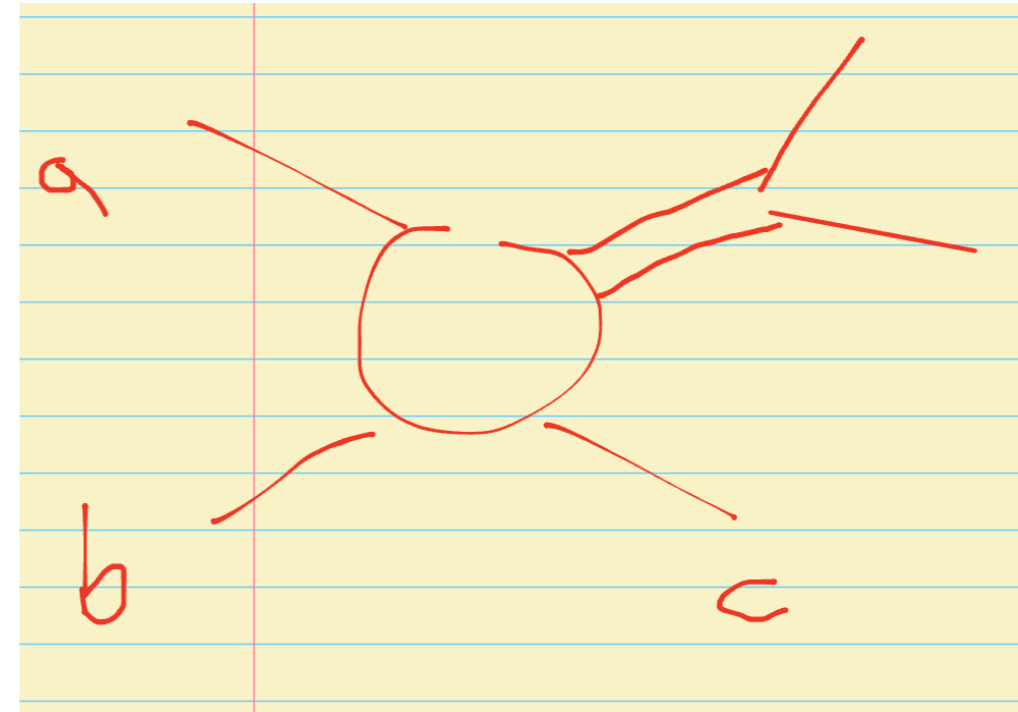


# Production of (strange) resonances (Importance of production)

- A few general remarks; amplitude analysis, interpretation of the experiment and theory.
- Physics of quasi-elastic photo-production :  $\pi/K$  exchange
- Consequences for spectroscopy

Adam Szczepaniak (IU/JLab)

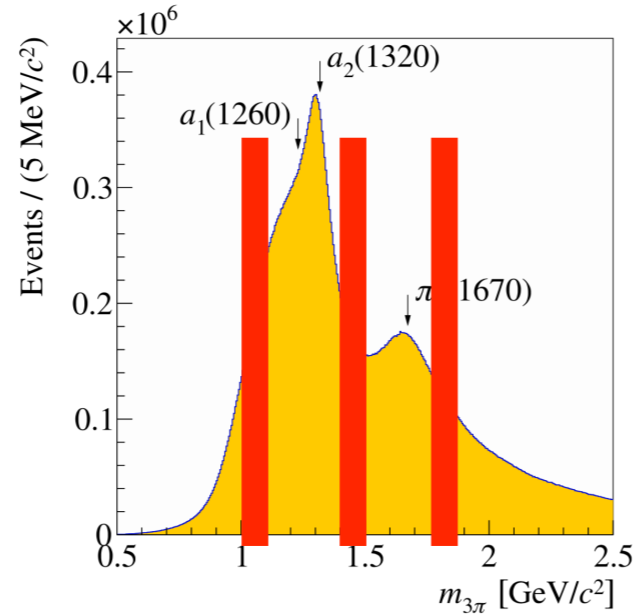


**J***PAC*

**EXOHAD**  
EXOTIC HADRONS TOPICAL COLLABORATION



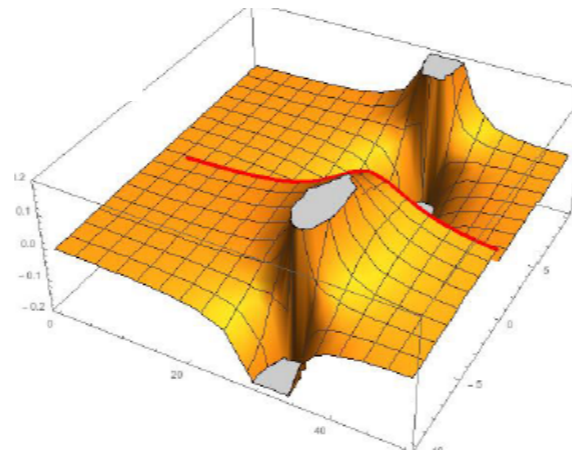
# Nature/QCD : kinematical variables are real



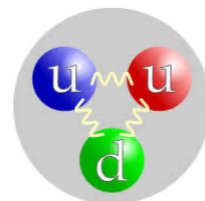
Less predictive power ✗  
Some physical interpretation ✗  
Minimally biased ✓

Dynamics : complex variables

Predictive power ✓  
Physical interpretation ✓  
(within the model! ✗)  
Biased by the input ✗



Physical interpretation: models



1. Covariant amplitudes fit data  $A(s, t)$ ,  
Physically: satisfy complicated unitarity relations
2. Partial waves  $A_l(s)$  satisfy simple unitarity relations, directly related to physics but have complicated analytical properties.
3. Ultimately physics is determined by singularities in energy **and in angular momentum**

