

Polarized predictions in diboson final states

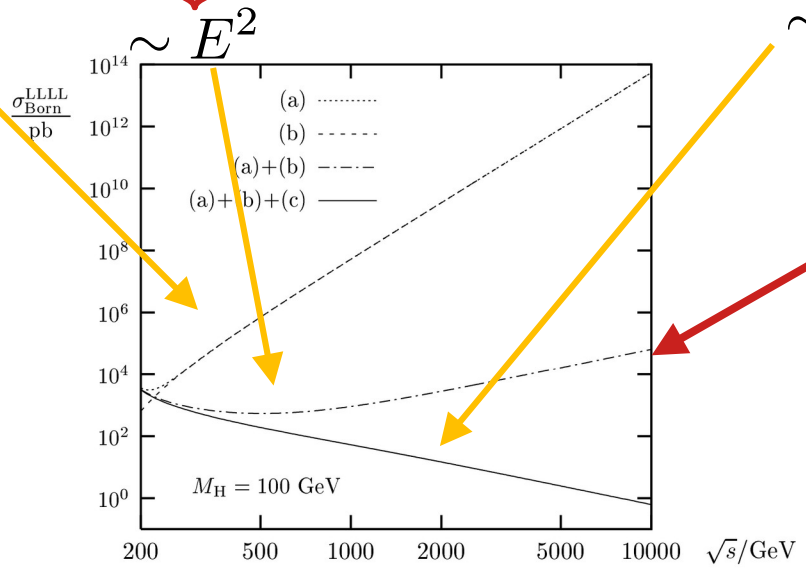
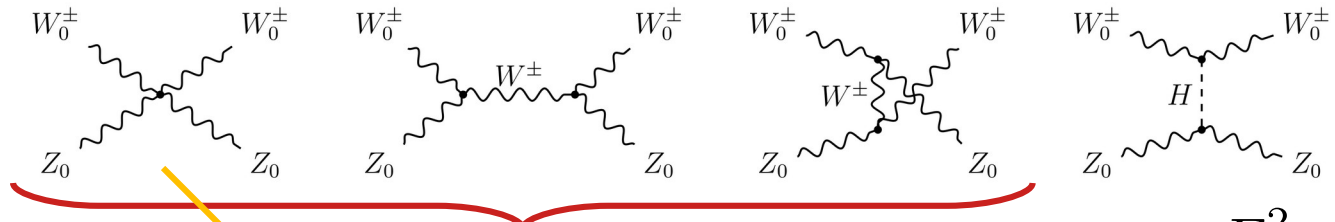
Rene Poncelet



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POLISH ACADEMY OF SCIENCES



Longitudinal Vector-Boson-Scattering (VBS)



Unitarity violation

Measurement of polarized boson scattering or production probes:

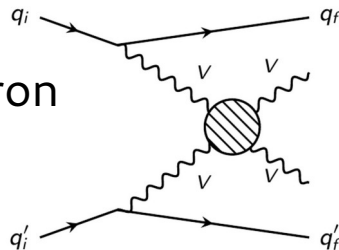
- EWSB mechanism
- Higgs and gauge sector
- New physics models

Radiative corrections to $W^+ W^- \rightarrow W^+ W^-$ in the electroweak standard model

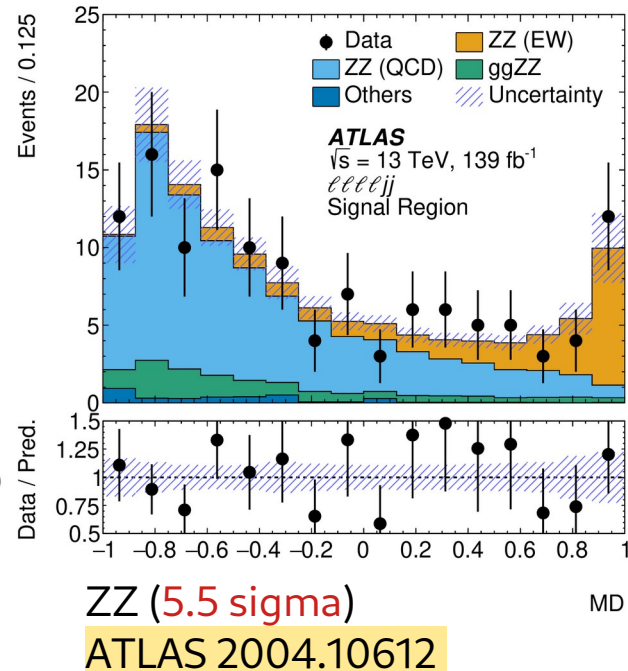
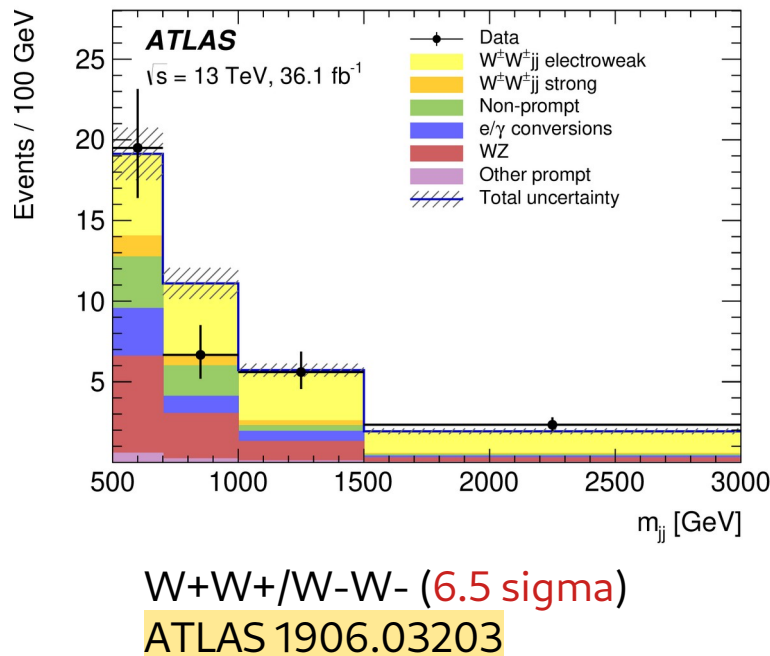
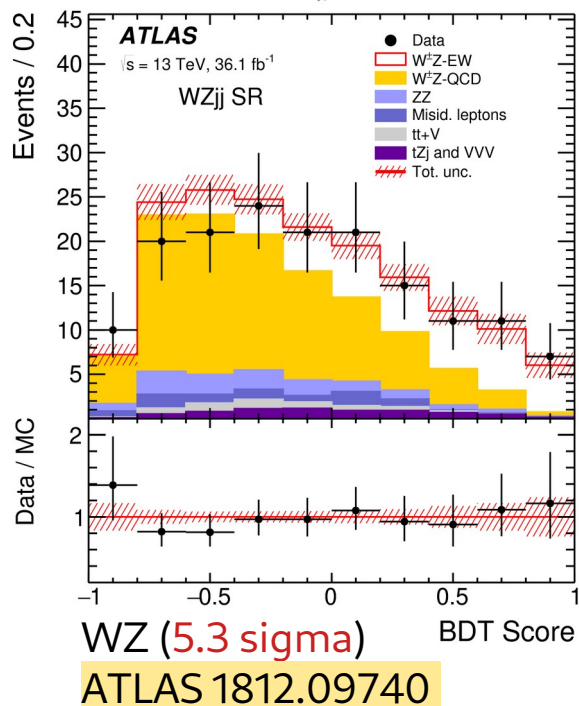
A. Denner, T. Hahn hep-ph/9711302

VBS at hadron colliders

VBS at hadron colliders

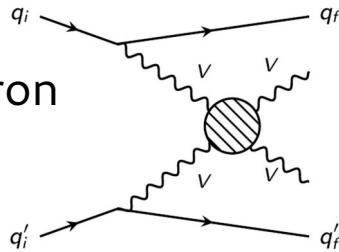


Separate from background processes through VBS topology
 → a rare process, but observed.

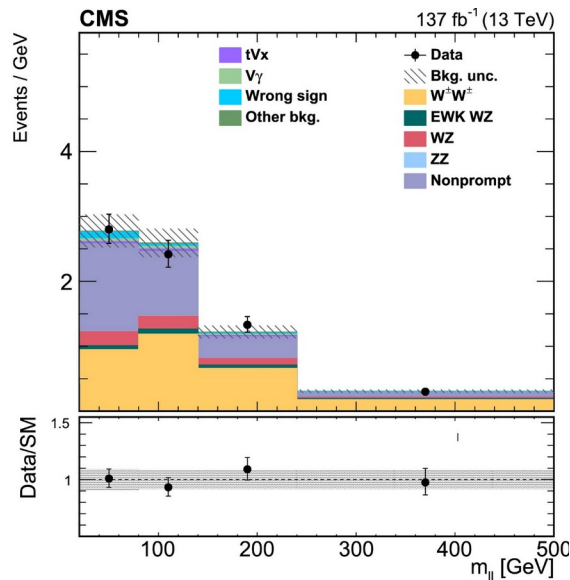
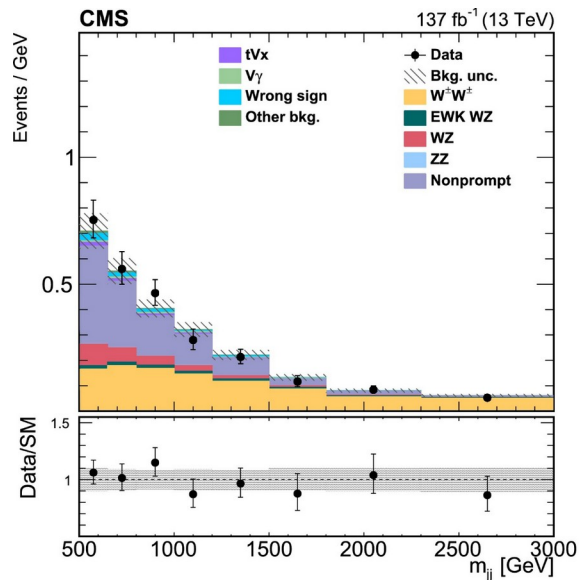


VBS at hadron colliders

VBS at hadron colliders

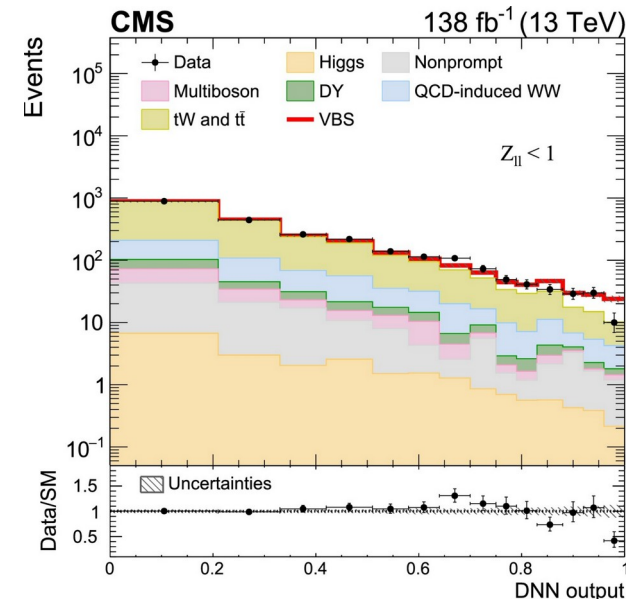


Separate from background processes through VBS topology
 → a rare process, but observed.



