

Electromagnetic probes in heavy-ion collisions



PROCEEDINGS
OF SCIENCE

Jean-François Paquet

February 14, 2024

**Electromagnetic probes in heavy-ion collisions:
progress and open questions**



Jean-François Paquet^{a,b}

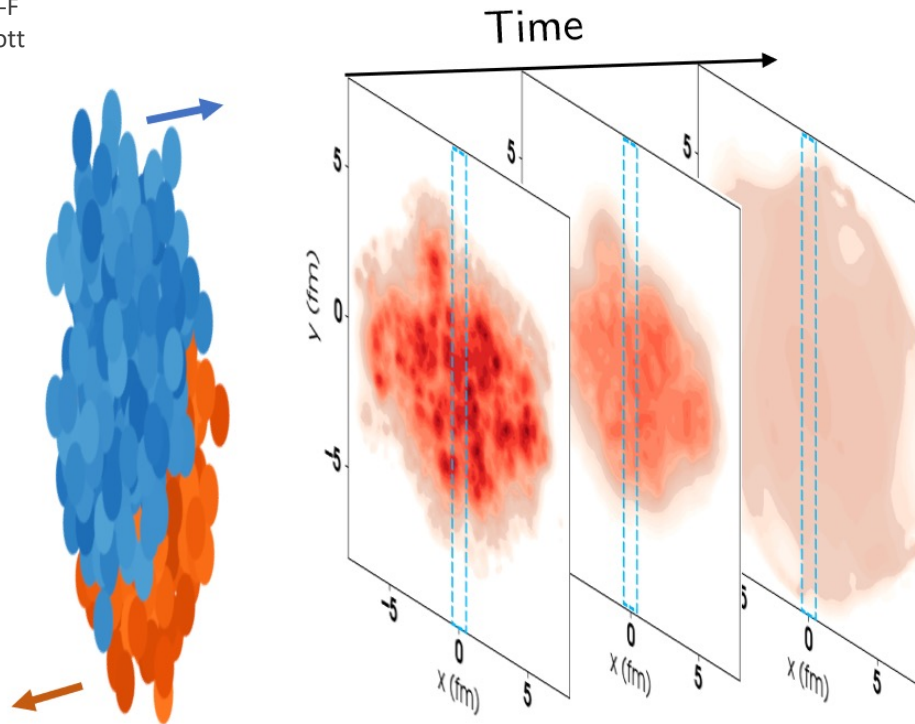
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Nashville TN 37212*

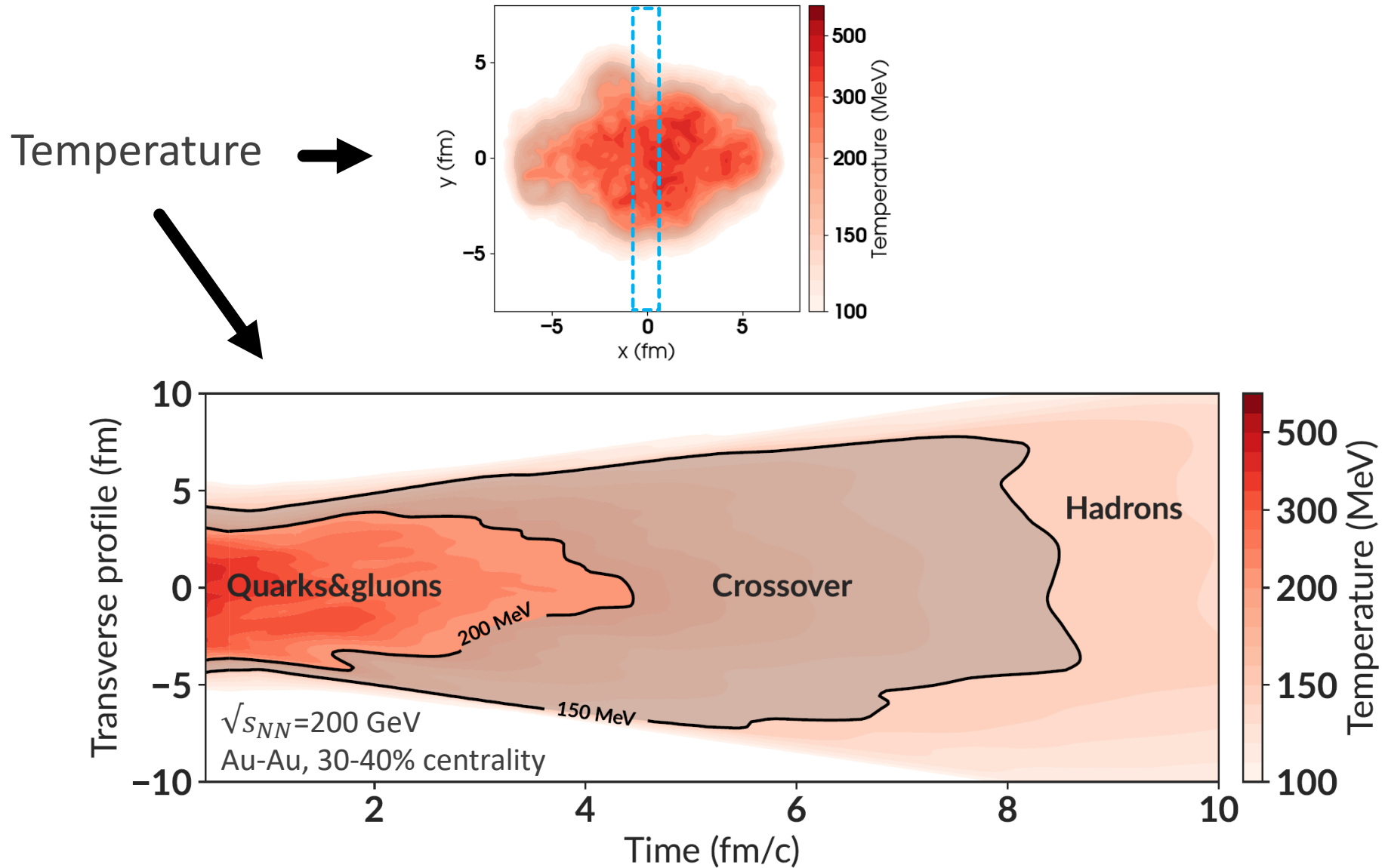
E-mail: jean-francois.paquet@vanderbilt.edu

Spacetime profile of heavy-ion collisions

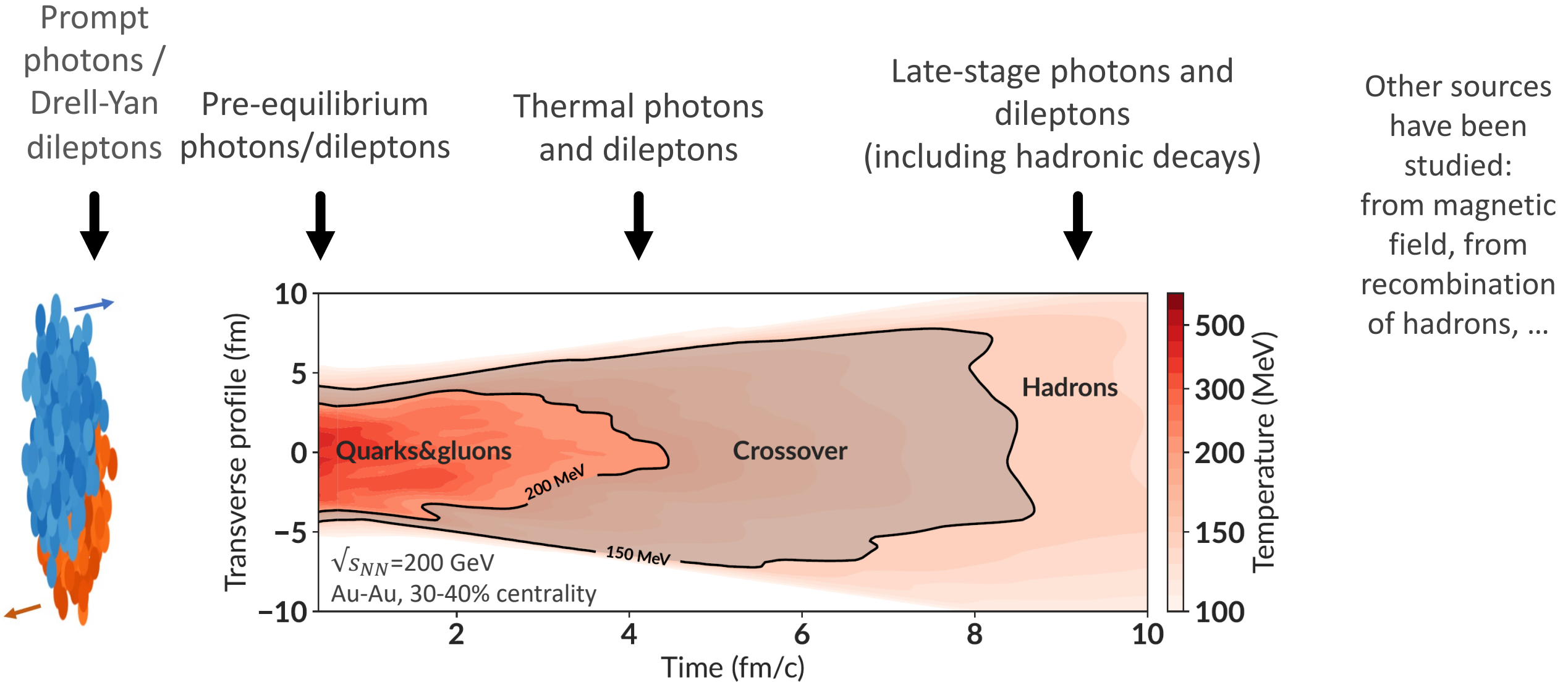
Figure credit: J-F
Paquet and Scott
Moreland

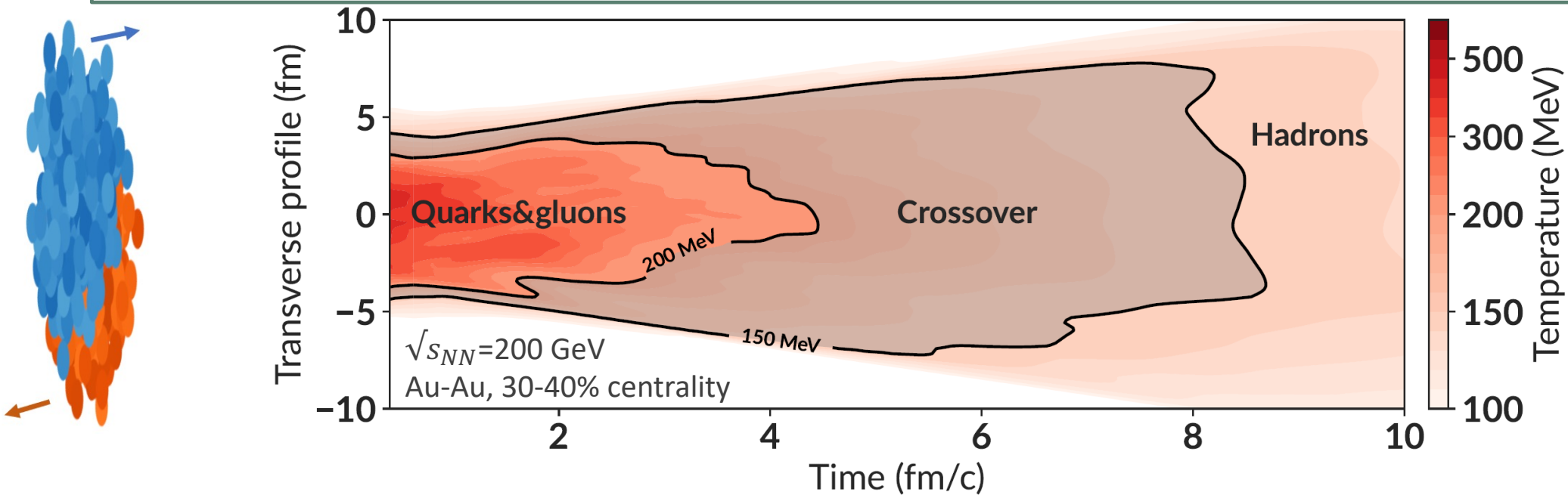
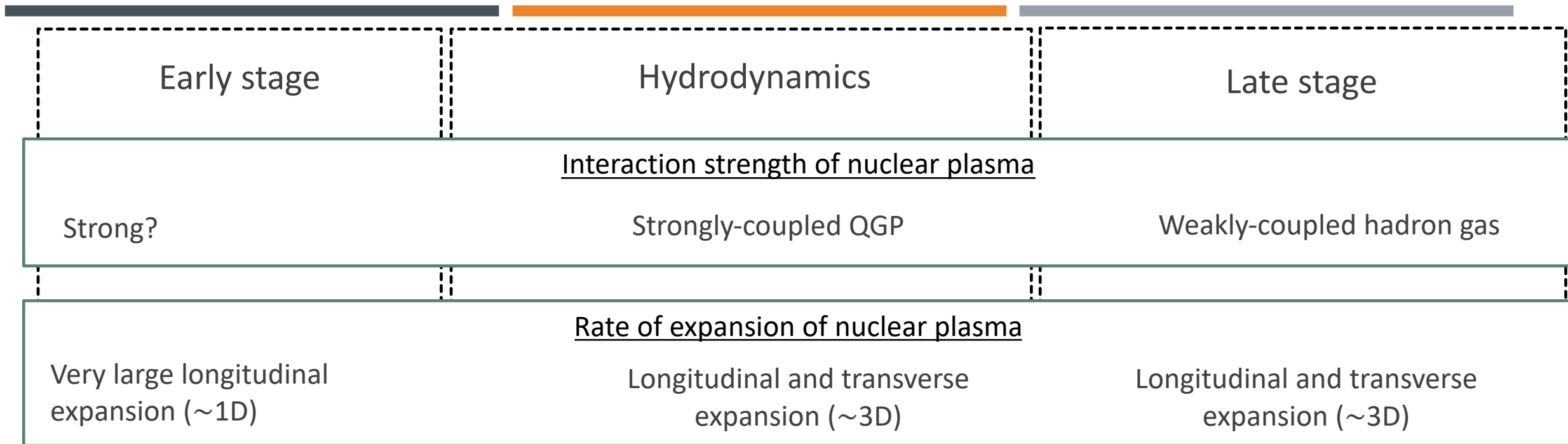


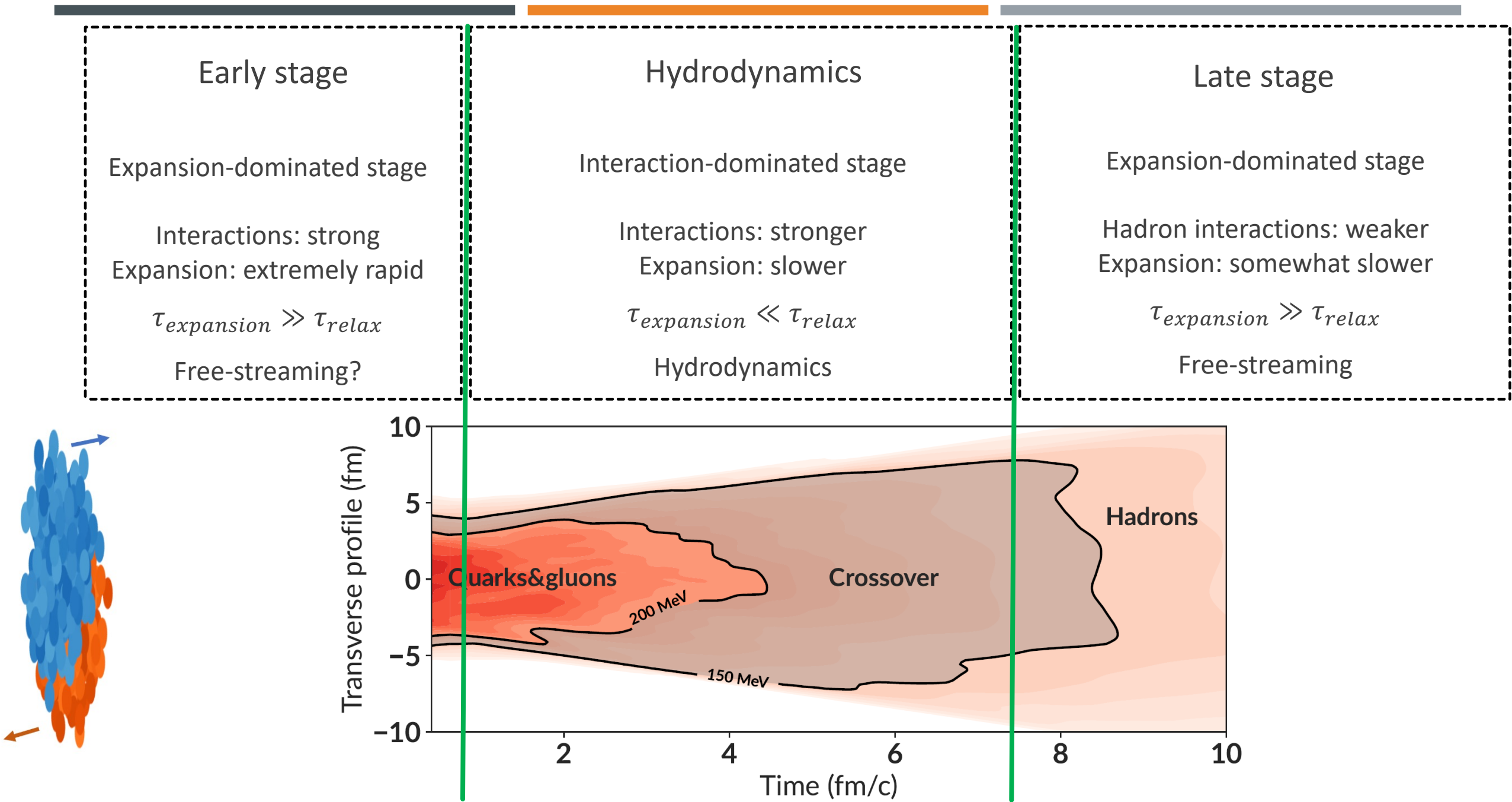
Spacetime profile of heavy-ion collisions

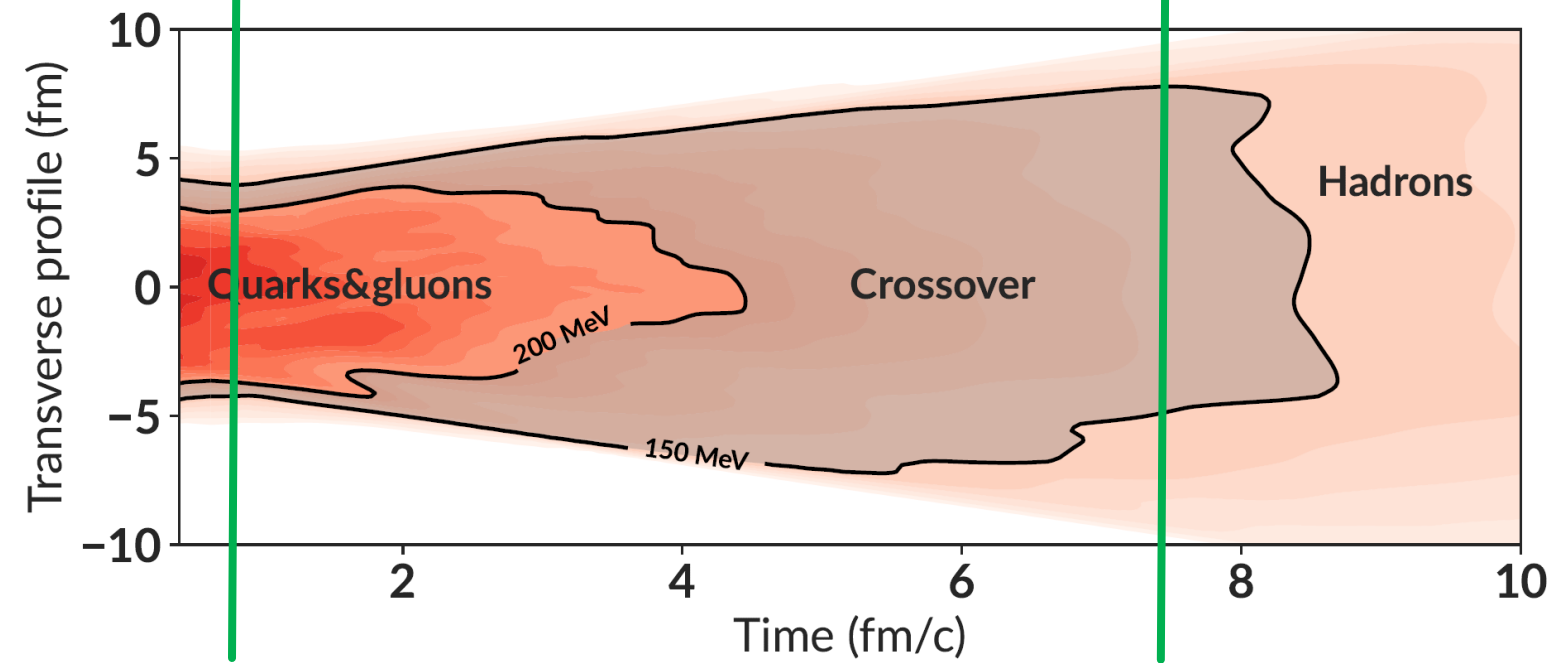
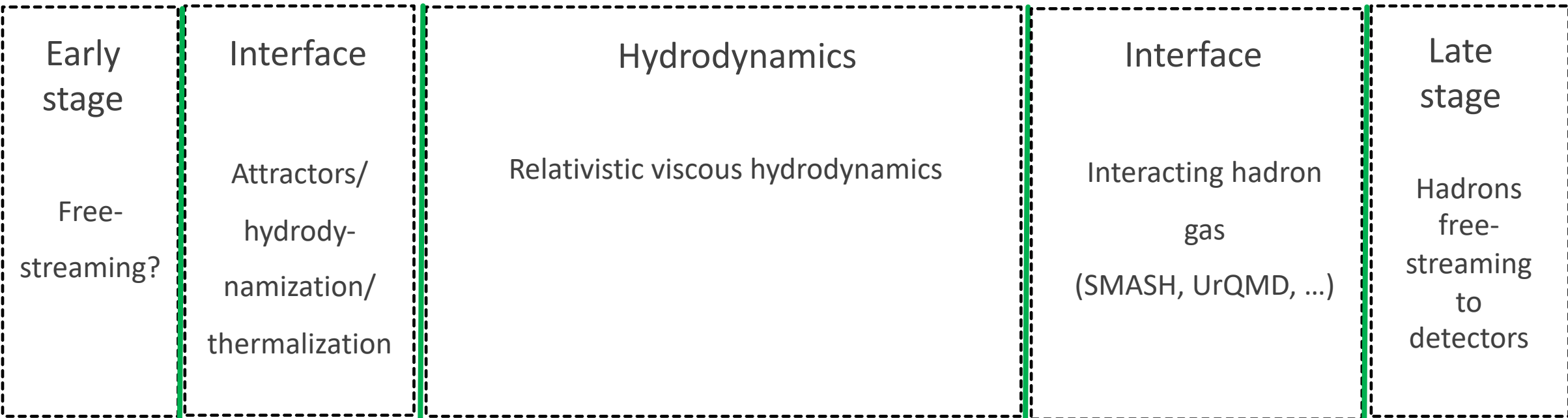


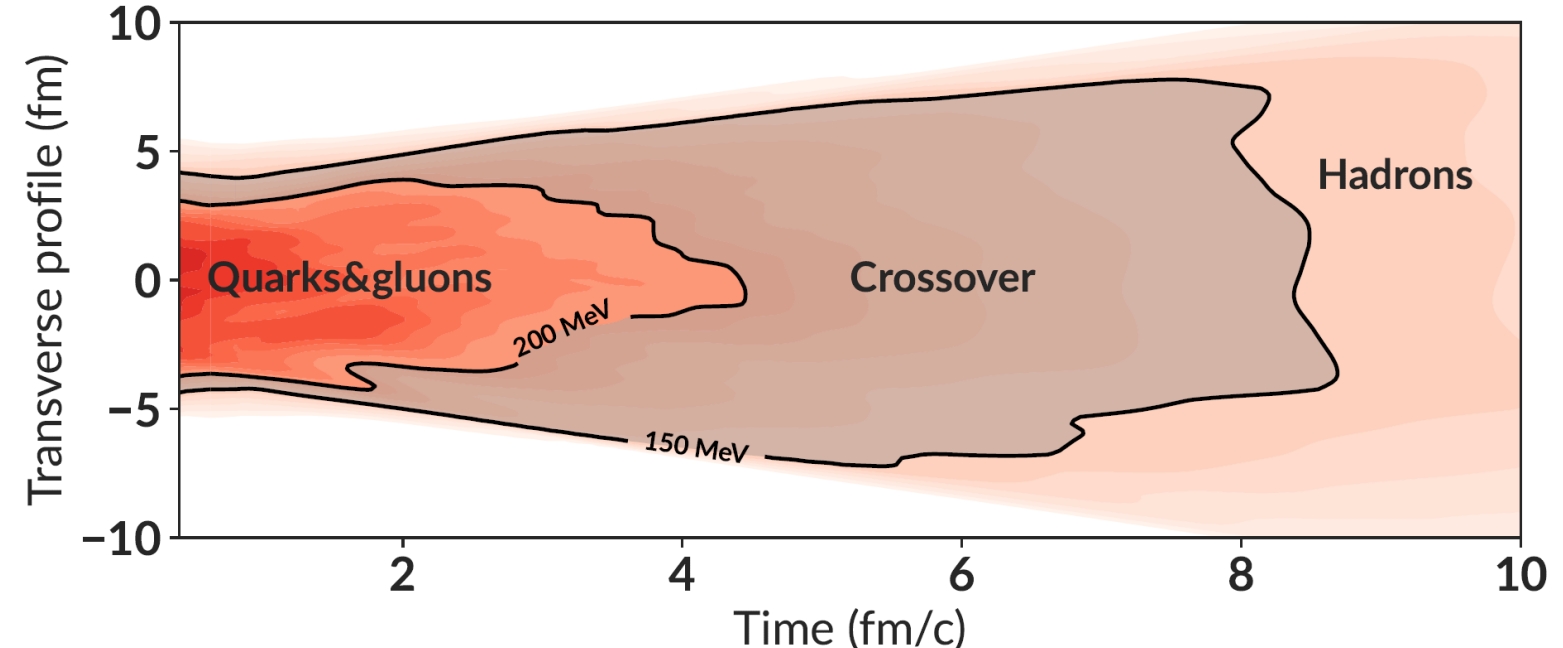
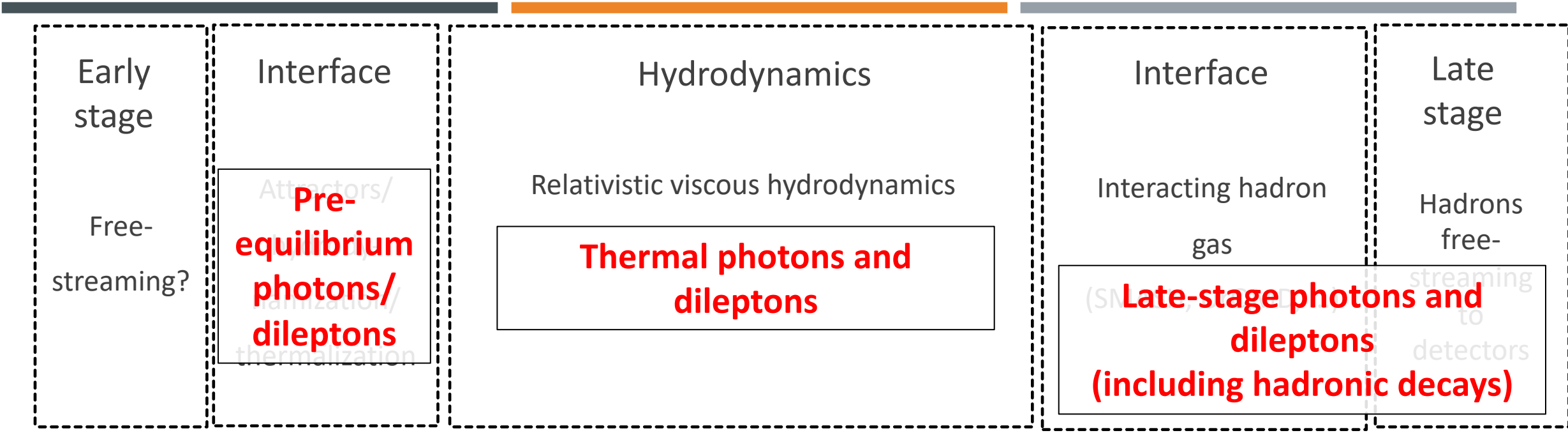
Electromagnetic probes in heavy-ion collisions



