



Particle Detectors

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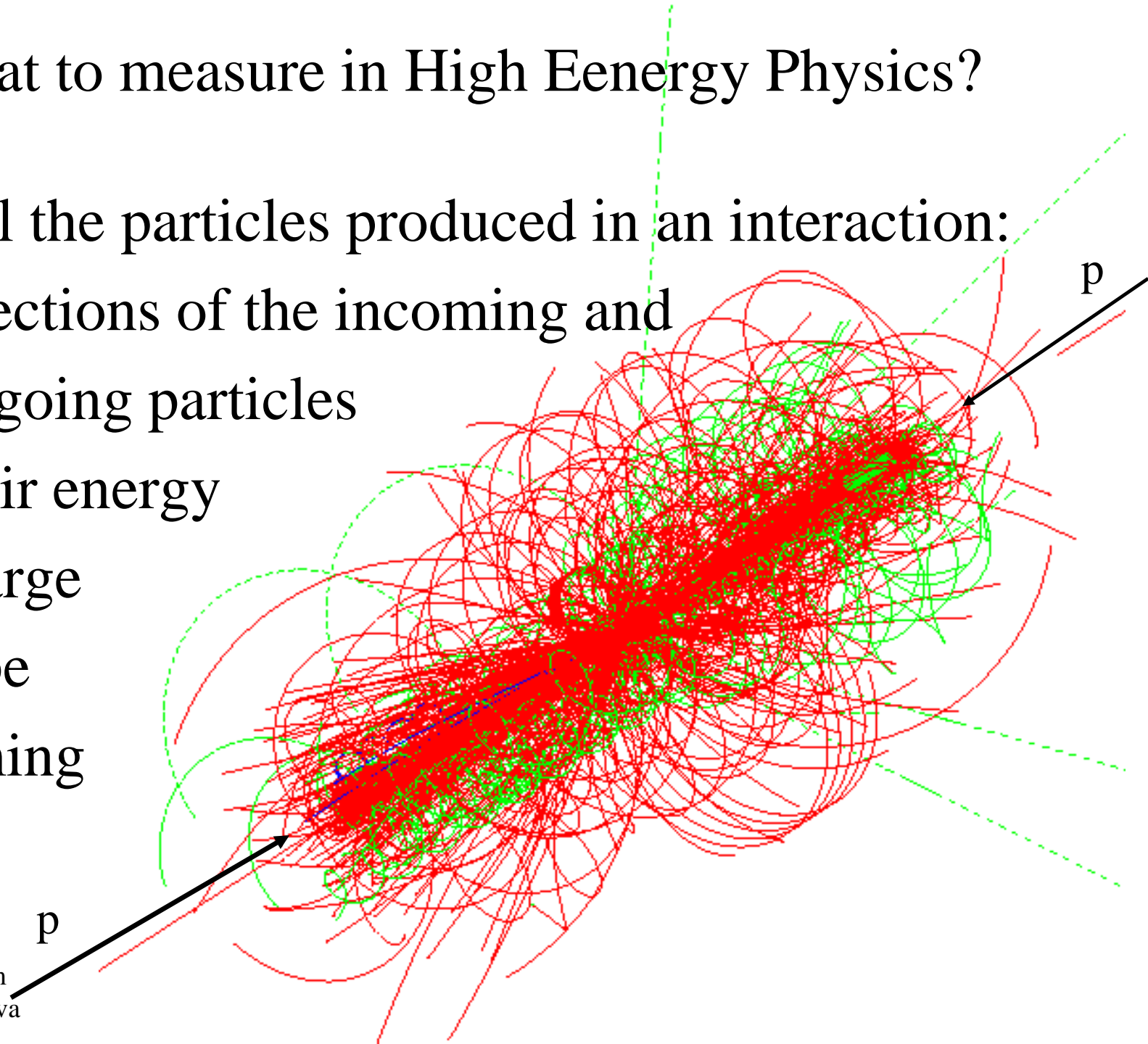
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What to measure in High Energy Physics?

For all the particles produced in an interaction:

- Directions of the incoming and outgoing particles
- Their energy
- Charge
- Type
- Timing



How to measure a radiation

- To measure radiation one needs to find out
 - **Type of radiation**
 - **Its losses in an assumed detector (will we get a signal, is it above the existing noise?)**
 - **Energy** : How much, in which form ?
 - **Interaction type and the** profile in time
 - **Position** : Interaction point, depth
- The ability of a particular instrument to measure the incident radiation can be measured by its
- **Efficiency** : of the detector we choose

Different types of radiation detectors

- Radiation detectors can be made on the basis of different materials
 - **Gas** : Gas filled detectors such as ionisation chambers, proportional counters and Geiger Muller (GM) tubes utilise a gas as the detection medium.
 - The above also includes the air in the atmosphere.
 - **Scintillation detectors** : Utilise either a liquid or solid state scintillator as the detection medium.
 - **Semiconductor detectors** : An elemental or compound semiconductor crystal is used as the detection medium.
- Every type has its own advantages
- A composite approach can often provide much more information about a radiation source.

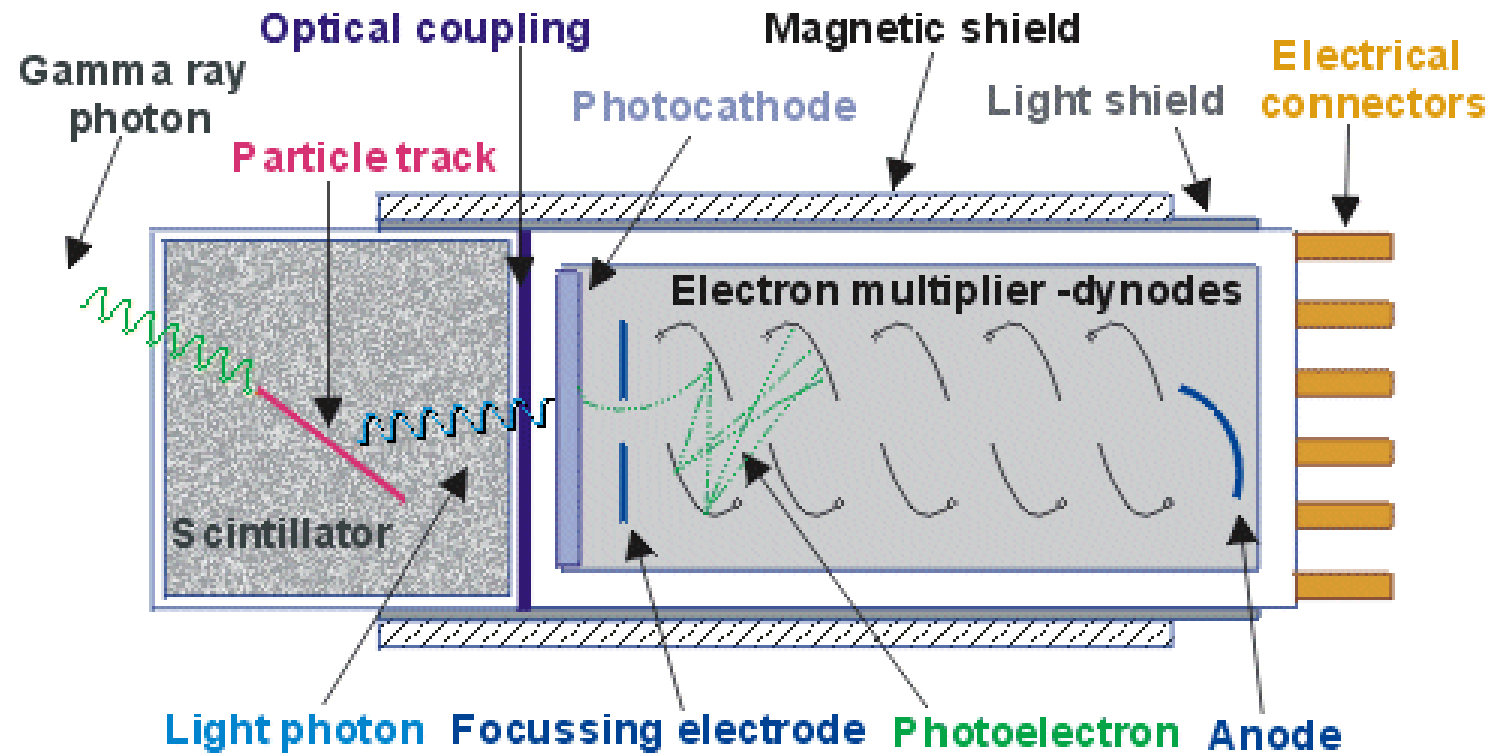
Gas detectors

- Gas detectors in general offer the following:
 - Poor energy resolution
 - Time resolution
 - Excellent position resolution (MWPC)
 - Low Efficiency:
 - Large volume possible
 - Low density/ Z – very low stopping power

Solid state detectors

- The use of a solid state detection medium is used in many radiation detection applications.
- For the measurement of high-energy electrons or gamma-rays detector dimensions can be kept much smaller than equivalent gas filled detectors because solid densities are some 1000 times greater than that for a gas.
- **Scintillation detectors** offer:
 - Average energy resolution.
 - Good/very good time resolution (sub ns)
 - Reasonable position resolution (~mm)
 - High efficiency

Scintillation detectors



The energy required to produce one information carrier is of the order of 100eV, the number of carriers created in a typical interaction is usually no more than a few thousand → statistical fluctuations → poor energy resolution.

A dream detector

Should respond to all types of particles, providing:

- particle identification
- measurement of energy/momentum
- measurement of trajectory (direction/origin)
- cover the full 4π solid angle
- time response

in addition:

- provide short dead-time (allow for high rate)
- wide dynamic range
- high radiation hardness
- long-time operation

As a rule, a real detector is a compromise



Ideal domestic animal -
Eierliegende Wollmilchsau

Different particle detection techniques

Particle can interact with matter by producing:

- Ionisation of atoms
- Bremsstrahlung and photon conversions
- Inelastic nuclear interactions
- Cherenkov or transition radiation
- Emission of scintillation or fluorescence light

How can we “visualise” these processes?