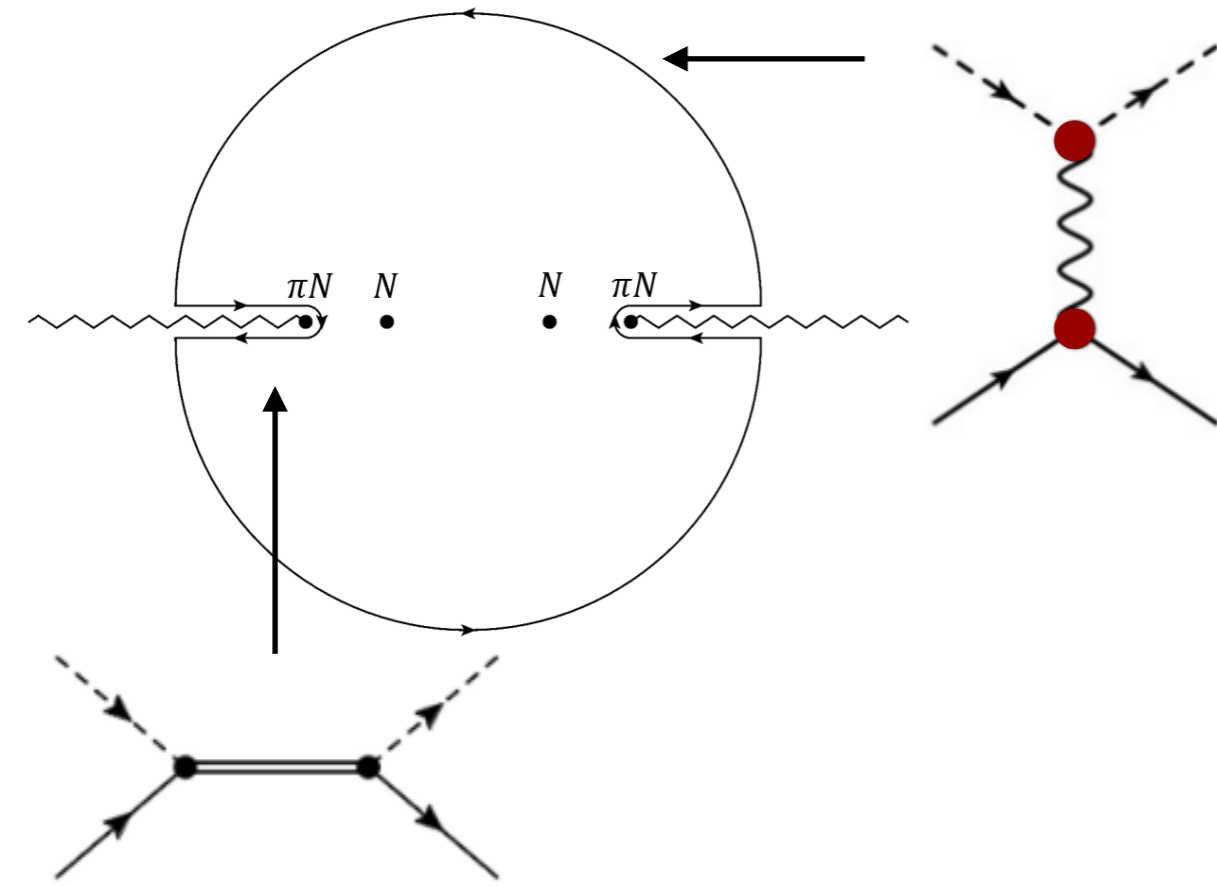
Experiment



Theory

$$A_l(s) = A(s, l)$$

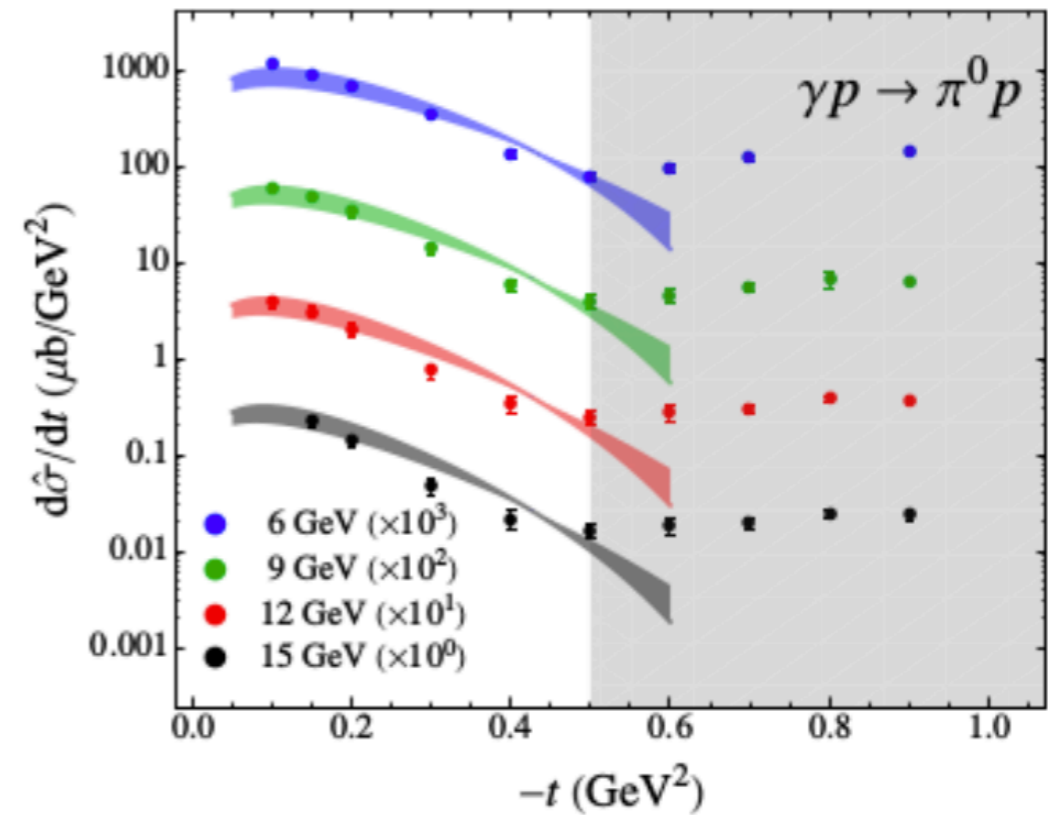
# Finite Energy Sum Rules



$$\int_0^{\Lambda} \text{Im } A_i(\nu, t) \nu^k d\nu = \beta_i(t) \frac{\Lambda^{\alpha(t)+k}}{\alpha(t) + k}$$

$$\beta_i(t) = \frac{\alpha(t) + k}{\Lambda^{\alpha(t)+k}} \int_0^{\Lambda} \text{Im } A_i(\nu, t) \nu^k d\nu$$

- High s-data calculated from s-channel p.w.a



V.Mathieu, et al. (JPAC) 2018

- Resonance parameters (low-s) can be constrained by high-s data

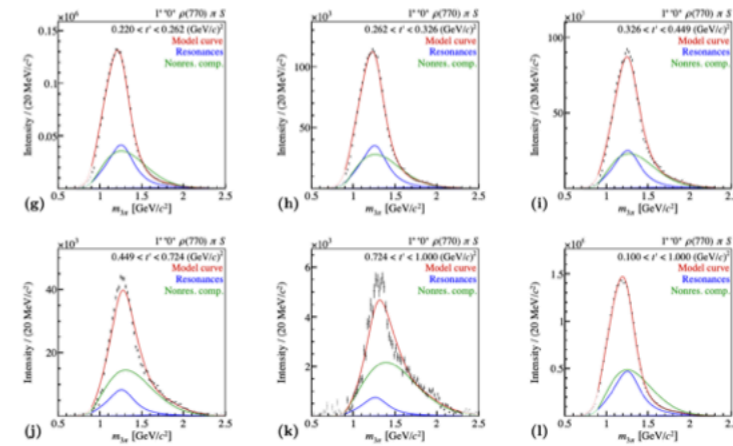
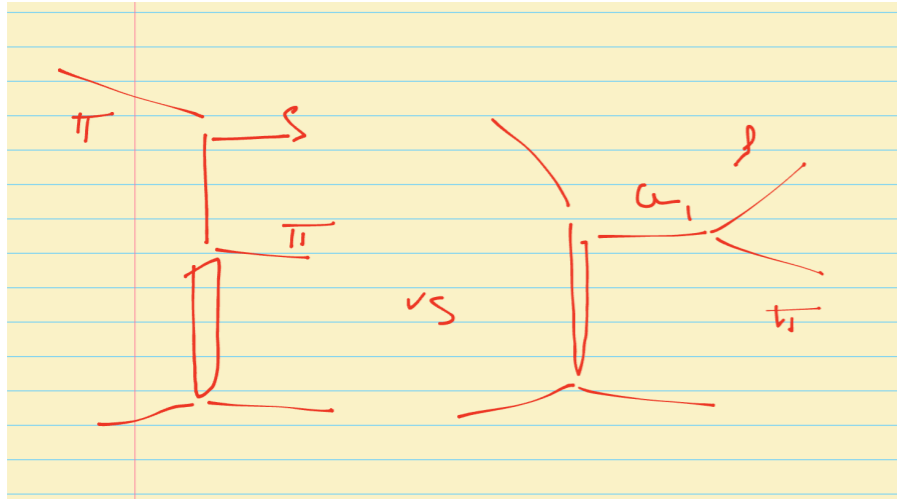
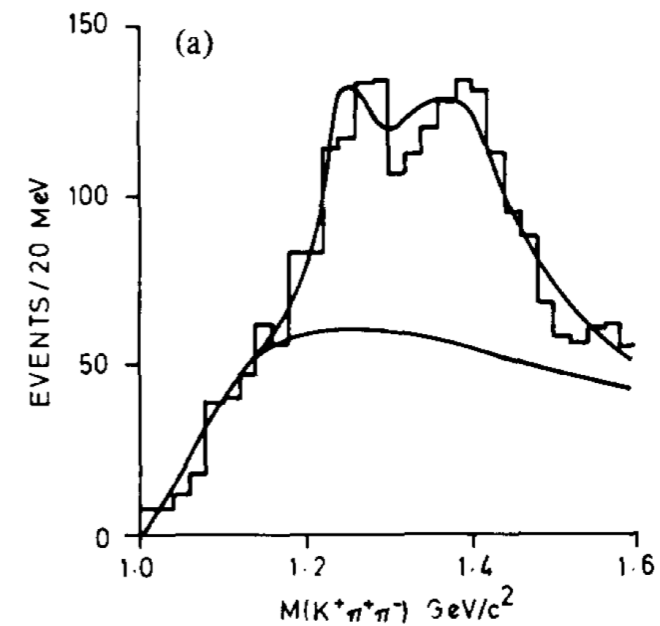


FIG. 30: (a) to (k): Intensity distributions of the  $1^{++}0^+\rho(770)\pi S$  wave in the 11  $t'$  bins. (l) The  $\rho(1640)$  component is so small that it is barely visible in linear scale.

- The  $a_1(1260)$  was for some time confused with the pion exchange

- Analogous  $K\pi\pi$  analysis: duality between  $K_1(1270), K_1(1400), K_1(1650)(?)$  K-exchange



M.G.Bowler et al. (CERN 1973)

Compare with 720 K events at COMPASS (see S.Wallner)

