

Standard Model 3/4

Andreas Weiler (TU Munich)

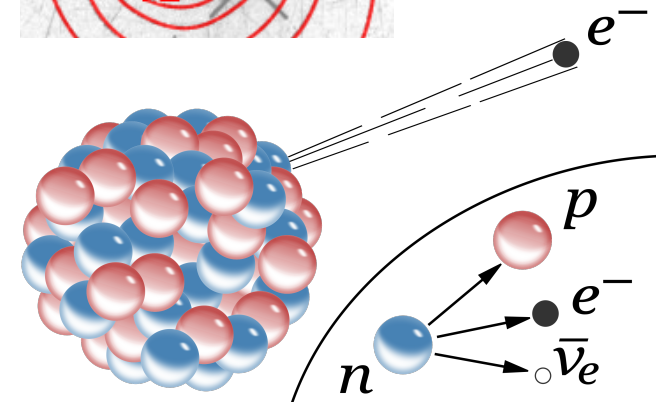
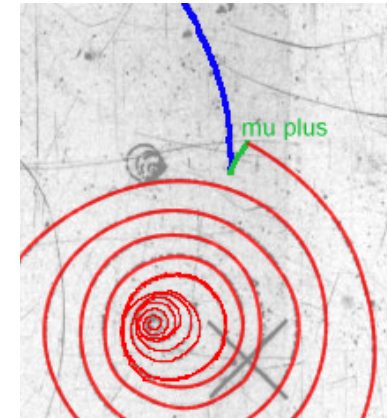
CERN, 7/2024



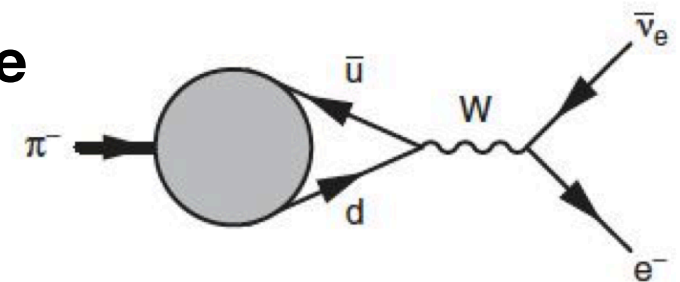
The weak interactions

We want to explain these processes

- muon decay (lifetime: $\tau \approx 10^{-6} s$)
- neutron decay (lifetime: $\tau \approx 877 s$)
- charged pion decay (lifetime: $\tau \approx 10^{-8} s$)



How can one interaction be responsible for such different life-times?



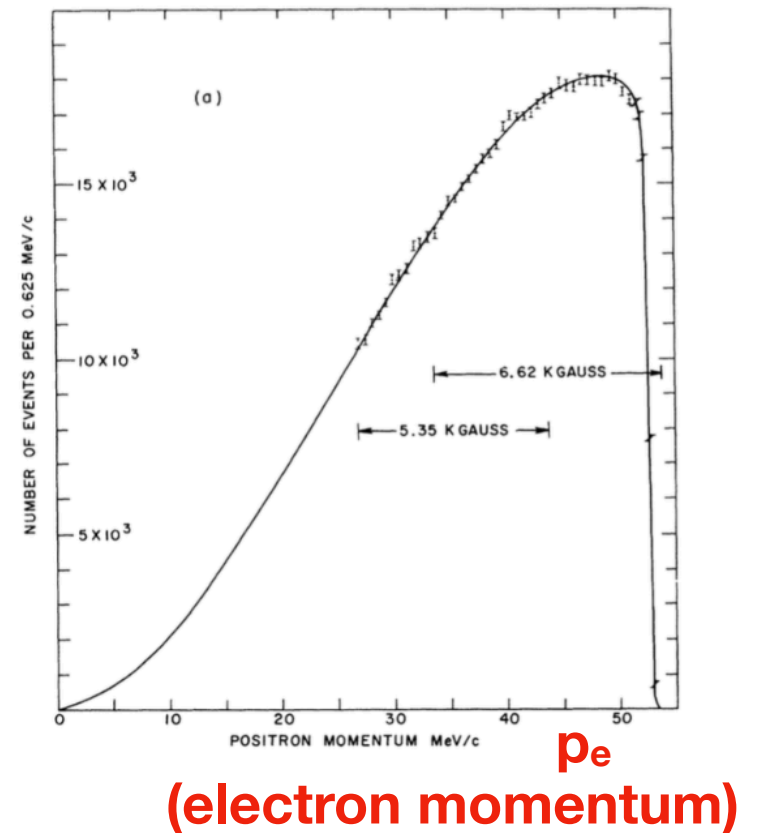
Muon decay

We observe a muon decaying into an electron:

X is something undetected.

On the right you see the electron spectrum
Can X be just one particle?

**No, because in two-body decays
 p_e would be fixed!**



47.4.2. *Two-body decays* :

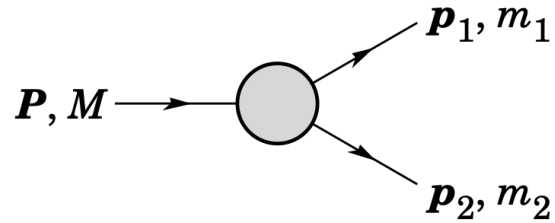


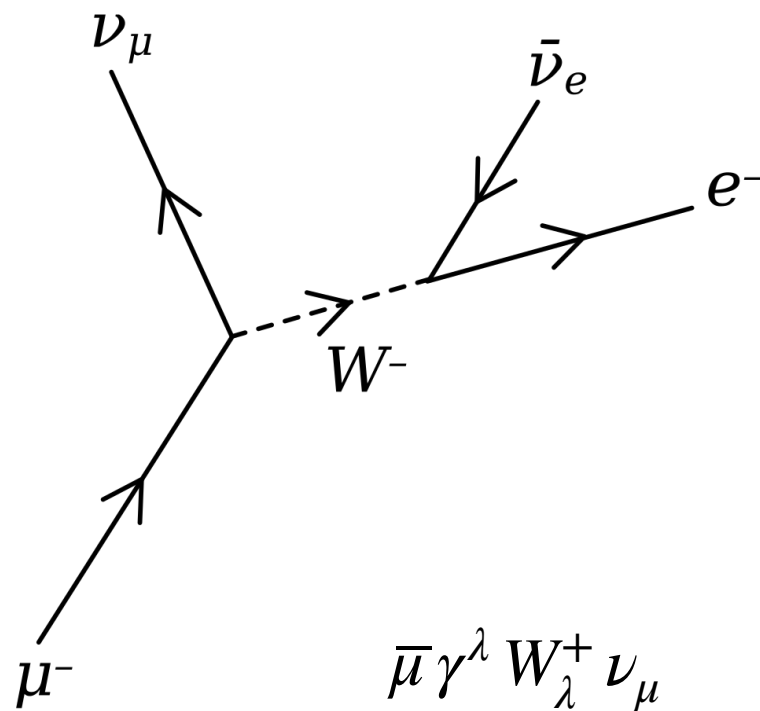
Figure 47.1: Definitions of variables for two-body decays.

