# Recent Results in Conformal Killing Gravity (CKG)

#### Alan Barnes

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• A 3<sup>rd</sup> order metric theory of gravitation due to Harada (2023) with field equations (HFEs):  $\tilde{G}_{(ab;c)} = \tilde{T}_{(ab;c)}$ where  $\tilde{G}_{ab} \equiv G_{ab} - \frac{1}{6}Gg_{ab}$  and  $\tilde{T}_{ab} \equiv T_{ab} - \frac{1}{6}Tg_{ab}$ .



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• Actually 
$$K_{ab} \equiv \tilde{G}_{ab} - \tilde{T}_{ab}$$
 is a Killing tensor and  $C_{ab} = K_{ab} + \frac{1}{2}Kg_{ab}$ .



#### Yet another theory – why study CKG?

• Harada considered the evolution of the scale factor *a*(*t*) of the Friedmann-Robertson-Walker-Lemaître (FRWL) metric:

$$ds^{2} = \mathrm{d}t^{2} - a^{2}(t) \left( \frac{\mathrm{d}r^{2}}{1 - kr^{2}} + r^{2}\mathrm{d}\theta^{2} + r^{2}\sin^{2}\theta\mathrm{d}\phi^{2} \right).$$

• He obtained a third order ODE which has the first integral:

$$2\left(\frac{\dot{a}}{a}\right)^2 - \frac{\ddot{a}}{a} + \frac{2k}{a^2} = \frac{4\pi G}{3}(5\rho + 3p) + \frac{\Lambda}{3}$$



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- He showed that even in the case of a matter-dominated universe (p = 0) with Λ = 0 there could be a transition from deceleration to accelerating expansion.
- Thus removing the need to assume the existence of dark energy.

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### Static black hole solutions

• Harada(2023) obtained a metric for a static spherically symmetric vacuum field in CKG:

$$\mathrm{d}s^{2} = \mathrm{e}^{\mathbf{a}(r)}\mathrm{d}t^{2} - \left(\mathrm{e}^{-\mathbf{a}(r)}\mathrm{d}r^{2} + r^{2}(\mathrm{d}\theta^{2} + \sin^{2}\theta\mathrm{d}\phi^{2})\right) \quad (1)$$

where  $e^{a(r)} = 1 - \frac{2m}{r} - \frac{1}{3}\Lambda r^2 - \frac{1}{5}\lambda r^4$  where *m*,  $\Lambda$  and  $\lambda$  are integration constants. Generalised Schwarzschild-Kottler solution.



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- However (1) is not the most general form for a spherically symmetric static metric which is

$$\mathrm{d}s^{2} = \mathrm{e}^{\mathbf{a}(r)}\mathrm{d}t^{2} - \left(\mathrm{e}^{\mathbf{b}(r)}\mathrm{d}r^{2} + r^{2}(\mathrm{d}\theta^{2} + \sin^{2}\theta\mathrm{d}\phi^{2})\right).$$



• The HFEs for the metric(2) with  $T_{ab} = 0$  lead to 2 ODEs of third order for a(r) and b(r). However, writing b = 2f - a one may deduce a 2nd order ODE  $rf'' - 2rf'^2 - f' = 0$ . Hence  $e^b = e^{-a}/(c + dr^2)$  where c and d are constants.



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- Eliminating b the remaining field equation may be linearised by writing y = e<sup>a</sup> to yield

$$(c+dr^{2})r^{3}y'''-(2c-dr^{2})r^{2}y''-(2c+dr^{2})ry'+8cy=8. (3)$$



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• If d = 0, c may be set to 1 by a rescaling of the *t*-coordinate. Similarly if c = 0, d may be set to 1. In either case the ODE becomes homogeneous and easy to solve. Otherwise we may set  $c = \pm 1$ . For Lorentzian signature  $c + dr^2 > 0$ .



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- For *c* = 0 there are strange vacuum & electrovac solutions (Barnes 2023):

$$y = e^a = \lambda - \Lambda r^2 / 3 + m \log r - 1 / (2r^2) + q^2 / (4r^4), \quad e^b = e^{-a} / r^2.$$

• Physical Interpretation ???!!!