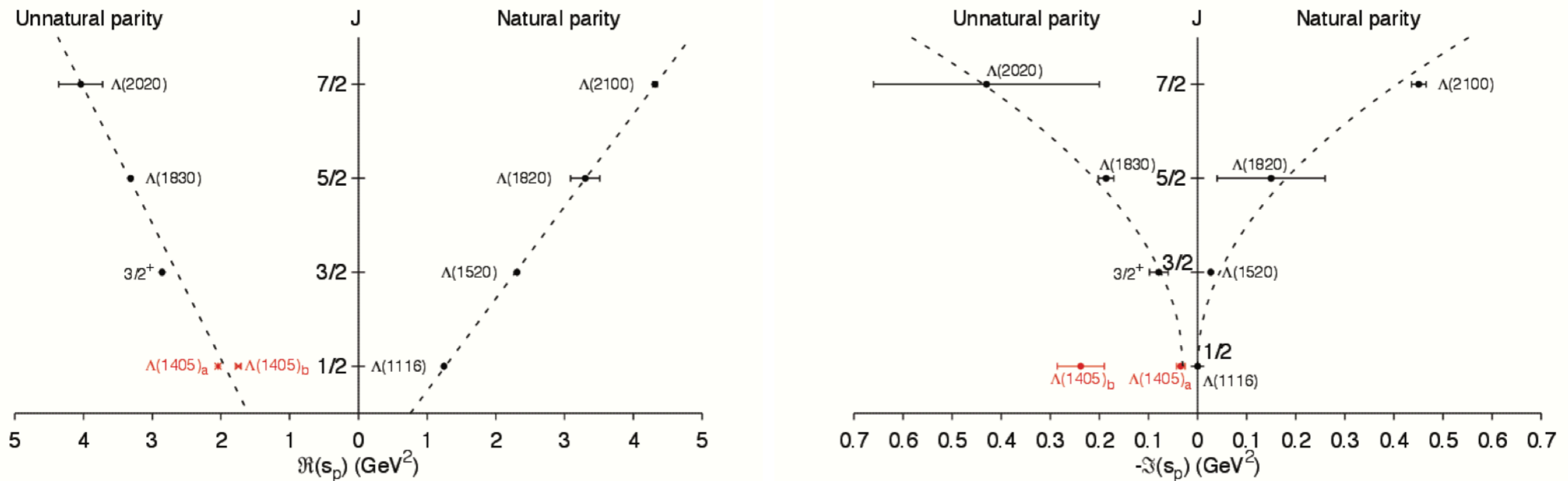
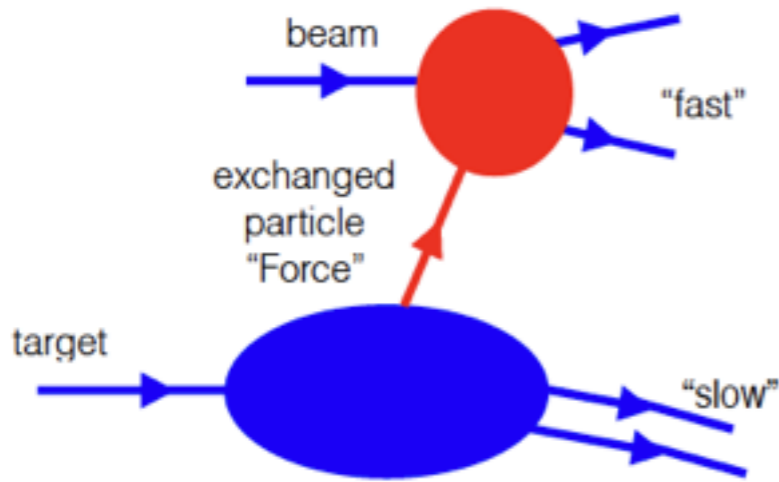


Figure 2: Leading Regge trajectories for the  $\Lambda$  resonances. Dashed lines are displayed to guide the eye.

C.Fernandez-Ramirez, et al. (JPAC) (2016)

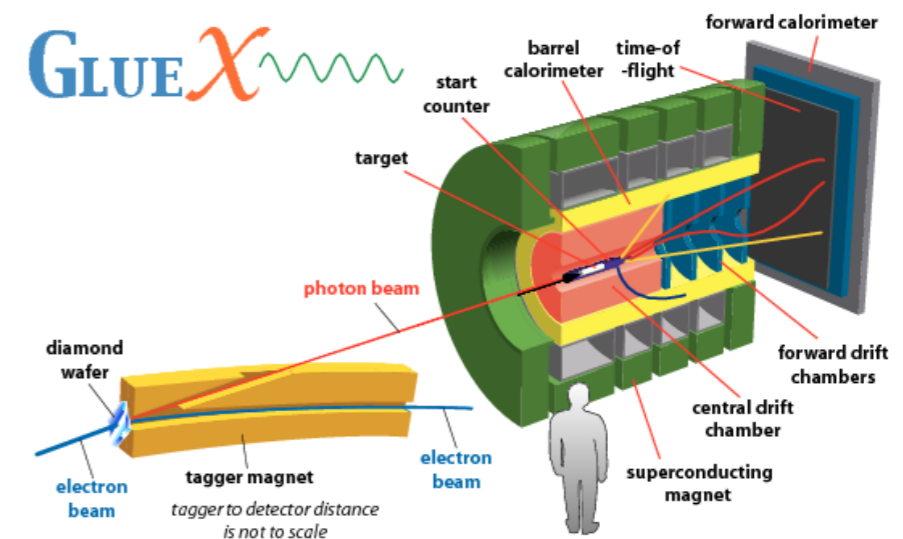


$$A(s, t) = \sum_l (2l + 1) A_l(s) P_l(t)$$

For large s,  $A_l(s)$  are small  
(tails of resonances)

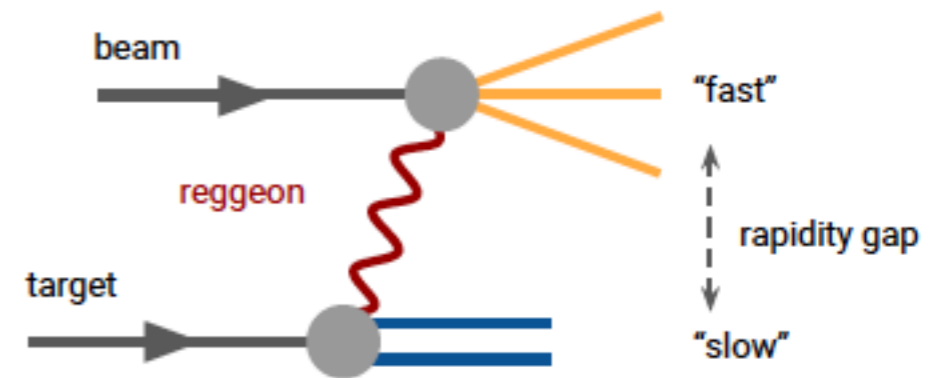
$$A(s, t) = \sum_l (2l + 1) A_l(t) P_l(s)$$

The t-channel p.w. with  
largest  $l_{eff} = \alpha(t)$   
dominates



# Global Regge analysis

- Test Regge pole hypothesis and estimate corrections (daughters, cuts)
- Regge poles : Factorizable residues  
t-dependent phase  
Shrinkage of the forward peak



$$A_{\lambda_i} = \sqrt{-t}^{|\lambda_1 - \lambda_3|} \sqrt{-t}^{|\lambda_2 - \lambda_4|} \beta_{\lambda_1 \lambda_3}^e \beta_{\lambda_2 \lambda_4}^a A_R(s, t)$$

