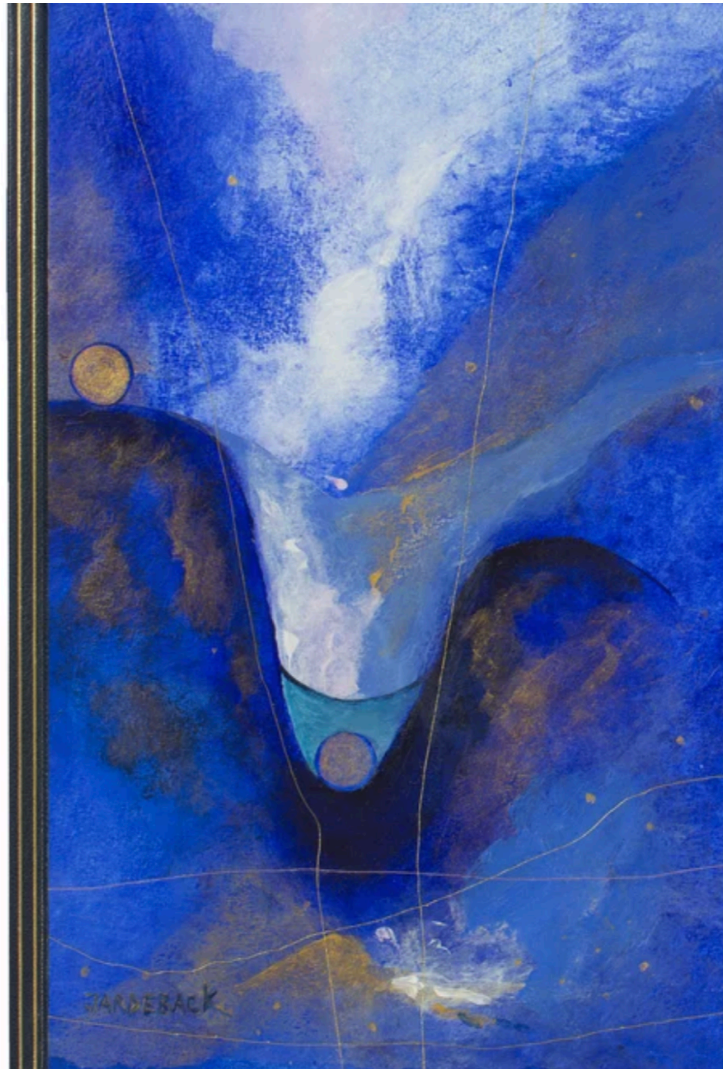


# Standard Model 4/4

Andreas Weiler (TU Munich)

CERN, 7/2024





Kungliga  
Svenska Vetenskapsakademien  
har den 8 oktober 2013 beslutat att med det

## NOBELPRIS

som detta är tillerkännas den  
som inom fysikens område gjort den  
viktigaste upptäckten eller uppfinningen  
gemensamt belöna

*Peter W Higgs*  
och *François Englert*

för den teoretiska upptäckten av en mekanism som  
bidrar till förståelsen av massans ursprung hos  
subatomära partiklar, och som nyligen, genom upp-  
täckten av den förutsagda fundamentala partikeln,  
bekräftats av ATLAS- och CMS-experimenten  
vid CERN:s accelerator LHC.

● STOCKHOLM DEN 10 DECEMBER 2013 ●

*Barbarasam*  *Staff W. Higgs* 

The certificate for 2013 physics laureate Peter Higgs — who shared the prize with Francois Englert — with original artwork by Susanne Jardeback.

Lovisa Engblom/Courtesy of and copyright (c) The Nobel Foundation

Susanne Jardeback's work on the certificate for 2013 physics laureate Peter Higgs -- a collage inspired partly by the divine, and partly by a **sombrero**.

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Barbara Larsson  Staff W. Higgs 

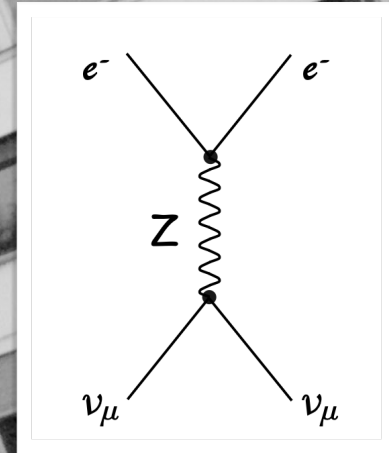
The certificate for 2013 physics laureate Peter Higgs — who shared the prize with Francois Englert — with original artwork by Susanne Jardeback.

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**First: some further discussion on the Gargamelle discovery from last lecture.**

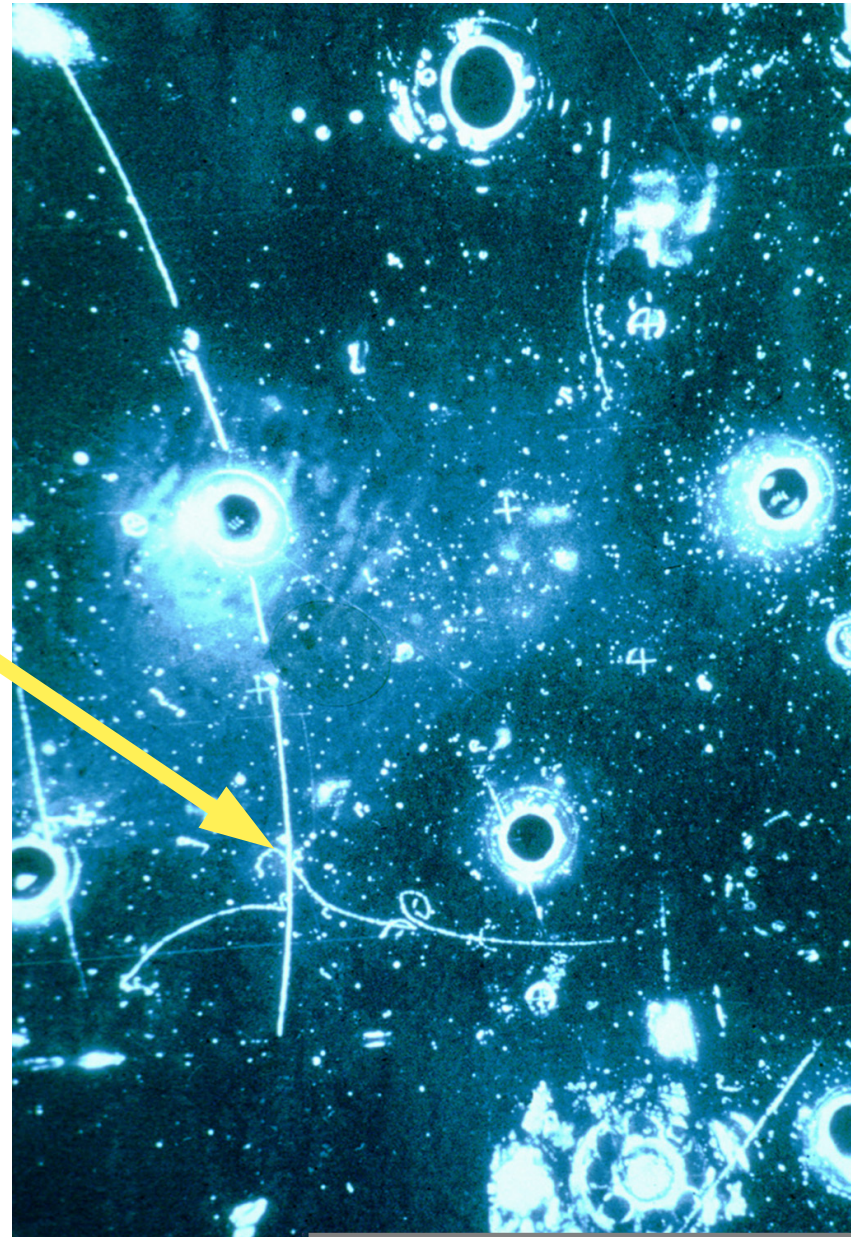
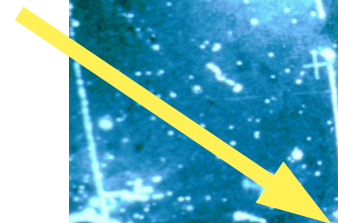
## Discovered at CERN: Gargamelle bubble chamber

$$\nu_{\mu} e^{-} \rightarrow \nu_{\mu} e^{-}$$



In front of  
the CERN  
theory  
group.

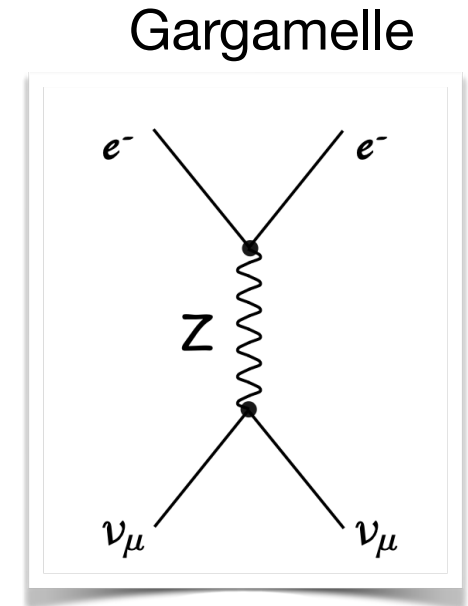
- Obviously weak currents



# Let us understand the Gargamelle experiment

## Shows existence of neutral current

- How did they know it is a muon-neutrino?!



- Why is not a charge current process? (mediated by W-boson)

# Let us understand the Gargamelle experiment

## Shows existence of neutral current

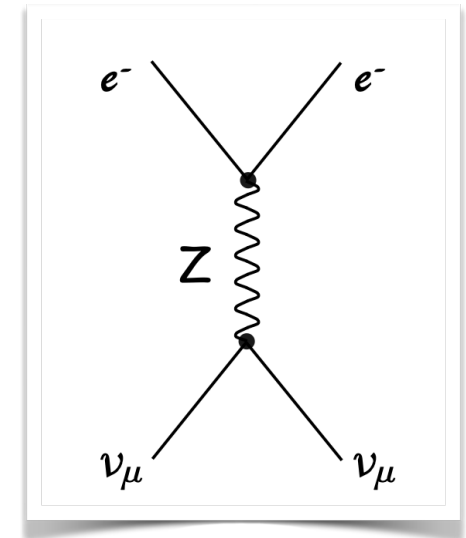
- How did they know it is a muon-neutrino?!

Do you know of a particle that prefers to decay to muon neutrinos? \*\*\*cough\*\*\*  $\pi^+$  \*\*\*cough\*\*\*

$$\frac{\Gamma(\pi^- \rightarrow \mu^- \bar{\nu}_\mu)}{\Gamma(\pi^- \rightarrow e^- \bar{\nu}_e)} \Big|_{EXP} \approx 10^4$$

- Why is not a charge current process? (mediated by W-boson)

Gargamelle





# Let us understand the Gargamelle experiment

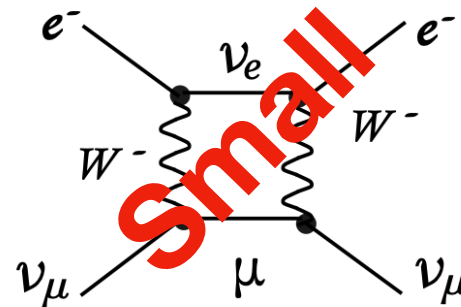
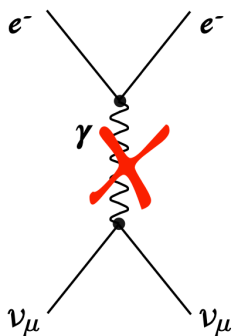
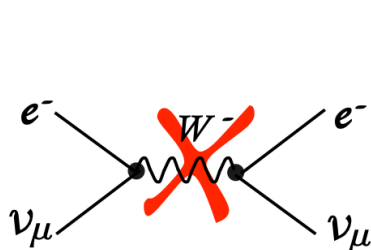
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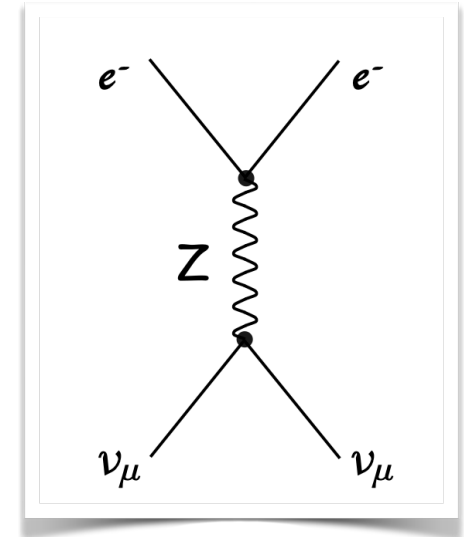
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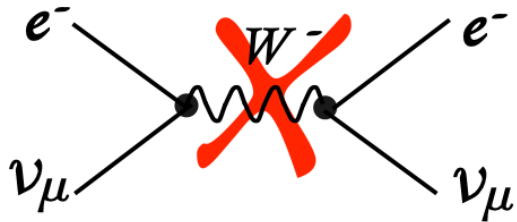
- Why is not a charge current process? (mediated by W-boson)



suppressed since  
higher order in expansion  
 $\sim g^4/(16 \pi^2) \sim 0.001$

Gargamelle





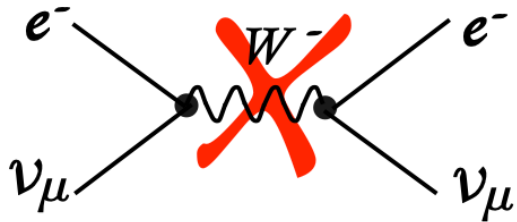
This would jump generations (change the “flavor” from electron to muon)

We know this transition must be possible, since we know that neutrinos change flavor in neutrino oscillations!

