

Symmetry reduction of gravitational Lagrangians

based on: G. Frausto, I. Kolář, TM, Ch. Torre, (soon on arXiv)

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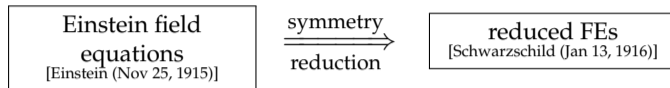
Outline

- 1 Motivation: Weyl trick
- 2 Rigorous treatment: Principle of symmetric criticality
- 3 Systematic study
- 4 Examples

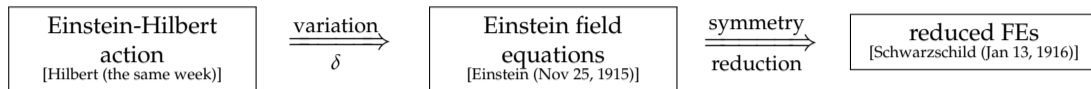
Weyl trick to derive the Schwarzschild solution [Weyl (1917)]

Einstein field
equations
[Einstein (Nov 25, 1915)]

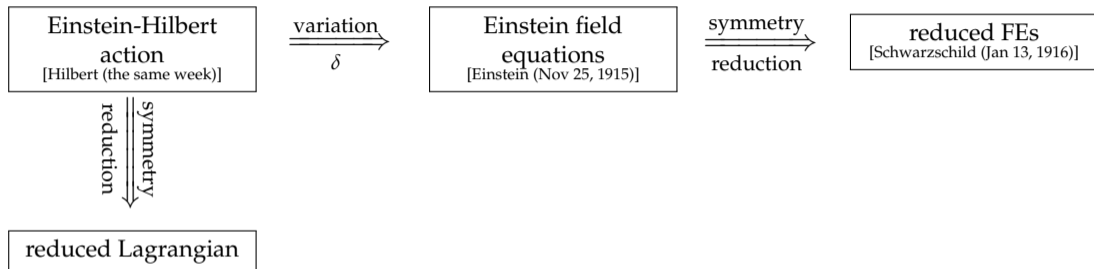
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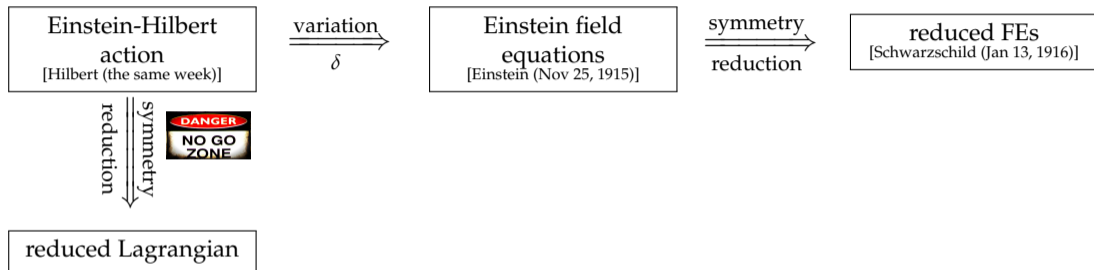
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1 symmetry reduction of Lagrangian

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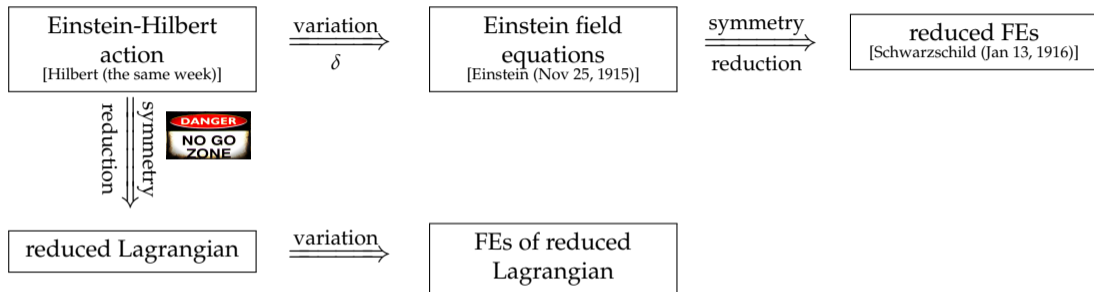
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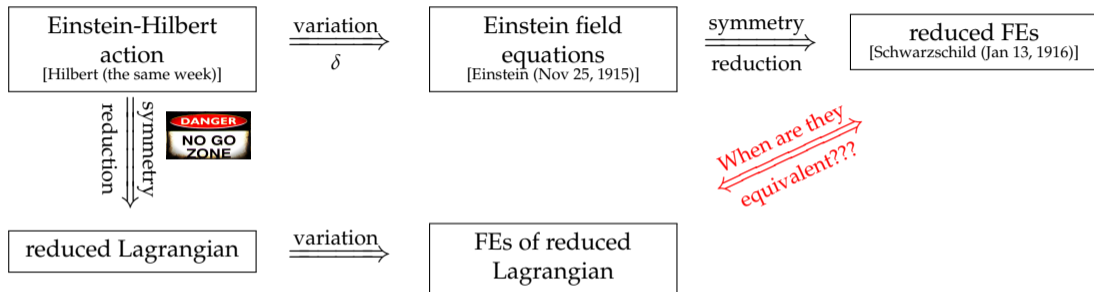
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$$b' = 0 \quad \text{and} \quad (r(a-1))' = 0 \quad \implies \quad a = 1 - \frac{c_1}{r}, \quad b = c_2$$