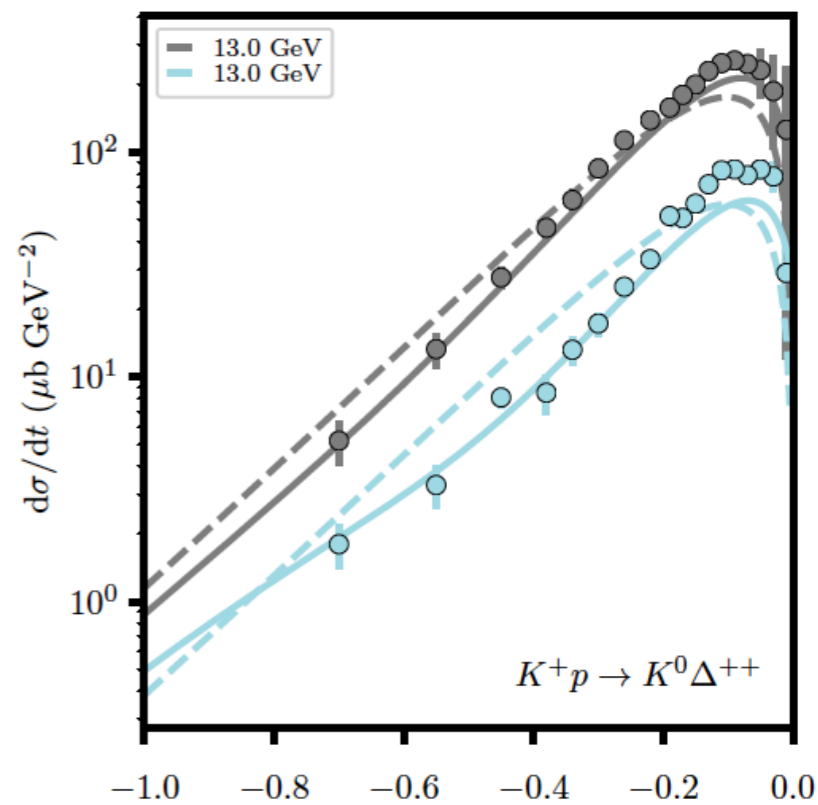
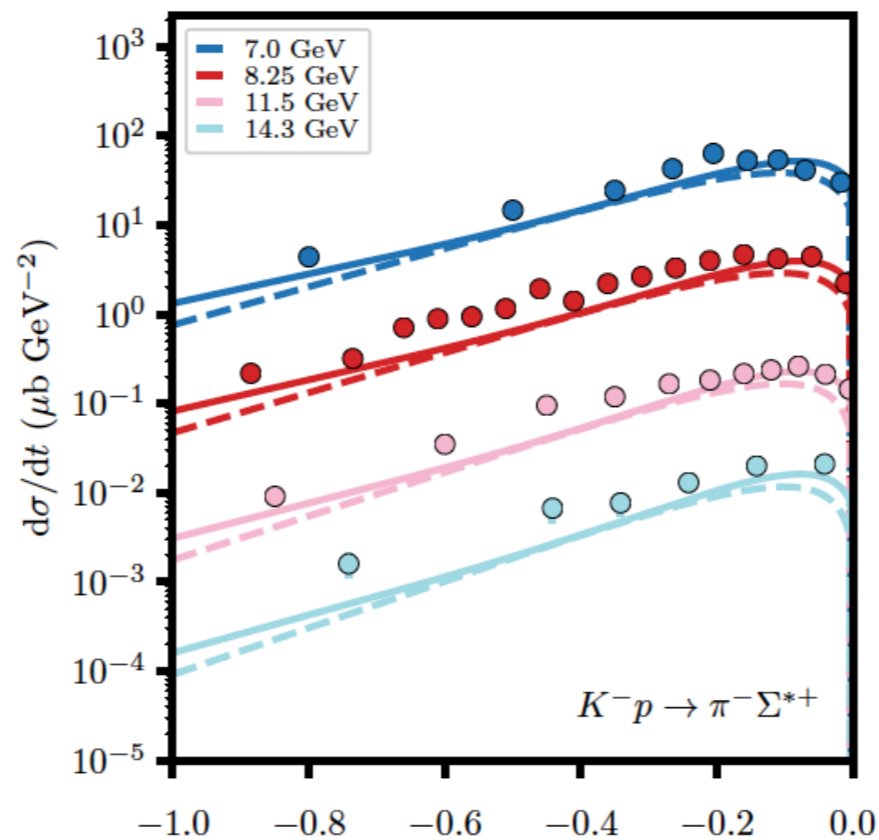
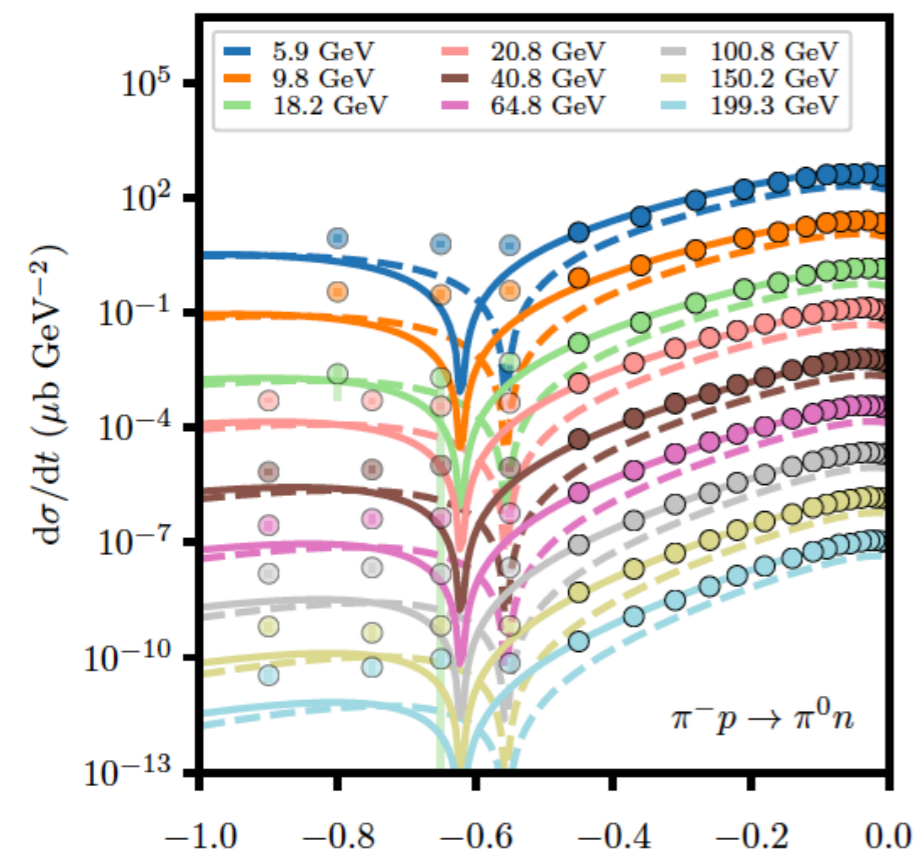
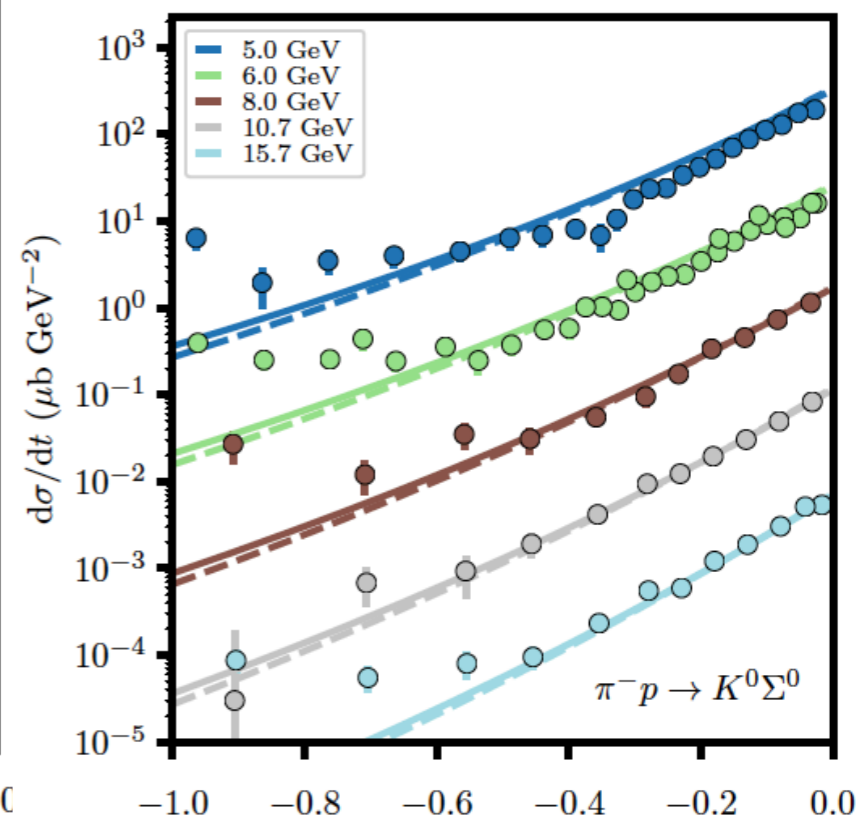
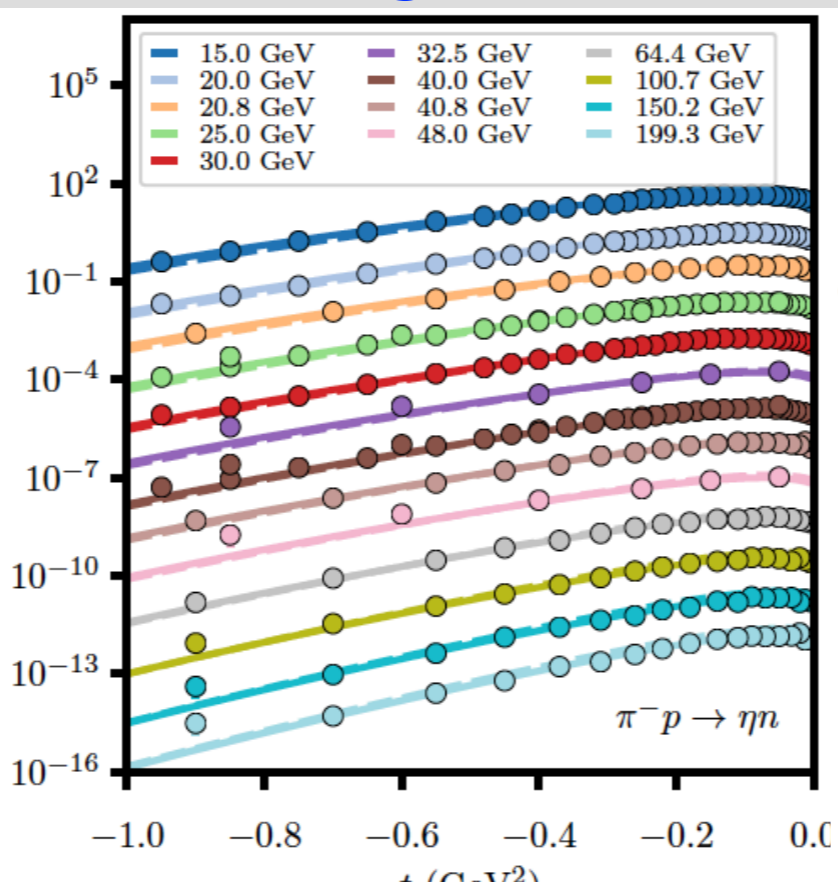
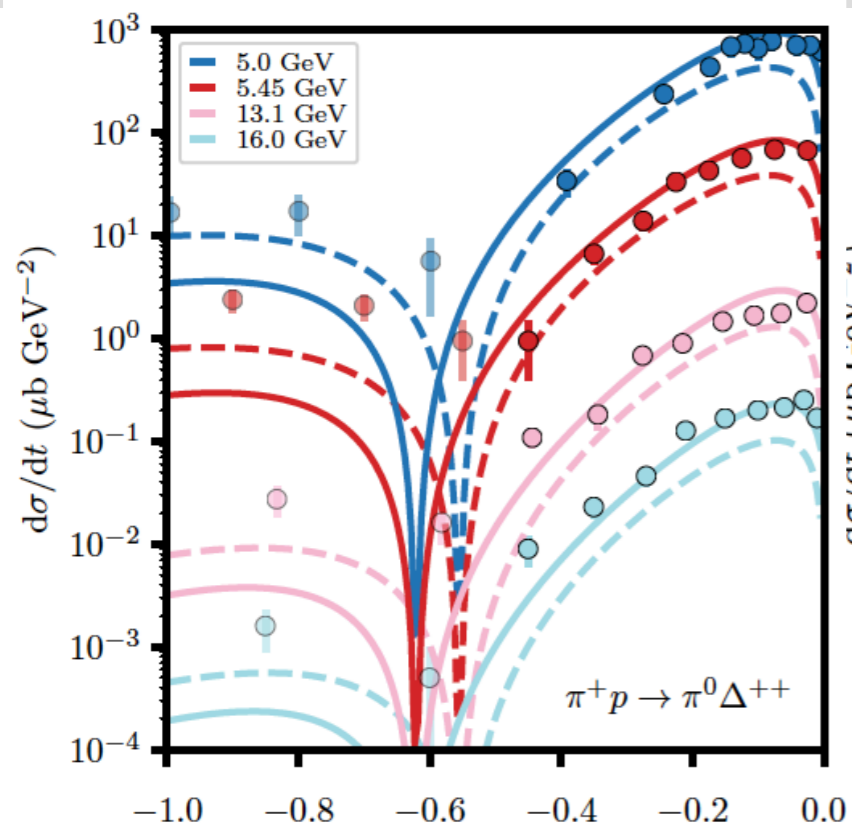
$\lambda_i = t$ -channel helicities

