

Standard Model 2/4

Andreas Weiler (TU Munich)

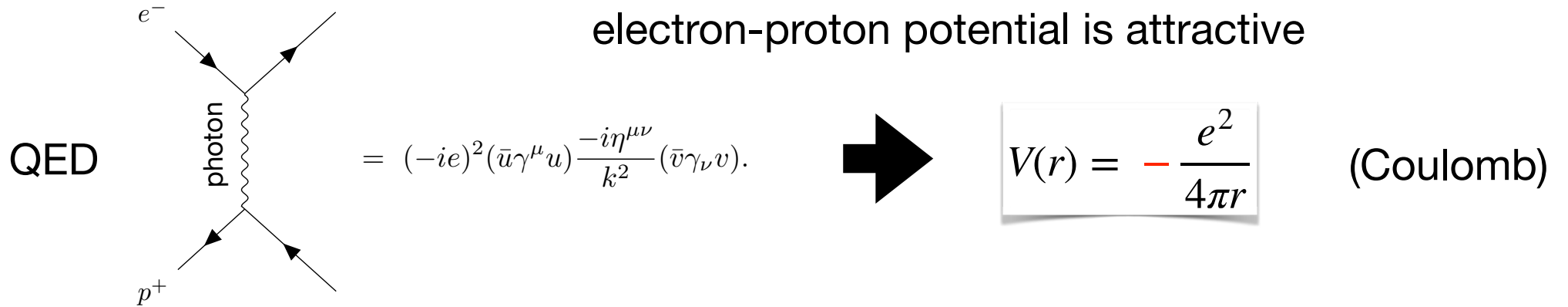
CERN, 7/2024



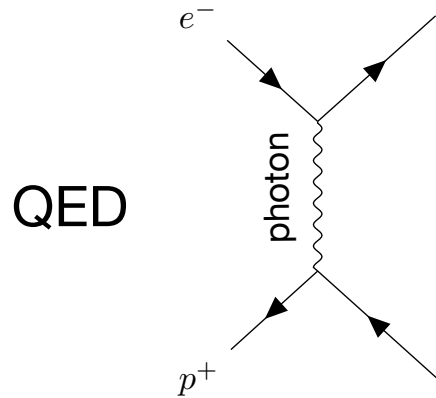
Recap

- Can you show why the photon (or the gluon) turns out to be massless in a gauge invariant theory?
- How many polarizations does a photon or a gluon have? (*Hint*: it has spin 1 and travels with the speed of light). How many entries are in the photon field $A^\mu(x)$, $\mu = 0,1,2,3$?

Example: Coulomb potential

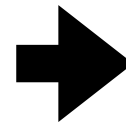


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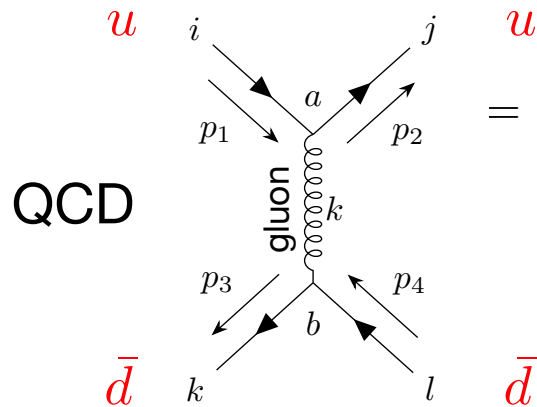
electron-proton potential is attractive

$$= (-ie)^2 (\bar{u} \gamma^\mu u) \frac{-i \eta^{\mu\nu}}{k^2} (\bar{v} \gamma_\nu v).$$



$$V(r) = -\frac{e^2}{4\pi r} \quad (\text{Coulomb})$$

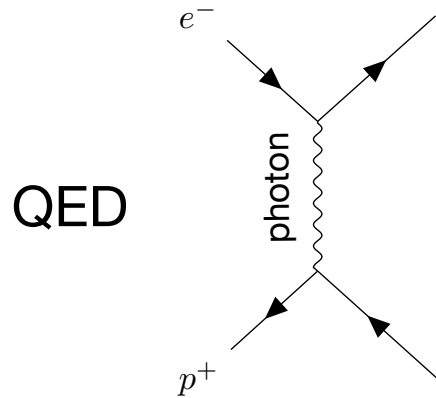
quark-quark potential is only attractive for color neutral combinations *



$$= (ig_s)^2 T_{ji}^a T_{kl}^a \times \bar{u}_j \gamma^\mu u_i \frac{-i \eta_{\mu\nu}}{k^2} \bar{v}_k \gamma^\nu v_l$$

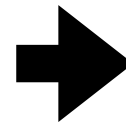
* QCD is strongly coupled at low energies, perturbative calculations are not reliable

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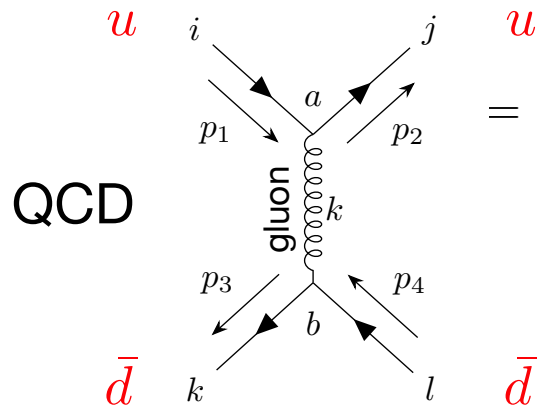
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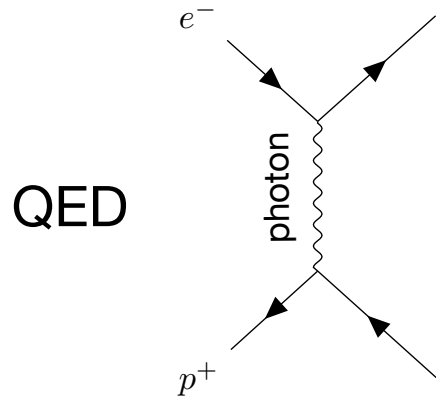


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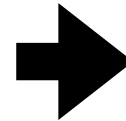
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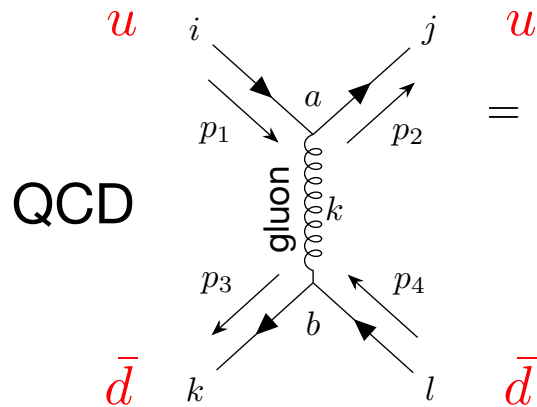
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$$V(r) = \frac{1}{6} \frac{g_s^2}{4\pi r} \quad (\text{color octet})$$

$$V(r) = -\frac{4}{3} \frac{g_s^2}{4\pi r}. \quad (\text{color singlet})$$

* QCD is strongly coupled at low energies, perturbative calculations are not reliable

$$u(r) \bar{d}(g) \rightarrow u(r) \bar{d}(g)$$

$$u(r) \bar{d}(r) \rightarrow u(b) \bar{d}(b)$$

Kinetic term for SU(N) gauge boson

We cannot recycle the Maxwell action. The Lagrangian would not be invariant under a local SU(N) transformation

$$A_\mu(x) \rightarrow U(x)A_\mu(x)U(x)^\dagger - \frac{i}{g}(\partial_\mu U(x))U(x)^\dagger$$

Field strength now contains a **non-abelian contribution**

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu + ig[A_\mu, A_\nu]$$

It transforms homogeneously

$$F_{\mu\nu} \rightarrow U(x) F_{\mu\nu} U^{-1}(x)$$

and we can build an invariant Lagrangian

$$\mathcal{L} = -\frac{1}{4}\text{Tr}(F_{\mu\nu}F^{\mu\nu}) = \dots + gAAA + g^2AAAA$$

Note: Gluons carry **colour** charge and do interact with themselves.

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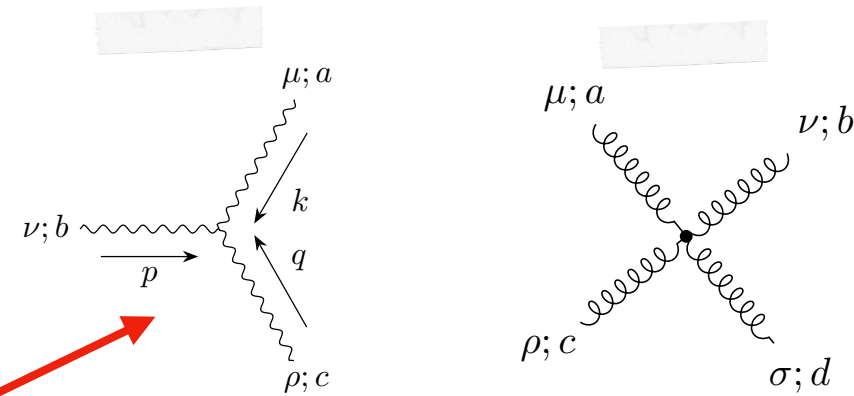
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$$\mathcal{L} = ?$$

Which theory is realized in nature? $SU(N)$? $U(1)$? Which particles?

How can we discover the Lagrangian of the universe?

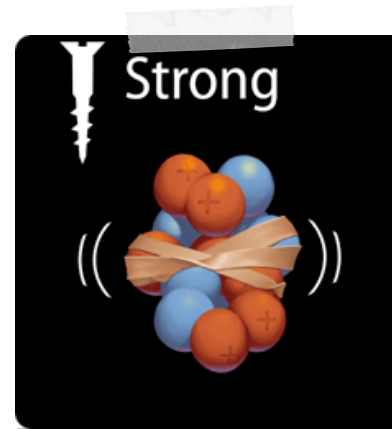
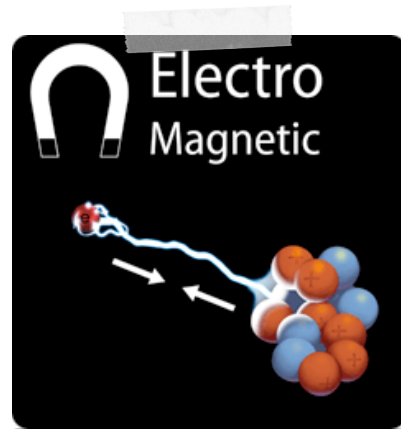
We need experiments!

The strong and the electromagnetic interactions

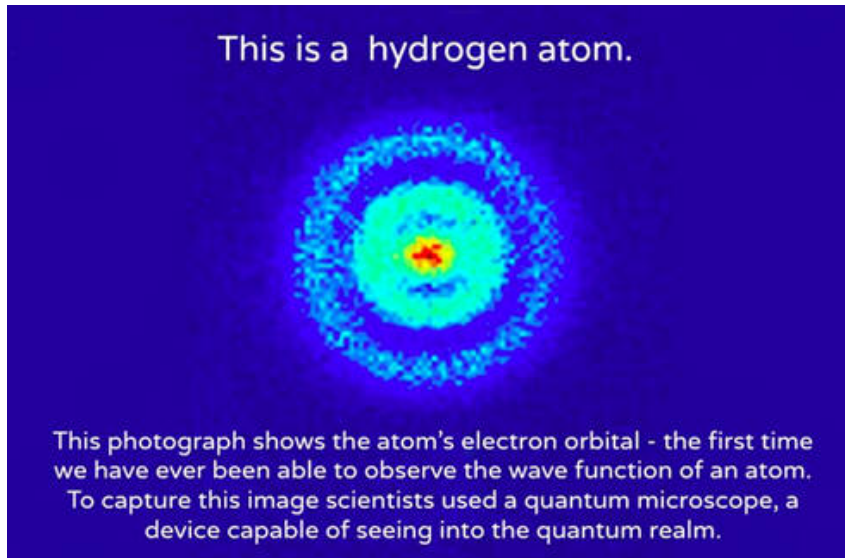
$U(1)_{em}$

$SU(3)_c$

- Why was evidence of electromagnetic interactions discovered before evidence of strong interactions?
- Why have we never seen a free quark, unlike electrons or protons?
- How can we test predictions about quarks if we don't observe them as free particles?



