
Primordial **B**lack **H**ole formation from a nonspherical density profile with a misaligned deformation tensor

Chulmoon Yoo(Nagoya Univ.)

CY arXiv: 2403.11147

DYNAMICS OF PRIMORDIAL BLACK HOLE FORMATION II

@NAGOYA, JAPAN

Oct. 7 - 10

INVITED SPEAKERS

Gabriele Franciolini (CERN, Switzerland)

Marcos Flores (LPENS, France)

***John T. Giblin (Kenyon Coll.)**

Juan Carlos Hidalgo (Universidad Nacional Autonoma de Mexico)

Kazunori Kohri (NAOJ, KEK, KIPMU)

Eugene A. Lim (King's Coll. London)

Shi Pi (Chinese Academy of Sciences, China)

*:to be confirmed



***To access the participants and invited speakers list page, you will need the following password:**

dpbhf2

Nagoya?



Why you should attend this workshop

©Objectives

To take this opportunity to promote international mutual understanding, build cooperative relationships, and lead to future constructive and cutting-edge international joint research.

©It should be a good enough reason, but you may have additional reasons, e.g., ...

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Schedule

OCTOBER 2024						
Sun	Mon	Tue	Wed	Thu	Fri	Sat
29	30	1	2	3	4	5
6	7	8	9	10	11	12
	← workshop →					
13	14	15	16	17	18	19
20	21	22	23	24	25	26
	← COSMO'24@Kyoto →					
27	28	29	30	31	1	2

Main organizers

Albert Escrivà

(Nagoya)

Tomohiro Harada

(Rikkyo)

Yuichiro Tada

(Nagoya)

Takahiro Terada

(KMI, Nagoya),

Shuichiro Yokoyama

(KMI, Nagoya)

Chulmoon Yoo

(Nagoya)

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Primordial Black Hole

◎ Remnants of primordial non-linear inhomogeneity

◎ BHs not produced by late time stellar collapse

◎ Reliable formation scenario:

Collapse of rarely dense regions generated by quantum fluctuation during inflation
It's rare, but has a finite probability!!

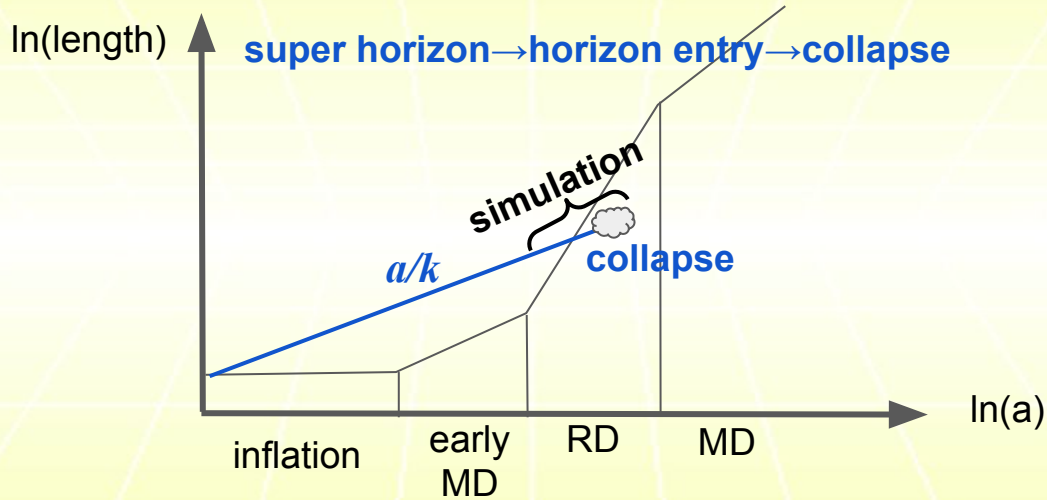
◎ If you accept inflation, you should be able to accept the **PBH** formation

◎ **PBH** is an appealing candidate for dark matter and BH binaries

- BHs “exist” in our universe
- BHs behave as dark matter in a cosmological scale
- Reliable scenario of **PBH** formation