Data = 1271 points,  $N_{\text{par}} = 6$  SU(3) couplings, 1 mixing angle, 2 exp. slopes )

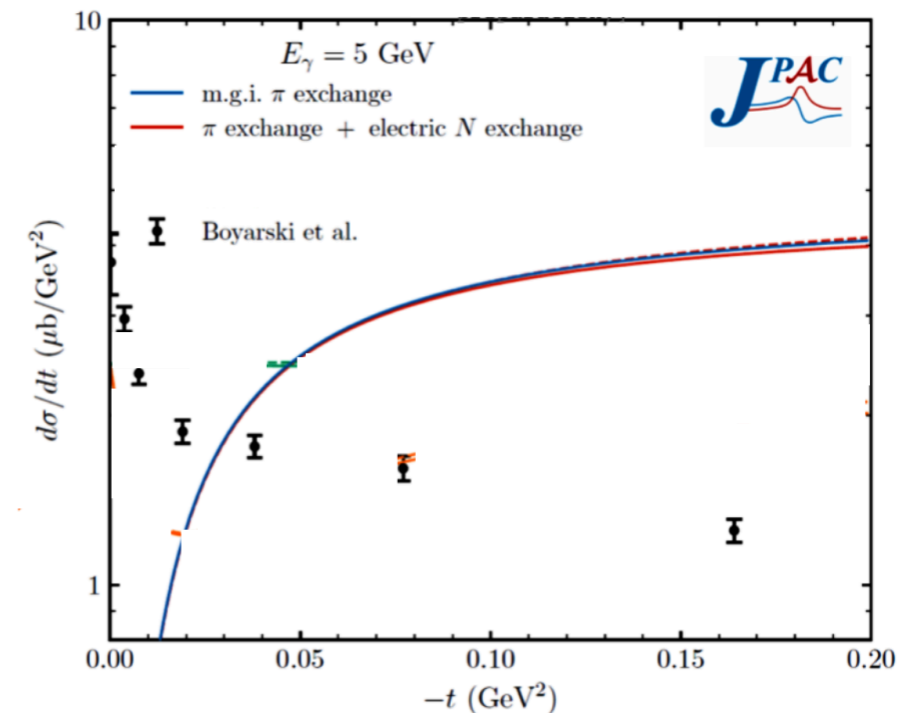
Except pion, pomeron exchange.



# Global Regge pole analysis



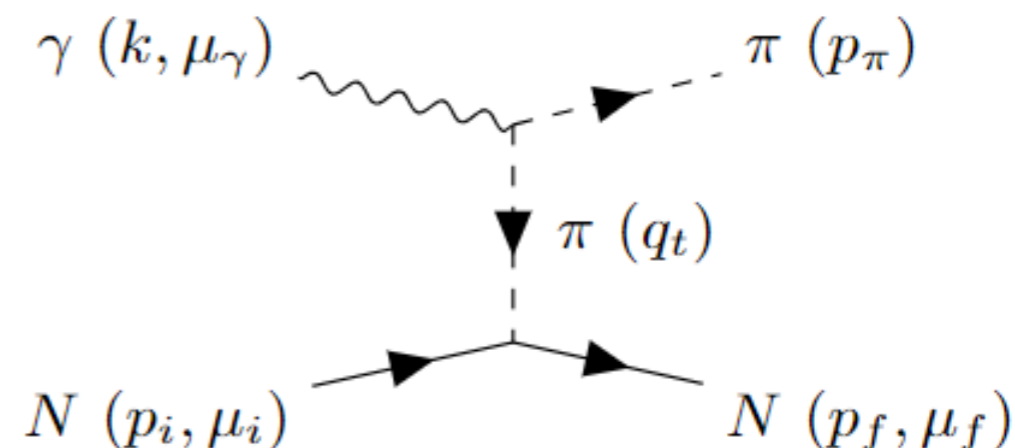
- At low- $t$  charge exchanges reactions should be dominated by  $\pi(K)$  exchanges
- Constraints from QCD, chiral, SU(3) symmetry.
- Relevant for light and heavy flavor pheno. (e.g. XYZ's)
- But there are still open issues
  - What is pion exchange ?
  - What happens at  $t \sim 0$



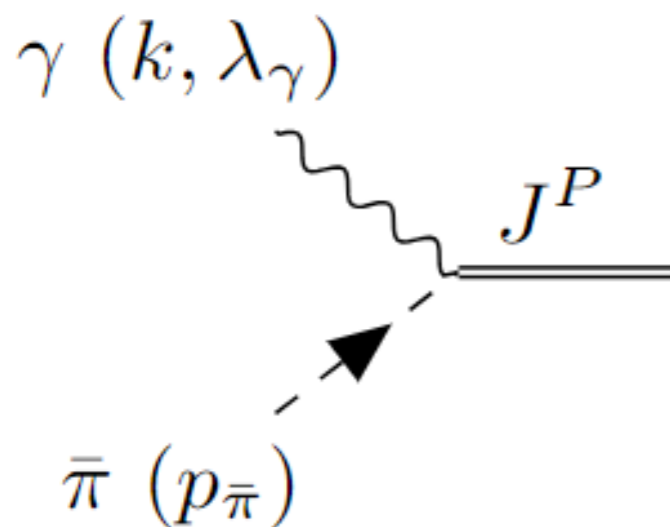
G.Montana et al. (JPAC) in preparation

$$A_{\lambda_i}(s, t) = \sum_J (2J + 1) A_{\lambda_i}^J(t) d_{\lambda_2 - \lambda_4, \lambda_1}^J(\theta_t)$$

In  $\gamma p \rightarrow \pi^+ n$  ( $K^+ Y$ ), t-channel p.w. with  $J^P = 0^-$  has the singularity closest to the physical region, but ...

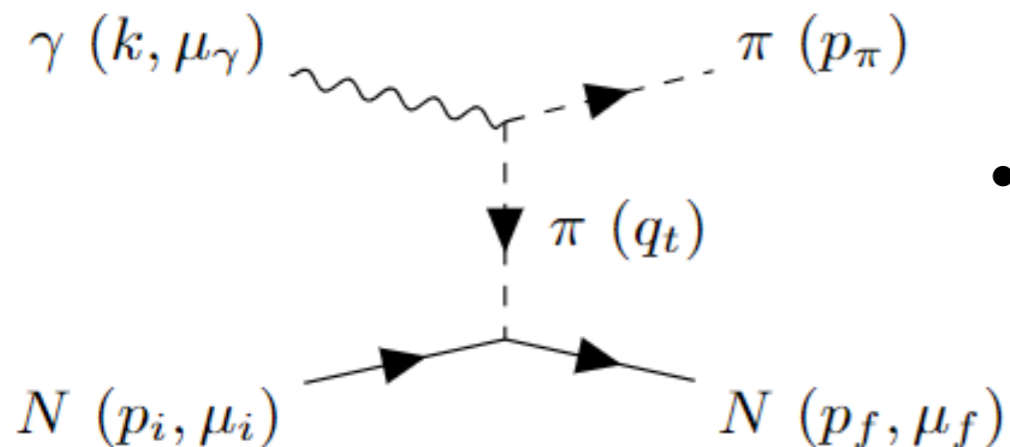


$\pi(K)$  contribution to t-channel vanishes !