QED binds electrons and nuclei inside atoms and molecules



But quarks are fundamental objects, *not* the composite nuclei

$$p = (uud), \quad n = (udd), \dots$$

Charges:

$$\text{up-quark:} \quad + 2/3$$

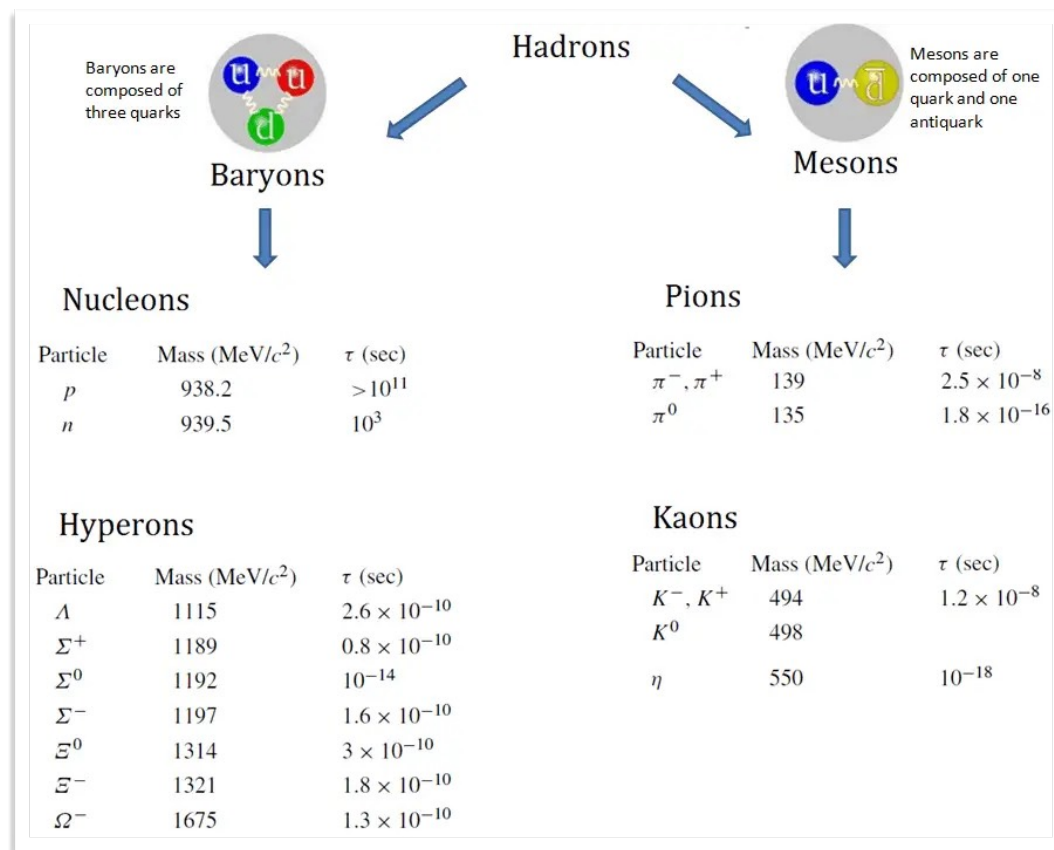
$$\text{down-quark:} \quad -1/3$$

$$\text{electron:} \quad -1$$

$$u(x) \rightarrow e^{i\frac{2}{3}e\alpha(x)} u(x), \quad d(x) \rightarrow e^{-i\frac{1}{3}e\alpha(x)} d(x), \quad e(x) \rightarrow e^{-ie\alpha(x)} e(x)$$

QCD binds quarks into hadrons

$p = (uud)$, $n = (udd)$, $\pi^+ = (\bar{d}u)$, ... and hundreds more.



Coupling “constants” : QED

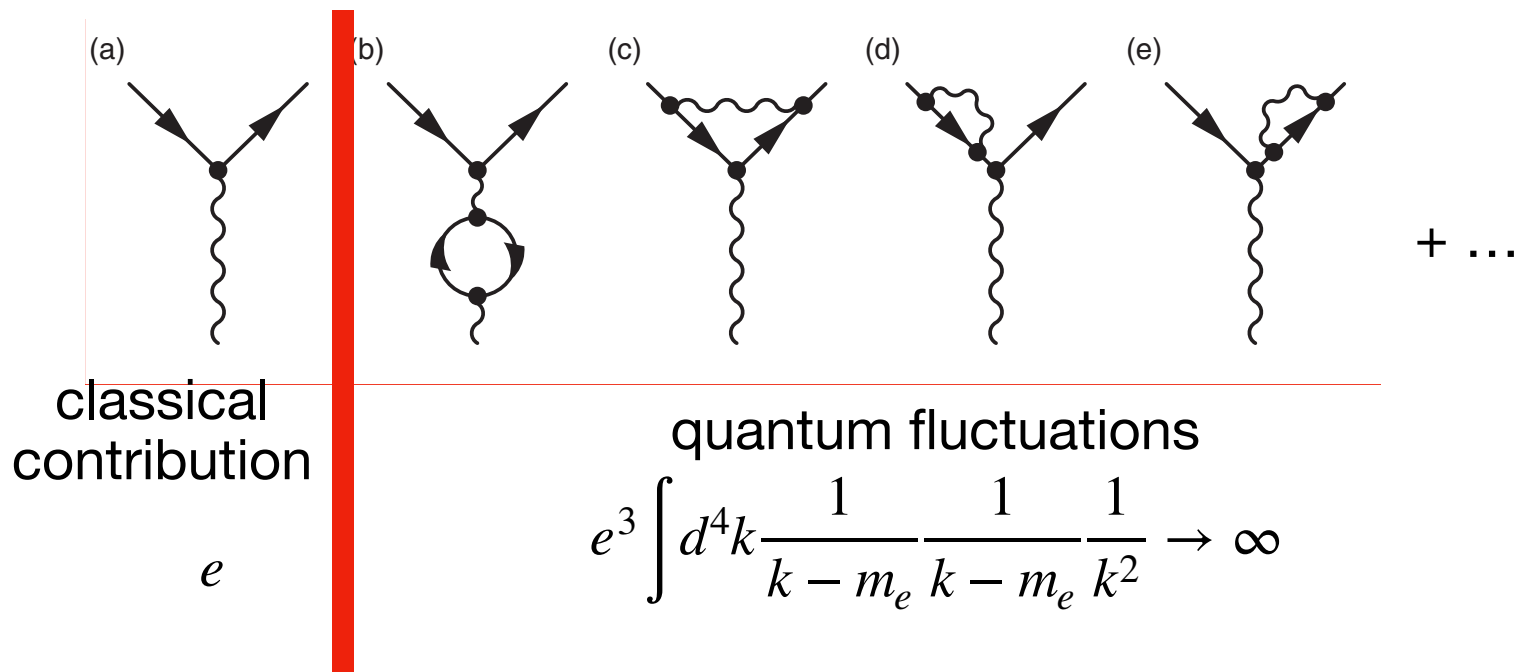
Classical physics: forces depend on distances

Quantum field theory: **charges** also depend on distances

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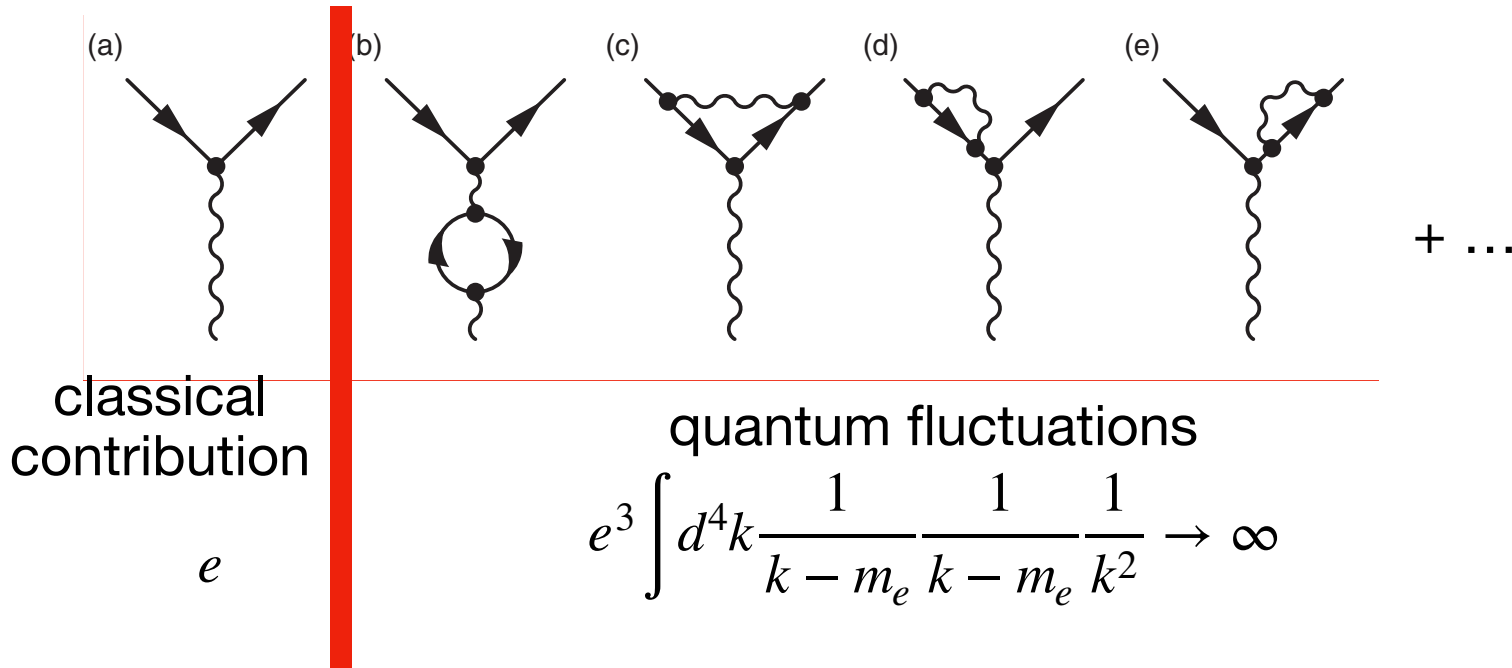
Quantum field theory: **charges** also depend on distances



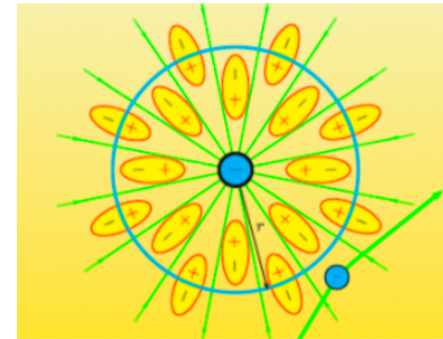
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intuitive picture