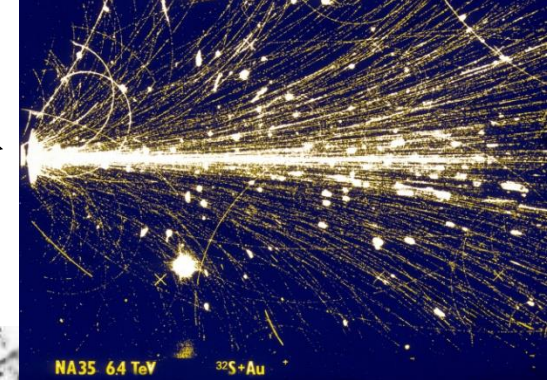
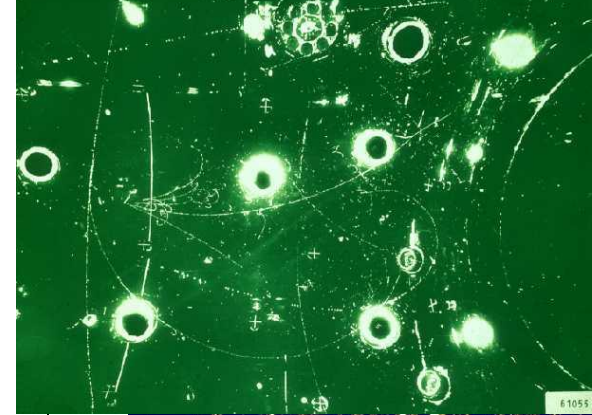
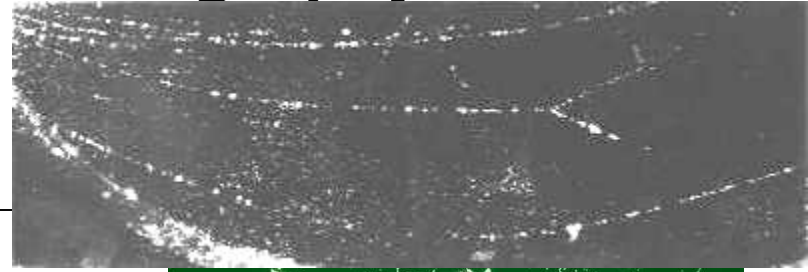
- By taking snapshots
- By collection of induced by ionisation charge
- By detecting of photons

Detection technique: Photography

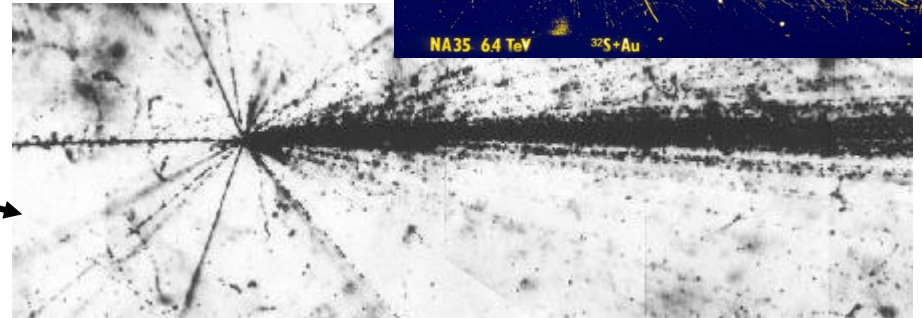
Charged particles ionise atoms along their trajectory.

(I) Ions act as seeds for:

- condensation in super saturated gas (Wilson chamber)
- bubble-formation in super-heated liquid
- electrical discharge or plasma formation



(II) Ionisation can also be made visible chemically in emulsion targets



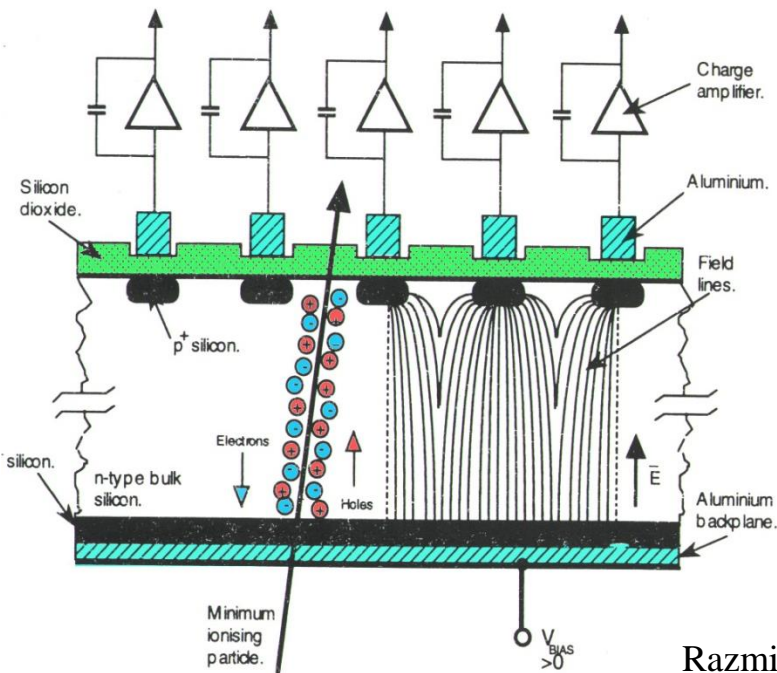
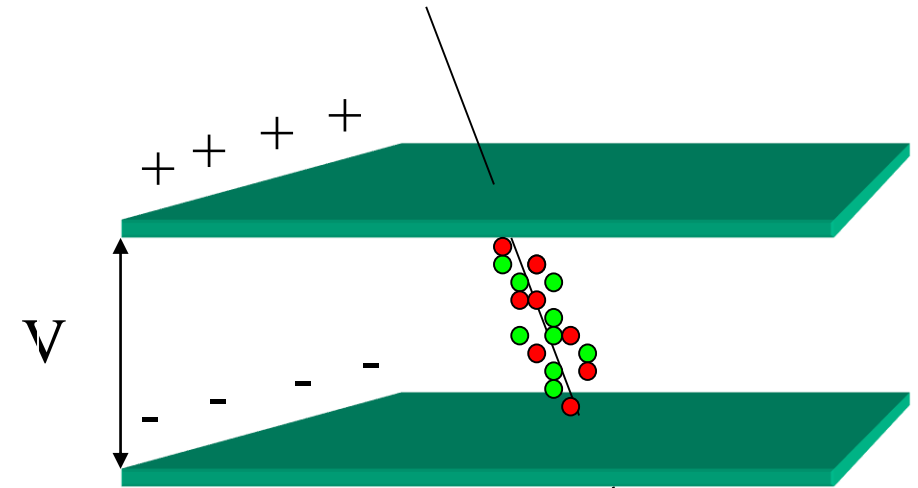
Detection technique: Collect electrical charge

Particle causes ionisation in a material.

Charge is separated/collected by an electric field.

Requirement on material:

- no/few free charge carriers (non-conducting)
- mechanism for transport of charge



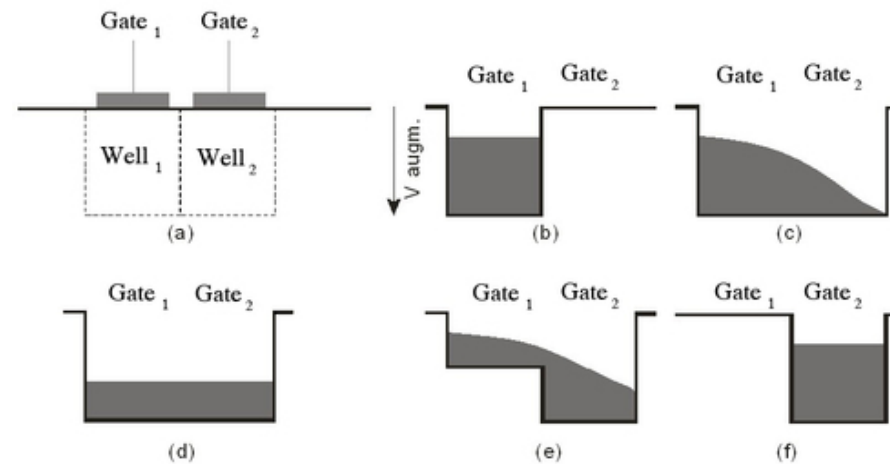
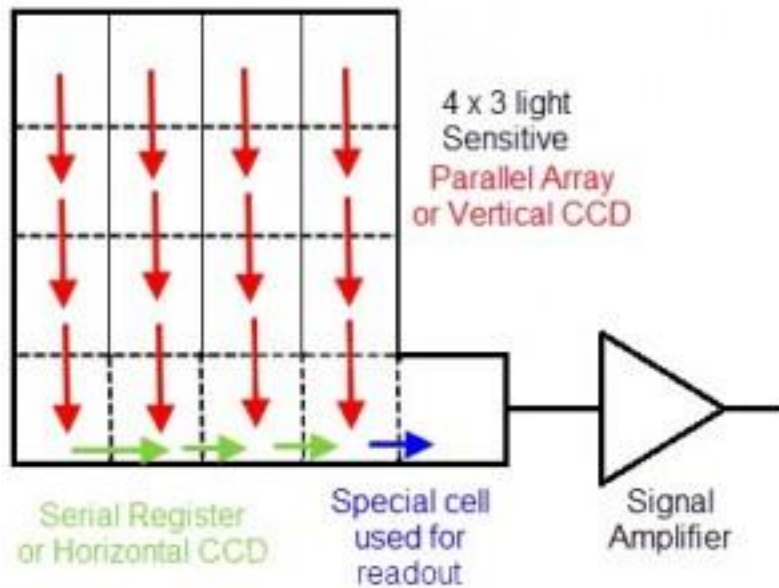
Proportional chambers, Drift chambers, ..

Insulating gas/liquid between anode and cathode (transport through drift). Sometimes also low conductivity solids.

Silicon strip detectors, CCDs, ..

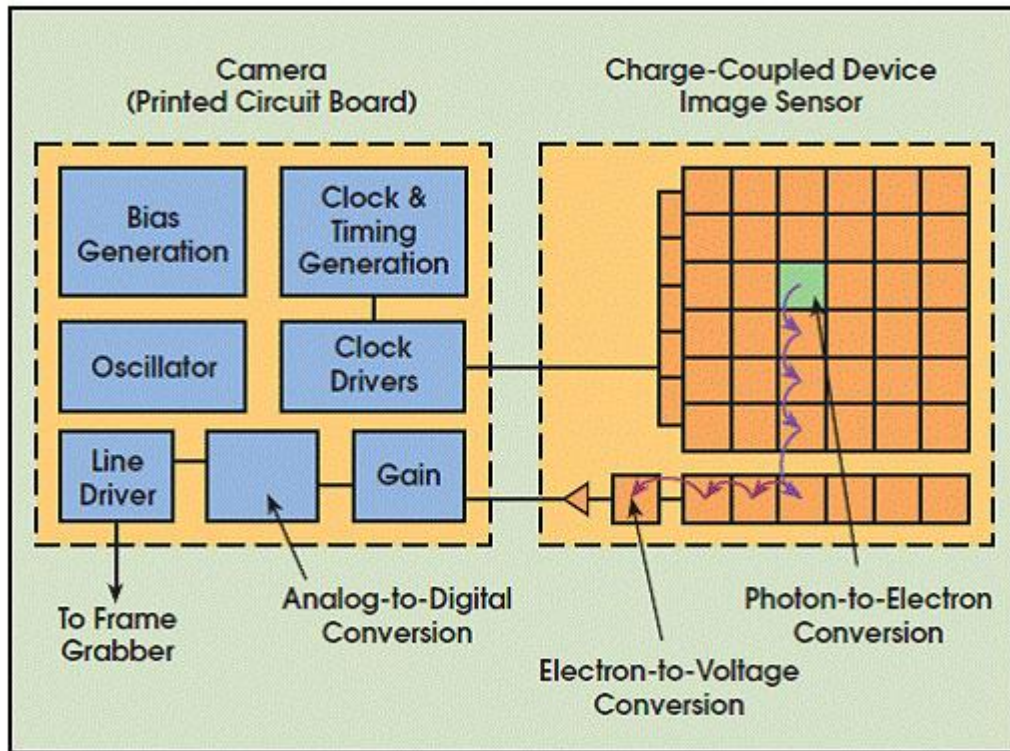
Using a semi-conducting material: Mostly in the form of a reverse-biased pn-junction diode.