In the rest frame of a particle of mass M , decaying into 2 particles labeled 1 and 2,

$$E_1 = \frac{M^2 - m_2^2 + m_1^2}{2M} ,$$

$$|\mathbf{p}_1| = |\mathbf{p}_2| = \frac{[(M^2 - (m_1 + m_2)^2) (M^2 - (m_1 - m_2)^2)]^{1/2}}{2M} ,$$

Muon decay (1st draft)



$$\mu^- \rightarrow e^- + \nu_\mu + \nu_e$$

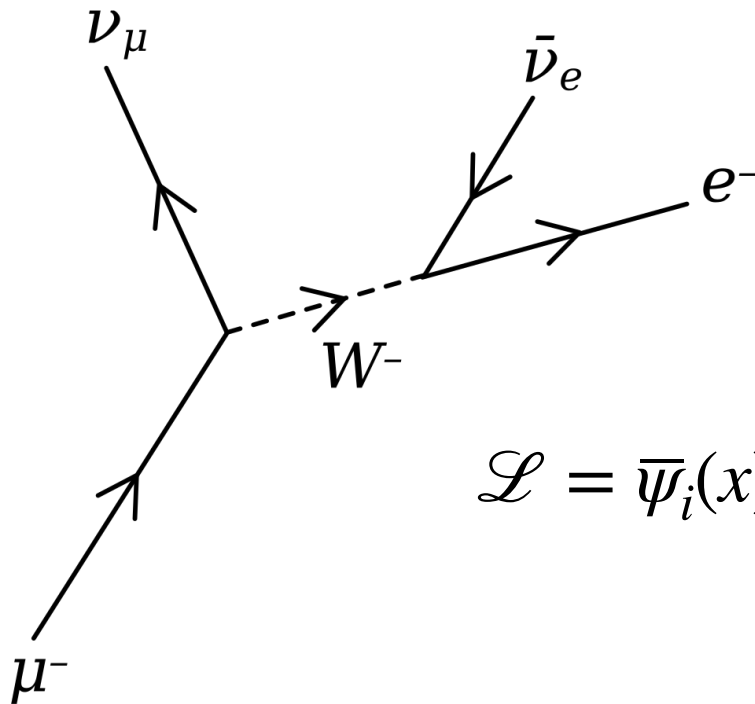
$$\bar{\mu} \gamma^\lambda W_\lambda^+ \nu_\mu$$

$$\begin{pmatrix} \mu \\ \nu_\mu \end{pmatrix}^\dagger \gamma_0 (i\partial_\mu - g \begin{pmatrix} 0 \\ \nu_\mu \end{pmatrix}) \begin{pmatrix} \mu \\ \nu_\mu \end{pmatrix}$$

Muon decay (1st draft)*

* chiral structure will come later.

$$\mu^- \rightarrow e^- + \nu_\mu + \bar{\nu}_e$$



Use non-abelian gauge theory template to describe this:

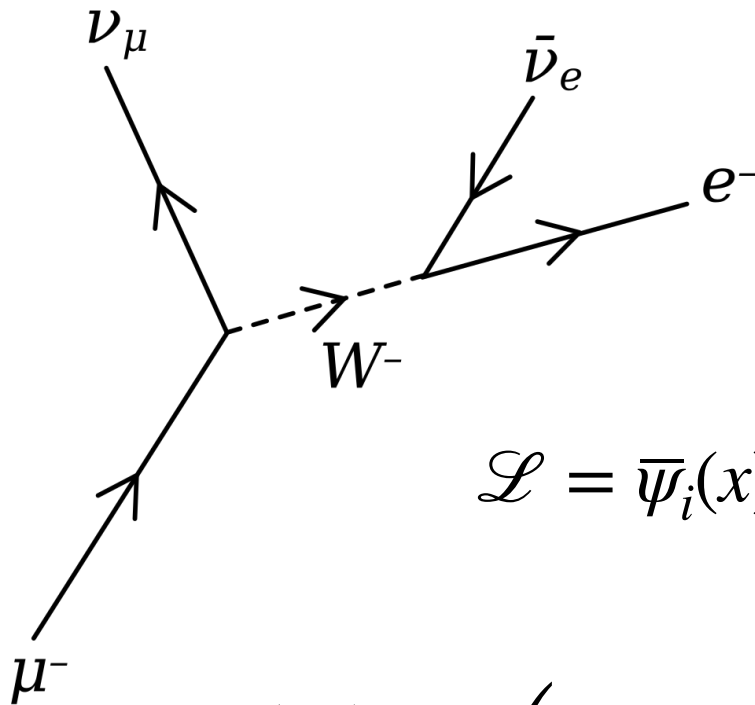
Plan: replace colors (r, g, b) with (ν_μ, μ)

$$\mathcal{L} = \bar{\psi}_i(x) \gamma^\mu (i\partial_\mu - g_s A_\mu^a(x) T_{ij}^a - m) \psi_j(x)$$

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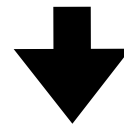
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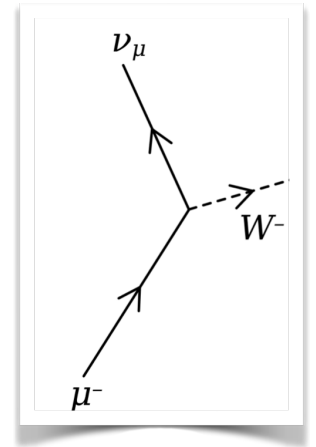
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$$\mathcal{L} = \begin{pmatrix} \bar{\nu}_\mu \\ \bar{\mu} \end{pmatrix} \gamma_L^\lambda \left(i\partial_\lambda - g \begin{pmatrix} 0 & W_\lambda^+ \\ W_\lambda^- & 0 \end{pmatrix} + \dots \right) \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix} = g \bar{\nu}_\mu \gamma_L^\lambda W_\lambda^+ \mu + \dots$$

Muon decay (1st draft)

“Charged current” (the W boson carries electric charge) interactions for the muon and electron

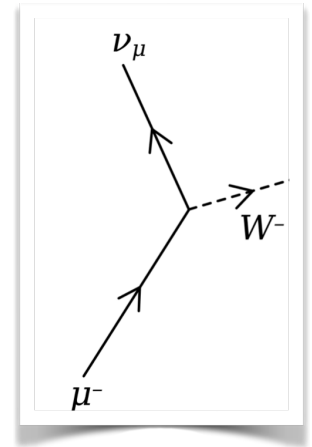


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Since this is a non-abelian theory, this is **universal!**

