

EFT approaches to flavour physics beyond the Standard Model

Wolfgang Altmannshofer
waltmann@ucsc.edu

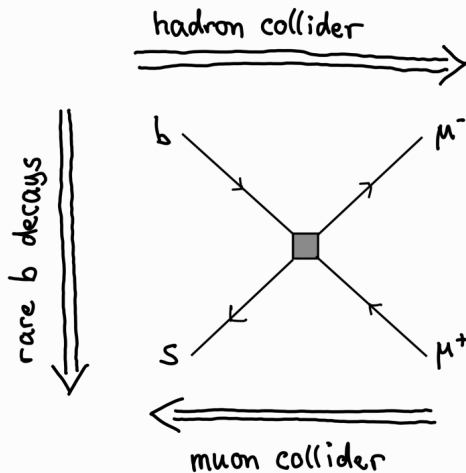


The Flavour Path to New Physics
June 5 - 7, 2024, University of Zurich

Part 1:
Collider Probes of
 $b \rightarrow s\mu\mu$

based on 2306.15017 with A. Gadam and S. Profumo

Collider Probes of $b \rightarrow s\mu\mu$



Non-Standard $\mu^+ \mu^- \rightarrow bs$ at a Muon Collider

$$\frac{d\sigma(\mu^+ \mu^- \rightarrow b\bar{s})}{d\cos\theta} = \frac{3}{16} \sigma(\mu^+ \mu^- \rightarrow bs) \left(1 + \cos^2\theta + \frac{8}{3} A_{\text{FB}} \cos\theta \right)$$

$$\frac{d\sigma(\mu^+ \mu^- \rightarrow \bar{b}s)}{d\cos\theta} = \frac{3}{16} \sigma(\mu^+ \mu^- \rightarrow bs) \left(1 + \cos^2\theta - \frac{8}{3} A_{\text{FB}} \cos\theta \right)$$

Total cross section **increases with the center of mass energy**
(unless the contact interaction is resolved)

$$\sigma(\mu^+ \mu^- \rightarrow bs) = \frac{G_F^2 \alpha^2}{8\pi^3} |V_{tb} V_{ts}^*|^2 s \left(|\Delta C_9|^2 + |\Delta C_{10}|^2 \right)$$

Non-Standard $\mu^+ \mu^- \rightarrow bs$ at a Muon Collider

$$\frac{d\sigma(\mu^+ \mu^- \rightarrow b\bar{s})}{d\cos\theta} = \frac{3}{16} \sigma(\mu^+ \mu^- \rightarrow bs) \left(1 + \cos^2\theta + \frac{8}{3} A_{\text{FB}} \cos\theta \right)$$

$$\frac{d\sigma(\mu^+ \mu^- \rightarrow \bar{b}s)}{d\cos\theta} = \frac{3}{16} \sigma(\mu^+ \mu^- \rightarrow bs) \left(1 + \cos^2\theta - \frac{8}{3} A_{\text{FB}} \cos\theta \right)$$

Total cross section **increases with the center of mass energy**
(unless the contact interaction is resolved)

$$\sigma(\mu^+ \mu^- \rightarrow bs) = \frac{G_F^2 \alpha^2}{8\pi^3} |V_{tb} V_{ts}^*|^2 s \left(|\Delta C_9|^2 + |\Delta C_{10}|^2 \right)$$

Forward backward asymmetry is sensitive to the **chirality structure**

$$A_{\text{FB}} = \frac{-3\text{Re}(\Delta C_9 \Delta C_{10}^*)}{2(|\Delta C_9|^2 + |\Delta C_{10}|^2)}$$

Non-Standard $\mu^+ \mu^- \rightarrow bs$ at a Muon Collider

$$\frac{d\sigma(\mu^+ \mu^- \rightarrow b\bar{s})}{d\cos\theta} = \frac{3}{16} \sigma(\mu^+ \mu^- \rightarrow bs) \left(1 + \cos^2\theta + \frac{8}{3} A_{\text{FB}} \cos\theta \right)$$

$$\frac{d\sigma(\mu^+ \mu^- \rightarrow \bar{b}s)}{d\cos\theta} = \frac{3}{16} \sigma(\mu^+ \mu^- \rightarrow bs) \left(1 + \cos^2\theta - \frac{8}{3} A_{\text{FB}} \cos\theta \right)$$

Total cross section **increases with the center of mass energy**
(unless the contact interaction is resolved)

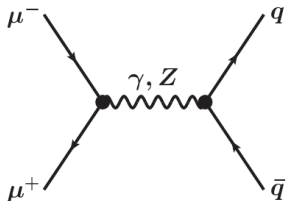
$$\sigma(\mu^+ \mu^- \rightarrow bs) = \frac{G_F^2 \alpha^2}{8\pi^3} |V_{tb} V_{ts}^*|^2 s \left(|\Delta C_9|^2 + |\Delta C_{10}|^2 \right)$$

Forward backward asymmetry is sensitive to the **chirality structure**

$$A_{\text{FB}} = \frac{-3\text{Re}(\Delta C_9 \Delta C_{10}^*)}{2(|\Delta C_9|^2 + |\Delta C_{10}|^2)}$$

Need **charge tagging** to measure the forward backward asymmetry

Main Background



- ▶ Mistagged dijets

$$\sigma_{bg}^{jj} = \sum_{q=b,c,s,d,u} 2\epsilon_q(1 - \epsilon_q)\sigma(\mu^+\mu^- \rightarrow q\bar{q})$$

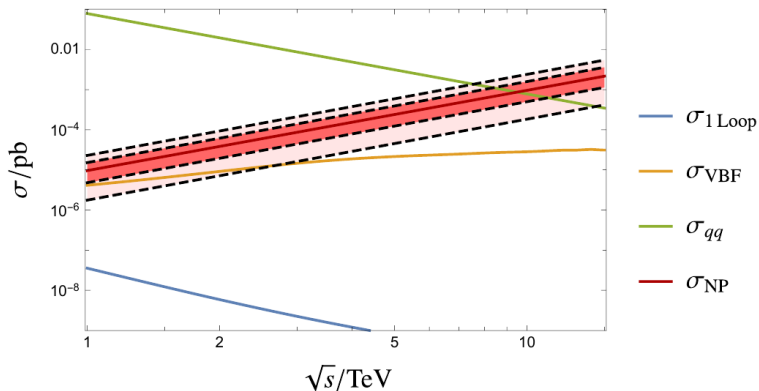
- ▶ Assume b tagging comparable to current LHC performance

$$\epsilon_b = 70\% , \quad \epsilon_c = 10\% , \quad \epsilon_u = \epsilon_d = \epsilon_s = 1\%$$

- ▶ Turns out to be the dominant background.

Signal vs. Background

WA, Gadam, Profumo 2203.07495, 2306.15017



- ▶ Main background falls with \sqrt{s} ; new physics signal increases.
- ▶ Signal/Background ~ 1 for $\sqrt{s} \sim 10$ TeV.