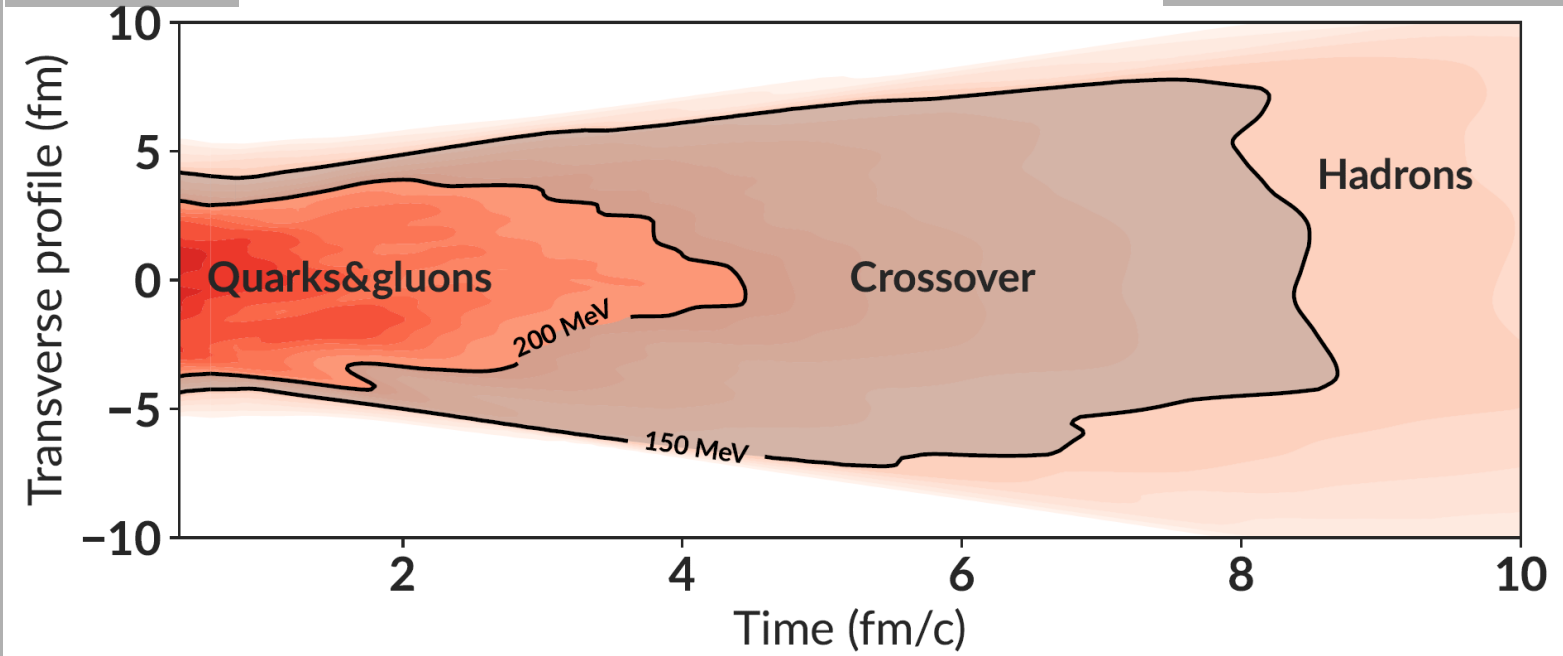
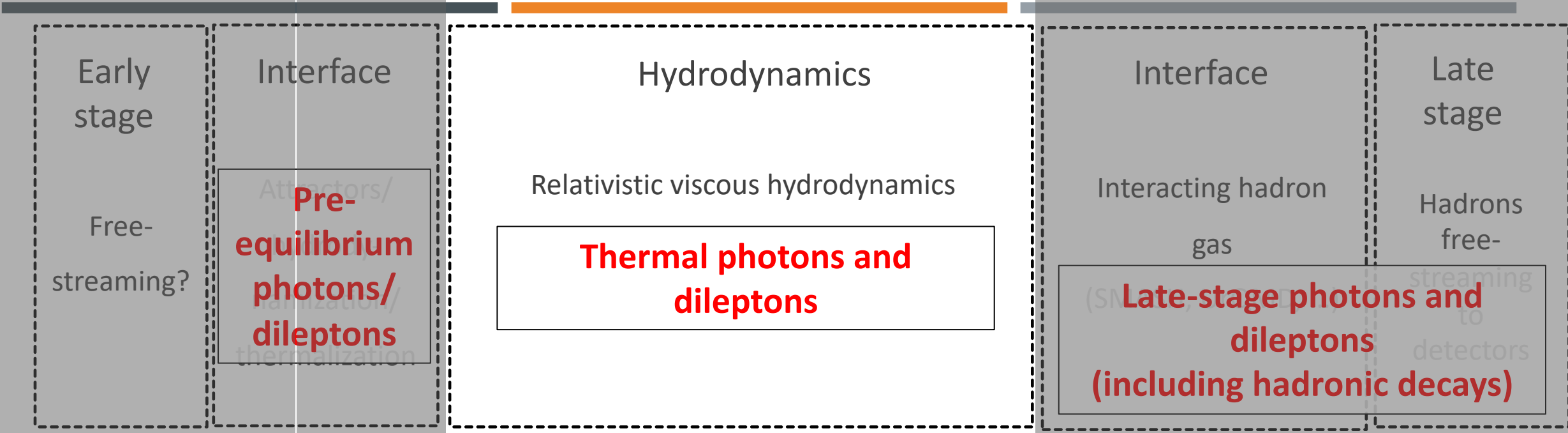








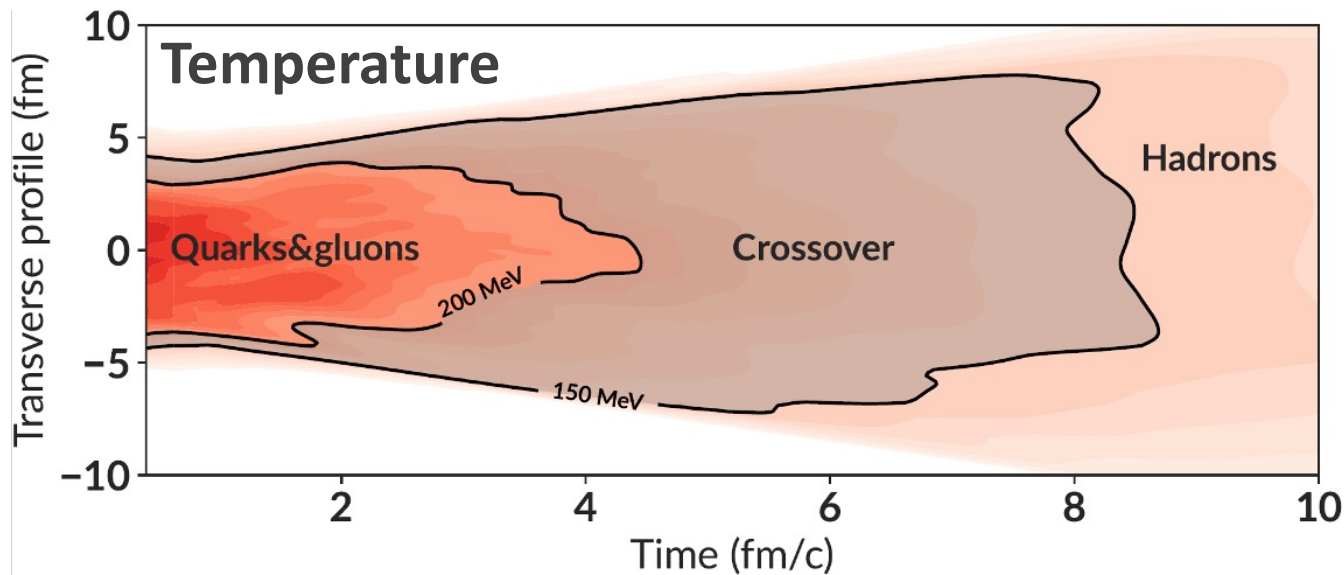
Other sources...



Other sources...

Photons from deconfined plasma

- What is the spacetime profile of quarks/gluons/hadrons?



Spacetime profile of plasma

- Photon production:
$$\frac{dN_\gamma}{d^3p} = \int d^4X \frac{d\Gamma_\gamma}{d^3p} (p, \overbrace{T(X), u^\mu(X), \dots})$$

Photon emission rate

Photon emission rate

Spacetime profile of plasma

▪ Photon production:
$$\frac{dN_\gamma}{d^3p} = \int d^4X \frac{d\Gamma_\gamma}{d^3p} (p, T(X), u^\mu(X), \dots)$$

Photon emission rate

State of matter/Temperatures

Gas of hadrons below $T \approx 160$ MeV

Deconfinement for $T \approx 160 - 200$ MeV

Strongly-coupled quark/gluons
for $T \sim 200 - 500$ MeV

Weakly-coupled QGP at $T \gg 1$ GeV

Photon emission rate

Effective hadronic models

Extrapolated rates from low/high
temperatures

Lattice QCD, holography, effective
models

Perturbative QCD



Photon emission rate

Spacetime profile of plasma

▪ Photon production:
$$\frac{dN_\gamma}{d^3p} = \int d^4X \frac{d\Gamma_\gamma}{d^3p} (p, T(X), u^\mu(X), \dots)$$

Photon emission rate

State of matter/Temperatures

**Photon emissivity of the quark-gluon plasma:
A lattice QCD analysis of the transverse
channel**

Marco Cè (U. Bern, AEC and Bern U. and CERN), Tim
Harris (Edinburgh U.), Ardit Krasniqi (U. Mainz, PRISMA), Harvey B.
Meyer (U. Mainz, PRISMA and Helmholtz Inst., Mainz and
Darmstadt, GSI), Csaba Török (U. Mainz, PRISMA)

May 5, 2022

26 pages

weakly coupled QCD at $T \gg T_c$

Photon emission rate

Effective hadronic models

Extrapolated rates from low/high
temperatures

Lattice QCD, holography, effective
models

Perturbative QCD

Photon emission rate

Spacetime profile of plasma

▪ Photon production: $\frac{dN_\gamma}{d^3p} = \int d^4X \frac{d\Gamma_\gamma}{d^3p}(p, T(X), u^\mu(X), \text{viscous part (shear and bulk)})$

Photon emission rate

State of matter/Temperatures

Equilibrium rate

Non-equilibrium corrections?

Gas of hadrons

Effective hadronic models

Partly known

Deconfinement

Extrapolated rates from low/high temperatures

Very limited information

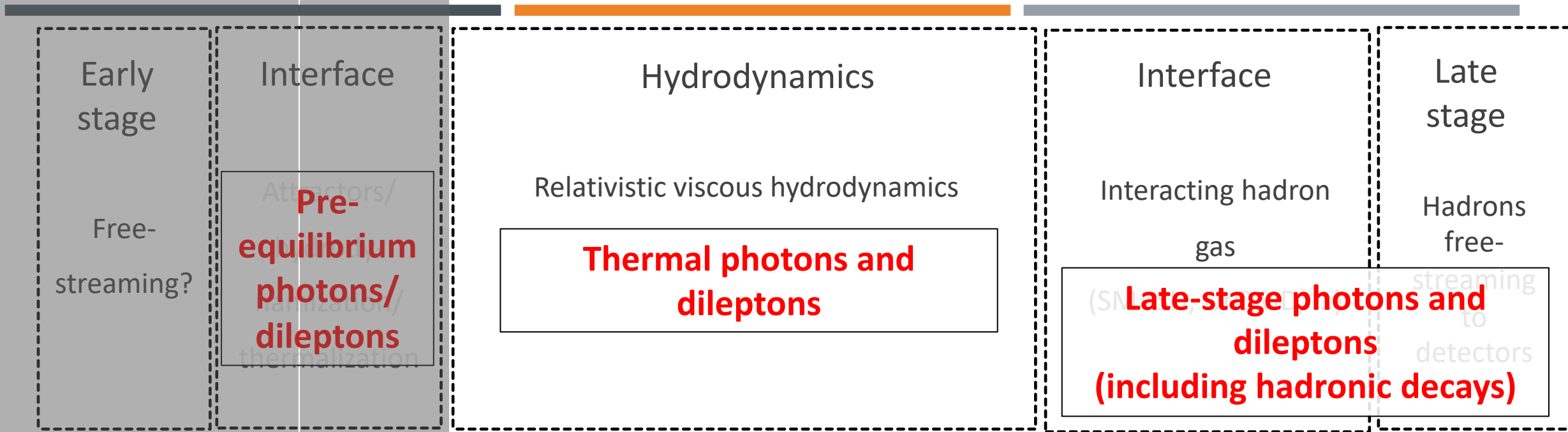
Strongly-coupled QGP

Lattice QCD, holography, effective models

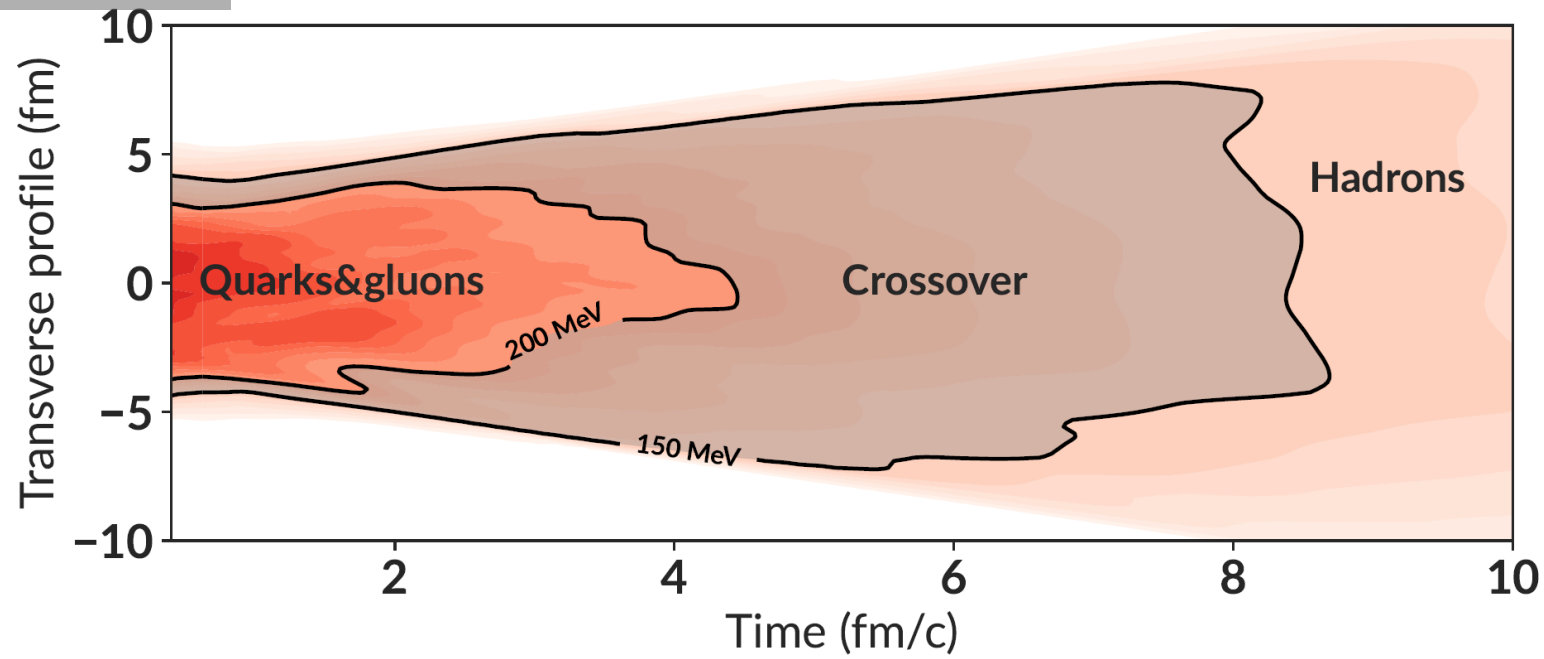
Weakly-coupled QGP

Perturbative QCD

Partly known



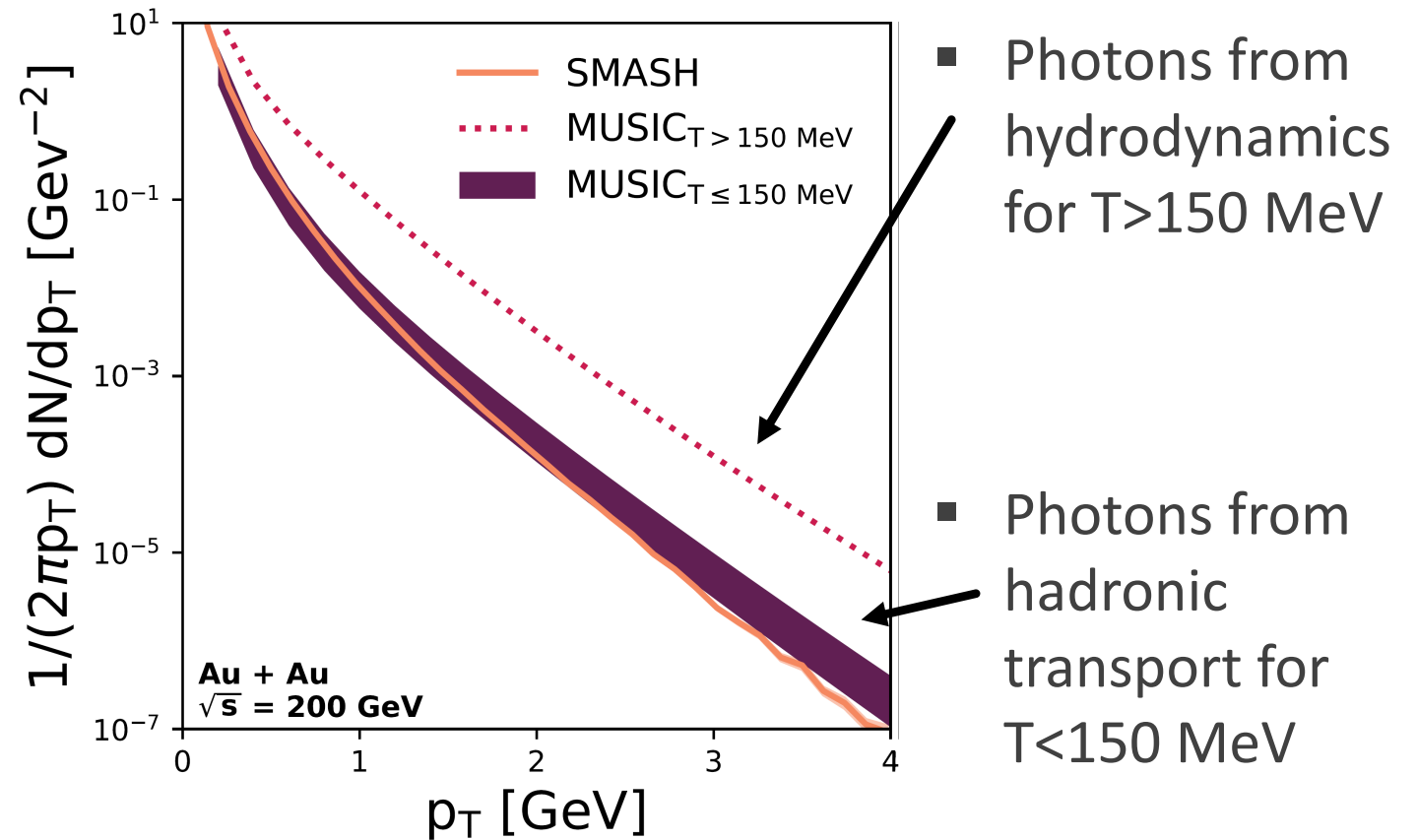
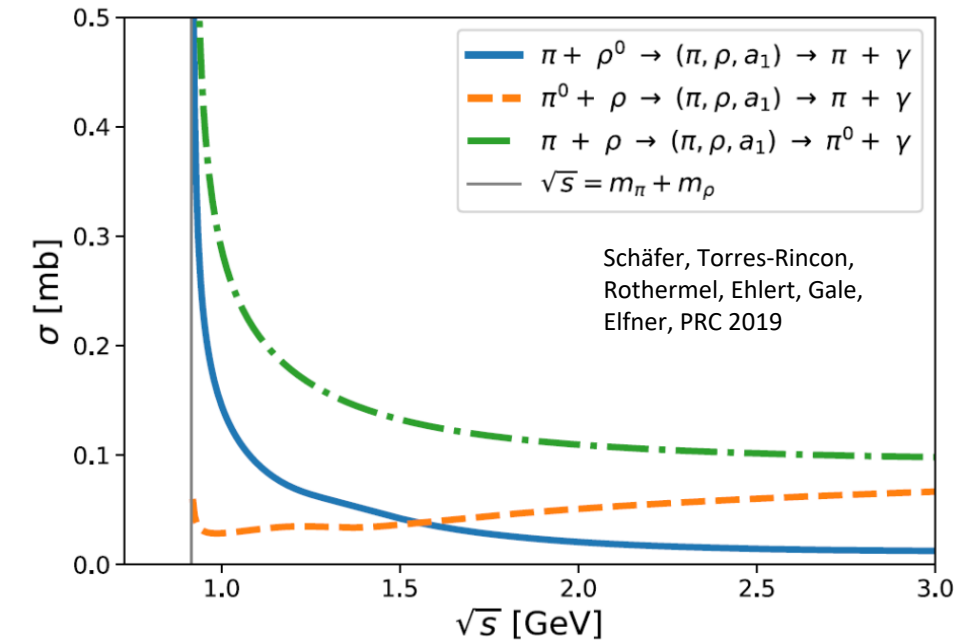
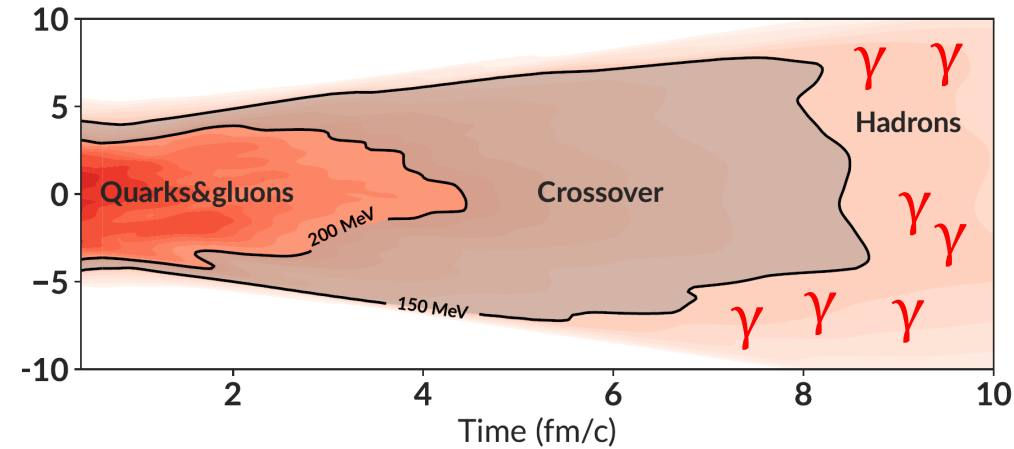
Prompt photons / Drell-Yan dileptons

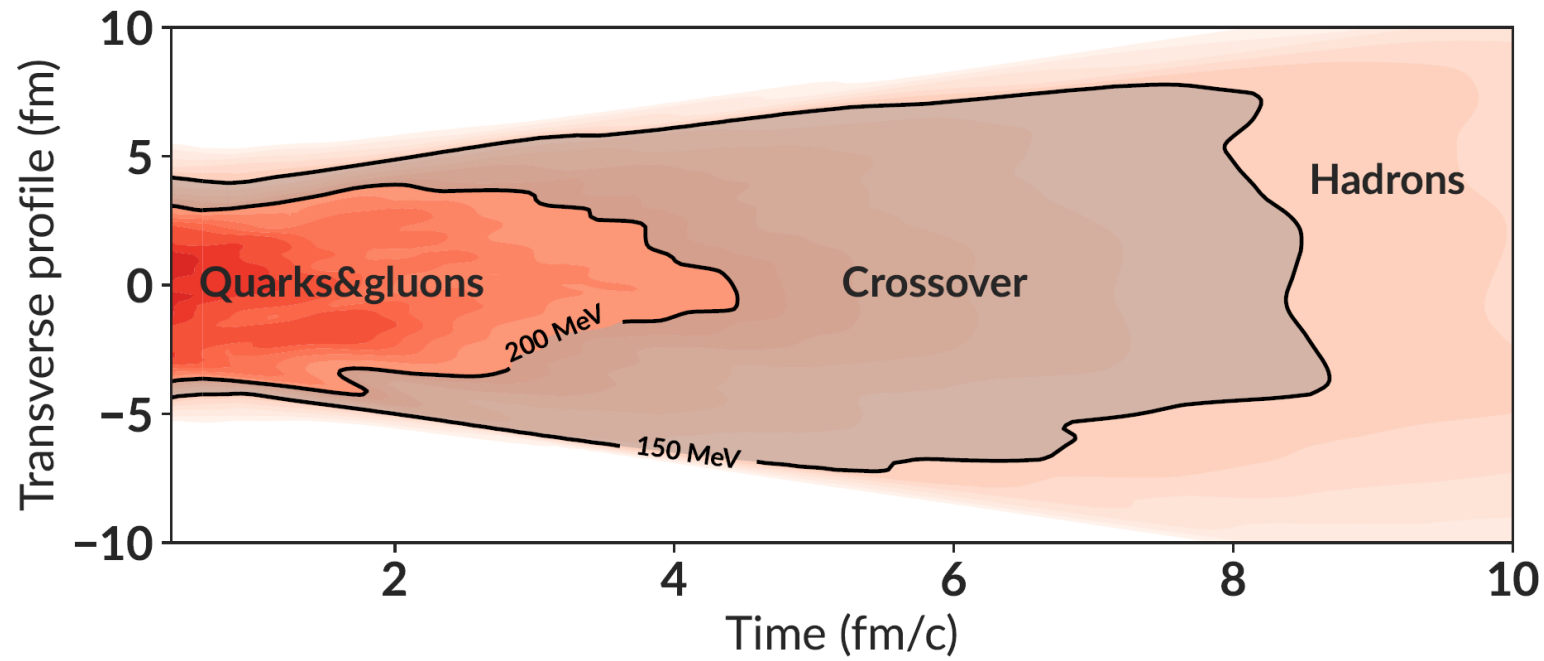
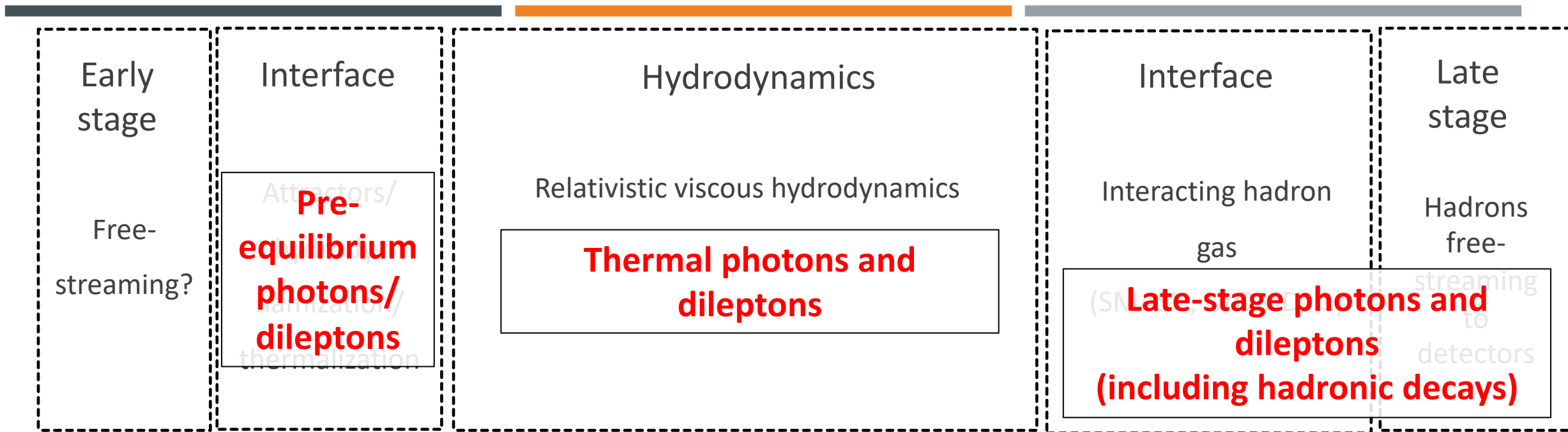


Other sources...