How Functions a Charge Coupled Device (CCD)

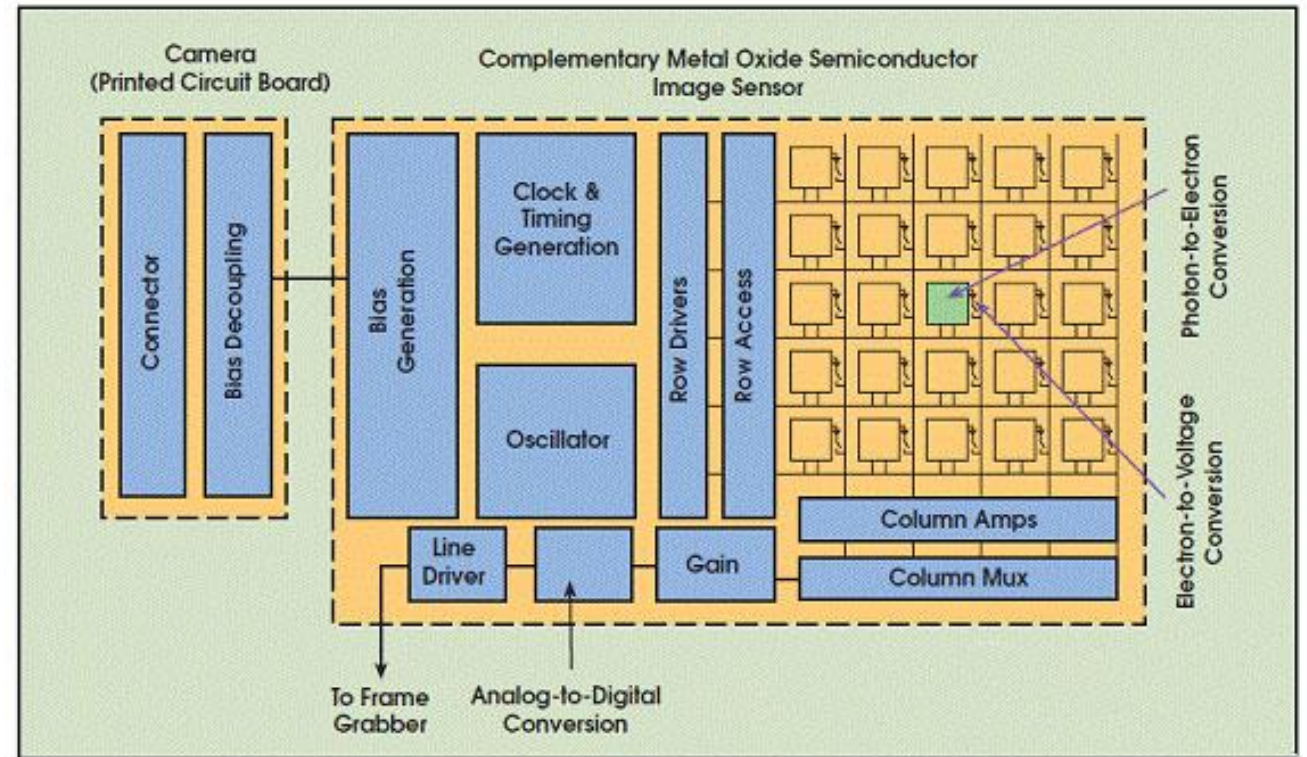


CCD versus CMOS sensor (these are what you have in your pocket, in your mobile phone)

CCD



CMOS



Detection technique: Photo-detection

Charged particles can produce photons via scintillation, Cherenkov or transition-radiation effects. One can detect these by using:

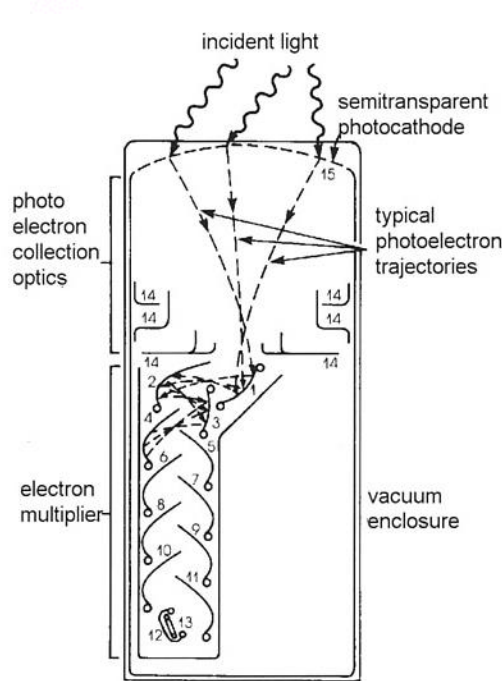


Photo-Multiplier Tube (PMT):

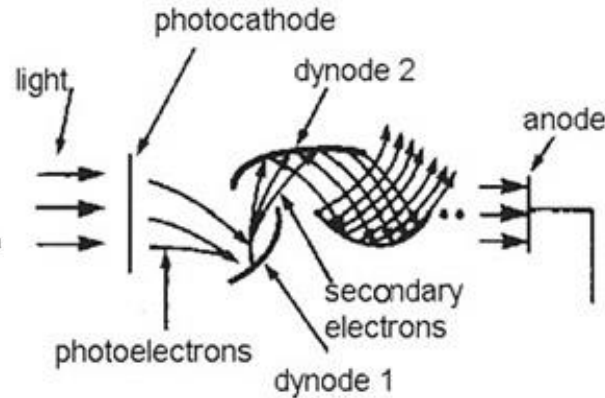
Electrons from photo-electric effect

“Electron multiplier” provides charge cascade

Very sensitive vacuum device



1 - 12 dynodes 14 focussing electrodes
13 anode 15 photocathode



Hamamatsu



Semi-conductor devices:
Photo-diodes or CCD's

R-12992-100

7 dynode



mat window

ETE D573KFLSA

7 dynode



clear window

ETE D569/3SA

8 dynode

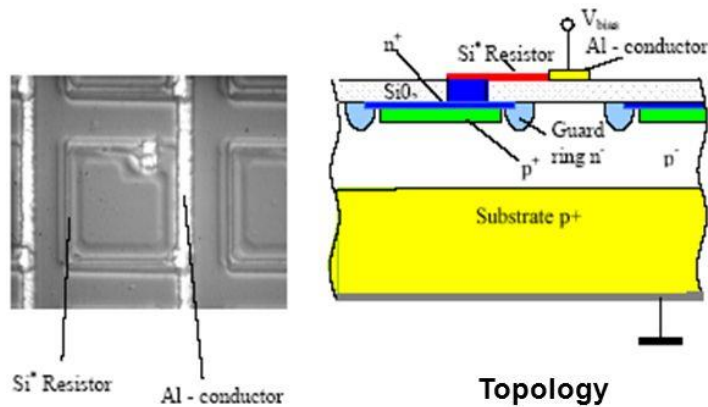


Scale [cm]

- Classical bialkali PMTs developed by us together with Electron Tubes (Great Britain) and Hamamatsu (Japan)
- The PMT on the top is currently worldwide the best

Silicon Photo Multiplier (SiPM)

SiPM - Principle of Operation



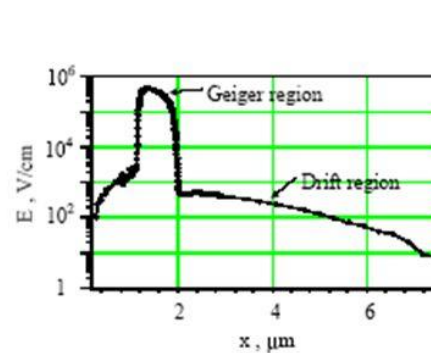
Si⁺ Resistor Al - conductor

Topology

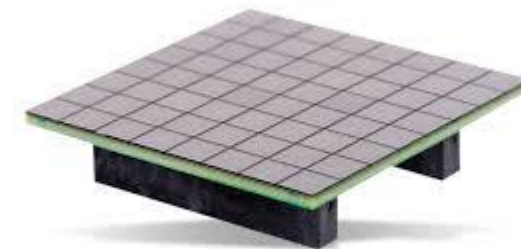
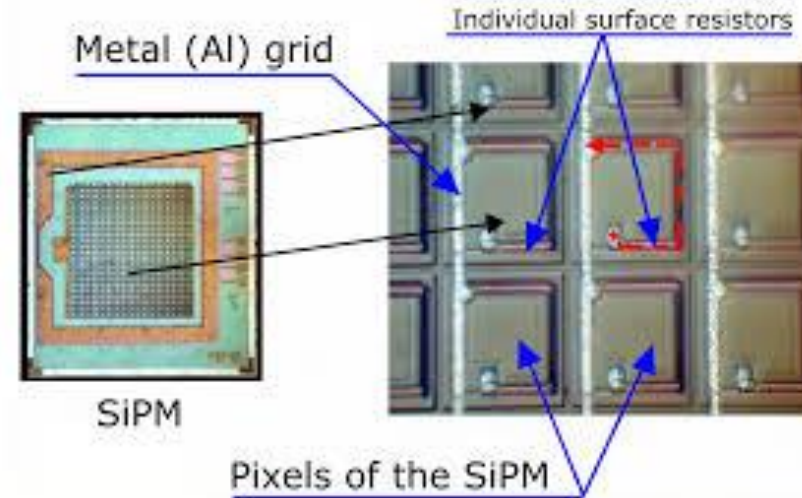
The pixel size – 32x32 μm²

Total number of pixels – 576 for 1x1 mm photodetector

Real topology is patented and different from it



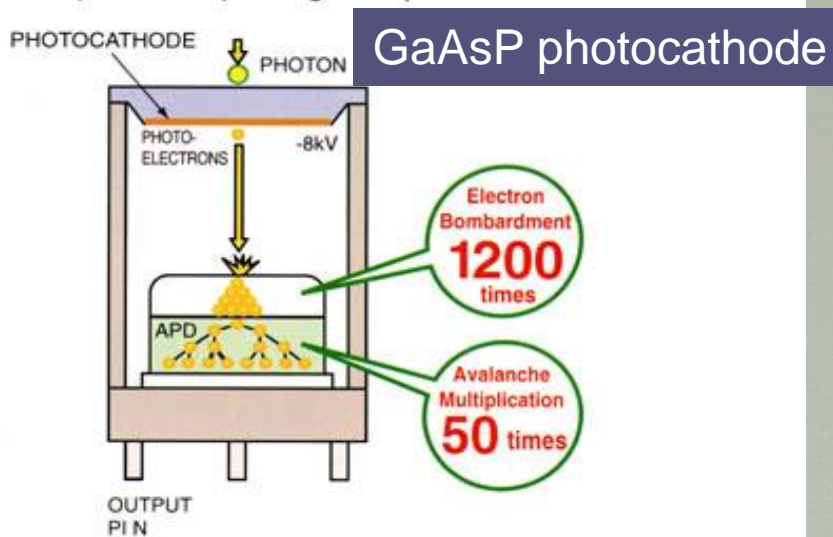
Electric field distribution in epitaxial layer



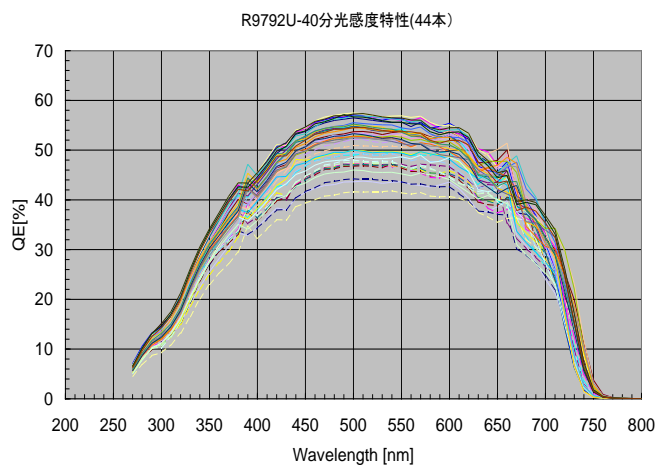
Vacuum detector HPD R9792U-40 developed for MAGIC