Forward Backward Asymmetry and Charge Tagging

$$\frac{d\sigma(\mu^+\mu^- \rightarrow b\bar{s})}{d\cos\theta} = \frac{3}{16}\sigma(\mu^+\mu^- \rightarrow bs) \left(1 + \cos^2\theta + \frac{8}{3}A_{\text{FB}}\cos\theta\right)$$

$$\frac{d\sigma(\mu^+\mu^- \rightarrow \bar{b}s)}{d\cos\theta} = \frac{3}{16}\sigma(\mu^+\mu^- \rightarrow bs) \left(1 + \cos^2\theta - \frac{8}{3}A_{\text{FB}}\cos\theta\right)$$

Need **charge tagging** to measure the forward backward asymmetry

Forward Backward Asymmetry and Charge Tagging

$$\frac{d\sigma(\mu^+\mu^- \rightarrow b\bar{s})}{d\cos\theta} = \frac{3}{16}\sigma(\mu^+\mu^- \rightarrow bs)\left(1 + \cos^2\theta + \frac{8}{3}A_{\text{FB}}\cos\theta\right)$$

$$\frac{d\sigma(\mu^+\mu^- \rightarrow \bar{b}s)}{d\cos\theta} = \frac{3}{16}\sigma(\mu^+\mu^- \rightarrow bs)\left(1 + \cos^2\theta - \frac{8}{3}A_{\text{FB}}\cos\theta\right)$$

Need **charge tagging** to measure the forward backward asymmetry

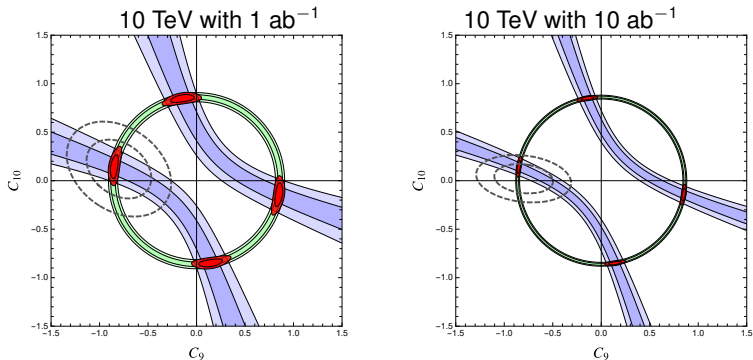
Imperfect charge tagging dilutes the forward backward asymmetry

$$A_{\text{FB}}^{\text{obs}} = (2\epsilon_{\pm} - 1) \left(\frac{N_{\text{sig}}}{N_{\text{tot}}} A_{\text{FB}} + \frac{N_{\text{bg}}}{N_{\text{tot}}} A_{\text{FB}}^{\text{bg}} \right)$$

As a benchmark, we assume charge tagging power as at LEP $\epsilon_{\pm} \simeq 70\%$

Sensitivity Projections

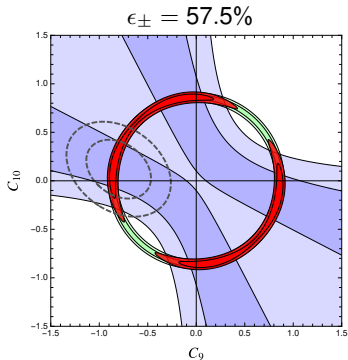
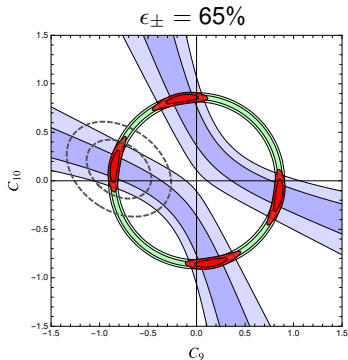
WA, Gadam, Profumo 2203.07495 and 2306.15017



- ▶ Branching ratio (green) and A_{FB} (blue) are complementary.
- ▶ In dashed: our global rare B decay fit.
- ▶ If there is new physics in $b \rightarrow sll$, a 10 TeV muon collider would clearly see it, and one does not need to worry about long distance QCD.

(see also Huang et al. 2103.01617; Asadi et al. 2104.05720; Azatov et al. 2205.13552)

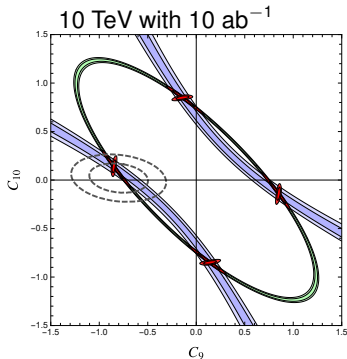
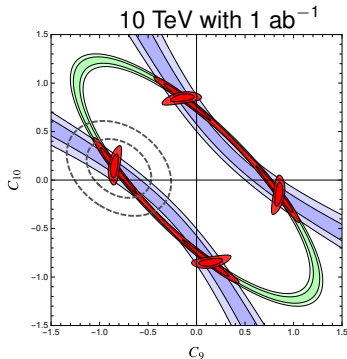
Impact of Charge Tagging



- ▶ The forward backward asymmetry gives useful information for charge tagging as low as $\sim 60\%$.
- ▶ For $\epsilon_{\pm} \lesssim 57.5\%$ two of the four red regions start to merge.

Impact of Beam Polarization

WA, Gadam, Profumo 2203.07495 and 2306.15017

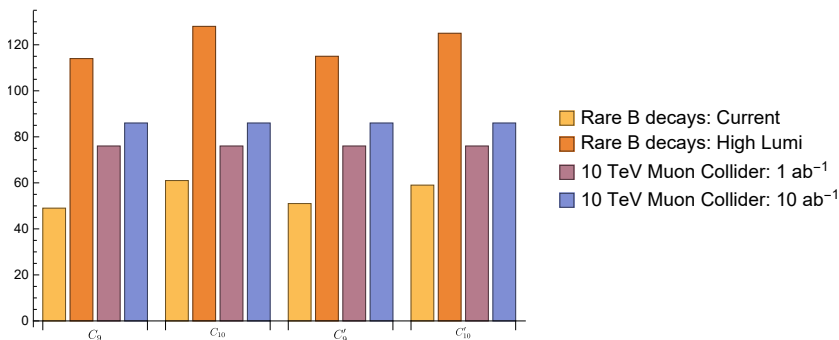


- ▶ So far had assumed that muon beams are unpolarized.
- ▶ Can expect a typical residual polarization of $\sim 20\%$ from pion decay. Higher polarization could be obtained at the cost of luminosity.
- ▶ Plots show the case of 50% polarization.

