

Istituto Nazionale di Fisica Nucleare



# Physics channels at LHCspin

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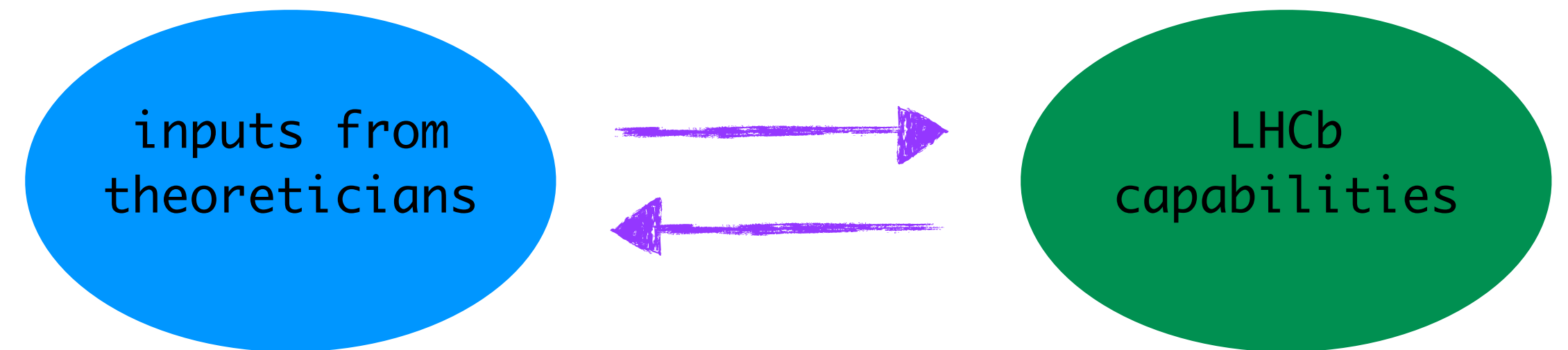
LHCspin kick-off meeting, 18/12/2023



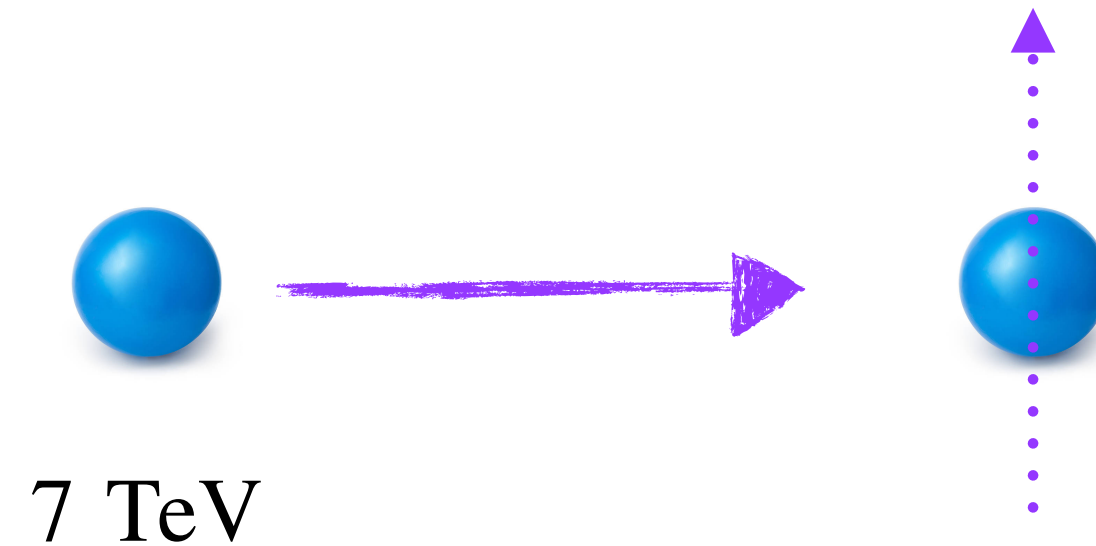
# Introduction

In this talk:

1. Kinematic acceptance & reconstruction efficiency
2. Expected yields for selected channels
3. An example measurement: the gluon Sivvers function



- The results are based on **full LHCb simulations** with p-H FT collisions
- Plenty of channels available, will focus on  $J/\psi \rightarrow \mu^+\mu^-$  as a benchmark



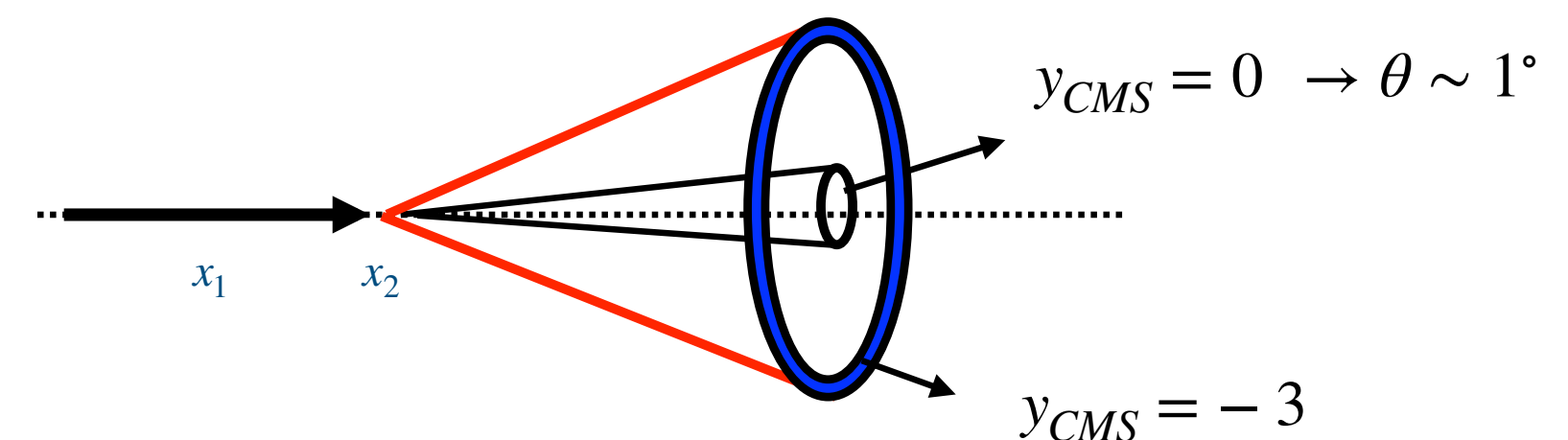
$$\sqrt{s} = \sqrt{2m_N E_p} = 115 \text{ GeV}$$

$$\gamma = \frac{\sqrt{s}}{2m_p} \sim 60$$

1: beam, 2: target

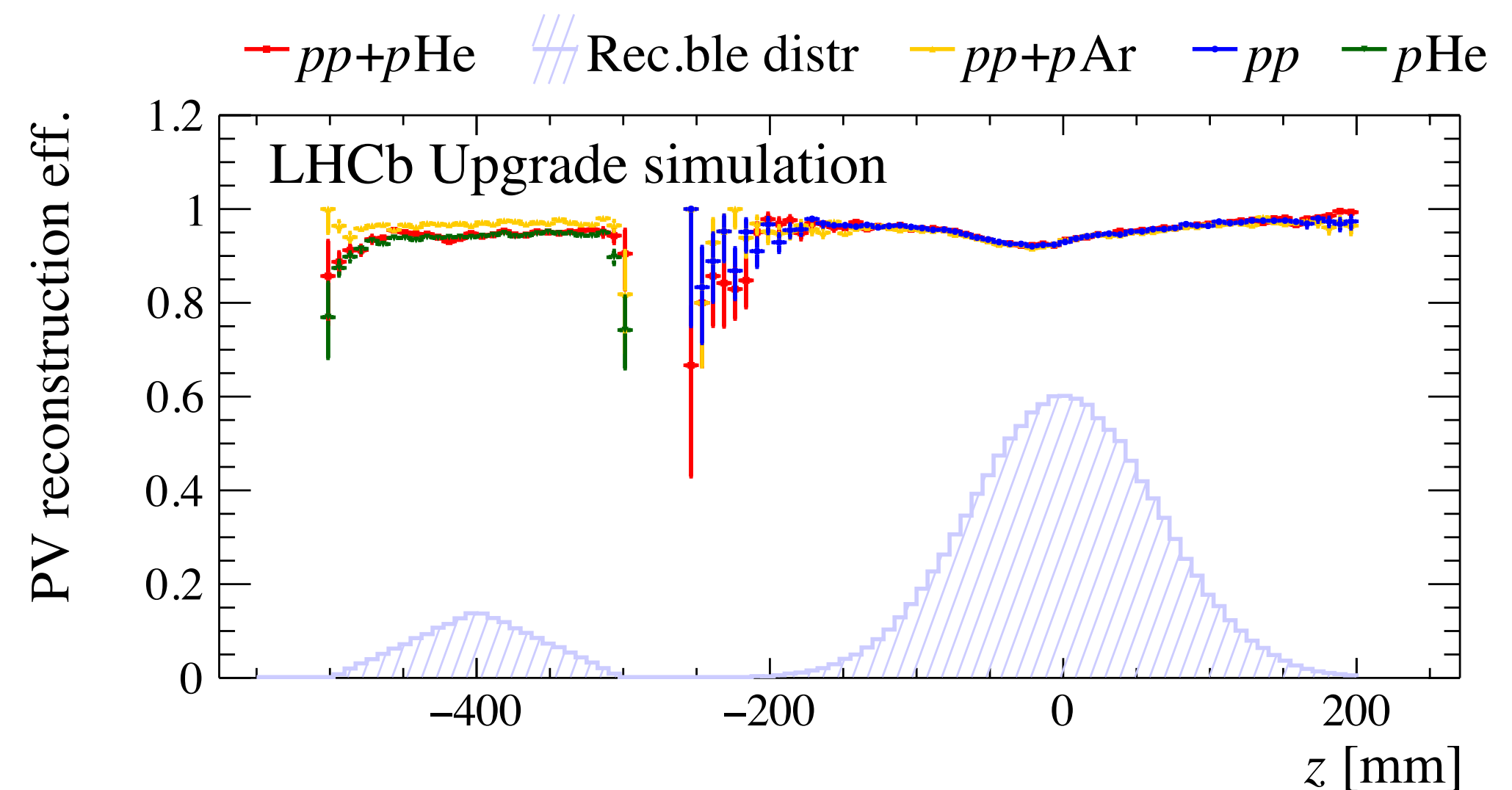
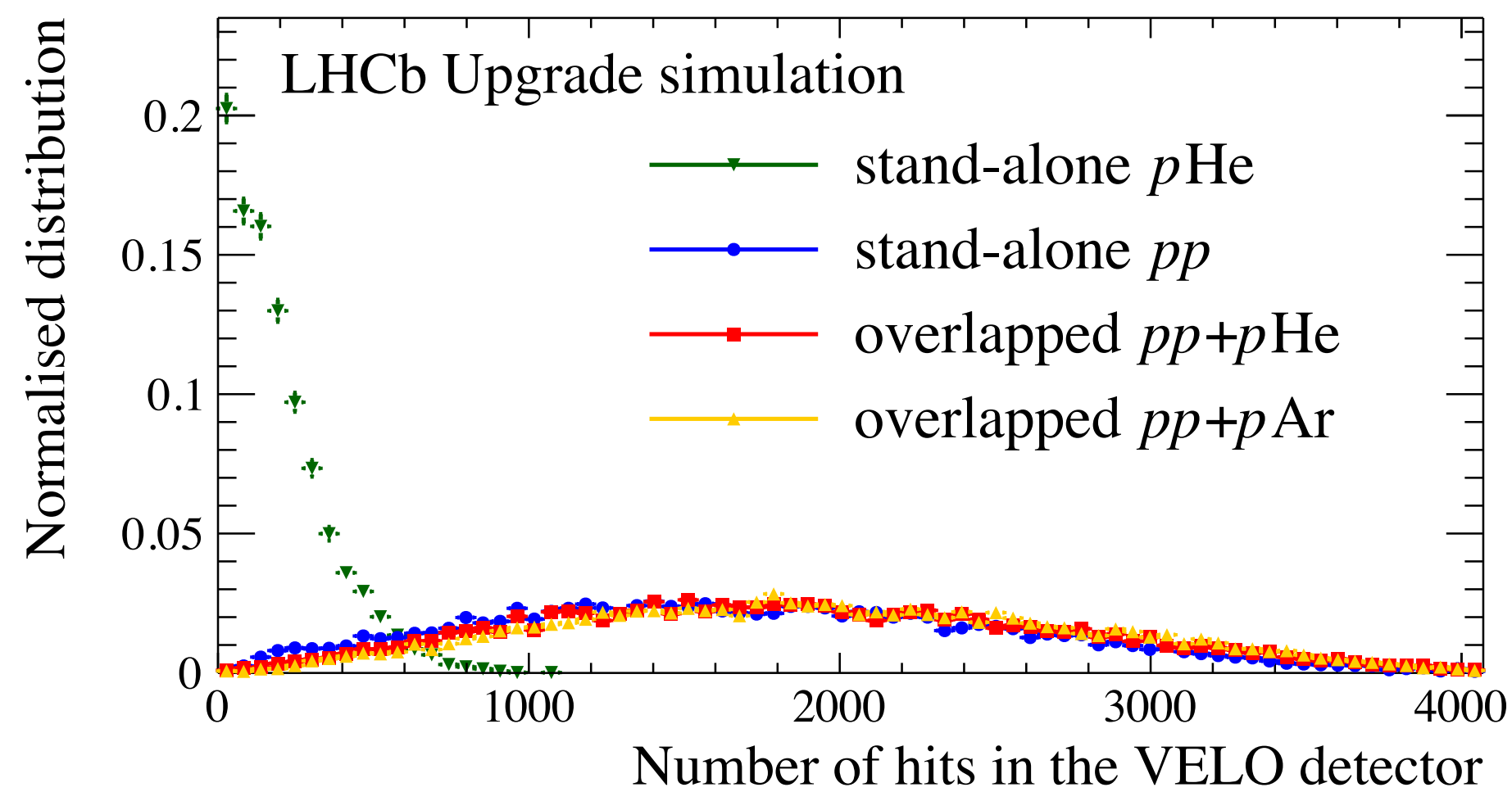
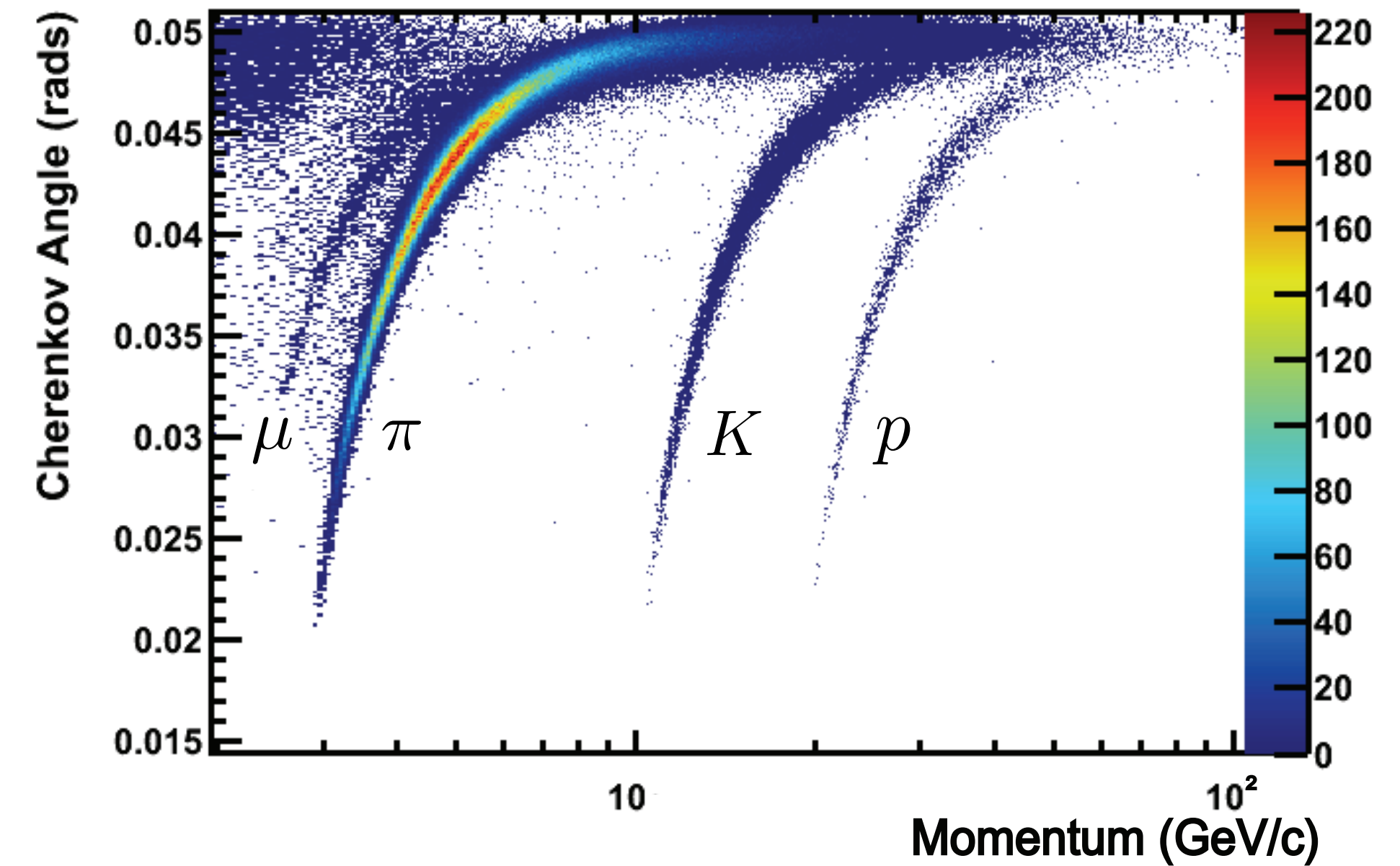
Large CM boost  $\rightarrow$  **large**  $x_2$  values