18mm GaAsP HPD by us with Hamamatsu

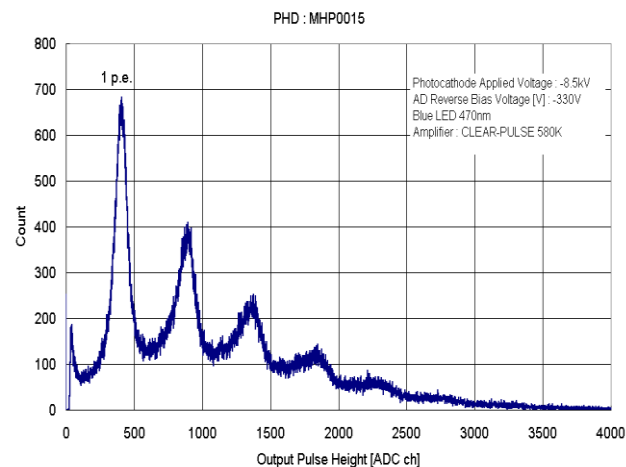
Compact HPD Operating Principle



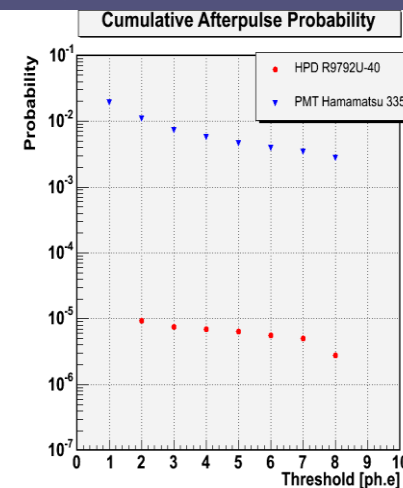
High Q.E.

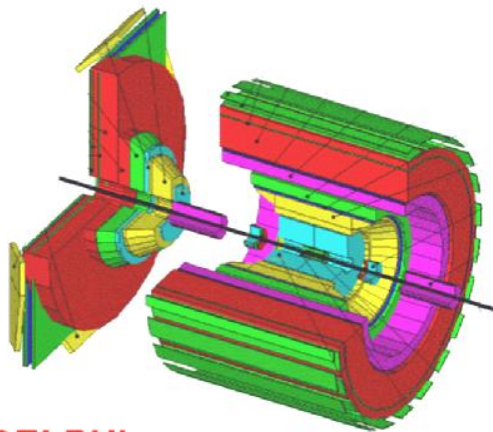


Good Charge Resolution



Very low after pulse rate

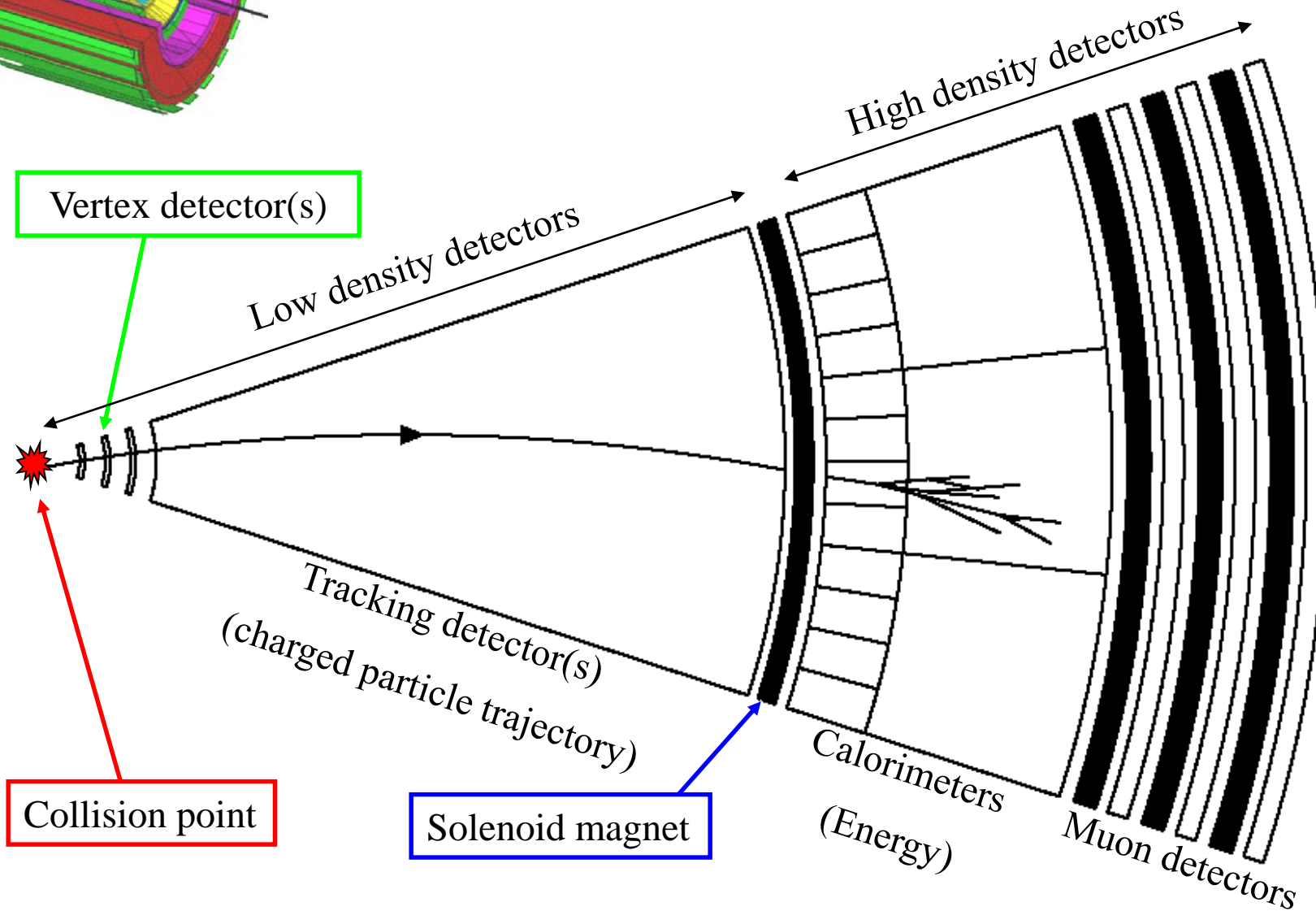




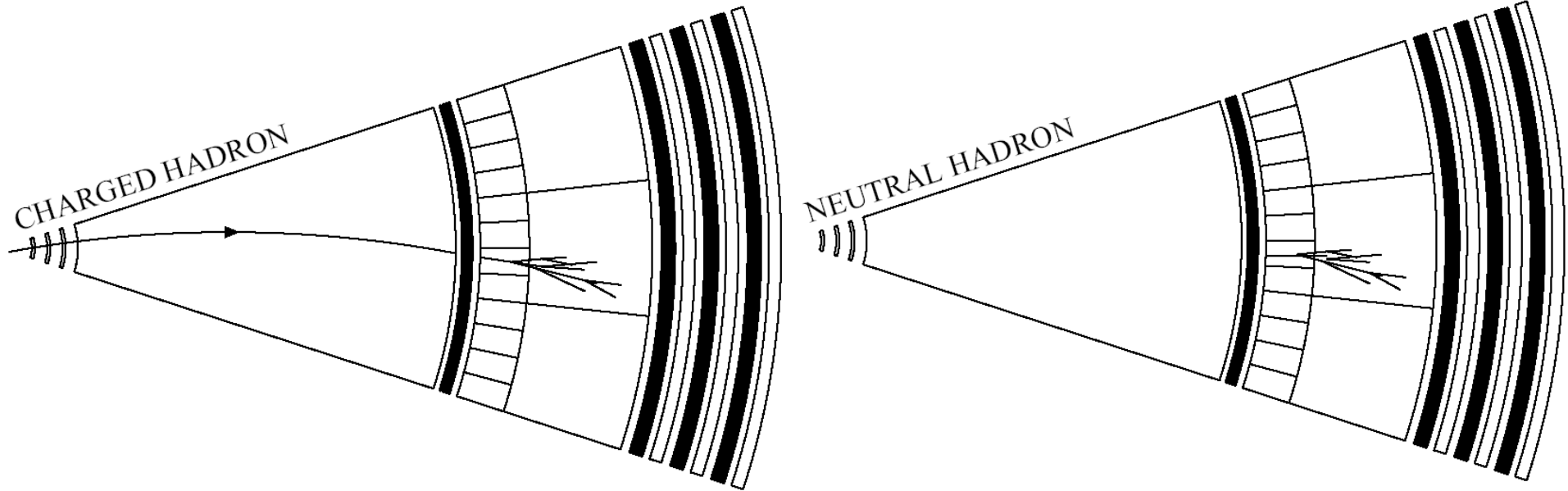
Global detector layout:

- barrel-shape surrounding beam-pipe
- 2 cone- or wheel-shaped end-caps

Nearly 4π coverage and good accessibility!



Particle signatures (first glance)



Charged hadrons:

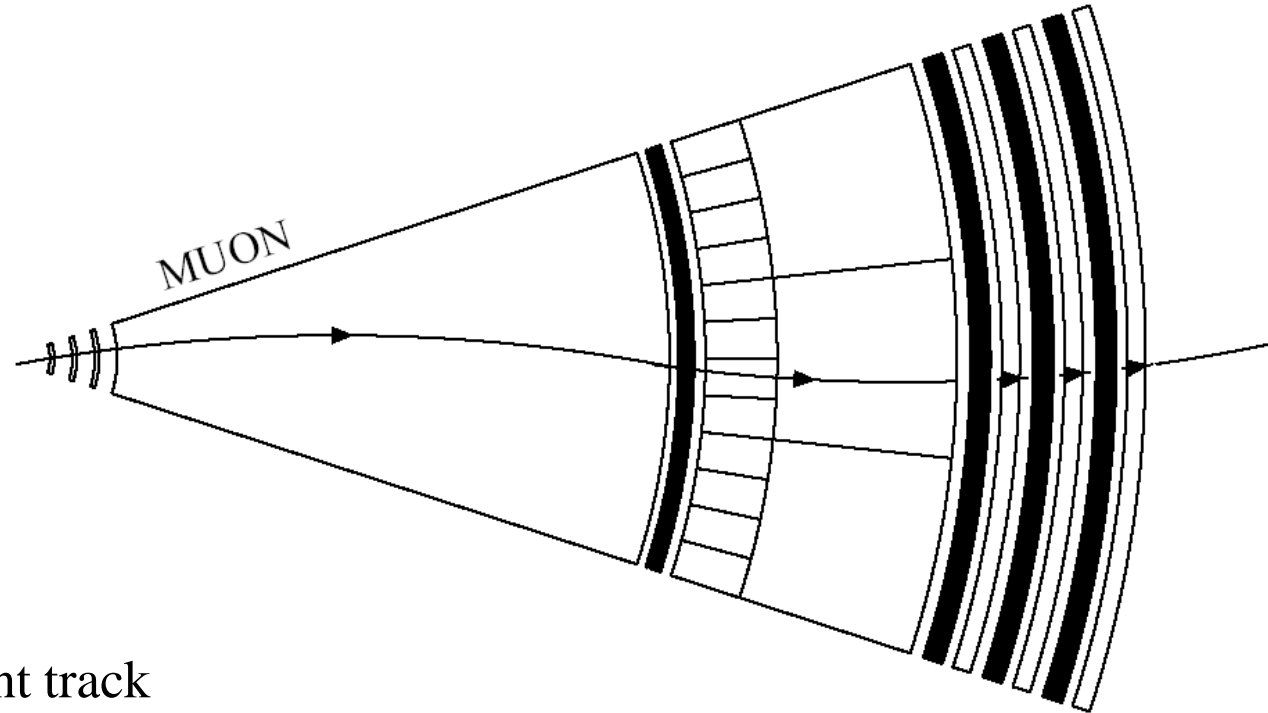
- leave a bent track
- stopped deep in calorimeter