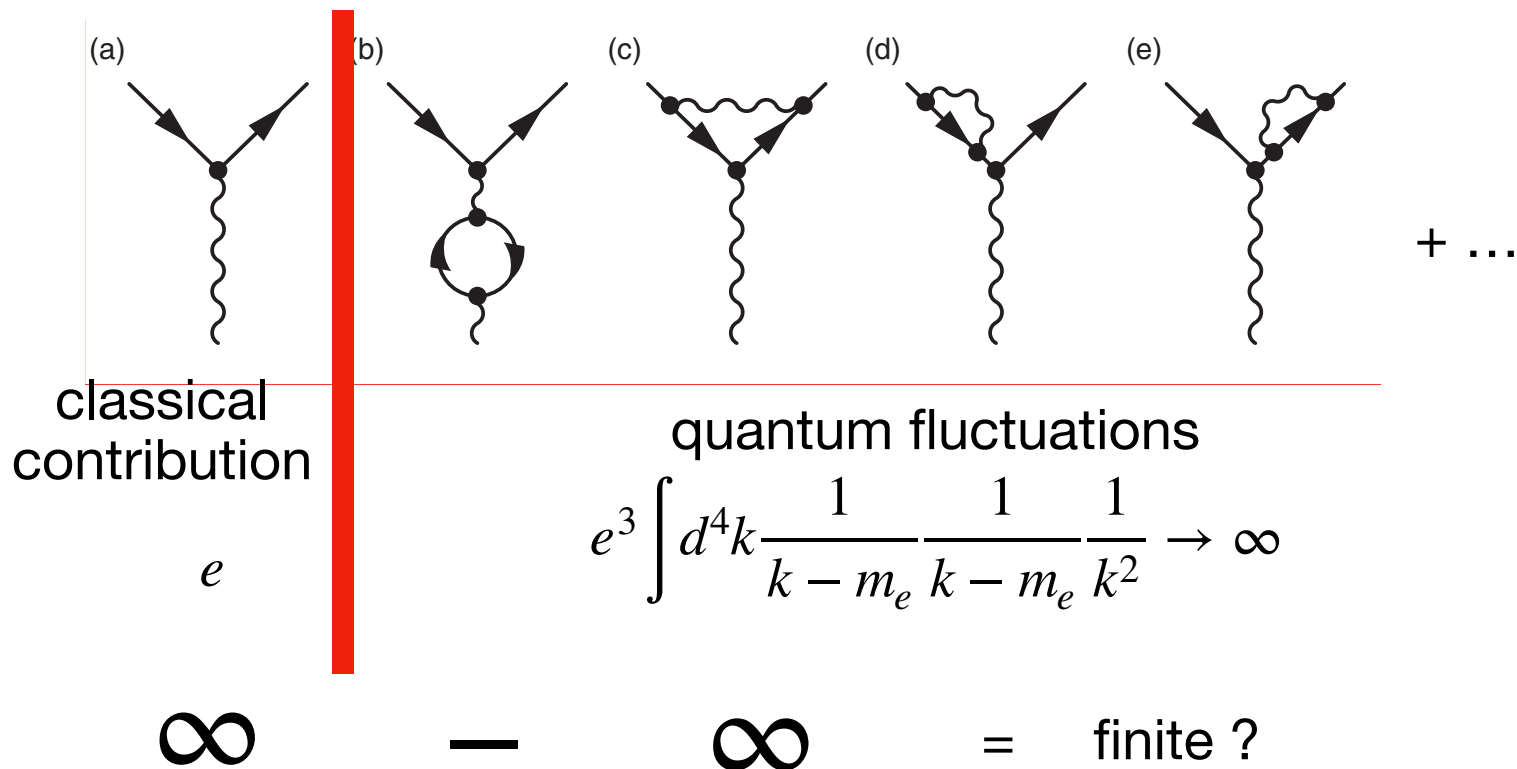
The vacuum screens the electric charge
 -> **infrared free**

charge weaker at lower E
 at larger r

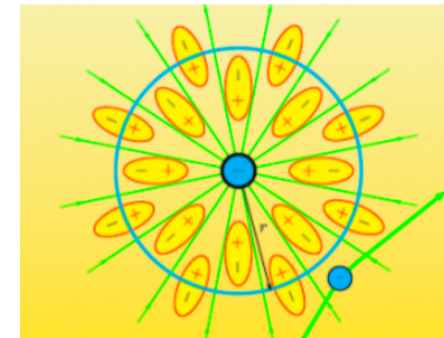
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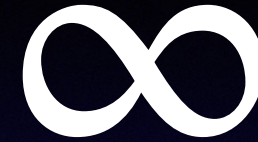
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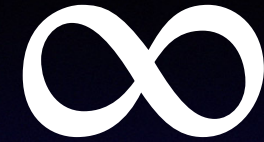
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The development of quantum electrodynamics.
1937 (colourised).



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Modern view: “**Ignorance is no shame**”.

We can't trust our QFT up to infinite energy, so we should not include virtual particles up to infinite energy. We introduce a maximum energy (a cut-off) to **regularize** the theory. We compare with the measurement to determine the value of classical + regularized virtual (= **renormalize**).

This is a good thing: for example, Feynman, Schwinger, Tomonaga, and others who developed quantum electrodynamics did not have to know about the top quark.

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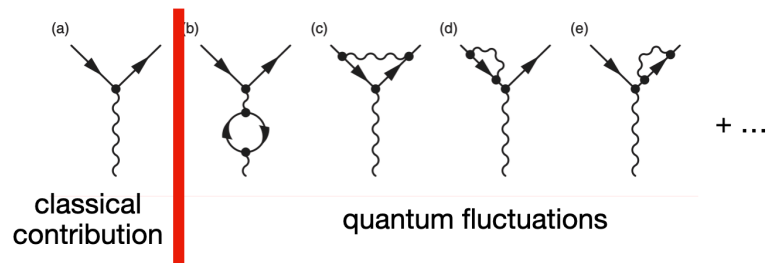
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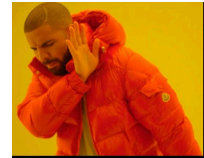
Coupling “constants” : QED

Classical physics: forces depend on distances

Quantum field theory: charges also depend on distances



Fine structure constant: $\alpha_{QED} = \frac{e^2}{4\pi}$



42

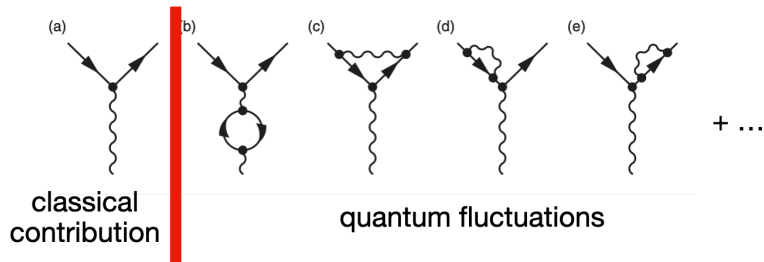


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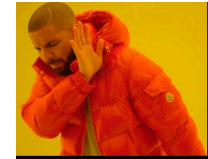


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$$\frac{1}{\alpha(0)} = 137.035999074(44)$$
$$\frac{1}{\alpha(90\text{GeV})} = 127.950(17)$$



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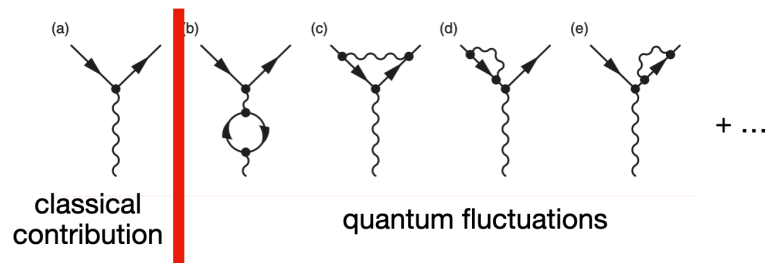
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1/137

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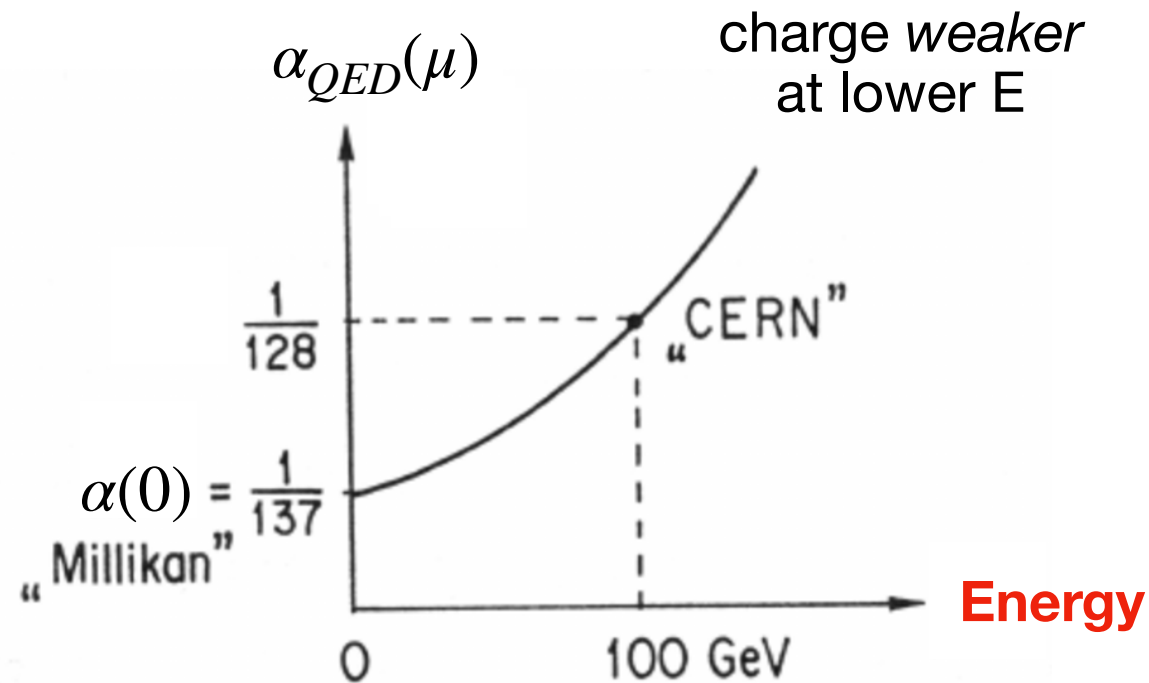
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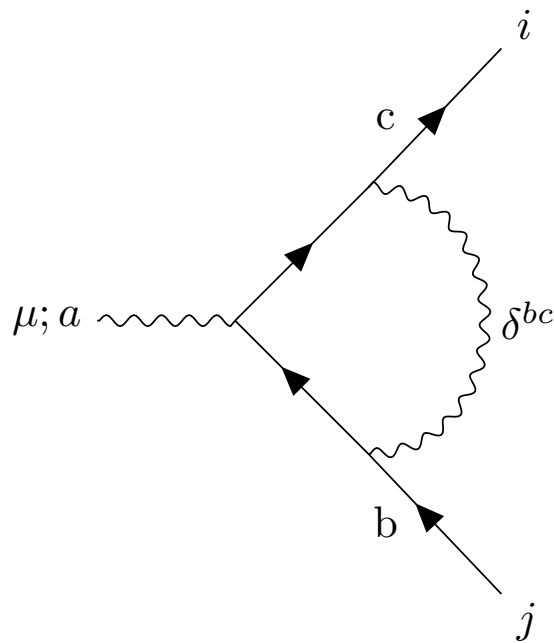
Q: Should we expect the virtual particle contribution to be the same as in QED?

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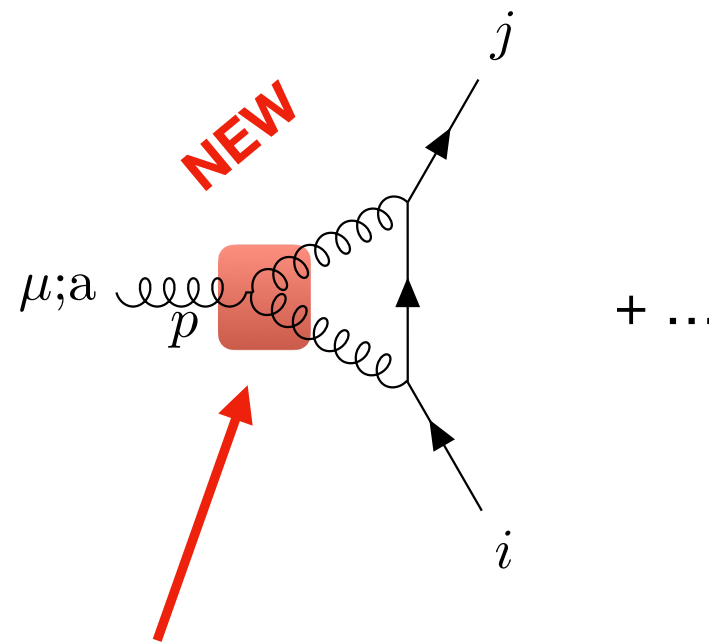
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QED like

