# Study of EFT effects in $n T G C$ and $a Q G C$ in ElectroWeak processes 

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## nTGC? aQGC? Why?

- SM provides gauge boson coupling of:



Triple Gauge Coupling (TGC): WWZ, WW $\gamma$

Quartic Gauge Coupling (QGC): WWWW, WWZZ, WWZ $\gamma$,
WW $\gamma \gamma$
-What if prohibited vertices exist?


A direct hint to

Neutral TGC (nTGC):
anomalous QGC (aQGC)


## nTGC? aQGC? How?

- Standard Model Effective Field Theory (SMEFT) based on Taylor expansion in local operators with mass dimension > 4

$$
\begin{aligned}
& \mathcal{L}_{S M E F T}=\mathcal{L}_{S M}+\sum_{i} \frac{c_{i}^{d=6}}{\Lambda_{i}^{2}} \mathcal{O}_{i}^{d=6}+\sum_{j} \frac{c_{j}^{d=8}}{\Lambda^{4}} \mathcal{O}_{j}^{d=8}+\cdots \\
& \text { Wilson Coefficient } \\
& \text { Onergy Scale } \quad \text { Operators }
\end{aligned}
$$

- The nTGCs and aQGCs (without aTGCs counterpart), which is today's focus, are described in diemension-8. While dimension-5 has one operators for neutrino mass and aTGCs arises from dimension-6.
- When the energy scale parameter $\Lambda \gg \sqrt{s}$ the expansion term can be truncated.
- Growth of amplitude with $\sqrt{s}$ can violate unitarity



## nTGC? aQGC? How? - The Basic Way

- Cross-section with single operator

$$
\sigma_{S M E F T}=\sigma_{S M}+\left(\frac{c^{d=8}}{\Lambda^{4}}\right) \sigma_{\text {int }}+\left(\frac{c^{d=8}}{\Lambda^{4}}\right)^{2} \sigma_{E F T}
$$

- Cross-section of interference term is proportional to coefficient.
- Cross-section of pure EFT contribution is proportional to the square of coefficient
- With cross-section of $\sigma_{i n t}$ and $\sigma_{E F T}$ from MC when $c=1$, a likelihood test can be performed to the measured crosssection:
$\mathcal{L}=\frac{1}{\sqrt{(2 \pi)^{k}|C o v|}} \exp \left(-\frac{1}{2}\left(\vec{\sigma}_{\text {data }}-\vec{\sigma}_{\text {SMEFT }}-\sum_{i} \theta \cdot \vec{e}_{\theta}\right)^{T} \operatorname{Cov}^{-1}\left(\vec{\sigma}_{\text {data }}-\vec{\sigma}_{\text {SMEFT }}-\sum_{i} \theta \cdot \vec{e}_{\theta}\right)\right) \times \prod_{i} \mathcal{N}\left(\theta_{i}\right)$


## $z^{2 / v}$ <br> $\sin _{z / v}$

nTGCs

## nTGCs - Parameters



Neutral TGC (nTGC):
$\mathrm{ZZZ}, \mathrm{ZZ} \gamma, \mathrm{Z}_{\gamma \gamma}$

$$
\begin{align*}
i e \Gamma_{Z Z V}^{\alpha \beta \mu}\left(\mathrm{q}_{1}, \mathrm{q}_{2}, \mathrm{q}_{3}\right) & =\frac{-e\left(\mathrm{q}_{3}^{2}-m_{V}^{2}\right)}{M_{Z}^{2}}\left[f_{4}^{V}\left(\mathrm{q}_{3}^{\alpha} g^{\mu \beta}+\mathrm{q}_{3}^{\beta} g^{\mu \alpha}\right)-f_{5}^{V} \epsilon^{\mu \alpha \beta \rho}\left(\mathrm{q}_{1}-\mathrm{q}_{2}\right)_{\rho}\right]  \tag{1.1}\\
i e \Gamma_{Z \gamma V}^{\alpha \beta \mu}\left(\mathrm{q}_{1}, \mathrm{q}_{2}, \mathrm{q}_{3}\right) & =\frac{-e\left(\mathrm{q}_{3}^{2}-m_{V}^{2}\right)}{M_{Z}^{2}}\left\{h_{1}^{V}\left(\mathrm{q}_{2}^{\mu} g^{\alpha \beta}-\mathrm{q}_{2}^{\alpha} g^{\mu \beta}\right)+\frac{h_{2}^{V}}{M_{Z}^{2}} \mathrm{q}_{3}^{\alpha}\left[\left(\mathrm{q}_{3} \mathrm{q}_{2}\right) g^{\mu \beta}-\mathrm{q}_{2}^{\mu} \mathrm{q}_{3}^{\beta}\right]\right. \\
& \left.-h_{3}^{V} \epsilon^{\mu \alpha \beta \rho} q_{2 \rho}-\frac{h_{4}^{V}}{M_{Z}^{2}} \mathrm{q}_{3}^{\alpha} \epsilon^{\mu \beta \rho \sigma} \mathrm{q}_{3 \rho} q_{2 \sigma}\right\} \tag{1.2}
\end{align*}
$$