Neutral hadrons:

- leave no track
- stopped deep in calorimeter

Second (+) layers of calorimeter: “*Hadron calorimeter*”

Particle signatures (at first glance)

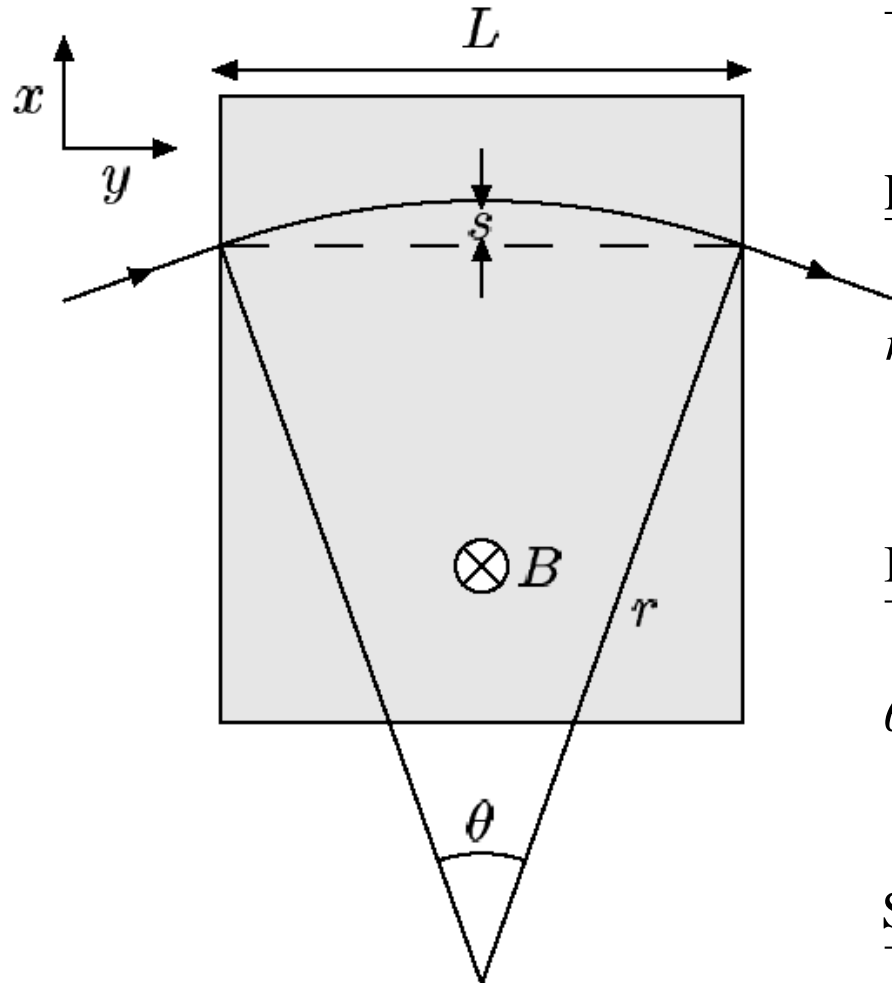


Muons:

- leave a bent track
- not stopped in calorimeter
- track in muon detectors

(Calorimeter, tracking and muon-detector information)

Charged particle in a magnetic field



Magnetic force: $\vec{F} = q(\vec{v} \times \vec{B})$ or $|F| = qv_{\perp}B$

Centrifugal force: $|F| = \frac{mv_{\perp}^2}{r}$

Radius: $r = \frac{mv_{\perp}}{qB} = \frac{P_{\perp}}{qB}$,

$r = \frac{P_{\perp}}{0.3B}$ (P_{\perp} in GeV/c and B in Tesla)

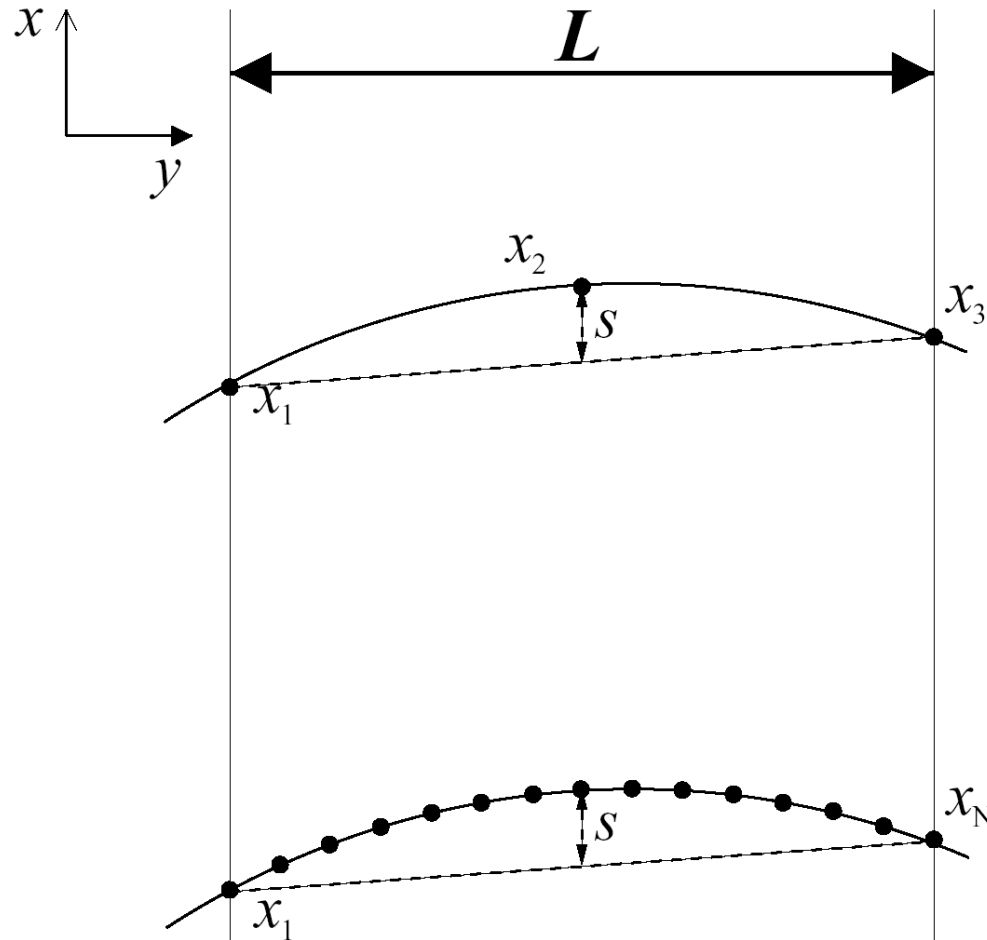
Bending angle: $\frac{\theta}{2} \approx \sin \frac{\theta}{2} = \left(\frac{L/2}{r} \right) \Rightarrow$

$\theta \approx \frac{0.3BL}{P_{\perp}}$

Sagitta :

$s = r - r \cos \frac{\theta}{2} = r \left(1 - \cos \frac{\theta}{2} \right) \approx r \frac{\theta^2}{8} \approx \frac{0.3}{8} \frac{BL^2}{P_{\perp}}$

Measurement of the sagitta



For 3 measured points :

x_1, x_2, x_3 at $y = 0, L/2, L$

$$s = x_2 - \frac{x_1 + x_3}{2}$$

$$\sigma(s) = \sqrt{\frac{3}{2}} \sigma(x)$$

$$\frac{\sigma(P_{\perp})}{P_{\perp}} = \frac{\sigma(s)}{s} = \frac{\sqrt{\frac{3}{2}} \sigma(x) 8P_{\perp}}{0.3BL^2}$$

For N equidistant measurements :

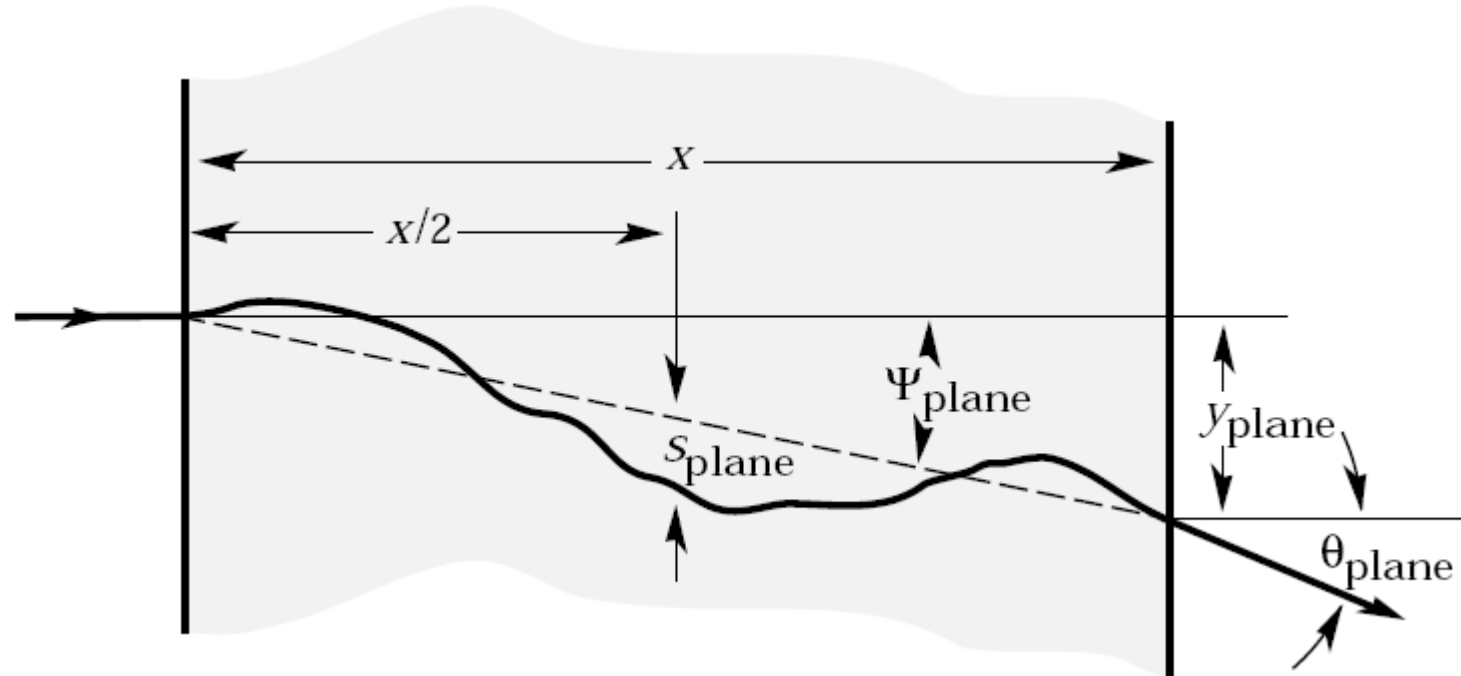
$$\frac{\sigma(P_{\perp})}{P_{\perp}} = \frac{\sigma(x) P_{\perp}}{0.3BL^2} \sqrt{\frac{720}{N+4}}$$

(R.L. Gluckstern, NIM 24 (1963) 381)

Thus for precise measurement momentum we need: high B field, large volume tracking detector, many measurements along the trajectory. But ...