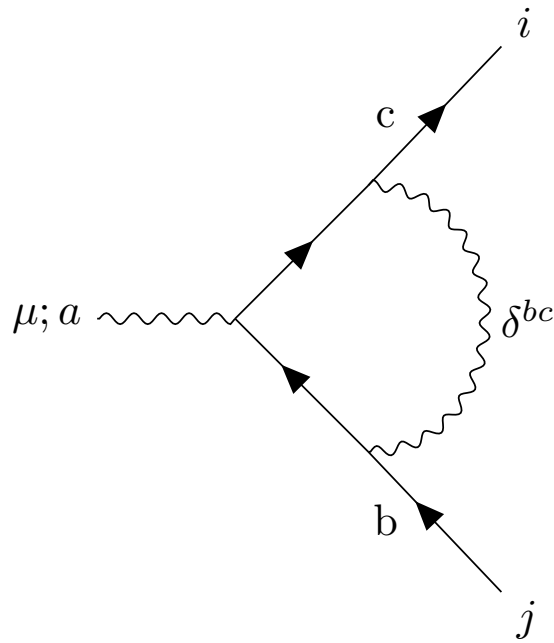


gluons interact with themselves

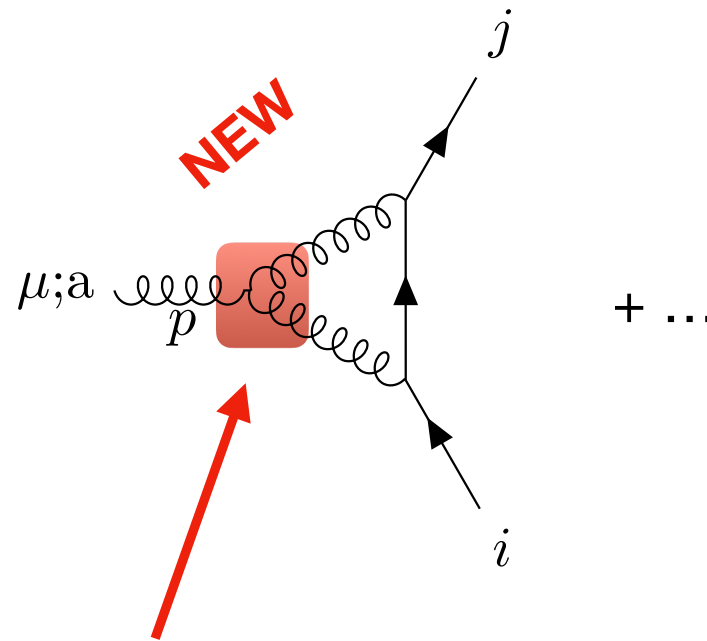
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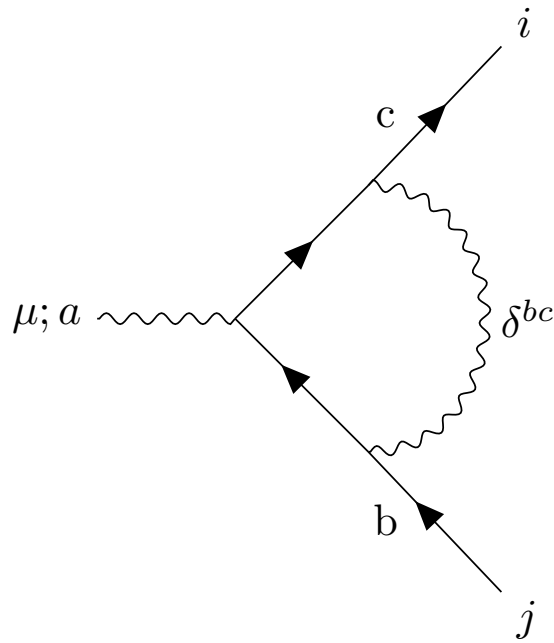


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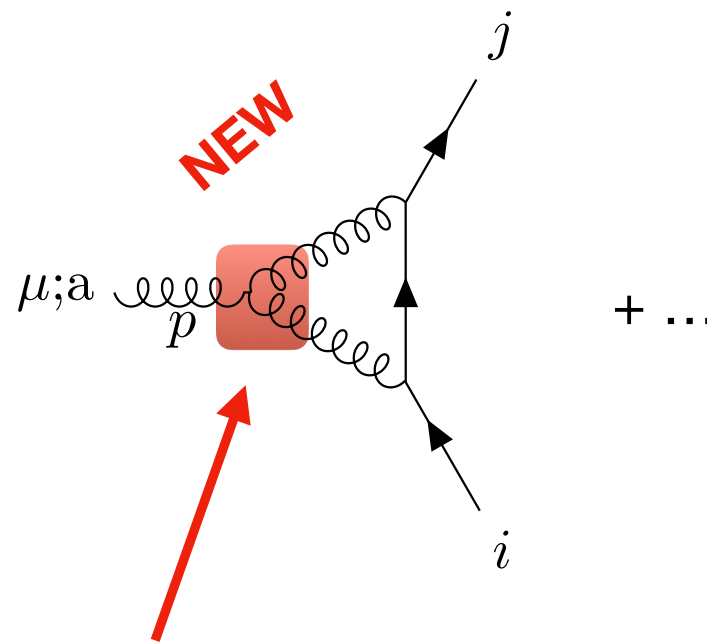
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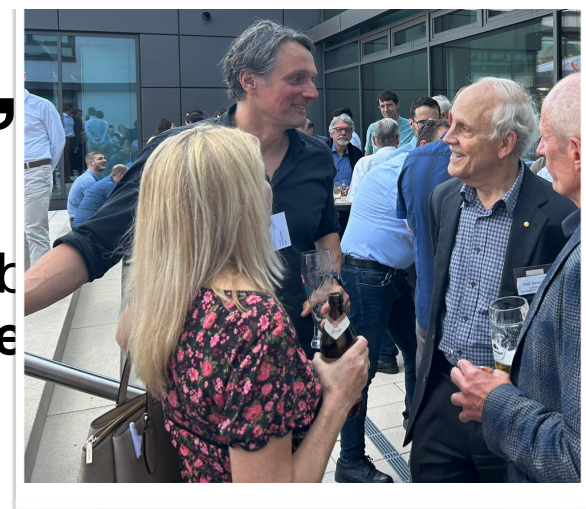
QED like



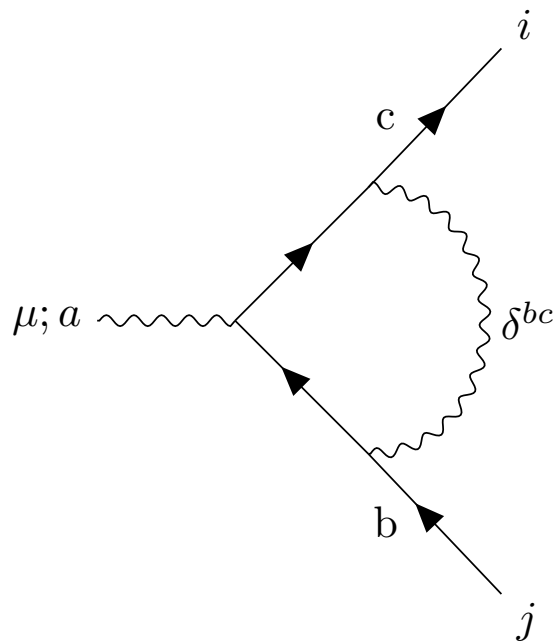
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Coupling “constants”

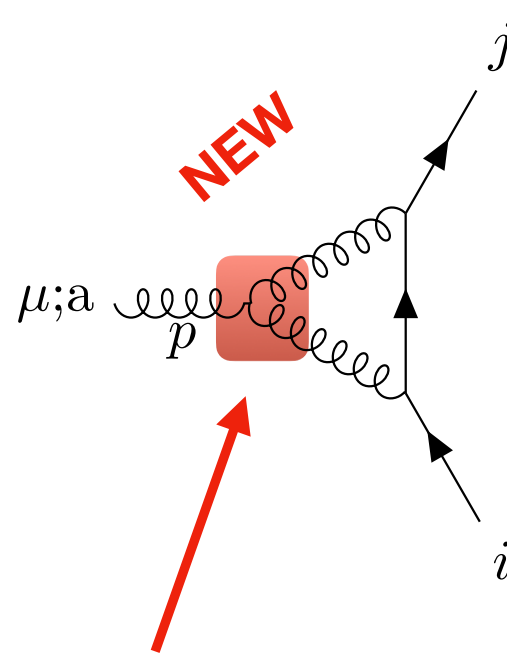
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fact



QED like



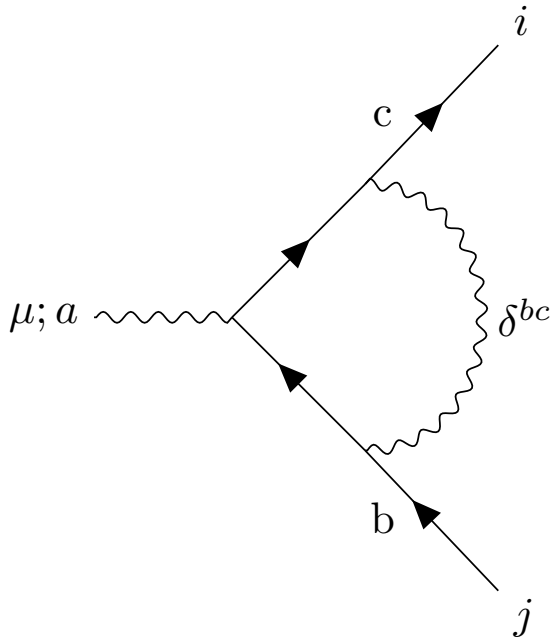
gluons interact with themselves

+ ...

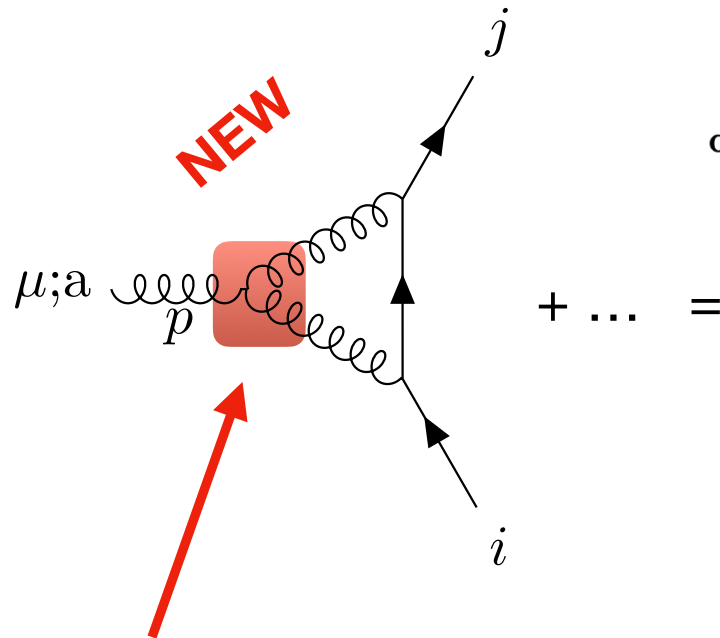
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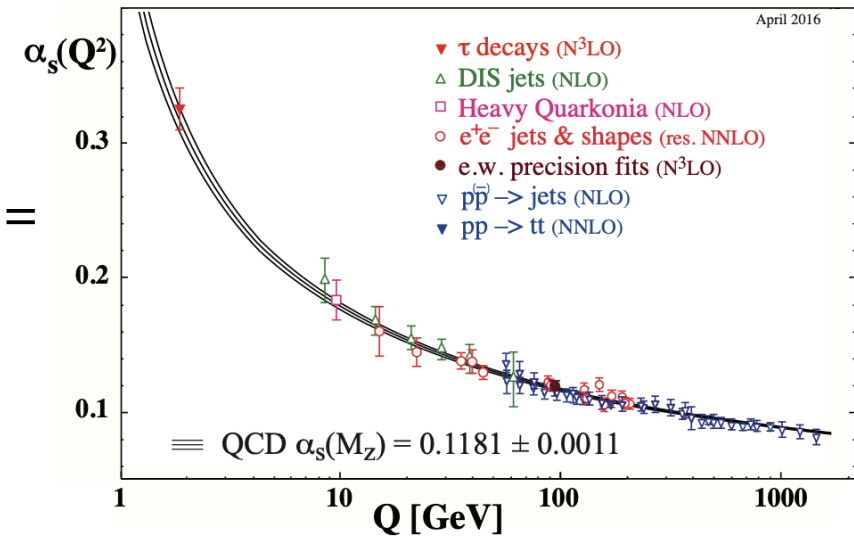


QED-like



gluons interact with themselves

stronger at lower E



Evolution of coupling constants

QED: virtual particles screen charge, *weaker* longer distances
= “infrared freedom”

infinite range $q_1 \bullet \text{---}\gamma\text{---} \bullet q_2$ $V = \frac{q_1 q_2}{r}$

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infinite range $q_1 \bullet \text{---} \gamma \text{---} \bullet q_2$ $V = \frac{q_1 q_2}{r}$

QCD: virtual particles anti-screen charge, *stronger* longer distances
= “asymptotic* freedom”

(* asymptotic means at high energies)



Cannot separate color charges!

