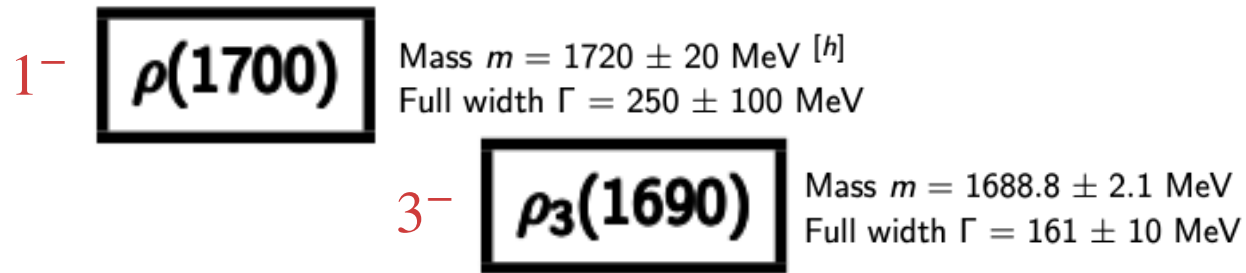


J^{--} meson resonances in QCD

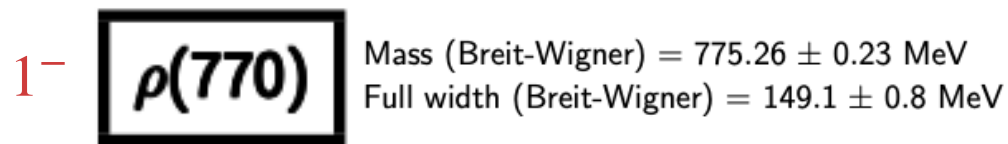
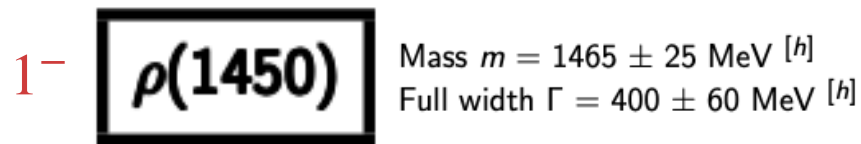
Jozef Dudek

... with focus on those containing strange quarks ...

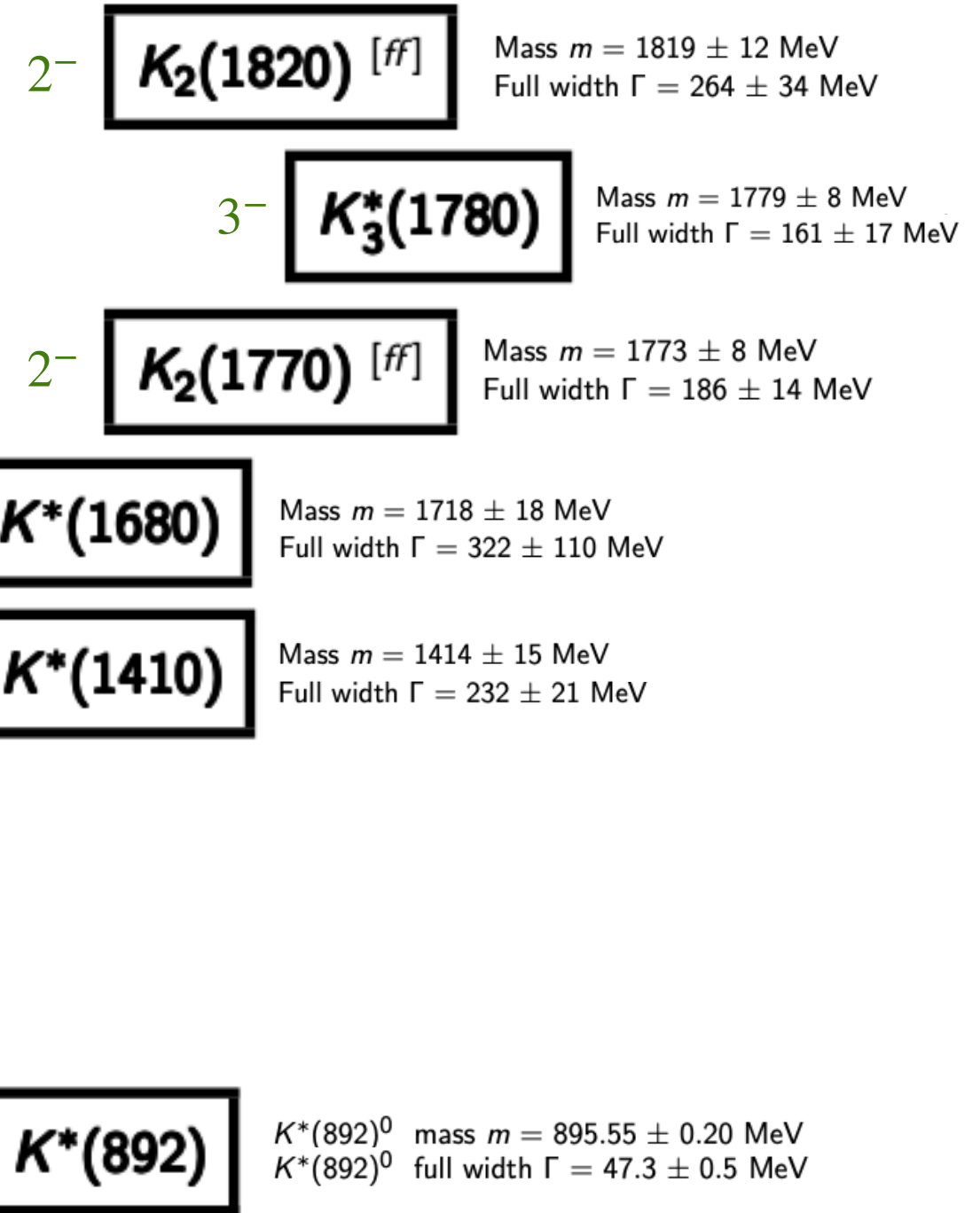
isovectors



no 2^- candidates



strange mesons



isoscalars

3^- **$\phi_3(1850)$** Mass $m = 1854 \pm 7$ MeV
Full width $\Gamma = 87^{+28}_{-23}$ MeV

3^- **$\omega_3(1670)$** Mass $m = 1667 \pm 4$ MeV
Full width $\Gamma = 168 \pm 10$ MeV

1^- **$\phi(1680)$** Mass $m = 1680 \pm 20$ MeV ^[h]
Full width $\Gamma = 150 \pm 50$ MeV ^[h]

no 2^- candidates

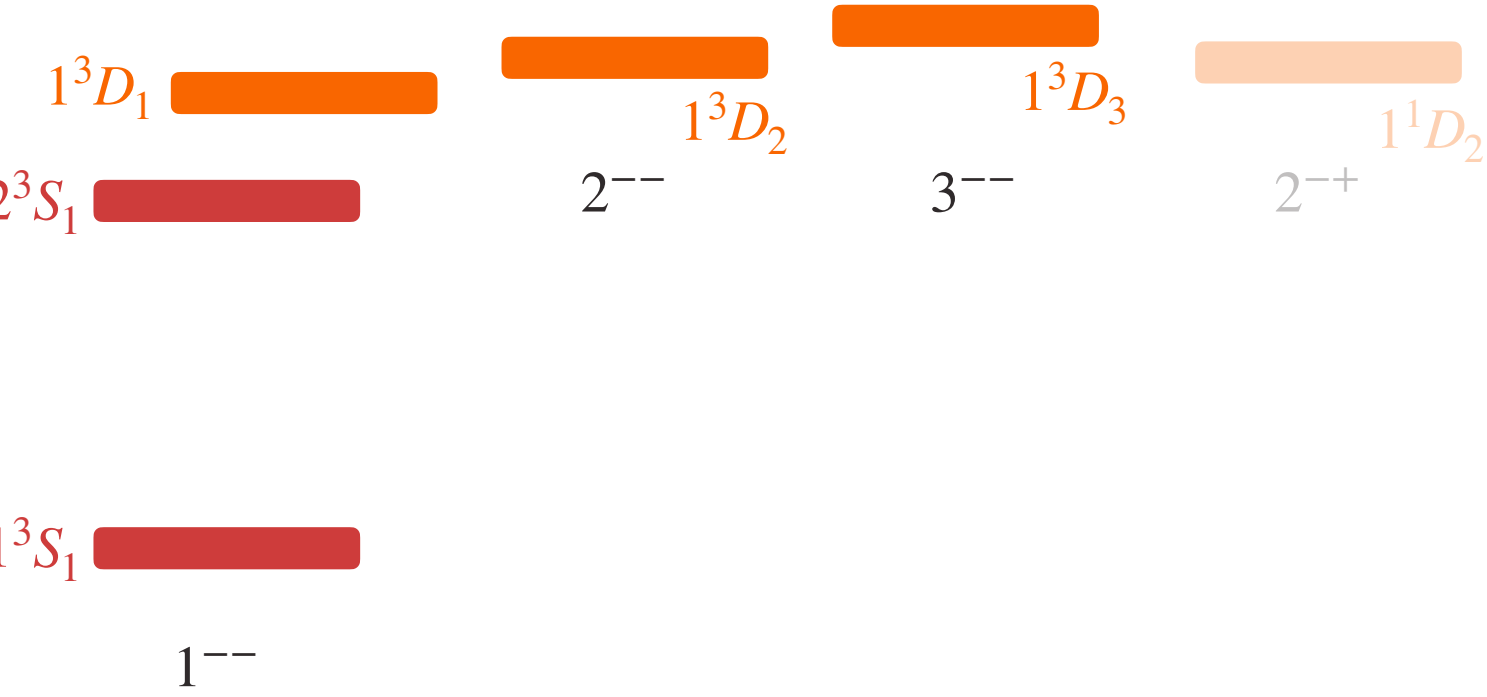
1^- **$\omega(1650)$ [k]** Mass $m = 1670 \pm 30$ MeV ^[h]
Full width $\Gamma = 315 \pm 35$ MeV ^[h]

1^- **$\omega(1420)$ [j]** Mass $m = 1410 \pm 60$ MeV ^[h]
Full width $\Gamma = 290 \pm 190$ MeV ^[h]

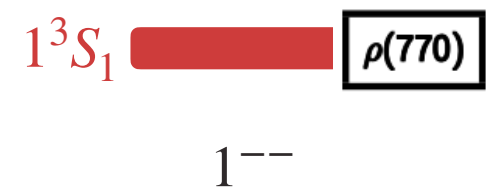
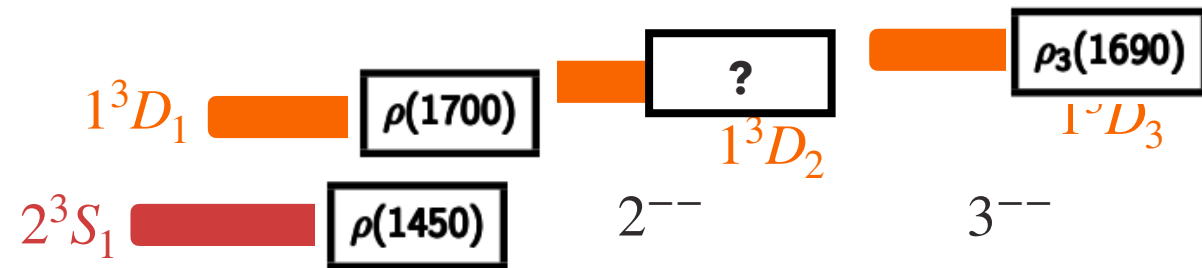
1^- **$\phi(1020)$** Mass $m = 1019.461 \pm 0.016$ MeV
Full width $\Gamma = 4.249 \pm 0.013$ MeV

1^- **$\omega(782)$** Mass $m = 782.66 \pm 0.13$ MeV
Full width $\Gamma = 8.68 \pm 0.13$ MeV

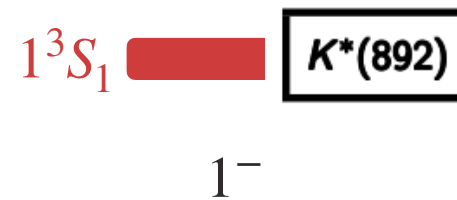
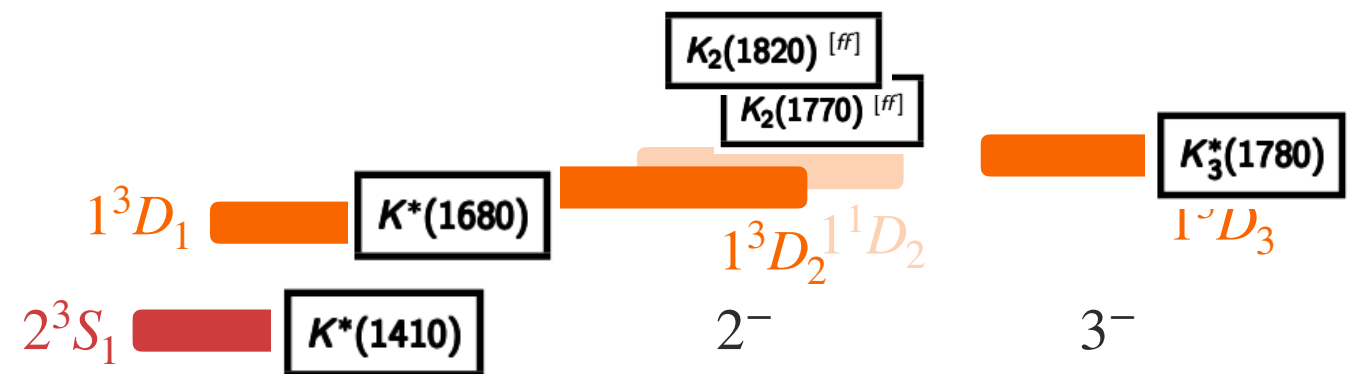
$q\bar{q}$ quark model picture is simple



isovectors



strange mesons



1^- states could be $2^3S_1, 1^3D_1$ admixtures in this model

ω, ϕ names typically assigned assuming **OZI rule** holds

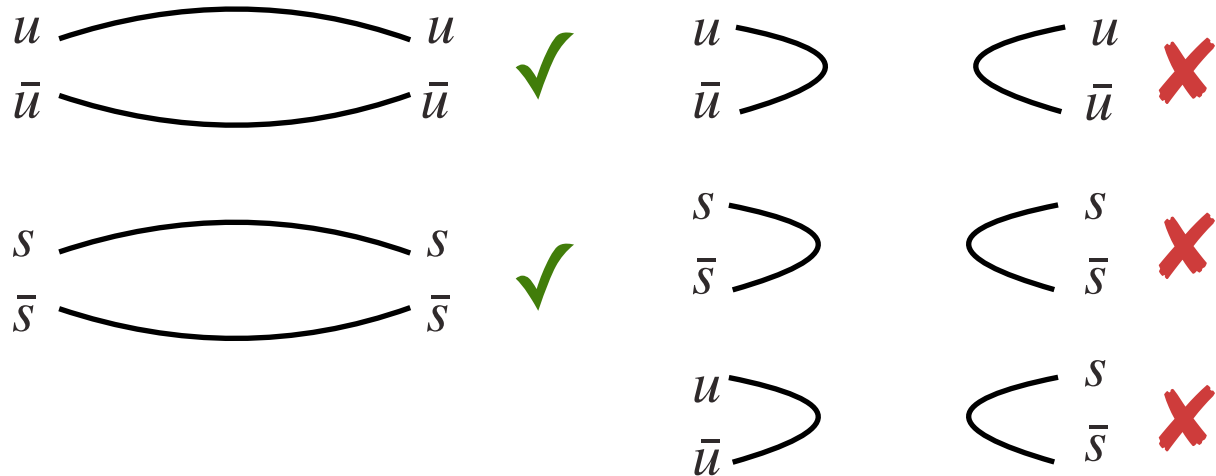
ideal flavor mixing

$$\omega \sim u\bar{u} + d\bar{d}$$

$$\phi \sim s\bar{s}$$

OZI rule (generalized)

"connected diagrams favored over disconnected"



ω, ϕ names typically assigned assuming **OZI rule** holds

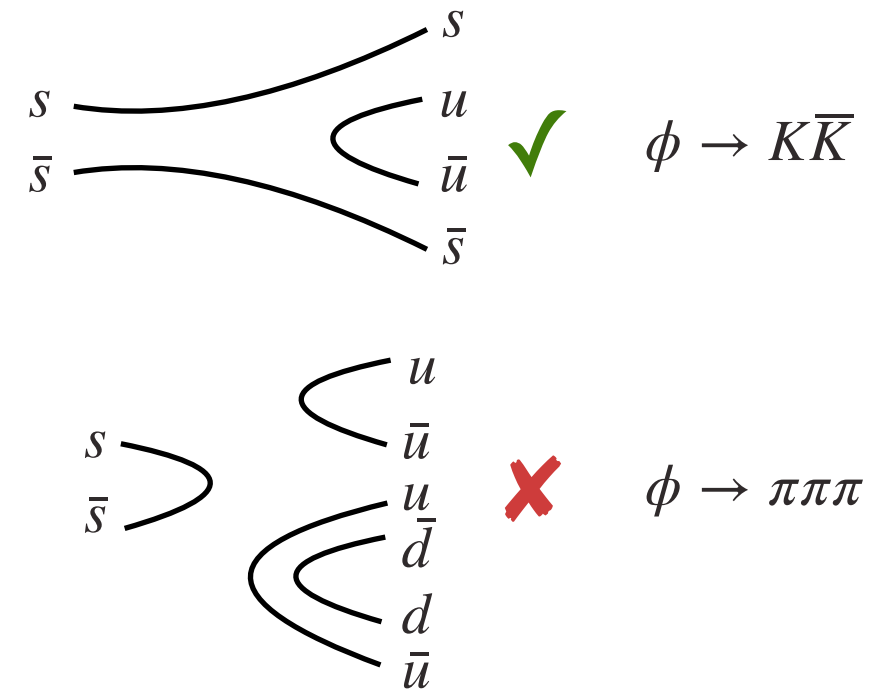
ideal flavor mixing

$$\omega \sim u\bar{u} + d\bar{d}$$

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OZI rule (generalized)

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ω, ϕ names typically assigned assuming **OZI rule** holds

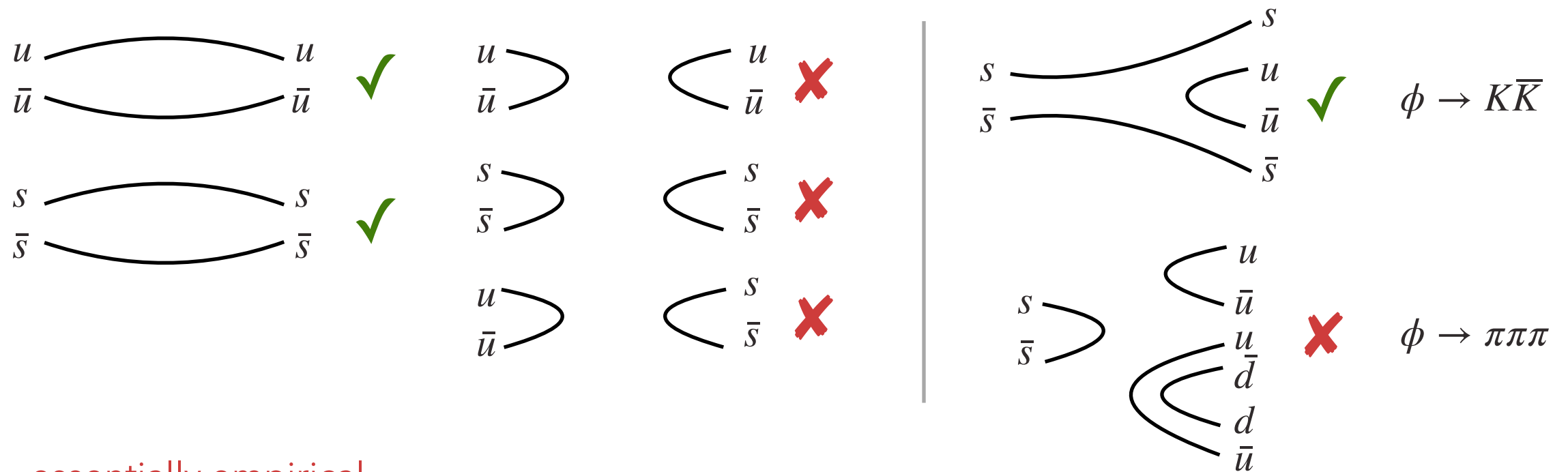
ideal flavor mixing

$$\omega \sim u\bar{u} + d\bar{d}$$

$$\phi \sim s\bar{s}$$

OZI rule (generalized)

"connected diagrams favored over disconnected"



essentially empirical

famous strong evidence: $\Gamma(\phi(1020) \rightarrow K\bar{K}) \gg \Gamma(\phi(1020) \rightarrow \pi\pi\pi)$

$$m_{\rho(770)} \approx m_{\omega(782)}$$

but doesn't work **at all** for pseudoscalars

$$\eta \sim u\bar{u} + d\bar{d} - 2s\bar{s}$$

$$\eta' \sim u\bar{u} + d\bar{d} + s\bar{s}$$

maybe axials not ideally mixed either

ω, ϕ names typically assigned assuming OZI rule holds

ideal flavor mixing

$$\omega \sim u\bar{u} + d\bar{d}$$

$$\phi \sim s\bar{s}$$

OZI rule (generalized)

“connected diagrams favored over disconnected”



essentially empirical

famous strong evidence: $\Gamma(\phi(1020) \rightarrow K\bar{K}) \gg \Gamma(\phi(1020) \rightarrow \pi\pi\pi)$

$$m_{\rho(770)} \approx m_{\omega(782)}$$

we shouldn't assume 'strangeonium' is how QCD chooses to manifest isoscalar mesons in all J^P , and at all energies

but doesn't work at all for pseudoscalars

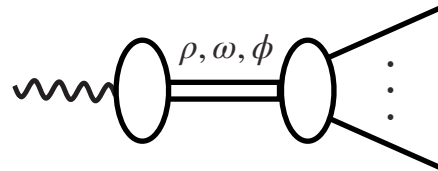
$$\eta \sim u\bar{u} + d\bar{d} - 2s\bar{s}$$

$$\eta' \sim u\bar{u} + d\bar{d} + s\bar{s}$$

it might, but we should establish it ...

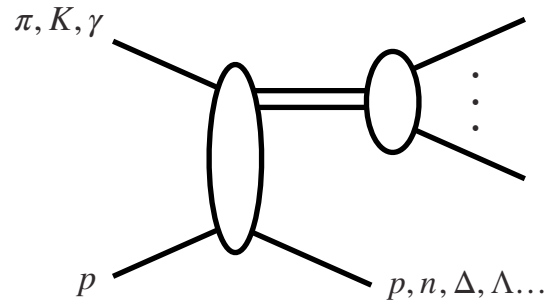
(this talk is not going to achieve this)

e^+e^- annihilation



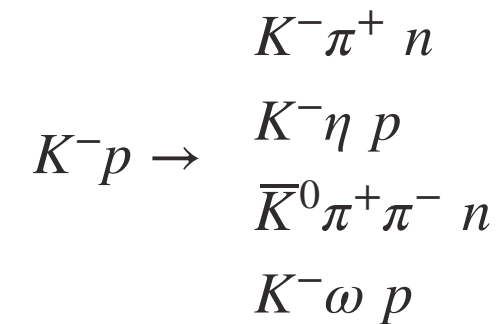
1^{--} only, absolute normalization (cross-section)

peripheral production

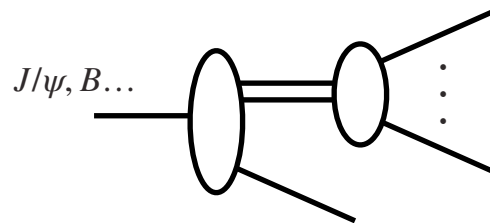


all partial-waves, have to be separated

e.g. **LASS** (1980s)



heavy flavor decays



Dalitz plot analysis

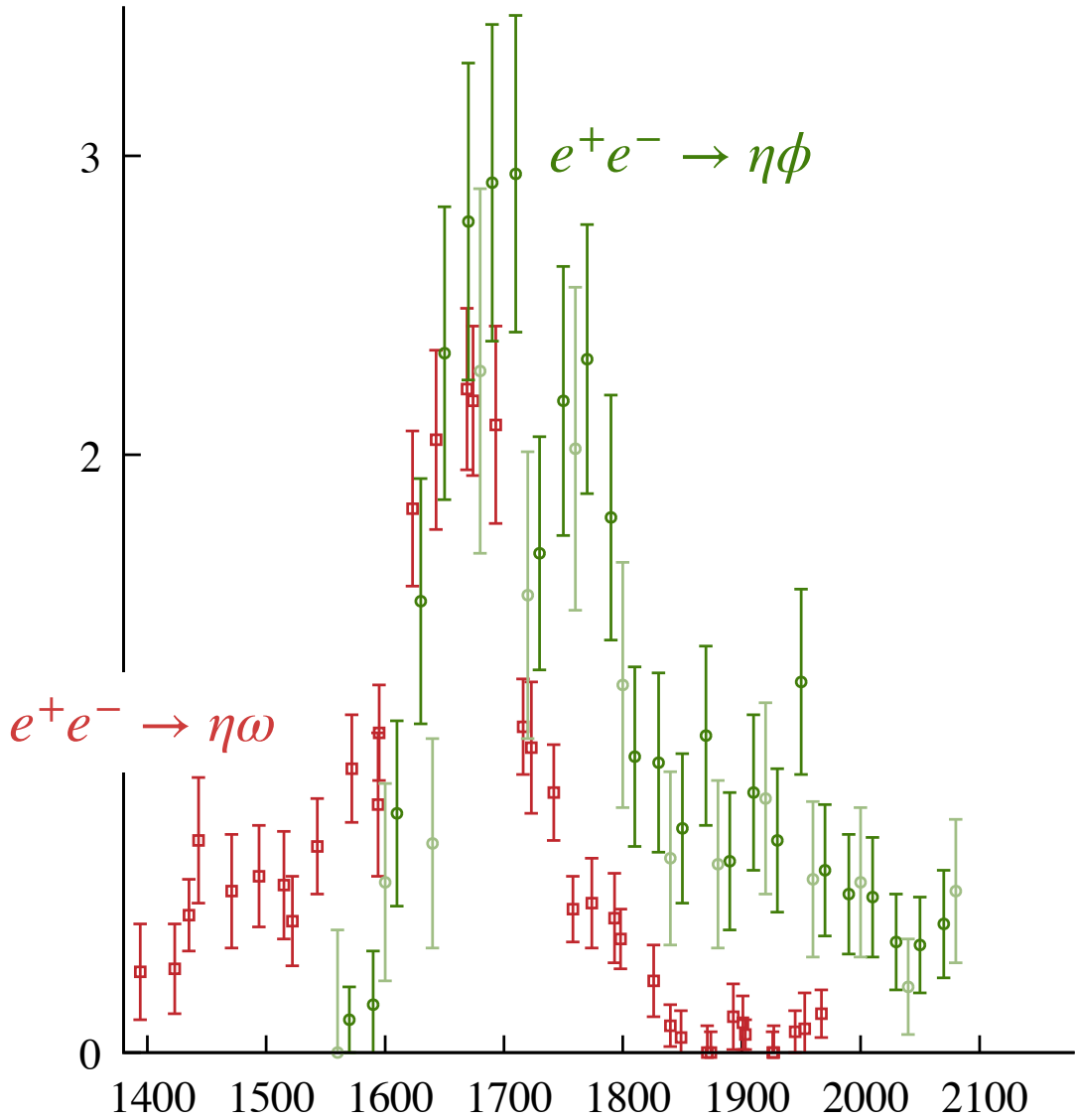
model multiple pairs simultaneously

e.g. **LHCb**



put in all PDG K^* resonances, assuming they are well established and consistent

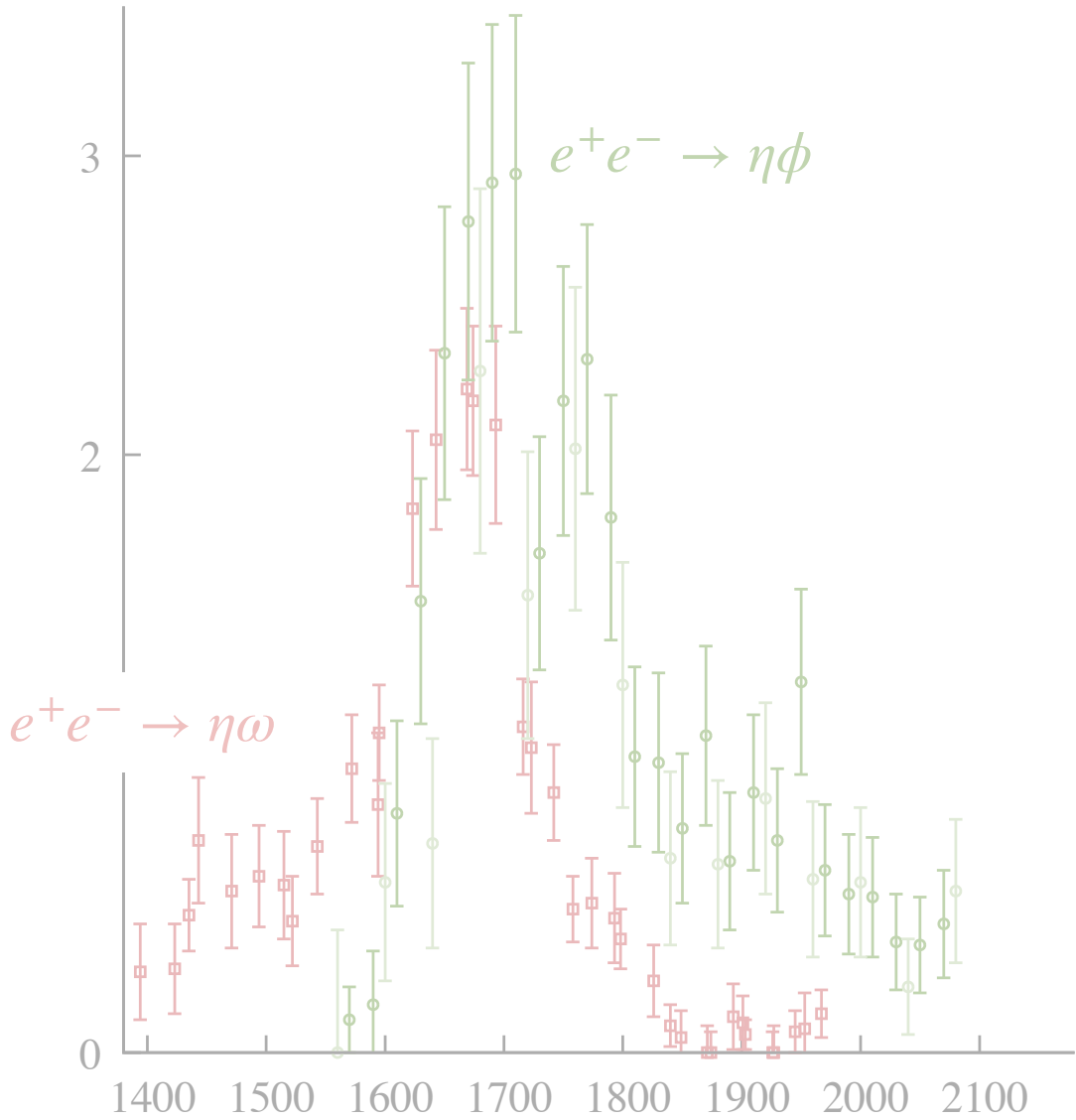
but even the 'simple' cases prove to be tricky:



published analyses of these data have **no resonances in common !**

in principle expect **four** resonances could contribute to each of these final states

but even the 'simple' cases prove to be tricky:



and it's not clear resonance identifications are really consistent, e.g.

$\omega(1420)$ WIDTH

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	COMMENT
290 ± 190 OUR ESTIMATE				
440 ± 125	267	1 ACHASOV	20B SND	$e^+e^- \rightarrow \omega\eta \rightarrow \eta\pi^0\gamma$
$104 \pm 35 \pm 10$	824	2 AKHMETSHIN	17A CMD3	$1.4-2.0 e^+e^- \rightarrow \omega\eta$
880 ± 170	13.1k	3 AULCHENKO	15A SND	$1.05-1.80 e^+e^- \rightarrow \pi^+\pi^-\pi^0$
480 ± 180		4 ACHASOV	10D SND	$1.075-2.0 e^+e^- \rightarrow \pi^0\gamma$
$130 \pm 50 \pm 100$		AUBERT	07AU BABR	$10.6 e^+e^- \rightarrow \omega\pi^+\pi^-\gamma$
$450 \pm 70 \pm 70$		AUBERT,B	04N BABR	$10.6 e^+e^- \rightarrow \pi^+\pi^-\pi^0\gamma$

$\omega(1650)$ WIDTH

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	COMMENT
315 ± 35 OUR ESTIMATE				
110 ± 16	267	1 ACHASOV	20B SND	$e^+e^- \rightarrow \omega\eta \rightarrow \eta\pi^0\gamma$
$194 \pm 8^+_{-7} 15$	183k	2 ABLIKIM	19AQ BES	$J/\psi \rightarrow K^+K^-\pi^0$
95 ± 11		ACHASOV	19 SND	$e^+e^- \rightarrow \pi^+\pi^-\pi^0\eta$
$113 \pm 9 \pm 10$	824	3 AKHMETSHIN	17A CMD3	$1.4-2.0 e^+e^- \rightarrow \omega\eta$
110 ± 20	898	4 ACHASOV	16B SND	$1.34-2.00 e^+e^- \rightarrow \omega\eta$
310 ± 30	13.1k	5 AULCHENKO	15A SND	$1.05-1.80 e^+e^- \rightarrow \pi^+\pi^-\pi^0$
$222 \pm 25 \pm 20$		AUBERT	07AU BABR	$10.6 e^+e^- \rightarrow \omega\pi^+\pi^-\gamma$
114 ± 14	13	AUBERT	06D BABR	$10.6 e^+e^- \rightarrow \omega\eta\gamma$
$230 \pm 30 \pm 20$		AUBERT,B	04N BABR	$10.6 e^+e^- \rightarrow \pi^+\pi^-\pi^0\gamma$
$490^+_{-150} 200 \pm 130$	1.2M	6 ACHASOV	03D RVUE	$0.44-2.00 e^+e^- \rightarrow \pi^+\pi^-\pi^0$

lattice QCD

... using the physical light quark mass very challenging here

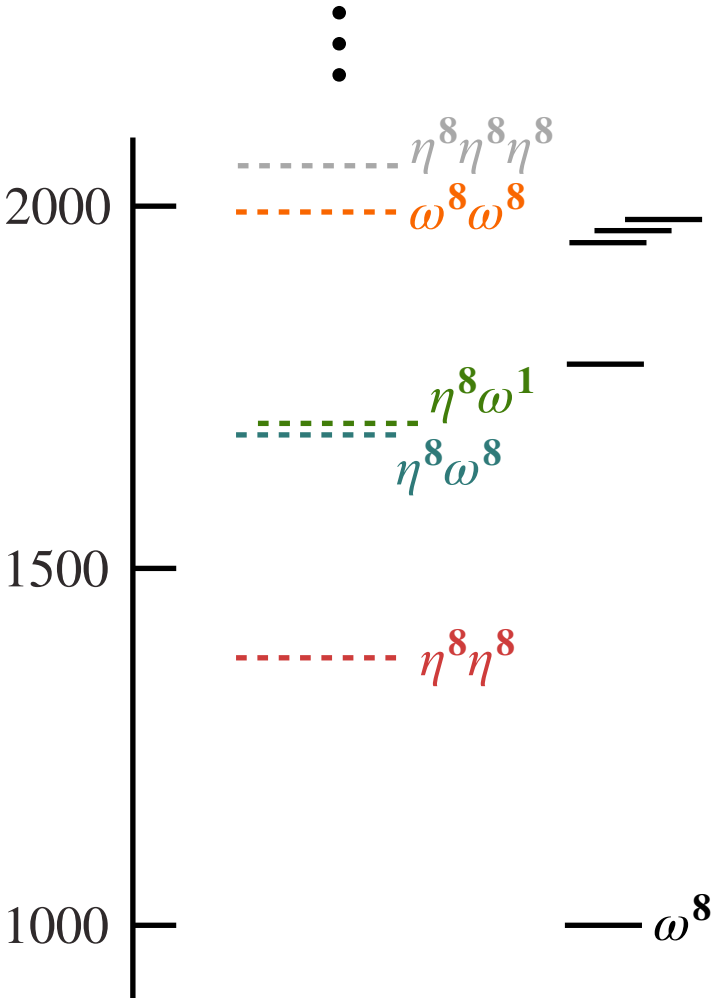
begin with a simpler case, still QCD but take $m_{u,d} \rightarrow m_s$

ensures an **exact** $SU(3)$ flavor symmetry – mesons in either **octet (8)** or **singlet (1)** representations