WZ (6.8 sigma) + W+W+/W-W- (diff. xsec)

CMS 2005.01173



W+W- (5.6 sigma)

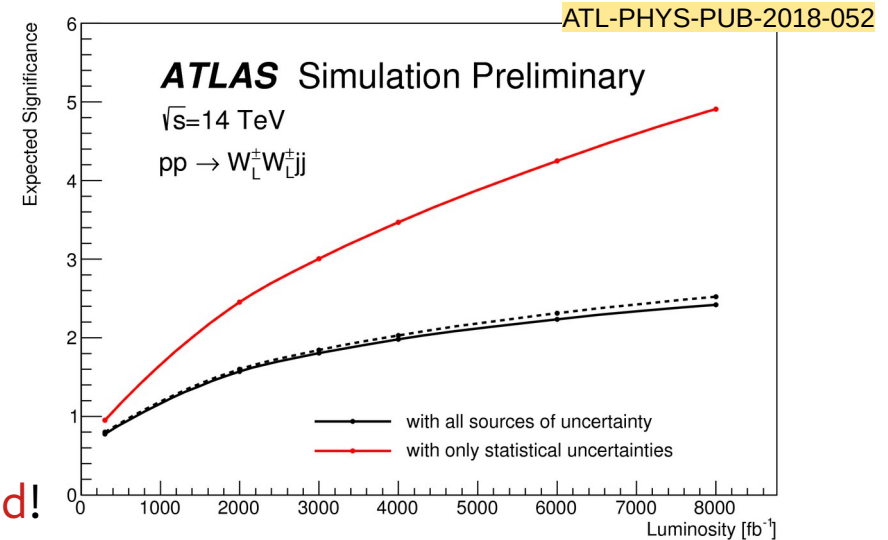
CMS 2205.05711

Polarised VBS at HL-LHC

If we want to study unitarisation/EWSB we need to **extract the longitudinal component**

- only 5-10 % of the total rate
→ **very challenging**
(remember: $130\text{fb}^{-1} \rightarrow \sim 5\text{-}7$ sigma
→ naive improvement by factor 10 necessary for observation)
- Requires CMS/ATLAS combination and/or new techniques at HL-LHC
→ **improvement of systematic uncertainties needed!**

ATLAS HL-LHC projection

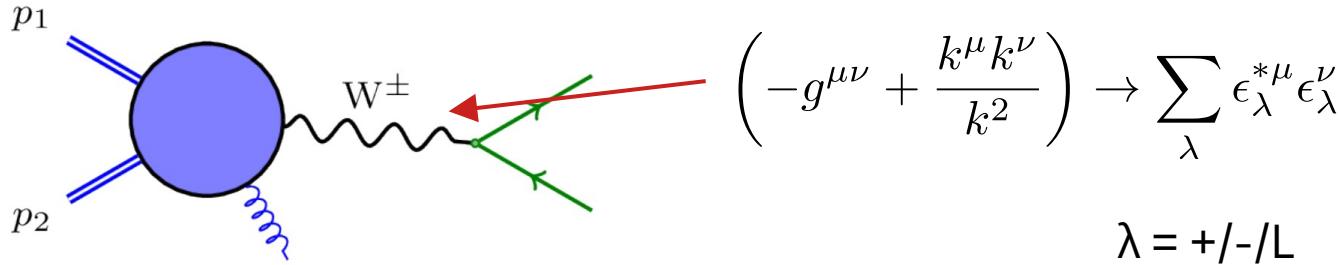


How to improve on the (theory) systematics?

→ Improved signal and background (i.e. transverse part)

→ Effective separation of boson polarisation

Polarised boson production



Can we extract the longitudinal component?

Measurements of longitudinal polarisation fractions:

Measurement of the Polarization of W Bosons with Large Transverse Momenta in W+Jets Events at the LHC,

CMS 1104.3829

Measurement of the polarisation of W bosons produced with large transverse momentum in pp collisions at $\sqrt{s}=7$ TeV with the ATLAS experiment,

ATLAS 1203.2165

Measurement of WZ production cross sections and gauge boson polarisation in pp collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector,

ATLAS 1902.05759

Measurement of the inclusive and differential WZ production cross sections, polarization angles, and triple gauge couplings in pp collisions at $\sqrt{s} = 13$ TeV,

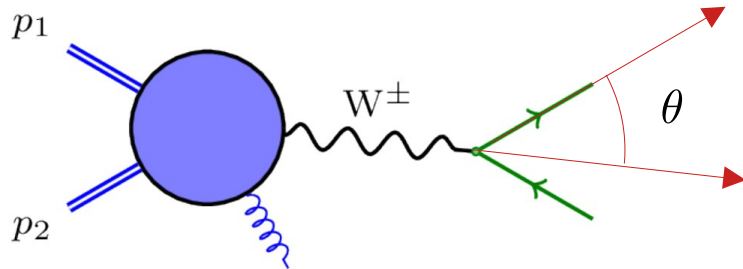
CMS 2110.11231

Observation of gauge boson joint-polarisation states in WZ production from pp collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector

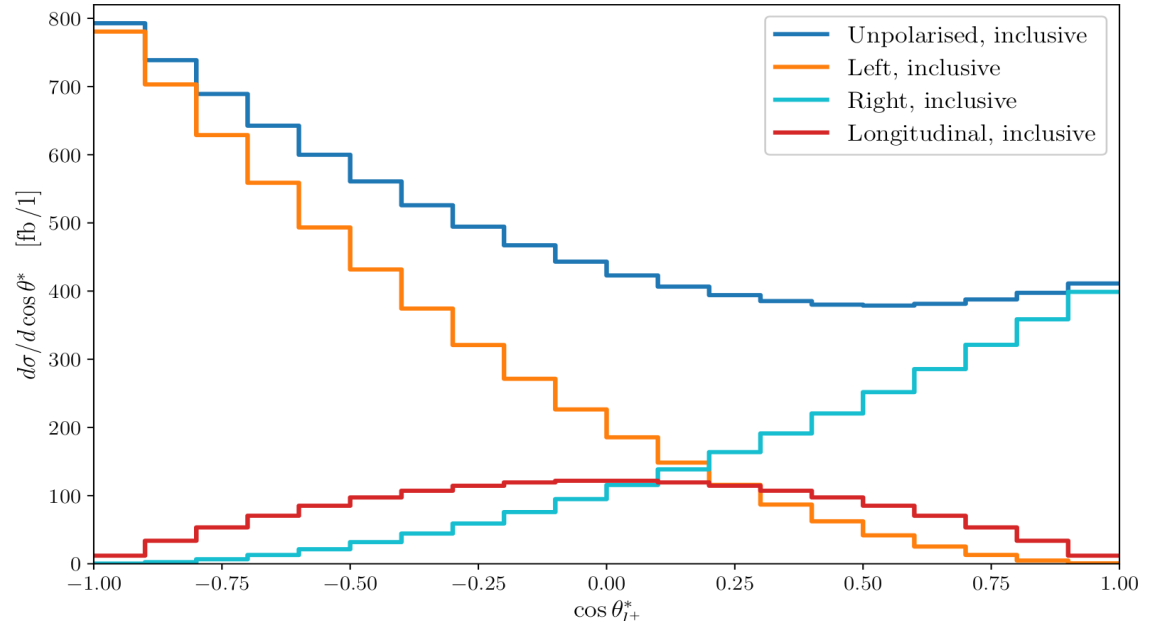
ATLAS 2211.09435

How to measure polarized bosons?

- We can't measure boson polarization directly.
- Luckily decay products can be used as a “polarimeter”:

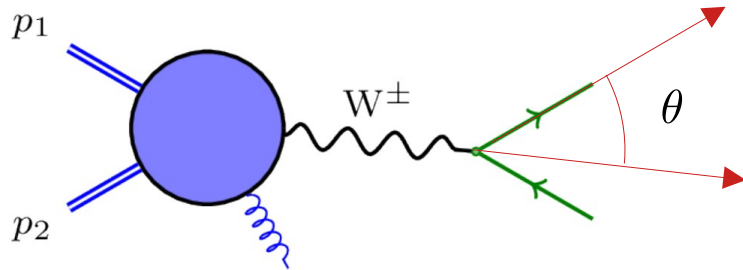


W^+ decay (W^- mirrored around 0)

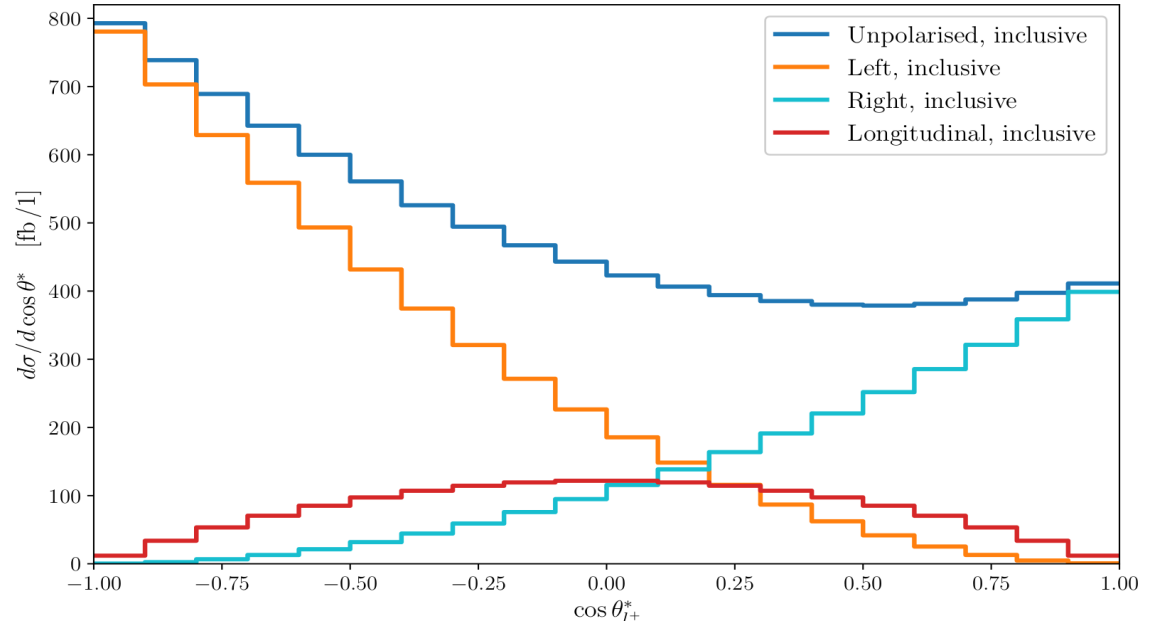


How to measure polarized bosons?