Photons from hadronic interactions

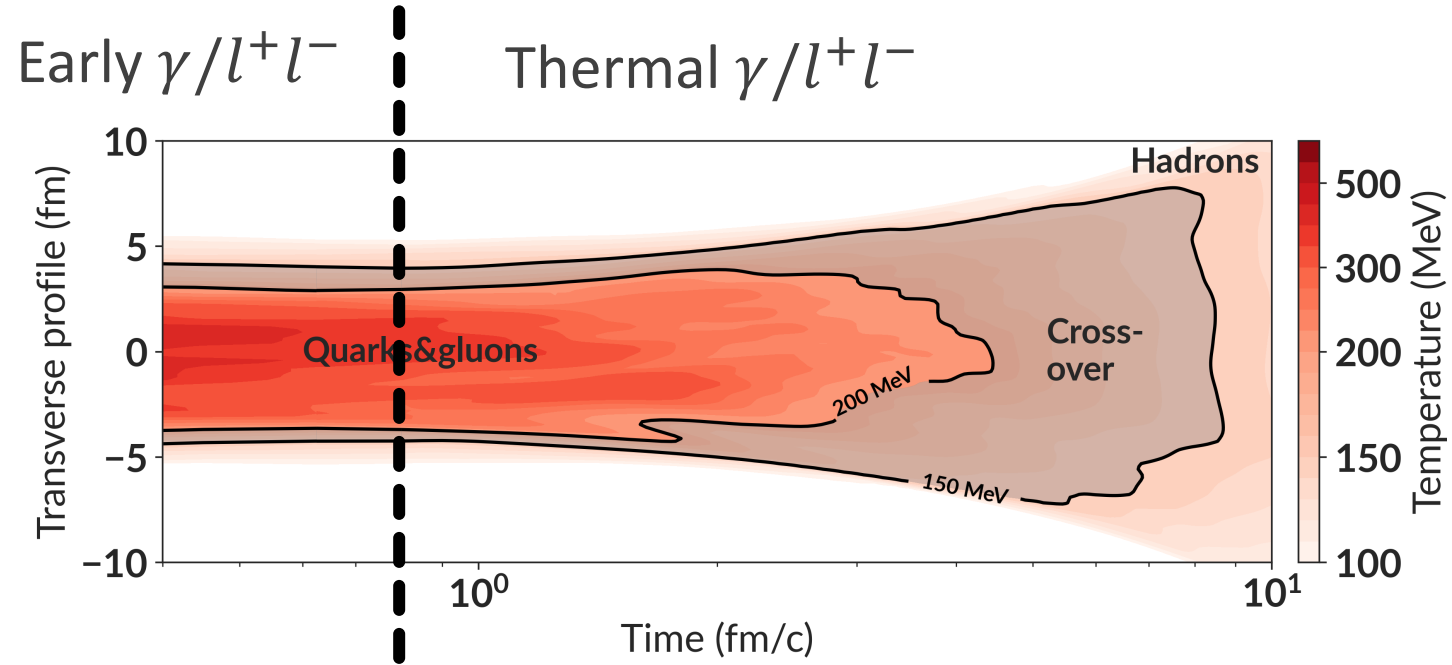
$$\frac{d^4 N_\gamma}{d^4 k} = \int d^4 X \frac{d^4 \Gamma_{\gamma/l+l^-}}{d^3 p} (p, T(X), u^\mu(X), \dots)$$





Other sources...

Early-stage emission

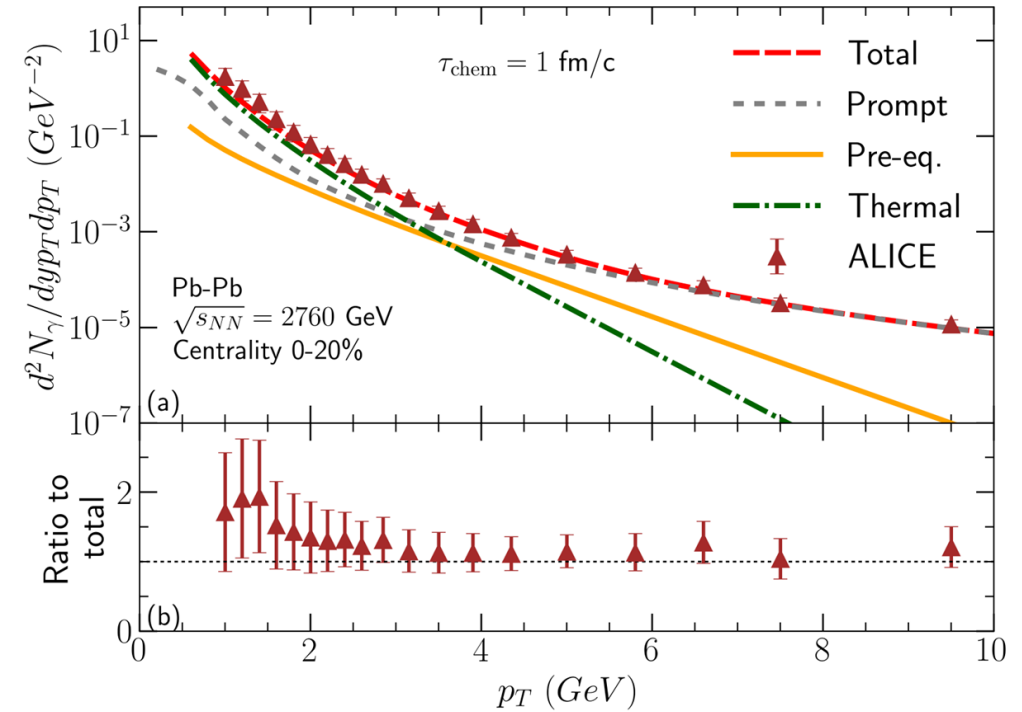
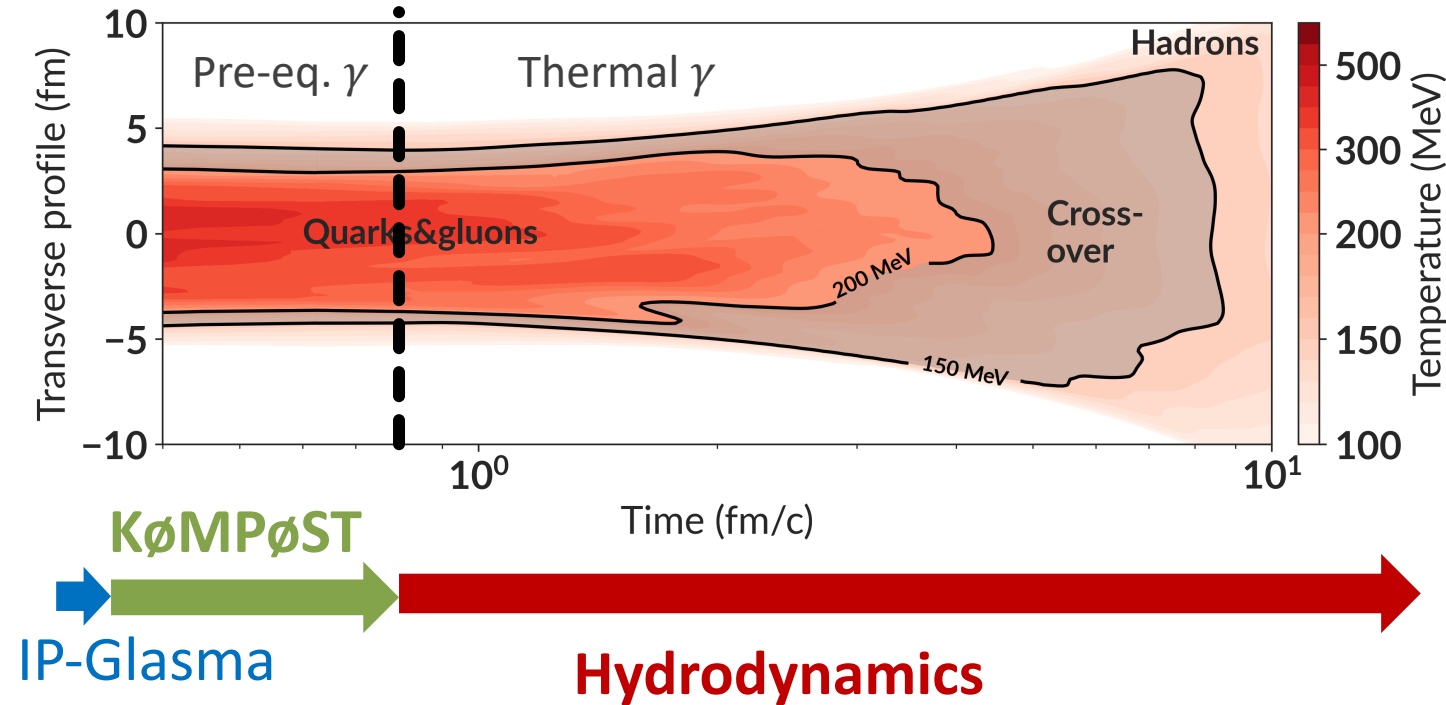


Emission from the soft bath

- What is the spatial distribution?
- What is the **rate** of photon and dilepton emission at early times?
 - **Rate** determined by quark/gluon ratio and momentum distributions
 - e.g. thermal distributions = equilibrium emission rate ($e^{-\text{energy}/T}$)

Early-stage photons

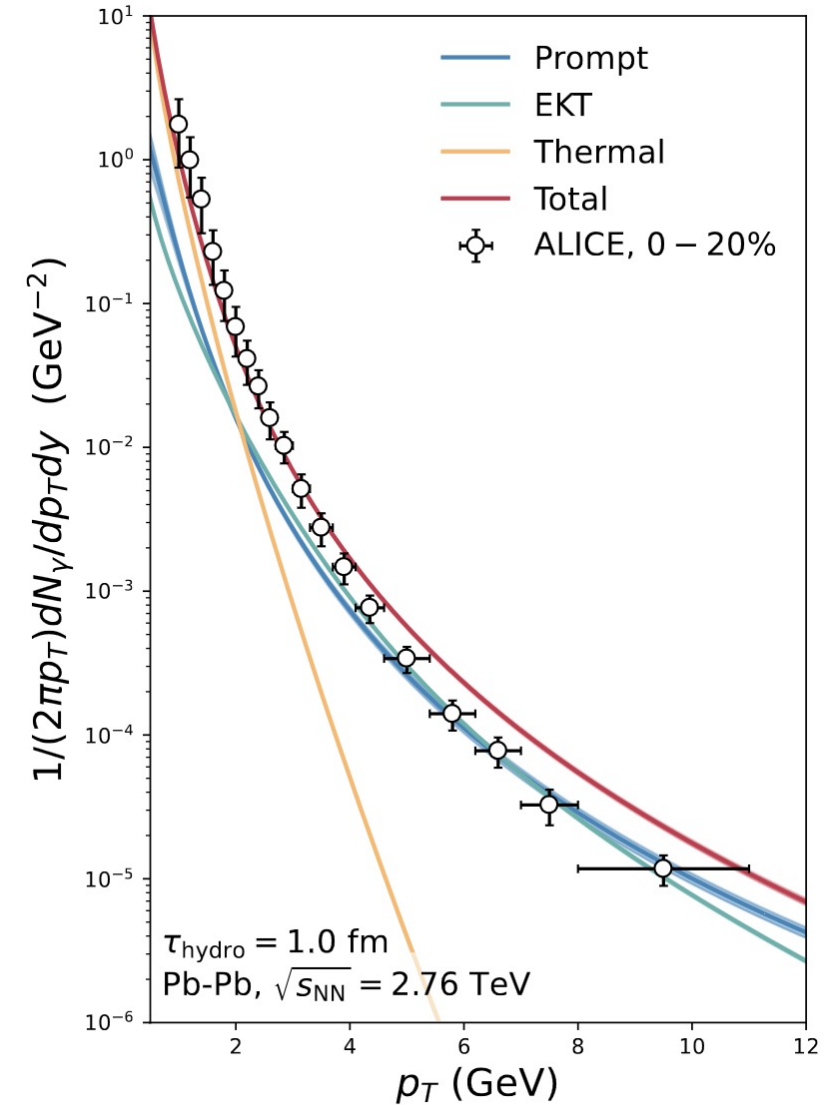
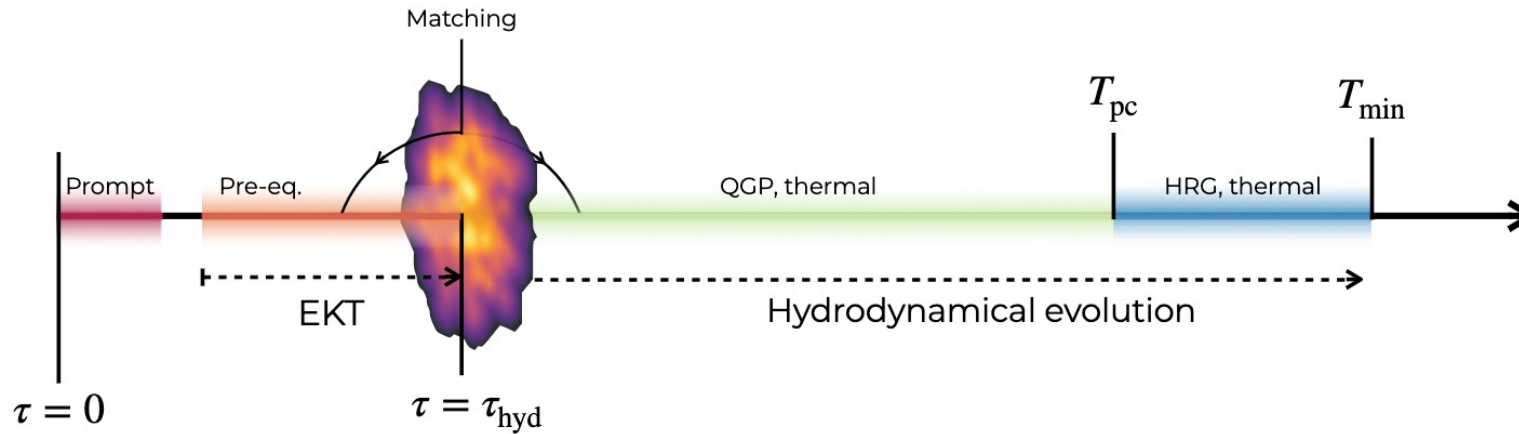
Gale, Paquet, Schenke,
Shen (2022) PRC

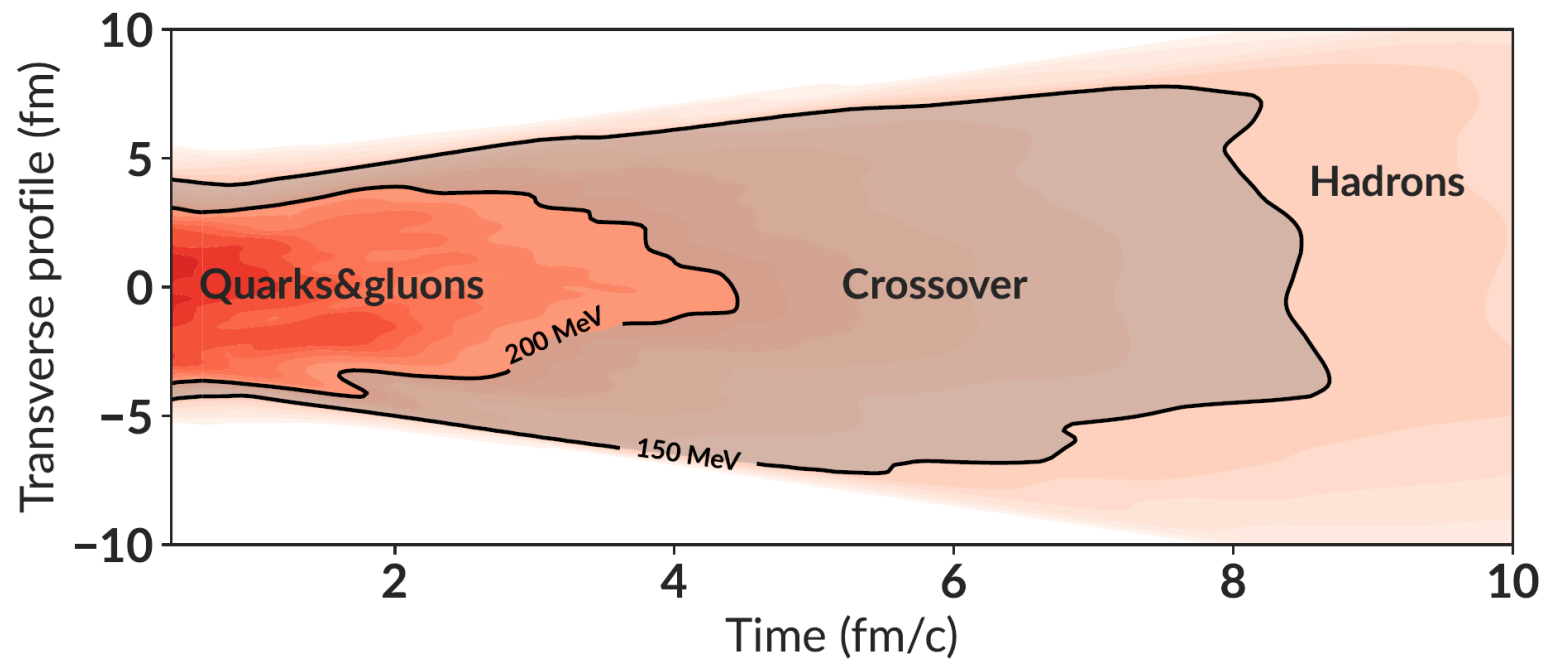
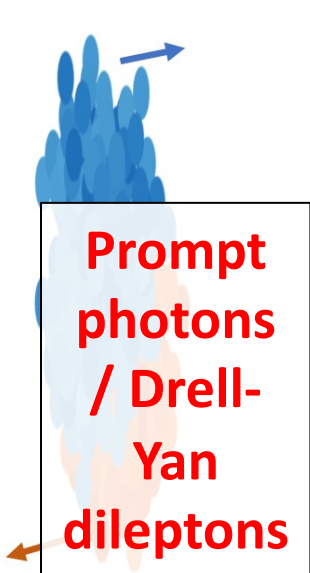
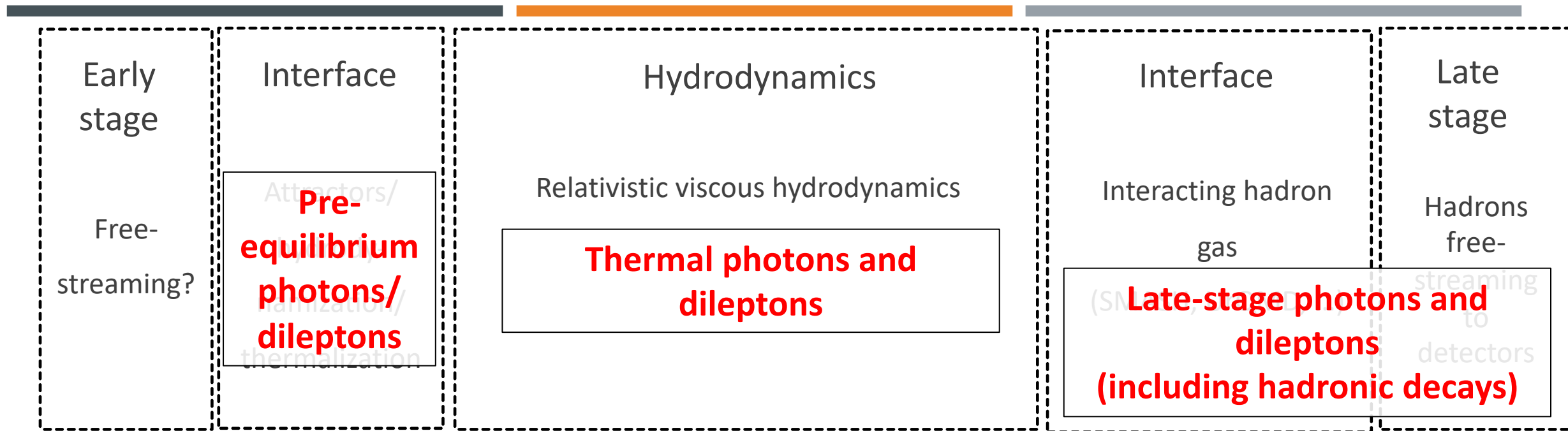


- What is the spatial distribution? IP-Glasma+KøMPøST
- What is the rate of photon emission at early time?
 - Thermal rate w/ viscous corrections + rate suppression factor for chemistry

Early-stage photons

Garcia-Montero, Mazeliauskas, Plaschke, Schlichting
[arXiv:2308.09747]

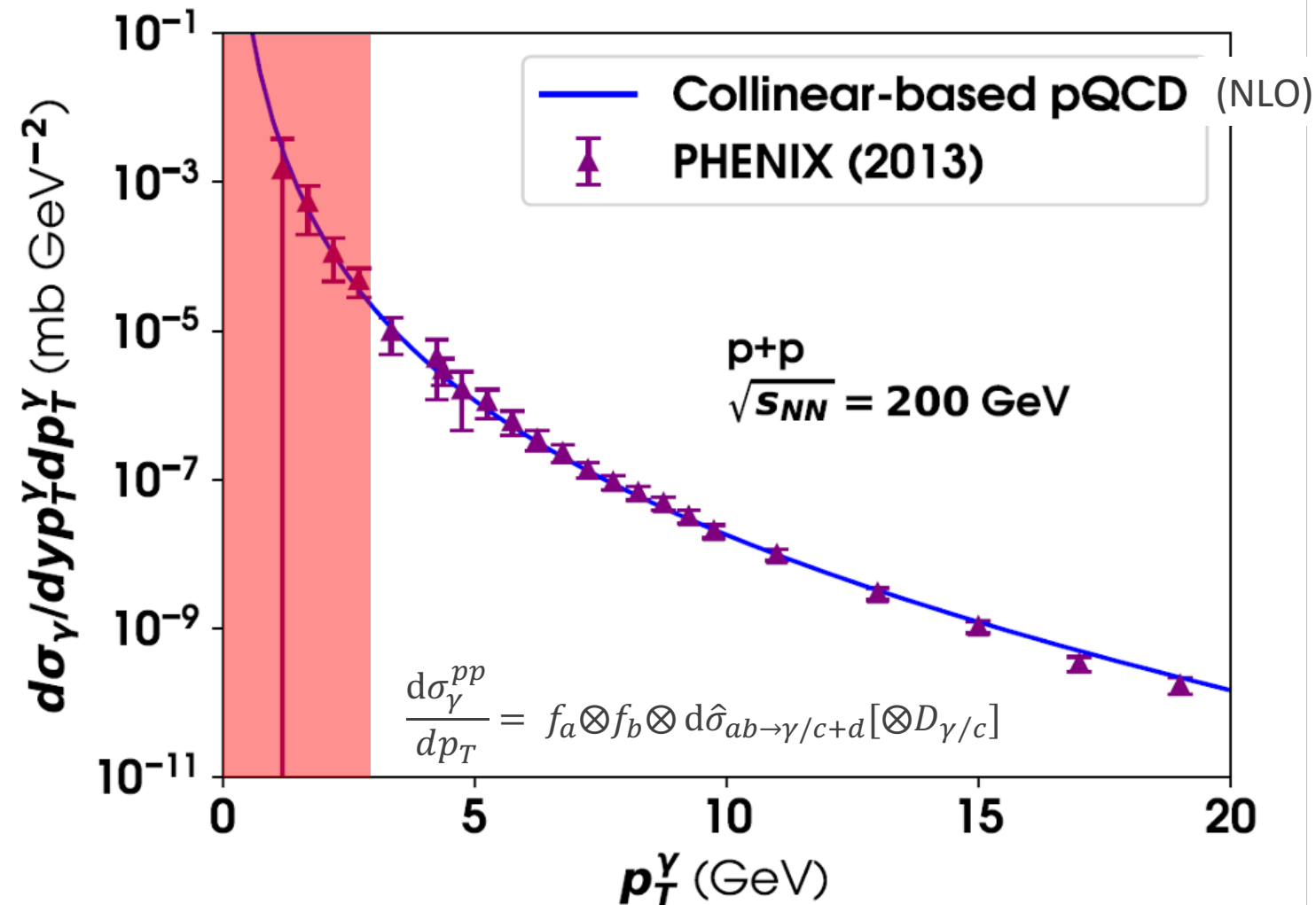




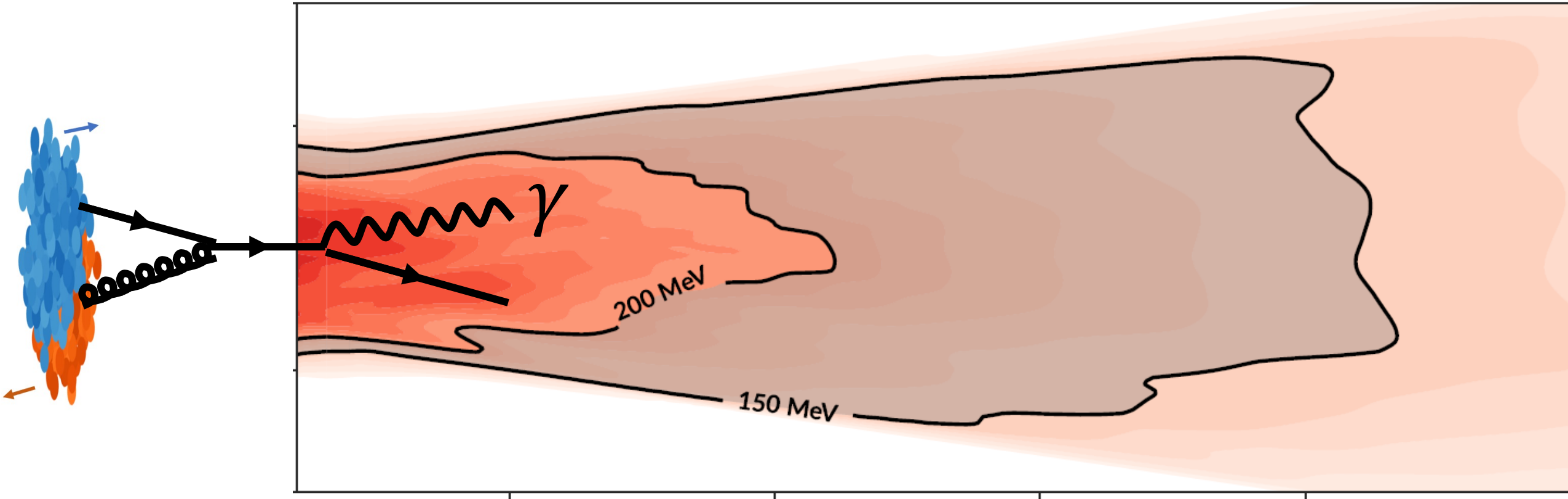
Other sources...