But the effect is tiny! It is proportional to the neutrino mass

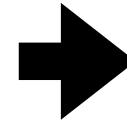
$$m_\nu < 1 \text{ eV} \ll m_e = 511 \text{ keV}$$

(We will see that for quarks these flavor changes can be much bigger!)



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We know this transition must be possible, since we know that neutrinos change flavor in neutrino oscillations!



THURSDAY, 25 JULY

09:15 → 10:10 **Neutrino Physics** ¶

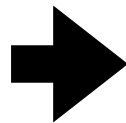
Speaker: Joachim Kopp (CERN)

[QR Code - Neutrino ...](#) [Survey - Neutrino Ph...](#)

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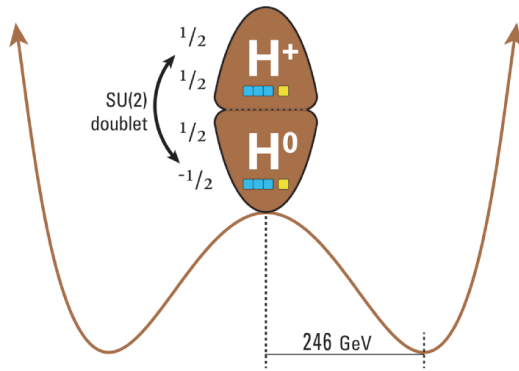
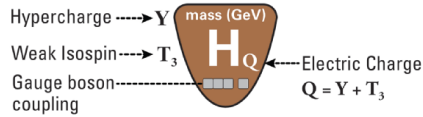


10:25 → 11:20 **Flavour Physics 1/3** ¶

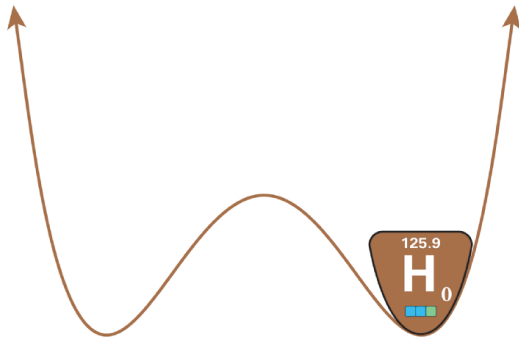
Speaker: Yasmine Sara Amhis (IJCLab/CERN)

# Higgs

Spin 0  
(Higgs Boson)

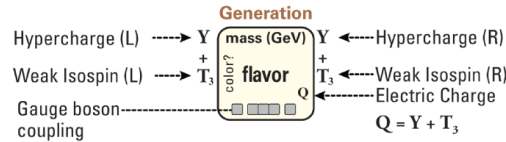


Unbroken Symmetry  
Broken Symmetry



# Matter

Spin 1/2  
(Fermions)

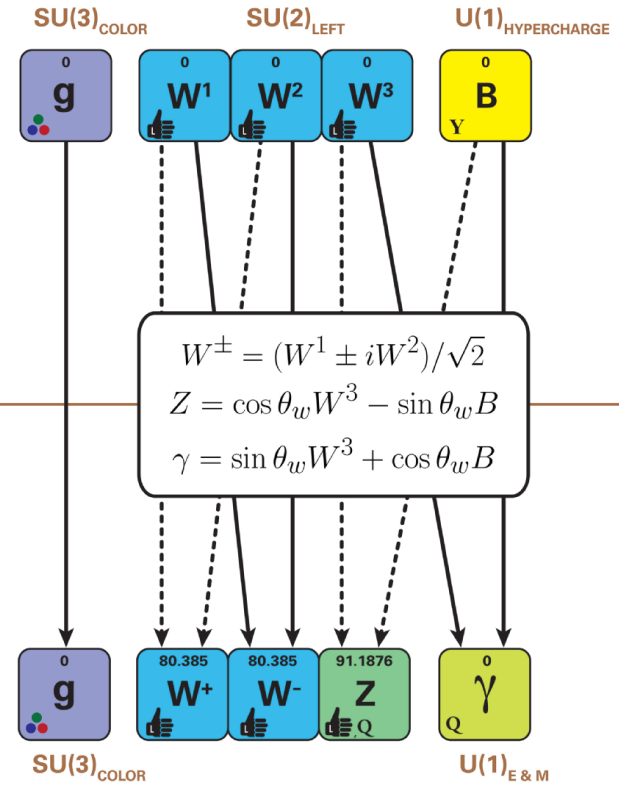
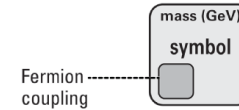


	1st	2nd	3rd	
Left handed SU(2) doublet	$1/6$ $u$ $2/3$	$1/6$ $c$ $2/3$	$1/6$ $t$ $2/3$	Quarks
	$1/6$ $d$ $-1/3$	$1/6$ $s$ $-1/3$	$1/6$ $b$ $-1/3$	
Left handed SU(2) doublet	$-1/2$ $\nu_e$ $0$	$-1/2$ $\nu_\mu$ $0$	$-1/2$ $\nu_\tau$ $0$	Leptons
	$-1/2$ $e$ $-1$	$-1/2$ $\mu$ $-1$	$-1/2$ $\tau$ $-1$	

	1st	2nd	3rd
$0.0023$ $u$ $2/3$	$1.275$ $c$ $2/3$	$173.07$ $t$ $2/3$	
$0.0048$ $d$ $-1/3$	$0.095$ $s$ $-1/3$	$4.18$ $b$ $-1/3$	
$m_1$ $\nu_e$ $0$	$m_2$ $\nu_\mu$ $0$	$m_3$ $\nu_\tau$ $0$	
$0.000511$ $e$ $-1$	$0.105658$ $\mu$ $-1$	$1.77682$ $\tau$ $-1$	

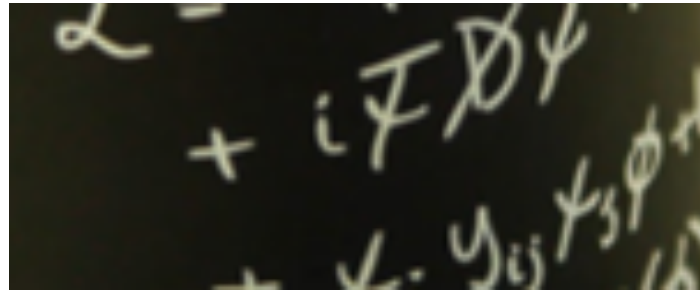
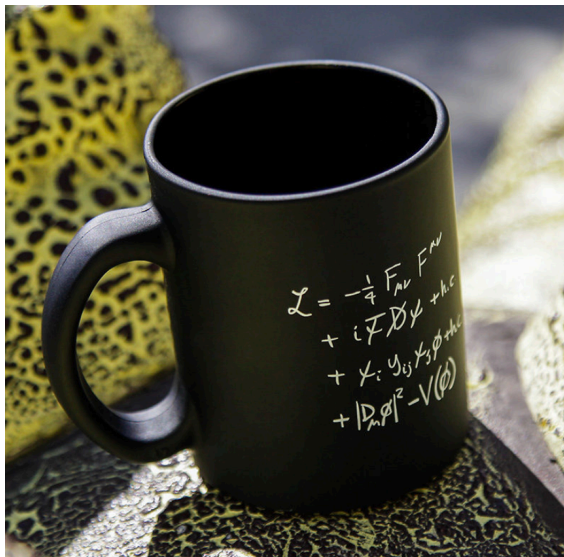
# Forces

Spin 1  
(Gauge Bosons)



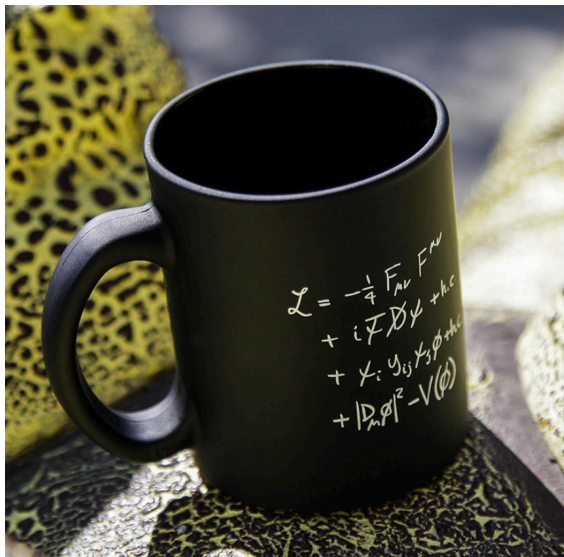
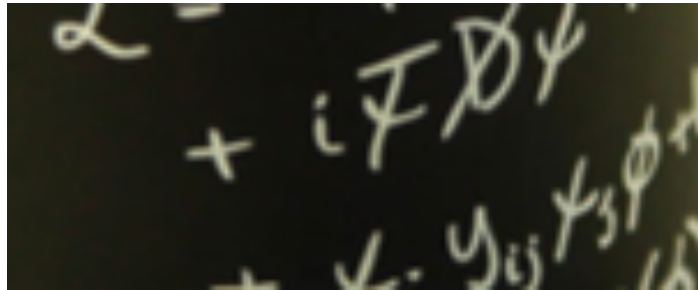
# Matter quantum fields: full form

Meaning of the mug?



# Matter quantum fields: full form

Meaning of the mug?



Left-handed quark field

weak SU(2) (1,2)      color SU(3) (1,2,3)

$$\psi(x) = \psi_{\alpha f}^{ib}(x)$$

Dirac-index (1,2)

generations (1,2,3)

# Weak interactions paradoxes

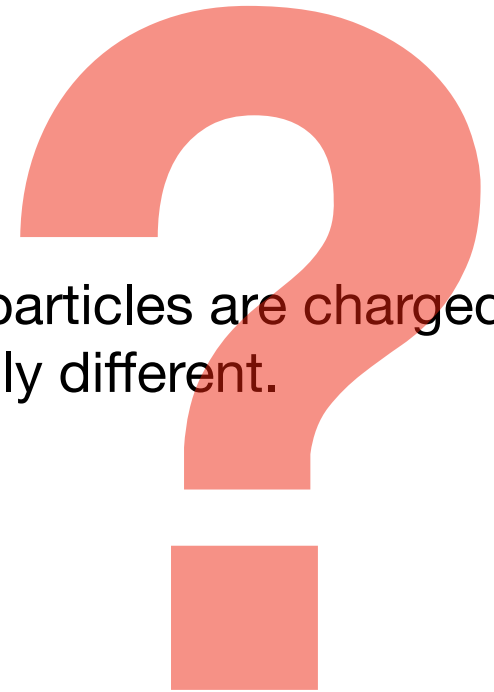
We need non-abelian gauge theory, but gauge bosons need to be massive!  
How can we save SU(2) gauge invariance?

$$M_W^2 W_\mu^+ W^{\mu-} + \frac{1}{2} M_Z^2 Z_\mu Z^\mu$$

Since the weak interactions are chiral (only the **left-handed** particles are charged), the **left-handed** and **right-handed** particles are fundamentally different.  
How can they have a mass term?

