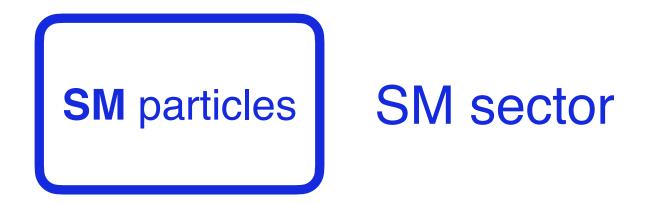
# Physics potential of beam dump experiments at future accelerators

Corfu Workshop on Future Accelerators 25 May 2024

**Daiki Ueda (Technion)** 

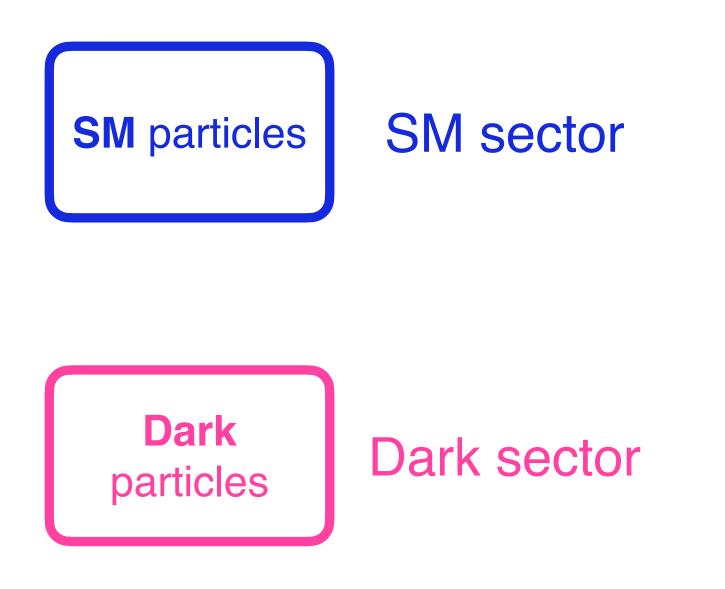


· Unsolved problems, e.g., neutrino mass, hierarchy problem, and dark matter(DM), remain in the SM



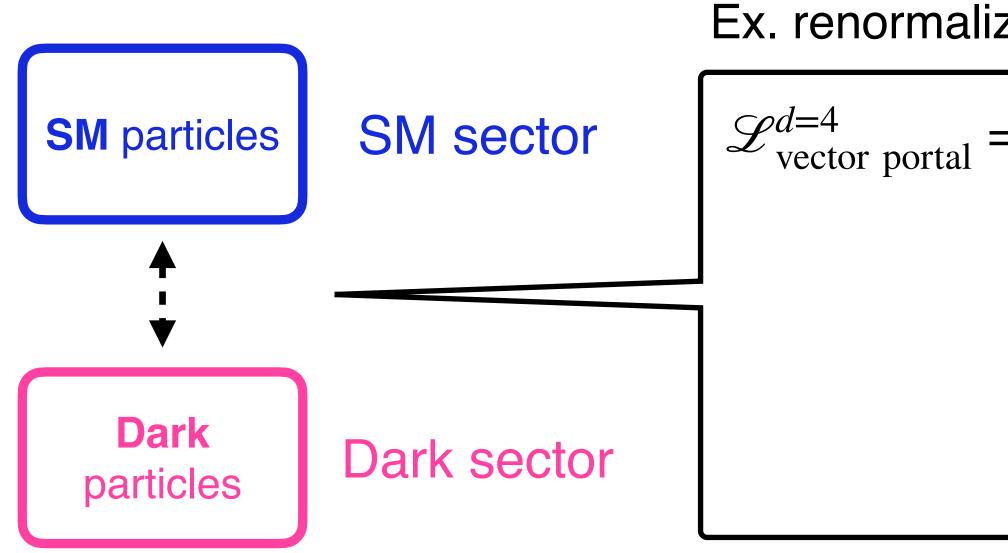
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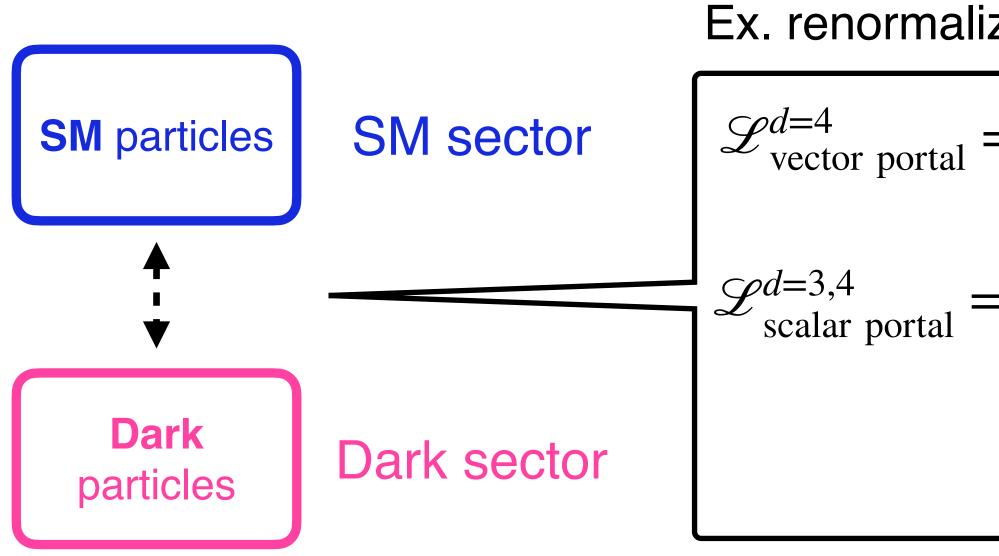


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Ex. renormalizable interaction b/w SM and dark sector

 $-\frac{1}{2\cos\theta_W}B_{\mu\nu}F'_{\mu\nu}$  $B_{\mu
u}$  hypercharge field strength,  $F_{\mu
u}' U(1)'$  filed strength



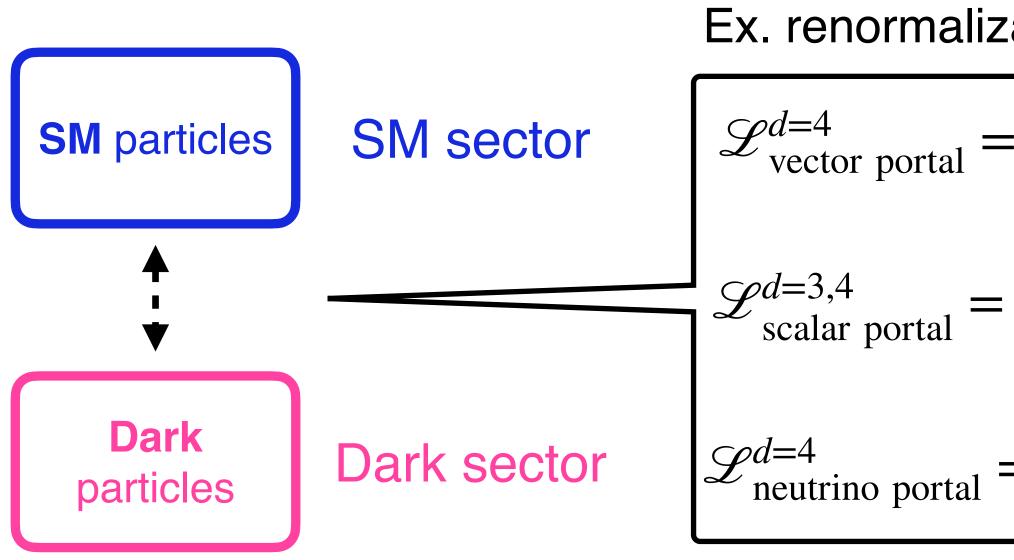


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 $\frac{1}{2\cos\theta_W} \frac{B_{\mu\nu}F_{\mu\nu}}{F_{\mu\nu}}$  $B_{\mu\nu}$  hypercharge field strength,  $F'_{\mu\nu}$  U(1)' filed strength  $= -H^{\dagger}H(AS + \lambda S^2)$ H SM Higgs field, S SM singlet scalar field



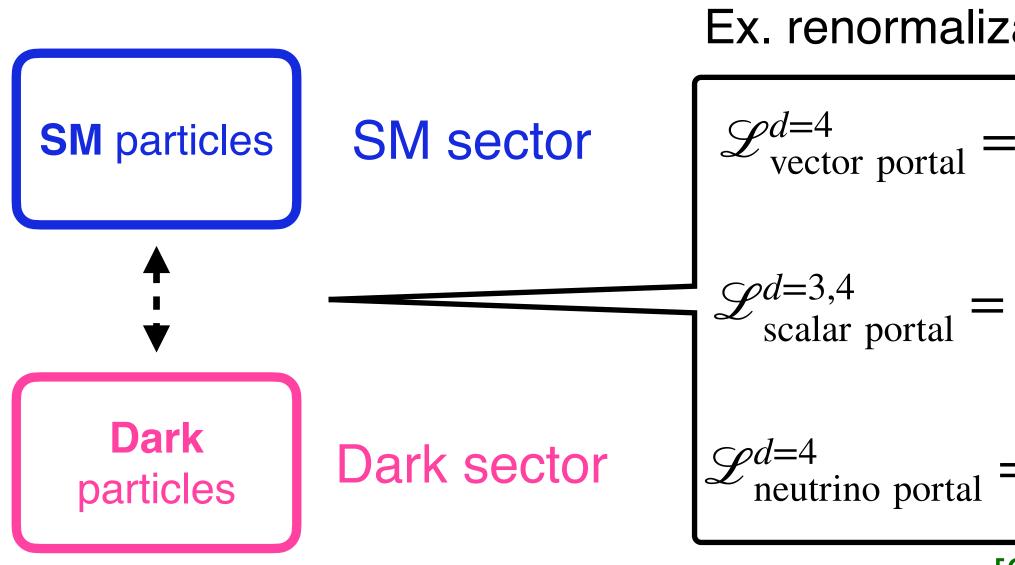


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$$= -H^{\dagger}H (A S + \lambda S^2) \qquad H \text{ SM Higgs field, } S \text{ SM singlet scalar field}$$
$$= -\sum y_{\nu}^{\alpha I} (\bar{L}_{\alpha} H) N_I \qquad L \text{ SM lepton field, } N \text{ SM singlet fermion field}$$



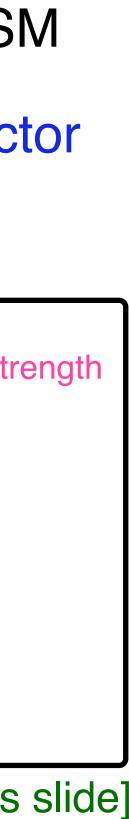


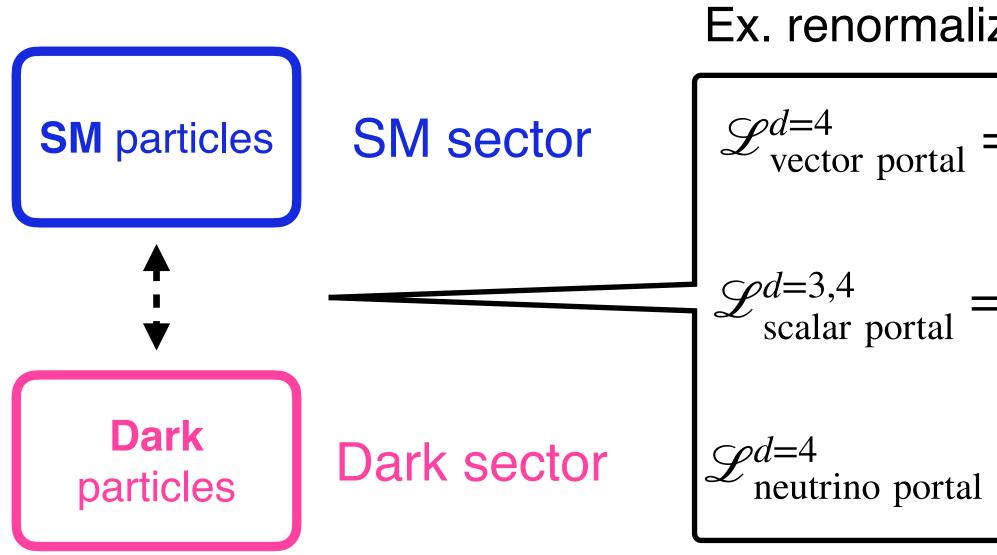
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[See Sebastian Trojanowski's, Bhupal Dev's, and Vedran Brdar's slide]





Neutrino portal can yield neutrino mass

Higgs portal, e.g., Twin Higgs models, can solve hierarchy problem

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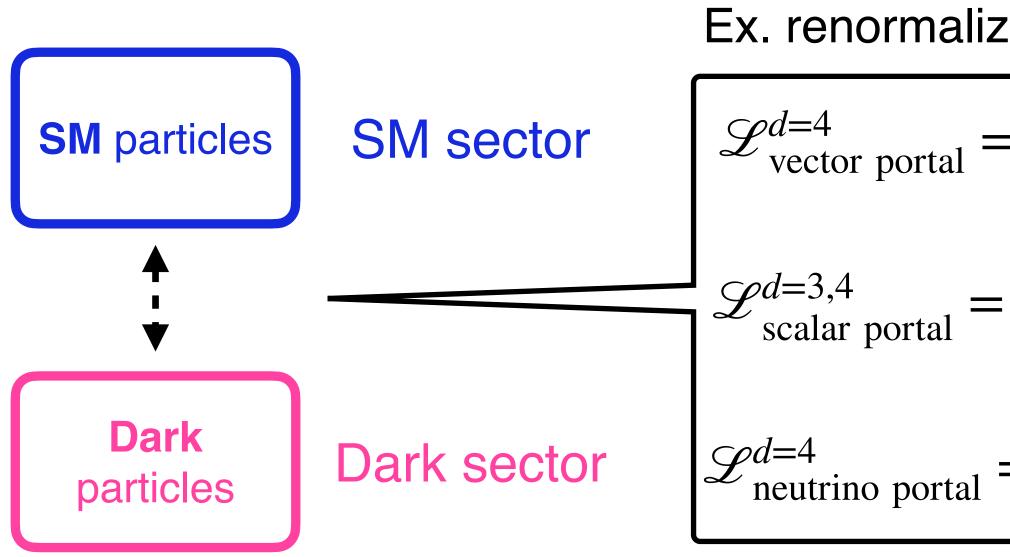
 $\mathscr{L}_{\text{vector portal}}^{d=4} = -\frac{1}{2\cos\theta_W} \frac{B_{\mu\nu}F'_{\mu\nu}}{B_{\mu\nu}F'_{\mu\nu}}$  $B_{\mu\nu}$  hypercharge field strength,  $F'_{\mu\nu}$  U(1)' filed strength  $\mathscr{L}_{\text{scalar portal}}^{d=3,4} = -H^{\dagger}H(AS + \lambda S^2)$ H SM Higgs field, S SM singlet scalar field

 $\mathscr{L}_{\text{neutrino portal}}^{d=4} = -\sum y_{\nu}^{\alpha I} (\bar{L}_{\alpha} H) N_{I}$ L SM lepton field, N SM singlet fermion field

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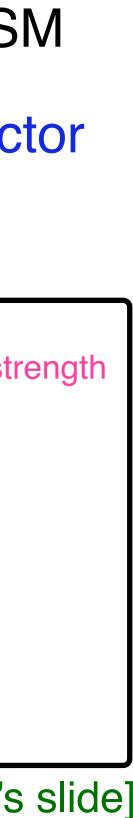
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  - [See Sebastian Trojanowski's, Bhupal Dev's, and Vedran Brdar's slide]
- Neutrino portal can yield neutrino mass

Dark sector potentially includes DM candidates



[attractive features of thermal DM]





[attractive features of thermal DM]

- Freeze-out mechanism can yield DM abundance  $\Rightarrow$  DM-SM reaction cross section can be large





[attractive features of thermal DM]

- Viable DM mass range is limited  $\Rightarrow$  Thermal mass window (~1 MeV to ~100 TeV)

#### • One of the DM candidates is thermal DM, i.e., DM is thermalized with SM particles in the early universe

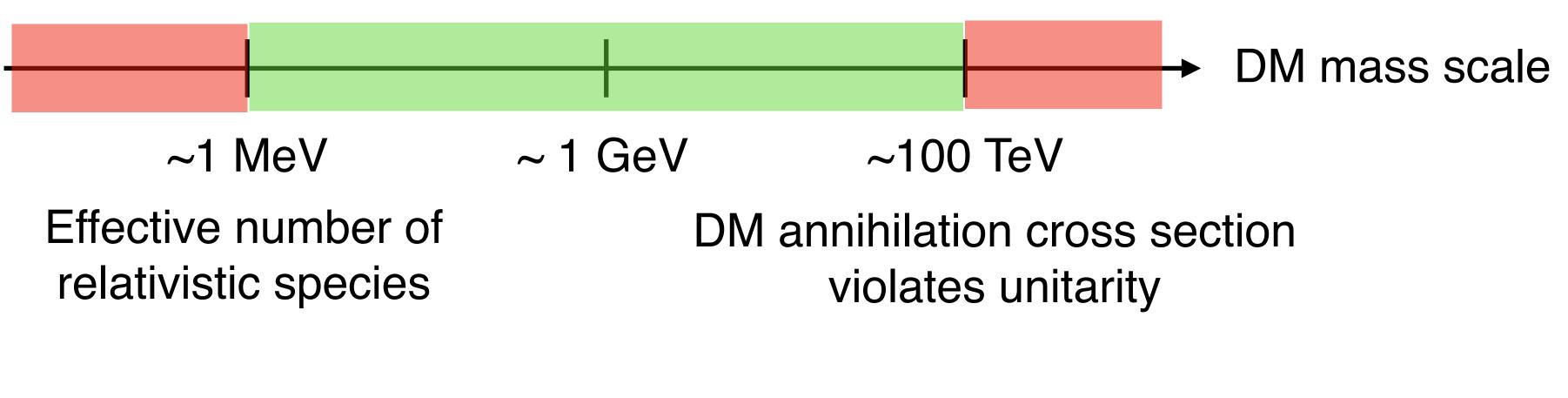
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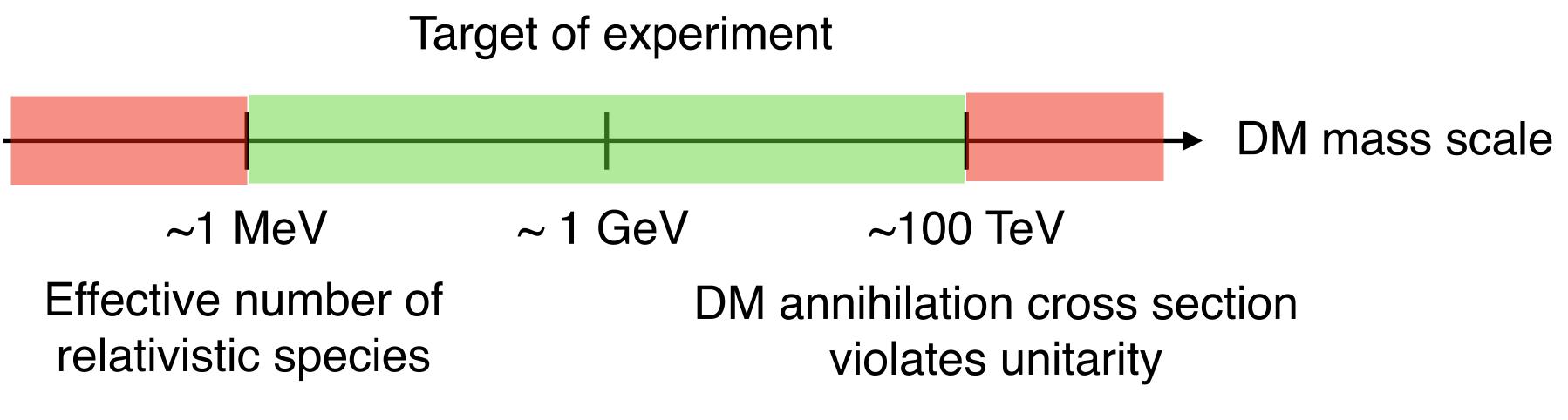






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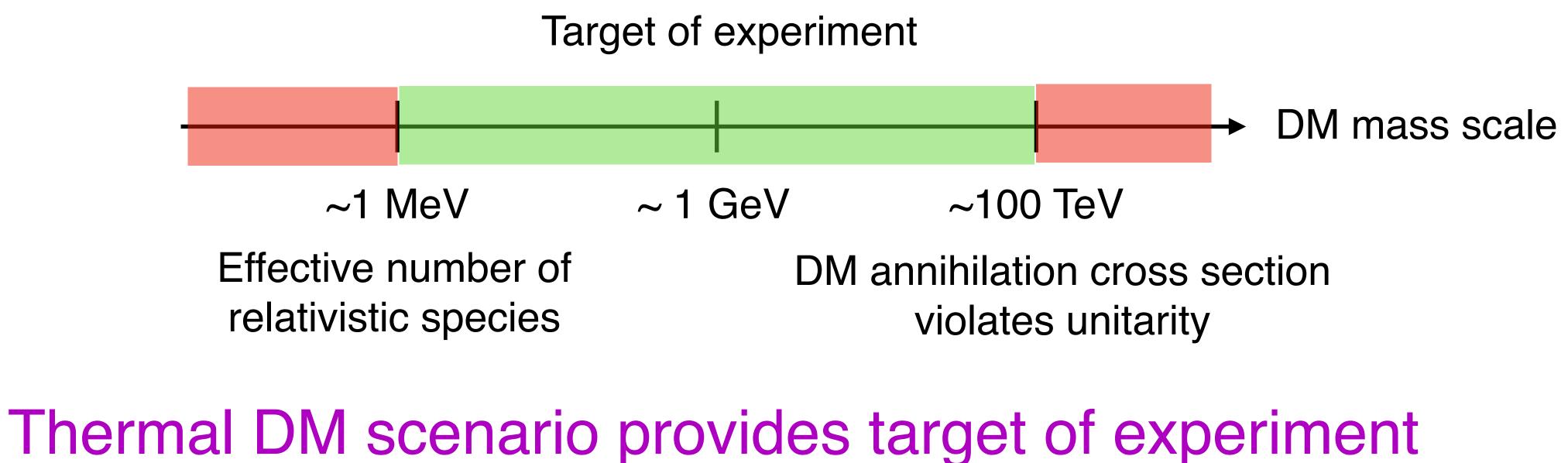






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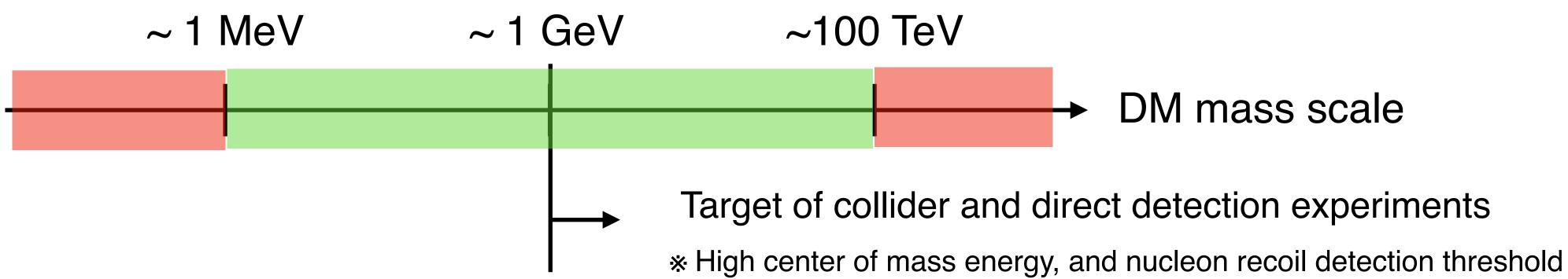






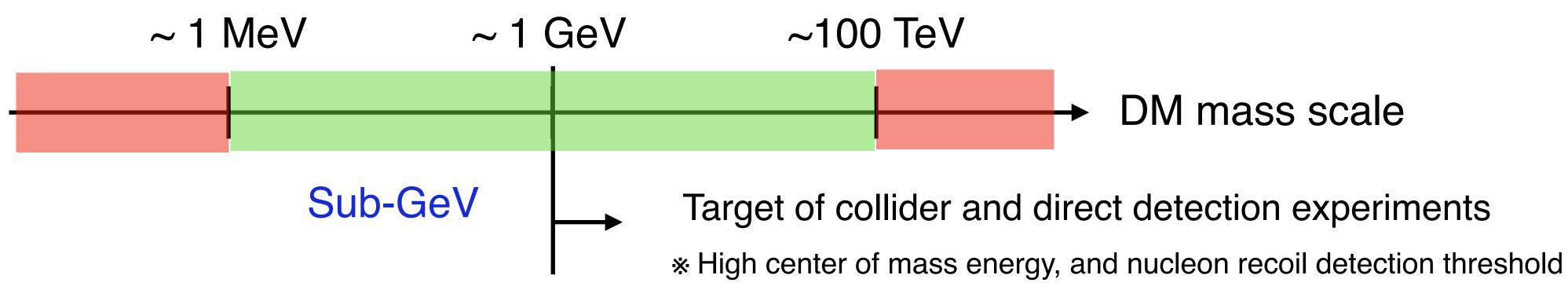
#### Sub-GeV dark matter

#### ~1 MeV



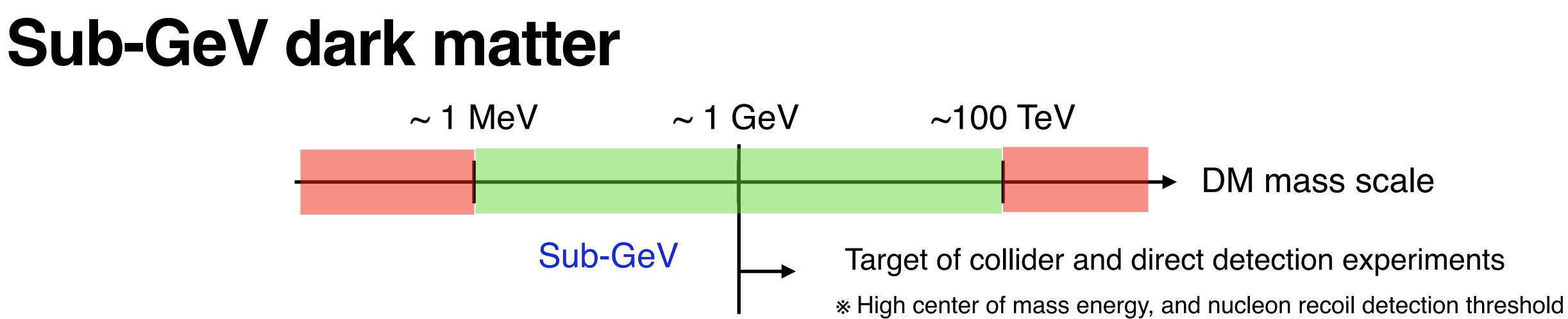


## Sub-GeV dark matter



Sub-GeV DM is also a DM candidate but is feebly coupled with SM particles

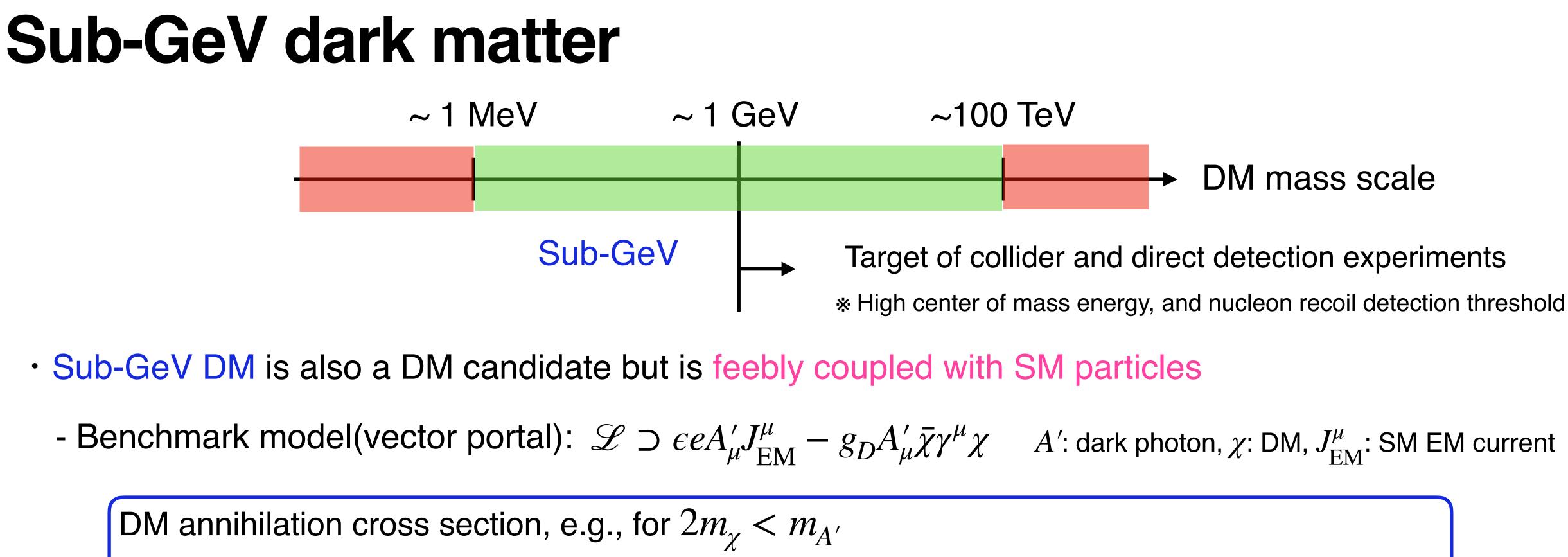


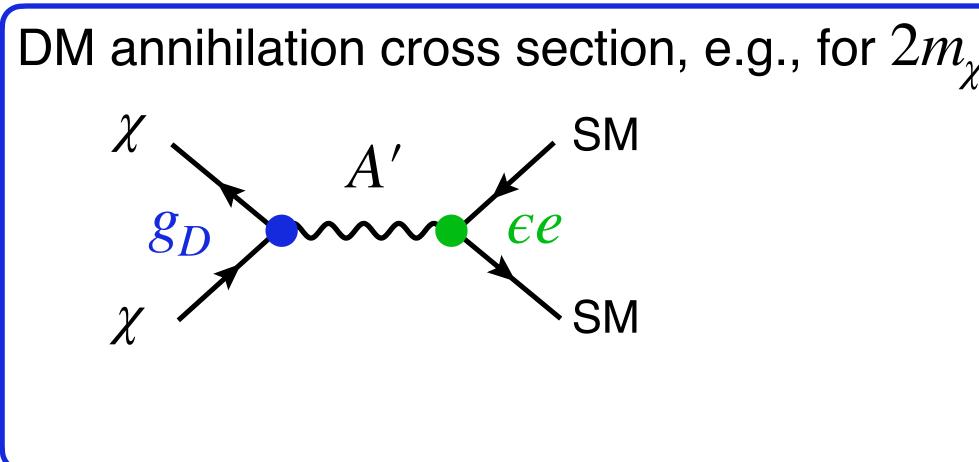


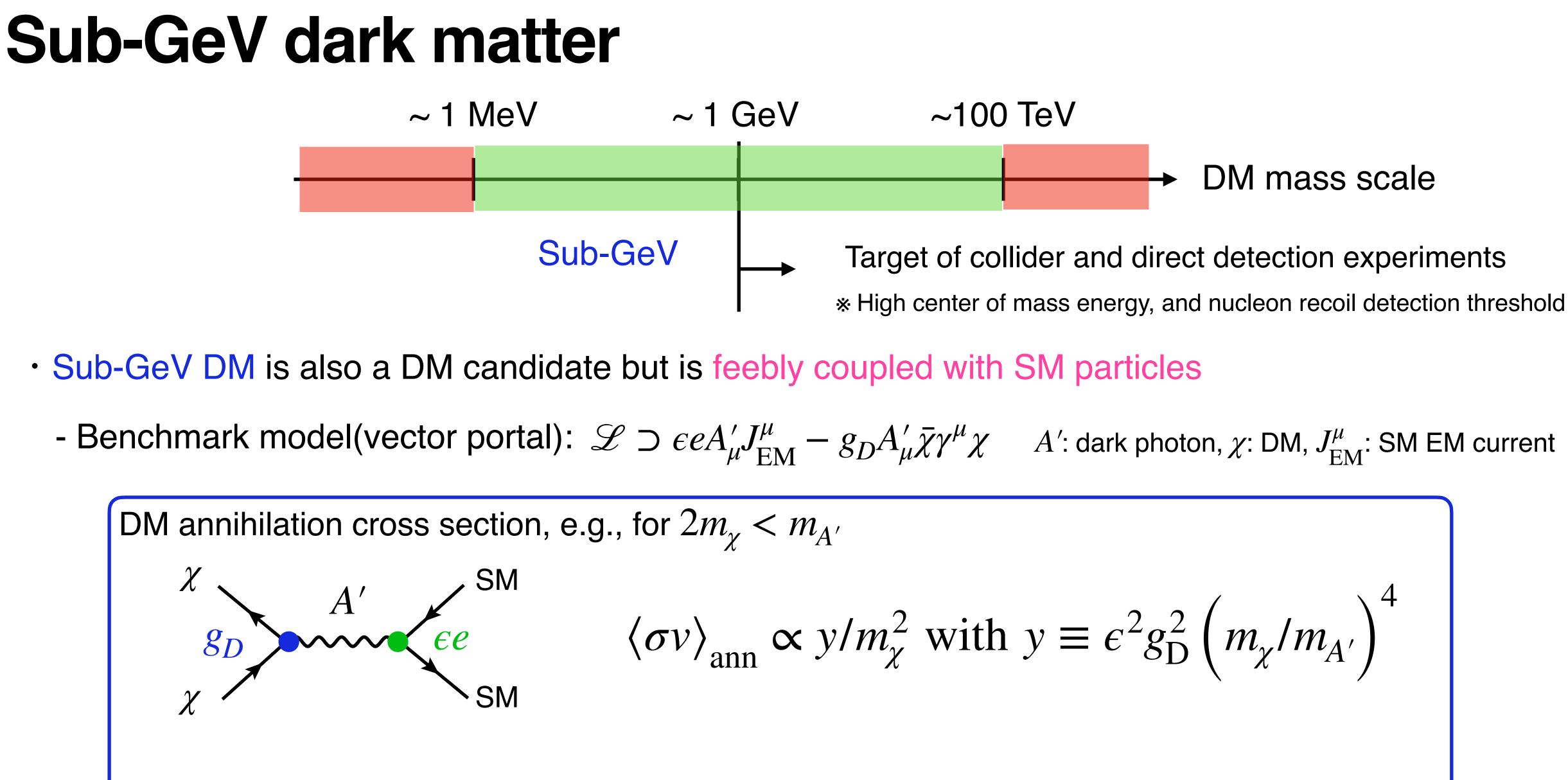
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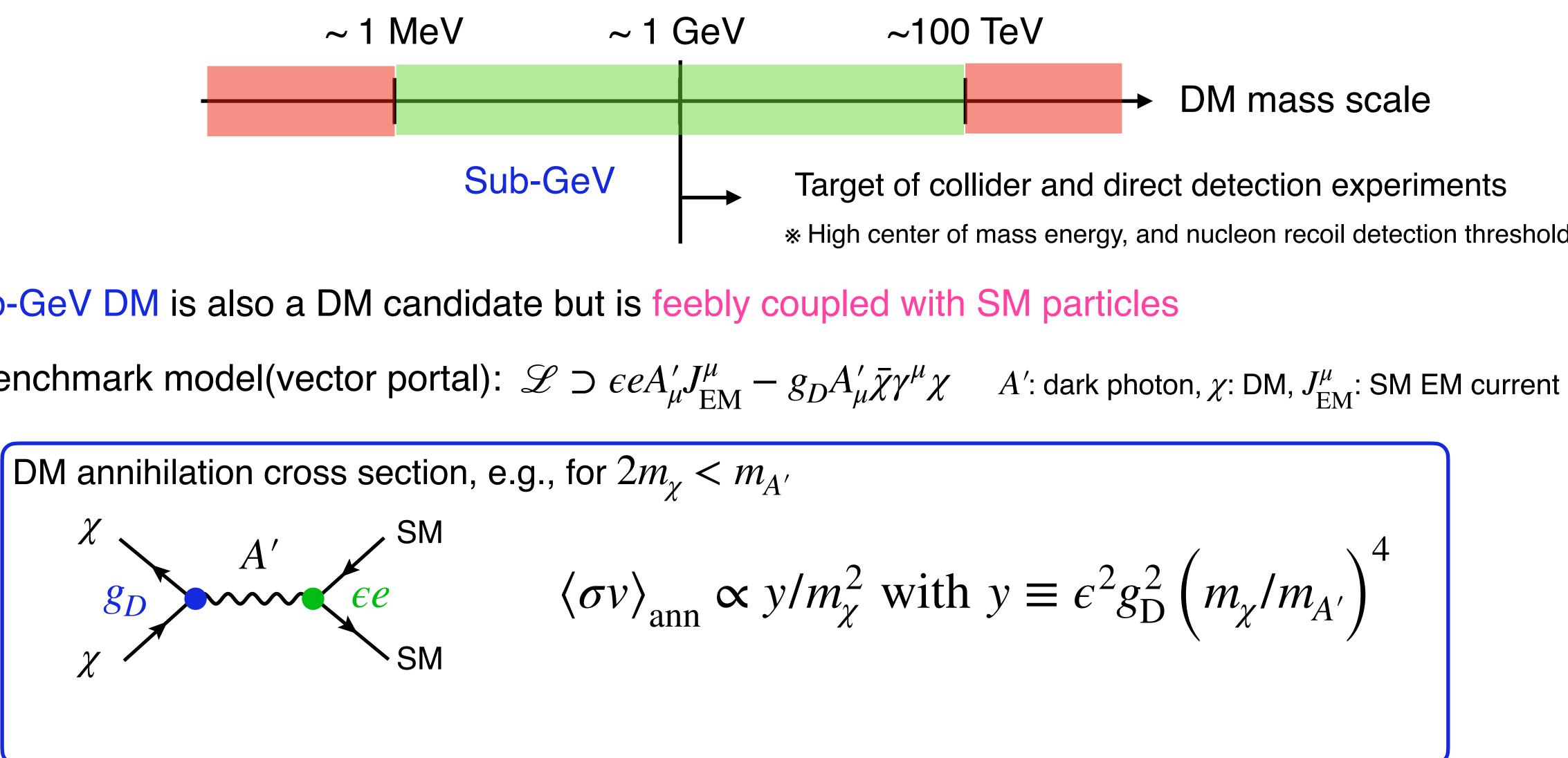
- Benchmark model (vector portal):  $\mathscr{L} \supset \epsilon e A'_{\mu} J^{\mu}_{EM} - g_D A'_{\mu} \bar{\chi} \gamma^{\mu} \chi$  A': dark photon,  $\chi$ : DM,  $J^{\mu}_{EM}$ : SM EM current