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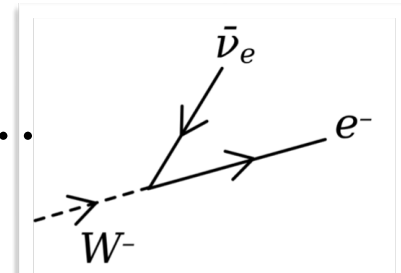
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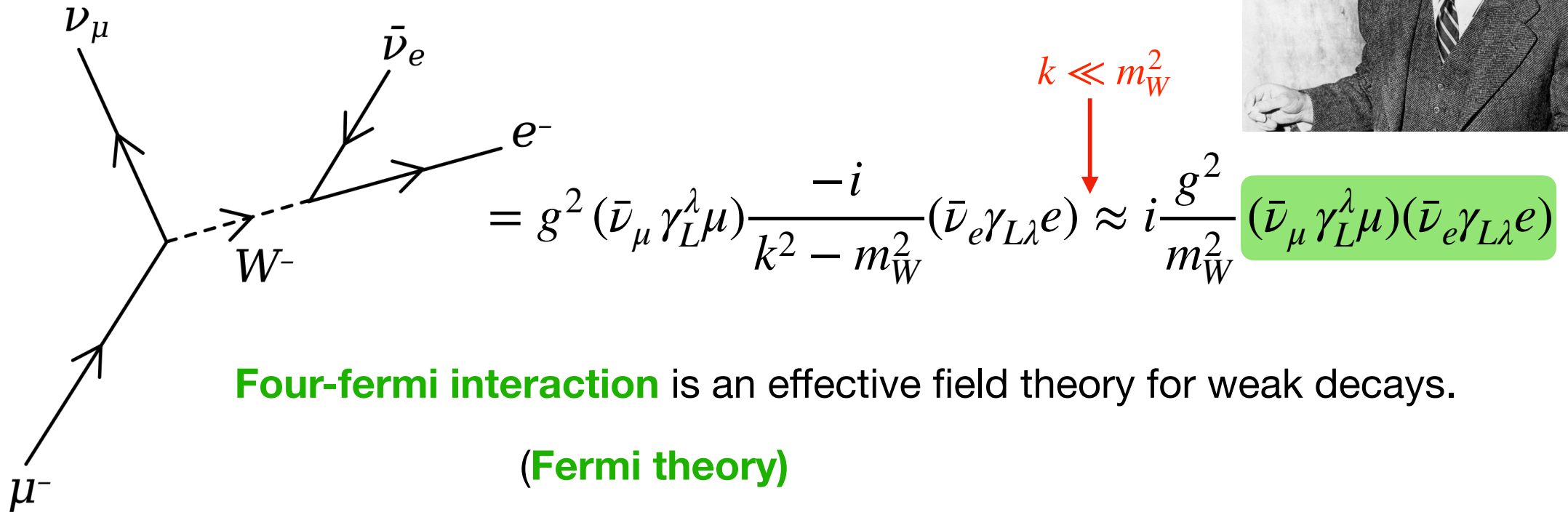
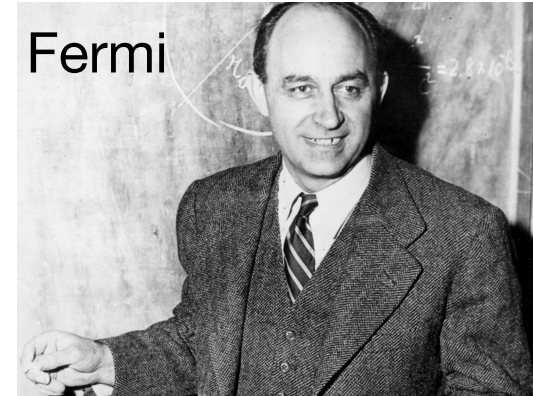
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$$\mathcal{L} = \begin{pmatrix} \bar{\nu}_e \\ \bar{e}^- \end{pmatrix} \gamma_L^\lambda (i\partial_\lambda - g \begin{pmatrix} 0 & W_\lambda^+ \\ W_\lambda^- & 0 \end{pmatrix}) \begin{pmatrix} \nu_e \\ e^- \end{pmatrix} = g \bar{\nu}_e \gamma_L^\lambda W_\lambda^+ e^- + \dots$$



Muon decay “calculation”



Four-fermi interaction is an effective field theory for weak decays.

(Fermi theory)

We can now use *dimensional analysis* to estimate the decay width Γ .

Dimensional analysis estimate of muon life-time

$\Gamma = 1/\tau_{lifetime}$ has dimensions of Energy

$$\Gamma \propto \left| \begin{array}{c} \begin{array}{c} \nu_{\mu} \\ \mu^{-} \end{array} \\ \text{---} \\ W^{-} \\ \text{---} \\ \begin{array}{c} \bar{\nu}_e \\ e^{-} \end{array} \end{array} \right|^2 \approx \frac{g^4}{m_W^4}$$

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Full calculation:

$$\Gamma(\mu) = \left(\frac{g^4}{32m_W^4} \right) \frac{m_\mu^5}{192\pi^3} = G_F^2 \frac{m_\mu^5}{192\pi^3}$$

$= G_F^2$

Muon lifetime:

$$\Gamma \sim 10^{-19} GeV \quad \text{or} \quad \tau = 1/\Gamma \sim 10^{-6} s$$

Muon decay vs. neutron decay

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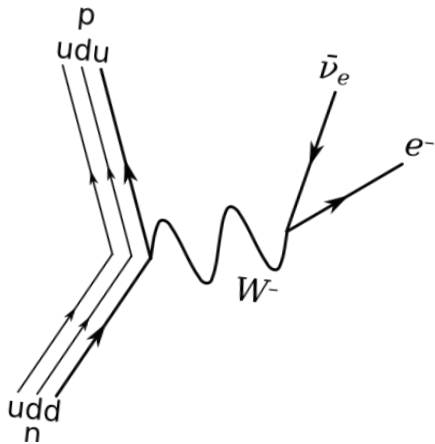
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What replaces m_μ ?

The energy released in neutron decay!



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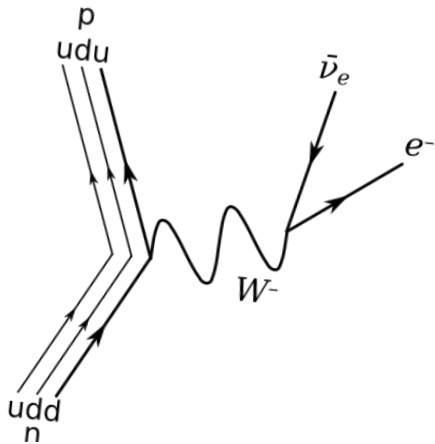
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$$\Gamma(n) \approx \xi G_F^2 \frac{(m_n - m_p)^5}{\pi^3} \approx 10^{-28} \text{ GeV} \quad \tau = 1/\Gamma \sim 10^3 \text{ s}$$

