

# New experimental techniques and new analysis ideas for probes of quark Yukawa interactions (excluding top)

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on behalf of the ATLAS and CMS collaborations

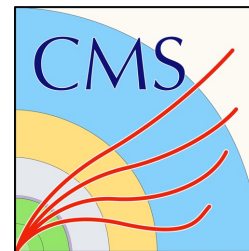
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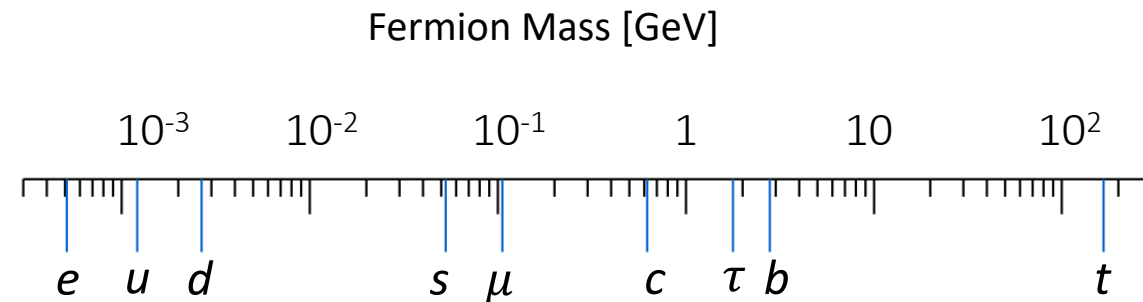
Standard Model at the LHC 2024

9<sup>th</sup> May 2024



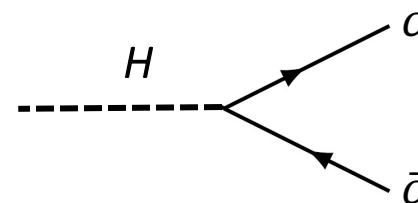
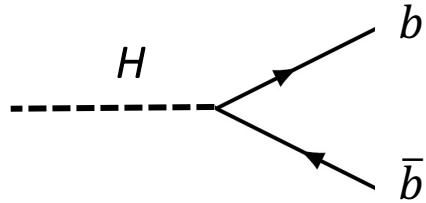
# Introduction

- Couplings of Higgs field to quarks and leptons – Yukawa couplings - are a potential source of the fermion masses
- Interaction so far only observed for 3<sup>rd</sup> generation of fermions (top, bottom and tau) and evidence found for coupling with muons
- Measurement of Yukawa couplings probe Standard Model expectation
- There is no guarantee all Higgs fermion couplings behave in a similar way
- Deviations could give insight into origin of the fermion mass hierarchy
- **Of utmost importance to measure all Higgs couplings to fermions with best possible precision!**



# Heavy quark Yukawa couplings

- Probability of Higgs boson decays to bottom/charm quarks of 58.2%/3.9% in Standard Model
- **Standard Model Higgs Yukawa coupling** to charm quarks is rather **small** ( $y_c = \sqrt{2} m_c(\mu=m_H)/v \approx 0.2 \times y_b$ )
- **One of largest contributions to  $\Gamma_H$**  (by SM expectations) **yet to be established experimentally**
- Both decay modes are susceptible to **significant modifications** in some **new physics** scenarios (e.g. two Higgs doublet models, EFTs)
- $H \rightarrow b\bar{b}$  analyses in **precision measurement** phase of cross-section and Yukawa coupling
- $H \rightarrow c\bar{c}$  analyses in search phase, **improving sensitivity** to decay and coupling



# Probe options

- Different probes of the  $Hb\bar{b}/c\bar{c}$  coupling have been proposed and investigated by the ATLAS and CMS experiments

- **Inclusive  $H \rightarrow b\bar{b}/c\bar{c}$  decays**

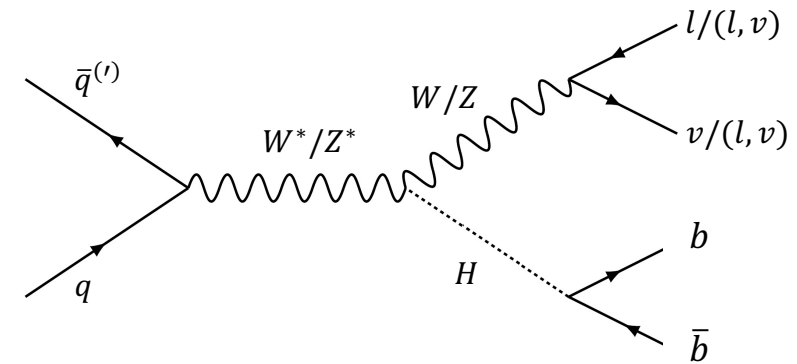
- Direct access to coupling
- Usually targeting **VH production** ( $V = W, Z$  boson)
  - **Enhanced Signal over Background ratio** w.r.t to inclusive Higgs production
  - Built around the use of **b/c-jet tagging algorithms**

### Leptonic mode

- W/Z boson decays into leptons allow for a **convenient trigger strategy**
- **Suppression of multi-jet backgrounds**

### All hadronic mode

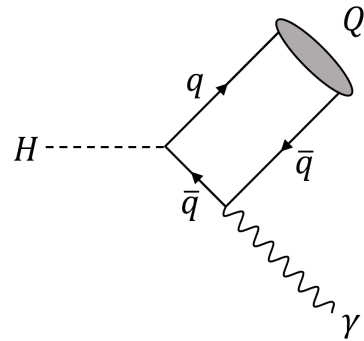
- W/Z boson decays into quarks, **increasing statistics**
- **Dominant multi-jet background**



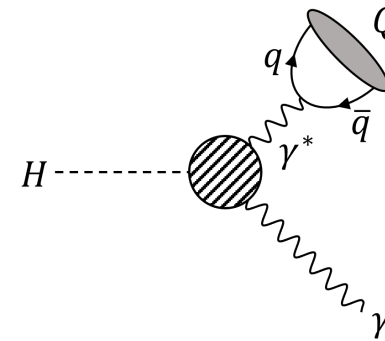


# Probe options II

- **Exclusive  $H \rightarrow J/\psi \gamma$  decays** (other rare decay modes can also probe lighter quark couplings)
  - Rare decay but **experimentally clean** probe ([Phys. Rev. D 88 \(2013\) 053003](#))
  - Direct amplitude sensitive to the **magnitude and sign of the Higgs-charm coupling**
  - **Indirect amplitude** constitutes **dominant** contribution to the Higgs boson width
  - Production rate is dominated by the indirect contribution, hindering the sensitivity to the coupling



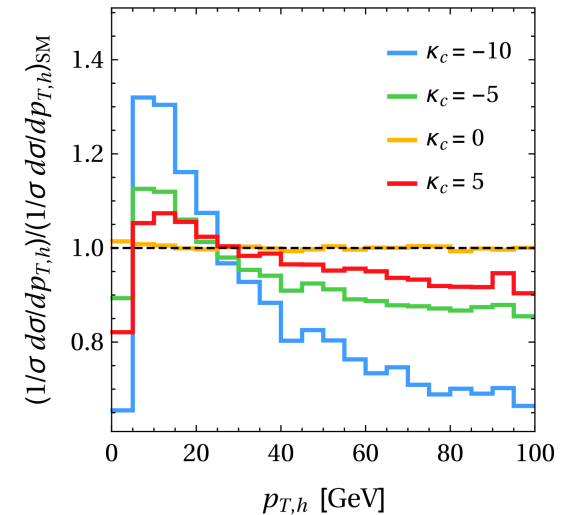
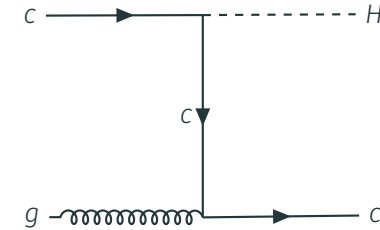
Direct contribution



Indirect contribution

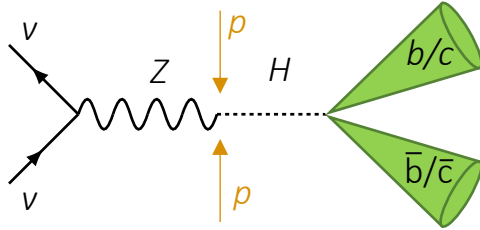
# Probe options III

- Loop induced gluon-gluon fusion and quark initiated Higgs boson production
  - Higgs production with quark contribution, e.g., via  $gg \rightarrow Hg$ ,  $gb/gc \rightarrow Hb/Hc$ ,  $b\bar{b}/c\bar{c} \rightarrow Hg$  ([Phys. Rev. Lett. 118 \(2017\) 121801](#))
  - Sensitivity to coupling via total production **differential cross-section measurements**
  - $p_T^H$  shape and normalisation will be changed in case of Yukawa couplings different to SM prediction
  - Normalisation effects coupled to Higgs width changes
  - Obtained **constraints on  $Hb\bar{b}$  ( $Hc\bar{c}$ ) coupling not as stringent (comparable)** as the ones set by inclusive Higgs decays searches
  - Subject to **different assumptions and sources of uncertainty**
  - Complementary **sensitivity to the sign** of the coupling modifiers



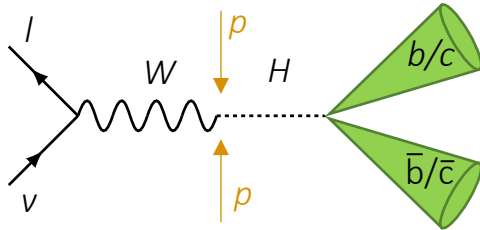
# VH Production: “Resolved” regime

Note:  
 $l = e, \mu$



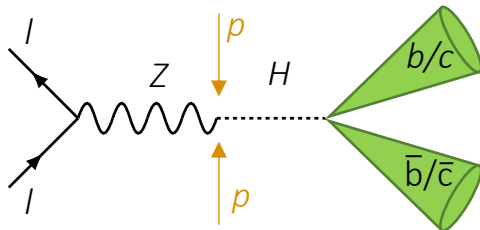
## 0 lepton channel

- Targets  $Z(\nu\nu)H(b\bar{b}/c\bar{c})$  production
- Large missing transverse energy in events (due to neutrinos)



## 1 lepton channel

- Targets  $W(l\nu)H(b\bar{b}/c\bar{c})$  production
- Large missing transverse energy
- One electron or muon



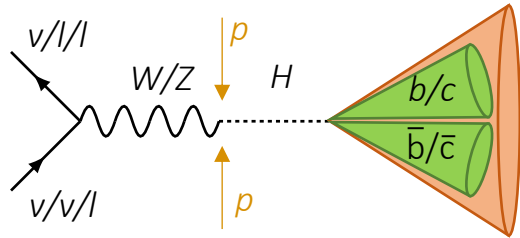
## 2 lepton channel

- Targets  $Z(l l)H(b\bar{b}/c\bar{c})$  production
- $e^+e^-$  or  $\mu^+\mu^-$  pairs

Drives sensitivity up to  $p_T^V$  of 250-400 GeV (where  $p_T^V$  is  $p_T$  of associated vector boson produced)

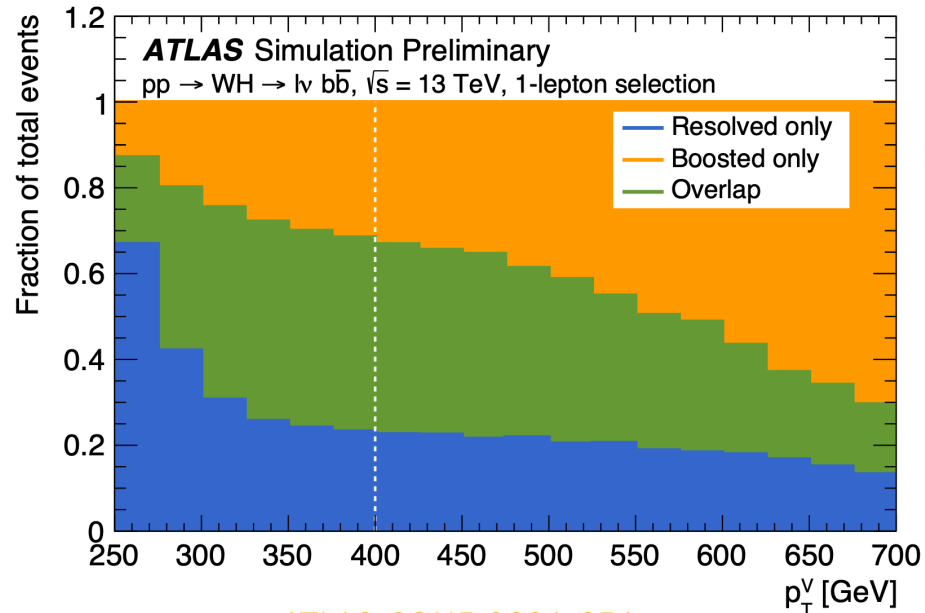
# VH Production: “Boosted” leptonic regime

Note:  
 $l = e, \mu$



## 0/1/2 lepton channel

- Targets  $Z(\nu\nu)H(b\bar{b}/c\bar{c})$ ,  $W(l\nu)H(b\bar{b}/c\bar{c})$ ,  $Z(ll)H(b\bar{b}/c\bar{c})$  production
- Large radius jet, with substructure exploited



## Resolved/boosted overlap

- Slightly different treatments between ATLAS and CMS
- Common factor – priority in event allocation given to resolved regime in overlap region as improves expected sensitivity and decreases uncertainties
- Events used only once