The range of the strong interactions determined by the exchange of the **lightest colorless hadron** (= pion)

$$V = \frac{g_1 g_2}{r} e^{-m_\pi r}$$

$$m_\pi = 125 \text{ MeV}, \quad \frac{1}{m_\pi} \approx 1 \text{ Fermi} = 10^{-13} \text{ cm}$$

Measuring the quark charge

Can we measure the electric charge of the quarks?

Can we test that there are $N_C = 3$ colors?

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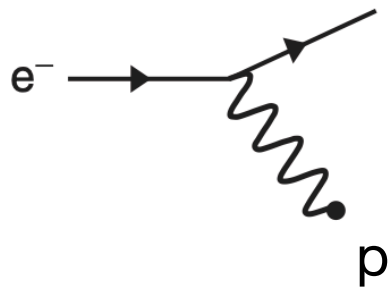
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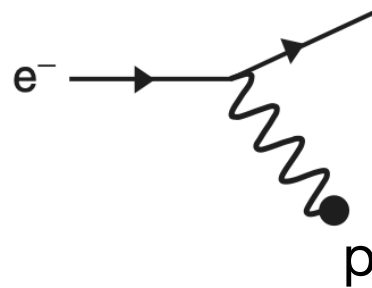
p appears pointlike

p has geometric features

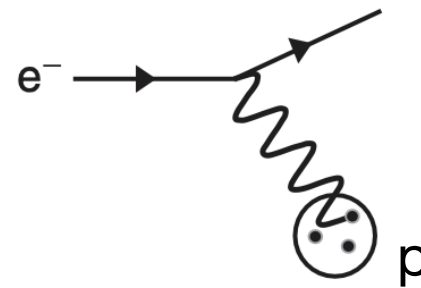
we see quarks !



$$\lambda \gg r_p$$



$$\lambda \sim r_p$$



$$\lambda < r_p$$

r_p : proton radius



Measuring the quark charge

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Let's collide an electron and a positron

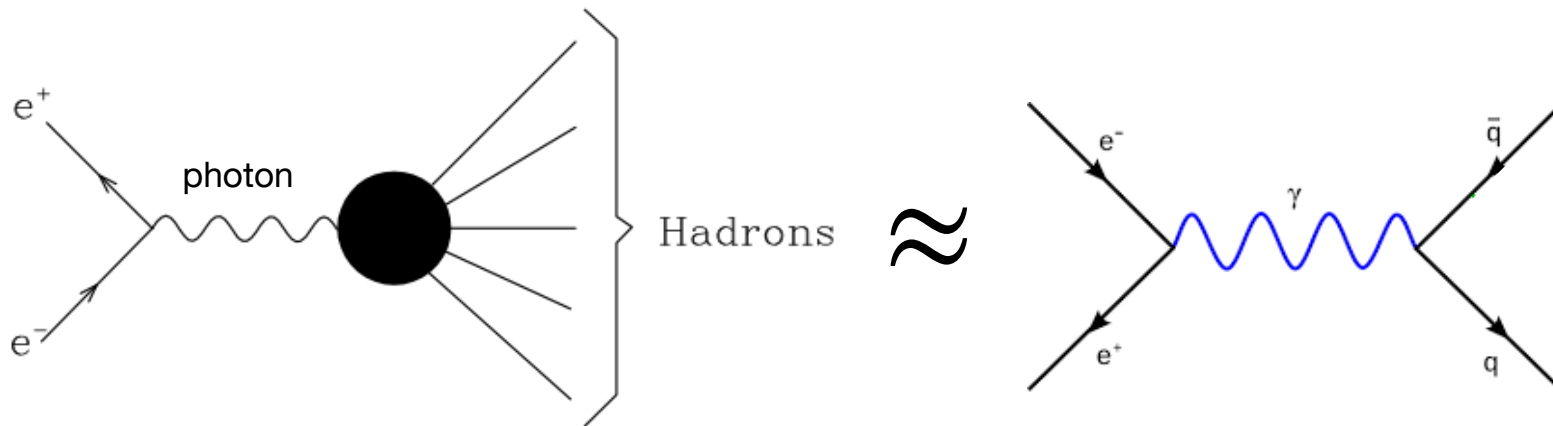
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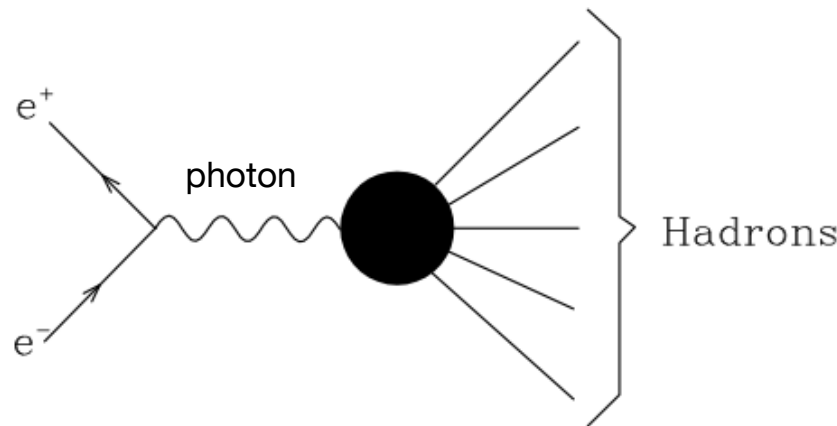
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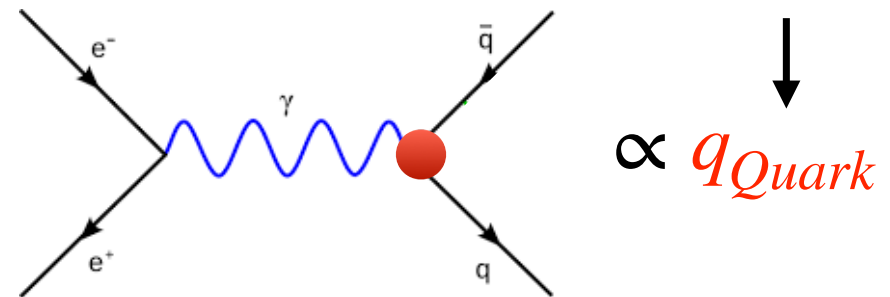
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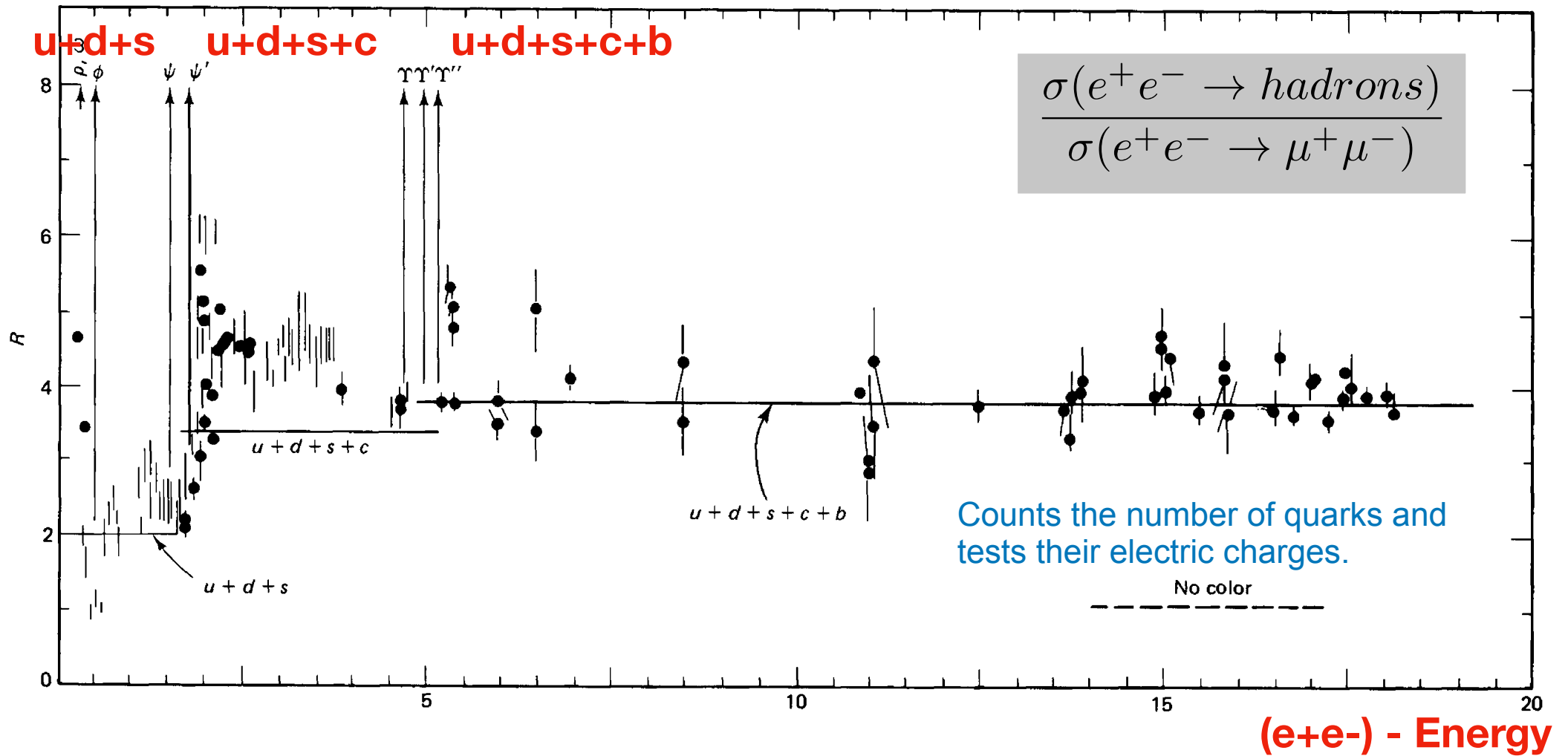
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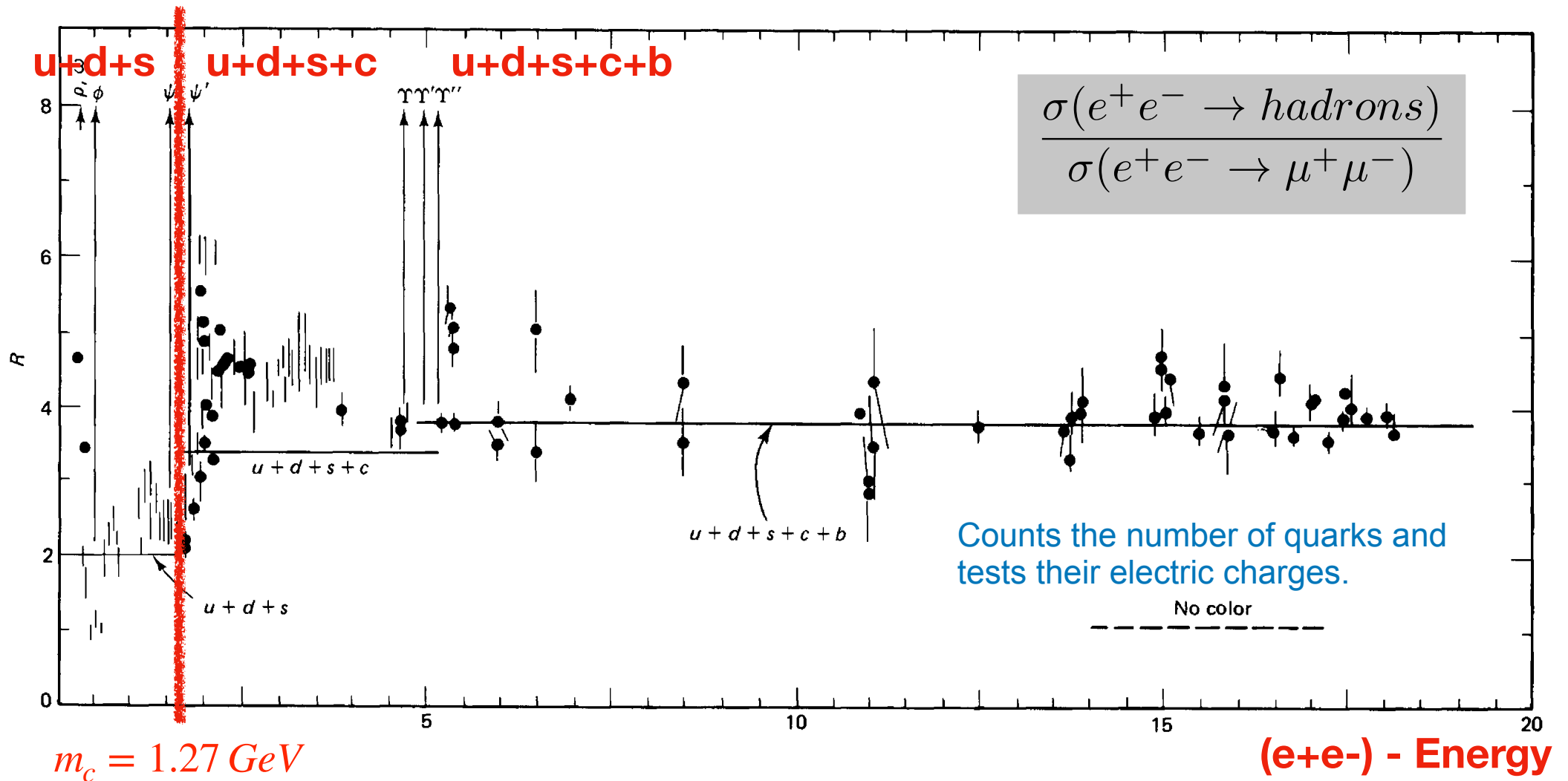
up, charm, top: $+2/3$
down, strange, bottom: $-1/3$