- $f_{i}^{V}$ require on shell ZZ, while $h_{i}^{V}$ require on shell $Z \gamma$.
- A recent paper point out that extra operators and form factor should be introduced in nTGCs.-> PRD 107 (2023) 035005
Basis of dim 8 operators for nTGCs:

$$
\begin{aligned}
\mathcal{O}_{\widetilde{B} W} & =i H^{\dagger} \widetilde{B}_{\mu \nu} W^{\mu \rho}\left\{D_{\rho}, D^{\nu}\right\} H, \\
\mathcal{O}_{B \widetilde{W}} & =i H^{\dagger} B^{\mu \nu} \widetilde{W}_{\mu \rho}\left\{D_{\rho}, D^{\nu}\right\} H, \\
\mathcal{O}_{\widetilde{W} W} & =i H^{\dagger} \widetilde{W}_{\mu \nu} W^{\mu \rho}\left\{D_{\rho}, D^{\nu}\right\} H, \\
\mathcal{O}_{\widetilde{B} B} & =i H^{\dagger} \widetilde{B}_{\mu \nu} B^{\mu \rho}\left\{D_{\rho}, D^{\nu}\right\} H .
\end{aligned}
$$



## CMS - ZZ $(\rightarrow 4 l)$

- Sensitive to two nTGCs: ZZZ, ZZ $\gamma$.
- Test predictions at next-to-next-to-leading order (NNLO) in QCD.
- Low background contribution (~3\%) due to the requirement for four well-reconstructed and isolated leptons.



- Cross section measurement:

| Year | Total cross section, pb |
| :--- | :--- |
| 2016 | $18.1 \pm 0.6$ (stat) $)_{-0.5}^{+0.6}$ (syst) $\pm 0.4$ (theo) ${ }_{-0.4}^{+0.5}$ (lumi) |
| 2017 | $17.0 \pm 0.5$ (stat) ${ }_{-0.5}^{+0.6}$ (syst) $\pm 0.4$ (theo) $\pm 0.4$ (lumi) |
| 2018 | $17.1 \pm 0.4$ (stat) $\pm 0.5$ (syst) $\pm 0.4$ (theo) $\pm 0.4$ (lumi) |
| Combined | $17.4 \pm 0.3$ (stat) $\pm 0.5$ (syst) $\pm 0.4$ (theo) $\pm 0.3$ (lumi) |

- Consistent with the NNLO prediciton


## CMS - ZZ $(\rightarrow 4 l)$

- 2-D constraints, set limit to two parameters simultaneously. Predicted cross section:


## $\sigma_{\text {SMEFT }}$

$$
=\sigma_{S M}+c_{1} \sigma_{i n t 1, S M}+c_{2} \sigma_{i n t 2, S M}+c_{1} c_{2} \sigma_{i n t 1,2}+c_{1}^{2} \sigma_{E F T, 1}+c_{2}^{2} \sigma_{E F T, 2}
$$

- Constraints are set on $m_{Z Z}$, CP-even variable. Hence CP-odd parameters $\left(f_{4}^{V}\right)$ interference term are vanished.
- Overflow contribution are included in the last bin.



