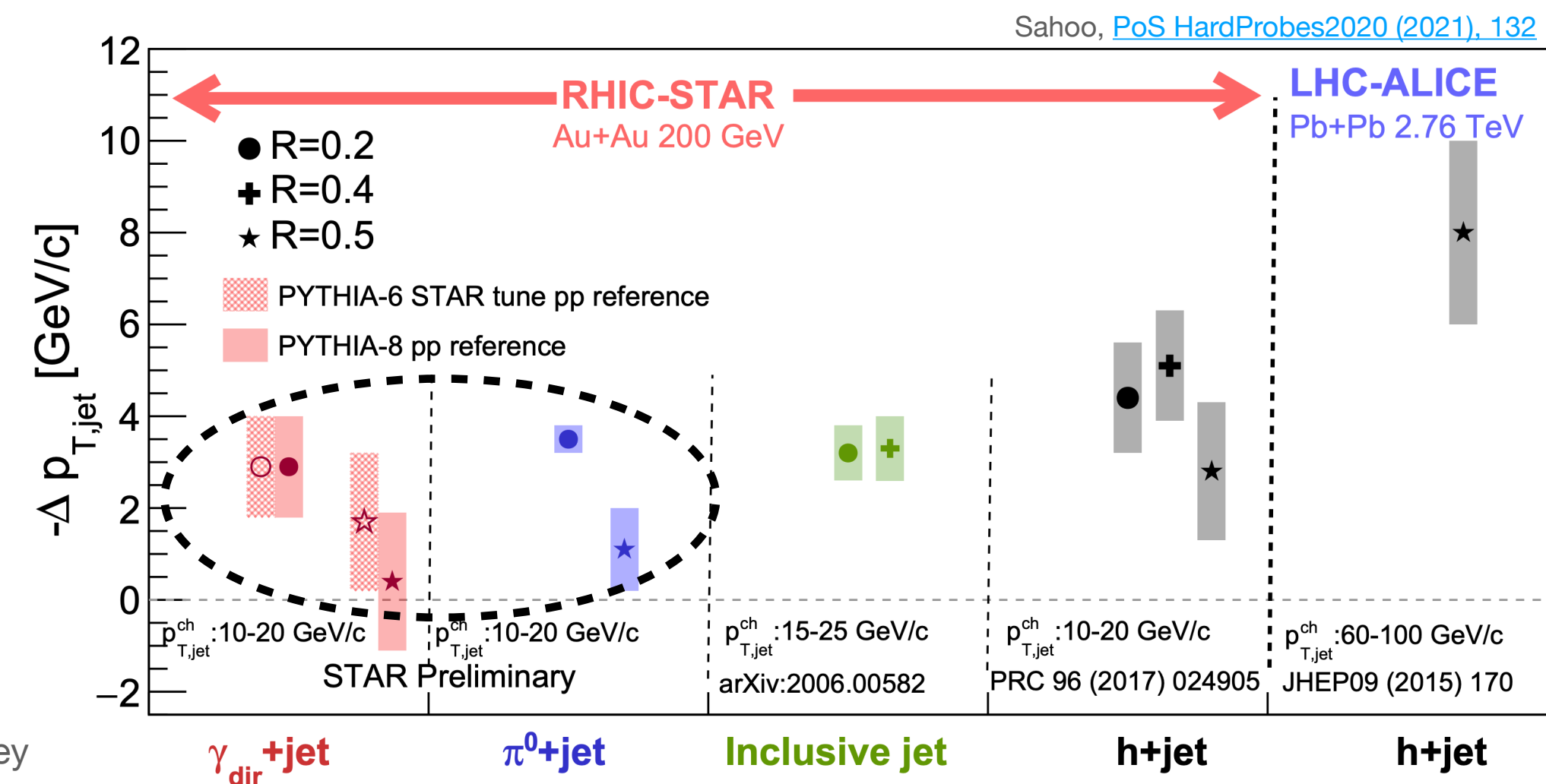


Semi-inclusive yield modification

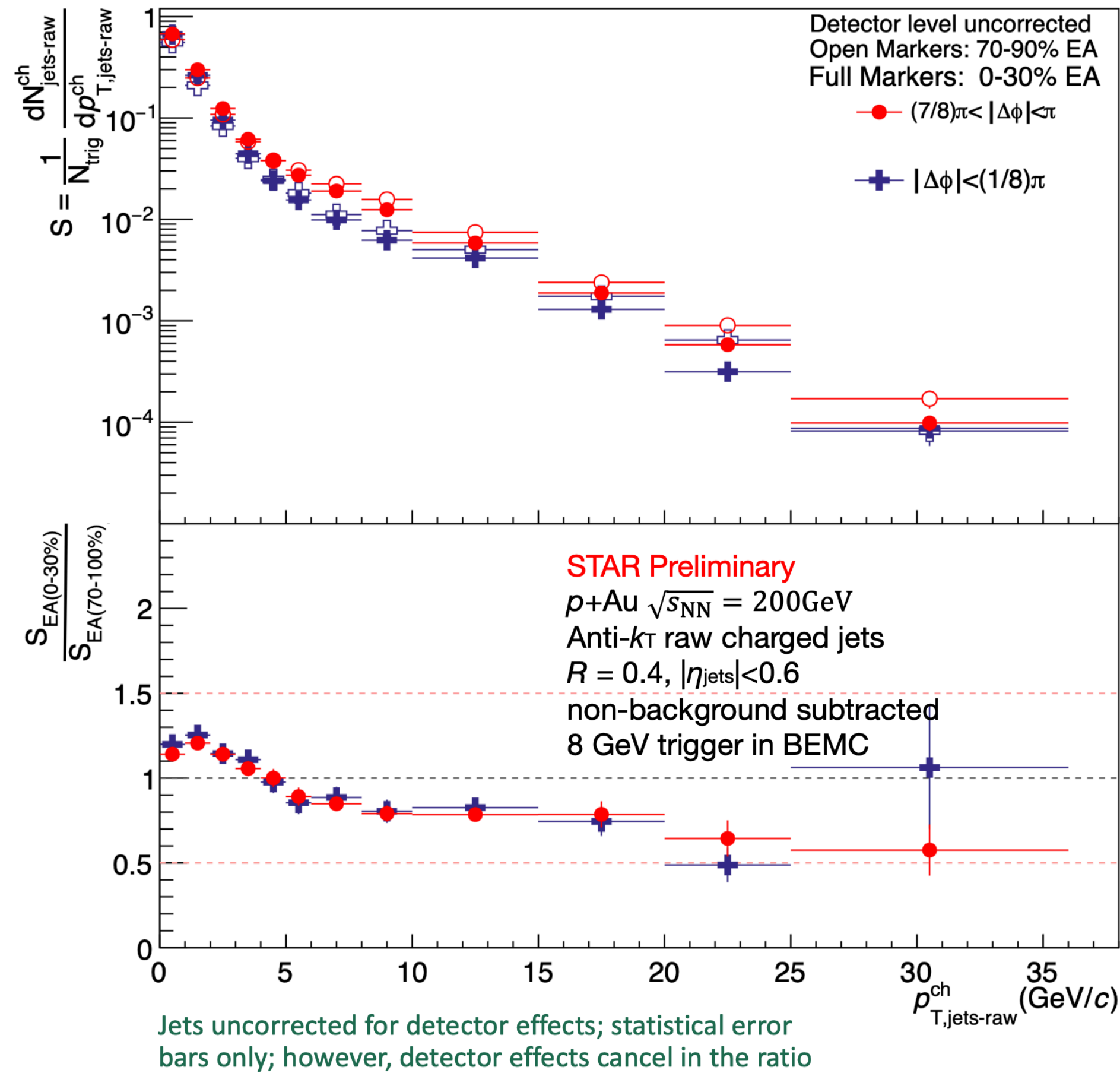
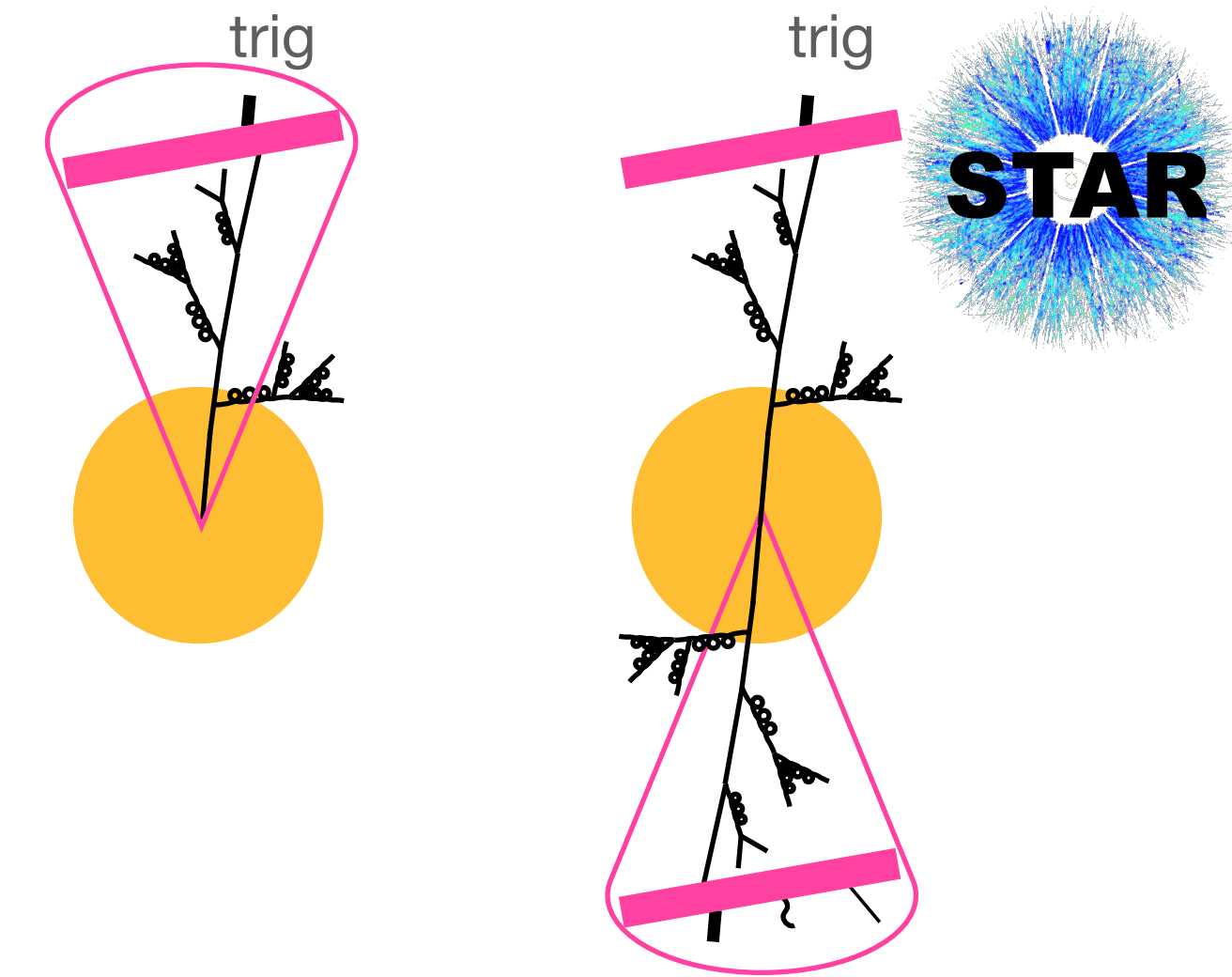
$$I_{AA} = Y^{AA}(p_{T,jet}^{ch}, R) / Y^{pp}(p_{T,jet}^{ch}, R), Y(p_{T,jet}^{ch}, R) = \frac{1}{N_{trig}} \int_{3\pi/4}^{5\pi/4} d\Delta\phi \left[\frac{d^2 N_{jet}(R)}{dp_{T,jet}^{ch} d\Delta\phi} \right]_{E_T^{trig} \in [E_T^{min}, E_T^{max}]}$$