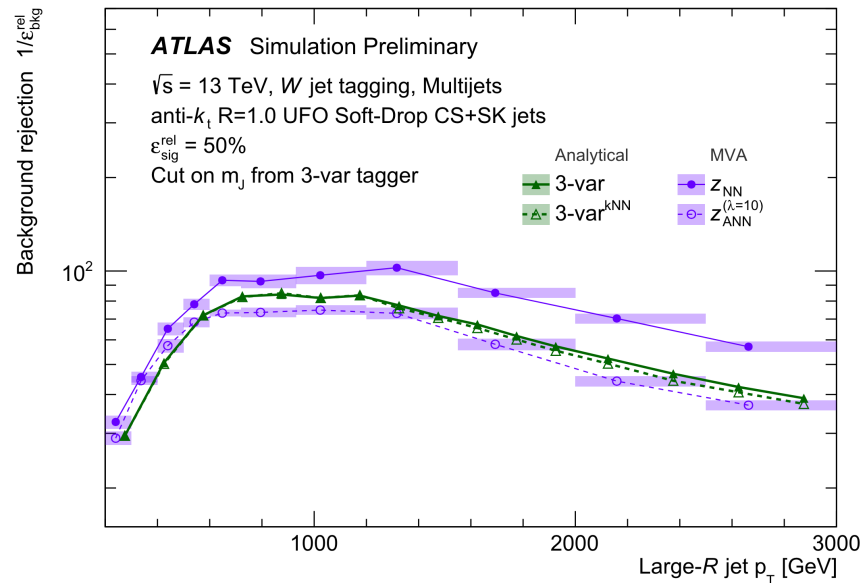
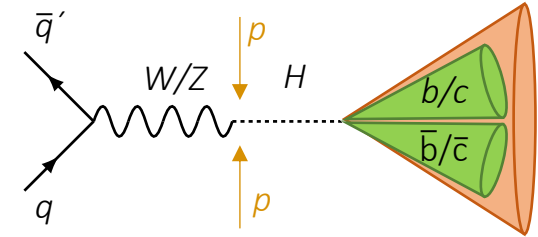
# VH Production: “Boosted” hadronic regime

## Hadronic channel

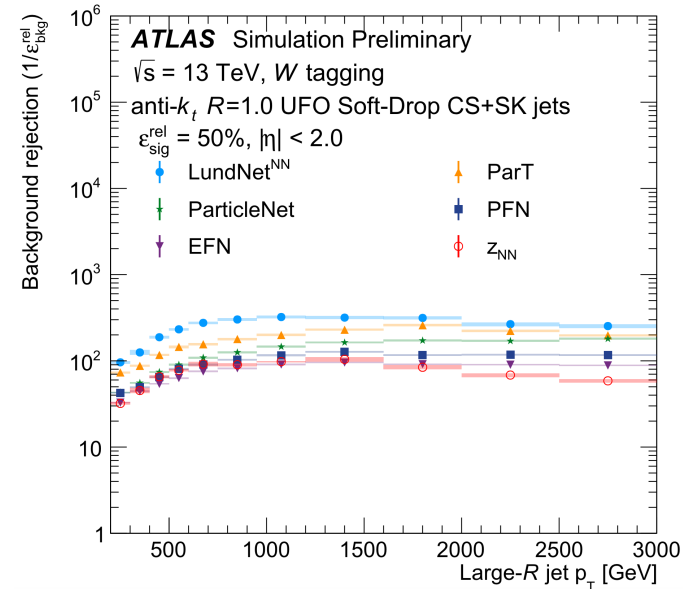
- Targets  $V(qq)H(b\bar{b}/c\bar{c})$  production
- Vector boson tagged via dedicated taggers

## Hadronically decaying W/Z boson tagging

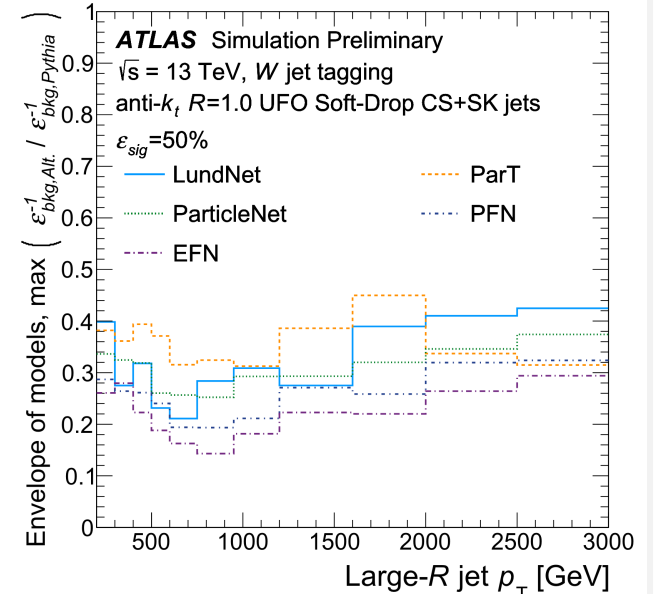
- First ATLAS hadronic VH,  $H(b\bar{b})$  analysis ([PhysRevLett.132.131802](https://arxiv.org/abs/1311.1802)) made use of a cut-based tagger, exploiting jet mass, number of tracks and substructure variables as energy correlation function ratios
- Performance improvements found in **DNN/GNN and transformers** architectures using jet constituents' energy and kinematics



JETM-2022-006

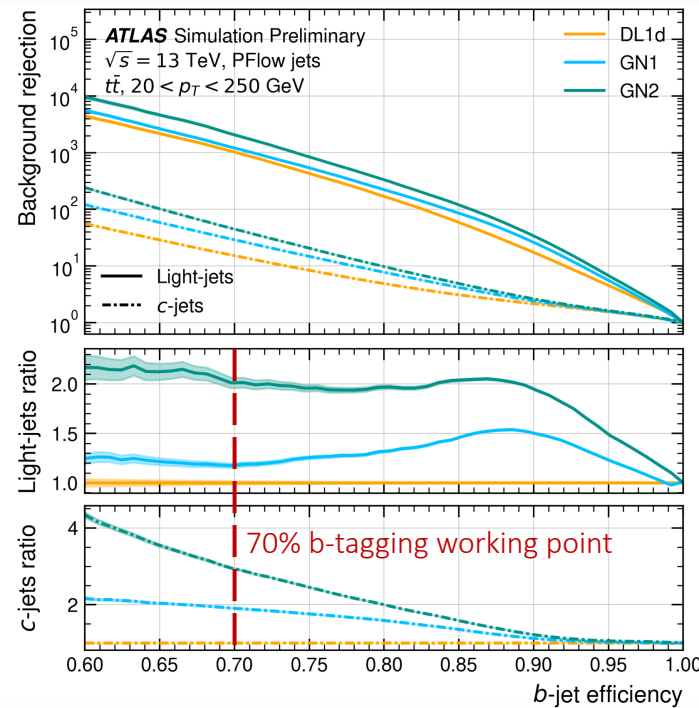


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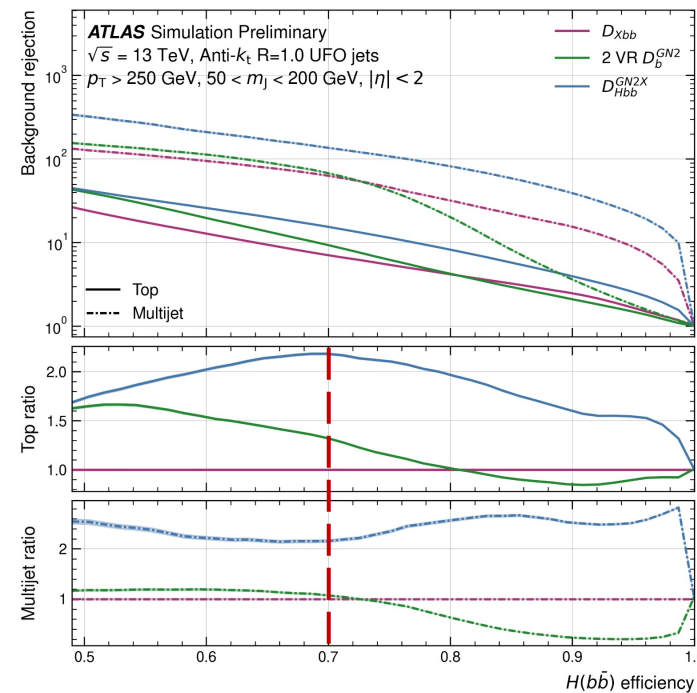


# VH Production: Jet Heavy Flavour Tagging

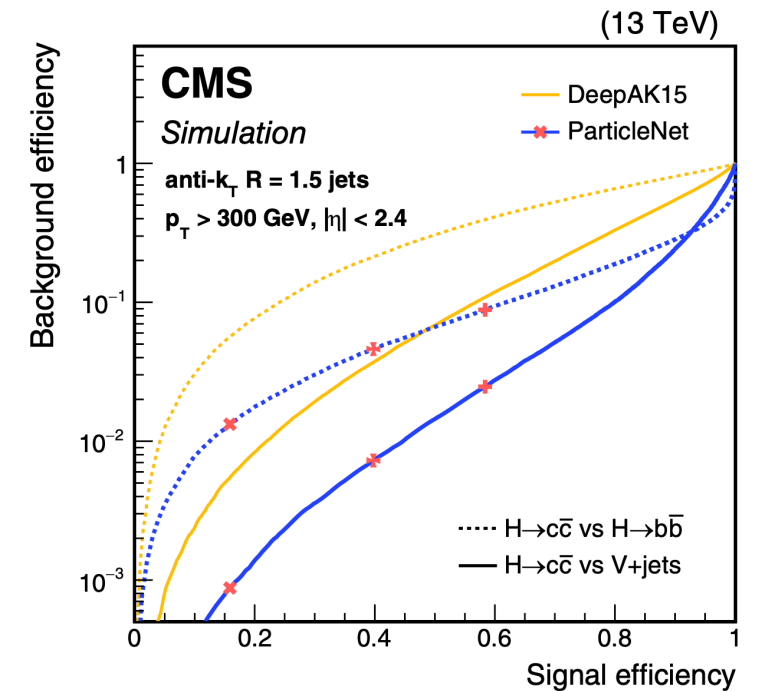
- Bottom or charm Higgs candidate jets identified via jet flavour tagging algorithms
- Taggers optimised separately for resolved or boosted regime
- Targeting both Higgs to bottom and charm decays
- DNN/CNN-based taggers being used in most of latest analyses
- Further improvements found in graph and transformer neural network architectures, by better exploiting jet internal structure



[ATLAS GNN Flavour Tagging](#)



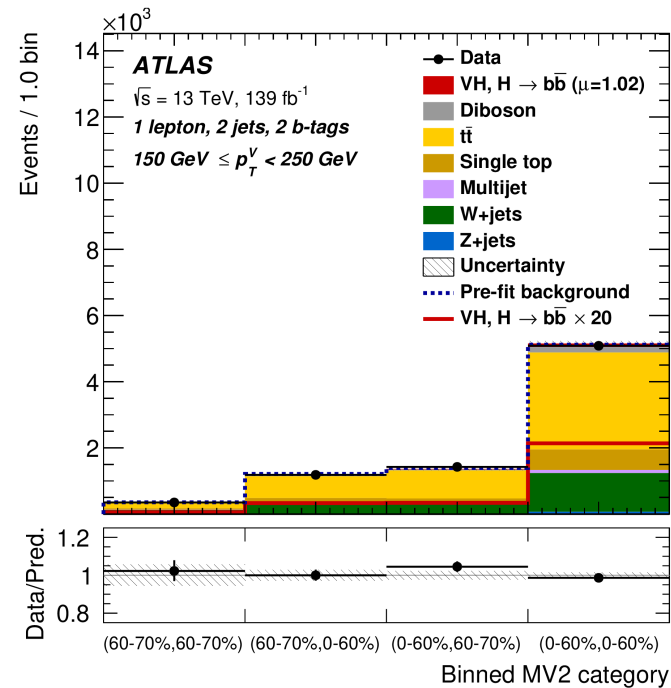
[ATL-PHYS-PUB-2023-021](#)



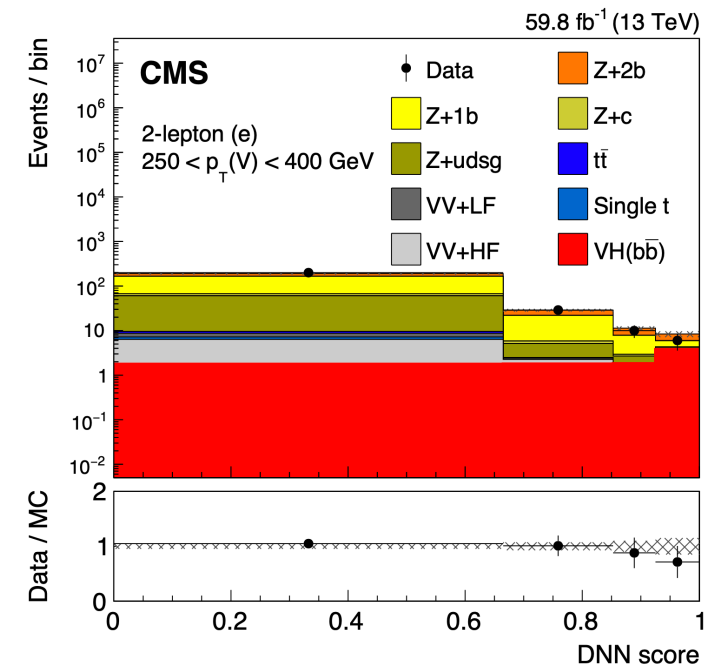
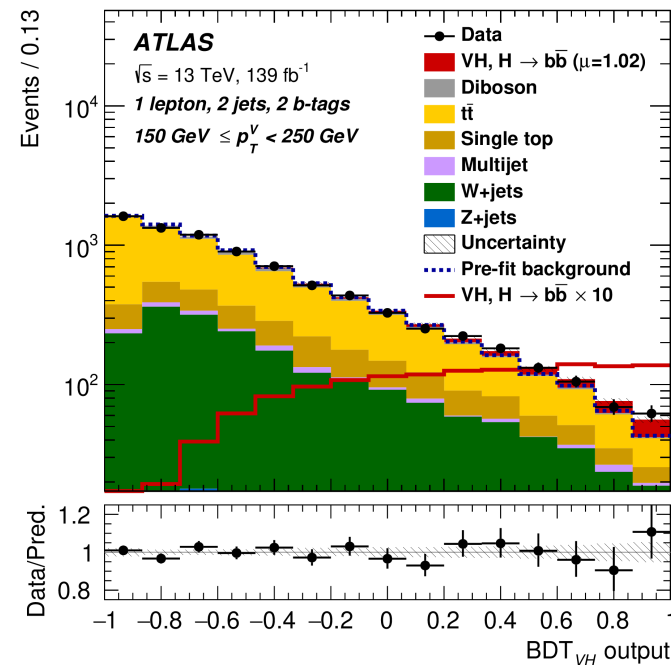
[PhysRevLett.131.061801](#)