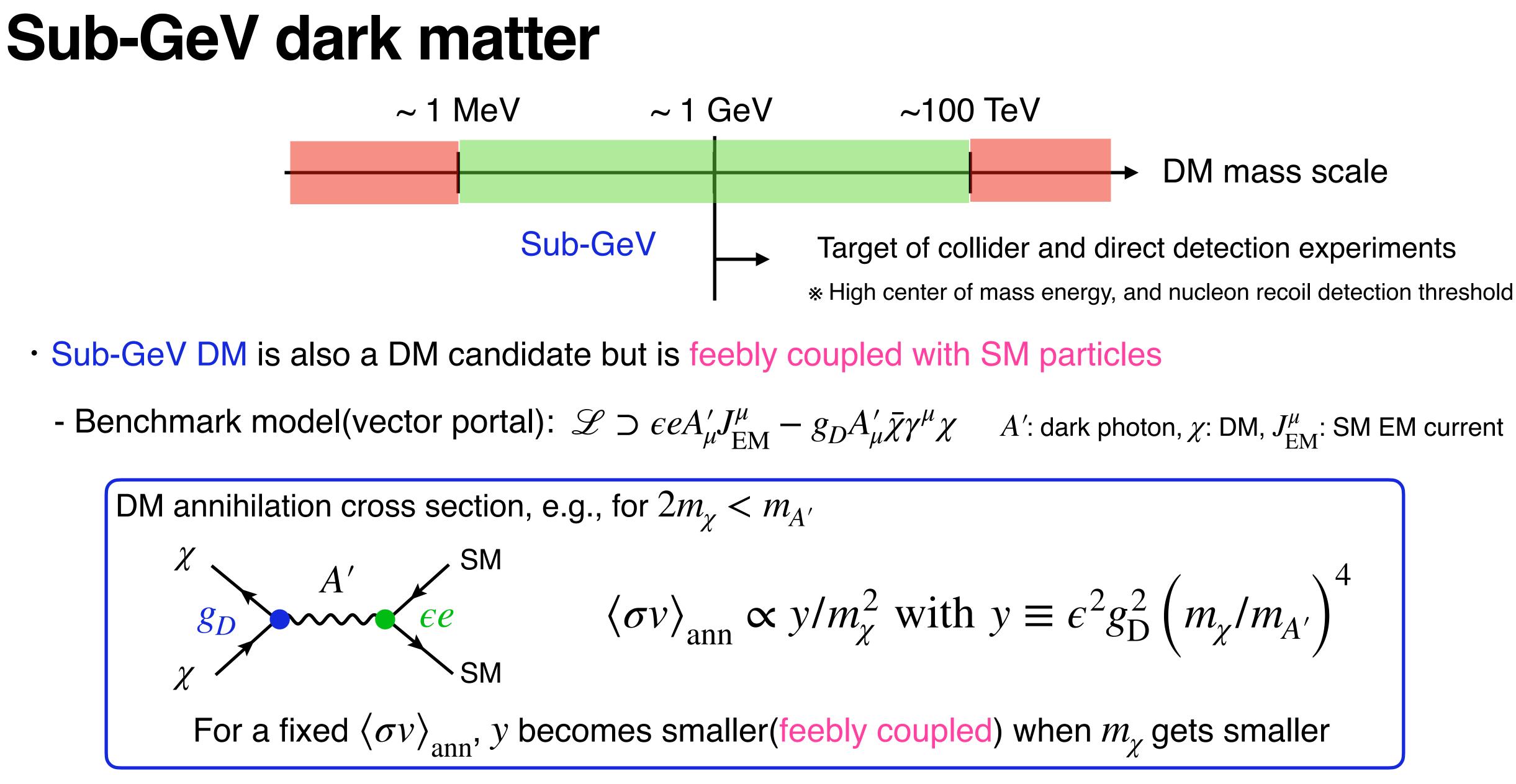


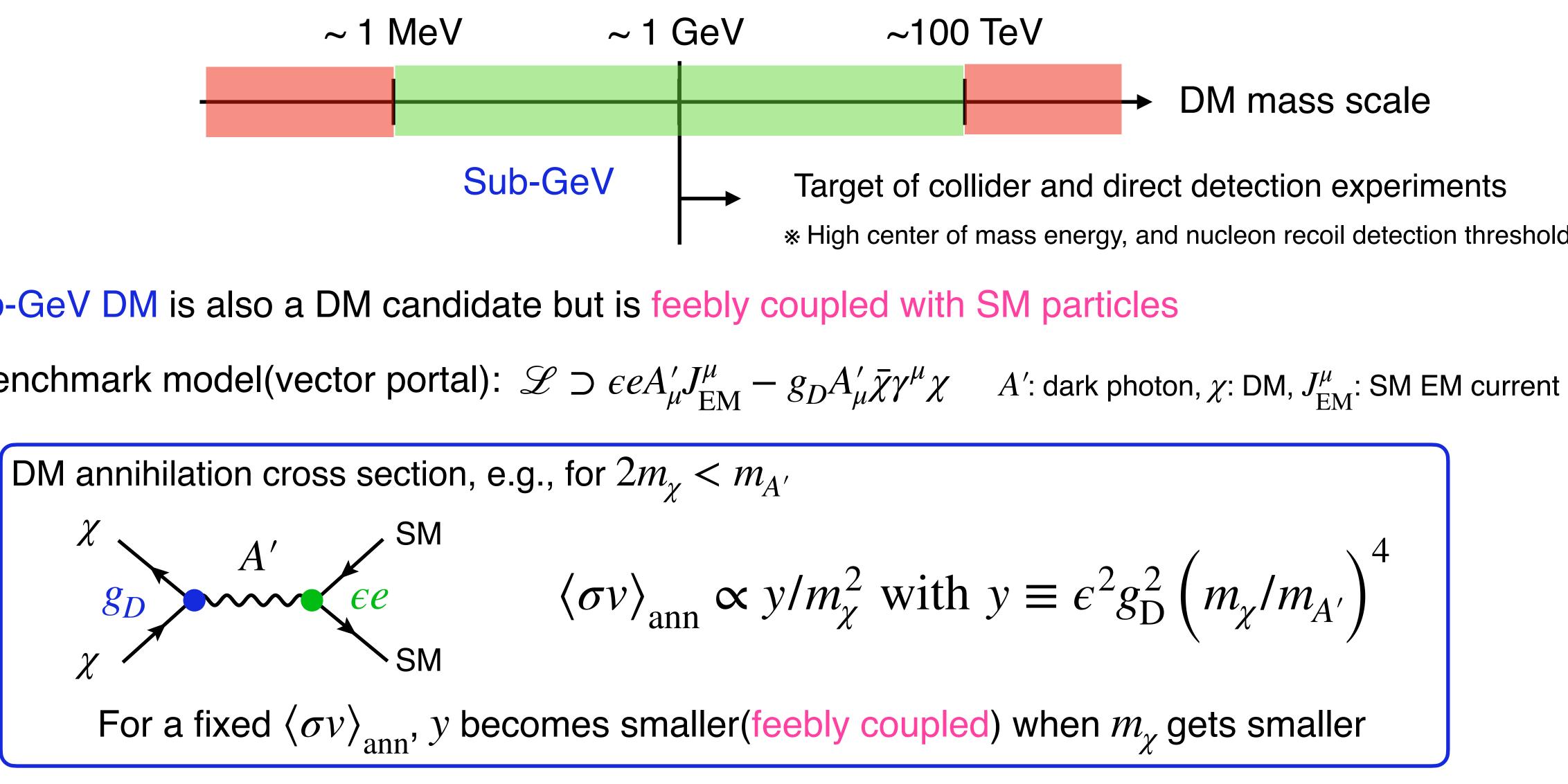


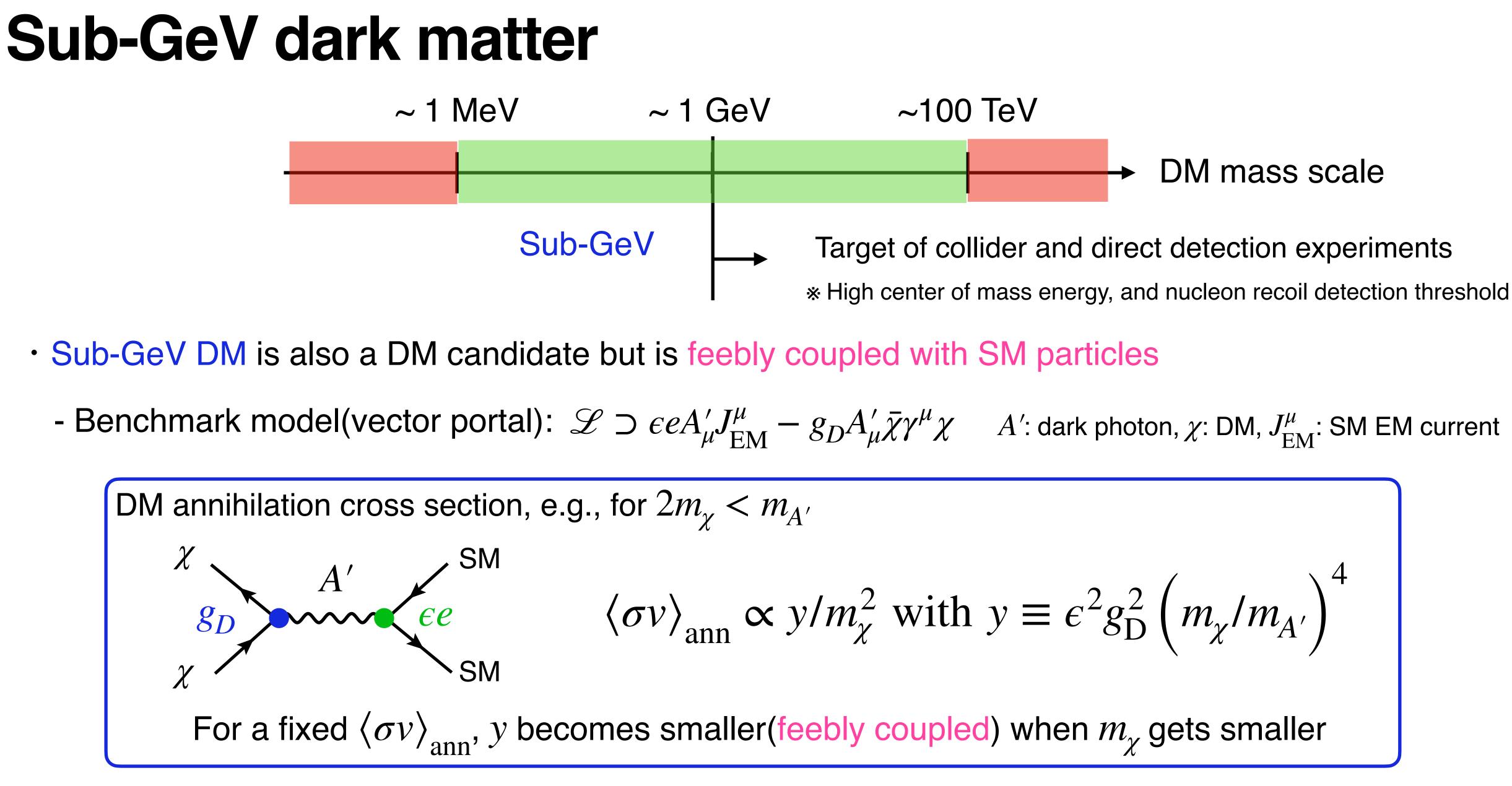


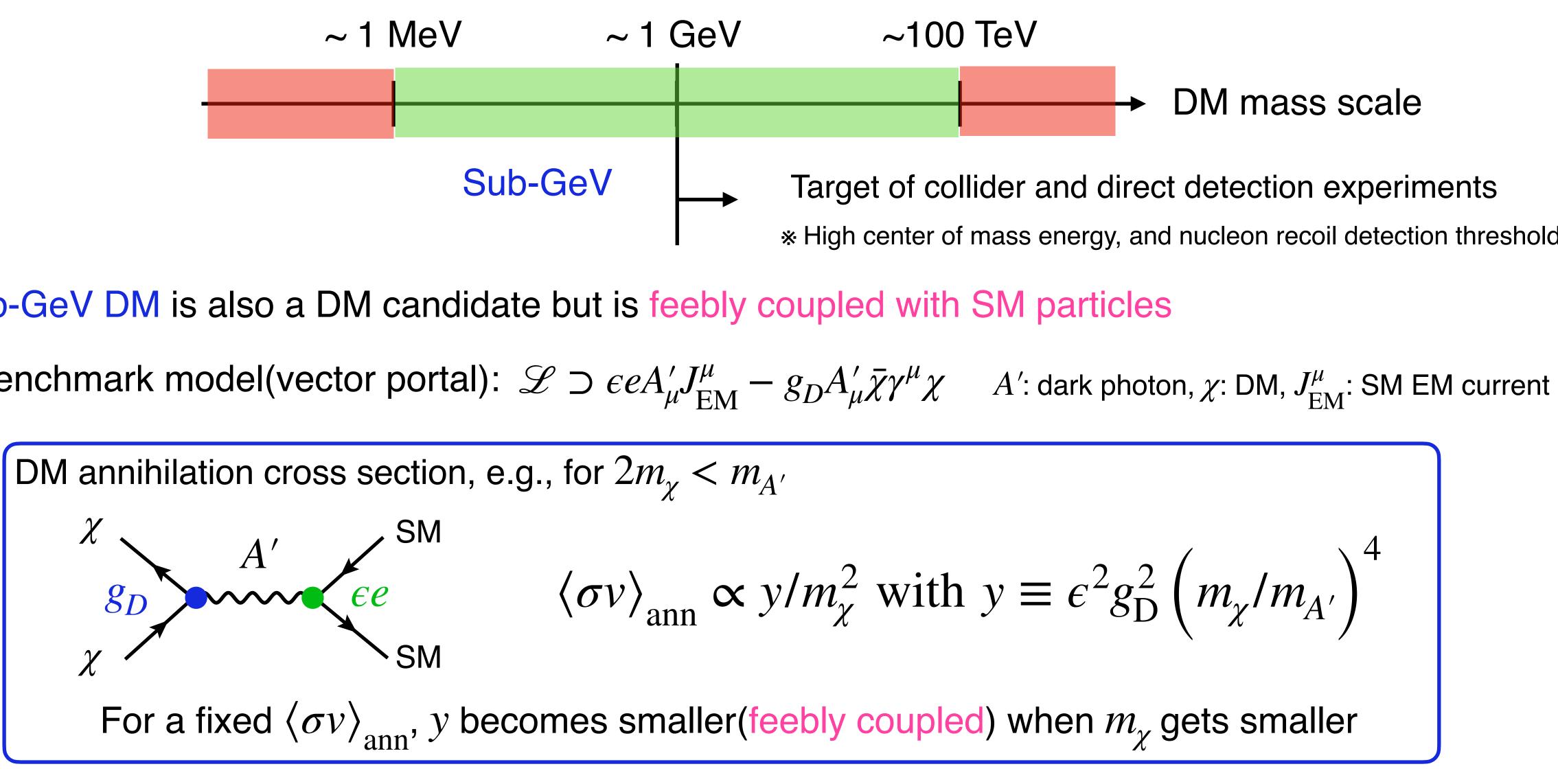












#### Hight intensity experiments are needed to search for Sub-GeV dark states

Beam dump experiments are high-intensity experiments and are sensitive to Sub-GeV particles

- - three components of beam dump experiment:

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Beam particle

e.g., electron, proton, and muon

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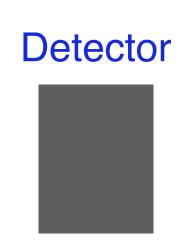
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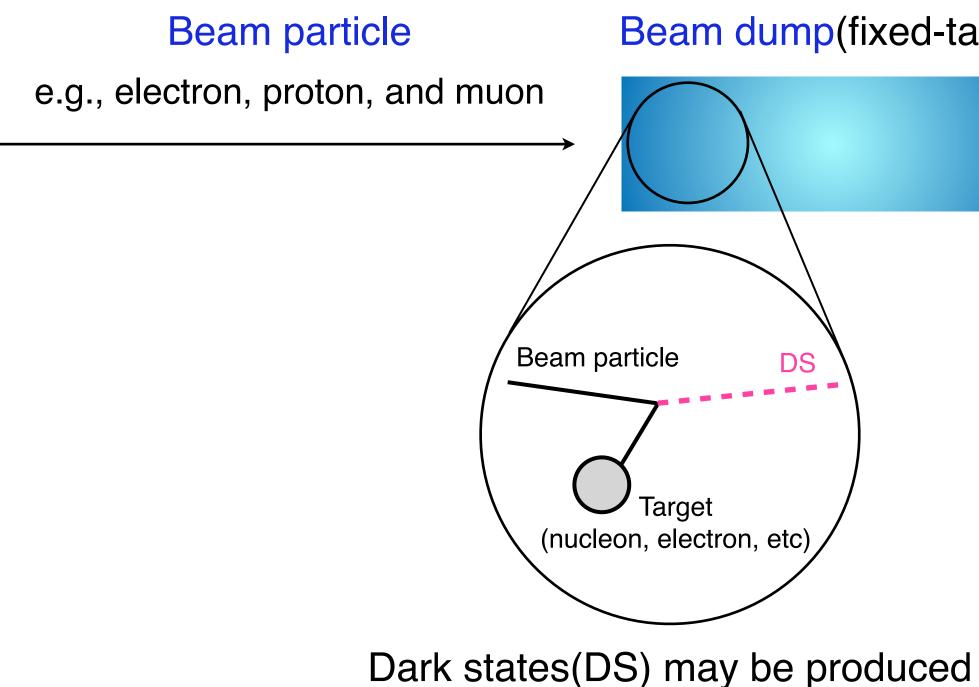
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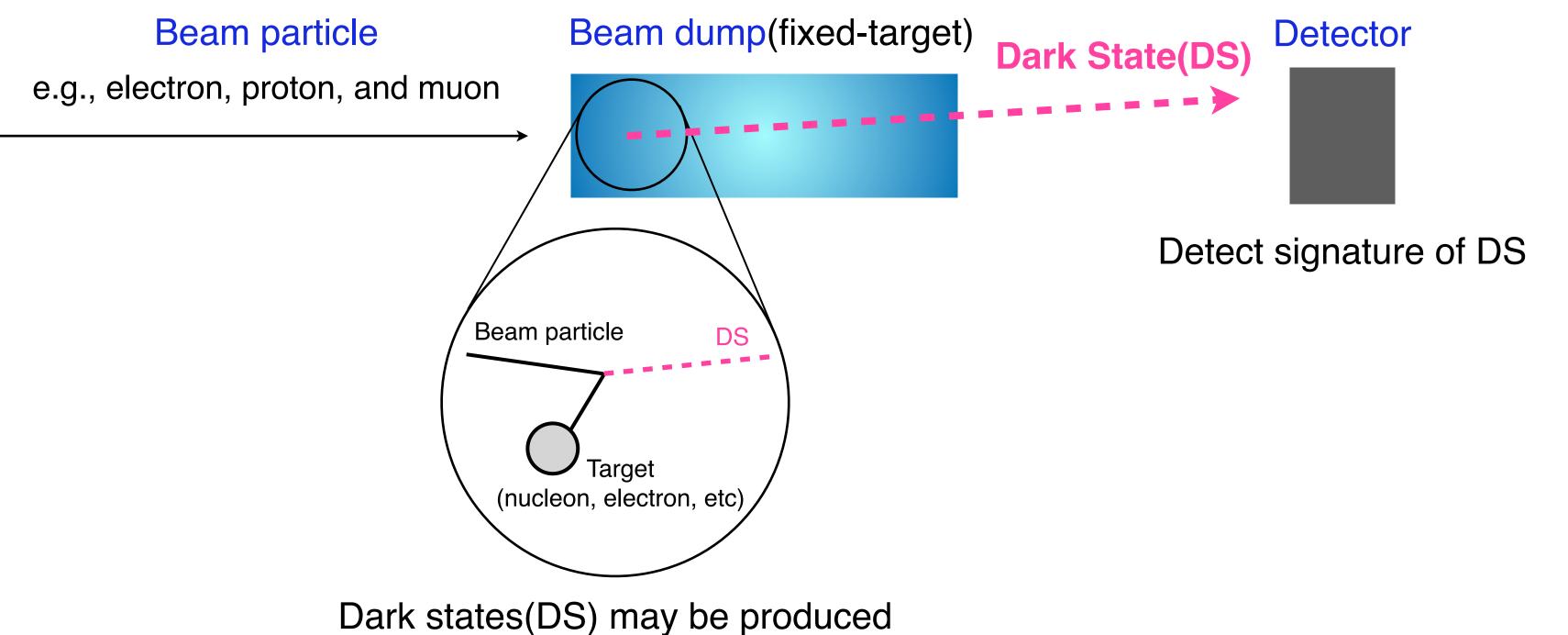


Beam dump experiments are high-intensity experiments and are sensitive to Sub-GeV particles

Detector

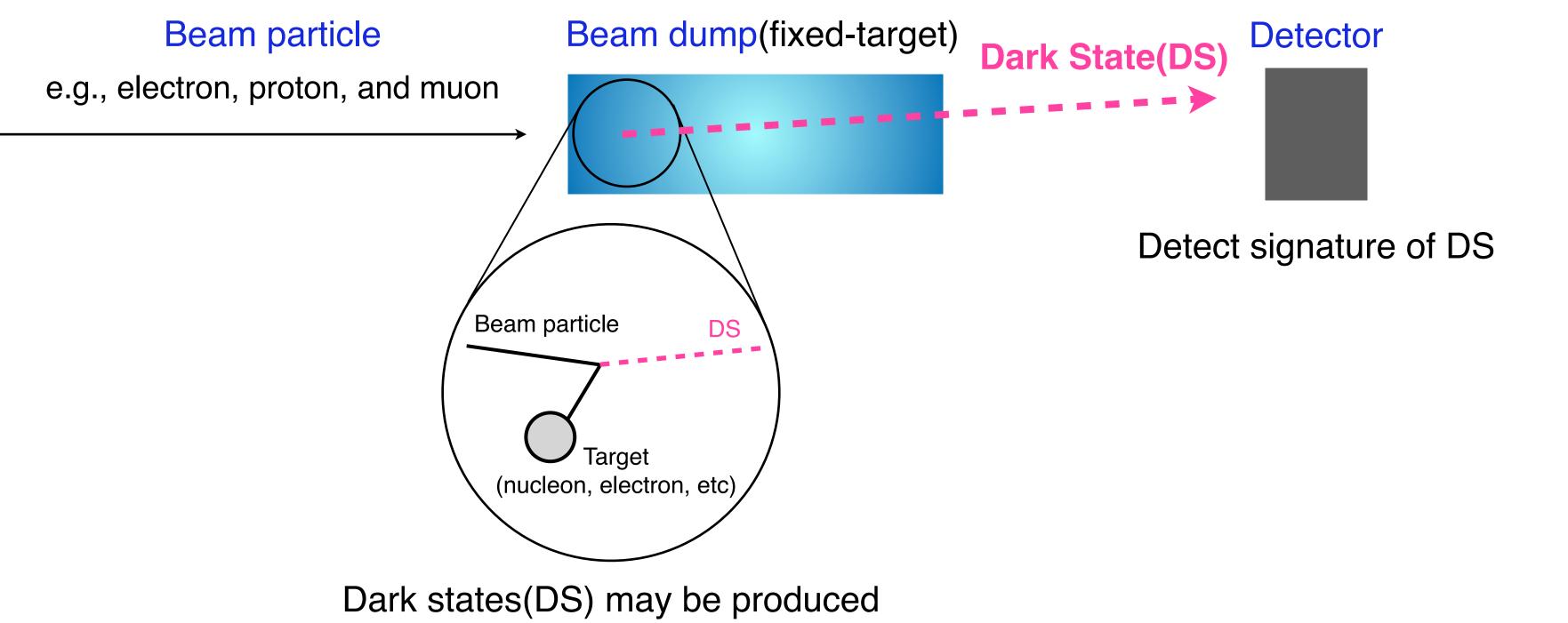
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- - three components of beam dump experiment:



[My talk's focus]

Detection of dark state signatures produced by beam-target collision

· Beam dump experiments are high-intensity experiments and are sensitive to Sub-GeV particles

## Outline

- Introduction
  - dark sector and Sub-GeV dark matter
  - beam dump(fixed-target) experiment
- Key features of beam dump experiment
- A classification of beam dump experiment
- Sensitivity of beam dump experiments at future accelerators
- Summary

#### Key features of beam dump experiment (1)

[Center of mass energy]

$$\sqrt{s} = \sqrt{m_{\text{beam}}^2}$$

where  $m_{\text{beam}}$  is mass of beam particle,  $m_{\text{target}}$  is mass of target particle, and  $E_{\text{beam}}$  is beam energy

 $+ m_{\text{target}}^2 + 2E_{\text{beam}}m_{\text{target}}$ 

\* This feature is determined only by beam and target properties



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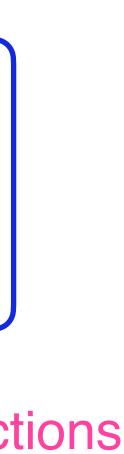
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 Center of mass energy is smaller than collid are kinematically allowed

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· Center of mass energy is smaller than collider energy scale, but Sub-GeV dark particle productions

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Ex. Beam = electron (
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 = 0.5 MeV),

$$\sqrt{s} = \sqrt{m_{\text{beam}}^2} +$$

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where  $m_{\text{beam}}$  is mass of beam particle,  $m_{\text{target}}$  is mass of target particle, and  $E_{\text{beam}}$  is beam energy

• Center of mass energy is smaller than collider energy scale, but Sub-GeV dark particle productions

target = nucleon ( $m_{target}$  = 1 GeV),  $E_{beam}$  = 10 GeV

 $+ m_{\text{target}}^2 + 2E_{\text{beam}}m_{\text{target}} \simeq 5 \text{ GeV}$ 

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[Center of mass energy]

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Ex. Beam = electron (
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### Boosted Sub-GeV particles can be produced in beam dump

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(# of produced DS) = (DS production cross section  $[L^2]$ )

X



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Beam dump experiment is high luminosity frontier

(# of produced DS) = (DS production cross section  $[L^2]$ )

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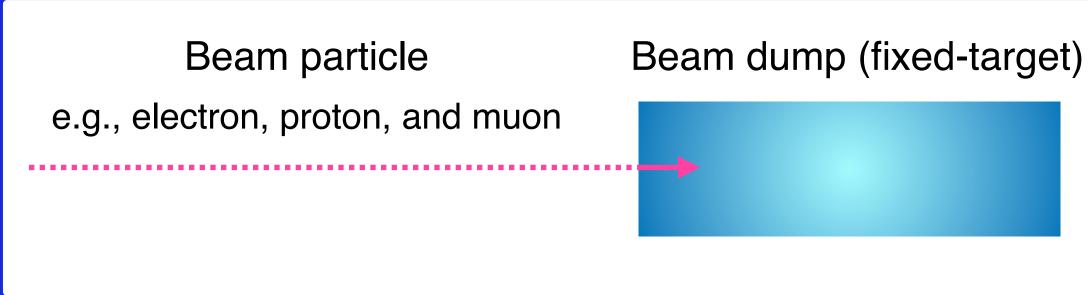
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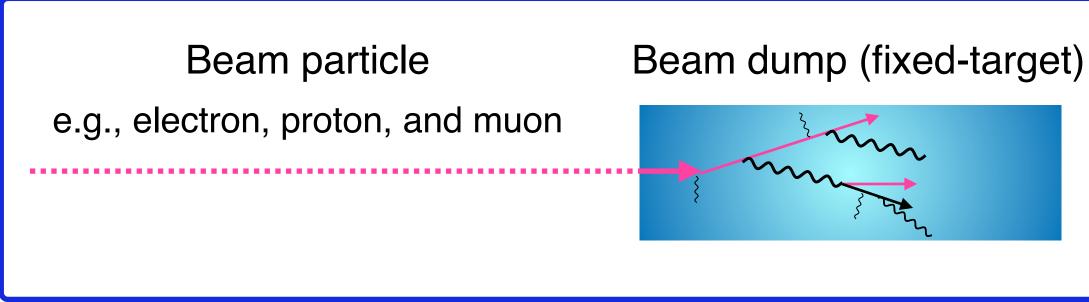
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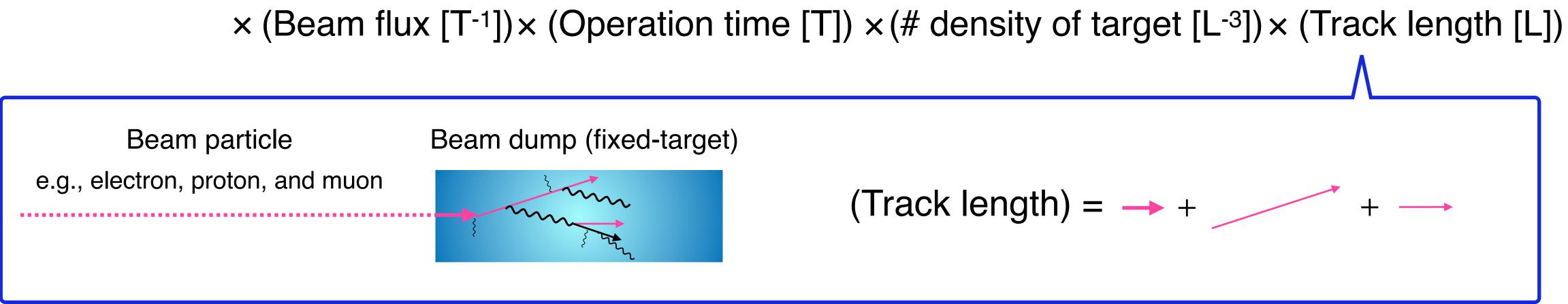
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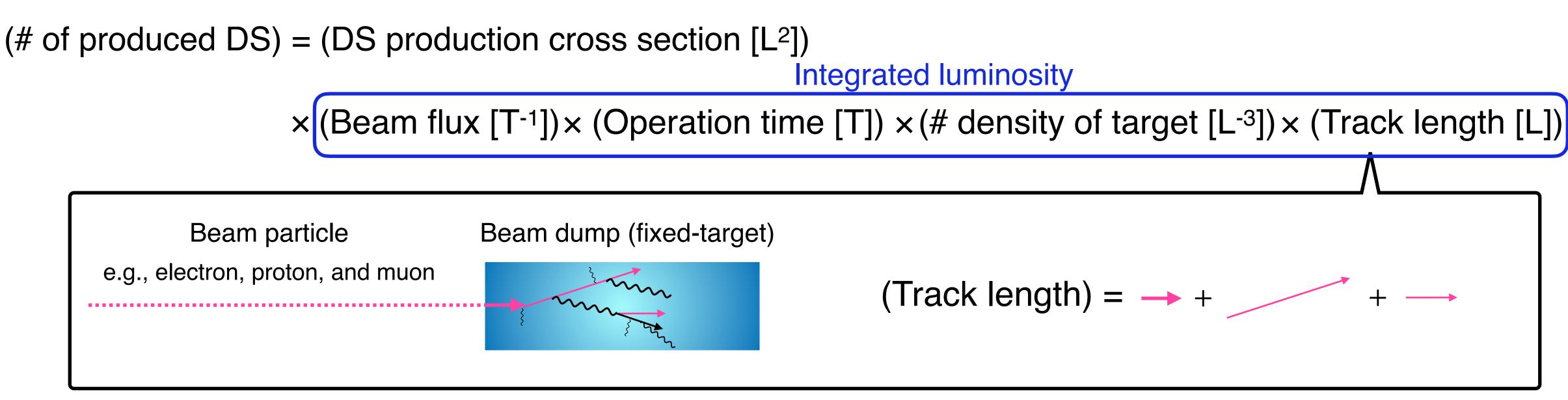
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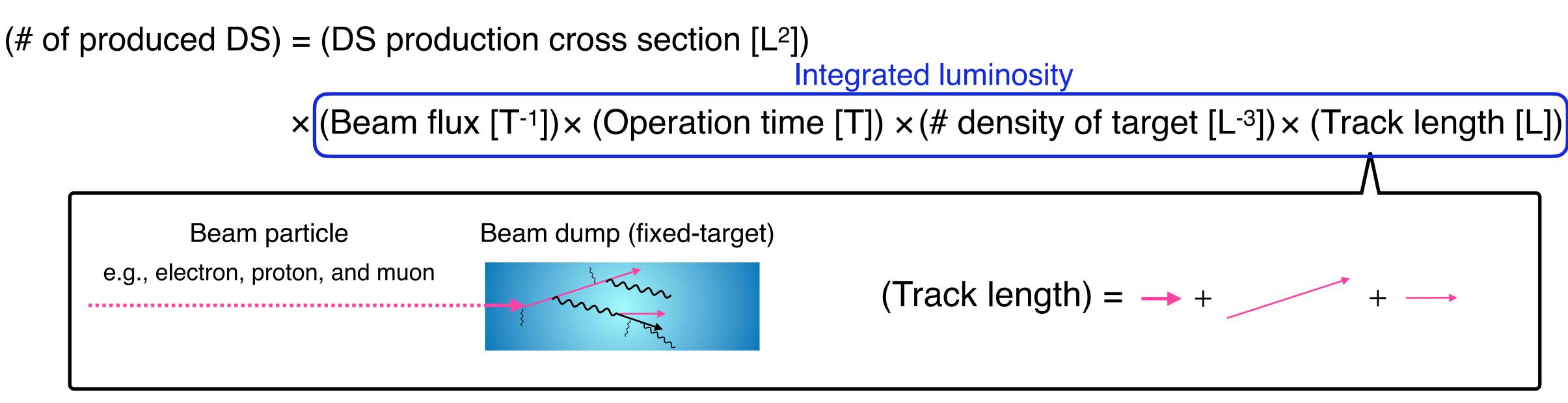
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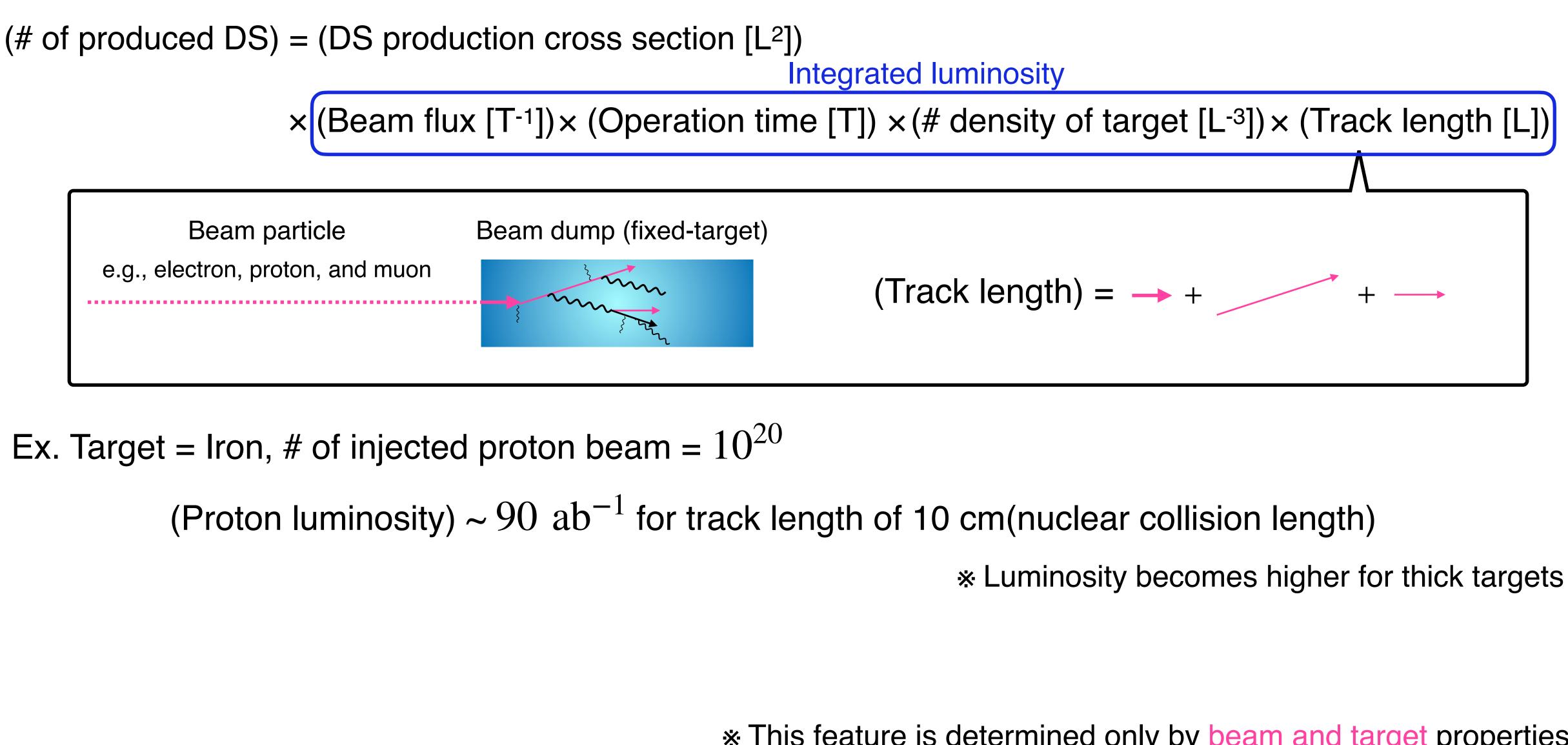


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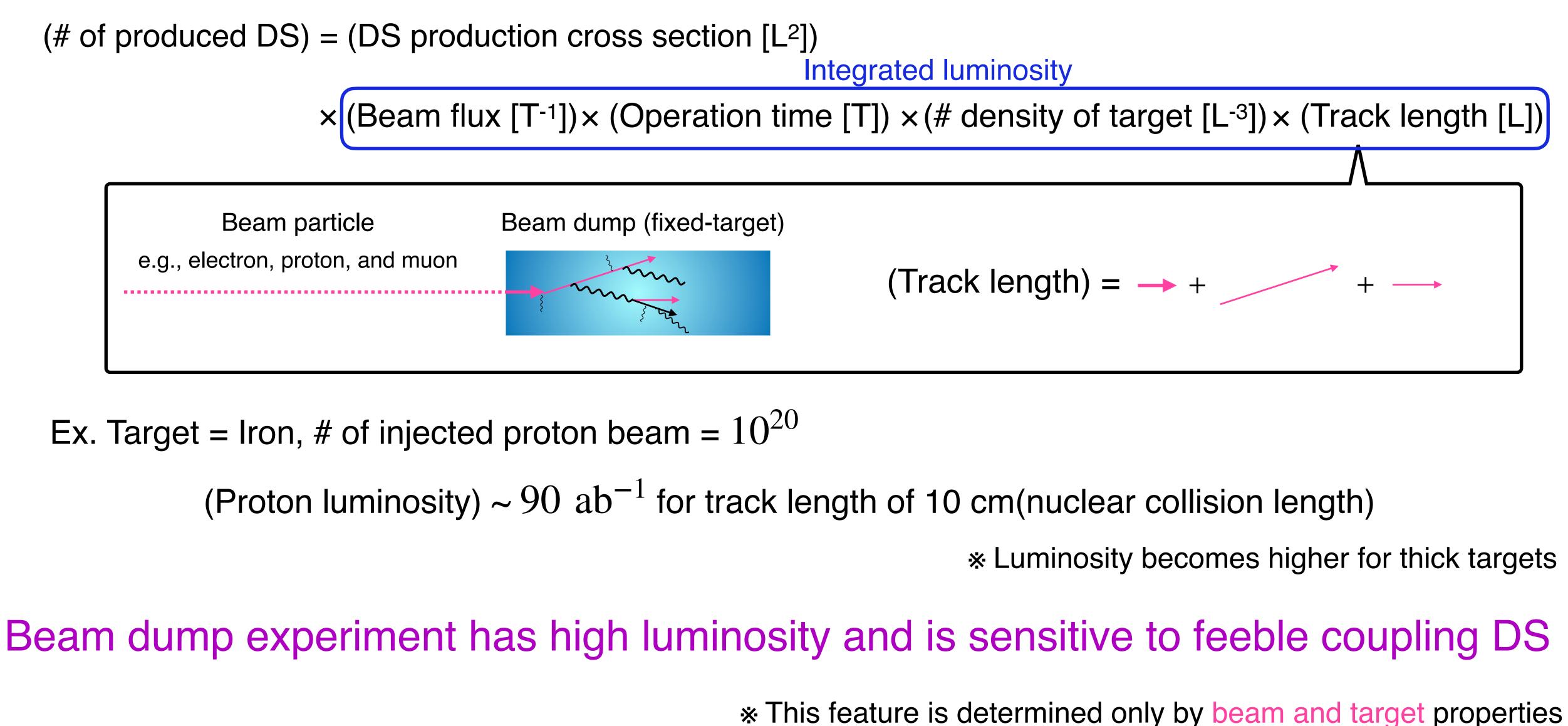