Multiple Scattering

- Particle



MS Theory

- Average scattering angle is roughly Gaussian for small deflection angles

$$\theta_0 = \frac{13.6 \text{ MeV}}{\beta c p} z \sqrt{\frac{x}{X_0}} \left[1 + 0.038 \ln \left(\frac{x}{X_0} \right) \right]$$

- With

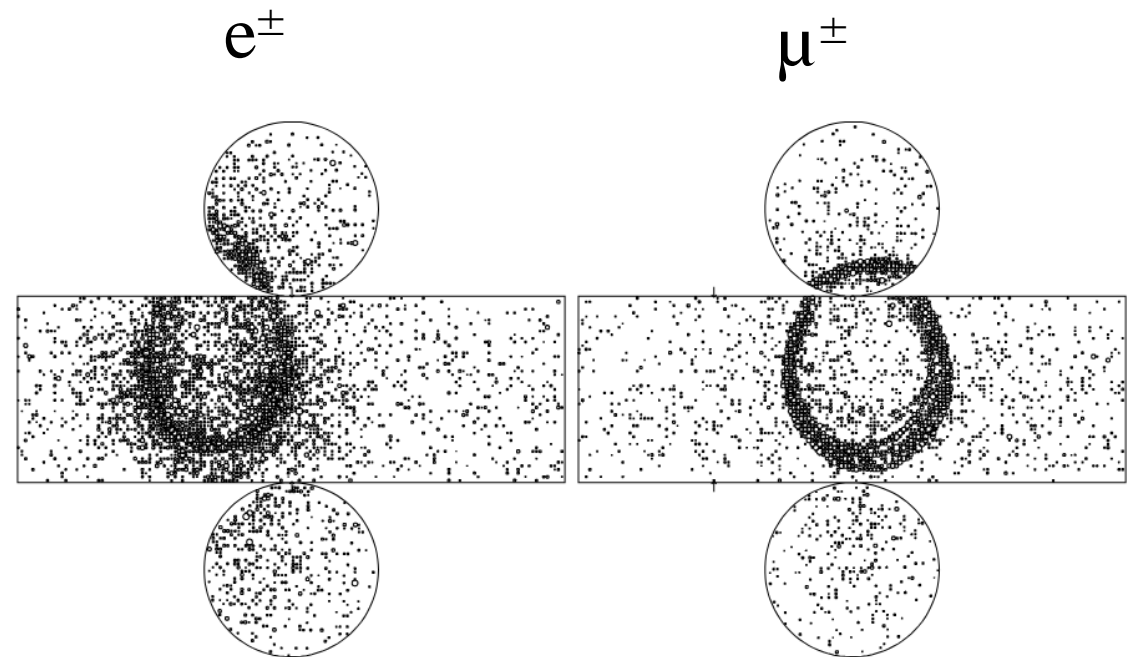
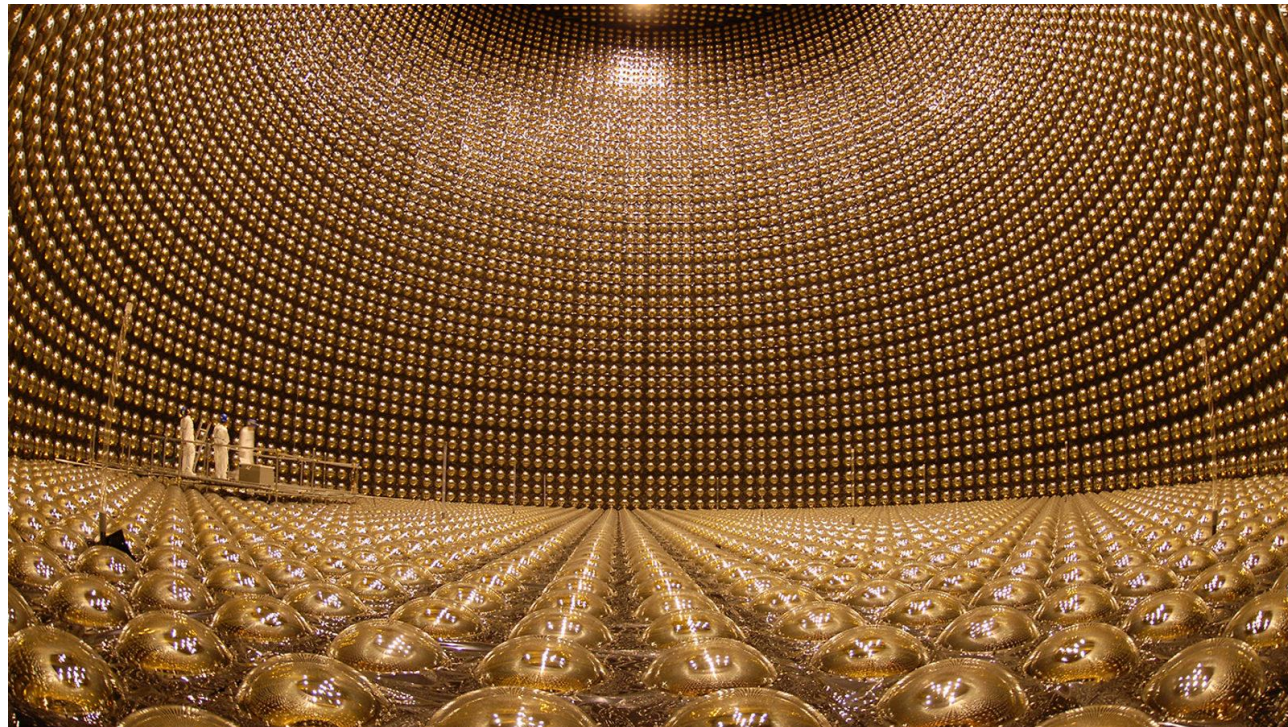
$X_0 \equiv$ radiation length

- Angular distributions are given by $\frac{dN}{d\Omega} \propto \frac{1}{2\pi\theta_0^2} \exp\left(-\frac{\theta_{space}^2}{2\theta_0^2}\right)$

$$\frac{dN}{d\theta_{plane}} \propto \frac{1}{\sqrt{2\pi}\theta_0} \exp\left(-\frac{\theta_{plane}^2}{2\theta_0^2}\right)$$

Electron and muon images in Superkamiokande detector demonstrate the effect of multiple scattering

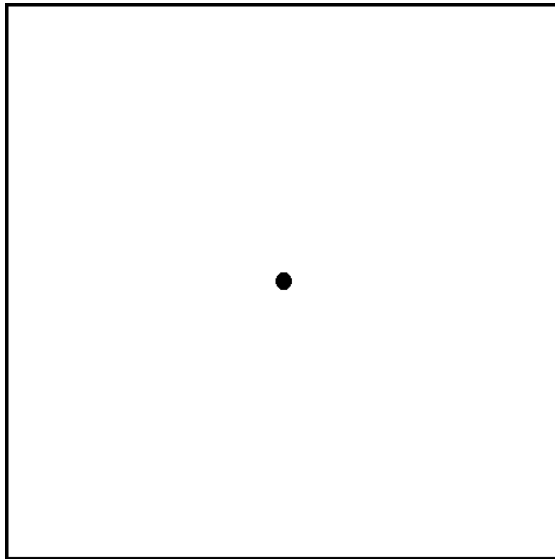
Kajita, 2004



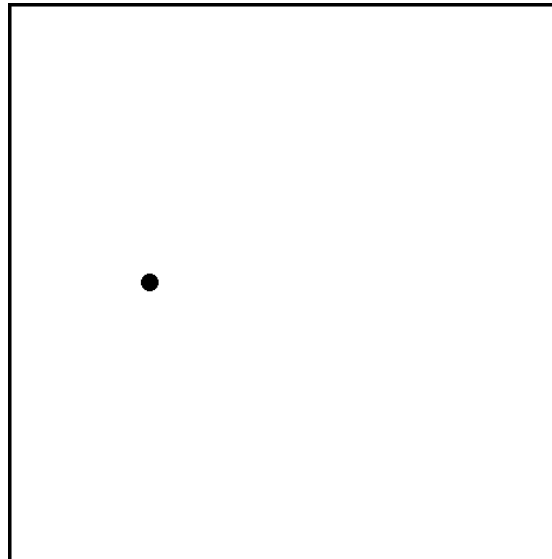
The Superkamiokande detector in Japan
uses 11200 PMTs of 0.5m size

Cherenkov Radiation

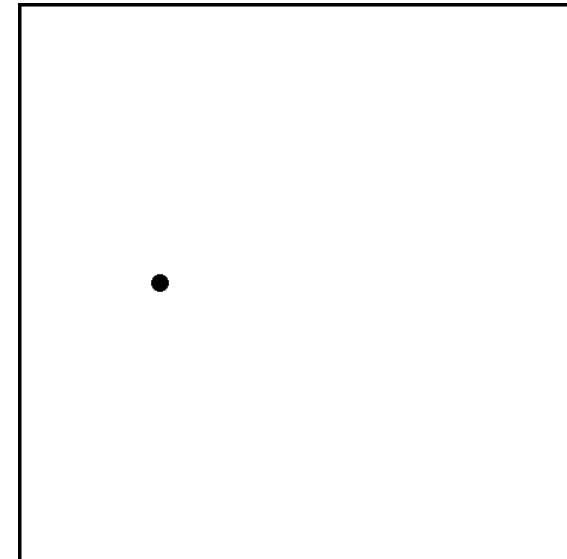
- Charge motion in a transparent dielectric medium



at rest



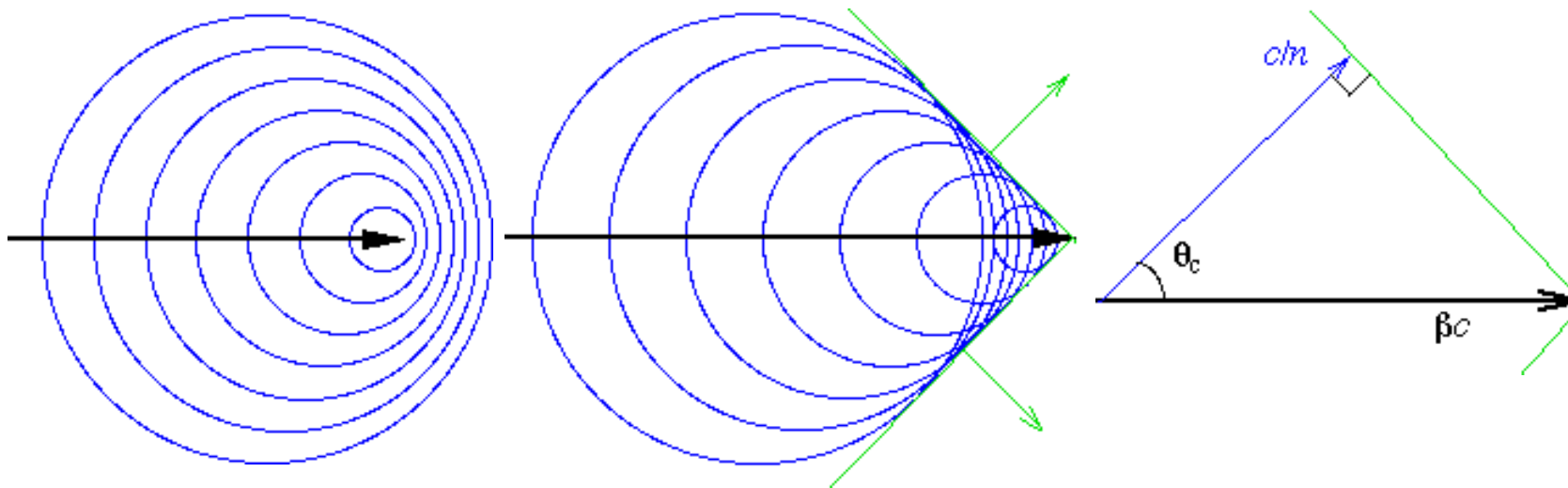
Slow $< c$



Fast $> c$

Cherenkov Radiation (2)