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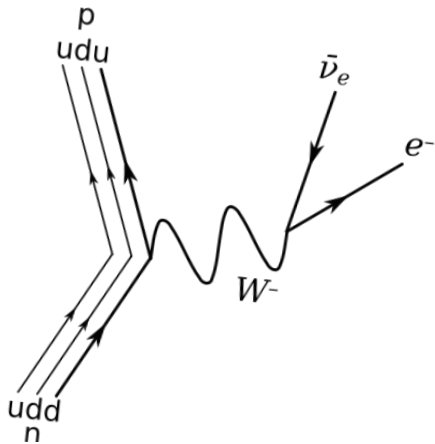
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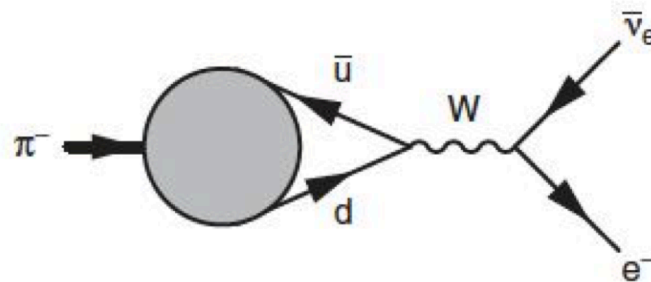
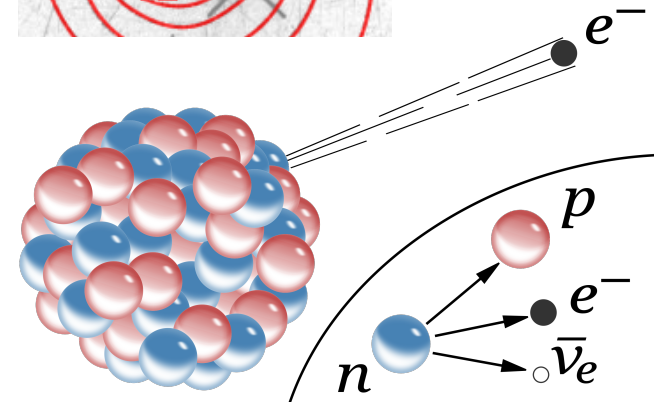
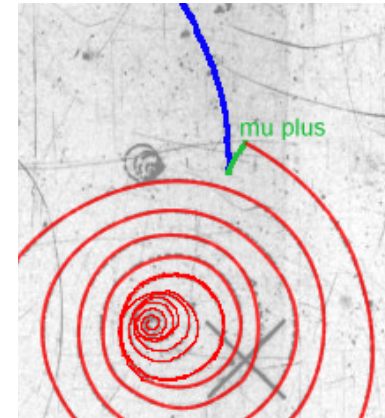
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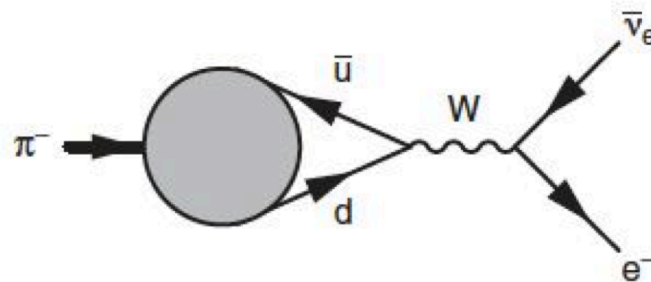
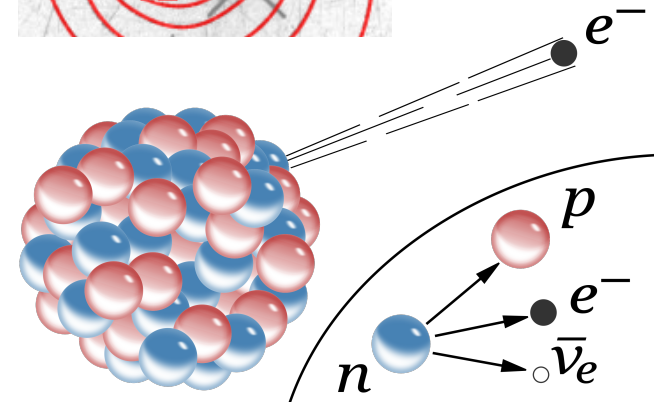
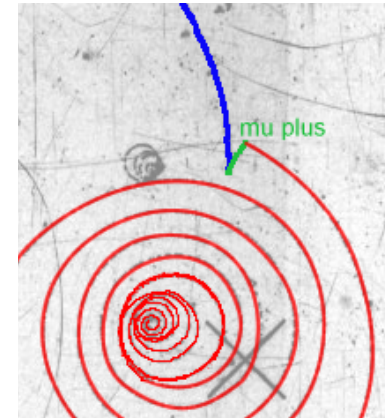
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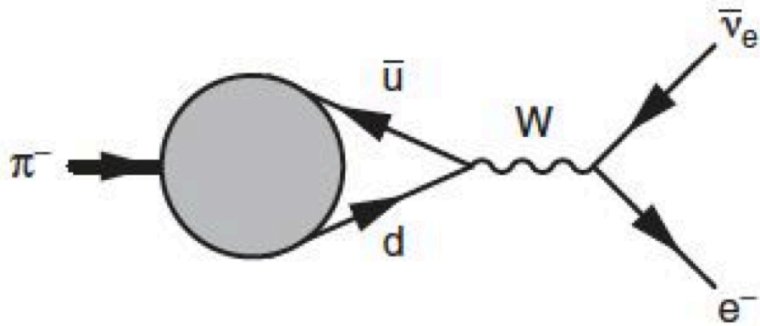
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Charged pion decay

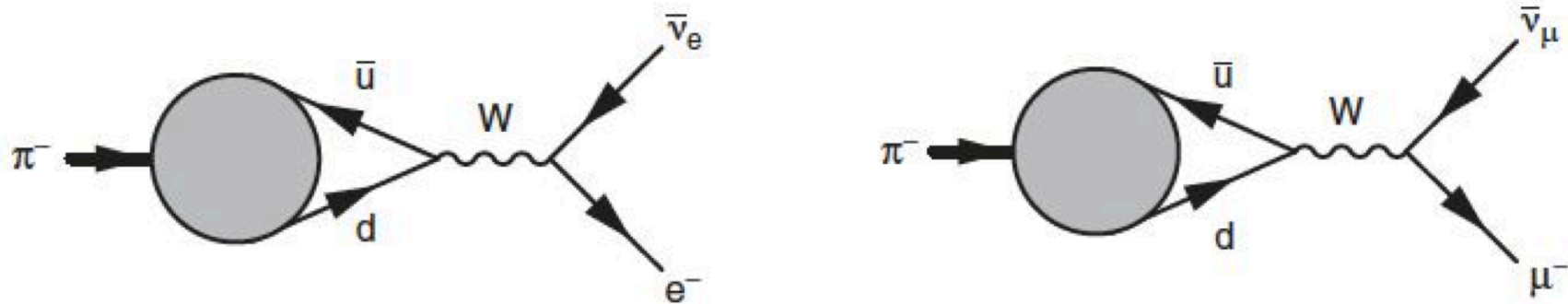


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The paradox: charged pion decay

We measure a lifetime: $\tau \approx 10^{-8} s$

Two contributions:



How important do we expect each one to be?

Using dimensional analysis, we predict

$$\frac{\Gamma(\pi^- \rightarrow e^- \bar{\nu}_e)}{\Gamma(\pi^- \rightarrow \mu^- \bar{\nu}_e)} = \frac{(m_\pi - m_e)^5}{(m_\pi - m_\mu)^5} \sim 10^3 \quad \text{vs. experiment} \quad \frac{\Gamma(\pi^- \rightarrow e^- \bar{\nu}_e)}{\Gamma(\pi^- \rightarrow \mu^- \bar{\nu}_e)} \Big|_{EXP} \approx 10^{-4}$$