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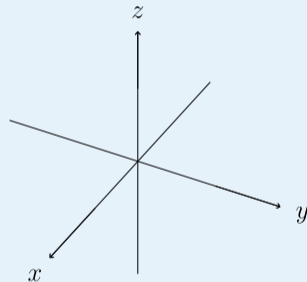
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Isometries

Infinitesimal group action Γ on M

given by d -dim Lie algebra of isometry generators $X \in \Gamma$ (Killing vectors)

Example: symmetries of S^2

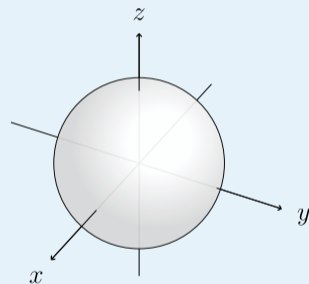


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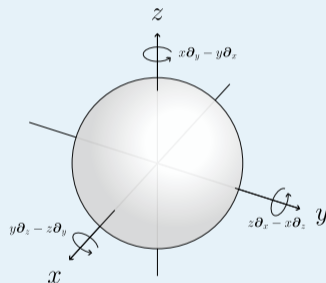


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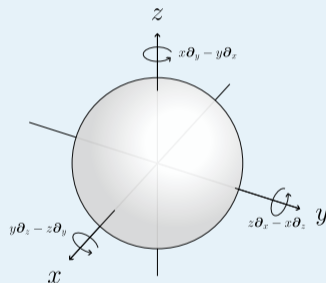
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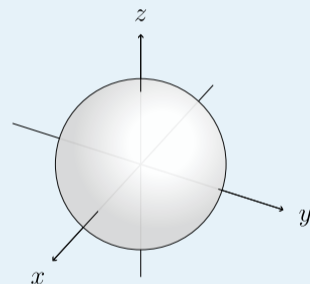
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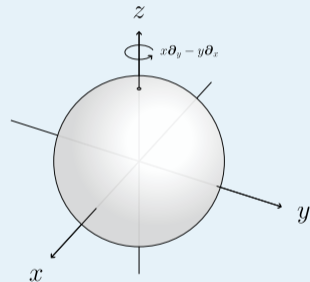
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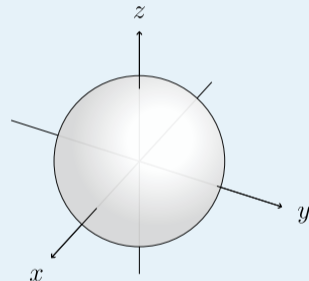
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relation to dim p of Γ_x : $l = d - p$

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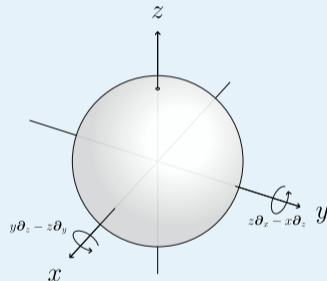
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 orbit of $(0,0,1)$ is unit sphere

Symmetry reduction of Lagrangian

Purely gravitational theory on a 4-dimensional spacetime

$$S = \int_{\mathcal{M}} \underline{\epsilon}(\mathbf{g}) L[\mathbf{g}]$$

- Levi-Civita tensor $\underline{\epsilon}(\mathbf{g})$ defines the volume element $\sqrt{-g}d^4x$
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- accomplished by contraction with Γ -invariant antisymmetric l -chain $\chi = \chi^{i_1 \dots i_l} \mathbf{X}_{i_1} \cdots \mathbf{X}_{i_l}$

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- reduced Lagrangian $\hat{\underline{L}} = \chi \bullet \underline{L}[\hat{\mathbf{g}}] = \hat{\underline{\epsilon}}(\hat{\mathbf{g}})L[\hat{\mathbf{g}}]$, (where $\hat{\underline{\epsilon}}(\hat{\mathbf{g}}) = \chi \bullet \underline{\epsilon}(\hat{\mathbf{g}})$)

Principle of symmetric criticality

Variation of Lagrangian 4-form

$$\delta \underline{L} = \underline{E}(\underline{L}) \cdot \delta \mathbf{g} + \underline{d}\underline{\eta}(\delta \mathbf{g})$$

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Principle of symmetric criticality [Palais (1979), M. E. Fels, C. G. Torre (2002)]

Variation of Lagrangian commutes with symmetry reduction for all possible theories:

$$\forall \underline{L} : \underline{E}(\underline{L})[\hat{\mathbf{g}}] = 0 \iff \underline{E}(\hat{\underline{L}})[\hat{\mathbf{g}}] = 0$$

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Two conditions imposed solely on Γ are necessary and sufficient for validity of PSC.