In the Absence of New Physics

WA, Gadam, Profumo 2203.07495 and 2306.15017

Λ/TeV

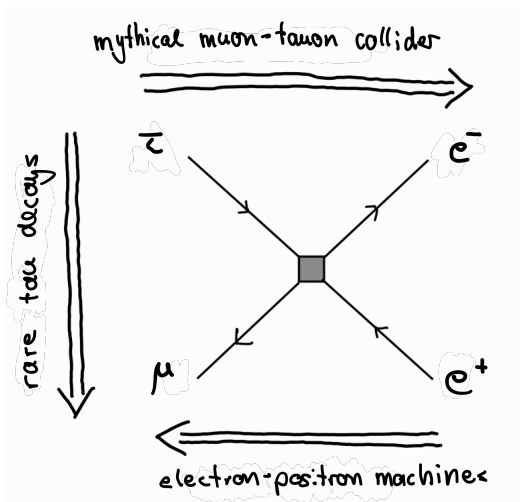


- ▶ In the absence of new physics, rare B decays and a 10 TeV muon collider have comparable sensitivity.
- ▶ Rare B decays have the advantage that a small new physics amplitude can interfere with the SM.
- ▶ At a muon collider one has to look for $|\text{new physics}|^2$.

Part 2:
Collider Probes of
Lepton Flavor Violation

based on 2305.03869 with P. Munbodh and T. Oh

Collider Probes of Lepton Flavor Violation



Lepton Flavor Violation

- ▶ In the SM, charged lepton flavor violation is suppressed by the tiny neutrino mass splittings

$$\text{e.g. } \text{BR}(\mu \rightarrow 3e) \sim \text{BR}(\mu \rightarrow e\nu_e\nu_\mu) \left| \frac{g^2}{16\pi^2} \frac{\Delta m_\nu^2}{m_W^2} \right|^2 \sim 10^{-50}$$

- ▶ Any observation in the foreseeable future would be an **unambiguous sign of new physics**.

Lepton Flavor Violation

- ▶ In the SM, charged lepton flavor violation is suppressed by the tiny neutrino mass splittings

$$\text{e.g. } \text{BR}(\mu \rightarrow 3e) \sim \text{BR}(\mu \rightarrow e\nu_e\nu_\mu) \left| \frac{g^2}{16\pi^2} \frac{\Delta m_\nu^2}{m_W^2} \right|^2 \sim 10^{-50}$$

- ▶ Any observation in the foreseeable future would be an **unambiguous sign of new physics**.
- ▶ Can search for lepton flavor violation in many different ways:
 - 1) At low energies in **lepton or hadron decays**: $\mu \rightarrow e\gamma$, $B_s \rightarrow \tau\mu$, ...

Lepton Flavor Violation

- ▶ In the SM, charged lepton flavor violation is suppressed by the tiny neutrino mass splittings

$$\text{e.g. } \text{BR}(\mu \rightarrow 3e) \sim \text{BR}(\mu \rightarrow e\nu_e\nu_\mu) \left| \frac{g^2}{16\pi^2} \frac{\Delta m_\nu^2}{m_W^2} \right|^2 \sim 10^{-50}$$

- ▶ Any observation in the foreseeable future would be an **unambiguous sign of new physics**.
- ▶ Can search for lepton flavor violation in many different ways:
 - 1) At low energies in **lepton or hadron decays**: $\mu \rightarrow e\gamma$, $B_s \rightarrow \tau\mu$, ...
 - 2) At high energies in **decays of heavy resonances**: $Z \rightarrow \mu e$, $h \rightarrow \tau\mu$, ...

Lepton Flavor Violation

- ▶ In the SM, charged lepton flavor violation is suppressed by the tiny neutrino mass splittings

$$\text{e.g. } \text{BR}(\mu \rightarrow 3e) \sim \text{BR}(\mu \rightarrow e\nu_e\nu_\mu) \left| \frac{g^2}{16\pi^2} \frac{\Delta m_\nu^2}{m_W^2} \right|^2 \sim 10^{-50}$$

- ▶ Any observation in the foreseeable future would be an **unambiguous sign of new physics**.
- ▶ Can search for lepton flavor violation in many different ways:
 - 1) At low energies in **lepton or hadron decays**: $\mu \rightarrow e\gamma$, $B_s \rightarrow \tau\mu$, ...
 - 2) At high energies in **decays of heavy resonances**: $Z \rightarrow \mu e$, $h \rightarrow \tau\mu$, ...
 - 3) At high energies in **non-resonant production**: $e^+e^- \rightarrow \tau\mu$, ...

- Generic scaling of a new physics effect with the flavor changing coupling g_{NP} and the new physics scale Λ_{NP}

$$\frac{\text{BR}(\mu \rightarrow 3e)}{\text{BR}(\mu \rightarrow e\nu_\mu\bar{\nu}_e)} \sim g_{\text{NP}}^2 \left(\frac{v}{\Lambda_{\text{NP}}} \right)^4 \lesssim 10^{-12}$$

$$\frac{\text{BR}(\tau \rightarrow 3\mu)}{\text{BR}(\tau \rightarrow \mu\nu_\mu\bar{\nu}_\tau)} \sim g_{\text{NP}}^2 \left(\frac{v}{\Lambda_{\text{NP}}} \right)^4 \lesssim 10^{-8}$$

New Physics Sensitivity of LFV at Low Energies

- ▶ Generic scaling of a new physics effect with the flavor changing coupling g_{NP} and the new physics scale Λ_{NP}

$$\frac{\text{BR}(\mu \rightarrow 3e)}{\text{BR}(\mu \rightarrow e\nu_\mu\bar{\nu}_e)} \sim g_{\text{NP}}^2 \left(\frac{v}{\Lambda_{\text{NP}}} \right)^4 \lesssim 10^{-12}$$

$$\frac{\text{BR}(\tau \rightarrow 3\mu)}{\text{BR}(\tau \rightarrow \mu\nu_\mu\bar{\nu}_\tau)} \sim g_{\text{NP}}^2 \left(\frac{v}{\Lambda_{\text{NP}}} \right)^4 \lesssim 10^{-8}$$

- ▶ For O(1) couplings, this corresponds to new physics scales of

$$\Lambda_{\text{NP}} \gtrsim 100 \text{ TeV} \quad \text{for muons}$$

$$\Lambda_{\text{NP}} \gtrsim 10 \text{ TeV} \quad \text{for taus}$$

New Physics Sensitivity of Heavy Resonance Decays

- Consider LFV decays of the Z boson, the Higgs, the top in the presence of generic new physics

$$\frac{\text{BR}(Z \rightarrow \mu e)}{\text{BR}(Z \rightarrow \mu\mu)} \sim g_{\text{NP}}^2 \left(\frac{v}{\Lambda_{\text{NP}}} \right)^4, \quad \frac{\text{BR}(H \rightarrow \tau\mu)}{\text{BR}(H \rightarrow \tau\tau)} \sim g_{\text{NP}}^2 \left(\frac{v}{\Lambda_{\text{NP}}} \right)^4$$

$$\frac{\text{BR}(t \rightarrow c\mu e)}{\text{BR}(t \rightarrow Wb)} \sim \frac{g_{\text{NP}}^2}{16\pi^2} \left(\frac{v}{\Lambda_{\text{NP}}} \right)^4$$