SU(2)  $\begin{pmatrix} u \\ d \end{pmatrix}_L$   $u_R$   $d_R$

vs.  $m(\bar{u}_L u_R + \bar{u}_R u_L)$  **not SU(2) invariant !**



# Avengers of the SM assemble



Peter Higgs



Francois Englert



Robert Brout



4th of July 2012  
Genf (CH)

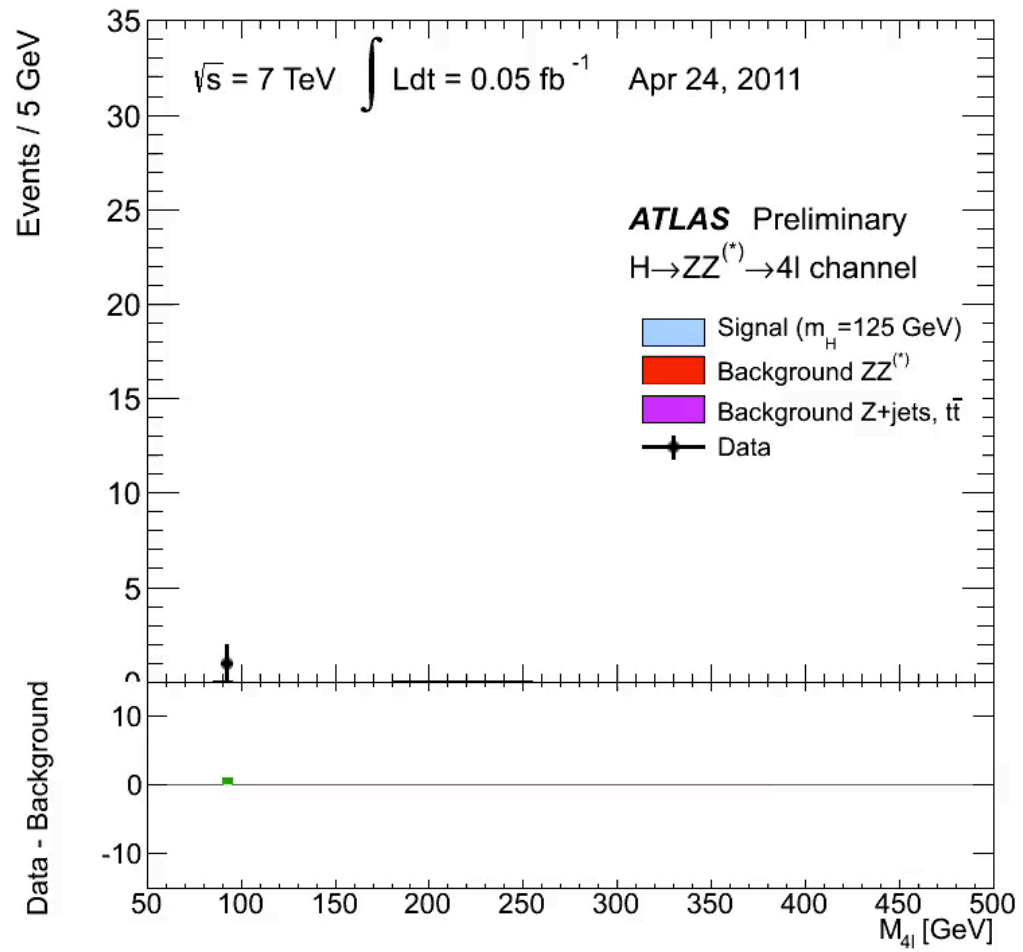


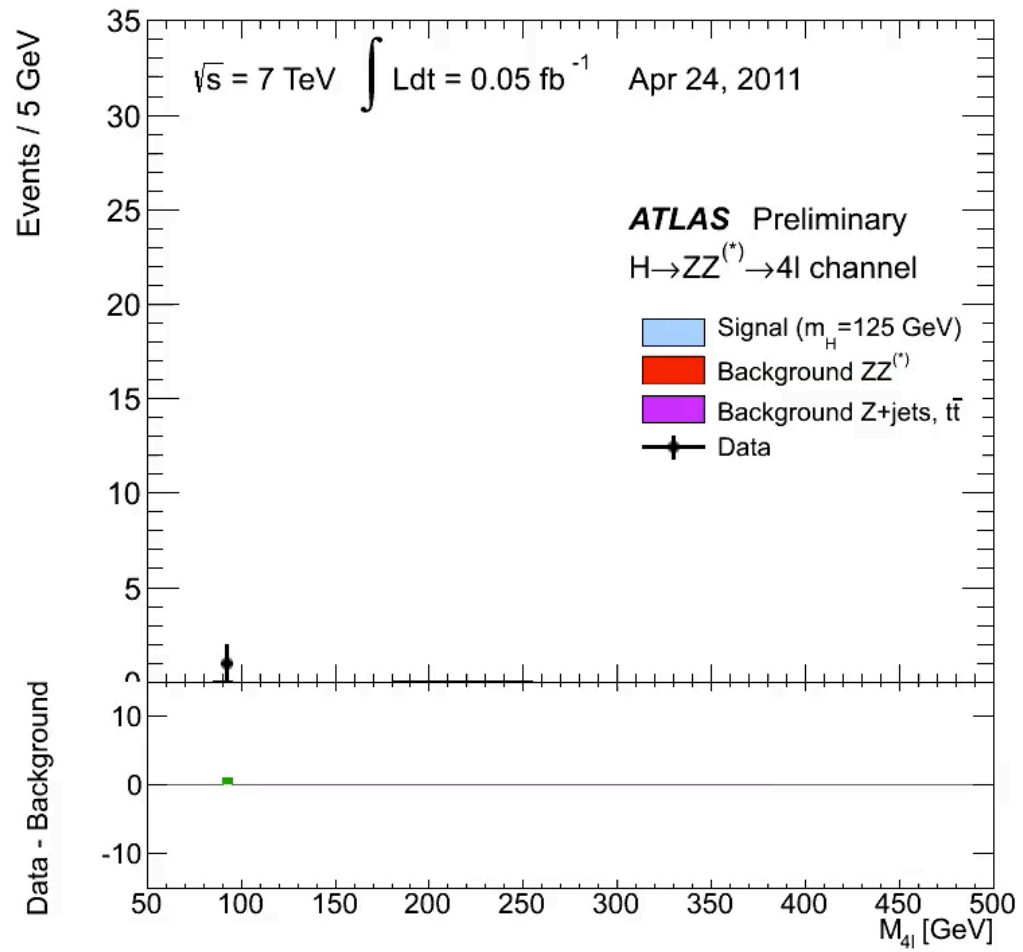




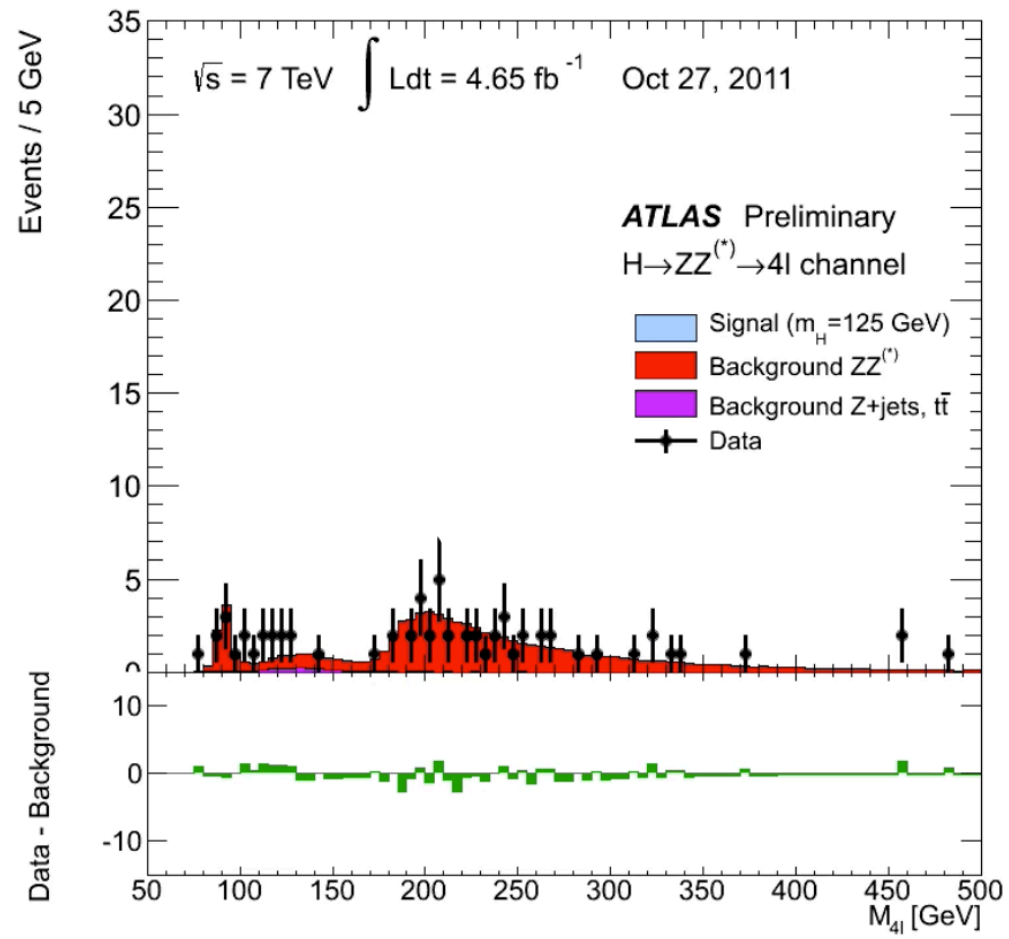


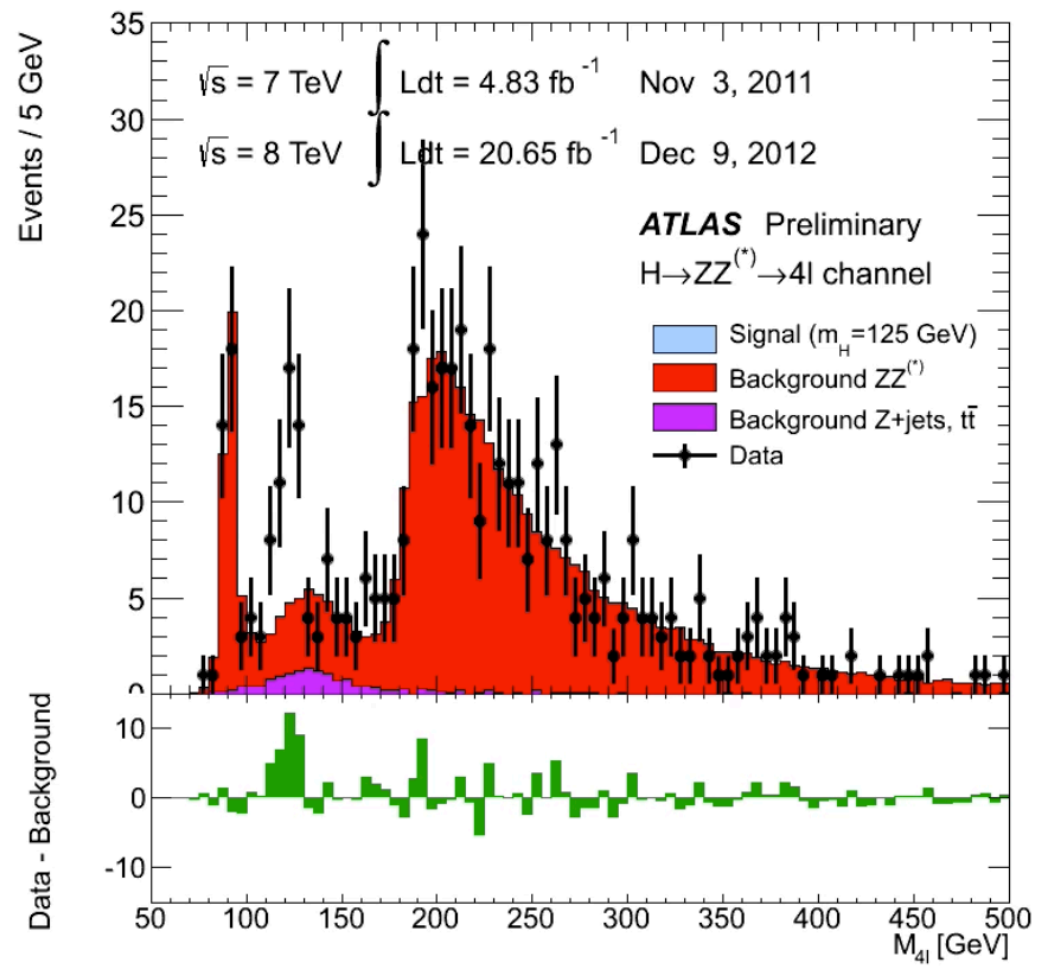


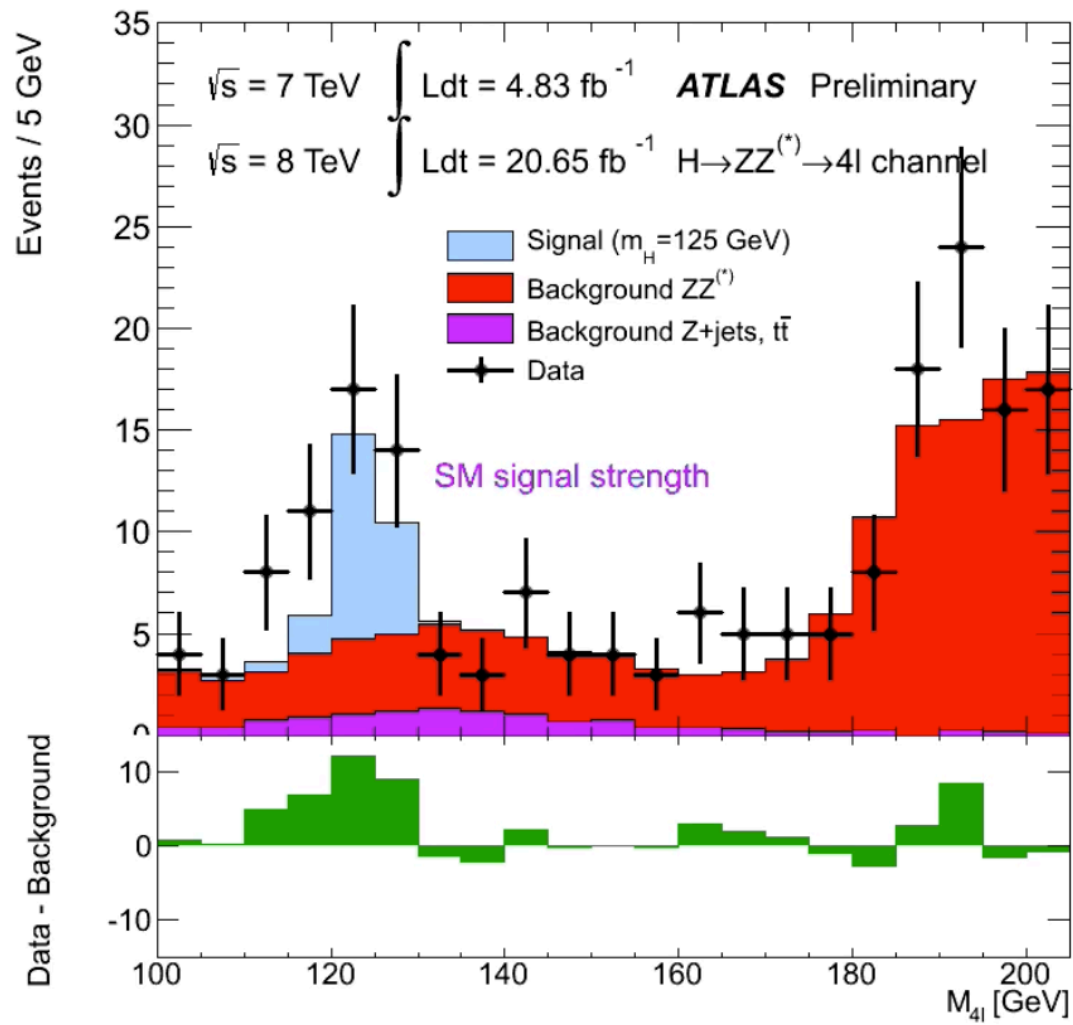


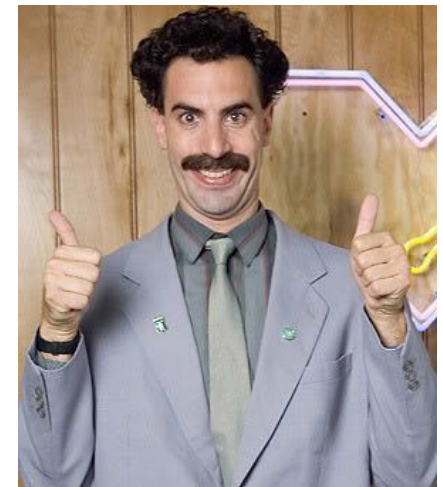
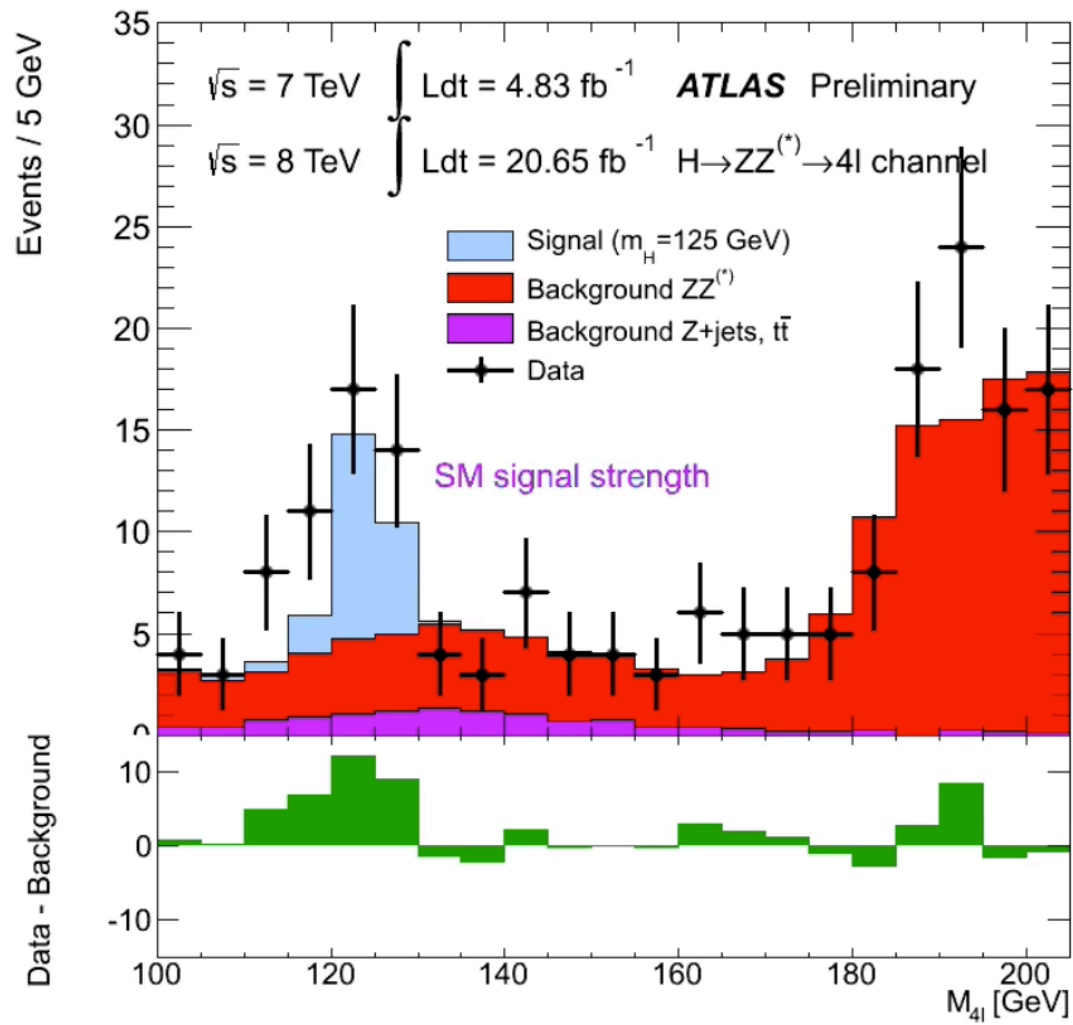














4/7/2012  
discovery of the  
Higgs

theory : 1964

design : 1984

construction : 1998

One year later in the theory  
group at CERN...









## The Nobel Prize in Physics 2013

François Englert, Peter Higgs



The Nobel Prize in Physics 2013 was awarded jointly to François Englert and Peter W. Higgs "for the theoretical discovery of a mechanism that contributes to our understanding of the origin of mass of subatomic particles, and which recently was confirmed through the discovery of the predicted fundamental particle, by the ATLAS and CMS experiments at CERN's Large Hadron Collider."