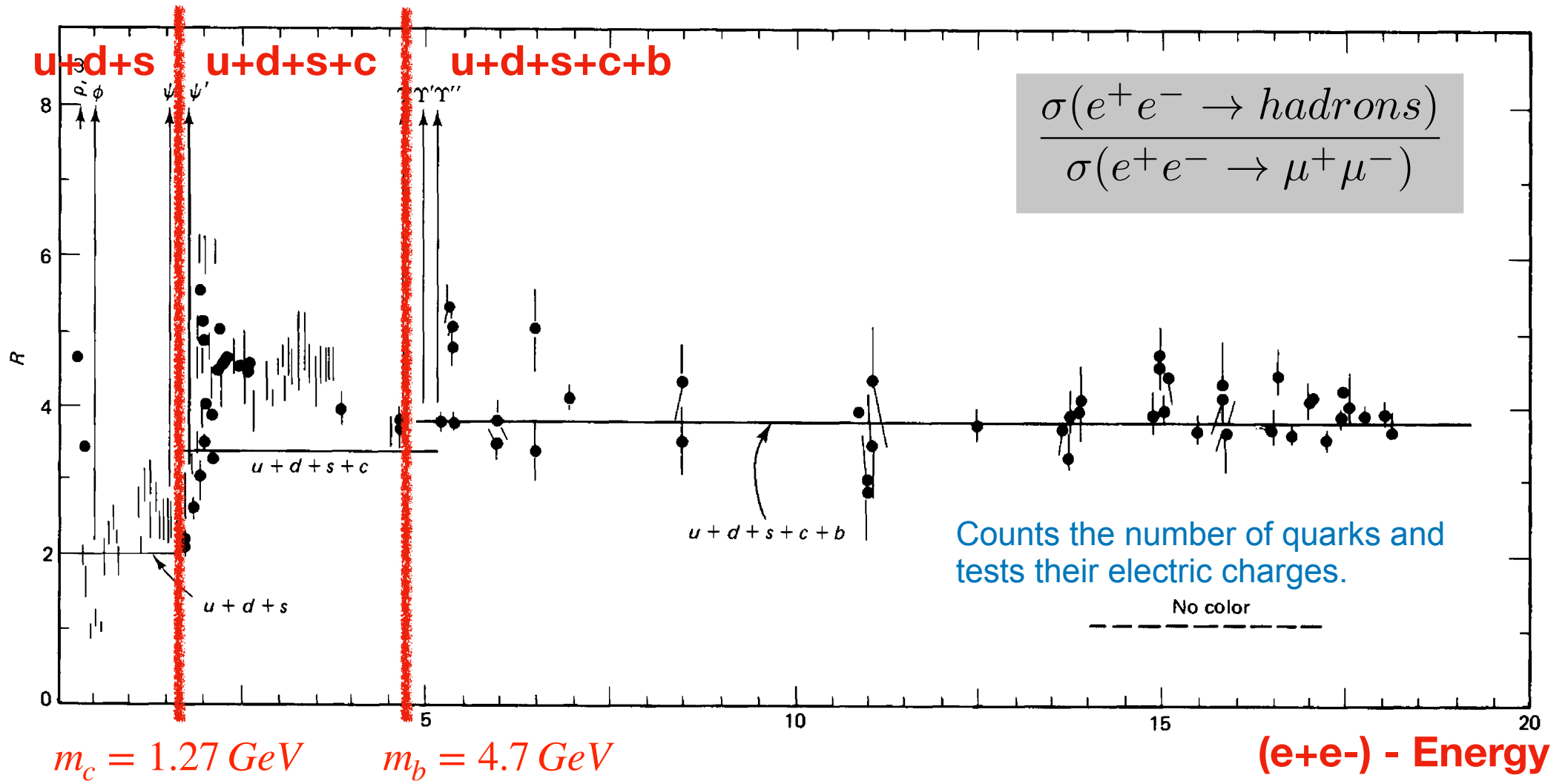
Testing the quark charges and N_c



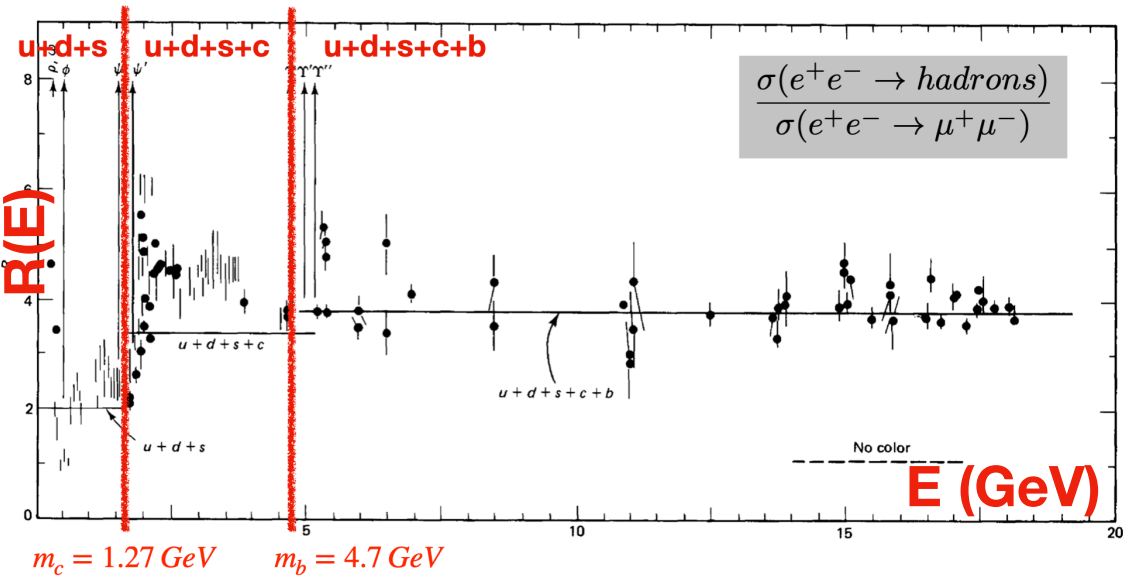
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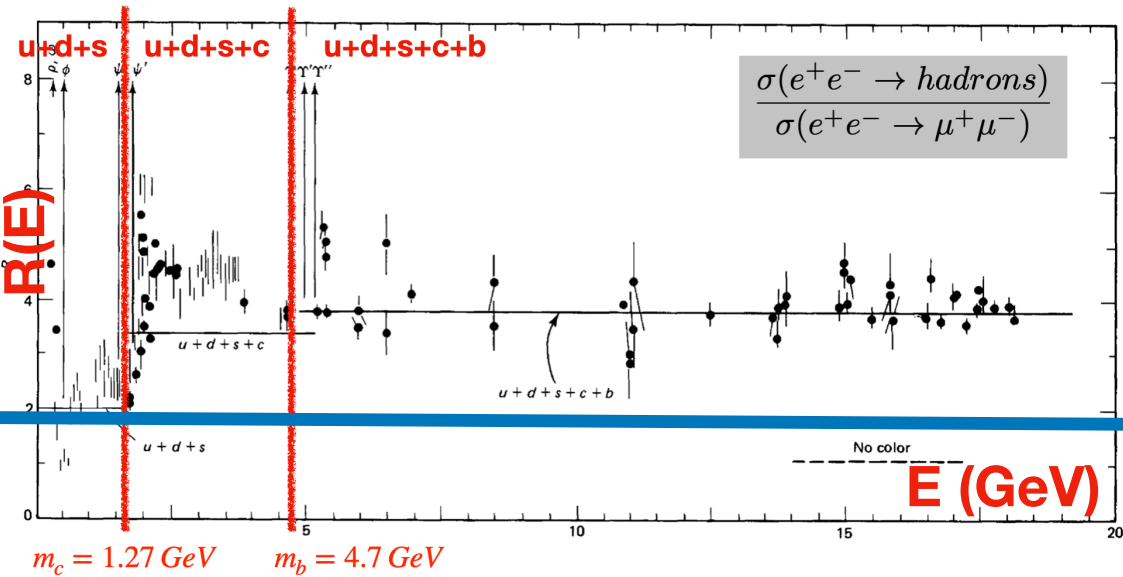


Testing the quark charges and N_c



$$R(E) = \sum_{m_i < E} Q_i^2 = \text{“sum over (quark charges)}^2 \text{ accessible at energy } E\text{”}$$

Testing the quark charges and N_c

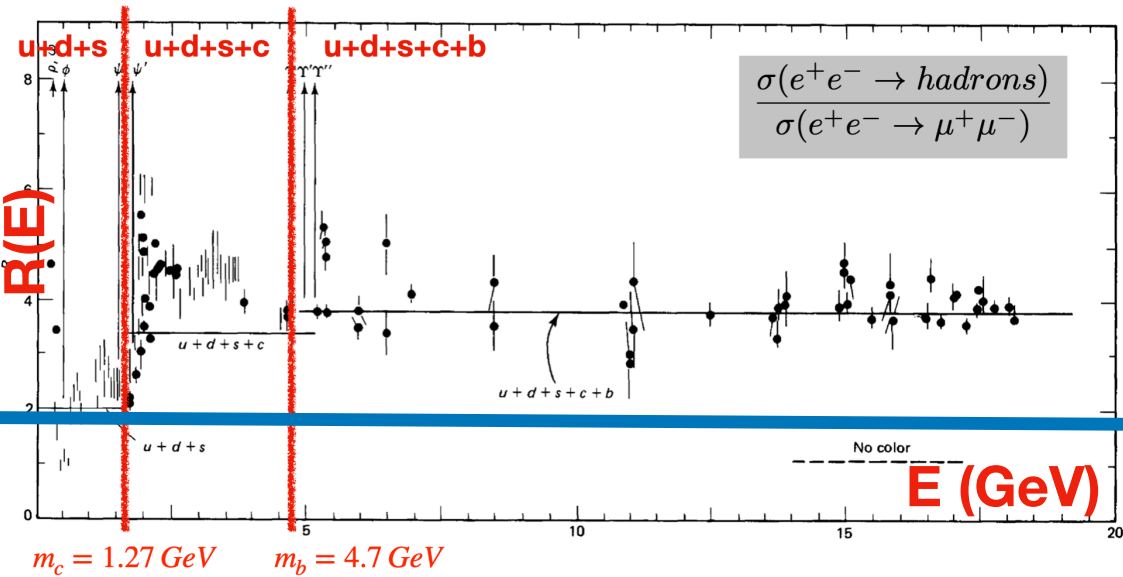


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2

$$R(E = 1 \text{ GeV}) = 3 \times \left(\left(\frac{2}{3} \right)^2 + \left(-\frac{1}{3} \right)^2 + \left(-\frac{1}{3} \right)^2 \right) = 2$$