- Wave front comes out at certain angle

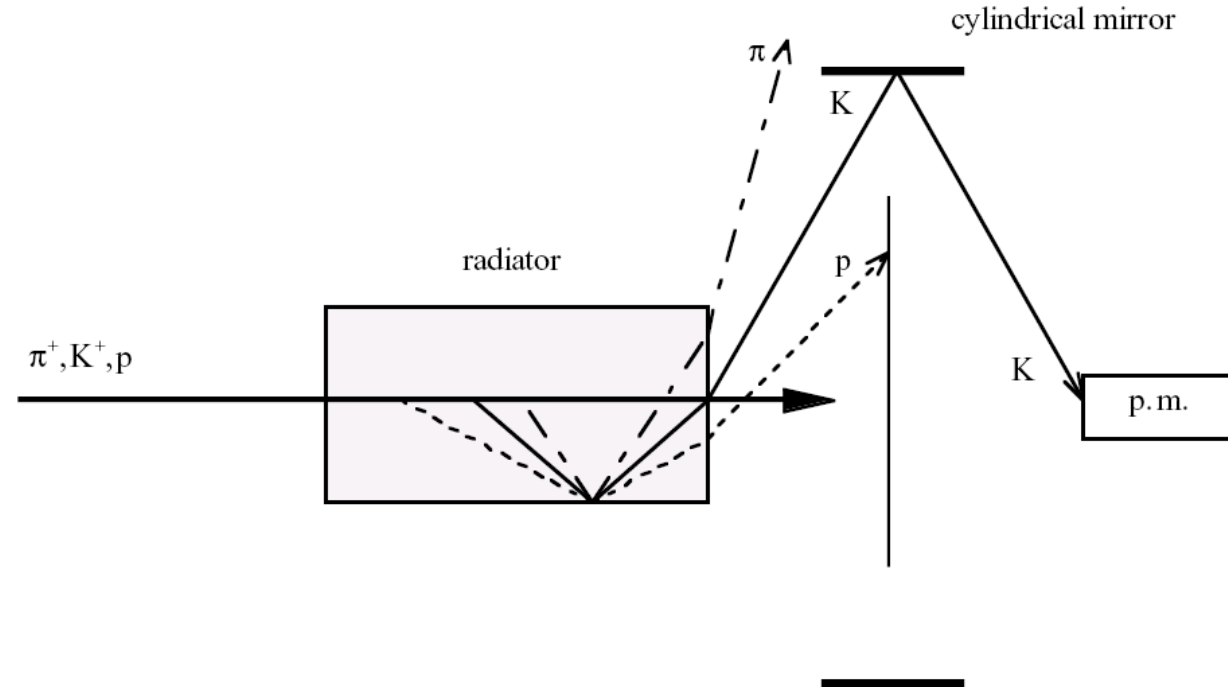


$$\cos \theta_c = \frac{1}{\beta n}$$

Different Types of Cherenkov Detectors

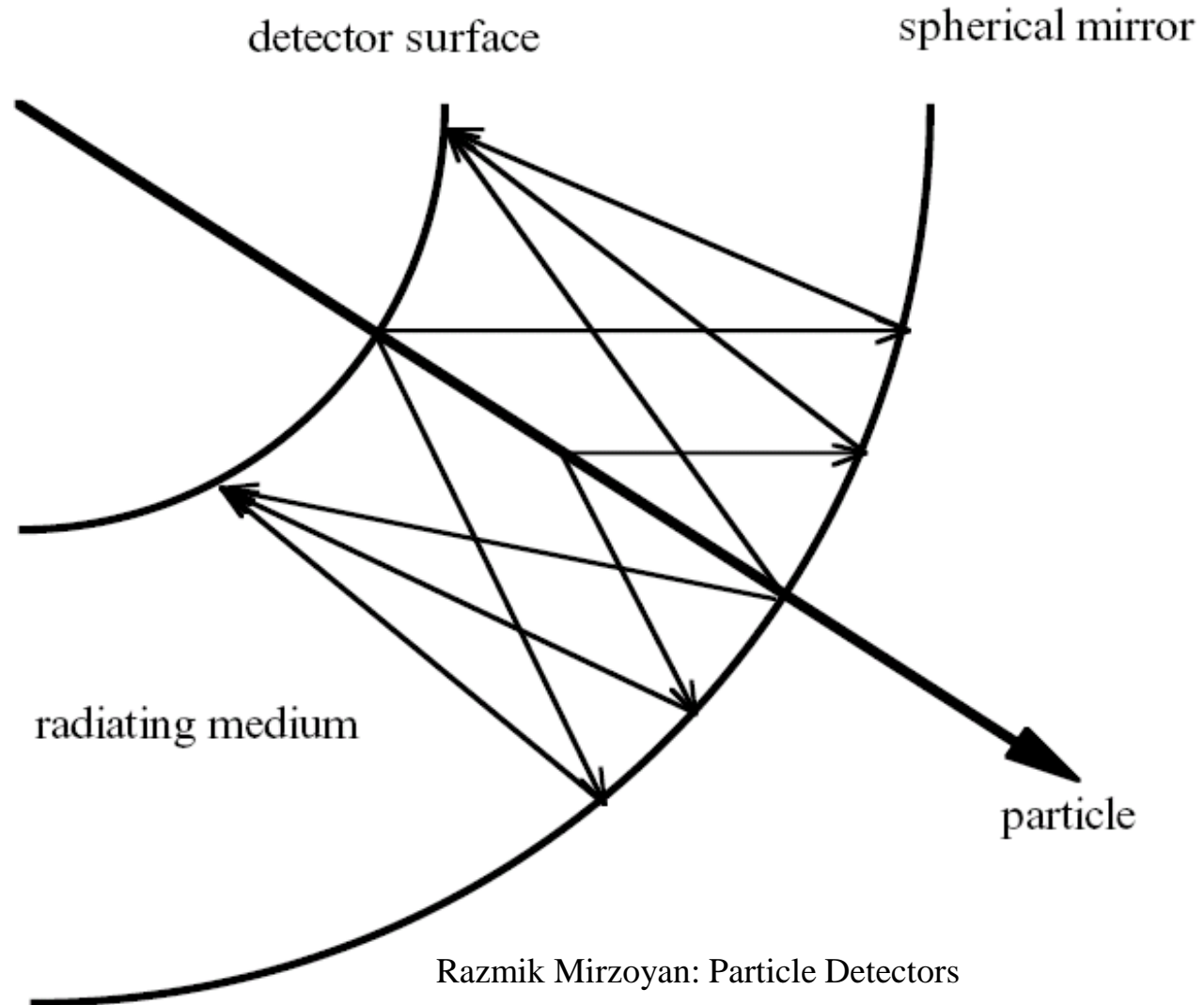
- Threshold Detectors
 - Yes/No on whether the speed is $\beta > 1/n$
- Differential Detectors
 - $\beta_{\max} > \beta > \beta_{\min}$
- Ring-Imaging Detectors
 - Measure β

Differential Detectors



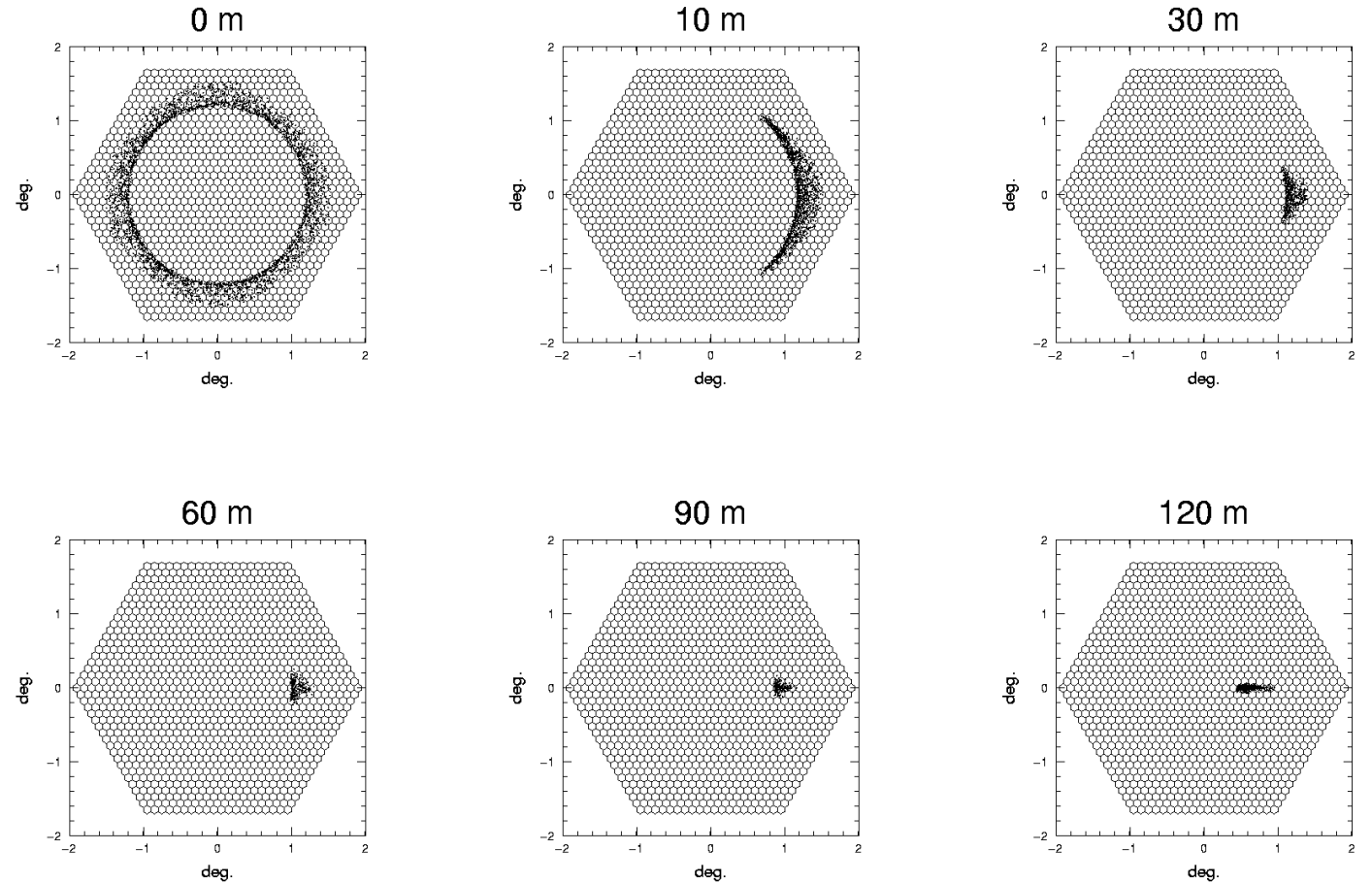
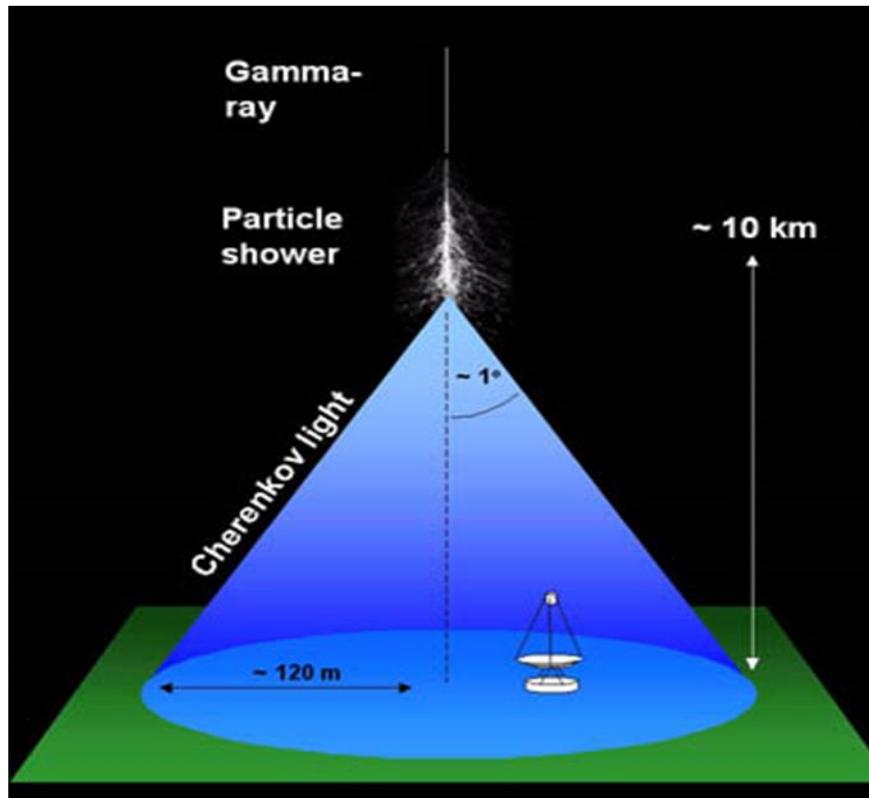
- Will reflect light onto PMT for certain angles only $\Leftrightarrow \beta$ Selecton