PBH formation

© Comoving scale of an inhomogeneity $\sim 1/k$



© GR simulation starting from a **super-horizon non-linear initial data**

© We assume the linear equation of state $p=w\rho$ with $w=1/3$ and $1/5$

Super-horizon scale metric

◎ Leading order metric with gradient expansion

$$ds^2 = -dt^2 + a^2 \Psi^4 f_{ij} dx^i dx^j$$

◎ The next-leading order perturbations

$$\psi(t, x^k) = \Psi(x^k)(1 + \xi(t, x^k)), \quad \delta = \frac{\rho - \rho_b}{\rho_b}, \quad h_{ij} = \tilde{\gamma}_{ij} - f_{ij}, \quad \chi = \alpha - 1, \quad \kappa = \frac{K - K_b}{K_b}, \quad \tilde{A}_{ij}$$

$$\xi = -\frac{1}{6w+6} f \left(\frac{1}{aH_b} \right)^2 + \mathcal{O}(\epsilon^4)$$

$$\chi = -\frac{3w+1}{3w+3} f \left(\frac{1}{aH_b} \right)^2 + \mathcal{O}(\epsilon^4) \quad f(x^k) := -\frac{4}{3} \frac{\bar{\Delta}\Psi}{\Psi^5}$$

$$\delta = f \left(\frac{1}{aH_b} \right)^2 + \mathcal{O}(\epsilon^4) \quad p_{ij}(x^k) := \frac{1}{\Psi^4} \left[-\frac{2}{\Psi} (\bar{\mathcal{D}}_i \bar{\mathcal{D}}_j \Psi - \frac{1}{3} f_{ij} \bar{\Delta}\Psi) + \frac{6}{\Psi^2} (\bar{\mathcal{D}}_i \Psi \bar{\mathcal{D}}_j \Psi - \frac{1}{3} f_{ij} \bar{\mathcal{D}}^k \Psi \bar{\mathcal{D}}_k \Psi) \right]$$

$$h_{ij} = -\frac{4}{(3w+5)(3w-1)} p_{ij} \left(\frac{1}{aH_b} \right)^2 + \mathcal{O}(\epsilon^4)$$

$$\tilde{A}_{ij} = \frac{2}{3w+5} p_{ij} H_b \left(\frac{1}{aH_b} \right)^2 + \mathcal{O}(\epsilon^4)$$

$\Psi(x^k)$ fixes everything

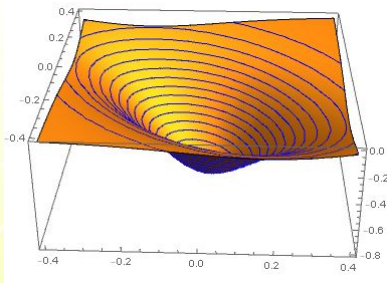
Initial condition

©Initial curvature perturbation

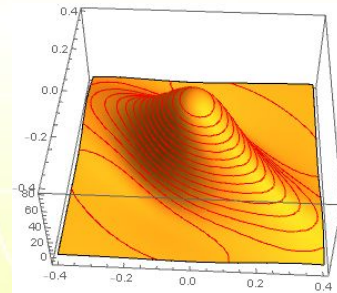
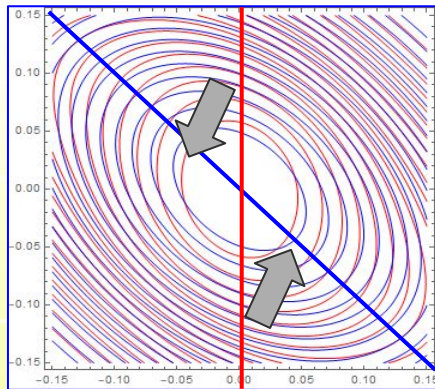
$$\frac{\zeta}{\mu} := -\frac{2}{\mu} \ln \Psi \simeq -1 + \frac{1}{2} (k_1^2 (x+y)^2 / 2 + k_2^2 (x-y)^2 / 2 + k_3^2 z^2) + \mathcal{O}(r^4)$$

$$\frac{\Delta \zeta}{\mu k^2} \simeq 1 - \frac{1}{2} (\kappa_1^2 x^2 + \kappa_2^2 y^2 + \kappa_3^2 z^2) + \mathcal{O}(r^4)$$

$\zeta \sim$ gravitational potential on (x,y) plane



$\Delta \zeta \sim$ energy density on (x,y) plane



tidal torque \Rightarrow angular momentum transfer \Rightarrow spinning PBH

Parameter setting

◎Initial curvature perturbation

$$\frac{\zeta}{\mu} \simeq -1 + \frac{1}{2}(k_1^2(x+y)^2/2 + k_2^2(x-y)^2/2 + k_3^2 z^2) + \mathcal{O}(r^4)$$

$$k_1^2 = \frac{1}{3}(\xi_1 + 3\xi_2 + \xi_3) \quad k_2^2 = \frac{1}{3}(\xi_1 - 3\xi_2 + \xi_3) \quad k_3^2 = \frac{1}{3}(\xi_1 - 2\xi_3)$$