Ex. Target = Iron, # of injected proton beam =  $10^{20}$ 

Beam dump experiment is high luminosity frontier



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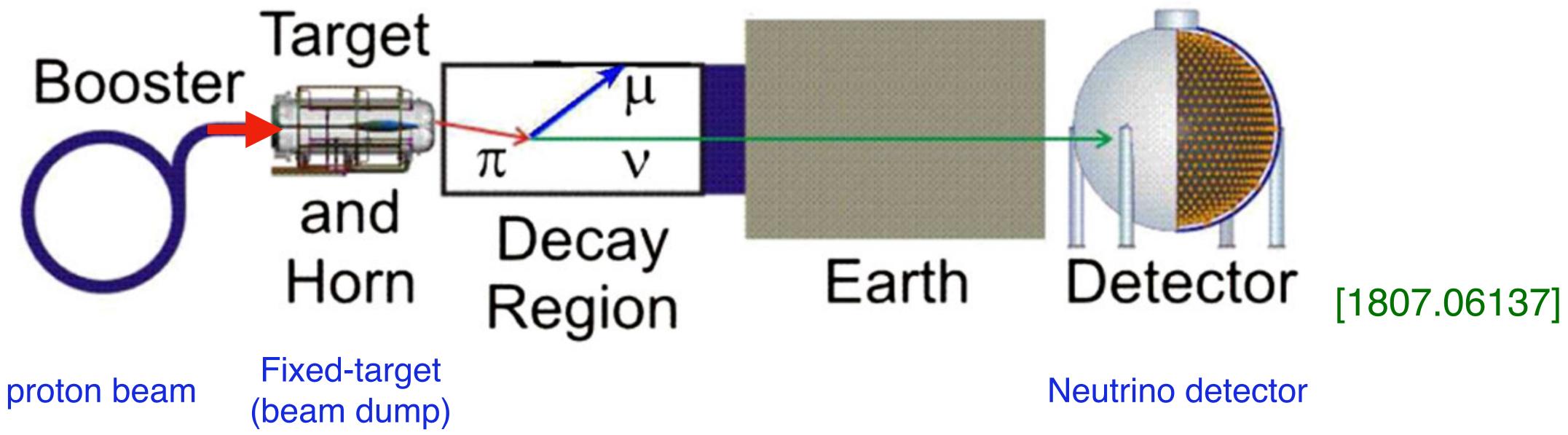
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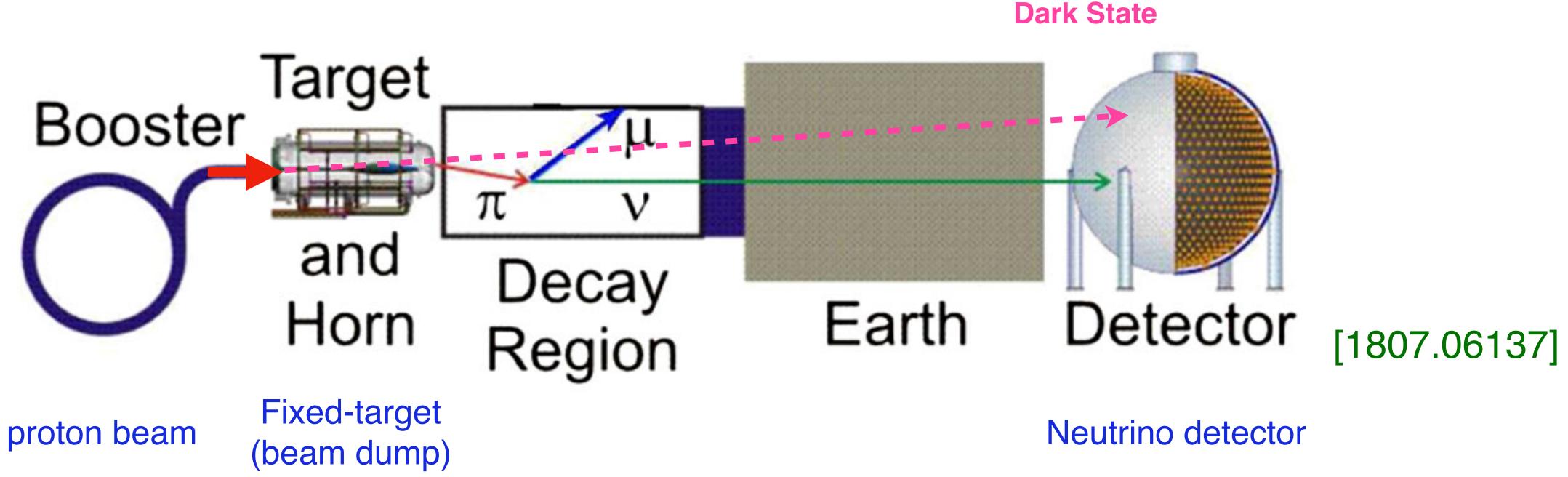
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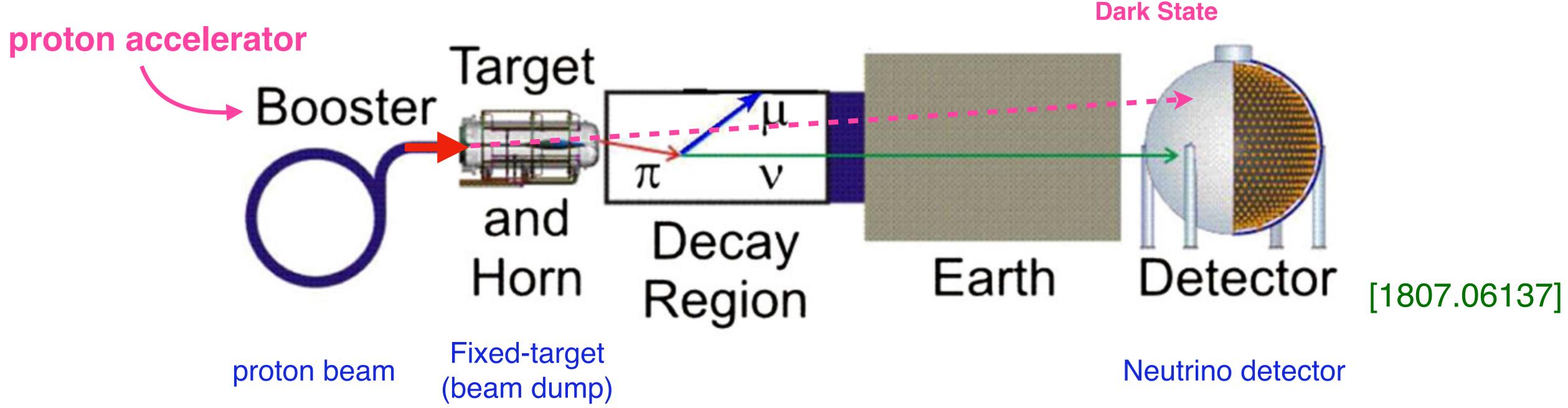




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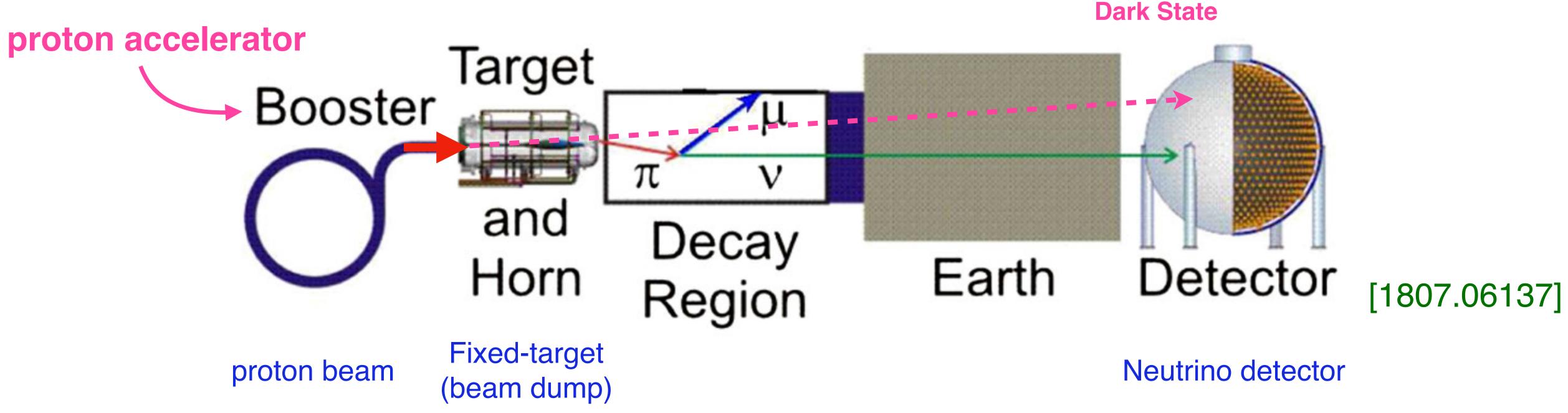




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Beam dump experiments can use accelerator facilities of neutrino experiments





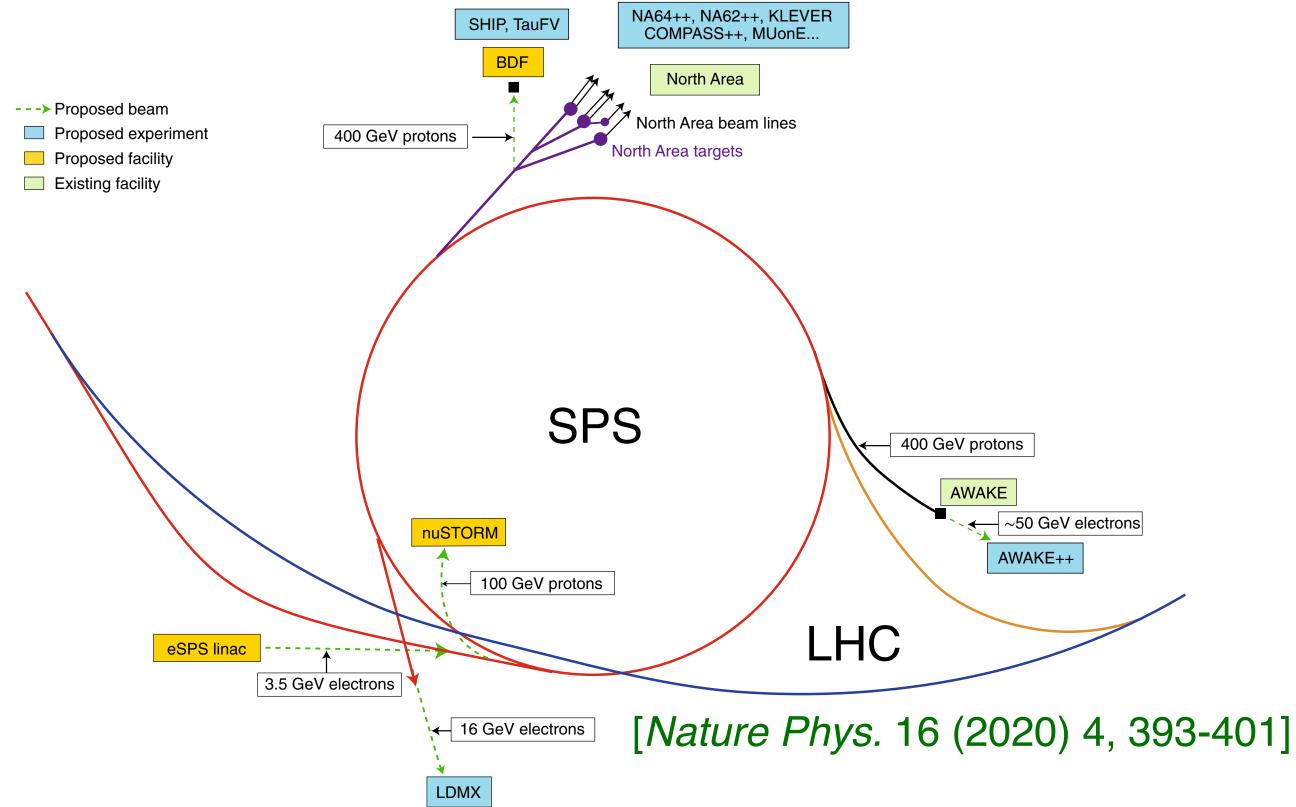


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NA62, NA64, etc  $\Rightarrow$  leverage extracted beam from CERN Super Proton Synchrotron(SPS) accelerator



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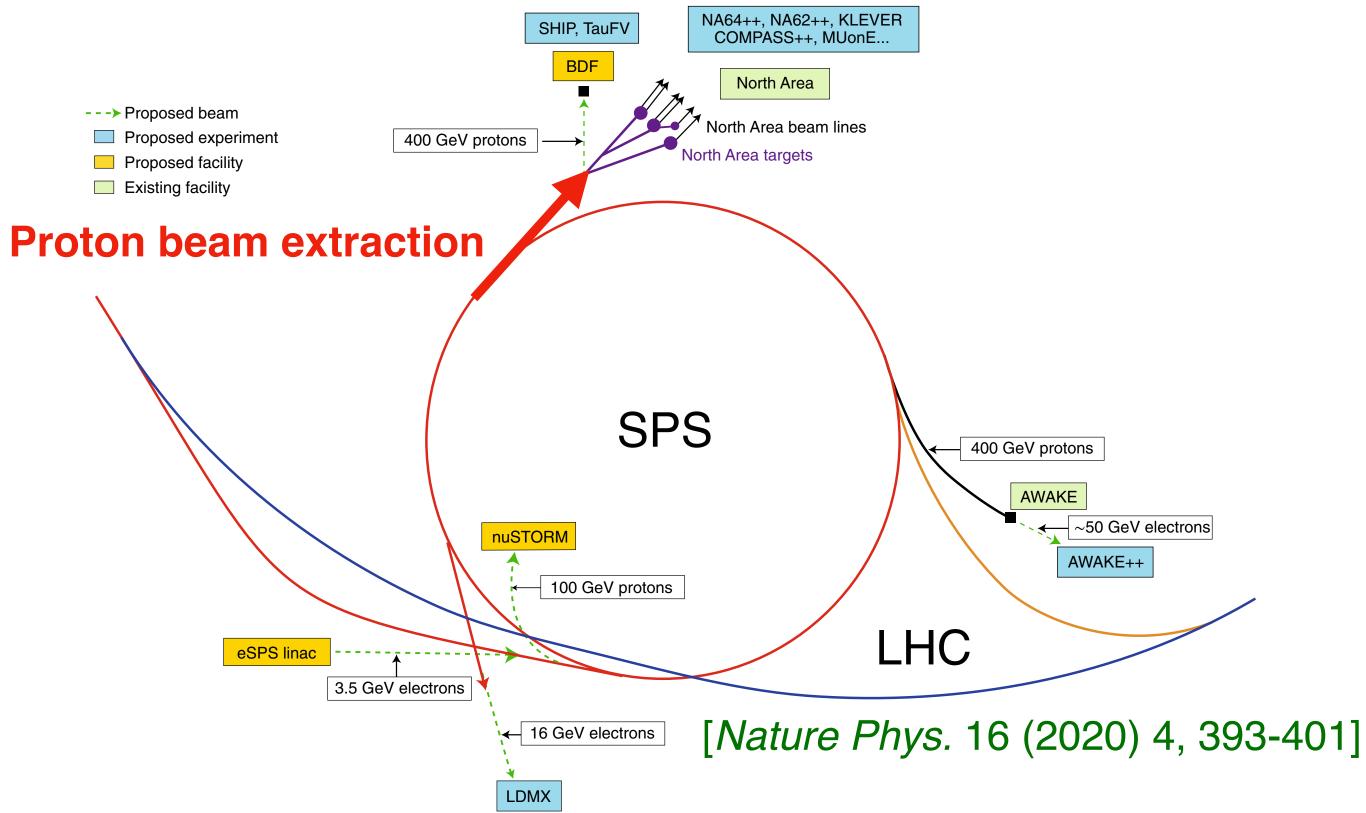


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~50 GeV electrons



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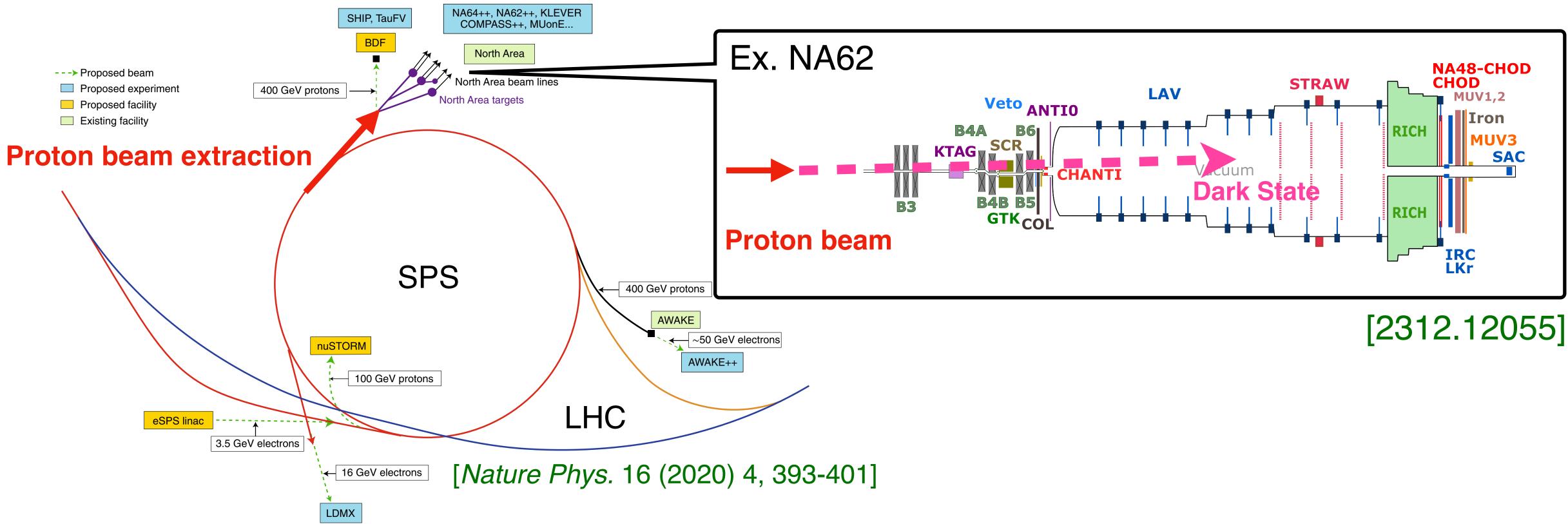


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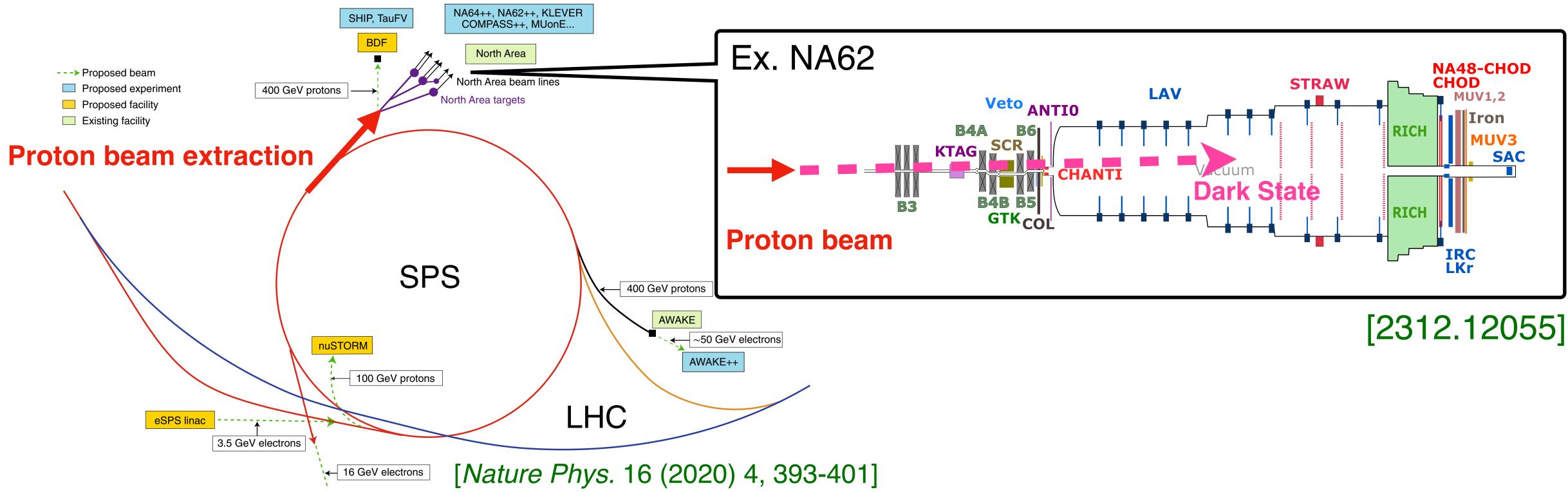
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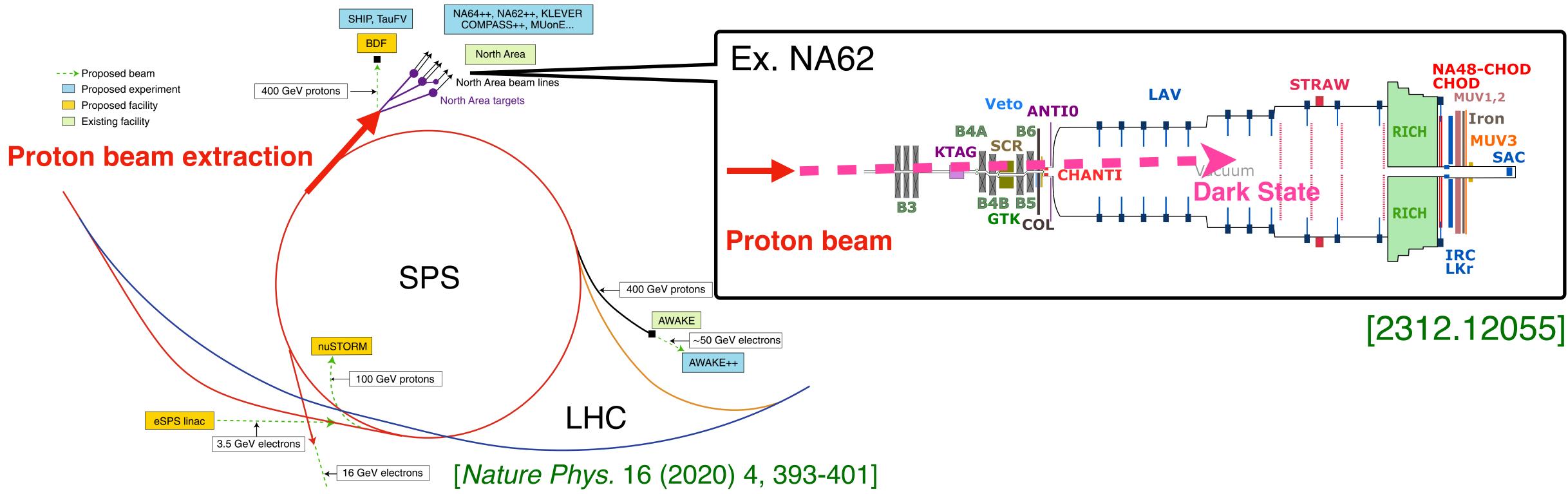
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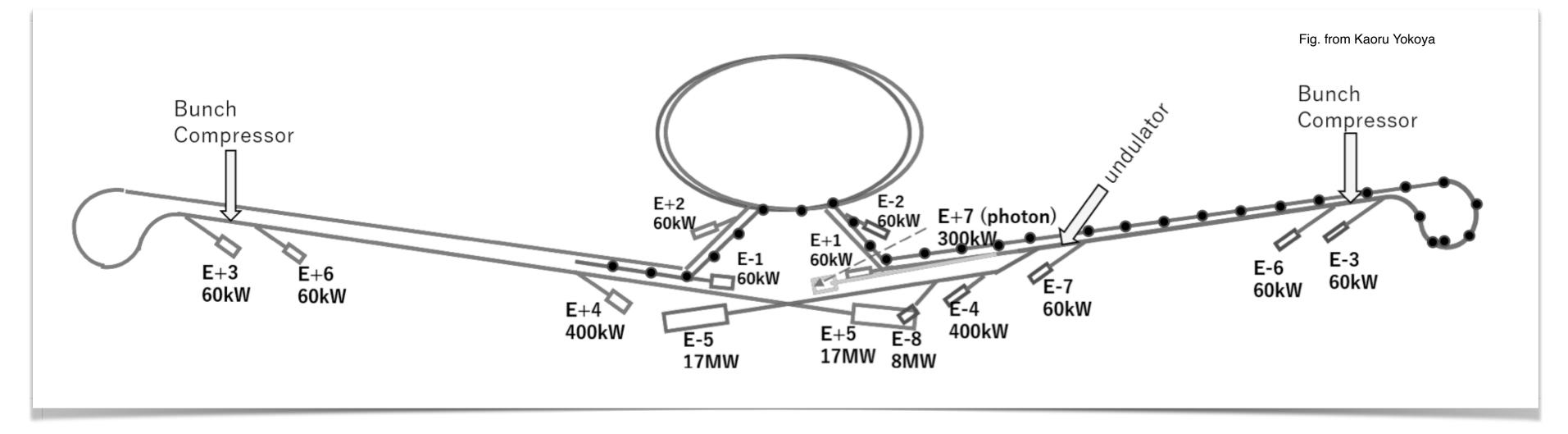
 $\Rightarrow$  beam dump experiments will be conducted at **future accelerators** 

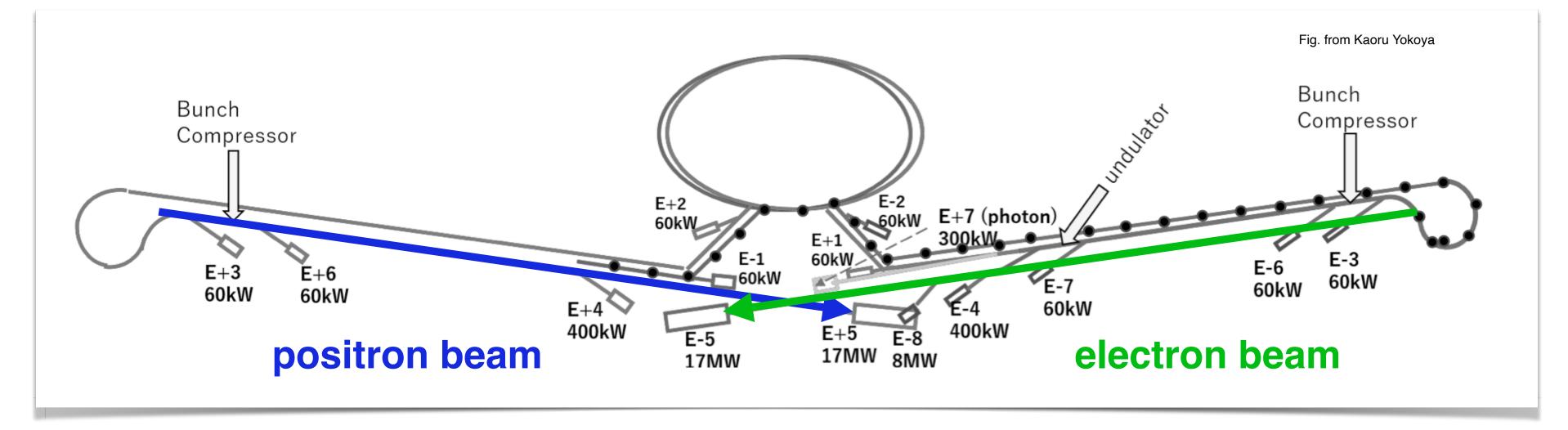
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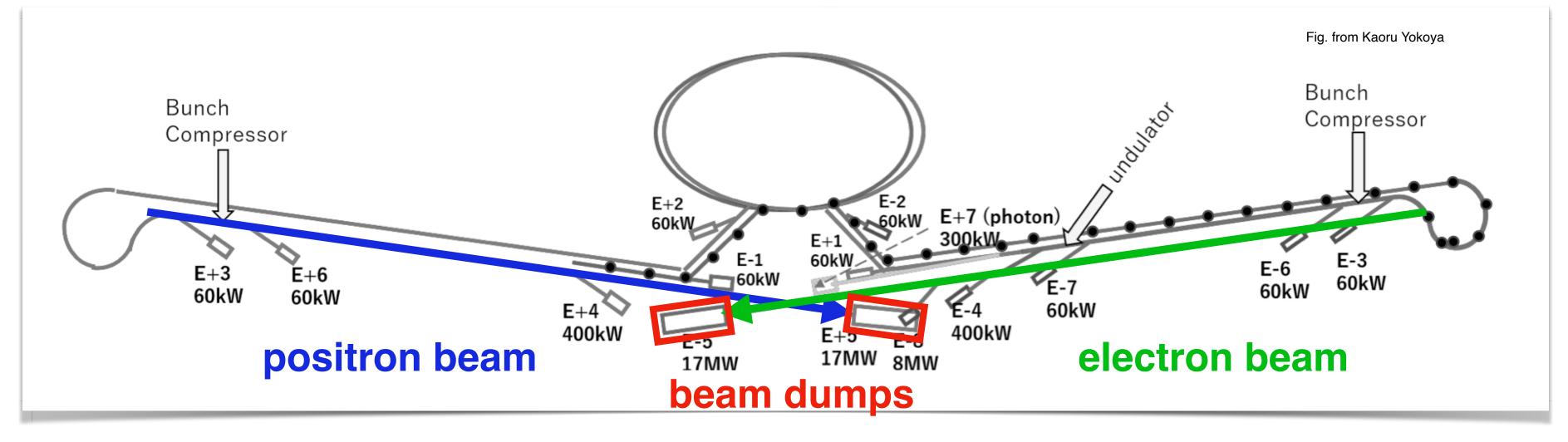
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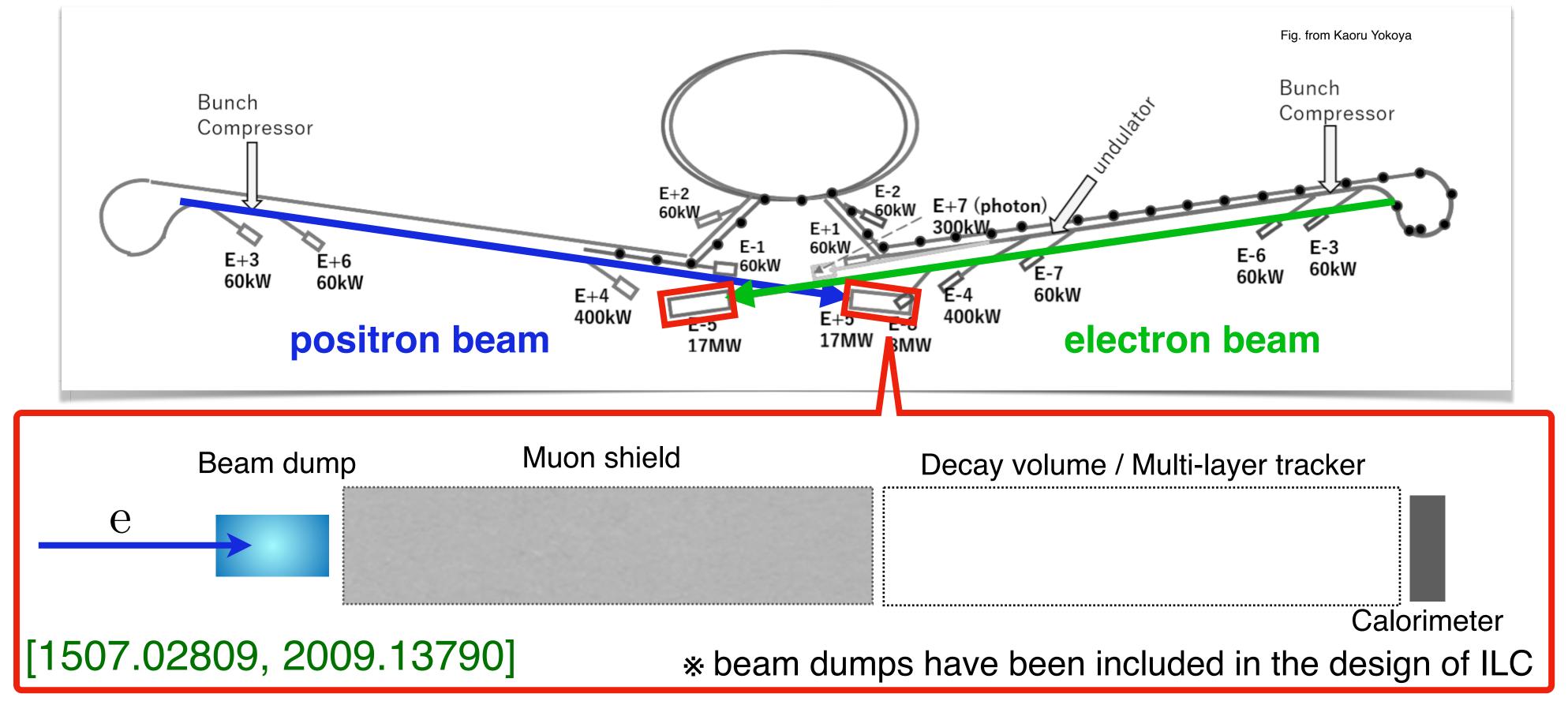


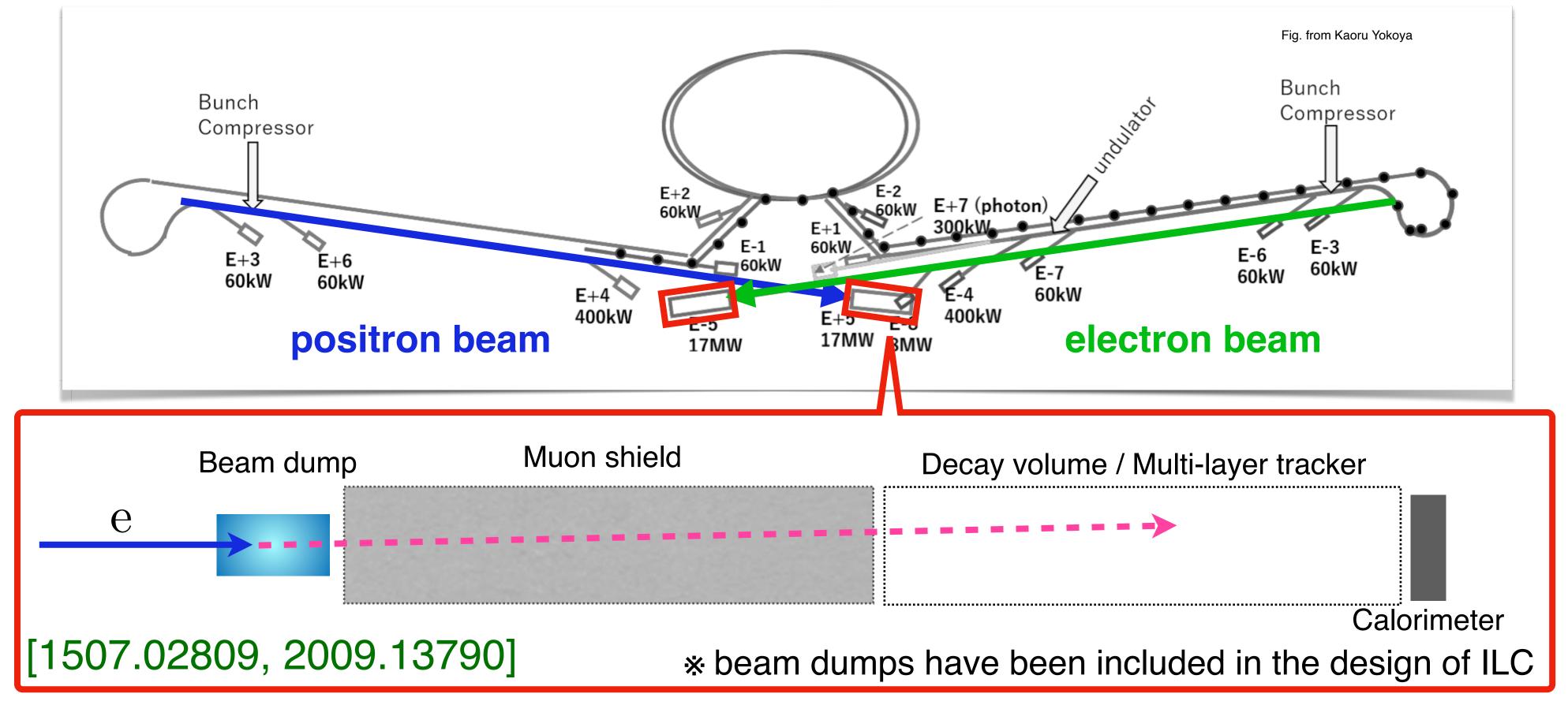


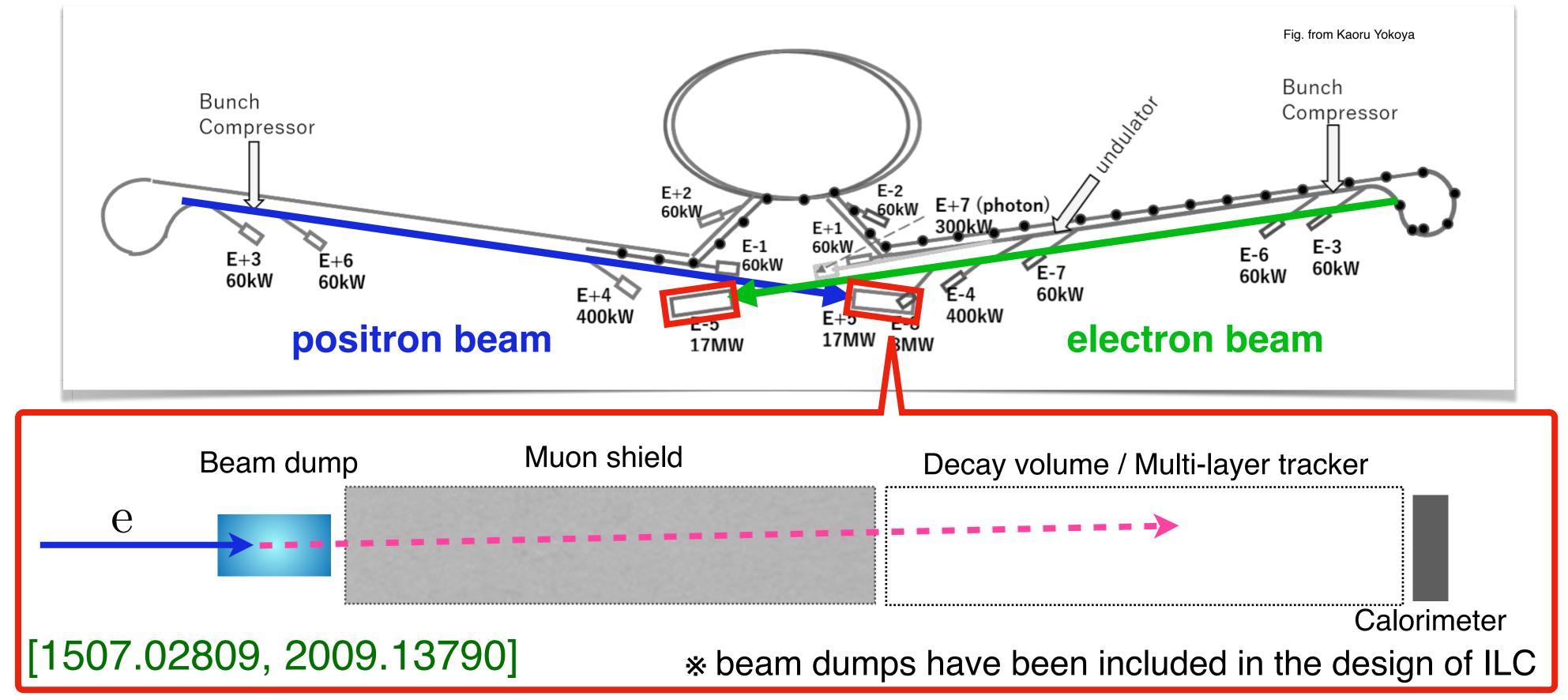








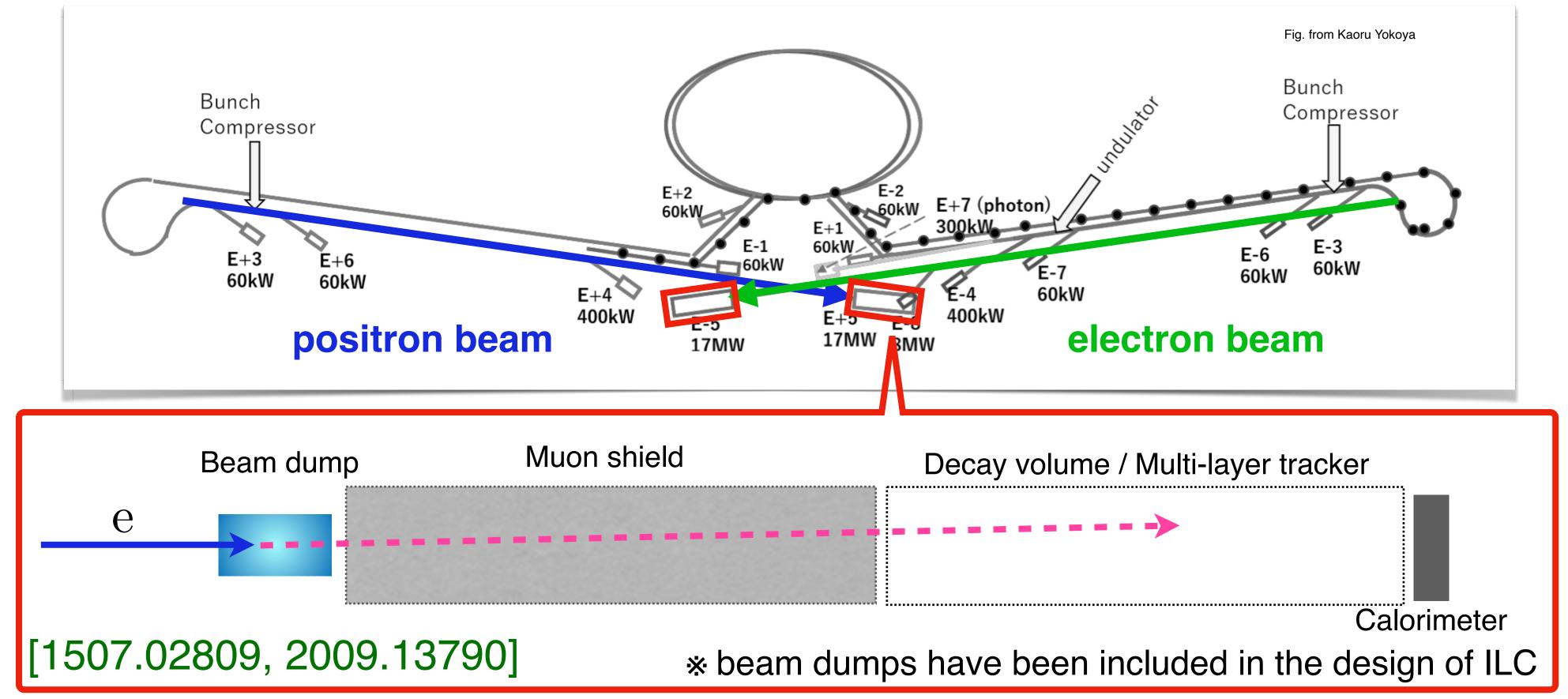




- space for muon shield is enough and beam dump facilities could be constructed if ILC is approved
- beam dump experiments in CLIC and C<sup>3</sup> would be similarly performed