- Recoil jet yield suppression in AuAu, stronger in small R jets
- Clear observation of *intra-jet broadening*
- Models unable to quantitatively describe the effect



Semi-inclusive yield modification in pAu collisions



¹CMS, [JHEP 09 \(2010\), 091](#)
²CMS, [PLB 718 \(2013\), 795](#)
³ALICE, [PLB 719 \(2013\), 29](#)
⁴ATLAS, [PLB 748 \(2015\), 392](#)
⁵PHENIX, [PRL 116 \(2016\), 122301](#)
⁶ALICE, [Nat. Phys. 13 \(2017\), 535](#)

- Hot nuclear matter effects in pAu collisions?^{1,2,3,4,5,6,...}
- Jet yield suppression, but on both *near* and *away* side → not surface bias as typical in AA with high p_{T} trigger...
- Jet substructure*, dijet p_{T} balance A_{J}^* also unmodified
- Anti-correlation of event activity at large rapidity with jet p_{T} at mid-rapidity* suggests $t \sim 0$ kinematics^{7,8}

*not shown

Precision QCD; exploring the Lund plane
with *multi-dimensional jet substructure*

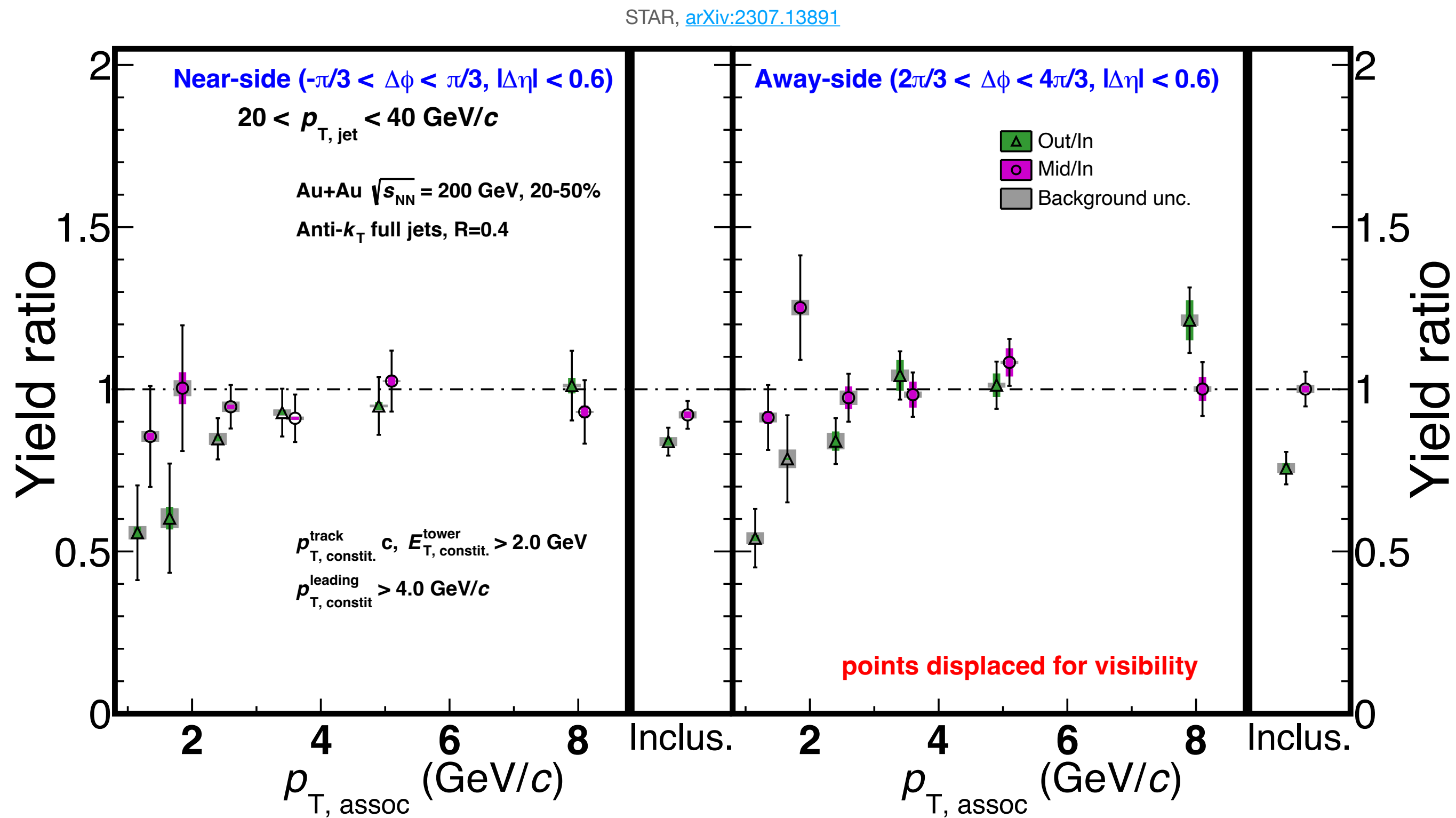
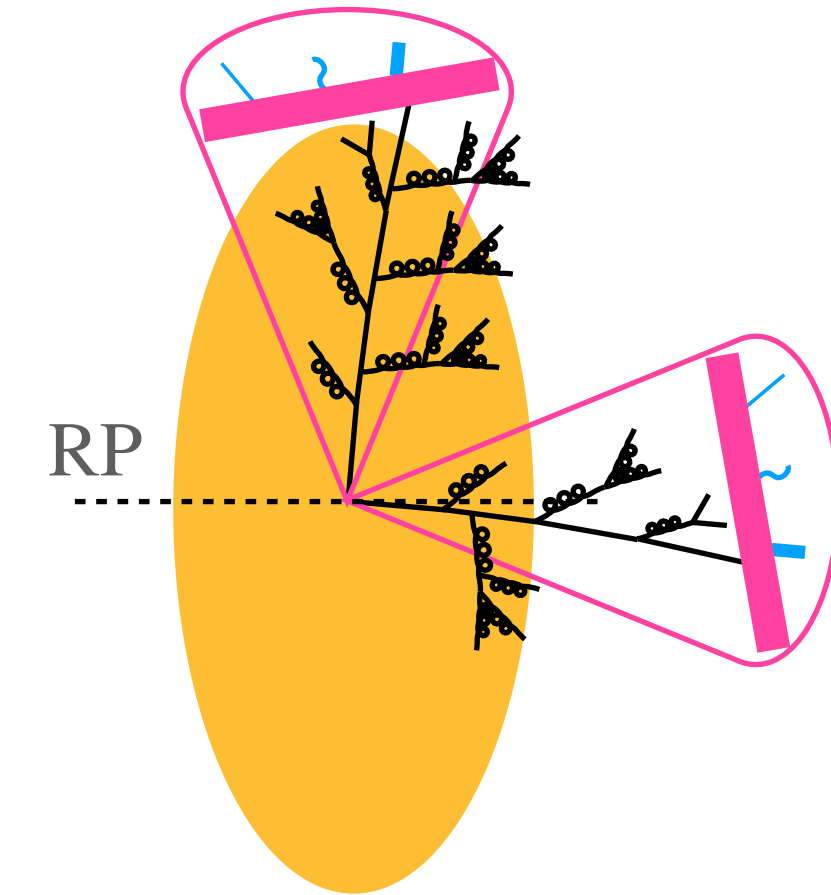
Separating p-QCD and np-QCD
with *energy correlators*

Energy flows

Path-length dependence of jet energy loss in medium
with *jet anisotropies (with respect to event plane)*

Energy-density dependence of jet energy loss in medium;
angular distribution of radiation in quenched jets
with *inclusive/semi-inclusive jet & high- p_T hadron yields*

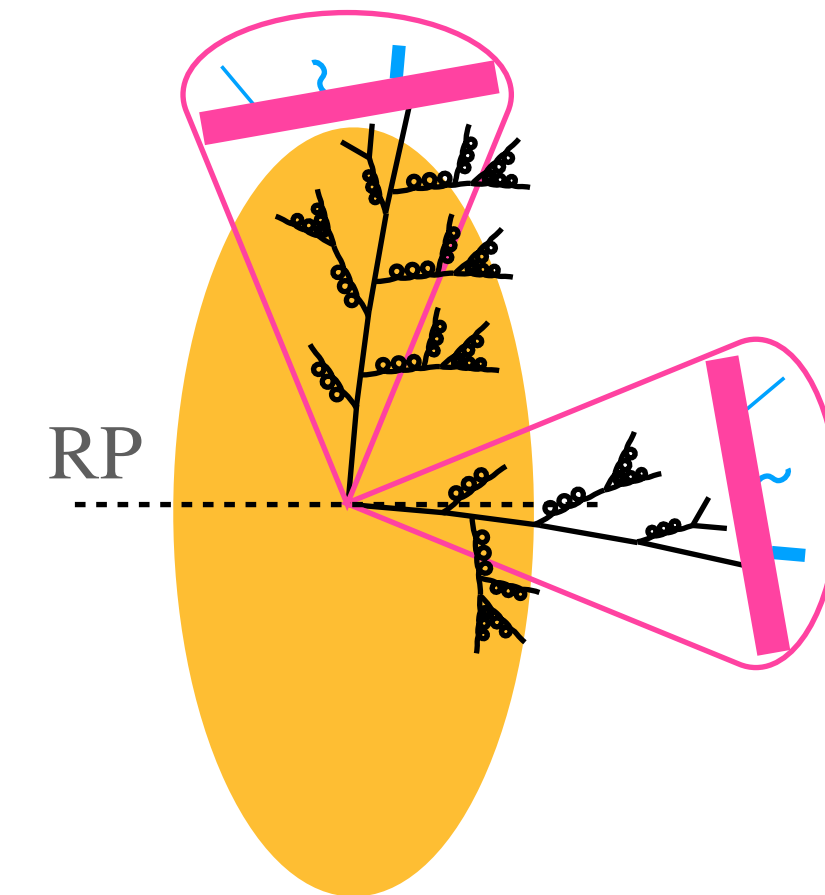
Event plane (EP) dep. of associated hadron yields