$$\frac{\Delta\zeta}{\mu k^2} \simeq 1 - \frac{1}{2}(\kappa_1^2 x^2 + \kappa_2^2 y^2 + \kappa_3^2 z^2) + \mathcal{O}(r^4)$$

$$\kappa_1^2 = \frac{1}{3}(\tilde{\xi}_1 + 3\tilde{\xi}_2 + \tilde{\xi}_3) \quad \kappa_2^2 = \frac{1}{3}(\tilde{\xi}_1 - 3\tilde{\xi}_2 + \tilde{\xi}_3) \quad \kappa_3^2 = \frac{1}{3}(\tilde{\xi}_1 - 2\tilde{\xi}_3)$$

◎With reference to the peak theory

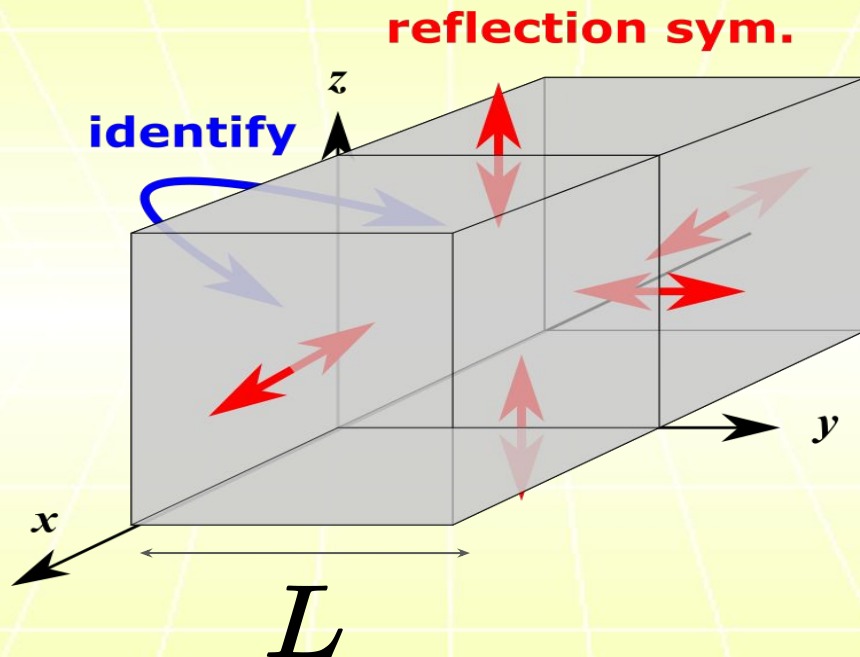
$$\vec{k}^2 = \vec{\kappa}^2 = \xi_1 = \tilde{\xi}_1 = 100/L^2$$

$$\xi_3 = \tilde{\xi}_3 = 0 \quad \text{The most probable values}$$

$$\xi_2 = 10/L^2 \quad \text{Too large to be statistically expected}$$

$$\tilde{\xi}_2 = 15/L^2 \quad \Rightarrow \text{Too much idealized for spin generation}$$

Numerical domain



About the numerical code

©Originally provided by Hirotada Okawa (for E-eqs and real scalar field w/ periodic BC)

©COSMOS(秋桜) code by C++ [CY, Hirotada Okawa, Ken-ichi Nakao(1306.1389),
Hirotada Okawa, Helvi Witek, Vitor Cardoso(1401.1548)]

©Basically follows the SACRA(桜) code by Fortran [Tetsuro Yamamoto, Masaru Shibata, Keisuke Taniguchi(arXiv:0806.4007)]

©Independently developed by CY and Hirotada Okawa

©In the CY side, it is mainly dedicated to PBH formation as follows

- **Inhomogeneous (scale-up) coordinate system** has been implemented [CY, Taishi Ikeda,Hirotada Okawa(arXiv:1811.00762)]
- **Fluid evolution code** has been implemented [CY, Tomohiro Harada,Hirotada Okawa(arXiv:2004.01042)]
- **1+1 code for spherical systems** has been developed based on COSMOS [CY, Harada, Hirano, Okawa, Sasaki(arXiv:2112.12335)]
- Recently, **a mesh refinement procedure** has been implemented

Summary of resolution difference

◎ Resolution in the previous simulation 2004.01042 CY, Tomohiro Harada, Hirotada Okawa