Chirality of weak interactions

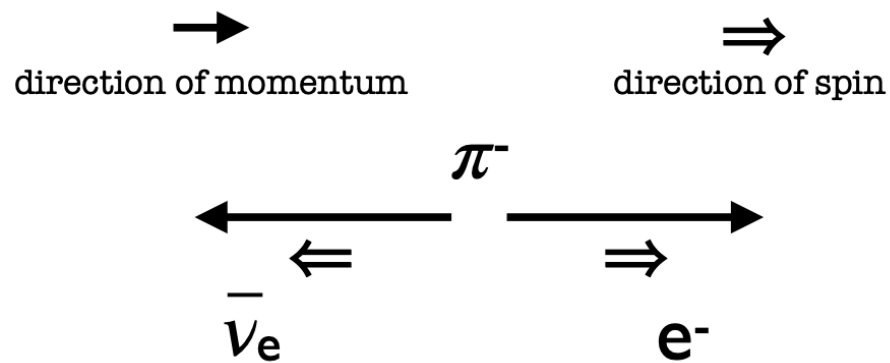
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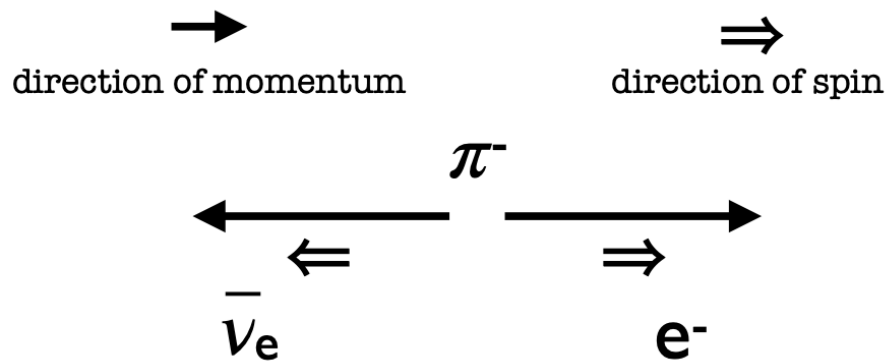
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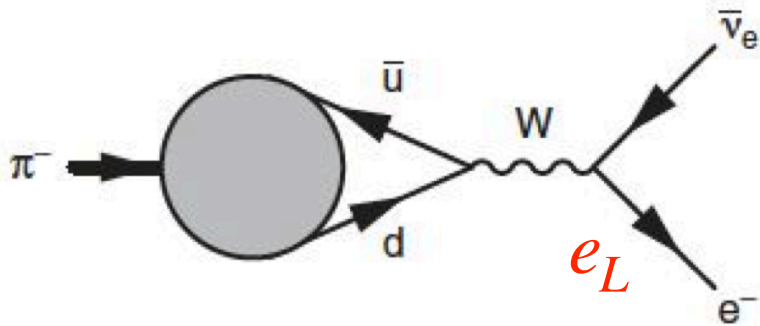
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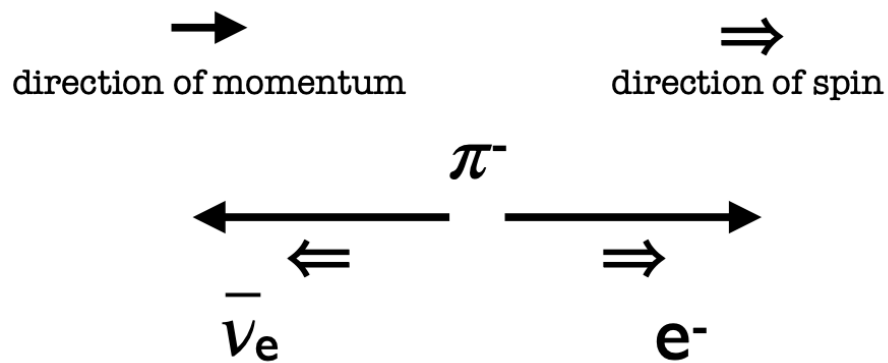
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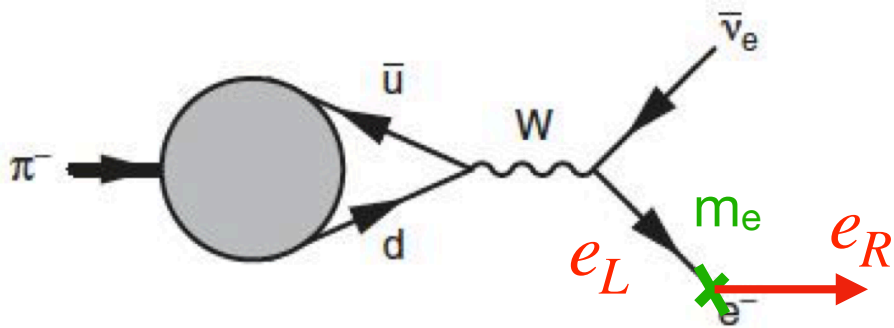
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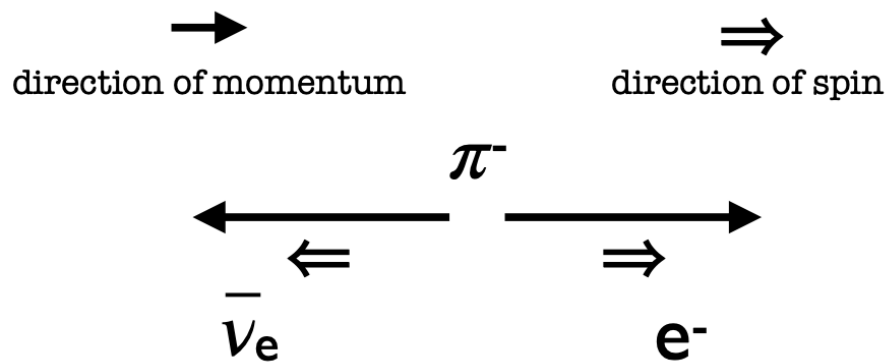
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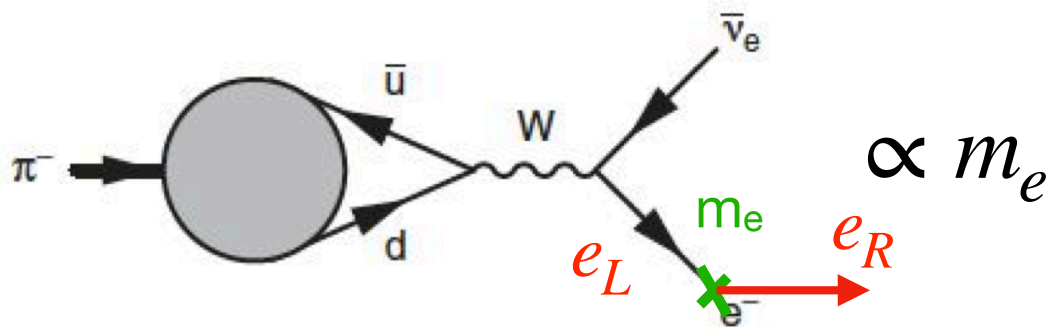
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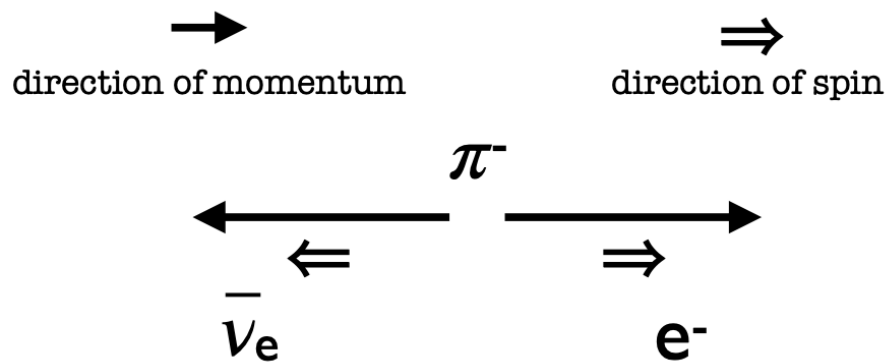
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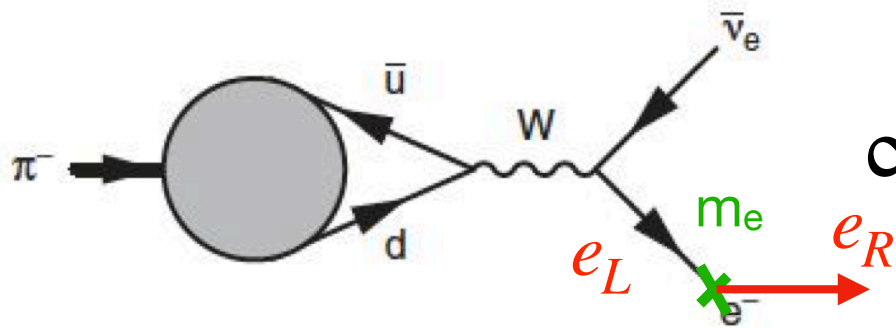
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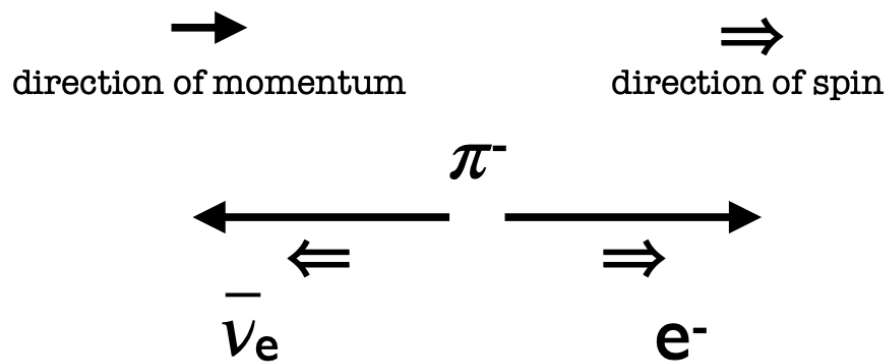
$$\propto m_e \rightarrow \frac{\Gamma(\pi^- \rightarrow e^- \bar{\nu}_e)}{\Gamma(\pi^- \rightarrow \mu^- \bar{\nu}_\mu)} \propto \frac{m_e^2}{m_\mu^2} \sim 2 \times 10^{-5} \sim 10_{\text{obs}}^{-4}$$