New Physics Sensitivity of Heavy Resonance Decays

- ▶ Consider LFV decays of the Z boson, the Higgs, the top in the presence of generic new physics

$$\frac{\text{BR}(Z \rightarrow \mu e)}{\text{BR}(Z \rightarrow \mu\mu)} \sim g_{\text{NP}}^2 \left(\frac{v}{\Lambda_{\text{NP}}} \right)^4, \quad \frac{\text{BR}(H \rightarrow \tau\mu)}{\text{BR}(H \rightarrow \tau\tau)} \sim g_{\text{NP}}^2 \left(\frac{v}{\Lambda_{\text{NP}}} \right)^4$$

$$\frac{\text{BR}(t \rightarrow c\mu e)}{\text{BR}(t \rightarrow Wb)} \sim \frac{g_{\text{NP}}^2}{16\pi^2} \left(\frac{v}{\Lambda_{\text{NP}}} \right)^4$$

- ▶ Same dependence on new physics as the low energy probes, but typically much **less Z, Higgs, top available in experiments.**
- ▶ Note: these are extremely generic/naive expectations; situation can be very different in concrete models.

[for a review see WA, Caillol, Dam, Xella, Zhang 2205.10576]

Example: LFV Z Decays

- ▶ Results from the LHC: ATLAS (139 fb^{-1})

Phys.Rev.Lett. 127 (2022) 271801; Nature Phys. 17 (2021) 7, 819-825; ATLAS-CONF-2021-042

$$\text{BR}(Z \rightarrow \mu e) < 3.04 \times 10^{-7}$$

$$\text{BR}(Z \rightarrow \tau e) < 5.0 \times 10^{-6}$$

$$\text{BR}(Z \rightarrow \tau \mu) < 6.5 \times 10^{-6}$$

- ▶ Slightly better than LEP bounds for all decay modes.
- ▶ In all searches there are backgrounds \Rightarrow expect sensitivities to improve with $\sqrt{\mathcal{L}}$, i.e. \sim factor of 5 at the HL-LHC.

Expected Sensitivities at Proposed Z Pole Machines

based on FCC-ee study Dam 1811.09408 (see also the FCC-ee whitepaper 2203.06520)

$Z \rightarrow \mu e$

- ▶ background from $Z \rightarrow \tau\tau \rightarrow \mu\nu\nu e\nu\nu$ is under control. Momentum resolution of 10^{-3} and Z mass constraint implies background rate of $\sim 10^{-11}$.
- ▶ main background: $Z \rightarrow \mu\mu$ where one muon suffers from “catastrophic” bremsstrahlung and is identified as electron.
- ▶ mis-id probability $\sim 10^{-7}$ limits the sensitivity to $\text{BR}(Z \rightarrow \mu e) \sim 10^{-8}$.
- ▶ With improved e/μ separation (dE/dx) might be able to go down to $\text{BR}(Z \rightarrow \mu e) \sim 10^{-10}$.

Expected Sensitivities at Proposed Z Pole Machines

based on FCC-ee study Dam 1811.09408 (see also the FCC-ee whitepaper 2203.06520)

$Z \rightarrow \mu e$

- ▶ background from $Z \rightarrow \tau\tau \rightarrow \mu\nu\nu e\nu\nu$ is under control. Momentum resolution of 10^{-3} and Z mass constraint implies background rate of $\sim 10^{-11}$.
- ▶ main background: $Z \rightarrow \mu\mu$ where one muon suffers from “catastrophic” bremsstrahlung and is identified as electron.
- ▶ mis-id probability $\sim 10^{-7}$ limits the sensitivity to $\text{BR}(Z \rightarrow \mu e) \sim 10^{-8}$.
- ▶ With improved e/μ separation (dE/dx) might be able to go down to $\text{BR}(Z \rightarrow \mu e) \sim 10^{-10}$.

$Z \rightarrow \tau e$
and
 $Z \rightarrow \tau\mu$

- ▶ minimize τ vs μ , e mis-id \rightarrow focus on hadronic taus
- ▶ background from $Z \rightarrow \tau_{\text{had}}\tau \rightarrow \tau_{\text{had}}\ell\nu\nu$
- ▶ limits sensitivity to $\text{BR}(Z \rightarrow \tau\ell) \sim 10^{-9}$

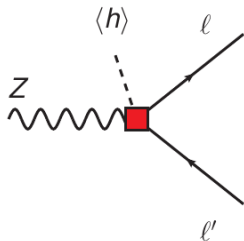
LFV Z Decays in the EFT Framework

- ▶ Parameterize New Physics in a systematic and controlled way: in terms of dim-6 operators of the SMEFT

dipoles

$$\mathcal{O}_{dW} = (\bar{\ell}\sigma^{\mu\nu}\tau^a P_R \ell') H W_{\mu\nu}^a$$

$$\mathcal{O}_{dB} = (\bar{\ell}\sigma^{\mu\nu} P_R \ell') H B_{\mu\nu}$$

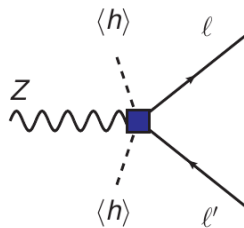


Higgs currents

$$\mathcal{O}_{hl}^{(3)} = (H^\dagger i \overleftrightarrow{D}_\mu^a H) (\bar{\ell} \gamma^\mu \tau^a P_L \ell')$$

$$\tilde{\mathcal{O}}_{hl}^{(1)} = (H^\dagger i \overleftrightarrow{D}_\mu H) (\bar{\ell} \gamma^\mu P_L \ell')$$

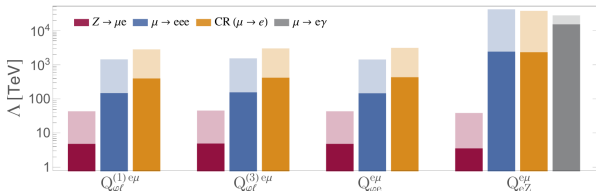
$$\mathcal{O}_{he} = (H^\dagger i \overleftrightarrow{D}_\mu H) (\bar{\ell} \gamma^\mu P_R \ell')$$



Comparison with Low Energy Probes

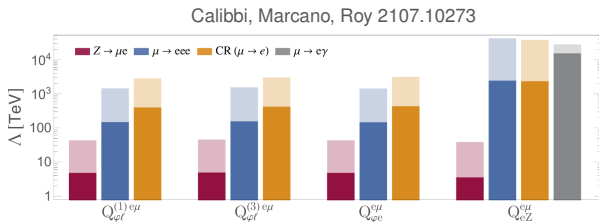
- ▶ Many flavor violating **low energy processes** will be affected as well.
- ▶ Severe indirect constraints on $Z \rightarrow \mu e$ from $\mu \rightarrow e\gamma$, $\mu \rightarrow 3e$, $\mu \rightarrow e$ conversion (barring accidental cancellations).

