



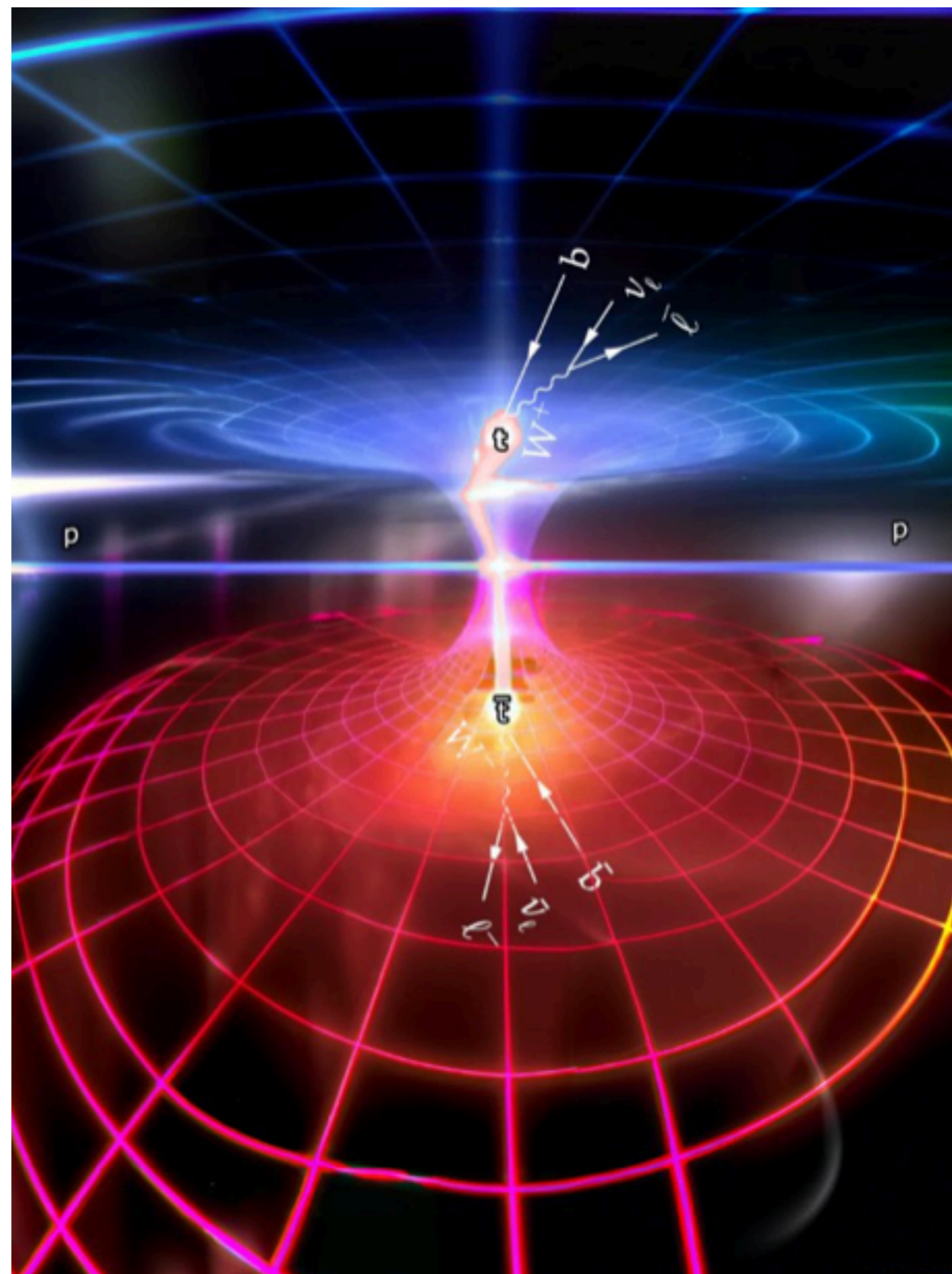
# Probing entanglement in top quark production with the CMS detector

**Giulia Negro**

on behalf of the CMS Collaboration

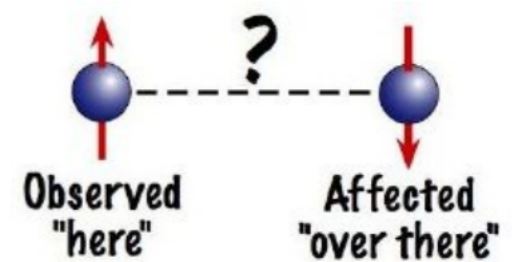
**Standard Model at the LHC 2024**

*10 May 2024*



# Entanglement at the LHC

- **Fundamental predictions of Quantum Mechanics:**
  - entangled states cannot be described by independent superpositions
  - measuring particle spin in an entangled system immediately reveals the spin state of the second particle



- **A lot of measurements with electrons and photons already performed**

Nobel Prize in 2022 for Aspect, Clauser, and Zeilinger

- **First observation of entanglement in  $t\bar{t}$  by ATLAS at the end of last year**

[arXiv:2311.07288](https://arxiv.org/abs/2311.07288)

- **Now also with CMS!**

## [CMS-PAS-TOP-23-001](#)

Available on the CERN CDS information server CMS PAS TOP-23-001

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**CMS Physics Analysis Summary**

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Contact: [cms-pag-conveners-top@cern.ch](mailto:cms-pag-conveners-top@cern.ch) 2024/04/01

Probing entanglement in top quark production with the CMS detector

The CMS Collaboration

**Observation of quantum entanglement in top-quark pair production using  $pp$  collisions of  $\sqrt{s} = 13$  TeV with the ATLAS detector**

The ATLAS Collaboration

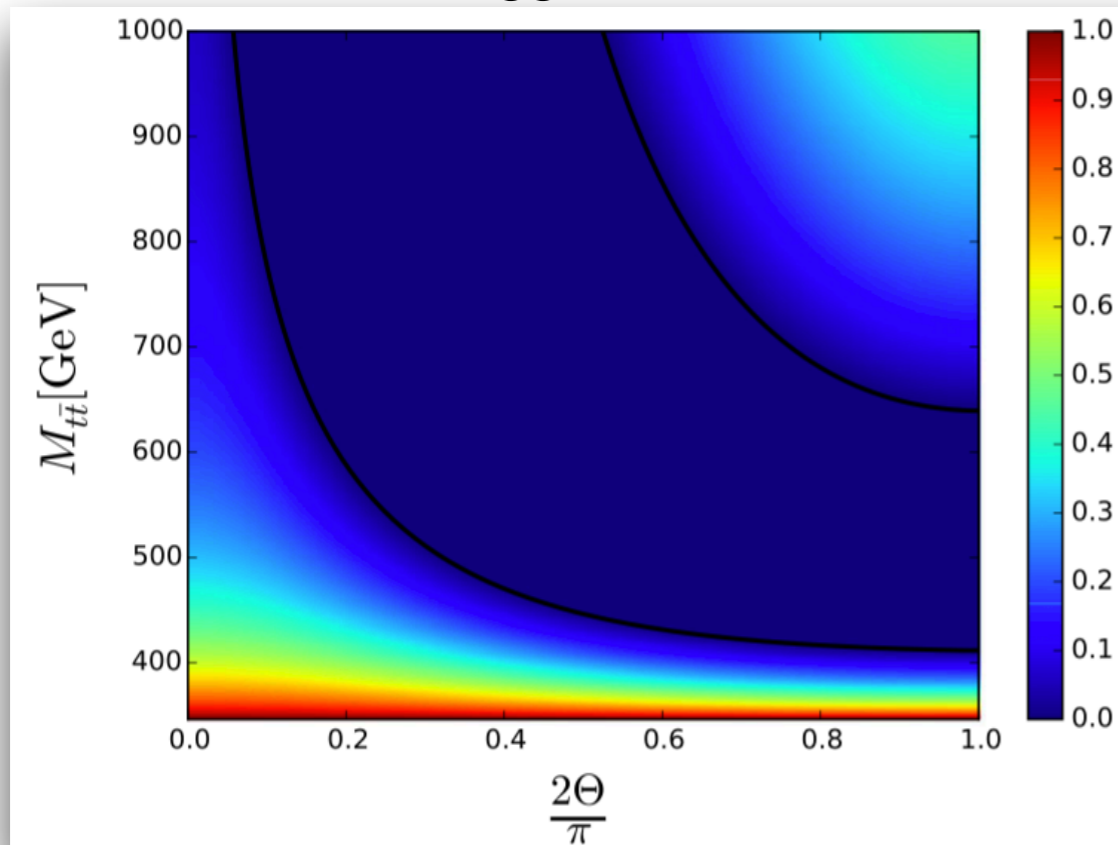
# Entanglement of top quarks

- Top quark = ideal candidate for spin measurements:
  - **extremely short lifetime** allows measuring polarization and spin correlation in  $t\bar{t}$  production
  - **spin information is preserved** in the angular distribution of its decay products