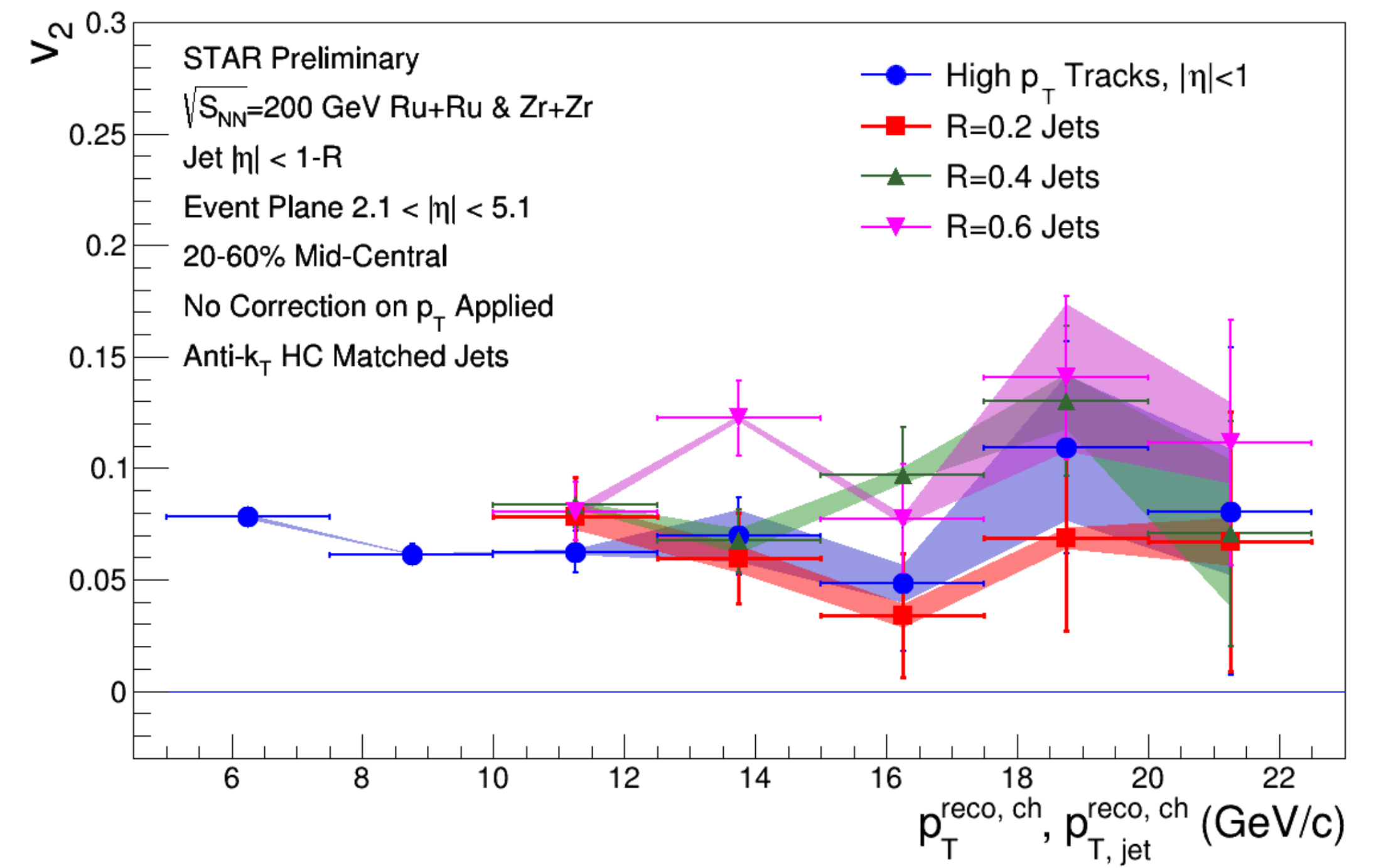
- *Expectation*: high (low)- p_T suppression (enhancement) for out-of-plane (OOP) vs. in-plane (IP) jets: path-length dependent quenching
- No significant deviation from unity within uncertainties
- Jet energy loss / medium density fluctuations spoiling effect?

Jet v_2

$$v_n(p_T, y) = \langle \cos(n(\phi - \Psi_{RP})) \rangle$$



- New forward detector at STAR, *EPD*, gives improved reaction plane (RP) resolution, no autocorrelation with mid-rapidity measurement
- v_2 in this context linked to **path-length dependent quenching**, not flow
- Clear v_2 signal, independent of jet R , p_T , in high-statistics isobar data



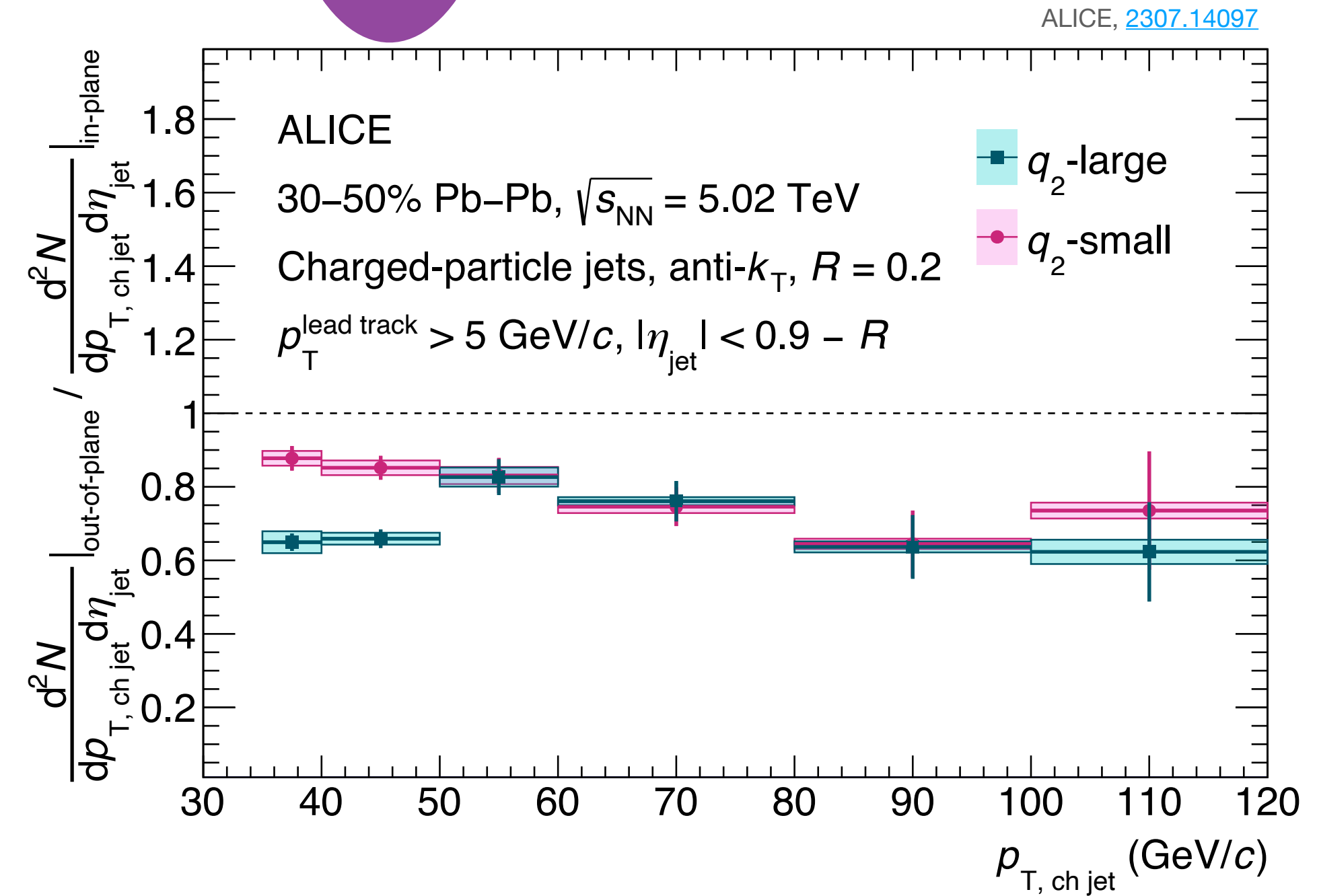
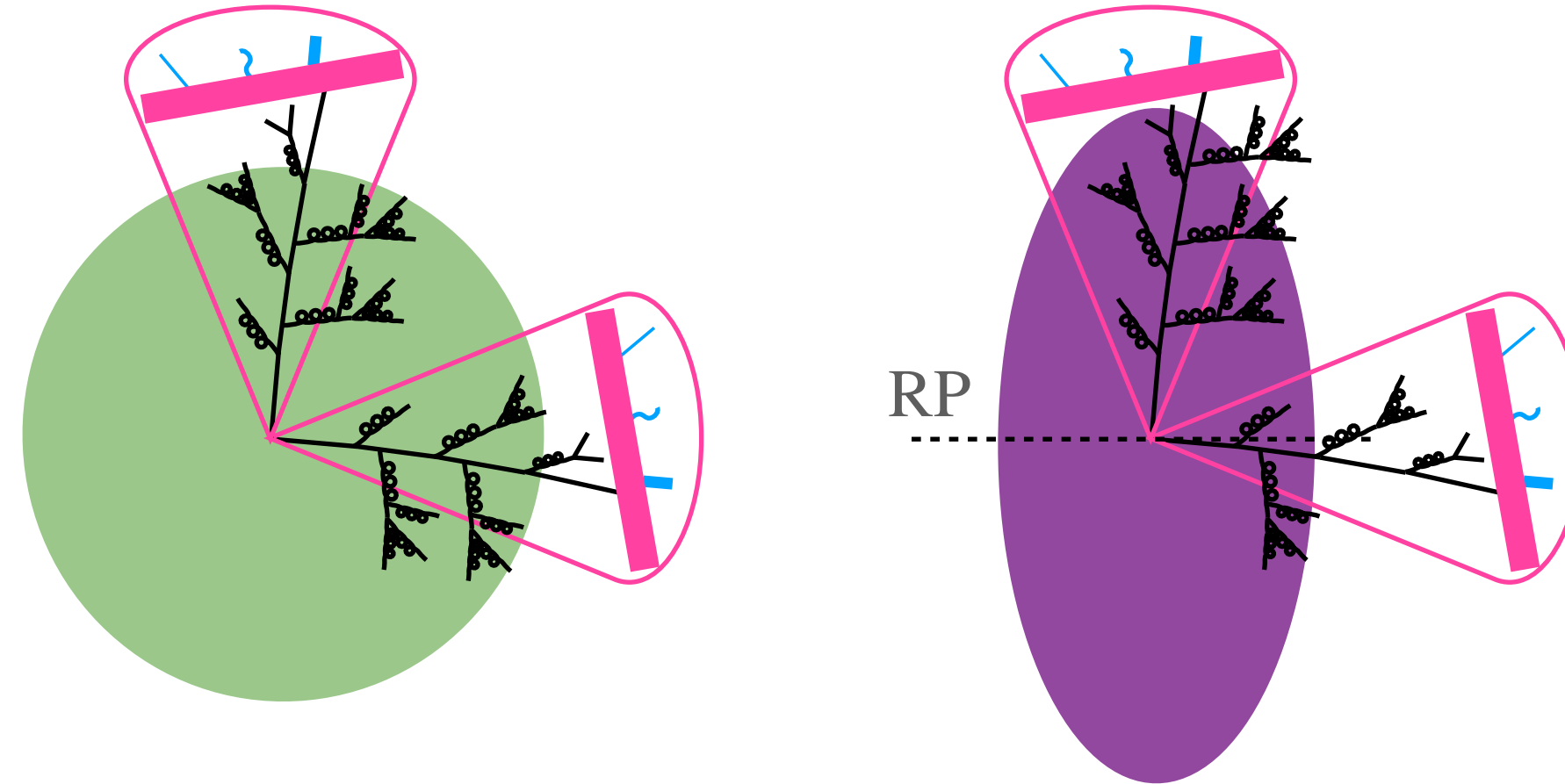
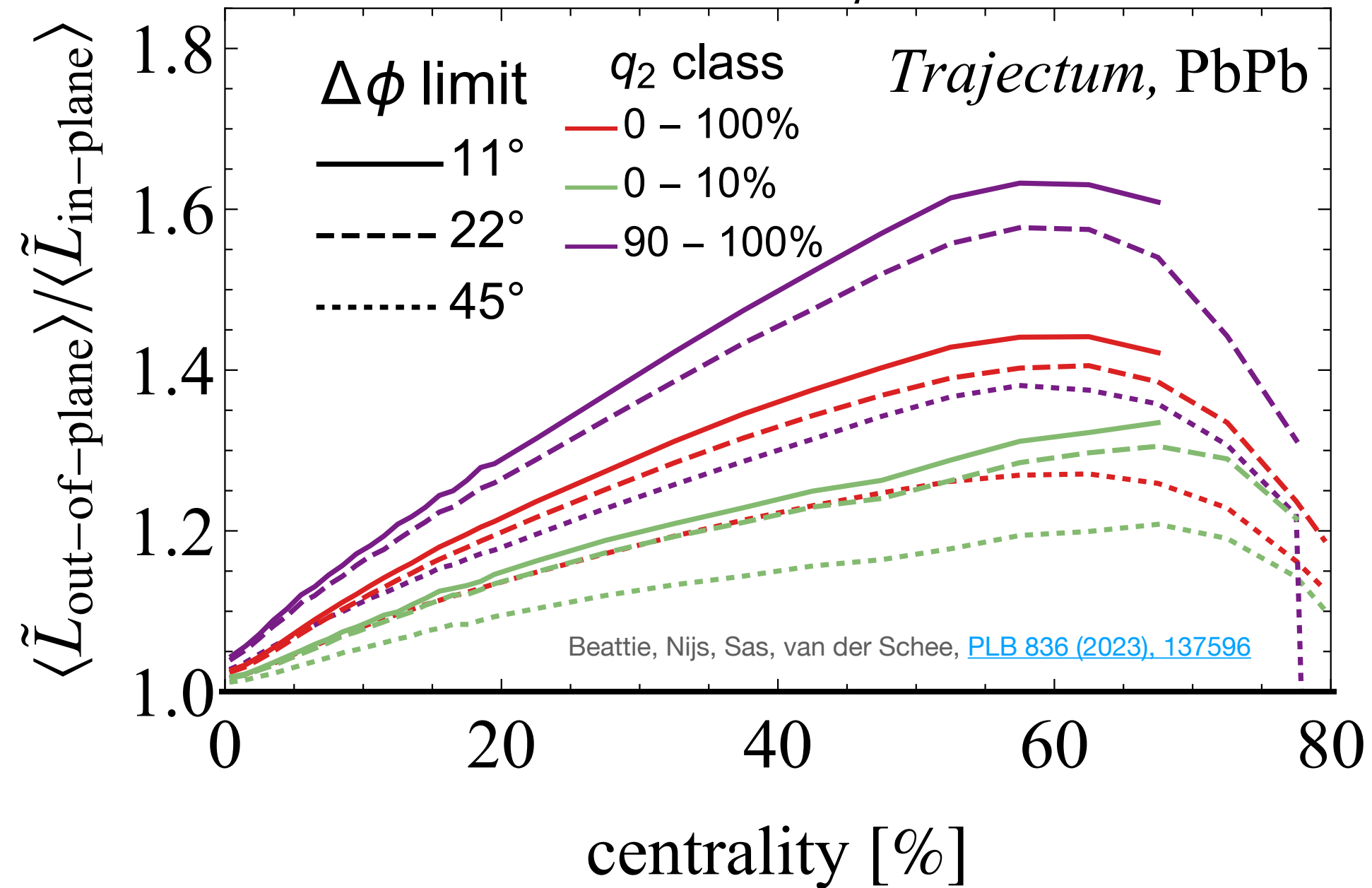
Event-shape engineering¹