## Key aspects

- Categorisation of events for optimised selections and signal sensitivity
- Data-driven methods/control regions to estimate  $t\bar{t}$ , W/Z+jets and multijet backgrounds
- Flavour tagging discriminants used along with jet mass and transverse momentum to discriminate VH production from backgrounds via BDTs/NNs



[Eur. Phys. J. C \(2021\) 81:178](#)



[arXiv:2312.07562](#)

# VH Production: Main uncertainties

## V(lep/had)H(bb̄)

- Starting to be dominated by **systematic uncertainties in resolved regime**
- Statistical uncertainties** still very relevant in boosted regime

## V(lep)H(c̄c̄)

- Statistical and systematic uncertainties of the same magnitude

## Leading systematic uncertainties

- Background modelling:** tt̄bar, W/Z+jets, multijet
- Signal modelling**
- Jet flavour tagging uncertainties**
- Jet energy scale and resolution uncertainties**
- Statistical uncertainty from limited size of MC samples**

→ More data will improve results directly and indirectly (better modelling, reduced uncertainties)

[arXiv:2312.07562](https://arxiv.org/abs/2312.07562)

CMS V(lep)H(bb̄)

	$\Delta\mu$
Background (theory)	+0.043 -0.043
Signal (theory)	+0.088 -0.059
MC sample size	+0.078 -0.078
Simulation modeling	+0.059 -0.059
b tagging	+0.050 -0.046
Jet energy resolution	+0.036 -0.028
Int. luminosity	+0.032 -0.027
Jet energy scale	+0.025 -0.025
Lepton ident.	+0.008 -0.007
Trigger ( $\vec{p}_T^{\text{miss}}$ )	+0.002 -0.001

[PhysRevLett.132.131802](https://arxiv.org/abs/1312.13180)

ATLAS boosted V(had)H(bb̄)

Uncertainty source	$\delta\mu$
Signal modeling	+0.10 -0.02
MC statistical uncertainty	+0.13 -0.13
Instrumental (pileup, luminosity)	+0.012 -0.004
Large-R jet	+0.13 -0.14
Top-quark modeling	+0.14 -0.15
Other theory modeling	+0.05 -0.03
H → b̄b̄ tagging	+0.52 -0.23
Multijet estimate (TF uncertainty)	+0.52 -0.41
Multijet modeling (TF vs BDT)	+0.14 -0.18
Total systematic uncertainty	+0.80 -0.61
Signal statistical uncertainty	+0.60 -0.60
Z + jets normalization	+0.42 -0.20
Total statistical uncertainty	+0.63 -0.63
Total uncertainty	+1.02 -0.88

[Eur. Phys. J. C \(2021\) 81:178](https://arxiv.org/abs/2103.13032)

ATLAS resolved V(lep)H(bb̄)

Source of uncertainty	VH
Total	0.177
Statistical	0.115
Systematic	0.134
Statistical uncertainties	
Data statistical	0.108
t̄t̄ eμ control region	0.014
Floating normalisations	0.034
Experimental uncertainties	
Jets	0.043
E <sub>T</sub> <sup>miss</sup>	0.015
Leptons	0.004
b-tagging	b-jets 0.045 c-jets 0.035 light-flavour jets 0.009
Pile-up	0.003
Luminosity	0.016
Theoretical and modelling uncertainties	
Signal	0.072
Z + jets	0.032
W + jets	0.040
t̄t̄	0.021
Single top quark	0.019
Diboson	0.033
Multi-jet	0.005
MC statistical	0.031



# H(b $\bar{b}$ ) Overview

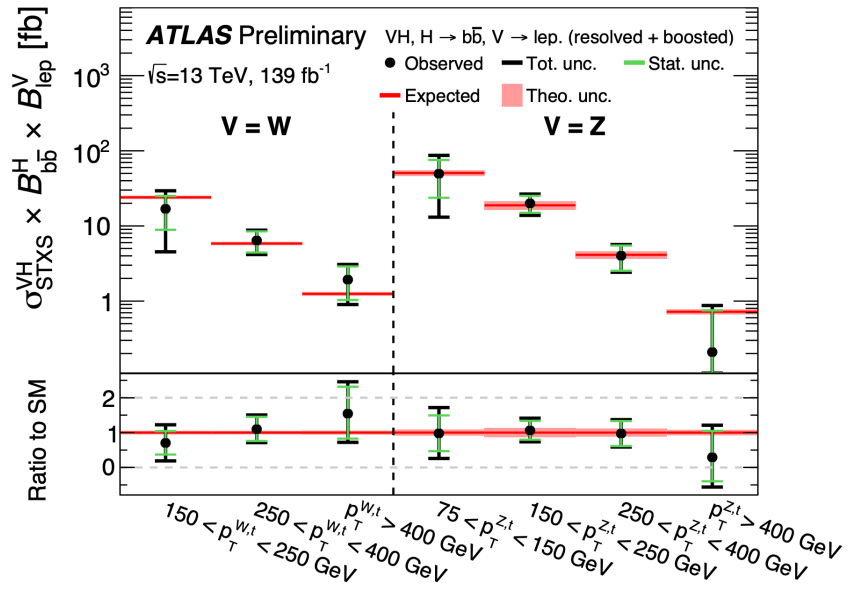
## VH(b $\bar{b}$ ) analyses

Analysis	Fitted signal strength $\mu_{VH} (H \rightarrow b\bar{b})$
ATLAS resolved V(l $\bar{e}$ p)H(b $\bar{b}$ )	$1.02^{+0.18}_{-0.17} = 1.02^{+0.12}_{-0.11}(\text{stat.})^{+0.14}_{-0.13}(\text{syst.})$
ATLAS boosted V(l $\bar{e}$ p)H(b $\bar{b}$ )	$0.72^{+0.39}_{-0.36} = 0.72^{+0.29}_{-0.28}(\text{stat.})^{+0.26}_{-0.22}(\text{syst.})$
ATLAS resolved+boosted V(l $\bar{e}$ p)H(b $\bar{b}$ )	$1.00^{+0.18}_{-0.17} = 1.02^{+0.12}_{-0.11}(\text{stat.})^{+0.14}_{-0.13}(\text{syst.})$
CMS resolved+boosted V(l $\bar{e}$ p)H(b $\bar{b}$ )	$1.15^{+0.22}_{-0.20}$
ATLAS boosted V(had)H(b $\bar{b}$ )	$1.4^{+1.0}_{-0.9} = 1.4^{+0.6}_{-0.6}(\text{stat.})^{+0.8}_{-0.7}(\text{syst.})$

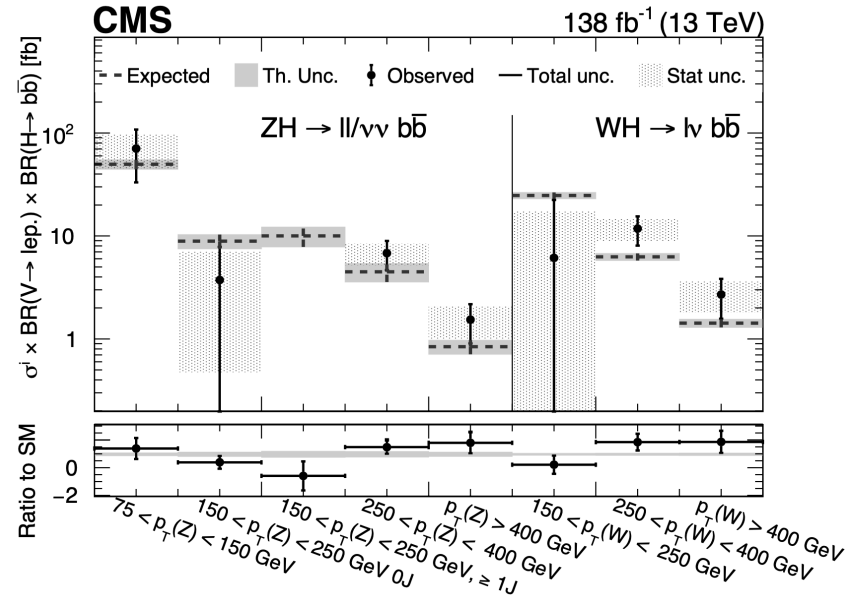
- Moving towards precision measurements
  - SXTS and differential cross-section measurements being undertaken

# STXS Measurements

- Simplified Template Cross-Sections (STXS) involve categorisation of events in bins designed to separate different production modes



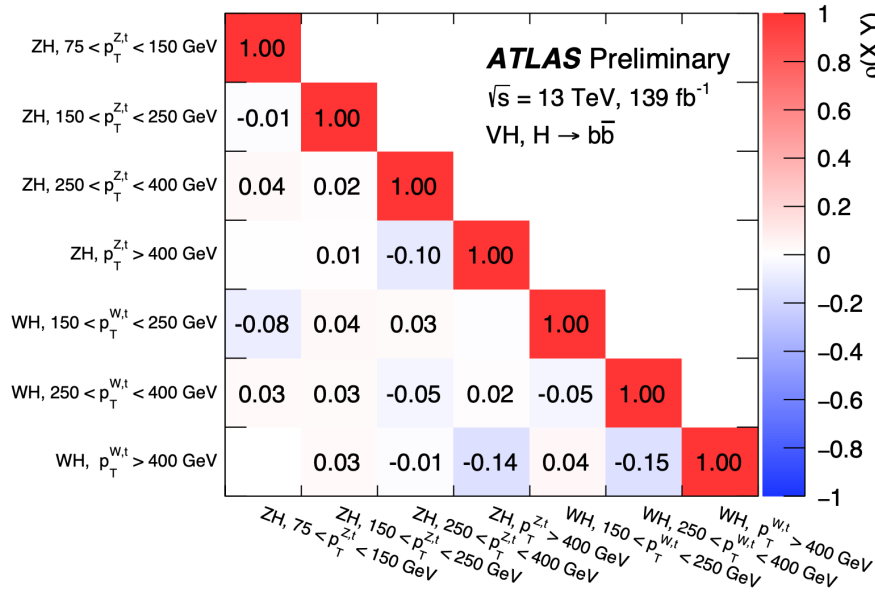
[ATLAS-CONF-2021-051](#)



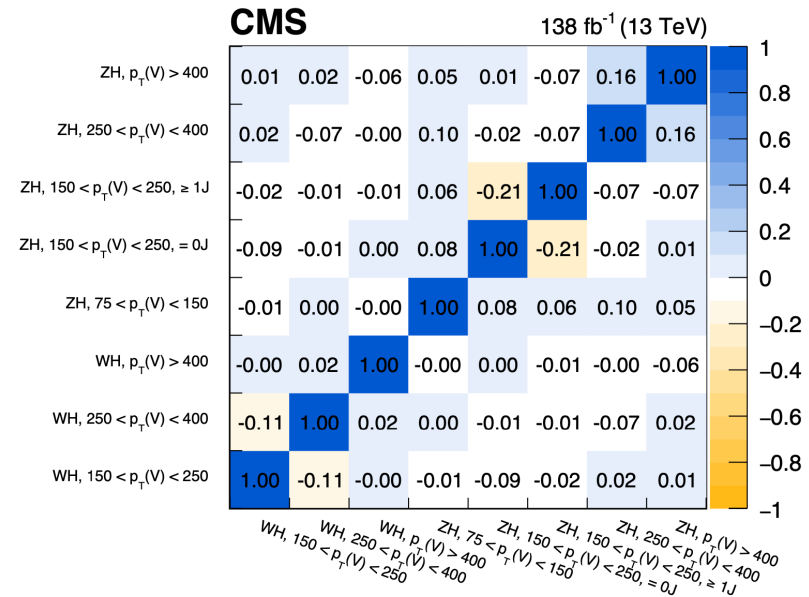
[arXiv:2312.07562](#)

# STXS Measurements II

- Cross-contamination generally negligible between different STXS production bins
- Some WH contribution in ZH bins via 0 lepton channel



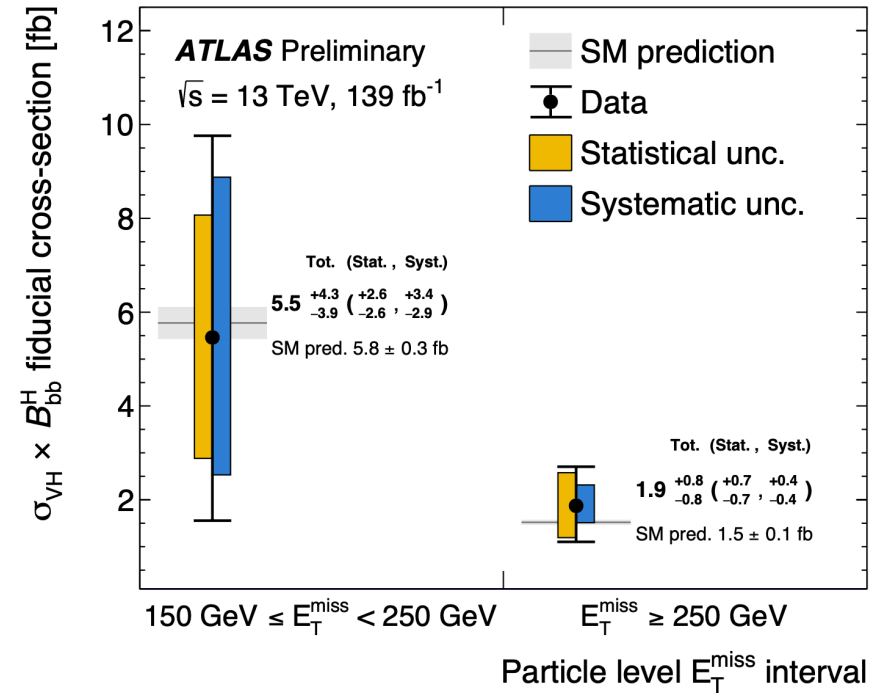
[ATLAS-CONF-2021-051](#)



[arXiv:2312.07562](#)