# Higgs is at the center of SM

- the only particle that **talks to everybody**\*
- the only elementary particle that **doesn't spin**
- the only particle that is **condensed** in the universe
- the **source of all masses** of elementary particles
- We **don't know why** all of this is the case
- important & special particle!

\* if we ignore the graviton

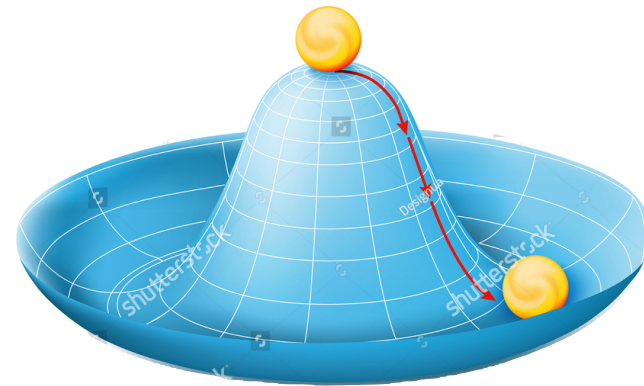
# Plan

## Model building the SM – requirements

- Explain the 3 massive weak bosons ( $W^+$ ,  $W^-$ ,  $Z$ ), **SU(2)** structure
- Two types of massless gauge bosons (**photon**, 8 gluons)
- Explain why the left-handed and right-handed particles can have different quantum numbers (= are different particles) while being massive
- Make some predictions we can test.



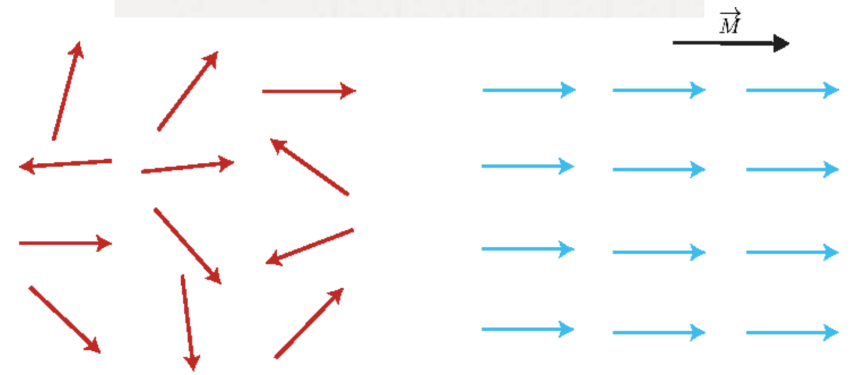
# Higgs mechanism!



# Spontaneous symmetry breaking



# Spontaneous symmetry breaking

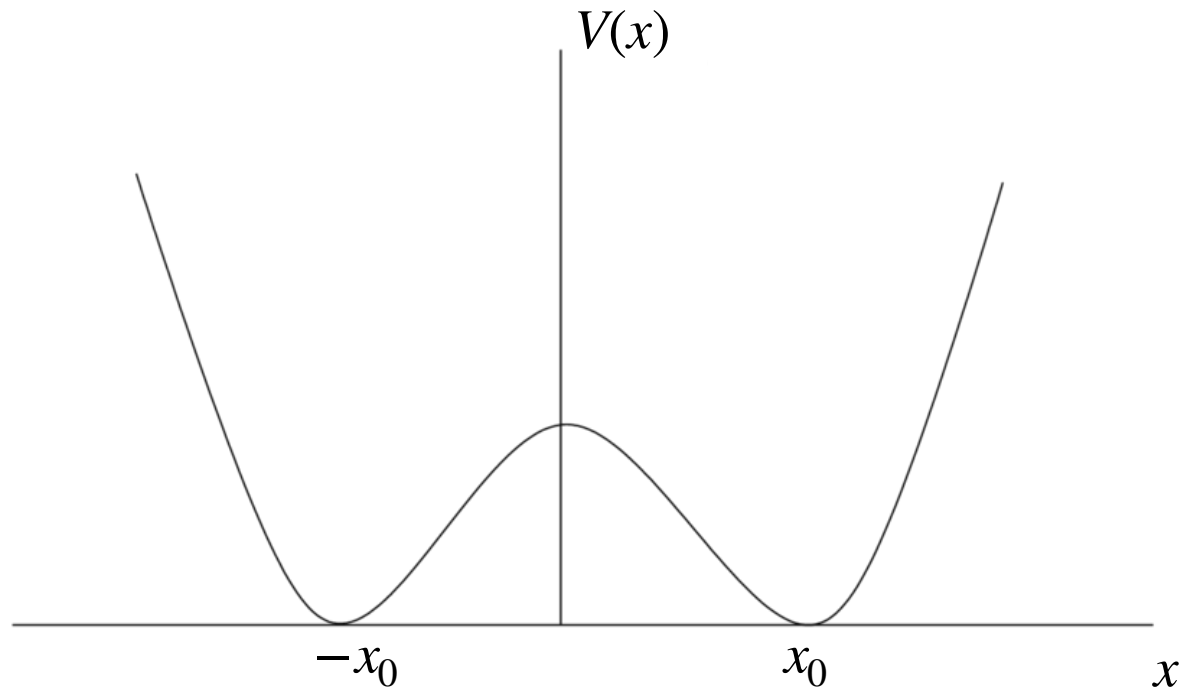


(a) The ferromagnet for  $T > T_c$ .  
The global rotational symmetry is unbroken.

(b) The ferromagnet for  $T < T_c$ .  
The global rotational symmetry is spontaneously broken.

# Classical mechanics

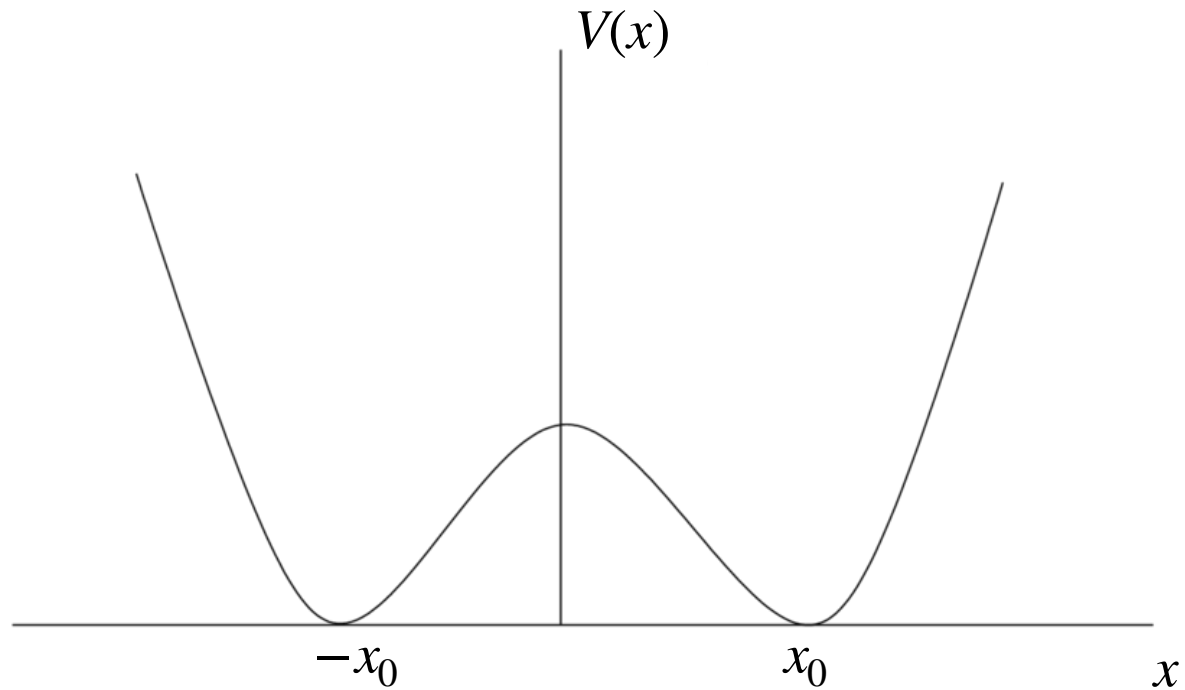
Symmetries? what is the ground state?