PSC1 “Lie algebra condition”

- PSC1 ensures that the reduction of the boundary term $\underline{d}\eta$ is a boundary term $\underline{d}\hat{\eta}$ for the reduced Lagrangian

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- if PSC1 satisfied then Euler-Lagrange equations of the reduced Lagrangian always yield at least a subset of the reduced equations

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Let S_x and S_x^* denote the vector space of Γ_x -invariant $\binom{0}{2}$ and $\binom{2}{0}$ tensors at x , respectively. Denote by V_x^0 the vector space of $\binom{2}{0}$ tensors which have a vanishing scalar contraction with all elements of S_x . Then in the neighborhood of x :

$$S_x^* \cap V_x^0 = \{0\}$$

i.e. there is no Γ_x -invariant $\binom{2}{0}$ tensor that contracts to zero with all Γ_x -invariant $\binom{0}{2}$ tensors

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- PSC2 satisfied iff the isotropy algebra contains no null-rotation subalgebra

Classification of infinitesimal group actions

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Hicks classification [Hicks, Ph.D. thesis (2016)]

based on classifying isometry algebra and isotropy subalgebra pairs (Γ, Γ_x)

- isotropy subalgebras Γ_x can be identified with subalgebras of the Lorentz algebra
- cases denoted by $[d, l, c]$
 - 1 d is dim of Γ
 - 2 l is dim of orbits ($l = d - p$)
 - 3 c enumerates possible cases of given dimensions
- explicit infinitesimal generators given for each case

PSC-compatible infinitesimal group actions

Hicks #	PSC1	PSC2	PSC
3,2,1	✓	✓	✓
3,2,2	✓	✓	✓
3,2,3	✓	✓	✓
3,2,4	✓	✓	✓
3,2,5	✓	✓	✓
3,3,1	✗	✓	✗
3,3,2	✓	✓	✓
3,3,3	✓	✓	✓
3,3,4	✗	✓	✗
3,3,5	✗	✓	✗
3,3,6	✗	✓	✗
3,3,7	✗	✓	✗
3,3,8	✓	✓	✓
3,3,9	✓	✓	✓
4,3,1	✓	✓	✓
4,3,2	✓	✓	✓
4,3,3	✓	✓	✓
4,3,4	✓	✓	✓
4,3,5	✓	✓	✓
4,3,6	✓	✓	✓
4,3,7	✗	✓	✗
4,3,8	✓	✓	✓
4,3,9	✓	✓	✓

Hicks #	PSC1	PSC2	PSC
4,3,10	✓	✓	✓
4,3,11	✓	✓	✓
4,3,12	✗	✓	✗
4,3,13	✓	✗	✗
4,3,14	✓	✗	✗
4,3,15	✓	✗	✗
4,3,16	✓	✗	✗
4,3,17	✓	✗	✗
4,3,18	✓	✗	✗
4,3,19	✓	✗	✗
4,3,20	✓	✗	✗
4,4,1	✓	✓	✓
4,4,2	✓	✓	✓
4,4,3	✗	✓	✗
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4,4,6	✗	✓	✗
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4,4,8	✗	✓	✗
4,4,9	✓	✓	✓
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4,4,11	✗	✓	✗
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4,4,17	✗	✓	✗
4,4,18	✓	✓	✓
4,4,19	✗	✓	✗
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4,4,21	✗	✓	✗
4,4,22	✓	✓	✓
4,4,23	✗	✓	✗
5,4,-1	✓	✗	✗
5,4,-2	✗	✗	✗
5,4,-3	✗	✗	✗
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5,4,-6	✗	✗	✗
5,4,1	✓	✓	✓
5,4,2	✓	✓	✓
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5,4,11	✓	✗	✗
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6,3,2	✓	✓	✓
6,3,3	✓	✓	✓
6,3,4	✓	✓	✓
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3,3,5	✗	✓	✗
3,3,6	✗	✓	✗
3,3,7	✗	✓	✗
3,3,8	✓	✓	✓
3,3,9	✓	✓	✓
4,3,1	✓	✓	✓
4,3,2	✓	✓	✓
4,3,3	✓	✓	✓
4,3,4	✓	✓	✓
4,3,5	✓	✓	✓
4,3,6	✓	✓	✓
4,3,7	✗	✓	✗
4,3,8	✓	✓	✓
4,3,9	✓	✓	✓

Hicks #	PSC1	PSC2	PSC
4,3,10	✓	✓	✓
4,3,11	✓	✓	✓
4,3,12	✗	✓	✗
4,3,13	✓	✗	✗
4,3,14	✓	✗	✗
4,3,15	✓	✗	✗
4,3,16	✓	✗	✗
4,3,17	✓	✗	✗
4,3,18	✓	✗	✗
4,3,19	✓	✗	✗
4,3,20	✓	✗	✗
4,4,1	✓	✓	✓
4,4,2	✓	✓	✓
4,4,3	✗	✓	✗
4,4,4	✗	✓	✗
4,4,5	✗	✓	✗
4,4,6	✗	✓	✗
4,4,7	✗	✓	✗
4,4,8	✗	✓	✗
4,4,9	✓	✓	✓
4,4,10	✗	✓	✗
4,4,11	✗	✓	✗
4,4,12	✗	✓	✗

Hicks #	PSC1	PSC2	PSC
4,4,13	✗	✓	✗
4,4,14	✗	✓	✗
4,4,15	✗	✓	✗
4,4,16	✗	✓	✗
4,4,17	✗	✓	✗
4,4,18	✓	✓	✓
4,4,19	✗	✓	✗
4,4,20	✗	✓	✗
4,4,21	✗	✓	✗
4,4,22	✓	✓	✓
4,4,23	✗	✓	✗
5,4,-1	✓	✗	✗
5,4,-2	✗	✗	✗
5,4,-3	✗	✗	✗
5,4,-4	✗	✗	✗
5,4,-5	✗	✗	✗
5,4,-6	✗	✗	✗
5,4,1	✓	✓	✓
5,4,2	✓	✓	✓
5,4,3	✓	✓	✓
5,4,4	✗	✓	✗
5,4,5	✗	✓	✗
5,4,6	✓	✓	✓