## ATLAS - ZZ $(\rightarrow 4 l)$ Angular

- Search for CP-violation and nTGCs in ZZ(4l) on-shell events (dim-8 EFT)
- Measure the ZZ polarization in 4 l channel (extract the LL component)
- Measure the spin correlation between ZZ bosons




## ATLAS - ZZ $(\rightarrow 4 l)$ Angular

- To improve sensitivity, the two CP sensitive angles are combined as:

$$
T_{y z, 1(3)}=\sin \phi_{1(3)} \cos \theta_{1(3)}
$$

- An Optimal Observable (OO) is defined from the 2D distribution of $T_{y z, 1} V . S . T_{y z, 3}$ to maximise the sensitivity for the fourlepton system.

| aNTGC parameter | Interference only |  | Full |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | Expected | Observed | Expected | Observed |  |
| $f_{Z}^{4}$ | $[-0.16,0.16]$ | $[-0.12,0.20]$ | $[-0.013,0.012]$ | $[-0.012,0.012]$ |  |
| $f_{\gamma}^{4}$ | $[-0.30,0.30]$ | $[-0.34,0.28]$ | $[-0.015,0.015]$ | $[-0.015,0.015]$ |  |

- A BDT is used to determine the three different ZZ polarisation pairs: $Z_{L} Z_{L}$ (Signal) \| $Z_{T} Z_{L} Z_{T} Z_{T}$ (Background)
- Fiducial cross section (4.3 $\sigma$ for $Z_{L} Z_{L}$ ):
- $\sigma_{Z_{L} Z_{L}}^{o b s .}=2.45 \pm 0.56$ (stat.) $\pm 0.21$ (syst.) fb
- $\sigma_{Z_{L} Z_{L}}^{p r e d .}=2.10 \pm 0.09 \mathrm{fb}$

CMS -Z $(\rightarrow v \bar{v}) \gamma$

- Invisible Z decay has a higher branching fraction (20\%) compared to the leptonic $Z$ channel (10\%) and cleaner signature compared to both leptonic $Z$ decay and hadronic decay.
- Measurement divided into barrel and endcaps due to different detector response on fake backgrounds:



ZZ $\gamma$ Vertex

| Parameter | Expected | Observed |
| :--- | :--- | :--- |
| $h_{3}^{\gamma} \times 10^{4}$ | $(-2.8,2.9)$ | $(-3.4,3.5)$ |
| $h_{4}^{\gamma} \times 10^{7}$ | $(-5.9,6.0)$ | $(-6.8,6.8)$ |
| $h_{3}^{Z} \times 10^{4}$ | $(-1.8,1.9)$ | $(-2.2,2.2)$ |
| $h_{4}^{Z} \times 10^{7}$ | $(-3.7,3.7)$ | $(-4.1,4.2)$ |

The sensitivities to CP-conserving and CP- violating couplings are comparable in the probed $p_{T}$ regime.


## aQGCs - Parameters

- The Eboli Model:
- tensor (T): EWK field strength tensors derivatives

anomalous QGC (aQGC)
- scalar (S): Higgs doublet derivatives
- mixed (M): both

|  | WWWW | WWZZ | ZZZZ | WWAZ | WWAA | ZZZA | ZZAA | ZAAA | AAAA |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathcal{L}_{S, 0}, \mathcal{L}_{S, 1}$ | X | X | X | O | O | O | O | 0 | O |
| $\mathcal{L}_{M, 0}, \mathcal{L}_{M, 1}, \mathcal{L}_{M, 6}, \mathcal{L}_{M, 7}$ | X | X | X | X | X | X | X | 0 | 0 |
| $\mathcal{L}_{M, 2}, \mathcal{L}_{M, 3}, \mathcal{L}_{M, 4}, \mathcal{L}_{M, 5}$ | O | X | X | X | X | X | X | 0 | 0 |
| $\mathcal{L}_{T, 0}, \mathcal{L}_{T, 1}, \mathcal{L}_{T, 2}$ | X | X | X | X | X | X | X | X | X |
| $\mathcal{L}_{T, 5}, \mathcal{L}_{T, 6}, \mathcal{L}_{T, 7}$ | O | X | X | X | X | X | X | X | X |
| $\mathcal{L}_{T, 9}, \mathcal{L}_{T, 9}$ | O | O | X | O | O | X | X | X | X |
| he analyses, it is essential to rtial-wave unitarity is satisfied |  |  |  |  |  |  |  | und |  |
|  |  |  |  |  |  |  |  | erator |  |
|  |  |  | Wilson coefficient |  |  |  | For $\sqrt{s}<1.5(3) \mathrm{TeV}$ |  |  |
|  |  |  | $\left\|\frac{f_{S_{0}}}{\Lambda^{4}}\right\|$ |  |  | $32 \pi s^{-2}$ |  | $20(1.2) \mathrm{TeV}^{-4}$ |  |
|  |  |  | $\left\|\frac{f_{s, 1}}{\Lambda^{4}}\right\|$ |  |  | $\frac{96}{7} \pi s^{-2}$ |  | $8.5(0.53) \mathrm{TeV}^{-4}$ |  |
|  |  |  | $\left\|\frac{f_{S_{2}}}{\Lambda^{4}}\right\|$ |  |  | $\frac{96}{5} \pi \mathrm{~s}^{-2}$ |  | $8.5(0.53) \mathrm{TeV}^{-4}$ |  |