Direct photons in proton-proton collisions: “low” energy

- Low p_T photons:
 - Few measurements (in proton-proton collisions)
 - Difficult to compute from first principles
 - Non-perturbative effects likely significant



Medium-modified prompt photons



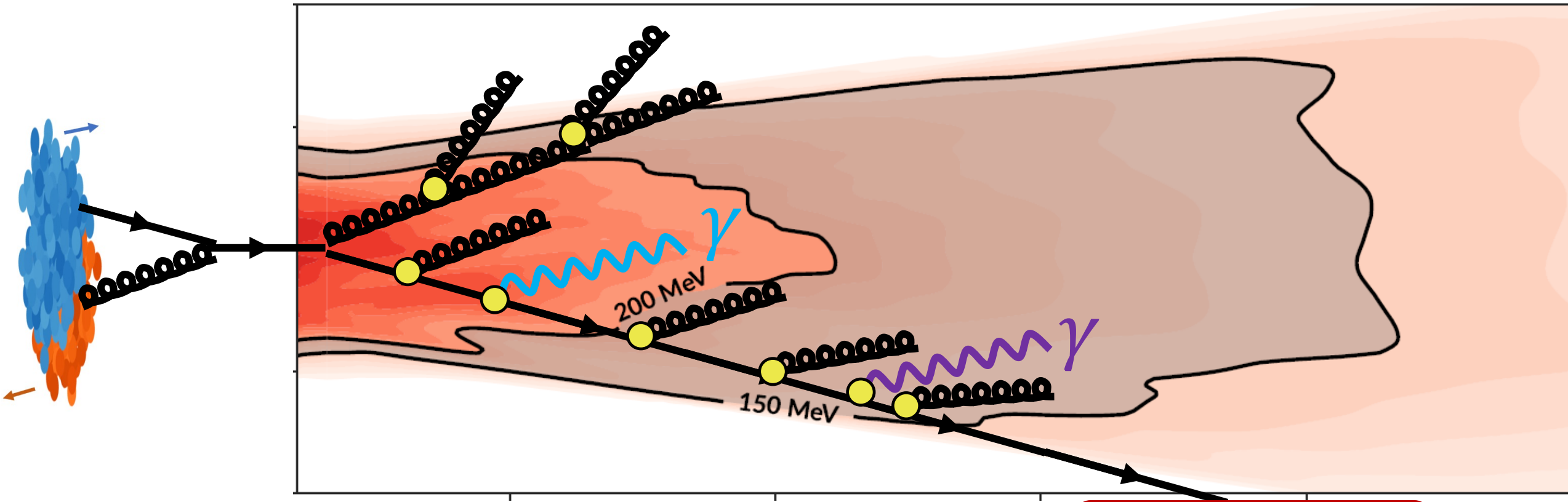
No medium effects on Compton scattering
and $q \bar{q}$ annihilation

$$q + \bar{q} \rightarrow g + \gamma$$

$$q + g \rightarrow q + \gamma$$

$$q + g \rightarrow q + g + \gamma ?$$

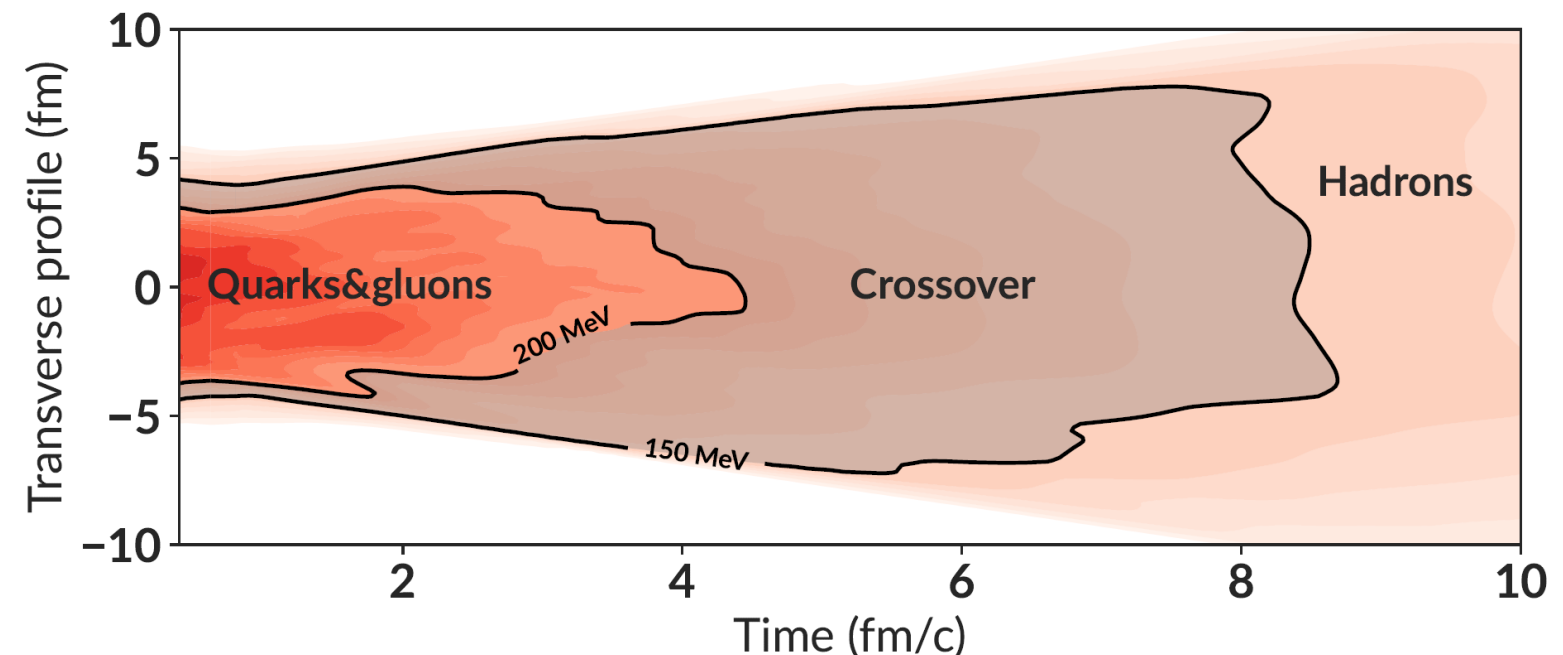
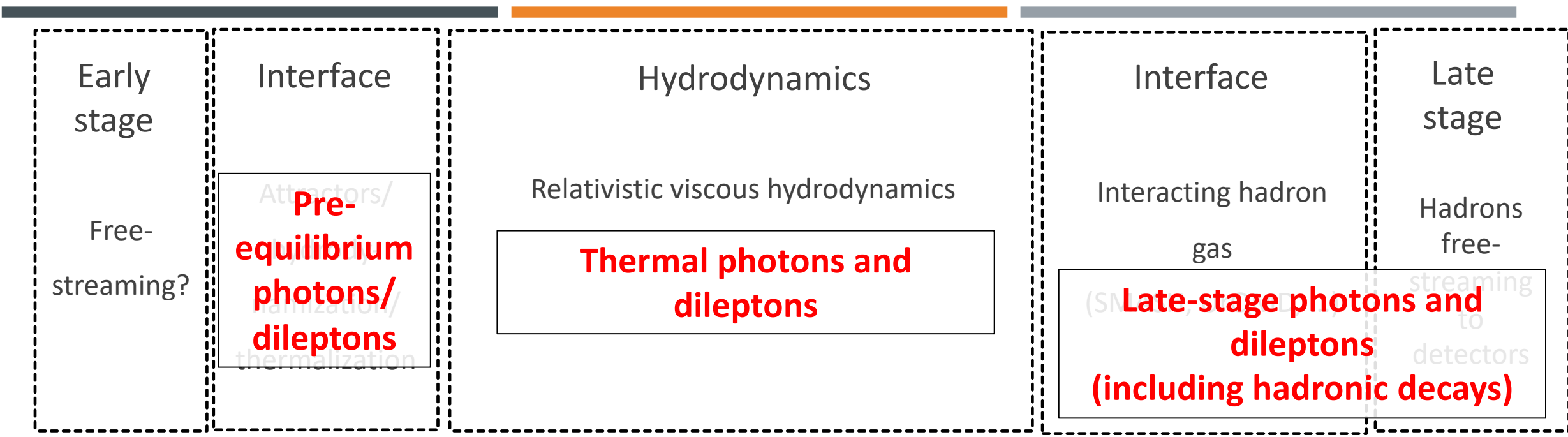
Medium-modified prompt photons



Medium-modified DGLAP-like radiation
 +medium-induced photons (“jet-medium”) [also l^+l^-]
 +non-perturbative fragmentation

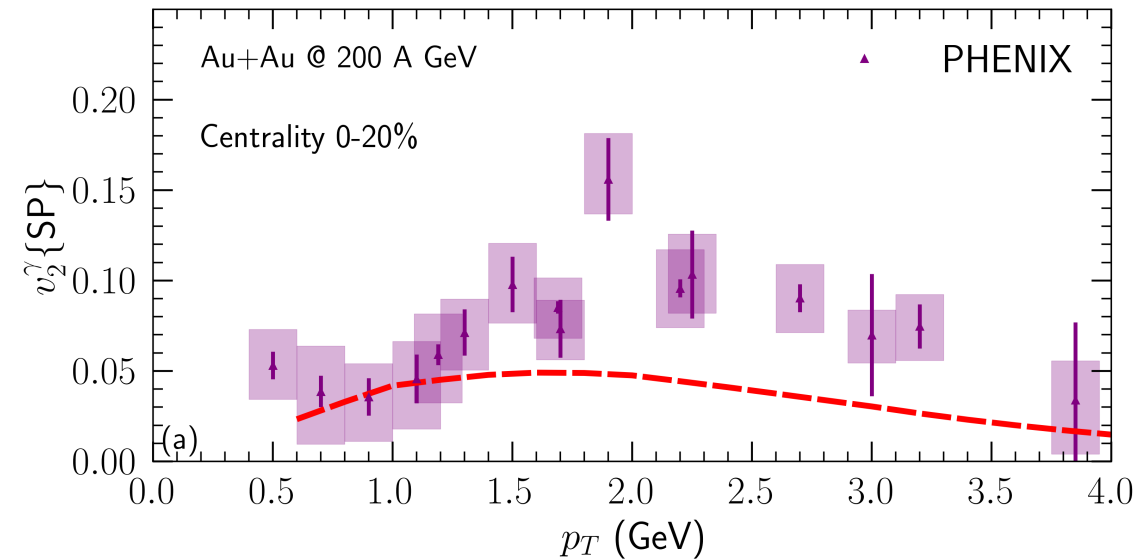
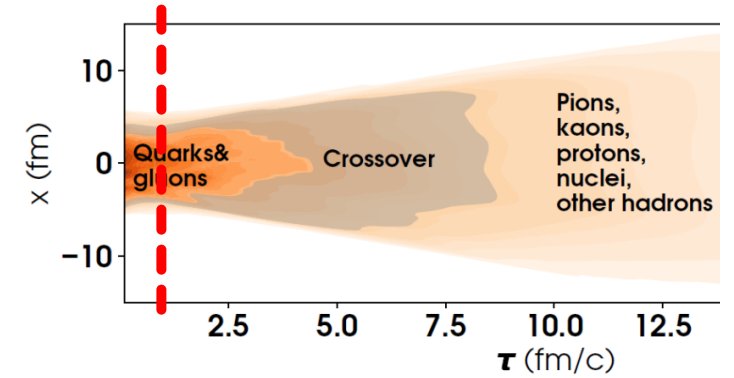
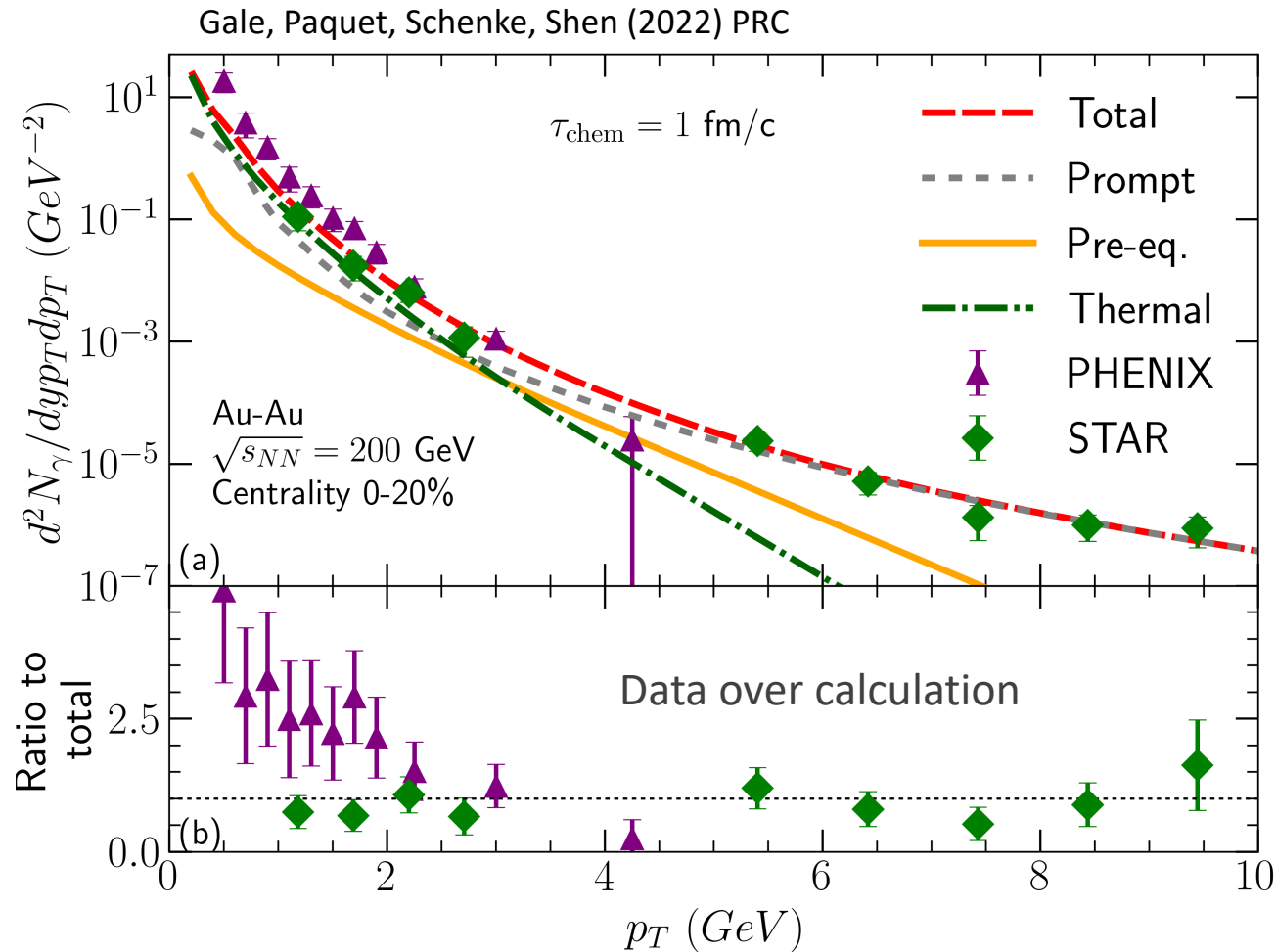
Fragmentation

Recent work: Shi, Modarresi Yazdi, Gale, Jeon (2022)

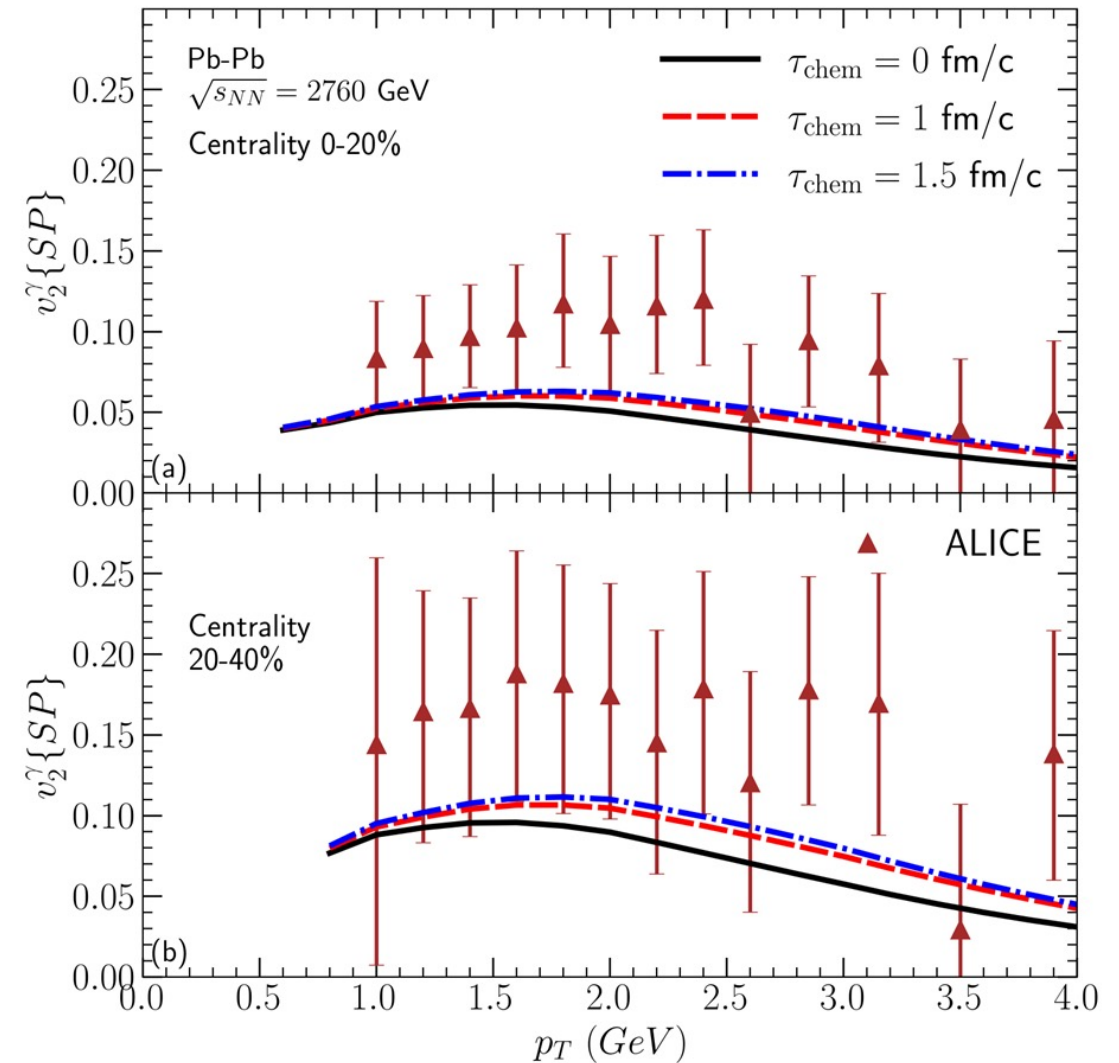
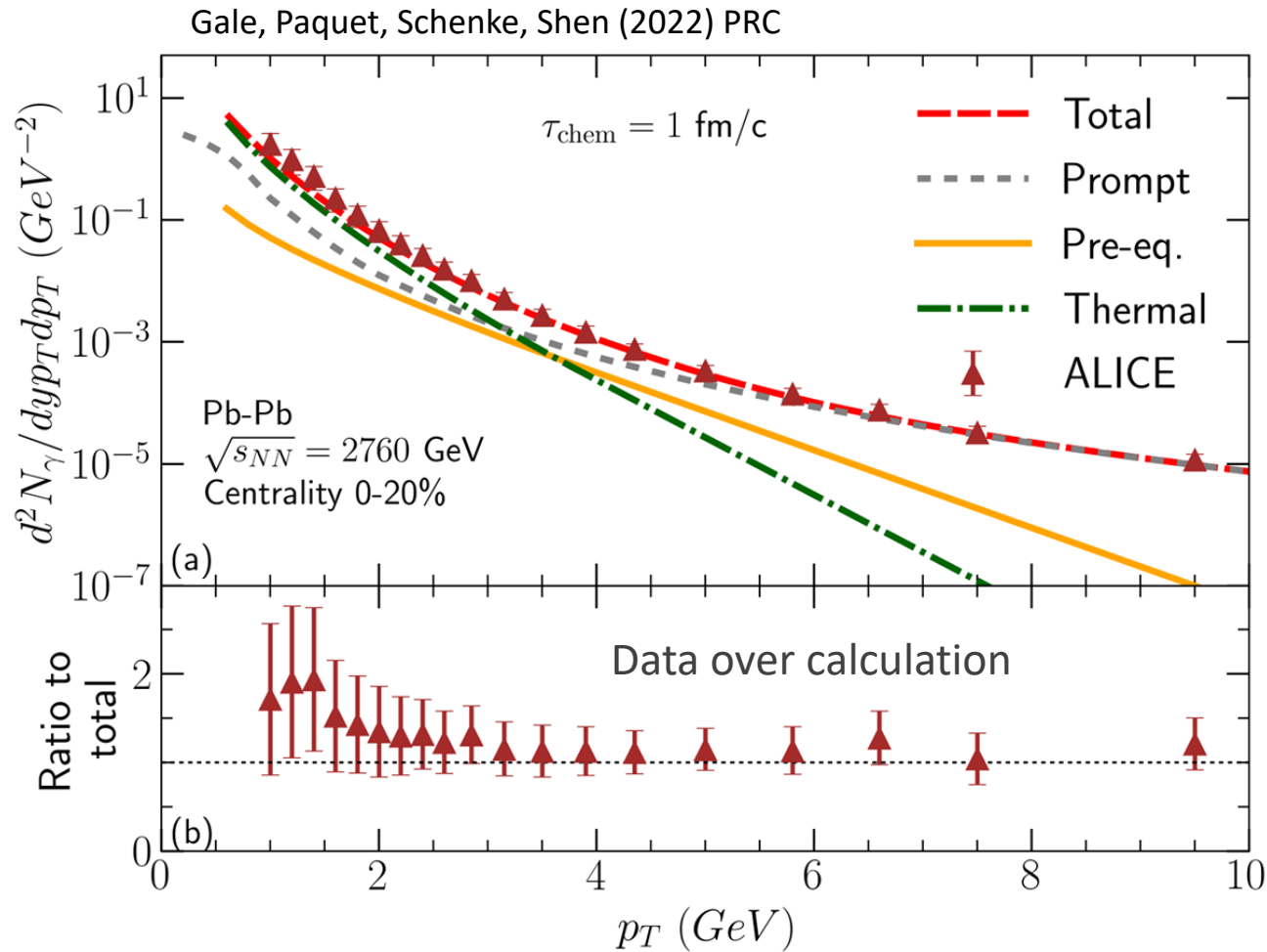


Other sources...