Hicks #	PSC1	PSC2	PSC
5,4,7	✓	✓	✓
5,4,8	✗	✓	✗
5,4,9	✗	✓	✗
5,4,10	✗	✗	✗
5,4,11	✓	✗	✗
6,3,1	✓	✓	✓
6,3,2	✓	✓	✓
6,3,3	✓	✓	✓
6,3,4	✓	✓	✓
6,3,5	✓	✓	✓
6,3,6	✓	✓	✓
6,4,-1	✓	✗	✗
6,4,1	✓	✓	✓
6,4,2	✓	✓	✓
6,4,3	✓	✓	✓
6,4,4	✓	✓	✓
6,4,5	✓	✓	✓
6,4,6	✓	✗	✗
7,4,1	✓	✓	✓
7,4,2	✓	✓	✓
7,4,3	✓	✓	✓
7,4,4	✓	✓	✓
7,4,5	✓	✗	✗

- number of cases

total	PSC1	PSC2	PSC
92	57	71	44

PSC-compatible infinitesimal group actions

Hicks #	PSC1	PSC2	PSC
3,2,1	✓	✓	✓
3,2,2	✓	✓	✓
3,2,3	✓	✓	✓
3,2,4	✓	✓	✓
3,2,5	✓	✓	✓
3,3,1	✗	✓	✗
3,3,2	✓	✓	✓
3,3,3	✓	✓	✓
3,3,4	✗	✓	✗
3,3,5	✗	✓	✗
3,3,6	✗	✓	✗
3,3,7	✗	✓	✗
3,3,8	✓	✓	✓
3,3,9	✓	✓	✓
4,3,1	✓	✓	✓
4,3,2	✓	✓	✓
4,3,3	✓	✓	✓
4,3,4	✓	✓	✓
4,3,5	✓	✓	✓
4,3,6	✓	✓	✓
4,3,7	✗	✓	✗
4,3,8	✓	✓	✓
4,3,9	✓	✓	✓

Hicks #	PSC1	PSC2	PSC
4,3,10	✓	✓	✓
4,3,11	✓	✓	✓
4,3,12	✗	✓	✗
4,3,13	✓	✗	✗
4,3,14	✓	✗	✗
4,3,15	✓	✗	✗
4,3,16	✓	✗	✗
4,3,17	✓	✗	✗
4,3,18	✓	✗	✗
4,3,19	✓	✗	✗
4,3,20	✓	✗	✗
4,4,1	✓	✓	✓
4,4,2	✓	✓	✓
4,4,3	✗	✓	✗
4,4,4	✗	✓	✗
4,4,5	✗	✓	✗
4,4,6	✗	✓	✗
4,4,7	✗	✓	✗
4,4,8	✗	✓	✗
4,4,9	✓	✓	✓
4,4,10	✗	✓	✗
4,4,11	✗	✓	✗
4,4,12	✗	✓	✗

Hicks #	PSC1	PSC2	PSC
4,4,13	✗	✓	✗
4,4,14	✗	✓	✗
4,4,15	✗	✓	✗
4,4,16	✗	✓	✗
4,4,17	✗	✓	✗
4,4,18	✓	✓	✓
4,4,19	✗	✓	✗
4,4,20	✗	✓	✗
4,4,21	✗	✓	✗
4,4,22	✓	✓	✓
4,4,23	✗	✓	✗
5,4,-1	✓	✗	✗
5,4,-2	✗	✗	✗
5,4,-3	✗	✗	✗
5,4,-4	✗	✗	✗
5,4,-5	✗	✗	✗
5,4,-6	✗	✗	✗
5,4,1	✓	✓	✓
5,4,2	✓	✓	✓
5,4,3	✓	✓	✓
5,4,4	✗	✓	✗
5,4,5	✗	✓	✗
5,4,6	✓	✓	✓

Hicks #	PSC1	PSC2	PSC
5,4,7	✓	✓	✓
5,4,8	✗	✓	✗
5,4,9	✗	✓	✗
5,4,10	✗	✗	✗
5,4,11	✓	✗	✗
6,3,1	✓	✓	✓
6,3,2	✓	✓	✓
6,3,3	✓	✓	✓
6,3,4	✓	✓	✓
6,3,5	✓	✓	✓
6,3,6	✓	✓	✓
6,4,-1	✓	✗	✗
6,4,1	✓	✓	✓
6,4,2	✓	✓	✓
6,4,3	✓	✓	✓
6,4,4	✓	✓	✓
6,4,5	✓	✓	✓
6,4,6	✓	✗	✗
7,4,1	✓	✓	✓
7,4,2	✓	✓	✓
7,4,3	✓	✓	✓
7,4,4	✓	✓	✓
7,4,5	✓	✗	✗

- number of cases

total	PSC1	PSC2	PSC
92	57	71	44

- only 42 qualitatively different cases (the answer to the ultimate question of life, the universe, and everything.)