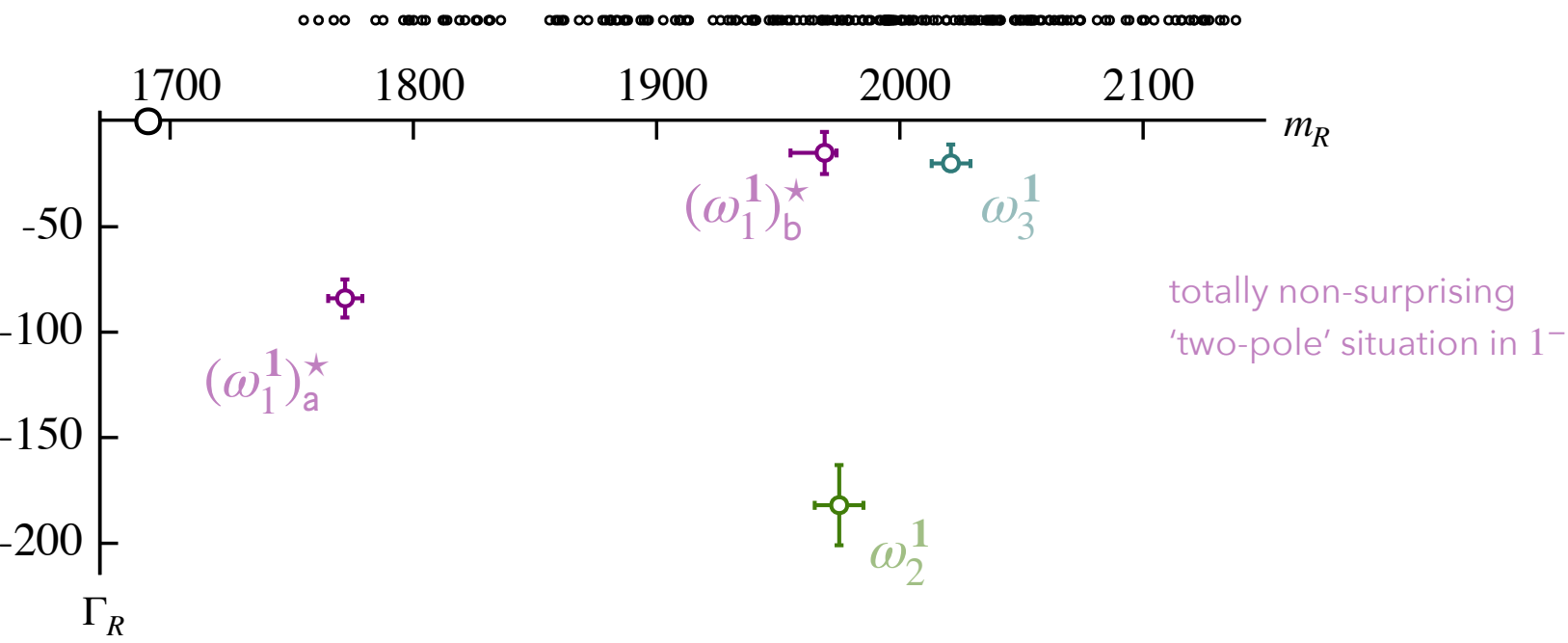
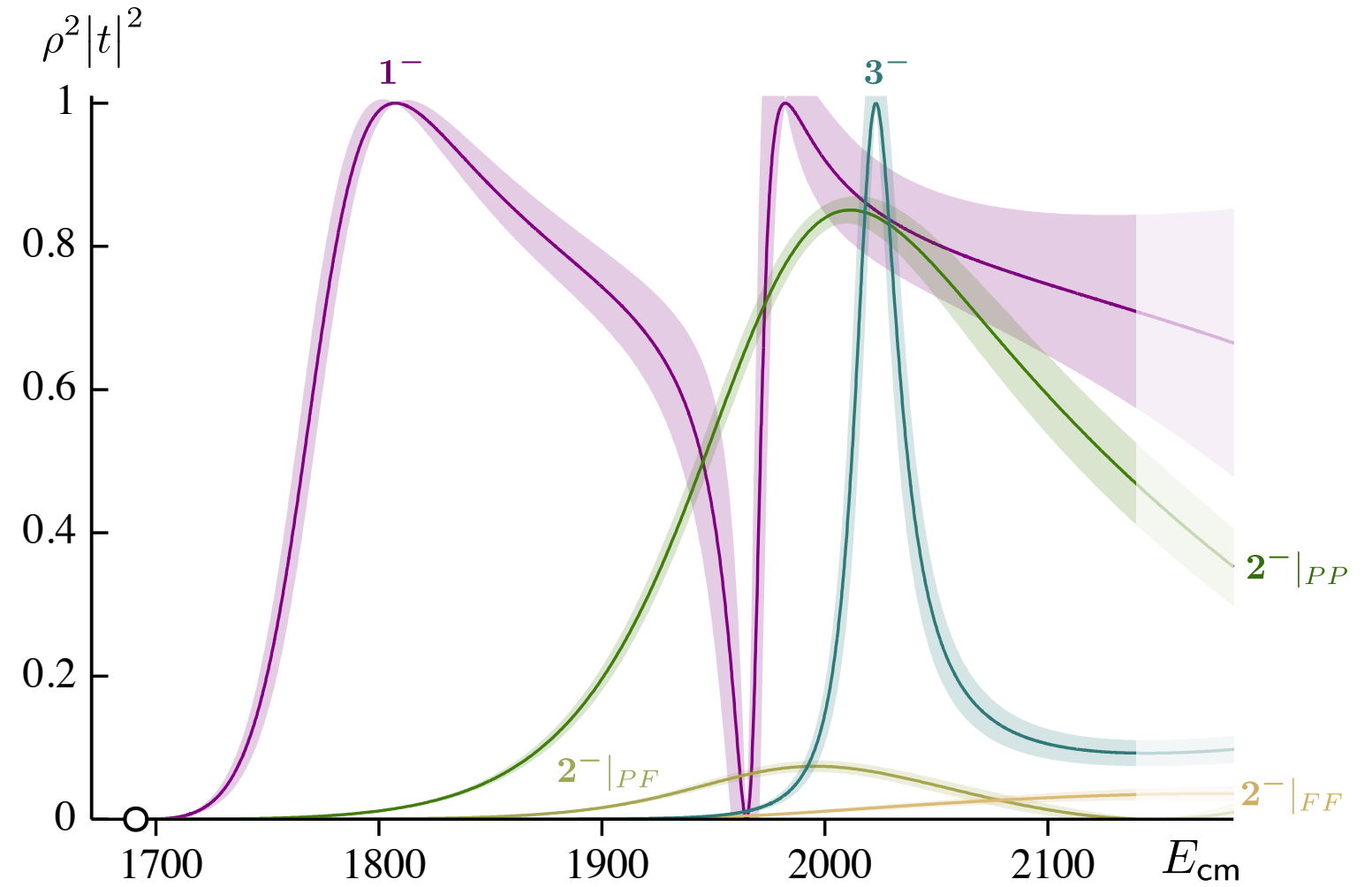
ρ, ω_8, K^* ω_1

lightest pseudoscalar becomes heavy, $m(\eta^8) \sim 700$ MeV

and only meson-meson decays kinematically accessible to the resonances of interest



singlet case computed in 2021
just elastic $\eta^8 \omega^8$ scattering



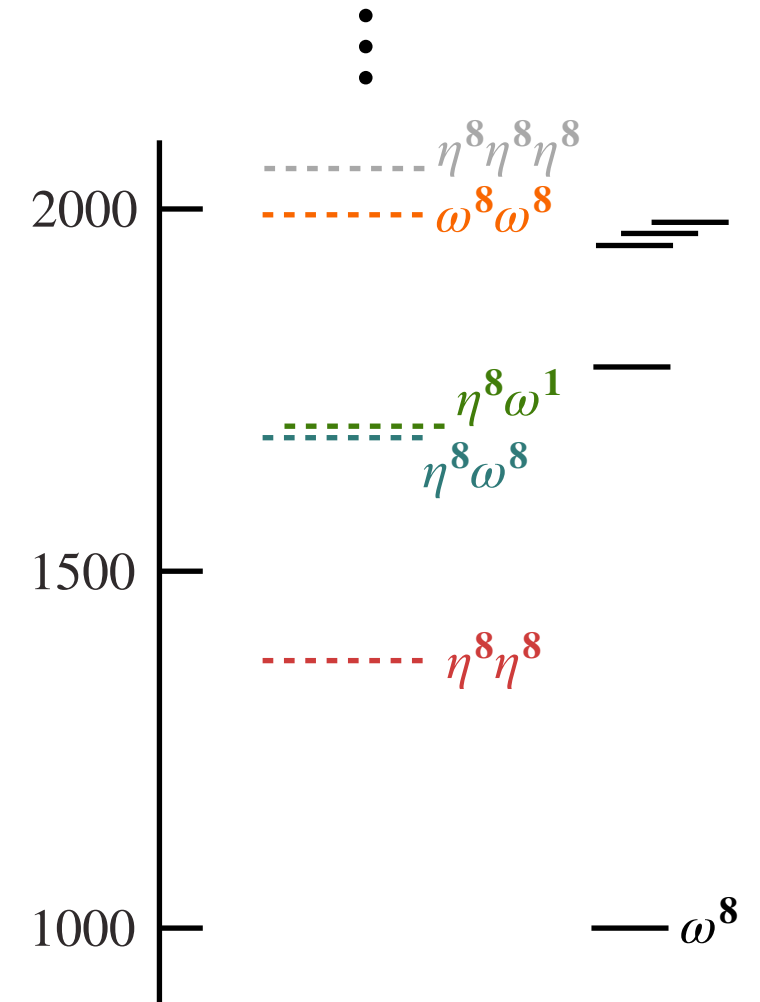
PHYSICAL REVIEW D **103**, 074502 (2021)

Excited J^{--} meson resonances at the SU(3) flavor point from lattice QCD

Christopher T. Johnson^{*,†} and Jozef J. Dudek[†]

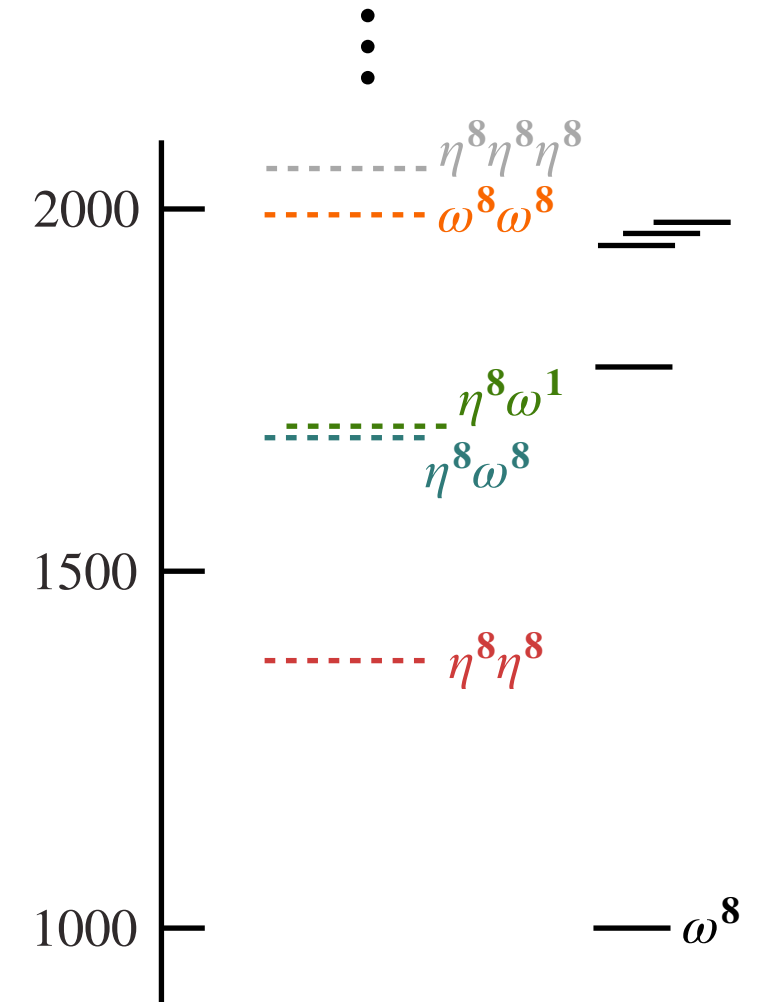
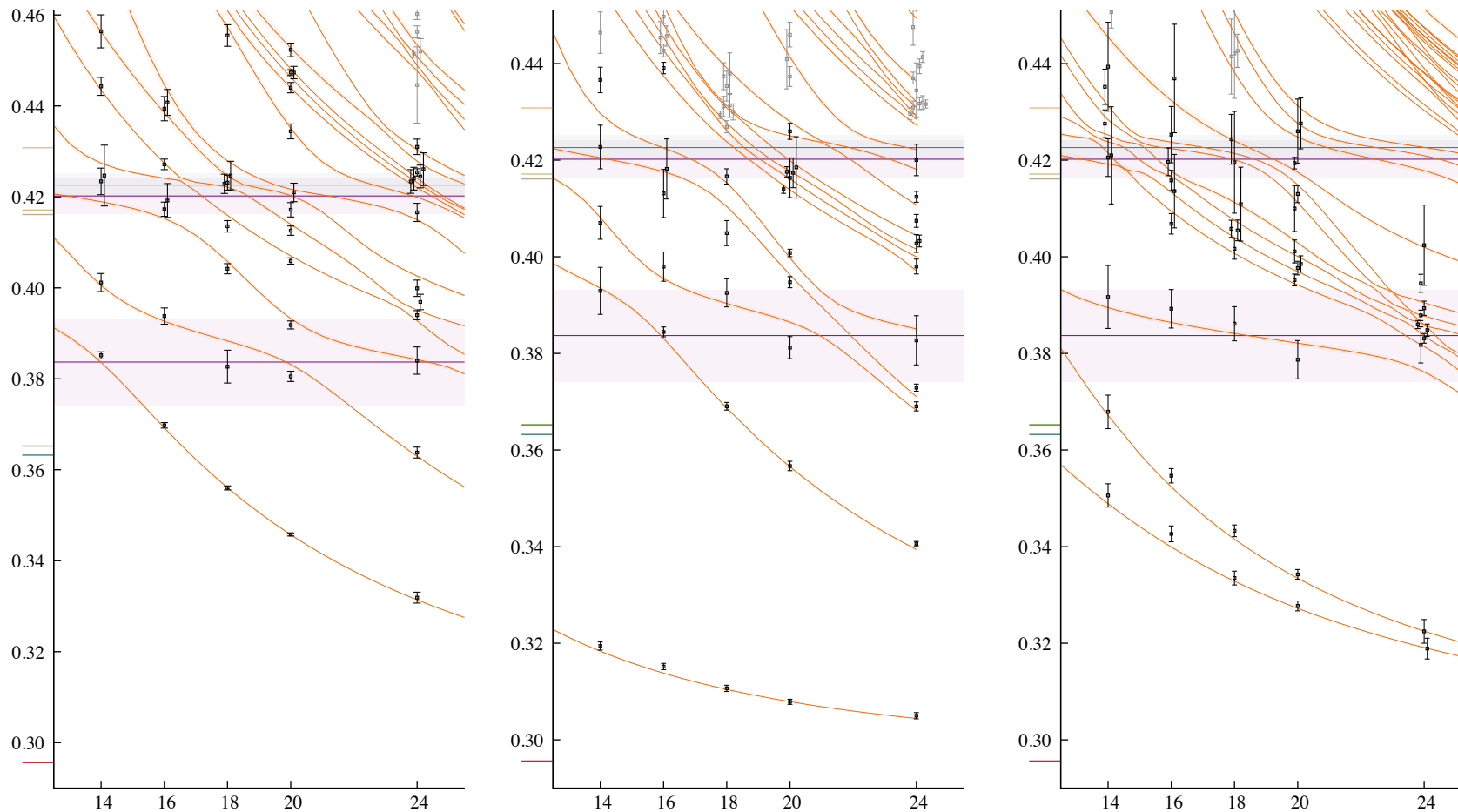
(for the Hadron Spectrum Collaboration)

octet case is new, presenting here for the first time
 coupled-channel $\eta^8\eta^8, \eta^8\omega^8, \eta^8\omega^1$ scattering

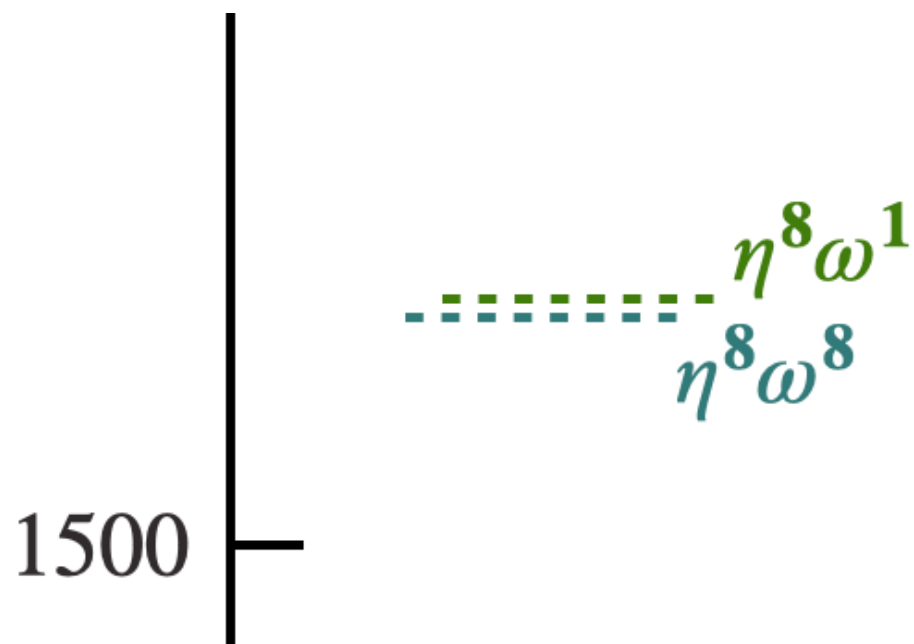
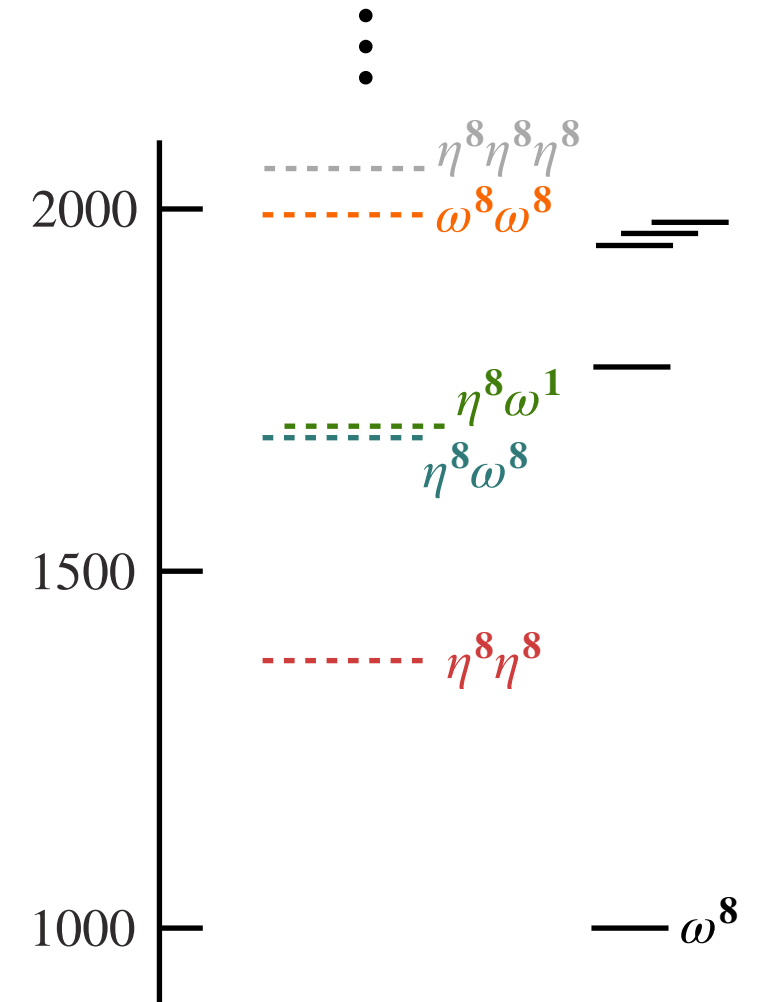


octet case is new, presenting here for the first time
 coupled-channel $\eta^8\eta^8, \eta^8\omega^8, \eta^8\omega^1$ scattering

computed in five lattice volumes,
 hundreds of energy levels constraining
 the scattering amplitudes

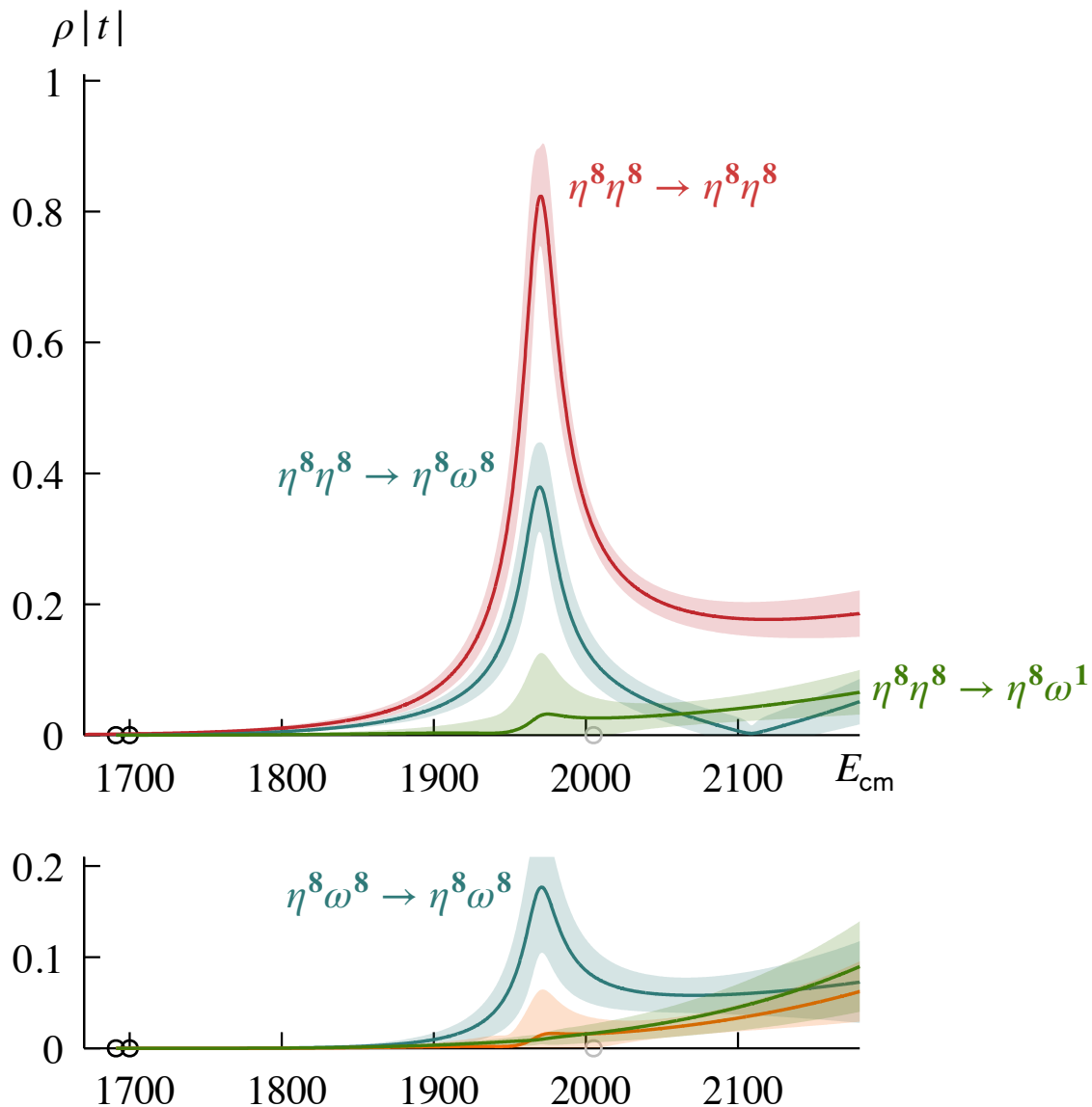


octet case is new, presenting here for the first time
 coupled-channel $\eta^8\eta^8, \eta^8\omega^8, \eta^8\omega^1$ scattering



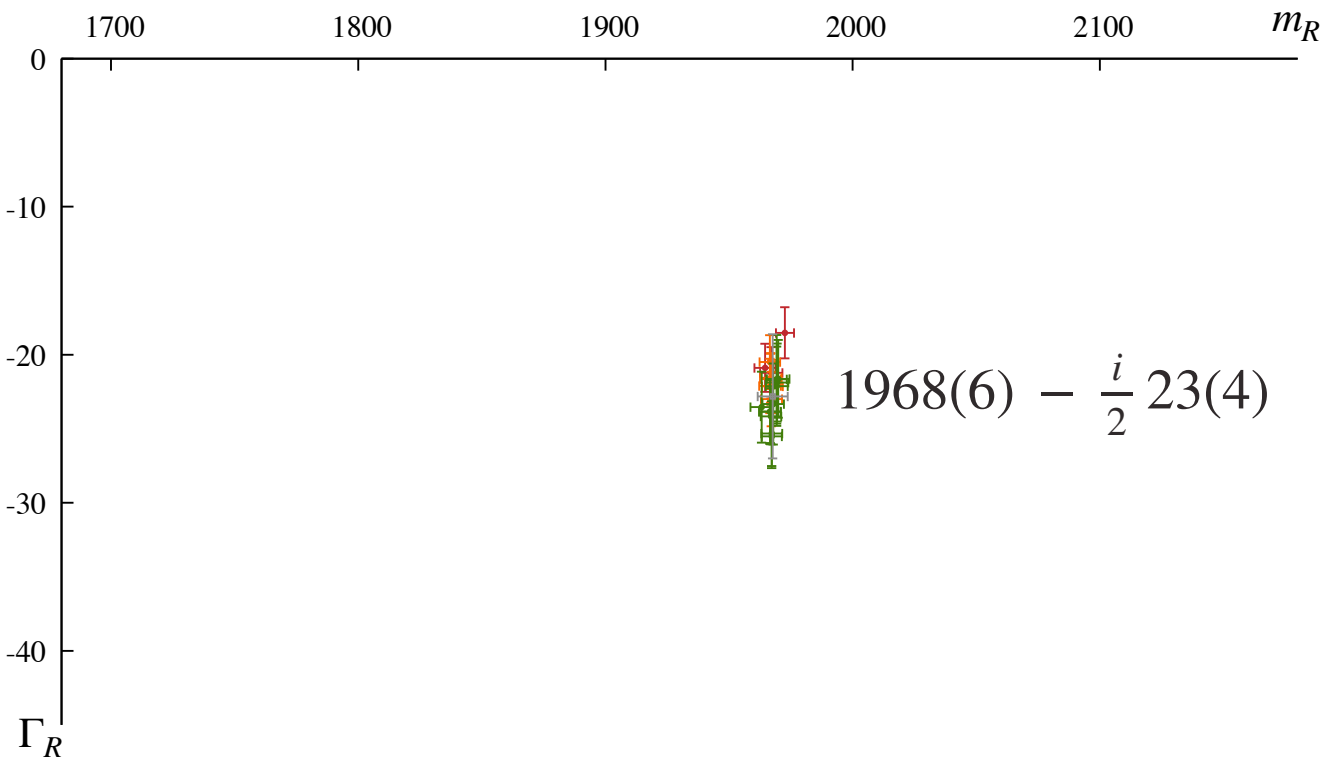
very similar masses for octet and singlet vector mesons
 makes separation of these channels challenging in finite volume

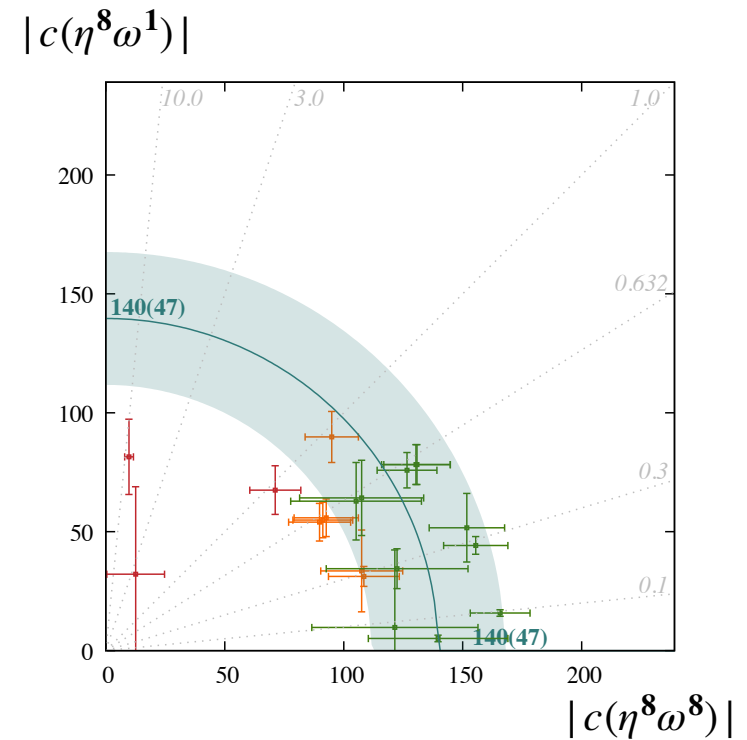
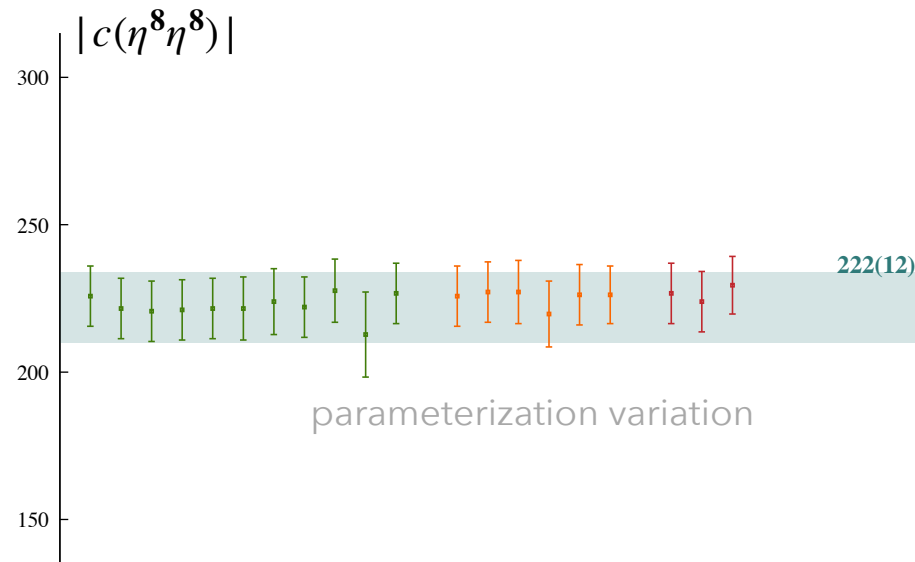
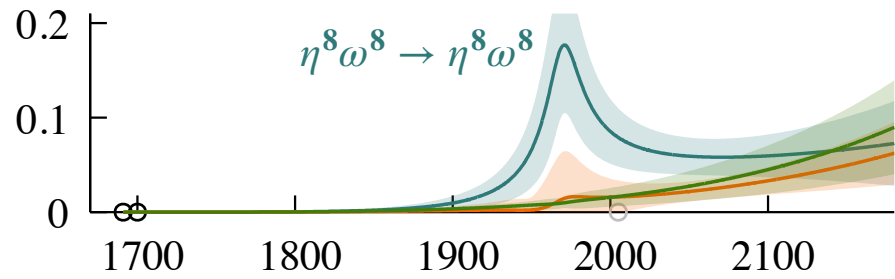
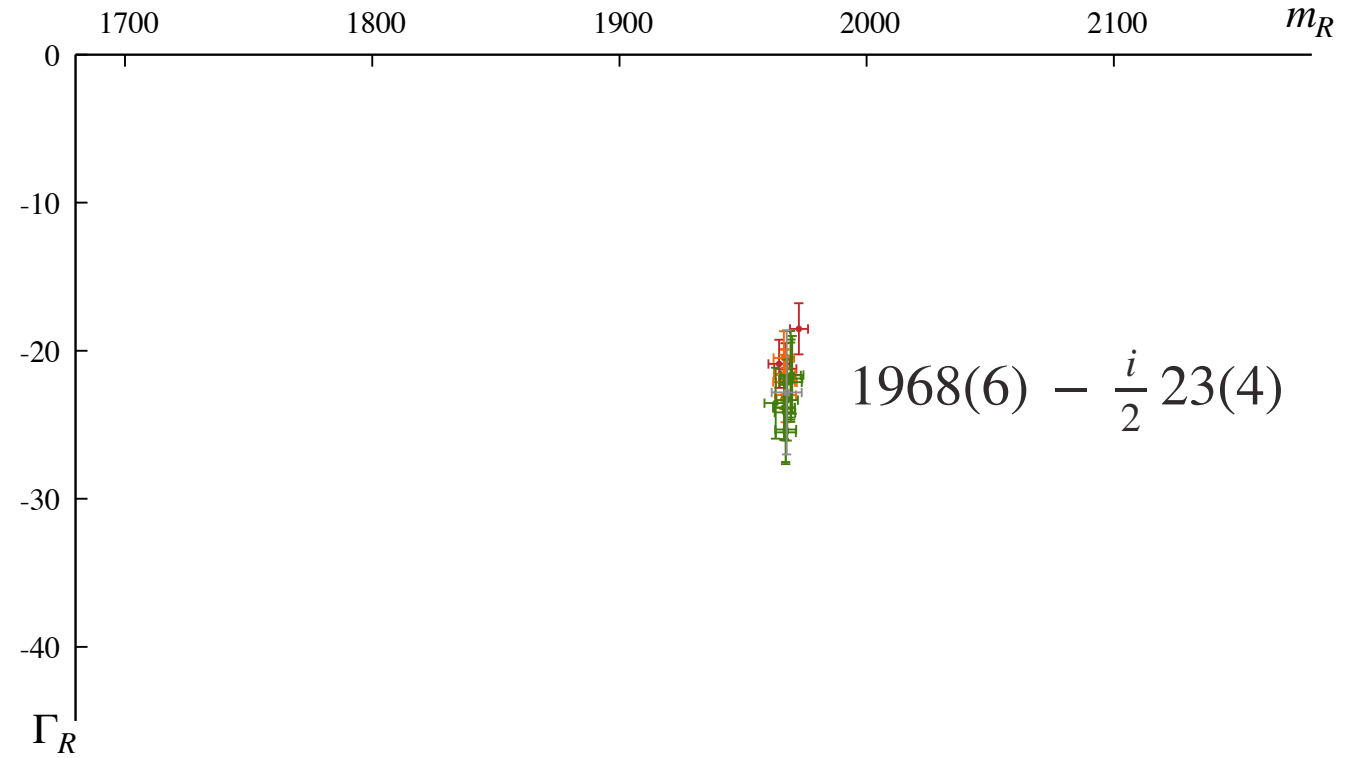
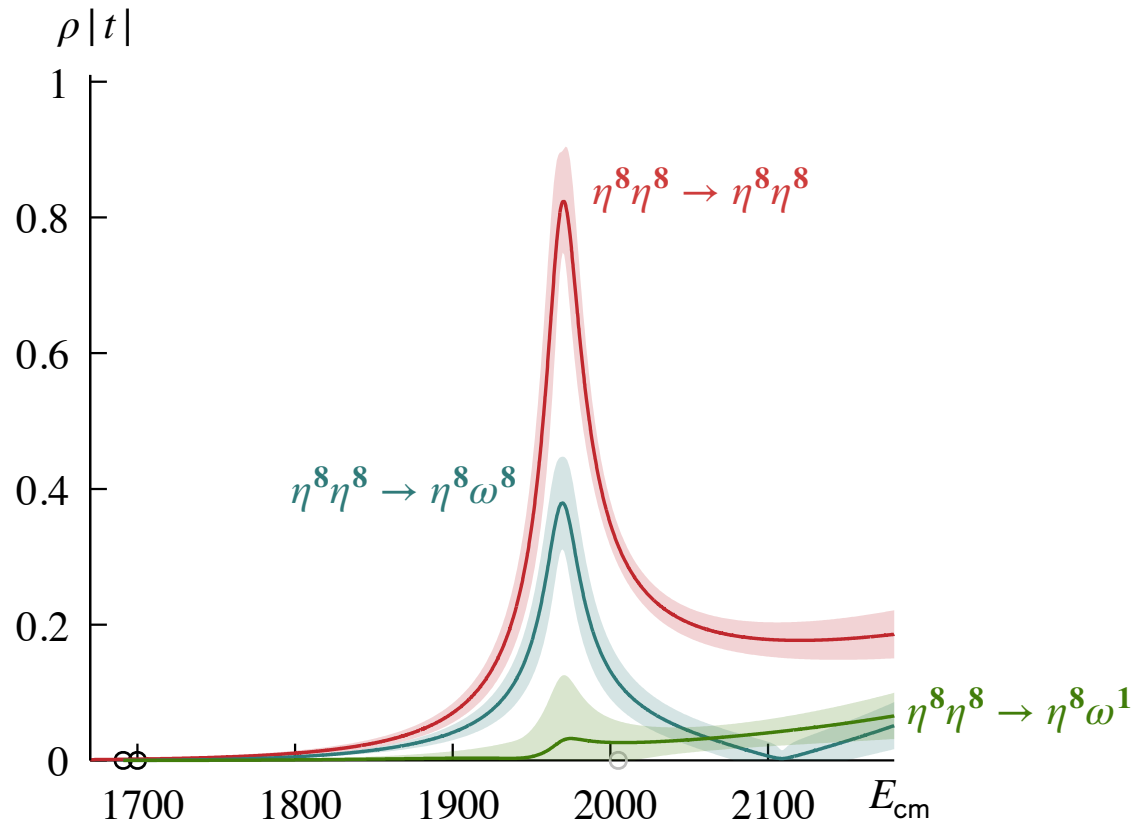
finite-volume approach works via energy shifts