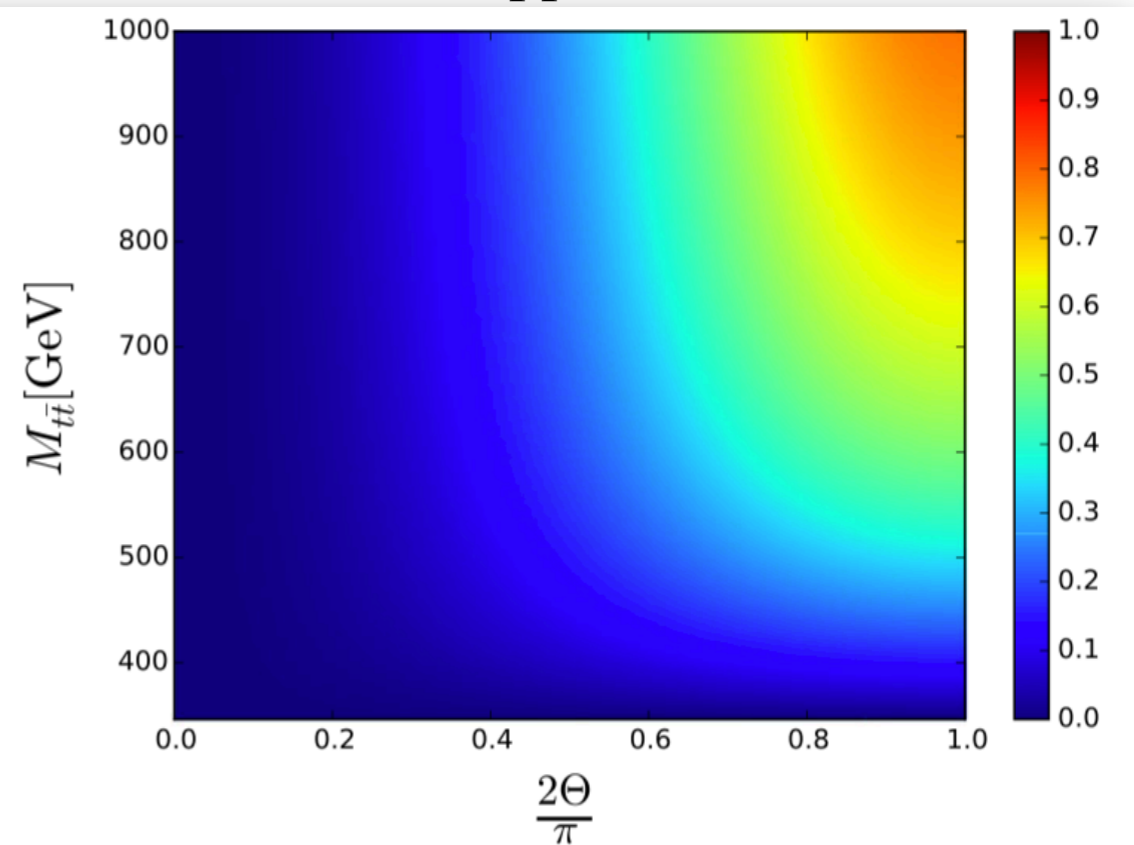
$$\underbrace{\frac{1}{m_t}}_{\substack{\text{production} \\ 10^{-27} \text{ s}}} < \underbrace{\frac{1}{\Gamma_t}}_{\substack{\text{lifetime} \\ 10^{-25} \text{ s}}} < \underbrace{\frac{1}{\Lambda_{\text{QCD}}}}_{\substack{\text{hadronization} \\ 10^{-24} \text{ s}}} < \underbrace{\frac{m_t}{\Lambda^2}}_{\substack{\text{spin-flip} \\ 10^{-21} \text{ s}}}$$

- Entanglement present in top quark pairs can be measured using **spin correlations variables**
- Entanglement depends on production mode,  $m_{t\bar{t}}$ , scattering angle of the top quark ( $\Theta$ )

$$gg \rightarrow t\bar{t}$$



$$q\bar{q} \rightarrow t\bar{t}$$



# How to probe entanglement

- At the LHC, top quarks are produced in a mixed state  
→ can be represented as a density operator:

$$\rho = \frac{I_4 + \sum_i (B_i^+ \sigma^i \otimes I_2 + B_i^- I_2 \otimes \sigma^i) + \sum_{i,j} C_{ij} \sigma^i \otimes \sigma^j}{4}$$

- $B^{+/-}$  = 3-vectors characterizing degree of top quark/antiquark **polarization**
- $C$  = 3x3 matrix characterizing top quark and antiquark **spin correlations**

- **Peres-Horodecki criterion:**

Peres, [Phys. Rev. Lett. 77, 1413](#)  
Horodecki, [Phys. Lett. A 232, 5](#)

if a state is separable (i.e., non-entangled), the transpose with respect to a subspace of the density operator is positive definite

→ a state is non-separable (i.e., entangled) if this condition doesn't hold

→ **top quarks are entangled in a certain phase space if at least one eigenvalue is < 0**

# How to probe entanglement

- **Peres-Horodecki criterion:**  
using simpler observables, a **sufficient condition to observe entanglement in top quarks is:**

$$\Delta = C_{33} + |C_{11} + C_{22}| - 1 > 0 \quad \text{Eur. Phys. J. Plus } \mathbf{136}, 907$$

- At low  $m_{t\bar{t}}$ ,  $C_{11} > 0$  and  $C_{22} > 0 \rightarrow \Delta + 1 = \text{tr}[C] > 1$
- $\text{tr}[C]$  can be probed from a single-differential cross section:

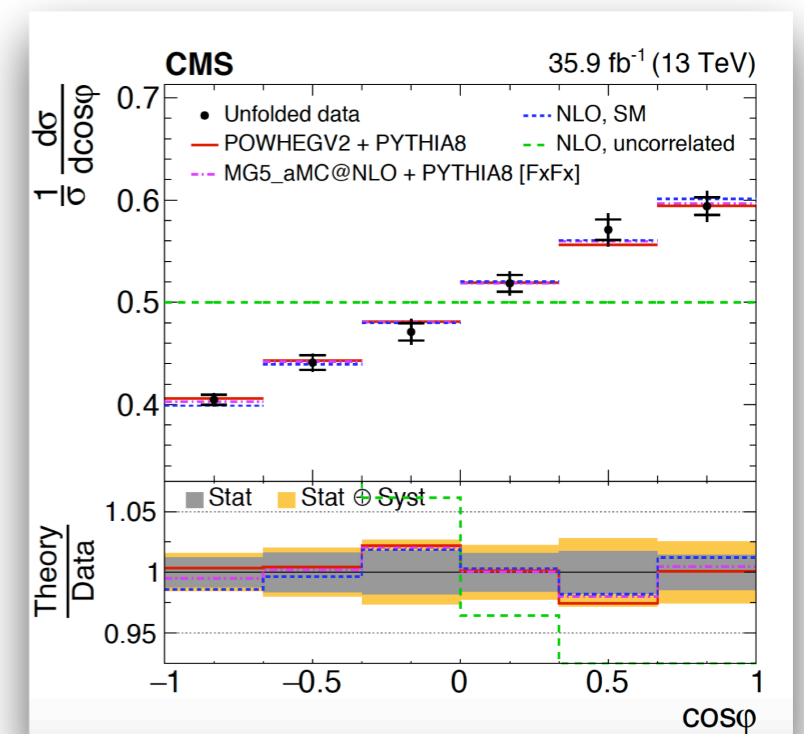
$$\frac{1}{\sigma} \frac{d\sigma}{d\cos\varphi} = \frac{1}{2}(1 - D \cos\varphi) \quad D = -\frac{\text{tr}[C]}{3} \rightarrow D < -1/3$$

**Sufficient condition for entanglement !**

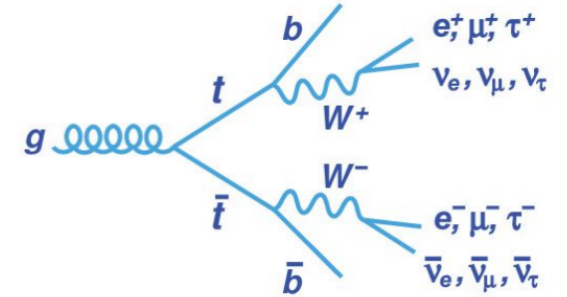
→ measure D to access entanglement information in top quark events!

- $\cos\varphi = \hat{\ell}_1 \cdot \hat{\ell}_2$  is the opening angle between leptons in parent top rest frame  
→ most sensitive and experimentally well measured observable  
→ **focus of entanglement measurement**

Phys. Rev. D 100  
(2019) 072002



# Analysis strategy

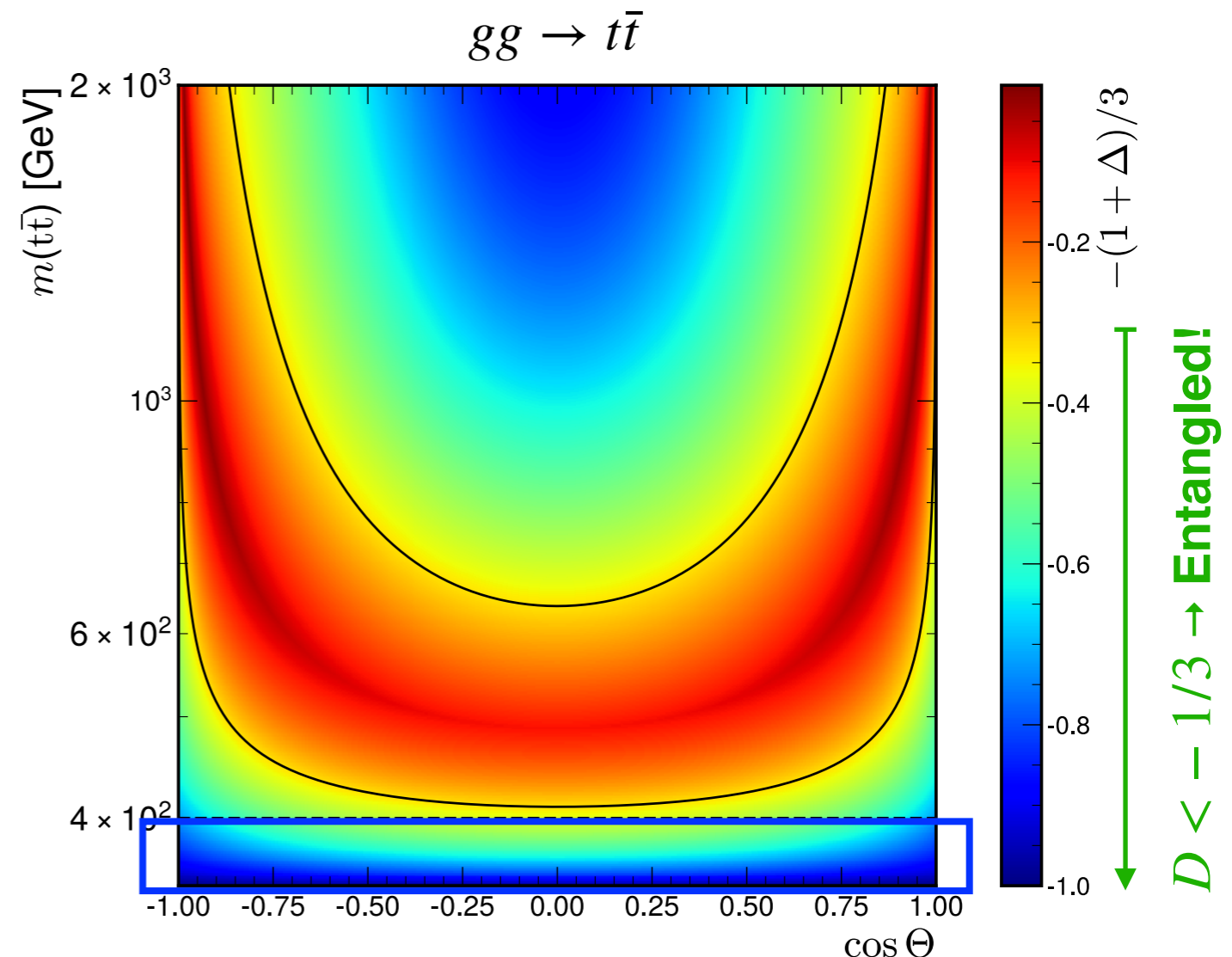


- The degree of entanglement is highly phase space-dependent
  - scan of  $\cos \Theta$  vs  $m_{t\bar{t}}$  to determine most sensitive phase space while minimizing expected total uncertainties
- Focus on **low-mass region** ( $345 < m_{t\bar{t}} < 400$  GeV) to increase entanglement
  - threshold region dominated by  $gg$
  - maximal sensitivity with high statistics
- Cut on velocity along the beam line of the  $t\bar{t}$  system to increase  $gg/q\bar{q}$  fraction:

Aguilar-Saavedra,  
Casas  
[arXiv:2205.00542](https://arxiv.org/abs/2205.00542)

$$\beta = \left| \frac{p_z^t + p_z^{\bar{t}}}{E^t + E^{\bar{t}}} \right| < 0.9$$

- Use **leptonic final states** to measure the helicity angle  $\cos \varphi = \hat{\ell}_1 \cdot \hat{\ell}_2$ 
  - fully encapsulates the spin correlations information for  $gg$  fusion production at low mass
- Perform a **profile maximum likelihood fit of the  $\cos \varphi$  distribution** in the  $m_{t\bar{t}} - \beta$  signal region



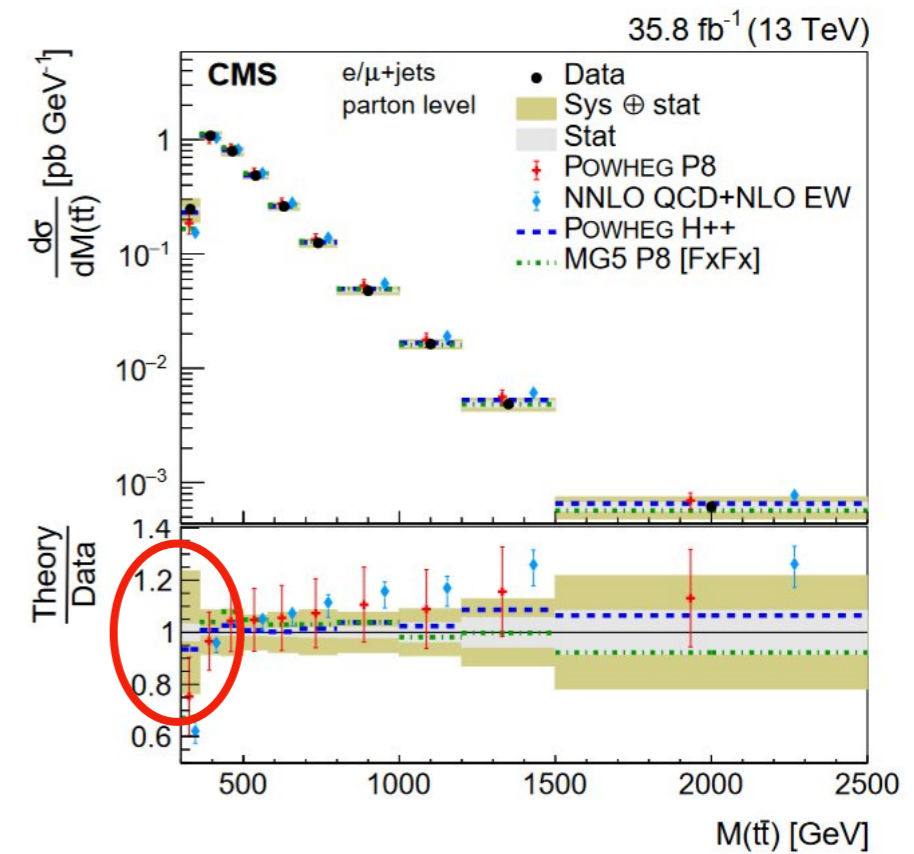
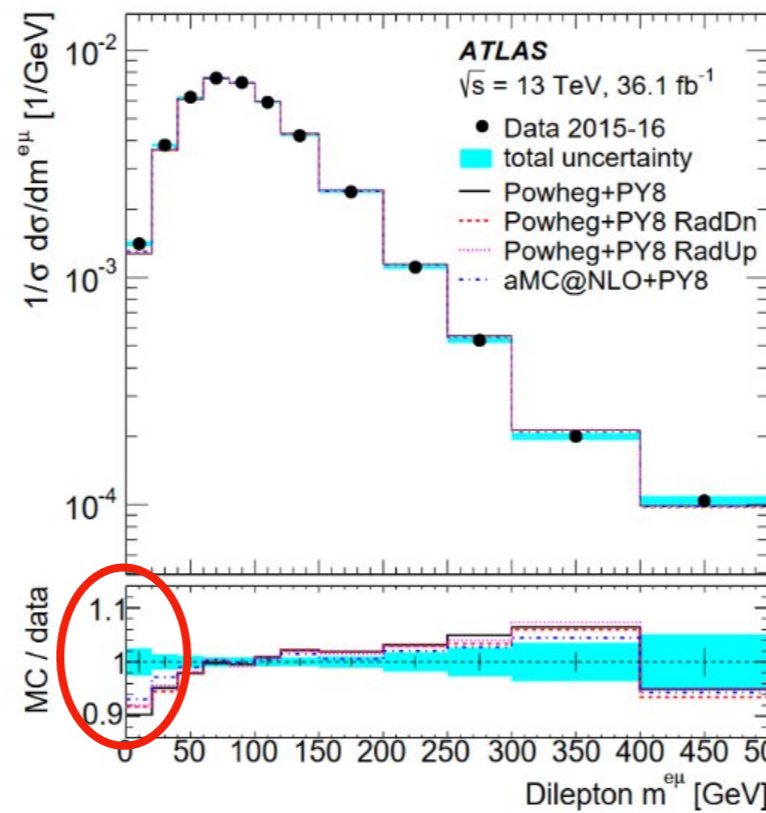
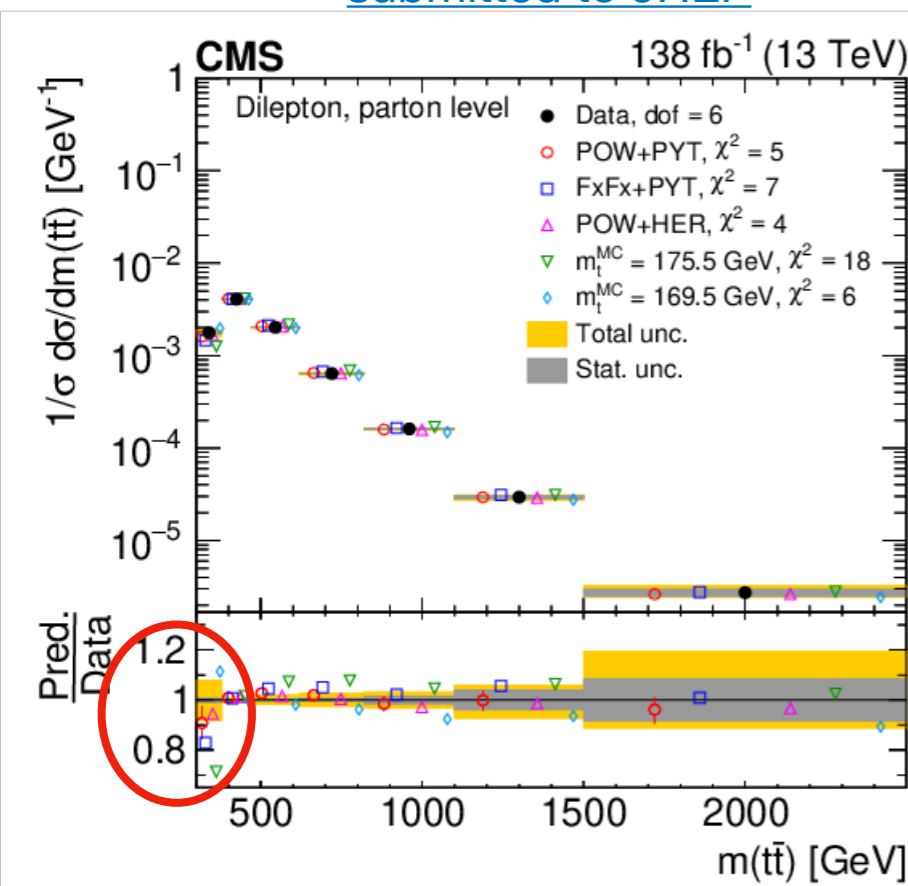
# Threshold region

- Mis-modeling at a level of  $\sim 10\%$  seen for  $m_{t\bar{t}} \sim 345$  GeV ( $m_{e\mu} < 50$  GeV)
- Consistent between dilepton and lepton+jets analyses in both CMS and ATLAS

[arXiv:2402.08486](https://arxiv.org/abs/2402.08486),  
submitted to JHEP

[EPJ C 80, 6](https://arxiv.org/abs/1703.07501)

[Phys. Rev. D 97, 112003](https://arxiv.org/abs/1703.07501)