PSC-compatible infinitesimal group actions

Hicks #	PSC1	PSC2	PSC
[3,2,1]	✓	✓	✓
[3,2,2]	✓	✓	✓
[3,2,3]	✓	✓	✓
[3,2,4]	✓	✓	✓
[3,2,5]	✓	✓	✓
[3,3,1]	✗	✓	✗
[3,3,2]	✓	✓	✓
[3,3,3]	✓	✓	✓
[3,3,4]	✗	✓	✗
[3,3,5]	✗	✓	✗
[3,3,6]	✗	✓	✗
[3,3,7]	✗	✓	✗
[3,3,8]	✓	✓	✓
[3,3,9]	✓	✓	✓
[4,3,1]	✓	✓	✓
[4,3,2]	✓	✓	✓
[4,3,3]	✓	✓	✓
[4,3,4]	✓	✓	✓
[4,3,5]	✓	✓	✓
[4,3,6]	✓	✓	✓
[4,3,7]	✗	✓	✗
[4,3,8]	✓	✓	✓
[4,3,9]	✓	✓	✓

Hicks #	PSC1	PSC2	PSC
[4,3,10]	✓	✓	✓
[4,3,11]	✓	✓	✓
[4,3,12]	✗	✓	✗
[4,3,13]	✓	✗	✗
[4,3,14]	✓	✗	✗
[4,3,15]	✓	✗	✗
[4,3,16]	✓	✗	✗
[4,3,17]	✓	✗	✗
[4,3,18]	✓	✗	✗
[4,3,19]	✓	✗	✗
[4,3,20]	✓	✗	✗
[4,4,1]	✓	✓	✓
[4,4,2]	✓	✓	✓
[4,4,3]	✗	✓	✗
[4,4,4]	✗	✓	✗
[4,4,5]	✗	✓	✗
[4,4,6]	✗	✓	✗
[4,4,7]	✗	✓	✗
[4,4,8]	✗	✓	✗
[4,4,9]	✓	✓	✓
[4,4,10]	✗	✓	✗
[4,4,11]	✗	✓	✗
[4,4,12]	✗	✓	✗

Hicks #	PSC1	PSC2	PSC
[4,4,13]	✗	✓	✗
[4,4,14]	✗	✓	✗
[4,4,15]	✗	✓	✗
[4,4,16]	✗	✓	✗
[4,4,17]	✗	✓	✗
[4,4,18]	✓	✓	✓
[4,4,19]	✗	✓	✗
[4,4,20]	✗	✓	✗
[4,4,21]	✗	✓	✗
[4,4,22]	✓	✓	✓
[4,4,23]	✗	✓	✗
[5,4,-1]	✓	✗	✗
[5,4,-2]	✗	✗	✗
[5,4,-3]	✗	✗	✗
[5,4,-4]	✗	✗	✗
[5,4,-5]	✗	✗	✗
[5,4,-6]	✗	✗	✗
[5,4,1]	✓	✓	✓
[5,4,2]	✓	✓	✓
[5,4,3]	✓	✓	✓
[5,4,4]	✗	✓	✗
[5,4,5]	✗	✓	✗
[5,4,6]	✓	✓	✓

Hicks #	PSC1	PSC2	PSC
[5,4,7]	✓	✓	✓
[5,4,8]	✗	✓	✗
[5,4,9]	✗	✓	✗
[5,4,10]	✗	✗	✗
[5,4,11]	✓	✗	✗
[6,3,1]	✓	✓	✓
[6,3,2]	✓	✓	✓
[6,3,3]	✓	✓	✓
[6,3,4]	✓	✓	✓
[6,3,5]	✓	✓	✓
[6,3,6]	✓	✓	✓
[6,4,-1]	✓	✗	✗
[6,4,1]	✓	✓	✓
[6,4,2]	✓	✓	✓
[6,4,3]	✓	✓	✓
[6,4,4]	✓	✓	✓
[6,4,5]	✓	✓	✓
[6,4,6]	✓	✗	✗
[7,4,1]	✓	✓	✓
[7,4,2]	✓	✓	✓
[7,4,3]	✓	✓	✓
[7,4,4]	✓	✓	✓
[7,4,5]	✓	✗	✗

- number of cases

total	PSC1	PSC2	PSC
92	57	71	44

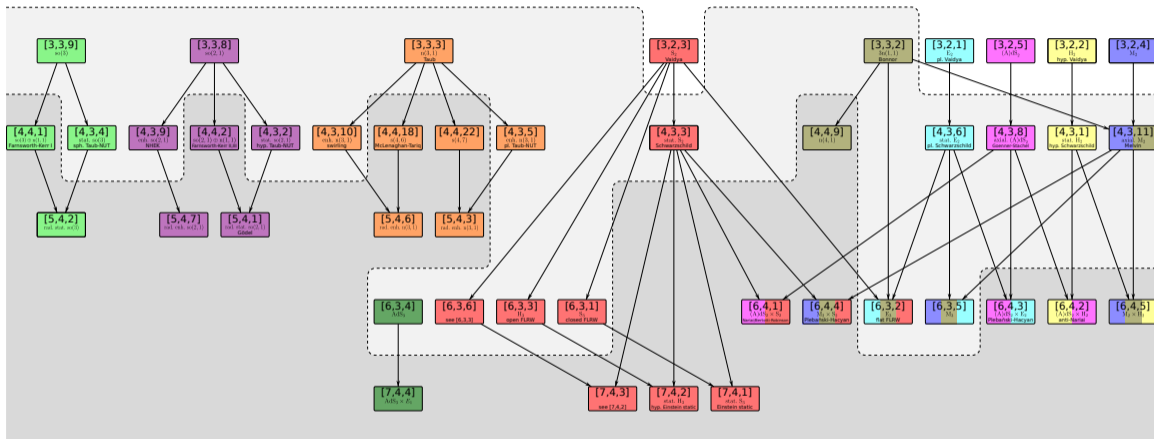
- only 42 qualitatively different cases (the answer to the ultimate question of life, the universe, and everything.)