$$Q_2 = \left(\sum_{i=1}^M w_i \cos(2\phi_i), \sum_{i=1}^M w_i \sin(2\phi_i) \right), \quad q_2 = |Q_2|/\sqrt{M},$$

w_i : nMIP weight, M : multiplicity

$$v_2 = \langle \cos(2(\phi - \Psi_2)) \rangle$$

$$\tilde{L} = \int 1/\gamma u_\mu dL^\mu$$



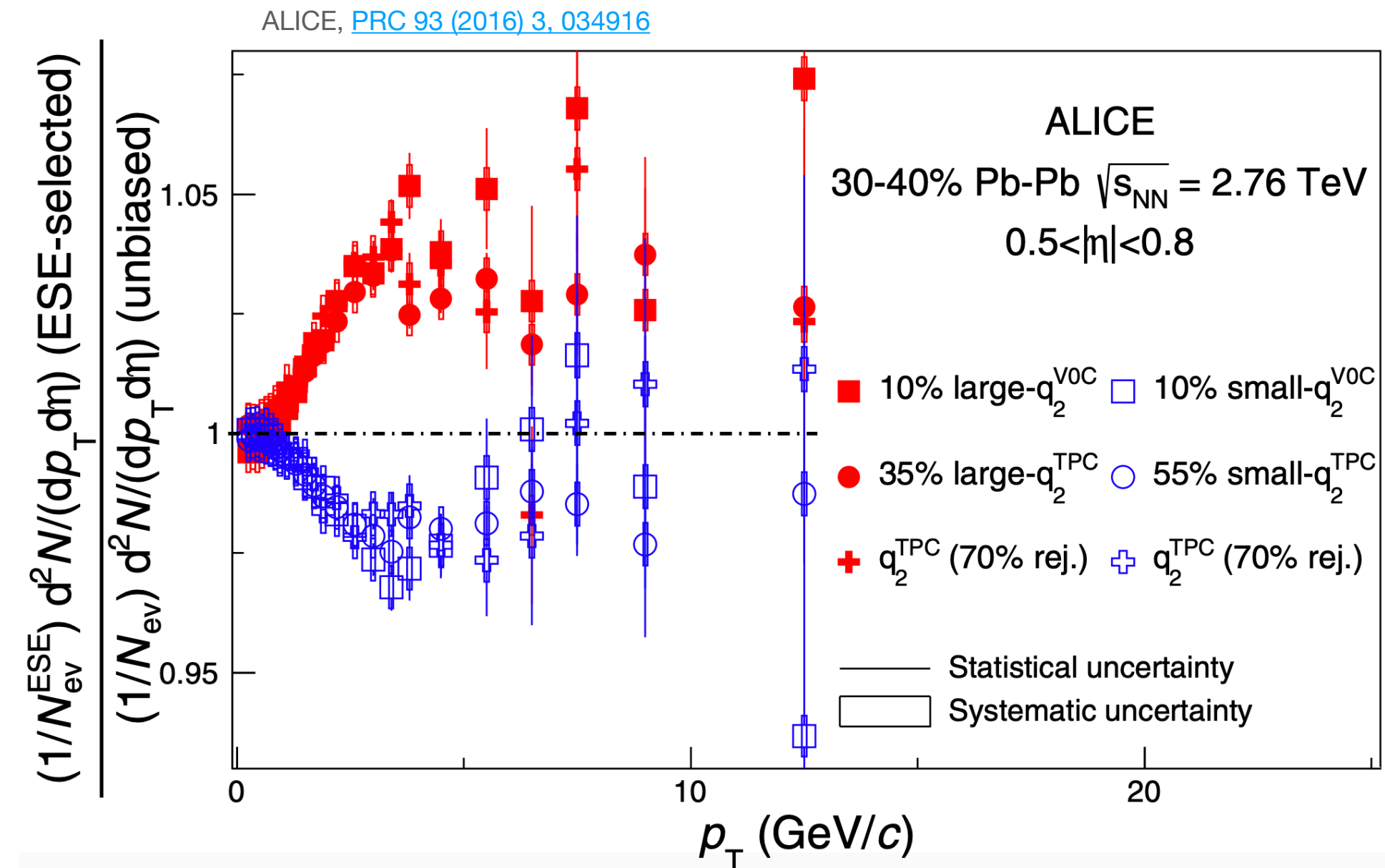
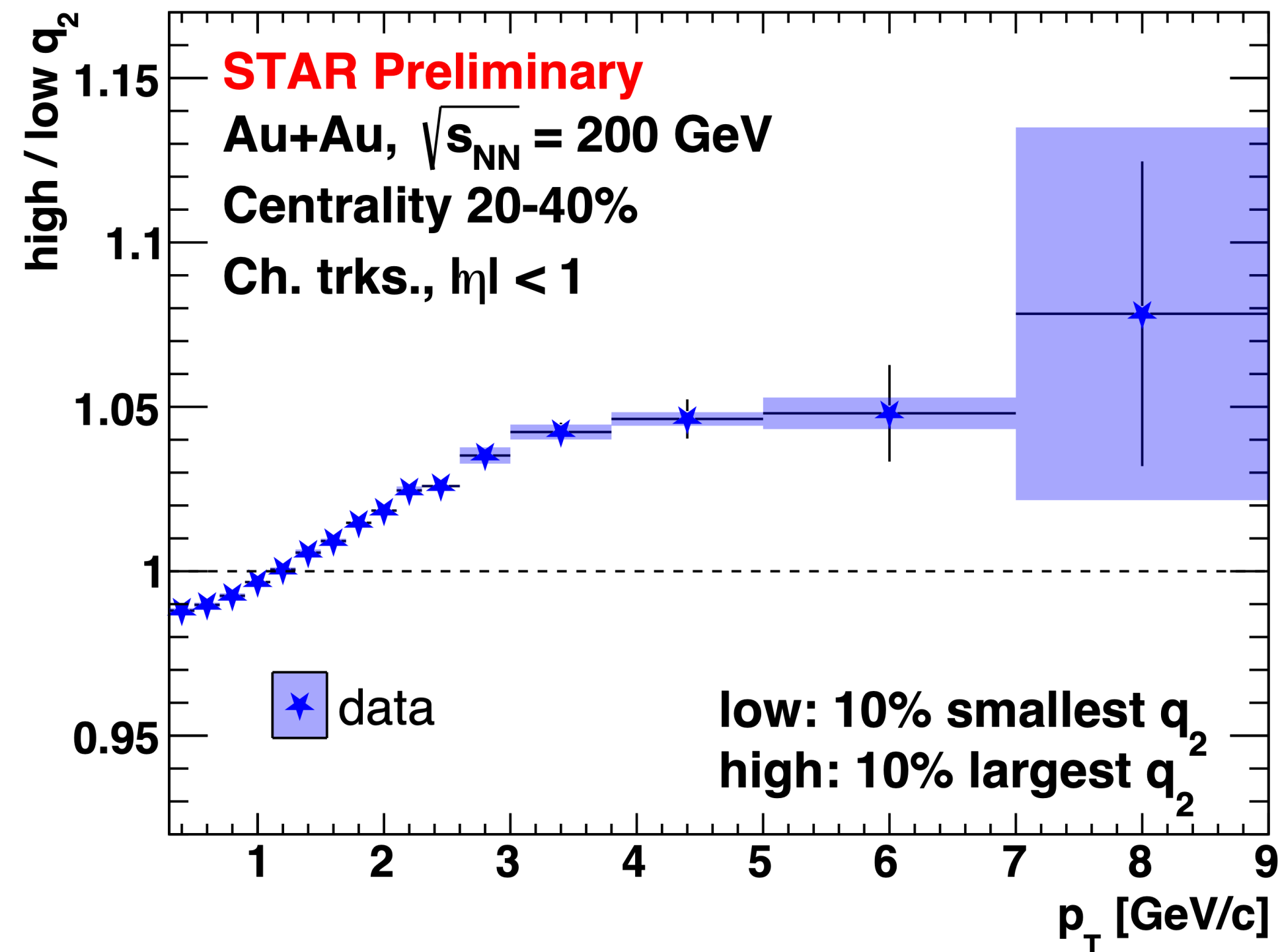
- Select on the shape using reduced flow vector q_2
→ average path length difference in Trajectum model
- ALICE: low- p_T difference in IP, OOP yields for highly elliptical events