### Ex. ILC-BDX $\Rightarrow$ parasitic running of International Linear Collider(ILC) experiment



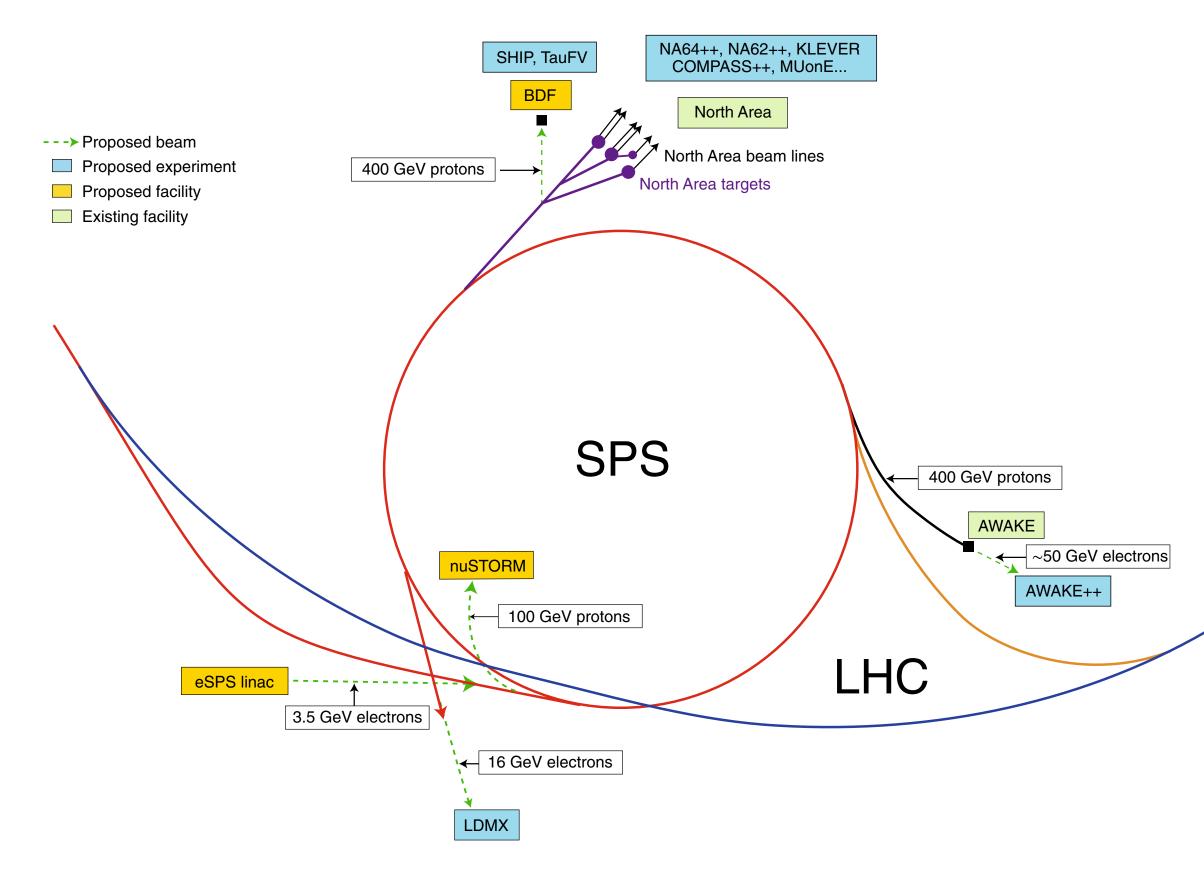
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Beam dump experiments will run in parallel with future accelerator facilities

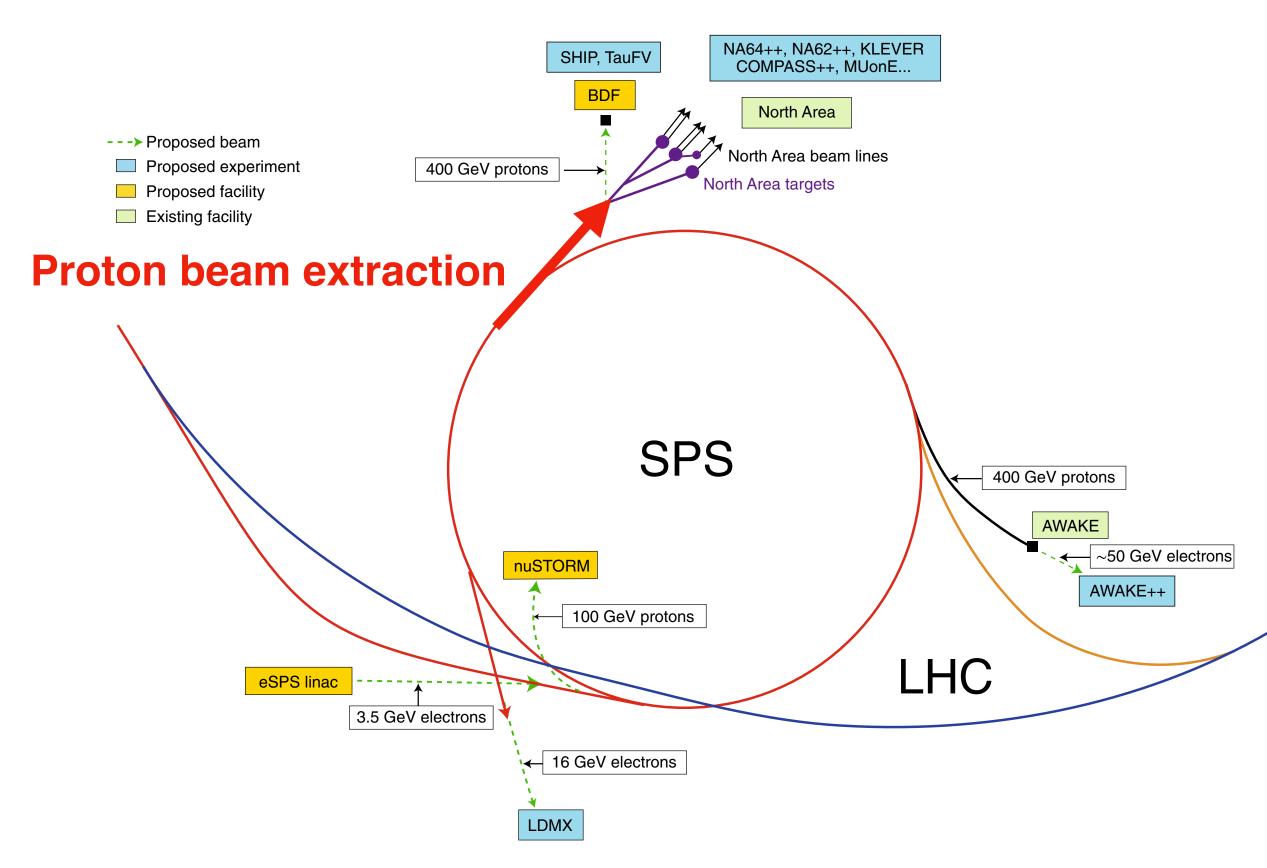


Ex. SHiP, LDMX, HIKE, SHADOWS  $\Rightarrow$  parasitic running of HL-LHC, etc

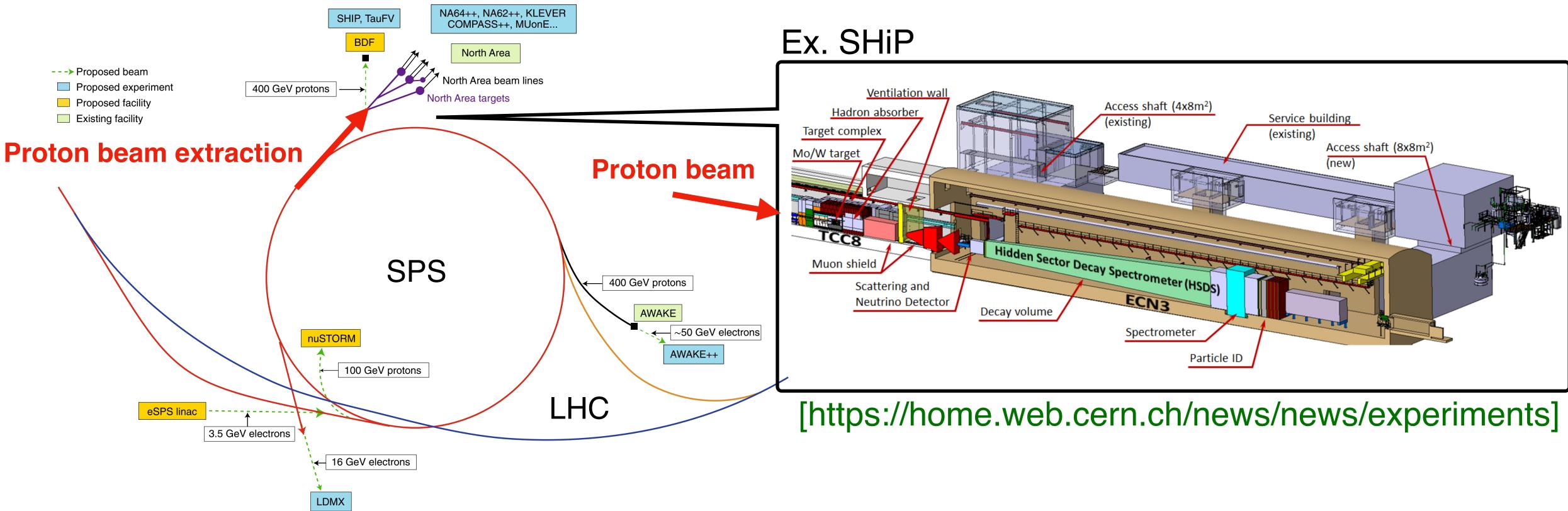
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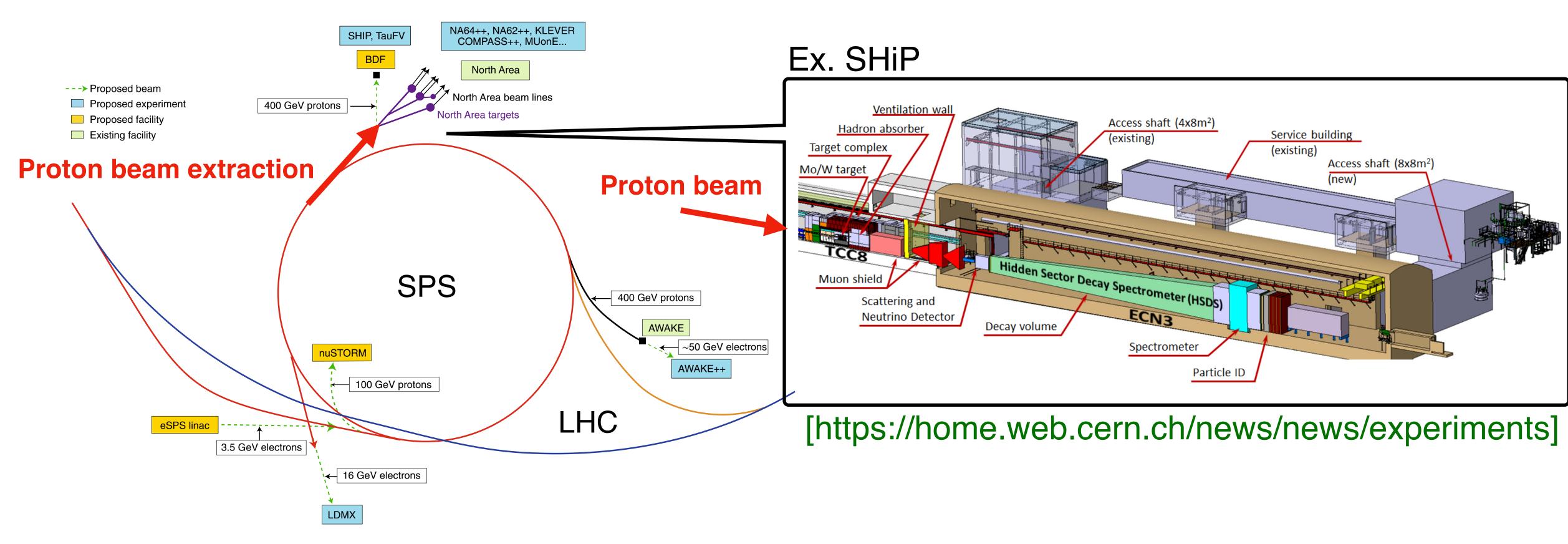
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- LDMX@eSPS, HIKE, SHADOWS, etc are proposed at CERN

Beam dump experiments will run in parallel with future accelerator facilities

- SHiP has been approved recently by CERN and will start to explore in 2031 in the North Area's ECN3 hall





- Examples of proposed beam dump experiments at future accelerator
  - SHiP, LDMX<sup>\*</sup>, HIKE, SHADOWS, etc [1901.09966, 2211.16586, etc]

 $\Rightarrow$  parasitic running of HL-LHC, etc

- ILC-BDX [1507.02809, 2009.13790, etc]
  - $\Rightarrow$  parasitic running of International Linear Collider(ILC) experiment

\* LDMX@SLAC received pre-project funds from Dark Matter New Initiatives and awaits construction funding

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- [2002.03005, 2203.08079] - DUNE, PIP2-BD, etc
  - $\Rightarrow$  parasitic running of accelerator facilities of future neutrino experiment

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  - $\Rightarrow$  parasitic running of muon collider experiments

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Beam dump experiments will run in parallel with future accelerator facilities

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# Outline

- Introduction
  - dark sector and Sub-GeV dark matter
  - beam dump(fixed-target) experiment
- Key features of beam dump experiment
- A classification of beam dump experiment
- Sensitivity of beam dump experiments at future accelerators
- Summary

• Past, current, and future beam dump experiments use the following detection methods:

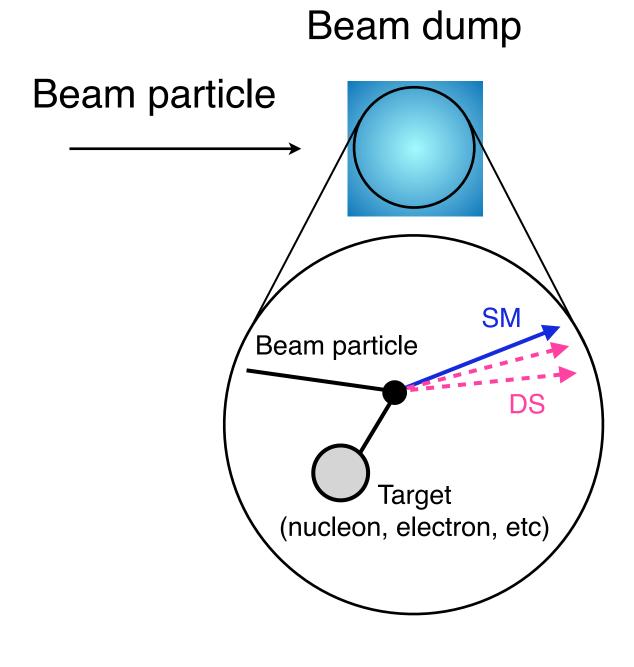
Beam dump

Beam particle

#### Detector



• Past, current, and future beam dump experiments use the following detection methods:

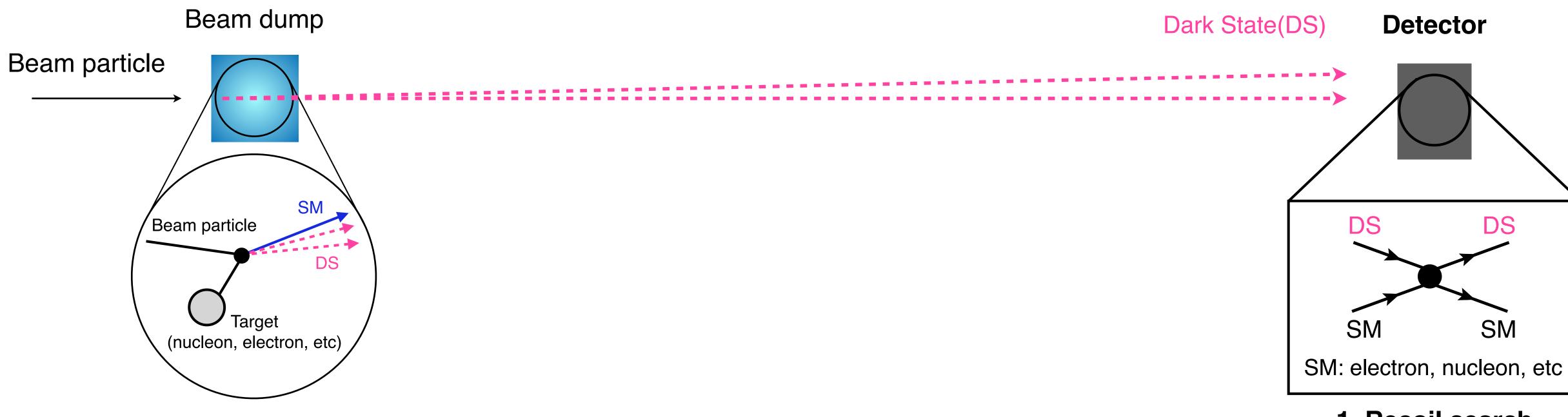


#### Detector



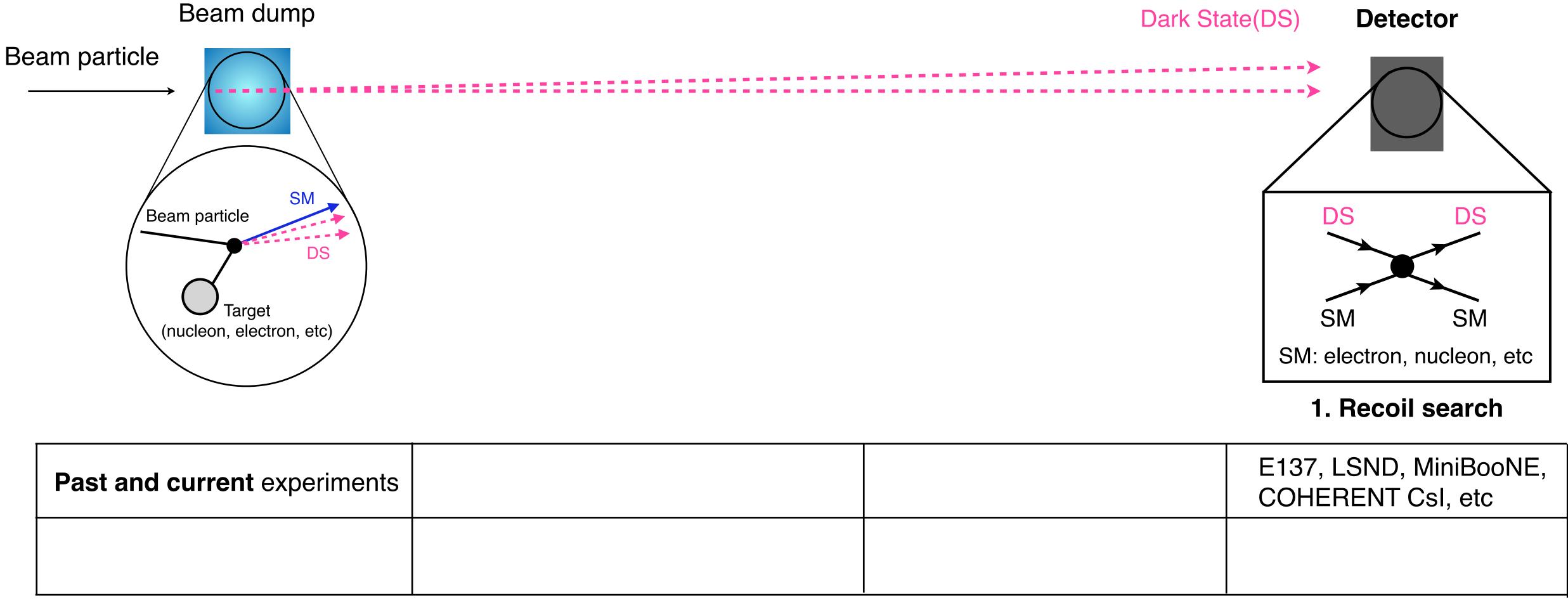


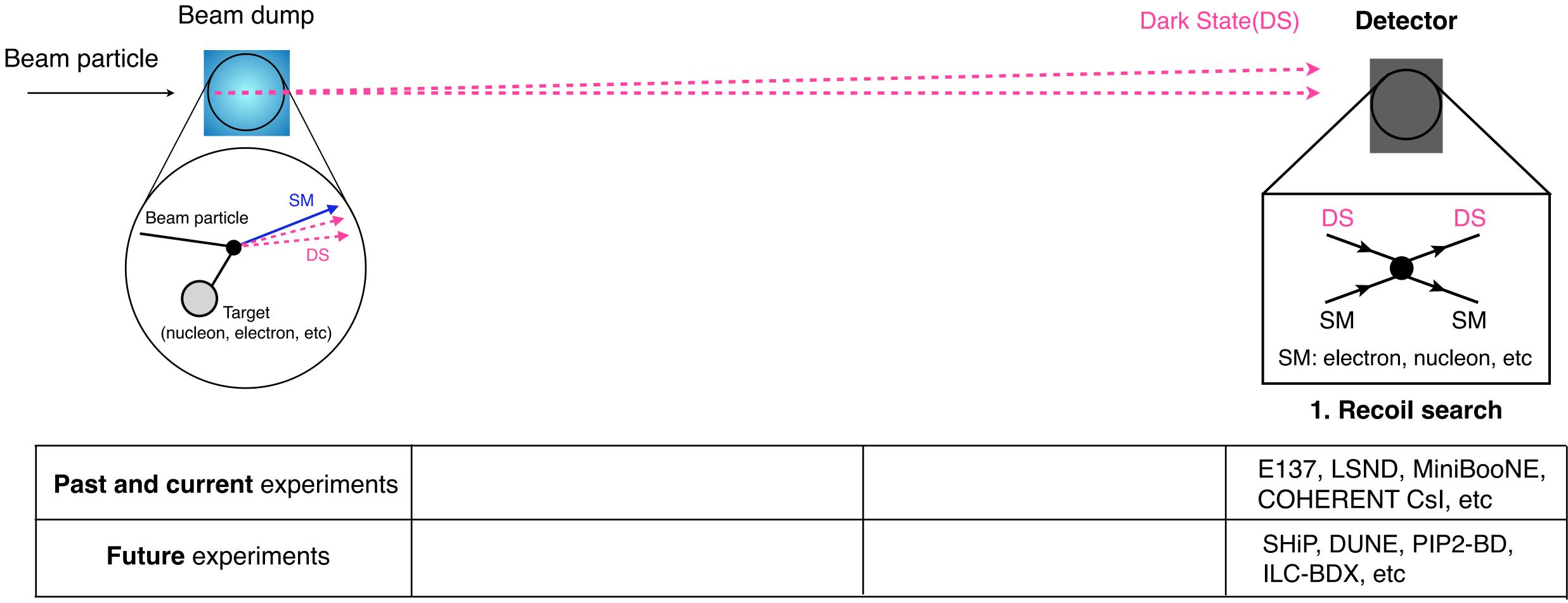
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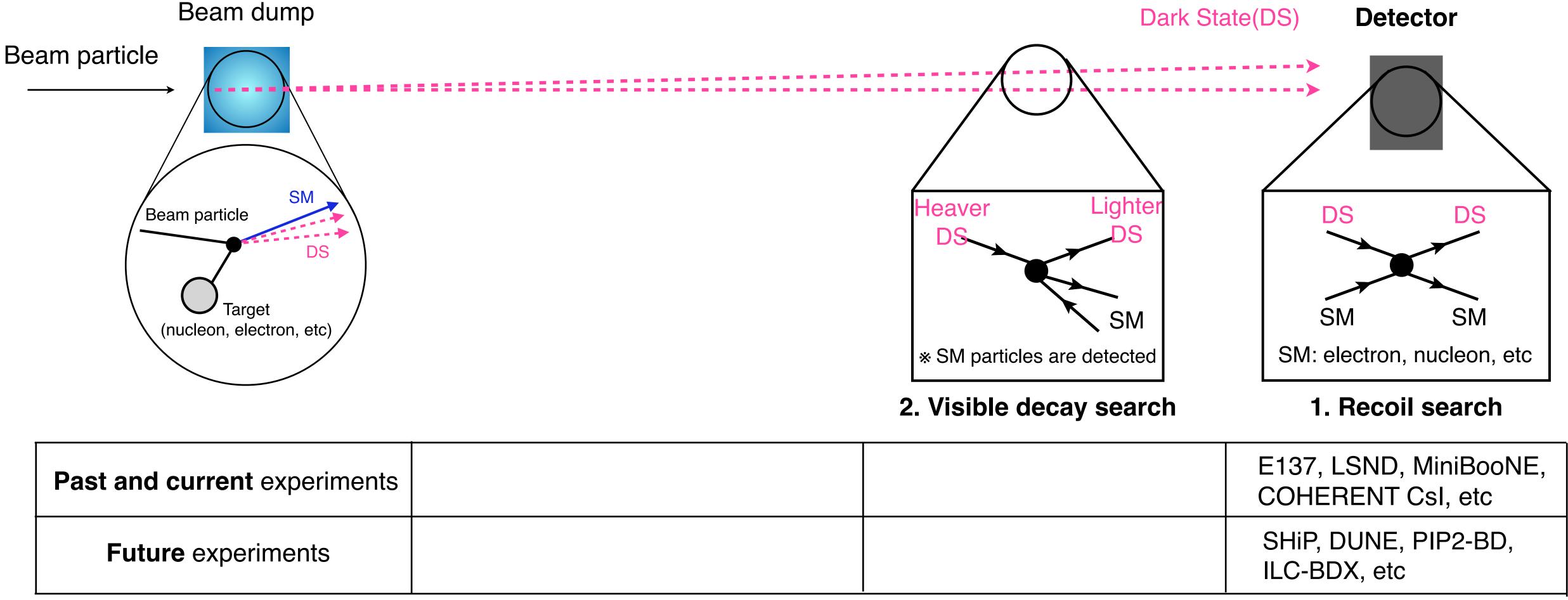
**1. Recoil search** 



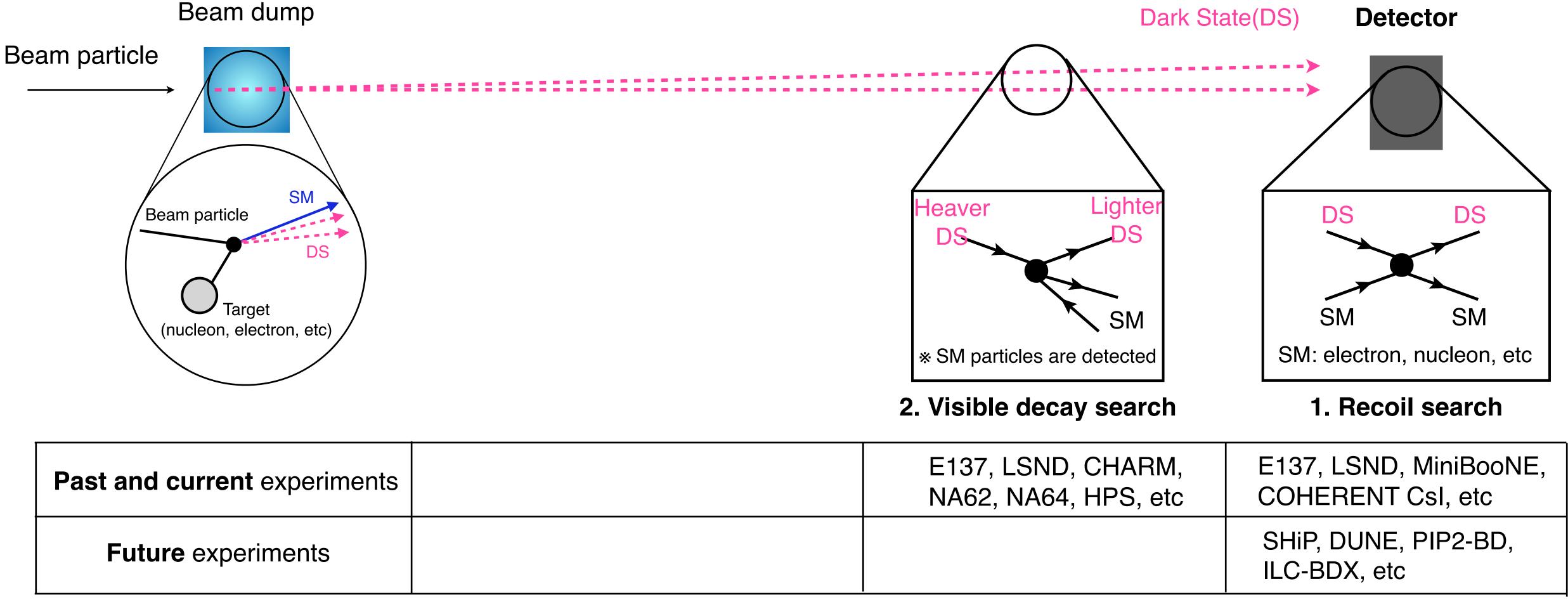




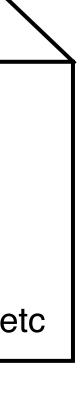
Past and current experiments	
Future experiments	

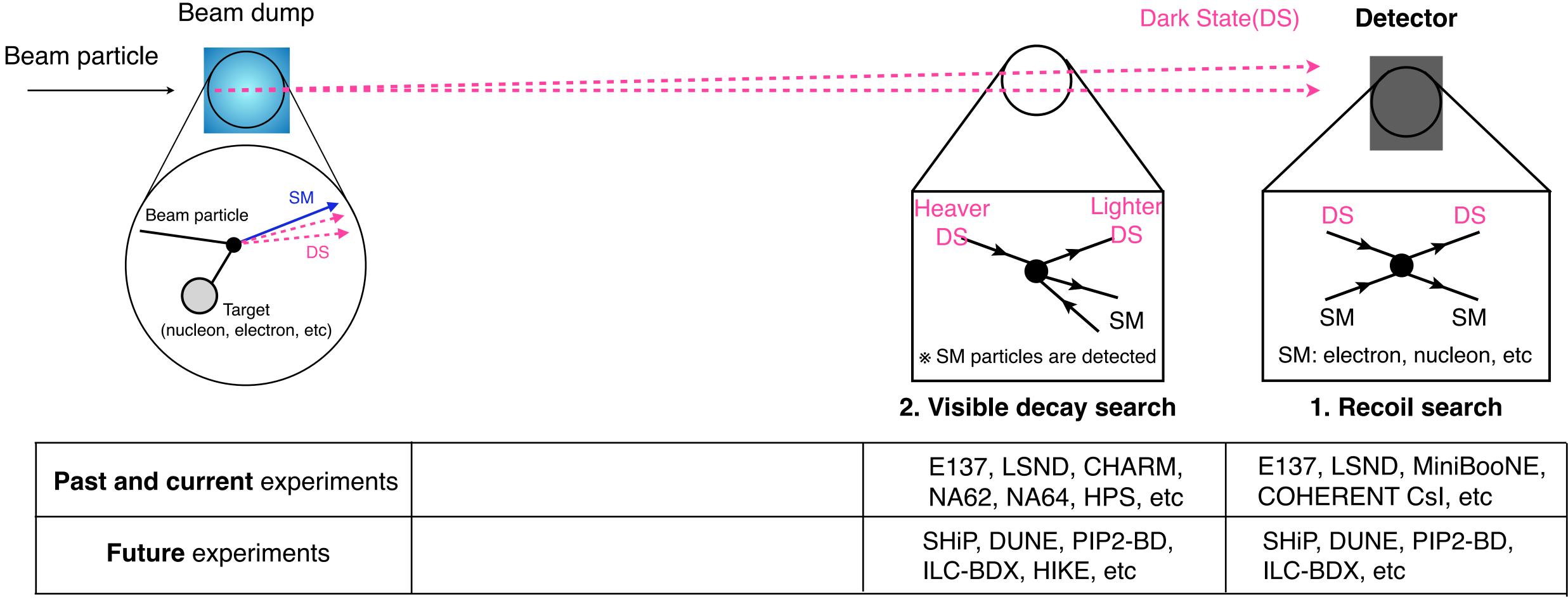


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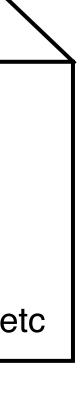


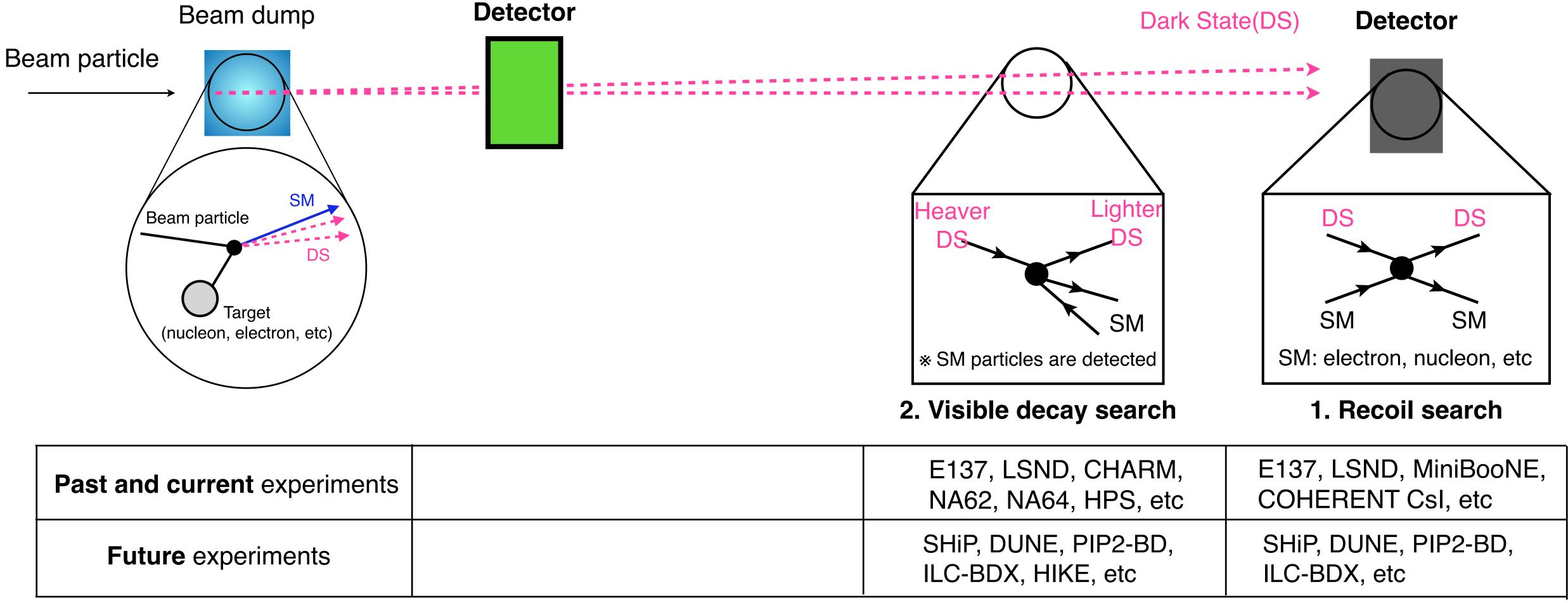
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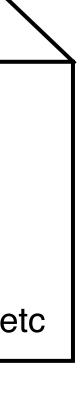