# Classical mechanics

Symmetries? what is the ground state?

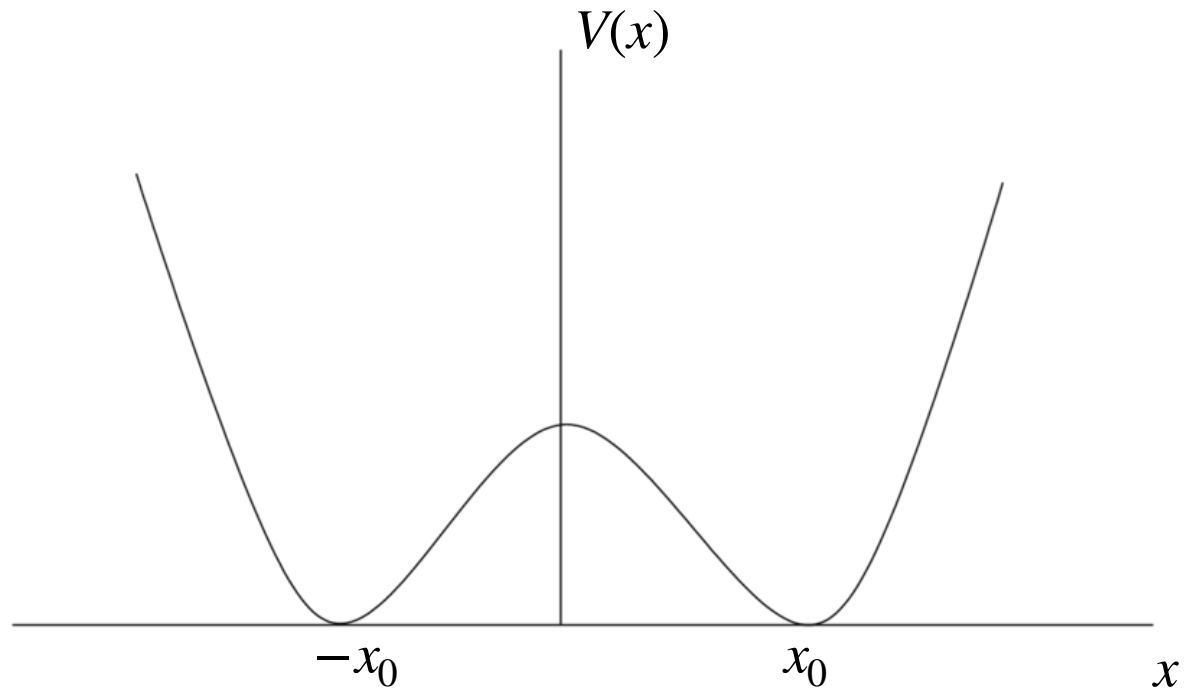
$$x \rightarrow -x$$





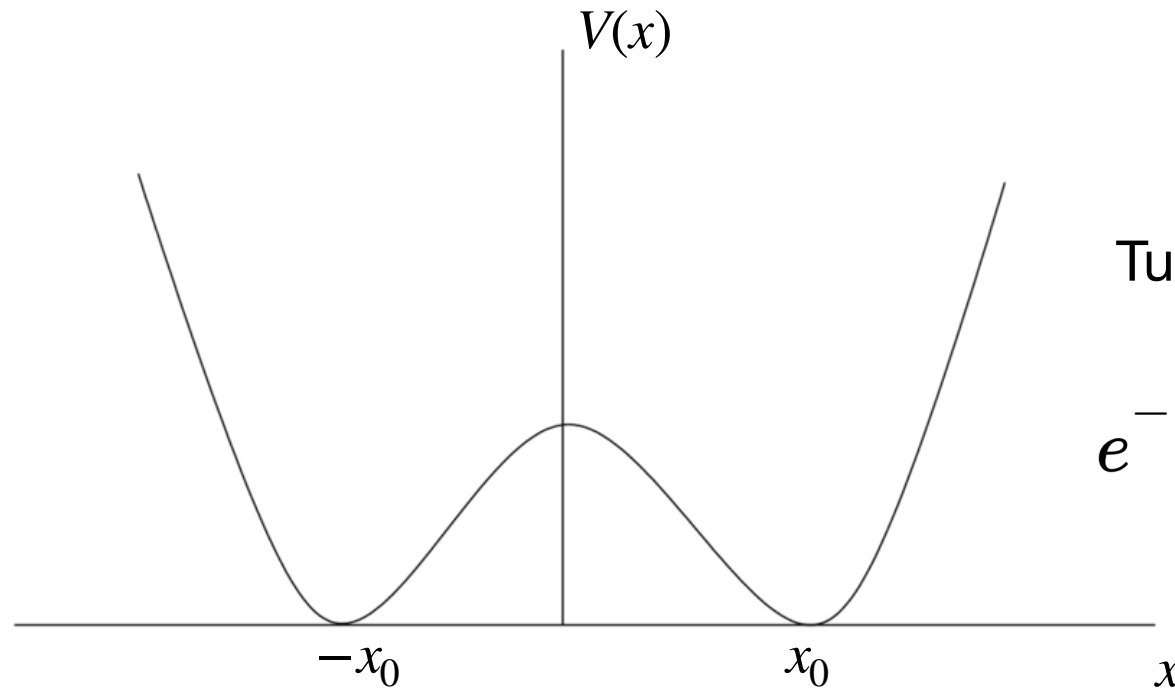
# Quantum mechanics

Symmetries? what is the ground state?



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Symmetries? what is the ground state?

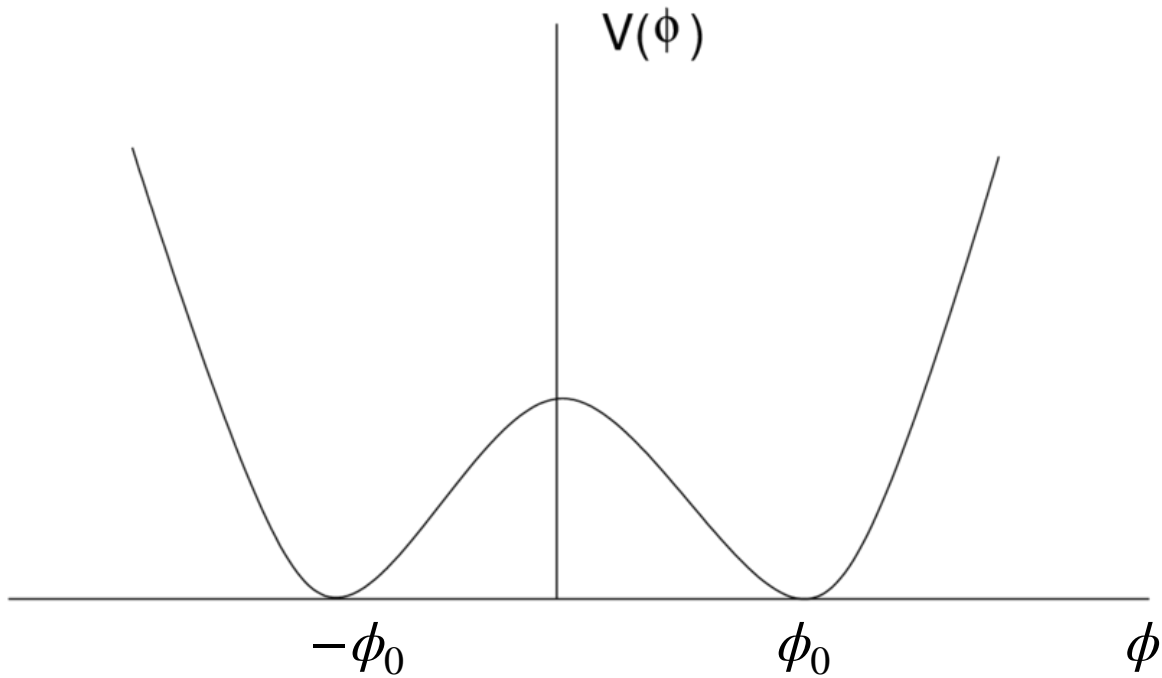


Tunneling amplitude

$$e^{-\int dx \sqrt{\frac{2m}{\hbar^2} (V(x) - E)}}$$

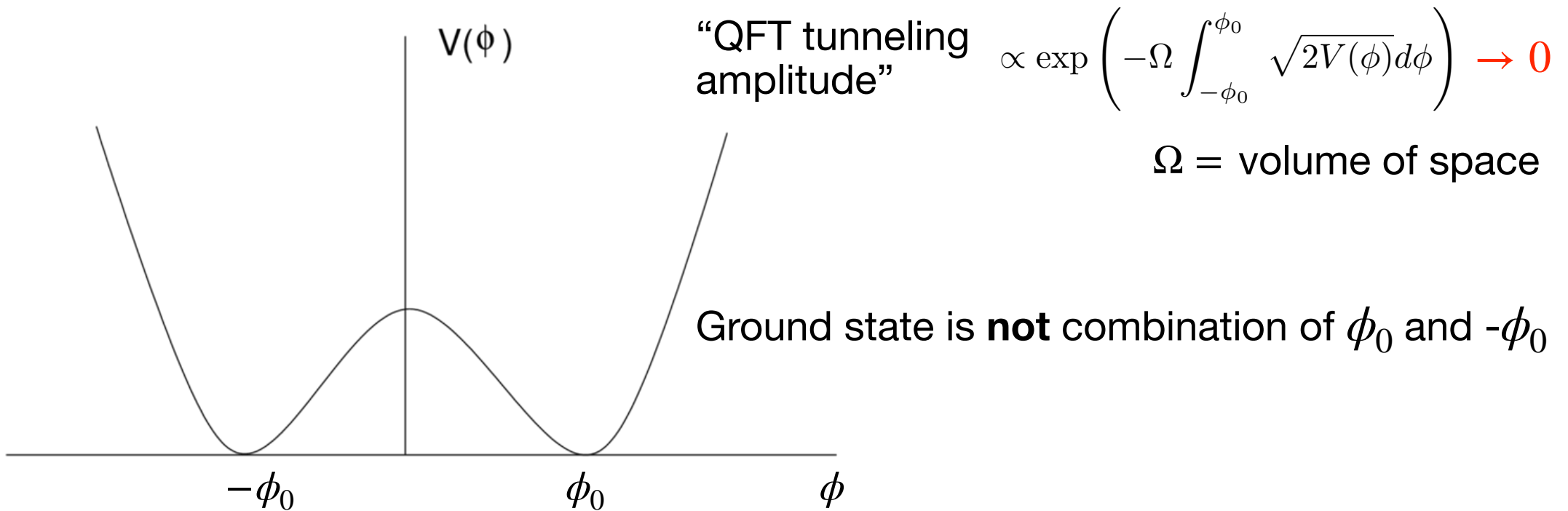
# Quantum Field Theory

Symmetries? what is the ground state?