Calibbi, Marcano, Roy 2107.10273

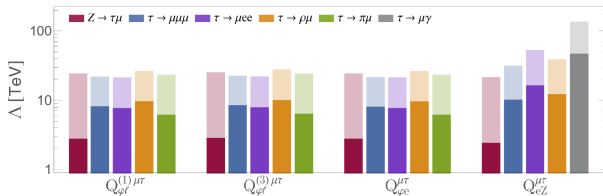


Comparison with Low Energy Probes

- ▶ Many flavor violating **low energy processes** will be affected as well.
- ▶ Severe indirect constraints on $Z \rightarrow \mu e$ from $\mu \rightarrow e\gamma$, $\mu \rightarrow 3e$, $\mu \rightarrow e$ conversion (barring accidental cancellations).



- ▶ **Complementary** sensitivity in the case of taus.



- ▶ The scaling of LFV cross sections with the center of mass energy depends on the type of operator:

$$\frac{\sigma(e^+e^- \rightarrow \tau\mu)}{\sigma(e^+e^- \rightarrow \tau^+\tau^-)} \sim$$

New Physics Sensitivity of Non-Resonant LFV

- ▶ The scaling of LFV cross sections with the center of mass energy depends on the type of operator:

$$\frac{\sigma(e^+e^- \rightarrow \tau\mu)}{\sigma(e^+e^- \rightarrow \tau^+\tau^-)} \sim g_{\text{NP}}^2 \left(\frac{v^4}{\Lambda_{\text{NP}}^4} \right),$$

New Physics Sensitivity of Non-Resonant LFV

- ▶ The scaling of LFV cross sections with the center of mass energy depends on the type of operator:

$$\frac{\sigma(e^+e^- \rightarrow \tau\mu)}{\sigma(e^+e^- \rightarrow \tau^+\tau^-)} \sim g_{\text{NP}}^2 \left(\frac{V^4}{\Lambda_{\text{NP}}^4} \right), \quad g_{\text{NP}}^2 \left(\frac{SV^2}{\Lambda_{\text{NP}}^4} \right),$$

New Physics Sensitivity of Non-Resonant LFV

- ▶ The scaling of LFV cross sections with the center of mass energy depends on the type of operator:

$$\frac{\sigma(e^+e^- \rightarrow \tau\mu)}{\sigma(e^+e^- \rightarrow \tau^+\tau^-)} \sim g_{\text{NP}}^2 \left(\frac{V^4}{\Lambda_{\text{NP}}^4} \right), g_{\text{NP}}^2 \left(\frac{SV^2}{\Lambda_{\text{NP}}^4} \right), g_{\text{NP}}^2 \left(\frac{S^2}{\Lambda_{\text{NP}}^4} \right)$$

- ▶ For some operators one will have **enhanced sensitivity at high energies**. (Assuming one does not resolve the higher dimensional operators.)
- ▶ How sensitive is one to $\tau\mu$ production at future e^+e^- colliders?

New Physics Sensitivity of Non-Resonant LFV

- ▶ The scaling of LFV cross sections with the center of mass energy depends on the type of operator:

$$\frac{\sigma(e^+e^- \rightarrow \tau\mu)}{\sigma(e^+e^- \rightarrow \tau^+\tau^-)} \sim g_{\text{NP}}^2 \left(\frac{V^4}{\Lambda_{\text{NP}}^4} \right), g_{\text{NP}}^2 \left(\frac{SV^2}{\Lambda_{\text{NP}}^4} \right), g_{\text{NP}}^2 \left(\frac{S^2}{\Lambda_{\text{NP}}^4} \right)$$

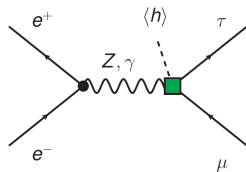
- ▶ For some operators one will have **enhanced sensitivity at high energies**. (Assuming one does not resolve the higher dimensional operators.)
- ▶ How sensitive is one to $\tau\mu$ production at future e^+e^- colliders?
- ▶ In **WA, Munbodh, Oh 2305.03869** we show that high-energy runs of FCC-ee/CEPC have sensitivity that is comparable and complementary to other probes.

Systematic SMEFT Parameterization of New Physics

dipoles

$$\mathcal{O}_{dW} = (\bar{\tau} \sigma^{\alpha\beta} T^a P_R \mu) H W_{\alpha\beta}^a$$

$$\mathcal{O}_{dB} = (\bar{\tau} \sigma^{\alpha\beta} P_R \mu) H B_{\alpha\beta}$$

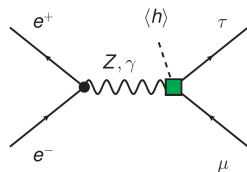


Systematic SMEFT Parameterization of New Physics

dipoles

$$\mathcal{O}_{dW} = (\bar{\tau} \sigma^{\alpha\beta} T^a P_R \mu) H W_{\alpha\beta}^a$$

$$\mathcal{O}_{dB} = (\bar{\tau} \sigma^{\alpha\beta} P_R \mu) H B_{\alpha\beta}$$

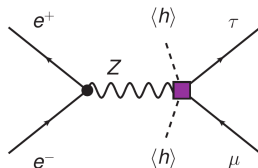


Higgs currents

$$\mathcal{O}_{hl}^{(3)} = (H^\dagger i \overleftrightarrow{D}_\alpha^a H) (\bar{\tau} \gamma^\alpha T^a P_L \mu)$$

$$\mathcal{O}_{hl}^{(1)} = (H^\dagger i \overleftrightarrow{D}_\alpha H) (\bar{\tau} \gamma^\alpha P_L \mu)$$

$$\mathcal{O}_{he} = (H^\dagger i \overleftrightarrow{D}_\alpha H) (\bar{\tau} \gamma^\alpha P_R \mu)$$

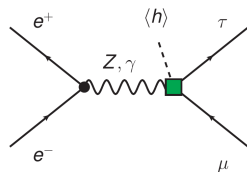


Systematic SMEFT Parameterization of New Physics

dipoles

$$\mathcal{O}_{dW} = (\bar{\tau} \sigma^{\alpha\beta} T^a P_R \mu) H W_{\alpha\beta}^a$$

$$\mathcal{O}_{dB} = (\bar{\tau} \sigma^{\alpha\beta} P_R \mu) H B_{\alpha\beta}$$