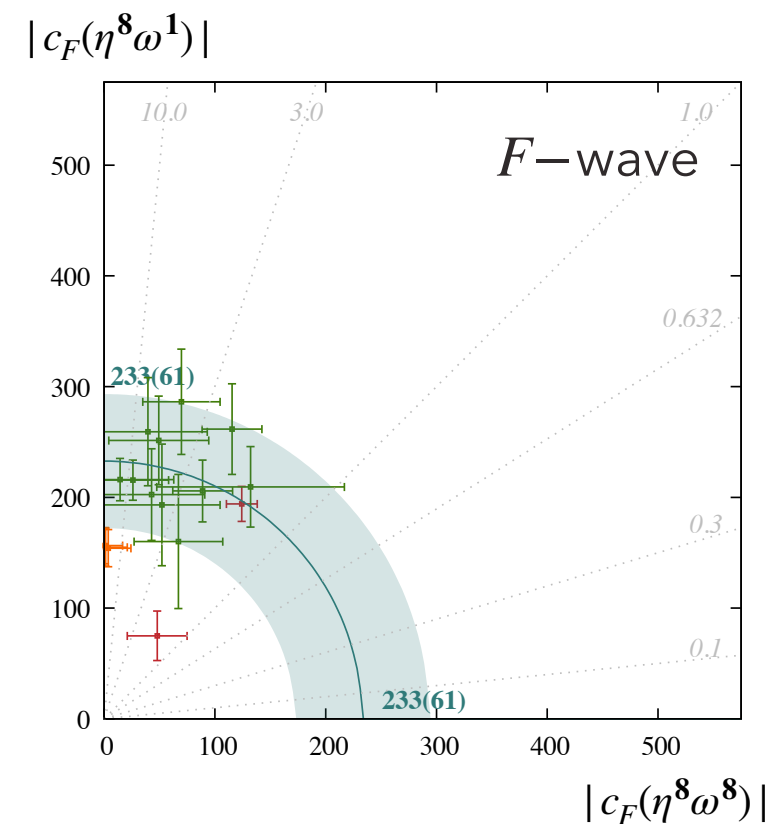
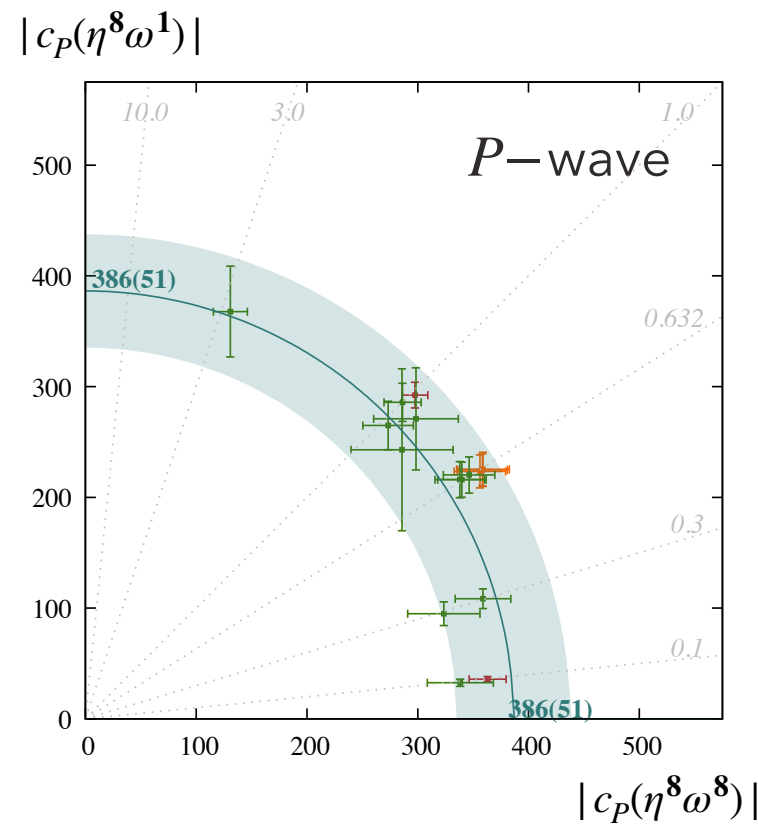
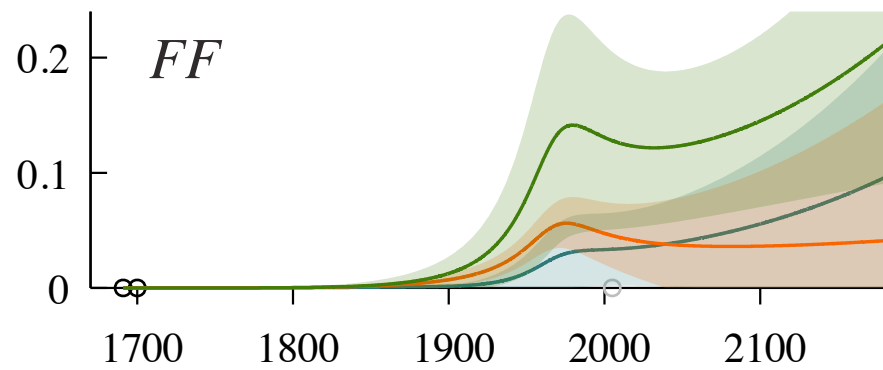
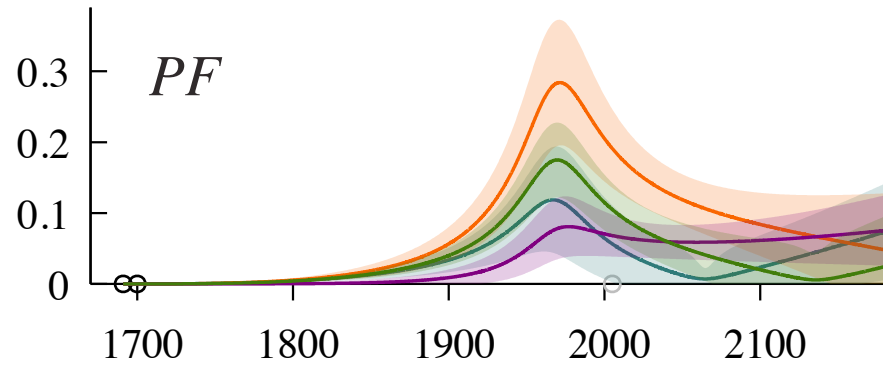
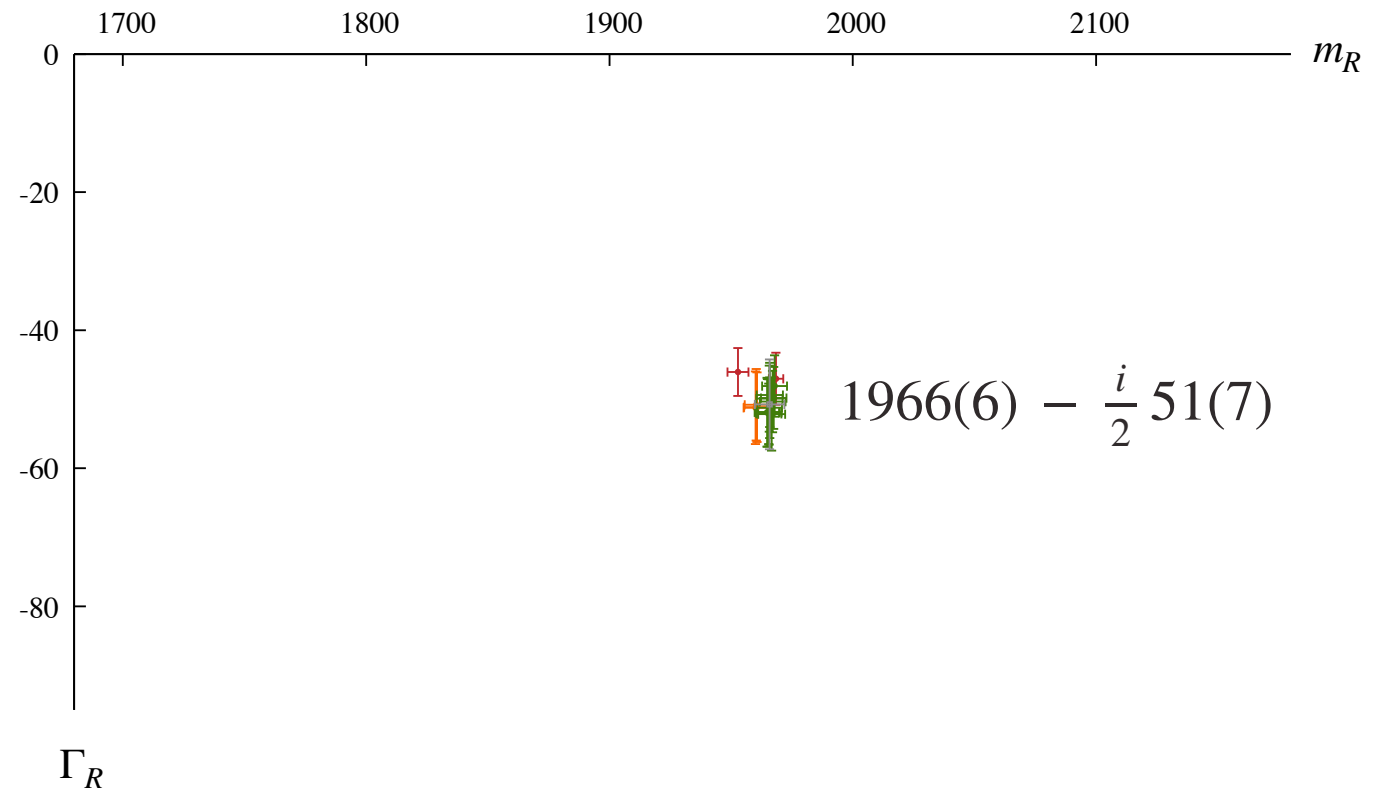
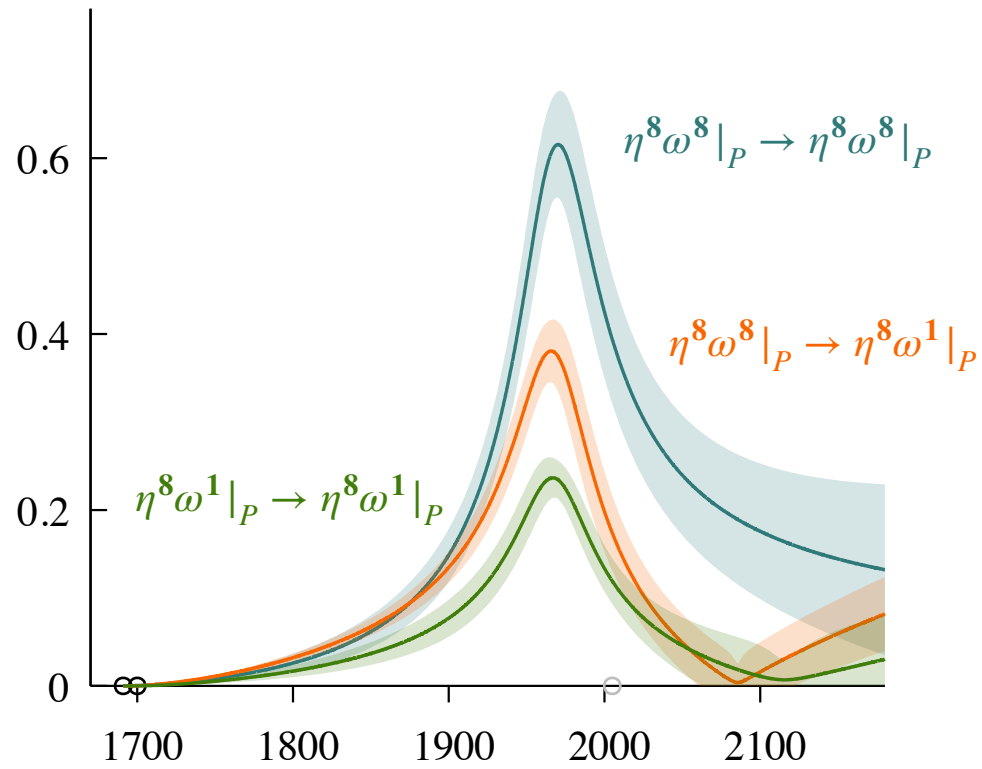


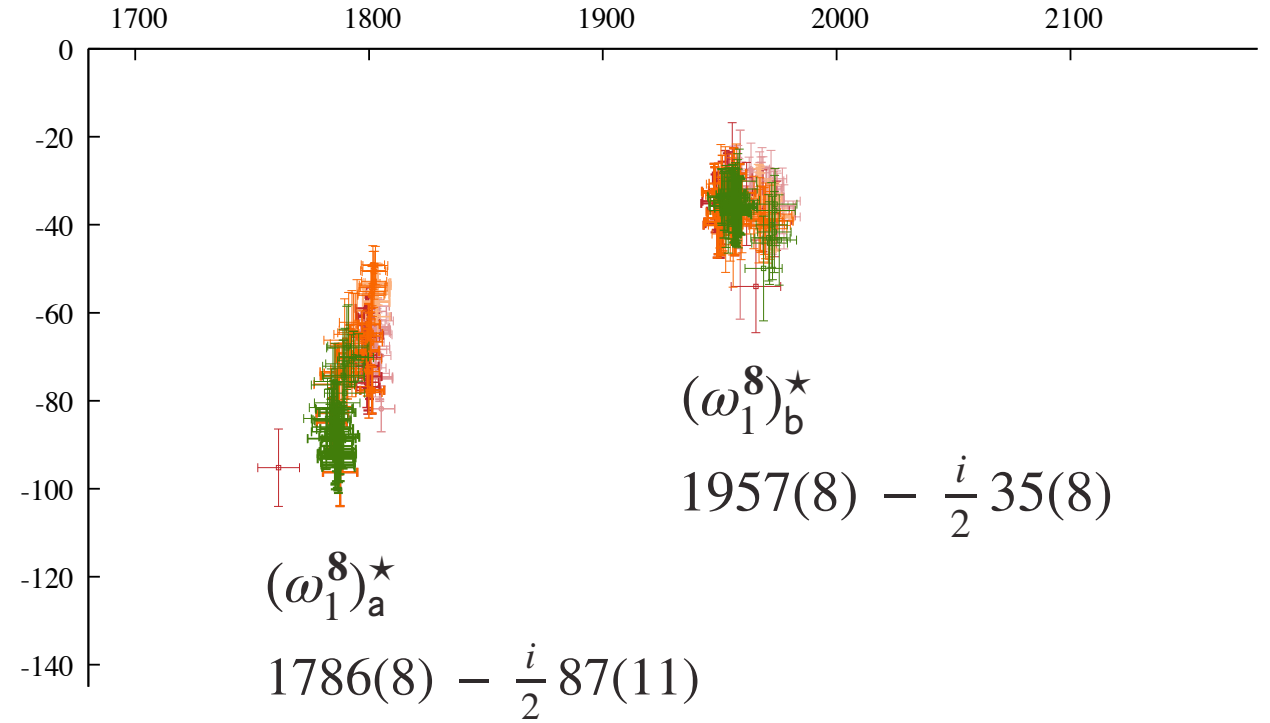
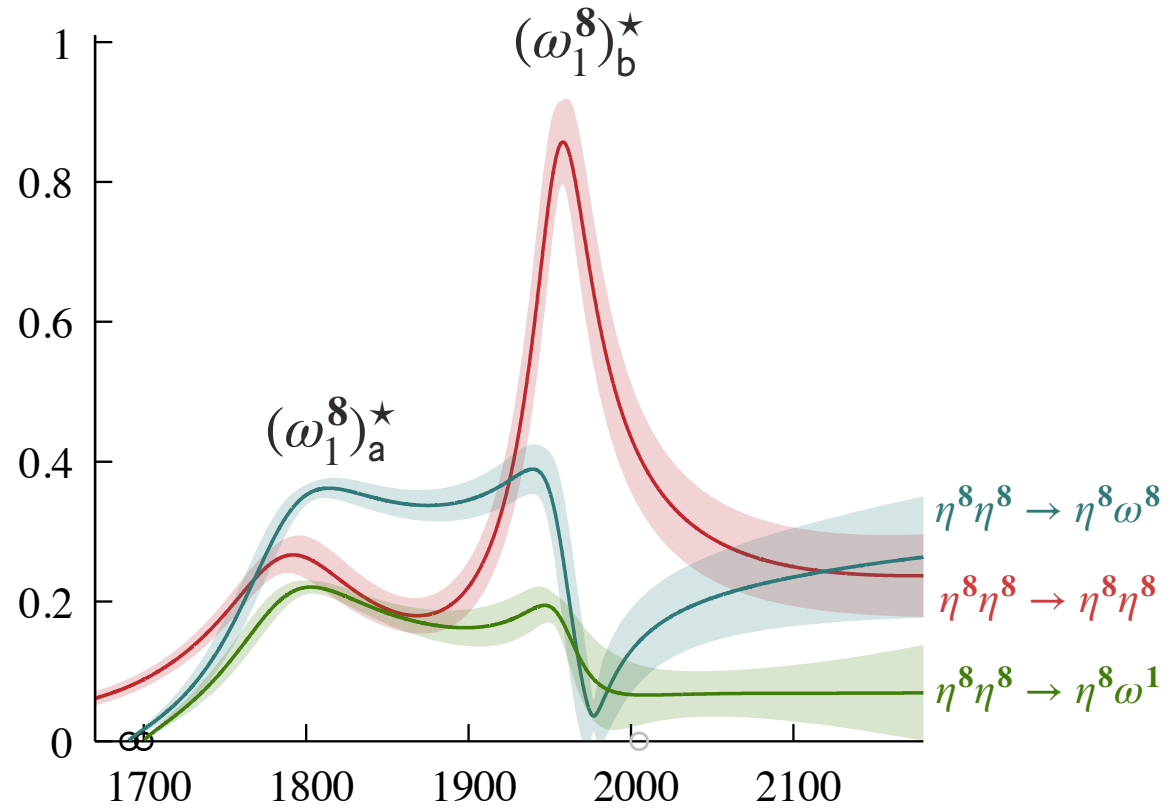
$$t_{ij}(s \sim s_R) \sim \frac{c_i c_j}{s_R - s}$$

$$\sqrt{s_R} = m_R - \frac{i}{2} \Gamma_R$$



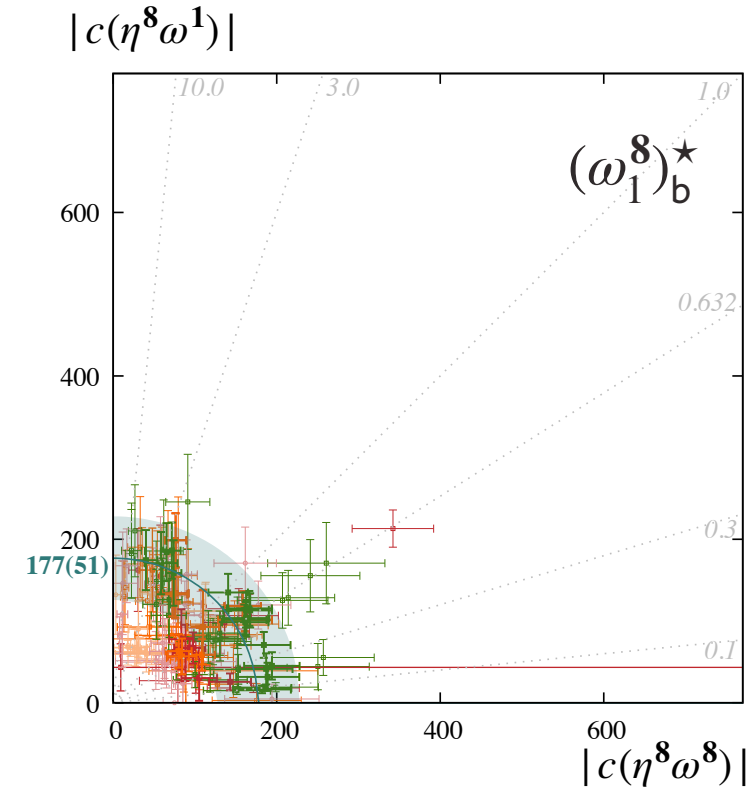
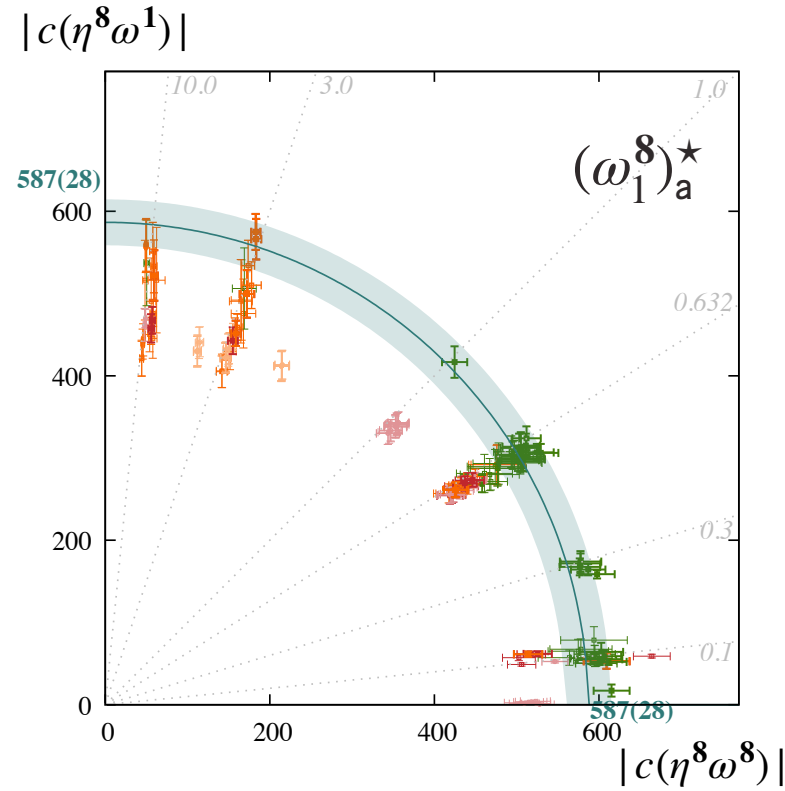
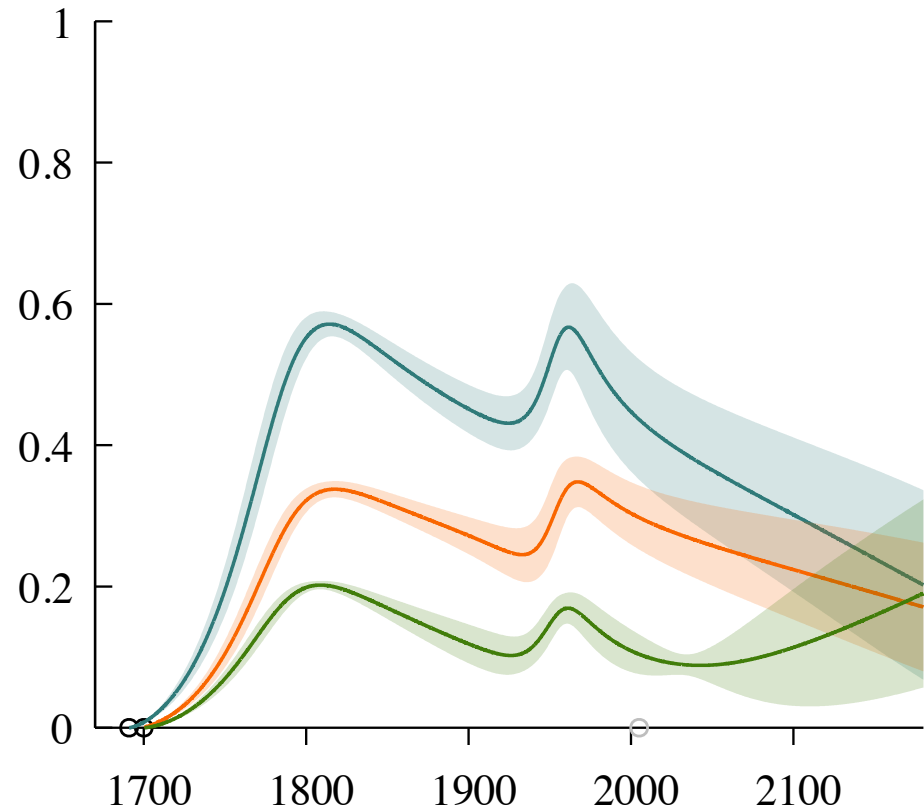


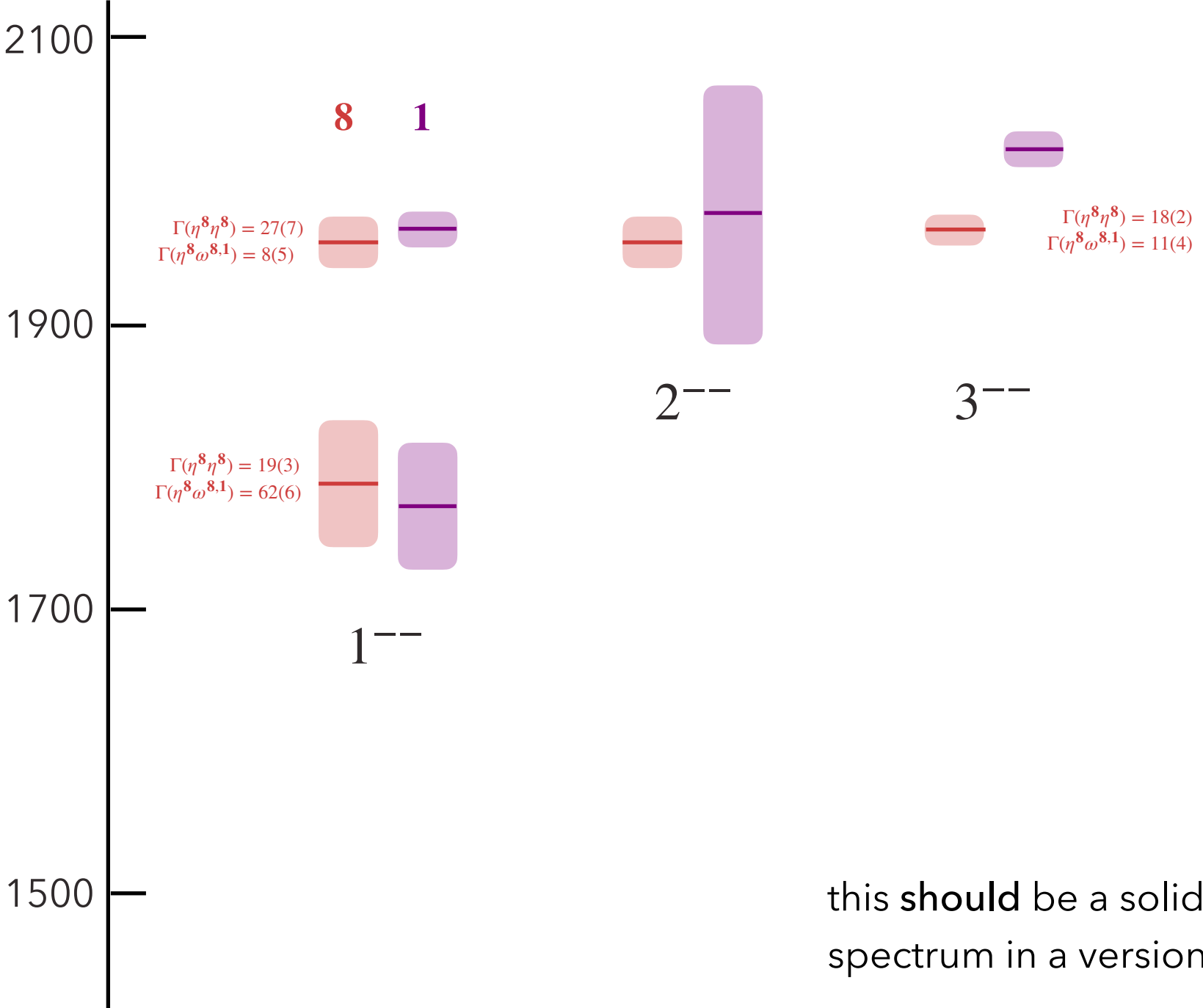




$|c(\eta^8\eta^8)_a| = 233(19)$

$|c(\eta^8\eta^8)_b| = 270(33)$





this **should** be a solid prediction for the spectrum in a version of QCD with $m_u = m_d = m_s \approx m_s^{\text{phys}}$

up to probably modest discretization effects or blunders

decompose $SU(3)$ into (I, I_z, S) states

$$\text{e.g. } |\mathbf{8}(I=0, I_z=0, S=0)\rangle = \frac{1}{\sqrt{20}} (K^+ \bar{K}^{*-} + K^- \bar{K}^{*+} - K^0 \bar{K}^{*0} - \bar{K}^0 K^{*0}) - \frac{1}{\sqrt{5}} (\pi^+ \rho^- + \pi^- \rho^+ - \pi^0 \rho^0 - \eta_8 \omega_8)$$

assume ideally mixed ω, ϕ
assume $\eta = \eta_8$

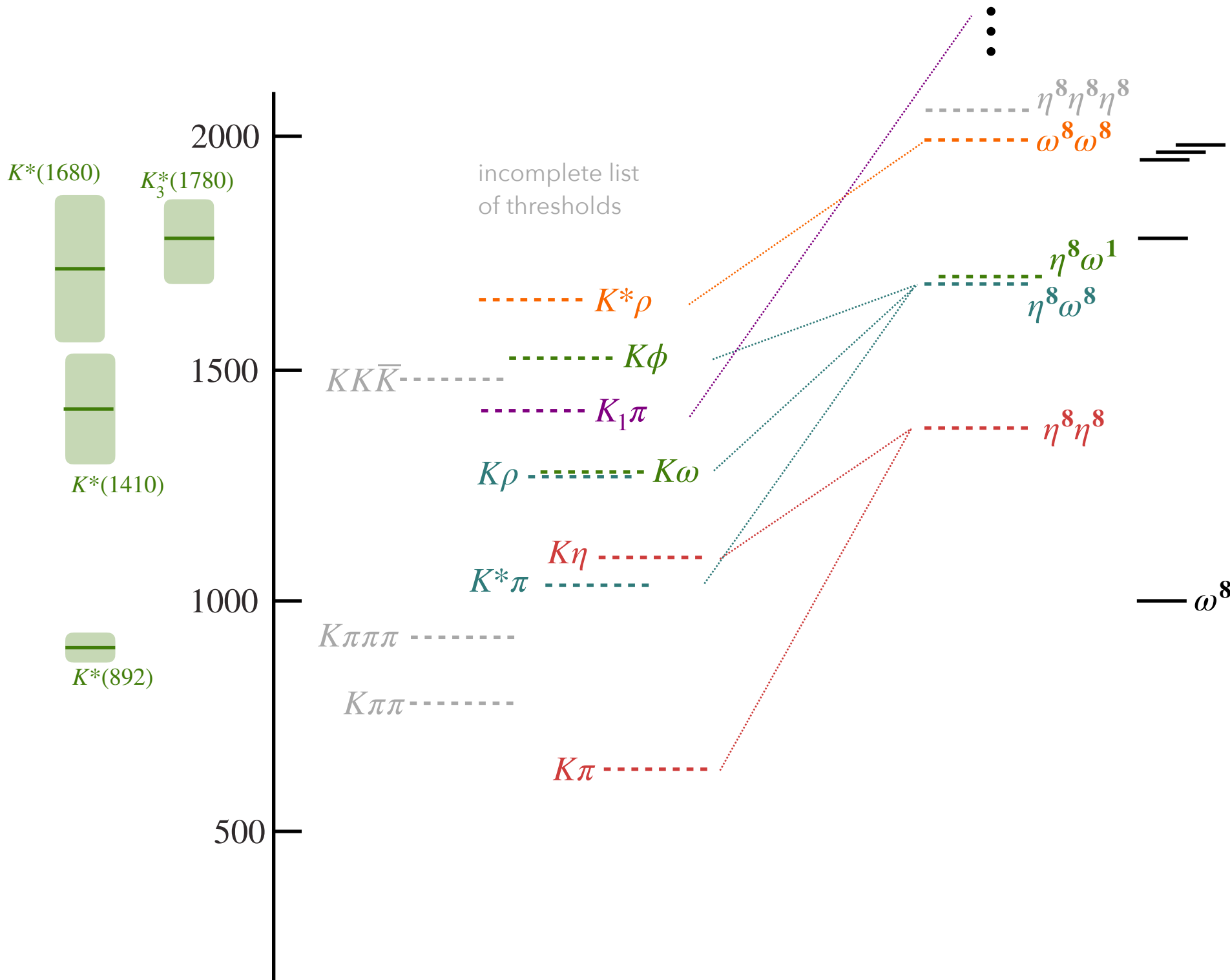
assume couplings are constant with changing light quark mass
apart from angular-momentum barrier:

$$|c_{\text{phys}}| = \frac{k_{\text{phys}}^\ell}{k_{\text{lat}}^\ell} |c_{\text{lat}}|$$

use the PDG masses to set the kinematics

construct partial widths via $\Gamma = \frac{\rho}{M} |c|^2$

everything shown after this point is a MODEL,
not a direct result of a QCD calculation



$K_3^*(1780)$

$$\Gamma(K\pi) \sim 30 \quad \Gamma(K\eta) \sim 13$$

$$\Gamma(K\rho) \sim 9 \quad \Gamma(K^*\pi) \sim 14 \quad \Gamma(K^*\eta) \sim 0 \quad \Gamma(K\omega) \sim 3 \quad \Gamma(K\phi) \sim 1$$

$$\sum_i \Gamma_i \sim 70$$

$K_3^*(1780)$

Mass $m = 1779 \pm 8$ MeV
Full width $\Gamma = 161 \pm 17$ MeV

" $K_2(1780)$ "

$$\Gamma(K\rho) \sim 55 \quad \Gamma(K^*\pi) \sim 74 \quad \Gamma(K^*\eta) \sim 3 \quad \Gamma(K\omega) \sim 17 \quad \Gamma(K\phi) \sim 9$$

$$\sum_i \Gamma_i \sim 157$$

$K_2(1820)$ [ff]

Mass $m = 1819 \pm 12$ MeV
Full width $\Gamma = 264 \pm 34$ MeV

$K_2(1770)$ [ff]