$$2 \leq y_{lab} \leq 5 \rightarrow -3.0 \leq y_{CMS} \leq 0$$



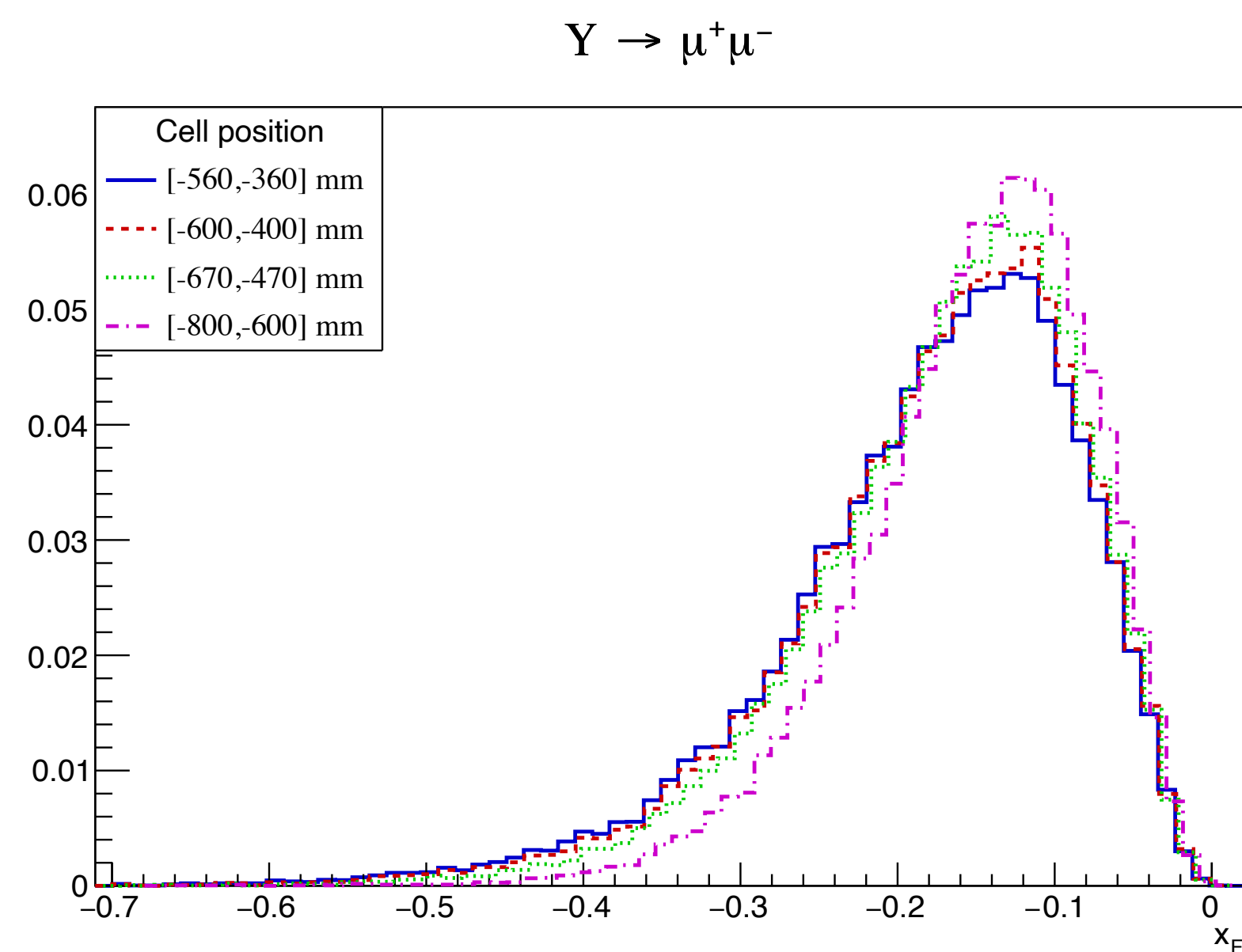
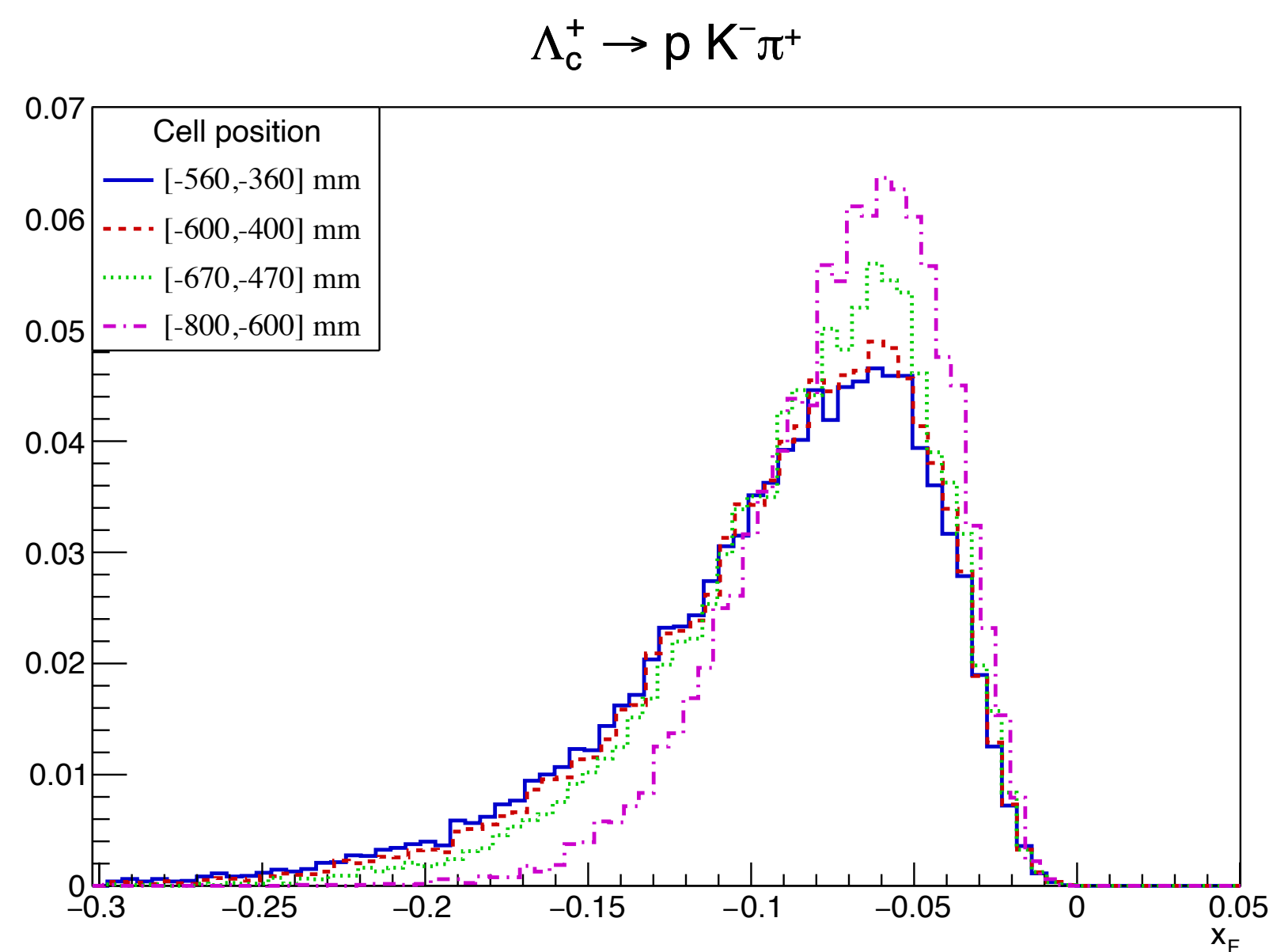
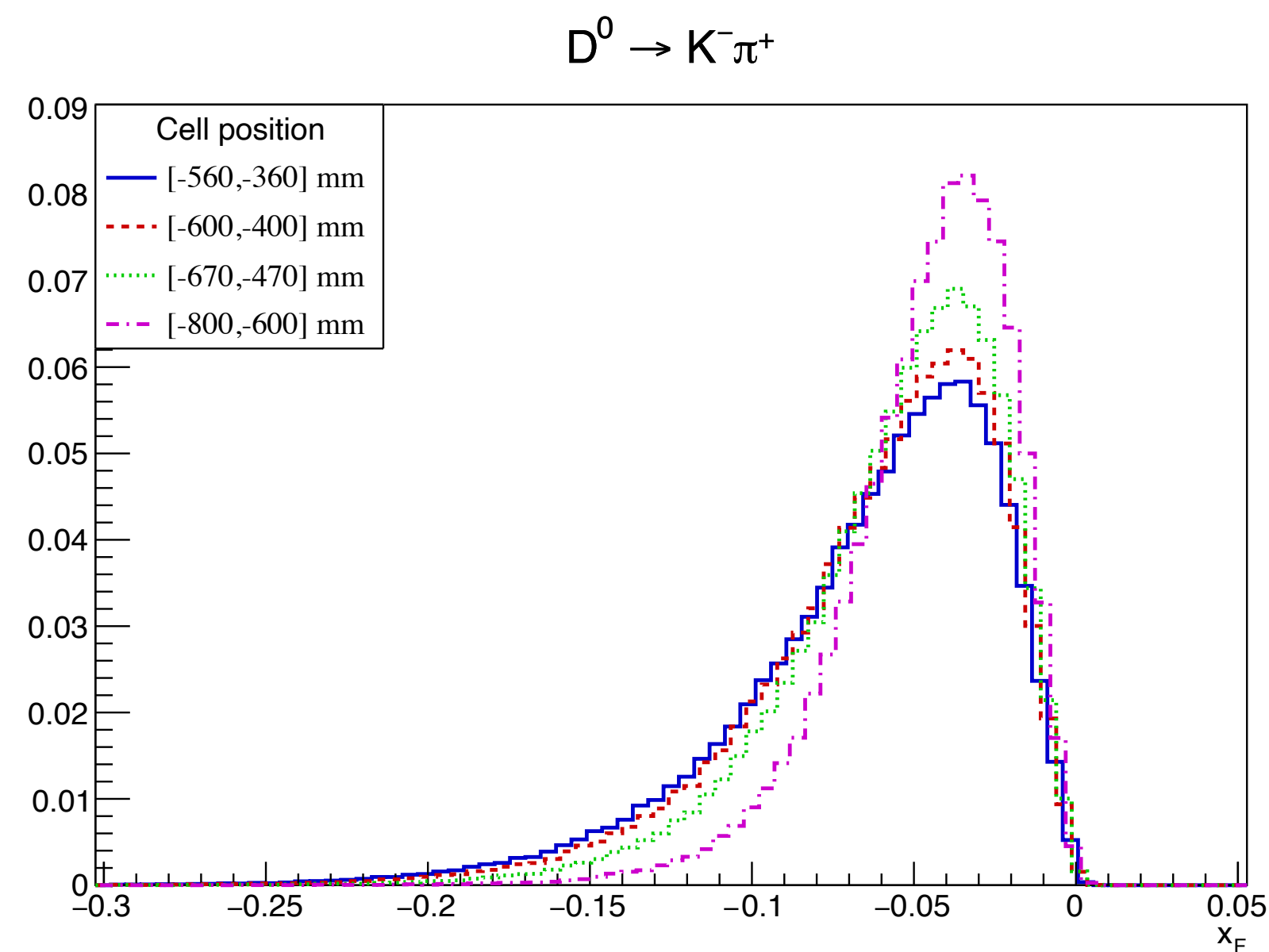
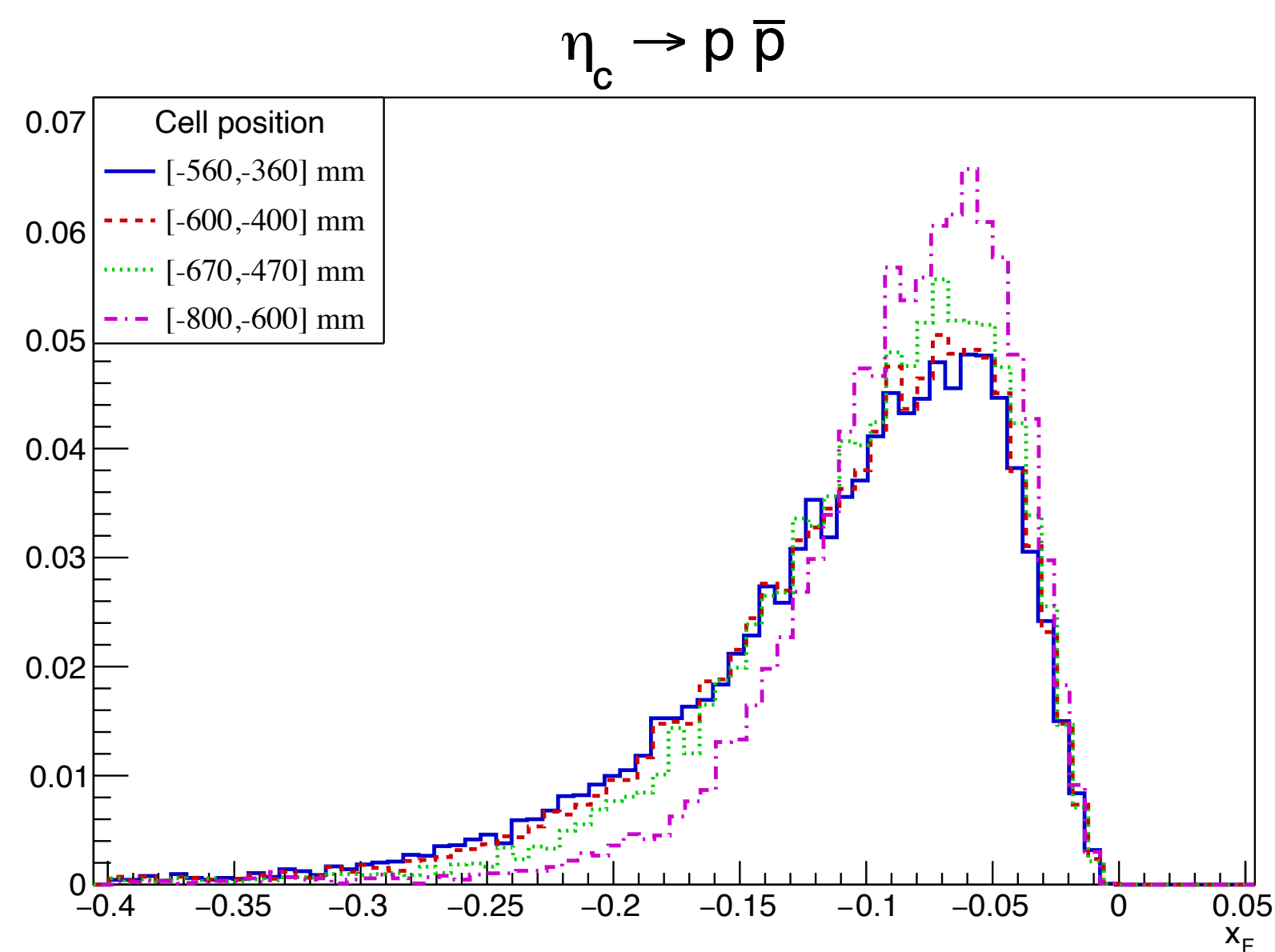
# LHCb capabilities for FT