Event-shape engineering

$$Q_2 = \left(\sum_{i=1}^M w_i \cos(2\phi_i), \sum_{i=1}^M w_i \sin(2\phi_i) \right), \quad q_2 = |Q_2|/\sqrt{M},$$

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$$v_2 = \langle \cos(2(\phi - \Psi_2)) \rangle$$



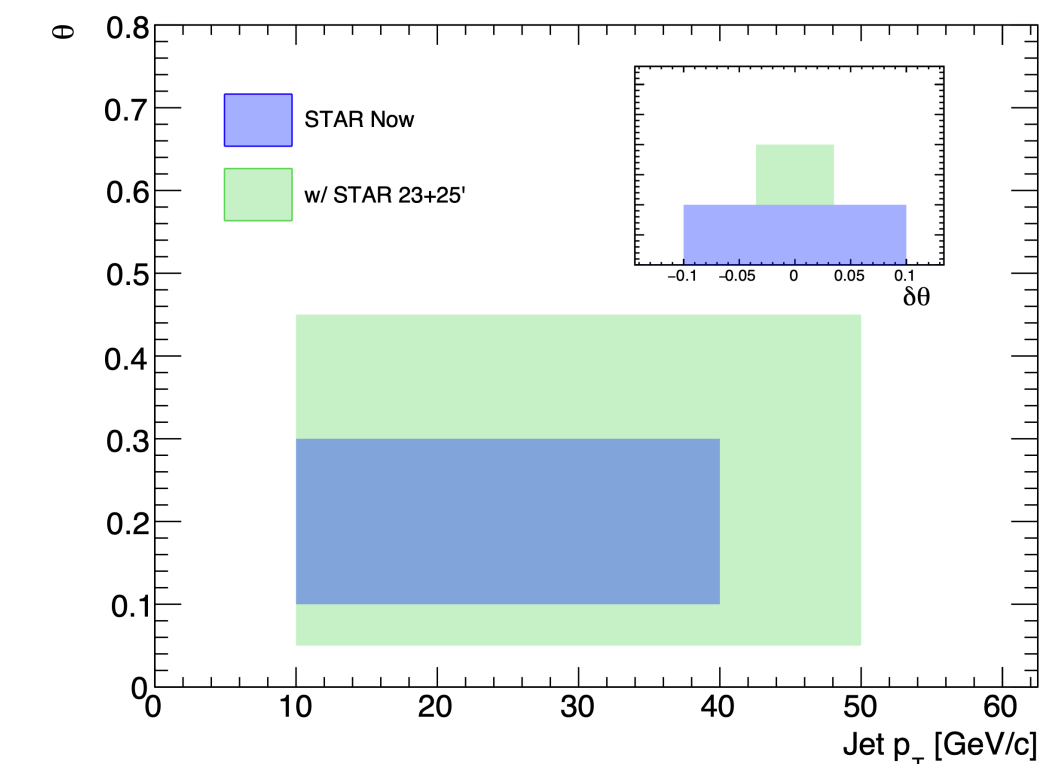
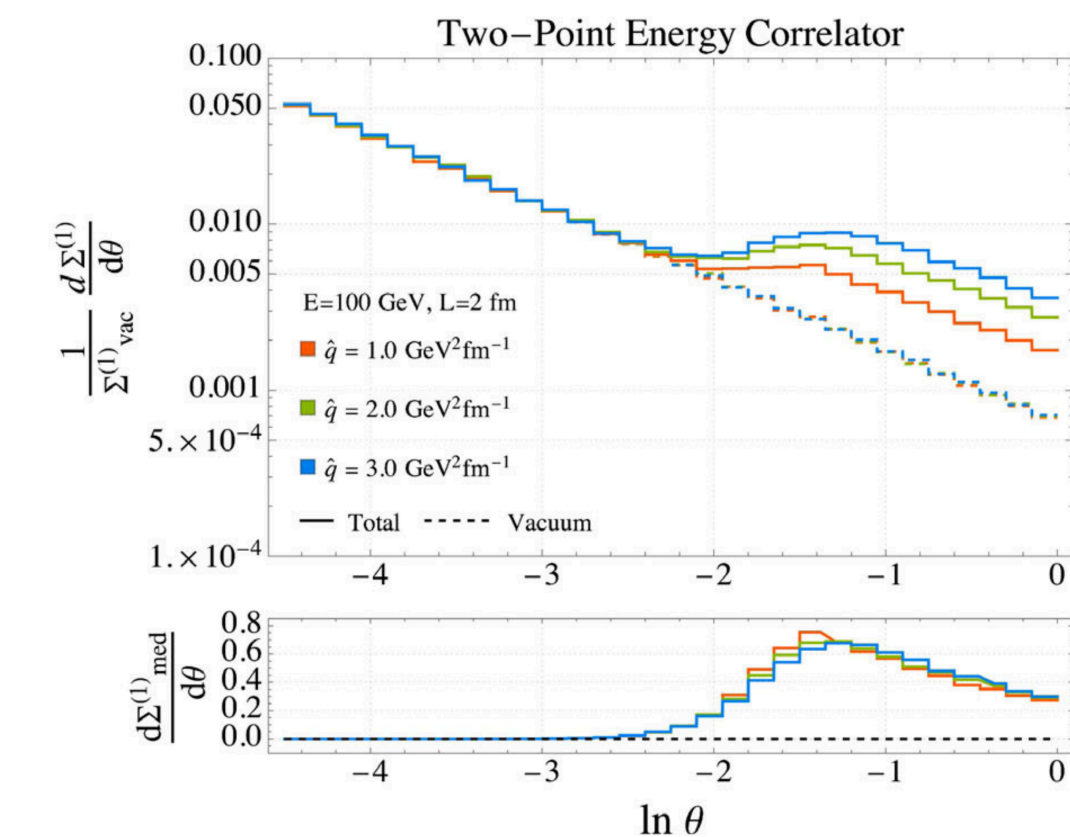
- STAR analysis ongoing — without selecting on EP angle, see enhancement at mid- p_T of charged track yields, for high vs. low q_2 events
- Interplay between eccentricity/density, elliptic/radial flow. Also observed by ALICE

Future prospects

Energy flows

Andres, Dominguez, Kunnawalkam Elayavalli, Holguin, Marquet, Moul, [PRL 130 \(2023\) 26, 262301](https://arxiv.org/abs/2208.12301)

- *Generalized angularities*: less conservative systematic uncertainties, extension to jet momentum profile $\rho(r)$
- *EECs*: higher orders; charge-dependent; in heavy-ion collisions
- R_{AA} : analyzing R_{pAu}
- *Jet v_2* : extended to OO collisions, studying non-flow contribution
- *Event shape engineering*: event-plane angle dependence study in progress
- *Runs 23+25^{1,2}*: expected $\sim 3x$ increase in statistics relative to current AA analyses w/ Run 14 \rightarrow improved uncertainties e.g. for $\gamma_{dir}+jet$ I_{AA} , and **kinematic reach / overlap with LHC**



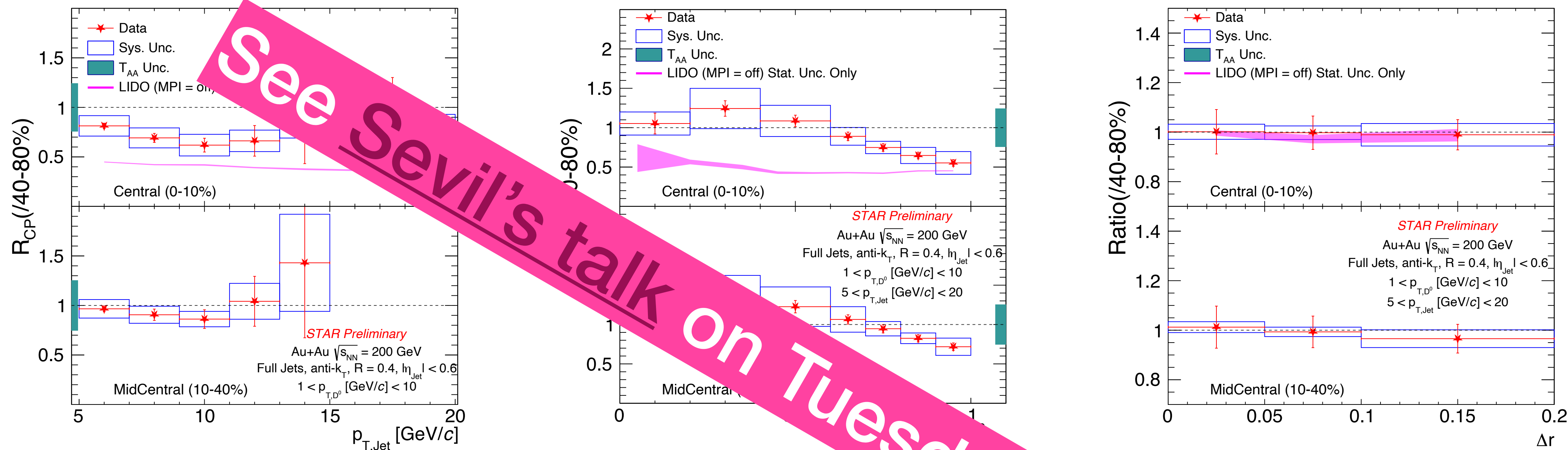
Charm quark energy loss, diffusion,
fragmentation modification in medium
with *charmed-jet yields*