Mass $m = 1773 \pm 8$ MeV
Full width $\Gamma = 186 \pm 14$ MeV

don't yet have the 2^{-+} state to admix

$K^*(1410)$

$$\Gamma(K\pi) \sim 19 \quad \Gamma(K\eta) \sim 9$$

$$\Gamma(K\rho) \sim 41 \quad \Gamma(K^*\pi) \sim 95 \quad \Gamma(K^*\eta) \sim 0 \quad \Gamma(K\omega) \sim 12$$

$$\sum_i \Gamma_i \sim 176$$

$K^*(1410)$

Mass $m = 1414 \pm 15$ MeV
Full width $\Gamma = 232 \pm 21$ MeV

$K^*(1680)$

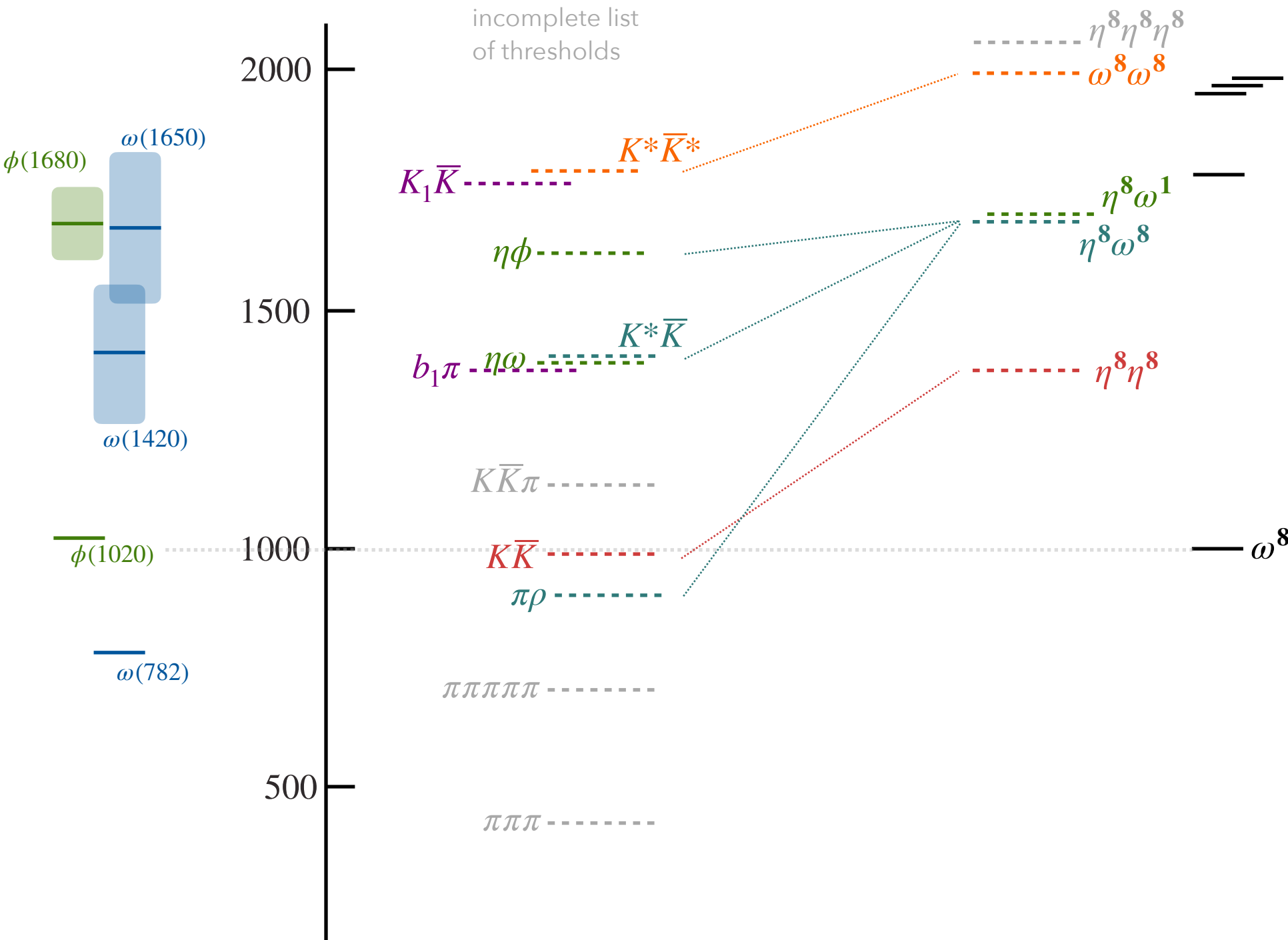
$$\Gamma(K\pi) \sim 24 \quad \Gamma(K\eta) \sim 16$$

$$\Gamma(K\rho) \sim 5 \quad \Gamma(K^*\pi) \sim 7 \quad \Gamma(K^*\eta) \sim 0 \quad \Gamma(K\omega) \sim 2 \quad \Gamma(K\phi) \sim 1$$

$$\sum_i \Gamma_i \sim 56$$

$K^*(1680)$

Mass $m = 1718 \pm 18$ MeV
Full width $\Gamma = 322 \pm 110$ MeV



in this case states are superpositions of **octet** and **singlet**

$SU(3)$ point calculation can't tell us the mixing angle – leave as a free parameter

$$|\omega^*\rangle = \cos \theta |\omega_8\rangle - \sin \theta |\omega_1\rangle$$

$$|\phi^*\rangle = \sin \theta |\omega_8\rangle + \cos \theta |\omega_1\rangle$$

ideal flavor $\theta = -54.74^\circ$,
 $\cos \theta = \sqrt{\frac{1}{3}}$, $\sin \theta = -\sqrt{\frac{2}{3}}$

ideal flavor mixing:

$$\Gamma(\omega^* \rightarrow K\bar{K}) \sim 6$$

$$\Gamma(\omega^* \rightarrow \pi\rho) \sim 61$$

$$\Gamma(\omega^* \rightarrow K\bar{K}^*) \sim 2$$

$$\Gamma(\omega^* \rightarrow \eta\omega) \sim 1$$

$$\Gamma(\omega^* \rightarrow \eta\phi) \sim 0$$

$$\Gamma(\phi^* \rightarrow K\bar{K}) \sim 29$$

$$\Gamma(\phi^* \rightarrow \pi\rho) \sim 0$$

$$\Gamma(\phi^* \rightarrow K\bar{K}^*) \sim 20$$

$$\Gamma(\phi^* \rightarrow \eta\omega) \sim 0$$

$$\Gamma(\phi^* \rightarrow \eta\phi) \sim 3$$

$$\sum_i \Gamma_i \sim 70$$

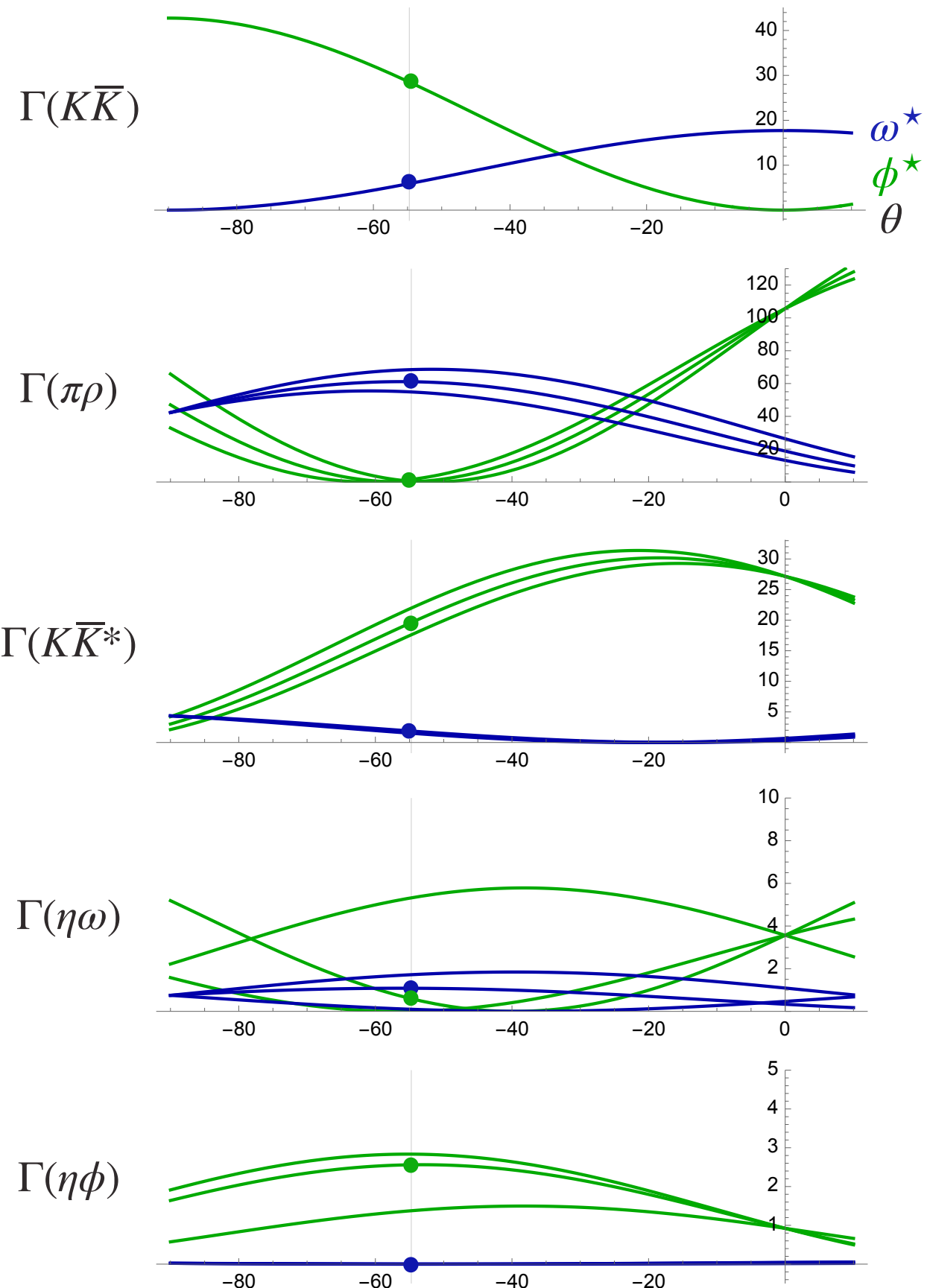
$$\sum_i \Gamma_i \sim 51$$

$\omega_3(1670)$

Mass $m = 1667 \pm 4$ MeV
Full width $\Gamma = 168 \pm 10$ MeV

$\phi_3(1850)$

Mass $m = 1854 \pm 7$ MeV
Full width $\Gamma = 87^{+28}_{-23}$ MeV



ideal flavor mixing:

P-wave, F-wave

$$\Gamma(\omega^* \rightarrow \pi\rho) \sim 237,110$$

$$\Gamma(\phi^* \rightarrow \pi\rho) \sim 6,6$$

$$\Gamma(\omega^* \rightarrow K\bar{K}^*) \sim 38,2$$

$$\Gamma(\phi^* \rightarrow K\bar{K}^*) \sim 102,35$$

$$\Gamma(\omega^* \rightarrow \eta\omega) \sim 12,2$$

$$\Gamma(\phi^* \rightarrow \eta\omega) \sim 0,0$$

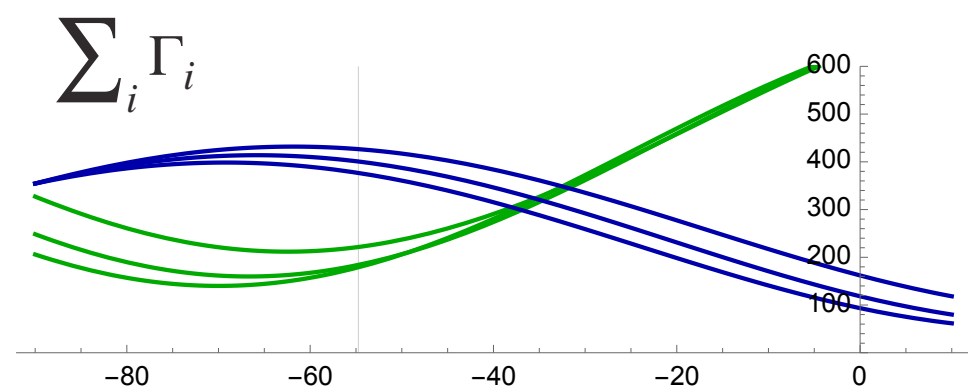
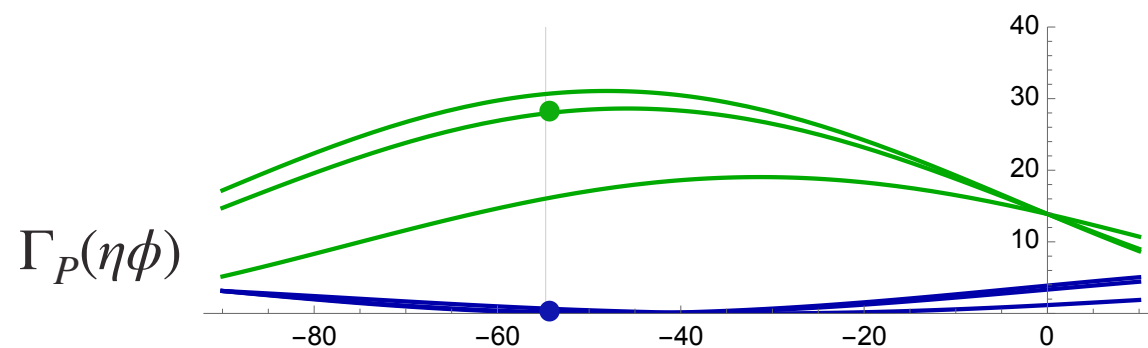
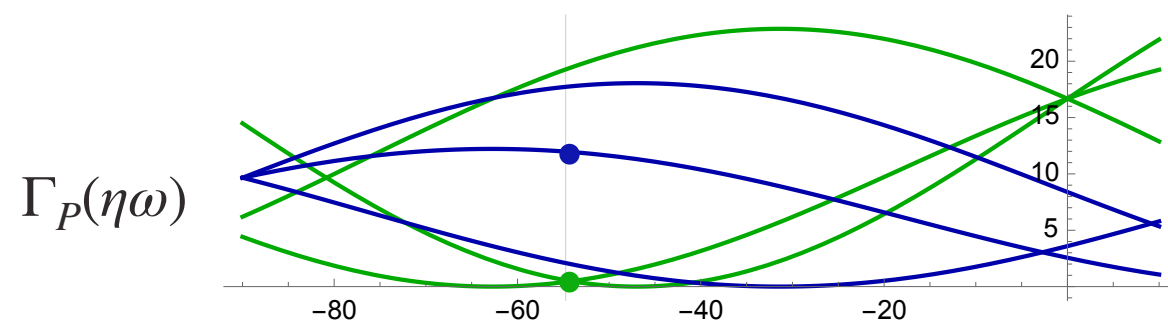
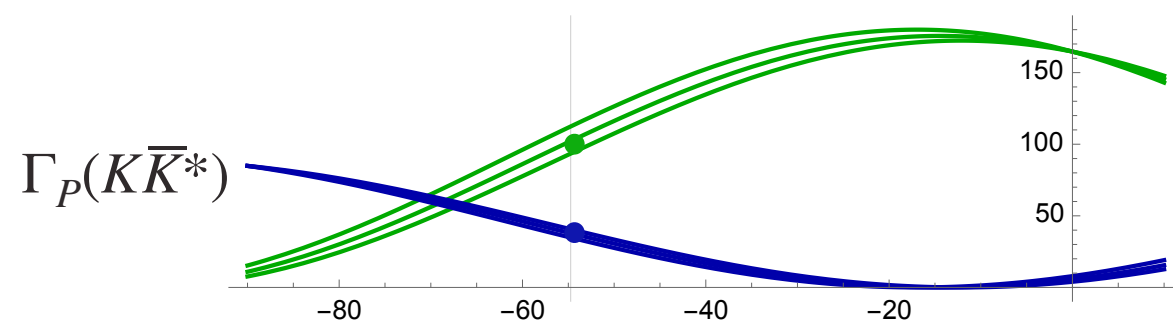
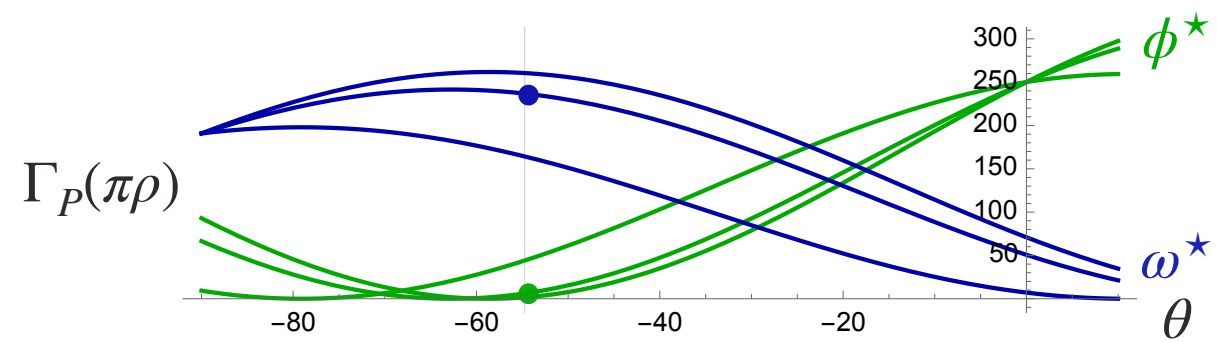
$$\Gamma(\omega^* \rightarrow \eta\phi) \sim 0$$

$$\Gamma(\phi^* \rightarrow \eta\phi) \sim 28,6$$

$$\sum_i \Gamma_i \sim 401$$

$$\sum_i \Gamma_i \sim 184$$

! likely to be very broad !



ideal flavor mixing:

$$\Gamma(\omega^* \rightarrow K\bar{K}) \sim 7$$

$$\Gamma(\omega^* \rightarrow \pi\rho) \sim 468$$

$$\Gamma(\omega^* \rightarrow K\bar{K}^*) \sim 2$$

$$\Gamma(\omega^* \rightarrow \eta\omega) \sim 6$$

$$\Gamma(\omega^* \rightarrow \eta\phi) \sim 0$$

$$\Gamma(\phi^* \rightarrow K\bar{K}) \sim 25$$

$$\Gamma(\phi^* \rightarrow \pi\rho) \sim 10$$

$$\Gamma(\phi^* \rightarrow K\bar{K}^*) \sim 191$$

$$\Gamma(\phi^* \rightarrow \eta\omega) \sim 0$$

$$\Gamma(\phi^* \rightarrow \eta\phi) \sim 37$$

$$\sum_i \Gamma_i \sim 482$$

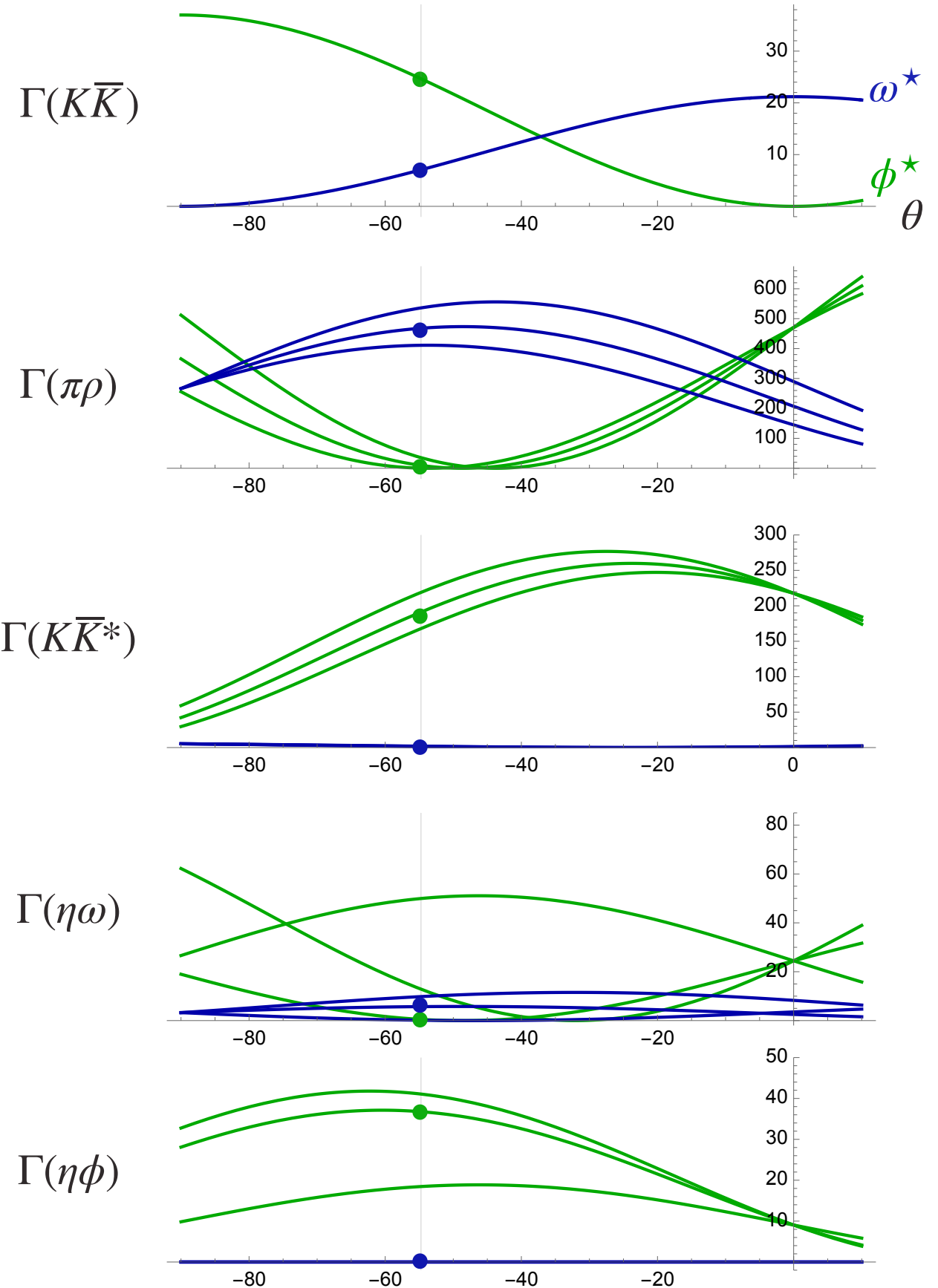
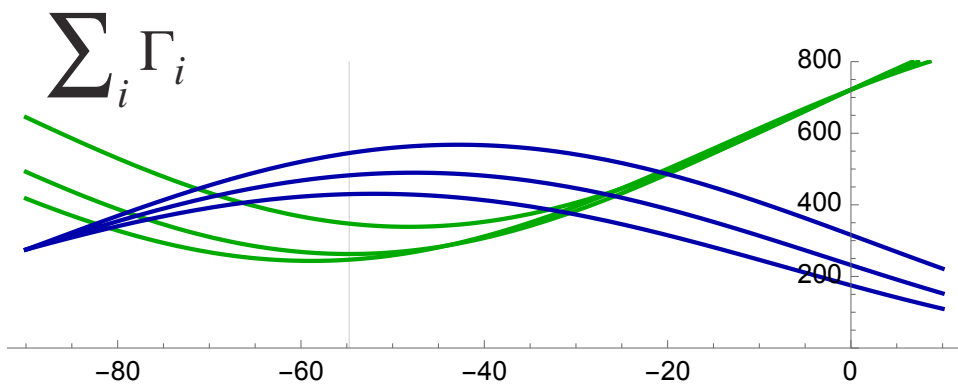
$$\sum_i \Gamma_i \sim 263$$

$\omega(1420)$ [j]

$\phi(1680)$

Mass $m = 1410 \pm 60$ MeV [h]
Full width $\Gamma = 290 \pm 190$ MeV [h]

Mass $m = 1680 \pm 20$ MeV [h]
Full width $\Gamma = 150 \pm 50$ MeV [h]



ideal flavor mixing:

$$\Gamma(\omega^* \rightarrow K\bar{K}) \sim 11$$

$$\Gamma(\phi^* \rightarrow K\bar{K}) \sim 28$$

$$\Gamma(\omega^* \rightarrow \pi\rho) \sim 28$$

$$\Gamma(\phi^* \rightarrow \pi\rho) \sim 0$$

$$\Gamma(\omega^* \rightarrow K\bar{K}^*) \sim 3$$

$$\Gamma(\phi^* \rightarrow K\bar{K}^*) \sim 12$$

$$\Gamma(\omega^* \rightarrow \eta\omega) \sim 1$$

$$\Gamma(\phi^* \rightarrow \eta\omega) \sim 0$$

$$\Gamma(\omega^* \rightarrow \eta\phi) \sim 0$$

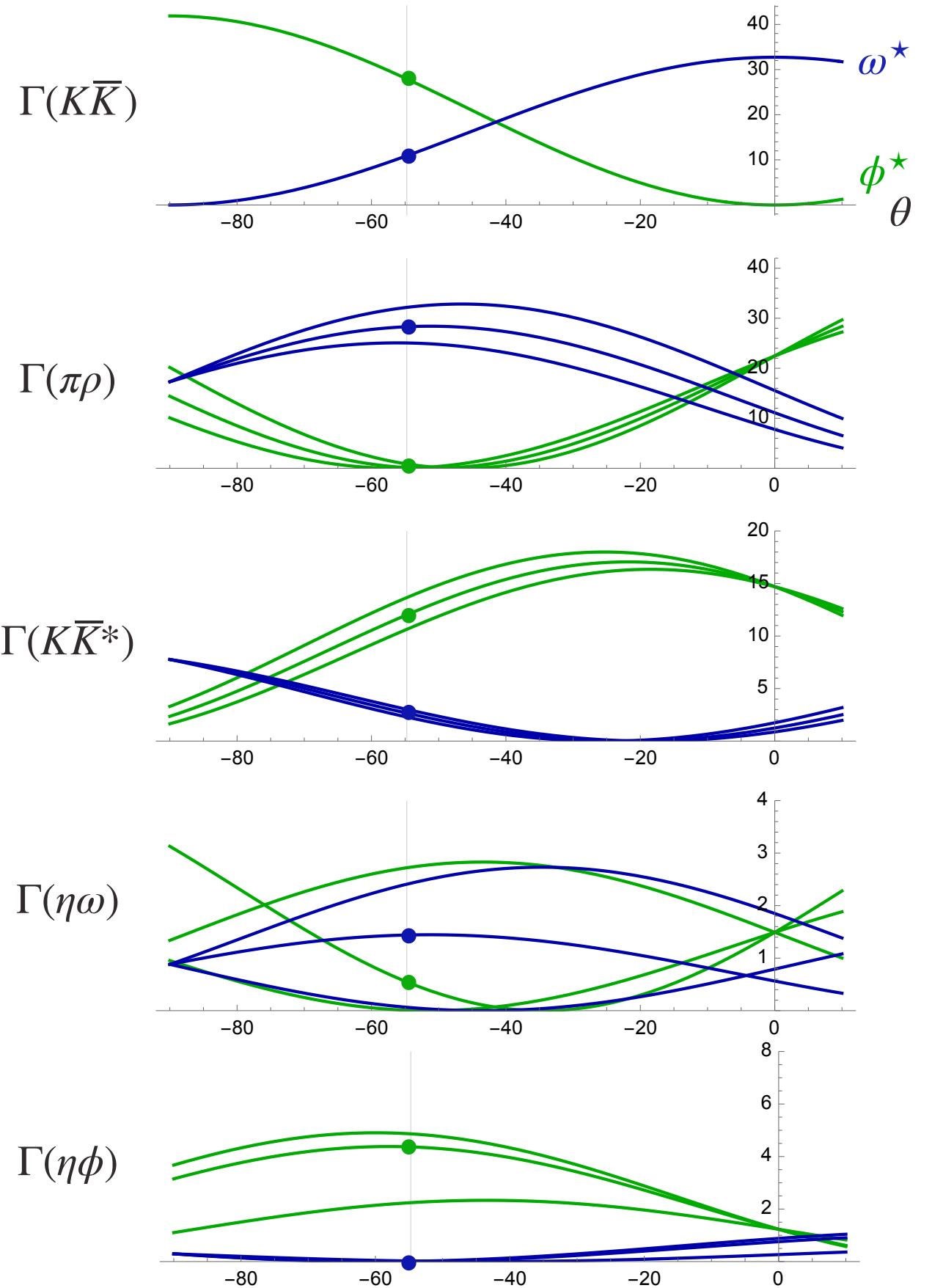
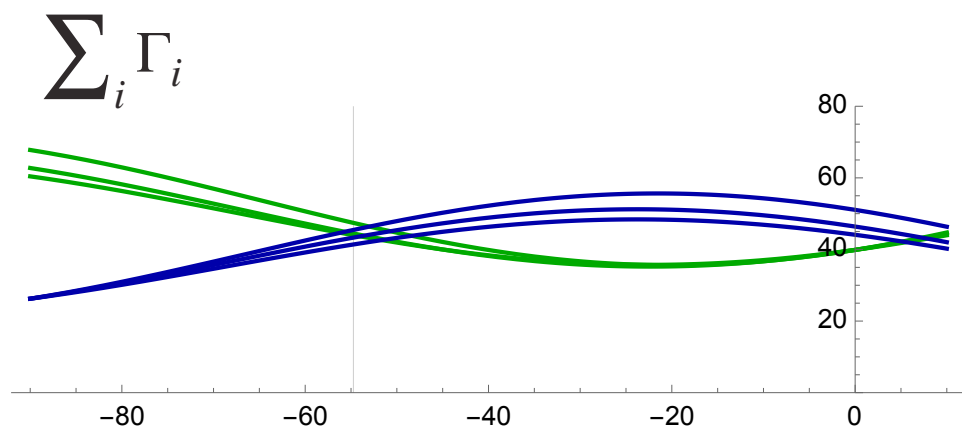
$$\Gamma(\phi^* \rightarrow \eta\phi) \sim 4$$

$$\sum_i \Gamma_i \sim 43$$

$$\sum_i \Gamma_i \sim 44$$

$\omega(1650)$ [k]

Mass $m = 1670 \pm 30$ MeV [h]
Full width $\Gamma = 315 \pm 35$ MeV [h]



not sure I have any yet ... but I can summarize some observations :

SU(3) point calculation doesn't indicate any large deviations from OZI expectations

2^{--} isoscalars likely to be **very broad** (even before 2^{+0-} decays are considered)

the **lighter** 1^- state has large 0^-1^- decay branches – can lead to a large total width

the **heavier** 1^- state has modest 0^-0^- , 0^-1^- decay branches

– if the PDG widths are right, there would need to be modes not considered here

depart from SU(3) symmetry point, by reducing $m_{u,d}$:

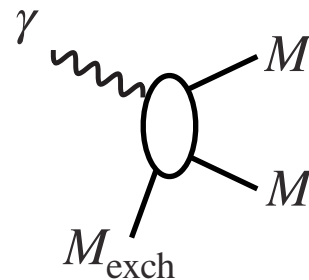
- establish pattern of octet, singlet mixing
- investigate reasonableness of extrapolation scheme

maturity of three-body analysis techniques will allow analysis at closer to physical light quark masses

compute amplitudes closer to those experimentally measured :

- in 1^{--} case can compute $\langle 0 | \bar{\psi} \gamma^\mu \psi | MM \rangle$ to get $A(e^+e^- \rightarrow MM)$

- can calculate contributions to photoproduction treating exchange meson as stable and on-shell



previously done for $\gamma\pi \rightarrow \pi\pi, \gamma K \rightarrow \pi K$

PHYSICAL REVIEW D **93**, 114508 (2016)

$\pi\pi \rightarrow \pi\gamma^*$ amplitude and the resonant $\rho \rightarrow \pi\gamma^*$ transition from lattice QCD

Raúl A. Briceño,^{1,2,*} Jozef J. Dudek,^{1,2} Robert G. Edwards,¹ Christian J. Shultz,²
Christopher E. Thomas,³ and David J. Wilson³
(Hadron Spectrum Collaboration)

PHYSICAL REVIEW D **106**, 114513 (2022)

Radiative decay of the resonant K^* and the $\gamma K \rightarrow K\pi$ amplitude from lattice QCD

Archana Radhakrishnan^{1,2,*} Jozef J. Dudek^{1,2,†} and Robert G. Edwards^{1,2,‡}

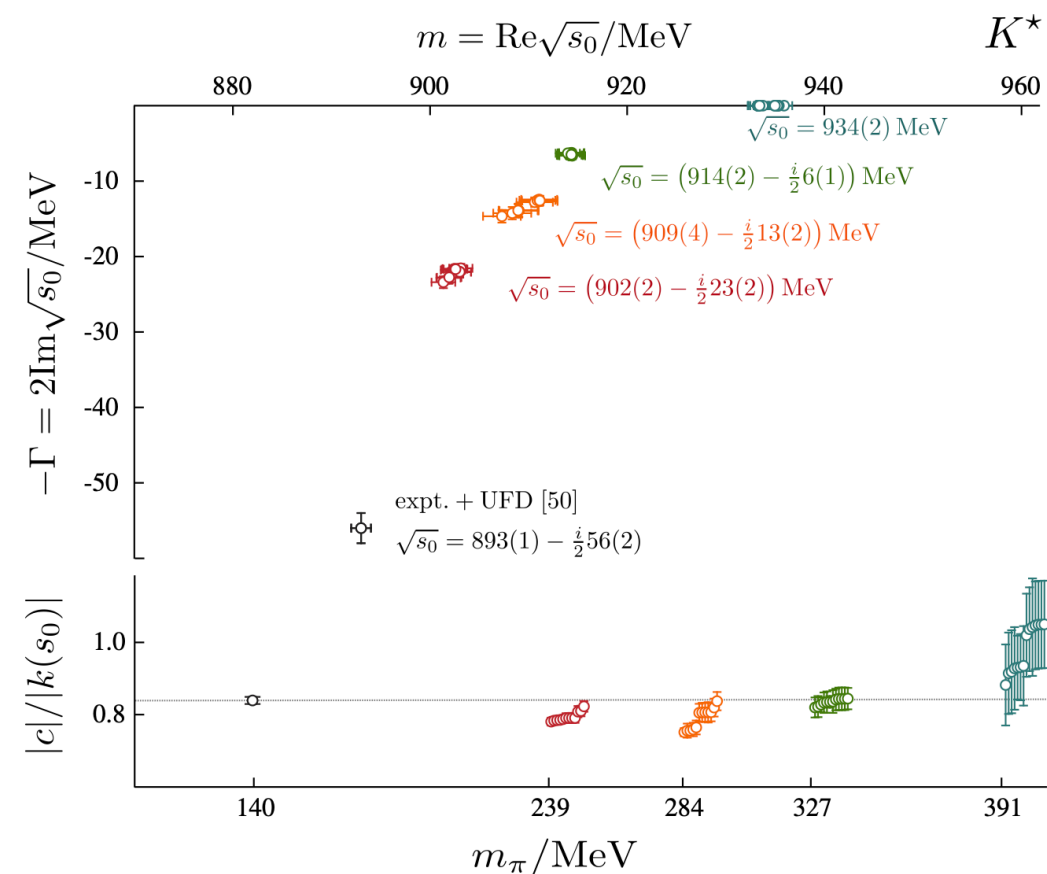
(for the Hadron Spectrum Collaboration)

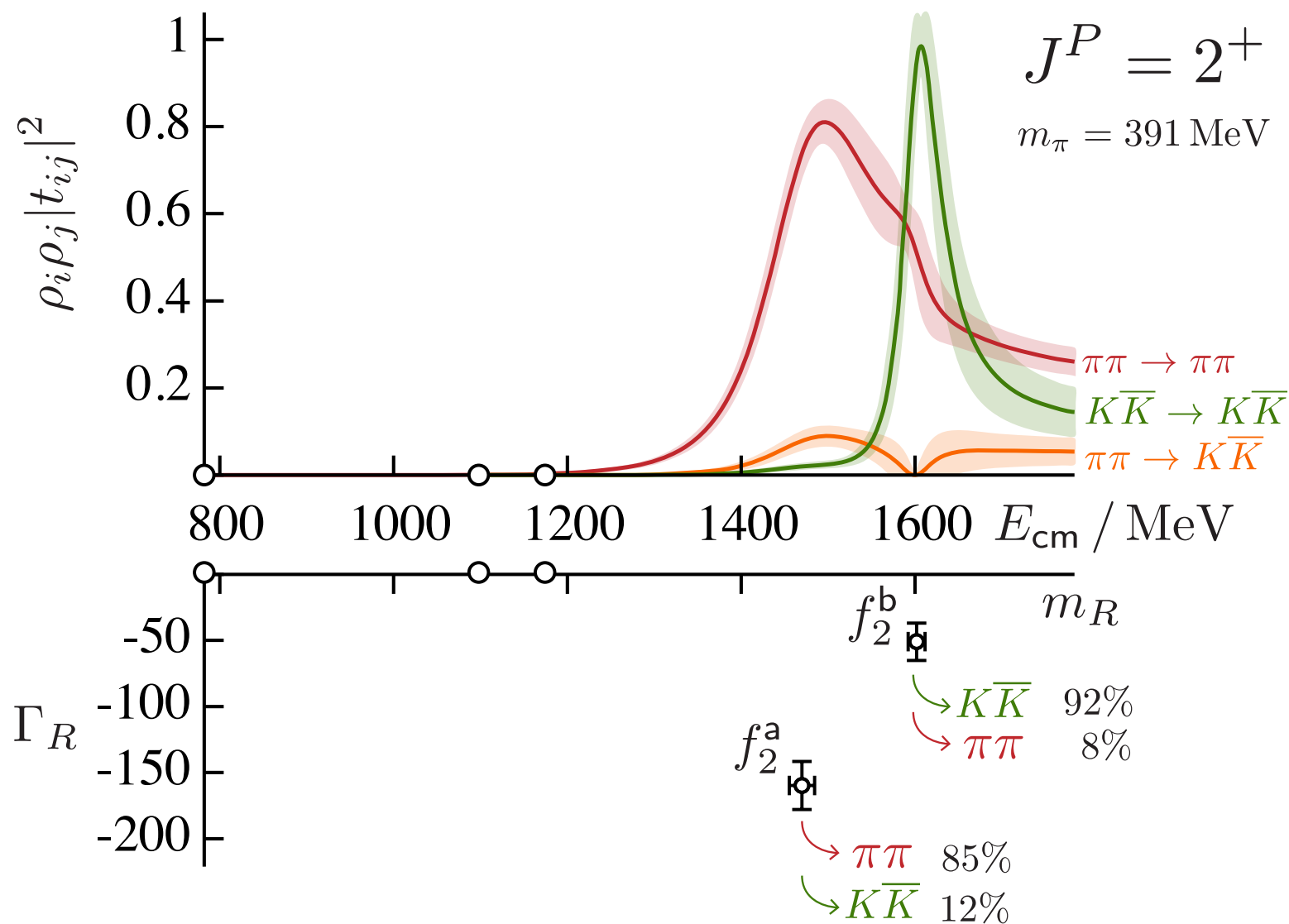
coupling extrapolation

This approach is motivated by observations made in lattice calculations of the decays of $b_1 \rightarrow \omega\pi$ dominantly in S -wave [68], ρ to $\pi\pi$ in P -wave [46], K^* to $K\pi$ in P -wave [93] and f_2, f'_2 decays to $\pi\pi$ and $K\bar{K}$ in D -wave [66], which appear to show quark-mass independence when treated this way. For example in the b_1 case, the coupling computed in [68] at $m_\pi \sim 391$ MeV is $|c| = 564(114)$ MeV, in good agreement with the coupling $|c|^{\text{phys}} = 556(17)$ MeV extracted from the experimental b_1 decay width. In the P -wave ρ decay, an explicit factor of k is required for the scaling to work, as presented in Ref. [46]. In addition, as shown in Fig. 4 of Ref. [93], the K^* coupling scaled in this way is approximately constant for four different light-quark masses corresponding to $m_\pi = 239$ to 391 MeV, even when the K^* is a shallow bound state, and is in agreement with the experimentally measured coupling. Scaling the f_2, f'_2 D -wave couplings computed at $m_\pi \sim 391$ MeV in [66] gives, in comparison to values extracted from the Particle Data Group (PDG) review [90],