- Strength points of LHCb:
  - Fully instrumented in  $2 < \eta < 5$  and optimised for b- and c-hadron detection
  - Excellent momentum resolution:  $\sigma_p/p = 0.5 - 1.0\%$  ( $p \in [2, 200]$  GeV)
  - Excellent particle identification with RICH+CALO+MUON, e.g.  $\epsilon_\mu \sim 98\%$  with  $\epsilon_{\pi \rightarrow \mu} \lesssim 1\%$ , hadron separation,  $\gamma/e$
  - Light hadrons ( $\pi, K, \dots$ ) are also abundantly produced and well reconstructed, but will not be considered in the following since are not unique to LHCspin
  - Can do simultaneous p-p and p-gas data-taking



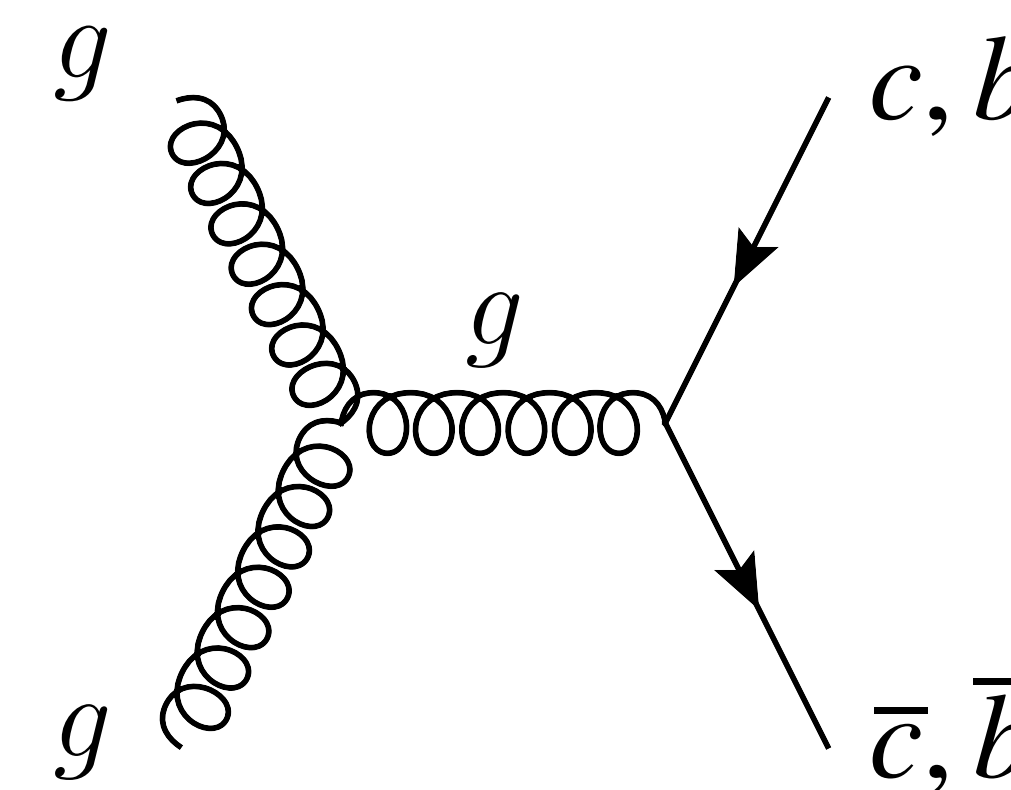
[LHCb-FIGURE-2022-002]

# Heavy flavour channels: examples



- c-hadrons will have the largest product of cross section and reconstruction efficiency
- Exotic probes and b-hadrons are also possible
- This is just a portion of the expected data (see later)

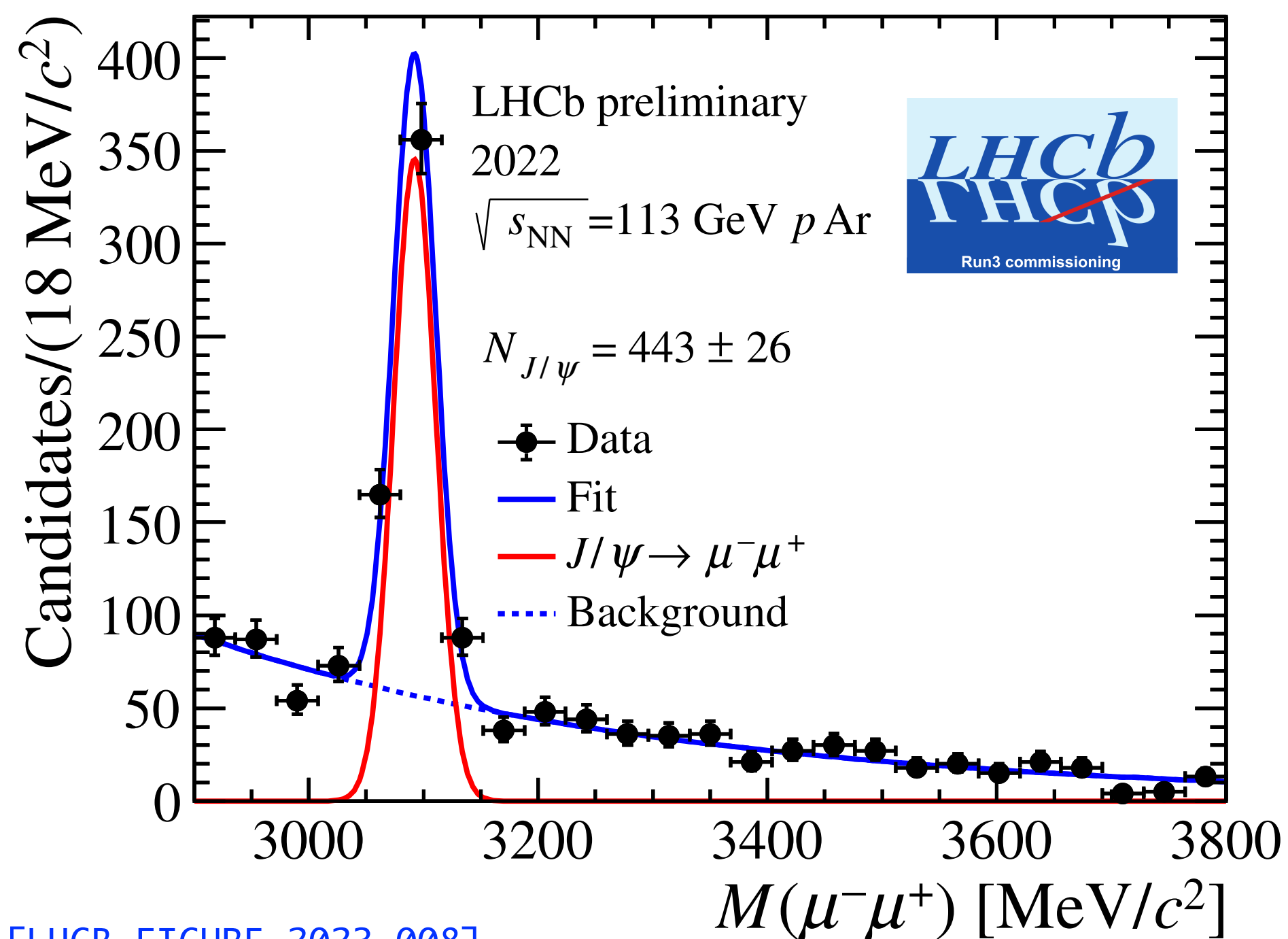
→ unique opportunity to probe gluon TMDs over a broad  $x$  range!





# Expected yields

- SMOG2 is performing above the expectation despite detectors are still under commissioning:



[LHCb-FIGURE-2023-008]

- Based on this milestone, we can project to  $\phi = 6.5 \times 10^{16}$  atoms/s (HERMES ABS) and SMOG2-like cell conductance  $\rightarrow \theta = 3.7 \times 10^{13}/\text{cm}^2$  for polarised hydrogen

- Assuming 84 hours of data-taking per week and 120-week Run... we can collect a huge amount of data!