Hadronization mechanism
with *flavor correlators*

Constituent identity

Hadrochemistry modification via
medium response
with *baryon-to-meson ratios*

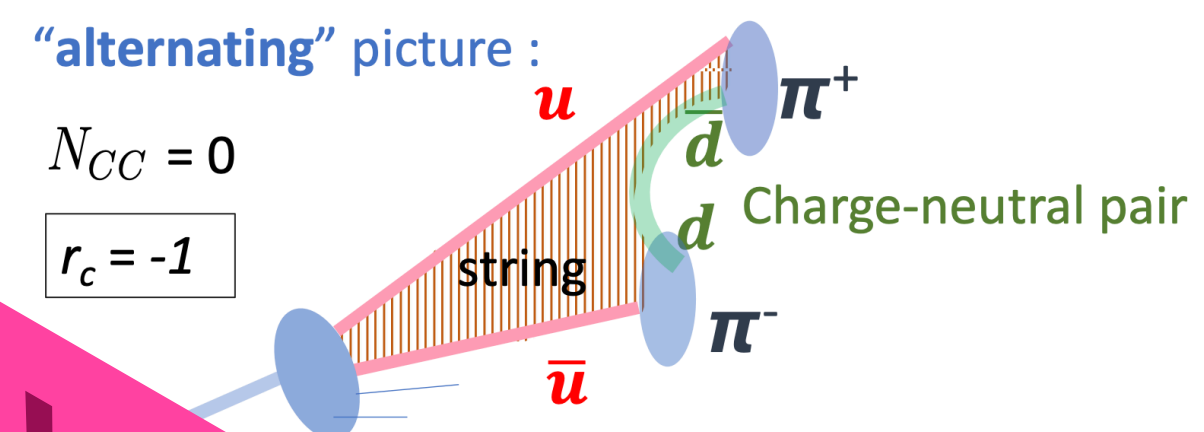
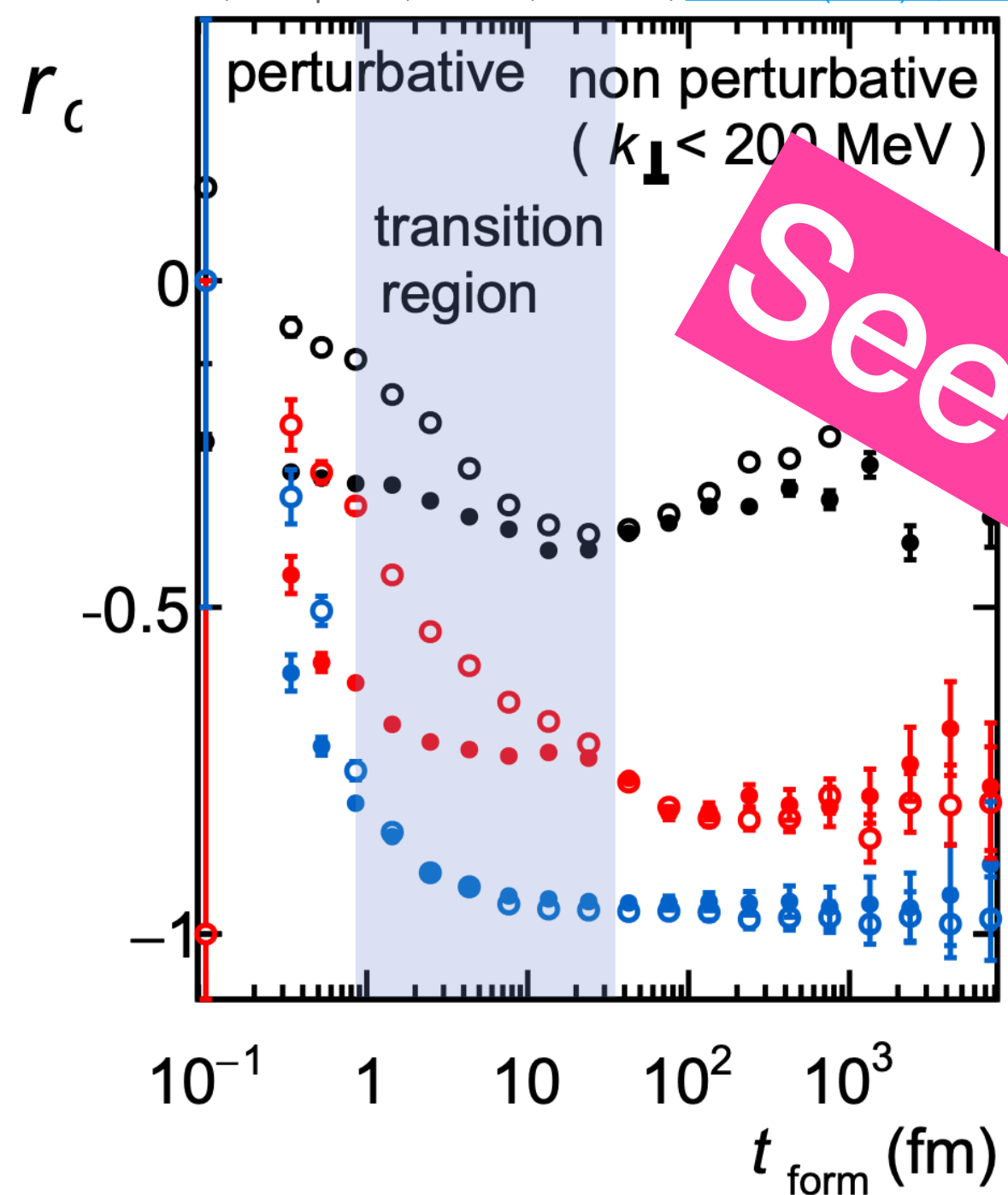
D⁰-jet spectra, profile, fragmentation



- Testing charm quark energy loss, diffusion, and fragmentation modification
- *Hint of suppression of yield at low- p_T . Hard-fragmenting charm quarks are suppressed. No diffusion.*
- Model including radiative and collisional energy loss during heavy quark propagation underpredicts central yields — MPI might be important for D⁰ p_T this low

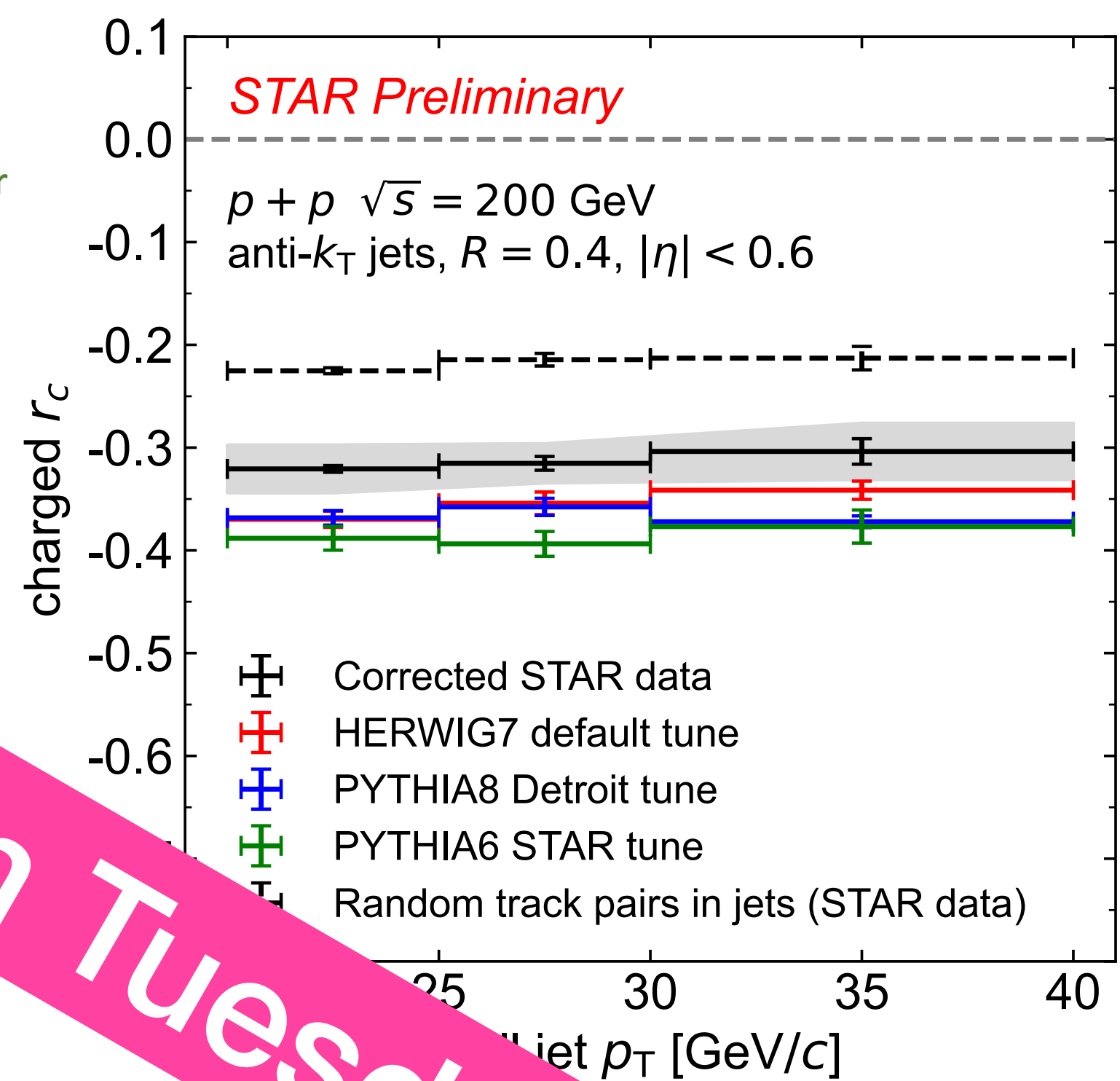
Flavor correlators in jets

Chien, Deshpande, Mondal, Sterman, [PRD 105 \(2022\) 5, L051502](#)



$$r_c(X) = \frac{d\sigma_{h_1 h_2} / dX - d\sigma_{h_1 \bar{h}_2} / dX}{d\sigma_{h_1 h_2} / dX + d\sigma_{h_1 \bar{h}_2} / dX}$$