	Scaled	PDG
$ c(f_2 \rightarrow \pi\pi) $	488(28)	453_{-4}^{+9} ,
$ c(f_2 \rightarrow K\bar{K}) $	139(27)	132(7),
$ c(f'_2 \rightarrow \pi\pi) $	103(32)	33(4),
$ c(f'_2 \rightarrow K\bar{K}) $	321(50)	389(12),

which is quite a reasonable agreement given the large extrapolation in quark mass.¹³





$$f_2^a \sim u\bar{u} + d\bar{d} \quad f_2^b \sim s\bar{s}$$

couplings from pole residue

	$\frac{a_t c_{\pi\pi} }{(a_t k_{\pi\pi})^2}$	$\frac{a_t c_{K\bar{K}} }{(a_t k_{K\bar{K}})^2}$
f_2^a	7.1(4)	4.8(9)
f_2^b	1.0(3)	5.5(8)

zero in 'OZI' limit

– requires $s\bar{s}$ annihilation

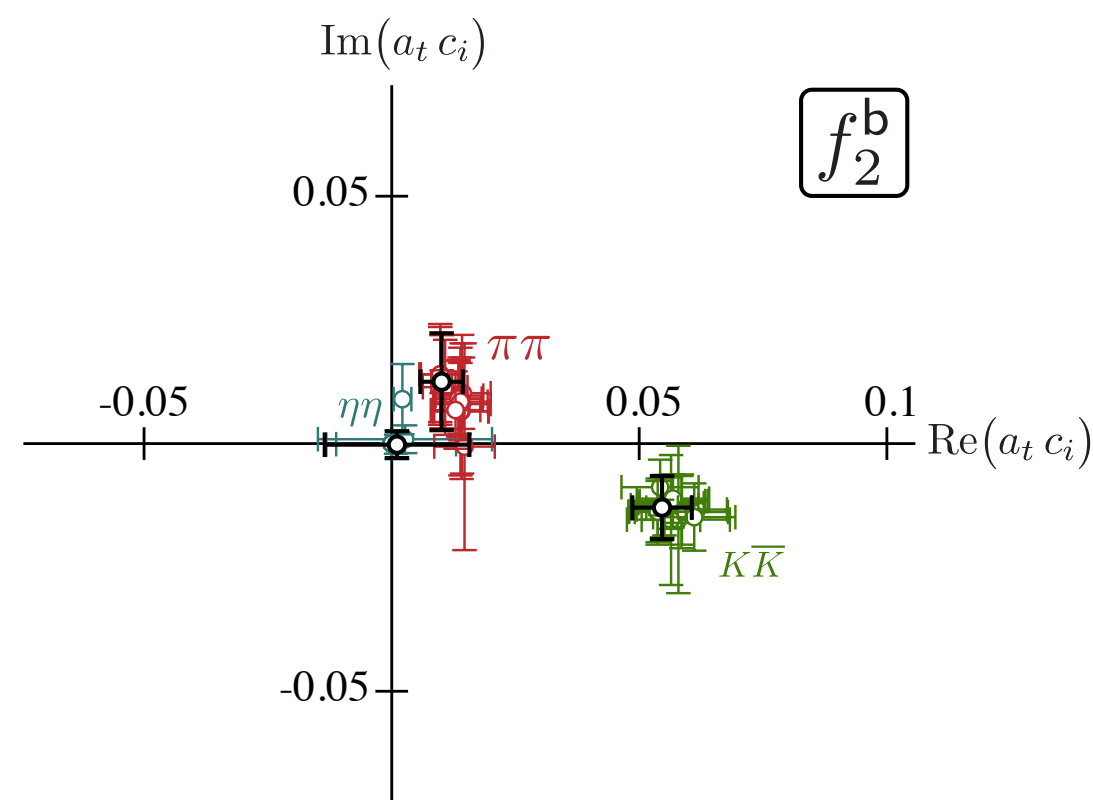
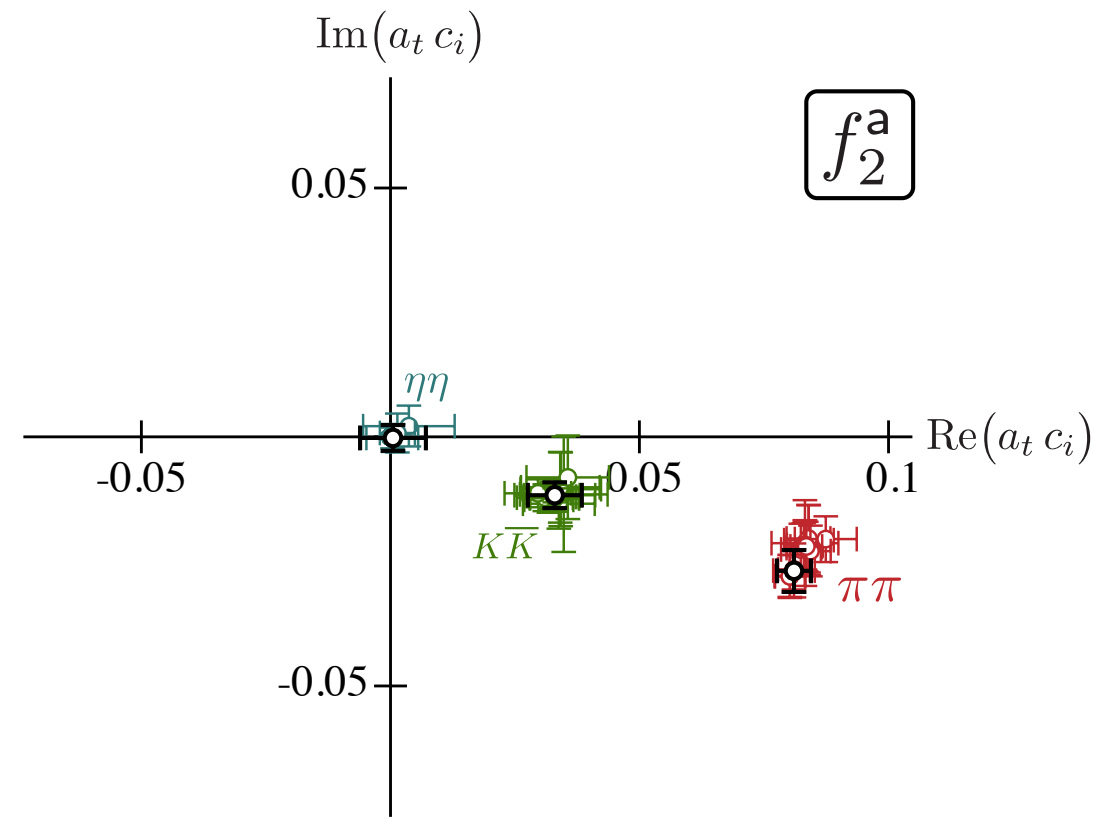
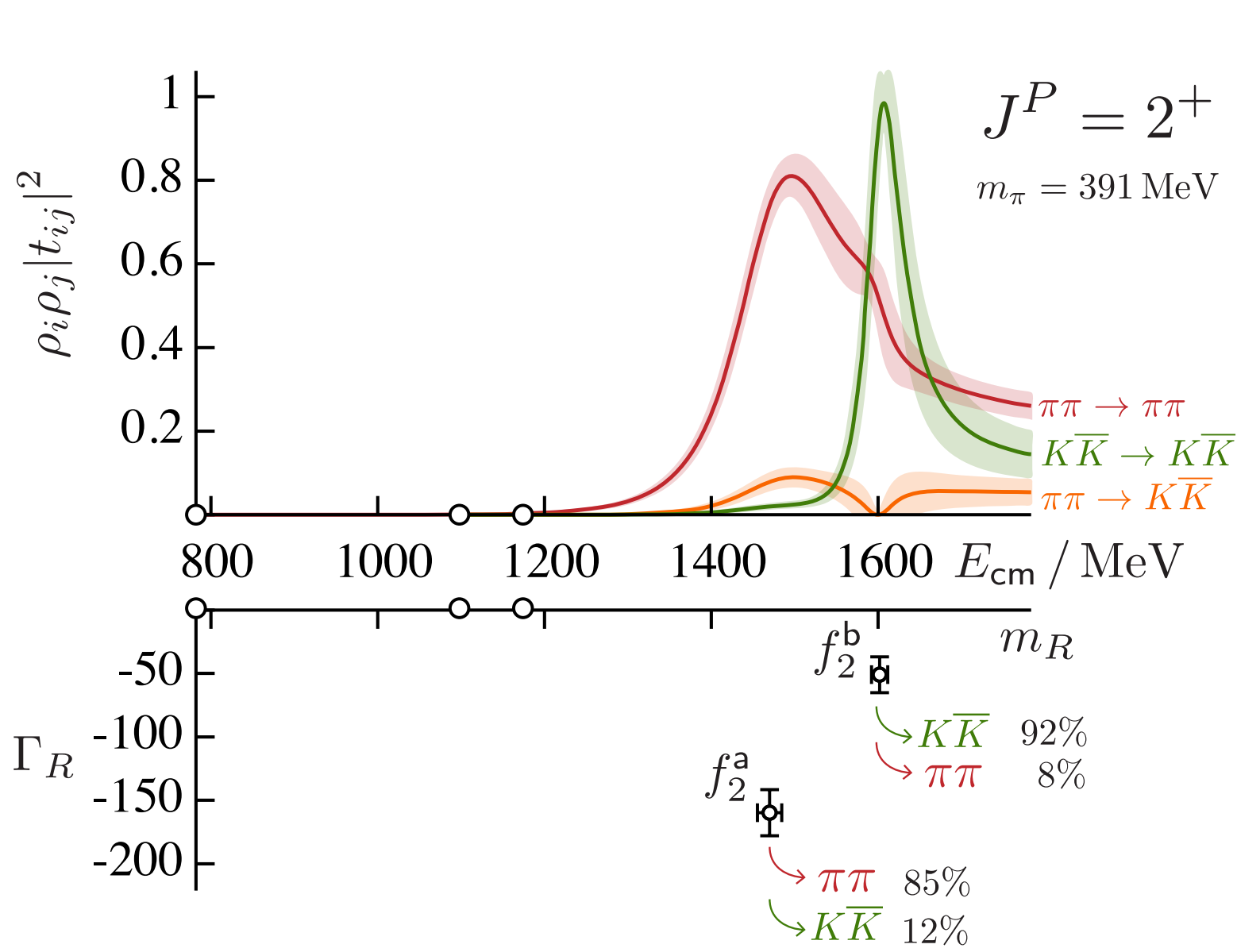
PHYSICAL REVIEW D 97, 054513 (2018)

Editors' Suggestion

Isoscalar $\pi\pi, K\bar{K}, \eta\eta$ scattering and the σ, f_0, f_2 mesons from QCD

Raul A. Briceño,^{1,2,*} Jozef J. Dudek,^{1,3,†} Robert G. Edwards,^{1,‡} and David J. Wilson^{4,§}

(for the Hadron Spectrum Collaboration)



PHYSICAL REVIEW D 97, 054513 (2018)

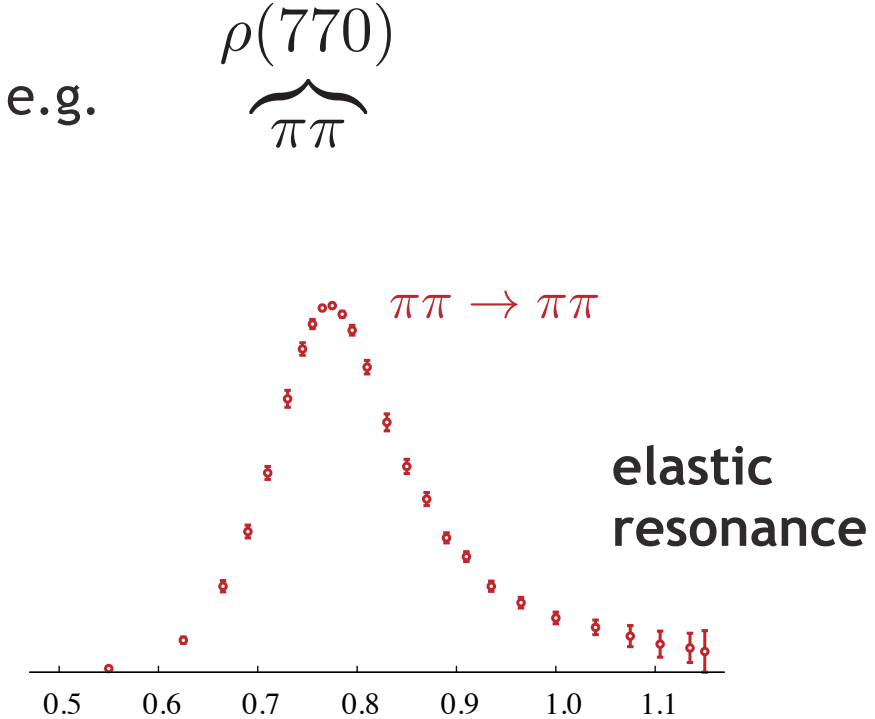
Editors' Suggestion

Isoscalar $\pi\pi, K\bar{K}, \eta\eta$ scattering and the σ, f_0, f_2 mesons from QCD

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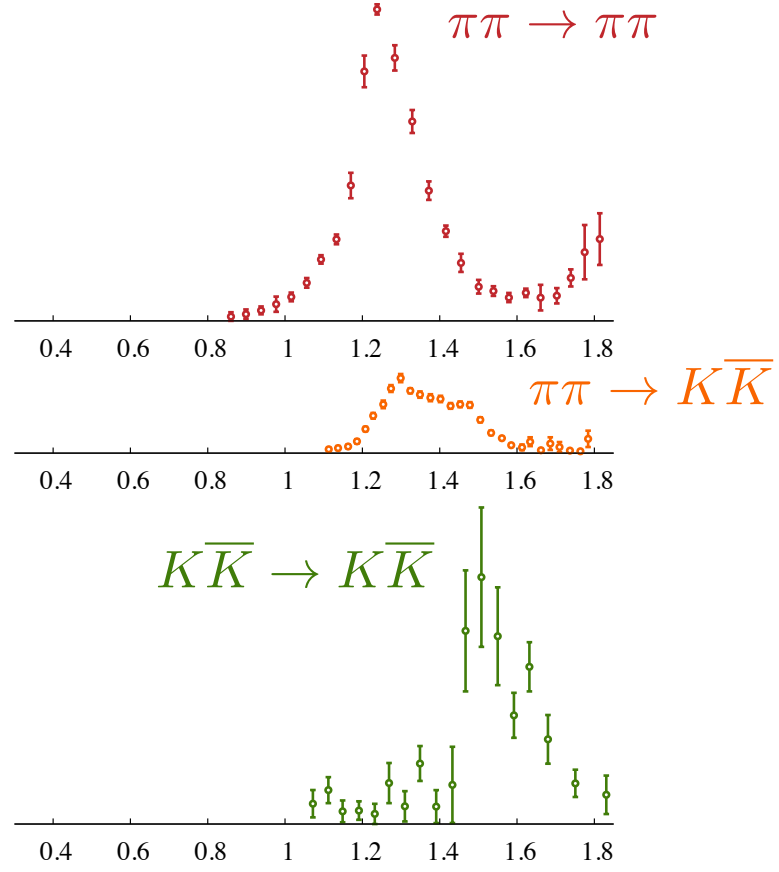
(for the Hadron Spectrum Collaboration)

most resonances can decay into more than one final state

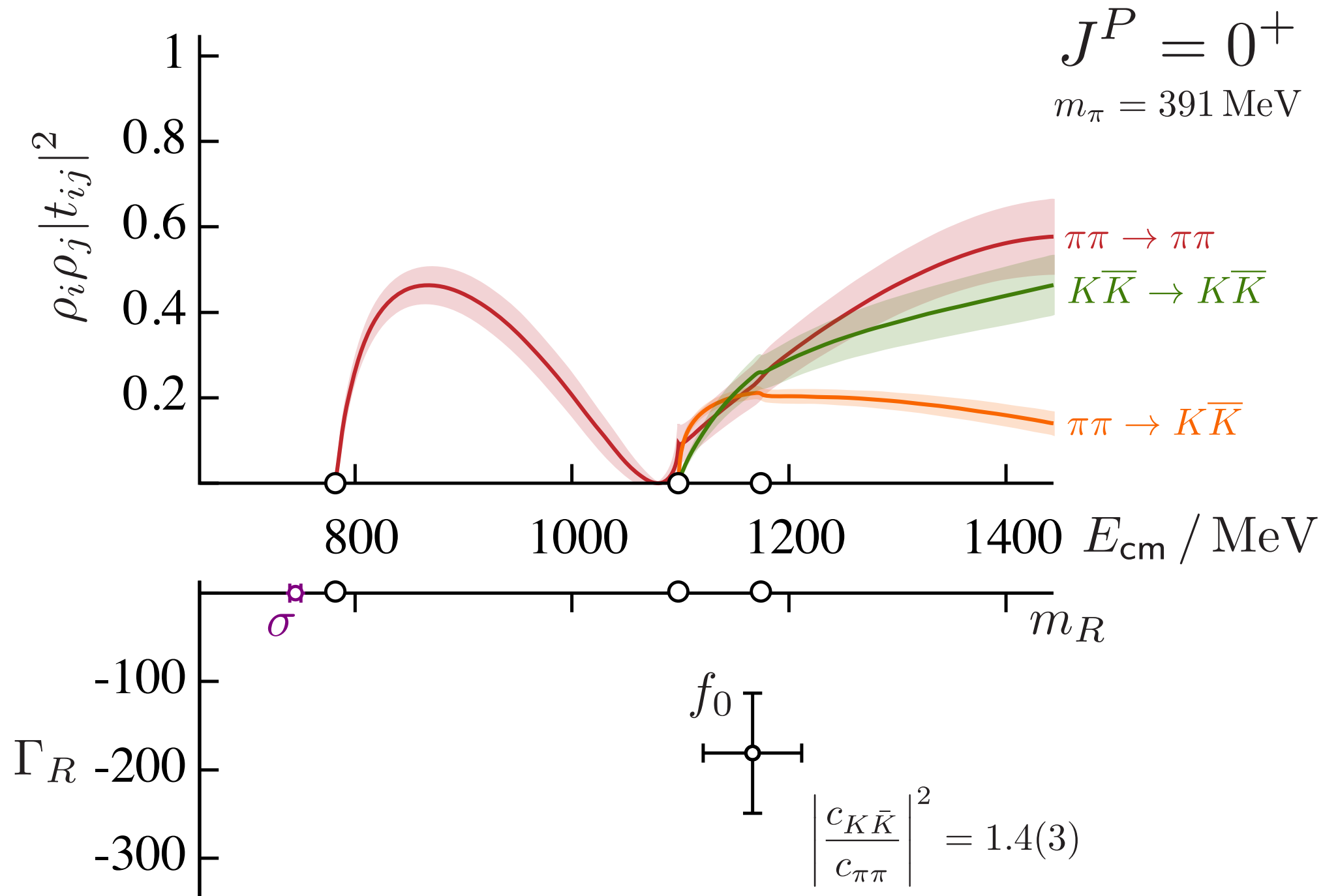


$f_2(1270)$
 $\pi\pi, \pi\pi\pi\pi, K\bar{K}, \eta\eta$

$f_2(1525)$
 $\pi\pi, K\bar{K}, \eta\eta$



coupled-channel resonances



PHYSICAL REVIEW D **97**, 054513 (2018)

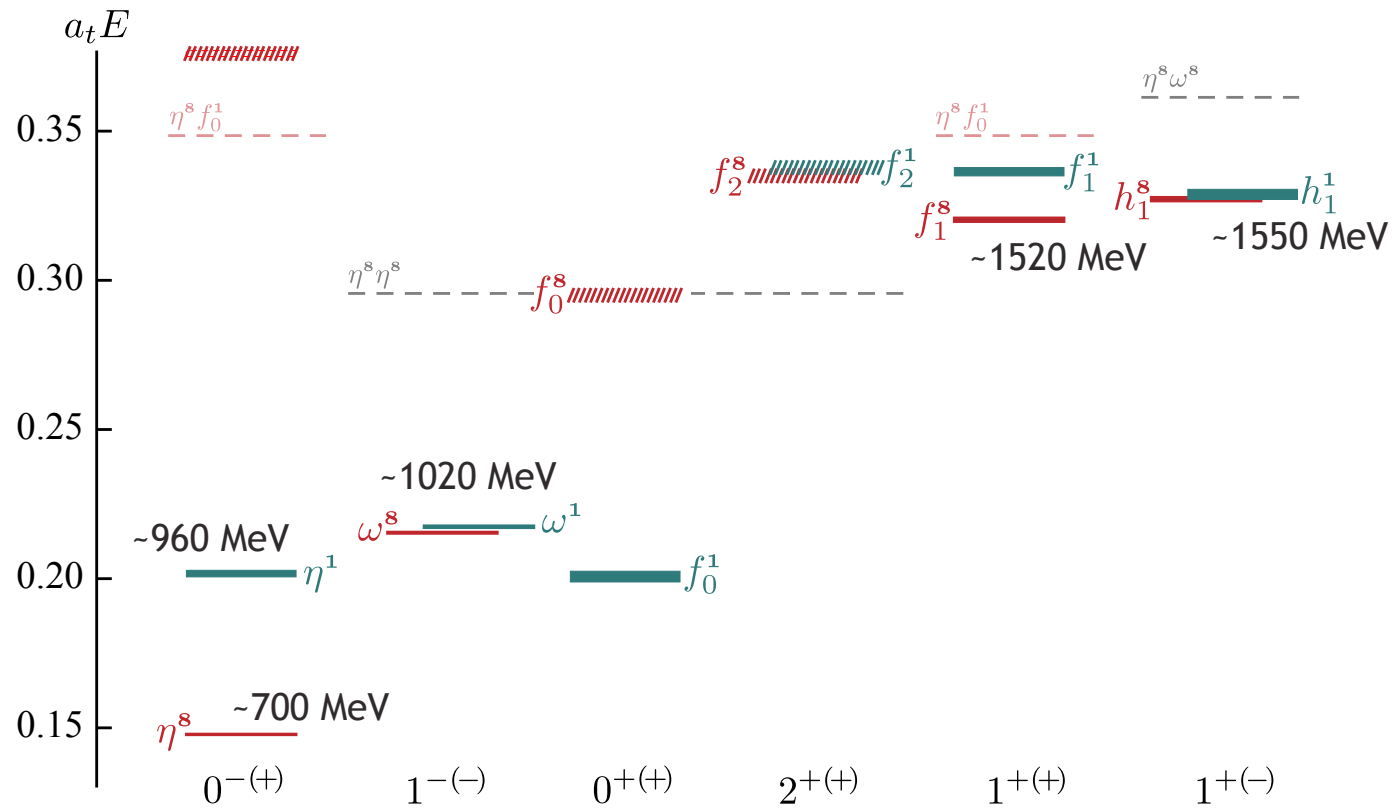
Editors' Suggestion

Isoscalar $\pi\pi, K\bar{K}, \eta\eta$ scattering and the σ, f_0, f_2 mesons from QCD

Raul A. Briceño,^{1,2,*} Jozef J. Dudek,^{1,3,†} Robert G. Edwards,^{1,‡} and David J. Wilson^{4,§}

(for the Hadron Spectrum Collaboration)

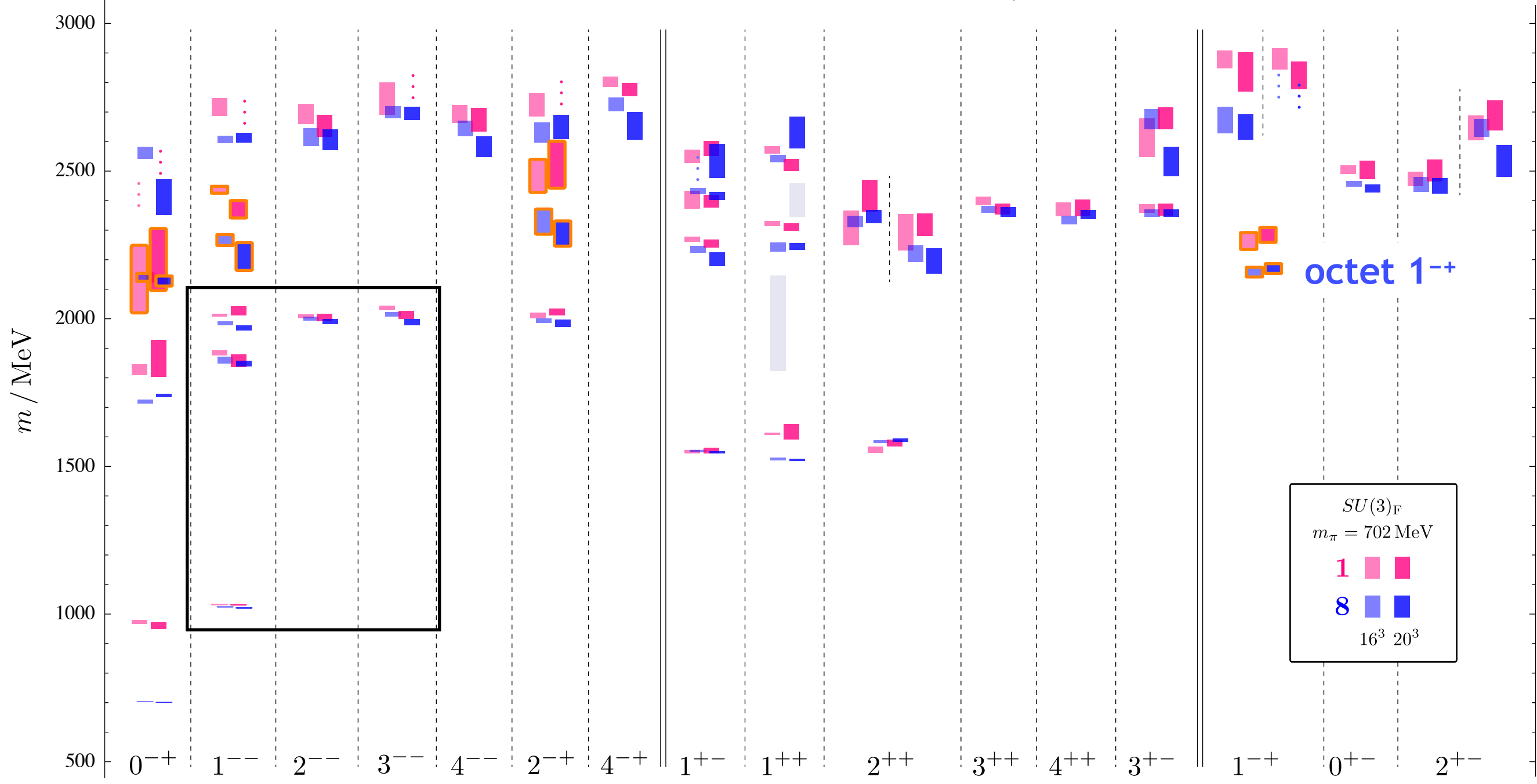
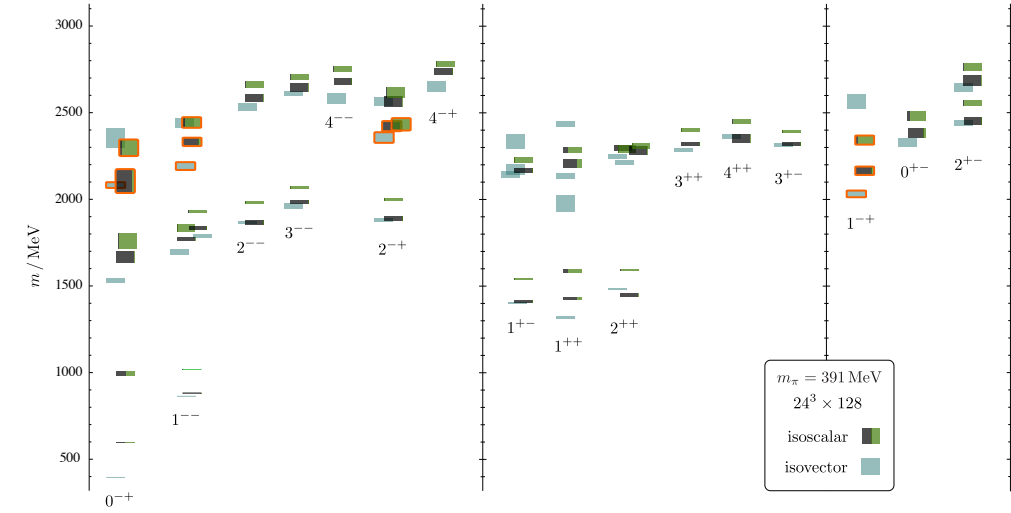
several stable mesons:



$m_u=m_d=m_s$ $SU(3)_F$ point

increase the light quark mass to the strange quark mass ...

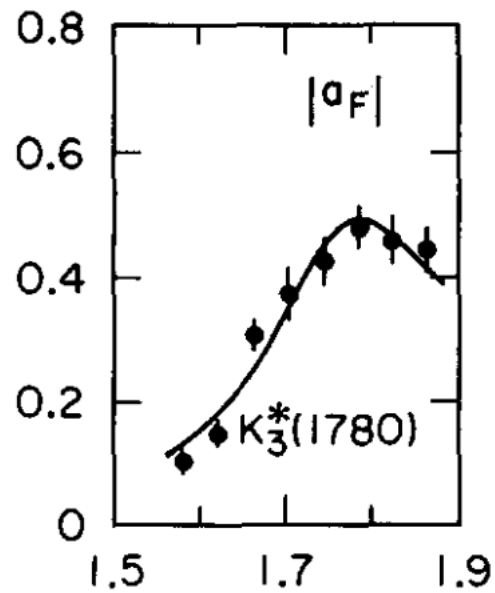
(incomplete) lattice spectrum calculation PRD 88 094505 (2013)



$K_3^*(1780)$

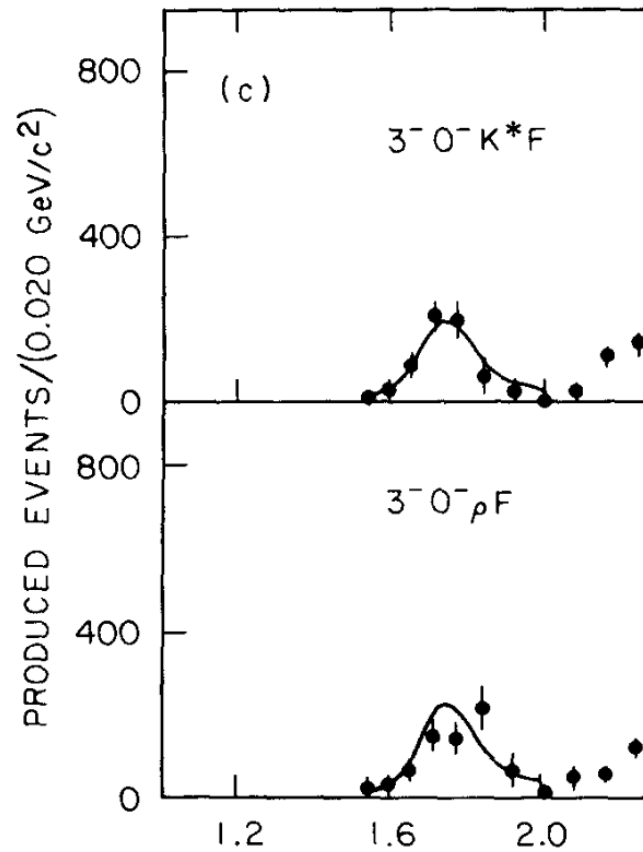
Mass $m = 1779 \pm 8$ MeV
Full width $\Gamma = 161 \pm 17$ MeV

$K\pi$



$\text{br}(K\pi) \sim 0.19(1)$

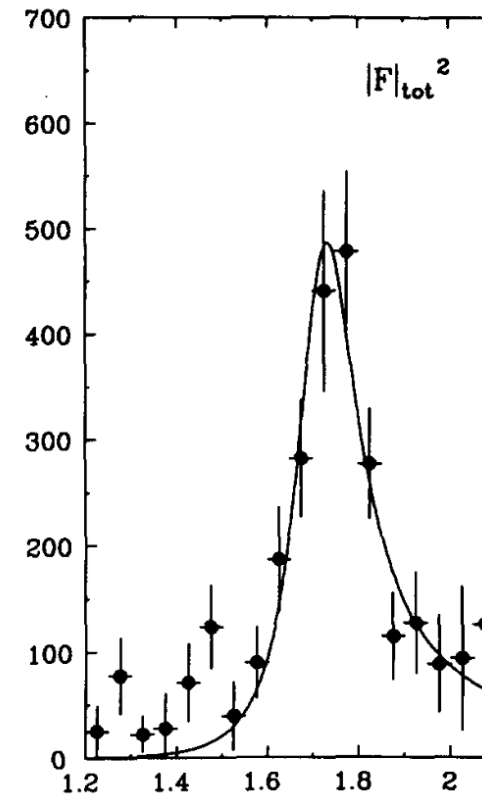
$K\pi\pi$



$$\frac{\Gamma(K\rho)}{\Gamma(K^*\pi)} \sim 1.5(2)$$

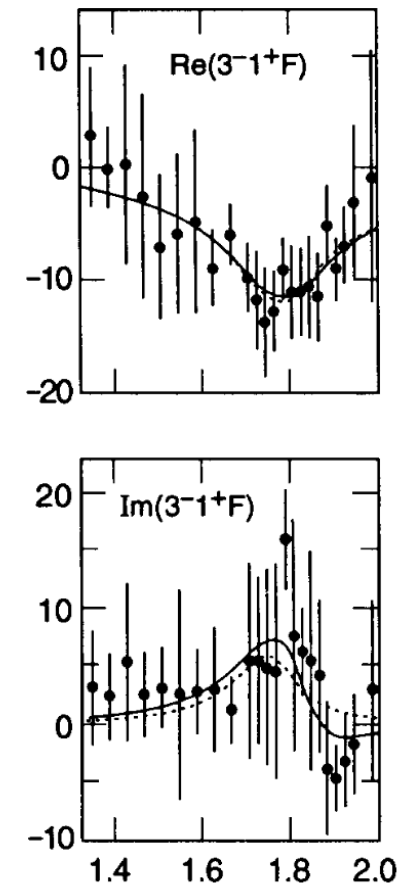
$$\frac{\Gamma(K_2^*\pi)}{\Gamma(K^*\pi)} \lesssim 0.8$$

$K\eta$



$$\frac{\Gamma(K\eta)}{\Gamma(K\pi)} \sim 0.5(2)$$

$K\omega$



$K_2(1820)$ [ff]

Mass $m = 1819 \pm 12$ MeV
Full width $\Gamma = 264 \pm 34$ MeV

$K_2(1770)$ [ff]

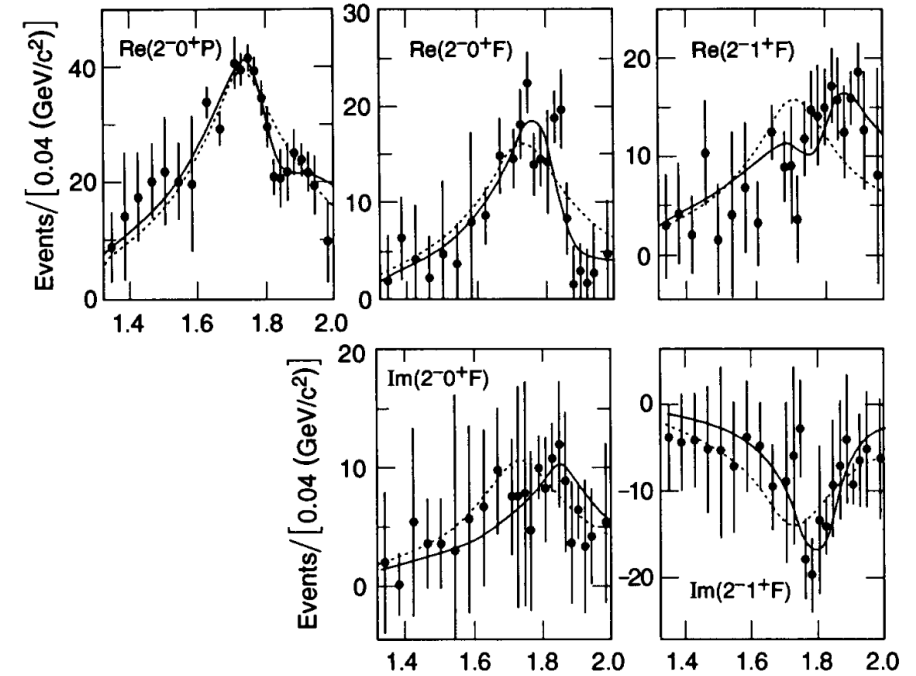
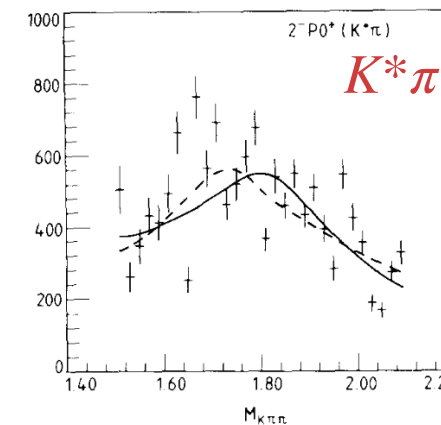
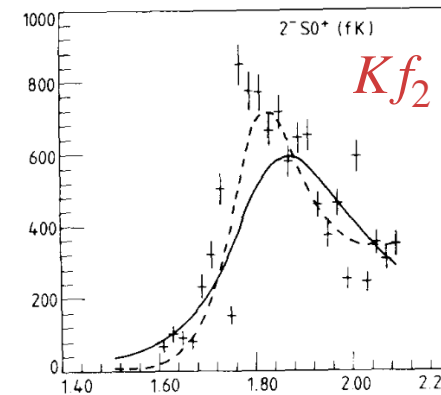
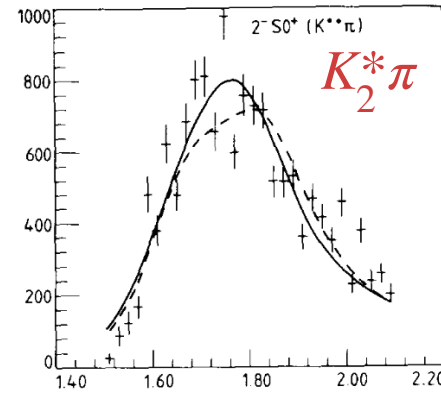
Mass $m = 1773 \pm 8$ MeV
Full width $\Gamma = 186 \pm 14$ MeV

$K\pi\pi$

2^- not seen in LASS
(dominance of pion exchange?)