# Differential Cross-section Measurements

- Current statistics and understanding of systematic uncertainties allow for differential cross-section measurements to be undertaken
- ATLAS performed such measurement in the 0 lepton channel ( $Z(\nu\nu)H(b\bar{b})$ )
  - Remaining channels constrain the SM background via one-binned regions
- Detector effects corrected for via profile-likelihood unfolding procedure
- Impacted by **usual systematic uncertainties**, with some reduced constraints due to one-binned regions
- Results also affected by migration of events from outside fiducial phase-space



[ATLAS-CONF-2022-015](#)

# H(c $\bar{c}$ ) Overview

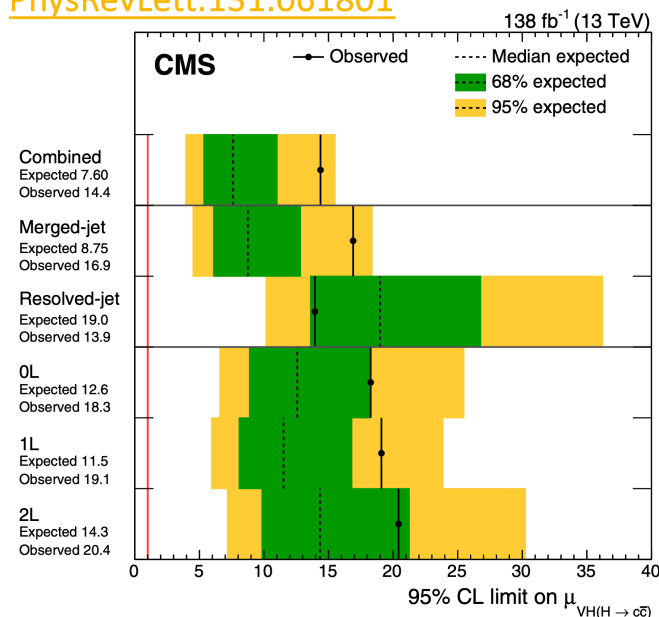
## H(c $\bar{c}$ ) analyses

Analysis	Observed (Expected) 95% CL limit on $\mu_{VH} (H \rightarrow c\bar{c})$
ATLAS resolved V(l $\bar{e}$ )H(c $\bar{c}$ ) H(b $\bar{b}$ )	26.0 $\times$ SM ( $31_{-8}^{+12}$ $\times$ SM)
CMS resolved+boosted V(l $\bar{e}$ )H(c $\bar{c}$ )	14.4 $\times$ SM ( $7.6_{-2.3}^{+3.4}$ $\times$ SM)
CMS boosted ggF H(c $\bar{c}$ )	47.0 $\times$ SM (39 $\times$ SM)

- Improving sensitivity

- Use of MVA techniques in analysis design and jet flavour tagging methods bringing significant enhancements in sensitivity
- Boosted regime can further exploit such tools to bring additional relevant contributions

[PhysRevLett.131.061801](https://arxiv.org/abs/1306.6580)



# $\kappa$ Framework

- Measure deviations of Higgs couplings from SM
- Inspired by leading order diagrams of Higgs couplings
- Assumes Higgs boson resonance at 125 GeV and narrow width approximation

$$\sigma \times \text{BR}(i \rightarrow H \rightarrow f) = \frac{\sigma_i \times \Gamma_f}{\Gamma_H}$$

+

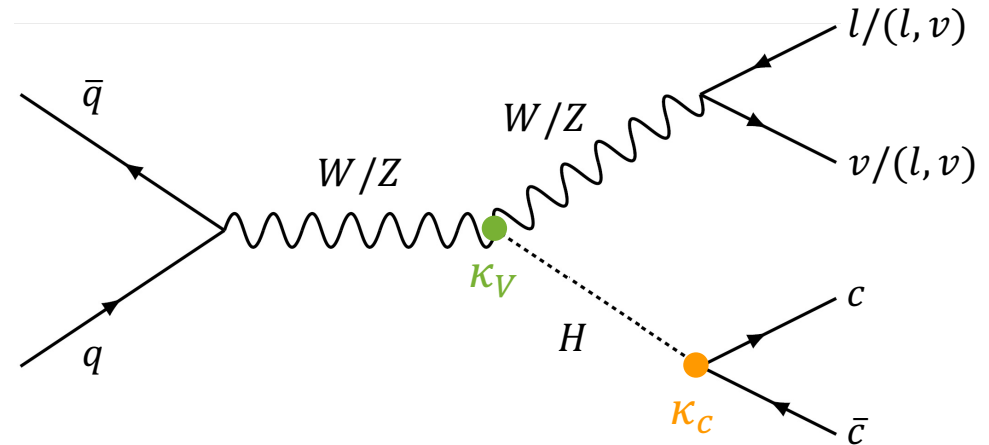
$$y_j \rightarrow \kappa_j y_j$$

↓

$$\sigma \times \text{BR}(i \rightarrow H \rightarrow f) = \frac{\sigma_i^{SM} \kappa_i^2 \times \Gamma_f^{SM} \kappa_f^2}{\Gamma_H^{SM} \kappa_H^2}$$

↓

$$\sigma \times \text{BR}(V \rightarrow H \rightarrow c\bar{c}) = \sigma_{SM} \times \text{BR}(V \rightarrow H \rightarrow c\bar{c})_{SM} \times \frac{\kappa_V^2 \kappa_c^2}{\kappa_H^2}, \quad \kappa_H^2 = \frac{\sum_j \Gamma_j^{SM} \times \kappa_j^2}{\Gamma_H^{SM}}$$

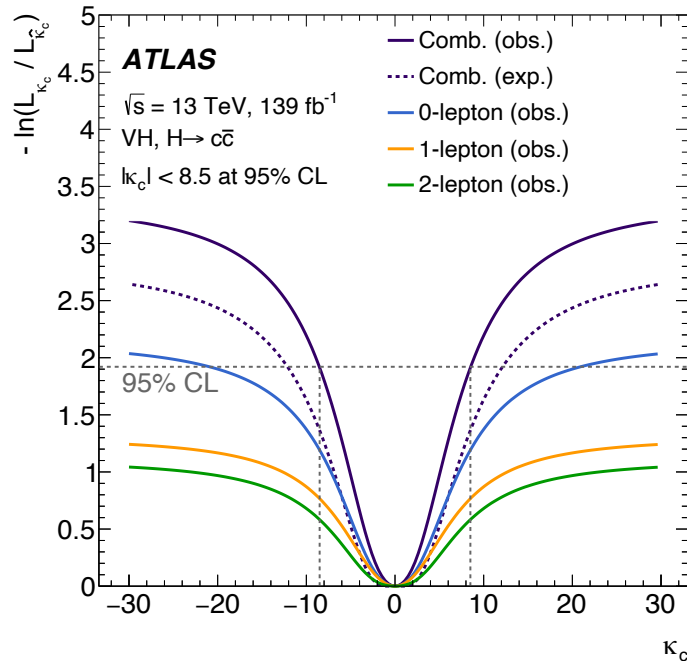


# $\kappa_C$ Interpretation

- $V(\text{lep})H(c\bar{c})$  analyses sensitive to exclude some parts of  $\kappa_C$  parameter space (parametrisation below used)

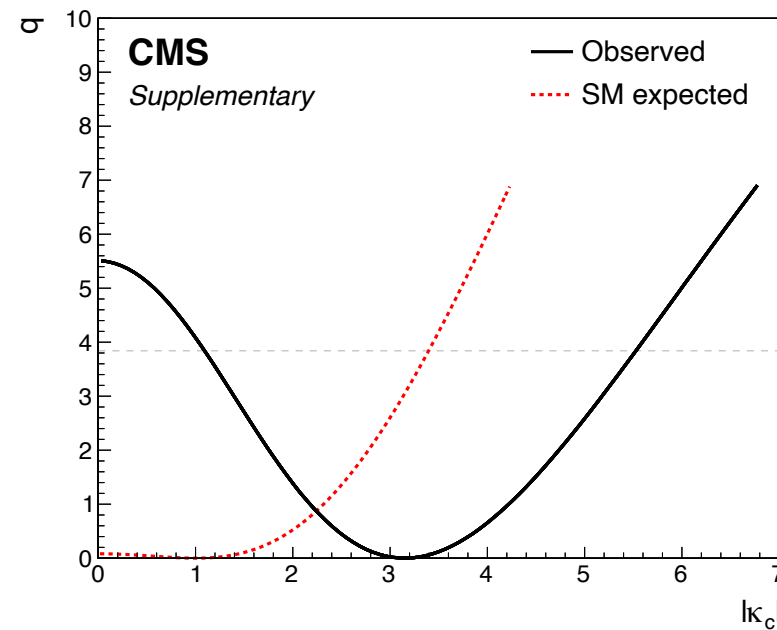
- $$\mu = \frac{\kappa_C^2}{[(1 - BR_{Hcc}) + BR_{Hcc} * \kappa_C^2]}$$
 (other coupling modifiers set to 1, no BSM contributions to Higgs width)

[Eur. Phys. J. C \(2022\) 82:717](#)



$|\kappa_C| \leq 8.5 \text{ at } 95\% \text{ CL}$

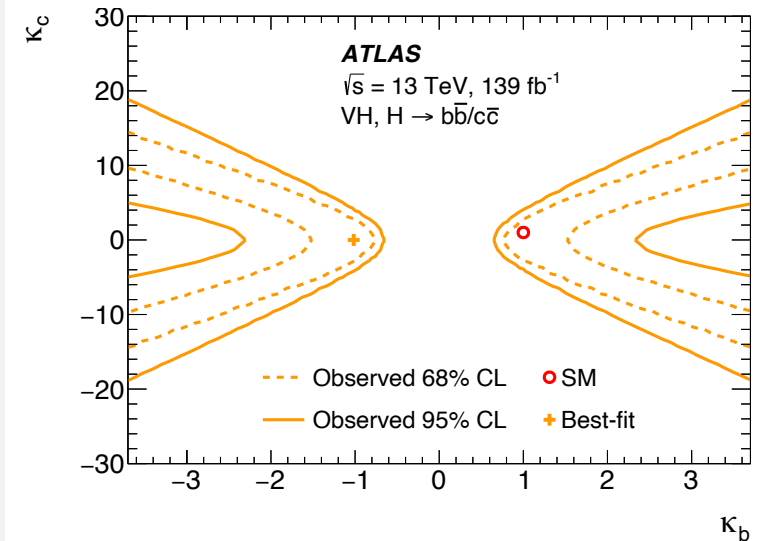
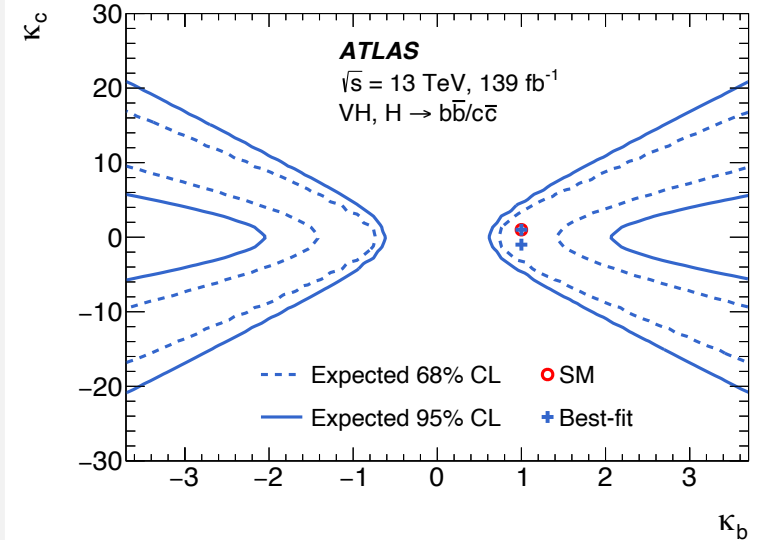
[PhysRevLett.131.061801](#)



$|\kappa_C| \leq 3.4 \text{ (} 1.1 < \kappa_C < 5.5 \text{) at } 95\% \text{ CL}$

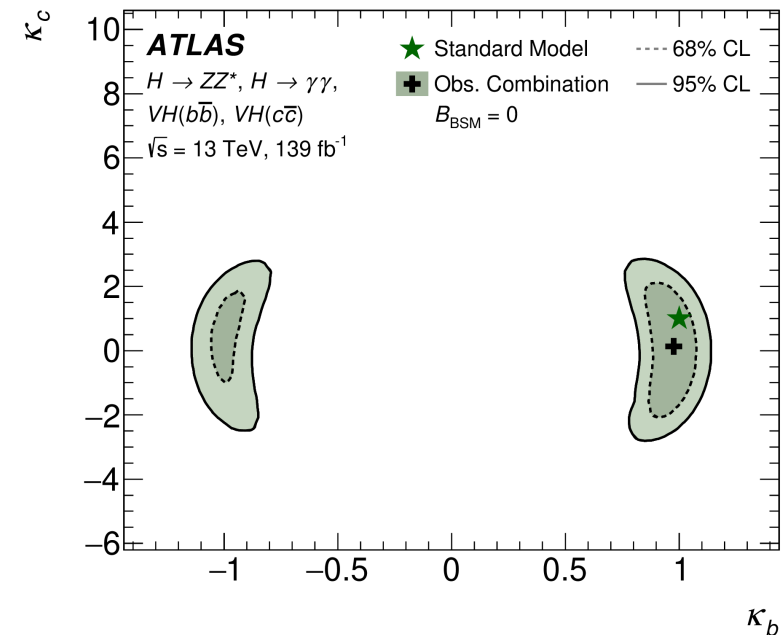
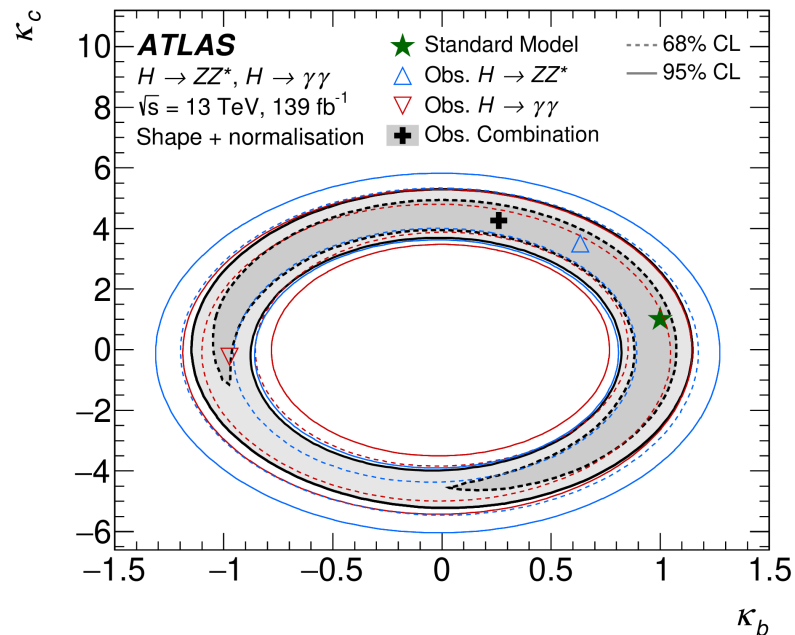
# $\kappa_c$ and $\kappa_b$ Interpretation