(Non-decay) photons in Au-Au $\sqrt{s_{NN}} = 200$ GeV, 0-20%



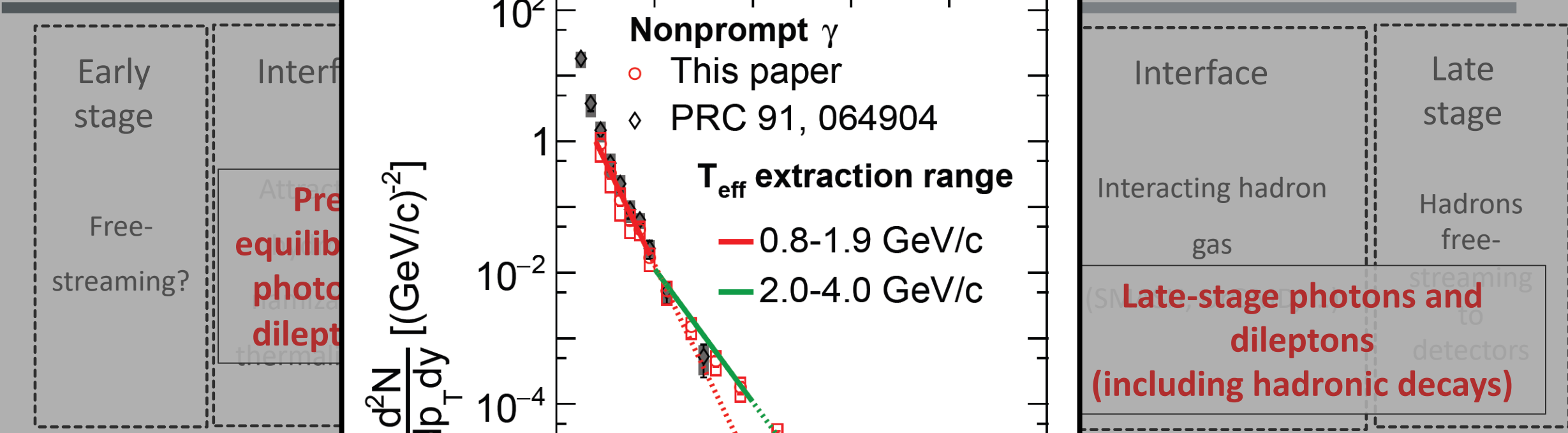
(Non-decay) photons in Pb-Pb $\sqrt{s_{NN}} = 2760$ GeV, 0-20%





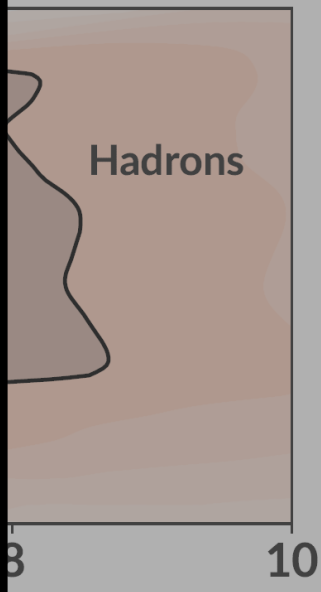
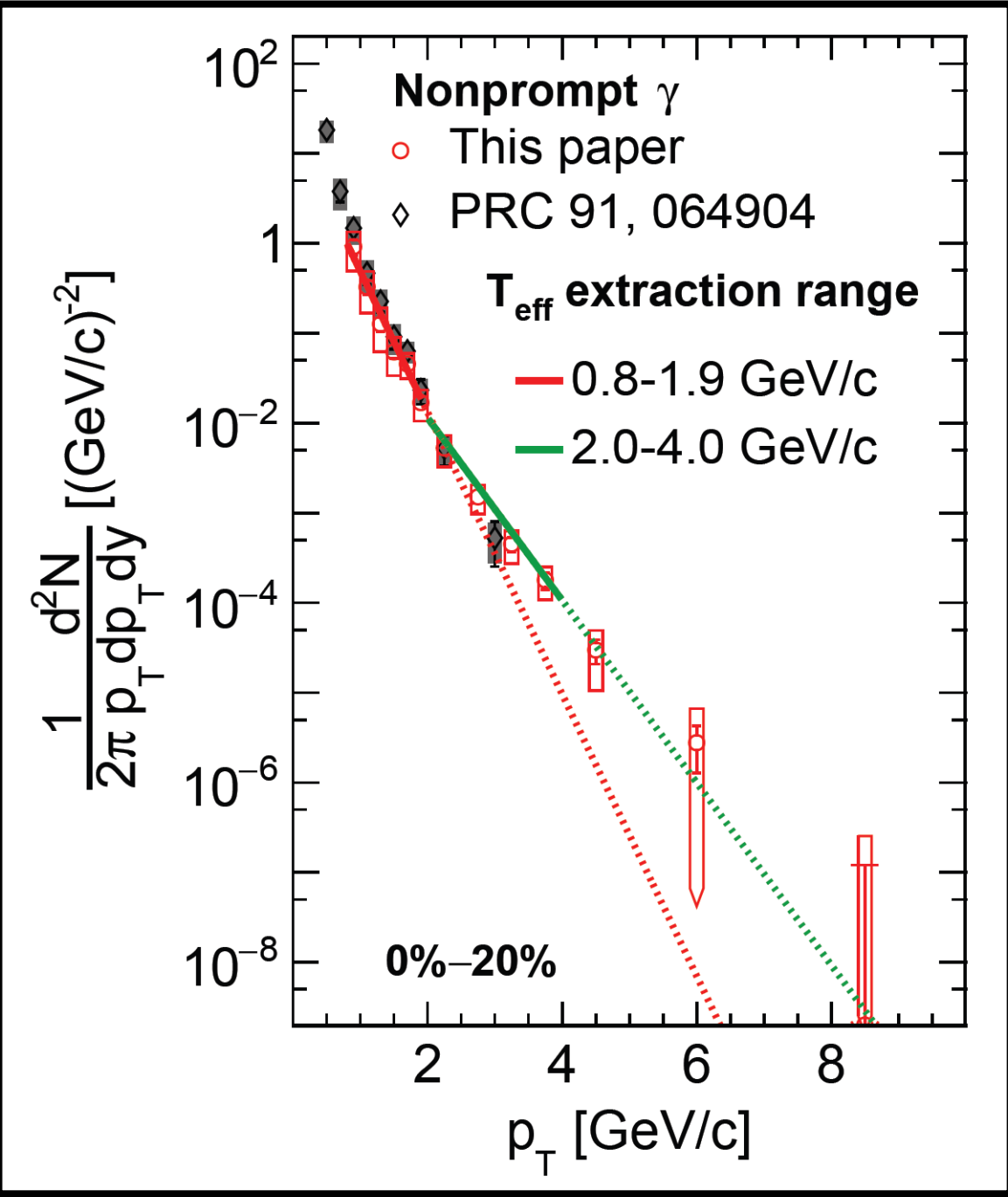
THE TEMPERATURE AND THE SLOPE





~~Prompt photons / Dileptons~~

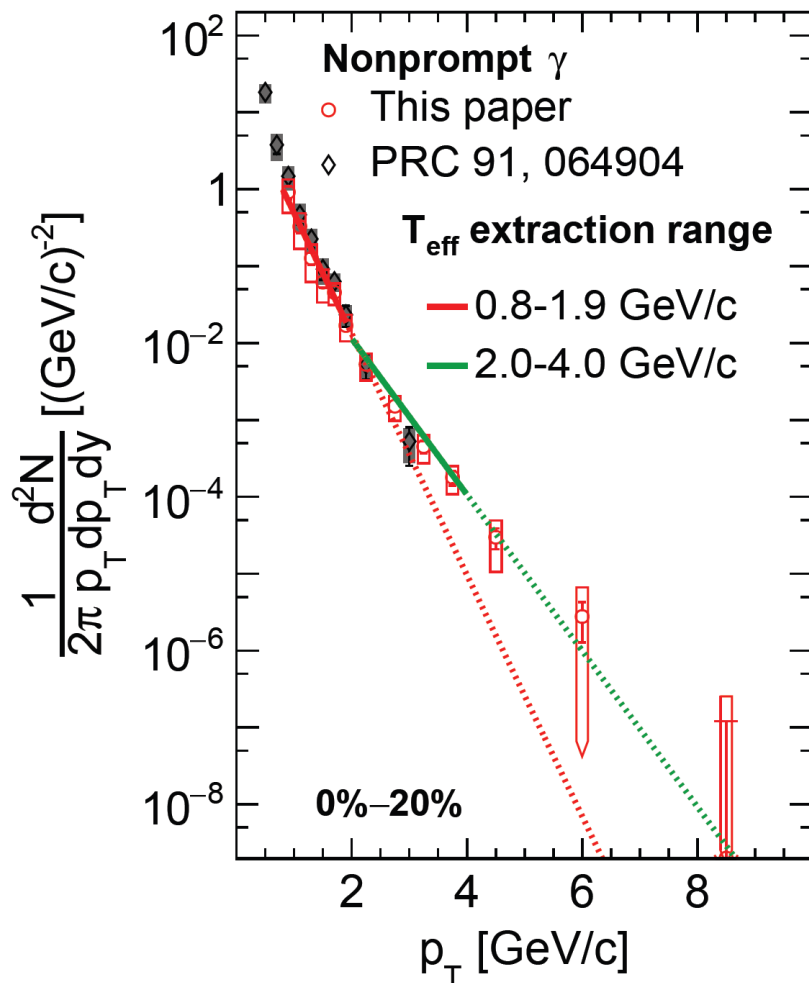
Transverse profile (fm)



Other sources...

Photon p_T spectrum and inverse slope

Ref.: PHENIX Collaboration [arXiv:2203.17187]



$$\frac{1}{2\pi p_T} \frac{dN_\gamma}{dp_T} \Big|_{y=0, p_{T,min} < p_T < p_{T,max}} \propto \exp\left(-\frac{p_T}{T_{eff}}\right)$$

centrality	T_{eff} (GeV/c)		T_{eff} (GeV/c)	
	0.8 < p_T < 1.9 GeV/c		2 < p_T < 4	
0%–20%	0.277 ± 0.017	+0.036 −0.014	0.428 ± 0.031	+0.031 −0.030
20%–40%	0.264 ± 0.010	+0.014 −0.007	0.354 ± 0.019	+0.020 −0.030
40%–60%	0.247 ± 0.007	+0.005 −0.004	0.392 ± 0.023	+0.022 −0.022
60%–93%	0.253 ± 0.011	+0.012 −0.006	0.331 ± 0.036	+0.031 −0.041

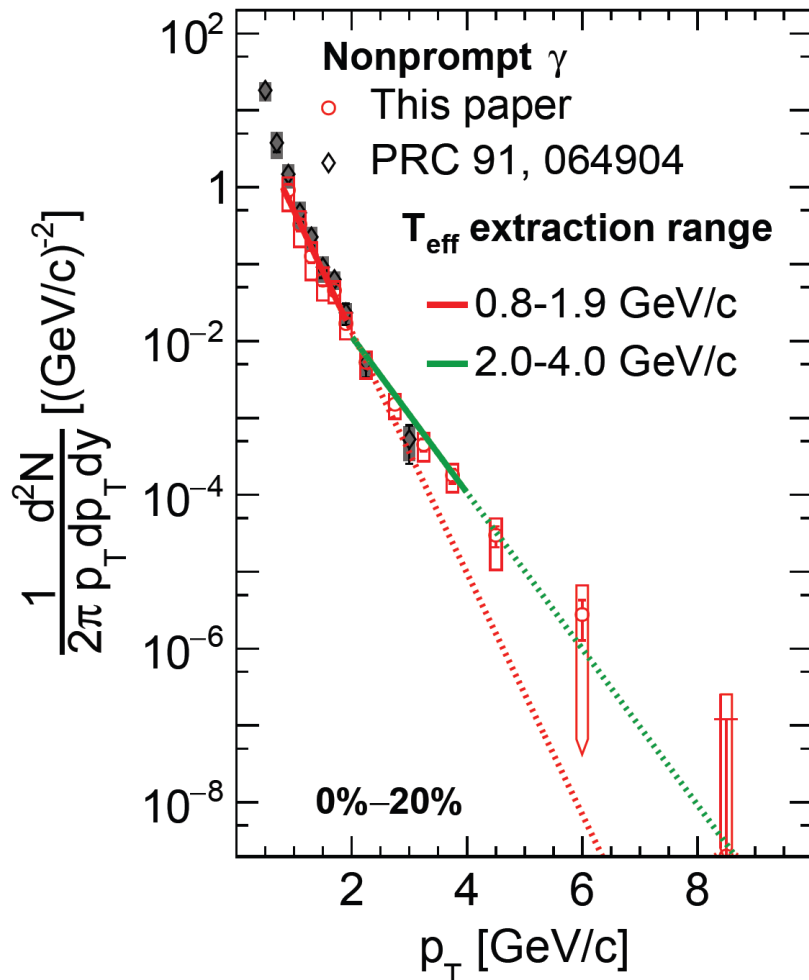
Results at the LHC by the ALICE Collaboration:

Centrality	T_{eff} (GeV) 0.9 < p_T < 2.1 GeV	T_{eff} (GeV) 1.1 < p_T < 2.1 GeV
0-20%	0.297	-
20-40%	-	0.410

Au-Au $\sqrt{s_{NN}} = 200$ GeV, 0-20%

Caveats: other sources of photons (e.g. pre-equilibrium), viscosity, ...

Ref.: PHENIX Collaboration [arXiv:2203.17187]



p_T cut	T_{eff}	$T_0 = \frac{T_{eff}}{1 - \frac{5 T_{eff}}{2 p_T}}$
$0.8 < p_T < 1.9$ GeV	277 MeV	570 MeV
$2 < p_T < 4$ GeV	428 MeV	670 MeV

From hydro fit to hadronic data: $T_0 \approx 530$ MeV

[from Gale, Paquet, Schenke, Shen (2022) PRC]

Partly explains large p_T -cut dependence of T_{eff}

Paquet and Bass [arXiv:2205.12299]; Paquet (2023) PRC

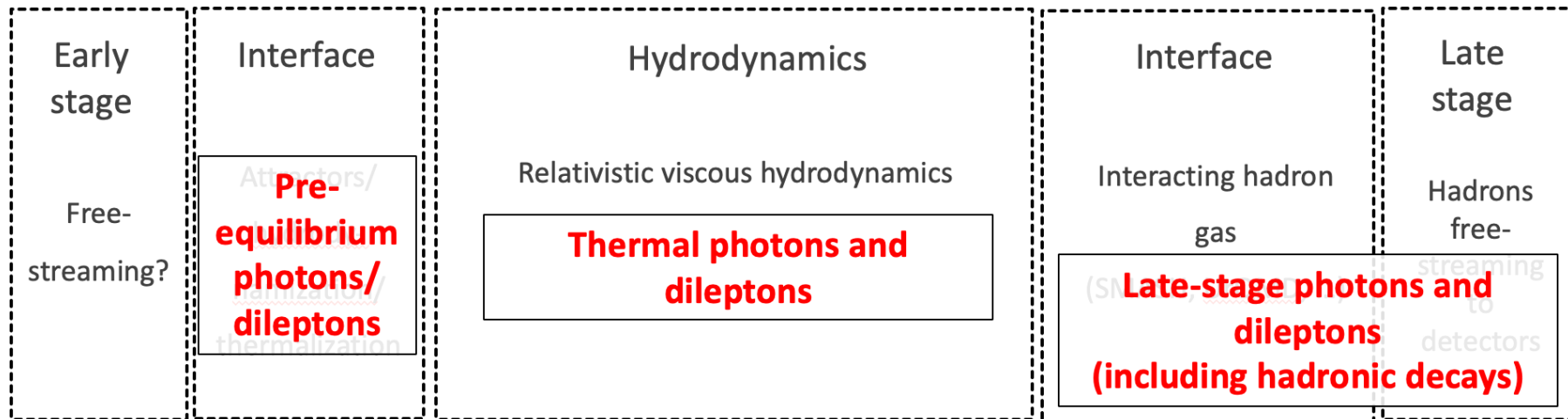


SUMMARY



Summary

- Electromagnetic probes are sensitive to all stages of heavy-ion collisions
- Progress at every stage, especially the interfaces; still more to do!



- Lower collisions energy or smaller system size = more interface



QUESTIONS?

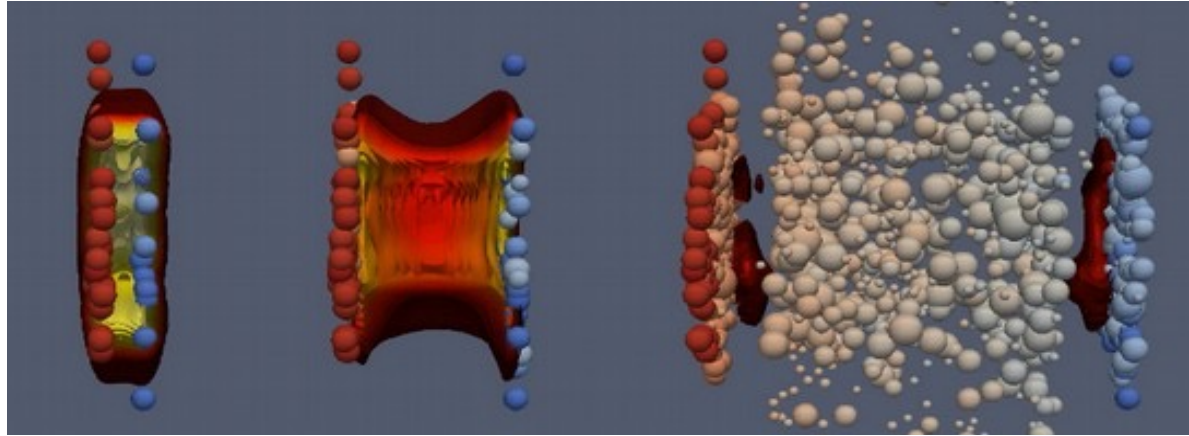
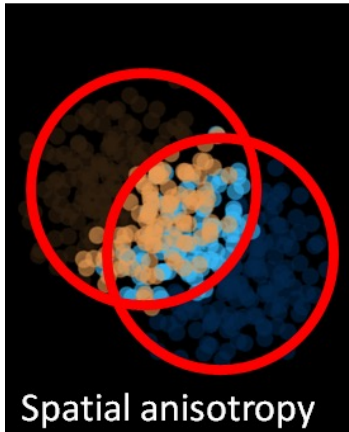




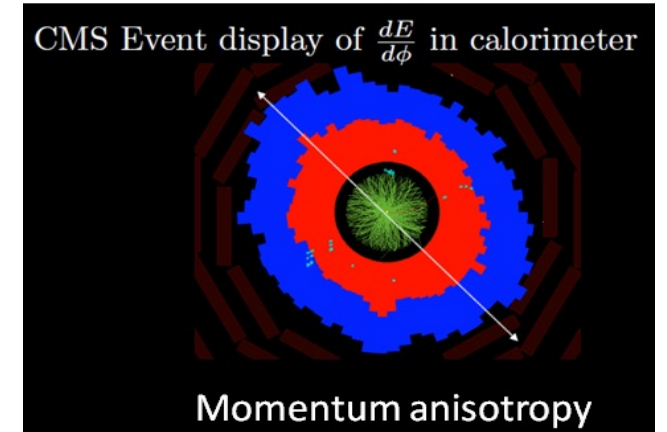
BACKUP



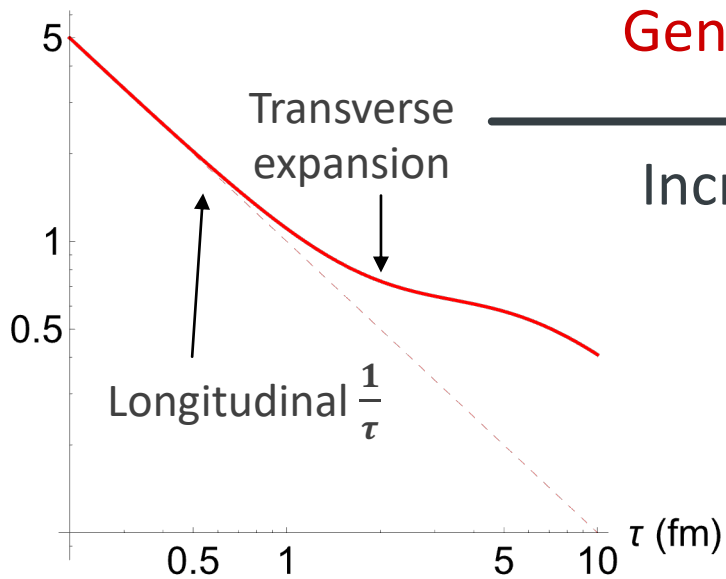
Interaction and expansion



Based on figures by Derek Teaney, CMS Coll., MADAI, H. Elfner and J. Bernhard



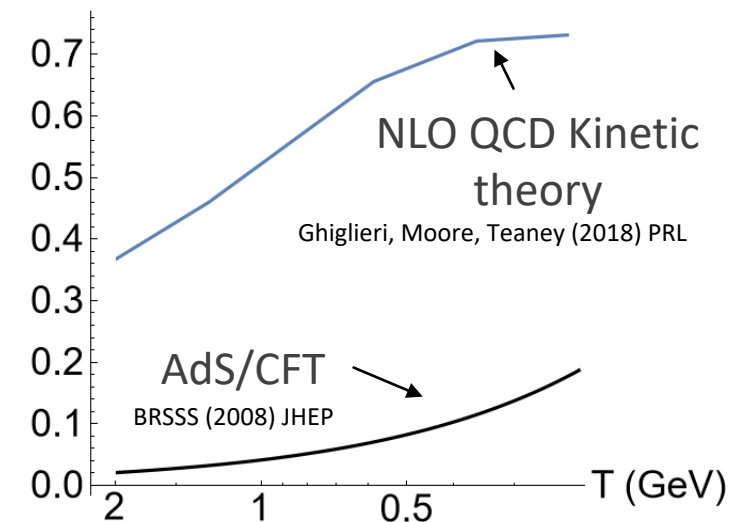
$$\theta = \partial_v u^v \text{ (1/fm)}$$



General decrease in expansion rate

Increase in local equilibration time (relaxation time)

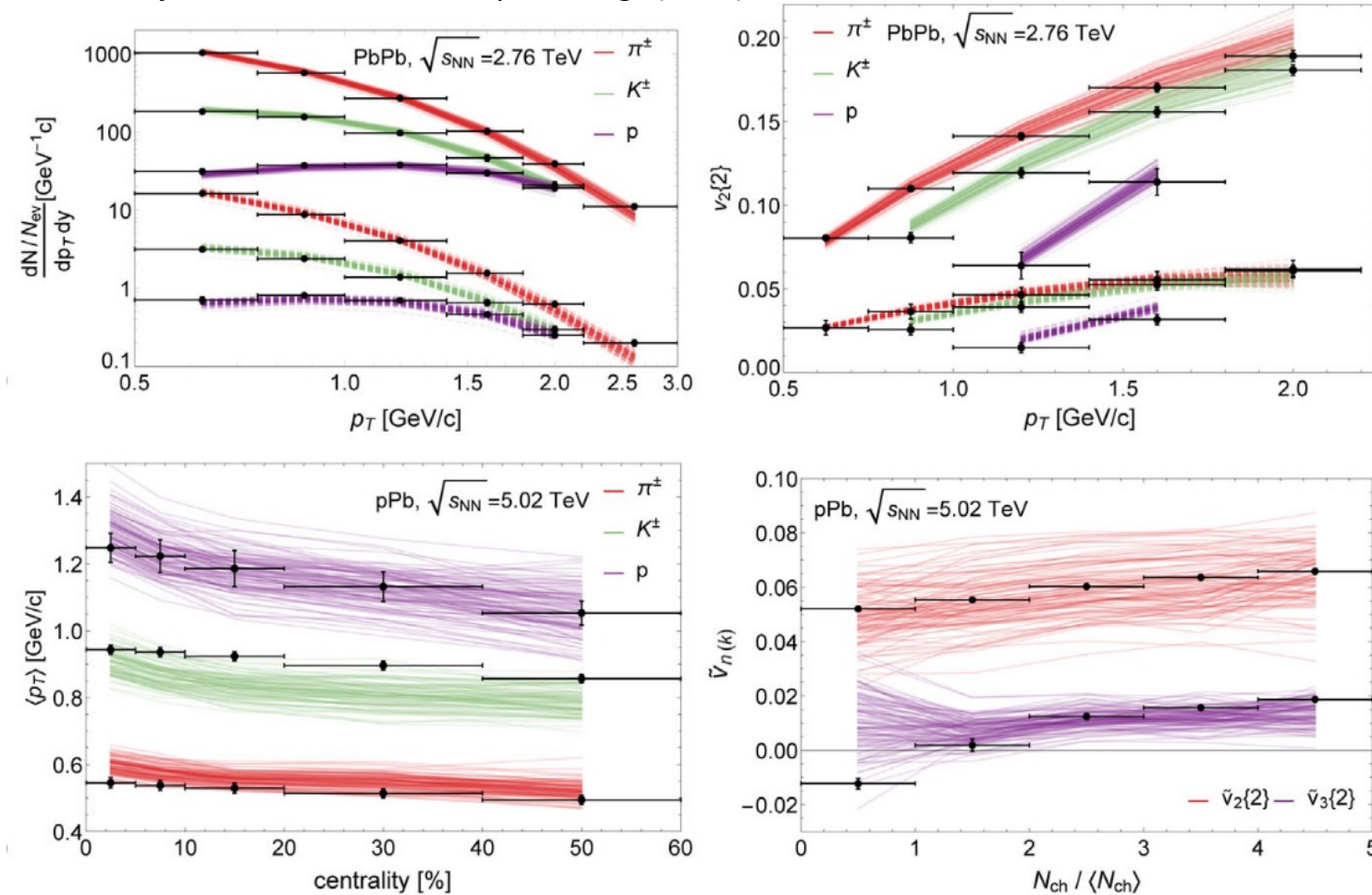
$$\tau_R \sim (\eta/s)/T \text{ (fm)}$$



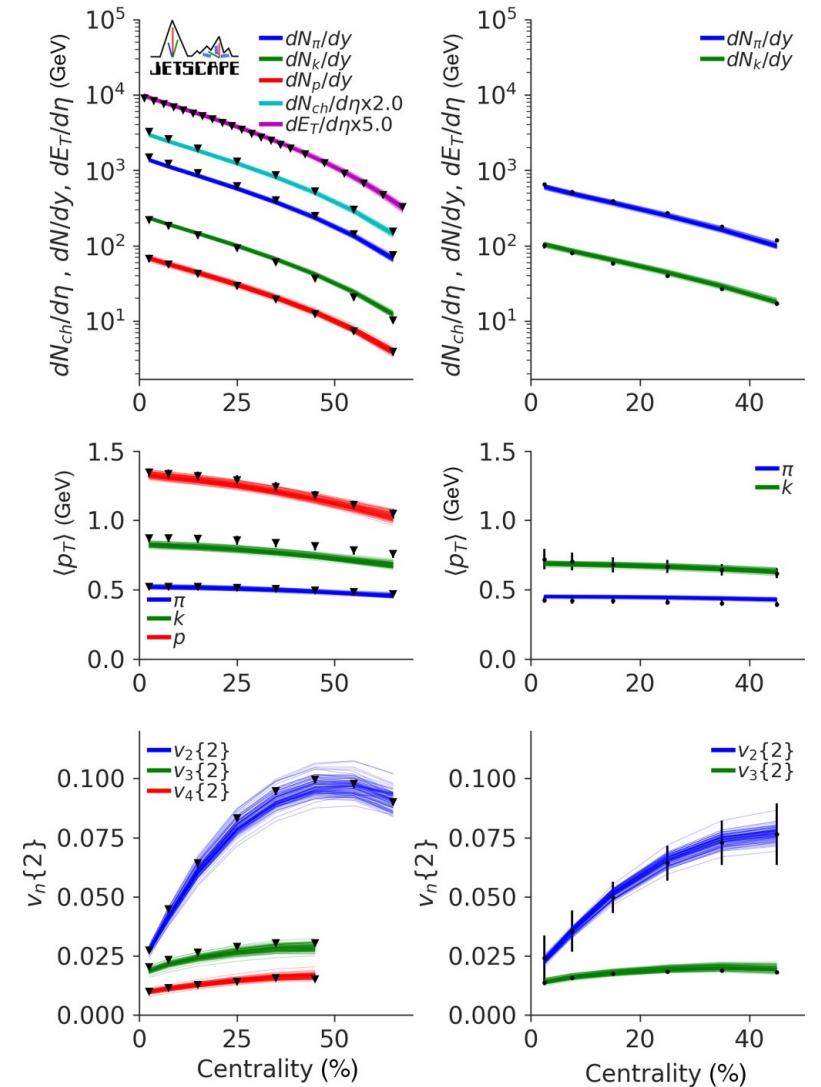
Hydrodynamic-based simulations of heavy ion collisions