$K\omega$

ACCMOR 1981
(higher beam energy)
 $K^\pm p \rightarrow K^\pm \pi^+ \pi^- p$



slight χ^2 improvement here is the main evidence for two 2^- resonances

ACCMOR data suggests significant branches into $K_2^* \pi$, Kf_2 which are closed at the SU(3) flavor point

(A signal in $2^- P(\rho K)$ is barely significant.)

	Mass (GeV)	Width (MeV)	Approximate branching ratios		
			$K^{**}\pi$	fK	$K^*\pi$
One resonance	1.82	~ 200	0.6	0.16	0.24
Two resonances	1.78	~ 210	0.03	0.74	0.23
	1.84	~ 230	0.77	0.18	0.05

COMPASS data to make two 2^- states firm?

$K^*(1680)$

Mass $m = 1718 \pm 18$ MeV
Full width $\Gamma = 322 \pm 110$ MeV

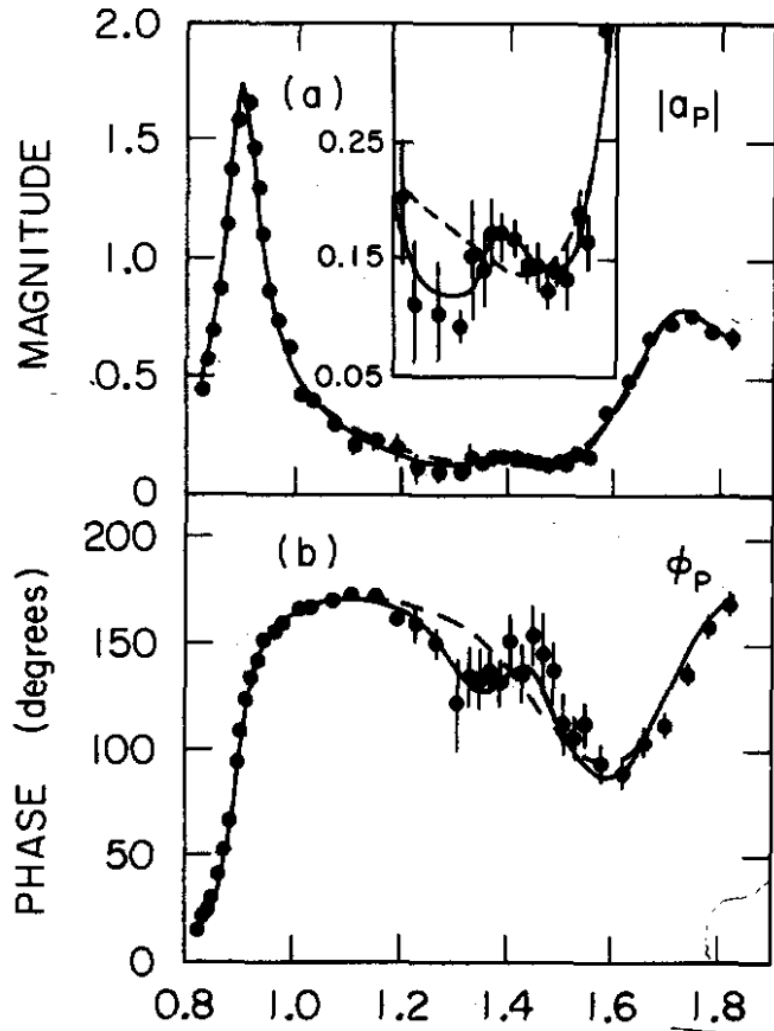
205(16)(34) in $K\pi$, 423(18)(30) in $K\pi\pi$ compatible ?

$K^*(1410)$

Mass $m = 1414 \pm 15$ MeV
Full width $\Gamma = 232 \pm 21$ MeV

176(52)(22) in $K\pi$, 240(18)(12) in $K\pi\pi$

$K\pi$

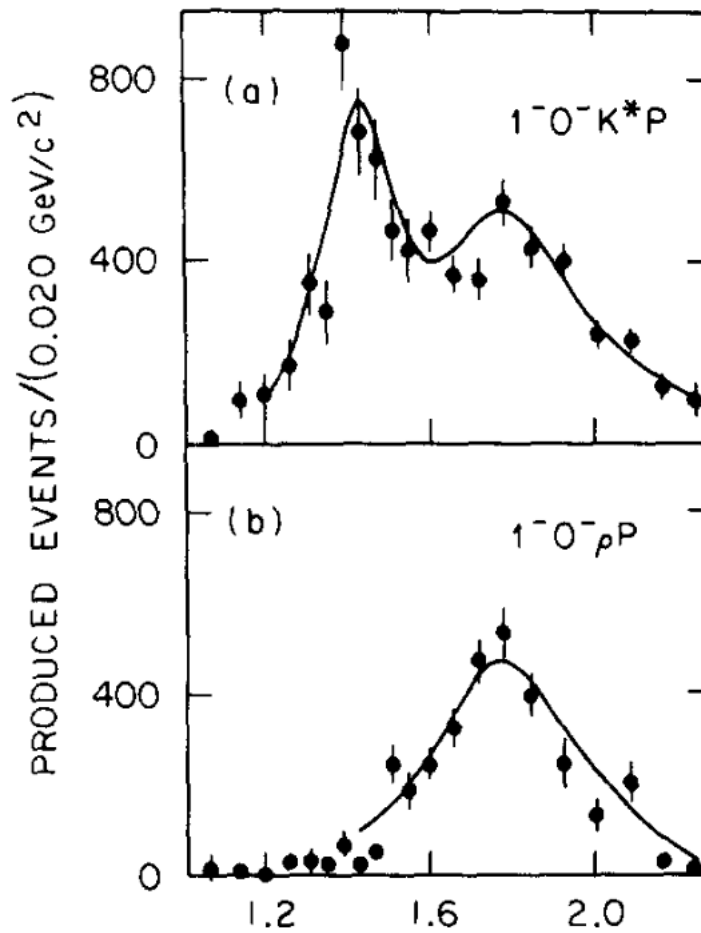


very weak contribution from $K^*(1410)$

$\text{br}(K^*(1410) \rightarrow K\pi) \sim 0.07(1)$

$\text{br}(K^*(1680) \rightarrow K\pi) \sim 0.39(3)$

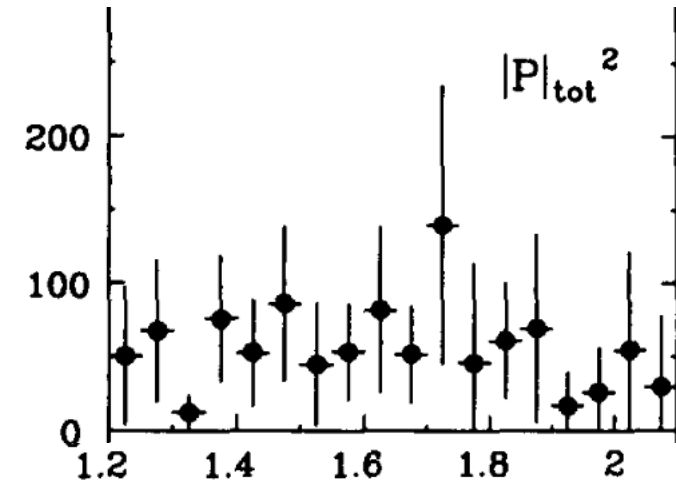
$K\pi\pi$



no $K^*(1410)$ in $K\rho$?

$$\frac{K^*(1680) \rightarrow \Gamma(K\rho)}{K^*(1680) \rightarrow \Gamma(K^*\pi)} \sim 0.97(10)$$

$K\eta$



no 1^- signal

$K\omega$

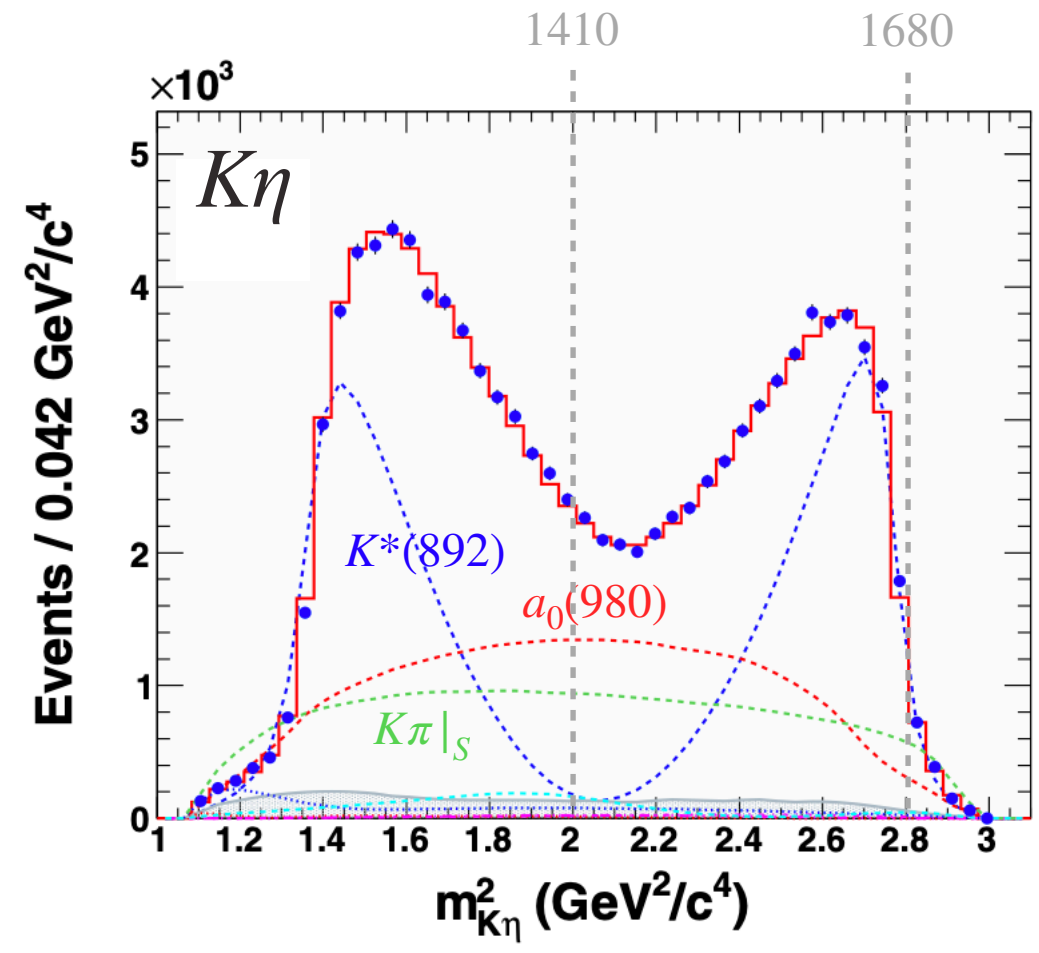
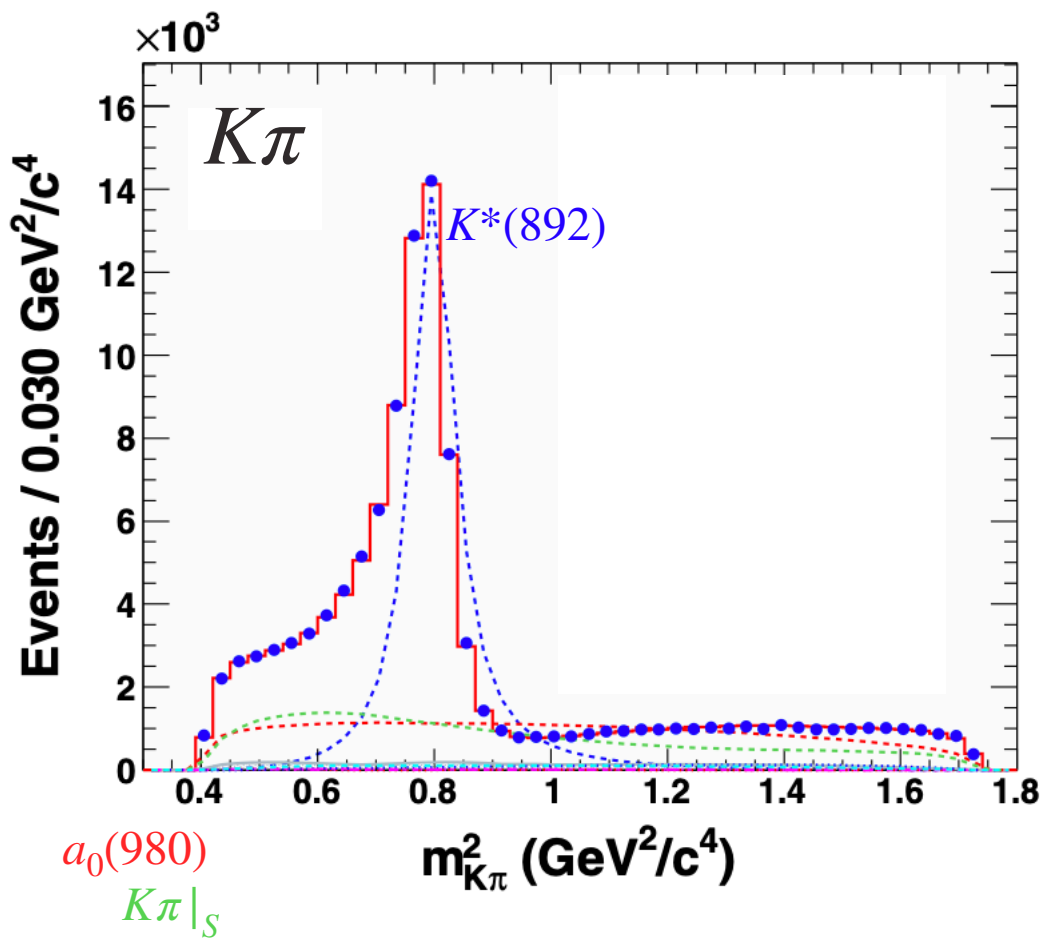
ambiguous 1^- wave
(only in a thesis)

ACCMOR 1981
(higher beam energy)

$K^\pm p \rightarrow K^\pm \pi^+ \pi^- p$

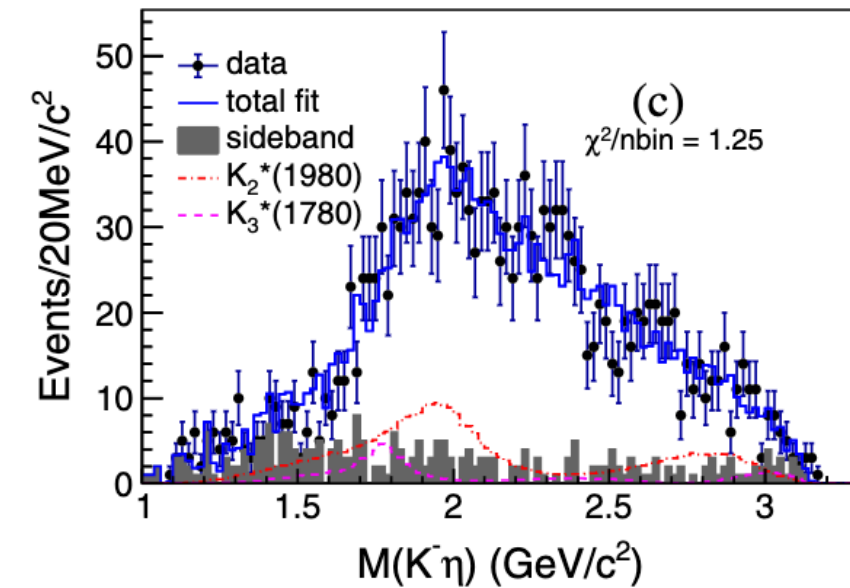
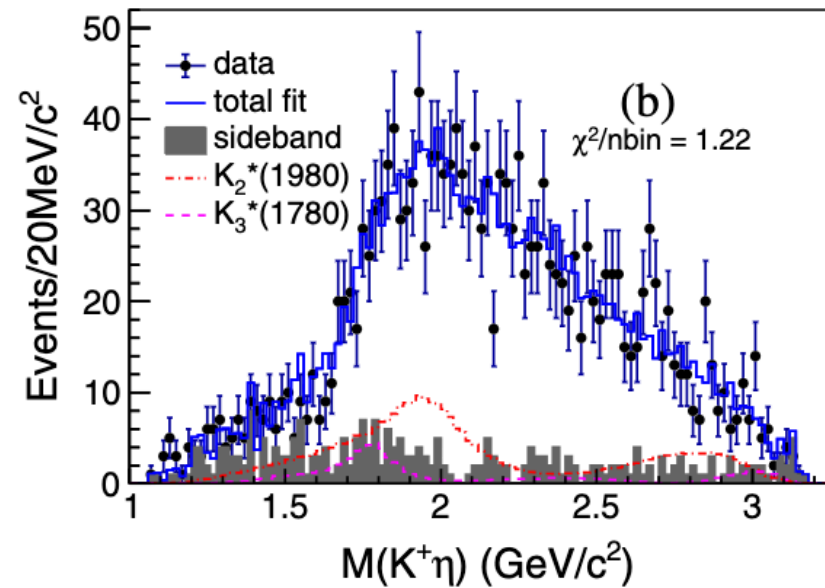
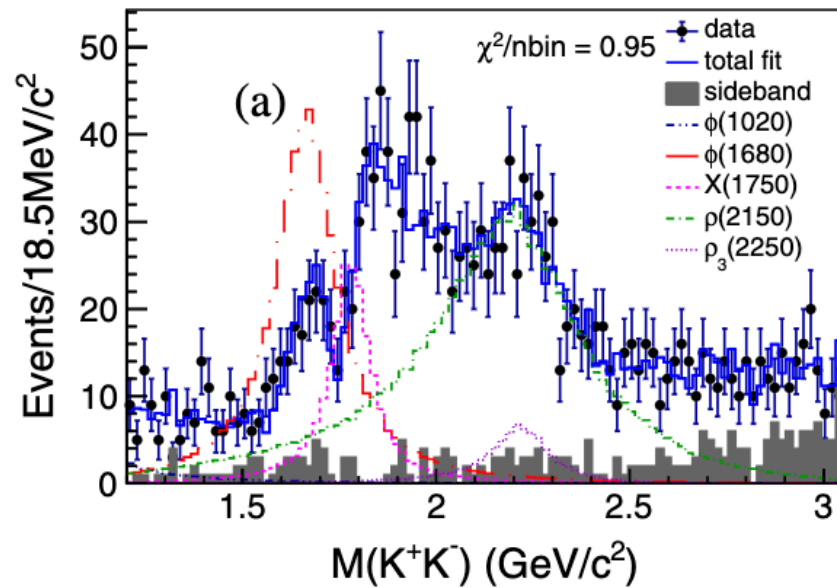
has no 1^- waves

Belle Dalitz plot analysis of $D^0 \rightarrow K^- \pi^+ \eta$



other contributions smaller than combinatoric background

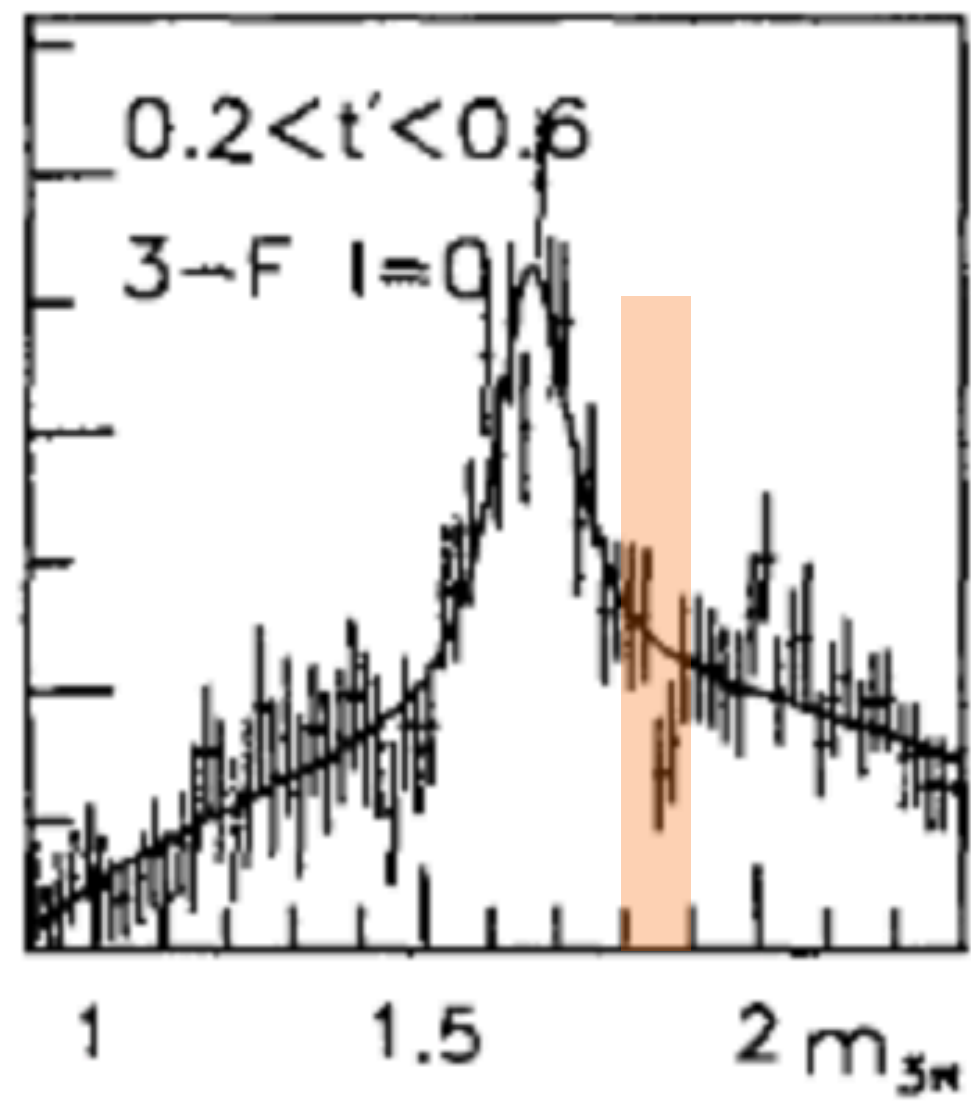
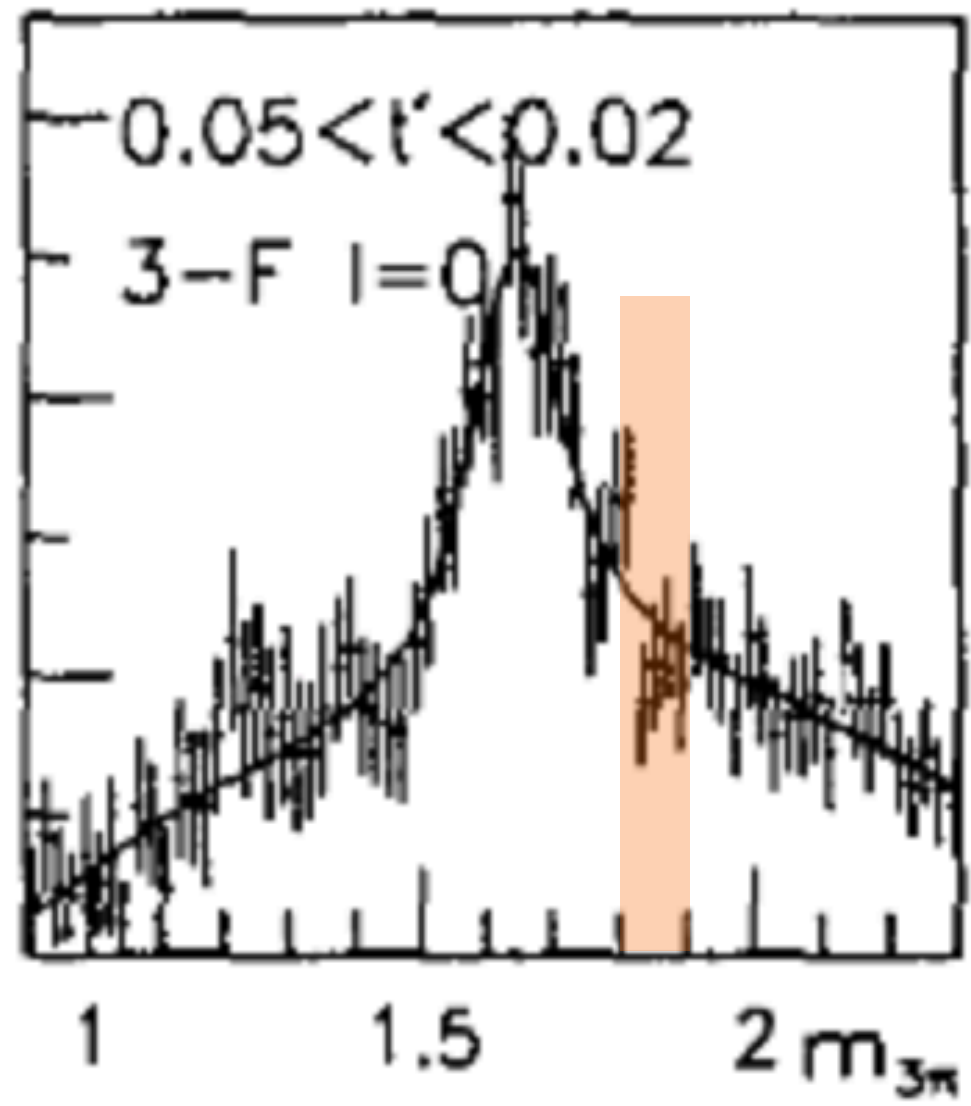
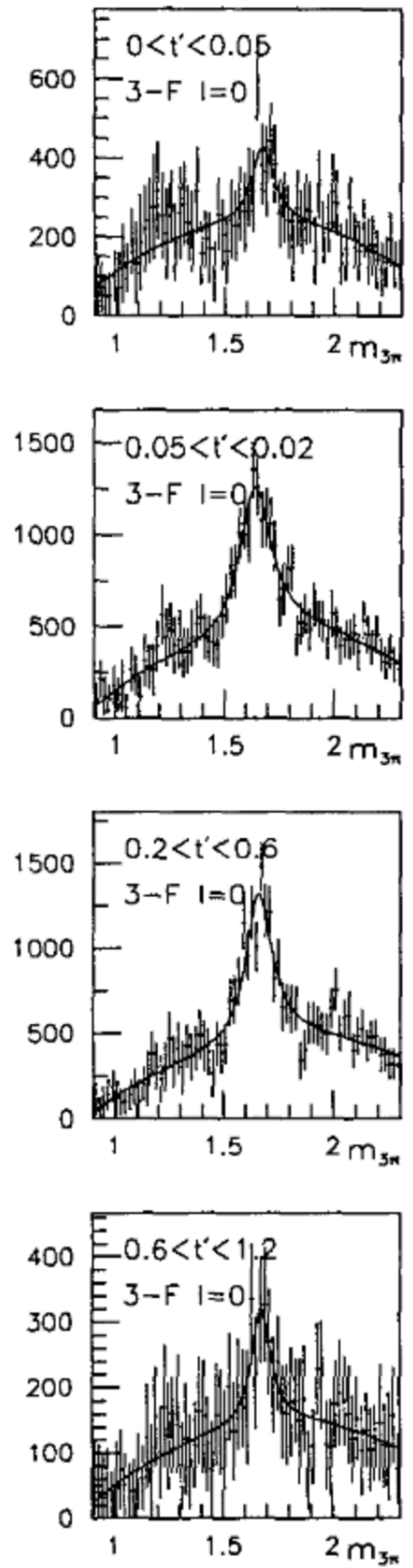
BES Dalitz plot analysis of $\psi' \rightarrow K^- K^+ \eta$



claim of rather small $K_3^*(1780)$ contributing in $K\eta$

$\omega_3(1670)$ in $\pi\rho$ (and $\phi_3(1850)$?)

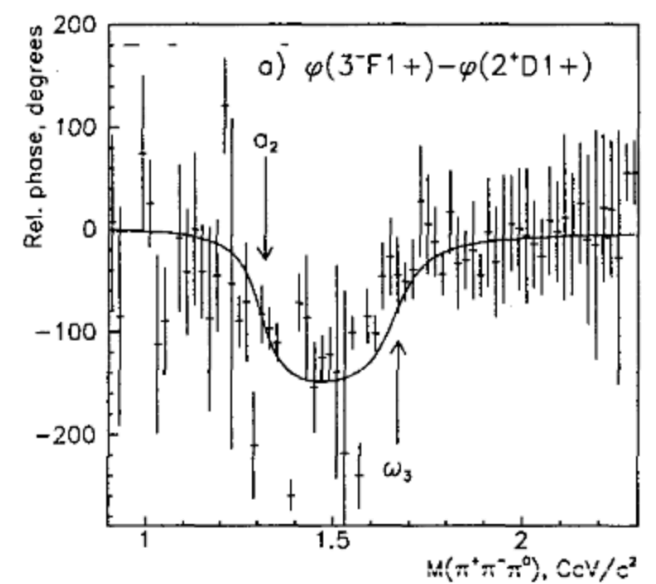
VES $\pi^-p \rightarrow \pi^+\pi^-\pi^0n$ @ 36 GeV only 'modern' high statistics evidence for this state

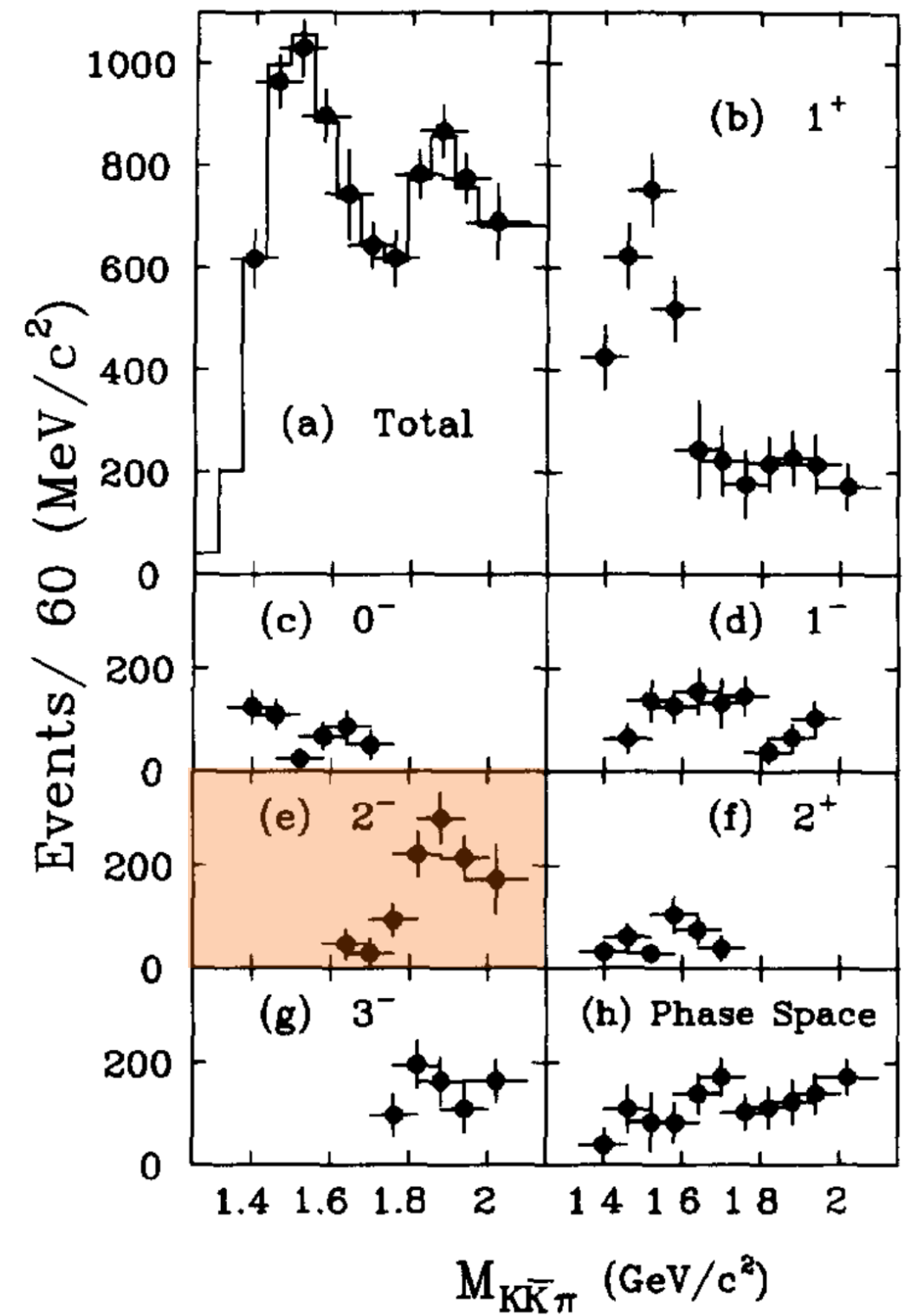


could there be a weak $\phi_3(1850)$ present in $\pi\rho$?

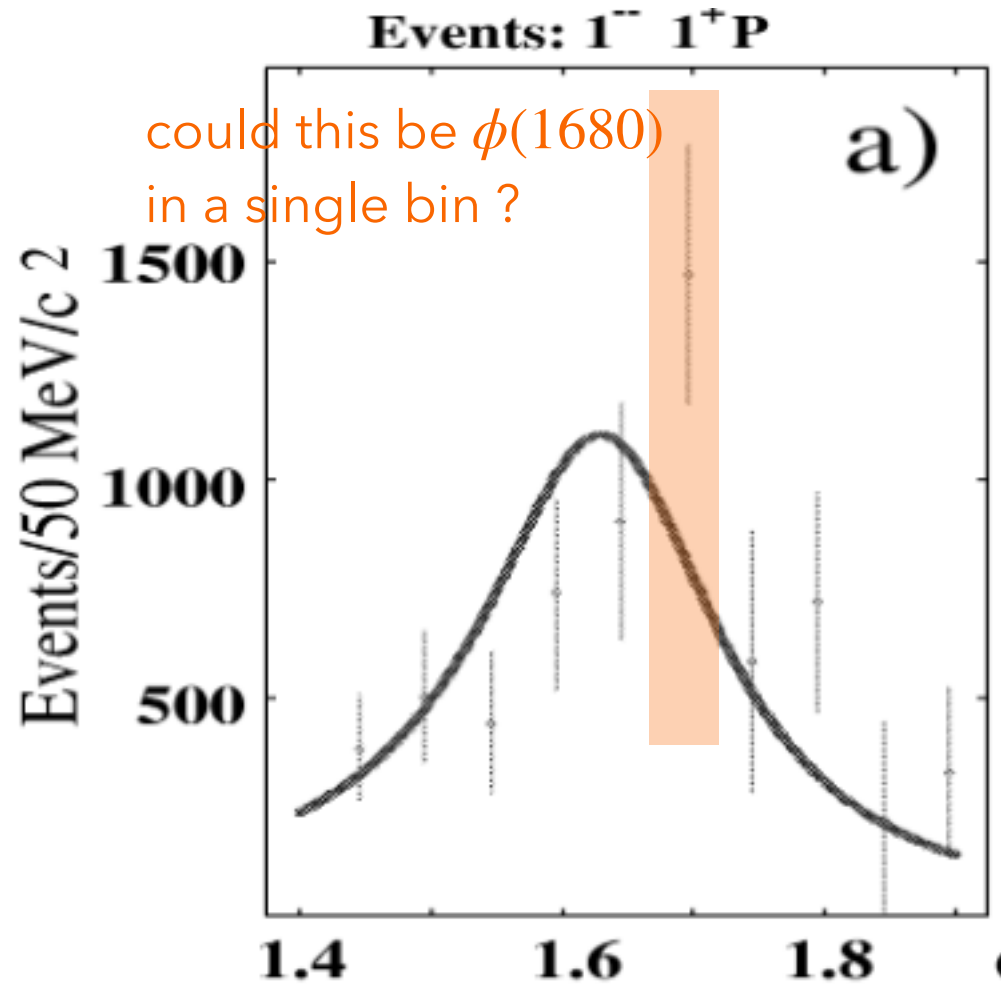
2^+ wave fit only has $a_2(1320)$ and not any $a_2(1700)$
so unclear what to conclude from this good fit to the phase difference

no sharp feature in the phase, but the ϕ_3 coupling could be small

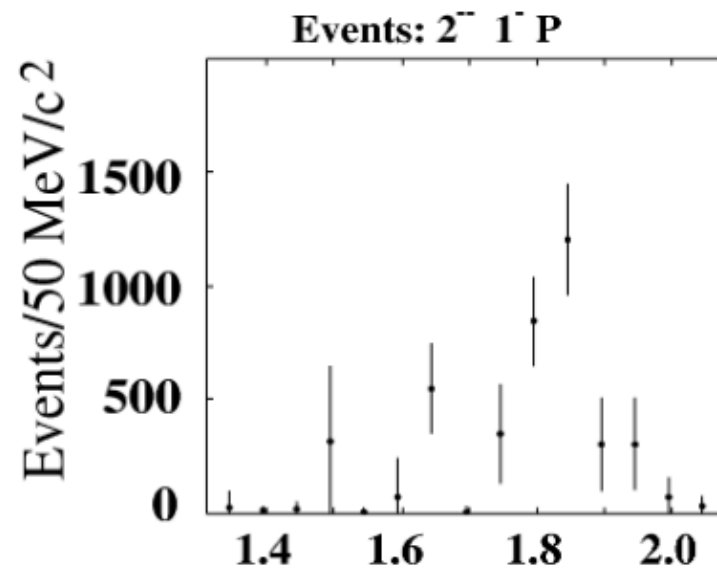


LASS $K^-p \rightarrow K_S^0 K^\pm \pi^\mp \Lambda$ 

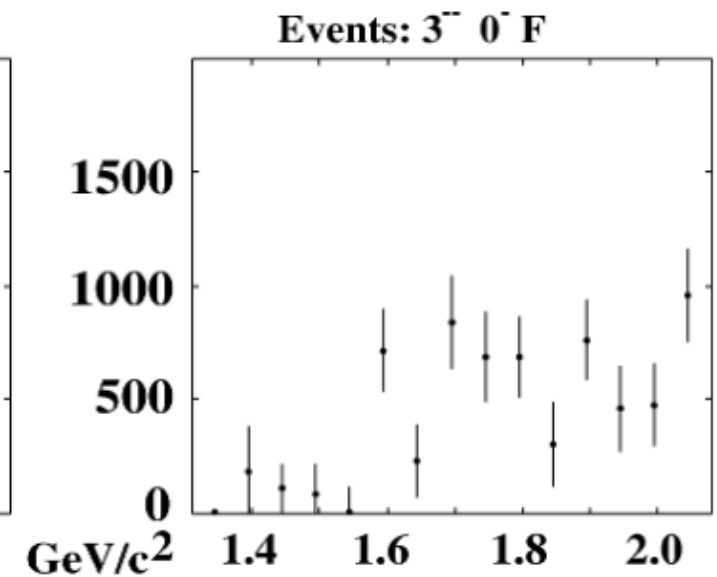
could this be a $\phi_2(\sim 1900)$ with $\Gamma \sim 200$?