- $V(\text{lep})H(b\bar{b})$  and  $V(\text{lep})H(c\bar{c})$  combined likelihood parametrised in terms of both  $\kappa_c$  and  $\kappa_b$
- Other coupling modifiers set to their SM predictions
- Gluon-initiated ZH production has no explicit parametrisation as function of  $\kappa_c$ 
  - Only taken into account for  $\kappa_b$
  - Only has an impact for large values of  $\kappa_b$  ( $\kappa_b \gtrsim 10$ )
- Non Drell-Yan-like  $q\bar{q} \rightarrow VH$  production also neglected given the low values of  $\kappa_c$  that the combination probes
- **Observed best fit value  $(\kappa_b, \kappa_c) = (-1.02, 0)$**
- Difference in value of the log-likelihood between best-fit value and  $(\kappa_b, \kappa_c) = (+1.02, 0)$  is 0.02
  - Analysis has little power to constrain the sign of  $\kappa_b$
  - Small likelihood asymmetry coming from b-quark loop contributions to  $g\bar{g} \rightarrow ZH$  production





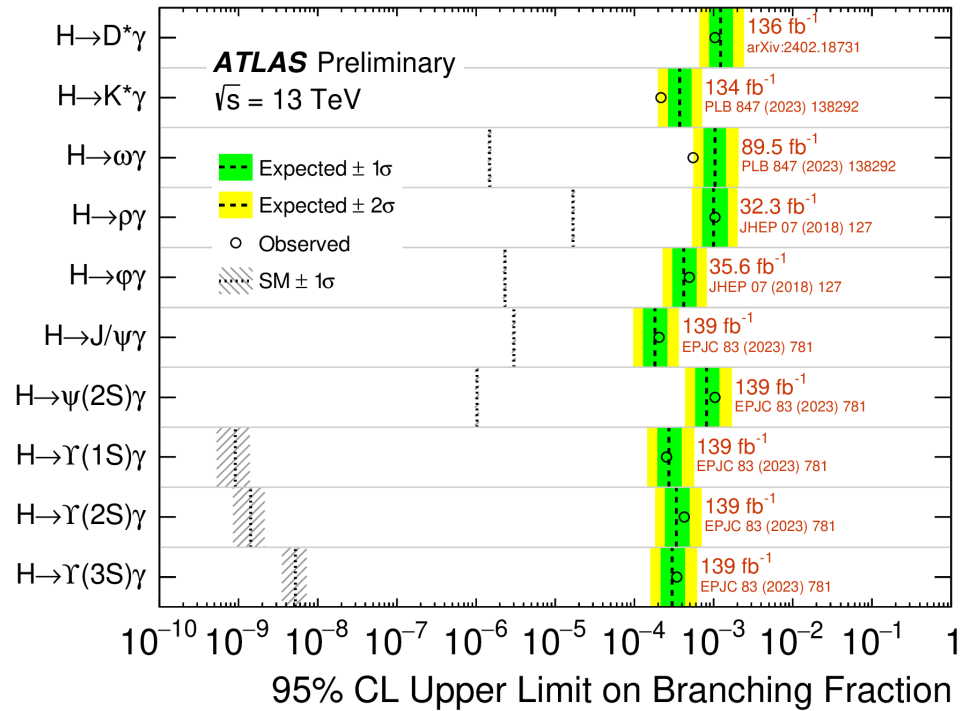
- Precision measurements of  $p_T^H$  cross-section in  $H \rightarrow ZZ^*$  and  $H \rightarrow \gamma\gamma$  channels provide **complimentary** information on couplings
- Individual analyses **dominated by statistical uncertainty**
- **Statistical combination between them** and in addition with  $V(\text{lep})H(b\bar{b})$  and  $V(\text{lep})H(c\bar{c})$  analyses allows for improved constraints
- Results for **shape and normalisation effects considered** shown here, different assumptions on backup



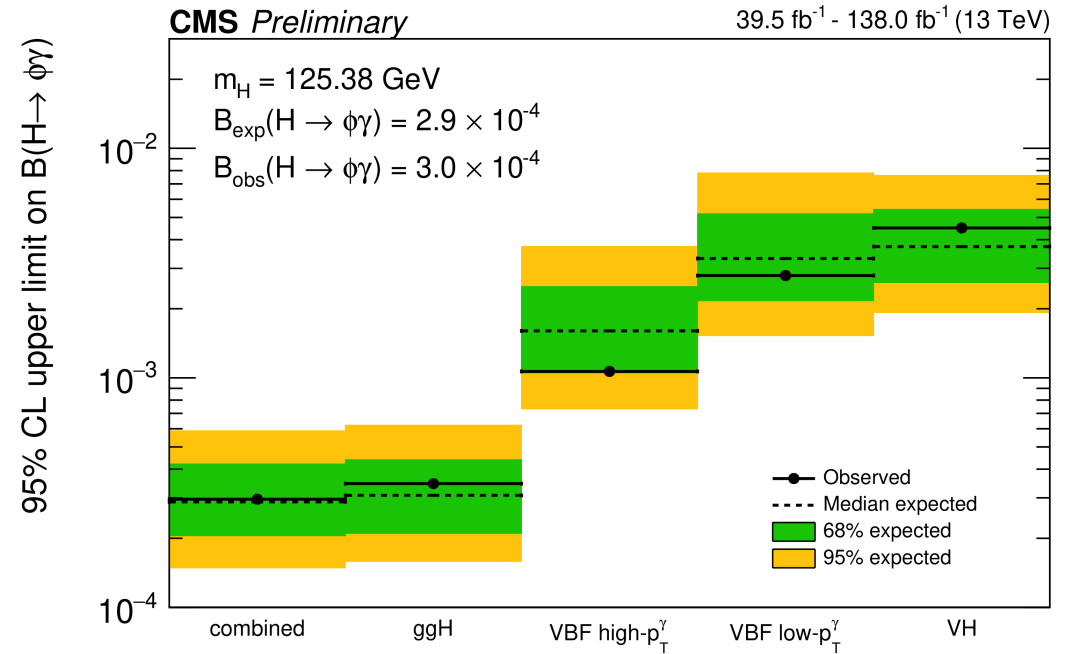
# Probing lighter quark Yukawas

$q\bar{q} \rightarrow H(\rightarrow WW) + \gamma$  differential XS measurement as alternative probe: [PhysRevLett.132.121901](https://arxiv.org/abs/2302.11901), in backup

- Rare **exclusive decays** of the Higgs boson into a meson and a photon allow **access to coupling to lighter flavour quarks** (up, down, strange) in addition to bottom and charm quarks
- Also used as probe of potential **flavour-violating Higgs boson interactions**
- **Dedicated triggers, different production modes, event selection with BDTs and background modelling enhancements** improving sensitivity



[ATL-PHYS-PUB-2023-004](https://arxiv.org/abs/2302.004)



[CMS-PAS-HIG-23-005](https://arxiv.org/abs/2302.005), more CMS results on rare decays in backup

# Summary

- **Different probes sensitive to Higgs quark Yukawa couplings are available**
  - Leading sensitivity to **heavy quark Yukawa** couplings currently via **Higgs production in association with W/Z bosons**
  - **Loop induced ggF and quark initiated Higgs production and exclusive Higgs boson decays into meson+photon** provide **complimentary information** on these couplings and give insight into **lighter quark couplings**
- **Machine-learning improvements in analysis design and jet tagging techniques leading to continuous enhancements in sensitivity**
  - Exploited both in **resolved** and **merged** jet topologies
  - Possible **tagging of Higgs candidates and/or W/Z bosons**
  - MVA discrimination of **Higgs signal** with respect to **backgrounds**
- **Main uncertainties related to modelling, flavour tagging and jet contributions or statistics**
- **Different suite of measurements currently possible, from STXS bins to coupling modifiers**

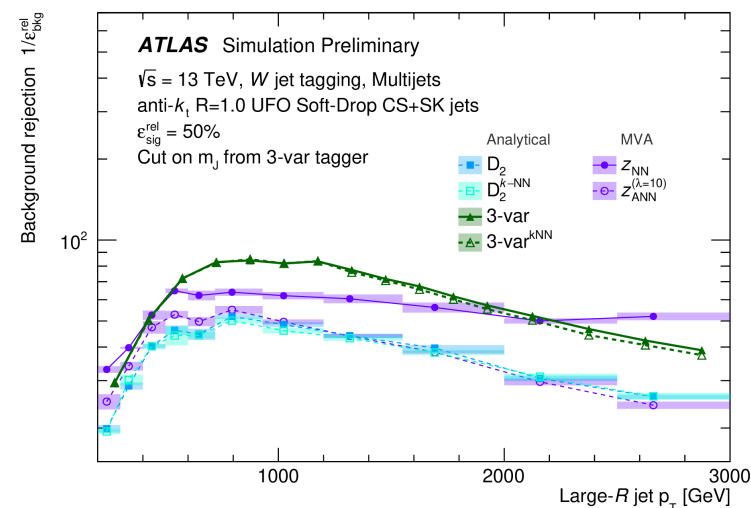
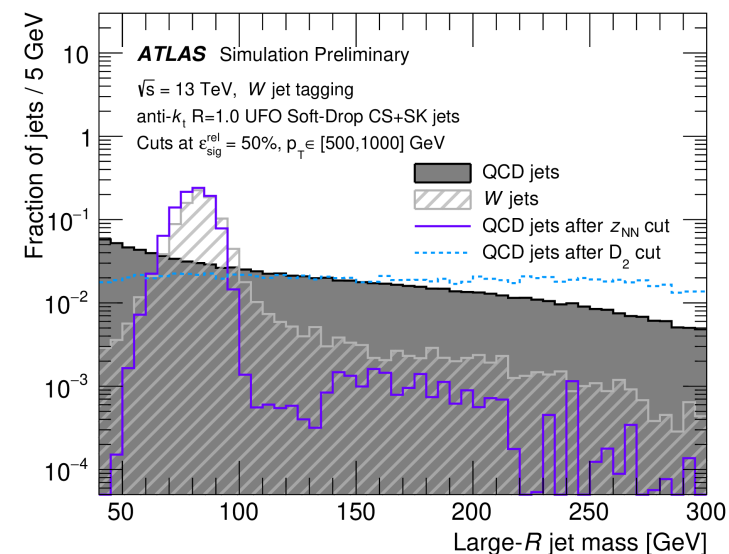
# Backup

# Vector boson DNN tagger

- DNN trained on high-level jet substructure quantities

Variable	Description
$D_2, C_2$	Energy correlation ratios
$\tau_{21}$	$N$ -subjettiness
$R_2^{\text{FW}}$	Fox-Wolfram moment
$\mathcal{P}$	Planar flow
$a_3$	Angularity
$A$	Aplanarity
$Z_{\text{cut}}, \sqrt{d_{12}}$	Splitting scales
$Kt\Delta R$	$k_t$ -subjettiness $\Delta R$

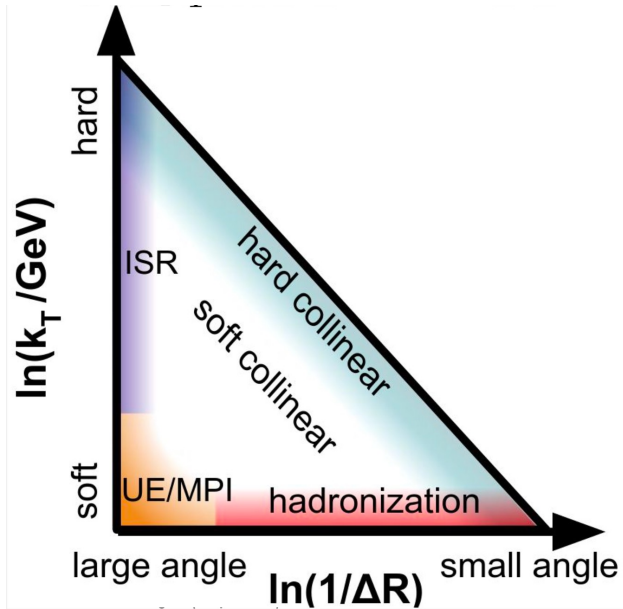
- Algorithm decorrelated from jet mass to prevent bias in background jet-mass distribution from being around the  $W$  boson mass
- Achieved via sequential training of jet classification DNN and jet mass inferring adversarial NN (ANN)
  - Classification DNN first optimised
  - DNN fixed, ANN optimised
  - Combined optimisation of both algorithms