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2

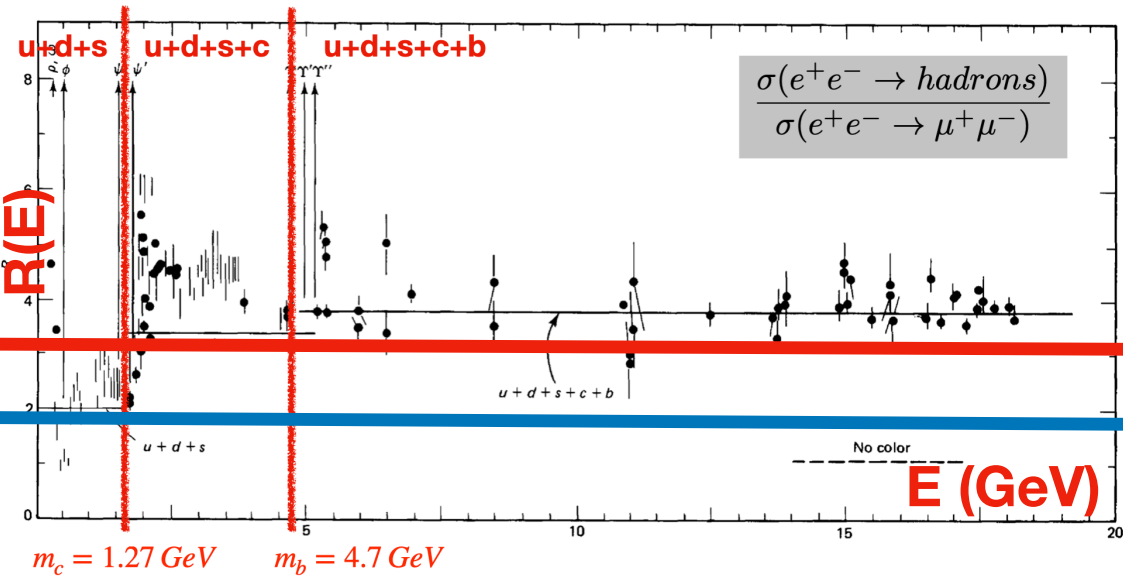
q_u q_d q_s

because there are 3 colors

$$R(E = 1\text{ GeV}) = 3 \times \left(\left(\frac{2}{3}\right)^2 + \left(-\frac{1}{3}\right)^2 + \left(-\frac{1}{3}\right)^2 \right) = 2$$

$SU(3)_c$

Testing the quark charges and N_c



$$R(E) = \sum_{m_i < E} Q_i^2 = \text{“sum over (quark charges)}^2 \text{ accessible at energy } E\text{”}$$

3.33

2

q_u

q_d

q_s

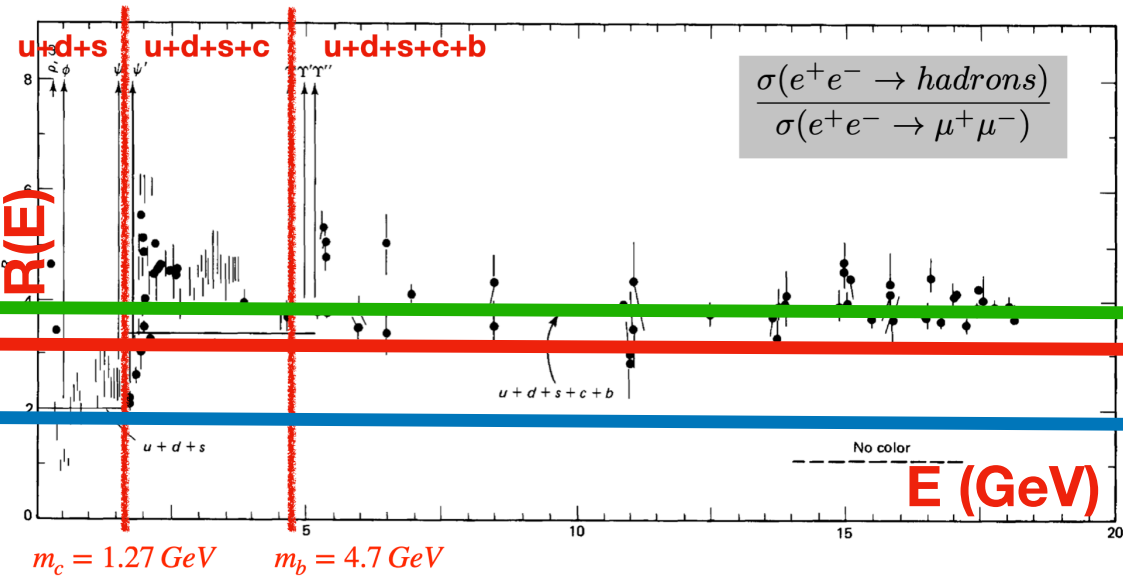
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$SU(3)_c$

$$R(1.2 < E < 4.6) = 3 \times \left(\left(\frac{2}{3} \right)^2 + \left(-\frac{1}{3} \right)^2 + \left(-\frac{1}{3} \right)^2 + \left(\frac{2}{3} \right)^2 \right) = 10/3 \approx 3.33$$

Testing the quark charges and N_c



$$R(E) = \sum_{m_i < E} Q_i^2 = \text{“sum over (quark charges)}^2 \text{ accessible at energy } E\text{”}$$

$$3.66$$

$$3.33$$

$$2$$

because there are 3 colors

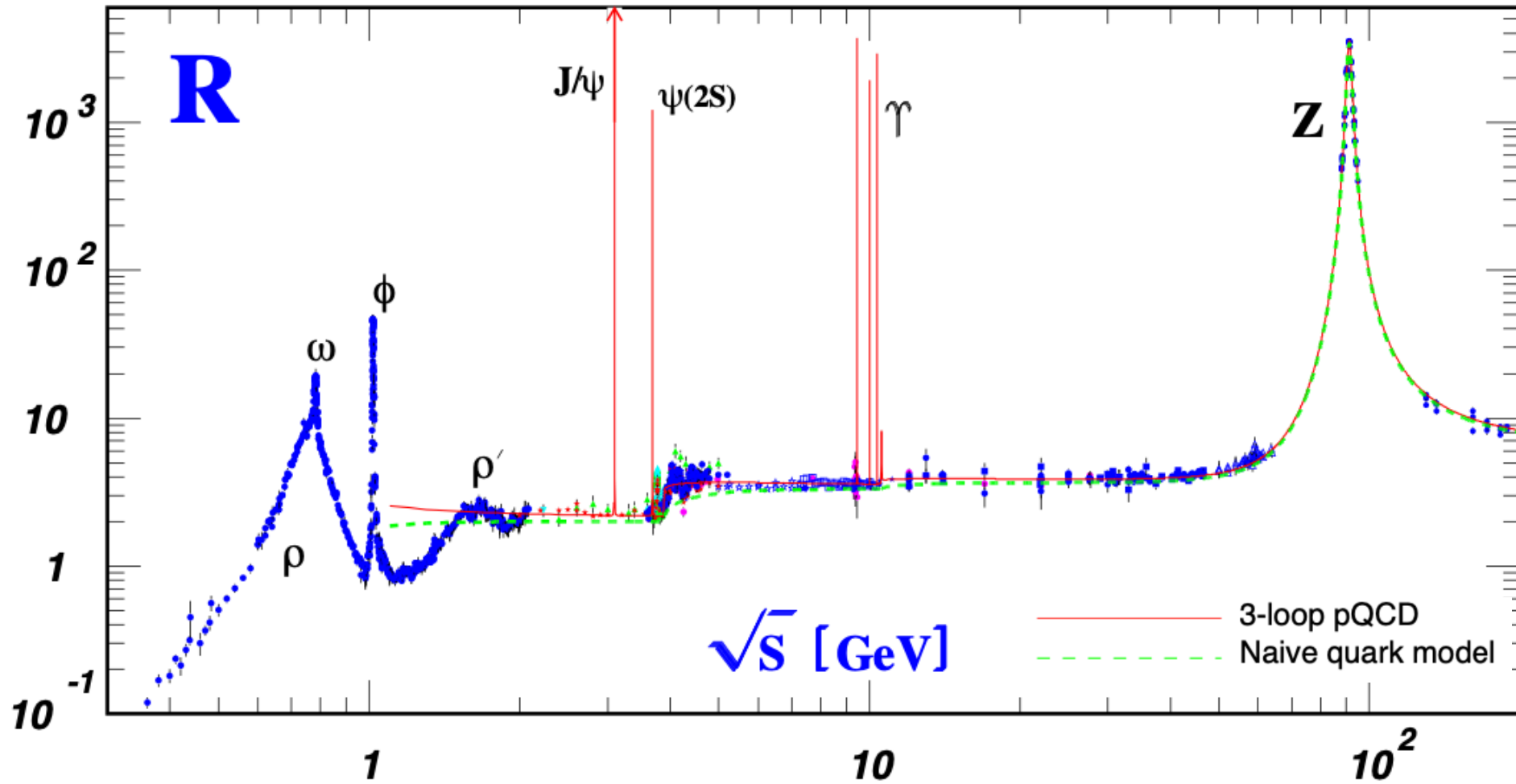
$SU(3)_c$

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$$R(1.2 < E < 4.6) = 3 \times \left(\left(\frac{2}{3}\right)^2 + \left(-\frac{1}{3}\right)^2 + \left(-\frac{1}{3}\right)^2 + \left(\frac{2}{3}\right)^2 \right) = 10/3 \approx 3.33$$

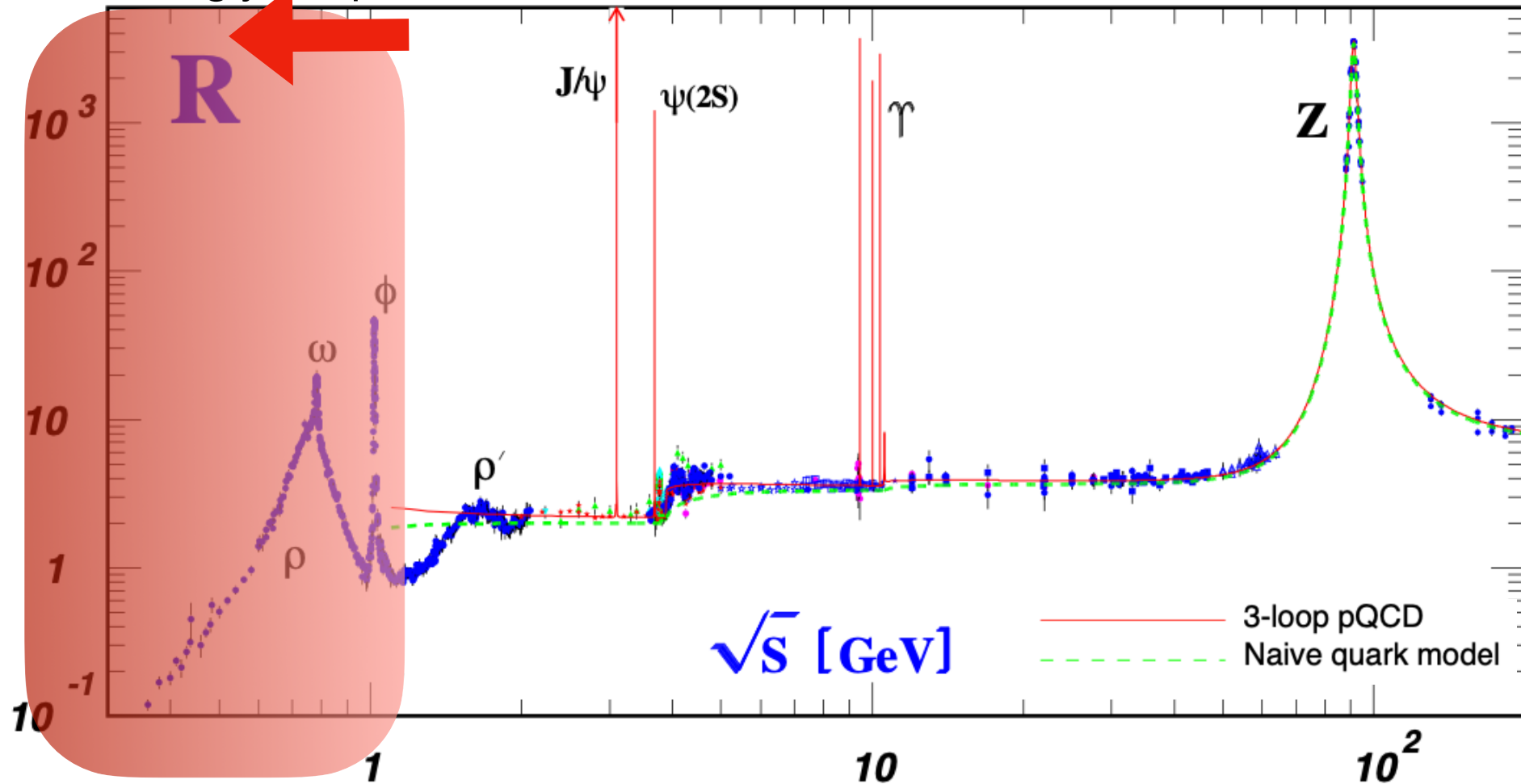
$$R(E > 4.7\text{GeV}) = 3 \times \left(\left(\frac{2}{3}\right)^2 + \left(-\frac{1}{3}\right)^2 + \left(-\frac{1}{3}\right)^2 + \left(\frac{2}{3}\right)^2 + \left(-\frac{1}{3}\right)^2 \right) = 11/3$$