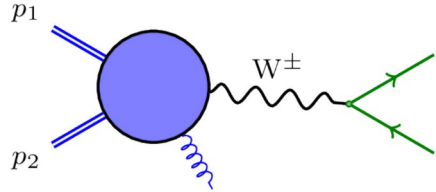
- We can't measure boson polarization directly.
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W^+ decay (W^- mirrored around 0)



Polarized cross sections



On-shell bosons: $\left(-g^{\mu\nu} + \frac{k^\mu k^\nu}{k^2}\right) \rightarrow \sum_{\lambda} \epsilon_{\lambda}^{*\mu} \epsilon_{\lambda}^{\nu}$
(DPA or NWA)

$$M = \mathbf{P}_{\mu} \cdot \frac{-g_{\mu\nu} + \frac{k^{\mu} k^{\nu}}{k^2}}{k^2 - M_V^2 + iM_V \Gamma_V} \cdot \mathbf{D}_{\nu}$$

$$|M|^2 = \underbrace{\sum_{\lambda} |M_{\lambda}|^2}_{\text{polarised x-sections}} + \underbrace{\sum_{\lambda \neq \lambda'} M_{\lambda}^* M_{\lambda'}}_{\text{Interferences}}$$

→ polarised x-sections Interferences

Create samples of fixed polarisation: $\frac{d\sigma}{dX} = f_L \frac{d\sigma_L}{dX} + f_R \frac{d\sigma_R}{dX} + f_0 \frac{d\sigma_0}{dX} \left(+ f_{int.} \frac{d\sigma_{int.}}{dX} \right)$

Template fit f_L, f_R, f_0 to measured $\frac{d\sigma^{exp.}}{dX}$

Polarized cross sections

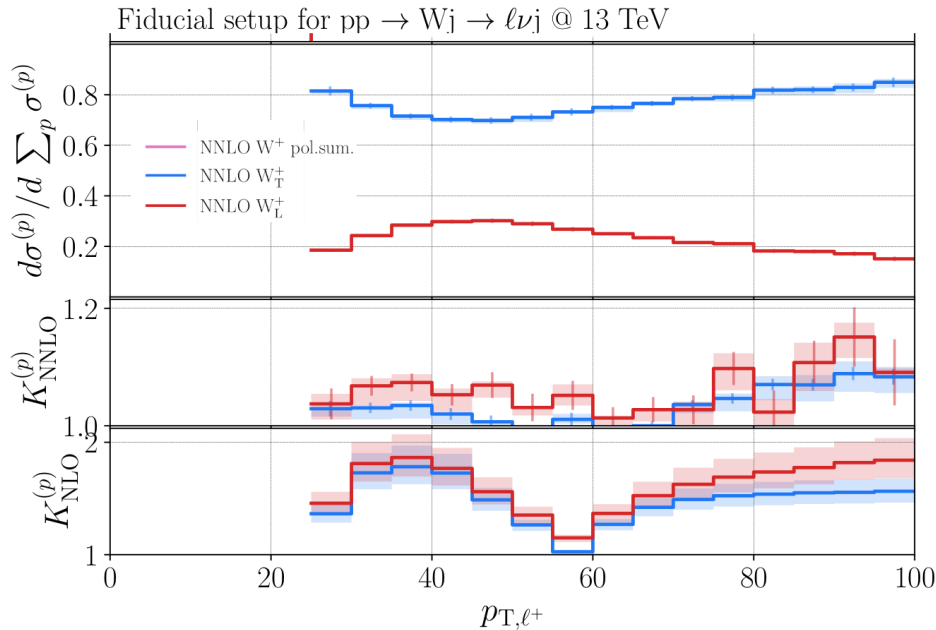
$$\frac{d\sigma}{dX} = f_L \frac{d\sigma_L}{dX} + f_R \frac{d\sigma_R}{dX} + f_0 \frac{d\sigma_0}{dX} \left(+ f_{int.} \frac{d\sigma_{int.}}{dX} \right)$$

- Interferences can be handled
- Does not rely on extrapolations to the full phase space
X can be any observable → lab frame observables
- $\frac{d\sigma_i}{dX}$ can be systematically improved

Higher-order QCD/EW corrections + PS
to minimize uncertainties from MHO (scale uncertainties)

Why do we need higher-order corrections?

Example: $pp \rightarrow W^\pm (\rightarrow l\nu) j$



Important

Just using single NNLO K-factors is not enough

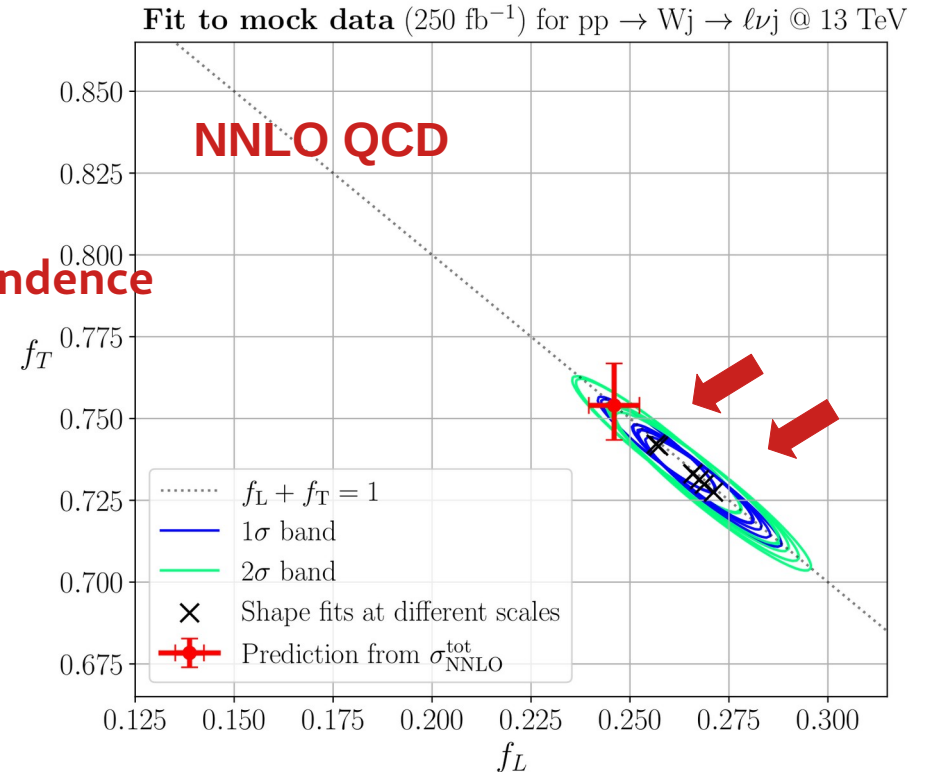
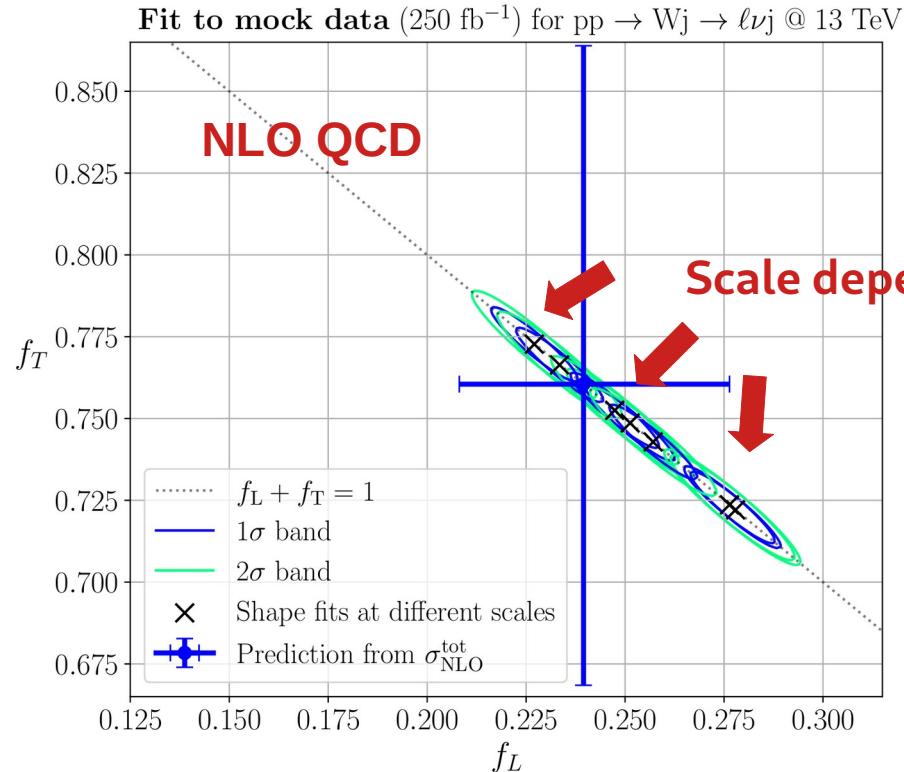
- 1) Differential polarization fractions have shapes (not just one number!)
- 2) Higher-order corrections dependent on polarization! Just using unpolarized K-factor would lead to distortion of spectrum.
- 3) NNLO QCD needed to reach percent-level scale-dependence \rightarrow MHO

Polarised W+j production at the LHC: a study at NNLO QCD accuracy,
Pellen, Poncelet, Popescu 2109.14336

W+jet: mock-data fit

Fit to mock-data (based on NNLO QCD and 250 fb⁻¹ stats):
→ extreme case to see effect of scale dependence reduction

Observable: $\cos(\ell, j_1)$



Status of polarization calculations

Process	LO	NLO	NLO EW	NNLO	+ PS
pp → WW	X	X	X	X	X
pp → ZZ	X	X	X		X
pp → WZ	X	X	X		X
pp → W/Z	X	X	X	Ang.	X
pp → W+j	X	X	(X)	X	
pp → Z+j	X	Ang.		Ang.	
VBS	X	X			

(Collection of papers in the backup)

Ang. = angular coefficients

Polarised nLO+PS: SHERPA

See also talk by Steffen

Polarised cross sections for vector boson production with SHERPA
Hoppe, Schönherr, Siegert 2310.14803

- New bookkeeping of boson polarizations in SHERPA for LO MEs
- Approximate NLO corrections: nLO+PS
 - Reals+matching are treated exact
 - loop matrix elements unpolarised (reweighted by pol. tree MEs)
- Comparison with multi-jet merged calculations

Comparison with fixed order

- nLO+PS approximation in fair agreement with full NLO
 - good for polarization fractions

W ⁺ Z	σ^{NLO} [fb]	Fraction [%]	K-factor	$\sigma_{\text{SHERPA}}^{\text{nLO+PS}}$ [fb]	Fraction [%]	K-factor
full	35.27(1)		1.81	33.80(4)		
unpol	34.63(1)	100	1.81	33.457(26)	100	1.79
Laboratory frame						
L-U	8.160(2)	23.563(9)	1.93	7.962(5)	23.796(25)	1.91
T-U	26.394(9)	76.217(34)	1.78	25.432(21)	76.01(9)	1.75
int	0.066(10) (diff)	0.191(29)	2.00	0.064(7)	0.191(22)	2.40(40)
U-L	9.550(4)	27.577(14)	1.73	9.275(16)	27.72(5)	1.72
U-T	25.052(8)	72.342(31)	1.83	24.156(18)	72.20(8)	1.81
int	0.028(10) (diff)	0.081(29)	-0.49	0.026(7)	0.079(22)	-0.471(34)