Channel	Events / week	Total yield
$J/\psi \rightarrow \mu^+ \mu^-$	$1.3 \times 10^7$ !!	$1.5 \times 10^9$
$D^0 \rightarrow K^- \pi^+$	$6.5 \times 10^7$	$7.8 \times 10^9$
$\psi(2S) \rightarrow \mu^+ \mu^-$	$2.3 \times 10^5$	$2.8 \times 10^7$
$J/\psi J/\psi \rightarrow \mu^+ \mu^- \mu^+ \mu^-$ (DPS)	8.5	$1.0 \times 10^3$
$J/\psi J/\psi \rightarrow \mu^+ \mu^- \mu^+ \mu^-$ (SPS)	$2.5 \times 10^1$	$3.1 \times 10^3$
Drell Yan ( $5 < M_{\mu\mu} < 9$ GeV)	$7.4 \times 10^3$	$8.8 \times 10^5$
$\Upsilon \rightarrow \mu^+ \mu^-$	$5.6 \times 10^3$	$6.7 \times 10^5$
$\Lambda_c^+ \rightarrow p K^- \pi^+$	$1.3 \times 10^6$	$1.5 \times 10^8$

- This is the number of fully-reconstructed and selected events based on our current capabilities
  - Event rate enhanced during HL-LHC (Upgrade II)
  - Fully-software trigger allows further improvements

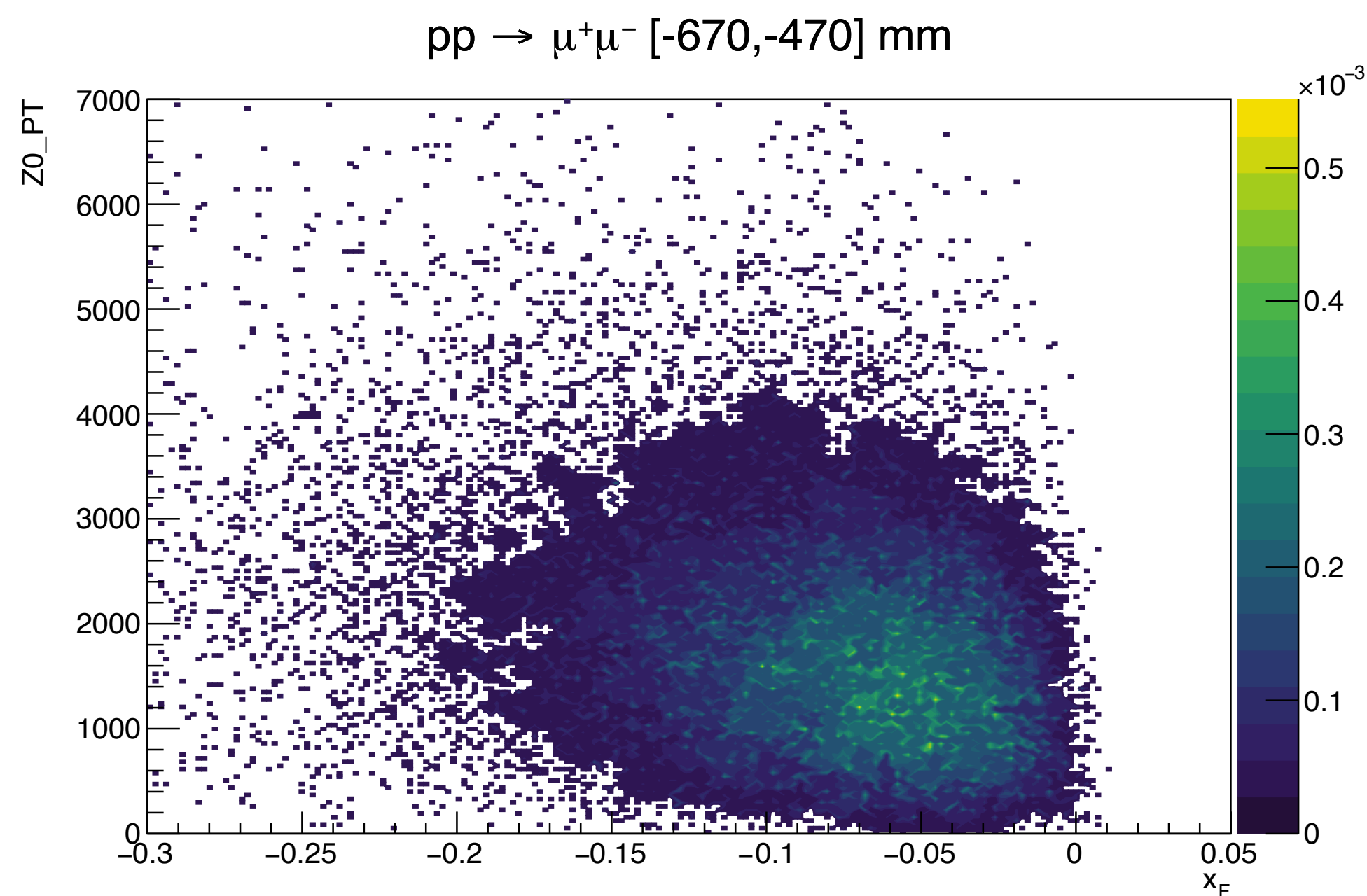


# Polarised Drell-Yan

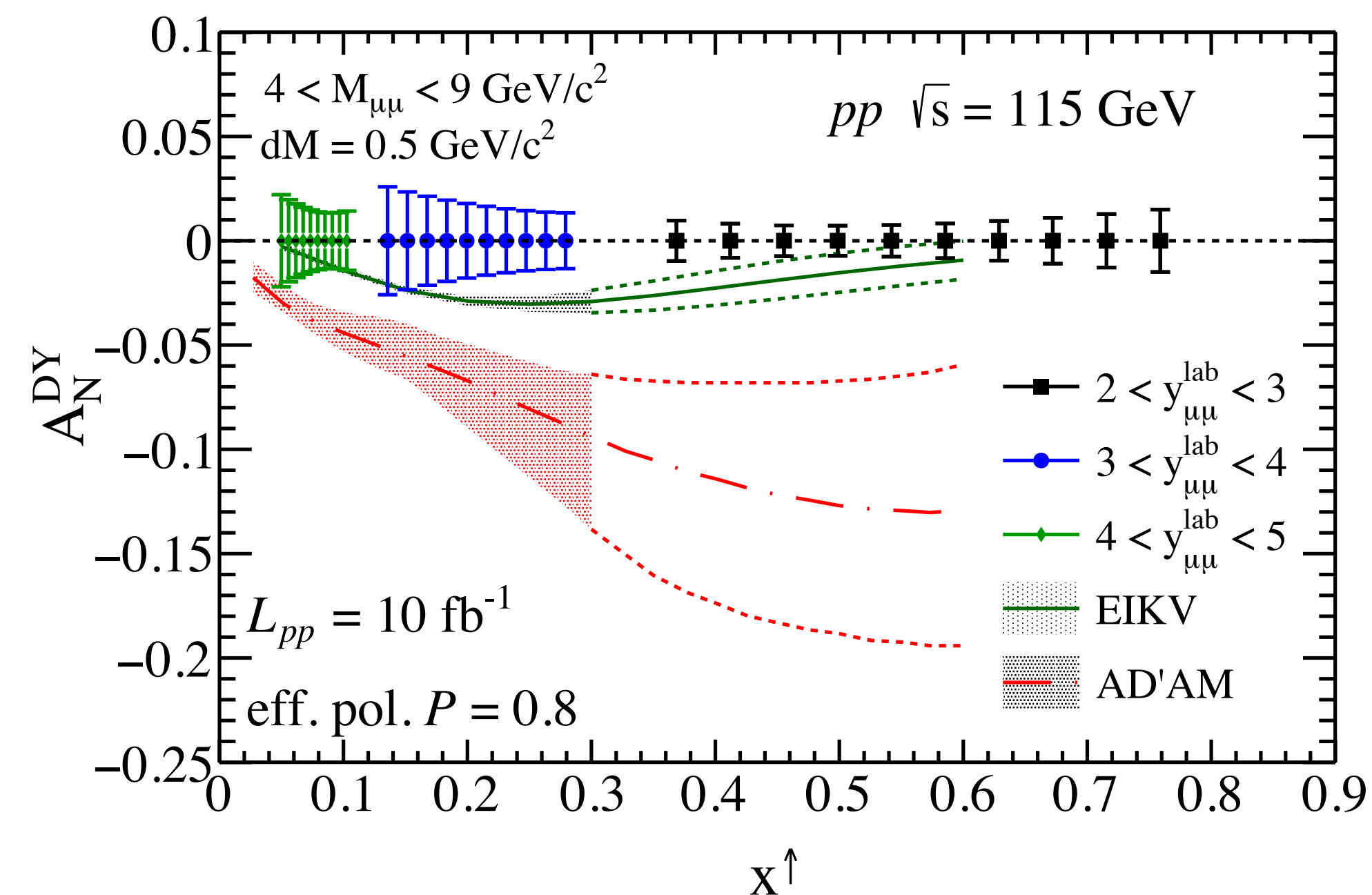
- To access the transverse motion of partons inside a polarised nucleon: measure TMDs via **TSSAs** at high  $x_2^\uparrow$  (and low  $x_1$ )

$$A_N = \frac{1}{P} \frac{\sigma^\uparrow - \sigma^\downarrow}{\sigma^\uparrow + \sigma^\downarrow} \longrightarrow A_N \sim \frac{f_1^q(x_1, k_{T1}^2) \otimes f_{1T}^{\perp\bar{q}}(x_2, k_{T2}^2)}{f_1^q(x_1, k_{T1}^2) \otimes f_1^q(x_2, k_{T2}^2)}$$

- Kinematics (~30k events):



- Projections of polarised DY with  $10 \text{ fb}^{-1}$  of data from [\[ArXiv:1807.00603\]](https://arxiv.org/abs/1807.00603) :



- Precise measurements but also unique features:
  - Verify the **sign change** of the Sivers TMD in DY wrt SIDIS:
 
$$f_{1T}^{\perp q}(x, k_T^2)_{\text{DY}} = -f_{1T}^{\perp q}(x, k_T^2)_{\text{SIDIS}}$$
  - + isospin effect with polarised deuterium



# More TMDs

- Can probe TMDs via azimuthal asymmetries of the dilepton pair:
- $h_q^1$  : transversity → difference in densities of quarks having T pol. ↑↑ or ↑↓ in T pol. nucleon
- $f_{1T}^{\perp q}$  : Sivers → dependence on  $p_T$  orientation wrt T pol. nucleon
- $h_1^{\perp q}$  : Boer-Mulders → dependence on  $p_T$  orientation wrt T pol. quark in unp. nucleon
- $h_{1T}^{\perp q}$  : pretzelosity → dependence on  $p_T$  and T. pol of both T pol. quark and nucleon
- $f_1^q$  : unpolarised TMD, always present at the denominator

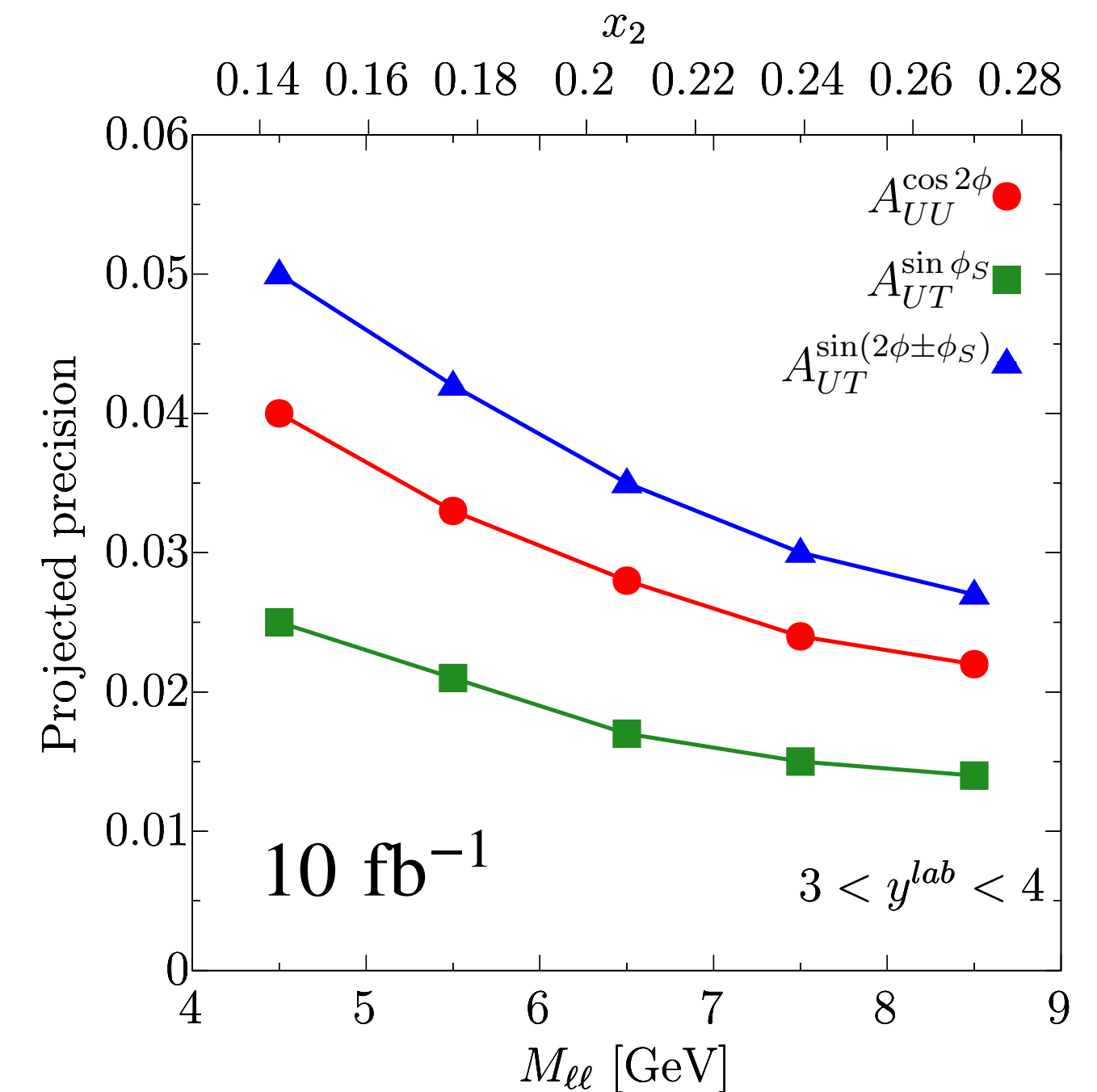
$$A_{UU}^{\cos 2\phi} \sim \frac{h_1^{\perp q}(x_1, k_{1T}^2) \otimes h_1^{\perp \bar{q}}(x_2, k_{2T}^2)}{f_1^q(x_1, k_{1T}^2) \otimes f_1^{\bar{q}}(x_2, k_{2T}^2)}$$