: constraints on each Wilson coefficient can be obtained after restricting EFT contribution within $\sqrt{S}<E_{C}$, and the unitarity bound of $E_{c}$ can be calculated and compared with the constraints.

## ATLAS - VBS ZZ $(\rightarrow 4 l)$

- Motivation and goal
- Sensitive to 3 and 4 -weak boson self-interactions
- Differential cross-sections can probe New Physics (aTGC, aQGC)
- Unfolded differential cross section measurement.
- Remove detector response


Physics distribution $y_{i}$

(Particle-level)


## ATLAS - VBS ZZ $(\rightarrow 4 l)$

- Unfolded cross-sections in agreement with predictions (some underestimation from MG5+PY8 strong production)
- Limits to dim-8 operators from a combined $m_{j j}+m_{4 \ell}$ fit with overflow contributions.

PRD 101, 113003 (2020)

- Clip scan is performed via $E_{c}=m_{4 l}$ to check the unitarity bound (if violated).

Clip scan is performed by estimating the limit of Wilson coefficient when clipping all the EFT event with the energy higher than the given $E_{c}$.



| Wilson <br> coefficient | $\left\|\mathcal{M}_{\mathrm{d} 8}\right\|^{2}$ <br> Included | 95\% confidence <br> Expected | Observed |
| :---: | :---: | :---: | :---: |
| $f_{\mathrm{T}, 0} / \Lambda^{4}$ | yes | $[-1.00,0.97]$ | $[-0.98,0.93]$ |
|  | no | $[-19,19]$ | $[-23,17]$ |
| $f_{\mathrm{T}, 1} / \Lambda^{4}$ | yes | $[-1.3,1.3]$ | $[-1.2,1.2]$ |
|  | no | $[-140,140]$ | $[-160,120]$ |
| $f_{\mathrm{T}, 2} / \Lambda^{4}$ | yes | $[-2.6,2.5]$ | $[-2.5,2.4]$ |
|  | no | $[-63,62]$ | $[-74,56]$ |
| $f_{\mathrm{T}, 5} / \Lambda^{4}$ | yes | $[-2.6,2.5]$ | $[-2.5,2.4]$ |
|  | no | $[-68,67]$ | $[-79,60]$ |
| $f_{\mathrm{T}, 6} / \Lambda^{4}$ | yes | $[-4.1,4.1]$ | $[-3.9,3.9]$ |
|  | no | $[-550,540]$ | $[-640,480]$ |
| $f_{\mathrm{T}, 7} / \Lambda^{4}$ | yes | $[-8.8,8.4]$ | $[-8.5,8.1]$ |
|  | no | $[-220,220]$ | $[-260,200]$ |
| $f_{\mathrm{T}, 8} / \Lambda^{4}$ | yes | $[-2.2,2.2]$ | $[-2.1,2.1]$ |
|  | no | $[-3.9,3.8] \times 10^{4}$ | $[-4.6,3.1] \times 10^{4}$ |
| $f_{\mathrm{T}, 9} / \Lambda^{4}$ | yes | $[-4.7,4.7]$ | $[-4.5,4.5]$ |
|  | no | $[-6.4,6.3] \times 10^{4}$ | $[-7.5,5.5] \times 10^{4}$ |