Polarised NLO+PS: POWHEG

Polarised-boson pairs at the LHC with NLOPS accuracy
Pelliccioli, Zanderighi 2311.05220

- NLO QCD + PS in POWHEG-BOX-RES framework
 - Study of PS (Pythia8) + hadronisation effects on fractions and differential distributions WW/WZ/ZZ
- 1-5% effect on distributions, but generally small impact on fractions (~1% effects)

state	σ [fb] LHE	ratio [/unp., %] LHE	σ [fb] PS+hadr	ratio [/unp., %] PS+hadr
Inclusive setup				
full off-shell	98.36(3) ^{+4.8%} _{-3.9%}	101.20	95.27(3) ^{+4.9%} _{-3.9%}	101.28
unpolarised	97.20(3) ^{+4.8%} _{-3.9%}	100	94.07(3) ^{+4.9%} _{-3.9%}	100
LL	4.499(2) ^{+2.8%} _{-2.3%}	4.63 ^{+0.13} _{-0.13}	4.359(2) ^{+2.8%} _{-2.2%}	4.63 ^{+0.13} _{-0.13}
LT	13.151(4) ^{+7.0%} _{-5.7%}	13.53 ^{+0.28} _{-0.27}	12.730(5) ^{+7.0%} _{-5.7%}	13.53 ^{+0.28} _{-0.28}
TL	12.724(4) ^{+7.3%} _{-5.9%}	13.09 ^{+0.32} _{-0.31}	12.314(5) ^{+7.4%} _{-5.9%}	13.09 ^{+0.31} _{-0.32}
TT	66.88(2) ^{+4.0%} _{-3.3%}	68.81 ^{+0.47} _{-0.51}	64.74(2) ^{+4.1%} _{-3.2%}	68.82 ^{+0.46} _{-0.51}
interference	-0.058	-0.06	-0.069	-0.06

NNLO QCD polarized WW production

NNLO QCD study of polarised W+W- production at the LHC,
Poncelet, Popescu 2102.13583



Technical aspects:

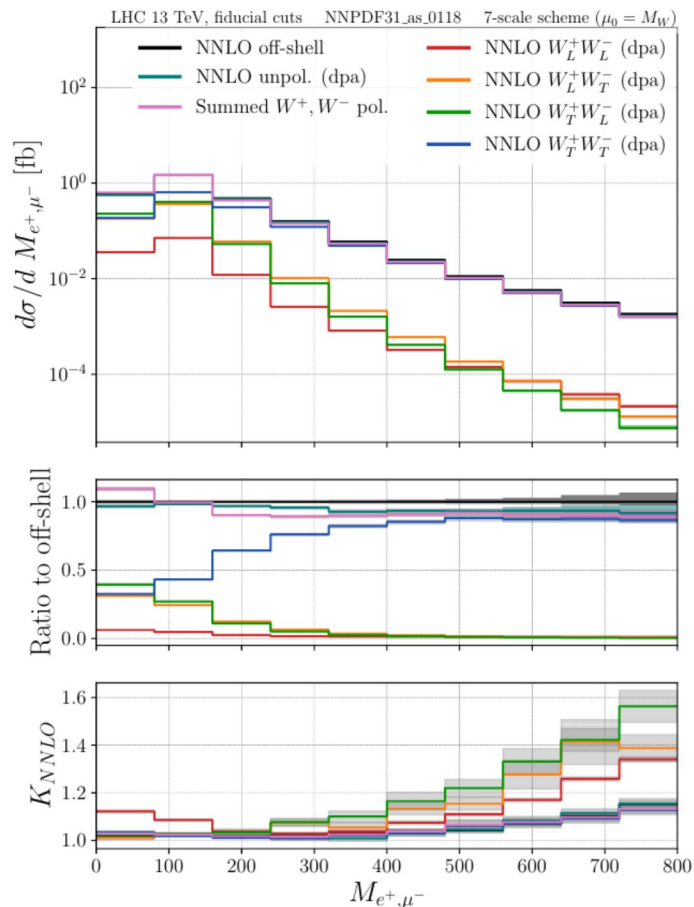
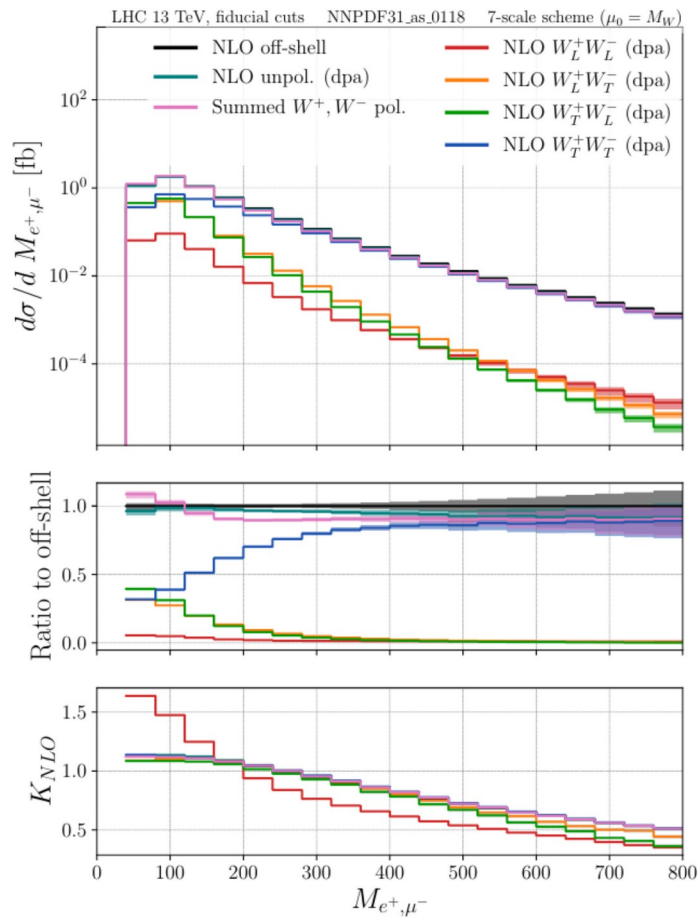
- Implementation of NNLO QCD in c++ sector-improved residue subtraction framework [1408.2500,1907.12911]
- Massive b-quarks → get rid of top production ($pp \rightarrow b\bar{b}W^+W^-$ enters at NNLO)
- NNPDF31 and a fixed renormalisation scale: $\mu_R = \mu_F = m_W$

Fiducial phase space

Measurement of fiducial and differential W+W- production cross-sections at $\sqrt{s} = 13$ TeV with the ATLAS detector
ATLAS 1905.04242

- Leptons: $p_T(\ell) \geq 27$ GeV $|y(\ell)| < 2.5$ $m(\ell\bar{\ell}) > 55$ GeV
- Missing transverse momentum: $p_{T,\text{miss}} = p_T(\nu_e + \bar{\nu}_\mu) \geq 20$ GeV
- Jet-veto: $p_T(j) > 35$ GeV $|y(j)| < 4.5$

Polarised di-boson production



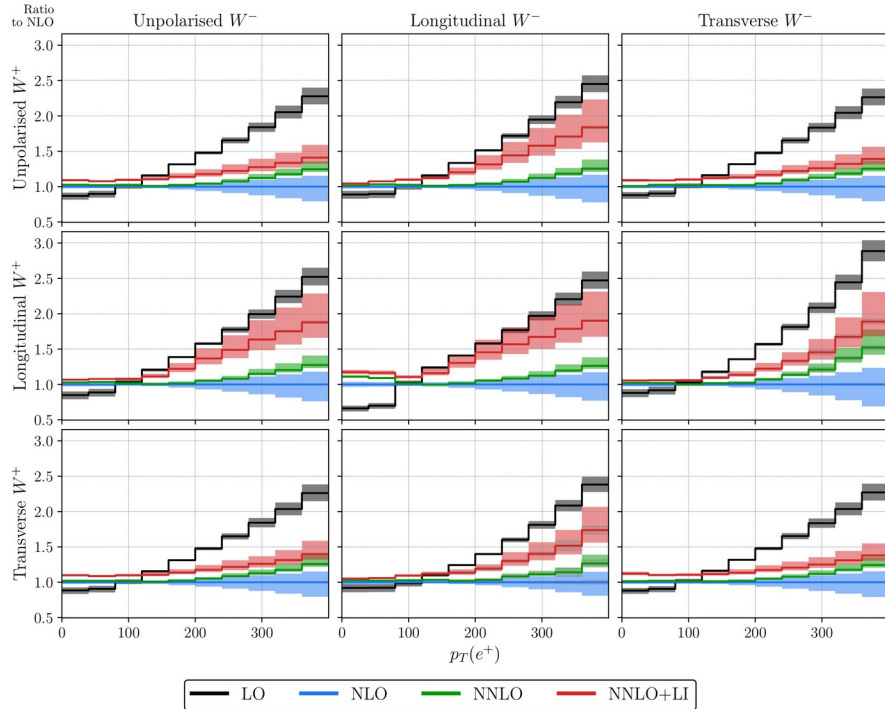
Doubly polarised cross sections

	NLO	NNLO	K_{NNLO}	LI	NNLO+LI
off-shell	220.06(5) ^{+1.8%} _{-2.3%}	225.4(4) ^{+0.6%} _{-0.6%}	1.024	13.8(2) ^{+25.5%} _{-18.7%}	239.1(4) ^{+1.5%} _{-1.2%}
unpol. (nwa)	221.85(8) ^{+1.8%} _{-2.3%}	227.3(6) ^{+0.6%} _{-0.6%}	1.025	13.68(3) ^{+25.5%} _{-18.7%}	241.0(6) ^{+1.5%} _{-1.1%}
unpol. (dpa)	214.55(7) ^{+1.8%} _{-2.3%}	219.4(4) ^{+0.6%} _{-0.6%}	1.023	13.28(3) ^{+25.5%} _{-18.7%}	232.7(4) ^{+1.4%} _{-1.1%}
W_L^+ (dpa)	57.48(3) ^{+1.9%} _{-2.6%}	59.3(2) ^{+0.7%} _{-0.7%}	1.032	2.478(6) ^{+25.5%} _{-18.3%}	61.8(2) ^{+1.0%} _{-0.8%}
W_L^- (dpa)	63.69(5) ^{+1.9%} _{-2.6%}	65.4(3) ^{+0.8%} _{-0.8%}	1.026	2.488(6) ^{+25.5%} _{-18.3%}	67.9(3) ^{+0.9%} _{-0.8%}
W_T^+ (dpa)	152.58(9) ^{+1.7%} _{-2.1%}	155.7(6) ^{+0.7%} _{-0.6%}	1.020	11.19(2) ^{+25.5%} _{-18.8%}	166.9(6) ^{+1.6%} _{-1.3%}
W_T^- (dpa)	156.41(7) ^{+1.7%} _{-2.1%}	159.7(6) ^{+0.5%} _{-0.6%}	1.021	11.19(2) ^{+25.5%} _{-18.8%}	170.9(6) ^{+1.7%} _{-1.3%}
$W_L^+ W_L^-$ (dpa)	9.064(6) ^{+3.0%} _{-3.0%}	9.88(3) ^{+1.3%} _{-1.3%}	1.090	0.695(2) ^{+25.5%} _{-18.8%}	10.57(3) ^{+2.9%} _{-2.4%}
$W_L^+ W_T^-$ (dpa)	48.34(3) ^{+1.9%} _{-2.5%}	49.4(2) ^{+0.9%} _{-0.7%}	1.021	1.790(5) ^{+25.5%} _{-18.3%}	51.2(2) ^{+0.6%} _{-0.8%}
$W_T^+ W_L^-$ (dpa)	54.11(5) ^{+1.9%} _{-2.5%}	55.5(4) ^{+0.6%} _{-0.7%}	1.025	1.774(5) ^{+25.5%} _{-18.3%}	57.2(4) ^{+0.7%} _{-0.7%}
$W_T^+ W_T^-$ (dpa)	106.26(4) ^{+1.6%} _{-1.9%}	108.3(3) ^{+0.5%} _{-0.5%}	1.019	9.58(2) ^{+25.5%} _{-18.9%}	117.9(3) ^{+2.1%} _{-1.6%}