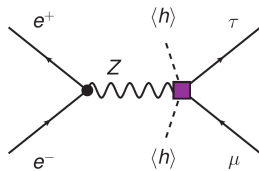


Higgs currents

$$\mathcal{O}_{hl}^{(3)} = (H^\dagger i \overleftrightarrow{D}_\alpha^a H) (\bar{\tau} \gamma^\alpha T^a P_L \mu)$$

$$\mathcal{O}_{hl}^{(1)} = (H^\dagger i \overleftrightarrow{D}_\alpha H) (\bar{\tau} \gamma^\alpha P_L \mu)$$

$$\mathcal{O}_{he} = (H^\dagger i \overleftrightarrow{D}_\alpha H) (\bar{\tau} \gamma^\alpha P_R \mu)$$



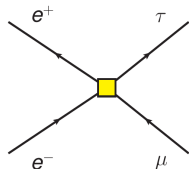
$$\mathcal{O}_{\ell\ell} = (\bar{e} \gamma^\alpha P_L e) (\bar{\tau} \gamma_\alpha P_L \mu)$$

$$\mathcal{O}_{ee} = (\bar{e} \gamma^\alpha P_R e) (\bar{\tau} \gamma_\alpha P_R \mu)$$

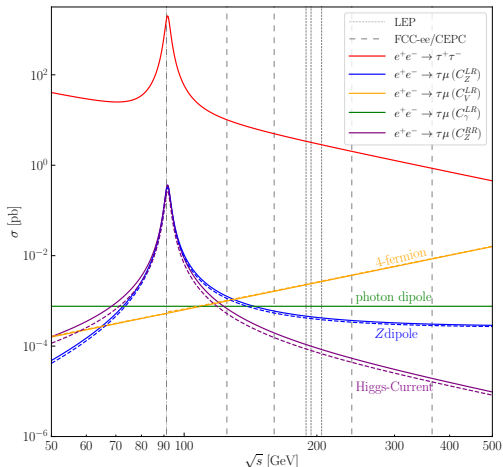
$$\mathcal{O}_{\ell e} = (\bar{e} \gamma^\alpha P_L e) (\bar{\tau} \gamma_\alpha P_R \mu)$$

$$\mathcal{O}_{e\ell} = (\bar{e} \gamma^\alpha P_R e) (\bar{\tau} \gamma_\alpha P_L \mu)$$

4-fermion contact interactions



Dependence on the Center of Mass Energy



WA, Munbodh, Oh 2305.03869
 (in the plot $\Lambda_{NP} = 3 \text{ TeV}$, $C_i = 1$)

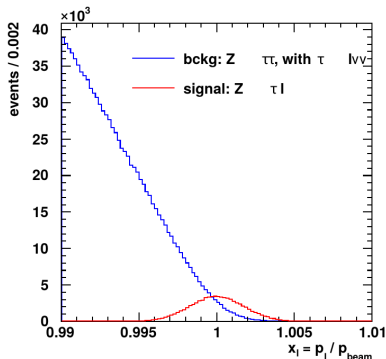
- ▶ $\tau^+\tau^-$ background falls like $1/s$
- ▶ $\tau\mu$ production increases linearly with s for 4-fermion operators
- ▶ $\tau\mu$ production is flat in s for dipole operators
- ▶ $\tau\mu$ production falls like $1/s$ for Higgs current operators
- ▶ resonance at $s = m_Z^2$ if Z-mediated

Signal and Most Important Background

signal: $e^+e^- \rightarrow \tau\mu$

bkg: $e^+e^- \rightarrow \tau^+\tau^- \rightarrow \tau\mu\nu\nu$

- ▶ **Signal** is a sharp peak at $x = p_\mu/p_{\text{beam}} = 1$
- ▶ **Background** is a smooth distribution with $x \lesssim 1$
- ▶ Width of the signal peak and spread of background to $x > 1$ is determined by the beam energy spread and the muon momentum resolution.

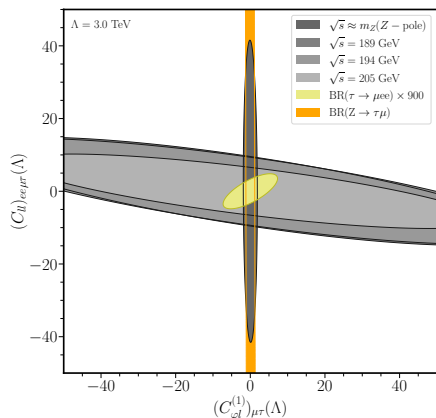


Dam 1811.09408
(study on the Z peak)

- ▶ Impact of initial state radiation? (work in progress with Munbodh)

Existing Constraints from LEP

WA, Munbodh, Oh 2305.03869



- ▶ LEP has searched for $e^+ e^- \rightarrow \tau \mu$ at the Z pole (e.g. OPAL Z.Phys.C 67 (1995) 555-564) and at $\sqrt{s} \sim 200 \text{ GeV}$ (OPAL PLB 519, (2001) 23-32).
- ▶ Z pole search mainly sensitive to the Higgs current operators.
- ▶ High \sqrt{s} search mainly sensitive to 4-fermion operators.
- ▶ LEP searches have sensitivity comparable to $Z \rightarrow \tau \mu$ at the LHC, but cannot compete with tau decays.

Projections for FCC-ee

machine and detector parameters from FCC-ee CDR vol. 2, 1909.12245, 2107.02686, 2203.06520

\sqrt{s} [GeV]	\mathcal{L}_{int} [ab $^{-1}$]	$\frac{\delta\sqrt{s}}{\sqrt{s}}$ [10 $^{-3}$]	$\frac{\delta p_T}{p_T}$ [10 $^{-3}$]	$\epsilon_{\text{bkg}}^{x_c}$ [10 $^{-6}$]	N_{bkg}	σ [ab]
91.2 (Z -pole)	75	0.93	1.35	1.55	9700 ± 100	45
87.7 (off-peak)	37.5	0.93	1.33	1.46	520 ± 20	21
93.9 (off-peak)	37.5	0.93	1.37	1.59	930 ± 30	28
125 (H)	20	0.03	1.60	1.44	12 ± 3	8
160 (WW)	12	0.93	1.89	2.44	6 ± 2	10
240 (ZH)	5	1.17	2.60	4.39	2 ± 1	18
365 ($t\bar{t}$)	1.5	1.32	3.78	8.61	0.5 ± 0.7	50

- ▶ Estimate background efficiency by imposing a cut $x > 1$. (could be further optimized)
- ▶ Expect sizable background on the Z -peak, very few background events at higher energies.
- ▶ Can achieve sensitivity to $e^+e^- \rightarrow \tau\mu$ cross sections of $\mathcal{O}(10 \text{ ab})$.