↑ phase space factor included

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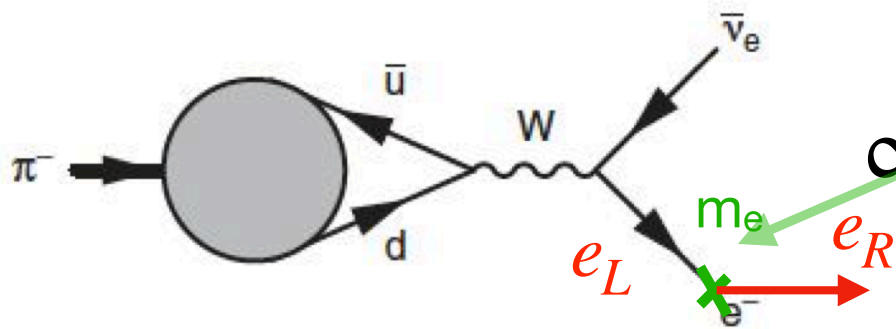
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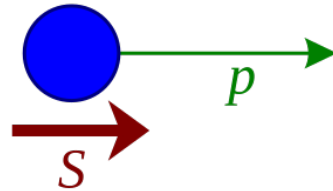
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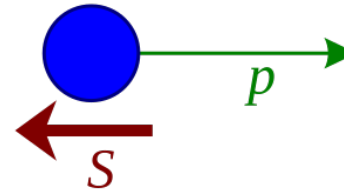
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Chirality and Dirac equation

Right-handed:



Left-handed:



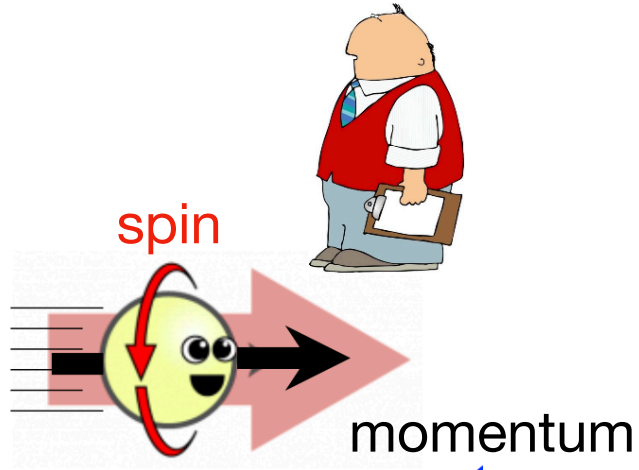
$$\mathcal{L} = \bar{\psi}(x) \gamma^\mu (i\partial_\mu - m) \psi(x)$$

$$\Psi(x) = \Psi_L(x) + \Psi_R(x)$$

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L and R chirality do **not** talk to each other if fermion is **massless (m=0)**!

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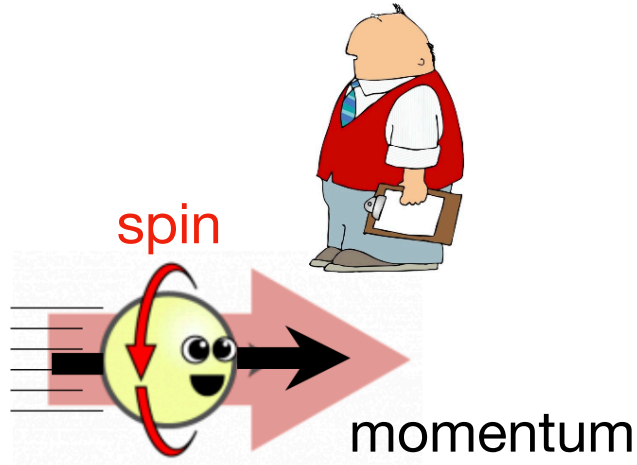


massive right-handed
spinor

(velocity < c)

picture courtesy Gian Giudice

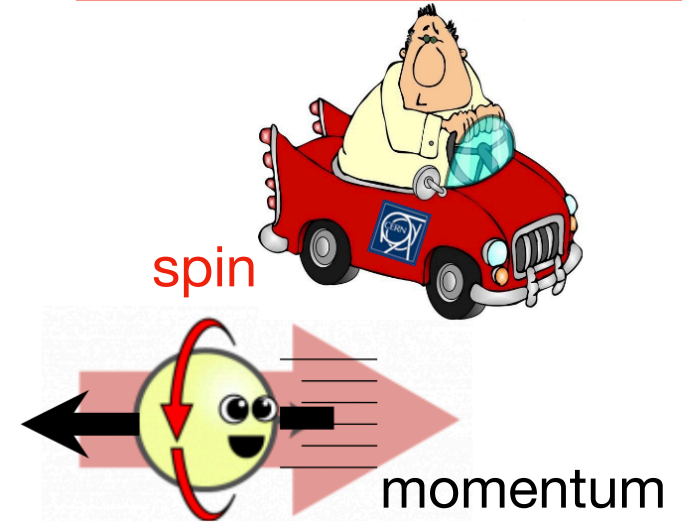
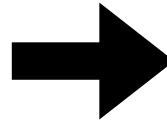
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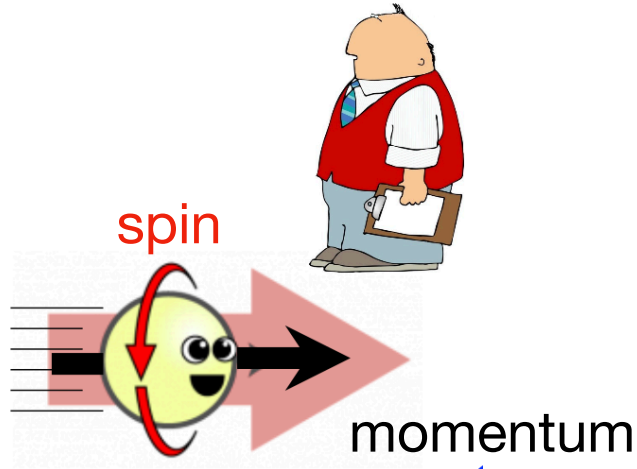
boost to
new frame