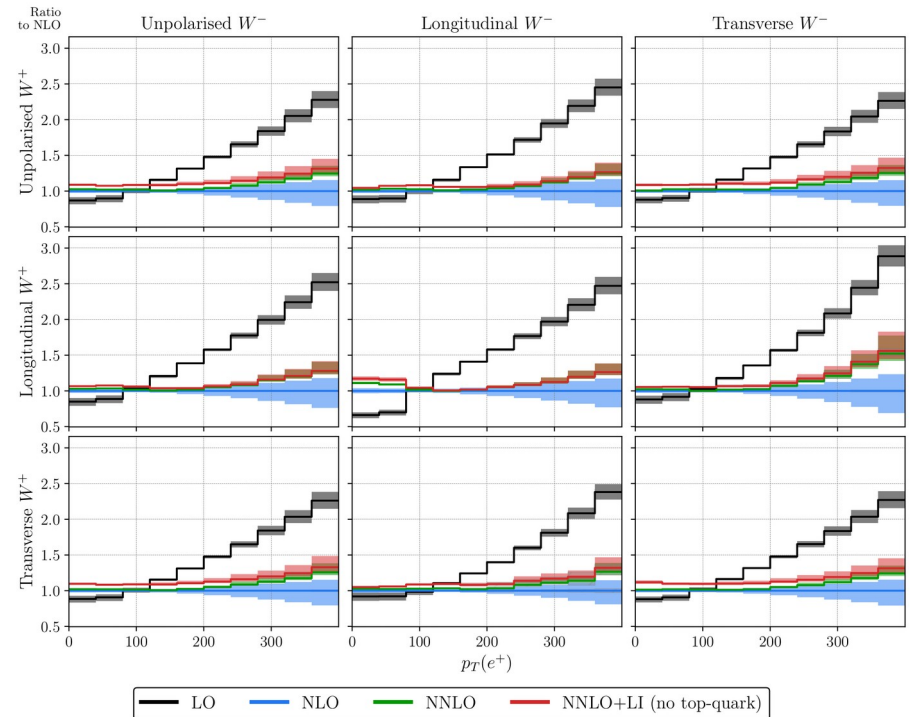
Small LL contribution, with large corrections (\rightarrow polarization frame)

Loop induced $gg \rightarrow WW$ contributions

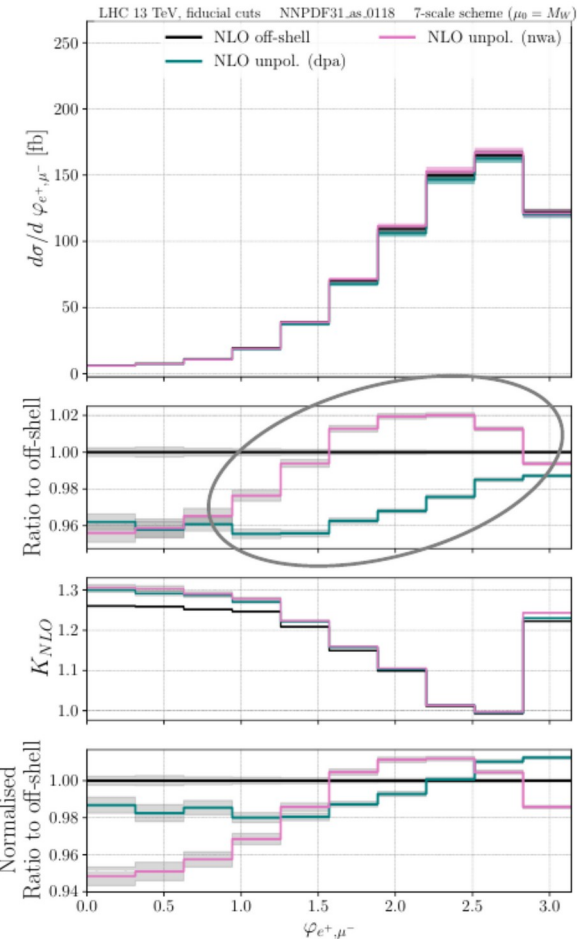
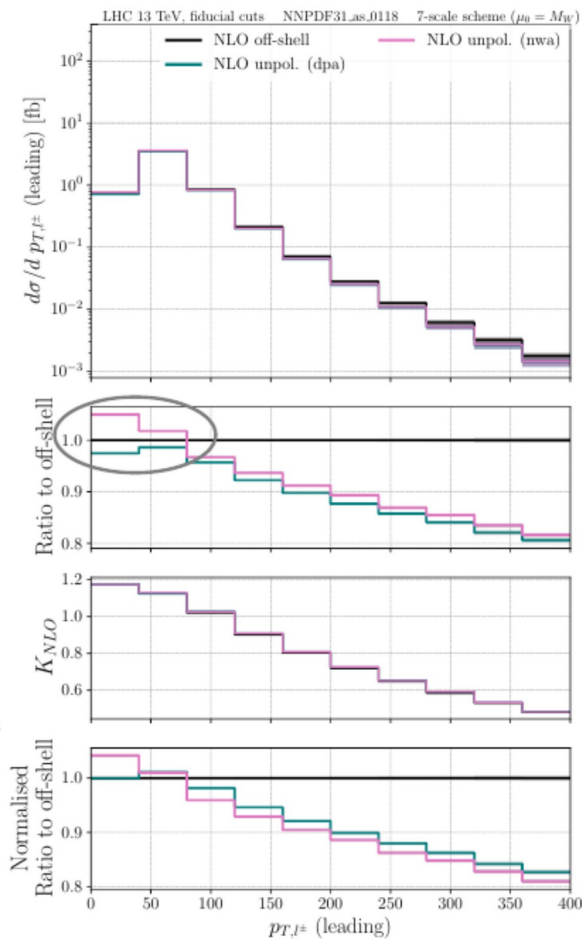
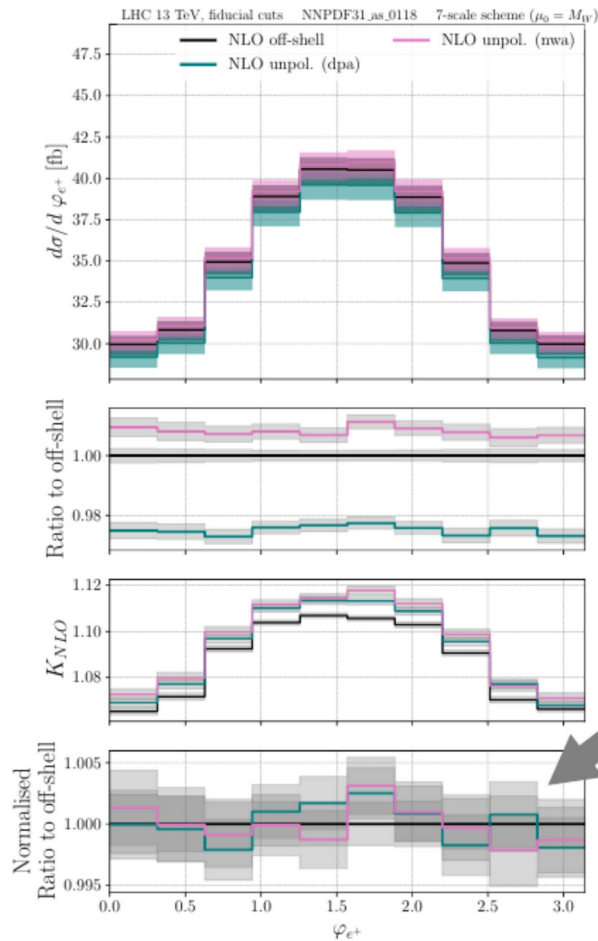
With top-quark loops in gg LI



Without top-quark loops in gg LI



NWA vs. DPA



Take home messages

- Precise (and accurate) SM predictions for polarized cross section are important to pin down the longitudinal component.
- NLO QCD/EW (+PS) are the state-of-the-art for polarized EW boson processes
→ few process are available at NNLO QCD
- Looking at higher-orders NNLO QCD
 - Scale dependence can mimic signal → NNLO QCD needed to reduce these effects
 - Loop-induced contributions: 'LO' at NNLO → needs partial N3LO QCD
- What's next? → Phenomenology, benchmark new tools (Powheg, SHERPA, Madgraph)
 - NNLO QCD for VV, (+ NLO EW), providing templates through *high tea*
 - + SMEFT



Comprehensive Multiboson Experiment-Theory Action

- WG1 - Theoretical framework, precision calculations and simulation
- WG2 - Technological innovation in data analysis
- WG3 - Experimental Measurements
- WG4 - Management and Event Organization
- WG5 - Inclusiveness and Outreach

Further information:

<https://www.cost.eu/actions/CA22130/> and <https://cometa.web.cern.ch/>

Backup

Polarized VV @ (N)NLO QCD / NLO EW

Fiducial polarization observables in hadronic WZ production: A next-to-leading order QCD+EW study,

Baglio, Le Duc 1810.11034

Anomalous triple gauge boson couplings in ZZ production at the LHC and the role of Z boson polarizations,

Rahama, Singh 1810.11657

Polarization observables in WZ production at the 13 TeV LHC: Inclusive case,

Baglio, Le Duc 1910.13746

Unravelling the anomalous gauge boson couplings in ZW⁺- production at the LHC and the role of spin-1 polarizations,

Rahama, Singh 1911.03111

Polarized electroweak bosons in W+W⁻ production at the LHC including NLO QCD effects,

Denner, Pelliccioli 2006.14867

NLO QCD predictions for doubly-polarized WZ production at the LHC,

Denner, Pelliccioli 2010.07149

NNLO QCD study of polarised W+W⁻ production at the LHC,

Poncelet, Popescu 2102.13583

NLO EW and QCD corrections to polarized ZZ production in the four-charged-lepton channel at the LHC,

Denner, Pelliccioli 2107.06579

Breaking down the entire spectrum of spin correlations of a pair of particles involving fermions and gauge bosons,

Rahama, Singh 2109.09345

Doubly-polarized WZ hadronic cross sections at NLO QCD+EW accuracy,

Duc Ninh Le, Baglio 2203.01470

Doubly-polarized WZ hadronic production at NLO QCD+EW: Calculation method and further results

Duc Ninh Le, Baglio, Dao 2208.09232

NLO QCD corrections to polarised di-boson production in semi-leptonic final states