- Scale-up reference coordinates x^i related to the Cartesian coord. X^i by

$$X^i = x^i - \frac{S}{1+S} \frac{L}{\pi} \sin\left(\frac{\pi}{L} x^i\right) \text{ with } S = 15$$

- Resolution at the center ($\Delta x = L/100$)

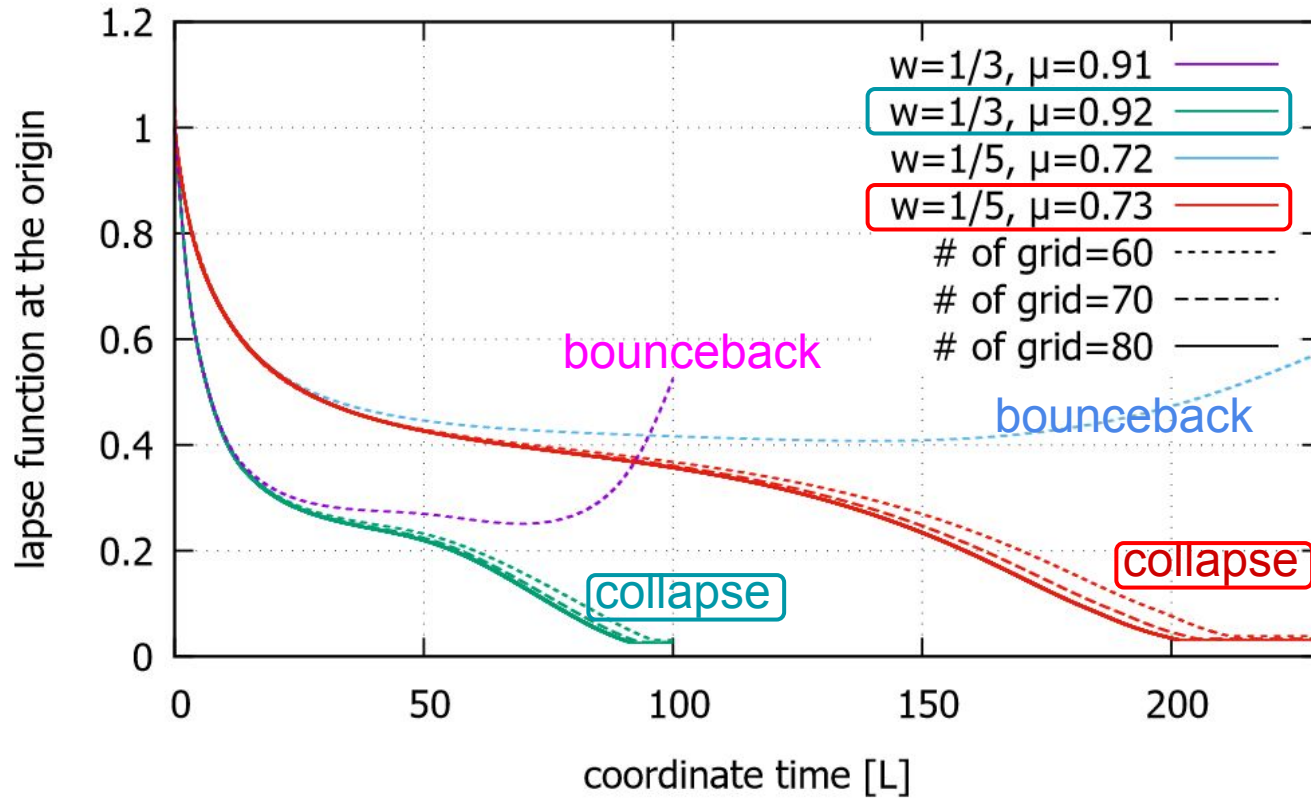
$$\Delta X|_{\text{center}} = \frac{1}{1+S} \Delta x = \frac{1}{16} \frac{L}{100} = \frac{L}{1600}$$