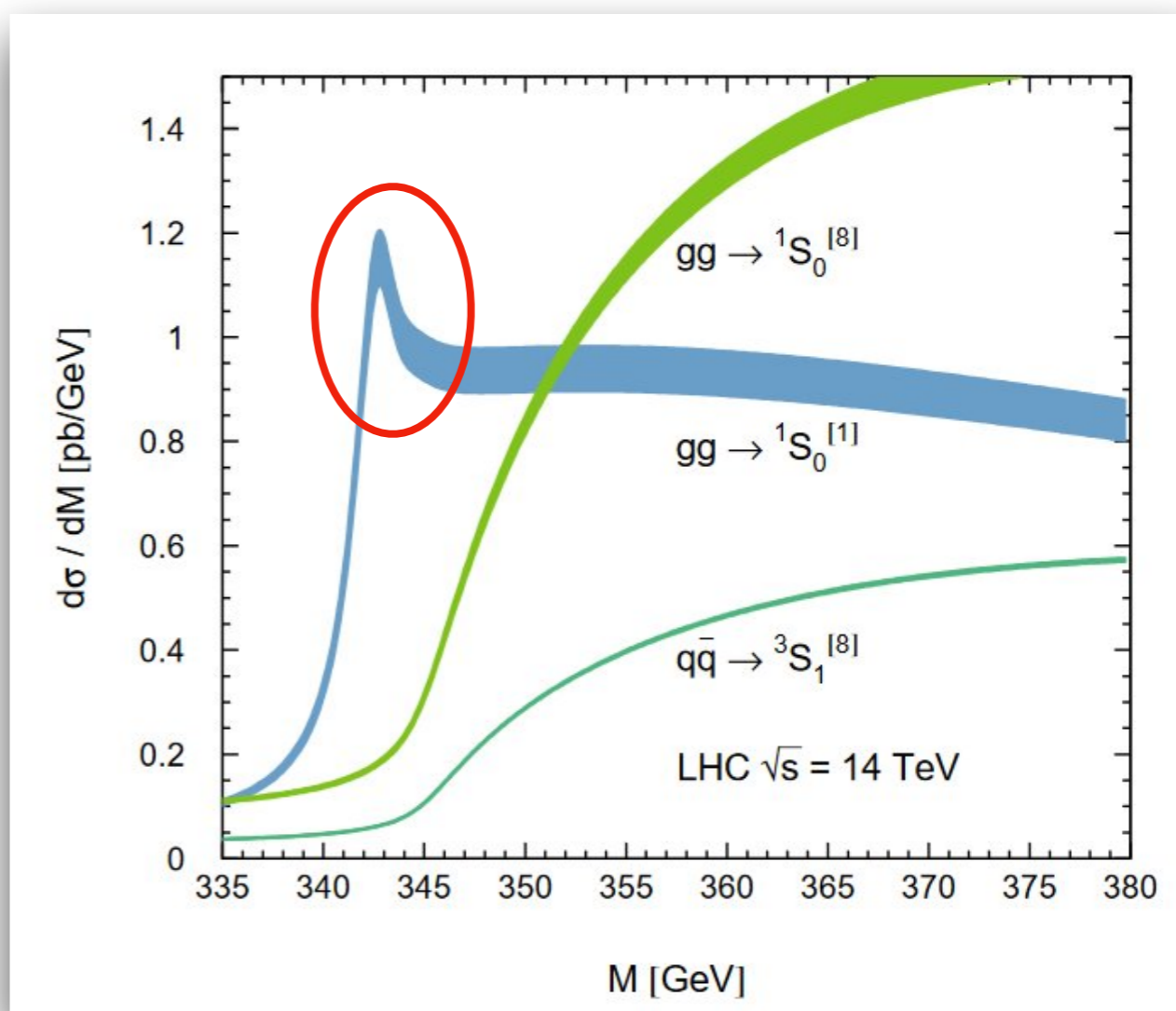
# Threshold region

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- Consistent between dilepton and lepton+jets analyses in both CMS and ATLAS
- NRQCD contributions close to threshold
  - spin and color singlet state ( $\eta_t$ ): maximally entangled **toponium**
- Excess seen could come from toponium ?

[JHEP 06, 158](#)

→ **inclusion of toponium ( $\eta_t$ ) contributions in our signal model**  
using simplistic model based on  
[Phys Rev D 104 034023](#)

Toponium = predicted top quark-antiquark quasi-bound state with a mass of 343 GeV and width of 7 GeV





# Dataset and signal model

- Current analysis = extension of 2016 top quark spin correlations analysis in dilepton events

- 35.9 fb<sup>-1</sup> of data @13 TeV collected in 2016

Phys. Rev. D 100 (2019) 072002

- **Combined signal model:  $t\bar{t}$  + toponium ( $\eta_t$ )**

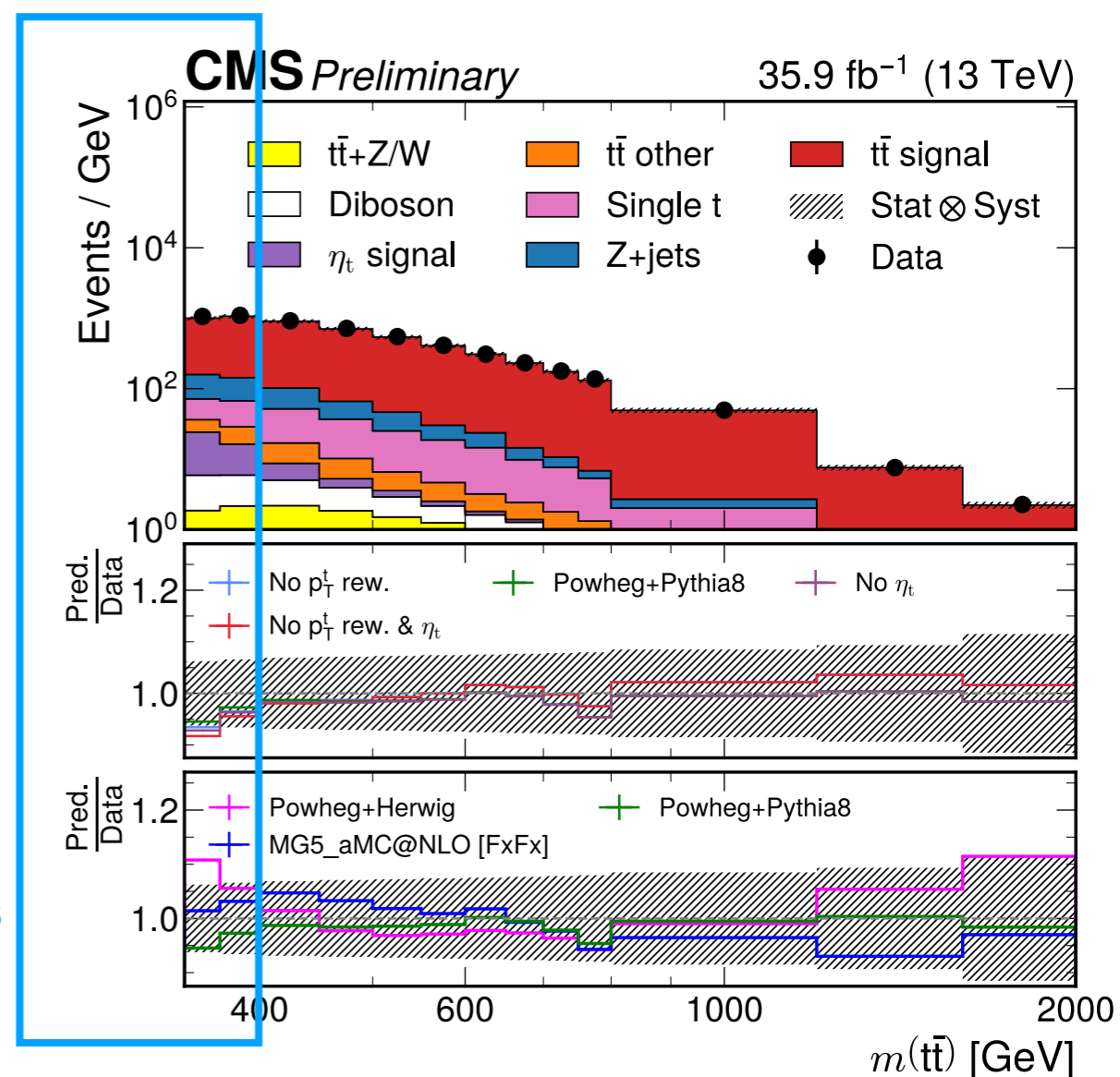
- PowhegBox+Pythia8 as nominal  $t\bar{t}$  sample
- PowhegBox+Herwig and MG5 aMC@NLO(+MadSpin) [FxFx] as alternative  $t\bar{t}$  samples

- $\eta_t$  improves data modeling in the threshold region

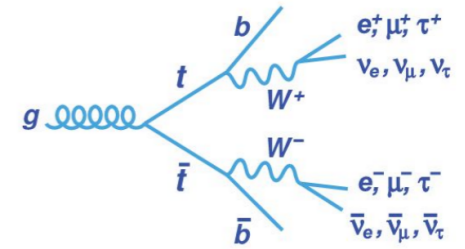
- only spin-0  $\eta_t$  accounted (colour singlet pseudoscalar state) [[PRD 104 \(2021\) 034023](#)]
- toponium model generated with MG5 aMC@NLO(LO)+Pythia8 with  $337 < m_{\eta_t} < 349$  GeV

- **Main background sources:**

- Z+jets (MG5\_aMC@NLO + data-driven corrections)
- single top (Powheg MC)
- diboson (Pythia8 MC)



# Event selection



- Current analysis = extension of 2016 top quark spin correlations analysis in dilepton events
  - same strategy for event selection, kinematic reconstruction, and background estimation
  - **optimized sensitivity for entanglement measurement**