- for each PSC-compatible Γ we determined the corresponding l -chains χ and Γ -invariant metrics \hat{g} in adapted coordinates

$$\hat{g} = \sum_{i=1}^s \phi_i q_i, \text{ where } s = \begin{cases} 2, & \text{for } [6, 3, \star], [6, 4, \star], [7, 4, \star] \\ 4, & \text{for } [3, 2, \star], [4, 3, \star], [5, 4, \star] \text{ and } \phi_i = \phi_i(x_1, x_2, \dots, x_{(4-i)}) \\ 10, & \text{for } [3, 3, \star], [4, 4, \star] \end{cases}$$

Relations among PSC-compatible infinitesimal group actions



$$\hat{g} = \sum_{i=1}^s \phi_i q_i, \text{ where } s = \begin{cases} 2, & \text{for } [6,3,*], [6,4,*], [7,4,*] \\ 4, & \text{for } [3,2,*], [4,3,*], [5,4,*] \text{ and } \phi_i = \phi_i(x_1, x_2, \dots, x_{(4-l)}) \\ 10, & \text{for } [3,3,*], [4,4,*] \end{cases}$$

Weyl trick revisited ([4,3,3]: stationary S_2)

- infinitesimal group action

$$\Gamma = \text{span}\{\cos \varphi \partial_{\vartheta} - \cot \vartheta \sin \varphi \partial_{\varphi}, \sin \varphi \partial_{\vartheta} + \cot \vartheta \cos \varphi \partial_{\varphi}, \partial_{\varphi}, \partial_t\}$$

Weyl trick revisited ([4,3,3]: stationary S_2)

- infinitesimal group action

$$\Gamma = \text{span}\{\cos \varphi \partial_\vartheta - \cot \vartheta \sin \varphi \partial_\varphi, \sin \varphi \partial_\vartheta + \cot \vartheta \cos \varphi \partial_\varphi, \partial_\varphi, \partial_t\}$$

- Γ -invariant metric

$$\hat{g} = -\phi_1(r) \mathbf{d}t^2 + \phi_2(r) (\mathbf{d}t \vee \mathbf{d}r) + \phi_3(r) \mathbf{d}r^2 + \phi_4(r) (\mathbf{d}\vartheta^2 + \sin^2 \vartheta \mathbf{d}\varphi^2)$$

Weyl trick revisited ([4,3,3]: stationary S_2)

- infinitesimal group action

$$\Gamma = \text{span}\{\cos \varphi \partial_\vartheta - \cot \vartheta \sin \varphi \partial_\varphi, \sin \varphi \partial_\vartheta + \cot \vartheta \cos \varphi \partial_\varphi, \partial_\varphi, \partial_t\}$$

- Γ -invariant metric

$$\hat{g} = -\phi_1(r) \mathbf{d}t^2 + \phi_2(r) (\mathbf{d}t \vee \mathbf{d}r) + \phi_3(r) \mathbf{d}r^2 + \phi_4(r) (\mathbf{d}\vartheta^2 + \sin^2 \vartheta \mathbf{d}\varphi^2)$$

- residual gauge freedom $t \rightarrow t + A(r), r \rightarrow B(r)$ do not invalidate PSC and allows us to fix $\phi_2 = 0$ and $\phi_4 = r^2$

Weyl trick revisited ([4,3,3]: stationary S_2)

- infinitesimal group action

$$\Gamma = \text{span}\{\cos \varphi \partial_\vartheta - \cot \vartheta \sin \varphi \partial_\varphi, \sin \varphi \partial_\vartheta + \cot \vartheta \cos \varphi \partial_\varphi, \partial_\varphi, \partial_t\}$$

- Γ -invariant metric

$$\hat{g} = -\phi_1(r) \mathbf{d}t^2 + \phi_2(r) (\mathbf{d}t \vee \mathbf{d}r) + \phi_3(r) \mathbf{d}r^2 + \phi_4(r) (\mathbf{d}\vartheta^2 + \sin^2 \vartheta \mathbf{d}\varphi^2)$$

- residual gauge freedom $t \rightarrow t + A(r), r \rightarrow B(r)$ do not invalidate PSC and allows us to fix $\phi_2 = 0$ and $\phi_4 = r^2$
- l -chain

$$\chi = \csc \vartheta \partial_t \wedge \partial_\vartheta \wedge \partial_\varphi$$

Weyl trick revisited ([4,3,3]: stationary S_2)

- infinitesimal group action

$$\Gamma = \text{span}\{\cos \vartheta \partial_\vartheta - \cot \vartheta \sin \varphi \partial_\varphi, \sin \varphi \partial_\vartheta + \cot \vartheta \cos \varphi \partial_\varphi, \partial_\varphi, \partial_t\}$$