◎ New simulation with mesh refinement

- $S = 10, \Delta x = L/80$
- Two additional layers for the mesh refinement

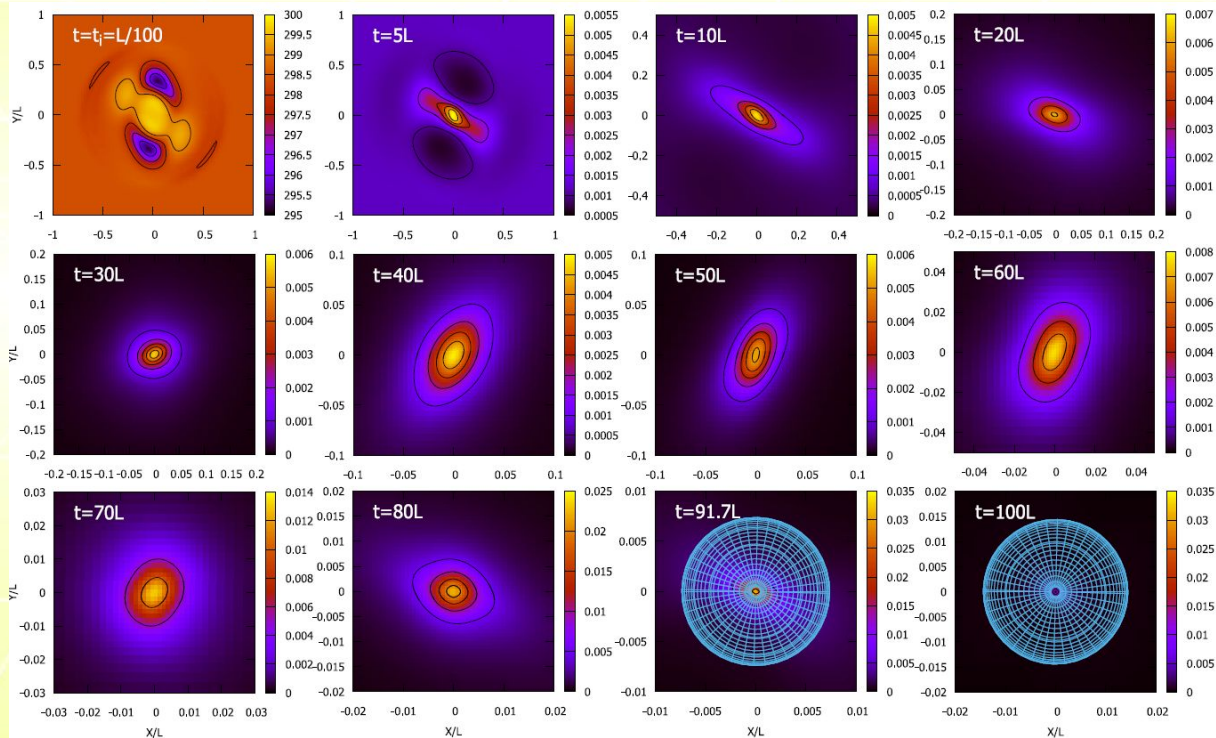
$$\Delta X|_{\text{center}} = \frac{1}{1+S} \times \frac{1}{2^2} \times \Delta x = \frac{1}{44} \times \frac{L}{80} = \frac{L}{3520}$$