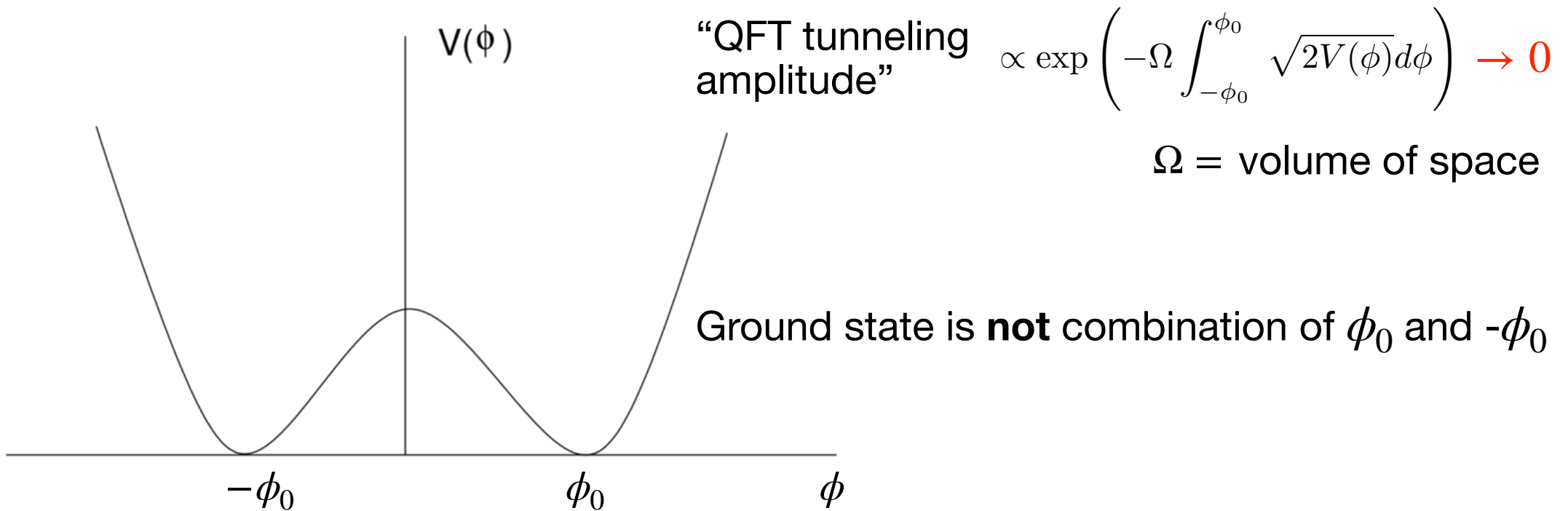
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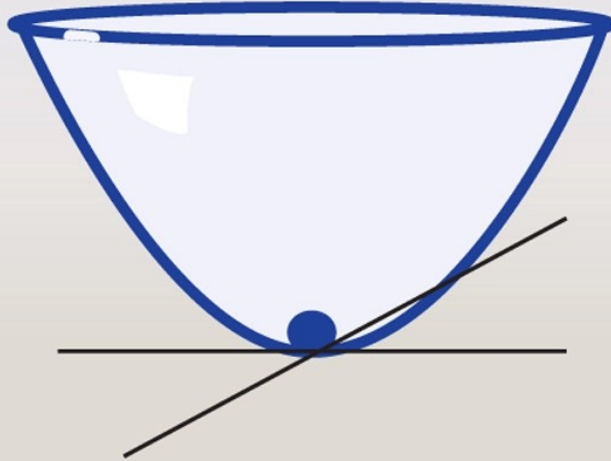
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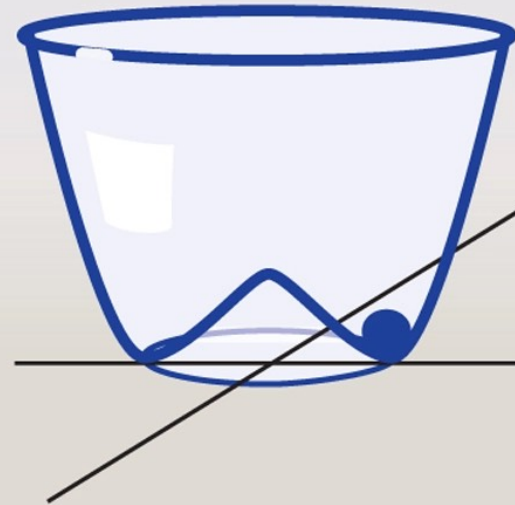


Ground state solution has **fewer symmetries** than the potential.

Unbroken symmetry



Spontaneously broken symmetry

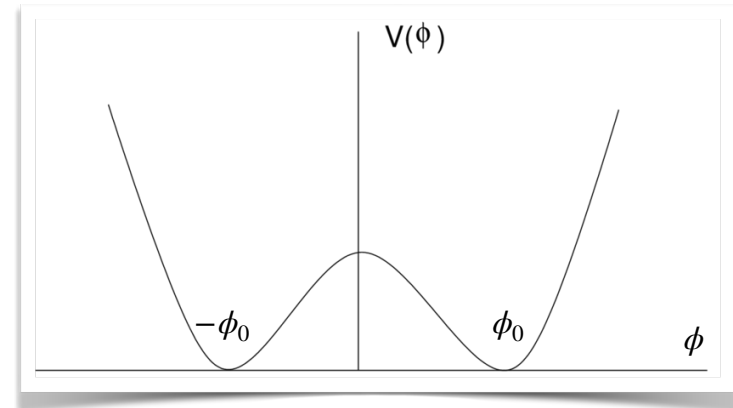


# Spontaneous symmetry breaking

Real scalar field (see Tim's lectures)

$$\mathcal{L} = \frac{1}{2} \partial^\mu \phi \partial_\mu \phi - V(\phi)$$

$$V(\phi) = -\frac{\mu^2}{2} \phi^2 + \frac{\lambda}{4} \phi^4$$

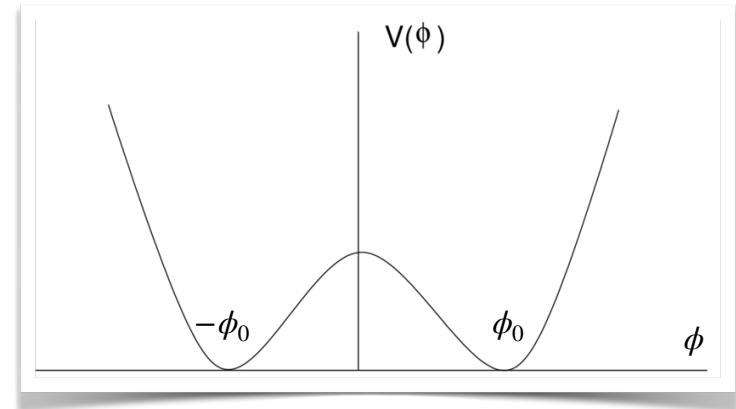


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Minimum at:  $\frac{\partial V}{\partial \phi} = 0 \Leftrightarrow -\mu^2 \phi + \lambda \phi^3 = 0$

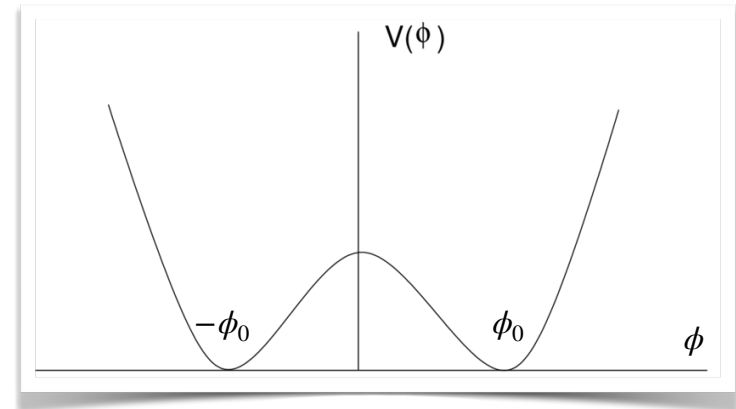


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➔  $\phi_0 = \pm \frac{\mu}{\sqrt{\lambda}} = \pm v.$

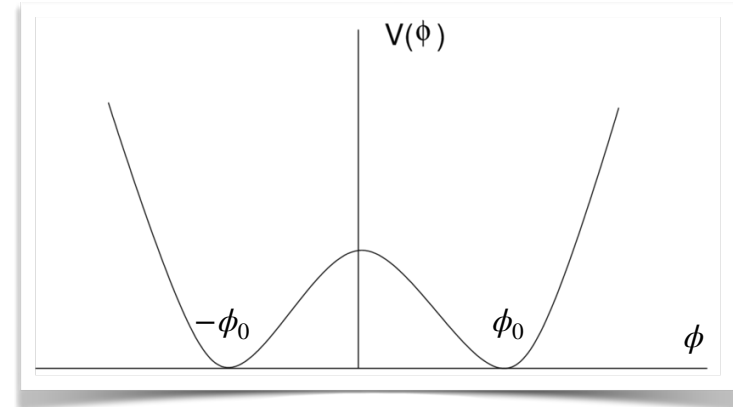
Scalar field develops “**vacuum expectation value**” ( $v = \text{vev}$ ).

# Spontaneous symmetry breaking

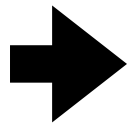
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Scalar field develops “**vacuum expectation value**” ( $v = \text{vev}$ ).

Pick one!

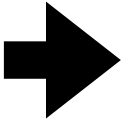
# Anderson-Higgs mechanism

Couple complex scalar field to a photon

gauge symmetry:

$$\phi(x) \rightarrow e^{i\alpha(x)} \phi(x)$$

$$A_\mu(x) \rightarrow A_\mu + \frac{1}{g} \partial_\mu \alpha(x).$$



# Anderson-Higgs mechanism

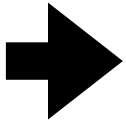
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$$\partial_\mu \phi \quad \text{becomes} \quad D_\mu \phi = (\partial_\mu - igA_\mu) \phi$$



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$$\rightarrow \mathcal{L} = D_\mu \phi^\dagger D_\mu \phi + \mu^2 \phi^\dagger \phi - \frac{\lambda}{2} (\phi^\dagger \phi)^2 - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}$$

# Anderson-Higgs mechanism

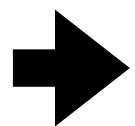
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$$\phi(x) \rightarrow e^{i\alpha(x)} \phi(x)$$

$$A_\mu(x) \rightarrow A_\mu + \frac{1}{g} \partial_\mu \alpha(x).$$

$$\partial_\mu \phi \text{ becomes } D_\mu \phi = (\partial_\mu - igA_\mu)\phi$$



$$\mathcal{L} = D_\mu \phi^\dagger D_\mu \phi + \mu^2 \phi^\dagger \phi - \frac{\lambda}{2} (\phi^\dagger \phi)^2 - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}$$

Groundstate  $A_\mu^{(0)} = 0$   $\phi^{(0)} = v = \frac{\sqrt{\mu}}{\lambda}$  (vev)

(Would it have been problematic if the photon developed a vacuum value?)

# Higgs-Anderson Mechanism

Particle spectrum

real scalar fields

vev



Expand around minimum:

$$\phi(x) = v + \frac{1}{\sqrt{2}} (\chi(x) + i\theta(x))$$

# Higgs-Anderson Mechanism

Particle spectrum

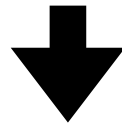
vev

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$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + D_{\mu} \phi^{\dagger} D_{\mu} \phi + \mu^2 \phi^{\dagger} \phi - \frac{\lambda}{2} (\phi^{\dagger} \phi)^2$$



$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu}^2 + e^2 v^2 A_{\mu} A^{\mu} + \frac{1}{2} (\partial_{\mu} \chi)^2 - \mu^2 \chi^2 + \dots$$



# Higgs-Anderson Mechanism

Particle spectrum

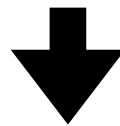
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Higgs particle

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu}^2 + e^2 v^2 A_{\mu} A^{\mu} + \frac{1}{2} (\partial_{\mu} \chi)^2 - \mu^2 \chi^2 + \dots$$

$m_A^2$

the gauge boson became **massive**.

# Higgs-Anderson Mechanism

Particle spectrum

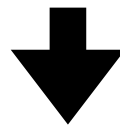
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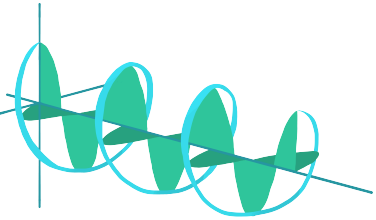
the gauge boson became **massive**.

What happened to the  $\theta(x)$  field? It got **eaten by the gauge boson!**

# Count physical degrees of freedom

$$\phi(x) = v + \frac{1}{\sqrt{2}} (\chi(x) + i\theta(x))$$

**massless photon**



**2 polarizations**

(2 transverse)

**complex scalar:**

**2 real degrees of freedom**

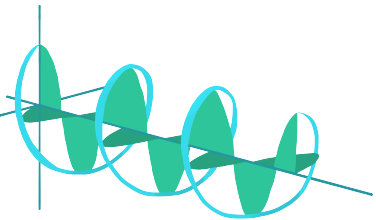
$$\Sigma \quad 2+2 = 4$$

# Count physical degrees of freedom

$$\phi(x) = v + \frac{1}{\sqrt{2}} (\chi(x) + i\theta(x))$$

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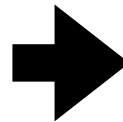
massive spin  $s=1$



**2 polarizations**

(2 transverse)

$$\langle \phi(x) \rangle = v$$



**(2s+1) = 3 polarizations**

(2 transverse + 1 longitudinal)

**complex scalar:**

**2** real degrees of freedom

$$\Sigma$$

$$2+2 = 4$$

*before SSB*

**real scalar**

**1** degree of freedom

$$\Sigma$$

$$3+1 = 4$$

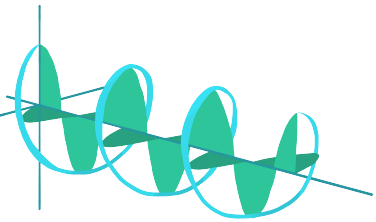
*after SSB*

# Count physical degrees of freedom

$$\phi(x) = v + \frac{1}{\sqrt{2}} (\chi(x) + i\theta(x))$$

massless photon

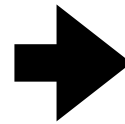
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**2 polarizations**

(2 transverse)

$$\langle \phi(x) \rangle = v$$



**(2s+1) = 3 polarizations**

(2 transverse + 1 longitudinal)