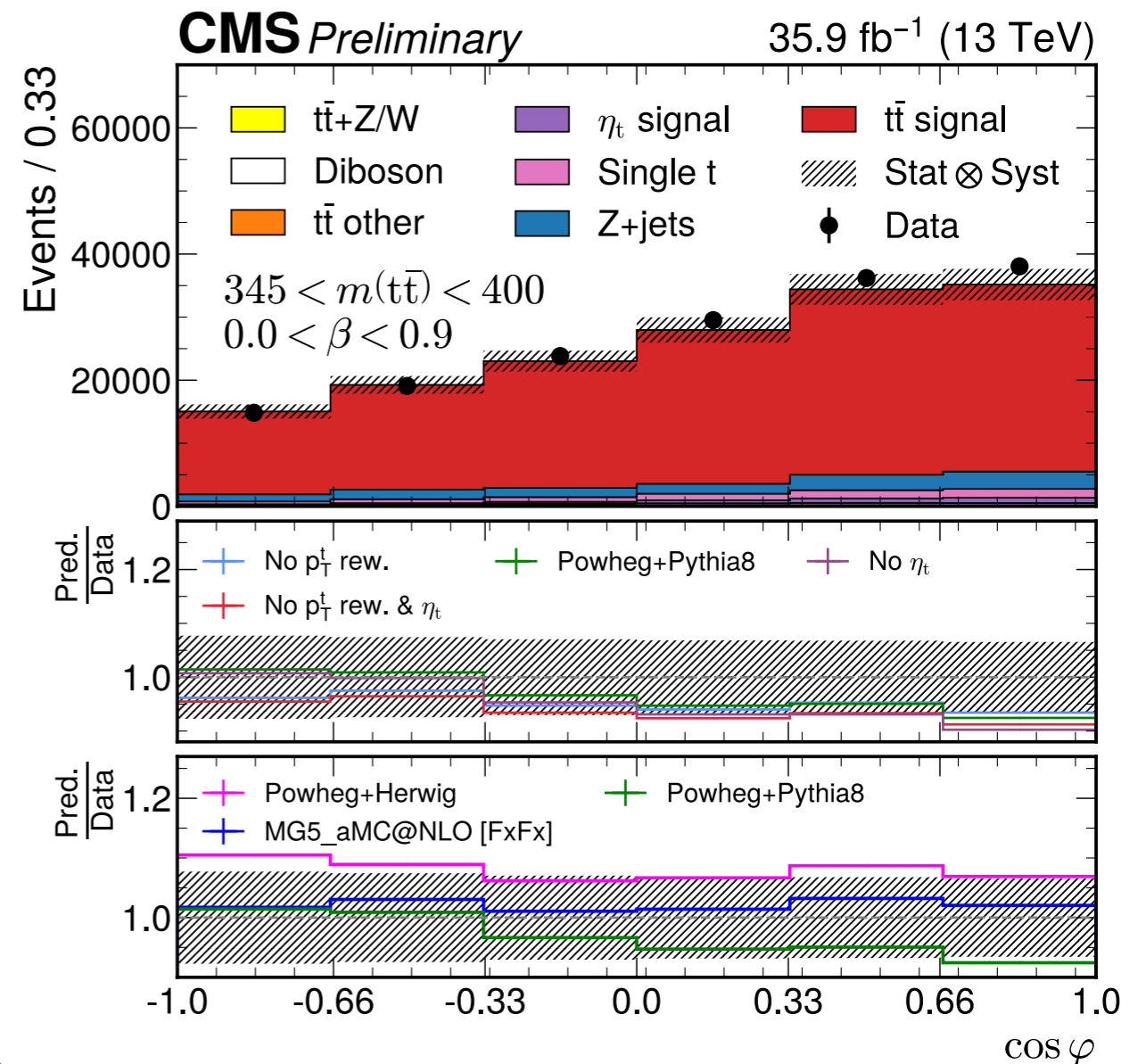
Phys. Rev. D 100 (2019) 072002

- 2 oppositely charged isolated leptons (ee, eμ and μμ)
  - **including also leptons from tau decays** (different from 2016 analysis)
  - $p_T > 25(20)$  GeV, for leading(trailing) lepton and  $|\eta| < 2.4$
  - veto events with more than two leptons
  - reject events with  $m_{\ell\bar{\ell}} < 20$  GeV
  - single lepton + dilepton triggers

- $\geq 2$  jets (R=0.4),  $\geq 1$  b jet
  - $p_T > 30$  GeV and  $|\eta| < 2.4$
  - jet cleaning:  $\Delta R(\ell, \text{jet}) > 0.4$

- ee, μμ channels:
  - $E_{\text{miss}}^T > 40$  GeV
  - Z veto:  $|m_Z - m_{\ell\bar{\ell}}| > 15$  GeV

- Top quark reconstruction with  $m_{\ell b}$  weighting method
  - take solution with smallest  $m_{t\bar{t}}$



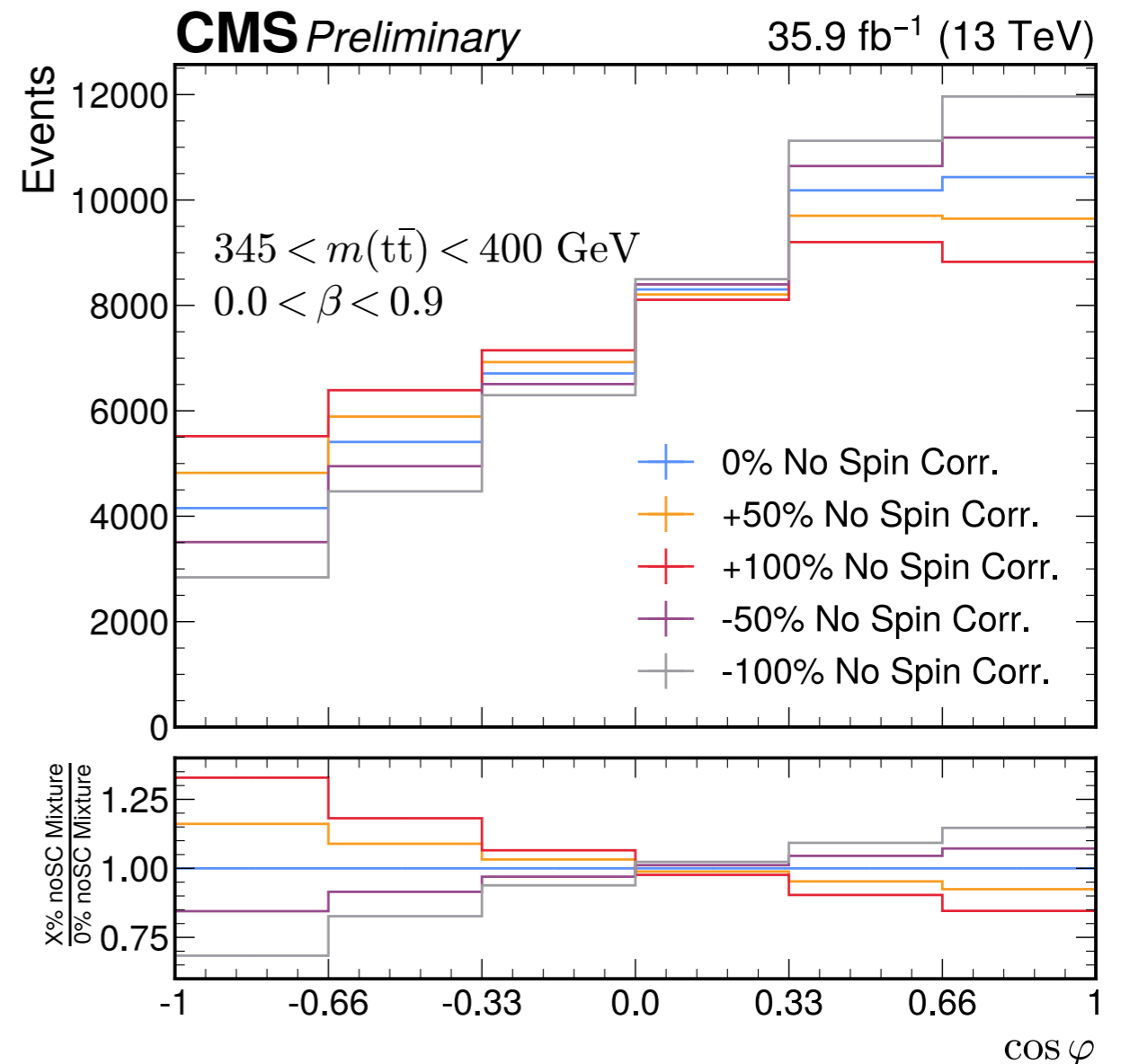
# Extraction of entanglement proxy

- The entanglement proxy  $D$  is extracted with a template fit
  - all systematic effects included as nuisances
- How can we create variations of  $D$  outside of SM?
  1. generate top quark pairs with no spin correlations  $\rightarrow D = 0$  (noSC samples)
  2. create new samples with mixtures of SM and noSC to obtain  $D \in [D_{SM}, 0]$
  3. extend the fit for variations of  $[-1, D_{SM}]$

- Use mixtures of SC and noSC to change fraction of  $t\bar{t}$  with aligned vs opposite spins  
 $\rightarrow$  any value of  $D$  between -1 and +1 can be reached

$$D \sim \frac{\sigma(\uparrow\uparrow) + \sigma(\downarrow\downarrow) - \sigma(\uparrow\downarrow) - \sigma(\downarrow\uparrow)}{\sigma(\uparrow\uparrow) + \sigma(\downarrow\downarrow) + \sigma(\uparrow\downarrow) + \sigma(\downarrow\uparrow)}$$