$\omega(1650)$ with fixed PDG (in 2001) params
 $m = 1649(24), \Gamma = 220(35)$



could this be a $\phi_2(1850)$?



nothing here

$\omega(1650) [k]$

Mass $m = 1670 \pm 30 \text{ MeV} [h]$
 Full width $\Gamma = 315 \pm 35 \text{ MeV} [h]$

can access $I = 1, I_z = 0 (\rho)$ and $I = 0, I_z = 0 (\omega, \phi)$

Measurement of the $e^+e^- \rightarrow \pi^+\pi^-\pi^0$ and $e^+e^- \rightarrow \omega\pi^+\pi^-$ reactions in the energy interval 1350–2400 MeV

DM2 Collaboration

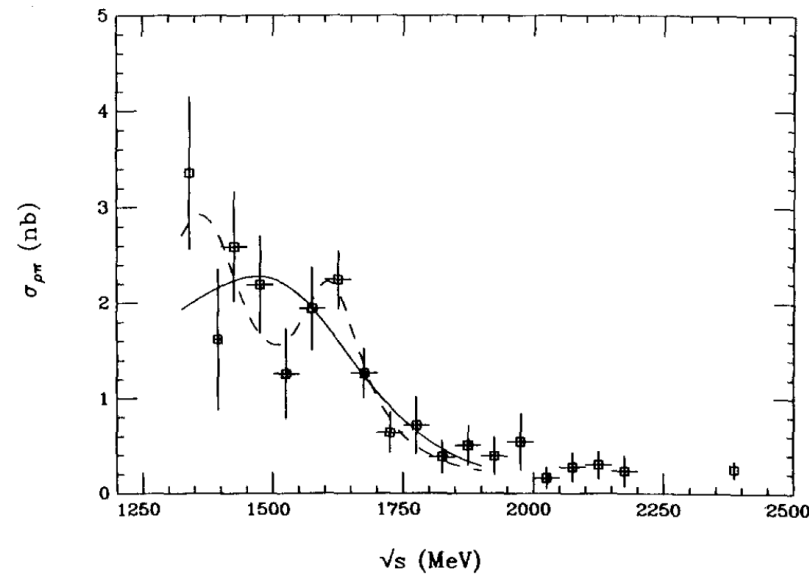
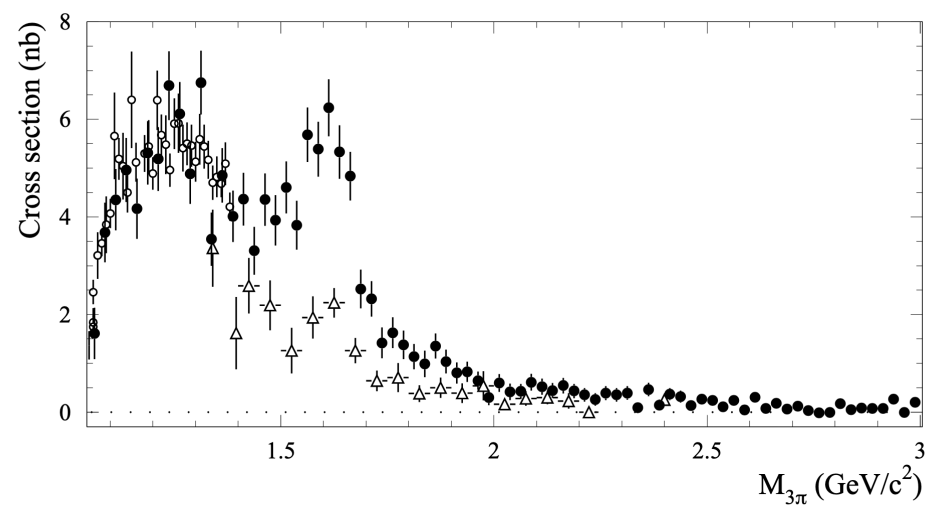


Fig. 3. $e^+e^- \rightarrow \rho\pi$ cross section. The solid and dotted lines refer respectively to the fits to a single or two Breit Wigner functions

$$G = - \implies I = 0$$

in principle expect **four** resonances could contribute here

Study of $e^+e^- \rightarrow \pi^+\pi^-\pi^0$ process using initial state radiation with *BABAR*



appears to be a significant tension between BaBar cross-section and DM2 above 1.5 GeV

FIG. 18: The $e^+e^- \rightarrow \pi^+\pi^-\pi^0$ cross section measured in this work (filled circles), by SND (open circles), and DM2 (open triangles).

The $e^+e^- \rightarrow 2(\pi^+\pi^-)\pi^0, 2(\pi^+\pi^-)\eta, K^+K^-\pi^+\pi^-\pi^0$ and $K^+K^-\pi^+\pi^-\eta$ cross sections measured with initial-state radiation

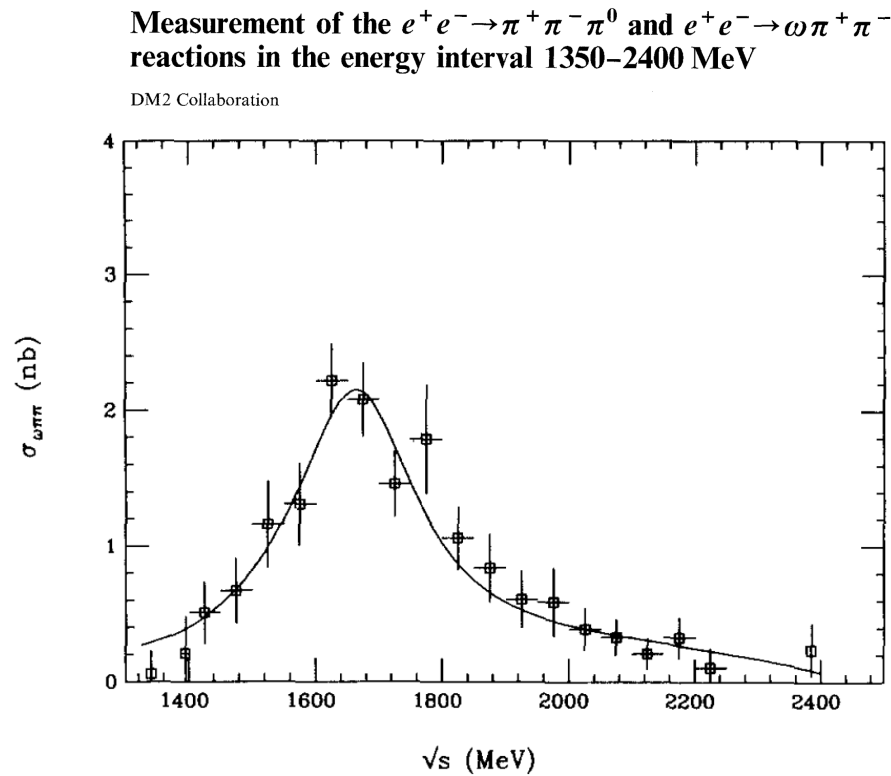


Fig. 6. $e^+e^- \rightarrow \omega\pi^+\pi^-$ cross section. The line refers to the single Breit Wigner fit

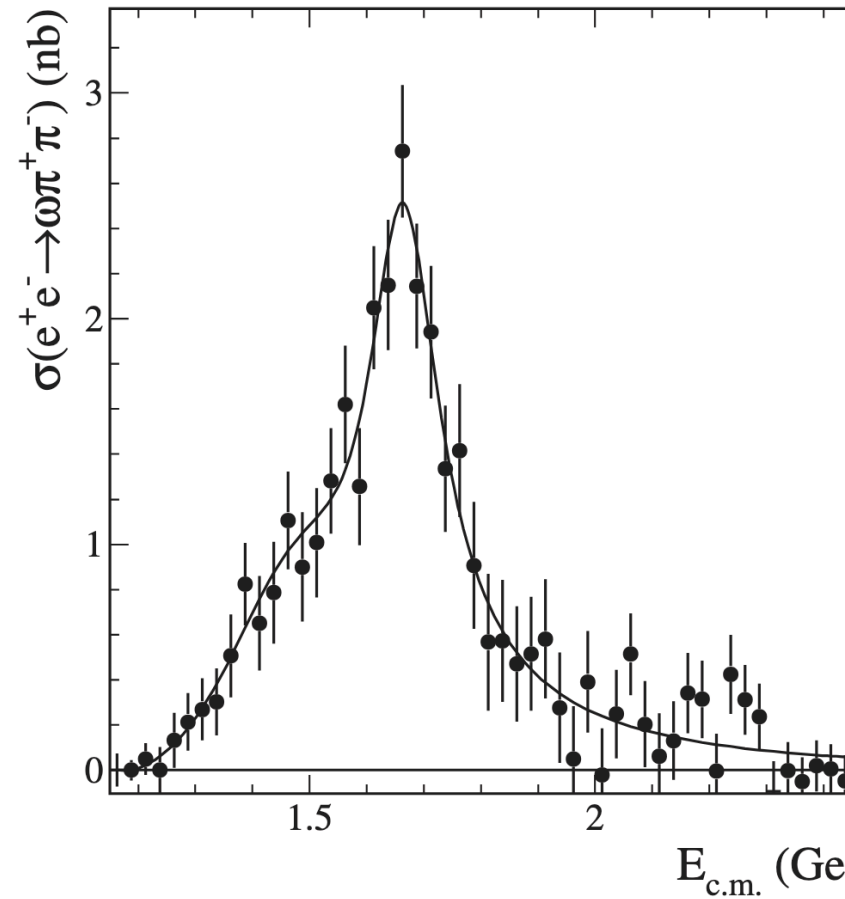


FIG. 16. The $e^+e^- \rightarrow \omega\pi^+\pi^-$ cross section excluding $\omega f_0(980)$ contribution (points). The curve shows the result the fit of the $\omega(1420)$ and $\omega(1650)$ resonances described in text (Fit 3 in Table VI).

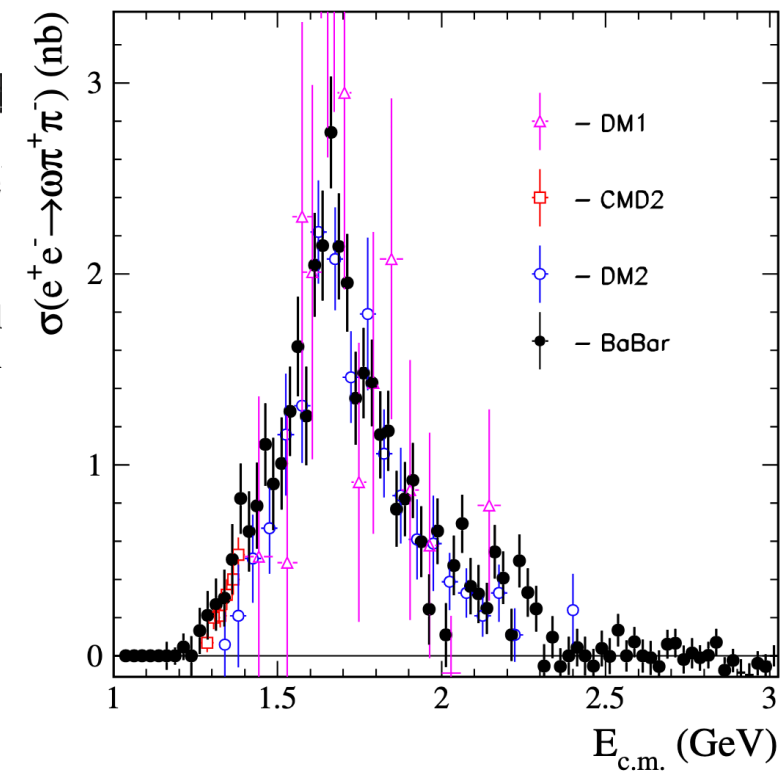


FIG. 14 (color online). The $e^+e^- \rightarrow \omega\pi^+\pi^-$ cross section as a function of c.m. energy obtained via ISR at BABAR. The direct measurements from DM1, DM2, and CMD2 are also shown. Only statistical errors are shown.

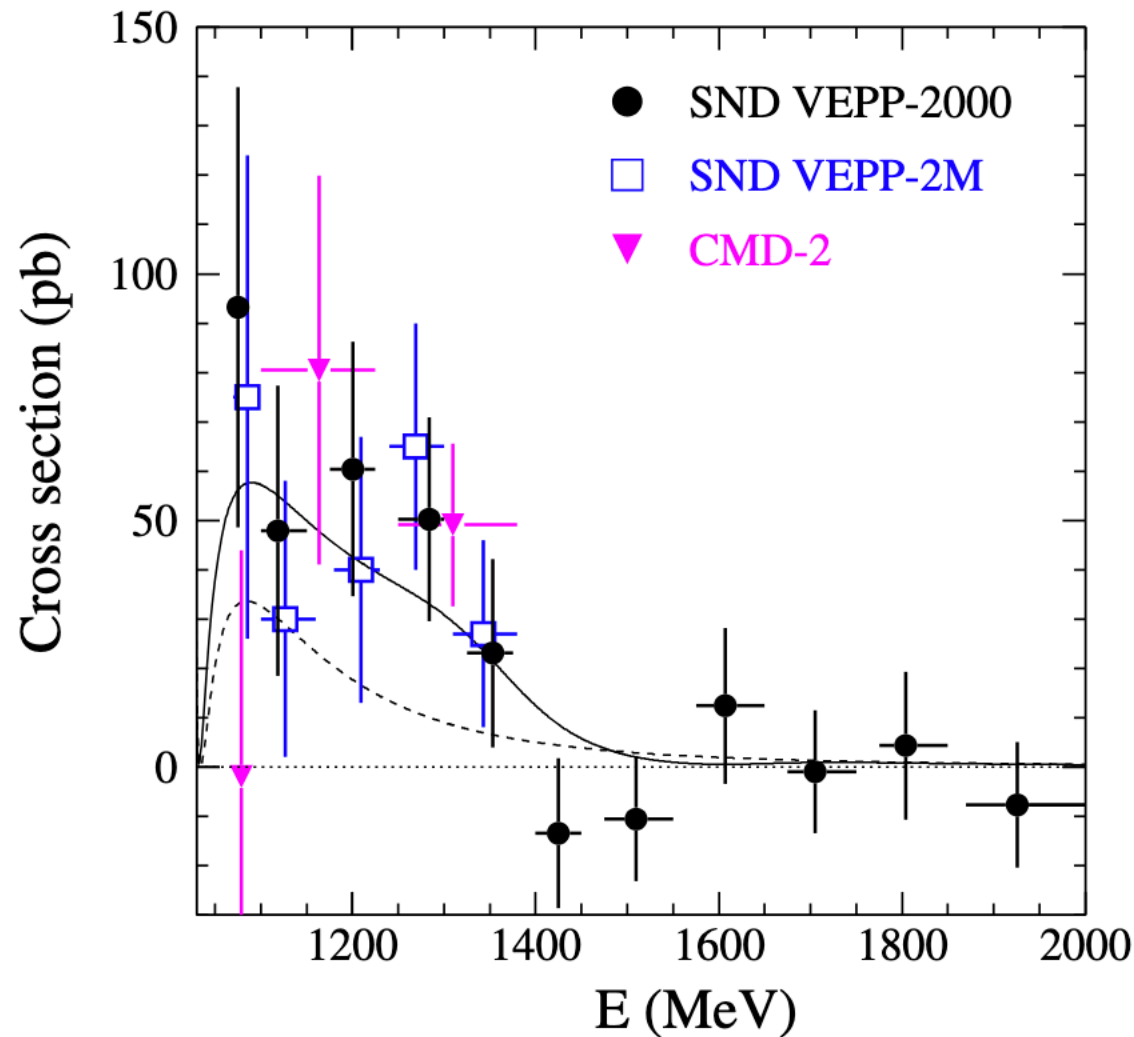
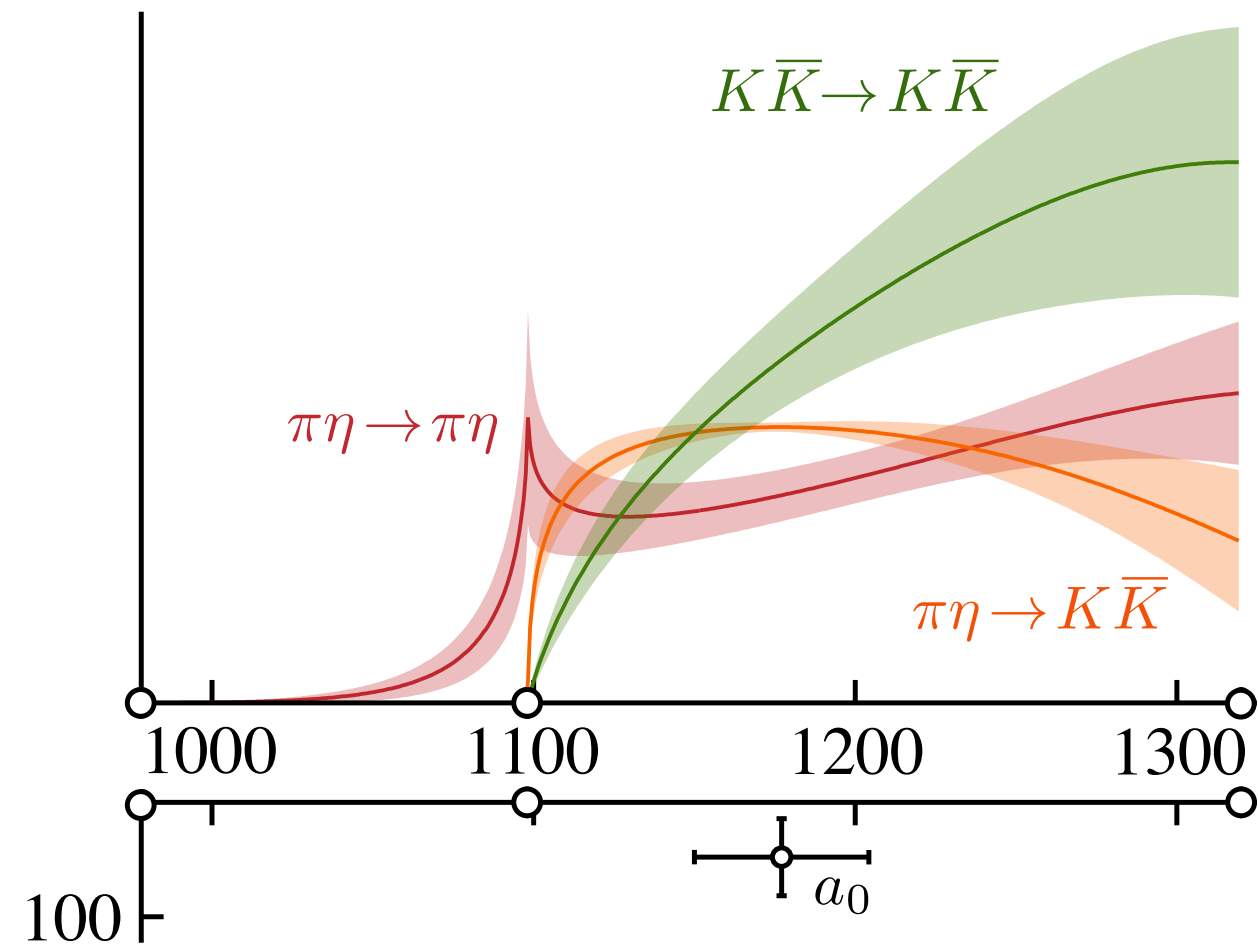
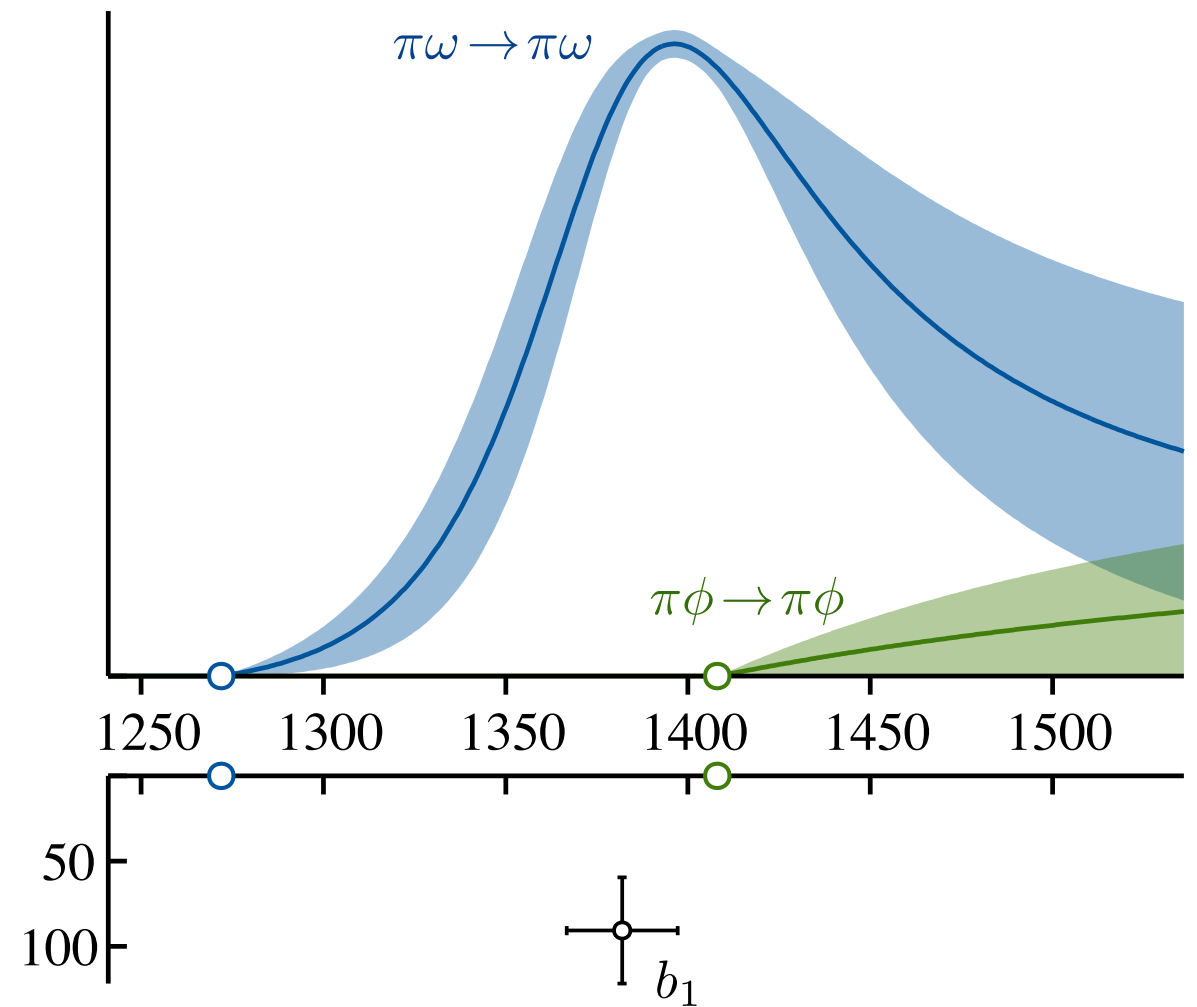


FIG. 5. The $e^+e^- \rightarrow \pi^0\gamma$ cross section measured in this work (SND@VEPP-2000) in comparison with the previous measurements: SND@VEPP-2M [4], and CMD-2 [5]. The solid curve is the result of the VMD fit with two excited resonances, V' and V'' . The dashed curve represents the result of the fit with $\sigma_{V'} = \sigma_{V''} = 0$.

$$J^P = 0^+ \quad I^G = 1^- \quad (\eta\pi, K\bar{K})$$



$$J^P = 1^+ \quad I^G = 1^- \quad (\pi\omega, \pi\phi)$$

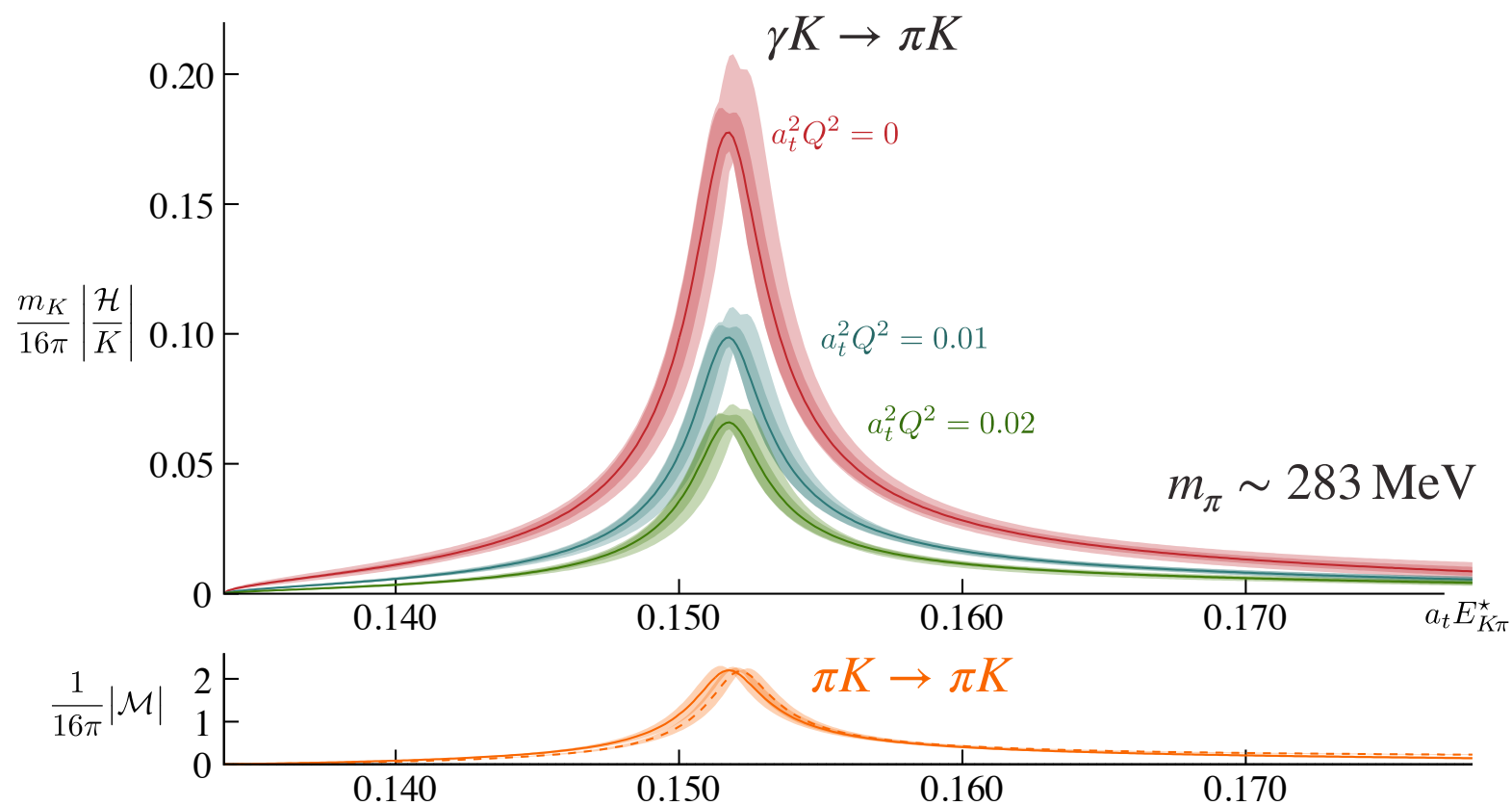


production amplitudes

can also couple to an external current, e.g. electroweak production

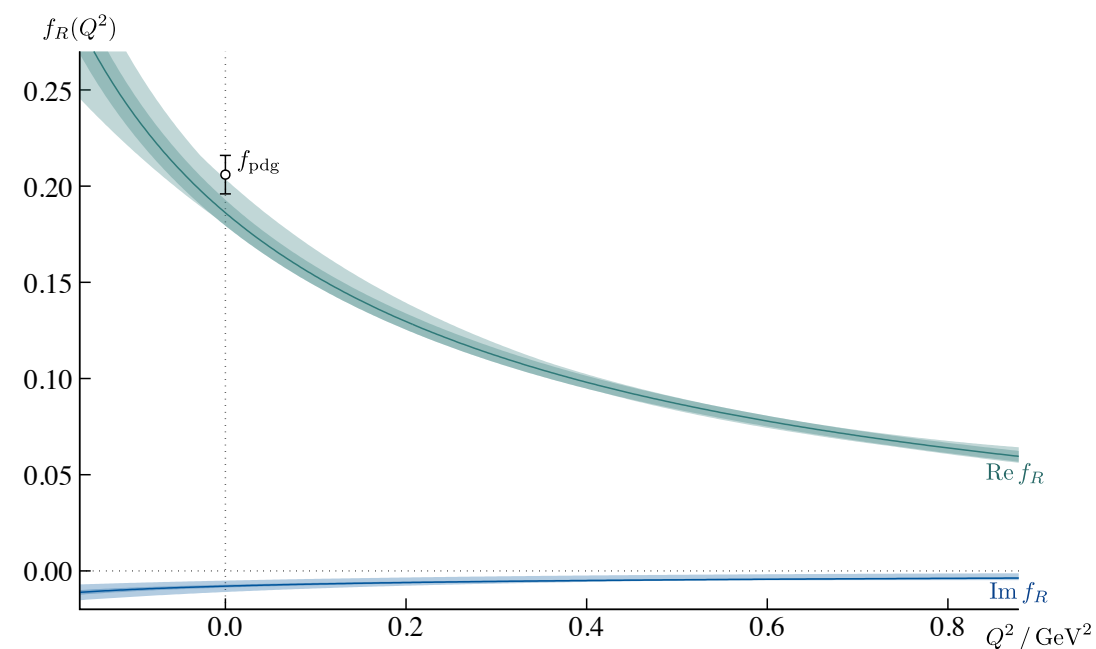
compute three-point functions with a current insertion
correct for the finite-volume

a recent example of the approach ...



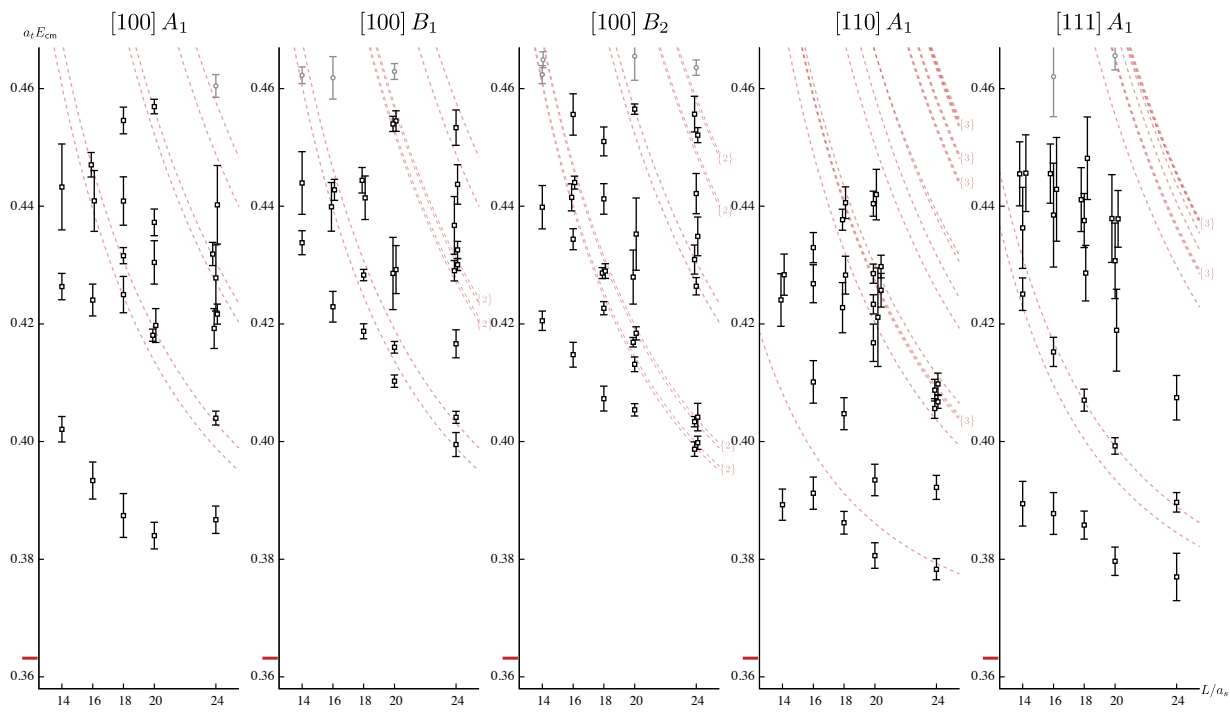
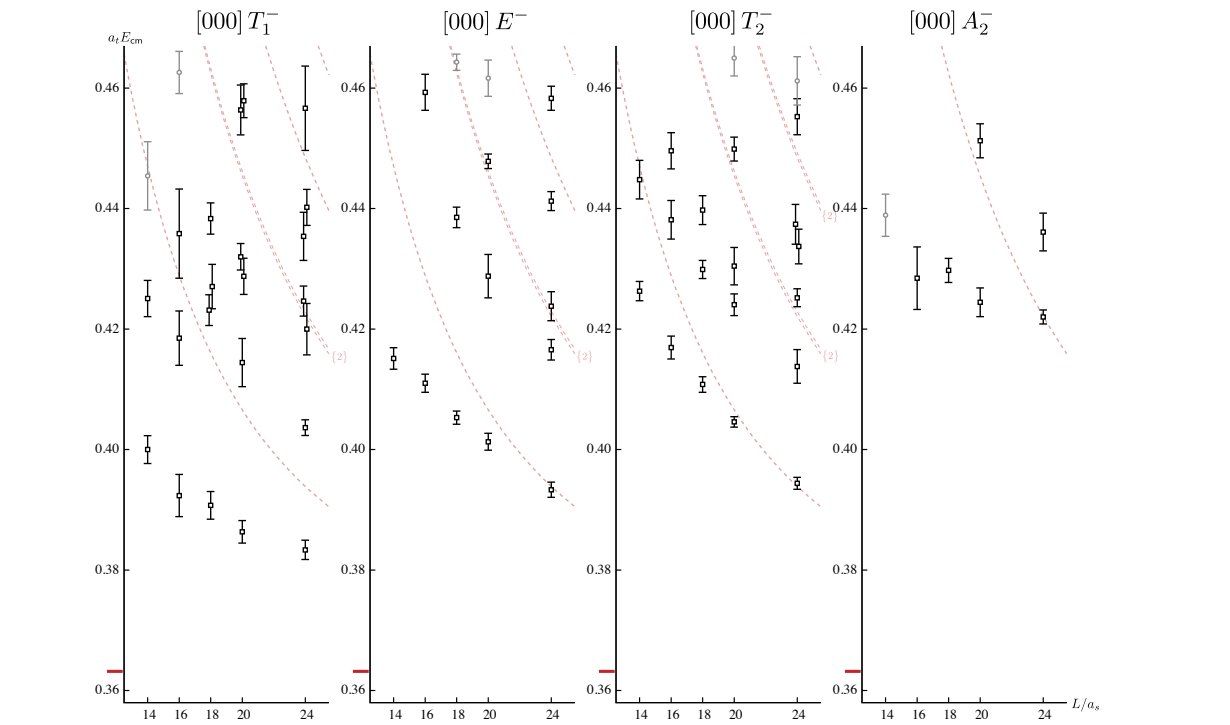
PRD 106 114513 (2022)

resonance transition form-factor $K^* \rightarrow K\gamma$

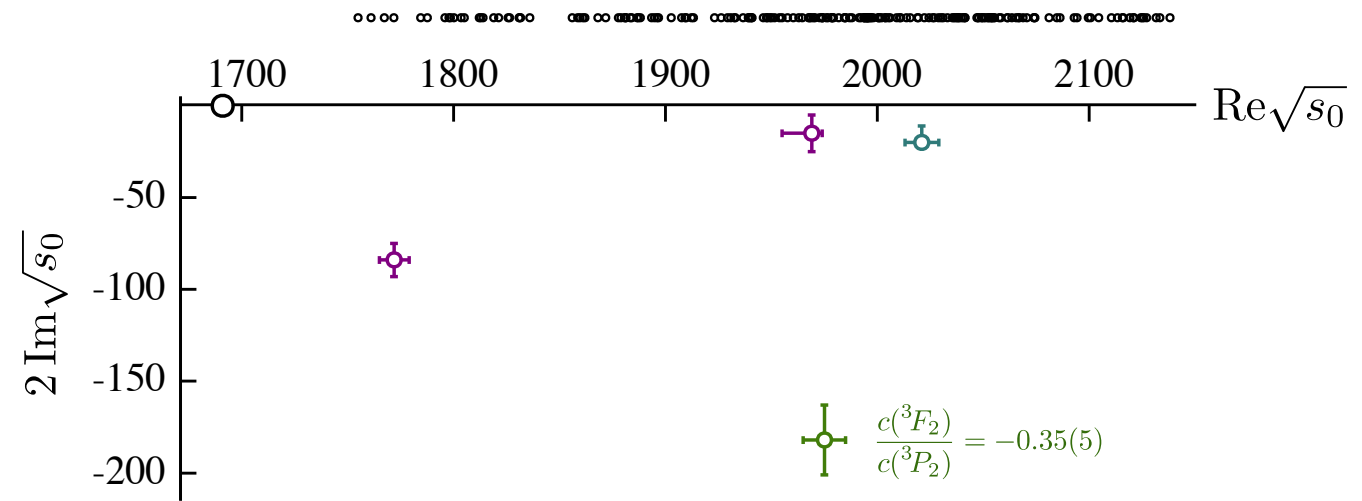
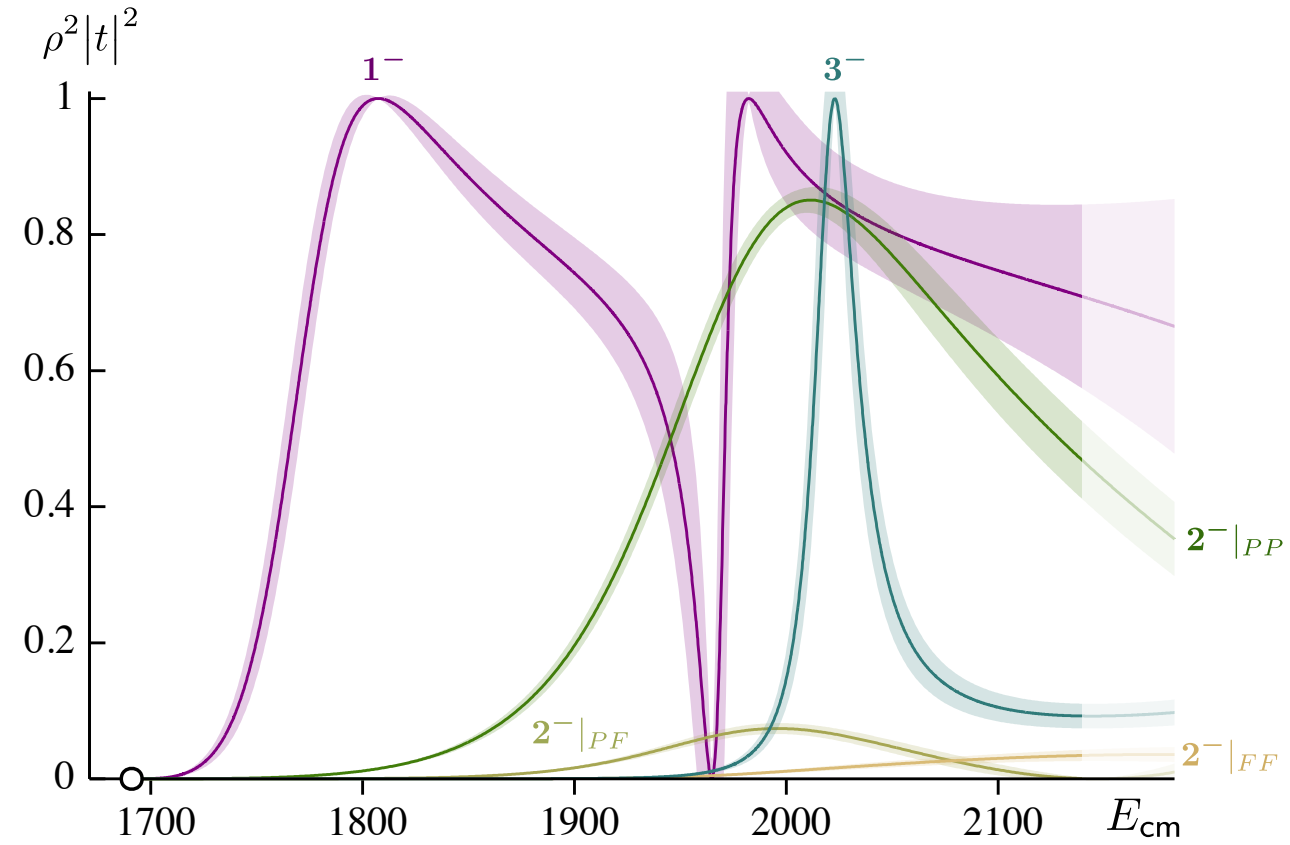


exact SU(3) flavor symmetry

$$\omega_J^1 \rightarrow \eta^8 \omega^8$$



unprecedented number of energy levels



operator basis: ‘single-meson’

$$\bar{\psi}\Gamma\psi$$

(& if you like,
tetraquark & ...)