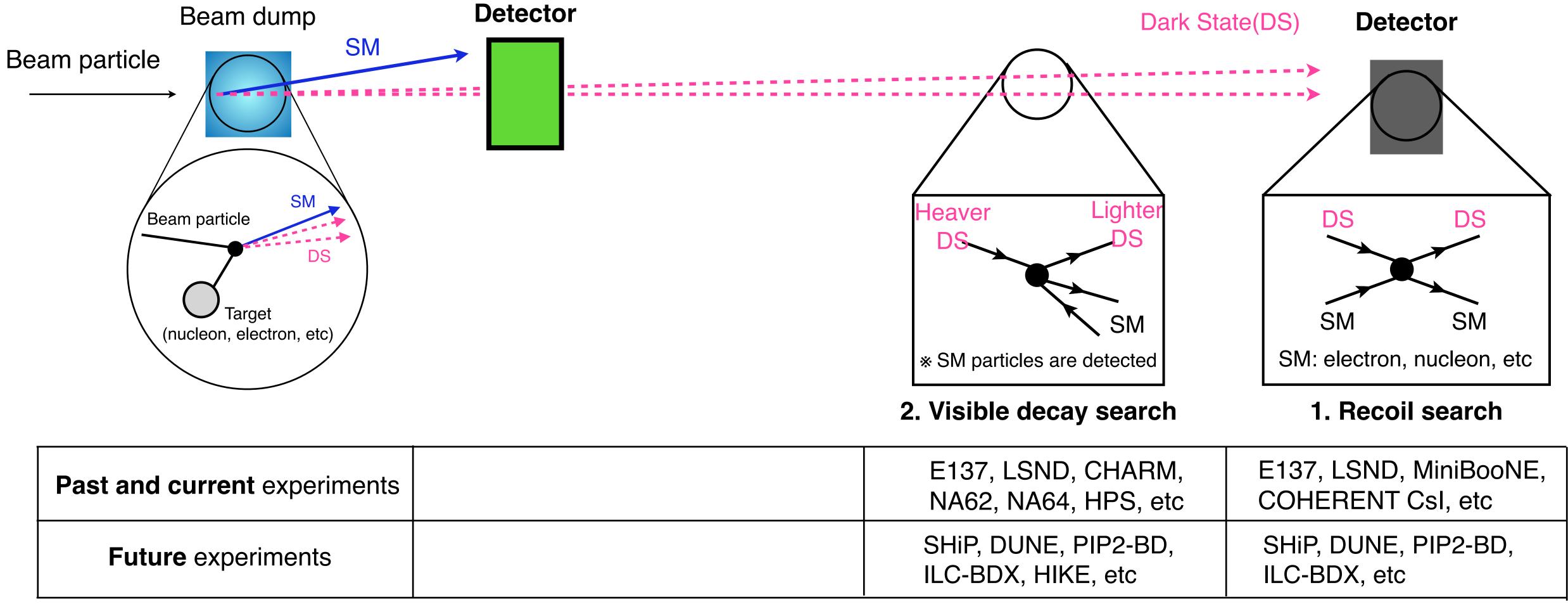
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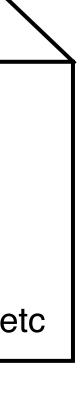


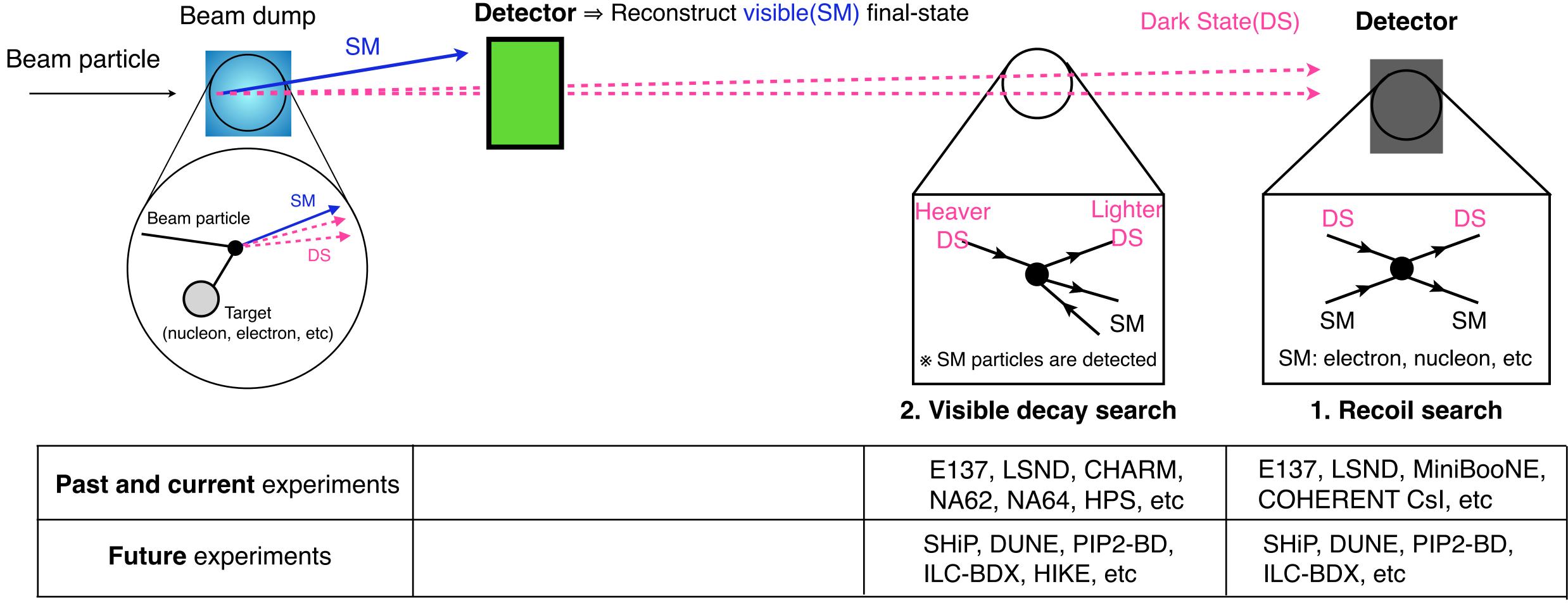
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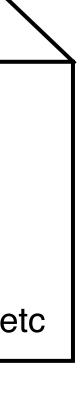


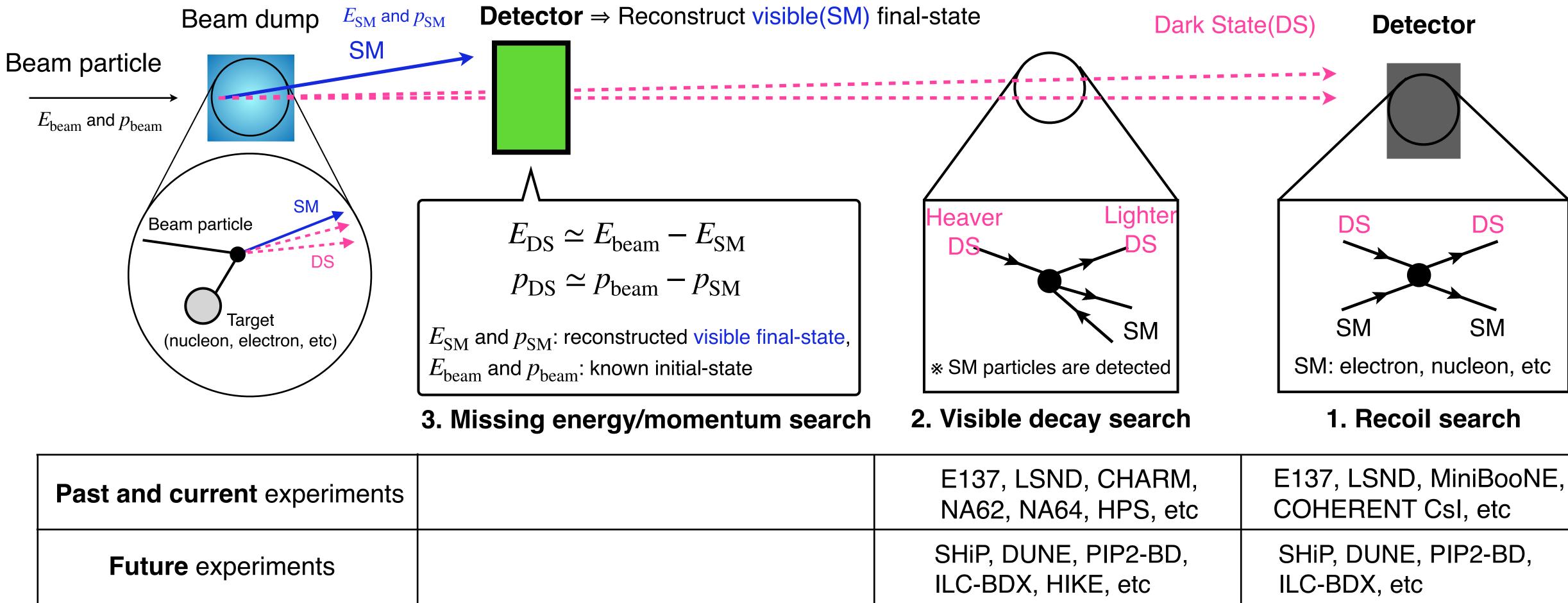
Past and current experiments	
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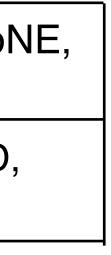
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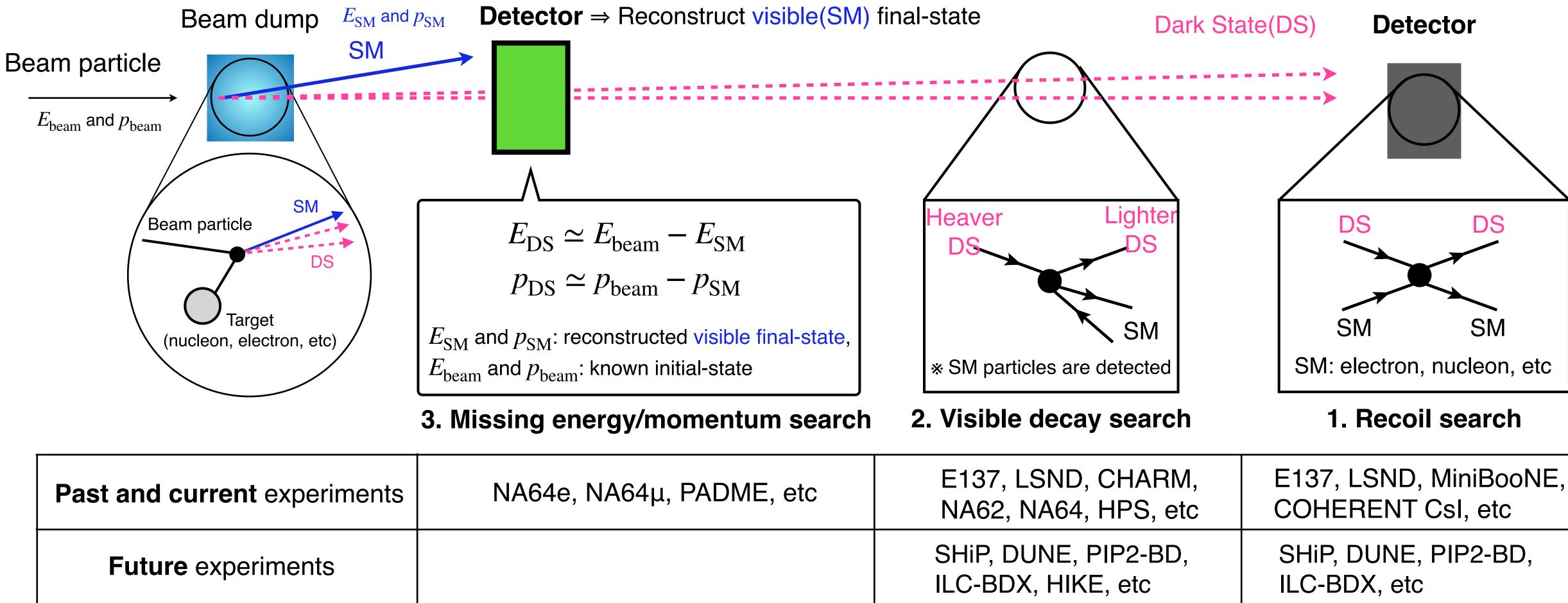




Past and current experiments	
Future experiments	

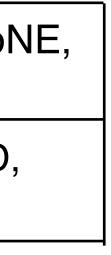


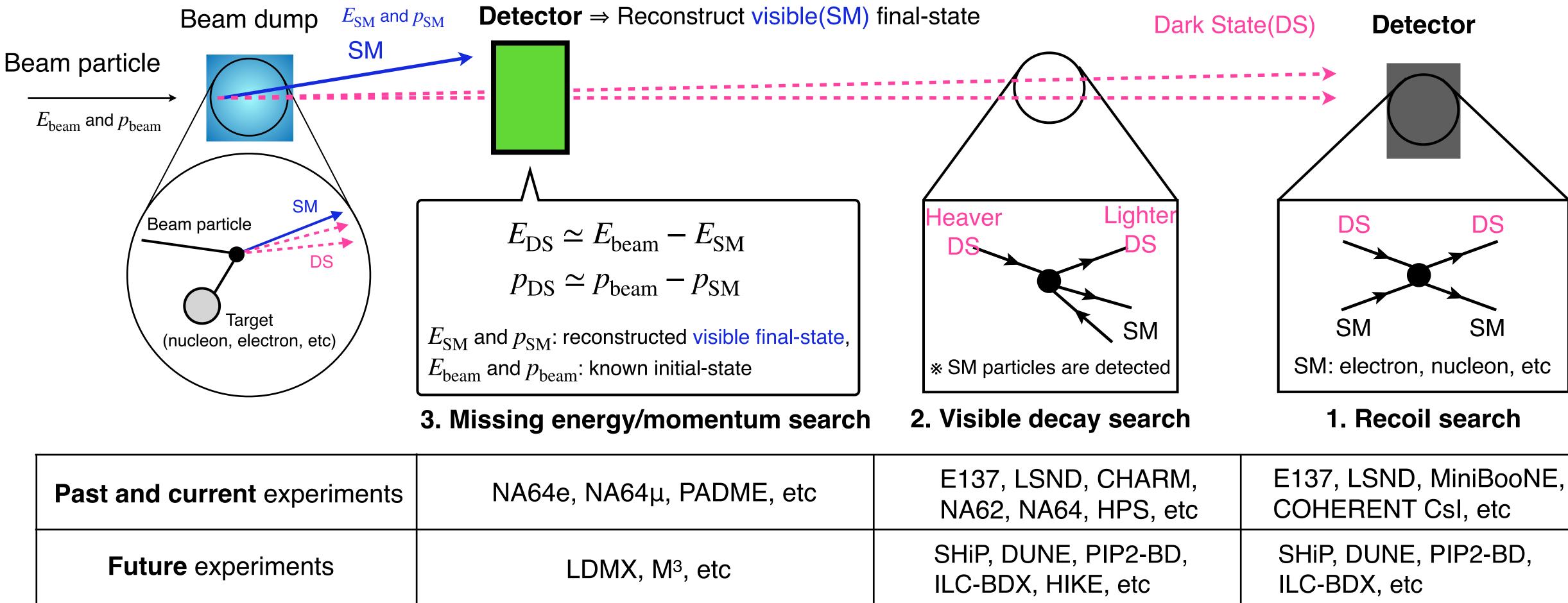




Past and current experiments	NA64e, NA64µ, PADN
Future experiments	

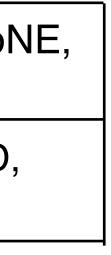




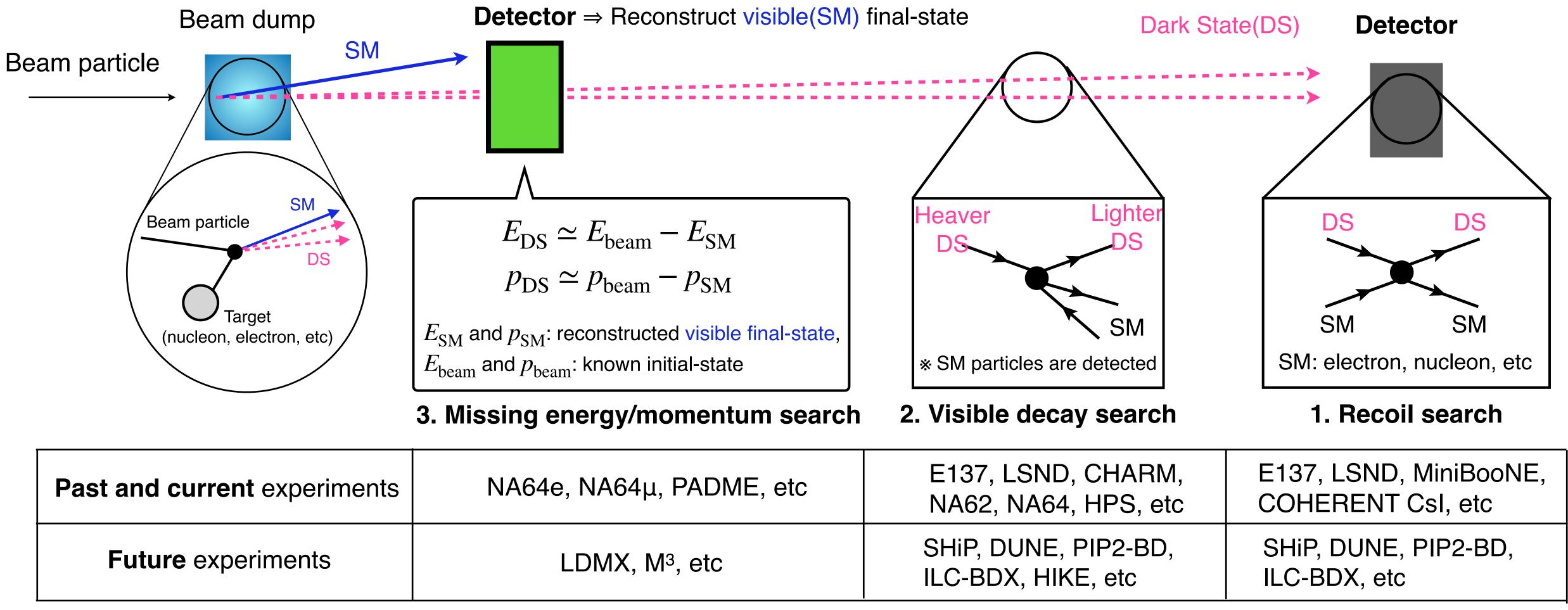


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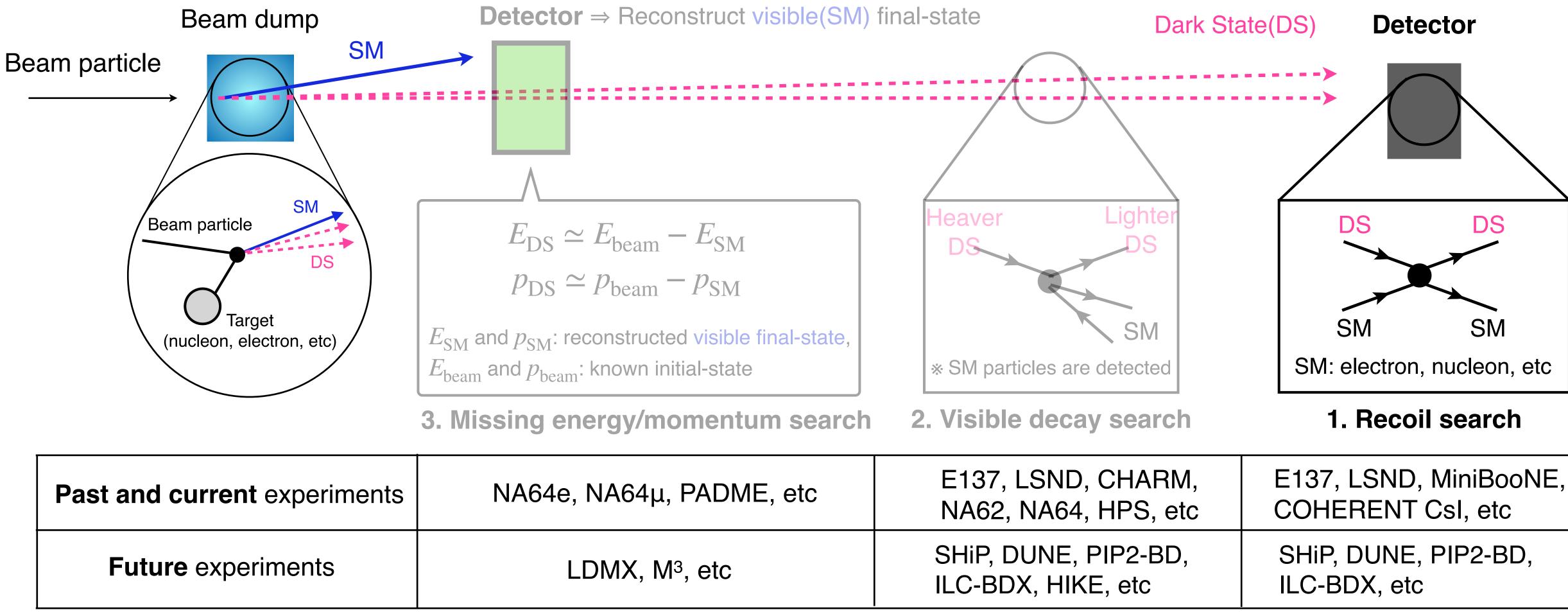
Past and current experiments	NA64e, NA64µ, PADME
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#### Beam dump experiments are divided into three detection approaches



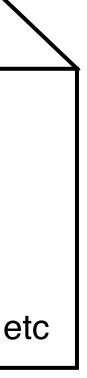


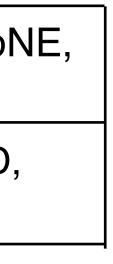
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Past and current experiments	NA64e, NA64µ, PADME
Future experiments	LDMX, M <sup>3</sup> , etc

#### Beam dump experiments are divided into three detection approaches







#### Typical setup:

Beam



#### Typical setup:

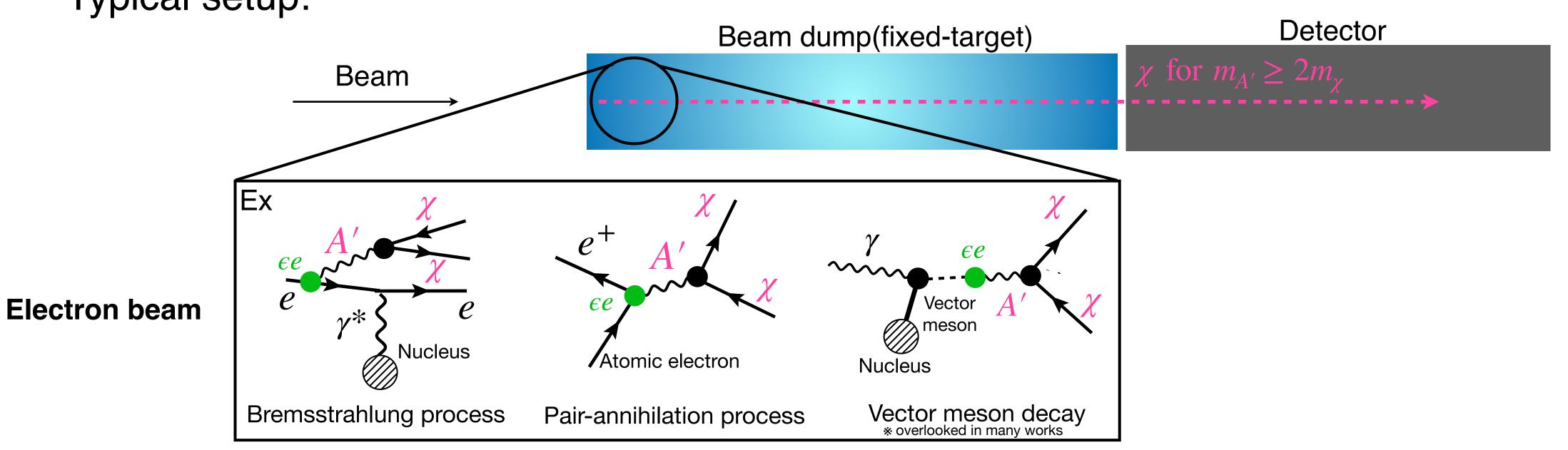
Beam

Benchmark model:  $\mathscr{L} \supset \epsilon e A'_{\mu} J^{\mu}_{EM} - g_D A'_{\mu} \bar{\chi} \gamma^{\mu} \chi A'$ : Dark photon,  $\chi$ : Dark Matter,  $J^{\mu}_{EM}$ : SM EM current



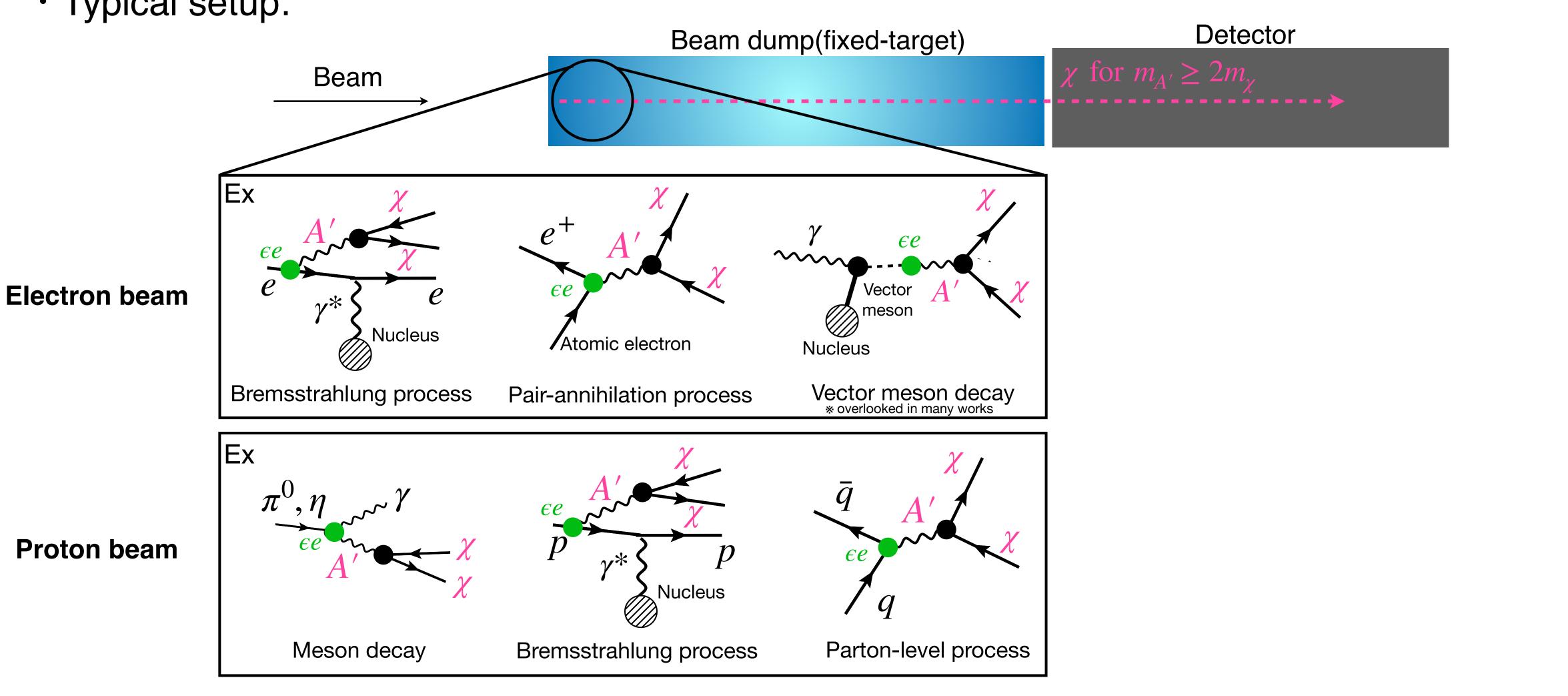
# rrent

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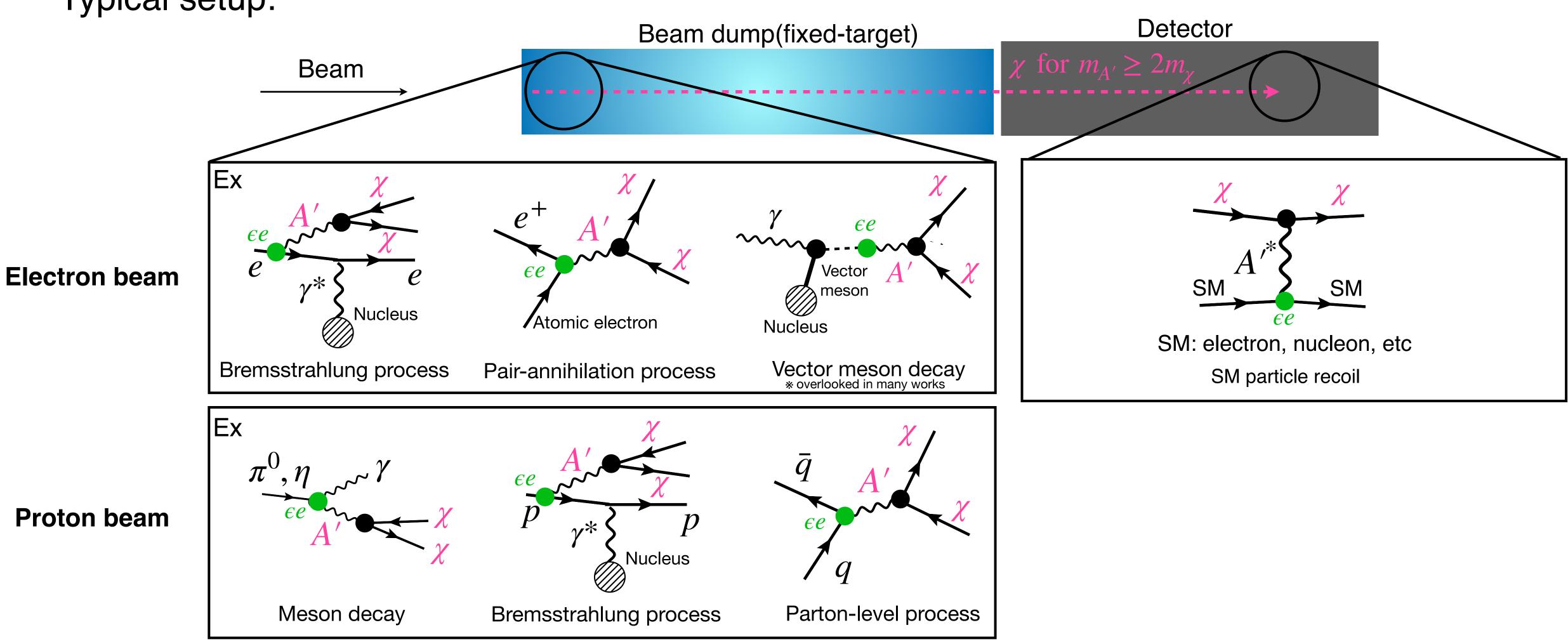
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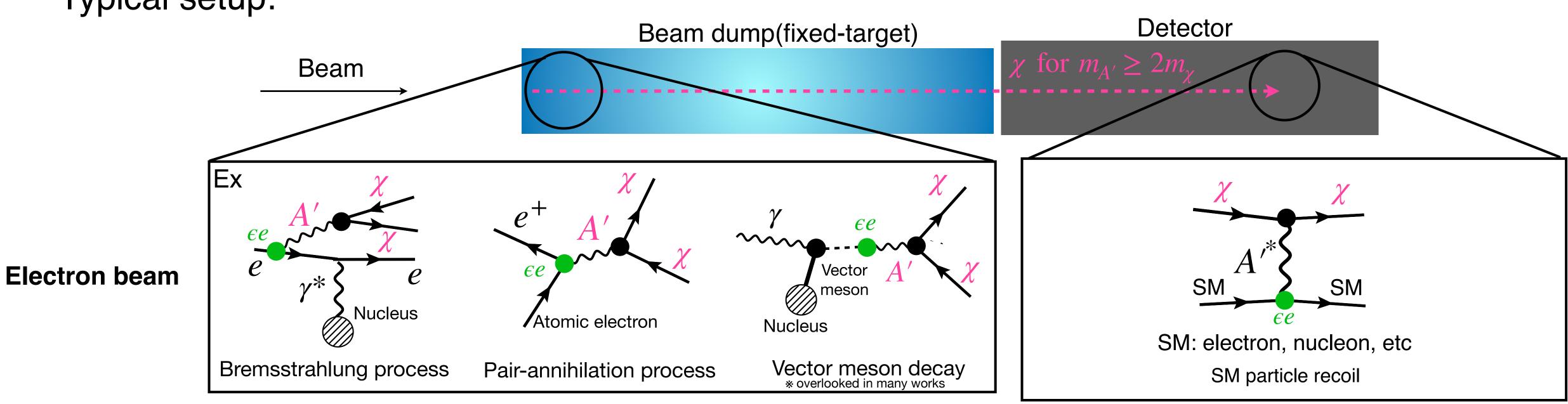
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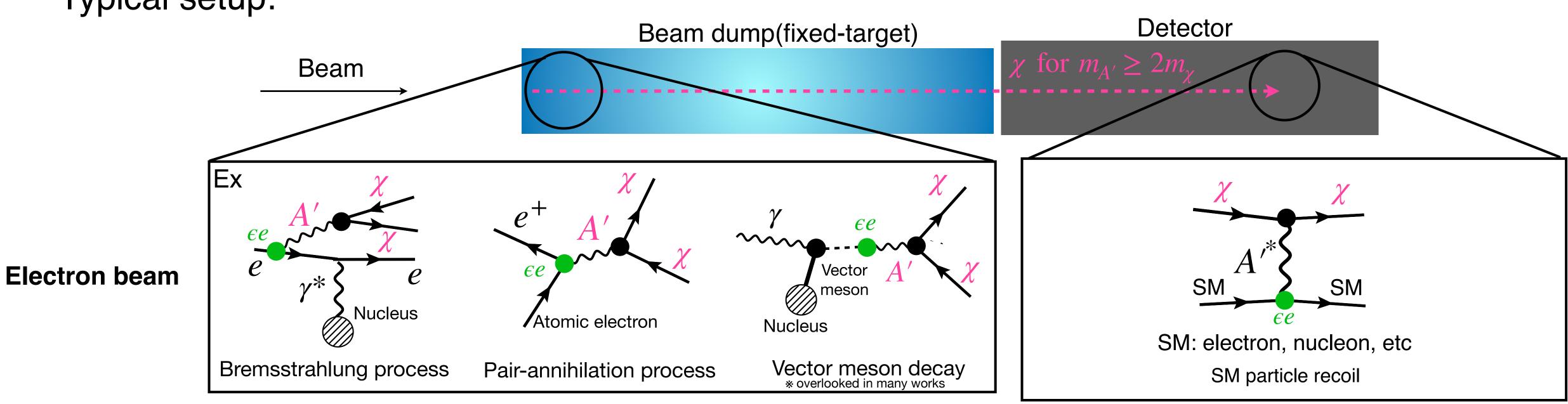
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Typical setup:



# of detected DS signature (signal events):

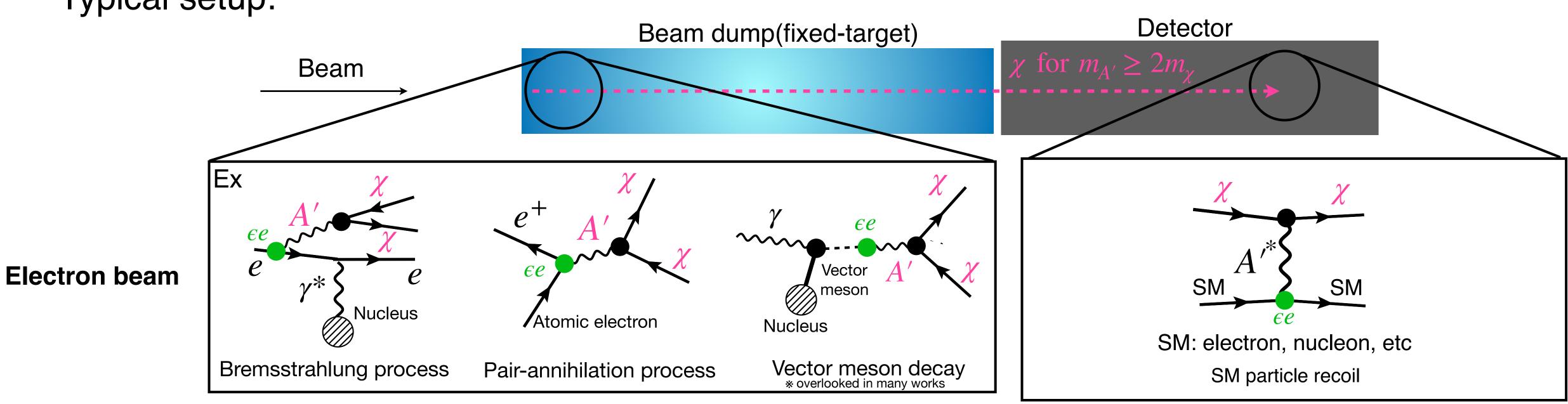
Typical setup:



# of detected DS signature (signal events):

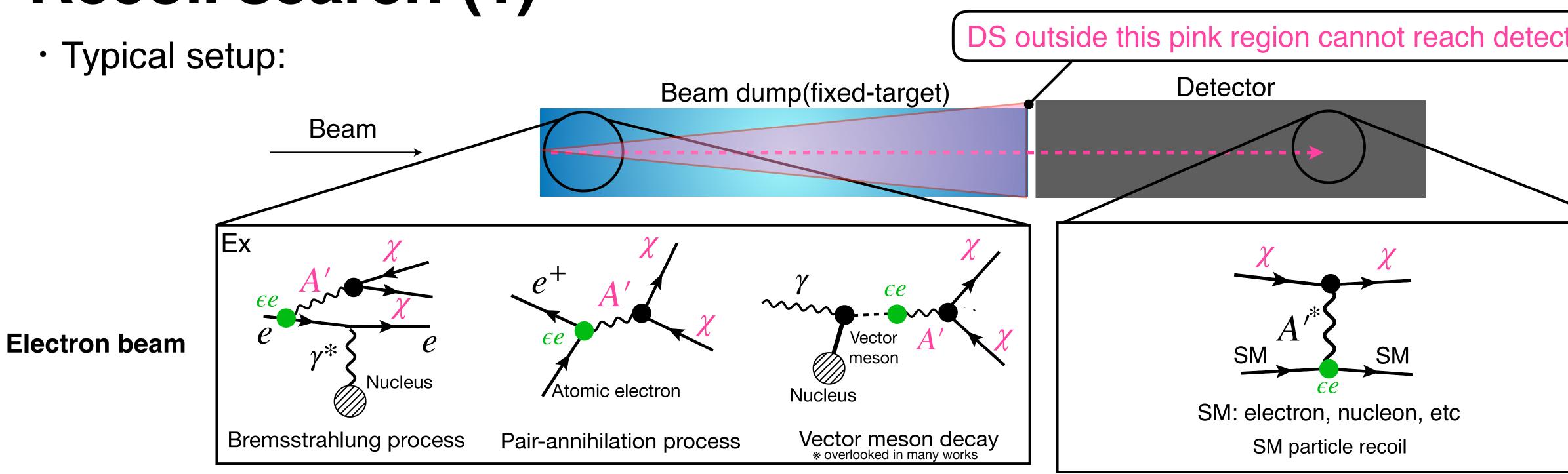
~ (# of produced DS) ×

• Typical setup:



# of detected DS signature (signal events):

 $\sim$  (# of produced DS)  $\times$  (Probability DS reaches detector)  $\times$ 

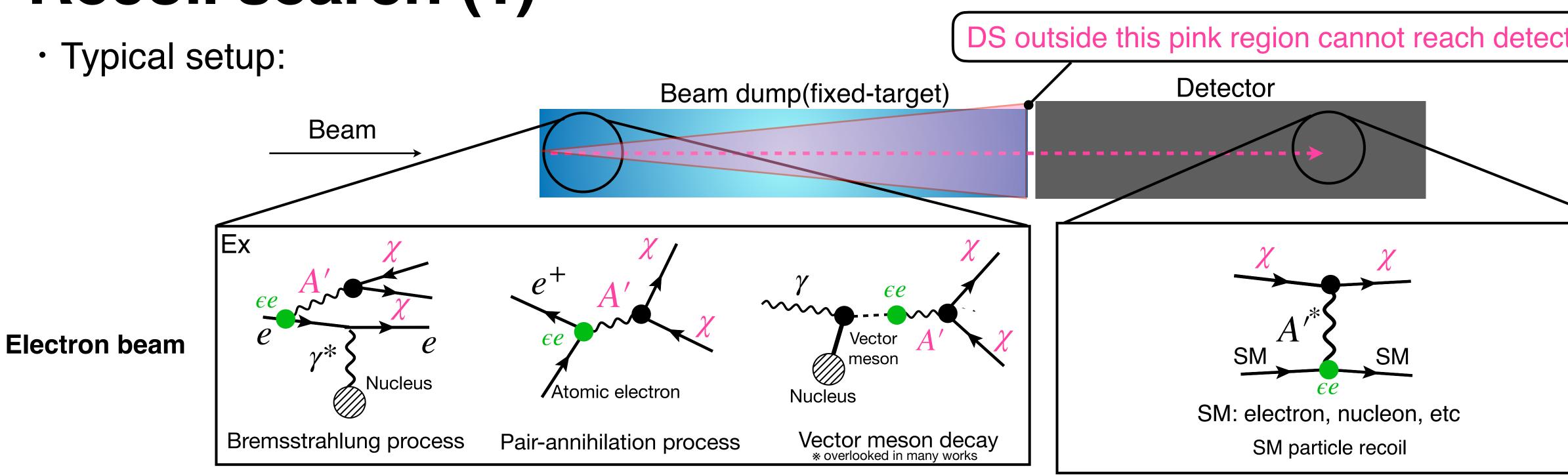


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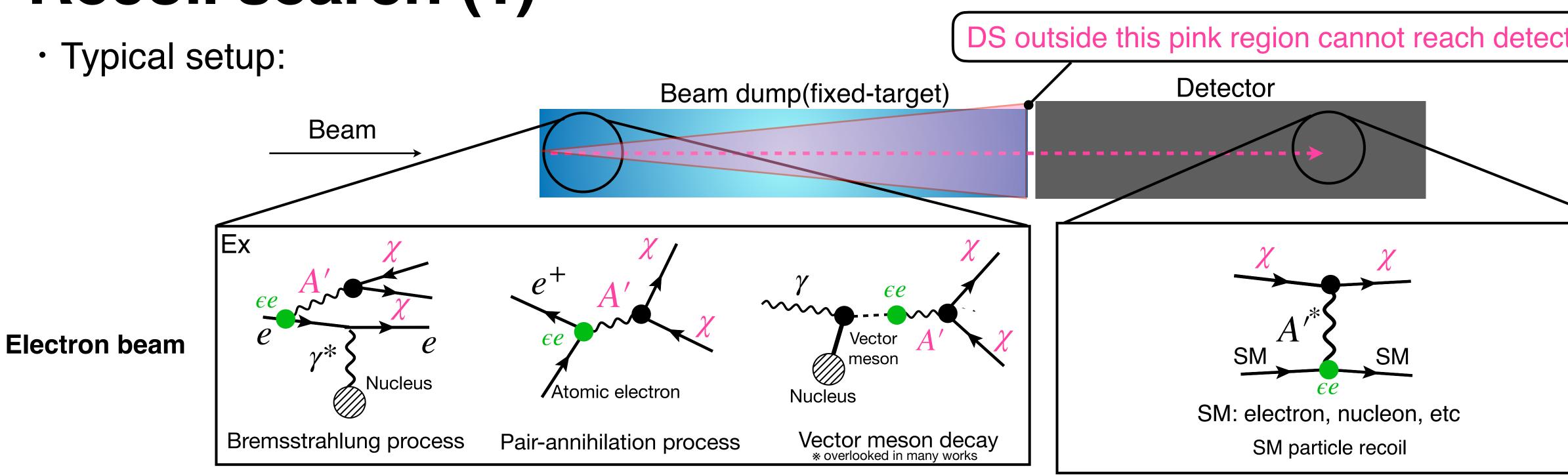
to	r

	_



# of detected DS signature (signal events):

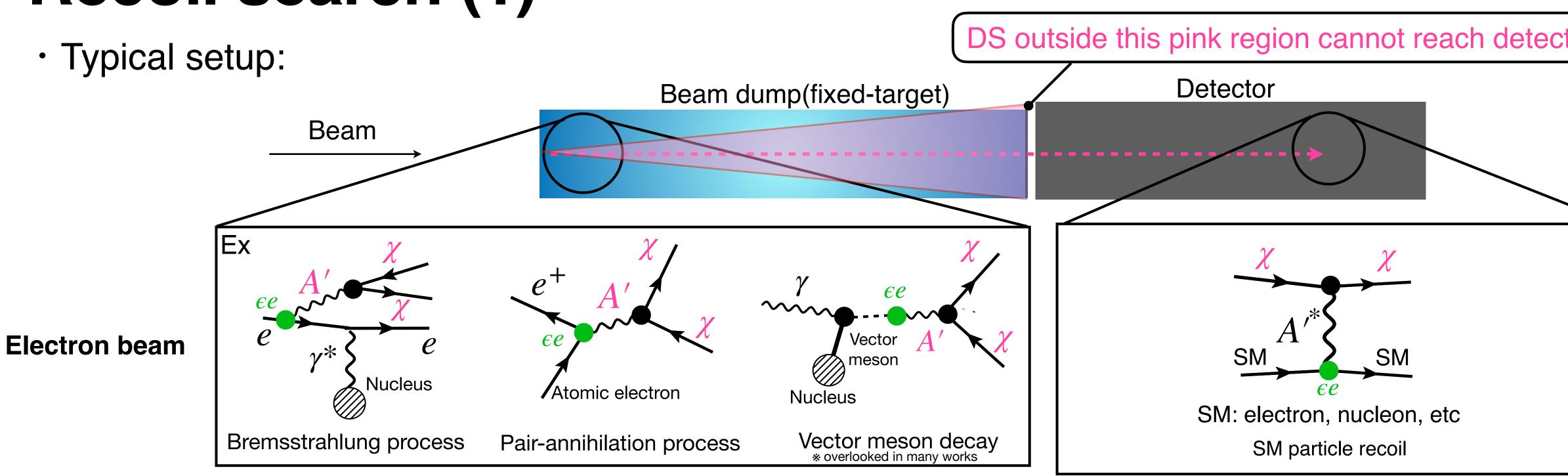
to	r



# of detected DS signature (signal events):

 $\propto$  (Beam flux)×( $\epsilon e$ )<sup>2</sup>

to	r

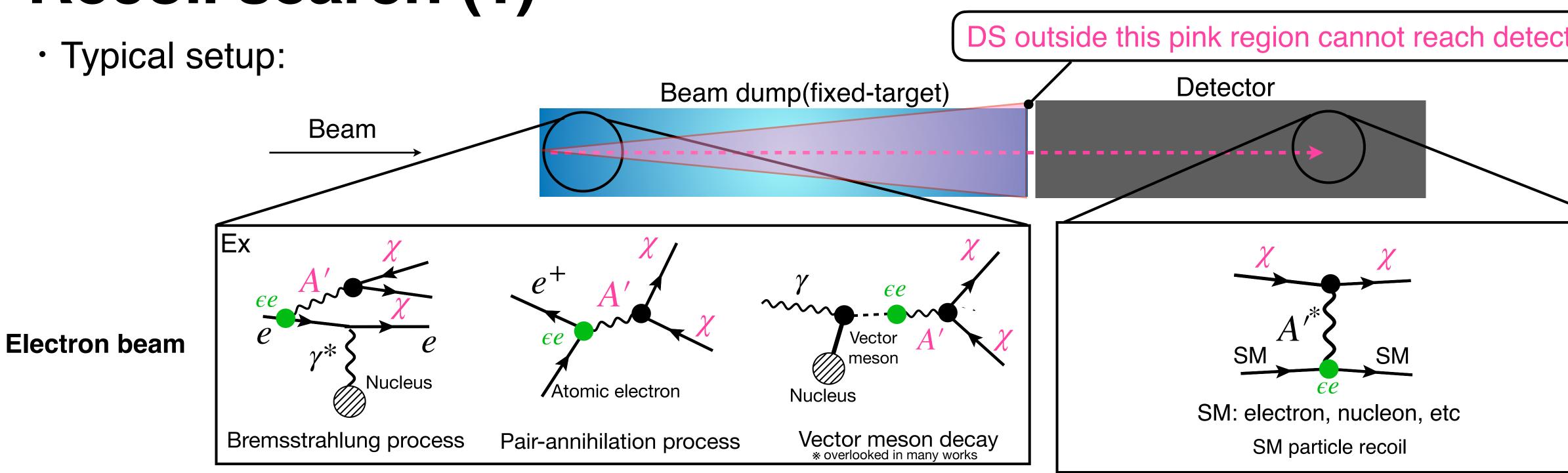


# of detected DS signature (signal events):

 $\propto$  (Beam flux)×( $\epsilon e$ )<sup>2</sup>

 $\propto$  (height of detector)<sup>2</sup>  $\times$  (length b/w beam dump and detector)<sup>-2</sup>

to	r



# of detected DS signature (signal events):

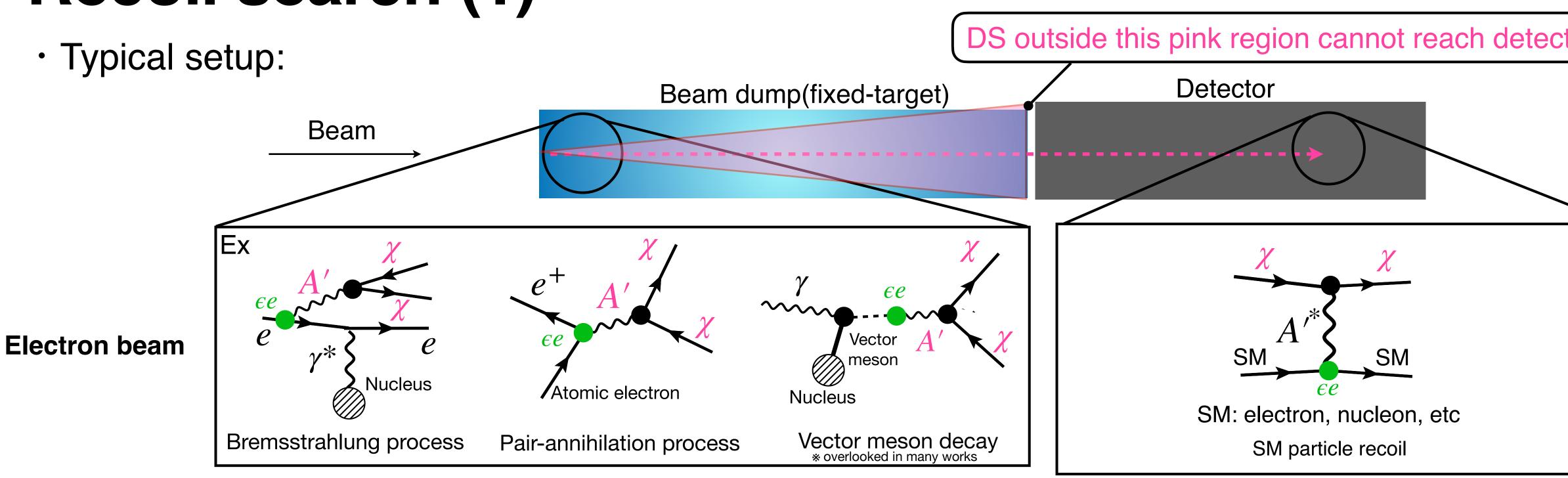
 $\propto$  (Beam flux)×( $\epsilon e$ )<sup>2</sup>

 $\propto$  (height of detector)<sup>2</sup>  $\times$  (length b/w beam dump and detector)<sup>-2</sup>

### $\sim$ (# of produced DS) × (Probability DS reaches detector) × (Probability DS is detected)

 $\propto$  (Length of detector)×( $\epsilon e$ )<sup>2</sup>

to	r



# of detected DS signature (signal events):

~ (# of produced DS) ×

 $\propto$  (Beam flux)×( $\epsilon e$ )<sup>2</sup>

(height of detector)<sup>2</sup>  $\propto$  $\times$  (length b/w beam dump and detector)<sup>-2</sup>

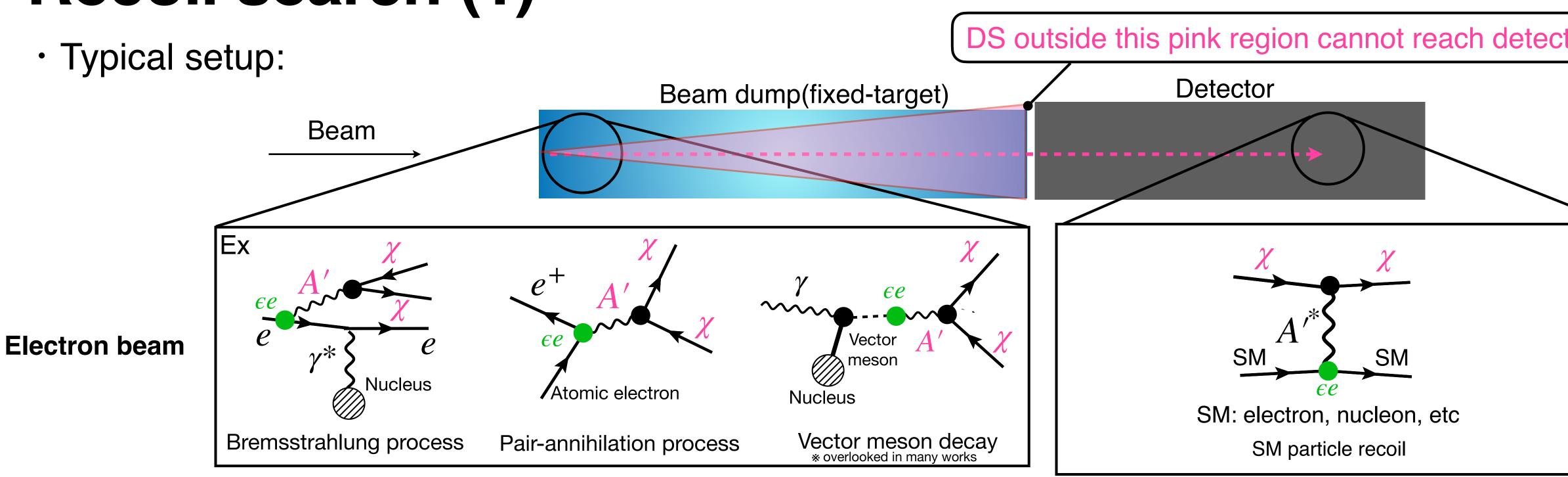
#### Acceptance

### (Probability DS reaches detector) $\times$ (Probability DS is detected)

 $\propto$  (Length of detector)×( $\epsilon e$ )<sup>2</sup>

to	r





# of detected DS signature (signal events):

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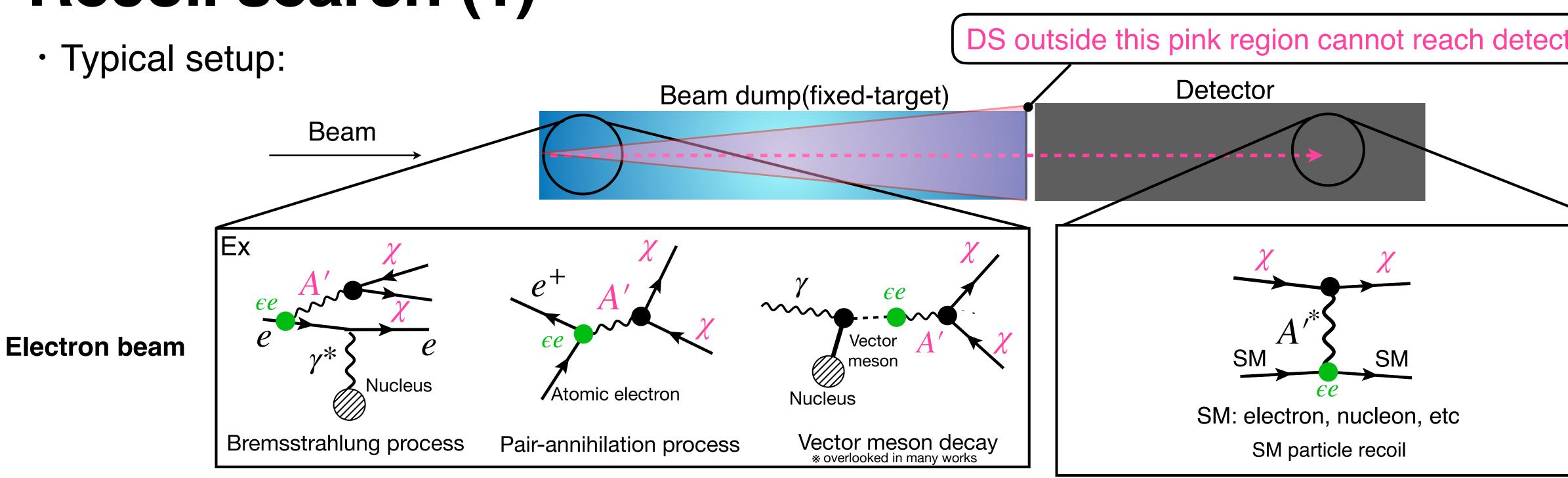
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### High flux beam, near\* and large detectors are suited for recoil search

to	r





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 $\propto$  (Beam flux)×( $\epsilon e$ )<sup>2</sup>

(height of detector)<sup>2</sup>  $\propto$  $\times$  (length b/w beam dump and detector)<sup>-2</sup>

\* Detector cannot be too near because large beam dump or shield is needed to reduce beam-induced BG

#### Acceptance

### (Probability DS reaches detector) $\times$ (Probability DS is detected)

 $\propto$  (Length of detector)×( $\epsilon e$ )<sup>2</sup>

### High flux beam, near\* and large detectors are suited for recoil search

to	r)

		ł







### • Typical setup:

Beam



#### Detector/Decay volume

### • Typical setup:

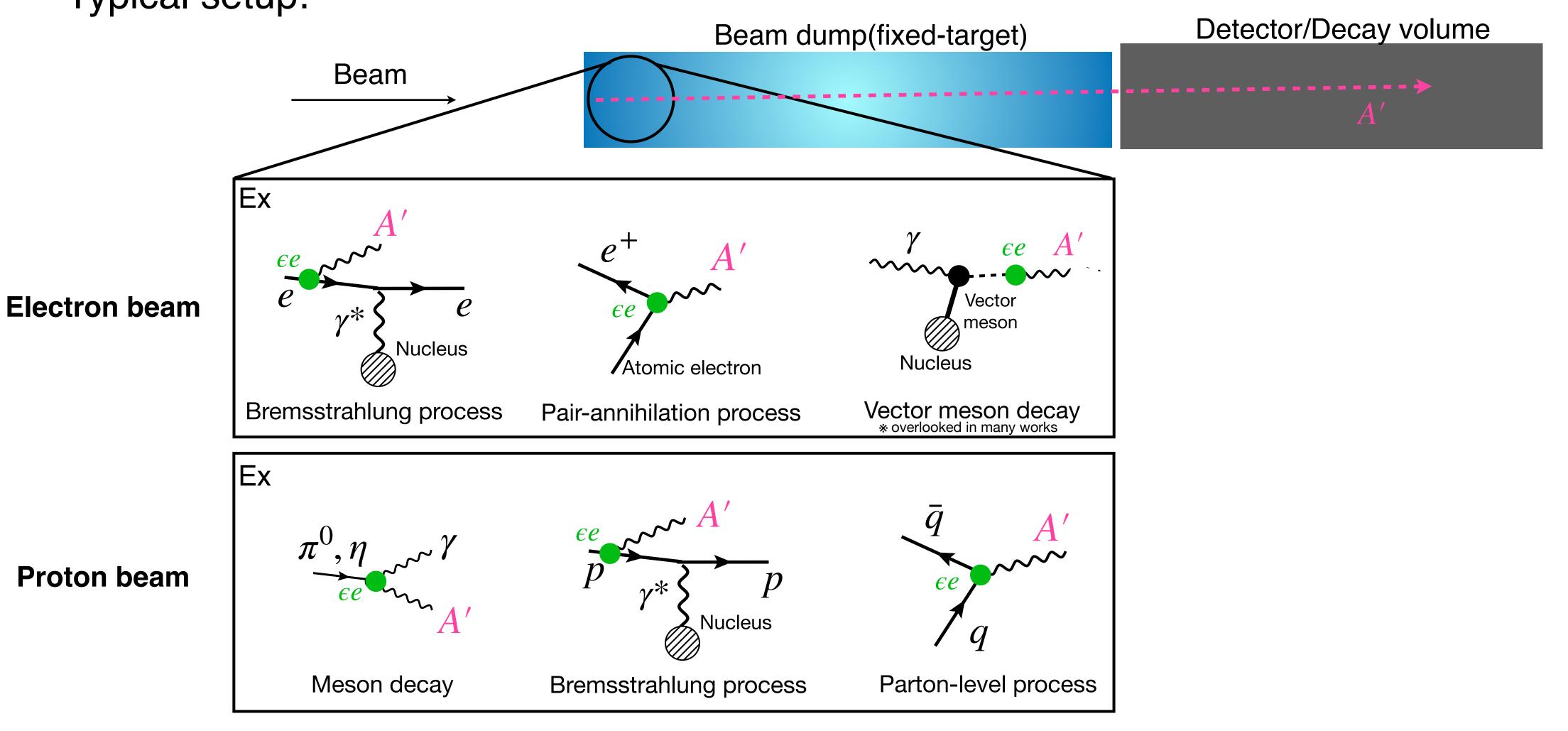
Beam



#### Detector/Decay volume

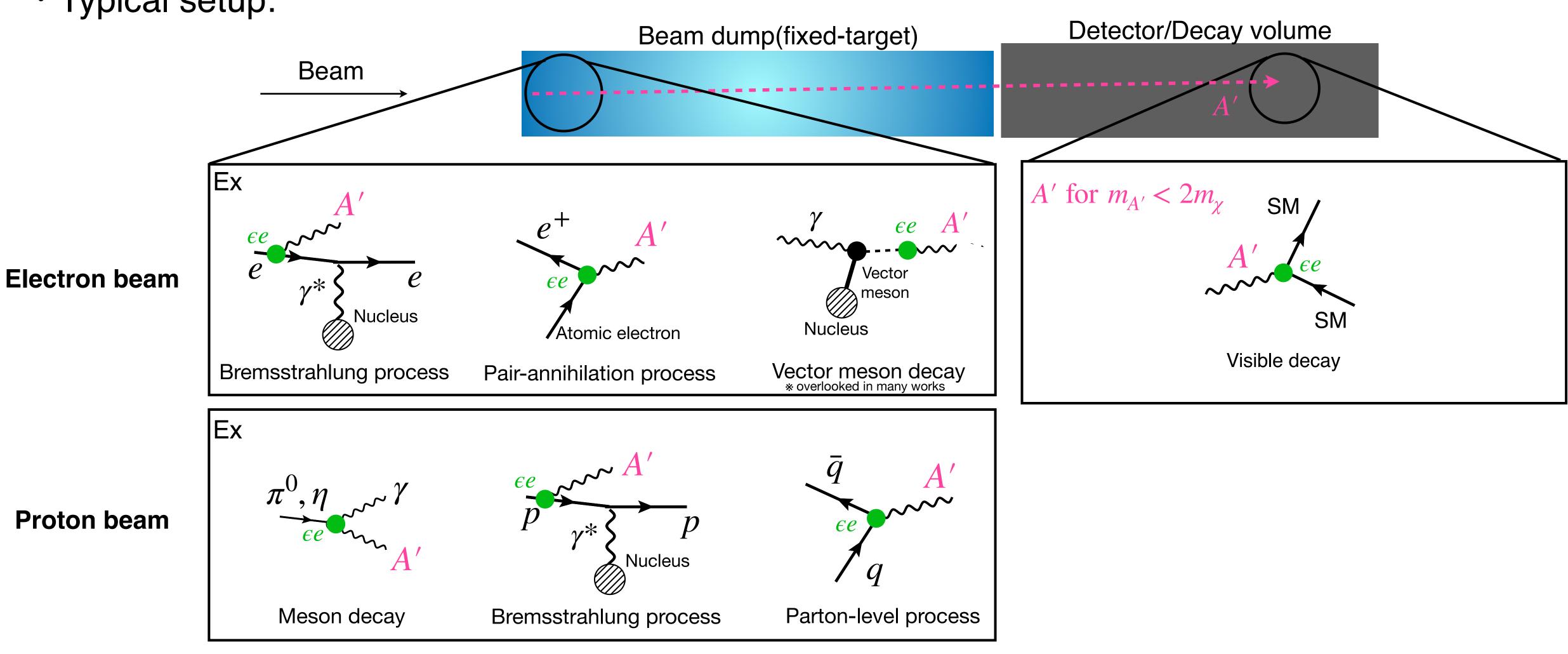
Benchmark model:  $\mathscr{L} \supset \epsilon e A'_{\mu} J^{\mu}_{EM}$  where A': Dark photon, and  $J^{\mu}_{EM}$ : SM EM current

### • Typical setup:



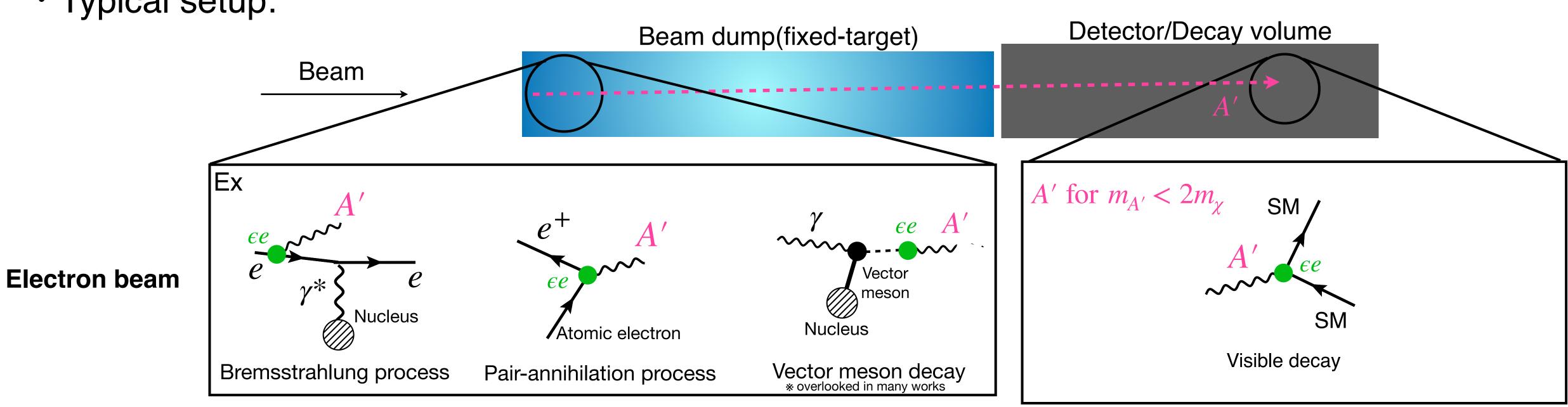
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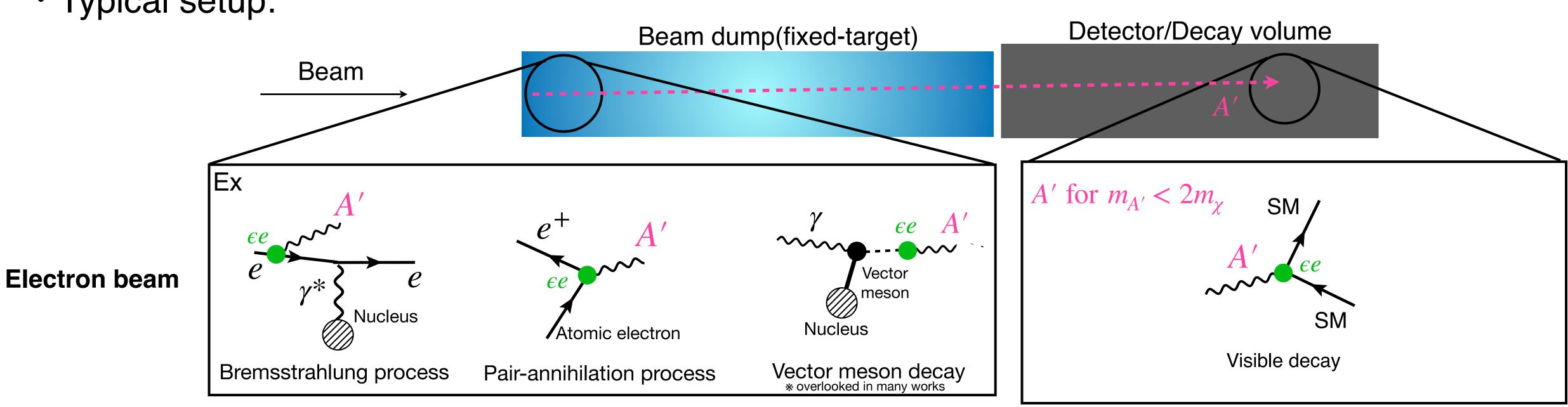
Benchmark model:  $\mathscr{L} \supset \epsilon e A'_{\mu} J^{\mu}_{EM}$  where A': Dark photon, and  $J^{\mu}_{EM}$ : SM EM current

### • Typical setup:



# of detected dark state(DS) signature:

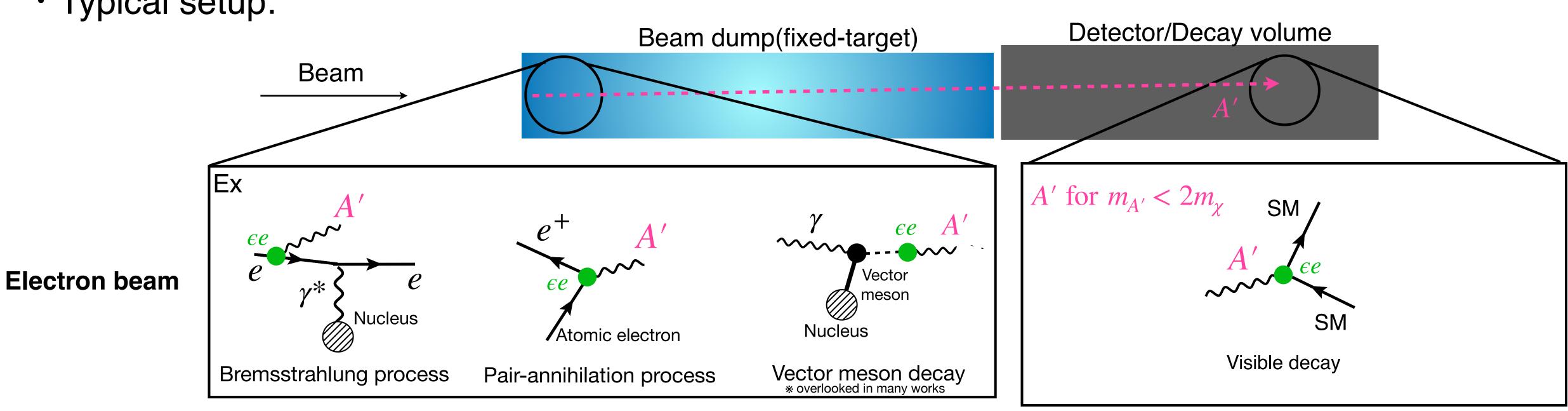




# of detected dark state(DS) signature:

~ (# of produced DS) ×

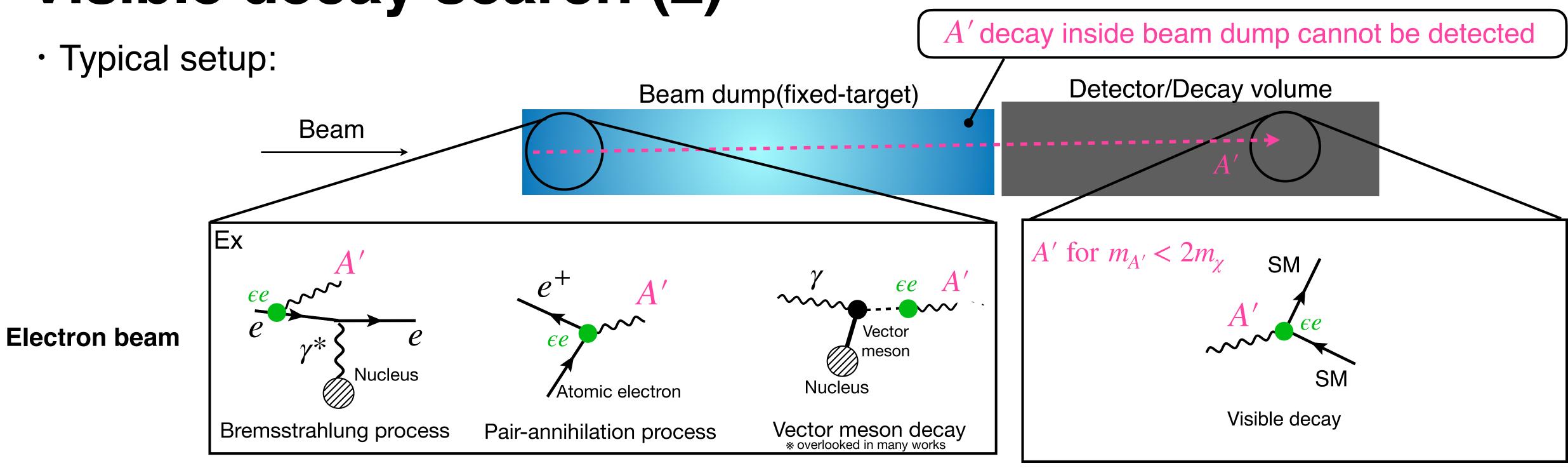




# of detected dark state(DS) signature:

 $\sim$  (# of produced DS)  $\times$  (Probability DS reaches detector)  $\times$ 

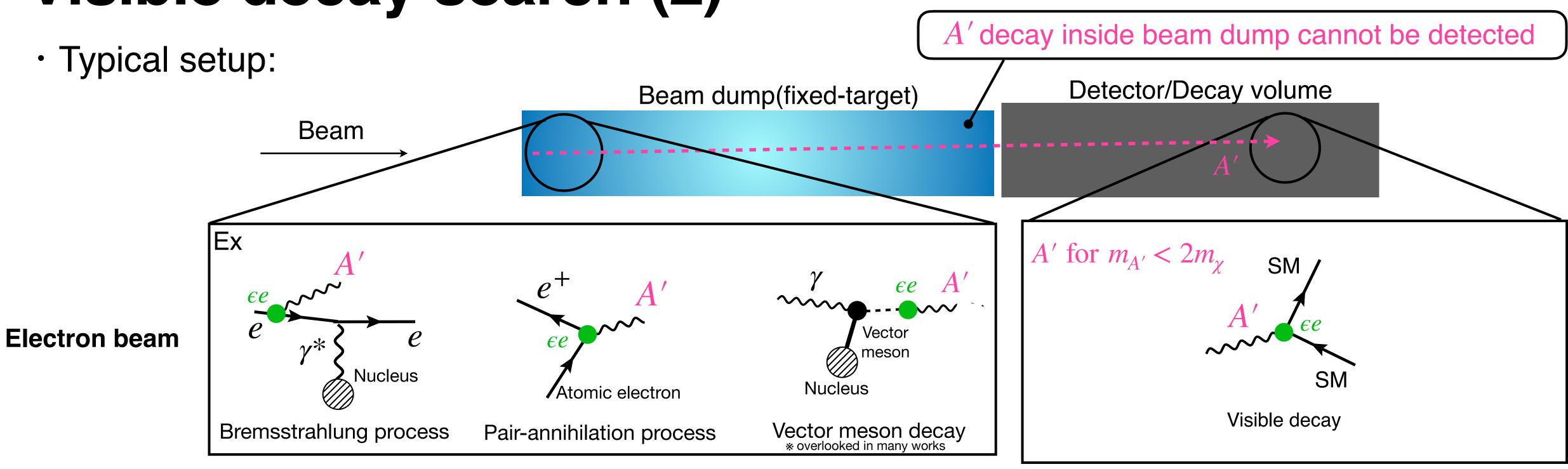




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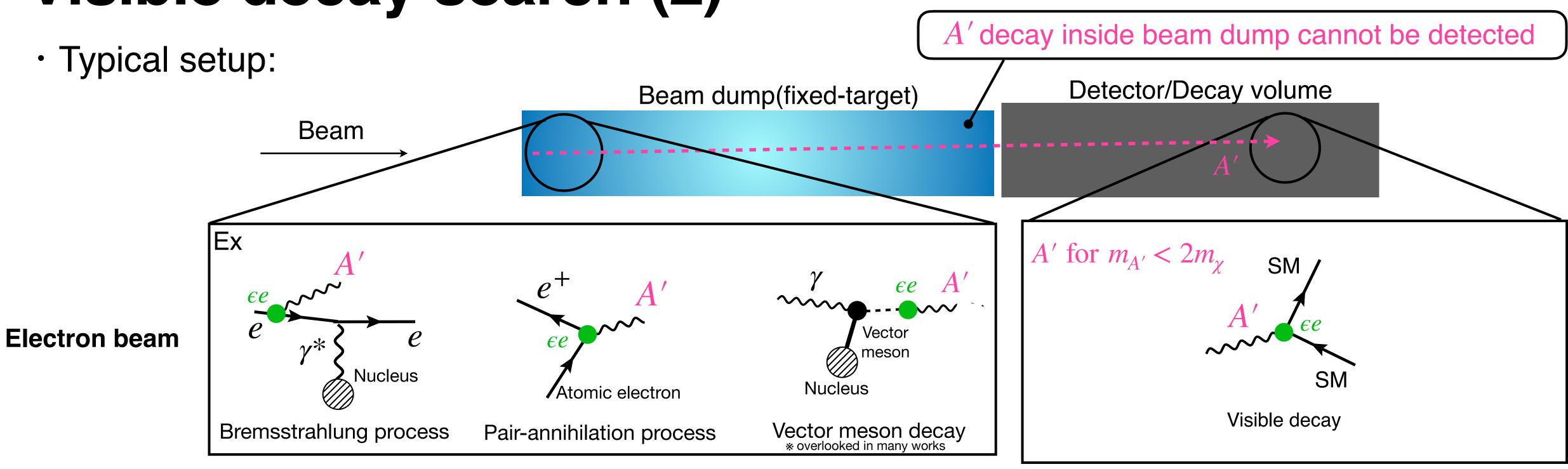
 $\sim$  (# of produced DS)  $\times$  (Probability DS reaches detector)  $\times$ 