Modern version of the measurement



Modern version of the measurement

QCD strongly coupled



More on QCD at colliders:



10:25

→ 11:20

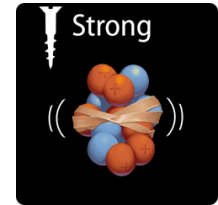
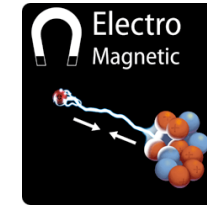
Making Predictions at Hadron Colliders 1/2

Speaker: Alexander Yohei Huss (CERN)

SM without weak interactions

Summary

- Symmetry: $SU(3) \times U(1)_{em}$
- Particles: $u = \mathbf{3}_{2/3}$, $d = \mathbf{3}_{1/3}$, $e = \mathbf{1}_{-1}$ (per generation)

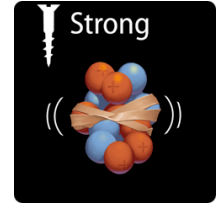
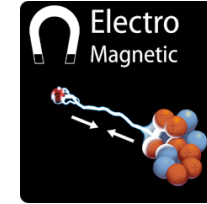


$$\mathcal{L}_{QCD+QED}(g_s, e, m_{u_i}, m_{d_i}, m_{e_i})$$

SM without weak interactions

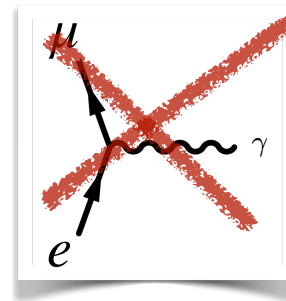
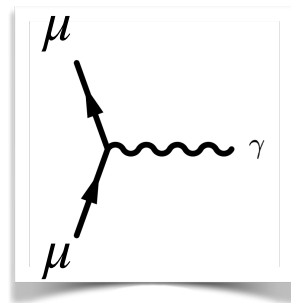
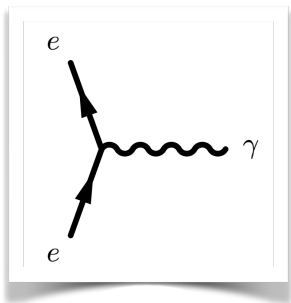
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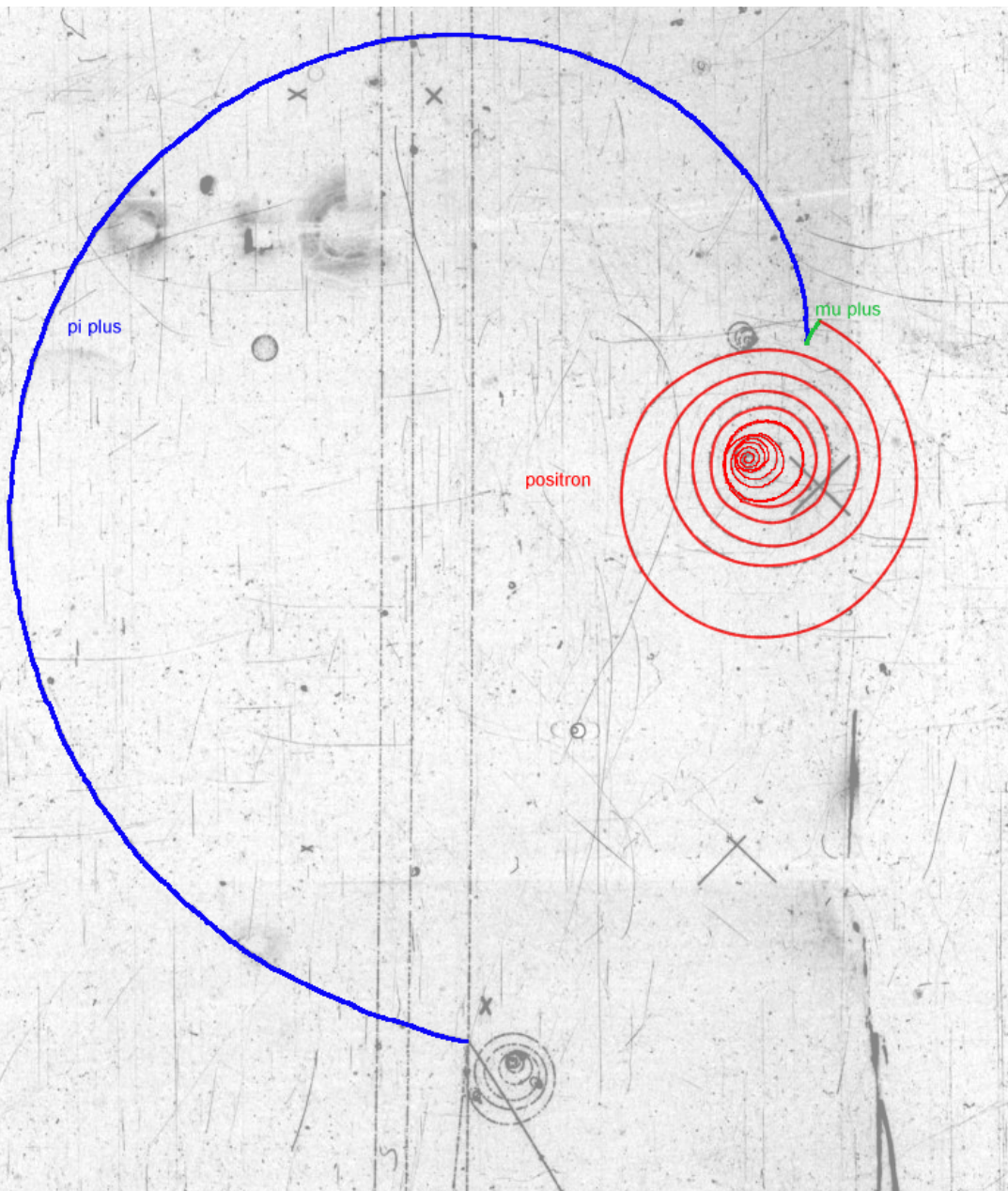


$$\mathcal{L}_{QCD+QED}(g_s, e, m_{u_i}, m_{d_i}, m_{e_i})$$

- One tiny problem: no way to violate individual quark or lepton number!



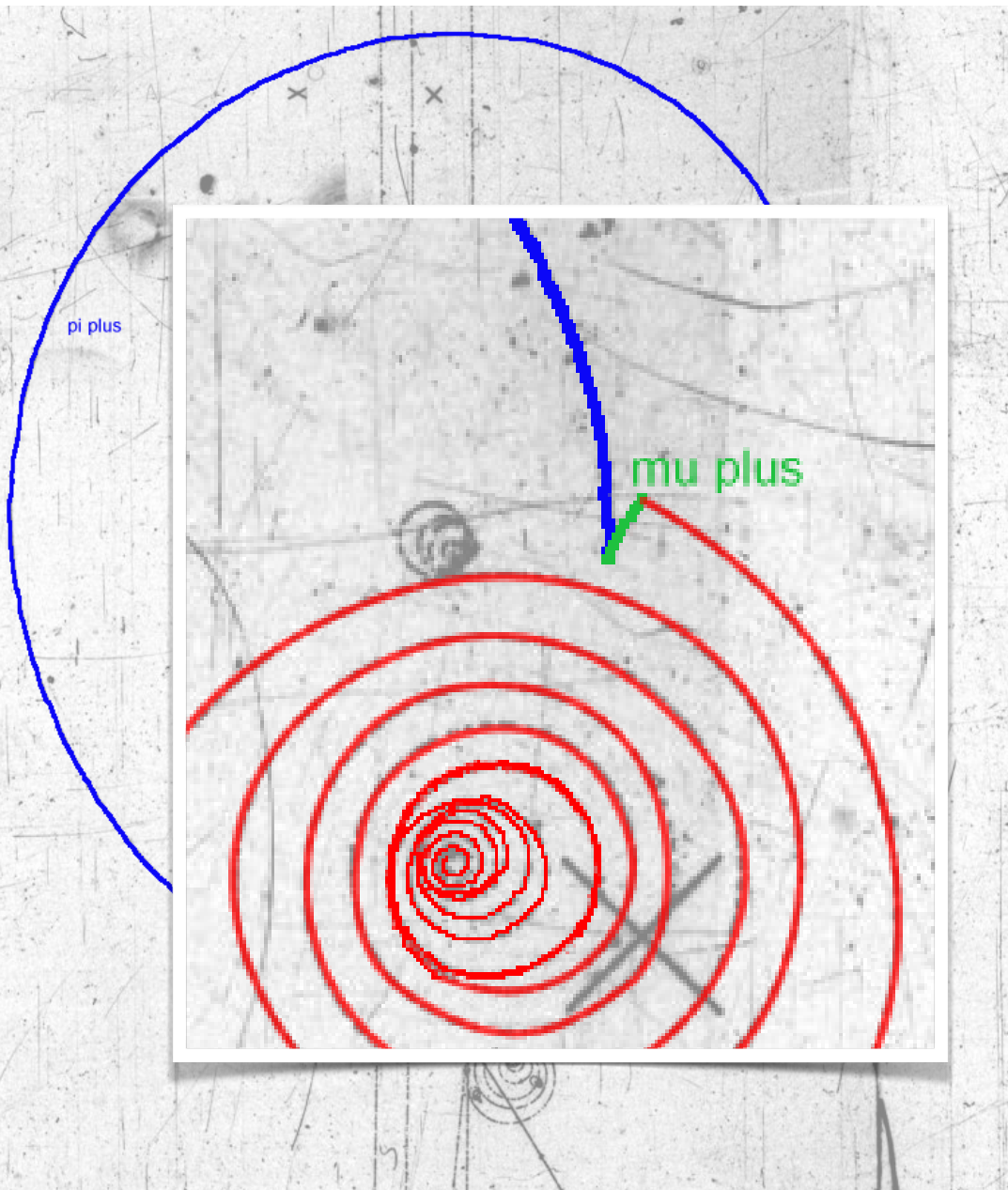
But: the muon decays!



CERN 2 metre hydrogen bubble chamber, exposed to a 10 GeV/c K^+ beam (from top of the picture).

Example of a **pion** stopping and then decaying into a **muon**. This a 'two-body decay' - the muon is accompanied by a neutrino moving in the opposite direction with equal and opposite momentum. Energy and momentum conservation force this momentum to be about 30 MeV/c. The **range of a muon** with this momentum in hydrogen is about a centimetre.

At the end of its short range, the muon itself decays into a **positron** which spirals characteristically.



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We are missing an interaction: the weak force!