Projections for CEPC

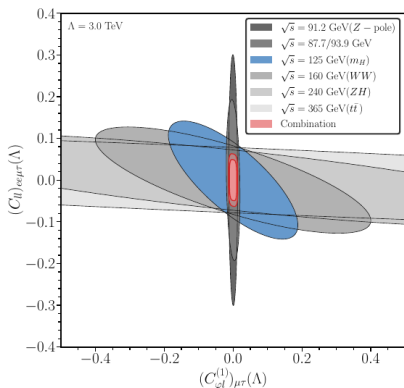
machine and detector parameters from 1809.00285, 1811.10545, 2203.09451, 2205.08553

\sqrt{s} [GeV]	\mathcal{L}_{int} [ab ⁻¹]	$\frac{\delta\sqrt{s}}{\sqrt{s}}$ [10 ⁻³]	$\frac{\delta p_T}{p_T}$ [10 ⁻³]	$\epsilon_{\text{bkg}}^{x_c}$ [10 ⁻⁶]	N_{bkg}	σ [ab]
91.2 (<i>Z</i> -pole)	50	0.92	1.35	1.53	6400 ± 80	55
87.7 (off-peak)	25	0.92	1.33	1.46	350 ± 20	27
93.9 (off-peak)	25	0.92	1.37	1.59	620 ± 25	35
160 (<i>WW</i>)	6	0.99	1.89	2.49	3 ± 2	17
240 (<i>ZH</i>)	20	1.20	2.60	4.42	7 ± 3	6.6
360 (<i>t\bar{t}</i>)	1	1.41	3.74	8.61	0.3 ± 0.5	72

- ▶ Estimate background efficiency by imposing a cut $x > 1$.
(could be further optimized)
- ▶ Expect sizable background on the *Z*-peak, very few background events at higher energies.
- ▶ Can achieve sensitivity to $e^+e^- \rightarrow \tau\mu$ cross sections of $\mathcal{O}(10 \text{ ab})$.

Complementarity of Different Observables (FCC-ee)

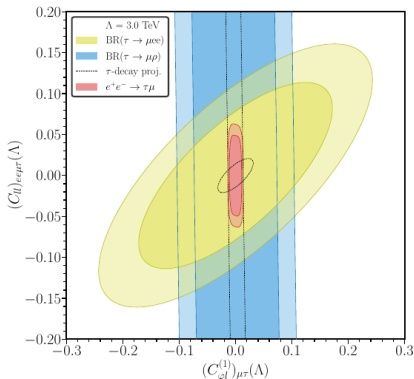
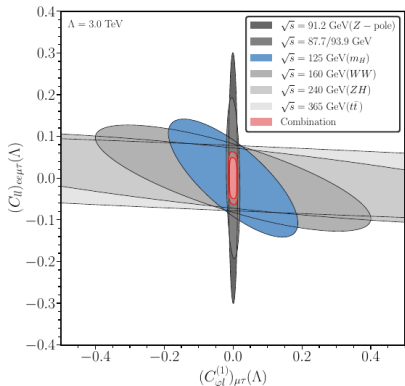
WA, Munbodh, Oh 2305.03869



- As in the case of LEP, the Z -pole searches and the high- \sqrt{s} searches are **complementary**.

Complementarity of Different Observables (FCC-ee)

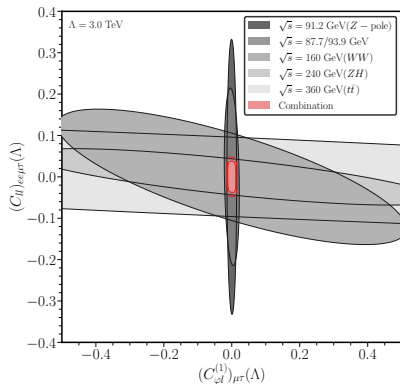
WA, Munbooth, Oh 2305.03869



- ▶ As in the case of LEP, the Z -pole searches and the high- \sqrt{s} searches are **complementary**.
- ▶ Expected **FCC-ee sensitivity** rivals the one from current and future searches for **LFV τ decays**.

Complementarity of Different Observables (CEPC)

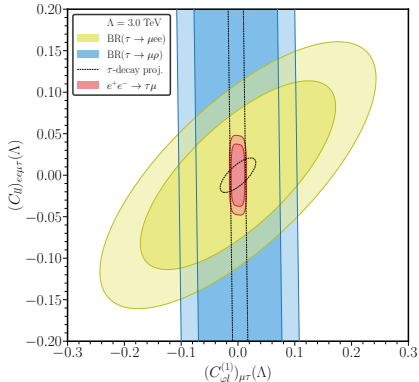
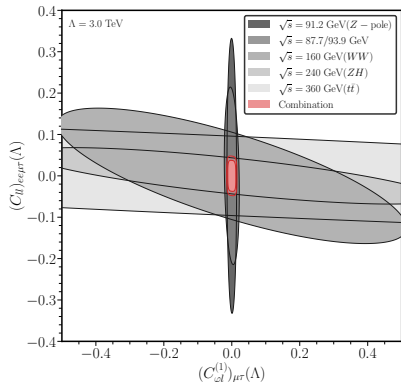
WA, Munbodh, Oh 2305.03869



- ▶ As in the case of LEP, the Z -pole searches and the high- \sqrt{s} searches are **complementary**.

Complementarity of Different Observables (CEPC)

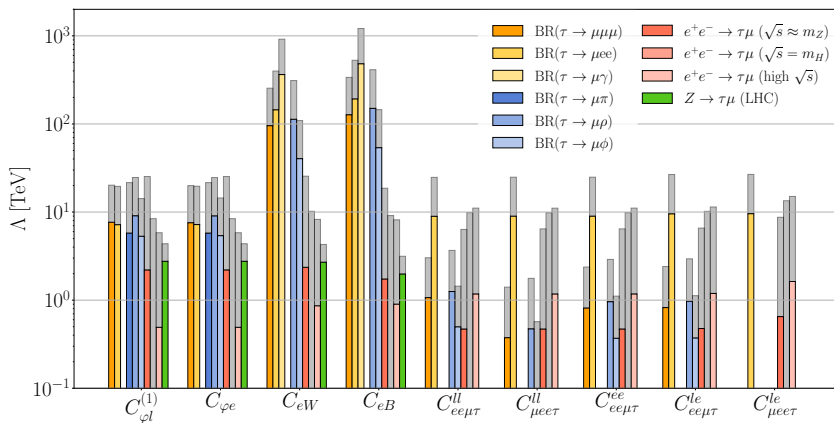
WA, Munbodh, Oh 2305.03869



- ▶ As in the case of LEP, the Z -pole searches and the high- \sqrt{s} searches are **complementary**.
- ▶ Expected **CEPC sensitivity** rivals the one from current and future searches for **LFV τ decays**.

Summary of Generic Sensitivities

WA, Munbodh, Oh 2305.03869



If a Signal is Seen ...

- ▶ If a signal is seen at one \sqrt{s} :
⇒ look at different \sqrt{s} to identify the operator class
(dipole, Higgs current, 4-fermion)

If a Signal is Seen ...