+ ‘meson-meson’

$$\sum_{\hat{\mathbf{p}}_1, \hat{\mathbf{p}}_2} C(\mathbf{p}_1, \mathbf{p}_2; \mathbf{p}) M_1(\mathbf{p}_1) M_2(\mathbf{p}_2)$$

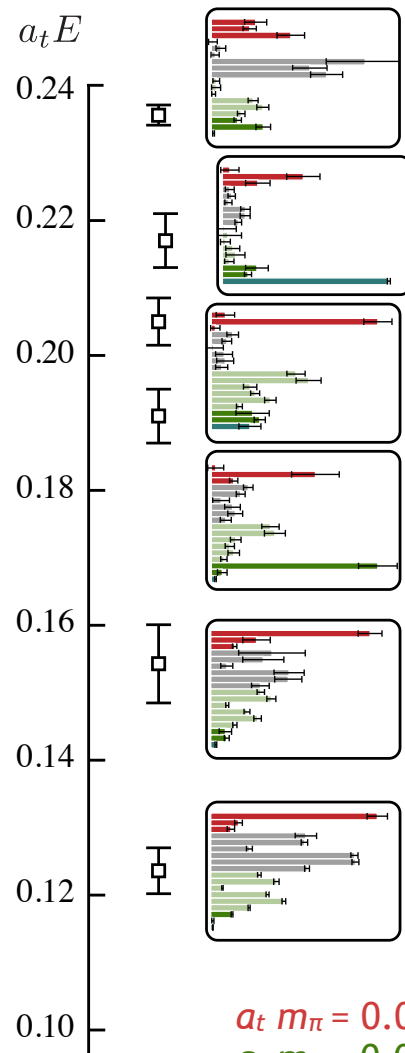
$$\mathbf{p} = \frac{2\pi}{L} [n_x, n_y, n_z]$$

maximum momentum
guided by non-interacting
energies

$$\sqrt{m_1^2 + \mathbf{p}_1^2} + \sqrt{m_2^2 + \mathbf{p}_2^2}$$

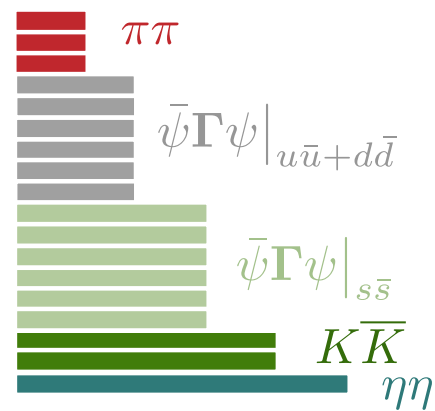
solutions of the det equation
when $t = 0$

[000] A_1^+ 24^3

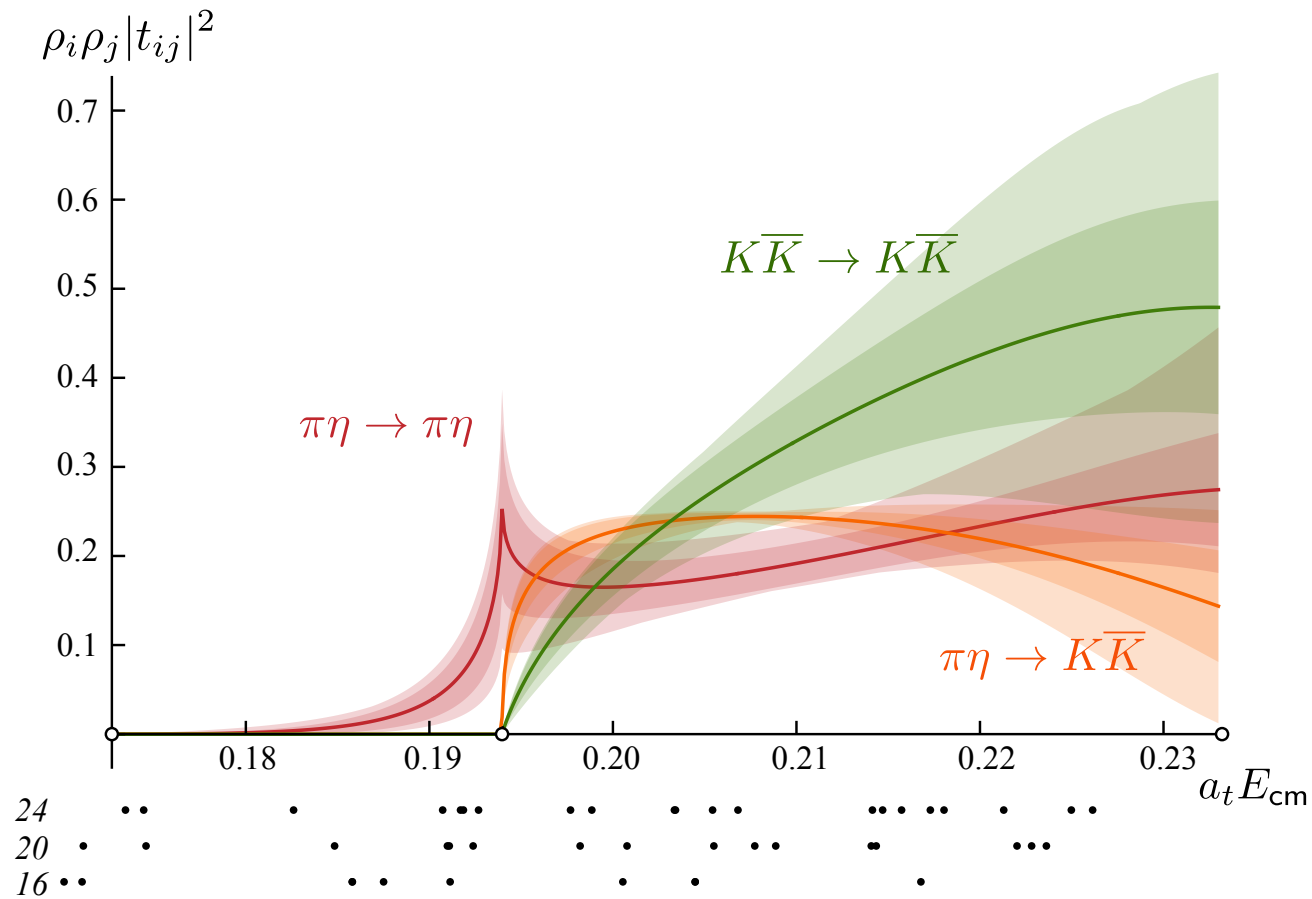


$a_t m_\pi = 0.069$
 $a_t m_K = 0.097$
 $a_t m_\eta = 0.104$

operator basis



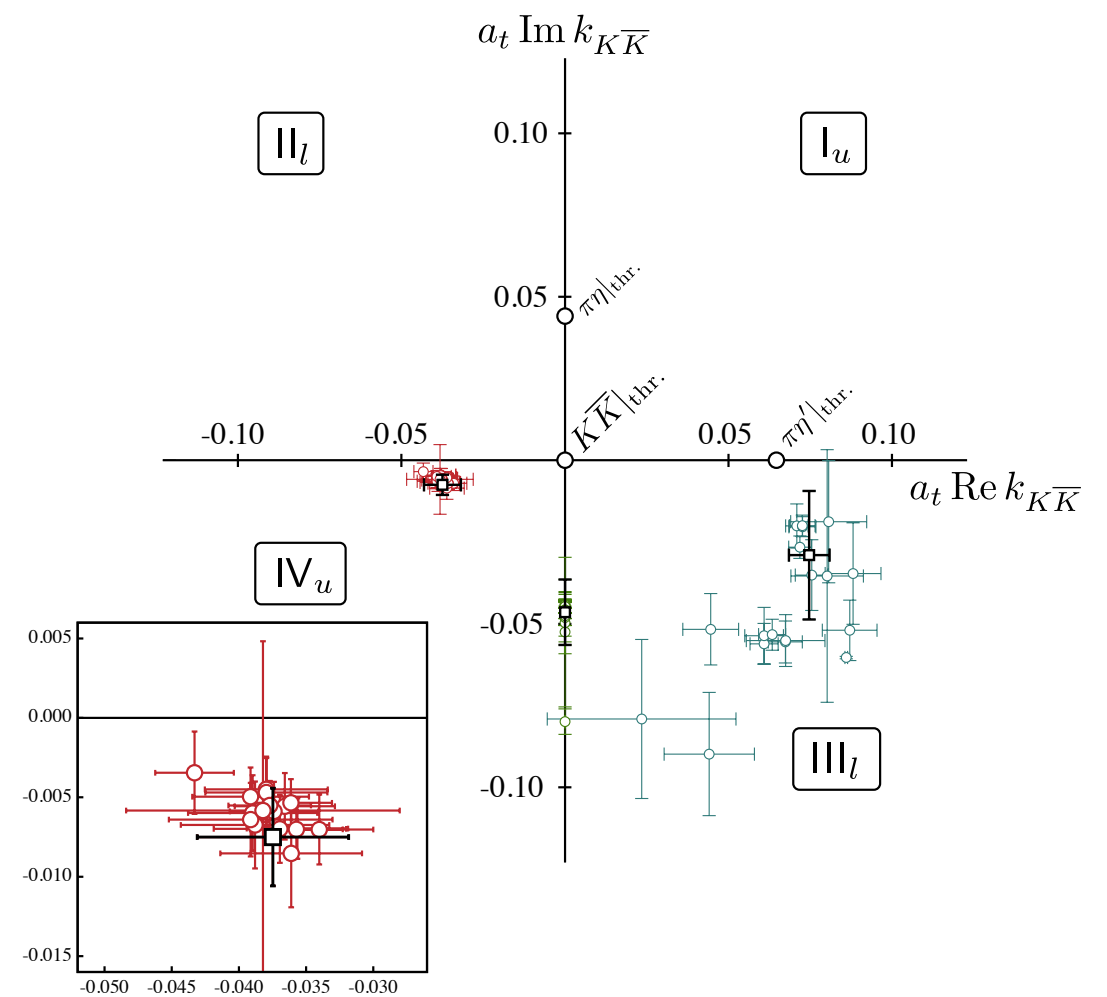
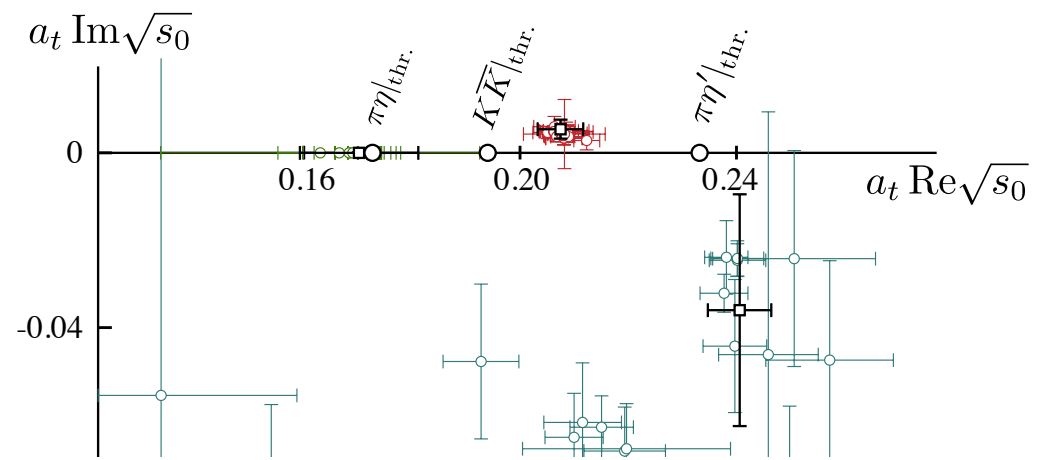
S-wave amplitudes



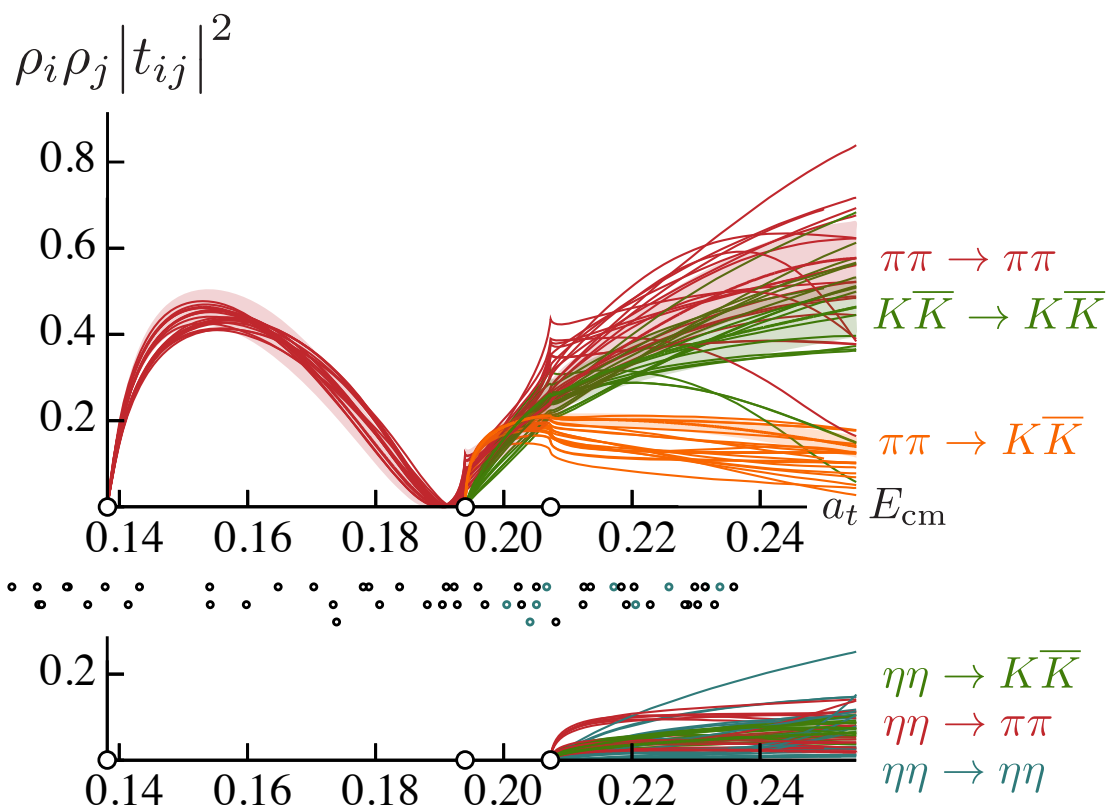
$a_0('980')$ as sheet IV pole

$$\left| \frac{c_{K\bar{K}}}{c_{\pi\eta}} \right| = 1.3(4)$$

pole singularities



S-wave amplitudes

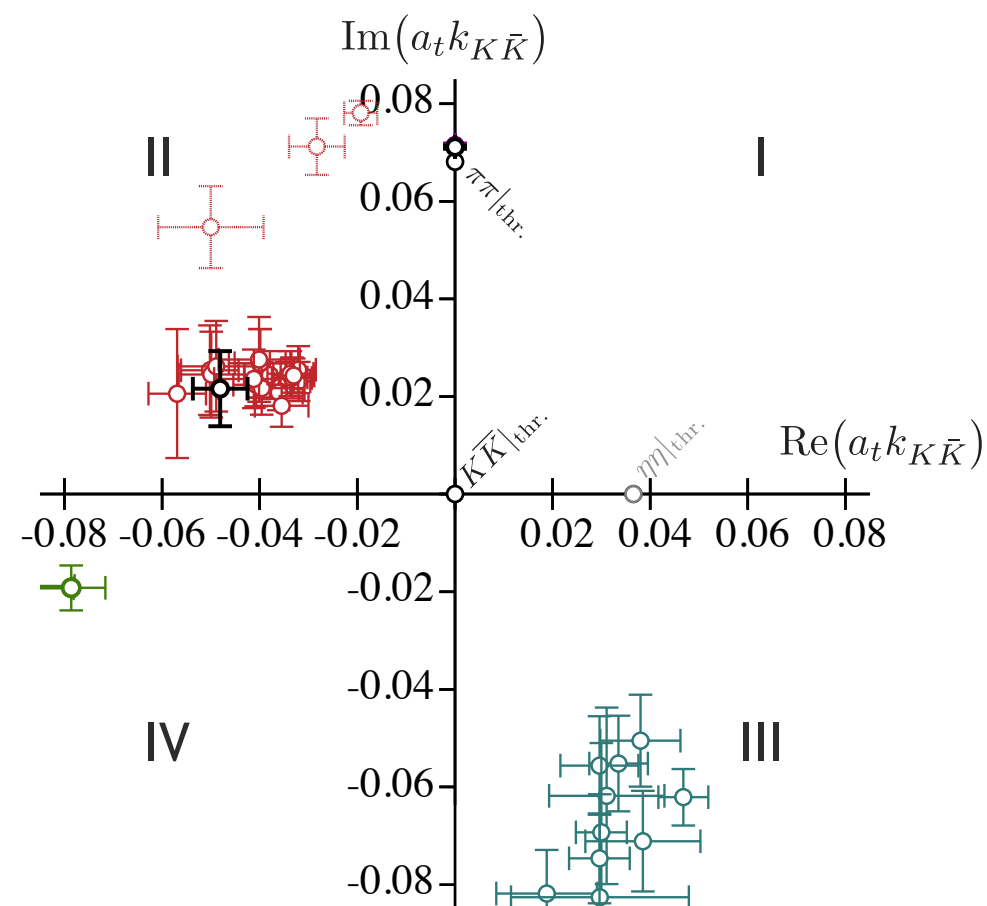
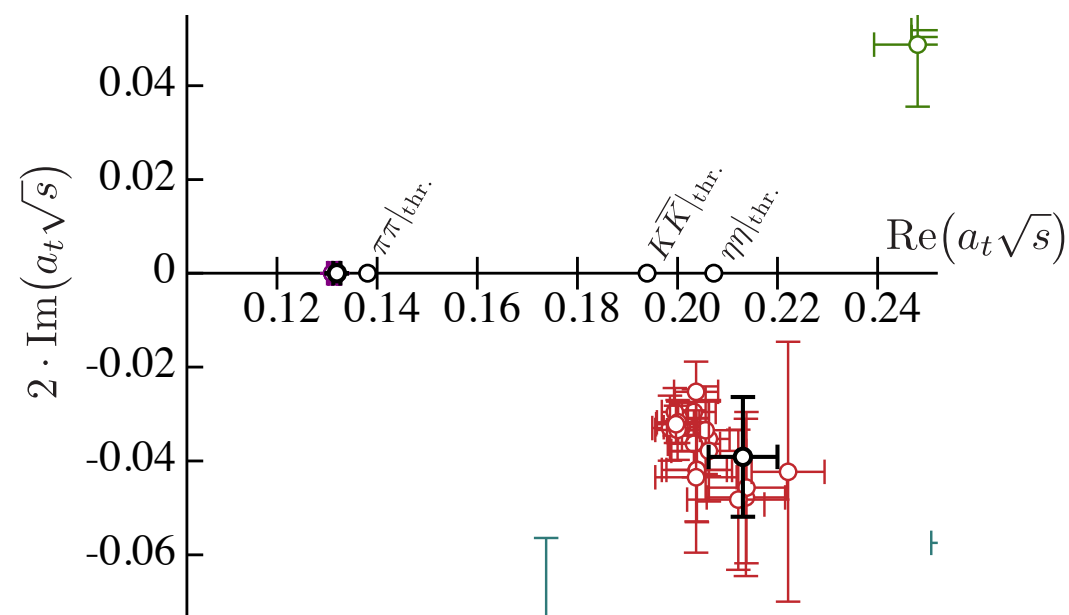


σ as a stable bound-state

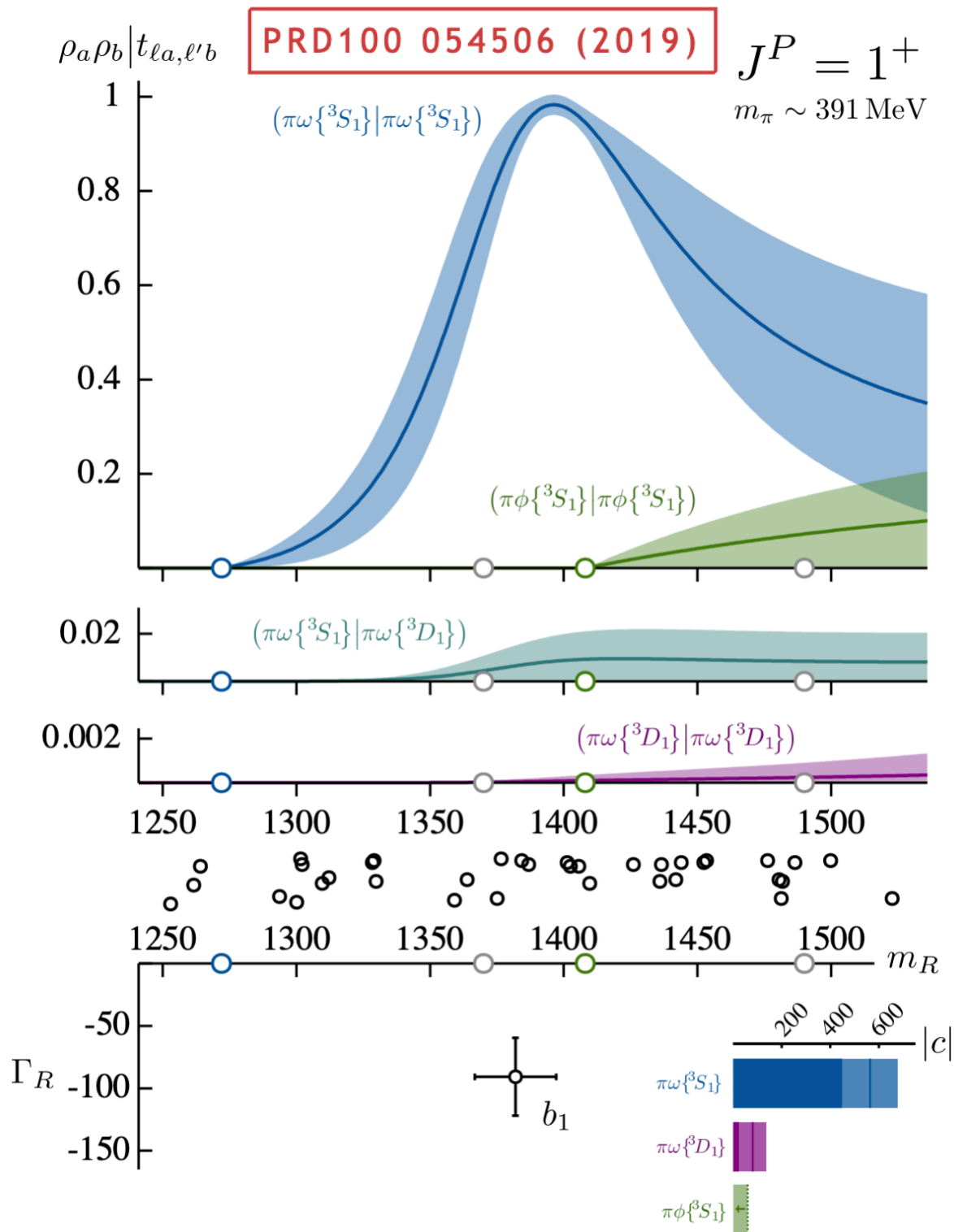
$f_0('980')$ as sheet II pole

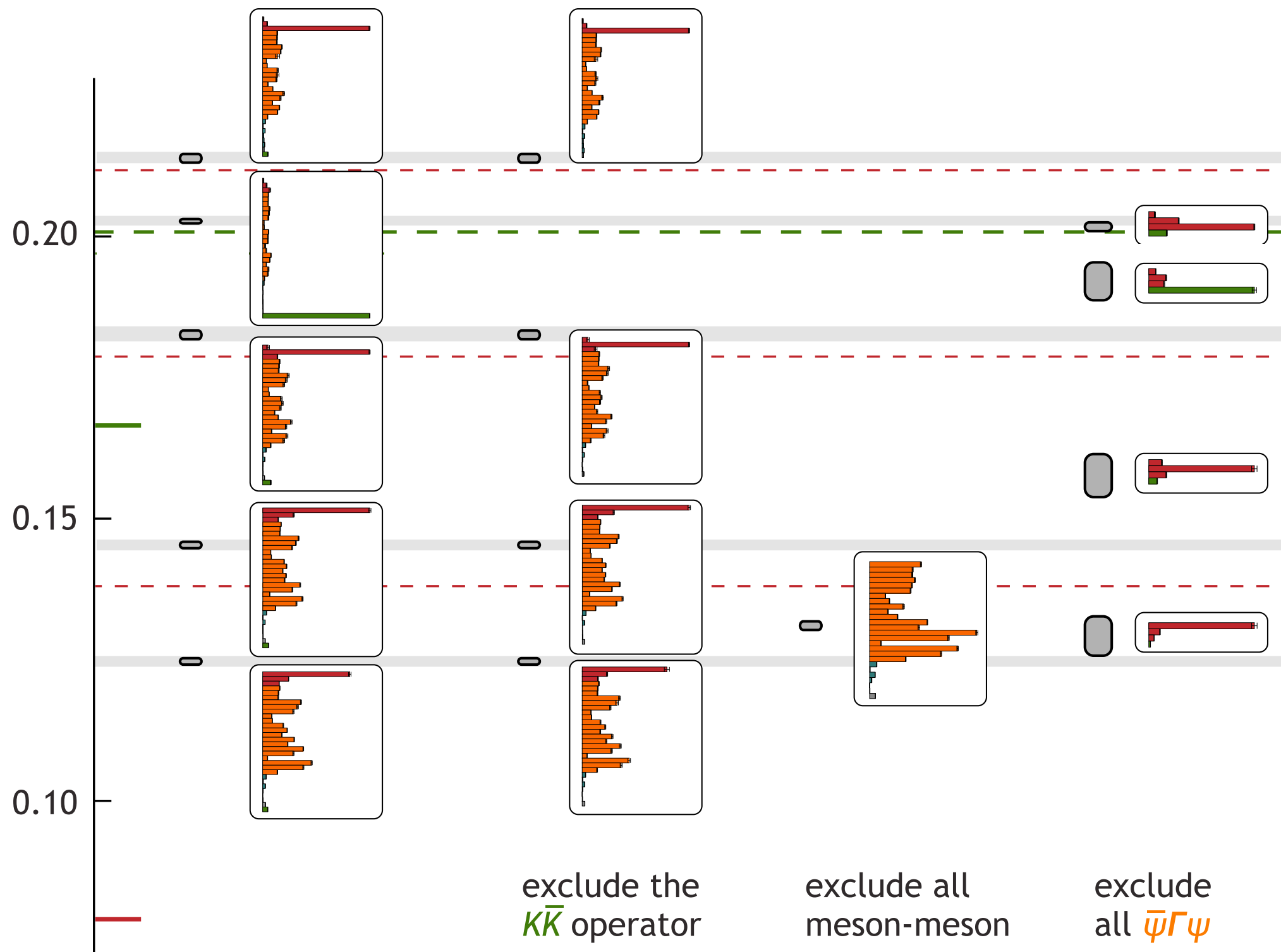
$$|c_{K\bar{K}}| \approx |c_{\pi\pi}|$$

pole singularities



$$b_1 \rightarrow \pi\omega, \pi\phi$$





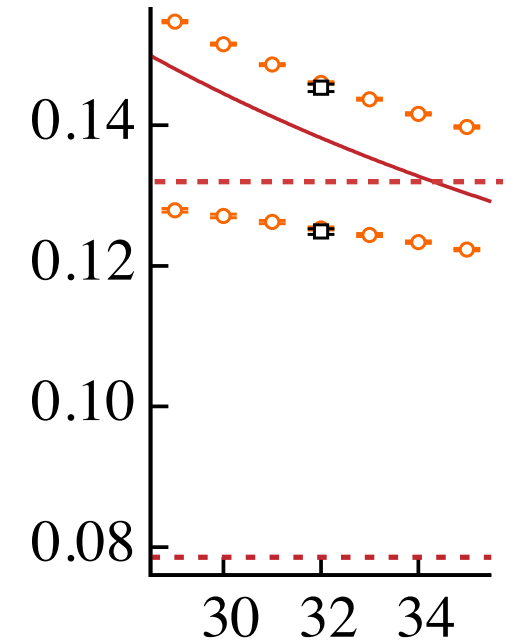
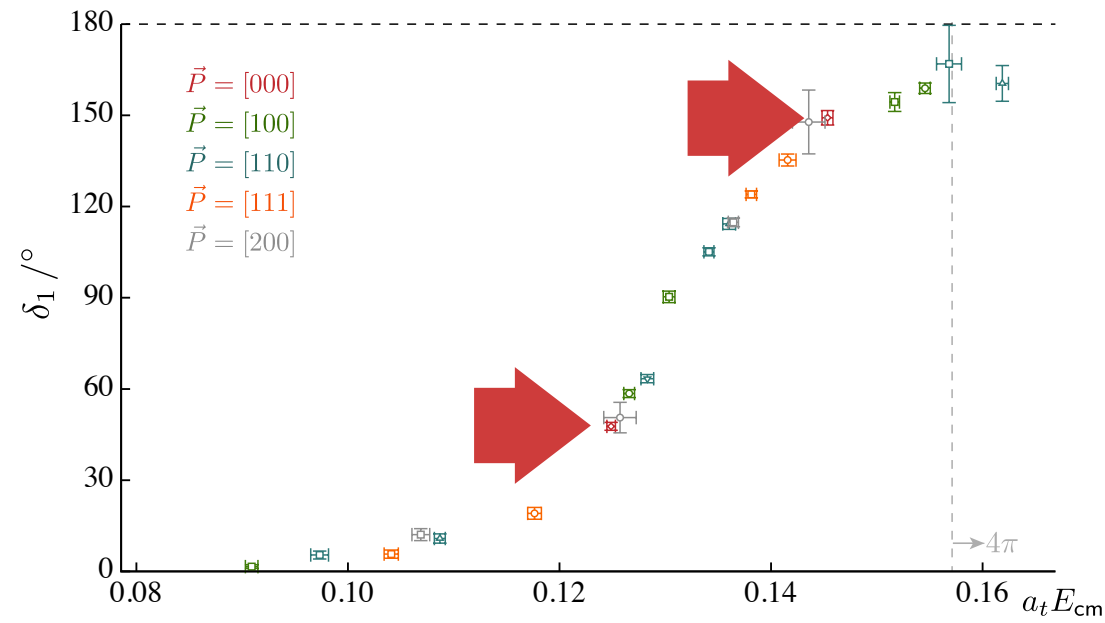
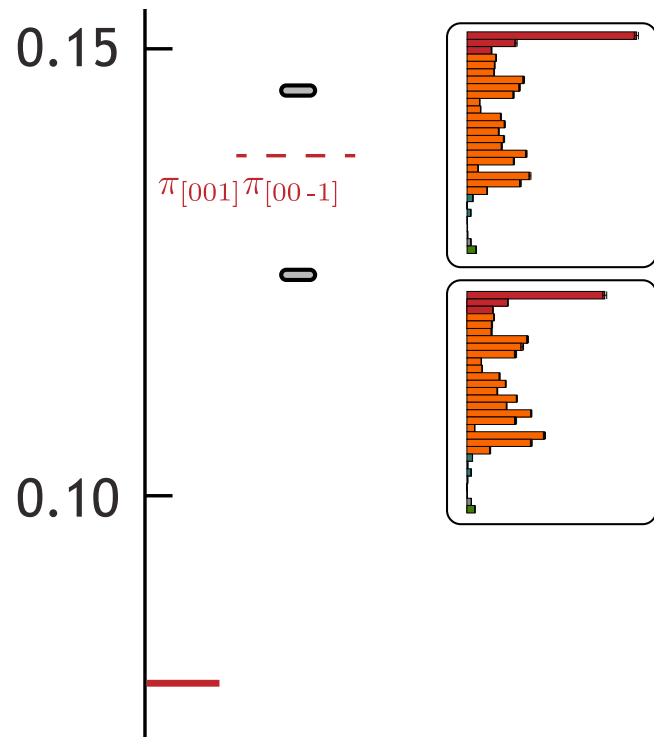
$m_\pi = 0.039$
 $m_K = 0.083$ $L \sim 3.8$ fm

exclude the $K\bar{K}$ operator

exclude all meson-meson

exclude all $\bar{\psi}\Gamma\psi$

focus on the lowest two states



an avoided level crossing