Thresholds



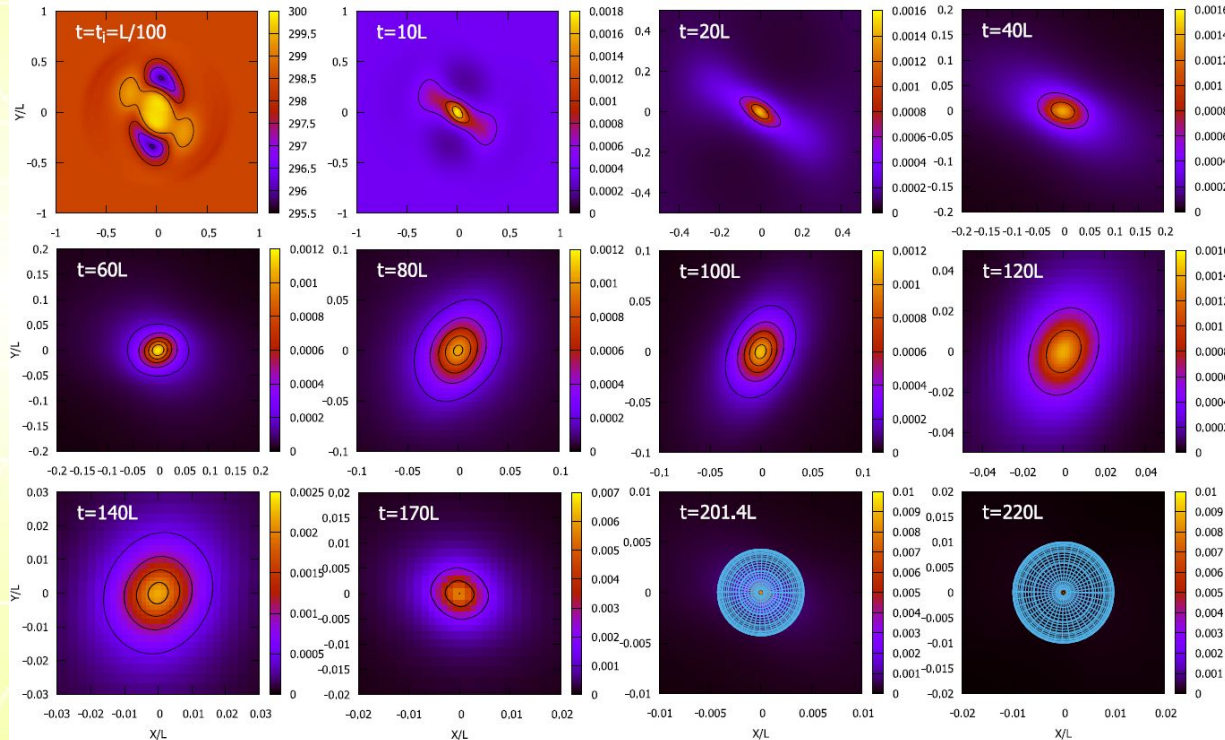
Effective spin parameter

⊙ $w=p/\rho=1/3$, amplitude $\mu=0.92$ case



Softer equation of state

© $p/\rho=w=1/5$, amplitude $\mu=0.73$ case



Effective spin parameter

©Kerr black hole case

$$\text{area : } A_{\text{Kerr}} = 8\pi(M^2 + \sqrt{M^4 - a^2 M^2})$$

$$\text{equatorial circumference : } d_{\text{Kerr}} = 4\pi M$$

$$\Rightarrow s_{\text{Kerr}}^2 := \frac{a^2}{M^2} = \frac{4\pi A_{\text{Kerr}}(d_{\text{Kerr}}^2 - \pi A_{\text{Kerr}})}{d_{\text{Kerr}}^4}$$