right-handed massive
spinor

picture courtesy Gian Giudice

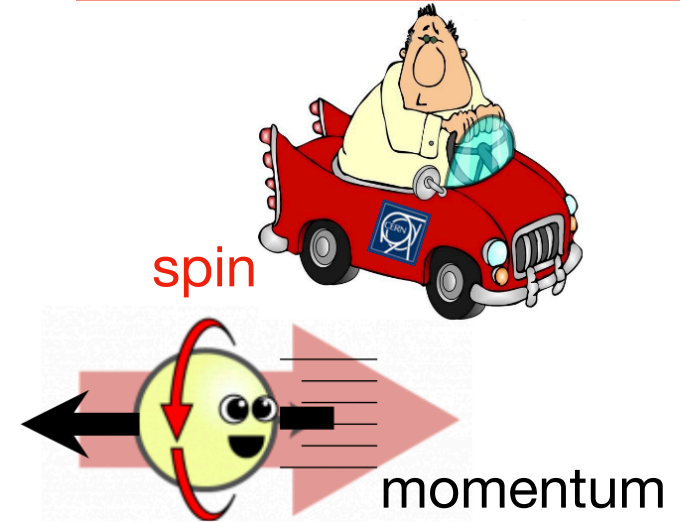
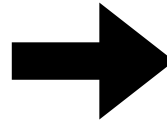
$$\mathcal{L}_{\text{Dirac}} = \bar{\psi}_L \gamma^\mu \partial_\mu \psi_L + \bar{\psi}_R \gamma^\mu \partial_\mu \psi_R + m(\bar{\psi}_L \psi_R + \bar{\psi}_R \psi_L)$$



massive right-handed
spinor

(velocity < c)

boost to
new frame



right-handed massive
spinor

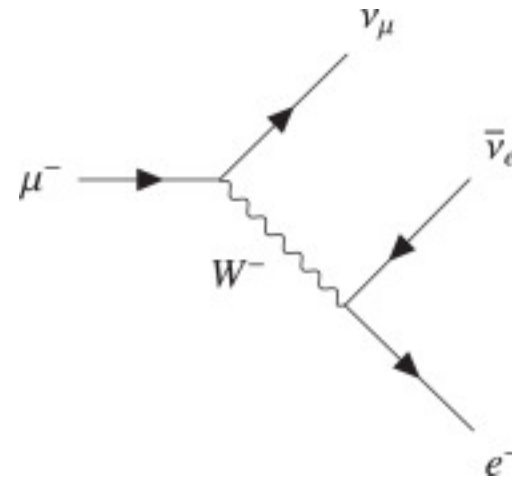
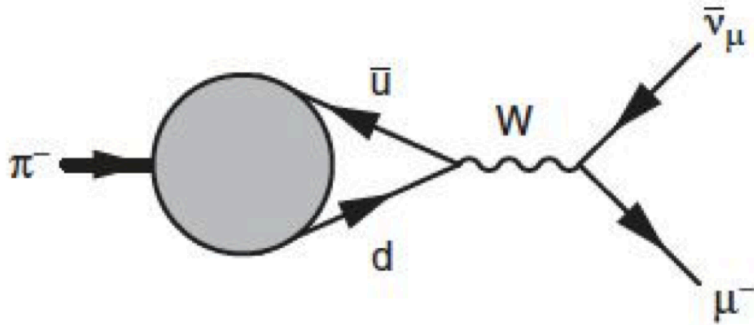
Conclusion: If our theory is able to distinguish left-handed and right-handed fermions, then we must have $m_e = 0$ (or give up on theory?).

picture courtesy Gian Giudice

Muons are polarized

Clever idea by Lederman/Garwin (57)

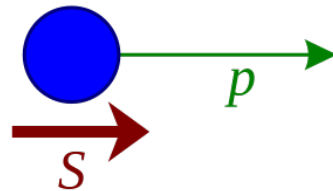
Decay pion to muon and investigate the muon decay products



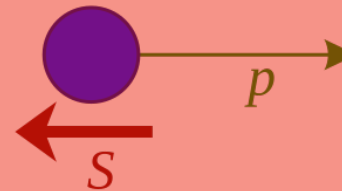
The angular distribution of the electrons reveals the chirality of the interaction.

Chirality and handedness

Right-handed:



Left-handed:



The weak interactions only talk to left-handed particles!

Historically found this in decay of Cobalt (C.S. Wu '56)



Weak interactions: neutral boson

W_3 is charge neutral

Let's extend this to SU(2)

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$
$$\sigma^\pm = \sigma^1 \pm i\sigma^2$$

$$\mathcal{L} = \begin{pmatrix} \bar{u} \\ \bar{d} \end{pmatrix} \gamma_L^\lambda (i\partial_\lambda - g \begin{pmatrix} W_\lambda^3 & W_\lambda^+ \\ W_\lambda^- & -W_\lambda^3 \end{pmatrix}) \begin{pmatrix} u \\ d \end{pmatrix}$$

Could the W_3 be the photon?

Weak interactions: neutral boson

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Could the W_3 be the photon?

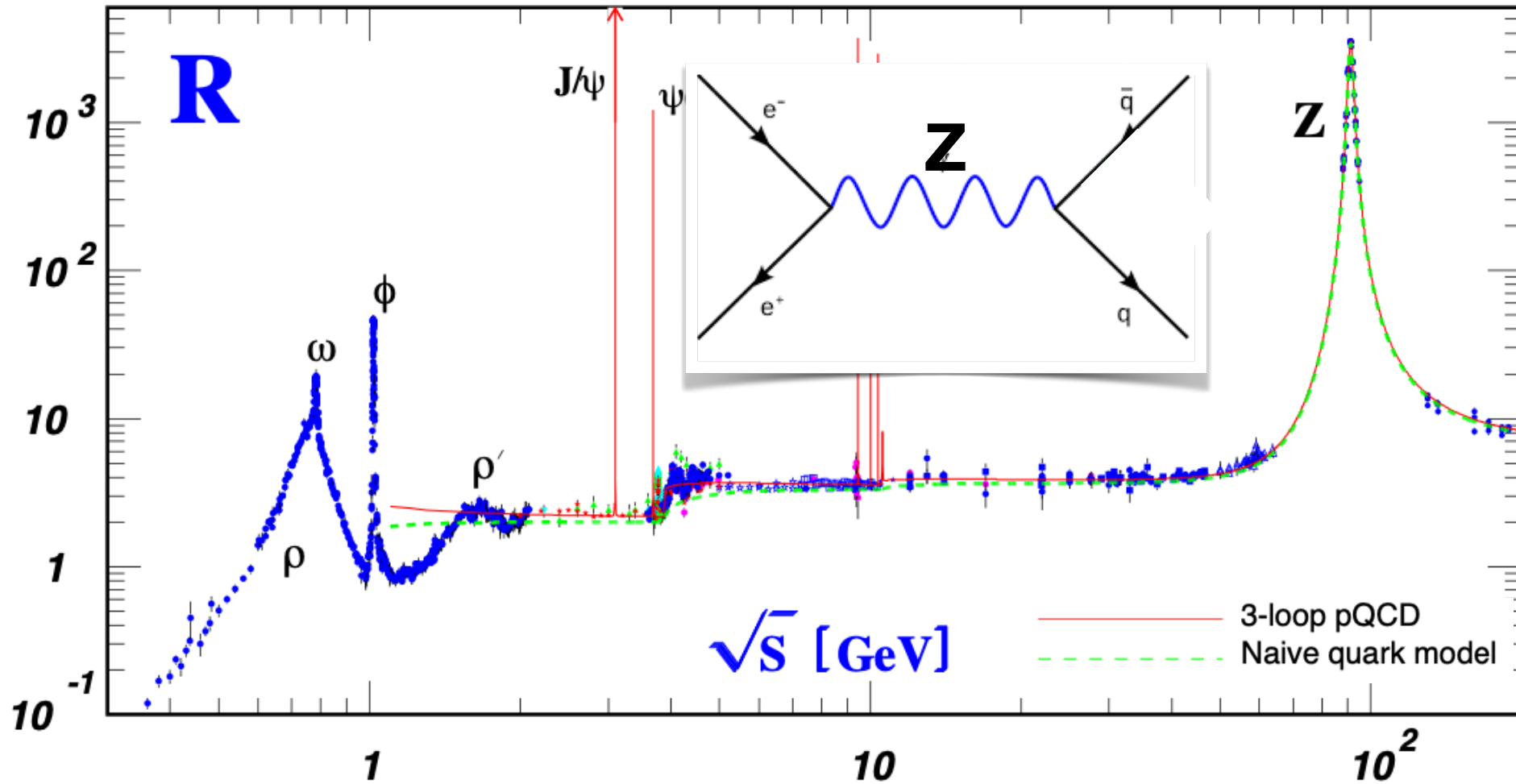
No. Would e.g. predict **+1** charge for up, **-1** charge for down.