Mixed samples with and without spin correlations



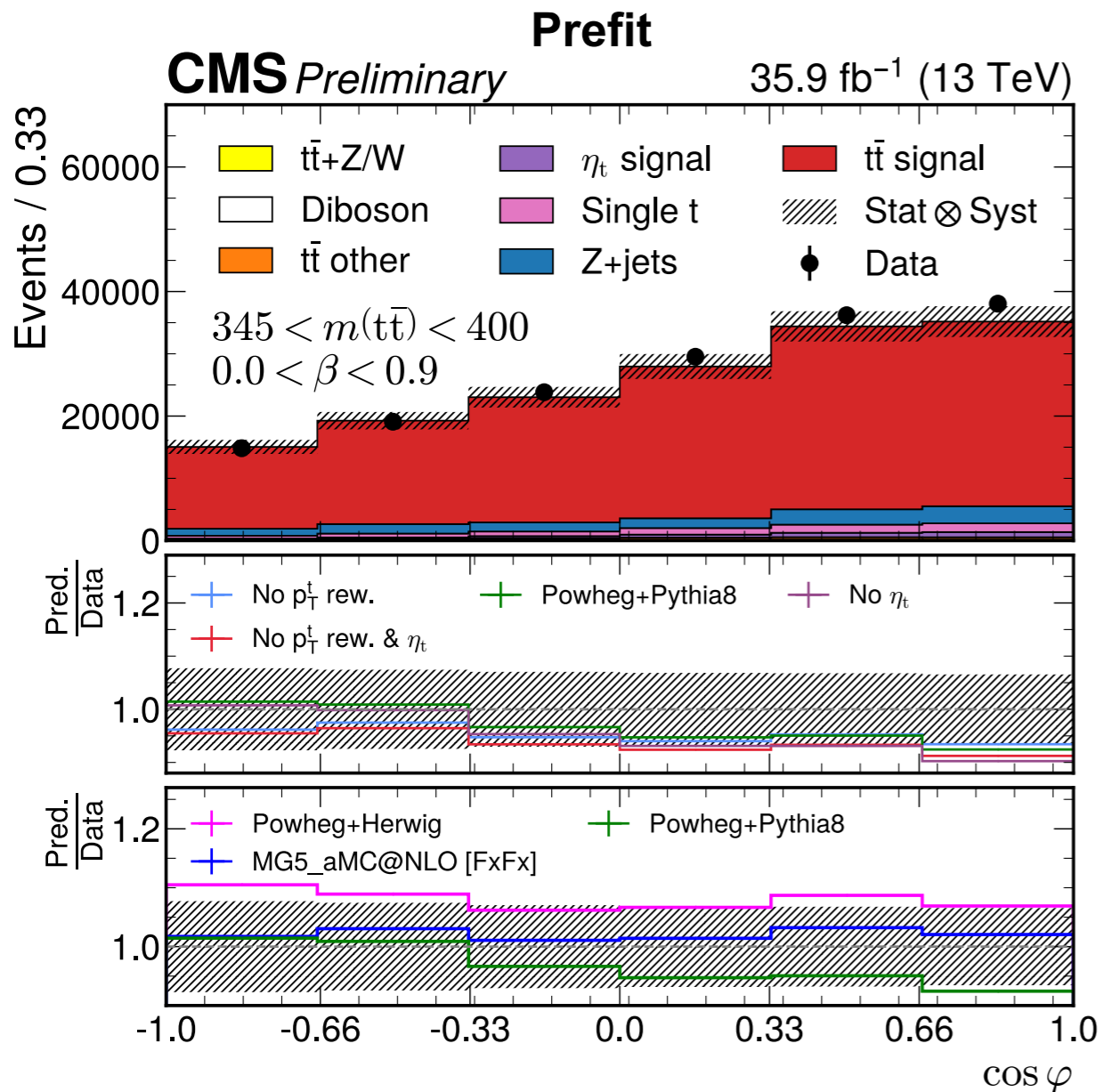
# Systematic uncertainties

- **Current analysis = extension of 2016 top quark spin correlations analysis in dilepton events**
  - same uncertainties considered + **additional ones for toponium**:
    - a flat uncertainty of 50% is applied on toponium
    - a binding energy uncertainty of  $\pm 0.5$  GeV is considered
- **Breakdown of leading syst. unc. in the entanglement proxy  $D$  at the post-fit level**
- **Leading experimental uncertainties:**
  - Jet energy scale and resolution
- **Leading theory-based uncertainties:**
  - Toponium normalization
  - Parton Shower

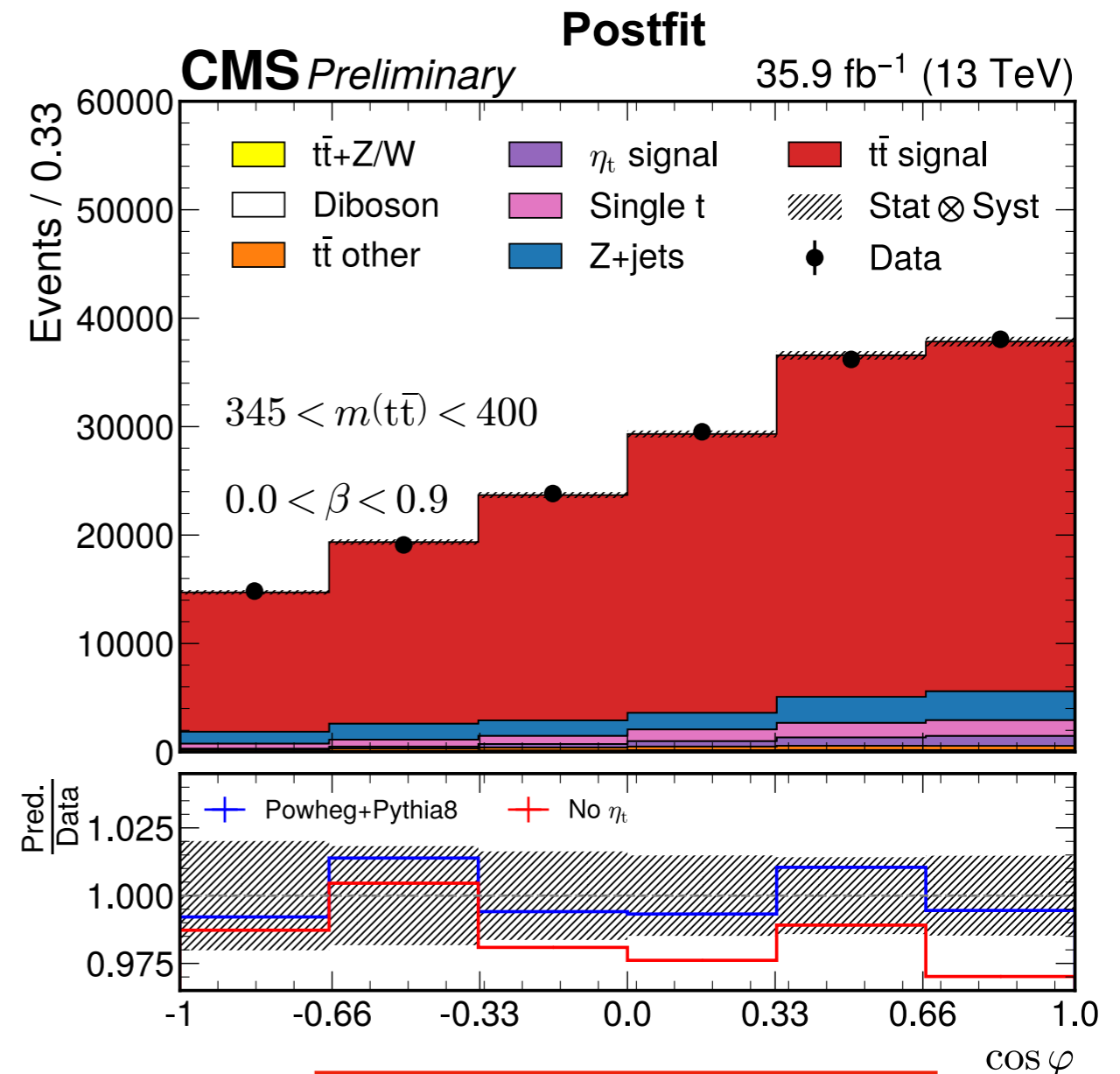
Source	Uncertainty $D$
JES	10.1%
Toponium normalization	10.1%
Parton Shower (ISR)	6.3%
Scale	1.8%
Parton Shower (FSR)	1.2%
JER	0.9%
Z+jets shape	0.8%
b quark fragmentation	0.4%
$t\bar{t}$ normalization	0.3%
PDF	0.3%

# Results

- Result of the binned profile likelihood fit of the  $\cos \varphi$  distribution
  - ~47500 signal candidates
- Good agreement with SM predictions



→



No  $\eta_t = \eta_t$  removed from signal without repeating the fit

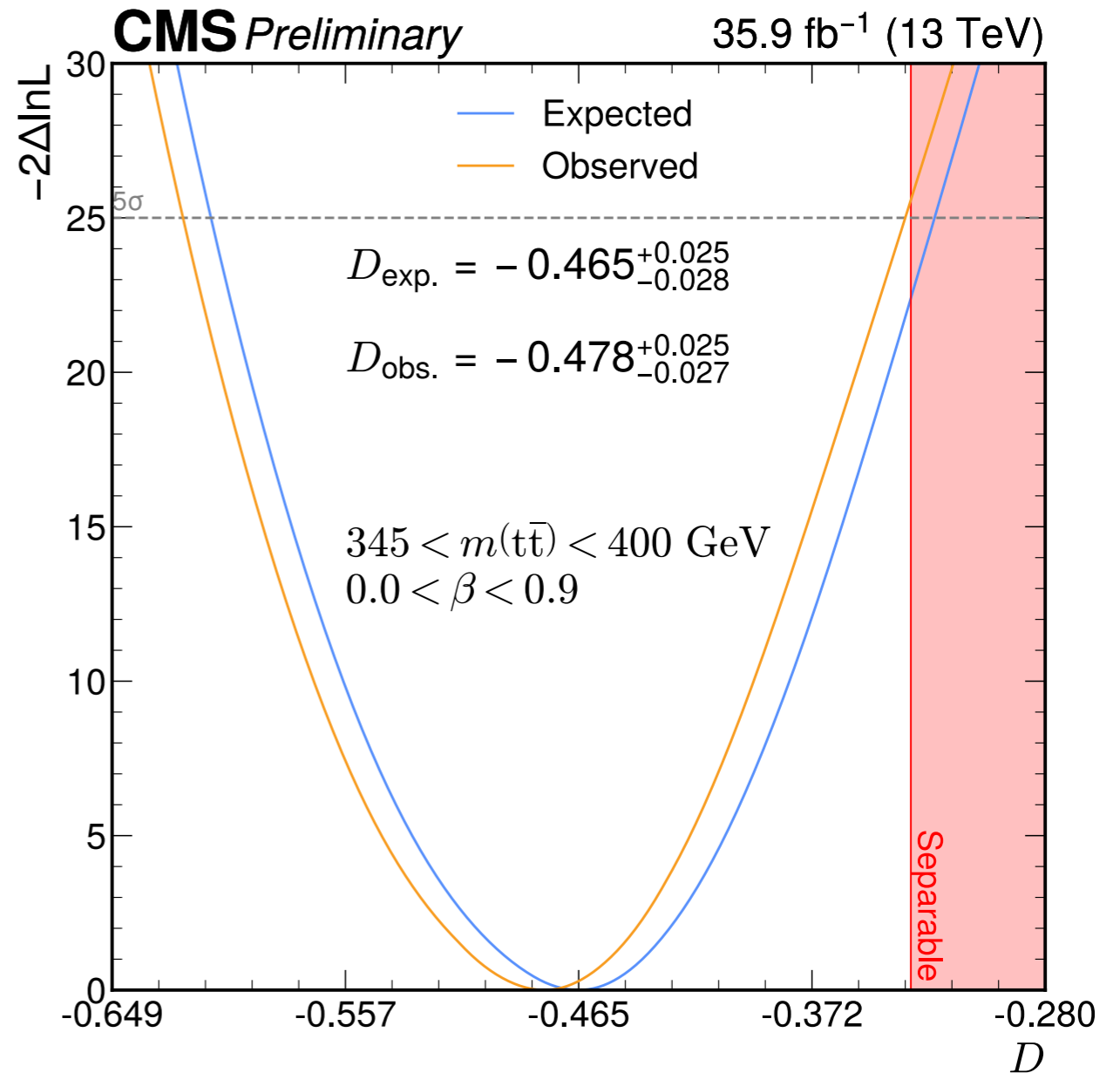
# Results

- Scan of the  $-2\Delta\ln L$  distribution yields  $D$  at parton level, accounting for all detector effects

$$D_{obs} = -0.478 \pm 0.017(\text{stat})_{-0.021}^{+0.018}(\text{syst})$$

$$D_{exp} = -0.465_{-0.017}^{+0.016}(\text{stat})_{-0.022}^{+0.019}(\text{syst})$$