**complex scalar:**

2 real degrees of freedom

$\Sigma$

$$2+2 = 4$$

*before SSB*

**real scalar**

1 degree of freedom

$\Sigma$

$$3+1 = 4$$

*after SSB*



# Warm-up: Higgs for a SU(2) gauge boson

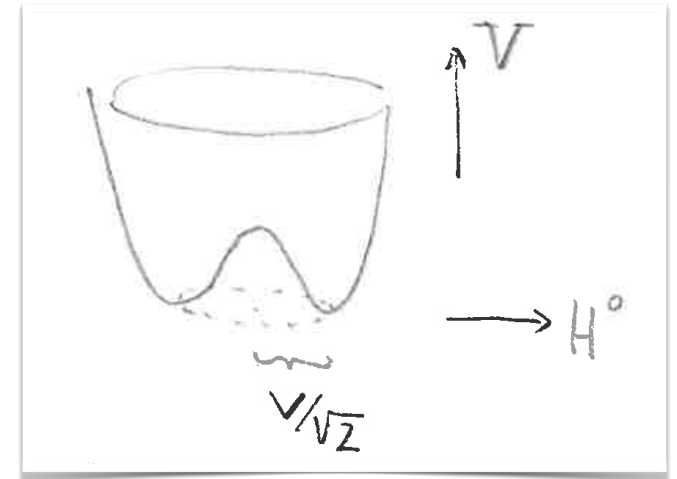
Higgs must couple to SU(2) bosons:

Complex 2 vector  $H(x) = \begin{pmatrix} H^+ \\ H^0 \end{pmatrix} \quad H(x) \rightarrow e^{i\sigma^a \alpha^a(x)} H(x)$

$$\mathcal{L}_H = D_\mu H^\dagger D^\mu H - \lambda \left( H^\dagger H - \frac{v^2}{2} \right)^2$$

$$H(x) = \begin{pmatrix} H^+ \\ H^0 \end{pmatrix} = \frac{1}{\sqrt{2}} e^{i\theta^a(x)\sigma^a} \begin{pmatrix} 0 \\ v + h(x) \end{pmatrix}$$

↑  
get eaten

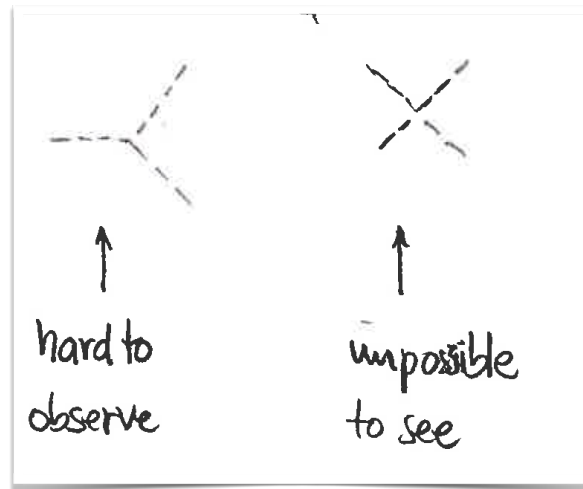


# Higgs potential

probe at LHC and future machines!

$$V \rightarrow + \lambda v^2 h^2 + \lambda v h^3 + \frac{\lambda}{4} h^4$$

$$\uparrow$$
$$m_h^2 = 2\lambda v^2$$



# Higgs potential

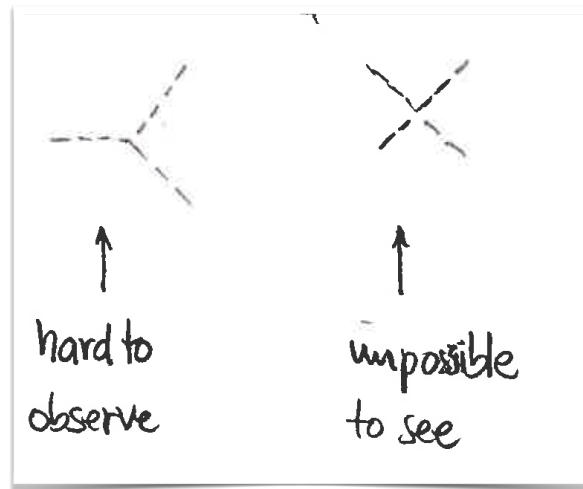
probe at LHC and future machines!

✓

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↑

$$m_h^2 = 2\lambda v^2$$





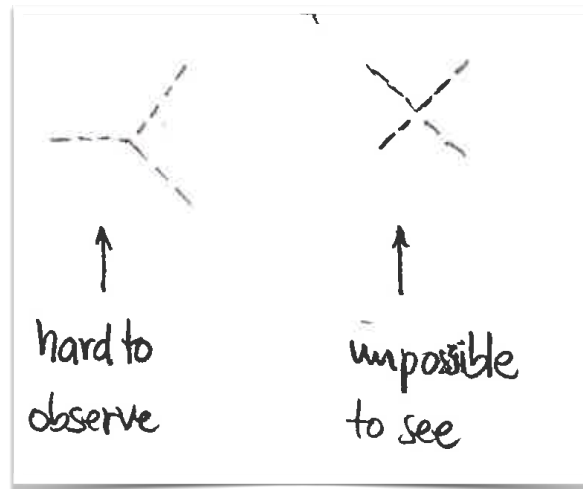
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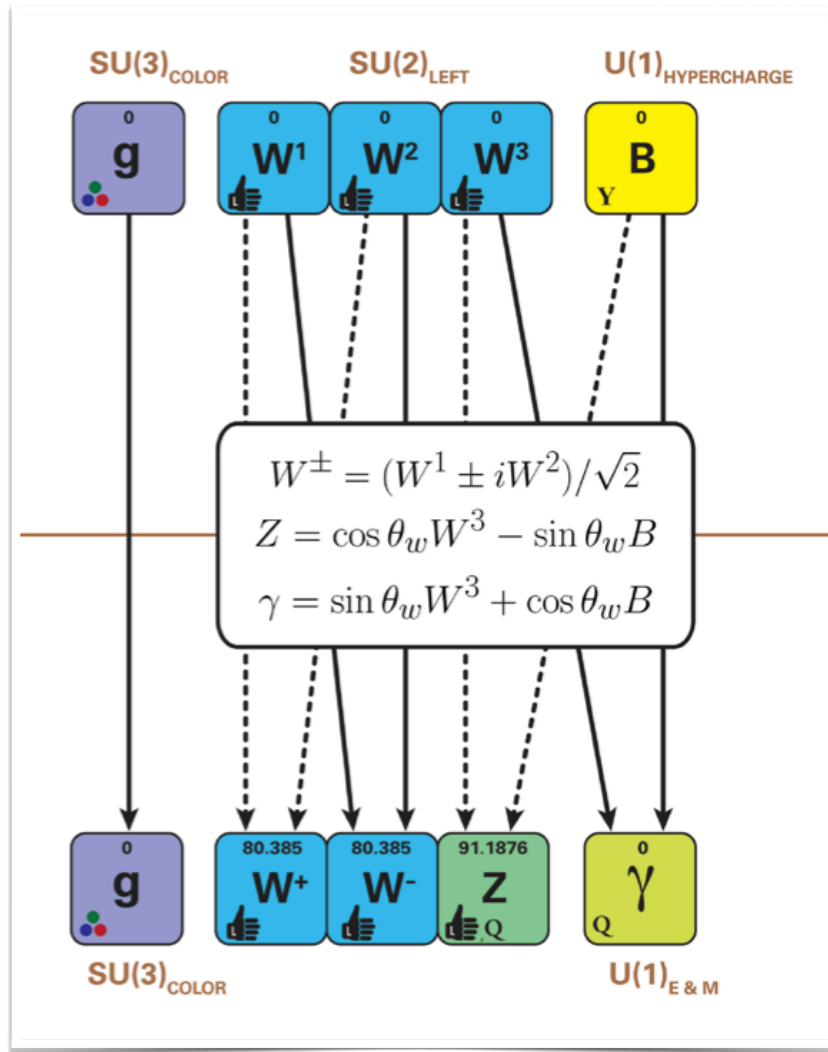
$$V \rightarrow + \lambda v^2 h^2 + \lambda v h^3 + \frac{\lambda}{4} h^4$$

↑

$$m_h^2 = 2\lambda v^2$$



# Higgs mechanism for W,Z masses



To arrive at  $W^+, W^-, Z$ , and  $\gamma$   
 propose:  $SU(2)_L \times U(1)_Y$



Glashow, Salam, Weinberg (Nobel '73)

# Higgs mechanism for W,Z masses

Gauge structure:  $H(x) \rightarrow e^{i\beta(x)} e^{i\sigma^a \alpha^a(x)} H(x)$   $\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

$$D_\mu H = \left( \partial_\mu - ig' \frac{1}{2} B_\mu - ig \frac{\sigma^a}{2} W_\mu^a \right) H$$

$U(1)_Y$        $SU(2)_L$

# Higgs mechanism for W,Z masses

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$$D_\mu H = \left( \partial_\mu - ig' \frac{1}{2} B_\mu - ig \frac{\sigma^a}{2} W_\mu^a \right) H$$

in components:

$$D_\mu H = \partial_\mu H - \frac{i}{2} \begin{pmatrix} gW_\mu^3 + g'B_\mu & \sqrt{2}gW_\mu^+ \\ \sqrt{2}gW_\mu^- & -gW_\mu^3 + g'B_\mu \end{pmatrix} H \quad \text{with} \quad W_\mu^\pm = \frac{1}{\sqrt{2}} (W_\mu^1 \mp W_\mu^2)$$

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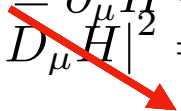
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$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix}$$



$$|D_\mu H|^2 = \frac{1}{4} g^2 v^2 W_\mu^+ W^{-\mu} + \frac{1}{8} (W_\mu^3 B_\mu) \begin{pmatrix} g^2 v^2 & -gg'v^2 \\ -gg'v^2 & g'^2 v^2 \end{pmatrix} \begin{pmatrix} W^{3\mu} \\ B^\mu \end{pmatrix}$$

# Higgs mechanism for W,Z masses

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$U(1)_Y$        $SU(2)_L$

$$D_\mu H = \left( \partial_\mu - ig' \frac{1}{2} B_\mu - ig \frac{\sigma^a}{2} W_\mu^a \right) H$$

Diagonalize

in components:

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# Higgs mechanism for W,Z masses

Diagonalize

$$A = \sin \theta_W W_3 + \cos \theta_W B$$

$$Z = \cos \theta_W W_3 - \sin \theta_W B$$

$$D_\mu H^\dagger D^\mu H \rightarrow \frac{g^2 v^2}{4} W_\mu^+ W^{-\mu} + \frac{(g^2 + g'^2) v^2}{4} \frac{1}{2} Z_\mu Z^\mu + 0 \cdot A_\mu A^\mu$$

$M_W^2$ 
 $M_Z^2$ 
 $M_A^2$

↑
↑
↓

↗

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix}$$

$$\sin \theta_w = \frac{g'}{\sqrt{g^2 + g'^2}}$$

# Higgs mechanism for W,Z masses

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$M_A^2$



$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix}$$



$M_W^2$



$M_Z^2$



$$v = 246 \text{ GeV}$$

$$g = 0.65, \quad g' = 0.36$$

$$M_W = 80.38 \text{ GeV}$$

$$M_Z = 90.19 \text{ GeV}$$

$$\sin^2 \theta_w = 0.231$$

$$\sin \theta_w = \frac{g'}{\sqrt{g^2 + g'^2}}$$



# Electro-weak unification

$$e = \frac{g g'}{\sqrt{g^2 + g'^2}}$$

The electro-magnetic coupling is derived from a more fundamental theory!