- Γ -invariant metric

$$\hat{g} = -\phi_1(r) \mathbf{d}t^2 + \phi_2(r) (\mathbf{d}t \vee \mathbf{d}r) + \phi_3(r) \mathbf{d}r^2 + \phi_4(r) (\mathbf{d}\vartheta^2 + \sin^2 \vartheta \mathbf{d}\varphi^2)$$

- residual gauge freedom $t \rightarrow t + A(r), r \rightarrow B(r)$ do not invalidate PSC and allows us to fix $\phi_2 = 0$ and $\phi_4 = r^2$

- l -chain

$$\chi = \csc \vartheta \partial_t \wedge \partial_\vartheta \wedge \partial_\varphi$$

- Levi-Civita tensor

$$\underline{\epsilon}(\hat{g}) = r^2 \sin \vartheta \sqrt{\phi_1 \phi_3} \mathbf{d}t \wedge \mathbf{d}r \wedge \mathbf{d}\vartheta \wedge \mathbf{d}\varphi$$

Weyl trick revisited ([4,3,3]: stationary S_2)

- infinitesimal group action

$$\Gamma = \text{span}\{\cos \vartheta \partial_\vartheta - \cot \vartheta \sin \varphi \partial_\varphi, \sin \varphi \partial_\vartheta + \cot \vartheta \cos \varphi \partial_\varphi, \partial_\varphi, \partial_t\}$$

- Γ -invariant metric

$$\hat{g} = -\phi_1(r) \mathbf{d}t^2 + \phi_2(r) (\mathbf{d}t \vee \mathbf{d}r) + \phi_3(r) \mathbf{d}r^2 + \phi_4(r) (\mathbf{d}\vartheta^2 + \sin^2 \vartheta \mathbf{d}\varphi^2)$$

- residual gauge freedom $t \rightarrow t + A(r), r \rightarrow B(r)$ do not invalidate PSC and allows us to fix $\phi_2 = 0$ and $\phi_4 = r^2$

- l -chain

$$\chi = \csc \vartheta \partial_t \wedge \partial_\vartheta \wedge \partial_\varphi$$

- Levi-Civita tensor

$$\underline{\epsilon}(\hat{g}) = r^2 \sin \vartheta \sqrt{\phi_1 \phi_3} \mathbf{d}t \wedge \mathbf{d}r \wedge \mathbf{d}\vartheta \wedge \mathbf{d}\varphi$$

- reduced Lagrangian 1-form

$$\underline{\hat{L}} = \chi \bullet \underline{\epsilon}(\hat{g}) L[\hat{g}] = r^2 \sqrt{\phi_1 \phi_3} L[\hat{g}] \mathbf{d}r$$

Symmetry reduction for flat FLRW cosmologies ([6,3,2]: E_3)

- infinitesimal group action

$$\Gamma = \text{span}\{\partial_x, \partial_y, \partial_z, x\partial_y - y\partial_x, y\partial_z - z\partial_y, z\partial_x - x\partial_z\}$$

Symmetry reduction for flat FLRW cosmologies ([6,3,2]: E_3)

- infinitesimal group action

$$\Gamma = \text{span}\{\partial_x, \partial_y, \partial_z, x\partial_y - y\partial_x, y\partial_z - z\partial_y, z\partial_x - x\partial_z\}$$

- Γ -invariant metric

$$\hat{g} = -\phi_1(t) dt^2 + \phi_2(t)(dx^2 + dy^2 + dz^2)$$

Symmetry reduction for flat FLRW cosmologies ([6,3,2]: E_3)

- infinitesimal group action

$$\Gamma = \text{span}\{\partial_x, \partial_y, \partial_z, x\partial_y - y\partial_x, y\partial_z - z\partial_y, z\partial_x - x\partial_z\}$$

- Γ -invariant metric

$$\hat{g} = -\phi_1(t) dt^2 + \phi_2(t)(dx^2 + dy^2 + dz^2)$$

- residual gauge freedom $t \rightarrow A(t)$ which would allow us to set $\phi_1 = 1$ breaks PSC

Symmetry reduction for flat FLRW cosmologies ([6,3,2]: E_3)

- infinitesimal group action

$$\Gamma = \text{span}\{\partial_x, \partial_y, \partial_z, x\partial_y - y\partial_x, y\partial_z - z\partial_y, z\partial_x - x\partial_z\}$$

- Γ -invariant metric

$$\hat{g} = -\phi_1(t) dt^2 + \phi_2(t)(dx^2 + dy^2 + dz^2)$$

- residual gauge freedom $t \rightarrow A(t)$ which would allow us to set $\phi_1 = 1$ breaks PSC
- l -chain

$$\chi = \partial_x \wedge \partial_y \wedge \partial_z$$

Symmetry reduction for flat FLRW cosmologies ([6,3,2]: E_3)

- infinitesimal group action

$$\Gamma = \text{span}\{\partial_x, \partial_y, \partial_z, x\partial_y - y\partial_x, y\partial_z - z\partial_y, z\partial_x - x\partial_z\}$$

- Γ -invariant metric

$$\hat{g} = -\phi_1(t) dt^2 + \phi_2(t)(dx^2 + dy^2 + dz^2)$$

- residual gauge freedom $t \rightarrow A(t)$ which would allow us to set $\phi_1 = 1$ breaks PSC
- l -chain

$$\chi = \partial_x \wedge \partial_y \wedge \partial_z$$

- Levi-Civita tensor

$$\sqrt{\phi_1\phi_2^3} dt \wedge dx \wedge dy \wedge dz$$

Symmetry reduction for flat FLRW cosmologies ([6,3,2]: E_3)