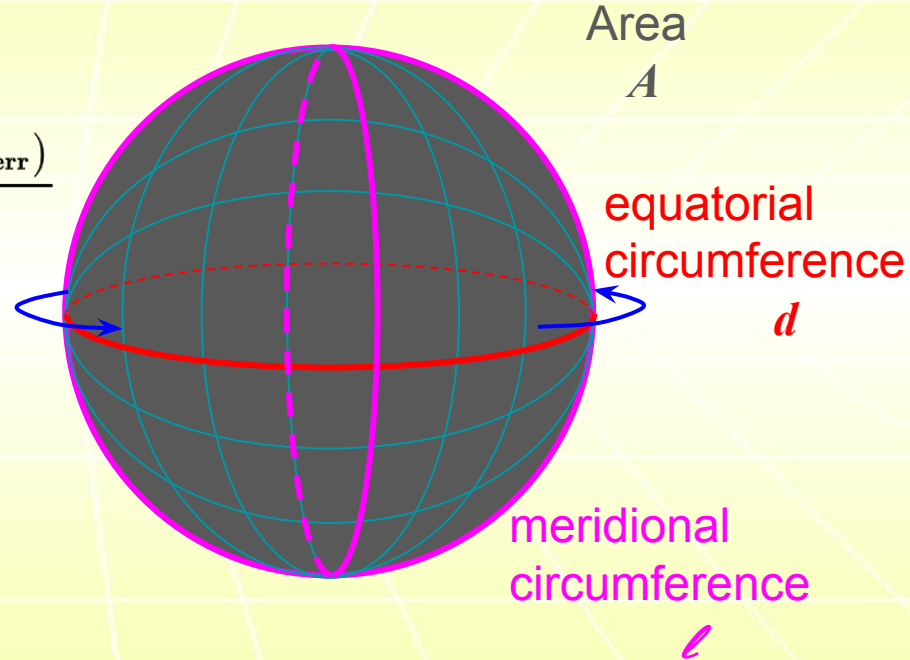
©Effective dimensionless spin

$$s^2 := \frac{4\pi A(d^2 - \pi A)}{d^4}$$

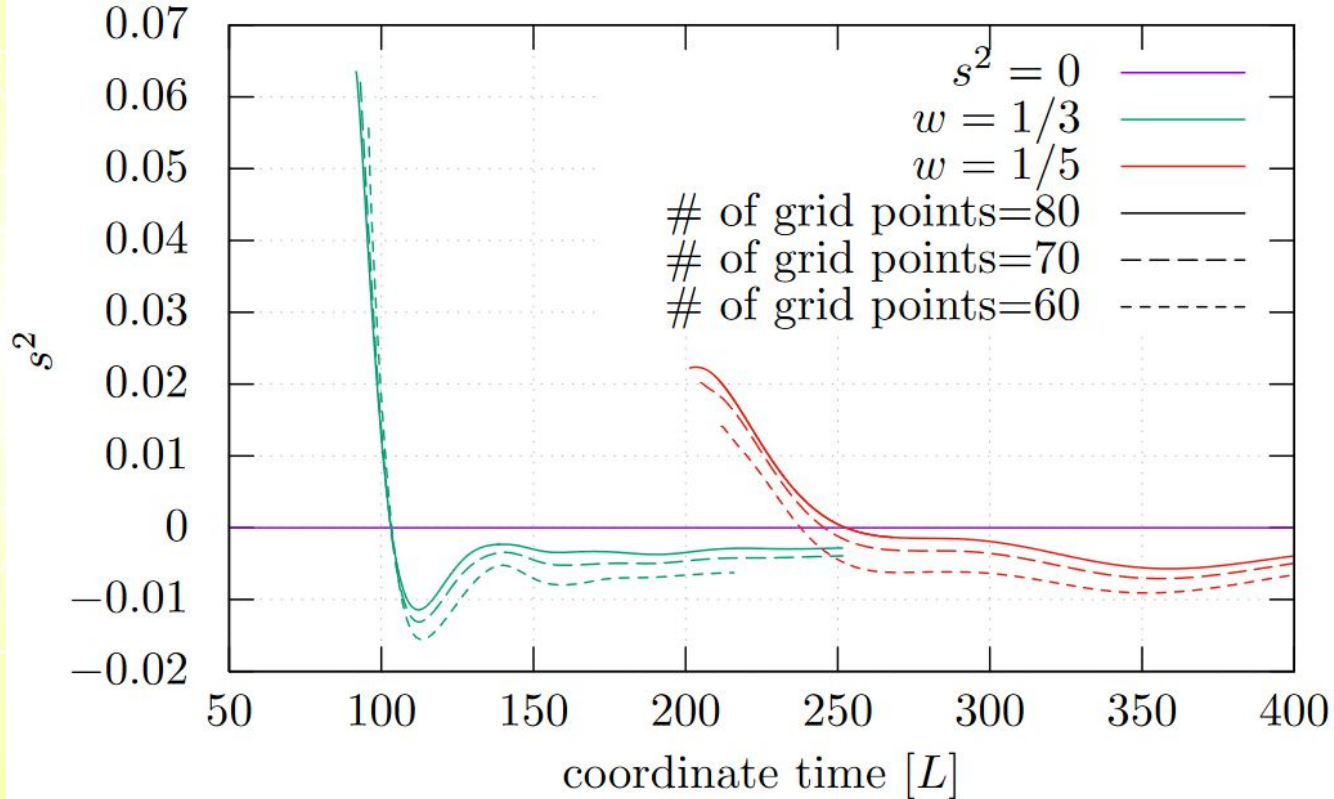
©Asphericity parameters

$$\text{meridional circumference : } \ell_{x=0, l_y=0}$$

$$\alpha := \frac{\ell_{x=0}}{d} \quad \beta := \frac{\ell_{y=0}}{d}$$



Effective spin parameter



Effective spin parameter

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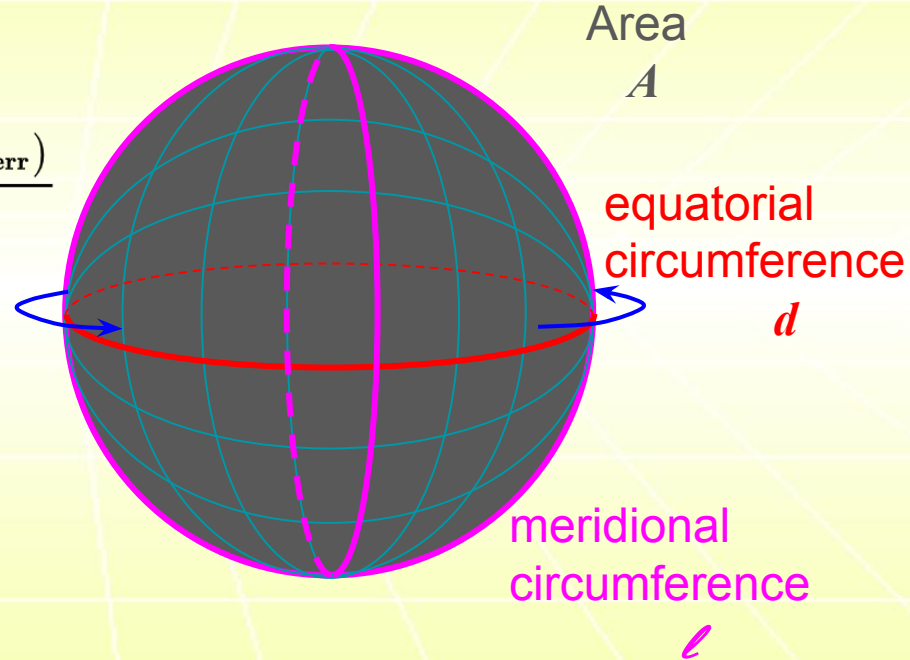