$$A_{UT}^{\sin \phi_S} \sim \frac{f_1^q(x_1, k_{1T}^2) \otimes f_{1T}^{\perp \bar{q}}(x_2, k_{2T}^2)}{f_1^q(x_1, k_{1T}^2) \otimes f_1^{\bar{q}}(x_2, k_{2T}^2)}$$

$$A_{UT}^{\sin(2\phi + \phi_S)} \sim \frac{h_1^{\perp q}(x_1, k_{1T}^2) \otimes h_{1T}^{\perp \bar{q}}(x_2, k_{2T}^2)}{f_1^q(x_1, k_{1T}^2) \otimes f_1^{\bar{q}}(x_2, k_{2T}^2)}$$

$$A_{UT}^{\sin(2\phi - \phi_S)} \sim \frac{h_1^{\perp q}(x_1, k_{1T}^2) \otimes h_1^{\bar{q}}(x_2, k_{2T}^2)}{f_1^q(x_1, k_{1T}^2) \otimes f_1^{\bar{q}}(x_2, k_{2T}^2)}$$

- **Polarised Drell-Yan** to access unpolarised TMDs of sea quarks and polarised TMDs in the valence region
- A good precision can be attained with ~1M events and 20-40 MeV mass resolution
- Also, dedicated trigger lines can be implemented e.g. tight muonID and relaxed p cuts, electron lines...



[ArXiv:1807.00603]



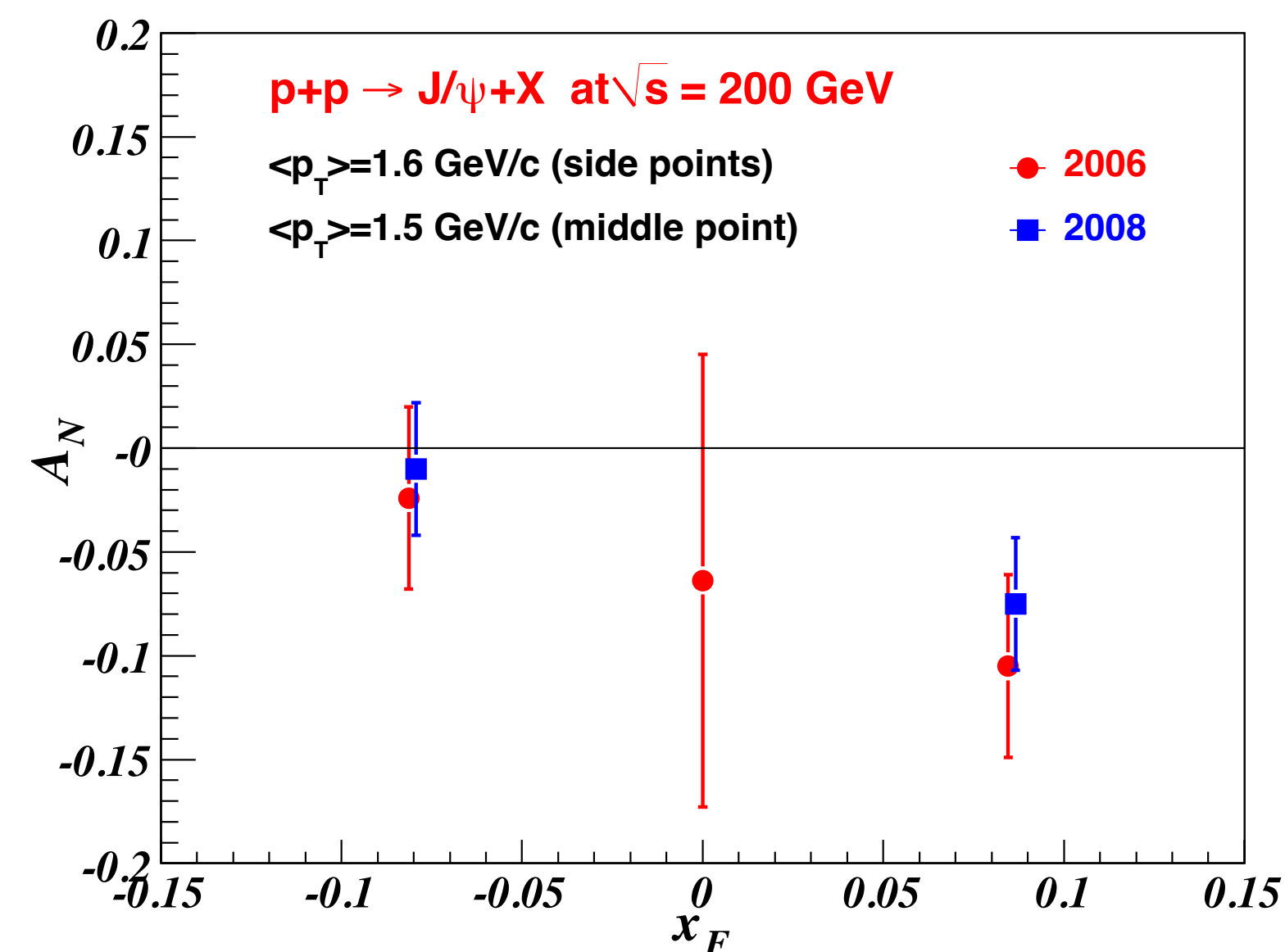
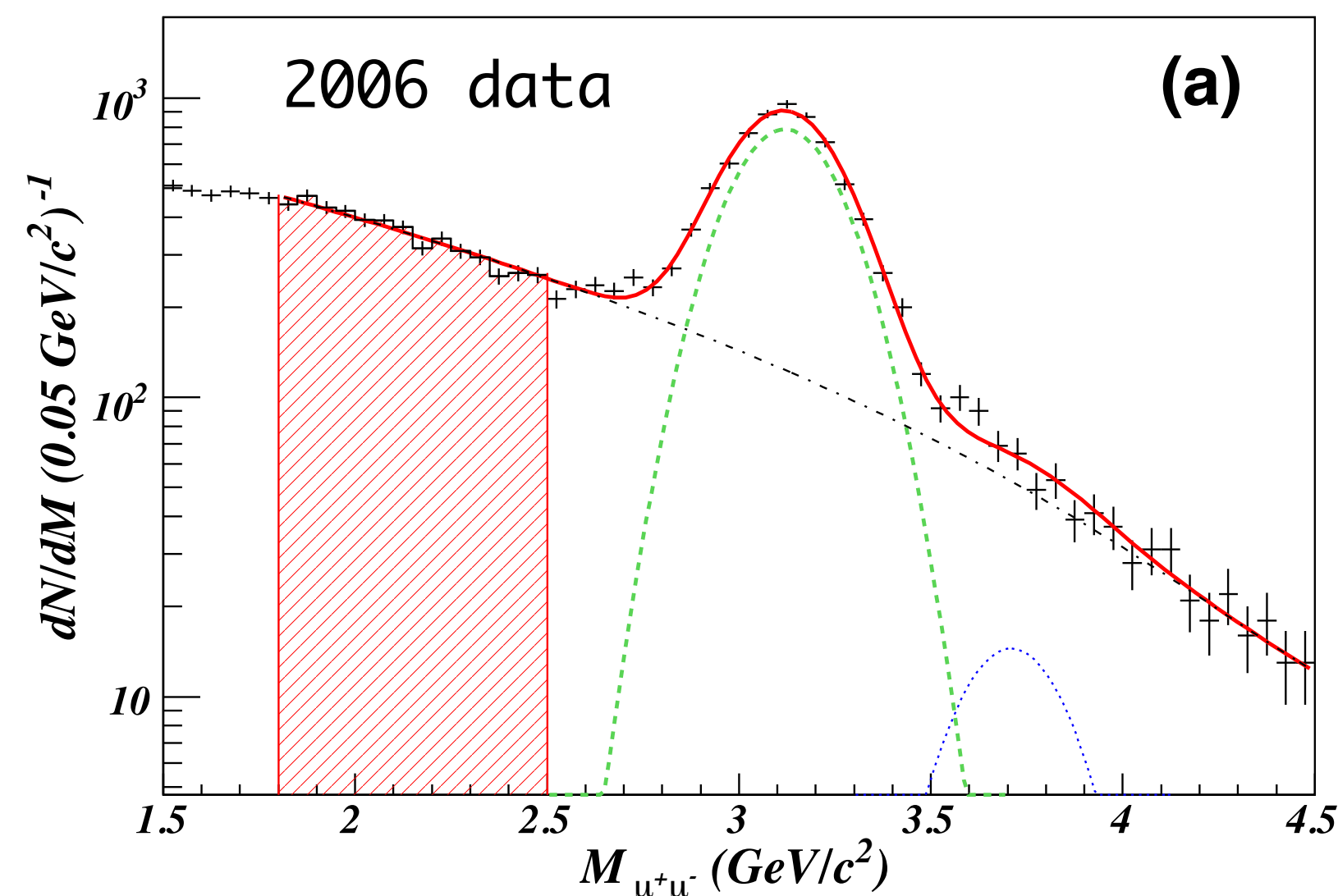
# Heavy-flavour and comparison with PHENIX

- LHCspin strength point and uniqueness will be heavy flavours, mostly unexplored by existing facilities
- The (only?) exception is  $J/\psi$ , for which measurements have been performed at PHENIX in polarised pp at  $\sqrt{s} = 200$  GeV. However, compared to LHCspin:
  - $\sim 21$ k signal candidates (2006 + 2008 data)  $\rightarrow$  can be collected in  $\sim 10$  minutes (cell) or  $\sim 7$  hours (jet)
  - Mass resolution  $\sim 150$  MeV  $\rightarrow \sigma_{\mu\mu} \approx 13$  MeV at the  $J/\psi$  mass and  $\sigma_{\mu\mu} \approx 42$  MeV at the  $\Upsilon$  mass
  - $\psi(2S)$  barely visible  $\rightarrow$  can measure excited states
- $\rightarrow$  we can greatly enrich these results with high precision measurements and much larger kinematic coverage!

- Also, with a few thousands of di- $J/\psi$  events we could measure the gluon Sivers  $f_{1T}^{\perp g}$ , transversity  $h_{1T}^g$  and pretzelosity  $h_{1T}^{\perp g}$
- Unique in FT! Challenging, but specific lines can be developed already in the Run 3

[PLB 784 (2018) 217-222]

[J. Bor @ DESY 2023]



[PRD 82 (2010) 112008]



# Expected precision on $A_N$

- Expected uncertainty on a TSSA at LHCspin:

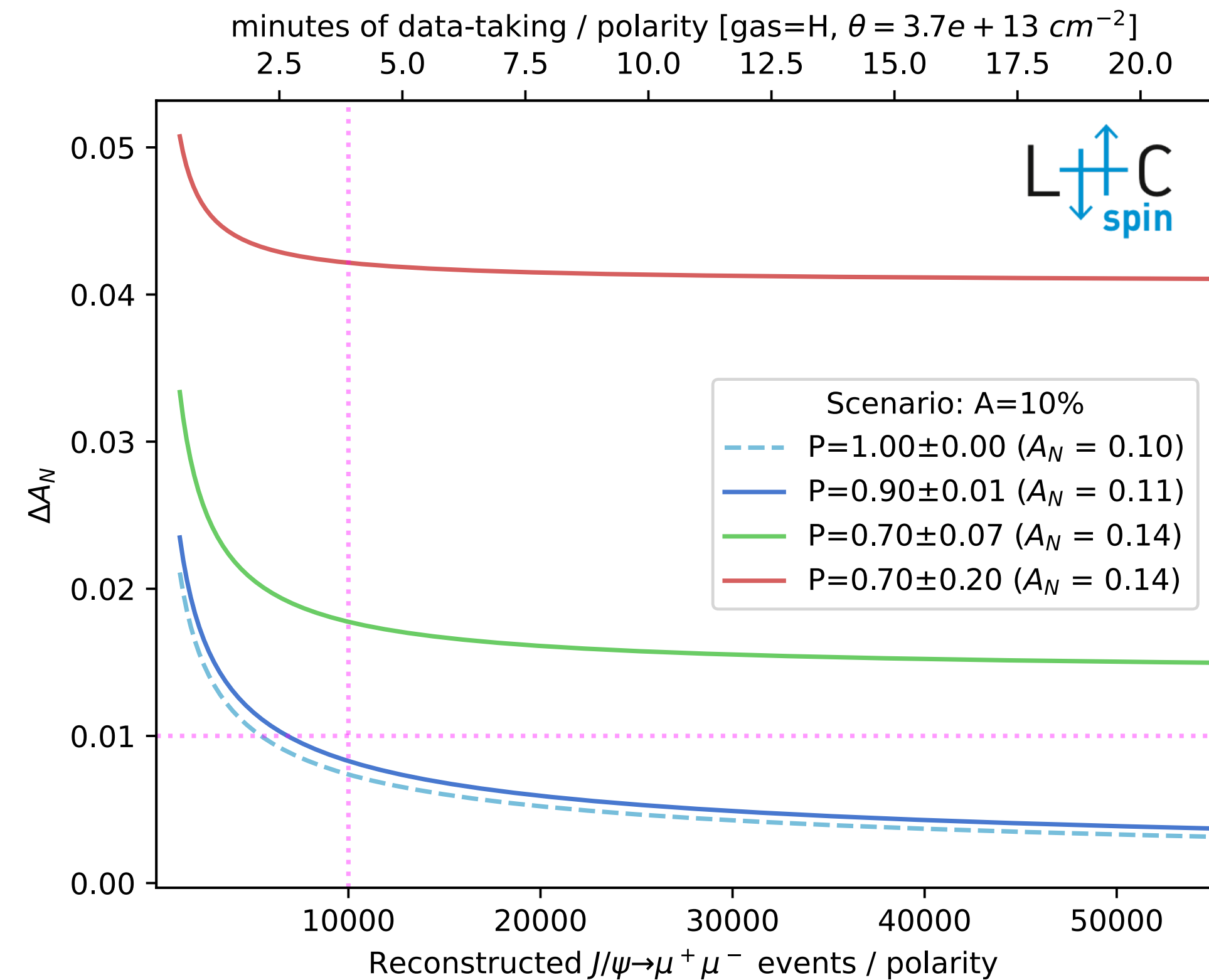
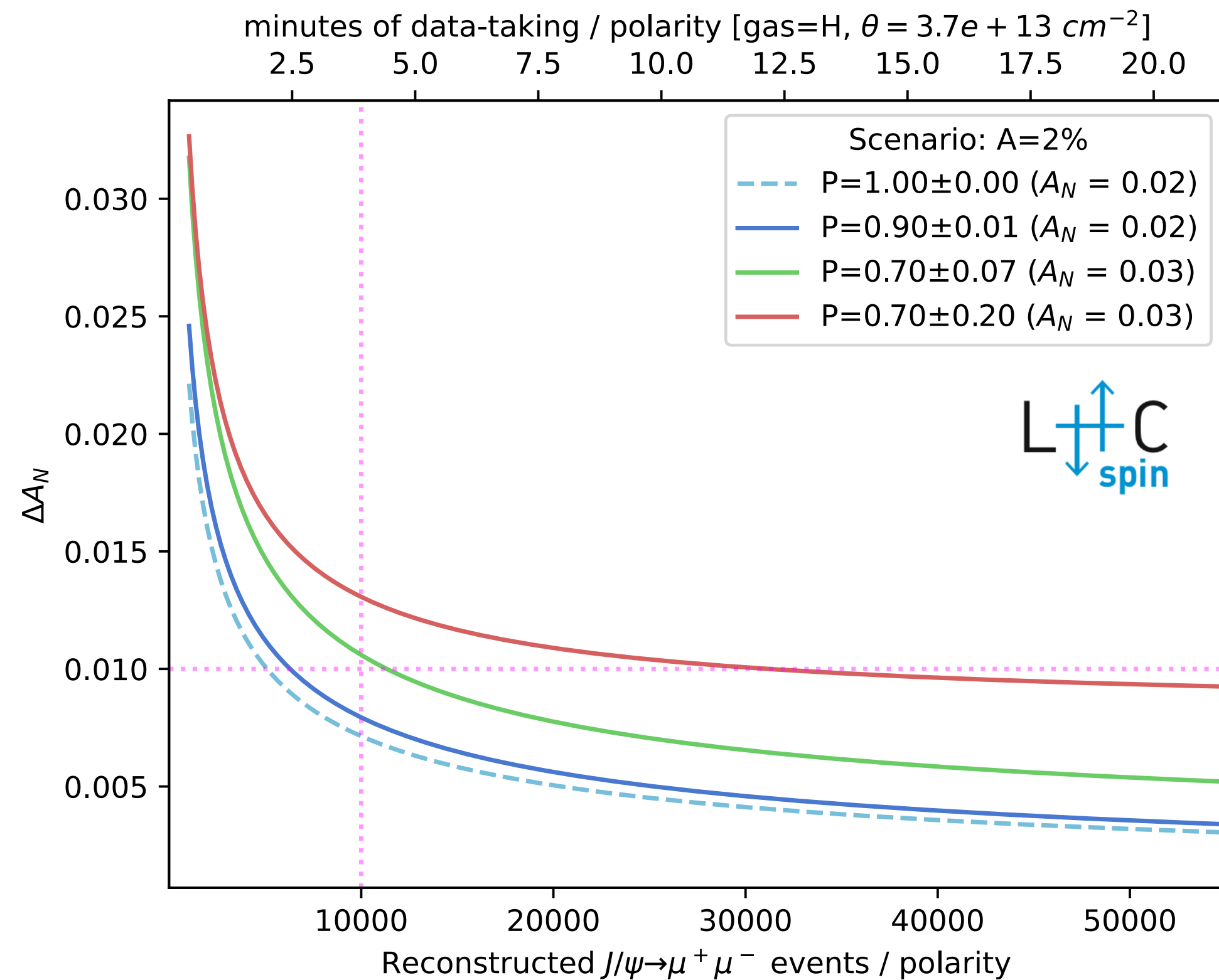
$$A_N = \frac{1}{P} \frac{N^\uparrow - N^\downarrow}{N^\uparrow + N^\downarrow} \rightarrow \Delta A \approx \frac{1}{\sqrt{2N^\uparrow}}$$

- $\Delta A_N$  showed for different polarisation degrees on two scenarios: small asymmetry  $A = 2\%$  (left) and large asymmetry  $A = 10\%$  (right)

- Systematic limit from P reached after few minutes for  $J/\psi \rightarrow \mu^+ \mu^-$ : precision TSSA measurements possible with very short  $pH^\uparrow$  runs!

- Cell target example:  $P = 0.70 \pm 0.07$ ,  $\theta = 3.7 \times 10^{13}/\text{cm}^2$  (used in the plots)

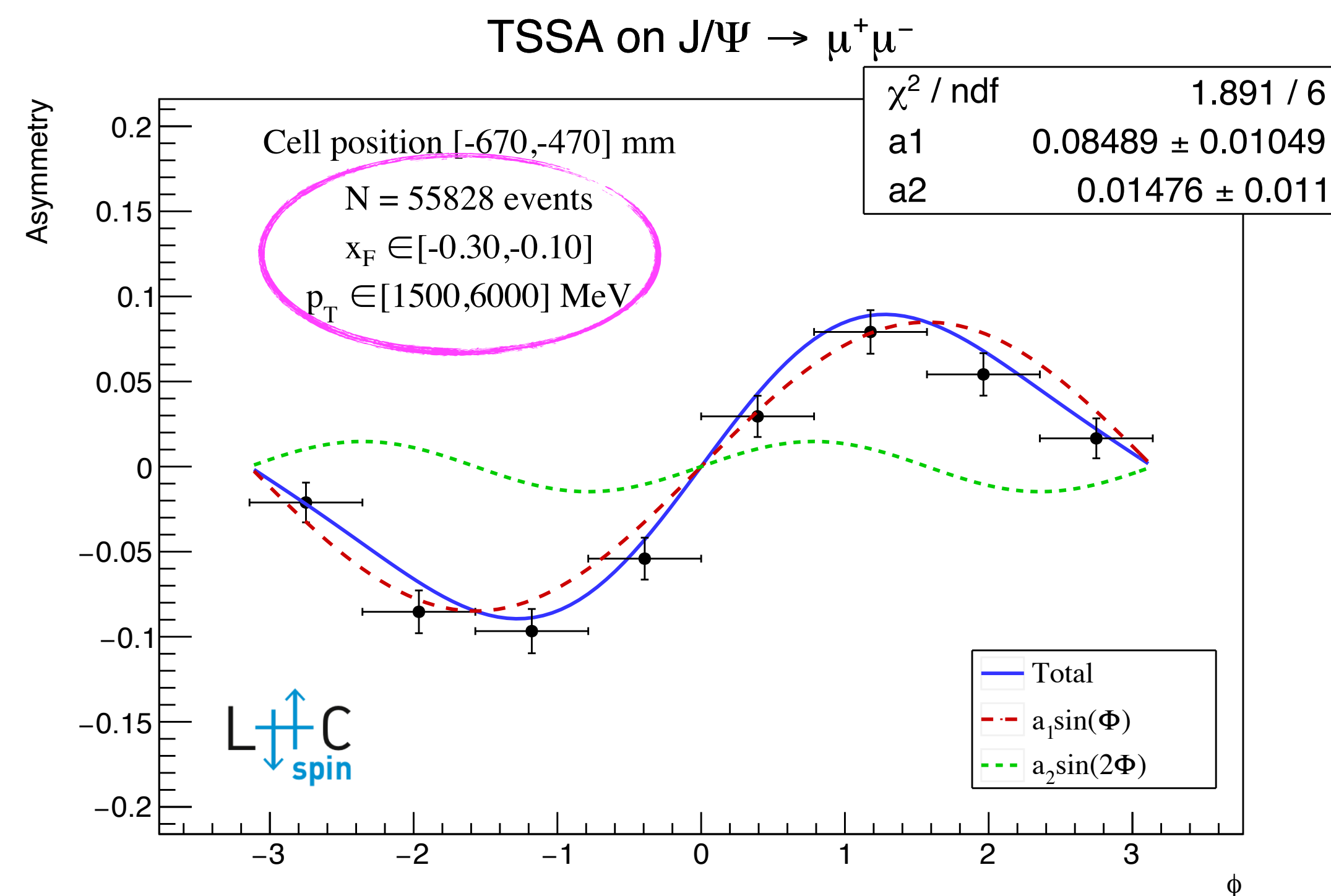
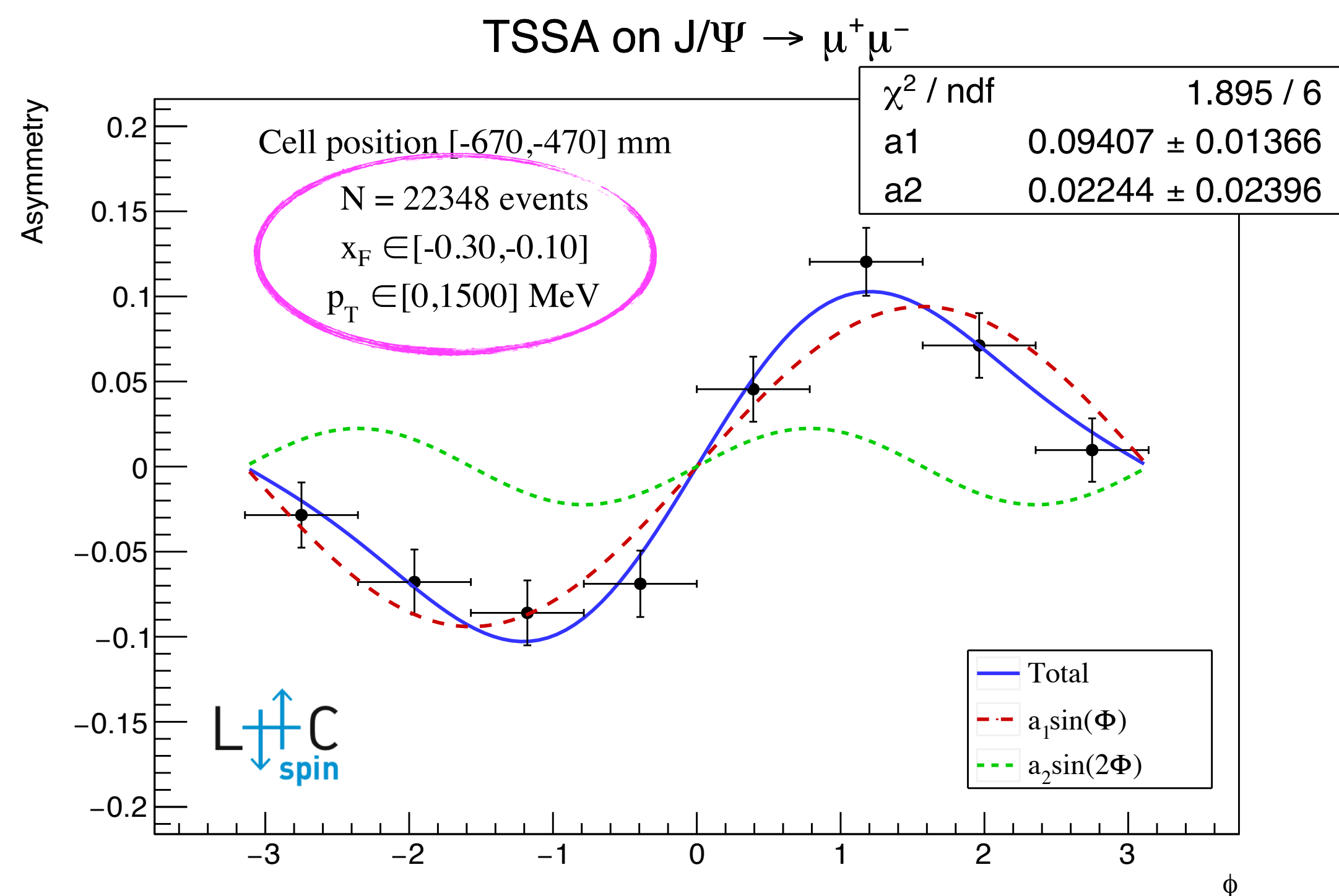
- Jet target example:  $P = 0.90 \pm 0.01$ ,  $\theta \approx 10^{12}/\text{cm}^2$