Denner, Haitz, Pelliccioli 2211.09040

Polarised cross sections for vector boson production with SHERPA

Hoppe, Schönherr, Siegert 2310.14803

Polarised-boson pairs at the LHC with NLOPS accuracy

Pelliccioli, Zanderighi 2311.05220

NLO EW corrections to polarised W+W⁻ production and decay at the LHC

Denner, Haitz, Pelliccioli 2311.16031

NLO electroweak corrections to doubly-polarized W+W⁻ production at the LHC

Thi Nhung Dao, Duc Ninh 2311.17027

Polarized ZZ pairs in gluon fusion and vector boson fusion at the LHC

Javurkova, Ruiz, Coelho, Sandesara 2401.17365

Other polarized cross section calculations

- Polarised VBS (so far LO):

W boson polarization in vector boson scattering at the LHC,

Ballestrero, Maina, Pelliccioli 1710.09339

Polarized vector boson scattering in the fully leptonic WZ and ZZ channels at the LHC,

Ballestrero, Maina, Pelliccioli 1907.04722

Automated predictions from polarized matrix elements

Buarque Franzosi, Mattelaer, Ruiz, Shil 1912.01725

Different polarization definitions in same-sign WW scattering at the LHC,

Ballestrero, Maina, Pelliccioli 2007.07133

- Single boson production

Left-Handed W Bosons at the LHC,

Z. Bern et. al. 1103.5445

Electroweak gauge boson polarisation at the LHC,

Stirling, Vryonidou 1204.6427

What Does the CMS Measurement of W-polarization Tell Us about the Underlying Theory of the Coupling of W-Bosons to Matter?,

Belyaev, Ross 1303.3297

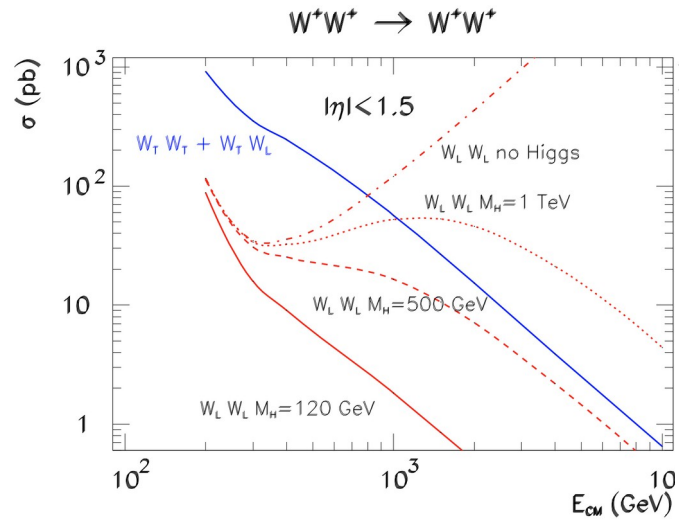
Polarised W+j production at the LHC: a study at NNLO QCD accuracy,

Pellen, Poncelet, Popescu 2109.14336

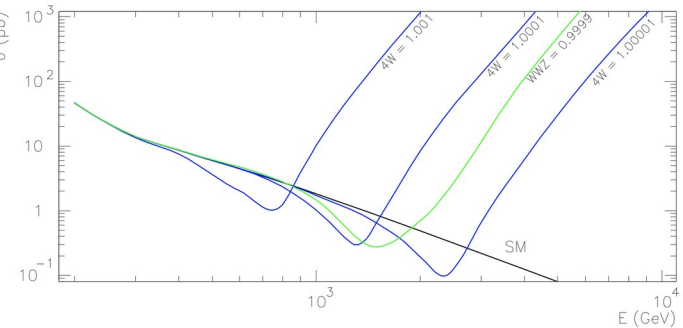
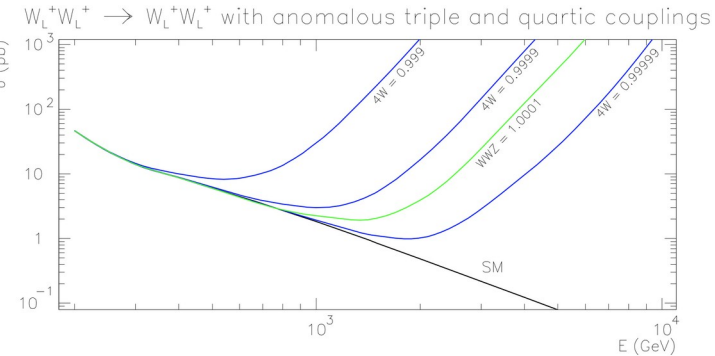
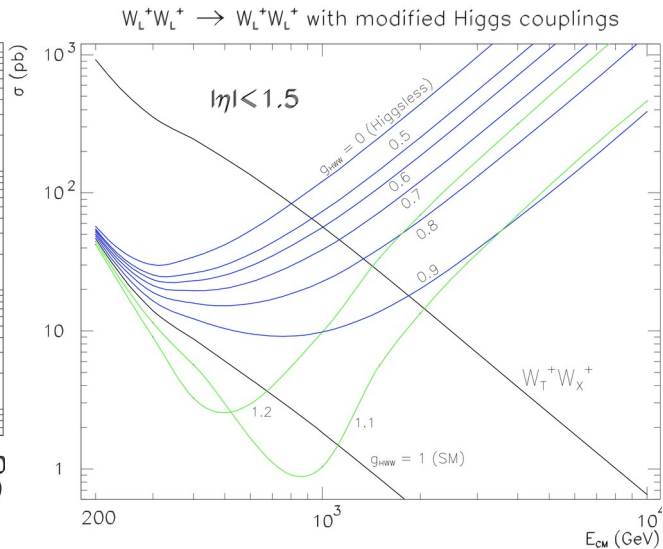
Longitudinal Vector-Boson-Scattering (VBS)

The Higgs boson and the physics of WW scattering before and after Higgs discovery
M. Szleper 1412.8367

Sensitivity to the Higgs mass



Modified HVV, VVV, VVVV couplings



EWSB

The reason is the EWSB in the SM:

$$\mathcal{L}_{EW} = -\frac{1}{4}(W_{\mu\nu}^i)^2 - \frac{1}{4}(B_{\mu\nu}^i)^2 + (D_\mu\phi)^2 - V(\phi^\dagger\phi)$$

- Higgs potential and minimum:

$$V(\phi^\dagger\phi) = -\mu^2(\phi^\dagger\phi)^2 + \lambda(\phi^\dagger\phi)^4 \quad \phi = U(\pi^i) \begin{pmatrix} 0 \\ \frac{v+H}{\sqrt{2}} \end{pmatrix} \quad \text{VEV: } \phi^\dagger\phi = \frac{\mu^2}{2\lambda} \equiv \frac{v^2}{2}$$

- Goldstone bosons can be absorbed via gauge transformation (unitary gauge).
This gives rise to massive gauge bosons:

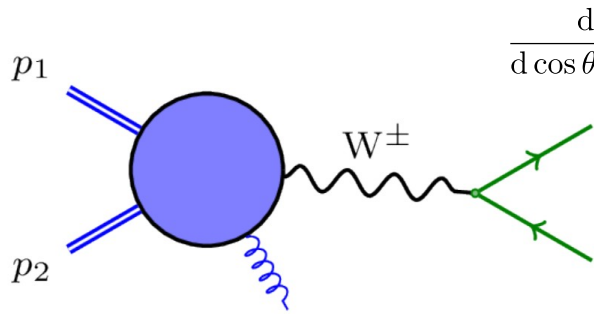
$$\phi = U^{-1}(\pi^i)\phi, \quad W_\mu = U^{-1}W_\mu U - \frac{i}{g_W}U^{-1}\partial_\mu U$$

$$|D_\mu\phi|^2 \ni \frac{v^2}{8} [2g_W^2 W_\mu^+ W^{-\mu} + (g_W W_\mu^3 - g'_W B_\mu)^2] \quad \longrightarrow \quad M_W = \frac{1}{2}vg_W, \quad M_Z = \frac{M_W}{\cos\theta_W}$$

- Restores renormalizability and unitarity

Angular coefficients

Angular decomposition of 2-body W decay:



$$\frac{d\sigma}{d\cos\theta d\phi dX} = \frac{d\sigma}{dX} \frac{3}{16\pi} \left[(1 + \cos^2\theta) + \frac{A_0}{2}(1 - 3\cos^2\theta) + A_1 \sin 2\theta \cos\phi + \frac{A_2}{2} \sin^2\theta \cos 2\phi \right. \\ \left. + A_3 \sin\theta \cos\phi + A_4 \cos\theta + A_5 \sin^2\theta \sin 2\phi + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \right]$$

After azimuthal integration:

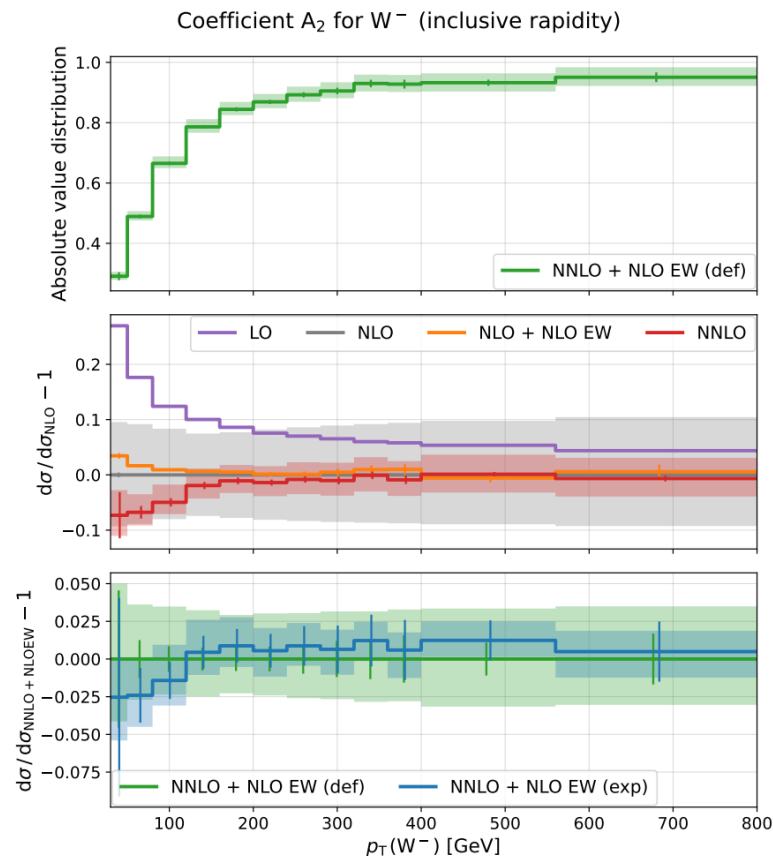
$$\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta} = \frac{3}{4} \sin\theta f_0 + \frac{3}{8} (1 - \cos\theta)^2 f_L + \frac{3}{8} (1 + \cos\theta)^2 f_R$$

Idea: Suitable projections (or fits) extract fractions of left, right and longitudinal components.