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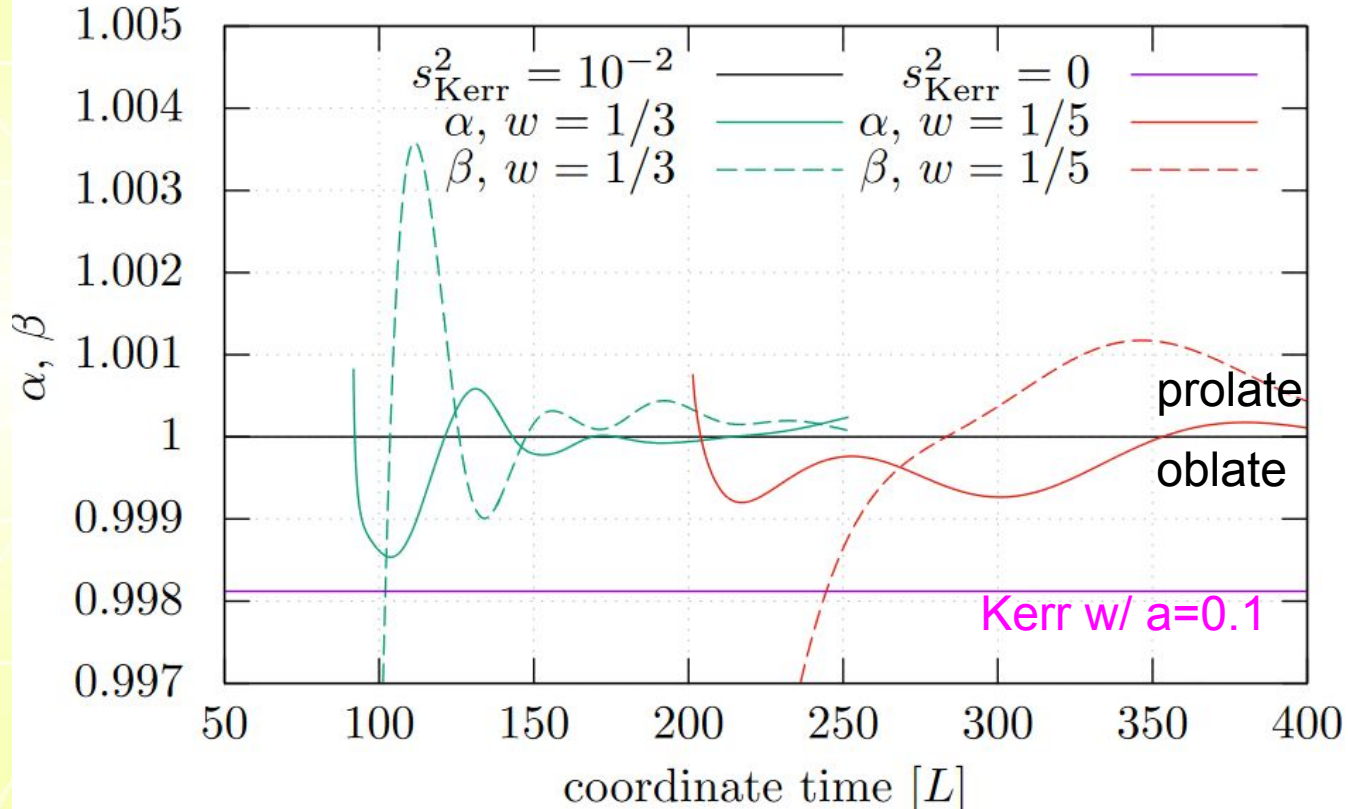
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Non-sphericity



Spin of **PBH** is very small

for the equation of states $p=w\rho$ with $w>1/5$

©Caveats

- Results only for a specific initial profile
- The case in which the amplitude is very close to the critical value
- etc.

Thank you for your attention