- Our MC is unpolarised. The polarisation can be emulated by assigning  $\pm 1$  state according to

$$\rho = \frac{1}{2} \left[ 1 + \left( a_1 + a_2 \frac{x - \bar{x}}{x_{max}} + a_3 \frac{p_T - \bar{p}_T}{p_{T\ max}} \right) \sin \phi + \left( b_1 + b_2 \frac{x - \bar{x}}{x_{max}} + b_3 \frac{p_T - \bar{p}_T}{p_{T\ max}} \right) \sin 2\phi \right]$$

- The resulting pseudo-data are fitted, in this examples with two amplitudes in  $4 x_F \times 2 p_T \times 8 \phi$  bins using  $A_N \sim 0.1 \pm 0.01$  and  $\Delta P = 5\%$  :

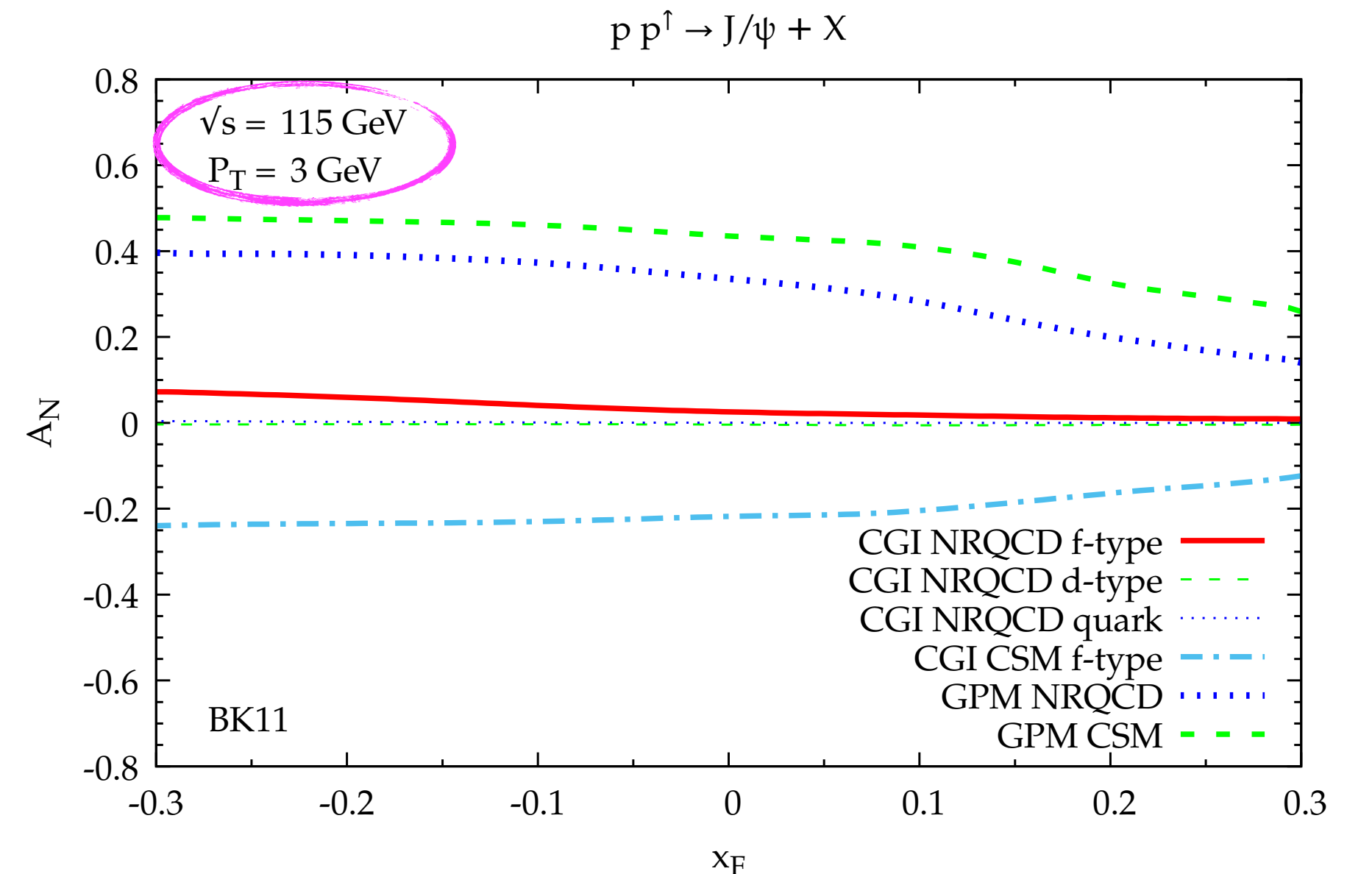
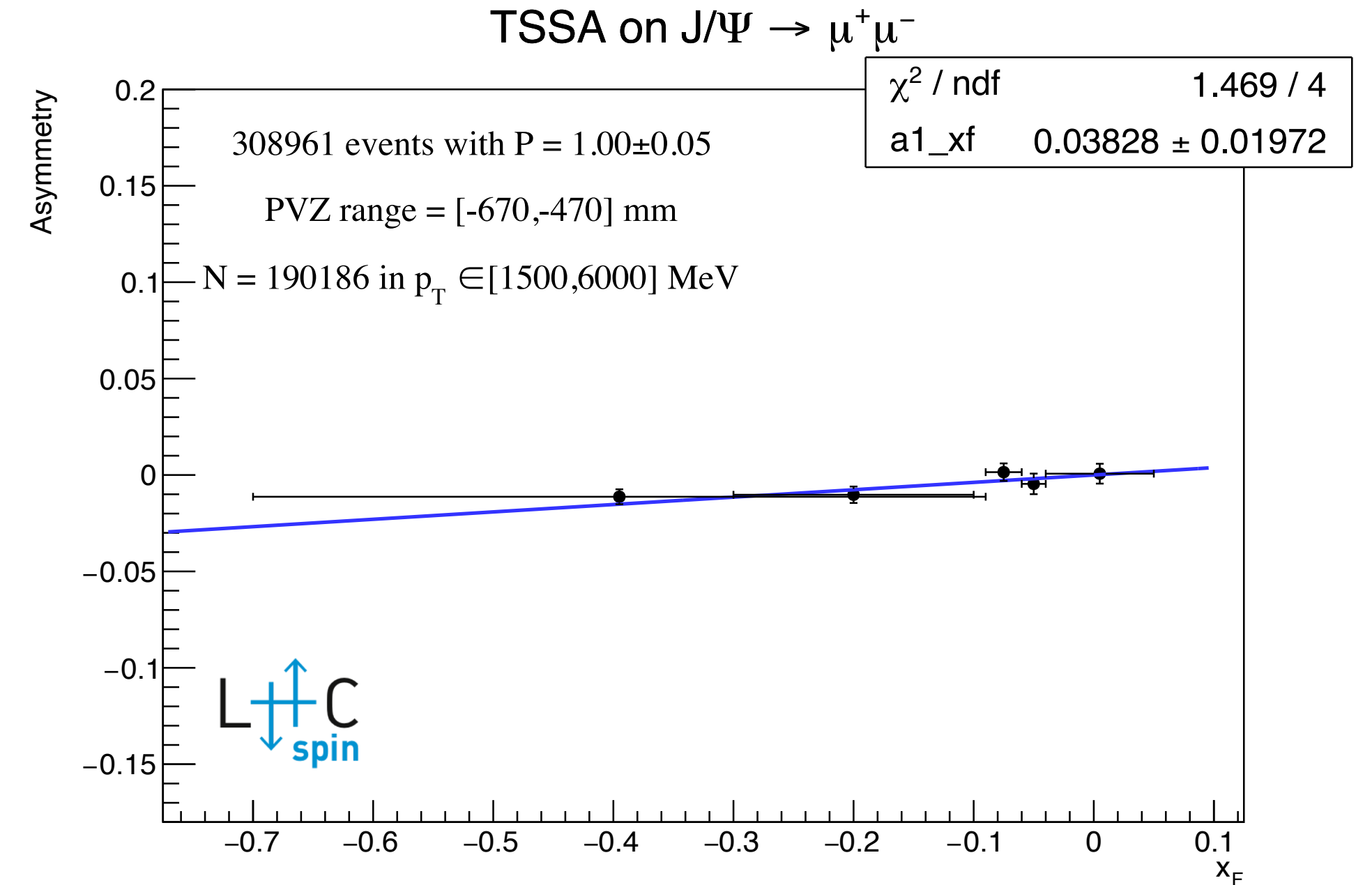
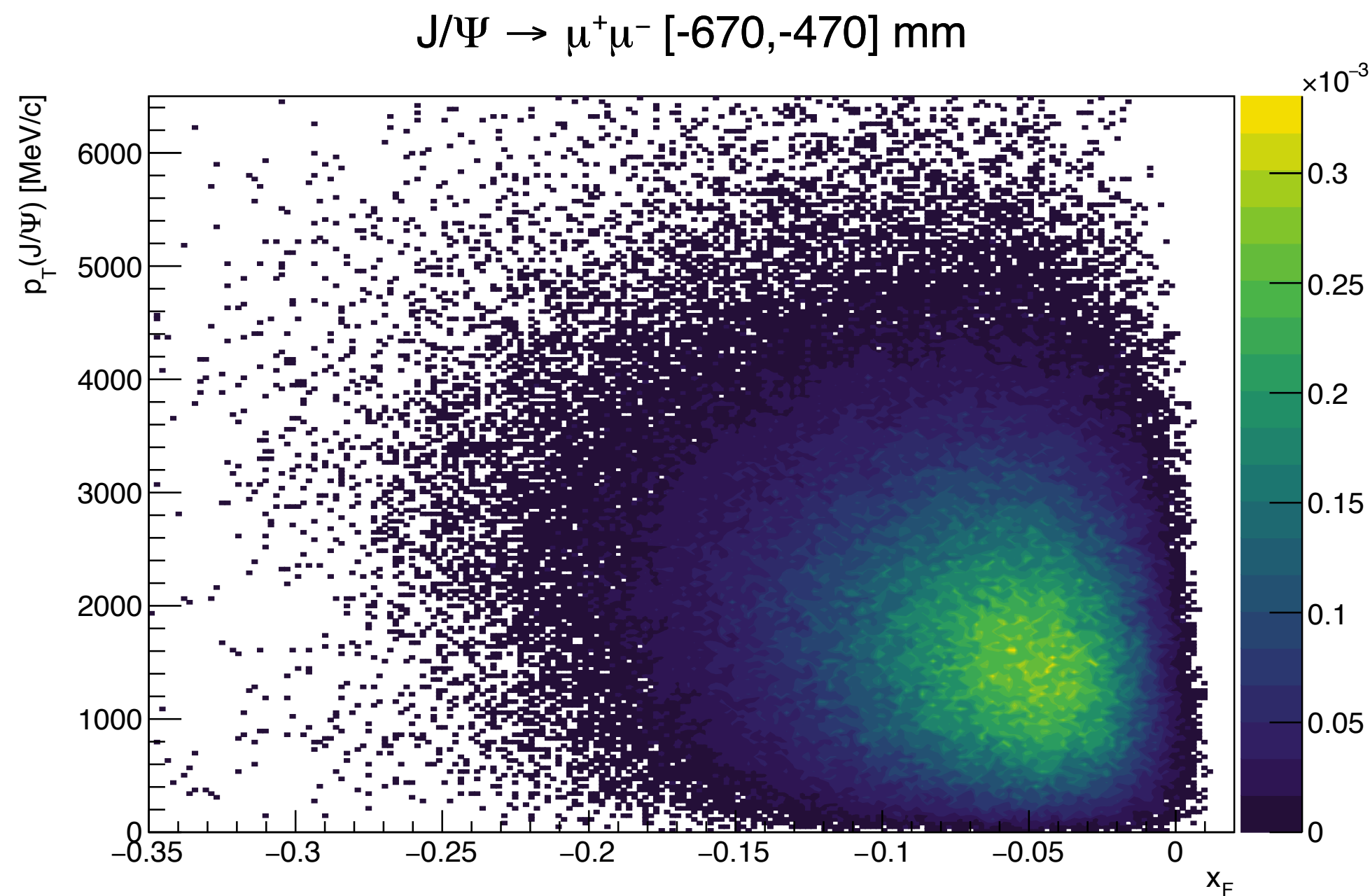


- Just a starting point: can refine & test any model of choice!