Ring Imaging Detectors (1)



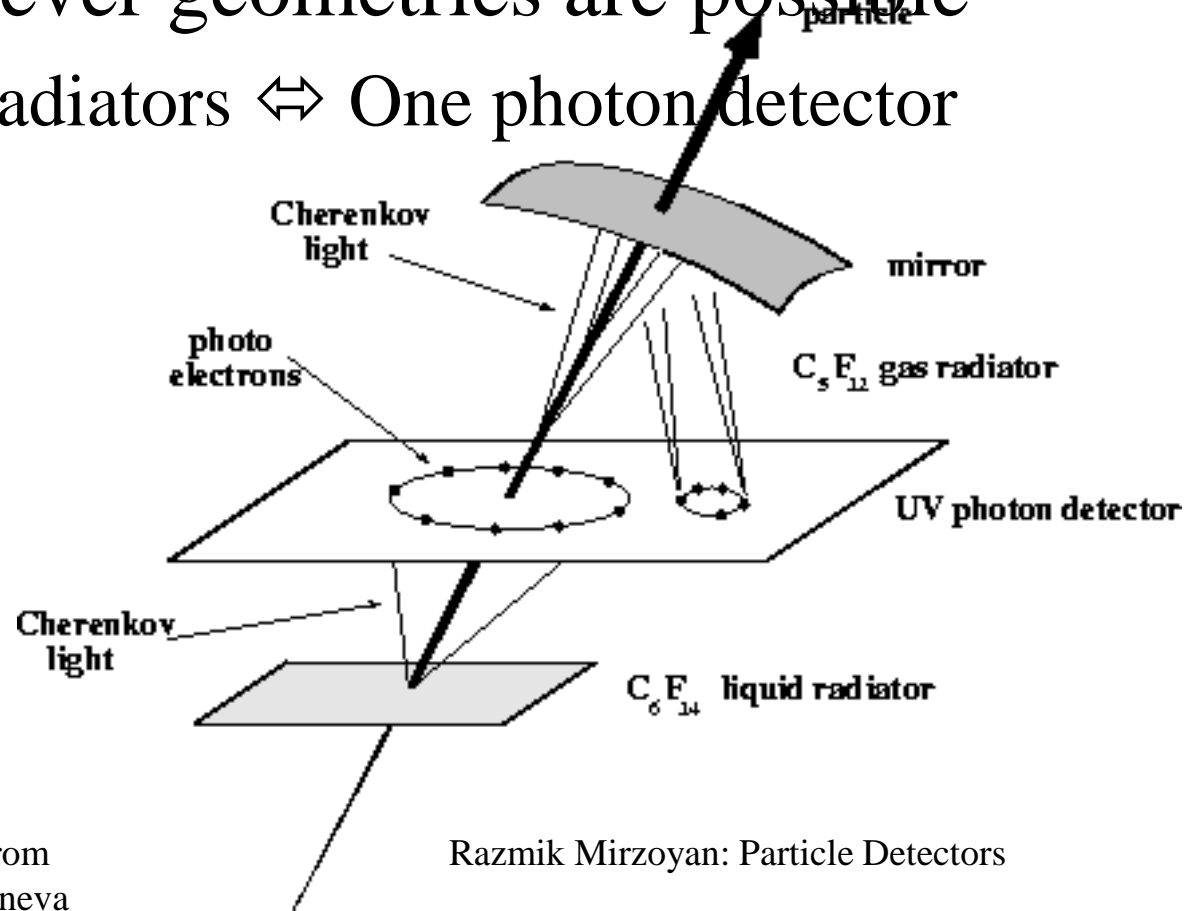
Shape of muon images versus the impact parameter in an Imaging Atmospheric Cherenkov Telescope

Imagine a downward going muon



Ring Imaging Detectors (3)

- More clever geometries are possible
 - Two radiators \Leftrightarrow One photon detector



Transition Radiation (3)

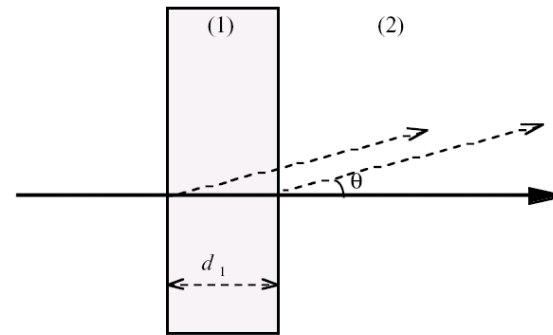
- Consider relativistic particle traversing a boundary from material (1) to material (2)

$$\frac{d^2 N}{d\omega d\Omega} = \frac{z^2 \alpha}{\pi^2 \omega} \phi^2 \times \left(\frac{1}{\omega_p^2 / \omega^2 + \phi^2 + 1/\gamma^2} - \frac{1}{\phi^2 + 1/\gamma^2} \right)^2$$

ω_p = plasma frequency

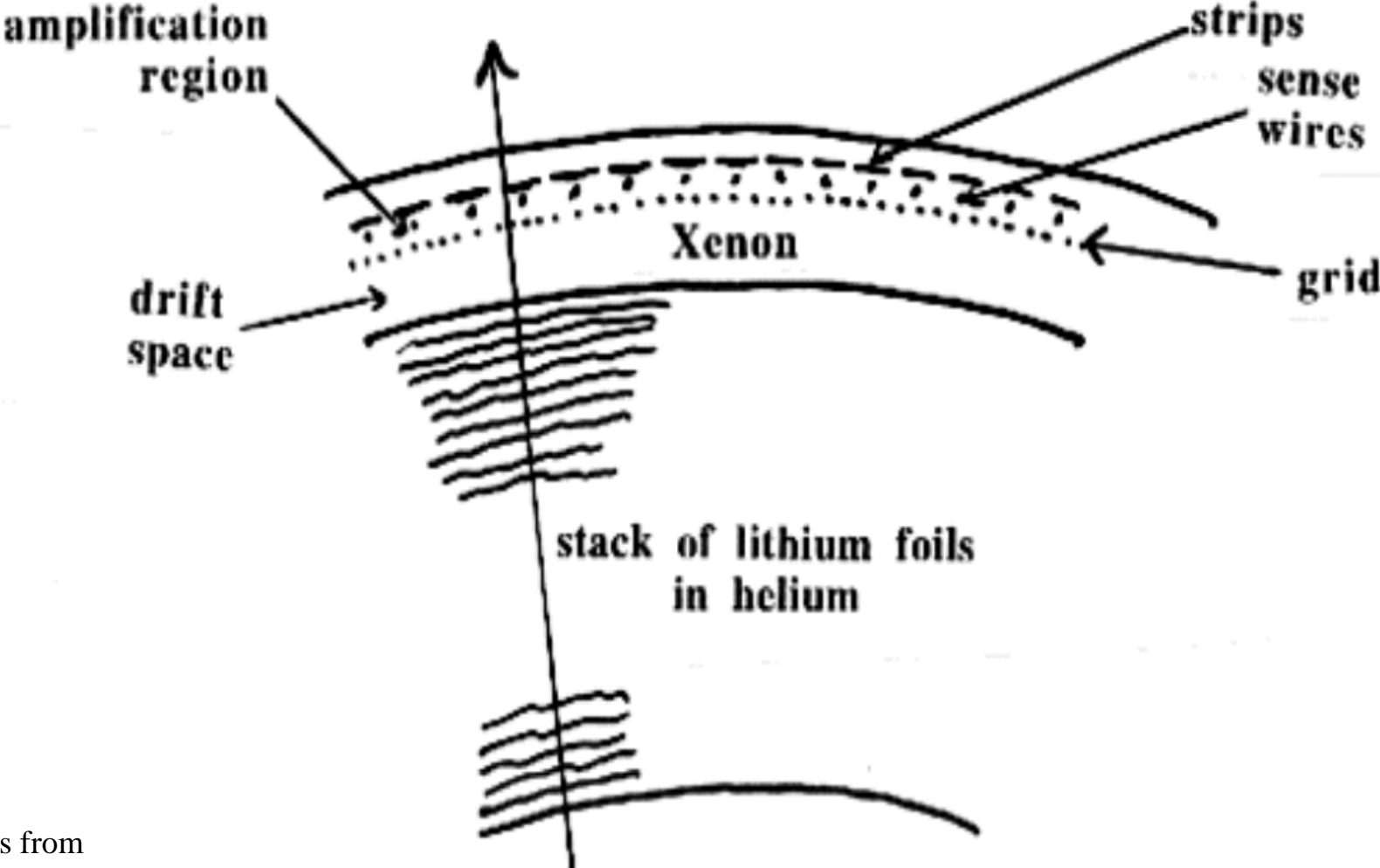
- Total energy radiated

$$E_{\text{TR}} = \frac{2}{3} \alpha h \omega_{p1} \gamma$$



- Can be used to measure γ

Transition Radiation Detector



Acknowledgements

- The author thanks for the sources:
 - Phys389 - Semiconductor Applications L11
 - 36. Particle Detectors for Non-Accelerator Physics
 - Particle Detection Techniques in HEP, Post-graduate lecture series, Joost Vosseveld
 - Interaction of Particles with Matter, Alfons Weber, CCLRC & University of Oxford Graduate Lecture 2004
 - New Journal of Physics 6 (20024) 194