- infinitesimal group action

$$\Gamma = \text{span}\{\partial_x, \partial_y, \partial_z, x\partial_y - y\partial_x, y\partial_z - z\partial_y, z\partial_x - x\partial_z\}$$

- Γ -invariant metric

$$\hat{g} = -\phi_1(t) dt^2 + \phi_2(t)(dx^2 + dy^2 + dz^2)$$

- residual gauge freedom $t \rightarrow A(t)$ which would allow us to set $\phi_1 = 1$ breaks PSC
- l -chain

$$\chi = \partial_x \wedge \partial_y \wedge \partial_z$$

- Levi-Civita tensor

$$\sqrt{\phi_1\phi_2^3} dt \wedge dx \wedge dy \wedge dz$$

- reduced Lagrangian 1-form

$$\underline{\hat{L}} = \chi \bullet \underline{\epsilon}(\hat{g})L[\hat{g}] = \sqrt{\phi_1\phi_2^3} L[\hat{g}] dt$$

Conclusion

- We established the essential ingredients for a successful symmetry reduction:
 - 1 identified all possible PSC-compatible infinitesimal group actions Γ
 - 2 determined corresponding Γ -invariant metrics and l -chains in adapted coordinates
 - 3 minimized the amount of unknown functions employing residual gauge freedom compliant with PSC
- As a by-product, we implemented the symmetry reduction of Lagrangians in MATHEMATICA employing the `xAct` package.

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Thank you! Obrigado!

List of Γ -invariant l -chains

Hicks #	l -chain χ
3,2,1	$\partial_{y_3} \wedge \partial_{y_4}$
3,2,2	$\operatorname{csch} y_3 \partial_{y_3} \wedge \partial_{y_4}$
3,2,3	$\operatorname{csc} y_3 \partial_{y_3} \wedge \partial_{y_4}$
3,2,4	$\partial_{y_1} \wedge \partial_{y_2}$
3,2,5	$\operatorname{sech} y_2 \partial_{y_1} \wedge \partial_{y_2}$
3,3,2	$\partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_4}$
3,3,3	$\partial_{y_1} \wedge \partial_{y_3} \wedge \partial_{y_4}$
3,3,8	$\operatorname{csch} y_3 \partial_{y_1} \wedge \partial_{y_3} \wedge \partial_{y_4}$
3,3,9	$\operatorname{csc} y_3 \partial_{y_1} \wedge \partial_{y_3} \wedge \partial_{y_4}$
4,3,1	$\operatorname{csch} y_3 \partial_{y_1} \wedge \partial_{y_3} \wedge \partial_{y_4}$
4,3,2	$\operatorname{csch} y_3 \partial_{y_1} \wedge \partial_{y_3} \wedge \partial_{y_4}$
4,3,3	$\operatorname{csc} y_3 \partial_{y_1} \wedge \partial_{y_3} \wedge \partial_{y_4}$
4,3,4	$\operatorname{csc} y_3 \partial_{y_1} \wedge \partial_{y_3} \wedge \partial_{y_4}$
4,3,5	$\partial_{y_1} \wedge \partial_{y_3} \wedge \partial_{y_4}$
4,3,6	$\partial_{y_1} \wedge \partial_{y_3} \wedge \partial_{y_4}$
4,3,8	$\operatorname{sech} y_2 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_4}$
4,3,9	$\operatorname{sech} y_2 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3}$
4,3,10	$\partial_{y_1} \wedge \partial_{y_3} \wedge \partial_{y_4}$
4,3,11	$\partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_4}$
4,4,1	$\operatorname{csc} y_3 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
4,4,2	$\operatorname{csch} y_3 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
4,4,9	$\partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
4,4,18	$\partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
4,4,22	$\partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
5,4,1	$\operatorname{csch} y_3 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$

5,4,2	$\operatorname{csc} y_3 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
5,4,3	$\partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
5,4,6	$\partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
5,4,7	$\operatorname{csch} y_3 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
6,3,1	$\operatorname{csc}^2 y_2 \operatorname{csc} y_3 \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
6,3,2	$\partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
6,3,3	$\operatorname{csc} y_3 \operatorname{csch}^2 y_2 \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
6,3,4	$\operatorname{csch} y_2 \operatorname{sech} y_2 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3}$
6,3,5	$\partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_4}$
6,3,6	$\operatorname{csc} y_3 \operatorname{csch}^2 y_2 \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
6,4,1	$\operatorname{csc} y_3 \operatorname{sech} y_2 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
6,4,2	$\operatorname{csch} y_3 \operatorname{sech} y_2 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
6,4,3	$\operatorname{sech} y_2 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
6,4,4	$\operatorname{csc} y_3 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
6,4,5	$\operatorname{csch} y_3 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
7,4,1	$\operatorname{csc}^2 y_2 \operatorname{csc} y_3 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
7,4,2	$\operatorname{csc} y_3 \operatorname{csch}^2 y_2 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
7,4,3	$\operatorname{csc} y_3 \operatorname{csch}^2 y_2 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$
7,4,4	$\operatorname{csch} y_2 \operatorname{sech} y_2 \partial_{y_1} \wedge \partial_{y_2} \wedge \partial_{y_3} \wedge \partial_{y_4}$