- Successful in describing broad sets of measurements

Nijs, van der Schee, Gürsoy, Snellings (2021) PRC, PRL

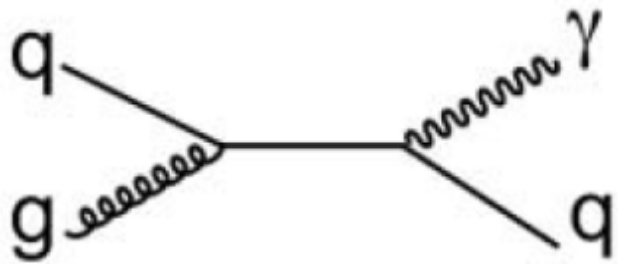


JETSCAPE Collaboration, (2021) PRC, PRL



Direct photons in proton-proton collisions: channels

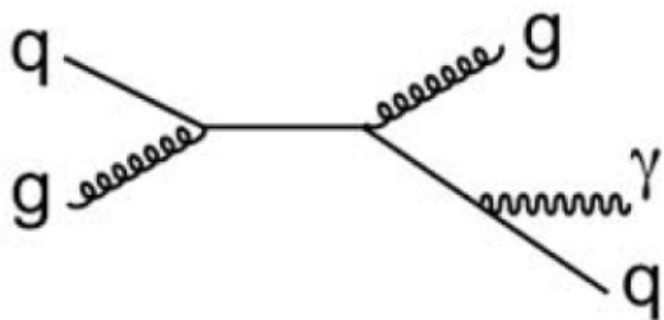
- Hard partonic collisions
 - “Isolated”



(Can be
calculated at
NNLO)

$$d\sigma_{\gamma}^{pp}/dp_T = f_a \otimes f_b \otimes d\hat{\sigma}_{ab \rightarrow \gamma/c+d}$$

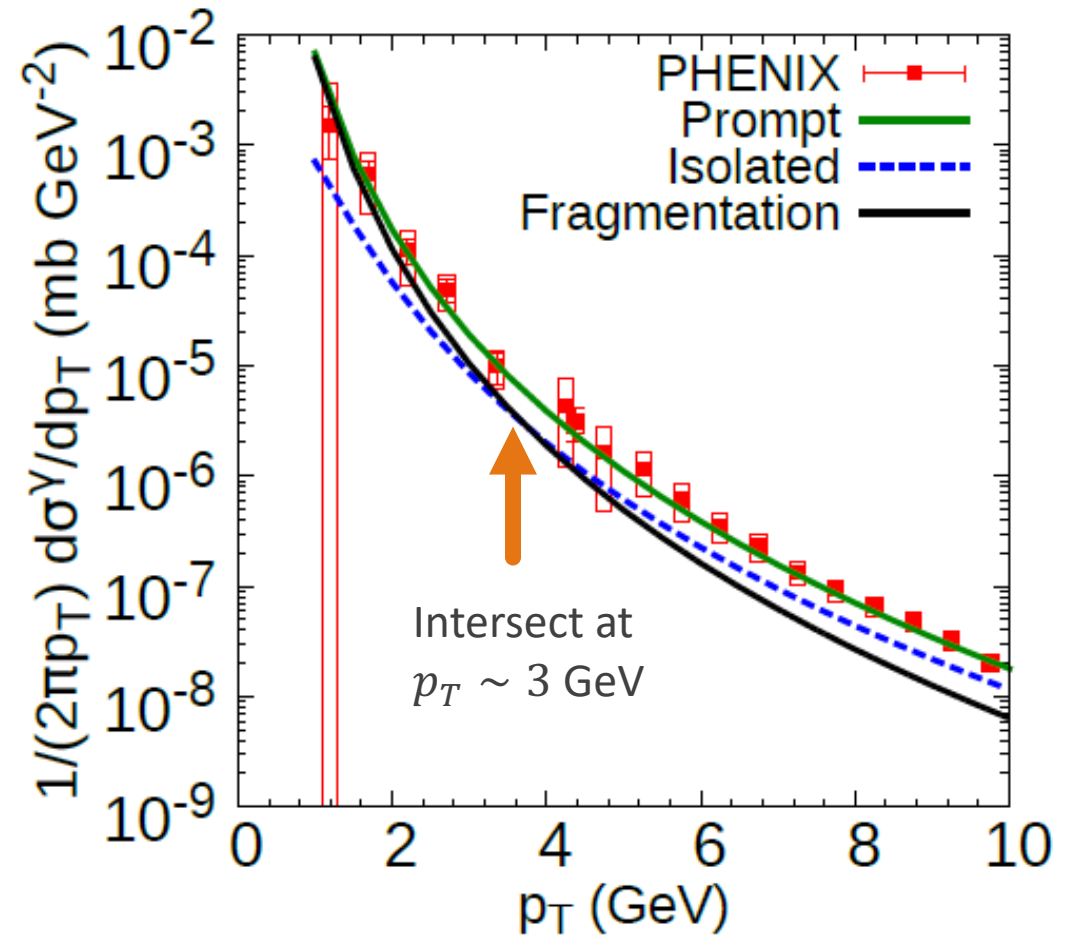
- Fragmentation



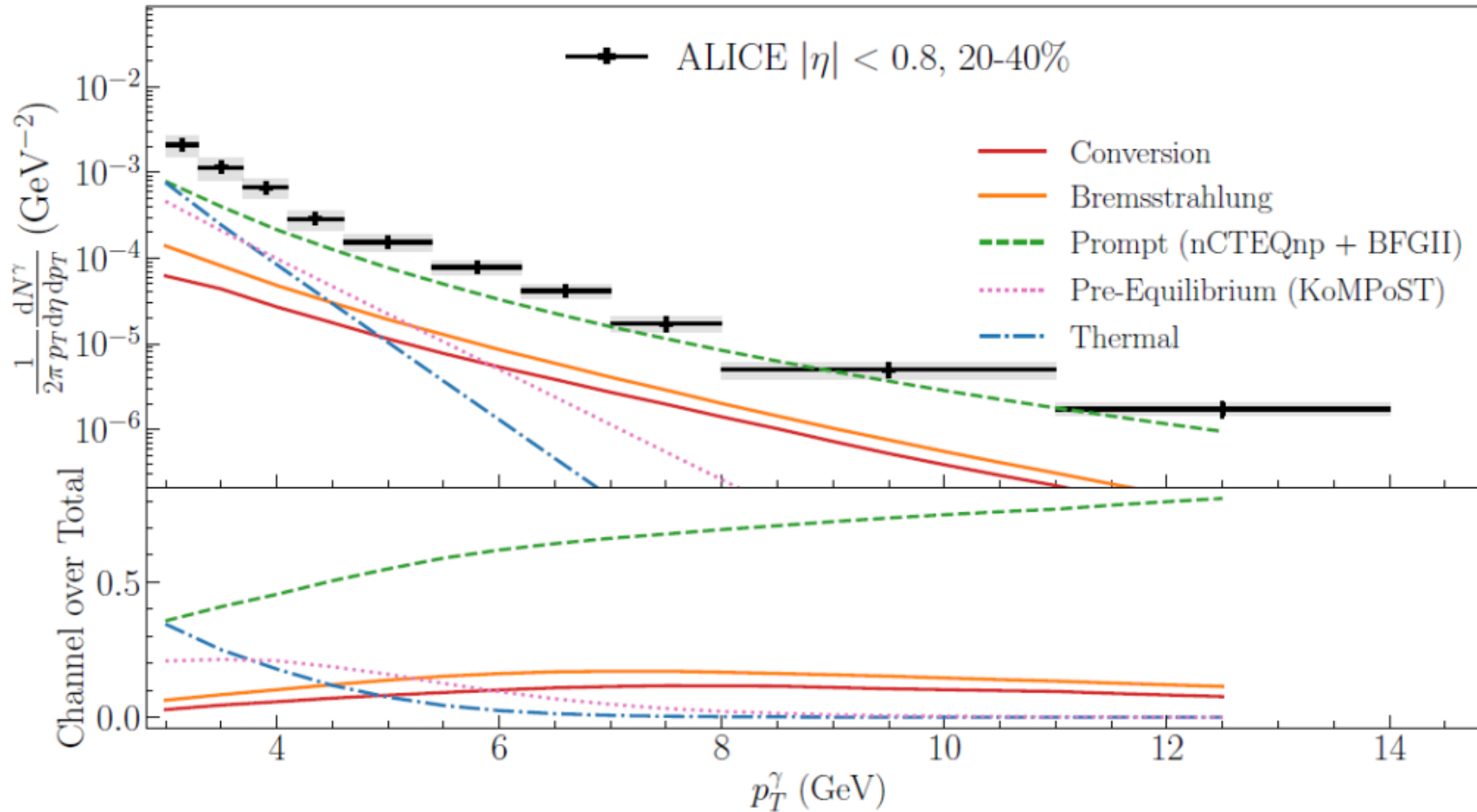
(Fragmentation
function
unmeasured at
NNLO and poorly
constrained at NLO)

$$d\sigma_{\gamma}^{pp}/dp_T = f_a \otimes f_b \otimes d\hat{\sigma}_{ab \rightarrow \gamma/c+d} \otimes D_{\gamma/c}$$

p+p $\sqrt{s}=0.2$ TeV

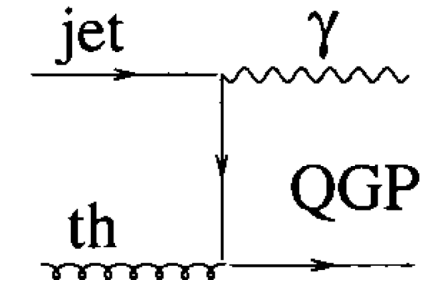


Jet-medium photons

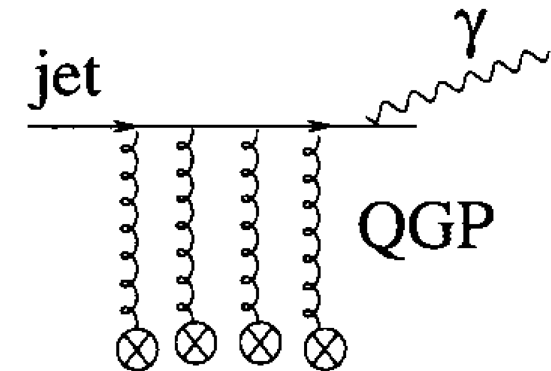


Shi, Modarresi Yazdi, Gale, Jeon (2022)

Conversion



Bremsstrahlung



Photon emission rate

Spacetime profile of plasma

▪ Photon production:
$$\frac{dN_\gamma}{d^3p} = \int d^4X \frac{d\Gamma_\gamma}{d^3p}(p, T(X), u^\mu(X), \dots)$$

Photon emission rate

Temperatures

$$T < 160 \text{ MeV}$$

$$T \approx 160 - 200 \text{ MeV}$$

$$T \sim 200 - 500 \text{ MeV}$$

$$T \gg 1 \text{ GeV}$$

Photon emission rate

Effective hadronic models: Texas A&M/McGill rates (“massive Yang-Mills Lagrangian”), Stony Brook rate (“chiral reduction”)

Lattice QCD:

Cè et al (2022) PRD;

Jackson&Laine (2019) JHEP;

Ghiglieri et al (2016) PRD

Perturbative QCD: Arnold, Moore, Teaney;
Ghiglieri, Moore, Teaney (2013) JHEP

Thermal photon spectrum: Doppler shift

$$\ln \left(\frac{1}{E} \frac{dN_\gamma}{d^3p} \right) = \ln \left(\int d^4X \frac{1}{E} \frac{d\Gamma_\gamma}{d^3p} (p, T(X), u^\mu(X), \dots) \right) \sim cte - \frac{E}{T_{eff}}?$$

Photon emission rate: $\frac{1}{E} \frac{d\Gamma_\gamma}{d^3p} \sim e^{-\frac{E}{T}}$

$$\ln \left(\frac{1}{E} \frac{dN_\gamma}{d^3p} \right) \approx \ln \left(\int d^4X e^{-\frac{P \cdot u(X)}{T(X)}} \right) + cte = \ln \left(\int d\phi d\eta_s dx_\perp e^{-\frac{P \cdot u(X)}{T(X)}} \right) + cte$$

Doppler shift

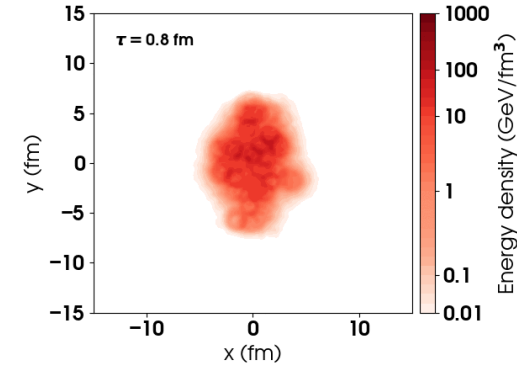


At midrapidity, $P \cdot u = p_T \left(\cosh(\eta_s) \sqrt{1 + u_\perp^2} - u_\perp \cos(\phi) \right)$

Thermal photon spectrum: Doppler shift

$$\ln \left(\frac{1}{E} \frac{dN_\gamma}{d^3p} \right) = \ln \left(\int d^4X \frac{1}{E} \frac{d\Gamma_\gamma}{d^3p} (p, T(X), u^\mu(X), \dots) \right) \sim cte - \frac{E}{T_{eff}} ?$$

Photon emission rate: $\frac{1}{E} \frac{d\Gamma_\gamma}{d^3p} \sim e^{-\frac{E}{T}}$



$$\begin{aligned} \ln \left(\frac{1}{E} \frac{dN_\gamma}{d^3p} \right) &\approx \ln \left(\int d^4X e^{-\frac{P \cdot u(X)}{T(X)}} \right) + cte = \ln \left(\int d\phi d\eta_s dx_\perp e^{-\frac{P \cdot u(X)}{T(X)}} \right) + cte \\ &\approx \ln \left(\int dx_\perp \exp \left(-\frac{E}{T \left(1 + \frac{u_\perp^2}{4E/T} (1 + (E/T - 2)(E/T)) \right)} \right) \right) + cte \end{aligned}$$

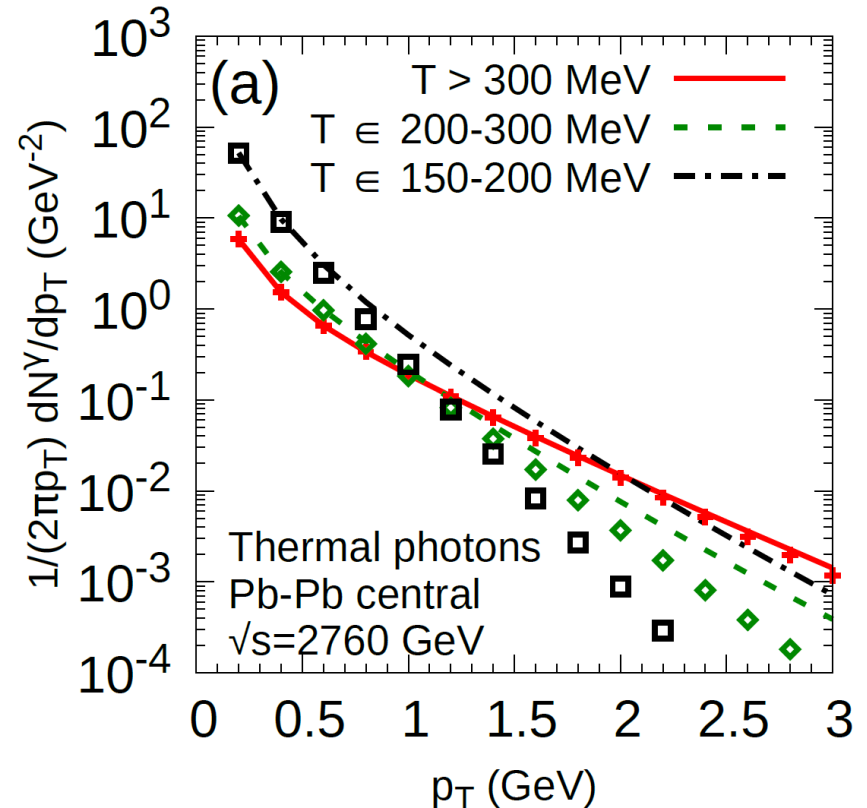
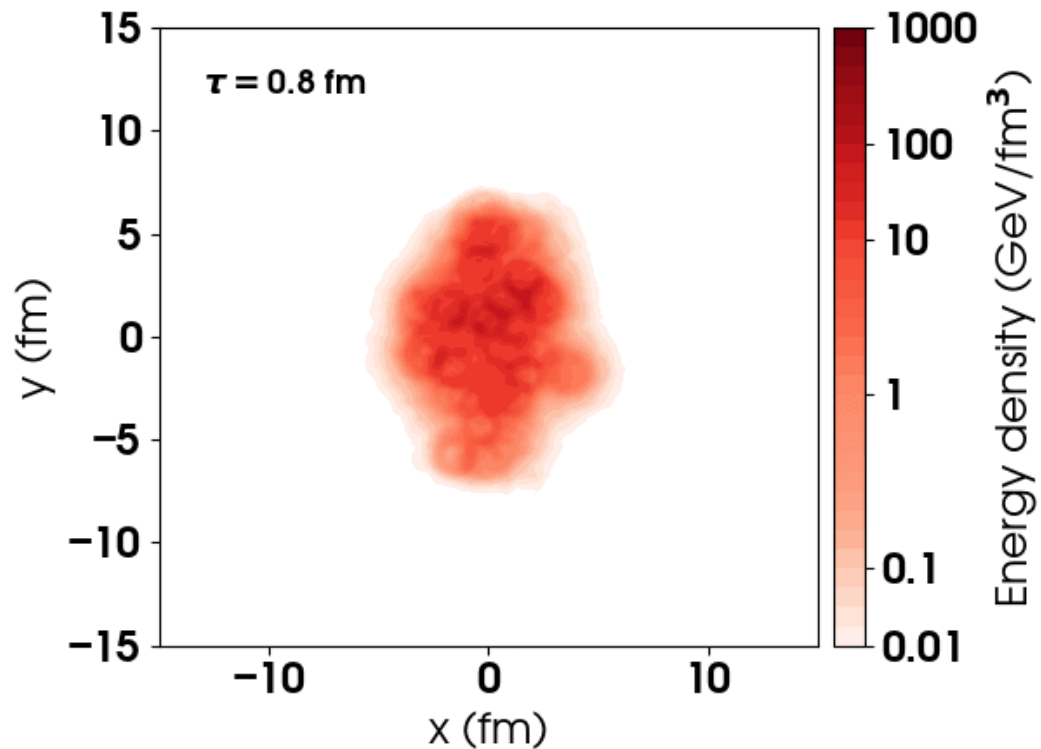
Doppler shift



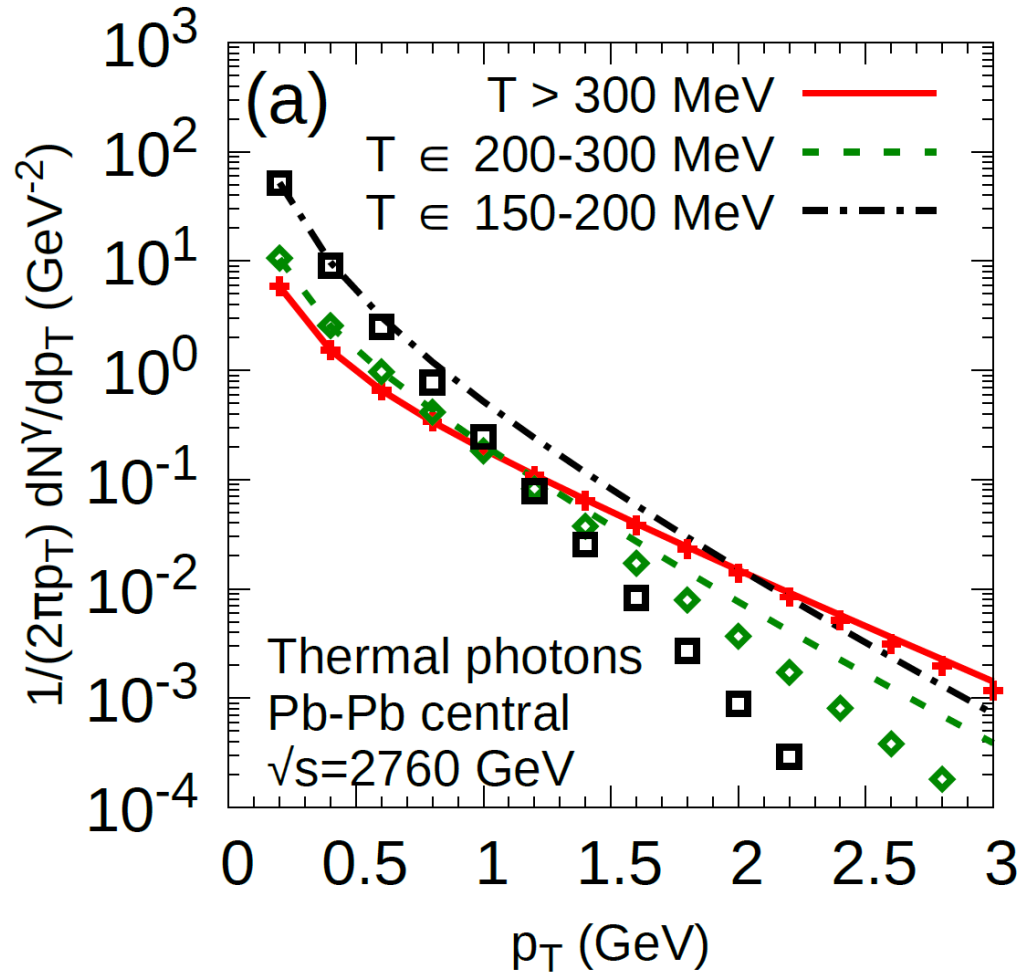
Thermal photon spectrum: Doppler shift

$$\ln \left(\frac{1}{E} \frac{dN_\gamma}{d^3p} \right) \approx \ln \left(\int dx_\perp \exp \left(- \frac{E}{T \left(1 + \frac{u_\perp^2}{4E/T} (1 + (E/T - 2)(E/T)) \right)} \right) \right) + cte$$