# An example measurement: the Gluon Sivers function

- A simpler approach for large asymmetry:  $x_F$  and  $p_T$  dependence of  $A_N$  (here  $\sim 0$ )  $\rightarrow$  this precision can be achieved with just few hours of data-taking!
- We can provide precision measurements well able to distinguish among theoretical predictions of the GSF
- Kinematic coverage given below for  $\sim 300k$  events (small sample)
- Simulations with the full LHCb detector shown  $\rightarrow$  to be performed also for the IR4 setup

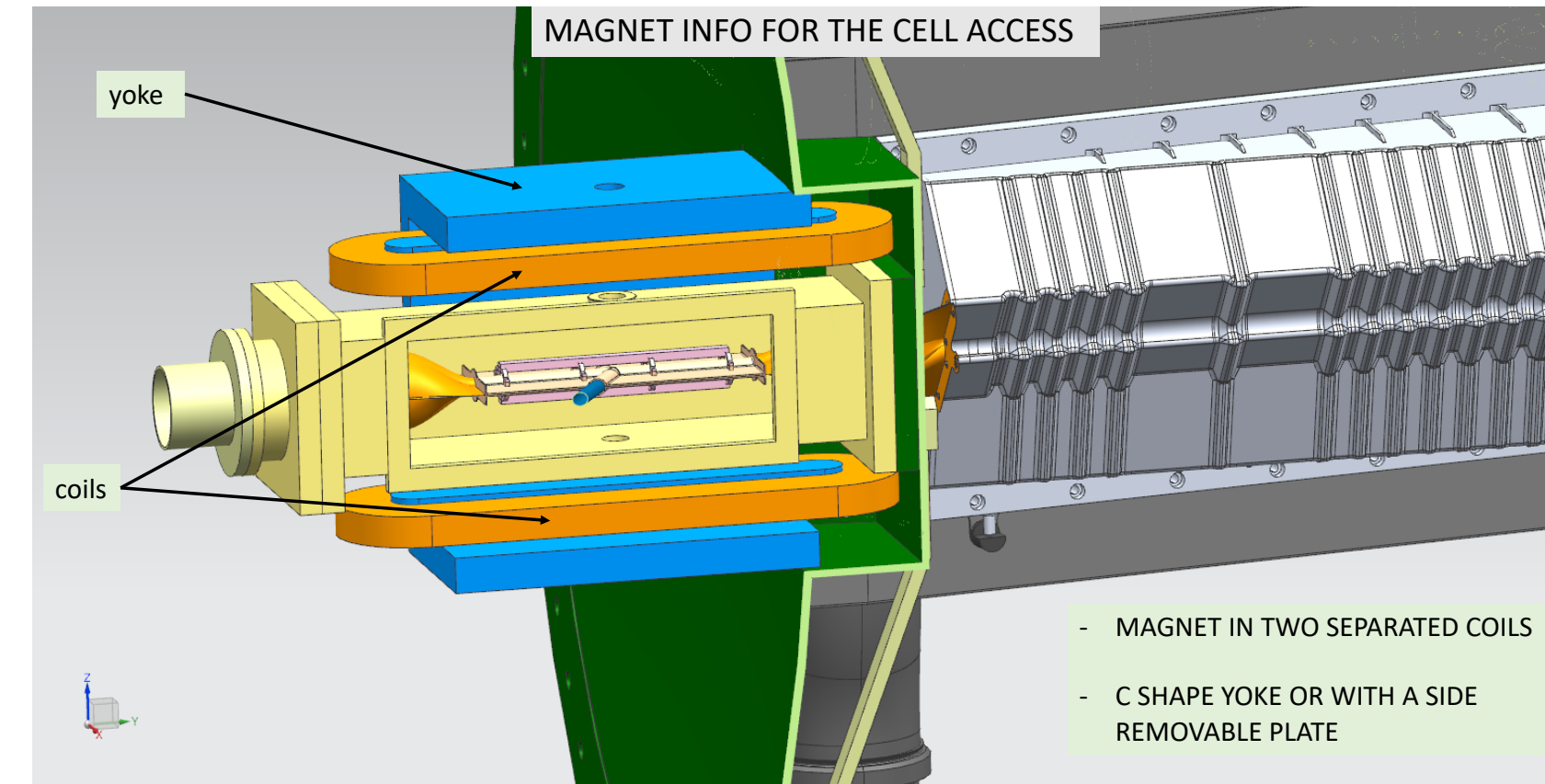


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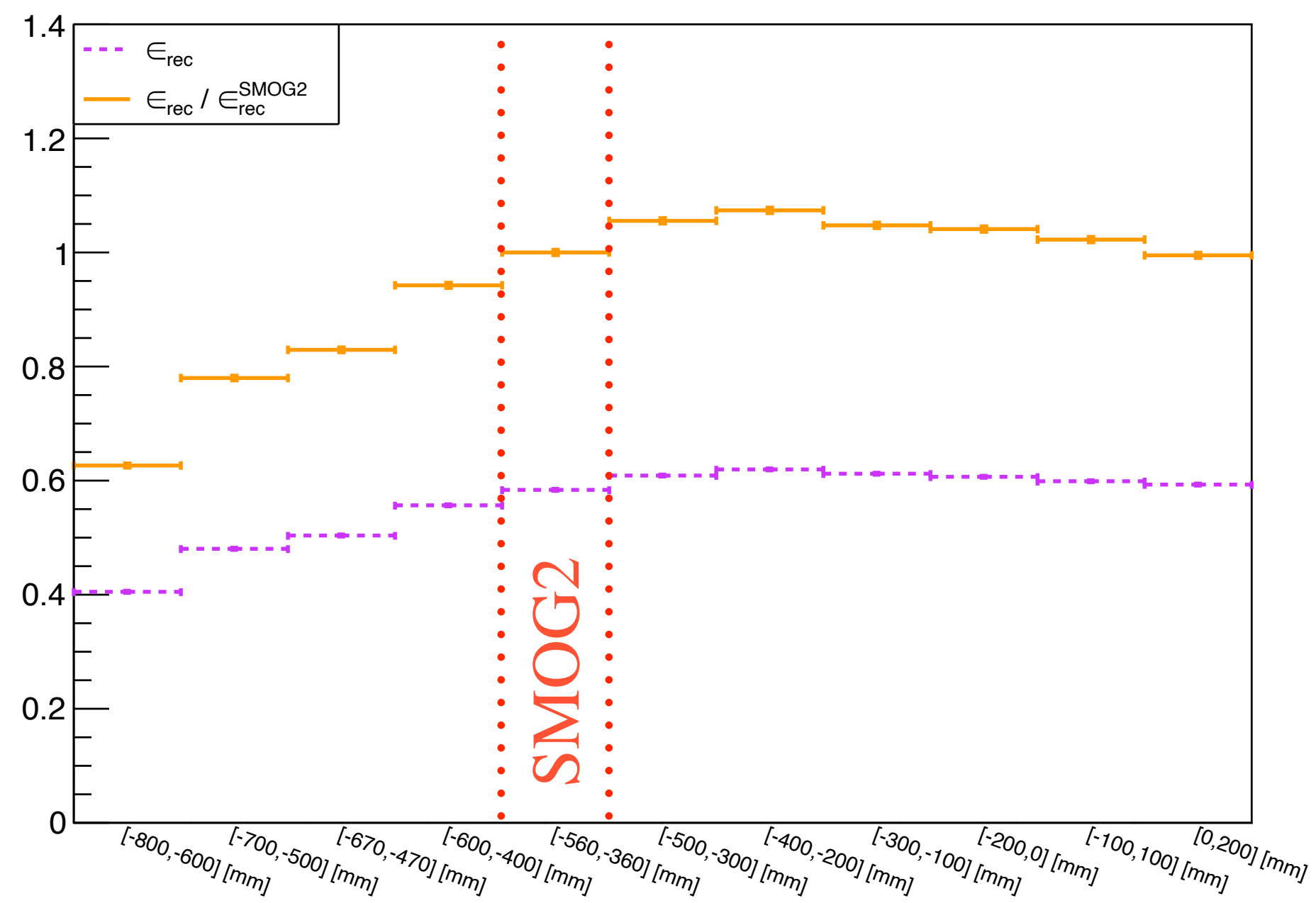


# Reconstruction efficiency

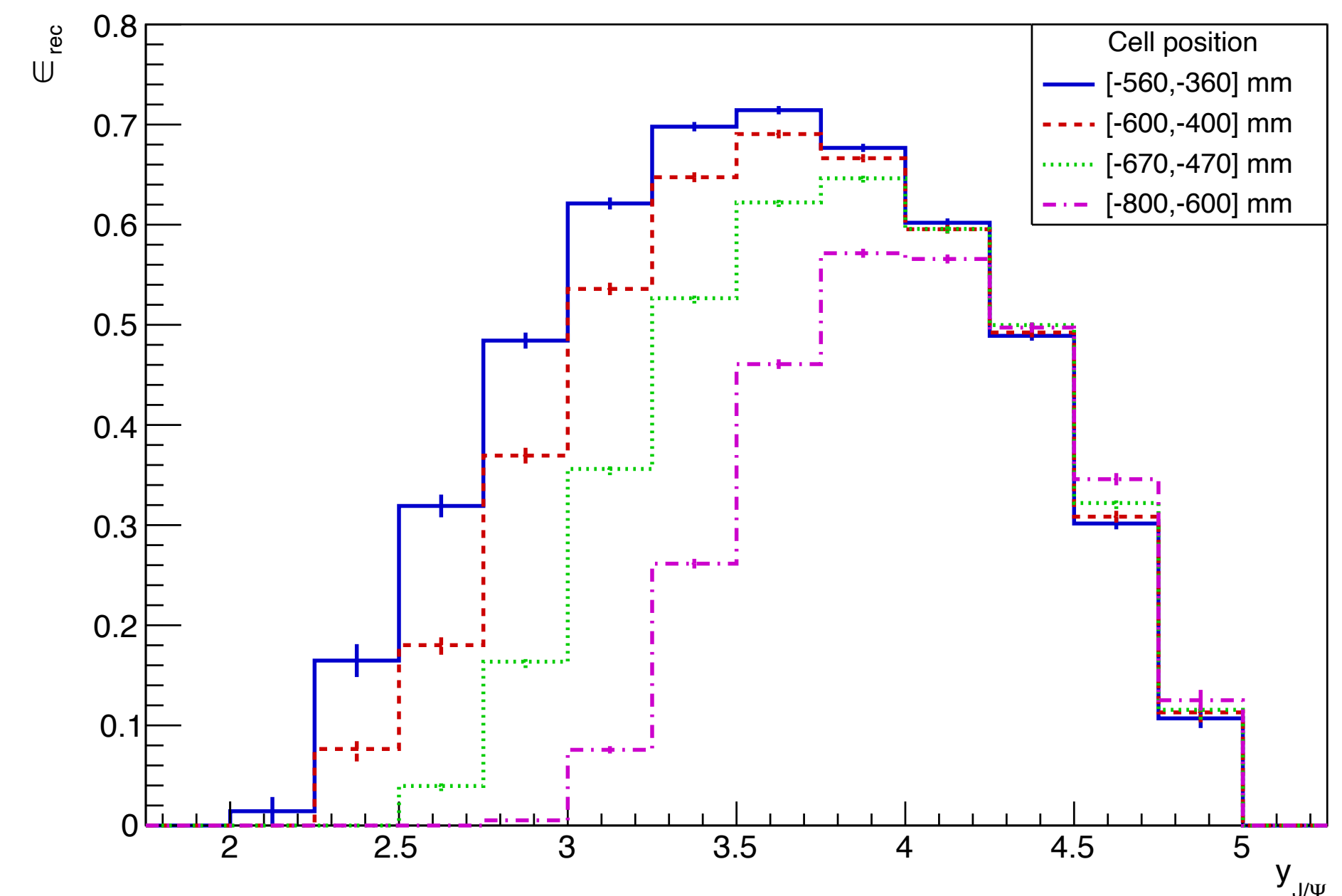
- LHCb simulations show broader kinematic acceptance & better reconstruction efficiency when the cell is close to the VELO, however:
  - Several reconstruction & trigger improvements have already been deployed for the Run 3
  - More reliable number will come with data next year, so here the **focus is on ratios of efficiencies**



$J/\Psi \rightarrow \mu^+\mu^-$  reconstruction efficiency vs cell position



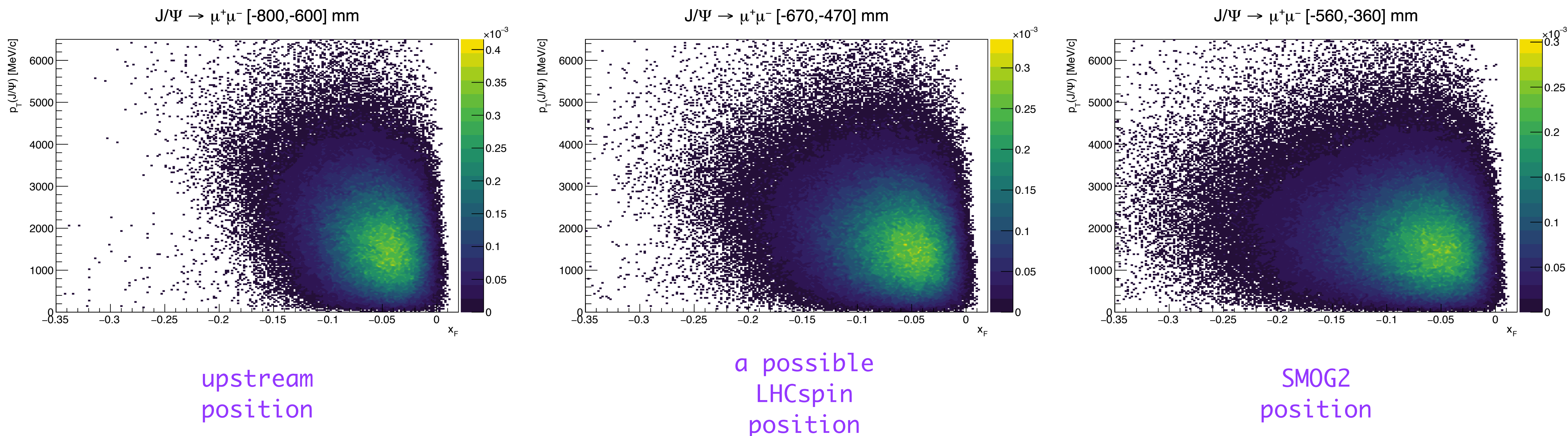
$J/\Psi \rightarrow \mu^+\mu^-$  reconstruction efficiency



# $J/\psi \rightarrow \mu^+ \mu^-$ kinematics vs cell position

- Using  $x_F = 2E_T / \sqrt{s_{NN}} \sinh(y^*)$  with  $E_T^2 = M^2 + P_T^2$
- Actual SMOG2 region  $[-560, -360]$  mm as a reference,  $[-670, -470]$  mm a possible solution to fit the LHCspin setup
- The kinematic coverage depends on the cell position  $\rightarrow p_T$  slightly affected,  $x$  range shrinks when moving upstream:

▶ VELO



- Here are  $\sim 300k$  events per plot, just a small fraction of the expected data