## ATLAS - VBS $\mathrm{Z}(\rightarrow \nu \bar{v}) \gamma$

- VBS $Z(\rightarrow v \bar{v}) \gamma$ is observed in low energy phase space ( $15<$ $E_{T}^{\gamma}<110 \mathrm{GeV}$ ) by ATLAS (EPJC 82 (2022) 105). But low energy phase space has no sensitivity to aQGCs.
- This analysis conduct the VBS $Z(\rightarrow v \bar{v}) \gamma$ in high energy phasespace ( $E_{T}^{\gamma}>150 \mathrm{GeV}$ ). Both phase-space can be combined to obtain higher sensitivity to aQGCs.
- Dominant background from QCD $Z(\rightarrow v \bar{v}) \gamma j j$ and $\mathrm{W}(\rightarrow l v) \gamma j j$.
- Combined measurement has found a $6.3 \sigma$ ( $6.6 \sigma$ ) significance on signal strengthen of VBS $Z(\rightarrow \nu \bar{v}) \gamma$ and the fiducial cross section of high energy phase space is measured:

$$
\left.\sigma_{Z \gamma \mathrm{EWK}}=0.77_{-0.30}^{+0.34} \mathrm{fb}=0.77_{-0.23}^{+0.25} \text { (stat.) }\right)_{-0.18}^{+0.22} \text { (syst.) fb. }
$$




## ATLAS - VBS $\mathbb{L}(\rightarrow 1$ $\cdot$ Probed for nQGCs via $E_{T}^{\gamma}$

- Clip scan performed by setting clip energy $E_{c}=$ $m_{Z_{\gamma}}$ (using particle-level information).
- The regime in which $E_{c}$ is less than 4 TeV is obtained with an $E_{T}^{\gamma}$ threshold of $600 \mathrm{GeV}(400$ GeV ) for $f_{T}\left(f_{M}\right)$.
- The regime in which $E_{c}$ exceeds 4 TeV is obtained with an $E_{T}^{\gamma}$ threshold of 900 GeV .

| Coefficient | $E_{\mathrm{c}}[\mathrm{TeV}]$ | Observed limit $\left[\mathrm{TeV}^{-4}\right]$ | Expected limit $\left[\mathrm{TeV}^{-4}\right]$ |
| :---: | :---: | :---: | :---: |
| $f_{T 0} / \Lambda^{4}$ | 1.7 | $[-8.7,7.1] \times 10^{-1}$ | $[-8.9,7.3] \times 10^{-1}$ |
| $f_{T 5} / \Lambda^{4}$ | 2.4 | $[-3.4,4.2] \times 10^{-1}$ | $[-3.5,4.3] \times 10^{-1}$ |
| $f_{T 8} / \Lambda^{4}$ | 1.7 | $[-5.2,5.2] \times 10^{-1}$ | $[-5.3,5.3] \times 10^{-1}$ |
| $f_{T 9} / \Lambda^{4}$ | 1.9 | $[-7.9,7.9] \times 10^{-1}$ | $[-8.1,8.1] \times 10^{-1}$ |
| $f_{M 0} / \Lambda^{4}$ | 0.7 | $[-1.6,1.6] \times 10^{2}$ | $[-1.5,1.5] \times 10^{2}$ |
| $f_{M 1} / \Lambda^{4}$ | 1.0 | $[-1.6,1.5] \times 10^{2}$ | $[-1.4,1.4] \times 10^{2}$ |
| $f_{M 2} / \Lambda^{4}$ | 1.0 | $[-3.3,3.2] \times 10^{1}$ | $[-3.0,3.0] \times 10^{1}$ |




## ATLAS - VBS WZ

- Boost Decision Tree (BDT) for separating QCD WZij and VBS WZ. 15 input variables are used, including jetkinematics variables, vector-bosons-kinematics variables, and variables related to both jets and leptons kinematics.
- Four bins in BDT score ( $[-1,-0.25,0.17,0.72,1]$ ) and five bins in $m_{T}^{W Z}([0,400,750,1050,1350, \infty] \mathrm{GeV})$ are used and arranged in a one-dimensional histogram of 20 statistically independent bins for EFT re-interpretation.





## ATLAS - VBS W $\gamma$

- Setting $f_{T}$ constraints via unfolded $p_{T}^{j j}$ distribution, $f_{M}$ constraints via unfolded $p_{T}^{l}$ distribution.
- Clip scan cut-off performed via $M_{W_{\gamma}}$
- A first measurement on $f_{T 3}$ and $f_{T 4}$ in LHC.