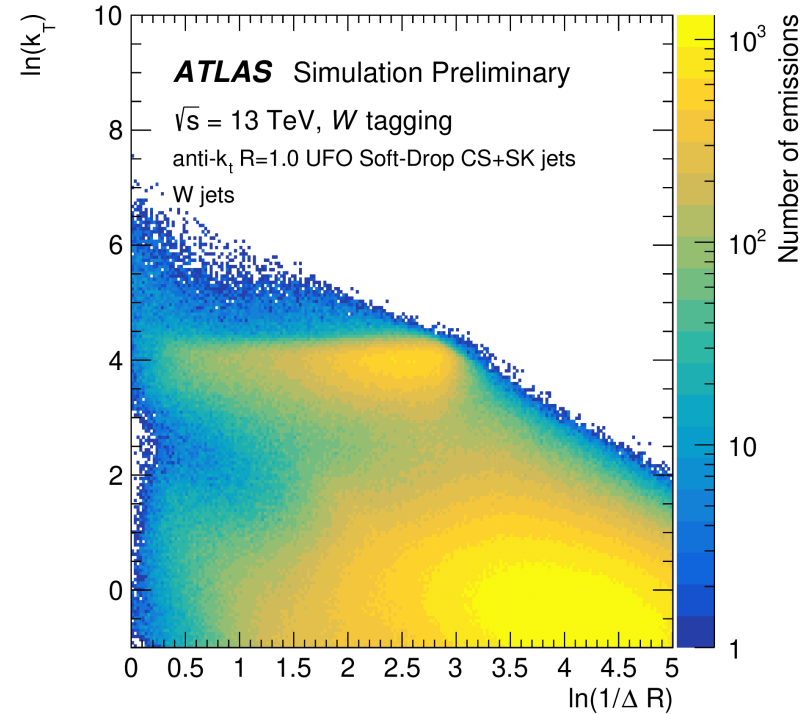
# Lund jet planes

$$\Delta R_{ij} = \sqrt{\Delta y_{ij}^2 + \Delta \phi_{ij}^2}, \quad z = \frac{p_T^j}{p_T^i + p_T^j}, \quad k_t = p_T^j \Delta R_{ij}$$

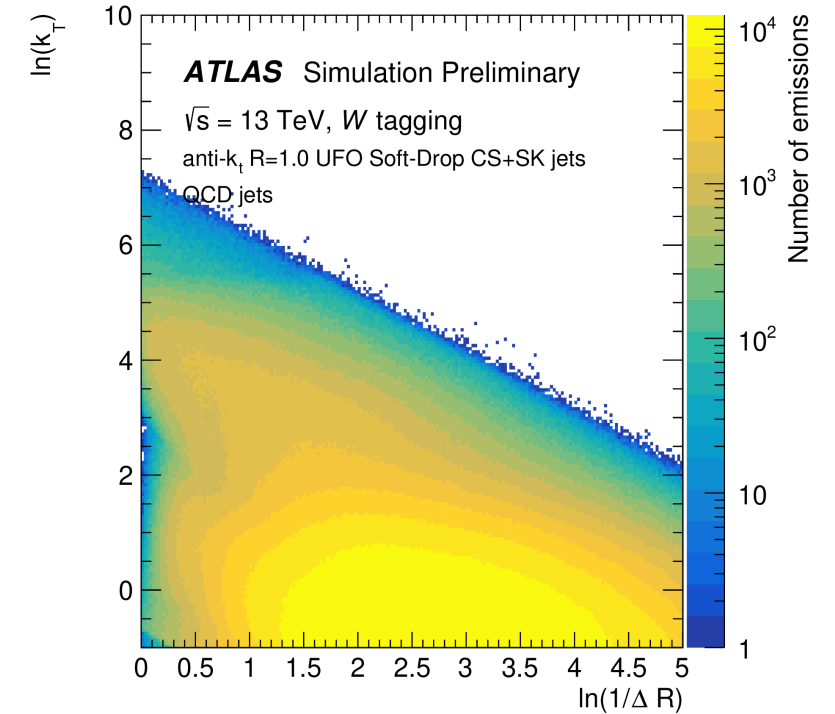
- Lund jet plane represents two-dimensional space built from the opening angle and the momentum fraction of a given gluon emission with respect to its emitter



[Jet substructure at high-pT](#)

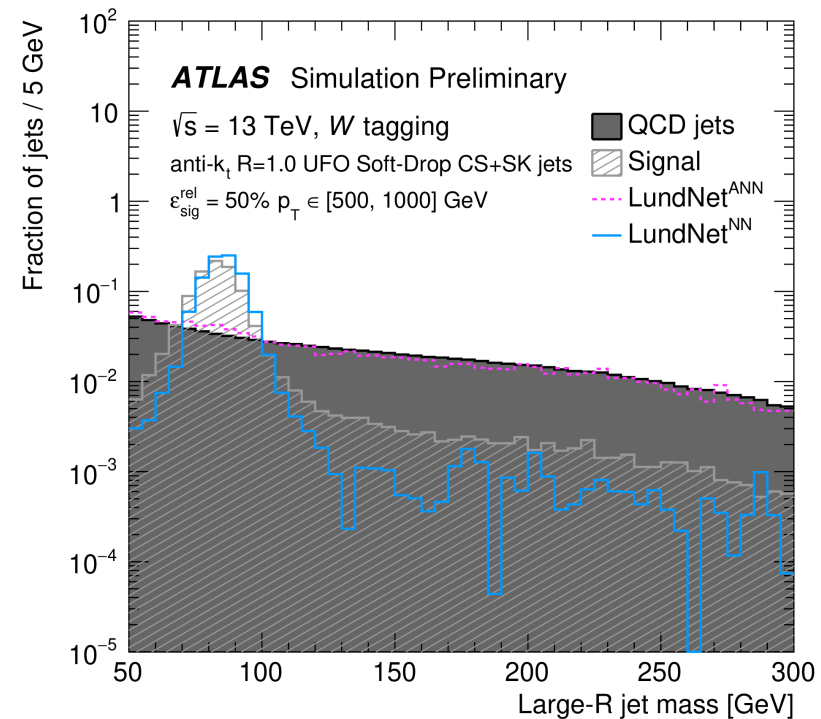
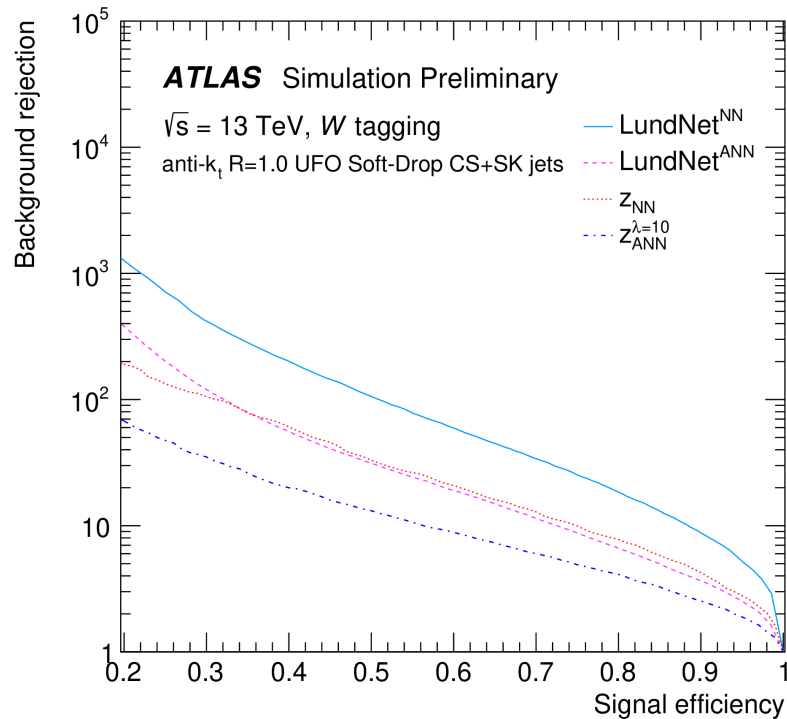


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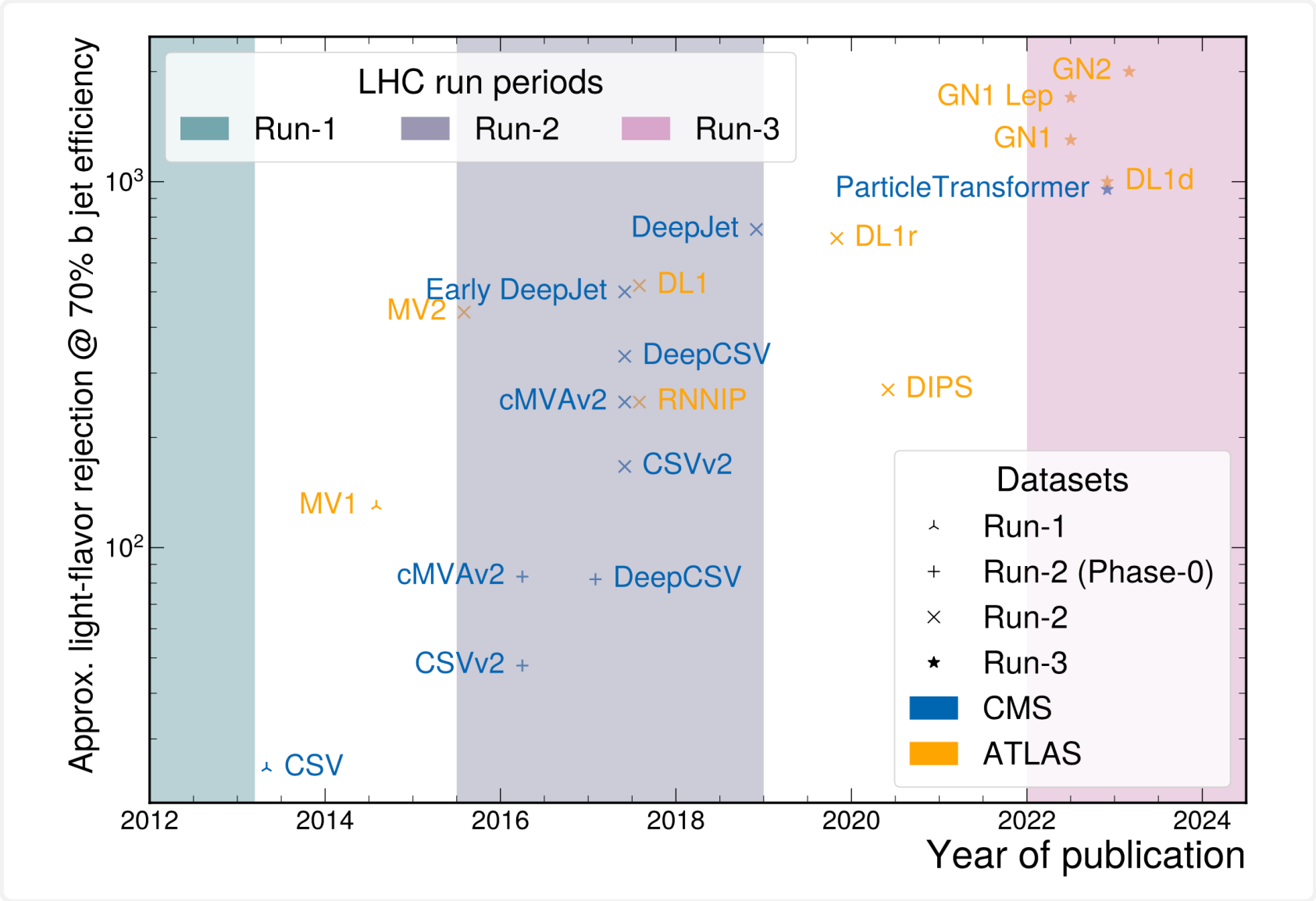


$$\Delta R_{ij} = \sqrt{\Delta y_{ij}^2 + \Delta \phi_{ij}^2}, \quad z = \frac{p_T^j}{p_T^i + p_T^j}, \quad k_i = p_T^i \Delta R_{ij}$$

- Sequence of jet emissions coming from the full declustering of large radius jet can be represented as a graph
- GNN inputs include  $\ln(1/\Delta R)$ ,  $\ln(k)$ ,  $\ln(1/z)$  and number of tracks
- Algorithm decorrelated from jet mass to prevent bias in background jet-mass distribution from being around the  $W$  boson mass
- Achieved via combined training of jet classification GNN and jet mass inferring adversarial NN (ANN)