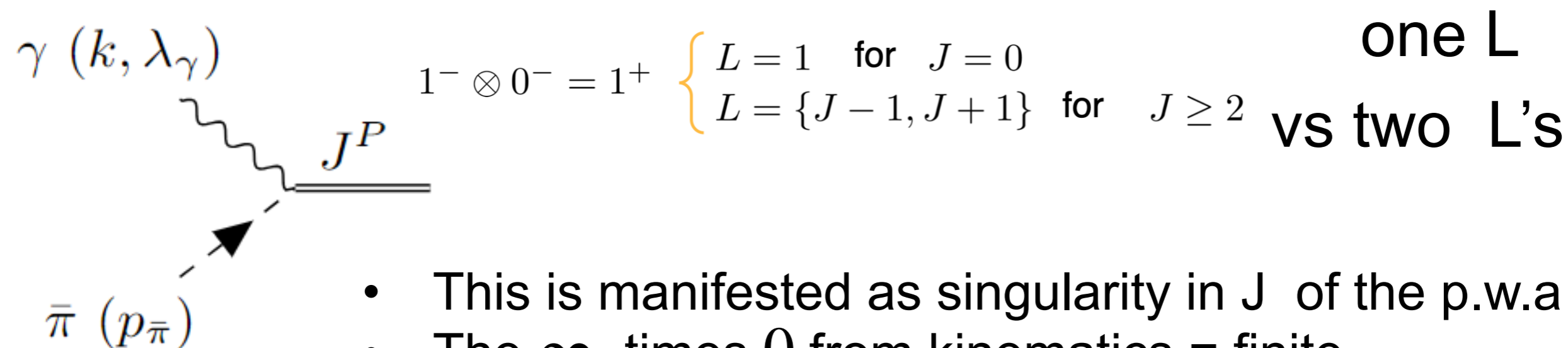
$$s_z^\gamma = \pm 1 \quad s_\pi = 0$$

$M_\pi$  would have to be  $\pm 1$

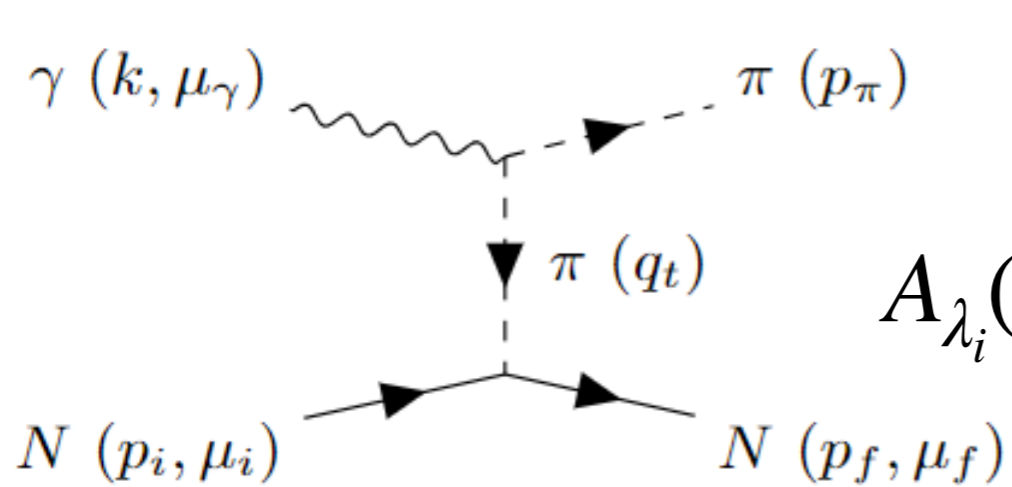


- The s-channel amplitude finite  $O(s^{J=0})$  : **Where does it come from ?**

- s/t-channel amplitudes are related (Wigner rotation) this means there has to be a Regge pole near  $J \sim 0$  in the t-channel p.w.



- $1^- \otimes 0^- = 1^+$ 
  - $L = 1$  for  $J = 0$  **one L**
  - $L = \{J - 1, J + 1\}$  for  $J \geq 2$  **vs two L's**
- This is manifested as singularity in  $J$  of the p.w.a
- The  $\infty$  times  $0$  from kinematics = finite



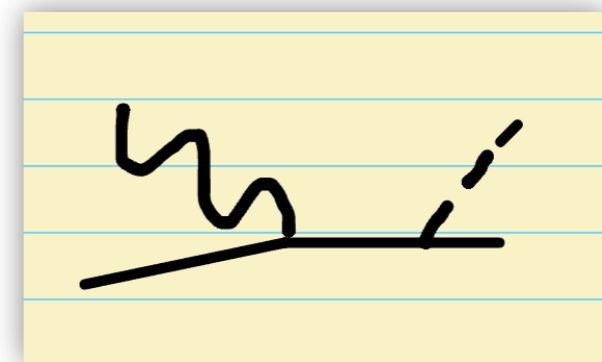
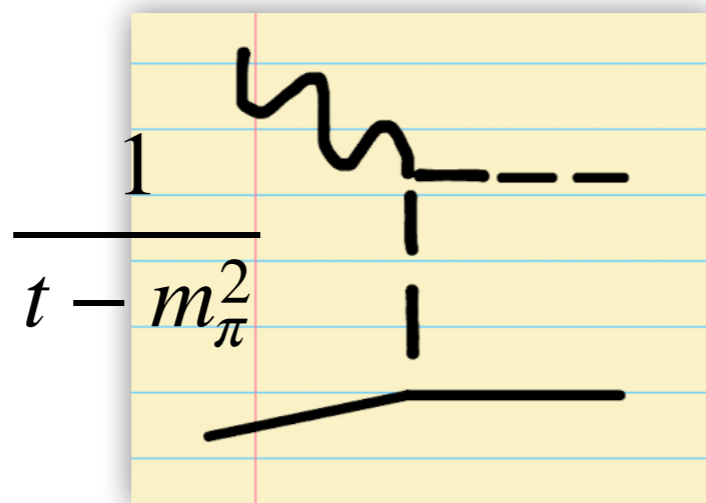
$$s + u = 2M^2 + m_\pi^2 - t$$

$$A_{\lambda_i}(s, t) = \sum_J (2J + 1) A_{\lambda_i}^J(t) d_{\lambda_2 - \lambda_4, \lambda_1}^J(\theta_t)$$

$$A_{\lambda_i}^J(t) \sim \frac{1}{J} \frac{1}{J - \alpha_\pi(t)}$$

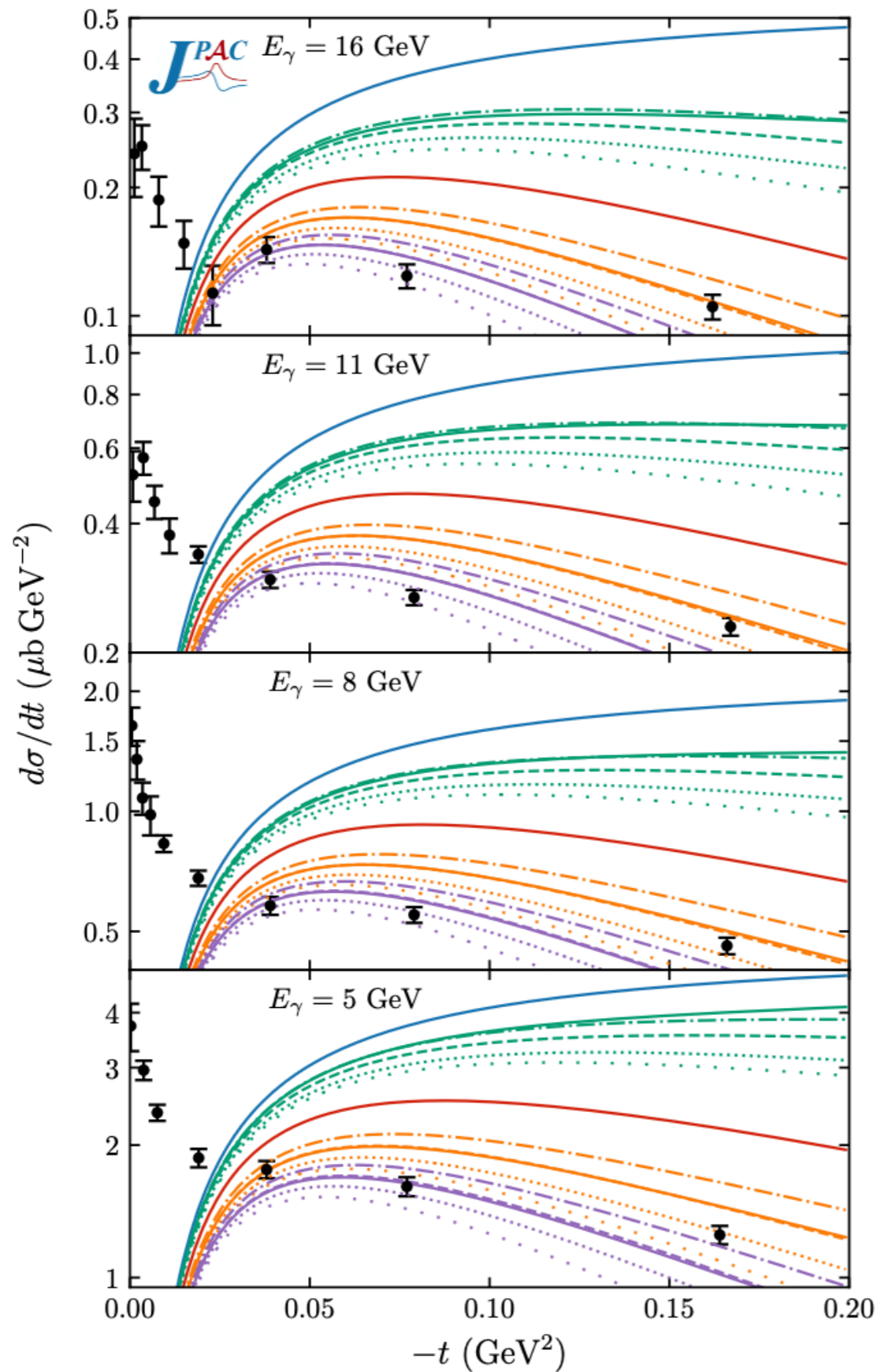
$$d_{\lambda_2 - \lambda_4, \lambda_1}^J(\theta_t) \rightarrow J \frac{t - m_\pi^2}{s - u}$$

t-channel = A x d = nucleon pole (s-channel) !