**>5 standard deviations observation  
of top quarks being entangled at  $t\bar{t}$  threshold !**



# Results

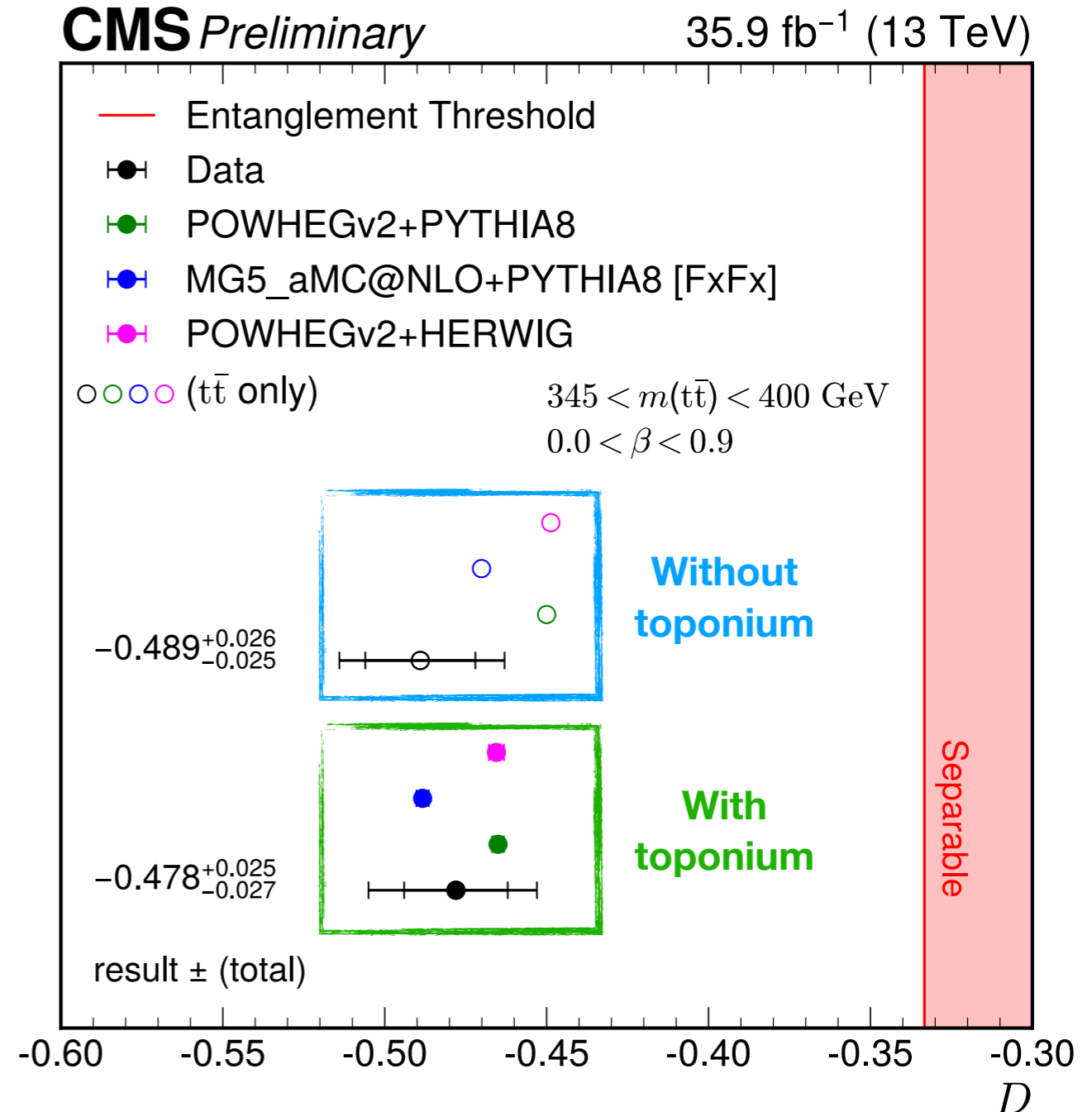
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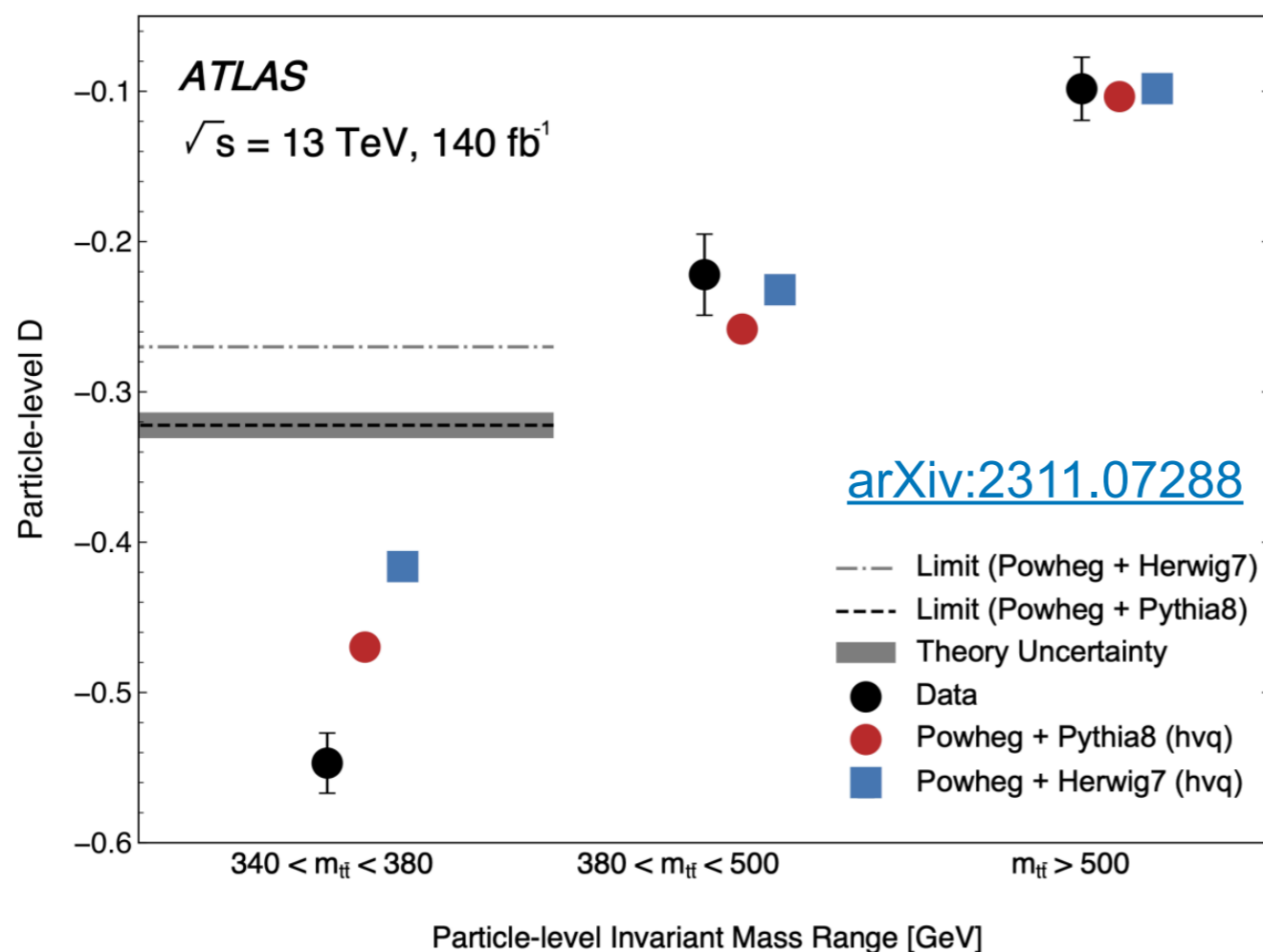
>5 standard deviations observation of top quarks being entangled at  $t\bar{t}$  threshold !

- Good agreement with SM predictions
  - significantly improved with  $\eta_t$  inclusion

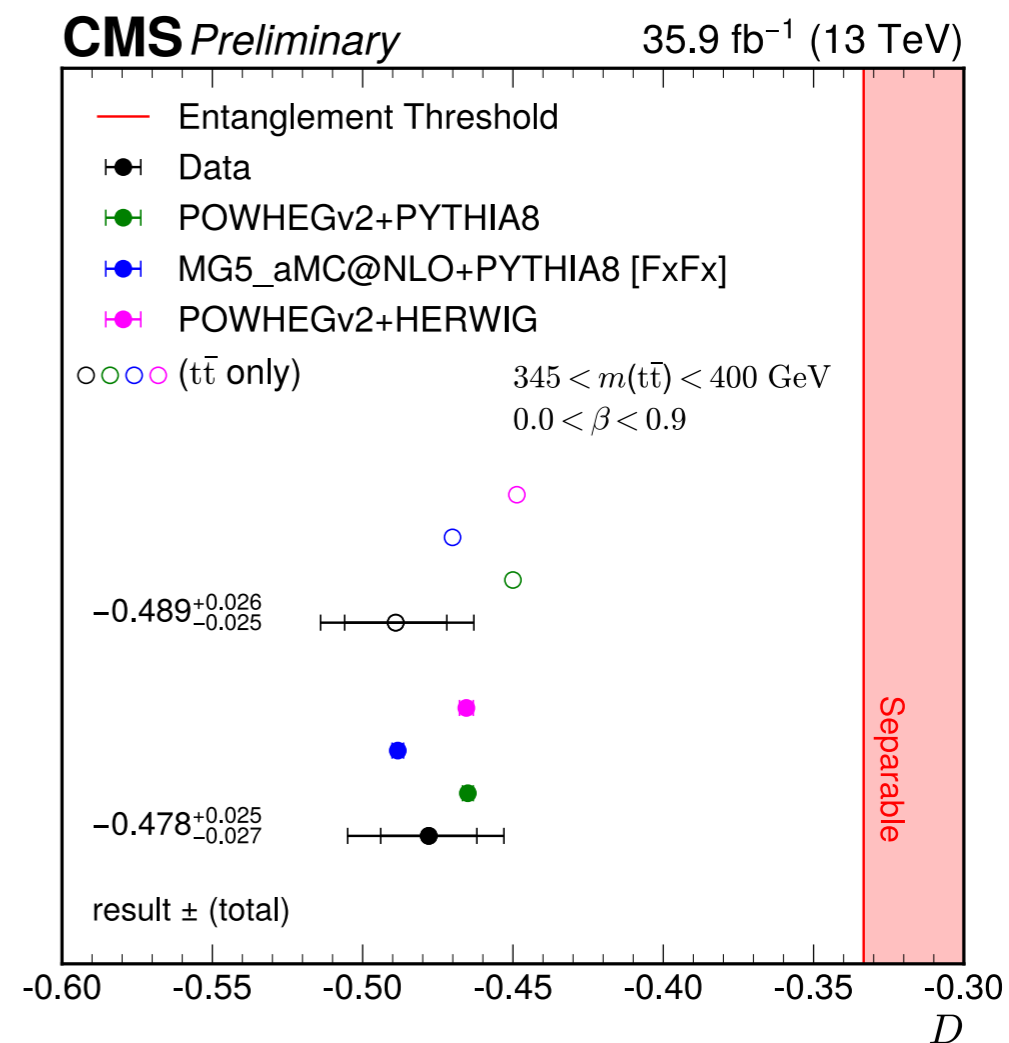


# Comparison with ATLAS

- **Entanglement in top quark observed by both ATLAS and CMS with >5 standard deviations!**
- No clear preference for a specific MC prediction
- **Both analyses are dominated by systematic uncertainty**
- Total (stat.) uncertainty is an order of magnitude larger in the CMS analysis
- Total (syst.) uncertainty is similar between ATLAS & CMS, but different systematics are considered