- ▶ If a signal is seen at one \sqrt{s} :
⇒ look at different \sqrt{s} to identify the operator class (dipole, Higgs current, 4-fermion)
- ▶ The signal can be further characterized by **angular distributions** (θ = angle between the beam axis and the outgoing muon) and **CP asymmetries** ($\tau^+ \mu^-$ vs. $\tau^- \mu^+$)

$$\frac{1}{\sigma_{\text{tot}}} \frac{d(\sigma + \bar{\sigma})}{d \cos \theta} = \frac{3}{8}(1 - F_D)(1 + \cos^2 \theta) + A_{\text{FB}} \cos \theta + \frac{3}{4}F_D \sin^2 \theta ,$$

$$\frac{1}{\sigma_{\text{tot}}} \frac{d(\sigma - \bar{\sigma})}{d \cos \theta} = \frac{3}{8}(A^{\text{CP}} - F_D^{\text{CP}})(1 + \cos^2 \theta) + A_{\text{FB}}^{\text{CP}} \cos \theta + \frac{3}{4}F_D^{\text{CP}} \sin^2 \theta ,$$

- ▶ For a sufficiently large signal, it might be possible to significantly narrow down the **chirality structure of the operator** that is responsible for $e^+ e^- \rightarrow \tau \mu$

- ▶ $\mu^+\mu^- \rightarrow bs$ at a 10 TeV muon collider is an interesting probe of new physics.
- Could test the B anomalies without having to worry about hadronic effects.
- In the absence of new physics, could probe $(\mu\mu)(bs)$ contact interactions at scales of ~ 80 TeV.

- ▶ $\mu^+ \mu^- \rightarrow bs$ at a 10 TeV muon collider is an interesting probe of new physics.
 - Could test the B anomalies without having to worry about hadronic effects.
 - In the absence of new physics, could probe $(\mu\mu)(bs)$ contact interactions at scales of ~ 80 TeV.
- ▶ $e^+ e^- \rightarrow \tau\mu$ offers interesting opportunities to probe lepton flavor violation at FCC-ee/CEPC.
 - Different LFV operators show characteristic dependence on the center of mass energy.
 - Estimated sensitivity rivals the one from rare tau decays.

Back Up

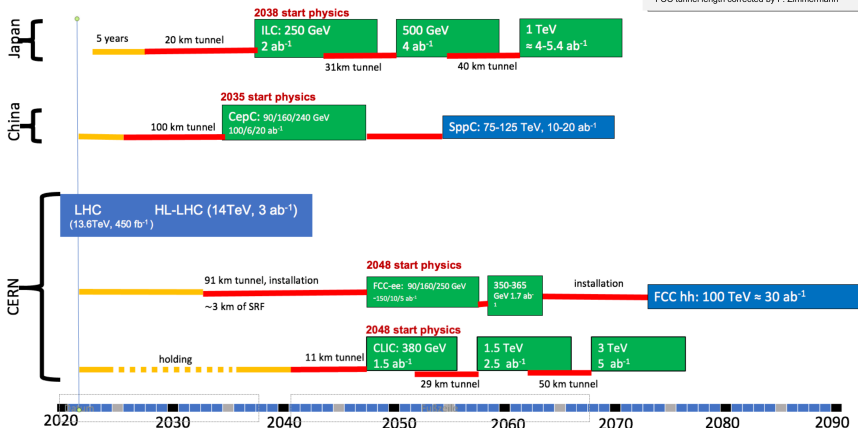
Future Colliders

Indicative scenarios of future colliders [considered by ESG]

- Proton collider
- Electron collider
- Muon collider

- Construction/Transformation
- Preparation / R&D

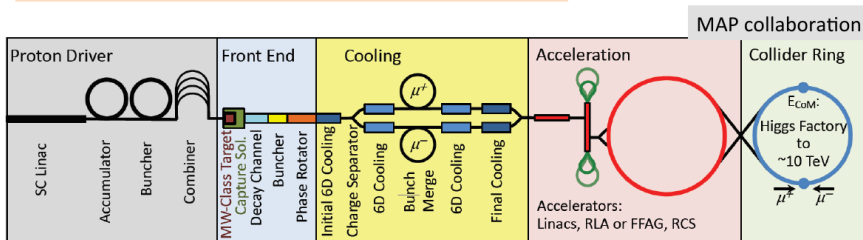
Original from ESPP by Ursula Bassler
 Updated July 25, 2022 by Meenkshi Narain
 FCC tunnel length corrected by F. Zimmermann



Karl Jacobs (ECFA chair) @ 2nd ECFA meeting on e^+e^- Higgs, electroweak, and top factories
 Oct 11-13, 2023, Paestum, Italy

A Muon Collider?

Muon collider design is driven by finite muon lifetime



Short, intense proton bunches to produce hadronic showers

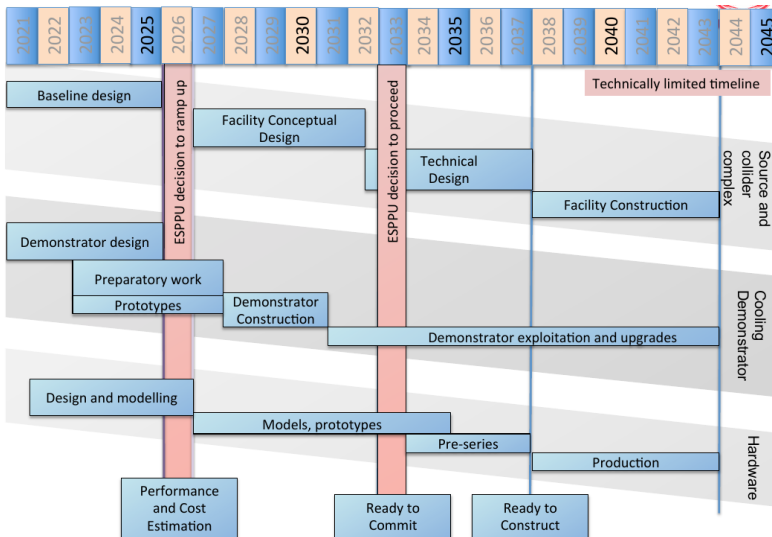
Muon are captured, bunched and then cooled by ionisation cooling in matter

Acceleration to collision energy

Protons produce pions
Pions decay to muons

talk by D. Schulte @ Muon Collider Agora, Feb 16 2022

A Muon Collider!



talk by D. Schulte @ Muon Collider Agora, Feb 16 2022

Another $\tau\mu$ Background at High Energies?

$$e^+e^- \rightarrow W^+W^- \rightarrow \tau\mu\nu\nu$$

- ▶ Muon momentum does not extend all the way to $x = 1$
- ▶ Decay kinematics is such that

$$x < \frac{1}{2} \left(1 + \sqrt{1 - \frac{4m_W^2}{s}} \right) < 1$$

- ▶ e.g. for $\sqrt{s} = 240$ GeV one has $x \lesssim 0.87$

⇒ this background is **not an issue**.