$$\frac{1}{s - M^2}$$





- Elementary exchange
- Regge  $\pi$  (VGL)
- - - Regge  $\pi + N$  (VGL)
- $\sum_J^{s \rightarrow \infty}, R = \sqrt{1/2s_0}$
- - -  $\sum_J, R = \sqrt{1/2s_0}, j_p = 1, j_z = 1$
- ⋯  $\sum_J, R = \sqrt{1/2s_0}, j_p = 2, j_z = 1$
- · -  $\sum_J, R = \sqrt{1/2s_0}, j_p = 1, j_z = 2$
- ⋯  $\sum_J, R = \sqrt{1/2s_0}, j_p = 1, j_z = 0.5$
- $\sum_J^{s \rightarrow \infty}, R = 1 \text{ fm}$
- - -  $\sum_J, R = 1 \text{ fm}, j_p = 1, j_z = 1$
- ⋯  $\sum_J, R = 1 \text{ fm}, j_p = 2, j_z = 1$
- · -  $\sum_J, R = 1 \text{ fm}, j_p = 1, j_z = 2$
- ⋯  $\sum_J, R = 1 \text{ fm}, j_p = 1, j_z = 0.5$
- $\sum_J^{s \rightarrow \infty}, R = 2 \text{ fm}$
- - -  $\sum_J, R = 2 \text{ fm}, j_p = 1, j_z = 1$
- ⋯  $\sum_J, R = 2 \text{ fm}, j_p = 2, j_z = 1$
- · -  $\sum_J, R = 2 \text{ fm}, j_p = 1, j_z = 2$
- ⋯  $\sum_J, R = 2 \text{ fm}, j_p = 1, j_z = 0.5$

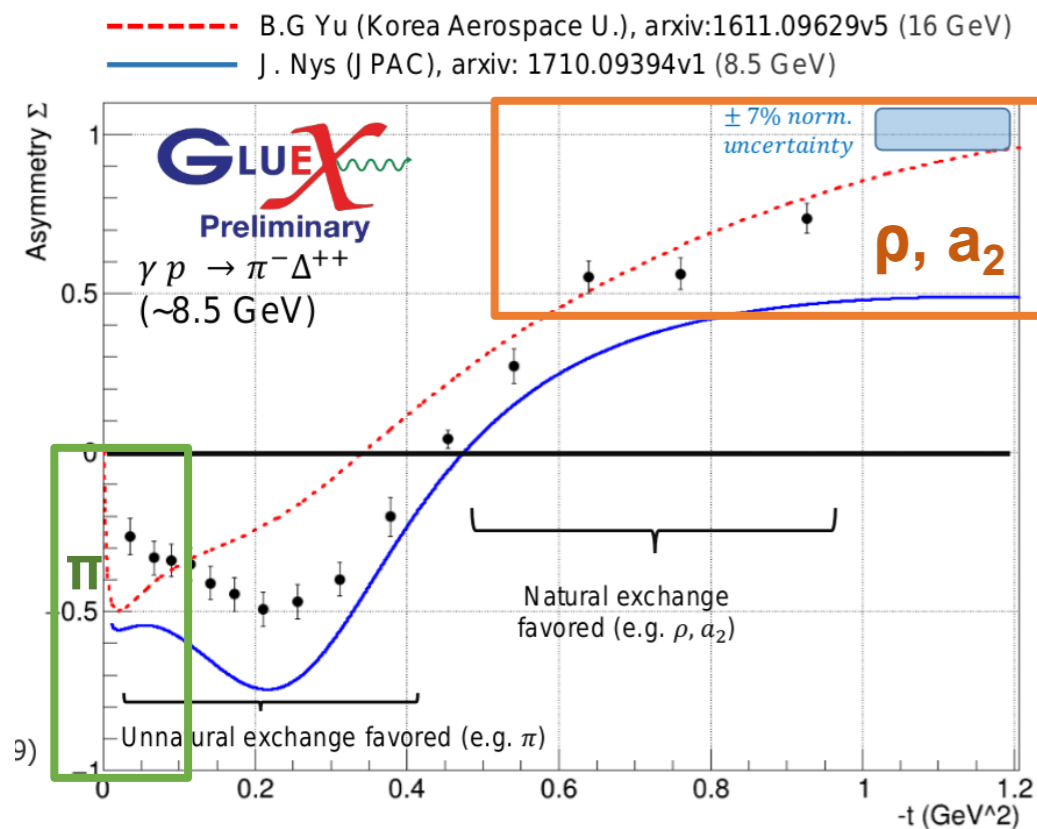
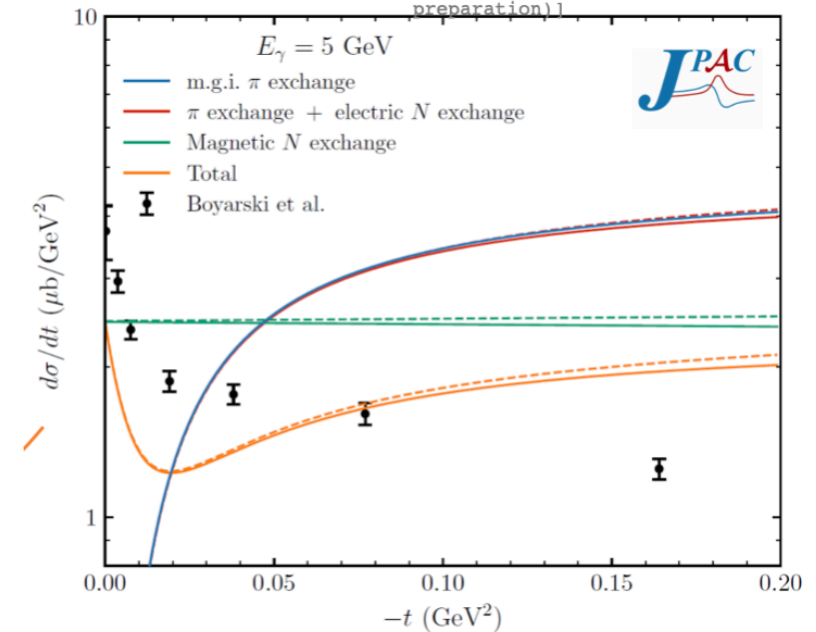


# Pion/Kaon exchange

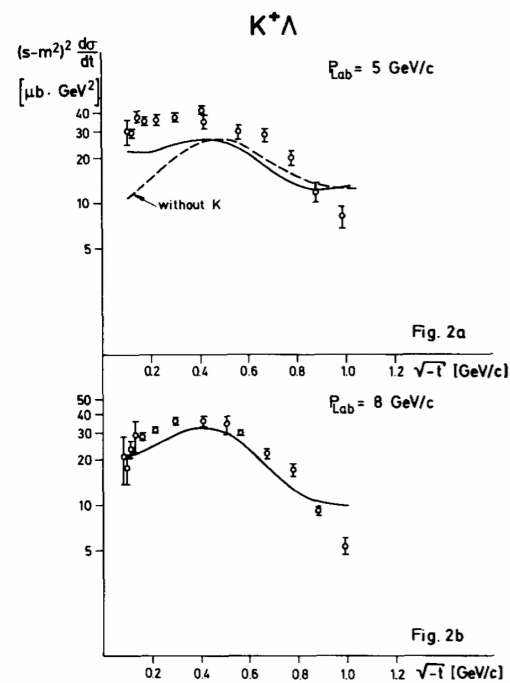
- The J=0 pole is equivalent to nucleon contribution (current conservation)
- As it should be : there is no need to mix s- and t-channel amplitudes

$$\frac{-t}{\mu^2 - t} = 1 + \frac{-\mu^2}{\mu^2 - t}$$

[G.Montana et al. (in preparation)]



- Photo-production is the “cleanest” probe of OPE.
- At higher-t natural exchanges dominate
- At t~0 other interesting phenomena: absorption, cuts, conspiracy between pion and nucleon poles etc. (stay tuned)



A.M.Boyarski et al. SLAC (1969) (exp)  
 N.Levy et al. (1973)

- Low-t vs. high-t physics : K vs K\* exchanges
- SU(3) relations
- Factorization breaking effects : need low-t,  $|-t'| < m_K^2 \sim 0.25 \text{ GeV}^2$