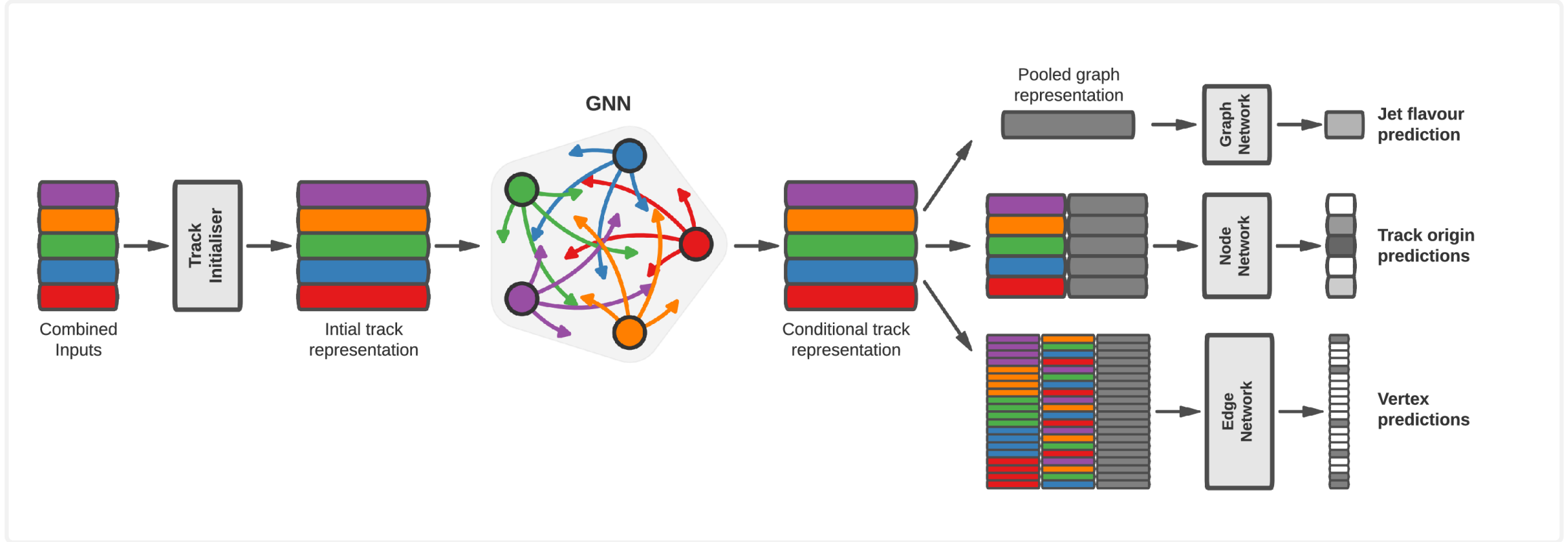


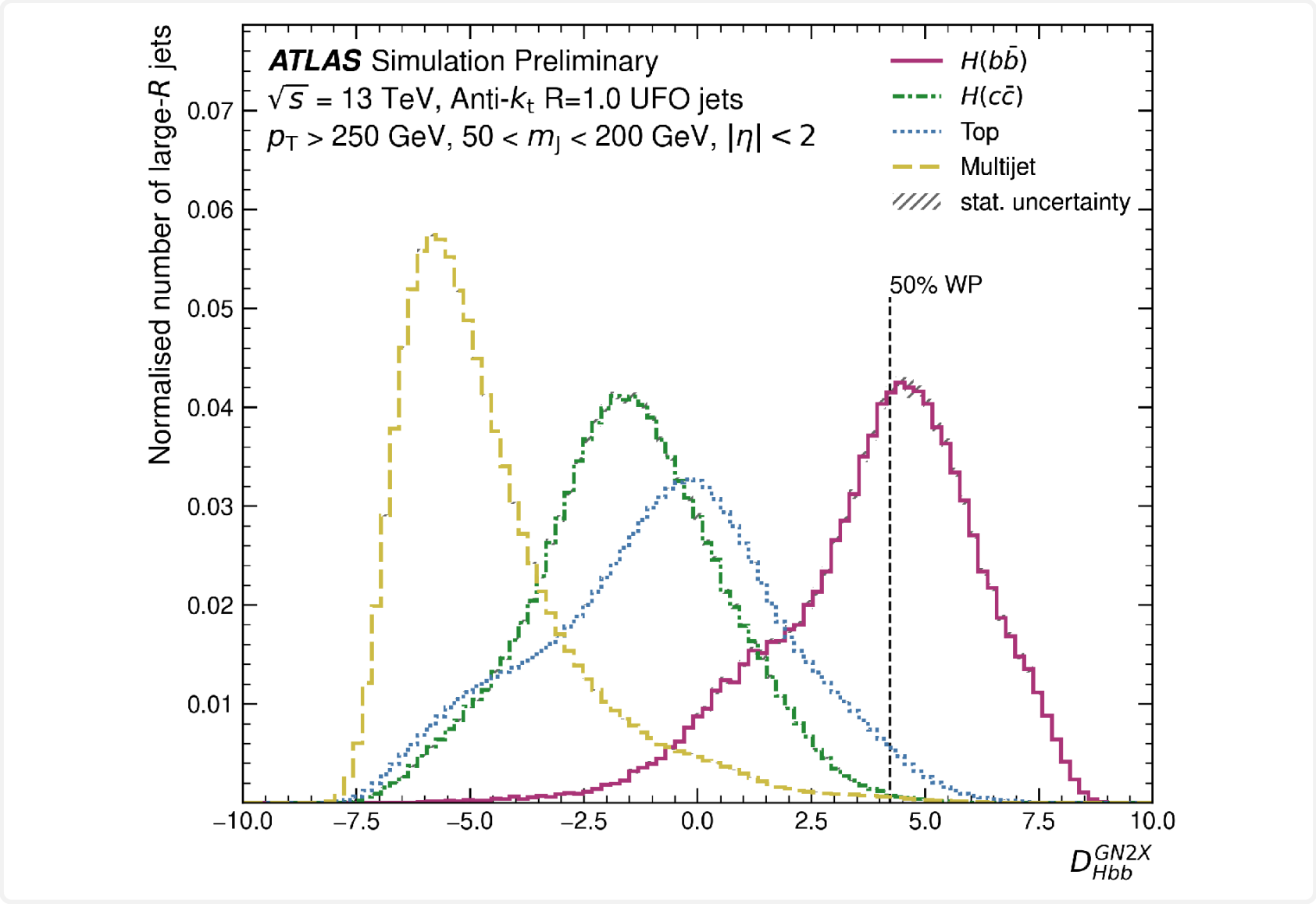
# Comparison of different jet taggers



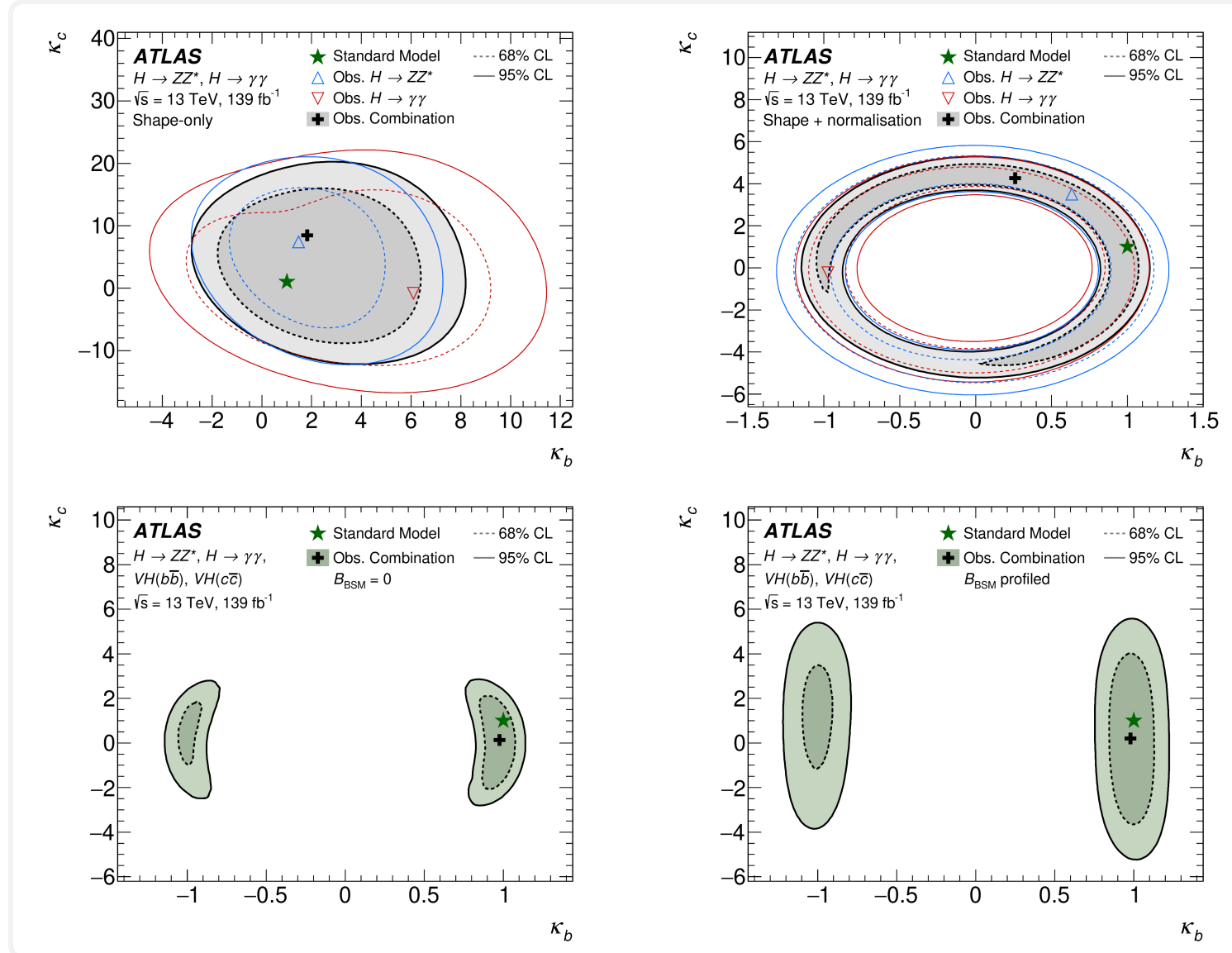


# GN1 Architecture

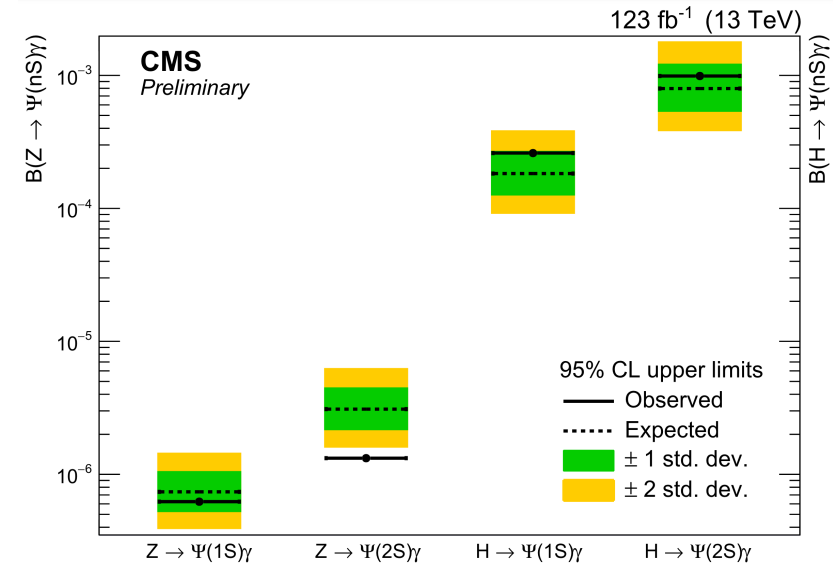
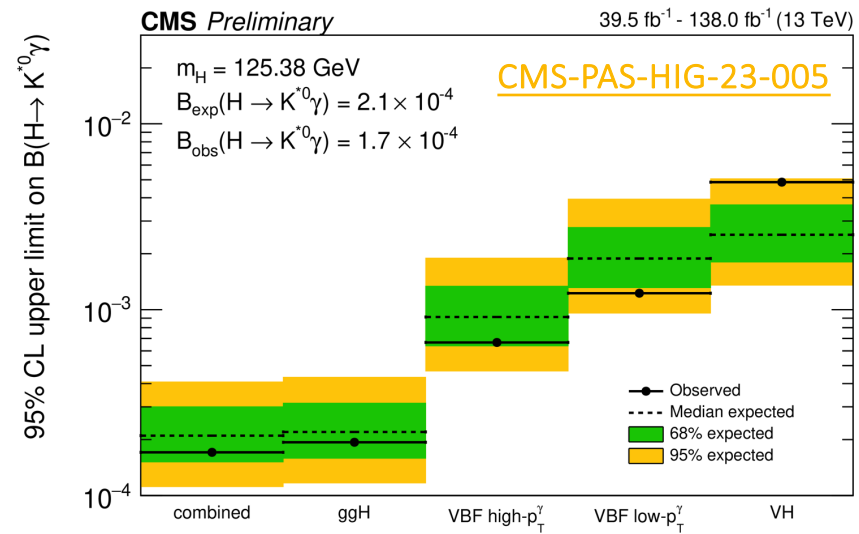
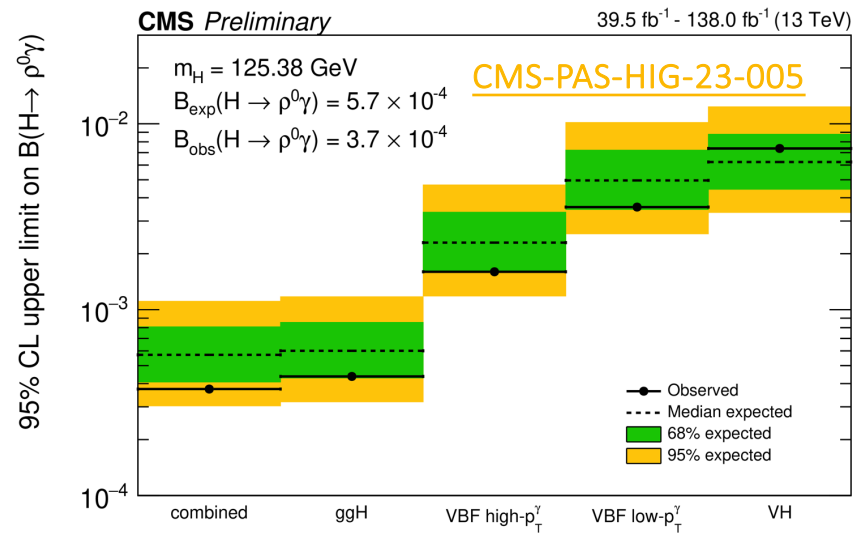




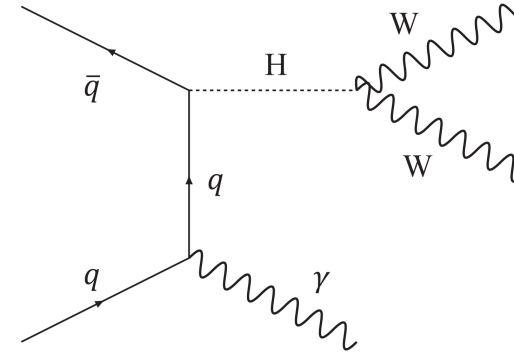
# $\kappa_C$ and $\kappa_b$ Interpretation



# CMS Higgs to meson + photon searches



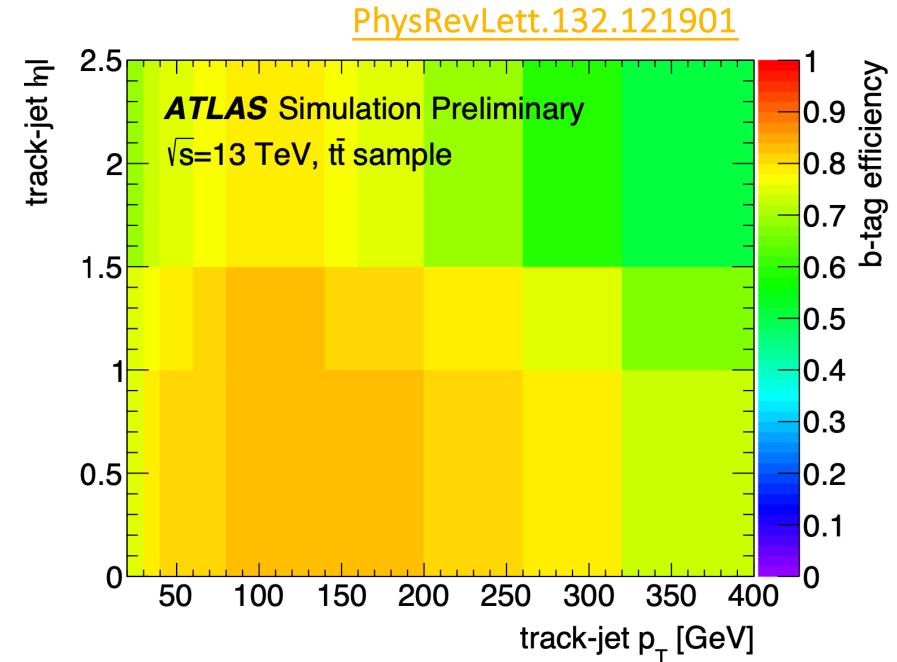
- Gluon initiated contribution disappears according to [Furry's theorem](#)
- Sensitive to Higgs quark Yukawas
- Cut-based analysis, **statistical and systematic uncertainties** with similar magnitude
- Other coupling modifiers set to their SM predictions