## CMS - VBS ssWW with hadronic $\tau$

$\mathcal{L}=138 \mathrm{fb}^{-1} @ 13 \mathrm{TeV}$

- VBS same-sign (ss) WW with one W decays to $e$ or $\mu$, another $W$ decays to hadronic $\tau$. Signal: $\tau v_{\tau} l v_{l} j j(l=e, \mu)$
- Significance of SM process at $2.7 \sigma$, signal strength: $1.44_{-0.56}^{+0.63}$
- First simultaneous extraction of dim-6 and dim-8 constraints


- 2-D constraints set via transverse mass $M_{o 1}$ :

$$
M_{o 1}^{2}=\left(p_{T}^{\tau}+p_{T}^{l}+p_{T}^{\text {miss }}\right)^{2}-\left|\vec{p}_{T}^{\tau}+\vec{p}_{T}^{l}+\vec{p}_{T}^{\text {miss }}\right|^{2}
$$

- Cross section for dim-6 + dim-8 operator:

$$
\sigma_{S M E F T}=\sigma_{S M}+c_{d-6} \sigma_{i n t}+c_{d-6}^{2} \sigma_{d-6}+c_{d-8} \sigma_{i n t}+c_{d-8}^{2} \sigma_{d-8}
$$

## CMS - VBS ssWW with hadronic $\tau$

$\mathcal{L}=138{f b^{-1} @ 13 \mathrm{TeV}}^{\text {@ }}$

- Also 1-D constraints are set via Deep Neural Network (DNNs) score

| Wilson coefficient |  | 68\% CL interval(s) |  | 95\% CL interval |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | Expected | Observed | Expected | Observed |
| dim-6 | $c_{l l}^{(1)}$ | $[-12.9,-8.03] \cup[-2.95,1.91]$ | [-11.6, 0.045] | [-14.6,3.53] | [-13.5,2.11] |
|  | $c_{q 9}^{(1)}$ | [-0.501, 0.576] | [-0.341, 0.416] | [-0.742, 0.818 ] | [-0.605, 0.681] |
|  | $c_{W}$ | [-0.681, 0.669] | [ $-0.513,0.481$ ] | [ $-0.987,0.974$ ] | -0.842,0.818] |
|  | $c_{\text {HW }}$ | [-7.00, 6.09] | [ $-5.48,4.31$ ] | [-9.99, 9.05] | [-8.68,7.60] |
|  | $c_{\text {HWB }}$ | [-41.7, 69.6] | [30.7, 89.2] | [-66.6, 96.4] | [-49.7,110] |
|  | $\mathcal{c}_{\text {H }}$ | [-16.6,18.1] | [-12.0, 14.0] | [-24.7,26.3] | [-20.9,22.7] |
|  | $c^{\text {HD }}$ | [-24.6,34.7] | [-15.3,31.5] | [-38.2, 48.8] | [-31.4,45.5] |
|  | $c_{\text {Hl }}^{(1)}$ | [-28.8, 29.9] | [-38.2,39.5] | [-49.4,49.7] | [-69.3,68.3] |
|  | $c_{\text {H1 }}$ | $[-1.43,2.23] \cup[5.88,9.54]$ | [-0.045, 8.58] | [-2.64,10.8] | [-1.59, 9.94] |
|  | $c^{(1)}$ | [-4.53, 4.42] | [-3.27, 3.44$]$ | [-6.56, 6.44] | [-5.55, 5.60] |
|  | $c_{\text {Hq }}$ | [-2.39, 1.37] | [-1.88, 0.705] | [-3.24, 2.16] | [-2.82, 1.61] |
| dim-8 | $f_{T 0}$ | [ $-1.02,1.08$ ] | [-0.774, 0.842] | [ $-1.52,1.58$ ] | [-1.32, 1.38] |
|  | $f_{T 1}$ | [ $-0.426,0.480]$ | [ $-0.319,0.381$ ] | [-0.640, 0.695] | [-0.552, 0.613] |
|  | $f_{T 2}$ | [-1.15, 1.37] | [-0.851, 1.12] | [ $-1.75,1.98$ ] | [-1.51,1.76] |
|  | $f_{\text {M0 }}$ | [-9.89,9.74] | [-8.07,7.70] | [-14.6, 14.5] | [-13.1, 12.8] |
|  | $f_{M 1}$ | [-12.5, 13.3] | [-9.54, 11.15] | [-18.7,19.6] | [-16.4,17.7] |
|  | $f_{M 7}$ | [-20.3,19.2] | [-17.6, 15.3] | [-29.9, 28.8] | [-27.6, 25.8] |
|  | $f_{S 0}$ | [-11.6, 12.0] | [-9.60,9.82] | [-17.4, 17.9] | [-15.9,16.1] |
|  | $f_{S 1}$ | [-37.4,38.8] | [-40.9, 41.3] | [-57.2,58.6] | [-60.9,61.8] |
|  | $f_{S 2}$ | [-37.4,38.8] | [-40.9, 41.3] | [-57.2,58.6] | [-60.9,61.8] |