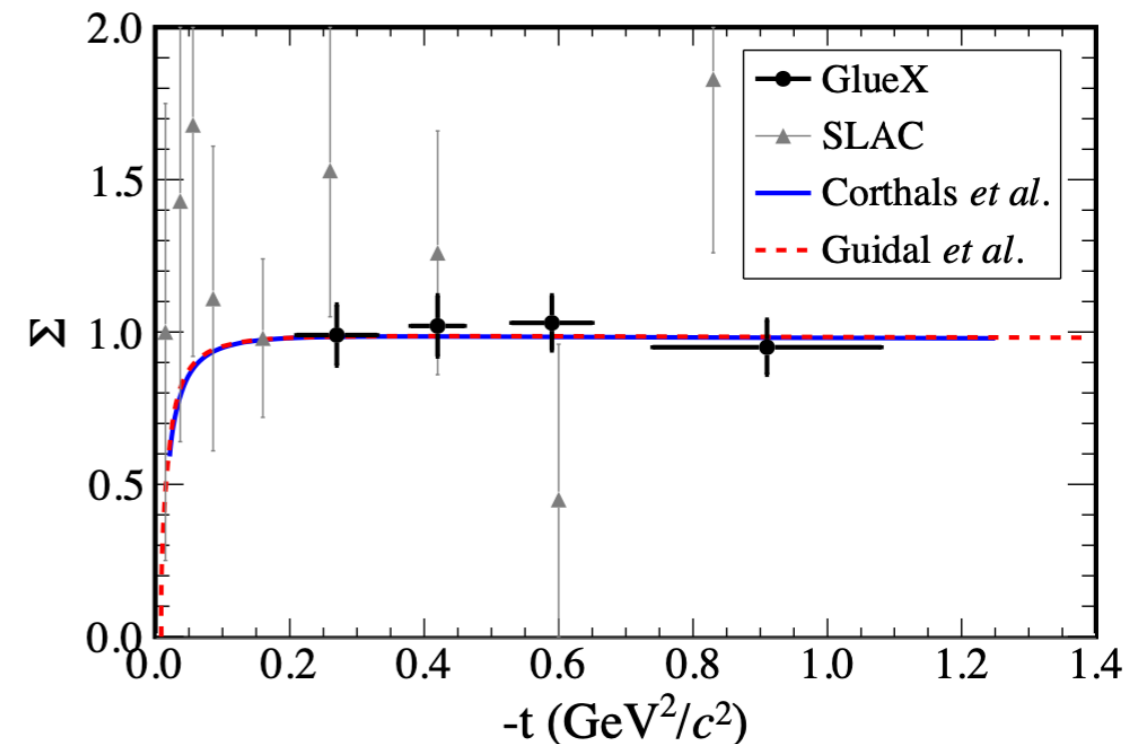
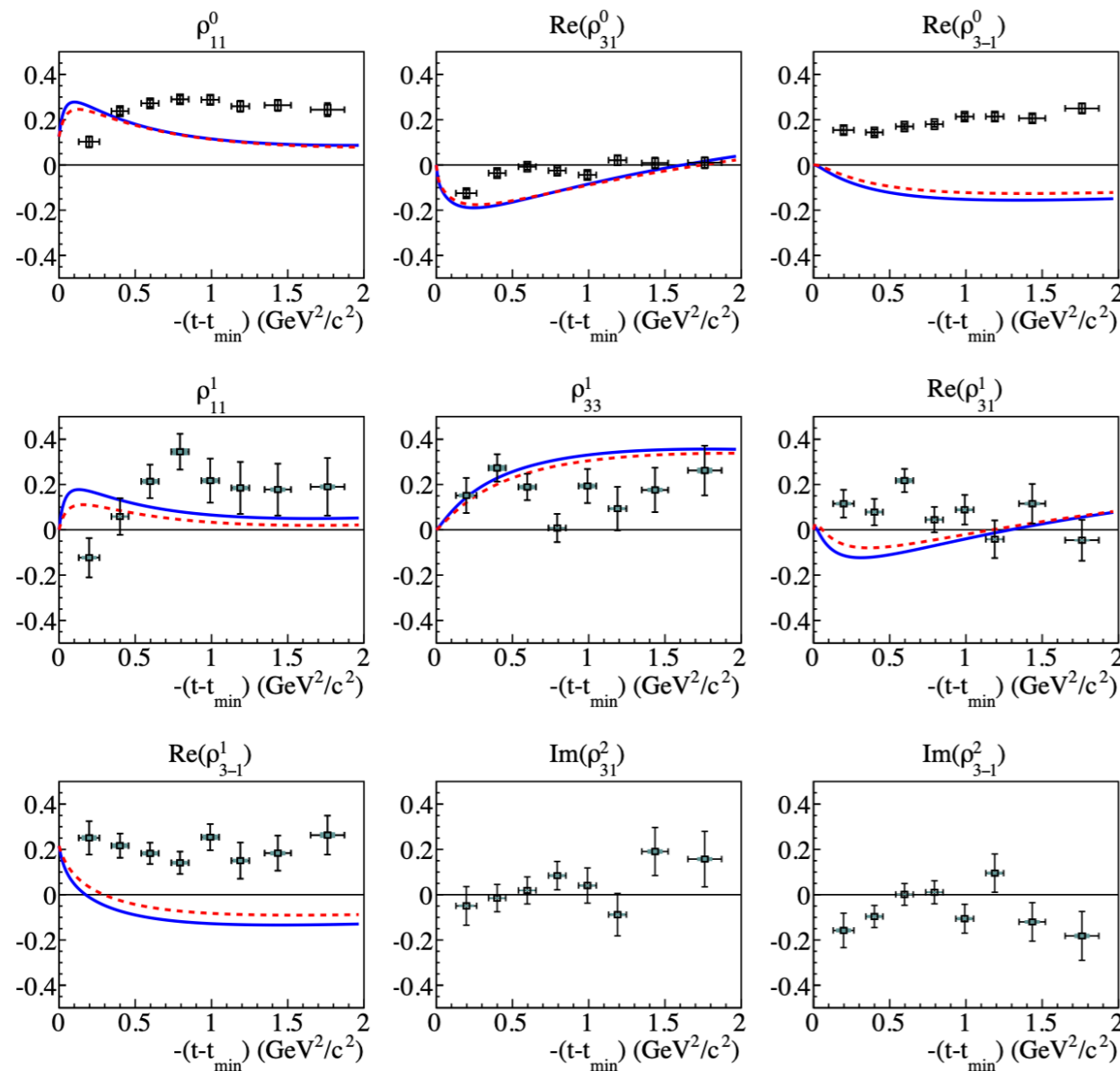


FIG. 10. The beam asymmetry  $\Sigma$  for  $\vec{\gamma}p \rightarrow K^+\Sigma^0$  as a function of  $-t$ . The results from the 0/90 and -45/45 data sets are averaged (solid circles) where horizontal error bars indicate the RMS widths of the  $t$  bins and vertical error bars represent statistical and systematic uncertainties added in quadrature. An additional 2.1% overall relative polarization uncertainty is not included. The triangles are previous SLAC results [2] at  $E_\gamma = 16 \text{ GeV}$ , the curves show predictions from RPR-2007 [6, 7] (solid) and Guidal *et al.* [5] (dashed) at  $E_\gamma = 8.5 \text{ GeV}$ .

J.Hernandez talk





- SDME's can be used to test Regge pole dominance, e.g.  $\rho^2$  vanishes if residues factorize

- Models need to be scrutinized

FIG. 7. Spin density matrix elements and predictions by Yu and Kong (based on Ref. [14]), using parameters based on data from CLAS [9] and LEPS [6, 7] (blue solid) and using parameters based on data from LAMP2 [5] and SLAC [4] (red dashed). The vertical error bars show the statistical uncertainty, the blue shaded boxes the scaling uncertainty from the polarization, and the black boxes the remaining systematic uncertainties combined in quadrature. The horizontal error bars show the RMS widths within the  $-(t - t_0)$  bins.

S.Adhikari, et al. GlueX 2022

# Summary

- (Quasi) elastic production:  
determines Regge exchanges
- Use Regge theory to correlate  
resonance production with  
resonance decays, e.g. FESR's
- One pole, two pole, etc. structures  
should be understood in terms of  
trajectories they belong to





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Arkaitz Rodas Bilbao  
Old Dominion University /  
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Astrid Hiller Blin  
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Mikhail Mikhasenko  
Ruhr-Universität Bochum



Robert Perry  
University of Barcelona



Sergi González-Solís  
University of Barcelona



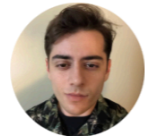
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Jefferson Lab



Vincent Mathieu  
University of Barcelona



Wyatt Smith  
George Washington  
University



Jinfeng Liao  
Indiana University



Andrew Jackura  
William & Mary



Derek Glazier  
University of Glasgow



Geoffrey Fox  
University of Virginia



Igor Danilkin  
JG University Mainz



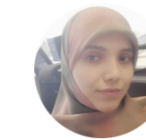
Jorge A. Silva-Castro  
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Kevin Quirion  
Indiana University



Michael Döring  
George Washington  
University



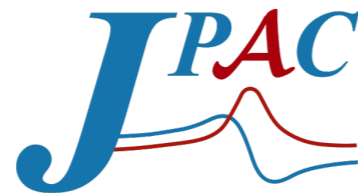
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Stephen Sharpe  
University of Washington



Eric Swanson  
University of Pittsburgh



Adam Szczepaniak  
Indiana University



The Exo(tic) Had(ron) Collaboration started in 2023 to explore all aspects of exotic hadron physics, from predictions within lattice QCD, through reliable extraction of their existence and properties from experimental data, to descriptions of their structure within phenomenological models.