New massive neutral gauge boson: this will turn into the Z-boson!

This was also called “a new neutral current process”, discovered at CERN.

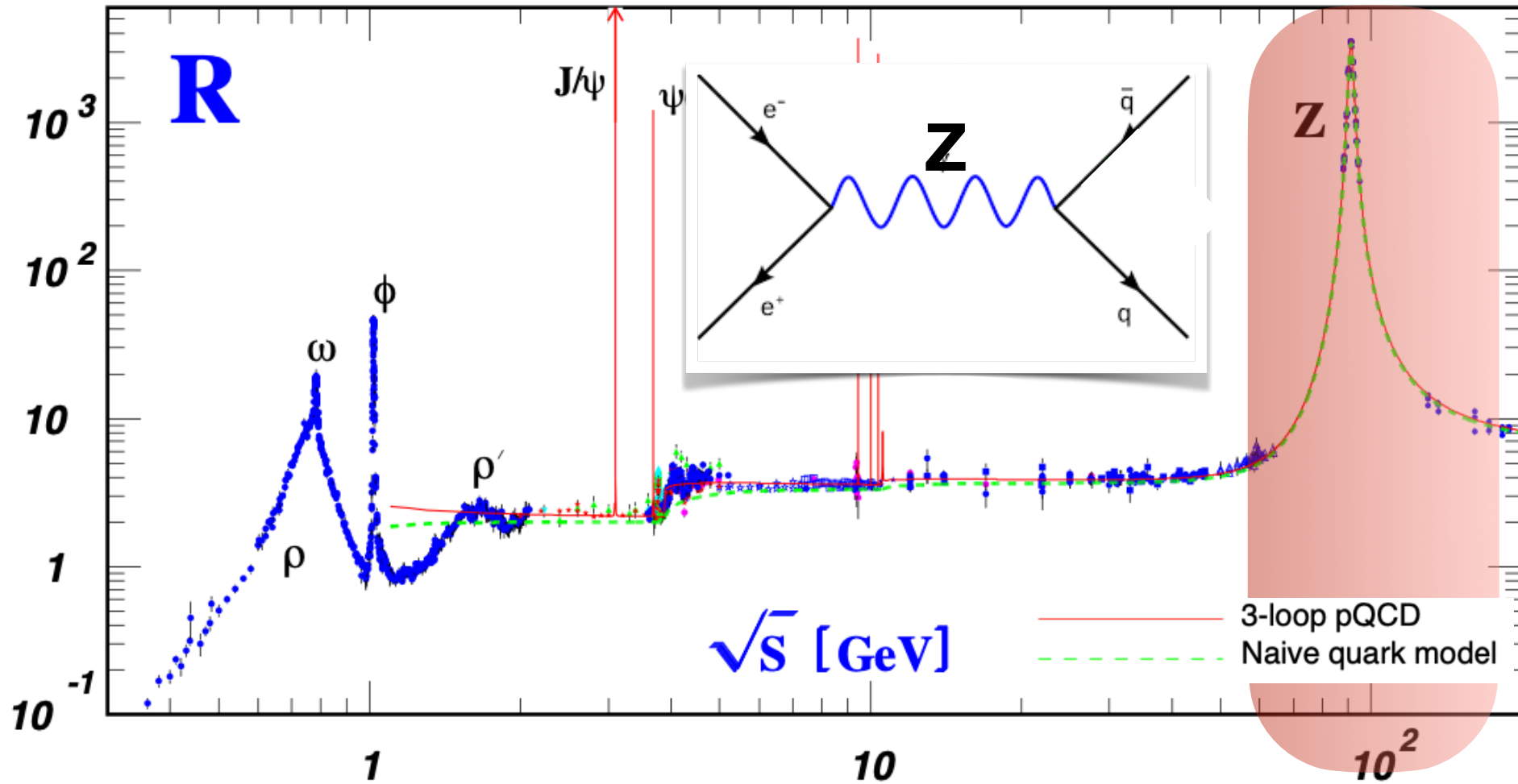
You have already seen the Z boson!

91.2 GeV



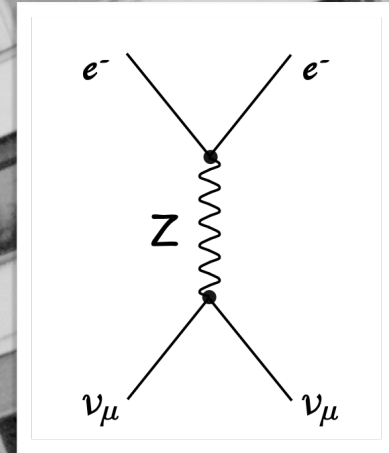
You have already seen the Z boson!

91.2 GeV



Discovered at CERN: Gargamelle bubble chamber

$$\nu_{\mu} e^{-} \rightarrow \nu_{\mu} e^{-}$$



In front of
the CERN
theory
group.

Weak interactions paradoxes

We need non-abelian gauge theory, but gauge bosons need to be massive!
Not SU(2) gauge invariant?

$$M_W^2 W_\mu^+ W^{\mu-} + \frac{1}{2} M_Z^2 Z_\mu Z^\mu$$

Since the weak interactions are chiral (only the **left-handed** particles are charged), the **left-handed** and **right-handed** particles are fundamentally different.
How can they have a mass term?

$$\text{SU}(2) \quad \begin{pmatrix} u \\ d \end{pmatrix}_L \quad u_R \quad d_R$$

vs. $m(\bar{u}_L u_R + \bar{u}_R u_L)$ **not** SU(2) invariant !

Next lecture