Angular coefficients as function of V kinematics

Keeping azimuthal dependence & boson kinematics:

$$\frac{d\sigma}{dp_{T,W} dy_W dm_{\ell\nu} d\Omega} = \frac{3}{16\pi} \frac{d\sigma^{U+L}}{dp_{T,W} dy_W dm_{\ell\nu}} \left((1 + \cos^2 \theta) + A_0 \frac{1}{2} (1 - 3 \cos^2 \theta) \right. \\ \left. + A_1 \sin 2\theta \cos \phi + A_2 \frac{1}{2} \sin^2 \theta \cos 2\phi + A_3 \sin \theta \cos \phi + A_4 \cos \theta \right. \\ \left. + A_5 \sin^2 \theta \sin 2\phi + A_6 \sin 2\theta \sin \phi + A_7 \sin \theta \sin \phi \right),$$



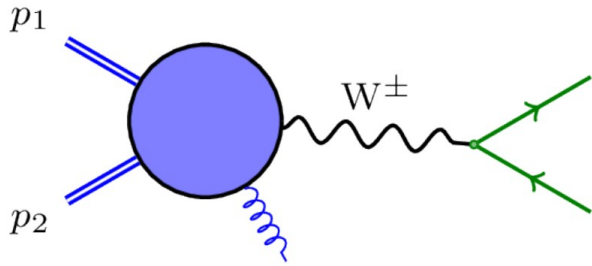
Angular coefficients in $W+j$ production at the LHC with high precision
 Pellen, Poncelet, Popescu, Vitos, 2204.12394

Angular coefficients, practical considerations

This simple idea suffers from:

- Fiducial phase space requirements on the leptons:
 - Interferences do not cancel
 - Correspondence between fractions (f_0, f_L, f_R) and angular distributions broken.
- Higher order corrections to decay (QED radiation or QCD in hadronic decays)
 - Decomposition in $\{A_i\}$ does not hold any more
- Angles in boson rest frame
 - Z rest frame accessible, but W more difficult to reconstruct

Polarised W+jet cross sections



Why looking at polarised W+jet with leptonic decays?

- The EW part is simple:
 - no non-resonant backgrounds
 - neutrino momentum approx. accessible (missing ET)
- Large cross section → precise measurements

Goals:

- Use W+j data to **extract the longitudinal polarisation fraction** (done before by exp.)
→ understand impact of NNLO QCD corrections (reduced scale dependence)
- Study **inclusive** (in terms of W decay products) and **fiducial** phase spaces
→ How does the sensitivity to longitudinal Ws depend on this?
Which observables have **small interference/off-shell** effects?
- Are there any differences between W+ and W-?
From PDFs and the fact that we cut on the charged lepton?

Setup W+jet: LHC @ 13 TeV

Polarised W+j production at the LHC: a study at NNLO QCD accuracy, Pellen, Poncelet, Popescu 2109.14336

Inclusive phase space:

- At least one jet with $|y(j)| \leq 2.4$ and $p_T(j) \geq 30$ GeV

Fiducial phase space:

Measurement of the differential cross sections for the associated production of a W boson and jets in proton-proton collisions at $\sqrt{s}=13$ TeV, CMS 1707.05979

- Lepton cuts: $p_T(\ell) \geq 25$ GeV, $|\eta(\ell)| \leq 2.5$ and $\Delta R(\ell, j) > 0.4$
- Transverse mass of the W: $M_T(W) = \sqrt{m_W^2 + p_T^2(W)} \geq 50$ GeV

Technical aspects:

- NNPDF31 and dynamical scale choice: $\mu_R = \mu_F = \frac{1}{2} \left(m_T(W) + \sum p_T(j) \right)$
- Implementation in STRIPPER framework (NNLO QCD subtractions) [1408.2500]
 - Narrow-Width-Approximation and OSP/Pole-Approximation
 - Matrix elements from: AvH [1503.08612], OpenLoops2 [1907.13071] (cross checks with Recola [1605.01090]) and VVamp [1503.04812]

Extraction of polarisation fractions

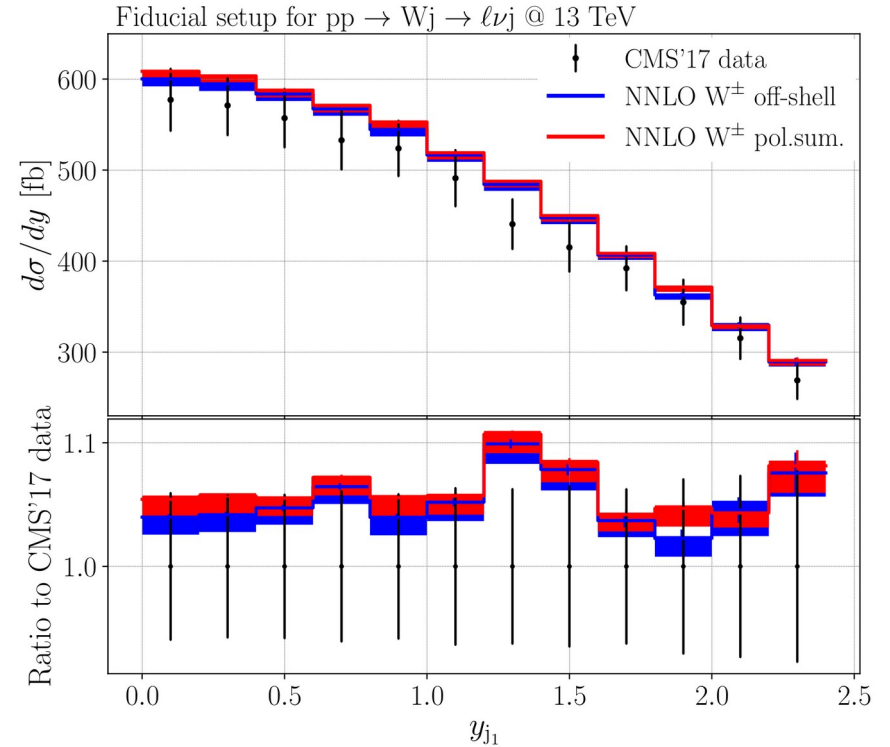
Identified 4 observables (ranges) with

→ Small interference effects (<2%)

→ Small off-shell effects (<2%)

→ Shape differences between L and T

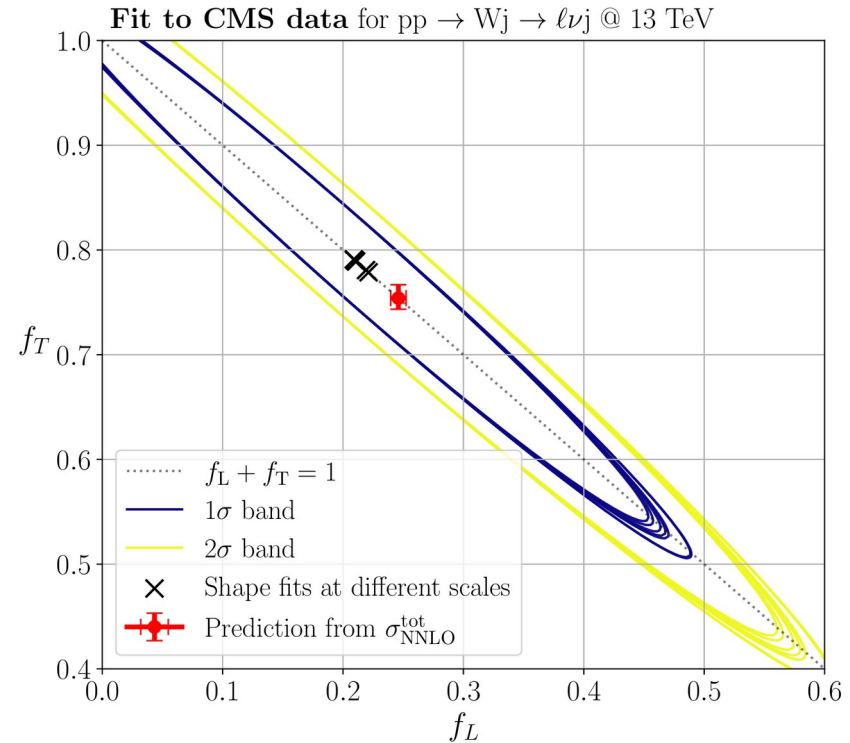
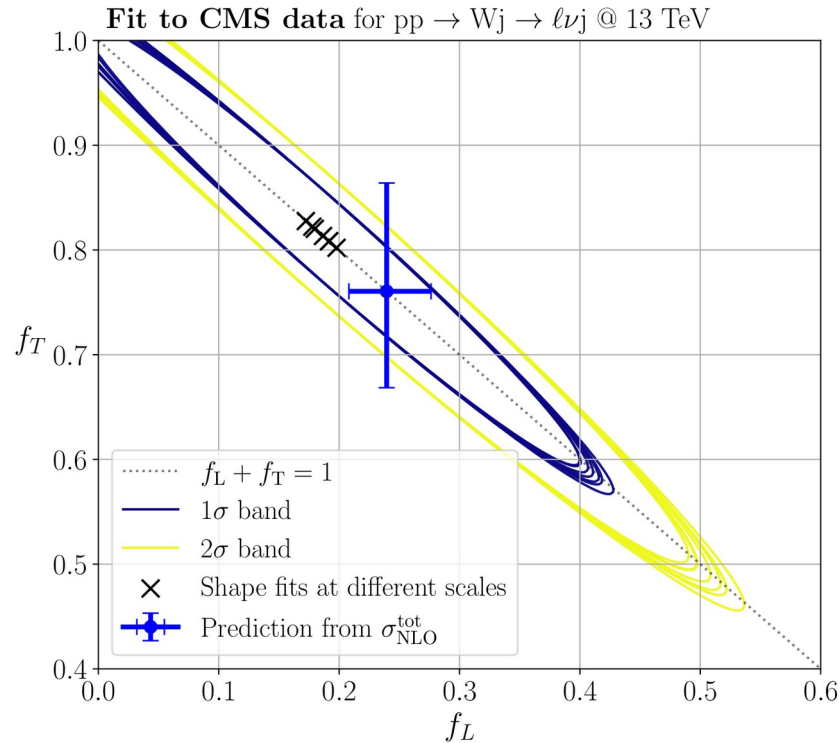
- $\Delta\phi(\ell, j_1) \geq 0.3$
- $25 \text{ GeV} \leq p_T(\ell) < 70 \text{ GeV}$
- $\cos(\theta_\ell^*) \geq -0.75$
- $|y(j_1)| \leq 2$