$$r_c = \frac{N_{SS} - N_{OS}}{N_{SS} + N_{OS}}$$



bath of charges with no net charge

jet with ~no net charge

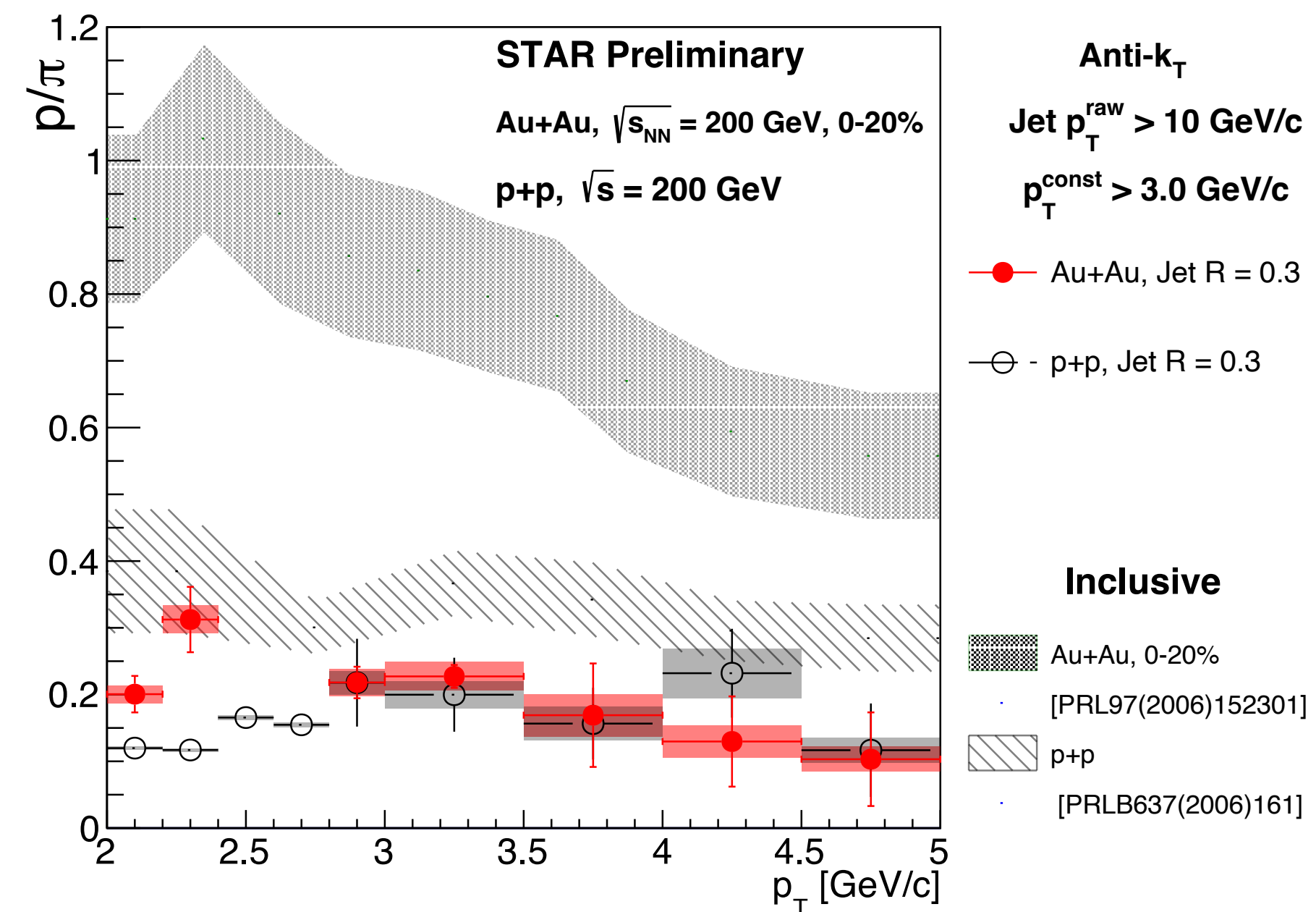
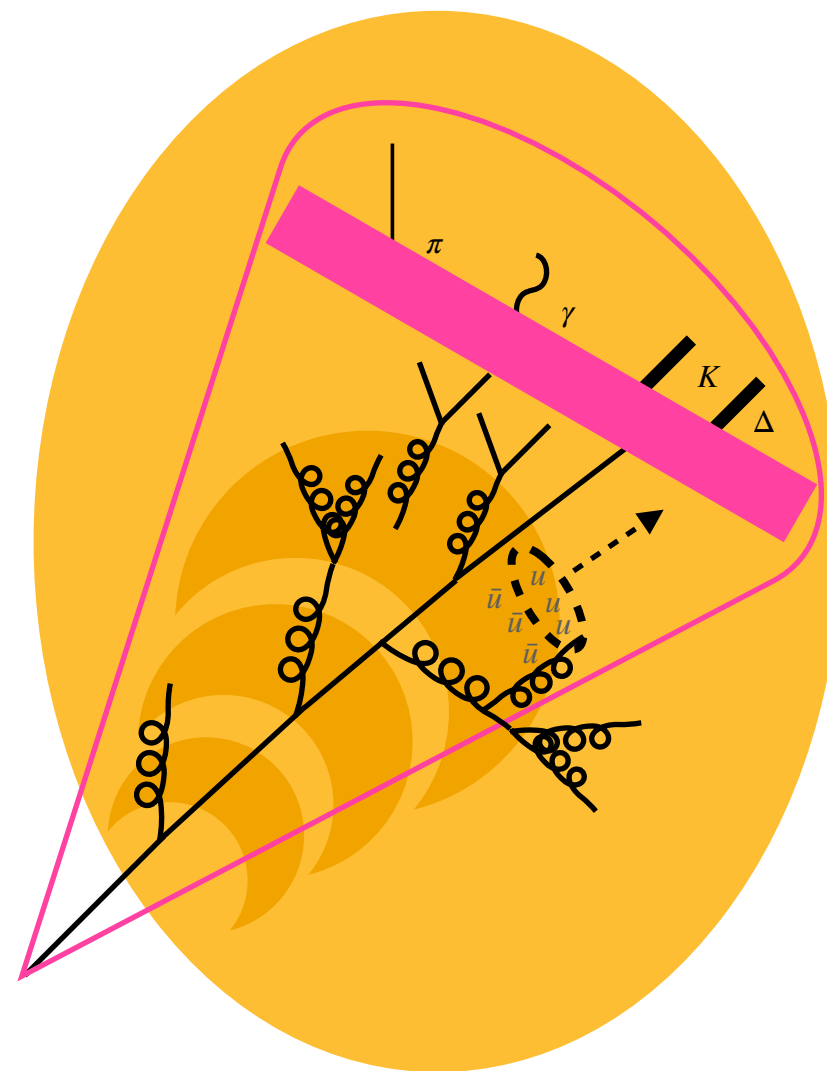
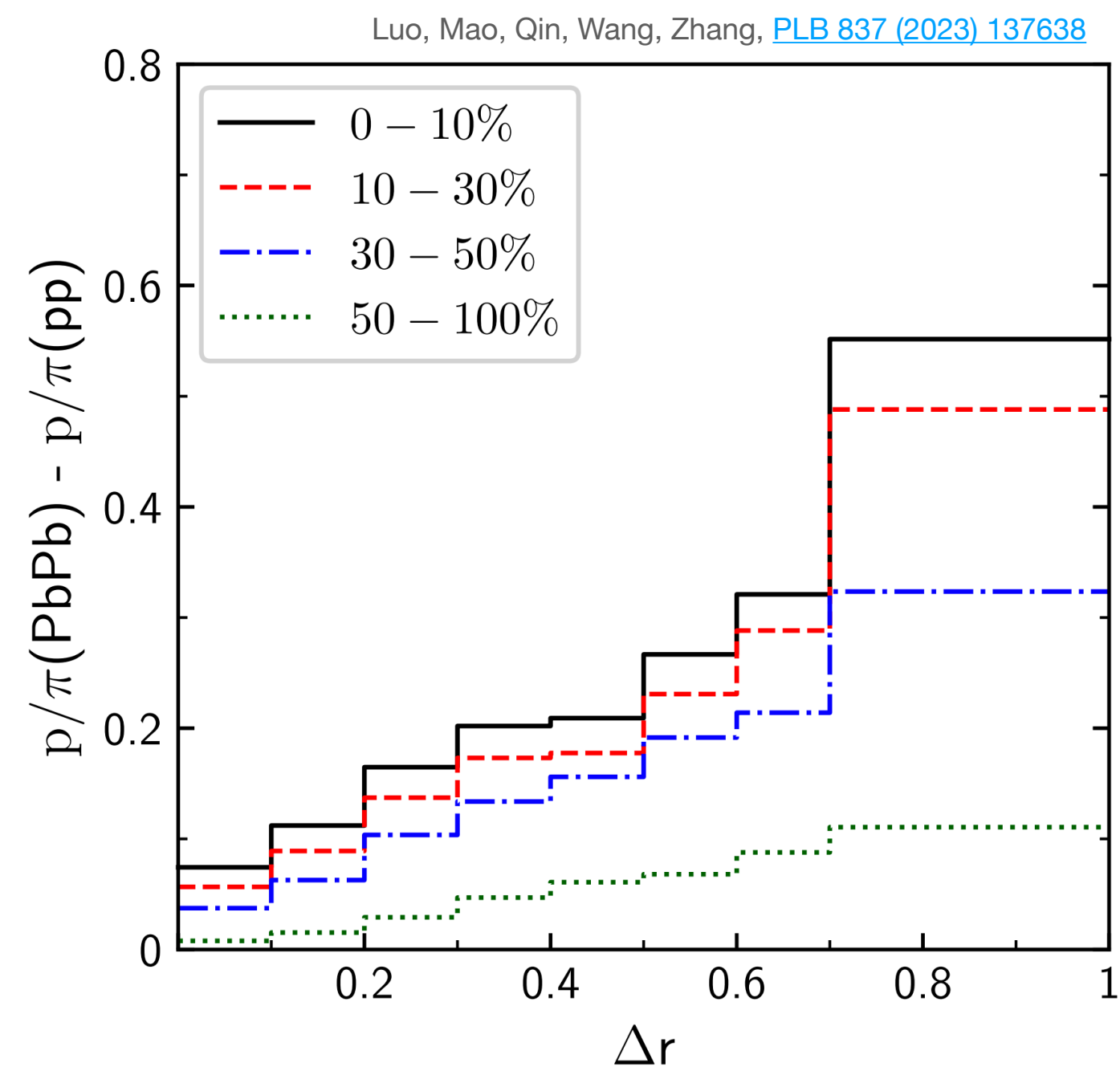
perfectly correlated

See Youqi's talk on Tuesday at 11 AM!

- r_c can probe contribution of string-like fragmentation
- First measurement in pp: Pythia predicts more string-like fragmentation than supported by data, but difficult to conclude with default tune

Baryon-to-meson ratios

Signature of medium response?