ATLAS: limit of  $D = -1/3$  is folded from parton to **particle-level**



CMS: limit of  $D = -1/3$  is shown at **parton-level**

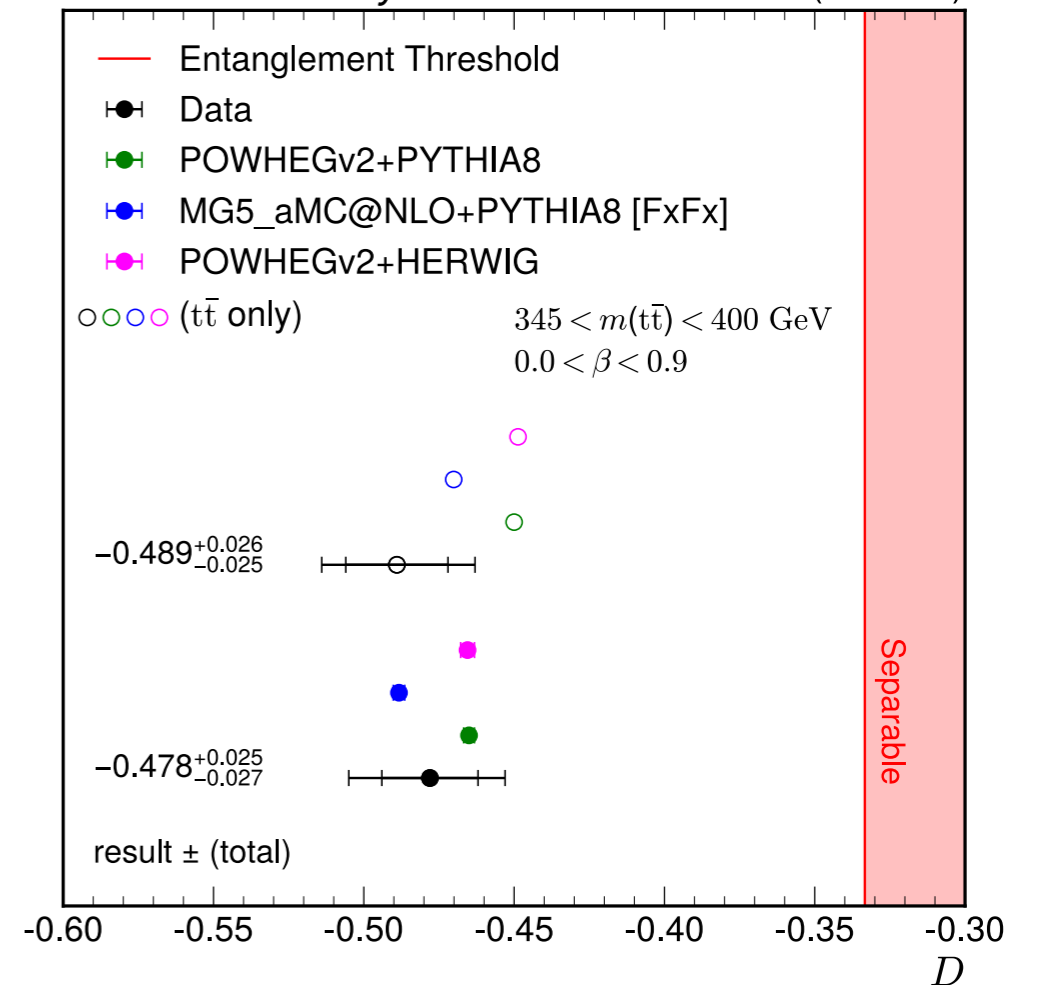


# Conclusions

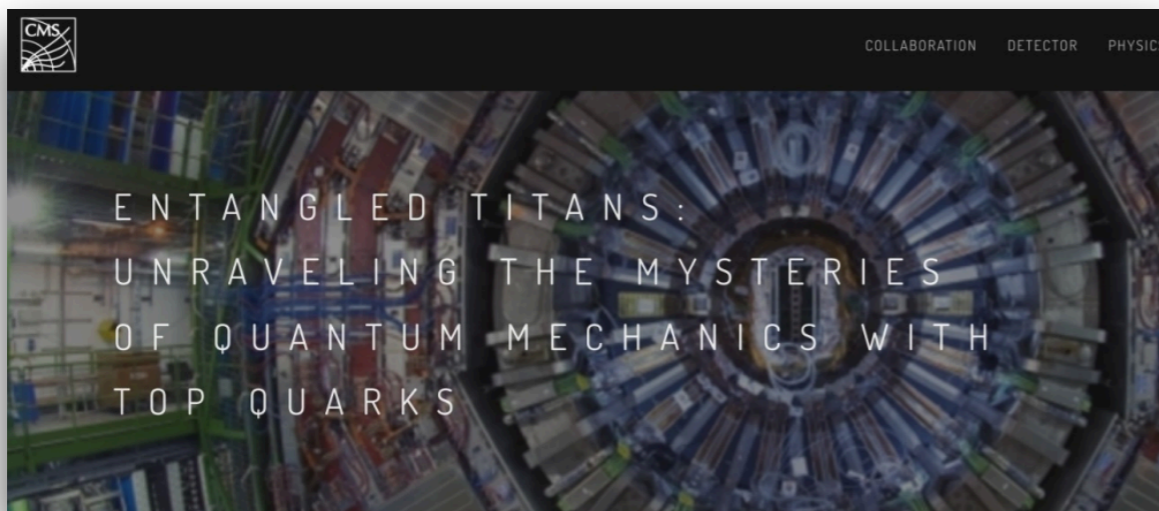
- **First observation of entanglement between top quarks with CMS data**
- One of few quantum information studies in high energy physics
- Even in presence of a “toponium” bound state, **we confirm the existence of entanglement in the  $t\bar{t}$  system with  $> 5$  standard deviations**
- A better modeling next to the production threshold is required  $\rightarrow$  theory community is working on improving the prediction of mainstream generators for precision measurements



CMS Preliminary 35.9 fb<sup>-1</sup> (13 TeV)



[CMS briefing](#)



**BACKUP**

# Top quark reconstruction

- Use algebraic method to solve for neutrino 3-vectors
- Results in quartic equation for neutrino momenta
- Pick solution with lowest  $m_{t\bar{t}}$
- Repeat process 100x for leptons and b jets smeared within resolution
- Weight solutions by the  $m_{\ell b}$  distribution

$$0 = \sum_{i=0}^4 c_i(m_t, p_{\ell^+}, p_{\ell^-}, p_b, p_{\bar{b}}) p_x(\bar{\nu})^i$$

$$\begin{aligned} H_x &= p_{\nu_x} + p_{\bar{\nu}_x} \\ H_y &= p_{\nu_y} + p_{\bar{\nu}_y} \end{aligned}$$

$$\begin{aligned} m_{W^+}^2 &= (E_{\ell^+} + E_{\nu})^2 - (p_{\ell_x^+} + p_{\nu_x})^2, \\ &\quad - (p_{\ell_y^+} + p_{\nu_y})^2 - (p_{\ell_z^+} + p_{\nu_z})^2, \\ m_{W^-}^2 &= (E_{\ell^-} + E_{\bar{\nu}})^2 - (p_{\ell_x^-} + p_{\bar{\nu}_x})^2, \\ &\quad - (p_{\ell_y^-} + p_{\bar{\nu}_y})^2 - (p_{\ell_z^-} + p_{\bar{\nu}_z})^2, \\ m_t^2 &= (E_b + E_{\ell^+} + E_{\nu})^2 - (p_{b_x} + p_{\ell_x^+} + p_{\nu_x})^2, \\ &\quad - (p_{b_y} + p_{\ell_y^+} + p_{\nu_y})^2 - (p_{b_z} + p_{\ell_z^+} + p_{\nu_z})^2, \\ m_{\bar{t}}^2 &= (E_{\bar{b}} + E_{\ell^-} + E_{\bar{\nu}})^2 - (p_{\bar{b}_x} + p_{\ell_x^-} + p_{\bar{\nu}_x})^2, \\ &\quad - (p_{\bar{b}_y} + p_{\ell_y^-} + p_{\bar{\nu}_y})^2 - (p_{\bar{b}_z} + p_{\ell_z^-} + p_{\bar{\nu}_z})^2. \end{aligned}$$

# Mixtures of SC and noSC

- In order to have templates implementing an alternative value of the entanglement proxy  $D$ , we employ the noSC sample and “mix” it in steps ranging from  $-100\%$  to  $100\%$  with the combined signal model SM template
- The negative mixtures are created mirroring the corresponding positive mixtures around the  $0\%$  noSC mixture, i.e., the nominal combined signal model
- Any particular mixture of combined SC and noSC signal corresponds to a certain value of  $D$  at the parton level by means of calculating a 2-bin asymmetry:

$$A_D = (N(\cos \varphi > 0) - N(\cos \varphi < 0)) / (N(\cos \varphi > 0) + N(\cos \varphi < 0))$$

yields  $D$  as  $-2 \cdot A_D$ , with  $N$  always being the sum of  $t\bar{t}$  and  $\eta_t$ .