# Higgs mechanism for W,Z masses

$$D_\mu H = \left( \partial_\mu - ig' \frac{1}{2} B_\mu - ig \frac{\sigma^a}{2} W_\mu^a \right) H$$

$$D_\mu H^\dagger D^\mu H \rightarrow \frac{g^2 v^2}{4} W_\mu^+ W^{-\mu} + \frac{(g^2 + g'^2) v^2}{4} \frac{1}{2} Z_\mu Z^\mu$$

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix}$$

$$M_W^2$$

$$M_Z^2$$

$$\sin \theta_W = \frac{g}{\sqrt{g^2 + g'^2}}$$

$$W^\pm = \frac{W_1 \mp iW_2}{\sqrt{2}}$$

$$A = \sin \theta_W W_3 + \cos \theta_W B$$

$$Z = \cos \theta_W W_3 - \sin \theta_W B$$

# Fermion masses

The SM is chiral. Left- and right-handed electrons are different particles. But m

$$m_e(\bar{e}_L e_R + \bar{e}_R e_L)$$

Not gauge invariant!

$$\begin{aligned} \begin{pmatrix} \nu_e \\ e \end{pmatrix}_L &\rightarrow e^{i\sigma^a \alpha^a(x)} \begin{pmatrix} \nu_e \\ e \end{pmatrix}_L \\ e_R &\rightarrow e^{-i\beta(x)} e_R \end{aligned}$$

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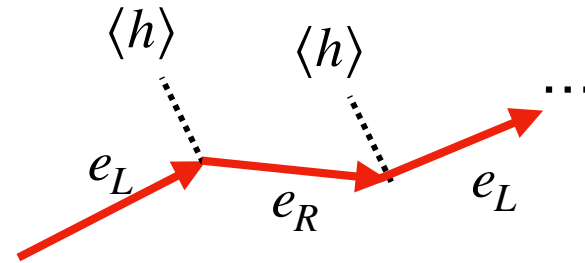
$$e_R \rightarrow e^{-i\beta(x)} e_R$$



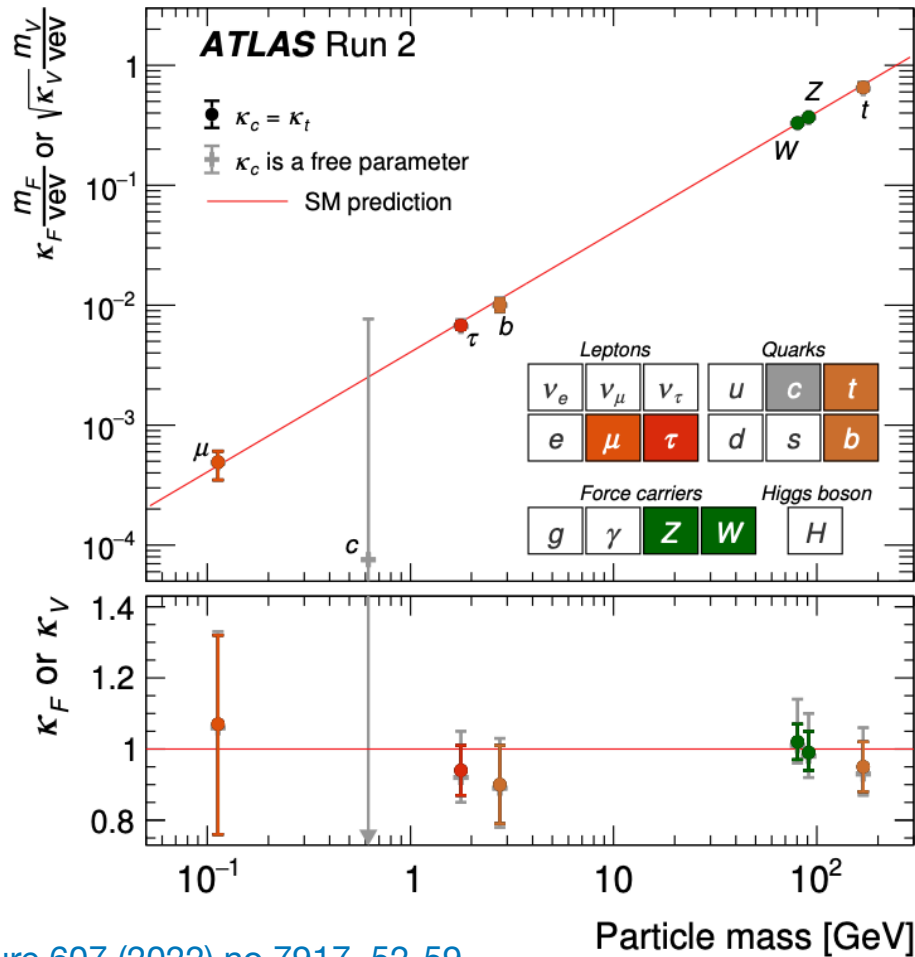
Higgs to the rescue! Fermion masses are emergent and are a result of interacting with the Higgs vacuum expectation value.

$$y_e \begin{pmatrix} \bar{\nu}_L \\ \bar{e}_L \end{pmatrix} \cdot \begin{pmatrix} H^+ \\ H^0 \end{pmatrix} e_R = \frac{y_e v}{\sqrt{2}} \left( \bar{e}_L e_R + \frac{1}{v} \bar{e}_L e_R h \right)$$

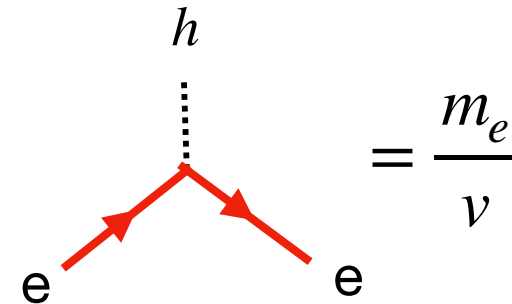
$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + h \end{pmatrix}$$



# Is the Higgs coupling proportional to mass of particle?



$$\frac{y_e v}{\sqrt{2}} \left( \bar{e}_L e_R + \frac{1}{v} \bar{e}_L e_R h \right) = m_e \left( \bar{e}_L e_R + \frac{1}{v} \bar{e}_L e_R h \right)$$



?

Nature 607 (2022) no.7917, 52-59

# Flavor in the SM

Higgs matter interactions are matrices, introduce generation hopping interactions (flavor change)

$$\mathcal{L}_{yukawa} = Y_L^{ij} \begin{pmatrix} \bar{\nu}_L \\ \bar{l}_L \end{pmatrix}^i H l_R^j + Y_U^{ij} \begin{pmatrix} \bar{u}_L \\ \bar{d}_L \end{pmatrix}^i \tilde{H} u_R^j + Y_D^{ij} \begin{pmatrix} \bar{u}_L \\ \bar{d}_L \end{pmatrix}^i H d_R^j$$

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$$Y_U \approx \begin{pmatrix} 6 \cdot 10^{-6} & -0.001 & 0.008 + 0.004i \\ 1 \cdot 10^{-6} & 0.004 & -0.04 + 0.001i \\ 8 \cdot 10^{-9} + 2 \cdot 10^{-8}i & 0.0002 & 0.98 \end{pmatrix}$$

$$Y_D \approx \text{diag} (2 \cdot 10^{-5} \quad 0.0005 \quad 0.02)$$



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What is the reason for this ?  
compare to:

$$Y_D \approx \text{diag} (2 \cdot 10^{-5} \quad 0.0005 \quad 0.02)$$

$$g_s \sim 1, \quad g \sim 0.6, \quad g' \sim 0.3, \quad \lambda_{\text{Higgs}} \sim 1$$

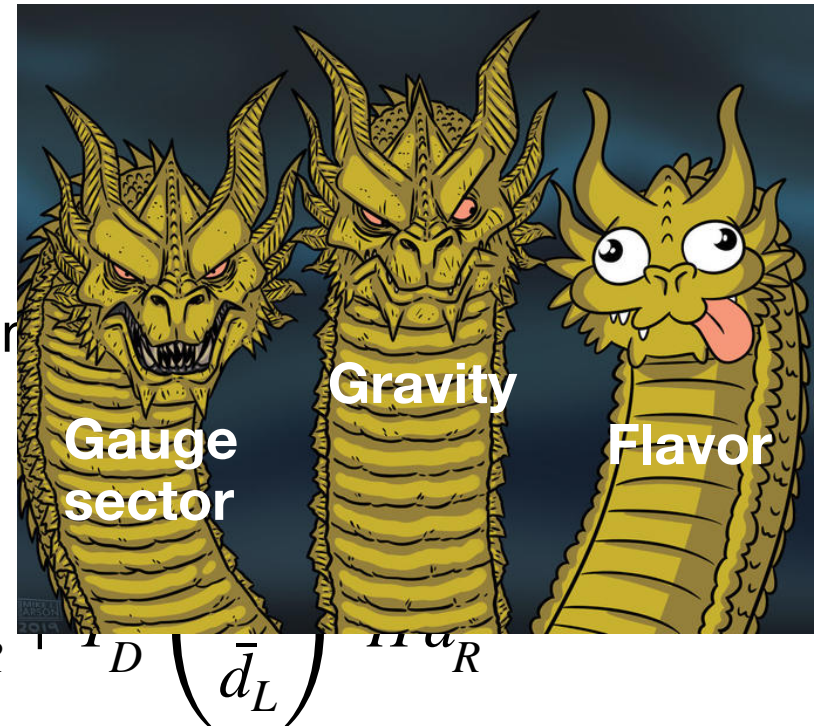
# Flavor in the SM

Higgs matter interactions are matrices, in  
interactions (flavor change)

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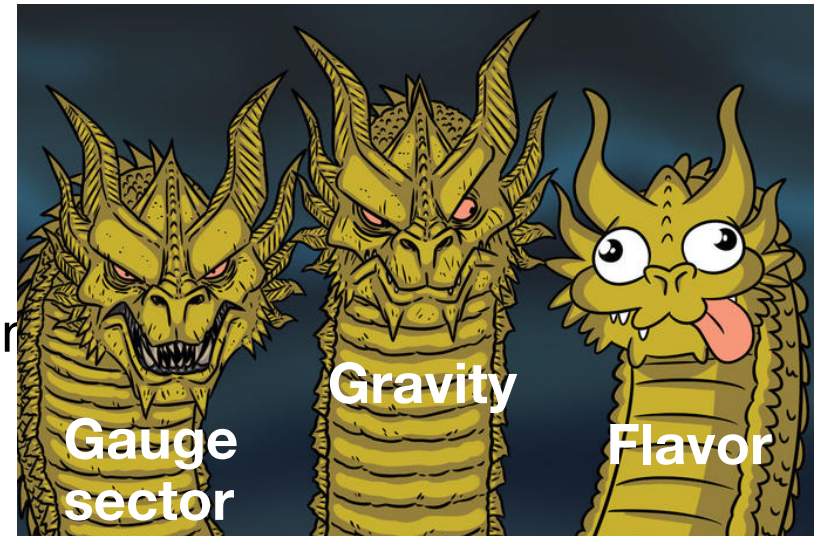


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# Flavor in the SM

Higgs matter interactions are matrices, in interactions (flavor change)



10:25

→ 11:20

Flavour Physics 1/3 ¶

Speaker: Yasmine Sara Amhis (IJCLab/CERN)

$$Y_U \approx \begin{pmatrix} 6 \cdot 10^{-6} & -0.001 & 0.008 + 0.004i \\ 1 \cdot 10^{-6} & 0.004 & -0.04 + 0.001i \\ 8 \cdot 10^{-9} + 2 \cdot 10^{-8}i & 0.0002 & 0.98 \end{pmatrix}$$

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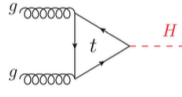
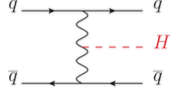

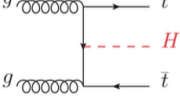
$$g_s \sim 1, \quad g \sim 0.6, \quad g' \sim 0.3, \quad \lambda_{\text{Higgs}} \sim 1$$

# LHC is pinning down the Higgs (> 8 Million produced)

## Nano Overview of Main Higgs Analyses at (HL) LHC

Most channels already covered at the Run 2 with only 5% (~150 fb<sup>-1</sup>) of full HL-LHC dataset!

72

Channel categories		Br	ggF  ~8 M vets produced	VBF  ~600 k vets produced	VH  ~400 k vets produced	ttH  ~80 k evts produced
Cross Section 13 TeV (8 TeV)			48.6 (21.4) pb*	3.8 (1.6) pb	2.3 (1.1) pb	0.5 (0.1) pb
Observed modes	$\gamma\gamma$	0.2 %	✓	✓	✓	✓
	ZZ	3%	✓	✓	✓	✓
	WW	22%	✓	✓	✓	✓
	$\tau\tau$	6.3 %	✓	✓	✓	✓
	bb	55%	✓	✓	✓	✓
Remaining to be observed	Z $\gamma$ and $\gamma\gamma^*$	0.2 %	✓	✓	✓	✓
	$\mu\mu$	0.02 %	✓	✓	✓	✓
Limits	Invisible	0.1 %	✓ (monojet)	✓	✓	✓

courtesy Marumi Kado

**much more on the Higgs searches here:**

**11:35**

→ 12:30

**Experimental Physics at Hadron Colliders 1/4**

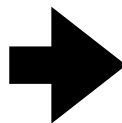
**Speaker:** Markus Klute (Karlsruhe Inst. of Technology (GER))

# Good news: SM is incomplete

At least five missing pieces in the SM

- non-baryonic dark matter
- neutrino mass
- dark energy
- inflation
- baryon asymmetry

We don't know their energy scales.



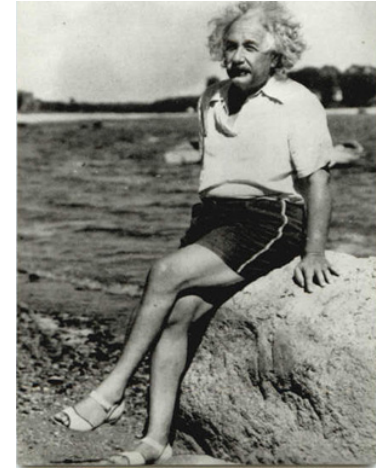
TUESDAY, 30 JULY

09:15 → 10:10 **Beyond the Standard Model Physics 1/4** ¶

Speaker: Tevong You (King's College London)



– *Henri Poincaré's recommendation letter for  
A. Einstein, 1911*



“I do not mean to say that all these anticipations will withstand the test of experiment on the day such a test would become possible. Since he seeks in all directions one must, on the contrary, expect most of the trails [...] to be blind alleys.



– *Henri Poincaré's recommendation letter for  
A. Einstein, 1911*



“I do not mean to say that all these anticipations will withstand the test of experiment on the day such a test would become possible. Since he seeks in all directions one must, on the contrary, expect most of the trails [...] to be blind alleys.

But one must hope at the same time that one of the directions he has indicated may be the right one, and that is enough. This is indeed how one should proceed. The role of mathematical physics is to ask the right questions, and experiment alone can resolve them.”





**Enjoy the rest of your summer  
at CERN!**

**Office hours: 2:00-3:00 (today)**

**Office: 4/2-026**

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