# of detected dark state(DS) signature:

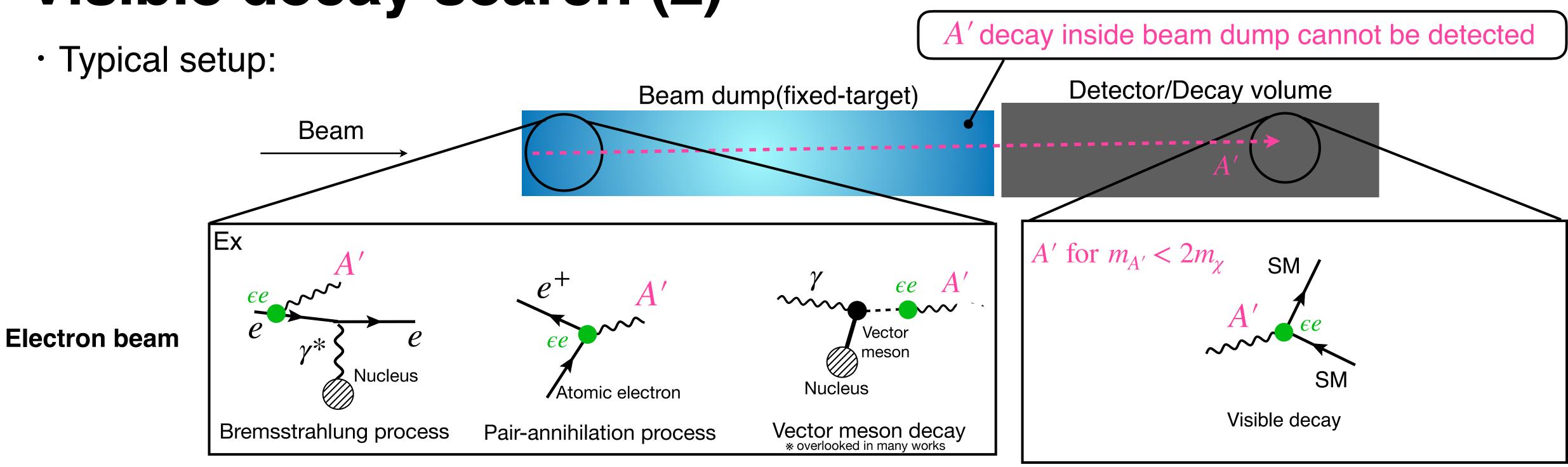




# of detected dark state(DS) signature:

#### Acceptance





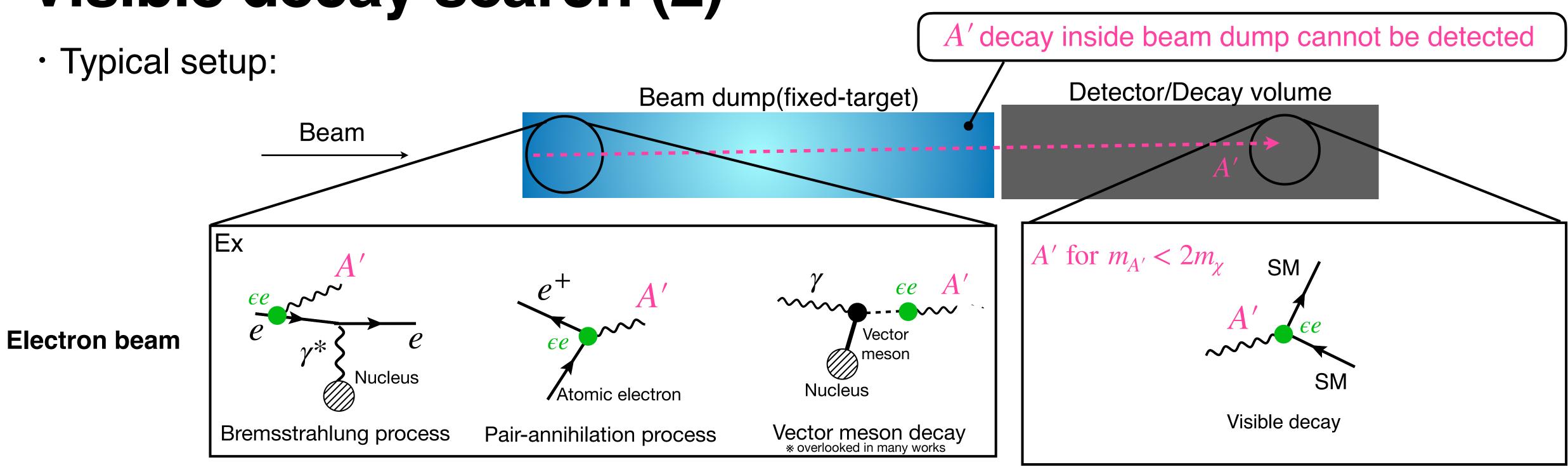
• # of detected dark state(DS) signature:

 $\propto$ 

#### Acceptance

 $\sim$  (# of produced DS) × (Probability DS reaches detector) × (Probability DS is detected)

decay length of  $A' \ge$  length of beam dump 0 decay length of A' < length of beam dump



• # of detected dark state(DS) signature:

~ (# of produced DS) × (Probability DS read  

$$\propto \begin{cases} 1 \\ 0 \end{cases}$$

### Thick(Thin) target experiments can be sensitive to long(short) lifetime dark states

\* In thin target experiments, e.g, HPS, NA64(visible decay search), continuous beam(low-intensity) is used to distinguish signal signatures

#### Acceptance

ches detector)  $\times$  (Probability DS is detected)

- decay length of  $A' \ge$  length of beam dump
- 0 decay length of A' < length of beam dump



Missing energy search, e.g., NA64:

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EM and hadron calorimeter, i.e., active target



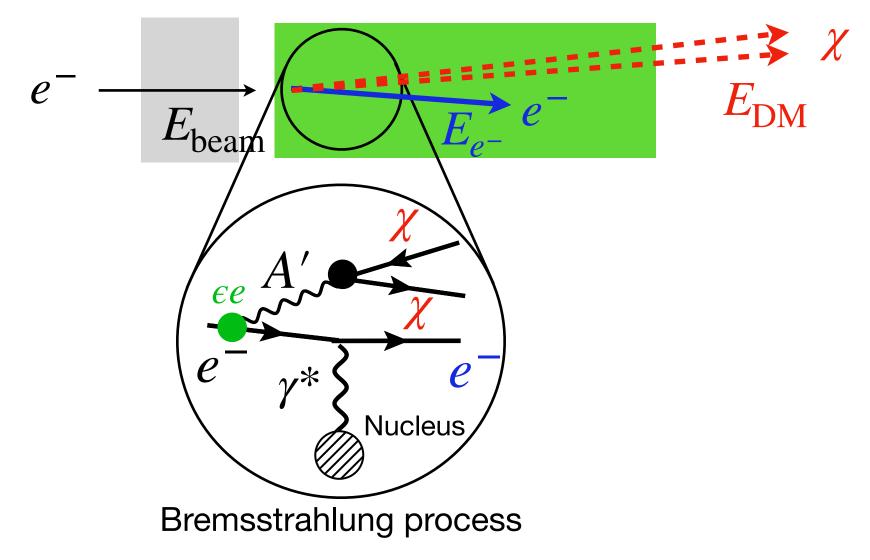
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Tracker EM and hadron calorimeter, i.e., active target



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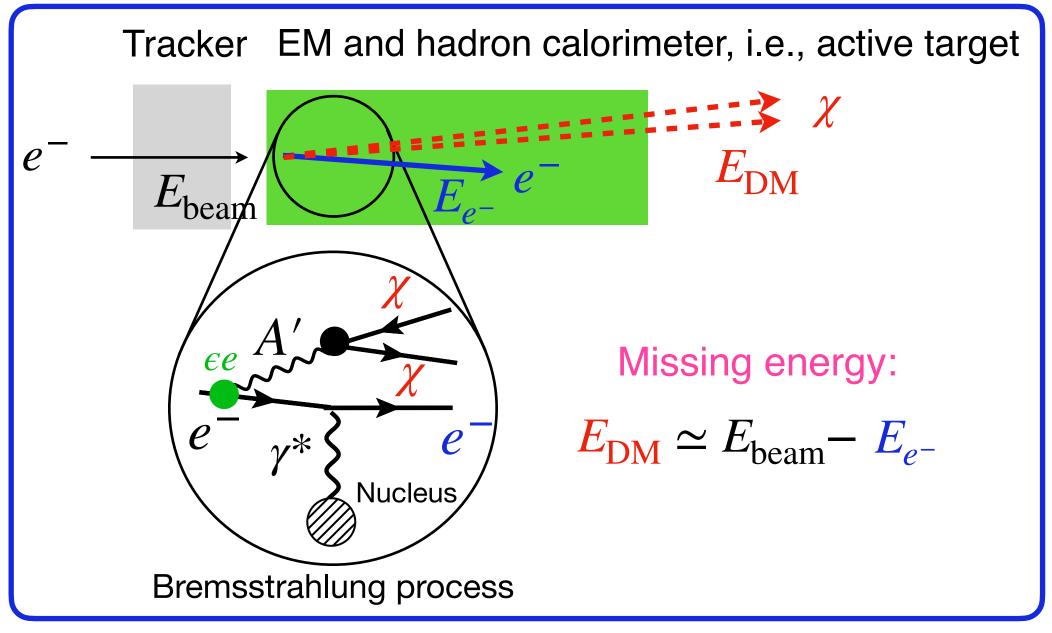
Tracker EM and hadron calorimeter, i.e., active target



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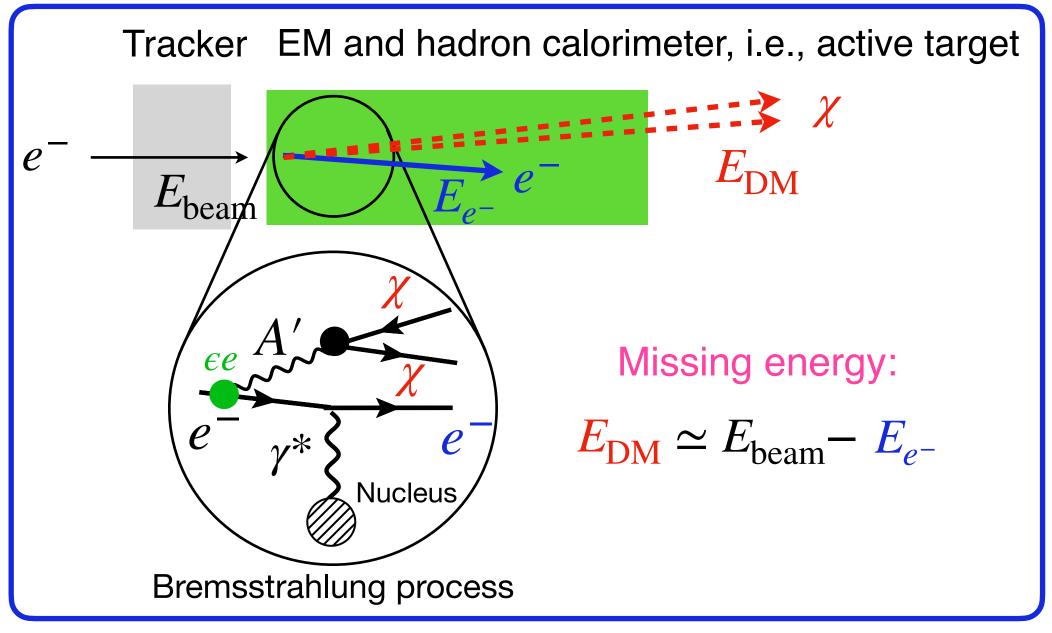


\* Energy transfer to nucleus is modest in Bremsstrahlung process

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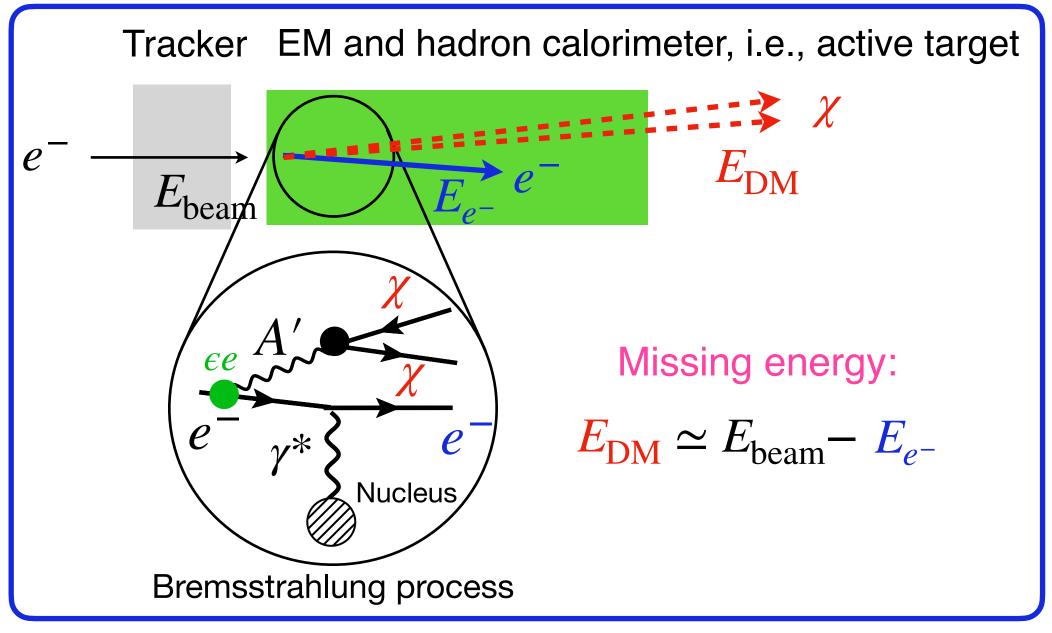
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Missing momentum search, e.g., LDMX:



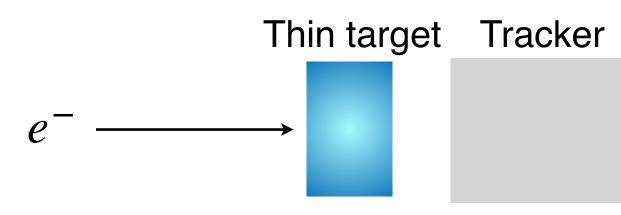
#### Missing energy search, e.g., NA64:



\* Energy transfer to nucleus is modest in Bremsstrahlung process

Benchmark model:  $\mathscr{L} \supset \epsilon e A'_{\mu} J^{\mu}_{EM} - g_D A'_{\mu} \bar{\chi} \gamma^{\mu} \chi A'$ : Dark photon,  $\chi$ : Dark Matter,  $J^{\mu}_{EM}$ : SM EM current

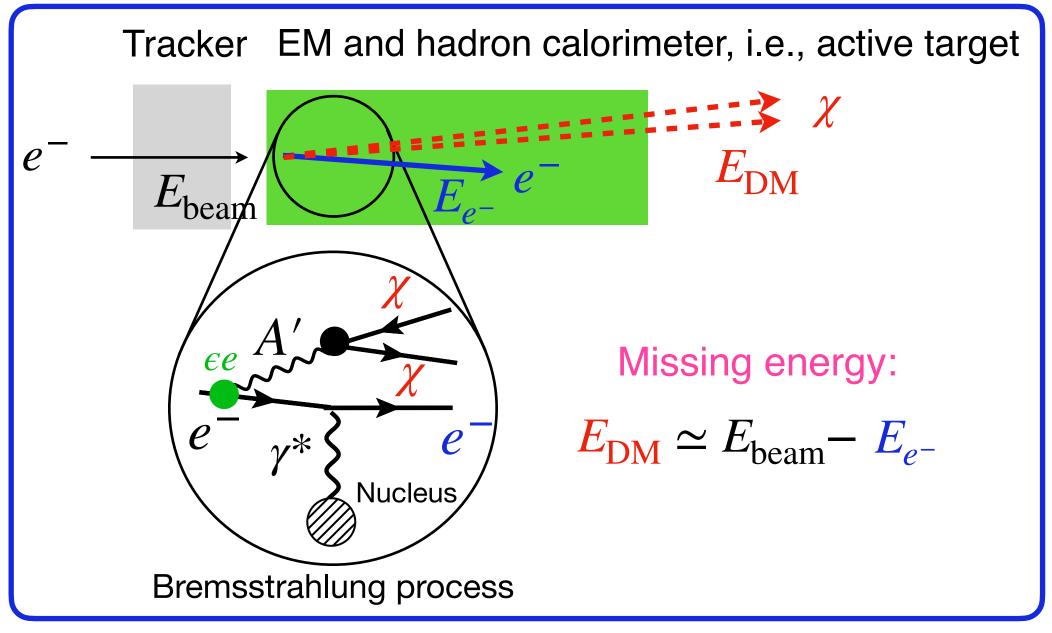
#### Missing momentum search, e.g., LDMX:



EM and Hadron calorimeter



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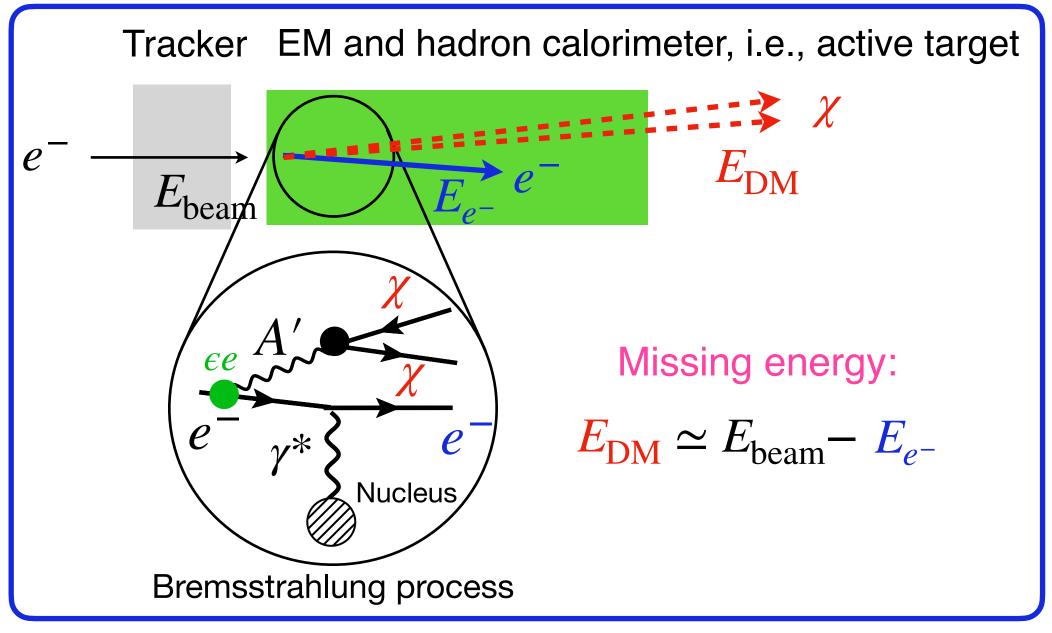
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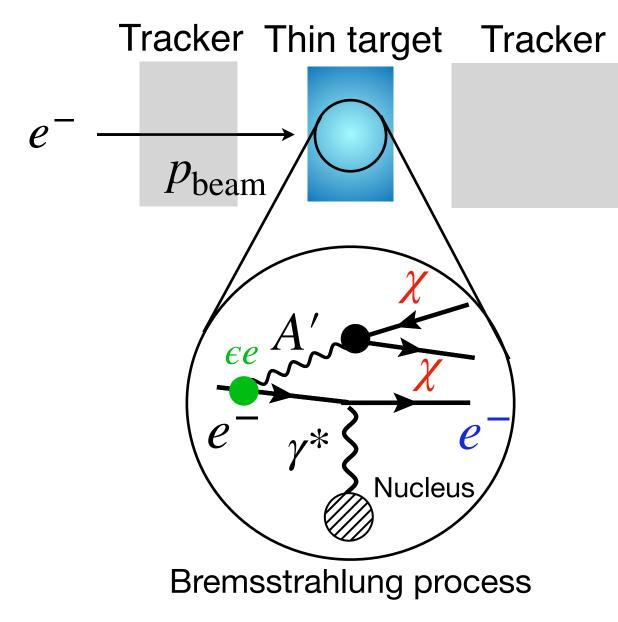
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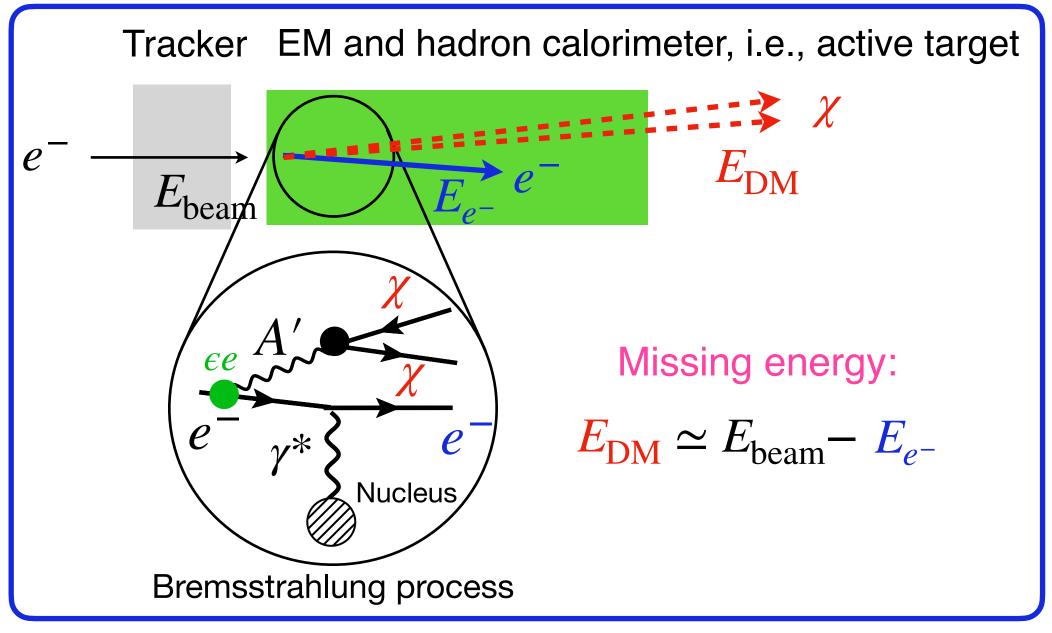
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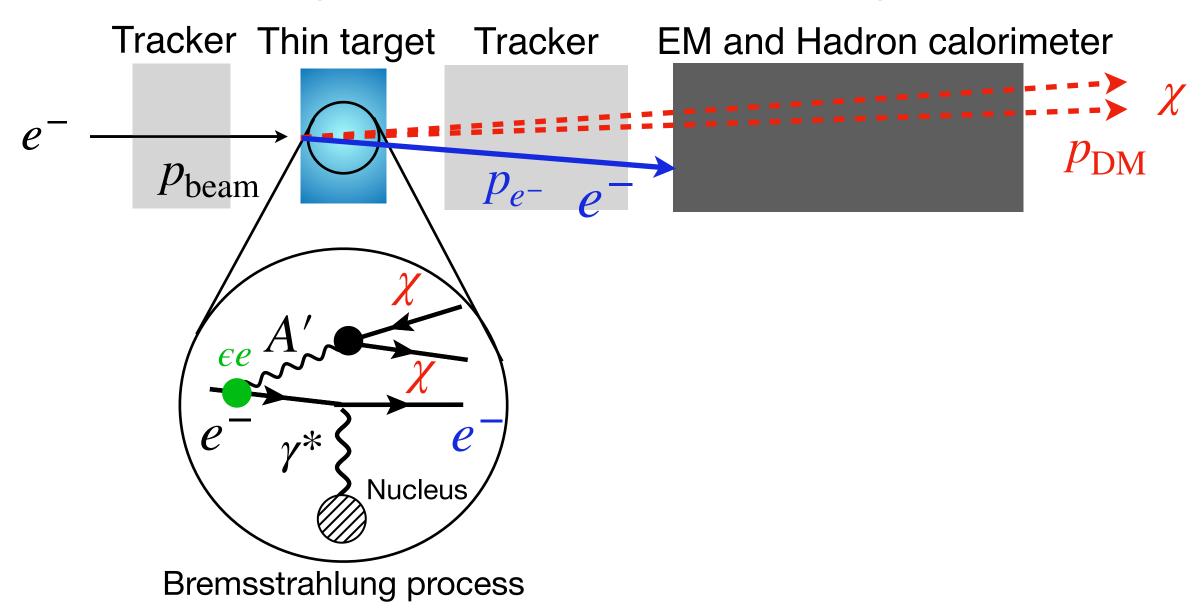


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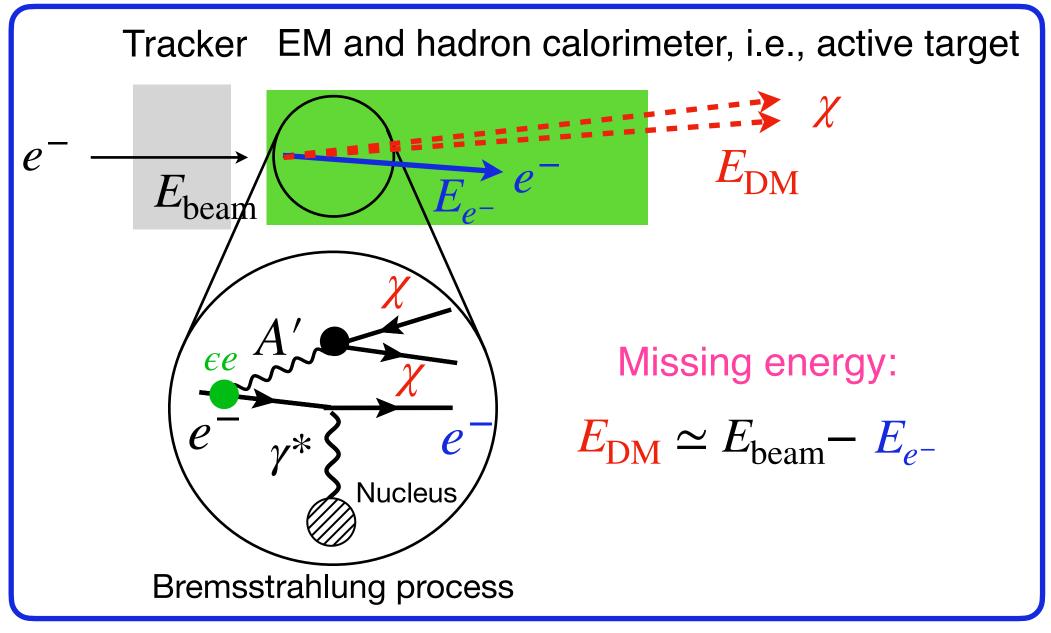
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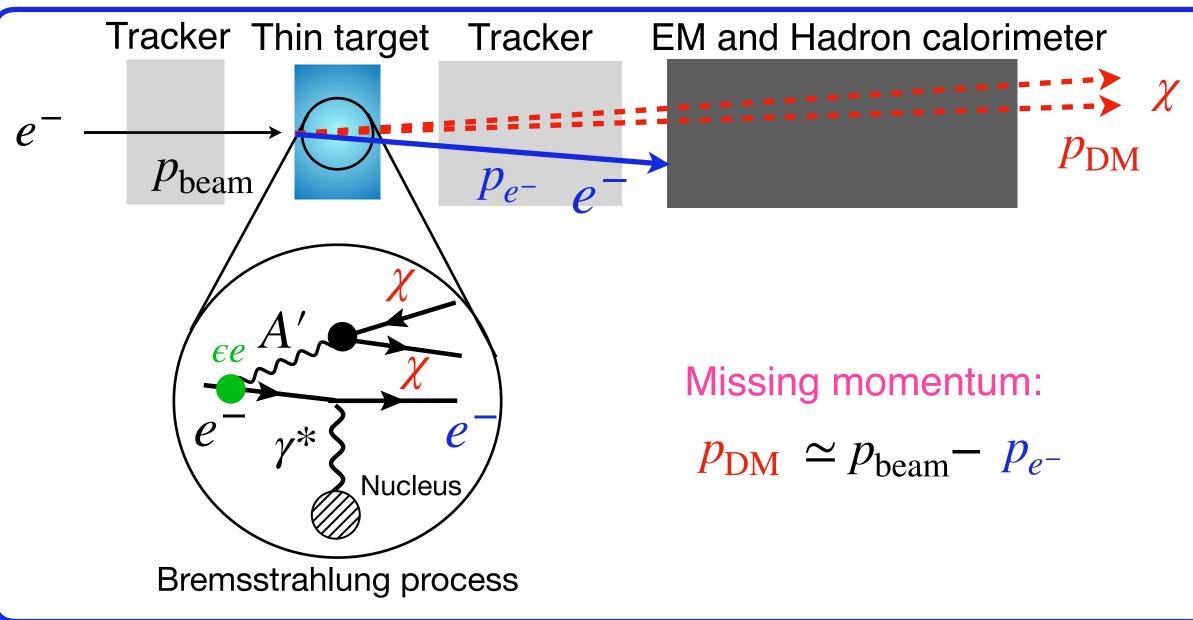
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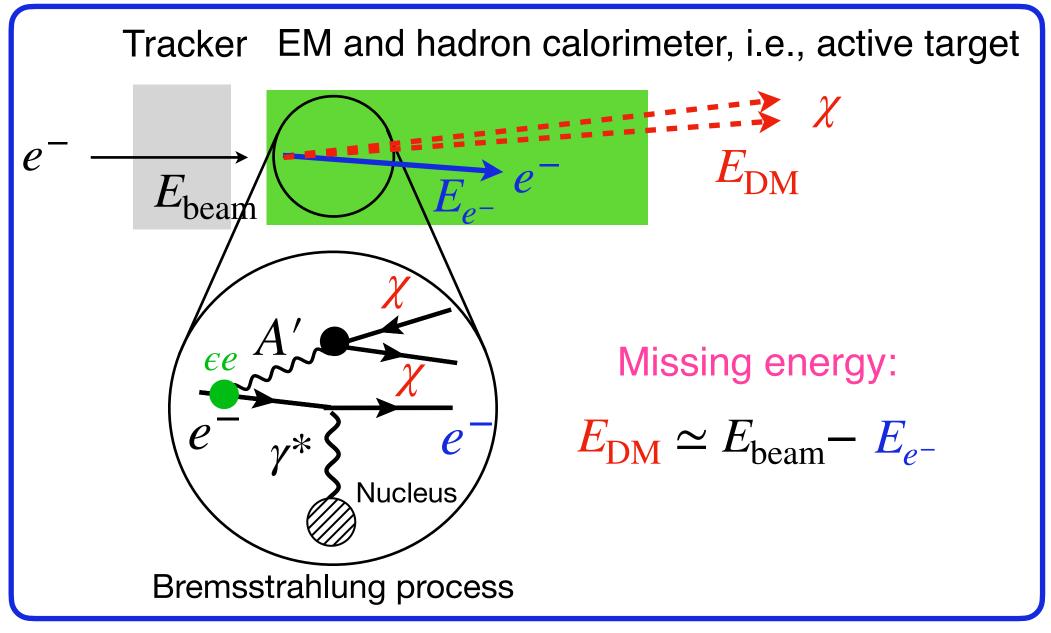
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\* Target is thin to reconstruct final state electron

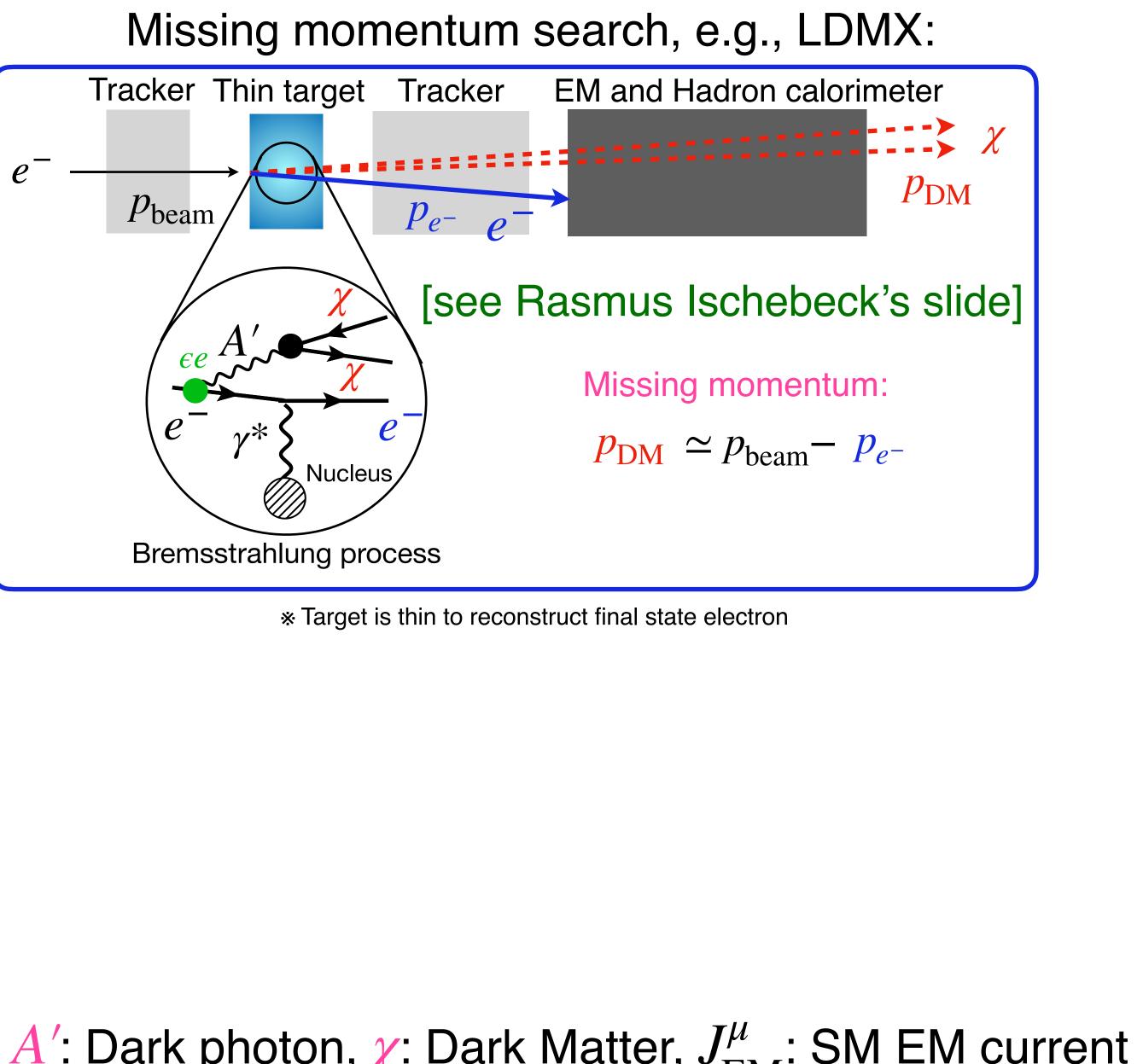


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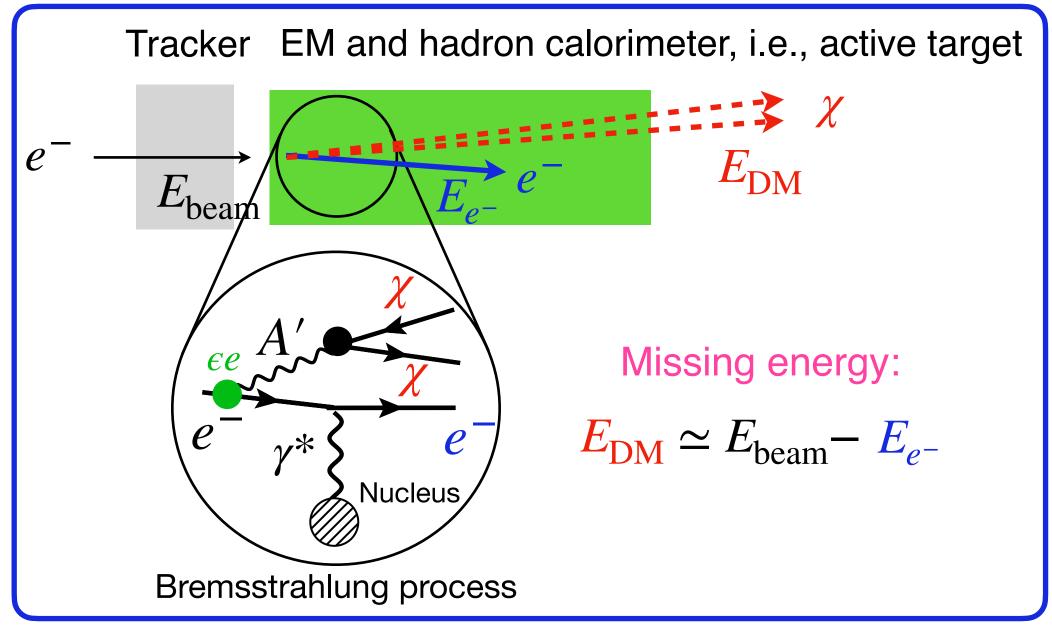


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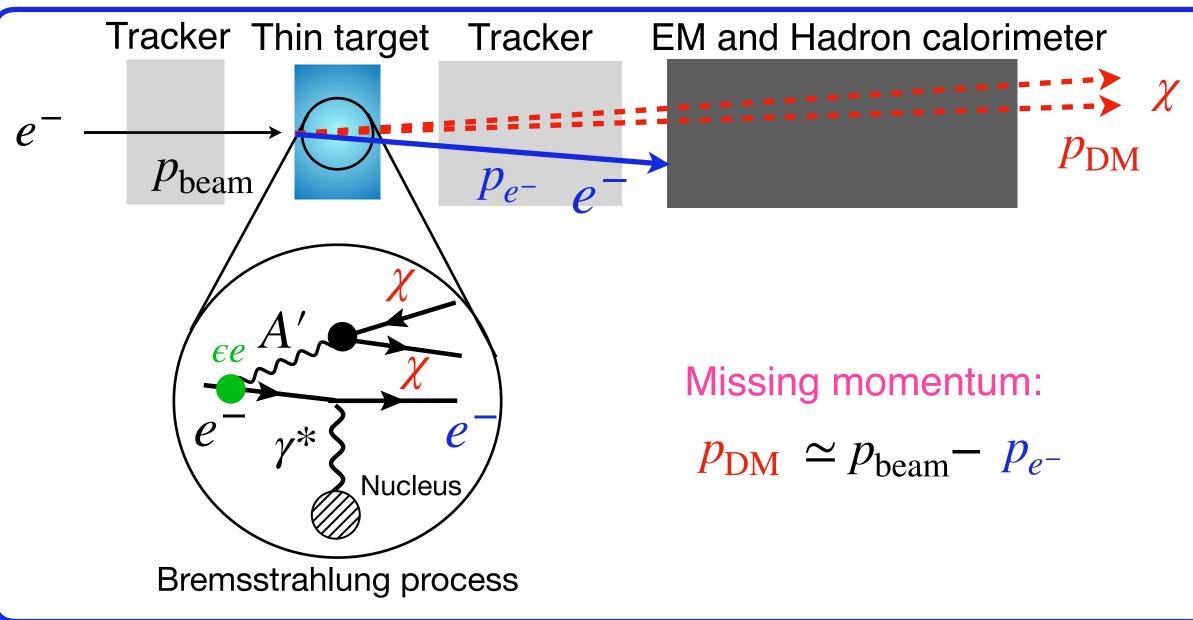


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# of missing events:

Benchmark model: 
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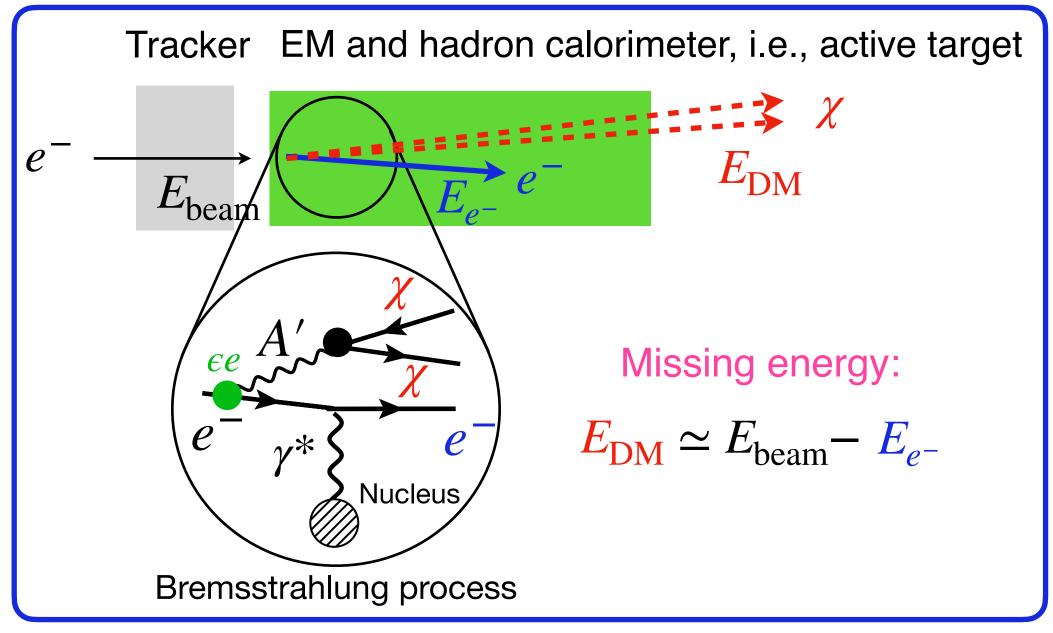


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A': Dark photon,  $\chi$ : Dark Matter,  $J_{\rm EM}^{\mu}$ : SM EM current



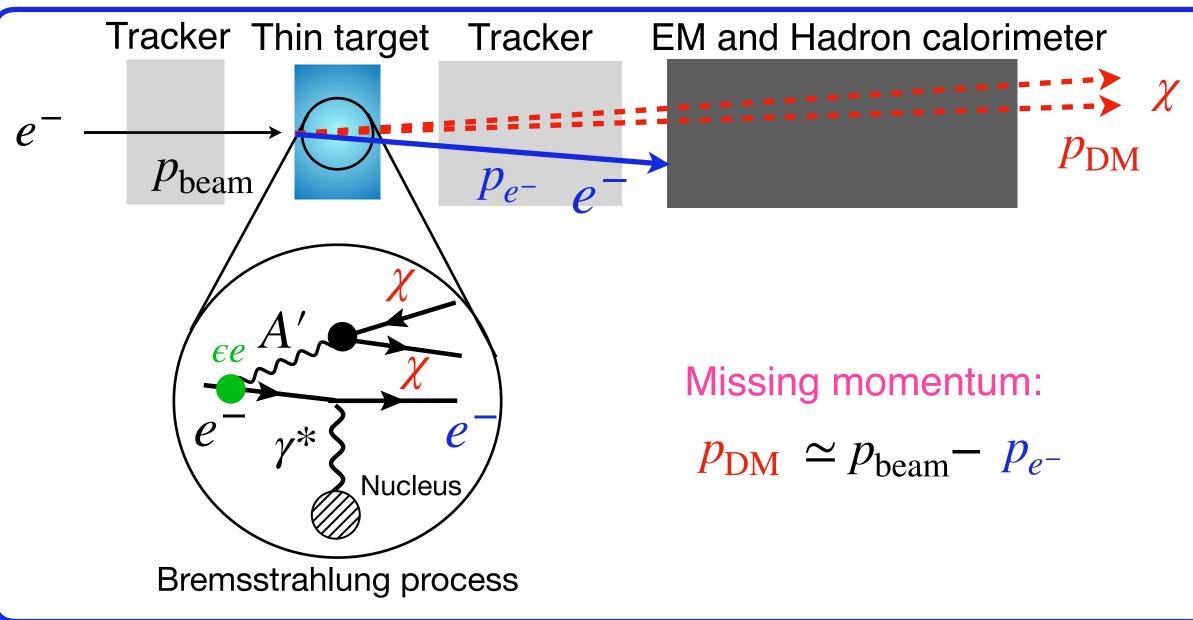
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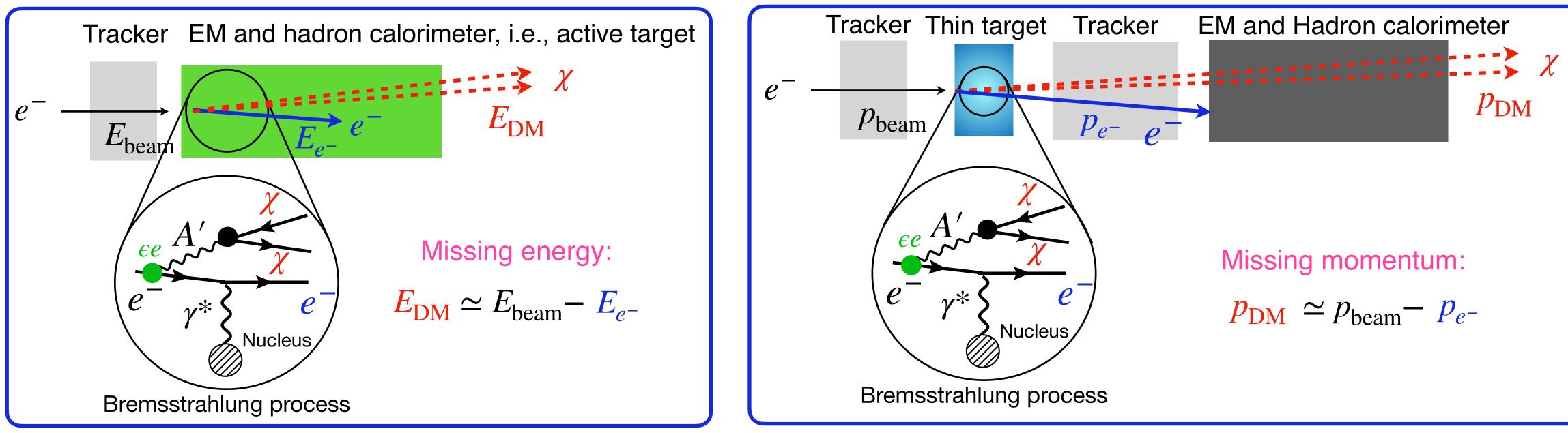


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#### Missing energy search, e.g., NA64:



\* Energy transfer to nucleus is modest in Bremsstrahlung process

• # of missing events:

\* not proportional to  $(\epsilon e)^2$  in contrast to recoil and visible search

Benchmark model:  $\mathscr{L} \supset \epsilon e A'_{\mu} J^{\mu}_{EM} - g_D A'_{\mu} \bar{\chi} \gamma^{\mu} \chi A'$ : Dark photon,  $\chi$ : Dark Matter,  $J^{\mu}_{EM}$ : SM EM current

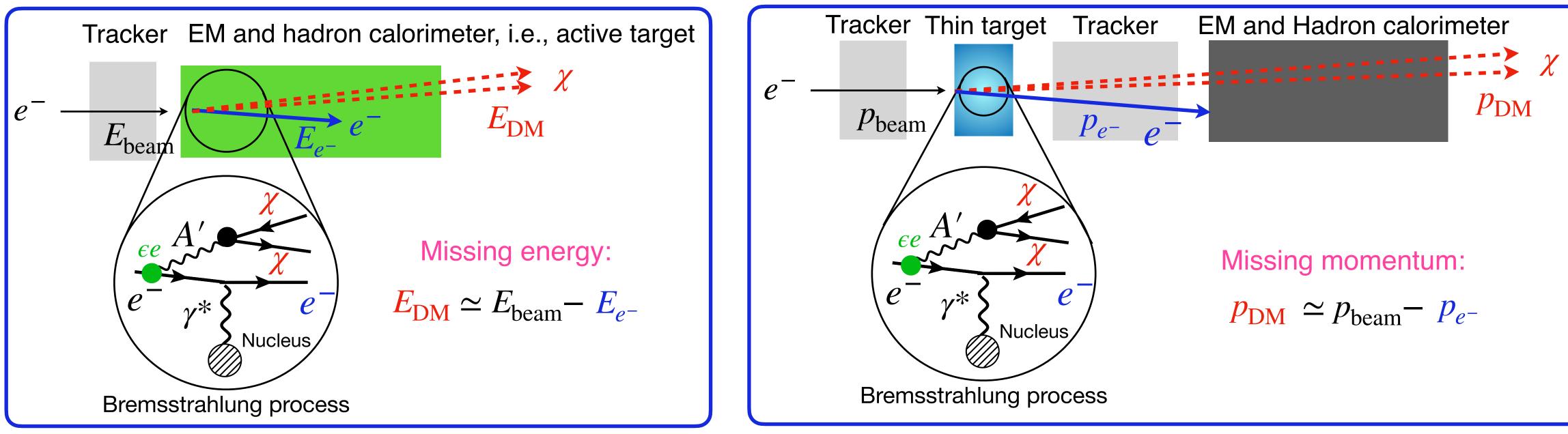
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\* Target is thin to reconstruct final state electron

#### Acceptance



#### Missing energy search, e.g., NA64:



\* Energy transfer to nucleus is modest in Bremsstrahlung process

• # of missing events:

~ (# of produced DM) × (Probability DM rea  

$$\propto$$
 (Beam flux)×( $\epsilon e$ )<sup>2</sup> \* not proportional

Missing momentum search, e.g., LDMX:

\* Target is thin to reconstruct final state electron

#### Acceptance

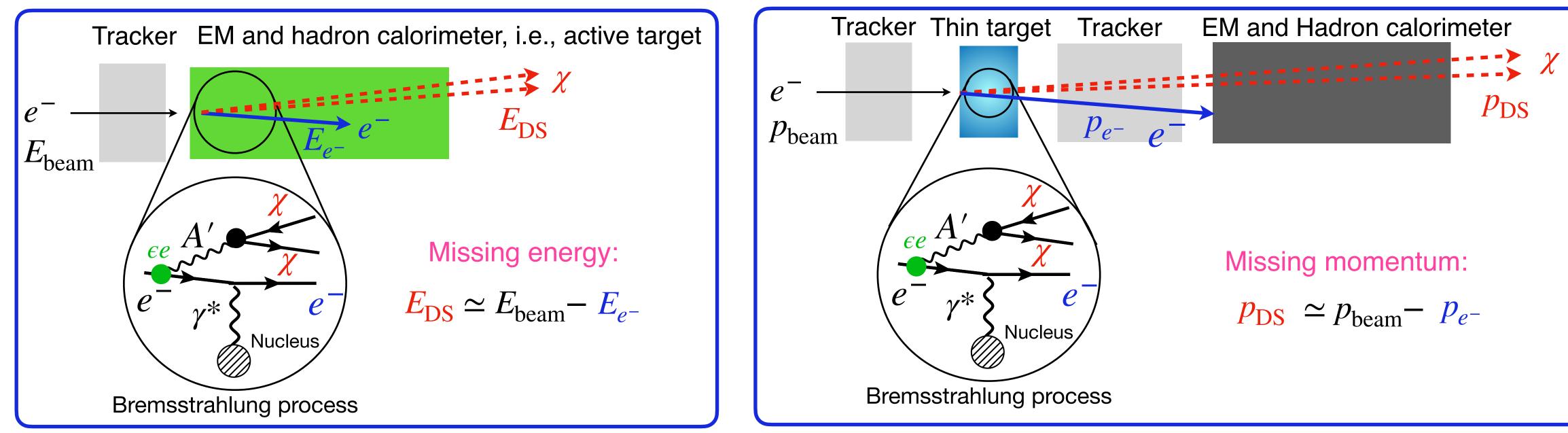
ches detector) × (Probability DM is detected)

Ito  $(\epsilon e)^2$  in contrast to recoil and visible search

Benchmark model:  $\mathscr{L} \supset \epsilon e A'_{\mu} J^{\mu}_{EM} - g_D A'_{\mu} \bar{\chi} \gamma^{\mu} \chi A'$ : Dark photon,  $\chi$ : Dark Matter,  $J^{\mu}_{EM}$ : SM EM current



Missing energy search, e.g., NA64:



• *#* of missing events:

 $\propto$  (Beam flux)×( $\epsilon e$ )<sup>2</sup>

Acceptance is good, but the continuous beam (low-intensity) is needed to reconstruct final state

Missing momentum search, e.g., LDMX:

\* Energy transfer to nucleus is modest in Bremsstrahlung process

#### Acceptance

~ (# of produced DM)  $\times$  (Probability DM reaches detector)  $\times$  (Probability DM is detected)

\* not proportional to  $(\epsilon e)^2$  in contrast to recoil and visible search

Benchmark model:  $\mathscr{L} \supset \epsilon e A'_{\mu} J^{\mu}_{EM} - g_D A'_{\mu} \bar{\chi} \gamma^{\mu} \chi A'$ : Dark photon,  $\chi$ : Dark Matter,  $J^{\mu}_{EM}$ : SM EM current

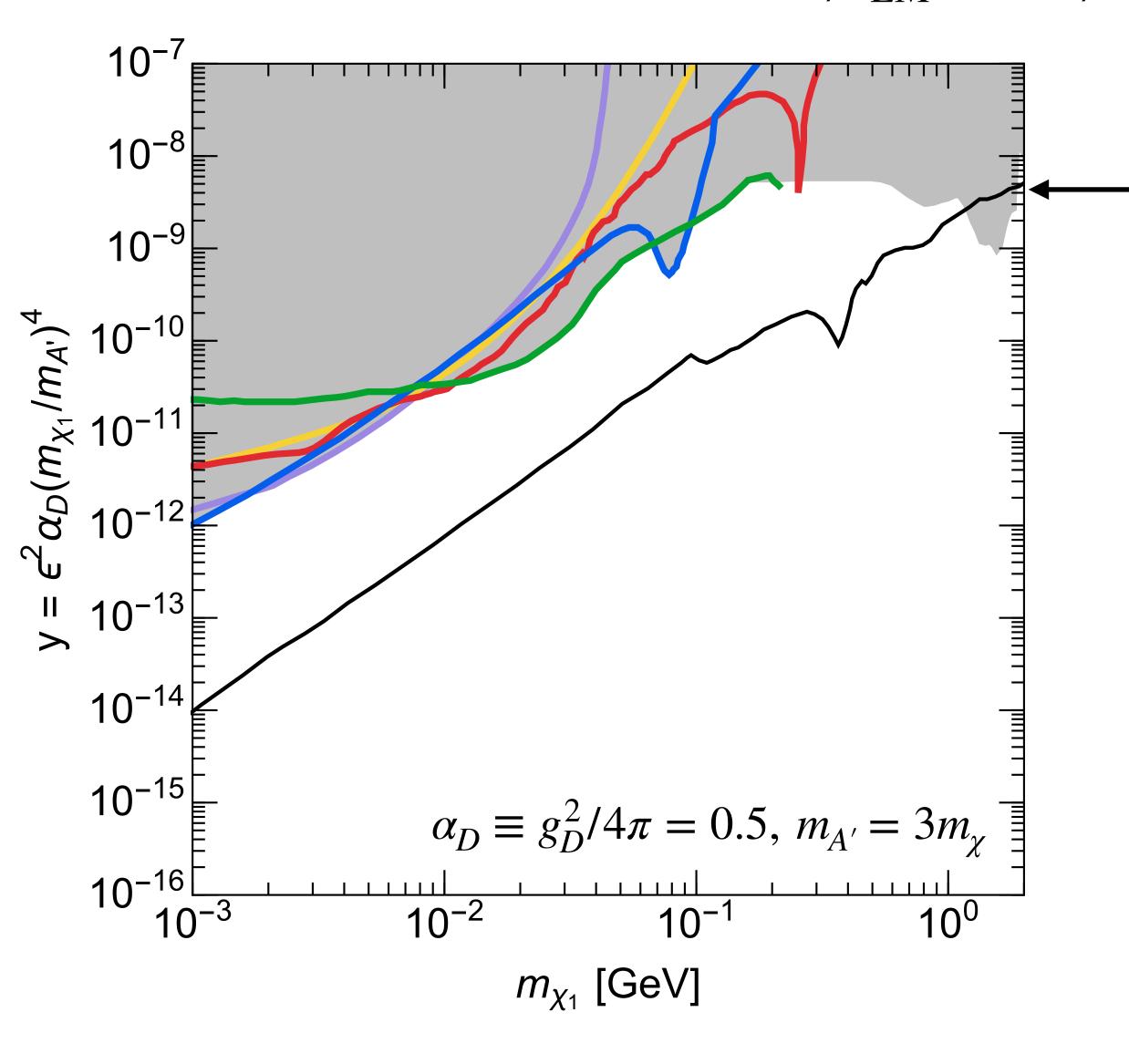


# Outline

- Introduction
  - dark sector and Sub-GeV dark matter
  - beam dump(fixed-target) experiment
- Key features of beam dump experiment
- A classification of beam dump experiment
- Sensitivity of beam dump experiments at future accelerators
- Summary

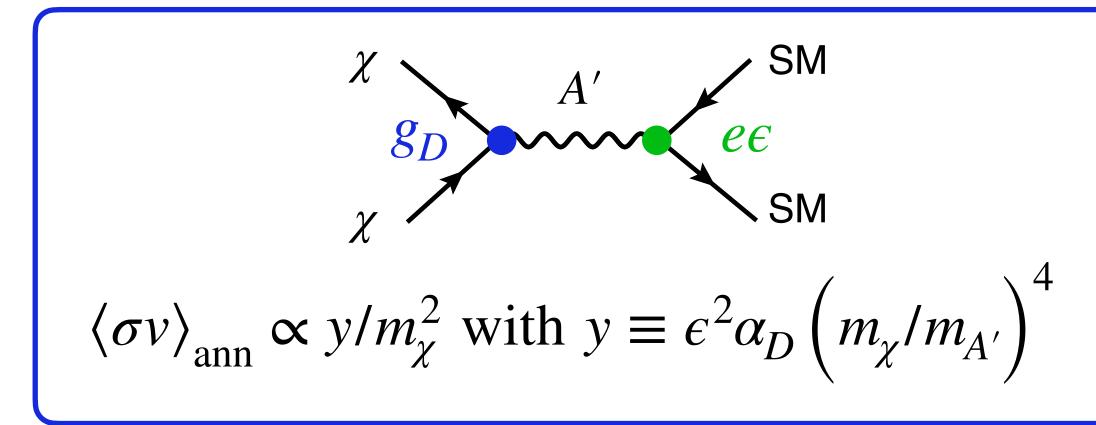
### Excluded regions by beam dump experiments (1)

• Benchmark model (1):  $\mathscr{L} \supset \epsilon \cdot eA'_{\mu}J^{\mu}_{EM} - g_D A'_{\mu}\bar{\chi}\gamma^{\mu}\chi$  where A': dark photon, and  $\chi$ : DM



 $\chi$  saturates observed DM abundance

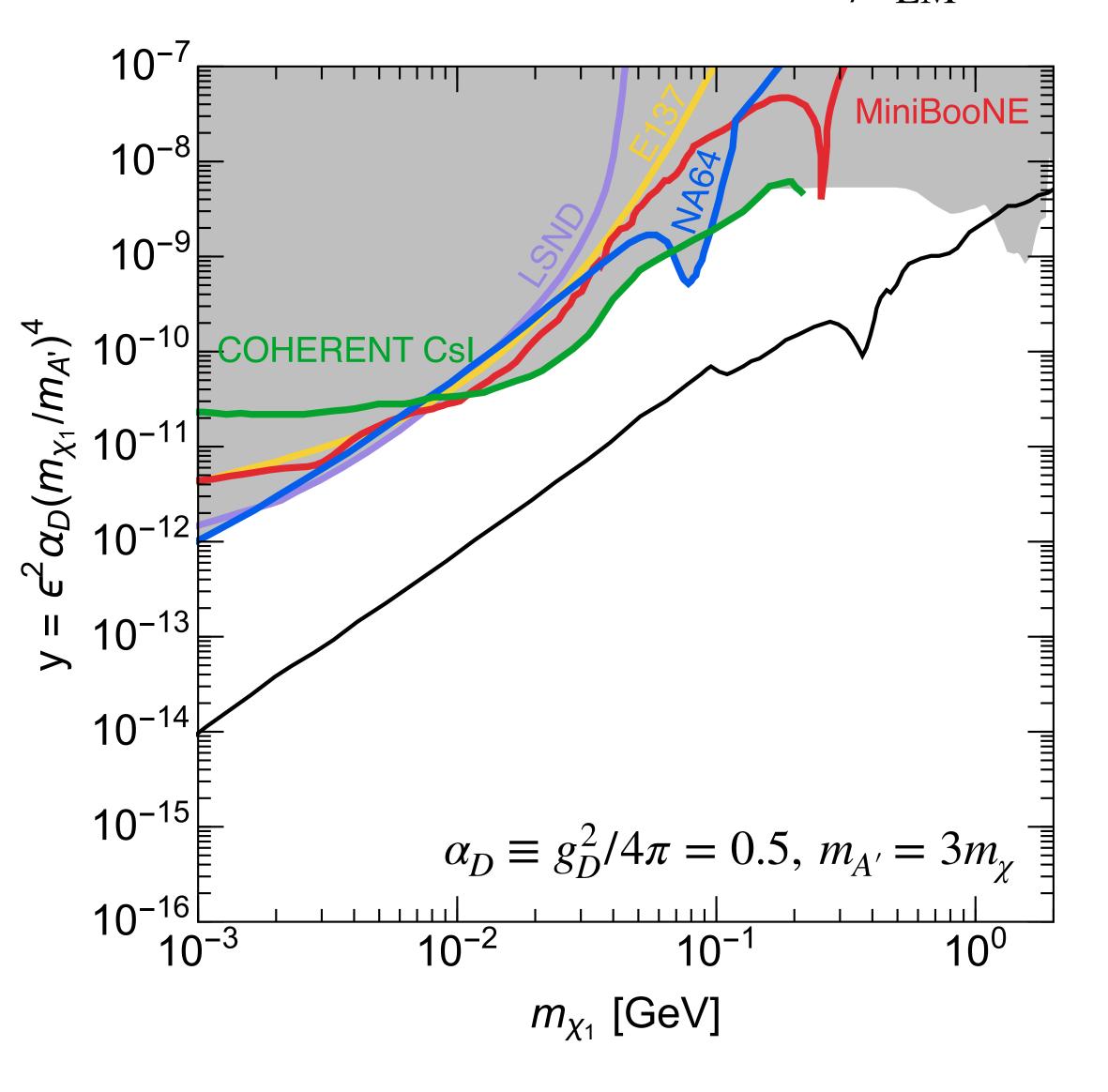
DM annihilation cross section

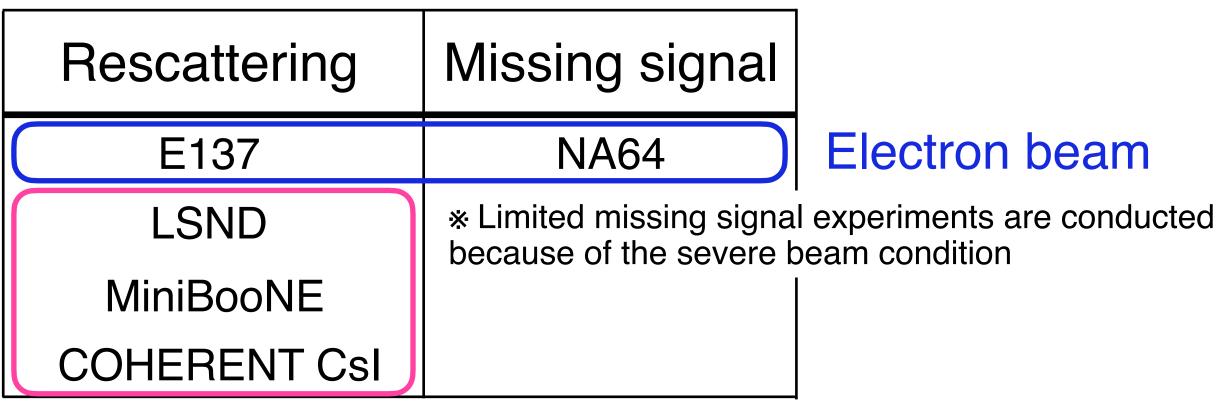




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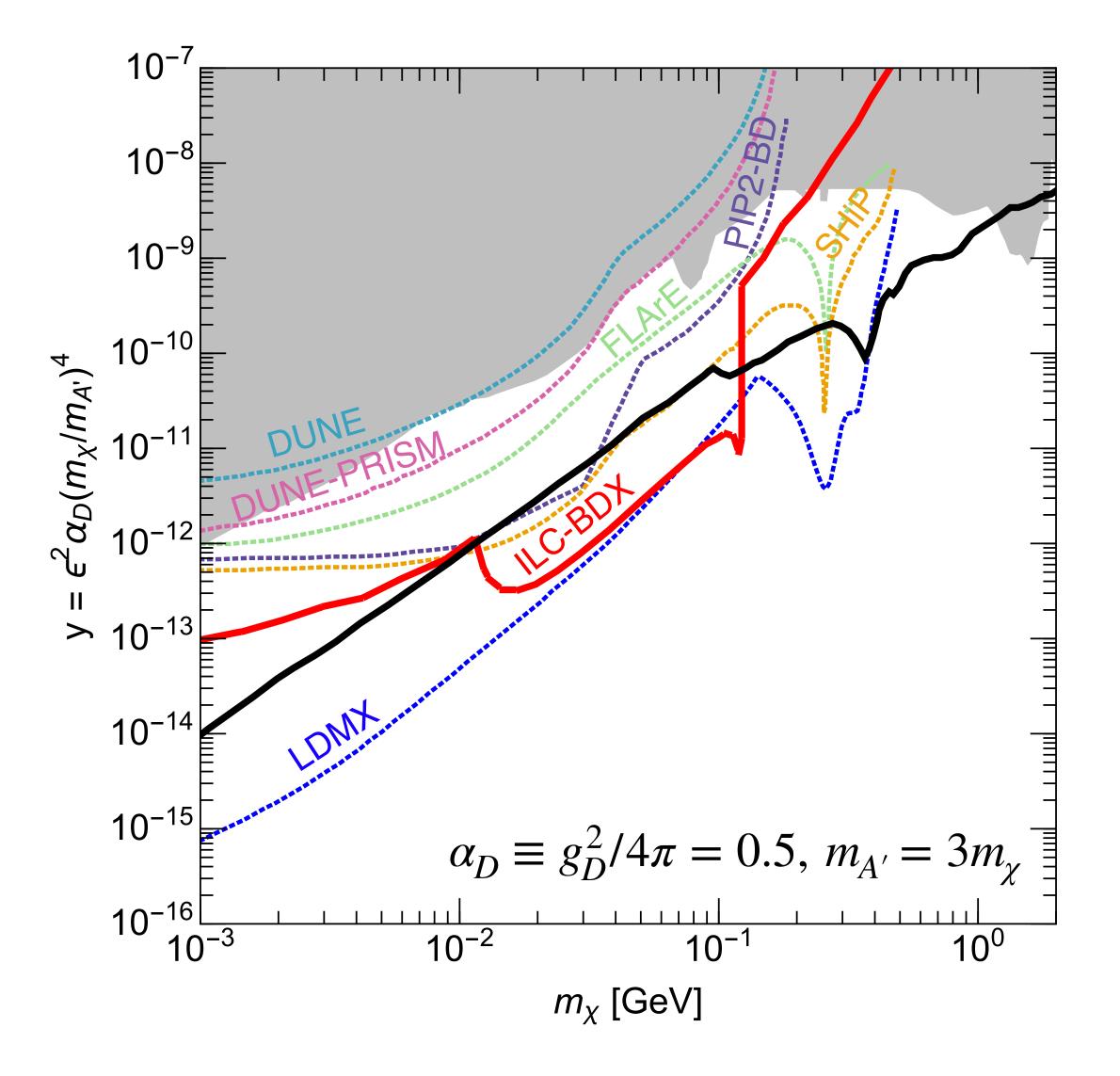


### Proton beam

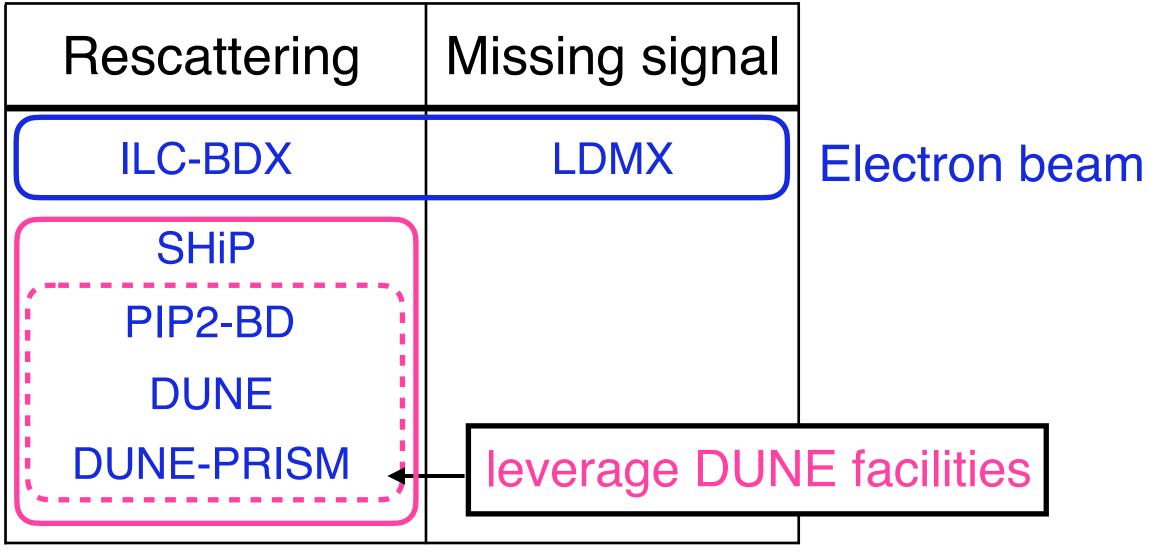
- $\Rightarrow$  parasitic running of neutrino experiment
  - \* MiniBooNE is off-target running to reduce neutrino BG

### Sensitivity of beam dump experiments at future accelerators (1)

• Benchmark model (1):  $\mathscr{L} \supset \epsilon \cdot eA'_{\mu}J^{\mu}_{EM} - g_D A'_{\mu}\bar{\chi}\gamma^{\mu}\chi$  where A': dark photon, and  $\chi$ : DM



Examples of BD at future accelerators



Proton beam

- LDMX is highly sensitive because of good acceptance
- ILC is sensitive because of high energy positron annihilation
- DUNE-PRISM(off-axis detector) is more sensitive than DUNE(on-axis detector) because of neutrino BG reductions

\* FLARE is LHC auxiliary detector experiments in HL-LHC phase



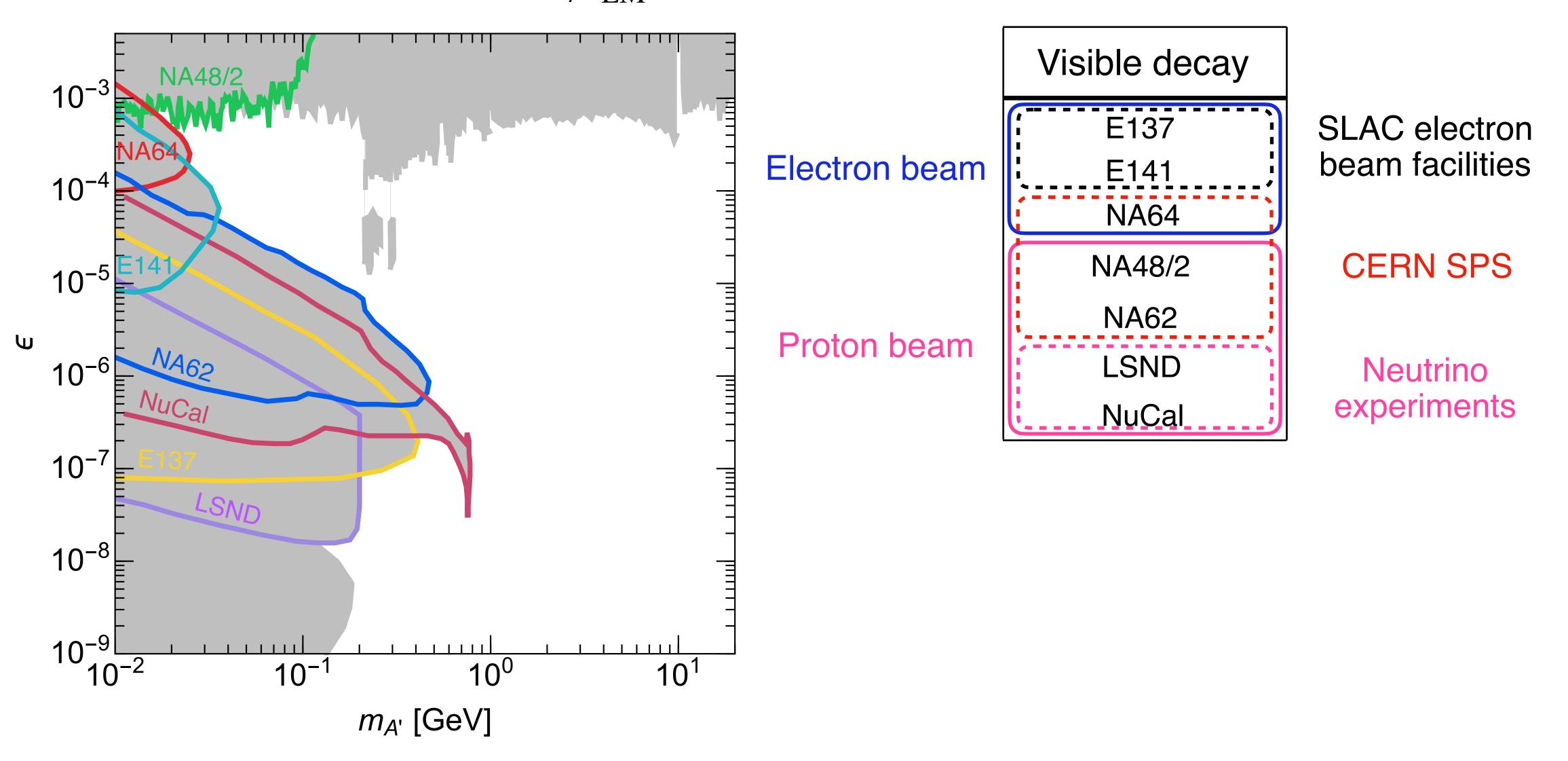






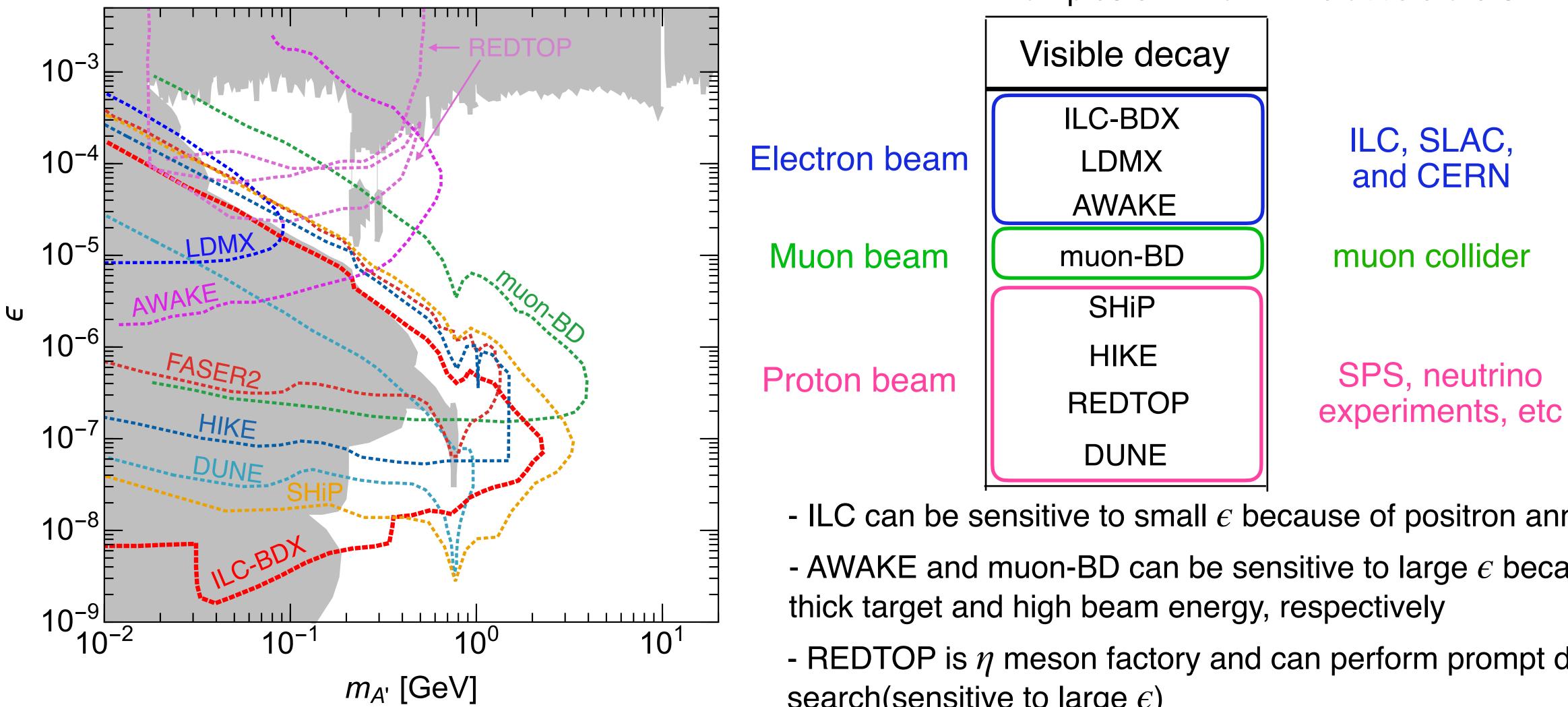
# **Excluded regions by beam dump experiments (2)**

• Benchmark model (2):  $\mathscr{L} \supset \epsilon \cdot eA'_{\mu}J^{\mu}_{EM}$  where A': dark photon



### Sensitivity of beam dump experiments at future accelerators (2)

• Benchmark model (2):  $\mathscr{L} \supset \epsilon \cdot eA'_{\mu}J^{\mu}_{EM}$  where A': dark photon



Examples of BD at future accelerators

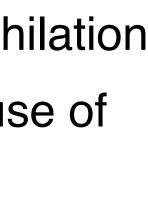
- ILC can be sensitive to small  $\epsilon$  because of positron annihilation

- AWAKE and muon-BD can be sensitive to large  $\epsilon$  because of

- REDTOP is  $\eta$  meson factory and can perform prompt decay search(sensitive to large  $\epsilon$ )

\* FASER2 is LHC auxiliary detector experiments in HL-LHC phase









# Summary

- Key features of beam dump(fixed target) experiments:
  - The beam dump experiment is high luminosity experiment sensitive to Sub-GeV scale  $\Rightarrow$  The physical potential is complemental to the other experiments, e.g., collider experiments, direct detection experiments
  - The beam dump experiment can run in parallel with accelerator-based experiments

e.g., HL-LHC, ILC, and muon collider.

beam dump experiments

beyond the SM

- $\Rightarrow$  The beam dump experiments are economical and would also run with future accelerators,
- $\Rightarrow$  High energy and high flux beams in the future accelerator lead to high sensitivity of the parasitic
- Regarding the dark sector search, the physical potential of the beam dump experiment depends on various factors, e.g., beam flux, beam energy, beam particle, acceptance, detection approach,...
- The beam dump experiment tandems with the future accelerators and potentially sheds light on the