• For c > 0 and  $d \neq 0$ , Barnes(2023) obtained solutions involving power series with radius of convergence  $1/\sqrt{|d|}$ .  $y = e^a = 1 + (1+2dr^2)q^2/r^2 - 2mp_1(r)/r - \lambda p_2(r)r^4/5 - \Lambda r^2/3$  $p_1(r) = 1 + dr^2/2 - d^2r^4 \dots$  $p_2(r) = 1 - 4dr^2/7 + 8d^2r^4/21\dots$  $e^b = e^{-a}/(1 + dr^2).$ 



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- When any of d, Λ or λ are non-zero, the solution is not asymptotically flat.
- When  $m = \lambda = 0$ , the power series disappear. We have type D 'cosmological' solutions.
- If, in addition,  $q = \Lambda = 0$  we have the static Einstein Universe which is a vacuum solution in CKG!



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- In these the scale factor *a* satisfies  $\dot{a}^2 \alpha a^4 \frac{\Lambda}{3}a^2 + k$ . The solutions in general involve Jacobi elliptic functions, but are elementary if  $\alpha = 0$  or  $\frac{\Lambda^2}{9} + 4k\alpha = 0$ .



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- Open question: does CKG admit spherically symmetric matter distributions emitting gravitational waves?



• Defining pp-waves as fields which admit a covariantly constant null bivector:  $W_{ab;c} = 0$  where  $W_{ab} = p_{[a}k_{b]}$  with  $p_ap^a = -1$  and  $k_ak^a = 0$ , it follows that  $k_{a;b} = 0$  and the metric may be written as

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- if  $H = c(x^2 + y^2)$  the metric is a non-flat vacuum conformally flat field in CKG.

#### Multiple Matter Sources in CKG

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- Thus different matter sources may produce the same metric. The metric does not uniquely determine the energy-momentum tensor.
- Thus, if  $\tilde{T}_{ab} \& \tilde{T}'_{ab}$  differ by the Killing tensor  $\lambda g_{ab}/3$ , ( $\lambda$  constant) then  $T_{ab} \& T'_{ab}$  differ by a dark energy term  $\lambda g_{ab}$ . Hence dark energy does not gravitate in CKG.



• If a metric admits symmetries generated by Killing vectors  $\xi_I^a, I = 1 \dots N$ , then  $K_{ab} = kg_{ab} + \sum_{I=1}^N \sum_{J=1}^N k_{IJ}\xi_{I(a}\xi_{|J|b)}$ , where k and the  $k_{IJ}$  are constants, is a Killing tensor.



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- All known exact solutions in CKG admit Killing vectors: Static spherically symmetric 4KVs, pp-waves at least 1KV plane waves at least 5KVs, FRWL metrics 6KVs Einstein Universe 7KVs, Minkowski, de Sitter & ADS 10 KVs. Thus there are a plethora of possible sources in these cases.



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- Even with no symmetries there is the ubiquitous dark energy ambiguity.



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- Dark energy does not gravitate in CKG.
- Many solutions have multiple possible matter sources.



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- Thank you for listening if you have been! Questions?



## Bibliography

- Harada J (2023) Phys. Rev. D 108, 044031 arXiv:2308.07634
- Harada J (2023) Phys. Rev. D 108, 104037 arXiv:2308.02115
- Mantica J C & Molinari L G (2023) Phys. Rev. D 108, 124029 arXiv:2308.06803 [gr-qc]
- Tarciso Junior, J S S et al. (2023) arXiv:2310.19508 [gr-qc]
- Barnes A (2023) arXiv:2309.05336 & arXiv:2311.09171 [gr-qc]
- Barnes A (2024) Class. Quantum Grav. 41, 155007
- Clément G and Nouicer K (2024) arXiv:2404.00328 [gr-qc]
- Barnes A (2024) arXiv:2404.09310 [gr-qc]