- Possible sign of parton coalescence in jet: enhanced baryon-to-meson ratio in AA
- No observed modification of *in-jet* p/π ratio for $R = 0.3$ jets

Future prospects

Constituent identity

- *D⁰-jet*: adding another dataset to increase statistics; adding generalized angularities; tightening D⁰ p_T threshold

- r_C : extension to heavy-ion collisions underway

- Herwig tune to RHIC kinematics ongoing

- *Baryon-to-meson ratios*: studying dependence on constituent p_T threshold

Esha, [Hard Probes 2023](#)

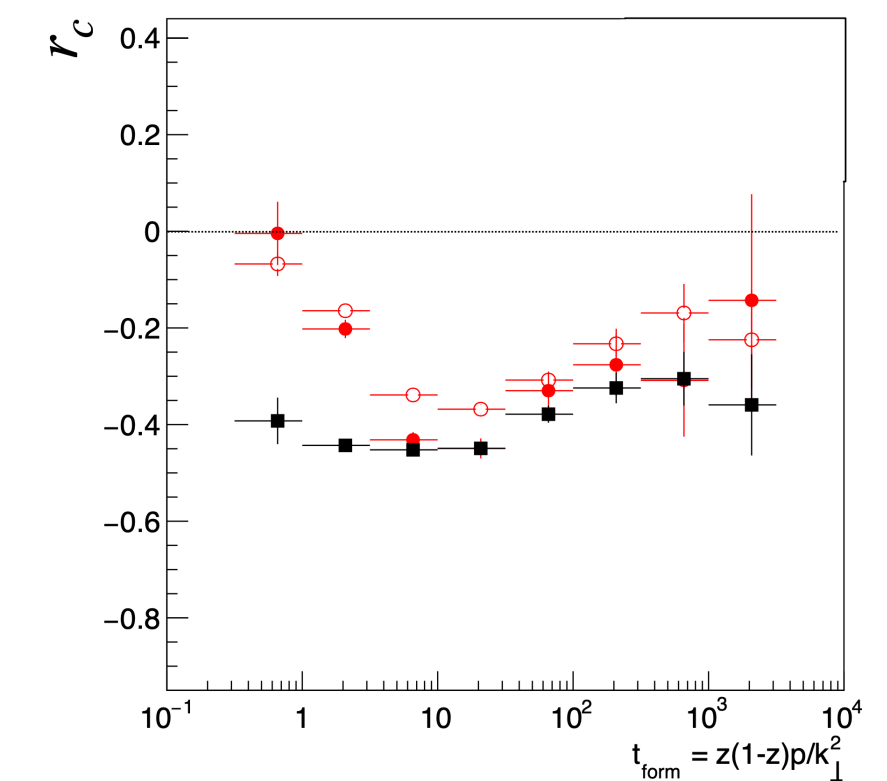
HIJING, Au+Au $\sqrt{s_{NN}} = 200$ GeV

π^\pm Quenched Unquenched

PYTHIA, p+p $\sqrt{s} = 200$ GeV

π^\pm ■

$R = 0.4, p_T^{\text{jet}} > 10$ GeV/c

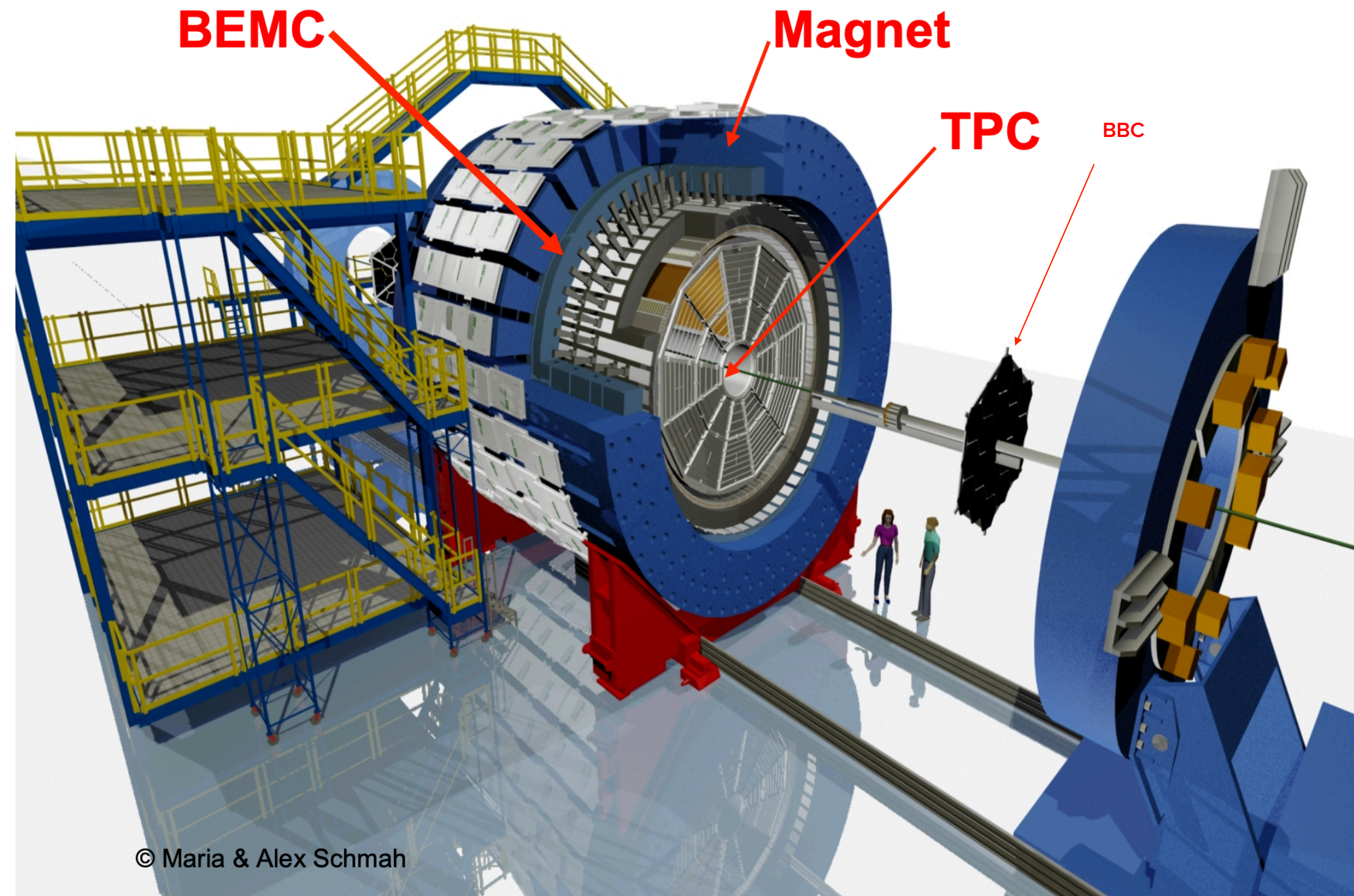


What we've learned

- **Precision era of jet substructure:** many-dimensional corrections and correlations, systematically mapping the phase space for QCD radiation in vacuum at lower \sqrt{s}
- **First measurements** of new observables **EECs and r_c** separate perturbative and non-perturbative physics cleanly for **improved theoretical control**
- Demonstrated **strong dependence of quenching** on N_{part} (**~similar energy density**) across collision species; **more energy lost at RHIC** than LHC, relative to jet p_T ; **jet profile broadening**, with radiation roughly recovered by ~ 0.5 radians; and **finite jet v_2** . **No quenching** observed in **pAu** collisions.
- **No medium-induced hadrochemistry** effect observed. **Suppression of jets with hard-fragmenting charm hadrons** but as yet no observed corresponding enhancement of soft-fragmented charm jets or diffusion to broader angles

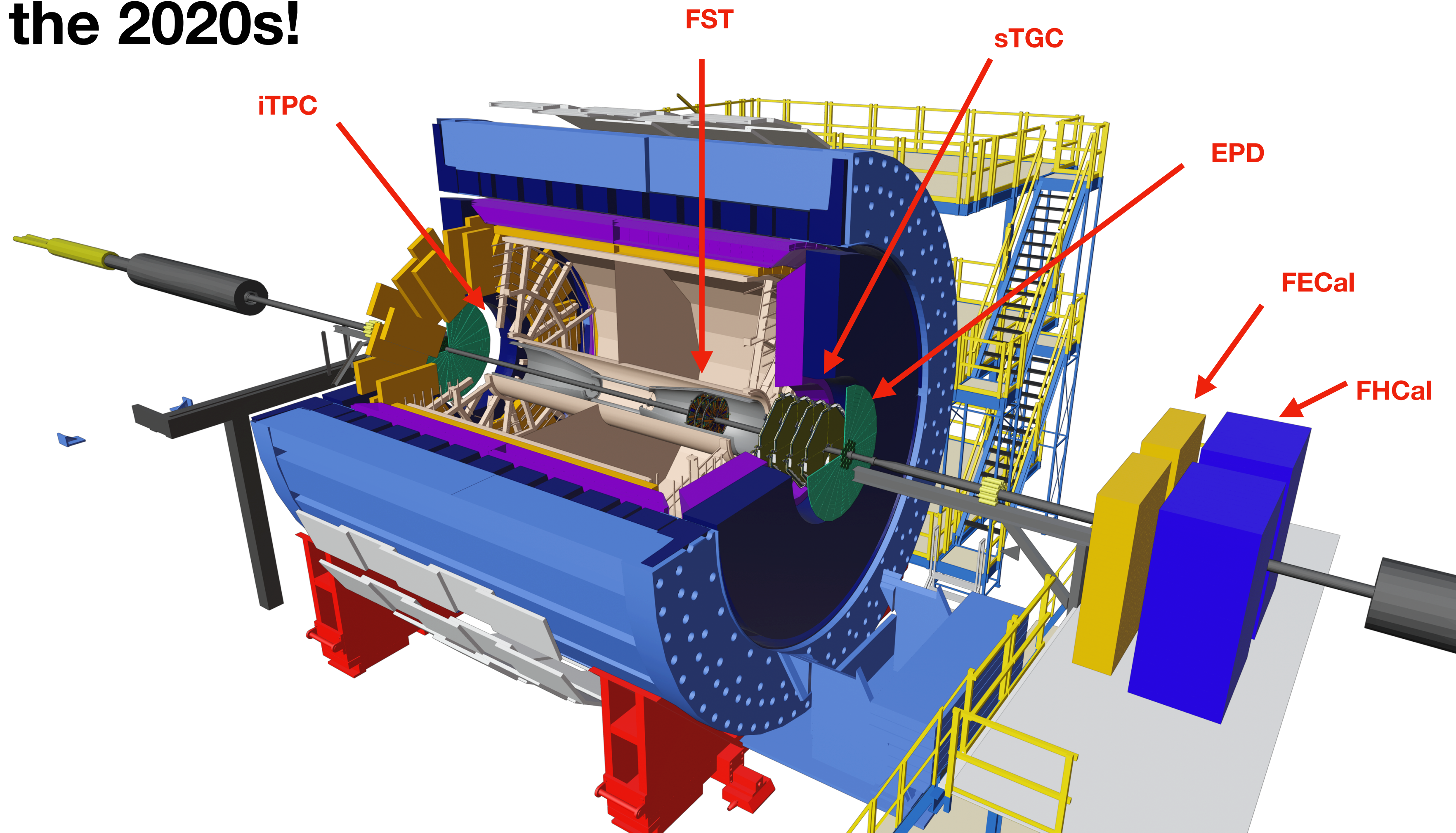
Jets at STAR

In the 2010s



Jets at STAR

In the 2020s!



Precision tracking

*Forward jets →
different x ; q v. g*

*Unbiased centrality/
EP determination*

Etc!



Backup

STAR Zero Degree Calorimeters