W+jet : fit to CMS data

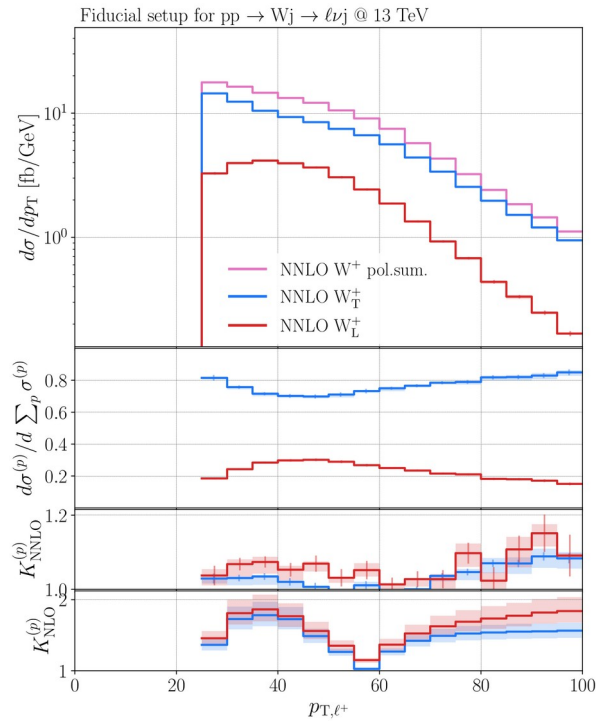
Fit to actual data, here $|y(j_1)|$

→ dominated by experimental uncertainties (no correlations available)

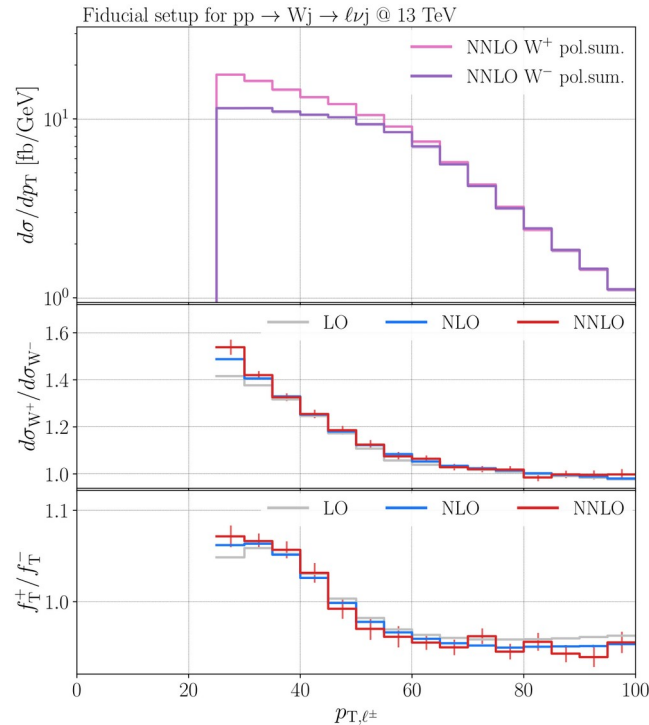


Example: lepton transverse momentum

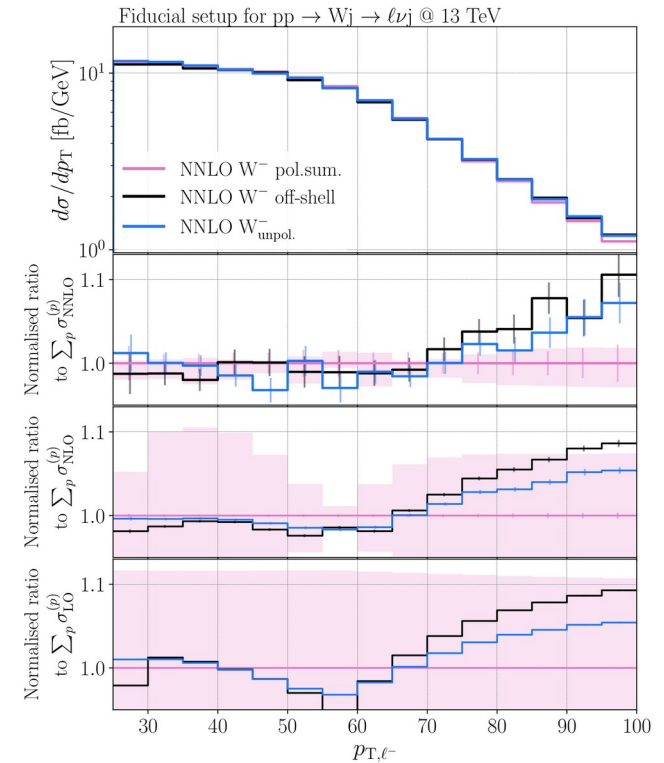
Perturbative corrections



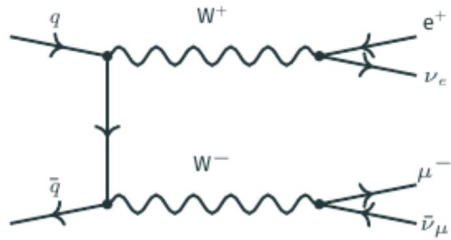
Charge differences



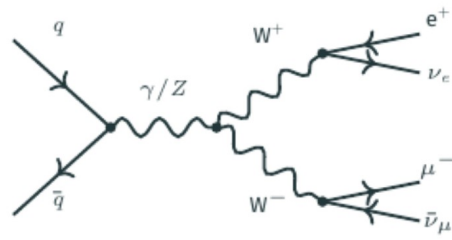
Off-shell/Interference effects



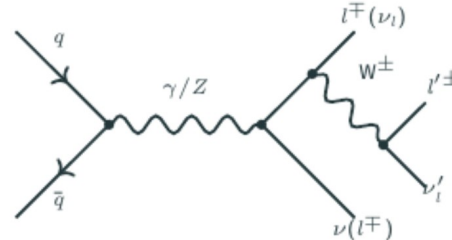
W-boson pair production



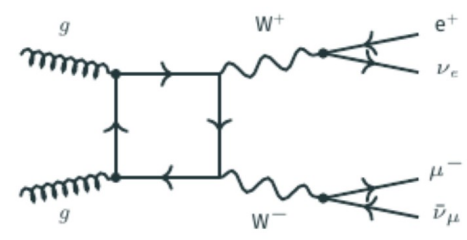
Double resonant (DR)



Double resonant (DR)

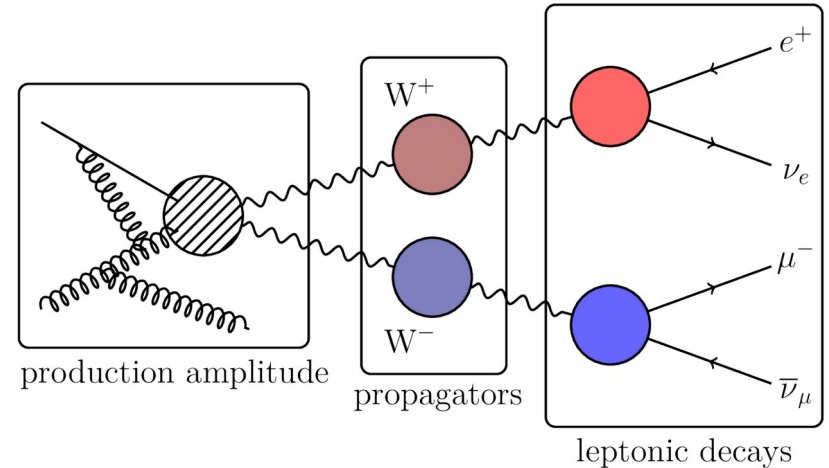


Single resonant (SR)



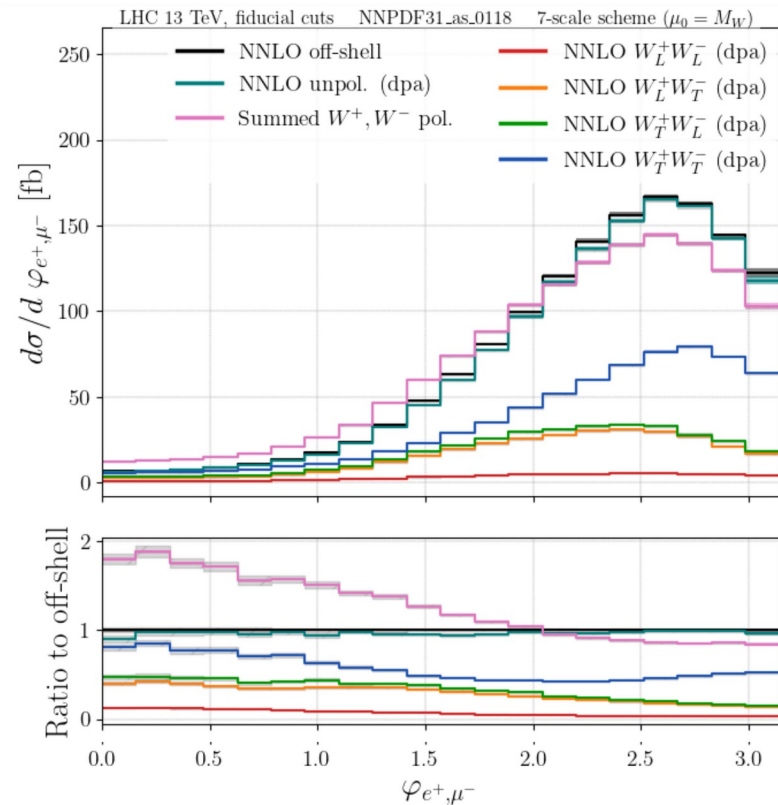
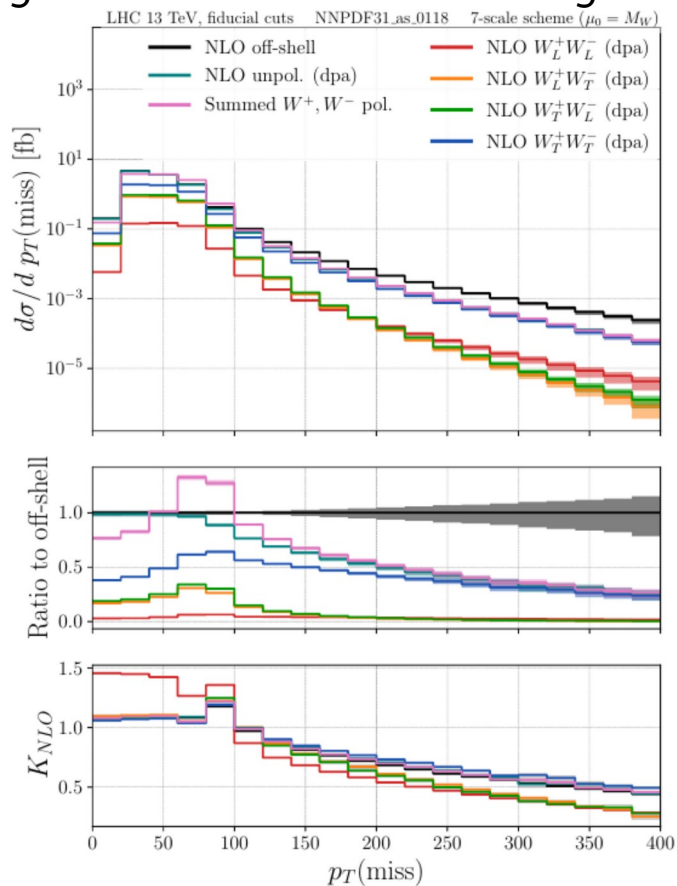
Loop-induced (LI)

- Single resonant backgrounds:
Definition of polarizations states in DPA [1710.09339] and NWA
- LI enters at NNLO \rightarrow large corrections



Interference and off-shell effects

Large off-shell effect from single-resonant contributions



Large interference effects through phase space constraints