Process	$\sigma$ upper limits obs. (exp.) [fb]	$\kappa_q$ limits obs. (exp.) at 95% C.L.
$u\bar{u} \rightarrow H + \gamma \rightarrow e\mu\nu_e\nu_\mu\gamma$	86 (67)	$ \kappa_u  \leq 16000$ (13000)
$d\bar{d} \rightarrow H + \gamma \rightarrow e\mu\nu_e\nu_\mu\gamma$	72 (58)	$ \kappa_d  \leq 17000$ (14000)
$s\bar{s} \rightarrow H + \gamma \rightarrow e\mu\nu_e\nu_\mu\gamma$	66 (49)	$ \kappa_s  \leq 1700$ (1300)
$c\bar{c} \rightarrow H + \gamma \rightarrow e\mu\nu_e\nu_\mu\gamma$	88 (66)	$ \kappa_c  \leq 190$ (110)

# Truth flavour jet tagging

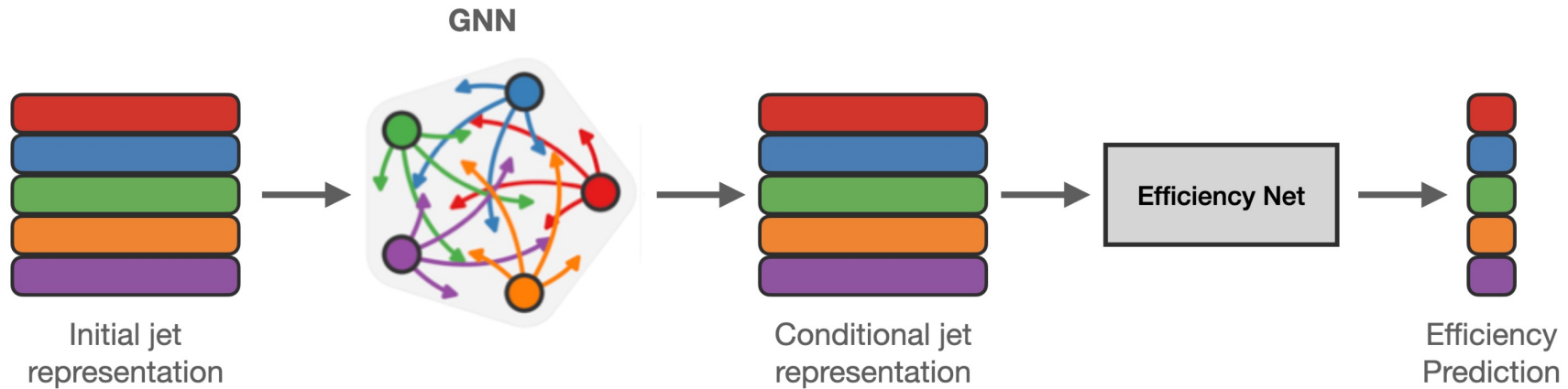
- Accepting/rejecting (“**direct tagging**”) events based on jet flavour tagging discriminants significantly **reduces available background** (e.g., V+light flavour jets processes) **statistics** due to **enhanced rejection** factors in latest algorithms
- ATLAS has implemented a “**truth-flavour tagging**” approach in recent  $V(\text{lep})H(b\bar{b})$  and  $V(\text{lep})H(c\bar{c})$  analyses
  - Consists in applying **tagging (in-)efficiency** of the Higgs candidate jets **as event weights**
  - **Efficiency maps** usually parametrised as **function of jet  $p_T$  and  $\eta$**
- Leads to **reduction of background statistical uncertainties**, since all available events are used
- Some **systematic uncertainties can arise** from differences to direct tagging approach (e.g. due to close-by jets) – **still resulting in better sensitivity**



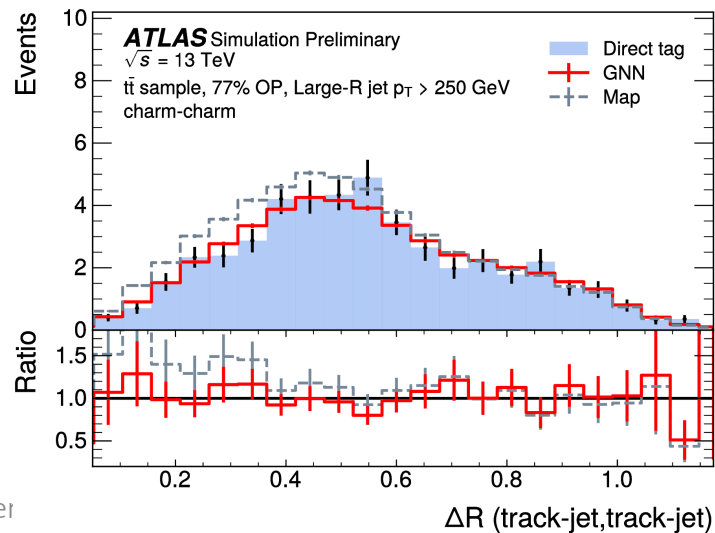
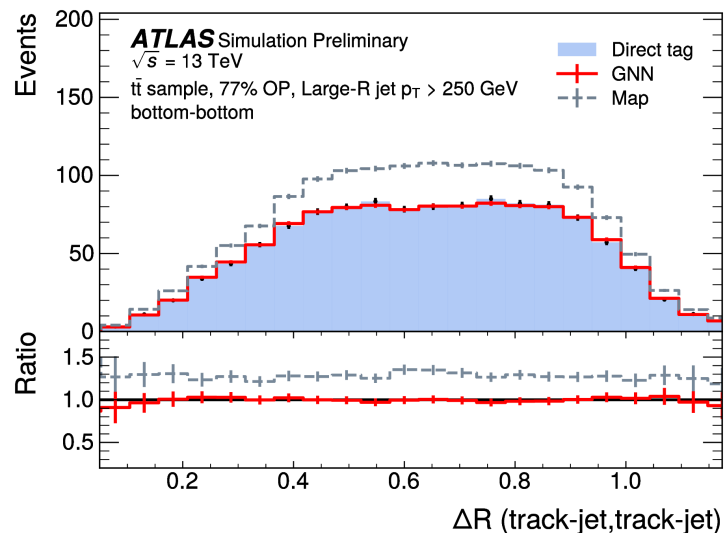
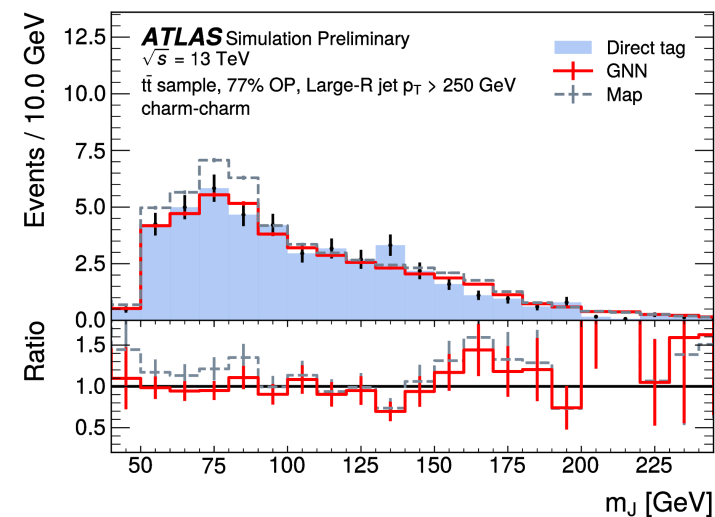
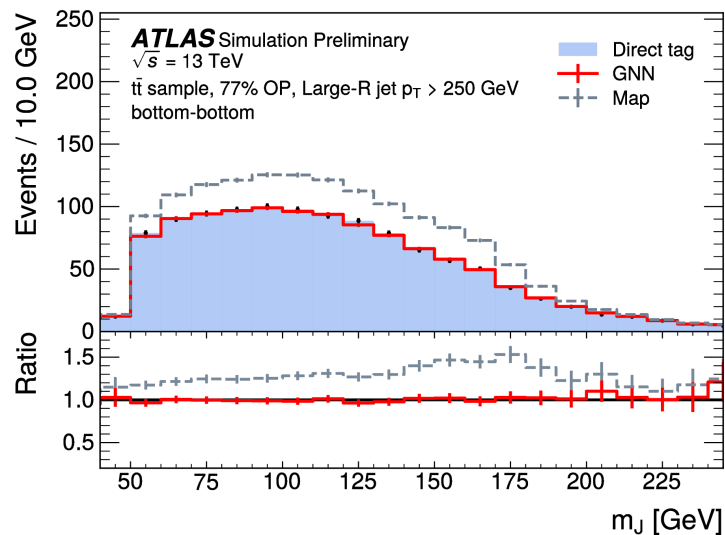
[Eur. Phys. J. C \(2022\) 82:717](#)

Source of uncertainty	$\mu_{VH(c\bar{c})}$
Total	21.5
Statistical	16.2
Systematics	14.0
Truth-flavour tagging	
$\Delta R$ correction	3.0
Residual non-closure	1.4

- Efficiency maps fail to capture all kinematic correlations, and can be affected by binning choices
- A GNN can potentially solve these problems, capturing high-dimensional correlations between the jets and leading to smooth output efficiency distributions
- GNN inputs include jet, event and jet-pair variables, trained with  $t\bar{t}$  events and large radius jets

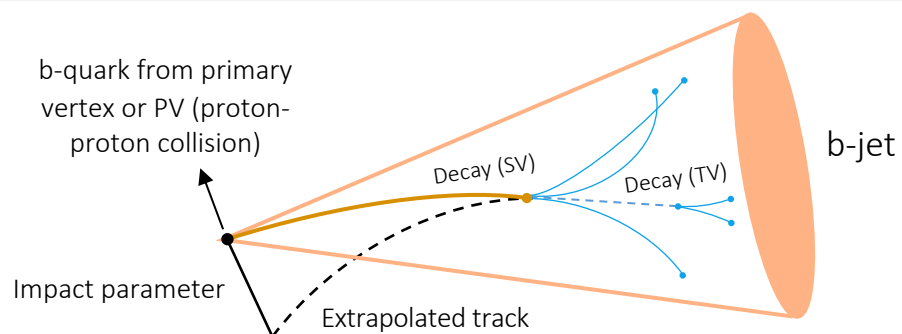


- GNN-based approach generally with similar or better agreement with direct tagging approach when compared to efficiency maps approach, particularly for very close-by jets

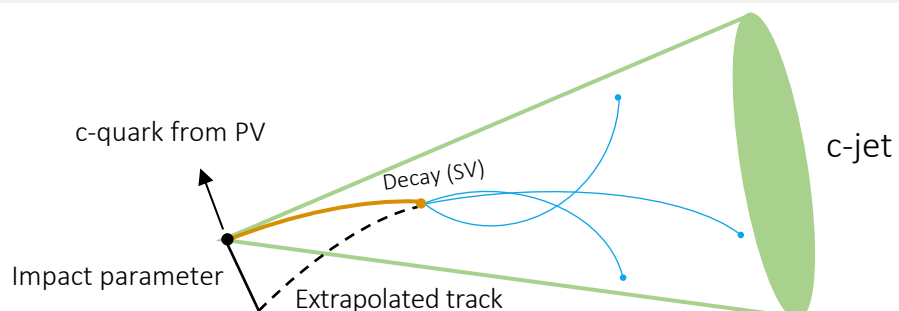




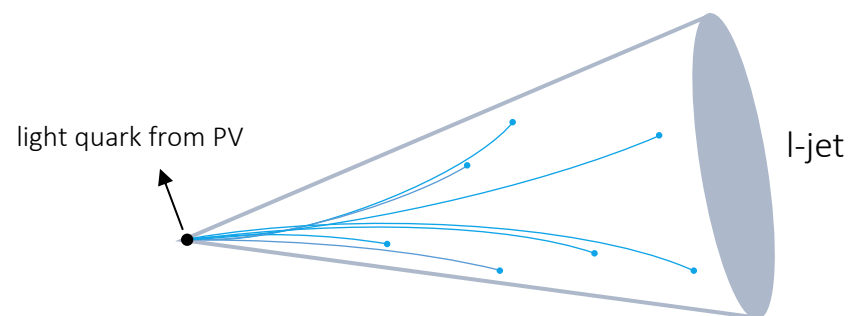
# Flavoured Jets



- b-quark fragments into **b-hadron** carrying around **80% of the jet energy**
- **High b-hadron decay product multiplicities** (around 5 charged particles per decay)
- Most b-hadrons ( $\approx 90\%$ ) **decay into c-hadrons**
- b(c)-hadron **decay vertex often displaced** from the PV(SV) by a few mm
- Tracks from both of these vertices often have **large impact parameters**



- c-quark fragments into **c-hadron** carrying around **55% of the jet energy**
- 2 to 3 times **lower c-hadron decay product multiplicities** than for b-hadrons (around 2 charged particles/decay)
- c-hadron **decay vertex often displaced** from the PV by a few mm
- Tracks from this vertex can often have **large impact parameters**



- Light quark hadronises into **many light hadrons sharing the jet energy**
- Tracks from this vertex can often have **impact parameters consistent with zero**
- Long-lived light hadrons can be produced, but more likely to decay cms away from PV