Di-Boson Interaction

$$
\mathrm{qq} \rightarrow \mathrm{ZZ} \rightarrow 4 \mathrm{I}
$$

$$
\mathrm{gg} \rightarrow \mathrm{ZZ} \rightarrow 4 \mid
$$

$$
\mathrm{gg} \rightarrow \mathrm{H} \rightarrow \mathrm{ZZ} \rightarrow 4 \mid
$$



What Else?

## CMS - TriboSon VYY $\quad \mathcal{L}=137 \mathrm{fb}^{-1} @ 13 \mathrm{TeV}$

- e/ $\mu$ channels are used and combined in this measurement
- Background dominated by misid-jet and misid electrons
- Sensitive to the dim-6 and dim-8 operators, but lower statistics than di-bosons results in much weaker limits
- EFT constraints set via $p_{T, \gamma \gamma}$

|  | $\mathrm{W} \gamma \gamma\left(\mathrm{TeV}^{-4}\right)$ |  | $\mathrm{Z} \gamma \gamma\left(\mathrm{TeV}^{-4}\right)$ |  |
| :---: | :---: | :---: | :---: | :---: |
| Parameter | Expected | Observed | Expected | Observed |
| $f_{\mathrm{M} 2} / \Lambda^{4}$ | $[-57.3,57.1]$ | $[-39.9,39.5]$ | - | - |
| $f_{\mathrm{M} 3} / \Lambda^{4}$ | $[-91.8,92.6]$ | $[-63.8,65.0]$ | - | - |
| $f_{\mathrm{T} 0} / \Lambda^{4}$ | $[-1.86,1.86]$ | $[-1.30,1.30]$ | $[-4.86,4.66]$ | $[-5.70,5.46]$ |
| $f_{\mathrm{T} 1} / \Lambda^{4}$ | $[-2.38,2.38]$ | $[-1.70,1.66]$ | $[-4.86,4.66]$ | $[-5.70,5.46]$ |
| $f_{\mathrm{T} 2} / \Lambda^{4}$ | $[-5.16,5.16]$ | $[-3.64,3.64]$ | $[-9.72,9.32]$ | $[-11.4,10.9]$ |
| $f_{\mathrm{T} 5} / \Lambda^{4}$ | $[-0.76,0.84]$ | $[-0.52,0.60]$ | $[-2.44,2.52]$ | $[-2.92,2.92]$ |
| $f_{\mathrm{T} 6} / \Lambda^{4}$ | $[-0.92,1.00]$ | $[-0.60,0.68]$ | $[-3.24,3.24]$ | $[-3.80,3.88]$ |
| $f_{\mathrm{T} 7} / \Lambda^{4}$ | $[-1.64,1.72]$ | $[-1.16,1.16]$ | $[-6.68,6.60]$ | $[-7.88,7.72]$ |
| $f_{\mathrm{T} 8} / \Lambda^{4}$ | - | - | $[-0.90,0.94]$ | $[-1.06,1.10]$ |
| $f_{\mathrm{T} 9} / \Lambda^{4}$ | - | - | $[-1.54,1.54]$ | $[-1.82,1.82]$ |



## Summary \& Outlook

- As nTGCs and aQGCs are direct hints to BSM physics, the re-interpretation is become one main part of bosonic electroweak analysis.
- nTGCs limits are set by diboson $Z Z$ or $Z \gamma$ production. aQGCs limits are obtain by VBS and Tri-boson production. All results are compatible with SM so far.
- Unfolded analysis allows test of new models in the future.


## Thank you!

- Challenges:
- Current analysis set constraints on one parameter / two parameters in simultaneously. How about more parameters and even full model?
- Unitarity violation when including higher energy overflow contribution
- Higher order correction of BSM model is absent, current analysis uses EFT model generated in tree-level
- LHC Run3 is on-going. Higher statistic and higher enegy $\rightarrow$ Higher sensitivity to BSM physics! Moreover, new global fit of Run 2 is await to be conducted.


## Backup

## nTGCs summary in 2020



## nTGCs summary in 2018



## aQGCs summary in 2020



## aQGCs summary in 2020



## aQGCs summary in 2020