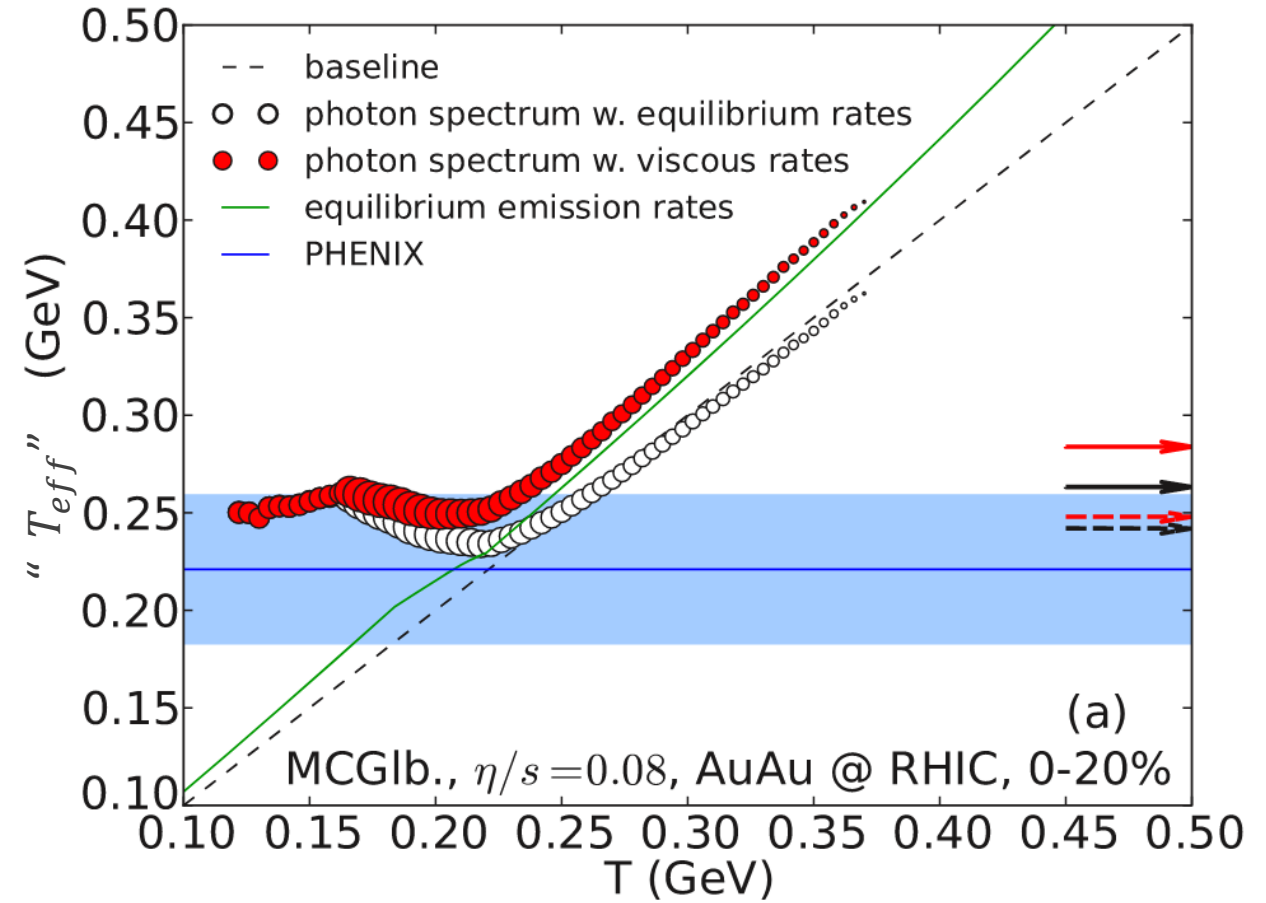
← Transverse Doppler shift



Effect of transverse Doppler shift

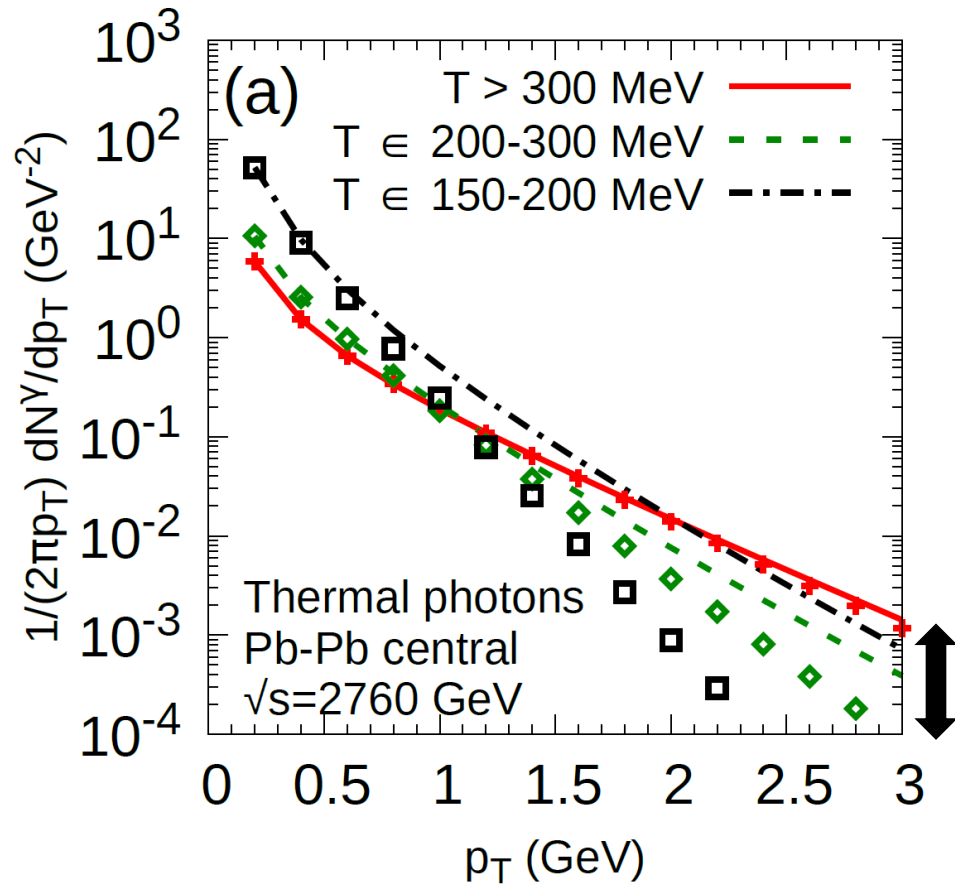


Local effect of Doppler shift

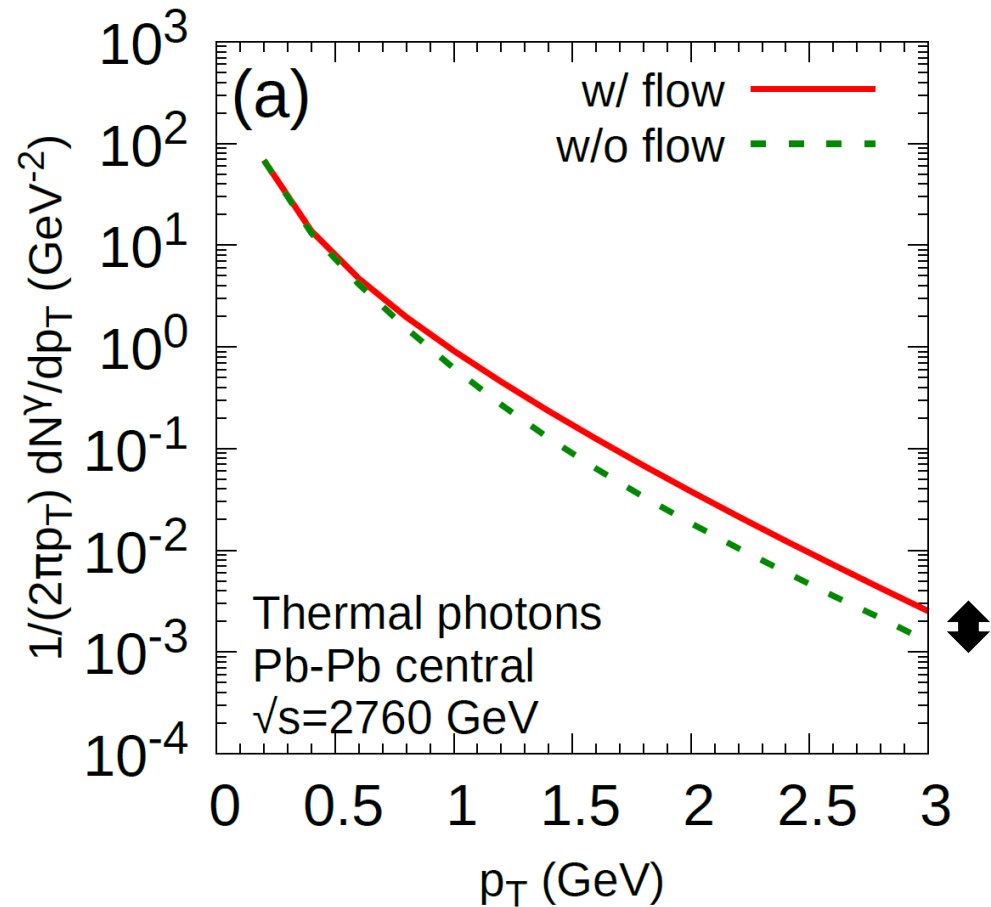


Ref.: Shen, Heinz, Paquet, Gale (2014) PRC;
See also van Hees, Gale, Rapp (2011) PRC

Effect of transverse Doppler shift



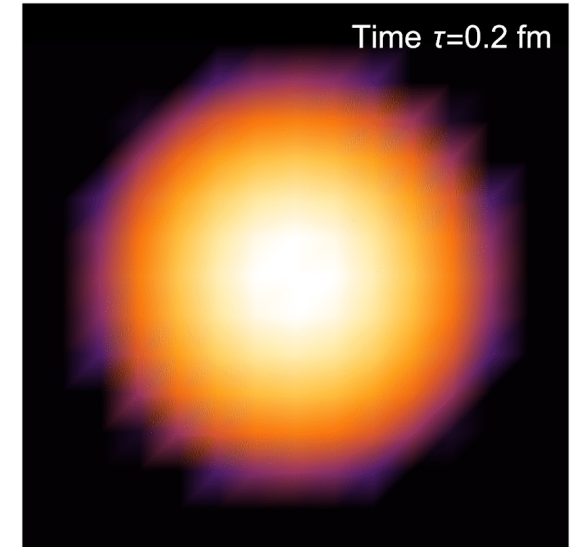
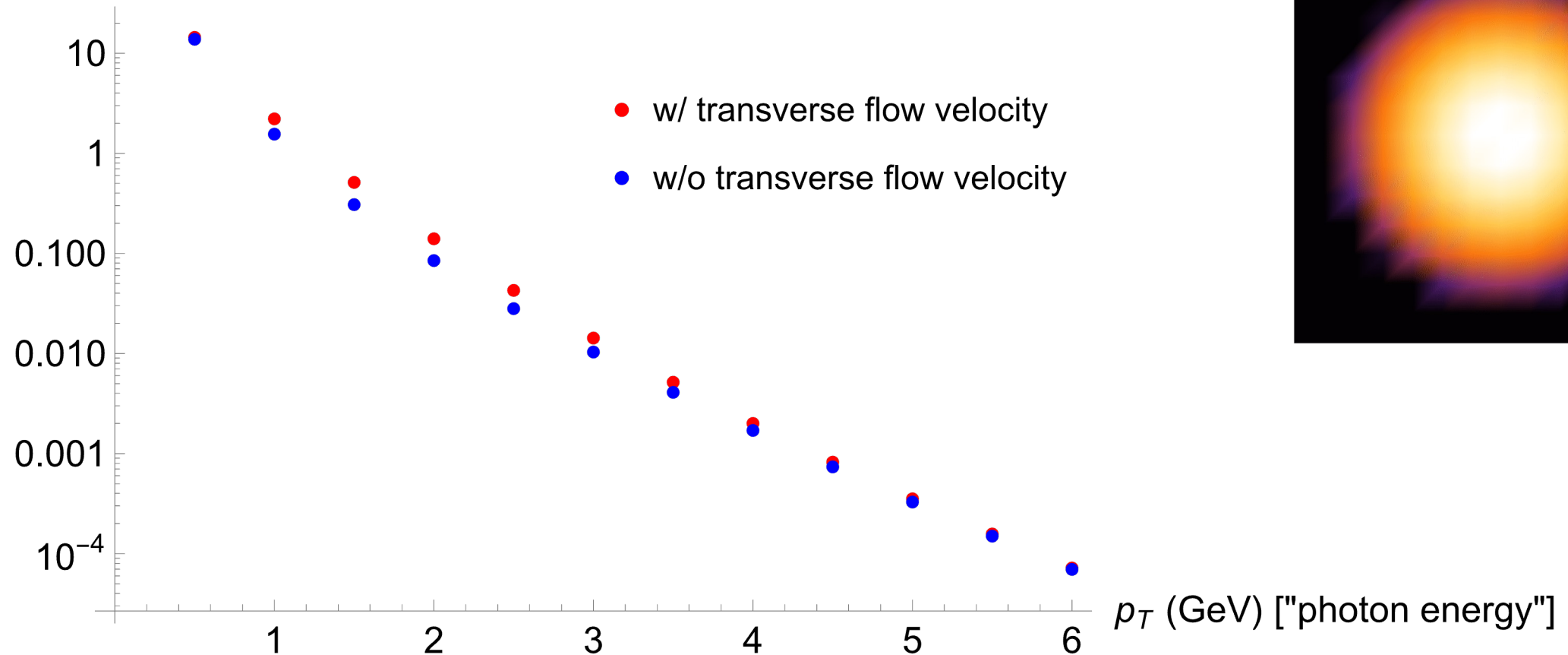
Local effect of Doppler shift



Global effect of Doppler shift

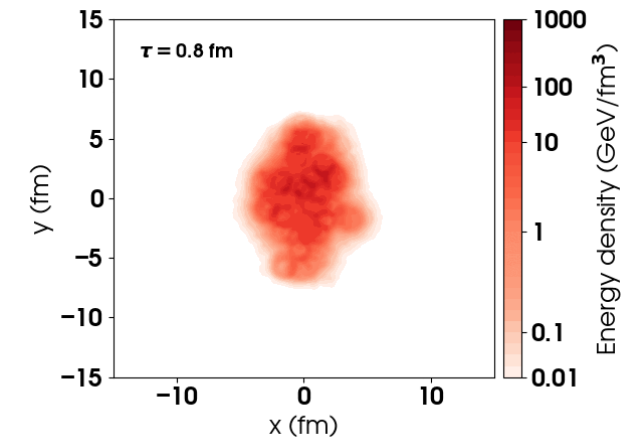
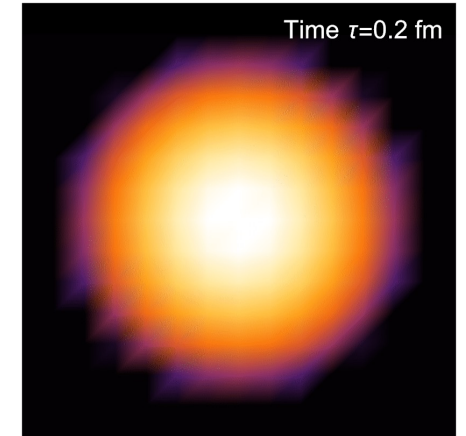
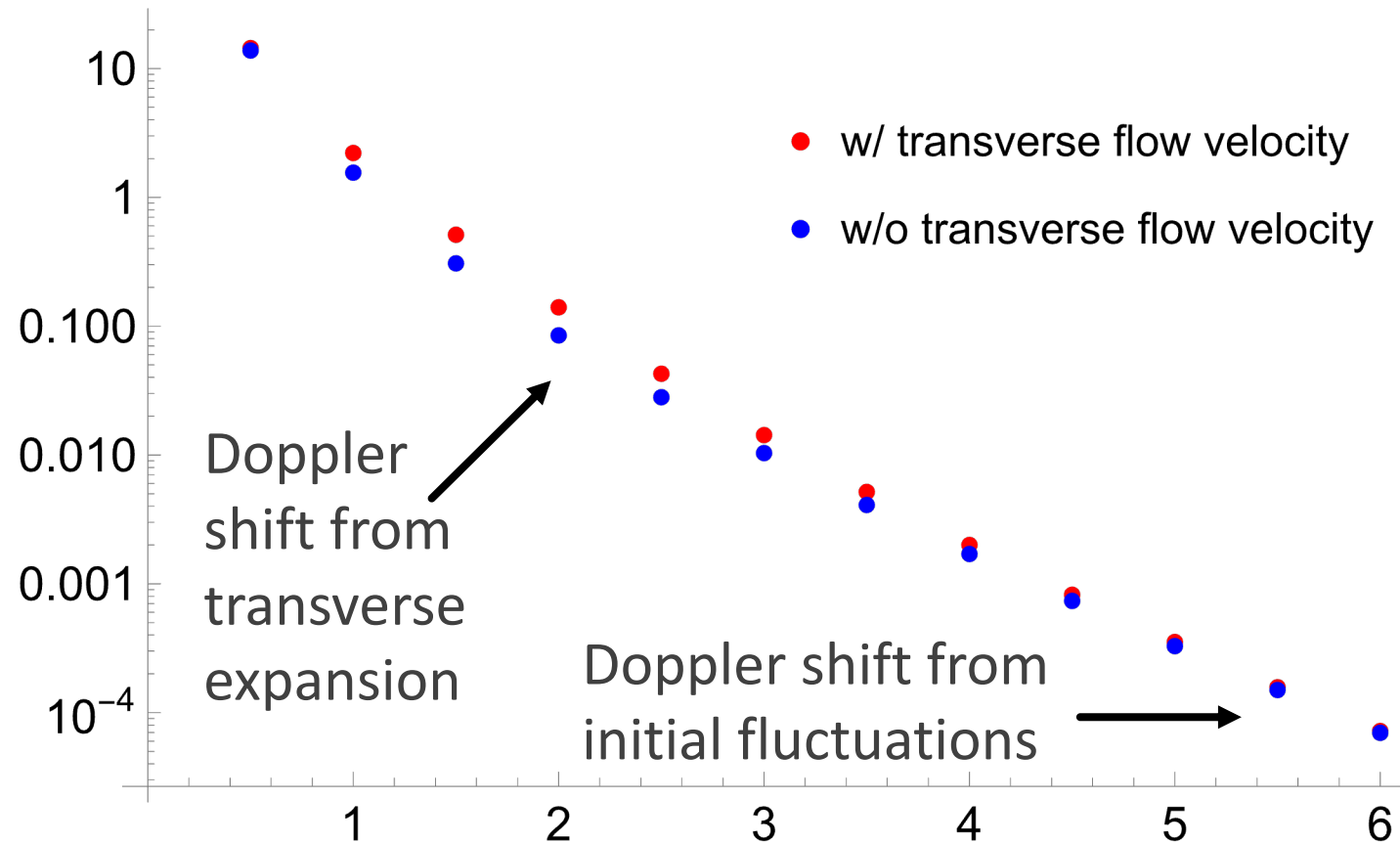
Not all Doppler shifts are equal

$$1/(2\pi p_T) dN/dp_T dy$$



Different origins of the Doppler shift

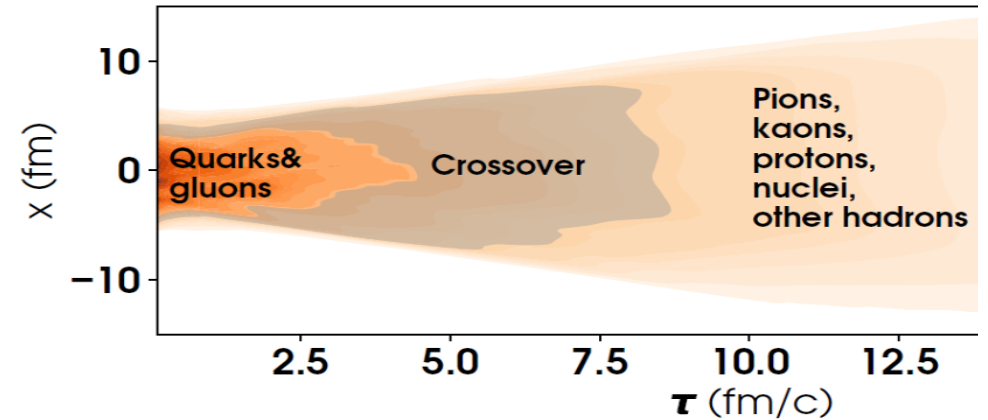
$$1/(2\pi p_T) dN/dp_T dy$$



p_T (GeV) ["photon energy"]

Thermal photon spectrum

$$\ln \left(\frac{1}{E} \frac{dN_\gamma}{d^3p} \right) \approx \ln \left(\int d\phi d\eta_s dx_\perp e^{-\frac{P \cdot u(X)}{T(X)}} \right) + cte$$

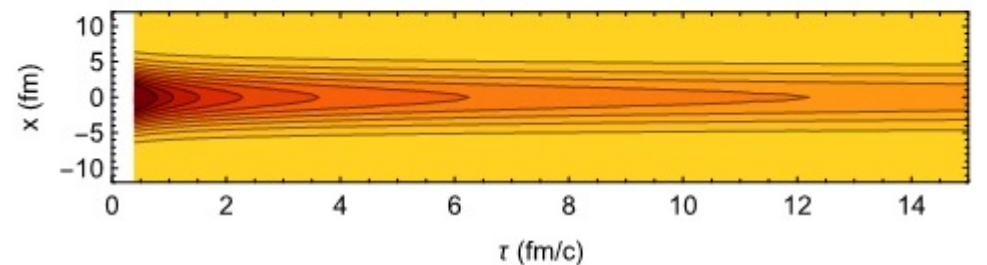
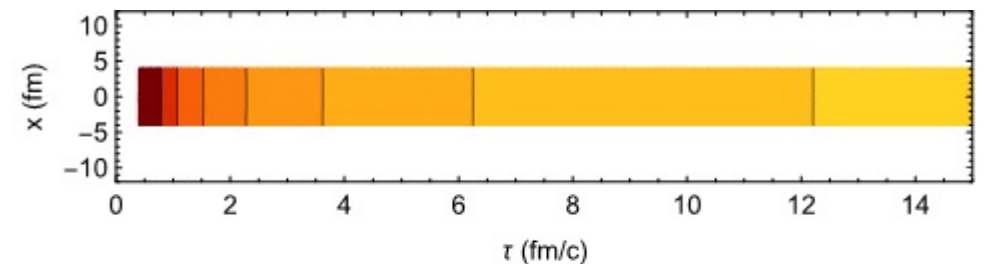


Spacetime profile of plasma: complicated, but can look at simple models

Bjorken hydrodynamics for longitudinal-dominated expansion: $T(\tau) = T_0 \left(\frac{\tau_0}{\tau} \right)^{c_s^2}$

➔ Black disk approx: $T(\tau, r < \sigma) = T_0 \left(\frac{\tau_0}{\tau} \right)^{c_s^2}$

➔ Gaussian approx: $T(\tau, r) = T_0 e^{-\frac{r^2}{2\sigma^2}} \left(\frac{\tau_0}{\tau} \right)^{c_s^2}$



Paquet and Bass [arXiv:2205.12299]

Thermal photon spectrum

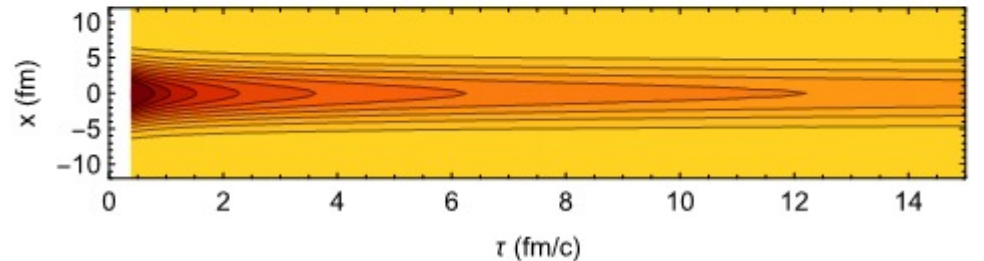
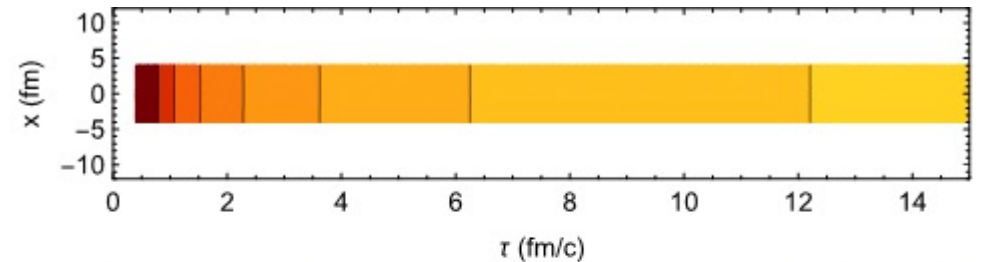
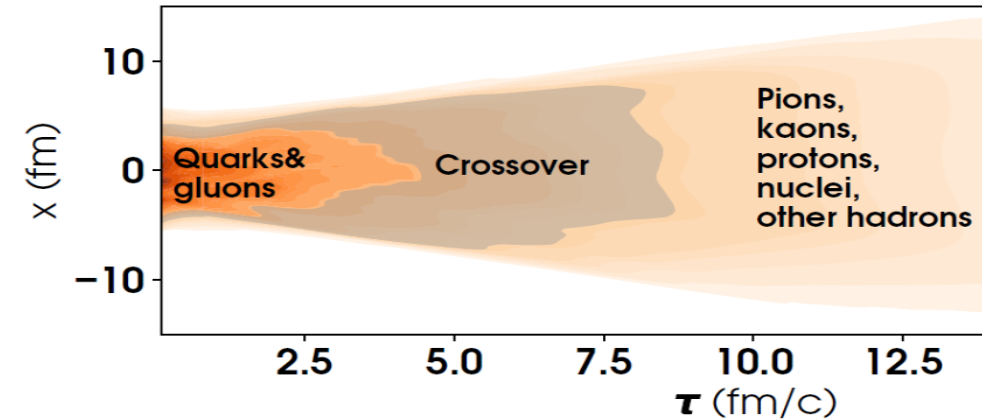
$$\ln \left(\frac{1}{E} \frac{dN_\gamma}{d^3p} \right) \approx \ln \left(\int d\phi d\eta_s dx_\perp e^{-\frac{P \cdot u(X)}{T(X)}} \right) + cte$$

$$\ln \left(\frac{1}{E} \frac{dN_\gamma}{d^3p} \right) \approx -\frac{E}{T_0} + \frac{3}{2} \log \left(\frac{T_0}{E} \right) + cte + O \left(\frac{T_0}{E} \right)$$

Paquet and Bass [arXiv:2205.12299]

$$\ln \left(\frac{1}{E} \frac{dN_\gamma}{d^3p} \right) \approx -\frac{E}{T_0} + \frac{5}{2} \log \left(\frac{T_0}{E} \right) + cte + O \left(\frac{T_0}{E} \right)$$

$$\ln \left(\frac{1}{E} \frac{dN_\gamma}{d^3p} \right) \approx -\frac{E}{T_0} + \mu \log \left(\frac{T_0}{E} \right) + cte \approx -\frac{E}{T_{eff}} + cte$$



$$T_0 \approx \frac{T_{eff}}{1 - \frac{T_{eff}}{E} \mu \ln \mu}$$

Matching to 0+1D (boost-invariant) hydrodynamics

- In 0+1D hydro, we can characterize $T^{\mu\nu}$ with single component: energy density
- 0+1D dynamical models with smooth transition to hydrodynamics:
 - Kinetic theory (gluons, QCD, RTA) or AdS/CFT

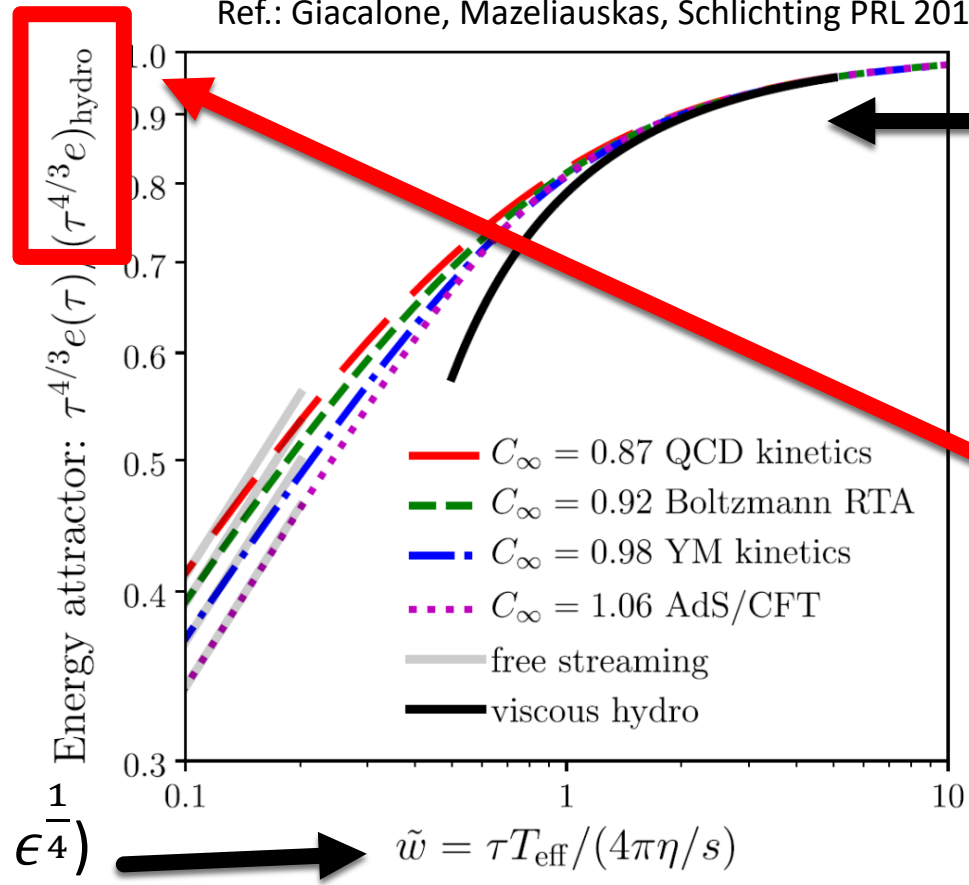
Conclusion:

Properly scaled 0+1D systems approach hydro similarly

Timescale necessary to converge to hydro depends:

- Strength of interaction $\left(\frac{\eta}{s} \sim \frac{1}{\alpha_s^2}\right)$
- Energy density of the system
(or “effective temperature” $T_{eff} \propto \epsilon^{\frac{1}{4}}$)

Ref.: Giacalone, Mazeliauskas, Schlichting PRL 2019



All systems converge to 0+1D viscous hydro

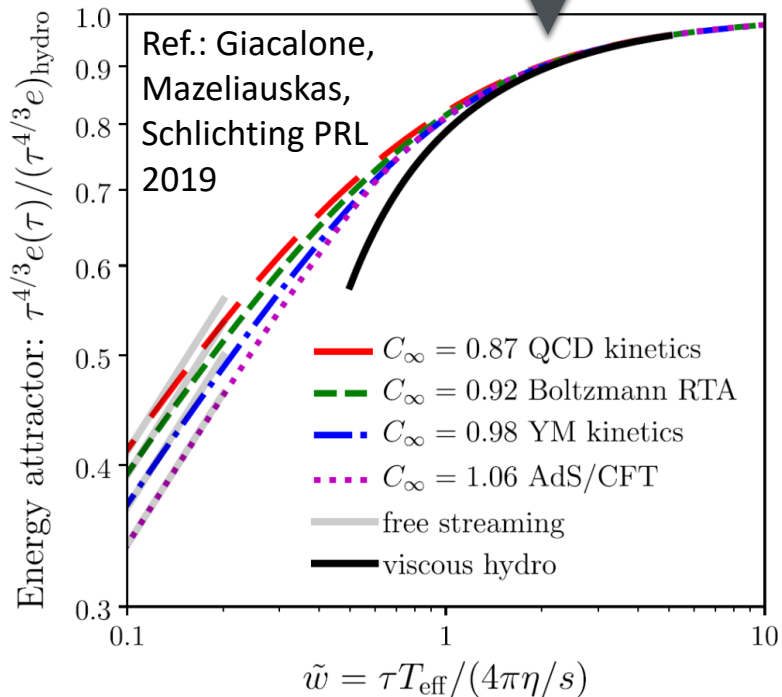
(note: not all systems have the same final energy density, but this can be rescaled)

Matching to 2+1D hydrodynamics: “KøMPøST”

- Take a 2+1D pre-hydro system: how does it approach hydrodynamics?
- Better approximation [KøMPøST]: decompose $T^{\mu\nu}$ in 0+1D background + **linear**

$$T^{\mu\nu}(\tau_{\text{hydro}}, \mathbf{x}) = \boxed{\bar{T}_{\mathbf{x}}^{\mu\nu}(\tau_{\text{hydro}})} + \frac{\bar{T}_{\mathbf{x}}^{\tau\tau}(\tau_{\text{hydro}})}{\bar{T}_{\mathbf{x}}^{\tau\tau}(\tau_{\text{EKT}})} \int d^2\mathbf{x}' G_{\alpha\beta}^{\mu\nu}(\mathbf{x}, \mathbf{x}', \tau_{\text{hydro}}, \tau_{\text{EKT}}) \boxed{\delta T_{\mathbf{x}}^{\alpha\beta}(\tau_{\text{EKT}}, \mathbf{x}')}$$

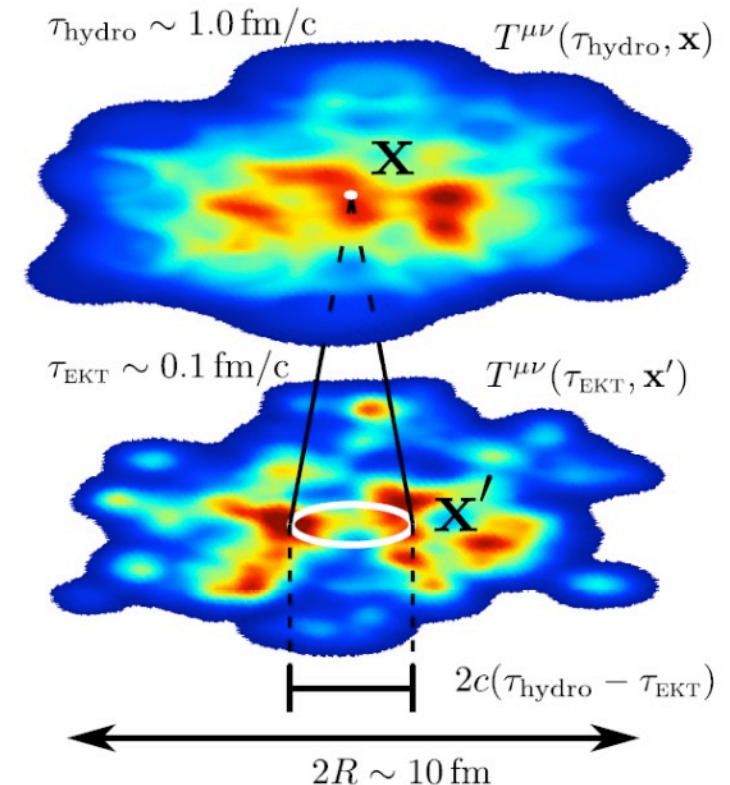
Ref.: Kurkela, Mazeliauskas, Paquet, Schlichting, Teaney PRL2019, PRC2019



Response functions describing evolution of perturbations

- Can be evaluated with QCD kinetic theory
- Also exhibits scaling with interaction strength and energy density of system

[See also Kamata, Martinez, Plaschke, Ochsenfeld, Schlichting, PRD 2020]



Dilepton rate

Figure by Gojko Vujanovic

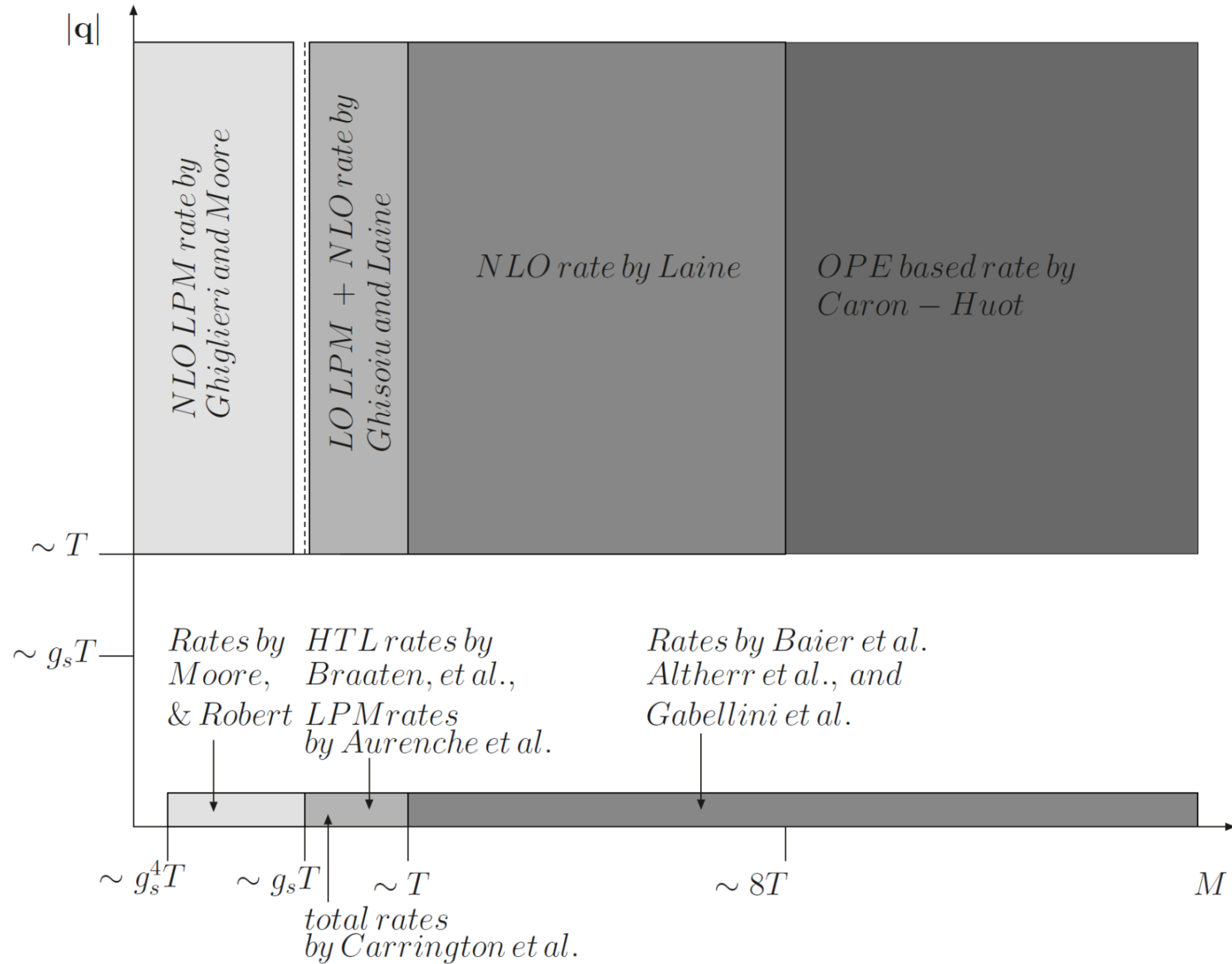
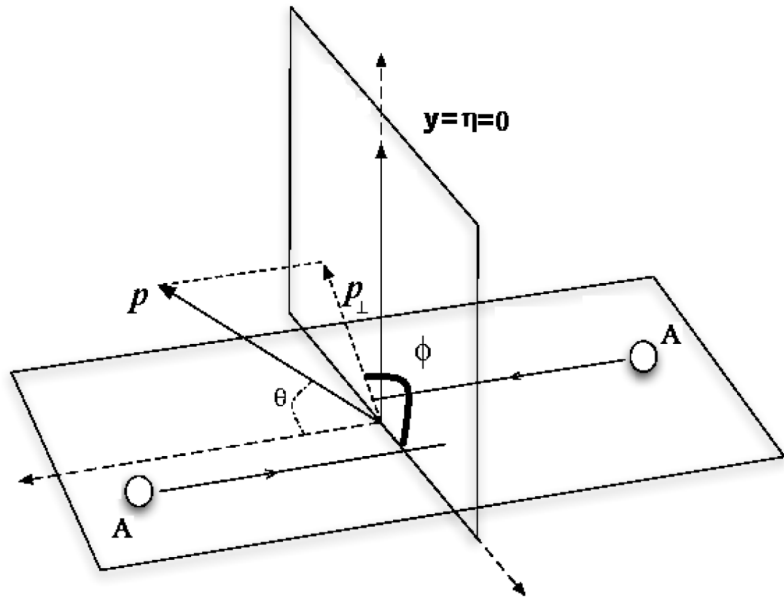
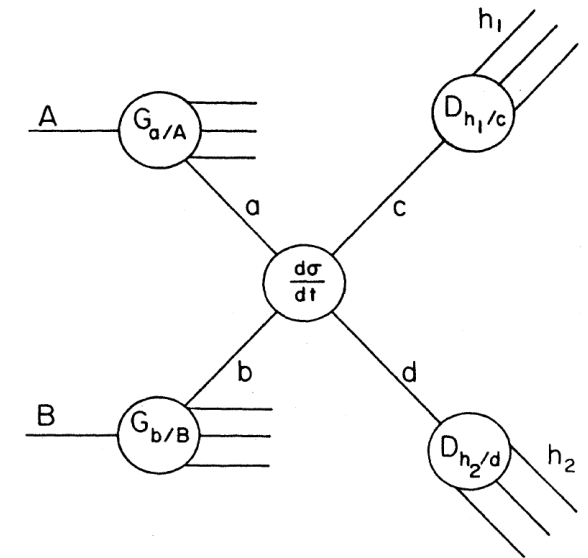


Figure adapted from K. Tuchin (2013) AHEP



PROTON-PROTON COLLISIONS



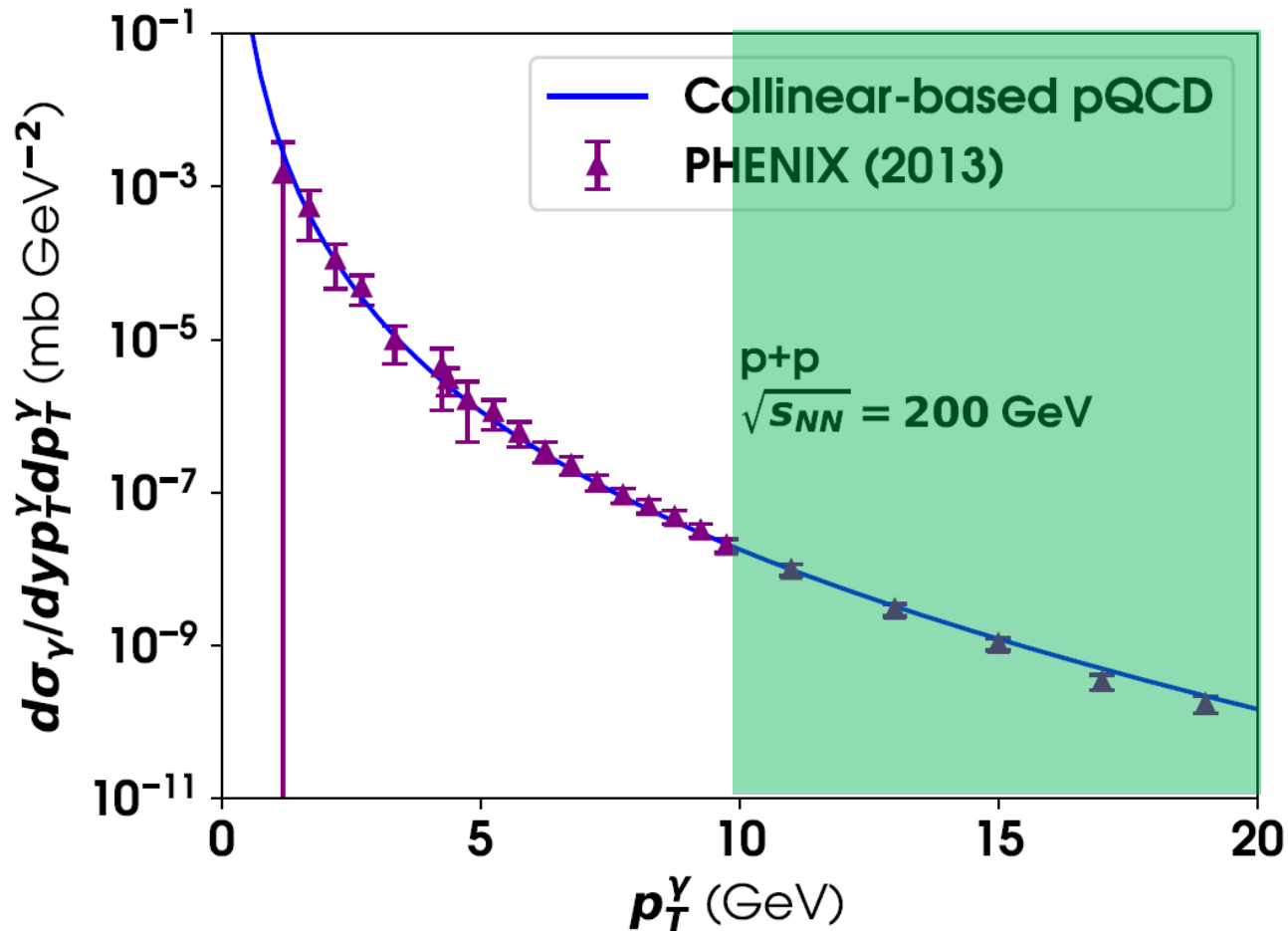
Ref: Owens (1987) RMP

Direct photons in p-p collisions: high energy

Nuclear Physics B327 (1989) 105–143
North-Holland, Amsterdam

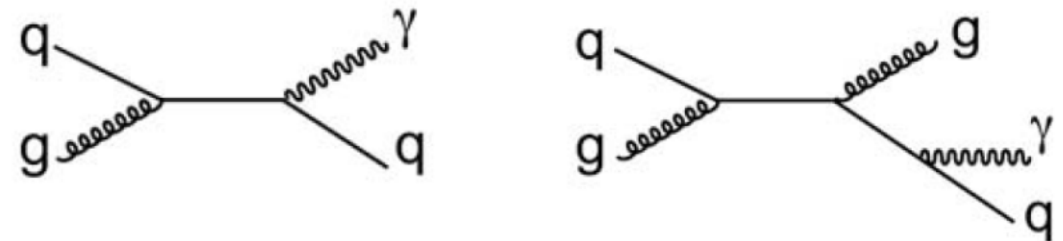
QCD CORRECTIONS TO PARTON-PARTON SCATTERING PROCESSES

F. AVERSA*, P. CHIAPPETTA, M. GRECO*, J.Ph. GUILLET**



- Can be calculated in collinear-factorization based perturbative QCD, up to next-to-leading order

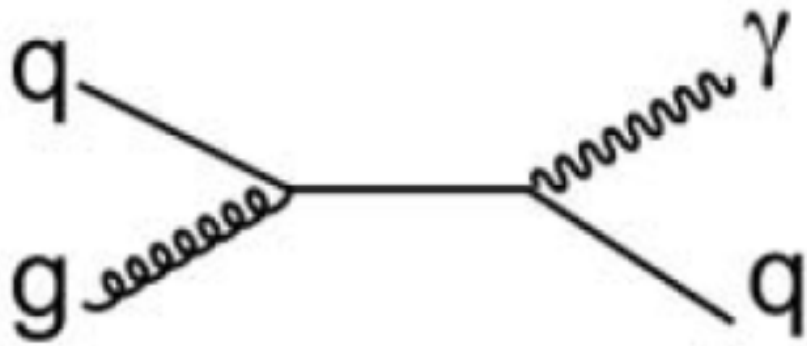
$$\frac{d\sigma_{\gamma}^{pp}}{dp_T} = f_a \otimes f_b \otimes d\hat{\sigma}_{ab \rightarrow \gamma/c+d} [\otimes D_{\gamma/c}]$$



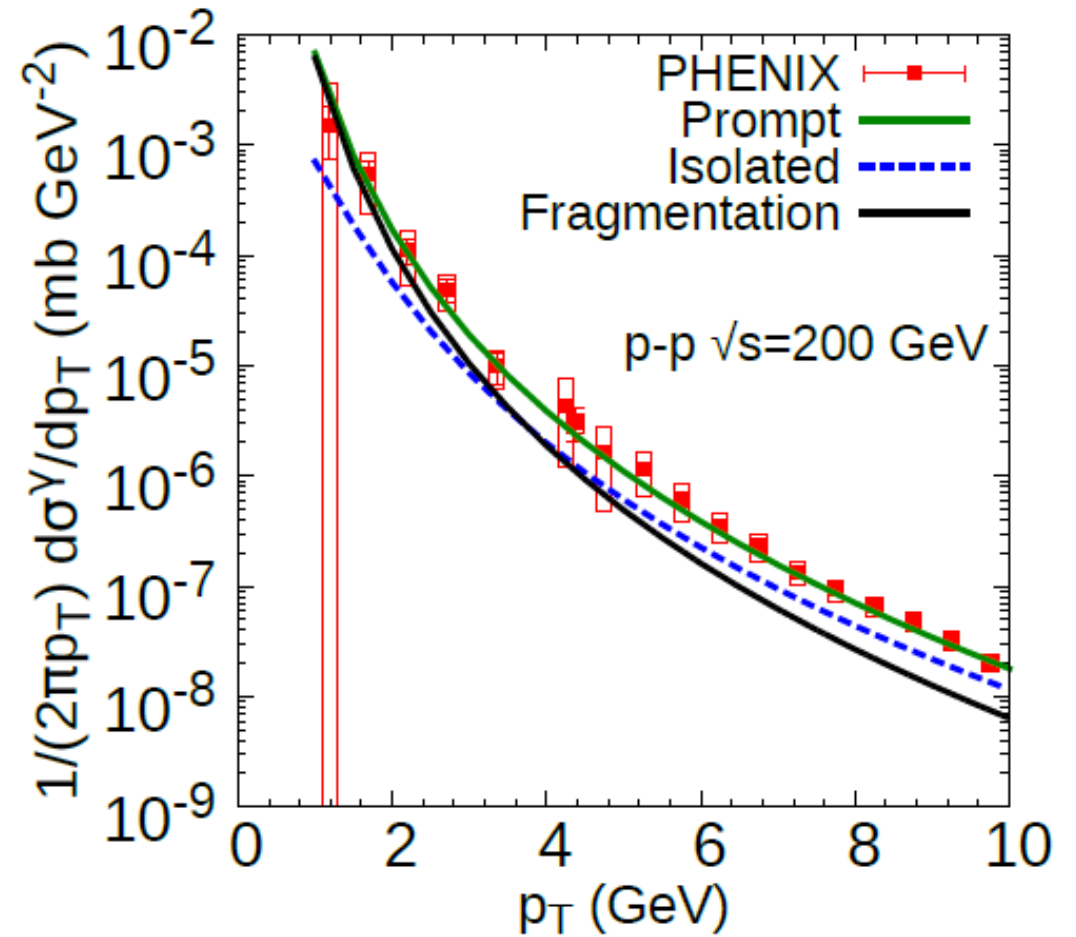
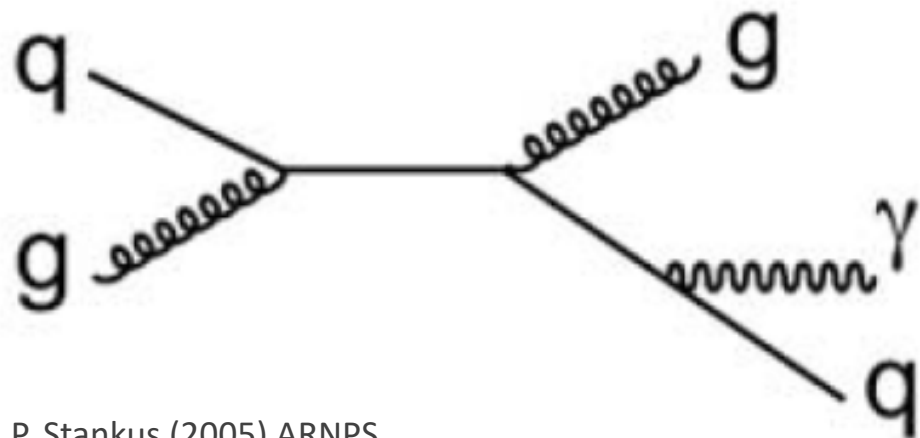
Frag fct: Bourhis, Fontannaz, Guillet (1998) EPJ

Direct photons in proton-proton collisions: channels

- Hard partonic collisions
 - “Isolated”

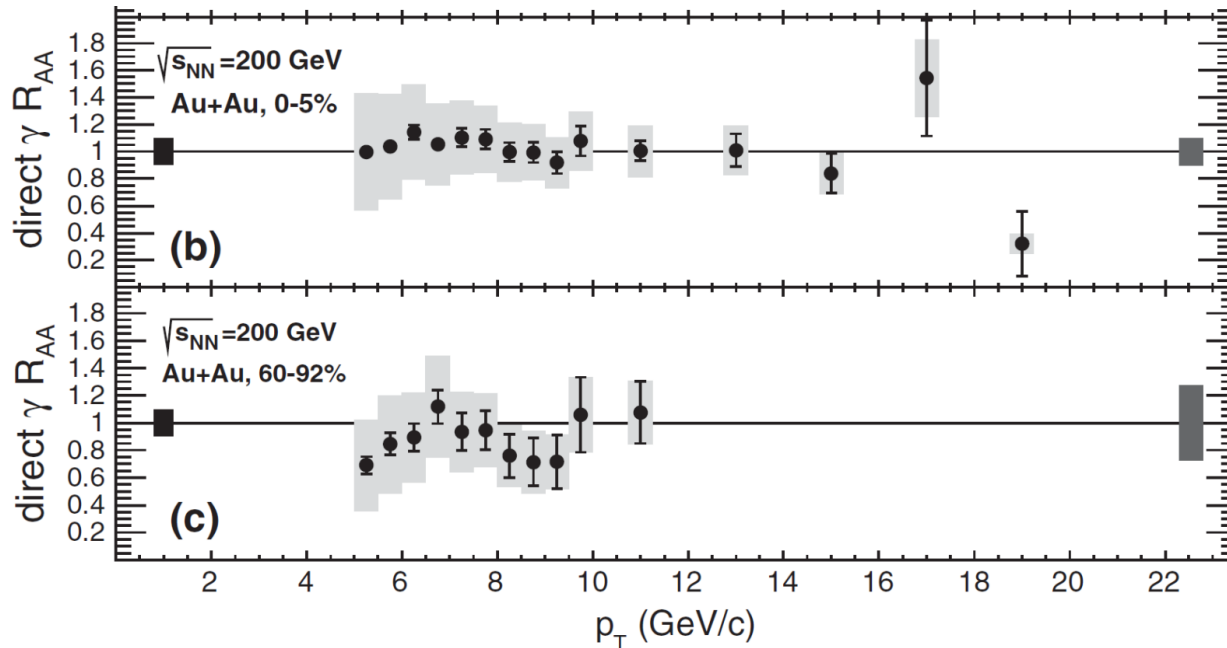


- Fragmentation



Photons in heavy-ion collisions: high p_T

- **Prompt photons** produced as superposition of nucleon-nucleon collisions (“binary scaling”)



Ref.: PHENIX Collaboration (2012) PRL

$$R_{AA}^{\gamma} = \frac{\frac{dN_{\gamma}^{AA}}{dp_T}}{\left(\frac{N_{binary}}{\sigma_{pp}^{inel}}\right) \frac{d\sigma_{\gamma}^{pp}}{dp_T}} \approx 1 \quad (\text{at high } p_T)$$

Deviations from $R_{AA}^{\gamma} = 1$ originate from:

- Isospin effect (parton content of n vs p)
- Nuclear effects on parton distribution functions
- Parton energy loss [more about this later]